

**2DBTOR -- A TOROIDAL GEOMETRY
NEUTRON DIFFUSION CODE**

A Thesis

by

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August 1990

Major Subject: Nuclear Engineering

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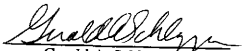
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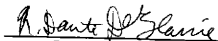
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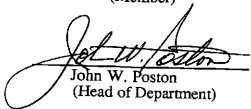
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ABSTRACT

2DBTOR -- A Toroidal Geometry Neutron Diffusion Code (August 1990)

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The objective of the research performed here was to produce a scoping code that could be used for fusion reactor blanket design. To this end, the present research initially explored a technique proposed by Pomraning and Stevens, in which the toroidal diffusion problem in toroidal geometry is cast into cylindrical (r - θ) form by a spatially dependent redefinition of the diffusion coefficient, absorption cross-section, and extraneous source function. This idea was explored but was later abandoned in favor of the direct finite differencing of the toroidal diffusion equation.

The direct finite differencing approach was programmed into an existing two-dimensional(x - y , r - z , r - θ , triangular), multi-group neutron diffusion code, 2DB, that had previously been converted to execute on the IBM-AT. Neutronic scoping calculations relevant to fusion reactor design were then performed in a micro-computer environment. The modified code was renamed 2DBTOR.

To verify that 2DBTOR was operating correctly, comparisons were made to both analytical and numerical solutions for several types of problems. Both ANISN and 2DB were used to verify and compare the solutions obtained from 2DBTOR. It was also shown that as the aspect ratio approached infinity (i. e., the major radius became large) the 2DBTOR solution approached the solution for that of 1-D cylindrical geometry. After verifying the solution for a large major radius, the errors associated

with using a non-toroidal scoping code were examined versus using 2DBTOR. To accomplish this, neutron cross-sections for a benchmark problem were input to 2DBTOR and the output was compared to that from ANISN. A method proposed by Price and Chapin, that used volume correction factors to compute the reaction rates in the benchmark blanket, was utilized to provide a means of checking 2DBTOR's results versus those given by a Monte Carlo code. It is also worth noting that 2DBTOR makes possible the calculation of material depletion in the fusion blanket, which is a unique advantage of the new program, 2DBTOR.

In future versions of the 2DBTOR program, it is recommended that the central vacuum should be modelled through an internal boundary condition. A separate void streaming calculation should be used to define the internal boundary condition by specifying the neutron flux to current ratio as a function of position along the vacuum wall. Improved modelling of the central void region will be required if 2DBTOR is to prove to be an attractive program for Tokamak blanket scoping calculations.

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TABLE OF CONTENTS

	Page
I. INTRODUCTION.....	1
I.A The Fusion Process.....	2
I.B The Need for Scoping Codes to Study the Fusion Process.....	5
I.C Previous Studies.....	6
I.D Computer Codes Used in This Work.....	7
I.E Objectives of This Work.....	8
II. THEORY.....	10
II.A Introduction.....	10
II.B Pomraning and Stevens' Method.....	22
II.C Finite Difference Approximation to the Diffusion Equation in Cylindrical Coordinates.....	24
II.D Finite Difference Approximation to the Diffusion Equation in Toroidal Coordinates.....	33
II.E Analytical Solutions.....	37
III. RESULTS.....	40
III.A 2DBTOR Verification and Infinite Cylinder Problems.....	40
III.B 2DBTOR Verification and Toroidal Problems.....	43
III.C Standard Blanket Solutions.....	46
III.C.1 Infinite Cylinder.....	46
III.C.2 Comparison of Torus and Infinite Cylinder Models.....	49
IV. SUMMARY AND CONCLUSIONS.....	64

TABLE OF CONTENTS (CONT.)

	Page
REFERENCES.....	66
APPENDIX A. Pomraning and Stevens' Analytical Solution.....	69
APPENDIX B. Average D_k Solution.....	72
APPENDIX C. XPROC2DB Listing.....	75
APPENDIX D. 2DB Errata.....	79
APPENDIX E. Plasma Source.....	81
APPENDIX F. Equivalent Source to the Plasma Source at the Inner Edge of the Blanket.....	83
APPENDIX G. 2DBTOR Manual.....	84
APPENDIX H. 2DBTOR Listing.....	147
VITA.....	237

LIST OF FIGURES

Figure	Page
1-1 Reactivity Curves.....	4
2-1 Toroidal Coordinate Description.....	20
2-2 Finite Difference Coordinates Used in 2DB and 2DBTOR.....	25
3-1 Flux versus Poloidal Angle at a Radius of 150 cm for Several Aspect Ratios.....	44
3-2 Flux versus Poloidal Angle at a Radius of 250 cm for Several Aspect Ratios.....	45
3-3 Standard Blanket Configuration.....	47

LIST OF TABLES

Table	Page
2-1 Area and Volume Elements in Cylindrical and Toroidal Coordinates.....	38
3-1 Standard Blanket Constituents.....	48
3-2 Radial Averaged Tritium Breeding Ratio's for 2DB and ANISN, Infinite Cylinder Case.....	50
3-3 Radial Averaged Tritium Breeding Ratio's for 2DBTOR Versus 2DB, Torus Case (Aspect Ratio = 3).....	52
3-4 Radial Averaged Tritium Breeding Ratio's for 2DBTOR Versus 2DB, Torus Case (Aspect Ratio = 5).....	53
3-5 Zone by Zone Averaged T6 for 2DBTOR Versus 2DB, Torus Case (Aspect Ratio = 3).....	54
3-6 Zone by Zone Averaged T7 for 2DBTOR Versus 2DB, Torus Case (Aspect Ratio = 3).....	55
3-7 Zone by Zone Averaged T6 for 2DBTOR Versus 2DB, Torus Case (Aspect Ratio = 5).....	56
3-8 Zone by Zone Averaged T7 for 2DBTOR Versus 2DB, Torus Case (Aspect Ratio = 5).....	57
3-9 T6 for 2DBTOR Versus Volume Corrected T6 for 2DB, Torus Case (Aspect Ratio = 3).....	60
3-10 T7 for 2DBTOR Versus Volume Corrected T7 for 2DB, Torus Case (Aspect Ratio = 3).....	61

Table	Page
3-11 T6 for 2DBTOR Versus Volume Corrected T6 for 2DB, Torus Case (Aspect Ratio = 5).....	62
3-12 T7 for 2DBTOR Versus Volume Corrected T7 for 2DB, Torus Case (Aspect Ratio = 5).....	63

I. INTRODUCTION

In order to sustain the current level of world civilization man requires energy. However, at the present pace of energy demand, it is predicted that coal and U-235 resources will be depleted in the next 25 to 100 years.¹ To meet the world's energy demand as the above resources become more scarce, new options will have to be employed. Of the methods available to produce energy, only solar energy, fission breeders, and fusion are thought to be capable of meeting the world's long-term needs. Of these three energy resources, fission breeders produce and use special nuclear materials, which makes them undesirable without stringent safeguards. This leaves solar energy and fusion as the most viable energy resources to meet future energy demands. Both will have to be developed, since a stable society needs alternate sources of energy to call upon in an uncertain future.

Because of progress in plasma confinement, fusion has warranted increased attention in the past decade. What makes fusion of most interest is the fact that one of the fuels required, deuterium, is essentially an inexhaustible resource. Deuterium, which has an average abundance of 0.015% in elemental hydrogen, can be separated relatively easily and cheaply from water.² A secondary, but perhaps more important advantage, as regards public perception, is fusion's inherent safety and reduced radioactivity hazard relative to fission. Tritium, one of the more radioactive of the fusion fuels, has a relatively short half-life(12.36y) and decays by emitting low energy β -rays. Unfortunately, activation of structural materials, such as the first wall in a fusion reactor, presents a possible hazard, but this can be controlled by choosing

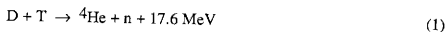
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suitable materials. In regard to safety, with magnetic confinement fusion, a sudden increase in power is likely to be counteracted by altering the conditions that are necessary to position and heat the plasma. In addition to the above advantages, use of exotic fuels which only produce charged particles, might lead to energy conversion efficiencies approaching 100%, since charged particles can possibly be directly converted to electricity.

I.A The Fusion Process

Fusion is essentially the process of two nuclei coming together to form one or more nuclei with an accompanying release of energy. Since very high temperatures (over 10^8 °C) are required to overcome the Coulomb repulsion between the reacting nuclei, a plasma must be produced. The nuclei, which are stripped of electrons at such temperatures, must be confined long enough to fuse. As a consequence of the high temperatures required for fusion, the reactants can not be contained within physical walls, since interactions with the wall material would likely cool the nuclei down below their required temperature for fusion. By virtue of the nuclei being charged particles, however, they can be contained by magnetic fields in various plasma confinement configurations.

Several reactions are considered to be possible for producing power from fusion break-even conditions, but the deuterium-tritium (D-T) reaction has the best chance of being the first to reach the required plasma conditions. The D-T reaction is described below:

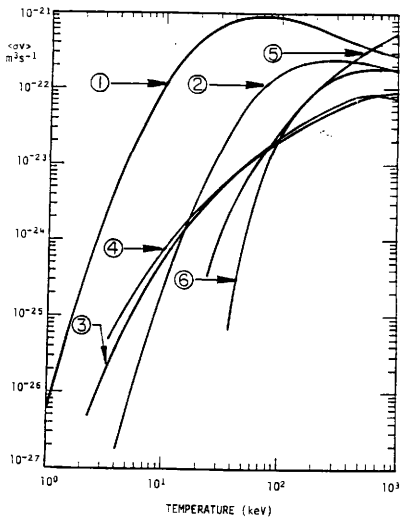


This reaction requires the use of tritium, which is extremely rare in nature and must be artificially produced. Tritium production is expected to be accomplished by surrounding the plasma with a lithium containing blanket. Tritium is produced from lithium according to the neutron reactions that follow:



The reaction with Li-7 has a threshold of 2.7 MeV, while the Li-6 reaction takes place with neutrons down to thermal energies. Although the above reactions require consumption of lithium, the known U.S. reserves of lithium are thought to be able to last 600 years.³ The D-T reaction discussed above has great potential for fusion due to its large reactivity at relatively low temperatures (<100 keV; 1 keV = 10⁹ °C). The reactivity, $\langle\sigma v\rangle_{jk}$, is defined to be the average of the fusion cross-section multiplied by the relative speed between the reacting particles j and k averaged over Maxwellian distributions. Curves of the reactivity versus energy for several fuels are shown in Figure 1-14. As can be seen, the D-T reactivity is by far the highest at temperatures below 100 keV. Since the reaction rate per unit volume for two nuclei can be shown to be given by

$$RR_{jk} = a_{jk} n_j n_k \langle\sigma v\rangle_{jk} \quad (4)$$



Reaction rate parameters for mixtures of Maxwellian distributions at the same temperature. (1) $\text{D} + \text{T} \rightarrow \text{n} + {}^4\text{He}$, (2) $\text{D} + {}^3\text{He} \rightarrow \text{H} + {}^4\text{He}$, (3) $\text{D} + \text{D} \rightarrow \text{H} + \text{T}$, (4) $\text{T} + \text{T} \rightarrow {}^4\text{He} + 2\text{n}$, (5) $\text{T} + {}^3\text{He} \rightarrow$ (various products), (6) $\text{H} + {}^{11}\text{B} \rightarrow 3({}^4\text{He})$.

Figure 1-1. Reactivity Curves.

where RR_{jk} = reaction rate per unit volume for nuclei j and k ,

$$a_{jk} = 1 \text{ for } j \neq k \text{ and } 1/2 \text{ for } j=k,$$

$$n_j = \text{atom density of nuclei } j,$$

$$\text{and } n_k = \text{atom density of nuclei } k,$$

it is apparent that the D-T reaction for fixed reactant densities will provide the greatest number of fusions per cm^3 per sec at the lowest temperatures.⁵ Since lowering of the temperature needed for fusion makes the required plasma conditions more achievable, it is likely that the first fusion reactions will use D-T fuel.

I.B The Need for Scoping Codes to Study the Fusion Process

In the past decade, as plasma confinement experiments have advanced nearer to break-even, there has been a concurrent interest in the engineering problems associated with power production from nuclear fusion. Since a magnetically confined plasma burning D-T fuel should be a reality within the next decade, calculations of the transport of 14 MeV neutrons through the first wall and blanket of a fusion reactor have become particularly pertinent. Accurate computations of neutron transport are currently feasible; however, these calculations can be labor intensive and expensive. To evaluate the myriad of possibilities for the first wall and blanket, scoping codes to estimate reaction rates and leakage could be advantageous for design purposes.

Engineers who do neutronic computations for fission reactors commonly use neutron transport and diffusion codes to determine the flux profiles within a nuclear reactor. By using a neutron diffusion code, scoping runs can be made to provide an

initial estimate for the actual flux profiles. Since diffusion calculations can be performed more easily than the more exact calculations based on transport theory, neutron diffusion codes are used to decrease the computational effort involved in the trial and error process that takes place in designing the cores of nuclear reactors. In the same manner, a need has arisen to perform flux computations for the nuclear design of Tokamak-type fusion reactors which have a toroidal or doughnut shape.

I.C Previous Studies

Several computational models have been applied to estimate the flux profile in toroidally shaped blankets. One approach is to perform diffusion or transport calculations in one-dimensional cylindrical geometry and to ignore the curvature associated with the major radius. In particular, the geometry of the toroidal reactor blanket can be approximated as an infinite right circular cylinder, so that a one-dimensional discrete ordinates code, such as ANISN⁶, can be used to solve for the neutron flux in the radial direction. This method is advantageous in that one-dimensional codes are readily available, easy to use, and relatively inexpensive to run. On the other hand, because one-dimensional cylindrical diffusion and transport codes can not model the curvature associated with the major radius present in Tokamaks, such codes are probably not suitable for tori with small aspect ratios.

To take into account the curvature of the major radius as well as the minor radius, two-dimensional cylindrical calculations in r - z geometry have been used to solve for the neutron flux. A typical code for this method is TWOTRAN-II⁷, a two-dimensional discrete ordinates transport code which can be run in r - z geometry. Although it is possible to model the curvature of the major radius adequately using such a code,

small mesh spacings are required to correctly model curved surfaces in the r-z plane; this can lead to costly and time consuming computations. More recently, computer programs have been developed to perform neutron transport calculations that are more suitable for toroidal geometry. One example of such a code is TRISM⁸, a two-dimensional discrete ordinates transport code developed at Los Alamos National Laboratory. Finally, full three-dimensional and Monte Carlo⁹ programs have also been applied to blanket problems. While these more advanced methods allow fluxes to be calculated accurately, they can be quite cumbersome and expensive to use.

I.D Computer Codes Used in This Work

2DB, a code written primarily for use in fast reactor calculations, provides two options that are potentially useful for fusion reactor design.¹⁰ First, a fixed source option will enable modeling of fusion blanket problems. Second, material depletion in the blanket of a Tokamak can potentially be calculated by using 2DB's material burn-up option.

ANISN, as stated previously, is a neutron transport code commonly used for fusion blanket analysis in one dimension. It can also be used to provide for a benchmark comparison to the one-dimensional results obtained with 2DB. In addition, ANISN enables one to reduce the number of groups used in the calculations by averaging over the fine groups to produce broad group cross sections. Reducing the number of groups will prove useful in a neutron diffusion code, since run times can be greatly decreased, thus facilitating its use on a micro-computer.

Since ANISN is a transport code, the results for the flux distribution will be more accurate than those obtained based on diffusion theory. In particular, if the validity of

diffusion theory in the vacuum region of a Tokamak inordinately affects the flux distribution in the rest of the Tokamak, then results from ANISN will provide a means to check this in 2DB.

The neutron cross section data used in the above codes will be taken from the CLAW-IV¹¹ library. Most of the materials suggested for fusion applications are contained in this multi-group cross section set. Each material in the CLAW-IV library has four cross section matrices ($P_0 - P_3$), with 30 neutron groups and 12 photon groups coupled in each. In the diffusion calculations, only the P_0 matrix is used.

I.E Objectives of This Work

The objective of the research performed here is to produce a scoping code that can be used for fusion blanket design. Pomraning and Stevens¹² have previously explained a technique in which a toroidal geometry problem can be cast into cylindrical ($r-\theta$) form by a spatially dependent redefinition of the diffusion coefficient, absorption cross-section, and extraneous source function. Use of this method will first be attempted in an existing two-dimensional ($x-y, r-z, r-\theta$, triangular), multi-group neutron diffusion code, 2DB, that already executes on the IBM-AT. This method will also be assessed and compared to direct finite differencing of the toroidal diffusion equation. The method with the most advantages will then be implemented in a new version of 2DB that will be called 2DBTOR.

To verify that 2DBTOR is operating correctly, comparisons will be made to both analytical and numerical solutions for several types of problems. ANISN and 2DB

will be used to verify the solutions obtained from 2DBTOR.

The modifications that were made to 2DB, a description of the problems and analytical solutions used to verify 2DBTOR, and the further modifications that were necessary to use 2DBTOR for fusion blanket calculations are presented in the following sections of this thesis.

II. THEORY

II.A. Introduction

The design of tokamak reactors requires an accurate calculation of the flux profile in the toroidally-shaped blanket. As stated previously, an r - θ neutron diffusion code will be used, with the appropriate modifications, to obtain a solution for the flux while taking into account the curvature of the torus. The solution from this modified neutron diffusion code might then be used as a first approximation for later runs in a more accurate transport code, such as TRISM⁸.

In order to make the appropriate changes so that a two-dimensional, r - θ , diffusion code can be made to model toroidal geometry, an understanding of the form of the diffusion equation in toroidal geometry is necessary. The diffusion equation in toroidal geometry can be developed from the equation for the time rate of change of the number of neutrons at energies, E , in an arbitrary differential volume, dV . Neutrons of energies between E and $E+dE$, within dV , can be lost or gained by a variety of processes including: (1) production directly from a source, (2) absorption, (3) leakage and (4) scattering. The time rate of change of the number of neutrons in dV and between E and $E+dE$ can be obtained by integrating the neutron density ($n(r,E,t)$) over dV , and balancing this with the gains and losses as follows:

$$\frac{\partial}{\partial t} \int_V \frac{\Phi(\vec{r}, E, t)}{v} dV = \left[\begin{array}{c} \text{source neutron} \\ \text{production rate} \\ \text{in } V \text{ at } E \end{array} \right] - \left[\begin{array}{c} \text{absorption} \\ \text{rate in } V \\ \text{at } E \end{array} \right] - \left[\begin{array}{c} \text{change due} \\ \text{to leakage} \\ \text{from } V \text{ at } E \end{array} \right] \\ - \left[\begin{array}{c} \text{neutron scattering} \\ \text{rate out of} \\ \text{E in } V \end{array} \right] + \left[\begin{array}{c} \text{neutron scattering} \\ \text{rate into} \\ \text{E in } V \end{array} \right] \quad (5)$$

where $\Phi(r,E,t)$ = flux of neutrons in r at E and t = $n(r,E,t)v$

and v = the speed of the neutrons at E .

Equation (5) is known as the neutron continuity equation. Since the energy dependence of the neutron cross sections vary, equation (5) is usually solved for discrete energy groups (groups denoted by g in this case); thus equation (5) can be written as

$$\frac{\partial}{\partial t} \int_V \frac{\Phi_g(r,t)}{v_g} dV = \left[\begin{array}{c} \text{source neutron} \\ \text{production rate} \\ \text{in } V \text{ for} \\ \text{group, } g \end{array} \right] - \left[\begin{array}{c} \text{absorption} \\ \text{rate in } V \\ \text{for group, } g \end{array} \right] - \left[\begin{array}{c} \text{change due} \\ \text{to leakage} \\ \text{from } V \\ \text{for group, } g \end{array} \right] \\ - \left[\begin{array}{c} \text{neutron scattering} \\ \text{rate out of} \\ \text{group, } g \text{ in } V \end{array} \right] + \left[\begin{array}{c} \text{neutron scattering} \\ \text{rate into} \\ \text{group, } g \text{ in } V \end{array} \right] \quad (6)$$

Before substituting expressions for the terms in equation (6), it is necessary to consider the flow of neutrons in a medium. Since the neutron current vector, $J(r,E,t)$ is a vector quantity, then

$$J \cdot n = \text{net rate of flow across a surface, } A \text{ (} n = \text{outward unit normal to a surface, } A \text{).} \quad (7)$$

To determine the total net rate of flow out of a closed surface, S , integration of equation (7) over S yields

$$\int_S \vec{J}(\vec{r}, E, t) \cdot \hat{n} \, dA,$$

where n is the unit normal to the surface dA . Using the Divergence Theorem, then

$$\int_S \vec{J}(\vec{r}, E, t) \cdot \hat{n} \, dA = \int_V \vec{\nabla} \cdot \vec{J}(\vec{r}, E, t) \, dV. \quad (8)$$

Since equation (6) will be used to solve for the neutron flux, it is necessary to relate $J(\vec{r}, E, t)$ to the flux. In the diffusion approximation the relation between current and flux is assumed to be as follows:

$$\vec{J}(\vec{r}, E, t) \equiv -D(\vec{r}) \vec{\nabla} \Phi(\vec{r}, E, t) \quad (9)$$

where $D(\vec{r}) =$ the diffusion coefficient.

Equation (9), known as Fick's Law for Diffusion, describes the flow of neutrons as being proportional to the negative of the density(flux) gradient, since particles tend to flow from a region of higher density to a region of lower density. Fick's Law is only valid for:

- 1.) points away from a vacuum boundary,
- 2.) points away from sources (by a few mean free paths),
- 3.) isotropic scattering (e.g., equal probability of scattering in any direction,
- and 4.) a slowly varying flux (absorption \ll scattering).¹³

The above conditions can be somewhat relaxed, depending on the accuracy required for the solution of equation (6).

Using Fick's Law it is now possible to express the balance equation only in terms of the neutron flux. Substituting the right hand side (RHS) of equation (9) into the leakage term found from the Divergence Theorem (equation (8)), then

$$\left[\begin{array}{l} \text{change due} \\ \text{to leakage} \\ \text{from V for} \\ \text{group, } g \end{array} \right] = \int_V \vec{\nabla} \cdot (-D_g(\vec{r}) \vec{\nabla} \Phi_g(\vec{r}, t)) dV \quad (10)$$

The use of Fick's Law in the neutron balance equation, equation (6), is known as the diffusion approximation, and the resulting equations are known as neutron diffusion equations.

The source term in equation (6) can be expressed by defining a source density, $S_g(\vec{r}, t)$, for the group, g . If $S_g(\vec{r}, t)$ is integrated over V , then

$$\left[\begin{array}{l} \text{source neutron} \\ \text{production rate} \\ \text{in V for} \\ \text{group, g} \end{array} \right] = \int_V S_g(\vec{r}, t) dV . \quad (11)$$

If source neutrons are allowed to be produced by either fission or extraneous (independent of the neutron flux) sources, then let

$$S_g(\vec{r}, t) = S_g^{\text{ext}}(\vec{r}, t) + S_g^{\text{fis}}(\vec{r}, t) \quad (12)$$

where $S_g^{\text{ext}}(\vec{r}, t)$ = extraneous source density rate
and $S_g^{\text{fis}}(\vec{r}, t)$ = fission source density rate.

The fission source density rate is given by

$$S_g^{\text{fis}}(\vec{r}, t) = \left[\begin{array}{l} \text{average number of} \\ \text{fission neutrons emitted} \\ \text{within group, g} \end{array} \right] \sum_{g'=1}^G \left\{ \left[\begin{array}{l} \text{fission rate in} \\ \text{group, g' per} \\ \text{unit volume} \end{array} \right] \left[\begin{array}{l} \text{average number of} \\ \text{neutrons released} \\ \text{into group, g, that} \\ \text{occur from fission} \\ \text{in group, g'} \end{array} \right] \right\} \quad (13)$$

where G = total number of groups.

It is common to denote the average number of neutrons released into group, g , that occur from fission in group, g' , by $\nu_{g'g}$ and the average number of fission neutrons emitted within group, g , by χ_g .¹⁴ Thus

$$S_g^{\text{fis}}(\vec{r}, t) = \chi_g \sum_{g'=1}^G v_{g'} \left[\begin{array}{l} \text{fission rate} \\ \text{in group, } g' \\ \text{per unit volume} \end{array} \right] \quad (14)$$

Since

$$\left[\begin{array}{l} \text{reaction rate} \\ \text{per unit} \\ \text{volume} \end{array} \right] = \Sigma \Phi \quad (15)$$

where Σ = macroscopic cross section,

the expression for the fission source can be written as

$$S_g^{\text{fis}}(\vec{r}, t) = \chi_g \sum_{g'=1}^G v_{g'} \Sigma_{fg'}(\vec{r}) \Phi_{g'}(\vec{r}, t). \quad (16)$$

The rest of the terms in the neutron balance equation can be expressed analogously to the above. For the absorption rate term, the absorption rate per unit volume is integrated over V to give

$$\left[\begin{array}{l} \text{absorption rate} \\ \text{in } V \text{ for group, } g \end{array} \right] = \int_V \Sigma_{ag}(\vec{r}) \Phi_g(\vec{r}, t) dV \quad (17)$$

where $\Sigma_{ag}(\vec{r})$ = the macroscopic absorption cross-section in group g .

Assuming that neutrons can not scatter to groups of higher energy (upscatter), then the neutron outscattering rate term is given by (decreasing group numbers correspond to increasing neutron energy)

$$\left[\begin{array}{l} \text{neutron scattering rate} \\ \text{out of group } g \text{ in } V \end{array} \right] = \int_V \sum_{g'=1}^g \Sigma_s(g \rightarrow g') \Phi_g(\vec{r}, t) dV \quad (18)$$

and the neutron in-scattering rate term is given by

$$\left[\begin{array}{l} \text{neutron scattering rate} \\ \text{into group } g \text{ in } V \end{array} \right] = \int_V \sum_{g'=1}^{g-1} \Sigma_s(g \rightarrow g') \Phi_{g'}(\vec{r}, t) dV \quad (19)$$

where $\Sigma_s(g \rightarrow g')$ = macroscopic scattering cross-section from g to g'

and $\Sigma_s(g \rightarrow g')$ = macroscopic scattering cross-section from g' to g .

Now that all of the terms in the neutron balance equation have been determined, equation (6) can be rewritten. Assuming steady state ($\partial/\partial t$ term = 0) and no upscatter, the neutron balance equation becomes, combining terms,

$$0 = \int_V \left(S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f,g'} \Phi_{g'}) - \Sigma_{a,g} \Phi_g - \sum_{g'=g+1}^G (\Sigma_{s(g \rightarrow g')} \Phi_{g'}) \right. \quad (20)$$

$$\left. + \vec{\nabla} \cdot (D_g \vec{\nabla} \Phi_g) + \sum_{g'=1}^{g-1} (\Sigma_{s(g \rightarrow g')} \Phi_{g'}) \right) dV$$

where the (r,t) has been dropped for clarity. Since the volume, V , was arbitrarily chosen, equation (20) reduces to the form that follows:

$$-\vec{\nabla} \cdot (D_g \vec{\nabla} \Phi_g) + \Sigma_{a,g} \Phi_g + \sum_{g'=g+1}^G (\Sigma_{s(g \rightarrow g')} \Phi_{g'}) = \quad (21)$$

$$S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f,g'} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_{s(g \rightarrow g')} \Phi_{g'})$$

Since the removal of neutrons from group g is caused by both downscattering and absorption, the removal cross section for group g is defined as shown below:

$$\Sigma_g^r \equiv \Sigma_{a,g} + \sum_{g'=g+1}^G (\Sigma_{s(g \rightarrow g')})$$

$$= \Sigma_{a,g} + \Sigma_{s(g \rightarrow g)} + \left[\sum_{g'=g+1}^G \Sigma_{s(g \rightarrow g')} \right] - \Sigma_{s(g \rightarrow g)}$$

$$= \Sigma_{tr,g} - \Sigma_{s(g \rightarrow g')}$$

where Σ_{tr_g} = the macroscopic transport cross section = $1/(3D_g)$,¹⁴

Thus the removal rate/cm³ is

$$\Sigma_g^r \Phi_g = \Sigma_{a_g} \Phi_g + \sum_{g'=g+1}^G (\Sigma_s(g \rightarrow g') \Phi_{g'}). \quad (22)$$

Finally, substituting equation (22) into equation (21) gives

$$-\vec{\nabla} \cdot (D_g \vec{\nabla} \Phi_g) + \Sigma_g^r \Phi_g = S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f_{g'}} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_s(g \rightarrow g') \Phi_{g'}) \quad (23)$$

To apply the above equation to practical problems, the ∇ operator must be written in the coordinate system of interest. In general orthogonal curvilinear coordinates

$$\vec{\nabla} \Phi = \frac{1}{h_1} \frac{\partial \Phi}{\partial u_1} \hat{e}_1 + \frac{1}{h_2} \frac{\partial \Phi}{\partial u_2} \hat{e}_2 + \frac{1}{h_3} \frac{\partial \Phi}{\partial u_3} \hat{e}_3$$

and

$$\vec{\nabla} \cdot \vec{A} = \frac{1}{h_1 h_2 h_3} \left[\frac{\partial}{\partial u_1} (h_2 h_3 A_1) + \frac{\partial}{\partial u_2} (h_3 h_1 A_2) + \frac{\partial}{\partial u_3} (h_1 h_2 A_3) \right]$$

where $u_1, u_2,$ and u_3 are the coordinates;

$e_1, e_2,$ and e_3 are the corresponding coordinate vectors;

and $h_1, h_2,$ and h_3 are scale factors that depend on the coordinates.

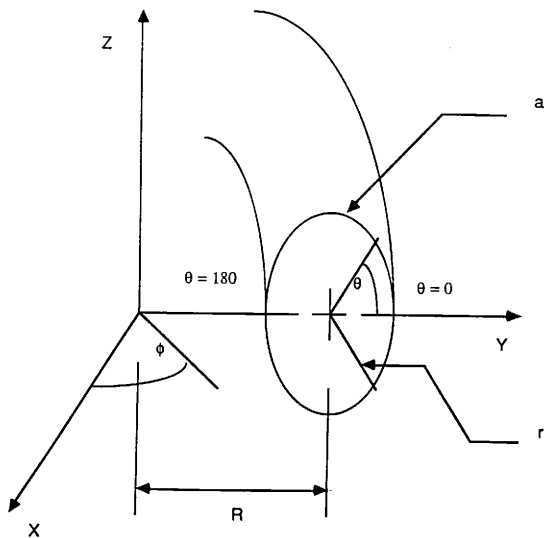
For cylindrical coordinates (r, θ, z)

$$\begin{array}{lll} h_1 = 1 & u_1 = r & e_1 = r \\ h_2 = r & u_2 = \theta & e_2 = \theta \\ h_3 = 1 & u_3 = z & e_3 = z. \end{array}$$

For toroidal coordinates (r, θ, ϕ) (see Figure 2-1)

$$\begin{array}{lll} h_1 = 1 & u_1 = r & e_1 = r \\ h_2 = r & u_2 = \theta & e_2 = \theta \\ h_3 = R+r\cos\theta & u_3 = \phi & e_3 = \phi.^{15} \end{array}$$

Consider the diffusion equation in cylindrical coordinates. From the above,



R = major radius, a = minor radius, r = radial coordinate,
 θ = poloidal angle, and ϕ = toroidal angle.

Figure 2-1: Toroidal Coordinate Description.

$$D\vec{\nabla}\Phi = D\frac{\partial\Phi}{\partial r}\hat{r} + \frac{D}{r}\frac{\partial\Phi}{\partial\theta}\hat{\theta} + D\frac{\partial\Phi}{\partial z}\hat{z} \quad (24)$$

Also,

$$\vec{\nabla} \cdot (D\vec{\nabla}\Phi) = \frac{1}{r} \left[\frac{\partial}{\partial r} (rD\frac{\partial\Phi}{\partial r}) + \frac{\partial}{\partial\theta} \left(\frac{D}{r} \frac{\partial\Phi}{\partial\theta} \right) + \frac{\partial}{\partial z} (rD\frac{\partial\Phi}{\partial z}) \right] \quad (25)$$

Analogously for toroidal coordinates,

$$D\vec{\nabla}\Phi = D\frac{\partial\Phi}{\partial r}\hat{r} + \frac{D}{r}\frac{\partial\Phi}{\partial\theta}\hat{\theta} + \frac{D}{R+r\cos\theta}\frac{\partial\Phi}{\partial\phi}\hat{\phi} \quad (26)$$

and

$$\vec{\nabla} \cdot (D\vec{\nabla}\Phi) = \frac{1}{r(R+r\cos\theta)} \left[\frac{\partial}{\partial r} (r(R+r\cos\theta)D\frac{\partial\Phi}{\partial r}) + \frac{\partial}{\partial\theta} ((R+r\cos\theta)\frac{D}{r}\frac{\partial\Phi}{\partial\theta}) + \frac{\partial}{\partial\phi} \left(\frac{rD}{(R+r\cos\theta)} \frac{\partial\Phi}{\partial\phi} \right) \right] \quad (27)$$

The diffusion equation in a cylindrical coordinate system, assuming axial symmetry (i.e., the flux, $\Phi(r, \theta, z)$ depends only on the r and θ coordinates), is given by

$$\begin{aligned}
& -\frac{1}{r} \left[\frac{\partial}{\partial r} (r D_g \frac{\partial \Phi_g}{\partial r}) + \frac{\partial}{\partial \theta} \left(\frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} \right) \right] + \Sigma_g^r \Phi_g \\
& = S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f'} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_s(g \rightarrow g') \Phi_{g'})
\end{aligned} \tag{28}$$

Similarly the diffusion equation in a toroidal coordinate system with axisymmetry (i.e., the flux, $\Phi(r, \theta, \phi)$ depends only on the r and θ coordinates), is given by

$$\begin{aligned}
& -\frac{1}{r(R+r\cos\theta)} \left[\frac{\partial}{\partial r} (r(R+r\cos\theta) D_g \frac{\partial \Phi_g}{\partial r}) + \frac{\partial}{\partial \theta} \left((R+r\cos\theta) \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} \right) \right] \\
& + \Sigma_g^r \Phi_g = S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f'} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_s(g \rightarrow g') \Phi_{g'})
\end{aligned} \tag{29}$$

II.B. Pomraning and Stevens' Method

If the toroidal form of the diffusion equation (equation (25)) is multiplied by $(R + r\cos\theta)/R$, then

$$-\frac{1}{r} \frac{\partial}{\partial r} \left(r \left(1 + \frac{r}{R} \cos \theta \right) D_g \frac{\partial \Phi_g}{\partial r} \right) - \frac{1}{r^2} \frac{\partial}{\partial \theta} \left(\left(1 + \frac{r}{R} \cos \theta \right) D_g \frac{\partial \Phi_g}{\partial \theta} \right) + \left(1 + \frac{r}{R} \cos \theta \right) \Sigma_g^r \Phi_g = \left(1 + \frac{r}{R} \cos \theta \right) \bar{S}_g$$

where $S_g =$ three terms representing sources on the right hand side of equation (29).

If now one defines

$$D_g' = \left(1 + r/R \cos \theta \right) D_g$$

$$\Sigma_g^r = \left(1 + r/R \cos \theta \right) \Sigma_g^r$$

$$S_g' = \left(1 + r/R \cos \theta \right) S_g$$

Then the above equation becomes

$$-\frac{1}{r} \frac{\partial}{\partial r} \left(r D_g' \frac{\partial \Phi_g}{\partial r} \right) - \frac{1}{r^2} \frac{\partial}{\partial \theta} \left(D_g' \frac{\partial \Phi_g}{\partial \theta} \right) + \Sigma_g^r \Phi_g = \bar{S}_g'$$

This equation is just the cylindrical diffusion equation with a modified diffusion coefficient, removal cross section, and source function. Pomraning and Stevens¹² proposed the possibility of using an existing r - θ neutron diffusion code and redefining the diffusion coefficient, removal cross section, and source function to model the

curvature of the torus naturally. This method was the original idea that suggested the topic of this research ; however, this approach was later abandoned in favor of direct finite differencing of the toroidal diffusion equation.

II.C. Finite Difference Approximation to the Diffusion Equation in Cylindrical Coordinates

To develop a finite difference approximation for the cylindrical diffusion equation (with axial symmetry), it is first necessary to integrate equation (28) over a small, arbitrary volume ΔV (see Figure 2-2). Thus,

$$\begin{aligned}
 & - \int_{\Delta V} \frac{1}{r} \left[\frac{\partial}{\partial r} \left(r D_g \frac{\partial \Phi_g}{\partial r} \right) + \frac{\partial}{\partial \theta} \left(\frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} \right) \right] dV + \int_{\Delta V} \left[\Sigma_g^r \Phi_g \right] dV \\
 & = \int_{\Delta V} \left[S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{r_s} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_s(g \rightarrow g') \Phi_{g'}) \right] dV
 \end{aligned} \tag{30}$$

where the first term on the LHS of the equation is the leakage term, the second term on

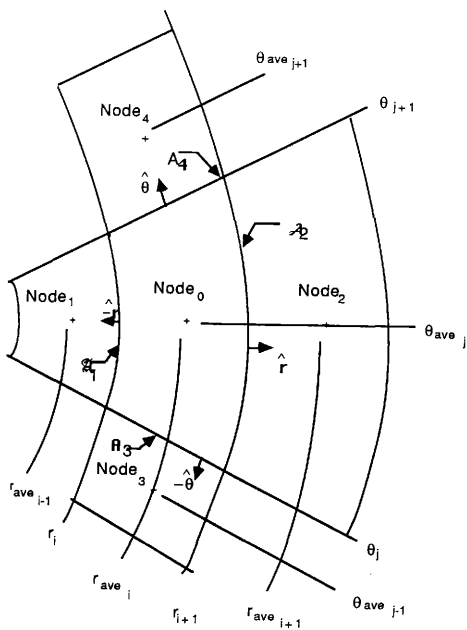


Figure 2-2. Finite Difference Coordinates Used in 2DB and 2DBTOR.

the LHS of the equation is the removal term, and the RHS represents the source terms including fission and scatter. Thus for the removal term,

$$\begin{aligned} \int_{\Delta V} \Sigma_g^r \Phi_g dV &= \Sigma_{g_0}^r \Phi_{g_0} (r_{i+1} - r_i) (\theta_{j+1} - \theta_j) r_{ave_1} \\ &= \Sigma_{g_0}^r \Phi_{g_0} V_0 \end{aligned} \quad (31)$$

where Φ_{g_0} = flux associated with meshpoint 0
and $\Sigma_{g_0}^r$ = removal cross section associated with meshpoint 0.

Just as above, the source term on the RHS can be shown to have a V_0 given by

$$V_0 = (r_{i+1} - r_i) (\theta_{j+1} - \theta_j) r_{ave_1}. \quad (32)$$

The leakage term is changed to an integral over the surface area of the volume element, thus from the Divergence Theorem

$$- \int_{\Delta V} \vec{\nabla} \cdot D_g \vec{\nabla} \Phi_g dV = - \int_A D_g \vec{\nabla} \Phi_g \cdot \hat{n} dA .$$

Using equation (24) for the $D_g \vec{\nabla} \Phi_g$ term (where $d\Phi/dz = 0$)

$$- \int_A D_g \vec{\nabla} \Phi_g \cdot \hat{n} dA = - \int_A \left(D_g \frac{\partial \Phi_g}{\partial r} \hat{r} + \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} \hat{\theta} \right) \cdot \hat{n} dA .$$

The normal vector, n , is r or θ for those area elements having normals in the positive r or θ , or increasing r or θ , respectively. In the same manner, n is $-r$ or $-\theta$ for those elements in the $-r$ or $-\theta$ directions. The area element corresponding to a normal in the $\pm r$ direction is $r d\theta$, while the area element for a normal in the $\pm\theta$ direction is dr . Thus for the 4 area elements

$$\begin{aligned}
 & - \int_A \left(D_g \frac{\partial \Phi}{\partial r} \hat{e}_r + \frac{D_g}{r} \frac{\partial \Phi}{\partial \theta} \hat{e}_\theta \right) \cdot \hat{n} dA \\
 = & - \int_{A_1} \left(D_g \frac{\partial \Phi}{\partial r} \hat{e}_r + \frac{D_g}{r} \frac{\partial \Phi}{\partial \theta} \hat{e}_\theta \right) \cdot \hat{n} dA_1 \\
 & - \int_{A_2} \left(D_g \frac{\partial \Phi}{\partial r} \hat{e}_r + \frac{D_g}{r} \frac{\partial \Phi}{\partial \theta} \hat{e}_\theta \right) \cdot \hat{n} dA_2 \\
 & - \int_{A_3} \left(D_g \frac{\partial \Phi}{\partial r} \hat{e}_r + \frac{D_g}{r} \frac{\partial \Phi}{\partial \theta} \hat{e}_\theta \right) \cdot \hat{n} dA_3 \\
 & - \int_{A_4} \left(D_g \frac{\partial \Phi}{\partial r} \hat{e}_r + \frac{D_g}{r} \frac{\partial \Phi}{\partial \theta} \hat{e}_\theta \right) \cdot \hat{n} dA_4
 \end{aligned}$$

Substituting for n and the dA 's (see Figure 2-2) then the RHS becomes

$$\begin{aligned}
&= - \int_{\theta_j}^{\theta_{j+1}} \left(D_g \frac{\partial \Phi_g}{\partial r} \hat{r} + \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} \hat{\theta} \right) \cdot \hat{r} r_1 d\theta \Big|_{A_1} \\
&\quad - \int_{\theta_j}^{\theta_{j+1}} \left(D_g \frac{\partial \Phi_g}{\partial r} \hat{r} + \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} \hat{\theta} \right) \cdot \hat{r} r_{i+1} d\theta \Big|_{A_2} \\
&\quad - \int_{r_i}^{r_{i+1}} \left(D_g \frac{\partial \Phi_g}{\partial r} \hat{r} + \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} \hat{\theta} \right) \cdot -\hat{\theta} dr \Big|_{A_3} \\
&\quad - \int_{r_i}^{r_{i+1}} \left(D_g \frac{\partial \Phi_g}{\partial r} \hat{r} + \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} \hat{\theta} \right) \cdot \hat{\theta} dr \Big|_{A_4}
\end{aligned}$$

Since $\mathbf{r} \cdot \hat{\theta} = \hat{\theta} \cdot \mathbf{r} = 0$ and $\mathbf{r} \cdot \mathbf{r} = \hat{\theta} \cdot \hat{\theta} = 1$, then the RHS simplifies to

$$= \int_{\theta_j}^{\theta_{j+1}} D_g \frac{\partial \Phi_g}{\partial r} r_1 d\theta \Big|_{A_1} - \int_{\theta_j}^{\theta_{j+1}} D_g \frac{\partial \Phi_g}{\partial r} r_{i+1} d\theta \Big|_{A_2} + \int_{r_i}^{r_{i+1}} \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} dr \Big|_{A_3} - \int_{r_i}^{r_{i+1}} \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} dr \Big|_{A_4} \quad (33)$$

Since the partial derivatives of the flux will be obtained by differencing the two neighboring flux values, then letting k be the adjacent mesh point to mesh point o gives

$$\left. \frac{d\phi}{dx} \right| \approx \frac{\phi_k - \phi_0}{\Delta x}$$

where $x = r$ or θ depending on the derivative being considered.

Then

$$\int_{\theta_i}^{\theta_{j+1}} D_g \frac{\partial \Phi_g}{\partial r} r_i d\theta \Big|_{A_i} = \bar{D}_{g1} \frac{\partial \Phi_g}{\partial r} \Big|_{r_i} r_i (\theta_{j+1} - \theta_j)$$

$$\approx \bar{D}_{g1} \frac{\phi_{g1} - \phi_{g0}}{r_{ave_i} - r_{ave_{i+1}}} r_i (\theta_{j+1} - \theta_j)$$

where D_{gk} is defined to be (see Appendix B)

$$\bar{D}_{gk} = \frac{D_{g0} D_{gk} (\Delta r_o + \Delta r_k)}{(D_{g0} \Delta r_k + D_{gk} \Delta r_o)} \quad \text{or} \quad \frac{D_{g0} D_{gk} (\Delta \theta_o + \Delta \theta_k)}{(D_{g0} \Delta \theta_k + D_{gk} \Delta \theta_o)} \quad \text{for } k = 1, 2, 3, \text{ or } 4$$

and D_{g0} = average D_g for volume element o ,

D_{gk} = average D_g for volume element k ,

$\Delta r_o = \Delta r$ for volume element o ,

$\Delta r_k = \Delta r$ for volume element k ,

$\Delta\theta_o = \Delta\theta$ for volume element o, and

$\Delta\theta_o = \Delta\theta$ for volume element o.

Let

$$\begin{aligned} A_1 &= r_i (\theta_{j+1} - \theta_j) \\ L_1 &= r_{ave_i} - r_{ave_{i-1}} \end{aligned}$$

Then

$$\int_{\theta_j}^{\theta_{j+1}} D_g \frac{\partial \Phi_g}{\partial r} r_i d\theta \approx -\bar{D}_{g1} \frac{\phi_{g1} - \phi_{g0}}{L_1} A_1 \quad (34)$$

For r_i close to r_{i+1} , $\ln(r_{i+1}/r_i) \approx (r_{i+1} - r_i)$. Thus,

$$\ln \left(\frac{r_{i+1}}{r_i} \right) \approx r_{i+1} - r_i \quad ,$$

which when used in the third term in equation (33) gives

$$\int_{r_i}^{r_{i+1}} \frac{D_g \partial \Phi_g}{r \partial \theta} dr = \bar{D}_{g3} \frac{\partial \Phi_g}{\partial \theta} \Big|_{\theta_j} \ln r \Big|_{r_i}^{r_{i+1}}$$

$$\approx -\bar{D}_{g3} \frac{\phi_{g3} - \phi_{g0}}{\theta_{ave_j} - \theta_{ave_{j-1}}} (r_{i+1} - r_i) .$$

Let

$$A_3 = (r_{i+1} - r_i)$$

$$L_3 = \theta_{ave_j} - \theta_{ave_{j-1}} ,$$

then

$$\int_{r_i}^{r_{i+1}} \frac{D_g \partial \Phi_g}{r \partial \theta} dr \approx -\bar{D}_{g3} \frac{\phi_{g3} - \phi_{g0}}{L_3} A_3 . \quad (35)$$

The above process can be applied to the remaining terms of equation (33) to give

$$\begin{aligned}
-\int_{\theta_j}^{\theta_{j+1}} D_g \frac{\partial \Phi_g}{\partial r} r_{i+1} d\theta &\approx \bar{D}_{g2} \frac{\phi_{g2} - \phi_{g0}}{r_{ave_{i+1}} - r_{ave_i}} r_{i+1} (\theta_{j+1} - \theta_j) \\
&= -\bar{D}_{g2} \frac{\phi_{g2} - \phi_{g0}}{L_2} A_2
\end{aligned} \tag{36}$$

and

$$\begin{aligned}
-\int_{r_i}^{r_{i+1}} \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} dr &\approx \bar{D}_{g4} \frac{\phi_{g4} - \phi_{g0}}{\theta_{ave_{j+1}} - \theta_{ave_j}} (r_{i+1} - r_i) \\
&= -\bar{D}_{g4} \frac{\phi_{g4} - \phi_{g0}}{L_4} A_4
\end{aligned} \tag{37}$$

where $L_2 = r_{ave_{i+1}} - r_{ave_i}$,

$$A_2 = r_{i+1} (\theta_{j+1} - \theta_j),$$

$$L_4 = \theta_{ave_{j+1}} - \theta_{ave_j},$$

and $A_4 = r_{i+1} - r_i$.

Combining all of the above terms back into equation (30)

$$\begin{aligned}
& - \sum_{k=1}^4 \left[\bar{D}_{g_k} \left(\frac{\phi_{g_k} - \phi_{g_0}}{L_k} \right) A_k \right] + \Sigma_{g_0}^r \Phi_{g_0} V_0 \\
& = \left[S_{g_0}^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_g \cdot \Sigma_{f_{g'}} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_{s_g}(g \rightarrow g') \Phi_{g'}) \right] V_0 .
\end{aligned} \tag{38}$$

II.D. Finite Difference Approximation to the Diffusion Equation in Toroidal Coordinates

In a similar manner as was used in the previous section, the diffusion equation in a toroidal coordinate system can be written in finite difference form. Using the Divergence Theorem to transpose the leakage term in the toroidal diffusion equation (equation (29)) from a volume to a surface integral and substituting for the $D_g \nabla \Phi_g$ term from equation (26), then

$$\begin{aligned}
& - \int_A \left[D_g \frac{\partial \Phi_{g\hat{r}}}{\partial r} + \frac{D_g}{r} \frac{\partial \Phi_{g\hat{\theta}}}{\partial \theta} + \frac{D_g}{(R+r\cos\theta)} \frac{\partial \Phi_{g\hat{\phi}}}{\partial \phi} \right] \cdot \hat{n} \, dA + \int_V \left[\Sigma_g^r \Phi_{g\hat{r}} \right] dV \quad (39) \\
& = \int_V \left[S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{r_g'} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_s(g \rightarrow g') \Phi_{g'}) \right] dV
\end{aligned}$$

The leakage term in toroidal coordinates becomes (assuming axisymmetry)

$$- \int_A 2\pi(R+r\cos\theta) \left(D_g \frac{\partial \Phi_{g\hat{r}}}{\partial r} + \frac{D_g}{r} \frac{\partial \Phi_{g\hat{\theta}}}{\partial \theta} \right) \cdot \hat{n} \, dA .$$

where the $(R+r\cos\theta)$ term appears from multiplying by $(R+r\cos\theta)$ before taking $d\Phi/d\phi = 0$. As in the cylindrical form of the diffusion equation, n and the dA 's are the same, so the only difference in evaluating the toroidal form of the diffusion equation is the $(R+r\cos\theta)$ terms. Thus from the above, the leakage term becomes

$$\begin{aligned}
& 2\pi \int_A (R+r\cos\theta) \left(D_g \frac{\partial \Phi_g}{\partial r} + \frac{D_g}{r} \frac{\partial \Phi_g}{\partial \theta} \right) \cdot \hat{n} \, dA = \\
& 2\pi \int_{\theta_j}^{\theta_{j+1}} D_g(R+r\cos\theta) \frac{\partial \Phi_g}{\partial r} r_i \, d\theta \Big|_{A_1} - 2\pi \int_{\theta_j}^{\theta_{j+1}} D_g(R+r\cos\theta) \frac{\partial \Phi_g}{\partial r} r_{i+1} \, d\theta \Big|_{A_2} \\
& + 2\pi \int_{r_i}^{r_{i+1}} \frac{D_g}{r} (R+r\cos\theta) \frac{\partial \Phi_g}{\partial \theta} \, dr \Big|_{A_3} - 2\pi \int_{r_i}^{r_{i+1}} \frac{D_g}{r} (R+r\cos\theta) \frac{\partial \Phi_g}{\partial \theta} \, dr \Big|_{A_4} \\
& \approx -2\pi R \bar{D}_{g1} \frac{\phi_{g1} - \phi_{g0}}{L_1} \left[r_i \left[(\theta_{j+1} - \theta_j) + \frac{r_i}{R} (\sin\theta_{j+1} - \sin\theta_{j+1'}) \right] \right] \\
& - 2\pi R \bar{D}_{g2} \frac{\phi_{g2} - \phi_{g0}}{L_2} \left[r_{i+1} \left[(\theta_{j+1} - \theta_j) + \frac{r_{i+1}}{R} (\sin\theta_{j+1} - \sin\theta_{j+1'}) \right] \right] \\
& - 2\pi R \bar{D}_{g3} \frac{\phi_{g3} - \phi_{g0}}{L_3} \left[(r_{i+1} - r_i) \left[\frac{\cos\theta_{ave_{j+1}}}{R} + 1 \right] \right] \\
& - 2\pi R \bar{D}_{g4} \frac{\phi_{g4} - \phi_{g0}}{L_4} \left[(r_{i+1} - r_i) \left[\frac{\cos\theta_{ave_j}}{R} + 1 \right] \right]
\end{aligned}$$

For the removal term

$$\begin{aligned}
\int_V \Sigma_g^r \Phi_g dV &= \Sigma_{g_0}^r \Phi_{g_0} \int_{r_i}^{r_{i+1}} \int_{\theta_j}^{\theta_{j+1}} 2\pi r (R + r \cos \theta) dr d\theta \\
&= \Sigma_{g_0}^r \Phi_{g_0} 2\pi R \left[\left(\frac{r_{i+1}^2 - r_i^2}{2} \right) (\theta_{j+1} - \theta_j) + \frac{r_{i+1}^3 - r_i^3}{3R} (\sin \theta_{j+1} - \sin \theta_j) \right] \\
&= \Sigma_{g_0}^r \Phi_{g_0} 2\pi R \left[(r_{i+1} - r_i) r_{\text{ave}} (\theta_{j+1} - \theta_j) + \frac{r_{i+1}^3 - r_i^3}{3R} (\sin \theta_{j+1} - \sin \theta_j) \right] \\
&= \Sigma_{g_0}^r \Phi_{g_0} 2\pi R V_0
\end{aligned}$$

As in the cylindrical coordinate system diffusion equation, V_0 for the source term is the same as the above term for V_0 .

Let

$$\begin{aligned}
A_1 &= \left[r_i \left[(\theta_{j+1} - \theta_j) + \frac{r_i}{R} (\sin \theta_{j+1} - \sin \theta_j) \right] \right] \\
A_2 &= \left[r_{i+1} \left[(\theta_{j+1} - \theta_j) + \frac{r_{i+1}}{R} (\sin \theta_{j+1} - \sin \theta_j) \right] \right] \\
A_3 &= \left[(r_{i+1} - r_i) \left[\frac{\cos \theta_{\text{ave}_{j+1}}}{R} + 1 \right] \right] \\
A_4 &= \left[(r_{i+1} - r_i) \left[\frac{\cos \theta_{\text{ave}_j}}{R} + 1 \right] \right]
\end{aligned}$$

Using the above, the toroidal diffusion equation in finite difference form after using

equation (39) and dividing through by $2\pi R$ becomes

$$\begin{aligned}
 & - \sum_{k=1}^4 \left[\bar{D}_{gk} \left(\frac{\phi_{gk} - \phi_{g0}}{L_k} \right) A_k \right] + \sum_{g_0}^r \Phi_{g_0} V_0 \\
 & = \left[S_{g_0}^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f_{g'}} \Phi_{g'_0}) + \sum_{g'=1}^{g-1} (\Sigma_{s_0}(g \rightarrow g') \Phi_{g'_0}) \right] V_0 .
 \end{aligned}$$

Thus, the finite difference approximation to the diffusion equation in any coordinate system results when appropriate area and volume elements are used. The differences between the cylindrical and toroidal coordinate system area and volume elements are listed in Table 2-1. By changing the area and volume elements to those for toroidal coordinates, a standard neutron diffusion code can be generalized to solve the multi-group diffusion equations in toroidal coordinates and in practice this is easier than the redefinition of cross sections proposed by Pomraning and Stevens.

II.E. Analytical Solutions

In order to insure that the above methodology was implemented correctly, analytical solutions for the flux profiles of an elementary case were obtained. Pomraning and Stevens¹² presented the analytical solution for the diffusion of neutrons due to a line source in a homogeneous cylindrical medium of radius, a (see Appendix A). This cylinder did not have a central void. Approximate solutions were then derived for the

Table 2-1. Area and Volume Elements in Cylindrical and Toroidal Coordinates.

	Cylindrical	Toroidal
A ₁	$r_i(\theta_{j+1} - \theta_j)$	$[r_i[(\theta_{j+1} - \theta_j) + \frac{r_i}{R}(\sin\theta_{j+1} - \sin\theta_j)]]$
A ₂	$r_{i+1}(\theta_{j+1} - \theta_j)$	$[r_{i+1}[(\theta_{j+1} - \theta_j) + \frac{r_{i+1}}{R}(\sin\theta_{j+1} - \sin\theta_j)]]$
A ₃	$(r_{i+1} - r_i)$	$[(r_{i+1} - r_i)[\frac{\cos\theta_{ave}}{R}j + 1]]$
A ₄	$r_{i+1} - r_i$	$[(r_{i+1} - r_i)[\frac{\cos\theta_{ave}}{R}j + 1]]$
V _o	$(r_{i+1} - r_i)(\theta_{j+1} - \theta_j)r_{ave}$	$[(r_{i+1} - r_i)r_{ave}(\theta_{j+1} - \theta_j) + \frac{r_{i+1}^3 - r_i^3}{3R}(\sin\theta_{j+1} - \sin\theta_j)]$

flux profile within a torus of major radius, R , and minor radius, a . The analytical result for the cylindrical problem is

$$\Phi(r) = \frac{1}{2\pi} \left[K_0\left(\frac{k r}{a}\right) - \frac{K_0(k)}{I_0(k)} I_0\left(\frac{k r}{a}\right) \right] \quad (40)$$

and the approximate result for the toroidal problem is

$$\Phi(r, \theta) \approx \frac{1}{2\pi} \left[K_0\left(\frac{k r}{a}\right) - \frac{K_0(k)}{I_0(k)} I_0\left(\frac{k r}{a}\right) \right] \left(1 - \frac{r R \cos \theta}{2} \right) + O\left(\left(\frac{a}{R}\right)^2\right) \quad (41)^{12}$$

where $k = (\Sigma_a/D)^{1/2}$; K_0 and I_0 are the zero order modified Bessel functions of the first and second kind, respectively; and $O((a/R)^2)$ denotes an error term of order $(a/R)^2$. For the above solutions, the only source is a line source of unit strength at $r=0$ ¹².

III. RESULTS

III.A 2DBTOR Verification and Infinite Cylinder Problems

In order to develop and verify 2DBTOR, modifications had to be made to 2DB, as described previously in the theory section. In addition, 2DB was tested to insure that it was operating correctly. This was done by comparing 2DB's results to both analytical and numerical flux solutions for several types of problems. ANISN was used to obtain numerical results to verify one-dimensional results from 2DB; Pomraning and Stevens' solution for a solid, homogeneous torus with a line source in the center provided an analytical verification in the limit as the aspect ratio of the torus became large (e.g., the torus approximated an infinite cylinder).¹²

The first part of the investigation was to compute, using 2DB, the flux in an infinite, homogeneous cylinder due to a line source. The radius of the cylinder was 300cm. Graphite was chosen as the material for the cylinder, since graphite has good neutron scattering properties ($\Sigma_s \gg \Sigma_a$) and thus should prove to give good results for the flux profile when diffusion theory is used. A source density given by

$$S_v = \frac{D}{\pi r_v^2} \quad (42)^{16}$$

where $S_v =$ source density (n/cm^3),

$D = 1/3\Sigma_{tr} =$ diffusion coefficient,

$\Sigma_{tr} =$ macroscopic transport cross section,

and $r_v =$ radius of the area for the source,

was input into 2DB for a small radius ($r_v = 6\text{cm}$) which would approximate the line source of equation (40) from the theory (see Appendix A¹⁶). ANISN was used to collapse the 30 group graphite cross sections to one group using S_2 angular quadrature with the P_0 matrix from the CLAW-IV neutron cross section library. After collapsing to one group, 2DBXPROC, a cross section processor to manipulate the cross sections from ANISN's format into 2DB's format (see Appendix C for a listing of 2DBXPROC), was used so that 2DB could employ the one group cross sections for graphite. Solutions for the flux away from the source for both the analytical calculation and the 2DB run gave results that were in excellent agreement, after errors for both the volume and area elements were corrected (see Appendix D). The one group cross section set for graphite was then input to ANISN with the same conditions as above. Solutions of the flux from ANISN were also in excellent agreement with 2DB.

After studying the infinite, homogeneous cylinder, computations were made to determine the flux due to a uniformly distributed source of 14 Mev neutrons (see Appendix E) in an infinite homogeneous cylinder with a central void. The source radius was 150 cm and the inner and outer radii of the medium were 200cm and 300cm, respectively. ANISN was again used to collapse the 30 group cross section

set for graphite to one group. In order to model the central void region, fictitious cross sections were first input into 2DB such that the absorption cross section was zero and the scattering cross section was equal to the transport cross section of graphite. This method proved to give somewhat erroneous results, however, since the value for the diffusion coefficient in the vacuum region was arbitrary. To alleviate this problem, a sufficiently thin source equivalent to the plasma source was placed at the inner edge of the annulus (see Appendix F) and the problem was re-run. This quasi-albedo boundary condition overcame the need for a fictitious scattering cross section in the void. Re-runs of the above problem gave 2DB and ANISN results which were in good agreement.

To complete the test of 2DB, some of the above problems were re-run with more than one energy group. This served to verify that 2DB's use of the downscattering cross sections was being performed correctly. Each problem was solved as before, except that 9 groups were used instead of one. Again the cross sections were obtained by collapsing the 30 group P_0 matrix using ANISN. Each solution of the flux was correct for each problem done previously for graphite. Since fusion problems would involve $(n,2n)$ and (n,n') reactions, it was decided to further test 2DB by using Nb-93, which has a significant $(n,2n)$ reaction at high energy for the problems previously studied. The results obtained using Nb-93 in 2DB for the two cases above (with and without central voids) were discovered to be quite different from those of ANISN. An investigation of 2DB's source code was therefore performed to identify the reason for the differences. It was then discovered that 2DB erroneously computed the downscattering (see Appendix D). After correcting 2DB to compute the proper downscatter contribution, the Nb-93 problem was re-run with a 9-group cross section

set, and the results were in good agreement with those obtained using ANISN.

III.B 2DBTOR Verification and Toroidal Problems

After making appropriate changes, as described earlier in this thesis (i.e, change the volume and area elements), to change 2DB to compute the flux for toroidal geometry (see Appendices G and H), a new code was produced which was called 2DBTOR. Using the same conditions stated previously for an infinite homogeneous graphite cylinder with one-group cross sections and without a central void, 2DBTOR was run for aspect ratios of 3 and 5 and compared to Pomraning and Stevens' solutions¹² for the same aspect ratio. The solutions for the flux were in excellent agreement. These solutions for both aspect ratios were then compared to the infinite cylinder solutions at radii of 150cm and 250cm from $\theta = 0$ to 2π . The results are shown in Figures 3-1 and 3-2 for 150cm and 250cm, respectively. As can be seen from Figures 3-1 and 3-2, the flux has a minimum at the outside part of the torus ($\theta = 0$) and increased up to the inner part ($\theta = \pi$), where the flux was a maximum. This was an expected result since the area of the inner portion of the torus is smaller than the outer portion. In addition, as the aspect ratio increases the solution approaches that of the infinite cylinder.

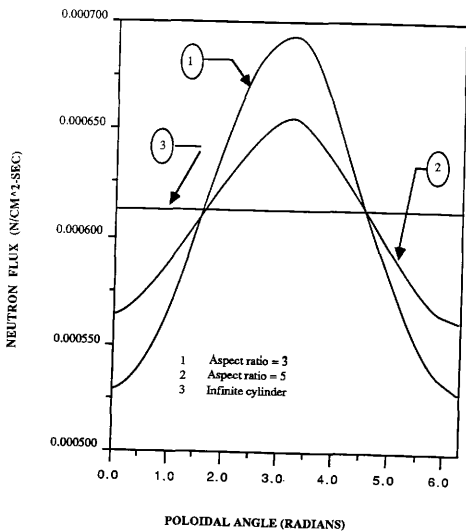


Figure 3-1. Flux Versus Poloidal Angle at a Radius of 150 cm for Several Aspect Ratios.

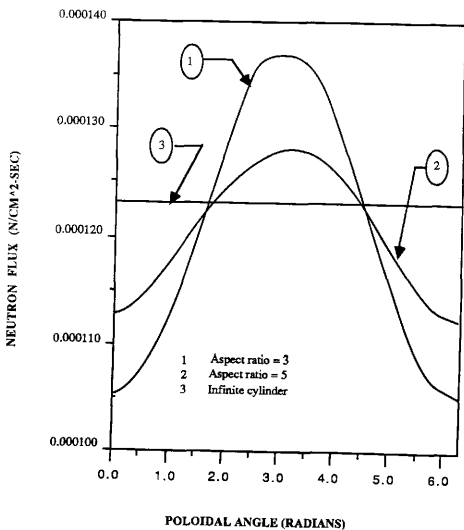


Figure 3-2. Flux Versus Poloidal Angle at a Radius of 250 cm for Several Aspect Ratios.

III.C Standard Blanket Solutions

III.C.1 Infinite Cylinder

The blanket used for this analysis is shown in Figure 3-3, and it is the so-called 'standard' blanket that was formulated as a benchmark for neutronic calculations.¹⁷ The materials used and their atom densities are given in Table 3-1.¹⁷ From Figure 3-3 it can be seen that the blanket consists of 10 zones, or regions, which contain one of 3 mixtures (except for the first 2 zones which are the vacuum and plasma). These are denoted A, B, or C. Each mixture has one or more materials including Li-6, Li-7, C-12, and Nb-93. When lithium is present in a mixture, the medium is assumed to be homogeneous with volume fractions of 94% lithium and 6% niobium. The lithium serves as both a coolant and tritium producer, while the niobium provides the structural function. The plasma has a radius of 150cm and occupies zone 1. A vacuum region is located from 150cm to 200cm and occupies zone 2. The blanket extends from 200cm to 300cm consisting of the first wall (200-200.5cm), a tritium production region (200.5-203.5cm), the second wall (203.5-204cm), 3 tritium production regions of 20cm thickness each (204-264cm), a carbon reflector (264-294cm), and a final tritium production region (294-300cm).

After setting up the standard blanket problem to run on ANISN, the 30-group cross sections in CLAW-IV for all the materials were collapsed to 9-group cross sections, with a P_0 Legendre expansion and S_2 quadrature to represent a diffusion theory calculation. 50 mesh intervals were used in this computation (see Figure 3-3). The 9-group cross section set was unable to produce satisfactory results for the tritium

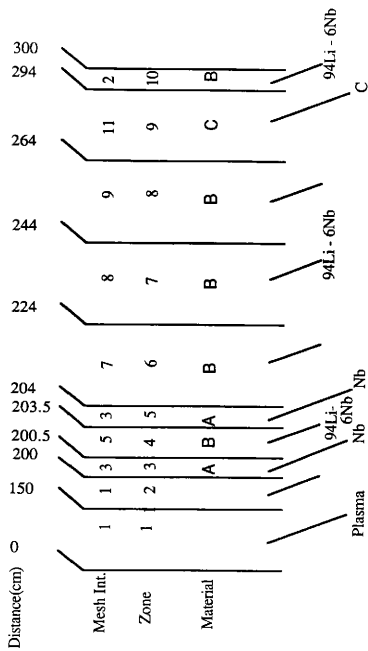


Figure 3-3. Standard Blanket Configuration.

Table 3-1. Standard Blanket Constituents.

Material Code Letter (from Figure 3-3)	Constituents	Atom density (atoms/cm ³ x 10 ⁻²⁴)
A	Nb-93	0.05556
B	Nb-93	0.003334
	Li-6	0.003234
	Li-7	0.04038
C	C-12	0.0804

breeding ratio (TBR = tritium production rate/ source neutron rate) for both T6 (TBR from Li-6) and T7 (TBR from Li-7). When the 9-group cross section set was input back into ANISN for comparison to the results of the 30-group set, differences in the corresponding zone TBR's were detected which were deemed to be unacceptable. To overcome this problem, it was decided to use the 30-group, uncollapsed, cross section set in 2DB and 2DBTOR for TBR computations. Again as in ANISN, 50 mesh intervals were used in both 2DB and 2DBTOR. The TBR results for 2DB and ANISN were compared in the zones of interest (i.e., zones 4,6,7,8, and 10 where lithium was present). The results of the 2DB and ANISN computations are given in Table 3-2 for the above 5 zones. It is seen from Table 3-2 that there is agreement between the results from the codes, especially in the T6 values. The T7 values are acceptable in the inner zones but have significant error in the outer zones. Likewise, the total TBR in the entire blanket is in closest agreement for T6, and differs by less than 5% for T7. A comparison of the individual group fluxes showed a similarly acceptable agreement between 2DB and ANISN. All of the above implies that 2DB can now be used successfully to perform fusion scoping calculations in one-dimensional geometry.

III.C.2 Comparison of Torus and Infinite Cylinder Models

Since 2DB was found to produce adequate results for the standard blanket compared to the ANISN results, a comparison of the results from both 2DBTOR and

Table 3-2. Radially Averaged Tritium Breeding Ratio's for 2DB and ANISN,
Infinite Cylinder Case.

Zone Number	Zone Radii (cm)	T6, reaction rate per source neutron rate from Li-6		T7, reaction rate per source neutron rate from Li-7	
		ANISN	2DB	ANISN	2DB
4	200.5-203.5	0.05043	0.05776	0.08191	0.10119
6	204.0-224.0	0.30399	0.34143	0.29809	0.33643
7	224.0-244.0	0.24598	0.25690	0.12189	0.09031
8	244.0-264.0	0.29518	0.27286	0.05050	0.02108
10	294.0-300.0	0.04923	0.03955	0.00093	0.00009
Total		0.94482	0.96850	0.55331	0.54910

2DB was made next. The 2DBTOR analysis was run in both toroidal (aspect ratios = 3 and 5) and cylindrical geometry. To begin the analysis, zones were chosen as in the previous infinite cylinder computation, so that radially averaged TBR's could be computed. In Tables 3-3 and 3-4, the TBR results for aspect ratios of 3 and 5 versus the infinite cylinder results were compared for T6 and T7, respectively. As can be seen there, the total T6 and T7 were about the same for both the toroidal and infinite cylinder calculations. This suggested the idea that if one were only interested in global reaction rates (e.g., not the outside versus inside reaction rates for the torus), then a one-dimensional neutron transport or diffusion code would be adequate for determining these quantities.

To investigate the local effects introduced by the toroidal geometry, the standard blanket was further sub-divided into ten zones of 36° each in the angular or poloidal direction. This resulted in 100 zones for the entire torus cross section (10 radially x 10 poloidally). The infinite cylinder T6 and T7 results are presented in Tables 3-5 and 3-6 for an aspect ratio of 3 and in Tables 3-7 and 3-8 for an aspect ratio of 5. Only the 5 zones from 0 to π are presented since the torus is symmetric about its axis (see Figure 2-1). From these tables, it was noted that the TBR was largest at the outer part of the torus ($\theta = 0^\circ$ to 36°) and decreased to a minimum value at the inner part of the torus ($\theta = 144^\circ$ to 180°). The effect of different aspect ratios was also seen by comparing the corresponding zone rates -- e.g., the 0° to 36° zone rate for an aspect ratio of 3 was about 7% larger than the 0° to 36° zone rate for an aspect ratio of 5, while the 144° to 180° zone rate was about 10% smaller for an aspect ratio of 3. This is a desirable result in a practical sense since the tritium on the outside part of the

Table 3-3. Radially Averaged Tritium Breeding Ratio's for 2DBTOR Versus 2DB,
Torus Case (Aspect Ratio = 3).

Zone Number	Zone Radii (cm)	T6, reaction rate per source neutron rate from Li-6		T7, reaction rate per source neutron rate from Li-7	
		2DBTOR	2DB	2DBTOR	2DB
4	200.5-203.5	0.05767	0.05776	0.10119	0.10119
6	204.0-224.0	0.34095	0.34143	0.33643	0.33643
7	224.0-244.0	0.25667	0.25690	0.09031	0.09031
8	244.0-264.0	0.27238	0.27286	0.02107	0.02108
10	294.0-300.0	0.03950	0.03955	0.00009	0.00009
Total		0.96717	0.96850	0.54909	0.54910

Table 3-4. Radially Averaged Tritium Breeding Ratio's for 2DBTOR Versus 2DB, Torus Case (Aspect Ratio = 5).

Zone	Zone	T6, reaction rate per source		T7, reaction rate per source	
Number	Radii (cm)	neutron rate from Li-6		neutron rate from Li-7	
		2DBTOR	2DB	2DBTOR	2DB
4	200.5-203.5	0.05781	0.05776	0.10114	0.10119
6	204.0-224.0	0.34167	0.34143	0.33619	0.33643
7	224.0-244.0	0.25690	0.25690	0.09029	0.09031
8	244.0-264.0	0.27286	0.27286	0.02107	0.02108
10	294.0-300.0	0.03950	0.03955	0.00009	0.00009
Total		0.96874	0.96850	0.54878	0.54910

Table 3-5. Zone by Zone Averaged T6 for 2DBTOR Versus 2DB,
Torus Case (Aspect Ratio = 3). Angle Measured in Degrees.

Zone Number	T6						
	2DB	0-36	36-72	72-108	108-144	144-180	0-360
4	0.05776	0.00688	0.00644	0.00503	0.00459	0.00459	0.05504
6	0.34143	0.04091	0.03824	0.03393	0.02960	0.02691	0.33914
7	0.25690	0.03110	0.02898	0.02555	0.02209	0.01994	0.25530
8	0.27286	0.03348	0.03107	0.02714	0.02319	0.02073	0.27123
10	0.03955	0.00495	0.00456	0.00394	0.00331	0.00292	0.03935
Total	0.96850	0.11732	0.10929	0.09559	0.08278	0.07509	0.96006

Table 3-6. Zone by Zone Averaged T7 for 2DBTOR Versus 2DB,
Torus Case (Aspect Ratio = 3). Angle Measured in Degrees.

Zone Number	T7						
	2DB	0-36	36-72	72-108	108-144	144-180	0-360
4	0.10119	0.01221	0.01141	0.01012	0.00884	0.00804	0.11240
6	0.33643	0.04074	0.03805	0.03364	0.02926	0.02655	0.33648
7	0.09031	0.01103	0.01027	0.00904	0.00780	0.00703	0.09034
8	0.02108	0.00260	0.00241	0.00211	0.00181	0.00162	0.02108
10	0.00009	0.00001	0.00001	0.00001	0.00001	0.00001	0.00009
Total	0.54910	0.06659	0.06215	0.05492	0.04772	0.04325	0.54924

Table 3-7. Zone by Zone Averaged T6 for 2DBTOR Versus 2DB,
 Torus Case (Aspect Ratio = 5). Angle Measured in Degrees.

Zone Number	T6						
	2DB	0-36	36-72	72-108	108-144	144-180	0-360
4	0.05776	0.00642	0.00617	0.00576	0.00535	0.00510	0.05758
6	0.34143	0.03812	0.03655	0.03402	0.03152	0.03000	0.34043
7	0.25690	0.02888	0.02762	0.02560	0.02360	0.02237	0.25612
8	0.27286	0.03095	0.02950	0.02719	0.02491	0.02347	0.27204
10	0.03955	0.00454	0.00431	0.00395	0.00358	0.00335	0.03946
Total	0.96850	0.11732	0.10929	0.09559	0.08278	0.07509	0.96563

Table 3-8. Zone by Zone Averaged T7 for 2DBTOR Versus 2DB,
Torus Case (Aspect Ratio = 5). Angle Measured in Degrees.

Zone Number	T7						
	2DB	0-36	36-72	72-108	108-144	144-180	0-360
4	0.10119	0.01136	0.01089	0.01012	0.00935	0.00887	0.10117
6	0.33643	0.03786	0.03624	0.03362	0.03100	0.02938	0.33619
7	0.09031	0.01023	0.00977	0.00903	0.00829	0.00783	0.09030
8	0.02108	0.00240	0.00229	0.00211	0.00193	0.00182	0.02108
10	0.00009	0.00001	0.00001	0.00001	0.00001	0.00001	0.00009
Total	0.54910	0.06186	0.05920	0.05489	0.05058	0.04800	0.54882

torus is probably more accessible than that of the inner part. In addition, the TBR computed in the top zone (72° to 108°) was about the same for both aspect ratios. It also indicates that the toroidal TBR depended on volumetric effects, since an area on the outer part of the torus corresponds to a larger volume than an area on the inner part.¹⁸

D. L. Chapin of the Princeton Plasma Physics Laboratory presented the above idea concerning volumetric effects in a topical report.¹⁸ From Table 2-1, the volume occupied by a specified zone is

$$V_t = V_c \left[1 + \frac{r_{i+1}^3 - r_i^3}{3R(r_{i+1} - r_i) r_{ave}(\theta_{j+1} - \theta_j)} (\sin\theta_{j+1} - \sin\theta_j) \right] \quad (43)$$

where V_t = volume of the torus,

and $V_c = (r_{i+1} - r_i) r_{ave}(\theta_{j+1} - \theta_j)$ = volume of the cylinder with the same zone as the torus.

Thus, a volume factor, F_v , can be defined as

$$F_v = \frac{\text{volume of the torus}}{\text{volume of the cylinder}} = \left[1 + \frac{r_{i+1}^3 - r_i^3}{3R(r_{i+1} - r_i) r_{ave}(\theta_{j+1} - \theta_j)} (\sin\theta_{j+1} - \sin\theta_j) \right] \quad (44)^{18}$$

When F_v is multiplied by the TBR given by the cylinder, good agreement is expected between the torus TBR values and the volume corrected TBR values of the cylinder according to Chapin.¹⁸ Tables 3-9 through 3-12 show the differences between using volume corrections on the cylinder T6 and T7 for aspect ratios of 3 and 5 and then comparing to the torus's T6 and T7. This provides another check on the validity of 2DBTOR for doing toroidal scoping calculations, since Chapin's results were produced using a Monte Carlo code, which provides very accurate analysis for toroidal geometry. Since the TBR results from Tables 3-9 through 3-12 were in good agreement, it is apparent that 2DBTOR worked quite well for doing toroidal geometry scoping computations. Of course, the real test and usefulness of 2DBTOR will be for those designs which are not poloidally symmetric like the standard blanket used here.

Table 3-9. T6 for 2DBTOR Versus Volume Corrected T6 for 2DB,
 Torus Case (Aspect Ratio = 3). Angle Measured in Degrees.

Zone Number	T6						
	2DBTOR	0-36	36-72	72-108	108-144	144-180	0-360
4	0.05504	0.00699	0.00653	0.00578	0.00503	0.00456	0.05776
6	0.33914	0.04174	0.03884	0.03414	0.02945	0.02654	0.34143
7	0.25530	0.03194	0.02955	0.02569	0.02183	0.01944	0.25691
8	0.27123	0.03449	0.03174	0.02729	0.02283	0.02008	0.27286
10	0.03935	0.00518	0.00471	0.00396	0.00320	0.00273	0.03955
Total	0.96006	0.12034	0.11137	0.09685	0.08233	0.07336	0.96850

Table 3-10. T7 for 2DBTOR Versus Volume Corrected T7 for 2DB,
 Torus Case (Aspect Ratio = 3). Angle Measured in Degrees.

Zone Number	T7						
	2DBTOR	0-36	36-72	72-108	108-144	144-180	0-360
4	0.11240	0.01224	0.01143	0.01012	0.00881	0.00799	0.10119
6	0.33648	0.04113	0.03827	0.03364	0.02901	0.02615	0.33643
7	0.09034	0.01123	0.01039	0.00903	0.00767	0.00683	0.09031
8	0.02108	0.00267	0.00245	0.00211	0.00176	0.00155	0.02108
10	0.00009	0.00001	0.00001	0.00001	0.00001	0.00001	0.00009
Total	0.54924	0.06728	0.06256	0.05491	0.04726	0.04254	0.54910

Table 3-11. T6 for 2DBTOR Versus Volume Corrected T6 for 2DB,
 Torus Case (Aspect Ratio = 5). Angle Measured in Degrees.

Zone Number	T6						
	2DBTOR	0-36	36-72	72-108	108-144	144-180	0-360
4	0.05758	0.00650	0.00623	0.00578	0.00533	0.00505	0.05776
6	0.34043	0.03870	0.03696	0.03414	0.03133	0.02958	0.34143
7	0.25612	0.02944	0.02801	0.02569	0.02337	0.02194	0.25691
8	0.27204	0.03161	0.02996	0.02729	0.02461	0.02296	0.27286
10	0.03946	0.00469	0.00441	0.00396	0.00350	0.00322	0.03955
Total	0.96563	0.11095	0.10556	0.09685	0.08814	0.08275	0.96850

Table 3-12. T7 for 2DBTOR Versus Volume Corrected T7 for 2DB,
 Torus Case (Aspect Ratio = 5). Angle Measured in Degrees.

Zone Number	T7						
	2DBTOR 0-36	36-72	72-108	108-144	144-180	0-360	
4	0.10117	0.01139	0.01091	0.01012	0.00933	0.00884	0.10119
6	0.33619	0.03814	0.03642	0.03364	0.03089	0.02915	0.33643
7	0.09030	0.01035	0.00985	0.00903	0.00822	0.00771	0.09031
8	0.02108	0.00244	0.00231	0.00211	0.00190	0.00177	0.02108
10	0.00009	0.00001	0.00001	0.00001	0.00001	0.00001	0.00009
Total	0.54882	0.06233	0.05950	0.05491	0.05032	0.04749	0.54910

IV. SUMMARY AND CONCLUSIONS

The objective of the research performed here was to produce a scoping code that could be used for fusion reactor design. To this end, the present research initially explored a technique proposed by Pomraning and Stevens¹², in which a toroidal geometry diffusion problem is cast into cylindrical (r - θ) form by a spatially dependent redefinition of the diffusion coefficient, absorption cross-section, and extraneous source function. This idea suggested the approach of a direct finite differencing of the toroidal diffusion equation. Direct finite differencing proved to be more advantageous for incorporation into a computer program and to allow the curvature of the torus to be accounted for naturally.

The direct finite differencing approach was programmed into an existing two-dimensional(x - y , r - z , r - θ , triangular), multi-group neutron diffusion code, 2DB¹⁰, that had previously been converted to execute on the IBM-AT. Neutronic scoping calculations relevant to fusion reactor design were then performed in a micro-computer environment. This modified code was called 2DBTOR.

To verify that 2DBTOR was operating correctly, comparisons were made to both analytical and numerical solutions for several types of problems. ANISN and 2DB were used to verify and compare the solutions obtained from 2DBTOR. It was also shown that as the aspect ratio approached infinity (e. g., the major radius became large) the 2DBTOR solution approached the solution for that of 1-D cylindrical geometry. After verifying the solution for a large major radius, the errors associated with using a non-toroidal scoping code were examined versus 2DBTOR. Neutron

cross-sections for a benchmark problem were input into 2DBTOR and the output was compared to that from ANISN. A method proposed by Price and Chapin¹⁸, that used volume correction factors to compute the reaction rates in the benchmark blanket, was utilized to provide a means of checking 2DBTOR's results versus those given by a Monte Carlo code. It was also noted that 2DBTOR would also make possible the calculation of depletion in the fusion blanket, which was a unique advantage of the new program, 2DBTOR.

In future versions of the 2DBTOR program, it is anticipated that the central vacuum should be modelled through an internal boundary condition. A separate void streaming calculation will be used to define the internal boundary condition by specifying the neutron flux to current ratio as a function of position along the vacuum wall.¹⁹ Improved modelling of the central void region will be required if 2DBTOR is to prove to be an attractive program for blanket scoping calculations.

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APPENDIX A
POMRANING AND STEVENS' ANALYTICAL SOLUTION

For an infinite line source at the center of a non-fissioning medium with radius, a , the monoenergetic neutron diffusion equation is given by the following (assuming that both the diffusion coefficient does not vary with position and the solution is only radially dependent):

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \phi}{\partial r} \right) - k^2 \phi = Q \quad (\text{A-1})$$

where

$$k^2 = \frac{\Sigma_a}{D} ,$$

$$Q = \frac{S}{D} = 0 \text{ for } r > 0 ,$$

ϕ = the flux,

Σ_a = the macroscopic absorption cross section,

and

D = the diffusion coefficient .

Solving equation (A-1) gives the following:

$$\phi(r) = AK_0(kr) + CI_0(kr) \quad (\text{A-2})$$

where both A and C are constants. Using the boundary condition at $r=a$ where $\phi(a)=0$ allows C to be given as

$$C = -A \frac{K_o(ka)}{I_o(ka)}$$

which when substituted into equation (A-2) gives the following:

$$\phi(r) = A \left[K_o(kr) - \frac{K_o(ka)}{I_o(ka)} I_o(kr) \right] \quad (A-3)$$

In order to solve for A, one must use a source condition given by

$$\lim_{r \rightarrow 0} 2\pi r J(r) = S \quad (A-4)$$

where $J(r) = -D \frac{\partial \phi(r)}{\partial r}$.

It can be shown that when equation (A-3) is substituted into equation (A-4) the following results:

$$\lim_{r \rightarrow 0} r J(r) = \frac{S}{2\pi D} = \frac{Q}{2\pi} \quad (A-5)$$

Since equation (A-4) can also be written as

$$\lim_{r \rightarrow 0} r J(r) = \lim_{r \rightarrow 0} A k r K_1(kr) \quad (A-6)$$

which from Meghreblian and Holmes, Reactor Analysis, page 185, is equal to A, this implies that

$$A = \frac{Q}{2\pi} \quad (A-7)$$

When equation (A-7) is substituted into equation (A-3) the following results:

$$\phi(r) = \frac{Q}{2\pi} \left[K_0(kr) - \frac{K_0(ka)}{I_0(ka)} I_0(kr) \right] . \quad (\text{A-8})$$

In order to employ equation (A-8) for comparison to 2DB, it is necessary to determine a volumetric source that is equivalent to the line source used in deriving the flux solution. Thus, the line source is given by

$$S_l = S_v \pi r^2 \quad (\text{A-9})$$

where S_l = the line source (n/cm-sec)

and S_v = the volumetric source (n/cm³-sec).

For Q=1 and since Q=S_l/D, then this implies that S_l = D. When S_l = D is substituted into equation (A-9), the resulting equation for S_v is given by the following:

$$S_v = \frac{D}{\pi r^2} . \quad (\text{A-10})$$

This is the S_v value that will be input to 2DB so that comparisons to the analytic solution can be made.

APPENDIX B
AVERAGE D_k SOLUTION

From continuity of current for two adjacent mesh points, k and $k+1$ (see Figure B-1), it can be shown that

$$\frac{D_k}{\frac{\delta r_k}{2}} (\phi^{1/2} - \phi_k) = \frac{D_{k+1}}{\frac{\delta r_{k+1}}{2}} (\phi_{k+1} - \phi^{1/2}) \quad (\text{B-1})$$

where $\frac{\delta r_k}{2}$ = distance from mesh point k to mesh point $1/2$

and $\frac{\delta r_{k+1}}{2}$ = distance from mesh point $k+1$ to mesh point $1/2$.

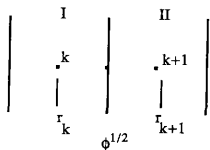


Figure B-1. Mesh points

Rearranging equation (B-1) gives the following:

$$\left[\frac{D_k}{\frac{\delta r_k}{2}} + \frac{D_{k+1}}{\frac{\delta r_{k+1}}{2}} \right] \phi^{1/2} = \frac{D_{k+1}}{\frac{\delta r_{k+1}}{2}} \phi_{k+1} + \frac{D_k}{\frac{\delta r_k}{2}} \phi_k \quad (\text{B-2})$$

which when solved for $\phi^{1/2}$ yields

$$\phi^{1/2} = \frac{\frac{D_{k+1}}{\delta r_{k+1}} \phi_{k+1} + \frac{D_k}{\delta r_k} \phi_k}{\left[\frac{D_k}{\delta r_k} + \frac{D_{k+1}}{\delta r_{k+1}} \right]} \quad (\text{B-3})$$

Substituting the RHS of equation (B-3) for the left-hand term in equation (B-1) gives the following:

$$J_k = \frac{D_{k+1}}{\delta r_{k+1}} \frac{D_k}{\delta r_k} \frac{1}{\left[\frac{D_k}{\delta r_k} + \frac{D_{k+1}}{\delta r_{k+1}} \right]} (\phi_{k+1} - \phi_k) \quad (\text{B-4})$$

Further simplifying equation (B-4) gives

$$J_k = \frac{D_{k+1} D_k (\delta r_k + \delta r_{k+1})}{[D_k \delta r_{k+1} + D_{k+1} \delta r_k]} \left(\frac{\phi_{k+1} - \phi_k}{r_{k+1} - r_k} \right) \quad (\text{B-5})$$

where $r_{k+1} - r_k = \frac{\delta r_k}{2} + \frac{\delta r_{k+1}}{2}$

Thus, D can be defined to be

$$\bar{D} = \frac{D_{k+1}D_k(\delta r_k + \delta r_{k+1})}{[D_k \delta r_{k+1} + D_{k+1} \delta r_k]} , \quad (\text{B-6})$$

to yield

$$J_k = \bar{D} \left(\frac{\phi_{k+1} - \phi_k}{r_{k+1} - r_k} \right) . \quad (\text{B-7})$$

APPENDIX C
XPROC2DB LISTING

```

DIMENSION A(30,46,4,4),B(30,46,4,4)
C  OPEN(UNIT=9,FILE='STDXS.C.DAT',STATUS='OLD')
C  OPEN(UNIT=10,FILE='STDXS.C2.DAT',STATUS='UNKNOWN')
  OPEN(UNIT=9,FILE='XPROC30.DAT',STATUS='OLD')
  OPEN(UNIT=10,FILE='XPROC30.OUT',STATUS='UNKNOWN')

C  K = NUMBER OF MATERIALS
C  I = GROUP #
C  J = XSC TABLE #
C  NMAT=NUMBER OF MATERIALS
C  NGRUP=# OF GROUPS
C  NTAB=XSEC TABLE LENGTH
C  NSCAT=SCATTERING ORDER

  NSCAT=3
  NMAT=4
  NGRUP=30
  NTAB=46

DO 30 K=1,NMAT
  DO 25 L=1,NSCAT+1
    DO 20 I=1,NGRUP
      READ(9,10) (A(I,J,K,L),J=1,NTAB)
10     FORMAT(1X/,9(5(E13.5)),E13.5)
      A(I,14,K,L)=A(I,8,K,L)

```

```
16          FORMAT(78X,I2)
20    CONTINUE
      READ(9,13)
13    FORMAT(1X)
25    CONTINUE
      READ(9,13)
30    CONTINUE
      DO 50 K=1,NMAT
        L=1
          DO 40 I=1,NGRUP
            WRITE(10,10) (A(I,J,K,L),J=1,NTAB)
            WRITE(10,99)
99          FORMAT(3X,' T',2X)
40        CONTINUE
            WRITE(10,16)K
            WRITE(10,13)
50      CONTINUE
        K=4
        L=1
          DO 60 I=1,NGRUP
            A(I,17,K,L)=A(I,16,K,L)
60      CONTINUE
        K=4
```

```
L=1
      DO 70 I=1,NGRUP
        DO 65 J=1,15
          A(I,J,K,L)=0.0
65      CONTINUE
70     CONTINUE
      DO 170 I=1,NGRUP
        DO 165 J=18,NTAB
          A(I,J,K,L)=0.0
165    CONTINUE
170   CONTINUE

K=4
L=1
      DO 260 I=1,NGRUP
        WRITE(10,10) (A(I,J,K,L),J=1,NTAB)
        WRITE(10,99)
260   CONTINUE
        WRITE(10,16)K
        WRITE(10,13)

STOP
END
```


APPENDIX D
2DB ERRATA

I. The area elements for r- θ geometry can be shown to be given by

$$A_r = r_{ave0} \Delta\theta_0 \quad (D-1)$$

$$A_\theta = \Delta r_0 \quad (D-2)$$

where A_r = radial area element and A_θ = axial area element for a volume element, 0.

In the original 2DB code, the above area elements included a 2π factor. Since this is incorrect, the 2π terms were deleted (see Appendix H, subroutine INIT). For x-y geometry, the axial area element stays the same as above, but the radial element does not have a r_{ave0} term and thus, is given by

$$A_r = \Delta\theta_0 \quad (D-3)$$

Again, the area elements in 2DB included a 2π term. The 2π terms were deleted to cause 2DB to correctly solve the diffusion equation (see Appendix H, subroutine INIT).

II. For the downscattering term in the diffusion equation, 2DB uses the following:

$$XD = \Sigma_{tr} - \Sigma_s - \Sigma_a \quad (D-4)$$

where XD = the sum of the downscattering cross sections.

It can be shown that

$$\Sigma_{tr} = \Sigma_a + \Sigma_s + XD - \Sigma_{n,2n} \quad (D-5)$$

where $\Sigma_{n,2n}$ = (n,2n) cross section.

Substituting equation (D-5) into equation (D-4) gives the following:

$$XD = XD - \Sigma_{n,2n} \quad (D-6)$$

Since XD cannot equal $XD - \Sigma_{n,2n}$, this implies that 2DB has incorrectly calculated the downscattering component for those problems which have a significant amount of (n,2n) scattering reactions taking place. Therefore, 2DB was changed so that the downscattering term was calculated by summing over all the downscattering cross sections for a specified group of neutrons (see Appendix H, subroutine S860).

APPENDIX E

PLASMA SOURCE

The problem to be solved is the determination of the uniformly distributed source of 14.1 MeV neutrons in the plasma (radius = R_{pl}) impinging upon the blanket's inner wall at R_w . It is assumed that 10 MW/m^2 will be the maximum wall load, thus, the current on the inner wall is given by

$$J = \frac{10\text{MW}}{\text{m}^2} \left(\frac{10^6 \text{W}}{\text{MW}} \right) \left(\frac{\text{J/s}}{\text{W}} \right) \left(\frac{\text{m}^2}{10^4 \text{cm}^2} \right) \left(\frac{\text{MeV}}{1.6 \times 10^{-13} \text{J}} \right) \left(\frac{1}{14.1 \text{ MeV/n}} \right) \quad (\text{E-1})$$

$$= 4.43 \times 10^{14} \frac{\text{n}}{\text{cm}^2 \cdot \text{sec}} .$$

Equating the number of neutrons emitted in the plasma to the number of neutrons impinging upon the inner wall, then for a unit width in the toroidal direction the balance is given by

$$J 2\pi R_w = \text{FRV} \pi R_{pl}^2 \quad (\text{E-2})$$

where FRV = the uniformly distributed source of 14.1 MeV neutrons in the plasma.

Rearranging equation (E-2) to solve for FRV gives the following:

$$\text{FRV} = \frac{J 2\pi R_w}{R_{pl}^2} . \quad (\text{E-3})$$

Substituting for J from equation (E-1), while letting $R_w = 200\text{cm}$ and $R_{pl} = 150\text{cm}$, the uniformly distributed source of 14.1 MeV neutrons in the plasma is equal to $7.88 \times 10^{12} \text{ n/cm}^2\text{-sec}$.

APPENDIX F
EQUIVALENT SOURCE TO THE PLASMA SOURCE
AT THE INNER EDGE OF THE BLANKET

The problem to be solved is the determination of the equivalent source at the inner edge of the blanket (radius = R_w) to the of the uniformly distributed source of 14.1 MeV neutrons in the plasma (radius = R_{pl}). Equating the number of neutrons emitted from the plasma to the number of neutrons emitted from a thin source (width = Δr) at the inner wall (radius = R_w), then for a unit width in the toroidal direction the balance is given by

$$\frac{FRV \pi R_{pl}^2}{2\pi R_{pl}} = \frac{S_e \pi ((R_w + \Delta r)^2 - R_w^2)}{2\pi R_w} \quad (F-1)$$

where FRV = the uniformly distributed source of 14.1 MeV neutrons in the plasma and S_e = equivalent source at the inner wall. Rearranging and simplifying equation (F-1) to solve for S_e gives the following:

$$S_e = \frac{FRV R_{pl} R_w}{((R_w + \Delta r)^2 - R_w^2)} \quad (F-2)$$

APPENDIX G
2DBTOR MANUAL

2DBTOR MANUAL

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Acknowledgements

Much of this manual for 2DBTOR was influenced by both the previous 2DB manual (2DB User's Manual written by W. W. Little, Jr. and R. W. Hardie) and the ANISN/PC manual (ANISN/PC Manual written by D. Kent Parsons). Some of the sections of both manuals were incorporated into the 2DBTOR manual, although the style was based for the most part on the ANISN/PC manual.

Contents

ACKNOWLEDGEMENTS.....	86
ABSTRACT.....	89
1. INTRODUCTION.....	90
2. THEORETICAL FOUNDATIONS.....	91
2.1 Discretization of the Diffusion Equation.....	91
2.1.1 Energy Discretization.....	91
2.1.2 Spatial Discretization -- The Finite Difference Method.....	94
2.2 Discussion of Boundary Conditions Used in 2DBTOR.....	97
2.2.1 Zero Flux Gradient Boundary Condition.....	97
2.2.2 Zero Flux Boundary Condition.....	98
2.2.3 Periodic Boundary Condition.....	99
2.3 Discussion of Triangular Mesh Option.....	100
2.4 Iterative Solution Methods of 2DBTOR.....	101
2.4.1 Outer Iterations.....	102
2.4.2 Inner Iterations.....	103
3. USER'S GUIDE.....	105
3.1 2DBTOR Input Data Description (Logical Unit 5).....	105
3.2 Supplemental Input Information for 2DBTOR.....	121
3.2.1 Cross Section Considerations.....	121
3.2.2 Search Considerations.....	123
3.2.2.1 Time Absorption (α Calculation).....	123
3.2.2.2 Material Concentration (C Calculation).....	124

3.2.2.3 Zone Dimensions (δ Calculation).....	126
3.2.2.4 Buckling (B^2 Calculation).....	126
3.2.3 Burnup Model.....	126
3.2.3.1 Remarks on Burnup Equations.....	128
3.2.4 Source Problems.....	128
3.2.5 Remarks on Code Operation.....	129
3.3 2DBTOR Format Description.....	130
3.3.1 Operators.....	130
3.4 Printed Output Description.....	131
4. PROGRAMMER'S GUIDE TO 2DBTOR.....	133
4.1 2DBTOR Files Description.....	133
4.2 2DBTOR Subroutines.....	134
4.2.1 Subroutine Calling Sequences.....	136
4.3 Selected Definitions of Variables and Arrays Used.....	138
5. REFERENCES.....	146

2DBTOR MANUAL -- REVISION OF 2DB MANUAL (REVISION 1)

ABSTRACT

1. Program Identification: 2DBTOR is a revised version of 2DB¹.
2. Description of Problem: 2DBTOR is a two-dimensional (X-Y, R-Z, R- θ , triangular, toroidal), multi-group neutron diffusion code for use in fast reactor criticality and burnup analysis. In addition, 2DBTOR may also be used to study fusion blanket problems in toroidal geometry. 2DBTOR solves the multi-group diffusion theory eigenvalue, adjoint, time absorption, fixed source and criticality search (concentration, zone thickness, and buckling) problems.
3. Method of Solution: Multi-group finite difference neutron diffusion equations are solved iteratively in 2DBTOR. The power method, accelerated by a fission source over-relaxation factor calculated in the code, is used for the outer iterations. Inner iterations are accelerated by use of an over-relaxation factor input by the user.

1. INTRODUCTION

A number of significant additions and alterations (e.g., a toroidal geometry option and an activity cross section option) have been made to the 2DB¹ diffusion-burnup code. In addition, some bugs were discovered in the old 2DB version which have been corrected in the current code, 2DBTOR. This manual gives a complete description of the code including all modifications. A description of both the mathematical model and user instructions are given in the body of the report; a sample problem is included in the appendix.

2DBTOR is designed for use in both fast reactor and fusion analysis. Eigenvalues are computed by standard source-iteration techniques. Group rebalancing and successive over-relaxation with line inversion are used to accelerate convergence. Adjoint solutions are obtained by inverting the input data and redefining the source terms.

Variable dimensioning is used to make maximum use of the available fast memory. Since only one energy group is in the fast memory at any given time, the storage requirements are insensitive to the number of energy groups.

Criticality searches can be performed on buckling, time absorption, material concentrations, and region dimensions. Alpha and k_{eff} can be used as parametric eigenvalues. Criticality searches can be performed during burnup to compensate for fuel depletion.

2. THEORETICAL FOUNDATIONS

2.1. Discretization of the Diffusion Equation

2.1.1. Energy Discretization. The diffusion equation can be developed from the equation for the time rate of change of the number of neutrons within dE of energy, E , in an arbitrary differential volume, dV . Neutrons of energies between E and $E+dE$, within dV , can be lost or gained by a variety of processes including: (1) production directly from a source, (2) absorption, (3) leakage and (4) scattering. The time rate of change of the number of neutrons in dV and between E and $E+dE$ can be obtained by integrating the neutron density ($n(r,E,t)$) over dV , and balancing this with the gains and losses as follows:

$$\frac{\partial}{\partial t} \int_V \frac{\Phi(\vec{r}, E, t)}{\nu} dV = \left[\begin{array}{c} \text{source neutron} \\ \text{production rate} \\ \text{in } V \text{ at } E \end{array} \right] - \left[\begin{array}{c} \text{absorption} \\ \text{rate in } V \\ \text{at } E \end{array} \right] - \left[\begin{array}{c} \text{change due} \\ \text{to leakage} \\ \text{from } V \text{ at } E \end{array} \right] - \left[\begin{array}{c} \text{neutron scattering} \\ \text{rate out of} \\ E \text{ in } V \end{array} \right] + \left[\begin{array}{c} \text{neutron scattering} \\ \text{rate into} \\ E \text{ in } V \end{array} \right] \quad (1)$$

where $\Phi(r,E,t)$ = flux of neutrons at r , E and t = $n(r,E,t)\nu$

and ν = the speed of the neutrons at E .

Equation (1) is known as the neutron continuity equation. Since the energy

dependence of the neutron cross sections vary, equation (1) is usually solved for discrete energy groups (groups denoted by g in this case); thus equation (1) can be written as

$$\frac{\partial}{\partial t} \int_V \frac{\Phi_g(\vec{r}, t)}{v_g} dV = \left[\begin{array}{c} \text{source neutron} \\ \text{production rate} \\ \text{in V for} \\ \text{group, } g \end{array} \right] - \left[\begin{array}{c} \text{absorption} \\ \text{rate in V} \\ \text{for group, } g \end{array} \right] - \left[\begin{array}{c} \text{change due} \\ \text{to leakage} \\ \text{from V} \\ \text{for group, } g \end{array} \right] \\ - \left[\begin{array}{c} \text{neutron scattering} \\ \text{rate out of} \\ \text{group, } g \text{ in V} \end{array} \right] + \left[\begin{array}{c} \text{neutron scattering} \\ \text{rate into} \\ \text{group, } g \text{ in V} \end{array} \right] \quad (2)$$

Assuming steady state ($\partial/\partial t$ term = 0) and no upscattering, equation (2) becomes upon substituting the corresponding mathematical expressions for the RHS terms,

$$-\vec{\nabla} \cdot (D_g \vec{\nabla} \Phi_g) + \Sigma_{a_g} \Phi_g + \sum_{g'=g+1}^G (\Sigma_s(g \rightarrow g') \Phi_{g'}) = \\ S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f_{g'}} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_s(g \rightarrow g') \Phi_{g'}) \quad (3)$$

where the (\vec{r}, t) arguments have been dropped for clarity. This is the multi-group neutron diffusion equation.

Equation (3) can be further simplified by noting that the removal of neutrons from group g is caused by both downscattering and absorption. The removal cross section for group g is defined as shown below:

$$\begin{aligned}
\Sigma_g^r &\equiv \Sigma_{a_g} + \sum_{g'=g+1}^G (\Sigma_s(g \rightarrow g')) \\
&= \Sigma_{a_g} + \Sigma_s(g \rightarrow g) + \left[\sum_{g'=g+1}^G \Sigma_s(g \rightarrow g') \right] - \Sigma_s(g \rightarrow g) \\
&= \Sigma_{tr_g} - \Sigma_s(g \rightarrow g')
\end{aligned}$$

where Σ_{tr_g} = the macroscopic transport cross section = $1/(3D_g)$.¹⁴

Thus the removal rate/cm³ is

$$\Sigma_g^r \Phi_g = \Sigma_{a_g} \Phi_g + \sum_{g'=g+1}^G (\Sigma_s(g \rightarrow g') \Phi_{g'}). \quad (4)$$

Substituting equation (4) into equation (3) gives

$$-\vec{\nabla} \cdot (D_g \vec{\nabla} \Phi_g) + \Sigma_g^r \Phi_g = S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f_{g'}} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_s(g \rightarrow g') \Phi_{g'}) \quad (5)$$

the form of the multi-group neutron diffusion equation used in 2DBTOR.

2.1.2. Spatial Discretization - The Finite Difference Method. To develop a finite difference approximation for the diffusion equation (with axial symmetry), it is first necessary to integrate equation (5) over a small, arbitrary volume ΔV (see Figure 2-1) where the mesh points are considered to be in the center of the homogeneous mesh interval. Thus,

$$\begin{aligned}
 & - \int_{\Delta V} \vec{\nabla} \cdot D_g \vec{\nabla} \Phi_g dV + \int_{\Delta V} [\Sigma_g^r \Phi_g] dV \\
 & = \int_{\Delta V} \left[S_g^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f,g'} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_s(g \rightarrow g') \Phi_{g'}) \right] dV
 \end{aligned} \tag{6}$$

where the first term on the LHS of the equation is the leakage term, the second term on the LHS of the equation is the removal term, and the RHS represents the source terms including fission and scatter. Thus for the removal term,

$$\int_{\Delta V} \Sigma_g^r \Phi_g dV = \Sigma_{g_0}^r \Phi_{g_0} V_o \tag{7}$$

where Φ_{g_0} = flux associated with meshpoint o

and $\Sigma_{g_0}^r$ = removal cross section associated with meshpoint o.

The source term on the RHS is done similarly to the above. The leakage term is changed to an integral over the surface area of the volume element, thus from the Divergence Theorem

$$-\int_{\Delta V} \vec{\nabla} \cdot D_g \vec{\nabla} \Phi_g dV = -\int_A D_g \vec{\nabla} \Phi_g \cdot \hat{n} dA .$$

The flux partial derivatives will be obtained by differencing the two neighboring flux values. Thus, volume integration of equation (5) for mesh point 0 (see Figure 2-1; where r stands for x or r, and θ stands for y, z, or θ , depending on the geometry) leads to the expression

$$-\sum_{k=1}^4 \left[\bar{D}_{gk} \left(\frac{\phi_{gk} - \phi_{g0}}{L_k} \right) A_k \right] + \Sigma_{g0}^r \Phi_{g0} V_0 \quad (8)$$

$$= \left[S_{g0}^{\text{ext}} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f_{g'}} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_{s_0}(g \rightarrow g') \Phi_{g'}) \right] V_0 .$$

where

$\Sigma_{g_0}^r$ = removal cross section associated with mesh point 0,

$S_{g_0}^{ext}$ = extraneous source rate associated with mesh point 0,

V_0 = volume associated with mesh point 0,

Φ_{g_k} = flux associated with mesh point 0,

L_k = distance between mesh point k and mesh point 0,

A_k = area of boundary between mesh point k and mesh point 0,

and

\bar{D}_{g_k} = effective diffusion coefficient mesh point k and mesh point 0

$$\left(= \frac{D_{g_0} D_{g_k} (\Delta r_0 + \Delta r_k)}{(D_{g_0} \Delta r_k + D_{g_k} \Delta r_0)} \text{ or } \frac{D_{g_0} D_{g_k} (\Delta \theta_0 + \Delta \theta_k)}{(D_{g_0} \Delta \theta_k + D_{g_k} \Delta \theta_0)} \right)$$

where $\Delta r_k = \Delta r$ for volume element k,
 $\Delta \theta_k = \Delta \theta$ for volume element k

Equation (8) can be recast into a form more convenient for performing flux iterations by rearranging equation (8) to that given below:

$$\phi_0 = \frac{\left[S_{g_0}^{ext} + \chi_g \sum_{g'=1}^G (v_{g'} \Sigma_{f_{g',g'}} \Phi_{g'}) + \sum_{g'=1}^{g-1} (\Sigma_{s_{g',g}} (g-g') \Phi_{g'}) \right] V_0 + \sum_{k=1}^4 C_k \phi_k}{C_5} \quad (9)$$

$$\text{where } C_k = \frac{\bar{D}_{g_k} A_k}{L_k} \quad k=1, \dots, 4 \quad (10)$$

$$\text{and } C_5 = \Sigma_{s_0}^r V_0 + \sum_{k=1}^4 C_k \quad (11)$$

2.2. Discussion of Boundary Conditions Used in 2DBTOR

Three boundary conditions are available in 2DBTOR: reflective, extrapolated vacuum, and periodic. The reflective boundary condition is used on boundaries where $\nabla\Phi = 0$; the extrapolated vacuum boundary condition is used on boundaries where the flux is assumed to be zero at the extrapolated boundary; and the periodic boundary condition is used on boundaries where material conditions are repeating. The above mentioned boundary conditions are described in more detail below.

2.2.1. Zero Flux Gradient Boundary Condition. Consider the left-hand boundary of the one-dimensional reactor shown in Figure 2-2. Imagine that a

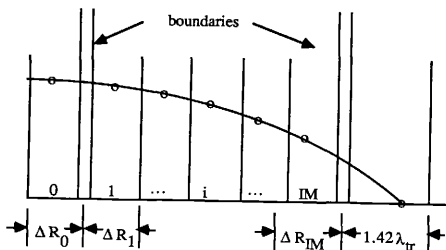


Figure 2-2. Schematic Diagram of 1-D Reactor

pseudo-mesh interval, interval 0, has been added on the left-hand side of the boundary with the same composition and thickness of interval 1. Clearly, then, if $\nabla\Phi = 0$ at the boundary, $\phi_0 = \phi_1$. Therefore, since $(\phi_0 - \phi_1) = 0$, the coefficient of $(\phi_0 - \phi_1)$, C_1 (see equations (8) and (10)), is immaterial -- hence C_1 can be set equal to zero. The calculation is therefore performed assuming that ϕ_0 does not exist and $C_1 = 0$.

2.2.2. Zero Flux Boundary Condition. Again, imagine that a pseudo-mesh interval with the same composition as interval IM has been added to the right hand side of the boundary in Figure 2-2. Now, since $\phi_{IM} \neq 0$ and $\phi_{IM+1} = 0$, the coefficient of $(\phi_{IM} - \phi_{IM+1})$ in equation (8) cannot be disregarded. In fact, from equation (10), it is clear that

$$C_k = \frac{D_k A_k}{.5 \Delta R_{IM} + .71 \lambda_{tr}}$$

where λ_{tr} is assumed to equal $1/\Sigma_{tr}$

Note, as in the reflective boundary condition case, that there is no contribution of the pseudo-flux in equation (9). For a zero flux gradient, $C_k = 0$; whereas for a zero flux, $\phi_k = 0$.

2.2.3. Periodic Boundary Condition. Periodic boundary conditions are only available for the top and bottom boundaries(e.g., boundaries in the y, z, or θ direction). In this option(see Figure 2-3),

$$\phi_{IM} = \phi_0$$

$$\phi_{IM+1} = \phi_1$$

and

$$C_{k(1 \rightarrow IM)} = \frac{\bar{D}_k A_k}{.5(\Delta R_1 + \Delta R_{IM})}$$

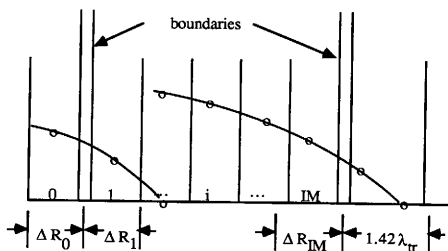


Figure 2-3. Schematic Diagram of 1-D Reactor

It should be stressed that the pseudo-mesh intervals discussed above are not in any way a part of the code. They are mentioned here only for heuristic purposes.

2.3. Discussion of Triangular Mesh Option

Since most fast reactors are composed of hexagonal assemblies, 2DBTOR includes a triangular mesh option. Hexagons are formed by appropriate grouping of six triangular mesh intervals.

In the triangular mesh option, the (i,j) coordinate grid is composed of a rectangular array of triangles. As in the other geometry options, the mesh points are placed in the center of each interval, or triangle. See Figure 2-4 for a simple 3 x 4 example. In contrast to the other geometry options, however, the mesh boundaries must be equally spaced. In fact, the radial (RB_i) and axial (θB_j) mesh boundaries must be computed by the expressions

$$RB_i = (i-1) \frac{FTF}{2\sqrt{3}}, \quad i=1, \dots, IM+1 \quad (12)$$

$$\theta B_j = (j-1) \frac{FTF}{2}, \quad j=1, \dots, JM+1 \quad (13)$$

where FTF is the flat-to-flat hexagon width.

Only vacuum and reflective boundary conditions are available with the triangular mesh option. The user is cautioned against using reflective left and right boundaries since this implies no surface leakage from each mesh interval on the left and right border.

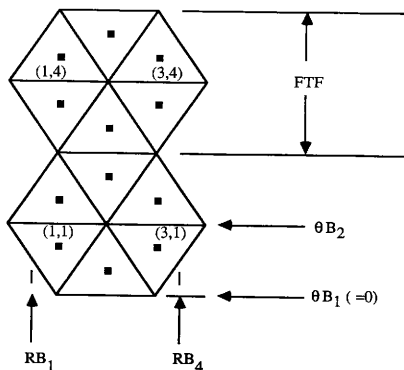


Figure 2-4. Triangular Mesh Example (3 x 4)

2.4. Iterative Solution Methods of 2DBTOR

Within the 2DBTOR code, two distinct levels of iteration may be found for general problems. The top level of iteration (i.e., outer iterations) is for the spatially and energy-group summed fission sources. The second level of iteration (i.e., inner iterations) is for the individual group fluxes that result from a given source.

The following sections describe both the inner and outer iteration procedures used in 2DBTOR and discuss the methods used to accelerate those procedures.

2.4.1. Outer Iterations. Outer iterations in 2DBTOR are based on the power iteration method. That is, at each outer iteration, a total fission source is calculated. Upon convergence, the ratio of the latest fission source to the previous fission source is the eigenvalue. Thus, the eigenvalue is used to renormalize the fission sources between outer iterations, and the ratio between normalized fission source iterates approaches 1.0.

For search problems, the eigenvalue is defined to be the value of the search quantity (e.g., time absorption, zone thickness, etc.) that produces criticality. In these problems, the eigenvalue is used to change the problem at each search step so that the fission source ratio still approaches 1.0 at convergence.

For each outer iteration, the inner iteration procedure starts with group 1 and sweeps through the groups in order of decreasing energy. The downscattering component of the source for the current group is calculated from the latest values of the higher energy fluxes.

Fission source over-relaxation is employed in 2DBTOR to accelerate convergence. The procedure is as follows: After the new fission source rate profile, F^{v+1}_1 , is calculated, a second "new" value, F^{v+1}_2 , is computed by magnifying the difference between the new fission source rate and the old fission source rate. Thus,

$$F_2^{v+1} = F^v + \beta'(F_1^{v+1} - F^v),$$

where β' is the fission source over-relaxation factor. F^{v+1}_2 is then normalized to give the same total source as F^{v+1}_1 .

2.4.2. Inner Iterations. Inner iterations are computed using successive line over-relaxation (SLOR). That is, the fluxes on each vertical (or horizontal) line are simultaneously computed (by the familiar Crout reduction technique) and then over-relaxed using

$$\phi^{v+1} = \phi^v - \beta(\phi^{v+1} - \phi^v) ,$$

where β is the over-relaxation factor. In r-q problems or problems involving periodic boundary conditions, direct inversion is performed on vertical lines beginning at the left boundary and proceeding by column to the right boundary. In triangular problems, direct inversion is performed along horizontal lines beginning at the bottom boundary and proceeding by row to the top boundary. In all other situations, direct inversion is performed along the dimension with the most mesh points. One mesh sweep is defined as one inner iteration.

The flux over-relaxation factor, β , is an input parameter. The fission source over-relaxation factor, β' , is computed internally from the expression

$$\beta' = 1.0 + 0.6(\beta - 1) .$$

As in the original version of 2DB, the flux in each group is normalized (by balancing the total source rate and loss rate) immediately before each group-flux calculation. Thus, a one-region problem with zero-gradient boundary conditions

would be solved in exactly one outer iteration.

It should be mentioned that an alternating direction SLOR scheme (using line inversion for rows and then columns in alternation) is included as an option to enhance convergence for problems involving tight mesh spacing in both directions.

3. USER'S GUIDE

3.1. 2DBTOR Input Data Description (Logical Unit 5)

This section describes the input data for the 2DBTOR code. The input has the following general structure:

<u>Section</u>	<u>Description</u>
A	Title Card
B	Single Integer and Real Numbers
C	Cross Section Data
D	Fixed Source Data (if needed)
E	Miscellaneous Data
F	Burnup Data

Each input section is begun on a new line and ended on a later line by a terminate marker. A "T" anywhere on a nontitle card or on a line by itself constitutes a terminate marker.

Each of the input sections will be described below. Locations of the terminate markers will also be given. The length of each data section is denoted by the number or variable in slashes "/" by the input description. The expression in braces "{}" by the input description is the condition that requires the input to be present. If no condition is given, the input is always needed.

The data format conventions used by 2DBTOR are described fully in Section 3.3. Succinctly described, however, 2DBTOR uses free format augmented by an operator notation, which conserves space in the input.

Further details about the various input options for 2DBTOR may be found in Section 3.2, which immediately follows the input data description.

A. Title Card - (4A20 Format)

B. Single Integer and Real Numbers

B.1 Integer control parameters /13/

	Variable	
	<u>Name</u>	<u>Description</u>
1.	A02	0 - regular calculation 1 - adjoint calculation
2.	I04	Eigenvalue type 0 - distributed source (D.) 1 - eigenvalue calculation 2 - time absorption (α) search calculation (E.7.) 3 - concentration search (C) calculation

Variable		
	<u>Name</u>	<u>Description</u>
		4 - zone thickness search (δ) calculation (E.11.-14.)
		5 - buckling (B^2) calculation (E.5.)
3.	S02	Parametric eigenvalue type 0 - no effect 1 - eigenvalue (k_{eff}) 2 - α (E.7.)
4.	IGM	Number of energy groups
5.	IHT	Position of σ_{tr}
6.	NXCM	Number of downscattering terms
7.	MCR	Number of cross section sets to be read from cards/tape (+N/-N) (+N go to C.)
8.	G07	Maximum number of inner iterations for each group per outer iteration (suggested value ≥ 20)

Variable		
	<u>Name</u>	<u>Description</u>
9.	D05	Maximum number of outer iterations(suggested value ≥ 50 for general usage but D05 = 1 fixed source calculations without fission)
10.	MAXT	Maximum run time (minutes)
11.	NPRT	Print option 0 - mini 1 - midi 2 - cross sections 3 - fluxes
12.	M07	Flux guess 0 - no effect 1 - flux guess from file FOR14.DAT
13.	NPUN	Flux dump 0 - no effect 1 - flux dump to file FOR16.DAT

Terminate Marker

B.2 Integer control parameters /14/

Variable		
	<u>Name</u>	<u>Description</u>
1.	IGE	Geometry parameter 0 - X-Y 1 - R-Z 2 - R- θ or toroidal (B.2.2.) 3 - triangular 4 - toroidal
2.	ITOR	Toroidal specifier (0/1 = R- θ /toroidal)
3.	NACT	Number of activities (E.15.) 0 - no effect >0 - read table positions for N activities
4.	IM	Number of radial fine mesh intervals
5.	JM	Number of axial fine mesh intervals
6.	IZM	Number of zones (or regions)
7.	MT	Total number of material (MCRI + mixtures formed in

Variable		
	<u>Name</u>	<u>Description</u>
		mixing table)
8.	M01	Mixing table length
9.	B01	Left boundary condition
		0 - vacuum
		1 - reflection
10.	B02	Right boundary condition
		0 - vacuum
		1 - reflection
11.	B03	Top boundary condition
		0 - vacuum
		1 - reflection
		2 - periodic
12.	B04	Bottom boundary condition
		0 - vacuum
		1 - reflection
		2 - periodic
13.	IZ	Number of radial zones
		0 - no effect
		>0 - only if I04 = 4
14.	JZ	Number of axial zones
		0 - no effect
		>0 - only if I04 = 4

Terminate Marker

B.3 Floating point control parameters /6/

	Variable	
	<u>Name</u>	<u>Description</u>
1.	EV	First eigenvalue guess
2.	EVM	Initial eigenvalue modifier for search problems, zero otherwise
3.	S03	parametric eigenvalue
4.	BUCK	Buckling (cm^{-2}) (E.5)
5.	LAL	Lower limit for λ in search calculations, zero otherwise
6.	LAH	Upper limit for λ in search calculations, zero otherwise

Terminate Marker

B.4 Floating point control parameters [6]

Variable		
	<u>Name</u>	<u>Description</u>
1.	EPS	Eigenvalue relative convergence criterion (suggested value, EPS = 0.0001)
2.	EPSA	Parametric convergence criterion
3.	G06	Pointwise flux relative convergence criterion, zero otherwise (suggested value, EPSA = 2.0*EPS)
4.	POD	Parameter oscillation damper (suggested value, POD = 1.0)
5.	ORF	Over relaxation factor (suggested value, $1.0 \leq \text{ORF} \leq 2.0$)
6.	S01	S01 < 0, power (MWT) for R-Z geometry, (power/height (MWT/cm) for all but R-Z geometry) S01 > 0, neutron source rate (n/cm ³) (0.0 for source without fission)

Terminate Marker

B.4 Floating point control parameters // {ITOR = 1}

Variable	
Name	Description
1. BGR	Major radius (cm)

Terminate Marker

C. Cross Section Data {MCR>0}

Variable	
Name	Description
1. HOLN(MCR)	Name of isotope
2. ATW(MCR)	Atomic weight of isotope

Terminate Marker

$ITL = NXCM + IHT + 1 =$ Cross section table length

3. C(ITL,IGM,MT) Read cross sections
for first group /ITL/.

4. *Terminate Marker*

5. Repeat C.3 and C.4 for all groups /IGM-1/

6. Repeat C.1 - C.5 for all materials /MT-1/

D. Fixed Source Data (I04 = 0)

	Variable	
	<u>Name</u>	<u>Description</u>
1.	S2(IM*JM)	Source in first group

2. *Terminate Marker*

3. Repeat D.1 - D.2 for all groups /IGM-1/

E. Miscellaneous Data

	<u>Description</u>
1.	Radial mesh line coordinates defining the IM fine mesh intervals /IM+1/ (should be strictly ascending in order)

Terminate Marker

Description

2. Axial mesh line coordinates defining the JM fine mesh intervals /JM+1/ (should be strictly ascending in order)

Terminate Marker

3. Zone numbers by fine mesh interval /IM*JM/

Terminate Marker

4. Material numbers by zone /IZM/

Terminate Marker

5. Buckling coefficients by zone /IZM/ {I04=5 or BUCK>0}

Terminate Marker

6. Fission spectrum data /IGM/ (the sum of the entries should equal 1.0 for eigenvalue calculations (I04=1))

Terminate Marker

Description

7. Neutron velocities by group /IGM/ {I04=2 or S02=2}

Terminate Marker

8. Mixture material numbers in mixing table /M01/ {M01>0}

Terminate Marker

9. Component material numbers of mixtures in mixing table /M01/
{M01>0}

Terminate Marker

10. Atom densities of component materials in mixing table /M01/
{M01>0}

Terminate Marker

11. Delta option radial zone numbers by fine mesh interval /IM/
{I04=4}

Terminate Marker

12. Delta option radial zone modifiers /IZ/ {I04=4}

Terminate Marker

13. Delta option axial zone numbers by fine mesh interval /JM/
{I04=4}

Terminate Marker

14. Delta option axial zone modifiers /JZ/ {I04=4}

Terminate Marker

15. Cross section table position for activities /NACT/ {NACT>0}

Terminate Marker

16. End of problem identifier (NCON)

0 - End of problem (only if no burnup)

>0 - Take time step of DELT and read burnup data for
N isotopes (F.)

<0 - Take time step of DELT (F.)

(DELT is in Burnup Data)

Terminate Marker

F. Burnup Data (NCON≠0)

F.1 Integer control parameters /11/

	Variable	
	<u>Name</u>	<u>Description</u>
1.	DELT	Time step (days)

*Terminate Marker**F.2 Integer control parameters /12/ (NCON>0)*

1.	MATN(NCON)	Material sequence number of burnable isotope
2.	NBR(NCON)	0 - No effect 1 - Fertile isotope 2 - Fissile isotope
3.	LD(NCON)	0 - No decay source N - Decay source from burnable isotope N

Variable		
	<u>Name</u>	<u>Description</u>
4.	LCN(NCON,1)	0 - No capture source N - Capture source from burnable isotope N
5.	LCN(NCON,2)	0 - No capture source N - Capture source from burnable isotope N
6.	LFN(NCON,1)	0 - No fission source N - Fission source from burnable isotope N
7.	LFN(NCON,2)	0 - No fission source N - Fission source from burnable isotope N
8.	LFN(NCON,3)	0 - No fission source N - Fission source from burnable isotope N
9.	LFN(NCON,4)	0 - No fission source N - Fission source from burnable isotope N
10.	LFN(NCON,5)	0 - No fission source N - Fission source from burnable isotope N

Variable		
	<u>Name</u>	<u>Description</u>
11.	LFN(NCON,6)	0 - No fission source N - Fission source from burnable isotope N
12.	LFN(NCON,7)	0 - No fission source N - Fission source from burnable isotope N

Terminate Marker

F.3 Integer control parameters /1/ {NCON>0}

1.	ALAM(ITEMP)	Decay constant for decay of burnable isotope N (days ⁻¹) (0.0 for no decay)
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Terminate Marker

3.2. Supplemental Input Information for 2DBTOR

3.2.1. **Cross Section Considerations.** Cross sections input into 2DBTOR are ordered for each group as shown in Table 2.

Table 2. Order of cross sections in 2DBTOR

<u>Cross Section</u>	<u>Group</u>	<u>Position Description</u>
σ activity 1 (optional)	g	
.	.	
.	.	
.	.	
σ activity N (optional)	g	(N = IHT - 4)
σ fission	g	
σ absorption	g	
$v\sigma$ fission	g	
σ transport	g	IHT
σ selfscatter	g \rightarrow g	IHS
σ downscatter	g - 1 \rightarrow g	
.	g	
.	g	
.	g	
σ downscatter	g - NXCM \rightarrow g	ITL

If activity cross sections are not present, then IHT = 4.

The absorption cross section is used only for editing purposes. If a removal cross section is to be calculated, then 2DBTOR computes this by subtracting the self scatter cross section from the transport cross section.

Material numbers in 2DBTOR start at 1 and go through MT (the user specified total number of materials). Materials entered by cards or tape start at 1 and run through MCR (the user-specified number of materials from cards or tape). Materials formed in the mixing table are numbered from MCR+1 through MT.

The cross section mixing table is controlled by three input arrays: I0, I1, and I2. The length of the cross section mixing table (M01) is specified by the user. It is important to remember to initialize each array to zero before performing a mix in the mixing table since 2DBTOT does not do this internally in the program.

For each row of entries in the mixing table, three operations are possible. First, all of the cross sections in a given material number may be multiplied by a constant. This option is useful in number density variations. Second, a set of cross sections may be multiplied by a constant and added into another material. This option is useful in mixing cross sections. Finally, all of the cross sections of a material may be multiplied by the eigenvalue. This option is useful in concentration searches (I04 = 3).

Table 3 illustrates the three options available from the mixing table:

Table 3. 2DBTOR mixing table options

Options	Material Number (I0)	Component Number (I1)	Concentration or Numeric Constant (I2)
1	M	0	X
1	M	0	X
1	M	0	X

3.2.2. Search Considerations. The 2DBTOR code computes implicit eigenvalue searches on time absorption, material composition, zone thickness, and material buckling. In contrast to a k_{eff} calculation, the fission spectrum is not multiplied by $1/\lambda$ after each outer iteration. Instead, after a converged λ has been obtained ($|\lambda^{V+1} - \lambda^V| < \epsilon$) by a sequence of outer iterations, the desired parameter is perturbed to make λ approach unity. That is, first a converged λ is calculated for the initial system. The system is then altered by the amount specified in the input (the eigenvalue modifier) and a second converged λ is calculated. Subsequent parameter changes are determined using either linear or parabolic interpolation procedures. The iteration is continued until $|\lambda - 1| < \epsilon$.

3.2.2.1. Time Absorption (α calculation). For simplicity, consider the one group, time dependent neutron diffusion equation

$$\frac{1}{\nu} \frac{\partial \phi(\vec{r}, t)}{\partial t} = DV^2 \phi(\vec{r}, t) - \Sigma_a \phi(\vec{r}, t) + \nu \Sigma_f \phi(\vec{r}, t) . \quad (14)$$

If one now assumes that

$$\phi(\vec{r}, t) = \phi(\vec{r}) e^{\alpha t} . \quad (15)$$

then equation (14) can be rewritten in the form

$$DV^2 \phi(\vec{r}) - (\Sigma_a + \frac{\alpha}{\nu} \phi(\vec{r})) + \nu \Sigma_f \phi(\vec{r}) = 0 . \quad (16)$$

In a time absorption calculation, the parameter α , as defined and used in equations (15) and (16), is computed as the eigenvalue. Note that α/ν is effectively an absorption cross section -- hence the name "time absorption."

3.2.2.2 Material Concentration (C calculation). 2DBTOR can perform an extremely flexible and comprehensive criticality search on material composition. Any number of materials can simultaneously be added, depleted, or interchanged in any number of zones.

The format for specifying concentration searches can best be described by a simple example. Suppose that a zone mixture, say Mix 10, is to be composed of Materials 2 and 4. Further, assume that Material 2, with an initial density of 0.02 (atoms/barn-cm), shall be varied to obtain criticality, and Material 4 shall have a fixed

density of 0.04 (atoms/barn-cm).

The I0, I1, and I2 arrays would then be set up as shown in the following tabulation.

<u>Mix Number (I0)</u>	<u>Material Number(I1)</u>	<u>Density(I2)</u>
10	0	0
10	2	0.02
10	10	0
10	4	0.04

The first row (10,0,0) instructs the code to clear a storage area for Mix 10. The second row (10,2,0.02) causes Material 2 to be added to Mix 10 with a density of 0.02. The third row (10,10,10) causes the current contents of Mix 10 to be multiplied by the eigenvalue. Finally, the last row (10,4,0.04) instructs the code to add Material 4 to Mix 10 with a density of 0.04.

All of the foregoing can be summarized by the expression

$$\Sigma_{10} = 0.02 \sigma_2 EV + 0.04\sigma_4 \quad (17)$$

where:

Σ_{10} = macroscopic cross section for Mix 10,

σ_2 = microscopic cross section for Material 2,

σ_4 = microscopic cross section for Material 4,

EV = the eigenvalue.

3.2.2.3 *Zone Dimensions (δ calculation).* 2DBTOR searches on reactor dimensions by varying the dimensions of each axial and radial mesh interval. Each mesh width, δr^i , is computed from the expression

$$\delta r^i = \delta r_0^i [1 + (\text{mesh modifier})^i EV] , \quad (18)$$

where δr^i is the initial mesh spacing and EV is the eigenvalue. Different mesh modifiers can be specified for each axial and radial mesh interval.

3.2.2.4 *Buckling (B^2 calculation).* In a buckling search, the quantity $D_i B^2$, where D_i is the zone dependent diffusion constant for group i , is added to the i th group absorption cross section. The in-group scattering cross section, σ_{gg}^i , is reduced by the same amount so that the calculated total cross section remains equal to the input total cross section. The buckling is then computed as the eigenvalue.

3.2.3. **Burnup Model.** The basic burnup equation for each zone has the form

$$\frac{dN^i}{dt} = -\lambda^i N^i - \sigma_a^i \phi N^i + \lambda^k N^k + \sum_{j=1}^{2} \sigma_c^{j-} \phi N^j + \sum_{m=1}^7 \sigma_f^{m-} \phi N^m \quad (19)$$

where:

- N^i = density of nuclide i ,
 λ^i = decay constant for nuclide i ,
 σ_a^i = spectrum averaged absorption cross section for nuclide i ,
 σ_f^i = spectrum averaged fission cross section for nuclide i ,
 σ_c^i = spectrum averaged capture cross section for nuclide i ,
 $\bar{\phi}$ = total average flux.

The last two terms in equation (19) allow provision for two capture and seven fission sources. The latter option, for example, could be used to compute the fission product buildup.

Each input time step is arbitrarily subdivided into 10 smaller time steps. Equation (19) is then solved as a march-out problem using the subdivided time intervals. If one rewrites equation (19) in the form

$$\frac{d\vec{N}}{dt} = \vec{f}(\vec{N}, t) \quad , \quad (20)$$

then the particular march-out algorithm used can be written as

$$\vec{N}_{j+1} = \vec{N}_j + \frac{\delta t}{2} (\vec{f}_j + \vec{f}_{j+1}) \quad , \quad (21)$$

where j is the index on time and δt is the fine-step time interval.

Observe that equation (21) is implicit in the sense that N_{j+1} must be known in

order to compute f_{j+1} . One must therefore iterate on N at each time point. This procedure leads to the algorithm

$$\vec{N}_{j+1}^{v+1} = \vec{N}_j + \frac{\delta t}{2}(\vec{f}_j + \vec{f}_{j+1}) , \quad (22)$$

where v is the iteration index.

3.2.3.1. Remarks on Burnup Equations. The zone averaged flux and cross sections appearing in equation (19) are computed before each time step. The total reactor power (from the burnable isotopes) and flux profile (relative zone fluxes) are held constant during the fine-step march-out described by equation (22).

It should be clear from the mathematical model presented that relatively short time steps should be employed if rapid variations in isotopic concentrations or flux profiles are anticipated.

3.2.4. Source Problems. 2DBTOR will compute the flux profiles resulting from an extraneous (in space and energy) source distribution. The following suggestions will assist the user in running source problems:

1. A source problem is meaningless (and will not converge) if $k > 1.0$.
2. Convergence can be accelerated by giving the code an estimate of k (Card 5, Word 1).
3. A good estimate of the initial total neutron production rate (Card 6, Word 6)

will enhance convergence. This value can be estimated using the simple expression

$$N = \frac{kS}{1 - k}, \quad (23)$$

where:

N = total neutron production rate from fission,

S = total neutron source rate from extraneous source,

k = multiplication constant.

3.2.5. Remarks on Code Operation.

1. Since the input data is inverted for transposed calculations, all group indicies in the output of adjoint cases are transposed. Furthermore, the balance tables in adjoint calculations do not have a direct physical interpretation.
2. The material inventory tables are inapplicable for a mixture specification more complex than a mix in a mix (e.g., a mix in a mix in a mix).
3. An isotope cannot be mentioned more than once in the same mix in burnup calculations. If mentioned more than once in other calculations, the printed inventory will be incorrect.
4. Although the new eigenvalue and material densities are computed and printed after the last time step, the zone averaged cross sections and reaction rates are not. These can easily be obtained, however, by simply taking 1 extra burnup step of

zero length. Similarly, the zone averaged cross sections and reaction rates can be obtained in non-burnup runs by simply calling for 1 (dummy) burnup step of zero length.

5. A flux dump is given only when:

- 1) A dump is called for, and
- 2) The burnup time is zero.

Thus, if a dump is called for in a burnup calculation, only one dump (the initial flux) is given.

6. Tight mesh spacing in the dimension perpendicular to line inversion can cause excessive running time. Thus, if tight mesh spacing is used, it should be along the dimension containing the most mesh intervals. For the same reason, the "dummy" dimension in one-dimensional problems should contain large mesh intervals.

3.3 2DBTOR Format Description

3.3.1 Operators. Generalized Format = OPnn xx

1. nn = first subfield with integer entry
2. OP = operator = second subfield with character entry
3. xx = third subfield with real or integer entry
4. There must not be a delimiter (blank or end of card) between the first and second subfields.

5. There must be a delimiter between the second and third subfields.
6. There must be a delimiter between the third subfield of one operator and the first subfield of the following operator.

	<u>Operator</u>	<u>Description</u>
1.	C	Continue the current array nn times with the previous xx entries
2.	F	Fill the remainder of the current array with xx
3.	I	Linear interpolation; generate nn entries between xx

	<u>Operator</u>	<u>Description</u>
		and the previous number
4.	R	The value xx is generated nn times
5.	T	Terminate the current array

3.4. Printed Output Description

The first output section of 2DBTOR is a brief edit of the first 37 input data preceded by the title card. The size of the array required to run the problem is printed after the above edit. If more or less than the required 37 entries is read than an error message is printed and the calculation stops.

The next section is an edit of the cross section , source, mesh interval, and zone data. First, a cross section edit is printed with a list of the cross sections read followed by a consistency check. If the problem involves a great amount of inelastic scattering,

then this check can be ignored. Second, if a source problem is run, then the source distribution is printed for each group. Each group source distribution is preceded by the required number of entries. Third, an edit of the mesh intervals is performed for both the radial and axial points. Again, each edit is preceded by the required number of entries. Finally, an edit of the zone numbers by mesh interval followed by an edit of the material numbers by zone is printed. Each of the above edits is preceded by the required number of entries.

The next section consists of an edit of the fission spectrum and mixing table data. The fission spectrum is printed for all groups. Next, the mixing table is printed for the I0, I1, and I2 arrays. Again, the required number of entries is printed.

The next section consists of a map of both the zone numbers by mesh interval and the material numbers by fine mesh interval. This provides a means to get a picture of the problem. If more than approximately 50 mesh intervals are used in the radial direction, then the printed inventory will leave off the excess due to problems with printing off the page. This will cause an error, which will not stop further running of the problem.

The next section is a brief edit of the time in days that the problem was started, of the mixing table in easy to read format, and of the cross sections ($NPRT \geq 2$) input into the problem for each group.

After printing out the input edits above, the number of inner iterations per outer iteration with the associate eigenvalue, eigenvalue slope, and lambda after the outer iteration are printed. This is followed by a balance table, which lists the number of fissions, in-scattering neutrons, out-scattering neutrons, absorptions, and leakage for each group. The total for each neutron process follows the above. The neutron

multiplication constant (not k_{eff}) is printed, followed by an edit of the radii, average radii, ax_{ii} , and average ax_{ii} dimensions.

Next, if $NPRT > 2$, then the flux by mesh interval for each group followed by the total flux by mesh interval is printed. This is followed by the power density (MWT/liter) for each mesh interval.

To end the problem (if no burnup is required), a brief edit of both the mass (kilograms) and volume (liters) inventory for each zone is printed followed by their totals. This concludes the problem.

If a burnup run is called for, then the amount of days in the burnup followed by another eigenvalue edit and balance table edit is printed. Next, a brief edit of the burnup input dat is printed in easy to read format. Finally, the absorption and fission rates for each material burned by zone is printed with the number density of each material called in the burnup.

4. PROGRAMMER'S GUIDE TO 2DBTOR

4.1. 2DBTOR Files Description

Logical

<u>Unit</u>	<u>Name</u>	<u>Format</u>	<u>Usage</u>
3	FOR3.DAT	Unformatted	Cross section storage
4	FOR4.DAT	Unformatted	Scratch storage
5	TORACT5.DAT	Formatted	Input

6	TORACT5.OUT	Formatted	Output
8	FOR8.DAT	Unformatted	Flux storage
9	FOR9.DAT	Unformatted	Source storage
10	FOR10.DAT	Unformatted	Scratch storage
11	FOR11.DAT	Unformatted	Scratch storage
12	FOR12.DAT	Unformatted	Scratch storage
14	FOR14.DAT	Unformatted	Input of a flux dump
15	FOR15.DAT	Unformatted	Scratch storage
16	FOR16.DAT	Unformatted	Output of a flux dump

4.2 2DBTOR Subroutines

<u>Name</u>	<u>Function</u>
CALC	Main program
INP	Controls reading and printing of all input dat, computes variable dimension pointers, and computes program constants
ERR02	Prints error messages
S860	Reads cross sections from cards, performs adjoint reversals if required, and writes cross section tape
S862	Reads input fluxes and prepares a flux tape
S864	Reads input source and prepares a source tape
REAG2	Reads floating point data

<u>Name</u>	<u>Function</u>
REAI2	Reads integer data
RREAG2	Reads toroidal data (major radius)
MAPR	Produces a picture by zone and by material
INIT	Performs adjoint reversals, mixes cross sections, modifies geometry, and calculates areas, volumes, and fission neutrons
CLEAR	Sets an array of a specified length to a constant
FISCAL	Calculates fission sums and performs outer iteration normalization
S8830	Prints time, eigenvalue, lambda, etc. after each outer iteration
OUTER	Performs a complete outer iteration
INNER	Calculates flux in a specified group
INNER1	Calculates coefficients for the flux equation
INNERT	Calculates coefficients for the flux equation in triangular geometry
INNER2	Calculates flux in a specified group
INNERP	Calculates flux in a specified group for periodic boundary conditions
IFLUXN	Normalizes flux before each group flux calculation
IFLUXL	Normalizes flux before each group flux calculation (used for toroidal geometry source problems)

<u>Name</u>	<u>Function</u>
CNNP	Performs convergence tests and computes a new eigenvalue for search calculations
S8850	Prints the monitor line, group fluxes, total flux, power density, and fission source rate
S8847	Computes and prints group totals
PRT	Prints any IM*JM array
GRAM	Calculates and print the zone volume and the mass of each material in each zone
INPB	Reads and prints the input burnup data
AVERAG	Calculates zone averaged fluxes, fission cross sections, absorption cross sections, and breeding ratio
MARCH	Calculates the time dependent isotopic concentrations

4.2.1. Subroutine Calling Sequences.

<u>Subroutine</u>	<u>Called By</u>	<u>Calls</u>
CALC	---	INP, INIT, FISCAL, S8830, ERR02, OUTER, CNNP, S8850, GRAM, INPB, AVERAG, MARCH
INP	CALC	S860,S862, S864, REAG2, REA12, RREAG2, MAPR, ERR02

<u>Subroutine</u>	<u>Called By</u>	<u>Calls</u>
ERR02	CALC, INP, REAI2, REAG2, INIT, CNNP	--
S860	INP	--
S862	INP	REAG2
S864	INP	REAG2
REAG2	INP, S862, S864	ERR02
REAI2	INP	ERR02
RREAG2	INP	--
MAPR	INP	--
INIT	CALC	CLEAR, ERR02
CLEAR	INIT, GRAM	--
FISCAL	CALC	--
S8830	CALC, S8850	--
OUTER	CALC	INNER1, INNER, INNERP
INNER	OUTER	IFLUXN, IFLUXL
INNER1	OUTER	--
INNERT	OUTER	--
INNER2	OUTER	IFLUXN, IFLUXL
INNERP	OUTER	IFLUXN
IFLUXN	INNER, INNER2, INNERP	--
IFLUXL	INNER, INNER2	--

<u>Subroutine</u>	<u>Called By</u>	<u>Calls</u>
CNNP	CALC	ERR02, CLEAR
S8850	CALC	PRT, S8830, S8847
S8847	S8850	--
PRT	S8850	--
GRAM	CALC	CLEAR
INPB	CALC	--
AVERAG	CALC	--
MARCH	CALC	--

4.3. Selected Definitions of Variables and Arrays Used

<u>Variable</u>	<u>Description</u>
** INTERNAL VARIABLES **	
NINP	Input tape
NOUT	Output tape
NCR1	Cross section tape
NFLUX1	Flux tape
NSCRAT	Scratch tape
NSORCE	Source tape
ALA	Lambda
B07	Used for internal computation in FISCAL and INIT
CNT	Convergence trigger for lambda
CVT	Convergence trigger
DAY	Burnup time in days

Variable	Description
DELT	Length of time step (days)
E0(IGP)	Fission rate
E1(IGP)	Fission source
E2(IGP)	In-scatter and extraneous source
E3(IGP)	Out-scatter
E4(IGP)	Absorptions
E5(IGP)	Left leakage
E6(IGP)	Right leakage
E7(IGP)	Top leakage
E8(IGP)	Bottom leakage
E9(IGP)	Total leakage
E01	Temporary
E02	Temporary
E03	Temporary
EQ	Temporary for S852(CNNP)
EVP	Previous eigenvalue
EVPP	Eigenvalue for two iterations back
FEF	Energy released /fission (215MeV)
GBAR	Group indicator
GLH	Initial clock time in seconds
IGEP	IGE + 1
IGP	IGM + 1
IGV	Group indicator for inner and outer

Variable	Description
IHS	Position of sigma self scatter
IHT	Position of sigma transport
II	Inner iteration count for a single group
IMJM	IM*JM
IP	IM + 1
ITEMP	Temporary
ITEMP1	Temporary
ITEMP2	Temporary
ITL	Cross section table length
IZP	IZM + 1
JP	JM + 1
K07	Not used
KPAGE	Page counter for monitor print
LAP	Lambda for previous eigenvalue
LAPP	Lambda for two iterations back
LAR	Lambda for previous iteration
LC	Loop count (total II in a single outer iteration)
ML	MCR + MTP
NCON	-/0/+ = take time step of DELT/ end of problem/ take time step of DELT and read burnup data
NGOTO	Temporary
ORFP	ORF for $(1-\lambda) < 10 \cdot \text{EPS}$
P02	Outer iteration count

<u>Variable</u>	<u>Description</u>
PBAR	Temporary
SBAR	Temporary
SK7	Sum of K7 over all groups
T06	0/1 = no effect/delta calculation
T7	Alpha/velocity
T11	Previous fission total
TEMP	Temporary
TEMP1	Temporary
TEMP2	Temporary
TEMP3	Temporary
TEMP4	Temporary
TI	Time
TSD	(MW-sec)/(fissions)
V11	Total source for the group
INPUT VARIABLES	
ID(20)	ID card
A02	0/1 = flux calc./adjoint calc.
I04	Eigenvalue type (0/1/2/3/4 = source/keff/alpha/ concentration/delta/buckling)
S02	Parametric eigenvalue type (0/1/2 = none/keff/alpha)
IGM	Number of groups
NXCM	Number of downscattering terms
MCR	Number of materials from cards

<u>Variable</u>	<u>Description</u>
MTP	Number of materials from tape
G07	Inner iteration max per group (if neg., alt dir)
D05	Max number of outer iterations
MAXT	Max run time (minutes)
NPRT	Print option (0/1/2/3 = mini/midi/Xsec/fluxes)
M07	Flux guess (0/1 = none/flux from FOR14.DAT)
NPUN	Flux dump (0/1 = none/ flux dump to FOR16.DAT)
IGE	Geometry parameter (0/1/2/3/4 = X-Y/R-Z/R- θ /triangular/ toroidal)
IM	Number of radial intervals
JM	Number of axial intervals
IZM	Number of zones
MT	Total number of materials including mixes
M01	Number of mixture specifications
B01	Left B. C. (0/1 = vacuum/reflective)
B02	Right B. C. (0/1 = vacuum/reflective)
B03	Top B. C. (0/1/2 = vacuum/reflective/periodic)
B04	Bottom B. C. (0/1/2 = vacuum/reflective/periodic)
IZ	Radial zones (delta option only)
JZ	Axial zones (delta option only)
NACT	Number of activations
EV	Eigenvalue
EVM	Eigenvalue modifier

Variable	Description
S03	Parametric eigenvalue
BUCK	Buckling
LAL	Lambda lower
LAH	Lambda upper
EPS	Eigenvalue convergence criterion
EPSA	Pointwise convergence criterion
G06	Inner iteration test (if 0 no test)
POD	Parameter oscillation damper
ORF	Over-relaxation factor
S01	-/+ = power(MWT)/neutron source rate
ARRAY VARIABLES	
ATW(ML)	Material atomic weight
HOLN(ML,2)	Material name
ALAM(ML)	Decay constant (days ⁻¹)
C0(ITL,MT)	Cross section array for current group
N0(IM,JM)	Total flux (old)
N2(IM,JM)	Total flux (new)
A0(IP)	Radial area element
A1(IM)	Axial area element
F0(IM,JM)	Fissions (old)
F2(IM,JM)	Fissions (new)
I0(M01)	Mix number
I1(M01)	Material number for mix

<u>Variable</u>	<u>Description</u>
I2(M01)	Material density
I3(M01)	Material densities for gram calc.
K6(IGM)	Fission spectrum (effective)
K7(IGM)	Fission spectrum (input)
M0(IM,JM)	Zone numbers
M2(IM,JM)	Material numbers by zone
R0(IP)	Initial radii
R1(IP)	Current radii
R2(IM)	Radial zone numbers (delta calc. only)
R3(IZ)	Radial zone modifiers (delta calc. only)
R4(IM)	Average radii
R5(IM)	Delta-R
S2(IM,JM)	Fixed source
V0(IM,JM)	Volume elements
V7(IGM)	Neutron velocities
Z0(JP)	Initial axii
Z1(JP)	Current axii
Z2(JM)	Axial zone numbers (delta calc. only)
Z3(JZ)	Axial zone modifiers (delta calc. only)
Z4(JM)	Average axii
Z5(JM)	Delta-Z
CXS(IM,JM,3)	Constants involving cross sections for flux calculation
VOL(IZM)	Zone volume (liters)

<u>Variable</u>	<u>Description</u>
MASS(ML,IZM)	Material inventory in each zone
MATN(ML)	Material number for burnable isotopes
NBR(ML)	0/1/2 = none/fertile/fissile
LD(ML)	Source isotope for decay
LCN(ML,2)	Source isotopes for capture
LFN(ML,7)	Source isotope for fission
PHIB(IZM)	Zone averaged flux
AXS(ML,IZM)	Spectrum averaged absorption cross section
FXS(ML,IZM)	Spectrum averaged fission cross section
MASSP(ML,IZM)	Material inventory in each zone (previous)
CXR(JM)	Constants for right boundary
CXT(IM)	Constants for top boundary
HA(IM OR JM)	Temp storage for line inversion
PA(IM OR JM)	Temp storage for line inversion

5. REFERENCES

1. W. W. Little, Jr., and R. W. Hardie. 2DB User's Manual, BNWL - 831
Revision 1, Batelle Pacific Northwest Laboratory, Richland, Washington
(Unpublished).

APPENDIX H
2DBTOR LISTING

TLIST2.FOR

Wednesday, February 21, 1990 12:42 pm

Page 1

C	PROGRAM ZDB		
C	2	TAPE5=INPUT,TAPE6=OUTPUT,TAPE3=MCR1,TAPE4=MSCRAT,	CALC <
C	3	TAPE8=MFLUX1,TAPE9=MSORCE,	CALC <
C	4	TAPE14=MFLX1,TAPE16=MFLXD,TAPE15=MXS)	CALC <
C	*****	DESCRIPTION OF SUBROUTINES *****	CALC 9
C			CALC 10
C	CALC	MAIN PROGRAM--SETS UP TAPE UNITS AND CALLS INP, INIT,	CALC 11
C		FISCAL, SB830, ERRO2, OUTER, CNMP, SB850, GRAM,	CALC 12
C		INPB, AVERAG, AND MARCH.	CALC 13
C			CALC 14
C	INP	SUBROUTINE TO CONTROL THE READING AND PRINTING OF ALL	CALC 15
C		INPUT DATA, COMPUTE VARIABLE DIMENSION POINTERS AND	CALC 16
C		PROGRAM CONSTANTS. INP IS CALLED BY CALC AND CALLS	CALC 1<
C		SB60, SB62, SB64, REAG2, REA12, MAPR, AND ERRO2.	CALC 18
C			CALC 19
C	ERRO2	ERRO2 IS USED TO PRINT AN ERROR MESSAGE. IT IS CALLED	CALC 20
C		BY CALC, INP, REA12, REAG2, INIT, AND CNMP.	CALC 21
C			CALC 22
C	SB60	SUBROUTINE TO READ CROSS SECTIONS FROM CARDS, PERFORM	CALC 23
C		ADJOINT REVERSALS IF REQUIRED, AND WRITE CROSS SECTION	CALC 24
C		TAPE. SB60 IS CALLED BY INP.	CALC 25
C			CALC 26
C	SB62	SB62 READS INPUT FLUXES AND PREPARES A FLUX TAPE. IT IS	CALC 27
C		CALLED BY INP AND CALLS REAG2.	CALC 28
C			CALC 29
C	SB64	SB64 READS INPUT SOURCE AND PREPARES A SOURCE TAPE. IT	CALC 30
C		IS CALLED BY INP AND CALLS REAG2.	CALC 31
C			CALC 32
C	REAG2	SUBROUTINE TO READ FLOATING POINT DATA. REAG2 IS CALLED	CALC 33
C		BY INP, SB62, AND SB64. REAG2 CALLS ERROR.	CALC 34
C			CALC 35
C	REA12	SUBROUTINE TO READ INTEGER DATA. REA12 IS CALLED BY INP	CALC 36
C		AND CALLS ERRO2.	CALC 37
C			CALC 38
C	MAPR	SUBROUTINE TO PRODUCE A PICTURE BY ZONE AND	CALC 39
C		MATERIAL. MAPR IS CALLED BY INP.	CALC 40
C			CALC 41
C	INIT	INIT PERFORMS ADJOINT REVERSALS(SB06), MIXES CROSS	CALC 42

TLIST2.FOR

Wednesday, February 21, 1990 12:42 pm

Page 2

C	SECTIONS(SB07), MODIFIES GEOMETRY(SB10), AND CALCULATES	CALC 43
C	AREAS AND VOLUMES(SB11), AND FISSION NEUTRONS(SB21).	CALC 44
C	INIT IS CALLED BY CALC AND CALLS CLEAR AND ERRO2.	CALC 45
C		CALC 46
C	CLEAR SETS AN ARRAY OF A SPECIFIED LENGTH TO A GIVEN	CALC 47
C	CONSTANT. THE SUBROUTINE IS CALLED BY INIT AND GRAM.	CALC 48
C		CALC 49
C	FISCAL CALCULATES FISSION SUMS(SB22) AND PERFORMS	CALC 50
C	NORMALIZATION(SB23). FISCAL IS CALLED BY CALC.	CALC 51
C		CALC 52
C	SB830 SB830 IS THE MONITOR PRINT SUBROUTINE--PRINTS TIME,	CALC 53
C	EIGENVALUE, LAMBDA, ETC. AFTER EACH OUTER ITERATION.	CALC 54
C	IT IS CALLED BY CALC AND SB850.	CALC 55
C		CALC 56
C	OUTER PERFORMS A COMPLETE OUTER ITERATION. CALLS INNER1,	CALC 57
C	INNER, AND INNERP. OUTER IS CALLED BY CALC.	CALC 58
C		CALC 59
C	INNER CALCULATES THE FLUX IN SPECIFIED GROUP. IT IS CALLED	CALC 60
C	BY OUTER AND CALLS IFLUXN.	CALC 61
C		CALC 62
C	INNER1 CALCULATES COEFFICIENTS FOR THE FLUX EQUATION. INNER1	CALC 63
C	IS CALLED BY OUTER.	CALC 64
C		CALC 65
C	INNER2 INNET CALCULATES COEFFICIENTS FOR THE FLUX EQUATION FOR	CALC 66
C	TRIANGULAR GEOMETRY. INNER2 IS CALLED BY OUTER.	CALC 67
C		CALC 68
C	INNER2 CALCULATES THE FLUX IN SPECIFIED GROUP. IT IS CALLED	CALC 69
C	BY OUTER AND CALLS IFLUXN.	CALC 70
C		CALC 71
C	INNERP CALCULATES THE FLUX IN SPECIFIED GROUP FOR PERIODIC B. C.	CALC 72
C	IT IS CALLED BY OUTER AND CALLS IFLUXN.	CALC 73
C		CALC 74
C	IFLUXN SUBROUTINE TO NORMALIZE THE FLUXES BEFORE EACH	CALC 75
C	GROUP FLUX CALCULATION. IT IS CALLED BY INNER, INNER2,	CALC 76
C	AND INNERP.	CALC 77
C		CALC 78
C	CNMP PERFORMS CONVERGENCE TESTS(SB51) AND COMPUTES A NEW	CALC 79
C	EIGENVALUE FOR SEARCH OPTIONS(SB52). CNMP IS CALLED	CALC 80

TLIST2.FOR

Wednesday, February 21, 1990 12:42 pm

Page 3

C		BY CALC AND CALLS ERRO2 AND CLEAR.	CALC 81
C			CALC 82
C	S8850	FINAL PRINT SUBROUTINE--PRINTS THE MONITOR LINE,	CALC 83
C		GROUP FLUXES, TOTAL FLUX, POWER DENSITY, AND FISSION	CALC 84
C		SOURCE RATE. IT IS CALLED BY CALC AND CALLS PRT, S8830,	CALC 85
C		AND S8847.	CALC 86
C			CALC 87
C	S8847	SUBROUTINE TO COMPUTE AND PRINT GROUP TOTALS. S8847 IS	CALC 88
C		CALLED BY S8850.	CALC 89
C			CALC 90
C	PRT	SUBROUTINE TO PRINT ANY IM*JM ARRAY. IT IS CALLED BY	CALC 91
C		S8850.	CALC 92
C			CALC 93
C	GRAM	CALCULATES AND PRINTS THE MASS OF EACH MATERIAL IN EACH	CALC 94
C		ZONE AND THE ZONE VOLUME. IT IS CALLED BY CALC AND	CALC 95
C		CALLS CLEAR.	CALC 96
C			CALC 97
C	INPB	SUBROUTINE TO READ AND PRINT THE INPUT BURNUP DATA. IT	CALC 98
C		IS CALLED BY CALC.	CALC 99
C			CALC 100
C	AVERAG	AVERAG CALCULATES ZONE AVERAGED FLUXES, FISSION CROSS	CALC 101
C		SECTIONS, ABSORPTION CROSS SECTIONS, AND BREEDING RATIO.	CALC 102
C		THE SUBROUTINE IS CALLED BY CALC.	CALC 103
C			CALC 104
C	MARCH	SUBROUTINE TO CALCULATE THE TIME DEPENDENT ISOTOPIC	CALC 105
C		CONCENTRATIONS. MARCH IS CALLED BY CALC.	CALC 106
C			CALC 107
C		***** INTERNAL VARIABLES *****	CALC 108
C			CALC 109
C	NIMP	INPUT TAPE	CALC 110
C	NOUT	OUTPUT TAPE	CALC 111
C	NCR1	CROSS SECTION TAPE	CALC 112
C	NFLUX1	FLUX TAPE	CALC 113
C	NSCRAT	SCRATCH TAPE	CALC 114
C	NCR1	SOURCE TAPE	CALC 115
C	ALA	LAMBDA	CALC 116
C	B07	USED FOR INTERNAL COMPUTATION IN FISCAL AND INIT	CALC 117
C	CNT	CONVERGENCE TRIGGER FOR LAMBDA	CALC 118

TL1S12.FOR

Wednesday, February 21, 1990 12:42 pm

Page 4

C	CVT	CONVERGENCE TRIGGER	CALC 119
C	DAY	BURNUP TIME IN DAYS	CALC 120
C	DELT	LENGTH OF TIME STEP (DAYS)	CALC 121
C	E0(1GP)	FISSION RATE	CALC 122
C	E1(1GP)	FISSION SOURCE	CALC 123
C	E2(1GP)	IN-SCATTER (AND EXTRANEOUS SOURCE)	CALC 124
C	E3(1GP)	OUT-SCATTER	CALC 125
C	E4(1GP)	ABSORPTIONS	CALC 126
C	E5(1GP)	LEFT LEAKAGE	CALC 127
C	E6(1GP)	RIGHT LEAKAGE	CALC 128
C	E7(1GP)	TOP LEAKAGE	CALC 129
C	E8(1GP)	BOTTOM LEAKAGE	CALC 130
C	E9(1GP)	TOTAL LEAKAGE	CALC 131
C	E01	TEMPORARY	CALC 132
C	E02	TEMPORARY	CALC 133
C	E03	TEMPORARY	CALC 134
C	E0	TEMPORARY FOR S852 (CNMP)	CALC 135
C	EQ	PREVIOUS EIGENVALUE	CALC 136
C	EVP	EIGENVALUE FOR TWO ITERATIONS BACK	CALC 137
C	FEP	ENERGY RELEASED PER FISSIION (=215 MEV)	CALC 138
C	GBAR	GROUP INDICATOR FOR TAPE MOTION IN S824 (OUTER)	CALC 139
C	GLN	INITIAL CLOCK TIME IN SECONDS (INTEGER)	
C	IGEP	IGE + 1	CALC 141
C	IGV	IGH + 1	CALC 142
C	IGV	GROUP INDICATOR FOR INNER AND OUTER	CALC 143
C	IHS	POSITION OF SIGMA SELF SCATTER	CALC 144
C	IHT	POSITION OF SIGMA TRANSPORT	CALC 145
C	II	INNER ITERATION COUNT FOR A SINGLE GROUP	CALC 146
C	IMJM	IM*JM	CALC 147
C	IP	IN + 1	CALC 148
C	ITEMP	TEMPORARY	CALC 149
C	ITEMP1	TEMPORARY	CALC 150
C	ITEMP2	TEMPORARY	CALC 151
C	ITL	CROSS SECTION TABLE LENGTH	CALC 152
C	I2P	I2M + 1	CALC 153
C	JP	JM + 1	CALC 154
C	K07	NOT USED	CALC 155
C	KPAGE	PAGE COUNTER FOR MONITOR PRINT	CALC 156

TLIST2.FOR

Wednesday, February 21, 1990 12:42 pm

Page 5

C	LAP	LAMBDA FOR PREVIOUS EIGENVALUE	CALC 157
C	LAPP	LAMBDA FOR TWO ITERATIONS BACK	CALC 158
C	LAR	LAMBDA FOR PREVIOUS ITERATION	CALC 159
C	LC	LOOP COUNT (TOTAL II IN A SINGLE OUTER ITERATION)	CALC 160
C	ML	MCR + MTP	CALC 161
C	NCON	NEG/ZERO/POS=TAKE TIME STEP OF DELT/END OF PROBLEM/	CALC 162
C		TAKE TIME STEP OF DELT AND READ BURNUP DATA	CALC 163
C	NGOTO	TEMPORARY	CALC 164
C	NPRT	0/1/2/3=MINI/MIDI/CROSS SECTION/FLUX PRINT	
C	ORFP	ORF FOR 1 - LAMBDA LESS THAN 10*EPS	CALC 166
C	PO2	OUTER ITERATION COUNT	CALC 167
C	PBAR	TEMPORARY	CALC 168
C	SBAR	TEMPORARY	CALC 169
C	SK7	SUM OF K7 OVER ALL GROUPS	CALC 170
C	TO6	0/1=NOT DELTA/DELTA CALCULATION	CALC 171
C	T7	ALPHA/VELOCITY	CALC 172
C	T11	PREVIOUS FISSION TOTAL	CALC 173
C	TEMP	TEMPORARY	CALC 174
C	TEMP1	TEMPORARY	CALC 175
C	TEMP2	TEMPORARY	CALC 176
C	TEMP3	TEMPORARY	CALC 177
C	TEMP4	TEMPORARY	CALC 178
C	TI	TIME	CALC 179
C	TSD	(MW-SEC)/(FISSIONS)	CALC 180
C	V11	TOTAL SOURCE FOR THE GROUP	CALC 181
C			CALC 182
C	*****	INPUT VARIABLES (CARDS 1-5) *****	CALC 183
C			CALC 184
C	ID(20)	IDENTIFICATION CARD	CALC 185
C	A02	0/1=FLUX CALCULATION/ADJOINT CALCULATION	CALC 187
C	I04	EIGENVALUE TYPE (1/2/3/4/5=KEFF/ALPHA/CONCENTRATION/CALC	CALC 188
C		DELTA/BUCKLING)	CALC 189
C	S02	PARAMETRIC EIGENVALUE TYPE (0/1/2=NONE/KEFF/ALPHA)	CALC 190
C	IGH	NUMBER OF GROUPS	CALC 191
C	NXCM	NUMBER OF DOWNSCATTERING TERMS	CALC 192
C	MCR	NUMBER OF MATERIALS FROM CARDS/TAPE (+N/-N)	
C	MTP	NUMBER OF MATERIALS FROM TAPE	CALC 194
C	G07	INNER ITERATION MAX PER GROUP (IF NEG, ALT DIR)	

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 6

C	S04	INVERSION DIRECTION (0/1=NO EFFECT/ALTERNATE DIRECTI	CALC 199
C	D05	MAXIMUM NUMBER OF OUTER ITERATIONS	
C	MAXT	MAXIMUM TIME (MINUTES)	
C	NPRT	PRINT OPTION (0/1/2=MINI/MIDI/MAXI)	
C	M07	FLUX GUESS (0/1=NONE/INPUT FROM TAPE 14)	
C	NPUN	FLUX DUMP (0/1=NONE/DUMP TO TAPE 16)	
C	IGE	GEOMETRY (0/1/2/3=X-Y/R-Z/R-THETA/TRIANGULAR)	CALC 202
C	IM	NUMBER OF RADIAL INTERVALS	CALC 203
C	JM	NUMBER OF AXIAL INTERVALS	CALC 204
C	I2M	NUMBER OF MATERIAL ZONES	CALC 205
C	MT	TOTAL NUMBER OF MATERIALS INCLUDING MIXES	CALC 206
C	M01	NUMBER OF MIXTURE SPECIFICATIONS	CALC 207
C	B01	LEFT BOUNDARY CONDITION (0/1=VACUUM/REFLECTIVE)	CALC 208
C	B02	RIGHT BOUNDARY CONDITION (0/1=VACUUM/REFLECTIVE)	CALC 209
C	B03	TOP BOUNDARY CONDITION (0/1/2=VAC/REFL/PERIODIC)	CALC 210
C	B04	BOTTOM BOUNDARY CONDITION (0/1/2=VAC/REFL/PERIODIC)	CALC 211
C	I2	RADIAL ZONES (DELTA-OPTION ONLY)	CALC 212
C	J2	AXIAL ZONES (DELTA-OPTION ONLY)	CALC 213
C	NACT	NUMBER OF ACTIVATIONS	CALC 214
C	EV	FIRST EIGENVALUE GUESS	CALC 220
C	EVM	EIGENVALUE MODIFIER	CALC 221
C	S03	PARAMETRIC EIGENVALUE	CALC 222
C	BUCK	BUCKLING	CALC 223
C	LAL	LAMBDA LOWER	CALC 224
C	LAH	LAMBDA UPPER	CALC 225
C	EPS	EIGENVALUE CONVERGENCE CRITERIA	CALC 226
C	EPSA	POINTWISE CONVERGENCE CRITERIA	CALC 227
C	G06	INNER ITERATION TEST (IF ZERO, NO TEST)	CALC 228
C	POD	PARAMETER OSCILLATION DAMPER	CALC 229
C	ORF	OVER-RELAXATION FACTOR	CALC 230
C	S01	NEG/POS=POWER (MWT)/NEUTRON SOURCE RATE	CALC 231
C			CALC 232
C	*****	SUBSCRIPTED VARIABLES *****	CALC 233
C			CALC 234
C	ATW(ML)	MATERIAL ATOMIC WEIGHT	CALC 235
C	HOLM(ML,2)	MATERIAL NAME	CALC 236
C	ALAM(ML)	DECAY CONSTANT (DAYS-1)	CALC 237
C	CO(ITL,MT)	CROSS SECTION ARRAY FOR CURRENT GROUP	CALC 238

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 7

C	N0(IM,JM)	TOTAL FLUX (OLD)	CALC 239
C	N2(IM,JM)	TOTAL FLUX (NEW)	CALC 240
C	A0(IP)	RADIAL AREA ELEMENT	CALC 241
C	A1(IM)	AXIAL AREA ELEMENT	CALC 242
C	F0(IM,JM)	FISSIONS (OLD)	CALC 243
C	F2(IM,JM)	FISSIONS (NEW)	CALC 244
C	I0(M01)	MIX NUMBER	CALC 245
C	I1(M01)	MATERIAL NUMBER FOR MIX	CALC 246
C	I2(M01)	MATERIAL DENSITY	CALC 247
C	I3(M01)	MATERIAL DENSITIES FOR GRAM CALCULATION	CALC 248
C	K6(IGM)	FISSION SPECTRUM (EFFECTIVE)	CALC 249
C	K7(IGM)	FISSION SPECTRUM (INPUT)	CALC 250
C	M0(IM,JM)	ZONE NUMBERS	CALC 251
C	M2(I2M)	MATERIAL NUMBERS BY ZONE	CALC 252
C	R0(IP)	INITIAL RADII	CALC 253
C	R1(IP)	CURRENT RADII	CALC 254
C	R2(IM)	RADIAL ZONE NUMBERS (DELTA CALCULATION ONLY)	CALC 255
C	R3(I2)	RADIAL ZONE MODIFIERS (DELTA CALCULATION ONLY)	CALC 256
C	R4(IM)	AVERAGE RADII	CALC 257
C	R5(IM)	DELTA-R	CALC 258
C	S2(IM,JM)	FIXED SOURCE	CALC 259
C	V0(IM,JM)	VOLUME ELEMENTS	CALC 260
C	V7(IGM)	NEUTRON VELOCITIES	CALC 261
C	Z0(JP)	INITIAL AXII	CALC 262
C	Z1(JP)	CURRENT AXII	CALC 263
C	Z2(JM)	AXIAL ZONE NUMBERS (DELTA CALCULATION ONLY)	CALC 264
C	Z3(J2)	AXIAL ZONE MODIFIERS (DELTA CALCULATION ONLY)	CALC 265
C	Z4(JM)	AVERAGE AXII	CALC 266
C	Z5(JM)	DELTA-Z	CALC 267
C	CXS(IM,JM,3)	CONSTANTS INVOLVING CROSS SECTIONS FOR FLUX CALC.	CALC 268
C	VOL(I2M)	ZONE VOLUME (LITERS)	CALC 269
C	MASS(ML,I2M)	MATERIAL INVENTORY IN EACH ZONE	CALC 270
C	MATN(ML)	MATERIAL NUMBER FOR BURNABLE ISOTOPES	CALC 271
C	NBR(ML)	0/1/2=NO EFFECT/FERTILE/FISSILE ISOTOPE	CALC 272
C	LD(ML)	SOURCE ISOTOPE FOR DECAY	CALC 273
C	LCM(ML,2)	SOURCE ISOTOPES FOR CAPTURE	CALC 274
C	LFN(ML,7)	SOURCE ISOTOPES FOR FISSION	CALC 275
C	PHIB(I2M)	ZONE AVERAGED FLUX	CALC 276

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 8

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C   AXS(ML, IZM)  SPECTRUM AVERAGED ABSORPTION CROSS SECTION      CALC 277
C   FXS(ML, IZM)  SPECTRUM AVERAGED FISSION CROSS SECTION        CALC 278
C   MASSP(ML, IZM) MATERIAL INVENTORY IN EACH ZONE (PREVIOUS)    CALC 279
C   CXR(JM)       CONSTANTS FOR RIGHT BOUNDARY                   CALC 280
C   CXT(IM)       CONSTANTS FOR TOP BOUNDARY                      CALC 281
C   HA(IM OR JM)  TEMP STORAGE FOR LINE INVERSION                CALC 282
C   PA(IM OR JM)  TEMP STORAGE FOR LINE INVERSION                CALC 283
C
C   INCLUDE 'ABC.FOR'

COMMON/PAKED/A(50000)
OPEN(UNIT=3, STATUS='SCRATCH', FORM='UNFORMATTED')
OPEN(UNIT=4, STATUS='SCRATCH', FORM='UNFORMATTED')
C   USE BELOW ON A VAX                                             HRA2
C   OPEN(UNIT=3, FILE='FOR3.DAT', STATUS='SCRATCH', FORM='UNFORMATTED') HRA2
C   OPEN(UNIT=4, FILE='FOR4.DAT', STATUS='SCRATCH', FORM='UNFORMATTED') HRA2
C   OPEN(UNIT=5, FILE='toract5.DAT', STATUS='OLD', FORM='FORMATTED')
C   OPEN(UNIT=6, FILE='toract5.DAT', STATUS='UNKNOWN', FORM='FORMATTED')
C   OPEN(UNIT=8, STATUS='SCRATCH', FORM='UNFORMATTED')
C   OPEN(UNIT=9, STATUS='SCRATCH', FORM='UNFORMATTED')
C   USE BELOW ON A VAX                                             HRA2
C   OPEN(UNIT=8, FILE='FOR8.DAT', STATUS='SCRATCH', FORM='UNFORMATTED') HRA2
C   OPEN(UNIT=9, FILE='FOR9.DAT', STATUS='SCRATCH', FORM='UNFORMATTED') HRA2
C   OPEN(UNIT=10, FILE='FOR10.DAT', STATUS='UNKNOWN', FORM='UNFORMATTED')
C   OPEN(UNIT=11, FILE='FOR11.DAT', STATUS='UNKNOWN', FORM='UNFORMATTED')
C   OPEN(UNIT=12, FILE='FOR12.DAT', STATUS='UNKNOWN', FORM='UNFORMATTED')
C   OPEN(UNIT=14, FILE='FOR14.DAT', STATUS='UNKNOWN', FORM='UNFORMATTED')
C   OPEN(UNIT=15, FILE='FOR15.DAT', STATUS='UNKNOWN', FORM='UNFORMATTED')
C   OPEN(UNIT=16, FILE='FOR16.DAT', STATUS='UNKNOWN', FORM='UNFORMATTED')

1   CONTINUE
    REWIND 3                                                         CALC 291
    REWIND 4                                                         CALC 292
    REWIND 8                                                         CALC 293
    REWIND 9                                                         CALC 294
    CALL IMP(BIGR)                                                  HRA2 295
102 CALL INIT(A(LK6), A(LK7), A(LI0), A(LI1), A(LI2), A(LM0), A(LM2), CALC 296
    1   A(LN0), A(LR0), A(LR1), A(LR2), A(LR3), A(LR4), A(LR5), CALC 297

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TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 9

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2      A(LZ0), A(LZ1), A(LZ2), A(LZ3), A(LZ4), A(LZ5), A(LA0), CALC 298
3      A(LA1), A(LF0), A(LC0), A(LV0), ITL, IM, A(LV7), CALC 299
4      JM, MT, ML, A(LGAM), A(LHOLN))
CALL FISCAL(A(LM0),A(LF0),A(LV0),A(LC0),A(LK6),          CALC 301
2      A(LM0),A(LM2),ITL,MT)                             CALC 302
C      CALL MONITOR PRINT                                 CALC 303
101     CALL S8850                                         CALC 304
GO TO (100, 106, 106, 107), NGOTO
106     CALL ERRO2('MONPR',106,1)                         HVX
C      PERFORM AN OUTER ITERATION                       CALC 307
107     CALL OUTER(A(LA0), A(LA1), A(LC0), A(LF0), A(LK6), CALC 308
1      A(LM0), A(LM2), A(LM0), A(LM2),                 CALC 309
2      A(LS2), A(LV0), A(LV7), A(LZ5),                 CALC 310
3      A(LF2), ITL, MT, A(LCX5), IM, JM, A(LR5), A(LR4), CALC 311
4      A(LZ4), A(LCXR), A(LCXT), A(LHA), A(LPA),A(LR1),A(LZ1)) HRA2
C      PERFORM FISSION CALCULATION                     CALC 313
CALL FISCAL(A(LM0),A(LF0),A(LV0),A(LC0),A(LK6),          CALC 314
2      A(LM0),A(LM2),ITL,MT)                             CALC 315
C      PERFORM CONVERGENCE AND NEW PARAMETER CALCULATIONS CALC 316
CALL CNWP(A(LF2), A(LK6))                                CALC 317
GO TO (100, 101, 102), NGOTO
C      100/101/102=FINAL PRINT/MONITOR PRINT/SEARCH CALCULATION
100     CALL S8850(A(LF2),A(LM2),A(LR1),A(LZ1),A(LR4),A(LZ4),A(LV7), CALC 320
1      IM,JM,A(LN2),A(LC0),A(LM0),A(LM0),A(LM2),A(LF0),ITL,MT)
CALL GRAM(A(LMASS), A(LVOL), A(LATW), A(LHOLN), IM, JM,   CALC 323
1      A(LM0), A(LM2), A(LV0), A(LI0), A(LI1), A(LI2), ML,   CALC 324
2      A(LI3))                                                CALC 325
CALL INPB(A(LMATN),A(LNBR),A(LLD),A(LLCN),A(LLFN),A(LALAM), CALC 326
1      A(LHOLN),ML,A(LI2))                                     CALC 327
IF (NCON) 170,1,170
170     CALL AVERAG(A(LPHIB),A(LAXS),A(LFXS),A(LMATN),A(LMASS),A(LATW), WLP
1      A(LVOL),A(LC0),A(LN2),A(LM0),A(LV0),A(LHOLN),ML,ITL,   CALC 330
C      2      A(LNBR),A(LT6),A(LT8),A(LNTWON))                HRA2
2      A(LNBR),A(LACT),NACT,A(LACPOS))                       HRA3
IF (DELT) 180,1,180
180     CONTINUE                                             WLP
CALL MARCH(A(LPHIB),A(LMATN),A(LFXS),A(LAXS),A(LVOL),A(LMASS), WLP
1      A(LMASSP),A(LALAM),A(LLD),A(LLCN),A(LLFN),ML,          CALC 334

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 10

2	A(L10),A(L11),A(L12),A(LM2))	CALC 335
	GO TO 102	CALC 336
	END	CALC 337
	SUBROUTINE OUTER (A0, A1, C0, F0, K6, M0, M2, N0, N2,	OUTE 2
1	S2, V0, V7, Z5, F2, JTL, JMT, CXS,	OUTE 3
2	JIM, JJM, R5, R4, Z4, CKR, CXT, HA, PA, R1, Z1)	HRA2
	DIMENSION XR(50),XD(50)	HRA2
	DIMENSION AO(1), A1(1), F0(1), K6(1), M0(1), M2(1),	OUTE 5
1	N0(1), N2(1), R1(1), Z1(1), S2(1),	HRA2 6
2	V0(1), V7(1), Z5(1), F2(1), C0(JTL, JMT), HA(1), PA(1),	OUTE 7
3	CXS(JIM, JJM, 3), R5(1), R4(1), Z4(1), CKR(1), CXT(1)	OUTE 8
	INTEGER GBAR, PBAR, SBAR	OUTE 9
	INCLUDE 'ABC.FOR'	
	IGV=1	OUTE 13
C	SOURCE CALCULATION	OUTE 14
	REWIND 12	HRA2
10	CONTINUE	OUTE 15
	READ(NCR1) ((CO(I,M), I=1, ITL), M=1, MT)	OUTE 16
C	READ(11) (XR(M), M=1, MT)	HRA2
	READ(11) (XD(M), M=1, MT)	HRA2
	DO 1110 M=1, MT	HRA2
1110	CONTINUE	HRA2
	IF (I04) 15, 12, 15	OUTE 17
12	READ (NSORCE) (S2(I), I = 1, INJM)	OUTE 18
	GO TO 30	OUTE 19
15	DO 20 I=1, INJM	OUTE 20
20	S2(I)=0.	OUTE 21
30	IF(A02) 60, 40, 60	OUTE 22
40	DO 50 I=1, INJM	OUTE 23
50	S2(I)=S2(I)+K6(IGV)*F0(I)	OUTE 24
	GO TO 80	OUTE 25
60	DO 70 I=1, INJM	OUTE 26
	ITEMP=MO(1)	OUTE 27
	ITEMP1=M2(ITEMP1)	OUTE 28
70	S2(I)=S2(I)+C0(INT-1, ITEMPI)*F0(I)	OUTE 29
80	GBAR=IGV+INS-ITL	OUTE 30
	IF(GBAR-1) 90, 100, 100	OUTE 31

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 11

90	GBAR=1	OUTE 32
100	PBAR = IHS + IGV - 1	OUTE 33
	IF(PBAR - ITL) 115,115,110	OUTE 34
110	PBAR = ITL	OUTE 35
115	IF(GBAR - IGV) 120, 140, 140	OUTE 36
120	READ (NSCRAT) (N2(I),I=1,INJM)	OUTE 37
	DO 130 I=1,INJM	OUTE 38
	ITEMP1=N0(I)	OUTE 39
	ITEMP1=N2(ITEMP1)	OUTE 40
	ITEMP=ITEMP1	OUTE 41
	TEMP=C0(PBAR,ITEMP)	OUTE 42
130	S2(I)=S2(I)+N2(I)*TEMP	OUTE 43
	GO TO 150	OUTE 44
140	READ (NFLUX1) (N2(I),I=1,INJM)	OUTE 45
150	GBAR=GBAR+1	OUTE 46
	PBAR=PBAR-1	OUTE 47
	IF(GBAR - IGV) 120, 140, 160	OUTE 48
160	IF(IGV - IGM) 180, 170, 180	OUTE 49
170	REWIND NCR1	OUTE 50
	REWIND 11	HRA2
	REWIND 12	HRA2
180	V11=0.	OUTE 51
	DO 190 I=1,INJM	OUTE 52
	S2(I)=S2(I)*V0(I)	OUTE 53
190	V11=V11+S2(I)	OUTE 54
	E2(IGV) = V11 - E1(IGV)	OUTE 55
C	SOURCE-ALPHA	OUTE 56
200	IF(I04 - 2) 210, 240, 210	OUTE 57
210	IF(S02 - 2) 230, 220, 230	OUTE 58
220	T7 = S03/V7(IGV)	OUTE 59
	GO TO 250	OUTE 60
230	T7 = 0.0	OUTE 61
	GO TO 270	OUTE 62
240	T7 = EV/V7(IGV)	OUTE 63
250	DO 260 K = 1, 12M	OUTE 64
	ITEMP1 = N2(K)	OUTE 65
260	C0(IHS, ITEMP1) = C0(IHS,ITEMP1) - T7	OUTE 66
270	CONTINUE	OUTE 67

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 13

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          IF(V11) 450, 500, 450
450  IF(A02) 460, 480, 460
460  ED(IGV)=0.0
          DO 470 I=1,INJM
          ITEMP1=N0(I)
          ITEMP1=N2(ITEMP1)
          ED(IGV)=ED(IGV) + C0(IHT-3,ITEMP1)*N2(I)*V0(I)
470  F2(1)=F2(1)+K6(IGV)*N2(1)
          GO TO 500
480  ED(IGV)=0.0
          DO 490 I=1,INJM
          ITEMP1=N0(I)
          ITEMP1=N2(ITEMP1)
          ED(IGV)=ED(IGV) + C0(IHT-3,ITEMP1)*N2(I)*V0(I)
490  F2(1)=F2(1)+C0(IHT-1,ITEMP1)*N2(1)
500  CONTINUE
          IGV=IGV+1
          IF(IGV-IGN) 10, 10, 510
510  T11 = E1(IGP)
C    SWITCH TAPE DESIGNATIONS
          REWIND NCR1
          REWIND NSCRAT
          REWIND NFLUX1
          REWIND 11
          REWIND 12
          ITEMP = NSCRAT
          NSCRAT = NFLUX1
          NFLUX1 = ITEMP
          IF (I04) 514,512,514
512  REWIND NSORCE
514  CONTINUE
C
C    OVER-RELAX FISSION SOURCE
          ORFF = 1.0 + .6*(ORF - 1.0)
          E02 = .0
          IF(A02) 520,580,520
520  E1(IGP) = .0
C    FOR ADJOINT CALCULATION, S2(1) STORES ORFF F2(1)

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          QUTE 112
          QUTE 113
          QUTE 114
          QUTE 115
          QUTE 116
          QUTE 117
          HRA2 118
          QUTE 119
          QUTE 120
          QUTE 121
          QUTE 122
          QUTE 123
          QUTE 124
          HRA2 125
          QUTE 126
          QUTE 127
          QUTE 128
          QUTE 129
          QUTE 130
          QUTE 131
          QUTE 132
          QUTE 133
          QUTE 134
          HRA2
          HRA2
          QUTE 135
          QUTE 136
          QUTE 137
          QUTE 138
          QUTE 139
          QUTE 140
          QUTE 141
          QUTE 142
          QUTE 143
          QUTE 144
          QUTE 145
          QUTE 146
          QUTE 147

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 14

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DO 522 I=1,IMJM                                QUTE 148
522 S2(I) = F0(I) + ORFF*(F2(I) - F0(I))        QUTE 149
DO 540 IIG = 1,IGM                              QUTE 150
READ(NCR1) ((CO(I,J), I=1,ITL), J=1,MT)        QUTE 151
E1(IIG) = .0                                     QUTE 152
DO 530 I=1,IMJM                                QUTE 153
ITEMP = NO(I)                                   QUTE 154
ITEMP = M2(ITEMP)                              QUTE 155
E1(IIG) = E1(IIG) + CO(IHT-1,ITEMP)*F2(I)*VO(I) HRA2 156
530 E02 = E02 + CO(IHT-1,ITEMP)*S2(I)*VO(I)    HRA2 157
540 E1(IGP) = E1(IGP) + E1(IIG)                QUTE 158
TEMP1 = E1(IGP)/E02                             QUTE 159
DO 550 I=1,IMJM                                QUTE 160
550 F0(I) = TEMP1*S2(I)                         QUTE 161
REWIND NCR1                                     QUTE 162
REWIND 11                                       HRA2
REWIND 12                                       HRA2
GO TO 620                                       QUTE 163
580 E01 = 0.0                                    QUTE 164
DO 590 I=1,IMJM                                QUTE 165
E01 = E01 + VO(I)*F2(I)                         QUTE 166
F2(I) = F0(I) + ORFF*(F2(I) - F0(I))          QUTE 167
590 E02 = E02 + VO(I)*F2(I)                   QUTE 169
TEMP1=0.                                         QUTE 170
IF(E02.NE.0.0)TEMP1=E01/E02                    QUTE 171
DO 600 I=1,IMJM                                QUTE 172
600 F0(I) = TEMP1*F2(I)                        QUTE 173
DO 610 IIG = 1,IGM                              QUTE 174
610 E1(IIG) = K6(IIG)*E01                      QUTE 175
IF(104) 620,609,620                             QUTE 176
609 TEMP1 = .0
IF(E01.EQ. 0.0) GO TO 613                       QUTE 178
C ACCELERATION FOR EXTRANEIOUS SOURCE PROBLEMS QUTE 179
611 TEMP1 = (1.0 - EV*T11/E01)/(1.0 - EV)       QUTE 180
IF (T11/E01 - .01) 620,620,612                 QUTE 181
612 IF (T11/E01 - 1./(EV + .0001)) 613,613,620 QUTE 182
613 DO 614 I = 1,IMJM                           QUTE 183
614 F0(I) = TEMP1 * F0(I)                       QUTE 184

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 15

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DO 616 IIG = 1,IGM                                QUTE 185
EO(IIG) = TEMP1*EO(IIG)                            QUTE 186
616 E1(IIG) = TEMP1*E1(IIG)                          QUTE 187
620 E1(IGP) = 0.0                                    QUTE 188
EO(IGP) = 0.0                                        QUTE 189
DO 640 IIG = 1,IGM                                QUTE 190
EO(IGP) = EO(IGP) + EO(IIG)                       QUTE 191
640 E1(IGP) = E1(IGP) + E1(IIG)                   QUTE 192
RETURN                                              QUTE 193
END                                                  QUTE 194
SUBROUTINE PRT (JIM,JJM, N2, Z4, NOUT)              PRT 2
DIMENSION N2(JIM,JJM), Z4(1)                       PRT 3
REAL N2                                             PRT 4
C DATA XRR/6HXRR  ',YZT/6HYZT  /                  PRT 5
CHARACTER*6 XRR, YZT                               HVX
DATA XRR/'XRR  ',YZT/'YZT  '/                     HVX
DATA LINES/O/                                       PRT 6
JM = JIM                                           PRT 7
JJ = JJM                                           PRT 8
DO 50 I=1,IM,6                                     PRT 10
I1=I
I2=I+5
IF(I2-IM) 20, 20, 10                              PRT 12
10 I2=IM                                           PRT 13
20 WRITE(NOUT,30) YZT,(XRR,JJ,JJ=11,12)           PRT 14
30 FORMAT(1X,A3,6(3X,A3,14,1X))
DO 50 JJ=1,JM                                       PRT 16
J=JJ                                               PRT 17
40 FORMAT (14, 6E11.4, F9.3)
IF(J.EQ.1)GO TO 45                                  PRT 19
DO 42 K=11,12                                       PRT 20
IF(N2(K,J).NE.N2(K,J-1)) GO TO 43                 PRT 21
42 CONTINUE                                         PRT 22
LINES=LINES+1                                       PRT 24
IF(J-JM) 50,43,43                                    PRT 25
43 IF(LINES.EQ.0)GO TO 45                            PRT 26
WRITE(NOUT,44) LINES                                PRT 27
44 FORMAT('      NEXT',15,' LINES SAME AS PRECEDING LINE')
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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 16

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      LINES=0                                PRT 29
      IF(J-JM) 45,50,45                       PRT 30
45  WRITE(NOUT,40)J,(N2(K,J),K=11,12), 24(J)
50  CONTINUE                                  PRT 34
      RETURN                                  PRT 35
      END                                      PRT 36
      SUBROUTINE AVERAG(PHIB,AXS,FXS,MATN,MASS,ATW,VOL,C0,N2,M0,V0,
C   1  HOLN, JML, JTL, NBR,T6,TB,MTWON)      HRA2
      1  HOLN, JML, JTL, NBR,ACT,JNACT,IACPOS) HRA3
      DIMENSION PHIB(1), AXS(JML,1), FXS(JML,1), MATN(1), MASS(JML,1),
      1  ATW(1), VOL(1), C0(JTL,1), N2(1), M0(1), V0(1),
      2  HOLN(JML,1), NBR(1)
C   DIMENSION T6(JML,1),TB(JML,1)           HRA2
      DIMENSION ACT(JML,1,JML,1),IACPOS(JNACT),T1(20,20),ACTIV(20,20) HRA3
      INCLUDE 'ABC.FOR'
C   THIS SUBROUTINE CALCULATES ZONE AVERAGED FLUXES, FISSION CROSS
C   SECTIONS, AND ABSORPTION CROSS SECTIONS.
      RL = 0.0                                AVER 8
      RC = 0.0                                AVER 9
      DO 10 KZ=1,12M                           AVER 10
      PHIB(KZ) = 0.0                            AVER 11
      DO 10 KN =1,NCOM                           AVER 12
      AXS(KN,KZ) = 0.0                          AVER 13
      FXS(KN,KZ) = 0.0                          AVER 14
      T6(KN,KZ) = 0.0                            AVER 15
C   T6(KN,KZ) = 0.0                            HRA2
C   TTOTAL = 0.0                               HRA2
C   TB(KN,KZ) = 0.0                            HRA2
      DO 1000 I=1,NACT                           HRA3
      K=IACPOS(I)                                HRA3
      LN = MATN(KN)                              HRAH
      ACT(LN,K,KN,KZ) = 0.0                     HRA3
1000 CONTINUE                                  HRA3
      LN = MATN(KN)                              AVER 17
      IF (MASS(LN,KZ) .EQ. 0) GO TO 10
      MASS(LN,KZ) = (MASS(LN,KZ)*.6023)/(ATW(LN)*VOL(KZ))
      AVER 18
10  CONTINUE
      DO 100 IIG=1,IGM                           AVER 19
      READ(NCR1) ((CO(I1,J), I1=1,ITL),J=1,MT)  AVER 20

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 17

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      READ(NFLUX1) (N2(I), I=1,IMJM)
      DO 11 J=1,MCR
C      CALL REAG2(' RO',AC0(1,MT),3)
11  CONTINUE
      DO 100 I=1,IMJM
      KZ = N0(I)
      PHIB(KZ) = PHIB(KZ) + N2(I)*V0(I)
      DO 100 KN=1,MCON
      LN = MATN(KN)
C      CO(5,1 AND 2) ARE ACTIVITIES FOR TPROD
C      IF(LN.EQ.1)THEN
C      T6(KN,KZ) = T6(KN,KZ) + CO(5,LN)*N2(I)*V0(I)
C      ENDF
C      IF(LN.EQ.2)THEN
C      T8(KN,KZ) = T8(KN,KZ) + CO(5,LN)*N2(I)*V0(I)
C      ENDF
      DO 2000 K=1,NACT
      J=IACPOS(K)
      ACT(LN,J,KN,KZ) = ACT(LN,J,KN,KZ) + CO(J,LN)*N2(I)*V0(I)
C      WRITE(6,*)J,ACT(J,KN,KZ),IACPOS(K)
2000 CONTINUE
      AXS(KN,KZ) = AXS(KN,KZ) + CO(IHT-2,LN)*N2(I)*V0(I)
100  FXS(KN,KZ) = FXS(KN,KZ) + CO(IHT-3,LN)*N2(I)*V0(I)
C      DO 200 KZ=1,IZM
      DO 209 KZ=1,IZM
C      TRIT6=0.0
C      TRIT7=0.0
      DO 2500 KI=1,NACT
      DO 2480 LL=1,MCON
      MM=MATN(LL)
      ACTIV(MM,KI)=0.0
2480 CONTINUE
2500 CONTINUE
      TEMP3 = PHIB(KZ)
      IF (PHIB (KZ) .EQ. 0) GO TO 105
      PHIB(KZ) = PHIB(KZ)/(VOL(KZ)*1000.)
105  CONTINUE
      WRITE(NOUT,110) KZ, PHIB(KZ), VOL(KZ)

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AVER 21
 HRA2
 HRA2
 HRA2
 AVER 22
 AVER 23
 AVER 24
 AVER 25
 AVER 26
 HRA2
 HRA2
 HRA2
 HRA2
 HRA2
 HRA2
 HRA3
 HRA3
 HRA3
 HRA3
 HRA2 27
 HRA2 28
 HRA2 29
 HRA2 29
 HRA2
 HRA2
 HRA3
 HRAH
 HRAH
 HRA3
 HRAH
 HRAH
 HRA3
 AVER 30
 AVER 31
 AVER 32

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 18

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110  FORMAT(1H1,2X,9H Z O N E ,13,5X,7H FLUX =,1PE10.4,5X,9H VOLUME =,
      1 1PE10.4,7H LITERS/)
      WRITE(NOUT,120)
      AVER 34
120  FORMAT (1X,'BURNABLE MAT.   NAME   ATOM '
      1,'  FISSON  ABSORPTION SIGMA  SIGMA ' /
      3 1X,'ISOTOPE NO. ', 15X,'DENSITY',4X,'RATE',8X,'RATE',4X,
      3 'FISSION',2X,'ABSORPTION' /)
      DO 200  KN=1,NCON
      LN = MATH(KN)
      AVER 41
      TEMP1 = AXS(KN,KZ)*MASS(LN,KZ)
      AVER 42
      TEMP2 = FXS(KN,KZ)*MASS(LN,KZ)
      AVER 43
      AVER 44
      C  T1 = T6(KN,KZ)*MASS(LN,KZ)
      HRA2
      C  T2 = T8(KN,KZ)*MASS(LN,KZ)
      HRA2
      C  T6(KN,KZ) = T6(KN,KZ)/TEMP3
      HRA2
      C  T8(KN,KZ) = T8(KN,KZ)/TEMP3
      HRA2
      DO 3000 K=1,NACT
      HRA3
      J=IACPOS(K)
      HRA3
      T1(LN,K) = ACT(LN,J,KN,KZ)*MASS(LN,KZ)
      HRA3
      ACT(LN,J,KN,KZ) = ACT(LN,J,KN,KZ)/TEMP3
      HRA3
3000  CONTINUE
      HRA3
      AXS(KN,KZ) = AXS(KN,KZ)/TEMP3
      AVER 45
      FXS(KN,KZ) = FXS(KN,KZ)/TEMP3
      AVER 46
130  FORMAT (1X, 13, 5X, 13, 4X, 2A4, 5E11.4)
      HRA2
      WRITE(NOUT,130) KN, LN, (HOLN(LN,K),K=1,2), MASS(LN,KZ), TEMP2,
      1 TEMP1,FXS(KN,KZ), AXS(KN,KZ)
      C  TRIT6=TRIT6+T1
      HRA2
      C  TRIT7=TRIT7+T2
      HRA2
      C  TTOTAL=TTOTAL+T1+T2
      HRA2
      DO 4000 KKK=1,NACT
      HRA3
      ACTIV(LN,KKK)=ACTIV(LN,KKK)+T1(LN,KKK)
      HRA3
4000  CONTINUE
      ITEMP = NBR(KN)
      AVER 50
      IF(ITEMP - 1) 200, 140, 160
      AVER 51
140  RC = RC + TEMP1 - TEMP2
      AVER 52
      GO TO 200
      AVER 53
160  RL = RL + TEMP1
      AVER 54
200  CONTINUE
      AVER 55
      C  WRITE(6,134)TRIT6,TRIT7
      HRA2

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 19

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C134  FORMAT(' T6=',E11.4,3X,'T7=',E11.4)          HRA2
      DO 5500 JJ=1,NCON                             HRAH
      LN=MATN(JJ)                                   HRAH
      DO 5000 KKK=1,NACT                             HRA3
      WRITE(6,1340)KKK,LN,ACTIV(LN,KKK)             HRA3
1340  FORMAT(' ACTIVITY',I2,' FOR MAT. NO.',I2,' ',E11.4) HRA3
5000  CONTINUE                                       HRA3
5500  CONTINUE                                       HRAH
      WRITE(6,*)
      WRITE(6,*)
      WRITE(6,*)
209   CONTINUE                                       HRA2
      IF(RL.EQ.0.0)THEN                             HRA2
          TEMP=0.0                                   HRA2
          GO TO 340                                   HRA2
      ENDIF                                          HRA2
      TEMP = RC/RL                                   AVER 56
      C  WRITE(NOUT,350) TEMP                         AVER 5
7
340  WRITE(NOUT,350) TEMP                           HRA2
350  FORMAT(1H '///' BREEDING RATIO                 =' ,F7.4)
      C  WRITE(6,401) TTOTAL                         HRA2
C401  FORMAT(' TTOTAL=',E11.4)                   HRA2
      REWIND NCR1                                    AVER 59
      REWIND NFLUX1                                  AVER 60
      RETURN                                         AVER 61
      END                                            AVER 62
      SUBROUTINE CLEAR (X,Y,N)                      CLER 2
      DIMENSION Y(1)                                CLER 3
      DO 1 I=1,N                                     CLER 4
1     Y(I)=X                                         CLER 5
      RETURN                                         CLER 6
      END                                            CLER 7
      SUBROUTINE CNMP (F2,K6)                       CNMP 2
      DIMENSION F2(1), K6(1)                       CNMP 3
      INCLUDE 'ABC.FOR'

      IF (MAXT) 25,25,10

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TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 20

10	JUMP=2	CNMP	5
	CALL TCHEK(GLH,JUMP)	CNMP	6
	GO TO (15,25),JUMP	CNMP	#
15	WRITE(NGOUT,20)	CNMP	8
20	FORMAT(53H1 * * RUNNING TIME EXCEEDED--FORCED CONVERGENCE * *//)	CNMP	9
	GO TO 90	CNMP	10
25	CONTINUE	CNMP	11
30	E01=1.0-ALA	CNMP	12
	IF(ABS (E01)-10.0*EPS) 40, 40, 45	CNMP	13
40	ORF = ORFP	CNMP	14
45	CONTINUE	CNMP	15
	E02=ABS(E01)	CNMP	16
50	IF(E1(1GP)) 55, 130, 55	CNMP	17
55	IF (E02 - EPS) 60, 60, 70	CNMP	18
60	CVT=1	CNMP	19
70	CALL CLEAR (0.0, F2, INJM)	CNMP	23
	GO TO 105	CNMP	24
80	EV=EV+POD*EQ*E01	CNMP	25
	GO TO 170	CNMP	26
C	FINAL PRINT	CNMP	27
90	NGOTO=1	CNMP	28
	IF (104 - 1) 135, 95, 80	CNMP	29
95	EV=0.0	CNMP	30
	DO 100 I=1,1GM	CNMP	31
100	EV=EV+K6(I)	CNMP	32
	EV=SK7/EV	CNMP	33
	GO TO 135	CNMP	34
105	IF(CVT-1) 110, 90, 110	CNMP	35
110	IF(104-1) 115, 120, 140	CNMP	36
C	MONITOR PRINT	CNMP	37
115	NGOTO=2	CNMP	38
	GO TO 135	CNMP	39
120	EV=0.	CNMP	40
	DO 125 I=1,1GM	CNMP	41
125	EV=EV+K6(I)	CNMP	42
	EV=SK7/EV	CNMP	43
	GO TO 115	CNMP	44
130	IF(104.EQ.0)GO TO 55	CNMP	45

TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 21

CALL ERRO2(6H**CNNP,130,1)	CNNP 46
135 RETURN	CNNP 47
140 CONTINUE	CNNP 48
C	CNNP 49
C CALCULATE NEW PARAMETERS FOR SEARCH CALCULATIONS	CNNP 50
145 EQ3=ABS (ALA-LAR)	CNNP 51
IF (LAPP) 270, 150, 270	CNNP 52
150 IF (LAP) 230, 155, 230	CNNP 53
155 IF (EQ) 200, 160, 200	CNNP 54
160 IF (EQ3-EPSA) 175, 175, 165	CNNP 55
C MONITOR PRINT.	CNNP 56
165 NGOTO=2	CNNP 57
RETURN	CNNP 58
C FINAL PRINT EXIT.	CNNP 59
170 NGOTO=1	CNNP 60
RETURN	CNNP 61
175 LAP=ALA	CNNP 62
EVP=EV	CNNP 63
IF (E01) 185,185,180	CNNP 64
180 EV=EV-EVM	CNNP 65
GO TO 190	CNNP 66
185 EV=EV+EVM	CNNP 67
190 IF (104-2) 195, 165, 195	CNNP 68
C MIX X-SECS.	CNNP 69
195 NGOTO=3	CNNP 70
RETURN	CNNP 71
200 IF (CVT) 170, 205,170	CNNP 72
205 EV=EV+POD*EQ*E01	CNNP 73
210 IF ((LAPP-1.0)/(LAP-1.0)) 215, 190, 190	CNNP 74
215 TEMP1=AMIN1(EVP,EVPP)	CNNP 75
IF (EV-TEMP1) 220, 225, 225	CNNP 76
220 EV=(EVPP+EVP)/2.	CNNP 77
GO TO 190	CNNP 78
225 TEMP1=AMAX1(EVP,EVPP)	CNNP 79
IF (EV-TEMP1) 190, 220, 220	CNNP 80
230 IF (E03-EPSA) 235, 235, 165	CNNP 81
235 EQ=(EVP-EV)/(LAP-ALA)	CNNP 82
240 IF (CNT) 260, 245, 260	CNNP 83

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 22

245	IF (E02-LAL) 265, 265, 250	CNNP 84
250	IF (E02-LAH) 260, 260, 255	CNNP 85
255	E01=SIGN (LAH,E01)	CNNP 86
260	LAPP=LAP	CNNP 87
	LAP=ALA	CNNP 88
	EVPP=EVPP	CNNP 89
	EVP=EV	CNNP 90
	GO TO 205	CNNP 91
265	CNT=1	CNNP 92
	LAP=0.0	CNNP 93
	LAPP=0.0	CNNP 94
	GO TO 205	CNNP 95
270	IF (E03-EPSA) 275, 275, 165	CNNP 96
C	CALCULATE QUADRATIC COEFFICIENTS.	CNNP 97
275	TEMP1=EVPP-EV	CNNP 98
	TEMP2=EVPP-EV	CNNP 99
	TEMP3=EVPP-EVPP	CNNP 100
	TEMP4=TEMP1*(EVP+EV)	CNNP 101
	TEMP5=-TEMP2*(EV+EVPP)	CNNP 102
	TEMP6=TEMP3*(EVPP+EVP)	CNNP 103
	DENOM=TEMP3*TEMP2*TEMP1	CNNP 104
	EQA=((LAPP-1.0)*TEMP1*EVP*EV-(LAP-1.0)*TEMP2	CNNP 105
	1*EV*EVPP+(ALA-1.0)*TEMP3*EVPP*EVP)/DENOM	CNNP 106
	EQB=- (LAPP*TEMP4+LAP*TEMP5+ALA*TEMP6)/DENOM	CNNP 107
	EQC=(LAPP*TEMP1-LAP*TEMP2+ALA*TEMP3)/DENOM	CNNP 108
	DISCR=EQB*EQB-4.0*EQA*EQC	CNNP 109
	IF (DISCR) 235, 280, 280	CNNP 110
280	IF (E02-LAL) 265, 265, 285	CNNP 111
285	TEMP1=EQC+EQC	CNNP 112
	TEMP=SQRT (DISCR)	CNNP 113
	EQ=1.0/(EQB+EV*TEMP1)	CNNP 114
	LAPP=LAP	CNNP 115
	LAP=ALA	CNNP 116
	EVPP=EVPP	CNNP 117
	EVP=EV	CNNP 118
	EV1=(TEMP-EQB)/TEMP1	CNNP 119
	EV2=- (TEMP+EQB)/TEMP1	CNNP 120
	EVA=ABS (EV-EV1)	CNNP 121

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 23

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      EVB=ABS (EV-EV2)
      IF (EVA-EVB) 290, 290, 295
290  EV=EV1
      GO TO 210
295  EV=EV2
      GO TO 210
      END
      SUBROUTINE ERROZ( HOL,JSUBR,I)
      COMMON NSORCE, NINP, NOUT, NCR1, NFLUX1, NSCRAT
      CHARACTER*6 HOL
      DATA NERR/0/
      NERR=NERR+1
      WRITE (NOUT,1)          HOL,JSUBR
1    FORMAT(2H */9H ERROR IN,A6,3H AT,16/2H */2H *)
      IF(NERR.EQ.100)GO TO 3
      GO TO (3,4),I
3    STOP
4    RETURN
      END
      SUBROUTINE FISCAL (NO, FO, VO, CO, K6, MO, M2, JTL,JMT)
      INCLUDE 'ABC.FOR'

      DIMENSION NO(1), FO(1), VO(1), CO(JTL,1),K6(1), MO(1), M2(1)
      LAR = ALA
      C    FISSION SUMS
      IF(B07) 90,90,10
10     IF(A02) 20, 40, 20
20     DO 30 IIG=1,1GM
      READ (NCR1) ((CO(I,J),I=1,ITL),J=1,NT)
      E1(IIG)=0.
      DO 30 I=1,IMJM
      ITEMF=MO(I)
      ITEMF=M2(ITEMF)
30     E1(IIG)=E1(IIG)+CO(IHT-1,ITEMF)*FO(I)*VO(I)
      REWIND NCR1
      GO TO 70
40     E01=0.
      DO 50 I=1,IMJM

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CNMP 122

CNMP 123

CNMP 124

CNMP 125

CNMP 126

CNMP 127

CNMP 128

ERR2 2

ERR2 3

HVX

ERR2 4

ERR2 7

ERR2 8

ERR2 9

ERR2 1#

ERR2 12

ERR2 13

ERR2 14

FISC 2

FISC 5

FISC 6

FISC 7

FISC 8

FISC 9

FISC 10

FISC 11

FISC 12

FISC 13

FISC 14

FISC 15

FISC 16

FISC 17

FISC 18

FISC 19

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 24

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50  E01=E01+V0(I)*F0(I)          FISC 20
    DO 60 IIG=1,IGM             FISC 21
60  E1(IIG)=K6(IIG)*E01        FISC 22
70  E1(IGP)=0.                 FISC 23
    EO(IGP)=0.                 FISC 24
    DO 80 IIG=1,IGM             FISC 25
    EO(IGP)=EO(IGP)+EO(IIG)    FISC 26
80  E1(IGP)=E1(IGP)+E1(IIG)    FISC 27
    IF(B07) 140, 90, 140       FISC 28
90  IF(T11.EQ. 0.0)GO TO 95
    ALA=E1(IGP)/T11
    TEMP=1.0/ALA                FISC 31
95  IF(I04-1) 230,100,140
100 DO 110 IIG=1,IGM           FISC 32
    E1(IIG)=E1(IIG)*TEMP       FISC 33
110 K6(IIG)=K6(IIG)*TEMP      FISC 34
    E1(IGP)=E1(IGP)*TEMP      FISC 35
    IF(A02) 120, 140, 120     FISC 36
120 DO 130 I=1,IMJM           FISC 37
130 F0(I)=F0(I)*TEMP          FISC 38
140 CONTINUE                  FISC 39
C                               FISC 40
C  NORMALIZATION              FISC 41
    B07=0                      FISC 42
150 IF(S01) 160, 230, 170     FISC 43
160 E01 = ABS(S01)/(EO(IGP)*TSD) FISC 44
    GO TO 180                  FISC 45
170 E01=S01/E1(IGP)           FISC 46
180 DO 190 IIG=1,IGP          FISC 47
190 E1(IIG)=E01*E1(IIG)      FISC 48
    DO 200 I=1,IMJM           FISC 49
200 F0(I)=E01*F0(I)          FISC 50
230 RETURN                    FISC 51
    END                        FISC 52
    SUBROUTINE GRAM(MASS, VOL,  ATW, HOLN,J1M,JJM, M0, M2, V0,  GRAM 2
1      10, I1, I2,JML, I3)    GRAM 3
    INCLUDE 'ABC.FOR'

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TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 25

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        DIMENSION MASS(JML,1), VOL(1), ATW(1), HOLN(JML,1), MO(JIM,JJM),
1      MZ(1), VO(JIM,JJM), IO(1), I1(1), I2(1), I3(1)      GRAM 6
C      THIS SUBROUTINE CALCULATES THE MASS OF THE VARIOUS MATERIALS      GRAM 7
      WRITE(NOUT,10) (IO(1), I=1,20)      GRAM 8
10     FORMAT(1H1,20A4///)      GRAM 9
      WRITE(NOUT, 20)      GRAM 10
20     FORMAT(45H MATERIAL INVENTORY (KILOGRAMS) FOR EACH ZONE / )      GRAM 11
      CALL CLEAR(0.0,VOL,IZM)      GRAM 12
      ITEMP = ML*IZM      GRAM 13
      CALL CLEAR(0.0,MASS,ITEMP)      GRAM 14
      DO 30 J = 1, JM      GRAM 15
      DO 30 I = 1, IM      GRAM 16
      K = MO(I,J)      GRAM 17
30     VOL(K) = VOL(K) + VO(I, J)*.001      GRAM 18
      DO 39 M=1,MO1      GRAM 19
      I3(M) = I2(M)      GRAM 20
      IF(I0(M) - I1(M)) 39,35,39      GRAM 21
35     IF(I2(M)) 39,36,39      GRAM 22
36     DO 38 MM=1,M      GRAM 23
      IF(I0(M) - I0(MM)) 38,37,38      GRAM 24
37     I3(MM) = I2(MM)*EV      GRAM 25
38     CONTINUE      GRAM 26
39     CONTINUE      GRAM 27
      DO 190 N = 1, IZM      GRAM 28
      NN = MZ(N)      GRAM 29
      DO 190 M = 1,MO1      GRAM 30
      IF(I0(M) - NN) 190, 40, 190      GRAM 31
40     L = I1(M)      GRAM 32
      IF(L - ML) 170, 170, 50      GRAM 33
50     NAAA = L      GRAM 34
      IF(L - I0(M)) 130,190, 130      GRAM 35
130    DO 160 MAA = 1, MO1      GRAM 36
      IF(I0(MAA) - NAAA) 160, 140, 160      GRAM 37
140    L = I1(MAA)      GRAM 38
      IF(L) 160, 160, 150      GRAM 39
150    EO1 = I3(MAA)*I3(M)      GRAM 40
      MASS(L,N) = ((EO1*ATW(L)*VOL(N))/ .6023) + MASS(L,N)      GRAM 41
160    CONTINUE      GRAM 42

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 26

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      GO TO 190                                     GRAM 43
170 IF(L) 190, 190, 180                           GRAM 44
180 E01 = I3(M)                                    GRAM 45
      MASS(L,N) = ((E01*ATM(L)*VOL(N))/ .6023) + MASS(L,N) GRAM 46
190 CONTINUE                                       GRAM 47
      DO 260 L = 1, IZM, 5                          GRAM 49
      LL = L + 4                                    GRAM 50
      IF(LL - IZM) 210, 210, 200                   GRAM 51
200 LL = IZM                                       GRAM 52
210 WRITE(NOUT,220) ((K), K=L, LL)
220 FORMAT (///' MATERIAL      ATOMIC WT.      ',2X,5(' ZONE',I3,3X)/)
      DO 240 K = 1, ML                               GRAM 57
      DO 233 I=L,LL                                  UPD1 3
      IF(MASS(K,I) .NE. 0.) GO TO 238               UPD1 4
233 CONTINUE                                       UPD1 5
      GO TO 240                                     UPD1 6
238 WRITE(NOUT,250) K,(HOLN(K,N),N=1,2),ATW(K), (MASS(K,I), I=L,LL)

240 CONTINUE                                       UPD1 8
250 FORMAT (1X,I3,1X,2A4, F12.2, 1X, 5E11.3)
      IF(LL - IZM) 260, 270, 270
260 CONTINUE                                       GRAM 60
      GRAM 61
C COMPUTE TOTAL MASSES                             UPD1 9
270 WRITE (NOUT,275)
275 FORMAT (///' MATERIAL      ATOMIC WT.      TOTAL'//)
      DO 310 K=1,ML                                  UPD1 16
      TEMP=0.0
      DO 280 L=1,IZM
280 TEMP=TEMP+MASS(K,L)
310 WRITE(NOUT,250) K,(HOLN(K,N),N=1,2),ATW(K),TEMP

      WRITE (NOUT,350)
350 FORMAT (///,' ZONE NUMBER      VOLUME (LITERS)'//)
      DO 400 L=1,IZM
      WRITE (NOUT,360) L,VOL(L)
360 FORMAT (6X,I4,6X,1PE12.3)
400 CONTINUE
      RETURN                                       UPD1 18

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TL1STZ.FOR

Wednesday, February 21, 1990 12:31 pm

Page 27

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END
SUBROUTINE IFLUXN (N2, CO, VO, CXS, MO, M2, JTL,JIM,JJM, CXR, CXT HRAZ 2
1 ,XR,XD) HRAZ 2
INCLUDE 'ABC.FOR'

DIMENSION N2(1), CO(JTL,1), VO(1),CXS(JIM,JJM,3),MO(1), M2(1), IFLU 4
1 CXR(1), CXT(1) IFLU 5
DIMENSION XD(50),XR(50) HRAZ 2

C THIS SUBROUTINE NORMALIZES FLUXES BEFORE EACH INNER ITERATION IFLU 6
C ABSORPTION AND OUT-SCATTER IFLU 7
E3(IGV) = 0.0 IFLU 8
E4(IGV) = 0.0 IFLU 9
DO 10 I=1, IMJM IFLU 10
TEMP = VO(I)*N2(I) IFLU 11
ITEMP = MO(I) IFLU 12
ITEMP = M2(ITEMP) IFLU 13
E3(IGV) = E3(IGV) + (XD(ITEMP))*TEMP HRAZ 14
10 E4(IGV) = E4(IGV) + CO(INT-2,ITEMP)*TEMP HRAZ 15
C LEFT LEAKAGE IFLU 16
IF(B01) 20, 20, 40 IFLU 17
20 E5(IGV) = 0.0 IFLU 18
DO 30 KJ = 1, JM IFLU 19
I = (KJ - 1)*IM + 1 IFLU 20
30 E5(IGV) = E5(IGV) + CXS(1,KJ,1)*N2(I) IFLU 21
GO TO 50 IFLU 22
40 E5(IGV) = .0 IFLU 23
C RIGHT LEAKAGE IFLU 24
50 IF(B02) 60, 60, 80 IFLU 25
60 E6(IGV) = 0.0 IFLU 26
DO 70 KJ = 1, JM IFLU 27
I = KJ*IM IFLU 28
70 E6(IGV) = E6(IGV) + CXR(KJ)*N2(I) IFLU 29
GO TO 90 IFLU 30
80 E6(IGV) = 0.0 IFLU 31
C TOP LEAKAGE IFLU 32
90 IF(B03-1) 120, 140, 100 IFLU 33
100 E7(IGV) = .0 IFLU 34
DO 110 K1 = 1,IM IFLU 35

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TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 28

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      I = IMJM - IM + KI                                IFLU 36
110  E7(IGV) = E7(IGV) + CXS(K1,1,2)*(N2(1) - N2(K1)) IFLU 37
      EB(IGV) = - E7(IGV)                               IFLU 38
      GO TO 190                                         IFLU 39
120  E7(IGV) = 0.0                                     IFLU 40
      DO 130 KI = 1, IM                                IFLU 41
      I = IMJM - IM + KI                                IFLU 42
130  E7(IGV) = E7(IGV) + CXT(K1)*N2(1)                IFLU 43
      GO TO 150                                         IFLU 44
140  E7(IGV) = 0.0                                     IFLU 45
C     BOTTOM LEAKAGE                                   IFLU 46
150  IF(B04) 160, 160, 180                             IFLU 47
160  EB(IGV) = 0.0                                     IFLU 48
      DO 170 KI = 1, IM                                IFLU 49
170  EB(IGV) = EB(IGV) + CXS(K1,1,2)*N2(K1)          IFLU 50
      GO TO 190                                         IFLU 51
180  EB(IGV) = 0.0                                     IFLU 52
190  E9(IGV) = E5(IGV) + E6(IGV) + E7(IGV) + EB(IGV) IFLU 53
      TEMP = (E1(IGV) + E2(IGV))/(E3(IGV) + E4(IGV) + E9(IGV)) IFLU 54
      DO 200 I = 1, IMJM                                IFLU 55
200  N2(I) = TEMP*N2(1)                                 IFLU 56
      E3(IGV) = TEMP*E3(IGV)                            IFLU 57
      E4(IGV) = TEMP*E4(IGV)                            IFLU 58
      E5(IGV) = TEMP*E5(IGV)                            IFLU 59
      E6(IGV) = TEMP*E6(IGV)                            IFLU 60
      E7(IGV) = TEMP*E7(IGV)                            IFLU 61
      EB(IGV) = TEMP*EB(IGV)                            IFLU 62
      E9(IGV) = TEMP*E9(IGV)                            IFLU 63
      RETURN                                             IFLU 64
      END                                               IFLU 65
SUBROUTINE INIT (K6, K7, I0, I1, I2, M0, M2, N0, R0, R1, R2,
1      R3, R4, R5, Z0, Z1, Z2, Z3, Z4, Z5, A0, A1,
2      F0, C0, V0, JTL, JIM, V7, JJM, JMT, JML, GAM, HOLN)
      INCLUDE 'ABC.FOR'

DIMENSION K6(1), K7(1), I0(1), I1(1), I2(1), R0(1), R1(1),
1      R2(1), R3(1), R4(1), R5(1), Z0(1), Z1(1), Z2(1),
2      Z3(1), Z4(1), Z5(1), A0(1), A1(1), C0(JTL, JMT),
INIT 6
INIT 7
INIT 8

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 29

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3          VO(JIM,JJM), M0(1), M2(1), N0(1), F0(1), V7(1)      INIT  9
4          ,GAM(1), HOLN(JML,1)
      DIMENSION XD(50),XR(50)                                     HRAZ
      IF (P02) 20, 10, 20                                         INIT 11
10 WRITE(NOUT,15) DAY                                           INIT 12
15 FORMAT(1H1,24X,' T I M E  =' ,F8.3,8H D A Y S///)
20 CONTINUE                                                     INIT 14
C  ADJOINT REVERSALS                                           INIT 15
      IF(A02) 25, 45, 25                                         INIT 16
25 IF(P02) 45, 30, 45                                         INIT 17
30 IF(NCON) 45, 35, 45                                         INIT 18
35 IIG=1                                                         INIT 19
      IGBAR=IGM                                                  INIT 20
40 TEMP=K7(IIG)                                                 INIT 21
      K7(IIG)=K7(IGBAR)                                         INIT 22
      K7(IGBAR)=TEMP                                           INIT 23
      TEMP=V7(IIG)                                              INIT 24
      V7(IIG)=V7(IGBAR)                                         INIT 25
      V7(IGBAR)=TEMP                                           INIT 26
      IIG=IIG+1                                                 INIT 27
      IGBAR=IGBAR-1                                             INIT 28
      IF(IIG-IGBAR) 40, 45, 45                                  INIT 29
45 CONTINUE                                                     INIT 30
C  MIX CROSS-SECTIONS                                          INIT 32
      B07=1                                                      INIT 33
      IF(P02) 50, 55, 50                                         INIT 34
50 GO TO (245,245,85,245,185), 104                              INIT 3#
55 IF(M01) 70, 70, 60                                          INIT 36
60 WRITE(NOUT,61)
61 FORMAT(1H0,4X,16H MIXTURE NUMBER ,18H MIX COMMAND ,
      124H MATERIAL ATOMIC DENSITY /)
      DO 67 J=1,M01
      WORD = 0
      WORD2 = WORD
      I=1(J)
      UPD1 22
      UPD1 23
      UPD1 24
      UPD1 25
      UPD1 27
      IF(1 .GT. 0 .AND. I .LE. NL) WORD=HOLN(I,1)
      IF(1 .GT. 0 .AND. I .LE. NL) WORD2=HOLN(I,2)
      HVX
      WRITE(NOUT,63) J, I0(J), I1(J), I2(J), WORD, WORD2

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 30

63	FORMAT(15,19,115,E28.8,12X,2A4)	
	IF(J.EQ.M01.OR.I0(J).EQ.I0(J+1)) GO TO 67	UPD1 31
66	FORMAT(1X,15(5H-----))	
	WRITE(NOUT,66)	UPD1 33
67	CONTINUE	UPD1 34
70	IF(NPRT-1) 85, 85, 75	
75	WRITE(NOUT,80)	INIT 41
80	FORMAT(/19H1CROSS-SECTION EDIT)	INIT 42
85	REWIND NCR1	INIT 43
	DO 180 IIG=1,IGM	INIT 44
	READ(NCR1) ((CO(I,J),I=1,ITL),J=1,MT)	INIT 45
C	READ(11)(XR(J),J=1,MT)	HRA2
	READ(12)(XD(J),J=1,MT)	HRA2
	IF(M01) 90, 145, 90	INIT 46
90	DO 140 M=1,M01	INIT 47
	IF(10(M)-MT) 100, 100, 95	INIT 48
95	CALL ERRO2('**INIT',95,1)	HVX
100	IF(11(M)-MT) 105, 105, 95	INIT 50
105	N=10(M)	INIT 51
	L=I1(M)	INIT 52
	E01=I2(M)	INIT 53
	IF(L) 125, 125, 110	INIT 54
110	IF(E01) 125, 115, 125	INIT 55
115	IF(N-L) 125, 120, 125	INIT 56
120	E01 = EV	INIT 57
	L = 0	INIT 58
125	DO 140 I=1,ITL	INIT 59
	IF(L) 130, 135, 130	INIT 60
130	CO(I,N)=CO(I,N)+CO(I,L)*E01	INIT 61
	GO TO 140	INIT 62
135	CO(I,N)=CO(I,N)*E01	INIT 63
140	CONTINUE	INIT 64
	IF(M01) 900, 145, 900	HRA2 46
900	DO 1140 M=1,M01	HRA2 47
	IF(10(M)-MT) 1000, 1000, 950	HRA2 48
950	CALL ERRO2('**INIT',950,1)	HRA2 HVX
1000	IF(11(M)-MT) 1050, 1050, 950	HRA2 50
1050	N=10(M)	HRA2 51

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 34

	Z4(J)=(Z1(J+1)+Z1(J))*0.5	WLP
	Z5(J)=Z1(J+1)-Z1(J)	WLP
	IF(Z5(J)) 350, 350, 355	WLP
350	CALL ERRO2 ('*Z5(J)',350,1)	HVX
355	CONTINUE	WLP
	DO 370 I=1,IM	WLP
	GO TO (360,365,367,360), IGP	HRA2
360	V0(I,J)=R5(I)*Z5(J)	WLP
	GO TO 370	WLP
365	V0(I,J)=P12*R5(I)*Z5(J)*R4(I)	HRA2
	GO TO 370	HRA2
367	V0(I,J)=(R5(I)*Z5(J)*R4(I)+((R1(I+1))**3-(R1(I))**3)/3.0	HRA2
	1 /BIGR*(SIN(Z1(J+1))-SIN(Z1(J))))	HRA2
370	CONTINUE	WLP
375	CONTINUE	WLP
	C	INIT 170
	C MATERIAL ADDRESSES	INIT 171
380	IF(P02) 405, 385, 405	INIT 172
385	SK7=0.	INIT 173
	DO 400 IIG=1,IGM	INIT 174
	IF(S02-1) 395, 390, 395	INIT 175
390	K6(IIG)=K7(IIG)/S03	INIT 176
	GO TO 400	INIT 177
395	K6(IIG)=K7(IIG)	INIT 178
400	SK7=SK7+K7(IIG)	INIT 179
405	CONTINUE	INIT 180
	C	INIT 181
	C FISSION NEUTRONS	INIT 182
	T11=E1(IGP)	INIT 183
410	CALL CLEAR(0.0,F0,IMJM)	INIT 184
	DO 425 IIG=1,IGM	INIT 185
	E0(IIG) = .0	INIT 186
	READ (NFLUX1) (NO(I),I=1,IMJM)	INIT 187
	READ (MCR1) ((CO(I,J),I=1,ITL),J=1,MT)	INIT 188
	DO 425 J = 1, JM	INIT 189
	DO 425 K = 1, IM	INIT 190
	I = K + (J-1)*IM	INIT 191

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 35

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ITEMP=NO(1) INIT 192
ITEMP=M2(ITEMP) INIT 193
EO(IIG) = EO(IIG) + VO(K,J)*NO(1)*CO(IHT-3,ITEMP) HRA2 194
IF(AO2) 415, 420, 415 INIT 195
415 FO(1)=FO(1)+K7(IIG)*NO(1) INIT 196
GO TO 425 INIT 197
420 FO(1)=FO(1)+CO(IHT-1,ITEMP)*NO(1) INIT 198
425 CONTINUE INIT 199
REWIND NFLUX1 INIT 200
REWIND NCR1 INIT 201
RETURN INIT 202
END INIT 203
SUBROUTINE INNER(N0, N2, CXS, S2, M0, M2, V0, CO, JIM, JJM, JTL, INNR 2
1 CXR, CXT, HA, PA, XR, XD) HRA2
INCLUDE 'ABC.FOR'

DIMENSION NO(1), N2(1), CXS(JIM, JJM, 3), S2(1), M0(1), M2(1), INNR 6
1 V0(1), CO(JTL, 1), CXR(1), CXT(1), HA(1), PA(1) INNR 7
DIMENSION XD(50), XR(50) HRA2
CALL IFLUXN (N2, CO, V0, CXS, M0, M2, ITL, IM, JM, CXR, CXT, XR, XD) HRA2 9
2 DO 4 I=1, IMJM INNR 18
4 NO(1) = M2(1) INNR 19
C BEGIN FLUX CALCULATION INNR 20
IKB = IM - 1 INNR 21
JKB = JM - 1 INNR 22
C FLUX CALCULATION USING SOR WITH LINE INVERSION INNR 23
C INNR 24
C CALCULATION OF LEFT BOUNDARY FLUX INNR 25
KI = 1 INNR 26
KJ = 1 INNR 27
I = KI + (KJ - 1)*IM INNR 28
HA(KJ) = CXS(KI, KJ+1, 2)/CXS(KI, KJ, 3) INNR 29
PA(KJ) = (S2(I) + CXS(KI+1, KJ, 1)*N2(I+1))/CXS(KI, KJ, 3) INNR 30
DO 5 KJ = 2, JKB INNR 31
I = KI + (KJ - 1)*IM INNR 32
HA(KJ) = CXS(KI, KJ+1, 2)/(CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1)) INNR 33
5 PA(KJ) = (S2(I) + CXS(KI+1, KJ, 1)*N2(I+1) + CXS(KI, KJ, 2)*PA(KJ-1))/INNR 34
1 (CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1)) INNR 35

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TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 36

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KJ = JM                                INNRR 36
I = KI + (KJ - 1)*IM                    INNRR 37
N2(I) = (S2(I) + CXS(KI+1,KJ,1)*N2(I+1) + CXS(KI,KJ,2)*PA(KJ-1))/INNRR 38
1 (CXS(KI,KJ,3) - CXS(KI,KJ,2)*HA(KJ-1)) INNRR 39
DO 10 KJ = 2, JM                          INNRR 40
KJ = JM - KJJ + 1                          INNRR 41
I = KI + (KJ - 1)*IM                    INNRR 42
10 N2(I) = PA(KJ) + HA(KJ) * N2(I+IM)      INNRR 43
DO 15 KJ = 1, JM                          INNRR 44
I = KI + (KJ - 1)*IM                    INNRR 45
N2(I)=NO(I)+ORF*(N2(I)-NO(I))           INNRR 46
15 IF(N2(I).LE.0)N2(I)=ABS(NO(I)+(N2(I)-NO(I))/ORF) INNRR 47
C PRINCIPAL FLUX LOOP                      INNRR 48
DO 40 KI = 2, IKB                          INNRR 49
KJ = 1                                      INNRR 50
I = KI + (KJ - 1)*IM                    INNRR 51
HA(KJ)= CXS(KI,KJ+1,2)/CXS(KI,KJ,3)      INNRR 52
PA(KJ)= (S2(I) + CXS(KI,KJ,1)*N2(I-1) + CXS(KI+1,KJ,1)*N2(I+1))/INNRR 53
1 CXS(KI,KJ,3)                             INNRR 54
DO 25 KJ = 2, JKB                          INNRR 55
I = KI + (KJ - 1)*IM                    INNRR 56
HA(KJ) = CXS(KI,KJ+1,2)/(CXS(KI,KJ,3)- CXS(KI,KJ,2)*HA(KJ-1)) INNRR 57
PA(KJ) = (S2(I) + CXS(KI,KJ,1)*N2(I-1) + CXS(KI+1,KJ,1)*N2(I+1) + INNRR 58
1 CXS(KI,KJ,2)*PA(KJ-1))/(CXS(KI,KJ,3) - CXS(KI,KJ,2)*HA(KJ-1)) INNRR 59
KJ = JM                                    INNRR 60
I = KI + (KJ - 1)*IM                    INNRR 61
N2(I) = (S2(I) + CXS(KI,KJ,1)*N2(I-1) + CXS(KI+1,KJ,1)*N2(I+1) + INNRR 62
1 CXS(KI,KJ,2)*PA(KJ-1))/(CXS(KI,KJ,3) - CXS(KI,KJ,2)*HA(KJ-1)) INNRR 63
DO 30 KJ = 2, JM                          INNRR 64
KJ = JM - KJJ + 1                          INNRR 65
I = KI + (KJ - 1)*IM                    INNRR 66
10 N2(I) = PA(KJ) + HA(KJ) * N2(I+IM)      INNRR 67
DO 35 KJ = 1, JM                          INNRR 68
I = KI + (KJ - 1)*IM                    INNRR 69
N2(I)=NO(I)+ORF*(N2(I)-NO(I))           INNRR 70
35 IF(N2(I).LE.0)N2(I)=ABS(NO(I)+(N2(I)-NO(I))/ORF) INNRR 71
40 CONTINUE                                INNRR 72
C CALCULATION OF RIGHT BOUNDARY FLUX      INNRR 73

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TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 37

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KI = IM                                INNRR 74
KJ = 1                                  INNRR 75
I = KI + (KJ - 1)*IM                    INNRR 76
HA(KJ)= CXS(KI,KJ+1,2)/CXS(KI,KJ,3)     INNRR 77
PA(KJ)= (S2(I) + CXS(KI,KJ,1)*N2(I-1))/CXS(KI,KJ,3) INNRR 78
DO 45 KJ = 2,JKB                          INNRR 79
I = KI + (KJ - 1)*IM                    INNRR 80
HA(KJ) = CXS(KI,KJ+1,2)/(CXS(KI,KJ,3)- CXS(KI,KJ,2)*HA(KJ-1)) INNRR 81
45 PA(KJ) = (S2(I) + CXS(KI,KJ,1)*N2(I-1) + CXS(KI,KJ,2)*PA(KJ-1))/ INNRR 82
1 (CXS(KI,KJ,3) - CXS(KI,KJ,2)*HA(KJ-1)) INNRR 83
KJ = JM                                  INNRR 84
I = KI + (KJ - 1)*IM                    INNRR 85
N2(I) = (S2(I) + CXS(KI,KJ,1)*N2(I-1) + CXS(KI,KJ,2)*PA(KJ-1))/ INNRR 86
1 (CXS(KI,KJ,3) - CXS(KI,KJ,2)*HA(KJ-1)) INNRR 87
DO 50 KJJ = 2,JM                          INNRR 88
KJ = JM - KJJ + 1                        INNRR 89
I = KI + (KJ - 1)*IM                    INNRR 90
50 N2(I) = PA(KJ) + HA(KJ) * N2(I+1M)     INNRR 91
DO 55 KJ = 1,JM                          INNRR 92
I = KI + (KJ - 1)*IM                    INNRR 93
N2(I)=NO(I)+ORF*(N2(I)-NO(I))           INNRR 94
55 IF(N2(I).LE.0)N2(I)=ABS(NO(I)+(N2(I)-NO(I))/ORF) INNRR 95
TEMP1 = .0                                INNRR 96
DO 90 I = 1,INJM                          INNRR 97
TEMP2 = ABS (1.0 - NO(I)/N2(I))          INNRR 98
IF (TEMP1 - TEMP2) 80,90,90              INNRR 100
80 TEMP1 = TEMP2                          INNRR 101
90 CONTINUE                               INNRR 103
C                                          INNRR 104
C INNER ITERATION CONTROL                INNRR 105
133 LC = LC + 1                           INNRR 106
II = II + 1                               INNRR 107
IF (II - G07) 533, 1033, 1033           INNRR 117
533 IF(TEMP1-EPS) 633,633,2              INNRR 118
633 IF(G06) 733, 1033, 733              INNRR 119
733 IF(TEMP1-G06)1033,1033,2            INNRR 120
C1033 CONTINUE
1033 WRITE(NOUT,213)                      HRA2

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TL1ST2.FOR                Wednesday, February 21, 1990 12:31 pm                Page 38

213  FORMAT('          GROUP   IN. IT.   ')          HRA2
WRITE(NOUT,2133)IGV,II          HRA2
2133 FORMAT('          ',I3,'          ',I4)          HRA2
IF(104.NE.0)GO TO 1133          HRA2
CALL IFLUXL (N2,CO,VO,CXS,M0, M2, ITL, IM, JM, CXR,CXT,XR,XD) HRA2
1133 CONTINUE                    HRA2
RETURN                            INNR 125
END                                INNR 129
SUBROUTINE INNER1(M0, M2, CXS, VO, CO, AO, Z5, R5, R4, Z4, A1, INN1 2
2                                JIM,JJM,JTL,CXR,CXT,XR,XD,R1,Z1) HRA2 3
DIMENSION XR(50),XD(50)          HRA2
DIMENSION M0(1), M2(1),CXS(JIM,JJM,3),VO(1), CO(JTL,1), INN1 4
1                                AO(1), Z5(1), R5(1), R4(1), Z4(1), A1(1),CXR(1), CXT(1)INN1 5
2                                ,R1(1),Z1(1)          HRA2
INCLUDE 'ABC.FOR'

C  THIS SUBROUTINE CALCULATES COEFFICIENTS FOR THE FLUX EQUATION INN1 7
P12 = 6.28318                    INN1 8
DO 45 KJ = 1, JM                 INN1 9
DO 45 KI = 1, IM                 INN1 10
TEMPA=AO(KI)                     HRA2
TEMPB=AO(IP)                     HRA2
TEMPC=A1(KI)                     HRA2
GO TO (10,10, 5), IGEP          INN1 1#
5  TEMP = (Z4(KJ) - Z4(KJ-1))*R4(KI) HRA2 12
AO(KI)=AO(KI)*(1.0+R1(KI)/B1GR*(SIN(Z1(KJ+1))-SIN(Z1(KJ))))/Z5(KJ)HRA2
AO(IP)=AO(IP)*(1.0+R1(IP)/B1GR*(SIN(Z1(KJ+1))-SIN(Z1(KJ))))/Z5(KJ)HRA2
A1(KI)=A1(KI)*(1.0+COS(Z4(KJ))/B1GR) HRA2
GO TO 15                          INN1 13
10  TEMP = Z4(KJ) - Z4(KJ-1)      INN1 14
15  I = KI + (KJ-1)*IM           INN1 15
ITEMP = M0(I)                    INN1 16
ITEMP = M2(ITEMP)                INN1 17
CXS(KI,KJ,3) = VO(1)*(CO(IHT,ITEMP) - CO(IHS,ITEMP)) INN1 18
IF(I - 1) 44,44,18              HRA2 19
18  ITEMPI = M0(I-1)            INN1 20
ITEMPI = M2(ITEMPI)             INN1 21
IF (ITEMP - ITEMPI) 25,20,25    INN1 22

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 39

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20  CXS(KI,KJ,1)=AO(KI  )*Z5(KJ)/(3.*CO(IHT,ITEMP)*(R4(KI)-R4(KI-1))) INN1 23
    GO TO 30 INN1 24
25  CXS(KI,KJ,1) = AO(KI)*Z5(KJ)*(R5(KI-1)+R5(KI))/((R4(KI)-R4(KI-1))*INN1 25
1  (3.*(R5(KI-1)*CO(IHT,ITEMP1) + R5(KI)*CO(IHT,ITEMP)))) INN1 26
30  IF(I - IM) 44,44,32 INN1 27
32  ITEMPS = M0(I - IM) INN1 28
    ITEMPS = M2(ITEMPS) INN1 29
    IF (ITEMP - ITEMPS) 40,35,40 INN1 30
35  CXS(KI,KJ,2) = A1(KI)/(3.*CO(IHT,ITEMP)*TEMP) INN1 31
    GO TO 44 HRA2 32
40  CXS(KI,KJ,2) = A1(KI)*(Z5(KJ-1) + Z5(KJ))/(TEMP* INN1 33
1  (3.*(Z5(KJ-1)*CO(IHT,ITEMPS) + Z5(KJ)*CO(IHT,ITEMP)))) INN1 34
44  AO(KI)=TEMPA HRA1
    AO(IP)=TEMPB HRA1
    A1(KI)=TEMPC HRA1
45  CONTINUE HRA1 35
    DO 200 KJ = 1, JM INN1 36
    DO 200 KI = 1, IM INN1 37
        TEMPA=AO(KI) HRA2
        TEMPB=AO(IP) HRA2
        TEMPC=A1(KI) HRA2
        GO TO (55,55,50), IGEP INN1 38
50  TEMP = .5*Z5(KJ)*R4(KI) HRA2 39
    AO(KI)=AO(KI)*(1.0+R1(KI)/BIGR*(SIN(Z1(KJ+1))-SIN(Z1(KJ)))/Z5(KJ))HRA2
    AO(IP)=AO(IP)*(1.0+R1(IP)/BIGR*(SIN(Z1(KJ+1))-SIN(Z1(KJ)))/Z5(KJ))HRA2
    A1(KI)=A1(KI)*(1.0+COS(Z4(KJ))/BIGR) HRA2
    GO TO 60 INN1 40
55  TEMP = .5*Z5(KJ) INN1 41
60  I = KI + (KJ-1)*IM INN1 42
    ITEMPS = M0(I) INN1 43
    ITEMPS = M2(ITEMPS) INN1 44
    TEMP1 = CXS(KI+1,KJ,1) INN1 45
    TEMP2 = CXS(KI,KJ+1,2) INN1 46
    IF(KJ - 1) 65,65,100 INN1 47
65  IF(B04 - 1) 90,95,70 INN1 48
70  GO TO ( 80, 80, 75), IGEP INN1 48
75  TEMP3 = P12*.5*(Z5(KJ) + Z5(JM)) HRA2 50
    GO TO 85 INN1 51

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 40

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80  TEMP3 = .5*(Z5(KJ) + Z5(JM))          INN1 52
85  ITEMP3 = I + IM*(JM - 1)             INN1 53
    ITEMP3 = M0(ITEMP3)                  INN1 54
    ITEMP3 = M2(ITEMP3)                  INN1 55
    CXS(K1,KJ,2) = A1(K1)*(Z5(JM) + Z5(KJ))/(TEMP3*
1    (3.*(Z5(JM)*CO(IHT,ITEMP3) + Z5(KJ)*CO(IHT,ITEMP3))) INN1 56
    GO TO 125                             INN1 58
90  CXS(K1,KJ,2) = A1(K1)/(3.*CO(IHT,ITEMP3)*(TEMP +.71/
1  CO(IHT,ITEMP3)))                      INN1 59
    GO TO 125                             INN1 61
95  CXS(K1,KJ,2) = .0                    INN1 62
    GO TO 125                             INN1 63
100 IF (KJ - JM) 125,105,105             INN1 64
105 IF (B03 - 1) 115,120,110            INN1 65
110 TEMP2 = CXS(K1,1,2)                  INN1 66
    CXT(K1) = TEMP2                       INN1 67
    GO TO 125                             INN1 68
115 TEMP2 = A1(K1)/(3.*CO(IHT,ITEMP3)*(TEMP +.71/
1  CO(IHT,ITEMP3)))                      INN1 69
    CXT(K1) = TEMP2                       INN1 71
    GO TO 125                             INN1 72
120 TEMP2 = .0                           INN1 73
    CXT(K1) = TEMP2                       INN1 74
125 IF (K1 - 1) 130,130,145             INN1 75
130 IF (B01) 135,135,140                INN1 76
135 CXS(K1,KJ,1) = A0(K1) *Z5(KJ)/(3.*CO(IHT,ITEMP3)*
1 (.5*R5(K1) + .71/CO(IHT,ITEMP3))) INN1 77
    GO TO 165                             INN1 79
140 CXS(K1,KJ,1) = .0                    INN1 80
    GO TO 165                             INN1 81
145 IF (K1 - IM) 165,150,150            INN1 82
150 IF (B02) 155,155,160                INN1 83
155 TEMP1 = A0(K1+1)*Z5(KJ)/(3.*CO(IHT,ITEMP3)*
1 (.5*R5(K1) + .71/CO(IHT,ITEMP3))) INN1 84
    CXR(KJ) = TEMP1                       INN1 86
    GO TO 165                             INN1 87
160 TEMP1 = .0                           INN1 88
    CXR(KJ) = TEMP1                       INN1 89

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TLIST2.FOR                               Wednesday, February 21, 1990 12:31 pm           Page 41

165  CXS(KI,KJ,3) = CXS(KI,KJ,3) + CXS(KI,KJ,1) + CXS(KI,KJ,2)      INN1 90
1   + TEMP1 + TEMP2                                                INN1 91
   AO(KI)=TEMPA                                                    HRA1
   AO(IP)=TEMPS                                                    HRA1
   A1(KI)=TEMPC                                                    HRA1
200  CONTINUE                                                       INN1 92
   RETURN                                                         INN1 93
   END                                                            INN1 94
   SUBROUTINE INNERZ(MO, M2, CXS, S2, MO, M2, VO, CO, JIM, JJM, JTL,
1   CXR, CXT, HA, PA)                                             INN2 3
   INCLUDE 'ABC.FOR'

   DIMENSION NO(1), N2(1), CXS(JIM, JJM, 3), S2(1), MO(1), M2(1),
1   VO(1), CO(JTL, 1), CXR(1), CXT(1), HA(1), PA(1)             INN2 6
   CALL IFLUXN (M2, CO, VO, CXS, MO, M2, ITL, IM, JM, CXR, CXT, XR, XD) INN2 7
2   DO 4 I=1, IMJM                                                INN2 8
4   NO(I) = N2(I)                                                 INN2 19
C   BEGIN FLUX CALCULATION                                         INN2 20
   IKB = IM - 1                                                  INN2 21
   JKB = JM - 1                                                  INN2 22
C   FLUX CALCULATION USING SOR WITH LINE INVERSION                INN2 23
C                                                                 INN2 24
C   CALCULATION OF BOTTOM BOUNDARY FLUX                            INN2 25
   KI = 1                                                         INN2 26
   KJ = 1                                                         INN2 27
   I = KI + (KJ - 1)*IM                                          INN2 28
   HA(KI) = CXS(KI+1, KJ, 1) / CXS(KI, KJ, 3)                   INN2 29
   PA(KI) = (S2(I) + CXS(KI, KJ+1, 2) * N2(I+IM)) / CXS(KI, KJ, 3) INN2 30
   DO 5 KI = 2, IKB                                              INN2 31
   I = KI + (KJ - 1)*IM                                          INN2 32
   HA(KI) = CXS(KI+1, KJ, 1) / (CXS(KI, KJ, 3) - CXS(KI, KJ, 1) * HA(KI-1)) INN2 33
5   PA(KI) = (S2(I) + CXS(KI, KJ+1, 2) * N2(I+IM) + CXS(KI, KJ, 1) * PA(KI-1)) / INN2 34
1   (CXS(KI, KJ, 3) - CXS(KI, KJ, 1) * HA(KI-1))                 INN2 35
   KI = IM                                                         INN2 36
   I = KI + (KJ - 1)*IM                                          INN2 37
   N2(I) = (S2(I) + CXS(KI, KJ+1, 2) * N2(I+IM) + CXS(KI, KJ, 1) * PA(KI-1)) / INN2 38
1   (CXS(KI, KJ, 3) - CXS(KI, KJ, 1) * HA(KI-1))                 INN2 39
   DO 10 KII = 2, IM                                             INN2 40

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TLISTZ.FOR

Wednesday, February 21, 1990 12:31 pm

Page 42

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KI = IM - KII + 1                                INN2 41
I = KI + (KJ - 1)*IM                             INN2 42
10 N2(I) = PA(KI) + HA(KI) * N2(I+1)              INN2 43
DO 15 KI = 1,IM                                  INN2 44
I = KI + (KJ - 1)*IM                             INN2 45
N2(I)=NO(I)+ORF*(N2(I)-NO(I))                   INN2 46
15 IF(N2(I).LE.0)N2(I)=ABS(NO(I)+(N2(I)-NO(I))/ORF) INN2 47
C PRINCIPAL FLUX LOOP                            INN2 48
DO 40 KJ = 2,JKB                                 INN2 49
KI = 1                                           INN2 50
I = KI + (KJ - 1)*IM                             INN2 51
HA(KI) = CXS(KI+1,KJ,1)/CXS(KI,KJ,3)            INN2 52
PA(KI) = (S2(I) + CXS(KI,KJ,2)*N2(I-IM) + CXS(KI,KJ+1,2)*N2(I+IM))/ INN2 53
1 CXS(KI,KJ,3)
DO 25 KI = 2,IKB                                 INN2 54
I = KI + (KJ - 1)*IM                             INN2 56
HA(KI) = CXS(KI+1,KJ,1)/(CXS(KI,KJ,3) - CXS(KI,KJ,1)*HA(KI-1)) INN2 57
25 PA(KI) = (S2(I) + CXS(KI,KJ,2)*N2(I-IM) + CXS(KI,KJ+1,2)*N2(I+IM) + INN2 58
1 CXS(KI,KJ,1)*PA(KI-1))/(CXS(KI,KJ,3) - CXS(KI,KJ,1)*HA(KI-1)) INN2 59
KI = IM                                          INN2 60
I = KI + (KJ - 1)*IM                             INN2 61
N2(I) = (S2(I) + CXS(KI,KJ,2)*N2(I-1-IM) + CXS(KI,KJ+1,2)*N2(I+1-IM) + INN2 62
1 CXS(KI,KJ,1)*PA(KI-1))/(CXS(KI,KJ,3) - CXS(KI,KJ,1)*HA(KI-1)) INN2 63
DO 30 KII = 2,IM                                 INN2 64
KI = IM - KII + 1                                INN2 65
I = KI + (KJ - 1)*IM                             INN2 66
30 N2(I) = PA(KI) + HA(KI) * N2(I+1)            INN2 67
DO 35 KI = 1,IM                                  INN2 68
I = KI + (KJ - 1)*IM                             INN2 69
N2(I)=NO(I)+ORF*(N2(I)-NO(I))                   INN2 70
35 IF(N2(I).LE.0)N2(I)=ABS(NO(I)+(N2(I)-NO(I))/ORF) INN2 71
40 CONTINUE                                      INN2 72
C CALCULATION OF TOP BOUNDARY FLUX               INN2 73
KJ = JM                                          INN2 74
KI = 1                                           INN2 75
I = KI + (KJ - 1)*IM                             INN2 76
HA(KI) = CXS(KI+1,KJ,1)/CXS(KI,KJ,3)            INN2 77
PA(KI) = (S2(I) + CXS(KI,KJ,2)*N2(I-1-IM))/CXS(KI,KJ,3) INN2 78

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 45

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      KJ = JM - KJJ + 1                                INNP 43
      I = KI + (KJ - 1)*IM                             INNP 44
      KI1 = (JM-1)*IM + KI                             INNP 45
10    N2(I) = PA(KJ) + HA(KJ) * N2(I+IM) + N2(I) * N2(KI1) INNP 46
      DO 15 KJ = 1, JM                                  INNP 47
      I = KI + (KJ - 1)*IM                             INNP 48
15    N2(I) = N0(I) + ORF*(N2(I) - N0(I))             INNP 49
C    PRINCIPAL FLUX LOOP                              INNP 50
      DO 40 KI = 2, 1KB                                INNP 51
      KJ = 1                                           INNP 52
      I = KI + (KJ - 1)*IM                             INNP 53
      HA(KJ) = CXS(KI, KJ+1, 2) / CXS(KI, KJ, 3)       INNP 54
      N2(I) = CXS(KI, 1, 2) / CXS(KI, KJ, 3)          INNP 55
      TEMP1 = N2(I)                                    INNP 56
      TEMP = HA(I)                                     INNP 57
      PA(KJ) = (S2(I) + CXS(KI, KJ, 1)*N2(I-1) + CXS(KI+1, KJ, 1)*N2(I+1)) / INNP 58
1    CXS(KI, KJ, 3)                                    INNP 59
      TEMP2 = PA(KJ)                                   INNP 60
      DO 25 KJ = 2, JKB                                INNP 61
      I = KI + (KJ - 1)*IM                             INNP 62
      HA(KJ) = CXS(KI, KJ+1, 2) / (CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1)) INNP 63
      N2(I) = CXS(KI, KJ, 2) * N2(I-IM) /              INNP 64
1    (CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1))      INNP 65
      TEMP1 = TEMP1 + TEMP*N2(I)                       INNP 66
      PA(KJ) = (S2(I) + CXS(KI, KJ, 1)*N2(I-1) + CXS(KI+1, KJ, 1)*N2(I+1) + INNP 67
1    CXS(KI, KJ, 2)*PA(KJ-1)) / (CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1)) INNP 68
      TEMP2 = TEMP2 + TEMP*PA(KJ)                     INNP 69
25    TEMP = TEMP*HA(KJ)                               INNP 70
      KJ = JM                                          INNP 71
      I = KI + (KJ - 1)*IM                             INNP 72
      TEMP1 = (TEMP1 + TEMP)*CXS(KI, 1, 2) + CXS(KI, KJ, 2)*N2(I-IM) INNP 73
      N2(I) = (S2(I) + CXS(KI, KJ, 1)*N2(I-1) + CXS(KI+1, KJ, 1)*N2(I+1) + INNP 74
1    CXS(KI, 1, 2)*TEMP2 +                             INNP 75
1    CXS(KI, KJ, 2)*PA(KJ-1)) / (CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1) - INNP 76
2    TEMP1)                                           INNP 77
      DO 30 KJJ = 2, JM                                INNP 78
      KJ = JM - KJJ + 1                                INNP 79
      I = KI + (KJ - 1)*IM                             INNP 80

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 46

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      KII = (JM-1)*IM + KI
30  N2(I) = PA(KJ) + HA(KJ) * N2(I+IM) + N2(I) * N2(KII)
      DO 35 KJ = 1, JM
      I = KI + (KJ - 1)*IM
35  N2(I) = NO(I) + ORF*(N2(I) - NO(I))
40  CONTINUE
C   CALCULATION OF RIGHT BOUNDARY FLUX
      KI = IM
      KJ = 1
      I = KI + (KJ - 1)*IM
      HA(KJ) = CXS(KI, KJ+1, 2) / CXS(KI, KJ, 3)
      N2(I) = CXS(KI, 1, 2) / CXS(KI, KJ, 3)
      TEMP1 = N2(I)
      TEMP = HA(I)
      PA(KJ) = (S2(I) + CXS(KI, KJ, 1)*N2(I-1)) / CXS(KI, KJ, 3)
      TEMP2 = PA(KJ)
45  DO 45 KJ = 2, JKB
      I = KI + (KJ - 1)*IM
      HA(KJ) = CXS(KI, KJ+1, 2) / (CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1))
      N2(I) = CXS(KI, KJ, 2) * N2(I-IM) /
1      (CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1))
      TEMP1 = TEMP1 + TEMP*N2(I)
      PA(KJ) = (S2(I) + CXS(KI, KJ, 1)*N2(I-1) + CXS(KI, KJ, 2)*PA(KJ-1)) /
1      (CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1))
      TEMP2 = TEMP2 + TEMP*PA(KJ)
45  TEMP = TEMP*HA(KJ)
      KJ = JM
      I = KI + (KJ - 1)*IM
      TEMP1 = (TEMP1 + TEMP)*CXS(KI, 1, 2) + CXS(KI, KJ, 2)*N2(I-IM)
      N2(I) = (S2(I) + CXS(KI, KJ, 1)*N2(I-1) + CXS(KI, KJ, 2)*PA(KJ-1)
1      + CXS(KI, 1, 2)*TEMP2) /
1      (CXS(KI, KJ, 3) - CXS(KI, KJ, 2)*HA(KJ-1) - TEMP1)
      DO 50 KJJ = 2, JM
      KJ = JM - KJJ + 1
      I = KI + (KJ - 1)*IM
      KII = (JM-1)*IM + KI
50  N2(I) = PA(KJ) + HA(KJ) * N2(I+IM) + N2(I) * N2(KII)
      DO 55 KJ = 1, JM

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INNP 81

INNP 82

INNP 83

INNP 84

INNP 85

INNP 86

INNP 87

INNP 88

INNP 89

INNP 90

INNP 91

INNP 92

INNP 93

INNP 94

INNP 95

INNP 96

INNP 97

INNP 98

INNP 99

INNP 100

INNP 101

INNP 102

INNP 103

INNP 104

INNP 105

INNP 106

INNP 107

INNP 108

INNP 109

INNP 110

INNP 111

INNP 112

INNP 113

INNP 114

INNP 115

INNP 116

INNP 117

INNP 118

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 47

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      I = KI + (KJ - 1)*IM                INNP 119
55  N2(1) = NO(1) + ORF*(M2(1) - NO(1))   INNP 120
C    CALCULATION OF ERROR CRITERION      INNP 121
      TEMP1 = .0                          INNP 122
      DO 90 I = 1,IM,JM                   INNP 123
      TEMP2 = ABS (1.0 - NO(I)/M2(1))     INNP 124
      IF (TEMP1 - TEMP2) 80,90,90        INNP 125
80  TEMP1 = TEMP2                         INNP 126
90  CONTINUE                             INNP 128
C                                         INNP 129
C    INNER ITERATION CONTROL            INNP 130
133 LC = LC + 1                          INNP 131
      II = II + 1                         INNP 132
      IF (II - G07) 533, 1033, 1033      INNP 139
533 IF (TEMP1 - EPS) 633,633,2          INNP 140
633 IF(G06) 733, 1033, 733             INNP 141
733 IF (TEMP1 - G06) 1033, 1033, 2     INNP 142
1033 CONTINUE
      RETURN                              INNP 144
      END                                 INNP 145
SUBROUTINE INNERT(M0, M2, CXS, V0, C0, A0, Z5, R5, R4, Z4, A1,
2      JIM,JJM,JTL,CXR,CXT)             INNT 2
      DIMENSION M0(1), M2(1),CXS(JIM,JJM,3),V0(1), C0(JTL,1),
1      A0(1), Z5(1), R5(1), R4(1), Z4(1), A1(1),CXR(1), CXT(1)INNT 5
      INCLUDE 'ABC.FOR'

C    THIS SUBROUTINE CALCULATES COEFFICIENTS FOR TRIANGULAR GEOMETRY INNT 7
DO 45 KJ = 1, JM                        INNT 8
DO 45 KI = 1, IM                        INNT 9
TEMP = KI - 2* (KI/2) - (KJ - 2* (KJ/2)) INNT 10
TEMP = ABS(TEMP)                        INNT 11
I = KI + (KJ-1) *IM                    INNT 12
ITEMP = NO(1)                          INNT 13
ITEMP = M2(ITEMP)                       INNT 14
CXS(KI,KJ,3) = V0(1)*(C0(IHT,ITEMP) - C0(IHS,ITEMP)) INNT 15
IF(I - 1) 45, 45, 18                   INNT 16
18  ITEMPI = NO(I-1)                    INNT 17
      ITEMPI = M2(ITEMPI)                INNT 18

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 48

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      IF(ITEMP - ITEMP1) 25, 20, 25                                INNT 19
20  CXS(KI,KJ,1) = A0(KI)/(2.*CO(IHT,ITEMP))*25(1))              INNT 20
      GO TO 30                                                    INNT 21
25  CXS(KI,KJ,1) = A0(KI)/((CO(IHT,ITEMP1) + CO(IHT,ITEMP))*25(1)) INNT 22
30  IF(I - IM) 45, 45, 32                                         INNT 23
32  ITEMP3 = M0(I - IM)                                           INNT 24
      ITEMP3 = M2(ITEMP3)                                         INNT 25
      IF(ITEMP - ITEMP3) 40, 35, 40                               INNT 26
35  CXS(KI,KJ,2) = A1(KI)*TEMP/(2.* CO(IHT,ITEMP))*25(1))      INNT 27
      GO TO 45                                                    INNT 28
40  CXS(KI,KJ,2) = A1(KI)*TEMP/((CO(IHT,ITEMP3)+CO(IHT,ITEMP))*25(1)) INNT 29
45  CONTINUE                                                      INNT 30
      DO 200 KJ = 1, JM                                           INNT 31
      DO 200 KI = 1, IM                                           INNT 32
      TEMP = KI - 2*(KI/2) - (KJ-2*(KJ/2))                       INNT 33
      TEMP = ABS(TEMP)                                             INNT 34
      I = KI + (KJ-1)*IM                                          INNT 35
      ITEMP = M0(I)                                               INNT 36
      ITEMP = M2(ITEMP)                                           INNT 37
      TEMP1 = CXS(KI+1,KJ,1)                                       INNT 38
      TEMP2 = CXS(KI,KJ+1,2)                                       INNT 39
      IF(KJ-1) 65, 65, 100                                         INNT 40
65  IF(B04-1) 90, 95, 95                                         INNT 41
90  CXS(KI,KJ,2) = A1(KI)*TEMP/(3.*CO(IHT,ITEMP))*25(1)/3. + .71/ INNT 42
      I CO(IHT,ITEMP)))                                           INNT 43
      GO TO 125                                                  INNT 44
95  CXS(KI,KJ,2) = .0                                             INNT 45
      GO TO 125                                                  INNT 46
100 IF(KJ - JM) 125, 105, 105                                     INNT 47
105 IF(B03 - 1) 115, 120, 120                                     INNT 48
115 TEMP = KI - 2*(KI/2) - (KJ + 1 - 2*((KJ+1)/2))              INNT 49
      TEMP = ABS(TEMP)                                             INNT 50
      TEMP2=A1(KI)*TEMP/(( 3.*CO(IHT,ITEMP))*25(1)/3.+ .71/CO(IHT,ITEMP))) INNT 51
      CXT(KI) = TEMP2                                             INNT 52
      GO TO 125                                                  INNT 53
120 TEMP2 = .0                                                    INNT 54
      CXT(KI) = TEMP2                                             INNT 55
125 IF(KI-1) 130, 130, 145                                       INNT 56

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 49

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130 IF(B01) 135, 135, 140                                INNT 57
135 CXS(KI,KJ,1) = A0(KI)/(3.*CO(IHT,ITEMP))*(Z5(1)/3.   INNT 58
1 + .71/CO(IHT,ITEMP))                                INNT 59
GO TO 165                                              INNT 60
140 CXS(KI,KJ,1) = .0                                  INNT 61
GO TO 165                                              INNT 62
145 IF(KI - IM) 165, 150, 150                          INNT 63
150 IF(B02) 155, 155, 160                              INNT 64
155 TEMP1 = A0(KI+1)/(3.*CO(IHT,ITEMP))*(Z5(1)/3.+    INNT 65
.71/CO(IHT,ITEMP))
CXK(KJ) = TEMP1                                       INNT 66
GO TO 165                                              INNT 67
160 TEMP1 = .0                                          INNT 68
CXK(KJ) = TEMP1                                       INNT 69
165 CXS(KI,KJ,3) = CXS(KI,KJ,3) + CXS(KI,KJ,1) + CXS(KI,KJ,2) INNT 70
1 + TEMP1 + TEMP2                                    INNT 71
200 CONTINUE                                           INNT 72
RETURN                                                INNT 73
END                                                    INNT 74
SUBROUTINE INP                                         INP 2
INCLUDE 'ABC.FOR'

COMMON/PACKED/A(50000)
DIMENSION IDUM(25), DUM(12)
C THIS SUBROUTINE CONTROLS THE READING OF ALL INPUT DATA INP 8
NCR1 = 3
NSCRAT = 4                                           INP 10
NINP = 5                                             INP 11
NOUT = 6                                             INP 12
NFLUX1 = 8                                           INP 13
NSORCE = 9                                           INP 14
WRITE(NOUT,9)                                        INP 15
9 FORMAT(1H1)
WRITE(NOUT,10)
10 FORMAT (24X, '***** P 2 D B *****')
READ(NINP,20,END=14) (ID(1),I=1,20)
GO TO 15
14 STOP
15 CONTINUE

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TLIST2.FOR          Wednesday, February 21, 1990 12:31 pm          Page 50

20  FORMAT(20A4)
    WRITE(NOUT,30) (ID(I),I=1,20)
30  FORMAT (1X, 20A4//)
    I=13
    CALL REA12 (' INP', IDUM(1),I)
    I=14
    CALL REA12 (' INP', IDUM(14),I)
    I=6
    CALL REAG2 (' INP', DUM(1),I)
    CALL REAG2 (' INP', DUM(7),I)
    WRITE (NSCRAT) (IDUM(I), I=1,27)
    WRITE (NSCRAT) (DUM(I), I=1,12)
    REWIND NSCRAT
    READ (NSCRAT) A02, I04, S02, IGM,IHT, NXCM, MCR, G07, D05, MAXT,
1  NPRT,M07, NPUN, IGE, ITOR, NACT, IM, JM, IZM, MT, M01, B01,
2  B02, B03, B04,IZ, JZ
    WRITE(NOUT,60)  A02, I04, S02, IGM,IHT, NXCM, MCR
60  FORMAT(
1  16,' A02  0/1=REGULAR CALCULATION/ADJOINT CALCULATION'/
2  16,' I04  EIGENVALUE TYPE (0/1/2/3/4/5=SOURCE/K/ALPHA/CONC/'
3  , 'DELTA/BUCK)'/
5  16,' S02  PARAMETRIC EIGENVALUE TYPE (0/1/2=NONE/K/ALPHA)'/
7  16,' IGM  NUMBER OF GROUPS'/
4  16,' IHT  POSITION OF SIGMA TRANSPORT'/
9  16,' NXCM NUMBER OF DOWNSCATTERING TERMS'/
2  16,' MCR  NUMBER OF MATERIALS FROM CARDS/TAPE (+N/-N)')
    WRITE (NOUT,70) G07, D05, MAXT, NPRT, M07, NPUN
70  FORMAT(
1  16,' G07  INNER ITERATION MAX PER GROUP (IF NEG, ALT DIR)'/
3  16,' D05  MAXIMUM NUMBER OF OUTER ITERATIONS'/
5  16,' MAXT MAXIMUM TIME IN MINUTES (IF 0, NO EFFECT)'/
7  16,' NPRT PRINT OPTION (0/1/2/3=MINI/MIDI/XS/FLUXES)'/
9  16,' M07  FLUX GUESS (0/1=NONE/FLUX FROM TAPE 14)'/
2  16,' NPUN FLUX DUMP (0/1=NONE/FLUX DUMP TO TAPE 16)'/
    WRITE(NOUT,80) IGE, ITOR, NACT, IM, JM, IZM, MT, M01
80  FORMAT(
1  16,' IGE  GEOMETRY (0/1/2/3=X-Y/R-Z/R-THETA/TRIANGULAR)'/
4  16,' ITOR TOROIDAL SPECIFIER (0/1=R-THETA/TOROIDAL)'/

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HRA2

HRA2

HRA2

HRA2

HRA2

HRA2

HRA2

HRA2

HRA2 60

INP 61

HRA2

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 51

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6 16,' NACT ACTIVITIES (0/>0/=NO EFFECT/READ TABLE POSITIONS' HRA2
8,' FOR ACTIVITIES)'/ HRA2
3 16,' IM NUMBER OF RADIAL INTERVALS'/
5 16,' JM NUMBER OF AXIAL INTERVALS'/
7 16,' IZM NUMBER OF ZONES'/
9 16,' MT TOTAL NUMBER OF MATERIALS INCLUDING MIXES'/
2 16,' M01 NUMBER OF MIXTURE SPECIFICATIONS')
WRITE(NOUT,90) B01, B02, B03, B04, IZ, JZ INP 74
90 FORMAT( INP 75
1 16,' B01 LEFT BOUNDARY CONDITION (0/1=VACUUM/REFLECTIVE)'/
3 16,' B02 RIGHT BOUNDARY CONDITION (0/1=VACUUM/REFLECTIVE)'/
5 16,' B03 TOP BOUNDARY CONDITION (0/1/2=VACUUM/REFLECT'/
6,'PERIODIC)'/
7 16,' B04 BOTTOM BOUNDARY CONDITION (0/1/2=VACUUM/REFL'/
8,'PERIODIC)'/
9 16,' IZ NUMBER OF RADIAL ZONES (DELTA OPTION ONLY)'/
2 16,' JZ NUMBER OF AXIAL ZONES (DELTA OPTION ONLY)'/)
READ(NSCRAT) EV, EVM, S03, BUCK, LAL, LAH, EPS, EPSA, G06,
1 POD, ORF, S01 INP 99
REWIND NSCRAT
WRITE(NOUT,110) EV, EVM, S03, BUCK, LAL, LAH INP 101
110 FORMAT( INP 102
1 1X,1PE11.4,' EV FIRST EIGENVALUE GUESS'/
3 1X,1PE11.4,' EVM EIGENVALUE MODIFIER'/
5 1X,1PE11.4,' S03 PARAMETRIC EIGENVALUE'/
7 1X,1PE11.4,' BUCK BUCKLING (CM-2)'/
9 1X,1PE11.4,' LAL LAMBDA LOWER'/
2 1X,1PE11.4,' LAH LAMBDA UPPER'/)
WRITE (NOUT,120) EPS, EPSA, G06, POD, ORF, S01
120 FORMAT( INP 116
1 1X,1PE11.4,' EPS EIGENVALUE CONVERGENCE CRITERION'/
3 1X,1PE11.4,' EPSA PARAMETER CONVERGENCE CRITERION'/
5 1X,1PE11.4,' G06 INNER ITERATION TEST (IF ZERO, NO TEST)'/
7 1X,1PE11.4,' POD PARAMETER OSCILLATION DAMPER'/
9 1X,1PE11.4,' ORF OVER-RELAXATION FACTOR'/
2 1X,1PE11.4,' S01 -N/*N=POWER (MWT)/NEUTRON SOURCE RATE')
IF(ITOR.EQ.1)THEN HRA2
I=1 HRA2

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 52

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CALL RREAG2 (' INP', BIGR,1)          HRA2
WRITE(6,11)BIGR                      HRA2
11  FORMAT(2X,'BIGR=',E13.6)        HRA2
ELSE                                  HRA2
BIGR=1.0E+20                          HRA2
ENDIF
S04=0
IF (G07.LE.0) S04=1
G07=ABS(G07)
MTP=0
IF (MCR .LE. 0) MTP=-MCR
IF (MCR .LE. 0) MCR=0
IF(I2 + J2) 230, 210, 230            IMP 129
210 IF(I04 - 4) 230, 220, 230        IMP 130
220 CALL ERRO2 ('***I04',220,1)      HVX
230 CONTINUE                          IMP 132
IF(S02) 240, 260, 240                IMP 133
240 IF(S03) 260, 250, 260            IMP 134
250 CALL ERRO2 ('***S03',250,1)
260 CONTINUE                          IMP 136
FEF = 200.0                          IMP 137
TSD = FEF*1.602*10.**(-19)           IMP 138
CCCC COMMENT OUT CALL SETTIN ON A VAX HRA2
C   CALL SETTIN (0,0,0,0)             HRA2
KPAGE = 100                          IMP 140
IHS = IHT+1                           HRA 141
ITL = MXCM + IHS                       HRA 142
C   IHT = 4                            HRA 14
3
IZP = IZM + 1                          IMP 144
IP = IH + 1                            IMP 145
JP = JN + 1                            IMP 146
ML = MCR + MTP
IGP = IGM + 1                          IMP 148
IGEP = IGE + 1                         IMP 149
INJM = IMP*JM                          IMP 150
EQ = .0                                 IMP 151
LAP = .0                                 IMP 152

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TL1S2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 53

	LAPP = .0	INP 153
	LAR = 0.0	INP 154
	DAY = 0.0	INP 155
	ALA = .0	INP 156
	LC = 0	INP 157
	PO2 = 0	INP 158
	CVT = 0	INP 159
	CNT = 0	INP 160
	NCON = 0	INP 161
	TO6 = 0	INP 163
	IF(104-4) 310, 300, 310	INP 164
300	TO6 = 1	INP 165
310	CONTINUE	INP 166
	ORFP =1.0*(ORF - 1.0) + 1.0	INP 167
C	COMPUTE DIMENSION POINTERS	
	LHOLN = 1	HVX
	LATW = LHOLN + 2*ML	
	LALAM = LATW + ML	HVX
	LCO = LALAM + ML	
	LNO = LCO + ITL*MT	INP 174
	LN2 = LNO + IMJM	INP 175
	LA0 = LN2 + IMJM	INP 176
	LA1 = LA0 + IP	INP 177
	LFO = LA1 + IM	INP 178
	LF2 = LFO + IMJM	INP 179
	L10 = LF2 + IMJM	INP 180
	L11 = L10 + MD1	INP 181
	L12 = L11 + MD1	INP 182
	L13 = L12 + MD1	INP 183
	LK6 = L13 + MD1	INP 184
	LK7 = LK6 + IGM	INP 185
	LM0 = LK7 + IGM	INP 186
	LM2 = LM0 + IMJM	INP 187
	LR0 = LM2 + IZM	INP 188
	LR1 = LR0 + IP	INP 189
	LR2 = LR1 + IP	INP 190
	LR3 = LR2 + TO6*IM	INP 191
	LR4 = LR3 + TO6*I2	INP 192

TL1S12.FOR

Wednesday, February 21, 1990 12:31 pm

Page 54

	LRS = LR4 + IM	INP 193
	LS2 = LR5 + IM	INP 194
	LVO = LS2 + IMJM	INP 195
	LV7 = LVO + IMJM	INP 196
	LZ0 = LV7 + IGM	INP 197
	LZ1 = LZ0 + JP	INP 198
	LZ2 = LZ1 + JP	INP 199
	LZ3 = LZ2 + JM*TO6	INP 200
	LZ4 = LZ3 + JZ*TO6	INP 201
	LZ5 = LZ4 + JM	INP 202
	LCKS = LZ5 + JM	INP 203
	LVOL = LCKS + IMJM*3	INP 204
	LMASS = LVOL + IZM	INP 205
	LMATN = LMASS + ML*IZM	INP 206
	LNBR = LMATN + ML	INP 207
	LLD = LNBR + ML	INP 208
	LLCN = LLD + ML	INP 209
	LLFN = LLCN + ML*2	INP 210
	LPHIB = LLFN + ML*7	INP 211
	LAXS = LPHIB + IZM	INP 212
C	LT6=LAXS+ML*IZM	HRA
C	LT8=LT6+ML*IZM	HRA
C	LNTWON=LT8+ML*IZM	HRA
CC	LFXS = LAXS + ML*IZM	INP 2
13		
C	LFXS = LNTWON + ML*IZM	HRA 21
3		
	LACT=LAXS+ML*IZM	HRA3
	LACPOS=LACT+ML*IZM*NACT*ML	HRA3
	LFXS=LACPOS+NACT	HRA3
	LMASSP = LFXS + ML*IZM	INP 214
	L CXR = LMASSP + ML*IZM	INP 215
	L CXT = L CXR + JM	INP 216
	LHA = L CXT + IM	INP 217
	LPA = LHA + MAX0(IM, JM)	INP 218
	LGAM=LPA + MAX0(IM, JM)	INP 219
	LAST = LGAM + IZM	
	ITEMP = 1 + ML*(3+IGM*ITL) + IABS(MTP) + MT*ITL	

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 55

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      LAST=MAX0(LAST,ITEMP)
      WRITE(NOUT,316) LAST
315  IF (LAST.GT.50000) CALL ERRO2 ('  INP',315,1)
316  FORMAT(7H0LAST = ,I7)
      DO 317 I=1,ITEMP
317  A(I)=0.0
C    READ CROSS SECTIONS AND WRITE CROSS SECTION TAPE                INP 237
      CALL SB60(A(LNO),A(LCO),ITL,IGM,MT,ML,A(LATW),A(LHOLN),A(LALAM))
      DO 325 I=LCO, LAST                                             WLP

325  A(I) = .0                                                       INP 240
C    READ FLUXES AND WRITE FLUX TAPE                                  INP 241
      CALL SB62(A(LNO), A(LRO), A(LZO))                             INP 242
C    READ EXTERNAL SOURCE                                           INP 243
      IF (104) 328,326,328                                          INP 244
326  CALL SB64 (A(LS2))
328  CONTINUE                                                       INP 246
      WRITE(NOUT,330)                                               INP 247
330  FORMAT(51H0MESH BOUNDARIES (R0/Z0=RADIAL POINTS/AXIAL POINTS)) INP 248
C    READ RADIAL INTERVALS                                          INP 249
      CALL REAG2('  R0',A(LRO),IP)                                  HVX
C    READ AXIAL INTERVALS                                           INP 251
      CALL REAG2('  Z0',A(LZO),JP)                                  HVX
C    READ ZONE NUMBERS                                              INP 253
      IF (NPRT .GT. 1) GO TO 335
      CALL REAR12 ('  INP',A(LMO),IMJM)                             HRA2
      GO TO 345
335  WRITE(NOUT,340)                                               INP 254
340  FORMAT(30H0ZONE NUMBERS BY MESH INTERVAL)                     INP 255
      CALL REAR12 ('  M0',A(LMO),IMJM)                             HRA2
C    READ MATERIAL NUMBERS                                          INP 257
345  WRITE(NOUT,350)                                               INP 258
350  FORMAT(25H0MATERIAL NUMBERS BY ZONE)                          INP 259
      CALL REAR12('  M2',A(LM2),IZM)                               HRA2
      IF(104-5) 351,352,351                                         INP 261
351  IF(BUCK) 352,358,352                                          INP 262
352  WRITE(NOUT,354)                                               INP 263
354  FORMAT(30H0BUCKLING COEFFICIENTS BY ZONE)                    INP 264

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TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 56

	CALL REAG2(' GAM', A(LGAM),IZM)	HVX
358	CONTINUE	INP 266
C	READ FISSION FRACTIONS	INP 267
	WRITE(NOUT,360)	INP 268
360	FORMAT(17HOFISSION SPECTRUM)	INP 269
	CALL REAG2(' K7',A(LK7),1GM)	HVX
C	READ VELOCITIES	INP 271
	IF (104 .EQ. 2 .OR. S02 .EQ. 2) GO TO 365	
	GO TO 375	
365	WRITE(NOUT,370)	INP 272
370	FORMAT(17HONEUTRON VELOCITY)	INP 273
	CALL REAG2(' V7',A(LV7),1GM)	HVX
375	IF(M01) 400, 400, 380	INP 275
380	WRITE(NOUT,390)	INP 276
390	FORMAT ('ONMIXTURE SPECIFICATIONS (10/I1/I2=MIX NUMBER/MAT. NUMBER 1FOR MIX/MAT. DENSITY)')	
	CALL REAR12(' 10',A(L10),M01)	HRA2
	CALL REAR12(' 11',A(L11),M01)	HRA2
	CALL REAG2(' 12',A(L12),M01)	HVX
400	CONTINUE	INP 282
C	CHECK FOR DELTA CALCULATION	INP 283
	IF(104 - 4) 440, 410, 440	INP 284
410	WRITE(NOUT,420)	INP 285
420	FORMAT (1H0,'DELTA OPTION DATA (R2/R3=RADIAL ZONE NUMBERS/ZONE MOD IFIERS)')	
	CALL REAR12(' R2',A(LR2),1M)	HRA2
	CALL REAG2(' R3',A(LR3),1Z)	HVX
	WRITE(NOUT,430)	HRA2
430	FORMAT (1H0,'DELTA OPTION DATA (Z2/Z3=AXIAL ZONE NUMBERS/ZONE MOD IFIERS)')	
	CALL REAR12(' Z2',A(LZ2),JM)	HRA2
	CALL REAG2(' Z3',A(LZ3),JZ)	HVX
440	CONTINUE	INP 292
	IF(NACT.GT.0)THEN	HRA3
	WRITE(NOUT,490)	HRA3
490	FORMAT (1H0,'ACTIVITY POSITION DATA')	HRA3
	CALL REAR12(' ACPOS',A(LACPOS),NACT)	HRA3
	WRITE(NOUT,*)	HRA3

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 57

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ENDIF
CALL MAPR(A(LM0),A(LM2),IM,JM,A(LC0))
RETURN
END
SUBROUTINE INPB(MATN,NBR,LD,LCN,LFN,ALAM,HOLN,JML,12)
INCLUDE 'ABC.FOR'

DIMENSION MATN(1), NBR(1), LD(1),LCN(JML,1),LFN(JML,1), ALAM(1), INPB 4
1 HOLN(JML,1), 12(1)
DIMENSION IDUM(12)
C THIS SUBROUTINE READS AND PRINTS THE BURNUP DATA INPB 6
CALL REA12 (' INPB',ITEMP,1)
DELT = .0
IF (ITEMP .NE. 0) CALL REAG2 (' INPB',DELT,1)
DAY = DAY + DELT INPB 9
CVT = 0 INPB 10
CNT = 0 INPB 11
POZ = 0 INPB 12
ALA = 0.0 INPB 13
LAP = 0.0 INPB 14
LAPP = 0.0 INPB 15
LAR = 0.0 INPB 16
KPAGE = 100 INPB 17
IF(ITEMP) 100, 15, 20
15 NCON = ITEMP INPB 19
GO TO 100 INPB 24
20 NCON = ITEMP INPB 25
DO 40 M = 1, NCON INPB 26
REWIND NSCRAT
CALL REA12 (' INPB',IDUM(1),12)
WRITE (NSCRAT) (IDUM(I),I=1,12)
REWIND NSCRAT
READ (NSCRAT) MATN(N),NBR(N),LD(N),(LCN(N,K),K=1,2),(LFN(N,K), INPB 28
1 K=1,7) INPB 29
REWIND NSCRAT
ITEMP=MATN(N)
CALL REAG2 (' INPB',ALAM(ITEMP),1)
40 CONTINUE

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 59

	FMT2(3)=PICT(N1)	MAR 31
	FMT3(3)=PARK(N1)	MAR 32
	MARK1=MARK(N1)	MAR 33
	NN=1	MAR 34
	NN=60	MAR 35
	IF (NI.EQ.2) NN=40	MAR 36
120	IF (NM.GT.IM) NN=1M	MAR 37
	WRITE (NOUT,190) (ID(1),I=1,20)	MAR 38
	DO 130 JJ=1,JM	MAR 39
	J=JM-JJ+1	MAR 40
130	WRITE (NOUT,FMT1) J,(MO(I,J),I=NN,MM)	MAR 41
	WRITE (NOUT,FMT3) (MARK1,I=NN,MM)	MAR 42
	WRITE (NOUT,FMT2) (I,I=NN,MM)	MAR 43
	WRITE (NOUT,200)	MAR 44
	IF (NM.EQ.IM) GO TO 140	MAR 45
	NN=NN+1	MAR 46
	NN=NN+NN-1	MAR 47
	GO TO 120	MAR 48
140	CONTINUE	MAR 49
	NN=1	MAR 50
	NN=60	MAR 51
	IF (NI.EQ.2) NN=40	MAR 52
	IF (NM.GT.IM) NN=1M	MAR 53
150	WRITE (NOUT,190) (ID(1),I=1,20)	MAR 54
	DO 170 JJ=1,JM	MAR 55
	J=JM-JJ+1	MAR 56
	DO 160 L=NN,MM	MAR 57
	N=MO(L,J)	MAR 58
160	K(L)=IABS(N2(N))	MAR 59
170	WRITE (NOUT,FMT1) J,(K(L),L=NN,MM)	MAR 60
	WRITE (NOUT,FMT3) (MARK1,I=NN,MM)	MAR 61
	WRITE (NOUT,FMT2) (I,I=NN,MM)	MAR 62
	WRITE (NOUT,200)	MAR 63
	IF (NM.EQ.IM) GO TO 180	MAR 64
	NN=NN+1	MAR 65
	NN=NN+NN-1	MAR 66
	GO TO 150	MAR 67
180	RETURN	MAR 68

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 60

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190 FORMAT (1H1,20A4///) MARC 70
200 FORMAT (2H A/2H X/2H 1/2H A/2H L//BH RADIAL) MARC 71
END MARC 72-
SUBROUTINE MARCH(PHIB,MATN,FXS,AXS,VOL,MASS,MASSP,ALAM,LD,LCN, MARC 2
1 LFN,JML,10,11,I2,M2) MARC 3
DIMENSION PHIB(1),MATN(1),FXS(JML,1),AXS(JML,1),VOL(1), MARC 4
1 MASS(JML,1),MASSP(JML,1),ALAM(1),LD(1),LCN(JML,1), MARC 5
2 LFN(JML,1),10(1),11(1),I2(1),M2(1) MARC 6
INCLUDE 'ABC.FOR'

C THIS SUBROUTINE COMPUTES THE TIME DEPENDENT ISOTOPIC CONCENTRATIONMARC 8
TEMP = DELT * 24. * 3600. / 10. MARC 9
TEMP1 = .0 MARC 10
DO 5 KZ = 1,I2M MARC 11
PHIB(KZ) = PHIB(KZ) * 10.**(-24) MARC 12
DO 5 KN = 1,NCOM MARC 13
LN = MATN(KN) MARC 14
5 TEMP1 = TEMP1 + FXS(KN,KZ)*PHIB(KZ)*MASS(LN,KZ)*VOL(KZ) MARC 15
DO 200 KT = 1,10 MARC 16
TEMP3 = .0 MARC 17
DO 20 KZ = 1,I2M MARC 18
DO 20 KN = 1,NCOM MARC 19
LN = MATN(KN) MARC 20
20 MASSP(LN,KZ) = MASS(LN,KZ) MARC 21
DO 100 KZ = 1,I2M MARC 22
DO 50 KKK = 1,5 MARC 23
DO 50 KN = 1,NCOM MARC 24
LN = MATN(KN) MARC 25
CCC WARNING WARNING HRA2
TEMP2=-((MASS(LN,KZ)+MASSP(LN,KZ))*((ALAM(LN)+AXS(KN,KZ)*PHIB(KZ))) MARC 26
CC ONLY USE ABOVE WHEN SIGMA ABS DOES INCLUDE SIGMA FIS HRA2
C TEMP2=-((MASS(LN,KZ)+MASSP(LN,KZ))*((ALAM(LN)+ HRA2
C 1 (AXS(KN,KZ)+FXS(KN,KZ))*PHIB(KZ))) HRA2 2
6
IF (LD(KN)) 30, 30, 28 MARC 27
28 KK = LD(KN) MARC 28
KK = MATN(KK) MARC 29
TEMP2 = TEMP2 + ALAM(KK)*(MASS(KK,KZ) + MASSP(KK,KZ)) MARC 30

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 61

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30 DO 32 K = 1,2                                MARC 31
    KK = LCN(KN,K)                              MARC 32
    KL = MATN(KK)                              MARC 33
    IF (KK) 32,32,31                            MARC 34
CCC WARNING WARNING                            HRA2
31 TEMP2 = TEMP2 + (AXS(KK,KZ) - FXS(KK,KZ))*PHIB(KZ)* MARC 35
CC ONLY USE ABOVE WHEN SIGMA ABS DOES INCLUDE SIGMA FIS HRA2
C31 TEMP2 = TEMP2 + (AXS(KK,KZ) )*PHIB(KZ)*     HRA2 3
5
1 (MASS(KL,KZ) + MASSP(KL,KZ))                MARC 36
32 CONTINUE                                    MARC 37
DO 36 K = 1,7                                  MARC 38
    KK = LFN(KN,K)                              MARC 39
    KL = MATN(KK)                              MARC 40
    IF (KK) 36,36,34                            MARC 41
CCC THE BELOW IS FOR YIELDS OF FISSION PRODUCT POISONS HRA2
C34 YIELD = 1.0
C IF(KL.EQ.1)THEN
C IF(LN.EQ.7)YIELD=0.061
C IF(LN.EQ.8)YIELD=0.003
C IF(LN.EQ.9)YIELD=0.0113
C IF(LN.EQ.11)YIELD=1.0
C ENDDIF
C IF(KL.EQ.6)THEN
C IF(LN.EQ.7)YIELD=0.055
C IF(LN.EQ.8)YIELD=0.000
C IF(LN.EQ.9)YIELD=0.019
C IF(LN.EQ.11)YIELD=1.0
C ENDDIF
C34 TEMP2=TEMP2+YIELD+FXS(KK,KZ)*PHIB(KZ)*(MASS(KL,KZ)+MASSP(KL,KZ)) HRA2 4
2
CCC THE ABOVE IS FOR YIELDS OF FISSION PRODUCT POISONS
34 TEMP2 = TEMP2 + FXS(KK,KZ)*PHIB(KZ)*(MASS(KL,KZ)+MASSP(KL,KZ)) MARC 42
36 CONTINUE                                    MARC 43
50 MASS(LN,KZ) = MASSP(LN,KZ) + .5*TEMP*TEMP2 MARC 44
DO 100 KN = 1,NC00N                            MARC 45
    LN = MATN(KN)                              MARC 46
100 TEMP3 = TEMP3 + FXS(KN,KZ)*PHIB(KZ)*MASS(LN,KZ)*VOL(KZ) MARC 47

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 62

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                IF(TEMP3) 200,200,110
110 DO 120 KZ = 1,12M
120 PHIB(KZ) = PHIB(KZ) * TEMP1/TEMP3
200 CONTINUE
    DO 500 KZ = 1,12M
500 PHIB(KZ) = PHIB(KZ)*10.**(24)
    DO 540 KZ=1,12M
    DO 540 M=1,M01
    IF(10(M) - M2(KZ)) 540,520,540
520 DO 530 KN=1,MCON
    LN = MATN(KN)
    IF(LN - 11(M)) 530,525,530
525 I2(M) = MASS(LN,KZ)
530 CONTINUE
540 CONTINUE
    RETURN
    END
    SUBROUTINE REAG2(HOLL,ARRAY,MCOUNT)
CCC REPLACE ALL FORMAT STATEMENTS WITH * ON A VAX
    COMMON NSORCE,MIMP, NOUT
    DIMENSION ARRAY (1), HOL(80),NE(40),LE(40)
    CHARACTER*6 HOLL
    CHARACTER*1 HOL
    CHARACTER*20 HE
    J=0
    1 READ (MIMP, 10) (HOL(I),I=1,80)
    10 FORMAT (80A1)
    DO 20 L=1,40
    20 LE(L)=0
        I=0
        L=0
    30 L=L+1
    40 I=I+1
        IF(I.LE.80) GO TO 50
        L=L-1
        GO TO 100
    50 IF (HOL(I).EQ.' ') GO TO 60
        IF (HOL(I).EQ.'T') GO TO 55

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MARC 48
MARC 49
MARC 50
MARC 51
MARC 52
MARC 53
MARC 54
MARC 55
MARC 56
MARC 57
MARC 58
MARC 59
MARC 60
MARC 61
MARC 62
MARC 63
MARC 64
REAG2001
HRAZ

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      IF (HOL(I).EQ.'C') GO TO 55
      IF (HOL(I).EQ.'F') GO TO 55
      IF (HOL(I).EQ.'R') GO TO 55
      IF (HOL(I).EQ.'I') GO TO 55
      LE(L)=LE(L)+1
      HE(L)(LE(L);LE(L))=HOL(I)
      GO TO 40
55  IF (LE(L).GT.0) L=L+1
      LE(L)=1
      HE(L)(1:1)= HOL(I)
      IF (HOL(I).EQ.'R' .OR. HOL(I).EQ. 'I') GO TO 30
      GO TO 100
60  IF (LE(L) .EQ. 0) GO TO 40
      GO TO 30
100 LL=L
      L=0
110 L=L+1
      IF (L .GT. LL) GO TO 1
      IF (HE(L)(1:1).EQ.'T') GO TO 150
      IF (HE(L)(1:1).EQ.'C') GO TO 115
      IF (HE(L)(1:1).EQ. 'F') GO TO 120
      IF (HE(L)(1:1).EQ. 'R') GO TO 130
      IF (HE(L)(1:1).EQ. 'I') GO TO 140
      J=J+1
      READ (HE(L)(1:LE(L)),112) ARRAY(J)
C   DECODE (LE(L),112,HE(L)) ARRAY(J)
112 FORMAT (E20.2)
      GO TO 110
C   CYCLE
115 READ (HE(L+1)(1:LE(L+1)),132) J1
C115 DECODE (LE(L+1),132,HE(L+1)) J1
      READ (HE(L+2)(1:LE(L+2)),132) J2
C   DECODE (LE(L+2),132,HE(L+2)) J2
      J0 = J
      DO 119 K1=1,J1
      DO 119 K2=1,J2
      J=J+1
119 ARRAY(J)=ARRAY(J0 - J2 + K2)

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 64

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      L=L+2
      GO TO 110
C     FILL
120  DO 125 JJ=J+1,NCOUNT
125  ARRAY(JJ)=ARRAY(J)
      J=NCOUNT
      GO TO 150
C     REPEAT
130  READ (HE(L+1))(1:LE(L+1)),132) J1
C130 DECODE (LE(L+1),132,HE(L+1)) J1
132  FORMAT (I20)
      READ (HE(L+2))(1:LE(L+2)),112) T1
C     DECODE (LE(L+2),112,HE(L+2)) T1
      DO 135 JJ=J+1,J+J1
135  ARRAY(JJ)=T1
      J=J+J1
      L=L+2
      GO TO 110
C     INTERPOLATE
140  READ (HE(L+1))(1:LE(L+1)),132) J1
C140 DECODE (LE(L+1),132,HE(L+1)) J1
      READ (HE(L+2))(1:LE(L+2)),112) ARRAY(J+J1+1)
C     DECODE (LE(L+2),112,HE(L+2)) ARRAY(J+J1+1)
      T1= (ARRAY(J+J1+1) - ARRAY(J))/(J1+1)
      DO 145 JJ=J+1,J+J1
145  ARRAY(JJ)= ARRAY(JJ-1) + T1
      J=J +J1+1
      L=L+2
      GO TO 110
150  IF (HOLL .EQ. ' INP') GO TO 155
      IF (HOLL .EQ. ' INPB') GO TO 155
      IF (HOLL .EQ. ' SB60') GO TO 155
      WRITE (NOUT,160) HOLL,J,(ARRAY(I),I=1,J)
155  IF (J-NCOUNT) 170,180,170
160  FORMAT (6X,A6,16/(6E12.5))
170  CALL ERRO2 ( ' REAG2',170,1)
180  RETURN
      END

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 65

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SUBROUTINE RREAG2(HOLL,A,NCOUNT)
CCC  REPLACE ALL FORMAT STATEMENTS WITH * ON A VAX
COMMON NSORCE,NINP,NOUT
DIMENSION ARRAY (1), HOL(80),HE(40),LE(40)
CHARACTER*6 HOLL
CHARACTER*1 HOL
CHARACTER*20 HE
J=0
1  READ (NINP, 10) (HOL(I),I=1,80)
10  FORMAT (80A1)
DO 20 L=1,40
20  LE(L)=0
    I=0
    L=0
30  L=L+1
40  I=I+1
    IF(I,LE.80) GO TO 50
    L=L-1
    GO TO 100
50  IF (HOL(I).EQ.' ') GO TO 60
    IF (HOL(I).EQ.'T') GO TO 55
    IF (HOL(I).EQ.'C') GO TO 55
    IF (HOL(I).EQ.'F') GO TO 55
    IF (HOL(I).EQ.'R') GO TO 55
    IF (HOL(I).EQ.'I') GO TO 55
    LE(L)=LE(L)+1
    HE(L)(LE(L):LE(L))=HOL(I)
    GO TO 40
55  IF (LE(L).GT.0) L=L+1
    LE(L)=1
    HE(L)(1:1)= HOL(I)
    IF (HOL(I).EQ.'R' .OR. HOL(I).EQ. 'I') GO TO 30
    GO TO 100
60  IF (LE(L) .EQ. 0) GO TO 40
    GO TO 30
100 LL=L
    L=0
110 L=L+1

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REAG2001
HRA2

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      IF (L .GT. LL) GO TO 1
      IF (HE(L)(1:1).EQ.'T') GO TO 150
      IF (HE(L)(1:1).EQ.'C') GO TO 115
      IF (HE(L)(1:1).EQ.'F') GO TO 120
      IF (HE(L)(1:1).EQ.'R') GO TO 130
      IF (HE(L)(1:1).EQ.'I') GO TO 140
      J=J+1
      READ (HE(L)(1:LE(L)),112) A
C     DECODE (LE(L),112,HE(L)) A
112  FORMAT (E20.2)
      GO TO 110
C     CYCLE
115  READ (HE(L+1)(1:LE(L+1)),132) J1
C115 DECODE (LE(L+1),132,HE(L+1)) J1
      READ (HE(L+2)(1:LE(L+2)),132) J2
C     DECODE (LE(L+2),132,HE(L+2)) J2
      J0 = J
      DO 119 K1=1,J1
      DO 119 K2=1,J2
      J=J+1
119  ARRAY(J)=ARRAY(J0 - J2 + K2)
      L=L+2
      GO TO 110
C     FILL
120  DO 125 JJ=J+1,NCOUNT
125  ARRAY(JJ)=ARRAY(J)
      J=NCOUNT
      GO TO 150
C     REPEAT
130  READ (HE(L+1)(1:LE(L+1)),132) J1
C130 DECODE (LE(L+1),132,HE(L+1)) J1
132  FORMAT (I20)
      READ (HE(L+2)(1:LE(L+2)),112) T1
C     DECODE (LE(L+2),112,HE(L+2)) T1
      DO 135 JJ=J+1,J+J1
135  ARRAY(JJ)=T1
      J=J+J1
      L=L+2

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TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 67

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      GO TO 110
C     INTERPOLATE
140  READ (HE(L+1)(1:LE(L+1)),132) J1
C140 DECODE (LE(L+1),132,HE(L+1)) J1
      READ (HE(L+2)(1:LE(L+2)),112) ARRAY(J+J1+1)
C     DECODE (LE(L+2),112,HE(L+2)) ARRAY(J+J1+1)
      T1= (ARRAY(J+J1+1) - ARRAY(J))/(J1+1)
      DO 145 JJ=J+1,J+J1
145  ARRAY(JJ)= ARRAY(JJ-1) + T1
      J=J +J1+1
      L=L+2
      GO TO 110
150  IF (HOLL .EQ. ' INP') GO TO 155
      IF (HOLL .EQ. ' INPB') GO TO 155
      IF (HOLL .EQ. ' SB60') GO TO 155
      WRITE (NOUT,160) HOLL,J,(ARRAY(I),I=1,J)
155  IF (J-NCOUNT) 170,180,170
160  FORMAT (6X,A6,16/(6E12.5))
170  CALL ERRO2 (' REAG2',170,1)
180  RETURN
      END
      SUBROUTINE REA12(HOLL,IARRAY,NCOUNT)
      COMMON NSORCE, NINP, NOUT
      DIMENSION IARRAY (1), HOL(80),HE(40),LE(40)
      CHARACTER*6 HOLL
      CHARACTER*1 HOL
      CHARACTER*20 HE
      J=0
1   READ (NINP, 10) (HOL(I),I=1,80)
10  FORMAT (80A1)
      DO 20 L=1,40
20  LE(L)=0
      I=0
      L=0
30  L=L+1
40  I=I+1
      IF (HOL(I) .EQ. '/') GO TO 45
      IF(I.LE.80) GO TO 50

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REA12001

TL1ST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 68

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45  L=L-1
    GO TO 100
50  IF (HOL(I).EQ.' ') GO TO 60
    IF (HOL(I).EQ.'T') GO TO 55
    IF (HOL(I).EQ.'F') GO TO 55
    IF (HOL(I).EQ.'R') GO TO 55
    IF (HOL(I).EQ.'I') GO TO 55
    IF (HOL(I) .EQ. 'C') GO TO 55
    LE(L)=LE(L)+1
    HE(L)(LE(L):LE(L))=HOL(I)
    GO TO 40
55  IF (LE(L).GT.0) L=L+1
    LE(L)=1
    HE(L)(1:1)= HOL(I)
    IF (HOL(I) .EQ. 'C') GO TO 30
    IF (HOL(I).EQ.'R' .OR. HOL(I).EQ. 'I') GO TO 30
    GO TO 100
60  IF (LE(L) .EQ. 0) GO TO 40
    GO TO 30
100 LL=L
    L=0
110 L=L+1
    IF (L .GT. LL) GO TO 1
    IF (HE(L)(1:1).EQ.'T') GO TO 150
    IF (HE(L)(1:1).EQ. 'C') GO TO 115
    IF (HE(L)(1:1).EQ. 'F') GO TO 120
    IF (HE(L)(1:1).EQ. 'R') GO TO 130
    IF (HE(L)(1:1).EQ.'I') GO TO 140
    J=J+1
    READ (HE(L)(1:LE(L)),112) IARRAY(J)
C   DECODE (LE(L),112,HE(L))  IARRAY(J)
112 FORMAT (I20)
    GO TO 110
C   CYCLE
115 READ (HE(L+1)(1:LE(L+1)),112) J1
C115 DECODE (LE(L+1),112,HE(L+1)) J1
    READ (HE(L+2)(1:LE(L+2)),112) J2
C   DECODE (LE(L+2),112,HE(L+2)) J2

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 69

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      J0 = J
      DO 119 K1=1,J1
      DO 119 K2=1,J2
      J=J+1
119  IARRAY(J)=IARRAY(J0 - J2 + K2)
      L=L+2
      GO TO 110
C     FILL
120  DO 125 JJ=J+1,NCOUNT
125  IARRAY(JJ)=IARRAY(J)
      J=NCOUNT
      GO TO 150
C     REPEAT
130  READ (HE(L+1):(1:LE(L+1)),112) J1
C130 DECODE (LE(L+1),112,HE(L+1)) J1
      READ (HE(L+2):(1:LE(L+2)),112) I1
C     DECODE (LE(L+2),112,HE(L+2)) I1
      DO 135 JJ=J+1,J+J1
135  IARRAY(JJ)=I1
      J=J+J1
      L=L+2
      GO TO 110
C     INTERPOLATE
140  READ (HE(L+1):(1:LE(L+1)),112) J1
C140 DECODE (LE(L+1),112,HE(L+1)) J1
      READ (HE(L+2):(1:LE(L+2)),112) IARRAY(J+J1+1)
C     DECODE (LE(L+2),112,HE(L+2)) IARRAY(J+J1+1)
      I1= (IARRAY(J+J1+1) - IARRAY(J))/(J1+1)
      DO 145 JJ=J+1,J+J1
145  IARRAY(JJ)= IARRAY(JJ-1) + I1
      J=J +J1+1
      L=L+2
      GO TO 110
150  IF (HOLL.EQ.' INP') GO TO 155
      IF (HOLL .EQ. ' INPB') GO TO 155
      IF (HOLL .EQ. ' SB60') GO TO 155
      WRITE (NOUT,160) HOLL,J,(IARRAY(1),I=1,J)
155  IF (J-NCOUNT) 170,180,170

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 70

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160  FORMAT (6X,A6,I6/(10I6))
170  CALL ERRO2 (' REA12',170,1)
180  RETURN
      END
      SUBROUTINE REAR12(HOLL,MIARRAY,NCOUNT)
                                                    REA12001
      COMMON NSORCE, NINP, NOUT
      DIMENSION MIARRAY (1), HOL(80),HE(40),LE(40)
      CHARACTER*6 HOLL
      CHARACTER*1 HOL
      CHARACTER*20 HE
      J=0
1    READ (NINP, 10) (HOL(I),I=1,80)
10   FORMAT (80A1)
      DO 20 L=1,40
20   LE(L)=0
      I=0
      L=0
30   L=L+1
40   I=I+1
      IF (HOL(I) .EQ. '/') GO TO 45
      IF (I .LE. 80) GO TO 50
45   L=L-1
      GO TO 100
50   IF (HOL(I) .EQ. ' ') GO TO 60
      IF (HOL(I) .EQ. 'T') GO TO 55
      IF (HOL(I) .EQ. 'F') GO TO 55
      IF (HOL(I) .EQ. 'R') GO TO 55
      IF (HOL(I) .EQ. 'I') GO TO 55
      IF (HOL(I) .EQ. 'C') GO TO 55
      LE(L)=LE(L)+1
      HE(L)(LE(L):LE(L))=HOL(I)
      GO TO 40
55   IF (LE(L) .GT. 0) L=L+1
      LE(L)=1
      HE(L)(1:1)= HOL(I)
      IF (HOL(I) .EQ. 'C') GO TO 30
      IF (HOL(I) .EQ. 'R' .OR. HOL(I) .EQ. 'I') GO TO 30
      GO TO 100

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 71

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60  IF (LE(L) .EQ. 0) GO TO 40
    GO TO 30
100  LL=L
    L=0
    L=L+1
110  IF (L .GT. LL) GO TO 1
    IF (HE(L)(1:1).EQ.'T') GO TO 150
    IF (HE(L)(1:1).EQ.'C') GO TO 115
    IF (HE(L)(1:1).EQ.'F') GO TO 120
    IF (HE(L)(1:1).EQ.'R') GO TO 130
    IF (HE(L)(1:1).EQ.'I') GO TO 140
    J=J+1
    READ (HE(L)(1:LE(L)),112) MIARRAY(J)
C    DECODE (LE(L),112,HE(L)) RIARRAY(J)
112  FORMAT (I20)
    GO TO 110
C    CYCLE
115  READ (HE(L+1)(1:LE(L+1)),112) J1
C115 DECODE (LE(L+1),112,HE(L+1)) J1
    READ (HE(L+2)(1:LE(L+2)),112) J2
C    DECODE (LE(L+2),112,HE(L+2)) J2
    J0 = J
    DO 119 K1=1,J1
    DO 119 K2=1,J2
    J=J+1
119  MIARRAY(J)=MIARRAY(J0 - J2 + K2)
    L=L+2
    GO TO 110
C    FILL
120  DO 125 JJ=J+1,NCOUNT
125  MIARRAY(JJ)=MIARRAY(J)
    J=NCOUNT
    GO TO 150
C    REPEAT
130  READ (HE(L+1)(1:LE(L+1)),112) J1
C130 DECODE (LE(L+1),112,HE(L+1)) J1
    READ (HE(L+2)(1:LE(L+2)),112) J1
C    DECODE (LE(L+2),112,HE(L+2)) J1

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 72

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DO 135 JJ=J+1,J+J1
135 MIARRAY(JJ)=I1
    J=J+J1
    L=L+2
    GO TO 110
C INTERPOLATE
140 READ (HE(L+1)(1:LE(L+1)),112) J1
C140 DECODE (LE(L+1),112,HE(L+1)) J1
    READ (HE(L+2)(1:LE(L+2)),112) MIARRAY(J+J1+1)
C DECODE (LE(L+2),112,HE(L+2)) MIARRAY(J+J1+1)
    I1= (MIARRAY(J+J1+1) - MIARRAY(J))/(J1+1)
DO 145 JJ=J+1,J+J1
145 MIARRAY(JJ)= MIARRAY(JJ-1) + I1
    J=J +J1+1
    L=L+2
    GO TO 110
150 IF (HOLL.EQ.' IMP') GO TO 155
    IF (HOLL .EQ. ' IMPB') GO TO 155
    IF (HOLL .EQ. ' SB60') GO TO 155
    WRITE (NOUT,160) HOLL,J,(MIARRAY(1),1=1,J)
155 IF (J-NCOUNT) 170,180,170
160 FORMAT (6X,A6,16/(10I6))
170 CALL ERRO2 (' REA12',170,1)
180 RETURN
    END
SUBROUTINE SB60 (C,CO,JTL,JGM,JMT,JML,ATW,HOLH,ALAM)
    INCLUDE 'ABC.FOR'

    DIMENSION C(JTL,JGM,JMT), CO(JTL,JMT), ATW(1), HOLN(JML,1),ALAM(1)
    DIMENSION XSD(50,50),XSR(50,50),XD(50),XR(50)
    THIS SUBROUTINE READS CROSS SECTIONS, PERFORMS ADJOINT SB60 6
C REVERSALS IF REQUIRED, AND WRITES CROSS SECTION TAPE SB60 7
C WRITE(NOUT,5) (ID(I), I=1,20) SB60 8
5 FORMAT(1H1,20A4 ///) SB60 9
10 WRITE (NOUT, 20 ) SB60 11
20 FORMAT (55H CROSS SECTIONS ARE READ-IN FOR THE FOLLOWING MATERIALSSB60 12
1/) SB60 13
DO 50 I=1,ML

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 73

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      IF (MTP) 40,40,30
30  READ (15) (HOLN(I,K),K=1,2), ATW(I)
      READ(15) ((C(L,IIG,I), L=1,ITL), IIG=1,IGM)
      GO TO 48
      READ (NINP,42) (HOLN(I,K),K=1,2)
42  FORMAT (2A4)
      CALL REAG2 (' S860', ATW(I), 1)
      DO 45 IIG=1,IGM
      CALL REAG2 (' S860',C(I,IIG,I), ITL)
45  CONTINUE
48  WRITE (NOUT,55) I, (HOLN(I,K),K=1,2)
50  CONTINUE
55  FORMAT (I3, 6X, 2A4)
C   CHECK ON CROSS SECTION CONSISTENCY AND ORDER
      IF(MCR) 70,70,90
70  REWIND 15
90  CONTINUE
      DO 140 J=1,ML
      DO 140 I=1,IGM
      G = C(2,I,J) + C(5,I,J)
      DO 110 K = 1, NXCM
      KK = I + K
      M = 5 + K
      IF(KK - IGM) 100, 100, 110
100  G = G + C(M,KK,J)
110  CONTINUE
      XSR(I,J)=G-C(IHS,I,J)
      XSD(I,J)=G-C(IHT-2,I,J)-C(IHS,I,J)
      IF(C(4,I,J).EQ. 0.0) GO TO 130
      G=ABS((G-C(4,I,J))/C(4,I,J))
      IF(G<.0001) 140,130,130
130  WRITE(NOUT,135) J,I,G
135  FORMAT (' CHECK MATERIAL ',I2, 5X, ' GROUP ',I2,G10.4)
140  CONTINUE
C
C   A02=0/1=FLUX CALCULATION/ADJOINT CALCULATION
160  IF(A02) 170, 280, 170
170  DO 190 IIG=1,IGM

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S860 19

S860 21

S860 26

S860 27

S860 28

S860 29

S860 30

S860 31

S860 32

S860 33

S860 34

S860 35

S860 36

S860 37

S860 38

HRA2

HRA2

S860 39

S860 40

S860 41

S860 43

S860 44

S860 45

S860 46

S860 47

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 74

	IGBAR=IGM-IIG+1	S860 48
	DO 180 M=1,ML	
	DO 180 L = IHT-3,IHS	HRA2 50
	TEMP=C(L,IIG,M)	S860 51
	C(L,IIG,M)=C(L,IGBAR,M)	S860 52
180	C(L,IGBAR,M)=TEMP	S860 53
	IF (IGBAR - IIG -1) 200, 200, 190	S860 54
190	CONTINUE	S860 55
200	CONTINUE	S860 56
	KK = ITL - IHS	S860 57
	IF (KK) 280, 280, 210	S860 58
210	CONTINUE	S860 59
	DO 240 M=1,ML	
	DO 240 IIG = 1,IGM	S860 61
	IGBAR = IGM - IIG + 1	S860 62
	DO 240 L = 1,KK	S860 63
	IF (L - IIG) 220, 240, 240	S860 64
220	I = L + IHS	S860 65
	ITEMP = IGBAR + L	S860 66
	IF (IIG - ITEMP) 230, 230, 240	S860 67
230	TEMP = C(I, IIG, M)	S860 68
	C(I,IIG,M) = C(I,ITEMP,M)	S860 69
	C(I,ITEMP,M) = TEMP	S860 70
240	CONTINUE	S860 71
	C WRITE CROSS SECTION TAPE	
280	DO 300 IIG=1,IGM	S860 72
	DO 295 M=1,MT	S860 73
	IF(M .LE. ML) GO TO 288	
	DO 284 L=1,ITL	
284	C0(L,M)=0.0	
	GO TO 295	
288	CONTINUE	
	DO 290 L=1,ITL	S860 75
290	C0(L,M)=C(L,IIG,M)	S860 76
	XR(M)=XSR(IIG,M)	HRA2
	XD(M)=XSD(IIG,M)	HRA2
295	CONTINUE	
	WRITE (11) (XR(M),M=1,MT)	HRA2

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 75

	WRITE (12) (XD(M),M=1,MT)	HRA2
	DO 88 M=1,MT	HRA2
88	CONTINUE	HRA2
300	WRITE (NCR1) ((C0(L,M),L=1,ITL),M=1,MT)	S860 77
	REWIND NCR1	S860 78
	REWIND 11	HRA2
	REWIND 12	HRA2
	RETURN	S860 79
	END	S860 80
	SUBROUTINE SB62(N0,RF,ZF)	S862 2
	INCLUDE 'ABC.FOR'	
	DIMENSION N0(1), RF(1), ZF(1)	S862 4
C	THIS SUBROUTINE READS THE INPUT FLUXES AND PREPARES A FLUX TAPE	S862 5
	WRITE(NOUT,5)	S862 6
5	FORMAT(1H1)	S862 7
C	N07=0/1=NO FLUX INPUT/FLUX FROM TAPE 14	
	DO 1000 IIG = 1, IGM	S862 14
	IF (N07 .GT. 0) GO TO 120	
	DO 59 I=1,IM	
	DO 59 J=1,JM	
	ITEMP= (J-1)*IN + 1	
59	N0(ITEMP) = 1.0	
	GO TO 200	
120	READ(14) (N0(I), I=1, IMJM)	S862 34
200	WRITE(NFLUX1) (N0(I), I=1, IMJM)	S862 35
1000	CONTINUE	S862 36
	REWIND 14	
	REWIND NFLUX1	
	RETURN	S862 40
	END	S862 41
	SUBROUTINE SB64 (S2)	
	INCLUDE 'ABC.FOR'	
	DIMENSION S2(1)	
C	THIS SUBROUTINE READS THE EXTERNAL SOURCE AND PREPARES A SOURCE TAP	S864 6
	DO 50 IIG = 1,IGM	S864 32
	CALL REAG2 (' S0',S2,INJM)	HVX

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 76

	WRITE (NSORCE) (S2(1), I = 1, INJM)	S864 59
50	CONTINUE	S864 60
	REWIND NSORCE	S864 61
	RETURN	S864 62
	END	S864 63
	SUBROUTINE S8830	S883 2
	INCLUDE 'ABC.FOR'	
	INTEGER*2 IT1, IT2, IT3, IT4	
C	MONITOR PRINT	
C	CHANGE GETTIM TO CLOCK(IT1, IT2, IT3) ON A VAX	HRA2
	CALL GETTIM (IT1, IT2, IT3, IT4)	
	TI = FLOAT (3600*IT1 + 60*IT2 + IT3)/60.	
	KPAGE = KPAGE + 1	S883 7
	IF(KPAGE - 40) 220, 160, 160	S883 8
160	KPAGE = 0	S883 9
210	WRITE(NOUT, 213)	S883 10
213	FORMAT (1H1, ' TIME OUTER IN. IT. EIGENVALUE	
	1 EIGENVALUE LAMBDA')	
	WRITE(NOUT, 215)	S883 13
215	FORMAT (' (MINUTES) ITERATIONS PER LOOP SLOPE'/)	
220	WRITE(NOUT, 225) TI, P02, LC, EQ, EV, ALA	S883 16
225	FORMAT (2X, F6.2, 7X, I4, 7X, I4, 6X, 3(E14.7))	
230	P02=P02+1	S883 18
	LC=0	S883 19
	IF(P02-D05)430,430,330	S883 20
330	NGOTO = 1	S883 21
	GO TO 630	S883 22
430	NGOTO = 4	S883 23
630	RETURN	S883 24
	END	S883 25
	SUBROUTINE S8847	S884 2
	INCLUDE 'ABC.FOR'	
C	THIS SUBROUTINE PRINTS THE FINAL NEUTRON BALANCE TABLE	
	E2(1GP) = .0	S884 4
	E3(1GP) = .0	S884 5
	E4(1GP) = .0	S884 6
	E5(1GP) = .0	S884 7

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 77

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E6(IGP) = .0                      S884  8
E7(IGP) = .0                      S884  9
E8(IGP) = .0                      S884 10
E9(IGP) = .0                      S884 11
DO 10 I = 1, IGM                  S884 12
E2(IGP) = E2(IGP) + E2(I)        S884 13
E3(IGP) = E3(IGP) + E3(I)        S884 14
E4(IGP) = E4(IGP) + E4(I)        S884 15
E5(IGP) = E5(IGP) + E5(I)        S884 16
E6(IGP) = E6(IGP) + E6(I)        S884 17
E7(IGP) = E7(IGP) + E7(I)        S884 18
E8(IGP) = E8(IGP) + E8(I)        S884 19
10 E9(IGP) = E9(IGP) + E9(I)      S884 20
WRITE(NOUT,20)                   S884 21
20 FORMAT (1H1, 2BH FINAL NEUTRON BALANCE TABLE///
1 ' GROUP FISSION IN-SCAT OUT-SCAT ABSORB '
2, ' L.L. R.L. T.L. B.L. '///)   S884 26
DO 30 I = 1, IGM                  S884 26
25 FORMAT (14,3X, 1PBE9.2)
30 WRITE(NOUT,25) I,E1(I),E2(I),E3(I),E4(I),E5(I),E6(I),E7(I),
1 E8(I)                            S884 30
WRITE(NOUT,35)                   S884 30
35 FORMAT (1H )                  S884 31
I = IGM + 1                       S884 32
WRITE(NOUT,25) I,E1(I),E2(I),E3(I),E4(I),E5(I),E6(I),E7(I),
1 E8(I)                            S884 33
XX=E1(I)/(E4(I)+E9(I))
WRITE(NOUT,70) XX
70 FORMAT (1H0/5X, 'NEUTRON MULTIPLICATION CONSTANT = ',F10.6)
RETURN                            S884 35
END                                S884 36
SUBROUTINE S8850(F2,M2,R1,Z1,R4,Z4,V7,J1M,JJM,FN2,
1 CO,NO,MO,M2,F0,JTL,JMT)
INCLUDE 'ABC.FOR'

DIMENSION F2(J1M,JJM), M2(J1M,JJM), R1(1), Z1(1), R4(1), Z4(1), S885 5
1 FLUX(6), FN2(1), CO(JTL,JMT), NO(J1M,JJM), MO(J1M,JJM), S885 6
2 M2(1), F0(J1M,JJM),V7(1)        S885 7

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 78

C	SB850 FINAL PRINT	
	ICARD = 1	S885 9
	CALL SB830	S885 11
	CALL SB847	S885 13
	IF (NPRT-1) 90,90,15	
15	J=IP	S885 14
	IF(IP-JP) 30, 30, 20	S885 15
20	J=JP	S885 16
30	WRITE (NOUT, 40) (I,R1(I),R4(I),Z1(I),Z4(I),I=1,J)	S885 17
40	FORMAT (1H1, 16X, 'RADII', 9X, 'AVG RADII', 11X, 'AXII',	
	1 14X, 'AVG AXII'/(14, 4F18.4))	
	J=J+1	S885 20
	IF(IP-JP) 50, 90, 70	S885 21
50	WRITE (NOUT, 60) (I,Z1(I),Z4(I),I=J,JP)	S885 22
60	FORMAT(14,36X,2F18.4)	S885 23
	GO TO 90	S885 24
70	WRITE (NOUT, 80) (I,R1(I),R4(I),I=J,IP)	S885 25
80	FORMAT(14,2F18.4)	S885 26
90	CONTINUE	S885 27
	DO 100 I=1, IM	S885 28
	DO 100 J=1, JM	S885 29
	NO(I,J) = 0.0	S885 30
100	F2(I,J) = 0.0	S885 31
	DO 220 IIG=1, IGM	S885 32
	IF (NPRT .GT.2) WRITE (NOUT, 110) IIG	
110	FORMAT(1H1, 20X,14HFLUX FOR GROUP,13)	S885 34
	READ (NFLUX1)((N2(I,J),I=1,IM),J=1,JM)	S885 35
	READ(NCR1)((CO(I,J), I1 = 1, ITL), J = 1, MT)	S885 36
	DO 120 I=1, IM	S885 37
	DO 120 J=1, JM	S885 38
	NO(I,J) = NO(I,J) + N2(I,J)	S885 39
	ITEMP = NO(I,J)	S885 40
	ITEMP = N2(ITEMP)	S885 41
120	F2(I,J) = F2(I,J) + CO(IHT-3,ITEMP)*N2(I,J)*1000.*TSD	HRA2 42
	IF(NPJUN) 210, 210, 205	S885 43
205	WRITE(16) ((N2(I,J),I=1,IM),J=1,JM)	S885 63
210	IF (NPRT .GT.2) CALL PRT (IM,JM,N2,Z4,NOUT)	
220	CONTINUE	S885 65

TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 79

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      IF (NPRT .LE. 0) GO TO 250
      WRITE(NOUT, 230 )
230  FORMAT(1H1//, 19X,11H TOTAL FLUX//)
      CALL PRT(IM,JM,MO,Z4,NOUT)
      WRITE(NOUT, 240 )
240  FORMAT(1H1//, 19X, 26HPOMER DENSITY (MMT/LITER))
      CALL PRT(IM,JM,F2,Z4,NOUT)
250  CONTINUE
      IF(NPUM - 1) 270, 260, 260
260  END FILE 16
      WRITE(NOUT,265)
265  FORMAT(1H0,50X,'***** FLUXES, ETC. DUMPED TO TAPE *****')
270  REWIND NCR1
      REWIND NFLUX1
      RETURN
      END
      SUBROUTINE TCHEK(LGH,JUMP)
      INTEGER*2 IT1,IT2,IT3,IT4
CC CHANGE GETTIM TO CLOCK(IT1,IT2,IT3) ON A VAX
      CALL GETTIM (IT1,IT2,IT3,IT4)
      ISEC = 3600*IT1 + 60*IT2 + IT3
      IF(ISEC.GT.60*MAXT.AND.MAXT.GT.0) JUMP=1
      RETURN
      TCHE 6
      END
      TCHE 7
CC THIS SUBROUTINE MUST BE USED ON THE VAX
CCC  SUBROUTINE CLOCK (IT1,IT2,IT3)
C    INTEGER*2 IT1,IT2,IT3
C    IT1=0
C    IT2=0
C    IT3=0
C    RETURN
C    END
SUBROUTINE IFLUXL (N2, CO, VO, CXS, MO, M2, JTL,JIM,JJM, CXR, CXT HRA2 2
1  ,XR,XD)
      HRA2
      INCLUDE 'ABC.FOR'

      DIMENSION M2(1), CO(JTL,1), VO(1),CXS(JIM,JJM,3),MO(1), M2(1), IFLU 4
1      CXR(1), CXT(1)
      IFLU 5

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 80

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DIMENSION XD(50),XR(50)                                HRAZ
C THIS SUBROUTINE NORMALIZES FLUXES BEFORE EACH INNER ITERATION  IFLU 6
C ABSORPTION AND OUT-SCATTER                                IFLU 7
E3(IGV) = 0.0                                           IFLU 8
E4(IGV) = 0.0                                           IFLU 9
DO 10 I=1, IMJM                                         IFLU 10
TEMP = VO(I)*N2(1)                                       IFLU 11
ITEMP = NO(I)                                           IFLU 12
ITEMP = N2(ITEMP)                                       IFLU 13
E3(IGV) = E3(IGV) + (XD(ITEMP))*TEMP                    HRAZ 14
10 E4(IGV) = E4(IGV) + CO(1HT-2,ITEMP)*TEMP             IFLU 15
C LEFT LEAKAGE                                           IFLU 16
IF(B01) 20, 20, 40                                       IFLU 17
20 E5(IGV) = 0.0                                         IFLU 18
DO 30 KJ = 1, JM                                         IFLU 19
I = (KJ - 1)*IM + 1                                     IFLU 20
30 E5(IGV) = E5(IGV) + CXS(1,KJ,1)*N2(1)                IFLU 21
GO TO 50                                                 IFLU 22
40 E5(IGV) = .0                                          IFLU 23
C RIGHT LEAKAGE                                          IFLU 24
50 IF(B02) 60, 60, 80                                     IFLU 25
60 E6(IGV) = 0.0                                         IFLU 26
DO 70 KJ = 1, JM                                         IFLU 27
I = KJ*IM                                               IFLU 28
70 E6(IGV) = E6(IGV) + CXR(KJ)*N2(1)                   IFLU 29
GO TO 90                                                 IFLU 30
80 E6(IGV) = 0.0                                         IFLU 31
C TOP LEAKAGE                                           IFLU 32
90 IF(B03-1) 120, 140, 100                               IFLU 33
100 E7(IGV) = .0                                         IFLU 34
DO 110 KI = 1, IM                                       IFLU 35
I = IMJM - IM + KI                                     IFLU 36
110 E7(IGV) = E7(IGV) + CXS(KI,1,2)*(N2(I) - N2(KI))   IFLU 37
E8(IGV) = - E7(IGV)                                     IFLU 38
GO TO 190                                               IFLU 39
120 E7(IGV) = 0.0                                        IFLU 40
DO 130 KI = 1, IM                                       IFLU 41
I = IMJM - IM + KI                                     IFLU 42

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TLIST2.FOR

Wednesday, February 21, 1990 12:31 pm

Page 81

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130  E7(IGV) = E7(IGV) + CXT(KI)*N2(I)          IFLU 43
      GO TO 150                                IFLU 44
140  E7(IGV) = 0.0                             IFLU 45
C    BOTTOM LEAKAGE                           IFLU 46
150  IF(B04) 160, 160, 180                     IFLU 47
160  EB(IGV) = 0.0                             IFLU 48
      DO 170 KI = 1, IM                       IFLU 49
170  EB(IGV) = EB(IGV) + CXS(KI,1,2)*N2(KI)   IFLU 50
      GO TO 190                                IFLU 51
180  EB(IGV) = 0.0                             IFLU 52
190  E9(IGV) = E5(IGV) + E6(IGV) + E7(IGV) + EB(IGV) IFLU 53
      RETURN                                   IFLU 64
      END                                       IFLU 65

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PROGRAM ABC.FOR

COMMON NSORCE

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COMMON NINP,NOUT,MCR1,HFLUX1,NSCRAT,ALA,B07,
1CNT,CVT,DAY,DELTA,E0(S1),E1(S1),E2(S1),E3(S1),E4(S1),E5(S1),
2E6(S1),E7(S1),E8(S1),E9(S1),
3E01,E02,E03,EQ,EVP,EVPP,FEF,GBAR,GLH,IGEP,IGP,IGV,IHS,IHT,II,
4INJN,IP,ITEMP,ITEMP1,ITEMP2,ITL,IZP,JP,K07,KPAGE,LAP,LAPP,LAR,
5LC,ML,NCON,NGOTO,NPRT,ORFP

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COMMON P02,PBAR,SBAR,SK7,T06,T7,T11,TEMP,TEMP1,TEMP2,TEMP3,TEMP4,
7TI,TSO,V11

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COMMON ID(20),MAXT,A02,I04,S02,IGH,NXCM,MCR,MTP,M07,D05,G07,S04,
1NPUN,IGE,IM,JM,IZM,MT,M01,B01,B02,B03,B04,I2,JZ,EV,EVM,S03,BUCK,
2LAL,LAH,EPS,EPSA,G06,POD,ORF,S01,NACT,NFD

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COMMON LATW,LNOLN,LALAM,LCD,LMO,LN2,LA0,LA1,LF0,LF2,L10,L11,L12,
1L13,LK6,LK7,LMO,LN2,LRO,LR1,LR2,LR3,LR4,LR5,LS2,LV0,LV7,L20,L21,
2L22,L23,L24,L25,LCXS,LVOL,LMASS,LMATN,LNBR,LLD,LLCN,LLFN,LPH1B,

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C    3LAXS,LFXS,LMASS,LCCR,LCXT,LHA,LPA,LT6,LT8,Intwon,bigr
3LAXS,LFXS,LMASS,LCCR,LCXT,LHA,LPA,lact,lacpos,bigr
INTEGER A02,B01,B02,B03,B04,B07,CNT,CVT,D05,G07,P02,S02,S04,T06,
1R2,ZZ,GLH
REAL I2,I3,K6,K7,LAH,LAL,LAP,LAPP,LAR,NO,N2,MASS,MASSP

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