ANALYSIS OF OPTIMAL CONTROL OF A

FOUR-GIMBAL SYSTEM

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ANALYSIS OF OPTIMAL CONTROL OF A

FOUR-GIMBAL SYSTEM

by

Michael Andrew Gennert

Submitted to the Department of Electrical Engineering and Computer Science on May 9, 1980 in partial fulfillment of the requirements for the Degrees of Bachelor of Science and Master of Science.

ABSTRACT

This thesis investigates modelling and control of a four-gimbal inertial system. The system under study is used to stabilize an inertial platform and to isolate the platform from vibration and rotation of the vehicle in which the system is mounted.

A few simplifying assumptions are made about the gimbal system. Using these assumptions and Euler's torque equations for a rotating body, a set of linear equations is developed relating angular acceleration of the gimbal elements to torque motor voltage. Taking a state-space approach, a set of nonlinear differential equations is used to compute the orientations of the gimbal elements from the torque motor voltages. A novel approach to the incorporation of static friction is presented, which leads to a simplified set of equations in the presence of static friction. Coulomb friction is also taken into account.

Modern optimal control techniques are applied to a linearized discrete-time version of the state equations to produce an optimal control scheme. The gimbal system and controller are simulated on a digital computer using the FORTRAN programming language. A listing of the program is included in the appendix. Comparisons are made with an earlier control strategy showing the reduction of platform misorientation, reduction of required torque, and elimination of switching transients. THESIS SUPERVISOR: Nils R. Sandell, Jr. TITLE: Adjunct Professor of Electrical Engineering

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I. Introduction

This thesis investigates modelling and control of a four-gimbal system. Gimbals are generally used for precise orientation and/or stabilization. Typical applications include: attachment of a rocket engine so that the engine may be aimed, suspension of a ship's compass in a horizontal position despite pitch and roll, mounting a radar to rapidly track a target, stabilization of an inertial platform and isolation of the platform from vibration. It is this last application that will be of concern to us in this paper.

Inertial guidance and navigation systems generally use gyroscopes and accelerometers as sensing devices. High performance inertial guidance systems usually have these sensors mounted on an inertial platform and a series of concentric gimbals connecting the platform to the case. Gyros on the platform sense rotations of the platform with respect to inertial space, and are used in feedback loops to maintain an inertial reference.

The inertial platform and gimbals are housed in the inertial measurement unit case. The case is rigidly affixed to a vehicle whose rotation rate will be changing with time. The rotation rate is not measured directly; it can be calculated from other quantities, as will be shown. The rotation can be viewed as an input to the gimbal system, uninfluenced by the behavior of the system. As such, the vehicle's motion provides a set of boundary conditions for the kinematic equations describing the behavior of the gimbals.

To fully isolate the inertial platform from vehicle motion requires a minimum of three gimbals, providing three degrees of freedom. It is possible for two of the gimbal axes to become "gimbal lock" is then said to occur and one degree of parallel: freedom is lost. If all three axes lie in one plane, rotation axis perpendicular to this plane is impossible. about an Clearly, gimbal lock must be avoided. However, it is not sufficient that the system stay out of gimbal lock; it must not even get close because, as gimbal lock is approached. increasingly high torque levels are required to keep the platform If the required torque should exceed the maximum inertial[5]. available torque, then the inertial platform may lose its inertial reference.

There are basically two strategies available for dealing with the gimbal lock problem. The simplest solution is to restrict the vehicle's motion so that gimbal lock cannot occur. Early guidance systems used exactly this restricted attitude scheme. The drawbacks are obvious. A present state-of-the-art all-attitude guidance system avoids gimbal lock by adding a fourth gimbal (Figure 1.1). The extra degree of freedom ensures that it will always be possible to avoid gimbal lock. If two gimbal axes are aligned there will still be three degrees of freedom. However, if the system is not properly controlled it is possible for all four axes to lie in one plane, a second degree of freedom will be lost, and gimbal lock will result. The problem then is one of allocation of control among the four gimbals to stabilize the inertial platform while avoiding gimbal lock given the vehicle's rotation rate.

Control is effected through torque motors mounted on the outer three gimbals and the case. The torquers are driven by saturating amplifiers, limiting the maximum available torque. Information on the state of the system is available from three sources. Gyroscope outputs indicate any deviation of the platform attitude from inertial, resolvers mounted on each gimbal indicate the angles between gimbals, and tachometers measure angular velocities.

Presently, the inner two gimbals are driven directly by gyroscopes, and control is switched between the two outer gimbals, depending on the two middle angles. The control law takes the form of decision rules, so that control is allocated based upon the zone in which the middle two angles reside. Although the zone control does avoid gimbal lock, it is not optimal. Large attitude errors and torque transients may occur when switching zones. The maximum torque requirements are excessive; by reducing them it will be possible to improve torque motor performance and/or reduce the torquer size, weight and cost. Furthermore, reductions in attitude errors resulting from optimization will contribute to overall system accuracy.

The approach taken is as follows. The mechanics of the gimbal system are discussed first. Simplifying assumptions and approximations are presented and justified. Based upon Euler's torque equations a set of equations are derived that characterize the system. We examine friction and its effects. Modern optimal control techniques are applied to a linearized discrete-time version of the torque equations to yield an optimal control scheme. Various methods of implementing the controller are suggested. The controller is realized as a simulation on a digital computer using the FORTRAN programming language. Results of the simulation are analyzed and compared with an earlier control strategy.

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Figure 1.1 Gimbal Configurations

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II. Nomenclature

A	Continuous-time dynamics matrix
A#	Discrete-time dynamics matrix
В	Continous-time matrix from control signals to state derivative
B [#]	Discrete-time matrix from control signals to state derivative
С	Case
टै	Continuous-time constant vector
č *	Discrete-time constant vector
с¥	Coordinate transformation from j to k system
D	State information compression matrix
ē	State error vector
E	Elevation gimbal = Inertial platform
Ē	Optimal next state
^{fik} j	Angular momentum of gimbal j in the k frame
I	Inner gimbal
I	3 x 3 identity matrix
J	Cost function
J_k^1	Inertia tensor of gimbal k in the l reference frame
Jkv	Moment of inertia of gimbal k about its v-axis in the k frame
L	Matrix transforming acceleration to torques
M	Middle gimbal
M	Matrix transforming torques to accelerations = Γ'
0	Outer gimbal
Q	Symmetric state weight matrix in cost function
R	Symmetric torque weight matrix in cost function
S	Inertial space

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īţ	Total torque on gimbal j in the k frame					
\vec{T}_{kj}^{l}	Torque on gimbal j supplied by gimbal k in the l frame					
l Tjkv	Component of T in the v direction					
បិ	Control vector					
<u>↓</u> Wkj	Rotation of gimbal j with respect to gimbal k in the l frame					
x	State vector					
Ŷ	Vector composed of torques and torque-like terms					
Ż	Angular acceleration vector					
I	Angle between E and I					
В	Angle between I and M					
0	Angle between M and O					
Ø	Angle between 0 and C					
λ	Gimbal lock angle					

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III. System Description

The four-gimbal system is shown schematically in Figure 3.1. Pictured are the case (C), outer gimbal (O), middle gimbal (M), inner gimbal (I) and elevation gimbal (E). The terms "elevation gimbal" (1) and "inertial platform" refer to the same thing and will be used interchangably. "Case" and "vehicle" will also be used interchangably in the context of rotation and acceleration, although they do not refer to the same thing. The case is securely bolted to the vehicle and thus experiences the same velocity and acceleration.

The outer, middle and inner gimbals look much the same except for size. Two slipring assemblies connect each gimbal to the next innermost and next outermost gimbals. The slipring assemblies contain resolvers, tachometers and torque motors. The relative position and velocity of each gimbal pair may be directly observed (after filtering to remove noise). The torque motors are the sole actuators present in the system.

The elevation gimbal is totally different from the others. It is essentially a platform laden with sensors. The only sensors of concern to us here will be the gyroscopes. The

(1) The phrase "elevation gimbal" is carried over from three-gimbal system days when the elevation angle I was exactly equal to the elevation of the vehicle with respect to the earth's surface. What is now the inner gimbal was then called the "azimuth gimbal." It is still occaisionally referred to by the older name. We will stick with "inner gimbal." The letter "B" used for the angle between the inner and middle gimbals reflects the fact that this angle equalled the bearing of the vehicle in the three-gimbal system. gyroscopes will be treated as though there were three single degree of freedom (SDF) gyros. In fact, two two degree of freedom (TDF) gyros may be used, one degree of freedom being redundant. The gyros are aligned so that their input axes lie along XE, YE and ZE. Any rotation of the inertial platform will be sensed by one or more gyros. Any misalignment of the gyroscopes with respect to the inertial platform will be subject to compensation elsewhere in the guidance system and will not concern us.

Six different Cartesian coordinate systems may be defined. Four of these coordinate systems are fixed to the four gimbals, the fifth and sixth coordinate systems are associated with the case and inertial space (S). One may restate the purpose of the controller as being to keep the elevation gimbal coordinate frame and the inertial space coordinate frame as closely aligned as possible given the rotation rate of the case coordinate frame. The rotation rates of the case and gimbals with respect to inertial space coordinatized in the case and gimbal frames may be defined as follows:

$$\vec{W}_{SC}^{C} \triangleq \begin{vmatrix} W_{CX} \\ W_{CY} \\ W_{CZ} \end{vmatrix} \qquad \vec{W}_{SO}^{O} \triangleq \begin{vmatrix} W_{OX} \\ W_{OY} \\ W_{OZ} \end{vmatrix}$$

. ..

$$\vec{W}_{SM} \triangleq \begin{vmatrix} W_{MX} \\ W_{MY} \\ W_{MZ} \end{vmatrix} \qquad \vec{W}_{SI} \triangleq \begin{vmatrix} W_{IX} \\ W_{IY} \\ W_{IZ} \end{vmatrix} \qquad \vec{W}_{EX} \\ \vec{W}_{EX} \\ W_{EX} \\$$

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The above vectors are interpreted as the rotation rate of the coordinate system denoted by the the right subscript with respect to the coordinate system denoted by the left subscript as seen from the coordinate system denoted by the superscript. This convention is discussed in more detail in Britting[3].

In order to relate the various coordinate frames it is necessary to define the angle between adjacent gimbals.

Angle name	<u>Between</u>	<u>Also</u> called
I	E and I	Elevation Angle
В	I and M	Inner Angle
0	M and O	Middle Angle
Ø	O and C	Outer Angle

That only a single degree of freedom exists between gimbals simplifies the direction cosine matrices. Specifically:

$$C_{C}^{0} = \begin{bmatrix} \cos \emptyset & 0 & -\sin \emptyset \\ 0 & 1 & 0 \\ \sin \emptyset & 0 & \cos \emptyset \end{bmatrix}$$
(3.1)

 $C_0^{\mathsf{M}} = \begin{bmatrix} 1 & 0 & 0 \\ 0 & \cos\theta & \sin\theta \\ 0 & -\sin\theta & \cos\theta \end{bmatrix}$ (3.2)

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$$C_{M}^{I} = \begin{vmatrix} \cos B & \sin B & 0 \\ -\sin B & \cos B & 0 \\ 0 & 0 & 1 \end{vmatrix}$$
(3.3)
$$C_{I}^{E} = \begin{vmatrix} \cos I & 0 & -\sin I \\ 0 & 1 & 0 \\ \sin I & 0 & \cos I \end{vmatrix}$$
(3.4)

The above matrices are interpreted as a linear transform the coordinate system denoted by the subscript to the from coordinate system denoted by the superscript. Direction cosine matrices are treated in more detail in Appendix A. The special form of the direction cosine matrices is due to the fact that $\mathcal D$ is measured around the outer gimbal and case y-axis, Q is measured around the middle and outer z-axis, B is measured around the inner and middle x-axis, and Y is measured around the elevation and inner y-axis. These definitions are entirely arbitrary but survive for historical reasons. The time derivatives of these angles are nothing but the relative rotation rates. That is:

$$\vec{W}_{CO}^{C} = \vec{W}_{CO}^{O} \triangleq \vec{\emptyset} \qquad \vec{W}_{OM}^{O} = \vec{W}_{OM}^{M} \triangleq 0$$

$$\vec{w}_{MI}^{M} = \vec{w}_{MI}^{I} \stackrel{2}{=} \begin{vmatrix} 0 \\ 0 \\ \dot{B} \end{vmatrix} \qquad \vec{w}_{IE}^{I} = \vec{w}_{IE}^{E} \stackrel{2}{=} \begin{vmatrix} 0 \\ \dot{I} \\ 0 \\ 0 \end{vmatrix}$$

It is now possible to relate the rotation of any gimbal to inertial space. This is necessary to express the torque equations later. Starting with the outer gimbal and applying equations (A.2) and (A.6) we have:

$$\vec{W}_{SO}^{O} = \vec{W}_{SC}^{O} + \vec{W}_{CO}^{O} = C_{C}^{O}\vec{W}_{SC}^{C} + \vec{W}_{CO}^{O}$$

$$= \begin{vmatrix} \cos \emptyset & W_{CX} - \sin \emptyset & W_{CZ} \\ & W_{CY} + \dot{\emptyset} & & \\ & \sin \emptyset & W_{CX} + \cos \emptyset & W_{CZ} \end{vmatrix}$$
(3.5)

$$\vec{W}_{SM}^{M} = C_{O}^{M} \vec{W}_{SO}^{O} + \vec{W}_{OM}^{M}$$

$$= \begin{vmatrix} W_{OX} + \dot{\Theta} \\ \cos \Theta W_{OY} + \sin \Theta W_{OZ} \\ -\sin \Theta W_{OY} + \cos \Theta W_{OZ} \end{vmatrix}$$
(3.6)

$$\vec{W}_{SI}^{T} = C_{M}^{I} \vec{W}_{SM}^{M} + \vec{W}_{MI}^{I}$$

$$= \begin{vmatrix} \cos B W_{MX} + \sin B W_{MY} \\ -\sin B W_{MX} + \cos B W_{MY} \\ W_{MZ} + \dot{B} \end{vmatrix}$$
(3.7)

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$$\vec{W}_{SE}^{E} = C_{I}^{E}\vec{W}_{SI}^{I} + \vec{W}_{IE}^{E}$$

$$= \begin{vmatrix} \cos I W_{IX} - \sin I W_{IZ} \\ W_{IY} + \dot{I} \\ \sin I W_{IX} + \cos I W_{IZ} \end{vmatrix}$$
(3.8)

Equations (3.5) through (3.8) and (A.2) may be combined to compute the case rates.

$$\vec{w}_{SE}^{E} = c_{C}^{E}\vec{w}_{SC}^{C} + c_{O}^{E}\vec{w}_{CO}^{O} + c_{M}^{E}\vec{w}_{OM}^{M} + c_{I}^{E}\vec{w}_{MI}^{I} + \vec{w}_{IE}^{E}$$
(3.9)

Rearranging terms and multiplying by $C_{\underline{E}}^{\underline{C}}$ yields:

$$\vec{\mathbf{W}}_{SC}^{C} = c_{E}^{C}\vec{\mathbf{w}}_{SE}^{E} - c_{E}^{C}c_{O}^{E}\vec{\mathbf{w}}_{CO}^{O} - c_{E}^{C}c_{M}^{E}\vec{\mathbf{w}}_{OM}^{M} - c_{E}^{C}c_{I}^{E}\vec{\mathbf{w}}_{MI}^{I} - c_{E}^{C}\vec{\mathbf{w}}_{IE}^{E}$$

$$= c_{E}^{C}\vec{\mathbf{w}}_{SE}^{E} - c_{O}^{C}\vec{\mathbf{w}}_{CO}^{O} - c_{M}^{C}\vec{\mathbf{w}}_{OM}^{M} - c_{I}^{C}\vec{\mathbf{w}}_{MI}^{I} - c_{E}^{C}\vec{\mathbf{w}}_{IE}^{E}$$

$$\vec{\mathbf{w}}_{SC}^{C} = c_{E}^{C}\vec{\mathbf{w}}_{SE}^{E} - \vec{\mathbf{w}}_{CO}^{C} - c_{O}^{C}\vec{\mathbf{w}}_{OM}^{O} - c_{M}^{C}\vec{\mathbf{w}}_{MI}^{M} - c_{I}^{C}\vec{\mathbf{w}}_{IE}^{I}$$
(3.10)

The left hand side is the rotation rate of the case, which is to be determined; the right hand side is dependent only upon measurable quantities. We will want to relate torque to acceleration in the next section, so we may apply equation (A.8) to equations (3.5) through (3.8).

$$\vec{w}_{SO}^{O} = c_{SC}^{O} \vec{w}_{SC}^{C} - \vec{w}_{CO}^{O} \times c_{C}^{O} \vec{w}_{SC}^{C} + \vec{w}_{CO}^{O}$$
(3.11)

$$\mathbf{\hat{w}}_{SM}^{M} = c_{O}^{M} \mathbf{\hat{w}}_{SO}^{O} - \mathbf{\hat{w}}_{OM}^{M} \times c_{O}^{M} \mathbf{\hat{w}}_{SO}^{O} + \mathbf{\hat{w}}_{OM}^{M}$$
(3.12)

$$\vec{\mathbf{w}}_{SI}^{I} = c_{M}^{I} \vec{\mathbf{w}}_{SM}^{M} - \vec{\mathbf{w}}_{MI}^{I} \times c_{M}^{I} \vec{\mathbf{w}}_{SM}^{M} + \vec{\mathbf{w}}_{MI}^{I} \qquad (3.13)$$

$$\vec{\mathbf{w}}_{SE}^{E} = C_{ISI}^{E} \vec{\mathbf{w}}_{SI}^{I} - \vec{\mathbf{w}}_{IE}^{E} \times C_{ISI}^{E} \vec{\mathbf{w}}_{SI}^{I} + \vec{\mathbf{w}}_{IE}^{E}$$
(3.14)

Unfortunately, equation (3.11) contains \overline{W}_{SC}^{C} , the acceleration of the case, and a difficult quantity to measure.

It will be desirable to know $\dot{\overline{w}}_{SC}^{C}$ in order to predict the trajectory of \overline{w}_{SC}^{C} and thereby optimize the performance of the gimbal system at some time in the future. For a massive vehicle such as the one under consideration here the rotation rate cannot change rapidly. Unable to measure the vehicle's acceleration directly to predict its behavior, a reasonable approach is to assume that it does not change at all. Therefore, throughout this paper it will be assumed that $\dot{\overline{w}}_{SC}^{C} = \overline{0}$. This is not such a bad assumption over a short time interval. Thus, equation (3.11) reduces to

$$\vec{W}_{SO}^{O} = -\vec{W}_{CO}^{O} \times c_{C}^{O} \vec{W}_{SC}^{C} + \vec{W}_{CO}^{O}$$
(3.15)

In theory it is possible to predict the vehicle's acceleration knowing the generated thrust and mass. It is preferable, though, to keep the four-gimbal controller as decoupled as possible from all other vehicular systems, including propulsion.

The vector angular acceleration equations, although compact, are of limited utility by themselves[10]. They need to be expressed in terms of scalar quanties. To this end, equations (3.12) through (3.15) will be expanded using equations (3.5) through (3.8).

$$\vec{W}_{SO} = \vec{D} \quad W_{OZ}$$

$$\vec{W}_{OX} = \vec{D} \quad (3.16)$$

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These last four equations are second order differential equations. Note that the only place where second time derivatives appear is on angles. This is quite a propitious occurrence because, in a later section the state variables will be specified, and the angles will be among the state variables. We will want to express the highest order derivatives of the state variables as functions of lower order derivatives and other known quantities, and to do this we must separate the highest order derivatives from all other factors. Equations (3.16) through (3.19) show where the high order derivatives lie and this is a great help.

We now introduce three variables \triangle SV, \triangle J and \triangle SR. They represent the tilt (rotation) of the inertial platform with respect to inertial space. \triangle SR is measured about the x-axis, \triangle J is measured about the y-axis and \triangle SV is measured about the z-axis of the elevation gimbal. The tilts equal the angular displacement of the inertial platform as sensed by the gyros about the relevant axes. They may be described by differential equations by noting that the rate of change of the tilts must equal the rotation rate of the inertial platform. The rotation rate of the platform is merely \vec{W}_{SE}^{E} . Applying equation (3.8) we have:

$$\Delta SR = \cos \mathbf{I} W_{IX} - \sin \mathbf{I} W_{IZ} \qquad (3.20)$$

$$\Delta J = W_{IY} + I \qquad (3.21)$$

 $\Delta SV = \sin I W_{IX} + \cos I W_{IZ} \qquad (3.22)$

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Now to define the moments of inertia of the gimbals. Let J_k^l be the inertia tensor of gimbal k in the l coordinate system. The matrix representation of an inertia tensor transforms under similarity transformations, i.e., $J_k^m = C_1^m J_k^l C_m^l$. Because the gimbals are symmetric and have been evenly balanced, and because the gimbal-fixed coordinate systems are aligned with the principal axes, the inertia matrix will have zeros off the diagonal when coordinatized in the reference frame of that gimbal. Thus:

$$J_{k}^{k} = \begin{bmatrix} J_{kx} & 0 & 0 \\ 0 & J_{ky} & 0 \end{bmatrix}$$

In general, this is true only when the inertia is coordinatized in the reference frame of that gimbal, and not true in most other reference frames. Thus, J_k^1 , $l \neq k$ will, in general, have nonzero terms off the diagonal. Furthermore, the elevation gimbal is almost symmetric, so we may approximate JEX = JEY = JEZ = JEXYZ The other three gimbals take the shape of bands, each having two roughly equal moments of inertia and a third distinct moment of inertia, the distinct inertia corresponding to the gimbal axis passing through the "hole" in the gimbal. For the given geometry:

 $J_{IY} \stackrel{=}{=} J_{IZ} \stackrel{\triangleq}{=} J_{IYZ}$ $J_{MX} \stackrel{=}{=} J_{MZ} \stackrel{\triangleq}{=} J_{MXZ}$ $J_{OX} \stackrel{=}{=} J_{OY} \stackrel{\triangleq}{=} J_{OXY}$

These approximations will greatly simplify the torque equations. As shown in Table 3.1 the approximations are good ones. The largest error introduced is 8% for the E gimbal, 2.5% for the M gimbal and 0% for the other gimbals. The 8% E gimbal error will have a negligible effect because that gimbal should remain inertial and the exact value of its moment of inertia ought not to matter much.

It should be noted here that the symmetry of the E gimbal gives rise to some useful results.

$$J_{E}^{E} = \begin{bmatrix} J_{EX} & 0 & 0 \\ 0 & J_{EY} & 0 \\ 0 & 0 & J_{EZ} \end{bmatrix} = \begin{bmatrix} J_{EXYZ} & 0 & 0 \\ 0 & J_{EXYZ} & 0 \\ 0 & 0 & J_{EXYZ} \end{bmatrix} = J_{EXYZ} I$$
(3.23)

Thus in any coordinate system k,

$$J_{E}^{k} = C_{E}^{k} J_{E}^{E} C_{E}^{E} = C_{E}^{k} (J_{E}^{I}) C_{E}^{k-1} = J_{EXYZ}^{I}$$
(3.24)

$$J_{E}^{E} = J_{E}^{I} = J_{E}^{M} = J_{E}^{O}$$
 (3.25)

Equation (3.25) has the following interpretation: J_E^E , J_E^I , J_E^M and J_E^O are all different tensors; they just happen to share the same matrix representation.

Table 3.1

Moments of Inertia

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<u>Gimbal</u>	<u>Axis</u>	Name	<u>Value</u>
Elevation	X	J _{EX}	1.3
	Y	JEY	1.2
	Z	JEZ	1.1
Inner	X	JIX	1.7
	Y	JIY	1.3
	Z	JIZ	1.3
Middle	X	JMX	2.2
	Y	JMY	3.0
	Z	JMZ	2.3
Outer	X	JOX	3.0
	Y	Joy	3.0
	2	JOZ	3.9

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Figure 3.1



IV. Derivation of Torque-Acceleration Equations

Torque is the rate of change of angular momentum. For a four-gimbal system there will be four angular momentum vectors to consider, one for each gimbal. This will lead to four torque These four torque equations will be solved for the equations. (Y. B. Ö, Ø). four angular accelerations The angular accelerations can be integrated twice to solve for the angular velocities and the angles themselves, thus completely characterizing the system.

Torques are applied to gimbals through their pivot assemblies. Torques may be applied either along a torque motor axis or normal to a torque motor axis or both. Torques normal to a motor axis are coupled through the bearings; these forces are not controlled directly. Control is exerted directly only on the components of torque along the motor axes. There are four sources in all of torques about a motor axis. They are control voltage, back-emf, coulomb friction and static friction.

Let's examine the relationship between angular momentum and torque. The angular momentum of gimbal j with respect to inertial space (1) is

(1) Strictly speaking, angular momentum is only defined with respect to inertial space. Nonetheless it will be convenient to treat angular momentum like any other vector, especially as regards coordinate transformations. As long as we remember that torque is the rate of change of angular momentum in inertial space there will be no problem.

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$$\vec{H}_{j}^{S} = J_{j}^{S} \vec{W}_{Sj}^{S}$$
(4.1)

In the j coordinate system equation (4.1) becomes:

•

$$\vec{H}_{j} = c_{S}^{j}\vec{H}_{j}^{S} = c_{S}^{j}J_{j}^{S}\vec{W}_{Sj}^{S} = c_{S}^{j}J_{j}^{S}c_{S}^{S}c_{S}^{j}\vec{W}_{Sj}^{S}$$

$$= J_{j}^{j}\vec{W}_{Sj}^{j}$$

$$(4.2)$$

Torque is the time derivative of angular momentum in an inertial coordinate frame. Differentiating equation (4.2) and applying (A.2) and (A.8)

$$\begin{split} \vec{T}_{j}^{j} &= c_{S}^{j}\vec{T}_{j}^{S} = c_{S}^{j}d/dt(\vec{H}_{j}^{S}) = c_{S}^{j}d/dt(c_{j}^{S}\vec{H}_{j}^{j}) \\ &= c_{S}^{j}[c_{j}^{S}\vec{H}_{j}^{i} - \vec{W}_{Sj}^{S} \times (c_{j}^{S}\vec{H}_{j}^{j})] \\ &= \hat{H}_{j}^{j} + c_{S}^{j}[\vec{W}_{Sj}^{S} \times (c_{j}^{S}\vec{H}_{j}^{j})] \\ &= \hat{H}_{j}^{j} + (c_{S}^{j}\vec{W}_{Sj}^{S}) \times (c_{S}^{j}c_{j}^{S}\vec{H}_{j}^{j}) \\ &= \hat{H}_{j}^{j} + \vec{W}_{Sj}^{j} \times \vec{H}_{j}^{j} \end{split}$$
(4.3)

For a rigid body such as a gimbal, $d/dt(J_j^j) = 0$, so equation (4.3) becomes:

$$\vec{T}_{j}^{j} = J_{j}^{j} \vec{W}_{Sj}^{j} + \vec{W}_{Sj}^{j} \times (J_{j}^{j} \vec{W}_{Sj}^{j})$$
(4.4)

For the O gimbal we have:

$$\vec{T}_{0} = \vec{T}_{C0} + \vec{T}_{M0} = J_{0}^{0} = J_{0}^{0} = \vec{W}_{S0} + \vec{W}_{S0} \times (J_{0}^{0} = \vec{W}_{S0})$$
(4.5)

$$\overline{T}_{CO}^{O} = J_{O}^{O} \overline{W}_{SO}^{O} + \overline{W}_{SO}^{O} \times (J_{O}^{O} \overline{W}_{SO}^{O}) + C_{M}^{O} \overline{T}_{M}^{M}$$
(4.6)

 \overline{T}_{CO}^0 represents the torque transmitted from the case to the outer gimbal as seen from the outer gimbal. The form of equation (4.6) will prove most useful. Similar equations can be written for the other gimbals.

$$\overline{\mathbf{T}}_{OM}^{M} = \mathbf{J}_{\underline{M}}^{M} \overline{\mathbf{M}}_{M}^{M} + \overline{\mathbf{W}}_{\underline{S}}^{M} \mathbf{X} \left(\mathbf{J}_{\underline{M}}^{M} \overline{\mathbf{W}}_{\underline{N}}^{M} \right) + \mathbf{C}_{\underline{I}}^{M} \overline{\mathbf{M}}_{\underline{I}}^{I}$$

$$(4.7)$$

$$\overline{\mathbf{T}}_{\mathbf{MI}}^{\mathbf{I}} = \mathbf{J}_{\mathbf{I}}^{\mathbf{I}} \overline{\mathbf{W}}_{\mathbf{I}}^{\mathbf{I}} + \overline{\mathbf{W}}_{\mathbf{I}}^{\mathbf{I}} \mathbf{X} \left(\mathbf{J}_{\mathbf{I}}^{\mathbf{I}} \overline{\mathbf{W}}_{\mathbf{I}}^{\mathbf{I}} \right) + \mathbf{C}_{\mathbf{E}}^{\mathbf{I}} \overline{\mathbf{T}}_{\mathbf{E}}^{\mathbf{I}}$$

$$\overline{\mathbf{T}}_{\mathbf{E}}^{\mathbf{E}} = \mathbf{I}_{\mathbf{E}}^{\mathbf{E}} \overline{\mathbf{E}}_{\mathbf{I}}^{\mathbf{E}} \mathbf{U}_{\mathbf{E}}^{\mathbf{E}} \mathbf{U}_{\mathbf{E}}^{\mathbf{E}}$$

$$(4.8)$$

$$\mathbf{\tilde{I}}_{E} = J_{E}^{E} \mathbf{W}_{SE}^{E} + \mathbf{W}_{SE}^{E} \mathbf{X} (J_{E}^{E} \mathbf{W}_{SE}^{E})$$
(4.9)

Because the elevation gimbal is assumed to be symmetric, and \vec{W}_{SE}^E is to be kept small, equation (4.9) reduces to:

$$\hat{T}_{IE}^{E} \simeq J_{E}^{E} \hat{W}_{SE}^{E}$$
 (4.10)

The torque motor force from the inner to the elevation gimbal is along the y-axis of both gimbals.

$$T_{IEY}^{E} = J_{EXYZ} \dot{W}_{EY}$$

= $J_{EXYZ} \{\dot{\mathbf{Y}} - \sin B (-\dot{\emptyset} W_{OZ} + \ddot{\Theta}) + \cos B (\cos \Theta \ddot{\emptyset} + \sin \Theta \dot{\emptyset} W_{OX} + \dot{\Theta} W_{MZ}) - \dot{B} W_{IX}\}$ (4.11)

Equation (4.11) can be rewritten as:

$$Y_{E} = J_{EXYZ}\ddot{Y} - \sin B J_{EXYZ}\ddot{\Theta} + \cos B \cos \Theta J_{EXYZ}\ddot{\Theta}$$
(4.12)
Where $Y_{E} \triangleq$
$$T_{IEY}^{E} + J_{EXYZ}(-\sin B \dot{D} W_{OZ} - \cos B \sin \Theta \dot{D} W_{OX} - \cos B \dot{\Theta} W_{MZ} + \dot{B} W_{IX})$$
(4.13)

 Y_E is a quantity that contains all of the terms of the torque equation for T^E_{IEY} that do not contain an angular acceleration. Similar definitions will be made for the other gimbals. Proceeding in a parallel manner with the inner gimbal we repeat equation (4.8).

$$\overline{\mathbf{T}}_{MI}^{I} = J_{I}^{I} \overline{\mathbf{W}}_{SI}^{I} + \overline{\mathbf{W}}_{SI}^{I} \times (J_{I}^{I} \overline{\mathbf{W}}_{SI}^{I}) + C_{E}^{I} \overline{\mathbf{T}}_{E}^{E}$$
(4.8)

Applying equations (3.14) and (4.13) to (4.8):

$$\vec{T}_{MI}^{I} = J_{I}^{I} \vec{w}_{SI}^{I} + \vec{w}_{SI}^{I} \times (J_{I}^{I} \vec{w}_{SI}^{I}) + c_{EJE}^{I} (c_{I}^{E} \vec{w}_{SI}^{I})
- \vec{w}_{IE}^{E} \times (c_{I}^{E} \vec{w}_{SI}^{I}) + \vec{w}_{IE}^{E})
= (J_{I}^{I} + c_{EJECE}^{I}) \vec{w}_{I}^{I} + \vec{w}_{SI}^{I} \times (J_{I}^{I} \vec{w}_{SI}^{I})
+ c_{EJECE}^{I} (-\vec{w}_{IE}^{I} \times \vec{w}_{SI}^{I} + \vec{w}_{IE}^{I})
= (J_{I}^{I} + J_{E}^{I}) \vec{w}_{SI}^{I} + \vec{w}_{SI}^{I} \times (J_{I}^{I} \vec{w}_{SI}^{I}) + J_{E}^{I} (-\vec{w}_{IE}^{I} \times \vec{w}_{IE}^{I} + \vec{w}_{IE}^{I})
= (J_{I}^{I} + J_{E}^{I}) \vec{w}_{SI}^{I} + \vec{w}_{SI}^{I} \times (J_{I}^{I} \vec{w}_{SI}^{I}) + J_{E}^{I} (-\vec{w}_{IE}^{I} \times \vec{w}_{IE}^{I} + \vec{w}_{IE}^{I})
= (4.14)$$

Recalling equation (3.25) and expanding (4.14):

$$\overline{T}_{MI}^{I} = \begin{vmatrix} J_{IX} + J_{EXYZ} & 0 & 0 \\ 0 & J_{IYZ} + J_{EXYZ} & 0 \\ 0 & 0 & J_{IYZ} + J_{EXYZ} & 0 \\ 0 & 0 & J_{IYZ} + J_{EXYZ} & 0 \\ W_{IY} & W_{IZ} & (J_{IYZ} - J_{IYZ}) \\ W_{IZ} & W_{IX} & (J_{IX} - J_{IYZ}) \\ W_{IX} & W_{IY} & (J_{IYZ} - J_{IX}) & 0 & J_{EXYZ} & 0 \\ W_{IX} & W_{IY} & (J_{IYZ} - J_{IX}) & 0 & 0 & J_{EXYZ} \\ \end{vmatrix} \begin{vmatrix} (J_{IX} + J_{EXYZ}) & W_{IX} - J_{EXYZ} & W_{IZ} \\ (J_{IYZ} + J_{EXYZ}) & W_{IY} + J_{EXYZ} & Y \\ (J_{IYZ} + J_{EXYZ}) & W_{IZ} + J_{EXYZ} & Y \\ \end{vmatrix} \begin{vmatrix} (J_{IYZ} + J_{EXYZ}) & W_{IZ} + J_{EXYZ} & Y \\ W_{IX} & W_{IY} & (J_{IYZ} - J_{IY}) & W_{IY} & (J_{IYZ} - J_{IY}) \end{vmatrix}$$

When the elevation gimbal is inertial both W_{IX} and W_{IZ} should be small; ideally they will be zero. The product of such small terms will certainly be negligible. Therefore the term W_{IX} W_{IZ} ($J_{IX} - J_{IYZ}$) has been dropped from equation (4.15). Motor torque from the M to the I gimbal is along the z-axis.

$$T_{MIZ}^{I} = (J_{IYZ} + J_{EXYZ}) \dot{W}_{IZ} + W_{IX} W_{IY} (J_{IYZ} - J_{IX}) + J_{EXYZ} \dot{I} W_{IX}$$
$$= (J_{IYZ} + J_{EXYZ}) (-\sin\theta \ddot{\theta} + \cos\theta \dot{\theta} W_{OX} - \dot{\theta} W_{MX}$$
$$+ \ddot{B}) + W_{IX} W_{IY} (J_{IYZ} - J_{IX}) + J_{EXYZ} \dot{I} W_{IX}$$
(4.16)

Equation (4.16) can be rewritten in a similar fashion to (4.11):

$$Y_{I} = (J_{IYZ} + J_{EXYZ}) (\ddot{B} - \sin \varphi \ddot{p})$$
(4.17)
Where $Y_{I} \stackrel{c}{=} T_{MIZ}^{I} + (J_{IYZ} + J_{EXYZ}) (\dot{\varphi} W_{MX} - \cos \varphi \dot{p} W_{OX})$

+
$$W_{IX} W_{IY} (J_{IX} - J_{IYZ}) - \dot{Y} J_{EXYZ} W_{IX}$$
 (4.18)

Proceeding to the middle gimbal:

$$\vec{T}_{OM}^{M} = J_{M}^{M} \vec{w}_{SM}^{M} + \vec{w}_{SM}^{M} \times (J_{M}^{M} \vec{w}_{SM}^{M}) + C_{I}^{M} \vec{T}_{MI}^{I}$$
(4.7)

$$= \begin{vmatrix} J_{MXZ} & 0 & 0 \\ 0 & J_{MY} & 0 \\ 0 & 0 & J_{MXZ} \end{vmatrix} \begin{vmatrix} \dot{W}_{MX} \\ \dot{W}_{MY} \\ \dot{W}_{MZ} \end{vmatrix} + \begin{vmatrix} W_{MY} & W_{MZ} & (J_{MXZ} - J_{MY}) \\ W_{MZ} & W_{MX} & (J_{MXZ} - J_{MXZ}) \\ W_{MX} & W_{MY} & (J_{MY} - J_{MXZ}) \end{vmatrix}$$
$$+ \begin{vmatrix} \cos B & -\sin B & 0 \\ \sin B & \cos B & 0 \\ 0 & 0 & 1 \end{vmatrix} \begin{vmatrix} (J_{IX} + J_{EXYZ}) & \dot{W}_{IX} - J_{EXYZ} & W_{IZ} \\ (J_{IYZ} + J_{EXYZ}) & \dot{W}_{IY} + J_{EXYZ} & \ddot{Y} \\ (J_{IYZ} + J_{EXYZ}) & \dot{W}_{IZ} + W_{IX} & W_{IY} & (J_{IYZ} - J_{IX}) \\ + J_{EXYZ} & \dot{Y} & W_{IX} & (4.19) \end{vmatrix}$$

Torque from the O to the M gimbal is along the x-axis. Using equations (3.17) and (3.18) the x component of (4.19) can be expanded as follows:

$$T_{OMX}^{M} = J_{MXZ} \overset{V}{}_{MX} + \overset{W}{}_{MY} \overset{W}{}_{MZ} (J_{MXZ} - J_{MY}) + \cos B \{ (J_{IX} + J_{EXYZ}) \overset{V}{}_{IX} - J_{EXYZ} \overset{W}{}_{IZ} \overset{I}{}_{I} \} - \sin B \{ (J_{IYZ} + J_{EXYZ}) \overset{V}{}_{IY} + J_{EXYZ} \overset{V}{}_{I} \}$$

$$= J_{MXZ} \overset{\tilde{W}_{MX}}{} + \overset{W}_{MY} \overset{W}_{MZ} (J_{MXZ} - J_{MY}) + \cos B (J_{IX} + J_{EXYZ}) (\cos B \overset{W}{W}_{MX} + \sin B \overset{W}{W}_{MY} + \dot{B} W_{IY}) - \cos B J_{EXYZ} \overset{W}{IZ} \overset{\tilde{X}}{} - \sin B (J_{IYZ} + J_{EXYZ}) (-\sin B \overset{W}{W}_{MX} + \cos B \overset{W}{W}_{MY} - \dot{B} W_{IX}) - \sin B J_{EXYZ} \overset{\tilde{Y}}{} = (J_{MXZ} + \cos^{2}B J_{IX} + \sin^{2}B J_{IYZ} + J_{EXYZ}) (\ddot{\Theta} - \dot{D} W_{OZ}) + \overset{W}{}_{MY} \overset{W}{}_{MZ} (J_{MXZ} - J_{MY}) + \sin B \cos B (J_{IX} - J_{IYZ}) (\cos \Theta \overset{D}{D} + \sin \Theta \overset{D}{D} W_{OX} + \dot{\Theta} W_{MZ}) + \dot{B} \{\cos B (J_{IX} + J_{EXYZ}) (-\sin B W_{MX} + \cos B W_{MY}) + \sin B (J_{IYZ} + J_{EXYZ}) (\cos B W_{MX} + \sin B W_{MY}) \} - \cos B J_{EXYZ} \overset{W}{}_{IZ} \overset{T}{} - \sin B J_{EXYZ} \overset{T}{}_{IX}$$

$$(4.20)$$

Equation (4.20) can be rearranged like this:

$$Y_{M} = -\sin B J_{EXYZ} \ddot{T} + (J_{MXZ} + \cos^{2}B J_{IX} + \sin^{2}B J_{IYZ} + J_{EXYZ}) \ddot{\Theta} + \sin B \cos B \cos \Theta (J_{IX} - J_{IYZ}) \ddot{\Theta}$$
(4.21)

Where
$$Y_{M} \triangleq T_{OMX}^{M}$$

+ $(J_{MXZ} + \cos^{2B} J_{IX} + \sin^{2B} J_{IYZ} + J_{EXYZ}) \overset{\circ}{D} W_{OX}$
+ $\sin^{B} \cos^{B} (J_{IYZ} - J_{IX}) (\sin^{O} \overset{\circ}{D} W_{OX} + \overset{\circ}{\Theta} W_{MZ} - \overset{\circ}{B} W_{MX})$
- $(\cos^{2B} J_{IX} + \sin^{2B} J_{IYZ} + J_{EXYZ}) \overset{\circ}{B} W_{MY}$
+ $\cos^{B} J_{EXYZ} W_{IZ} \overset{\circ}{Y}$ (4.22)

Finally, for the outer gimbal:

$$\vec{T}_{CO}^{O} = J_{O} \vec{W}_{SO}^{O} + \vec{W}_{SO}^{O} \times (J_{O} \vec{W}_{SO}^{O}) + C_{M}^{OTM}$$
(4.6)

The component of interest here is along the y-axis since that is where the torque motor is. The algebra required is extremely tedious, and little insight is obtained. We will not go through the entire derivation. A rigorous derivation is given in [6]. The resulting equation for Y_O (which is really all we want) is:

$$\begin{split} \mathbf{Y}_{0} &= \cos B \ \cos \theta \ \mathbf{J}_{\mathrm{EXYZ}} \ \mathbf{\ddot{Y}} \\ &= \sin \theta \ (\mathbf{J}_{\mathrm{IYZ}} + \mathbf{J}_{\mathrm{EXYZ}}) \ \mathbf{\ddot{B}} \\ &+ \sin B \ \cos B \ \cos \theta \ (\mathbf{J}_{\mathrm{IX}} - \mathbf{J}_{\mathrm{IYZ}}) \ \mathbf{\ddot{\theta}} \\ &+ (\mathbf{J}_{0\mathrm{XY}} + \sin^{2} \theta \ \mathbf{J}_{\mathrm{MXZ}} + \cos^{2} \theta \ \mathbf{J}_{\mathrm{MY}} + \sin^{2} \theta \ \mathbf{J}_{\mathrm{IYZ}} \\ &+ \sin^{2} B \ \cos^{2} \theta \ \mathbf{J}_{\mathrm{IX}} + \cos^{2} B \ \cos^{2} \theta \ \mathbf{J}_{\mathrm{IYZ}} + \mathbf{J}_{\mathrm{EXYZ}}) \ \mathbf{\ddot{\theta}} \qquad (4.23) \end{split}$$

Where $\mathbf{Y}_{0} \triangleq \mathbf{T}_{\mathrm{COY}}^{0} \\ &+ \mathbf{W}_{0\mathrm{X}} \ \mathbf{W}_{0\mathrm{Y}} \ (\mathbf{J}_{0\mathrm{Z}} - \mathbf{J}_{0\mathrm{XY}}) \\ &+ \sin B \ \cos B \ \cos \theta \ (\mathbf{J}_{\mathrm{IX}} - \mathbf{J}_{\mathrm{IYZ}}) \ (\mathbf{\dot{\theta}} \ \mathbf{W}_{0\mathrm{Z}} - \mathbf{\dot{B}} \ \mathbf{W}_{\mathrm{MY}}) \\ &+ \sin \theta \ \cos \theta \ (\mathbf{J}_{\mathrm{MXZ}} - \mathbf{J}_{\mathrm{MY}} + \sin^{2} \theta \ \mathbf{J}_{\mathrm{IYZ}} - \mathbf{J}_{\mathrm{IX}}) \ (\mathbf{\dot{\theta}} \ \mathbf{W}_{0\mathrm{X}} - \mathbf{\dot{\theta}} \ \mathbf{W}_{0\mathrm{Y}}) \\ &- (\sin^{2} \theta \ \mathbf{J}_{\mathrm{MXZ}} + \cos^{2} \theta \ \mathbf{J}_{\mathrm{MY}} + \sin^{2} \theta \ \mathbf{J}_{\mathrm{IYZ}} + \sin^{2} B \ \cos^{2} \theta \ \mathbf{J}_{\mathrm{IX}} \\ &+ \cos^{2} B \ \cos^{2} \theta \ \mathbf{J}_{\mathrm{IYZ}} + \mathbf{J}_{\mathrm{EXYZ}}) \ \mathbf{\dot{\theta}} \ \mathbf{W}_{0\mathrm{Z}} \\ &+ \sin \theta \ \mathbf{W}_{\mathrm{MY}} \ \mathbf{W}_{\mathrm{MX}} \ (\mathbf{J}_{\mathrm{MY}} - \mathbf{J}_{\mathrm{MXZ}}) \\ &+ \sin \theta \ \mathbf{W}_{\mathrm{MY}} \ \mathbf{W}_{\mathrm{MX}} \ (\mathbf{J}_{\mathrm{IYZ}} - \mathbf{J}_{\mathrm{IX}}) \\ &+ \sin B \ \cos \theta \ \mathbf{J}_{\mathrm{EXYZ}} \ \mathbf{\dot{Y}} \ \mathbf{W}_{\mathrm{IX}} \qquad (4.24) \end{split}$

Equations (4.12), (4.17), (4.21) and (4.23) may be combined into a single matrix equation.

Where $\vec{Y} \cong (Y_E, Y_I, Y_M, Y_O)^T$ $\vec{Z} \cong (\vec{Y}, \vec{B}, \vec{O}, \vec{D})^T$	
$L = \begin{bmatrix} L_{11} & L_{12} & L_{13} & L_{14} \\ L_{21} & L_{22} & L_{23} & L_{24} \\ L_{31} & L_{32} & L_{33} & L_{34} \\ L_{41} & L_{42} & L_{43} & L_{44} \end{bmatrix}$	
With $L_{11} = J_{EXYZ}$	(4.26)
$L_{12} = L_{21} = 0$	(4.27)
$L_{13} = L_{31} = -sinB J_{EXYZ}$	(4,28)
$L_{14} = L_{41} = \cos B \cos \theta J_{EXYZ}$	(4.29)
$L_{22} = J_{IYZ} + J_{EXYZ}$	(4.30)
$L_{23} = L_{32} = 0$	(4.31)
$L_{24} = L_{42} = -\sin \Theta \left(J_{IYZ} + J_{EXYZ} \right)$	(4.32)
$L_{33} = J_{MXZ} + \cos^2 B J_{IX} + \sin^2 B J_{IYZ} + J_{EXYZ}$	(4.33)
$L_{34} = L_{43} = sinB cosB cosP (J_{IX} - J_{IYZ})$	(4.34)
$L_{44} = J_{0XY} + \sin^2 \theta J_{MXZ} + \cos^2 \theta J_{MY} + \sin^2 \theta J_{IYZ}$	
+ $\sin^{2}B \cos^{2}\Theta J_{IX} + \cos^{2}B \cos^{2}\Theta J_{IYZ} + J_{EXYZ}$	(4.35)

A term lying on the diagonal of L, L_{ii} , is the effective moment of inertia of gimbal i and those gimbals inside it as seen looking into the pivot axis of gimbal i. For example, if all four gimbals are treated as a single unit, then the inertia along the y-axis of the outer gimbal is just L_{44} . Similarly, L_{33} is the inertia of the three innermost gimbals along the x-axis of the middle gimbal. The off-diagonal terms of L are a consequence of the fact that an inertia matrix may no longer be diagonal if

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 $\vec{Y} = L \vec{Z}$

(4.25)
coordinatized in a coordinate system not attached to the appropriate gimbal.

L contains information about the geometry of the gimbals. We have already assumed that the elevation gimbal is symmetric. This implies that the orientation of the elevation gimbal is not relevant to the overall geometry of the system and therefore we would not expect the elevation angle \mathbf{X} to appear in equations (4.26) through (4.35). The outer angle Ø also should not affect the gimbal geometry, so we would not expect Ø to appear in equations (4.26) through (4.35) either. These expectations are realized. The gimbal configuration as defined by the matrix L is only a function of B and Θ_{c}

Note that L is symmetric. This is an instance of a reciprocity relationship between torque and angular acceleration. A torque applied at angle i will produce a response at angle j equal to the response at angle i to a torque at angle j.

The actual torque values are nestled into the Y vector together with a great many other terms having the same dimensions as torque. These other terms for the most part resemble Coriolis forces, although their exact interpretation is not always obvious. In any event, for reasonable gimbal rates and reasonable torque levels the torque terms will dominate the Coriolis forces.

Equation (4.25) allows the computation of torque given acceleration. In actuality we know the torque since the controller will be supplying the control signals; it is the acceleration we wish to compute. So we may take the inverse of equation (4.25) to come up with:

$$\overline{Z} = L^{-1}\overline{Y} = M \overline{Y}$$
 where $M \triangleq L^{-1}$ (4.36)

M will of course be symmetric since L is. The computation of M is aided by repeated application of the following matrix identity:

Before presenting the terms of M it is helpful to define a quantity called DENOM. DENOM is the determinant of $(L_{44}-L_{43}L_{33}L_{34})$ obtained when using formula (4.37). Since this quantity appears in each element of M it will be much easier to define DENOM once than to write it out in full each time.

$$DENOM = [J_{OXY} + \sin^{2}\Theta J_{MXZ} + \cos^{2}\Theta (J_{MY} + \cos^{2}B J_{IYZ} + \sin^{2}B \{J_{IX} + J_{EXYZ}\})] [J_{MXZ} + \cos^{2}B (J_{IX} + J_{EXYZ}) + \sin^{2}B J_{IYZ}]$$

$$- \cos^{2}B \sin^{2}B \cos^{2}\Theta [J_{IYZ} - J_{IX} - J_{EXYZ}]^{2} (4.38)$$

$$M_{11} = \{[J_{OXY} + \sin^{2}\Theta J_{MXZ} + \cos^{2}\Theta (J_{MY} + J_{IX} + J_{EXYZ})]$$

$$[J_{EXYZ} + \sin^{2}B J_{IYZ} + \cos^{2}B J_{IX} + J_{MXZ}]$$

$$- \cos^{2}B \cos^{2}\Theta [J_{IX} - J_{IYZ}] [J_{MXZ} + J_{IX} + J_{EXYZ}]\}$$

$$/ DENOM / J_{EXYZ} (4.39)$$

$$M_{12} = -\cos B \sin \Theta \cos \Theta [J_{EXYZ} + J_{IX} + J_{MXZ}] / DENOM (4.40)$$

$$M_{13} = \sin B [J_{OXY} + \sin^{2}\Theta J_{MXZ} + J_{EXYZ})] / DENOM (4.41)$$

$$M_{14} = -\cos B \cos \Theta [J_{MXZ} + J_{IX} + J_{EXYZ}] / DENOM (4.42)$$

$$M_{22} = \{\cos^{2B} \cos^{2\Theta} [J_{IYZ} - J_{IX} - J_{EXYZ}] [J_{EXYZ} + J_{IX} + J_{MXZ}] + [J_{MXZ} + \cos^{2B} (J_{IX} + J_{EXYZ}) + \sin^{2B} J_{IYZ}] \\ [J_{0XY} + \cos^{2\Theta} (J_{MY} + J_{IX}) + \sin^{2\Theta} (J_{MXZ} + J_{IYZ}) \\ + J_{EXYZ}] \} / [J_{IYZ} + J_{EXYZ}] / DENOM$$
(4.43)
$$M_{23} = \sin B \cos B \sin \Theta \cos \Theta [J_{IYZ} - J_{IX} - J_{EXYZ}] / DENOM (4.44) \\ M_{24} = \sin \Theta [J_{MXZ} + \cos^{2B} (J_{IX} + J_{EXYZ}) + \sin^{2B} J_{IYZ}] / DENOM$$
(4.45)

$$M_{33} = [J_{0XY} + \sin^2 \theta J_{MXZ} + \cos^2 \theta (J_{MY} + \cos^2 B J_{IYZ} + \sin^2 B \{J_{IX} + J_{EXYZ}\})] / DENOM \qquad (4.46)$$

$$M_{34} = \cos B \sin B \cos \theta \left[J_{IYZ} - J_{IX} - J_{EXYZ} \right] / DENOM \qquad (4.47)$$
$$M_{44} = \left[J_{MXZ} + \cos^2 B \left(J_{IX} + J_{EXYZ} \right) + \sin^2 B J_{IYZ} \right] / DENOM \qquad (4.48)$$

The above equations are rather difficult to manipulate and verify. By writing a computer program to numerically multiply L amd M it was found that M is indeed the inverse of L.

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V. Torque and Friction

The torque produced by a given torque motor is proportional to the current through it. The constant of proportionality is Kt/r. This current will equal the applied voltage, in this case the control signal, minus the back-emf generated by the motor, divided by the resistance of the motor windings. Back-emf is created when a torque motor acts like an electric generator, putting out a voltage proportional to the relative rotation rate of the rotor and stator, tending to cancel any rotation of the gimbals. The constant of proportionality is denoted by Kv. These torque motor parameters will differ from gimbal to gimbal. Inductive effects in the motors are negligible.

Anathema to designers of precision guidance equipment, friction is nonetheless a force to be reckoned with, or at least accounted for. It is a major factor in the four-gimbal system; much of the torque supplied by the torque motors is used to overcome friction. In fact, in the absence of friction there would be almost no forces acting to perturb the inertial platform except in the neighborhood of gimbal lock.

There are essentially two types of friction: static friction and Coulomb friction. Although they originate in the same intermolecular forces the analysis of the two types of friction is substantially different. We deal first with Coulomb friction.

Gimbals in relative motion will be subject to Coulomb friction. We will use a very simple model for friction in the simulation.

Tcoulomb = -sgn (relative gimbal rate) X Tcoulomb-limit(5.1) Where sgn(x) = $\begin{cases} -1 & x < 0 \\ 0 & x = 0 \\ 1 & x > 0 \end{cases}$

This simple model has adequately predicted Coulomb friction in earlier simulations. It has the advantage of requiring only a single parameter for each gimbal. Gully[7] goes into more sophisticated models. The net torque at each pivot can now be determined in terms of control signals and rotation rates.

$$T_{IEY}^{E} = (Kt/r)_{E} \{U_{E} - (Kv)_{E} \dot{\mathbf{I}}\} - sgn(\dot{\mathbf{I}}) Tcl_{E}$$
(5.2)

$$T_{MIZ} = (Kt/r) \{U - (Kv) B\} - sgn(B) Tcl [5.3]
 T_{M} = (Kt/r) \{U - (Kv) \dot{\Theta}\} - sgn(\dot{\Theta}) Tcl [5.4]
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$$T_{COY}^{0} = (Kt/r)_{0} \{U_{0} - (Kv)_{0}\dot{p}\} - sgn(\dot{p}) Tcl_{0}$$
 (5.5)

Static friction or stiction as it is often called, is the force tending to prevent adjacent bodies from moving at all relative to one another once they have stopped moving. Static friction is in general stronger than Coulomb friction, the latter being effective only after the onset of relative motion. Static friction is quite annoying from the viewpoint of the four-gimbal controller. It means that a comparatively large torque must be applied to get a stuck gimbal pair unstuck.

The model used for static friction here is extremely simple. Others are certainly possible and ought to be analyzable in the same framework. The model used here is characterized by a single parameter, the static friction torque limit. The static friction torque limit will differ from gimbal to gimbal. The model works as follows:

Whenever two adjacent gimbals are not in relative motion (i.e. their relative rotation rate is zero) they will be considered stuck until the magnitude of the torque supplied by a torque motor from one gimbal to the other exceeds the static friction torque limit. If a greater amount of torque is applied, then the gimbals will be free to rotate subject to Coulomb friction. If the relative rotation rate is nonzero, no matter how small in magnitude, then the gimbals will not be stuck.

This may cause some difficulty in the computer simulation of the system. Because of numerical considerations it is unlikely that the relative rate of any gimbal pair will exactly equal zero in the simulation. The approach taken then is to check if the relative rotation rate about any axis has recently passed through zero (i.e. changed sign). If so, then a comparison of applied torque with the static friction limit is made as though the rotation rate were exactly zero, and the system is treated accordingly.

When two gimbals are stuck they will travel together. Neither a relative velocity nor a relative acceleration will be experienced, despite any applied torque up to the static friction torque limit. This causes problems in applying equation (4.25). We no longer know the net torque being supplied between the stuck gimbals. The motor torque is known, but not the amount of stiction. Static friction will be just adequate to prevent motion along the affected axis, but it is not possible to predict a priori. Calculation of angular acceleration by means of equation (4.36) is thereby rendered impossible. Some other method is required.

The method used is to go back to equation (4.25). If all net torques were known then (4.25) could be inverted as was done in (4.36). But the net torque will not be known at a stuck gimbal pair. So a constraint will have been lost from equation (4.25) and the system will be indeterminate. However, another constraint may be added, namely that the acceleration of the affected angle will be zero. This can be best expressed by rearranging and partitioning the elements of equation (4.25)

Let \vec{Y}_1 be a vector containing those elements of \vec{Y} not affected by stiction. \vec{Y}_1 can be computed since the net torque is readily computable in the absence of stiction. Let \vec{Z}_1 be a vector containing the angular accelerations in \vec{Z} not affected by stiction. These are the values we wish to compute. Similarly, let \vec{Y}_2 be a vector containing the elements of \vec{Y} that are affected by static friction. Even though the torque motor contributions to \vec{Y}_2 will be known, the static friction contributions will not, as was discussed above. Lastly, let \vec{Z}_2 be a vector containing those angular accelerations that are affected by static friction. \vec{Z}_2 will be identically zero. Introduce a new matrix L' whose elements are permuted elements of L such that:

$$\begin{vmatrix} \overline{Y}_1 \\ -\overline{Y}_2 \\ \overline{Y}_2 \end{vmatrix} = L^* \begin{vmatrix} \overline{Z}_1 \\ -\overline{Z}_2 \\ \overline{Z}_2 \end{vmatrix}$$
(5.6)

L' can be partitioned like so:

$$L' = \begin{vmatrix} L'_{11} & L'_{12} \\ L'_{21} & L'_{22} \end{vmatrix}$$
(5.7)

Equations (5.6) and (5.7) can be combined as follows:

$$\vec{Y}_{1} = L_{11}^{\prime} \vec{Z}_{1} + L_{12}^{\prime} \vec{Z}_{2}$$

$$= L_{11}^{\prime} \vec{Z}_{1} \qquad (5.8)$$

$$\vec{Z}_{1} = (L_{11}^{\prime})^{-1} \vec{Y}_{1} \qquad (5.9)$$

$$\vec{Z} = \begin{vmatrix} \vec{Z} \\ -\vec{1} \\ \vec{Z} \\ 2 \end{vmatrix} = \begin{vmatrix} (L_1) - 1 & \vec{Y} \\ -\vec{1} \\ \vec{Z} \\ 2 \end{vmatrix}$$
(5.10)

This is what we wanted. The presence of stiction leads to a smaller set of equations to solve. The exact contribution of static friction was not needed. If stiction is present and equation (5.10) is used, or stiction is absent and equation (4.36) is used, the angular accelerations, and thus the angles can be correctly determined. VI. Optimal Control of the Four-Gimbal System

Up to this point, differential equations have been derived that relate acceleration of the gimbals to control signals. These have all been scalar equations although they may be considered selected components of a set of vector equations of the type exemplified by (4.3). Rearranging the scalar equations into state-space form will aid the application of modern optimal control theory to the four-gimbal problem. Define an 11-dimensional state vector \vec{X} and a 4-dimensional control vector \vec{U} by:

$$\overline{\mathbf{X}} = (\mathbf{Y}, \mathbf{B}, \mathbf{\Theta}, \emptyset, \dot{\mathbf{I}}, \dot{\mathbf{B}}, \dot{\mathbf{\Theta}}, \dot{\emptyset}, \Delta SR, \Delta J, \Delta SV)^{\mathrm{T}}$$
$$\overline{\mathbf{U}} = (\mathbf{U}_{\mathrm{E}}, \mathbf{U}_{\mathrm{I}}, \mathbf{U}_{\mathrm{M}}, \mathbf{U}_{\mathrm{O}})^{\mathrm{T}}$$

 \vec{X} is composed of the gimbal angles and velocities plus the inertial platform tilts ΔSV , ΔJ , ΔSR . The entire dynamics of the four-gimbal system can be compressed into a single nonlinear vector differential equation by writing:

$$\dot{\vec{X}} = \vec{f} (\vec{X}, \vec{U}, \vec{W}_{SC}^{C})$$
(6.1)

Explicitly, \overline{X} may be expanded using equations (3.20)-(3.22) and (4.36) to yield:

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$$\dot{X}_1 = X_5$$
 (6.2)

$$\dot{x}_2 = x_6$$
 (6.3)

$$\dot{x}_3 = x_7$$
 (6.4)

$$X_{4} = X_{8} \tag{6.5}$$

$$X_5 = M_{11}Y_1 + M_{12}Y_2 + M_{13}Y_3 + M_{14}Y_4$$
 (6.6)

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$$X_6 = M_{21}Y_1 + M_{22}Y_2 + M_{23}Y_3 + M_{24}Y_4$$
 (6.7)

$$\dot{X}_7 = M_{31}Y_1 + M_{32}Y_2 + M_{33}Y_3 + M_{34}Y_4$$
 (6.8)

$$\dot{\mathbf{X}}_{8} = M_{41}\mathbf{Y}_{1} + M_{42}\mathbf{Y}_{2} + M_{43}\mathbf{Y}_{3} + M_{44}\mathbf{Y}_{4}$$
(6.9)

$$\dot{X}_{9} = \cos X_{1} W_{IX} - \sin X_{1} W_{IZ}$$
(6.10)

$$\dot{X}_{10} = W_{IY} + X_5$$
 (6.11)

$$\dot{X}_{11} = \sin X_1 W_{IX} + \cos X_1 W_{IZ}$$
 (6.12)

The M's are functions of B and Θ , or X₂ and X₃. The W's are functions of case rates, gimbal angles and angular rates, and so can be expressed in terms of $\overrightarrow{W}_{SC}^C$ and X's. The Y's are also functions of X's and U's. Only the states, controls and case rates appear on the right hand side of equations (6.2)-(6.12) in accordance with the formulation (6.1).

One advantageous aspect of this formulation of the system relates to sensors. Each state variable has a unique sensor associated with it. X_1 through X_4 are measured by resolvers, X_5 through X_8 are measured by tachometers and X_9 through X_{11} are measured by gyros. There can be no question as to whether or not the system is observable. Measurement noise does complicate the picture somewhat, but filtering of the sensor data should suffice to provide accurate estimates of the state variables. The oft-quoted Separately from issues of control for linear systems. The system under study is not linear, but as we will shortly see, it can be approximated by linear equations. Henceforth we will not be concerned with estimation of state except insofar as it relates to the validity of simulation studies. The nonlinear equations embodied in (6.1) are fine for numerical analysis and simulation. They allow for numerical integration of the dynamical equations given any inputs to predict the trajectory of the system. As far as optimal control is concerned, equation (6.1) is horrendous. The theory of nonlinear optimal control is difficult to apply to actual real-time processes. For this reason linear quadratic optimal control will be applied to a linearized discrete-time version of the state equations.

Start by looking at the system at time to and at short intervals thereafter. Over a short enough interval, tens of milliseconds for example, the system will not change state much and the dynamics may be faithfully described by linear equations.

It is necessary to choose a nominal operating point about which to perform the linearization. One could choose $\vec{X} = \vec{X}(t0)$, $\vec{U} = \vec{U}(t0)$ and $\vec{W} = \vec{W}_{SC}^{C}(t0)$. This is valid if \vec{X} , \vec{U} , and \vec{W} are slowly time-varying. It has already been assumed that \vec{W} is. \vec{X} is also slowly changing on the time scale of interest here. But \vec{U} need not be so constrained. \vec{U} , the control vector, is a quantity that ultimately will be minimized. Since U ideally will be near zero we will use $\vec{X} = \vec{X}(t0)$, $\vec{U} = \vec{0}$ and $\vec{W} = \vec{W}_{SC}^{C}(t0)$ as a nominal operating point. Assuming constant case velocity, equation (6.1) can be approximated by

$$\Delta \vec{x} = (d\vec{r}/d\vec{x}) \Delta \vec{x} + (d\vec{r}/d\vec{u}) \Delta \vec{u}$$
 (6.13)

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$$\dot{\vec{X}}(t) = \dot{\vec{X}}(t0) \approx (d\vec{r}(\vec{X}, \vec{U}, \vec{W})/d\vec{X}) | \{\vec{X}(t) = \vec{X}(t0)\} \\ | \vec{X} = \vec{X}(t0) \\ | \vec{U} = \vec{U}_{C} \\ | \vec{W} = \vec{W}_{SC} \\ + (d\vec{r}(\vec{X}, \vec{U}, \vec{W})/d\vec{U}) | \vec{U}(t) \\ | \vec{X} = \vec{X}(t0) \\ | \vec{U} = \vec{O}_{C} \\ | \vec{W} = \vec{W}_{SC} \end{cases}$$
(6.14)

Equation (6.14) can be rewritten as:

$$\vec{X}(t) = A \vec{X}(t) + B \vec{U}(t) + \vec{X}(t0) - A \vec{X}(t0)$$
 (6.15)

Where
$$A = d\vec{f}(\vec{X}, \vec{0}, \vec{W})/d\vec{X}$$
 (6.16)

$$B = df(\overline{X}, \overline{U}, \overline{W})/d\overline{U}$$
 (6.17)

Equation (6.15) is a linear continuous-time approximation to the four-gimbal system. Computation of A and B is extremely complex. Unfortunately, we do not have at our disposal a computer that can exactly simulate in a finite amount of time the continuous behavior of the system that is implicit in (6.15). It is appropriate to ask what the state of the system will be at time t0 + Δ t given the state and control at time t0. Simulating samples of the state will relieve the computational burden required for a continuous solution. Assuming $\vec{U}(t)$ to be constant in the interval [t0, t] and A to be nonsingular, the solution to the dynamical equation (6.15) is:

$$\vec{X}(t) = e^{A} (t-t0) \vec{X}(t0) + A^{-1} [e^{A} (t-t0) - I] [B \vec{U} + \vec{X}(t0) - A \vec{X}(t0)]$$
(6.18)

Equation (6.18) can be differentiated versus time to check that it does solve the dynamical equation. Plugging in t0 for t allows us to check the initial conditions, too.

$$\dot{\vec{X}}(t) = Ae^{A(t-t0)} \vec{X}(t0) + A^{-1}Ae^{A(t-t0)} [B \vec{U} + \vec{\vec{X}}(t0) - A \vec{X}(t0)]$$

$$= Ae^{A(t-t0)} \vec{X}(t0) + e^{A(t-t0)} [B \vec{U} + \vec{\vec{X}}(t0) - A \vec{X}(t0)]$$

$$= [B \vec{U} + \vec{\vec{X}}(t0) - A \vec{X}(t0)] + [B \vec{U} + \vec{\vec{X}}(t0) - A \vec{X}(t0)]$$

$$= A\{e^{A(t-t0)} \vec{X}(t0) + A^{-1}[e^{A(t-t0)} - I][B \vec{U} + \vec{\vec{X}}(t0)$$

$$= A \{\vec{X}(t0)\}\} + B \vec{U} + \vec{\vec{X}}(t0) - A \vec{X}(t0)$$

$$= A \vec{X}(t0) + B \vec{U} + \vec{\vec{X}}(t0) - A \vec{X}(t0)$$
(6.19)

$$X(t0) = e^{X(t0-t0)}X(t0) + A^{-1}[e^{X(t0-t0)} - I][B \overline{U}$$

+ $\dot{\overline{X}}(t0) - A \overline{X}(t0)]$
= $\overline{X}(t0)$ (6.20)

Denoting t by t0 + Δ t, equation (6.18) can be rewritten as:

$$\vec{X}(t0+\Delta t) = e^{A\Delta t}\vec{X}(t0) + A^{-1}[e^{A\Delta t} - I][B\vec{U} + \vec{X}(t0) - A\vec{X}(t0)]$$

= $e^{A\Delta t}\vec{X}(t0) + A^{-1}[e^{A\Delta t} - I]B\vec{U}$
+ $A^{-1}[e^{A\Delta t} - I][\vec{X}(t0) - A\vec{X}(t0)]$ (6.21)

Equation (6.21) can be put in discrete form as:

$$-\vec{X}[n+1] = A^{*}\vec{X}[n] + B^{*}\vec{U}[n] + \vec{C}^{*}$$
(6.22)
$$A^{*} = e^{A\Delta t}$$
(6.23)

Where

$$A^{T} = e^{A\Delta t}$$
(6.23)
$$A^{T} = e^{A\Delta t}$$
(6.23)

$$B^{-} = A^{-1} [e^{A \Delta c} - I]B \qquad (6.24)$$

$$C = A^{-1}[e^{A\Delta t} - I][\vec{X}(t0) - A \vec{X}(t0)]$$
 (6.25)

Because Δt is assumed small equations (6.23) through (6.25) may be approximated to second order:

$$A^{\#} = I + A_{\Delta}t/11 + A^{2}_{\Delta}t^{2}/21$$
 (6.26)

$$B^{*} = A^{-1}[I + A_{\Delta}t/1! + A^{2}_{\Delta}t^{2}/2! - I]B$$

= [I_{\Delta}t + A_{\Delta}t^{2}/2!]B (6.27)
$$C^{*} = [I_{\Delta}t + A_{\Delta}t^{2}/2!][\vec{X}(t0) - A \vec{X}(t0)]$$
(6.28)

A cost function will be used as a measure of performance. The performance index will be a quadratic function of those parameters to be minimized by the controller. They are motor torques, inertial platform misorientation as sensed by the gyros, inertial platform rotation rate, and proximity to gimbal lock. Only proximity to gimbal lock remains to be expressed mathematically.

The angle from the inner gimbal x-axis, X_I , to the outer gimbal plane defined by X_O and Y_O is a convenient measure of gimbal lock. This angle is called λ . λ can be shown to obey the following equation:

$$sin\lambda = sinB sin\Theta$$
 (6.29)

Gimbal lock occurs when λ equals ± 90 degrees. Equation (6.29) requires that both B equal ± 90 degrees and θ equal ± 90 degrees for this to happen. In keeping with a philosophy of linearizing and sampling the equations, the gimbal lock contribution to performance is approximately:

 $sin\lambda[n+1] = cosB sin\Theta \Delta B + sinB cos\Theta \Delta \Theta + sin\lambda[n+1] (6.30)$

The inertial platform rotation rates can be handled in similar fashion.

$$W_{EX}[n+1] = W_{EX}[n] + \sum_{i} (dW_{EX}/dX_{i}) \Delta X_{i}$$
 (6.31)

$$W_{EY}[n+1] = W_{EY}[n] + \sum_{i} (dW_{EY}/dX_{i}) \Delta X_{i}$$
 (6.32)

$$W_{EZ}[n+1] = W_{EZ}[n] + \sum_{i} (dW_{EZ}/dX_{i}) \Delta X_{i}$$
 (6.33)

Equations (6.30) through (6.33) can be combined into a single equation.

$$\vec{e}[n+1] = D[n] \vec{X}[n+1] + \vec{E}[n]$$
 (6.34)

Where $\overline{e}[n+1] = (\sin\lambda[n+1], \Delta SR[n+1], \Delta J[n+1], \Delta SV[n+1],$

 $W_{EX}[n+1], W_{EY}[n+1], W_{EZ}[n+1])$ (6.35) D[n] is a matrix of derivatives with respect to state $\vec{E}[n]$ is a vector containing those terms in (6.30)

through (6.33) not explicitly dependent on X[n+1]

A quadratic cost function was chosen because of a desire to penalize large misorientations of the inertial platform over small ones. Perhaps it would be more appropriate to minimize the maximum torque rather than minimize the RMS torque, but the latter approach is compatible with a quadratic cost function and is certainly more tractable. The one-step performance index is given by:

$$J[n] = \vec{e}[n+1]^{T}Q \vec{e}[n+1] + \vec{U}[n]^{T}R \vec{U}[n]$$
 (6.36)

R is a positive definite symmetric matrix reflecting the cost associated with any control

The cost function in equation (6.36) can be rewritten using matrix trace. Equations (6.34) and (6.22) can then be used to express the cost in terms of \vec{U} .

$$J[n] = \vec{e}[n+1]^{T}Q \vec{e}[n+1] + \vec{U}[n]^{T}R \vec{U}[n]$$
(6.36)
= Tr{Q $\vec{e}[n+1] \vec{e}[n+1]^{T} + R \vec{U}[n] \vec{U}[n]^{T}$ }
= Tr{Q $(D[n](A^{*}\vec{X}[n] + B^{*}\vec{U}[n] + \vec{C}^{*}) + \vec{E}[n])$
 $(D[n](A^{*}\vec{X}[n] + B^{*}\vec{U}[n] + \vec{C}^{*}) + \vec{E}[n])^{T} + R \vec{U}[n]\vec{U}[n]^{T}$ }
(6.37)

Applying the Matrix Minimum Princ_ple[1,2] and taking the gradient of equation (6.37) with respect to \overline{U} yields:

$$dJ/d\overline{U} = 2(B^{*T}D[n]^{T}Q D[n] B^{*} + R) \overline{U}$$

+ $2B^{*T}D[n]^{T}Q (D[n](A^{*}\overline{X}[n] + \overline{C}^{*}) + \overline{E}[n])$ (6.38)

Setting equation (6.38) to $\overline{0}$ and solving for $\overline{10}$ while keeping in mind that things are really dependent on n gives:

$$\vec{v}_{opt} = -(B^{*T}D^{T}Q D B^{*} + R)^{-1}B^{*T}D^{T}Q (D(A^{*}\vec{X} + \vec{C}) + \vec{E}) (6.39)$$

This can be expressed as:

$$\vec{U}_{opt} = K_1 \vec{X} + \vec{K}_2$$
 (6.40)

Where
$$K_1 = -(B^{*T}D^{T}Q D B^{*} + R)^{-1}B^{*T}D^{T}Q D A^{*}$$
 (6.41)

$$\vec{K}_{2} = -(B^{*T}D^{T}Q D B^{*} + R)^{-1}B^{*T}D^{T}Q (D\vec{C}^{*} + \vec{E})$$
 (6.42)

Equations (6.40) through (6.42) immediately suggest an implementation like that depicted in Figure 6.2. Here the state vector is multiplied by gain matrix K_1 to produce an intermediate control signal. The intermediate control is corrected by \vec{K}_2 before driving the actuators. K_1 and \vec{K}_2 are functions of the state, so there are two feedback loops operating here.

Alternatively equation (6.40) can be written as:

$$\vec{v}_{opt} = K_3(\vec{K}_4 - \vec{X})$$
 (6.43)

Where
$$K_3 = (B^*T D^T Q D B^* + R)^{-1} B^* D^T Q D A^*$$
 (6.44)

 $\vec{K}_{4} = -A^{*-1} (D^{T} (D D^{T})^{-1} \vec{E} + \vec{C}^{*})$ (6.45)

These equations, although representing the same system as (6.40) through (6.42) suggest a different implementation shown in Figure 6.3. The controller should behave the same way regardless of which implementation is chosen. It is obvious that a great deal of effort is required to compute $K_1, \vec{K_2}$ or $K_3, \vec{K_4}$ since they are complicated functions of complicated functions. Their calculation poses an immense computational burden. Some way should be found to reduce the amount of work necessary.

One method is to update $K_1, \vec{K_2}$ or $K_3, \vec{K_4}$ less often. The relatively simple calculation of the control vector could be performed very frequently whereas it might be possible to update the gain matrix at a lower rate without sacrificing either performance or stability. Such an analysis has yet to be undertaken.

Another strategy for coping with the complexity of

computation would be to simplify the torque equations by ignoring high order effects. This would hopefully not degrade performance, but might enable more frequent calculation of the gain matrix and offset vector. Carried to an extreme one could ignore everything in the Y's except for torque, and approximate L and M by constant matrices. Simplifications will get propagated through A, B, A^* , B^* , \overline{C}^* etc. leading to more tractable formulas for the K's. In practice, some combination of both strategies may be most feasible.



Figure 6.1

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Gimbal Lock Angle





Proposed Controller Configuration 1





Proposed Controller Configuration 2

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VII. Results

Embedded in the R and Q matrices of the previous section four parameters called TORQWT, LOCKWT, TILTWT and RATEWT. are They are the weights assigned to torque motor control signals. gimbal lock proximity, inertial platform tilt and inertial platform rate respectively in the cost function. These weights were not assigned in any specific fashion. Rather, a trial and error approach was taken to get results that look good. The simulation was run with various values for the weights and performance was judged on the basis of low control voltage, gimbal lock avoidance, small inertial platform tilts and rates, and stability of the controller. The four parameters were tweaked until the controller exhibited the desired behavior. It may be possible to further improve performance by further refining the weights but it is not clear that any significant amelioration will result. In any event, the cost function weights were not chosen in any formal way.

Before examining the performance of the optimal gimbal controller let see what it replaces. The currently us implemented controller uses a zone control scheme. In this the B-O plane is divided into 16 regions (Figure 7.1). scheme Torque motor control signals are generated based on the current The idea is to steer clear of gimbal lock by staying zone. within the numbered zones and avoiding those that include the gimbal lock condition. This is done by driving the elevation and inner gimbals from two of the gyros, and using the third gyro to control either the middle or outer gimbal depending on the zone. The remaining redundant gimbal is used to assist in some sensible fashion. Essentially it is a three-gimbal controller modified for an extra gimbal. Additionally, two of the physical gyros are replaced in the controller by "computed" gyros. The computed gyros, $\triangle R$ and $\triangle V$, lie in the same plane as $\triangle SR$ and $\triangle SV$. However, they point in the same direction as X_I and Z_I respectively. Equation (3.4) can be used to show that:

$$\Delta R = \Delta SR \cos \mathbf{I} - \Delta SV \sin \mathbf{I} \tag{7.1}$$

$$\Delta V = \Delta SR \sin I + \Delta SV \cos I \qquad (7.2)$$

The zone control works fairly well until a zone switch is necessary. When a zone switch occurs, large transient effects arise. Large torque levels may be required to keep the platform inertial. Inertial platform misorientations are greatest immediately following zone changes. The decision rules are:

- Zones 1-4 AR drives the middle gimbal (Iactual - Icommanded) drives the outer gimbal
- Zones 5-8 AR drives the middle gimbal sinB drives the outer gimbal
- Zones 9-12 AR drives the outer gimbal sin0 drives the middle gimbal
- All zones $\triangle J$ drives the elevation gimbal $\triangle V$ drives the inner gimbal

The next several pages compare the optimal controller with the zone control over a variety of orientations and case rates. In all examples the optimal controller exhibits much smaller gyro errors, plus lower RMS and peak torques while avoiding gimbal lock at least as well as the zone control. It wouldn't be optimal otherwise! Much of the apparent advantage of the optimal controller stems from the elimination of zone switch transients. Examples provided courtesy of H. M. Jones. For all examples the time between control updates is 5 milliseconds for the optimal controller, whereas the zone control is simulated as a continous system using a fourth order Runge-Kutta numerical integration technique with a time interval of 1 millisecond.



Zone Control Zones

Figure 7.1

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Table 7.1

Optimal vs. Zone Control Run 1

Case Rates (deg/sec)	Initial Angles (deg)	
Roll -30.0	<i>b</i> 0.0	
Pitch 0.0	θ 0.0	
Yaw -90.0	B 60.0	
	Y 0.0	
<u>Peak Torque</u> (ft-lbs)	Optimal	<u>Zone</u>
Ø	0.328	0.601
0	0.203	0.460
B	0.172	0.160
I .	0.223	0.121
RMS Torque (ft-1bs)		
Ø	0.111	0.218
θ	0.112	0.135
В	0.100	0.100
I	0.090	0.090
<u>Peak Gyro Errors</u> (mill:	iradians)	
⊾SR	0.03	0.51
ΔJ	0.26	0.42
∆SV	0.06	0.38

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Figure 7.2

Optimal vs. Zone Control Trajectory 1

Table 7.2

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Optimal vs. Zone Control Run 2

<u>Case Rates</u> (deg/sec)	<u>Initial Angles</u> (deg)		
Roll 0.0	Ø 0.0		
Pitch 0.0	0 45.0		
Yaw -90.0	B 90.0		
	Y 0.0		
<u>Peak Torque</u> (ft-lbs)	<u>Optimal</u>	Zone	
Ø	0.0	0.0	
0	0.258	0.763	
В	0.0	0.0	
I	0.196	0.117	
RMS Torque (ft-lbs)			
Ø	0.0	0.0	
G	0.119	0.413	
В	0.0	0.0	
Σ.	0.090	0.090	
<u>Peak Gyro Errors</u> (milliradians)			
∆SR	0.0	0.0	
۷Ţ	0.098	0.436	
asv	0.0	0.0	

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Figure 7.3

Optimal vs. Zone Control Trajectory 2

Table 7.3

Optimal vs. Zone Control Run 3

<u>Case</u> Rate	<u>es</u> (deg/sec)	<u>Initial Angles</u> (deg)	
Roll	0.0	Ø 135.0	
Pitch	0.0	9 4.1	
Yaw	90.0	B 41.2	
		¥ 0.0	
Peak Toro	<u>lue</u> (ft-lbs)	<u>Optimal</u>	<u>Zone</u>
Ø		0.395	1.13
θ.		0.133	0.673
В		0.322	0.209
Y .		0.131	0.115
RMS Torqu	ue (ft-lbs)		
Ø		0.150	0.532
0		0.110	0.194
В	•	0.101	0.101
I		0.090	0.090
<u>Peak Gyro Errors</u> (milliradians)			
⊿SR		0.133	0.460
$ abla_{j} $		0.049	0.260
∆ SV		0.117	0.518



----- Optimal

Figure 7.4

Optimal vs. Zone Control Trajectory 3

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Table 7.4

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Optimal vs. Zone Control Run 4

<u>Case Rates</u> (deg/sec)	<u>Initial Angles</u> (deg)		
Roll -90.0	Ø 0.0		
Pitch 0.0	θ 0.0		
Yaw -90.0	Β 0.0		
	¥ 0.0		
<u>Peak Torque</u> (ft-lbs)	<u>Optimal</u>	<u>Zone</u>	
Ø	0.75	1.63	
e	0.36	1.23	
B	0.22	0.138	
I	0.26	0.157	
<u>RMS Torque</u> (ft-lbs)			
Ø	0.169	0.238	
e	0.075	0.192	
В	0.082	0.100	
I ·	0.015	0.085	
<u>Peak Gyro Errors</u> (milliradians)			
⊿SR	0.05	1.06	
$ abla_1 $	0.18	0.84	
∆ SV	0.05	0.08	

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— Optimal

Figure 7.5

Optimal vs. Zone Control Trajectory 4

Table 7.5

Optimal vs. Zone Control Run 5

<u>Case Rates</u> (deg/sec)	<u>Initial Angles</u> (deg)		
Roll 0.0	Ø 180.0		
Pitch 0.0	0 -105.0		
Yaw 90.0	B 43.5		
	Y 0.0		
<u>Peak Torque</u> (ft-lbs)	<u>Optimal</u>	Zone	
Ø	0.939	1.64	
9	1.24	1.23	
B	0.234	0.184	
Y .	0.415	0.177	
<u>RMS Torque</u> (ft-1bs)			
Ø	0.276	0.503	
9	0.174	0.173	
В _	0.100	0.099	
I	0.093	0.091	
<u>Peak Gyro Errors</u> (milliradians)			
∆SR	0.029	0.411	
۵J	0.305	0.902	
∆SV	0.086	1.440	





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Optimal vs. Zone Control Trajectory 5

Table 7.6

Optimal vs. Zone Control Run 6

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<u>Case Rat</u>	<u>tes</u> (deg/sec)	Initial Angles (deg)	
Roll	0.0	Ø 135.0	
Pitch	0.0	θ 0.0	
Yaw	90.0	B -45.0	
		I 0.0	
Peak To	<u>rque</u> (ft-lbs)	<u>Optimal</u>	<u>Zone</u>
Ø		0.637	1.29
Ð		0.230	1.23
B		0.301	0.138
I		0.183	0.111
RMS Tor	<u>que</u> (ft-lbs)		
Ø		0.252	0.390
e		0.110	0.300
В	-	0.099	0.087
I		0.089	0.081
<u>Peak Gyro Errors</u> (milliradians)			
⊿SR		0.029	0.887
∠J		0.139	0.236
₽SA		0.117	0.232


Figure 7.7

Optimal vs. Zone Control Trajectory 6

Table 7.7

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Optimal vs. Zone Control Run 7

<u>Case</u>	<u>Rates</u> (deg/sec)	<u>Initial Angles</u> (deg)	·
Roll	0.0	Ø 180.0	
Pitch	0.0	θ 0.0	
Yaw	90.0	B -45.0	
		I 0.0	
<u>Peak</u>	<u>Torque</u> (ft-lbs)	<u>Optimal</u>	<u>Zone</u>
Ø		0.637	1.10
e		0.230	0.401
В		0.301	0.158
Y		0.183	0.149
<u>RMS 1</u>	<u>forque</u> (ft-1bs)		
Ø		0.252	0.430
9		0.110	0.123
В	-	0.099	0.085
Y		0.089	0.081
<u>Peak</u>	<u>Gyro Errors</u> (mil	liradians)	
⊾SR		0.029	0.278
ΔJ		0.139	0.846
⊿ SV		0.117	0.312
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Figure 7.8

Optimal vs. Zone Control Trajectory 7

VIII. Conclusions

Modern optimal control provides a useful framework in which to analyze and improve the performance of feedback systems. Many applications exist for this untapped powerful theory. Unfortunately, it is not always used to advantage. This thesis has attempted to relieve this situation for one particular Simulation studies indicate great success. system. The optimal controller for a four-gimbal system potentially far outperforms an earlier nonoptimal controller.

This improvement in performance does not come free. A significant computational burden is imposed by optimization. Some techniques for reducing the load have been suggested. Work remains to be done actually implementing the proposed controller. Final judgement on its feasibility awaits.

There is no reason to be content even with an optimal controller. Under different optimality criteria it is conceivable that a controller could be designed with more desirable operating characteristics. A bang-bang controller is one worth considering. By applying full torque in short pulses it may be possible to further reduce platform tilts.

Leaving such speculation aside, the fact remains that with a suitable model developed, optimal control can be applied to components of inertial guidance equipment. One can only hope that deployment precludes actual use.

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Appendix A. Coordinate Transformations and Notation

The notation used here is based on work by Britting[3]. This notation is helpful for representing orientations, rotations and coordinate transformations. The reference frame of a vector is indicated by a superscript. \vec{r}^j is a vector coordinatized in the j reference frame. Any vector in the j frame can be expressed in the k frame by premultiplying the original vector by a coordinate transformation C_j^k . The subscript indicates the original reference frame and the superscript denotes the new reference frame. Thus, for the example given:

$$\vec{r}^{k} = C_{j}^{k} \vec{r}^{j}$$
 (A.1)

Note that the original superscript has been canceled by the subscript of C_j^k . For Cartesian coordinate systems, in which the basis vectors are orthonormal, the entries of a coordinate transformation matrix are direction cosines. Direction Cosine Matrix (DCM) is a term often used to describe such a matrix. The direction cosine from the m-axis of reference frame j to the n-axis of frame k is the mnth entry of C_j^k . DCM's exhibit many interesting properties. Some follow:

$$c_{k}^{l}c_{j}^{k} = c_{j}^{l} \text{ but } c_{j}^{k}c_{k}^{l} \neq c_{j}^{l}$$
(A.2)

$$C_{j}^{J} = I \tag{A.3}$$

$$C_{j}^{R} = (C_{k}^{J})^{-1}$$
(A.4)

$$C_{j}^{n} = (C_{k}^{j})^{T}$$
(A.5)

Rotations satisfy the same superscript convention as other vectors. In addition rotation vectors have two subscripts. The sense of rotation is from the left subscript to the right subscript. To be precise, coordinate systems rotate, not subscripts. \overline{W}_{kj}^1 would be the rotation rate of system j with respect to system k as seen from the l reference frame. Rotations add vectorially. When they do, subscripts cancel.

$$\vec{W}_{ki}^{l} = \vec{W}_{kj}^{l} + \vec{W}_{ji}^{l}$$
(A.6)

It follows that:

$$\vec{W}_{kj}^{l} = -\vec{W}_{jk}^{l}$$
(A.7)

The superscripts must be the same for these relations to hold. Differention of vectors is no longer simple in rotating reference frames. For any vector \vec{r}^i we have the following equivalent expressions:

$$\vec{r}^{i} = c_{j}^{i}\vec{r}^{j} - c_{j}^{i}(\vec{w}_{ji}^{j} \times \vec{r}^{j})$$

$$= c_{j}^{i}\vec{r}^{j} - c_{j}^{i}\vec{w}_{ji}^{j} \times c_{j}^{i}\vec{r}^{j}$$

$$= c_{j}^{i}\vec{r}^{j} - \vec{w}_{ji}^{i} \times \vec{r}^{i}$$

$$= c_{j}^{i}\vec{r}^{j} + \vec{w}_{ij}^{i} \times \vec{r}^{i} \qquad (A.8)$$

Appendix B. Summary of Computer Routines Used

Main Program

1. Calls INITLZ routine

2. Calls DERIVE routine

3. Calls OUTPUT routine

4. Calls UPDATE routine

5. Loops to 2.

INITLZ (X, DXDT, W, OLDRATE, TOTALT, DELTAT)

1. Clears out storage areas

2. Initializes state, case rates and other parameters

OUTPUT (X, DXDT, U, W, I, TIME)

1. Prints output 1 out of J invocations else returns

2. Prints state, derivative, control and case rate vectors

UPDATE (X, DXDT, U, W, OLDRATE, DELTAT)

1. Updates state via 4th order Runge-Kutta Integration

2. Calls DERIVE during computation

DERIVE (X, DXDT, U, W, OLDRATE)

1. Computes friction as described in section V.

2. Derives torque-acceleration equations as per section VI.

3. Solves for angular accelerations using SIMQ

4. Returns state derivative in DXDT

SIMQ (A, B, N, KS)

1. Solves system of equations of form AX=B

2. Returns solution in B

MINV (A, N, D, L, M)

1. Inverts a matrix

2. Returns result in A

MATMPY (A, B, C)

1. Computes C=AB

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CONTRL (X, U, W, OLDRATE, DELTAT)

Computes linear discrete-time equations as in section VII
 Calls MINV and MATMPY to perform matrix manipulations

3. Returns control in U

Appendix C. Computer Simulation of the Four-Gimbal System

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000100 SIMULAG -- 4 GIMBAL SYSTEM SIMULATION MICHAEL A. GENNERT 004200 IMPLICIT DOUBLE PRECISION(A-H,0-Z) 00030 PARAHETER IXX=11, IDR=7 DIMENSION X(11), DXDT(11), U(4), W(3), OLDRATE (4) 00040 00050 00060C DELTAT = TIME INCREMENT SIZE 001700 TOTALT = TOTAL SIMULATION TIME X = STATE VECTOR DXDT = TIME RATE OF CHANGE OF STATE VECTOR 0003080 809900 ÖÖIDÖČ = CONTROL VECTOR 001100 11 W = CASE RATE VECTOR OLDRATE= VECTOR CONTAINING PREVIOUS VALUES OF GIMBAL ANGLE RATES 00120C 001380 00140C USED TO DETERMINE FRICTION EFFECTS IN DERIVE ROUTINE 001500 00160C 00170C X(1) X(2) X(3) = PSI(E) = BETA Ū 001800 = THETA (\aleph) X(4) X(5) 00190C = PHI **(f)** = PSIDOT (dE/dt) 002000 00210C 00220C X(6) X(7) = BETADOT (dI/dt) = THETADOT (dH/dt) 002300 X(8) = PHIDOT (d0/dt) 002400 X(9) = DELTASR X(10) = DELTAJX(11) = DELTASV002500 00260C 002700 002800 U(f) = CONTROL ON E GIHBAL = CONTROL ON I GINGAL = CONTROL ON M GINBAL 002900 Ū(2) **90300**C Ū(3) = CONTROL ON O GINBAL 003100 U(4) 093280 003300 W(1) W(2) W(3) = X AXIS CASE RATE (VCX) = Y AXIS CASE RATE (VCX) = Z AXIS CASE RATE (VCZ) 00350Č #0360C 00370C 00380C INITIALIZE 003966 CALL INITLZ (X, DXDT, ₩, OLDRATE, TOTALT, DELTAT) N = TOTALT/DELTAT DO 100 I = 1, N+2 00400 00410 00420 004300 00440C DETERMINE CONTROL SIGNAL TO BE APPLIED 00450C 00460 CALL CONTRL (X, U, W, OLDRATE, DELTAT) 00470C 09480C COMPUTE STATE DERIVATIVE VECTOR PRIOR TO DUTPUT 00490C 00500 CALL DERIVE (X, DXDT, U, W, OLDRATE) 005100 00520C 00530C 00540 PRINT STATE AND CONTROL INFORMATION II = I - iCALL OUTPUT (X, DXDT, U, W, II, DBLE(FLOAT(II))*DELTAT) 00550 00560C 00570C UPDATE STATE EQUATIONS AND INTEGRATE 005800 00590 CALL UPDATE (X, DXDT, U, W, OLDRATE, DELTAT) 80698 180 CONTINUE 00650 00620 STOP END 006300 00640C 00650C INITIALIZE SUBROUTINE SUBROUTINE INITLZ (X, DXDT, W, OLDRATE, TOTALT, DELTAT) . IMPLICIT DOUBLE PRECISION(A-H,O-Z) 00660 00670 PARAMETER IXX=11 **\$**068**\$**

00690 DIMENSION X (1), DXDT (1), W (1), OLDRATE (1) XK = 3.14159265358979323846D0/180.D0 09700 00710C CLEAR STATE VECTOR 007200 09730C 00740 DO 100 I = 1, IXX X(I) = 0.D0 100 CONTINUE 00750 00760 007700 007800 E, I, M, O FOLLOW IN DEGREES 80790C 00800 X (1) = 0.D02XK 00810 X(2)=60.DD*XK X(3)=0.D0#XK X(4) = 0.D0#XK 00820 00830 008400 dE/dt, dI/dt, dH/dt, dD/dt FOLLOW IN DEGREES/SECOND 008500 008600 00876 X (5) = 0.D0XKX(6)=30,D0#XK 00880 X(7)=90.D0#XK 00890 00900 X(8) = 0.D0#XX 009100 00920C SET OLDRATE TO ANGLEDOT FOR FRICTION COMPUTATION 009300 $\begin{array}{l} DO \ ii0 \] = i, \ 4 \\ OLDRATE \ (I) \ = \ X \ (I \ + \ 4) \end{array}$ 08940 00950 00960 110 CONTINUE 00970C 00980Č CALCULATE CASE RATES IN KADIANS/SECOND 009900 SB = DSIN (X (2)) CB = DCOS (X (2)) ST = DSIN (X (3))11001 01010 01020 01030 CT = DCOS(X(3))SF = DSIN (X (4))CF = DCOS (X (4))WIY = -X (5)01040 01050 01060 WHX = -SEXWIY 01070 WMY = CB*WIYWMZ = -X (6) 01080 01090 01100 WOX = WHX - X (7) WOY = CT#WHY-ST#WHZ 01110 01129 HOZ = STARHY+CTARHZ 01130 01140 $W(1) = CF \pm WO X + SF \pm WO Z$ W(2) = WO Y - X(8)W(3) = -SF WOX + CF WOZ01150 01160C 01170C SET UP TIME PARAMETERS IN SECONDS 01180C 01198 TOTALT = 1.DODELTAT=1.D0/3000.D0 01200 01210 01220 RETURN END 01230C 01240C 01250C 01260 OUTPUT SUBROUTINE SUBROUTINE OUTPUT (X, DXDT, U, W, I, TINE) IMPLICIT DOUBLE PRECISION(A-H, 0-Z) 01270 01280 DIMENSION X (1), DXDT (1), U (1), W(1) J=30IPRINT = 6 01290 01300 01310 XK = 3.14159265358979323846D8/180.D0 013210 PRINT EVERY Jth TIME, RETURN THE OTHER J-1 OCCURANCES 01339C 013400 01350 IF (I.NE.(I/J)#J) RETURN 01360 WRITE (IPRINT, 900) TIME

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900 FORMAT (//" TIME =", F8.3, " SECONDS")
WRITE (IPRINT,901)
901 FORMAT (18X, "DEG", 11X, "DEG/SEC", 7X, "DEG/SEC/SEC", 7X, "CONTROL")
WRITE (IPRINT,902) "E", X (1)/XK, X (5)/XK, DXDT (5)/XK, U (1)
902 FORMAT (6X, 1A1, 4(6X, F8.3))
WRITE (IPRINT,902) "1", X (2)/XK, X (6)/XK, DXDT (6)/XK, U (2)
WRITE (IPRINT,902) "H", X (3)/XK, X (7)/XK, DXDT (6)/XK, U (2)
WRITE (IPRINT,902) "0", X (4)/XK, X (7)/XK, DXDT (7)/XK, U (3)
WRITE (IPRINT,902) "0", X (4)/XK, X (8)/XK, DXDT (8)/XK, U (4)
WRITE (IPRINT,903)
903 FORMAT (18X, "MRAD", 9X, "MRAD/SEC", 25X, "DEG/SEC")
WRITE (IPRINT,904) "SR", X (9)*1E3, DXDT (9)*1E3, "WCX", W (1)/XK
904 FORMAT (6X, "DELTA", A2, F10.3, 6X, F10.3, 16X, A3, 5X, F8.3)
WRITE (IPRINT,904) "J", X (10)*1E3, DXDT (10)*1E3, "WCZ", W (2)/XK
WRITE (IPRINT,904) "SV", X (11)*1E3, DXDT (11)*1E3, "WCZ", W (3)/XK
RETURN 01370 900 FORMAT (//" TIME =", F8.3, " SECONDS") 01380 01390 01400 01410 01420 01430 01440 01450 01460 01478 01480 01490 01580 01510 RETURN 01520 END 81530C UPDATE SUBROUTINE 015400 ŎĨŚŚOC CALLS DERIVE WHICH COMPUTES DERIVATIVE, THEN EMPLOYS RUNGE-KUTTA 4th ORDER INTEGRATION 015600 01570C 01580 SUBROUTINE UPDATE (X, DXDT, U, W, OLDRATE, DELTAT) IMPLICIT DOUBLE PRECISION(A-H,O-Z) 01590 PARAMETER IXX=11 DIMENSION X (1), DXDT (1), U (1), W (1), OLDRATE (1) DIMENSION Q (IXX), XSTOR (IXX) 01600 01610 01620 016300 STORE STATE VECTOR IN XSTOR 016400 01650C DO 100 I = 1, IXX XSTOR (I) = X (I) 01660 01670 100 CONTINUE 01680 01690C 01700C 01710C COMPUTE DERIVATIVE AND MAKE 1st APPROXIMATION 01720 CALL DERIVE (XSTOR, DXDT, U, W, OLDRATE) DO iiO I = i, IXX Q (I) = DXDT (I)01730 01740 01750 XSTOR (I) = X (I) + .5D0 # DELTAT # DXDT (I) 110 CONTINUE 01760 01770C 01780C 01790C 2nd APPROXIMATION 01800 CALL DERIVE (XSTOR, DXDT, U, W, DLDRATE) $\begin{array}{l} D G \ 120 \ I \ = \ 1, \ IXX \\ Q(I) \ = \ Q \ (I) \ + \ 2.D0 \ \ DXDT \ (I) \\ XSTOR \ (I) \ = \ X \ (I) \ + \ .SD0 \ \ DELTAT \ \ DXDT \ (I) \\ \end{array}$ 01810 01820 01830 01840 120 CONTINUE 018500 018600 3rd APPROXIMATION 01870C 01880 CALL DERIVE (XSTOR, DXDT, U, W, OLDRATE) DO 130 I = 1, IXX Q (I) = Q (I) + 2.DO \ddagger DXDT (I) XSTOR (I) = X (I) + DELTAT \ddagger DXDT (I) 01890 01900 01918 01920 130 CONTINUE 01930C 01940C FINAL APPROXIMATION 019500 81968 CALL DERIVE (XSTOR, DXDT, U, W, OLDRATE) 01978C 01980C STORE OLD VALUES OF ANGLE RATES FOR FRICTION COMPUTATION 81999C DD 140 I = 1, 4 $OLDRATE_{(I)} = X (I+4)$ 02000 02010 02020 140 CONTINUE DO 150 I = 1, IXX DXDT (I) = (Q (I) + DXDT (I)) / 6.D0 02030 12040

02850 02060 02070 02080	150	X (I) = X (I) + DELTAT ‡ DXDT (I) CONTINUE RETURN END
D2190C D2190C D2110C		DERIVATIVE SUBROUTINE COMPUTES DXDT GIVEN X, U AND W
22130 32130 32140 32150 32150 32150 32150 32150 32150 32150 32150 32170 32180 32190 32219 32220 32220 32240 32250 32240 32250 32280C 32300C 32300C 32300C	i	SUBROUTINE DERIVE (X. DXDT, U. W, OLDRATE) IMPLICIT DOUBLE PRECISION(A-Z) LOGICAL STUCK INTEGER I,II,FLAG,KS DIMENSION X (1), DXDI (1), U (1), W (1), OLDRATE (1) DIMENSION STUCK (4), TORQUE (4), KTR (4), KV (4), FSTATC (4) DIMENSION FCOULM (4) DIMENSION FCOULM (4) DIMENSION FCOULM (4) DATA JE, JIX, JIYZ, JMXZ, JMY, JOXY, JOZ (12D-2, 1.7D-2, 1.3D-2, 2.25D-2, 3.0D-2, 3.0D-2, 3.9D-2/ DATA FSTATC /.09D0, 1D0, 11D0, 165D0 / DATA FCOULM /.09D0, 1D0, 11D0, 165D0 / DATA KTR /1.9D-2, 1.9D-2, 4.1D-2, 5.13D-2/ DATA KTR /1.9D-2, 1.9D-1, 1.25D0/ THIS ROUTINE SOLVES THE FOLLOWING MATRIX EQUATON. Y1 IS A FUNCTION ONLY OF GIMBAL ANGLES, GIMEAL RATES, AND TIEY. Y2 IS A FUNCTION ONLY OF GIMBAL ANGLES, GIMEAL RATES, AND THIZ. Y3 IS A FUNCTION ONLY OF GIMBAL ANGLES, GIMEAL RATES, AND TOMX. Y4 IS A FUNCTION ONLY OF GIMBAL ANGLES, GIMBAL RATES, AND TOMX. Y4 IS A FUNCTION ONLY OF GIMBAL ANGLES, GIMBAL RATES, AND TOMX. Y4 IS A FUNCTION
02330C 02330C 02340C 02350C 02350C		i Yi i i Lii Li2 Li3 Li4 i i PSIDOUBLEDOT i i Y2 i i L2i L22 L23 L24 i i BETADOUBLEDOT i i Y3 i = i L3i L32 L33 L34 i i THETADOUBLEDOT i i Y4 i i L4i L42 L43 L44 i i PHIDOUBLEDOT i
02370C 02380C 02390C		STUCK = ONE FLAG FOR EACH GIMBAL. VALUE IS TRUE IF THE SPECIFIED GIMBAL IS STUCK TO THE NEXT OUTER GIMBAL DUE TO STATIC
02410C 02420C 02420C		FSTATC = STATIC FRICTION TORQUE LIMIT, SPECIFIES THE STATIC FRICTION LEVELS THAT MUST BE OVERCOME TO FREE A STUCK GIMPAL. FCOMM M = COMMONS FRICTION TORQUE LIMIT, SPECIFIES THE FRICTION
02430C 02450C 02450C 02450C 02478C 02478C 02480C 02500C 02510C 02510C 02520C 02530C		TOUGLINCONSTANTSFROM THE GINBALS ARE UNSTUCK.KTR= CONVERSION CONSTANTS FROM TORQUE MOTOR VOLTAGES TO TORQUESKV= PROPORTIONALITY CONSTANTS FROM ANGLEDOTS TO BACK ENFSJE= INERTIA ABOUT ANY AXIS OF THE ELEVATION GIMBALJIX= INERTIA ABOUT THE X AXIS OF THE INNER GIMBALJIX= INERTIA ABOUT THE X AND Z AXES OF THE INNER GIMBALJMXZ= INERTIA ABOUT THE X AND Z AXES OF THE MIDDLE GIMBALJMY= INERTIA ABOUT THE X AND Z AXES OF THE MIDDLE GIMBALJOXY= INERTIA ABOUT THE Y AND Y AXES OF THE OUTER GIMBALJOZ= INERTIA ABOUT THE X AXIS OF THE OUTER GIMBAL
02550C 02560C 02560C		ASSOCIATE VARIABLES WITH ARRAY ELEMENTS COMPUTE REQUIRED TRIGONOMETRIC FUNCTIONS
02580 02590 02590 02610 02610 02610 02620 02630 02650 02650 02650 02680 02680 02680 02700 02710 02710		PSI = X (1) BETA = X (2) THETA = X (3) PHI = X (4) SP = DSIN (PSI) CP = DCOS (PSI) SB = DSIN (BETA) CB = DCOS (BETA) ST = DSIN (THETA) CT = DCOS (THETA) SF = DSIN (PHI) CF = DCOS (PHI) SB2 = SB \$ SB CB2 = CB \$ CB SI2 = ST \$ ST

02730 $CT2 = CT \ddagger CT$ 02740C DEFINE L11 THROUGH L44 FOR THE SIGNIFICANCE OF THESE QUANTIES REFER TO THE GIMBAL TORQUE EQUATION DERIVATIONS 027500 02760C 02770C 02780C L (1,1) = JE L (1,2) = 0.D0 L (1,3) = $-SB \times JE$ L (1,4) = CB $\times CT \times JE$ L (2,1) = 0.D0 L (2,2) = JIYZ + JE L (2,3) = 0.D0 L (2,4) = $-ST \times (JIYZ + JE)$ L (3,1) = $-SR \times JE$ L (3,2) = 0.D0 L (3,3) = (JE + SB2 $\times JIYZ + CB2 \times JIX + JHXZ)$ L (3,4) = SB $\times CB \times CI \times (JIX - JIYZ)$ L (4,4) = SB $\times CB \times CI \times (JIX - JIYZ)$ L (4,4) = SB $\times CB \times CI \times (JIX - JIYZ)$ L (4,4) = J0XY + CI2 $\times JHY + SI2 \times (JHXZ + JIYZ) + SB2 \times CI2 \times JIX$ $\times + CB2 \times CI2 \times JIYZ + JE$ 02790 02808 02810 02820 02830 02840 02850 02860 02870 02880 02890 02909 02910 02920 02930 02940 02950 + CB2 # CT2 # JIYZ + JE Ł 029600 02970C 02970C 02980C 02990 03000 ROUTINE TO CONVERT CONTROL SIGNALS TO TORQUES, INCLUDING FRICTION DO 120 I = 1, 4 TORQUE (I) = KTR (I) * (U (I) - KV (I) * X (I+4)) 03010C IF THE MAGNITUDE OF THE APPLIED TORQUE DOES NOT EXCEED THE STATIC FRICTION LIMIT AND THE GIMBAL RATE IS PASSING THROUGH ZERO (ie. ANGLEDOT CHANGES SIGN) THEN THE GIMBALS WILL BE STUCK TOGETHER 030200 030300 030400 03050C 03060 IF (ABS (TORGUE (I)).LE.FSTATE (I).AND.((X (I+4)*OLDRATE (I)). & LT.0.D0.OR.X(I+4).EQ.0.D0)) GOTO 100 83070 030800 030900 GIMBALS NOT STUCK TOGETHER -- CLEAR STUCK FLAG, SUBTRACT FRICTION 031000 STUCK (I) = .FALSE. TORQUE (I) = TORQUE (I)-SIGN (FCOULM (I), X (I+4)) GOTD 120 03110 03120 03130 03140C GINBALS STUCK TOGETHER -- SET STUCK FLAG, SET ANGLEDOT TO ZERO SET 1th ROW AND 1th COLUMN OF (L) TO ZERO, SET L (I,I) = 1. 031500 031600 031700 03180 100 STUCK (I) = .TRUE. 03190 03200 03210 X (I+4) = 0.00DO 110 II = 1, L (I,II) = 0.D8 L (II,I) = 0.D0 83220 03230 110 CONTINUE 03240C 03250C SET L (I SET L (I,I) = 1. SO AS NOT TO HAVE A SINGULAR MATRIX 032500 SETL(I, 03260C 03270 L(I,I)=i. 03280 120 CONTINUE 03290C 03300C DEFINE GI L(I,I)=1.D0 DEFINE GIMBAL RATES 033100 03320 03330 PSIDOT = X (S) $\begin{array}{l} \text{RETADOT} = X (6) \\ \text{THETADOT} = X (7) \end{array}$ 83340 03350 PHIDOT = X (8) WCX = W (1) 03360 WCZ = W (2) WCZ = W (3) WCX = CF * WCX - SF * WCZ 03371 03390 03370 0**3400** NOY = NCY + PHIDOT

NOZ = SF # NCX + CF # NCZ NHX = NOX + THETADOT 03410 03420 WHY = CT * WOY + ST * WOZ 03430 WHZ =-ST * NOY + CT * NOZ 03440 WIX = CB * WMX + SB * WMY WIY =-SB * WMX + CB * WMY 03450 03469 WIZ = WHZ + BETADOT 03470 WEX = CP * WIX - SP * WIZ WEY = WIY + PSIDOT WEZ = SP * WIX + CP * WIZ 83480 03490 03500 035100 035200 DEFINE TORQUES 835380 03540 TIEY = TORGUE (1) THIZ = TORQUE (2) 03550 TOHX = TORQUE (3) 03560 TCOY = TORQUE (4) 03570 03580C 03596C TORQUE EQUATIONS FOR THE FOUR GIMBALS 03600C Y (1) = TIEY + JE * (-SB * PHIDOT * WOZ - CB * ST * PHIDOT * WOX - CB * THETADOT * WHZ + BETADOT * WIX) Y (2) = THIZ + (JIYZ + JE) * (THETADOT * WHY - CT * PHIDOT * WOX) 03610 03620 ł 03630 + WIX * WIY * (JIX-JIYZ) - PSIDOT * JE * WIX 03640 è Y (3) = TOHX + (JHXZ + CF2 * JIX + SR2 * JIYZ + JE) * PHIDOT * WOZ - SB * CB * (JIX-JIYZ) * (ST * PHIDOT * WOX + THETADOT * WHZ - BETADOT * WHX) + WHY * WHZ * (JHY - JHXZ) - (CB2 * JIX + SB2 * JIYZ + JE) * BETADOT * WHY + CB * JE * PSIDOT 03650 å 03660 03670 Ł 03680 ł \$ # WIZ Y (4) = TCDY - WOX * WOZ * (JOXY - JOZ) - (CT2 * (JHY + SB2 * JIX) * ST2 * (JHXZ + JIYZ) + JE) * THETADOT * WOZ + ST * CT * ST2 * (JHXZ + JIYZ) + JE) * (THETADOT * WOY 03690 83700 03710 * (JHY - JMXZ + SB2 * (JIX-JIYZ)) * (THETADOT * WOY 03720 å - PHIDOT # WOX) + SB * CB * CT * (JIX - JIYZ) * (PHIDOT * WOZ - BETADOT * WAY) + (CT * (SB2 * JIX + CB2 * JIYZ + JE) * BETADOT + ST * WAY * (JMY- JHXZ)) * WAX + SB * CT * JE 03730 å 03740 Ł 03750 Ł * PSIDDT * WIZ + ST * (JE * PSIDDT + WIY * (JIYZ - JIX)) 03760 Ł ¥ WIX 03770 å 03780C 03790C CALL SING TO SOLVE FOR ACCELERATIONS 03800C CALL SING (L, Y, 4, KS) 03810 038200 SET TO ZERO THE ACCELERATION OF ANY GINBAL THAT IS STUCK 03830C 03840C D0 130 II = 1, 4 IF (STUCK (II)) Y (II) = 0.00 03850 03860 530 CONTINUE 83878 038800 SET DXDT TO THE COMPUTED DERIVATIVE 138900 03988C DXDT (1) = PSIDOT 03910 DXDT (2) = BETADOT DXDT (3) = THETADOT 83920 03930 03940 DXDT (4) = PHIDOT DXDT (5) = Y (1) 03950 DXDT (6) = Y (2) DXDT (7) = Y (3) DXDT (8) = Y (4) 03960 = Y (2). 03970 03980 DXDT (9) = WEX DXDT (10) = WEY DXDT (11) = WEZ RETURN 03990 04000 04010 04020 04039 END 048400 SUBROUTINE TO SOLVE SYSTEMS OF SIMULTANEOUS LINEAR EQUATIONS 04050C 84050C SUBROUTINE SING (A, B, N, KS) 04970 IMPLICIT DOUBLE PRECISION(A-H, 0-Z) 84680

DIMENSION A (1), B (1) 04100C SOLVE SET OF EQUATIONS AX=B A = MATRIX OF COEFFICIENTS STORED COLUMNWISE. THESE ARE DESTROYED IN THE COMPUTATION. THE SIZE OF A IS N BY N. B = VECTOR OF ORIGINAL CONSTANTS (LENGTH N). THESE ARE REPLACED BY FINAL SOLUTION VALUES, VECTOR X. N = NUMBER OF EQUATIONS AND VARIABLES. N = NUMBER OF EQUATIONS AND VARIABLES. 04110C 04140C KS= OUTPUT DIGIT. 0 FOR NORMAL SOLUTION, 1 FOR A SINGULAR SYSTEM. TOL = 0.090KS = 0JJ = -NDO 65 J = 1, N JY = J + 1 JJ = JJ + N + .1BIGA = 0.DO IT = JJ - J $\overline{D}\overline{D}$ $\overline{3}\overline{D}$ $\overline{1}$ = \overline{J} , N IJ = II + I04300 IF (ABS (BIGA) - ABS (A (IJ))) 20, 30, 30 20 BIGA = A (IJ) IMAX = I30 CONTINUE IF (ABS (BIGA) - TOL) 35, 35, 40 35 KS = 1 04340 RETURN 40 II = $J + N \neq (J - 2)$ $\begin{array}{l} IT = IMAX - J\\ DO 50 K = J, N\\ I1 = I1 + N \end{array}$ $\begin{array}{l} 12 = 11 + 11 \\ 12 = 11 + 11 \\ SAVE = A (11) \\ A (11) = A (12) \\ A (12) = SAVE \\ \end{array}$ 04420 50 A (II) = A (II) / BIGA SAVE = B (IMAX) B (IMAX) = B (J) B (1MAX) = B (J) B (J) = SAVE / BIGA IF (J - N) 55, 70, 55 55 IQS = N \ddagger (J - 1) D0 65 IX = JY, N IXJ = IQS + IX IT = J - IX -D0 60 JX = JY, N IXJX = N \ddagger (JX - 1) + IX IY = IYIX + IT 04500 $\begin{array}{r} IXJX = N \pm (JX - 1) + IX \\ JJX = IXJX + IT \\ 60 A (IXJX) = A (IXJX) - A (IXJ) \pm A (JJX) \\ 65 B (IX) = B (IX) - B (J) \pm A (IXJ) \\ 70 NY = N - 1 \\ IT = N \pm N \\ D0 80 J = 1, NY \\ IA = IT - J \\ IB = N - J \\ TC = N \end{array}$ ĨČ = ₩ DO BO K = 1, J B (IB) = B (IB) - A (IA) * B (IC) IA = IA - N B0 IC = IC - IRETURN 84700C END SUBROUTINE TO INVERT A MATRIX 04740 SUBROUTINE MINV (A, N, D, L, H) INPLICIT DOUBLE PRECISION(A-H, D-Z) DIMENSION A (1), L (1), H (1) 04760C

INVERT A MATRIX A = INPUT ARRAY, DESTROYED IN COMPUTATION AND REPLACED BY INVERSE N = ORDER OF MATRIX A D = RESULTANT DETERMINANT. A ZERO DETERMINANT INDICATES A 04770C 04760C 047900 048000 SINGULAR MATRIX L = WORK VECTOR OF LENGTH N M = WORK VECTOR OF LENGTH N 048100 048200 048300 04849C 14850 D=1.DO $\frac{D}{D} = \frac{1}{N}$ $\frac{D}{D} = \frac{1}{N}$ $\frac{D}{D} = \frac{1}{N}$ 04860 84870 04881 04890 NK = NK + NL(K) = KH(K) = K04900 04910 KK = NK + K BIGA = A (KK) DO 20 J = K, N IZ = N \$ (J - i) DO 20 I = K, N IJ = IZ + I i0 IF (ABS (BIGA) - ABS (A (IJ))) i5, 20, 20 i5 BIGA = A (IJ) L (K) = I M (K) = T KK = NK + K04920 04930 04940 04950 84960 04970 04980 04990 L (K) = 1 H (K) = J 20 CONTINUE J = L (K) IF (J - K) 35, 35, 25 25 KI = K - N D0 30 1 = 1, N KI = KI + N VOI D = - (KT) 05000 05010 05020 05030 05040 05050 05060 HOLD = -A (KI) JI = KI -K + JA (KI) = A (JI) 30 A (JI) = HOLD 05070 05080 05090 05100 $\begin{array}{l} 30 \ A \ (J1) = HULD \\ 35 \ I = H \ (K) \\ IF \ (I - K) \ 45, \ 45, \ 38 \\ 38 \ JP = N \ (I - 1) \\ DD \ 40 \ J = 1, \ N \\ JK = NK + J \\ JI = JP + J \\ HOLD = -A \ (JK) \\ A \ (JK) = A \ (JI) \\ 40 \ A \ (JI) = HOLD \end{array}$ 05119 05120 05130 05140 05150 05160 05170 05180 40 A (JI) = HOLD 45 IF (DIGA) 48, 46, 48 05190 05200 05210 05220 05230 05240 05240 05260 05260 05260 46 - D = 0.00RETURN 48 DO 55 I = 1, N IF (I - K) 58, 55, 50 50 IK = NK + I A (IK) = A (IK) / (-BIGA) 55 CONTINUE 55 CONTINUE 05288 05290 DO 65 I = 1, NIK = NK + I05300 05310 HOLD = A (IK)IJ = I - N DO 65 J = 1, N 05320 UU 05 J = 1, H IJ = IJ + N IF (I - K) 68, 65, 60 68 IF (J - K) 62, 65, 62 62 KJ = IJ - I + K A (IJ) = HOLD \$ A (KJ) + A (IJ) 65 CONTINUE KT = K - M 05330 05340 05350 05360 05370 05380 05390 05400 KJ = K - N NJ - KJ = 1, N KJ = KJ + N IF (J - K) 70, 75, 78 70 A (KJ) = A (KJ) /BIGA 05418 05420 \$5430 05440 **75 CONTINUE**

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 $D = D \ddagger BIGA$ A(KX)=1.D0/BIGA 80 CONTINUE K = N 100 K = (K - 1) IF (K) 150, 150, 105 105 I = L (K) IF (I - K) 120, 120, 108 108 JQ = N # (K - 1) JR = N # (I - 1) DO 110 J = 1, N JK = JQ + J HOLD = A (JK) JI = JR + J A (JK) = -A (JI) 110 A (JI) = HOLD 120 J = M (K) IF (J - K) 100, 100, 125 125 KI = K - N DO 130 I = 1, N KI = KI + N HOLD = A (KI) JI = KI - K + J A (KI) = -A (JI) (KI = KI - K + J A (KI) = -A (JI) **80 CONTINUE** 05510 05540 05550 05580 05620 05650 05660 A(KI) = -A(JI)A (JT) = HOLD GO TO 100 150 RETURN END MATRIX MULTIPLICATION SUBROUTINE **95750C** 05760 05780 05790C SUBROUTINE MATMPY (A. B. C. L1, L2, M1, M2, N1, N2, IFLAG) IMPLICIT DOUBLE PRECISION(A-H, 0-Z) DIMENSION A (1), B (1), C (1) 05800C A = L1 BY L2 INPUT MATRIX B = M1 BY M2 INPUT MATRIX C = N1 BY M2 OUTPUT MATRIX 05810C 05820C 05840C 05850C IF IFLAG = 1 THEN $C = A \neq B$ = 2 = 3 = TRANSPOSE (A) & B = A & TRANSPOSE (B) = TRANSPOSE (A) & TRANSPOSE (B) 05870C = 4 _GOTO (100, 110, 120, 130), IFLAG 100 I1 = L1 I2 = M2 13 = 12GOTO 148 110 II = L2 I2 = M2 05940 13 = Li $\begin{array}{c} 13 - L1 \\ \text{GOTO} \quad 140 \\ 120 \quad 11 = L1 \\ 12 = M1 \\ 13 = L2 \\ \text{GOTO} \quad 140 \\ 130 \quad 11 = L2 \\ 2 = M1 \\ 130 \quad 2 = M1 \\ 1$ 05980 06020 I2 = hi13 = L1 D0 200 1 = 1, I1 D0 200 II = 1, I2 TEMP = 0.D0 DO 190 III = 1, I3 GOTO (150, 160, 170, 180), IFLAG. 150 IEMP = IEMP + A (I + L1 * (III - 1)) * B (M1 * (II - 1) + III) **GOTO 190** 160 TEMP = TEMP + A (L1 # (I - 1) + III) # B (M1 # (II - 1) + III)

86130 GOTO 199 TEHP = TEHP + A (I + Li * (III - i)) * B (II + Hi * (III - i))06140 170 GOTO 190 06150 180 TEHP = TEHP + A (Li * (I - 1) + III) * B (II + Hi * (III - i)) 06160 06170 190 CONTINUE \overline{C} (I + Ni t (II - i)) = TEMP 96180 200 CONTINUE 06190 06200 RETURN 06210 END 06220C CONTROL SIGNAL GENERATION SUBROUTINE 062300 06240C 06250 SUBROUTINE CONTRL (X, U, W, OLDRATE, DELTAT) IMPLICIT DOUBLE PRECISION(A-H,J-Z) 06260 06270 TAPLICIT DOUBLE PRECISION(A-H,J-Z) PARAMETER IXX=11,IDR=7 DIMENSION AAA(IXX,IXX),AAAA(IXX,IXX) DIMENSION AAA(IXX,IXX), BSTAR (IXX,4), CSTAR (IXX) DIMENSION ASTAR (IXX,IXX), BSTAR (IXX,4), CSTAR (IXX) DIMENSION ASTAR (IXX,IXX), BSTAR (IXX,4), CSTAR (IXX) DIMENSION A (IXX,IXX), B (IXX,4), XDOTO (IXX) DIMENSION AA (IXX,IXX), DIA (IXX,IXX), XDOTAX (IXX) DIMENSION ZERO (4), IWDRK1 (4), IWDRK2 (4) DIMENSION DE (IDR,IXX), E (IDR), W (IDR,IDR), R (4,4) DIMENSION DB (IDR,4), EDQ (4,IDR), BDGDB (4,4), REDGDB (4,4) DIMENSION AX (IXX), AXC (IXX), DAXC (IDR), DAXCE (IDR), BDQDAXCE (4) DIMENSION KTR (4), KU (4), FCOULM (4) DIMENSION ASTRX(IXX),XNEWIIXX) DIMENSION XTEMP(IXX) 06280 06290 06300 06310 06320 06330 06349 06350 06360 06370 06380 06390 DIMENSION XTEMP(IXX) DIALASION ATERCIAA DATA JE, JIX, JIYZ, JMXZ, JMY, JOXY, JOZ 4 /1.2D-2, 1.7D-2, 1.3D-2, 2.25D-2, 3.0D-2, 3.0D-2, 3.9D-2/ DATA FCOULM /.09D0,.1D0,.11D0,.165D0 / DATA KTR /1.9D-2,1.9D-2,4.1D-2,5.13D-2/ DATA KTR /1.9D-1,5.1D-1,7.9D-1,1.25D0/ DATA ZER0 /4#0.D0/ DATA ZER0 /4#0.D0/ 06400 06410 06420 06430 06440 06450 DATA TOROWT, LOCKWI, TILIWI, RATEWI /1.D-13, 1.D-13, 1.D8,1.D-13/ 06460 064700 ZERO = ZERO VECTOR USED TO COMPUTE DERIVATIVE WITH U = 0 TORGWT = WEIGHT ASSIGNED TO TOPSUE REQUIREMENT IN COST FUNCTION LOCKWT = WEIGHT ASSIGNED TO GIMBAL LOCK PROXIMITY IN COST FUNCTION TILTWT = WEIGHT ASSIGNED TO ELEVATION GIMBAL TILT IN COST FUNCTION RATEWT = WEIGHT ASSIGNED TO ELEVATION GIMBAL RATE IN COST FUNCTION FCOULM = COULOMB FRICTION TORGUE LIMIT. SPECIFIES THE FRICTION MAGNITUDE WHEN THE GIMBALS ARE UNSTUCK. KTR = CONVERSION CONSTANTS FROM TORGUE KOTOR VOLTAGES TO TORGUES KU = PROPORTIONALITY CONSTANTS FROM ANCIEDODS TO BACK FMES 064800 064900 06500C 065100 065200 06538C 06540C 065500 = CONVERSION CONSTANTS FROM TORSUE HOTOR VOLTAGES TO TOU = PROPORTIONALITY CONSTANTS FROM ANGLEDOTS TO BACK EMFS = INERTIA ABOUT ANY AXIS OF THE LEVATION GIMBAL = INERTIA ABOUT THE X AXIS OF THE INNER GIMBAL = INERTIA ABOUT THE X AND Z AXES OF THE INDER GIMBAL = INERTIA ABOUT THE X AND Z AXES OF THE MIDDLE GIMBAL = INERTIA ABOUT THE X AND Y AXES OF THE MIDDLE GIMBAL = INERTIA ABOUT THE X AND Y AXES OF THE OUTER GIMBAL = INERTIA ABOUT THE X AND Y AXES OF THE OUTER GIMBAL = INERTIA ABOUT THE Z AXIS OF THE OUTER GIMBAL 065600 χŲ ĴĖ JIX 065700 065900 JIYZ 06600C JHXZ 06610C 06620C JMY JÖXY 06630C JOZ 06640C THE LINEARIZED CONTINUOUS TIME SYSTEM EQUATIONS ARE OF THE FORM 06650C 06660C X(t) = A(t0) (X(t) - X(t0)) + B(t0) U(t) + X(t0) 066700 066800 THE LINEARIZED DISCRETE TIME SYSTEM EQUATIONS ARE OF THE FORM 066900 06798C X[n+i] = A[n] X[n] + B[n] U[n] + C[n] 067100 06720C 06730C 06740C THE COST FUNCTION TAKES THE FORM 06750C J = (D[n] X[n+1] + E[n]) Q (D[n] X[n+1] + E[n]) + U[n] R U[n]86760C 05770C 06780C SOLUTION IS #T T ¥-117 T 1 U[n] = -(R + B D Q D B) B D Q (D (A X + C) + E)067900 **66800C**

06810C 06820C	ASSOCIATE VARIABLES WITH ARRAY ELEMENTS COMPUTE REQUIRED TRIGONOMETRIC FUNCTIONS
06830C 86840	DO 98 I=1.1XX
06850	XTEMP(I)=X(I) 8 CONTINIE
06870	PSI = X (1)
06890	$\frac{dETH}{THETA} = X (3)$
06900 06910	PHI = X (4) SP = DSIN (PSI)
86920 86920	CP = DCOS (PSI) SR = DSIN (RETA)
06940	CB = DCOS (BETA)
06960	CT = DCOS (THETA)
06970 06980	SF = DSIN (PHI) CF = DCOS (PHI)
06990 07890	$SB2 = SB \pm SB$ $CB2 = CB \pm CB$
	ST2 = ST ¥ ST T2 = T ¥ ST
87830	PSIDOT = X (5)
07040 07050	$\frac{BETADOT}{THETADOT} = X (7)$
07060 07070	PHIDOT = X (8) NCX = W (1)
87090 N7890	HCY = H (2) HC7 = H (3)
07100	WOX = CF * WCX - SF * WCZ
07120	$HOZ = SF \neq WCX + CF \neq WCZ$
87148	WHX = WUX + THETADDF WHY = CT * WOY + ST * WOZ
07150 07160	WAZ =-SI X WUY + CI X WUZ WIX = CB X WAX + SB X WAY
07170 87180	WIY =-SB # WMX + CB # WMY WT7 = WM7 + BETAMAT
07198	WEX = CP * WIX - SP * WIZ
07210	WEZ = SP # WIX + CP # WIZ
07220L 07230C	i psidot i i mii miz miz mia i i yi i
07240C 07250C	1 BETADOT 1 1 M21 M22 M23 M24 1 1 Y2 1 -1 THETADOT 1 = 1 M31 M32 M33 M34 1 1 Y3 1
07260C 07270C	1 PSIDOT 1 1 M41 H42 H43 H44 1 1 Y4 1
07280C	TORQUE EQUATIONS FOR THE FOUR GINBALS EXCLUDING CONTROL SIGNALS
07310	Yi = $-KTR$ (i) * KV (i) * PSIDET + TF + (-CD + PUIDET + UCZ - CD + CT + PUIDET + UCZ
07320	4 - CB # THETADOT # WHZ + BETADOT # WIX)
07340	
07350 07360	5 + WIX # WIY # (JIX-JIYZ) ~ PSIDOT # JE # WIX Y3 = -KTR (3) # KV (3) # THETADOT
07370	$\begin{array}{c} \bullet \\ \bullet $
07390	& # WHZ - RETADOT # WHX) + WHY # WHZ # (JHY - JHXZ) - (CB2 # TTY + CD2 # TTY7 + TE) # PETADOT # UHY + CD # TE # PETDOT
07410	$\frac{1}{2} = \frac{1}{2} $
67430	$\delta = -\frac{1}{100} \times \frac{1}{100} \times \frac{1}{100} \times \frac{1}{100} - \frac{1}{100} \times \frac{1}{100} - \frac{1}{100} \times \frac{1}{100} \times$
#7440 07450	$ \begin{array}{c} \bullet \\ \bullet $
07468 07470	a - PHIDOT \$ NGX) + SB \$ CB \$ CT \$ (JIX - JIYZ) \$ (PHIDOT a NOZ - BETADOT \$ NNY) + (CT \$ (SB2 \$ JIX + CB2 \$ JIYZ + JE)
07480	& # BETADOT + ST # WHY # (JNY- JNXZ)) # WHX + SB # CT # JE

* PSIDOT * WIZ + ST * (JE * PSIDOT + WIY * (JIYZ - JIX)) * WIX 07490 07500 ł £ 075100 M MATRIX -- THIS IS THE INVERSE OF THE L MATRIX ABOVE BECAUSE M IS SYMMETRIC ONLY THE UPPER HALF NEED BE COMPUTED 175200 07530C 07540C 07550 DENOM = (JOXY+ST2*JHXZ+CT2*(JHY+CR2*JIYZ+SB2*(JIX+JE)))*(JHXZ +CB2*(JIX+JE)+SB2*JIYZ)-CB2*SB2*CT2*(JIYZ-JIX-JE)**2 H11 = ((JOXY+ST2*JHXZ+CT2*(JHY+JIX+JE))*(JE+SB2*JIYZ 07560 Ł 07570 +CB2#JIX+JHXZ)-CB2#CT2#(JIX-JIYZ)#(JHXZ+JIX+JE))/JE/DENOH H12 = -CB#ST#CT#(JE+JIX+IHXZ)/DENOH 07580 å 07590 M13 = SB*(JOXY+ST2*JHXZ+CT2*(JHY+JIX+JE))/DENOH 07600 07610 M14 = -CB±CT±(JHXZ+JIX+JE)/DENCH M22 = (CB2*CT2*(JIYZ-JIX-JE)*(JE+JIX+JMXZ)+(JHXZ+CB2*(JIX+JE)+SB2* JIYZ)*(J0XY+CT2*(JMY+JIX)+ST2*(JHXZ+JIYZ)+JE))/(JIYZ+JE)/DENOM 07620 07630 Ł 07640 H23 = SB*CB*ST*CT*(JIYZ-JIX-JE)/DENOM M24 = ST*(JMXZ+CB2*(JIX+JE)+SB2*JIYZ)/DENOM 07650 M33 = (JOXY+ST2*JHXZ+CT2*(JHY+CB2*JIYZ+SB2*(JIX+JE)))/DENOM M34 = CB*SB*CT*(JIYZ-JIX-JE)/DENOM 07669 07670 H44 = (JHXZ+CB2#(JIX+JE)+SB2#JIYZ)/DENOH 07680 076900 PARTIAL DERIVATIVES OF Y AND H WITH RESPECT TO X FOLLOW 877000 DY1DB = dY1/dBETA, DY1DFD = dY1/dPHIDOT ETC. 077100 077200 DY10B = JE\$(-CB*PHIDOT*WOZ+SB\$ST\$PHIDOT*WOX+SB\$THETADOT*WHZ+ 07730 07749 BETADOT #WIY) å = JE#CB*(THETADOT#WHY-CT#PHIDOT#WOX)+SB#PETADOT#WHZ#JE 07750 DYIDT DY1DF = JE*(-SB*PHIDOT*WOX+CB*ST*PHIDOT*WOZ-CB*CT*THETADUT*WOX 07760 07770 +PETALOT*(SB*ST*WDX-CB*WOZ)) $\delta_{DYiDPD} = -KTR(i) \pm KV(i)$ 07790 DYIDED = JETWIX DYIDTD = -JEXCBXWHZ+JEACBABETADOT 07800 DY1DFD = JE\$(-SB*W0Z-CB*ST*W0X+CB*ST*THETADOT+SB*CT*BETADOT) DY2DB = (WIY**2-WIX**2)*(JIX-JIYZ)-PSIDOT*JE*WIY 07810 07820 07830 DY2DI=WHZ#((JIYZ+JE)*THETADOT+(SB*WIY+CB*WIX)#(JIX-JIYZ)-PSIDOT*SB*JE)+(JIYZ+JE)*ST*PHIDOT*WOX 07840 Ł 07850 07860 DY2DF=(JIYZ+JE)*(ST*THETADOT*HOX+CT*PHIDOT*HOZ)+(-CB*HOZ+SB*ST *WOX)*(WIY*(JIX-JIYZ)-PSIDOT*JE)+WIX*(SB*WOZ+CB*CT#WOX) ŝ 07878 å **X**(JIX-JIYZ) DY2DPD = -JEXWIX 07880 DY2DBD = -KTR (2)*KV (2) DY2DTD = (JIYZ+JE)*HMY+(JIX-JIYZ)*(CB*WIY-SB*WIX)-CB*JE*PSID0T CB*JE*PSID0T 07890 07900 DY2DFD=CT#(JIYZ+JE)#(THETADOT-NOX)-CT#(JIX-JIYZ)#WMX-SB#CT#JE 07910 07920 #PSIDDT Ł DY3DB = 2*SB*CB*(JIYZ-JIX)*(PHIDOT*WOZ-BETADOT*WHY)+(SB2-CB2) *(JIX-JIYZ)*(ST \$7930 \$7940 *PHIDOT*WOX+THETADOT*WHZ-BETADOT*WHX)-SR*JE*PSIDOT*WIZ 07950 07960 DY3DT = SB*CB*(JIX-JIYZ)*(-CT*PHIDOT*WOX+THETADOT*NHY)+(WHZ**2 07570 -WHY\$\$2)\$(JHY-JHXZ)-(CB2\$JIX+SB2\$JIYZ+JE)*BETADOT\$WHZ-CB **\$JE**\$PSIDDT\$WHY £. DY3DF = (JHXZ+CB2#JIX+SB2#JIYZ+JE)#PHIDOT#HOX-SB#CB#(JIX-JIYZ) 07990 #(-ST#PHIDOT#WOZ+THETADOT#CT#WOX+BETADOT#WOZ)+(ST#WHZ 08000 ¥. +CT#WHY)#WOX#(JHY-JHXZ)-(CB2#JIX+SB2#JIYZ+JE)#BETADOT 08010 *ST*WDX+CB*JE*PSIDOT*CT*WDX 18920 88830 DY3DPD = CB#JE#WIZ DY3DBD = SB*CB*(JIX-JIYZ)*UHX-(CB2*JIX+SP2*JIYZ+JE)*WHY+CB*JE*PSIDOT 08040 DY3DTD = -KTR (3)*KV (3)+SB*CB*(JIX-JIYZ)*(SETADOT-WHZ) DY3DFD = (JHXZ+CB2*JIX+SB2*JIYZ+JE)*WOZ+SB*CB*SI*(JIX-JIYZ)*(THETADOT 08050 08060 -HOX)+(CT#HHZ-ST#HHY)#(JHY-JHXZ)-CT#(CB2#JIX+SB2#JIYZ+JE) 88970 L #BETADOT-CB#ST#JE#PSIDOT 08030 Ł DY4DB = 2\$\$B\$CB\$(-CT2*JIX*THETADOT*W0Z+(JIX-JIYZ)*(THETADOT*W0Y \$\$T\$CT-ST*CT*PHIDOT*K0X+CT*BETADOT*WMX))+(CB2-SB2)*CT \$(JIX-JIYZ)*(PHIDOT*W0Z-BETADOT*WMY)+CB*CT*JE*PSIDOT*WIZ 88898 08100 å 88110 Ł -ST\$WIX\$\$2\$(JIYZ-JIX)+ST*(JE\$FSIDOT+WIY*(JIYZ-JIX))*WIY 18120 1 DY4DT = 2*CT*ST*(JHY+SB2*JIX-JHXZ-JIYZ)*THETADOT*H0Z+(CT2-ST2) 08130 *(JHY-JHXZ+SF2*(JIX-JI/Z))*(T+ETA:OT*WOY-PHIDOT#WOX)+SR *CH*(JIX-JIYZ)*(-ST*(PHIDOT*WOZ-PETA:OT*WHY)-CT*RETADOT 08140 å 08150 é #WHZ)-ST#(SB2#JIX+CB2#JIYZ+JE)#BETADOT#WHX+(CT#WHY+ST#WHZ) 08160 å

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08170 ł #(JMY-JMXZ)#UMX-SP#JE#PSIDOT#(ST#WIZ+CT#WMY)+ET#(JE#PSIDOT 08180 ā +WIY#(JIYZ-JIX))#WIX+CE#ST#WHZ#(JIYZ-JIX)#WIX+ST#(JE *PSIDOT+WIY*(JIYZ-JIX))*SB*WHZ 08190 å DY4DF = (W0Z\$\$2-W0X\$\$2)\$(J0XY-J0Z)-(CT2\$(JWY+SB2\$JIX)+ST2\$(JWXZ 08200 08219 08220 +JIYZ)+JE)*THETADOT#WOX+ST#CT#(JHY-JHXZ+SB2#(JIX-JIYZ)) Ł ŝ #PHIDOT#HUZ+SB*CB*CT*(JIX-JIYZ)*(PHIDOT*HUX-BETADOT*ST 18230 #WDX)+ST2#WDX#(JHY-JHXZ)#WHX-(CT#(SB2#JIX+CB2#JIYZ+JE) ł 08240 #BETADOT+ST#WHY#(JHY-JHXZ))#W0Z+SB#CT2#JE#PSIDOT#W0X+ST #WIX&(JIYZ-JIX)*(CF*ST*HOX+SB*HOZ)+ST*(JE*PSIDOT+WIY *(JIYZ-JIX))*(SB*ST*WOX-CB*WOZ) 08250 08260 ŝ. 08270 DY4DPD = SB\$CT*JE*WIZ+ST*JE*WIX 08280 DY4DBD = -SB*CB*CT*(JIX-JIYZ)*WHY+CT*(SB2*JIX+CB2*JIYZ+JE)*WHX 08290 08300 +SBICT#JEIPSIDOT Ł DY4DTD = -(CT2*(JHY+SB2*JIX)+ST2*(JHXZ+JIYZ)+JE)*W0Z+ST*CT*(JHY 08310 Ł -JHXZ+SE2*(JIX-JIYZ))#U0Y+ST#UHY#(JHY-JHXZ)+ST#(JIYZ-JIX) 08320 #(CB#WIY-SB#WIX)+CT#(SB2#JIX+CB2#JIYZ+JE)#BETADOT+CB#ST å 08330 #JE#PSIDDT å 08340 DY4DFD = -KTR (4)*KU (4)+ST*CT*(JHY-JHXZ+SB2*(JIX-JIYZ))*(THETADOT -WOX)+SB*CB*CT*(JIX-JIYZ)*(WOZ-CT*BETADOT)+ST*CT*(JHY-JHXZ) 08350 08360 08370C 08380 £ ***WHX+ST*CT*(JIYZ-JIX)*(CB*WIX+SB*WIY)** DDENDB = 2xSBxCBx(JIYZ-JIX-JE)x(JOXY+ST2xJHXZ+CT2x(JHY-JHXZ)) DDENDT = 2*ST*CT*((JHXZ-JHY-CP2*JIYZ-SB2*(JIX+JE))*(JHXZ+CB2*(JIX+JE)) +SB2*JIYZ)+SB2*CB2*(JIYZ-JIX-JE)**2) DHi1DB = 2*SB*CB*(JOXY+ST2*JHXZ+CT2*(JHY-JHXZ))*(JIYZ-JIX)/JE/DENOM 08390 08400 Ł 08410 08420 Ł -DDENDE#H11/DENOH \$8430 DHiiDI = 2*ST*CT*((JHXZ-JHY-JIX-JE)*(JE+SB2*JIYZ+CB2*JIX+JHXZ) 08440 08450 +CB2#(JIX-JIYZ)#(JHXZ+JIX+JE))/JE/DENOH-DDENDT#H11/DENOH Ł DM12DB = SB*ST*CT*(JE+JIX+JHXZ)/DENOH-DDENDB*H12/DENOH DM12DT = CB*(ST2-CT2)*(JE+JIX+JHXZ)/DENOH-DDENDT*H12/DENOM 08460 DHI2DI = CB*(3)2-C12)*(3+21)*(3)2/2)/DEHCM-DDENDIAH12/DEHOM DHI3DB = CB*(J0XY+ST2*JHXZ+CT2*(JHY+JIX+JE))/DEHOM-DDENDBRH13/DEHOM DHI3DT = 2*SB*ST*CT*(JHXZ-JHY-JIX-JE)/DEHOM-DDENDT*H13/DEHOM DHI4DB = SB*CT*(JHXZ+JIX+JE)/DEHOM-DDENDT*H14/DEHOM DH14DT = CB*ST*(JHXZ+JIX+JE)/DEHOM-DDEHDT*H14/DEHOM DH22DB = 2*SB*CB*(JIYZ-JIX-JE)*(J0XY+CT2*(JHY-JHXZ)+ST2*(JIYZ 08470 18480 08490 08510 +JHXZ+JE))/(JIYZ+JE)/DENCH-DDENDR*H22/DENCH DH22DT = 2*CT*ST*(CB2*(JIX-JIYZ+JE)*(JE+JIX+JHXZ)+(JHXZ+CB2*(JIX +JE)+SB2*JIYZ)*(JHXZ+JIYZ-JHY-JIX))/(JIYZ+JE)/DENCH 08520 Ł 08530 08540 Ł 08550 -DDENDT*H22/DENON Ł DM23DB = (CP2-SB2)*ST*CT*(JIYZ-JIX-JE)/DENOM-DDENDB*M23/DENOM DM23DJ = SB*CF*(CT2-ST2)*(JIYZ-JIX-JE)/DENOM-DDENDB*M23/DENOM DM24DB = 2*SB*CB*ST*(JIYZ-JIX-JE)/DENOM-DDENDB*M24/DENOM DM24DT = CT*(JMXZ+CB2*(JIX+JE)+SB2*JIYZ)/DENOM-DDENDT*h24/DENOM DM33DB = 2*SB*CB*CT2*(JIX+JE-JIYZ)/DENOM-DDENDB*M33/DENOM -DM33DT = 2*ST*CT*(JMXZ-JMY-CB2*JIYZ-SB2*(JIX+JE))/DENOM 08560 08570 08580 08590 88666 08610 08620 -DDENDT#H33/DENON Ł DH34DB = (CB2-SB2)*CT*(JIYZ-JIX-JE)/DENOM-DDENDB*H34/DENOM DH34DT = SB*CB*ST*(JIX+JE-JIYZ)/DENOM-DDENDB*B380[#ENOM DH44DB = 2*SB*CB*(JIYZ-JIX-JE)/DENOM-DDENDB*H44/DENOM 08630 08640 08650 08660 DM44DT = -DDENDT*H44/DENOM08670C **#8680C** 08690C 08700C DERIVATIVES OF WE WITH RESPECT TO X FOLLOW DWXDFD = dWEX/dPHIDOT ETC. 98710C 08720 DWXD? = -WEZ DWXDB = CP * WIY DWXDT = CP * SB * WHZ + SP * WHY 08738 08749 08750 DWXDF = CP # SB * ST # W0X - CP # CB # W0Z - SP # CT # W0X 08760 08771 08780 DWYDB = -WIX 8779 DWYDT = CB * WMZ DWYDF = SB * WDZ + CB * ST * WDX 08800 08810 DWYDPD=1.DO 08820 18830 DWYDTD = -SB DWYDFD = CB # CT 88840

08850 88860 88870 08850 08890 08910 08910 08920C 08930C 08930C 08930C 08940C 08950C	DWZDP = WEX DWZDB = SP * WIX DWZDT = SP * SB * WHZ - CP * WMY DWZDF = SP * SB * ST * WOX - SP * CB * WOZ + CP * CT * WOX DWZDBD = CP DWZDTD = SP * CB DWZDFD = SP * SB * CT - CP * ST SET UP MATRIX A CLEAR OUT A
08960C 08970 08980 08990 09000 09000 09000 09010C 09020 09020	DO 100 I = 1, IXX DO 100 II = 1, IXX A (I,II) = 0,D0 100 CONTINUE DO 110 I = 1, 4 A (I, IAA)=5 D0
070340 07050 07050 07060 05050 05050 05100 05110 05120 05120	<pre>ii0 CONTINUE A (5,2) = DMiiDB * Yi + Kii * DYiDB + DMi2DB * Y2 + Mi2 * DY2DB b + DMi3DB * Y3 + Mi3 * DY3DB + DMi4DB * Y4 + Mi4 * DY4DB A (5,3) = DMiiDT * Yi + Mi1 * DYIDT + DMi2DT * Y2 + Mi2 * DY2DT b + DMi3DT * Y3 + Mi3 * DY3DT + DMi4DT * Y4 + Mi4 * DY4DT A (5,4) = Mii * DYIDF + Mi2 * DY2DF + Mi3 * DY3DF + Mi4 * DY4DF A (5,5) = Mii * DYIDFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,6) = Mii * DYIDTD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,7) = Hii * DYIDTD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DYIDTD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DYIDTD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD A (5,8) = Mii * DY1DFD + Mi2 * DY2DFD + Mi3 * DY3DFD + Mi4 * DY4DFD </pre>
09140 09150 09160 09170 09180 09190 09200 09210 09220 09220	A (6,2) = DM12DB x Y1 + H12 x DY1DB + DH22DB x Y2 + H22 x DY2DB b + DH23DB x Y3 + H23 x DY3DB + DH24DB x Y2 + H22 x DY2DB A (6,3) = DM12DT x Y1 + H12 x DY1DT + DH22DT x Y2 + H22 x DY2DT b + DH23DT x Y3 + H23 x DY3DT + DH24DT x Y4 + H24 x DY4DT A (6,4) = H12 x DY1DF + H22 x DY2DF + H23 x DY3DFD + H24 x DY4DF A (6,5) = H12 x DY1DF + H22 x DY2DFD + H23 x DY3DFD + H24 x DY4DF A (6,6) = H12 x DY1DFD + H22 x DY2DFD + H23 x DY3DFD + H24 x DY4DFD A (6,6) = H12 x DY1DFD + H22 x DY2DFD + H23 x DY3DFD + H24 x DY4DFD A (6,6) = H12 x DY1DFD + H22 x DY2DFD + H23 x DY3DFD + H24 x DY4DFD A (6,7) = H12 x DY1DFD + H22 x DY2DFD + H23 x DY3DFD + H24 x DY4DFD A (6,8) = H12 x DY1DFD + H22 x DY2DFD + H23 x DY3DFD + H24 x DY4DFD A (6,8) = H12 x DY1DFD + H22 x DY2DFD + H23 x DY3DFD + H24 x DY4DFD A (6,8) = H12 x DY1DFD + H22 x DY2DFD + H23 x DY3DFD + H24 x DY4DFD A (2,2) = DM13DFD x Y1 + H13 x DY1DR + DH23DFD x Y2 + H23 x DY2DR
09240 09250 09260 09270 09280 09280 09280 09280 09300 09310 09320	 t DH33DB \$ Y3 + M33 \$ DY3DB + DH34DB \$ Y4 + M34 \$ DY4DB A (7,3) = DH13DT \$ Y1 + M13 \$ DY1DT + DH23DT \$ Y2 + M23 \$ DY2DT b DH33DT \$ Y1 + M13 \$ DY1DT + DH23DT \$ Y2 + M23 \$ DY2DT c DH33DT \$ Y3 + M33 \$ DY3DT + DH34DT \$ Y4 + M34 \$ DY4DT A (7,4) = M13 \$ DY1DF + M23 \$ DY2DF + M33 \$ DY3DF + M34 \$ DY4DF A (7,5) = M13 \$ DY1DF + M23 \$ DY2DF + M33 \$ DY3DF + M34 \$ DY4DF A (7,5) = M13 \$ DY1DFD + M23 \$ DY2DPD + M33 \$ DY3DFD + M34 \$ DY4DFD A (7,6) = M13 \$ DY1DFD + M23 \$ DY2DFD + M33 \$ DY3DFD + M34 \$ DY4DFD A (7,7) = M13 \$ DY1DFD + M23 \$ DY2DFD + M33 \$ DY3DFD + M34 \$ DY4DFD A (7,8) = M13 \$ DY1DFD + M23 \$ DY2DFD + M33 \$ DY3DFD + M34 \$ DY4DFD A (7,8) = M13 \$ DY1DFD + M23 \$ DY2DFD + M33 \$ DY3DFD + M34 \$ DY4DFD A (8,2) = DM14DB \$ Y1 + M14 \$ DY1DB + DH24DB \$ Y2 + M24 \$ DY2DB
07330 07340 07350 07360 87370 07380 87370 07380 87370 07400 07410 89410	A (8,3) = DH34DB x 13 + H34 x DY3DB + DH44DB x 14 + H44 x DY4DB A (8,3) = DH24DT x Y1 + H14 x DY1DT + DH24DT x Y2 + H24 x DY2DT b + DH34DT x Y3 + H34 x DY3DT + DH44DT x Y4 + H44 x DY4DT A (8,4) = H14 x DY1DF + H24 x DY2DF + H34 x DY3DF + H44 x DY4DF A (8,5) = H14 x DY1DFD + H24 x DY2DFD + H34 x DY3DFD + H44 x DY4DFD A (8,6) = H14 x DY1DFD + H24 x DY2DFD + H34 x DY3DFD + H44 x DY4DFD A (8,7) = H14 x DY1DFD + H24 x DY2DFD + H34 x DY3DFD + H44 x DY4DFD A (8,8) = H14 x DY1DFD + H24 x DY2DFD + H34 x DY3DFD + H44 x DY4DFD A (8,8) = H14 x DY1DFD + H24 x DY2DFD + H34 x DY3DFD + H44 x DY4DFD A (8,8) = H14 x DY1DFD + H24 x DY2DFD + H34 x DY3DFD + H44 x DY4DFD A (8,9) = H14 x DY1DFD + H24 x DY2DFD + H34 x DY3DFD + H44 x DY4DFD A (9, 1) = DW2DP A (9, 2) = DW2DP
09430 09440 09450 09450 09460 89478 89480 05450 89580 89510 89528	A $(9, 3) = DWXDT$ A $(9, 4) = DWXDF$ A $(9, 4) = DWXDFD$ A $(9, 7) = DWXDFD$ A $(9, 8) = DWXDFD$ A $(10, 2) = DWYDFD$ A $(10, 3) = DWYDF$ A $(10, 4) = DWYDF$ A $(10, 5) = DWYDFD$ A $(10, 7) = DWYDFD$

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19530 19540 19550 19568 19580 19580 19580 19580 19600 19610 19610		A (10, 8) = DWYDFD A (11, 1) = DWZDP A (11, 2) = DWZDB A (11, 3) = DWZDT A (11, 4) = DWZDF A (11, 6) = DWZDFD A (11, 7) = DWZDFD A (11, 8) = DWZDFD
9620C		SET UP ARRAY B
09640C		CLEAR OUT B
096500 09660 09670 09680 09690	120	DO 120 I = 1, IXX DO 120 II = 1, 4 B (I, II) = 0.D0 CONTINUE
97000 9710 9720 9730 9740 9750 9750		B (5, 1) = KTR (1) * H11 B (5, 2) = KTR (2) * H12 B (5, 3) = KTR (2) * H13 B (5, 4) = KTR (4) * H14 B (6, 1) = KTR (1) * H12 B (6, 2) = KTR (2) * H22 B (6, 3) = KTR (3) * H23
05760 05750 05810 07810 07810 07810 07820 07830 07850 07850 07850		B (6, 4) = KTR (4) \pm M24 B (7, 1) = KTR (1) \pm M13 B (7, 2) = KTP (2) \pm M23 B (7, 3) = KTR (3) \pm M33 B (7, 4) = KTR (4) \pm M34 B (8, 1) = KTR (4) \pm M34 B (8, 3) = KTR (2) \pm M24 B (8, 3) = KTR (3) \pm M34 B (8, 4) = KTR (4) \pm M34
09870C 09880C		SET UP VECTOR XD070
09890C 09900		CALL DERIVE (XTEMP, XDOTO, ZERO, W, OLDRATE)
09910C 09920C		4 GINBAL SYSTEH DYNAMIC EQUATIONS ARE NOW COMPLETELY LINEARIZED
19930C 19940C		THE DISCRETE TIME APPROXIMATIONS FOLLOW
09950C 09960C		$\begin{array}{c} 2 \\ A = I + DELTAT * A + DELTAT * A / 2! \end{array}$
09970C 09980 09990 10000 10010		DELTA2 = .5D0 * DELTAT ** 2 DELTA3=DELTAT*DELTA2/3.D0 DELTA4=DELTAT*DELTA3/4.D0 DELTAS=DELTAT*DELTA4/5.D0
10030 10040 10050 10060 10060		CALL MATMPY (A, A, AA, IXX, IXX, IXX, IXX, IXX, IXX,
10080 10090 10100 10110 10110	i30	ASTAR (I, II) = DELTAT * A (I, II) + DELTA2 * AA (I, II) +DELTAJ#AAA(I,II)+DELTA4#AAAA(I,II) IF (I .EQ .II) ASTAR (I, II) = ASTAR (I, II) + 1.DO CONTINUE
		B = (DELTAT # I + DELTAT # A / 2!) B
10150 10160 10170 10180 10180	ł	DO 149 I = 1, IXX DO 149 II = 1, IXX DIA (I, II) = DELTA2 # A (I, II) S+DELTA3#AA(I,II)+DELTA4#AAA(I,II)+DELTA5#AAAA(I,II) IF (I .EQ. II) DIA (I, II) = DIA (I, II) + DELTAT
10200	140	CONTINUE

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10210
10220C
                   CALL MATHPY (DIA, B, BSTAR, ii, ii, ii, 4, ii, 4, i)
                     1
10230C
10240C
                    C = (DELTAT \ddagger I + DELTAT \ddagger A / 2!) (X (t0) - A X (t0))
                   D0 160 I = 1, IXX
TEMP = XD0T0 (I)
D0 150 II = 1, IXX
TEMP = TEMP - A (I, II) $ XTEMP (II)
10250
10260
19270
10280
            150 CONTINUE
10290
10300
                    XDOTAX (I)=TEMP
10310
           160 CONTINUE
                    CALL HATMPY (DIA, XDOTAX, ESTAR, IXX, IXX, IXX, 1, IXX, 1, 1)
10320
10330C
10340C
                    THE MATRIX R REFLECTS THE WEIGHT OF THE CONTROL SIGNALS IN THE COST FUNCTION
103500
10360C
10370
                   DO 170 I = 1, 4
DO 170 II = 1, 4
R (I, II) = 0.DO
IF (I .EQ. II) R (I, II) = TORQWT
10380
10390
10400 IF (1 .EC
10410 170 CONTINUE
10420C
                    THE MATRIX & REFLECTS THE WEIGHT OF THE STATE IN THE COST FUNCTION
10430C
10440C
10450
10469
                    DO 180 I = 1, IDR
DO 180 II = 1, IDR
            Q (I, II) = 0.D0
IF (I.EQ. II.AND. I.GE. 2) Q (I, II) = TILTWT
IF(I.EQ.II.AND.I.GE.5)Q(I,II)=RATEWT
100 CONTINUE
10470
10480
10490
10500
10510
10520C
10530C
                    Q (1, 1) = LOCKWT
                    THE MATRIX D COMPRESSES THE STATE INFORMATION AND LINEARIZES THE GIMBAL LOCK COST
10540C
105500
            D6 190 I = 1, IDR
D0 190 II = 1, IXX
D (I, II) = 0.D0
190 CONTINUE
10560
10570
10580
           D (1, 11) = 0.D0

190 CONTINUE

DO 200 I = 1, 3

D (1 + 1, I + 8) = 1.D0

200 CONTINUE

D(1,2)=CB*ST

D(5,1)=DWXDP

-D(5,2)=DWXDF

D(5,2)=DWXDFD

D(5,3)=DWXDFD

D(5,4)=DWXDFD

D(5,6)=DWXDFD

D(5,6)=DWXDFD

D(5,2)=DWYDFD

D(5,2)=DWYDFD

D(5,2)=DWYDFD

D(5,2)=DWYDFD

D(6,2)=DWYDFD

D(6,3)=DWYDFD

D(6,3)=DWYDFD

D(6,3)=DWYDFD

D(6,3)=DWYDFD

D(6,3)=DWYDFD

D(6,3)=DWYDFD

D(6,3)=DWYDFD

D(6,3)=DWYDFD

D(7,1)=DWZDF

D(7,4)=DWZDFD

D(7,8)=DWZDFD
10590
10600
10610
10620
10630
10640
10650
10670
10680
18698
10700
10710
10720
10730
10740
10750
10760
10770
10790
10800
10810
10820
18830
10840
                    D(7,8)=DWZDFD
108500
                    THE MATRIX E EXPRESSES THE OPTIMAL LINEARIZED NEXT STATE
108500
108700
10880
                    E(1)=SB$S1-BETA$CB$ST-THETA$SB$CT
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-97-
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10870		DO 210 I = 2, 4
10900		$E(I) = 0.00^{\circ}$
10910	210	CONTINUE
10920		E(5)=WEX-PSI*DWYDP-BETA*DWXDB-THETA*DWXDT-PHI*DWXDF
10930		BETADOT#DW\DBD-THETADOT#DWXDTD-PHIDOT#DWXDFD
10940		E(6)=WEY-BETA*DWYDB-THETA*DWYDT-PHI*DWYDF-PSIDOT#
10950		& DWYDPD-THETADOT*DWYDTD-PHIEOT*DWYDFD
10960		E(7)=WEZ-PSI\$DUZDP-BETA\$DUZDB-THETA\$DUZDT-PHI\$DUZDF-
10970	4	BETADOT#DWZDED-THETADOT#DWZDTD-PHIDOT#DWZDFD
109800		
10990C		COMPUTE U
110000		
11010		CALL MATHPY (D, BSTAR, DB, IDR, IXX, IXX, 4, IDR, 4, 1)
11020		CALL MATMPY (DB, Q, BDQ, IDR, 4, IDR, IDR,4, IDR, 2)
11030		CALL MATMPY (BDQ, DB, BDQDB, 4, IDR, IDR, 4, 4, 4, 1)
11040		DO 220 I = 1, 4
11050		DO 220 II = 1, 4
11960		RBDQDB (I, II) = R (I, II) + BDQDB (I, II)
11070	220	CONTINUE
11080		CALL MINV (RBDQDB, 4, DETERN, INORK1, INORK2)
11090		CALL MATMPY (ASTAR, XTEMP, AX, IXX, IXX, IXX, 1, IXX, 1, 1)
11190		DO 230 I = 1, IXX
11110		AXC (I) = AX (I) + CSTAR (I)
11120	239	CONTINUE
11130		CALL MATMPY (D, AXC, DAXC, IDR, XXX, IXX, 1, IDR, 1, 1)
11140		DO 240 I = 1, IDR
11150		DAXCE (I) = DAXC (I) + E (I)
11160	249	CONTINUE
11170		CALL MATMPY (BDD, DAXCE, BDQDAXCE, 4, IDR, IDR, 1, 4, 1, 1)
11180		CALL MATMPY (RBDQDB, BDQDAXCE, U, 4, 4, 4, 1, 4, 1, 1)
11190		DO 250 I = 1, 4
11200		$U(\mathbf{I}) = -U(\mathbf{I})$
11210	250	CONTINUE
11220		KE IUKN
11250		ENU

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