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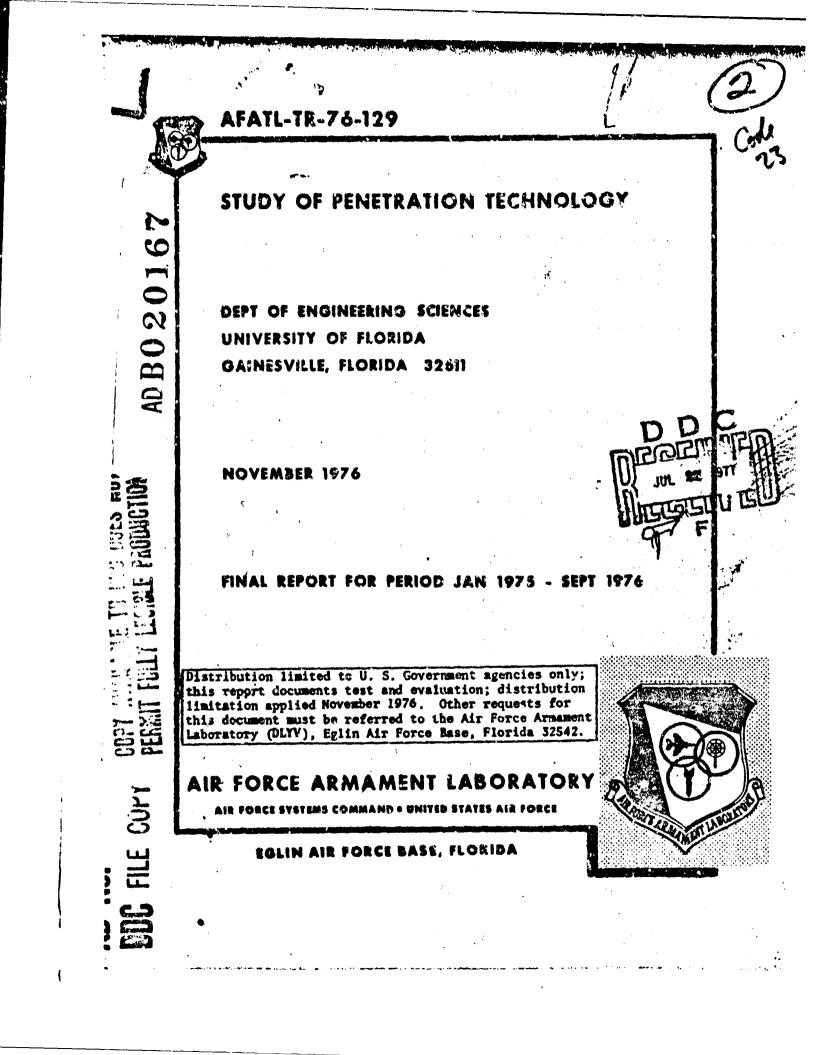
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## PREFACE

This report represents the results of a study on the soil penetration performance of terradynamic impactors. This study was conducted Jan 1, 1975 to Sept 29, 1976 by the Engineering Sciences Department, University of Floridi, Gainesville, Florida, 32611, under Contract No. F08635-75-C-0054 with the Air Force Armament Laboratory, Eglin Air Force Base, Florida. Mr. John Collins served as program manager for the Armament Laboratory.

This technical report has been reviewed and is approved for publication.

FOR THE COMMANDER

J. R. MURRAY Chief, Weapon Systems Analysis Division

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UNCLASSIFIED SECURITY CLASSIFICATION OF THIS PADER to Date Entered and other pressure transducers, strain gages on test chamber walls, and cont. breaking-wire sensors. Data analysis was performed by using classical Poncelet predictive techniques, empirical Sandia penetration equations, semi-analytical Cavity Expansion Theory, and a three dimensional code for trajectory analysis developed under this contract for use with an assumed three dimensional force law. To obtain essential input to the aforementioned models, a study of the acoustic wave velocity in a sand medium was made, and results are included in this report. . . **.** هده م الح الجان A. • • **..**. . UNCLASSIFIED SECURITY JLASS FIGATION OF THIS PASE The Date Balance أأسط والمردفية كاردا والأقعد وم الاحجا لمانحم صادرتهم المرمصان

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#### SECTION I

#### INTRODUCTION AND BACKGROUND

This report presents the results of a joint investigation conducted by the Vulnerability Assessment Branch (ELYV) of the Air Force Armament Laboratory and the University of Florida. The University gave analytical support to an experimental program conducted by DLYV/AFATL at Eglin Air Force Base.

The support included reviewing the proposed soil penetration experiments, recommending changes, participating in some of the experiments at Eglin, making independent laboratory invostigations at the University of several types of sensors and of ultrasonic wave speeds in sand, extensive data analysis of the Eglin Experiments, study of existing terradynamic penetration models, modification of the models and application of them to the interpretation of the Eglin experiments.

The study of the mechanics of high speed earth penetrators, including predictions of trajectory, depth of penetration, cavity formation, stability, and target interaction has in recent years been given the name of terradynamics. While this area of study has been investigated since the early 18th century, technological barriers have hindered experimental programs in assessing models advanced for characterizing penetrator performance. The principal difficulty encountered has been the unavailability of experimental tools for examining the sequential motion of a vehicle pausing through opaque loose and/or semicohesive media.

A recent review of the State of the Art of Earth Penetration Technology by Triandafilidis (Reference 1) has categorized predictive penetration techniques as semi-analytical, analytical, theoretical, and empirical models. The first technique, which includes the classical penetration models based upon Newtonian mechanics, such as Poncelet (Reference 2), requires experimental data for evaluation of the important penetration constants. So-called analytical techniques, which include the Cavity Expansion (References 3 through 7) and Differential Force Law Models (Reference 8), rely upon knowledge of constitutive target material properties. The theoretical models proposed (References 9 through 11) are based upon continuum mechanics formulations describing the penetrator and target, and rely upon finite difference and finite element computer codes as solution techniques. Finally, empirical techniques based upon extensive laboratory and field testing have been introduced with the most extensive work in this area developed at Sandia Laboratories (References 12.13). Additional background on the experimental program is presented in Section II.

The purpose of the experimental program at Eglin was to obtain more complete transient records of the penetration events than previous investigators had obtained in order to provide insight into the actual physical mechanisms involved, which could lead to better cerradynamic penetration models for predicting trajectories, penetration depths, and the three acting on the projectile. In the test program, five consecutively space X-ray units have been used to visually record the transient position of several penetrators. Nonspinning projectiles of stable configuration with various nose shapes have been tested in dry and saturated sand at three impact

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velocities with near zero impact obliquity. This is believed to be the most extensive use ever made of flash radiography in terradynamic research. In addition to the X-ray units, velocity coil sets is have been used as monitoring devices in conjunction with a magnetic tape recording system.

The experimental setup at Eglin and some of the experiments at the University on sensors are described in Section II after a short background account of previous experimental studies. Data from the Eglin tests are described in Section III, with details tabulated in Appendix A. Analysis of the data by classical semi-analytical penetration models and empirical methods is presented in Section IV. The analytical technique based on the spherical cavity expansion technique is discussed in Section V and applied to the Eglin experiments. In Section VI a three-dimensional terradynamic model is developed and applied. Sound speed measurements are reported in Section VII and a summary of the conclusions is given in Section VIII.

#### SECTION II

#### EXPERIMENTAL EQUIPMENT AND PROCEDURES

#### 2.1 INTRODUCTION

Penetration experiments were performed by firing projectiles horizontally into sand targets contained in specially designed test chambers. After some preliminary tests with 0.50 caliber and 20mm standard rounds, the major part of the investigation used modelled 20mm projectiles fabricated both at the AFATL and at the University of Florida. These projectiles were cylinders 0.02 meter in diameter by 0.22 to 0.24 meter in length. Three specific nose shapes were investigated: biconic, flat ended, and step-tier. Some of the biconic and step-tier projectiles had a hollow afterbody, but the majority of the results were obtained using solid projectiles.

Various sensing methods were investigated to determine as much as possible about the projectile's position and orientation, the shape of the cavity formed around the projectile, deformation patterns and force distributions in the sand, and shock waves ahead of the projectile. The most successful sensing method was flash radiography. In the simmary test program five X-ray heads were fired sequentially with delay times set to record projectile position as it moved through a 1.2-meterlong test chamber. The primary test program was planned to include firings of two projectile types (flat and step-tier projectiles) at three different velocities (approximately 210,320 and 400 m/sec) in dry sand and in saturated sand, with four replications of each type of shot, and five Xray pictures taken in each shot. This program was completed successfully. Results of these tests are presented in Section III, along with a few examples of other projectile types.

Besides giving a more certain indication of projectile trajectory and attitude than any other sensing method, the X-rays give a good indication of the position on the projectiles where the sand separates to form a cavity, and can show also the reattachment point on the afterbody as the projectile slows down. In the primary test program the X-rays showed that reattachment seldom occurred in the 1.2 meters of the trajectory observed.

The X-rays also revealed a detached bow wave in some cases (notably the higher-speed impacts in dry sand). The bow wave is a density discontinuity moving with the projectile, resembling the detached shock wave ahead of a supersonic aircraft. The X-ray method was emphasized because it was the only method known that could give transient information about separation and about the shock wave shape and density gradients. Other types of sensors envisioned for use in the test program were investigated to complement the X-ray technique or to be used in case the X-ray equipment was not available.

Some of the sensing methods investigated at the University were microwaves, breaking wires, and magnetic sensors. The microwave technique was considered as an alternative to the X-rays for continuous position monitoring, but it was not used in the experiments at Eglin, since the X-ray equipment was available. Various breaking-wire sensors and velocity screens were used at Eglin, and the magnetic sensing method was used extensively both in the primary lest program and in the preliminary testing before the X-ray system was fully developed. A pressure transducer in the floor of the test chamber and strain gages on the walls were also used in atrempting to build a complete data base.

The general set-up for the primary test program at Eglin and the flash X-ray method are described in paragraph 2.3 after a brief review in paragraph 2.2 of some previous terradynamic experiments. Other sensors used in or examined for the test program are discussed in paragraph 2.4.

#### 2.2 BACKGROUND

Until fairly recently the only experimental data available on ballistic penetration of soils consisted of tabulations of striking velocity  $V_0$  versus final penetration distance S. Comparisons of the plots of S versus  $V_0$  with integration of assumed force laws, e.g., of the form

 $-dV/dt = cV^2 + \beta V + \gamma$ 

could in principle determine the coefficients for such laws. The scatter in the data because of variations of <u>in situ</u> soil properties or because of tumbling or other unstable projectile behavior made conclusions from S versus  $V_0$  data difficult to draw.

In 1957 Allen, Mayfield, and Morrison (Reference 14) reported what were apparently the first laboratory investigations to record projectile transient motion. They used a photographic-electronic chronograph to record the successive breaking of copper grid wires located 0.1 meter apart along the trajectory and were able to obtain better determination of force law coefficients than could be obtained from final penetration depths alone.

This brief discussion will not attempt a complete historical account of penetration experiments, but will mention a few of the more recent investigations that have obtained transient data. Some additional historical information is given in References 1 and 14 through 16 and in a 1972 survey f the state of the art by McNeill (Reference 17), which also gives a bibliography. A more extensive bibliography has been prepared by Triandafilidis (Reference 1), and a 1974 annotated bibliography (Reference 18) lists Sandia Laboratories Publications related to Terradynamics.

According to McNeill, significant strides in penetrator system technology began at Sandia in 1961 with penetrators 2.4 to 3 meters in length and 0.23 to 0.46 meter in diameter, with masses on the order of 450 kilograms, delivered by ground-launched rockets or by airplanes. Some of these tests used on-board accelerometers and telemetered data. Since that time, the accelerometer-carrying air-dropped penetrometer has been developed into a practical tool for rapid survey of subsurface soil properties. Wood (Reference 19) has discussed instrumentation and telemetry. Trailing wires have also been used for air gun projectiles at speeds up to 120 m/sec (Reference 19). Murff and Coyle (Reference 20) have obtained deceleration-time records for impact at speeds up to 90 m/sec into three soil types (compacted kaolin clay, dense Ottawa sand and a mixture of kaolin clay and sand). Projectiles varied from 38 to 76 millimeters in diameter and had masses ranging from 1.4 to 52 kilograms.

A microwave monitoring system was developed at the University of New Mexico. Its use was reported in a Ph.D dissertation in 1965 by Hakala (Reference 15). The technique showed considerable promise, although questions of how long a path could be monitored and whether the technique could be used in moist soil were not addressed.

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Successful use of flash radiography in soil beginning in 1974 was reported by Culp et al (References 21, 22). Their later work in clay showed the existence of a detached shock wave. Color enhancement techniques of the X-rays revealed density variations. An automatic scanning and image storing and processing technique was used. One significant result of the scanning technique was the discovery that the soil cavity around the projectile seemed to be larger than it had appeared in visual inspection of the radiographs. Flash radiography in soil had been used earlier (Reference 23), but few details about it have been made public.

Although transient trajectory measurements were not made, the share of the trajectory was revealed by post-test excavation in a 1973 Master's Thesis by Biele (Reference 24), which investigated the stability of scaled model projectiles of various nose types. Initial angles of impact were revealed by yaw cards, and plots of lateral deflection versus penetration distance were made for various initial angles. Even with quite small initial angles (1-2 degrees) lateral deflections of as much as 0.15 meter were observed in a penetration distance of 1.06 meters.

#### 2.3 SETUP OF PENETRATION EXPERIMENTS AT EGLIN

The test setup used in collecting the data base for analysis was developed in an evolutionary manner. Several sensing devices, projectile shapes, and velocity regimes were studied before the basic elements of the primary test matrix were investigated. Details of these techniques are given elsewhere in this section and principal attention is focused upon the test assemblage as used in the March and April 1975 test program. The primary test matrix is shown below with the complete matrix described in paragraph 3.2.

Projectile	Target		Velocity		
Туре		210 m/sec	320 m/sec	400 m/sec	
Flat Nose	Dry Sand	15,16,17, 18,19	20,22,23,24	14,25,26,27,29	
Flat Nose	Wet Sand	70,71,72,73	36,37,38,74,81	76,82,83,84	
Step Tier	Dry Sand	52,53,54,55,57	56,58,59,61	62,63,64,65	
Step Tier	Wet Sand	42,43,44,45	39,40,41,49	50,51,68,69	

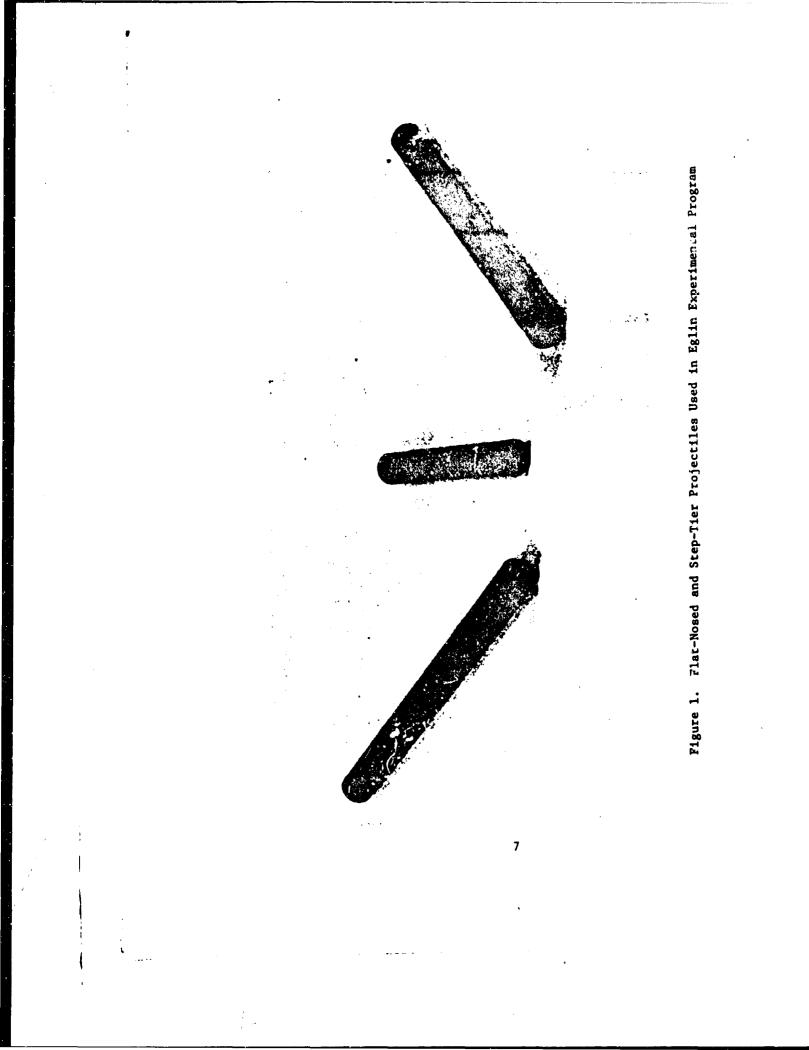
The flat-nose projectiles used in these experiments were solid cylinders 0.0198 meter in diameter by 0.225 meter long. For the stepcier projectiles the afterbody was a cylinder of 0.0198 meter diameter and 0.232 meter length, with a cylindrical nose 0.0095 meter in diameter and 0.0065 meter long. The material used for the projectiles was a high carbon content steel drill rod, supplied in rod form with nominal dimensions of 0.02 meter in diameter and 1 meter long. For specimens made at the University an AISI-WI water quenched bar stock was used, while for specimens fabricated at the AFATL, AISI-01 oil quenched bar stock was used. Three of the projectiles used in the Eglin penetration experiments are shown in Figure 1. In addition to the two projectiles described above, the photograph shows a shorter flat-nosed projectile, length 0.152 meter, used for some later tests.

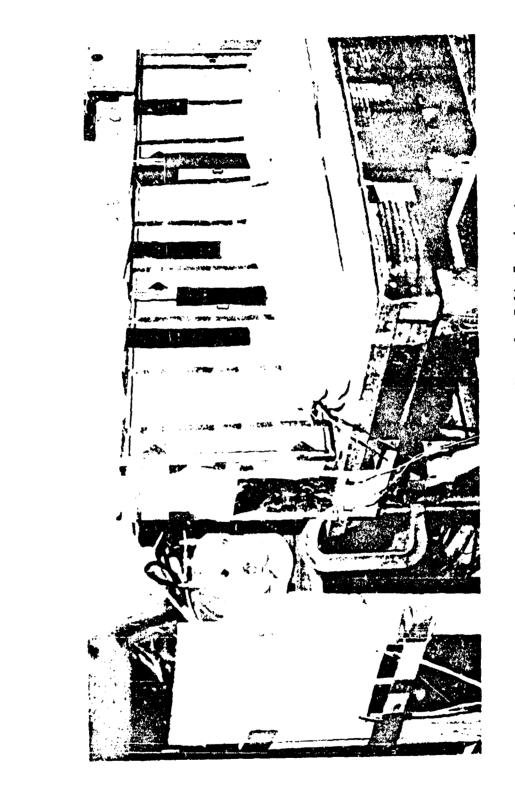
In the Eglin experiments the projectiles were fired horizontally into a test chamber consisting of an open-top box of nominal dimensions 0.15 meter wide by 0.40 meter high and 1.2 meters long. The side walls and floor of the box were made of 0.0023 meter aluminum sheet framed by steel brackets and mounted on a flat wooden table platform as show. in Figure 2. The ends of the test chamber were closed by fiber board that was easily penetrated by the projectiles. The test chamber was backed up by a large open-topped wooden box fitted with vertical slots to accommodate partitions. The partitions were used to fill the box with varying amounts of sand in order to contain the projectiles for re-use in the various velocity regimes tested.

The boxes were filled with Eglin sand that had been sieved with a U.S. Standard Sieve Series No. 25 sieve to remove large debris, but not sieved to a controlled size range. For the dry sand tests the sand was poured slowly into the test chamber from a bucket assembly attached to an overhead crane. The wet sand tests were for the fully saturated condition. For the wet sand tests the sand was first mixed with water in a container and then shoveled into the test chamber. It was maintained in a fully saturated condition by adjusting a flow of water into the open top to compensate for leakage and maintain an essentially constant water level.

Standard triaxial tests were performed on two samples of the Eglin sand. For these tests the sand was first carefully dried following procedures as described in Reference 25. Each sample was tested at three different constant values of the lateral confining pressure  $\sigma_3$  (0.1962, 0.392, and 0.589 MPa) with axial compressive stress  $\sigma_1$  increasing until failure occurred (significant increase of axial strain at constant load). The two samples were a loose sand and one compacted by vibration before testing. Table 2 lists the initial density  $\rho$  and the angle of friction  $\phi$ determined for each sample by analysis of the triaxial data as well as the value  $(\sigma_1 - \sigma_3)_f$  of the stress difference at failure for each of the confining pressures.

The curve of  $\sigma_1 - \sigma_3$  versus axial strain  $\epsilon_1$  for the loose sand at the highest confining pressure will be given in Section V, where it is used to determine the deviatoric properties for the penetration analysis by the spherical cavity expansion theory method. Several confined uniaxial strain tests were also performed on dry Eglin sand. These will discussed in Sections V and VII.





#### TABLE 2. TRIAXIAL DATA FOR DRY EGLIN SAND

	٥з	(01-03) f
Loose Sand	0.1962 MPa	0.538 MPa
ი <sub>0</sub> = 1519 kg/m <sup>3</sup>	0.392	0.983
φ = 33.4 <sup>°</sup>	0.589	1.447
Compacted Sand	0.1962	0.763
$\rho_{0} = 1698 \text{ kg/m}^{3}$	0.392	1.423
$\phi = 39.7^{\circ}$	0.589	2.02

The projectiles were fired into the test chamber with a 20mm gun. Firing velocity was controlled by varying the powder load in a primed 20mm case. Striking velocity was measured by timing the interval between the breaking of two paper back velocity screens with a Terminal Ballistics Data Acquisition System. The start screen was 1.22 meters from the front of the target, and the stop screen was 0.61 meter from the target. The timing signals were also recorded on a magnetic tape system and later transferred to a pape. oscillograph record. Also recorded on the tape was the signal from a break wire on the gun muzzle 3.54 meters from the target.

To monitor the projectile flight path through the sand in the 1.2meter-long box, five Hewlett-Packard flash X-ray units were used: one 150 KV soft X-ray unit and four 300 KV hard X-ray units. They were spaced sequentially along the horizontal, with the first unit (150 KV) located at 0.038 meter from the front of the box and the four 300 KV units spaced 0.38 meter between centers. Standoff distance for the 300 KV units was nominally 0.55 meter. The front ends of the i50 KV unit (small cylinder) and of one 300 KV unit (large cylinder) are visible in Figure 2.

Several types of X-ray film were evaluated in the course of the test program, with the majority of the data recorded on Dupont Lightning Plus X-ray film. The film cassettes (not shown in Figure 2) were mounted on the outside of the box opposite from the X-ray units. The film plane was positioned at 0.08 meter from the centerline. In the main test program a series of metal letters (A through Q) were taped along the box, separated horizontally by approximately 0.07 meter along a line 0.20 meter from the top of the box, to serve as markers for locating the projectile position in the X-ray pictures. Some of the earlier tests used fewer markers in the form of metal arrows. In some of the tests displacement of the soil medium was observed by suspending 0.0015-meter steel markers that moved with the sand. Preshot and postshot X-ray records were made to locate the initial and final positions of the markers.

Figure 3 is a photograph of the three prints made from the X-ray negatives for Shot No. 26. The two panels on the left show the nose of the flat-ended projectile in four successive praitions at the times of the sequential firing of the X-rays. The fifth one in the third panel does not show up well in the reproduction but could be seen in the negative. The separation angles and the cavity around the afterbody are clearly shown in three of the positions illustrated. The aft end of the projectile is not usually visible, since it is out of the main X-ray beam when the firing is correctly timed to show the nose.

A magnetic system was used to furnish supplementary velocity information after preliminary laboratory investigations at the University had established the feasibility of the method. The steel projectiles were magnetized to a strength of about 150 gauss, as measured at the center of the nose with a Hall-effect gaussmeter. When this magnetized projectile passed through a 0.15-meter-diameter coil mour ted inside the test chamber a voltage signal was generated. Four such coils .ere used in most of the Eglin tests, one on the front of the box and three inside at distances of approximately 0.49, 0.80 and 1.09 meters from the front of the box. The positions are recorded in Appendix A for the 26 shots for which the magnetic sensor data were analyzed. The projectile can be seen passing through two of them in the X-ray picture of Figure 3. Each sensing coil as formed with 40 turns of copper wire, forming a rim about 0.004 meter this. The voltage signals were recorded without any preamplification on the magn cic tape recording system and later transcribed to an oscillograph record. he records indicated voltage peaks of the order of 40 to 80 mv in tests with initial impact velocities around 200 m/sec. Some recovered projectiles showed a residual magnetic strength at the nose of around 20 percent of the value before firing.

Laboratory tests were made at the University on smaller diameter projectiles fired from an air gun. The time when the projectile nose arrived at the plane of the coil was precisely determined with a light beam. Comparison with the time of the peak voltage output showed that in tests with a coil formed by two parallel wires 0.025 meter apart the peak voltage occurred precisely as the nose passed through the loop. With a 0.165-meter-diameter coil a discrepancy was noted, indicating that the nose of the 0.215-meter-long by 0.0095-meter-diameter projectile has advanced approximately 0.021 to 0.027 meter beyond the coil plane when the maximum voltage was observed for the lowspeed shots in air (20 to 35 m/sec). Since a comparable direct check could not conveniently be made with the larger diameter projectiles used in the Eglin experiments, an indirect check was made by statically mapping the radial component of the magnetic field.

The mapping was first made for the laboratory projectiles to see if it agreed with the laboratory dynamic measurements. At a radial distance of 0.09 meter from the projectile axis the peak radial magnetic field occurred at a distance of 0.021 meter back of the plane of the nose, in approximate agreement with the discrepancies noted above. Similar mappings of one flat-nosed and one step-tier solid projectile of the type used in the Eglin experiments showed that the maximum radial component of the magnetic field occurred at distances of 0.020 meter and 0.022 meter respectively, back of the nose tip plane when measured at a radial distance of 0.086 meter from the projectile axis. This indicates that the maximum response should occur when the projectiles have penetrated some 0.02 meter through the plane of the sensor coil,





assuming a circular undeformed coil and a straight horizontal flight path through the center of the coil. Possible sources of additional error are imprecise measurement of the coil locations and especially of the time of the peak response, since with coils this large the response curve does not show a very sharp peak.

2.4 OTHER SENSORS USED OR TESTED

#### 2.4.1 Breaking-Wire Sensors

Both wire-grid and coated paper or plastic velocity screens are widely used to time the airborne part of a ballistic test. They have also been used buried in soil targets or sandwiched between slabs of rock or concrete. Because it was believed that the standard wire-grid screens might disturb the deformation patterns and force fields in the target, an attempt was made to develop wire sensors that would interfere less, by using finer wires in parallel arrays, less closely spaced than the screens. A developmental investigation at the University tested single wires impacted by projectiles fired from an air gun. Of particular interest was a method of verifying how much lateral motion of the wires occurred before they broke to give a signal. The 0.0095-meter-diameter projectile was 0.15 meter long. Two pinholes in the air gun barrel near the muzzle transmitted light to a photo-multiplier timing system for measuring projectile velocity. The dual-beam oscilloscope was triggered when the aft end of the projectile passed the first pinhole (farther from the muzzle). The wire sensor was placed 0.15 meter from the second pinhole, so that the projectile nose impacted the first wire just as the aft end passed the second pinhole. The time difference between the two signals (from the second pinhole and from the breaking wire) determined the time delay (or advance) of the breaking-wire signal. With the known projectile velocity, the position error that would be caused by assuming that the wire broke instantaneously in its orginal position could be determined.

Several kinds of wire were tested. The wires were stretched between supports 0.15 meter apart. The first tests were performed in air. Ductile wires of copper and stainless steel stretched so much that the projectile traveled almost 0.05 meter before the wire broke in impacts at 32 m/sec. Brittle wires gave better results. A brittle 0.0001-meter-diameter tungsten wire broke after about 0.0025 meter of travel.

Tests were then performed with the stretched tungsten wires buried in sand. The wires broke before the projectile reached them, because of the sand pushed ahead of the flat-ended projectile. At 39 m/sec the distance was about 0.006 meter and at 65 m/sec about 0.009 meter from the projectile to the initial position of the wire when it broke. In all cases the breaking wire gave a good sharp step on the oscilloscope trace. The last group of tests used 0.0002-meter-diameter steel music wire (static breaking force 89 N as compared to 17.8 N for the tungsten wire). In these last tests the wire sometimes did not break, but was deflected to one side of the projectile. A small perturbation in the voltage trace was noted at about the time the projectile reached the wire's initial position. Such a perturbation may be usable for timing purposes, although it lacks the sharp step that occurs when the wire breaks.

Two of the standard wire grid velocity screens were checked in the test apparatus at the University. With a 12.7 mm-diameter projectile impacting the screen in air at a speed 25.4 m/sec, the screen bent to allow about 0.0025 meter of travel before it broke. When z similar test was performed in sand, the travel appeared to be about 0.0075 meter before the break.

The figures quoted for distances from the initia<sup>1</sup> wire position to the projectile position when the break occurred apply, of course, only for the specific projectiles, wires and configurations tested. Sensors will have to be calibrated in conditions similar to those they are to be used in.

The first test firings at Eglin on 6 June 1975 were planned to test breaking-wire sensors and capacitor sensors prepared at Eglin. Tungsten wires and steel music wires of the types previously tested at the University were strung between supports 0.76 meter apart, and in addition two standard wire grid velocity screens were placed in the sand near the front end of the target. Signals were to be recorded both by counters and on magnetic tape recorders. No signals were obtained from any of the breaking-wire sensors. Post-test checks showed that the velocity screens were broken but the two other wire sensors were not broken by the 0.50 caliber projectiles, Similar wire systems could be used in sand, especially with the larger 20mm projectiles, but it would be necessary to check them out carefully with each projectile and test configuration. The X-ray method could be used as a check. Little further use of breaking-wire sensors in sand was made in the Eglin experiments because the magnetic sensors were so much better, and later the X-ray method gave still better results.

2.4.2 Sensors Responding to Pressure or Deformation: Capacitors, Pressure Transducers, and Strain Gages

1-

Although the major effort in the experimental program was directed toward recording trajectory information and the bow wave formation and cavity formation, several types of sensors were tried that could give some additional information about arrival times and intensities of the stress and deformation waves in the target medium.

A capacitive transducer was developed at AFATL, consisting of two thin metal foils separated by a layer of foam rubber and encased in a flexible electrical insulating material. When the sensor was compressed along with the surrounding sand a voltage change occurred across the charged capacitor. This furnished timing information about the arrival of the pressure wave. With suitable calibration it could also furnish quantitative information about pressure and deformation. It also served as a good antenna for detecting and recording the actual firing times of the flash X-rays.

A pressure cell in the bottom of the test chamber 0.127 meter from the front of the box also gave information about the arrival time and intensity of the pressure waves and furnished a good signal.

One to five strain gages were also mounted on the aluminum plates at the sides of the box. Good strong signals were obtained from the gages. The interpretation of these signals depends on the interaction between the pressure wave in the sand and a flexural wave in the plate.

#### 2.4.3 Microwaves

A microwave monitoring system was reported on in 1965 by Hakala (Reference 15). Its operation gives output depending on the interference between a transmitted signal and a signal reflected from the moving projectile. The transmitter and receiver were at the opposite end of the sand target from the impact point. The interference frequency is a function of the projectile velocity. Since few details about power requirements for penetrating various distances in dry and moist sand were available, an experimental program to determine some of this informatic, was undertaken at the University early in 1975.

A microwave oscillator of maximum power 1 mw fed a variable gain amplifier at a frequency of about 10 GHz through a coupler in a microwave horn into the sand contained in a 1.2-meter-long box. For these static experiments the signal generator carrier frequency was modulated by a 1 kHz square wave. The signal was reflected from the target, which was the end of s metal rod inserted into the opposite end of the box from the horn. Portions of the mixed incident and reflected signals were detected by a crystal detector. The detector output (DC with amplitude varying at the modulation frequency) was fed to a Standing-Wave Ratio Meter (which contained an internal amplifier with a narrow pass band around 1 kHz).

When the rod was moved axially by one quarter wave length the round trip path from coupler to rod was shortened by one half wave length. The reflected and incident waves interfered and a half wave length reduction in path was required for the detected mixed signal to go from a maximum to a minimum.

Preliminary tests showed a strong signal response at a distance of 0.30 meter with 3 mw power output from the microwave amplifier. At 0.60 meter the difference between the maximum and minimum response was down about 5 dB from the difference at 0.30 meter indicating a power transmission drop by about a factor of one-third, but the signal was still clearly distinguishable. In fact it was still clear at 1.0 meter. Precise attenuation factors could not be obtained with the preliminary test set-up, because of reflections from the sides of the box containing the sand.

Additional microwave studies attempted to repeat with naturally moist Eglin sand (approximately 5 percent by weight moisture content) the kind of measurements previously made in dry sand. With the target at a distance of only 0.15 meter however, the attenuation was so great that the alternate constructive and destructive interference by the reflected wignal as the target advanced a quarter wave length was barely perceptible.

Some additional measurements with a tuned microwave horn pickup replacing the target were made verifying that detectable microwave signals were transmitted through 0.30 meter of the moist sand even with the low power signal source.

It may be possible to increase the transmitted power and to obtain a coupler that will pick up a smaller fraction of the transmitted signal to mix with the reflected signal from the target in order to enhance the interference. For projectile velocity monitoring, the carrier wave would not be modulated by the 1 kHz square wave. A projectile advancing at constant speed at 300 m/sec will produce an interference frequency of approximately 20 kHz without any modulation of the original signal. This frequency will decrease as the projectile slows. The amplified signal could be recorded both on an oscilloscope and on one or two channels of a Biomation transient recorder. It would be recorded as a quasi-sinusoidal 'gnal of decreasing frequency. The time between a maximum and a minimum of this signal is the time for the projectile to advance one quarter wave length (of the order of 0.0075 meter although precise values would have to be estabilished by calibration). At a projectile velocity of 300 m/sec the time between maximum and minimum is about  $25 \times 10^{-6}$  sec, increasing as the projectile slows.

The microwave system was not actually used in the Eglin experiments since the X-ray equipment was available, but it is a possible option for future use if the power requirements can be met.

In Section III the data collected in the Eglin experiments will be described.

#### SECTION III

#### RESULTS OF EGLIN PENETRATION EXPERIMENTS

#### 3.1 INTRODUCTION

During the period from 22 January to 24 May 1976 the Eglin penetration experiments included a total of 91 shots in 17 missions. X-ray data from two or more stations were obtained in 74 shots (No's 14 to 91 except for Shots 21, 28, 60, and 75). Appendix A lists data obtained in Shots 14 through 91, except for the four shots for which X-ray data were not obtained. One page is used for each shot, and they are listed in order. A description of the various kinds of data in Appendix A, both the experimentally measured data and several kinds of information calculated in the data analysis, is given in paragraph 3.3 after an overview of the primary and secondary test programs in paragraph 3.2.

3.2 TEST PROGRAM MATRICES

#### 3.2.1 Primary Test Program Matrix

The primary test program at Eglin was planned to test two projectile configurations at three impact speeds and two target moisture conditions (dry Eglin sand and saturated Eglin sand). With four replications of each cest the plan called for 48 shots. Four extra replications brought the total to 52 shots as summarized in Table 3. The two projectile configurations were both solid cylinders of nominal diameter 0.0198 meter, one with a flat nose and the other with a step-tier nose, as described in paragraph 2.3. The wet sand was fully saturated. The letters after the shot numbers indicate that special analysis was made of those shots. The letter V indicates that veloc'ty data from the X-rays were fitted to a Poncelet force law as described in paragraph 4.4, with information about the fitting tabulated in Appendix A. The letter M indicates that the velocity results were compared with the magnetic sensor data obtained with the velocity coils, with results of the comparison listed at the end of each tabulation in Appendix A. The letter B indicates that a bow wave was observed in front of the projectile in one or more of the X-rays (see paragraph 3.5.1).

#### 3.2.2 Secondary Test Program Matrix

Shot numbers not included in the primary test program are listed in Table 4.

#### 3.3 DESCRIPTION OF EXPERIMENTAL DATA

For each of the 74 shots listed in Appendix A the position and cavity separation angle information obtained from the X-ray records is given in the first data group. An example is shown in Table 5.

Projectile Type and	Velocity Range			
Sand Condition	210 m/sec	320 m/sec	400 m/sec	
			14 B	
	15	20 VB	25 VB	
Solid	16	22 VB	26 VMB	
Flat Nose	17 V	23 VB	27 VB	
Dry Sand	18 V	24 VMB	29 VM	
	19 V			
	70 VM	36		
Solid	71 VM	37. VK	76 VM	
Flat Nose	72 VM	38 V	82 V	
Wet Sand	73 VM .	74 V	83 VM	
		81 VM	84 VM	
	52 V	56 VM	62 VB	
Solid	53	58 VM	63 VB	
Step-Tier Nose	54 VM	59 VM	64 VMB	
Dry Sand	55 VM	61 VM	65 VB	
	57 ¥			
	42	39 V	50 M	
Solid	43	40	51	
Step-Tier Nose	44	41	63 VM	
let Sand	45 V	49 VM	69 VM	
let Sand	45 V alculated veloci	49 VM ties of nose and	69 VM center of gravity are	

## TABLE 3. SHOT NUMBERS OF EXPERIMENTAL MATRIX FOR PRIMARY TEST PROGRAM

17

Projectile	Target	Impact Velocity Ranges			
Type	Medium	Shot Numbers			
Solid	Dry	320 m/sec			
Biconic	Sand	30 V, 31 V, 32 V			
Hollow	Dry	350 m/sec 400 m/sec			
Biconic	Sand	35 V 33, 34			
Solid Flat Nose 0.152 meter long	Dry Sand	77,	240 m/sec 78 V, 79 VM		
Hollow	Dry	230 m/sec			
Step-Tier	Sand	88, 89 V			
Hollow	Wet	440 m/sec			
Step-Tier	Sand	90			
Hollow Step-Tier	Water	230 m/sec 85, 86 VM, 87 VM, 91 V			
Special	Wet Sand	250 m/sec 46	550 m/sec 47, 48		
Model	Dry Send		66, 67		

TABLE 4. SHOT NUMBERS OF EXPERIMENTAL MATRIX OF SECONDARY TEST PROGRAM

pavison with the magnetic sensor data.)

#### TABLE 5. EXAMPLE OF FIRST DATA GROUP

1

SHOT 26 ( 12 MARCH, 1976, ND. 3 )

SAND: DRY, DENSITY:1538 Prujectile: Solid Flat N	KG/H++1 DSE MASS	1 APPRU	KG, D=0	CLUCITY:	406 M/S
X-RAY STATION	N0.1	5, 0N	N0.3	N0.4	N0.5
TIME (SECOND) CENTEN OF GRAVITY POSITI Horigontal	-000108 C.H. (M)	.000575	.001687	.003003	+004582
HORIZONTAL VERTICAL	-0.085 0.134	0.119 0.149	0.420	0,732	1.035
INCLINATION ANGLE (DEG). Separation Angle (Degree)	0.0	1.0	0.5	-1.5	-6.0
BELON	**** **** 0650.0	5.0 4.52 0.0228	3.5 4250	2,5 0.0230	0,5
MURIZUNTAL AAAAAAAAAA	0.027 0.134	0.232	0.533 0.147	0.843 0.150	1.147
INPUT NOSE PÜŚITIÓN (M) Morizontal Vertical	0.024 =0.081	0,232 -0,074	0,533 -0,068	0,844 =0,064	-0,067

The first line below the shot number and date gives the target medium: dry sand, wet sand or H-OH (meaning H<sub>2</sub>O for shots into water), the density in  $kg/m^3$ , the projectile's approaching velocity in meters per second as measured by the counter start and stop velocity screens as described in paragraph 2.3 and/or by the time from the break-wire on the gun muzzle to the X-ray trigger foil switch located on the front of the box as recorded on the oscillograph strip chart. The second line gives the projectile type, mass in kg, nominal diameter and length in meters. Below each X-ray station number is listed the time in seconds from the foil switch trigger to the firing of that X-ray. The firing times were determined from the delay settings and also by noting signals appearing on the strip chart records for various sensors. The next two lines give the calculated center of gravity position coordinates in meters measured from the front and bottom of the target box. Note that the first horizontal coordinate is negative, since the projectile is still outside the box.

The projectile's angle of inclination to the horizontal and the cavity separation angles at the nose, measured with respect to the projectile axis, are listed in the next two lines of data. A row of asterisks ( $\frac{1}{2}$ indicates missing data. These angles were measured on the X-ray negatives, as was the measured nose width in meters. The next two lines are calculated nose position coordinates, corrected for X-ray beam divergence. The last two lines give the raw data on the nose position in the X-ray picture, with the vertical position measured from the row of letter markers on the wall of the box, so that a negative coordinate indicates distance below the letters. For 22 shots (those not marked with a V or M in the experimental matrices of paragraph 3.2, no further data are tabulated in Appendix A. The other 52 shots have additional groups of calculated data to be discussed in paragraph 4.3, and the 26 marked with an M in the experimental matrices of paragraph 3.2 include magnetic sensor data in the last data group, as shown in Table 6.

#### TABLE 6. EXAMPLE OF LAST DATA GROUP (SHOT NO. 26)

RECORDED TIME OF MAXIMUM/MINIMUH CU	L VOLTAGE (8)
MAX         000099         001394           MAX         000099         001394           HIN         000625         002307           CDHPUTED NUSE POSITION AT MAX/MIN CO         0.509           AT MAX         0.024         0.509           AT MIN         0.214         0.509	002864 003947 005650 011 VULTAGE (M)
COMPUTED NUSE POSITION AT MAX/MIN CO AT MAX 0.024 0.509 RECORDED COIL POSITION (M)	
RECORDED CUIL POSITIUN (D/	
DIFFERENCE BETWEEN COIL AND NOSE AT AT MAX 0.024 0.023 AT MIN 0.214 0.03	0.778 MXX/MIN VOLTAGE (M) 0.035 0.251 0.273
AT HAX 0,024 0,023	0.035 0.042

The first line of data in this last data group (see Table 6) lists the time in seconds from the X-ray trigger coil switch time to the maximum voltage from the four magnetic sensor velocity coils. The second line lists the times of the minimum voltage at each of the coils. These times were transcribed from the strip chart. The next two lines are computed nose positions, as will be described in paragraph 4.4. The next line lists the actual positions of the four coils as recorded in the log book, and the last two lines record the differences between the two computed nose positions and the actual positions. The significance of these differences will be discussed in paragraph 4.4.

#### 3.4 PRELIMINARY ANALYSIS OF TABULATED DATA

#### 3.4.1 Nose Positions

Nose positions as measured on the X-ray photos were recorded as INPUT NOSE POSITION in the first data group of the tabulations of Appendix A (see Table 5) and also the apparent nose width. This apparent nose width as compared to the known actual nose width provides a first-order correction for the divergence of the X-ray beams. A simple computer program, based on similar triengles with apex at the X-ray source, was used to correct all apparent horizontal and vertical distances in proportion to the known correction for nose width. The corrected nose positions are tabulated immediately above the raw data input nose positions.

#### 3.4.2 Center of Gravity Position

The center of gravity position was calculated from the corrected nose position and the (uncorrected) inclination angle by using the known distance from the nose of the projectile to its center of gravity. This correction did not account for projectile yaw. Yaw was believed negligible because of the straightness and lateral stability of the trajectory. Further data analysis is given in Section IV.

#### 3.5 DATA NOT ANALYZED

#### 3.5.1 Bow Waves

Several of the X-rays showed a detached shock wave ahead of the projectile, revealed by a density discontinuity. This occurred notably in the higher speed impacts in dry sand. In the primary test matrix of Table 3 shot numbers marked with a B showed well defined shock waves. Thus the flat-nosed projectiles showed shock waves in the intermediate velocity range also.

The bow shock wave appeared as a roughly parabolic curve (almost a circular arc near the vertex) with vertex at a distance ahead of the projectile nose of the order of magnitude of the projectile diameter. The X-ray pictures have been retained for possible use in future theoretical analysis of the deformation. Figure 4 shows tracings of two of the bow waves ahead of the projectile nose in two positions in Shot 26. The distance between the two positions is not to scale in the figure, but the position of each bow wave is shown relative to the nose.

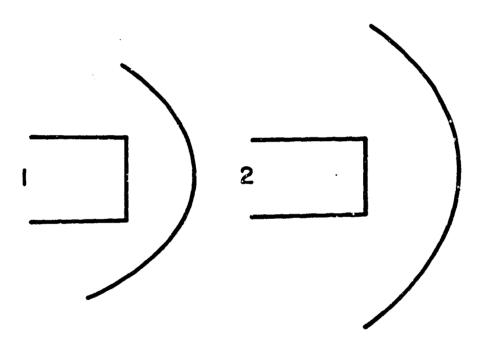


Figure 4. Tracing of a Bow Wave for Shot No. 26

Similar shock waves were also observed in some of the preliminary shots with 0.50 caliber projectiles.

### 3.5.2 Separation Angles

In some shots the separation angles above and below appeared to be approximately symmetric with respect to the nose velocity vector, which was slightly different from the projectile heading as given by the recorded inclination angle. In almost all of the cases recorded in the primary test program only the nose was in contact with the sand. In future analyses it may be possible to relate the separation angles to the shape of the false nose of sand formed in front of the flat-nose projectiles and/or to the lift forces exerted on the nose.

#### 3.5.3 Marker Movements

The preshot and postshot X-rays showing movement of the small steel markers have not been analyzed. This information may be useful for evaluating future theoretical analyses of target medium deformation.

#### 3.5.4 Pressure and Strain-Gage Measurements

The pressure transducer strip chart records may furnish useful data to compare with the observed bow waves and/or with future theoretical analyses of stress and deformation wave propagation in the target. The strain gage measurements on the aluminum test chamber walls were also recorded on the strip chart. These strain pulses are also related to the pressure wave in the sand, but are strongly influenced by the response of the aluminum plate to a traveling and varying dynamic load.

#### SECTION IV

#### CLASSICAL AND EMPIRICAL ANALYSIS OF EXPERIMENTAL DATA

#### 4.1 INTRODUCTION

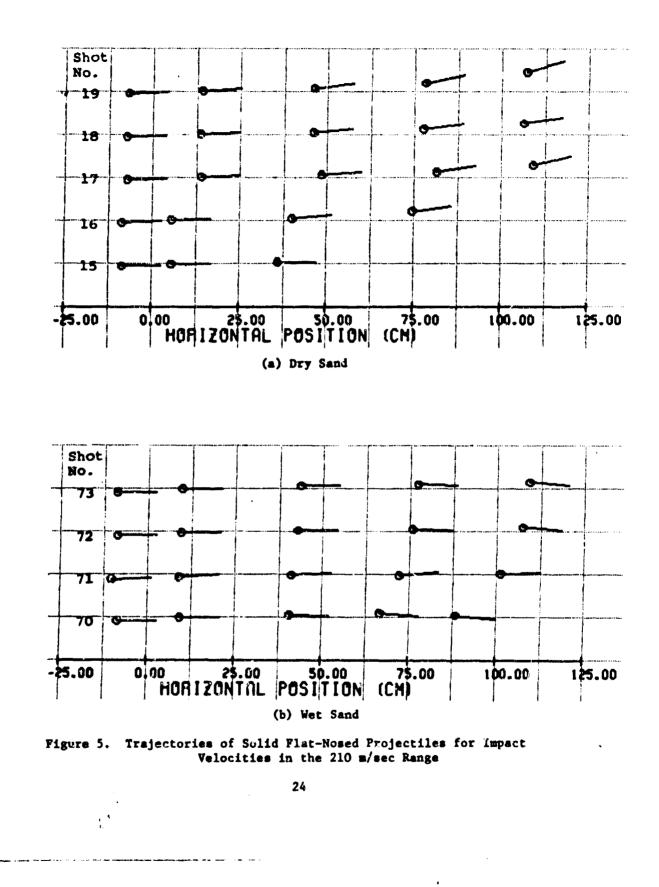
Trajectory plots for the 52 shots of the primary test program are given in paragraph 4.2. Since these trajectories are very nearly straight and horizontal throughout the 1.2-meters-long region of observation, analysis by one-dimensional penetration models is feasible. In paragraph 4.3.1 computer plots of horizontal positon versus time and of velocity versus position are given for 21 of the shots. These were obtained by first fitting a cubic polynominal interpolation formula to the position data and then fitting a Poncelet force law to each shot, as described in paragraph 4.3.2, which also contains a comparison of the values obtained for the Poncelet drag coefficients of the 41 shots of the primary matrix that have been analyzed by this method. Results of magnetic sensing are compared with the X-ray data in paragraph 4.4. This data analysis was performed at the University.

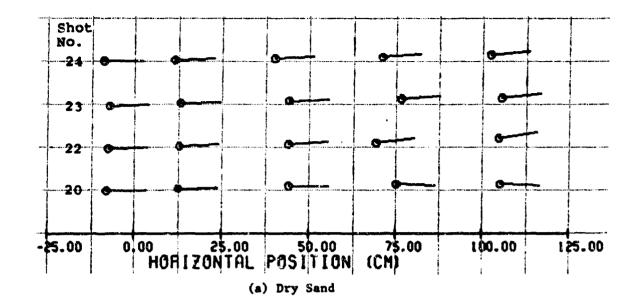
Results of a different method of determing the Poncelet coefficient for each shot and also results of empirical analysis by methods similar to those developed at Sandia Laboratories (References 12,13) are given in paragraph 4.5. Variation of the drag coefficient within a shot is discussed in paragraph 4.6 by considering separately different segments of several trajectories. These last two data analyses were performed at the AFATL.

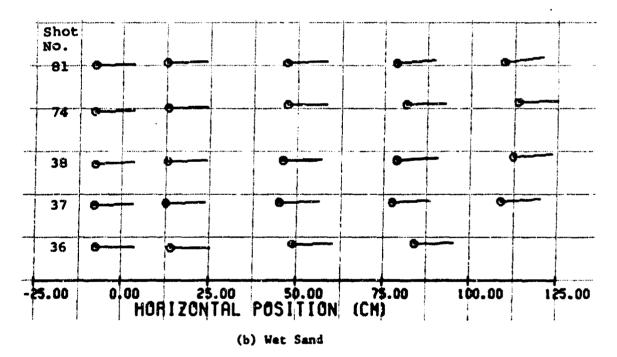
#### 4.2 TRAJECTORIES OF PRIMARY TEST PROGRAM

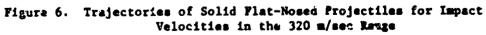
Computer-plotted trajectories for the 52 shots of the primary test program matrix listed in Table 3 are shown in Figures 5 through 10 based on the X-ray data for the positions and inclination angles. In each plot the circles mark center of gravity positions and the other end of the line from the circle is the nose positon. Shot number is shown at the left end of each trajectory. The horizontal and vertical scales are the same, but each successively numbered trajectory in a figure is plotted displaced upward one square (12.5 cm) from the preceding one. The plots give a pictorial summary of the trajectory data. Precise positions are given in the tabulations of Appendix A.

The most remarkable feature of the trajectories is their straightness, following in most cases a nearly horizontal straight line through the 1.2-meters-long target box. All but one of trajectories have a slight upward trend. The greatest rise, 6.2 cm, occurred for Shot 19 in Figure 5(a). Shots 16 to 19 of this group for the solid flat-nose projectile impacting dry sand at about 210 m/sec all show a continuously increasing angle of inclination, reaching 16.5 degrees in Shot 19. This was the largest inclination angle recorded. Positive final inclination angles were recorded for 31 shots, negative for 19 shots and zero for one, Shot 76 in Figure 7(b). In dry sand the flat-nose projectiles showed 9 positive and 4 negative final inclination angles while the step-tier



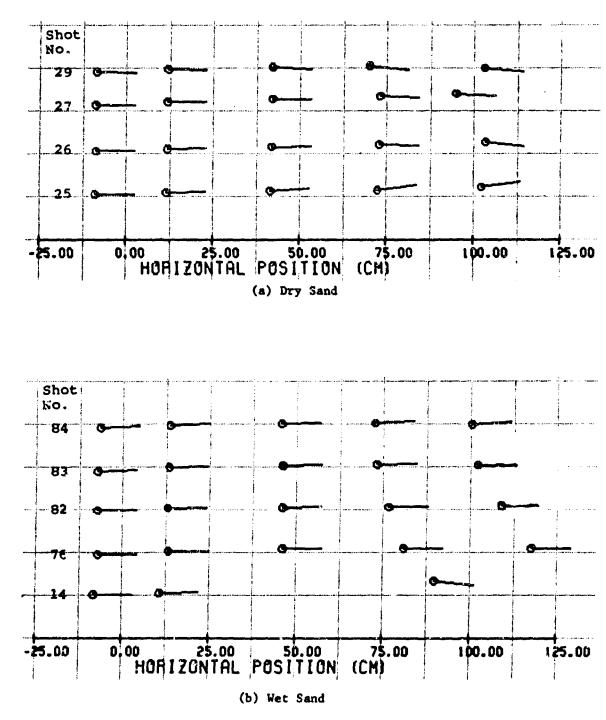


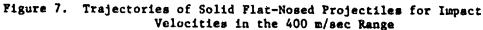


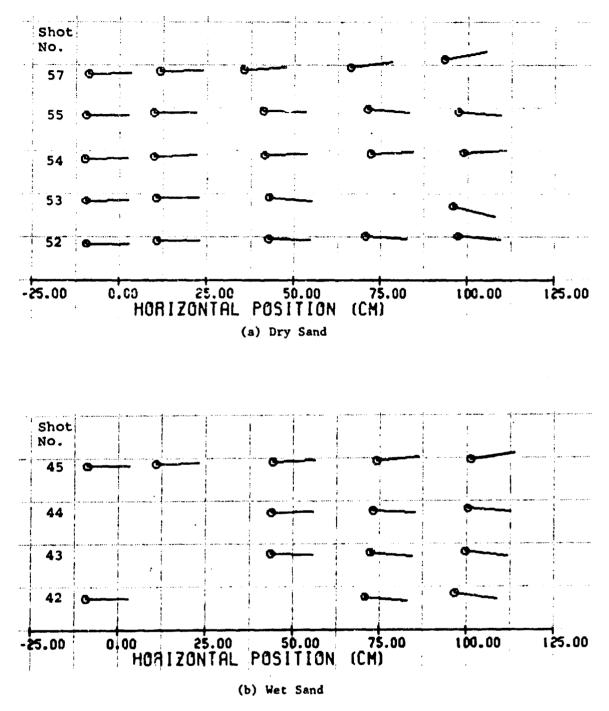


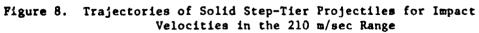
•

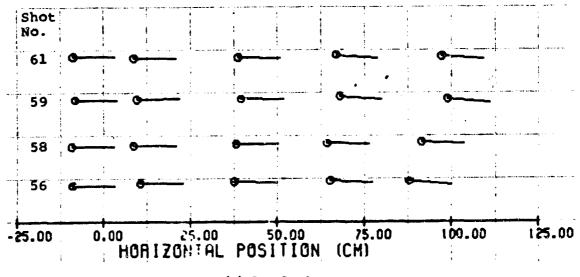
· 25











(a) Dry Sand

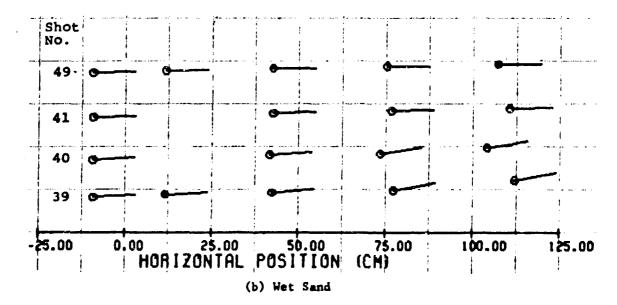
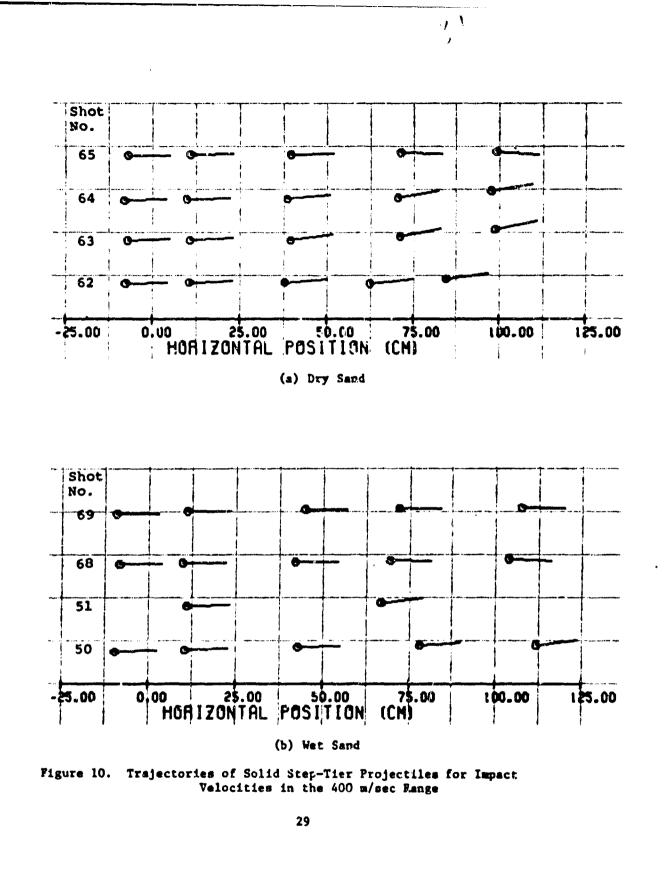


Figure 9. Trajectories of Solid Step-Tier Projectiles for Impact Velocities in the 320 m/sec Range



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projectiles showed 5 positive and 4 negative. In saturated sand the fiatnoses showed 11 positive and 2 negative final inclination angles while the step-tier noses showed 7 positive and 5 negative. The largest negative angle was -14.5 degrees in Shot 53 of Figure 8(a). This was also the only trajectory that did not rise.

Some of the trajectories show a continued rise, even with a negative angle. This is most evident in Shots 27 and 29 of Figure 7(a) for the flat nose impacting dry sand at 400 m/sec. The X-ray pictures show that during the part of the trajectory observed only the flat nose was in contact with the sand, so that no forces were acting on the afterbody surface.

Since the trajectories are so straight in the region of observation, analysis by one-dimensional penetration models is reasonable, for example the Poncelet force law to be discussed in paragraphs 4.3 and 4.5 and the Cavity Expansion Theory penetration model in Section V.

4.3 ANALYSIS AND DISCUSSION OF TABULATED POSITION-TIME RESULTS BASED ON X-RAY DATA

4.3.1 Cubic Interpolation, X, t-plocs and V, x-plots

Velocity analysis has been carried out for 52 shots for which complete X-ray data (5 stations) were available. This includes 41 shots from the primary test program and 11 from the secondary program (those marked with a V in Tables 3 and 4.) In the data reduction at the University of Florida, the velocity analysis was performed first by fitting a cubic interpolation formula to the data, and later improved results for the x-component of the center of gravity velocity were obtained by fitting the data to a formula derived from the Poncelet force-law penetration model. The Poncelet law fittings will be discussed in paragraph 4.3.2.

The coefficients of each cubic polynomial are listed in the tabulations of Appendix A.

For example, for Shot 26, the tabulated coefficients for C.G. VELOCITY X-COMP. imply the polynomial

$$\mathbf{x} = -0.1277 + 393.3t - 46,190t^2 + 3,432,000t^3 \tag{1}$$

for  $\pi$  in meters and t in seconds. The coefficients were determined by a least-squares fit. At the end of each set of tabulated coefficients is listed the standard deviation in maters [square root of the sum of the squares of the differences between the measured x-positions and those calculated by the cubic at the times of firing of the X-rays]. For Shot 26 the standard deviation is 0.0014 meter for the x-component, indicating a very good fit.

Velocities were calculated from the fitted cubics. For example, for Shot 26 the center of gravity x-component velocity is given by

$$V = \frac{dx}{dt} = 393.3 - 92,380t + 10,296,000t^2 \text{ m/sec}$$
(2)

This should give a good approximation to the velocity near the center of the interval, but larger errors would be expected at the ends. Computer plots of: (a) the calculated x,t-curve and (b) the calculated  $\forall,x$ -curve for 21 shots of the primary test program are shown in Figures 11 through 52. On each x,t-plot the five experimental data points are marked by squares. The solid curve is the fitted cubic, and the curve marked with vertical strokes is a curve based on the Poncelet force law, to be discussed in paragraph 4.3.2. It is seen that the fitted cubic x,t-curves agree very well with the experimental data. The cubic and Poncelet x,t-curves are also close to each other through the whole interval. Their slopes begin to differ at the ends of the intervals of observation. The V,xplots by the two methods therefore show considerable differences at the ends. The cubic interpolation would give completely unreasonable results outside the interval of observation (0 to 1.2 meters).

4.3.2 One-Dimensional Analysis of Velocities by Fitted Poncelet Force Law

The Poncelet force law (Reference 2) takes the following form, after dividing through by the mass m of the projectile,

$$-\frac{\mathrm{d}V}{\mathrm{d}t} = A + BV^2 \tag{3}$$

where A and B are parameters depending on the target material as well as on m. For the high velocities in the interval of observation in the present program, the contribution of A is negligible, and the V,x-curves have been fitted by taking A equal to zero and determining a best fit for B by a nonlinear regression procedure that minimizes the standard deviation from the experimental data of the x,t-curve obtained by integrating Equation (3).

Equation (3) can be integrated explicitly for given initial data  $(V_{\alpha}, x_{\alpha}, t_{\alpha})$  to obtain

$$\nabla = \left[ \left( \frac{A}{B} + \nabla_{0}^{2} \right) e^{-2B(x-x_{0})} - \frac{A}{B} \right]^{1/2}$$
(4)

or, with A = 0

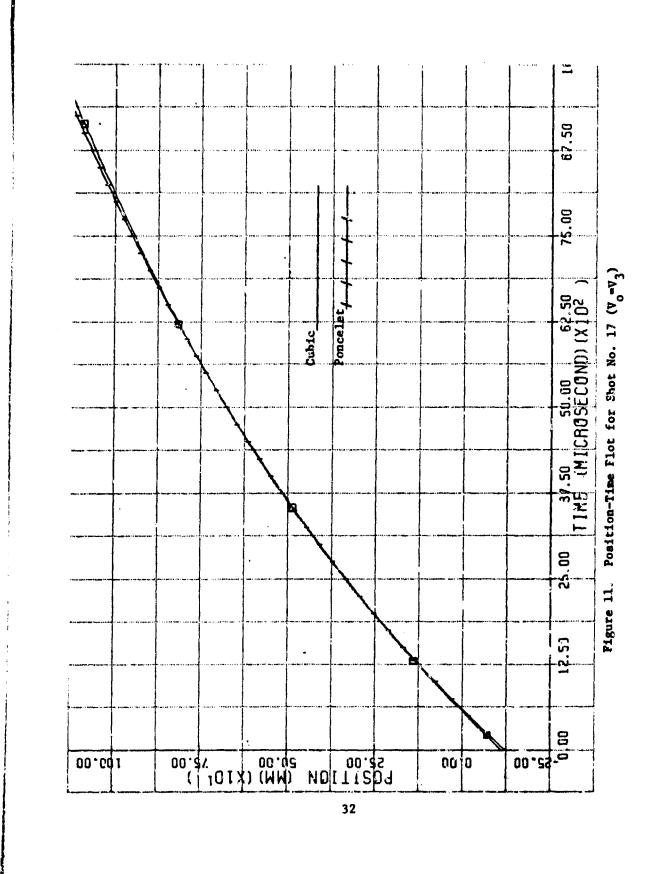
$$\mathbf{V} = \mathbf{V}_{o} e^{-\mathbf{B}(\mathbf{x} - \mathbf{x}_{o})}$$
(5)

With V = dx/dt a second integration of Equation (5) gives then

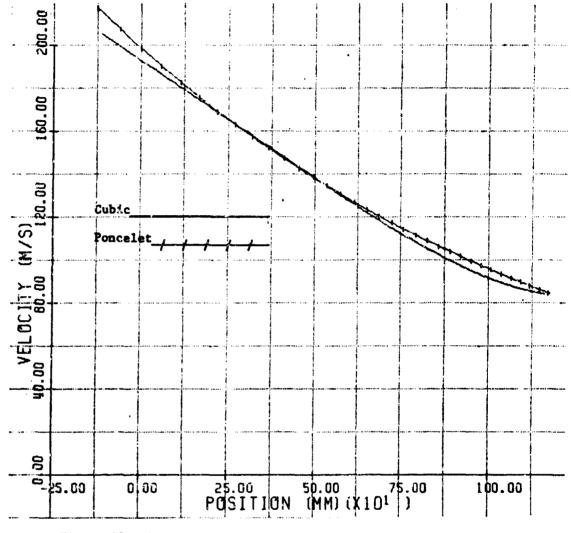
$$x - x_{o} = \frac{1}{B} \ln \left[ 1 + E V_{o} (t - t_{o}) \right]$$
 (6)

The more complicated case, with  $A \neq 0$ , is discussed in Section V on the Cavity Expansion Theory.

Equation (6) was fitted to the experimental x,t-data. A nonlinear regression is required. The procedure followed in this section was to take initial conditions  $x_0$ ,  $V_0$ , and  $t_0$  from the experimental data and the cubic

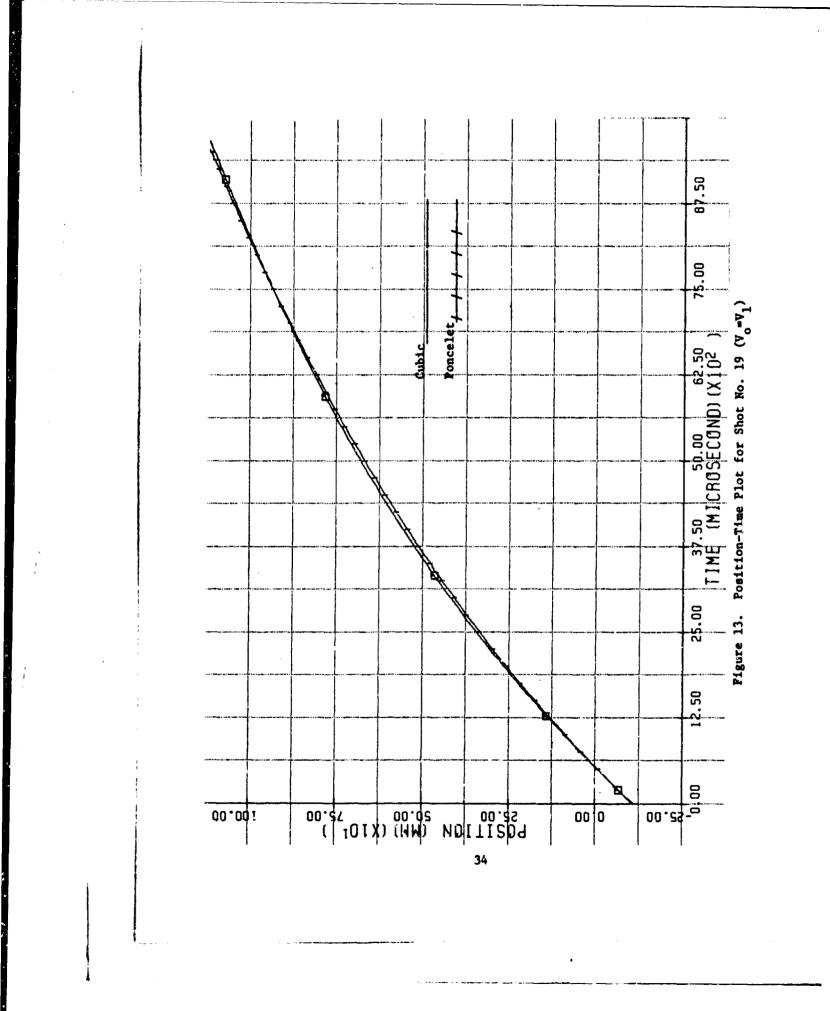


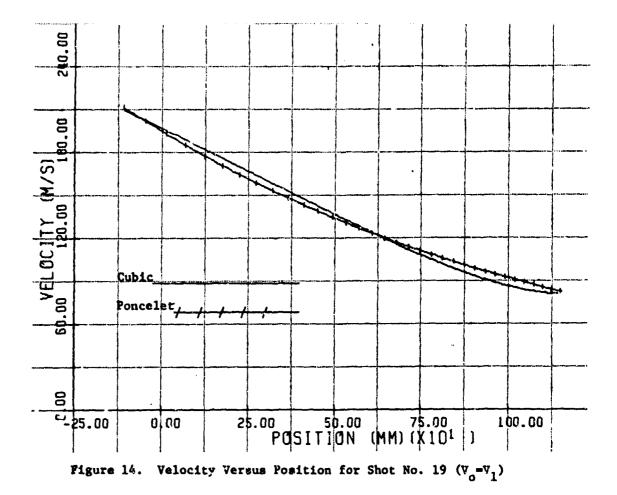
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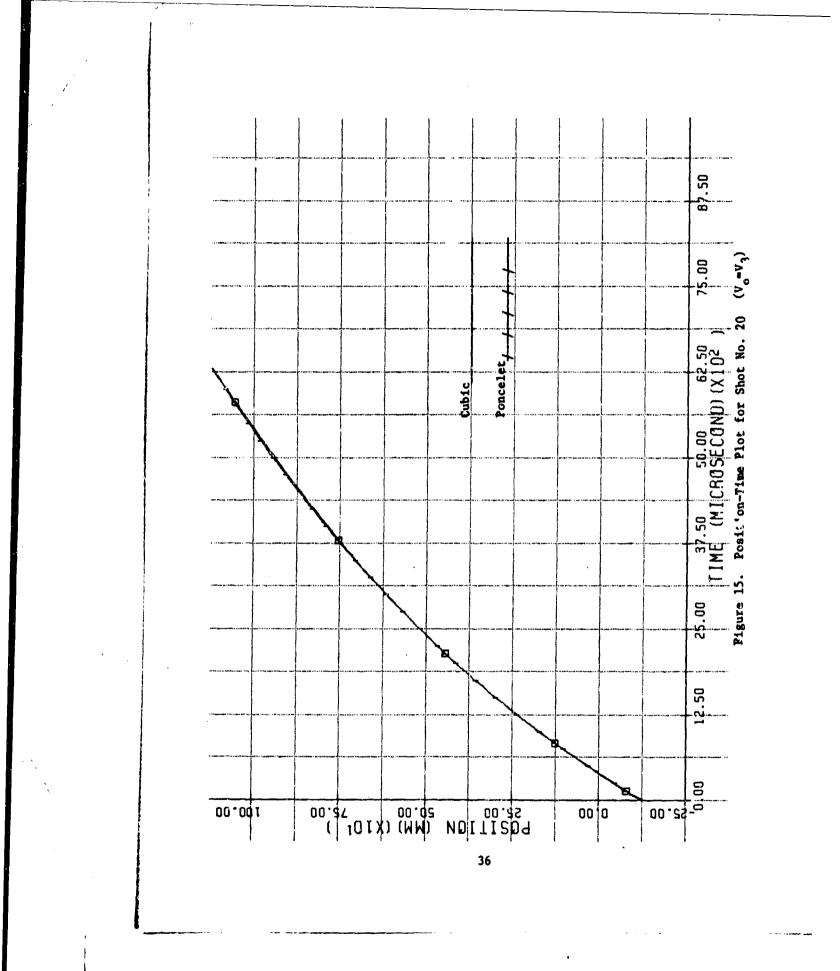


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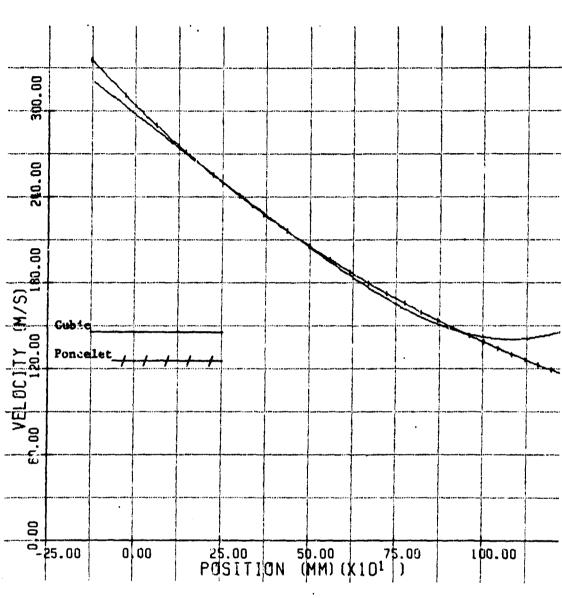




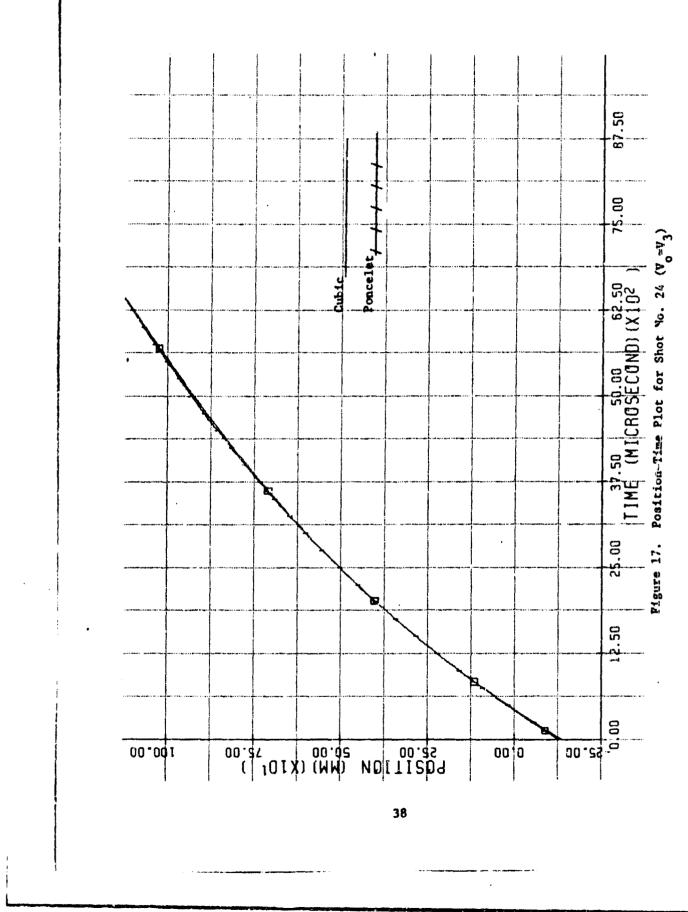


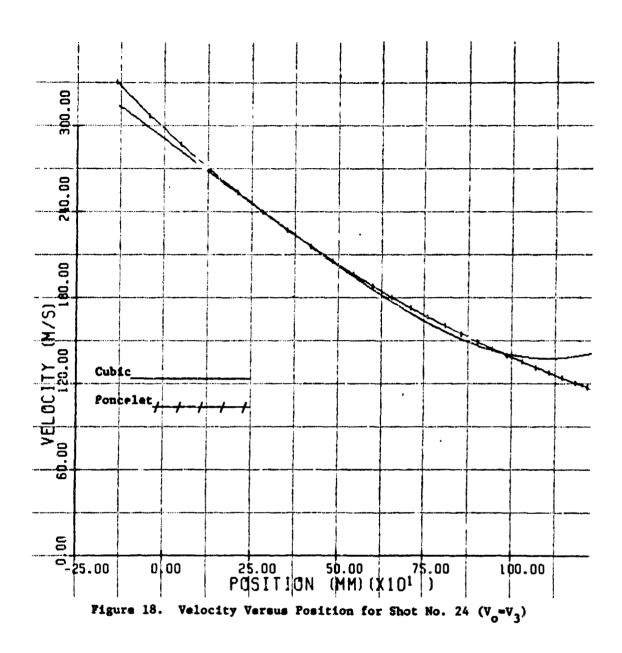


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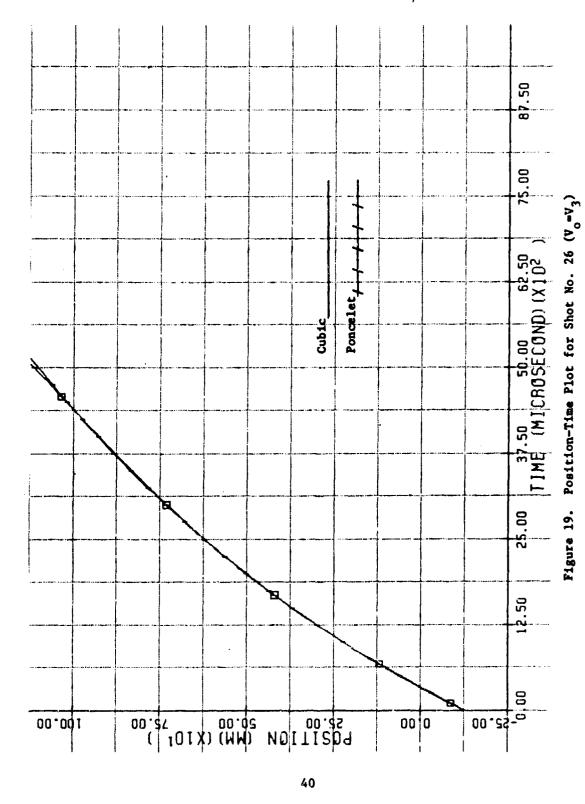


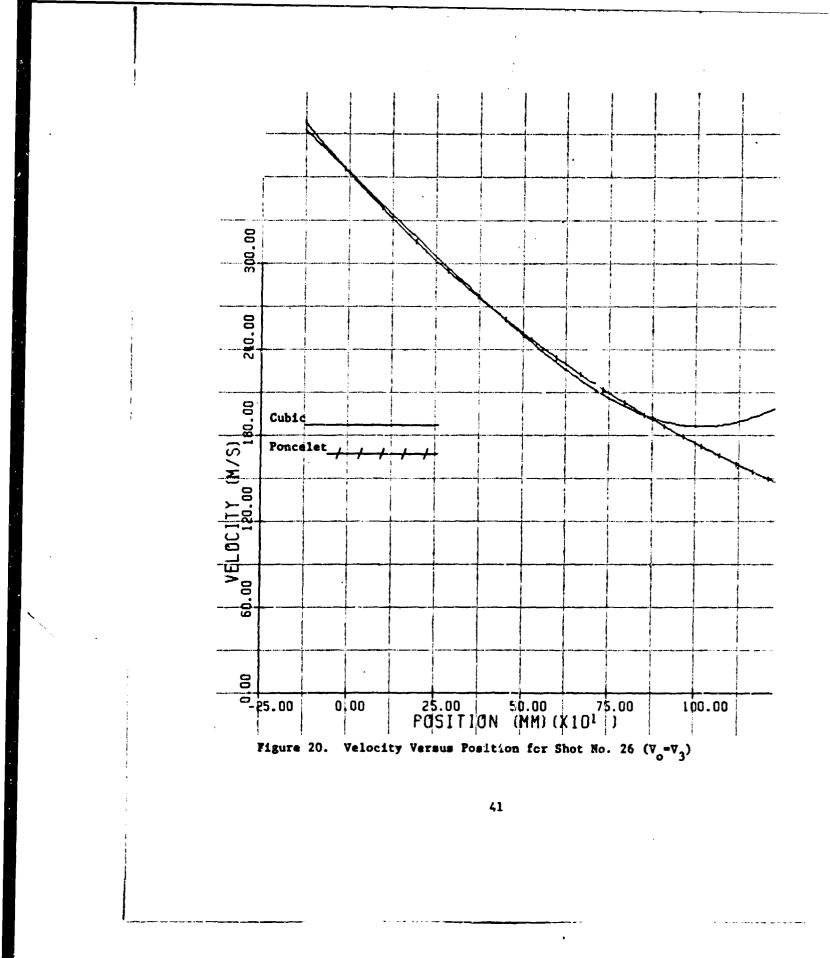


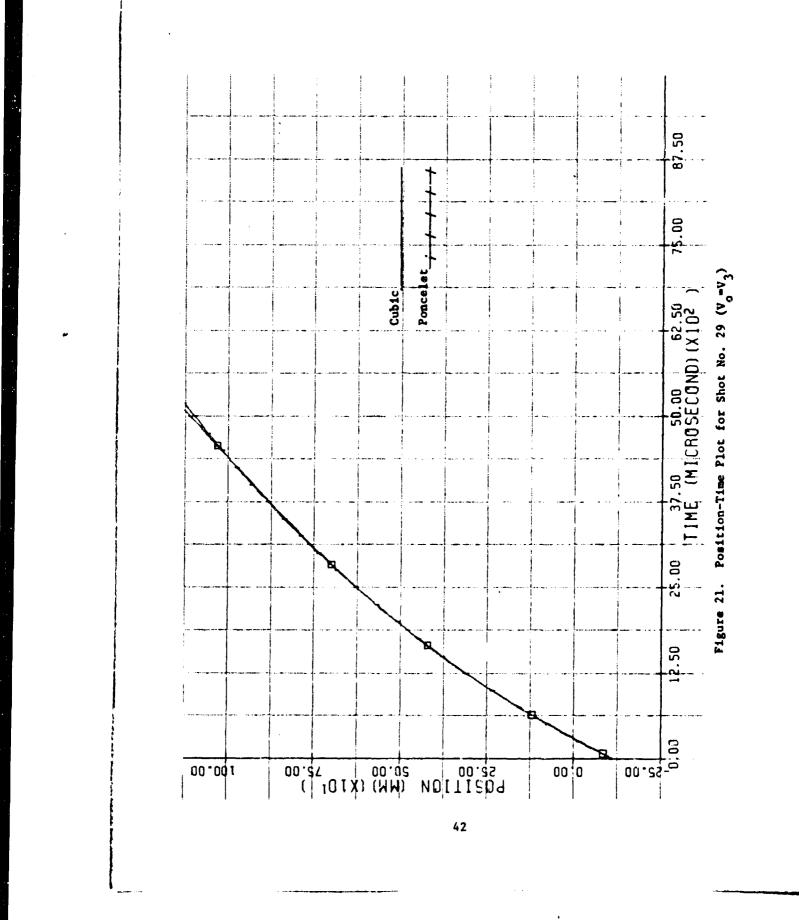


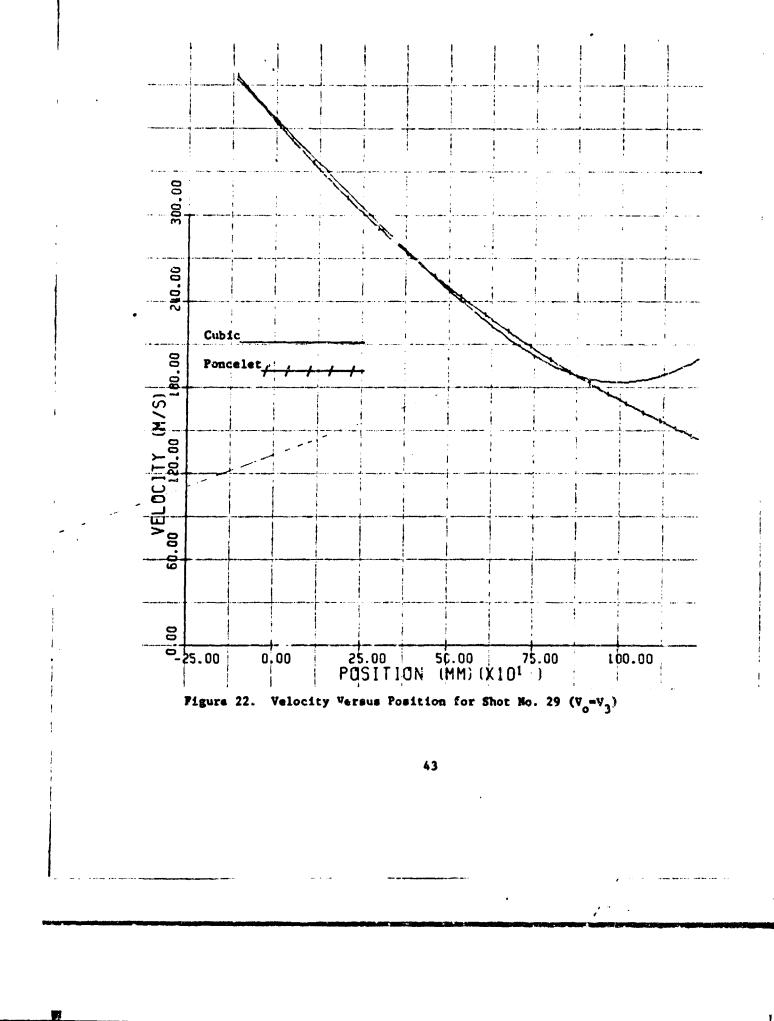




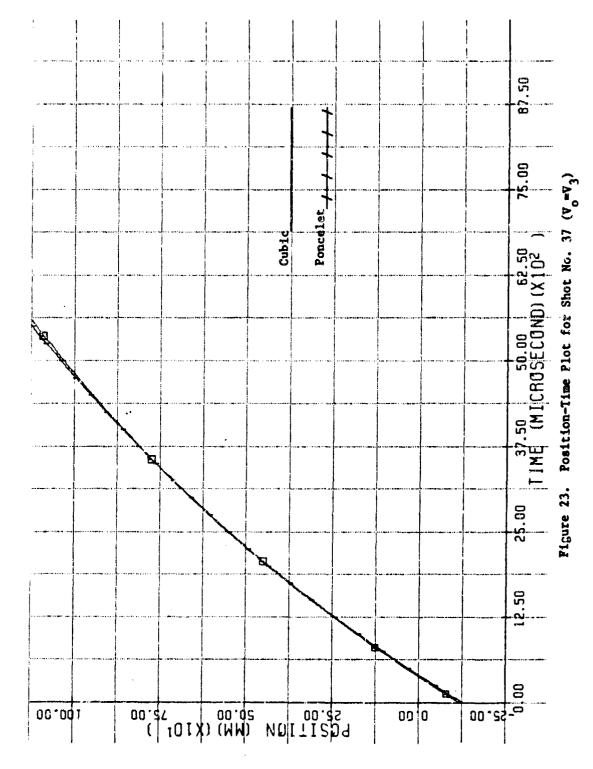


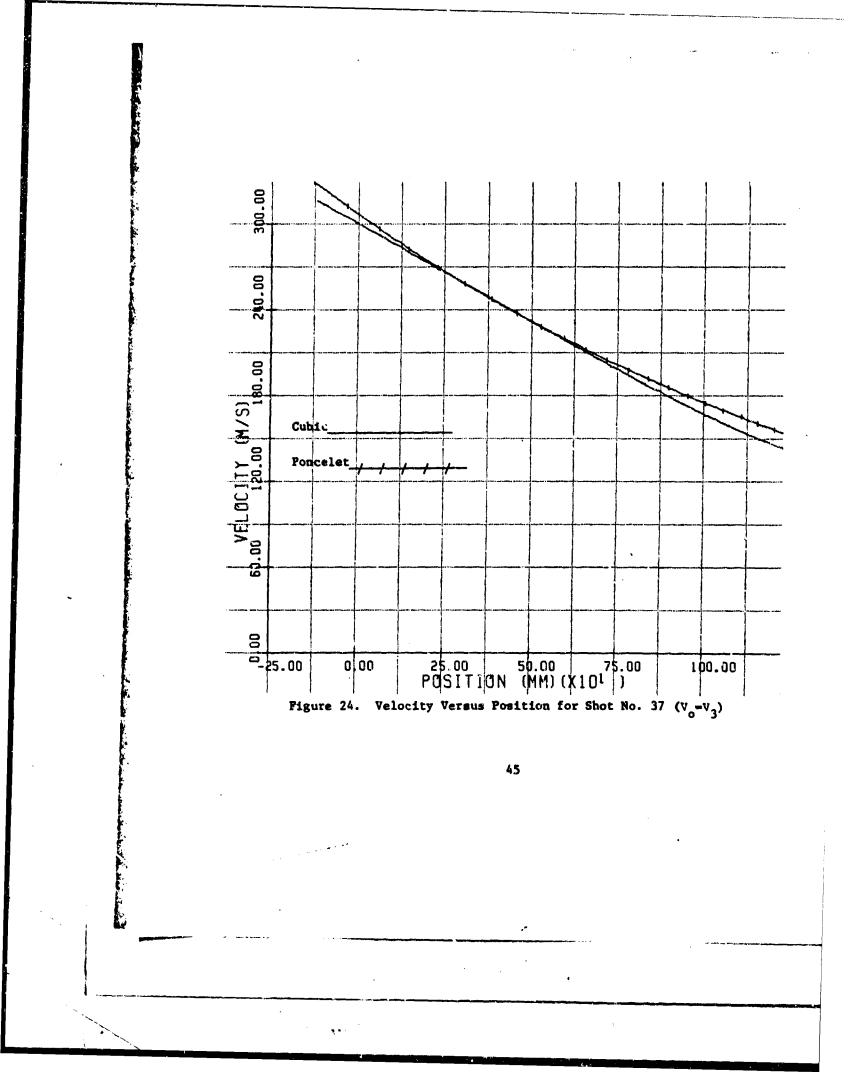


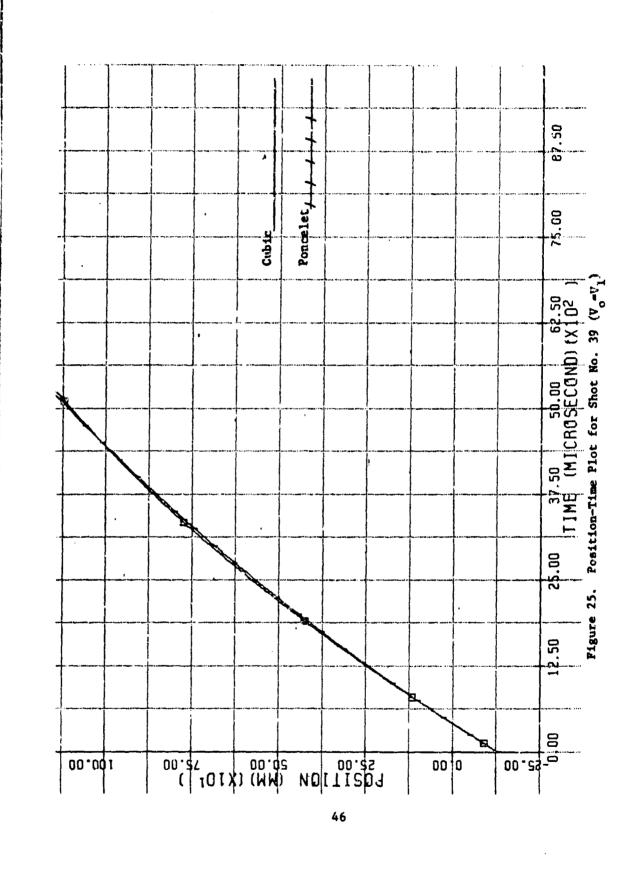


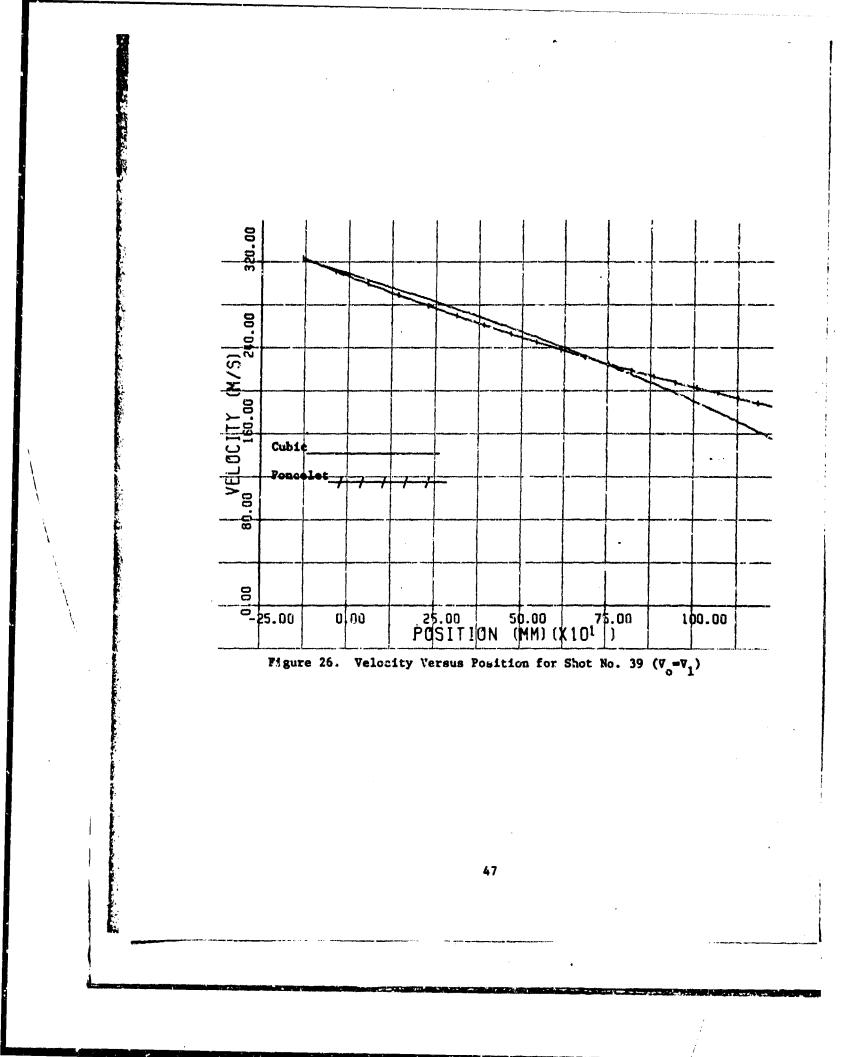


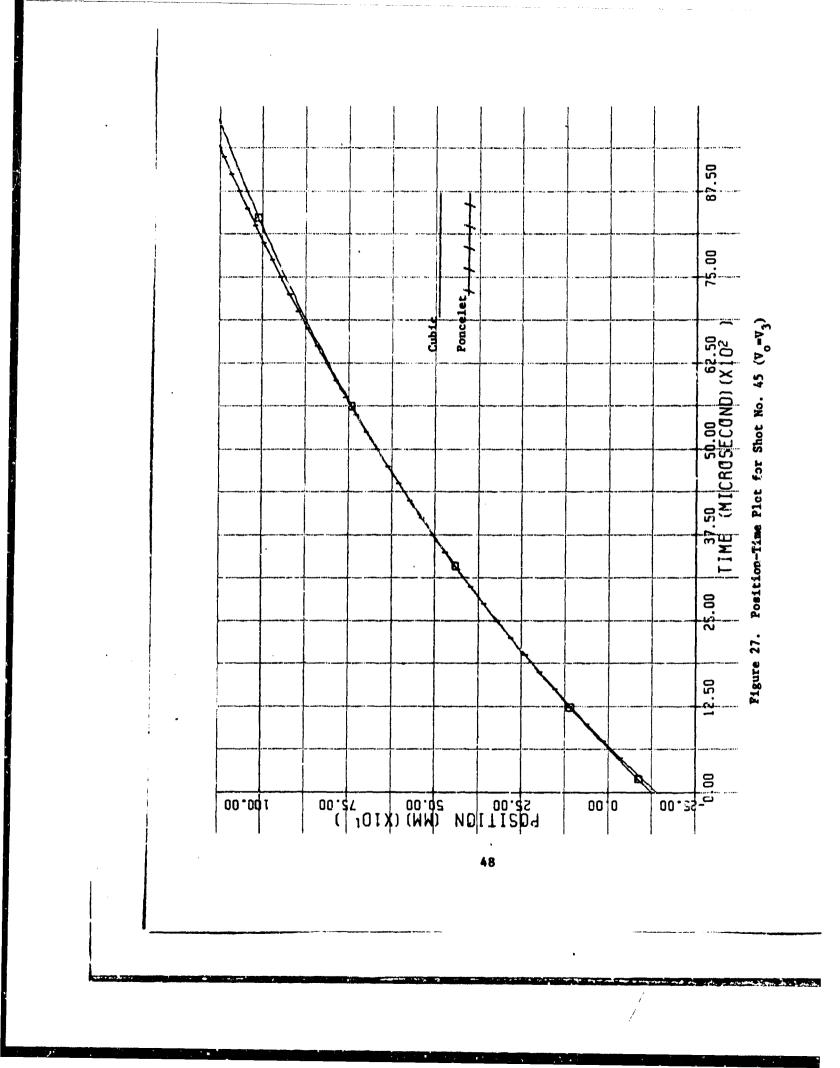
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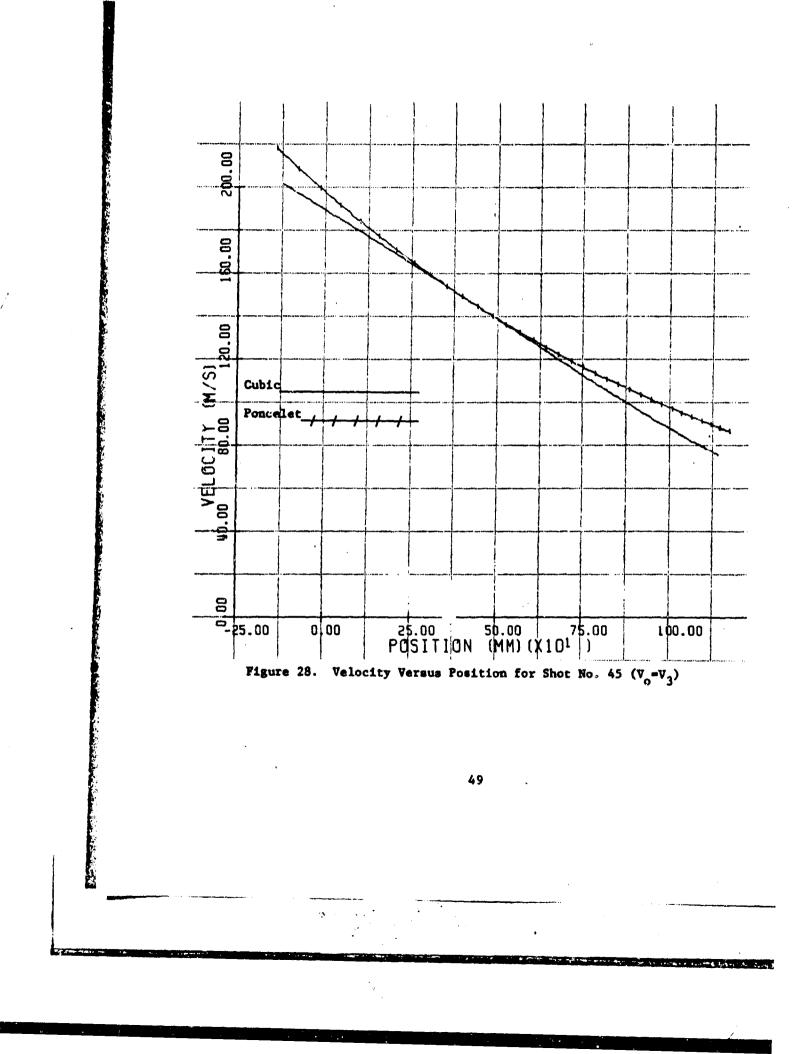


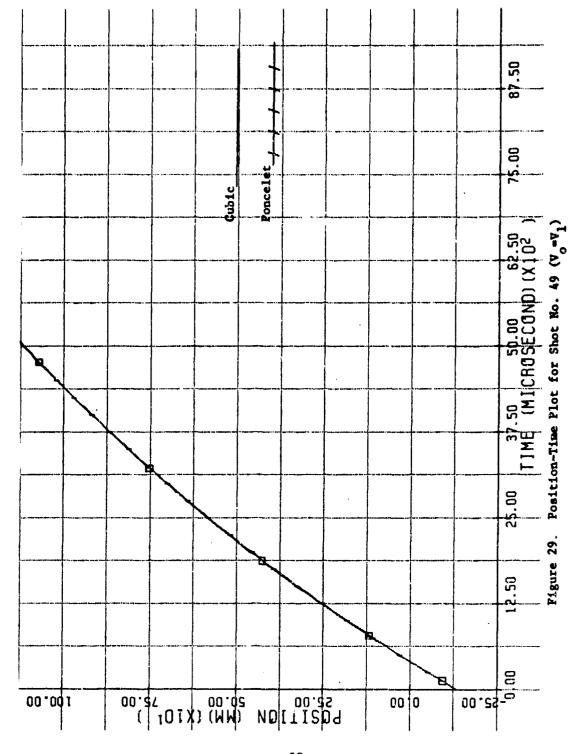


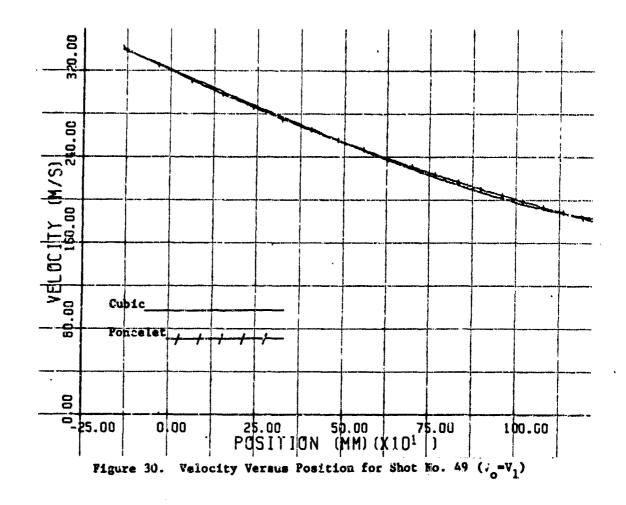




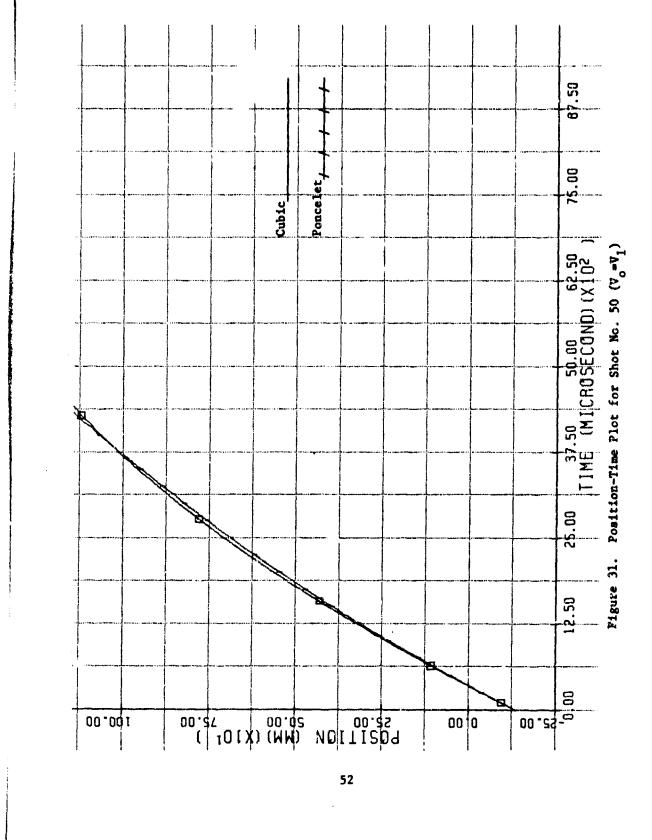




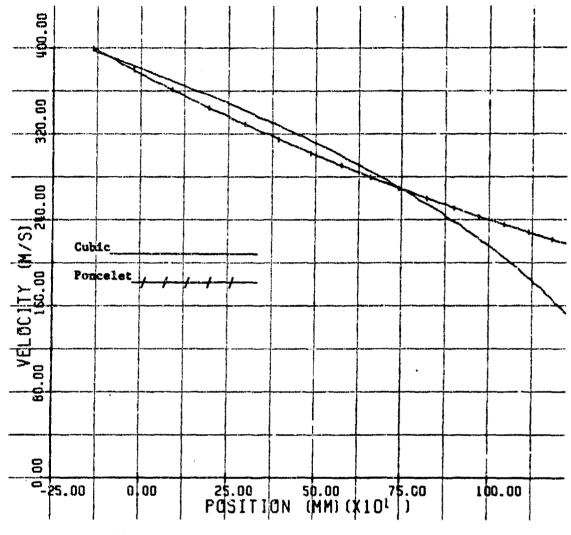




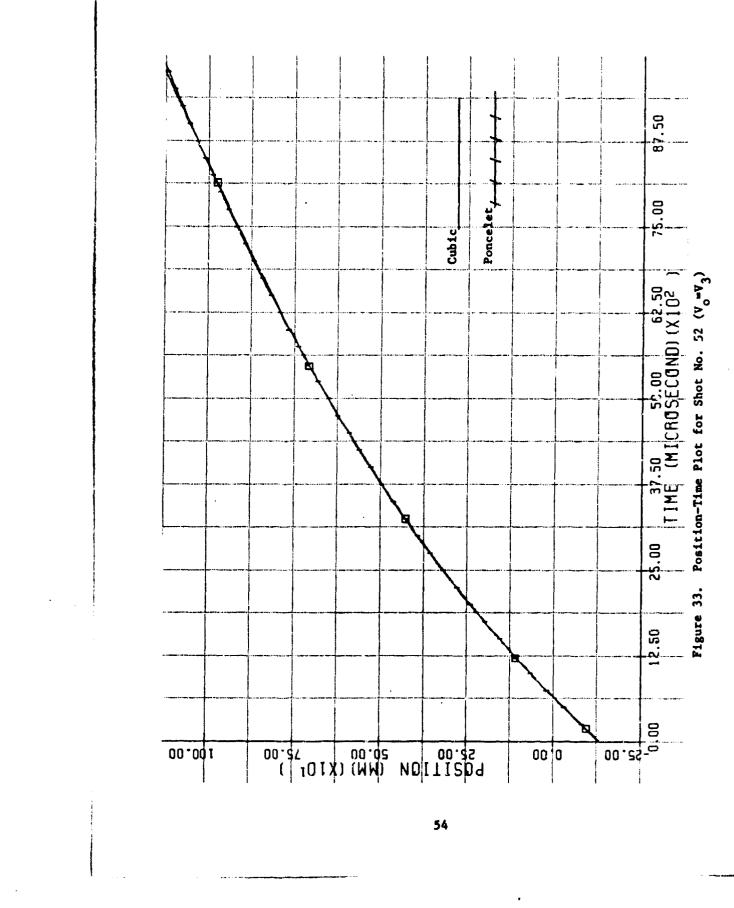


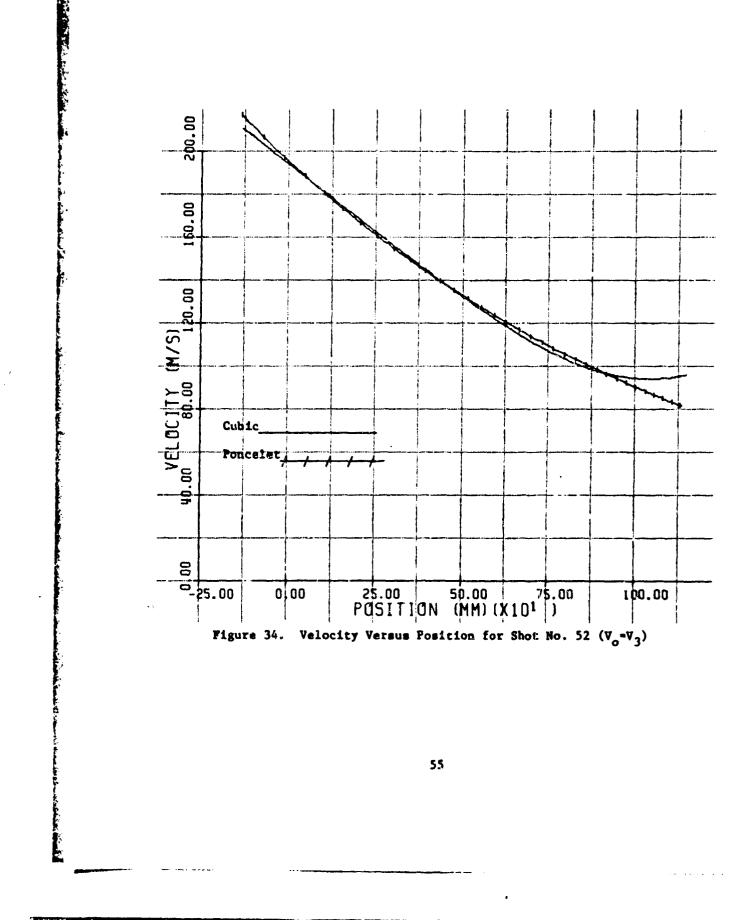


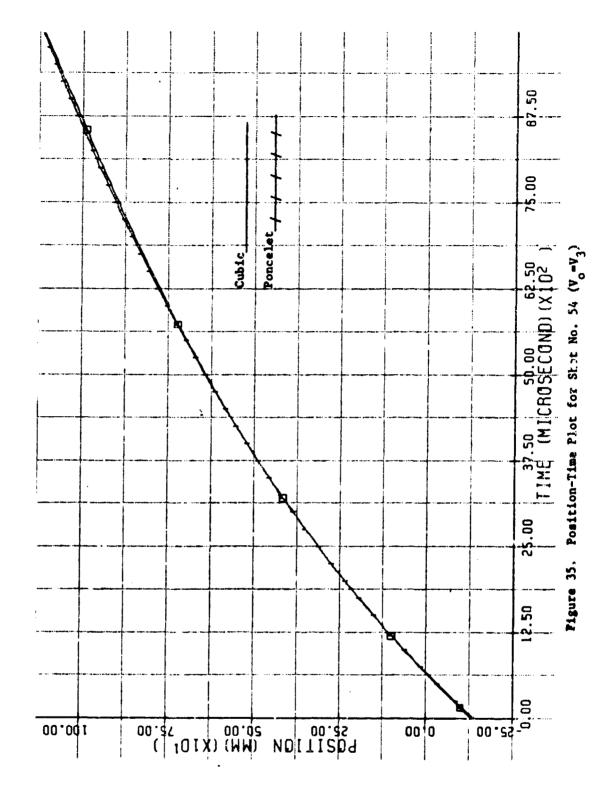
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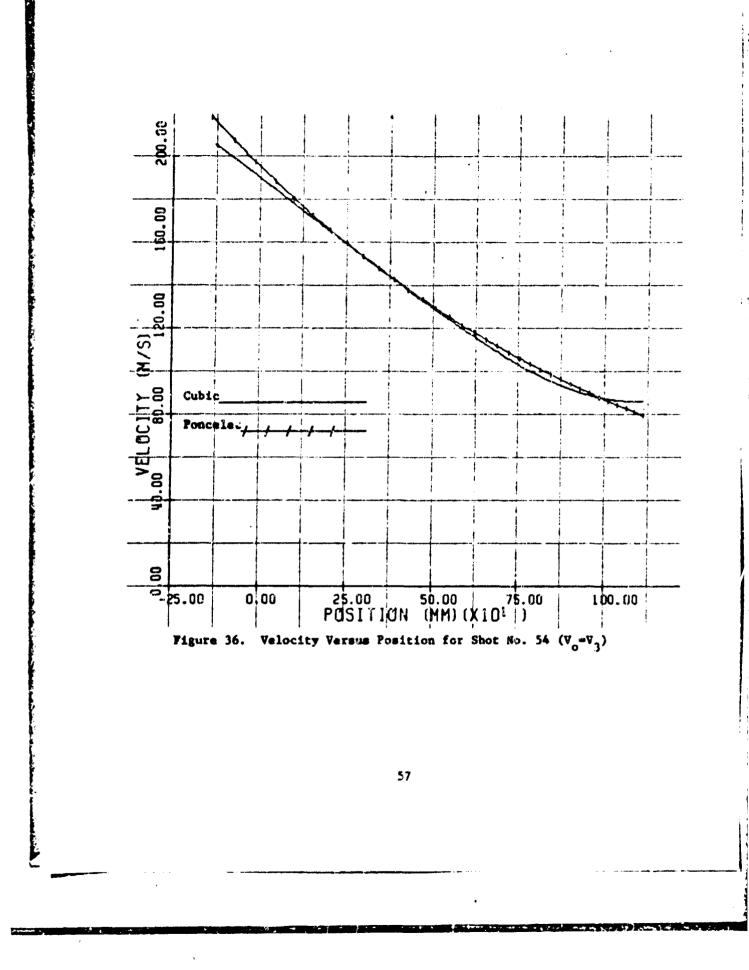


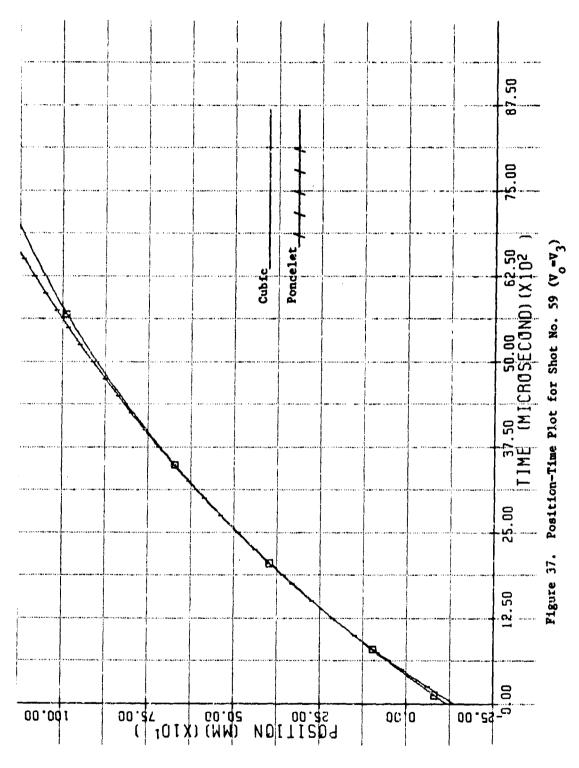


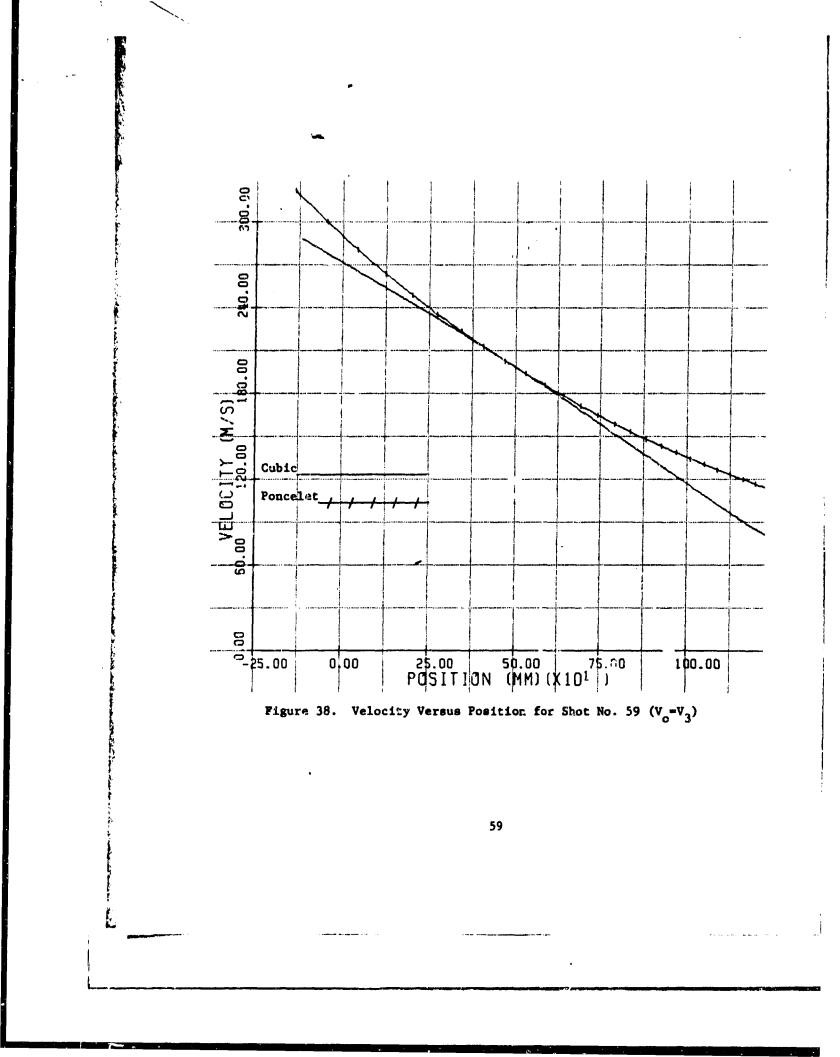


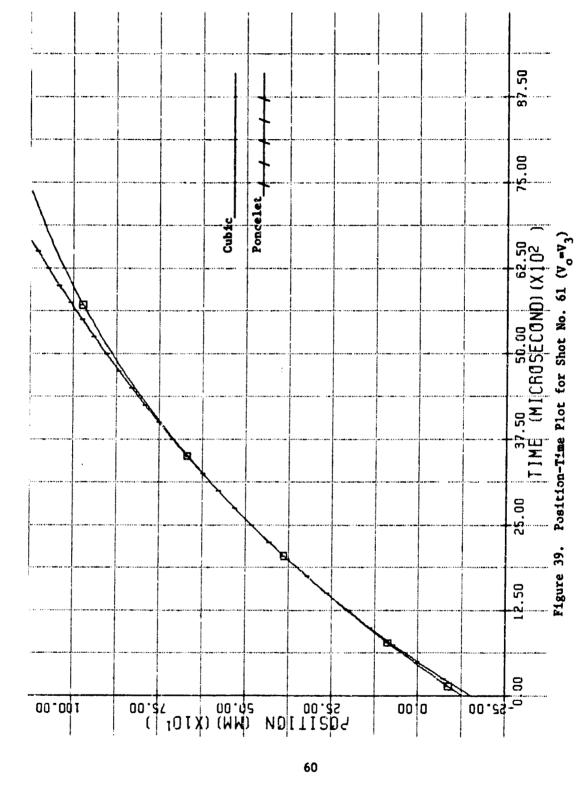


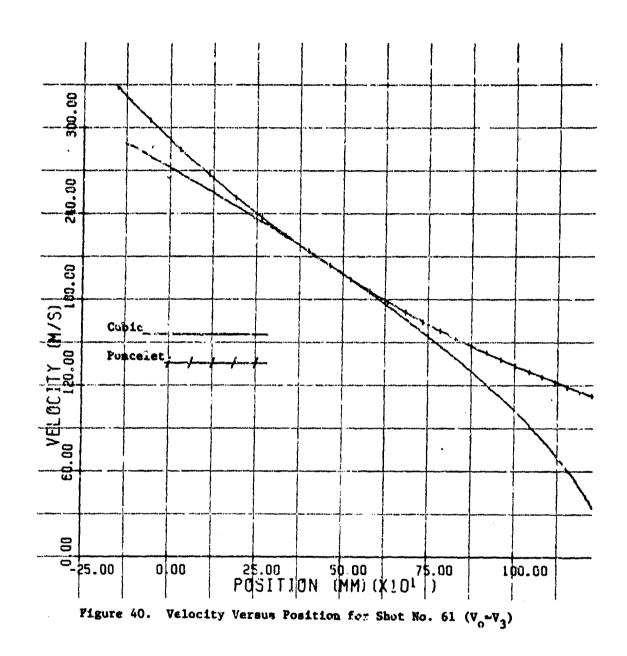


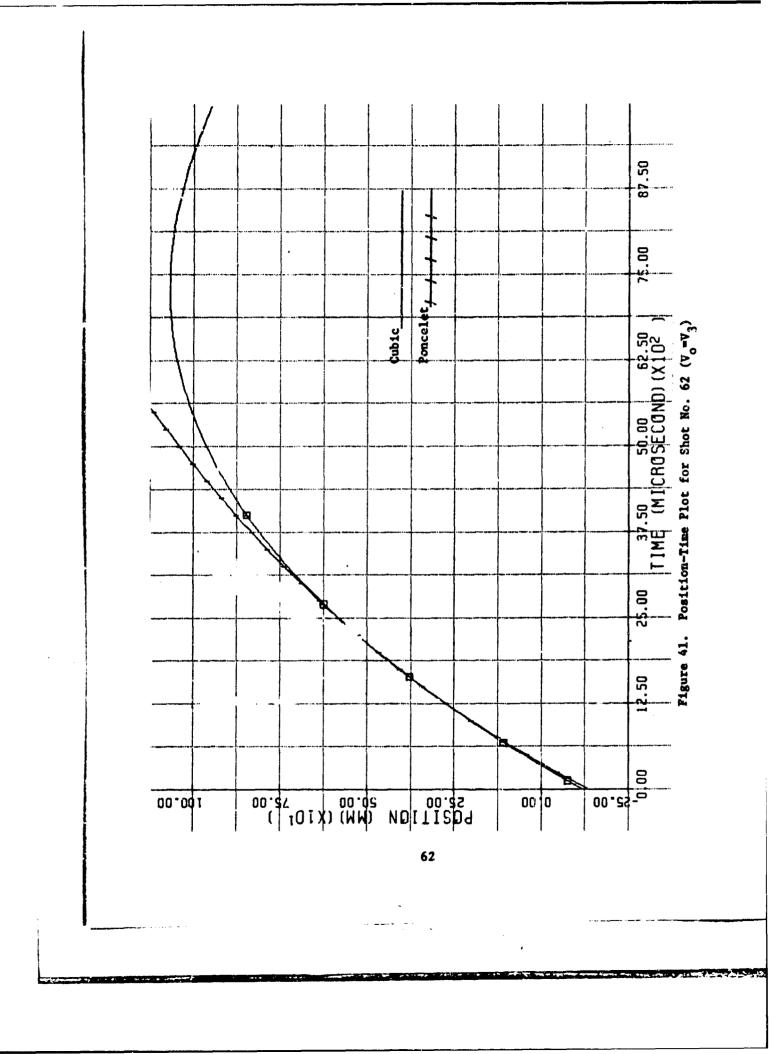


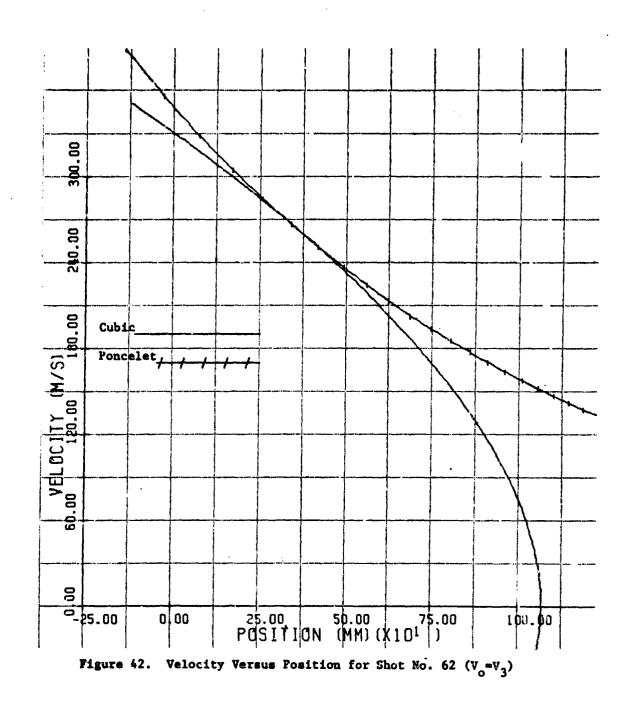


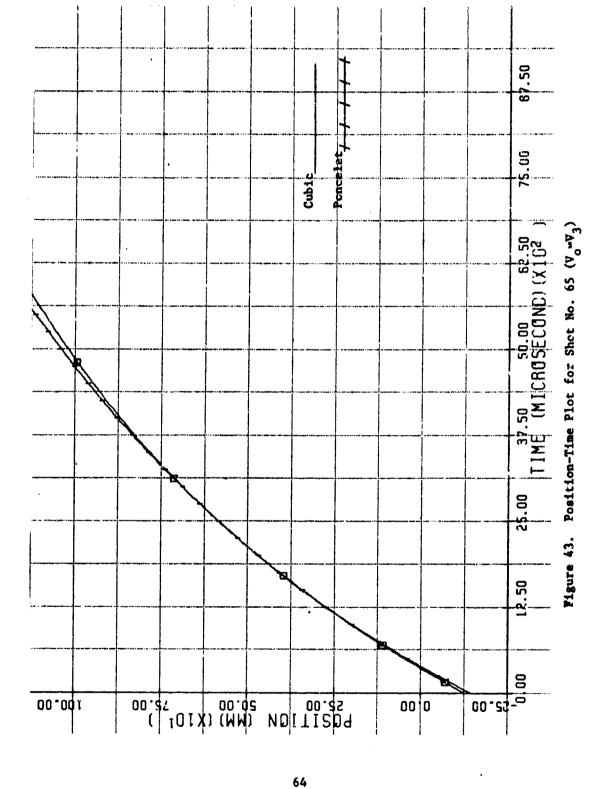




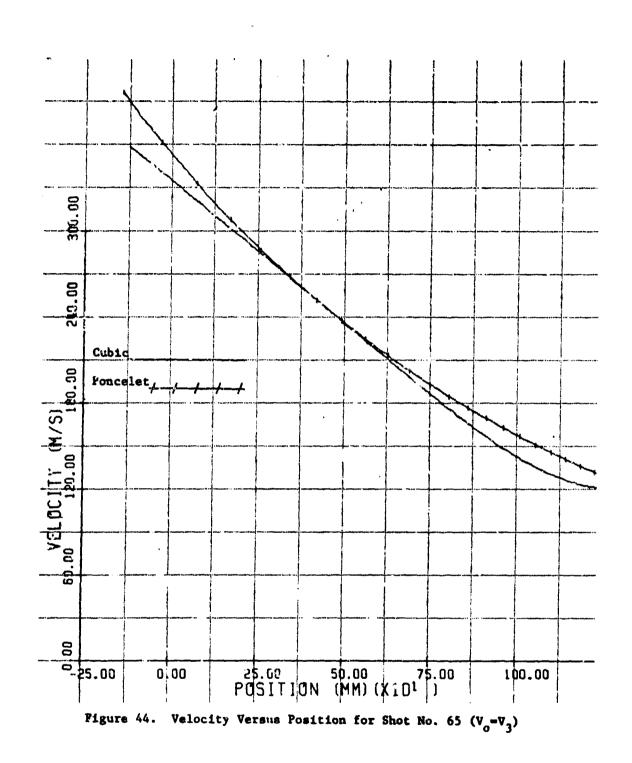


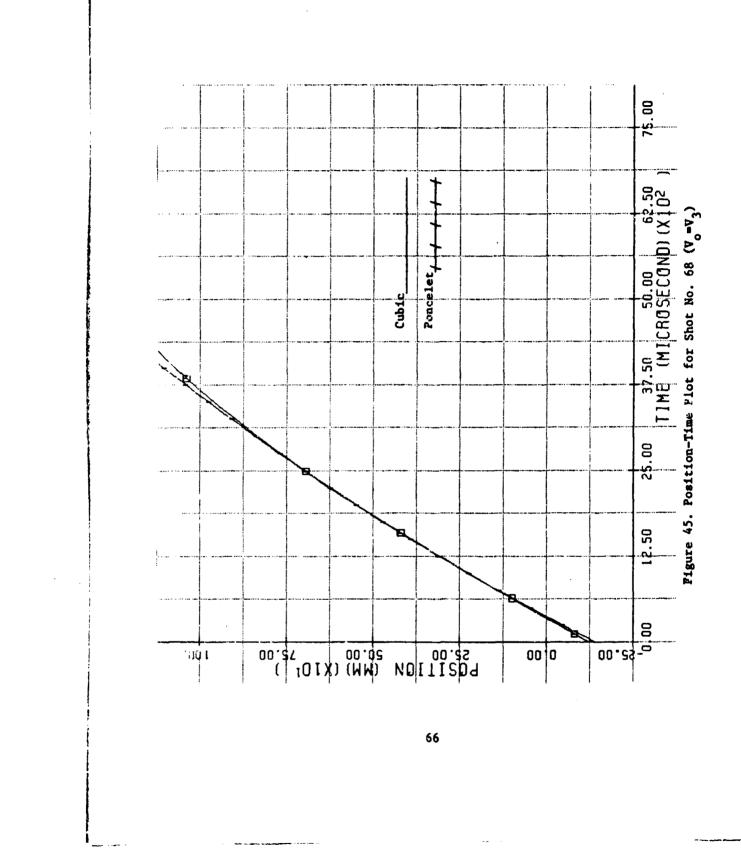




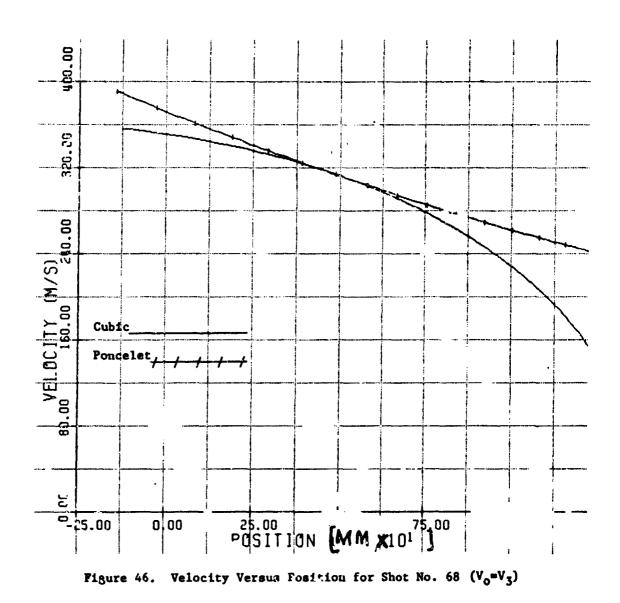


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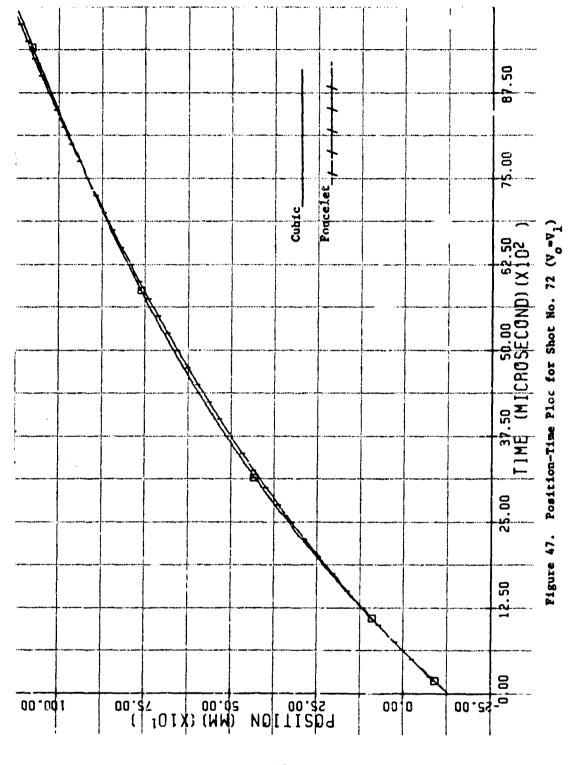


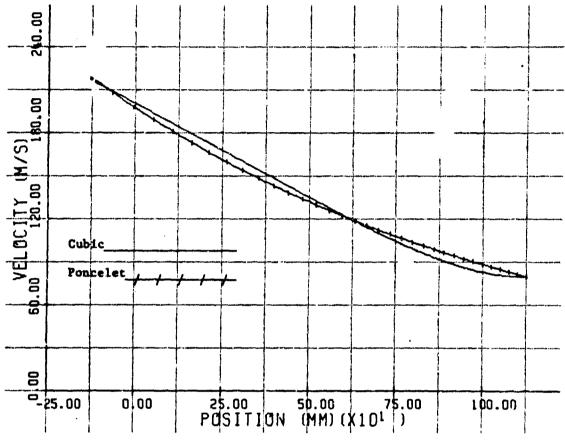
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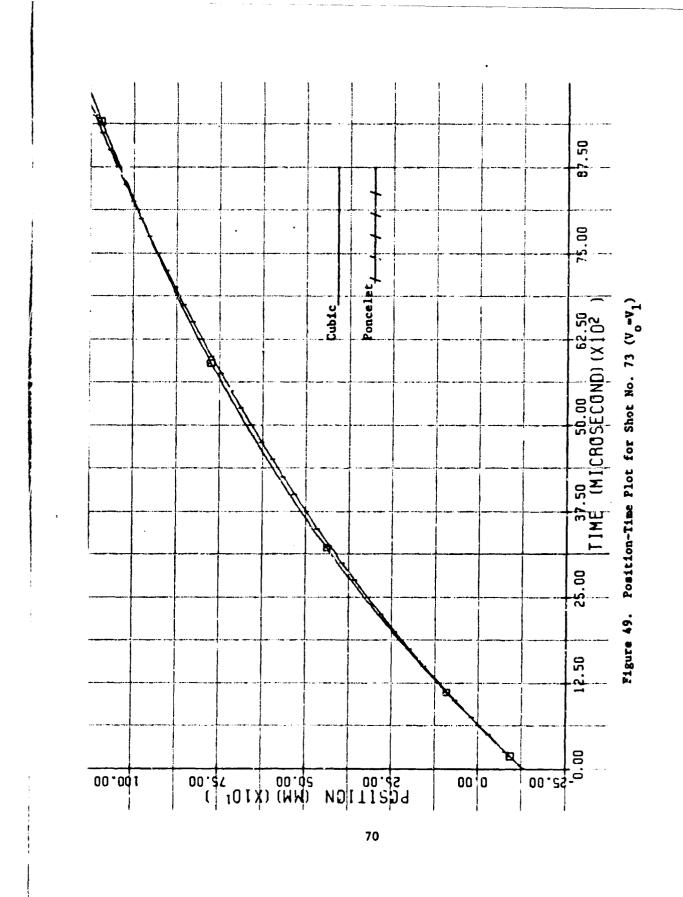


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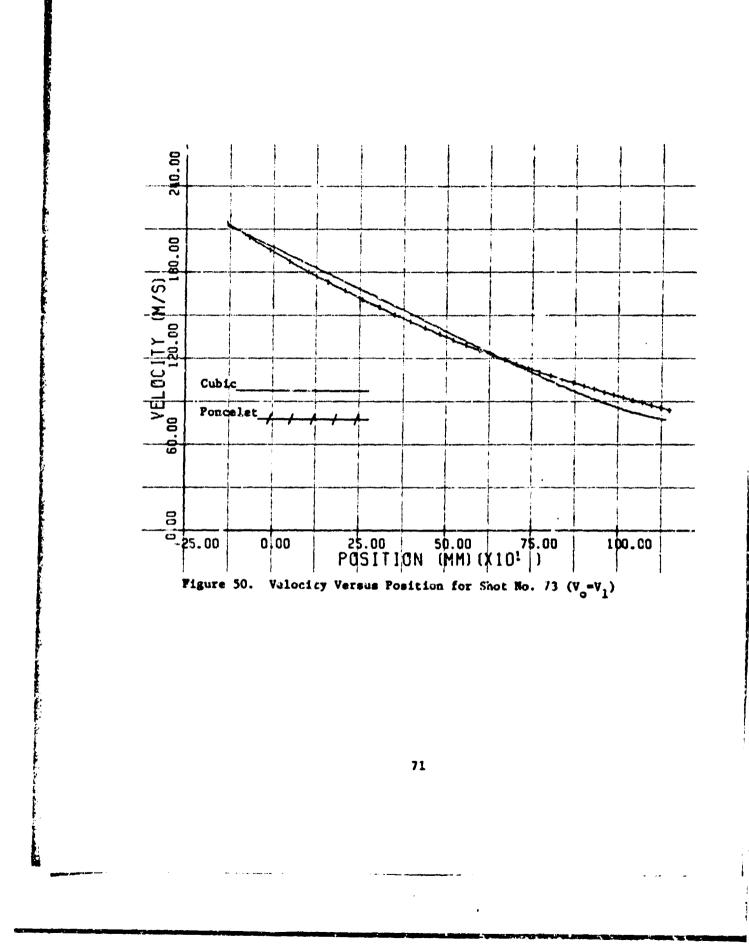


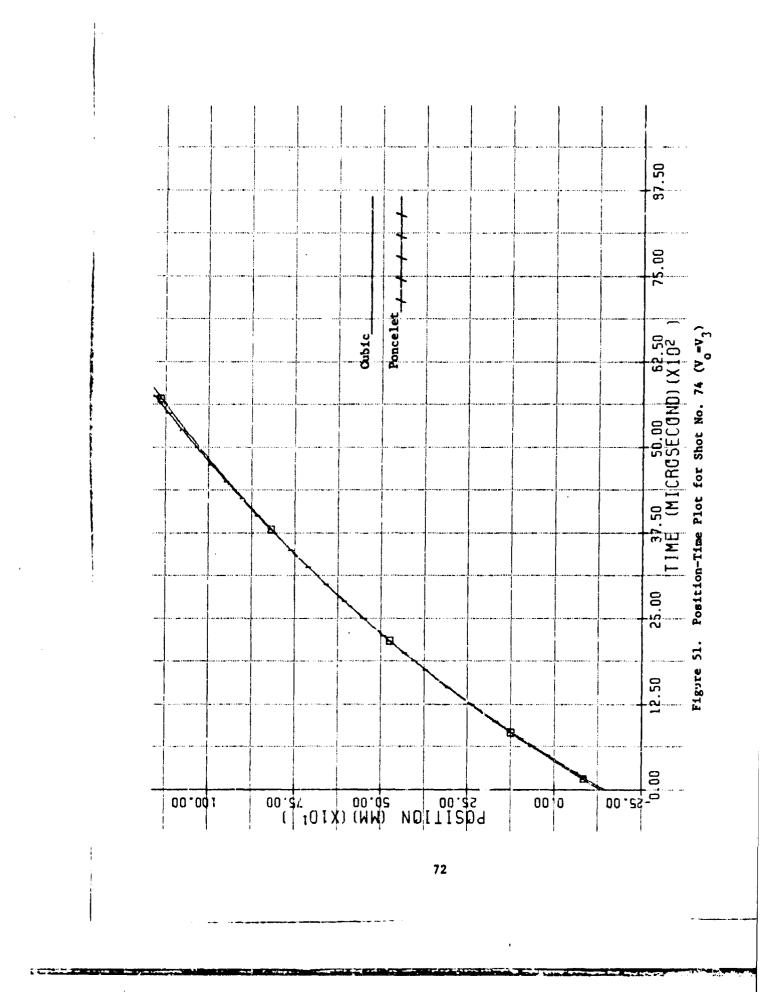


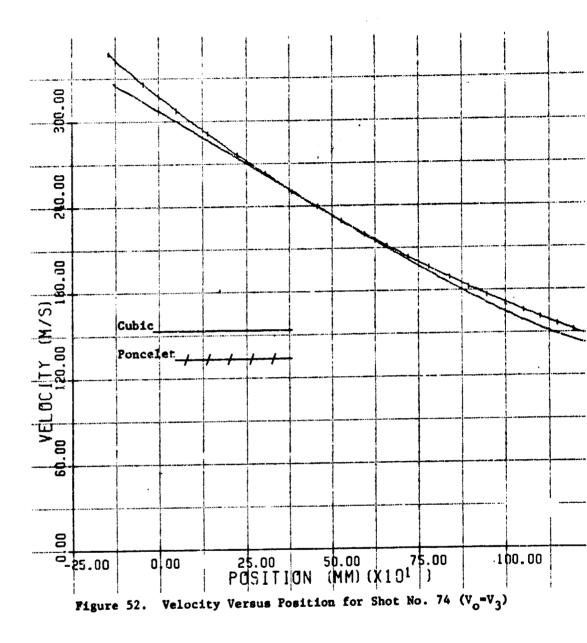


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interpolation results of paragraph 4.3.1. A different procedure will be used in paragraphs 4.5 and 4.6. Measured approach velocities were available but there was some uncertainty about the exact time the nose of the projectile first impacted the target. Positions at the times of the X-ray firings were more precisely known. The majority of the aralyses (33 shots) were therefore made with Station 3 as initial-value point, with  $V_0 = V_3$  as given by the cubic interpolation, since that point is near the middle of the interval where differentiation of the cubic interpolation formula should be most accurate. Calculated velocities at Stations 1 and 2 were then obtained for negative values of t-t, and x-x. For 19 shots that procedure led to unreasonably large calculated values for V at x = -0.115 meter where V should be equal to the approach velocity. For these cases the analysis was therefore made with Station 1 selected as initial point.

In the tabulations of Appendix A the value of  $V_0$  is tabulated as VO on the line below the tabulated Poncelet drag coefficient. For example, for Shot 26 the value VO = 267 is listed. Comparision with C.G.Vel.X-Comp. values four lines above this entry shows that in this case  $V_0 = V_3$ . For the 21 cases plotted in paragraph 4.3.1, the captions include either the statement  $V_0 = V_1$  or  $V_0 = V_3$ , and it is also easy to tell from the V, x-plots where they were made to agree.

All the fitted Poncelet curves (those marked by vertical strokes) now give a reasonable agreement at x = -11.5 cm with the measured approach velocity.

The standard deviation of the x-positions calculated by Equation (6) from the experimental values is tabulated for each case in Appendix A immediately following the tabulated  $V_{\alpha}$ .

The Poncelet drag coefficient tabulated on the same line is related to the coefficient B of Equation (1) as follows. In aerodynamics a dimensionless drag coefficient  $C_D$  is defined such that the drag force on an object of projected area  $A_1$  on a plane perpendicular to the velocity is given by

Inertial Drag Force = 
$$\rho A_1 c_D V^2 / 2$$
 (7)

where  $\rho$  is the density of the medium being traversed. Comparison with Equation (3), neglecting A, shows

$$B = \rho A_1 C_0 / 2m \quad \text{or } C_0 = 2m B / \rho A_1 \tag{8}$$

The projectile mass in kilograms and the target sand density in  $kg/m^3$  are tabulated for each shot in Appendix A.

Table 7 lists the Poncelet drag coefficients  $C_D$  for each of the 41 shots of the primary test program whose velocities were analyzed. In dry sand there is little variation with impact velocity of the value of  $C_D$ required to fit the X-ray position time data in the region observed. These dry sand results are well characterized by the Poncelet force law with a

Projectile Type		It	mpact Vel	ocity Range	•	
and Sand Condition		m/sec		m/sec		w/sec
	Shot	c <sub>D</sub>	Shot	с <sub>р</sub>	Shot	° CD
Solid	17	1.64	20	1.77	25	1.68
	18	1.62	22	2.30 <sup>4</sup>	26	1.65
Fiat Nose	19	1.69	23	1.72	27	1.77
Dry Sand			24	1.72	29	1.71
	Avg	1.65	Avg	1.74	Avg	1.70
Solid	70	1.59	37	0.95	76	0.96
Flat Nose	71	1.55	38	0.94	82	0.73
FIGE NUBE	72	1.36	74	1.02	83	0.82
Wet Sand	73	1.23	81	0.96	84	0.78
	Avg	1.43	Avg	<u>U.97</u>	Avg	0.82
Solid	52	1.81	56	1.78	62	1.83
Step-Tier	54	1.90	58	1.91	63	1.81
	55	1.89	59	1.67	64	1.92
Dry Sand	57	1.62	61	1.82	65	0.54 #
	Avg	1.80	Avg	1.80	Avg	1.85
Solid	45	1.24	39	0.72	50	0.79
Step-Tier			49	0.82	68	0.61
- 1				1	69	0.71
Wet Sand	Avg	1.24	Avg	0.77	Avg	0.70
	(Valu	es marked	with e	xcluded fro	on avera	2e)

# TABLE 7 . PONCELET/DRAG COEFFICIENTS CALCULATED FOR PRIMARY TEST PROGRAM

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value of  $C_D = 1.7$  for the flat-nose projectile and  $C_D = 1.8$  for the steptier projectile. In paragraph 4.6, a different method of determining  $C_D$  is presented, however, which shows up variations in  $C_D$  along a trajectory.

In the wet sand (fully saturated) there is a downward trend of  $C_D$ as the impact velocity increases, indicating that the penetration is not well characterized by a Poncelet force law with coefficients independent of velocity. There is also more scatter in the  $C_D$  values obtained for some of the velocity regimes in the wet sand.

The saturated sand  $C_D$  values are all lower than any of the dry sand values (except for Shot 65, which is believed to be in error). Shot No. 87 into a water target at 241 m/sec was fitted by  $C_D = 0.51$ . Apparently the fully saturated sand tends to respond somewhat like a water target.

Magnetic sensor response is compared with X-ray data in the following paragraph.

4.4 COMPARISON OF MAGNETIC SENSOR RESULTS WITH X-RAY DATA

For the 26 shots marked with an M in Tables 3 and 4 the cubic interpolation formula described in paragraph 4.3.1 was used to compute the nose position of the projectile at the time of maximum and minimum coil voltage from each magnetic sensing coil. These are compared to the coil position in the last data group tabulated from each of these shots in Appendix A. The differences between these two positions (tabulated in two last lines) give the apparent distance back from the nose to where the maximum outward radial component of the projectile's magnetic field cut the sensing coil (or the maximum inward component for the minimum voltage case).

In Shot 26, for example, the distances were 0.024, 0.023, 0.035, 0.042 back to the apparent maxima. The first two of these agree quite closely with the laboratory mapping of the magnetic field of the projectile reported in paragraph 2.3, which indicated that the maximum should occur about 0.02 meter back from the nose of the flat-nosed projectile if the magnetized projectile passed through the center of the coil. There was some discrepancy at the last two stations, but an error of 0.02 meter is quite small compared to the position coordinate of 1.076 meters at the last station. Shot 26 had one of the straightest trajectories.

Larger discrepancies are recorded for several shots. The largest positive difference noted is for Station 2 of Shot 50 where the apparent maximum was 0.063 meter back of the nose as compared with the laboratory measurement of 0.022 meter for a solid step-tier projectile. Since this occurred at Station 2 (position 0.486 meter ) it would represent about an 8 percent error in position indication.

For several stations a negative difference was noted. The largest was -0.027 meter for Station 3 in Shot 58, also a solid step-tier projectile, indicating an apparent maximum 0.027 meter ahead of the nose or 0.049 meter ahead of the maximum position according to the laboratory measurement of the field, leading to a 6 percent position error at Station 3.

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In the majority of the cases the apparent errors were smaller than these. No laboratory measurement was made for the hollow step-tier projectiles, but differences between computed nose positions and coil positions for them in Shots 79, 86, and 87 were in the range of the differences for the solid projectiles.

A possible source of the errors could be off-center projectile paths, although both Shots 50 and 58 had quite straight trajectories according to the X-ray records. Other possible sources of error are imprecise measurement of coil positions and/or maximum time on the strip chart and possibly distortions of the magnetic field of the projectile caused by the impact.

It is believed that the X-ray data are generally more precise than the magnetic sensor data. Nevertheless the investigation has indicated that the magnetic sensing method can give quite good results, and it is certainly an economical method.

A different procedure for determining the Poncelet drag coefficients will be presented in paragraph 4.5, which also considers application of the empirical penetration formula developed at Sandia Laboratories (References 12,13).

#### 4.5 COMPARISON OF PONCELET AND SANDIA EMPIRICAL FORMULA RESULTS

The Sandia empirical formulas (References 12,13) are rewritten here in SI units. Thus, according to Young (Reference 13) the total depth of penetration D is given in terms of the initial impact velocity  $V_0$  by an equation of the form

$$D = 0.0117 \text{ KSN}(W/A_1)^{1/2}(V_0 - 31.5) \text{ for } V_0 > 61 \text{ m/sec}$$
(9)

or by

$$D = 2KSN(W/A_{1})2n[1 + 2V_{0}^{2}(10^{-4})] \text{ for } V_{0} < 61 \text{ m/sec}$$
(10)

W is projectile weight A<sub>1</sub> is cross sectional area

Since all impacts in the present study had  $V_0 > 61$  m/sec, a procedure based on a method used by Young (Reference 13) to modify Equation (9) for use with layered media was used to analyze the Eglin experiments. In Equation (9), N is a nose coefficient, S is a soil coefficient, and K is an independently determined parameter. Since K, S, and N appear only as the product KSN, the procedure followed was to determine the best value of KSN to fit the experimental velocity versus position data. The projectiles for the present experiments were considerably smaller and lighter than those for which the empirical formulations and scaling laws had been shown to give good results.

According to Young (Reference 13), for  $V_o > 61$  m/sec the constant A in the Poncelet Equation (3) of paragraph 4.3.2 is negligible, so that

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Equation (5) is the appropriate form of to position, namely (with  $x_0 = 0$ )

mediat inw to relate velocity

$$\mathbf{v} = \mathbf{v}_{o} \mathbf{e}^{-\mathbf{B}^{*}}$$
(11)

In the procedure of this section both the Poncelet parameter B and the initial velocity  $V_0$  at x = 0 are considered as unknowns to be determined from the experimental data of Appendix A as follows. The velocity  $V_1$  at four positions  $x_1$ , each located midway between the positions of the projectile in two successive X-ray pictures, is calculated as  $V_1 = \Delta x_1 / \Delta t_1$ , where  $\Delta x_1$  is the distance traveled during the time between the firings of the two x-rays. Then the regression procedure gives

$$B = [(\Sigma x_{i})(\frac{1}{4}\Sigma \ell n V_{i}) - \Sigma x_{i} \ell n V_{i}] / [\Sigma x_{i}^{2} - \frac{1}{4}(\Sigma x_{i})(\Sigma x_{i})]$$
(12)

and

$$V_{o} = Exp \left[ \Sigma \ln V_{i} + \frac{1}{4} B(\Sigma x_{i}) \right]$$
(13)

where 1 = 1 to 4.

The Poncelet drag coefficient  $C_D$  is then given by  $C_D = 2mB/\rho A_1$  as in Equation (8).

The basis for the application of the Sandia empirical formula to layered media (Reference 13) was the assumption of constant deceler "ion through each layer. Thus if a is the acceleration magnitude nondimusionalized with respect to g, the acceleration of gravity (so that a g is the dimensional deceleration), then the velocities  $V_{n+1}$  at the beginning and the end of a layer were related (Reference 13) by

$$V_{n+1}^{2} = V_{n}^{2} - 2g[(a_{n-1} + a_{n})\frac{L}{2} + a_{n}(t_{n} - L)]$$
(14)

where  $t_n$  is the thickness of the layer and it is assumed that  $t_n >> L$ . For the flat-nosed projectile, with L =0, Equation (14) reduces to

$$v_{n+1}^2 = v_n^2 - 2a_n gt_n$$
 (15)

which is the basis for the following regression procedure to determine the Sandia parameter KSN. Let

$$G_{i} = v_{1}^{2} / [0.0117(v_{1} - 30.5)(v_{1}^{2} - v_{1}^{2}) (W/A)^{1/2}]$$
(16)

Then

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$$KSN = \frac{\Sigma(x_{1} - x_{1})(\frac{1}{3}\Sigma G_{1}) - \Sigma(x_{1} - x_{1})(G_{1})}{\Sigma(x_{1} - x_{1})^{2} - \frac{1}{3}[\Sigma(x_{1} - x_{1})][\Sigma(x_{1} - x_{1})]}$$
(17)  
= 2 to 4

The different regression analyses presented in this paragraph and in paragraph 4.3.2 gave for the classical Poncelet equations in most cases almost

the same results for  $C_D$ . The results for  $C_D$  as a function of initial velocity for the solid flat-nosed projectiles in dry sand and in wet sand are summarized in Figure 53. When the two regression results differed significantly the result of paragraph 4.3.2 was used. As shown by the solid triangles the value of  $C_D$  is almost independent of the initial velocity in dry sand, while the values for saturated sand (marked by the solid circles) are lower than for dry sand and show a downward trend with increasing initial velocity.

A similar plot of the fitted value of the Sandia parameter KSN for the same shots is shown in Figure 54. The greater scatter in the fitted values of KSN than in those of  $C_D$  indicated that these penetration events are not well characterized by a single value of KSN for each shot. The greater discrepancies with the Sandia equation can be explained in part by the assumption of a constant deceleration pegnitude in each segment, in contrast with the Poncelet prediction which does fit the dry sand experimental data very well.

In paragraph 4.6 the pollible variation of the Poncelet coefficient with position along a tribulary is examined by fitting separace values for different portions the trajectory for each shot. Some of the values  $\operatorname{clip}$  and KSN for the shots analyzed by the regression methods described in this section will also be given in Table 8 for dry sand and Table 9 for wet sand.

#### 4.6 DRAG COEFFICIENT VARIATION WITH POSITION ALONG A TRAJECTORY

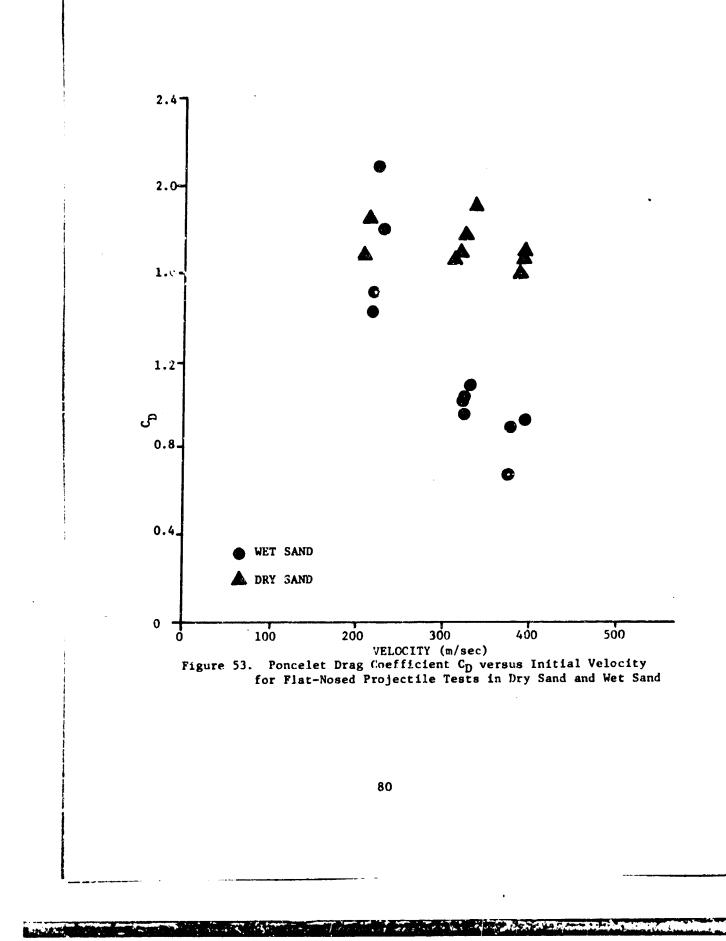
The favorable agreement with the experimental data of the positiontime and velocity-position curves calculated by the Poncelet force law suggested further analysis by this model. The variation in the drag coefficient along a trajectory was examined for some of the solid flat-nose projectile trajectories, for 9 shots in dry sand and 10 in wet sand, including examples from each of the three impact velocity regimes of the primary test program. These calculations were made by separately evaluating the Poncelet parameter B for three different segments of each trajectory by the following equation.

$$B = \rho C_{D} A_{1} / 2m = [\ell_{n} (V_{n} / V_{n+1})] / (x_{n+1} - x_{n})$$
(18)

where subscripts n and n + 1 identify values at the beginning and the end of a segment.

Tables 8 and 9, for dry sand and wet sand, respectively, exhibit in the third column the resulting drag coefficients calculated by Equation (18) for three segments of the trajectory for each shot. The tables also list in the last two columns some of the values of  $C_D$  and the Sandia parameter KSN fitted to the whole trajectory by the methods of paragraphs 4.3.2 and 4.5. When two values of  $C_D$  were listed for a shot, the first was calculated by the method of paragraph 4.3.2 and the second by the method of paragraph 4.5.

As was pointed out at the end of paragraph 4.3 the  $C_D$  values fitted to the dry sand shots are consistently higher than those for saturated sand. Moreover the wet sand values show a downward trend with increasing impact



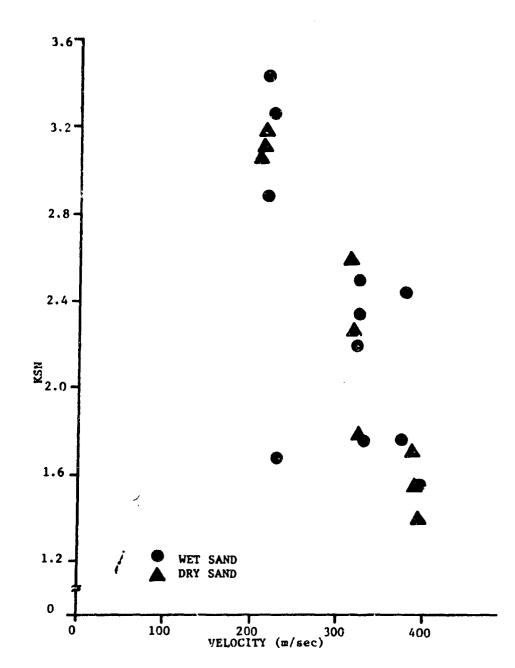


Figure 54. Sandia Penetration Coefficient KSN versus Initial Velocity for Flat-Nosed Projectile Tests in Dry Sand and Wet Sand

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TABLE 8.	DRAG COEFFICIENTS C	AND KSN	VALUES F	FOR SELECTED	SHOTS I	N DRY SAND
	(SOL	ID FLAT-	NOSE PROJ	ECTILES)		

Shot	Striking	Seguent	Distance	Average Velocity	Shot	Shot
No.	Velocity (m/sec)	C <sub>D</sub>	(m)	(m/sec)	C <sub>D</sub> *	KSN
17	212.1	1.779	0.320	174.1 139.3	1.643	3.05
		2.049	0.901	108.6	1.1/3	
18	213.4	1.887	0.307	174.6	1.624	3.20
		1.673	0.570	139.3	2.230	1
		2.413	0.826	106.5		
19	210.6	1.849	0.311	176.5	1.692	3.17
		1.880	0.573	139.6	1.973	
		2.149	0.880	106.5		
20	329.2	1.678	0.295	268.1	1.769	2.26
		1.910	0.552	214.8	1.786	
		1.687	0.862	168.3	<u></u>	
23	328.2	2.014	0.300	269.5	1.725	1.77
		1.578	0.562	216.0	1.867	1
		2.152	0.863	168.7		
24	327.4	1.643	0.269	265.9	1.723	2.59
		1.774	0.525	217.8	1.765	1
		1.814	0.826	172.0		
25	406.0	1.842	0.277	329.1	1.685	1.38
		1.400	0.534	264.2	1.561	[
		1.615	0.832	211.3		
26	406.3	1.783	0.280	329.6	1.649	1.70
		1.760	0.538	266.5	1.670	
		1.515	0.810	214.0		
29	405.0	1.933	0.281	330.8	1.706	1.54
		1.761	0.522	266.3	1.735	
		1.625	0.839	214.0		

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# TABLE 9. DRAG COEFFICIENTS CD AND KSN VALUES FOR SELECTED SHOTS IN WET SAND (SOLID FLAT-NOSE PROJECTILES)

Shot No.	Striking Velocity (m/sec)	Segment CD	Distance (m)	Average Velocity (m/sec)	Shot C <sub>D</sub> *	Shot KSN
71	207.9	2.167	0.268	182.3	1.545	1.614
		1.560	(	134.2	1.802	1
		1.921	0.826	98.8		1
72	214.0	1.302	0.285	181.1	1.359	3.40
	2.4.0	1.561	0.541	142.3	1.564	1 3.40
		1.673	0.865	104.9	1	
73	212.8	1.410	0.287	180.9	1.233	2.92
1.5	212.0	1.288	0.549	143.6	1.474	2.52
		1.697	0.878	108.3	1.4/4	
37	336.5	1.134	0.298	279.6	0.947	2.214
57	550.5	0.824	0.559	237.0	1.023	2.214
		1.218	0.883	196.6	1.025	
38	333.1	1.032	0.300	284.0	0.937	2.36
		0.924	0.587	240.2	0.969	
		0.997	0.917	200.1		ļ
74	334.0	1.240	0.309	282.7	1.016	
		0.923	0.580	233.4		1
		1.230	0.917	189.7		
81	333.8	1.023	0.309	278.7	0.963	
		0.885	0.570	236.3		Į –
		1.308	0.896	193.6		ļ
82	404.8	0.985	0.308	343.2	0.728	
		0.645	0.562	299.3		ſ
		0.499	0.889	269.2		l
83	419.4	1.228	0.307	347.0	0.819	[
		0.697	0.547	298.5		[
		1.148	0.854	258.7		ļ
34	405.7	0.978	0.308	333.3	0.777	
	1	0.584	0.545	294.0		
	1	1.322	0.843	253.0		1

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velocity, while for dry sand the values were essentially independent of striking velocity.

Of particular interest here is the change in  $C_D$  along a trajectory as shown by the two or three different values listed in the third column for each shot. In many cases the first segment  $C_D$  is higher than the second and then the trend is reversed to give a third segment  $C_D$  higher than the second. This pattern is followed in 8 of the 10 wet sand cases and in 4 of the 9 dry sand cases, for which three segments were calculated.

The dry sand coefficients show less variation along the trajectory than the wet sand coefficients, a variation of the order of 30 percent between the maximum and minimum values in dry sand and two or three times this much variation in wet sand.

This variation might imply that  $C_D$  is velocity dependent instead of being a constant. It seems that in some cases the drag coefficient variation can be fitted to a power law in the velocity. An example of this is given in Section VI.

The apparent variation in the coefficient may, however, actually be a result of assuming an incorrect form for the force law. If the force law contains a term linearly dependent on velocity in addition to the term depending on the square of the velocity, both with constant coefficients, this would lead to an apparent variation in the  $C_D$  determined by Equation (18), which is based on a law where the force is proportional to the square of the velocity.

The high drag at the beginning of the trajectory may also be related to shock effects involving not only the velocity but also the appulse duration and the acoustic impedance of the target medium. The minimum  $C_D$ appears to be related to momentum transfer to the target medium, where the predominant drag is proportional to the square of the velocity. Finally, as the projectile slows and the cavity collapses onto it, friction and shear resistance in the target medium become important, giving rise to an increase in the apparent drag coefficient.

#### SECTION V

#### CAVITY EXPANSION THEORY PENETRATION CALCULATION

In 1975 Bernard and Hanagud (Reference 6) published a report e-cending the approximate penetration calculation method for projectiles with a hemispherical nose, based on the theory of expansion of a spherical cavity (CET), to projectiles with conical and ogival noses and showing how it could be extended to an arbitrary axially symmetric projectile. The first  $n_{\rm e}=$  of CET methods for dynamic penetration was by Goodter (Reference 26) for a spherical projectile impacting an incompressible strain hardening target. Hanagud and Ross (Reference 3) modified the method to account approximately for target compressibility by treating the target material as a locking medium. The method has been applied to penetration calculations for flatnose projectiles by Rohani (Reference 4), with the implicit assumption that a false hemispherical nose of target material is formed and carried by the projectile along a stable straight path.

It should be remembéred that sev il quite important assumptions are made in applying the cavity expansion method to penetration calculations, so that extensive experimental verification is necessary to check on the range of approximate validity of predictions by the method. Nevertheless it has achieved some remarkable success in predicting penetration depths from measured soil properties (Reference: 4,6,7). Bernard and Hanagud (Reference 6) defined a dimensionless parameter  $R_g$ , which they call the solid Reynolds number

$$R_{g} = \frac{\rho V^{2}}{Y}$$
(19)

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where  $\rho$  is target density, Y is target yield strength in a unimial strain test, and V is projectile velocity. It was concluded that final penetration depth was reasonably well predicted for  $R_g$  between zero and about 100. They considered the upper bound of 100 as a conservative one, since results of experiments at high values of  $R_g$  are needed in order to establish a more realistic range. They remarked that accurate prediction of details of the complete deceleration history might demand a much stronger limitation on  $R_g$ .

According to the spherical cavity expansion theory for an infinite locking compressible medium the compressive normal stress p at the cavity surface is

$$p = p_{a} + p_{T} = p_{a} + \rho_{p} (B_{1}ab + B_{2}a^{2})$$
 (20)

where p and  $p_1$  are the separate contributions of the material deformation (shear) and inertia, which Bernard and Hanagud (Reference 6) call the shear resistance and the dynamic pressure, respectively. In this equation  $\rho_p$  is the locked plastic density in the region behind the expanding spherical plastic locking shock wave, a is the instantaneous cavity radius, Å, and ä are the radial velocity and acceleration of the cavity surface and  $p_8$ ,  $B_1$ , and  $B_2$  are parameters related to properties of the material. The way these parameters are calculated will be indicated later in this section.

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In applying this theory to penetration by a projectile with a hemispherical nose two very important assumptions are made (References 3,4) in order to get a simple theory.

(1). The parts  $p_g$  and  $p_I$  of the normal pressure at the tip of a hemispherical nose of radius a on a projectile traveling at speed V and acceleration  $\hat{V}$  are assumed to be equal to the values of  $p_g$  and  $p_I$ on a spherical cavity surface of the same instantaneous radius a, but expanding with  $\dot{a} = V$  and  $\ddot{a} = \dot{V}$ .

(2). The entire hemispherical nose is assumed to be in contact with the target material, and the dynamic pressure on the projectile's hemispherical nose is assumed to vary from the stagnation point value at the nose tip to zero at the shoulder as the cosine of the polar angle measured from tip (colatitude), while  $p_s$  is uniform over the nose.

If friction on the nose and all afterbody forces are neglected this leads in a straightforward manner to the following equation of motion for the projectile of mass M

$$(M + \frac{2}{3}\pi a^{3}\rho_{p}B_{1})\frac{dV}{dt} = -\pi a^{2}(\rho_{s} + \frac{2}{3}\rho_{p}B_{2}V^{2})$$
(21)

The term  $\frac{2}{3\pi a} {}^3\rho_{\rm p} B_1$  is an added mass term resulting from the acceleration term  $\rho_{\rm B_1}$  at in the expression for the dynamic pressure in Equation (20). In most cases that have been treated the added-mass term was negligible in comparison to the projectile mass M.

In modifying the method to other exisymmetric nose shapes, Bernard and Hanagud replaced the assumption on the variation of  $p_T$  over the nose surface by an assumption on the variation along the nose of the tangential component  $V_t$  of the target absolute velocity. For a fully embedded nose, the material was assumed to be in contact all along the nose surface (no separation before the base of the nose), an assumption which they recognized was not generally strictly correct. The normal component  $V_n$  of target material velocity was therefore required to be equal to the normal component of the velocity of the projectile nose surface. For a conical nose of half apex angle  $\phi$ , it was assumed that the tangential component of target velocity varies from V cos  $\phi$  at the nose tip to zero at the base of the cone, according to the law

$$V_{t} = [1 - (z/L)^{2}]^{1/2} V \cos \phi \qquad 0 \le z \le L \qquad (22)$$

where L is the nose length and z is the axial coordinate measured back from the tip of the nose. The velocity dependent part of the dynamic pressure  $p_{\rm I}$  was then assumed to depend on the local resultant particle velocity magnitude  $V_{\rm p} = [V_{\rm n}^2 + V_{\rm t}^2]^{1/2}$  in the same way that  $p_{\rm I}$  depends on a in the spherical cavity expansion theory. Similar assumptions could be made about the dependence on the local particle acceleration on the nose surface, but because that term is usually much smaller than the velocity-dependent term, Bernard and Hanagud (Reference 6) chose to use the nose tip acceleration  $\dot{V}$ , so that  $p_{\rm T}$  is given by

$$p_{I} = \rho_{V} B_{1} a \dot{V} + \rho_{p} B_{2} V_{p}^{2}$$
(23)

where  $V_p$  varies over the nose while  $\dot{V}$  does not. Integration of the axial force component over the nose then gives the following equation of motion, replacing Equation (21)

$$(M + \pi a^{3} \rho_{p} B_{1}) \frac{dV}{dt} = -\pi a^{2} (p_{s} + \rho_{p} B_{2} f_{n} V^{2})$$
(24)

where the dimensionless nose-shape factor  $f_n$  is given for a conical nose of length-to-diameter ratio L/D by

$$f_{n} = 1 - \frac{2}{3} \frac{4(I/D)^{2}}{4(I/D)^{2} + 1}$$
(25)

Besides containing the factor  $f_n$  in place of the factor 3/2, Equation (24) differs from Equation (21) by lacking the factor 3/2 in the term containing  $S_1$ .

Bernard and Hanagud (Reference 6) observed that with the assumption listed above the variation in  $f_n$  as L/D varies from zero (flat mose) to infinity (long pointed nose) produces a variation in predicted final penetration depth by a factor of three. For L/D = 0.5,  $f_n$  is 3/2 and the prediction is the same as for the hemispherical nose, when the contribution of the added-mass term containing  $B_1$  is negligible.

For other fully embedded convex axisymmetrical nose shapes, Bernard and Hanagud (Reference 6) gave a method for estimating  $V_n$  and  $V_t$  at any position on the nose by considering the circumscribed cone tangent tr the nose at the point. The base of the cone was in the same plane as the actual nose base, but the tip was forward of the actual tip. The components  $V_n$ and  $V_t$  at the point of tangency were assumed to be equal to the values of  $V_n$ and  $V_t$  that would be assumed on the circumscribed cone at the point of tangency if it were an actual conical nose, calculated by the procedure described above for conical noses. This gave a continuous variation of  $V_t$ over the nose, dropping to zero at the base of the nose. Explicit formulas for  $V_p$  and  $p_I$  as a function of position on the nose were given for ogives (Reference 6). For the special case where the ogive is a hemisphere, the distribution of  $p_I$  over the nose as calculated by this procedure differs slightly from that given by the previous procedure assuming  $p_g$  to vary as the cosine of the polar angle. But the resultant axial force is the same when the added-mass term containing  $B_1$  is negligible.

Because of the formation of a false nose of target material, it does not seem reasonable to apply these assumptions to actual flat-nose projectiles. Neither the X-ray pictures nor post-test examination had shown the actual shape of the false noses formed in the Eglin penetration experiments. In the following analysis of Shots 20 and 21 of the Eglin experiments two kinds of assumptions were made for the chape of the false nose. The first assumption was a hemispherical nose, leading to Equation (21) for the equation of motion. The second kind of assumption was a conical nose, leading to Equation (24). The second assumption was applied for L/D values of 0.5, 0.4, 0.2, and 0.

The material properties for the dry sand target were determined as follows. The shear or deviatoric properties were based on a triaxial test performed in the Civil Engineering Laboratories at the University. Figure 55 shows a plot of  $\sigma_1 - \sigma_3$  versus  $\varepsilon_1$ , where  $\sigma_1$  and  $\varepsilon_1$  are axial compressive stress and strain and  $\sigma_3$  is the constant lateral confining pressure of

0.589 MPa. The dots denote the experimental curve while the two straight lines are the bilinear fit to it. From the bilinear fit the values of

$$E = 54 MPa$$
,  $E_{\perp} = 1.39 MPa$ ,  $Y = 1.4 MPa$  (26)

were determined for the elastic modulus  $E_v$  tangent modulus  $E_t$  in the plastic regime, and yield stress Y (at the intersection of the two straight lines). The bilinear approximation, required for the simple cavity expansion theory, seems to be quite a reasonable one for the triaxial test curve.

The required approximation to the compressibility properties, determined from a uniaxial strain test as shown in Figure 56 is more extreme and more arbitrary. It is necessary to approximate the curve by two vertical (incompressible) lines joined by a horizontal jump representing the change in density from the locked elastic density at uniaxial strain  $\varepsilon_1$ , to the locked plastic density at uniaxial strain  $\varepsilon_p$  upon the passage of the plastic shock wave. The approximations assumed correspond to uniaxial strain values of

Elastic  $c_1 = 0.03$  Plastic  $c_p = 0.094$  (27) The initial density was  $\rho_c = 1540 \text{ kg/m}^3$ . Then

$$\rho_{\rm p} = 1700 \, \rm kg/m^3$$
 and  $\rho_{\rm o}/\rho_{\rm p} = 0.906$  (28)

leading to the following numerical values for constants in the theory

$$\beta = \frac{Y}{2E} - \frac{1}{3}c_1 = 0.00296 \qquad a_p = 1 - (\rho_0/\rho_p) = 0.094$$

$$\sigma = 1 - (\rho_0/\rho_p)e^{-3\beta} = 0.1020 \quad B_1 = 1 - 6^{1/3} = 0.533$$

$$B_2 = 1.5 + (1+a_p)6^{1/3} + 0.56^{4/3} = 1.013$$

$$P_8 = \frac{4}{9}E(1-e^{-3\beta}) - \frac{2}{3}Y \ln 6 + \frac{2}{27}\pi^2 E_t - \frac{4}{9}E_t(\frac{\pi}{h=1}(6^n/n^2)) = 3.396 \text{ MPa}$$
(29)

The formulas for  $\beta$ ,  $\alpha_p$ ,  $B_1$ ,  $B_2$ , and  $p_2$  are as given in Hanagud and Ross (Reference 3), and differ slightly from the versions in Bernard and Hanagud (Reference 6). The differences have to do with inclusion of various terms that contribute little to the actual numerical values obtained. With this value of  $B_1$ , the added-mass term containing  $B_1$  in Equations (21) or (24) is 0.00285 kg, which is negligible in comparison to the projectile mass m (0.5451 kg in Shot 20 and 0.5443 kg in Shot 25). The equation then takes the form

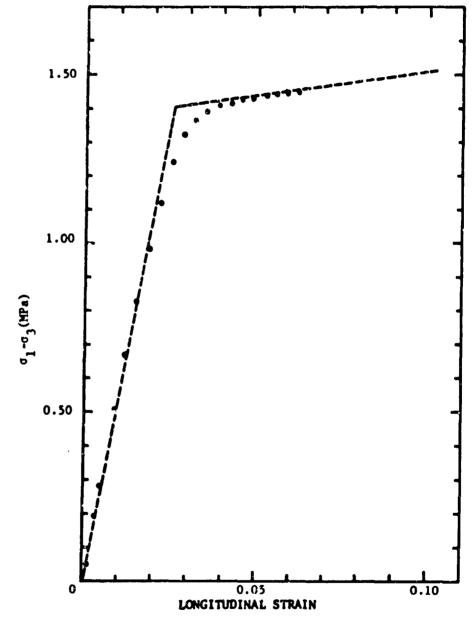
$$\frac{dV}{dt} = - (A + BV^2)$$
(30)

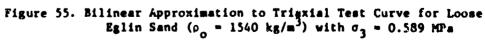
as in the Poncelet force law. Equation (30) is integrated to give

$$x = x_{o} + \frac{1}{B} ln \{ cos(\sqrt{AB} (t-t_{o})) + \sqrt{B/A} V_{o}sin(\sqrt{AB} (t-t_{o})) \}$$
(31)

and

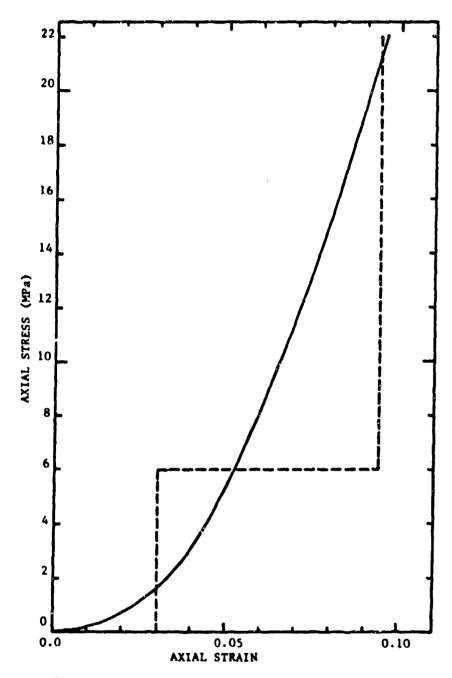
$$v = \left\{ \left( \frac{A}{B} + v_o^2 \right) e^{-2B(x-x_o)} - \frac{A}{B} \right\}^{1/2}$$
(32)



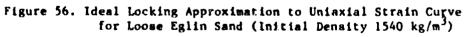


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These results are plotted in Figure 57 for Shot 20 and Figure 58 for Shot 25. The coefficients A and B for the two shots have the values shown in Table 10.

Shot 20			Shot 25				
Curve	J./D	A(m/s <sup>2</sup> )	B(m <sup>-1</sup> )	Curve	L/D	A(m/s <sup>2</sup> )	B(m <sup>-1</sup> )
1	0.5	1900	0.6614	1	0.5	1903	0.6624
2	0.4	1900	0.7340	2	0.4	1903	0.7351
3	0.2	1900	0.9009	3	0.2	1903	0.9023
4	0	1900	0.9922	4	0	1903	0.9936

TABLE 10. COEFFICIENTS FOR CAVITY EXPANSION THEORY PENETRATION CALCULATIONS

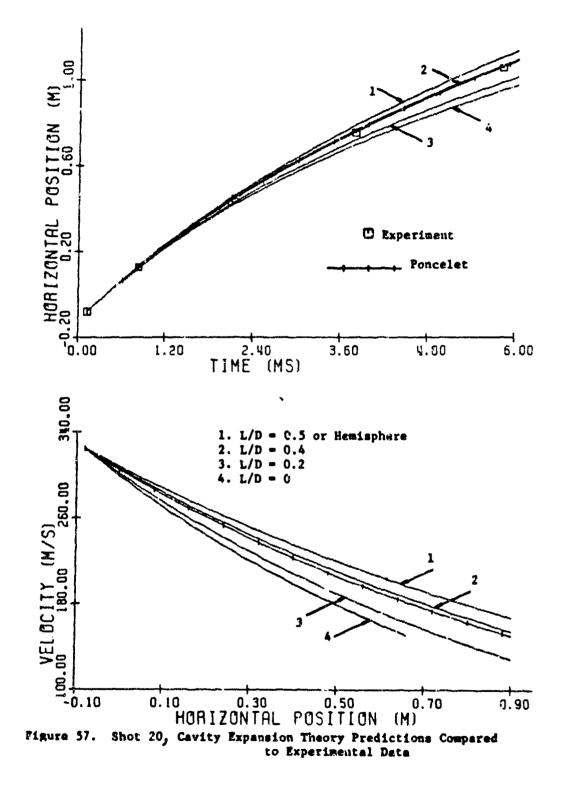
Also plotted in each figure is the fitted Poncelet curve through the experimental points. The agreement is fairly good for the hemispherical nose (Curves 1, L/D = 0.3), overestimating the penetration by 3 to 5 percent and the final velocity by about 11 percent in Shot 20. The xt-curve for L/D = 0.4 (Curve 2) is essentially coincident with the experimental curve in each case. The Vx-curve is also coincident with the experimental curve in Shot 25 and overestimates the final velocity by only 3 percent for Shot 20. It is emphasized that the parameter values used for the prediction were determined from the two static curves as shown in Figures 55 and 56. These parameters were determined before calculation of Figures 57 and 58. No adjustments were made to the parameter values to get a better agreement with the experimental results.

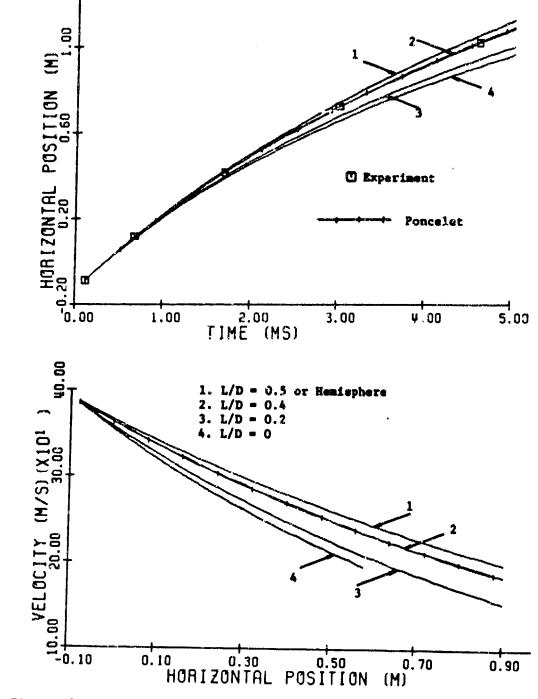
Thus, with an assumed false conical nose with L/D = 0.4 the details of the deceleration history are remarkably well predicted for Shots 20 and 25 in the region of observation. The prediction with an assumed spherical nose (or a cone with L/D = 0.5) is also not bad. The solid Reynolds number for these two shots is  $R_g = 127$  for Shot 20 and  $R_g = 181$  for Shot 25, both above the range of  $R_g$  in which the cavity expansion theory was known to give good results. (See the discussion following Equation (19).)

Any attempt to apply the theory directly to the flat-nosed projectiles by using the limiting case of L/D = 0 as in the Curves No. 4 in Figures 57 and 58 would greatly overpredict the drag and underpredict the penetration for Shots 20 and 25.

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It is concluded that the cavity expansion theory method has considerable merit despite the strong assumptions involved in its application to penetration theory. For projectiles with actual conical or ogival noses or other nonflat axisymmetric shapes it should prove profitable to consider oblique impacts and attempt to modify the theory to apply locally to a surface area element somewhat in the manner of the assumed differential area force laws to be discussed in Section VI. This has not been done yet for a





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complete trajectory, although Bernard and Hanagud (Reference 6) discussed the embedment process at the beginning of an oblique impact for a projectile with a conical nose.

It might also be possible to make similar calculations for the flat-nosed projectile, but this would severely test any assumed false nose shape. It may be possible to obtain some hints about the false nose shape from a study of the separation angles as shown in the X-rays.

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# SECTION VI

# RIGID BODY MOTION IN A SOIL MEDIUM

### 6.1 EQUATIONS OF MOTION

Soil penetration prediction techniques as classified by Triandafilidis (Reference 1) are considered to fall into two broad categories i.e., mathematical and experimental. Under the broad heading of mathematical further suggested subdivisions are semi-analytical, analytical and theoretical. The semi-analytical techniques listed in Reference 1 were all restricted to completely normal penetration. Even the cavity expansion models, classified as analytical, are also restricted to normal impact. Oblique impact may be analyzed using an analytical Differential Area Force Law (Reference 8). However, a computer program based on this type of analysis has limited access because of proprietary restrictions, as well as being expensive to operate. In light of the above, a study was initiated to develop a simple multidegree-of-freedom set of equations of motion for bodies of revolution in a scil medium. The ground rules for this development are as follows:

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- 1. The projectile was to be a body of revolution with zero rotation rate about the longitudinal axis.
- Classical six-degree-of-freedom equations of motion would be used with force and moment terms from assumed or empirical force expressions.
- 3. Force expressions would be selected by joint agreement between the contractor and the project engineer.
- Results were to be projectile position-time tabulation using Cartesian coordinates for center of mass position and Euler angles for body rotations.

For the derivation a set of body fixed axes x, y, z with unit vectors  $\hat{i}$ ,  $\hat{j}$ ,  $\hat{k}$  and an inertial frame x', y', z', with unit vectors  $\hat{i}'$ ,  $\hat{j}'$ ,  $\hat{k}'$  fixed in the soil, were selected and are whown schematically in Figure 59.

The generalized six-degree-of-freedom equations of motion for a rigid nonsymmetrical body written relative to the body axes are given as (Reference 27):

 $F_{u} = m(\dot{U} + QW - RV)$ (33)

$$F_{y} = m(\dot{V} + RU - PW)$$
(34)

$$F_{-} = m(\dot{w} + PV - QU)$$
(35)

$$L = I_{xx} \dot{P} + I_{xy} (PR - \dot{Q}) - I_{xz} (\dot{R} + PQ) + RQ (I_{zz} - I_{yy}) + I_{yz} (R^2 - Q^2)$$
(36)

$$M = -I_{xy}(P+RQ) + I_{yy}(PQ-R) + RP(I_{xx} - I_{zz}) + I_{xz}(P^2 - R^2)$$
(37)

$$N = I_{xz}(QR-P) - I_{yz}(Q+PR) + I_{zz}R + I_{xy}(Q^2 - P^2) + QP(I_{yy} - I_{xx})$$
(38)

where:

 $F_{u}$ ,  $F_{u}$ ,  $F_{z}$  Applied forces in x, y, z directions respectively.

L, M, N Applied torques or moments about x, y, z axes respectively.

- U,V,W Projectile or body translational velocities relative to x, y, z axes respectively.
- P,Q,R Projectile or body rotational velocities about x, y, z axes respectively.
- m Projectile mass.

 $I_{xx}$ ,  $I_{yy}$ , Conventional moments of inertia and products of inertia for body axem x, y, 2  $I_{xx}$ ,  $I_{xy}$ ,

I .... Ivz

÷

The dot above any term represents the time derivative or time rate of change of that term.

For a symmetrical body the products of inertia are zero and I =I, and Equations (33) through (38) reduce to

$$\mathbf{F}_{\mathbf{u}} = \mathbf{m}(\mathbf{U} + \mathbf{Q}\mathbf{W} - \mathbf{R}\mathbf{V}) \tag{39}$$

$$\mathbf{F}_{\mathbf{v}} = \mathbf{m}(\mathbf{\dot{v}} + \mathbf{R}\mathbf{U} - \mathbf{P}\mathbf{W}) \tag{40}$$

$$\mathbf{F}_{\mathbf{v}} = \mathbf{n}(\mathbf{W} + \mathbf{P}\mathbf{V} - \mathbf{Q}\mathbf{U}) \tag{41}$$

$$\mathbf{L} = \mathbf{I}_{\mathbf{u}} \mathbf{P} \tag{42}$$

$$M = I_{yy} \dot{Q} + RP(I_{xx} - I_{yy})$$
(43)

$$\mathbf{M} = \mathbf{I}_{\mathbf{X}\mathbf{X}} - \mathbf{Q}\mathbf{P}(\mathbf{I}_{\mathbf{X}\mathbf{X}} - \mathbf{I}_{\mathbf{Y}\mathbf{Y}})$$
(44)

## 6.2 FORCE EXPRESSIONS

 $\begin{array}{c} \mathbf{A}_{\mathbf{x}}, \ \mathbf{B}_{\mathbf{x}}, \ \mathbf{C}_{\mathbf{x}}, \\ \mathbf{A}_{\mathbf{y}}, \ \mathbf{B}_{\mathbf{y}}, \ \mathbf{C}_{\mathbf{y}}, \\ \mathbf{A}_{\mathbf{z}}, \ \mathbf{B}_{\mathbf{z}}, \ \mathbf{C}_{\mathbf{z}}, \end{array}$ 

n, n, n

The force exerted on the projectile by the soil was assumed to be of the form

$$\frac{d\bar{P}}{dA} = n_{\chi} (A_{\chi} + B_{\chi} | v | + C_{\chi} v^{2}) \hat{1} + n_{\chi} (A_{\chi} + B_{\chi} | v | + C_{\chi} v^{2}) \hat{j}$$
(45)  
+ n\_{\chi} (A\_{\chi} + B\_{\chi} | w | + C\_{\chi} w^{2}) \hat{k}

where:

are components of outward unit vector normal to surface of projectile.

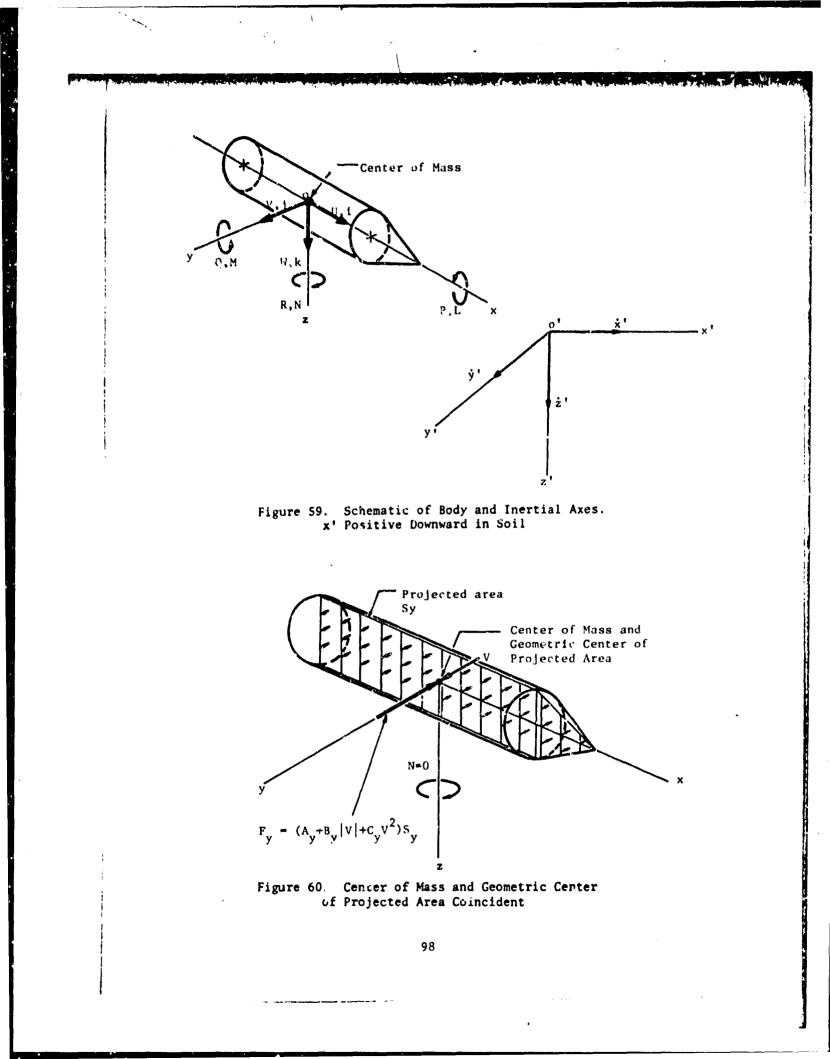
The forces  $F_{x}$ ,  $F_{y}$ , and  $F_{z}$  may then be determined by the projected wetted sreas normal to the velocities U, V, W respectively. A force is assured to exist only when the watted surface has an outward velocity component normal to the projected area. This type of force distribution applicable only to axisymmetric bodies, assumes a uniformly distributed pressure over the projected wetted area, giving rise to a resultant force passing brough the geometric center of the projected wetted area. When the greatric center of the projected wetted area coincides with the center of the mass of the projectils then no applied moment L. M. N. exists the projectile. If the projectile is hollow or if the geometric center of the projected area and the center of mass do not coincide then an applied moment exists and is equal to the applied force times the distance between the geometric center of the projected wetted area and the center of mass. For a body of revolution completely submerged in the medium the force distribution due to a lateral translational velocity V is shown in Figura 60 for the case of zero moment. In a case where the geometric center of the area and the center of mass do not coincide the force distribution is shown in Figure 61. The projected area used to determine the force Fy for a completely submerged axisymmetric projectile is the same as for  $F_{\pi}$ . The projected area for  $F_{\pi}$  for the completely submargad projectile is simply the cross section of the projectile normal to the x axis.

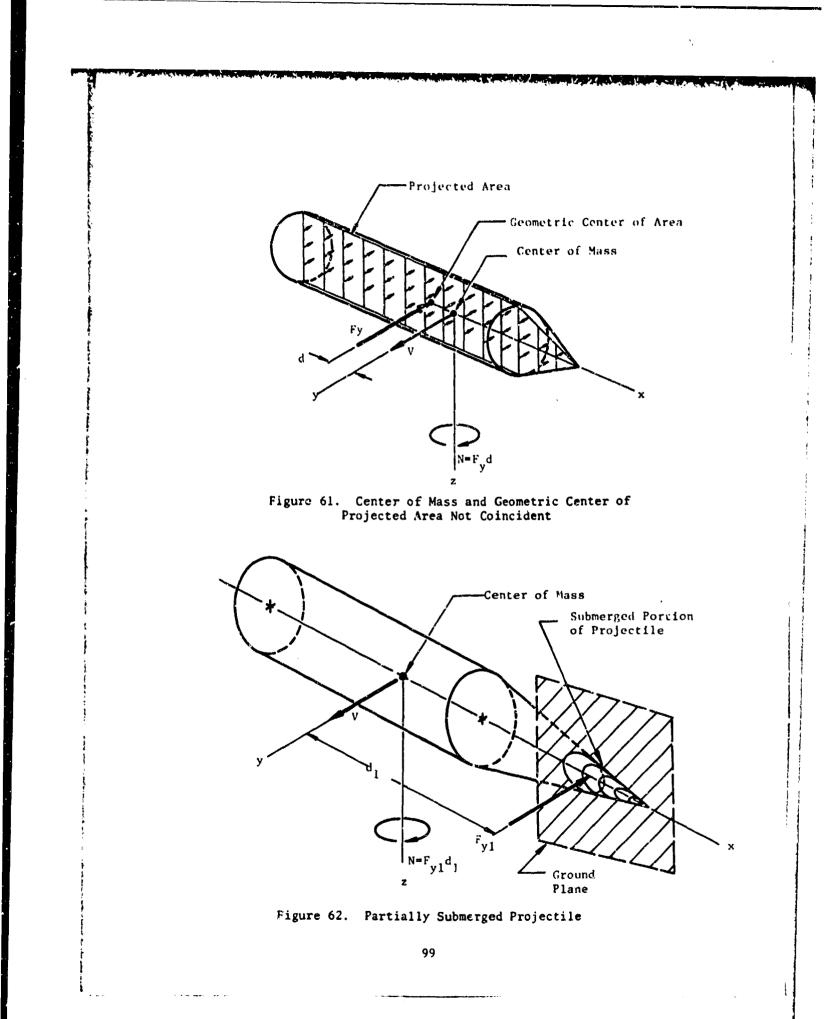
If the projectile is only partially submerged then only a portion of the total projected area is in contact with the soil medium and a resultant force and morent will exist as shown in Figure 62. These moments and forces are clanging with depth of penetration and become functions of depth. For this case the forces and moments are not simple expressions due to the complicated expressions required to calculate the areas. The derivation of the forces and moments required for partial penetration of a conical nose are given in Appendix B. These equations are not included in the main body of the report as time did not permit complete computer modeling of these equations and therefore are simply included as information.

### 6.3 COORDINATE TRANSFORMATIONS

Translational and angular positions relative to an inertial frame fixed in the soil may be expressed in terms of angular and translational velocities of the body fixed axes by use of coordinate transformations relating the two systems. The coordinate rotations required for these transformations are given in Reference 28 as:

- Start with body axis x, aligned with inertial axis x' and rotate about body axis z, through an azimuthal angle Y. This produces a new set of body axes X<sub>2</sub>, Y<sub>2</sub>, Z<sub>2</sub>.
- 2. Rotate about  $Y_2$  through a pitch angle  $\Theta$ . This produces a new set of body axes  $X_2$ ,  $Y_2$ ,  $Z_3$ .
- 3. Finally, rotate about X<sub>3</sub> through a roll angle  $\phi$ , which brings the body into its final being axis system X, Y, Z. (This rotation is not important for a body possessing complete symmetry about the X axis.





The transformation equation based on these angles is given as

$$\overline{A}_{I} = T_{BI} \overline{A}_{B}, \qquad (46)$$

which transforms an arbitrary vector  $\overline{A}_{p}$  given in the body system to a vector  $\overline{A}_{p}$  in the inertial system. The transformation matrix  $T_{BI}$  is given as

where:  $c\theta = \cos \theta, s\theta = \sin \theta, etc.$ 

For orthogonal transformation such as this the inverse of  $T_{BI}$  is equal to the transpose of  $T_{BI}$ ; therefore the inverse relation of Equation (46) is simply

$$\overline{A}_{B} = [T_{BI}]^{T} \overline{A}_{I} .$$
(48)

By using Equation (46) the velocity vectors U, V, W may be transformed to give the velocities x', y', z', in the inertial system. These velocities,

$$\begin{array}{c} \dot{x}' \\ \dot{y}' \\ \dot{z}' \end{array} - \begin{bmatrix} T_{BI} \end{bmatrix} \begin{cases} U \\ V \\ W \end{cases}$$

$$(49)$$

may be integrated to determine the position x', y', z' of the center of mass in the inertial frame.

The angular velocities  $\dot{\Psi}, \dot{\Theta}, \dot{\phi}$  are related to the angular velocities P, Q, R through a transformation matrix ROT defined by

 $\begin{array}{c} \mathbf{P} \\ \mathbf{Q} \\ \mathbf{R} \end{array} = \begin{bmatrix} \mathbf{ROT} \end{bmatrix} \begin{pmatrix} \dot{\mathbf{\Phi}} \\ \dot{\mathbf{\Theta}} \\ \dot{\mathbf{\Psi}} \end{pmatrix}$  (50)

where:

$$[ROT] = \begin{bmatrix} 1 & 0 & -\sin\theta \\ 0 & \cos\phi & \sin\phi\cos\theta \\ 0 & -\sin\phi & \cos\phi\cos\theta \end{bmatrix}$$
(51)

Due to the non-orthogonality of [ROT] the inverse of the [ROT] transformation is not [ROT]<sup>T</sup>. It is given by

sinétan0 cosétan0

 $[ROT]^{-1} = [1]$ 

and

(52)

Equations (33) through (38),(49) and (53) represent the necessary equations. When solved simultaneously with the proper initial conditions they will yield the position and orientation of the body as functions of time in the inertial frame.

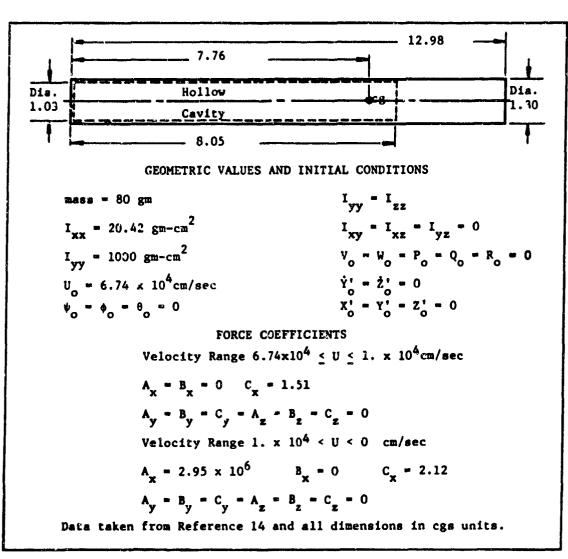
# 6.4 CALCULATIONS

The solution of Equations (33) through (38), (49) and (53) may be obtained numerically provided expressions for the forces and moments are available. In the general case where impact is not normal the expressions are rather complicated and require considerable computer technique and programming ability. However, a completely submerged projectile could be handled very easily for the force distribution as given by Equation (45). Also if the assumption is made that only the nose of the projectile is in contact with the soil and submersion of the nose is instantaneous, a normal impact can be handled very easily.

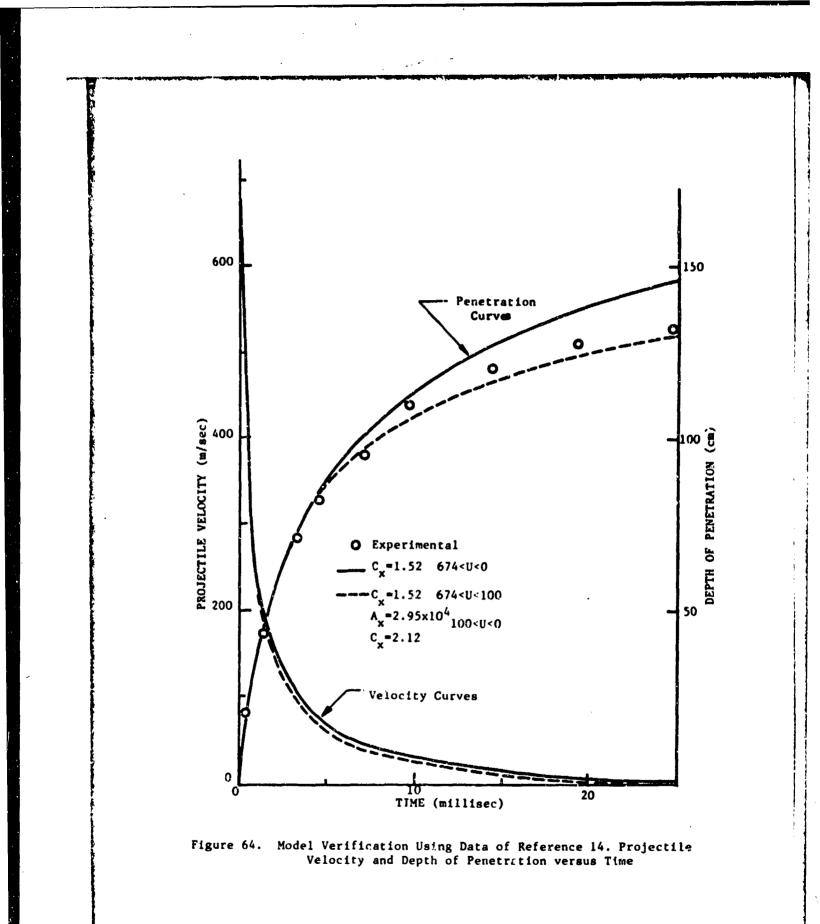
The case of complete submersion was programmed using a MIMIC source language program. This program, available on the CPC 6600 in the Mathematical Laboratory at Eglin AFB, is essencially a fourth order Runge-Kutta numerical method for solving simultaneous differential equations. A program shown as Computer Program I of Appendix C was developed for solution of Equations (33) through (38), (49) and (53), for a blunt nosed cylinder, a conical nosed cylinder and a hemispherical nosed cylinder. For this case the force coefficients A through C are assumed to be constents.

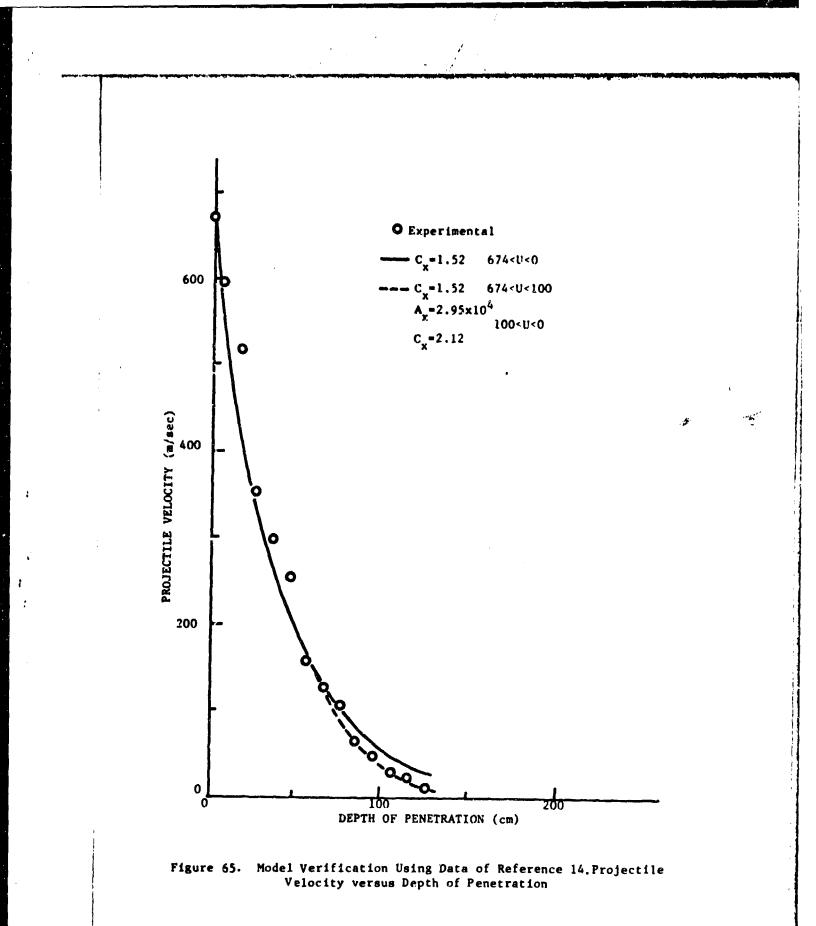
The MIMIC program allows for naming constants or variable parameters and for Program I, Appendix C, all the initial conditions, soil force coefficients, and geometric properties were given a parameter status. Only integration time and print frequency were named constants. A check ca this program was accomplished using the data for the blunt nosed cylinder listed in Figure 63. The assumption in this case is that, for normal penetration, nose submersion is instantaneous and only the nose is in contact with the soil.

For this case only the forces parallel to the x axis of the projectile are operative; therefore all other coefficients are set to zero. It is important to note here that the coefficients used for this case were obtained from test date of Reference 14. The results of this case are shown in Figures 64 and 65. Both plots show very









good agreement between model and experimental values. Both figures show a comparison between using a constant value of C for the whole velocity range and using values for A and C for velocities below the critical velocity of 100 m/sec. The overall effect of using • a value for A is to reduce the depth of penetration at the lower velocities and bring the projectile to rest at some finite time. The significant difference of the two cases is shown in the reduction of depth of penetration of Figure 64.

Further verification of the analytical model was obtained by use of experimental data of the Eglin experiments. As discussed in Section II velocity measurements of solid blunt nosed cylinders were made using an X-ray technique. The reduction in velocity during the trajectory was assumed to be attributed to a drag term based on the expression

$$F_{x} = \frac{1}{2} C_{D^{\mu} s} A(\dot{x}')^{2}$$
(54)

where  $C_{D}$  is defined as the drag coefficient in the x' direction,  $\rho$  is the undisturbed soil density and A is the cross sectional area of the projectile. By using this assumption and analyzing the data as in paragraph 4.6, it was found that C had a variation with velocity as given in Table 11. Also, if x'and x are assumed to be collinear, then  $\dot{x}'$  and U are equal.

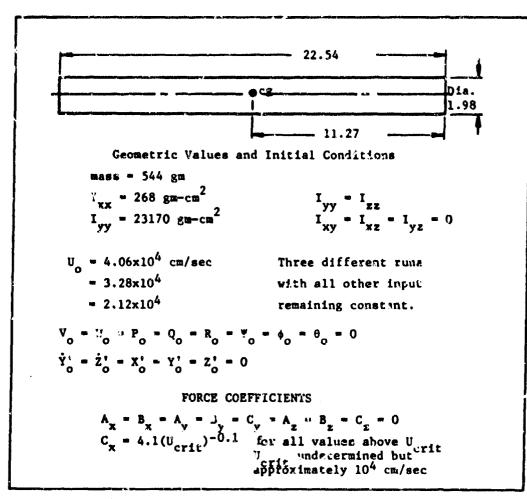
DRY	SAND
x' m/sec	с <sub>р</sub>
0	4.0
12.7	3.0
25.4	2.2
76.2	2.0
152.4	1.9
228.6	1.25
304.8	1.25
2540.0	1.7

TABLE 11. VARIATION OF C WITH VELOCITY \*'.

The data of Table 11 fit a power law of the form

$$C_{D} = aU^{D}$$

(55)





where a = 3.12 and b = -0.1 when U is expressed in cm/sec. if equation (54) is compared to Equation (45) then the relation between C and C is given as

$$C_{\mathbf{x}} = \frac{1}{2} \rho_{\mathbf{g}} C_{\mathbf{p}}.$$
 (56)

The density of the soil p used in the experiment was 1.6 gm/cm<sup>3</sup>; theref re  $C_y = 0.8C_p$ .

Modifying the Computer Program I to include Equations (54)to (56) results in the Computer Program II given in Appendix C. The data used as input for this case are given in Figure 66.

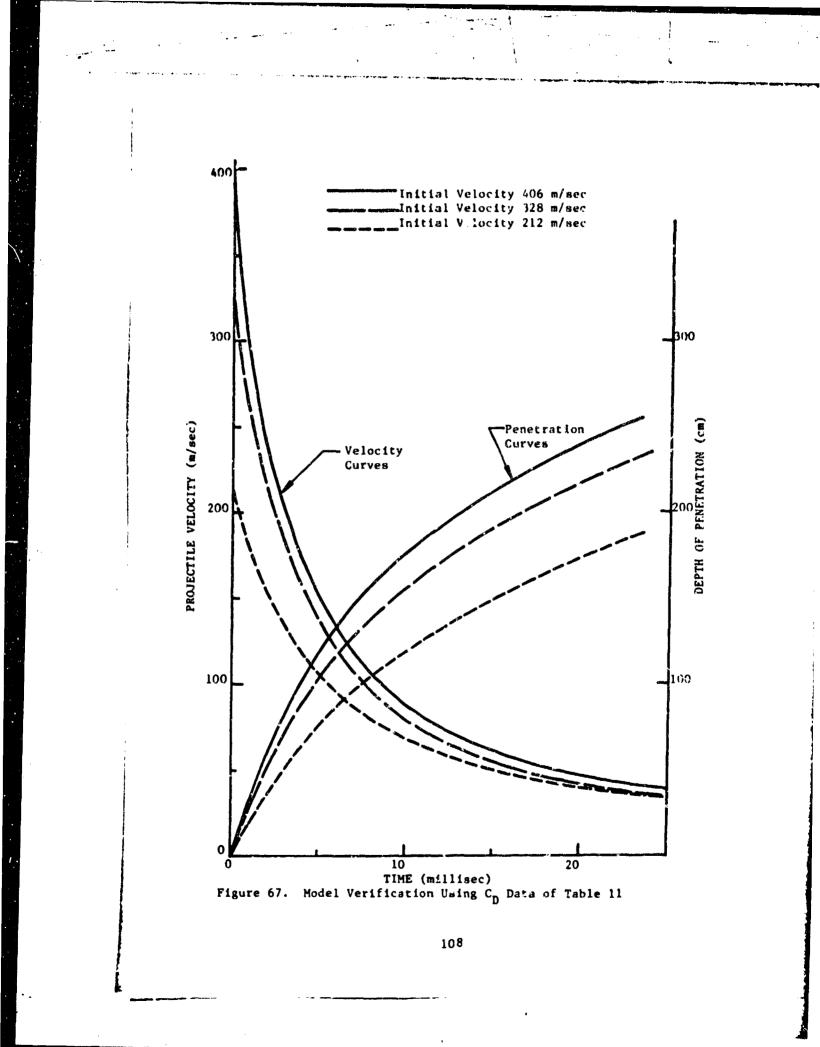
The results of these runs are given in the graphs of Figures 67 and 68. Figure 57 shows the variation of depth of penetration with variations in impact or initial velocity. Figure 68 shows correlation of model prediction with experimentally determined data.

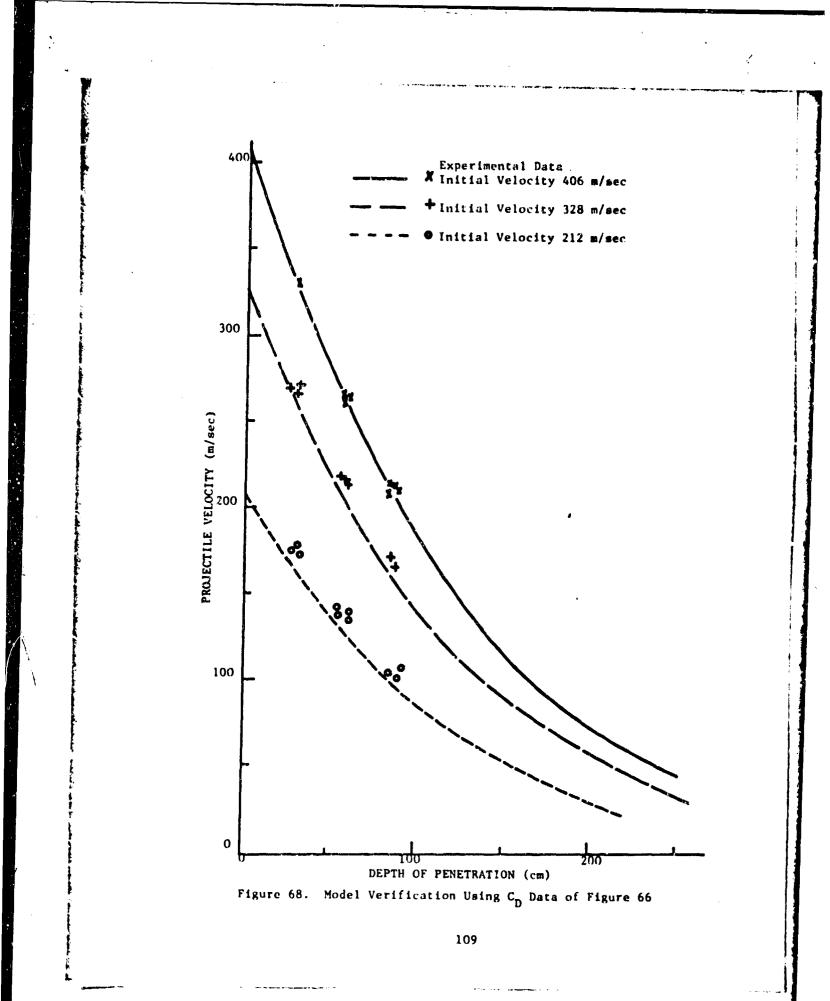
### 6.5 CONCLUSIONS

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The results of the preceding paragraph show that a simple basic terradynamic approach using experimentally determined force coefficients will yield reasonable results. However, it must be emphasized that this method is highly dependent on good data obtained by experiment. The lack of force coefficients for angle of attack and other than normal impact prevents model verification for a general case.





### SECTION VII

### SONIC AND ULTRASONIC WAVE SPEED MEASUREMENTS

## 7.1 INTRODUCTORY REMARKS

The preceding sections have dealt with the testing and performance characteristics of terradynamic vehicles, including trajectory, cavity shape, separation, reattachment and stability. Predictive techniques which allow for a quantitative assessment of earth penetrating vehicle performance require information on certain parameters. Important among these parameters as inputs into terradynamic models are the density and the acoustic impedance of the target medium (Reference 29). The acoustic impedance, which is related to the wave velocity in the material, has particular importance in delineating regimes of application of penetration equations as well as in perhaps predicting bow wave speeds observed in the X-ray studies.

One technique which has been explored in the current test program for obtaining information on the above parameters is ultrasonic wave speed measurement. The use of ultrasound as a diagnostic and measurement tool is well documented. Pohlman (Reference 30), for example, has catalogued ultrasonic research topics in a series of volumes with frequent updating of the literature. Specific descriptions of ultrasonics in diagnostic applications have also been discussed in such references as (References 31,32), while techniques for measurement of mechanical parameters are contained in References 33 through 35. Hany mechanical properties measurements by ultrasound have been directed toward obtaining information on the elastic properties of either solid, liquid, or gaseous media using the pulse-echo or through-transmission techniques, as described for example by McSkimin (Reference 36), Papadakis (Reference 37) and others (References 33 through 35). Some recent properties measurements on solid heterogeneous media such as fibrous composite materiels and rock media have been reported on in References 40 through 42. These media, while dispersive in nature, remain amenable to conventional ultrasonic testing procedures because of the retention of specimen shape during machining. For granular or solid media, in which three distinct phases are present (solid, liquid and gas), and for which confined samples are not readily produced, measurement of mechanical properties becomes more involved. Some data collection related to soil media has been reported in References 43 through 46. Because of the phase inhomogeneity, properties measurements for soil types such as dry or saturated sands are difficult tasks. For example, dilatational wave speeds through dry sands by sonic radiation provide reasonable properties data, while similar measurements through saturated sand have proved unsatisfactory (Reference 43).

The filled pores allow rapid propagation of dilatational waves, so that measurements of the waves transmitted by the skeletal phase generally require the measurement of shear waves instead. These waves appear to reflect a better standard measurement of the skeletal stiffness and are frequently used in testing soils properties at ultrasonic frequencies. For this reason many of the tests reported in the literature report data on the propagation of shear waves (Reference 44).

In obtaining such data it is important to note that considerable disparity in the reported magnitude of acoustic waves in soils appears in the literature. This discrepancy is partially due to the measuring technique used, type of pulse disturbance used for generating the transmitted signal, and amplitude of resulting disturbances. In any event it is important to note that current analytical models appear inadequate to predict wave velocities and it is necessary to obtain quantitative measures of the wave velocities in real soils by experimental procedures (Reference 44).

It the studies reported in this section, wave speeds in any or moist Eglin sand (5 to 15 percent moisture by weight) have been investigated with three purposes in mind: (1) to obtain input for existing penetration codes requiring this information or for delineating bunds on the usefulness of terradynamic equations (Reference 29), (2) for poter tal relationship to observed how shock wave speeds (Reference 21), and the possible use as a tool for establishing the compaction state of sand.

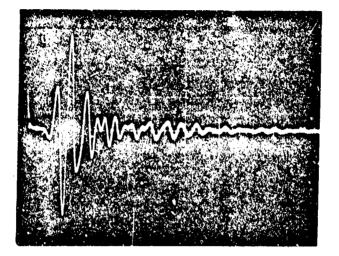
The major part of the investigation has been concerned with ultrasonic wave speed measurements as a function of compaction and testing pressure. Several difficulties were encountered with the testing program, and it still has had only limited success. The difficulties were associated with the dispersive nature of the sand medium, especially at high frequencies and low testing pressures. The ultrasonic wave speed measurement techniques and results will be described in paragraph 7.2.

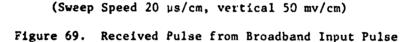
A brief discussion will be given in paragraph 7.3 of some low frequency field measurements of sound wave speeds.

#### 7.2 EXPERIMENTAL PROCEDURES AND RESULTS FOR ULTRASONIC WAVE SPEEDS

A Panametrics Ultrasonic Intervalometer system was available at the University that could be used either in a pulse-echo-overlap method with the same transducer used both for sending and for receiving the reflected signal or in a through-transmission method with separate mending and receiving transducers. Because of the dispersive nature of the medium, effort was concentrated on the through-transmission method. The Panameterics system can be used alone with broad band single pulses, or in conjunction with a pulsed radio-frequency (RF) oscillator it can be used with a burat of RF oscillations.

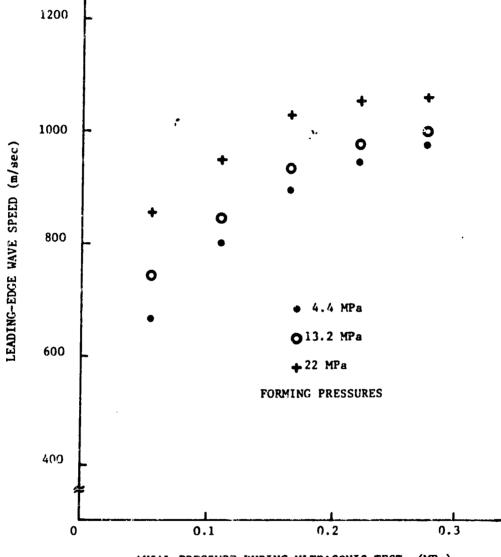
The first testing of sand samples used the single broad band pulse from the Ultrasonic Pulsing Module of Panametrics system as input signal to a Panametrics Type V201 5 MHz longitudinal transducer. The sand was firsc compacted under conditions of uniaxial strain in a steel cylinder 0.05 meter in diameter under axial pressure of 4.4 to 22 MPa. After unloading, end plates each containing one of the transducers were mounted on the cylinder and the broad band pulse applied to one end. The broad band pulse contains all frequency components, but the received signal resembled a distorted sine wave with the first few oscillations at a frequency 0.02 times the 5 MHz transducer resonant frequency. An example is shown in Figure 69.

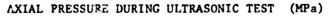


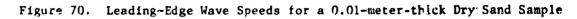


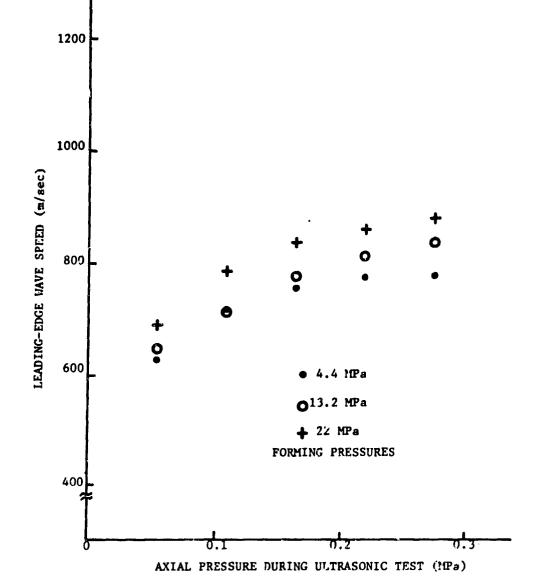
A portion of the input broadband pulse is seen at the beginning of the oscilloscope trace, driving the signal off screen. The received pulse is amplified and mixed with the broadband pulse internally in the Panametrics unit and then displayed on an oscilloscope. Precise timing can be accomplished in a manner similar to that for the RF bursts as will be described later in this section. More details on the operation of the Panametrics system are given in Reference 45.

The pulse displayed in Figure 69 was transmitted through a sand sample 0.0094 meter thick while under an axial pressure of about 0.2 MPa after compaction by an axial pressure of 4.4 MPa in the steel cylinder. The attenuation of the pulse by the sand was so great that this equipment and procedure could not be used with samples much thicker than 0.05 meter. Also the signal became erratic when the applied axial pressure during the wave speed measurement was about 0.1 MPa(approximately 15 psi) and disappeared altogether at some lower pressure. But the most disappointing feature of the results was that the measured wave speed appeared to depend on the thickness of the sample. This is a result of the change in shape of the pulse as it propagates through the dispersive medium. What was measured in these tests was the speed of propagation of the leading edge of the received pulse, which would be the group velocity of the transmitted wave packet of oscillations if the transmitted packet did not change its shape so much.









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Figure 71. Leading-Edge Wave Speeds for a 0.023-meter-thick Dry Sand Sample

Figure 70 shows the leading-edge wave speed for a 0.01-meter-thick dry sand sample versus axial testing pressure varying from about 0.05 MPa to about 0.30 MPa for specimens previously compacted at three different axial pressures. Figure 71 shows the same kind of plot for a 0.023-meter-thick dry sand sample. Figure 72 is a photograph of the steel cylinder in place in a fixture mounted in a Tinius Olsen universal testing machine to provide the axial force during testing. The two transducer leads can be seen coming out of the fixutre. In this setup the transducers (not visible) are inside the end plates in direct contact with the sand.

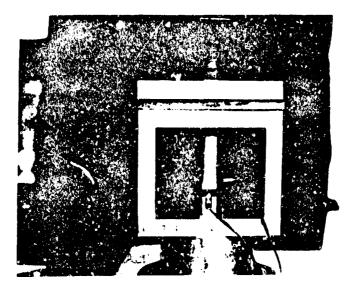


Figure 72. Fixture for Ultrasonic Wave Speed Measurements in Sand Contained in Cylinder Under Axial Load

Some additional tests of this type were performed on samples with moisture contents of 5, 10, and 15 percent by weight. At the highest forming pressure the 15 percent sample was saturated. The other samples were not saturated. The results indicated that increasing the moisture content increases the attenuation and decreases the wave speed, but the leading-edge speed results were again imprecise because of the changing shape of the transmitted pulse.

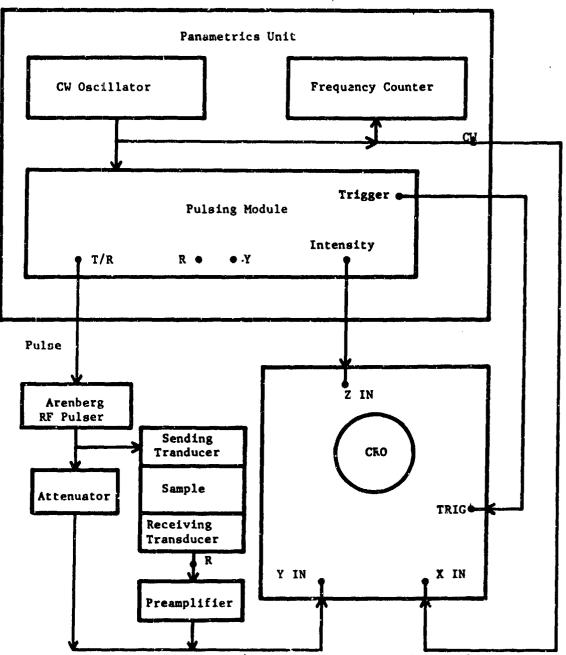
It was proposed then to obtain an RF pulser, since it was believed that the RF bursts, containing a single dominant RF frequency could be used to measure group velocity of the RF burst in a manner that would not appear to depend upon specimen thickness. It was also proposed to use the new equipment to study the effect of various amounts of air and water in the three-phase sand medium. In particular it was proposed to compare results with the predictions of an equation derived by Liahov (Reference 46) and discussed by Cristescu (Reference 47). An Arenberg Ultrasonic Laboratory oscillator, Model II (PG-65-2-C) with 400 watts peak power was obtained. It was furnished with three coils for the RF frequency ranges of 4.3 to 7.5 MHz, 0.81 to 1.1 MHz, and 0.45 to 0.62 MHz. Additional coils can extend the range to operate anywhere in the 0.13 to 190 MHz frequency range with some loss of power at the lower frequencies.

Figure 73 shows a block diagram of the Arenberg oscillator (RF pulser) connections with the Panametrics unit. The single pulse from the Pulsing Module of the Panametrics unit triggers the Arenberg RF Pulser which then emits an RF frequency burst of variable length (e.g. 5 to 20 cycles at 0.5 MHz). This input pulse travels to the sending transducer which sends a mechanical stress wave burst through the sample to the receiver and preamplifier and then to the oscilloscope (CRO). The input pulse is also attenuated and fed directly to the CRO y-input terminal for comparison with the received signal.

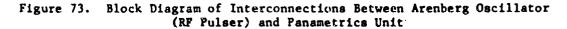
If the transmitted signal has the same wave form as the input signal, very precise timing can be obtained as follows. After the attenuated input signal and amplified output signal have been displayed on the oscilloscope (CRO) using the internal sweep of the CRO, the final precise measurement is made by switching to a sweep provided by the variable-frequency CW oscillator of the Panametrics system. The sweep frequency is adjusted so that the transmitted and received signals are made to overlap. The time interval between the two signals is then measured by the frequency counter of the Panametrics system.

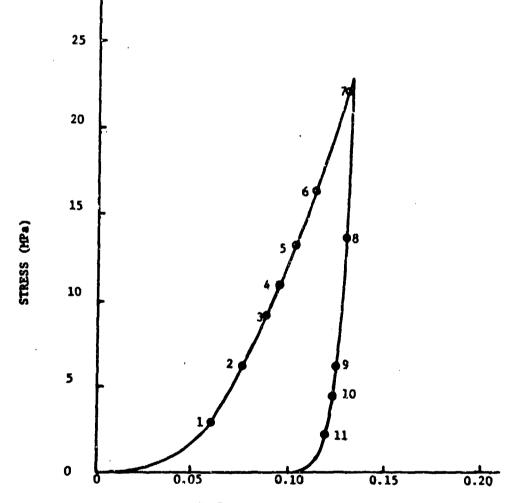
In actuality the transmitted wave pulse is modulated by the transducers and the transmission through the sand, so that the output resembles the received signal from the broadband input (see Figure 69) more than it resembles the input RF signal, which has an essentially square envelope. But the oscillations within the pulse are at the frequency of the RF, and the overlapping technique can still be applied, except at low testing pressures where there is excessive distortion. For best results the RF burst should contain at least 20 cycles and the overlapping should be made to coincide at the middle of the burst, since there is some distortion at the beginning and at the end of the burst. Because the great attenuation made it necessary to use a very short specimen path, it was not possible to use such a long pulse, and usually the overlapping was performed on the second or third peak in the burst. This use of the RF technique did succeed in removing the apparent dependence of the measured wave speed on the thickness of the sample when the test was performed at axial pressures above about 0.5 MPa. Differences in wave speed between the different specimens was within 3 percent which is considered very close agreement for different sand samples. The sand specimen holder was redesigned so that the transducers were not in direct contact with the sand but transmitted the signal through the steel end plates without themselves being subjected to the static axial loads. They could thus be left in place during the axial loading to compact the sand and could measure wave speed at various times during the loading and unloading.

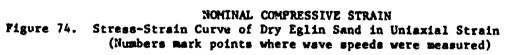
Figure 74 shows a stress-strain curve for a uniaxial strain test carried out in the steel cylinder. At the points marked on the curve, wave speed measurements were performed, with the results shown in Table 12.



For use with broad band pulse, the T/R terminal is connected directly to the sending transducer, and the receiving transducer is connected to the R terminal on the Fanametrics unit for internal amplification and mixing with the input pulse.







Point	Stress (MPa)	Density Ratic $\rho/\rho_0$	Wave Speed (m/sec)
Loading			
1	3.08	1.065	910
2	6.14	1.081	1195
3	9.22	1.097	, 1280
4	10.98	1.106	1340
5	13.17	1.114	1390
6	16.48	1.128	1470
7	22.17	1.152	1600
Unloading			
8	13.61	1.150 1420	
9	8.78	1.145 1280	
10	6.14	1.142 116	
11	4.39	1.140 1070	
12	2.19	1.136	960

TABLE 12. WAVE SPEEDS DURING UNIAXIAL STRAIN TEST

These wave speeds are group velocities of the RF bursts. They show a very strong dependence on the testing pressure. It is remarkable how little difference was measured between the velocities during unlosding and those during loading, despite the different compaction states and the different slopes of the loading and unloading curves. It seems clear that ultrasonic wave speed measurement will not be a good tool for determining the compaction condition.

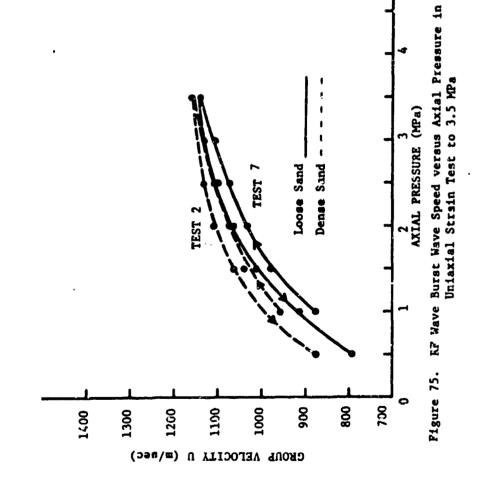
For the last series of tests the 5 MHz transducer was replaced by a 1 MHz Panametrics Type V103 transducer, and testing was again performed with bursts of 0.5 MHz. By operating nearer the transducer resonance and by using the larger 1 MHz transducers with 4 times as much frontal area a greater power could be transmitted. It was hoped that this would permit testing at much lower axial pressures more like the ambient pressures in the penetration experiments. A signal was in fact received at pressures well below the previous minimum testing pressure of 0.5 MPs. but the transmitted signals were so badly distorted that the pulse overlap technique could not be used. The oscillations in the transmitted pulse appeared to be at a frequency about 20 percent below the RF frequency of the input signal. The reason for this is not clear, but it is believed to be related to the fact that at these low pressures the phase velocity is so low that the wave length approaches the order of magnitude of the sand particle size. Thus the medium no longer responds as a continuum. Because of the distortion, only leading-edge wave speeds could be measured at the low pressures and these with the same kind of errors and apparent dependence of the wave speed on specimen thickness as were previously observed with the broadband pulses at all testing pressures.

At higher testing pressures, however, good clean RF bursts were transmitted by the 1 MHz transducers operating at 0.5 MHz, and the group velocities could be measured quite accurately by the overlap technique.

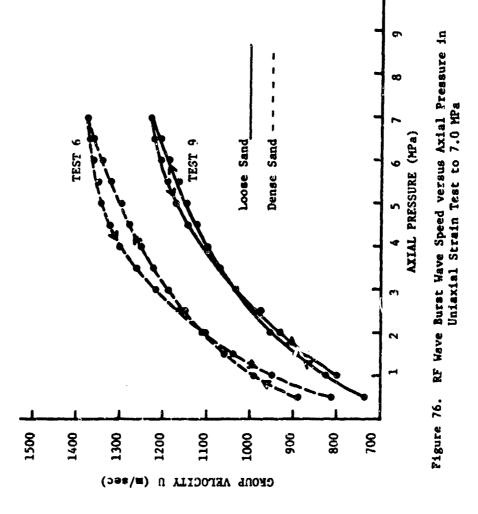
The last series of tests examined the variation of the group velocity with testing pressure for pairs of dry sand specimens all of about 0.025 meter thickness with different initial densities. The two different initial compaction states at the same pressure were achieved by shaking one of the two specimens (by tapping the side of the steel container with a hammer) to compact the sand instead of by initially compressing it under axial load. Results for three such pairs of specimens, tested by loading to three different maximum values of axial pressure are shown in Figures 75 to 77. Each figure thus gives one curve for an initially loose sand (solid curve) and one for an initially dense sand tested over the same range of pressures. The so-called dense sand had an initial density 4 to 7 percent greater than that for the initially loose sand. The additional density increase during the test varied from about 1 percent for the dense sand in Figure 75 to about 5 percent for the loose sand tested to a higher pressure in Figure 77. The arrows on each curve indicate the direction of loading or unloading. Although the loading and unloading curves are not identical, the wave speeds at any testing pressure do not vary a great deal from the loading curve to the unloading curve. The extreme values from the three curves are listed in Table 13. The last-point densities upon unloading were not recorded, but they are approximately equal to the maximum densities.

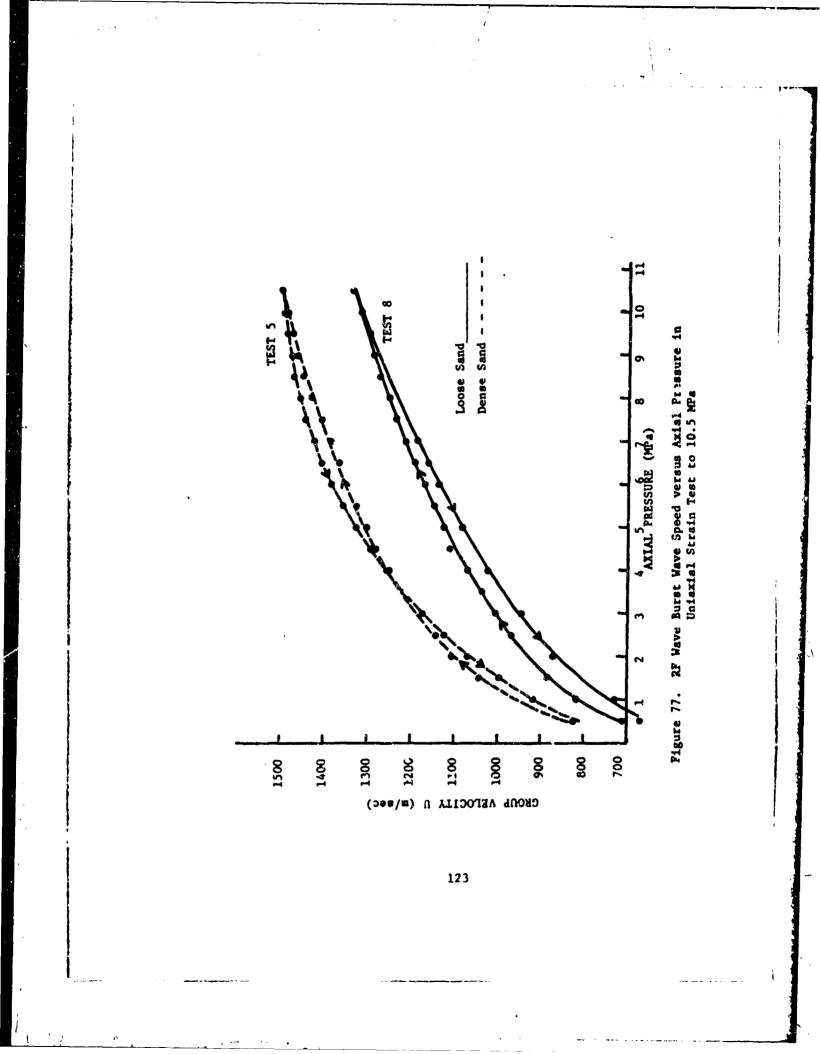
Evidently the wave speeds are much more dependent on testing pressure than on the compaction state, so that wave speed measurements are not a good measure of compaction state.

Because of the difficulties encountered in the program, time did not permit further testing of wave speeds and attenuation as a function of moisture content. The Liahov equation (References 46,47) predicts significant differences in sound wave speeds in almost saturated sands for very small changes in the air content. Variations in the air volume fraction from 0.005 to 0.04 would change the wave speed by a factor of a third. It was not possible to attempt any verification of this theoretical prediction in the present program because of the difficulty in controlling the air content as well as the difficulty in measuring ultrasonic velocities at low pressures in sand.









		Density (kg/m <sup>3</sup> )	Pressure (MPa)	Wave Speed (m/sec)
Figure 75				
Loose Sand	First Point	1638	1.0	875
Test 7	Maximum	1666	3.5	1140
	Last Point		0.5	790
Dense Sand	First Point	1726	1.0	960
Test 2	Maximum	1742	3.5	1170
	Last Point		0.5	875
Figure 76				
Loose Sand	First Point	1573	1.0	735
Test 9	Maximum	1636	7.0	1230
	Last Point		0.5	810
Dense Sand	First Point	1677	1.0	890
Test 6	Maximum	1723	7.0	1370
	Last Point		0.5	810
Figure 77				
Loose Sand	First Point	1585	0.5	710
Test 8	Maximum	1653	10.5	1335
	Last Point		0.5	670
Dense Sand	First Point	1644	0.5	835
Test 5	Maximum	1723	10.5	1500
	Last Point		0.5	810

TABLE 13. EXTREME VALUES OF PRESSURE AND WAVE SPEED IN FIGURES 75 to 77

The difficulty at the lower pressures is now believed to be caused by the fact that the lower wave speeds at the lower confining pressures lead to wave lengths of the order of magnitude of particle size. At low pressures, wave speeds of the order of 250 m/sec have been reported (Reference 14) although not at ultrasonic frequencies. The relationship between phase velocity c, frequency f, and wave length  $\lambda$  is (Reference 48)

 $c = f\lambda$  (57)

At f = 1 MHz, and c = 1000 m/sec (the order of magnitude of the group velocities observed in the present program under high pressures) the wave length is  $10^{-3}$  meters while the sand grain size is of the order of  $10^{-4}$ to 2 x  $10^{-4}$  meters. If at low pressures a speed of c = 250 m/sec could be expected, this would give a wave length at 1 MHz of 2.5 x  $10^{-4}$  meters, about the same as the grain size. If the same speed 250 m/sec prevailed at lower frequencies, it would give  $\lambda = 5 \times 10^{-4}$  meters at 0.5 MHz and  $\lambda = 10^{-3}$  meters at 0.25 MHz.

This suggests that the high frequencies are not suitable for use in the sand. A brief discussion of some lower frequency testing is given in paragraph 7.3. The discussion of Equation (57) applies to phase velocities, while the measurements reported in this section have all been group velocities. In a dispersive medium, the phase velocity of a dilatational plane wave is a function of wave length, say  $c = c(\lambda)$ . A wave packet, such as the RF bursts of the experiments described, contains a spectrum of phase velocities with a dominant mean phase velocity, say  $c_c$  at wave length  $\lambda_c$ . The packet containing wave lengths predominantly near  $\lambda_c$  travels at a group velocity U, which is related to  $c_c$  by the following equation (Reference 48)

$$U = c_0 - \lambda_0 \left(\frac{dc}{d\lambda}\right)_0$$
 (58)

If the medium is nondispersive  $dc/d\lambda = 0$ . Then  $U = c_0$  and U is independent of wave length and frequency. But if the medium is dispersive,  $dc/d\lambda \neq 0$ , and the group velocity may differ markedly from the phase velocity. The phase velocities for the sand are not known at these frequencies, so an evaluation of U by Equation (18) is not feasible. It was observed in some preliminary tests that the group velocity of the RF was somewhat frequency dependent, but time and available equipment did not permit a determination of the frequency dependence over a wide frequency range.

The pulsed RF measurements can be used to determine phase velocity by using a specimen of thickness equal to one wave length (Reference 39), but this is not easy to do with the short wave lengths of the ultrasound. The technique is similar to that used with continuous waves. Some preliminary tests with continuous waves at lower frequencies are described in paragraph 7.3.

#### 7.3 SOUND WAVE PHASE VELOCITIES

In July 1976 some preliminary field tests of low-frequency sonic phase velocities were performed at Eglin. The source for the sound waves was a 50-pound dynamic force MB vibration-testing shaker and power amplifier belonging to the University. A 0.1-meter-diameter aluminum plate was fabricated at the University and mounted on a shaft attached to the shaker armature and extending 0.15 meter outside a wooden box containing the shaker. The box was partly buried in the ground.

The tryout of the equipment was conducted jointly with the Mines Branch AFATL/DLJM, which provided geophones, recording equipment, and Apectrum analysis. The setup worked well. A good strong continuous wave signal was obtained at a distance of 3.7 meters from the shaker, even when the power amplifier driving the shaker was operating at very low power. Higher power sometimes led to distortion of the sine-wave signal. The technique involved recording the signal at two stations. Frequency was increased slowly until the two signals were in phase, indicating that the two stations were one wave length  $\lambda$  apart. The dilatational phase velocit! c is then given by  $c = f\lambda$ .

The preliminary tests indicate some dependence of wave speed on frequency, varying from about 108 m/sec at 59 Hz to 170 to 180 m/sec at around 100 Hz for two receiving stations 1.85 meters apart. The phase velocity at 100 Hz was comparable to the speed of 168 m/sec determined from the leading edge wave speeds of pulses produced by hammer blows. Higher frequency tests (up to about 19,000 Hz) were also recorded on magnetic tape for later analysis by methods not requiring the two signals to be in phase. Further tests of this type should be performed. It seems to be a good method for low-frequency sound-wave speed measurements, although care must be exercised in interpreting the results, which may be affected by reflections from the free surface and/or from internal boundaries in stratified media, especially when the two geophones are more than one wave length apart.

#### SECTION VIII

### SUMMARY AND CONCLUSIONS

The results and conclusions of each phase of the investigation have been reported in previous sections. This final section summarizes them and indicates where more detail about them may be found. Section II described the Eglin experimental program and the various types of sensors used in it or evaluated for possible use. The sequential flash X-ray technique was judged to be the most successful method investigated, since it not only gave more complete and precise information about the trajectory and the projectile's position and attitude at various times than did any other method but also gave information on cavity formation and separation points and, in some cases, showed a shock wave ahead of the projectile. This investigation is believed to be the most extensive use ever made of flash radiography in terradynamic research. The magnetic sensors also provided good information about horizontal velocity.

The results of the experiments were described in Section III and interpreted in Section IV. The trajectory plots of paragraph 4.2 for the primary test program showed that the flat-nosed and step-tier projectiles had followed remarkably straight and stable horizontal paths through the 1.2-meters-long test chamber, although most of them exhibited a slight rise. Because the paths were so nearly straight, analysis by one-dimensional terradynamic models was feasible. A cubic interpolation formula gave a very accurate representation of the horizontal position-time data, and of the velocity near the middle of the interval. A classical Poncelet forcelaw penetration model, discussed in paragraph 4.3.2, gave an excellent account of the observed parts of the trajectories in dry sand, with a drag coefficient essentially independent of the striking velocit in the range of velocities observed. In saturated sand, each shot could be fitted by the Poncelet model, but the drag coefficient appeared to depend on the striking velocity, which showed that the Poncelet model does not really apply.

Drag coefficient variation along the trajectory was exhibited in paragraph 4.6. Although the velocity calculations of that section, each based on average velocity between only two stations, tend to magnify an error at one of the stations, they do show a trend of variation along the path, more pronounced in the wet sand cases than in the dry. The classical Poncelet force law gave more consistent results than a modification of the Sandia empirical method. In all of the analyses of Section IV, force law coefficients were determined to fit observed penetration data, and the success of a model was judged on the basis of agreement between the coefficient values fitted to the different shots.

The cavity-expansion penetration model of Section V, on the other hand, attempts to predict the penetration behavior from statically measured soil properties. Despite the rather strong assumptions involved in this simple analytical model, it gave very good results in predicting the behavior for two flat-nosed projectiles in dry mand. It was necessary to assume a shape for the false nose of sand carried along by the flat-mosed projectile. An assumed hemispherical nose or a conical nose with length-to-diameter ratio of 0.4 to 0.5 led to very close agreement between the predicted and observed position-time and velocity-position curves, even for shots in a velocity range higher than the range for which previous investigations had validated the method. The success of this model suggests that it should be considered further, possibly for oblique impacts.

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A three-dimensional trajectory analysis based on an assumed three-dimensional differential force law was presented in Section VI. The procedure was carried through for a case of a straight trajectory with a drag coefficient varying with velocity according to a power law, with reasonable results. It could be applied to a trajectory with an angle of attack or an oblique impact if suitable force coefficients could be determined or estimated.

fection VII reported on an independent investigation of ultrasonic wave speeds as a function of sand compaction and testing pressure.' Several difficulties were encountered in the investigation. Pulse shape changes made it impossible to establish single broadband pulse propagation speeds that were independent of path length. This difficulty was overcome by using RF bursts instead of single pulses. Group velocities measured for these bursts gave (with an RF frequency of 0.5 MHz) consistent results for ambient testing pressures greater than about 0.5 MPa, but at lower pressures the signals were too badly distorted to give consistent results. This is believed to be a result of the fact that the lower pressures lead to lower wave speeds and a wave length of the order of sand grain dimensions. It is recommended that in further studies of sound wave speeds in sand attention be concentrated on lower frequencies. The group velocities showed a greater dependence on the testing pressure than on the compaction state.

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### APPENDIX A

### DATA FROM EGLIN PENETRATION EXPERIMENTS

This appendix tabulates data from 7% shots of the Eglin experimental program, using one page for each shot, and listed in order by Shot Numbers. Descriptions of the various kinds of entries in the tables, both of experimentally measured data and of several kinds of information calculated in the data analysis, are given in paragraph 3. Further explanations of some of the calculated data groups may be found in paragraphs 3.3 and 4.4.

### SHOT 14 ( 16 MARCH 1976, NO. 3 )

SAND: URV. DENSITY:1538. KG/M##3: APPROACHING VELOCITY: 420. M/S Projectile: Solid Flat Nose Mass:3.5451 kg. D=0.02 M. L=0.225 M

A-RAY STATION	NU . 1	NO. 2	NO+3	NO • 4	NQ.5
TIME (SECOND)		+001304	.003401	.005621	.009161
CENTER OF GRAVITY POSIT					0.900
HUNIZONTAL	-6.031	0.109	18.515	18.820	
VENTICAL		0•1Jž	18.232	18.232	0.167
INCLINATION ANGLE (DEG).	. U.U	1.0	*****	****	-6.5
SEP ARATION ANGLE (DEGREE		****	****	****	****
ABOVE	****				
8ELON	***	****	****	****	****
NOSE # IDTH (# ON FILM).		0.0230	*****	*****	0.0230
NCSE PUSITION (M)					
	6.032	0.222	18.535	16.839	1.013
HCRIZONTAL					0.154
VERTICAL	S. 127	0.134	18.121	18.121	0.104
INPUT NOSE POSITIUN (M)					
		A 1. 14	******	*****	0.993
HCRIZONTAL	0-030	0 +221			
VERTICAL	-0.096	-4.063	*****	*****	-0.067

## SHOT 15 ( 10 NARCH, 1970, ND. 4 )

SANDI DRY: DENSITY:1538 Projectile: Solid Flat					
X-RAY STATION	NG + 1	NU+2	N0+3	NO+4	NO.5
TIME (SECOND) CENTER OF GRAVITY POSIT		.001273	.003323	.005932	•009099
HCRIZONTAL	-0.084	0.058 0.124	0.361 0.130	18.820 18.232	19.125 18.232
INCLINATION ANGLE(DEG). Separation Angle(Degree	0.5 }	0.5	0.0	*****	****
ABOVE	****	****	**** ****	**** ****	****
NOSE WIDTH (M ON FILM). NOSE POSITION (N)	6.0245	¢+0225	0.0225	*****	***** .
HORIZONTAL	U•029 U•126	0.171 6.125	0.474 U.13U	18.839 18.121	19.144 18.121
INPUT NOSE POSITION (M) Horizontal	0 . U 27 -0 . 0 95	00164 -000€0-	-0.068	******	*****

# SHUT 16 ( 11 MARCH, 1976, NO. 1 )

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SANDI DRY. DENSITY:1538. KG/M Prujectile: Solid Flat Nose M	##3: APPRO \$5:0.5451	ACHING VI KG. D#0	ELOCITY:	212. N/S 0.225 M
X-RAY STATION NO.1	NO. 2	N0.3	NQ . 4	NG . S
TIME (SECOND)		•003434	.005308	******
HERIZONTAL	6.057	0.40J J.132	0.748 U.153	19-125
INCLINATION ANGLE(DEG). 2.0	1. Ú	5.0	7.0	*****
SEPERATION ANGLE (DEGREE) ABOVE	4+++	7.0	****	****
NOSE BIDTH (M ON FILM) + 0+0250	**** G•023G	2•0 0•0240	**** 0.0250	**** *****
NOSE PUSITION (M) Hok (ZUNTAL		0.515	0.860 0.167	19.144
VERTICAL		0.512	0.866	******
VENTICAL		-0.072	-0.040	*****

## SHOT 17 ( 1) MARCH. 1976. NO. 2 }

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SANC: DRY. DENSI Projectile: Soli	TY:1538. KG/H++3 D FLAT NOSE MASS	: APPROACHING VE :U.5451 KG. D=U.	LOCITY: 212. M/S 02 M. L=0.225 M
X-RAY STATION	· ···· NO.1	NU.2 NO.3	NO.4 NO.5
TIME (SECOND) CENTER OF GRAVIT	Y POSITION (M)		.006220 .009140
MGRIZONTAL Vertical	••••••• -0•071 ••••••• 0•126	0.146 0.486 6.128 J.135	0.816 1.092 0.143 0.162
INCLINATION ANGLE SEPARATION ANGLE	(DEGREE)	2.3 4.6	9.0 12.5
ABOVE		4.0 5.0 2.5 1.8	9.0 9.5
NOSE WIDTH (M ON NOSE POSITION (M)	FILM). 0.0240	0.0230 0.0230	1.0 0.0230 0.02.J
HERIZONTAL		U.252 0.599	0.928 1.202
INPUT NOSE POSITI	IGN (M)	0.132 0.143	
MCRIZONTAL		∪•256 0•609 -J•685 -0•072	0.941 1.211 -0.052 -0.022
NOSE VEL. Y-COMP. Coefficients of (	CUBIC PULYNOMIAL	STANDARD DEVIAT	ION
0.122UE 00 0.1	7758E 01 -0.6430	E U3 0.6267E 05	/ 0.0013 (M)
NOSE VEL. X-COMP Coefficients of (	(M/S): 201. Cubic Polynomial	178. 139. STANDARD DEVIAT	106. 86. Ion
-0.4578E-03 0.4	2054E 03 -0.1116	E 05 0.3377E 95	/ U.0043 (M)
NOSE VEL. DIRECTI Separation Angle(	LON(DEG) 2.1 (Degree). Relativ	2.1 2.3	3.8 7.8
ABOVE		3.8 3.3	3.8 4.8
BELUW 381888888	****	2.7 3.5	6.2 5.7
C.G. VEL. Y-COMP. COEFFICIENTS OF C			4. 10.
0.11858 00 0.8	457E 01 -6.1420	04 9.1117E 06	/ 0.0002 (M)
C.G. VEL. X-COMP. COEFFICIENTS OF C	(M/S): 200. Weic Pulynumial	178. 139. Standard Deviati	106. 66. Ion
-0.1133E 00 0.2			
PONCELET DRAG COL			
VO = 139. STAND. C.G. VEL. X-COMP.	DEVIA. = 0.6087 (H/S): 210.	(M) 180• 139•	109. 89.

## SHUT 18 ( 11 MARCH. 1976, NO. 3 )

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SAND: DRY, DENSITY:1538. KG/M# Projectile: Sulid Flat Nuse ma	+3; APPROACHING V&LOCITY: 213. M/ SS:0.5451 Kg. D=0.02 M. L=0.225 M	1
X-RAY STATION NG.L	NG+2 NQ+3 NQ+4 ND+5	
CENTER OF GRAVITY PUSITION (M)		
HCRIZUNTAL	v+128 0+133 0+143 0+159	
INCLINATION ANGLE(DEG). 1.6 SFPARATION ANGLE(DEGREE)	2.u 4.6 7.4 8.5	
ABOVE	5.5 5.0 6.5 9.5	
BELOW	4.u 1.4 1.0 1.0 U.0235 J.0230 G.0240 0.0255	
NOSE POSITION (M)	UNUEJJ UNUEJU UNUE4U UNUEJJ	
HERIZONTAL		
VERTICAL	0.132 0.142 0.157 0.175	
HCRIZONTAL Q.U41		
VERTICAL ************* -0+045	-ů.v85 -v.c73 -0.c54 -c.j28	
NOSE VEL. Y-COMP. (M/S): 7.		
COEFFICIENTS OF CUBIC POLYNONIA 0.1217E 00 0.7488E 01 -0.433	AL/STANDARD DEVIATION 2JE UJ U.27878 05 / 0.0014 (M)	
U. 1217E UU U. 7480E UI -0.434	KUE UJ U.2787E VJ / U.UVI4 (M)	
NUSE VEL . X-COMP. (M/S): 200.		
COEFFICIENTS OF CUBIC POLYNOMIA	AL/STANDARD DEVIATION 14e 05  0.3059e 06 / 0.0044 (m)	
-0.30012-02 0.20512 03 -0.111	14E 02 013034E 00 / 010044 (M)	
NOSE VEL . DIRECTION (DEG) 2.1	2.1 2.3 2.9 4.8	
SEPARATION ANGLE(DEGREE), RELAT	5.6 2.7 2.0 5.8	
BELUW	3.9 3.7 5.5 4.7	
C.G. VEL. Y-COMP. (M/S): 7. COEFFICIENTS OF CUBIC POLYNOMIA		
J.1187E 00 0.6934E 01 -0.925	SOE 03 0.7026E 05 / 0.0019 (M)	
C.G. VEL. X-COMP. (M/S): 200. CDEFFICIENTS OF CUBIC POLYNOMIA		
-U.1165E UO 0.205GE 03 -0.410		
PONCELET DRAG COEFF. = 1.624		
VO + 200. STAND. DEVIA. = 0.01		
C.G. VEL. X-COMP. (N/S): 200.	173. 137. 109. 87.	

## SHUT 19 ( 11 MARCH+ 1976: ND+ 4 )

SAND: DRY, DENSITY:1538. KG/M+#3: APPRDACHING V Phojectile: Soliu flat nose massiu.5459 kg. d=0	
X-RAY STATION	NO.4 NO.5
TIME (SECOND)	.005932 .009099
HÜRIZONTAL ******** ~6.068 G*141 0*464 VERTICAL ********* 0*121 0*126 0*136	
INCLINATION ANGLE(DEG). 2.0 3.3 7.0 Separation angle(Degree)	12.0 16.5
	11.0 11.0
BELOW	0.5 Ú.5 0.0220 V.C225
HORIZONTAL	0.893 1.180
VERTICAL	0.174 0.215
HCHIZONTAL	0.899 1.185
VERTICAL	-9.039 0.008
NOSE VEL. Y-COMP. (H/S): 9. 8. 8.	10. 16.
CUEFFICIENTS OF CUBIC PULYNOMIAL/STANDARD DEVIA	TION
0.1232E 00  0.8648E 01 -0.3090E 63  0.6144E 09	5 / 0.0005 (N)
NOSE VEL. X-COMP. (M/S): 205. 180. 14C. CDEFFICIENTS OF CUBIC POL/NUMIAL/STANDARD DEVIA	
0.4456E-02 0.2097E 03 -0.1249E 05 0.3988E 00	
NOSE VEL. DIRECTION (DEG) 2.4 2.6 3.4	5.7 11.0
SEPARATION ANGLE (DEGRÉÉ), RÉLATIVE TU NOSE VELOC Above 4.3 2.9	1177 407 545
BELOW	
C.G. VEL. Y-COMP. (M/S): 6. 5. 5.	7. 14.
COEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIAT 0.11975 00 0.66485 01 -0.74505 03 0.85475 09	
C.G. VEL. X-COMP. (N/S): 2050 180. 1400 Coefficients of Cubic Polynomial/Standard Deviat	TUN
-0 +1084E 06 0+2096E 03 -0+1236E 05 0+3930E 36	/ G.0029 (M)
PONCELET DRAG CULFF. = 1.694	
VD = 205. STANL. (EVIA. = 0.0098 (M) C.G. Vel. X-CUMP. (M/S): 205. 176. 138.	109. 87.
Area rest of fight a function of the fight fight	

## SHUT 20 ( 31 MARCH, 1976, ND. 5 )

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**COMPANY** 

				OCITY: 340. M/S 2 M. L=0.225 M
X-HAY STATIUN	* * * * * * * * * *	ND-1 10-2	N0.3	NO.4 NO.5
TIME (SECOND) Center of Grav			2 .002143 .	003795 .065867
HERIZUNTAL .				0.753 1.054 0.144 Grl44
INCLINATION AN	GLE(DEG).			-1.7 -2.3
SEPARATION ANG Abuve Below			4.0 5.7	2.5 2.5
NOSE WIDTH (M NOSE POSITION	ON FILMIS	0.0225	0.0225 0	•0225 <b>v</b> •C220
HCHIZONTAL . VERTICAL		0.125 0.132	0 • 557 3 • 1 38	0.866 1.167 0.140 0.140
INPUT NOSE POS HCRIZONTAL • VERTICAL •••		0	V.566	0.870 1.169 0.676 -0.077
NOSE VEL. Y-CC Coefficients 0	F CUBIC POL		ARD DEVIATI	DN
MOSE VEL. X-CO	MP. (M/5):	383. 274.	215.	165. 141.
COEFFICIENTS O -0.8087E-02		_YNUMIAL/STAND. -G.3511E (5 (		
NUSE VEL. DIRE				-0.0 0.3
		**** 6.8	5.4	4.2 5.1
BELOW 000000	••••	**** 6.6	4.9	5.3 / 4.2
C.G. VEL # Y-CO	F CUBIC POL	YNONIAL/STANDA	ARD DEVIATIO	)N
0.12391 UU		-U.9086E 03 (		
COEFFICIENTS OF	F CUSIC POL	YNUMIAL/STAND	ARD DEVIATIO	2N
PONCELET DRAG				
C.G. VEL. X-CO	MP. (N/S):	325. 276.	215. 1	68. 133.

SHUT 22 ( 11 MARCH. 1976, NO. 7 )

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SANC: DRY. DENSITY:1538. KG/M4#3; APPROACHING VE Projectile: Solid flat Nose Mass:0.5451 kg. D#0.	
X-RAY STATION NU.1 NU.2 MOL3	N0+4 N0+5
TIME (SECOND)	•003855 •005820
HGRIZGNTAL	0.693 1.049 0.139 0.152
INCLINATION ANGLE(DEG). 1.5 2.5 4.0 Separation angle(degree)	7.0 9.0
ABOVE ************************************	
	0.0246 0.0236
	0.805 1.160 0.152 0.169
INPUT NUSE POSITION (M) HURIZONTAL ************************************	0.799 1.163
	-0.060 -0.042
NOSE VEL. Y-COMP. (M/3): 13. 10. 5. Cuefficients of cubic polynomial/standard deviat	
U-1247E 00 0-1428E 02 -0-3347E 04 0-3809E 06	
NDSE VEL. X-COMP. (M/S): 372. 300. 167. CDEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIAT	
-0.1326E-01 C.J953L 03 -0.7622E 05 0.7393E 07	
NOSE VEL. DIRECTION(DEG) 2.0 1.9 1.7 Separation Angle(degree), relative to nose veloc	2•3 3•1
ABUVE	5.3 5.1
BELÜW ************************************	6.7 7.4
C.G. VEL. Y-COMP. (M/S): 12. 8. 2. COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIAT	3. 13.
0.1215E 00 0.1304E 02 -0.4045E 04 0.4636E 06	
C.G. VEL. X-COMP. (M/S): 372. 300. 167. Coefficients of Cubic Polynumial/Standard Deviat:	
-U.1262E OG U.J951E UJ -O.7611E OS U.7381E UT	
PONCELET DRAG COEFF. = 2.498 VO = 372. Stand. Devia. = 0.0503 (m)	
C.G. VEL. X-CONP. (M/5): 372. 304. 202.	147. 112.

140

## SHOT 23 ( 11 MARCH. 1975. NO. 8 )

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SAND: DRY, DENSITY:1538, KG/M##3; APPRUACHING VELOCITY: Projectile: Solid flat nose mass:g.5451 kg, d#g.02 m, l	
X-RAY STATEGN	ND . 5
TINE (SECOND)	•005789
HERIZUNTAL	
INCLINATION ANGLEIDEG) 200 1.5 2.2 3.7	6.5
SEPARATIUN ANGLÉ(DEGRÉE) Above ++++++++++++++++++++++++++++++++++++	9.0
BELGW ************************************	1.5
NOSE POSITION (A)	
VERTICAL	1 • 1 70 0 • 1 57
INPUT NOSE POSITION (M) HERIZUNTAL	1.175
VENTICAL	-0.054
NOSE VEL. Y-COMP. (M/S): 10. 8. 6. 4.	3.
COEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIATION 0.12462 00 0.10032 02 -0.10692 04 0.54022 05 / 0.00	13 (N)
	131.
COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION	
Ů.4990Ě-U2 0.3155E UJ -3.2726E 05 0.1305E 07 / 0.068	53 (M)
NOSE YEL. DIRECTIC (DEG) 1.8 1.8 1.6 1.5	1.3
SEPARATION ANGLE(DEGREE), RELATIVE TO NOSE VELOCITY Above	3.0
ABOVE	6.7
C.G. VEL. Y-COMP. (M/S): 11. 9. 5. 2. CDEFFICIENTS OF CUBIC POLYNDHIAL/STANDARD LEVIATION	Ŭ•
0.1206E 00 C.1153E 02 - 0.1924E 04 0.1119E 06 / 0.00	18 (M)
C.G. VEL. X-COMP. (M/S): 309. 274. 216. 166. CDEFFICIENTS OF GUBIC PULYNUMIAL/STANDARD DEVIATION	132.
-0.1079E 00 C.3154E 03 -0.272JE 05 0.1309E 07 / 0.00	55 (M)
PUNCELET DRAG COEFF. = 1.723	
VO = 218. STAND. DEVIA. = G.(106 (M) C.G. Vel. X-COMP. (M/S): 326. 279. 218. 171.	135.

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# SHOT 24 ( 12 MARCH. 1976. NO. 1 )

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SAND: DRY. DE Projectile: S	INSIGY:1538. Dilo flat n	KG/N++31 APPRD. 1056 HASS10.5443	ACHING VELOCIT KG. D=G.02 M.	Y: 327. M/S L=0.225 M
X-RAY STATION		NO.1 NO.2	NO.3 NO.4	NL 35
CENTER OF GRA	VITY POSITI	.4460136 .000636	.002019 .0036	16 .605696
HCRIZONTAL		-0.089 U.115 C.126 C.130	0.402 0.71 ucl35 0.13	0 1-024 9 6-145
INCLINATION A SEPARATION AN			2.3 4.0	5.C
BELOW		**** 4.0	5.0 5.6 J.J 2.5	L 🖬 🐪
NOSE WIDTH ( H NOSE POSITION HCRIZONTAL	I (M)	6.02*0 0.0223 6.024 0.228		
VERTICAL INPUT NOSE PO	SITION (M)	0.126 0.134	0.140 C.14	7 0.154
HORIZONTAL Ventical +•	•••••	0.021 0.228 -0.091 -0.083	0.513 0.82 -0.077 -0.06	1 1.136 9 -J.061
NUSE VEL . Y-C			5. 3.	5.
		-0.1841E 04 U		0015 (M)
COLPFICIENTS	OF CUBIC POL	37. 270. Lynumial/standaf	RD DEVIATION	140.
-0,1628E-01	0.3142E 03	-0.2816E 05 0.	150SE 07 / 0.0	0016 (M)
NOSE VEL. DIR Separation an	GLE(DEGREE)	, RELATIVE TO NO	1.3 1.1 DSF VELOCITY	1 . 9
ABOVE Below		**** 4.6 **** 4.4	3.8 2.1 4.2 5.4	1.9 4.6
Colle VELo V-C				3.
COEFFICIENTS	0F CUBIC PUL 0.59425 01	LYNDHIAL/STANDAR -0.8590E UJ 0.	D DEVIATION 7040E 05 / 0.0	002 (M)
		307. 270. YNOMI AL/STANDAN		140.
-0.1293E 00	0.3102E U3	-0.28145 05 0.	1500E 07 / 0.0	0010 (M)
VO N 219. ST C.G. VEL: X-C	AND. DEVIA.	= 0.0048 (M)	219. 173.	135.
MAX •( MIN •(	00 MAXIMUM/ 000126 -0 000712 -0	/MINIMUH COIL VO 01925 - 00352 002771 - 00487	19 005418 9 005418	
COMPUTED NOSE	POSITION AT 0.022	0.494 0.80	OLTAGE (N) 8 1.095	
AT MIN Recorded Coil	PUSITION IM	6.670 0.93 1) 0.486 0.77		
DIFFERENCE BET	WEEN CULL A	ND NUSE AT MAX/ G.GOB G.23	MIN VOLTAGE (M	))
AT MIN	0.194	0.184 Vo21	3 0,230	

## SHUT 25 ( 12 MARCH. 1976. NO. 2 )

SAND: DRY. DENSITY:1538. KG/M##3: APPRUA Projectile, solid flat nose massig.6443	
X-RAY STATIUN	NO.3 NJ.4 ND.5
TIME (SECOND)	.001703 .093009 .004613
	3.5 7.0 9.0
SEP ON ATTOM ANGLE LOEGREE )	
ABOVE	6.5 7.0 10.0 3.0 1.5 0.0
NOSE WIDTH (M ON FILMI: 6.0240 0.0230 NGSE POSITION (M)	
HCRIZONTAL	0.529 0.838 1.136 0.150 0.159 0.171
INPUT NOSE POSITIUN (M) HORIZUNTAL	
VERTICAL	
NOSE VEL. Y-COMP. (M/S): 16. 12. Cuefficients of Cubig Polynomial/Standar	D DEVIATION
C+1305E 00 0+176GE 02 -0+3958E 04 0+	
NOSE VELS X-COMPS (M/S): 381. 335. Crefficients of Cubic Polynumial/Standar	266. 206. 176. D DEVIATION
-0+1964E-01 C+3917E 63 -0++494E 05 0+	J121E 07 / 0.0052 (M)
NUSE VEL. DIRECTION(DEG) 2.4 2.1	1.7 1.7 3.5
SEPARATION ANGLE(JEGREE), RELATIVE TO NO	SE VELOCITY 4.7 1.7 4.5
BELOW	
C.G. VEL. Y-COMP. (M/S): 12. 8. Coefficients of Cubic Polynumial/s7andar	D DEVIATION
0.1305E 00 0.1206E 02 -0.4363E 14 0.	-
C.G. VEL. X-COMP. (M/S): 381. 335. CREFFICIENTS OF CUBIC POLYNOMIAL/STANDAS	266. 207. 177. D DEVIATION
-U.1326E 00 0.3915E US -0.4469E 05 0.	3092E 07 / 0.0054 (M)
PÜNCELET DRAG CUEFF: # 1.684 VD = 266. Stand: Devia: = 0.6039 (M)	-
C.G. VELA X-COMP. (M/S): 389. 334.	266. 211. 169.

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## SHOT 26 ( 12 MARCH. 1976, NO. 3 )

				CITY: 406. M/S M. L=0.225 M
X-RAY STATIO	••••••	NG.1 NO.2	NQ.3 N	0+4 NC+5
CENTER OF GRA	AVITY POSITI	-6.086 0.11	9 3.423 0	03003 •064582 •730 1•035 •153 0•160
INCLINATION A	NGLË (DEG) . IGLE (DE GREE )	ŭ.ŭ 1.U	0.5 -	1.5 -6.0
ABOVE BELOW NOSE WIDTH (J NOSE POSITION		+*** 4.5	4.5	2.5 0.5 3.5 8. 0236 6.0125
HCRIZONTAL VENTICAL	SITION (M)	0.027 0.23 0.134 0.14	2 0.147 0	•843 1•147 •150 0•148
HORIZONTAL VERTICAL ••	•••••	-0.081 -0.07		•844 1•148 •064 -0•067
NOSE VEL • V-C Coefficients 0 • 13322 00	OF CUBIC PU	13. 9. Lynumial/stan -0.35646 04	DARO DEVIATIO	-01. N 9.0013 (M)
NOSE VEL • X-C COEFFICIENTS -0 • 1469E-01	OF CUBIC PU	383. 336. Lynumial/stan -0.4603e 65	DARD DEVLATIO	09. 185. N 0.0015 (m)
NOSE VEL. DIR	ECTIUN (DEG)	1.9 1.6	0.9 -(	7.0 -0.4
ABOVE		• RELATIVE TO ++++ 5.6 ++++ 3.9	3.9	
C.G. VEL. Y-C Coefficients U.1334e 00	OF CUBIC PO	9. 8. Lynumi al/stani -0.1439e 64	6. DARD DEVIATIO 0.1306E 96 /	5. 5. N 0.0007 (M)
C.G. VEL . X-C	OMP. (M/S); Of Cubic Poi		267. 20 DARD DEVIATION	)9. 187. N
PONCELET DRAG VO = 267. ST	COEFF. # 1 AND. DEVIA.	•649 = 0°0025 (M)		
RECORDED TIME	OF MAXIMUM	386. 333. /MINIMUM COIL	VOLTAGE (S)	120 170.
MAX MIN COMPUTED NOSE AT MAX	000099 .0 000625 .0 Pusition A1 0.024	01594 .002 02307 .003 T MAX/MIN COIL 0.509 0	. VOLTAGE (N)	_
AT MIN RECORDED COIL	0+214 POSITION () 0+0	0.689 1. M} D.486 0.	029 1.34 778 1.07	9 76
DIFFERENCE BE At Max At Min	TWEEN COIL / G.624 0.214		X/MIN VOLTAGE 635 0.04 251 0.27	2

### SHUT 27 ( 12 MARCH. 1976. ND. 4 )

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SANG: DRY. DENS	SITY:1530. KG/###3	APPRUACHING VE	LOCITY: 469. M/S
	LID FLAT NUSE MASS		
	NGel		
CENTER OF GRAVI			•
MCRIZONTAL	•••••••• -0.087 •••••••	0.119 0.422 0.152 0.161	
INCLINATION AND		0.0 '-0.5	-2.5 -2.5
SEPARATION ANGL	LE(DEGHEE)	A.5 2.5	1.0 2.5
56LU#		4.5 2.5 6.u 7.0	7.0 7.0
NUSE WIDTH (M C NGSE POSITION (			0.0230 0.0230
HCRIZONTAL			0.845 1.063
VERTICAL	• • • • • • • • • • • • • • • • • • •	v.152 C.160	
HERIZUNTAL	C.023	0.233 0.535 -0.062 -0.053 -	0.846 1.051 -J.048 -0.0-0-0
		••••	
NOSE VEL . Y-CUM	MP. (M/5): 18.	12. 5.	2. 8.
COEFFICIENTS OF 0.1409É ÚG Ú	F CUBIC FOLYNUMIALS G.1959E G2 -G.6071E	04 U.7017E 06	/ G.UU05 (M)
NOSE VEL. X-COM	MF. (M/S): 371.	3363 273.	189. 44.
COEFFICIENTS OF	CUBIC FOLYNOMIAL/ 3.37742 C3 -C.3029E	STANUARD DEVIATI	(UN / G.OG72 (M)
••••••			
NOSE VEL . DIREC	TIGN(DEG) 2.8 Le(Degree), relativ		0.6 5.8
ABOVE		6.6 4.1	4.1 10.8
BELUN		3.5 5.4	3.9 -1.3
C.C. MEL . Y-COM	49. (M/S): 16.	13. 8.	4. 6.
COFFFICIENTS OF	F CUBIC PULYNUMIAL/	STANDARD DEVIATE	(ON
	0.1705E 02 -0.3688E		/ 0+0010 (M)
C.G. VEL . X-COM	AP. (M/S): 371. ".Cuhic Polynumial/	337. 273. Standard Deviati	189. 84. Ion
-U.1269E UU U	.3781E 03 -0.3029E	05 -0.2130E 06	/ 0.0073 (M)
PONCELET DRAG C	06FF. = 1.764		
VD = 273. STAN C.G. VEL. X-COM	D. DEVIA. = 0.0350 4. (M/5): 413.	349. 273.	213. 168.

#### SHUT 29 ( 12 MARCH, 1976, NO. 6 )

		KGZHODI APPRUACHING ( Se massiu 5447 kg. D=(	
X-RAY STATIO	• • • • • • • • • •	NO+1 NO+2 NO+3	N034 N0+5
CENTER OF GRA HERIZONTAL	AVITY PUSITIO		0.701 1.031
INCLINATION A	NGLE(DEG).	-1.0 -1.5 -4.0	-0.5 -6.G
BELOW NOSE WIDTH (M NGSE POSITION	4 ON FILM). U N (M)	**** 9°0 10°630 **** 9°0 10°630	
VERTICAL INPUT NUSE PO	SITIUN (M)	0.025 0.233 0.534 0.108 -0.100 -0.096	
NOSE VEL • Y-C CGEFFICIENTS U • 1114E UU	OF CUBIC PUL	15. 80. Ynumial/standard devia -0.6907e u4 c.7803e 0	TION
COEFFICIENTS	OF CUBIC POLY	387. 337. 265. Ynumial/standard Devia -0.4837e 05 0.3711e 0	TIUN
NUSE VEL . DIR Separation an	GLE(NEGRÉE).	RELATIVE TO NOSE VELO	-1.2 0.5 IC ( TY
		**** 6.9 6.9 **** 6.1 5.6	6.6 7.5 5.7 3.5
C.G. VEL, Y-C COEFFICIENTS U.1132E 00	OF CUBIC POLY	16. Nomial/Standard Devia -0.459je 64 6.33288 0	-14. TION 6 / 0.0003 (M)
COEFFICIENTS	OF CUBIC POLY	J87• J38• 265• Yndmial/standard devia -0•4866e u5 0•367ce 0	TION
PONCELET DRAG	CGEFF. = 1.7	706	
C.G. VEL. X-C	OMP. (M/S): 3	388• 333• 265•	214. 167.
MAX • Min •	006062 .61 606557 .00	MINIMUM COIL VOLTAGE ( 01564 -002864 -00 02276 -003762 -00	04474 05542
AT HAX AT HIN	PÚSITION AT 0.023 0 0.204 0	MAX/MIN CDIL VOLTAGE 5+511 - 0+818 5+689 - 0+995	(M)
RECORDED COIL DIFFLRENCE BE	TWEEN CUIL AN	) )+486	1.076 TAGE (M)
AT MEN		).025 0.04) ).203 9.217	9•051 9•254

## SHUT 30 ( 12 MAHCH, 1976, NO. 7 )

			CHING VELOCITY: 352. M/5 Kg. D=0.02 M. L=0.226 M
X-RAY STATIO	)N	NU.1 NU.2	NG.3 NO.4 NO.5
			.001733 .002879 .004613'
HERIZONTAL	RAVITY PUSITI	-6-092 6-100	
	ANGLE(CEG). NGLE(DEGHÉE)		2+5 6+0 11+0
ABLVE		**** ****	4.5 5.0 6.5
NOSE NIDTH (	M DN FLLMA	\$\$\$\$ \$\$\$\$ ∪•4250 U•0230 (	2•0 1•0 1•5 0•0225 6•0220 C•0203
			0.529 0.039 1.214
INPUT NOSE P	OSITION (M)		0.130 0.143 J.185
HCRIZONTAL VERTICAL 0	• • • • • • • • • • • • • • • • • • •	0.028 0.221 -0.106 -0.097 -	0.529 0.839 1.215 -0.688 -0.074 -C.C33
NOSE VEL . Y-			
		LYNUMIAL/STANDARD -0.45502 04 U+1	019E 07 / 0.0007 (M)
		349. 324.	
		LYNOMIAL/STANDARL -0.2290E 05 0.7	/ DEVIATION /905E 06 / 0.0019 (M)
NOSE VEL . DI			1.7 3.3 11.1
	NGLE(NEGREE);		
BELOW	0 <b>0 0 0 0 0 0 0 0 0 0 0 0 0</b>	**** ****	2.6 3.7 1.4
		12. 7.	
0,1127E 00	0F CUBIC POL 0.1368E 02	LYNGMIAL/STANDARD -0.6010E 04 0.1	159E 07 / 0.0001 (M)
		349. 324. YNDHIAL/STANDARD	
			028E 06 / 0.0019 (M)
PONCELET DRAG			
		349• 321•	286. 244. 206.

## SHUT 31 ( 12 MARCH. 1976. NC. 6 )

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SAND: DRY. DENSITY:1538. KG/M+#J; APPHDACHING VELDCITY: Projectile: Solio Biconic - Massig.4964 kg. D=0.02 m. L=	352. M/S ().226 M
X-HAY STATION ******* NU+1 NO+2 NO+3 NO+4	NO. 5
TIMÉ (SECOND)	.004303
HCRIZONTAL ********* -0*081 0*100 0*406 0*697 Vertical ********** 0*110 0*114 0*123 0*136	1.019 0.190
INCLINATION ANGLE(UEG): 1.5 3.0 5.5 12.0 Separation Angle(degree)	
ABCVE	
NOSE #10TH (N ON FILM). 0.0245 0.0235 0.0225 0.0220 ( NOSE POSITION (M)	0.0210
HCRIZONTAL ******** 0.0041 0.222 0.528 0.816 Ventical ********** 0.113 0.121 0.135 0.162	1+134
INPUT NOSE POSITION (M)	
HCRIZONTAL ******** 0.042 C.221 0.527 0.814 VERTICAL ******** -0.107 -0.098 -0.082 -0.053	1.134 Q.016
NOSE VEL. Y-COMP. (M/S): 13. 13. 18. 31. COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION	62+
0.1117E 00 0.1331E 02 -0.1442E 04 0.1094E 07 / 0.0061	7 (M)
NOSE VEL. X-COMP. (M/S): 346. 323. 201. 239. Cuefficients of cubic polynumial/standard deviation	185.
-0.6516E-C2 0.3524E C3 -U.2206E C5 0.4055E 06 / 0.0026	5 (M)
NOSE VEL: DIRECTION(DEG) 2:1 2:3 3:6 7:4 Separation Angle(degree): Relative to nose velocity	18.4
	9.4
ABOVE **** **** **** 4*1 3*4 BELUM ************************************	****
C.G. VEL. Y-COMP. (M/S): 12. B. B. 20. CUEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIATION	54.
0.1079E 00 0.1393E 02 -0.5891E 04 0.1642E 07 / 0.0012	C M 3
C.G. VEL. X-COMP. (M/S): 346. 323. 282. 241. Coefficients of cubic polynumial/standard deviation	189.
-0.1283E 00 0.3518E 03 -0 2146E 05 0.3850E 06 / 0.0628	6 ( M )
PONCELET DRAG CUEFF. = 0.937	
VO = 346. STAND. DEVIA. = 0.0083 (M) C.G. Vel. X-COMP. (M/S): 346. 319. 278. 244.	209.

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### SHOT 32 ( 1" MARCH. 1976. NO. 1 )

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				ELOCITY: 334. M/S 62 M. L=0.226 M
X-HAY STA	TION	NU-1 N	0.2 NG.3	N0.4 N0.5
	OND)		03672 .001694	.002981 .004288
	TAL			0.739 1.007 0.140 0.183
	UN ANGLE(DEG). N Angle(Degree	2.0	2.5 6.5	15.0 29.0
ABUVE .	3 * * * * * * * * * * * * * * * * * * *	**** *	*** ****	4444 4444 4444
NELUW .		****** *:	*** ****	
NCSE POSI	H (M ON FALM). Tion (M)			
	TAL	0.048 0	•227 0•529	0.857 1.114 0.171 0.242
INPUT NUS	L			0.1/1 0.242
HCRIZON	TAL	ບັບບໍລິພັ ປີ.	.227 0.528	0.861 1.104
VENT ICAL		-0.115 -0	125 -J.Q87	-9.037 0.063
NOSE VEL.	Y-COMP. (M/S) NTS OF CUBIC P	: 10. CHANGALAST	130 220 TANDARD DEVIAT	48+ 69+
0.1002E	00 J.4629E 0	1 0.1641E (	04 J.6218E 06	/ Q.0015 (N)
	X-COMP. (M/S) NTS OF CUBIC P			
	-02 0.3361E 0			
	DIRECTION (DEG			10.3 22.6
	N ANGLE (DEGREE			
		이 가 가 가 다 다 다 다 다 다 다 다 다 다 다 다 다 다 다 다		**** ****
C.G. VEL.	Y-COMP. ( N/5)	: 9.	8. 12.	24. 45.
	NTS OF CUBIC P 00 0.9176E 0			
C.G. VEL.	X-COMP. (M/S)	: 333. 31	6. 281.	232. 178.
	NTS OF CUBIC PI			
-0+1151E	00 0.033696 0	3 -C.1531E 0	)5 -0.4915E 06	/ U00032 (M)
	RAG COEFF. =			
	STANUS DEVIA			216. 210.
veve vikë		, JJJJ 4 JV	·₩₩ 45134	2784 CIV4

## SHOT 33 ( 13 MARCH, 1976, NO. 2 )

SANC: DRY. DENSITY: 1548 PROJECTILE: HOLLUW BICO	• KG/R44)	SS APPRO	ACHING VI	ELOCITY:	400. M/S
PRUSELILES MULLUM DIEM					
X-RAY STATION	NG + 1	ND+2	NG+3	NO • 4	NG • 5
TIME (SECUND)		.000786	.001703	-002e17	•004303
CENTER OF GRAVITY POSIT HORIZONTAL	-0.051	U.190 6.104	0.458 0.000	18.737	19.342
VERTICAL					
INCLINATION ANGLE(LEG).	, - <b>4</b> .J	-11-0	-16.0	0 • U	0.0
SEPARATION ANGLE(DEGREE ABOVE	****	1.0	****	****	****
	****	9.0		4844	***
NOSE WIDTH (M ON FILM).	4.0250	0.0236	0.0230	******	******
NOSE POSITION (M)		1 200	0.556	18.839	19.144
HURIZONTAL	0.051	Ú • 29Ú			
VERTICAL	じゅしなち	0.084	0.032	18.121	18.121
INPUT NOSE PUSITION (M)					•
HCRIZONTAL	0.054	0.299	0.560	*****	*****
				******	*****
VERTICAL	-0.130	-6.140	-0.200	******	*****

## SHOT 34 ( 13 MARCH. 1976, NO. 3 )

SAND: DRY, DENSITY:1538, KG/MS Prujectile: Hollow Bicunic Ma	131 APPRGAC 5510-3434 K	HING VEL G. D=0.6	CCITY: 2 M, L=	413. M/3 0.226 M
X-RAY STATION	N0+2	NO.3	NO = 4	NQ . 5
TIME (SECOND)	9 .000774 ·	001703 •	002811	.004319
CENTER OF GRAVITY PUSITION (M) HURIZONTAL			8.737 8.121	19.042 18.121
INCLINATION ANGLE (DEG) - 1.7	-5.0	-8.5	0.0	0
SEPARATION ANGLE (DEGREE) Above	1.J 6.5 0.0235 J	1.0 6.0 .0250	**** **** ****	**** **** *****
NOSE POSITION (A) HCRIZONTAL			8.839	19.144 18.121
INPUT NOSE POSITIUN (A) HERIZONTAL			*****	*****

### SHGT 35 ( 13 MARCH, 1976, NO. 4 )

				LOCITY: 302. M/S 02 M. L=G.226 M
X-RAY STATION	••••••	NU-1 NG-2	N0+3	N0+4 N0+5
TIME (SECOND) Center of Grav	VITY POSITI	+600139 +006796 [On (M]	. 91919	003514 .005418
HERIZONTAL Vertical		-0.063 0.146 0.096 0.192		
INCLINATION AN		1.5 4.0	6.0	7.5 14.0
ABOVE	*********	***	5.0	7.0 ****
HELOW Nose fidth (n		*### #### U=D245 U=U230	1.G J.0240 (	4.C ++++ ).U250 V.U31.
NOSE POSITION		V VZ45 V V VZJU		/+VZ34 V+V315
HCRIZONTAL		0.039 0.247	0.541	0.876 1.191
VERTICAL		0.095	0.126	0.157 C.215
HORIZONTAL		0.639 0.250	0.543	0.885 1.218
VERTICAL		-0.125 -0.111	-0.092 -	0.053 0.049
		16. 15.		
		LYNUMIAL/STANDA		
U + 9659E-01	V.10335 02	-4.1343E 04 0	11434YE 40	/ 0.0008 (M)
		335. 295.		
		LYNUMIAL/STANDA		
-0.10806-02	0.34372 03	-0+3338E 05 0	1991E 07	/ 0.0038 (M)
NOSE VEL. GIRE	ECTION(DEG)	2.7 2.9 . RELATIVE TU N	3.9	7+2 14+3
		**** ****		
BELDW				
C.G. VEL. Y-CO	3MP. (M/S):	8. 10.	14.	21. 29.
COEFFICIENTS G	OF CUBIC PO	LYNOMIAL/STANDA	RD DEVIATI	ON
0.9512E-01	0.7102E U1	0.1731E 04 0	.J738E 05	/ 0+0097 (M)
		335. 295.		
		LYNOMIAL/STANDA		
-0.10922 00	4.34442 03	~0.3368E 05 0	.ZU42E U7	V 00038 (M)
PONCELET DRAG	COEFF. = 1	.011		
VO # 238. STA	ND. DEVIA.	= 0.0021 (M) 340. 294.	~ 7 0	107 140
L.U. VEL. X-CO	MMA (W121:	그국상소 《상학》	Z30.	10/0 1424

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## SHOT 36 ( 13 MARCH. 1976, NO. 6 )

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SANC: WET. CENSITY:2050. KG/M## Projectile: Solid Flat Nose Mas	3; APPKO 5;450	ACHING VE KG. D=6	02 X L	336. M/S =0.225 M
X-RAY STATION	N0.2	N0.3	NÜ+4	ND . 5
TIME (SECOND)	••••0836	.062217	-003901	.006094
CENTER OF GRAVITY POSITION (M) HCRIZONTAL	0.142	0.490 U.106	0.844 J.166	19.031 18:121
INCLINATION ANGLE(DEG) - 0.0	U.U	0.5	0.5	0.0
SEPARATION ANGLEVDEGREED ABOVE	17.0	7.5	5.0	****
BELOW	17.0 6.0236	8.5 v.0230	8.C Q.06.30	**** *****
NOSE MOSITION (M) HCRIZONTAL	0.255	0.693 0.107	0.957 0.109	19.144
VENTICAL AND	J.259	0.613	0.975	*****
VERTICAL	-0.125	-0.114	-0.111	*****

## SHUT 37 ( 24 MARCH+ 1976+ ND+ 1 )

		KG/M##3; APPROACHIN J56 Mass:J.5449 Kg. (	G VELOCITY: 336. M/S D=0.02 M, L=0.225 M
X-RAY STATION		NU.1 ND.2 NO.	3 NO.4 NO.5
TIME (SECUND) Center of Gra	ATTA DUSITI	.000136 .000805 .002	074 .003560 .005365
HCRIZONTAL		-0.079 0.126 0.49 0.095 0.100 0.10	
INCLINATION A SEPARATION AN		1.0 1.0 2.0	6 2.5 4.0
ABOVE		20.6 **** 11.6 20.0 **** 9.6	0 6.5 4.0
NOSE WIDTH (M NGSE POSITION	I ON FILM). ( ( (m)		
HER IZUNTAL		6.034 0.239 0.50 0.097 0.102 0.10	63 0.888 1.203 06 0.166 0.113
INPUT NOSE PO HURIZONTAL	SITION (M)	4.0033 0.240 0.50 0.127 -0.119 -0.13	58 J.896 1.212 15 -J.113 -C.107
NOSE VEL. Y-C CUEFFICIENTS	OF CUBIC POL	YNUMIAL/STANDARD DEV	VIATION
		-U.2090E 64 0.29548	
COEFFICIENTS	OF CUBIC POL	.311. 283. 238. .Ynumial/standard den -g.2123e 05 g.7992e	195. 157. Viation
-0.55565-02			
NOSE VEL. DIR Separation an	GLE(DEGREE).	RELATIVE TO NOSE VE	0.3 2.1 ELOCITY
A80VE	• • • • • • • • • • • •	20.5 **** 9.4	5+8 7+6
C.G. VEL. Y-C	OMB . (	7. 4. 1.	0. 4.
	OF CUBIC POL	YNUMIAL/STANGARD DEV -0.2613E 04 0.2755E	/IATION
C.G. VEL. X-C	OMP. (M/S):	311. 284. 238.	195. 157.
COEFFICIENTS -U.1183E 00	OF CUBIC PUL 0.31616 03	YNUMIAL/STANDARD DEV -C.2123E G5 U.8008E	/IATIUN   Og / Ceuc44 (M)
PONCELET DRAG			
		322. 287. 238.	199. 166.
		MINIMUM COIL VOLTAGE	
MAX + MIN +	030133 .u 006697 .u	U1811 .063229 62666 .004235 Max/MIN CUIL VOLTAG	. 00 50 68
AT MAX	0.030	U-502 U-821 U-702 I-013	
RECORDED COIL	PÚSITIÚN (M U.U	1 .456 0.794	1.092
AT MAX		ND NOSE AT MAX/MIN V U.016 U.U27	OLTAGE (H) 0+035
AT MIN	0.245	0.216 0.219	0.218

#### SHUT 30 ( 24 MARCH, 1976, NO. 2 )

SAND: WET, DENSITY:2050, KG/M3#3; APPROACHING VELOCITY Projectile: Solid flat nose mass:3.5446 kg, D=0.02 m.	: 333. M/S L¤G.225 M
K-RAY STATION NO.1 NO.2 NO.3 NO.4	NDc 5
TIME (SECOND)	0.005418
HCH1ZUNTAL ********* -G*076 0+131 0*461 0*788 Ventical ********** 0*089 0*097 0*106 0*100	
INCLINATION ANGLE(DEG): 2.5 1.0 1.5 2.6 Separation anglé(degrée)	3.5
ABCVÉ ************************************	
NOSE WIDTH (M ON FILM). 0.0250 0.0236 0.6230 0.0230 NOSE POSITION (M)	
HURIZUNTAL	
VERTICAL ********** 4*094 4*099 4*103 0*104 Input Nose Positiun (M)	6 - 17
HCRIZUNTAL	
NUSE VEL. Y-COMP. (M/S): 10. 6. 1. 2. Cuefficients of cubic Polynomial/Standard Deviation	14.
	006 (N)
NOSE VEL, X-COMP, (M'S); 319, 268, 240, 198. Cuefficients of Cubic Polynomial/Standard Seviation	168.
-0.4028E-02 0.3250E 03 -C.2426E 05 0.1237E 07 / 0.00	019 (M)
NOSE VELS DIRECTION(DEG) 1.9 1.1 0.2 0.6 Separation angle(degree). Relative to nose velocity	4,6
ABOVE ************************************	11-1
8ELOW ************************************	3.9
C.G. VEL. Y-COMP. (M/S): 14. 7. C. C. O. Cgefficients of cubic polyngmial/standard deviation	13.
0.8757E-01 0.154 JE 02 -0.5785E 04 0.6853E 06 / 0.00	007 (M)
C.G. VEL. X-COMP. (M/5): 319. 288. 240. 198. CJEFFICIENTS OF CUBIC POLYNOMIAL/STANDARD DEVIATION	169.
-U-1169E 00 0.3249E 63 -G.2421E 05 U.1207E 07 / 0.00	19 (M)
PONCELET DRAG CORFF. = 0.937 Vo = 243. Stand. Devia. = 0.0021 (n)	
C.G. VEL. X-COMF. (M/S): 324. 289. 240. 200.	166.

SHOT 39 ( 24 MAHCH, 1976, NO. 3 )

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SAND: WET, DENSITY:2050, KG/MODI; APPROACHIN Projectile: Solij Step-Tier Mass:0.5640 Kg,	
X-RAY STATION	d <sup>1</sup> 4 NDaJ
	1904 348 .065108 25 0.774 1.121 115 0.123 0.152
SEPARATION ANGLE(DEGREE) Above	5 9.5 10.0 5 13.0 ****
BELOW	
NCRIZONTAL +++++++ 0+032 4+238 4+5	546 C+894 1+241 28 0+144 G+174
HGRIZONTAL	
NOSE VEL. Y-COMP. (N/S): 11. 10. 10 Coefficients of Cubic Polynumial/standard de U.1087E 00 0.1186 02 -0.1478E 04 0.3259	VIATION
NOSE VEL: X-COMP. (M/S): 318. 297. 253 COEFFICIENTS OF CUBIC POLYMONIAL/STANDARD DE -0.9964E-02 G.3227E G3 -0.1043E G5 0.2335	VIATION
NOSE VEL . DIRECTION (DEG) 200 1.9 2.	1 3,3 7.3
SEPĂRATION ANGLE(LEGREE), ŘELATIVE TO NUSE V Above ••••••••••••••••••••••••••••••••••••	1 6.5 ****
C.G. VEL. Y-COMP. (M/S): 2J. B. 4 COEFFICIENTS OF CUBIC POLYNOMIAL/STANDARD DE U.1021E 00 0.1446E 02 -0.5314E 04 C.8614	
C.G. VEL. X-COMP. (M/5): 318. 297. 263 COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DE	• 221• 173• VIATIUN
-0.1317E 00 0.3221E 43 -0.1590E 05 0.1651 PONCELET DRAG CUEFF. # 0.720	E 46 / G.O006 (M)
VO = 318. STANU. DEVIA. # 0.0092 (M) C.G. VEL. X-COMP. (M/S)2 318. 292. 258	₂ 224∎ 193•

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#### SHOT 49 ( 25 NARCH, 1976, ND. 2 )

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SAND: WET, DENSITY:2050, KG/MAA Projectile: Solid Step-tier Mas		
X-RAY STATION NEUGODOLO NO.1	NQ.2 NO-3	N0+4 N0+5
TIME (SECOND)	na***** "201850	.003251 .004836
HCRIZUNTAL	18.216 0.410	0.735 1.041
VERTICAL	18.241 00102	0.105 0.122
INCLINATION ANGLE (DEG) - 3.0	#÷5## 4.0	8.5 9.2
SEPARATION ANGLE(DEGREE)		
ABOVE	**** 9+0	16.0 12.0
86400 *****	#### <b>4</b> 9 <b>5</b>	5+û 1+5
NUSE WIDTH (M CN FILM). 0.0260	≠≠≈≠≠≠ 0°°61539	0.0230 0.0230
NOSE POSITION (M)		
HLRIZÜNTAL ******** 0** 0**34	18-231 0-540	Q.855 1.161
VERTICAL	18.121 v.110	0.123 0.142
INPUT NOSE POSITION (P)		
HORIZONTAL	****** 0.541	0.858 1.164
VERTICAL	****** .0.110	-0.095 -0.073
ARMITEUR BEABBOJEEEEE _AE795	****** .A411A	-78423 -01013

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#### SHUT 41 ( 25 MARCH. 1976, ND. 3 )

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SAND: WET. DENSITY:2050. KG/M##3; APPROACHING VELOCITY: 320. M/5 Projectile: Solid Step-tier Mass:0.5645 kg. D=0.02 M. L=0.238 M X-RAY STATION ..... NC • 1 NO.2 NO.3 NO.4 N0.5 INCLINATION ANGLE(DEG). Separation Angle(Degree) 1.0 \*\*\*\* 1.0 1.0 . 2.0 ABOVE ............ \*\*\*\* \*\*\*\* 5.0 11.0 13.0 13. BELOW \*\*\*\* \*\*\*\* 4.0 10.0 \*\*\*\*\* U. ÚŽĴJ 0.0230 0. 550 0.889 1.228 HCRIZONTAL ....... 0.031 18.231 0.040 18.121 0.102 0-107 0.116 0.553 0.897 0.029 \*\*\*\*\*\* 1.241 -9.103 \*\*\*\*\* -0.139 --.120 -0.114

## SHOT 42 ( 25 MARCH, 1976, NO. 6 )

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SANC: WET. DENSITY:2080, KG/Mt4: Projectile: Sulid Step-tier Mas:				
X-RAY STATION	NO+ 2	N0.3	ND.4	NU . 5
TIME (SECOND)	******	.004536	.005263	.067740
HCRIZONTAL	18.210	10.511	0.709	0.968
/ERTICAL	16+541	18.241	0.096	0.107
INCLINATION ANGLE(LEG). U+0 SEPARATION ANGLE(DEGREE)	****	****4	-4.5	-8.0
ABUVE		****	****	****
BELOW	****	****	40.40	****
NOSE WIDTH (M GN FILM), 0.0260 NGSE POSITION (H)	*****	*****	0.0230	0.0230
HCRIZONTAL 0.031	18.231	18.535	0.831	1.089
VERTICAL	18.121	18.121	0.085	0.090
INPUT NOSE POSITION (M)		100161	44463	<b>M + M M</b>
	******			
HERIZONTAL	******	******	0.830	1.081
VERTICAL	<b>大卡女法亲</b> 教	******	-0.138	-0.133

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## SHUT 43 ( 25 MARCH+ 1976+ NO+ 5 )

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SANCI WET, DENSITY:20 Projectile: Solid Ste	50. KG/M##JI APPKU P-TIEñ Massju.5048	ACHING VELOCITY: Kg, D=0.02 M. L	210. M/5 #0.238 M
X-RAY STATION	•• NO.1 NO.2	ND.3 NO.4	N0.5
TIME (SECOND)		au03198 .005461	\$4180°
HORIZONTAL	18.019 18.210	0.437 0.725 0.698 6.101	0.997
INCLINATION ANGLE(DEG		-1-9 -5-0	+7.5
SEPARATION ANGLE DE GR		3.5 4444	****
BELOW	■■ 你追你来 食產業者	5.0 #### Us0230 0.0230	\$\$ 0 = 0 ≥ .30
NOSE POSITION (N) HCRIZONTAL	•••••••••••••••••••••••••••••••••••••••	0.559 0.847	1.118
VERTICAL	. 18.121 16.121	0.096 0.090	0.089
HCRIZONTAL	- + + + + + + + + + + + + + + + + + + +	0.563 0.848	1.114

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#### SHUT 44 ( 25 NARCH. 1976. NO. 6 )

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SANC: WET. DENSITY:2050. KG/M443: APPROACHING VELOCITY: 213. M/S PROJECTIE: SOLID STEP-TIER MASS:0.0659 kg. D=0.02 M. L=0.238 M NOc4 NG.2 NU<sub>1</sub>3 N0.5 X-RAY STATION ..... NU . I .006173 \*\*\*\*\*\*\* .JC3297 .005650 .048421 1.004 0.732 18.210 3.438 0.693 VERTICAL ........ 18.121 10.241 9.0398 G.104 INCLINATION ANGLE(DEG). SEPARATION ANGLE(DEGREE) AEGVE BELOW NUSE WIDTH (M ON FILM). NOSE MOSITION (M) C ... \*\*\*\*\* 1.0 -2.5 -4.5 4.0 Q. Q 3.0 .... \*\*\*\* 5.0 6.C 0.023C \*\*\*\* \*\*\*\* 4.Ú 6.U230 ...... \*\*\*\*\*\* 0.854 1.120 HLAIZONTAL ...... 18.040 0.560 18-231 0.095 0.693 18.121 10.121 0.096 0.856 1.123 \*\*\*\*\* 2.564 \*\*\*\*\* -0.127 -0.130 -0 128 . .

### SHUT 45 ( US APR.L. 1976, NO. 1 )

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SANC: WET. DENSITY:2050. KG/M##3; Projectile: Solid Step-Tier Mass:	: APPROACHING VELOCITY: 210. M/S 80.5662 kg. D=0.02 m. L=0.238 m
X-RAY STATION NO.1	NG+2 NO+3 NO+4 NO+5
TIME (SECOND)	
NCRIZUNTAL	C.110 0.442 0.742 1.013 U.110 0.115 0.118 (.122
SEPARATION ANGLEIUEGREE)	1.0 2.5 5.6 6.0
ABUVE	7•0 4•5 8•5 10•0 5∗5 3•5 5•0 1•0
HURIZONTAL	
INPUT NÜSE POSITIUN (M) HURIZGNTAL	
NOSE VEL. Y-COMP. (N/S): 6. COEFFICIENTS OF CUBIC POLYNOMIAL/	
	93 U-8368E 04 / C.J003 (M)
NOSE VEL. X-COMP. (M/S): 198. Cuefficients of Lubic Polynumial/	
-0.3201E-02 0.2010E U3 -6.5825E	04 G.2332E 06 / 0.0G11 (M)
NUSE VEL. DIRECTION(DEG) 1.6 Separation Angle(Gegree). Relative	
AHOVE	
C.G. VEL. Y-CUMP. (M/S): 6. CDEFFICIENTS OF CUDIC PULYNUMIAL/S	
0.1029E 00 0.6875E 01 -6.1156E	
C.G. VEL. X-COMP. (M/5): 198. 1 COEFFICIENTS OF CUBIC POLYNOMIAL/S	STANUARD DEVIATION
-0.12522 00 0.2017E 03 -0.4772E	04 0.2296E 06 / 0.CO12 (M)
PUNCELET DRAG CUEFF. = 1.241 VO = 145. STAND. DEVIA. = C.UIJ4	
C.G. VEL. X-COMP. (N/5): 212. 1	183. 145. 117. 95.

## SHUT 46 ( US APRIL: 1976, NO. 2 )

SANG: WET. DENSITY:2080. Projectile: A.M.H.		J: AP(?RD 5:6.0311			251. M/S =J.J70 M
X-RAY STATION	NQ+1	NO• 2	N0.3	NQ.4	NO.5
TIME (SECOND)	.U00699 UN (A)	•001901	.005000	.019317	.063951
HERIZONTAL	4.024	J.152	18.535	18.839	19.144
MERTICAL	0.111	0.118	18.121	18,121	18.121
INCLINATION ANGLE(DEG). Separation Angle(Degree)	3.5	13.0	****	****	*****
ABUVE	****		****	****	****
#ELUW **************	****	1.0	***	****	****
NOSE WIDTH (M ON FILM). ( NCSE POSITION (M)	6.0250	0.0230	*****	******	*****
HCRIZONTAL	0.024	U.152	18.535	18.839	19,144
VERTICAL "	<b>U-111</b>	6.118	18.121	18.121	18.121
INPUT NOSE POSITION (M)					
HCRIZUNIAL	0.021	6.141	******	*****	******
VENTICAL	-4.116	-0.101	******	******	Q#\$P\$\$

# SHUT 47 ( US APRIL: 1978, NO. 3 )

SANDI WET. DENSITY:2050. KG/M483 Projectile: A.M.N. Mass	5:3.00000 5:3.0012	KG. D=V	ELOCITY: 02 M. L.	552. M/S 10,070 N
X-RAY STATION NU.L	NQ+2	NQ.3	NQ • A	ND.5
TIME (SECUND)	.006970	.005000	-019192	•063951
CENTER OF GRAVITY POSITION (M) HERIZONTAL	G.254 G.121	0,398 9,164	18.839 18.121	19,144 18,121
INCLINATION ANGLE (DEG) . 8.0	20.0	152.0	*****	*****
SEP ARATION ANGLE (DEGHEE) ABOVE	24.U .	****	****	****
NOSE SIDTM (M ON FILM). 0.6256 NOSE PUSITION (M)	0.1235	ù• V 22U	*****	****
HCRIZUNTAL	G.264 J.121	<b>398 • ت</b> 164 • ب	18.839 18.121	19.144
INPUT NOSE PUSITION (A) HORIZONTAL	0.270	0.385 -0.051	******	*****
VERTICAL	-6.698	-0.091	*****	*****

#### SHOT 48 ( 06 APRIL. 1976: ND. 4 )

SAND: WEI. DENSITY:2050. KG/MO+: Prujectile: A.M.M. Mass	31 APPROA 5:0.0311 1	CHING VE KG• D≖U•	LOCITY:	538. M/S 0.070 M
X-RAY STATION NO.1	NU . 2	N0 . 3	NO • 4	NO.5
TIME (SECOND)	• 001200	.008044	.019969	•100000
CENTER OF GRAVITY POSITION (M) HCRIZONTAL	18.231 19.121	0.423 0.076	18.839 18.121	19.144
INCLINATION ANGLE(DEG)1.0		****	*****	*****
SEPARATION ANGLEIDEGREED Above		****	****	****
62100	****	**** 0=0230	**** ******	****
NOSE PUSITION (M) Hoxizontal	18.231 18.121	u • 423 0 • 6 76	18.839 18.121	19.144 13.121
INPUT NOSE POSITION (M) HCRIZUNTAL	*****	0.406	*****	*****
VERTICAL	******	-ú • 1 5u	*****	****

#### SHUT 49 ( U9 APRIL, 1976, NO. 1 )

SANC: WET. D Projectile:	DENSITY:2080 . KG/N## Solid Step-tier Mas	3: APPROACHING V 5:J+566J KG+ D≠0	ELOCITY: 327. M/S .U2 M. L=U.238 M
X-RAY STATIO	IN	N0+2 N0+3	N0.4 N0.5
TIME (SECOND CENTER OF GR	))	.000780 .001876	.003220 .004768
HCRIZONTAL	0.093	0.117 0.426 0.098 6.105	
INCLINATION A	ANGLE(DEG). 0.5 NGLE(DEGREE)	0•5 C•5	•••
ABOVE Below		11.0 10.5 11.0 7.5	9.0 8.0 7.0 10.0
NOSE POSITIO		0.0230 0.0235	0.0230 0.0250
VENTICAL .	••••••••••••••••••••••••••••••••••••••		
HCRIZONTAL VERTICAL	USITION (M) 	0.241 0.550	
			-08109 -04105
COEFFICIENTS	COMP. (M/S): 8. OF CUBIC PULYNOMIAL	STANDARD DEVIA	
	0.8645E 91 -0.1072		
COEFFICIENTS	COMP. (M/S): 334. Of Cubic Pulynumial 0.3401e 03 -g.2382	STANDARD DEVIAT	223• 192• Tion
	0004012 00 -002002		/ / 00010 (H)
SEPARATION AN	RECTIUN(DEG) 1.4 NGLE(DEGREE), RELATI	I.J I.I VE TO HOSE VELOG	0.8 C.4 ;17y
ABOVE Beluw	••••••••••• **** ••••••••••	11.8 11.1 10.2 6.9	9•3 8•9 6•7 9•1
C.G. VFL . Y-C	COMP. (N/S): 9.	7. 5.	<b>4. 4.</b>
	OF CUBIC POLYNOMIAL	STANDARD DEVIAT	ION
	COMP. (M/S): 334.	305. 263.	223. 193.
-0.1351E 00	OF CUBIC POLYNOMIAL G.34U1E ú3 -G.2384	E 05 0.1171E 07	'ION ' / G+0017 (m)
PONCELET DRAG	G CUEFF. = 0.822 Fand. Devia. = 0.002	5 (M)	
C.G. VEL. X-C	COMP. (M/S): 334.		225. 194.
RECORDED TIME	OF MAXIMUN/MINIMUN	COLL VOLTAGE (S	
MAX . NIN . COMPUTED NCSE	000155 .001894 000705 .002773 Pusitiun at Max/MI	•003137 •00 ####### • •00 N CDIL VOLTAGE (	4658 5621 M)
AT MAX AT MIN	0.039 0.554 0.215 0.772	Ú.855 1	•172 •353
	. POSITION (M) Ú-ú ú-487	0.792	•085
DIFFERENCE BE AT MAX AT MIN	TWEEN CUIL AND NOSE G.U39 G.067 0.215 G.265	0.063 0	•086
	VINIU VIE03		• 257

## SHUT 50 ( 09 APRIL: 1976, NO. 2 )

A STATISTICS

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SAND: #ET. DENSITY:2050. KG/N##31 APPROACHING VELOCITY: 395. Projectile: Solid Step-tier Mass:0.5560 kg. D=0.02 m. L=0.23	M/S 0 M
X-RAY STATION NO.1 NU.2 NO.3 NO.4 HO.	5
TIME (SECOND)	19
INCLINATION ANGLE(DEG). 1.5 1.5 2.0 3.5 7.5 SEPARATION ANGLE(DEGREE) ABUVE	5
BELOJ ************************************	30
HCR120NTAL	28 55
NOSE VEL. Y-COMP. (M/S): 12. 10. 8. 6. 6. 6. CUEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION	
J.9567E-01 U.1207E U2 -U.1727E U4 U.1531E 06 / U.OCO3 (M)	
NOSE VEL: X-COMP: (M/S]: J92: 368: 322: 263: 183: CDEFFICIENTS OF CUGIC PULYNOMIAL/STANDARD DEVIATION -U:1501E-01 U:3971E U:3 -0:2302E 05 -0:2949E 06 / 0:0021 (M)	
NOSE VEL. DIRECTION(DEG) 1.7 1.6 1.4 1.3 1.6 SEP ARATION ANGLE(DEGREE), RELATIVE TO NOSE VELOCITY	-
ABOVE	
C.G. VEL. Y-COMP. (M/S): 12. 10. 6. 22. Coefficients of cubic polynomial/standard deviation 0.9248E-01 0.1230E 02 -0.2173E 04 0.8346E 05 / 0.0002 (M)	
C.G. VEL. X-COMP. (M/5): 392. 368. 322. 263. 184. CDEFFICIENTS OF CUBIC POLYNOMIAL/STANDARD DEVIATION -0.137je 03 C.3972e 03 -0.2304e 05 -0.2785e 06 / 0.0023 (M)	
PONCELET DRAG CUEFF. = 0.788 VD = 392. Stand. UEVIA. = 0.0156 {m} C.G. VEL. X-COMP. (M/S): 392. 359. 312. 267. 226.	,
RECORDED TIME OF MAXIMUM/MINIMUM COLL V/LTAGE (5) Max ====================================	
MIN       +000606       +002236       +003137       +004665         COMPUTED NOSE PUSITION AT MAX/MIN CULL VOLTAGE (N)         AT MAX       0+034       0+550       0+825       1+146         AT MIN       0+217       0+755       0+995       1+307	
RECORDED COIL POSITION (A)	
DIFFERENCE BETWEEN COIL AND NOSE AT MAX/MIN VOLTAGE (M) AT MAX 0.034 0.063 0.033 0.060 AT MIN 0.217 0.208 0.203 0.221	

### SHOT 51 ( US APRIL: 1976, NO. 3 )

SANC: WET. DENSITY:2050. KG/M44 Frojectile: Sulio Step-tier Mas	3; APPRO	KG. D=U	ELOCITY: •02 M, L	402; M/S =0.238 M
X-RAY STATION NU-1	N0+ 2	HO.3	NQ.4	NU . 5
TIME (SECOND)		+001591	+002495	•0C3848
CENTER OF GRAVITY PUSITION (M) HCKIZONTAL	Ú.111 Ú.102	18.514	0.667	19.124 18.241
INCLINATION ANGLE(DEG). ****	1.5	****	6.0	*****
SEPARATION ANGLE (DE GREE ) Abuve	14.0	****	14.0	****
BELDU 000000000000000000000000000000000000	11.0 0.0230	**** *****	6.5 0.0230	#\$## <b>#####</b> #
NOSE POSITION (M) MGRIZONTAL	0.233	18,535	0.788 V:124	19.144 18.121
VERTICAL	ú•1ú5 0•234	18.121	0.781	******
HERIZONTAL	-0.116	*****	-0.094	*****

## SHUT 52 ( 10 APRIL: 1976, NO. 1 )

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SAND: DRY, DENSITY:1538, KG/M##3: APPRDACHING VELOCITY Projectile: Solid Stef-tier Mass:0.5682 kg, D#0.02 M.	
X-RAY STATION NO.1 NG.2 NO.3 NO.4	NG . 5
TIME (SECOND) ******* *900%66 *JG1217 *003251 *00547 Center of gravity position (M)	71 eC08136
HENIZUNTAL	
INCLINATION ANGLE(UEG). 0.0 0.0 0.0 -1.0 -3.5 Separation Angle(Degree)	-5.5
ABOVE	
NOSE NIDTH (M ON FILM). 0.0250 0.0230 0.0230 0.0230 NUSE POSITION (M)	
HCRIZONTAL	
VERTICAL	
HURIZONTAL ******** 0*425 0*231 0*552 0*828 Vertical ************************************	1.089 -0.102
NUSE VEL. Y-COMP. (M/S): 7. 5. 11.	-1 -
CUEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION	• •
NCSE VEL. X-CGMP. (M/S): 206. 181. 140. 110.	94.
COEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIATION -0.9796E-02 0.2107E 03 -0.1327E 05 0.5007E 06 / 0.0	044 (N)
NOSE VEL. DIRECTION(DEG) 2.0 1.6 0.6 -0.3 Separation angle(degree). Relative to nose velocity	-0.5
ABOVE	5.0 -2.4
C.G. VEL. Y-COMP. (M/S): 6. 5. 3. 2. CDEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION	-0.
0.1658E 00 0.6473E 01 -0.4980E 03 0.6752E C4 / 0.0	026 (M)
C.G. VEL. X-COMP. (M/S): 206. 181. 140. 111. COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION	94.
-0.1318E 00 0.2106E 03 -0.1326E 05 0.5005E 06 / 0.00	045 (M)
PONCELET DRAG COEFF. = 1.011 Vo = 146. stanus devias = 0.0020 (m)	
C.G. VEL. X-COMF. (M/S); 210. 180. 140. 113.	92 •

#### SHUT 5J ( 10 APRIL: 1976, NO. 2 )

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SAND: DRY. DENSITY:1538. KG/N##J; APPRUACHING VELOCITY: 212. M/S PROJECTILE: SOLID STEP-TIER MASS:0.5662 kg. D=0.402 M. L=0.238 M N0.3 NOAC NO.5 NO.L N0+2 X-RAY STATION ..... .000186 .001227 .003233 .005575 .008168 6.108 0.429 18.819 0.960 VERTICAL ....... 0.146 0.114 6.116 18-241 0.091 INCLINATION ANGLE(DEG). Separation angle(degree) Agove ...... 4.5 -14.5 -5.0 \*\*\*\* 1.4 \*\*\*\* 3.4 2+0 \*\*\*\* \*\*\*\* 7.0 v.0235 \*\*\*\* NUSE WIDTH (M GN FILM). 0.0240 NOSE POSITION (M) 3.0 \*\*\*\* \*\*\*\*\* 0.0250 J . 551 J . 105 1.078 HGRIZONTAL ....... 0.027 6.230 18.839 4.115 4.100 18-151 0.061 4.425 0.230 0.554 \*\*\*\*\* 1.062 -0.1.5 VERTICAL ........ -0.113 \*\*\*\* -0.173 -0.116

## SHUT 54 | 10 APRIL: 1976. NO. 3 3

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# SHUT 55 ( 10 APH 54, 1976, NO. 4 )

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			ING VELOCITY: 206. M/S . C=0.02 M, L=0.238 M
X-RAY STATION	N	. L ND. 2 N	0.3 NO.4 NO.5
	VITY PUSITION		G3283 .005814 .008618
	** ** ** **** -0 *	095 0+100 0	•411 0•713 0•975 •118 0•122 0•115
INCLINATION A	VGLE ( DE GHEE )		2.0 -5.0 -5.5
BELOW		** 5.0	
NOSE PUSITION HERIZONTAL			0230 0.0220 0.022C .533 0.834 1.097
VERTICAL INPUT NOSE PO	SITIÚN (M)	105 3.112 0	-114 0+112 C+1C4
HERIZONTAL Vertical ••		124 0.221 0 117 -0.108 -0	•533 0•834 1•092 •106 -0•108 -6•117
	COMP. (M/S): ( Of Cubic Polyne		G33.
0.1046E 00			83E 05 / U.JO15 (M)
CUEFFICIENTS	CMP. (M/S): 197 Of CUBIC POLYNO	MIAL/STANUARD	DEVIATION
-0.12826-01	0.20242 03 -00	12322 95 0.43	62E (6 / 0.0025 (N)
NÚSE VEL OIR Separation an	ECTION(DEG) 1. IGLE(DEGREE), RE	LATIVE TO NOSE	
A80VE Below			8•6 **** 4•8 6•4 **** 3•2
C.G. VEL. Y-C	OMP. (M/S): 5 Of Cubic Polyno		3C5.
0.1047E 90			07E 05 / 0.0013 (M)
COEFFICIENTS	CMP. (M/S): 197 OF CUBIC PULYNO	MIAL/ST/NDARD (	
	9.2023E 03 -6. CGEFF. = 1.894		19E 06 / C+0026 (M)
VU = 136. ST	ANL. LEVIA. = U DMP. (M/5): 205		36. 106. 86.
RECORDED TIME	OF MAXIMUM/NIN		
HIN .	COCIAL .0C31 001242 .048 PUSITILN AT MA	37 .005730 45 .007685 X/MIN COIL VOLT	•008882 •011211 [AGF (M)
AT MAX At Min	0.015 U.5 0.226 0.7	14 0.825	1.119
	PUSITION (M) 0.0 0.0		1.086
AT MAX AT MIN		NCSE AT MAX/MIR 26 0.034 42 0.222	

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# SHOT 56 ( 10 APRIL. 1976. NO. 5 )

SANU: DRY, DENSITY:1538, KG/M0#3; APPRUACHING VELOCITY: 324, M PhDJECTILE: SOLID STEP-TIER MASS:0.5653 KG, D=0.02 M, L=0.238	/5 M
X-RAY STATION ******* NU+1 NU+2 NO+3 NO+4 NO+5	
TIME (SECOND)	8
HORIZONTAL	
INCLINATION ANGLE(DEG). 0.0 0.5 -1.0 -2.5 -5.0	
SEPARATION ANGLE (DEGRÉE) ABOVE ************************************	
BÉLUM ************************************	
NOSE POSITION (M) HURIZONTAL	
VENTICAL	
HCHIZONTAL	
VERTICAL	
NOSE VEL. Y-COMP. (M/S): 10. 6. 134.	
CGEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIATION UNICATE 00 UNITION (M)	
NUSE VEL. X-COMP. (M/S): 314. 281. 228. 174. 126.	
CUEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIATION -0.24198-01 0.3243E 03 -0.2794E 05 0.1012E 07 / 0.0069 (M)	
NDSE VEL. DIRECTION(DEG) 1.8 1.3 0.3 -1.0 -2.0 Separation angle(degree). Relative to nose velocity	
ABOVE	
C.G. VEL. Y-COMP. (M/S): 9. 7. 4. 02.	
COEFFICIENTS OF CUBIC PULYNOMIAL/STANDARD DEVIATION U+1046E 00 - 0+9832E 01 -0+1919E 04 - 0+9958E 05 / 0+0006 (M)	
C.G. VEL. X-CÔMP. (M/S): 314. 280. 228. 174. 126.	
COEFFICIENTS OF CUBIC POLYNGMIAL/STANDARD DEVIATION -J.1462E 00 0.3243E 03 -0.2800E 05 0.1320E 07 / 0.0069 (M)	
PONCELET DRAG CUEFF. # 1.760	
VO = 314. STAND. DEVIA, = 0.0133 (M) C.G. VEL. X-COMP. (M/S): 314. 273. 223. 182. 150.	
erer neer a demonstration of the state the state state	
RECORDED TIME OF MAXIMUM/MINIMUM COLL VOLTAGE (S)	
MAR 000124 001957 003571 005839 MIN 000714 002950 0064752 007298	
COMPUTED NOSE PUSITION AT MAXIMIN COLL VULTAGE (M) AT MAX J.GIG J.511 J.B24 1.118	
AT MIN 3.194 0.715 0.995 1.248 Recorded Coll Position (M)	
0.0 U.486 0.791 1.086 DIFFERENCE BETWEEN CGIL AND NOSE AT MAX/MIN VOLTACE (N)	
AT MAX 0.016 0.025 0.033 0.032	
AT MIN 0+194 0+229 0+204 0+162	

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## SHUT 57 ( 16 APRIL: 1976; NO. 6 )

		31 APPROACHING VI S10+5651 KG+ D≖0+	ELDCITY: 198. M/S •02 N. L≈0.230 M
X-RAY STATION		NG+2 NG+3	NQ.4 NG.5
TIME (SECOND) Center of gravit		.001422 .003090	·005512 ·008429
HCHIZUNTAL	•••••••	0.117 00356 0.108 0.113	0n664 0.934 Call9 Cal42
INCLINATION ANGLE		1.0 3.G	8.0 10.5
ABOVE		3.5 2.5 4.5 5.6	#### 1+°
BELOW	FILMI	4.5 5.6 6.0236 J.0230	0.0240 0.025C
NOSE POSITION (M HCRIZONTAL,		0.239 0.477	0.785 1.054
VERTICAL		0.110 0.119	0.136 6.164
HCRIZONTAL	0.034	0.241 0.469 -0.110 -0.100	0.774 1.032
VERILGAL		-04110 -04100	
NOSE VEL . Y-COMP	• (M/S): 7.	6. 6.	
		L/STANDARD DEVIAT De 03 0.6165e 05	
NOSE VEL. X-COMP	• (M/S): 186.	162. 139.	169. 79.
-0.1329E-01 0.	1842E 03 -0.798	JSTANDARD DEVIAT BE 04 0.1388E 00	5 / C.OO95 (M)
NOSE VEL. DIRECT		2.1 2.4 IVE TU NOSE VELUC	400 809 :ITY
ABOVE	****	4.0 1.9	**** -0.1
BELU#		3:4 5:6	**** 3.1
		5. 2.	
COEFFICIENTS OF 0.9773E-01 0.	CUBIC POLYNONIAL 9566F 01 -0.2064	./STANDARD DEVIAT De 04 0.1834e 06	10N 2 V•0007 (M)
C.G. VEL. X-COMP	. (M/S): 179.	162. 139.	189. 79.
COEFFICIENTS OF -0.1351E 00 0.		./STANDARD DEVIAT Le Q4  051298e 06	
PONCELET DRAG CO			
VO = 139. STAND C.G. VEL. X-COMP	. UEVIA. = G.016 . (M/S): 19G.	57 (M) 165+ 139+	113. 92.

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## SHUT 56 ( 22 APRIL: 1976, NO. 1 )

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SANC: DRY. Projectile:	DENSITY:1536. Solid Step-t	KG/M+03; APPRUA Ilir Massiu.5647	ACHING VELOCIT Kg, D=0.02 N.	Y: 334. M/S L=0.238 M
X-RAV STATI	QN	140.1 NO.2	110+3 NO+4	N0.5
	D) -J Ravity positi	.000127 .000780	.061975 .0034	04 .005100
HCRIZÜNTA			0.381 0.64 0.102 0.10	
	ANGLE(DEG). ANGLE(DEGREE)		0.0 -1.0	-2.5
ABOVE			203 **** 400 ****	
	(M ON FILMI.		J.0240 0.024	
HCRIZONTA		0.031 4.208 0.055 0.698	0.503 0.76 0.102 0.10	
INPUT NOSE	POSITIUN (M)		0.497 0.75	U 1.015
VERTICAL	••••	0.029 C.205 -9.131 -0.124	-0.120 -0.12	6 -6.120
NUSE VEL . Y	-CUMP. (M/S):	7. 4.	10.	0.
0.9407E-0	5 DF CUBIC PO 1 0.7207E 01	LYNOMIAL/STÂNDAR -U.1977E L4 U.	1715E 06 / 0.	UC07 (M)
NOSE VEL. X	-COMP. (M/S): S GF CUBIC PO	JGJ. 266. LYNUMIAL/STANDAR	212. 170. D DEVIATION	153.
-0.1050E-0	1 0.3113E 03	-0.3136E 05 C.	2070E 07 / 0.	0C88 (M)
NOSE VEL . D	IRECTION (DEG)	1.3 1.0	0.4 -0.1	0+2
ABOVE		• KELATIVE TO NO **** 4.0 **** 5.5	2.9 ****	
		**** 5.5 5	3.6 ****	2 • J
C.G. VEL. V- COEFFICIENTS	-COMP. (M/S): 5 OF CUBIC PU	6. 4. LYNOMIAL/STANDAR	2. I. DEVIATION	3.
0 -9429E-01		-0+1423E C4 0.		GUCA (N)
C.G. VEL. X- CUEFFICIENTS	-COMP. (M/S): 5 OF CUBIC POU	303. 266. Lynghial/Standar	212. 170. D DEVIATION	153.
-J.1326E UL	. C.3112E 03	-0.3138E 05 0.	2074E 07 / 0.0	0068 (M)
VO = 212. S	AG CGEFF. = 1. Stand. Devia.	= 0.0088 (M)		a 13 <b>18</b>
C.G. VEL. X-	-CUMP. [M/S]:	311. 267.	212. 170.	137.
RECORDED TIM		MININUN CGIL VO 101916 -JU349		
MEN	.006745 .0	062929	1 .007075 DLTAGE (M)	
AT MAX At min	0.037 0.205	0.455 0.78 0.684 0.99	2 1.100	
	L PÚŠÍŤIÚN (N Uđu	Nj U•486 J•81		
AT MAK	6.037 -	AND NUSE AT MAX/ -0.001 -0.02	MIN VOLTAGE ()	4)
AT MIN	0.205	0.198 0.18	0 0.269	

#### SHUT 54 ( 22 APRIL. 1976, NO. 2 )

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SAND: DRY. DENSITY:1538. KG/M##3; APPRUACHING VELD Prujectile: Sqlid Step-tier Mass:0.5649 kg, D=0.02	CITY: 335. M/S M. L=0.238 M
X-RAY STATION ******* NO+1 NO+2 NO+3 NO	D.4 NO.5
TIME (SECOND)	03491 •005696 •684 c•989
	.115 C.IC8
INCLINATION ANGLE(DEG). 0.0 0.5 -1.0 -3 Separation Angle(Degree)	
BELOW	2•5 2•6 5•4 7:
NOSE POSITION (M)	.892 1.111
VERTICAL	107 U.U98
HCRIZONTAL	798 1.109 113 -0.122
	-26.
COEFFICIENTS OF CUBIC PULYNOMIAL/STANDARD DEVIATION 0.1059E 06 0.2160E 01 -0.4530E 03 -0.2995E 05 /	
NOSE VEL. X-COMP. (M/S): 277. 257. 219. 17 Coefficients of Cubic Polynomial/Standard Deviation	
0.5737E-02 0.2802E C3 -0.1459E C5 -0.9062E 05 /	
NOSE VEL. DIRECTION(DEG) ().4 U.3 -0.0 -0 Separation angle(degree). Relative to nose velocity	•7 -3•2
ABOVE #	• 3 3• 3 • 2 5• 7
C.G. VEL.3 Y-COMP. (N/S): -Q. 2. 3.	29.
COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION 0.1062E 00 -0.1004E 01 0.2105E 04 -0.3305E 06 /	
C.G. VEL. X-COMP. (M/S): 276. 257. 219. 17 Coefficients of Cubic Polynomial/Standard Deviation	5. 105.
-0.1163E 00 0.2601E 03 -0.1451E 05 -0.9933E 05 /	
PONCELET DRAG COEFF. # 1.667 VO # 219. STAND. DEVIA. # 0.0276 (M)	
C.G. VEL. X-COMP. (M/5): 314. 273. 219. 17	9. 140.
RECORDED TIME OF MAXIMUM/MININUM COIL VOLTAGE (S) MAX •006062 •001916 •093500 •005460	
MIN •000739 •00300G •004839 •007209 COMPUTED NOSE PUSITION AT MAX/MIN CUIL VOLTAGE (M) AT MAX 0•023 0•468 0•804 1•08	-
AT MIN 0.205 0.713 1.010 1.23 RECORDED COL PUSITION (M)	
0.0 0.486 0.810 1.080 DIFFERENCE BETWEEN COIL AND NOSE AT MAX/MIN VOLTAGE	(M)
AT MAX 0.023 0.002 -0.006 0.001 AT MIN 0.205 0.227 0.200 0.147	

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# SHUT 61 ( 22 APRIL, 1976, ND. 4 )

SAND: DRY. D Phujectile:	DENSITY:1538. KG/M##3; APPROACHING VELOCIT SULID STEP-TIER MASS:0.5636 KG. D=0.02 M.	Y: 336. M/S L=0.238 M
	064 + + + + + + + + + + + + + + + + + + +	
TIME (SECOND	))	09 .005714
HERIZONTAL		
SEPARATION A	ANCLE(DEG). U.U -0.5 -1.5 -4.5 NGLE(DEGREE)	-4.5
LELUI	••••••••••••••••••••••••••••••••••••••	
VERTICAL DE	C	2 6.099
	••••••••••••••••••••••••••••••••••••••	3 i.086 3 -0.124
CUEFFICIENTS	CGMP: (M/S): -8: -4: G: 1: DF CUBIC POLYNOMIAL/STANDARD DEVIATION -6:83716 01 0:31266 04 -0:34906 06 / 0:0	-6. )(17 (M)
NOSE VEL X-0	COMP+ (M/S): 284+ 266+ 215+ 149+	109.
	OF CUBIC PULYNUMIAL/STANDARD DEVIATION 0.28962 U3 -0.19392 U5 U.42212 06 / 0.0	G10 (M)
SEPARATION AN	RECTIUN(DEG) -1.5 -0.9 0.0 C.3 NGLE(DEGREE), PELATIVE TO NOSE VELICITY	-3.1
ABOVE Belgw	••••••••••• **** 7.6 7.5 8.8 ••••••••• **** 8.4 5.5 3.2	3.4 6.6
COEFFICIENTS	COMPS (M/S): -93. 4. 4. Of Cubic Polynumial/standard deviation -0.1.13e 02 0.5343e 04 -0.6252e 06 / 0.0	-10.
	COMP. (M/S): 284. 260. 216. 169.	109.
COEFFICIENTS -0.1272E 00	OF CUBIC POLYNOMIAL/STANUARD DEVIATION G.20952 03 -0.19262 05 C.40762 06 / 0.0	011 (M)
VO = 216. ST	S COEFF. = 1.815 Tang. Devia. = 0.0221 (m)	
C.G. VEL. X-C	(CHP+ (M/5): 317+ 274+ 218+ 173+	134.
MAX .	OF MAXIMUM/MINIMUM COIL VOLTAGE (S) 000109 001929 003863 005767	
	0000699 0002957 0004944 0007491 Pusition at Max/Min Cuil Voltage (M) 00026 00484 00805 10101	
AT MIN	0.100 0.693 1.004 1.254 . Positium (M)	
DIFFERENCE BE AT MAX AT MIN	Ú.L. U.486 U.81C I.086 TWEEN COIL ANÚ NUSE AT HAX/MIN VOLTAGE (M. 3.026 -J.0J2 -0.004 0.015 G.138 0.207 ú.194 0.168	)

# SHUT OF 1 22 APRIL, 1976, NO. 5 )

SANC: DRY. DENSITY:1538. KG/N##3: APPHDACHING VELOC Projectile: Solid Step-tier Mass:0.5633 kg. D=0.02	ITY: A. N. M/S M. L=Q.238 M
X-RAY STATIUN	la4 NDej
TIME (SECOND)	2702 .004005
SCOLIONTAL ANALANA -Q.077 GALUT VESTY VE	624 0.847 103 0.116
	.5 8.0
SEPARATION ANGLEIVEGREE	.0 ****
	.5 ****
	230 0.0226
NUSE BIDTH (M ON FILME GROUPS GROUPS CONTROL C	
HLRIZANTAL	746 6.968
VERTICAL	115 0.133
INPLT NOSE POSITION (M)	732 0.950
	105 -0.085
VERTICAL	
	5. 24.
NOSE VEL. Y-COMP. (M/S): 12. 5. 1. CUEFFICIENTS OF CUBIC PULYNDMIAL/STANDARD DEVIATION	
U-10422 00 0-1356E 02 -0-7442E 04 0-1455E 07 /	0.0015 (M)
	)3. 139.
- ARE RETAILSTO OF CONTERPONDED AND ARU DEVIATION	
U. 10332-02 U. 3510E 03 -0.2930E 05 U.4710E 06 /	010019 (M)
	.5 9.8
- RECIDATION ANGLE (LEGREE). RELATIVE TO NUSE VELUCIT	5.G ****
	200 <del></del>
BELUM	
	C. 22.
C.G. VEL. Y.COMP. (M/S): 12. 33. CUEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIATION	
U-1012E OU U-1389E 02 -0.9613E 04 U-1759E 07 /	G.CCCB (M)
	<b>140</b> ~
- COCCETTETENTE OF CURTE POLYMIAL /STANDARD DEVIALIUN	1
-U.1203E 06 U.3510E W3 -U.2918E 05 U.4669E U6 /	C.UC20 (M)
PONCELET DRAG COEFF. = 1.827	
- いっ こ つはた - たてんれい。 いたいてん こ エーじょりついじ 〔餘〕	~
C.G. VEL. X-COMP. (M/S): 374. 321. 259. 21	3. 175.

# SHOT 63 ( 22 APRIL: 1976. NO. 6 )

SAND& DRY& CENS Prujectile: Sou	SITY:1538. KG/M++. LID STEP-TIER MASS	3; APPROACHING VI 5:0:5533 KG; D=0	ELOCITY: 415, M/S .62 M, L#J.238 M
X-RAY STATION	NU.1	NG+2 ND+3	N0.4 N0.5
	ITY POSITION (M)	.000683 .001708	.003084 .004798
		u •109 0•396 U c l u • U • 104	0.711 C.989 0.114 0.136
INCLINATION AND SEPARATION ANGL		3.5 8.0	11.5 11.5
AUUVE		10.0 12.0 4.ú 1.5	
	DN FILMI. U.U256	6.6230 0.0230	0.0230 0.0230
HERIZUNTAL		0.231 0.516 0.111 0.121	0.831 1.108
VERTICAL	[TEUN (N)		0.138 0.161
HCKIZONTAL •• Vertical ••±•		u.231 U.514 -0.109 -0.098	0.630 1.103 -0.078 -0.052
NGSE VLL. Y-COM	H. (M/S): 6. Cubic Polynomial	8. 11. ASTANDARD DEVIAT	13. 13.
	-534 8E 01 4.2234		
	P. (M/S): 341. Cubic Polynumial		
	-3499E 03 -6.2920		
	TION(LEG) 1.0		
ALUVE		5.0 5.5	4.4 4.9
BELOF		6•u 7ºA	8.6 7.1
	( . ( M/S): -0.		9. 18.
	CUBIC POLYNUMIAL •41505 00 0.607J		
	P. (M/S): 342. Cubic Pulynomial		
	-JAHVE 63 -0.2662		
PONCELET DRAG CO	CEFF• = 1+811 L• DEVIA• = 0-8214	A (M)	
C.G. VEL. X-COM		327 <b>.</b> 259.	203. 160.

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# SHUT 64 ( 22 APRIL, 1976, NO. 7 )

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SAND: DRY, DENSITY:1538, KG/M##3; APPROACHING VELC Projectile: Solid Step-tier Mass:0.5633 kg, D=0.02	CITY: 409. H/S ! M. L=J.238 M
X-RAY STATION NO.1 NO.2 NO.3 N	10.4 NO.5
TIME (SECOND)	03106 0004861
MORIZONTAL	)•704 C•978 )•102 O•122
STRAGATION AGGLEIDEGREEN	0.0 9.5
DEFOR \$\$\$\$\$\$\$\$\$\$\$	3.0 11.0 2.0 1.5 0225 0.0220
NGSE PUSITION (H)	
HGRIZUNTAL	
HURIZONTAL ********* 0+043 0+220 0+504 0 Vertical ************************************	0.822 1.094 0.055 -0.075
NOSE VEL. Y-COMP. (M/S): 7. 8. 9. COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATIO	10. 11.
J.956UE-OL U.69048 UL U.8U9UE G3 -0.56J2E 05 /	0.00C2 (M)
COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATIO	91° 137• N
-0.8636E-02 (.3598E 63 -0.3429E 05 (.1536E 07 /	0+0034 (M)
NOSE VEL. DIRECTION(DEG) 1.2 1.5 2.1 Separation anglé(degree). Felative to nose velocit	3•1 4•5 Y
	6.1 6.0 8.9 6.5
C.G. VEL. Y-COMP. (M/S): 6. 3. 1.	5. 20.
COEFFICIENTS OF LUBIC POLYNUMIAL/STANDARD DEVIATIO U-9364E-01 U-7402E 01 -0-3646E 04 0-6699E 06 /	
C.G. VEL. X-COMP. (M/S): J49. 316. 258. 1 COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATIO	
-J.1305E 00 0.3590E 03 -0.3550E 05 0.1417E 07 /	0.0036 (M)
PONCELET DRAG CGEFF. = 1.917 VO = 258; Stande devia. = 0.0174 (m) C.C. Vel. X-comp. (m/5): 383. 329. 258. 1	99. 155.
RECURDED TIME OF MAXIMUM/MINIMUM CCIL VOLTAGE (S)           MAX         •000124         •001646         •002972         •0047           MIN         •000668         •002547         •004146         •0061	
COMPUTED NOSE PUSITION AT MAX/MIN CUIL VOLTAGE (M) AT MAX 0.035 0.497 0.798 1.00	
RECORDED COLL PUSITION (M)	61
JOG 00486 JOBIO LOD DIFFERENCE BETWEEN COIL AND NUSE AT MAX/MIN VOLTAGE AT MAX 40435 CO411 -00012 000	Ē"(M)
AT MIN 0.217 0.225 J.193 0.1	

#### SHUT 65 ( 22 APRIL, 1976, NO. 8 )

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SANC: DRY, DENSITY:5380. KG/N++3: APPRDACHING VELOC Projectile: Solid Step-tier Mass:0.6631 kg. D=0.02	
X-RAY STATION	1.4 NC.5
TIMÉ (SEGOND)	3121 .004804
	713 0.993
INCLINATION ANGLE(DEG). 0.0 0.0 0.5 -1	•0 -3•5
SEPARATION ANGLE(DEGREE) Above	.5 2.0
BELOW	
NOSE POSITION (M)	
INPUT NOSE PUSITION (M) MCHIZONTAL -200000000 U+U55 U+232 U+515 0+	835 1.111
VERTICAL	
NOSE VEL. Y-10MP. (M/S): 5. 3. 1	12.
CDEFFICIENTS OF CUBIC PULYNOMIAL/STANDARD DEVIATION 0.5963E-01 0.4867E D1 -0.1143E 04 0.5990E 05 /	
NOSE VEL. X~COMP. (M/S): J48. 314. 257. 19. Coefficients of Cubic Polynomial/Standard Deviation	
-0.2609E-02 C.J588E 43 -0.3408E 65 0.1622E 07 /	0.0022 (M)
NOSE VEL. DIRECTION (DEG) 0.7 0.6 0.3 -0	• 2 - 0 • 8
SEPARATION ANGLE (DEGREE), RELATIVE TO NOSE VELOCITY	
ABUVE ************************************	•3 ••7 •7 3•3
C.G. VEL. Y-COMP. (M/S): 2. 2. 2. 2. CUEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION	2. 2.
0.1003E 00 0.2379E 01 0.1000E 02 -0.1075E 05 / 0	0.0C25 (M)
C.G. VEL. X-COMP M/S): 348. 314. 257. 194 CDEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION	
-0.1246E 00 0.3569E 03 -0.3468E 05 0.1626E 07 / 3	
PONCELET DRAG CUEFF. # 0.536	
VO = 257. STAND. DEVIA. = 0.6133 (M) C.G. VEL. X-COMP. (M/S): 378. 325. 257. 199	9. 197.

SHOT 06 ( 22 APRIL. 1970. NO. 9 )

SAND: DRY, DENSITY:1538. K PROJECTILE: A.M.H.	G/N++3 MASS	: APPRUA :0 =0024	KG. D=0.	02 M. L.	588. #/5 10.076 M
	10 • 1	NU• 2	N0+3	NU.4	NO+5
TINE (SECOND)		.01158	.007983	.020000	-100000
CENTER OF GRAVITY PUSITION	•062 •088	0.267 U.U71	J-395 J-055	18.839 18.121	19-144 18-121
INCLINATION ANGLE(DEG)	6.0	-49.0	*****	*****	****
	***	**** ****	**** **** u.6240	**** **** *****	本(1) (1) (1) (1) (1) (1) (1) (1) (1) (1)
NOSE WIDTH (M ON FILM). 04 NOSE POSITION (M)	.026û ).662	0.0235 J.267	0.395	18.839	19.144
VERTICAL	.088	Ú.)71 0.274	0 • 0 5 5 0 • 3 6 7	18.121	18.121 ******
MERIZONTAL PRARECOVE (	).069 ).140	0.274 -0.156	-0.177	*****	******

SHOT 67 ( 22 APRIL, 1976, NO.10 )

SANGI DRY. DÉNSITVI1538 Projectile: A.M.H.		31 APPRO 310.0412			575. M/S =0.07J M
X-RAY STATION	NG • 1	ND.2	N0+3	NO . 4	NO+5
TIME (SECOND)		+¥01186	.J07981	.J20000	.100000
HURIZONTAL	0.000	N.248	0.364	18.839	19.144
VENTICAL	<b>U ~ V &amp; </b>	4.60	0.062	18.121	10.125
INCLINATION ANGLE(DEG), SEPARATION ANGLE(DEGREE	-7.5	- 58 - 5	*****	****	*****
	4	****	****	****	****
BELUN	14.0	****	****	****	****
NUSE WIDTH (M ON FILM). NOSE PUSITION (M)		0.0240	0.0240	*****	******
HCRIZONTAL	0.005	0.248	0.364	18.839	19.144
VENTICAL	0.086	0.069	0.662	18-121	18.121
INPUT NOSE PUSITION (M)	***	****		• • • • • • •	
HURIZONTAL	0 0 73	0.252	6.330	*****	*****
MEGTICAL			-0.168	*****	*****

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# SHOT 68 ( 23 APRIL. 1970. NO. 1 )

SAND: WET. DENSITY:2650. KG/M043; APPROACHING VELOCITY: 357. M/S Projectile: Solid Step-tier Massiu.5664 kg. D=0.02 m, L=0.238 m N0.5 NO .4 NG.J NO. 2 NO.1 X-RAY STATION ...... VERTICAL ........ -3.0 -1.0 0.0 INCLINATION ANGLE (UEG). 4.0 4.4 SEPARATION ANGLEIDEGHEET 7.0 12.0 14.0 13.0 12.0 12. 1.161 0.815 0.221 0.040 HEHIZONTAL ........ 0.197 0.105 0.107 0.099 0.103 1.164 0.812 0.544 0.220 6.041 HCHIZONTAL ...... -0-114 -0.110 -0.114 -6.118 NOSE VEL. Y-COMP. (M/S): 8. 5. 2. C. 1. COEFFICIENTS OF CUBIL POLYNUMIAL/STANDARD DEVIATION U.9832E-01 0.8637E 01 -U.3USUE 04 0.3592E 06 / 0.0011 (M) NOSE VEL. X-COMP. (M/S): 355. 346. 321. 288. 218. COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION -0.1462E-02 0.3561E 03 -0.6032E 04 -0.2060E 07 / 0.0013 (M) 0.3 0.0 13.0 10.3 8.7 11.0 C.G. VEL. Y-COMP. (M/S): 7. 5. 3. 3. 5. COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIATION 0.9858E-01 0.7234E 01 -0.210/E 04 0.3144E 06 / 0.0015 (M) C.G. VEL. X-COMP. (M/S): 354. 346. 321. 288. 218. COEFFICIENTS OF CUBIC PULYNOMIAL/STANDARD DEVIATION -0.1239E 00 0.3560E 03 ~0.6000E 04 -0.2064E 07 / 0.0014 (M) PONCELET DRAG CDEFF. = 0.614 VO = 321. STAND. DEVIA. = 0.6177 (M) C.G. VEL. X-COMP. (M/S): 385. 360. 256. 292. 321 . RECORDED TIME OF MAXIMUM/MINIMUM COIL VOLTAGE (S) .003615 .000084 .001494 .002457 .000565 .062149 ####### MAX .004394 COMPUTED NOSE PUSITIUN AT MAX/MIN COIL VOLTAGE (H) AT MAX 0.028 0.510 0.806 1.10 1.109 AT MAX 1.271 V.197 0.715 RECORDED COIL PUSITION (M) 1.086 L.486 0.810 DIFFERENCE BETWEEN CUIL AND NOSE AT MAX/MIN VOLTAGE (M) AT MAX 0.428 0.024 -0.004 0.023 AT MIN 0.147 4.226 ####### 0.185 AT MAX 0-197 4.225

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SHUT 65 ( 23 APRIL, 1976, NO. 2 )

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SANC: WET, DENSITY:2050, KG/M##3; APPRUACHING VELOCITY: 395, M/S Projectile: Solig Step-tiek Mass:0.5643 kg, d=0.62 m, l=0.238 m
X-RAY STATION NL.! NL.Z NG.3 NO.4 NO.5
TIME (SECUND)
MURIZUNTAL ********* -0.092 0.111 0.448 0.719 1.075 VERTICAL ************************************
INCLINATION ANGLE(UEG). U.5 0.0 0.0 0.0 -3.0 SEPARATION ANGLE(UEGREE)
ABUVE ************************************
NÜSE WIDTH (M ON FILM). 6.0266 0.0236 J.0240 0.0230 6.0240 NGSE POSITION (M) MERIZONTAL
VERTICAL
HGRIZONTAL ************************************
NOSE VEL. Y-COMP. (M/S): 10. 9. 5. 17. COEFFICIENTS OF CUBIC POLYNOMIAL/STANDARD DEVIATION
0.1206E 00 0.1079E 02 -C.1553E 04 -C.1359E 06 / 0.0016 (M)
NUSE VEL. X-COMP. (M/S); 401. 374. 325. 287. 235. CGEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIATION
-U.10242-01 V.40692 03 +0.2784E 05 G.9462E 06 / 0.0041 (M)
NDSE VEL. DIRECTIUN(DEG) 1.5 1.3 0.8 0.1 -1.8 Separation angle(degree). Relative to NDSE velocity
ABCVE ************************************
C+G•VEL•V-COMP•(M/5): 14• 10• 4• 2• 2• 2• Coefficients of Cubic Polynomial/Standard Deviation
0.1192E 00 0.147UE 02 -0.4452E 04 0.4882E 06 / 0.0016 (H)
C.G. VEL. X-COMP. (M/S): 401. 374. 325. 287. 235. CUEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIATION
-J.1322E 00 -U.4068E UJ -0.2782E 05 -0.9452E 06 / 0.0042 (M) Poncelet drag coeff. = 0.711
$\begin{array}{rcl} & \text{VO} = & \text{O} & $
RECORDED TIME OF MAXIMUM/NINIMUM COIL VOLTAGE (S) Max +000112 +001407 +002416 +003562
MIN +000559 +002028 +003158 +004464 Computed Nose Position at Max/Min Cuil Voltage (M) At Max 0+035 +0+516 +0+824 +1+129
AT MIN 0.205 0.706 1.027 1.322 Recorded Coil Position (M)
0.0 0.486 0.810 1.086 Difference between cuil and nuse at max/min voltage (m)
AT MAX 0.035 0.024 0.014 0.043 AT MIN 0.205 0.222 0.217 0.236

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SHCT 70 ( 23 APRIL: 1976, NO. 3 )

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#### SHUT 71 ( 23 APRIL: 1976, NO. 4 )

		KGZNAAJI APPROA		
X-RAY STATION		NÜ+1 ND+2	NO+3 NO+4	ND.S
CENTER OF GRA Herizontal	VITY PUSITI	-0.102 0.091	.003227 .00591 0.413 0.722 0.124 0.122	1.013
INCLINATION A SEPARATION AN ABOVE	GLE ( DE GREE )	2.5 3.5 · **** ****		
NOSE WIDTH (N NOSE POSITION	I ON FILM). ( ( (N)	0.0270 0.0230	0.0246 0.0240	0.0240
VENTICAL INPUT NOSE PL	SITION (M)	Ue116 Ue124	6.130 0.132	0.129
		0.002 0.200 -0.104 -0.094		
CGEFFICIENTS	OF CUBIC PU	8. 6. Lynumial/Standar -0.1285e 04 g.	RD DEVIATION	-0. C18 (m)
COEFFICIENTS	OF CUBIC PO	212. 185. Lynumial/Standar -6.1589£ 05 0.	O DEVIATION	
NOSE VEL. DIR	ECTION (DEG)	2.1 1.8 . HELATIVE TO NO	0.8 -0.4	
ABOVE		**** ****	4.3 **** 5.7 ****	****
COEFFICIENTS	OF CUBIC PU	9. 6. Lynomial/standar -0.1978e 64 0.	D DEVIATION	5. 104 (M)
CoGo VEL . X-C	OMP. (#/5);	212. 185.	135. 94.	
-U.1431E GU	0.2187E 03	LYNUNIAL/STANDAR -0.158JE 65 0.	5955E 06 / 0+G1	102 (M)
PONCELET DRAG VO = 212. ST C.G. VEL. X-C	AND. DEVIA.	(543) ≠ J.G125 (M) 212. 180.	134. 101.	76.
RECORDED TIME Max Min	0F MAXIMUM 000208 +0 001155 +0	/MINIMUH CUIL VO 063112 -00569 004621 -00773	LTAGE (S) 9 .009317 6 .012270	
COMPUTED NOSE AT MAX AT MIN	PUSITION A1 C.015 0.262	T MAX/MIN CUIL V 0.515 0.81 C.7u1 0.99	OLTAGE (M)	
RÉCORDED COIL	PLSITION (N	4) 0.486 0.81 AND NOSE AT MAX/	0 1.085	
AT MAX AT MIN	0.015	Ú.029 Ú.00	2 0.029	

# SHUT 72 ( 23 APRIL, 1976, NO. 5 )

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			HING VELOCITY: 214. M/5 5. D=0.02 M. L=0.225 M
K-RAY STATION		NO+1 NO+2 M	N0+3 N0+4 N0+5
CENTER OF GRA	VITY POSITI	UN (M)	03146 .005870 .009410
HCRIZONTAL VERTICAL ••			0.430 0.760 1.075 0.129 0.132 0.139
INCLINATION A	GLE (UEGREE)	••••	0.5 -1.5 -5.0
AGCVE Bélüw Nose Width (M		**** **** **** **** C.0250 U.0230 J.	6.0 3.0 4.0 5.5 8.0 5.0 ,0230 0.0240 0.0240
NOSE POSITION	(m)		
HCRIZONTAL	SITION (M)	3.024 4.206 0	.545 0.884 1.196
VERTICAL		-0.107 -0.097 -0	.687 -0.688 -0.088
COEFFICIENTS (	OF CUBIC PO	11. 7. LYNUMIAL/STANDARD	
		-0.2190ê 64 C.12 210. 187. 1	224E 06 / 0.0003 (M)
CUEFFICIENTS (	ÖF CUBIL PU	LYNOMIAL/STANDARD	
NG65 MEL 010.	SCTTON ( T.S. C.)	3.4. 3.3	0.6 -0.8 2.1
SEP ARATION AND	GLE(ŬEGRÉÉ)	3+0 2.2 • HELATIVE TU NOSE • + + + • • • •	VELOCITY 6.1 3.7 11.1
BELOW		****	6•1 3•7 11•1 5•4 7•3 v•9
CUEFFICIENTS (	OF CUBIC POL	10. 7. Lynumial/stanuard	2. C. 5. DEVIATION
			68E 06 / 0.JC06 (N)
COEFFICIENTS C	DF CUBIC POL	210. 167. 1 Lynomial/standard	43. 102. 81. Deviation 238 06 / Guocii (M)
PONCELET DRAG	CUEFF. = 1	.359	
VU = 210. STA C.G. VEL. X-CO	AND. LEVIA. DMP. (M/S):	≖ 0.0098 (M) ≖ 0.0098 (M)	40. 107. 82.
RECORDED TIME	OF MAXIMUNA	MINIMUM COIL VOLT	AGE (S)
	)041099 041099	203631 .005326 564375 .607174 7 Max/Min Coil Vol	• 008634 • 011431
AT MAX	0.621	U.527 U.816 U.707 U.998	1.124
RECORDED COIL	POSITION (P	0.707 0.998 4) 0.486 0.810	
DIFFERENCE BET	WEEN COIL A	AND NUSE AT MAX/MI	N VOLTAGE (M)
AT NIN	6.296	Ú•221 J•188	

# SHUT 73 ( 23 APRIL, 1976, NO. 6 )

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SHOT 74 ( 23 APRIL, 1976, NO. 7 )

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SANG: WET. DENSITY:2050. KG/M4#3; APPRUACHING VELOCITY: J34. M/S Projectile: Sglid flat Nuse Mass:0.5450 kg. D=0.02 N. L=0.225 M

# SHOT 76 ( 24 APRIL, 1976, NO. 2 )

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## SHOT 77 ( 24 APRIL, 1976, NG. 3 )

SAND: DRY. DENSITY:1536. KG/N##3; APPRUACHING VELUCITY: 242. M/S Projectile: Flat 6-inch cyl.Mass:L.3667 kg. D=0.02 M. L=U.152 M

X-RAY STATION	N0+1	NO. 2	NQ. 3	N0+4	N0.5
TIME (SECOND)		.000786	.002323	.004522	.008432
CENTER OF TRAVITY POSIT			0 741	G.882	19.132
HCRIZONTAL	-0.012	60681	0.341	V . 002	
VERTICAL	C • 1 2 U	u=120	ú•129	0.098	18.196
INCLINATION ANGLE(DEG). SEPARATION ANGLE(DEGREE	, -1.0	0.9	-4.0	-8.0	****1
					****
ABOVE	****	****	****	****	****
BELO#	****	****	****	****	***
NOSE WIDTH (M ON FILM).	じょじょうじ	9.6240	0.0235	J.J25J	******
NCSE POSITION (M)					
HCRIZONTAL	U.U64	0.157	0.417	0.957	19.144
VERTICAL SHARAAAAAAAA	6.119	v • 126	0.123	<b>℃</b> • 087	18+121
INPUT NOSE POSITION (M)					
HORIZONTAL	0.072	0.143	0.397	0.987	*****
VERTICAL		-0.092	-0.095	-0-140	******

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# SHUT 76 ( 24 APHIL. 1470. NO. 4 )

SANL: URY. CENSITY: 1538. KG/M443	5; APPHOACHING VELOCITY: 242. M/S 5:4.3687 kg. D=4.62 m. L=0.152 m
X-RAY STATION	
TIME (SECUND)	
CLNTER OF GRAVITY FESITION (M) HERIZUNTAL	C5191 0.404 0.748 1.690 0.130 0.129 0.144 0.176
INCLINATION ANGLE(DEG). Jou	2.5 6.0 11.6 16.5
SEPARATION ANGLE (DEGREE)           ABUVE         ****           BELOW         *****	**** 10.0 **** 11.C **** 2.0 **** 3.0
NUSE WIDTH (M GN FILM). U.U26U NUSE PUSITION (M)	0.6230 0.6230 0.0233 0.6215 0.177 6.486 0.822 1.169
HERIZONTAL	0.177 G.480 0.822 1.169 0.134 0.137 0.158 C.198
HCHIZUNTAL	J.169 G.472 G.820 1.171 -J.683 -J.679 -G.655 -9.715
NOSE VEL . Y-COMP . (M/SI: 3.	0. 5. 6. 12.
CEEFFICIENTS OF CUBIC PULYNUM AL U.1274E UG U.J.27E UI G.JUS	UE 63 0.1003E 05 / 6.0044 (M)
NOSE VEL. X-COMP. (M/S): 145. CLEFFICIENTS OF CUBIC FULYNUMIAL	
	7E 65 0.4047E 06 / 6.3146 (M)
NGSE VEL. DIHECTION(DEG) 0.9 Separation Angle(Deghee). Relayi	TAF IN WRZE AFFORTAL
	****         5.9         ****         10.5
C.G. VEL . Y-COMP. (M/S): 3.	3. 5. 12.
C.G. VEL. X-COMP. (M/S): 195. COEFFICIENTS OF CUBIC POLYN MIAL	
-u.7941E-03 0.1983E 03 -0.1201	₽E 05 U~4017E 3N / U+3146 (M)
PUNCELET DRAG CLEFF. = 1.311 VJ = 137. STANL. DEVIA. = 0.011 C.G. VEL. X-COMP. (M/S): 206.	39 (M) 181. 137. 103. 75.

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SHUT 79 ( 24 APHIL. 1976. NO. 5 )

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SANDI URY, DE	NS 1 TY : 1 538 .	KG/M++3; APFRUA ILEMMASSIU+4218	CHING VELOCITY	1: 2345 M/S L=0+238 M
Z-RAY STATION			N 3 NO.4	
		000124 .000876	. 002082	07 .010325
CENTES OF GRAV	/114 PUSSILI			
HCHIZGNTAL		U-084 U-073 U-120 U-126	C.39J 0.736 J.122 0.136	
VENTICAL				-5.5
INCLENATION AP	NGLE(DEG). GLE(DEGREE)	-0.5 0.0	••••	
ABUVE		0000 000 000 000 000 000 000 000 000 0	2.0 ****	10.5
NOSE WIDTH (M	ON FILMI. U		0.0230 0.023	5 0.0240
NOSE POSITION	(M)	UnC16 U0178	0.503 0.84	
VERTICAL		U.114 U.126	0.126 0.125	7 6.121
ILPUT NUSE PU		0.010 0.108	0.498 0.84	1 1.191
VENTIGAL		-U-104 - U-100 -U-104 - U-092	-0.092 -0.091	1 -03041
			11.	~ <b>0</b> •
NUSE VEL . Y-C			DOFVIATION	• •
0.1198E 00	0.4815E 01	-0.6016E 03 0	4165E 05 / U.(	0031 (M)
		10. 101.	141a 93a	81.
		LYNOMIAL/STANDAR -U.1654E US 04	ID DEVIATION	0069 (M)
-0.10432-02	0122042 03		, <u> </u>	
NOSE VEL. DIR	ECTION (DEG)	1.2 1.0	6.3 -0.7	-0.2
SEPARATION AN	GLE LLE GREE ) a	S RELATIVE TO THE	J+8 VELOCITY	6.8
85LUV		**** 7.4	4.7 ****	5.2
				_
C.G. VEL . Y-C	CMP. (N/S):	4. 3.	3. I.	-3.
CUEFFICIENTS J.IZUGE 00	OF CUBIC POL 0.3599E C1	-0.8300E 02 -0	1658E 05 / C.	0G26 (M)
		216. 193.	141. 93.	
			D DEVIATION	6667 (M)
-0.1120E 00		-4.1048E 65 0	103332 V0 / V+	
PUNCELET DRAG	CULFF. = 1	•789		
	ANG. GEVIA. Omp. (4/5):	= G.GIIU (M) 235. 199.	141. 99.	69.
RECERDED TIME	OF MAXIMUM	MININUM COIL VO	DLTAGE (5) 59 .009217	
MAX •	306127 •( 301047 •(	003005 00555 004637 00755	24 .012025	
COMPUTED NOSE	PUSITION A	T MAX/MIN CUIL	VOLTAGE (M) 16 1+117	
AT MAX AT MIN	0.020 0.206	U.722 U.9		
RECURDED COIL	POSITION (		1.086	
DIFFEHENCE BE	0.U TWEEN CUIL /	AND NUSE AT MAX	MIN VOLTAGE (	Mŷ
AT MAX AT MIN	0.026	0.037 0.00 0.236 0.10	00 0.031	

SHUT 66 1 24 APHIL: 1976. NO. 6 )

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SHUT 81 ( 24 APRIL: 1976, NO. 7 )

		G/M##31 APPRDA De Mass:J.5448		
X-RAY STATION	• • • • • • • • • •	10.1 NO.2	NU.3 NO.4	N0.5
	VITY FUSITION	CU155 .000839	.002174 .00363	• •005450
HCRIZUNTAL		.078 3.128		1.097 0.136
INCLINATION A SEPARATION AN			2.C 5c0	
BELQW	4 ON FILM) - U		9.0 9.0 6.0 3.5	10.0 1.5 0.0235
NOSE PUSITION HERIZUNTAL	i (m) 	•035 U•241	0.0240 0.0230 0.581 0.899	
INPUT NOSE PL	SITIUN (M)	6131 56136 6231 56136		0-151
VERTICAL			-3.677 -0.673	-0.062
		8. 5. NUMIAL/STANDARI		9.
0.1294E 00	0.0052E 01 -	6.2554E 04 0.	3187E 06 / 6.00	
CUEFFICIENTS	OF CUBIC PULY	10. 283. Numial/Standar( 0.2042e 05 0.4	DEVIATION	
			0.4 0.7	3.4
SEPARATION AN	*********	HELATIVE TO NOS	SE VELOCITY	5.4
8€L0W	*********	*** 8.0	7.6 7.8	5•1
COEFFICIENTS	OF CUBIC PULY	11. 5. NUMIAL/STANDAR(	DEVIATION	7.
		0.4691E 04 0.5 10. ∠83.		· •
CUEFFICIENTS	OF CUBIC PULY	NUMIAL/STANDARD	D DEVIATION	
PONCELET DRAG	CUEFF. = 0.9 ANU. LEVIA. =			
		25. 289.	237. 198.	164.
HECORDED TIME	OF MAXIMUM/M	INIMUM COIL VOL 1863 -003199 2671 -004193	TAGE (S)	
AT MAX	PUSITION AT	2671 •UG4193 Max/Min Coil VC •510	DLTAGE (M)	
AT MIN Recorded Coil	0.201 Č Pusition (m)	•696 1•601	1.303	
DIFFERENCE BE AT MAX AT MIN	TWEEN COIL AN	•486 U•B10 D NUSE AT MAX/N •024 U•001	IN VOLTAGE (M)	i -
A1 510	0-201 Ú	.212 0.191	J.217	

SHUT 62 ( 24 APRIL. 1976. WO. 8 )

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## SHUT 83 ( 25 APRIL, 1976, NO. 1 )

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# SHUT 84 ( 25 APRIL: 1976, NO. 2 )

X-RAY STATION NG.1 NG.2 NO.3	NQ.4 NO.5
TIME (SECOND)	.002702 .003929
HGRIZUNTAL	0.726 1.004 0.129 0.125
INCLINATION ANGLE(DEG): 2:5 2:5 2:0 Separation angle(degree)	2+5 3.8
ABOVE	15.0 12.0 9.0 4.0
NOSE WIDTH (M ON FILM). L.0255 U.0235 U.0230 RUSE POSITION (M)	0.0230 0.0230
HCRIZUNTAL	0.839 1.117 0.134 0.132
HCRIZONTAL	0.839 1.113 -0.083 -0.085
NGSE VEL. Y-COMP. (M/S): 12. 9. 3. CCEFFICIENTS OF CUBIC PULYNOMIAL/STANDARD DEVIA	-C2.
0.1183E UU 0.1337E 02 -0.3727E UA 0.3087E 0	6 / 0.0012 (M)
NUSE VEL. X-CCMP. (M/S): 362. J37. 293. COEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DEVIA -0.41206+03 0.36826 53 -0.22186 05 0.21306 0	
-V.41262-63 V.30022 53 -V.2210E US V.2130E U	0 / VIV94 (8/
NUSE VEL. DIRECTION(DEG) 2.5 1.5 0.6 Separation angle(degrée), relative to nose velo	-0+0 -C+5 CITY
ABUVE	12•5 7•7 11•5 8•3
C.G. VEL. Y-COMP. (M/S): 13. 9. 3. CREFFICIENTS OF CUBIC PULYNOMIAL/STANDARD DEVIA	-15. TION
0.1132E 00 0.1409E 02 -0.3062E 64 0.2008E 0	6 / G.QOCT (H)
C.G. VEL. X-COMP. (M/S): 362. 337. 293. CDEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD DEVIA	
-0.1133E 00 0.3602E 03 -0.2219E 05 0.2212E 0 Poncelet drag cueff. = 0.777	5 / C.OU54 (M)
VU = 293. STAND. DEVIA. = 0.0108 (M) C.G. VEL. X-COMP. (M/S): 374. 341. 293.	259. 226.
	03842
COMPUTED NOSE PUSITION AT MAX/MIN COIL VOLTAGE	04752 (m) 1•099
	1 • 271
UIFFENENCE BETWEEN CUIL AND NOSE AT MAX/MIN VOL	
	0+013 9+185

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## SHOT 85 ( 25 APRIL, 1976, NO. 3 )

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SAND:H-OH, DENSITY:1000. KG/M##3; APPRGACHING VELOCITY: 237. M/S PROJECTILE: HOLLOW STEP-TIERMASS:0.4230 KG. D=0.02 M. L=0.238 M PROVESTATION ADDACTOR NO.1 NO.2 NO.3 NO.4 NO.5

FRAY STATION	NO + 1	N0+2	N0+3	NO • 4	N0.5
TIME (SECUND)		.000885	.00229A	.004596	.008612
HORIZONTAL		∪•090 č•113	U.393 C.116	0.862 U.121	19•127 18•224
INCLINATION ANGLE(DEG). SEPARATION ANGLE(DEGREE	1-4	U•4	6.0	1.0	*****
ABOVE	****	电动电电	****	* * * *	***
BELOW	****	****	****	****	***
NOSE WIDTH (M ON FILM). NOSE POSITION (M)	6.0262	ü+0230	0.0231	0.0230	***~**
HORIZUNTAL	0.024	<b>U-195</b>	<b>U • 498</b>	0.967	19.144
VERTICAL	0 • 107	Ŭ•113	0.116	0.122	18+121
HORIZONTAL	6.020	0.190	0.493	0.986	******
VERTICAL	-6.116	-0.107	~0.104	-0.096	*****

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SHOT 60 ( 25 APRIL: 1976, NO. 4 )

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# SHUT 87 ( 25 APRIL, 1976, ND. 5 )

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SAND H-OH: DENS ITY: LUCU. KG/H++3: APPROACHI	NG VELOCITY: 241. M/S
PROJECTILE: MOLLOW STEP-TIERMASSIU +4227 KG+	
	• J NO• 4 NO• 5
TIME (SECOND)	2391 .004193 .006381
HCK120NTAL	406 0.775 1.122 114 0.120 C.123
INCLINATION ANGLE(UEG)1.5 -1.0 -0 Séparation angle(Uegnee)	•8 -1•7 -2•0
ABUVE	** 17.0 15.0 ** 20.0 20.0
NOSE WIDTH (M ON FILM). C.C260 U.C235 0.0	240 0.0245 0.0245
NCSE POSITION (M) Morizontal	511 C.886 1.227
HGHIZONTAL	112 U.117 0.119
HERIZUNTAL	201 00000 10240
VERTICAL	168 -0.103 -0.100
NOSE VELA Y-COMPA (M/S) BA 54	3. 1. 2.
NDSE VEL. Y-COMP. (M/S) 8. 5. COEFFILIENTS UF CUBIC ! JLYNUHIAL/STANDARD DO	EVIATION
0.1011E 00 0.0328E 1 - G.1679E C4 0.130	
NOSE VEL . X-COMP. (N/S : 227. 221. 21) CDEFFICIENTS OF CUBIC POLYNUMIAL/STANDARD DI	G. 194. 176. Eviation
-0.9338E-02 0.2181E 63 -6.3552E 64 -6.791	02 05 7 0.0052 (M)
NOSE VEL. DIRECTION(LEG) 200 1.4 0. Séparation anglé(degree). Relative to nose v	•7 0•3 C•8 VELOCITY
ABOVE	* 19.0 17.8
C.G. VEL. V-COMP. (M/S): 6. 5. :	3. 2. 2.
COEFFICIENTS OF CUBIC PULYNUMIAL/STANDARD D 0.10438 00 0.60358 C1 -C.80868 03 0.5197	
C.G. VEL. X-COMP. (M/5): 227. 221. 210	
COEFFICIENTS OF CUBIC PULYNUMIAL/STANUARD DE	EVIATION
-0.1143E 00  6.228LE UJ -0.352JE 64 -0.8422	(E U3 / 6.0651 (M)
PONCELET DRAG CUEFF. = 6.506 VD = 210. Stand. Devia. = 0.0064 (m)	
C.G. VEL. X-CUMP. (M/S): 236. 222. 210	. 196. 183.
RECORDED TIME OF MAXIMUM/MINIMUM COLL VOLTAC MAX .000149 .002376 .003929	SE (S) .005419
MAX .000145 .002376 .003929 MIN .001005 .003332 .004891 COMPUTED NOSE PUSITION AT MAX/MIN COIL VOLTA	A006503
AT MAX 0.025 0.511 0.527	1.110
AT MIN 0.217 0.708 1.012 Recorded Coil Pusition (M)	1.302
	1.086
DIFFERENCE BETWEEN CUIL AND NOSE AT MAX/MIN AT MAX 0.025 0.625 0.617	VOLTAGE (M) 0.024
AT MIN 0.217 0.222 0.202	0.216

-02

# SHUT 88 ( 24 MAY, 1976, NO. 8 )

SAND: DRY, DENSITY:1538. KG/M##3; APPROACHING VELOCITY: 233. M/S Projectile: Hollow Step-tiermass:0.4219 kg. D=0.02 M. L=0.238 M

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4555 A

A-RAY STATION	NG • 1	NO+ 2	NG3	NO.4	ND . 5
TIME (SECOND)			.062919	. 005792	.010028
CENTER OF GRAVITY FUSIT					
HGHIZONTAL	18.022	0.080	u • 40u	C.723	1.051
VERTICAL	18.224	<b>U.117</b>	V.122	0.124	0-108
INCLINATION ANGLE(DEG).	****	G.G '	~2.8	-7.3	-7.5
SEPARATION ANGLE (DE GREE	)				
ABOVE	****	****	****	1.0	2.5
BELUN	****	****	****	6.5	9.0
NOSE HIDTH (M ON FILM).	*****	u. U2.3u	0.0230	0.0230	0.0230
NUSE POSITIUN (M)					
HCHIZÜNTAL	18.046	0.185	0.565	0.828	1.155
VERTICAL	18.121	6.117	0.117	0-110	0.095
INPUT NOSE PUSITION (M)					
HCRICONTAL	* ** * **	0.178	0.501	0.026	1.157
VERTICAL	*****	-0.102	-0.102	-0.11U	-0.128

203

SHUT 89 ( 24 MAT. 1976, ND. 2 )

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SAND: DRY. DEN	SITY:1536. KG/M## Llow Step-tlermas	S; APPROACHING VE	ELUCITY: 227. M/S
	••••••••••••••••••••••••••••••••••••••		
CENTER OF GRAVI			
		0.008 0.393 0.119 0.125	
INCLINATION AND SEPARATION AND	LE(DEGRES)		
A80VE		6.0 **** 5.0 ****	**** ****
NOSE JIDTH (N C	ON FILMI. G.U200	6.0235 0.0230	0.0235 0.0230
HUR IZUNTAL		J.173 J.496	0.818 1.146
VERTICAL ++++ INPUT NUSE PUSI		3.118 0.110	0.095 0.064
HCRIZONTAL	LUUS		3.815 1.146
VENIIÇAL +++4			-9120 -9103
NOSE VEL . Y-CON	MP. (M/S): 2.	-64.	-7, -6.
Jelisse og s	F CUBIC POLYNUHIAL 3.2012E 61 -0.1297	/STANDARD DEVIAT 'E 04 J.5877E 05	10N 5 / 0.4027 (M)
NOSE VEL . X-COM	MP. (M/S): 215.	191. 137.	89. 77.
	F CUBIC POLYNUMIAL U-21995 03 -U-1093		
	CTIUN(GEG) G.4 Le(degree), relati		-4.6 -4.7
ABOVE			**** ****
HELOW		4.6 4.44	**** ****
C.G. VEL. Y-CCM	MP. (M/5): 9.	61.	-68.
	F CUBIC POLYNUMIAL 0.9566E u1 -0.02160		
	MP. (M/S): 215.		
COEFFICIENTS OF	F CULIC PULYNUMIAL	/STANDARD DEVIAT	ION
	0.2263E 03 -0.1693	E 05 V.6505E 06	) / U+9664 (M)
PUNCELET DRAG C	CUEFF. = 1.734 NU. DEVIA. = 0.612	6 (M)	
	4P. (M/5): 215.		97. 69.

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SHUT 90 1 24 MAY+ 1976, NO. 3 )

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X-RAY STATION ...... NU-1 NO-2 N0.3 NO.4 NQ.5 TIME (SECOND) · UU0146 .000700 .001708 .002800 .004354 18.213 0.577 0.885 1.174 6.136 VERTICAL ...... 0.137 0.135 4.121 18.224 INCLINATION ANGLE(DEG). SEPARATION ANGLE(DEGREE) ABOVE 0.4 -3.0 \*\*\*\*\* -3-0 -2-2 \*\*\*\* 7.0 2.0 \*\*\*\* \*\*\*\* 7.0 0.0235 .... 11.0 \*\*\*\* BELOW .... 0.0230 0.0220 \*\*\*\*\* 0.990 1,279 18-231 0.682 6.121 18.121 U.130 0-132 0.131

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u.704

-0.087

1.293

-0.087

1.017

-0.085

0.054

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SANC: WET, DENSITY:2053, KG/M++3; APPROACHING VELOCITY: 445, M/S Projectile: Hollow Step-Tiermass:0,4220 Kg, D=3.02 N; L=0.238 M

# SHUT 91 1 24 MAY. 1976, ND. 4 3

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SANCIN-UH. DENSITYIIQUG. KGZMAAGG APPROACHING VELUCITYI 230. M/S Projectile: Hollub Siep-tiermassiu.4222 kg. Dru.02 m. Leu.238 m.

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#### APPENDIX B

## DERIVATION OF FORCES AND MOMENTS ON A RIGID BODY IN A SOIL MEDILA

The purpose of this section is to develop the expressions for the forces and moments applied to a projectile in a soil medium. The forces are represented by  $F_X$ ,  $F_y$  and  $F_z$  along a body fixed set of axes x,y,z with the origin at the center of mass of the vehicle. The moments are about the x,y and z axes and are denoted L, M and N respectively and the velocity components are designated U, V, W respectively. It is assumed that the force exerted by the soil on a differential surface area of the body is of the form

$$\frac{dF}{dA} \propto (A_{x}+B_{x}|U|+C_{x}U^{2})n_{x}\hat{1} + (A_{y}+B_{y}|V|+C_{y}V^{2})n_{y}\hat{1} + (A_{z}+B_{z}|W|+C_{z}W^{2})n_{z}\hat{k} , \qquad (B-1)$$

where  $n_{x}$ ,  $n_{z}$  are components of the outward unit vector n normal to the surface of the body. The force is assumed to exist only when the velocity is directed toward the elemental area. The flow is assumed to be separated from the body aft of the nose.

The node shape that is treated is conical, with nomenclature as shown in Figure B-1. In the cylindrical coordinate system the equation for the surface of the conical nose is given by

$$\rho = \frac{r}{L_N} (\bar{x} - x) \quad . \tag{B-2}$$

The unit vector normal to the surface of the nose is

$$n = Sin \gamma i + Cos \gamma Cos \beta j + Cos \gamma Sin \beta k$$
. (B-3)

The area of a differential surface element is given by

$$dA = \frac{\rho \, dx \, d\beta}{\cos \gamma} \quad . \tag{B-4}$$

The inertial reference frame x'y'z', a rectangular Cartesian system, is choser with x' perpendicular to the surface of the target and pointing inward and y' and z' chosen in any convenient direction. In order to write the equation of the target surface prior to impact, in the body fixed axis system, a coordinate system  $x_1, y_1, z_1$  is chosen oriented parallel to the body axis system but whose origin is coincident with that of the inertial axis system. The equation of the target plane in the inertial system is x' = 0. Using the Euler transformation matrix introduced in Section VI the relation between the primed and sub-one system can be written as

$$\begin{pmatrix} x' \\ y' \\ z' \end{pmatrix} = T_{BI} \begin{pmatrix} x_1 \\ y_1 \\ z_1 \end{pmatrix} .$$
 (B.5)

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This yield?

$$x' = T_{11}x_1 + T_{12}y_1 + T_{13}Z_1 = 0, \qquad (B-6)$$

where

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$$T_{11} = \cos \Psi \cos \theta$$
  

$$T_{12} = \cos \Psi \sin \theta \sin \phi - \sin \Psi \cos \phi$$
  

$$T_{13} = \cos \Psi \sin \theta \cos \phi + \sin \Psi \sin \phi$$

Equation (B-6) can be written in the body axis system by a simple translation

$$x = x_1 + d$$
  
 $y = y_1$  (B-7)  
 $z = z_1$ ,

where d is the distance from the body axis origin to the inertial axis origin and is

$$d = \sqrt{(x')^2 + (y')^2 + (z')^2} ,$$

Substituting Equation (B-7) into Equation (B-6) results in

$$T_{11}x + T_{12}y + T_{13}z = T_{11}d$$
, (B-8)

the equation of the target plane in the body axis system. It is convenient to write this equation in a cylindrical body axis system as shown in Figure B-1.

$$T_{11}x + T_{12}\rho \cos\beta + T_{13}\rho \sin\beta = T_{11}d.$$
 (B-9)

Consider the case where the nose has impacted the target and is partially immersed in the soil as shown in Figure B-2. The locus of the intersection of the nose and the target plane will be given by the simultaneous solution of Equation (B-9) and Equation (B-2) which yields

$$x = \frac{T_{11}d - (r/L_N)\bar{x}[T_{12}\cos\beta + T_{13}\sin\beta]}{T_{11} - (r/L_N)[T_{12}\cos\beta + T_{13}\sin\beta]} = f_1(x',y',z',\Psi,\theta,\phi,\beta), \quad (B-10)$$

Using Equation (B-1) the differential force in the x direction becomes

$$dF_{x} = (A_{x} + B_{x}|U| + C_{x}U^{2})n_{x} dA$$

and realizing that a positive value of U will produce a negative force on the projectile, this force is rewritten as

$$dF_{x} = -\frac{U}{|U|} (A_{x} + B_{x}|U| + C_{x}U^{2})n_{x} dA$$
.

By using Equations (B-2), (B-3) and (B-4) and  $\frac{r}{L_N}$  = tany the above expression may be written as

$$dF_{x} = -\frac{U}{|U|} (A_{x} + B_{x} |U| + C_{x} U^{2}) (r/L_{N})^{2} (\bar{x} - x) dx d\beta,$$
  
integration wields

and upon integration, yields

$$F_{x} = -\frac{U}{|U|} (A_{x} + B_{x} |U| + C_{x} U^{2}) (r/L_{N})^{2} \int \int (\bar{x} - x) dx d\beta$$

where the limits for the integrals are chosen so that the integration is over the portion of the submerged area having a velocity component toward that surface. The integration with respect to x is performed first. The upper limit for the x integral is simply  $\bar{x}$ . The lower limit for x is a bit more complicated as shown by Figure B-3 which shows the x-z plane for a case where  $\bar{x} = \phi = 0$ . As long as the projectile has not penetrated very far the lower limit is simply given by the equation of the intersection of the cone and the target plane, that is  $f_1$  of Equation (B-10). As shown in Figure B-3 when the penetration reaches the base of the nose this limit is no longer correct since integration where  $x < L_1$  gives a contribution over the dashed part of the cone that does not exist. For this reason the lower limit becomes  $f_1$ , unless  $f_1 < L_1$  in which case the lower limit becomes  $L_1$ .

The limit on the  $\beta$  integral will depend upon the magnitude of the angle of attack,  $\alpha$  given below, as compared with the cone half angle  $\gamma$ . If the angle of attack is less than  $\gamma$  the limit will be from 0 to  $2\pi$ . This is expressed as

$$\alpha = \cos^{-1}\left[\frac{|u|}{\sqrt{u^2 + v^2 + w^2}}\right] < \gamma = \tan^{-1}\left[\frac{r}{L_N}\right].$$

To summarize

$$F_{x} = -\frac{U}{|U|} (A_{x} + B_{x} |U| + C_{x} U^{2}) \left[\frac{r}{L_{N}}\right]^{2} \int_{0}^{2\pi} f_{1}^{\overline{x}} (\overline{x} - x) dx d\beta , \qquad (B-11)$$

for  $\alpha < \gamma$  and for  $\overline{x} > \hat{x}_1 > L_1$ . For the case where  $\alpha > \gamma$  one side of the cone will be blanketed and the  $\beta$  integral will be integrated only over 180 degrees. In order to determine the proper range of  $\beta$  consider the y-z coordinates (x into the plane of the paper) as shown in Figure B-4. The V and W components determine the projected or wetted area in contact with the soil. Let  $\delta_t = \tan^{-1} \frac{W}{V}$ , and limits on  $\beta$  will be  $\pm \frac{\pi}{2}$  radians from the angle  $\delta_t$ or from  $-\frac{\pi}{2} + \delta_t$  to  $\frac{\pi}{2} + \delta_t$ . Collecting this we write

$$F_{x} = -\frac{U}{|U|} (A_{x} + B_{x} |U| + C_{x} U^{2}) \left[\frac{t}{L_{N}}\right]^{2} \int_{-\frac{\pi}{2} + \delta_{t}}^{\frac{\pi}{2} + \delta_{t}} \int_{f_{1} \ge L_{1}}^{\frac{\pi}{2} + \delta_{t}} \int_{f_{1} \ge L_{1}}^{\frac{\pi}{2} + \delta_{t}} (x - x) dx d\beta , \qquad (B-12)$$

for

 $\alpha > \gamma$ ,  $\overline{x} \ge f_1 \ge L_1$  and  $V \ge 0$ .

If V<0 a different set of limits on  $\beta$  are applicable. In this case

$$F_{x} = \frac{U}{|U|} (A_{x} + B_{x} |U| + C_{x} U^{2}) \left[\frac{r}{L_{N}}\right]^{2} \int_{\frac{\pi}{2} + \delta_{t}}^{\frac{3\pi}{2} + \delta_{t}} \int_{f_{1} \ge L_{1}}^{\frac{\pi}{2}} (\bar{x} - x) dx d\beta , \qquad (B-13)$$

for  $\alpha > \gamma$ ,  $x \ge f_1 \ge L_1$  and V < 0.

Once the nose has penetrated so that it is totally submerged the force becomes

$$F_{x} = -\frac{U}{|U|} (A_{x} + B_{x}|U| + C_{x}U^{2}) \frac{\pi r^{2}}{n}$$
(B-14)

where n = 1 for  $\alpha < \gamma$  and n = 2 for  $\alpha > \gamma$ .

In developing expressions for  $F_y$  and  $F_z$  the velocity components in the y and z directions would be made up of a velocity term due to translation plus a term due to rotation. Specifically, for a point on the axis of symmetry the velocity component in the y direction is V + Rx and the velocity component in the y direction is V + Rx and the velocity component in the z direction of x. If these velocity terms are used, the angle  $\delta_t$  becomes a function of x, namely  $\delta_t = \tan^{-1}[(V+Rx)/(W-Qx)]$ . Since  $\delta_t$  appears in the limits of integration with respect to  $\beta$ , these limits now become functions of x, and the integration is extremely cumbersome.

One approach to this problem is to treat the two velocity terms separately and add the resulting forces. This neglects the cross product terms 2VRx and -2WQx in the velocity squared terms of the differential force in Equation (B-1). Another approach would be to sum the two velocity terms first and replace the longitudinal integration with respect to x by integrations over several strips of finite length  $\Delta x$ . Then sum these segmental integrations to obtain the resultant forces. Only the first method is presented here, and further consideration of this problem is left for later research in which the possiblity of modifying the assumed force law should also be considered.

The expression for  $F_y$  is written as the sum of two parts. The first is due to translation,  $(F_y)_T$ , and is obtained in a manner similar to the detailed procedure outlined above for  $F_x$  and the other is a component of force due to a rotation rate R radians per second about the z-axis. This component is denoted  $(F_y)_T$ . The rotation rate will give the differential surface element a velocity component in the y direction equal to Px. This component is treated in the same manner as the translational velocity V in order to determine the  $F_y$  that it produces. The results are as follows:

$$F_y = (F_y)_T + (F_y)_r$$
 (B-15)

$$(\mathbf{F}_{\mathbf{y}})_{\mathrm{T}} = -\frac{\mathbf{V}}{|\mathbf{V}|} (\mathbf{A}_{\mathbf{y}} + \mathbf{B}_{\mathbf{y}} |\mathbf{V}| + \mathbf{C}_{\mathbf{y}} \mathbf{V}^{2}) \left[\frac{\mathbf{r}}{\mathbf{L}_{\mathrm{N}}}\right] \int_{-\frac{\pi}{2} + \delta_{\mathrm{T}}}^{\frac{\pi}{2} + \delta_{\mathrm{T}}} \int_{\mathbf{f}_{1} \ge \mathbf{L}_{1}}^{\frac{\pi}{2} + \delta_{\mathrm{T}}} \int_{\mathbf{f}_{1} \ge \mathbf{L}_{1}}^{\frac{\pi}{2} + \delta_{\mathrm{T}}} (\mathbf{x} - \mathbf{x}) d\mathbf{x} \cos\beta d\beta \qquad (B-16)$$

for V>0 and  $\bar{x} \ge f_1 \ge L_1$ . For the case where V<0 the limits on the  $\beta$  integral become  $\frac{\pi}{2} + \delta_t$  to  $\frac{3\pi}{2} + \delta_t$ .

$$(F_{y})_{x} = -\frac{R}{|R|} \left[ \frac{r}{L_{N}} \right] \int_{\frac{\pi}{2} + \delta_{x}}^{\frac{\pi}{2} + \delta_{x}} \int_{f_{1} \ge L_{1}}^{\frac{\pi}{2} + \delta_{x}} (A_{y} + B_{y} \times |R| + C_{y} \times^{2} R^{2}) (\bar{x} - x) dx \cos \beta d\beta \quad (B-17)$$

for  $R \ge 0$ ,  $\bar{x} \ge f_1 \ge l_1$  and  $\delta_r = -\tan \frac{-1Q}{R}$ . For R < 0 the limits on the  $\beta$  integral become  $\frac{\pi}{2} + \delta_r$  to  $\frac{3\pi}{2} + \delta_r$ .

Once the projectile has penetrated so that the nose is totally submerged the lower limit on all of the x integrals becomes  $L_1$ .

In a similar manner 
$$F_z$$
 is composed of  $(F_z)_T$  and  $(F_z)_T$ .  
 $F_z = (F_z)_T + (F_z)_T$ . (B-18)

Where

$$(\mathbf{F}_{z})_{T} = - \frac{W}{|W|} \begin{bmatrix} \mathbf{x} \\ \mathbf{L}_{N} \end{bmatrix} (\mathbf{A}_{z} + \mathbf{B}_{z} |W| + C_{z} W^{2}) \int_{-\frac{\pi}{2} + \delta_{t}}^{\frac{\pi}{2} + \delta_{t}} \int_{\mathbf{x}}^{\frac{\pi}{2} + \delta_{t}} (\bar{\mathbf{x}} - \mathbf{x}) d\mathbf{x} \sin \beta d\beta \qquad (B-19)$$

for  $V \ge 0$  and  $\bar{x} \ge f_1 \ge L_1$ . For V<0 the limits on the  $\beta$  integral become  $\frac{\pi}{2} + \delta_t$  to  $\frac{3\pi}{2} + \delta_t$ .

The component due to rotation is

$$(F_{z})_{r} = \left| \frac{Q}{Q} \right| \left[ \frac{r}{L_{N}} \right] \int_{-\frac{\pi}{2}+6}^{\frac{\pi}{2}+6} \int_{r}^{\bar{x}} \int_{1-L_{1}}^{(A_{z}+B_{z}x)} |Q| + C_{z} x^{2} Q^{2} (\bar{x}-x) dx \sin \theta d\theta \qquad (B-20)$$

for R>9 and  $\overline{x} \ge f_1 \ge L_1$ . For R<0 the limits on the  $\beta$  integral become  $\frac{\pi}{2} + \delta_r$  to  $\frac{3\pi}{2} + \delta_r$ .

Once the projectile has penetrated so that the nose is totally submerged in the target the lower limit on the x integrals becomes  $L_1$ .

The moments follow from the fact that there is a force on the differential element of area which is located a distance x from the center of mass. See Figure B-6. The force  $dF_y$  applied to the area dA will produce a moment about the z-axis, denoted N, whose magnitude will be x  $dF_y$ .

The moment N will be made up of two parts, one due to translation,  $(N)_T$  and one due to rotation rate  $(N)_r$ .

$$N = (N)_{T} + (N)_{T}$$
 (B-21)

The results for  $(N)_T$  are as follows:

$$(N)_{T} = -\frac{V}{|V|}(A_{y}+B_{y}|V|+C_{y}V^{2}) \left[\frac{r}{L_{N}}\right] \int_{-\frac{\pi}{2}+\delta_{t}}^{\frac{\pi}{2}+\delta_{t}} \int_{f_{1}\geq L_{1}}^{\frac{\pi}{2}+\delta_{t}} x(\bar{x}-x) dx \cos\beta d\beta \qquad (B-22)$$

for  $V \ge 0$  and  $x \ge f_1 \ge L_1$ . For V<0 the limits on the  $\beta$  integral become

 $\frac{\pi}{2}$  to  $\frac{3\pi}{2}$  to  $\frac{3\pi}{2}$ .

The values of moment due to rotation rate are

$$(N)_{r} = -\frac{R}{|R|} \left[\frac{r}{L_{N}}\right] \int_{-\frac{\pi}{2}+\delta}^{\frac{\pi}{2}+\delta} \int_{r}^{\overline{x}} \int_{f_{1}\geq L_{1}}^{\frac{\pi}{2}+\delta} (A_{y}+B_{y}x|R|+C_{y}x^{2}R^{2})x(\overline{x}-x) dx \cos\beta d\beta \quad (B-23)$$

for R>0 and  $\bar{x} \ge f_1 \ge L_1$ . For R<0 the limits on the  $\beta$  integral become

 $\frac{\pi}{2}+\delta_r$  to  $\frac{3\pi}{2}+\delta_r$ .

When the complete nose is submerged in the target the lower limit on the x integrals becomes  $L_1$ .

In a similar manner the moment about the y-axis,

$$M = (M)_{T} + (M)_{r}$$
 (B-24)

$$(\mathbb{M})_{T} = \frac{W}{|W|} \left[ \frac{r}{L_{N}} \right] (\mathbb{A}_{z} + \mathbb{B}_{z} |W| + \mathbb{C}_{z} W^{2}) \int_{-\frac{\pi}{2} + \delta_{t}}^{\frac{\pi}{2} + \delta_{t}} \int_{f_{1} \ge L_{1}}^{\tilde{x}} x(\tilde{x} - x) dx \operatorname{Sing} d\beta$$
 (B-25)

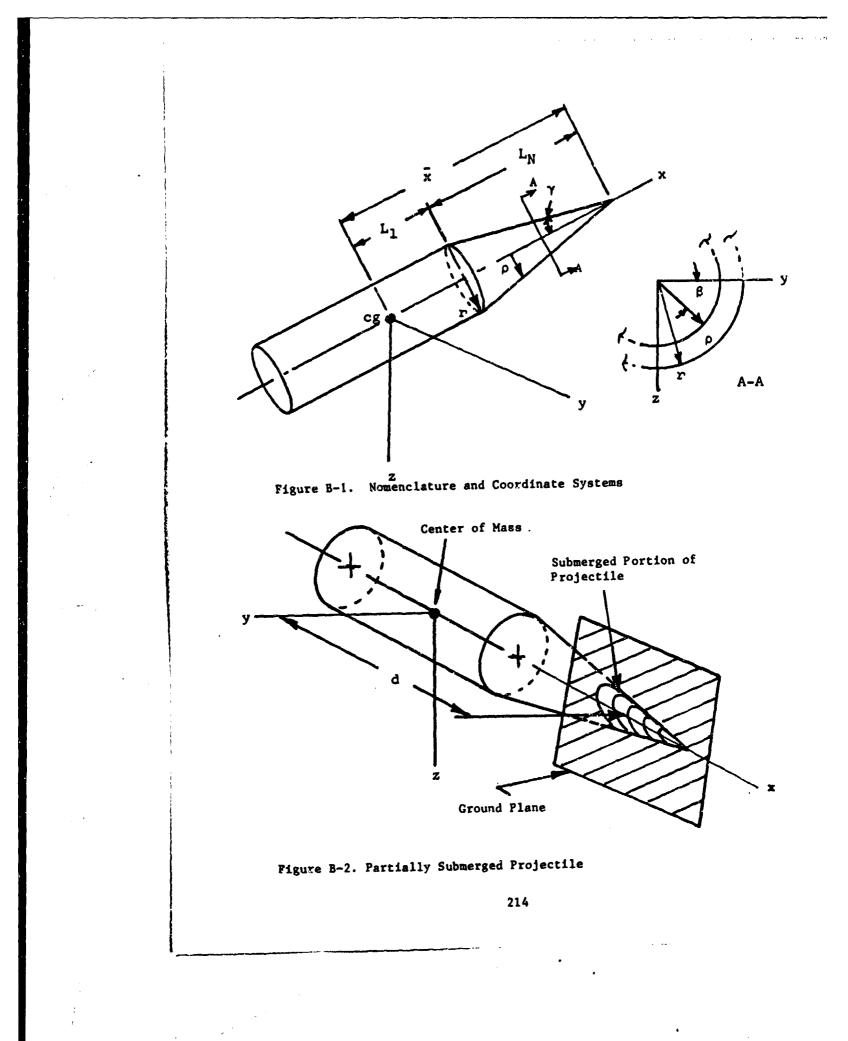
when  $V \ge 0$ ,  $\overline{x} \ge f_1 \ge L_1$ . For V<O the limits on the  $\beta$  integral become

when  $R\geq 0$  ,  $\bar{x}\geq f_1\geq L_1$  . For R<0 the limits on the  $\beta$  integral become

 $\frac{\pi}{2} + \delta_r$  to  $\frac{3\pi}{2} + \delta_r$ .

When the nose is completely submerged the lower limit on the  ${\bf x}$  integrals becomes  ${\bf L}_{j}$  .

Under the assumptions of this force model the moment about the xaxis, L, will be equal to zero for reasonable angles of attack and obliquity.



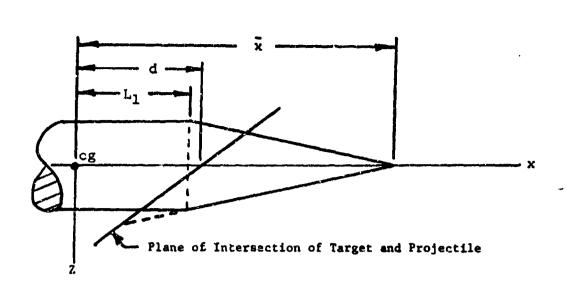
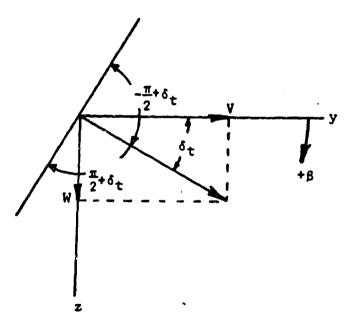


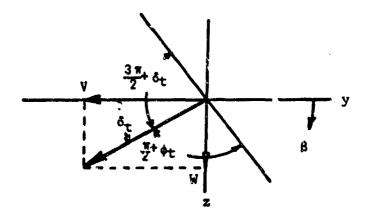
Figure B-3. Schematic Showing Plane of Intersection of Target and Projectile



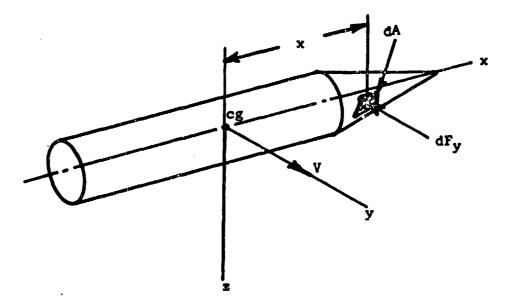


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APPENDIX C

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1.	LIST OF CO	IMPUTER SYMBOLS
	AX	Force coefficient, A <sub>x</sub>
	AY	Force coefficient, A
	AZ	Force coefficient, A
	BX	Force coefficient, B
	BY	Force coefficient, B
	BZ	Force coefficient, B
	œ	Drag coefficient, C <sub>D</sub>
	CMX	x coordinate center of mass
	<b>CM</b>	y coordinate center of mass
	CMZ	z coordinate center of mass
	CX	Force coefficient, C <sub>x</sub>
	CY	Force coefficient, C
	CZ	Force coefficient, C <sub>z</sub>
	DT	Print time frequency
	DTMAX	Maximum integration time
	DTMIN	Minimum integration time
	FIN	MIMIC finish statements
	FX	Total force component x direction, $F_{x}$
	fy	Total force component y direction, Fy
	FYCO	Force component y direction for cone only
	FYCY	Force component y direction for cylinder only
	Fyhe	Force component y direction for hemisphere only
	FZ	Total force component z direction, Fg
	FZCO	Force component z direction for cone only
	FZCY	Force component a direction for cylinder only
	YZHE	Force component z direction for hemisphere only
	HDR	Output headings
	INT	MIMIC integration symbol
	IXX	Moment of inertia, I
	IYY	Moment of incitia, I <sub>yy</sub>
	K1,K2,K3	Constants used to select proper force terms
	<b>K4,K5,K6</b>	Constants used to select proper moment terms
	L	Total applied moment about x axis

LCO	Applied moment about x axis due to cone only
LCX	Naturai logarithm of C <sub>X</sub>
LCY	Applied moment about x axis due to cylinder only
LHE	Applied moment about x axis due to hemisphere only
LNCY	Length of cylinder
LNCO	Length of nose cone
м	Total applied moment about y axis
MAS	Mass of projectile
MCO	Applied moment about y axis due to cone only
MCY	Applied moment about y axis due to cylinder only
MHE	Applied moment about y axis due to hemisphere only
N	Total applied moment about z axis
OUT	Output to be listed
NCO	Applied moment about z axis due to cone only
NCY	Applied moment about z axis due to cylinder only
P	Rotation velocity about x axis
PD	Rotation acceleration about x axis, P
PI	Universal constant m = 3.14159
PHI	Euler angle ¢
PHID	Euler angular velocity, ø
PHIO	Initial Euler angle, ¢ <sub>o</sub>
PO	Initail rotational velocity, P <sub>o</sub>
PRESU	Pressure exerted by velocity U
PRESV	Pressure exerted by velocity V
PRESW	Pressure exerted by velocity W
PSI	Euler angle, ψ
PSID	Euler angular velocity, $\psi$
PSIO	Initial Euler angle, ψ <sub>o</sub>
Q	Rotational velocity about y axis
QD	Rotational acceleration about y axis, Q
QO	Initial angular velocity about y axis, Q <sub>o</sub>
R	Rotational velocity about z axis
RAD	Radius of cylinder, base of nose cone, hemisphere
RD	Rotational acceleration about z axis, R
RO	Initial angular velocity about z axis, R <sub>o</sub>
Т	Time

TERM	1-u/u <sub>0</sub>
тн	Euler angle 0
THO	Euler angular velocity, $\hat{\theta}$
тно	Initial Euler angle, θ <sub>ο</sub>
U	Velocity relative to x axis
UD	Acceleration relative to x axis, U
UO	Initial velocity relative to x axis, Uo
v	Velocity relative to y axis
VD	Acceleration relative to y axis, V
vo	Initial velocity relative to y axis, $V_{o}$
W	Velocity relative to z axis
WD	Acceleration relative to z axis, W
WO	Initial velocity relative to z axis, W <sub>o</sub>
XP	x' axis
XPD	Velocity relative to x' axis
YP	y' axis
YPD	Velocity relative to y' axis
ZP	z' axis
ZPD	Velocity relative to z' axis

#### 2. INTRODUCTION

The necessary format and language for programming is contained in Control Data Corp. reference manual entitled Control Data MIMIC and only the details for datainput are included here. All input data either constants (CON)  $\gamma$ r parameters (PAR) are placed at the end of the program and correspond exactly to the manner in which they are called for by the CON and PAR cards preceding the program. Details of input data are shown for Computer Frogram II.

In the first case shown, Computer Program I, the force coefficients  $A_x$  through  $C_z$  are assumed to be independent of the velocity terms, therefore all these values are shown as parameters. However, in the second case, Computer frogram II, the value of  $C_x$  is then given a variable status included in the body of the program.

Parameters K1 through K6 are used to select forces and moments depending on type of nose cone used as well as the option not to use the forces and moments on the after body. These parameters take on values of zero or unity and are defined as follows:

Configuration 1	KI	K2	K3	K4	K5	<u>K6</u>
	1	0	0	1	0	0
Cylinder only, all moments and forces						
Cylinder plus cone, all moments and forces	1	1	0	1	1	<u>)</u>
Cylinder plus cone, only nose moments and forces	0	1	0	1	0	0
Cylinder plus hemisphere, all moments and forces	1	0	1	0		1
Cylinder plus hemisphere, only nose moment and forces	0	0	1	0	0	1

3. COMPUTER PROGRAM I

SOIL PENET	CCN(D., DTHIN, DTHAX) PAR(AX, BX, CX, AY, BY, CY) FAR(AZ, UZ, CZ) FAR(AZ, UZ, CZ) FAR(AS, IXX, IYY, CHX, CFY, CMZ) FAR(LAUY, LNCO, RAU) FAR(UU, JU, HU, PU, QU, PO) FAR(FHIU, THU, PSIO, XPO, YPO, ZPO) FAR(K1, K2, K3, K4, K5, KE) RATLUN CUMPLETELY SUBME=GED
FURCE FERM PRESU PRESU PRESU FRZUCO FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE FRZUCE F	3.14127. - (AX+UX+A3S(U)+CX+U+U)+U/ABS(U) - (AY+UY+A3S(V)+CY+V+V)+V/ABS(V) - (AZ+UZ+ABS(W)+CZ+W+W)+W/ABS(W) FI+RAC+PAJ+PRESU 2.*RAC+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+PRESU 2.*RAU+L4CY+R2+FZCO+K3+FYHE K1+FZCY+K2+FZCO+K3+FZHE

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FORCE EGUATIONS JD FX/MAS-Q+W+R+V
VO FY/NAS-RYU+PHN
V TNT(VC,VO) ND FZ/MAS-F+V+Q+U
INT (HU-HO)
MONENT EQUATIONS
PO L/1XX INT(PO,PO)
$\frac{1}{2}$
EULE ANGLE EQUAIIONS
PHID1 F+Q+SIN(PH1)+SIN(TH)/CO5(TH)
PHID2 R-COS(PHI)+SIN(TH)/COS(TH) PHID PHID1+PHID2
PHĪ INĪ (ĒHIO.PHIO)
THĂ Ā*CÓS(PHÍ)-(*SIN(PHI)
ΪΗ΄ ΙΝΤ(ΤΗΟ, ΤΗΟ) PS 10 G +SIN (PHI)/COS(TH) +R + CQS (PHI)/COS(TH)
PSI INT(PSID, PSID)
INERTIAL FRAME EQUATIONS
xFJ1     U*CUS(PSI)*COS(TH)       xFU2     V*(CUS(PSI)*SIN(TH)*SIN(PHI)*SIN(PSI)*COS(PHI))
XF02 V+(CUS(PSI)+SIN(TH)+SIN(PHI)-SI4(PSI)+COS(PHI)) XF03 N+(CUS(FSI)+SIN(TH)+CUS(PHI)+SIN(PSI)+SIN(PHI))
XPJ XPD1-XPU2+XPD3
<pre>YP01 U*(SIN(PSI)*COS(TH)) YP02 V*(SIN(PSI)*COS(TH)) </pre>
<b>YPJ3 H+(SIN(PSI)+SIN(TH)+CUS(PHI)+CUS(PSI)+SIN(PHI))</b>
YPD YF01+YPC2+YP03
VP         INT(VPU, VPO)           ZFD         -U*SIN(TH)+V*(COS(TH)*SIN(PHI))+W*(COS(TH)*CUS(PHI))
ZFD = -U*SIN(FH)+V+(COS(TH)+SIN(PHI))+W+(COS(TH)+CUS(PHI)) ZY INT(ZFD,ZPO)
HUL ( ; ,XP),XP) Out ( ; ,XP) . XP)
FIH(T, U3)
END

4. COMPUTER PROGRAM II

Sec. Sec. Sec. 1

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Sec.

CCN(CT, DTHIN, OTHAX)	
F & R ( A Z, B Z, A Y, E Y, C Y) F & R ( A Z, B Z, C Z)	an an an an an an an an an an an an an a
<u>FAR (PAS, IXX, IYY, CHX, CPY, CH7)</u>	
FAR(LNCY,LNCO,RAD) FAR(UO,VO,HO,FG,GO,RG)	
SOIL PENETRATION COMPLETELY SUBHERGED	

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FORCE TERM	5 3 - 1 4 1 5 9
LCX CX	1.4114LOG(U) Exp(LCX)
PRESU PRESV	- (AX + EX + ABS(U) + CX + U + U) + U/ABS(U) - (AY + EY + ABS(V) + CY + V + V) + V/ABS(V)
FX FX EYCY	- (A2+82+A85(W)+C2+W+W)+W/A85(W) FI+RAD+RAD+PRESU 2.+RAD+LNCY+PRESU
FZČY FYCO	2.+RAC+LNCY+PRESW FAD+LNCO+PRESV -RAD+LNCO+PRESW
FYHE FZHE FY	F I+RAD+RAD+PRESV/2, F I+RAD+RAD+PRESW/2, K I+F YCY+K2+F YCO+K3+F YHE
FZ MOMENT_TERI	K1+F2CY+K2+F2CO+K3+F2FE
	K4+(FZGY+CHY-FYCY+CHZ)       K5+(FZC0+CHY-FYC0+CHZ)       HC+(FZHC+CHY-FYHC+CHZ)
4CV 4C0 4HE	- K4*FZGY*(CMX-LNCY/2.) - K5*FZCO*(CMX-2.*LNCO/3.) - K5*FZHE*(CMX-5.*RAU/2.)
NCO NCO NHE	K4#FYCY#(CMX+LNCY/2.) K5#FYC0#/CMX-2.#LNCO/3.) KE#FYHE#(CMX+5.#RAD/8.)
	LCY+LCO+LHE #EY+HEO+HHE-F##CHZ
FORCE ECUAT	FJ/MAS-UPWARTV
	INT(LD,UO) FY/NAS_R+UAP+H INT(VC,VO)
MO MOMENT, EQUA	FZ/HAS-P#V+Q+U INT(HDyHO) TIONS
<u> </u>	(/IXX IVI(FC,PO) H/IYY-R*P*(IXX-IYY)/IYY
	<u>INI(0[_00)</u> N/IYY+Q=P+(IXX-IYY)/IYY
-EULER ANGLE	P+Q+SIN(PHI)+SIN(TH)/COS(TH)
PHID PHI	R +CCS (PHI) +SIN(TH)/COS (TH) PFID (#PHID2 INT (PHID, PHID)
	C 4 C A S ( PH 1) - R # SIN ( PH I ) INT ( THO, THO) G #SIN ( PHI) / COS ( TH) + R # COS ( PHI ) / COS ( TH )
	##14F610y#510}

INERTIAL FFAME EQUATIONS XP01 U+COS(PSI)+COS(TH) XP02 V+(CCSTPSI)+SIN(TH)+SIN(PHI)+SIN(PSI)+COS(PHI) XP03 H+(CGS(PSI)+SIN(TH)+CCS(PHI)+SIN(PSI)+SIN(PHI)) XP0 1 V+(SIN(PSI)+SIN(TH)+CCS(PHI)+SIN(PHI)) YP01 U+(SIN(PSI)+COS(TH)) YP02 V+(SIN(PSI)+SIN(TH)+COS(PHI)+COS(PSI)+SIN(PHI)) YP03 H+(SIN(PSI)+SIN(TH)+COS(PHI)-COS(PSI)+SIN(PHI)) YP0 YFD1+YP02+YP03 YP IN(TYP0,YP0) YP0 -L+SIN(TH)+V+(COS(TH)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 -L+SIN(TH)+V+(COS(TH)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP IN(TYP0,YP0) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI)+COS(PHI))+H+(COS(TH)+COS(PHI)) YP0 +COS(PSI)+SIN(PHI)+COS(PHI)+COS(PHI))+H+(COS(PHI))+COS(PHI)) YP0 +COS(PSI)+SIN(PHI)+COS(PHI)+COS(PHI)+COS(PHI))+COS(PHI)) YP0 +COS(PSI)+CO) COS(PSI)+SIN(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI))+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI)+COS(PHI

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5. INPUT DATA FOR COMPUTER PROGRAM II  $4.06 \times 10^4$  $DT = 5.0 \times 10^{-4}$ υ<sub>o</sub> DTMIN 1.0 x 10<sup>-5</sup> 0.0 ٧ð DTMAX 1.0 x 10<sup>-3</sup> 0.0 WO PO 0.0 AX 0.0 0.0 **Q**<sub>0</sub> BX 0.0 0.0 RO AY 0.0 PHIO 0.0 BY 0.0 THO 0.0 CY 0.0 0.0 **PSIO** AZ 0.0 0.0 XPO BZ 0.0 YPO 0.0 cz 0.0 MAS 5.44 x  $10^2$ 0.0 ZPO IXX 2.687 x  $10^2$ 1.0 **K1** IYY 2.317 x  $10^4$ K2 0.0 0.0  $CHOX 1.127 \times 10^{1}$ K3 K4 1.0 CMY 0.0 0.0 K5 CMZ 0.0 LNCY 2.254 x  $10^1$ K6 0.0 LNCO 0.0 RAD 0.992

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