

# PHYSICS LETTERS

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## I. OVERVIEW

This review is an updating through November 1985 of the Review of Particle Properties [Particle Data Group (1984)], a compilation of experimental results on the properties of particles studied in elementary particle physics. These properties include masses, widths or lifetimes, branching ratios, and other experimentally determined properties. Where feasible, we provide a suggested "best" value of each parameter based on our own judgment, using the best available data. A discussion of some of the procedures that we apply, and a brief review of the historical performance of averages of measurements, may be found below (Section IV Part D).

The results of this compilation are presented in two sections, the "Summary Tables of Particle Properties" and the

"Full Listings." The Summary Tables give our estimates of the properties of those states whose existence we consider well established. Our opinion of whether or not a particle is well established can change as new data become available. We attempt to be conservative, so particles awaiting confirmation are not included, even if they may be theoretically well understood.

All data used for the numerical estimates in the Summary Tables are included in the Full Listings, with references and our comments, if any. Those measurements considered recent enough or important enough to mention, but which for some reason were not used in the averaging, appear in parentheses. The Full Listings also contain information on unconfirmed particles and unsuccessful particle searches, as well as short "mini-reviews" about subjects of particular interest or data that have particular problems.

In the past, we have attempted to use the Full Listings as an archive of all reported data on particles of interest. This is no longer possible because the growth of information would require a 5 to 10% per year expansion in this Review. Therefore we refer interested readers to previous editions for references to data considered obsolete.

This edition we are implementing our new particle naming conventions [Barnett (1985) and Wohl (1984)], which primarily affect meson names. A few baryon states are renamed as well. In the Summary Tables of Particle Properties and the Full Listings each particle is listed by its new name, with the old name, if different, given below it. It is our hope that these new conventions, described in Section III below, if adopted by the community will bring order to the chaos of particle names and facilitate discussion and understanding. Since there will doubtless be a transition period during which the literature may contain a mixture of both old and new names, we will continue to list the old names with the new for several editions.

We categorize the particles into types, intended to correspond roughly to the different types of data and problems encountered:

**STABLE PARTICLES** — All particles stable under the strong interaction. These include the truly stable particles as well as those which decay weakly or electromagnetically, including the  $\eta$ ,  $D$ ,  $D_s$  (formerly called the  $F$ ),  $\Lambda_c$ ,  $W$ ,  $Z^0$ , and so on.

**MESONS** — All meson resonances that decay strongly, including the  $\psi$ ,  $\chi$ , and  $\Upsilon$  families.

**BARYONS** — All baryon resonances that decay strongly, including the resonant  $N$  and  $\Delta$  families, dibaryon candidates, and so on.

This classification scheme is used to organize the Summary Tables and the Full Listings.

We include a section of "Miscellaneous Tables, Figures, and Formulae." These provide a quick reference for the practicing elementary particle physicist. They normally presuppose some understanding of the subject matter, and do not attempt to serve as a textbook. We welcome all suggestions and comments regarding topics for inclusion or deletion, any errors or confusing passages, etc.

A pocket-sized Particle Properties Data Booklet is available. This contains the complete Summary Tables of Particle Properties and the most frequently used parts of the Miscellaneous Section, but not the Full Listings. For North and South America, Australia, and the Far East, write to Technical Information Department, Lawrence Berkeley Laboratory, Berkeley, CA 94720, USA. For all other areas, write to CERN Scientific Information Service, CH-1211 Geneva 23, Switzerland.

In 1984 we began a multiyear effort aimed at modernization and reorganization. In this edition, we have added Greek letters and larger fonts for headings in the Full Listings, as well as numerous more minor improvements in format. We are also modernizing our internal procedures, and some of these improvements are already in place.

## II. AUTHORS AND CONSULTANTS

The primary responsibilities of the authors are as follows:

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(3) *Baryon resonances*: R.L. Crawford, G.P. Gopal, G. Höhler, D.M. Manley, L.D. Roper, and C.G. Wohl.

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The usefulness of this compilation depends in large part on interaction between the users and the authors and consultants. We appreciate comments, criticisms, and suggestions for improvements of all stages of data retrieval, evaluation, and presentation.

### III. A NEW NAMING SCHEME FOR HADRONS

*"Young man, if I could remember the names of these particles, I would have been a botanist."*

*Enrico Fermi*

#### A. The need for a new scheme; guiding principles

We introduce in this edition a new naming scheme for the hadrons. Anyone who doubts the desirability of a better naming scheme is invited to give, without looking at the Meson Summary Table, the quantum numbers  $I, J, P, C$ , and  $G$  of the following established nonstrange mesons:

$S(975), \delta(980), H(1190), B(1235), D(1285), \epsilon(1300), \pi(1300), E(1420), \iota(1440), \rho(1600), \omega(1670), A(1680), \phi(1680), g(1690), \theta(1690), h(2030).$

There is no rhyme or reason to this alphabet soup of symbols – they convey nothing about the properties of the particles they name. Nor is the use of five different symbols,  $K, K^*, Q, L$ , and  $\kappa$ , to name just nine strange mesons informative or economical. The symbols for mesons containing heavy quarks and for ordinary baryons are fairly sensible, but it seems wise, while in the grip of reformist zeal, to make some rules regarding names of particles yet to be discovered, such as the whole spectrum of baryons containing one or more heavy quarks.

There are several obvious virtues any rational naming scheme ought to embody. The symbols ought to be as few and as simple as possible, with those already in common use retained where possible; the symbols ought to convey unambiguously the important quantum numbers of the particles they name; and the quark model ought to guide the whole scheme. There are, however, constraints: it is not practical, for example, to now rename the  $\phi$  meson the  $\omega'$ , or to call the  $K$  meson containing an  $s$  quark (as opposed to an  $\bar{s}$ ) a  $\bar{K}$  instead of a  $\bar{K}$ . Some compromise between simplicity and long-established usage is unavoidable.

The new scheme adopted here has evolved over the last two years in response to much discussion both within the Particle Data Group and with the larger community. Preliminary versions of the scheme were presented at the 1984 Santa Fe Meeting of the Division of Particles and Fields [Wohl (1984)] and at the 1985 International Conference on Hadron Spectroscopy [Barnett (1985)]. Several thousand copies of the proposal, with an invitation to comment, were distributed in the spring of 1985. A *Physics Today* news report discussed the proposal [Schwarzschild (1985)], and it

has been discussed in the CERN *Courier* (November 1985).

As indicated above, many of the mesons have been renamed. The Meson Summary Table in this edition gives both the new and old names, and a table of equivalent names will appear in foreseeable editions. Only two particles in the Stable Particle and Baryon Summary Tables are renamed (the  $F$  and the  $A^+$  become the  $D_s$  and the  $\Xi_c^+$ ).

#### B. "Neutral-flavor" mesons ( $S = C = B = T = 0$ )

Table I shows the naming scheme for mesons having the strangeness and all heavy-flavor quantum numbers equal to zero. The mesons are assumed to be quark-antiquark states. The rows of the table give the possible  $q\bar{q}$  content. The columns give the possible parity/charge-conjugation states,  $PC = -+, +-, --, ++$ ; these combinations correspond one-to-one with the angular-momentum state  $2S+1L_J$  of the  $q\bar{q}$  system being  $^1(L \text{ even})_J, ^1(L \text{ odd})_J, ^3(L \text{ even})_J$ , or  $^3(L \text{ odd})_J$ .<sup>\*</sup> In addition, the spin  $J$  is added to the main symbol as a subscript except for pseudoscalar and vector mesons ( $L=0$  states in the quark model), and the mass is given for any meson that decays strongly.

Experimental determination of the mass, quark content, and quantum numbers  $I, J, P$ , and  $C$  (or  $G$ ) of a meson thus fixes its symbol. Conversely, these properties may be inferred unambiguously from the symbol.

If the main symbol cannot be assigned, because the quantum numbers (other than  $J$ ) are unknown, the symbol  $X$  is used temporarily. Sometimes it is not known whether a meson is mainly the isospin-0 mix of  $u\bar{u}$  and  $d\bar{d}$  or is mainly  $s\bar{s}$ ; the prime (or symbol  $\phi$ ) may be used to distinguish two such mixing states.

Names have been assigned for the anticipated  $t\bar{t}$  mesons. No suggestion is made here for names for mesons (should any be found) with the "exotic" quantum numbers that a  $q\bar{q}$  system cannot have, namely  $J^{PC} = 0^{--}, 0^{+-}, 1^{-+}, 2^{+-}, 3^{-+}, \dots$ . Glueballs or other mesons that are not  $q\bar{q}$  states would (if the quantum numbers are *not* exotic) be named just as if they were  $q\bar{q}$  states, since they will probably be difficult to distinguish from such states and will likely mix with them.

The results of all this are as follows. None of the lowest

Table I. Symbols for mesons with the strangeness and all heavy-flavor quantum numbers equal to zero.

$q\bar{q}$ content	$J^{PC} = \begin{cases} 0^{-+} & 1^{+-} & 1^{--} & 0^{++} \\ 2^{-+} & 3^{+-} & 2^{--} & 1^{++} \\ \vdots & \vdots & \vdots & \vdots \end{cases}$				
	$2S+1L_J =$	$^1(L \text{ even})_J$	$^1(L \text{ odd})_J$	$^3(L \text{ even})_J$	$^3(L \text{ odd})_J$
$u\bar{d}, d\bar{d} - u\bar{u}, d\bar{u} \ (I=1)$	$\pi$	$b$	$\rho$	$a$	
$d\bar{d} + u\bar{u}$ and/or $s\bar{s}$	$\eta, \eta'$	$h, h'$	$\omega, \phi$	$f, f'$	
$c\bar{c}$	$\eta_c$	$h_c$	$\psi^\dagger$	$\chi$	
$b\bar{b}$	$\eta_b$	$h_b$	$\Upsilon$	$\chi_b$	
$t\bar{t}$	$\eta_t$	$h_t$	$\theta$	$\chi_t$	

<sup>†</sup>The  $J/\psi$  remains the  $J/\psi$ .

<sup>\*</sup> The relations between the quantum numbers are

$$P = (-1)^{L+1} \quad C = (-1)^{L+S} \quad G = (-1)^{L+S+I},$$

where of course the  $C$  quantum number (charge conjugation) is only relevant to charge-zero mesons.

mass pseudoscalar or vector mesons ( $\pi$ ,  $\eta$ , and  $\eta'$ ;  $\rho$ ,  $\omega$ , and  $\phi$ ) change names, nor do any of the  $c\bar{c}$  or  $b\bar{b}$  mesons. Established mesons whose names change slightly are:

Old name	New name	Old name	New name
$H(1190)$	$h_1(1190)$	$A_2(1320)$	$a_2(1320)$
$B(1235)$	$b_1(1235)$	$f'(1525)$	$f_2(1525)$
$f(1270)$	$f_2(1270)$	$\omega(1670)$	$\omega_3(1670)$
$A_1(1270)$	$a_1(1270)$	$\phi(1850)$	$\phi_f(1850)$

Established mesons whose names change completely are:

Old name	New name	Old name	New name
$S(975)$	$f_0(975)$	$\iota(1440)$	$\eta(1440)$
$\delta(980)$	$a_0(980)$	$A_3(1680)$	$\pi_2(1680)$
$D(1285)$	$f_1(1285)$	$g(1690)$	$\rho_3(1690)$
$\epsilon(1300)$	$f_0(1300)$	$\theta(1690)$	$f_2(1720)$
$E(1420)$	$f_1(1420)$	$h(2030)$	$f_4(2030)$

The  $S(975)$ ,  $D(1285)$ ,  $\epsilon(1300)$ ,  $E(1420)$ ,  $\theta(1690)$ , and  $h(2030)$  all become  $f$  mesons; the new scheme reveals that all have  $PC=++$  and are  ${}^3(L \text{ odd})_J$  states.

### C. Mesons with nonzero $S$ , $C$ , $B$ , and/or $T$

Since the strangeness or a heavy flavor is nonzero, none of the mesons here are eigenstates of charge conjugation, and in each of them one of the quarks must be heavier than the other. The rules are:

(1) The main symbol is an upper-case Roman letter indicating the heavier quark as follows:<sup>‡</sup>

$$s \rightarrow \bar{K} \quad c \rightarrow D \quad b \rightarrow \bar{B} \quad t \rightarrow T$$

(2) If the lighter quark is not a  $u$  or a  $d$  quark, its identity is given by a subscript.

(3) If the spin-parity is in the "normal" series  $J^P = 0^+, 1^-, 2^+, \dots$ , a superscript "\*" is added.

(4) The spin is added as a subscript unless the meson is a pseudoscalar or a vector ( $L=0$  states in the quark model).

Thus the pseudoscalar and vector  $K$ ,  $K^*$ ,  $D$ ,  $D^*$ , and  $B$  mesons do not change names. Established mesons whose names do change are:

Old name	New name	Old name	New name
$Q_1(1280)$	$K_1(1280)$	$L(1770)$	$K_2(1770)$
$\kappa(1350)$	$K_0^*(1350)$	$K^*(1780)$	$K_3^*(1780)$
$Q_2(1400)$	$K_1^*(1400)$	$K^*(2060)$	$K_4^*(2060)$
$K^*(1430)$	$K_2^*(1430)$	$F$	$D_s$

<sup>‡</sup>Two different conventions exist in the literature for the sign of the flavor of  $b$  quarks. We have adopted the convention that the sign of the flavor of a quark is the same sign as its charge, which is true for all flavors. Thus the strangeness of the  $s$  quark is negative, the charm of the  $c$  quark is positive, and the bottom of the  $b$  quark is negative. In addition,  $I_3$  of the  $u$  and  $d$  quarks is positive and negative, respectively. The effect of this convention is as follows: Any flavor carried by a charged meson has the same sign as its charge. Thus the  $K^+$ ,  $D^+$ , and  $B^+$ , have positive strangeness, charm, and bottom, respectively, and all have positive  $I_3$ . The  $D_s^+$  (formerly the  $F^+$ ) has positive charm and strangeness. Furthermore, the  $\Delta(\text{flavor}) = \Delta Q$  rule, which is best known for the kaons, applies to every flavor.

The most notable change is that of the  $F$  (the  $c\bar{s}$  state) to a  $D_s$ . However, with the prospect of  $B_s$ ,  $B_c$ ,  $T_s$ , and similar mesons, there is no consistent and economical alternative. The rules can lead to cumbersome symbols, such as a  $D_s^*2$ , but such particles are unlikely to be often seen.

### D. Baryons

The symbols  $N$ ,  $\Delta$ ,  $\Lambda$ ,  $\Sigma$ ,  $\Xi$ , and  $\Omega$  have been used for 20 years for the baryons made of light quarks ( $u$ ,  $d$ , and  $s$  quarks), and no change is made to these symbols here. They tell the isospin and quark content, and the same information ought to be conveyed by the symbols used for the baryons containing one or more heavy quarks ( $c$ ,  $b$ , and  $t$  quarks). The following system was invented earlier and independently by Hendry and Lichtenberg (1978) and by Samios (1980). The rules are (see also Fig. 1):

(1) Baryons with three  $u$  and/or  $d$  quarks are  $N$ 's (isospin 1/2) or  $\Delta$ 's (isospin 3/2).

(2) Baryons with two  $u$  and/or  $d$  quarks are  $\Lambda$ 's (isospin 0) or  $\Sigma$ 's (isospin 1). If the third quark is a heavy quark (not an  $s$  quark) its identity is given by a subscript. This nomenclature is already used for the  $\Lambda_c(2281)$ ,  $\Sigma_c(2450)$ , and  $\Lambda_b(5500)$ .

(3) Baryons with one  $u$  or  $d$  quark are  $\Xi$ 's (isospin 1/2). One or two subscripts are used if one or both of the remaining quarks are heavy: thus  $\Xi_c$ ,  $\Xi_{cc}$ ,  $\Xi_b$ , etc. The possible but not established  $A(2460)$  is renamed the  $\Xi_c(2460)$ .

(4) Baryons with no  $u$  or  $d$  quarks are  $\Omega$ 's (isospin 0), and subscripts indicate any heavy-quark content. The possible but not established  $T(2740)$  is renamed the  $\Omega_c(2740)$ .

In short, the total number of  $u$  and  $d$  quarks together with the isospin determine the main symbol, and subscripts indicate any content of heavy quarks. A  $\Sigma$  always has isospin 1, an  $\Omega$  always has isospin 0, etc. The only baryons whose names change are the  $A$  and the  $T$ .

Note in Fig. 1 that the SU(4) 20-plet that contains the basic SU(3) octet has an  $\Omega_c$  and an  $\Omega_{cc}$  although it has no  $\Omega$ . It has two  $\Xi_c$ 's, which would be distinguished by mass (they might also be distinguished by a prime on the heavier of the two).

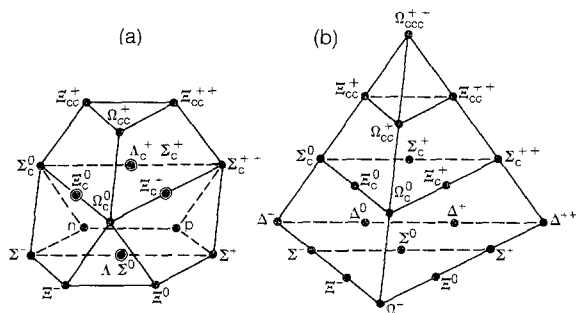


Fig. 1. SU(4) multiplets of baryons made of  $u$ ,  $d$ ,  $s$ , and  $c$  quarks. (a) The 20-plet with an SU(3) octet. (b) The 20-plet with an SU(3) decuplet.

## IV. PROCEDURES

### A. Selection and treatment of data

The Full Listings contain a complete record of all *relevant* data known to us. As a general rule, we do not include results from preprints or conference reports. It is our experience that preprinted results often change before publication. In some cases, such results may be cited but not used in computing the estimates given in the Summary Tables. There are a few exceptions to this exclusion, which we decide on a case-by-case basis after consultation with the experimenters.

As mentioned earlier, we no longer attempt to maintain an archival record of data of historical importance only. We do, however, quote the references of discoveries, even when the data are no longer useful.

If data are included in the Full Listings but not used in the final average given in the Summary Tables, they are enclosed in parentheses. We give explanatory comments in many such cases. If no comment is given, the reason the data were excluded is one or more of the following:

- The data are superseded or included in later results.
- No error was given.
- The data were contained in a preprint or conference report.
- The result involves some assumptions we do not wish to incorporate.
- The measurement has poor signal-to-noise ratio, low statistical significance, or is otherwise of much poorer quality than other data available.
- The measurement is clearly inconsistent with other results which appear to be highly reliable (see discussion in Section IV Part D below).
- The measurement is not independent of other results, e.g., it is from one of several partial-wave analyses, all of which use the same data, rendering averaging meaningless.

In some cases, *none* of the measurements is entirely reliable and no statistically meaningful average is quoted. For example, the masses of many of the baryon resonances, obtained from partial-wave analyses, are quoted as a range thought to probably include the true value rather than as an average with error. This is discussed in more detail in some of the mini-reviews in the Baryon Full Listings.

For upper limits, we normally quote in the Summary Tables the strongest limit available from a single experiment. We do not average or combine upper limits except in a very few cases where they may be re-expressed as measured numbers with Gaussian errors.

For quantum number assignments, we indicate in the Summary Tables those which are either well established or probable. In the Meson Summary Table, we underline those we consider well established; the others are inferred from whatever experimental evidence is available. In the Stable Particle Summary Table, nearly all quantum numbers are well established and we do not underline; those which are not well established are indicated by a footnote.

As is customary, we assume that antiparticles are the result of operating with *CPT* on particles, so both share the same spins, masses, and mean lives. There is an entry in the Miscellaneous Section, Tests of Conservation Laws, listing tests of *CPT* and other conservation laws.

### B. Criteria for new states

An experimentalist who sees indications of a new state will of course want to know what has been seen in that region in the past. Hence, we include in the Full Listings all reported states which have not been, in our opinion, disproved by better (e.g., more reliable) data.

For the Summary Tables we are much more conservative. We include only those reported states which we feel have a large chance of survival. One's betting odds for survival are of course subjective; therefore no precise criteria can be defined. For more detailed discussions, see the mini-reviews in the Full Listings. In what follows we shall attempt to specify some guidelines.

(a) When energy-independent partial-wave analyses are available (mostly for  $\pi N$  resonances), approximate Breit-Wigner behavior of the amplitude appears to us to be the most satisfactory test for a resonance. We can check that the Argand plot follows roughly a left-hand circle, and that the "speed" of the amplitude also shows a maximum near the resonance energy; further, there should be data well above the resonance, showing that the speed again decreases. Indeed, proper behavior of the partial-wave amplitude often establishes a resonance even if its elasticity is too small to make a noticeable peak in the cross section.

(b) When there are insufficient data to perform energy-independent analyses, one often resorts to energy-dependent partial-wave analyses. In this case Breit-Wigner behavior is an input. We usually require that resonance solutions be found by several different analyses, preferably in different channels ( $\bar{K}N \rightarrow \bar{K}N, \pi\Sigma$ , etc.), before putting the claim in the Summary Tables.

(c) Stable particles, most meson resonances,  $\Xi$  resonances, and high-mass  $N, \Delta, \Lambda$ , and  $\Sigma$  resonances fall into a category for which no partial-wave analyses exist. In general, we accept such states if they are experimentally reliable, of high statistical significance ( $4.5\sigma$  or better), or observed in several different production processes.

(d) Partial-wave analyses of three-body final states ( $\pi N \rightarrow \pi\pi N$ ) are also available. While these analyses are based on the isobar model ( $\pi N \rightarrow \rho N, \pi\Delta$ , etc.) and are subject to theoretical objections of varying importance, they provide increasingly reliable information on inelastic decay modes of otherwise-established resonances.

### C. Statistical Procedures

We divide this discussion on obtaining averages and errors into two sections:

1. The unconstrained case, or "simple averaging;" and
2. The constrained case.

In what follows, the term "error" means one standard deviation ( $1\sigma$ ); that is, for central value  $\bar{x}$  and error  $\delta\bar{x}$ , the range  $\bar{x} \pm \delta\bar{x}$  constitutes a 68.3% confidence interval.

#### 1. Unconstrained averaging

We use a standard Gaussian procedure with a "scale factor" applied to the errors as our method of averaging the data. The Student's t-distribution, the basis of an earlier experiment of ours in data averaging, would give more conservative (and perhaps more realistic) errors at the two-standard-deviation ( $2\sigma$ ) and higher level, but we do not choose to quote such errors. It is worth bearing in mind,

however, that a  $2\sigma$  error might more realistically be somewhat larger than twice a  $1\sigma$  error, owing to the non-Gaussian character of some sets of real measurements. This is a persistent problem in data averaging arising from the existence of mildly discrepant measurements.

We begin by assuming that measurements of a given quantity are uncorrelated, and thus we calculate a weighted average and error

$$\bar{x} \pm \delta\bar{x} = \left[ \frac{\sum_i w_i x_i / \sum_i w_i}{\sum_i w_i} \right] \pm \left[ \frac{\sum_i w_i}{\sum_i w_i} \right]^{-1/2},$$

$$w_i = [1/(\delta x_i)^2], \quad (1)$$

where  $x_i$  and  $\delta x_i$  are the value and error, respectively, reported by the  $i$ th experiment, and the sums run over  $N$  experiments. We also calculate  $\chi^2$  and compare it with its expectation value; assuming that the measurements obey a Gaussian distribution, this is  $N - 1$ .

If  $\chi^2/(N - 1)$  is less than or equal to 1, and there are no known problems with the data, we accept the above results.

If  $\chi^2/(N - 1)$  is very large, or if there is prior knowledge of extremely large inconsistencies among experiments, we may choose not to average the data at all. Alternatively, we may quote the calculated average, but then give an educated guess as to the error; such a guess is generally a quite conservative estimate designed to take into account known problems with the data.

Finally, if  $\chi^2/(N - 1)$  is greater than 1, but not greatly so, we still average the data, but then also do the following:

(a) We plot an ideogram to display the pattern of the data. Sometimes only one or two data points lie apart from the main body; other times the data split into two or more roughly equal-sized groups. The reader may use this information in deciding upon an alternative average, but caution is urged, as "outlying" data points are sometimes the "correct" ones. An example of such an ideogram is given in Fig. 2 below. Each experiment appearing in the plot is

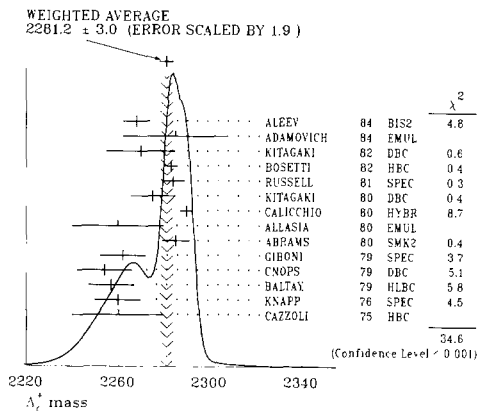


Fig. 2. Ideogram of measurements of the  $\Lambda_c^+$  mass. The "data point" at the top shows the position of the weighted average, while the width of the error bar (and the shaded pattern beneath it) shows the error in the average after scaling by the SCALE factor. Only those experiments indicated by + error flags were precise enough to be accepted in the calculation of the SCALE factor; the column on the far right gives the  $\chi^2$  contribution of each of these experiments. The less precise experiments were included in the calculation of the weighted average, but not SCALE; they have  $\perp$  error flags.

represented by a Gaussian with central value  $x_i$ , error  $\delta x_i$ , and area proportional to  $1/\delta x_i$ . The choice of area is somewhat arbitrary; it assumes that an experimenter will work to reduce the systematic errors until they are slightly smaller (but seldom much smaller) than the statistical errors. Thus, as a physicist collects more events, he or she will use them both to reduce the statistical errors and to study the biases. Our confidence that a significant systematic error has not been made in a given experiment, as compared with other contradictory experiments, then tends to go up as  $1/\delta x_i$ .

But why not assign a weight  $1/(\delta x_i)^2$ , as is done when computing a weighted average? We feel that this assignment is equivalent to assuming that large systematic errors are as infrequent as large statistical fluctuations, and that this assumption is unrealistic.

We emphasize the difference between least-squares averaging (where the weighting factor is the inverse square of the error) and the ideograms prepared for visual display. The former arithmetic is of course best if one has unbiased data whose errors are well understood. In particular, the error analysis assumes that the true error on each datum is sampled from a Gaussian whose width is correctly reported. Then we obtain a narrow Gaussian distribution centered at the weighted mean for the answer. The ideogram (often multi-peaked and certainly not Gaussian) is based on the opposite hypothesis that some of the input is systematically in error. The idea behind least-squares averaging is that experiments 1, 2, 3, etc., are *all* valid (so we should multiply their probabilities). Our *ideograms* are based on the assumption that 1 or 2 or 3, etc., is valid, "hedged" with  $1/\delta x_i$  betting odds; we then add their probabilities. Both approaches cannot simultaneously be right; we allow the reader to choose. However, we quote the least-squares result in the Summary Tables. This is the most precise value if the data satisfy the appropriate assumptions. A glance at the ideogram will show that the difference between the two approaches is usually not severe.

(b) The second way in which we try to take account of  $\chi^2/(N - 1)$  being greater than 1 is to scale up our quoted error  $\delta\bar{x}$  in Eq. (1) by a factor

$$\text{SCALE} = [\chi^2/(N - 1)]^{1/2}. \quad (2)$$

Our reasoning is as follows. Since we do not know which of the experiments are more than one standard deviation away from the correct value, we assume that all experimentalists underestimated their errors by the same scale factor (2). If we scale up all input errors by this factor,  $\chi^2$  becomes  $N - 1$ , and of course the output error scales up by the same factor.

If we are to combine experiments with widely varying errors, we modify this procedure slightly. This is because it is the more precise experiments that most influence not only the average value  $\bar{x}$ , but also the error  $\delta\bar{x}$ . Now, on the average, the low-precision experiments each contribute about unity to both the numerator and the denominator of SCALE, hence the  $\chi^2$  contribution of the sensitive experiments is diluted, i.e., reduced. Therefore, we evaluate SCALE by using *only* experiments for which the errors are not much greater than those of the more precise experiments, i.e., only those experiments with errors less than  $\delta_0$ , where the ceiling  $\delta_0$  is (arbitrarily) chosen to be

$$\delta_0 = 3N^{1/2}\delta\bar{x}.$$

Here  $\delta\bar{x}$  is the unscaled error of the mean of all the experiments. Note that if each experiment had the same error  $\delta x_i$ , then  $\delta\bar{x}$  would be  $\delta x_i/N^{1/2}$ , so each individual experiment would be well under the ceiling on SCALE.

This scaling approach has the property that if there are two values with comparable errors separated by much more than their stated errors (with or without a number of other experiments of lower accuracy), the error on the mean value  $\delta\bar{x}$  is increased so that it is approximately half the interval between the two discrepant values.

We wish to emphasize the fact that our scaling procedures for errors in no way affect central values. In addition, if one wishes to recover the unscaled error  $\delta\bar{x}$ , one need only divide the given error by the SCALE factor for that error.

## 2. Systematic errors and correlated measurements

Many experimental groups have now adopted the convention of presenting results with statistical and systematic errors explicitly indicated. Because of a lack of space in the Full Listings, we usually do not quote the two errors separately and at present we combine them.<sup>§</sup> In general we add the statistical and the systematic errors in quadrature. A comment is printed whenever this is done. When averages are calculated as described above in Eq. (1), the weight  $w_i$  is based upon the combined error.

It may happen that two measurements have correlated errors. For example, a group may improve the statistical or systematic errors by further data-taking or analysis. In this case we use only the improved result for averaging. The earlier result may still appear in the Listings (in parentheses), but in general we omit such obsolete entries.

A second case of correlated measurements is the occurrence of a common systematic error in experiments which are statistically independent. If two results  $A_1 \pm \sigma_1 \pm S$  and  $A_2 \pm \sigma_2 \pm S$  have completely correlated systematic errors  $S$ , one must first average  $A_1 \pm \sigma_1$  and  $A_2 \pm \sigma_2$  and then combine the resulting statistical error with  $S$ . One obtains, however, the same result by a second procedure, averaging  $A_1 \pm X\sigma_1$  and  $A_2 \pm X\sigma_2$  where

$$X = \left[ 1 + S^2/\sigma_1^2 + S^2/\sigma_2^2 \right]^{1/2}. \quad (3)$$

The second procedure has the advantage that the modified entries  $A_i \pm X\sigma_i$  may be averaged with further independent data as in Eq. (1). We therefore adopt this second procedure when appropriate.

## 3. Constrained fits

Except for trivial cases, all branching ratios and rate measurements are analyzed by making a simultaneous least-squares fit to all the data and extracting the partial decay fractions  $P_i$ , the width  $\Gamma$ , the partial widths  $\Gamma_i$ , and the associated error matrix.

Assume, for a simple example, that a state has only three partial decay fractions,  $P_1$ ,  $P_2$ , and  $P_3$  ( $\sum P_i = 1$ ), which have been measured in four different ratios,  $R_1, \dots, R_4$ ,

where, e.g.,  $R_1 = P_1/P_2$ ,  $R_2 = P_1/P_3$ , etc.\*\* Further assume that each ratio  $r$  has been measured by  $N_r$  experiments (we designate each experiment with a subscript  $x$ , e.g.,  $R_{1x}$ ). We then find the best values of  $P_1$ ,  $P_2$ , and  $P_3$  by minimizing  $\chi^2$ :

$$\chi^2 = \sum_{r=1}^4 \left[ \sum_{x=1}^{N_r} \left( \frac{R_{rx} - R_r(P_1, P_2, P_3)}{\delta R_{rx}} \right)^2 \right]. \quad (3)$$

In addition to the fitted values  $\bar{P}_i$ , we calculate an error matrix  $\langle \delta\bar{P}_i \delta\bar{P}_j \rangle$ . We tabulate the diagonal elements of  $\delta\bar{P}_i = \langle \delta\bar{P}_i \delta\bar{P}_i \rangle^{1/2}$  (except that some errors are scaled as discussed below). In the Full Listings we give the complete error matrix; we also calculate the fitted value of each ratio, for comparison with the input data, and list it below the relevant input, along with a simple unconstrained average of the same input.

Two further comments on the example above:

(1) There was no connection between measurements of the width and the branching ratios. But often we also have information on partial widths  $\Gamma_i$  as well as total width  $\Gamma$ . In this case we must introduce  $\Gamma$  as a parameter into the fit, along with the relations  $\Gamma_i = \Gamma P_i$ ,  $\sum \Gamma_i = \Gamma$ . When appropriate, we tabulate the  $\Gamma_i$  along with the  $P_i$ , and give error matrices in the Full Listings.

(2) We do *not* allow for correlations between input data. We *do* try to pick those ratios and widths which are as independent and as close to the original data as possible.

For *asymmetric* errors, we use a continuous function of  $\delta(P)^+$  and  $\delta(P)^-$  in the fitting. When no errors are reported, we merely list the data for inspection.

*Inconsistent constrained data.* According to Eq. (3), the double sum for  $\chi^2$  is first summed over experiments  $x = 1$  to  $N_r$ , leaving a single sum over ratios

$$\chi^2 = \sum_r \chi_r^2.$$

We test for SCALE factors after the fit. Knowing the fitted  $\chi_r^2$  and its expectation value  $\langle \chi_r^2 \rangle$ , we form SCALE factors (just as before), i.e.,

$$(\text{SCALE})_r^2 = \chi_r^2 / \langle \chi_r^2 \rangle,$$

and if any  $(\text{SCALE})_r$  is greater than 1, all  $N_r$  of the measurements of that particular ratio are equally penalized by having their errors increased by  $(\text{SCALE})_r$ . We then recycle the full fit, yielding new values  $\delta\bar{P}'_i$  for the errors in the partial decay modes, as well as new central values  $\bar{P}'_i$ .

Because of the constraint ( $\sum P_i = 1$ ), some of the new SCALE factors may still be greater than 1. If this is so, the whole procedure (i.e., increasing errors by the new SCALE factors and recycling through the fit) is repeated until the process converges.

At the end, we have final estimated errors  $\delta\bar{P}'_i$  for the  $\bar{P}'_i$ . If SCALE factors have been used, they normally will have caused a shift in the central fitted values  $\bar{P}'_i$ , as well as having given larger errors  $\delta\bar{P}'_i$ . Often we find that the shift  $|\bar{P}'_i - \bar{P}_i|$  due to the SCALE factors is the same size as (or

<sup>§</sup> We are considering a revision of the format of the Full Listings which would allow separation of these types of error, and also allow presentation of asymmetric errors.

\*\* We can handle any ratio  $R$  of the form  $\sum \alpha_i P_i / \sum \beta_i P_i$ , where  $\alpha_i$  and  $\beta_i$  are constants, usually 1 or 0. The forms  $R = P_i \cdot P_j$  and  $R = (P_i \cdot P_j)^{1/2}$  are also allowed.



greater than) the  $\delta\bar{P}'_i$ . We have decided to incorporate this shift into our errors as a reflection of the uncertainty due to the introduction of the SCALE factor; we tabulate an error

$$(\delta\bar{P}_i)_{\text{tab}} = [(\delta\bar{P}'_i)^2 + (\bar{P}_i - \bar{P}'_i)^2]^{1/2},$$

where  $\bar{P}_i$  is the fitted value of the  $i$ th partial decay mode before scaling,  $\bar{P}'_i$  is its value after all scaling, and  $\delta\bar{P}'_i$  is the error in  $\bar{P}'_i$ . The SCALE factors we finally list in such cases are defined by

$$(\text{SCALE})_i = (\delta\bar{P}_i)_{\text{tab}}/\delta\bar{P}'_i.$$

However, in line with our policy of not letting SCALE affect the central values, we quote the values of  $\bar{P}_i$  obtained from the original (unscaled) fit [which are always less than or equal to one standard deviation from  $\bar{P}'_i$ , by construction of  $(\delta\bar{P}_i)_{\text{tab}}$ ].

#### D. Discussion

The entire question of averaging data containing discrepant values is nicely discussed by Taylor (1982). He considers a number of algorithms which attempt to incorporate data which are not completely consistent into a meaningful average. Problems occur because it is very difficult to develop a procedure which handles simultaneously in a reasonable way two basic types of situations: (a) data which seem to lie apart from the main body of the data are incorrect (contain unreported errors); and (b) the opposite (the main body of the data is systematically wrong). Unfortunately, as Taylor shows, case (b) is not infrequent. His conclusion is that the choice of procedure is less significant than the initial choice of data to include or exclude.

We place a great emphasis on the choice of data to include or exclude. Unfortunately, the volume of data precludes spending as much time on the problem as we would like. We address this problem by soliciting the help of as many outside experts (consultants) as possible. In the final analysis, however, it is often impossible to determine which (if either) of two discrepant measurements is correct. Our SCALE factor technique is an attempt to address this ignorance by increasing the error above that suggested by least-squares analysis. In effect, we are saying that present experiments do not allow a precise determination of this constant because of unresolvable discrepancies, and one must await further measurements. The reader is warned of this situation by the size of the SCALE factor; he or she is then able to go back to the literature (via the Full Listings) and redo the average as desired.

Our situation with regard to discrepant data is easier to handle than most of the cases Taylor considers, such as estimates of the fundamental constants like  $\hbar$ , etc. Most of the errors in his case are dominated by systematic effects. In particle properties data, statistical effects are often at least as large as systematic effects, and statistical errors are usually easier to estimate. A notable exception occurs in partial-wave analyses, where different techniques applied to the same data yield different results. In this case, as stated earlier, we often do not attempt an average, but just quote a range of values.

A brief history of Particle Data Group averages is given in Rosenfeld (1975). Updated versions of some of Rosenfeld's figures are shown in Fig. 3. The least-squares error is shown by the thick portion of the error bars; the full

error bar exhibits the SCALE factor extension.

Some cases of rather wild fluctuation are shown; this usually represents the introduction of significant new data or the discarding of some older data. Older data are sometimes discarded in favor of more modern data if it is felt that the newer data had fewer systematic errors, had more checks on their systematic errors, made some corrections unknown at the time of the older experiments, or some such reason. Near the time at which a large jump takes place, the SCALE factor sometimes becomes large, reflecting the uncertainty introduced by the new existence of partly inconsistent data.

By and large, a full scan of our history plots shows a rather dull progression toward greater precision at a central value completely consistent with the first data point shown. These plots are available on request from the Berkeley Particle Data Group.

We conclude that the reliability of the combination of experimental data and Particle Data Group averaging procedures is usually good, but it is important to realize that fluctuations outside of the quoted errors can and do occur, perhaps with more frequency than expected for truly Gaussian errors.

#### ACKNOWLEDGMENTS

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# SUMMARY TABLES OF PARTICLE PROPERTIES

April 1986

## Particle Data Group

M. Aguilar-Benitez, R.M. Barnett, R.L. Crawford, R.A. Eichler, R. Frosch, G.P. Gopal, K.G. Hayes, J.J. Hernandez, I. Hinchliffe, G. Höhler, G.R. Lynch, D.M. Manley, L. Montanet, F.C. Porter, J. Primack, A. Rittenberg, M. Roos, L.D. Roper, R.H. Schindler, K.R. Schubert, T. Shimada, R.E. Shrock, N.A. Törnqvist, T.G. Trippe, W.P. Trower, C.G. Wohl, G.P. Yost, and B. Armstrong and G.S. Wagman (Technical Associates)

(Closing date for data: Dec. 1, 1985)

### Stable Particle Summary Table

(stable under strong decay)

For additional parameters, see Addendum to this table.

Quantities in italics are new or have changed by more than one (old) standard deviation since April 1984.

Particle	$J^G(J^{PC})^a$	Mass <sup>b</sup> (MeV)	Mean life <sup>b</sup>		Partial decay modes		<i>p</i> (MeV/c) <sup>c</sup>
			$\tau$ (sec)	$c\tau$ (cm)	Mode	Fraction <sup>b</sup>	
<b>GAUGE BOSONS</b>							
$\gamma$	0,1(1 <sup>-</sup> )	(< 3×10 <sup>-33</sup> )	-----		stable		
$W$	$J=1$	81.8 ±1.5 GeV	$\Gamma < 6.5$ GeV		$e\nu$ $\mu\nu$	seen seen	40.9 GeV/c 40.9
$Z$		92.6 ±1.7 GeV	$\Gamma < 4.6$ GeV		$e^+e^-$ $\mu^+\mu^-$	seen seen	46.3 GeV/c 46.3
<b>LEPTONS</b>							
$\nu_e$	$J=\frac{1}{2}$	< 46 eV <sup>d</sup>	stable		stable		
			[ $\tau > 3 \times 10^8 m_{\nu_e}$ sec ( $m_{\nu_e}$ in MeV)]				
$e$	$J=\frac{1}{2}$	0.5110034 ±0.0000014 MeV	stable		stable		
			(> 2×10 <sup>22</sup> years)				
$\nu_\mu$	$J=\frac{1}{2}$	< 0.25	stable		stable		
			[ $\tau > 1.1 \times 10^5 m_{\nu_\mu}$ sec ( $m_{\nu_\mu}$ in MeV)]				
$\mu$	$J=\frac{1}{2}$	105.65916 ±0.00030 S=1.1*	2.19703×10 <sup>-6</sup> ±0.00004		$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$ $e^- \bar{\nu}_e \nu_\mu \gamma$ $e^- \nu_e \bar{\nu}_\mu + e^-$ $e^- \gamma$ $e^- e^+ e^-$ $e^- \gamma \gamma$	(100)% (1.4 ± 0.4)% (<5)% (3.4 ± 0.4)×10 <sup>-5</sup> (<1.7)×10 <sup>-10</sup> (<2.4)×10 <sup>-12</sup> (<8.4)×10 <sup>-9</sup>	53 53 53 LF 53 53 53
			$c\tau = 6.5865 \times 10^4$		(or $\mu^+ \rightarrow$ chg. conj.)		

## Stable Particle Summary Table (cont'd)

Particle	$I^G(J^{PC})^a$	Mass <sup>b</sup> (MeV)	Mean life <sup>b</sup>		Partial decay modes		
			$\tau$ (sec)	$c\tau$ (cm)	Mode	Fraction <sup>b</sup>	$p$ (MeV/c) <sup>a</sup>
$\nu_\tau$	$J=\frac{1}{2}$	< 70					
$\tau$	$J=\frac{1}{2}$	1784.2 $\pm 3.2$	$(3.3 \pm 0.4) \times 10^{-13}$	$c\tau=0.010$	$\tau^- \rightarrow \nu_\tau$ (or $\tau^+ \rightarrow \text{chg. conj.}$ )		
					particle <sup>-</sup> neutrals	( 86.5 $\pm$ 0.3 ) %	
					$\mu^- \nu \nu$	( 17.6 $\pm$ 0.6 ) %	889
					$e^- \nu \nu$	( 17.4 $\pm$ 0.5 ) %	892
					hadron <sup>-</sup> $\geq 0\pi^0 \nu$	( 51.6 $\pm$ 0.7 ) %	
					hadron <sup>-</sup> $\nu$	( 10.8 $\pm$ 1.1 ) %	
					$\pi^- \nu$	( 10.1 $\pm$ 1.1 ) %	887
					$K^- \nu$	( 0.67 $\pm$ 0.17 ) %	824
					hadron <sup>-</sup> $\geq 1\pi^0 \nu$	( 40.8 $\pm$ 1.3 ) %	
					$\rho^- \nu$	( 21.8 $\pm$ 2.0 ) %	726
					$\pi^- \pi^0$ (non-res.) $\nu$	( 0.3 $\pm$ 0.3 ) %	881
					$\pi^- \pi^0 \pi^0 \nu$	( 6.0 $\pm$ 3.5 ) %	866
					$\pi^- \pi^0 \pi^0 \pi^0 \nu$	( 3.0 $\pm$ 2.7 ) %	840
					$K^- \geq 1\pi^0 \nu$	( 1.0 $\pm$ 0.3 ) %	
					$\pi^- \pi^- \pi^+ \geq 0\pi^0 \nu$	( 13.4 $\pm$ 0.3 ) %	
					$\pi^- \pi^- \pi^+ \geq 1\pi^0 \nu$	( 5.3 $\pm$ 0.8 ) %	
					$\pi^- \pi^- \pi^+ \nu$	( 8.1 $\pm$ 0.7 ) %	865
					$\pi^- \rho^0 \nu$	( 5.4 $\pm$ 1.7 ) %	718
					$\pi^- \pi^- \pi^+ \nu$ (non-res.)	( < 1.4 ) %	865
					$\pi^- \pi^- \pi^+ K^0 \geq 0\gamma \nu$	( < 0.27 ) %	
					$K^- 2\text{charged} \geq 0\pi^0 \nu$	( < 0.6 ) %	
					$K^- \pi^+ \pi^- \geq 0\pi^0 \nu$	( 0.2 $\pm$ 0.2 ) %	
					$K^+ K^- \pi^- \nu$	( 0.2 $\pm$ 0.2 ) %	690
					$3\pi^- 2\pi^+ \geq 0\pi^0 \nu$	( 0.15 $\pm$ 0.04 ) %	
					$3\pi^- 2\pi^+ \nu$	( 0.07 $\pm$ 0.03 ) %	
					$3\pi^- 2\pi^+ \pi^0 \nu$	( 0.07 $\pm$ 0.03 ) %	
					$\dagger [K^*(892)^- \nu$	( 1.7 $\pm$ 0.7 ) %	669
					$K_2^*(1430)^- \nu$	( < 0.9 ) %	323
Searches for massive neutrinos and lepton mixing $\nu$ bounds from astrophysics and cosmology Heavy lepton searches					} See Addendum and the Stable Particle Full Listings		
<b>LIGHT MESONS</b> <sup>a</sup>					$[\pi^+ = u\bar{d}, \pi^0 = (u\bar{u} - d\bar{d})/\sqrt{2}, \pi^- = \bar{u}d, \eta = c_1(u\bar{u} + d\bar{d}) + c_2(s\bar{s})]$		
$\pi^\pm$	$1^-(0^-)$	139.5685 $\pm 0.0010$ S=1.5*	$2.6030 \times 10^{-8}$ $\pm 0.0023$	$c\tau=780.4$	$\pi^\pm \rightarrow \nu$ (or $\pi^\pm \rightarrow \text{chg. conj.}$ )		
					$\mu^\pm \nu$	100%	30
					$e^\pm \nu$	( 1.228 $\pm$ 0.022 ) $\times 10^{-4}$	S=1.8* 70
					$\dagger [\mu^\pm \nu \gamma$	( 1.24 $\pm$ 0.25 ) $\times 10^{-4}$	30
					$e^\pm \nu \gamma$	( 5.6 $\pm$ 0.7 ) $\times 10^{-8}$	70
					$e^\pm \nu \pi^0$	( 1.025 $\pm$ 0.034 ) $\times 10^{-8}$	5
					$e^\pm \nu e^+ e^-$	( < 5 ) $\times 10^{-9}$	70
					$\mu^\pm \nu e$	( < 1.5 ) $\times 10^{-3}$	30
					$\mu^\pm \nu e$	( < 8 ) $\times 10^{-3}$	LF 30
$\pi^0$	$1^-(0^{++})$	134.9642 $\pm 0.0038$	$0.87 \times 10^{-16}$ $\pm 0.04$	S=1.8* $c\tau=2.6 \times 10^{-6}$	$\gamma\gamma$	( 98.799 $\pm$ 0.030 ) %	67
					$\gamma e^+ e^-$	( 1.198 ) %	67
					$e^+ e^- e^+ e^-$	( 3.24 ) $\times 10^{-5}$	67
					$\gamma\gamma\gamma$	( < 3.8 ) $\times 10^{-7}$	C 67
					$\gamma\gamma\gamma\gamma$	( < 4 ) $\times 10^{-6}$	67
					$e^+ e^-$	( 1.8 $\pm$ 0.7 ) $\times 10^{-7}$	67
					$\nu\nu$	( < 2.4 ) $\times 10^{-5}$	67
					$\mu^+ e^- + \mu^- e^+$	( < 7 ) $\times 10^{-8}$	LF 26

## Stable Particle Summary Table (cont'd)

Particle	$I^G(J^{PC})^a$	Mass <sup>b</sup> (MeV)	Mean life <sup>b</sup> $\tau$ (sec) $c\tau$ (cm)	Partial decay modes									
				Mode	Fraction <sup>b</sup>	$p$ (MeV/c) <sup>c</sup>							
$\eta$	$0^+(0^{-+})$	548.8 $\pm 0.6$ S=1.4*	$\Gamma=(1.05 \pm 0.15)$ keV S=1.7* Neutral decays (70.9 $\pm$ 0.6)%	$\left\{ \begin{array}{l} \gamma\gamma \\ 3\pi^0 \\ \pi^0\gamma\gamma \\ \pi^+\pi^-\pi^0 \\ \pi^+\pi^-\gamma \\ \pi^+e^-\gamma \\ \mu^+\mu^-\gamma \\ e^+e^-\gamma \\ \mu^+\mu^- \\ \pi^+\pi^-e^+e^- \\ \pi^+\pi^-\gamma\gamma \\ \pi^+\pi^-\pi^0\gamma \\ \pi^+\pi^- \\ \pi^0e^+e^- \\ \pi^0\mu^+\mu^- \\ \pi^0\mu^+\mu^-\gamma \end{array} \right.$	$\left\{ \begin{array}{l} (38.9 \pm 0.4)\% \\ (31.90 \pm 0.35)\% \\ (0.078 \pm 0.012)\% \\ (<5) \times 10^{-4} \\ (23.7 \pm 0.5)\% \\ (4.91 \pm 0.13)\% \\ (0.50 \pm 0.12)\% \\ (3.1 \pm 0.4) \times 10^{-4} \\ (<3) \times 10^{-4} \\ (6.5 \pm 2.1) \times 10^{-6} \\ (0.13 \pm 0.13)\% \\ (<0.21) \\ (<6) \times 10^{-4} \\ (<0.15) \\ (<5) \times 10^{-5} \\ (<5) \times 10^{-6} \\ (<3) \times 10^{-6} \end{array} \right.$	$\left\{ \begin{array}{l} 274 \\ 180 \\ 258 \\ C \\ 274 \\ 175 \\ 236 \\ 274 \\ 253 \\ 274 \\ 253 \\ 236 \\ 236 \\ 175 \\ P, CP \\ 236 \\ 258 \\ C \\ 211 \\ 211 \end{array} \right.$							
							<b>STRANGE MESONS<sup>a</sup></b>				$[K^+ = u\bar{s}, K^0 = d\bar{s}, \bar{K}^0 = \bar{d}s, K^- = \bar{u}s]$		
							$K^\pm$	$\frac{1}{2}(0^-)$	493.667 $\pm 0.014$  $m_{K^0} - m_{K^\pm} = 4.05$ $\pm 0.07$	1.2371 $\times 10^{-8}$ $\pm 0.0026$ S=1.9*  $c\tau=370.9$	$\left\{ \begin{array}{l} \mu^+\nu \\ \pi^+\pi^0 \\ \pi^+\pi^+\pi^- \\ \pi^+\pi^0\pi^0 \\ \pi^0\mu^+\nu \\ \pi^0e^+\nu \\ \dagger[\mu^+\nu\gamma \\ \mu^+\nu\gamma (SD+)^g \\ \pi^+\pi^0\gamma \\ \pi^+\pi^+\pi^-\gamma \\ \pi^0\mu^+\nu\gamma \\ \pi^0e^+\nu\gamma \\ \pi^0\pi^0e^+\nu \\ \pi^+\pi^-e^+\nu \\ \pi^+\pi^+e^+\nu \\ \pi^+\pi^-\mu^+\nu \\ \pi^+\pi^+\mu^+\nu \\ e^+\nu \\ e^+\nu\gamma (SD+)^g \\ e^+\nu\gamma (SD-)^g \\ \pi^+e^+e^- \\ \pi^-e^+e^+ \\ \pi^+\mu^+\mu^- \\ \pi^+\gamma\gamma \\ \pi^+\gamma\gamma\gamma \\ \pi^+\nu\nu \\ \pi^\mp e^+\mu^\pm \\ \pi^+e^-\mu^+ \\ e^+\nu\nu \\ \mu^+\nu\nu \\ \mu^+\nu e^+e^- \\ \mu^-\nu e^+e^+ \\ e^+\nu e^+e^- \\ \mu^+\nu_e \\ \mu^+\nu_e \\ \pi^0e^+\nu_e \end{array} \right.$	$\left\{ \begin{array}{l} (63.51 \pm 0.16)\% \\ (21.17 \pm 0.15)\% \\ (5.59 \pm 0.03)\% \text{ S}=1.1* \\ (1.73 \pm 0.05)\% \text{ S}=1.4* \\ (3.18 \pm 0.10)\% \text{ S}=1.9* \\ (4.82 \pm 0.05)\% \text{ S}=1.1* \\ e^i (5.4 \pm 0.3) \times 10^{-3} \\ (<3.0) \times 10^{-5} \\ h,e^i (2.75 \pm 0.16) \times 10^{-4} \\ e^i (1.0 \pm 0.4) \times 10^{-4} \\ e^i (<6) \times 10^{-5} \\ e^i (3.7 \pm 1.4) \times 10^{-4} \\ (1.8 \pm_{0.6}^{2.4}) \times 10^{-5} \\ (3.90 \pm 0.15) \times 10^{-5} \\ (<1.2) \times 10^{-8} \text{ SQ} \\ (1.4 \pm 0.9) \times 10^{-5} \\ (<3.0) \times 10^{-6} \text{ SQ} \\ (1.54 \pm 0.07) \times 10^{-5} \\ (1.52 \pm 0.23) \times 10^{-5} \\ (<1.6) \times 10^{-4} \\ (2.7 \pm 0.5) \times 10^{-7} \text{ FC} \\ (<1) \times 10^{-8} \text{ L} \\ (<2.4) \times 10^{-6} \text{ FC} \\ e^i (<8) \times 10^{-6} \\ e^i (<1.0) \times 10^{-4} \\ (<1.4) \times 10^{-7} \text{ FC} \\ (<7) \times 10^{-9} \text{ LF, or L} \\ (<5) \times 10^{-9} \text{ LF} \\ (<6) \times 10^{-5} \\ (<6) \times 10^{-6} \\ (11 \pm 3) \times 10^{-7} \\ (<2.0) \times 10^{-8} \text{ LF} \\ (2 \pm_{-1}^2) \times 10^{-7} \\ (<4) \times 10^{-3} \text{ LF} \\ (<3.3) \times 10^{-3} \text{ L} \\ (<3) \times 10^{-3} \text{ L} \end{array} \right.$	$\left\{ \begin{array}{l} 236 \\ 205 \\ 125 \\ 133 \\ 215 \\ 228 \\ 236 \\ 236 \\ 205 \\ 125 \\ 215 \\ 228 \\ 207 \\ 203 \\ 203 \\ 151 \\ 151 \\ 247 \\ 247 \\ 227 \\ 227 \\ 172 \\ 227 \\ 227 \\ 227 \\ 214 \\ 214 \\ 247 \\ 236 \\ 236 \\ 247 \\ 236 \\ 236 \\ 247 \end{array} \right.$

## Stable Particle Summary Table (cont'd)

Particle	$I(J^P)^a$	Mass <sup>b</sup> (MeV)	Mean life <sup>b</sup> $\tau$ (sec) $c\tau$ (cm)	Partial decay modes						
				Mode	Fraction <sup>b</sup>	$p$ (MeV/c) <sup>c</sup>				
$\frac{K^0}{\bar{K}^0}$	$\frac{1}{2}(0^-)$	497.72 $\pm 0.07$ S=1.4*		50% $K_S$ , 50% $K_L$						
$K_S^0$	$\frac{1}{2}(0^-)$		0.8923 $\times 10^{-10}$ $\pm 0.0022$  $c\tau=2.675$	$\pi^+\pi^-$	(68.61 $\pm$ 0.24)%	206				
				$\pi^0\pi^0$	(31.39 $\pm$ 0.20)%	S=1.1* 209				
				$\dagger[\pi^+\pi^-\gamma]$	$e^i(1.85 \pm 0.10)\times 10^{-3}$	206				
				$\mu^+\mu^-$	(<3.2) $\times 10^{-7}$	225				
				$e^+e^-$	(<3.4) $\times 10^{-4}$	FC 249				
				$\gamma\gamma$	(<4) $\times 10^{-4}$	249				
				$\pi^+\pi^-\pi^0$	(<4.9) $\times 10^{-5}$	133				
				$\pi^0\pi^0\pi^0$	(<3.7) $\times 10^{-5}$	139				
$K_L^0$	$\frac{1}{2}(0^-)$	$m_{K_L} - m_{K_S} = 0.5349 \times 10^{10} \text{ h sec}^{-1}$ $\pm 0.0022$  $= 3.521 \times 10^{-12} \text{ MeV}$ $\pm 0.014$	5.183 $\times 10^{-8}$ $\pm 0.040$  $c\tau=1554$	$\pi^0\pi^0\pi^0$	(21.5 $\pm$ 1.0)%	S=1.7* 139				
				$\pi^+\pi^-\pi^0$	(12.40 $\pm$ 0.20)%	S=1.3* 133				
				$\pi^\pm\mu^\mp\nu$	(27.1 $\pm$ 0.4)%	S=1.4* 216				
				$\pi^\pm e^\mp\nu$	(38.7 $\pm$ 0.6)%	S=1.5* 229				
				$\pi^+\pi^-$	$i^i(0.203 \pm 0.005)\%$	S=1.1* 206				
				$\pi^0\pi^0$	$i^i(0.094 \pm 0.019)\%$	S=1.5* 209				
				$\dagger[\pi e\nu\gamma]$	$e^i(1.3 \pm 0.8)\%$	229				
				$\pi^+\pi^-\gamma$	$e^i(4.41 \pm 0.32)\times 10^{-5}$	206				
				$\pi^0\gamma\gamma$	(<2.4) $\times 10^{-4}$	231				
				$\gamma\gamma$	(4.9 $\pm$ 0.4) $\times 10^{-4}$	249				
				$e\mu$	$j^i(<6)\times 10^{-6}$	LF 238				
				$\mu^+\mu^-$	(9.1 $\pm$ 1.9) $\times 10^{-9}$	FC 225				
				$\mu^+\mu^-\gamma$	(2.8 $\pm$ 2.8) $\times 10^{-7}$	FC 225				
				$\pi^0\mu^+\mu^-$	(<1.2) $\times 10^{-6}$	FC 177				
				$e^+e^-$	$j^i(<2.0)\times 10^{-7}$	FC 249				
				$e^+e^-\gamma$	(1.7 $\pm$ 0.9) $\times 10^{-5}$	FC 249				
				$\pi^0e^+e^-$	(<2.3) $\times 10^{-6}$	FC 231				
				$\pi^+\pi^-e^+e^-$	(<2.5) $\times 10^{-6}$	FC 206				
				$\pi^0\pi^\pm e^\mp\nu$	(6.2 $\pm$ 2.0) $\times 10^{-5}$	207				
				$(\pi\mu \text{ atom})\nu$	(1.05 $\pm$ 0.11) $\times 10^{-7}$	216				
				$\mu^+\mu^-e^+e^-$	(<5) $\times 10^{-6}$	FC 225				
				$e^+e^-e^+e^-$	(<2.6) $\times 10^{-6}$	FC 249				
				<b>CHARMED MESONS<sup>a</sup></b>				$[D^+ = c\bar{d}, D^0 = c\bar{u}, \bar{D}^0 = \bar{c}u, D^- = \bar{c}d; D_s^+ = c\bar{s}, D_s^- = \bar{c}s]$		
								$^k D^+ \rightarrow \text{---} \text{ (or } D^- \rightarrow \text{chg. conj.)}$		
				$D^\pm$	$\frac{1}{2}(0^-)$	1869.3 $\pm 0.6$  $m_{D^\pm} - m_{D^0} = 4.7$ $\pm 0.3$	(9.2 $^{+1.3}_{-1.0}$ ) $\times 10^{-13}$ $c\tau=0.028$	$e^+$ anything	(18.2 $\pm$ 1.7)%	
								$K^-$ anything	(16 $\pm$ 4)%	
								$K^+$ anything	(6.0 $\pm$ 3.3)%	
								$\bar{K}^0$ any + $K^0$ any	(48 $\pm$ 15)%	
$\eta$ anything	$e^i(<13)\%$									
$e^+\nu$	(<2.5)%	935								
$\mu^+\nu$	(<2)%	932								
$\pi^+\pi^0$	(<0.5)%	925								
$\pi^+\pi^+\pi^-$	(0.5 $\pm$ 0.2)%	908								
$\dagger[K^-\pi^+\pi^+]$	(11.4 $\pm$ 1.1)%	845								
$K^-\pi^+\pi^+\pi^0$	(6.4 $\pm$ 1.5)%	816								
$K^-\pi^+\pi^+\pi^-\pi^-$	(<6)%	772								
$\bar{K}^0\pi^+$	(4.1 $\pm$ 0.6)%	862								
$\bar{K}^0\pi^+\pi^0$	(13.4 $\pm$ 3.3)%	845								
$\bar{K}^0\pi^+\pi^+\pi^-$	(15.2 $\pm$ 5.8)%	814								
$\bar{K}^0K^+$	(1.2 $\pm$ 0.4)%	792								
$K^+K^-\pi^+$	(<1.6)%	744								
$K^+\pi^+\pi^-$	(<0.6)%	845								
$\dagger[\bar{K}^*(892)^0\pi^+]$	(<6.6)%	713								
$K^+K^-\pi^+$ (non-res.)	(0.7 $\pm$ 0.3)%	744								
$\phi\pi^+$	(1.0 $\pm$ 0.3)%	647								
$\bar{K}^*(892)^0K^+$	(0.5 $\pm$ 0.3)%	613								

## Stable Particle Summary Table (cont'd)

Particle	$I(J^P)^a$	Mass <sup>b</sup> (MeV)	Mean life <sup>b</sup> $\tau$ (sec) $c\tau$ (cm)	Partial decay modes		$\rho$ (MeV/c) <sup>c</sup>	
				Mode	Fraction <sup>b</sup>		
$D^0$ $\bar{D}^0$	$\frac{1}{2}(0^-)$	1864.6 $\pm 0.6$	$(4.3^{+0.5}_{-0.4}) \times 10^{-13}$ $c\tau=0.013$	<sup>k</sup> $D^0 \rightarrow \bar{D}^0$ (or $\bar{D}^0 \rightarrow \text{chg. conj.}$ )			
				$e^+$ anything	( 7.0 $\pm$ 1.1 ) %	S=1.3*	
				$K^-$ anything	( 44 $\pm$ 10 ) %		
				$K^+$ anything	( 8 $\pm$ 3 ) %		
				$\bar{K}^0$ any+ $K^0$ any	( 33 $\pm$ 10 ) %		
				$\eta$ anything	<sup>f</sup> ( < 13 ) %		
				$\pi^+\pi^-$	( 0.18 $\pm$ 0.05 ) %		922
				$\pi^+\pi^+\pi^-\pi^-$	( 1.5 $\pm$ 0.6 ) %		880
				$\mu^+\mu^-$	( < 3.4 ) $\times 10^{-4}$		FC 926
				$\pi^+\pi^-\pi^0$	( 1.1 $\pm$ 0.4 ) %		907
<sup>†</sup> $[K^*(892)^-\pi^+$	( 7.1 $\pm$ 2.5 ) %	711					
$\bar{K}^{*0}\pi^0$	( 2.3 $\pm$ 2.1 ) %	711					
$K^-\rho^+$	( 9.9 $\pm$ 3.5 ) %	679					
$\bar{K}^0\rho^0$	( 0.1 $\pm$ $^{0.9}_{0.1}$ ) %	677					
$K^-\pi^+\pi^0$ (non-res.)	( < 4.3 ) %	844					
$\bar{K}^0\pi^+\pi^-$ (non-res.)	( 2.0 $\pm$ 1.7 ) %	842					
$\bar{K}^*(892)^0\rho^0$	( 1.6 $\pm$ 1.6 ) %	423					
$K^-\pi^+\rho^0$	( 9.3 $\pm$ $^{1.5}_{2.6}$ ) %	613					
$\bar{K}^*(892)^0\pi^+\pi^-$	( < 2.9 ) %	685					
$K^-\pi^+\rho^+$ (1320) <sup>+</sup>	( < 1.9 ) %	198					
$\bar{K}^0\phi$	( 1.5 $\pm$ 0.4 ) %	520					
$\bar{K}^0K^+K^-$ (non-res.)	( 1.1 $\pm$ 0.5 ) % ]	544					
$D_s^\pm$	$0(0^-)^n$	1970.5 <sup>n</sup> $\pm 2.5$	$(2.8^{+1.6}_{-0.7}) \times 10^{-13}$	<sup>k</sup> $D_s^+ \rightarrow \bar{D}_s^-$ (or $D_s^- \rightarrow \text{chg. conj.}$ )			
was $F^\pm$	S=1.2*		$c\tau=0.008$	$\phi\pi^+$	( seen )	713	
				$\phi\pi^+\pi^+\pi^-$	( seen )	641	
				$\eta\pi^+$	( possibly seen )	903	
				$\eta\pi^+\pi^+\pi^-$	( possibly seen )	857	
				$\eta'\pi^+\pi^+\pi^-$	( possibly seen )	679	
				$\phi\rho^+$	( possibly seen )	411	
BOTTOM MESONS <sup>a</sup>				$[B^+ = u\bar{b}, B^0 = d\bar{b}, \bar{B}^0 = \bar{d}b, B^- = \bar{u}b]$			
$B^\pm$	$\frac{1}{2}(0^-)^p$	5271.2 $\pm 3.0$	$m_{B^0} - m_{B^\pm} = 4.0$ $\pm 3.4$	<sup>k</sup> $B^+ \rightarrow \bar{B}^0$ (or $B^- \rightarrow \text{chg. conj.}$ )			
				$\bar{D}^0\pi^+$	( 1.1 $\pm$ 0.6 ) %	2303	
				$D^*(2010)^-\pi^+\pi^+$	( 2.7 $\pm$ 1.7 ) %	2243	
				$J/\psi(3097)K^+$	( < 2.6 ) $\times 10^{-3}$	1678	
				$\rho^0\pi^+$	( < 6 ) $\times 10^{-4}$	2578	
$B^0$ $\bar{B}^0$	$\frac{1}{2}(0^-)^p$	5275.2 $\pm 2.8$		<sup>k</sup> $B^0 \rightarrow \bar{B}^0$ (or $\bar{B}^0 \rightarrow \text{chg. conj.}$ )			
				$\bar{D}^0\pi^+\pi^-$	( 7 $\pm$ 5 ) %	2299	
				$D^*(2010)^-\pi^+$	( 1.7 $\pm$ 0.7 ) %	2253	
				$D^*(2010)^-\rho^+$	( 8 $\pm$ $^{7}_{4}$ ) %	2180	
				$J/\psi(3097)K^+\pi^-$	( < 6.3 ) $\times 10^{-3}$	1649	
				$\pi^+\pi^-$	( < 5 ) $\times 10^{-4}$	FC 2634	
				$e^+e^-$	( < 3 ) $\times 10^{-4}$	FC 2638	
				$\mu^+\mu^-$	( < 2 ) $\times 10^{-4}$	FC 2635	
	$e^+\mu^- + \text{c.c.}$	( < 3 ) $\times 10^{-4}$	LF 2637				
$B^\pm, B^0, \bar{B}^0$	(not separated) <sup>q</sup>		$(14.2 \pm 2.7) \times 10^{-13}$	$e^\pm\nu$ hadrons		( 12.3 $\pm$ 0.8 ) %	
			$c\tau=0.043$	$\mu^\pm\nu$ hadrons		( 11.0 $\pm$ 0.9 ) %	
	$\frac{\Gamma(B \rightarrow \bar{B} \rightarrow \ell^- \text{ any})}{\Gamma(B \rightarrow \ell^\pm \text{ any})} < 0.12$			$D^0$ anything		( 80 $\pm$ 28 ) %	
	$\frac{\Gamma(B \rightarrow e^\pm\nu \text{ noncharm-hadrons})}{\Gamma(B \rightarrow e^\pm\nu \text{ hadrons})} < 0.04$			$K$ anything		( seen )	
				$\rho$ anything		( > 3.6 ) %	
				$\Lambda$ anything		( > 2.2 ) %	
				$e^+e^-$ anything		( < 0.6 ) %	
				$\mu^+\mu^-$ anything		( < 0.6 ) %	
				$J/\psi(3097)$ anything		( 1.2 $\pm$ 0.3 ) %	
				$D^*(2010)^\pm$ anything		( 23 $\pm$ 9 ) %	

## Stable Particle Summary Table (cont'd)

Particle	$I(J^P)^a$	Mass <sup>b</sup> (MeV)	Mean life <sup>b</sup> $\tau$ (sec) $c\tau$ (cm)	Partial decay modes		
				Mode	Fraction <sup>b</sup>	$p$ (MeV/c) <sup>c</sup>
<b>NUCLEONS<sup>a</sup></b> [ $p = uud, n = udd$ ]						
<b>p</b>	$\frac{1}{2}(\frac{1}{2}^+)$	938.2796 $\pm 0.0027$	stable ( $> 1.6 \times 10^{25}$ yr or $> 10^{31} - 3 \times 10^{32}$ yr) <sup>r</sup>	Partial mean lifetimes of protons (units $10^{30}$ yr)		Partial mean lifetimes of bound neutrons (units $10^{30}$ yr)
				$ q_p  -  q_e  < 10^{-21}  q_e ^s$	$e^+$ anything $> 0.6$	$\nu\gamma$ $> 11$
<b>n</b>	$\frac{1}{2}(\frac{1}{2}^+)$	939.5731 $\pm 0.0027$	898 $\pm$ 16 S=2.4* $c\tau=2.7 \times 10^{13}$	$\mu^+$ anything $> 12$	$e^+\pi^-$ $> 25$	
				$e^+\gamma$ $> 360$	$\mu^+\pi^-$ $> 38$	
				$\mu^+\gamma$ $> 280$	$e^+K^-$ $> 1.3$	
				$e^+\pi^0$ $> 250$	$\mu^+K^-$ $> 0.4$	
				$\mu^+\pi^0$ $> 100$	$\nu\pi^0$ $> 7$	
				$e^+K^0$ $> 77$	$\nu K^0$ $> 10$	
				$\mu^+K^0$ $> 40$	$e^+\rho^-$ $> 12$	
				$\nu\pi^+$ $> 5.8$	$\mu^+\rho^-$ $> 9$	
				$\nu K^+$ $> 9.6$	$\nu K^*(892)^0$ $> 7$	
				$e^+\omega$ $> 37$	$\nu\nu\nu$ $> 0.0005$	
				$\mu^+\omega$ $> 23$	$e^+\pi^0$ any $> 0.6$	
				$e^+\rho^0$ $> 17$	$\nu\eta$ $> 18$	
				$\mu^+\rho^0$ $> 16$	$\nu\rho^0$ $> 2$	
				$e^+e^+e^-$ $> 510$	$\nu\omega$ $> 16$	
				$\mu^+\mu^+\mu^-$ $> 190$	$e^+e^-\nu$ $> 26$	
				$\nu K^*(892)^+$ $> 9.6$	$\mu^+\mu^-\nu$ $> 19$	
				$e^+\pi^0$ any $> 0.6$	$e^-\pi^+$ $> 25$	
				$e^+\eta$ $> 200$	$\mu^-\pi^+$ $> 27$	
				$\mu^+\eta$ $> 46$	$e^-\rho^+$ $> 12$	
				$\nu\rho^+$ $> 8.4$	$\mu^-\rho^+$ $> 9$	
	$p\bar{e}^-\bar{\nu}$ 100%	1.2				
	$p\nu\bar{\nu}$ ( $< 9$ ) $\times 10^{-24}$ Q	1.3				
<b>STRANGENESS -1 BARYONS<sup>a</sup></b> [ $\Lambda = uds, \Sigma^+ = uus, \Sigma^0 = uds, \Sigma^- = dds$ ]						
<b><math>\Lambda</math></b>	$0(\frac{1}{2}^+)$	1115.60 $\pm 0.05$ S=1.2*	2.632 $\times 10^{-10}$ $\pm 0.020$ S=1.6* $c\tau=7.89$	$p\pi^-$ (64.2 $\pm$ 0.5)%		100
				$n\pi^0$ (35.8 $\pm$ 0.5)%		104
				$p\bar{e}^-\bar{\nu}$ (8.35 $\pm$ 0.14) $\times 10^{-4}$		163
				$p\mu^-\bar{\nu}$ (1.57 $\pm$ 0.35) $\times 10^{-4}$		131
				$\dagger [p\pi^-\gamma]$ (8.5 $\pm$ 1.4) $\times 10^{-4}$		100
$m_{\Sigma^0} - m_{\Lambda} = 76.86$ $\pm 0.08$						
<b><math>\Sigma^+</math></b>	$1(\frac{1}{2}^+)$	1189.37 $\pm 0.06$ S=1.8*	0.800 $\times 10^{-10}$ $\pm 0.004$ $c\tau=2.40$	$p\pi^0$ (51.64 $\pm$ 0.30)%		189
				$n\pi^+$ (48.36 $\pm$ 0.30)%		185
				$p\gamma$ (1.22 $\pm$ 0.10) $\times 10^{-3}$	S=1.1*	225
				$\dagger [n\pi^+\gamma]$ (4.5 $\pm$ 0.5) $\times 10^{-4}$		185
				$\Lambda e^+\nu$ (2.0 $\pm$ 0.5) $\times 10^{-5}$		71
				$n\mu^+\nu$ ( $< 3.0$ ) $\times 10^{-5}$	SQ	202
				$ne^+\nu$ ( $< 5$ ) $\times 10^{-6}$	SQ	224
$pe^+e^-$ ( $< 7$ ) $\times 10^{-6}$	FC	225				
$m_{\Sigma^-} - m_{\Sigma^+} = 7.97$ $\pm 0.07$ S=1.3*	$\frac{\Gamma(\Sigma^+ \rightarrow \ell^+ n \nu)}{\Gamma(\Sigma^- \rightarrow \ell^- n \nu)} < 0.04$					
<b><math>\Sigma^0</math></b>	$1(\frac{1}{2}^+)^t$	1192.46 $\pm 0.08$	5.8 $\times 10^{-20}$ $\pm 1.3$ $c\tau=1.7 \times 10^{-9}$	$\Lambda\gamma$ 100%		74
				$\Lambda e^+e^-$ (5.45 $\pm$ 0.3) $\times 10^{-3}$		74
				$\Lambda\gamma\gamma$ ( $< 3$ ) %		74
<b><math>\Sigma^-</math></b>	$1(\frac{1}{2}^+)$	1197.34 $\pm 0.05$	1.482 $\times 10^{-10}$ $\pm 0.011$ S=1.3* $c\tau=4.44$	$n\pi^-$ 100%		193
				$ne^-\bar{\nu}$ (1.022 $\pm$ 0.034) $\times 10^{-3}$		230
				$n\mu^-\bar{\nu}$ (4.5 $\pm$ 0.4) $\times 10^{-4}$		210
				$\Lambda e^-\bar{\nu}$ (5.74 $\pm$ 0.27) $\times 10^{-5}$		79
				$\dagger [n\pi^-\gamma]$ (4.6 $\pm$ 0.6) $\times 10^{-4}$		193
$m_{\Sigma^-} - m_{\Sigma^0} = 4.88$ $\pm 0.06$						

## Stable Particle Summary Table (cont'd)

Particle	$I(J^P)^a$	Mass <sup>b</sup> (MeV)	Mean life <sup>b</sup> $\tau$ (sec) $c\tau$ (cm)	Partial decay modes						
				Mode	Fraction <sup>b</sup>	$p$ (MeV/c) <sup>c</sup>				
<b>STRANGENESS -2 BARYONS<sup>a</sup></b>				$[\Xi^0 = u_{ss}, \Xi^- = d_{ss}]$						
$\Xi^0$	$\frac{1}{2}(\frac{1}{2}^+)^t$	1314.9 $\pm 0.6$	2.90 $\times 10^{-10}$ $\pm 0.10$  $c\tau=8.69$	$\Lambda\pi^0$	100%	135				
				$\Lambda\gamma$	( 0.5 $\pm$ 0.5 )%	184				
				$\Sigma^0\gamma$	( <7 )%	117				
				$p\pi^-$	( <3.6 ) $\times 10^{-5}$	$\Delta S$ 299				
				$pe^- \nu$	( <1.3 ) $\times 10^{-3}$	$\Delta S$ 323				
				$\Sigma^+ e^- \nu$	( <1.1 ) $\times 10^{-3}$	120				
				$\Sigma^- e^+ \nu$	( <0.9 ) $\times 10^{-3}$	SQ 112				
				$\Sigma^+ \mu^- \nu$	( <1.1 ) $\times 10^{-3}$	65				
				$\Sigma^- \mu^+ \nu$	( <0.9 ) $\times 10^{-3}$	SQ 49				
				$p\mu^- \nu$	( <1.3 ) $\times 10^{-3}$	$\Delta S$ 309				
				$\Xi^-$	$\frac{1}{2}(\frac{1}{2}^+)^t$	1321.32 $\pm 0.13$	1.642 $\times 10^{-10}$ $\pm 0.015$  $c\tau=4.92$	$\Lambda\pi^-$	100%	139
$\Lambda e^- \nu$	( 5.5 $\pm$ 0.6 ) $\times 10^{-4}$	S=2.0* 190								
$\Sigma^0 e^- \nu$	( 8.7 $\pm$ 1.7 ) $\times 10^{-5}$	123								
$\Lambda\mu^- \nu$	( 3.5 $\pm$ 3.5 ) $\times 10^{-4}$	163								
$\Sigma^0 \mu^- \nu$	( <8 ) $\times 10^{-4}$	70								
$n\pi^-$	( <1.9 ) $\times 10^{-5}$	$\Delta S$ 303								
$ne^- \nu$	( <3.2 ) $\times 10^{-3}$	$\Delta S$ 327								
$n\mu^- \nu$	( <1.5 ) $\times 10^{-2}$	$\Delta S$ 313								
$\Sigma^- \gamma$	( <1.2 ) $\times 10^{-3}$	118								
$p\pi^- \pi^-$	( <4 ) $\times 10^{-4}$	$\Delta S$ 223								
$p\pi^- e^- \nu$	( <4 ) $\times 10^{-4}$	$\Delta S$ 304								
$p\pi^- \mu^- \nu$	( <4 ) $\times 10^{-4}$	$\Delta S$ 250								
$\Xi^0 e^- \nu$	( <2.3 ) $\times 10^{-3}$	6								
<b>STRANGENESS -3 BARYON<sup>a</sup></b>								$[\Omega^- = s_{ss}]$		
$\Omega^-$	$0(\frac{3}{2}^+)^t$	1672.45 $\pm 0.29$	0.822 $\times 10^{-10}$ $\pm 0.013$  $c\tau=2.46$					$\Lambda K^-$	( 67.8 $\pm$ 0.7 )%	211
				$\Xi^0 \pi^-$	( 23.6 $\pm$ 0.7 )%	294				
				$\Xi^- \pi^0$	( 8.6 $\pm$ 0.4 )%	290				
				$\Xi^0 e^- \nu$	( 5.6 $\pm$ 2.8 ) $\times 10^{-3}$	319				
				$\Xi(1530)^0 \pi^-$	( 6.4 $\pm$ 5.1 ) $\times 10^{-4}$					
				$\Lambda\pi^-$	( <1.9 ) $\times 10^{-4}$	$\Delta S$ 449				
				$\Xi^- \gamma$	( <2.2 ) $\times 10^{-3}$	314				
				$\Xi^- \pi^+ \pi^-$	( 4.3 $\pm$ 3.4 ) $\times 10^{-4}$	190				
<b>CHARMED BARYON<sup>a</sup></b>				$[\Lambda_c^+ = udc]$						
$\Lambda_c^+$	$0(\frac{1}{2}^+)^t$	2281.2 $\pm 3.0$ $S=1.9^*$	(2.3 $^{+0.8}_{-0.5}$ ) $\times 10^{-13}$  $c\tau=0.007$	$pK^- \pi^+$	( 2.2 $\pm$ 1.0 )%	820				
				$p\bar{K}^0$	( 1.1 $\pm$ 0.7 )%	870				
				$p\bar{K}^0 \pi^+ \pi^-$	( <4, seen )%	751				
				$\Lambda$ anything	( 33 $\pm$ 29 )%					
				$\dagger[\Lambda\pi^+]$	( 0.6 $\pm$ 0.5 )%	861				
				$\Lambda\pi^+ \pi^+ \pi^-$	( < 3.1, seen )%	804				
				$\Sigma^0 \pi^+$	( seen )	822				
				$\dagger[pK^{*0}]$	( 0.48 $\pm$ 0.30 )%	681				
				$\Delta^+ K^-$	( 0.45 $\pm$ 0.27 )%	706				
				$pK^{*+} \pi^+$	( seen )	575				
				$e^+$ anything	( 4.5 $\pm$ 1.7 )%					
				$\dagger[pe^+ \text{ anything}]$	( 1.8 $\pm$ 0.9 )%					
				$\Lambda e^+$ anything	( 1.1 $\pm$ 0.8 )%					
$\Xi_c^+$ [was $A^+$ ]	}	Not established; see the Stable Particle Full Listings								
$\Omega_c^0$										
$\Lambda_b^0$										
Searches for top hadrons, free quarks, magnetic monopoles, axions, supersymmetric particles, and other stable particles				} See Addendum and the Stable Particle Full Listings						



ADDENDUM TO  
Stable Particle Summary Table

<b>Magnetic moment</b>			
$e^u$	1.001 159 652 209 $\pm 0.000\ 000\ 000\ 031$	$\frac{e\hbar}{2m_e c}$	
<b><math>\mu</math> decay parameters <sup>v</sup></b>			
$\mu^u$	1.001 165 923 $\pm 0.000\ 000\ 009$	$\frac{e\hbar}{2m_\mu c}$	$\rho = -0.752 \pm 0.003$ $\delta = 0.755 \pm 0.009$ $\xi'' = 0.65 \pm 0.36$
$\tau$	<b>Michel parameter</b> $\rho = 0.71 \pm 0.08$	$\alpha/A = (0.4 \pm 4.3) \times 10^{-3}$ $\beta/A = (3.9 \pm 6.2) \times 10^{-3}$ $\bar{\eta} = 0.016 \pm 0.076$	$\eta = -0.007 \pm 0.079$ $\xi' = 0.998 \pm 0.042$ $\xi P_\mu \delta / \rho > 0.9966^w$ $\alpha'/A = (-0.2 \pm 4.3) \times 10^{-3}$ $\beta'/A = (1.5 \pm 6.3) \times 10^{-3}$
$\eta$	<b>Mode</b> $\pi^+ \pi^- \pi^0$ $\pi^+ \pi^- \gamma$	<b>Left-right asymmetry</b> (0.12 $\pm$ 0.17)% (0.88 $\pm$ 0.40)%	<b>Sextant asymmetry</b> (0.19 $\pm$ 0.16)% <b>Quadrant asymmetry</b> (-0.17 $\pm$ 0.17)% $\beta = 0.047 \pm 0.062$ S=1.5*
$K$	<b>Slope parameters for <math>K \rightarrow 3\pi^x</math></b> $K^+ \rightarrow \pi^+ \pi^+ \pi^-$ $g = -0.215 \pm 0.004$ S=1.4* $K^- \rightarrow \pi^- \pi^- \pi^+$ $g = -0.217 \pm 0.007$ S=2.5* $K^\pm \rightarrow \pi^0 \pi^0 \pi^\pm$ $g = 0.607 \pm 0.030$ S=1.3* $K_L^0 \rightarrow \pi^+ \pi^- \pi^0$ $g = 0.670 \pm 0.014$ S=1.6* See Full Listings for quadratic coefficients.		<b>Form factors for <math>K_{e3}</math> decays <sup>v</sup></b> $K_{e3}^+ \left\{ \begin{array}{l} \lambda_+ = 0.029 \pm 0.004 \\  f_S/f_+  = 0.13 \pm 0.04 \\  f_T/f_+  = 0.22 \pm 0.14 \end{array} \right. S=1.3^* \quad K_{e3}^0 \left\{ \begin{array}{l} \lambda_+ = 0.0300 \pm 0.0016 \\  f_S/f_+  < 0.04 \\  f_T/f_+  < 0.23 \end{array} \right. S=1.2^* \\ K_{\mu 3}^+ \left\{ \begin{array}{l} \lambda_+ = 0.033 \pm 0.008 \\ \lambda_0 = 0.004 \pm 0.007 \\  f_T/f_+  = 0.02 \pm 0.12 \end{array} \right. S=1.6^* \quad K_{\mu 3}^0 \left\{ \begin{array}{l} \lambda_+ = 0.034 \pm 0.005 \\ \lambda_0 = 0.025 \pm 0.006 \\  f_T/f_+  = 0.12 \pm 0.12 \end{array} \right. S=2.3^*$
<b><math>\Delta S = -\Delta Q</math> in <math>K_{e3}^0</math> decay</b> Re $x = 0.009 \pm 0.020$ S=1.4* Im $x = -0.004 \pm 0.026$ S=1.1*		<b>CP-violation parameters <sup>v,i</sup></b> $ \eta_{+-}  = (2.275 \pm 0.021) \times 10^{-3}$ $\left  \frac{\epsilon'}{\epsilon} \right  = (-3 \pm 4) \times 10^{-3}$ $ \eta_{00}  = (2.299 \pm 0.036) \times 10^{-3}$ $\phi_{+-} = (44.6 \pm 1.2)^\circ$ $\phi_{00} = (54 \pm 5)^\circ$ Re $\epsilon = (1.621 \pm 0.088) \times 10^{-3}$ $ \eta_{+-0} ^2 < 0.12$ $ \eta_{000} ^2 < 0.1$ $\delta = (0.330 \pm 0.012)\%$	
<b>Magnetic moment</b> ( $e\hbar/2m_p c$ )	<b>Decay parameters <sup>z</sup></b>		
	<u>Measured</u>	<u>Derived</u>	<u>Coupling constant ratios</u>
	$\alpha$	$\phi$ (degree)	$\gamma$ $\Delta$ (degree)
$p^u$	2.7928444 $\pm 0.0000011$		
$n^u$	-1.91304308 $\pm 0.00000054$	$pe^- \nu$ $g_A/g_V = -1.254 \pm 0.006$ $\phi_{AV} = (180.11 \pm 0.17)^\circ$	
$\Lambda^u$	-0.613 $\pm 0.004$	$p\pi^-$ +0.642 $\pm$ 0.013 $n\pi^0$ +0.646 $\pm$ 0.044 $pe^- \nu$	-6.5 $\pm$ 3.5 0.76 7.7 $\pm$ 4.1 $g_A/g_V = -0.694 \pm 0.025$ S=1.3*
$\Sigma^+$	2.379 $\pm 0.020$	$p\pi^0$ -0.980 $\pm$ 0.015 $n\pi^+$ +0.068 $\pm$ 0.013 $p\gamma$ -0.72 $\pm$ 0.29	36 $\pm$ 34 167 $\pm$ 20 S=1.1* 0.16 187 $\pm$ 6 -0.97 -73 $^{+13}_-10$
$\Sigma^-$	-1.14 $\pm 0.05$ S=1.9*	$n\pi^-$ -0.068 $\pm$ 0.008 $ne^- \nu$ $\Lambda e^- \nu$	10 $\pm$ 15 0.98 249 $^{+12}_-116$ $ g_A/g_V  = 0.362 \pm 0.043$ S=1.7* $g_V/g_A = 0.01 \pm 0.10$ S=1.5* $g_{WM}/g_A = 2.4 \pm 1.7$
$\Xi^0$	-1.250 $\pm 0.014$	$\Lambda\pi^0$ -0.413 $\pm$ 0.022 S=2.0*	21 $\pm$ 12 0.85 218 $^{+12}_-18$
$\Xi^-$	-0.69 $\pm 0.04$	$\Lambda\pi^-$ -0.455 $\pm$ 0.015 S=1.8* $\Lambda e^- \nu$	4 $\pm$ 5 0.89 187 $\pm$ 9 $g_A/g_V = -0.25 \pm 0.05$
$\Omega^-$	$\Lambda K^-$ -0.024 $\pm$ 0.028 $\Xi^0 \pi^-$ +0.09 $\pm$ 0.14 $\Xi^- \pi^0$ +0.05 $\pm$ 0.21		

## ADDENDUM TO Stable Particle Summary Table

### SEARCH LIMITS

Search limits usually are critically dependent on assumptions. Further details and other limits are given in the Full Listings. For complete information, the original literature should be consulted. All limits given below are for CL=0.90, except those marked with an asterisk (\*), which are for CL=0.95.

#### Leptons

$L^+$ — charged lepton	$M > 22.5 \text{ GeV}^*$	if $\tau < 10 \text{ nsec}$
$E^0$ — neutral paralepton	$M > 22.5 \text{ GeV}^*$ $> 24.5 \text{ GeV}^*$	if $E^0 e W$ vertex is $V - A$ if $E^0 e W$ vertex is $V + A$
$e^{*\pm}$ — excited electron	$M > 72 \text{ GeV}^*$	if $\lambda = 1$
$\mu^{*\pm}$ — excited muon	$M > 25 \text{ GeV}^*$	if $\lambda = 1$
$\nu$ oscillation $\nu_\mu \rightarrow \nu_e$ ( $\theta =$ mixing angle)	$\sin^2 2\theta < 3.4 \times 10^{-3}$ $\Delta M^2 < 0.2 \text{ eV}^2$	if $\Delta M^2$ is large if $\sin^2 2\theta = 1$

#### Number of Light Neutrino Types (including $\nu_e, \nu_\mu,$ and $\nu_\tau$ )

$N < 7$	from $Z$ width in standard model (cosmological limits are lower)
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#### Additional $W$ Bosons

$W$ right-handed	$M \geq 475 \text{ GeV}^*$	assuming light right-handed neutrino
$W$ with standard couplings decaying to $e\nu$	$M \geq 210 \text{ GeV}$	

#### Additional $Z$ Bosons

[First four are not in Full Listings; see L.S. Durkin and P. Langacker, Phys. Lett. **166B**, 436 (1986), and V. Barger et al., Phys. Rev. Lett. **56**, 30 (1986).]

$Z'$ of $SU(2)_L \times SU(2)_R \times U(1)$	$M \geq 275 \text{ GeV}$	if magnitudes of $L$ & $R$ coupling constants are equal
$Z_\chi$ of $SO(10) \rightarrow SU(5) \times U(1)$	$M \geq 220 \text{ GeV}$	coupling constant derived from G.U.T.
$Z_\psi$ of $E_6 \rightarrow SO(10) \times U(1)_\psi$	$M \geq 114 \text{ GeV}$	coupling constant derived from G.U.T.
$Z_\eta$ of $E_6 \rightarrow SU(3) \times SU(2) \times U(1) \times U(1)_\eta$	$M \geq 126 \text{ GeV}$	coupling constant derived from G.U.T.; charges are $Q_\eta = \sqrt{3/8}Q_\chi - \sqrt{5/8}Q_\psi$
$Z$ with standard couplings decaying to $e^+e^-$	$M \geq 160 \text{ GeV}$	
$Z$ right-handed	$M \geq 150 \text{ GeV}^*$	Rizzo-Senjanovic formalism

#### Higgs Bosons (or Technipions)

$H^\pm$	$M \geq 17 \text{ GeV}^*$ $M \geq 13 \text{ GeV}^*$	if $B(H^\pm \rightarrow \tau\nu) \geq 0.25$ from hadronic decays, if $B(H^\pm \rightarrow bc)$ is small
$H^0$	No substantial limits	

#### Axions (including nonstandard axions)

$A$ — axion with mass $< 100 \text{ MeV}$	$M \geq 1 \text{ MeV}$	assuming $\tau_A > 10^{-11} \text{ sec}$
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#### Supersymmetric Particles

Assumptions include: 1) photino is lightest supersymmetric particle; 2)  $R$ -parity is conserved; 3) mass of exchanged particles is less than about 250 GeV; 4) unless otherwise stated,  $M(\tilde{J}_L) = M(\tilde{J}_R)$ , and  $u, d, s, c,$  and  $b$  scalar quarks are degenerate in mass.

When two limits are quoted, the parenthetical number has  $M(\tilde{J}_R) \gg M(\tilde{J}_L)$ .

$\tilde{\gamma}$ — photino ( $\tilde{\gamma}$ limits are from cosmology)	$M \leq 100 \text{ eV}$	
	or $\begin{cases} M \geq 0.5 \text{ GeV} \\ \geq 2 \text{ GeV} \\ \geq 5 \text{ GeV} \end{cases}$	if $M_{\text{heavy}} = \begin{cases} 20 \text{ GeV} \\ 40 \text{ GeV} \\ 100 \text{ GeV} \end{cases}$
$\tilde{\nu}$ — scalar neutrino	$[M(\tilde{\nu}_\tau) + M(\tilde{\nu}_{e,\mu})] \geq 1.7 \text{ GeV}$	
$\tilde{e}$ — scalar electron	$M \geq 51 (42) \text{ GeV}$ $\geq 33 (26) \text{ GeV}$ $\geq 22 (21) \text{ GeV}^*$	if $M_{\tilde{\gamma}} = \begin{cases} 0 \text{ GeV} \\ 10 \text{ GeV} \\ < 19 (15) \text{ GeV} \end{cases}$
$\tilde{\mu}$ — scalar muon	$M > 21 \text{ GeV}^*$	if $M_{\tilde{\gamma}} \leq 15 \text{ GeV}$
$\tilde{\tau}$ — scalar tau	$M > 18 \text{ GeV}^*$	if $M_{\tilde{\gamma}} < 13 \text{ GeV}$
$\tilde{q}$ — scalar quark	$M > 60\text{--}70 \text{ GeV}$	depending on $M_{\tilde{g}}$
$\tilde{g}$ — gluino	$M > 50\text{--}60 \text{ GeV}$	depending on $M_{\tilde{q}}$
	or $\begin{cases} M = 0\text{--}4 \text{ GeV} \\ = 2\text{--}4 \text{ GeV} \\ = 3\text{--}4 \text{ GeV} \\ \text{Low mass not allowed} \end{cases}$	if $M_{\tilde{q}} > \begin{cases} 330 \text{ GeV} \\ 300 \text{ GeV} \\ 150 \text{ GeV} \\ < 65 \text{ GeV} \end{cases}$
$\chi^\pm$ — chargino (mixtures of $\tilde{W}^\pm$ and $\tilde{H}_i^\pm$ )	$M \geq 22 \text{ GeV}^*$ $M > 16.5 \text{ GeV}^*$	Whether $M = 0\text{--}0.5 \text{ GeV}$ is allowed is the subject of dispute. if $M_{\tilde{\nu}} < 18 \text{ GeV}$ and if $B(\tilde{\chi}^\pm \rightarrow \ell\nu) \approx 100\%$ if hadronic decays dominate
$\chi^0$ — neutralino (mixtures of $\tilde{\gamma}, \tilde{Z}^0,$ and $\tilde{H}_i^0$ )	Limits are extremely model dependent. Reported limits based on a set of very specific assumptions may exclude $M \approx 6\text{--}33 \text{ GeV}$ .	

## Stable Particle Summary Table (cont'd)

- \*  $S = \text{Scale factor} = \sqrt{\chi^2/(N-1)}$ , where  $N \approx$  number of experiments.  $S$  should be  $\approx 1$ . If  $S > 1$ , we have enlarged the error of the mean,  $\delta\bar{x}$ ; i.e.,  $\delta\bar{x} \rightarrow S\delta\bar{x}$ . This convention is still inadequate, since if  $S \gg 1$  the experiments are probably inconsistent, and therefore the real uncertainty is probably even greater than  $S\delta\bar{x}$ . See the Introduction, and ideograms in the Full Listings.
- † Square brackets indicate subreactions of some previous unbracketed decay mode(s). Reactions in one set of brackets may overlap with reactions in another set of brackets. A radiative mode such as  $\pi \rightarrow \mu\nu\gamma$  is a subreaction of its parent mode  $\pi \rightarrow \mu\nu$ .
- a. The strangeness  $S$ , charm  $C$ , and bottom  $B$  of the hadrons which appear in the Summary Table are given in the Nonrelativistic Quark Model section.
- b. Quoted upper limits correspond to a 90% confidence level. Masses, mean lives, and partial rates are evaluated assuming equality for particles and antiparticles. Decays which are forbidden by a conservation law are indicated by the following abbreviations:  $LF \equiv$  lepton family number,  $L \equiv$  total lepton number,  $B \equiv$  baryon number,  $Q \equiv$  electric charge,  $C \equiv$  charge conjugation,  $P \equiv$  parity,  $\Delta S \equiv (\Delta S = 2)$ ,  $SQ \equiv (\Delta S = \Delta Q)$ ,  $FC \equiv$  flavor-changing neutral current. See the Tests of Conservation Laws section for further details.
- c. For a 2-body decay mode, this is the momentum of the decay products in the decay rest frame. For a 3-or-more-body mode, this is the maximum momentum any of the products can have in this frame.
- d. 99% confidence level. See footnote in the Full Listings.
- e. See the Full Listings for energy limits used in this measurement.
- f. Theoretical value; see also the Full Listings.
- g. Structure-dependent part with positive (SD+) and negative (SD-) photon helicity.
- h. The direct emission branching fraction is  $(1.56 \pm .35) \times 10^{-2}$ . For interference terms in  $K \rightarrow \mu\nu\gamma$ , see the Full Listings.
- i. The  $K_S^0 \rightarrow \pi\pi$  and  $K_L^0 \rightarrow \pi\pi$  branching fractions are from our branching-fraction and rate fits and do not include results of  $K_L^0 - K_S^0$  interference experiments. The  $\pi\pi$  rate results are combined with the interference results to obtain the  $|\eta_{+-}|$  and  $|\eta_{00}|$  values given in the Addendum.
- j. The stronger limit  $< 2 \times 10^{-9}$  of Clark et al., Phys. Rev. Lett. **26**, 1667 (1971) is not listed because of possible (but unknown) systematic errors; see the Full Listings.
- k. Many of the  $D^\pm$  and  $D^0$  hadronic branching fractions have been rescaled to reflect the cross section for charm production determined in a recent experiment. See the Full Listings for details.
- l. This is a weighted average of  $D^\pm$  (44%) and  $D^0$  (56%) branching fractions.
- m.  $D_1^0 - D_2^0$  limits inferred from limit on  $D^0 \rightarrow \bar{D}^0 \rightarrow \mu^-$  anything.
- n.  $D_s$  mass determined from  $\phi\pi$  mode. See note on conflicting  $D_s$  meson results in the Full Listings. Quantum numbers shown are favored but not yet established.
- p. Quantum numbers not measured. Values shown are quark model predictions.
- q. Except for the neutral-current decay modes ( $\ell^+\ell^-$  anything), only data from  $\Upsilon(10575)$  decays are used. Behrends et al. [Phys. Rev. Lett. **50**, 881 (1983)] estimate the  $\Upsilon(10575) \rightarrow B^+B^-$  and  $\Upsilon(10575) \rightarrow B^0\bar{B}^0$  branching fractions to be  $60 \pm 2$  and  $40 \pm 2\%$ .
- r. First limit is geochemical and independent of decay mode. Second limit assumes that the dominant decay modes are among those investigated. For antiprotons the best mean life limit, inferred from observation of cosmic ray  $\bar{p}$ 's, is  $\tau_{\bar{p}} > 10^7$  yrs, the cosmic ray storage time.
- s. Limit from neutrality-of-matter experiments. Assumes  $|q_n| = |q_p| - |q_e|$ .
- t.  $P$  for  $\Xi$ ,  $J^P$  for  $\Omega^-$  and  $\Sigma^0$ , and  $J$  for  $\Lambda_c^+$  not yet measured. Values shown are quark model predictions.
- u. For limits on electric dipole moment (forbidden by  $P$  and  $T$  invariance), see the Tests of Conservation Laws section.
- v. For more details and definitions of parameters, see the Full Listings.
- w.  $P$  is muon longitudinal polarization from  $\pi$  decay. In standard  $V-A$  theory,  $P = 1$  and  $\rho = \delta = 3/4$ .
- x. The definition of the slope parameter of the Dalitz plot is as follows [see also note in the Full Listings]:

$$|M|^2 = 1 + g(s_3 - s_0)/m_{\pi^+}^2.$$

- y. The definition for the CP-violation parameters is as follows [see also note in the Full Listings]:

$$\eta_{+-} = |\eta_{+-}| e^{i\phi_{+-}} = \frac{A(K_L^0 \rightarrow \pi^+\pi^-)}{A(K_S^0 \rightarrow \pi^+\pi^-)} \quad \eta_{00} = |\eta_{00}| e^{i\phi_{00}} = \frac{A(K_L^0 \rightarrow \pi^0\pi^0)}{A(K_S^0 \rightarrow \pi^0\pi^0)}$$

$$\delta = \frac{\Gamma(K_L^0 \rightarrow \ell^+) - \Gamma(K_L^0 \rightarrow \ell^-)}{\Gamma(K_L^0 \rightarrow \ell^+) + \Gamma(K_L^0 \rightarrow \ell^-)}, \quad |\eta_{+-0}|^2 = \frac{\Gamma(K_S^0 \rightarrow \pi^+\pi^-\pi^0)_{\text{CP viol.}}}{\Gamma(K_L^0 \rightarrow \pi^+\pi^-\pi^0)}, \quad |\eta_{000}|^2 = \frac{\Gamma(K_S^0 \rightarrow \pi^0\pi^0\pi^0)_{\text{CP viol.}}}{\Gamma(K_L^0 \rightarrow \pi^0\pi^0\pi^0)}.$$

- z. The definition of these quantities is as follows [for more details and sign convention, see note in the Full Listings]:

$$\alpha = \frac{2|s||p|\cos\Delta}{|s|^2 + |p|^2} \quad \beta = \sqrt{1 - \alpha^2} \sin\phi \quad g_A, g_V, g_{WM} \text{ defined by } \langle B_f | \gamma_\lambda (g_V - g_A \gamma_5) + (g_{WM}/m_{B_i}) \sigma^{\lambda\nu} q_\nu | B_i \rangle$$

$$\beta = \frac{-2|s||p|\sin\Delta}{|s|^2 + |p|^2} \quad \gamma = \sqrt{1 - \alpha^2} \cos\phi \quad \phi_{AV} \text{ defined by } g_A/g_V = |g_A/g_V| e^{i\phi_{AV}}$$

# Meson Summary Table

April 1986

In addition to the entries in the Meson Summary Table, the Meson Full Listings contain all substantial claims for meson resonances. See Contents of the Meson Full Listings at end of this Summary Table.

*Quantities in italics are new or have changed by more than one (old) standard deviation since April 1984.*

Particle	$J^{PC}$ <sup>a</sup> — <i>estab.</i>	Mass <i>M</i> (MeV)	Full width $\Gamma$ (MeV)	Partial decay modes		
				Mode	Fraction(%) [Upper limits (%) are 90% CL]	<i>p</i> <sup>b</sup> (MeV/c)
<b>NONFLAVORED MESONS</b>						
$\pi^\pm$	$1^-(0^{-+})$	139.57	0.0	See Stable Particle Summary Table		
$\pi^0$		134.96	7.57 $\pm 0.32$ eV			
$\eta$	$0^+(0^{-+})$	548.8 $\pm 0.6$	1.05 $\pm 0.15$ keV	Neutral Charged	70.9 29.1	See Stable Particle Summary Table
$\rho(770)$	$1^+(1^{--})$	770 $\pm 3^S$	153 $\pm 2$ MeV	$\pi\pi$ $\pi\gamma$ $\mu^+\mu^-$ $e^+e^-$ $\eta\gamma$	$\approx 100$ 0.046 $\pm$ 0.005 0.0067 $\pm$ 0.0012 <sup>d</sup> 0.0045 $\pm$ 0.0002 <sup>d</sup> seen	358 372 370 384 189
$\Gamma_{ee} = (6.9 \pm 0.3)$ keV				For upper limits, see footnote <i>e</i>		
$M$ and $\Gamma$ from neutral mode.						
$\omega(783)$	$0^-(1^{--})$	782.6 $\pm 0.2$ S=1.1*	9.8 $\pm 0.3$	$\pi^+\pi^-\pi^0$ $\pi^0\gamma$ $\pi^+\pi^-$ $\pi^0\mu^+\mu^-$ $e^+e^-$ $\eta\gamma$	89.6 $\pm$ 0.5 8.7 $\pm$ 0.5 1.7 $\pm$ 0.2 0.010 $\pm$ 0.002 0.0067 $\pm$ 0.0004 S=1.2* seen	327 380 366 349 391 199
$\Gamma_{ee} = (0.66 \pm 0.04)$ keV S=1.2*				For upper limits, see footnote <i>f</i>		
$\eta'(958)$	$0^+(0^{-+})$	957.57 $\pm 0.25$	0.24 $\pm 0.03$	$\eta\pi\pi$ $\rho^0\gamma$ $\omega\gamma$ $\gamma\gamma$ $3\pi^0$ $\mu^+\mu^-\gamma$	65.2 $\pm$ 1.6 30.0 $\pm$ 1.6 2.7 $\pm$ 0.5 1.9 $\pm$ 0.2 0.17 $\pm$ 0.04 0.009 $\pm$ 0.002	231 170 159 479 430 467
				For upper limits, see footnote <i>g</i>		
$f_0(975)$ was <i>S</i> (975)	$0^+(0^{++})$	975 <sup>c</sup> $\pm 4$ S=1.4*	33 <sup>c</sup> $\pm 6$	$\pi\pi$ $K\bar{K}$	78 $\pm$ 3 22 $\pm$ 3	467
$a_0(980)$ was $\delta$ (980)	$1^-(0^{++})$	983 <sup>h</sup> $\pm 2$	54 <sup>h</sup> $\pm 7$	$\eta\pi$ $K\bar{K}$	seen seen	320
$\phi(1020)$	$0^-(1^{--})$	1019.5 $\pm 0.1$ S=1.2*	4.22 $\pm 0.13$	$K^+K^-$ $K_L^0 K_S^0$ $\pi^+\pi^-\pi^0$ (incl. $\rho\pi$ ) $\eta\gamma$ $\pi^0\gamma$ $e^+e^-$ $\mu^+\mu^-$ $\pi^+\pi^-$	49.5 $\pm$ 1.5 34.3 $\pm$ 0.9 14.9 $\pm$ 1.4 1.3 $\pm$ 0.1 0.131 $\pm$ 0.013 0.031 $\pm$ 0.001 0.025 $\pm$ 0.003 0.02 $\pm$ 0.01	S=1.9* 127 S=1.3* 110 S=2.3* 462 S=1.1* 362 501 510 499 490
$\Gamma_{ee} = (1.31 \pm 0.06)$ keV				For upper limits, see footnote <i>i</i>		
$h_1(1190)$ was <i>H</i> (1190)	$0^-(1^{+-})$	1190 $\pm 60$	320 $\pm 50$	$\rho\pi$	seen	327

## Meson Summary Table (*cont'd*)

Particle	$f^G(J^{PC})^a$ — <i>estab.</i>	Mass $M$ (MeV)	Full width $\Gamma$ (MeV)	Partial decay modes		
				Mode	Fraction(%) [Upper limits (%) are 90% CL]	$p^b$ (MeV/c)
$b_1(1235)$ <i>was B(1235)</i>	$1^+(1^{+-})$	1233 $\pm 10^{\S}$	150 $\pm 10^{\S}$	$\omega\pi$ $\eta\rho$ [ $D/S$ amplitude ratio = $0.29 \pm 0.05$ ] For upper limits, see footnote <i>j</i>	dominant < 10	350
$f_2(1270)$ <i>was f(1270)</i>	$0^+(2^{++})$	1274 $\pm 5^{\S}$	176 $\pm 20^{\S}$	$\pi\pi$ $2\pi^+2\pi^-$ $K\bar{K}$ $\eta\eta$ $\gamma\gamma$ $\pi^+\pi^-2\pi^0$ For upper limits, see footnote <i>k</i>	84.3 $\pm$ 1.2 2.9 $\pm$ 0.4 2.9 $\pm$ 0.2 0.31 $\pm$ 0.08 0.0015 $\pm$ 0.0002 seen	622 559 398 323 637 562
$a_1(1270)$ <i>was A<sub>1</sub>(1270)</i>	$1^-(1^{++})$	1275 $\pm 28$	316 $\pm 45$	$\rho\pi$ $\pi(\pi\pi)_{S\text{-wave}}$ $\pi^+\gamma$	dominant < 0.7 <sup>§</sup> 0.20 $\pm$ 0.08	389 599 630
$f_1(1285)$ <i>was D(1285)</i>	$0^+(1^{+-})$	1283 $\pm 5^{\S}$	25 $\pm 3^{\S}$	$K\bar{K}\pi$ $\eta\pi\pi$ † [ $a_0(980)\pi$ ] $4\pi$ (prob. $\rho\pi\pi$ )	11 $\pm$ 3 49 $\pm$ 6 36 $\pm$ 7 40 $\pm$ 7	302 482 236 564
$f_0(1300)$ <i>was <math>\epsilon(1300)</math></i>	$0^+(0^{++})$	$\sim 1300$	150–400	$\pi\pi$ $K\bar{K}$ $\eta\eta$	$\sim 90$ $\sim 10$ seen	635 418 348
$\pi(1300)$	$1^-(0^{-+})$	1300 <sup>§</sup> $\pm 100^{\S}$	200–600	$\rho\pi$ $\pi(\pi\pi)_{S\text{-wave}}$	seen seen	407 612
Not a well-established resonance.						
$a_2(1320)$ <i>was A<sub>2</sub>(1320)</i>	$1^-(2^{++})$	1318 <sup>§</sup> $\pm 5^{\S}$	110 $\pm 5^{\S}$	$\rho\pi$ $\eta\pi$ $\omega\pi\pi$ $K\bar{K}$ $\pi\gamma$ $\gamma\gamma$ $\eta'\pi$	70.1 $\pm$ 2.2 14.5 $\pm$ 1.2 10.6 $\pm$ 2.5 4.9 $\pm$ 0.8 0.27 $\pm$ 0.06 0.0008 $\pm$ 0.0001 < 2 (CL=97%)	419 534 361 434 652 659 286
$f_1(1420)^{\ddagger}$ <i>was E(1420)</i>	$0^+(1^{+-})$	1422 <sup>§</sup> $\pm 10^{\S}$	56 $\pm 3$	$K\bar{K}\pi$ (incl. $K^*\bar{K} + K\bar{K}^*$ ) $\eta\pi\pi$ † [ $a_0(980)\pi$ ]	seen possibly seen possibly seen	423 565 348
$\eta(1440)^{\ddagger}$ <i>was <math>\iota(1440)</math></i>	$0^+(0^{-+})$	1440 <sup>§</sup> $\pm 20^{\S}$	76 $\pm 8$ S=1.3*	$K\bar{K}\pi$ (incl. $K^*\bar{K} + K\bar{K}^*$ ) $\eta\pi\pi$ † [ $a_0(980)\pi$ ]	<i>seen</i> <i>seen</i> <i>seen</i>	441 579 366
$f_2'(1525)$ <i>was f'(1525)</i>	$0^+(2^{++})$	1525 <sup>§</sup> $\pm 5^{\S}$	70 $\pm 10^{\S}$	$K\bar{K}$ $\pi\pi$ $\gamma\gamma$	dominant possibly seen 0.00016 $\pm$ 0.00007	578 750 763
$f_0(1590)$	$0^+(0^{++})$	1587 $\pm 16$	287 $\pm 50$	$\eta'\eta$ $\eta\eta$	<i>dominant</i> <i>large</i>	246 575
Seen by one group only.						
$\rho(1600)$	$1^+(1^{--})$	1590 <sup>§</sup> $\pm 20^{\S}$	260 <sup>§</sup> $\pm 100^{\S}$	$4\pi$ (incl. $\rho\pi^+\pi^-, a_1(1270)\pi$ ) $\pi\pi$ $K^*\bar{K} + \bar{K}^*K$ $\eta\pi\pi$ $K\bar{K}$ $e^-e^-$	60 $\pm$ 7 <sup>§</sup> 23 $\pm$ 7 <sup>§</sup> 9 $\pm$ 2 7 $\pm$ 2 1 $\pm$ 0.5 0.0029 $\pm$ 0.0011	733 783 377 669 623 795
$\Gamma_{ee} = (7.5 \pm 1.5) \text{ keV}$						

## Meson Summary Table (cont'd)

Particle	$I^G(J^{PC})^a$ _____ <i>estab.</i>	Mass $M$ (MeV)	Full width $\Gamma$ (MeV)	Partial decay modes		$p^b$ (MeV/c)
				Mode	Fraction(%) [Upper limits (%) are 90% CL]	
$\omega_3(1670)$ <i>was</i> $\omega(1670)$	$0^-(3^{--})$	1668 $\pm 5$	166 <sub>S</sub> $\pm 15^{\delta}$ S=1.1*	$3\pi$ $\dagger[\rho\pi]$ $5\pi$ $\dagger[\omega\pi\pi$ (prob. $b_1(1235)\pi$ )	seen seen seen seen	806 648 740 616
$\pi_2(1680)$ <i>was</i> $A_3(1680)$	$1^-(2^{-+})$	1680 <sub>S</sub> $\pm 30^{\delta}$	250 <sub>S</sub> $\pm 50^{\delta}$	$f_2(1270)\pi$ $\rho\pi$ $\pi(\pi\pi)_{S=wave}$ $K^*K + \bar{K}^*K$ For upper limits, see footnote $\ell$	53 $\pm$ 5 34 $\pm$ 6 9 $\pm$ 5 4 $\pm$ 1.4	336 656 813 459
$\phi(1680)$	$0^-(1^{--})$	1685 <sub>S</sub> $+75^{\delta}$ -15	130 <sub>S</sub> $\pm 50^{\delta}$	$K^*\bar{K} + \bar{K}^*K$ $\omega\pi\pi$ $K\bar{K}$ $e^+e^-$ $\pi^+\pi^-\pi^0$	<i>seen</i> <i>seen</i> <i>seen</i> <i>seen</i> <i>possibly seen</i>	466 624 683 842 814
<i>Not a well-established resonance.</i>						
$\rho_3(1690)$ <i>was</i> $g(1690)$	$1^+(3^{--})$	1691 <sub>S</sub> $\pm 5^{\delta}$	200 <sub>S</sub> $\pm 20^{\delta}$	$2\pi$ $4\pi$ (incl. $\pi\pi\rho, \rho\rho, a_2\pi, \omega\pi$ ) $K\bar{K}\pi$ (incl. $K^*\bar{K} + \bar{K}^*K$ ) $K\bar{K}$	23.8 $\pm$ 1.3 70.9 $\pm$ 1.9 3.8 $\pm$ 1.2 1.5 $\pm$ 0.3	834 787 625 684
$J^P, M,$ and $\Gamma$ from the $2\pi$ and $K\bar{K}$ modes.						
$f_2(1720)$ <i>was</i> $\theta(1690)$	$0^+(2^{++})$	1716 $\pm 8$	134 $\pm 19$	$\eta\eta$ $K\bar{K}$	<i>seen</i> <i>seen</i>	660 699
$\phi_J(1850)$ <i>was</i> $\phi(1850)$	$0^-(3^{--})$	1853 $\pm 10$	96 $\pm 32$	$K\bar{K}$ $K^*\bar{K} + \bar{K}^*K$	<i>seen</i> <i>seen</i>	784 601
$f_4(2030)$ <i>was</i> $h(2030)$	$0^+(4^{++})$	2026 $\pm 12$	200 $\pm 13$	$\pi\pi$ $K\bar{K}$ $\eta\eta$	17 $\pm$ 2 0.7 $^{+0.4}_{-0.2}$ 0.22 $\pm$ 0.10	1004 883 852
<b><math>c\bar{c}</math> MESONS</b>						
$\eta_c(2980)$	$0^+(0^{-+})$	2981 $\pm 2$	11 $\pm 4$	Decay modes into hadronic resonances		
				$\eta'\pi\pi$	4.1 $\pm$ 1.7	1321
				$K^*{}^0K^-\pi^+ + c.c.$	2.0 $\pm$ 0.7	1275
				$K^*\bar{K}^*$	0.9 $\pm$ 0.5	1194
				$\phi\phi$	0.8 $\pm$ 0.3	1087
				$a_0(980)\pi$	< 2.0	1324
				$a_2(1320)\pi$	< 2.0	1194
				$\rho\rho$	< 1.4	1277
				$f_2(1270)\eta$	< 1.1	1144
				$\omega\omega$	< 0.3	1269
				Decay modes into stable hadrons		
				$K\bar{K}\pi$	5.4 $\pm$ 1.8	S=1.1* 1377
				$\eta\pi\pi$	4.7 $\pm$ 1.5	1426
				$K^+K^-\pi^+\pi^-$	1.7 $\pm$ 0.6	1343
				$2(\pi^+\pi^-)$	1.30 $\pm$ 0.47	1458
				$p\bar{p}$	0.12 $\pm$ 0.06	1158
				$\eta K\bar{K}$	< 3.1	1261
				$\pi^+\pi^-\bar{p}\bar{p}$	< 1.2	1025
				Radiative decay modes		
				$\gamma\gamma$	< 0.06	1491

## Meson Summary Table (cont'd)

Particle	$J^G(J^{PC})^a$ — <i>estab.</i>	Mass $M$ (MeV)	Full width $\Gamma$ (MeV)	Partial decay modes		
				Mode	Fraction(%) [Upper limits (%) are 90% CL]	$p^b$ (MeV/c)
$J/\psi(3097)$	$0^-(1^{--})$	3096.9 $\pm 0.1$	0.063 $\pm 0.009$	$e^+e^-$ $\mu^+\mu^-$ hadrons + radiative	6.9 $\pm$ 0.9 6.9 $\pm$ 0.9 86.2 $\pm$ 2.0	1548 1545
$\Gamma_{ee} = (4.7 \pm 0.3) \text{ keV}$ (assuming $\Gamma_{ee} = \Gamma_{\mu\mu}$ )				Decay modes into hadronic resonances		
Decay modes into stable hadrons				$\dagger[\rho\pi$	1.27 $\pm$ 0.09	1449
$\dagger[2(\pi^+\pi^-)\pi^0$		3.4 $\pm$ 0.3	1496	$\omega 2\pi^+ 2\pi^-$	0.85 $\pm$ 0.34	1392
$3(\pi^+\pi^-)\pi^0$		2.9 $\pm$ 0.6	1433	$\rho a_2(1320)$	0.84 $\pm$ 0.45	1126
$\pi^+\pi^-\pi^0$		1.50 $\pm$ 0.15	1533	$K\bar{K}^*(892)$	0.75 $\pm$ 0.09	S=1.8* 1373
$\pi^+\pi^-\pi^0 K^+ K^-$		1.2 $\pm$ 0.3	1368	$\omega\pi^+\pi^-$	0.74 $\pm$ 0.12	1435
$4(\pi^+\pi^-)\pi^0$		0.9 $\pm$ 0.3	1345	$K^*(892)^0 \bar{K}_2^*(1430)^0 + \text{c.c.}$	0.67 $\pm$ 0.26	1011
$\pi^+\pi^- K^+ K^-$		0.72 $\pm$ 0.23	1407	$b_1(1235)^\pm \pi^\mp$	0.29 $\pm$ 0.07	1298
$K\bar{K}\pi$		0.61 $\pm$ 0.10	1440	$\omega f_2(1270)$	0.23 $\pm$ 0.08	S=1.2* 1143
$\rho\bar{\rho}\pi^+\pi^-$		0.60 $\pm$ 0.05	S=1.3* 1107	$\phi\pi^+\pi^-$	0.21 $\pm$ 0.09	1365
$\Sigma\bar{\Sigma}$		0.45 $\pm$ 0.07	988	$\omega\eta$	0.19 $\pm$ 0.04	1394
$2(\pi^+\pi^-)$		0.4 $\pm$ 0.1	1517	$\Delta(K\bar{K})$	0.18 $\pm$ 0.08	1176
$3(\pi^+\pi^-)$		0.4 $\pm$ 0.2	1466	$\phi(1232)^+ + \bar{p}\pi^-$	0.16 $\pm$ 0.05	1030
$n\bar{n}\pi^+\pi^-$		0.38 $\pm$ 0.36	1106	$\omega K\bar{K}$	0.16 $\pm$ 0.10	1265
$2(\pi^+\pi^-)K^+ K^-$		0.31 $\pm$ 0.13	1320	$\omega\rho\bar{\rho}$	0.13 $\pm$ 0.03	S=1.3* 768
$\Xi\bar{\Xi}$		0.25 $\pm$ 0.04	818	$\Delta(1232)^+ + \bar{\Delta}(1232)^{-}$	0.11 $\pm$ 0.03	938
$\rho\bar{\rho}\pi^+\pi^-\pi^0$		0.23 $\pm$ 0.09 <sup>m</sup>	S=1.9* 1033	$\eta\rho\bar{\rho}$	0.09 $\pm$ 0.04	S=1.7* 596
$\rho\bar{\rho}$		0.22 $\pm$ 0.01	1232	$\Sigma(1385)^\pm \bar{\Sigma}(1385)^\mp$	0.09 $\pm$ 0.03	692
$\rho\bar{\rho}\eta$		0.21 $\pm$ 0.02	948	$\bar{p}N(1440-1535)^+$	0.09 $\pm$ 0.04	979
$\rho\bar{n}\pi^-$ or $\bar{\rho}n\pi^+$		0.20 $\pm$ 0.01	1174	$\phi\eta$	0.067 $\pm$ 0.008	1320
$n\bar{n}$		0.18 $\pm$ 0.09	1231	$\omega\pi$	0.067 $\pm$ 0.013	1446
$\Lambda\bar{\Sigma}^\pm \pi^\mp$		0.15 $\pm$ 0.04	1332	$\rho K^-\Sigma(1385)^0$	0.05 $\pm$ 0.03	643
$\Lambda\Lambda$		0.13 $\pm$ 0.02	S=1.7* 1074	$\phi\eta'$	0.037 $\pm$ 0.006	1192
$\rho\bar{\rho}\pi^0$		0.11 $\pm$ 0.01	1176	$\omega\eta'$	0.040 $\pm$ 0.011	1279
$\rho K^-\bar{\Lambda}$		0.09 $\pm$ 0.02	876	$\phi f_2'(1525)$	0.037 $\pm$ 0.013	871
$2(K^+ K^-)$		0.07 $\pm$ 0.03	1131	$\Sigma(1385)^\pm \bar{\Sigma}^\mp$	0.03 $\pm$ 0.01	1234
$\rho K^-\bar{\Sigma}^0$		0.03 $\pm$ 0.01	820	$\phi f_0(975)$	0.026 $\pm$ 0.006	1184
$K^+ K^-$		0.024 $\pm$ 0.003	1468	$\rho\eta$	0.018 $\pm$ 0.004	1398
$\pi^+\pi^-$		0.015 $\pm$ 0.002	1542	$\pi^\pm a_2^\mp(1320)$	< 0.43	1263
$K_L^0 K_S^0$		0.010 $\pm$ 0.002	1466	$K\bar{K}_2^*(1430)$	< 0.4	1159
$\Lambda\bar{\Sigma}$		< 0.015]	1032	$K_2^*(1430)^0 \bar{K}_2^*(1430)^0$	< 0.29	606
Radiative decay modes				Radiative decay modes (cont'd)		
$\dagger[\gamma\eta\pi\pi$		0.62 $\pm$ 0.11	1486	$\gamma\eta_c(2980)$	1.27 $\pm$ 0.36	114
$\gamma 2(\pi^+\pi^-)$		0.49 $\pm$ 0.13	1517	$\gamma f_2'(1720)$	seen	1073
$\gamma\eta(1440) \rightarrow \gamma K\bar{K}\pi$		0.46 $\pm$ 0.07 <sup>n</sup>	1214	$\gamma\rho\bar{\rho}$	0.038 $\pm$ 0.010	1232
$\gamma\eta'$		0.42 $\pm$ 0.05	1400	$\gamma f_2'(1525)$	0.016 $\pm$ 0.005	1173
$\gamma\omega\omega$		0.18 $\pm$ 0.05	1336	$\gamma\pi^0$	0.004 $\pm$ 0.001	1546
$\gamma f_2(1270)$		0.16 $\pm$ 0.02	1286	$\gamma f_1(1285)$	< 0.6	1283
$\gamma\rho^0\rho^0$		0.13 $\pm$ 0.05	1344	$\gamma\rho\bar{\rho}\pi^+\pi^-$	< 0.08	1107
$\gamma\eta$		0.086 $\pm$ 0.008	1500	$2\gamma$	< 0.05	1548
				$3\gamma$	< 0.006]	1548

## Meson Summary Table (*cont'd*)

Particle	$J^{PC}$ <sup>a</sup> —— <i>estab.</i>	Mass <i>M</i> (MeV)	Full width $\Gamma$ (MeV)	Partial decay modes						
				Mode	Fraction(%) [Upper limits (%) are 90% CL]	$p$ <sup>b</sup> (MeV/c)				
$\chi_0(3415)$ <i>was</i> $\chi(3415)$	$0^+(0^{++})$	3414.9 $\pm 1.1$		$2(\pi^+\pi^-)$ (incl. $\pi\pi\rho$ )	$3.8 \pm 0.8$	1679				
				$\pi^+\pi^-K^+K^-$ (incl. $\pi K\bar{K}^*$ )	$2.9 \pm 0.8$	1580				
				$3(\pi^+\pi^-)$	$1.5 \pm 0.5$	1633				
				$\pi^+\pi^-$	$0.8 \pm 0.2$	1702				
				$\gamma J/\psi(3097)$	$0.7 \pm 0.2$	S=1.3* 303				
				$K^+K^-$	$0.7 \pm 0.2$	1635				
				$p\bar{p}\pi^+\pi^-$	$0.5 \pm 0.2$	1320				
				$\gamma\gamma$	$< 0.15$	1707				
				$p\bar{p}$	$< 0.09$	1427				
$\chi_1(3510)$ <i>was</i> $\chi(3510)$	$0^+(1^{++})$	3510.7 $\pm 0.5$ S=1.4*	$< 1.0$	$\gamma J/\psi(3097)$	$25.8 \pm 2.5$	389				
				$3(\pi^+\pi^-)$	$2.2 \pm 0.9$	1683				
				$2(\pi^+\pi^-)$ (incl. $\pi\pi\rho$ )	$1.7 \pm 0.5$	1727				
				$\pi^+\pi^-K^+K^-$ (incl. $\pi K\bar{K}^*$ )	$0.9 \pm 0.4$	1632				
				$\pi^+\pi^-p\bar{p}$	$0.14 \pm 0.09$	1381				
				$\pi^+\pi^- + K^+K^-$	$< 0.2$	1749				
				$\gamma\gamma$	$< 0.15$	1755				
				$p\bar{p}$	$< 0.12$	1483				
				$\chi_2(3555)$ <i>was</i> $\chi(3555)$	$0^+(2^{++})$	3556.3 $\pm 0.4$	2.9 $+1.8$ $-1.1$	$\gamma J/\psi(3097)$	$14.8 \pm 1.7$	429
								$2(\pi^+\pi^-)$ (incl. $\pi\pi\rho$ )	$2.2 \pm 0.5$	1750
$\pi^+\pi^-K^+K^-$ (incl. $\pi K\bar{K}^*$ )	$1.9 \pm 0.5$	1656								
$3(\pi^+\pi^-)$	$1.1 \pm 0.8$	1706								
$p\bar{p}$	$0.76 \pm 0.30$	1510								
$\pi^+\pi^-p\bar{p}$	$0.33 \pm 0.13$	1410								
$\pi^+\pi^-$	$0.19 \pm 0.10$	1772								
$K^+K^-$	$0.15 \pm 0.11$	1708								
$J/\psi\pi^+\pi^-\pi^0$	$< 1.5$	185								
$\gamma\gamma$	$< 0.06$	1778								
$\psi(3685)$	$0^-(1^{--})$	3686.0 $\pm 0.1$	0.215 $\pm 0.040$	$e^+e^- + \mu^+\mu^-$	$1.8 \pm 0.3$	1843				
				hadrons + radiative	$98.2 \pm 0.3$					
$\Gamma_{ee} = (2.1 \pm 0.2)$ keV				Decay modes into hadrons						
(assuming $\Gamma_{ee} = \Gamma_{\mu\mu}$ )				$\dagger [J/\psi\pi\pi$	$50 \pm 4$	477				
				$J/\psi\eta$	$2.7 \pm 0.4$	S=1.6* 196				
				$3(\pi^+\pi^-)\pi^0$	$0.35 \pm 0.16$	1746				
				$2(\pi^+\pi^-)\pi^0$	$0.31 \pm 0.07$	1799				
				$\pi^+\pi^-K^+K^-$	$0.16 \pm 0.04$	1726				
				$J/\psi\pi^0$	$0.10 \pm 0.03$	528				
				$p\bar{p}\pi^+\pi^-$	$0.08 \pm 0.02$	1491				
				$K^*(892)^0K^-\pi^+ + c.c.$	$0.067 \pm 0.025$	1674				
				$2(\pi^+\pi^-)$	$0.05 \pm 0.01$	1817				
				$\rho^0\pi^+\pi^-$	$0.042 \pm 0.015$	1751				
				$p\bar{p}$	$0.019 \pm 0.005$	1586				
				$3(\pi^+\pi^-)$	$0.015 \pm 0.010$	1774				
				$p\bar{p}\pi^0$	$0.014 \pm 0.005$	1543				
				$K^+K^-$	$0.010 \pm 0.007$	1776				
				$\pi^+\pi^-\pi^0$	$0.009 \pm 0.005$	1830				
				$\pi^+\pi^-$	$0.008 \pm 0.005$	1838				
				$\Delta\Delta$	$< 0.04$	1467				
				$\rho\pi$	$< 0.008$	1760				
				$K^+K^-\pi^0$	$< 0.003$	1754				
				$K^\pm K^*(892)^\mp$	$< 0.002$	1698				
$\psi(3770)$	$(1^{--})$	3769.9 $\pm 2.4$ S=1.8*	25 $\pm 3$	$e^+e^-$	$0.0011 \pm 0.0002$	1885				
				$D\bar{D}$	dominant	242				
$\Gamma_{ee} = (0.26 \pm 0.15)$ keV S=1.3*										



## Meson Summary Table (*cont'd*)

Particle	$J^{PC}$ <sup>a</sup> — <i>estab.</i>	Mass <i>M</i> (MeV)	Full width $\Gamma$ (MeV)	Partial decay modes		
				Mode	Fraction(%) [Upper limits (%) are 90% CL]	$p^b$ (MeV/c)
$\psi(4030)$	$(1^{--})$	$4030^{\S}$ $\pm 5^{\S}$	$52$ $\pm 10$	$e^+e^-$ $D\bar{D}$ $D\bar{D}^*(2010) + c.c.$ $D^*(2010)\bar{D}^*(2010)$	$0.0014 \pm 0.0004$ seen seen seen	2015 752 559 177
$\Gamma_{ee} = (0.75 \pm 0.15)$ keV						
$\psi(4160)$	$(1^{--})$	$4159$ $\pm 20$	$78$ $\pm 20$	$e^+e^-$	$0.0010 \pm 0.0004$	2079
$\Gamma_{ee} = (0.77 \pm 0.23)$ keV						
$\psi(4415)$	$(1^{--})$	$4415$ $\pm 6$	$43$ $\pm 20^{\S}$	$e^+e^-$	$0.0011 \pm 0.0006$	2207
$\Gamma_{ee} = (0.47 \pm 0.10)$ keV						
<b><math>b\bar{b}</math> MESONS</b>						
$\Upsilon(9460)$ or $\Upsilon(1S)$	$(1^{--})$	$9460.0$ $\pm 0.2$ $S=1.5^*$	$0.043$ $\pm 0.003$	$\tau^+\tau^-$ $\mu^+\mu^-$ $e^+e^-$ $\rho\pi$ $J/\psi + \text{anything}$	$3.2 \pm 0.4$ $2.8 \pm 0.2$ $2.8 \pm 0.3$ $< 0.21$ $< 2$	4381 4729 4730 4698
$\Gamma_{ee} = (1.22 \pm 0.05)$ keV						
$\chi_{b0}(9860)$ or $\chi_{b0}(1P)^o$	$(^{++})$	$9859.8$ $\pm 1.3$		$\Upsilon(9460)\gamma$	$< 6$	392
$\chi_{b1}(9895)$ or $\chi_{b1}(1P)^o$	$(^{++})$	$9891.9$ $\pm 0.7$		$\Upsilon(9460)\gamma$	$35 \pm 8$	422
$\chi_{b2}(9915)$ or $\chi_{b2}(1P)^o$	$(^{++})$	$9913.3$ $\pm 0.6$		$\Upsilon(9460)\gamma$	$22 \pm 4$	443
$\Upsilon(10023)$ or $\Upsilon(2S)$	$(1^{--})$	$10023.4$ $\pm 0.3$	$0.030$ $\pm 0.007$	$\mu^+\mu^-$ $\tau^+\tau^-$ $\Upsilon(9460)\pi^+\pi^-$ $\Upsilon(9460)\pi^0\pi^0$ $\Upsilon(9460)\eta$ $\chi_{b0}(9860)\gamma$ $\chi_{b1}(9895)\gamma$ $\chi_{b2}(9915)\gamma$	$1.8 \pm 0.4$ $1.7 \pm 1.6$ $18.7 \pm 1.0$ $8.6 \pm 1.1$ $< 0.2$ $4.3 \pm 1.0$ $6.7 \pm 0.9$ $6.6 \pm 0.9$	5011 4683 476 481 124 162 131 109
$\Gamma_{ee} = (0.54 \pm 0.03)$ keV						
$\chi_{b1}(10255)$ or $\chi_{b1}(2P)^o$	$(^{++})$	$10255$ $\pm 2$		$\gamma\Upsilon(9460)$ $\gamma\Upsilon(10023)$	seen seen	763 228
Seen in one experiment only.						
$\chi_{b2}(10270)$ or $\chi_{b2}(2P)^o$	$(^{++})$	$10271$ $\pm 2$		$\gamma\Upsilon(9460)$ $\gamma\Upsilon(10023)$	seen seen	779 245
Seen in one experiment only.						
$\Upsilon(10355)$ or $\Upsilon(3S)$	$(1^{--})$	$10355.5$ $\pm 0.5$	$0.012$ $+0.010$ $-0.004$	$\mu^+\mu^-$ $\Upsilon(9460)\pi^+\pi^-$ $\Upsilon(10023)\pi^+\pi^-$ $\chi_{b1}(10255)\gamma$ $\chi_{b2}(10270)\gamma$	$3.3 \pm 1.5$ $4.5 \pm 0.8$ $3 \pm 2$ $15.6 \pm 4.2$ $12.7 \pm 4.1$	5177 814 177 101 84
$\Gamma_{ee} = (0.40 \pm 0.03)$ keV						
$\Upsilon(10575)$ or $\Upsilon(4S)$	$(1^{--})$	$10577$ $\pm 4$	$24$ $\pm 2$			
$\Gamma_{ee} = (0.24 \pm 0.05)$ keV						
$\Upsilon(10860)$ or $\Upsilon(5S)$	$(1^{--})$	$10865$ $\pm 8$	$110$ $\pm 13$			
$\Gamma_{ee} = (0.31 \pm 0.07)$ keV						
$\Upsilon(11020)$ or $\Upsilon(6S)$	$(1^{--})$	$11019$ $\pm 9$	$79$ $\pm 16$			
$\Gamma_{ee} = (0.13 \pm 0.03)$ keV						

## Meson Summary Table (cont'd)

Particle	$I (J^P)$ — <i>estab.</i>	Mass $M$ (MeV)	Full width $\Gamma$ (MeV)	Partial decay modes		
				Mode	Fraction(%) [Upper limits (%) are 90% CL]	$p^b$ (MeV/c)
<b>STRANGE MESONS</b>				[ $K^+ = u\bar{s}, K^0 = d\bar{s}, \bar{K}^0 = \bar{d}s, K^- = \bar{u}s$ ]		
$K^+$ $K^0$	$1/2(0^-)$	493.67 497.72		See Stable Particle Summary Table		
$K^*(892)$	$1/2(1^-)$	892.1 $\pm 0.3$ S=1.4*	51.1 $\pm 0.8$ S=1.1*	$K\pi$ $K^0\gamma$ $K^\pm\gamma$ $K\pi\pi$	$\approx 100$ $0.23 \pm 0.02$ $0.10 \pm 0.01$ < 0.05	288 309 307 216
$M$ and $\Gamma$ from charged mode: $M^0 = 896.5 \pm 0.4$ S=1.6*						
$K_1(1280)$ was $Q(1280)$	$1/2(1^+)$	1270 <sup>S</sup> $\pm 10^S$	90 <sup>S</sup> $\pm 20^S$	$K\rho$ $K_0^*(1350)\pi$ $K^*(892)\pi$ $K\omega$ $Kf_0(1300)$	42 $\pm$ 6 28 $\pm$ 4 16 $\pm$ 5 11 $\pm$ 2 3 $\pm$ 2	45 298
$K_0^*(1350)$ was $\kappa(1350)$	$1/2(0^+)$	$\sim 1350$	$\sim 250$	$K\pi$	seen	574
$K_1(1400)$ was $Q(1400)$	$1/2(1^+)$	1406 $\pm 10$	184 $\pm 9$	$K^*(892)\pi$ $K\rho$ $Kf_0(1300)$ $K\omega$	94 $\pm$ 6 3 $\pm$ 3 2 $\pm$ 2 1 $\pm$ 1	403 299 285
$K_2^*(1430)$ was $K^*(1430)$	$1/2(2^+)$	1426 $\pm 2$ S=1.1*	99 $\pm 3$	$K\pi$ $K^*(892)\pi$ $K^*(892)\pi\pi$ $K\rho$ $K\eta$ $K\omega$ $K\gamma$	44.9 $\pm$ 2.7 23.6 $\pm$ 1.8 13.5 $\pm$ 2.7 8.6 $\pm$ 0.8 5.2 $\pm$ 2.9 4.3 $\pm$ 1.6 0.24 $\pm$ 0.05	S=3.2* 618 417 366 324 485 310 627
$M$ and $\Gamma$ from charged mode: $M^0 = 1423 \pm 2$ S=1.1*						
$K_2(1770)^\ddagger$ was $L(1770)$	$1/2(2^-)$	$\sim 1770^S$	$\sim 200^S$	$K_2^*(1430)\pi$ $K^*(892)\pi$ $Kf_2(1270)$ $K\phi$	dominant seen seen seen	286 651 438
$K_3^*(1780)^\ddagger$ was $K^*(1780)$	$1/2(3^-)$	1780 $\pm 4$	150 $\pm 17$ S=1.4*	$K\pi\pi$ $^+ [K\rho$ $^+ [K^*(892)\pi$ $K\pi$	large large large 17 $\pm 5^S$	796 620 657 815
$K_4^*(2060)$ was $K^*(2060)$	$1/2(4^+)$	2060 <sup>S</sup> $\pm 30^S$	210 <sup>S</sup> $\pm 40^S$	$K\pi$ $K^*(892)\pi\pi$ $\rho K\pi$ $\omega K\pi$ $K^*(892)\pi\pi\pi$	7 $\pm$ 1 seen seen seen seen	966 809 751 744 775
Not a well-established resonance.						
<b>CHARMED, NONSTRANGE MESONS</b>				[ $D^+ = c\bar{d}, D^0 = c\bar{u}, \bar{D}^0 = \bar{c}u, D^- = \bar{c}d$ ]		
$D^+$ $D^0$	$1/2(0^-)$	1869.3 1864.6		See Stable Particle Summary Table		
$D^*(2010)^+$	$1/2(1^-)$	2010.1 $\pm 0.7$	< 2.0	$D^0\pi^+$ $D^+\pi^0$ $D^+\gamma$	49 $\pm$ 8 34 $\pm$ 7 17 $\pm$ 11	39 38 136
$m_{D^{*+}} - m_{D^0} = 145.45 \pm 0.07$ MeV						
$D^*(2010)^0$	$1/2(1^-)$	2007.2 $\pm 2.1$	< 5	$D^0\pi^0$ $D^0\gamma$	51.5 $\pm$ 7.6 48.5 $\pm$ 7.6	44 137

## Meson Summary Table (cont'd)

Particle	$I (J^P)$ — <i>estab.</i>	Mass $M$ (MeV)	Full width $\Gamma$ (MeV)	Partial decay modes	
				Mode	Fraction(%) [Upper limits (%) are 90% CL] $p^b$ (MeV/c)
<b>CHARMED, STRANGE MESON</b>				$[D_s^+ = c\bar{s}, D_s^- = \bar{c}s]$	
$D_s^+$ <i>was</i> $F^+$	$0(0^-)$	1971		See Stable Particle Summary Table	
<b>BOTTOM MESON</b>				$[B^+ = u\bar{b}, B^0 = d\bar{b}, \bar{B}^0 = \bar{d}b, B^- = \bar{u}b]$	
$B^+$ $B^0$	$1/2(0^-)$	5271 5275		See Stable Particle Summary Table	

‡ See Meson Full Listings.

\* Quoted error includes scale factor  $S = \sqrt{\chi^2/(N-1)}$ . See footnote to Stable Particle Summary Table.

† Square brackets indicate a subreaction of the previous (unbracketed) decay mode(s).

§ This is only an educated guess; the error given is larger than the error on the average of the published values. (See the Meson Full Listings for the latter.)

a. Charge conjugation  $C$  applies only to neutral states.

b. For a 2-body decay mode, this is the momentum of the decay products in the decay rest frame. For a 3-or-more-body mode, this is the maximum momentum any of the products can have in this frame. The momenta have been calculated by using the averaged central mass values, without taking into account the widths of the resonances.

c. From pole position ( $M - i\Gamma/2$ ).

d. The  $e^+e^-$  branching fraction is from  $e^+e^- \rightarrow \pi^+\pi^-$  experiments only. The  $\omega\rho$  interference is then due to  $\omega\rho$  mixing only, and is expected to be small. The  $\mu^+\mu^-$  branching fraction is compiled from 3 experiments, each possibly with substantial  $\omega\rho$  interference. The error reflects this uncertainty; see notes in the Meson Full Listings. If  $e\mu$  universality holds,  $\Gamma(\rho^0 \rightarrow \mu^+\mu^-) = \Gamma(\rho^0 \rightarrow e^+e^-) \times 0.99785$ .

e. Empirical limits on fractions for other decay modes of  $\rho(770)$  are  $\pi^\pm\eta < 0.8\%$  (CL=84%),  $\pi^+\pi^+\pi^-\pi^- < 0.15\%$ .  $\pi^\pm\pi^+\pi^-\pi^0 < 0.2\%$  (CL=84%).

f. Empirical limits on fractions for other decay modes of  $\omega(783)$  are  $\pi^+\pi^-\gamma < 5\%$ ,  $\pi^0\pi^0\gamma < 1\%$ ,  $\eta + \text{neutral(s)} < 1.5\%$ ,  $\mu^+\mu^- < 0.02\%$ .

g. Empirical limits on fractions for other decay modes of  $\eta(958)$  are  $\pi^+\pi^- < 2\%$  (CL=84%),  $\pi^+\pi^-\pi^0 < 5\%$  (CL=84%),  $\pi^+\pi^+\pi^-\pi^- < 1\%$  (CL=95%),  $\pi^+\pi^+\pi^-\pi^-\pi^0 < 1\%$  (CL=84%),  $6\pi < 1\%$ ,  $\pi^+\pi^-e^+e^- < 0.6\%$ ,  $\pi^0e^+e^- < 1.3\%$  (CL=84%),  $\eta e^+e^- < 1.1\%$ ,  $\pi^0\rho^0 < 4\%$ ,  $\eta\mu^+\mu^- < 1.5 \times 10^{-5}$ ,  $\pi^0\mu^+\mu^- < 6 \times 10^{-5}$ .

h. The mass and width are from the  $\eta\pi$  mode only. If the  $K\bar{K}$  channel is strongly coupled, the width may be larger.

i. Empirical limits on fractions for other decay modes of  $\phi(1020)$  are  $\pi^+\pi^-\gamma < 0.7\%$ ,  $\omega\gamma < 5\%$  (CL=84%),  $\rho\gamma < 2\%$  (CL=84%),  $2\pi^+2\pi^-\pi^0 < 1\%$  (CL=95%),  $2\pi^+2\pi^- < 0.1\%$ .

j. Empirical limits on fractions for other decay modes of  $b_1(1235)$  are  $\pi\pi < 15\%$ ,  $K\bar{K} < 2\%$  (CL=84%),  $4\pi < 50\%$  (CL=84%),  $\phi\pi < 1.5\%$  (CL=84%),  $\eta\pi < 25\%$ .  $(\bar{K}K)^\pm\pi^0 < 8\%$ ,  $K_S K_S \pi^\pm < 2\%$ ,  $K_S K_L \pi^\pm < 6\%$ .

k. Empirical limits (CL=95%) on fractions for other decay modes of  $f_2(1270)$  are  $\eta\pi\pi < 1\%$ ,  $K^0 K^- \pi^+ + \text{c.c.} < 0.4\%$ .

l. Empirical limits on fractions for other decay modes of  $\pi_2(1680)$  are  $\eta\pi < 10\%$ ,  $5\pi < 10\%$ .

m. Includes  $\rho\bar{\rho}\pi^+\pi^-\gamma$  and excludes  $\rho\bar{\rho}\eta$ ,  $\rho\bar{\rho}\omega$ ,  $\rho\bar{\rho}\eta'$ .

n. See  $f_1(1420)$  mini-review.

o. Spectroscopic labeling for these states is theoretical, pending experimental information.

# Meson Summary Table (cont'd)

## Table of Contents of Meson Full Listings

● Indicates particle appears in Meson Summary Table above. We do not regard the other entries as established resonances.

Nonstrange ( $S = 0; C, B = 0$ )			Strange ( $ S  = 1; C, B = 0$ )				
entry	$I^G(J^{PC})^a$	entry	$I^G(J^{PC})^a$	entry	$I(J^P)$		
● $\pi$	$1^-(0^{-+})$	● $\rho$	(1600) $1^+(1^{--})$	$e^+e^-$	(1100–2200) ( $1^{--}$ )	● $K$	$1/2(0^-)$
● $\eta$	$0^+(0^{-+})$	● $\omega_3$	(1670) $0^-(3^{--})$	$\bar{N}N$	(1200–3600)	● $K^*$	(892) $1/2(1^-)$
● $\rho$	(770) $1^+(1^{--})$	● $\pi_2$	(1680) $1^-(2^{-+})$	$X$	(1900–3600)	● $K_1$	(1280) $1/2(1^+)$
● $\omega$	(783) $0^-(1^{--})$	● $\phi$	(1680) $0^-(1^{--})$	● $\eta_c$	(2980) $0^+(0^{-+})$	● $K_0^*$	(1350) $1/2(0^+)$
● $\eta'$	(958) $0^+(0^{-+})$	● $\rho_3$	(1690) $1^+(3^{--})$	● $J/\psi$	(3097) $0^-(1^{--})$	● $K_1$	(1400) $1/2(1^+)$
● $f_0$	(975) $0^+(0^{++})$	$X$	(1700) $+$	● $\chi_0$	(3415) $0^+(0^{++})$	$K^*$	(1410) $1/2(1^-)$
● $a_0$	(980) $1^-(0^{++})$	● $f_2$	(1720) $0^+(2^{++})$	● $\chi_1$	(3510) $0^+(1^{++})$	● $K_2^*$	(1430) $1/2(2^+)$
● $\phi$	(1020) $0^-(1^{--})$	$f_0$	(1730) $0^+(0^{++})$	● $\chi_2$	(3555) $0^+(2^{++})$	$K$	(1460) $1/2(0^-)$
● $h_1$	(1190) $0^-(1^{+-})$	$\pi$	(1770) $1^-(0^{-+})$	$\eta_c$	(3590) ( $+$ )	$K_2$	(1580) $1/2(2^-)$
● $b_1$	(1235) $1^+(1^{+-})$	$f_2$	(1810) $0^+(2^{++})$	● $\psi$	(3685) $0^-(1^{--})$	● $K_2$	(1770) $1/2(2^-)$
$f_0$	(1240) $0^+(0^{++})$	● $\phi_J$	(1850) $0$	● $\psi$	(3770) ( $1^{--}$ )	● $K_3^*$	(1780) $1/2(3^-)$
$\rho$	(1250) $1^+(1^{--})$	$X$	(1935)	● $\psi$	(4030) ( $1^{--}$ )	$K^*$	(1790) $1/2(1^-)$
● $f_2$	(1270) $0^+(2^{++})$	● $f_4$	(2030) $0^+(4^{++})$	● $\psi$	(4160) ( $1^{--}$ )	$K$	(1830) $1/2(0^-)$
● $a_1$	(1270) $1^-(1^{+-})$	$a_4$	(2040) $1^-(4^{++})$	● $\psi$	(4415) ( $1^{--}$ )	● $K_4^*$	(2060) $1/2(4^+)$
$\eta$	(1275) $0^+(0^{-+})$	$a_3$	(2050) $1^-(3^{++})$	● $T$	(9460) ( $1^{--}$ )	$K_2$	(2250) $1/2(2^-)$
● $f_1$	(1285) $0^+(1^{++})$	$\pi_2$	(2100) $1^-(2^{-+})$	● $\chi_{b0}$	(9860) ( $^{++}$ )	$K_3$	(2320) $1/2(3^+)$
● $f_0$	(1300) $0^+(0^{++})$	$\rho$	(2150) $1^+(1^{--})$	● $\chi_{b1}$	(9895) ( $^{++}$ )	$K_4$	(2500) $1/2(4^-)$
● $\pi$	(1300) $1^-(0^{-+})$	$f_2$	(2150) $0^+(2^{++})$	● $\chi_{b2}$	(9915) ( $^{++}$ )	Charmed ( $ C  = 1$ )	
● $a_2$	(1320) $1^-(2^{++})$	$X$	(2220) $0$ ( $^{++}$ )	● $T$	(10023) ( $1^{--}$ )	● $D$	$1/2(0^-)$
$f_2$	(1410) $0^+(2^{++})$	$f_2$	(2240) $0^+(2^{++})$	$\chi_{b0}$	(10235) ( $^{++}$ )	● $D^*$	(2010) $1/2(1^-)$
● $f_1$	(1420) $0^+(1^{++})$	$\rho_3$	(2250) $1^+(3^{--})$	● $\chi_{b1}$	(10255) ( $^{++}$ )	$D^*$	(2420) $1/2( )$
● $\eta$	(1440) $0^+(0^{-+})$	$f_4$	(2300) $0^+(4^{++})$	● $\chi_{b2}$	(10270) ( $^{++}$ )	● $D_s$	$0$ ( $0^-$ )
● $f_2'$	(1525) $0^+(2^{++})$	$\rho_5$	(2350) $1^+(5^{--})$	● $T$	(10355) ( $1^{--}$ )	$D_s^*$	(2110)
$f_1$	(1530) $0^+(1^{++})$	$a_6$	(2450) $1^-(6^{++})$	● $T$	(10575) ( $1^{--}$ )	Bottom ( $ B  = 1$ )	
● $f_0$	(1590) $0^+(0^{++})$	$f_6$	(2510) $0^+(6^{++})$	● $T$	(10860) ( $1^{--}$ )	● $B$	
				● $T$	(11020) ( $1^{--}$ )	$B^*$	(5325)
						Exotics	

# Baryon Summary Table

April 1986

The first, short table gives the name, the quantum numbers (where known), and the status of every entry in the Baryon Full Listings. Only the baryons with 3- or 4-star status are included in the main Baryon Summary Table. Due to insufficient data or uncertain interpretation, the other entries in the short table are not established as baryons.

$N(939) P_{11}$ ****	$\Delta(1232) P_{33}$ ****	$Z_0(1780) P_{01}$ *	$\Sigma(1193) P_{11}$ ****	$\Xi(1318) P_{11}$ ****
$N(1440) P_{11}$ ****	$\Delta(1550) P_{31}$ *	$Z_0(1865) D_{03}$ *	$\Sigma(1385) P_{13}$ ****	$\Xi(1530) P_{13}$ ****
$N(1520) D_{13}$ ****	$\Delta(1600) P_{33}$ **	$Z_1(1725) P_{11}$ *	$\Sigma(1480)$ *	$\Xi(1630)$ *
$N(1535) S_{11}$ ****	$\Delta(1620) S_{31}$ ****	$Z_1(1900) P_{13}$ *	$\Sigma(1560)$ **	$\Xi(1680)$ **
$N(1540) P_{13}$ *	$\Delta(1700) D_{33}$ ****	$Z_1(2150)$ *	$\Sigma(1580) D_{13}$ **	$\Xi(1820) 13$ ***
$N(1650) S_{11}$ ****	$\Delta(1900) S_{31}$ ***	$Z_1(2500)$ *	$\Sigma(1620) S_{11}$ **	$\Xi(1940)$ **
$N(1675) D_{15}$ ****	$\Delta(1905) F_{35}$ ****		$\Sigma(1660) P_{11}$ ***	$\Xi(2030) 1$ ***
$N(1680) F_{15}$ ****	$\Delta(1910) P_{31}$ ****	$\Lambda(1116) P_{01}$ ****	$\Sigma(1670) D_{13}$ ****	$\Xi(2120)$ *
$N(1700) D_{13}$ ***	$\Delta(1920) P_{33}$ ***	$\Lambda(1405) S_{01}$ ****	$\Sigma(1690)$ **	$\Xi(2250)$ **
$N(1710) P_{11}$ ***	$\Delta(1930) D_{35}$ ***	$\Lambda(1520) D_{03}$ ****	$\Sigma(1750) S_{11}$ ***	$\Xi(2370) 1$ **
$N(1720) P_{13}$ ****	$\Delta(1940) D_{33}$ *	$\Lambda(1600) P_{01}$ ***	$\Sigma(1770) P_{11}$ *	$\Xi(2500)$ *
$N(1960) ?$ *	$\Delta(1950) F_{37}$ ****	$\Lambda(1670) S_{01}$ ****	$\Sigma(1775) D_{15}$ ****	
$N(1990) F_{17}$ **	$\Delta(2000) F_{35}$ **	$\Lambda(1690) D_{03}$ ****	$\Sigma(1840) P_{13}$ *	$\Omega(1672) P_{03}$ ****
$N(2000) F_{15}$ **	$\Delta(2150) S_{31}$ *	$\Lambda(1800) S_{01}$ ***	$\Sigma(1880) P_{11}$ **	
$N(2080) D_{13}$ **	$\Delta(2200) G_{37}$ *	$\Lambda(1800) P_{01}$ ***	$\Sigma(1915) F_{15}$ ****	$\Lambda_c(2281)$ ****
$N(2090) S_{11}$ *	$\Delta(2300) H_{39}$ **	$\Lambda(1820) F_{05}$ ****	$\Sigma(1940) D_{13}$ ***	$\Sigma_c(2450)$ **
$N(2100) P_{11}$ *	$\Delta(2350) D_{35}$ *	$\Lambda(1830) D_{05}$ ****	$\Sigma(2000) S_{11}$ *	$\Xi_c(2460)$ *
$N(2190) G_{17}$ ****	$\Delta(2390) F_{37}$ *	$\Lambda(1890) P_{03}$ ****	$\Sigma(2030) F_{17}$ ****	$\Omega_c(2740)$ *
$N(2200) D_{15}$ **	$\Delta(2400) G_{39}$ **	$\Lambda(2000)$ *	$\Sigma(2070) F_{15}$ *	$\Lambda_b(5500)$ *
$N(2220) H_{19}$ ****	$\Delta(2420) H_{31}$ ****	$\Lambda(2020) F_{07}$ *	$\Sigma(2080) P_{13}$ **	
$N(2250) G_{19}$ ****	$\Delta(2750) I_{313}$ **	$\Lambda(2100) G_{07}$ ****	$\Sigma(2100) G_{17}$ *	<u>Dibaryons</u>
$N(2600) I_{111}$ ***	$\Delta(2950) K_{315}$ **	$\Lambda(2110) F_{05}$ ***	$\Sigma(2250)$ ***	$NN(2170) 1D2$ **
$N(2700) K_{113}$ **	$\Delta(\sim 3000)$	$\Lambda(2325) D_{03}$ *	$\Sigma(2455)$ **	$NN(2250) 3F3$ **
$N(\sim 3000)$		$\Lambda(2350)$ ***	$\Sigma(2620)$ **	$NN(?)$ *
		$\Lambda(2585)$ **	$\Sigma(3000)$ *	$\Lambda N(2130) 3S1$ **
			$\Sigma(3170)$ *	$\Xi N(?)$ *

- \*\*\*\* Good, clear, and unmistakable.  
 \*\*\* Good, but in need of clarification or not absolutely certain.  
 \*\* Not established; needs confirmation.  
 \* Evidence weak; likely to disappear.

Particle <sup>a</sup>	$J^P$	$L_{2I,2J}$	$P_{beam}^b$ (GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass <sup>c</sup> $M$ (MeV)	Full <sup>d</sup> width $\Gamma$ (MeV)	Decay modes		
						Mode <sup>e</sup>	Fraction <sup>f</sup> (%)	$p^g$ (MeV/c)
<b><math>N</math> RESONANCES (<math>S=0, I=1/2</math>)</b>						$[N^+ = uud, N^0 = udd]$		
$p$	$1/2^+$			938.3		See Stable Particle Table		
$n$				939.6				
$N(1440)$	$1/2^+$	$P_{11}$	$P \approx 0.61$ $\sigma = 31.0$	1400 to 1480	120 to 350 (200)	$N\pi$ $N\pi\pi$ $\Delta\pi$ $N\rho$ $N(\pi\pi)_S$ $p\gamma$ $n\gamma$	50-70 30-50 10-20 10-15 5-20 0.08-0.10 0.01-0.06	397 342 143 † 342 414 413
$N(1520)$	$3/2^-$	$D_{13}$	$P = 0.74$ $\sigma = 23.5$	1510 to 1530	100 to 140 (125)	$N\pi$ $N\eta$ $N\pi\pi$ $\Delta\pi$ $N\rho$ $N(\pi\pi)_S$ $p\gamma$ $n\gamma$	50-60 $\sim 0.1$ 40-50 20-30 15-25 < 5 0.43-0.57 0.34-0.51	456 149 410 228 † 410 470 470

## Baryon Summary Table (cont'd)

Particle <sup>a</sup>	$J^P$	$L_{2I,2J}$	$P_{\text{beam}}^b$ (GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass <sup>c</sup> $M$ (MeV)	Full <sup>d</sup> width $\Gamma$ (MeV)	Decay modes								
						Mode <sup>e</sup>	Fraction <sup>f</sup> (%)	$p^g$ (MeV/c)						
<b>N(1535)</b>	$1/2^-$	$S_{11}$	$P = 0.76$ $\sigma = 22.5$	1520 to 1560	100 to 250 (150)	$N\pi$	35-50	467						
						$N\eta$	45-55	182						
						$N\pi\pi$	$\sim 10$	422						
						$\Delta\pi$	$< 5$	242						
						$N\rho$	$\sim 5$	†						
						$N(\pi\pi)_S$	$\sim 5$	422						
						$p\gamma$	0.1-0.2	481						
						$n\gamma$	0.15-0.35	480						
						<b>N(1650)</b>	$1/2^-$	$S_{11}$	$P = 0.96$ $\sigma = 16.4$	1620 to 1680	100 to 200 (150)	$N\pi$	55-65	547
												$N\eta$	$\sim 1.5$	346
$\Delta K$	$\sim 8$	161												
$N\pi\pi$	20-35	511												
$\Delta\pi$	$< 10$	344												
$N\rho$	5-30	†												
$N(\pi\pi)_S$	$< 15$	511												
$p\gamma$	0.04-0.16	558												
$n\gamma$	0-0.17	557												
<b>N(1675)</b>	$5/2^-$	$D_{15}$	$P = 1.01$ $\sigma = 15.4$	1660 to 1690	120 to 180 (155)							$N\pi$	35-40	563
						$N\eta$	$\sim 1$	374						
						$\Delta K$	$\sim 0.1$	209						
						$N\pi\pi$	60-65	529						
						$\Delta\pi$	55-60	364						
						$N\rho$	$< 10$	†						
						$N(\pi\pi)_S$	$< 5$	529						
						$p\gamma$	$\sim 0.01$	575						
						$n\gamma$	0.07-0.12	574						
						<b>N(1680)</b>	$5/2^+$	$F_{15}$	$P = 1.01$ $\sigma = 15.2$	1670 to 1690	110 to 140 (125)	$N\pi$	55-65	567
$N\eta$	$< 1$	379												
$\Delta K$	not seen	218												
$N\pi\pi$	35-45	532												
$\Delta\pi$	10-15	369												
$N\rho$	10-20	†												
$N(\pi\pi)_S$	15-20	532												
$p\gamma$	0.21-0.30	578												
$n\gamma$	0.02-0.05	577												
<b>N(1700)</b>	$3/2^-$	$D_{13}$	$P = 1.05$ $\sigma = 14.5$	1670 to 1730	70 to 120 (100)							$N\pi$	5-15	580
						$N\eta$	$\sim 4$	400						
						$\Delta K$	$\sim 0.2$	250						
						$N\pi\pi$	80-90	547						
						$\Delta\pi$	15-70	385						
						$N\rho$	$< 20$	†						
						$N(\pi\pi)_S$	$< 70$	547						
						$p\gamma$	$\sim 0.01$	591						
						<b>N(1710)</b>	$1/2^+$	$P_{11}$	$P = 1.07$ $\sigma = 14.2$	1680 to 1740	90 to 130 (110)	$N\pi$	10-20	587
												$N\eta$	$\sim 25$	410
$\Delta K$	$\sim 15$	264												
$\Sigma K$	2-10	138												
$N\pi\pi$	$< 50$	554												
$\Delta\pi$	10-20	393												
$N\rho$	5-35	48												
$N(\pi\pi)_S$	5-35	554												
<b>N(1720)</b>	$3/2^+$	$P_{13}$	$P = 1.09$ $\sigma = 13.9$	1690 to 1800	125 to 250 (200)							$N\pi$	10-20	594
												$N\eta$	$\sim 3.5$	420
						$\Delta K$	$\sim 5$	278						
						$\Sigma K$	2-5	162						
						$N\pi\pi$	$< 75$	561						
						$\Delta\pi$	$< 15$	401						
						$N\rho$	$< 75$	104						
						$N(\pi\pi)_S$	$< 20$	561						

## Baryon Summary Table (cont'd)

Particle <sup>a</sup>	$J^P$	$L_{2I,2J}$	$P_{\text{beam}}^b$ (GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass <sup>c</sup> $M$ (MeV)	Full <sup>d</sup> width $\Gamma$ (MeV)	Decay modes		
						Mode <sup>e</sup>	Fraction <sup>f</sup> (%)	$p$ <sup>g</sup> (MeV/c)
<b>N(2190)</b>	$7/2^-$	$G_{17}$	$P = 2.07$ $\sigma = 6.21$	2120 to 2230	200 to 500 (350)	$N\pi$	~14	888
						$N\eta$	~3	790
						$\Delta K$	~0.3	712
<b>N(2220)</b>	$9/2^+$	$H_{19}$	$P = 2.14$ $\sigma = 5.97$	2150 to 2300	300 to 500 (400)	$N\pi$	~18	905
						$N\eta$	~0.5	811
						$\Delta K$	~0.2	732
<b>N(2250)</b>	$9/2^-$	$G_{19}$	$P = 2.21$ $\sigma = 5.74$	2130 to 2270	200 to 500 (300)	$N\pi$	~10	923
						$N\eta$	~2	831
						$\Delta K$	~0.3	754
<b>N(2600)</b>	$11/2^-$	$I_{111}$	$P = 3.12$ $\sigma = 3.86$	2580 to 2700	>300 (400)	$N\pi$	~5	1126
<b><math>\Delta</math> RESONANCES (<math>S=0, I=3/2</math>)</b>						[ $\Delta^{++} = uuu, \Delta^+ = uud, \Delta^0 = udd, \Delta^- = ddd$ ]		
<b><math>\Delta(1232)</math></b>	$3/2^+$	$P_{33}$	$P = 0.30$ $\sigma = 94.8$	1230 to 1234	110 to 120 (115)	$N\pi$	99.4	227
						$N\gamma$	0.56-0.66	259
<b><math>\Delta(1620)</math></b>	$1/2^-$	$S_{31}$	$P = 0.91$ $\sigma = 17.7$	1600 to 1650	120 to 160 (140)	$N\pi$	25-35	526
						$N\pi\pi$	65-75	488
						$\Delta\pi$	60-70	318
						$N\rho$	10-20	†
						$N\gamma$	~0.03	538
<b><math>\Delta(1700)</math></b>	$3/2^-$	$D_{33}$	$P = 1.05$ $\sigma = 14.5$	1630 to 1740	190 to 300 (250)	$N\pi$	10-20	580
						$N\pi\pi$	80-90	547
						$\Delta\pi$	50-90	385
						$N\rho$	<35	†
						$N\gamma$	0.14-0.33	591
<b><math>\Delta(1900)</math></b>	$1/2^-$	$S_{31}$	$P = 1.44$ $\sigma = 9.71$	1850 to 2000	130 to 300 (150)	$N\pi$	5-15	710
						$\Sigma K$	~10	410
<b><math>\Delta(1905)</math></b>	$5/2^+$	$F_{35}$	$P = 1.45$ $\sigma = 9.62$	1890 to 1920	250 to 400 (300)	$N\pi$	5-15	713
						$\Sigma K$	<3	415
						$N\pi\pi$	<75	687
						$\Delta\pi$	~25	542
						$N\rho$	<50	421
$N\gamma$	0.01-0.05	721						
<b><math>\Delta(1910)</math></b>	$1/2^+$	$P_{31}$	$P = 1.46$ <sup>†</sup> $\sigma = 9.54$	1850 to 1950	200 to 330 (220)	$N\pi$	15-25	716
						$\Sigma K$	2-20	421
						$N\pi\pi$	<75	691
						$\Delta\pi$	small	545
						$N\rho$	small	426
						$N(1440)\pi$	large	393
<b><math>\Delta(1920)</math></b>	$3/2^+$	$P_{33}$	$P = 1.48$ $\sigma = 9.38$	1860 to 2160	190 to 300 (250)	$N\pi$	15-20	722
						$\Sigma K$	~5	431
<b><math>\Delta(1930)</math></b>	$5/2^-$	$D_{35}$	$P = 1.50$ $\sigma = 9.21$	1890 to 1960	150 to 350 (250)	$N\pi$	5-15	729
						$\Sigma K$	<10	441
						$N\pi\pi$	not seen	704
<b><math>\Delta(1950)</math></b>	$7/2^+$	$F_{37}$	$P = 1.54$ $\sigma = 8.91$	1910 to 1960	200 to 340 (240)	$N\pi$	35-45	741
						$\Sigma K$	<1	460
						$N\pi\pi$	<40	716
						$\Delta\pi$	~30	574
						$N\rho$	<10	469
						$N\gamma$	0.08-0.17	749
<b><math>\Delta(2420)</math></b>	$11/2^+$	$H_{311}$	$P = 2.64$ $\sigma = 4.68$	2380 to 2450	300 to 500 (300)	$N\pi$	5-15	1023

## Baryon Summary Table (*cont'd*)

Particle <sup>a</sup>	$J^P$	$L_{I,2J}$	$P$ beam (GeV/c)	$\sigma = 4\pi\lambda^2$ (mb)	Mass <sup>c</sup> $M$ (MeV)	Full <sup>d</sup> width $\Gamma$ (MeV)	Decay modes		
							Mode	Fraction <sup>f</sup> (%)	$p$ <sup>g</sup> (MeV/c)
<b><math>\Lambda</math> RESONANCES (<math>S=-1, I=0</math>)</b>							[ $\Lambda^0 = uds$ ]		
$\Lambda$	$1/2^+$				1115.6		See Stable Particle Table		
$\Lambda(1405)$	$1/2^-$	$S_{01}$		Below $\bar{K}N$ threshold	1405 $\pm 5^h$	$40 \pm 10^h$	$\Sigma\pi$	100	152
$\Lambda(1520)$	$3/2^-$	$D_{03}$		$P = 0.395$ $\sigma = 82.3$	1519.5 $\pm 1.0^h$	15.6 $\pm 1.0^h$	$N\bar{K}$ $\Sigma\pi$ $\Lambda\pi\pi$ $\Sigma\pi\pi$ $\Lambda\gamma$	45 $\pm$ 1 42 $\pm$ 1 10 $\pm$ 1 0.9 $\pm$ 0.1 0.8 $\pm$ 0.2	244 267 252 152 351
$\Lambda(1600)$	$1/2^+$	$P_{01}$		$P = 0.58$ $\sigma = 41.6$	1560 to 1700	50 to 250 (150)	$N\bar{K}$ $\Sigma\pi$	15-30 10-60	343 336
$\Lambda(1670)$	$1/2^-$	$S_{01}$		$P = 0.74$ $\sigma = 28.5$	1660 to 1680	25 to 50 (35)	$N\bar{K}$ $\Sigma\pi$ $\Lambda\eta$	15-25 20-60 15-35	414 393 64
$\Lambda(1690)$	$3/2^-$	$D_{03}$		$P = 0.78$ $\sigma = 26.1$	1685 to 1695	50 to 70 (60)	$N\bar{K}$ $\Sigma\pi$ $\Lambda\pi\pi$ $\Sigma\pi\pi$	20-30 20-40 $\sim$ 25 $\sim$ 20	433 409 415 350
$\Lambda(1800)$	$1/2^-$	$S_{01}$		$P = 1.01$ $\sigma = 17.5$	1720 to 1850	200 to 400 (300)	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$ $N\bar{K}^*(892)$	25-40 seen seen seen	528 493 345 †
$\Lambda(1800)$	$1/2^+$	$P_{01}$		$P = 1.01$ $\sigma = 17.5$	1750 to 1850	50 to 250 (150)	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$ $N\bar{K}^*(892)$	20-50 10-40 seen 30-60	528 493 345 †
$\Lambda(1820)$	$5/2^+$	$F_{05}$		$P = 1.06$ $\sigma = 16.5$	1815 to 1825	70 to 90 (80)	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$	55-65 8-14 5-10	545 508 362
$\Lambda(1830)$	$5/2^-$	$D_{05}$		$P = 1.08$ $\sigma = 16.0$	1810 to 1830	60 to 110 (95)	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$	3-10 35-75 $>$ 15	553 515 371
$\Lambda(1890)$	$3/2^+$	$P_{03}$		$P = 1.21$ $\sigma = 13.6$	1850 to 1910	60 to 200 (100)	$N\bar{K}$ $\Sigma\pi$ $\Sigma(1385)\pi$ $N\bar{K}^*(892)$	20-35 3-10 seen seen	599 559 420 233
$\Lambda(2100)$	$7/2^-$	$G_{07}$		$P = 1.68$ $\sigma = 8.68$	2090 to 2110	100 to 250 (200)	$N\bar{K}$ $\Sigma\pi$ $\Lambda\eta$ $\Xi K$ $\Lambda\omega$ $N\bar{K}^*(892)$	25-35 $\sim$ 5 $<$ 3 $<$ 3 $<$ 8 10-20	751 704 617 483 443 514
$\Lambda(2110)$	$5/2^+$	$F_{05}$		$P = 1.70$ $\sigma = 8.53$	2090 to 2140	150 to 250 (200)	$N\bar{K}$ $\Sigma\pi$ $\Lambda\omega$ $\Sigma(1385)\pi$ $N\bar{K}^*(892)$	5-25 10-40 seen seen 10-60	757 711 455 589 524
$\Lambda(2350)$	$9/2^+$			$P = 2.29$ $\sigma = 5.85$	2340 to 2370	100 to 250 (150)	$N\bar{K}$ $\Sigma\pi$	$\sim$ 12 $\sim$ 10	915 867



## Baryon Summary Table (cont'd)

Particle <sup>a</sup>	$J^P$	$L_{I,2J}$	$P_{\text{beam}}^b$ (GeV/c) $\sigma = 4\pi\lambda^2$ (mb)	Mass <sup>c</sup> $M$ (MeV)	Full <sup>d</sup> width $\Gamma$ (MeV)	Decay modes		
						Mode	Fraction <sup>f</sup> (%)	$p^g$ (MeV/c)
<b><math>\Sigma</math> RESONANCES (<math>S=-1, I=1</math>)</b>						[ $\Sigma^+ = uus, \Sigma^0 = uds, \Sigma^- = dds$ ]		
$\Sigma^+$	$1/2^+$			1189.4		See Stable Particle Table		
$\Sigma^0$				1192.5				
$\Sigma^-$				1197.3				
$\Sigma(1385)^+$	$3/2^+$	$P_{13}$	Below $\bar{K}N$ threshold	$1382.8 \pm 0.4$ $S=2.0^i$	$36 \pm 1$	$\Delta\pi$ $\Sigma\pi$	$88 \pm 2$ $12 \pm 2$	208 127
$\Sigma(1385)^0$				$1383.7 \pm 1.0$ $S=1.4^i$	$36 \pm 5$			
$\Sigma(1385)^-$				$1387.2 \pm 0.6$ $S=2.2^i$	$39 \pm 2$ $S=1.7^i$			
$\Sigma(1660)$	$1/2^+$	$P_{11}$	$P = 0.72$ $\sigma = 29.9$	1630 to 1690	40 to 200 (100)	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$	10-30 seen seen	405 439 385
$\Sigma(1670)$	$3/2^-$	$D_{13}$	$P = 0.74$ $\sigma = 28.5$	1665 to 1685	40 to 80 (60)	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$	7-13 5-15 30-60	414 447 393
$\Sigma(1750)$	$1/2^-$	$S_{11}$	$P = 0.91$ $\sigma = 20.7$	1730 to 1800	60 to 160 (90)	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ $\Sigma\eta$	10-40 seen < 8 15-55	486 507 455 81
$\Sigma(1775)$	$5/2^-$	$D_{15}$	$P = 0.96$ $\sigma = 19.0$	1770 to 1780	105 to 135 (120)	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ $\Sigma(1385)\pi$ $\Delta(1520)\pi$	37-43 14-20 2-5 8-12 17-23	508 525 474 324 198
$\Sigma(1915)$	$5/2^+$	$F_{15}$	$P = 1.26$ $\sigma = 12.8$	1900 to 1935	80 to 160 (120)	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ $\Sigma(1385)\pi$	5-15 seen seen < 5	618 622 577 440
$\Sigma(1940)$	$3/2^-$	$D_{13}$	$P = 1.32$ $\sigma = 12.1$	1900 to 1950	150 to 300 (220)	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ $\Sigma(1385)\pi$ $\Delta(1520)\pi$ $\Delta(1232)\bar{K}$ $N\bar{K}^*(892)$	<20 seen seen seen seen seen seen	637 639 594 460 354 410 320
$\Sigma(2030)$	$7/2^+$	$F_{17}$	$P = 1.52$ $\sigma = 9.93$	2025 to 2040	150 to 200 (180)	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$ $\Xi K$ $\Sigma(1385)\pi$ $\Delta(1520)\pi$ $\Delta(1232)\bar{K}$ $N\bar{K}^*(892)$	17-23 17-23 5-10 < 2 5-15 10-20 10-20 < 5	702 700 657 412 529 430 498 438
$\Sigma(2250)$	?		$P = 2.04$ $\sigma = 6.76$	2210 to 2280	60 to 150 (100)	$N\bar{K}$ $\Delta\pi$ $\Sigma\pi$	<10 seen seen	851 842 803

## Baryon Summary Table (*cont'd*)

Particle <sup>a</sup>	$J^P$	$L_{2I-2J}$	Mass <sup>c</sup> $M$ (MeV)	Full <sup>d</sup> width $\Gamma$ (MeV)	Decay modes		
					Mode	Fraction (%)	$p^g$ (MeV/c)
<b><math>\Xi</math> RESONANCES (<math>S=-2, I=1/2</math>)</b>					$[\Xi^0 = uss, \Xi^- = dss]$		
$\Xi^0$	$1/2^+$		1314.9		See Stable Particle Table		
$\Xi^-$			1321.3				
$\Xi(1530)^0$	$3/2^+$	$P_{13}$	$1531.8 \pm 0.3$ $S = 1.3^i$	$9.1 \pm 0.5$	$\Xi\pi$	100	148
$\Xi(1530)^-$			$1535.0 \pm 0.6$	$10.1 \pm 1.9$			
$\Xi(1820)$	$3/2$		$1822 \pm 6^h$	$26^{+15}_{-10}{}^h$	$\Lambda\bar{K}$ $\Sigma\bar{K}$ $\Xi\pi$ $\Xi(1530)\pi$	$\sim 45$ $\sim 10$ small $\sim 45$	396 306 413 231
$\Xi(2030)$	?		$2025 \pm 6^h$	$20^{+15}_{-5}{}^h$	$\Lambda\bar{K}$ $\Sigma\bar{K}$ $\Xi\pi$ $\Xi(1530)\pi$	$\sim 20$ $\sim 80$ small small	587 524 573 418
<b>OTHER BARYONS</b>					$[\Omega^- = sss, \Lambda_c^+ = udc]$		
$\Omega^-$	$3/2^+$		1672.4		See Stable Particle Table		
$\Lambda_c^+$	$1/2^+$		2281		See Stable Particle Table		

Only the established baryons are included in this Baryon Summary Table. See the short table at the front of this main Table for a list of *all* the baryons for which there is evidence. See also the Notes on  $N$  and  $\Delta$  resonances, on  $\Lambda$  and  $\Sigma$  resonances, on  $\Xi$  resonances, and on dibaryons, introducing those sections of the Baryon Full Listings. In particular, there are Argand diagrams of all the  $\pi N$  and  $\bar{K}N$  elastic partial-wave amplitudes, and lengthy discussions of the main analyses of elastic and inelastic channels.

- f. This mode is energetically forbidden when the nominal mass of the decaying resonance (and of any resonance in the final state) is used, but is in fact allowed due to the nonzero width(s) of the resonance(s).
- a. The nominal mass here (in MeV) is used for identification; see column 5 for the actual mass.
- b. The quantities here are calculated using the nominal mass of column 1.
- c. Usually a conservatively large range of masses rather than a statistical average of the various determinations of the mass is given. In these cases, the mass determinations are nearly entirely from various phase-shift analyses of more or less the same data. It is thus not appropriate to treat the determinations as independent measurements or to average them together. The masses, widths, and branching fractions in this Table are Breit-Wigner parameters. The Baryon Full Listings also include pole parameters when they are available, and there is a table of pole parameters for  $N$  and  $\Delta$  resonances in the "Note on  $N$  and  $\Delta$  Resonances" in the Listings.
- d. Usually a conservatively large range of widths rather than a statistical average of the various determinations of the width is given (see note c for the reason). The nominal value in parentheses is then simply a best guess.
- e. The indented modes are subreactions of the  $N\pi\pi$  mode. The  $(\pi\pi)_S$  is the isospin-0,  $S$ -wave state of two pions (this used to be called the  $\epsilon$ ).
- f. Most of the inelastic branching fractions come from partial-wave analyses, and these determine  $(xx')^{1/2}$ , where  $x$  and  $x'$  are the elastic and inelastic branching fractions, not  $x'$  directly. Thus any uncertainty (and it is often considerable) in  $x$  carries over into  $x'$ . When  $x'$  so determined is really poorly known, we here simply note that the mode is seen. The values of  $(xx')^{1/2}$  are given in the Baryon Full Listings.
- g. For a 2-body decay mode, this is the momentum of the decay products in the decay rest frame. For a 3-or-more-body mode, this is the maximum momentum any of the products can have in this frame. The nominal mass of column 1 is used, as is the nominal mass of any resonance in the final state.
- h. The error given here is only an educated guess. It is larger than the error on the weighted average of the published values (the error on the weighted average is given in the Baryon Full Listings).
- i. The error given here has been scaled up by an "S factor" (see the \* footnote to the Stable Particle Summary Table for how S is defined) because the various measurements disagree more seriously than one would expect from statistics.

## PHYSICAL CONSTANTS\*

Quantity	Symbol, equation	Value	Uncert. (ppm)
speed of light	$c$	$2.997\,924\,58(1.2)\times 10^8\text{ m s}^{-1}$ (see note**)	0.004
Planck constant	$h$	$6.626\,176(36)\times 10^{-34}\text{ J s}$	5.4
Planck constant, reduced	$\hbar \equiv h/2\pi$	$1.054\,588\,7(57)\times 10^{-34}\text{ J s}$ $= 6.582\,173(17)\times 10^{-22}\text{ MeV s}$	5.4 2.6
electron charge magnitude	$e$	$1.602\,189\,2(46)\times 10^{-19}\text{ C} = 4.803\,242(14)\times 10^{-10}\text{ esu}$	2.9, 2.9
conversion constant	$\hbar c$	$197.328\,58(51)\text{ MeV fm}$	2.6
conversion constant	$(\hbar c)^2$	$0.389\,385\,7(20)\text{ GeV}^2\text{ mbarn}$	5.2
electron mass	$m_e$	$0.511\,003\,4(14)\text{ MeV}/c^2 = 9.109\,534(47)\times 10^{-31}\text{ kg}$	2.8, 5.1
proton mass	$m_p$	$938.279\,6(27)\text{ MeV}/c^2 = 1.672\,648\,5(86)\times 10^{-27}\text{ kg}$ $= 1.007\,276\,470(11)\text{ amu} = 1836.151\,52(70)\text{ }m_e^\dagger$	2.8, 5.1 0.011, 0.38
deuteron mass	$m_d$	$1875.628\,0(53)\text{ MeV}/c^2$	2.8
atomic mass unit (amu)	$(\text{mass }^{12}\text{C atom})/12 = (1\text{ g})/N_A$	$931.501\,6(26)\text{ MeV}/c^2 = 1.660\,565\,5(86)\times 10^{-27}\text{ kg}$	2.8, 5.1
permittivity of free space	$\epsilon_0$	$8.854\,187\,818(71)\times 10^{-12}\text{ C}^2\text{ N}^{-1}\text{ m}^{-2}$	0.008
permeability of free space	$\mu_0$	$4\pi\times 10^{-7} = 1.256\,637\,061\,4\times 10^{-6}\text{ N A}^{-2}$	—
		$\epsilon_0\mu_0 = 1/c^2$	
fine structure constant	$\alpha = e^2/4\pi\epsilon_0\hbar c$	$1/137.036\,04(11)^\dagger$	0.82
classical electron radius	$r_e = e^2/4\pi\epsilon_0 m_e c^2$	$2.817\,938\,0(70)\times 10^{-15}\text{ m}$	2.5
electron Compton wavelength	$\lambda_e = \hbar/m_e c = r_e\alpha^{-1}$	$3.861\,590\,5(64)\times 10^{-13}\text{ m}$	1.6
Bohr radius ( $m_{\text{nucleus}} = \infty$ )	$a_\infty = 4\pi\epsilon_0\hbar^2/m_e e^2 = r_e\alpha^{-2}$	$5.291\,770\,6(44)\times 10^{-11}\text{ m}$	0.82
Rydberg energy	$\hbar c R_\infty = m_e e^4/2(4\pi\epsilon_0)^2\hbar^2 = m_e c^2\alpha^2/2$	$13.605\,804(36)\text{ eV}$	2.6
Thomson cross section	$\sigma_T = 8\pi r_e^2/3$	$0.665\,244\,8(33)\text{ barn}$	4.9
Bohr magneton	$\mu_B = e\hbar/2m_e$	$5.788\,378\,5(95)\times 10^{-11}\text{ MeV T}^{-1}$	1.6
nuclear magneton	$\mu_N = e\hbar/2m_p$	$3.152\,451\,5(53)\times 10^{-14}\text{ MeV T}^{-1}$	1.7
electron cyclotron freq./field	$\omega_{\text{cycl}}^e/B = e/m_e$	$1.758\,804\,7(49)\times 10^{11}\text{ rad s}^{-1}\text{ T}^{-1}$	2.8
proton cyclotron freq./field	$\omega_{\text{cycl}}^p/B = e/m_p$	$9.578\,756(28)\times 10^7\text{ rad s}^{-1}\text{ T}^{-1}$	2.8
gravitational constant	$G_N$	$6.672\,0(41)\times 10^{-11}\text{ m}^3\text{ kg}^{-1}\text{ s}^{-2}$ $= 6.7065(41)\times 10^{-39}\text{ } \hbar c (\text{GeV}/c^2)^{-2}$	615 615
grav. accel., sea level, 45° lat.	$g$	$9.8062\text{ m s}^{-2}$	—
Fermi coupling constant	$G_F/(\hbar c)^3$	$1.166\,37(2)\times 10^{-5}\text{ GeV}^{-2}$	17
Avogadro number	$N_A$	$6.022\,045(31)\times 10^{23}\text{ mol}^{-1}\dagger$	5.1
Boltzmann constant	$k$	$1.380\,662(44)\times 10^{-23}\text{ J K}^{-1}$ $= 8.617\,35(28)\times 10^{-5}\text{ eV K}^{-1}$	32 32
molar volume, ideal gas at STP	$N_A k(273.15\text{ K})/(1\text{ atmosphere})$	$22.413\,83(70)\times 10^{-3}\text{ m}^3\text{ mol}^{-1}$	31
Stefan-Boltzmann constant	$\sigma = \pi^2 k^4/60\hbar^3 c^2$	$5.670\,32(71)\times 10^{-8}\text{ W m}^{-2}\text{ K}^{-4}$	125

$$\pi = 3.141\,592\,653\,589\,793\,238$$

$$e = 2.718\,281\,828\,459\,045\,235$$

$$\gamma = 0.577\,215\,664\,901\,532\,861$$

1 in $\equiv 0.0254\text{ m}$	1 barn $\equiv 10^{-28}\text{ m}^2$	1 eV $= 1.602\,189\,2\times 10^{-19}\text{ J}$	1 gauss (G) $\equiv 10^{-4}\text{ tesla (T)}$
1 Å $\equiv 10^{-10}\text{ m}$	1 dyne $\equiv 10^{-5}\text{ newton (N)}$	1 eV/c <sup>2</sup> $= 1.782\,676\times 10^{-36}\text{ kg}$	1 atmosphere $= 1.013\,25\times 10^5\text{ N/m}^2$
1 fm $\equiv 10^{-15}\text{ m}$	1 erg $\equiv 10^{-7}\text{ joule (J)}$	$2.997\,924\,58\times 10^9\text{ esu} = 1\text{ coulomb (C)}$	0° C $\equiv 273.15\text{ K}$

SI units take as their base: length (m), mass (kg), time (s), electric current (A), thermodynamic temperature (K), amount of a substance (mol), and luminous intensity (candela, cd), and the two supplementary units plane angle (rad) and solid angle (sr).

\* Revised 1985 by B.N. Taylor. Based mainly on the "1973 Least-Squares Adjustment of the Fundamental Constants," by E.R. Cohen and B.N. Taylor, J. Phys. Chem. Ref. Data 2, 663 (1973). The figures in parentheses after the values give the 1-standard-deviation uncertainties in the last digits; the uncertainties in parts per million (ppm) are given in the last column. The uncertainties of the values from a least-squares adjustment are in general correlated, and the laws of error propagation must be used in calculating additional quantities. The set of constants resulting from the 1973 adjustment has been recommended for international use by CODATA (Committee on Data for Science and Technology), and is the most up-to-date, generally accepted set currently available.

† Since 1973, new experiments have yielded better values for some of the constants:  $N_A = 6.022\,097\,8(63)\times 10^{23}\text{ mol}^{-1}$  (1.04 ppm);  $\alpha^{-1} = 137.035\,963(15)$  (0.11 ppm); and  $m_p/m_e = 1836.152\,701(100)$  (0.054 ppm – error increased by us pending final determination). However, since a change in the value of one constant usually leads to changes in the adjusted values of others, caution is required in using together the values from the 1973 adjustment and the results of more recent experiments.

\*\* In 1983, the Conf. Générale des Poids et Mesures adopted a new definition of the meter: it is the distance traveled by light in vacuum in  $1/299\,792\,458\text{ s}$ . Thus the speed of light is defined to be  $299\,792\,458\text{ m/s}$ . For a discussion, see B.W. Petley, Nature 303, 373 (1983).

## ASTROPHYSICAL CONSTANTS\*

Quantity	Symbol, equation	Value	Quantity	Symbol	Value
Planck mass	$M_{\text{Planck}}$ $= (\hbar c/G_N)^{1/2}$	$1.221\ 10(37)\times 10^{19}\ \text{GeV}/c^2$ $= 2.176\ 83(66)\times 10^{-8}\ \text{kg}$	cosmological constant	$\Lambda$	$ \Lambda  < 3\times 10^{-52}\ \text{m}^{-2}$
Hubble parameter <sup>1</sup>	$H_0$	$100h_0\ \text{km s}^{-1}\ \text{Mpc}^{-1}$ $= h_0\times 1.0\times 10^{-10}\ \text{year}^{-1}$	age of the universe <sup>1</sup>	$t_0$	$1.5(5)\times 10^{10}\ \text{years}$
normalized Hubble parameter <sup>1</sup>	$h_0$	$0.4 < h_0 < 1$	solar mass	$M_\odot$	$1.989(2)\times 10^{30}\ \text{kg}$
density parameter of the universe <sup>1</sup>	$\Omega_0 \equiv \rho_0/\rho_c$	$0.05 \leq \Omega_0 \leq 4$	solar luminosity	$L_\odot$	$3.826(8)\times 10^{26}\ \text{J s}^{-1}$
critical density of the universe <sup>1</sup>	$\rho_c = 3H_0^2/8\pi G_N$	$1.88\times 10^{-26}\ h_0^2\ \text{kg m}^{-3}$ $= 2.8\times 10^{11}\ h_0^2\ M_\odot\ \text{Mpc}^{-3}$	solar radius	$R_\odot$	$6.959\ 9(7)\times 10^8\ \text{m}$
			1 tropical year $\approx 3.155\ 69\times 10^7\ \text{s}$ 1 light year $= 9.460\ 528\times 10^{15}\ \text{m}$ 1 parsec (pc) $= 3.261\ 633\ \text{light years}$ 1 astro. unit $= 1.495\ 979\times 10^{11}\ \text{m}$		

\* Compiled with the help of K.A. Olive, J. Primack, and S. Rudaz. Some values are taken from C.W. Allen, *Astrophysical Quantities* (Athlone Press, London, 1973).

<sup>1</sup> Subscript 0 indicates present-day values.

## BIG BANG COSMOLOGY\*

All observational evidence to date indicates that our universe is very nearly homogeneous and isotropic. The most general space-time interval with these properties is the Friedmann-Robertson-Walker metric (with  $c = 1$ ):

$$ds^2 = dt^2 - R^2(t) \left[ \frac{dr^2}{1 - \kappa r^2} + r^2(d\theta^2 + \sin^2\theta d\phi^2) \right],$$

where  $\kappa = +1, -1, \text{ or } 0$  corresponds to closed, open, or spatially flat geometries;  $R(t)$  is a scale factor for distances in comoving coordinates. Einstein's equations lead to the Friedmann equation

$$H^2 \equiv \left( \frac{\dot{R}}{R} \right)^2 = \frac{8\pi G_N \rho}{3} - \frac{\kappa}{R^2} + \frac{\Lambda}{3},$$

as well as to

$$\frac{\ddot{R}}{R} = \frac{\Lambda}{3} - \frac{4\pi G_N}{3}(\rho + 3p),$$

where  $H(t)$  is the Hubble parameter,  $\rho$  is the total mass-energy density,  $p$  is the isotropic pressure, and  $\Lambda$  is the cosmological constant. (For limits on  $\Lambda$ , see the Table of Astrophysical Constants; we will assume here  $\Lambda = 0$ .) The Friedmann equation serves to define the density parameter  $\Omega_0$  (subscript 0 indicates present-day values):

$$\kappa/R_0^2 = H_0^2(\Omega_0 - 1), \quad \Omega_0 = \rho_0/\rho_c;$$

and the critical density is defined as

$$\rho_c \equiv \frac{3H_0^2}{8\pi G_N} = 1.88\times 10^{-26}\ h_0^2\ \text{kg m}^{-3},$$

with

$$H_0 = 100h_0\ \text{km s}^{-1}\ \text{Mpc}^{-1}.$$

Observational bounds give  $0.4 < h_0 < 1$ . The three possible values of  $\kappa$ ,  $+1, -1, \text{ and } 0$ , correspond to  $\Omega_0 > 1, < 1, \text{ and } = 1$ , i.e., to closed, open, and flat (critical) universes. The value of  $\Omega_0$  is inferred from velocity measurements on scales greater than 100 kpc, which are all consistent with  $0.1 \leq \Omega_0 \leq 0.4$ . Conservative bounds are  $0.05 \leq \Omega_0 \leq 4$ . The portion of  $\Omega$  in luminous matter is much smaller,  $0.005 \leq \Omega_{\text{lum}} \leq 0.02$ . The excess of  $\Omega_0$  over  $\Omega_{\text{lum}}$  leads to the inference that most of the matter in the universe is nonluminous "dark" matter.

Energy conservation implies that  $\dot{\rho} = -3(\dot{R}/R)(\rho + p)$ , so that for a matter-dominated ( $p = 0$ ) universe  $\rho \propto R^{-3}$ , while for a radiation-dominated ( $p = 1/3\rho$ ) universe  $\rho \propto R^{-4}$ . Thus the less singular curvature term  $\kappa/R^2$  in the Friedmann equation can be neglected at early times when  $R$  is small. Energy conservation also implies that the universe expands adiabatically,  $R^3 s = \text{constant}$ , where the entropy density  $s = (\rho + p)/T$  and  $T$  is temperature.

The energy density of radiation can be expressed as

$$\rho_r = \frac{\pi^2 k^4}{30} N(T) T^4,$$

with  $\hbar = 1$ , where  $N(T)$  counts the effectively massless degrees of freedom of bosons and fermions:

$$N(T) = \sum_B g_B + \frac{7}{8} \sum_F g_F.$$

For example, for  $m_\mu > kT > m_e$ ,  $N(T) = g_\gamma + 7/8(g_e + 3g_\nu) = 2 + 7/8[4 + 3(2)] = 43/4$ . For  $m_\pi > kT > m_\mu$ ,  $N(T) = 57/4$ .

In the early universe when  $\rho \sim \rho_r$ , then  $\dot{R} \sim 1/R$ , so that  $R \propto t^{1/2}$  and  $Ht \rightarrow 1/2$ ; the time-temperature relation then follows:

$$t = 2.4 [N(T)]^{-1/2} \left( \frac{1\ \text{MeV}}{kT} \right)^2 \text{ s}.$$

Today, the energy density in photons is  $\rho_\gamma = (\pi^2 k^4/15) T_0^4$ , where the present temperature of the microwave background is  $T_0 = 2.73 \pm 0.05\ \text{K}$ , and the number density of photons  $n_\gamma$  is  $400(T_0/2.7\ \text{K})^3\ \text{cm}^{-3}$ . For nonrelativistic matter (such as baryons) today, the energy density is  $\rho_B = m_B n_B$  with  $n_B \propto R^{-3}$ , so that for most of the history of the universe  $n_B/s$  is constant. Today, the entropy density is related to the photon density by  $s \approx 7n_\gamma$ . Big Bang nucleosynthesis calculations limit  $\eta = n_B/n_\gamma$  to  $3 \times 10^{-10} \leq \eta \leq 10^{-9}$ . The parameter  $\eta$  is also related to the portion of  $\Omega$  in baryons

$$\Omega_B = 3.6 \times 10^7 \eta h_0^{-2} (T_0/2.7\ \text{K})^3,$$

so that  $0.01 < \Omega_B h_0^2 < 0.04$  and hence the universe cannot be closed by baryons.

\* Written December 1985 by K.A. Olive and S. Rudaz.

## ATOMIC AND NUCLEAR PROPERTIES OF MATERIALS\*

Material	Z	A	Nuclear total cross section $\sigma_T$ [barn]	Nuclear <sup>b</sup> inelastic cross section $\sigma_I$ [barn]	Nuclear <sup>c</sup> collision length $\lambda_T$ [g/cm <sup>2</sup> ]	Nuclear <sup>c</sup> interaction length $\lambda_I$ [g/cm <sup>2</sup> ]	$dE/dx$ min <sup>d</sup>		Radiation length <sup>e</sup>		Density <sup>f</sup> [g/cm <sup>3</sup> ] ( ) is for gas [g/ℓ]	Refractive index $n^f$ ( ) is $(n-1) \times 10^6$ for gas
							$\frac{dE}{d\ell}$ [MeV/g/cm <sup>2</sup> ]	$\Delta E_{mp}$ l cm ( ) is for gas [keV]	$L_{rad}$ [cm] ( ) is for gas	[cm]		
H <sub>2</sub>	1	1.01	0.0387	0.033	43.3	50.8	4.12	(0.19)	61.28	865	0.0708(0.090)	1.112(140)
D <sub>2</sub>	1	2.01	0.073	0.061	45.7	54.7	2.07	(0.17)	122.6	757	0.162(0.177)	1.128
He	2	4.00	0.133	0.102	49.9	65.1	1.94	(0.16)	94.32	755	0.125(0.178)	1.024(35)
Li	3	6.94	0.211	0.157	54.6	73.4	1.58	0.70	82.76	155	0.534	—
Be	4	9.01	0.268	0.199	55.8	75.2	1.61	2.61	65.19	35.3	1.848	—
C	6	12.01	0.331	0.231	60.2	86.3	1.78	3.57	42.70	18.8	2.265 <sup>g</sup>	—
N <sub>2</sub>	7	14.01	0.379	0.265	61.4	87.8	1.82	(0.93)	37.99	47.0	0.808(1.25)	1.205(300)
O <sub>2</sub>	8	16.00	0.420	0.292	63.2	91.0	1.82	(1.31)	34.24	30.0	1.14(1.43)	1.22(266)
Ne	10	20.18	0.507	0.347	66.1	96.6	1.73	(0.75)	28.94	24.0	1.207(0.90)	1.092(67)
Al	13	26.98	0.634	0.421	70.6	106.4	1.62	3.81	24.01	8.9	2.70	—
Si	14	28.09	0.660	0.440	70.6	106.0	1.66	3.36	21.82	9.36	2.33	—
Ar	18	39.95	0.868	0.566	76.4	117.2	1.51	(1.30)	19.55	14.0	1.40(1.78)	1.233(283)
Fe	26	55.85	1.120	0.703	82.8	131.9	1.48	10.7	13.84	1.76	7.87	—
Cu	29	63.54	1.232	0.782	85.6	134.9	1.44	11.85	12.86	1.43	8.96	—
Sn	50	118.69	1.967	1.21	100.2	163	1.26	8.3	8.82	1.21	7.31	—
Xe	54	131.30	2.120	1.29	102.8	169	1.24	(3.57)	8.48	2.77	3.057(5.89)	(705)
W	74	183.85	2.767	1.65	110.3	185	1.16	21.1	6.76	0.35	19.3	—
Pb	82	207.19	2.960	1.77	116.2	194	1.13	11.7	6.37	0.56	11.35	—
U	92	238.03	3.378	1.98	117.0	199	1.09	19.3	6.00	≈0.32	≈18.95	—
Air, 20°C, 1 atm. (STP in paren.)					62.0	90.0	1.82	(1.12)	36.66	(30420)	0.001205(1.29)	1.000273(293)
H <sub>2</sub> O					60.1	84.9	2.03	1.72	36.08	36.1	1.00	1.33
Shielding concrete <sup>h</sup>					67.4	99.9	1.70	3.68	26.7	10.7	2.5	—
SiO <sub>2</sub> (quartz)					67.0	99.2	1.72	3.28	27.05	12.3	2.64	1.458
H <sub>2</sub> (bubble chamber 26°K)					43.3	50.8	4.12	0.20	61.28	≈1000	≈0.063 <sup>i</sup>	1.100
D <sub>2</sub> (bubble chamber 31°K)					45.7	54.7	2.07	0.22	122.6	≈900	≈0.140 <sup>i</sup>	1.110
H-Ne mixture (50 mole percent) <sup>j</sup>					65.0	94.5	1.84	0.59	29.70	73.0	0.407	1.092
Ilford emulsion G5					82.0	134	1.44	4.79	11.0	2.89	3.815	—
NaI					94.8	152	1.32	4.13	9.49	2.59	3.67	1.775
BaF <sub>2</sub>					92.1	146	1.35	5.72	9.91	2.05	4.89	1.56
BGO (Bi <sub>4</sub> Ge <sub>3</sub> O <sub>12</sub> )					97.4	156	1.27	8.07	7.98	1.12	7.1	2.15
Polystyrene, scintillator (CH) <sup>k</sup>					58.4	82.0	1.95	1.72	43.8	42.4	1.032	1.581
Lucite, Plexiglas (C <sub>5</sub> H <sub>8</sub> O <sub>2</sub> )					59.2	83.6	1.95	1.98	40.55	≈34.4	1.16–1.20	≈1.49
Polyethylene (CH <sub>2</sub> )					56.9	78.8	2.09	1.68	44.8	≈47.9	0.92–0.95	—
Mylar (C <sub>5</sub> H <sub>4</sub> O <sub>2</sub> )					60.2	85.7	1.86	2.24	39.95	28.7	1.39	—
Borosilicate glass (Pyrex) <sup>l</sup>					66.2	97.6	1.72	3.32	28.3	12.7	2.23	1.474
CO <sub>2</sub>					62.4	90.5	1.82	(1.92)	36.2	(18310)	(1.977)	(410)
Methane CH <sub>4</sub>					54.7	74.0	2.41	(0.91)	46.5	(64850)	0.423(0.717)	(444)
Isobutane C <sub>4</sub> H <sub>10</sub>					56.3	77.4	2.22	(3.43)	45.2	(16930)	(2.67)	(1270)
Freon 12 (CCl <sub>2</sub> F <sub>2</sub> ) gas, 26°C, 1 atm. <sup>m</sup>					70.6	106	1.62	4.49	23.7	4810	(4.93)	1.001080
Silica Aerogel <sup>n</sup>					65.5	95.7	1.83	0.28	29.85	≈150	0.1–0.3	1.0+0.25ρ
G10 plate <sup>o</sup>					62.6	90.2	1.87	2.7	33.0	19.4	1.7	—

\* Table revised April 1986 by W. Carithers.  $\sigma_T$ ,  $\sigma_I$ ,  $\lambda_T$ , and  $\lambda_I$  are energy dependent. Values quoted apply to high energy range given in footnote a or b, where energy dependence is weak.

a.  $\sigma_{total}$  at 80–240 GeV for neutrons ( $\approx \sigma$  for protons) from Murthy et al., Nucl. Phys. **B92**, 269 (1975). This scales approximately as  $A^{0.77}$ .

b.  $\sigma_{inelastic} = \sigma_{total} - \sigma_{elastic} - \sigma_{quasielastic}$ ; for neutrons at 60–375 GeV from Roberts et al., Nucl. Phys., **B159**, 56 (1979). For protons and other particles, see Carroll et al., Phys. Lett. **80B**, 319 (1979); note that  $\sigma_I(p) \approx \sigma_I(n)$ .  $\sigma_I$  scales approximately as  $A^{0.71}$ .

c. Mean free path between collisions ( $\lambda_T$ ) or inelastic interactions ( $\lambda_I$ ), calculated from  $\lambda = A/(N \times \sigma)$ , where  $N$  is the Avogadro number.

d. For minimum-ionizing protons and pions.  $\Delta E$  is energy loss per g/cm<sup>2</sup> from Barkas and Berger, *Tables of Energy Losses and Ranges of Heavy Charged Particles*, NASA-SP-3013 (1964). For electrons and positrons see: M.J. Berger and S.M. Seltzer, *Stopping Powers and Ranges of Electrons and Positrons* (2<sup>nd</sup> Ed.), U.S. National Bureau of Standards report NBSIR 82-2550-A (1982).  $\Delta E_{mp}$  is the most probable deposited energy in one cm, in MeV for solids and liquids, in keV for gases.  $E_{mp}$  varies with depth in a nonproportional manner. [See Sect. (1) of Passage of Particles Through Matter.] Parentheses refer to gaseous form at STP (0°C, 1 atm.).

e. From Y.S. Tsai, Rev. Mod. Phys. **46**, 815 (1974);  $L_{rad}$  data for all elements up to uranium may be found here. Corrections for molecular binding applied for H<sub>2</sub> and D<sub>2</sub>. Parentheses refer to gaseous form at STP (0°C, 1 atm.).

f. Values for solids, or the liquid phase at boiling point, except as noted. Values in parentheses for gaseous phase at STP (0°C, 1 atm.). Refractive index given for sodium D line.

g. For pure graphite; industrial graphite density may vary 2.1–2.3 g/cm<sup>3</sup>.

h. Standard shielding blocks, typical composition O<sub>2</sub> 52%, Si 32.5%, Ca 6%, Na 1.5%, Fe 2%, Al 4%, plus reinforcing iron bars. The attenuation length,  $\ell = 115 \pm 5$  g/cm<sup>2</sup>, is also valid for earth (typical  $\rho = 2.15$ ), from CERN-LRL-RHEL Shielding exp., UCRL-17841 (1968).

i. Density may vary about  $\pm 3\%$ , depending on operating conditions.

j. Values for typical working conditions with H<sub>2</sub> target: 50 mole percent, 29°K, 7 atm.

k. Typical scintillator; e.g., PILOT B and NE 102A have an atomic ratio H/C = 1.10.

l. Main components: 80% SiO<sub>2</sub> + 12% B<sub>2</sub>O<sub>3</sub> + 5% Na<sub>2</sub>O.

m. Used in Cerenkov counters. Values at 26°C and 1 atm. Indices of refraction from E.R. Hayes, R.A. Schluter, and A. Tamosaitis, ANL-6916 (1964).

n.  $n(\text{SiO}_2) + 2n(\text{H}_2\text{O})$  used in Cerenkov counters,  $\rho$  = density in g/cm<sup>3</sup>. From M. Cantin et al., Nucl. Instr. Meth. **118**, 177 (1974).

o. G10-plate, typical 60% SiO<sub>2</sub> and 40% epoxy.

## EXPECTED PARAMETERS OF FUTURE HIGH ENERGY COLLIDERS

The numbers shown here were received from authorized representatives of each collider in November 1985. All numbers are subject to change, and many are only estimates.

$H$   $\equiv$  horizontal direction,  $V$   $\equiv$  vertical direction, s.c.  $\equiv$  superconducting

	TEVATRON (Fermilab)	TRISTAN (KEK)	SLC (SLAC)	LEP (CERN)	HERA (DESY)	UNK (Serpukhov)	SSC (unknown)
Start date	Sept. 1986	Nov. 1986	March 1987	Early 1989	Spring 1990	1993	1994
Particles collided	$p\bar{p}$	$e^+e^-$	$e^+e^-$	$e^+e^-$	$ep$	$pp$	$pp$
Max. beam energy (TeV)	0.8-1.0	0.03	0.05	0.06	0.026/0.82 $e/p$	3 (Stage II)	20
Injection energy (TeV)	0.15	0.008	0.05	0.02	0.014/0.040 $e/p$	0.4-0.6	1
Luminosity ( $10^{30} \text{ cm}^{-2} \text{ s}^{-1}$ )	1	20	6 (0.6 1 <sup>st</sup> yr)	16	15	100	1000
Circumference (km)	6.28	3.02	1.45 + 1.47	26.66	6.336	20.77	82.9
No. of interaction regions	2 high $\mathcal{L}$ 2 low $\mathcal{L}$	4	1	4	3	4	4
No. of particles per bunch (units $10^{10}$ )	6	32	7.2 (5 1 <sup>st</sup> yr)	41.6	3.48/10 $e/p$	5	0.73
No. of bunches per ring per species	3	2	1	4	220	12,000	17,280
Average beam current per species (mA)	1.4	10	0.0014	3	58/163 $e/p$	1400	73
Beam-beam tune shift per crossing (units $10^{-4}$ )	17	300	-	300	250/20 $e/p$	2	9
Filling time (min)	2-3	20	-	0.25 mA/min	10/20 $e/p$	5	30
Luminosity lifetime (hr)	10-20	3-4	-	5	>3	1	73
Crossing angle ( $\mu$ rad)	0	0	0.00	0	0	0-500	75
Energy spread (units $10^{-3}$ )	0.12	1.64	0.2	0.82	0.91/0.4 $e/p$	0.13	0.5
Transverse emittance ( $10^{-9} \pi$ rad-m)	4.2	$H$ : 180	0.42	$H$ : 56 $V$ : 2.2	$H$ : 34.5/8.6 $e/p$ $V$ : 6.90/4.3 $e/p$	47	0.047
RF frequency (MHz)	53	508	-	352.2	499.7/208.2 $e/p$	200	375
Acceleration period (sec)	50	300	-	80		100	1000
Bunch length (cm)	50	1.2	0.1	1.6	0.78/15 $e/p$	50	7
Repetition rate (Hz)			180 (120 1 <sup>st</sup> yr)				
$\beta^*$ , amplitude function at interaction point (m)	1	$H$ : 0.8 $V$ : 0.05	0.01	$H$ : 1.75 $V$ : 0.07	$H$ : 2/10 $e/p$ $V$ : 0.70/1.0 $e/p$	3	$H$ : 0.5 $V$ : 0.5
Free space at interaction point (m)	$\pm 6.5$	$\pm 2.5$ $\pm 4.5$ (initial)	$\pm 2.2$	$\pm 3.5$	$\pm 5.5$	$\pm 9$	$\pm 20$
Beam radius ( $10^{-6}$ m)	65	$H$ : 367 $V$ : 23	1.7 (2.1 1 <sup>st</sup> yr)	$H$ : 312 $V$ : 12.5	$H$ : 263/293 $e/p$ $V$ : 69/66 $e/p$	375	4.8
No. of utility insertions	4	8	-	2	4	4	4
Length of standard cell (m)	59.5	16.1	5.2	79	23.5/47 $e/p$	91.8	192
Phase advance per cell (deg)	65	60	108	60	60/90 $e/p$	82.5	60
Magnetic length of dipole (m)	6.12	5.86	2.5	11.66/pair	9.2/8.9 $e/p$	5.8	16.5
No. of dipoles in ring	774	272	460+440	3280+24 inj. + 64 weak	400/416 $e/p$	2176	$H$ : 3840 $\times$ 2 (s.c.)
No. of quadrupoles in ring	216	400	-	520+288 + 8 s.c.	568/242 $e/p$	454	920 (1 ring) 1808 (2 r.)(s.c.)
Magnet type	s.c. $\cos \theta$ warm iron	room temp. C type	AGF in arcs cold bore cold iron nonsat.	iron domin. + 8 s.c. quads	C-shaped $e$ s.c., collared, cold iron $p$	s.c.	s.c., collared, cold iron
Peak magnetic field (T)	4.4	0.406	0.597	0.135	0.165/4.65 $e/p$	5	6.6
$\bar{p}$ source accum. rate ( $\text{hr}^{-1}$ )	$4 \times 10^{10}$	-	-	-	-	-	-
Max. no. $\bar{p}$ in accum. ring	$4 \times 10^{11}$	-	-	-	-	-	-

PERIODIC TABLE OF THE ELEMENTS

IA		VIII										VIIA										VIA	VA	IVA	III A	IIA	VIIIA	VIIA	VIA	VA	IVA	III A	IIA	IA																																																																																																																																																																											
1	H	21	Sc	22	Ti	23	V	24	Cr	25	Mn	26	Fe	27	Co	28	Ni	29	Cu	30	Zn	31	Ga	32	Ge	33	As	34	Se	35	Br	36	Kr	37	Rb	38	Sr	39	Y	40	Zr	41	Nb	42	Mo	43	Tc	44	Ru	45	Rh	46	Pd	47	Ag	48	Cd	49	In	50	Sn	51	Sb	52	Te	53	I	54	Xe																																																																																																																																								
1	H 1.00794	2	He 4.00260	3	Li 6.941	4	Be 9.01218	5	B 10.81	6	C 12.011	7	N 14.0067	8	O 15.9994	9	F 18.998403	10	Ne 20.179	11	Na 22.98977	12	Mg 24.305	13	Al 26.98154	14	Si 28.0855	15	P 30.97376	16	S 32.06	17	Cl 35.453	18	Ar 39.948	19	K 39.0983	20	Ca 40.08	21	Sc 44.9559	22	Ti 47.88	23	V 50.9415	24	Cr 51.996	25	Mn 54.9380	26	Fe 55.847	27	Co 58.9332	28	Ni 58.69	29	Cu 63.546	30	Zn 65.38	31	Ga 69.72	32	Ge 72.59	33	As 74.9216	34	Se 78.96	35	Br 79.904	36	Kr 83.80	37	Rb 85.4678	38	Sr 87.62	39	Y 88.9059	40	Zr 91.22	41	Nb 92.90634	42	Mo 95.94	43	Tc 98	44	Ru 101.07	45	Rh 102.9055	46	Pd 106.42	47	Ag 107.8682	48	Cd 112.41	49	In 114.82	50	Sn 117.60	51	Sb 121.75	52	Te 127.60	53	I 126.9045	54	Xe 131.29	55	Cs 132.9054	56	Ba 137.33	57	La 138.9054	58	Ce 140.12	59	Pr 140.9077	60	Nd 144.24	61	Pm 145	62	Sm 150.36	63	Eu 151.96	64	Gd 157.25	65	Tb 158.9254	66	Dy 162.50	67	Ho 164.9304	68	Er 167.26	69	Tm 168.9342	70	Yb 173.04	71	Lu 174.967	72	Hf 178.49	73	Ta 180.9479	74	W 183.85	75	Re 186.207	76	Os 190.2	77	Ir 192.22	78	Pt 195.08	79	Au 196.9665	80	Hg 200.59	81	Tl 204.383	82	Pb 207.2	83	Bi 208.9804	84	Po (209)	85	At (210)	86	Rn (222)	87	Fr 223	88	Ra 226.0254	89	Ac 227.0278	90	Th 232.0381	91	Pa 231.0359	92	U 238.0289	93	Np 237.0482	94	Pu 244	95	Am 243	96	Cm 247	97	Bk 247	98	Cf 251	99	Es 252	100	Fm 257	101	Md 258	102	No 259	103	Lr 260

The atomic number (top left) is the number of protons in the nucleus. The atomic mass (bottom) is weighted by isotopic abundance in earth's surface, relative to the mass of the carbon 12 isotope, which has been arbitrarily assigned a mass of 12.00000 atomic mass units (amu). Relative isotopic abundances vary considerably for many elements in naturally occurring specimens; commercially available samples may also vary due to undisclosed or inadvertent isotope separation. Numbers in parentheses are mass numbers (the whole number nearest the atomic mass, in amu) of most stable isotope of that element. Adapted from the *Handbook of Chemistry and Physics*, 66<sup>th</sup> Ed., 1985-1986.

## ELECTRONIC STRUCTURE OF THE ELEMENTS

At. no.	Name	Chem. Symbol	Electronic Configuration				
			K	L	M	N	O
			s	s p	s p d	s p d	s p
1	Hydrogen	H	1				
2	Helium	He	2				
3	Lithium	Li	2	1			
4	Beryllium	Be	2	2			
5	Boron	B	2	2	1		
6	Carbon	C	2	2	2		
7	Nitrogen	N	2	2	3		
8	Oxygen	O	2	2	4		
9	Fluorine	F	2	2	5		
10	Neon	Ne	2	2	6		
11	Sodium	Na	2	2	6	1	
12	Magnesium	Mg	2	2	6	2	
13	Aluminum	Al	2	2	6	2	1
14	Silicon	Si	2	2	6	2	2
15	Phosphorus	P	2	2	6	2	3
16	Sulfur	S	2	2	6	2	4
17	Chlorine	Cl	2	2	6	2	5
18	Argon	Ar	2	2	6	2	6
19	Potassium	K	2	2	6	2	6 ... 1
20	Calcium	Ca	2	2	6	2	6 ... 2
21	Scandium	Sc	2	2	6	2	6 1 2
22	Titanium	Ti	2	2	6	2	6 2 2
23	Vanadium	V	2	2	6	2	6 3 2
24	Chromium	Cr	2	2	6	2	6 5* 1
25	Manganese	Mn	2	2	6	2	6 5 2
26	Iron	Fe	2	2	6	2	6 6 2
27	Cobalt	Co	2	2	6	2	6 7 2
28	Nickel	Ni	2	2	6	2	6 8 2
29	Copper	Cu	2	2	6	2	6 10* 1
30	Zinc	Zn	2	2	6	2	6 10 2
31	Gallium	Ga	2	2	6	2	6 10 2 1
32	Germanium	Ge	2	2	6	2	6 10 2 2
33	Arsenic	As	2	2	6	2	6 10 2 3
34	Selenium	Se	2	2	6	2	6 10 2 4
35	Bromine	Br	2	2	6	2	6 10 2 5
36	Krypton	Kr	2	2	6	2	6 10 2 6
37	Rubidium	Rb	2	2	6	2	6 10 2 6 ... 1
38	Strontium	Sr	2	2	6	2	6 10 2 6 ... 2
39	Yttrium	Y	2	2	6	2	6 10 2 6 1 2
40	Zirconium	Zr	2	2	6	2	6 10 2 6 2 2
41	Niobium	Nb	2	2	6	2	6 10 2 6 4* 1
42	Molybdenum	Mo	2	2	6	2	6 10 2 6 5 1
43	Technetium	Tc	2	2	6	2	6 10 2 6 6 1
44	Ruthenium	Ru	2	2	6	2	6 10 2 6 7 1
45	Rhodium	Rh	2	2	6	2	6 10 2 6 8 1
46	Palladium	Pd	2	2	6	2	6 10 2 6 10* 0
47	Silver	Ag	2	2	6	2	6 10 2 6 10 1
48	Cadmium	Cd	2	2	6	2	6 10 2 6 10 2
49	Indium	In	2	2	6	2	6 10 2 6 10 2 1
50	Tin	Sn	2	2	6	2	6 10 2 6 10 2 2
51	Antimony	Sb	2	2	6	2	6 10 2 6 10 2 3
52	Tellurium	Te	2	2	6	2	6 10 2 6 10 2 4
53	Iodine	I	2	2	6	2	6 10 2 6 10 2 5
54	Xenon	Xe	2	2	6	2	6 10 2 6 10 2 6
55	Cesium	Cs	2	2	6	2	6 10 2 6 10 ... 2 6 ... 1
56	Barium	Ba	2	2	6	2	6 10 2 6 10 ... 2 6 ... 2
57	Lanthanum	La	2	2	6	2	6 10 2 6 10 ... 2 6 1 ... 2
58	Cerium	Ce	2	2	6	2	6 10 2 6 10 2* 2 6 ... 2
59	Praseodymium	Pr	2	2	6	2	6 10 2 6 10 3 2 6 ... 2
60	Neodymium	Nd	2	2	6	2	6 10 2 6 10 4 2 6 ... 2
61	Promethium	Pm	2	2	6	2	6 10 2 6 10 5 2 6 ... 2
62	Samarium	Sm	2	2	6	2	6 10 2 6 10 6 2 6 ... 2
63	Europium	Eu	2	2	6	2	6 10 2 6 10 7 2 6 ... 2
64	Gadolinium	Gd	2	2	6	2	6 10 2 6 10 7 2 6 1 ... 2
65	Terbium	Tb	2	2	6	2	6 10 2 6 10 9* 2 6 ... 2
66	Dysprosium	Dy	2	2	6	2	6 10 2 6 10 10 2 6 ... 2
67	Holmium	Ho	2	2	6	2	6 10 2 6 10 11 2 6 ... 2
68	Erbium	Er	2	2	6	2	6 10 2 6 10 12 2 6 ... 2
69	Thulium	Tm	2	2	6	2	6 10 2 6 10 13 2 6 ... 2
70	Ytterbium	Yb	2	2	6	2	6 10 2 6 10 14 2 6 ... 2
71	Lutetium	Lu	2	2	6	2	6 10 2 6 10 14 2 6 1 ... 2
72	Hafnium	Hf	2	2	6	2	6 10 2 6 10 14 2 6 2 ... 2
73	Tantalum	Ta	2	2	6	2	6 10 2 6 10 14 2 6 3 ... 2
74	Tungsten	W	2	2	6	2	6 10 2 6 10 14 2 6 4 ... 2
75	Rhenium	Re	2	2	6	2	6 10 2 6 10 14 2 6 5 ... 2
76	Osmium	Os	2	2	6	2	6 10 2 6 10 14 2 6 6 ... 2
77	Iridium	Ir	2	2	6	2	6 10 2 6 10 14 2 6 7 ... 2
78	Platinum	Pt	2	2	6	2	6 10 2 6 10 14 2 6 9 ... 1
79	Gold	Au	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 1
80	Mercury	Hg	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2
81	Thallium	Tl	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 1
82	Lead	Pb	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 2
83	Bismuth	Bi	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 3
84	Polonium	Po	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 4
85	Astatine	At	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 5
86	Radon	Rn	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 6
87	Francium	Fr	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 6 ... 1
88	Radium	Ra	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 6 ... 2
89	Actinium	Ac	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 6 1 2
90	Thorium	Th	2	2	6	2	6 10 2 6 10 14 2 6 10 ... 2 6 2 2
91	Protactinium	Pa	2	2	6	2	6 10 2 6 10 14 2 6 10 2* 2 6 1 2
92	Uranium	U	2	2	6	2	6 10 2 6 10 14 2 6 10 3 2 6 1 2
93	Neptunium	Np	2	2	6	2	6 10 2 6 10 14 2 6 10 4 2 6 1 2
94	Plutonium	Pu	2	2	6	2	6 10 2 6 10 14 2 6 10 6 2 6 ... 2
95	Americium	Am	2	2	6	2	6 10 2 6 10 14 2 6 10 7 2 6 ... 2
96	Curium	Cm	2	2	6	2	6 10 2 6 10 14 2 6 10 7 2 6 1 2
97	Berkelium	Bk	2	2	6	2	6 10 2 6 10 14 2 6 10 9* 2 6 ... 2
98	Californium	Cf	2	2	6	2	6 10 2 6 10 14 2 6 10 10 2 6 ... 2
99	Einsteinium	Es	2	2	6	2	6 10 2 6 10 14 2 6 10 11 2 6 ... 2
100	Fermium	Fm	2	2	6	2	6 10 2 6 10 14 2 6 10 12 2 6 ... 2
101	Mendelevium	Md	2	2	6	2	6 10 2 6 10 14 2 6 10 13 2 6 ... 2
102	Nobelium	No	2	2	6	2	6 10 2 6 10 14 2 6 10 14 2 6 ... 2
103	Lawrencium	Lr	2	2	6	2	6 10 2 6 10 14 2 6 10 14 2 6 1 2
104	---	---	2	2	6	2	6 10 2 6 10 14 2 6 10 14 2 6 2 2

\*Note irregularity



## PARTICLE DETECTORS\*

In this section we give various parameters for common detectors. The quoted numbers are usually based on some typical apparatus, and obviously should be regarded as rough approximations, valid only for preliminary design when applied to other cases. A more detailed introduction to detectors can be found in "A Consumer's Guide to Particle Detectors," by D.J. Miller, Ruthford Lab Report RL-76-072, July 1976.

(1) **Scintillators:** The photon yield in the frequency range of practical photomultiplier tubes is  $\approx 1\gamma$  per 100 eV of charged particle ionization energy loss in plastic scintillator<sup>1</sup> and  $\approx 1\gamma/25$  eV in NaI.<sup>1,2</sup>

(2) **Cerenkov:**<sup>3</sup> The half-angle  $\theta_c$  of the Cerenkov cone aperture in terms of the velocity  $\beta$  and the index of refraction  $n$  is:

$$\theta_c = \arccos \left( \frac{1}{\beta n} \right) \approx \left[ 2 \left( 1 - \frac{1}{\beta n} \right) \right]^{1/2}.$$

The threshold velocity is:  $\beta_t = 1/n$ ;  $\gamma_t = 1/(1 - \beta_t^2)^{1/2}$ . Therefore,  $\beta_t \gamma_t = 1/(2\delta + \delta^2)^{1/2}$ , where  $\delta = n - 1$ . Values of  $\delta$  for various commonly used gases are given as a function of pressure and wavelength in Ref. 4; for values at atmospheric pressure, see the Table of Atomic and Nuclear Properties.

The number of photons  $N$  per cm of path length is given by:

$$N = \frac{\alpha}{c} \int \left( 1 - \frac{1}{\beta^2 n^2} \right) 2\pi d\nu = \frac{\alpha}{c} \beta_t^2 \int \left( \frac{1}{\beta_t^2 \gamma_t^2} - \frac{1}{\beta^2 \gamma^2} \right) 2\pi d\nu$$

$$\approx 500 \sin^2 \theta_c / \text{cm (visible spectrum)}.$$

(3) **Photon collection:** In addition to the photon yield, one should take into account the light collection efficiency ( $\approx 10\%$  for typical 1-cm-thick scintillator), the attenuation length ( $\approx 1$  to 4 m for typical scintillators<sup>5</sup>), and the quantum efficiency of the photomultiplier cathode ( $\approx 25\%$ ).

### (4) Typical detector characteristics:

Detector Type	Accuracy (rms)	Resolution Time	Dead Time
Bubble chamber	$\approx \pm 10$ to $\approx \pm 150\mu\text{m}$	$\approx 1$ ms	$\approx 1/20$ s <sup>a</sup>
Streamer chamber	$\pm 300\mu\text{m}$	$\approx 2$ $\mu\text{s}$	$\approx 100$ ms
Proportional chamber	$\approx \pm 300\mu\text{m}$ <sup>b,c</sup>	$\approx 50$ ns	$\approx 200$ ns
Drift chamber	$\pm 50$ to $300\mu\text{m}$	$\approx 2$ ns <sup>d</sup>	$\approx 100$ ns
Scintillator	—	$\approx 150$ ps	$\approx 10$ ns
Emulsion	$\pm 1\mu\text{m}$	—	—
Silicon strip	$\pm 2.5\mu\text{m}$	e	e

<sup>a</sup> Multiple pulsing time.

<sup>b</sup>  $300\mu\text{m}$  is for 1 mm pitch.

<sup>c</sup> Delay line cathode readout can give  $\pm 150\mu\text{m}$  parallel to anode wire.

<sup>d</sup> For two chambers.

<sup>e</sup> Limited at present by noise and readout time of attached electronics.

(5) **Shower detectors:** We give below typical energy resolutions (FWHM) for an incident electron in the 1 GeV range;  $E$  is in GeV. For a fixed number of radiation lengths, FWHM in the last three detectors would be expected to be proportional to  $\sqrt{t}$  for  $t$  (= plate thickness)  $\geq 0.2$  radiation lengths.<sup>6</sup>

For all detectors, operational resolution may be up to 50% worse due to dead areas, non-normally incident tracks, and other effects.

$$\text{NaI (20 rad. lengths):}^7 \frac{2\%}{E^{1/4}}$$

$$\text{Lead glass (14 rad. lengths):}^8 \frac{10 - 12\%}{\sqrt{E}}$$

$$\text{Lead-liquid argon (15.75 rad. lengths):}^6 \frac{16\%}{\sqrt{E}}$$

(42 cells: 1.1 mm lead, 2 mm liquid argon, 2.3 mm lead-G10, 2 mm liquid argon)

$$\text{Lead-scintillator sandwich (12.5 rad. lengths):}^9 \frac{17\%}{\sqrt{E}}$$

(66 cells: 1 mm lead, 5 mm scintillator)

$$\text{Proportional wire shower chamber (17 rad. lengths):}^{10} \frac{40\%}{\sqrt{E}}$$

(36 cells: 0.474 rad. length type-metal + Al, 9.5 mm 80% Ar - 20% CH<sub>4</sub> gas)

(6)  **$dE/dx$  resolution in argon:** Particle identification (relativistic,  $Q = 1$  incident particles) by  $dE/dx$  is dependent on the width of the distribution:

Multiple-sample Ar gas counters (no lead):<sup>11</sup>

$$\text{FWHM} \left( \frac{dE}{dx} \Big|_{\text{most probable}} \right) = 0.96N^{-0.46}(tp)^{-0.32};$$

$$\frac{dE}{dx} \Big|_{\text{most probable}}$$

$N$  = no. samples,  $t$  = thickness per sample (cm),  $p$  = pressure (atm.); most commonly used chamber gases (except Xe) give approximately the same resolution.

(7) **Proportional chamber wire instability:** The limit on the voltage  $V$  for a wire tension  $T$ , due to mechanical effects when the electrostatic repulsion of adjacent wires exceeds the restoring force of wire tension, is given by (MSKA)<sup>12</sup>

$$V < \frac{S}{\ell C} \sqrt{4\pi\epsilon_0 T},$$

where  $s$ ,  $\ell$ , and  $C$  are the wire spacing, length, and capacitance per unit length. An approximation to  $C$  for chamber half-gap  $t$  and wire diameter  $d$  (good for  $s \approx t$ ) gives<sup>13</sup>

$$V \approx 59T^{1/2} \left[ \frac{t}{\ell} + \frac{s}{\pi\ell} \ell n \left( \frac{s}{\pi d} \right) \right],$$

where  $V$  is in kV, and  $T$  is in grams-weight equivalent.

(8) **Proportional and drift chamber potentials:** The potential distributions and fields in a proportional or drift chamber can usually be calculated with good accuracy from the exact formula for the potential around an array of parallel line charges  $q$  (coul/m) along  $z$  and located at  $y = 0$ ,  $x = 0, \pm s, \pm 2s, \dots$ ,

$$V(x,y) = -\frac{q}{4\pi\epsilon_0} \ell n \left\{ 4 \left[ \sin^2 \left( \frac{\pi x}{s} \right) + \sinh^2 \left( \frac{\pi y}{s} \right) \right] \right\}.$$

Errors from the presence of cathodes, mechanical defects, TPC-type edge effects, etc., are usually small and are beyond the scope of this review.

## PARTICLE DETECTORS (Cont'd)

(9) **Silicon strip detectors and photodiodes:** These are silicon diodes operated with a reverse bias voltage  $V$  (typically 30–300 volts) sufficient to deplete the sensitive volume of most mobile charge carriers (electrons and holes). The active (depletion layer) thickness  $t$  (cm) is given in a simple model by

$$t = \sqrt{\frac{2\epsilon V}{ne}} = \sqrt{2\rho\mu\epsilon V},$$

where

- $n$  = number of impurity centers/cm<sup>3</sup>
- $e$  = electron charge
- $\epsilon$  = dielectric constant  $\cong 1 pF/cm \cong 11.9 \epsilon_0$
- $\rho$  = resistivity  $\cong 1\text{--}20$  k $\Omega$ -cm
- $\mu$  = majority charge carrier mobility  
 $\cong 1300\text{--}1500$  cm<sup>2</sup>/volt-sec (electrons)  
 $\cong 450\text{--}600$  cm<sup>2</sup>/volt-sec (holes).

A minimum-ionizing particle has a Landau energy-loss distribution with average energy loss 39 keV/100  $\mu$ m, most probable energy loss 26 keV in 100  $\mu$ m (which scales within  $\sim \pm 10\%$  from  $\sim 20$  to  $\sim 300$   $\mu$ m), and full width at half-maximum of roughly  $0.1t/\beta^2$  keV, where  $t$  is the detector thickness in microns and  $\beta = v_{inc}/c$ . The width is usually increased further by electronic noise ( $\sigma \sim 1\text{--}10$  keV) and for thin layers by a Gaussian contribution due to atomic effects [ $\sigma \sim (0.3\text{--}0.4)\sqrt{t}$  keV]. The average energy required to produce an electron-hole pair is 3.6 eV, from which one can estimate total charge of either sign released. Silicon detectors can tolerate integrated charged-particle fluxes of up to  $\sim 10^{10}\text{--}10^{14}$ /cm<sup>2</sup> and still operate as efficient detectors.

Typical photodiodes are sensitive (quantum efficiencies greater than  $\sim 10\%$ ) to wavelengths from  $\sim 200$  nm to 1100 nm.

\* Updated April 1986 by Sherwood Parker, Ray Hagstrom, and Geoff Hall.

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## COSMIC RAY FLUXES\*

The fluxes of particles of different types depend at the  $\sim 10\%$  level on the latitude, their energy, and the conditions of measurement. Some typical sea-level values<sup>1</sup> for charged particles are given below:

- $I_v$  flux per unit solid angle per unit horizontal area about vertical direction  
 $\equiv j(\theta=0, \phi)$  [ $\theta$  = zenith angle,  $\phi$  = azimuthal angle];
- $J_1$  total flux crossing unit horizontal area from above  
 $\equiv \int_{\theta \leq \pi/2} j(\theta, \phi) \cos \theta d\Omega$  [ $d\Omega = \sin \theta d\theta d\phi$ ];
- $J_2$  total flux from above (impinging on a sphere of unit cross-sectional area)  
 $\equiv \int_{\theta \leq \pi/2} j(\theta, \phi) d\Omega$ .

	Total Intensity	Hard Component	Soft Component
$I_v$	$1.1 \times 10^2$	$0.8 \times 10^2$	$0.3 \times 10^2$ m <sup>-2</sup> sec <sup>-1</sup> sterad <sup>-1</sup>
$J_1$	$1.8 \times 10^2$	$1.3 \times 10^2$	$0.5 \times 10^2$ m <sup>-2</sup> sec <sup>-1</sup>
$J_2$	$2.4 \times 10^2$	$1.7 \times 10^2$	$0.7 \times 10^2$ m <sup>-2</sup> sec <sup>-1</sup>

Very approximately, about 75% of all particles at sea level are penetrating, and are muons (the dominant portion of the hard component at sea level). The sea-level vertical flux ratio for protons to muons (both charges together) is about  $3\frac{1}{2}$  at 1 GeV/c, decreasing to about  $\frac{1}{2}$  at 10 GeV/c.

The muon flux at sea level has a mean energy of 2 GeV and a differential spectrum falling as  $E^{-2}$ , steepening smoothly to  $E^{-3.6}$  above a few TeV. The angular distribution is  $\cos^2\theta$ , changing to  $\sec\theta$  at energies above a TeV, where  $\theta$  is the zenith angle at production. The  $+$ - charge ratio is 1.25–1.30. The mean energy of muons originating in the atmosphere is roughly 300 GeV at slant depths  $\geq$  a few hundred meters. Beyond slant depths of  $\sim 10$  km water-equivalent, the muons are due primarily to in-the-earth neutrino interactions (roughly 1/8 interaction ton<sup>-1</sup> year<sup>-1</sup> for  $E_\nu > 300$  MeV,  $\sim$  constant throughout the earth).<sup>2</sup> Muons from this source arrive with a mean energy of 20 GeV, and have a flux of  $2 \times 10^{-9}$  m<sup>-2</sup> sec<sup>-1</sup> sterad<sup>-1</sup> in the vertical direction and about twice that in the horizontal,<sup>3</sup> down at least as far as the deepest mines.

\* Updated April 1986.

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## PASSAGE OF PARTICLES THROUGH MATTER\*

(1) **Energy loss rates for heavy charged projectiles:** A heavy projectile (much more massive than an electron) of charge  $Z_{\text{inc}}e$ , incident at speed  $\beta c$  ( $\beta \gg 1/137$ ) through a slowing medium, dissipates energy principally via interactions with the electrons of the medium. The mean rate of such energy loss per unit path length  $x$ , called the stopping power, is given by the Bethe-Bloch equation:<sup>1</sup>

$$\left( \frac{dE}{dx} \right)_{\text{inc}} = \frac{D Z_{\text{med}} \rho_{\text{med}}}{A_{\text{med}}} \left( \frac{Z_{\text{inc}}}{\beta} \right)^2 \times \left[ \ell n \left( \frac{2m_e \gamma^2 \beta^2 c^2}{I} \right) - \beta^2 - \frac{\delta}{2} - \frac{C}{Z_{\text{med}}} \right] \left\{ 1 + \nu \right\},$$

where  $D = 4\pi N_A r_e^2 m_e c^2 = 0.3070 \text{ MeV cm}^2/\text{g}$  (see Physical Constants Table). Mean range and energy loss figures appear at the end of this section.

Here,  $Z_{\text{med}}$  and  $A_{\text{med}}$  are the charge and mass numbers of the medium and  $\rho_{\text{med}}$  is the mass density of the medium;  $I$ ,  $\delta$ ,  $C$ , and  $\nu$  are phenomenological functions. Frequently, the values of  $\delta$ ,  $C$ , and  $\nu$  are negligibly small; the parameter  $I$  characterizes the binding of the electrons of the medium. As a rule of thumb, we may estimate  $I$  for an idealized medium as  $I \approx 16 (Z_{\text{med}})^{0.9} \text{ eV}$  when  $Z_{\text{med}} > 1$ . For realistic media the value of  $I$  will vary at the 10% level from this estimate. Variations of this order occur due to atomic effects such as completion of a shell, also due to chemical binding, and even due to the phase of the substance. Hydrogen, perhaps the most sensitive, has  $I$  of about 15 eV in the atomic mode, rising to about 19.2 eV as  $\text{H}_2$  gas and to 21.8 eV as  $\text{H}_2$  liquid.<sup>2</sup> For many substances, the transition from gas to solid is accompanied by a 20-30% increase in  $I$ .<sup>2</sup> We may *approximately* treat media which are chemical mixtures or compounds by computing

$$\frac{dE}{dx} \approx \sum_{n=1}^N \left( \frac{dE}{dx} \right)_n,$$

with  $(dE/dx)_n$  appropriate to the  $n^{\text{th}}$  chemical constituent (using  $\rho_{\text{med}}^{(n)}$  as the partial density in the formula for  $dE/dx$ ).<sup>3</sup> For many chemical compounds, small corrections to this additivity rule may be found in Ref. 2.

The function  $\delta$  represents the density effect upon the energy loss rate; it is non-negligible only for highly relativistic projectiles in denser media.<sup>4</sup> For ultra-relativistic projectiles,  $\delta$  approaches  $2\ell n \gamma + \text{constant}$ , where the value of the constant depends upon the density of the medium as well as its chemical composition.

The function  $C$  represents shell corrections to the energy loss rate.<sup>1</sup> These effects are non-negligible only for projectiles with speeds not much faster than the speeds of the fastest electrons bound in the medium.

The function  $\nu$  represents corrections due to higher order electrodynamics.<sup>5</sup> These effects become important when  $|Z_{\text{inc}}/\beta|$  is comparable to 137. For relativistic unit-charge projectiles,  $|\nu|$  is of the order of 1%; positively charged projectiles lose energy more rapidly than do their charge conjugates.<sup>5,6</sup>

For nonrelativistic projectiles, our formulae above are inapplicable. At the very slowest speeds, total energy loss rates are believed to be proportional to  $\beta$ , rising through a peak at projectile speeds comparable to atomic speeds ( $\beta$  on the order of  $\alpha c$ ), after having passed through a smaller peak (due to elastic Coulomb collisions with the nuclei of the slowing medium<sup>7</sup>) at intermediate speeds. For example, for protons in Si,  $dE/dx = 61.23 \beta \text{ GeV}/(\text{gm cm}^{-2})$  for  $\beta < 0.005$ ; the peak occurs at  $\beta = 0.0126$  where  $dE/dx = 522 \text{ MeV}/(\text{gm cm}^{-2})$ . In some cases, energy loss rates depend significantly upon the relation of the projectile trajectory to the crystalline structure of the slowing medium.<sup>8</sup>

For relativistic projectiles,  $(dE/dx)_{\text{inc}}$  falls rapidly with increasing  $\beta$  until reaching a minimum around  $\beta = 0.96$  (almost indepen-

dent of medium), followed by a slow rise. Because of the density effect, the quantity in square brackets approaches  $\ell n \gamma + \text{constant}$  for large  $\gamma$ .

The quantity  $(dE/dx)_{\text{inc}} \delta x$  is the *mean* total energy loss via interactions with electrons of the medium in a layer of thickness  $\delta x$ . For any finite  $\delta x$ , Poisson fluctuations can cause the actual energy loss to deviate from the mean. For thin layers, the distribution is broad and skewed, being peaked below  $(dE/dx) \delta x$ , and having a long tail toward large energy losses.<sup>9</sup> Only for a very thick layer [ $(dE/dx) \delta x \gg 2m_e \beta^2 \gamma^2 c^2$ ] will the distribution of energy losses become nearly Gaussian. The large fluctuations of the total energy loss rate from the mean are due to a small number of collisions involving large energy transfers. The fluctuations are greatly reduced for the so-called restricted energy loss rate, described in Section (4).

(2) **Ionization yields:** Physicists frequently relate total energy loss to the number of ion pairs produced near the projectile's track. This relation becomes complicated for relativistic projectiles due to the wandering of energetic knock-on electrons whose ranges exceed the dimensions of the fiducial volume. For a qualitative appraisal of the nonlocality of energy deposition by such modestly energetic knock-on electrons in various media, see Ref. 10. Furthermore, the mean local energy dissipation per local ion pair produced,  $W$ , while essentially constant for relativistic projectiles, increases at slow projectile speeds.<sup>11</sup> The numerical value of  $W$  for gases can be surprisingly sensitive to trace amounts of various contaminants.<sup>11</sup> Of course, in addition to the preceding effects, practical ionization yields may be greatly influenced by subsequent recombinations and other factors.<sup>12</sup>

(3) **Energetic knock-on electrons:** For a relativistic point-charge projectile, the production of high energy (kinetic energy  $T \gg I$ ) electrons is given by:<sup>13</sup>

$$\frac{d^2 N}{dT dx} = \frac{1}{2} D \left( \frac{Z_{\text{med}}}{A_{\text{med}}} \right) \left( \frac{Z_{\text{inc}}}{\beta} \right)^2 \rho_{\text{med}} \frac{1}{T^2} F,$$

for  $I \ll T \leq T_{\text{max}}$ , where

$$T_{\text{max}} = \frac{2m_e \beta^2 \gamma^2 c^2}{1 + 2\gamma \frac{m_e}{M_{\text{inc}}} + \left( \frac{m_e}{M_{\text{inc}}} \right)^2},$$

$M_{\text{inc}}$  is the mass of the incident projectile, and all other quantities except  $F$  are as in Sec. (1).  $F (\approx 1 \text{ for } T \ll T_{\text{max}})$  is a factor dependent upon the spin of the projectile.

For spin-0 projectiles,

$$F = 1 - \beta^2 \frac{T}{T_{\text{max}}};$$

for spin-1/2 projectiles,

$$F = 1 - \beta^2 \frac{T}{T_{\text{max}}} + \frac{1}{2} \left( \frac{T}{T_{\text{inc}} + M_{\text{inc}} c^2} \right)^2,$$

where  $T_{\text{inc}}$  is the kinetic energy of the projectile; for electrons incident,

$$F = \beta^2 T^2 \left[ \frac{T_{\text{inc}}}{T(T_{\text{inc}} - T)} - \frac{1}{T_{\text{inc}}} \right]^2;$$

and for positrons incident,

$$F = \beta^2 \left[ 1 - \frac{T}{T_{\text{inc}}} + \left( \frac{T}{T_{\text{inc}}} \right)^2 \right]^2.$$

## PASSAGE OF PARTICLES THROUGH MATTER (Cont'd)

For incident electrons, the indistinguishability of projectile and target means that the range of  $T$  is only up to  $T_{\text{inc}}/2$ . For additional formulas see Ref. 14. Our formula is inaccurate for  $T$  close to  $I$ ; for  $2I \leq T \leq 10I$ , the  $1/T^2$  dependence above becomes  $\approx T^{-\eta}$  with  $3 \leq \eta \leq 5$ .<sup>15</sup>

(4) **Rates of restricted energy loss for relativistic charged projectiles:** The variability of energy loss for heavy projectiles is due primarily to the variability in the production of energetic knock-on electrons. Bremsstrahlung and pair-production processes make this variability even greater for electrons than for heavy particles as projectiles (see, e.g., the figure Fractional Energy Loss for Electrons and Positrons in Lead). If an instrument, such as a bubble chamber, is capable of isolating these high-energy-loss interactions, then it is appropriate to consider the rate of energy loss excluding them, i.e., a restricted energy loss rate. The mean energy loss rate via all collisions which have energy transfer  $T$  such that  $T \leq E_{\text{max}} \ll T_{\text{max}}$  is:<sup>1</sup>

$$\left( \frac{dE}{dx} \right)_{\leq E_{\text{max}}} = \frac{1}{2} D \frac{Z_{\text{med}} \rho_{\text{med}}}{A_{\text{med}}} \left( \frac{Z_{\text{inc}}}{\beta} \right)^2 \times \left[ \ell n \left( \frac{E_{\text{max}} T_{\text{max}}}{I^2} \right) - \beta^2 - \delta - \frac{2C}{Z_{\text{med}}} \right]$$

Notice the overall factor of  $1/2$ . See Sec. (1) above for definitions of the quantities in this equation.

The density effect causes the restricted energy loss rate to approach a constant, the Fermi plateau value, for the fastest projectiles.

(5) **Multiple scattering through small angles:** As a charged particle traverses a medium it is deflected by many small-angle elastic scatterings. The bulk of this deflection is due to elastic Coulomb scattering from the nuclei within the medium, hence the usual identification as multiple Coulomb scattering (note, however, that strong interactions do contribute to the total multiple scattering for hadronic projectiles). For both Coulomb and strong interactions, the Central Limit Theorem provides little useful guidance in establishing the precise nature of the distribution of the total deflections resulting from multiple scattering. The true distribution is roughly Gaussian only for small deflection angles, while it shows much greater probability for large-angle scatterings ( $\geq$  a few  $\theta_0$ , see below, depending on absorber) than the Gaussian would suggest. These tails on the distribution (a few per cent of peak height in the region where the Gaussian part becomes negligible) are more pronounced for hadrons than for muons as projectiles. The large-angle behavior of these distributions is best estimated by computing the exact distribution for the vectorial sum of the largest deflections based upon the true elastic scattering cross section of the projectile against the medium,<sup>16</sup> or, when applicable, by interpolation from tabular data.<sup>17</sup> An easier alternative which may suffice for noncritical applications would be to use a Gaussian approximation with the following width:<sup>18</sup>

$$\theta_0 = \frac{14.1 \text{ MeV}/c}{p\beta} Z_{\text{inc}} \sqrt{L/L_R} \left[ 1 + \frac{1}{9} \log_{10}(L/L_R) \right] \text{ (radians),}$$

where  $p$ ,  $\beta$ , and  $Z_{\text{inc}}$  are the momentum (in MeV/c), velocity, and charge number of the incident particle, and  $L/L_R$  is the thickness, in radiation lengths, of the scattering medium.  $L_R$  for certain materials is given in the Table of Atomic and Nuclear Properties of Materials. See also Sec. (7) below. The angle,  $\theta_0$ , is a fit to Moliere<sup>16</sup> theory, accurate to about 5% for  $10^{-3} < L/L_R < 10$  except for very light elements or low velocity where the error is about 10 to 20%. In this Gaussian approximation,  $\theta_0$  has the meaning

$$\theta_0 = \theta_{\text{plane, rms}} = \frac{1}{\sqrt{2}} \theta_{\text{space, rms}}$$

The nonprojected (space) and projected (plane) angular distributions are given approximately<sup>16</sup> by the Gaussian forms:

$$\frac{1}{2\pi\theta_0^2} \exp \left[ -\frac{\theta_{\text{space}}^2}{2\theta_0^2} \right] d\Omega,$$

$$\frac{1}{\sqrt{2\pi}\theta_0} \exp \left[ -\frac{\theta_{\text{plane}}^2}{2\theta_0^2} \right] d\theta_{\text{plane}},$$

where  $\theta$  is the deflection angle.

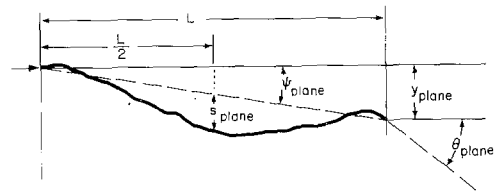
Other quantities are sometimes used to describe the amount of multiple Coulomb scattering: the auxiliary quantities  $\psi_{\text{plane}}$ ,  $y_{\text{plane}}$ , and  $s_{\text{plane}}$  (see the figure) obey:

$$\psi_{\text{plane, rms}} = \frac{1}{\sqrt{3}} \theta_{\text{plane, rms}} = \frac{1}{\sqrt{3}} \theta_0,$$

$$y_{\text{plane, rms}} = \frac{1}{\sqrt{3}} L \theta_{\text{plane, rms}} = \frac{1}{\sqrt{3}} L \theta_0,$$

and

$$s_{\text{plane, rms}} = \frac{1}{4\sqrt{3}} L \theta_{\text{plane, rms}} = \frac{1}{4\sqrt{3}} L \theta_0.$$



All the quantitative estimates in this section apply only in the limit of small  $\theta_{\text{plane, rms}}$  and in the absence of large-angle scatters.

(6) **Longitudinal distribution of electromagnetic showers:** A photon of energy  $E \geq 0.1$  GeV converting in a semi-infinite medium produces an electromagnetic cascade whose intensity initially increases with depth and then falls off. The average number of  $e^\pm$  with kinetic energy above 1.5 MeV, crossing a plane at a depth of  $L$  radiation lengths from the beginning of the medium, in a material of atomic number  $Z$ , calculated using the Monte Carlo program EGS,<sup>19</sup> can be fit by the empirical formula<sup>20</sup>

$$N = N_0 L^a e^{-bL},$$

where  $N_0 = 5.51 E(\text{GeV}) \sqrt{Z} b^{a+1} / \Gamma(a+1)$  and  $b = 0.634 - 0.0021 Z$ . For  $Z \geq 26$ ,  $a = 2.0 - Z/340 + (0.664 - Z/340)\ell n E$ . For  $Z = 13$ ,  $a = 1.77 - 0.52\ell n E$ . The maximum intensity,  $N_{\text{max}}$ , occurs at the depth  $L = a/b$ . The maximum error of the fit occurs in the vicinity of this depth and is less than  $0.15 N_{\text{max}}$ . The

integral of the tail,  $\int_{1.5a/b}^{\infty} N dL$  is fit to better than 2.5%. The total

longitudinally projected  $e^\pm$  path length,  $\int_0^{\infty} N dL = 5.51 E \sqrt{Z}$ , is less than the total  $e^\pm$  path length due primarily to multiple Coulomb scattering.

(7) **Radiation length:** For the passage of electromagnetically interacting particles through a medium it is convenient to measure

## PASSAGE OF PARTICLES THROUGH MATTER (Cont'd)

thickness in terms of radiation length.<sup>21</sup> For most electromagnetic processes (Bremsstrahlung, Coulomb scattering, showering, pair production, etc.), over large energy intervals, some or all of the dependence upon the medium is contained in the radiation length.

The radiation length may be defined as the distance  $L_R$  over which a high energy electron ( $\geq 1$  GeV for most materials) loses all but a fraction  $1/e$  of its energy to Bremsstrahlung, on average. For a homogeneous monoatomic medium,  $Z \geq 5$ ,

$$\frac{1}{L_R} = \frac{4\alpha r_e^2 N_A Z^2}{A} \left\{ \epsilon n \left( \frac{184.15}{Z^{1/3}} \right) + \frac{1}{Z} \epsilon n \left( \frac{1194}{Z^{2/3}} \right) - 1.202\alpha^2 Z^2 + 1.0369\alpha^4 Z^4 - \frac{1.008\alpha^6 Z^6}{1 + \alpha^2 Z^2} \right\} = \frac{Z^2 \left\{ \right\}}{716.405A},$$

where  $\alpha$ ,  $r_e$ , and  $N_A$  are found in the Physical Constants Table, and  $Z$  and  $A$  are the atomic number and weight of the medium. For  $Z < 5$ , a more complex numerical calculation is required. Radiation lengths for many substances are tabulated in the Table of Atomic and Nuclear Properties of Materials. For media which are chemical mixtures or compounds,

$$\frac{1}{L_R} \cong \sum_i \frac{f_i}{L_R^i},$$

where  $f_i$  is the fraction by mass of atoms of type  $i$ , radiation length  $L_R^i$ . Chemical binding can lower  $L_R$  from this, typically by a few per cent.

For electrons of energy below about one GeV, the average fractional energy loss per unit length decreases as the energy decreases (see Fractional Energy Loss for Electrons and Positrons in Lead figure). With distances measured in units of  $L_R$ , dependence of the Bremsstrahlung fractional energy loss upon  $Z$  of the medium in the low energy region ( $\geq 10$  MeV) is of order a few percent or less.

For photons of infinite energy, the total pair-production cross section is

$$\sigma = \frac{7}{9}(A/L_R N_A).$$

This is accurate to within a few per cent down to  $\sim 1$  GeV for most materials. For energies below about 1 GeV, the cross section varies in a manner which may be determined from the Photon Mass Attenuation figures. See also Contributions to Photon Cross Section in Carbon and Lead figure.

**(8) Electron practical range:** The electron "practical range" — a common measure of straight-line penetration distance — is shorter than the total path length because of multiple Coulomb scattering, which becomes increasingly important as the electron slows down. E.g., for a fast electron the rms projected angle due to multiple Coulomb scattering reaches 1 radian by the time the electron has slowed to 0.4 MeV in hydrogen, 1.5 MeV in carbon, 9 MeV in copper, and 24 MeV in lead. Electrons which have energy less than 0.2 MeV in Ar, 1.5 MeV in Cu, 3.5 MeV in Sn, and 5 MeV in Pb are likely to deposit 10% of their energy *behind* their starting plane. The practical range,  $R_p$ , is defined as that absorber thickness obtained by extrapolating to zero the linearly decreasing part of the curve of penetration probability vs. absorber thickness. Data for Al in the  $T$  range up to about 10 MeV are available, and fit (to  $\sim \pm 10\%$ )  $R_p = AT[1 - B/(1 + CT)]$  mg cm<sup>-2</sup>, a form suggested in Ref. 22, with  $A = 0.55$  mg cm<sup>-2</sup> keV<sup>-1</sup>,  $B = 0.9841$ , and  $C = 0.0030$  keV<sup>-1</sup>. At this penetration depth, 90 - 95% of the incident electrons have stopped. Data for other elements are sketchy, but suggest that higher- $Z$  ( $\leq 50$ ) elements have  $1 \leq R_p/R_p(\text{Al}) \leq 1.4$  below  $\sim 10$  keV, and  $0.6 \leq R_p/R_p(\text{Al}) \leq 1$  above  $\sim 100$  keV. The "critical energy" (above which the energy loss due to

bremsstrahlung exceeds that due to ionization, and showering becomes important) is 400 MeV for hydrogen, 100 MeV for carbon, 25 MeV for copper, and 10 MeV for lead. The mean positron range may differ from the mean electron range by several percent. See Refs. 23 and 24. Electron energy deposition and penetration probability vs. range are discussed in Refs. 10, 25, and 26.

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## COMMONLY USED RADIOACTIVE SOURCES

Nuclide	$t_{1/2}$ (y)	Decay mode	Principal emissions E(keV)/Intensity (%)		
			$\alpha$	$\beta$	$\gamma$
$^{22}\text{Na}$	2.602	$\beta^+$ , EC		546/90	(511) 1275/100
$^{54}\text{Mn}$	0.854	EC			5.4/22 $X$ 835/100
$^{55}\text{Fe}$	2.68	EC			5.89/25 $X$ 6.49/3.3 $X$
$^{57}\text{Co}$	0.742	EC			14.4/10 122/86 136/11
$^{60}\text{Co}$	5.271	$\beta^-$		318/100	1173/100 1332/100
$^{68}\text{Ge}$ ( $^{68}\text{Ga}$ )	0.742	EC $\rightarrow \beta^+$		1899/88	(511) 9.4/44 $X$
$^{90}\text{Sr}$ ( $^{90}\text{Y}$ )	28.5	$\beta^- \rightarrow \beta^-$		546/100 2284/100	
$^{106}\text{Ru}$ ( $^{106}\text{Rh}$ )	1.020	$\beta^- \rightarrow \beta^-$		39/100 3541/79	512/21 622/10
$^{109}\text{Cd}$	1.267	EC		62/42 84/44 87/10	$e^-$ 22/83 $X$ 25/17 $X$ 88/3.7
$^{113}\text{Sn}$	0.315	EC		364/28 388/7	$e^-$ 24/79 $X$ 392/65
$^{137}\text{Cs}$	30.17	$\beta^-$		512/95 1173/5	662/85
$^{133}\text{Ba}$	10.54	EC		45/48 75/7	$e^-$ 81/33 303/18 356/60
$^{207}\text{Bi}$	32.2	EC		482/2 976/7 1048/2	$e^-$ 570/98 1064/75 1770/7
$^{228}\text{Th}$ ( $^{224}\text{Ra}$ , $^{220}\text{Rn}$ , $^{216}\text{Po}$ , $^{212}\text{Pb}$ , $^{212}\text{Bi}$ , $^{212}\text{Po}$ )	1.913	$6\alpha$ , $2\beta^-$	(5341 -8785)	334/85 1794/18 2246/48	239/45 583/30 2615/36
$^{241}\text{Am}$	432.2	$\alpha$		5443/13 5486/85	14/13 $X$ 18/20 $X$ 60/36 $X$
$^{244}\text{Cm}$	18.10	$\alpha$		5763/24 5805/76	14/4 $X$ 18/4 $X$
$^{252}\text{Cf}$	2.638	0.12 fission neutron, 0.6 $\gamma$ ( $\leq 1$ MeV)/decay of Cf			
Am/Be	432.2	$6 \times 10^{-5}$ neutron, $4 \times 10^{-5}$ $\gamma$ (4.43 MeV)/decay of Am			

EC denotes electron capture,  $X$  an atomic  $X$ -ray. Maximum  $\beta^\pm$  energies are listed, unless followed by  $e^-$  indicating monoenergetic conversion electrons. (511) indicates annihilation radiation, where intensity depends on the number of stopped positrons. In some cases, the  $\gamma$ -ray values are approximate weighted averages of two or more close-together lines. Daughter isotopes, the actual sources of some lines, are listed in parentheses where appropriate.

E. Browne and R.B. Firestone, *Table of Radioactive Isotopes* (John Wiley & Sons, New York, 1986).

Half-lives from J.K. Tuli, *Nuclear Wallet Cards* (1985), National Nuclear Data Center.

Energies and intensities from D.C. Kocher, *Radioactive Decay Data Tables* (1981), DOE/TIC-11026.

Neutrons from *Neutron Sources for Basic Physics and Applications* (Pergamon Press, 1983).

## RADIOACTIVITY &amp; RADIATION PROTECTION

The International Commission on Radiation Units and Measurements (ICRU) recommends the use of SI units. Therefore we list SI units first, followed by cgs (or other common) units in parentheses, where they differ.

**Unit of activity** = becquerel (curie):

$$1 \text{ Bq} = 1 \text{ disintegration/sec} [= 1/(3.7 \times 10^{10}) \text{ Ci}].$$

**Unit of exposure**, the quantity of  $X$ - or  $\gamma$ - radiation at a point in space integrated over time, in terms of charge of either sign produced by showering electrons in a small volume of air about the point:

$$= 1 \text{ coul/kg of air (roentgen: } 1 \text{ R} = 2.58 \times 10^{-4} \text{ coul/kg)}$$

$$= 1 \text{ esu/cm}^3 = 87.8 \text{ erg released energy per g of air}; \text{ implied in the definition is the assumption that the small test volume is embedded in a sufficiently large uniformly irradiated volume that the number of secondary electrons entering the volume equals the number leaving.}$$

**Unit of absorbed dose** = gray (rad):

$$1 \text{ Gy} = 1 \text{ joule/kg} (= 10^4 \text{ erg/g} = 10^2 \text{ rad})$$

$$= 6.24 \times 10^{12} \text{ MeV/kg deposited energy.}$$

**Unit of dose equivalent** (for biological damage) = sievert [=  $10^2$  rem (roentgen equivalent for man)]:

Dose equivalent in Sv = grays  $\times Q$ , where  $Q$  (quality factor) expresses long-term risk (primarily cancer and leukemia) from low-level chronic exposure; it depends upon the type of radiation and other factors. For  $\gamma$  rays and  $\beta$  particles,  $Q \cong 1$ ; for protons,  $Q \cong 1$  at  $\sim 10$  MeV, rising gradually to  $\cong 2$  at  $\sim 1$  GeV; for thermal neutrons,  $Q \cong 3$ ; for fast neutrons,  $Q$  ranges up to 10; and for  $\alpha$  particles and heavy ions (assuming internal deposition — skin and clothing are usually sufficient protection against external sources),  $Q \cong 20$ .

**Natural annual background**, all sources: Most world areas,

whole-body dose equivalent rate  $\cong (0.4\text{--}4)$  mSv (40-400 millirems). Can range up to 50 mSv (5 rems) in certain areas.

U.S. average  $\cong 0.8$  mSv. The lungs receive an additional  $\cong 0.1$  mSv ( $\cong 10$  mrem) from inhaled natural radioactivity, mostly radon and radon daughters (good to  $\cong$  factor of 2 in open areas; can range an order of magnitude higher in buildings and up to  $1000\times$  in poorly ventilated mines).

**Cosmic ray background** in counters (Earth's surface):

$\sim 10^4$ /min/m<sup>2</sup>/ster. For more accurate estimates and more details, see Cosmic Rays section.

**Fluxes** (per m<sup>2</sup>) to deposit one Gy in one kg of matter, assuming uniform irradiation:

$\cong$  (**charged particles**)  $6.24 \times 10^{12} / (dE/dx)$ , where  $dE/dx$  (MeV m<sup>2</sup>/kg), the energy loss per unit length, may be obtained (after conversion of units) from the Mean Range and Energy Loss figures.

$\cong 3.5 \times 10^{13}$  minimum-ionizing singly charged particles in carbon.

$\cong$  (**photons**)  $6.24 \times 10^{12} / [E(\text{MeV})/(\lambda f)(\text{m}^2/\text{kg})]$ , for photons of energy  $E$ , attenuation length  $\lambda$  (see Photon Attenuation Length figures), and fraction  $f \leq 1$  expressing the fraction of the photon's energy deposited in a small volume of thickness  $\ll \lambda$  but large enough to contain the secondary electrons.

$\cong 2 \times 10^{15}$  photons of 1 MeV energy on carbon.

(Quoted fluxes good to about a factor of 2 for all materials.)

**U.S. maximum permissible occupational whole-body dose:**

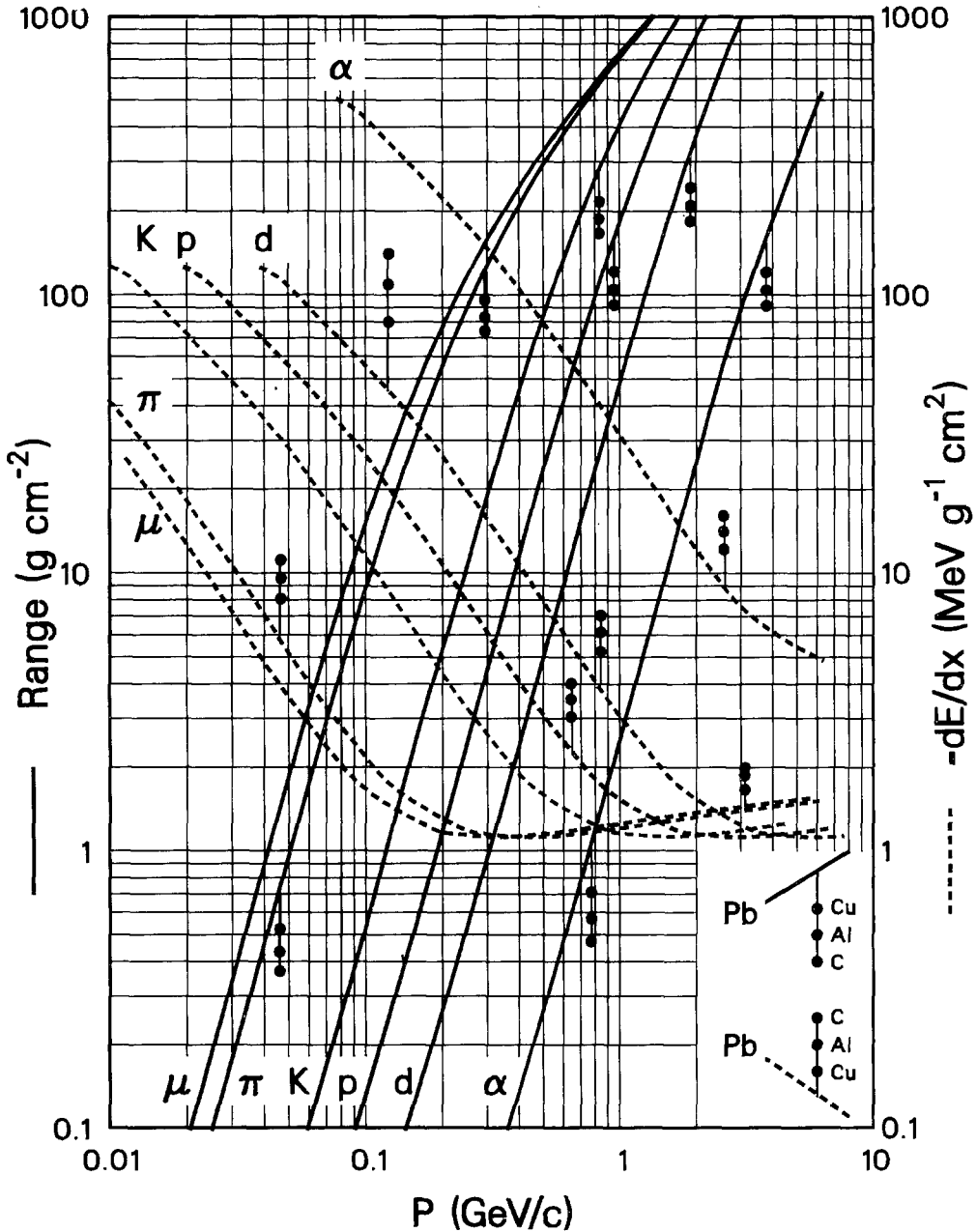
50 mSv/year (5 rem/year).

**Lethal dose:** Whole-body dose from penetrating ionizing radiation resulting in 50% mortality in 30 days (assuming no medical treatment), 2.5-3.0 Gy (250-300 rads) as measured internally on body longitudinal center line; surface dose varies due to variable body attenuation and may be a strong function of energy.

For a recent review, see E. Pochin, *Nuclear Radiation: Risks and Benefits* (Clarendon Press, Oxford, 1983).

## MEAN RANGE AND ENERGY LOSS

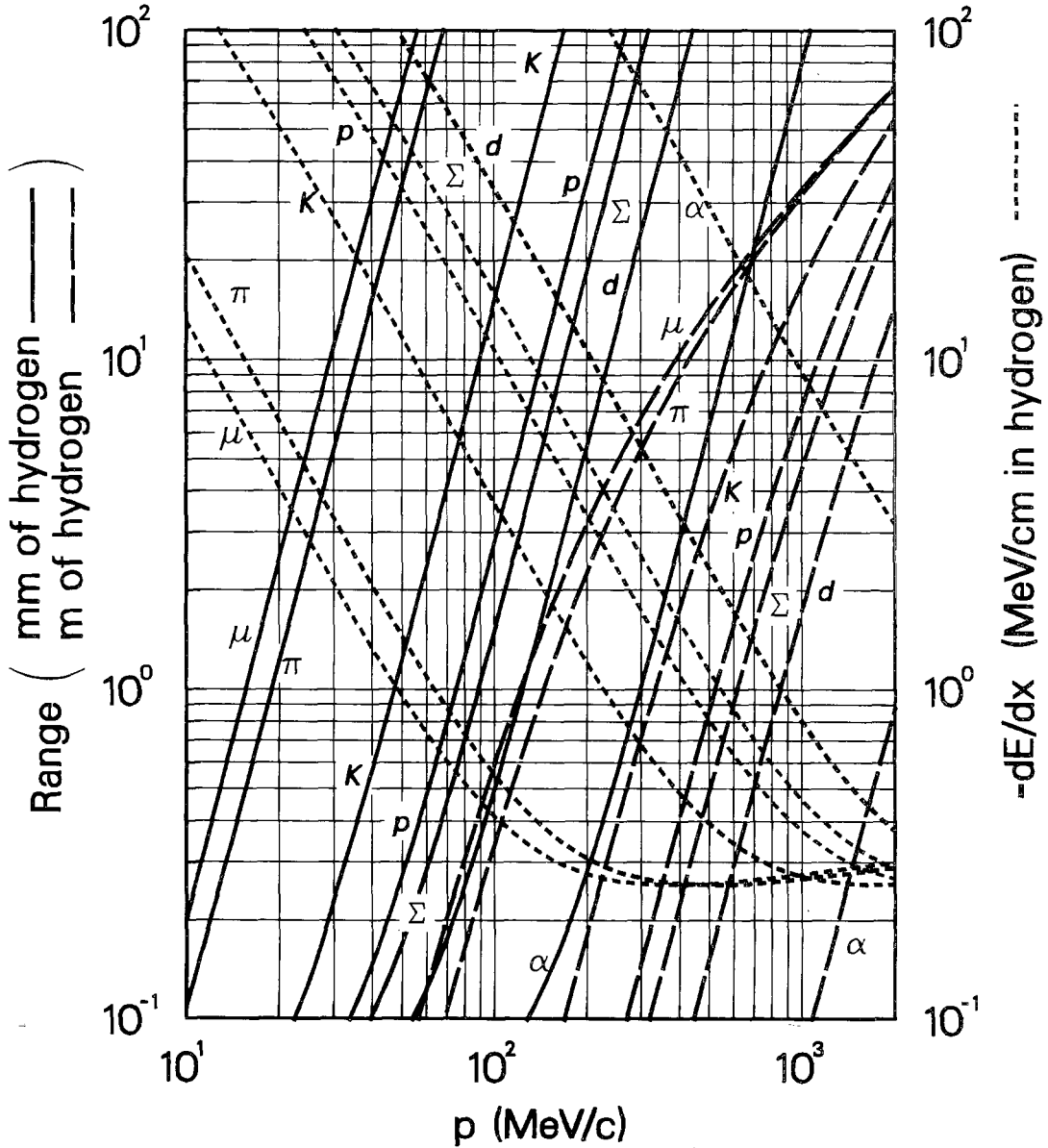
## Mean Range and Energy Loss in Lead, Copper, Aluminum, and Carbon



Mean range and energy loss due to ionization for the indicated particles in Pb, with scaling to Cu, Al, and C indicated, using Bethe-Bloch equation [See Sec. (1) of Passage of Particles Through Matter] with corrections. Calculated by M.J. Berger, using ionization potentials and density effect corrections as discussed in M.J. Berger and S.M. Seltzer, "Stopping Powers and Ranges of Electrons and Positrons," (2<sup>nd</sup> ed.), U.S. National Bureau of Standards Report NBSIR 82-2550-A (1982). The average ionization potentials ( $I$ ) assumed were: Pb (823 eV), Cu (322 eV), Al (166 eV), and C (78.0 eV). Figure indicates total path length; observed range may be smaller (by  $\sim 1\%$  -  $2\%$  in heavy elements) due to multiple scattering, primarily from small energy-loss collisions with nuclei. The functional forms have not been experimentally verified to better than roughly  $\pm 1\%$ . For higher energies refer to discussion by Cobb ["A Study of Some Electromagnetic Interactions of High Velocity Particles with Matter," University of Oxford Report HEP/T/55 (1973)] and by Turner ["Penetration of Charged Particles in Matter: A Symposium," National Academy of Sciences, Washington D.C. (1970), p. 48]. For lower energies both data and theory are not well understood. Scaling to other beam particles is, to a good approximation, described by the expression on the next page.

MEAN RANGE AND ENERGY LOSS (Cont'd)

Mean Range and Energy Loss in Liquid Hydrogen



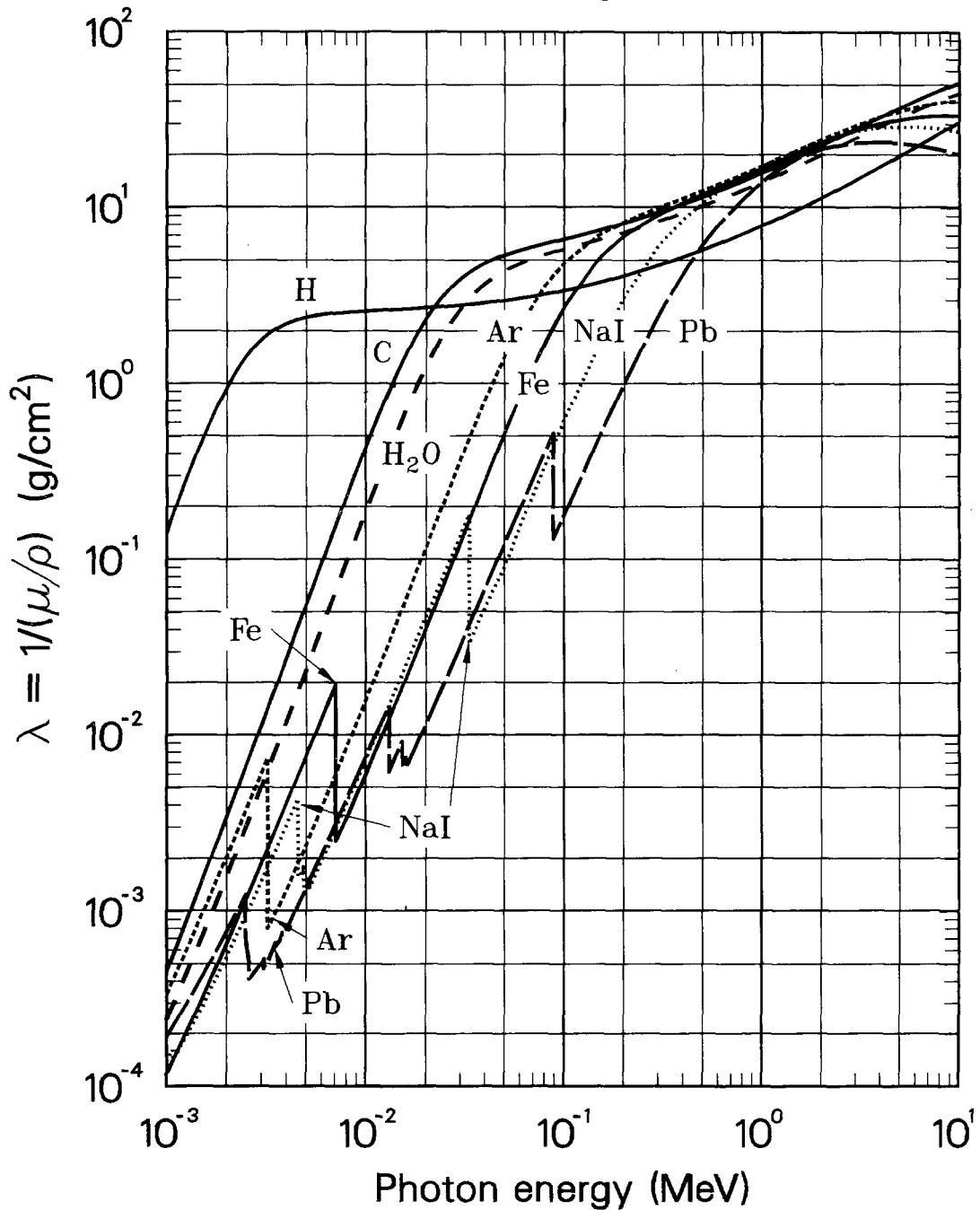
Range and energy loss in liquid hydrogen bubble chamber, based on Bethe-Bloch equation [See Sec. (1) of Passage of Particles Through Matter], using an average ionization potential for H<sub>2</sub> of  $I = 20.0$  eV, which is an approximate average of the experimental result of Garbincius and Hyman [Phys. Rev. A2, 1834 (1970)] and the theoretical result of Ford and Browne [Phys. Rev. A7, 418 (1973)]. Bubble chamber conditions are chosen to be those of Garbincius and Hyman: parahydrogen of density = 0.0625 g/cm<sup>3</sup> (note: range  $\propto 1/\text{density}$ ), with vapor-pressure 60.8 lb/in<sup>2</sup> (absolute) and temperature 26.2°K. The functional dependence of the Bethe-Bloch equation is not experimentally verified to better than about  $\pm 1\%$  over large momentum ranges. It should be noted that the number of bubbles per cm of a track in a bubble chamber is nearly proportional to  $1/\beta^2$ , not  $dE/dx$ . For the linear portions of the range curves,  $R \propto p^{3.6}$ . **Scaling law for particles of other mass or charge (except electrons):** for a given medium, the range  $R_b$  of any beam particle with mass  $M_b$ , charge  $z_b$ , and momentum  $p_b$  is given in terms of the range  $R_a$  of any other particle with mass  $M_a$ , charge  $z_a$ , and momentum  $p_a = p_b M_a / M_b$  (i.e., having the same velocity) by the expression:

$$R_b(M_b, z_b, p_b) = \left[ \frac{M_b / M_a}{z_b^2 / z_a^2} \right] R_a(M_a, z_a, p_a = p_b M_a / M_b).$$



## PHOTON AND ELECTRON ATTENUATION

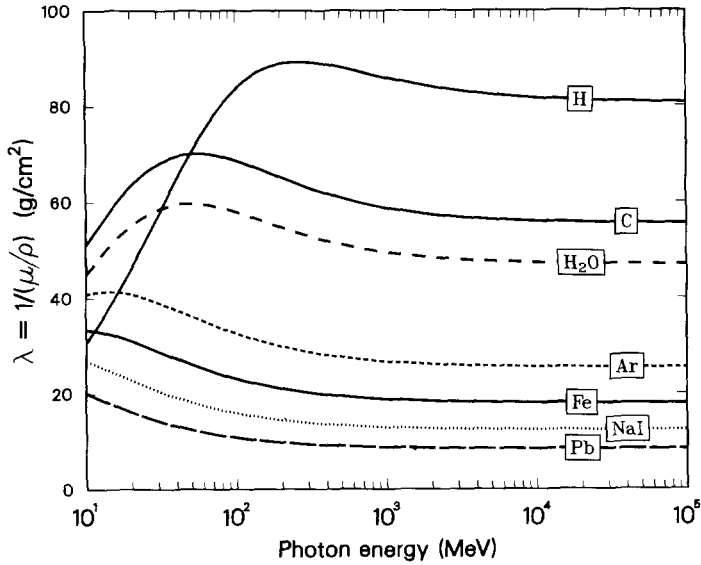
## Photon Attenuation Length



The photon mass attenuation length  $\lambda = 1/(\mu/\rho)$  (also known as mfp, mean free path) for various absorbers as a function of photon energy, where  $\mu$  is the mass attenuation coefficient. For a homogeneous medium of density  $\rho$ , the intensity  $I$  remaining after traversal of thickness  $t$  is given by the expression  $I = I_0 \exp(-t/\lambda)$ . The accuracy is a few percent. Interpolation to other  $Z$  should be done in the cross section  $\sigma = A/\lambda N_A \text{ cm}^2/\text{atom}$ , where  $A$  is the atomic weight of the absorber material in grams and  $N_A$  is the Avogadro number. For a chemical compound or mixture, use  $(1/\lambda)_{\text{eff}} \cong \sum w_i (1/\lambda)_i$ , accurate to a few percent, where  $w_i$  is the proportion by weight of the  $i^{\text{th}}$  constituent. See next page for high energy range. From Hubbell, Gimm, and Øverbø, *J. Phys. Chem. Ref. Data* 9, 1023 (1980). See also J.H. Hubbell, *Int. J. of Applied Rad. and Isotopes* 33, 1269 (1982). Data courtesy J.H. Hubbell.

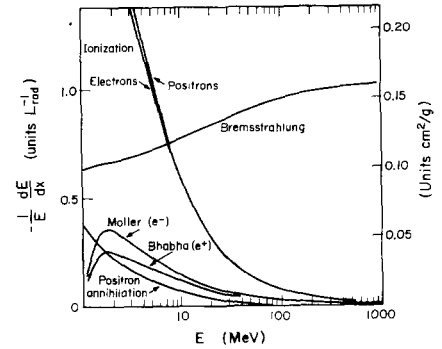
PHOTON AND ELECTRON ATTENUATION (Cont'd)

Photon Attenuation Length (High Energy)



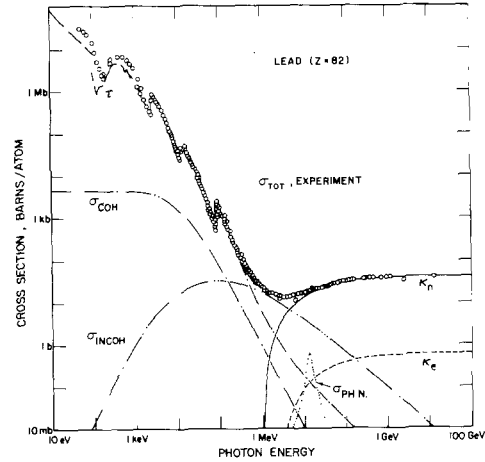
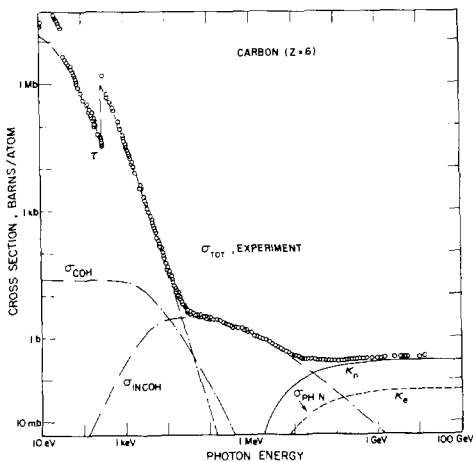
The photon mass attenuation length, high energy range (note that ordinate is linear scale). See caption on previous page for details.

Fractional Energy Loss for Electrons and Positrons in Lead



Fractional energy loss per radiation length in lead as a function of electron or positron energy. Electron (positron) scattering is considered as ionization when the energy loss per collision is below 0.255 MeV, and as Moller (Bhabha) scattering when it is above. Adapted from Fig. 3.2 from Messel and Crawford, *Electron-Photon Shower Distribution Function Tables for Lead, Copper, and Air Absorbers*, Pergamon Press, 1970. Messel and Crawford use  $L_r(\text{Pb}) = 5.82 \text{ g/cm}^2$ , but we have modified the figures to reflect the value given in the Table of Atomic and Nuclear Properties of Materials, namely  $L_r(\text{Pb}) = 6.4 \text{ g/cm}^2$ . The development of electron-photon cascades is approximately independent of absorber when the results are expressed in terms of inverse radiation lengths (i.e., scale on left of plot).

Contributions to Photon Cross Section in Carbon and Lead



Photon total cross sections as a function of energy in carbon and lead, showing the contributions of different processes.

- $\tau$  = Atomic photo-effect (electron ejection, photon absorption)
- $\sigma_{\text{COH}}$  = Coherent scattering (Rayleigh scattering — atom neither ionized nor excited)
- $\sigma_{\text{INCOH}}$  = Incoherent scattering (Compton scattering off an electron)
- $K_n$  = Pair production, nuclear field
- $K_e$  = Pair production, electron field
- $\sigma_{\text{PH.N.}}$  = Photonuclear absorption (nuclear absorption, usually followed by emission of a neutron or other particle)

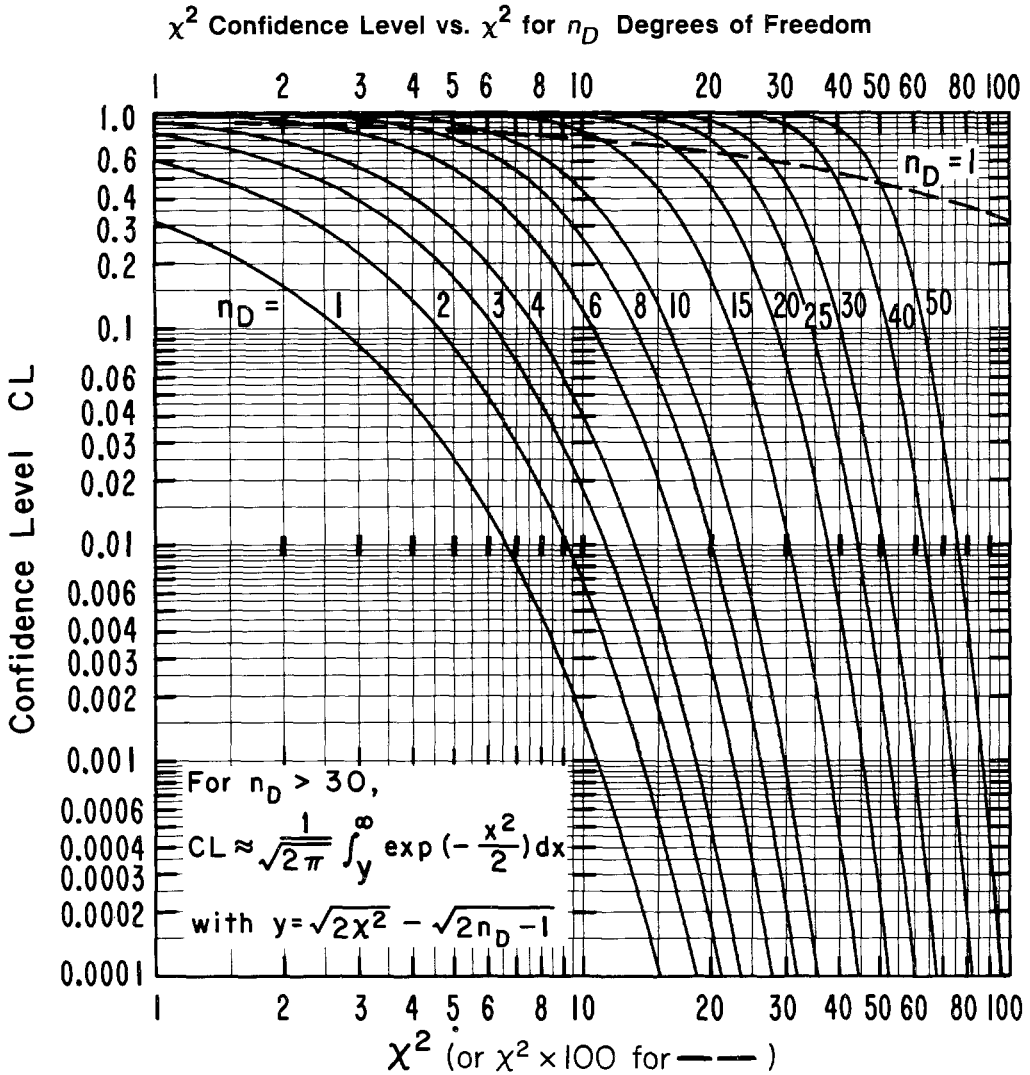
From Hubbell, Gimm, and Øverbø, *J. Phys. Chem. Ref. Data* 9, 1023 (1980). Figures courtesy J.H. Hubbell.

## PROBABILITY AND STATISTICS

### A. PROBABILITY DISTRIBUTIONS AND CONFIDENCE LEVELS

We give here properties of four probability distributions commonly encountered in high energy physics: normal (or Gaussian), chi-squared ( $\chi^2$ ), Poisson, and binomial. We warn the reader that

there is no universal convention for the term "confidence level"; thus, explicit definitions that correspond to common usage are given below. We explain below how confidence levels for the first three distributions may be extracted from the following figure.



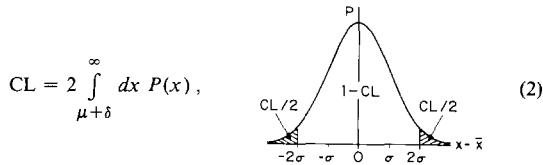
**PROBABILITY AND STATISTICS (Cont'd)**

**A.1 Normal distribution**

The normal distribution with mean  $\bar{x} = \mu$  and standard deviation  $\sigma$  (variance  $\sigma^2$ ) is:

$$P(x)dx = \frac{1}{\sigma\sqrt{2\pi}} e^{-(x-\mu)^2/2\sigma^2} dx \quad (1)$$

The confidence level associated with an observed deviation  $\pm\delta$  from the mean is the probability that  $|x-\mu| > \delta$ , i.e.,



since the distribution is symmetric about  $\mu$ . The small figure in Eq. (2) is drawn with  $\delta = 2\sigma$ . CL is given by the ordinate of the  $n_D = 1$  curve in the large figure at  $\chi^2 = (\delta/\sigma)^2$ . The odds against exceeding  $\delta$  in an observation,  $(1-CL)/CL$ , for  $\delta = 1\sigma$  are 2.15:1;  $2\sigma$ , 21:1;  $3\sigma$ , 370:1;  $4\sigma$ , 16,000:1;  $5\sigma$ , 1,700,000:1. Other measures of the width are related to  $\sigma$  as follows: probable error (CL = 0.5) =  $0.67\sigma$ ; mean absolute deviation =  $0.80\sigma$ ; RMS deviation =  $\sigma$ ; half width at half maximum =  $1.18\sigma$ .

The probability that a random observation of  $x$  will have  $|x - \mu| < \delta$  is  $(1 - CL)$ . Therefore, for a given measurement  $x_0$ , the interval  $x_0 - \delta$  to  $x_0 + \delta$  can be expected to contain  $\mu$  a fraction  $(1 - CL)$  of the time. Refer to Sec. B.1 below to combine multiple observations and treat them as a single measurement with appropriate variance. This interval is a  $(1 - CL)$  confidence interval for an unknown  $\mu$  when  $\sigma$  is known. Frequently used choices for CL and  $\delta$  are:

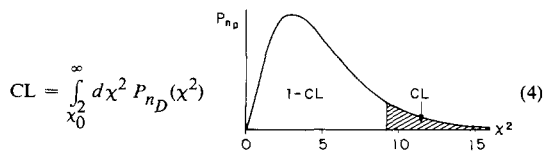
CL(%)	$\delta$	CL(%)	$\delta$
31.7	$1\sigma$	10	$1.64\sigma$
4.6	$2\sigma$	5	$1.96\sigma$
0.3	$3\sigma$	1	$2.58\sigma$
$6 \times 10^{-3}$	$4\sigma$	0.1	$3.29\sigma$
$6 \times 10^{-5}$	$5\sigma$	0.01	$3.89\sigma$

**A.2  $\chi^2$  distribution**

The  $\chi^2$  distribution for  $n_D$  degrees of freedom is:

$$P_{n_D}(\chi^2)d\chi^2 = \frac{1}{2^h\Gamma(h)} (\chi^2)^{h-1} e^{-\chi^2/2} d\chi^2 \quad (\chi^2 \geq 0) \quad (3)$$

where  $h$  (for "half") =  $n_D/2$ . The mean and variance are  $n_D$  and  $2n_D$  respectively. In evaluating Eq. (3) one may use Stirling's approximation:  $\Gamma(h) \approx 2.507 e^{-h} h^{(h-1/2)}(1 + 0.0833/h)$ , which is accurate to  $\pm 0.1\%$  for all  $h \geq 1/2$ . The confidence level associated with a given value of  $n_D$  and an observed value of  $\chi_0^2$  is the probability of the  $\chi^2$  in a random measurement exceeding the observed value, i.e.,



The small figure in Eq. (4) is drawn with  $n_D = 5$  and CL = 10%. CL is plotted as a function of  $\chi^2$  for several values of  $n_D$  in the large figure. For large  $n_D$ ,  $\chi^2$  becomes normally distributed about  $n_D$ . Thus,

$$y_1 = (\chi^2 - n_D)/\sqrt{2n_D} \quad (5)$$

becomes normally distributed with unit standard deviation and mean zero. A better approximation is that  $\chi$ , not  $\chi^2$ , becomes normally distributed; specifically,

$$y_2 = \sqrt{2}\chi^2 - \sqrt{2n_D-1} \quad (6)$$

approaches normality with unit standard deviation and mean zero. For small CL's in particular,  $y_2$  is much more accurate than  $y_1$ . Thus, for  $n_D = 50$  and  $\chi^2 = 80$ , the true CL = 0.45%, but  $y_1$  is 3.0 corresponding to a CL of 0.13%, while  $y_2$  is 2.7 corresponding to a CL of 0.35%.

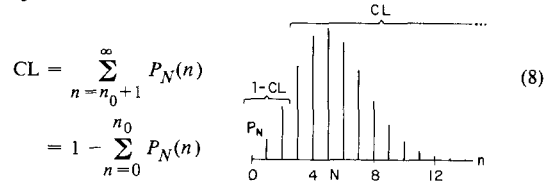
**A.3 Poisson distribution**

The Poisson distribution with mean  $\mu$  is:

$$P_\mu(n) = \frac{e^{-\mu}\mu^n}{n!} \quad (n = 0, 1, 2, \dots) \quad (7)$$

The observed result of a Poisson process is a non-negative integer  $n$ ; the parameter  $\mu$ , which is equal to both the mean and the variance, is any non-negative real number. When  $\mu$  is large ( $\geq 7$  or 8), one may often usefully approximate the distribution of  $n$  by a normal of mean  $\mu$  and variance  $\sigma^2 = \mu$ , as though  $n$  were a continuous variable. Two or more Poisson processes (e.g., signal plus background, with parameters  $\mu_S$  and  $\mu_B$  respectively) which independently contribute amounts  $n_S$  and  $n_B$  to a given measurement will produce an observed number  $n = n_S + n_B$  which is distributed according to a new Poisson distribution with parameter  $\mu = \mu_S + \mu_B$ .

When an observed value  $n_0$  of  $n$  is small, one is often interested in obtaining an upper limit  $N$  (not necessarily an integer) to the parameter  $\mu$  such that one may be, e.g., 90% or 95% confident that the true value of  $\mu$  lies below  $N$ . We define this  $N$  to be that value of  $\mu$  such that it would be exactly 90% (or 95%) probable that a random observation of  $n'$  would lie above the actual observation  $n_0$ . That is, if CL = 0.90 or 0.95,



The small figure in Eq. (8) illustrates the case with  $n_0 = 2$  and CL = 0.90, for which it may be shown that  $N = 5.3$ . We can obtain  $N$  for any desired  $n_0$  and CL from the large  $\chi^2$  figure, because of a particular relation between Poisson and  $\chi^2$  confidence levels: at the level  $1-CL$  on the curve  $n_D = 2(n_0 + 1)$ , the desired  $N$  is  $1/2$  the  $\chi^2$  value of that point. Some useful values are:

Poisson upper limits  $N$  for  $n_0$  observed events

$n_0$	CL=	CL=	$n_0$	CL=	CL=
	90%	95%		90%	95%
0	2.30	3.00	6	10.53	11.84
1	3.89	4.74	7	11.77	13.15
2	5.32	6.30	8	13.00	14.44
3	6.68	7.75	9	14.21	15.71
4	7.99	9.15	10	15.41	16.96
5	9.27	10.51			

Tables of confidence levels for all three of these distributions, the relation between Poisson and  $\chi^2$  confidence levels, and numerous other useful tables and relations may be found in Ref. 1.

## PROBABILITY AND STATISTICS (Cont'd)

### A.4 Binomial distribution

Any process with exactly two possible outcomes is a *Bernoulli* process. If the process is repeated  $n$  times independently, and if the probability of obtaining a certain outcome (a "success") in each trial is  $p$ , then the probability of obtaining exactly  $r$  successes is given by the binomial distribution:

$$P_{n,p} = \binom{n}{r} p^r q^{n-r} = \frac{n!}{r!(n-r)!} p^r q^{n-r}, \quad (9)$$

where  $q = 1 - p$  and the order in which the successes and failures come is assumed irrelevant. The mean of  $r$  is equal to  $np$  and the variance is  $npq$ . If  $r$  successes are observed in  $n_r$  Bernoulli trials with probability  $p$  of success, and if  $s$  successes are observed in  $n_s$  similar trials, then  $t = r + s$  is also binomial with  $n_t = n_r + n_s$ .

## B. STATISTICS

The probability density functions we have shown enable us to predict the frequency with which a random observation of the variable  $x$  (or  $n$  or  $\chi^2$ ) will lie in a certain region, once the parameters  $\mu$  and  $\sigma$  (or  $n_D$ ) are known. In statistics (*parametric* statistics) we have the opposite problem of estimating the parameters of the distribution from a set of actual observations.

Suppose one is presented with  $N$  independent data,  $y_n \pm \sigma_n$ , and it is desired to make some *inference* about the "true" value of the quantity represented by these data. For this purpose we interpret each datum  $y_n$  as a single sample point drawn randomly from a distribution having true mean  $\mu_n$  (which we wish to estimate) and variance  $\sigma_n^2$ . We do not require that they be normally distributed. (Identification of the true  $\sigma_n$  with the  $\sigma_n$  datum is often an *approximation* which may become seriously inaccurate when  $\sigma_n$  is an appreciable fraction of  $y_n$ .) Some methods of estimation commonly used in high energy physics are given below; see Ref. 2 for numerous applications. Section B.1 deals with the case in which all  $\mu_n$  are the same, i.e., several different measurements of the same quantity. Sec. B.2 deals with the case in which the true value of  $y_n$ ,  $\mu_n$ , is a variable dependent upon another variable  $x_n$ ; i.e.,  $\mu_n = f(x_n)$ . An example is the energy deposited in a calorimeter as a function of the energy of the incident particle,  $y_n = f(E_n)$ .

### B.1 Single mean and variance estimates

(1) If the  $y_n$  represent a set of values all supposedly drawn from a single distribution with mean  $\mu$  and variance  $\sigma^2$  (i.e., the  $\sigma_n$  are all the same, but their common value is unknown), then

$$\hat{y} = \frac{1}{N} \sum_{n=1}^N y_n \quad \text{and} \quad (10)$$

$$\hat{\sigma}^2 = \frac{1}{N-1} \sum_{n=1}^N (y_n - \hat{y})^2 = \frac{N}{N-1} \left[ \langle (y^2) \rangle - (\hat{y})^2 \right] \quad (11)$$

are unbiased *estimates* of  $\mu$  and  $\sigma^2$ ; the angular brackets denote an average over the data. The variance of  $\hat{y}$  is  $\sigma^2/N$ . If the parent distribution is normal and  $N$  is large, the variance of  $\hat{\sigma}^2$  is  $2\sigma^4/N$ .

(2) If the  $y_n$  are independent estimates of the same  $\mu$ , and the  $\sigma_n$  are known, then the weighted average

$$\hat{y} = \frac{1}{w} \sum_n w_n y_n, \quad (12)$$

where  $w_n = 1/\sigma_n^2$  and  $w = \sum w_n$ , is an appropriate unbiased estimate of  $\mu$ . This choice of weighting factors in Eq. (12) minimizes the variance of the estimate, which is  $1/w$ .

### B.2 Linear least-squares fit

We wish to determine the best fit of unbiased data  $y_n$ , measured at  $N$  points  $x_n$  (assumed known with negligible error), to the form  $y(x) = \sum a_i f_i(x)$ , where the  $f_i$  are known, linearly independent functions (e.g., Legendre polynomials) which are single-valued over the allowed range of  $x$ , and the sum runs from 1 to  $k$ . We require  $k \leq N$ , and at least  $k$  of the  $x_n$  must be distinct. We wish to estimate the linear coefficients  $a_i$ .

The Method of Least Squares assumes that each measured  $y_n$  is equal to this sum plus an error  $\epsilon_n$ . If the distribution of  $\epsilon_n$  has an expectation value of zero (unbiased) and has a finite, known variance which is fixed (does not depend on the parameters of the fit), then the estimates of  $a_i$  obtained by minimizing the sum of squares

$$SS = \sum_n [y_n - \sum a_i f_i(x_n)]^2 / \sigma_n^2$$

will be unbiased and have the smallest possible variance of all linear unbiased estimates (Gauss-Markov Theorem). If the point errors  $\epsilon_n$  are also Gaussian, then

$$SS = \sum_n \epsilon_n^2 / \sigma_n^2$$

will be distributed, over a large number of similar experiments, as a  $\chi^2$  with  $N - k$  degrees of freedom. If, more generally, the measured  $y_n$ 's are not independent, then the set of  $\sigma_n^2$ 's must be replaced by the  $N \times N$  variance matrix  $V_y$ . Then, in matrix notation, if  $H$  is the  $N \times k$  matrix with element  $H_{ni} = f_i(x_n)$ , the vector (all vectors are column vectors) of solutions  $\hat{a}$ , with  $k$  elements, is given by

$$\hat{a} = (H^T V_y^{-1} H)^{-1} H^T V_y^{-1} \bar{y} \equiv D \bar{y}, \quad (13)$$

where  $\bar{y}$  is the  $N$ -element vector of measured  $y_n$ 's. In terms of the  $k \times N$  matrix  $D$ , the standard variance matrix for the  $\hat{a}$  is estimated by

$$V_{\hat{a}} = D V_y D^T. \quad (14a)$$

If the measured  $y_n$ 's are independent,  $V_y$  is diagonal with  $n$ <sup>th</sup> element  $\sigma_n^2$  and

$$\left[ V_{\hat{a}}^{-1} \right]_{ij} = \sum_n f_i(x_n) f_j(x_n) / \sigma_n^2. \quad (14b)$$

The estimated variance of an interpolated or extrapolated value of  $y$  at a point  $x$ ,  $\hat{y} = \sum \hat{a}_i f_i(x)$ , is

$$\sigma^2(\hat{y}) = \sum_{ij} (V_{\hat{a}})_{ij} f_i(x) f_j(x). \quad (15)$$

The same results may be obtained by numerical or other techniques from the sum of squares,  $SS$ , directly, if we have a reasonable first guess  $\bar{a}_0$  for the solution vector:

$$\hat{a} = \bar{a}_0 - \left( \frac{\partial^2 SS}{\partial a^2} \right)_{\bar{a}_0}^{-1} \cdot \frac{\partial SS}{\partial a} \Big|_{\bar{a}_0} \quad (16a)$$

and

$$V_{\hat{a}} = 2 \left( \frac{\partial^2 SS}{\partial a^2} \right)_{\hat{a}}^{-1}, \quad (16b)$$

where  $\partial SS / \partial a$  is a  $k$ -element vector whose  $i$ <sup>th</sup> element is  $\partial SS / \partial a_i$ ,  $\partial^2 SS / \partial a^2$  is a  $k \times k$  matrix with  $ij$ <sup>th</sup> element  $\partial^2 SS / (\partial a_i \partial a_j)$ , and all derivatives are to be evaluated at the points indicated. If  $SS$  is a  $\chi^2$ , the second-derivative matrix is independent of  $\bar{a}$ , and the shape of the  $\chi^2$  as a function of  $\bar{a}$  is a parabola; otherwise one may need to iterate Eq. (16a) to arrive at a solution (Newton-Raphson method).

## PROBABILITY AND STATISTICS (Cont'd)

Note that the errors on the solution  $\hat{a}$  are independent of the value of SS or  $\chi^2$  at minimum – they depend only upon the shape about the minimum. Eq. (16b) implies that one-standard-deviation limits on the elements of  $\hat{a}$  are given by the set of  $\bar{a}'$  such that

$$SS(\bar{a}') = SS_{\min} + 1. \quad (17)$$

This equation, which defines a contour in  $\bar{a}$ -space, is often convenient for estimating errors in applications of least-squares techniques to nonlinear cases, where the second derivative [Eq. (16b)] may be a rapidly varying function of  $\bar{a}$ . In general, contours at  $m$  standard deviations may be found by replacing the 1 in Eq. (17) by  $m^2$ . If the problem is highly nonlinear, all such contours are at best only approximations to confidence regions with some given probability of covering the true value of  $\bar{a}$ .

If SS is a  $\chi^2$  with  $N - k$  degrees of freedom, its value at minimum may be used as a test of goodness-of-fit. If the model for  $y$  is correct and the measurements  $y_n$  are distributed as Gaussians with known variance matrix  $V_y$ , then the large figure above will give the probability that a similar experiment would be expected to yield the observed  $\chi^2_{\min}$  or larger. If this probability is smaller than some agreed-upon value, one should question the assumptions about the data or the validity of the model.

### Example — Straight-Line Fit

For the case of a *straight line fit*,  $y(x) = a + bx$ , one obtains, for independent measurements  $y_n$ , the following estimates of  $a$  and  $b$ ,

$$\hat{a} = (S_y S_{xx} - S_x S_{xy}) / D, \quad (18)$$

$$\hat{b} = (S_1 S_{xy} - S_x S_y) / D,$$

where

$$S_1, S_x, S_y, S_{xx}, S_{xy} = \sum (1, x_n, y_n, x_n^2, x_n y_n) / \sigma_n^2, \quad (19)$$

respectively, and

$$D = S_1 S_{xx} - S_x^2.$$

The covariance matrix of the fitted parameters is:

$$\begin{bmatrix} V_{aa} & V_{ab} \\ V_{ab} & V_{bb} \end{bmatrix} = \frac{1}{D} \begin{bmatrix} S_{xx} & -S_x \\ -S_x & S_1 \end{bmatrix}. \quad (20)$$

The estimated variance of an interpolated or extrapolated value of  $y$  at point  $x$  is:

$$(\hat{y} - y_{\text{true}})^2 |_{\text{est}} = \frac{1}{S_1} + \frac{S_1}{D} \left[ x - \frac{S_x}{S_1} \right]^2. \quad (21)$$

### B.3 Limits in the Case of Bounded Physical Regions

In certain statistical problems the true value of the parameter to be estimated,  $\mu$ , is constrained to lie within a bounded *physical region* (e.g., the mass of a neutrino is bounded from below by 0). However, due to random measurement error, real measured values may or may not occur inside the physical region. For this case no completely satisfactory approach exists, but here we suggest a technique for obtaining approximate limits within the physical region at specified confidence levels.

We assume a measurement  $x$ , which represents one observation (or the result of combining multiple measurements as in Sec. B.1) from a normal of true (but unknown) mean  $\mu$  and known, fixed, variance  $\sigma^2$ . We estimate  $\mu$  by  $\hat{\mu} = x$  and attempt to construct a confidence interval for  $\mu$  from the resultant Gaussian, as in A.1. If

$\hat{\mu}$  or a portion of the accepted region [unshaded in the figure with Eq. (2)] lies in the unphysical region, the result, while statistically perfectly correct as stated, is physically unsatisfactory.

If we assume  $\mu$  is bounded from below by  $\mu_{\min}$  (the argument for  $\mu$  bounded from above is similar), we may estimate an upper limit for  $\mu$  at the CL (e.g., 90% or 95%) level by the following procedure: 1) *renormalize* the normal probability distribution for  $x$  such that the integral of Eq. (1) with  $\mu = \hat{\mu}$  over  $x$  from  $\mu_{\min}$  to infinity (i.e., over the physical region) is equal to 1.0; 2) find the value  $\mu_1$  such that the integral over  $x$  of the renormalized distribution from  $\mu_{\min}$  to  $\mu_1$  is equal to the desired value of CL; 3) set  $\mu_1$  to be the desired upper limit with confidence CL. In fact, it can be shown that this is *conservative*, in the sense that the probability that this interval actually covers the true value of  $\mu$  is  $\geq$  CL.

For  $\mu - \mu_{\min} \gg \sigma$ , this technique, which may be applied for any measured  $x$  (physical or unphysical), converges smoothly to that of Sec. A.1, since  $x$  is then effectively confined to the physical region.

One should exercise caution for values of  $x$  which lie many standard deviations outside the physical region. It may be that the particular probability model (normal with variance  $\sigma^2$ ) may not be a correct description of the measurement process (e.g., the true variance may have unanticipated components and be  $> \sigma^2$ , or there may be a bias), in which case confidence levels of this sort will not be correct.

## C. ERROR PROPAGATION

Suppose one wishes to calculate the value and error of a function of some other quantities with errors, e.g., in a Monte Carlo program. Let  $\{y\}$  be a set of random variables with observed means  $\{\bar{y}\}$  and covariance matrix  $V$ . Then the mean and variance of a function  $f$  of these variables are approximately (to second order in  $\{y - \bar{y}\}$ ):

$$\bar{f} \approx f(\{\bar{y}\}) + \frac{1}{2} \sum_{mn} V_{mn} \left[ \frac{\partial^2 f}{\partial y_m \partial y_n} \right]_{\{y\} = \{\bar{y}\}}, \quad (22)$$

$$\overline{(f - \bar{f})^2} \approx \sum_{mn} V_{mn} \left[ \frac{\partial f}{\partial y_m} \right]_{\{y\} = \{\bar{y}\}} \left[ \frac{\partial f}{\partial y_n} \right]_{\{y\} = \{\bar{y}\}}. \quad (23)$$

E.g., the mean and variance of a function of a *single variable* with mean  $\bar{y}$  and variance  $\sigma^2$  are

$$\bar{f} \approx f(\bar{y}) + \frac{1}{2} \sigma^2 f''(\bar{y}), \quad (24)$$

$$\overline{(f - \bar{f})^2} \approx \sigma^2 f'(\bar{y})^2. \quad (25)$$

If, as is often the case, the observed mean  $\bar{y}_n$  differs from the true mean  $\mu_n$  by an amount of order  $\sqrt{V_{nn}}$ , then the second-order terms in Eqs. (22) and (24) may be small compared with the first-order errors introduced by the substitution.

1. M. Abramowitz and I. Stegun, eds., *Handbook of Mathematical Functions* (Dover, New York, 1972).
2. W.T. Eadie, D. Drijard, F.E. James, M. Roos, and B. Sadoulet, *Statistical Methods in Experimental Physics* (North Holland, Amsterdam and London, 1971); S.L. Meyer, *Data Analysis for Scientists and Engineers* (John Wiley and Sons, Inc., New York, 1975); A.G. Frodesen, O. Skjeggstad, and H. Tøfte, *Probability and Statistics in Particle Physics* (Universitetsforlaget, Oslo, Norway, 1979).

**ELECTROMAGNETIC RELATIONS**

Quantity	Gaussian CGS	MKSA
Units and conversions: Charge: Potential: Magnetic field: Electron charge:	2.99792×10 <sup>9</sup> esu (1/299.792) statvolt = (1/299.792) erg/esu 10 <sup>4</sup> gauss = 10 <sup>4</sup> dyne/esu e = 4.803 242×10 <sup>-10</sup> esu	= 1 coul = 1 amp-sec = 1 volt = 1 joule/coul = 1 tesla = 1 nt/amp-m = 1.602 189 2×10 <sup>-19</sup> coul
Lorentz force:	$\mathbf{F} = q(\mathbf{E} + \frac{\mathbf{v}}{c} \times \mathbf{B})$	$\mathbf{F} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B})$
Maxwell equations:	$\nabla \cdot \mathbf{D} = 4\pi\rho$ $\nabla \times \mathbf{E} = -\frac{1}{c} \frac{\partial \mathbf{B}}{\partial t}$ $\nabla \cdot \mathbf{B} = 0$ $\nabla \times \mathbf{H} = \frac{4\pi\mathbf{j}}{c} + \frac{1}{c} \frac{\partial \mathbf{D}}{\partial t}$	$\nabla \cdot \mathbf{D} = \rho$ $\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$ $\nabla \cdot \mathbf{B} = 0$ $\nabla \times \mathbf{H} = \mathbf{j} + \frac{\partial \mathbf{D}}{\partial t}$
Materials: Permittivity of free space: Permeability of free space:	$\mathbf{D} = \epsilon\mathbf{E}, \mathbf{B} = \mu\mathbf{H}$ $\epsilon_{\text{vac}} = 1$ $\mu_{\text{vac}} = 1$	$\mathbf{D} = \epsilon\mathbf{E}, \mathbf{B} = \mu\mathbf{H}$ $\epsilon_{\text{vac}} = \epsilon_0$ $\mu_{\text{vac}} = \mu_0$
Fields:	$\mathbf{E} = -\nabla V - \frac{1}{c} \frac{\partial \mathbf{A}}{\partial t}$ $\mathbf{B} = \nabla \times \mathbf{A}$	$\mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t}$ $\mathbf{B} = \nabla \times \mathbf{A}$
Static potentials: (coulomb gauge)	$V = \sum_{\text{charges}} \frac{q}{r}$ $\mathbf{A} = \frac{1}{c} \sum_{\text{currents}} \frac{\mathbf{I}}{r}$	$V = \frac{1}{4\pi\epsilon_0} \sum_{\text{charges}} \frac{q}{r}$ $\mathbf{A} = \frac{\mu_0}{4\pi} \sum_{\text{currents}} \frac{\mathbf{I}}{r}$
Relativistic transformations: (v is the velocity of primed system as seen in unprimed system)	$\mathbf{E}'_{\parallel} = \mathbf{E}_{\parallel}$ $\mathbf{E}'_{\perp} = \gamma(\mathbf{E}_{\perp} + \frac{1}{c} \mathbf{v} \times \mathbf{B})$ $\mathbf{B}'_{\parallel} = \mathbf{B}_{\parallel}$ $\mathbf{B}'_{\perp} = \gamma(\mathbf{B}_{\perp} - \frac{1}{c} \mathbf{v} \times \mathbf{E})$	$\mathbf{E}'_{\parallel} = \mathbf{E}_{\parallel}$ $\mathbf{E}'_{\perp} = \gamma(\mathbf{E}_{\perp} + \mathbf{v} \times \mathbf{B})$ $\mathbf{B}'_{\parallel} = \mathbf{B}_{\parallel}$ $\mathbf{B}'_{\perp} = \gamma(\mathbf{B}_{\perp} - \frac{1}{c^2} \mathbf{v} \times \mathbf{E})$
$4\pi\epsilon_0 = \frac{1}{c^2} \frac{10^7 \text{ coul}^2}{\text{nt sec}^2} = \frac{1}{8.987 55} \times 10^{-9} \frac{\text{coul}^2}{\text{nt m}^2}$ $\frac{\mu_0}{4\pi} = 10^{-7} \frac{\text{nt sec}^2}{\text{coul}^2}; \quad c = 2.997 924 58 \times 10^8 \text{ m sec}^{-1}$		

**Impedances (MKSA)**

$\rho$  = resistivity in 10<sup>-8</sup> Ωm:  
 ~ 1.7 for Cu ~ 5.5 for W  
 ~ 2.4 for Au ~ 73 for SS 304  
 ~ 2.8 for Al ~ 100 for Nichrome  
 (Al alloys may have double this value.)

For alternating currents, instantaneous current  $I$ , voltage  $V$ , angular frequency  $\omega$ :

$$V = V_0 e^{i\omega t} = ZI.$$

Impedance of self-inductance  $L$ :  $Z = i\omega L$ .

Impedance of capacitance  $C$ :  $Z = 1/i\omega C$ .

Impedance of free space:  $Z = \sqrt{\mu_0/\epsilon_0} = 376.7 \Omega$ .

Impedance per unit length of a flat conductor of width  $w$  (high frequency,  $\nu$ ):

$$Z = \frac{(1+i)\rho}{w\delta}, \text{ where } \delta = \text{effective skin depth};$$

$$\delta = \sqrt{\frac{\rho}{\pi\nu\mu}} \approx \frac{6.6 \text{ cm}}{\sqrt{\nu(\text{sec}^{-1})}} \text{ for Cu.}$$

**Capacitance  $\hat{C}$  and inductance  $\hat{L}$  per unit length (MKSA)**

Flat rectangular plates of width  $w$ , separated by  $d \ll w$ :

$$\hat{C} = \epsilon \frac{w}{d}; \quad \hat{L} = \mu \frac{d}{w};$$

$$\frac{\epsilon}{\epsilon_0} = 2 \text{ to } 6 \text{ for plastics; } 4 \text{ to } 8 \text{ for porcelain, glasses.}$$

Coaxial cable of inner radius  $r_1$ , outer radius  $r_2$ :

$$\hat{C} = \frac{2\pi\epsilon}{\ln(r_2/r_1)}; \quad \hat{L} = \frac{\mu}{2\pi} \ln(r_2/r_1).$$

Transmission lines (no loss):

$$\text{Impedance: } Z = \sqrt{\hat{L}/\hat{C}}.$$

$$\text{Velocity: } v = 1/\sqrt{\hat{L}\hat{C}} = 1/\sqrt{\mu\epsilon}.$$

**Motion of charged particles in a uniform, static, magnetic field**

The path of motion of a charged particle of momentum  $p$  is a helix of constant radius  $R$  and constant pitch angle  $\lambda$ , with the axis of the helix along  $\mathbf{B}$ :

$$p(\text{GeV}/c)\cos\lambda = 0.29979 q B(\text{tesla}) R(\text{m}),$$

where the charge  $q$  is in units of the electronic charge. The angular velocity about the axis of the helix is

$$\omega(\text{rad sec}^{-1}) = 8.98755 \times 10^7 q B(\text{tesla})/E(\text{GeV}),$$

where  $E$  is the energy of the particle.

## ELECTROMAGNETIC RELATIONS (Cont'd)

### Synchrotron radiation (CGS)

For a relativistic particle of charge  $e$ , velocity  $\beta$ ,  $\gamma$ , energy  $E$ , traveling in a circular orbit of radius  $R$ :

$$\begin{aligned} \text{Energy loss/revolution (MeV)} &= \frac{4\pi}{3} \frac{e^2}{R} \beta^3 \gamma^4 \\ &\cong 0.0885 [E(\text{GeV})]^4 / R(\text{m}) \text{ for } e^\pm \text{ if } \beta \cong 1. \end{aligned}$$

Energy spectrum: The energy radiated into the photon energy interval  $d(\hbar\omega)$  is

$$dI = \alpha \gamma F(\omega/\omega_c) d(\hbar\omega),$$

where  $\alpha = e^2/(\hbar c)$  is the fine-structure constant,

$F(y) = 2\sqrt{3}y \int_0^\infty dx K_{5/3}(x)$ , with  $K_{5/3}(x)$  a modified cylindrical

Bessel function of the third kind, and  $\omega_c = 3\gamma^3 c/R$  is a critical frequency;

$$\hbar\omega_c \text{ (keV)} \cong 4.44 [E(\text{GeV})]^3 / R(\text{m}) \text{ for } e^\pm \text{ if } \beta \cong 1.$$

In the limit  $\gamma \gg 1$ ,

for  $\omega \ll \omega_c$ :

$$\frac{dI}{d(\hbar\omega)} \cong 3.3\alpha \left( \frac{\omega R}{c} \right)^{1/3};$$

for  $\frac{\omega}{\omega_c} = (0.01, 0.1, 0.2, 1.0, 2.0)$ :

$$\frac{dI}{d(\hbar\omega)} \cong (1.0, 1.6, 1.6, 0.5, 0.08)\alpha\gamma, \text{ respectively;}$$

for  $\omega \gtrsim 2\omega_c$ :

$$\frac{dI}{d(\hbar\omega)} \cong \sqrt{3}\pi\alpha\gamma \left( \frac{\omega}{\omega_c} \right)^{1/2} e^{-2\omega/\omega_c}.$$

The radiation is confined to angles  $\lesssim 1/\gamma$  relative to the instantaneous direction of motion.

See J.D. Jackson, *Classical Electrodynamics*, 2nd edition (John Wiley & Sons, New York, 1975) for more formulae and details. (Prepared April 1974; revised April 1984.)

## CROSS SECTIONS, DECAYS, STRUCTURE FUNCTIONS, AND KINEMATICS

### A. LORENTZ TRANSFORMATIONS

The energy  $E$  and three-momentum  $\vec{p}$  of a particle form a four-vector  $p = (E, \vec{p})$ . If the particle has rest mass  $m$  and velocity  $\vec{v}$ , then  $\vec{p} = \gamma m \vec{v}$  and  $E = \gamma mc^2$ . It follows that the velocity expressed as a fraction of the speed of light is  $\beta \equiv \vec{v}/c = \vec{p}c/E$ . In what follows we adopt units where  $c = 1$ . Viewed from a second frame with velocity  $\beta\hat{z}$  relative to the original frame, the components of  $p$  are  $(E', \vec{p}')$ , where

$$E' = \gamma E - \beta\gamma p_z,$$

$$p'_z = \gamma p_z - \beta\gamma E,$$

$$p'_x = p_x; p'_y = p_y,$$

and where  $\gamma = (1 - \beta^2)^{-1/2}$ . It follows that the scalar product of two momenta,  $p_1 \cdot p_2 = E_1 E_2 - \vec{p}_1 \cdot \vec{p}_2$ , is invariant, that is, frame independent.

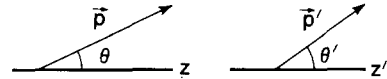
If the lab frame of two colliding particles of masses  $m_1$  and  $m_2$  is the frame in which  $m_2$  is at rest, then the lab velocity of the center of mass is

$$\beta_{\text{cm}} = p_1 / (E_1 + m_2).$$

The total energy available in the center of mass is

$$E_{\text{cm}}^{\text{tot}} = (m_1^2 + m_2^2 + 2E_1 m_2)^{1/2}.$$

If  $\vec{p}$  makes an angle  $\theta$  with the  $z$ -axis, then  $\vec{p}'$  makes an angle  $\theta'$  with the  $z'$ -axis,



where

$$\tan \theta' = \frac{|\vec{p}| \sin \theta}{\gamma |\vec{p}| \cos \theta - \beta \gamma E}.$$

In particular, if the unprimed frame is the center of mass and the primed frame is the lab, and if the velocity of the center of mass in the lab frame is  $\beta^* \hat{z}$ , we use  $\beta = -\beta^*$  above to find (denoting  $p_{\text{cm}} = |\vec{p}_{\text{cm}}|$ )

$$\tan \theta_{\text{lab}} = \frac{p_{\text{cm}} \sin \theta_{\text{cm}}}{\gamma^* p_{\text{cm}} \cos \theta_{\text{cm}} + \beta^* \gamma^* E_{\text{cm}}}.$$

If  $\beta^* > p_{\text{cm}}/E_{\text{cm}}$ , the particle is necessarily moving forward in the lab and

$$(\tan \theta_{\text{lab}})_{\text{max}} = \frac{p_{\text{cm}}}{\gamma^* E_{\text{cm}}} \frac{1}{\sqrt{\beta^{*2} - p_{\text{cm}}^2 / E_{\text{cm}}^2}}.$$

We denote  $p_{\perp} = p_{\perp}' = |\vec{p}| \sin \theta_{\text{cm}}$ . Then given a fixed  $p_{\text{cm}}$  and  $E_{\text{cm}}$ , as, for example, in a two-to-two scattering process, as  $\theta_{\text{cm}}$  varies from 0 to  $2\pi$  the lab momentum describes an ellipse:

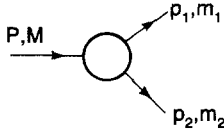
$$\frac{(p'_{\perp} - \beta^* \gamma^* E_{\text{cm}})^2}{\gamma^{*2} p_{\text{cm}}^2} + \frac{p_{\perp}^2}{p_{\text{cm}}^2} = 1.$$



## CROSS SECTIONS, DECAYS, STRUCTURE FUNCTIONS, AND KINEMATICS (Cont'd)

### B. DECAYS

#### B.1.a Two-body kinematics:



In the rest frame of the decaying particle,

$$E_1 = \frac{M^2 + m_1^2 - m_2^2}{2M},$$

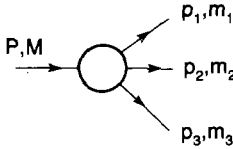
$$|\vec{p}_1| = \left[ \frac{[M^2 - (m_1 + m_2)^2][M^2 - (m_1 - m_2)^2]}{4M^2} \right]^{1/2}.$$

**B.1.b Two-body partial decay rate:** If  $\mathcal{M}$  is the Lorentz invariant matrix element (see Section D below), the partial decay rate in the rest frame of the decaying particle is

$$d\Gamma = \frac{1}{32\pi^2} |\mathcal{M}|^2 \frac{|\vec{p}_1| d\Omega}{M^2},$$

where  $d\Omega$  is the differential solid angle in the rest frame of the decaying particle.

#### B.2.a Three-body kinematics:



We denote

$$p_{12} = p_1 + p_2, \quad m_{12}^2 = p_{12}^2, \quad \text{etc.}$$

Then

$$m_{12}^2 + m_{23}^2 + m_{13}^2 = M^2 + m_1^2 + m_2^2 + m_3^2.$$

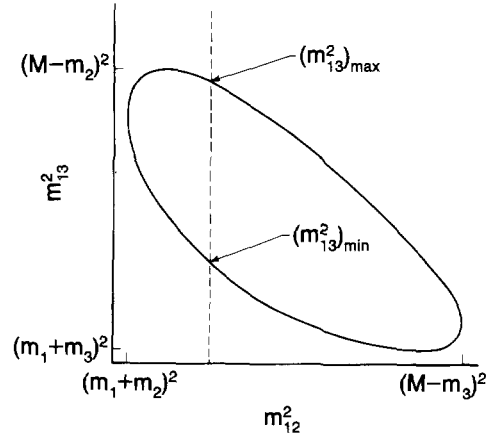
The invariant mass of the pair 1-2 is related to the energy of particle 3 in the rest frame of  $M$ ,

$$m_{12}^2 = (P - p_3)^2 = M^2 + m_3^2 - 2ME_3.$$

**B.2.b Dalitz plot:** If the orientation of the decaying particle is ignored, there are two kinematic variables, which may be chosen to be  $m_{12}^2$  and  $m_{13}^2$ . For fixed  $m_{12}^2$ , the range of  $m_{13}^2$  is determined by letting  $\vec{p}_1$  be parallel or antiparallel to  $\vec{p}_3$ . In the rest frame of  $(p_1 + p_2)$ , the energy of particle 3 is  $E_3^* = (M^2 - m_{12}^2 - m_3^2)/(2m_{12})$ , and that of particle 1 is  $E_1^* = (m_{12}^2 + m_1^2 - m_2^2)/(2m_{12})$ . Thus for a given  $m_{12}^2$ ,

$$(m_{13}^2)_{\max} = (E_1^* + E_3^*)^2 - \left( \sqrt{E_1^{*2} - m_1^2} - \sqrt{E_3^{*2} - m_3^2} \right)^2$$

$$(m_{13}^2)_{\min} = (E_1^* + E_3^*)^2 - \left( \sqrt{E_1^{*2} - m_1^2} + \sqrt{E_3^{*2} - m_3^2} \right)^2.$$



The scatter plot in  $m_{12}^2$  and  $m_{13}^2$  is called a Dalitz plot. Phase space density is uniform across the plot. See below.

**B.2.c Three-body phase space:** Fixing the energies  $E_1$  and  $E_2$  of two of the final state particles in the  $M$  rest frame determines the relative orientation of the three outgoing particles. Their momenta may then be regarded as a rigid body whose orientation with respect to the initial particle is specified by the Euler angles  $\alpha$ ,  $\beta$ , and  $\gamma$ . The partial decay rate in the  $M$  rest frame is

$$d\Gamma = \frac{(2\pi)^{-5}}{16M} |\mathcal{M}|^2 dE_1 dE_2 d\alpha d\cos\beta d\gamma.$$

If the angles are integrated out, we have the Dalitz plot form,

$$d\Gamma = \frac{(2\pi)^{-3}}{8M} |\mathcal{M}|^2 dE_1 dE_2 = \frac{(2\pi)^{-3}}{32M^3} |\mathcal{M}|^2 dm_{12}^2 dm_{23}^2.$$

An alternative expression is

$$d\Gamma = \frac{(2\pi)^{-5}}{16M^2} |\mathcal{M}|^2 |\vec{p}_1^*| |\vec{p}_3| dm_{12} d\Omega_1^* d\Omega_3,$$

where

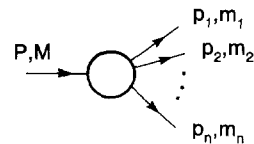
$$|\vec{p}_1^*| = \left[ \frac{[m_{12}^2 - (m_1 + m_2)^2][m_{12}^2 - (m_1 - m_2)^2]}{4m_{12}^2} \right]^{1/2}$$

is the momentum of particle 1 in the rest frame of  $m_{12}$ ,

$$|\vec{p}_3| = \left[ \frac{[M^2 - (m_{12} + m_3)^2][M^2 - (m_{12} - m_3)^2]}{4M^2} \right]^{1/2}$$

is the momentum of particle 3 in the  $M$  rest frame,  $d\Omega_1^*$  is the solid angle element for particle 1 in the 1-2 rest frame, and  $d\Omega_3$  is the solid angle element for particle 3 in the  $M$  rest frame.

#### B.3 $n$ -body phase space:



## CROSS SECTIONS, DECAYS, STRUCTURE FUNCTIONS, AND KINEMATICS (Cont'd)

The partial decay rate in the  $M$  rest frame is

$$d\Gamma = \frac{(2\pi)^4}{2M} |\mathcal{M}|^2 d\Phi_n(P; p_1, \dots, p_n),$$

where

$$d\Phi_n(P; p_1, \dots, p_n) = \delta^4(P - \sum_{i=1}^n p_i) \prod_{i=1}^n \frac{d^3 p_i}{(2\pi)^3 2E_i}.$$

In particular,

$$d\Phi_2(P; p_1, p_2) = (2\pi)^{-6} \frac{|\vec{p}_1^*|}{4M} d\Omega_1^*,$$

where  $|\vec{p}_1^*|$  is the momentum of particle 1 in the  $M$  rest frame and  $d\Omega_1^*$  is the solid angle element in the same frame.

Phase space for  $n$  particles can be related to that for  $n-1$  by treating particles 1 and 2 as a single system of momentum  $p_{12} = p_1 + p_2$  and mass squared  $m_{12}^2 = p_{12}^2$ . Thus

$$d\Phi_n(P; p_1, p_2, \dots, p_n) = d\Phi_{n-1}(P; p_{12}, p_3, \dots, p_n) \\ \times d\Phi_2(p_{12}; p_1, p_2) (2\pi)^3 dm_{12}^2.$$

### C. CROSS SECTIONS AND STRUCTURE FUNCTIONS

Throughout Section C, we set  $\hbar = 1$ ,  $c = 1$ . Use  $\hbar c = 197.3$  MeV fermi, and  $(\hbar c)^2 = 0.3894$  GeV<sup>2</sup> mb for conversions.

**C.1 Partial waves:** The amplitude in the center of mass for elastic scattering of spinless particles may be written in a partial wave expansion

$$f(k, \theta) = \frac{1}{k} \sum_{\ell} (2\ell + 1) a_{\ell} P_{\ell}(\cos \theta),$$

where  $k$  is the c.m. momentum,  $\theta$  is the c.m. scattering angle,  $a_{\ell} = (\eta_{\ell} e^{2i\delta_{\ell}} - 1)/2i$ ,  $0 \leq \eta_{\ell} \leq 1$ , and  $\delta_{\ell}$  is the phase shift of the  $\ell^{\text{th}}$  partial wave. For purely elastic scattering,  $\eta_{\ell} = 1$ . The differential cross section is

$$\frac{d\sigma}{d\Omega} = |f(k, \theta)|^2.$$

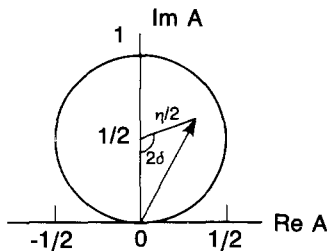
The optical theorem is

$$\sigma_{\text{tot}} = \frac{4\pi}{k} \text{Im} f(k, 0),$$

and the cross section in the  $\ell^{\text{th}}$  partial wave is

$$\sigma_{\ell} = \frac{4\pi}{k^2} (2\ell + 1) |a_{\ell}|^2 \leq \frac{4\pi (2\ell + 1)}{k^2}.$$

The partial-wave amplitude  $a_{\ell}$  can be displayed in an Argand plot.



The usual Lorentz invariant matrix element  $\mathcal{M}$  (see Section D below) for the elastic process is related to  $f(k, \theta)$  by

$$\mathcal{M} = -8\pi\sqrt{s} f(k, \theta),$$

so

$$\sigma_{\text{tot}} = -\frac{1}{2k\sqrt{s}} \text{Im} \mathcal{M}(t=0),$$

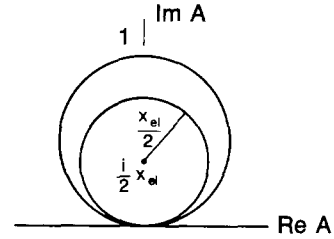
where

$s$  and  $t$  are the center-of-mass energy squared and momentum transfer squared, respectively (see Section C.3.a).

**C.2 Resonances:** The Breit-Wigner form for  $a_{\ell}$  with a resonance at c.m. energy  $E_R$ , elastic width  $\Gamma_{\text{el}}$ , and total width  $\Gamma_{\text{tot}}$  is

$$a_{\ell} = \frac{\frac{1}{2}\Gamma_{\text{el}}}{E_R - E - \frac{i}{2}\Gamma_{\text{tot}}},$$

where  $E$  is the c.m. energy. This gives a circle in the Argand plot with center  $ix_{\text{el}}/2$  and radius  $x_{\text{el}}/2$ , where  $x_{\text{el}} = \Gamma_{\text{el}}/\Gamma_{\text{tot}}$ . The quantity  $x_{\text{el}}$  is called the elasticity. The amplitude has a pole at  $E = E_R - i\Gamma_{\text{tot}}/2$ .

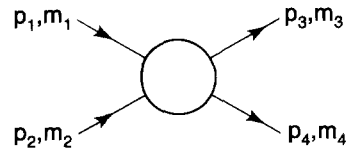


The Breit-Wigner cross section for a spin- $J$  resonance produced in the collision of particles of spin  $S_1$  and  $S_2$  is

$$\sigma_{BW}(E) = \frac{(2J+1)}{(2S_1+1)(2S_2+1)} \frac{\pi}{k^2} \frac{B_{\text{in}} B_{\text{out}} \Gamma_{\text{tot}}^2}{(E - E_R)^2 + \Gamma_{\text{tot}}^2/4},$$

where  $k$  is the c.m. momentum,  $E$  is the c.m. energy, and  $B_{\text{in}}$  and  $B_{\text{out}}$  are the branching fractions of the resonance into the entrance and exit channels. The  $2S+1$  factors are the multiplicities of the incident spin states, so they are replaced by 2 for photons, etc.

**C.3.a Two-body scattering kinematics:**



In the center of mass,

$$E_{\text{1cm}} = \frac{s + m_1^2 - m_2^2}{2\sqrt{s}}, \\ p_{\text{1cm}} = \left[ \frac{[s - (m_1 + m_2)^2][s - (m_1 - m_2)^2]}{4s} \right]^{1/2} \\ = \frac{p_{\text{1lab}} m_2}{\sqrt{s}},$$

## CROSS SECTIONS, DECAYS, STRUCTURE FUNCTIONS, AND KINEMATICS (Cont'd)

where  $\sqrt{s}$  is the total c.m. energy. The Lorentz invariant Mandelstam variables are

$$\begin{aligned} s &= (p_1 + p_2)^2 = (p_3 + p_4)^2 \\ &= m_1^2 + 2E_1E_2 - 2\vec{p}_1 \cdot \vec{p}_2 + m_2^2, \\ t &= (p_1 - p_3)^2 = (p_2 - p_4)^2 \\ &= m_1^2 - 2E_1E_3 + 2\vec{p}_1 \cdot \vec{p}_3 + m_3^2, \\ u &= (p_1 - p_4)^2 = (p_2 - p_3)^2 \\ &= m_1^2 - 2E_1E_4 + 2\vec{p}_1 \cdot \vec{p}_4 + m_4^2, \end{aligned}$$

and they satisfy

$$s + t + u = m_1^2 + m_2^2 + m_3^2 + m_4^2.$$

If  $\theta_{\text{cm}}$  is the c.m. scattering angle between particles 1 and 3, then (denoting  $p_{1\text{cm}} = |\vec{p}_{1\text{cm}}|$ ,  $p_{3\text{cm}} = |\vec{p}_{3\text{cm}}|$ )

$$t = (E_{1\text{cm}} - E_{3\text{cm}})^2 - (p_{1\text{cm}} - p_{3\text{cm}})^2 - 4p_{1\text{cm}}p_{3\text{cm}} \sin^2(\theta_{\text{cm}}/2).$$

For  $\theta_{\text{cm}} = 0$ ,  $-t$  is a minimum.

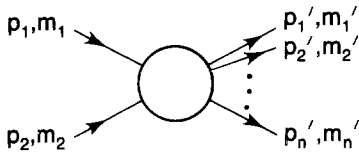
**C.3.b Two-body differential cross sections:** In the center of mass or lab,

$$\frac{d\sigma}{dt} = \frac{1}{64\pi s} \frac{1}{p_{1\text{cm}}^2} |\mathcal{M}|^2.$$

In the center of mass,

$$\frac{d\sigma}{d\Omega_{\text{cm}}} = \frac{p_{1\text{cm}}p_{3\text{cm}}}{\pi} \frac{d\sigma}{dt}.$$

**C.4  $n$ -body differential cross sections:**



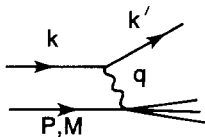
In the c.m. or lab

$$d\sigma = \frac{(2\pi)^4 |\mathcal{M}|^2}{4 \sqrt{(p_1 p_2)^2 - m_1^2 m_2^2}} d\Phi_n(p_1 + p_2; p_1', p_2', \dots, p_n'),$$

where  $n$ -body phase space,  $d\Phi_n$ , is described in Section B.3 above.

Note that  $\sqrt{(p_1 p_2)^2 - m_1^2 m_2^2} = p_{1\text{lab}} m_2 = p_{1\text{cm}} \sqrt{s}$ .

**C.5.a Leptonproduction kinematics:**



$q = k - k'$  is the four-momentum transferred to the target.

Invariant quantities:

$\nu = \frac{q \cdot P}{M} = E - E'$  is the lepton's energy loss in the lab (in earlier literature sometimes  $\nu = q \cdot P$ ). Here,  $E$  and  $E'$  are the initial and final lepton energies in the lab.

$Q^2 = -q^2 = 2(EE' - \vec{k} \cdot \vec{k}') - m_\ell^2 - m_{\ell'}^2$ , where  $m_\ell$  ( $m_{\ell'}$ ) is the initial (final) lepton mass. If  $EE' \sin^2(\theta/2) \gg m_\ell^2, m_{\ell'}^2$ , then

$$\approx 4EE' \sin^2(\theta/2), \text{ where } \theta \text{ is the lepton's scattering angle in the lab.}$$

$x = \frac{Q^2}{2M\nu}$  In the parton model,  $x$  is the fraction of the target nucleon's momentum carried by the struck quark. See section on QCD.

$y = \frac{q \cdot P}{k \cdot P} = \frac{\nu}{E}$  is the fraction of the lepton's energy lost in the lab.

$W^2 = (P + q)^2 = M^2 + 2M\nu - Q^2$  is the mass squared of the system recoiling against the lepton.

**C.5.b Leptonproduction cross sections:**

$$\frac{d^2\sigma}{dx dy} = 2M\nu E \frac{d^2\sigma}{d\nu dQ^2} = \frac{2\pi M\nu}{E'} \frac{d^2\sigma}{d\Omega_{\text{lab}} dE'} = 2xME \frac{d^2\sigma}{dx dQ^2}$$

**C.5.b.i Electroproduction structure functions:**

$$\begin{aligned} \frac{d^2\sigma}{dx dy} &= \frac{8\pi\alpha^2 ME}{Q^4} \left[ \frac{1+(1-y)^2}{2} 2xF_1^{\text{em}} \right. \\ &\quad \left. + (1-y)(F_2^{\text{em}} - 2xF_1^{\text{em}}) - \frac{M}{2E} xyF_2^{\text{em}} \right]. \end{aligned}$$

$F_1^{\text{em}}(x, Q^2)$  and  $F_2^{\text{em}}(x, Q^2)$  are the (unpolarized) structure functions, which are, in the naive parton model, independent of  $Q^2$ .

**C.5.b.ii Neutrino production structure functions:**

$$\begin{aligned} \frac{d^2\sigma^\nu}{dx dy} &= \frac{G_F^2 ME}{\pi} \left[ \left(1-y - \frac{M}{2E} xy\right) F_2^\nu + \frac{y^2}{2} 2xF_1^\nu \right. \\ &\quad \left. + \left(y - \frac{y^2}{2}\right) xF_3^\nu \right], \end{aligned}$$

$$\begin{aligned} \frac{d^2\sigma^{\bar{\nu}}}{dx dy} &= \frac{G_F^2 ME}{\pi} \left[ \left(1-y - \frac{M}{2E} xy\right) F_2^{\bar{\nu}} + \frac{y^2}{2} 2xF_1^{\bar{\nu}} \right. \\ &\quad \left. - \left(y - \frac{y^2}{2}\right) xF_3^{\bar{\nu}} \right]. \end{aligned}$$

The structure functions  $F_i^{\nu, \bar{\nu}}$  are related to quark distributions in the parton model (see section on QCD). There are separate  $F_i$ 's for neutral- and charged-current processes.

**C.6.a  $e^+e^-$  annihilation:** For pointlike spin-1/2 fermions in the c.m., the differential cross section for  $e^+e^- \rightarrow f\bar{f}$  via single photon annihilation is

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2}{4s} \beta \left[ 1 + \cos^2\theta + (1-\beta^2)\sin^2\theta \right] Q_f^2,$$

## CROSS SECTIONS, DECAYS, STRUCTURE FUNCTIONS, AND KINEMATICS (Cont'd)

where  $\beta$  is the velocity of the final state fermion in the center of mass and  $Q_f$  is the charge of the fermion in units of the proton charge. For  $\beta \rightarrow 1$ ,

$$\sigma = \frac{4\pi\alpha^2}{3s} Q_f^2 = \frac{86.8 Q_f^2 nb}{s(GeV^2)}.$$

At higher energies the  $Z^0$  (mass  $M_Z$  and width  $\Gamma_Z$ ) must be included, and the differential cross section for  $e^+e^- \rightarrow f\bar{f}$  becomes

$$\begin{aligned} \frac{d\sigma}{d\Omega} &= \frac{\alpha^2\beta}{4s} \left[ Q_f^2 [1 + \cos^2\theta + (1 - \beta^2)\sin^2\theta] \right. \\ &- 2Q_f\chi_1 \{ VV_f [1 + \cos^2\theta + (1 - \beta^2)\sin^2\theta] - 2a_f\beta\cos\theta \} \\ &+ \chi_2 \{ V_f^2(1 + V^2)[1 + \cos^2\theta + (1 - \beta^2)\sin^2\theta] \\ &\left. + \beta a_f^2(1 + V^2)[1 + \cos^2\theta] - 8\beta VV_f a_f \cos\theta \right], \end{aligned}$$

$$\chi_1 = \frac{1}{16\sin^2\theta_W \cos^2\theta_W} \frac{s(s - M_Z^2)}{(s - M_Z^2)^2 + \Gamma_Z^2 M_Z^2},$$

$$\chi_2 = \frac{1}{256\sin^4\theta_W \cos^4\theta_W} \frac{s^2}{(s - M_Z^2)^2 + \Gamma_Z^2 M_Z^2},$$

$$V = -1 + 4\sin^2\theta_W,$$

$$a_f = 2T_{3f},$$

$$V_f = 2T_{3f} - 4Q_f \sin^2\theta_W,$$

where the subscript  $f$  refers to the particular fermion and

$$T_3 = +1/2 \text{ for } \nu_e, \nu_\mu, \nu_\tau, u, c, t,$$

$$T_3 = -1/2 \text{ for } e^-, \mu^-, \tau^-, d, s, b.$$

**C.6.b  $e^+e^-$  two-photon process:** In the equivalent photon approximation, the cross section for  $e^+e^- \rightarrow e^+e^-X$  is related to the cross section for  $\gamma\gamma \rightarrow X$  by

$$d\sigma_{e^+e^- \rightarrow e^+e^-X}(s) = \eta^2 \int d\omega f(\omega) d\sigma_{\gamma\gamma \rightarrow X}(\omega s),$$

where

$$\eta \approx \frac{\alpha}{2\pi} \ell n \left[ \frac{s}{4m_e^2} \right]$$

and

$$f(\omega) = \frac{1}{\omega} \left[ (2 + \omega)^2 \ell n \frac{1}{\omega} - 2(1 - \omega)(3 + \omega) \right].$$

For the production of a resonance of mass  $m_R$  and spin  $J$ ,

$$\sigma(e^+e^- \rightarrow e^+e^-R) = \eta^2 \frac{(2J+1)8\pi^2 \Gamma(R \rightarrow \gamma\gamma)}{m_R s} f\left(\frac{m_R^2}{s}\right).$$

**C.7 Inclusive hadronic reactions:** A particle's momentum can be parametrized by selecting a particular direction for the  $z$ -axis and writing

$$(E = m_\perp \cosh y, p_z = m_\perp \sinh y, p_x, p_y),$$

where

$$m_\perp^2 = m^2 + p_x^2 + p_y^2,$$

$$y = \frac{1}{2} \ell n \left( \frac{E + p_z}{E - p_z} \right) = \ell n \left( \frac{E + p_z}{m_\perp} \right) = \tanh^{-1} \left( \frac{p_z}{E} \right).$$

The variable  $y$  is called the rapidity. A boost in the  $z$ -direction then modifies  $y$  by  $y \rightarrow y + \Delta$ , where  $\gamma = \cosh \Delta$ ,  $\beta = \tanh \Delta$ . Thus the shape of the distribution  $dN/dy$  is invariant under such a boost, and

$$E \frac{d^3\sigma}{d^3p} = \frac{d^3\sigma}{dy d^2p_\perp}.$$

A variable which may be easier to measure experimentally is the pseudorapidity  $\eta$

$$\eta = \ell n(\cot \theta/2)$$

where  $\theta$  is the angle between the particle and the beam. In the limit  $E \gg m$ , it follows that  $\eta = y$ .

Feynman's  $x$  variable is defined to be

$$x = \left( \frac{p_z}{p_{z \max}} \right)_{\text{cm}} \approx \frac{2p_{z \text{ cm}}}{\sqrt{s}} \approx \frac{2m_\perp \sinh y_{\text{cm}}}{\sqrt{s}}.$$

For  $y_{\text{cm}}$  not small ( $e^{-2y_{\text{cm}}} \ll 1$ )

$$x \approx \frac{m_\perp}{\sqrt{s}} e^{y_{\text{cm}}}$$

and

$$(y_{\text{cm}})_{\max} = \ell n \frac{\sqrt{s}}{m}.$$

## D. LORENTZ INVARIANT AMPLITUDES

The quantity  $-i\mathcal{M}$  is determined in perturbation theory by the Feynman rules. Our convention above is consistent with the Appendices of Bjorken and Drell except that fermion spinors are normalized so that  $\bar{u}u = 2m$ , etc. In particular, the  $S$ -matrix for two-body scattering is

$$\langle p'_1 p'_2 | S | p_1 p_2 \rangle = I - i(2\pi)^4 \delta^4(p_1 + p_2 - p'_1 - p'_2)$$

$$\times \frac{\mathcal{M}(p_1, p_2; p'_1, p'_2)}{(2E_1)^{1/2} (2E_2)^{1/2} (2E'_1)^{1/2} (2E'_2)^{1/2}},$$

where the states are normalized so

$$\langle p' | p \rangle = (2\pi)^3 \delta^3(\vec{p} - \vec{p}').$$

C.M. ENERGY AND MOMENTUM VS. BEAM MOMENTUM
(for scattering on a proton target)

E\_cm dE\_cm = m\_p dT\_beam = m\_p v\_beam dP\_beam ≈ m\_p dP\_beam

Table with 4 columns: P\_beam (GeV/c), E\_cm (GeV), Momentum in c.m. (GeV/c), and a corresponding set of values for each. Each column contains a list of values for different scattering angles and beam energies, with sub-headers for gamma-p, pi-p, and K-p interactions.

# CLEBSCH-GORDAN COEFFICIENTS, SPHERICAL HARMONICS, AND D FUNCTIONS

Note: A  $\sqrt{\quad}$  is to be understood over every coefficient; e.g., for  $-8/15$  read  $-\sqrt{8/15}$ .

Notation:  $\begin{matrix} J & J & \dots \\ M & M & \dots \end{matrix}$

$1/2 \times 1/2$   $\begin{matrix} 1 & 1 \\ +1/2 & +1/2 \\ 1 & 0 \\ 0 & 0 \end{matrix}$   $Y_1^0 = \sqrt{\frac{3}{4\pi}} \cos \theta$   $2 \times 1/2$   $\begin{matrix} 5/2 & 5/2 \\ +5/2 & 5/2 \\ 3/2 & 3/2 \\ +2 & 1 \\ 1/2 & -1/2 \\ 1/5 & 4/5 \\ +1 & -1 \end{matrix}$   $Y_1^1 = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}$   $Y_2^0 = \sqrt{\frac{5}{4\pi}} \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2}\right)$   $Y_2^1 = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}$   $Y_2^2 = \frac{1}{4} \sqrt{\frac{15}{2\pi}} \sin^2 \theta e^{2i\phi}$

$1 \times 1/2$   $\begin{matrix} 3/2 & 3/2 \\ +3/2 & 3/2 \\ 3/2 & 1/2 \\ +1 & 1/2 \\ 1/3 & 2/3 \\ 0 & 1/2 \end{matrix}$   $3/2 \times 1$   $\begin{matrix} 5/2 & 5/2 \\ +5/2 & 5/2 \\ 3/2 & 3/2 \\ +3/2 & 1 \\ 2/5 & 3/5 \\ +1/2 & 1 \end{matrix}$   $3/2 \times 1/2$   $\begin{matrix} 2 & 2 \\ +2 & 1 \\ 1 & 1 \\ -3/2 & -1/2 \\ 1/4 & 3/4 \\ +1/2 & 1/2 \end{matrix}$

$2 \times 1$   $\begin{matrix} 3 & 3 \\ +3 & 2 \\ 2 & 2 \\ +2 & 0 \\ 1/3 & 2/3 \\ 2/3 & -1/3 \end{matrix}$   $1 \times 1$   $\begin{matrix} 2 & 2 \\ +2 & 1 \\ 1 & 1 \\ 1/2 & 1/2 \\ 0 & 1/2 \end{matrix}$   $1 \times 1$   $\begin{matrix} 1 & 1 \\ +1 & 1 \\ 0 & 0 \\ 1/2 & 1/2 \\ 1/2 & -1/2 \end{matrix}$

$1 \times 1$   $\begin{matrix} 2 & 2 \\ +2 & 1 \\ 1 & 1 \\ 1/2 & 1/2 \\ 0 & 1/2 \end{matrix}$   $1 \times 1$   $\begin{matrix} 2 & 2 \\ +2 & 1 \\ 1 & 1 \\ 1/2 & 1/2 \\ 0 & 1/2 \end{matrix}$

$Y_l^{-m} = (-1)^m Y_l^m e^{-im\phi}$   $d_{m,0}^l = \sqrt{\frac{4\pi}{2l+1}} Y_l^m e^{-im\phi}$

$(j_1 j_2 m_1 m_2 | j_1 j_2 J M)$   
 $= (-1)^{J-j_1-j_2} (j_2 j_1 m_2 m_1 | j_2 j_1 J M)$

$d_{m',m}^{j'} = (-1)^{m-m'} d_{m,m'}^{j'}$   $d_{1/2,1/2}^{3/2} = \cos \frac{\theta}{2}$   $d_{1/2,-1/2}^{3/2} = -\sin \frac{\theta}{2}$   $d_{1,1}^{3/2} = \frac{1+\cos \theta}{2}$   $d_{1,0}^{3/2} = \frac{\sin \theta}{\sqrt{2}}$   $d_{1,-1}^{3/2} = \frac{1-\cos \theta}{2}$   $d_{0,0}^{3/2} = \cos \theta$

$2 \times 3/2$   $\begin{matrix} 7/2 & 5/2 \\ +7/2 & 5/2 \\ 5/2 & 3/2 \\ +3/2 & 1/2 \\ 3/7 & 4/7 \\ 4/7 & -3/7 \end{matrix}$   $2 \times 2$   $\begin{matrix} 4 & 3 \\ +4 & 3 \\ 3 & 2 \\ +2 & 1 \\ 1/2 & 1/2 \\ 1/2 & -1/2 \end{matrix}$   $3/2 \times 3/2$   $\begin{matrix} 3 & 2 \\ +3 & 2 \\ 1/2 & 1/2 \\ +1/2 & -1/2 \\ 1/5 & 1/5 \\ +1/2 & 1/2 \end{matrix}$

$d_{3/2,3/2}^{3/2} = \frac{1+\cos \theta}{2} \cos \frac{\theta}{2}$   $d_{3/2,1/2}^{3/2} = -\sqrt{3} \frac{1-\cos \theta}{2} \sin \frac{\theta}{2}$   $d_{3/2,-1/2}^{3/2} = \sqrt{3} \frac{1-\cos \theta}{2} \cos \frac{\theta}{2}$   $d_{3/2,-3/2}^{3/2} = -\frac{1-\cos \theta}{2} \sin \frac{\theta}{2}$

$d_{2,2}^{3/2} = \left(\frac{1+\cos \theta}{2}\right)^2$   $d_{2,1}^{3/2} = \frac{1-\cos \theta}{2} \sin \theta$   $d_{2,0}^{3/2} = \frac{\sqrt{6}}{4} \sin^2 \theta$   $d_{2,-1}^{3/2} = -\frac{1-\cos \theta}{2} \sin \theta$   $d_{2,-2}^{3/2} = \left(\frac{1-\cos \theta}{2}\right)^2$

$d_{1,1}^{3/2} = \frac{1+\cos \theta}{2} (2\cos \theta - 1)$   $d_{1,0}^{3/2} = -\sqrt{\frac{3}{2}} \sin \theta \cos \theta$   $d_{1,-1}^{3/2} = \frac{1-\cos \theta}{2} (2\cos \theta + 1)$   $d_{0,0}^{3/2} = \left(\frac{3}{2} \cos^2 \theta - \frac{1}{2}\right)$

Sign convention is that of Wigner (*Group Theory*, Academic Press, New York, 1959), also used by Condon and Shortley (*The Theory of Atomic Spectra*, Cambridge Univ. Press, New York, 1953), Rose (*Elementary Theory of Angular Momentum*, Wiley, New York, 1957), and Cohen (*Tables of the Clebsch-Gordan Coefficients*, North American Rockwell Science Center, Thousand Oaks, Calif., 1974). The signs and numbers in the current tables have been calculated by computer programs written independently by Cohen and at LBL. (Table extended April 1974.)

## SU(3) ISOSCALAR FACTORS AND REPRESENTATION MATRICES

The most commonly used isoscalar factors, corresponding to the singlet, octet, and decuplet content of  $8 \otimes 8$  and  $10 \otimes 8$ , are displayed at the right. The notation uses particle names to identify the coefficients, so that the pattern of relative couplings can be seen at a glance. We illustrate the use of the coefficients by example; see J.J de Swart, Rev. Mod Phys. 35, 916 (1963) for detailed explanation and phase conventions.

A  $\sqrt{\phantom{x}}$  is understood over every integer in the matrices; the exponent  $\frac{1}{2}$  is a reminder of this. For example, in de Swart's notation the  $\Xi \rightarrow \Omega K$  element of our  $10 \rightarrow 10 \otimes 8$  matrix reads

$$\begin{pmatrix} 10 & 8 & 10 \\ 0-2 & \frac{1}{2} 1 & \frac{1}{2} -1 \end{pmatrix} = \frac{-\sqrt{6}}{\sqrt{24}}.$$

Intramultiplet relative decay strengths can be read directly from our matrices. Thus, the partial widths for  $\Delta \rightarrow (N\pi)_{I=3/2}$  and  $\Omega^* \rightarrow (\Xi\bar{K})_{I=0}$  are in the ratio

$$\frac{\Gamma(\Omega^* \rightarrow (\Xi\bar{K})_{I=0})}{\Gamma(\Delta \rightarrow (N\pi)_{I=3/2})} = \frac{12}{6} \times (\text{phase space factors}).$$

Supplying isospin Clebsch-Gordan coefficients, one obtains, e.g.,

$$\frac{\Gamma(\Omega^{*-} \rightarrow \Xi^0 K^-)}{\Gamma(\Delta^+ \rightarrow p \pi^0)} = \frac{1/2}{2/3} \times \frac{12}{6} \times p.s.f. = \frac{3}{2} \times p.s.f.$$

Partial widths for  $8 \rightarrow 8 \otimes 8$  involve a linear superposition of  $8_1$  (symmetric) and  $8_2$  (antisymmetric) couplings. For example,

$$\Gamma(\Xi^* \rightarrow \Xi\pi) \sim \left[ -\sqrt{\frac{9}{20}}g_1 + \sqrt{\frac{3}{12}}g_2 \right]^2.$$

The relation between  $g_1, g_2$  (with de Swart's normalization) and the standard  $D, F$  couplings appearing in the interaction Lagrangian,

$$\mathcal{L} = -\sqrt{2}D \text{Tr}([\bar{B}, B]_+ M) + \sqrt{2}F \text{Tr}([\bar{B}, B]_- M),$$

is

$$D = \frac{\sqrt{30}}{40}g_1, \quad F = \frac{\sqrt{6}}{24}g_2.$$

Thus,

$$\Gamma(\Xi^* \rightarrow \Xi\pi) \sim (1-2\alpha)^2$$

where  $\alpha \equiv D/(D+F)$ .

When acting upon a representation of dimension  $d$ , the generators of SU(3) transformations,  $\lambda_a$  ( $a=1,8$ ), are  $d \times d$  matrices which obey the following relationships:

$$\begin{cases} [\lambda_a, \lambda_b] = 2if_{abc}\lambda_c \\ \{\lambda_a, \lambda_b\} = \frac{4}{3}\delta_{ab}I + 2d_{abc}\lambda_c \end{cases}$$

where  $I$  is a  $d \times d$  unit matrix. The  $f_{abc}$  are odd under the permutation of any pair of indices, while the  $d_{abc}$  are even. The nonzero elements are

$abc$	$f_{abc}$	$abc$	$d_{abc}$	$abc$	$d_{abc}$
123	1	118	$1/\sqrt{3}$	355	$1/2$
147	$1/2$	146	$1/2$	366	$-1/2$
156	$-1/2$	157	$1/2$	377	$-1/2$
246	$1/2$	228	$1/\sqrt{3}$	448	$-1/(2\sqrt{3})$
257	$1/2$	247	$-1/2$	558	$-1/(2\sqrt{3})$
345	$1/2$	256	$1/2$	668	$-1/(2\sqrt{3})$
367	$-1/2$	338	$1/\sqrt{3}$	778	$-1/(2\sqrt{3})$
458	$\sqrt{3}/2$	344	$1/2$	888	$-1/\sqrt{3}$
678	$\sqrt{3}/2$				

$1 \rightarrow 8 \otimes 8$

$$\begin{pmatrix} \Lambda \\ \Sigma \\ \Xi \\ \Omega \end{pmatrix}_1 \rightarrow \begin{pmatrix} N\bar{K} \Sigma\pi & \Lambda\eta & \Xi K \\ \Sigma\bar{K} \Sigma\pi & \Lambda\eta & \Xi K \\ \Xi\bar{K} \Sigma\pi & \Lambda\eta & \Xi K \\ \Omega\bar{K} \Sigma\pi & \Lambda\eta & \Xi K \end{pmatrix}_{8 \otimes 8} = \frac{1}{\sqrt{8}} \begin{pmatrix} 2 & 3 & -1 & -2 \end{pmatrix}^{1/2}$$

$8_1 \rightarrow 8 \otimes 8$

$$\begin{pmatrix} N \\ \Sigma \\ \Lambda \\ \Xi \end{pmatrix}_{8_1} \rightarrow \begin{pmatrix} N\pi & N\eta & \Sigma K & \Lambda K \\ N\bar{K} \Sigma\pi & \Delta\pi & \Sigma\eta & \Xi K \\ N\bar{K} \Sigma\pi & \Delta\pi & \Sigma\eta & \Xi K \\ \Sigma\bar{K} \Lambda\bar{K} & \Xi\pi & \Xi\eta \end{pmatrix}_{8 \otimes 8} = \frac{1}{\sqrt{20}} \begin{pmatrix} 9 & -1 & -9 & -1 \\ -6 & 0 & 4 & 4 & -6 \\ 2 & -12 & -4 & -2 \\ 9 & -1 & -9 & -1 \end{pmatrix}^{1/2}$$

$8_2 \rightarrow 8 \otimes 8$

$$\begin{pmatrix} N \\ \Sigma \\ \Lambda \\ \Xi \end{pmatrix}_{8_2} \rightarrow \begin{pmatrix} N\pi & N\eta & \Sigma K & \Lambda K \\ N\bar{K} \Sigma\pi & \Delta\pi & \Sigma\eta & \Xi K \\ N\bar{K} \Sigma\pi & \Delta\pi & \Sigma\eta & \Xi K \\ \Sigma\bar{K} \Lambda\bar{K} & \Xi\pi & \Xi\eta \end{pmatrix}_{8 \otimes 8} = \frac{1}{\sqrt{12}} \begin{pmatrix} 3 & 3 & 3 & -3 \\ 2 & 8 & 0 & 0 & -2 \\ 6 & 0 & 0 & 6 \\ 3 & 3 & 3 & -3 \end{pmatrix}^{1/2}$$

$10 \rightarrow 8 \otimes 8$

$$\begin{pmatrix} \Lambda \\ \Sigma \\ \Xi \\ \Omega \end{pmatrix}_{10} \rightarrow \begin{pmatrix} N\pi & \Sigma K \\ N\bar{K} \Sigma\pi & \Delta\pi & \Sigma\eta & \Xi K \\ \Sigma\bar{K} \Lambda\bar{K} & \Xi\pi & \Xi\eta \\ \Xi\bar{K} \Omega\eta \end{pmatrix}_{8 \otimes 8} = \frac{1}{\sqrt{12}} \begin{pmatrix} -6 & 6 \\ -2 & 2 & -3 & 3 & 2 \\ 3 & -3 & 3 & 3 \\ 12 \end{pmatrix}^{1/2}$$

$8 \rightarrow 10 \otimes 8$

$$\begin{pmatrix} N \\ \Sigma \\ \Lambda \\ \Xi \end{pmatrix}_8 \rightarrow \begin{pmatrix} \Delta\pi & \Sigma K \\ \Delta\bar{K} \Sigma\pi & \Sigma\eta & \Xi K \\ \Sigma\pi & \Xi K \\ \Sigma\bar{K} \Xi\pi & \Xi\eta & \Omega K \end{pmatrix}_{10 \otimes 8} = \frac{1}{\sqrt{15}} \begin{pmatrix} -12 & 3 \\ 8 & -2 & -3 & 2 \\ -9 & 6 \\ 3 & -3 & -3 & 6 \end{pmatrix}^{1/2}$$

$10 \rightarrow 10 \otimes 8$

$$\begin{pmatrix} \Delta \\ \Sigma \\ \Xi \\ \Omega \end{pmatrix}_{10} \rightarrow \begin{pmatrix} \Delta\pi & \Delta\eta & \Sigma K \\ \Delta\bar{K} \Sigma\pi & \Sigma\eta & \Xi K \\ \Sigma\bar{K} \Xi\pi & \Xi\eta & \Omega K \\ \Xi\bar{K} \Omega\eta \end{pmatrix}_{10 \otimes 8} = \frac{1}{\sqrt{24}} \begin{pmatrix} 15 & 3 & -6 \\ 8 & 8 & 0 & -8 \\ 12 & 3 & -3 & -6 \\ 12 & -12 \end{pmatrix}^{1/2}$$

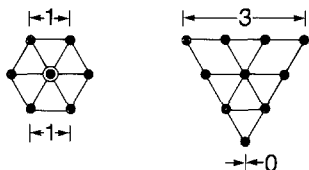
In the fundamental 3-dimensional representation, the  $\lambda_a$ 's are given by

$$\begin{aligned} \lambda_1 &= \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} & \lambda_2 &= \begin{pmatrix} 0 & -i & 0 \\ i & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} & \lambda_3 &= \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & 0 \end{pmatrix} \\ \lambda_4 &= \begin{pmatrix} 0 & 0 & 1 \\ 0 & 0 & 0 \\ 1 & 0 & 0 \end{pmatrix} & \lambda_5 &= \begin{pmatrix} 0 & 0 & -i \\ 0 & 0 & 0 \\ i & 0 & 0 \end{pmatrix} & \lambda_6 &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \\ \lambda_7 &= \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -i \\ 0 & i & 0 \end{pmatrix} & \lambda_8 &= \begin{pmatrix} 1/\sqrt{3} & 0 & 0 \\ 0 & 1/\sqrt{3} & 0 \\ 0 & 0 & -2/\sqrt{3} \end{pmatrix} \end{aligned}$$

## SU(N) MULTIPLETS AND YOUNG DIAGRAMS

This note tells how  $SU(n)$  particle multiplets are identified or labeled, how to find the number of particles in a multiplet from its label, how to draw the Young diagram for a multiplet, and how to use Young diagrams to determine the overall multiplet structure of a composite system, such as a 3-quark or a meson-baryon system.

(1) **Multiplet labels** — An  $SU(n)$  multiplet is uniquely identified by a string of  $(n-1)$  nonnegative integers:  $(\alpha, \beta, \gamma, \dots)$ . Any such set of integers specifies a multiplet. For an  $SU(2)$  multiplet such as an isospin multiplet, the single integer  $\alpha$  is the number of steps from one end of the multiplet to the other (i.e., it is one fewer than the number of particles in the multiplet). In  $SU(3)$ , the two integers  $\alpha$  and  $\beta$  are the numbers of steps across the top and bottom levels of the multiplet diagram. Thus the labels for the  $SU(3)$  octet and decuplet



are  $(1, 1)$  and  $(3, 0)$ . For larger  $n$ , the interpretation of the integers in terms of the geometry of the multiplets, which exist in an  $(n-1)$ -dimensional space, is not so readily apparent.

The label for the  $SU(n)$  singlet is  $(0, 0, \dots, 0)$ . In a flavor  $SU(n)$ , the  $n$  quarks together form a  $(1, 0, \dots, 0)$  multiplet, and the  $n$  anti-quarks belong to a  $(0, \dots, 0, 1)$  multiplet. These two multiplets are conjugate to one another, which means their labels are related by  $(\alpha, \beta, \dots) \leftrightarrow (\dots, \beta, \alpha)$ .

(2) **Number of particles** — The number of particles in a multiplet,  $N = N(\alpha, \beta, \dots)$ , is given as follows (note the pattern of the equations). In  $SU(2)$ ,  $N = N(\alpha)$  is

$$N = \frac{(\alpha+1)}{1}$$

In  $SU(3)$ ,  $N = N(\alpha, \beta)$  is

$$N = \frac{(\alpha+1)}{1} \cdot \frac{(\beta+1)}{1} \cdot \frac{(\alpha+\beta+2)}{2}$$

In  $SU(4)$ ,  $N = N(\alpha, \beta, \gamma)$  is

$$N = \frac{(\alpha+1)}{1} \cdot \frac{(\beta+1)}{1} \cdot \frac{(\gamma+1)}{1} \cdot \frac{(\alpha+\beta+2)}{2} \cdot \frac{(\beta+\gamma+2)}{2} \cdot \frac{(\alpha+\beta+\gamma+3)}{3}$$

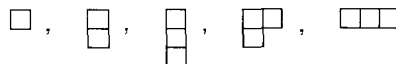
Note that there is no factor with  $(\alpha + \gamma + 2)$ ; only a consecutive sequence of the label integers appears in any factor. One more example should make the pattern clear for any  $SU(n)$ . In  $SU(5)$ ,  $N = N(\alpha, \beta, \gamma, \delta)$  is

$$N = \frac{(\alpha+1)}{1} \cdot \frac{(\beta+1)}{1} \cdot \frac{(\gamma+1)}{1} \cdot \frac{(\delta+1)}{1} \cdot \frac{(\alpha+\beta+2)}{2} \cdot \frac{(\beta+\gamma+2)}{2} \times \frac{(\gamma+\delta+2)}{2} \cdot \frac{(\alpha+\beta+\gamma+3)}{3} \cdot \frac{(\beta+\gamma+\delta+3)}{3} \cdot \frac{(\alpha+\beta+\gamma+\delta+4)}{4}$$

Multiplets that are conjugate to one another obviously have the same number of particles, but so can other multiplets. For example, the  $SU(4)$  multiplets  $(3, 0, 0)$  and  $(1, 1, 0)$  each have 20 particles.

(3) **Young diagrams** — A Young diagram consists of an array of boxes (or some other symbol) arranged in one or more left-justified

rows, with each row being at least as long as the row beneath. The correspondence between a diagram and a multiplet label is: The top row juts out  $\alpha$  boxes to the right past the end of the second row, the second row juts out  $\beta$  boxes to the right past the end of the third row, etc. A diagram in  $SU(n)$  has at most  $n$  rows. There can be any number of "completed" columns of  $n$  boxes buttressing the left of a diagram; these don't affect the label. Thus in  $SU(3)$  the diagrams



represent the multiplets  $(1, 0)$ ,  $(0, 1)$ ,  $(0, 0)$ ,  $(1, 1)$ , and  $(3, 0)$ . In any  $SU(n)$ , the quark multiplet is represented by a single box, the anti-quark multiplet by a column of  $(n-1)$  boxes, and a singlet by a completed column of  $n$  boxes.

(4) **Coupling multiplets together** — The following recipe tells how to find the multiplets that occur in coupling two multiplets together. To couple together more than two multiplets, first couple two, then couple the third with each of the multiplets obtained from the first two, etc.

First a definition: A sequence of the letters  $a, b, c, \dots$  is *admissible* if at any point in the sequence at least as many  $a$ 's have been reached as  $b$ 's, at least as many  $b$ 's have been reached as  $c$ 's, etc. Thus  $abcd$  and  $aabcb$  are admissible sequences and  $abb$  and  $acb$  are not. Now the recipe:

(a) Draw the Young diagrams for the two multiplets, but in one of the diagrams replace the boxes in the first row with  $a$ 's, the boxes in the second row with  $b$ 's, etc. The unlettered diagram forms the upper left-hand corner of all the enlarged diagrams constructed below.

(b) Add the  $a$ 's from the lettered diagram to the unlettered diagram to form all possible legitimate Young diagrams that have no more than one  $a$  per column. (All the  $a$ 's appear in each new diagram.)

(c) Use the  $b$ 's to further enlarge the diagrams already obtained, subject to the same rules. Throw away any diagram in which the sequence of letters formed by reading right to left in the first row, then the second row, etc., is not admissible.

(d) Proceed as in (c) with the  $c$ 's, etc.

Thus, for example, the calculation to find the multiplets that can occur in a system made up of two  $SU(3)$  octets (one might be the  $\pi$ -meson octet, the other the  $N$ -baryon octet) is as follows:

$$\begin{array}{|c|c|} \hline \square & \square \\ \hline \square & \square \\ \hline \end{array} \otimes \begin{array}{|c|c|} \hline a & a \\ \hline b & b \\ \hline \end{array} =$$

$$\begin{array}{|c|c|} \hline \square & a \\ \hline \square & b \\ \hline \end{array} a \oplus \begin{array}{|c|c|} \hline \square & a \\ \hline \square & a \\ \hline \end{array} a \oplus \begin{array}{|c|c|} \hline \square & a \\ \hline \square & b \\ \hline \end{array} a \oplus \begin{array}{|c|c|} \hline \square & a \\ \hline \square & a \\ \hline \end{array} a \oplus \begin{array}{|c|c|} \hline \square & a \\ \hline \square & b \\ \hline \end{array} a \oplus \begin{array}{|c|c|} \hline \square & a \\ \hline \square & a \\ \hline \end{array} a$$

where only the diagrams with admissible sequences and with fewer than four rows (since  $n = 3$ ) have been kept. In terms of multiplet labels, the above may be written

$$(1, 1) \otimes (1, 1) = (2, 2) \oplus (3, 0) \oplus (0, 3) \oplus (1, 1) \oplus (1, 1) \oplus (0, 0),$$

or in terms of numbers of particles,

$$8 \otimes 8 = 27 \oplus 10 \oplus \bar{10} \oplus 8 \oplus 8 \oplus 1.$$

The product of the numbers of the left is equal to the sum on the right. (See the section on the Nonrelativistic Quark Model for results for 3-quark systems.)



## TESTS OF CONSERVATION LAWS\*

In response to the current interest in tests of conservation laws, we have made a list of experimental limits on all weak and electromagnetic decays, mass differences, moments, and a few reactions, whose observation would violate conservation laws. The list is given only in the full Review of Particle Properties, not in the Data Booklet. For the benefit of Data Booklet readers, we have included the best limits from the list in the following text. The list is in two parts: "Discrete Space-Time Symmetries," i.e.,  $C$ ,  $P$ ,  $T$ ,  $CP$ , and  $CPT$ ; and "Number Conservation Laws," i.e., lepton, baryon, hadronic flavor, and charge conservation. The references for these data can be found in the Stable Particle Section of the Full Listings in this Review. A discussion of these tests follows.

### CPT INVARIANCE

General principles of relativistic field theory require invariance under the combined transformation  $CPT$ . The simplest tests of  $CPT$  invariance are the equality of the masses and lifetimes of a particle and its antiparticle. The best test comes from a limit on the mass difference between  $K^0$  and  $\bar{K}^0$ . Any such mass difference contributes to the  $CP$ -violating parameter  $\epsilon$ . In fact  $\epsilon$  can be explained by a  $CPT$ -conserving but  $CP$ -violating mixing of  $K^0$  and  $\bar{K}^0$ , which yields a prediction that  $\phi_{+-} \approx 44^\circ$ , while a  $K^0 - \bar{K}^0$  mass difference would yield  $\phi_{+-} \approx 44^\circ + 90^\circ$ . It is thus possible to deduce that  $|m(K^0) - m(\bar{K}^0)| < 10^{-4} |m(K_S) - m(K_L)| < 3 \times 10^{-10}$  eV. Also, an upper limit on  $|m(D^0) - m(\bar{D}^0)|$  can be derived from the bound  $|m(D_1^0) - m(D_2^0)| < 6.5 \times 10^{-10}$  MeV (inferred from bound on  $D^0 \rightarrow \bar{D}^0 \rightarrow \mu^-$  anything), given an input value of, or bound on, the  $CP$ -violation parameter  $\epsilon$  for  $D^0 - \bar{D}^0$  mixing.

### CP AND T INVARIANCE

Given  $CPT$  invariance,  $CP$  violation and  $T$  violation are equivalent. So far the only evidence for  $CP$  or  $T$  violation comes from the measurements of  $\eta_{+-}$ ,  $\eta_{00}$ , and the semileptonic decay charge asymmetry for  $K_L$ , e.g.,  $|\eta_{+-}| = |A(K_L^0 \rightarrow \pi^+ \pi^-) / A(K_S^0 \rightarrow \pi^+ \pi^-)| = (2.275 \pm 0.021) \times 10^{-3}$  and  $|\Gamma(K_L^0 \rightarrow \pi^- e^+ \nu) - \Gamma(K_L^0 \rightarrow \pi^+ e^- \bar{\nu})| / |\text{sum}| = (0.333 \pm 0.014)\%$ . Other searches for  $CP$  or  $T$  violation should be divided into (a) those that involve weak interactions or parity violation, and (b) those that involve processes allowed by the strong or electromagnetic interactions. In class (a) the most sensitive is probably the search for an electric dipole moment of the neutron, which is measured to be  $(2.3 \pm 2.3) \times 10^{-25}$  e cm. A nonzero value requires both  $P$  and  $T$  violation to be nonzero. Class (b) searches involve looking for  $C$  or  $T$  violation in strong or electromagnetic processes. Examples are the search for  $C$  violation in  $\eta$  decay, believed to be an electromagnetic process, e.g., as measured by  $\Gamma(\eta \rightarrow \mu^+ \mu^- \pi^0) / \Gamma(\eta \rightarrow \text{all}) < 5 \times 10^{-6}$ , and the search for  $T$  violation in a number of nuclear and electromagnetic reactions.

### CONSERVATION OF LEPTON NUMBERS

Present experimental evidence and the standard electroweak theory are consistent with the absolute conservation of three separate lepton numbers: electron number  $L_e$ , muon number  $L_\mu$ , and  $\tau$ -number  $L_\tau$ . Searches for violations are of the following types:

**a)  $\Delta L = 2$  for one type of lepton.** The best limit comes from the search for neutrinoless double beta decay  $(Z, A) \rightarrow (Z+2, A) + e^- + e^-$ . The best laboratory limit is  $t_{1/2} > 0.7 \times 10^{23}$  yr (CL=0.90) for  ${}^{76}\text{Ge}$ .

**b) Conversion of one lepton type to another.** For purely leptonic processes, the best limits are on  $\mu \rightarrow e \gamma$  and  $\mu \rightarrow 3e$ , measured as  $\Gamma(\mu \rightarrow e \gamma) / \Gamma(\mu \rightarrow \text{all}) < 1.7 \times 10^{-10}$  and  $\Gamma(\mu \rightarrow 3e) / \Gamma(\mu \rightarrow \text{all}) < 2.4 \times 10^{-12}$ . For semileptonic processes, the best limit comes from the coherent conversion process in a muonic atom  $\mu^- +$

$(Z, A) \rightarrow e^- + (Z, A)$ , measured as  $\Gamma(\mu^- \text{Ti} \rightarrow e^- \text{Ti}) / \Gamma(\mu^- \text{Ti} \rightarrow \text{all}) < 1.6 \times 10^{-11}$ . Of special interest is the case in which the hadronic flavor also changes, as in  $K_L \rightarrow e \mu$  and  $K^+ \rightarrow \pi^+ e^+ \mu^-$ , measured as  $\Gamma(K_L \rightarrow e \mu) / \Gamma(K_L \rightarrow \text{all}) < 6 \times 10^{-6}$  and  $\Gamma(K^+ \rightarrow \pi^+ e^+ \mu^-) / \Gamma(K^+ \rightarrow \text{all}) < 5 \times 10^{-9}$ . Limits on the conversion of  $\tau$  into  $e$  or  $\mu$  are found in  $\tau$  decay and are much less stringent than those for  $\mu \rightarrow e$  conversion, e.g.,  $\Gamma(\tau \rightarrow \mu \gamma) / \Gamma(\tau \rightarrow \text{all}) < 5.5 \times 10^{-4}$  and  $\Gamma(\tau \rightarrow e \gamma) / \Gamma(\tau \rightarrow \text{all}) < 6.4 \times 10^{-4}$ .

**c) Conversion of one type of lepton into another type of antilepton.** The case most studied is  $\mu^- + (Z, A) \rightarrow e^+ + (Z-2, A)$ , the strongest limit being  $\Gamma(\mu^- {}^{127}\text{I} \rightarrow e^+ {}^{127}\text{Sb}^{\text{stable}}) / \Gamma(\mu^- {}^{127}\text{I} \rightarrow \text{all}) < 3 \times 10^{-10}$ .

**d) Relation to neutrino mass.** If neutrinos have masses then it is expected even in the standard electroweak theory that separate lepton numbers are not conserved as a consequence of lepton mixing analogous to the Cabibbo quark mixing. However, in this case lepton-number-violating processes such as  $\mu \rightarrow e \gamma$  are expected to have extremely small probability. For small neutrino masses, the lepton-number violation would be observed first in neutrino oscillations, which have been the subject of extensive experimental searches. For example,  $\bar{\nu}_e \rightarrow \bar{\nu}_\mu$  gives measured limits  $\Delta(m^2) < 0.016$  eV<sup>2</sup> for  $\sin^2(2\theta) = 1$ , and  $\sin^2(2\theta) < 0.16$  for large  $\Delta(m^2)$ , where  $\theta$  is the neutrino mixing angle. Searches for  $\nu_\mu \rightarrow \nu_e$  set limits  $\Delta(m^2) < 0.2$  eV<sup>2</sup> for  $\sin^2(2\theta) = 1$ , and  $\sin^2(2\theta) < 0.0034$  for large  $\Delta(m^2)$ . For larger neutrino masses ( $\gg 1$  keV), lepton-number violation is searched for by looking for anomalous decays such as  $\pi \rightarrow e \nu_x$ , where  $\nu_x$  is a massive neutrino. If the  $\Delta L = 2$  type of violation occurs, it is expected that neutrinos will have a nonzero mass of the Majorana type.

### CONSERVATION OF HADRONIC FLAVORS

The conversion of quarks of a given charge,  $(d, s, b)$  or  $(u, c, t)$ , into one another is forbidden in strong and electromagnetic interactions by the conservation of hadron flavors: S (strangeness), C (charm), B (bottomness), and T (topness). The weak interactions violate these conservation laws as a result of the Cabibbo or Kobayashi-Maskawa mixing (see the section on the Kobayashi-Maskawa Mixing Matrix). The way in which these conservation laws are violated is tested as follows:

**a)  $\Delta S = \Delta Q$  rule.** In the semileptonic decay of strange particles, the strangeness change equals the change in charge of the hadrons. Tests come from limits on decay rates such as  $\Gamma(\Sigma^+ \rightarrow ne^+ \nu) / \Gamma(\Sigma^+ \rightarrow \text{all}) < 5 \times 10^{-6}$ , and from a detailed analysis of  $K_L \rightarrow \pi e \nu$ , which yields the parameter  $x$ , measured to be  $(\text{Re}x, \text{Im}x) = (0.009 \pm 0.020, -0.004 \pm 0.026)$ . A corresponding rule for charm decays is  $\Delta C = \Delta Q$ .

**b) Change of flavor by 2 units.** In the standard model this occurs only in second-order weak interactions. The one example for which this has been measured is the  $\Delta S = 2$   $K^0 - \bar{K}^0$  mixing, which is directly measured by  $m(K_S) - m(K_L) = (3.521 \pm 0.014) \times 10^{-12}$  MeV. A limit on the  $\Delta C = 2$   $D^0 - \bar{D}^0$  mixing,  $\Gamma(D^0 \rightarrow \bar{D}^0 \rightarrow \mu^- \text{ anything}) / \Gamma(D^0 \rightarrow \mu^+ \text{ anything}) < 0.044$ , provides a limit  $|m(D^0) - m(\bar{D}^0)| < 6.5 \times 10^{-10}$  MeV.

**c) Flavor-changing neutral currents.** In the standard model the neutral-current interactions do not change flavor. The low rate  $\Gamma(K_L \rightarrow \mu^+ \mu^-) / \Gamma(K_L \rightarrow \text{all}) = (9.1 \pm 1.9) \times 10^{-9}$  puts limits on such interactions; the nonzero value for this rate is attributed to a combination of the weak and electromagnetic interactions. The best test should come from a limit on  $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ , which occurs in the standard model only as a second-order weak process with a branching fraction of  $10^{-10}$  to  $10^{-11}$ . Limits for charm-changing or bottom-changing neutral currents are much less stringent:  $\Gamma(D^0 \rightarrow \mu^+ \mu^-) / \Gamma(D^0 \rightarrow \text{all}) < 3.4 \times 10^{-4}$  and  $\Gamma(B^0 \rightarrow \mu^+ \mu^-) / \Gamma(B^0 \rightarrow \text{all}) < 2 \times 10^{-4}$ .

\* Revised April 1986 by T.G. Trippe and L. Wolfenstein.

## TESTS OF CONSERVATION LAWS (Cont'd)

## Discrete Space-Time Symmetries

Quantity <sup>(a)</sup>	Value <sup>(b)</sup>	Symmetry tested or violated
$\pi^0 \rightarrow \gamma\gamma$ / all	$< 3.8 \times 10^{-7}$	C
$(e^+e^-)_J = 1 \rightarrow \gamma\gamma$ / all	$(5 \pm 3) \times 10^{-4(c)}$	C
$\eta \rightarrow \gamma\gamma$ / all	$< 7 \times 10^{-4}$	C
$\eta \rightarrow e^+e^-\pi^0$ / all	$< 5 \times 10^{-5}$	C (single photon process)
$\eta \rightarrow \mu^+\mu^-\pi^0$ / all	$< 5 \times 10^{-6}$	C (single photon process)
$\eta \rightarrow \pi^+\pi^-\pi^0$ parameters: left-right asymmetry	$(1.2 \pm 1.7) \times 10^{-3}$	C
sextant asymmetry	$(1.9 \pm 1.6) \times 10^{-3}$	C
quadrant asymmetry	$(-1.7 \pm 1.7) \times 10^{-3}$	C
$\eta \rightarrow \pi^+\pi^-\gamma$ parameters: left-right asymmetry	$(8.8 \pm 4.0) \times 10^{-3}$	C
$\beta$ (D-wave)	$0.047 \pm 0.062$	C
$\eta \rightarrow \pi^+\pi^-$ / all	$< 1.5 \times 10^{-3}$	P and CP
$e$ electric dipole moment	$< 3 \times 10^{-24}$ e cm	T and P
$\mu$ electric dipole moment	$(3.7 \pm 3.4) \times 10^{-19}$ e cm	T and P
$p$ electric dipole moment	$< 4 \times 10^{-21}$ e cm	T and P
$n$ electric dipole moment	$(2.3 \pm 2.3) \times 10^{-25}$ e cm	T and P
$\Delta$ electric dipole moment	$< 1.5 \times 10^{-16}$ e cm	T and P
$\alpha'/a$ from $\mu \rightarrow e\bar{\nu}\nu$	$(-0.2 \pm 4.3) \times 10^{-3}$	T
$\beta'/a$ from $\mu \rightarrow e\bar{\nu}\nu$	$(1.5 \pm 6.3) \times 10^{-3}$	T
$e$ pol. $\perp$ $\mu$ spin and $e^+$ mom. from $\mu \rightarrow e\bar{\nu}\nu$	$0.007 \pm 0.023$	T
$\text{Im } \xi$ in $K_{\mu 3}^0$ decay (from transverse $\mu$ pol.)	$-0.017 \pm 0.025$	T
$\text{Im } \xi$ in $K_{\mu 3}^0$ decay (from transverse $\mu$ pol.)	$-0.020 \pm 0.022$	T
$\phi(g_A) - \phi(g_V)$ for $n$	$(180.11 \pm 0.17)^\circ$	T ( $0^\circ$ or $180^\circ$ )
$n$ 3-vector corr. coeff.	$-0.0007 \pm 0.0014$	T
$K^\pm \rightarrow \pi^\pm\pi^+\pi^-$ rate difference / average	$(0.07 \pm 0.12)\%$	CP
$K^\pm \rightarrow \pi^\pm 2\pi^0$ rate difference / average	$(-0.03 \pm 0.55)\%$	CP
$K^\pm \rightarrow \pi^\pm\pi^0\gamma$ rate difference / average	$(0.9 \pm 3.3)\%$	CP
$K \rightarrow 3\pi^\pm$ slope $(g^+ - g^-)$ / sum	$(-0.7 \pm 0.5)\%$	CP
$ \eta_{+-0} ^2 = \Gamma(K_S^0 \rightarrow \pi^+\pi^-\pi^0) / \Gamma(K_L^0 \rightarrow \pi^+\pi^-\pi^0)$	$< 0.12$	CP
$ \eta_{000} ^2 = \Gamma(K_S^0 \rightarrow 3\pi^0) / \Gamma(K_S^0 \rightarrow 3\pi^0)$	$< 0.1$	CP
Charge asymm. $j$ in $K_L^0 \rightarrow \pi^+\pi^-\pi^0$	$0.0011 \pm 0.0008$	CP
$K_L^0 \rightarrow (\pi^-\mu^+\nu - \pi^+\mu^-\bar{\nu})$ / sum	$(0.319 \pm 0.038)\%$	CP (violated)
$K_L^0 \rightarrow (\pi^-\mu^+\nu - \pi^+\mu^-\bar{\nu})$ / sum	$(0.333 \pm 0.014)\%$	CP (violated)
$ \eta_{00}  =  A(K_L^0 \rightarrow \pi^0\pi^0) / A(K_S^0 \rightarrow \pi^0\pi^0) $	$(2.299 \pm 0.036) \times 10^{-3}$	CP (violated)
$ \eta_{+-}  =  A(K_L^0 \rightarrow \pi^+\pi^-) / A(K_S^0 \rightarrow \pi^+\pi^-) $	$(2.275 \pm 0.021) \times 10^{-3}$	CP (violated)
$ \epsilon'/\epsilon  = (1 -  \eta_{00}/\eta_{+-} ) / 3$	$(-3 \pm 4) \times 10^{-3}$	CP
$\phi_{+-}$ : phase of $\eta_{+-}$	$(44.6 \pm 1.2)^\circ$	CP (violated)
$\phi_{00}$ : phase of $\eta_{00}$	$(54 \pm 5)^\circ$	CP (violated)
$\text{Re } \epsilon$ in $K_L^0$ decay	$(1.621 \pm 0.088) \times 10^{-3}$	CP (violated)
$\alpha_\lambda / \alpha_\Lambda$	$-1.04 \pm 0.29$	CP
$(g_e^+ - g_e^-)$ / average	$(2.2 \pm 6.4) \times 10^{-11}$	CPT
$(g_e^+ - g_e^-)$ / average	$(-2.6 \pm 1.6) \times 10^{-8}$	CPT
$(\mu_p^+ - \mu_p^-)$ / average	$(-1 \pm 7) \times 10^{-3}$	CPT
$\pi^+ - \pi^-$ mass difference / average	$(2 \pm 5) \times 10^{-4}$	CPT
$K^+ - K^-$ mass difference / average	$(-0.6 \pm 1.8) \times 10^{-4}$	CPT
$ K^0 - \bar{K}^0 $ mass difference / average	$< 6 \times 10^{-19}$	CPT
$p - \bar{p}$ mass difference / average	$(7 \pm 4) \times 10^{-5}$	CPT
$\Lambda - \bar{\Lambda}$ mass difference / average	$(7 \pm 7) \times 10^{-5}$	CPT
$\Xi^- - \bar{\Xi}^+$ mass difference / average	$(1.1 \pm 2.7) \times 10^{-4}$	CPT
$\Omega^- - \bar{\Omega}^+$ mass difference / average	$(-0.7 \pm 4.7) \times 10^{-4}$	CPT
$\mu^+ - \mu^-$ mean life difference / average	$(3 \pm 8) \times 10^{-5}$	CPT
$\pi^+ - \pi^-$ mean life difference / average	$(5 \pm 7) \times 10^{-4}$	CPT
$K^+ - K^-$ mean life difference / average	$(1.1 \pm 0.9) \times 10^{-3}$	CPT
$\Lambda - \bar{\Lambda}$ mean life difference / average	$(4.4 \pm 8.5) \times 10^{-2}$	CPT
$\Xi^- - \bar{\Xi}^+$ mean life difference / average	$(0.02 \pm 0.18)$	CPT
$K^\pm \rightarrow \mu^\pm\nu$ rate difference / average	$(-0.54 \pm 0.41)\%$	CPT
$K^\pm \rightarrow \pi^\pm\pi^0$ rate difference / average	$(0.8 \pm 1.2)\%$	CPT <sup>(d)</sup>

a. Branching fractions are described by a shorthand notation. e.g., " $\mu^+ \rightarrow e^+\gamma$ /all" means  $\Gamma(\mu^+ \rightarrow e^+\gamma) / \Gamma(\mu^+ \rightarrow \text{all})$ .

b. Limits are given at 90% confidence level while errors are given as  $\pm 1$  standard deviation.

c. Orthopositronium data are from Liu and Roberts, Phys. Rev. Lett. **16**, 67 (1966).

d. Neglecting photon channels. See, e.g., A. Pais and S.B. Treiman, Phys. Rev. **D12**, 2744 (1975).

**TESTS OF CONSERVATION LAWS (Cont'd)**  
**Number Conservation Laws**

Quantity <sup>(a)</sup>	Value <sup>(b)</sup>	Conservation law tested
$\mu^+ \rightarrow e^+ \bar{\nu}_e \nu_\mu$ / all	$< 5 \times 10^{-2}$	Lepton family number <sup>(c,d)</sup>
$\rightarrow e^+ \gamma$ / all	$< 1.7 \times 10^{-10}$	Lepton family number <sup>(d)</sup>
$\rightarrow e^+ e^+ e^-$ / all	$< 2.4 \times 10^{-12}$	" " "
$\rightarrow e^+ \gamma \gamma$ / all	$< 8.4 \times 10^{-9}$	" " "
$\mu^- \text{ } ^{32}\text{S} \rightarrow e^- \text{ } ^{32}\text{S}$ / all	$< 7 \times 10^{-11}$	" " "
$\mu^- \text{Cu} \rightarrow e^- \text{Cu}$ / all	$< 1.6 \times 10^{-8}$	" " "
$\mu^- \text{Ti} \rightarrow e^- \text{Ti}$ / all	$< 1.6 \times 10^{-11}$	" " "
coupling for $(\mu^+ e^- \rightarrow \mu^- e^+)_{\text{bound}}$	$< 42 G_F$	" " "
$\tau^+ \rightarrow \mu^+ \gamma$ / all	$< 5.5 \times 10^{-4}$	" " "
$\rightarrow e^+ \gamma$ / all	$< 6.4 \times 10^{-4}$	" " "
$\rightarrow \mu^+ \mu^+ \mu^-$ / all	$< 4.9 \times 10^{-4}$	" " "
$\rightarrow e^+ \mu^+ \mu^-$ / all	$< 3.3 \times 10^{-4}$	" " "
$\rightarrow \mu^+ e^+ e^-$ / all	$< 4.4 \times 10^{-4}$	" " "
$\rightarrow e^+ e^+ e^-$ / all	$< 4.0 \times 10^{-4}$	" " "
$\rightarrow \mu^+ \pi^0$ / all	$< 8.2 \times 10^{-4}$	" " "
$\rightarrow e^+ \pi^0$ / all	$< 2.1 \times 10^{-3}$	" " "
$\rightarrow \mu^+ K^0$ / all	$< 1.0 \times 10^{-3}$	" " "
$\rightarrow e^+ K^0$ / all	$< 1.3 \times 10^{-3}$	" " "
$\rightarrow \mu^+ \rho^0$ / all	$< 4.4 \times 10^{-4}$	" " "
$\rightarrow e^+ \rho^0$ / all	$< 3.7 \times 10^{-4}$	" " "
$\pi^+ \rightarrow \mu^+ \nu_e$ / all	$< 8.0 \times 10^{-3(e)}$	" " "
$\pi^0 \rightarrow e \mu$ / all	$< 7 \times 10^{-8}$	" " "
$K^+ \rightarrow \pi^+ e^+ \mu^-$ / all	$< 7 \times 10^{-9}$	" " "
$\rightarrow \pi^+ e^- \mu^+$ / all	$< 5 \times 10^{-9}$	" " "
$\rightarrow \mu^+ \nu_e$ / all	$< 4 \times 10^{-3(e)}$	" " "
$\rightarrow \mu^- \nu_e e^+$ / all	$< 2 \times 10^{-8}$	" " "
$K_L^0 \rightarrow e \mu$ / all	$< 6 \times 10^{-6}$	" " "
$B^0 \rightarrow e \mu$ / all	$< 3 \times 10^{-4}$	" " "
$\nu$ oscillations		
$\Delta(m^2)$ for $\sin^2(2\theta)=1$		
$\bar{\nu}_e \leftrightarrow \bar{\nu}_e$	$< 0.016 \text{ eV}^2$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_e$	$< 0.2 \text{ eV}^2$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_\mu$	$< 0.9 \text{ eV}^2$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_\tau$	$< 3 \text{ eV}^2$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_\tau$	$< 2.2 \text{ eV}^2$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_\mu$	$< 0.23 \text{ eV}^2$ or $> 1500 \text{ eV}^2$	" " "
$\bar{\nu}_e \leftrightarrow \bar{\nu}_e$	$< 2.3 \text{ eV}^2$	" " "
$\sin^2(2\theta)$ for large $\Delta(m^2)$		
$\bar{\nu}_e \leftrightarrow \bar{\nu}_e$	$< 0.16$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_e$	$< 0.0034$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_e$	$< 0.013$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_\tau$	$< 0.013$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_\tau$	$< 0.044$	" " "
$\bar{\nu}_\mu \leftrightarrow \bar{\nu}_\mu$	$< 0.02 [\Delta(m^2) = 110 \text{ eV}^2]$	" " "
$\bar{\nu}_e \leftrightarrow \bar{\nu}_e$	$< 0.07$	" " "
$\bar{\nu}_e \leftrightarrow \bar{\nu}_\tau$	$< 0.7$	" " "
For other lepton mixing effects in particle decays, see Full Listings.		
$\mu^- \text{ } ^{32}\text{S} \rightarrow e^+ \text{ } ^{32}\text{S}^*$ / all	$< 9 \times 10^{-10}$	Total lepton number <sup>(f)</sup>
$\mu^- \text{ } ^{127}\text{I} \rightarrow e^+ \text{ } ^{127}\text{Sb}^{\text{stable}}$ / all	$< 3 \times 10^{-10}$	" " "
$\mu^- \text{Cu} \rightarrow e^+ \text{Co}$ / all	$< 2.6 \times 10^{-8}$	" " "
$\pi^+ \rightarrow \mu^+ \bar{\nu}_e$ / all	$< 1.5 \times 10^{-3(e)}$	" " "
$K^+ \rightarrow \pi^- e^+ e^+$ / all	$< 1 \times 10^{-8}$	" " "
$\rightarrow \pi^- e^+ \mu^+$ / all	$< 7 \times 10^{-9}$	" " "
$\rightarrow \mu^+ \bar{\nu}_e$ / all	$< 3.3 \times 10^{-3(e)}$	" " "
$\rightarrow e^+ \pi^0 \bar{\nu}_e$ / all	$< 3 \times 10^{-3(e)}$	" " "
neutrinoless double beta decay	See Full Listings	" " "

**TESTS OF CONSERVATION LAWS (Cont'd)**  
**Number Conservation Laws (Cont'd)**

Quantity <sup>(a)</sup>	Value <sup>(b)</sup>	Conservation law tested
$\tau_p / \text{BR}(p \rightarrow e^+ \pi^0)$	$> 2.5 \times 10^{32}$ years	Baryon number
$\tau_p / \text{BR}(p \rightarrow \mu^+ \pi^0)$	$> 1.0 \times 10^{32}$ years	" "
$\tau_p / \text{BR}(p \rightarrow e^+ K^0)$	$> 0.8 \times 10^{32}$ years	" "
$\tau_p / \text{BR}(p \rightarrow \mu^+ K^0)$	$> 0.4 \times 10^{32}$ years	" "
For other nucleon decay channels, see Stable Particle Summary Table.		
mean time for $n \rightarrow \bar{n}$ transition	$> 3$ years	" "
$e$ mean life	$> 2 \times 10^{22}$ years	Charge
$n \rightarrow p \bar{\nu} / p e^- \bar{\nu}$	$< 9 \times 10^{-24}$	" "
Re x from $K^0 \rightarrow \pi e \nu$	$0.009 \pm 0.020$	$\Delta S = \Delta Q$ (g)
Im x from $K^0 \rightarrow \pi e \nu$	$-0.004 \pm 0.026$	" "
$K^+ \rightarrow \pi^+ \pi^+ e^- \bar{\nu} / \text{all}$	$< 1.2 \times 10^{-8}$	" "
$\rightarrow \pi^+ \pi^+ \mu^- \bar{\nu} / \text{all}$	$< 3 \times 10^{-6}$	" "
$\Sigma^+ \rightarrow n e^+ \nu / \text{all}$	$< 5 \times 10^{-6}$	" "
$\rightarrow n \mu^+ \nu / \text{all}$	$< 3 \times 10^{-5}$	" "
$(\Sigma^+ \rightarrow n e^+ \nu) / (\Sigma^- \rightarrow n e^- \bar{\nu})$	$< 0.04$	" "
$\Xi^0 \rightarrow \Sigma^- e^+ \nu / \text{all}$	$< 9 \times 10^{-4}$	" "
$\rightarrow \Sigma^- \mu^+ \nu / \text{all}$	$< 9 \times 10^{-4}$	" "
$\rightarrow p e^- \bar{\nu} / \text{all}$	$< 1.3 \times 10^{-3}$	$\Delta S = 2$ forbidden <sup>(g)</sup>
$\rightarrow p \mu^- \bar{\nu} / \text{all}$	$< 1.3 \times 10^{-3}$	" "
$\Xi^- \rightarrow n e^- \bar{\nu} / \text{all}$	$< 3.2 \times 10^{-3}$	" "
$\rightarrow n \mu^- \bar{\nu} / \text{all}$	$< 1.5 \times 10^{-2}$	" "
$\rightarrow p \pi^- e^- \bar{\nu} / \text{all}$	$< 4 \times 10^{-4}$	" "
$\rightarrow p \pi^- \mu^- \bar{\nu} / \text{all}$	$< 4 \times 10^{-4}$	" "
$\Xi^0 \rightarrow p \pi^- / \text{all}$	$< 3.6 \times 10^{-5}$	" "
$\Xi^- \rightarrow n \pi^- / \text{all}$	$< 1.9 \times 10^{-5}$	" "
$\rightarrow p \pi^- \pi^- / \text{all}$	$< 4 \times 10^{-4}$	" "
$\Omega^- \rightarrow \Lambda \pi^- / \text{all}$	$< 1.9 \times 10^{-4}$	" "
$m_{K_L} - m_{K_S}$	$(3.521 \pm 0.014) \times 10^{-12}$ MeV	" "
$(D^0 \rightarrow \bar{D}^0 \rightarrow K^+ \pi^-) / (D^0 \rightarrow K \pi)$	$< 0.08$	$\Delta C = 2$ forbidden <sup>(g)</sup>
$(D^0 \rightarrow \bar{D}^0 \rightarrow \mu^- \text{ anything}) / (D^0 \rightarrow \mu^+ \text{ anything})$	$< 0.044$	" "
$ m_{D_1^0} - m_{D_2^0} $ (from previous limit)	$< 6.5 \times 10^{-10}$ MeV	" "
$(B^0 \rightarrow \bar{B}^0 \text{ via mixing}) / (B^0 \rightarrow \text{all})$	$< 0.12$	$\Delta B = 2$ forbidden <sup>(g)</sup>
$K_L^0 \rightarrow \mu^+ \mu^- / \text{all}$	$(9.1 \pm 1.9) \times 10^{-9}$	no flav. chng. neut. curr.
$\rightarrow e^+ e^- / \text{all}$	$< 2.0 \times 10^{-7}$	" " " " "
$\rightarrow \mu^+ \mu^- \gamma / \text{all}$	$(2.8 \pm 2.8) \times 10^{-7}$	" " " " "
$\rightarrow e^+ e^- \gamma / \text{all}$	$(1.7 \pm 0.9) \times 10^{-5}$	" " " " "
$\rightarrow \pi^0 \mu^+ \mu^- / \text{all}$	$< 1.2 \times 10^{-6}$	" " " " "
$\rightarrow \pi^0 e^+ e^- / \text{all}$	$< 2.3 \times 10^{-6}$	" " " " "
$\rightarrow \pi^+ \pi^- e^+ e^- / \text{all}$	$< 9 \times 10^{-6}$	" " " " "
$\rightarrow \mu^+ \mu^- e^+ e^- / \text{all}$	$< 4.9 \times 10^{-6}$	" " " " "
$\rightarrow e^+ e^- e^+ e^- / \text{all}$	$< 2.6 \times 10^{-6}$	" " " " "
$K_S^0 \rightarrow \mu^+ \mu^- / \text{all}$	$< 3.2 \times 10^{-7}$	" " " " "
$\rightarrow e^+ e^- / \text{all}$	$< 3.4 \times 10^{-4}$	" " " " "
$K^+ \rightarrow \pi^+ e^+ e^- / \text{all}$	$(2.7 \pm 0.5) \times 10^{-7}$	" " " " "
$\rightarrow \pi^+ \mu^+ \mu^- / \text{all}$	$< 2.4 \times 10^{-6}$	" " " " "
$\rightarrow \pi^+ \nu \bar{\nu} / \text{all}$	$< 1.4 \times 10^{-7}$	" " " " "
$D^0 \rightarrow \mu^+ \mu^-$	$< 3.4 \times 10^{-4}$	" " " " "
$B^0 \rightarrow e^+ e^- / \text{all}$	$< 3 \times 10^{-4}$	" " " " "
$\rightarrow \mu^+ \mu^- / \text{all}$	$< 2 \times 10^{-4}$	" " " " "
$\Sigma^+ \rightarrow p e^+ e^- / \text{all}$	$< 7 \times 10^{-6}$	" " " " "

- a. Branching fractions are described by a shorthand notation, e.g., " $\mu^+ \rightarrow e^+ \gamma / \text{all}$ " means  $\Gamma(\mu^+ \rightarrow e^+ \gamma) / \Gamma(\mu^+ \rightarrow \text{all})$ .  
b. Limits are given at 90% confidence level while errors are given as  $\pm 1$  standard deviation.  
c. Test of additive vs. multiplicative lepton family number conservation.  
d. Lepton family number conservation means separate conservation of  $e$ -number,  $\mu$ -number, and  $\tau$ -number.  
e. These limits are derived from the analysis of neutrino oscillation experiments.  
f. Violation of total lepton number conservation also implies violation of lepton family number conservation.  
g. Can be violated in second-order weak interactions.

## NONRELATIVISTIC QUARK MODEL

### A. QUANTUM NUMBERS

Each quark has spin 1/2. The additive quantum numbers (other than baryon number = 1/3) of the known (and presumed) quarks are shown in the table.

Quantum number	Quark type (flavor)					
	<i>d</i>	<i>u</i>	<i>s</i>	<i>c</i>	<i>b</i>	<i>t</i>
Q — electric charge	$-\frac{1}{3}$	$+\frac{2}{3}$	$-\frac{1}{3}$	$+\frac{2}{3}$	$-\frac{1}{3}$	$+\frac{2}{3}$
$I_z$ — z-component of isospin	$-\frac{1}{2}$	$+\frac{1}{2}$	0	0	0	0
S — strangeness	0	0	-1	0	0	0
C — charm	0	0	0	+1	0	0
B — bottomness	0	0	0	0	-1	0
T — topness	0	0	0	0	0	+1

With these conventions the strangeness *S* of the  $K^+$  is +1 and the bottomness *B* of the  $B^+$  is +1.

The *G*-parity operator is defined to be  $G = C e^{-i\pi I_y}$ , where *C* is the charge conjugation operator. The mesons with  $S = C = B = T = 0$  are eigenstates of *G*. If a meson is also an eigenstate of the charge conjugation operator with charge conjugation *C*, then  $G = C(-1)^I$ , where *I* is its isospin; all the other particles in the same isomultiplet have the same value of *G*:  $G(\pi^\pm) = G(\pi^0) = -1$ ,  $G(\rho^\pm) = G(\rho^0) = +1$ , etc.

### B. MESONS

Nearly all known mesons can be understood as bound states of a quark *q* and an antiquark  $\bar{q}'$  (the flavors of *q* and  $q'$  may be different). If the orbital angular momentum of the  $q\bar{q}'$  state is *L*, then the parity  $P = (-1)^{L+1}$ . A state  $q\bar{q}'$  of a quark and its own antiquark is also an eigenstate of charge conjugation with  $C = (-1)^{L+S}$ , where the spin  $S = 0$  or 1. The  $L = 0$  states are the pseudoscalars,  $J^P = 0^-$ , and the vectors,  $J^P = 1^-$ . See table below.

Standard quark model assignments for some of the known mesons. Some assignments, especially for  $0^{++}$ , are controversial. Note that only the states in the  $u\bar{u}$ ,  $d\bar{d}$ ,  $s\bar{s}$ ,  $c\bar{c}$ , and  $b\bar{b}$  columns and the neutral states in the  $I = 1$  column are eigenstates of charge conjugation *C*.

$2S+1L_J$	$J^{PC}$	$u\bar{d}, u\bar{u}, d\bar{d}$ <i>I</i> = 1	$u\bar{u}, d\bar{d}, s\bar{s}$ <i>I</i> = 0	$c\bar{c}$ <i>I</i> = 0	$b\bar{b}$ <i>I</i> = 0	$s\bar{u}, s\bar{d}$ <i>I</i> = 1/2	$c\bar{u}, c\bar{d}$ <i>I</i> = 1/2	$c\bar{s}$ <i>I</i> = 0	$b\bar{u}, b\bar{d}$ <i>I</i> = 1/2
$^1S_0$	$0^{-+}$	$\pi$	$\eta, \eta'$	$\eta_c$		<i>K</i>	<i>D</i>	$D_s$	<i>B</i>
$^3S_1$	$1^{--}$	$\rho$	$\phi, \omega$	$J/\psi$	$\Upsilon$	$K^*(892)$	$D^*(2010)$		
$^1P_1$	$1^{+-}$	$b_1(1235)$	$h_1(1190)$			$K_{1B}$			
$^3P_0$	$0^{++}$	$a_0(980)$	$f_0(975), f_0(1300)$	$\chi_0(3415)$	$\chi_{b0}(9860)$	$K_0^*(1350)$			
$^3P_1$	$1^{++}$	$a_1(1270)$	$f_1(1285), f_1(1420)$	$\chi_1(3510)$	$\chi_{b1}(9895)$	$K_{1A}$			
$^3P_2$	$2^{++}$	$a_2(1320)$	$f_2'(1525), f_2(1270)$	$\chi_2(3555)$	$\chi_{b2}(9915)$	$K_2^*(1430)$			
$^1D_2$	$2^{-+}$	$\pi_2(1680)$							
$^3D_1$	$1^{--}$			$\psi(3770)$					
$^3D_2$	$2^{--}$					$K_2(1770)$			
$^3D_3$	$3^{--}$	$\rho_3(1690)$	$\omega_3(1670)$			$K_3^*(1780)$			

States in the "normal" spin-parity series,  $P = (-1)^J$ , must, according to the above, have  $S = 1$  and hence  $CP = +1$ . Thus mesons with normal spin-parity and  $CP = -1$  are forbidden in the  $q\bar{q}'$  quark model. The  $J^{PC} = 0^{--}$  state is forbidden as well. Mesons with such  $J^{PC}$  could exist, but would lie outside the  $q\bar{q}'$  model.

States with the same  $J^P$  and additive quantum numbers can mix (if they are eigenstates of charge conjugation, they must also have the same value of *C*). Thus the physical  $J^P = 1^+$ , strangeness  $S = 1$  states,  $K_{1A}(1280)$  and  $K_{1B}(1400)$ , are mixtures of the pure quark model states  $K_{1A}$  and  $K_{1B}$ . The  $\psi(3770)$  is a mixture of  $^3S_1$  and  $^3D_1$ . The  $\eta$  and  $\eta'$  are mixtures of the SU(3) octet and singlet states.

For the pseudoscalar mesons, the Gell-Mann-Okubo formula is

$$m_\eta^2 = \frac{1}{3}(4m_K^2 - m_\pi^2),$$

assuming no octet-singlet mixing. However, the octet  $\eta_8$  and singlet  $\eta_1$  mix because of SU(3) breaking. The physical states  $\eta$  and  $\eta'$  are given by

$$\eta = \eta_8 \cos \theta_P - \eta_1 \sin \theta_P$$

$$\eta' = \eta_8 \sin \theta_P + \eta_1 \cos \theta_P.$$

These combinations diagonalize the mass-squared matrix

$$M^2 = \begin{pmatrix} M_{11}^2 & M_{18}^2 \\ M_{18}^2 & M_{88}^2 \end{pmatrix},$$

where  $M_{88}^2 = \frac{1}{3}(4m_K^2 - m_\pi^2)$ . It follows that

$$\tan^2 \theta_P = \frac{M_{88}^2 - m_\eta^2}{m_\eta^2 - M_{88}^2}.$$

The sign of  $\theta_P$  is meaningful in the quark model. If

$$\eta_1 = (u\bar{u} + d\bar{d} + s\bar{s})/\sqrt{3}$$

$$\eta_8 = (u\bar{u} + d\bar{d} - 2s\bar{s})/\sqrt{6},$$

then the matrix element  $M_{18}^2$ , which is due mostly to the strange

### NONRELATIVISTIC QUARK MODEL (Cont'd)

quark mass, is negative. From the relation

$$\tan \theta_P = \frac{M_{88}^2 - m_\eta^2}{M_{18}^2}$$

we find  $\theta_P < 0$ .

For the vector mesons we replace  $\pi \rightarrow \rho$ ,  $K \rightarrow K^*$ ,  $\eta \rightarrow \phi$ , and  $\eta' \rightarrow \omega$ , so

$$\phi = \omega_8 \cos \theta_V - \omega_1 \sin \theta_V$$

$$\omega = \omega_8 \sin \theta_V + \omega_1 \cos \theta_V$$

For "ideal mixing,"  $\phi = s\bar{s}$ ,  $\tan \theta_V = 1/\sqrt{2}$ , so  $\theta_V \approx 35.3^\circ$ . Experimentally,  $\theta_V$  is near  $35^\circ$ , the sign being determined by a formula analogous to that for  $\tan \theta_P$ . Following this procedure we find the mixing angles below. There are uncertainties of a few degrees arising from electromagnetic mass splittings and uncertainties in resonance masses.

Singlet-octet mixing for the pseudoscalar, vector, and tensor mesons. The sign conventions are as above. The value of  $\theta_{\text{quad}}$  is obtained from the equations above, and  $\theta_{\text{lin}}$  is obtained by replacing  $m^2 \rightarrow m$  throughout. Of the two isosinglets, the mostly octet one is listed first.

$J^{PC}$	Nonet Members	$\theta_{\text{quad}}$	$\theta_{\text{lin}}$
$0^{-+}$	$\pi, K, \eta, \eta'$	$-10^\circ$	$-23^\circ$
$1^{--}$	$\rho, K^*(892), \phi, \omega$	$39^\circ$	$36^\circ$
$2^{++}$	$a_2(1320), K_2^*(1430), f_2'(1525), f_2(1270)$	$28^\circ$	$26^\circ$
$3^{--}$	$\rho_3(1690), K_3^*(1780), \phi_J(1850), \omega_3(1670)$	$29^\circ$	$28^\circ$

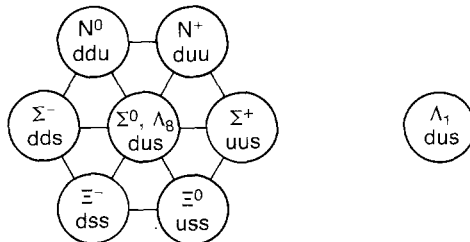
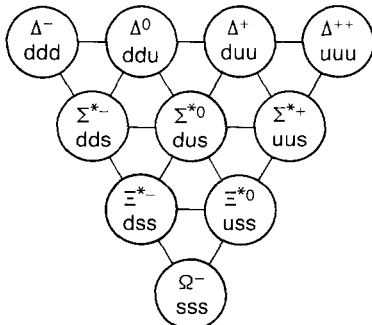
### C. BARYONS

All the established baryons are apparently 3-quark ( $qqq$ ) states, and each such state is an SU(3) color singlet, a completely antisymmetric state of the three possible colors. Since the quarks are fermions, the state function for any baryon must be antisymmetric under interchange of any two of its quarks. Thus the state is *symmetric* under interchange of the quantum labels other than color:

$$|qqq\rangle_A = |\text{color}\rangle_A \times |\text{space, spin, flavor}\rangle_S$$

where the subscripts  $S$  and  $A$  indicate symmetry or antisymmetry under interchange of any two of the quarks. Note the contrast with the state function for the three nucleons in  $^3\text{H}$  or  $^3\text{He}$ :

$$|NNN\rangle_A = |\text{space, spin, isospin}\rangle_A$$



This difference has major implications for internal structure, magnetic moments, etc. (For a nice discussion, see Ref. 1.)

Few of the baryons containing  $c$  or heavier quarks have yet been discovered, so we restrict further attention to baryons made up of just  $d, u,$  and  $s$  quarks. The three flavors imply a flavor SU(3), which requires that baryons made of these quarks belong to the multiplets on the right side of

$$3 \otimes 3 \otimes 3 = 10_S \oplus 8_M \oplus 8_M \oplus 1_A$$

(see the section on SU( $n$ ) Multiplets and Young Diagrams). Here the subscripts indicate symmetric, mixed-symmetry, or antisymmetric states under interchange of any two quarks. The figure shows particle assignments in these multiplets. States  $\Lambda_8$  and  $\Lambda_1$  that have the same spin and parity can mix; an example is the mainly octet  $D_{03} \Delta(1690)$  and mainly singlet  $D_{03} \Delta(1520)$ . The formalism is the same as for  $\eta-\eta'$  or  $\phi-\omega$  mixing (see above), except that for baryons the mass  $M$  instead of  $M^2$  is used. The section SU(3) Isoscalar Factors shows how relative decay rates in, say,  $10 \rightarrow 8 \otimes 8$  decays may be calculated. A summary of results of fits to the observed baryon masses and decay rates for the best-known SU(3) multiplets is given in Appendix II of our 1982 edition.<sup>2</sup>

Flavor and spin may be combined in a flavor-spin SU(6) in which the six basic states are  $d \uparrow, d \downarrow, \dots, s \downarrow$  ( $\uparrow, \downarrow =$  spin up, down). Then the baryons belong to the multiplets on the right side of

$$6 \otimes 6 \otimes 6 = 56_S \oplus 70_M \oplus 70_M \oplus 20_A$$

These SU(6) multiplets decompose into flavor SU(3) multiplets as follows:

$$56 = 4^{10} \oplus 2^8$$

$$70 = 2^{10} \oplus 4^8 \oplus 2^8 \oplus 2^1$$

$$20 = 2^8 \oplus 4^1$$

where the superscript ( $2S+1$ ) gives the net spin  $S$  of the quarks for each particle in the SU(3) multiplet. The  $J^P = 1/2^+$  octet containing the nucleon and the  $J^P = 3/2^+$  decuplet containing the  $\Delta(1232)$  together make up the "ground-state" 56-plet in which the orbital angular momenta between the quarks are zero (so that the spatial part of the state function is trivially symmetric). The 70 and 20 require some excitation of the spatial part of the state function in order to make the overall state function symmetric.

The quark model for baryons is extensively reviewed in Ref. 3.

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## QCD

### A. THE QCD COUPLING CONSTANT

The QCD coupling constant  $\alpha_s(Q^2)$  has a  $Q^2$  dependence given by

$$\frac{d\alpha_s(Q^2)}{d\ln(Q^2)} = -\frac{\alpha_s(Q^2)}{4\pi} \left[ \beta_0 + \frac{\alpha_s(Q^2)}{4\pi} \beta_1 + \dots \right]$$

where

$$\beta_0 = 11 - \frac{2n_f}{3}, \quad \beta_1 = 102 - \frac{38n_f}{3},$$

and  $n_f$  is the number of light quark flavors, i.e., those with masses much less than  $Q/2$ . This equation can be solved given a boundary condition: the value of  $\alpha_s(Q^2)$  at some arbitrary scale  $Q_0$ . This is the only parameter of QCD apart from quark masses. The solution to this equation in lowest order, i.e., retaining only the first term on the right-hand side of the equation, is

$$\alpha_s(Q^2) = \frac{12\pi}{(33 - 2n_f)\ell n(Q^2/\Lambda^2)}.$$

Here a parameter  $\Lambda$  has been introduced. It replaces the boundary condition  $\alpha_s(Q_0^2)$ .

To next order the solution is obtained by multiplying the right-hand side of the above equation by

$$\left\{ 1 - \frac{6(153 - 19n_f)\ell n[\ell n(Q^2/\Lambda^2)]}{(33 - 2n_f)^2\ell n(Q^2/\Lambda^2)} \right\}.$$

### B. THE QCD PARTON MODEL

The structure functions of deep inelastic scattering given in Secs. C.5.b.i and C.5.b.ii of Cross Sections, Decays, Structure Functions, and Kinematics can be written in terms of quark distribution functions in the QCD parton model.  $q(x, Q^2)dx$  is the probability that a parton (quark, antiquark, or gluon) carries a momentum fraction between  $x$  and  $x + dx$  of the nucleon's momentum in a frame where the nucleon's momentum is large. The energy scale  $Q$  is the invariant mass of the virtual probe (photon,  $W$ , or  $Z$ ) which hits the nucleon. The structure functions referred to in Cross Sections, Decays, Structure Functions, and Kinematics, Sec. C.5.b, are given by

$$F_2^{\mu\text{CC}} = 2x [d(x) + s(x) + \bar{u}(x) + \bar{c}(x)]$$

$$xF_3^{\mu\text{CC}} = 2x [d(x) + s(x) - \bar{u}(x) - \bar{c}(x)]$$

$$F_2^{\bar{\nu}\text{CC}} = 2x [u(x) + c(x) + \bar{d}(x) + \bar{s}(x)]$$

$$xF_3^{\bar{\nu}\text{CC}} = 2x [u(x) + c(x) - \bar{d}(x) - \bar{s}(x)]$$

$$F_2^{\text{em}} = x \left[ \frac{4}{9} [u(x) + \bar{u}(x)] + \frac{1}{9} [d(x) + \bar{d}(x)] + \dots \right]$$

$$F_2^{\mu\text{NC}} = 2\rho^2 x \left\{ \left[ \frac{1}{4} - \frac{2}{3} \sin^2 \theta_W + \frac{8}{9} \sin^4 \theta_W \right] [u(x) + \bar{u}(x)] \right. \\ \left. + \left[ \frac{1}{4} - \frac{1}{3} \sin^2 \theta_W + \frac{2}{9} \sin^4 \theta_W \right] [d(x) + \bar{d}(x)] \right\}$$

$$xF_3^{\mu\text{NC}} = 2\rho^2 x \left\{ \left[ \frac{1}{4} - \frac{2}{3} \sin^2 \theta_W \right] [u(x) - \bar{u}(x)] \right. \\ \left. + \left[ \frac{1}{4} - \frac{1}{3} \sin^2 \theta_W \right] [d(x) - \bar{d}(x)] \right\}$$

$F_2 = 2xF_1$  in all cases (This is the Callan-Gross relation, and ignores parton transverse momentum and higher order QCD corrections.)

$$F_i^{\bar{\nu}\text{NC}} = F_i^{\mu\text{NC}},$$

where  $\rho = M_W^2/(M_Z^2 \cos^2 \theta_W)$ .

The  $Q^2$  dependence has been suppressed for the above distributions and structure functions. The  $x$  dependence of the structure functions for  $10 < Q^2 < 30 \text{ GeV}^2$  is shown in the Structure Functions Section of the Plots of Cross Sections and Related Quantities. The  $Q^2$  dependence of the structure functions is predicted by QCD and is controlled by  $\alpha_s$ . The dependence is logarithmic and consequently slow. It is also shown in the Structure Functions plots. Over the limited  $Q^2$  range shown relatively little net change is seen. The coupling constant, or  $\Lambda$ , can be determined from these data.<sup>1</sup>

The QCD parton model can also be applied to purely hadronic collisions. For example, the cross section for the production of two jets in a hadronic collision is given by

$$\sigma(s) = \sum_{i,j} \int f_i(x_1, Q^2) f_j(x_2, Q^2) \hat{\sigma}_{ij}(\hat{s}) dx_1 dx_2.$$

Here  $\hat{s} = sx_1x_2$  is the invariant mass squared of the parton-parton system and  $s$  is that of the proton-proton system. In this formula  $f_i(x, Q^2)$  is the probability of a parton of type  $i$  being inside the proton with fraction  $x$  of the proton's momentum. The sums  $i$  and  $j$  run over all partons: quarks, for which  $f_i(x, Q^2) = q(x, Q^2)$ ; antiquarks, for which  $f_i(x, Q^2) = \bar{q}(x, Q^2)$ ; and gluons for which  $f_i(x, Q^2) = g(x, Q^2)$ . The quantity  $\hat{\sigma}_{ij}$  is the cross section for the scattering of the two partons  $i$  and  $j$ . The scale  $Q$  appearing in the distribution functions is not well defined. It is characteristic of the momentum transfers in the partonic process and can be determined operationally only after computing higher order corrections.

In order to describe one-particle inclusive production in  $e^+e^-$  annihilation or deep inelastic scattering, it is convenient to introduce a fragmentation function  $D_i^h(z, Q^2)/z$  which is the probability that a parton of type  $i$  and momentum  $p$  will fragment into a hadron of type  $h$  and momentum  $zp$ . The  $Q^2$  evolution is predicted by QCD and is similar to that of the parton distribution functions (see above). The  $D_i^h(z, Q^2)$  are normalized so that

$$\sum_h \int D_i^h(z, Q^2) dz = 1.$$

If the contributions of the  $Z$  boson and three-jet events are neglected, the cross section for producing a hadron  $h$  in  $e^+e^-$  annihilation is given by

$$\frac{1}{\sigma_{\text{had}}} \frac{d\sigma}{dz} = \frac{\sum_i e_i^2 D_i^h(z, Q^2)}{\sum_i e_i^2},$$

where  $e_i$  is the charge of quark-type  $i$ ,  $\sigma_{\text{had}}$  is the total hadronic cross section, and the momentum of the hadron is  $zE_{\text{cm}}/2$ .

In the case of deep inelastic muon scattering, the cross section for producing a hadron of energy  $E_h$  is given by

$$\frac{1}{\sigma_{\text{tot}}} \frac{d\sigma}{dz} = \frac{\sum_i e_i^2 q_i(x, Q^2) D_i^h(z, Q^2)}{\sum_i e_i^2 q_i(x, Q^2)},$$

where  $E_h = \nu z$ . (For the kinematics of deep inelastic scattering, see section C.5 of the Cross Sections, Decays, Structure Functions, and

## QCD (Cont'd)

Kinematics section of this Review.) The fragmentation functions for light and heavy quarks have different  $z$  dependence; the former peak near  $z = 0$ . They are illustrated in a figure in the section on Plots of Cross Sections and Related Quantities.

### C. TESTS OF QCD

When a process is calculated beyond leading order in perturbative QCD, the coupling constant must be renormalized (defined). In the case of QED, the coupling constant can be defined as the strength of the electron-electron-photon vertex in the limit of zero photon momentum. Such a definition is related directly to a physical process, namely Thompson scattering. Such a definition is not possible in perturbative QCD since the coupling constant is too large for perturbation theory to be applicable to such a process. Rather, the coupling constant must be defined by some more formal procedure. Two such renormalization schemes are in common use: modified minimal subtraction ( $\overline{\text{MS}}$ )<sup>2</sup> and momentum space subtraction (MOM).<sup>3</sup> They are related by

$$\alpha_{\text{MOM}}(Q^2) = \alpha_{\overline{\text{MS}}}(Q^2) \left[ 1 + (1.86 - 0.48n_f) \alpha_{\overline{\text{MS}}}(Q^2) + \dots \right].$$

The corresponding  $\Lambda$ 's are then obtained from the equations of Sec. A above. If a process is calculated in perturbation theory, the coefficient of the next-to-leading term depends upon the procedure used. This coefficient can also be altered by changing the value of the scale  $Q^2$  at which the coupling constant is evaluated. Care should therefore be exercised in estimating how well the perturbation series is converging.

The parameter  $\alpha_s(Q_0^2)$  or  $\Lambda$  is extracted by comparing a QCD prediction with some data. The values of  $\alpha_s$  extracted from different experiments will agree if and only if the following criteria are satisfied:

- 1) The same renormalization scheme is used.
- 2) The perturbation series is convergent. As indicated above, the series may appear to be more convergent depending upon the choice of scheme and of  $Q^2$ . The choice of scheme is arbitrary. The choice of  $Q^2$  is sometimes clearly indicated by the physics, for example in the case of deep inelastic scattering. In other cases it is not clear; for example, in the case of jet production, it could be either of the Mandelstam variables  $t$  or  $s$  appropriate to the partonic process or the transverse momentum of the jet. In these ambiguous cases, it should be taken to be the value which minimizes the next-to-leading term in the perturbation series. In this event, a statement concerning the convergence of the perturbation expansion is not possible until the next-to-next-to-leading-order term has been computed. The size of the next-to-leading corrections varies widely from process to process. The corrections are largest in processes with a small  $Q$ , where  $\alpha_s$  is larger. There are corrections of order 30% to single  $W$  or  $Z$

production in proton-antiproton collisions at  $\sqrt{s} = 540$  GeV. By contrast the corrections to the total hadronic cross section in  $e^+e^-$  annihilation are only a few percent at  $\sqrt{s} = 30$  GeV.

3) The process must be free of so-called higher twist corrections. These are terms proportional to  $M^2/Q^2$ , where  $M$  is some mass scale of order 1 GeV. These corrections consequently become irrelevant as  $Q^2$  is increased. These corrections cannot be calculated in perturbative QCD. They afflict attempts to extract  $\alpha_s$  from deep inelastic scattering and from 3-jet events in  $e^+e^-$  annihilation. In the latter case, they are manifested in the different  $\alpha_s$  values which are extracted using different Monte-Carlos to describe the materialization of quarks and gluons into hadrons.

4) If there are quark thresholds in the region of  $Q^2$  being used for the measurement, these thresholds should be correctly dealt with. This involves removing kinematic corrections and taking care to indicate how  $\Lambda$  depends on the thresholds.

Among the processes used to determine  $\alpha_s$  are: deep inelastic scattering,<sup>4</sup> the total hadronic cross section in  $e^+e^-$  annihilation,<sup>5</sup> the distribution of 3-jet events in  $e^+e^-$  annihilation,<sup>6</sup> the flow of energy in  $e^+e^-$  annihilation,<sup>7</sup> the branching fractions  $\Upsilon \rightarrow \gamma + X$  and  $\Upsilon \rightarrow \mu^+\mu^-$ ,<sup>8</sup> the behavior of form factors,<sup>9</sup> measurements of the photon structure functions,<sup>10</sup> and the hyperfine splittings in the  $\psi$  system.<sup>11</sup>

In view of all the problems both theoretical and experimental, it is remarkable that all these different measurement agree so well and all are consistent with a value of  $\Lambda_{\overline{\text{MS}}}$  of order 200 MeV.

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## THE KOBAYASHI-MASKAWA MIXING MATRIX\*

In the "standard model" with  $SU(2) \times U(1)$  as the gauge group of electroweak interactions, both the quarks and leptons are assigned to be left-handed doublets and right-handed singlets. The quark mass eigenstates are not the same as the weak eigenstates, and the matrix connecting them has become known as the Kobayashi-Maskawa matrix<sup>1</sup> since an explicit parametrization in the six-quark case was first published by them in 1973.

By convention, the three charge  $2/3$  quarks ( $u$ ,  $c$ , and  $t$ ) are unmixed, and all the mixing is expressed in terms of a  $3 \times 3$  unitary matrix  $V$  operating on the charge  $-1/3$  quarks ( $d$ ,  $s$ ,  $b$ ):

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}. \quad (1)$$

The values of individual K-M matrix elements can in principle all be determined from weak decays of the relevant quarks, or, in some cases, from deep inelastic neutrino scattering. Using the constraints discussed below, together with unitarity, and assuming only three generations, the 90% confidence limits on the magnitude of the elements of the complete matrix are:

$$\begin{pmatrix} 0.9742 \text{ to } 0.9756 & 0.219 \text{ to } 0.225 & 0 & \text{to } 0.008 \\ 0.219 & \text{to } 0.225 & 0.973 \text{ to } 0.975 & 0.037 \text{ to } 0.053 \\ 0.002 & \text{to } 0.018 & 0.036 \text{ to } 0.052 & 0.9986 \text{ to } 0.9993 \end{pmatrix}. \quad (2)$$

The ranges shown are for the individual matrix elements. The constraints of unitarity connect different elements, so choosing a specific value for one element restricts the range of the others.

There are several parametrizations of the K-M matrix. The form due to Maiani<sup>2</sup> has a number of convenient properties:

$$V = \begin{pmatrix} c_\beta c_\gamma & c_\beta s_\gamma & s_\beta \\ -s_\gamma c_\beta s_\theta e^{i\delta'} - s_\theta c_\gamma & c_\gamma c_\beta - s_\gamma s_\beta s_\theta e^{i\delta'} & s_\gamma c_\beta e^{i\delta'} \\ -s_\beta c_\gamma c_\theta + s_\gamma s_\beta e^{-i\delta'} & -c_\gamma s_\beta s_\theta - s_\gamma c_\beta e^{-i\delta'} & c_\gamma c_\beta \end{pmatrix} \quad (3)$$

where  $\theta$ ,  $\beta$ ,  $\gamma$ , and  $\delta'$  are angles and  $c_\beta = \cos \beta$ ,  $s_\beta = \sin \beta$ , etc. With  $\beta = \gamma = 0$ , the first two generations of quarks decouple from the third, and  $\theta$  is directly the Cabibbo angle.

In view of the need for a "standard" parametrization in the literature, we propose this form and request public comment.

Kobayashi and Maskawa<sup>1</sup> chose a parametrization involving the four angles,  $\theta_1$ ,  $\theta_2$ ,  $\theta_3$ ,  $\delta$ :

$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} c_1 & -s_1 c_3 & -s_1 s_3 \\ s_1 c_2 & c_1 c_2 c_3 - s_2 s_3 e^{i\delta} & c_1 c_2 s_3 + s_2 c_3 e^{i\delta} \\ s_1 s_2 & c_1 s_2 c_3 + c_2 s_3 e^{i\delta} & c_1 s_2 s_3 - c_2 c_3 e^{i\delta} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}, \quad (4)$$

where  $c_i = \cos \theta_i$  and  $s_i = \sin \theta_i$  for  $i = 1, 2, 3$ . In the limit  $\theta_2 = \theta_3 = 0$ , this also reduces to the usual Cabibbo mixing with  $\theta_1$  identified (up to a sign) with the Cabibbo angle. The angles  $\theta_1$ ,  $\theta_2$ ,  $\theta_3$  can all be made to lie in the first quadrant (so that all  $s_i$ ,  $c_i$  are positive) by an appropriate redefinition of quark field phases.

Slightly different forms of the Kobayashi-Maskawa parametrization are found in the literature. The K-M matrix used in the 1982 Review of Particle Properties is obtained by letting  $s_1 \rightarrow -s_1$  and  $\delta \rightarrow \delta + \pi$  in the matrix given above. An alternative used in another review<sup>3</sup> is to change Eq. (4) by  $s_1 \rightarrow -s_1$  but leave  $\delta$  unchanged. With this change in  $s_1$ ,  $\theta_1$  becomes the usual Cabibbo angle, with the "correct" sign (i.e.,  $d' = d \cos \theta_1 + s \sin \theta_1$ ), in the limit  $\theta_2 = \theta_3 = 0$ . The angles  $\theta_1$ ,  $\theta_2$ ,  $\theta_3$  can, as before, all be taken to lie in the first quadrant by adjusting quark field phases. Since all these parametrizations are referred to as "the" Kobayashi-Maskawa

form, some care about which one is being used is needed when the quadrant in which  $\delta$  lies is under discussion.

Another parametrization, which emphasizes the relative sizes of the matrix elements by expressing them in powers of the Cabibbo angle, was introduced by Wolfenstein.<sup>4</sup> Still other parametrizations<sup>5</sup> have come into the literature in connection with attempts to define "maximal CP violation." No physics can depend on which of the above parametrizations (or any other) is used as long as it is used consistently and care is taken to be sure that no other choice of phases is in conflict.

Our present knowledge of the matrix elements comes from the following sources:

(1) Nuclear beta decay, when compared to muon decay, gives<sup>6</sup>

$$|V_{ud}| = 0.9729 \pm 0.0012. \quad (5)$$

Recent refinements (wherein leading log radiative corrections are summed using the renormalization group and structure-dependent  $O(\alpha)$  terms are analyzed and estimated) have been included, thereby lowering  $|V_{ud}|$  by 0.13%.

(2) An analysis of hyperon and  $K_{e3}$  decays yields<sup>7</sup>

$$|V_{us}| = 0.221 \pm 0.002. \quad (6)$$

The isospin violation between  $K_{e3}^+$  and  $K_{e3}^0$  decays has been taken into account, bringing the values of  $|V_{us}|$  extracted from these two decays into agreement at the 1% level of accuracy. The hyperon data alone tend to give a higher value, but theoretically they have larger possible uncertainties because of first-order symmetry-breaking effects in the axial-vector couplings. A simultaneous fit to both data sets shows that the difference is not statistically significant and yields the mean value given above.

(3) From neutrino and antineutrino production of charm, the CDHS group has deduced<sup>8</sup>

$$|V_{cd}| = 0.24 \pm 0.03. \quad (7)$$

(4) Values of  $|V_{cs}|$  from such experiments are dependent on assumptions about the strange-quark density in the parton sea. Using the conservative assumption that the strange-quark sea does not exceed the value corresponding to an  $SU(3)$  symmetric sea, a bound on  $|V_{cs}|$  results<sup>8</sup> which is comparable to that given below. A different source of information on  $|V_{cs}|$  arises from comparing the experimental value for  $\Gamma(D \rightarrow \bar{K}e^+\nu_e)$  with the expression that follows from the standard weak interaction amplitude:

$$\Gamma(D \rightarrow \bar{K}e^+\nu_e) = |f_+^D(0)|^2 |V_{cs}|^2 (1.54 \times 10^{11} \text{ sec}^{-1}). \quad (8)$$

Here  $f_+^D[(p_D - p_K)^2]$  is the form factor for  $D_{e3}$  decay which is the analogue of  $f_+[(p_K - p_\pi)^2]$  for  $K_{e3}$  decay. With the parametrization  $f_+^D(t)/f_+^D(0) = M^2/(M^2 - t)$  and  $M = 2.1$  GeV from recent measurements,<sup>9</sup> its variation has been taken into account in deriving Eq. (8). From combining data on  $\text{BR}(D^+ \rightarrow \bar{K}^0 e^+ \nu_e)$  and  $\text{BR}(D^0 \rightarrow K^- e^+ \nu_e)$  with world-average values<sup>9</sup> of  $\tau_{D^+}$  and  $\tau_{D^0}$ , the resulting value of the left-hand side of Eq. (8) is  $0.79 \pm 0.11 \times 10^{11} \text{ sec}^{-1}$ . Therefore,

$$|f_+^D(0)|^2 |V_{cs}|^2 = 0.51 \pm 0.07. \quad (9)$$

With sufficient confidence in a theoretical calculation of  $|f_+^D(0)|$ , a value of  $|V_{cs}|$  follows;<sup>10</sup> but even with the very conservative assumption that  $|f_+^D(0)| < 1$ , it follows that

$$|V_{cs}| > 0.66. \quad (10)$$

The constraint of unitarity when there are only three generations gives a much tighter bound [see matrix (2)].

## THE KOBAYASHI-MASKAWA MIXING MATRIX (Cont'd)

(5) The ratio  $|V_{ub}/V_{cb}|$  is obtained from the semileptonic decay of  $B$  mesons by fitting to the lepton energy spectrum as a sum of contributions involving  $b \rightarrow u$  and  $b \rightarrow c$ . The relative overall phase-space factor between the two processes is calculated from the usual four-fermion interaction with one massive fermion ( $c$  quark or  $u$  quark) in the final state. The value of this factor is between 0.4 and 0.5, depending on the quark masses used. We use 0.45. The lack of observation of the higher momentum leptons characteristic of  $b \rightarrow u\ell\bar{\nu}_\ell$ , as compared to  $b \rightarrow c\ell\bar{\nu}_\ell$ , results in a limit which depends on the lepton energy spectrum assumed for each decay. As more data have accumulated, the inadequacy of previously used parametrizations has become clear.<sup>9</sup> Conservatively using only the lepton momentum region beyond the endpoint for  $b \rightarrow c\ell\bar{\nu}_\ell$  results in<sup>9</sup>

$$\frac{\Gamma(b \rightarrow u\ell\bar{\nu}_\ell)}{\Gamma(b \rightarrow c\ell\bar{\nu}_\ell)} < 0.08, \quad (11)$$

which translates to

$$|V_{ub}/V_{cb}| < 0.19. \quad (12)$$

Being slightly less conservative and including the last 200 MeV/ $c$  of the  $b \rightarrow c\ell\bar{\nu}_\ell$  spectrum gives a stronger limit<sup>9</sup>

$$\frac{\Gamma(b \rightarrow u\ell\bar{\nu}_\ell)}{\Gamma(b \rightarrow c\ell\bar{\nu}_\ell)} < 0.04, \quad (13)$$

which coincides with previous limits<sup>11</sup> and translates to

$$|V_{ub}/V_{cb}| < 0.14. \quad (14)$$

There are some theoretical uncertainties in this analysis stemming from the fact that the physical decays involve actual hadrons and not just quarks as is assumed in the calculations of the lepton spectra for  $b \rightarrow u\ell\bar{\nu}_\ell$  and  $b \rightarrow c\ell\bar{\nu}_\ell$ .

(6) The magnitude of  $V_{cb}$  itself can be determined if the measured semileptonic bottom hadron partial width is assumed to be that of a  $b$  quark decaying through the usual  $V-A$  interaction:

$$\Gamma(b \rightarrow c\ell\bar{\nu}_\ell) = \frac{BR(b \rightarrow c\ell\bar{\nu}_\ell)}{\tau_b} = \frac{G_F^2 m_b^5}{192\pi^3} F(m_b, m_c) |V_{cb}|^2, \quad (15)$$

where  $\tau_b$  is the  $b$  lifetime and  $F(m_b, m_c)$  is the phase-space factor chosen above as 0.45.

Using an average semileptonic branching ratio measured in the continuum of<sup>9</sup>  $12.1 \pm 0.8\%$  [which from Eq. (11) is  $BR(b \rightarrow c\ell\bar{\nu}_\ell)$  to within 8%], a world-average bottom hadron lifetime<sup>9</sup> of  $1.26 \pm 0.16 \times 10^{-12}$  sec. and  $m_b$  between 4.8 and 5.2 GeV, we get

$$0.037 < |V_{cb}| < 0.053. \quad (16)$$

where the range of  $m_b$  values used in extracting  $V_{cb}$  has been treated as a theoretical systematic error on top of the errors arising from experimental measurements.

The results for three generations of quarks, from Eqs. (5), (6), (7), (10), (12), and (16) plus unitarity, are given in matrix (2). The ranges there are different from those given in Eqs. (5)–(16) because of the constraint of unitarity, but are consistent with the quoted one-standard-deviation errors.

The data do not preclude there being more than three generations. Moreover, the entries deduced from unitarity might be altered when the K-M matrix is expanded to accommodate more generations. Conversely, the known entries restrict the possible values of additional elements if the matrix is expanded to account for additional generations. For example, unitarity and the known elements of the first row require that<sup>6</sup> any additional element in the first row have a magnitude  $|V_{ub'}| < 0.088$ . When there are more than three generations, the allowed ranges (at 90% C.L.) of the matrix elements connecting the first three generations are

$$\begin{pmatrix} 0.9710 & \text{to } 0.9748 & 0.218 & \text{to } 0.224 & 0 & \text{to } 0.01 & \dots \\ 0.192 & \text{to } 0.288 & 0.66 & \text{to } 0.98 & 0.037 & \text{to } 0.053 & \dots \\ 0 & \text{to } 0.14 & 0 & \text{to } 0.72 & 0 & \text{to } 0.999 & \dots \\ \vdots & & \vdots & & \vdots & & \vdots \\ \vdots & & \vdots & & \vdots & & \vdots \end{pmatrix} \quad (17)$$

where we have used unitarity (for the expanded matrix) and Eqs. (5), (6), (7), (10), (12), and (16).

Further information on the angles requires theoretical assumptions. In particular, as  $CP$ -violating amplitudes involve  $\sin \delta$ , assuming that observed  $CP$  violation is solely related to a nonzero value of  $\delta$  allows additional constraints to be brought to bear. While hadronic matrix elements whose values are imprecisely known now enter, the constraints from  $CP$  violation in the neutral kaon system are tight enough that there may be no solution at all for certain quark masses, values of  $\delta$ , etc. See the reviews in Ref. 12.

\* Prepared January 1986 by F.J. Gilman and K. Kleinknecht.

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## STANDARD MODEL OF ELECTROWEAK INTERACTIONS\*

The couplings of the photon,  $W^\pm$ , and  $Z$  to fundamental fermions are

$$\bar{\psi}\gamma^\mu \left[ eQA_\mu + \frac{e(1-\gamma_5)}{2\sqrt{2}\sin\theta_W} \left( T^+W_\mu^+ + T^-W_\mu^- \right) + \frac{e}{\sin\theta_W\cos\theta_W} \left( \frac{1}{2}(1-\gamma_5)T_3 - \sin^2\theta_W Q \right) Z_\mu \right] \psi, \quad (1)$$

where

$$\psi = \begin{pmatrix} u \\ d' \end{pmatrix}, \begin{pmatrix} c \\ s' \end{pmatrix}, \begin{pmatrix} t \\ b' \end{pmatrix}, \begin{pmatrix} \nu_e \\ e^- \end{pmatrix}, \begin{pmatrix} \nu_\mu \\ \mu^- \end{pmatrix}, \begin{pmatrix} \nu_\tau \\ \tau^- \end{pmatrix};$$

for mixing effects defining  $d'$ ,  $s'$ , and  $b'$ , see the section on the Kobayashi-Maskawa Mixing Matrix, Eq. (1);

$T^\pm$  = weak isospin raising operator ( $T^\pm$  act on left-handed fermions);

$T_3$  = third component of weak isospin (i.e.,  $1/2$  for  $\nu_e, \nu_\mu, \nu_\tau, u, c, t; -1/2$  for  $e^-, \mu^-, \tau^-, d, s, b$ );

$Q$  = electric charge operator, in units of proton charge;

$\theta_W$  = weak mixing angle;

$A$  = electromagnetic vector potential.

Thus, for example, the  $Wey$  coupling is

$$\left( \frac{e}{\sqrt{2}\sin\theta_W} \right) \left[ W_\mu^- \bar{e} \gamma^\mu \frac{1}{2}(1-\gamma_5)\nu + W_\mu^+ \bar{\nu} \gamma^\mu \frac{1}{2}(1-\gamma_5)e \right]$$

and the  $Zu\bar{u}$  coupling is

$$\left( \frac{e}{\sin\theta_W\cos\theta_W} \right) Z_\mu \bar{u} \gamma^\mu \left[ \frac{1}{4}(1-\gamma_5) - \frac{2}{3}\sin^2\theta_W \right] u.$$

The physical, neutral fields  $A$  and  $Z$  are mixtures of  $W_3$ , the partner of  $W^\pm$ , and another field  $B$ :

$$A = W_3 \sin\theta_W + B \cos\theta_W, \quad Z = W_3 \cos\theta_W - B \sin\theta_W.$$

The  $SU(2) \times U(1)$  gauge couplings  $g$  and  $g'$  appear as

$$gW_\mu \cdot \mathbf{T} + g'B_\mu \frac{Y}{2},$$

where electric charge  $Q$ ,  $T_3$ , and  $Y/2$  are connected by  $Q = T_3 + Y/2$ . The couplings and mixing angle are related by  $\tan\theta_W = g'/g$ ,  $\sin\theta_W = e/g$ .

Branching fractions of the  $W^\pm$  and  $Z$  are predicted to be roughly

$$\begin{aligned} \text{BF}(W^+ \rightarrow e^+\nu_e) &= 0.08, & \text{BF}(W^+ \rightarrow u\bar{d}) &= 0.24, \\ \text{BF}(Z \rightarrow \nu_e\bar{\nu}_e) &= 0.06, & \text{BF}(Z \rightarrow e^+e^-) &= 0.03, \\ \text{BF}(Z \rightarrow u\bar{u}) &= 0.10, & \text{BF}(Z \rightarrow d\bar{d}) &= 0.13, \end{aligned}$$

etc., and similarly for the other generations, assuming there is no suppression for phase space even for the  $t$  quark. The total widths are expected to be (with  $M_Z = 94$  GeV,  $M_{\text{top quark}} = 30$  GeV):  $\Gamma(W) = 2.8$  GeV and  $\Gamma(Z) = 2.8 \pm 0.1$  GeV.

In the standard model, the best fit<sup>1</sup> to current data yields  $\sin^2\theta_W = 0.226 \pm 0.004$ , which corresponds to  $M_Z = 92.5 \pm 0.5$  GeV, with theoretical uncertainties of about  $\pm 0.005$  for  $\sin^2\theta_W$  and  $\pm 0.7$  GeV for  $M_Z$ .

The minimal  $SU(2) \times U(1)$  model of electroweak interactions, which has only one Higgs doublet, has three fundamental parameters (aside from the masses of the fermions and the Higgs boson). In the Lagrangian they are the  $SU(2)$  coupling constant  $g$ , the  $U(1)$  coupling constant  $g'$ , and the vacuum expectation value (v.e.v.) of the Higgs  $SU(2)$  doublet field.

It is best to choose the three parameters so they are identified with physical quantities. Two of the parameters can be taken to be  $\alpha = 1/137.036$ , which is obtained from measurements of the Josephson effect, and  $G_F = 1.16637 \times 10^{-5}$  GeV<sup>-2</sup>, which is derived from the muon lifetime once QED corrections are taken into account, viz.:

$$\tau_\mu^{-1} = \frac{G_F^2 m_\mu^5}{192\pi^3} \left[ 1 + \frac{\alpha}{2\pi} \left( \frac{25}{4} - \pi^2 \right) \left( 1 + \frac{2\alpha}{3\pi} \ell n \frac{m_\mu}{m_e} \right) \right] \times \left[ 1 - \frac{8m_e^2}{m_\mu^2} \right]. \quad (2)$$

The third parameter will be taken to be  $M_Z$ , since it will soon be known accurately from measurements at LEP and SLC. The  $W$  mass  $M_W$  is then predicted to be

$$M_W = \frac{M_Z}{\sqrt{2}} \left[ 1 + \left( 1 - \frac{4\pi\alpha}{\sqrt{2}M_Z^2 G_F (1 - \Delta r)} \right)^{1/2} \right]^{1/2} \quad (3)$$

The quantity  $\Delta r$  is  $0.0696 \pm 0.0020$ . It arises from radiative corrections<sup>2-4</sup> and depends only slightly on the  $t$  quark mass (taken to be 36 GeV), the Higgs mass (taken equal to  $M_Z$ ), and the  $Z$  mass itself (taken to be 92 GeV). Historically the third parameter was taken to be the weak mixing angle  $\theta_W$ . This mixing angle can be defined by

$$\cos\theta_W \equiv \frac{M_W}{M_Z}. \quad (4)$$

In lowest order, this mixing angle is the same as that introduced in Eq. (1). Once radiative corrections are taken into account, this ceases to be true, and it would appear that the consequent confusion can be reduced by not regarding  $\theta_W$  as one of the fundamental parameters. These  $O(\alpha)$  radiative corrections are of three types:

- i) QED vacuum polarization;
- ii) QED bremsstrahlung and graphs involving virtual photons in loops — these are detector dependent and must be carefully calculated for each process, taking experimental cuts and energy resolution into account;
- iii) graphs involving weak particles in virtual loops.

At present precise measurements of electroweak interactions are made in a number of experiments and are conventionally expressed in terms of  $\sin^2\theta_W$  defined by Eq. (4). Among the experiments are:

1. Neutrino and antineutrino deep inelastic scattering from isosinglet targets.
2. Neutrino and antineutrino deep inelastic scattering by protons.
3. Elastic  $\nu_\mu p$  and  $\bar{\nu}_\mu p$  scattering.
4. Exclusive and inclusive  $\pi$  production in neutral-current events.
5. Neutrino disintegration of the deuteron:  $\bar{\nu}_e d \rightarrow \bar{\nu}_e np$ .
6. Polarized-electron deuteron deep inelastic scattering.
7. Forward-backward asymmetry in  $e^+e^- \rightarrow \mu^+\mu^-$ .
8. Elastic  $\nu_\mu e$  and  $\bar{\nu}_\mu e$  scattering.
9. Parity non-conservation in heavy atoms.

The data for these processes and a comparison with the predictions of the standard model are given in three extensive reviews.<sup>1,5-6</sup> The conclusion of these and all similar studies is that all available data are consistent with the standard model.

## STANDARD MODEL OF ELECTROWEAK INTERACTIONS (Cont'd)

As stated above, a global fit to the data gives<sup>1</sup>  $\sin^2\theta_W = 0.226 \pm 0.004$ , which corresponds to the value  $92.5 \pm 0.5$  GeV for the remaining fundamental parameter  $M_Z$ .

Table 1 shows the values of  $M_Z$ , and the derived quantity  $\sin^2\theta_W$ , extracted from different experiments (quoted uncertainties include both statistical and experimental systematic errors). The remarkable consistency is a strong confirmation of the standard electroweak model.

Table 1

Process	$M_Z$	$\sin^2\theta_W$	Ref.
$e^+e^- \rightarrow \mu^+\mu^-$	$103 \pm 4.8$	$0.17 \pm 0.02$	7
$\nu p \rightarrow \nu p$ $\bar{\nu} p \rightarrow \bar{\nu} p$ }	$91.8 \pm 2.8$	$0.23 \pm 0.02$	1,8
<i>ed</i> asymmetry	$93.3 \pm 2.1$	$0.220 \pm 0.014$	1,9
$\nu_\mu e \rightarrow \nu_\mu e$	$91.8 \pm 2.8$	$0.23 \pm 0.02^\dagger$	1,10
Parity violation in atoms	$98.5 \pm 7.9$	$0.19 \pm 0.04$	11
$\bar{\nu} N \rightarrow \mu X, \bar{\nu} X$ $\nu N \rightarrow \mu X, \nu X$ }	$92.4 \pm 0.6$	$0.226 \pm 0.004^\ddagger$	12

<sup>†</sup>Two recent experiments average to  $\sin^2\theta_W = 0.212 \pm 0.023$ , whereas seven measurements by four 1979-era experiments average to  $\sin^2\theta_W = 0.277 \pm 0.034$ .

<sup>‡</sup>These results from deep inelastic scattering also have theoretical uncertainties (largely from the imprecisely known *c*-quark mass) which are about  $\pm 0.005$  for  $\sin^2\theta_W$  and  $\pm 0.7$  GeV for  $M_Z$ . Two recent experiments have given us modified values (which we have included) due to discovery of an error in their radiative correction program; these are unpublished.

In the case of neutrino deep inelastic scattering, the charged- and neutral-current cross sections for the standard model are given by:

$$R_{\nu N} = \frac{\sigma_{\nu N}^{NC}}{\sigma_{\nu N}^{CC}} = \frac{\frac{1}{2} - \sin^2\hat{\theta}_W + \frac{20}{27} \sin^4\hat{\theta}_W + \epsilon \left( \frac{1}{6} - \frac{1}{3} \sin^2\hat{\theta}_W + \frac{20}{27} \sin^4\hat{\theta}_W \right)}{1 + \frac{1}{3}\epsilon},$$

$$R_{\bar{\nu} N} = \frac{\sigma_{\bar{\nu} N}^{NC}}{\sigma_{\bar{\nu} N}^{CC}} = \frac{\frac{1}{6} - \frac{1}{3} \sin^2\hat{\theta}_W + \frac{20}{27} \sin^4\hat{\theta}_W + \epsilon \left( \frac{1}{2} - \sin^2\hat{\theta}_W + \frac{20}{27} \sin^4\hat{\theta}_W \right)}{\frac{1}{3} + \epsilon}. \quad (5)$$

In lowest order, the quantity  $\sin^2\hat{\theta}_W$  is equal to  $\sin^2\theta_W$  defined above. When radiative corrections are taken into account,  $\sin^2\hat{\theta}_W = \sin^2\theta_W + \delta$ , with  $\delta$  of order 0.01. Its precise value is dependent upon the kinematics of a particular experiment.<sup>13</sup>

The parameter  $\epsilon \approx 0.2$  is the ratio of antiquark momentum to quark momentum in the nucleon. It is related to the ratio of the neutrino charged-current cross section to the antineutrino charged-current cross section:

$$\frac{\sigma_{\nu N}^{CC}}{\sigma_{\bar{\nu} N}^{CC}} = \frac{1 + \frac{1}{3}\epsilon}{\frac{1}{3} + \epsilon}.$$

The polarized-electron deuteron scattering experiment measured the asymmetry in the cross section for left-handed ( $\sigma_L$ ) and right-handed ( $\sigma_R$ ) electrons:

$$A = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}.$$

In the quark parton model, this has the form (with  $q^2 > 0$  for deep inelastic scattering)

$$\frac{A}{q^2} = a_1 + a_2 \left[ \frac{1 - (1-y)^2}{1 + (1-y)^2} \right], \quad (6)$$

where  $y$  is the fraction of the incident lepton's energy lost in the collision. For an isoscalar target like the deuteron, and ignoring the antiquarks, the standard model gives

$$a_1 = -\frac{G_F}{2\sqrt{2}\pi\alpha} \frac{9}{10} \left( 1 - \frac{20}{9} \sin^2\hat{\theta}_W \right)$$

$$a_2 = -\frac{G_F}{2\sqrt{2}\pi\alpha} \frac{9}{10} \left( 1 - 4 \sin^2\hat{\theta}_W \right). \quad (7)$$

In the case of the SLAC *ed* experiment,<sup>3</sup>  $\sin^2\hat{\theta}_W = \sin^2\theta_W + 0.006$ .

Neutrino electron elastic-scattering cross sections were calculated with the following formulae:

$$R_{\nu\bar{\nu}} = \frac{\sigma(\nu_\mu e \rightarrow \nu_\mu e)}{\sigma(\bar{\nu}_\mu e \rightarrow \bar{\nu}_\mu e)}$$

$$\approx \frac{3 - 12 \sin^2\hat{\theta}_W + 16 \sin^4\hat{\theta}_W}{1 - 4 \sin^2\hat{\theta}_W + 16 \sin^4\hat{\theta}_W} \quad (8)$$

$$R_{NC,CC} = \frac{\sigma(\nu_\mu e \rightarrow \nu_\mu e)}{\sigma(\nu_\mu e \rightarrow \nu_\mu e)}$$

$$\approx \frac{3 - 12 \sin^2\hat{\theta}_W + 16 \sin^4\hat{\theta}_W}{12} \left[ 1 - \frac{m_\mu^2}{2m_e E_\nu} \right]^{-2} \quad (9)$$

The standard model agrees very well with the data; however, there is no direct evidence concerning the Higgs structure. If the model contains Higgs representations other than doublets, the theory has an additional parameter. It is useful to take this parameter to be the  $W$  mass. It is traditional to parametrize such models by a parameter  $\rho$ , defined by

$$\rho = \frac{\sqrt{2}M_W}{M_Z \left[ 1 + \left( 1 - \frac{4\pi\alpha}{\sqrt{2}M_Z^2 G_F (1 - \Delta r)} \right)^{1/2} \right]^{1/2}}.$$

so that  $\rho \equiv 1$  in the standard model. The definition of the weak mixing angle becomes

$$\cos\theta_W = \frac{M_W}{\rho M_Z}.$$

## STANDARD MODEL OF ELECTROWEAK INTERACTIONS (Cont'd)

Data can then be fitted to the two parameters  $M_Z$  and  $M_W$  (or  $\sin^2\theta_W$  and  $\rho$ ). For example, Eq. (5) has an additional factor of  $\rho^2$  multiplying the right-hand side. A two-parameter fit (excluding some recent data) yields<sup>1</sup>

$$\sin^2\theta_W = 0.223 \pm 0.006$$

$$\rho = 1.006 \pm 0.008.$$

Note that  $\rho$  is consistent with the standard model value of 1.

It is useful to compare the values of  $M_Z$  (inferred from the data discussed above) and  $M_W$  [calculated from Eq. (3)] with the observed  $W$  and  $Z$  masses. Table 2 shows the comparison. Detailed tests of the model await higher precision data on the weak boson masses, lepton asymmetries at SLC and LEP, and elastic  $\nu e$  and  $\bar{\nu} e$ . Such data will permit a single test involving the four measured parameters  $\alpha$ ,  $G_F$ ,  $M_W$ , and  $M_Z$ , and two additional tests comparing the derived quantity  $\sin^2\theta_W$  with the values obtained in neutral-current experiments.

Table 2

	UA1 <sup>14</sup>	UA2 <sup>15</sup>	Values <sup>2</sup> obtained using Table 1
$M_W$ (GeV)	$83.5_{-1.0}^{+1.1} \pm 2.7$	$81.2 \pm 1.1 \pm 1.3$	$81.4 \pm 0.6$ (79.8 without radiative corr.)
$M_Z$ (GeV)	$93.0 \pm 1.4 \pm 3$	$92.5 \pm 1.3 \pm 1.5$	$92.5 \pm 0.5$ (90.2 without radiative corr.)

Should such a high-precision test reveal a discrepancy with the radiatively corrected theory, it could indicate that the model needs fundamental modification. Small deviations from the predictions described above could arise if the  $t$  quark mass or Higgs boson mass is far from the value assumed, if there are additional generations of fermions or scalars, or if additional currents beyond iso-

spin and electromagnetic exist.<sup>4</sup> At present there is no indication of the need for a modification of the theory.

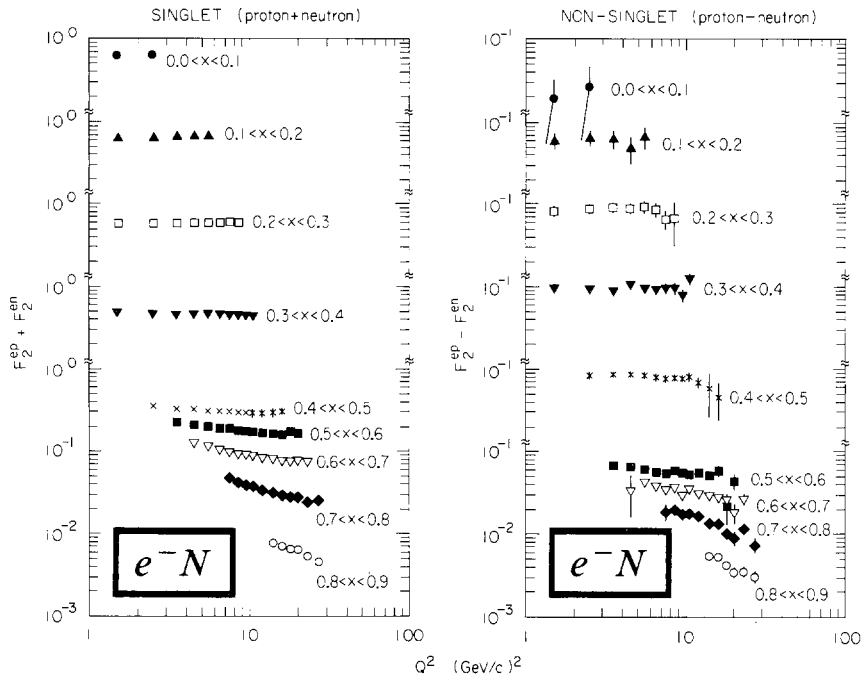
\* This section prepared with contributions from B. Lynn, L.S. Durkin, and P. Langacker.

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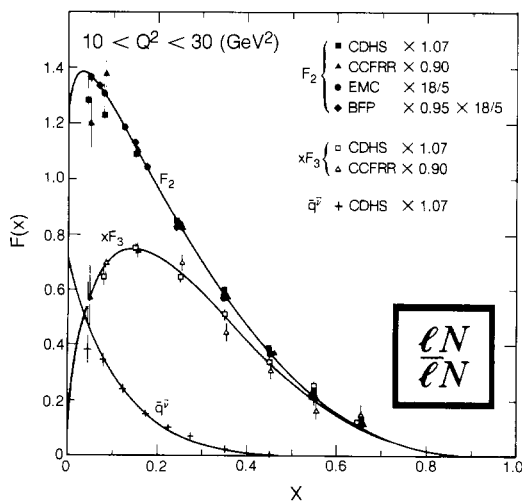
## PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES

NOTE: THE FIGURES IN THIS SECTION ARE INTENDED TO SHOW THE "BEST" OR "MOST REPRESENTATIVE" DATA IN THE OPINION OF THE COMPILER. THEY ARE NOT NECESSARILY COMPLETE COMPILATIONS OF ALL THE WORLD'S RELIABLE DATA.

### Structure Functions



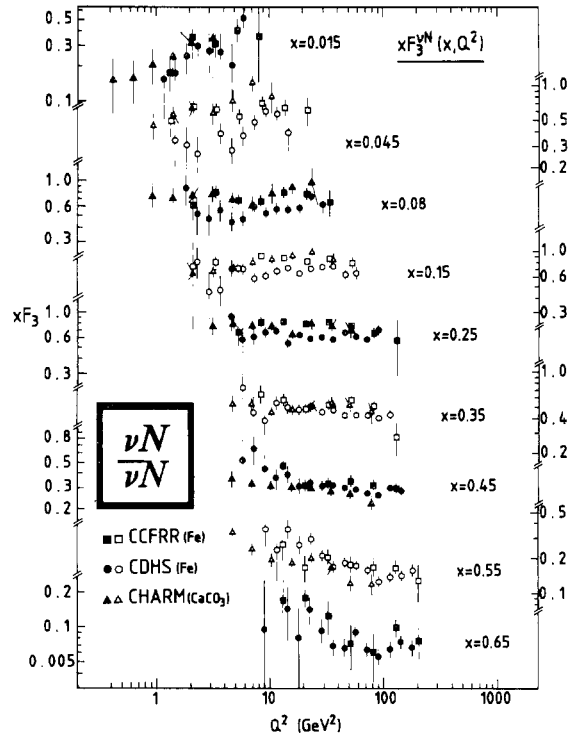
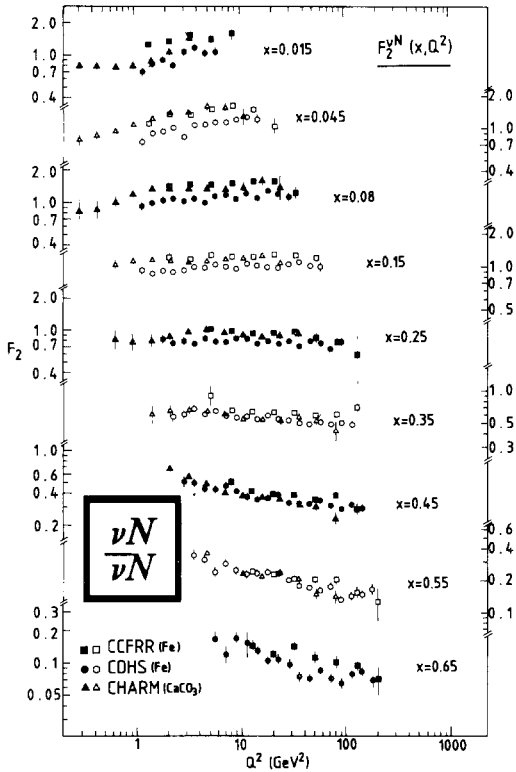
$F_2$  structure functions derived from inelastic electron-nucleon data taken at SLAC with recoil mass  $> 2$  GeV and four-momentum transfer squared  $Q^2 > 1$  (GeV/c) $^2$  are shown. For definitions of  $F_2$ ,  $x$ , and  $Q^2$ , see the Cross Sections, Decays, Structure Functions, and Kinematics Section.  $R \equiv \sigma_L/\sigma_T = 0.21^3$  was assumed. Systematic errors are comparable in size to the data point symbols. Corrections for nucleon motion in deuterium have been made. These corrections are small except for  $x > 0.7$ . No error was included to account for uncertainties in this correction. References: 1) A. Bodek et al., Phys. Rev. **D20**, 1471 (1979); 2) W.B. Atwood et al., SLAC Report No. 185 (1975); 3) M.D. Mestayer, SLAC Report No. 214 (1978); 4) S. Stein et al., Phys. Rev. **D12**, 1884 (1975). Courtesy W.B. Atwood, SLAC.



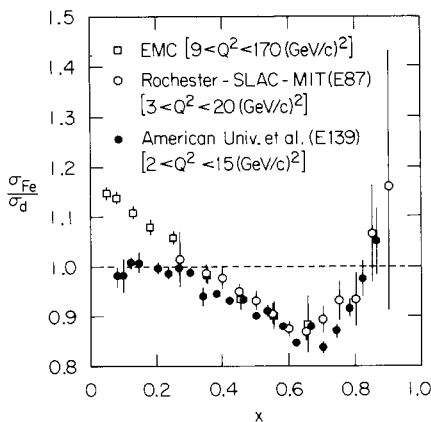
Structure functions  $F_2$ ,  $xF_3$ , and  $\bar{q}^{\bar{\nu}}$ , measured in different experiments, for fixed  $Q^2$  versus  $x$ , plotted assuming  $R = \sigma_L/\sigma_T = 0$ . The electromagnetic structure function  $F_2^{\mu N}$  measured by EMC and BFP is compared with the charged-current structure function  $F_2^{\nu N}$  using the 18/5 factor from the average charge squared of the quarks. No correction has been applied for the difference between the strange and charm sea quarks so the interpretation is  $F_2^{\nu N} = x[q + \bar{q} - \frac{1}{3}(s + \bar{s} - c - \bar{c})]$ . (In this  $Q^2$  range,  $F_2^{\nu N}$  is depleted by a similar amount due to charm threshold effects in the transition  $s \rightarrow c$ .) The antiquark distribution measured from antineutrino scattering is  $\bar{q}^{\bar{\nu}} = x(\bar{u} + \bar{d} + 2\bar{s})$ . The solid lines, valid for  $10 < Q^2 < 30$  GeV $^2$ , have the forms:  $F_2 = 3.9x^{0.55}(1-x)^{3.2} + 1.1(1-x)^8$ ,  $xF_3 = 3.6x^{0.55}(1-x)^{3.2}$ ,  $\bar{q}^{\bar{\nu}} = 0.7(1-x)^8$ . Relative normalization factors have been fitted to optimize agreement between the different data sets, and absolute changes have been arbitrarily chosen as indicated. References: CDHS — H. Abramowicz et al., Zeit. Phys. **C17**, 283 (1983); CCFRR — F. Sciulli, private communication; EMC — J.J. Aubert et al., Phys. Lett. **105B**, 322 (1981); and A. Edwards, private communication; BFP — A.R. Clark et al., Phys. Rev. Lett. **51**, 1826 (1983); and P. Meyers, Ph.D. Thesis, LBL-17108 (1983), Univ. of Calif., Berkeley. Courtesy J. Carr, Colorado.

PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)

Structure Functions



Structure functions  $F_2$  and  $xF_3$  for nucleons, measured in charged-current neutrino and antineutrino scattering from iron and marble targets with  $20 < E_\nu < 300$  GeV, versus  $Q^2$  for fixed bins of  $x$ , plotted assuming  $R = \sigma_L/\sigma_T = 0.1$  or  $R = 0.0$  for CHARM. The point-to-point systematic errors are generally smaller or of the same order as the statistical errors. In addition, CDHS quote an overall scale error of  $\pm 6\%$  for  $F_2$  and  $\pm 8\%$  for  $xF_3$ ; for the CCFRR data, the scale error is estimated to be  $\pm 4\%$ ; and for CHARM, these uncertainties are  $\pm 7\%$  for  $F_2$  and  $\pm 8\%$  for  $xF_3$ . References: CDHS — H. Abramowicz et al., *Zeit. Phys. C17*, 283 (1983); CCFRR — D.B. MacFarlane et al., *Zeit. Phys. C26*, 1 (1984); CHARM — F. Bergsma et al., *Phys. Lett. 123B*, 269 (1983); and F. Bergsma et al., *Phys. Lett. 141B*, 129 (1984). Courtesy K. Winter, CERN and J. Carr, Colorado. From F. Dydak, "Experimental Results from Lepton-Hadron and Photon-Hadron Scattering: Structure Functions and Final States," *Proceedings of the 1983 International Lepton/Photon Symposium* (Cornell, 1983), p. 634.

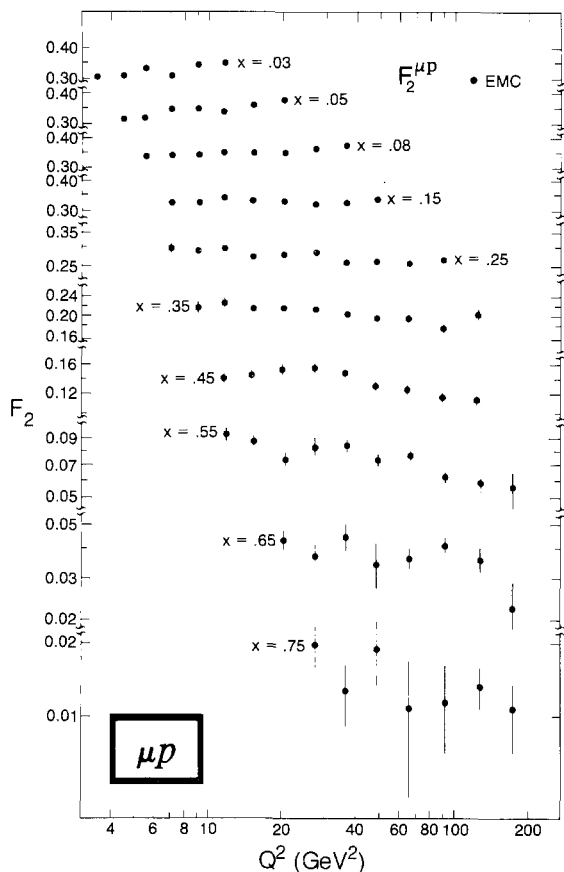


$\mu Fe / \mu d$

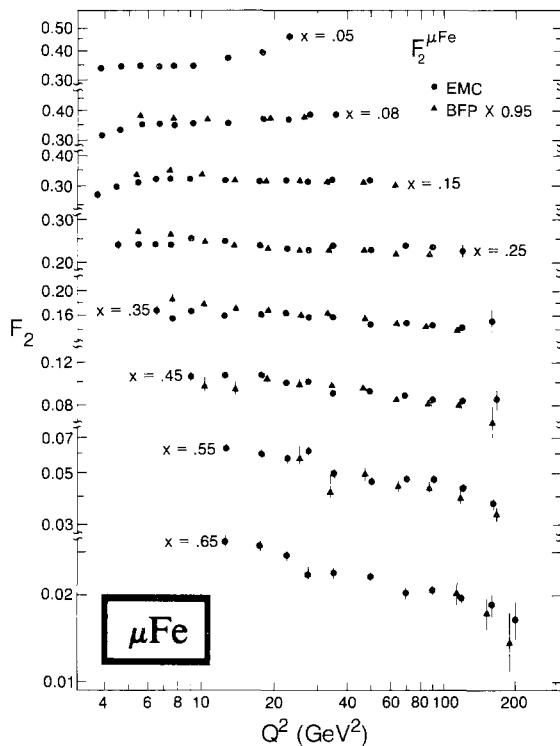
The "EMC" effect: the ratio of the differential cross section per nucleon,  $\sigma_{Fe}/\sigma_d$ , measured in electromagnetic deep inelastic scattering on iron and deuterium targets. For equal values of  $R = \sigma_L/\sigma_T$  on each target,  $\sigma_{Fe}/\sigma_d = F_2^{Fe}/F_2^d$ . The errors plotted are statistical only. References: American University et al. (electrons) — *Phys. Rev. Lett. 52*, 727 (1984); Rochester-SLAC-MIT (electrons) — *Phys. Rev. Lett. 50*, 1431 (1983); and EMC (muons) — *Phys. Lett. 123B*, 275 (1983). Forthcoming results from the BCDMS collaboration (CERN-EP/85-112) agree with EMC in the  $x$  range  $0.2 < x < 0.7$ . Courtesy D. Coward, SLAC.

## PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)

### Structure Functions

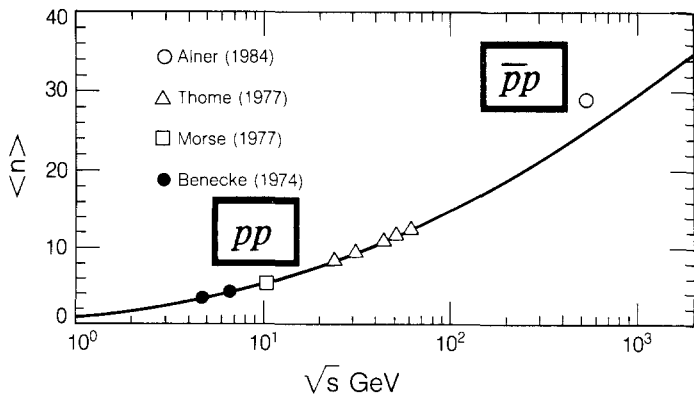


Structure function  $F_2^{\mu p}$ , measured in electromagnetic muon scattering from a hydrogen target with beam energies 120, 200, 240, 280 GeV, versus  $Q^2$  for fixed bins of  $x$ , plotted assuming  $R = \sigma_L/\sigma_T = 0$ . Reference: J.J. Aubert et al., Nucl. Phys. **B259**, 189 (1985). Courtesy J. Carr, Colorado.



Structure function  $F_2$  per nucleon, measured in electromagnetic muon scattering from iron targets with beam energies 120, 200, 250, 280 GeV (EMC) and 93, 215 GeV (BFP), versus  $Q^2$  for fixed bins of  $x$ , plotted assuming  $R = \sigma_L/\sigma_T = 0$ . A relative normalization factor has been fitted to optimize agreement between the different data sets and has been arbitrarily applied to one data set as indicated. References: EMC — J.J. Aubert et al., Phys. Lett. **105B**, 322 (1981); and A. Edwards, private communication; BFP — A.R. Clark et al., Phys. Rev. Lett. **51**, 1826 (1983); and P. Meyers, Ph.D. Thesis, LBL-17108 (1983), Univ. of Calif., Berkeley. Courtesy J. Carr, Colorado.

### Average $pp$ and $\bar{p}p$ Multiplicity

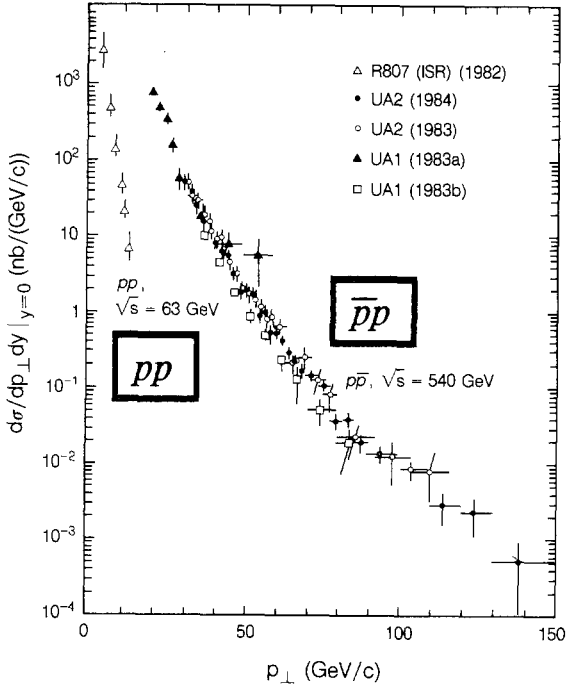


Average multiplicity as a function of  $\sqrt{s}$  for  $\bar{p}p$  at the  $S\bar{p}pS$  (open circles) and for  $pp$  at the ISR. Solid curve is a fit by Thomé et al. to their data (triangles) with the form  $\langle n \rangle = 0.88 + 0.44 \ell n s + 0.118 (\ell n s)^2$ . References:  $\bar{p}p$  — G.J. Alner et al., Phys. Lett. **138B**, 304 (1984);  $pp$  — W. Thomé et al., Nucl. Phys. **B129**, 365 (1977); W.M. Morse et al., Phys. Rev. **D15**, 66 (1977); and J. Benecke et al., Nucl. Phys. **B76**, 29 (1974).



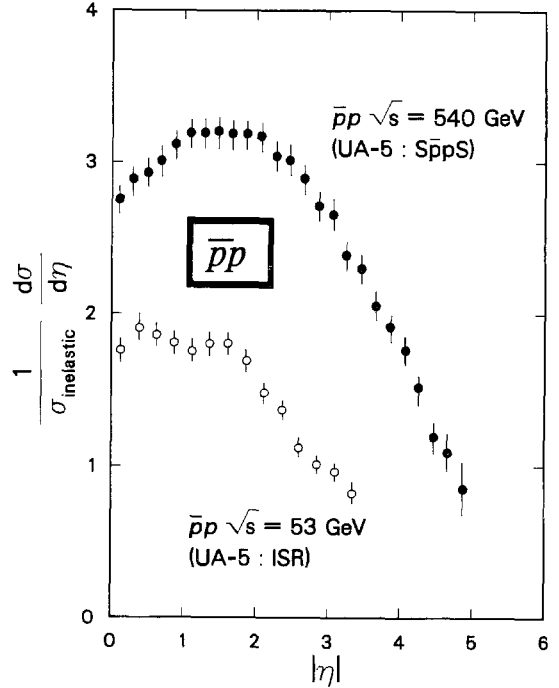
PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)

Jet Production in  $pp$  and  $\bar{p}p$  Interactions



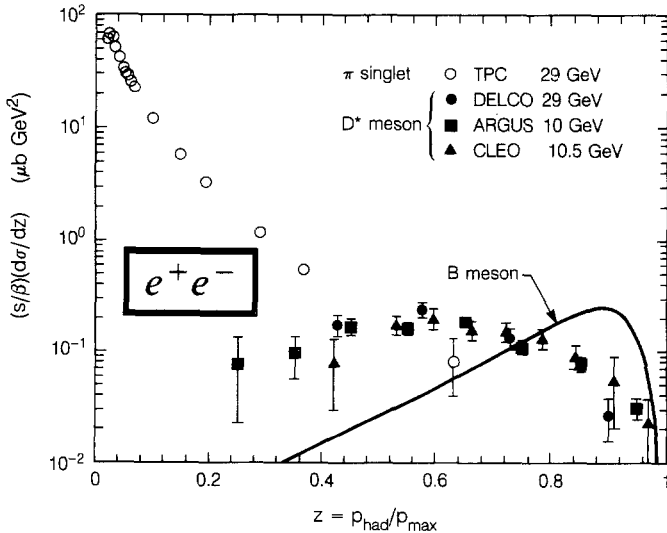
Differential cross sections for observation of a single jet of rapidity  $y = 0$  as a function of the jet transverse momentum. ISR ( $pp$ ) and  $S\bar{p}p$  collider ( $\bar{p}p$ ) data compared. Error bars include a contribution due to estimated systematic error in defining jet direction and  $p_T$ . References: ISR - T. Akesson et al., Phys. Lett. **118B**, 185 (1982); UA2 - P. Bagnaia et al., Phys. Lett. **138B**, 430 (1984); and P. Bagnaia et al., Z. Phys. **C20**, 117 (1983); UA1 - G. Arnison et al., Phys. Lett. **123B**, 115 (1983a); and G. Arnison et al., Phys. Lett. **132B**, 144 (1983b).

Pseudorapidity in  $\bar{p}p$  Interactions



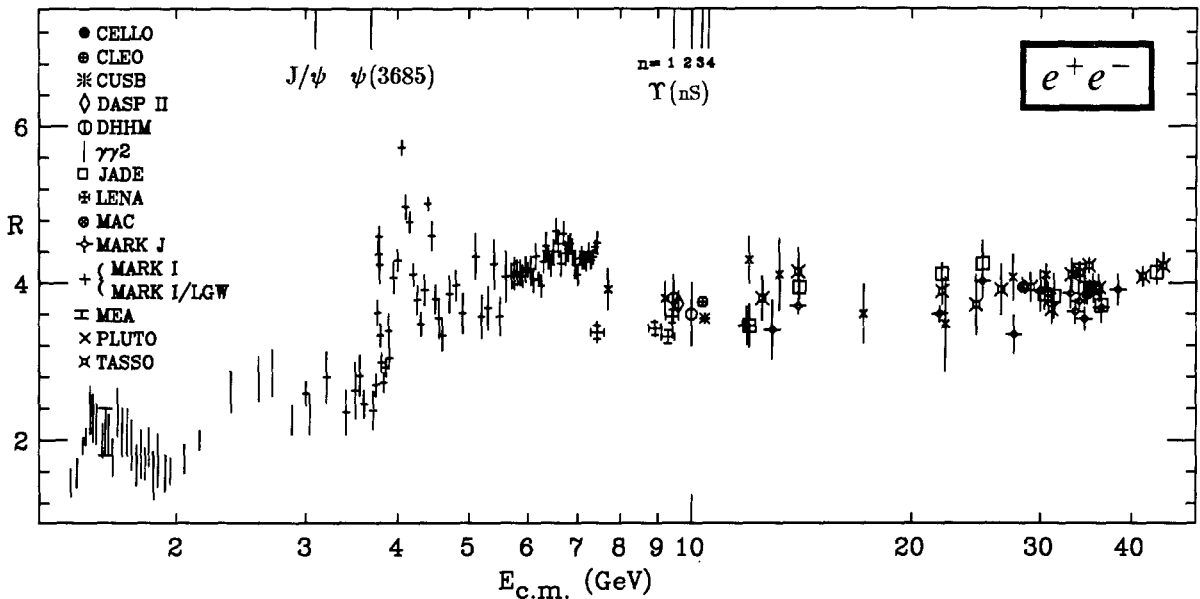
Comparison of the distribution of the pseudorapidity  $\eta = -\ln(\tan \theta_{cm}/2)$  for charged-particle production in proton-antiproton collisions at  $\sqrt{s} = 53$  GeV (1) and 540 GeV (2). References: (1) K. Alpgard et al., Phys. Lett. **112B**, 209 (1982); (2) UA5 Collaboration, presented by J. Rushbrooke in the *Proceedings of the XIV International Symposium on Multiparticle Dynamics*, eds. J.F. Gunion and P.M. Yager (World Scientific Publishing Co., Singapore, 1984).

Fragmentation Functions



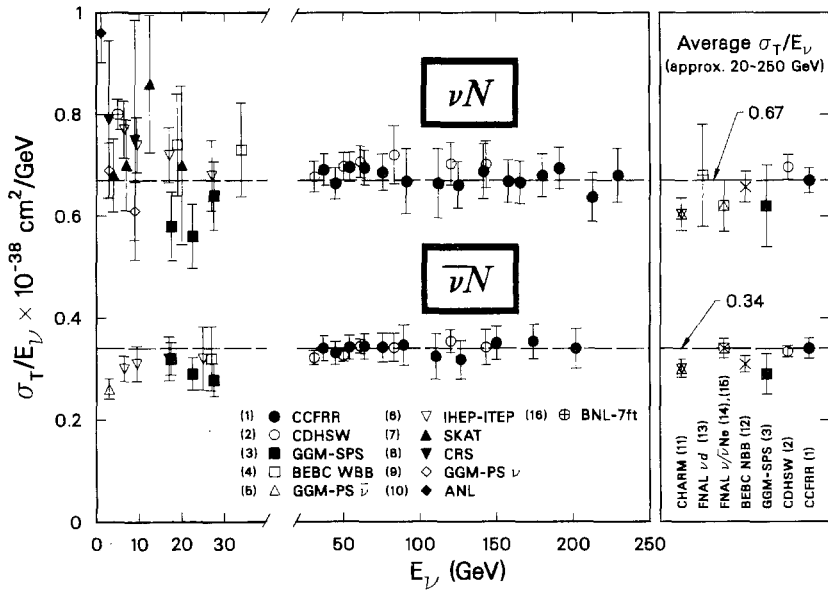
The cross section  $(s/\beta) d\sigma/dz$  versus  $z$  for producing a hadron  $h$  in  $e^+e^-$  annihilation, measured in different experiments, for fixed energies  $Q^2 = s$ . This quantity is closely related to the fragmentation function  $D_i^h(z, Q^2)$  as discussed in the QCD section. Note, however, that here we use the definition  $z \equiv p_{had}/(E_{beam}^2 - m_{had}^2)^{1/2}$ , whereas  $z \equiv E_{had}/E_{beam}$  is used by some experimenters, and theorists use  $z \equiv (E + p_{had})_{had}/(E + p_{had})_{quark}$ . The data are shown for pions (singlet term) and for  $D^*$  mesons (where  $b$ -quark contribution has been subtracted out). The data for heavy quarks are frequently parametrized by the Peterson et al. form,  $D(z) = Nz(1-z)^2/[(1-z)^2 + \epsilon_{had}z^2]^2$ . The  $\epsilon_c$  parameter ranges from 0.10 to 0.37 for the three  $D^*$  data sets shown. The  $B$ -meson curve corresponds to  $\epsilon_b = 0.014$ , which is an average derived from many experiments using an indirect method from  $B$  decays to  $e$  or  $\mu$  ( $N$  was chosen arbitrarily). References: C. Peterson et al., Phys. Rev. **D27**, 105 (1983); TPC - H. Aihara et al., Zeit. Phys. **C27**, 495 (1985); DELCO - H. Yamamoto, Ph.D. thesis (Cal Tech), report no. CALT-68-1318 (1985); ARGUS - H. Albrecht et al., Phys. Lett. **150B**, 235 (1985); and CLEO - C. Bebek et al., Phys. Rev. Lett. **49**, 610 (1982).

## PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)

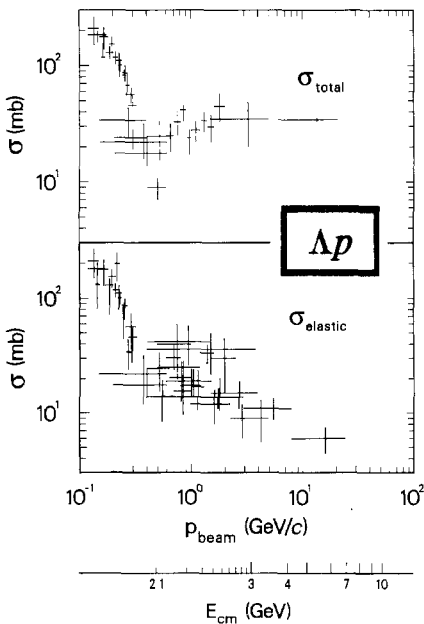


Selected measurements of  $R \equiv \sigma(e^+e^- \rightarrow \text{hadrons}) / \sigma(e^+e^- \rightarrow \mu^+\mu^-)$ , where the annihilation proceeds via one photon. The denominator is a calculated quantity; see the section on Cross Sections, Decays, Structure Functions, and Kinematics. Radiative corrections and, where important, corrections for two-photon processes and  $\tau$  production have been made. Note that the ADONE data ( $\gamma\gamma 2$  and MEA) is for  $\geq 3$  hadrons. The points in the  $\psi(3770)$  region are from the MARK I - Lead Glass Wall experiment. To preserve clarity only a representative subset of the available measurements is shown — references to additional data are included below. Also for clarity, some points have been combined or shifted slightly ( $< 4\%$ ) in  $E_{\text{cm}}$ , and some points with low statistical significance have been omitted. Systematic normalization errors are not included; they range from  $\sim 5 - 20\%$ , depending on experiment. Note the suppressed zero. The horizontal extent of the plot symbols has no significance. The positions of the  $J/\psi(3097)$ ,  $\psi(3685)$ , and the four lowest  $\tau$  vector-meson resonances are indicated at the top of the figure. References: **CELLO** — H.-J. Behrend et al., DESY 81-029; **CLEO** — R. Giles et al., Phys. Rev. **D29**, 1285 (1984); and D. Besson et al., Phys. Rev. Lett. **54**, 381 (1985); **CUSB** — E. Rice et al., Phys. Rev. Lett. **48**, 906 (1982); **DASP** — R. Brandelik et al., Phys. Lett. **76B**, 361 (1978); **DASP II** — Phys. Lett. **116B**, 383 (1982); **DHHM** — P. Bock et al. (DESY-Hamburg-Heidelberg-MPI München Collab.), Zeit. für Physik **C6**, 125 (1980);  $\gamma\gamma 2$  — C. Bacci et al., Phys. Lett. **86B**, 234 (1979); **JADE** — W. Bartel et al., Phys. Lett. **129B**, 145 (1983); and W. Bartel et al., Phys. Lett. **160B**, 337 (1985); **HRS** — D. Bender et al., Phys. Rev. **D31**, 1 (1985); **MAC** — E. Fernandez et al., Phys. Rev. **D31**, 1537 (1985); **MARK J** — B. Adeva et al., Phys. Rev. Lett. **50**, 799 (1983); and H. Newman, private communication; **MARK I** — J.L. Siegrist et al., Phys. Rev. **D26**, 969 (1982); **MARK I + Lead Glass Wall** — P.A. Rapidis et al., Phys. Rev. Lett. **39**, 526 (1977); and P.A. Rapidis, thesis, SLAC-Report-220 (1979); **MEA** — B. Esposito et al., Lett. Nuovo Cimento **19**, 21 (1977); **PLUTO** — A. Bäcker, thesis, Gesamthochschule Siegen, DESY F33-77/03 (1977); C. Gerke, thesis, Hamburg Univ. (1979); Ch. Berger et al., Phys. Lett. **81B**, 410 (1979); and W. Lackas, thesis, RWTH Aachen, DESY PLUTO-81/11 (1981); **TASSO** — R. Brandelik et al., Phys. Lett. **113B**, 499 (1982); and M. Althoff et al., Phys. Lett. **138B**, 441 (1984).

PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)

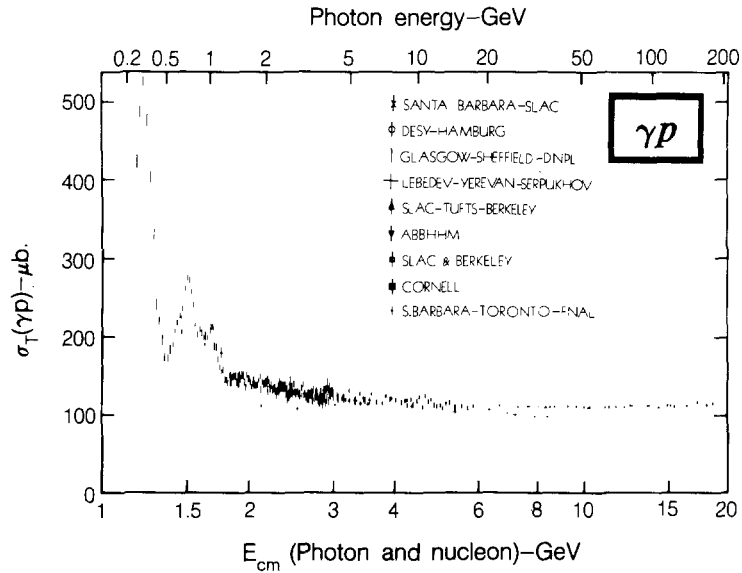


$\sigma_T/E_\nu$  for the muon neutrino and antineutrino charged-current total cross section as a function of neutrino energy. The error bars include both statistical and systematic errors. The straight lines are averages for the CCFRR measurement. Note the change in the energy scale between 30 and 50 GeV. The data points on the right give averages for other high energy measurements. References: (1) R. Blair et al., Phys. Rev. Lett. **51**, 343 (1983), and J.R. Lee, Ph.D. Thesis, Caltech (1981), "Measurements of  $\nu N$  Charged Current Cross Sections from  $E_\nu = 25$  GeV to  $E_\nu = 260$  GeV;" (2) H. Abramowicz et al., Zeit. fur Physik **C17**, 283 (1983); (3) J. Morfin et al., Phys. Lett. **104B**, 235 (1981); (4) D.C. Colley et al., Zeit. fur Physik **C2**, 187 (1979); (5) O. Erriquez et al., Phys. Lett. **80B**, 309 (1979); (6) A.S. Vovenko et al., Sov. J. Nucl. Phys. **30**, 527 (1979); (7) D.S. Baranov et al., Phys. Lett. **81B**, 255 (1979); (8) C. Baltay et al., Phys. Rev. Lett. **44**, 916 (1980); (9) S. Ciampolillo et al., Phys. Lett. **84B**, 281 (1979); (10) S.J. Barish et al., Phys. Rev. **D19**, 2521 (1979); (11) M. Jonker et al., Phys. Lett. **99B**, 265 (1981),  $E_\nu = 20-200$  GeV; (12) P. Bosetti et al., Phys. Lett. **110B**, 167 (1982),  $E_\nu = 20-200$  GeV; (13) T. Kitagaki et al., Phys. Rev. Lett. **49**, 98 (1982),  $E_\nu = 10-200$  GeV; (14) N.J. Baker et al., Phys. Rev. Lett. **51**, 735 (1983),  $E_\nu = 10-240$  GeV; (15) G.N. Taylor et al., Phys. Rev. Lett. **51**, 739 (1983),  $E_\nu = 5-250$  GeV; (16) N.J. Baker et al., Phys. Rev. **D25**, 617 (1982),  $E_\nu = 1.6-10$  GeV. Courtesy M.H. Shaevitz, Columbia University (Nevis Laboratory).

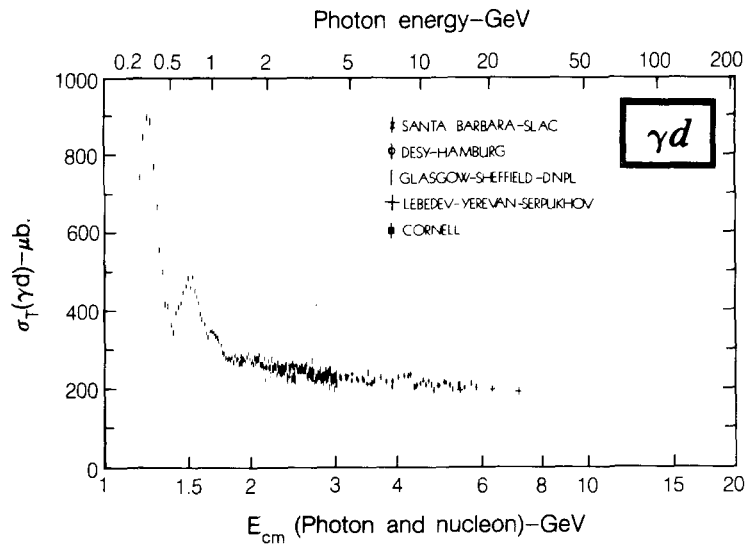


$\Delta p$  total and elastic cross sections vs. laboratory beam momentum  $p_{\text{beam}}$  and total center-of-mass energy  $E_{\text{cm}}$ . Data courtesy A. Baldini, V. Flaminio, W.G. Moorhead, D.R.O. Morrison, and N. Rivoire, CERN; and COMPAS Group, IHEP, Serpukhov, USSR. See CERN HERA Group compilations for references.

LOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)

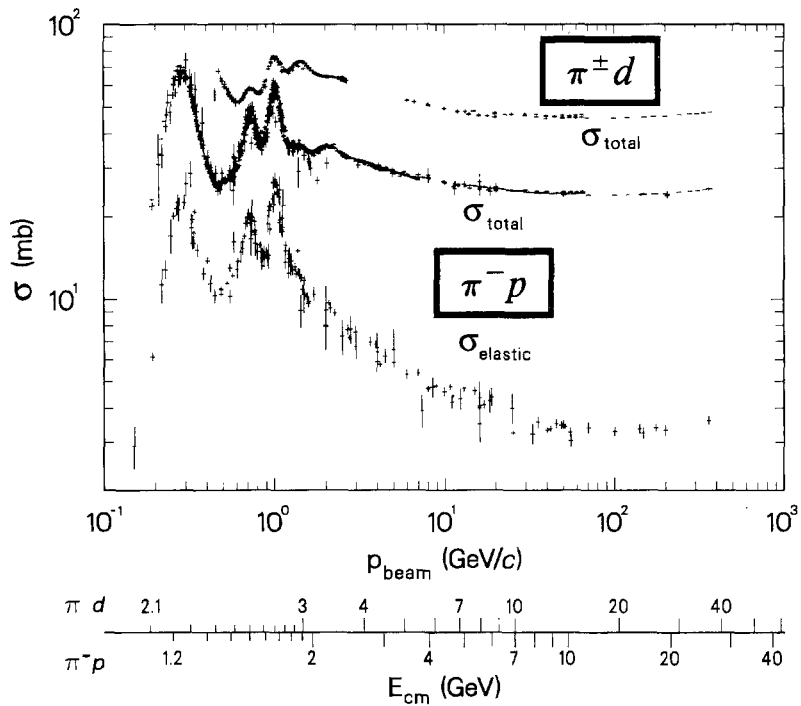
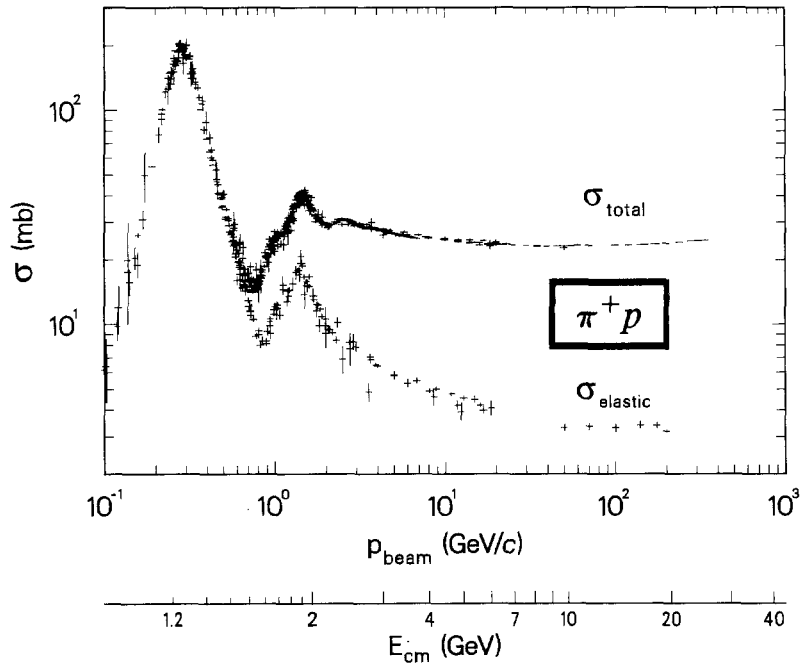


$\gamma p$  total cross section versus photon energy (top scale) and photon-plus-nucleon total center-of-mass energy (lower scale). References: **SANTA BARBARA-SLAC** — D.O. Caldwell et al., Phys. Rev. **D7**, 1362 (1973); **DESY-HAMBURG** — H. Meyer et al., Phys. Lett. **33B**, 189 (1970); **GLASGOW-SHEFFIELD-DNPL** — T.A. Armstrong et al., Phys. Rev. **D5**, 1640 (1972); **LEBEDEV-YEREVAN-SERPUKHOV** — A.S. Belousov et al., Preprint 19, Moscow, (1973); A. S. Belousov et al., Sov. Phys. Doklady **19**, 123 (1974); and A. S. Belousov et al., Sov. J. Nucl. Phys. **21**(3), 289 (1975); **SLAC-BERKELEY-TUFTS** — J. Ballam et al., Phys. Rev. **D5**, 545 (1972); **ABBHHM** — H.G. Hilpert et al., Phys. Lett. **27B**, 474 (1968); **SLAC and BERKELEY** — J. Ballam et al., Phys. Rev. Lett. **21**, 1544 (1968), and H.H. Bingham et al., Phys. Rev. **D8**, 1277 (1973); **CORNELL** — S. Michalowski et al., Phys. Rev. Lett. **39**, 737 (1977); **SANTA BARBARA-TORONTO-FNAL** — D.O. Caldwell et al., Phys. Rev. Lett. **40**, 1222 (1978). See, also, the ep data of E.D. Bloom et al., SLAC-PUB-653 (1969). Courtesy Gething M. Lewis, Glasgow.



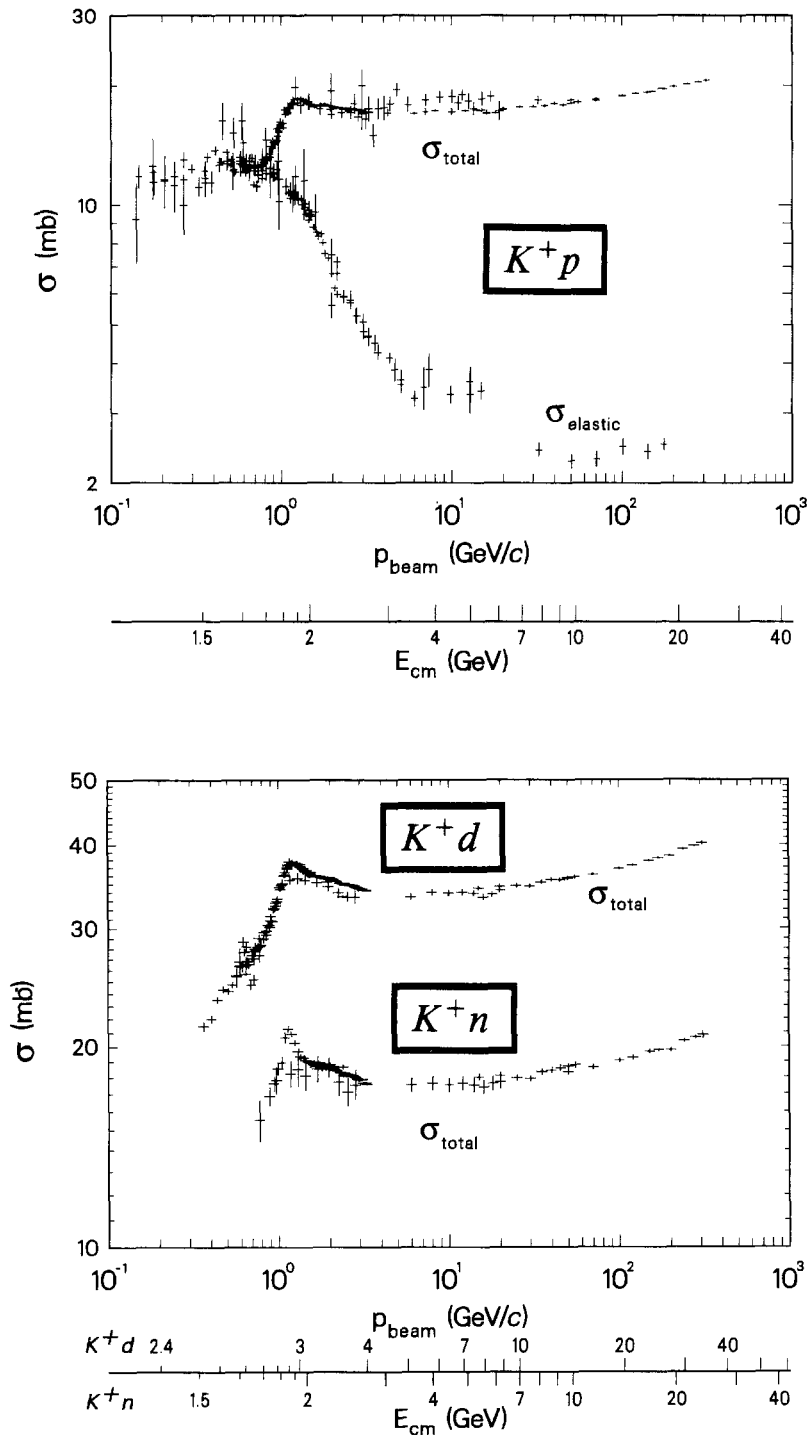
$\gamma d$  total cross section versus photon energy (top scale) and photon-plus-single-nucleon total center-of-mass energy (lower scale). References: **SANTA BARBARA-SLAC** — D.O. Caldwell et al., Phys. Rev. **D7**, 1362 (1973); **DESY-HAMBURG** — H. Meyer et al., Phys. Lett. **33B**, 189 (1970); **GLASGOW-SHEFFIELD-DNPL** — T.A. Armstrong et al., Nucl. Phys. **B41**, 445 (1972); **LEBEDEV-YEREVAN-SERPUKHOV** — A.S. Belousov et al., Sov. J. Nucl. Phys. **21**(3), 289 (1975); **CORNELL** — S. Michalowski et al., Phys. Rev. Lett. **39**, 737 (1977). Courtesy Gething M. Lewis, Glasgow.

## PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)



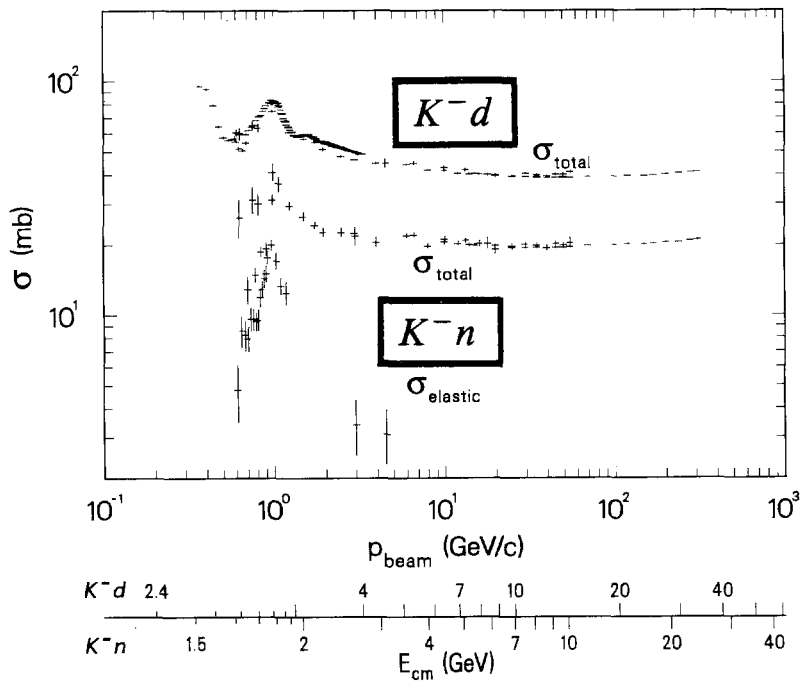
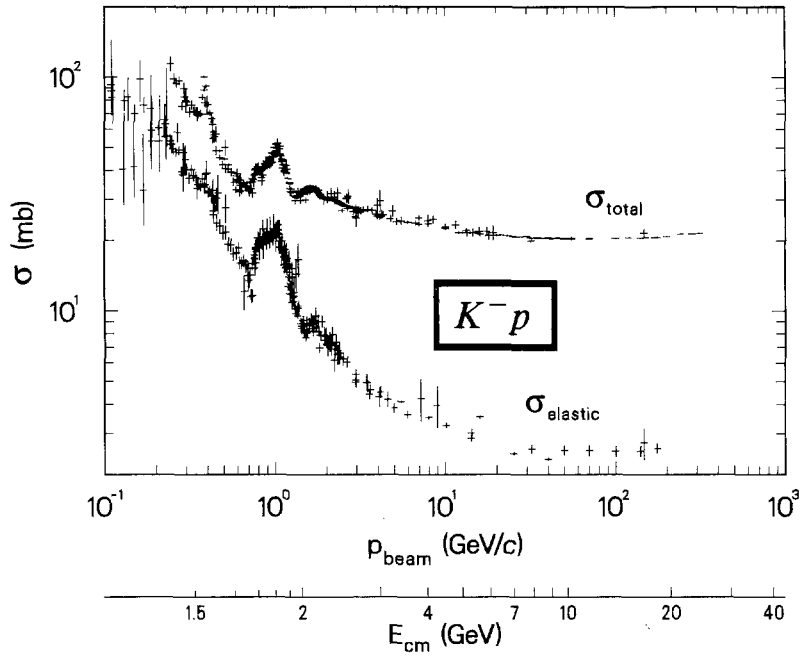
Hadronic total and elastic cross sections vs. laboratory beam momentum  $p_{\text{beam}}$  and total center-of-mass energy  $E_{\text{cm}}$ . Data courtesy A. Baldini, V. Flaminio, W.G. Moorhead, D.R.O. Morrison, and N. Rivoire, CERN; and COMPAS Group, IHEP, Serpukhov, USSR. See CERN HERA Group compilations for references.

## PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)



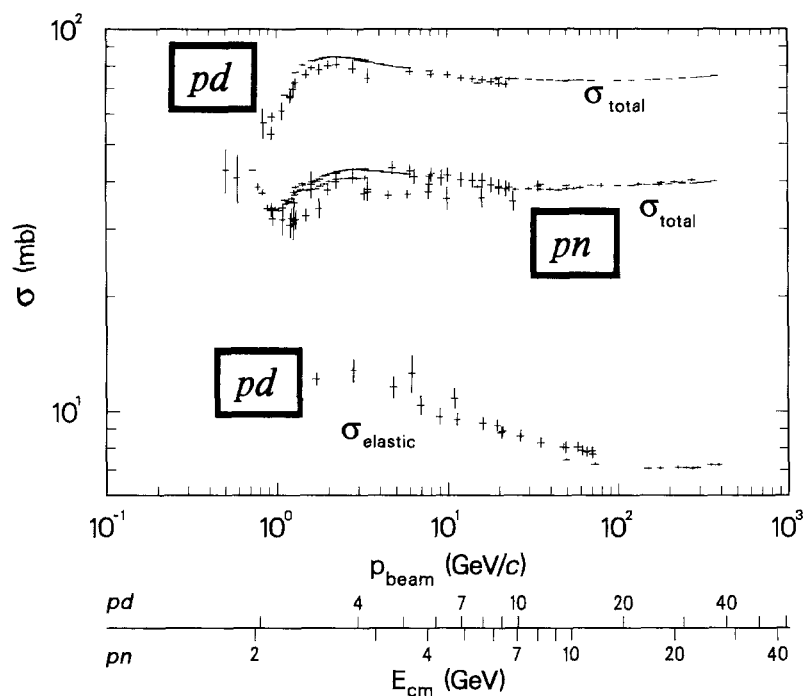
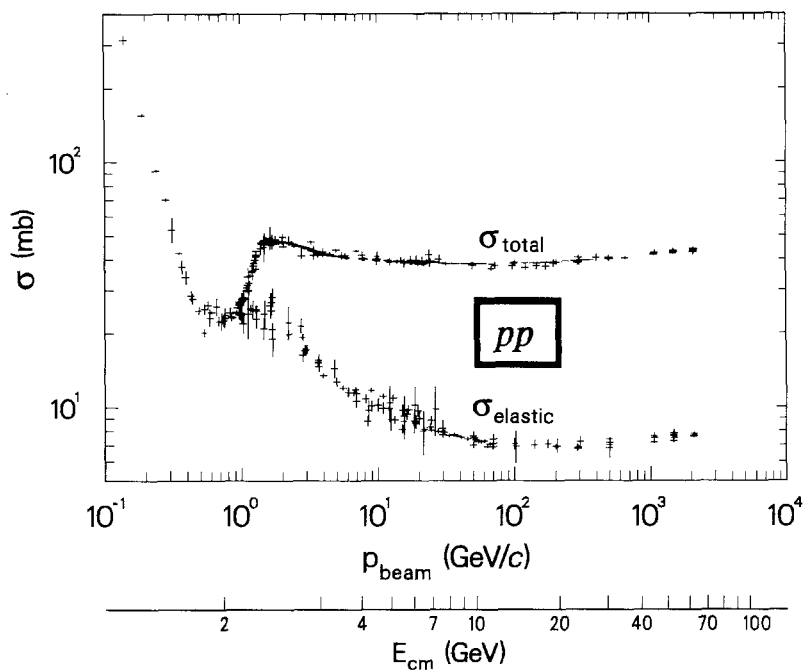
Hadronic total and elastic cross sections vs. laboratory beam momentum  $p_{\text{beam}}$  and total center-of-mass energy  $E_{\text{cm}}$ . Data courtesy A. Baldini, V. Flaminio, W.G. Moorhead, D.R.O. Morrison, and N. Rivoire, CERN; and COMPAS Group, IHEP, Serpukhov, USSR. See CERN HERA Group compilations for references.

## PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)



Hadronic total and elastic cross sections vs. laboratory beam momentum  $p_{\text{beam}}$  and total center-of-mass energy  $E_{\text{cm}}$ . Data courtesy A. Baldini, V. Flaminio, W.G. Moorhead, D.R.O. Morrison, and N. Rivoire, CERN; and COMPAS Group, IHEP, Serpukhov, USSR. See CERN HERA Group compilations for references.

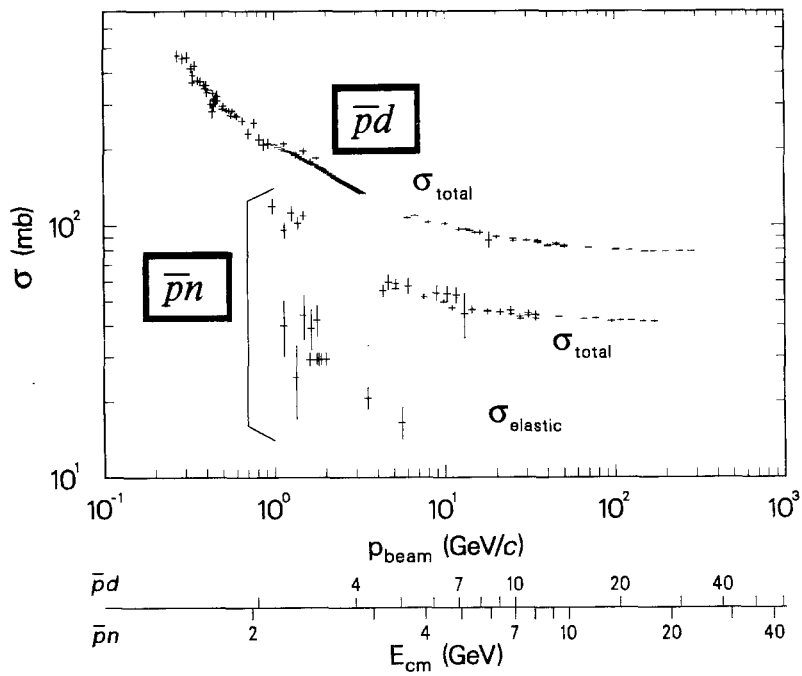
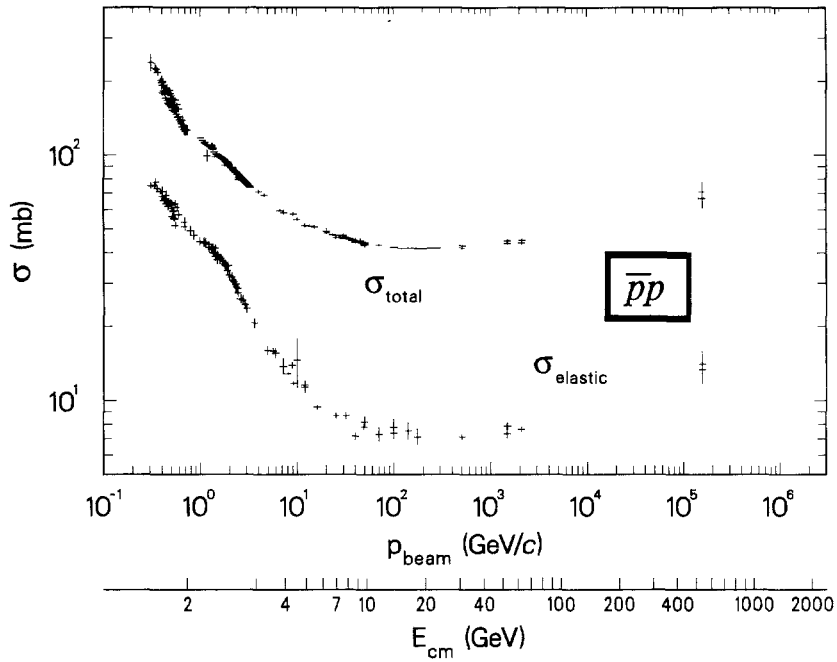
## PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)



Hadronic total and elastic cross sections vs. laboratory beam momentum  $p_{\text{beam}}$  and total center-of-mass energy  $E_{\text{cm}}$ . Data courtesy A. Baldini, V. Flaminio, W.G. Moorhead, D.R.O. Morrison, and N. Rivoire, CERN; and COMPAS Group, IHEP, Serpukhov, USSR. See CERN HERA Group compilations for references.



## PLOTS OF CROSS SECTIONS AND RELATED QUANTITIES (Cont'd)



Hadronic total and elastic cross sections vs. laboratory beam momentum  $p_{\text{beam}}$  and total center-of-mass energy  $E_{\text{cm}}$ . Data courtesy A. Baldini, V. Flaminio, W.G. Moorhead, D.R.O. Morrison, and N. Rivoire, CERN; and COMPAS Group, IHEP, Serpukhov, USSR. See CERN HERA Group compilations for references.

# Illustrative Key to the Full Listings

Name of particle ("old" name used before recent renaming scheme also given if different; see Introductory Text for details).

$a_0(1200)$   
was  $XX(1200)$

$I^G(J^{PC}) = 1^-(0^{+-})$

Particle quantum numbers (where known).

Indicates particle omitted from Particle Properties Summary Table, implying particle's existence is not confirmed.

OMITTED FROM SUMMARY TABLE

EVIDENCE NOT COMPELLING. MAY BE A KINEMATIC EFFECT.

General comments on particle.

Quantity tabulated below.

$a_0(1200)$  MASS (MeV)

Abbreviated reference for this result; full reference given below.

Code for quantity tabulated (M=mass, W=width, etc.).

M	1214.	11.	MERRILL	81	HBC	0	3.2	K-P
M	(150)(1192.)	(16.)	LYNCH	81	HBC	±	2.7	PI-P
M	LYNCH DATA HAS QUESTIONABLE BACKGROUND SUBTRACTION							
M	(1198.)	(10.)	PIERCE	83	ASPK	±	2.1	K-P
M	(1208.)		FENNER	83	HBC	0	4.2	PI+P
M	80 1210.	8.	SMITH	85	MMS	-	3.5	PI-P
M	SUPERSEDES EARLIER RESULT							
M	AVG	1206.9						
M								AVERAGE

Measurement technique (see abbreviations on next page).

Symbol used to key together data and related comments.

Number of events above background.

Charge(s) of particle(s) detected.

Measured value (parentheses indicate value not used in average; see Introductory Text for explanations).

$a_0(1200)$  WIDTH (MeV)

Reaction producing particle, or general comments.

± error in measured value (= field blank if error symmetric; parentheses on error only indicate data not used in average due to problems with error estimation).

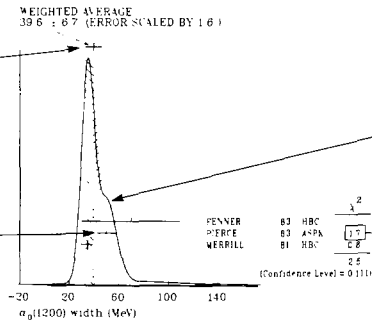
W	35.	5.	MERRILL	81	HBC	0	3.2	K-P
W	50.	8.	PIERCE	83	ASPK	±	2.1	K-P
W	70.	40.	FENNER	83	HBC	0	4.2	PI+P
W	(60.)	OR LESS	SMITH	85	MMS	-	3.5	PI-P
W	AVG	39.6						
W								AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)
W								(SEE IDEOGRAM BELOW)

"Change bar" indicates result added or changed since previous edition.

Scale factor > 1 indicates possibly inconsistent data.

Average value (and error) of quantity tabulated.

Top "data point" indicates average; width of error bar (and shaded pattern below) is ± error on average, scaled by "scale factor."



Ideogram to display possibly inconsistent data; curve is sum of Gaussians, one for each experiment (area of Gaussian = 1; error; width of Gaussian = ± error). See Introductory Text for discussion.

Value and error for each experiment.

Contribution of experiment to  $\chi^2$  (if no entry present, experiment not used in calculating  $\chi^2$  or scale factor because of very large error).

## $a_0(1200)$ PARTIAL DECAY MODES

Partial decay mode (labeled by  $P_i$ ).

P1  $a_0(1200) \rightarrow 3\pi$   
P2  $a_0(1200) \rightarrow K\bar{K}$

DECAY MASSES  
140+ 140+ 140  
494+ 494

Representative masses of decay products (used for calculating decay momentum in last column of Particle Property Summary Tables).

## $a_0(1200)$ BRANCHING RATIOS

Branching ratio (labeled by  $R_j$ ).

$a_0(1200) \rightarrow (3\pi)/total$								(P1)
R1	.66	.02	MERRILL	81	HBC	0	3.2	K-P
R1	L (.68)	(.03)	LYNCH	81	HBC	±	2.7	PI-P
R1	LYNCH DATA HAS QUESTIONABLE BACKGROUND SUBTRACTION							
R1	R1 FIT $0.675 \pm 0.012$ FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)							
$a_0(1200) \rightarrow (K\bar{K})/total$								(P2)
R2	.35	.05	PIERCE	83	ASPK	±	2.1	K-P
R2	R2 FIT $0.325 \pm 0.012$ FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)							
$a_0(1200) \rightarrow (K\bar{K})/(3\pi)$								(P2)/(P1)
R3	.50	.03	FENNER	83	HBC	0	4.2	PI+P
R3	.41	.04	SMITH	85	MMS	-	3.5	PI-P
R3	R3 AVG $0.468 \pm 0.043$ AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8)							
R3	R3 FIT $0.480 \pm 0.027$ FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)							

Branching ratio  $R_j$  in terms of partial decay mode fractions  $P_i$  above.

References, ordered by year, then author.

MERRILL 81 PRL 47 143  
LYNCH 81 PR 024 610  
PIERCE 83 PL 1238 230  
FENNER 83 NP 8213 372  
SMITH 85 PRL 55 14

A. MERRILL  
B. LYNCH  
N. PIERCE  
D. FENNER, B. BEANE  
J. SMITH

(SACLAY+CERN) [LJP]  
(BNL)  
(FNAL)  
(NYSE+AMEX)  
(SLAC)

Author(s).

Quantum number determinations in this reference.

Abbreviated reference form used on data entries above.

Journal, report, preprint, etc. (see abbreviations on next page).

Institution(s) of author(s) (see abbreviations on next page).

## Abbreviations Used in the Full Listings

### Journals

ADVP	Advances in Physics
AFIS	Anales de Fisica
ANP	Annals of Physics
APAH	Acta Phys. Acad. Hungarica
APJ	Astrophysical Journal
APP	Acta Physica Polonica
ARNPS	Annual Review of Nuclear & Particle Science
ARNS	Annual Review of Nuclear Science
BAPS	Bulletin of the American Physical Society
BASUP	Bulletin of the Academy of Science, USSR (Physics)
CJP	Canadian Journal of Physics
CNPP	Comments on Nuclear and Particle Physics
CZJP	Czechoslovak Journal of Physics
JAP	Journal of Applied Physics
JETP	English Translation of Soviet Physics JETP
JETPL	Letters of Soviet Physics JETP
JP	Journal of Physics (A,B,G)
JPSJ	Journal of the Physical Society of Japan
LNC	Letters to Nuovo Cimento
NAT	Nature
NC	Nuovo Cimento
NIM	Nuclear Instruments and Methods
NP	Nuclear Physics
PL	Physics Letters
PN	Particles and Nuclei
PPSL	Proc. of the Physical Society of London
PR	Physical Review
PRAM	Pramana
PRL	Physical Review Letters
PRPL	Physics Reports (Physics Letters C)
PRSE	Proc. of the Royal Society of Edinburgh
PRSL	Proc. of the Royal Society of London
PS	Physica Scripta
PTP	Progress of Theoretical Physics
RA	Radiochimica Acta
RMP	Reviews of Modern Physics
RNC	La Rivista del Nuovo Cimento
RPP	Reports on Progress in Physics
RRP	Revue Romaine de Physique
SCI	Science
SJNP	Soviet Journal of Nuclear Physics
SPU	Soviet Physics - Uspekhi
ZNAT	Zeitschrift fur Naturforschung
ZPHY	Zeitschrift fur Physik
GAM2	IHEP hodoscope Cerenkov $\gamma$ calorimeter GAMS-2000
GAM4	CERN hodoscope Cerenkov $\gamma$ calorimeter GAMS-4000
GOLI	CERN Goliath spectrometer
HBC	Hydrogen bubble chamber
HDDB	Hydrogen and deuterium bubble chambers
HEBC	Helium bubble chamber
HLBC	Heavy-liquid bubble chamber
HRS	SLAC high-resolution spectrometer
HYBR	Hybrid: bubble chamber + electronics
INDU	Magnetic induction
IPWA	Energy-independent partial-wave analysis
JADE	JADE detector at DESY
LASS	Large-angle superconducting solenoid spectrometer at SLAC
LENA	Nonmagnetic lead-glass NaI detector at DORIS
MAC	MAC detector at PEP/SLAC
MBR	Molecular beam resonance technique
MICA	Underground mica deposits
MMS	Missing mass spectrometer
MPS	Multiparticle spectrometer at BNL
MPSF	Multiparticle spectrometer at Fermilab
MPWA	Model-dependent partial-wave analysis
MRKJ	Mark-J detector at DESY
MRS	Magnetic resonance spectrometer
NEUL	Neuland large-angle neutrino spectrometer
OLYA	Detector at VEPP-4, Novosibirsk
OMEG	CERN OMEGA spectrometer
OSPK	Optical spark chamber
PBC	Propane bubble chamber
PLAS	Plastic detector
PLUT	DESY PLUTO detector
PWA	Partial-wave analysis
REDE	Resonance depolarization
RVUE	Review of previous data
SFM	CERN split-field magnet
SILI	Silicon detector
SMAG	SPEAR magnetic detector
SMK2	SLAC Mark-II detector
SMK3	SLAC Mark-III detector
SPEC	Spectrometer
SPRK	Spark chamber
STRC	Streamer chamber
TASS	DESY TASSO detector
THEO	Theoretical or heavily model-dependent result
TPC	TPC detector at PEP/SLAC
TPS	Tagged photon spectrometer at Fermilab
UA1	UA1 detector at CERN
UA2	UA2 detector at CERN
UA5	UA5 detector at CERN
WIRE	Wire chamber
XEBC	Xenon bubble chamber

### Conferences

Conferences are generally referred to by the location at which they were held (e.g., HAMBURG, TORONTO, CORNELL, BRIGHTON, etc.).

### Measurement Techniques (i.e., Detectors and Methods of Analysis)

AEMS	Argonne effective mass spectrometer
ARG	ARGUS detector at DORIS
ASP	Anomalous single-photon detector
ASPK	Automatic spark chambers
BDMP	Beam dump
BEBC	Big European bubble chamber at CERN
BIS2	BIS-2 spectrometer at Serpukhov
BONA	Bonanza nonmagnetic detector at DORIS
BPWA	Barrellet-zero partial-wave analysis
CALO	Calorimeter
CBAL	Crystal Ball detector at SLAC-SPEAR or DORIS
CC	Cloud chamber
CELL	CELLO detector at DESY
CHRM	CHARM neutrino detector at CERN
CIBS	CERN-IHEP boson spectrometer
CLEO	Cornell magnetic detector at CESR
CNTR	Counters
COSM	Cosmology and astrophysics
CUSB	Columbia U. - Stony Brook segmented NaI detector at CESR
DASP	DESY double-arm spectrometer
DBC	Deuterium bubble chamber
DLCO	DELCO detector at SLAC-SPEAR or SLAC-PEP
DMI	Detector at Orsay DMI collider
DPWA	Energy-dependent partial-wave analysis
ELEC	Electronic combination
EMC	European muon collaboration detector at CERN
EMUL	Emulsions
FBC	Freon bubble chamber
FIT	Fit to previously existing data
FRAB	ADONE $BB$ group detector
FRAG	ADONE $\gamma\gamma$ group detector
FRAM	ADONE MEA group detector
FRBC	Freon bubble chamber

### Institutions

AACH	Technische Univ. Aachen	Aachen, West Germany
AARH	Univ. of Aarhus	Aarhus, Denmark
ABO	Abo Akademi	Abo, Finland
ADEL	Adelphi Univ.	Garden City, NY, USA
AERE	Atomic Energy Res. Estab.	Harwell, Berks., England
AICH	Aichi Univ. of Education	Kariya, Aichi Pref., Japan
AIKO	Inst. Kernphys. Onderzoek	Amsterdam, Netherlands
ALAH	Univ. of Alabama at Huntsville	Huntsville, AL, USA
ALBA	State Univ. of New York at Albany	Albany, NY, USA
AMST	Univ. of Amsterdam	Amsterdam, Netherlands
ANIK	Amsterdam NIKHEF	Amsterdam, Netherlands
ANKA	Middle East Technical Univ.	Ankara, Turkey
ANL	Argonne National Lab.	Argonne, IL, USA
ARIZ	Univ. of Arizona	Tucson, AZ, USA
ARZS	Arizona State Univ.	Tempe, AZ, USA
ATEN	Nuclear Res. Centre Demokritos	Athens, Greece
ATHU	Univ. of Athens	Athens, Greece
AUCK	Univ. of Auckland	Auckland, New Zealand
BARC	Univ. de Barcelona	Barcelona, Spain
BARI	Univ. di Bari	Bari, Italy
BART	Bartol Research Foundation	Swarthmore, PA, USA
BASL	Univ. of Basel	Basel, Switzerland
BAYR	Univ. Bayreuth	Bayreuth, West Germany
BELG	Inst. Interuniv. des Sci. Nuc.	Bruxelles, Belgium
BELL	Bell Labs.	Murray Hill, NJ, USA
BERG	Univ. of Bergen	Bergen, Norway
BERL	Inst. Hochenergiephys. DAW	Berlin-Zeuthen, East Germany
BERN	Univ. Bern	Bern, Switzerland
BGNA	Univ. di Bologna	Bologna, Italy
BHAB	Bhabha Atomic Research Center	Bombay, India
BHEP	Inst. of High Energy Physics	Beijing, China
BIEL	Univ. Bielefeld	Bielefeld, West Germany
BING	State Univ. of New York at Binghamton	Binghamton, NY, USA
BIRM	Birmingham Univ.	Birmingham, England
BNL	Brookhaven National Lab.	Upton, L.I., NY, USA
BOHR	Niels Bohr Inst.	Copenhagen, Denmark

## Abbreviations Used in the Full Listings (*cont'd*)

### Institutions (*cont'd*)

BOIS	Boise State Univ.	BOIS, ID, USA	ILL	Univ. of Illinois	Urbana, IL, USA
BOMB	Univ. of Bombay	Bombay, India	ILLC	Univ. of Illinois at Chicago	Chicago, IL, USA
BONN	Univ. Bonn	Bonn, West Germany	ILLG	Inst. Laue-Langevin	Grenoble, France
BORD	Univ. de Bordeaux	Bordeaux, France	IND	Indiana Univ.	Bloomington, IN, USA
BOST	Boston Univ.	Boston, MA, USA	INEL	Idaho National Engineering Lab.	Idaho Falls, ID, USA
BRAN	Brandeis Univ.	Waltham, MA, USA	INFN	Ist. Nazionale di Fisica Nucleare	Roma, Italy
BRCO	Univ. of British Columbia	Vancouver, BC, Canada	INNS	Phys. Inst., Univ. Innsbruck	Innsbruck, Austria
BRIS	H. H. Wills Phys. Lab., U. of Bristol	Bristol, England	INRM	Inst. for Nuclear Research	Moscow, USSR
BROW	Brown Univ.	Providence, RI, USA	INUS	Inst. for Nuclear Study at Tokyo Univ.	Tokyo, Japan
BRUX	Univ. Libre de Bruxelles	Bruxelles, Belgium	IOFF	Ioffe Inst. of Physics and Tech.	Leningrad, USSR
BUCH	Bucharest State Univ.	Bucharest, Romania	IOWA	Univ. of Iowa	Iowa City, IA, USA
BUDA	Central Research Inst. of Physics	Budapest, Hungary	IPCR	Inst. of Physical and Chemical Research	Saitama-ken, Japan
BUFF	State Univ. of New York at Buffalo	Buffalo, NY, USA	IPN	Inst. de Phys. Nucleaire	Orsay, France
BURE	Inst. des Hautes Etudes Sci.	Bures-sur-Yvette, France	IPNP	Inst. de Physique Nucleaire	Paris, France
CAEN	Lab. de Phys. Corpusculaire	Caen, France	IPPC	Inst. for Particle Physics of Canada	Montreal, PQ, Canada
CAGL	Cagliari Univ.	Cagliari, Italy	IRAD	Inst. du Radium	Paris, France
CAIW	Carnegie Inst. of Washington	Washington, DC, USA	ISU	Iowa State Univ.	Ames, IA, USA
CAMB	Cambridge Univ.	Cambridge, England	ITEP	Inst. for Theor. and Exp. Phys.	Moscow, USSR
CANB	Australian National Univ.	Canberra, Australia	ITHA	Ithaca College	Ithaca, NY, USA
CARL	Carleton Univ.	Ottawa, ON, Canada	ITPU	Inst. for Theoretical Physics	Utrecht, Netherlands
CARN	Carnegie-Mellon Univ.	Pittsburgh, PA, USA	IUJU	Indiana U. - Purdue U. at Indianapolis	Indianapolis, IN, USA
CASE	Case Western Reserve Univ.	Cleveland, OH, USA	JAGL	Jagellonian Univ.	Cracow, Poland
CATH	Catholic Univ. of America	Washington, DC, USA	JHU	Johns Hopkins Univ.	Baltimore, MD, USA
CAVE	Cavendish Lab., Cambridge Univ.	Cambridge, England	JINR	Joint Inst. for Nucl. Research	Dubna, USSR
CCAC	Community College of Allegheny County	Pittsburgh, PA, USA	KAGO	Kagoshima Univ.	Kagoshima, Japan
CDEF	College de France	Paris, France	KANS	Univ. of Kansas	Lawrence, KS, USA
CEA	Cambridge Electron Accel.	Cambridge, MA, USA	KARL	Univ. Karlsruhe	Karlsruhe, West Germany
CENG	CEN, Grenoble	Grenoble, France	KAZA	Kazakh Academy of Science	Alma-Ata, USSR
CERN	European Org. for Nuclear Research	Geneva, Switzerland	KEK	Nat. Lab for High Energy Phys., Japan	Tsukuba-gun, Japan
CHIC	Univ. of Chicago	Chicago, IL, USA	KENT	Kent Univ. at Canterbury, Kent	Canterbury, England
CINC	Univ. of Cincinnati	Cincinnati, OH, USA	KEYN	Open Univ.	Milton Keynes, England
CIT	Calif. Inst. of Technology	Pasadena, CA, USA	KHAR	Phys.-Tech. Inst., Acad. Sci., Ukr.SSR	Kharkov, USSR
CLEV	Cleveland State Univ.	Cleveland, OH, USA	KIAE	Kurchatov Inst. of Atomic Energy	Moscow, USSR
CNRC	Canadian National Research Council	Ottawa, ON, Canada	KIEL	Kiel Univ.	Kiel, West Germany
COLO	Univ. of Colorado	Boulder, CO, USA	KIEV	Physical-Technical Inst.	Kiev, USSR
COLU	Columbia Univ.	New York, NY, USA	KINK	Kinki Univ.	Osaka, Japan
CORN	Cornell Univ.	Ithaca, NY, USA	KNTY	Univ. of Kentucky	Lexington, KY, USA
COSU	Colorado State Univ.	Fort Collins, CO, USA	KOBE	Kobe Univ.	Kobe, Japan
CRAC	Inst. for Nuclear Research	Cracow, Poland	KONS	B. P. Konstantinov Inst. of Nucl. Phys.	USSR
CUNY	City Univ. of New York	New York, NY, USA	KYOT	Kyoto Univ.	Kyoto, Japan
CURI	Laboratoire Joliot-Curie	Paris, France	LALO	Linear Accelerator Lab, Orsay	Orsay, France
DALH	Dalhousie Univ.	Halifax, NS, Canada	LANC	Lancaster Univ.	Lancaster, England
DARE	Daresbury Nuclear Physics Lab.	Daresbury, England	LANL	U.C. Los Alamos National Lab.	Los Alamos, NM, USA
DELH	Univ. of Delhi	Delhi, India	LAPP	Lab. d'Annecy de Phys. des Particules	Annecy, France
DESY	Deutsches Elektronen-Synchrotron	Hamburg, West Germany	LASL	U.C. Los Alamos Scientific Lab.	Los Alamos, NM, USA
DOE	U.S. Department of Energy	Washington, DC, USA	LAUS	Univ. of Lausanne	Lausanne, Switzerland
DORT	Univ. Dortmund	Dortmund, West Germany	LBL	U.C. Lawrence Berkeley Lab.	Berkeley, CA, USA
DUKE	Duke Univ.	Durham, NC, USA	LCGT	Lab. di Cosmo-Geofisica del CNR	Torino, Italy
DURH	Univ. of Durham	Durham, England	LEBD	Lebedev Physics Inst.	Moscow, USSR
DUUC	University College	Dublin, Ireland	LEED	Univ. of Leeds	Leeds, England
EDIN	Univ. of Edinburgh	Edinburgh, Scotland	LEHI	Lehigh Univ.	Bethlehem, PA, USA
EFT	Enrico Fermi Inst. for Nucl. Studies	Chicago, IL, USA	LEHM	Herbert H. Lehman College	Bronx, NY, USA
ELMT	Elmhurst College	Elmhurst, IL, USA	LEID	Inst. Lorentz	Leiden, Netherlands
EPOL	Ecole Polytechnique	Palaiseau, France	LEMO	Le Moyne College	Syracuse, NY, USA
ERLA	Univ. Erlangen-Nurnberg	Erlangen, West Germany	LENI	Inst. of Nucl. Phys., USSR Acad. Sci.	Leningrad, USSR
ETH	Swiss Federal Inst. of Technology	Zurich, Switzerland	LIBH	Lab. Interuniv. Belge High Eng.	Bruxelles, Belgium
FIRZ	Univ. di Firenze	Firenze, Italy	LINZ	Linz Inst. fur Physik, Kepler Hoch.	Linz, Austria
FISK	Fisk Univ.	Nashville, TN, USA	LISB	Univ. de Lisboa	Lisboa, Codex, Portugal
FLOR	Univ. of Florida	Gainesville, FL, USA	LIVP	Liverpool Univ.	Liverpool, England
FNAL	Fermi National Accelerator Lab.	Batavia, IL, USA	LJUB	Univ. of Ljubljana	Ljubljana, Yugoslavia
FOM	Found. for Fundamental Res. on Matter	Utrecht, Netherlands	LLL	Lawrence Livermore Lab.	Livermore, CA, USA
FRAS	Lab. Nazionali del C.N.E.N.	Frascati, Italy	LOIC	Imperial Col. of Sci. and Tech.	London, England
FREI	Univ. of Freiburg	Freiburg, West Germany	LOQM	Queen Mary College	London, England
FSU	Florida State Univ.	Tallahassee, FL, USA	LOUC	University College	London, England
GENO	Univ. di Genova	Genova, Italy	LOWC	Westfield College	London, England
GESC	General Electric Res. and Dev. Center	Schenectady, NY, USA	LPNP	Lab. de Phys. Nucl. et Hautes Energies	Paris, France
GEVA	Univ. de Geneve	Geneva, Switzerland	LPTP	Lab. de Phys. Theor. et Hautes Energies	Paris, France
GLAS	Univ. of Glasgow	Glasgow, Scotland	LRL	U.C. Lawrence Berkeley Lab.	Berkeley, CA, USA
GMAS	George Mason Univ.	Fairfax, VA, USA	LSU	Louisiana State Univ.	Baton Rouge, LA, USA
GRAZ	Univ. Graz	Graz, Austria	LUND	Univ. I Lund	Lund, Sweden
GREN	Inst. des Sci. Nuc., Univ. de Grenoble	Grenoble, France	LVLN	Univ. Catholique de Louvain	Louvain-La-Neuve, Belgium
GSCO	Geological Survey of Canada	Ottawa, ON, Canada	LYON	Univ. de Lyon	Villeurbanne, France
GUEL	Guelph Univ.	Guelph, ON, Canada	MADR	Junta de Energia Nuclear	Madrid, Spain
HAIF	Technion - Israel Inst. of Technology	Haifa, Israel	MADU	Univ. Autonome de Madrid	Madrid, Spain
HAMB	Univ. Hamburg	Hamburg, West Germany	MANI	Univ. of Manitoba	Winnipeg, MB, Canada
HARV	Harvard Univ.	Cambridge, MA, USA	MANZ	Univ. Mainz	Mainz, West Germany
HAWA	Univ. of Hawaii	Honolulu, HI, USA	MARS	Center National de la Recherche Sci.	Marseille, France
HEBR	Hebrew Univ.	Jerusalem, Israel	MASA	Univ. of Massachusetts	Amherst, MA, USA
HEID	Univ. Heidelberg	Heidelberg, West Germany	MCGI	McGill Univ.	Montreal, PQ, Canada
HELS	Helsingin Yliopisto	Helsinki, Finland	MCHS	Univ. Manchester	Manchester, England
HIRO	Hiroshima Univ.	Hiroshima, Japan	MCMS	McMaster Univ.	Hamilton, ON, Canada
HOUH	Univ. of Houston	Houston, TX, USA	MEIS	Meisel Univ.	Hino, Tokyo, Japan
HPC	Hewlett-Packard Corp.	Cupertino, CA, USA	MELB	Univ. of Melbourne	Parkville, Australia
IAS	Inst. for Advanced Study	Princeton, NJ, USA	MHCO	Mount Holyoke College	South Hadley, MA, USA
IBAR	Ibaraki Univ., Mito	Ibaraki-ken, Japan	MICH	Univ. of Michigan	Ann Arbor, MI, USA
IBM	International Business Machines	Palo Alto, CA, USA	MILA	Univ. di Milano	Milano, Italy
IFRJ	Inst. de Fisica, Rio de Janeiro	Rio de Janeiro, Brazil	MINN	Univ. of Minnesota	Minneapolis, MN, USA
IIT	Illinois Inst. of Tech.	Chicago, IL, USA	MIT	Massachusetts Inst. of Technology	Cambridge, MA, USA

## Abbreviations Used in the Full Listings (*cont'd*)

### Institutions (*cont'd*)

MONP	Univ. de Montpellier	Montpellier, France	SIN	Swiss Inst. of Nuclear Research	Villigen, Switzerland
MONS	Univ. de l'Etat, Mons	Mons, Belgium	SLAC	Stanford Linear Accel. Center	Stanford, CA, USA
MONT	Univ. de Montreal	Montreal, PQ, Canada	SMAS	Southeastern Massachusetts Univ.	North Dartmouth, MA, USA
MOSU	Moscow State Univ.	Moscow, USSR	SOFI	Bulgarian Acad. of Sci.	Sofia, Bulgaria
MPIH	Max Planck Inst. fur Kernphysik	Heidelberg, West Germany	STAN	Stanford Univ.	Stanford, CA, USA
MPIM	Max Planck Inst. fur Phys.-Astrophys.	Munich, West Germany	STEV	Stevens Inst. of Tech.	Hoboken, NJ, USA
MSU	Michigan State Univ.	East Lansing, MI, USA	STLO	St. Louis Univ.	St. Louis, MO, USA
MTHO	Mt. Holyoke College	South Hadley, MA, USA	STOH	Stockholm Univ.	Stockholm, Sweden
MULH	Centre Univ. du Haut-Rhin	Mulhouse, France	STON	State Univ. of New York at Stony Brook	Stony Brook, L.I., NY, USA
MUNI	Univ. of Munich	Munich, West Germany	STRB	Centre des Res. Nucleaires	Strasbourg, France
MURA	Midwestern Univ. Research Assoc.	Stroughton, WI, USA	SUSS	Univ. of Sussex	Falmer, Brighton, England
NAGO	Nagoya Univ.	Nagoya, Japan	SYDN	Univ. of Sydney	Sydney, Australia
NAPL	Univ. di Napoli	Napoli, Italy	SYRA	Syracuse Univ.	Syracuse, NY, USA
NASA	NASA, Goddard Space Flight Center	Greenbelt, MD, USA	TAMU	Texas A and M Univ.	College Station, TX, USA
NBS	U.S. National Bureau of Standards	Washington, DC, USA	TBLI	Tbilisi State Univ.	Tbilisi, USSR
NDAM	Univ. of Notre Dame	Notre Dame, IN, USA	TELA	Univ. of Tel-Aviv	Tel-Aviv, Israel
NEAS	Northeastern Univ.	Boston, MA, USA	TEMP	Temple Univ.	Philadelphia, PA, USA
NEUC	Univ. de Neuchatel	Neuchatel, Switzerland	TENN	Univ. of Tennessee	Austin, TN, USA
NIJM	R. K. Univ. Nijmegen	Nijmegen, Netherlands	TEXA	Univ. of Texas	Austin, TX, USA
NMSU	New Mexico State Univ.	Las Cruces, NM, USA	THES	Univ. of Thessaloniki	Thessaloniki, Greece
NORD	Nordisk Inst. for Teor. Atomfys.	Copenhagen, Denmark	TIFR	Tata Inst. of Fundamental Research	Bombay, India
NOVO	Inst. of Nucl. Phys.	Novosibirsk, USSR	TINT	Tokyo Inst. of Technology	Tokyo, Japan
NPOL	Northern Polytechnic	London, England	TMSK	Nucl. Phys. Inst., Tomsk Polytech Inst.	Tomsk, USSR
NRL	Naval Research Laboratory	Washington, DC, USA	TMU	Tokyo Metropolitan Univ.	Tokyo, Japan
NSF	U.S. National Science Foundation	Washington, DC, USA	TNTO	Univ. of Toronto	Toronto, ON, Canada
NTUA	National Technical Univ.	Athens, Greece	TOHO	Tohoku Univ.	Sendai, Japan
NWES	Northwestern Univ.	Evanston, IL, USA	TOKY	Univ. of Tokyo	Tokyo, Japan
NYU	New York Univ.	New York, NY, USA	TORI	Univ. di Torino	Torino, Italy
OHIO	Ohio Univ.	Athens, OH, USA	TRIK	Rikkyo Univ.	Tokyo, Japan
OKAY	Okayama Univ.	Okayama, Japan	TRIN	Trinity College	Dublin, Ireland
OKLA	Univ. of Oklahoma	Norman, OK, USA	TRIU	TRIUMF, Univ. of British Columbia	Vancouver, BC, Canada
OKSU	Oklahoma State Univ.	Stillwater, OK, USA	TRST	Univ. di Trieste	Trieste, Italy
OREG	Univ. of Oregon	Eugene, OR, USA	TSUK	Univ. of Tsukuba	Tsukuba, Japan
ORNL	Oak Ridge National Lab.	Oak Ridge, TN, USA	TTAM	Tamagawa Univ.	Tokyo, Japan
ORSA	Univ. de Paris, Fac. des Sci.	Orsay, France	TUFT	Tufts Univ.	Medford, MA, USA
OSAK	Osaka Univ.	Osaka, Japan	TWAS	Waseda Univ.	Tokyo, Japan
OSKC	Osaka City Univ.	Osaka, Japan	UBEL	Univ. of Belgrade	Belgrade, Yugoslavia
OSLO	Oslo Univ.	Oslo, Norway	UCB	Univ. of Calif. at Berkeley	Berkeley, CA, USA
OSU	Ohio State Univ.	Columbus, OH, USA	UCD	Univ. of Calif. at Davis	Davis, CA, USA
OTTA	Univ. of Ottawa	Ottawa, ON, Canada	UCI	Univ. of Calif. at Irvine	Irvine, CA, USA
OXF	Oxford Univ.	Oxford, England	UCLA	Univ. of Calif. at Los Angeles	Los Angeles, CA, USA
PADO	Univ. di Padova	Padova, Italy	UCND	Union Carbide Nuclear Division	Oak Ridge, TN, USA
PATR	Univ. of Patras	Patras, Greece	UCR	Univ. of Calif. at Riverside	Riverside, CA, USA
PAVI	Univ. di Pavia	Pavia, Italy	UCSB	Univ. of Calif. at Santa Barbara	Santa Barbara, CA, USA
PENN	Univ. of Pennsylvania	Philadelphia, PA, USA	UCSC	Univ. of Calif. at Santa Cruz	Santa Cruz, CA, USA
PGIA	Univ. di Perugia	Perugia, Italy	UCSD	Univ. of Calif. at San Diego	La Jolla, CA, USA
PHIL	Philippis Univ.	Marburg, West Germany	UDCF	Univ. de Clermont-Ferrand	Aubiere, France
PISA	Univ. di Pisa	Pisa, Italy	UMD	Univ. of Maryland	College Park, MD, USA
PITT	Univ. of Pittsburgh	Pittsburgh, PA, USA	UNCS	Union College	Schenectady, NY, USA
PNL	Pacific Northwest Lab.	Richland, WA, USA	UNM	Univ. of New Mexico	Albuquerque, NM, USA
PPA	Princeton-Penn. Proton Accel.	Princeton, NJ, USA	UOEH	Univ. of Occup. and Environ. Health	Kitakyushu, Japan
PRAG	Inst. of Physics, CSAV	Prague, Czechoslovakia	UPNJ	Uppsala College	East Orange, NJ, USA
PRIN	Princeton Univ.	Princeton, NJ, USA	UPPS	Gustaf Werner Inst.	Uppsala, Sweden
PSLL	Physical Science Lab.	Las Cruces, NM, USA	USC	Univ. of Southern California	Los Angeles, CA, USA
PUCB	Pontificia Univ. Catolica	Rio de Janeiro, Brazil	USCC	Univ. of South Carolina	Columbia, SC, USA
PURD	Purdue Univ.	Lafayette, IN, USA	USTL	Univ. Sci. et Tech. du Languedoc	Montpellier, France
QUKI	Queens Univ.	Kingston, ON, Canada	UTAH	Univ. of Utah	Salt Lake City, UT, USA
RAL	Rutherford Appleton Lab. (formerly RL)	Chilton, Did., Berks., England	UTRE	Univ. of Utrecht	Utrecht, Netherlands
REGE	Univ. Regensburg	Regensburg, West Germany	VAND	Vanderbilt Univ.	Nashville, TN, USA
REHO	Weizmann Inst. of Sci.	Rehovoth, Israel	VICT	Univ. of Victoria	Victoria, BC, Canada
RHEL	Rutherford High Energy Lab.	Chilton, Did., Berks., England	VIEN	Inst. for High Energy Physics, A. A. S.	Vienna, Austria
RICE	William Marsh Rice Univ.	Houston, TX, USA	VIRG	Univ. of Virginia	Charlottesville, VA, USA
RISO	Research Estab. Riso	Roskilde, Denmark	VPI	Virginia Polytechnic Inst./State Univ.	Blacksburg, VA, USA
RL	Rutherford Lab. (formerly RHEL)	Chilton, Did., Berks., England	VRUJ	Vrije Univ.	Amsterdam, Netherlands
RMCS	Royal Military College of Science	Shrivenham, England	WARS	Univ. of Warsaw	Warsaw, Poland
ROCH	Univ. of Rochester	Rochester, NY, USA	WASH	Univ. of Washington	Seattle, WA, USA
ROCK	Rockefeller Univ.	New York, NY, USA	WAYN	Wayne State Univ.	Detroit, MI, USA
ROMA	Univ. di Roma	Roma, Italy	WIEN	Univ. Wien	Wien, Austria
ROSE	Rose Polytechnic Inst.	Terre Haute, IN, USA	WILL	College of William and Mary	Williamsburg, VA, USA
RPI	Rensselaer Polytechnic Inst.	Troy, NY, USA	WINR	Warsaw Inst. of Nuclear Research	Warsaw, Poland
RUTG	Rutgers Univ.	New Brunswick, NJ, USA	WISC	Univ. of Wisconsin	Madison, WI, USA
SACL	Cntr. d'Etudes Nucl. Saclay	Gif-sur-Yvette, France	WITW	Univ. of the Witwatersrand	Johannesburg, S. Africa
SAGA	Saga Univ.	Saga, Japan	WMIU	Western Michigan Univ.	Kalamazoo, MI, USA
SANI	Ist. Superiore di Sanita	Roma, Italy	WOOD	Woodstock College	Woodstock, MD, USA
SBER	San Bernardino State College	San Bernardino, CA, USA	WUPG	Gesamthochschule Wuppertal	Wuppertal, West Germany
SCUC	Univ. of South Carolina	Columbia, SC, USA	WUPP	Univ. Wuppertal	Wuppertal, West Germany
SEAT	Seattle Pacific College	Seattle, WA, USA	WURZ	Univ. Wurzburg	Wurzburg, West Germany
SEIB	Research Center Seibersdorf	Vienna, Austria	WUSL	Washington Univ.	St. Louis, MO, USA
SEOU	Korea Univ.	Seoul, Korea	WYOM	Univ. of Wyoming	Laramie, WY, USA
SERP	Inst. of High Energy Physics	Serpukov, USSR	YALE	Yale Univ.	New Haven, CT, USA
SETO	Seton Hall Univ.	South Orange, NJ, USA	YERE	Yerevan Physics Inst.	Yerevan, Armenia, USSR
SFLA	Univ. of South Florida	Tampa, FL, USA	YOKO	Yokohama Univ.	Yokohama, Japan
SFSU	San Francisco State Univ.	San Francisco, CA, USA	YORK	York Univ.	Toronto, ON, Canada
SHEF	Univ. of Sheffield	Sheffield, England	ZAGR	Inst. Rudjer Boskovic	Zagreb, Yugoslavia
SHMP	Univ. of Southampton	Southampton, England	ZARA	Univ. of Zaragoza	Zaragoza, Spain
SIBE	Inst. of Nucl. Phys., USSR Acad. Sci.	Siberia, USSR	ZEEM	Zeevan Lab., Univ. of Amsterdam	Amsterdam, Netherlands
SIEG	Gesamthochschule Siegen	Huttenental, West Germany	ZURI	Univ. Zurich	Zurich, Switzerland



# Stable Particle Full Listings

## Z. Neutrinos

Z → (μ<sup>+</sup>μ<sup>-</sup>)/total  
R2 1 SEEN ARNISON 83 UA1 P PBAR ECM=546 GEV (P2)

Z → (e<sup>+</sup>e<sup>-</sup>γ)/total  
R4 1983 RADIATIVE Z0 DECAY EVENTS (Z0→e<sup>+</sup>e<sup>-</sup>γ, Z0→μ<sup>+</sup>μ<sup>-</sup>γ) (P3)  
R4 ARE IN EXCESS OF EXPECTED BREMSSTRAHLUNG RATE. IF  
R4 Z0→X GAMMA WITH X→e<sup>+</sup>e<sup>-</sup>, THE TWO e<sup>-</sup> GAMMA EVENTS GIVE  
R4 Mx=40-50 GEV. ELECTRON WIDTH OF X ABOUT 1 MEV. SEE ALSO SUBSECTIONS  
R4 ON W→e<sup>-</sup>NU GAM AND ON X HYPOTH IN OTHER STABLE PARTICLE SEARCHES.  
R4 A 1 SEEN IN 4 e<sup>+</sup>e<sup>-</sup> EVTS ARNISON 83 UA1 P PBAR ECM=546 GEV  
R4 B 1 SEEN IN 8 e<sup>+</sup>e<sup>-</sup> EVTS BAGNAIA 83 UA2 P PBAR ECM=546 GEV  
R4 C 1 SEEN IN 13 e<sup>+</sup>e<sup>-</sup> EVTS APPEL 86 UA2 P PBAR ECM=546+630  
R4 A ARNISON 83 EVENT HAS NEGATIVE TRACK OF 9-1 GEV WITH CALORIMETER  
R4 B BAGNAIA 83 EVENT HAS 3 BODY TOPOLOGY WITH GAMMA AND e<sup>-</sup> HIT SEPARATE  
R4 B CELLS.  
R4 C INCLUDES EVT OF BAGNAIA 83. PROB OF INTERNAL BREM = 19 PERCENT.

Z → (μ<sup>+</sup>μ<sup>-</sup>γ)/total  
R5 A 1 SEEN ARNISON 84 UA1 P PBAR ECM=546 GEV (P4)  
R5 B 2 SEEN IN 5 MU-MU- EVENTS ARNISON 84 UA1 P PBAR ECM=546 GEV  
R5 A ARNISON 84 EVENT HAS HARD PHOTON WITH ENERGY FRACTION OF 0.35 AND  
R5 A ANGLE OF 9 DEGREES FROM MU-. HARD TO INTERPRET AS INTERNAL BREM.

**CHARGE ASYMMETRY IN e<sup>+</sup>e<sup>-</sup> → μ<sup>+</sup>μ<sup>-</sup> (including rad. corrections)**

Z0A A BARTEL 82 JADE ECM=25.-36.86EV  
Z0A B FERNANDEZ 85 MAJ ECM=29 GEV  
Z0A C ALTHOFF 84 THORP 84 THORP ECM=34.4 GEV  
Z0A D ADEVA 85 MRKJ ECM=14.-46.8 GEV  
Z0A E ASH 85 MAC ECM=29 GEV  
Z0A F BARTEL 85 JADE ECM=12.-46. GEV  
Z0A G DERRICK 85 HRS ECM=29 GEV

Z0A A BARTEL 82 OBTAINED A=-0.118+-0.038 AFTER RAD. COR. AND FOR  
Z0A A COS(THETA) WITHIN (-0.8). CONSISTENT WITH STANDARD MODEL  
Z0A A PREDICTION A=-0.078).

Z0A B FERNANDEZ 83 GOT A=-.076+.018-.003.ST.MODEL(INCL.RC) GIVES A=-.060.  
Z0A C ALTHOFF 84 OBTAINED A=-.098--.023--.005, SIN(THETA/W)\*\*2-.27--.07.

Z0A D ADEVA 85 OBSERVE ASYMMETRY (STD MODEL) VS S\*\*1/2  
Z0A D -5.3+-5.0 (-1.5) 14.0 GEV  
Z0A D -11.7+-1.7 (-8.5) 34.6 GEV  
Z0A D -16.7+-4.0 (-16.0) 44.6 GEV

Z0A E ASH 85 OBSERVE -0.063+-0.008--0.002. STD MODEL GIVES -0.063  
Z0A E FROM 16000 EVENTS.

Z0A F BARTEL 85 OBSERVE ASYMMETRY (STD MODEL) VS S\*\*1/2  
Z0A F +2.4+-4.3 (-1.2) 13.9 GEV  
Z0A F -10.3+-1.6+-0.9 (-7.6) 34.4 GEV  
Z0A F -17.7+-3.8+-1.6 (-12.5) 42.4 GEV

Z0A G DERRICK 85 OBSERVE A=-0.049-0.015--0.005 CONSISTENT WITH THEORY  
Z0A G VALUE (WITH RAD. CORR'S) OF 0.059.

**CHARGE ASYMMETRY IN e<sup>+</sup>e<sup>-</sup> → τ<sup>+</sup>τ<sup>-</sup> (including rad. corrections)**

Z0T A BARTEL 85 OBSERVE ASYMMETRY (STD MODEL) VS S\*\*1/2  
Z0T A -6.0+-2.5+-1.0 (-8.8) 34.6 GEV  
Z0T A -11.8+-4.6+-1.0 (-14.8) 43.0 GEV

### REFERENCES FOR Z

BARTEL 82 PL 1088 140	JADE C. (DESY+HAMB+HEID+LANC+MCHS+RL+TOKY)
ARNISON 83 PL 1268 398	UA1 COLLAB.(AACH+LAPP+BIRM+CERN+HEL5+LOO+)
ARNISON2 83 PL 1268 273	UA1 COLLAB.(AACH+LAPP+BIRM+CERN+HEL5+LOO+)
BAGNAIA 83 PL 1298 150	UA2 COLLAB.(BERN+CERN+BOHR+ALD+PAVI+SACL)
FERNANDEZ 83 PRL 50 1238	+ (COLD+FRAS+HOUS+NEAS+STAN+SLAC+UTAH+WISC)
ALTHOFF 84 ZPHY C22 13	TASSO C. (AACH+BOHN+DESY+HAMB+LOIC+QFX+RL+)
ARNISON 84 PL 1358 290	UA1 COLLAB.(AACH+LAPP+BIRM+CERN+HEL5+LOO+)
ARNISON2 84 PL 1478 241	UA1 COLLAB.(AACH+LAPP+BIRM+CERN+HARV+HEL5+)
BAGNAIA 84 ZPHY C24 1	UA2 COLLAB.(BERN+CERN+BOHR+ORSA+PAVI+SACL)
ADEVA 85 PRL 55 665	MARK-J C (AACH+BNL+GIT+DESY+MIT+MADR+ANIK+)
ALSD 82 PRL 48 1701	MARK-J C (AACH+DESY+MIT+MADR+ANIK+DHEP)
ASH 85 PRL 55 183*	MAC C. (COLD+FRAS+HOUS+NEAS+SLAC+UTAH+WISC)
BARTEL 85 ZPHY C26 507	JADE C. (DESY+HAMB+HEID+LANC+MCHS+UMD+RHEL+)
BARTELZ 85 PL 1618 188	JADE C. (DESY+HAMB+HEID+LANC+MCHS+UMD+RHEL+)
DERRICK 85 PR 031 2352	HRS C. (ANL+INDI+MICH+PURD+LBL+SLAC)
APPEL 86 ZPHY C TO BE PUB	UA2 COLLAB.(BERN+CERN+BOHR+HEID+ORSA+PAVI+)
ARNISON 86 PL B TO BE PUB	UA1 COLLAB.(AACH+ANIK+LAPP+BIRM+CERN+HARV+)

### LEPTONS

#### NOTE ON NEUTRINOS

(by R.E. Shrock, State Univ. of New York, Stony Brook)

With the 1982 edition of this Review, the section on neutrino properties was expanded and reorganized. As before, there are listings which deal specifically with ν<sub>e</sub>,

ν<sub>μ</sub>, and ν<sub>τ</sub>. In addition, in the category of searches near the end of the Stable Particle Full Listings, we include sections which deal with correlated bounds on neutrino masses and lepton mixing but which do not pertain to any one weak eigenstate individually. Furthermore, we include constraints from cosmological and astrophysical data. (Since this Review is a compendium of data traditionally derived more or less directly from particle and nuclear physics, we have been somewhat less comprehensive in our coverage of astrophysical data.)

To summarize the current (December 1985) status, there is no undisputed experimental evidence for nonzero neutrino masses or mixing. The ITEP reports of a nonzero m(ν<sub>e</sub>) observed in the decay <sup>3</sup>H → <sup>3</sup>He + e<sup>-</sup> + ν̄<sub>e</sub> (LUBIMOV 80, LUBIMOV 83, BORIS 85) remain the only definite claims not directly contradicted by other experiments. However, from independent theoretical analysis of these data, it has been argued (SIMPSON 84, BERKVI1 85, BERKVI2 85) that, contrary to the ITEP claims, the ITEP data do not imply a nonzero value of m(ν<sub>e</sub>).

In contrast to the other particles in this Review, the neutrinos ν<sub>e</sub>, ν<sub>μ</sub>, and ν<sub>τ</sub> are defined as weak eigenstates (that is, states which couple weakly with unit strength to e, ν, and τ) and are not, in general, states of definite mass. In the standard model, where all neutrinos are assumed to be massless and hence degenerate, it is possible to define the weak eigenstates to be simultaneously mass eigenstates. However, in the general case of possibly massive (nondegenerate) neutrinos, the weak eigenstates have no well-defined masses, but instead are linear combinations of mass eigenstates. Let us denote the charged leptons as the set {ℓ<sub>a</sub>}, a = 1, ..., n, where n ≥ 3 is the number of generations, with ℓ<sub>1</sub> ≡ e, ℓ<sub>2</sub> ≡ μ, ℓ<sub>3</sub> ≡ τ. In the standard SU(2)<sub>L</sub> × U(1) electroweak theory<sup>1</sup> the mixing of the left-handed components of the mass eigenstates (ν<sub>j</sub>)<sub>L</sub> to form the weak gauge-group eigenstates (ν<sub>ℓ<sub>a</sub></sub>)<sub>L</sub> is specified by the transformation

$$(\nu_{\ell_a})_L = \sum_{j=1}^n U_{aj}(\nu_j)_L,$$

where U<sup>†</sup> = U<sup>-1</sup>. (In the case of Dirac neutrinos there are right-handed components of the ν<sub>j</sub>, but they are singlets under the gauge group; in the case of Majorana neutrinos in the standard theory there are no independent right-handed components.) The ordering of the mass eigenbasis is defined such that U is as nearly diagonal as possible, i.e., |U<sub>ij</sub>| (no sum on j) ≥ |U<sub>jk</sub>|, k ≠ j. This does not imply that m(ν<sub>j</sub>) > m(ν<sub>k</sub>) if j >

For notation, see key on page 91.

## Stable Particle Full Listings

Neutrinos

$k$ , although this ordering might be regarded as natural in view of the similar one that obtains in the quark sector. The virtue of this convention is that a mass limit on “ $m(\nu_{\ell_a})$ ” can

be used as a definite limit on  $\nu_j$ ,  $j = a$ , the dominantly coupled mass eigenstate in  $\nu_{\ell_a}$ .

Thus, in this general case of  $n$  possibly massive (Dirac or Majorana) neutrinos, decays such as  ${}^3\text{H} \rightarrow {}^3\text{He} + e^- + \bar{\nu}_e$  and  $\pi^+ \rightarrow \mu^+ + \nu_\mu$ , which have been used to set the best bounds on the respective neutrino masses, really consist of incoherent sums of the separate decay modes  ${}^3\text{H} \rightarrow {}^3\text{He} + e^- + \bar{\nu}_j$  and  $\pi^+ \rightarrow \mu^+ + \nu_k$ , where the  $\nu_j, \nu_k$  are mass eigenstates, and the indices  $j$  and  $k$  range over the subset  $\{1, \dots, n\}$  allowed by phase space in these two respective decays.<sup>2</sup> The coupling strengths for the  $j^{\text{th}}$  modes are given for the two decays by the factors  $|U_{1j}|^2$  and  $|U_{2j}|^2$ , respectively. There are, in addition, certain kinematic factors depending on the  $m(\nu_j)$  which enter in determining the branching ratio for the  $j^{\text{th}}$  decay mode. Assuming that the off-diagonal elements of the lepton mixing matrix  $U$  are small relative to the diagonal elements, the dominantly coupled decays are the ones with coupling strength  $|U_{aj}|^2$ ,  $a = j$ , i.e.,  ${}^3\text{H} \rightarrow {}^3\text{He} + e^- + \bar{\nu}_1$  and  $\pi^+ \rightarrow \mu^+ + \nu_2$ .

It follows that the old neutrino mass limits quoted in the literature for “ $m(\nu_e)$ ”, “ $m(\nu_\mu)$ ” and “ $m(\nu_\tau)$ ” are meaningful only insofar as they are reinterpreted as limits on the corresponding mass eigenstates. Specifically, a bound such as the Bergkvist limit,<sup>3</sup> “ $m(\nu_e)$ ”  $< 60$  eV (90% CL), really constitutes a weighted limit on each of the mass eigenstates  $\nu_j$  in the weak eigenstate  $\nu_e$  which are kinematically allowed to occur in tritium decay and which are coupled with strength  $|U_{1j}|^2$  sufficiently large to make a significant contribution to the observed spectrum. It is thus certainly a limit on  $\nu_1$ . If leptonic mixing is hierarchical as quark mixing is known to be (at least for the first three generations), i.e.,  $|U_{jj}|^2 \gg |U_{jk}|^2$ ,  $j \neq k$ , then  $\nu_1$  is the only mass eigenstate significantly constrained by a bound on “ $m(\nu_e)$ .” Furthermore, a neutrino mass limit cannot be stated in isolation; it always contains some implicit dependence on the relevant lepton mixing angles. Fortunately, this dependence is relatively unimportant for the dominantly coupled decay modes, i.e.,  $e\bar{\nu}_1$ ,  $\mu\bar{\nu}_2$ , and  $\tau\bar{\nu}_3$ . Since these modes were the ones responsible for the mass limits given previously, the latter can be reinterpreted without significant complication as proper limits on  $m(\nu_j)$ ,  $j = 1, 2$ , and 3, respectively.

In addition to mass and lifetime limits, we have added data on neutrino magnetic dipole moments.

These are of interest because a massless, purely chiral (empirically, left-handed) Dirac neutrino cannot have a magnetic (or electric) dipole moment. The same is true for a Majorana neutrino, whether massless or massive, because of its defining property of being self-conjugate.

If one considers the possibility of nonzero masses for neutrinos, for consistency one must also consider the leptonic mixing which would in general occur concomitantly. Accordingly we have devoted one category in the searches section to correlated bounds on neutrino masses and lepton mixing angles. These can be divided into two types. First, there are those due to decays involving neutrinos in the final state, which must be recognized to have the possible multimode structure pointed out above. In the two most sensitive cases suggested as tests for neutrino masses and mixing,<sup>2</sup> one obtains a limit on  $m(\nu_j)$  and  $|U_{aj}|^2$  individually for each  $j$ . Second, there are those due to processes involving the propagation and subsequent interaction of neutrinos. The latter are often called neutrino “oscillation”<sup>3</sup> limits, although this term is correct only if the differences in neutrino masses are sufficiently small relative to their momenta that the propagation is effectively coherent in a quantum mechanical sense; otherwise, the individual  $\nu_j$  from a given decay such as  $\pi_{\mu 2}$  or  $K_{\mu 2}$  propagate in a measurably incoherent manner and there is no “oscillation.” Experimentalists usually present their results in terms of a simplifying model in which mixing is assumed to occur only between two neutrino species. Then the transformation equation becomes

$$\begin{pmatrix} \nu_{\ell_a} \\ \nu_{\ell_b} \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \begin{pmatrix} \nu_i \\ \nu_j \end{pmatrix}$$

Let the distance between the source of the neutrinos and their point of interaction be labeled as  $x$ , and their energy as  $E$ . Assume furthermore that the  $m(\nu_j)$  are such that the coherence assumption is valid. Then, the probability of an initial  $\nu_{\ell_a}$  being equal to  $\nu_{\ell_b}$  at time  $t$ , or equivalently (given the above assumption) at distance  $x = t$ , is

$$|\langle \nu_{\ell_b}(0) | \nu_{\ell_a}(t) \rangle|^2 = \sin^2 2\theta \sin^2 \left[ \frac{\Delta m^2 x}{4E} \right],$$

where

$$\Delta m^2 = m(\nu_i)^2 - m(\nu_j)^2.$$



# Stable Particle Full Listings

Neutrinos,  $\nu_e$

Thus, neutrino oscillation experiments cannot measure individual neutrino masses, but only differences of masses squared, and indeed these are generally weighted in a more complicated way by mixing-matrix coefficients than in the two-species model. Experimental results are presented as allowed regions on a plot, the axes of which are  $|\Delta m^2|$  and  $\sin^2 2\theta$ . These are often summarized in terms of the asymptotic limits  $|\Delta m^2|_{\max}$  for  $\sin^2 2\theta = 1$ , and  $\sin^2 2\theta$  for "large"  $|\Delta m^2|$ , i.e., sufficiently large  $|\Delta m^2|$  that the detector averages over many cycles of oscillation (or there ceases to be any coherence). We refer the reader to the original papers for the two-dimensional plots; for the purpose of these Full Listings, we shall give only the asymptotic limits.

An important question has to do with whether neutrinos are Dirac or Majorana (self-conjugate) particles. In the former case neutrinoless double beta decay,  $(Z, A) \rightarrow (Z+2, A) + e^- + e^-$ , is forbidden from occurring.<sup>4</sup> In the Majorana case it may occur if (a) neutrinos are massive and/or (b) there are right-handed leptonic currents. In the light-neutrino case an upper limit on neutrinoless double beta decay yields a correlated upper bound on the quantity

$$\bar{m} \equiv \left| \sum_{j=1}^n U_{1j}^2 m(\nu_j) \right|$$

and  $\eta$ , the fractional admixture of right-handed leptonic current.

The correlated limits given in the section on Massive Neutrinos and Lepton Mixing are in digital form. For recent compendia of limits in convenient graphical form, see Refs. 5 and 6, and Figs. 1 and 2 (pp. 332-333) of Ref. 7.

Further explanatory notes are included in the Full Listings.

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$\nu_e$

$$J = \frac{1}{2}$$

NOT IN GENERAL A MASS EIGENSTATE

$\nu_e$  "MASS" (eV)

APPLIES TO  $\nu_e(1)$ , THE PRIMARY MASS EIGENSTATE IN  $\nu_e(e)$ . WOULD ALSO APPLY TO ANY OTHER  $\nu_e(j)$  WHICH MIXES STRONGLY IN  $\nu_e(e)$  AND HAS SUFFICIENTLY SMALL MASS THAT IT CAN OCCUR IN THE RESPECTIVE DECAYS. THE NEUTRINO MASS MAY BE OF DIRAC OR MAJORANA TYPE; THE FORMER CONSERVES TOTAL LEPTON NUMBER WHILE THE LATTER VIOLATES IT. IN GENERAL, EITHER WOULD VIOLATE LEPTON FAMILY NUMBER SINCE NOTHING FORCES THE NEUTRINO MASS EIGENSTATES TO COINCIDE WITH THE NEUTRINO INTERACTION EIGENSTATES. FOR LIMITS ON MAJORANA  $\nu_e(e)$  MASS, SEE THE SECTION ON MASSIVE NEUTRINOS AND LEPTON MIXING, PART 5(C), ENTITLED SEARCHES FOR NEUTRINOLESS DOUBLE BETA DECAY.

NOTE -- THE ABBREVIATION ANU IS USED BELOW FOR NUBAR

M	(250.)	OR LESS	LANGER	52	CNTR	ANU(E), TRITIUM								
M	(500.)	OR LESS	HAMILTON	53	CNTR	ANU(E), TRITIUM								
M	(550.)	(280.)	FRIEDMAN	58	CNTR	ANU(E), TRITIUM								
M	(4100.)	OR LESS	CL=.67	BECK	68	CNTR	NU, SODIUM 22							
M	(500.)	OR LESS	CL=.90	DARIS	69	CNTR	ANU(E), TRITIUM							
M	(320.)	OR LESS	CL=.90	SALGO	69	CNTR	ANU(E), TRITIUM							
M	(60.)	OR LESS	CL=.90	BERGKVIST	72	CNTR	ANU(E), TRITIUM							
M	(86.)	OR LESS	CL=.90	RODE	72	CNTR	ANU(E), TRITIUM							
M	(100.)	OR LESS	CL=.90	PIEL	73	CNTR	ANU(E), TRITIUM							
M	(4.5E5)	OR LESS	CL=.90	CLARK	74	ASPK	KE3 DECAY							
M	(35.)	OR LESS	CL=.90	TRETYAKOV	76	SPEC	ANU(E), TRITIUM							
M	(14.)	TO 46.	CL=.99	LUBIMOV	80	SPEC	ANU(E), TRITIUM							
M	(65.)	OR LESS	CL=.95	SIMPSON	81	CNTR	ANU(E), TRITIUM							
M	(1300.)	OR LESS	CL=.90	ANDERSEN	82	CNTR	NU, HOLMIUM 163							
M	(50.)	OR LESS	CL=.90	DERBIN	83	CNTR	ANU(E), TRITIUM							
M	(500.)	OR LESS	CL=.90	JONSON	83	CNTR	NU, PLATINIUM 193							
M	(20.)	OR MORE	CL=.95	LUBIMOV	83	CNTR	ANU(E), TRITIUM							
M	(1250.)	OR LESS	CL=.95	YASUMI	83	CNTR	NU, HOLMIUM 163							
M	(20.)	TO 45.	NO CL	BORIS	85	CNTR	ANU(E), TRITIUM							
M	DARIS	69	VALUE 75E5(CL=.67)	DISAGREES	WITH	THEIR	FIG.6. WE USE							
M	D	FIG.6.												
M	L	TRETYAKOV	76	DATA	INCLUDED,	AT	LEAST	IN	PART,	IN	LUBIMOV	80.		
M	L	NOTE	THAT	LUBIMOV	83	REMARKS	THAT	THE	14	EV	LOWER	LIMIT	GOES	TO
M	L	ZERO	IF	THE	INTRINSIC	RESOLUTION	OF	THE	CONVERSION	LINE	USED	FOR		
M	L	CALIBRATION	ARE	TAKEN	INTO	ACCOUNT.	A	DETAILED	DISCUSSION	IS	GIVEN			
M	L	BY	SIMPSON	84.	SEE	ALSO	THE	DISCUSSION	OF	THE	LUBIMOV	80	RESULT	
M	L	BY	BERGKVIST	80.	WE	CONTINUE	TO	USE	UPPER	LIMIT	FROM	LUBIMOV	80	
M	L	IN	THE	STABLE	PARTICLE	SUMMARY	TABLE.							
M	Y	LIMIT	OBTAINED	BY	YASUMI	83	ASSUMES	UPPER	LIMIT	ON	Q-VALUE			
M	Y	REPORTED	BY	ANDERSEN	82.									
M	P	PRELIMINARY	RESULT	FROM	BRIGHTON	CONF.	SEE	SIMPSON	84	AND	BERGKVIST	81		
M	P	WHOSE	INDEPENDENT	THEORETICAL	ANALYSES	OF	LUBIMOV	83	DATA	DISAGREE				
M	P	WITH	LUBIMOV	83	CLAIM	OF	NONZERO	MASS.						
M	B	INDEPENDENT	THEORETICAL	ANALYSIS	BY	BERGKVIST	85	OF	BORIS	85	DATA			
M	B	(LUBIMOV	DATA)	DISAGREES	WITH	BORIS	85	CLAIM	OF	NONZERO	MASS.			



# Stable Particle Full Listings

$\nu, \mu$

M L LU 80 COMBINES DAUM 79 PI+ --> MU+ NU(MU) MEASUREMENT WITH NEW LUBO  
 M L PI+ MASS AND REPLACES DAUM 79  
 M A ABELA 84 USE PDG84 VALUE FOR PI+- MASS, IN CONJUNCTION WITH MU  
 M A MOMENTUM MEASUREMENT IN PI -> MU NU(MU) DECAY.

## $\nu_2 - \bar{\nu}_2$ MASS DIFFERENCE (MeV)

TEST OF CPT FOR A DIRAC NEUTRINO (NOT A VERY STRONG TEST)

DM (0.45) OR LESS CL=.90 CLARK 74 ASPK KMU5 DECAY

## $\nu_2$ (MEAN LIFE)/MASS (sec/eV)

T B 0 (3. E-3) OR MORE CL=.90 BELLOTTI 76 HLBC NU, CERN GGM  
 T B 1 (1.3E-2) OR MORE CL=.90 BELLOTTI 76 HLBC NUBAR, CERN GGM  
 T B 0 (2.2E-3) OR MORE CL=.90 BARNES 77 DBC NU, ANL 12F1.  
 T B 0 (1.0E-2) OR MORE CL=.90 BLIETSCHA 78 HLBC NU(MU) CERN GGM  
 T B 0 (1.7E-2) OR MORE CL=.90 BLIETSCHA 78 HLBC NUBAR(MU) CERN GGM  
 T B 0 0.11 OR MORE CL=.90 FRANK 81 CNTR NU, NUBAR LAMPF  
 T B THESE EXPERIMENTS LOOK FOR NU(MU) --> NU(E)+GAMMA OR NUBAR(MU) --> ANU(E)+GAMMA.

$$\left| \frac{v-c}{c} \right| \quad (v \approx \nu_2 \text{ VELOCITY}) \text{ (units } 10^{-4}\text{)}$$

EXPECTED TO BE ZERO FOR MASSLESS NEUTRINO

V 77 (2.0) OR LESS CL=.99 ALSPECTOR 76 SPEC >50GEV NU  
 V 26 (4.0) OR LESS CL=.99 ALSPECTOR 76 SPEC >50GEV NU  
 V 980 (0.4) OR LESS CL=.95 KALBFLEIS 79 SPEC

## $\nu_2$ MAGNETIC MOMENT (eV/gauss)

MUST VANISH FOR MAJORANA NEUTRINO OR PURELY CHIRAL MASSLESS DIRAC NEUTRINO

MM K (4.7 E-17) OR LESS KIM 74  
 MM K KIM 74 IS A THEORETICAL ANALYSIS OF NUBAR(MU) REACTION DATA.

## REFERENCES FOR $\nu_\mu$

BARKAS 56 PR 101 778 W H BARKAS, W BIRNBAUM, F M SMITH (LRL)  
 DUDZIAK 59 PR 114 356 W F DUDZIAK, R SARGENT, J VEDDER (LRL)  
 FEINBERG 63 ARNS 13 431 G FEINBERG, L M LEIDERMAN (COLUMBIA)  
 ALLCOCK 65 PPSL 85 875 G R ALLCOCK (LIVERPOOL)  
 BARDON 65 PRL 14 469 BARDON, NORTON, PEOPLES + (COOL-STONY BROOK)  
 SHAFER 65 PRL 14 923 R E SHAFER, CROWE, JENKINS (LRL)  
 BOOTH 67 PL 268 39 BOOTH, JOHNSON, WILLIAMS, WORMALD (LIVERPOOL)  
 HYMAN 67 PL 258 376 -LOKEN, PEWITT, MCKENZIE- (ANL+CORN+AMES)  
 BACKENST 71 PL 368 203 BACKENSTOSS, DANIEL, KOCH+ (CERN, KARL+HEID)  
 SHRUM 71 PL 378 114 E V SHRUM, K O H ZIOCK (UNIV OF VIRGINIA)  
 BACKENST 73 PL 438 559 BACKENSTOSS, DANIEL, KOCH+ (CERN-KARL+MUNICH)  
 CLARK 74 PR D9 533 +ELIOFF, FRISCH, JOHNSON, KERTH, SHEN + (LBL)  
 KIM 74 PR D9 3050 J.E. KIM, V.S. MATHER, S. OKUBO (ROCK)  
 ALSPECTO 76 PRL 36 837 ALSPECTOR + (BNL+PURD+CIT+FNAL+ROCK)  
 BELLOTTI 76 LNC 17 552 +CAVALLI, FIORINI, ROLLIER (MILA)  
 BARNES 77 PRL 38 1049 +GARMONY, DAUVE, FERNANDEZ + (PURD+ANL)  
 BLIETSCH 78 NP B135 205 BLIETSCHAU-(AACH+LIRH-CERN-EPOL-MILA+ORSA+)  
 KALBFLEI 79 PRL 43 1361 KALBFLEISCH, BAGGETT, FOWLER-(FNAL+PURD+BELL)  
 DAUM 79 PR 920 2692 -EATON, FROSCH, HIRSCHMANN, MCCULLOCH+ (SIN)  
 ALSO 76 PL 608 380 DAUM, DUBAL, EATON, FROSCH, MCCULLOCH+(SIN-ETH)  
 ALSO 78 PL 748 126 DAUM, EATON, FROSCH, HIRSCHMANN, + (SIN)  
 LU 80 PRL 45 1066 +DELKER, DUGAN, WU, CAFFREY+ (YALE+COLU+JHU)  
 FRANK 81 PR D24 2001 BURMAN-(LASL+YALE-MIT+SAFL+SIN-CNRC+BERN)  
 ANDERHUB 82 PL 1148 76 +BOECKLIN, HOFER, KOTTMANN+ (ETH-SIN)  
 ABELA 84 PL 1468 431 +DAUM, EATON, FROSCH, JOST, KETTLE- (SIN)

$\mu$

$$J = \frac{1}{2}$$

## $\mu$ MASS (MeV)

M (105.659) (0.002) FEINBERG 63 RVUE  
 M (105.6599) (0.0014) TAYLOR 69 RVUE USING NEW E/H  
 M C (105.6597) (0.0005) CRANE 71 CNTR INCLUDED IN COHEN73  
 M D (105.6594) (0.0004) CROWE 72 CNTR INCLUDED IN COHEN73  
 M (105.65943) (0.00035) COHEN 73 RVUE  
 M A (105.65945) (0.00033) CASPERSON 77 CNTR -  
 M K (105.65933) (0.00029) KLEMP 82 CNTR + INCL IN MARIAM 82  
 M M 105.65928 0.00029 MARIAM 82 CNTR +  
 C CRANE 71 GIVES MU/ME=206.7687(85). WE USE ME=.5110041(16)MEV.  
 M D CROWE 72 GIVES MU/ME=206.7682(5) AND USES ME=.5110034(14)MEV.  
 M A CASPERSON 77 GIVES MU/ME=206.7685(29). WE USE ME=.5110034(14)MEV.  
 M K KLEMP 82 GIVES MU/ME=206.76835(11). WE USE ME=.5110034(14)MEV.  
 M M MARIAM 82 GIVES MU/ME=206.768259(62). WE USE ME=.5110034(14)MEV.  
 M FIT 105.65916 0.00030 FROM FIT (ERROR INCL. SCALE FACTOR 1.1)

## $\mu$ MEAN LIFE ( $\mu$ sec)

T 2.198 0.001 0.001 FARLEY 62 CNTR  
 T 2.203 0.004 LUNDY 62 CNTR CONLEY=.98  
 T 2.202 0.003 0.003 ECKHAUSE 63 CNTR  
 T 2.197 0.005 0.002 MEYER 63 CNTR +  
 T 2.198 0.002 0.002 MEYER 65 CNTR -  
 T W (2.20026) (0.00081) WILLIAMS 72 CNTR +  
 T 2.1973 0.0003 DUCLOS 73 CNTR +  
 T 2.19711 0.00008 BALANDIN 74 CNTR +  
 T 2.1948 0.0010 BAILVEZ 77 CNTR + STORAGE RINGS  
 T 2.1966 0.0020 BAILVEZ 77 CNTR + STORAGE RINGS  
 T (2.197182) (0.000121) BARDIN 81 CNTR REPL. BY BARDIN 84  
 T 2.197078 0.000073 BARDIN 84 CNTR +  
 T 2.197025 0.000155 BARDIN 84 CNTR +  
 T 2.19695 0.00006 GIOVANETTI 84 CNTR +  
 T W WILLIAMS 72 MEAN LIFE MEASUREMENT WAS NOT THE PRIMARY PURPOSE OF THEIR EXPERIMENT AND DISAGREES STRONGLY WITH LATER EXPTS. NOT AVGD.  
 T W  
 T AVG 2.197033 0.000039 0.000038 AVERAGE

## $\mu^+/\mu^-$ MEAN LIFE RATIO

DT TEST OF CPT  
 DT 1.000 0.001 MEYER 63 CNTR MEAN LIFE MU+/MU-  
 DT 1.0008 0.0010 BAILVEZ 77 CNTR STORAGE RING  
 DT 1.000024 0.000078 BARDIN 84 CNTR  
 DT  
 DT AVG 1.000029 0.000078 AVERAGE

## $\mu$ ANOMALOUS MAGNETIC MOMENT (units $10^{-6} e/2m\mu$ )

MM FOR REVIEWS OF THEORY AND EXPERIMENTS, SEE HUGHES 85, KINOSHITA 84,  
 MM COMBLEY 81, FARLEY 79, AND CALMET 77.  
 MM (1162.03) (5.34) CHAPK 62 CNTR +  
 MM (1166.16) (0.31) BAILLEY 68 CNTR + STOR. RINGS  
 MM IA (1165.922) (0.009) BAILLEY 77 CNTR + STORAGE RING  
 MM I (1165.910) (0.011) BAILLEY 79 CNTR + STORAGE RING  
 MM I (1165.937) (0.012) BAILLEY 79 CNTR + STORAGE RING  
 MM I 1165.923 0.0085 BAILLEY 79 CNTR + STORAGE RING  
 MM I BAILLEY 79 IS FINAL RESULT. INCLUDES BAILLEY 77 DATA.  
 MM I WE USE MUON/PROTON MAGNETIC MOMENT RATIO = 5.183 345 2 AND  
 MM I RECALCULATE THE BAILLEY 79 VALUES  
 MM I THIRD BAILLEY 79 RESULT IS FIRST TWO COMBINED.

MWR MU+/MU- G-FACTOR RATIO MINUS ONE, (G-/G-)-1  
 MWR TEST OF CPT  
 MWR -2.6E-8 BAILLEY 79

## $\mu$ ELECTRIC DIPOLE MOMENT (units $10^{-19} e-cm$ )

FORBIDDEN BY BOTH T INVARIANCE AND P INVARIANCE

EDM B (8.6) (4.5) BAILLEY 78 CNTR + STORAGE RINGS  
 EDM B (0.8) (4.5) BAILLEY 78 CNTR + STORAGE RINGS  
 EDM B 3.7 3.4 BAILLEY 78 CNTR + STORAGE RING  
 EDM B BAILLEY 78 YIELDS EDM < 1.05\*10\*\*18 WITH CL=.95. THIRD RESULT IS  
 EDM B FIRST TWO COMBINED ASSUMING CPT.

## $\mu/p$ MAGNETIC MOMENT RATIO

MPR THIS RATIO IS USED TO OBTAIN A PRECISE VALUE OF THE MUON MASS.  
 MPR (3.1865) (0.0022) COFFIN 58 CNTR + SPIN RESONANCE  
 MPR (3.1850) (0.0011) LUNDY 58 CNTR + PRECESSION STROB  
 MPR (3.176) (0.013) LUNDY 58 CNTR + PRECESSION STROB  
 MPR (3.1834) (0.0002) GARWIN 60 CNTR + PRECESSION PHASE  
 MPR (3.18336) (0.00007) BINGHAM 63 CNTR + PRECESSION STROB  
 MPR (3.1808) (0.0004) BINGHAM 63 CNTR + PRECESSION STROB  
 MPR (3.18338) (0.00004) HUTCHINS 63 CNTR + PRECESSION PHASE  
 MPR C (3.183351) (0.000016) EHRLLICH 69 CNTR HFS SPLITTING  
 MPR C (3.183314) (0.000034) THOMPSON 69 CNTR HFS SPLITTING  
 MPR C (3.183330) (0.000046) HUTCHINS 70 CNTR + PRECESSION PHASE  
 MPR H (3.183347) (0.00009) HAGUE 70 CNTR + PRECESSION PHASE  
 MPR C (3.183336) (0.000013) CRANE 71 CNTR HFS SPLITTING  
 MPR D (3.183349) (0.000015) DEVOE 71 CNTR HFS SPLITTING  
 MPR F (3.183326) (0.000013) FAVART 71 CNTR HFS SPLITTING  
 MPR H (3.183347) (0.000082) CROWE 72 CNTR + PRECESSION PHASE  
 MPR R THE RESULTS THROUGH 1972 ARE INCLUDED IN COHEN 73.  
 MPR R 3.1833402 .0000072 COHEN 73 RVUE  
 MPR E (3.1833299 .0000025) CASPERSON 75 CNTR  
 MPR C (3.1833403 .0000044) CASPERSON 77 CNTR + HFS SPLITTING  
 MPR (3.1833448 .0000029) CAMANI 78 CNTR + REPL. BY KLEMP82  
 MPR 3.1833441 .0000017 KLEMP 82 CNTR + PRECESSION STROB  
 MPR 3.1833461 .0000011 MARIAM 82 CNTR + HFS SPLITTING  
 MPR D DEVOE 71 SUPERCEDES EHRLLICH 69. THIS IS NOT A DIRECT MEASUREMENT.  
 MPR D WE GIVE A NEW VALUE WHICH CONTAINS A THEORETICAL CORRECTION OF  
 MPR D -7.8+-2.3 PPM, AS DISCUSSED IN FOOTNOTE 35A OF CROWE 72.  
 MPR C CRANE 71 SUPERCEDES THOMPSON 69. THIS IS NOT A DIRECT MEASUREMENT.  
 MPR H CROWE 72 SUPERCEDES HAGUE 70.  
 MPR F FAVART 71 ASSUMES A ZERO VALUE FOR THE PROTON POLARIZABILITY.  
 MPR E USES INCORRECT THEO. EXPRESSION FOR NU(HFS). SEE KLEMP 82, TBL.X1.  
 MPR AVG 3.1833452 .0000010 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)

## $\mu$ PARTIAL DECAY MODES

DECAY MASSES

P1  $\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$  .511+ 0+ 0  
 P2  $\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu \gamma$  .511+ 0+ 0+ 0  
 P3  $\mu^- \rightarrow e^- \nu_e \bar{\nu}_\mu$  .511+ 0+ 0  
 P4  $\mu^- \rightarrow e^- \gamma$  .511+ 0+ 0  
 P5  $\mu^- \rightarrow 3e^-$  .511+.511+.511  
 P6  $\mu^- \rightarrow e^- 2\gamma$  .511+ 0+ 0  
 P7  $\mu^- \rightarrow e^- e^+ e^- \bar{\nu}_e \nu_\mu$  .511+.511+.511+ 0+

For notation, see key on page 91.

# Stable Particle Full Listings

$\mu$

## $\mu$ BRANCHING RATIOS

$\mu^- \rightarrow (e^- \bar{\nu}_e \nu_\mu \gamma)/\text{total}$  (P2)

R1 27 EVENTS SEEN ASHKIN 59 CNTR  
R1 1.4E-2 0.4E-2 CRITTENDE 61 CNTR T(GAM) GT 10 MEV  
R1 (3.3E-3) (1.3E-3) CRITTENDE 61 CNTR T(GAM) GT 20 MEV  
R1 862 EVENTS SEEN BDGART 67 CNTR T(GAM) GT 14.5 MEV

$\mu^+ \rightarrow (e^+ \bar{\nu}_e \nu_\mu)/\text{total}$  (P3)

R2 FORBIDDEN BY ADDITIVE CONSERVATION LAW FOR LEPTON FAMILY NUMBER.  
R2 MULTIPLICATIVE LAW PREDICTS THIS BRANCHING RATIO TO BE 1/2.  
R2 FOR A REVIEW SEE NEMETHY 81.  
R2 (0.25) OR LESS CL=.90 EICHEN 73 HLBC +  
R2 (0.13) (0.15) BLIETSCHA 78 HLBC +- AVG. OF 4 VALUES  
R2 (0.09) OR LESS CL=.90 JONKER 80 CALD REPL. BY BERGSM 83  
R2 (-0.001) (0.061) WILLIS 80 CNTR +  
R2 0.05 OR LESS CL=.90 BERGSM 83 CALO ANUMU E --> MU- ANUE  
R2 A BERGSM 83 GIVES LIMIT ON INVERSE MUON DECAY CROSS SECTION RATIO  
R2 A SIG(NUBAR(MU) E --> MU- NUBAR(E)) / SIG(NUMU) E --> MU- NU(E)),  
R2 A WHICH IS ESSENTIALLY EQUIVALENT TO R2 FOR SMALL VALUES LIKE THAT  
R2 A QUOTED.

$\mu^- \rightarrow (e \gamma)/\text{total (units } 10^{-8})$  (P4)

R3 FORBIDDEN BY LEPTON FAMILY NUMBER CONSERVATION  
R3 (4.5) OR LESS CL=.90 FRANKEL 1 63 OSPK  
R3 (2.9) OR LESS CL=.90 PARKER 64 OSPK  
R3 (2.9) OR LESS CL=.90 KORENCH 71 OSPK + DUBNA  
R3 (0.10) OR LESS CL=.90 SCHAAF 80 ELEC + SIN  
R3 0.017 OR LESS CL=.90 KUMJONG 82 SPEC + LAMPF  
R3 (0.10) OR LESS CL=.90 AZUELOS 83 CNTR +

$\mu^- \rightarrow (3e)/\text{total (units } 10^{-10})$  (P5)

R4 FORBIDDEN BY LEPTON FAMILY NUMBER CONSERVATION  
R4 F (5000.) OR LESS CL=.90 PARKER 62 CNTR  
R4 F (1300.) OR LESS CL=.90 ALIKHANOV 62 OSPK  
R4 F (1500.) OR LESS CL=.90 FRANKEL 2 63 CNTR  
R4 F (1200.) OR LESS CL=.90 BABAEV 63 OSPK  
R4 K (62.) OR LESS CL=.90 KORENCH 71 OSPK DUBNA  
R4 K (19.) OR LESS CL=.90 KORENCHEN 76 SPEC + DUBNA  
R4 K (1.4) OR LESS CL=.90 BERTL 84 SPEC + SINDRUM  
R4 K (1.3) OR LESS CL=.90 BOLTON 84 CNTR LANL  
R4 K 0.024 OR LESS CL=.90 BERTL 85 SPEC SINDRUM  
R4 F FOUR ABOVE EXPERIMENTS EVALUATED UPPER LIMITS ASSUMING A SECOND  
R4 F ORDER V-A NEUTRINO LOOP DIAGRAM. LIMITS NOT SIGNIFICANTLY CHANGED  
R4 F BY ASSUMING A CONSTANT MATRIX ELEMENT.  
R4 K THESE EXPERIMENTS ASSUME A CONSTANT MATRIX ELEMENT.

$\mu^- \rightarrow (e 2\gamma)/\text{total (units } 10^{-7})$  (P6)

R5 FORBIDDEN BY LEPTON FAMILY NUMBER CONSERVATION  
R5 (150.) OR LESS CL=.90 FRANKEL 1 63 OSPK +  
R5 P (40.) OR LESS CL=.90 POUTISSOU 74 CNTR + LBL  
R5 A (0.5) OR LESS CL=.90 BOWMAN 78 CNTR DEPOMMIER 77 DATA  
R5 B 0.084 OR LESS CL=.90 AZUELOS 83 CNTR +  
R5 P POUTISSOU 74 LIMIT APPLIES TO SUM OF ALL NEUTRINOLESS MU+ DECAYS.  
R5 A BOWMAN 78 ASSUMES INT. LAGRANG. LOCAL ON SCALE OF INVERSE MU MASS.  
R5 B AZUELOS 83 USES PHASE SPACE DISTRIBUTION OF BOWMAN 78. SEE ABOVE.

$\mu^- \rightarrow (e^- e^+ e^- \bar{\nu}_e \nu_\mu)/\text{total (units } 10^{-5})$  (P7)

R6 L 3 (1.5) (1.0) LEE 59 HBC +  
R6 G 1 (2.) GUREVICH 60 EMUL +  
R6 C 7 (2.2) (1.5) CRITTENDE 61 HLBC + E(E+ E-) > 10MEV  
R6 B7443 5.4 0.4 BERTL 85 SPEC + SINDRUM  
R6 L IN THE THREE LEE 59 EVENTS, THE SUM OF ENERGIES S=(E+)+(E-)+E(E+)  
R6 L WAS S=51 MEV, 35 MEV, AND 35 MEV.  
R6 G GUREVICH 60 INTERPRET THEIR EVENT AS EITHER VIRTUAL OR REAL PHOTON  
R6 G CONVERSION. E+ AND E- ENERGIES NOT MEASURED.  
R6 C CRITTENDEN 61 COUNT ONLY THOSE DECAYS WHERE TOTAL ENERGY OF EITHER  
R6 C (E+, E-) COMBINATION IS > 10 MEV  
R6 B BERTL 85 HAS TRANSVERSE MOMENTUM CUT PT> 17 MEV/C.  
R6 B STAT ERROR (0.2) AND SYST ERROR (INCREASED BY US) ADDED IN QUADR.

## LIMIT ON $\mu^- \rightarrow e^-$ CONVERSION

FORBIDDEN BY LEPTON FAMILY NUMBER CONSERVATION

$\sigma(\mu^- 32S \rightarrow e^- 32S)/\sigma(\mu^- 32S \rightarrow \nu_\mu 32P^*)$   
RE (4. E-10) OR LESS CL=.90 BADERTSCH 77 STRC SIN  
RE 0.7E-10 OR LESS CL=.90 BADERTSCH 80 STRC SIN

$\sigma(\mu^- Cu \rightarrow e^- Cu)/\sigma(\mu^- Cu \rightarrow \text{capture})$   
RF 1.6E-8 OR LESS CL=.90 BRYMAN 72 SPEC

$\sigma(\mu^- Ti \rightarrow e^- Ti)/\sigma(\mu^- Ti \rightarrow \text{capture})$   
RG 1.6E-11 OR LESS CL=.90 BRYMAN 85 SPEC TRIUMF

## LIMIT ON $\mu^- \rightarrow e^+$ CONVERSION

FORBIDDEN BY TOTAL LEPTON NUMBER CONSERVATION

$\sigma(\mu^- 32S \rightarrow e^+ 32S^*)/\sigma(\mu^- 32S \rightarrow \nu_\mu 32P^*)$   
RP1 (1.5E-9) OR LESS CL=.90 BADERTSCH 78 STRC SIN  
RP1 0.9E-9 OR LESS CL=.90 BADERTSCH 80 STRC SIN

$\sigma(\mu^- 127I \rightarrow e^+ 127I^*)/\sigma(\mu^- 127I \rightarrow \text{anything})$   
RP2 A 0.3E-9 OR LESS CL=.90 ABELA 80 CNTR RADIOCHEMICAL TECH.  
RP2 A ABELA 80 IS UPPER LIMIT FOR MU- E+ CONVERSION LEADING TO PARTICLE-  
RP2 A STABLE STATES OF SB127. LIMIT FOR TOTAL CONVERSION RATE IS HIGHER  
RP2 A BY A FACTOR LESS THAN 4 (G. BACKENSTOSS, PRIVATE COMM.)

$\sigma(\mu^- Cu \rightarrow e^+ Co)/\sigma(\mu^- Cu \rightarrow \nu_\mu Ni)$   
RP3 (2.2E-7) OR LESS CL=.90 CONFORTO 62 OSPK  
RP3 2.6E-8 OR LESS CL=.90 BRYMAN 72 SPEC

## LIMIT ON $(\mu^+, e^-)$ BOUND STATE CONVERSION TO $(\mu^-, e^+)$

FORBIDDEN BY LEPTON FAMILY NUMBER CONSERVATION

$R_g = g_C/g_F$

MC WHERE  $g_F = 1.16637E-5 \text{ GEV}^{*-2}$  IS THE FERMI CONSTANT AND  
MC  $g_C$  IS AN EFFECTIVE COUPLING (DIMENSIONS  $\text{GEV}^{*-2}$ ) FOR A  
MC FOUR-FERMION INTERACTION ASSUMED TO BE RESPONSIBLE FOR THE  
MC CONVERSION OF THE  $(\mu^+, e^-)$  BOUND STATE TO  $(\mu^-, e^+)$ .  
MC  
MC 42 OR LESS CL=.95 MARSHALL 82 CNTR

## NOTE ON MUON DECAY PARAMETERS

(by F. Scheck, University of Mainz, West Germany and K. Mursula, Nordita, Copenhagen, Denmark)

The muon decay parameters describe the momentum spectrum ( $\rho$  and  $\eta$ ), the asymmetry ( $\xi$  and  $\delta$ ), and the polarization of the electron ( $\xi', \xi'', \alpha, \beta, \alpha', \beta'$ ) in the process  $\mu^- \rightarrow e^- + \nu_\mu + \bar{\nu}_e$ . Assuming a local, lepton-number-conserving, derivative-free, four-fermion interaction, the matrix element in charge-changing form may be written as<sup>1</sup>

$$\frac{G_F}{\sqrt{2}} \left\{ \sum_1 h_{ik} \langle \bar{e} | 1 + (-)^i \gamma_5 | \nu_e \rangle \right. \\ \times \langle \bar{\nu}_\mu | 1 + (-)^k \gamma_5 | \mu \rangle \\ + \sum_1 g_{ik} \langle \bar{e} | \gamma^\alpha [1 + (-)^i \gamma_5] | \nu_e \rangle \\ \times \langle \bar{\nu}_\mu | \gamma_\alpha [1 + (-)^k \gamma_5] | \mu \rangle \\ + \sum_1 f_{ii} \frac{1}{2} \langle \bar{e} | \sigma^{\alpha\beta} [1 + (-)^i \gamma_5] | \nu_e \rangle \\ \left. \times \langle \bar{\nu}_\mu | \sigma_{\alpha\beta} [1 + (-)^j \gamma_5] | \mu \rangle \right\} \quad (1)$$

The definitions of covariants and the sign conventions are the ones of Sachs and Sirlin (1975)<sup>2</sup> and Scheck (1978).<sup>3</sup> The connection to other charge-changing and charge-retention forms is worked out in Mursula and Scheck (1985).<sup>4</sup> Note that for massless particles the covariants chosen above project onto states of definite helicity. In the standard model,  $g_{22} = 1$  and all other coupling constants vanish.

All electron observables can be expressed, in a model-independent way, as functions of the ten standard real constants (see references above)  $a, b, c, a', b', c', \alpha, \beta, \alpha',$  and  $\beta'$ . The rate is proportional to  $A = a + 4b + 6c$ . The above decay parameters depend on nine of these constants only:

## Stable Particle Full Listings

$\mu$

$$\rho - \frac{3}{4} = \frac{3}{4}[-a + 2c]/A,$$

$$\eta = [\alpha - 2\beta]/A,$$

$$\delta - \frac{3}{4} = \frac{9}{4} \frac{[a' - 2c']/A}{1 - [a + 3a' + 4(b + b') + 6c - 14c']/A},$$

$$1 - \xi \frac{\delta}{\rho} = 4 \frac{[(b + b') + 2(c - c')]/A}{1 - [a - 2c]/A}, \quad (2)$$

$$1 - \xi' = [a + a' + 4(b + b') + 6(c + c')]/A,$$

$$1 - \xi'' = [-2a + 20c]/A,$$

$$\alpha/A, \beta/A, \alpha'/A, \beta'/A,$$

the last four of which are obtained from the transverse components of the electron polarization. These real constants are easily related to bilinear combinations of coupling constants in any form of the interaction. For the case of the form (1), they are given by [note the scale factor  $G_F/\sqrt{2}$  in Eq. (1)],

$$\left. \begin{array}{l} a \\ a' \end{array} \right\} = 16(|g_{12}|^2 \pm |g_{21}|^2) \pm |h_{11} + 6f_{11}|^2 \\ + |h_{22} + 6f_{22}|^2, \quad (3)$$

$$\left. \begin{array}{l} b \\ b' \end{array} \right\} = 4(|g_{11}|^2 \pm |g_{22}|^2) \pm |h_{12}|^2 + |h_{21}|^2, \quad (4)$$

$$\left. \begin{array}{l} c \\ c' \end{array} \right\} = \frac{1}{2} \left[ \pm |h_{11} - 2f_{11}|^2 + |h_{22} - 2f_{22}|^2 \right], \quad (5)$$

$$\left. \begin{array}{l} \alpha \\ \alpha' \end{array} \right\} = \left. \begin{array}{l} \text{Re} \\ \text{Im} \end{array} \right\} 8[g_{21}(h_{22}^* + 6f_{22}^*) \pm g_{12}(h_{11}^* + 6f_{11}^*)], \quad (6)$$

$$\left. \begin{array}{l} \beta \\ \beta' \end{array} \right\} = \left. \begin{array}{l} \text{Re} \\ \text{Im} \end{array} \right\} (-4)[g_{22}h_{21}^* \pm g_{11}h_{12}^*]. \quad (7)$$

As the decay parameters (2) depend on  $(b + b')$  but not on  $(b - b')$ , the constant  $h_{12}$  appears only in  $\beta$  and  $\beta'$ , Eq. (7). However,  $g_{11}$  is found to be compatible with zero, so the decay parameters do not determine  $h_{12}$ . (The coupling constant  $h_{12}$  does occur in the rate parameter  $A$  and may, in principle, be obtained by comparing  $\mu$  decay to other data.) By using Eqs. (2) and the experimental determinations of  $\rho$ ,  $\eta$ ,  $\xi\delta/\rho$ ,  $\delta$ ,  $\xi'$ ,  $\xi''$ ,  $\alpha$ ,  $\beta$ ,  $\alpha'$ , and  $\beta'$ , limits can be placed on the nine parameters  $a/A$ ,  $a'/A$ ,  $(b + b')/A$ ,  $c/A$ ,  $c'/A$ ,  $\alpha/A$ ,  $\alpha'/A$ ,  $\beta/A$ , and  $\beta'/A$ . These are given in the Listings. These limits are easily translated into limits on specific coupling constants in Eq. (1), depending on what kind of extension of the standard model one wishes to test. Examples

such as tests for right-handed interactions or for scalar/pseudoscalar effective couplings are given in Mursala and Scheck (1985).<sup>4</sup>

The limits on  $a$ ,  $a'$ ,  $\dots$ , can be recast into limits on the effective charge-retention coordinates  $g_S/g_V$ ,  $g_P/g_V$ ,  $g_A/g_V$ ,  $g_T/g_V$ ,  $\phi_{VA}$ , and  $\psi_{VA}$ , which were used in the earlier literature (cf., e.g., DERENZO 69). The most recent values are found in BURKARD2 85.

Note that the radiative corrections are unambiguous only if  $h_{11}$ ,  $h_{22}$ ,  $g_{12}$ ,  $g_{21}$ ,  $f_{11}$ , and  $f_{22}$  vanish.

### References

1. F. Scheck, in *Leptons, Hadrons, and Nuclei* (North Holland, Amsterdam, 1983).
2. A.M. Sachs and A. Sirlin, in *Muon Physics II*, eds. C.S. Wu and V. Hughes (Academic Press, New York, 1975), p. 49.
3. F. Scheck, *Phys. Rep.* **44**, 187 (1978).
4. K. Mursala and F. Scheck, *Nucl. Phys.* **B253**, 189 (1985); and K. Mursala et al., *Nucl. Phys.* **B219**, 321 (1983).

### $\mu$ DECAY PARAMETERS

$\rho$ PARAMETER	(V-A theory predicts $\rho=0.75$ )
RHO C (0.741) (0.027)	DUZIAK 59 CNTR + 20-53 MEV E+
RHO P9215 (0.745) (0.025)	PLANO 60 HBC + WHOLE SPECTRUM
RHO P TWO PARAMETER FIT TO RHO AND ETA	
RHO C 2276 (0.751) (0.034)	BLOCK 62 HBC - WHOLE SPECTRUM
RHO D (0.64) (0.04)	BARLOW 64 CNTR - WHOLE SPECTRUM
RHO D (0.661) (0.016)	BARLOW 64 CNTR + WHOLE SPECTRUM
RHO D (0.867) (0.035)	PONTECORV 64 CC -
RHO D RESULTS IN DOUBT.	
RHO C 800K (0.7503) (0.0026)	PEOPLES 66 ASPK + 20-53 MEV E-
RHO C 280K (0.760) (0.009)	SHERWOOD 67 ASPK + 25-53 MEV E-
RHO C 170K (0.762) (0.008)	FRYBERGER 68 ASPK + 25-53 MEV E-
RHO C ETA CONSTRAINED =0. THESE VALUES INCORPORATED INTO A TWO PARAMETER FIT TO RHO AND ETA BY DERENZO 69.	
RHO C 0.7518 0.0026	DERENZO 69 RVUE
RHO AVG 0.7517 0.0026	AVERAGE
$\eta$ PARAMETER	(V-A theory predicts $\eta=0$ )
ETA P 9213 (-2.0) (0.9)	PLANO 60 HBC - WHOLE SPECTRUM
ETA P TWO PARAMETER FIT TO RHO AND ETA	ETA-PLANO 60 DISCOUNTS VALUE FOR ETA
ETA C 800K (0.05) (0.5)	PEOPLES 66 ASPK + 20-53 MEV E+
ETA C 280K (-0.7) (0.6)	SHERWOOD 67 ASPK + 25-53 MEV E+
ETA C 170K (-0.7) (0.5)	FRYBERGER 68 ASPK + 25-53 MEV E+
ETA C RHO CONSTRAINED =0.75.	
ETA 6346 -0.12 0.21	DERENZO 69 HBC - 1.6-6.8 MEV E+
ETA B5.3M 0.011 0.085	BURKARD2 85 CNTR + 9-53 MEV E+
ETA D95.3M (-0.012) (0.016)	BURKARD2 85 CNTR + 9-53 MEV E-
ETA B STATISTICAL (0.081) AND SYSTEMATIC ERRORS ADDED IN QUADRATURE.	
ETA B MEASURED VALUE OF ETA IS AN ENERGY AVERAGE.	
ETA D ALPHA-PRIME=0 ASSUMED.	
ETA AVG -0.007 0.079	AVERAGE
$\xi''$ PARAMETER	
XPP 8326K 0.65 0.36	BURKARD1 85 CNTR + BHABHA + ANNIHIL
XPP B BURKARD1 85 MEASURE (X1''-X1*X1')/X1 AND X1'	
XPP B AND SET X1-1.	
( $\xi$ PARAMETER)*( $\mu$ LONGITUDINAL POLARIZATION)	
X1 (0.9959)OR MORE CL=90 CARR 83 SPEC + 11 KGAUSS	(V-A THEORY PREDICTS X1=1, LONG.POL=1)
X1 9K (0.97) (0.05)	BARDON 59 CNTR BROMFORM TARGET
X1 8354 (0.93) (0.06)	PLANO 60 HBC + 8.8 KGAUSS
X1 A (0.903) (0.027)	ALI-ZADE 61 EMUL + 27 KGAUSS
X1 A DEPOLARIZATION BY MEDIUM NOT KNOWN SUFFICIENTLY WELL.	
X1 66K (0.975) (0.015)	CUREVICH 64 EMUL REPL BY AKHMANOV 68
X1 (0.975) (0.015)	AKHMANOV 68 EMUL 140 KGAUSS
( $\xi$ PARAMETER)*( $\mu$ LONGITUDINAL POLARIZATION)* $\delta/\rho$	
XID S (0.9959)OR MORE CL=90 CARR 83 SPEC + 11 KGAUSS	
XID S 0.9966 OR MORE CL=90 STOKER 85 SPEC + MU-SPIN ROT	
XID S STOKER 85 FIND (X3*PMU*DELTA/RHO)=0.9959 AND >0.9966, WHERE FIRST	
XID S LIMIT IS FROM NEW MU-SPIN ROTATION DATA AND SECOND IS FROM	
XID S COMBINATION WITH CARR 83 DATA. (DELTA/RHO)=1.0 IN V-A THEORY.	
$\delta$ PARAMETER	(V-A theory predicts $\delta=0.75$ )
DEL 8354 0.78 0.05	PLANO 60 HBC + WHOLE SPECTRUM
DEL 0.782 0.031	KRUGER 61
DEL 490K 0.752 0.009	FRYBERGER 68 ASPK + 25-53 MEV E-
DEL VOSSLER 69 HAS MEASURED THE ASYMMETRY BELOW 10 MEV	
DEL AVG 0.7551 0.0085	AVERAGE



Stable Particle Full Listings

$\mu, \nu_\tau, \tau$

CUMBLEY 74 PRPL 14 1 F.CUMBLEY,E.PICASSO (CERN)
CALMET 77 RMP 49 21 J.CALMET,S.NARISON,M.PERROTET. (MARS)
KINOSHITA 78 TOKYO HEP P,571 T. KINOSHITA (COORN)

$\tau$  MEAN LIFE (units 10<sup>-13</sup> sec)

Table with columns for tau decay mode, branching ratio, and mean life. Includes entries like P1: tau+ -> mu+ nu\_mu nu\_tau, P2: tau+ -> e+ nu\_e nu\_tau, etc.

$\nu_\tau$

J = 1/2

EXISTENCE INDIRECTLY ESTABLISHED FROM TAU DECAY DATA COMBINED WITH NU REACTION DATA. SEE FOR EXAMPLE FELDMAN 81. KIRKSY 79 RULE OUT J=3/2 USING TAU -> PI NU(TAU) BRANCHING RATIO.

$\tau$  PARTIAL DECAY MODES

Table of partial decay modes for tau, labeled P1 to P30. Includes modes like P1: tau+ -> mu+ nu\_mu nu\_tau, P2: tau+ -> e+ nu\_e nu\_tau, P3: tau+ -> 3 pi+ nu\_tau, etc.

NOT IN GENERAL A MASS EIGENSTATE. SEE NOTE ON NEUTRINOS IN THE ELECTRON NEUTRINO SECTION ABOVE.

$\nu_\tau$  "MASS" (MeV)

APPLIES TO NU(3), THE PRIMARY MASS EIGENSTATE IN NU(TAU). WOULD ALSO APPLY TO ANY OTHER NU(i) WHICH MIXES STRONGLY IN NU(TAU) AND HAS SUFFICIENTLY SMALL MASS THAT IT CAN OCCUR IN THE RESPECTIVE DECAYS. (THIS WOULD BE NONTRIVIAL ONLY FOR A HYPOTHETICAL J.GE. 4, GIVEN THE NU(i) AND NU(j) "MASS" LIMITS ABOVE.)

Table listing various tau decay channels and their corresponding mass eigenstates. Includes channels like BA 594 (250.) OR LESS CL= .95 BACINO, M 3 (250.) OR LESS CL= .95 BLOCKER, etc.

$\nu_\tau$  BOUNDS ON MASS AND MIXING

LIMITS ON |U\_e tau|^2 AS FUNCTION OF nu\_tau MASS. Table with columns for UET, CL, BERGSM, BOMP, and M(NU(TAU)) values.

LEPTON FAMILY NUMBER VIOLATING MODES.

Table of lepton family number violating decay modes, labeled P31 to P44. Includes modes like P31: tau+ -> mu+ gamma, P32: tau+ -> e+ gamma, etc.

REFERENCES FOR nu\_tau

Table listing references for tau neutrino, including LNC 25 404, LNC 25 404, W.ALLES (BGMNJ), etc.

$\tau$

J = 1/2

TAU DISCOVERY PAPER IS PERL 75. E+ E- -> TAU-TAU- CROSS SECTION THRESHOLD BEHAVIOR AND MAGNITUDE CONSISTENT WITH POINTLIKE SPIN 1/2 DIRAC PARTICLE. BRANDELIX 78 RULES OUT POINTLIKE SPIN 0 OR SPIN 1 PARTICLE.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P\_i, as follows: The diagonal elements are P\_i + delta P\_i, where delta P\_i = sqrt((delta P\_i)^2 + (delta P\_j)^2). For the definitions of the individual P\_i see the listings above; only those P\_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table of tau mass (MeV) measurements and fits. Includes entries like M A 692 1783., M R 299 1787., M 1807., etc.

Matrix of fitted partial decay mode branching fractions. Rows and columns labeled P 1, P 2, P 7, P 12, P 16, P 17, P 18, P 19, P 20.

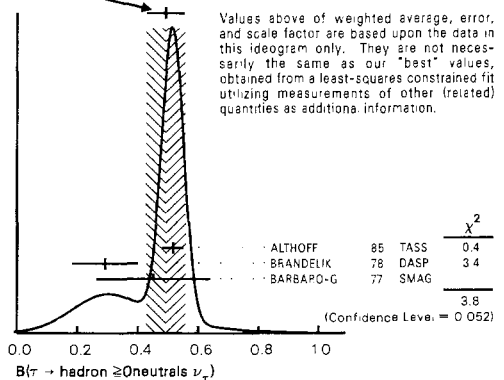
For notation, see key on page 91.

Stable Particle Full Listings

Table listing various tau branching ratios (BR) and statistical parameters. Columns include particle name, BR value, error, and a list of contributing experiments (e.g., BURMEST1, PERL, CAVALLISS, SMITH, BRANDELIK, BERGER, BEHREND, ALTHOFF, ASH, BALTRUSA, BERGER).

Table listing hadron and neutral tau branching ratios. Columns include particle name, BR value, error, and contributing experiments (e.g., WAGNER, BLOCKER, BEHREND, WAGNER, BLOCKER).

WEIGHTED AVERAGE 0.492 ± 0.065 (ERROR SCALED BY 1.9)





# Stable Particle Full Listings

T

$\tau^\pm \rightarrow (K^\pm (\geq 0 \text{ neutrals}) \nu_\tau)/\text{total}$  (P11)

R26 B SMALL BRANDELIX 77 DASP 3.6-5.2ECM E+E-  
R26 53 (0.0121) 0.0029 JAROS 78 SMAG E+E- ECM=29 GEV  
R26 B BRANDELIX 77 FINDS 0.07+-0.06 K+- PER EVT IN E+E- --> E+- PRONG+-.

$\tau^\pm \rightarrow (2\pi^\pm \pi^\mp \nu_\tau)/\text{total}$  (P17)

R28 DECAY MODES WITH KAONS ARE MEASURED TO BE SMALL, SO ALL HADRONS  
R28 ARE ASSUMED TO BE PIONS. BEHREND 84 SUBTRACTS KAONS BY HAND.  
R28 J 15 (0.07) (0.05) JAROS 78 SMAG E+E- ECM > 6 GEV  
R28 0.09 0.06 BRANDELIX 80 TASS E+E- ECM=30GEV  
R28 0.097 0.024 BRANDELIX 84 CELL E+E- ECM=14,22 GEV  
R28 1255 0.081 0.008 FERNADE 85 MAC E+E- ECM=29 GEV  
R28 J JAROS 78 EVENTS CONSISTENT WITH BEING RHO PI OR A1.  
R28 R28 AVG 0.0827 0.0075 AVERAGE  
R28 FIT 0.0810 0.0073 FROM FIT

$\tau^\pm \rightarrow (2\pi^\pm \pi^\mp (\text{non-resonant}) \nu_\tau)/\text{total}$  (P13-P9)

R29 0.014 OR LESS CL=.95 WAGNER 80 PLUT E+E- 4-5 GEV ECM

$\tau^\pm \rightarrow (3\pi^\pm \pi^0 \nu_\tau)/\text{total}$  (P25)

R30 WE FIT THIS AS P18 (TAU+- INTO 3 HADRONS+- GAMMA(S) NU(TAU)) AND  
R30 ASSUME THAT THE MULTI-PI0 FRACTION IS SMALL. (P17+P18)  
R30 (0.11) (0.07) JAROS 78 SMAG E+E- ECM > 6 GEV  
R30 (0.35) (0.11) BRANDELIX 78 DASP ASSUMES V-A DECAY  
R30 0.062 0.029 BEHREND 84 CELL E+E- ECM=14,22 GEV  
R30 R30 AVG 0.075 0.031 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)  
R30 FIT 0.0527 0.0078 FROM FIT

$\tau^\pm \rightarrow (3\pi^\pm (\geq 0 \pi^0) \nu_\tau)/\text{total}$  (P17+P18)

R31 DECAY MODES WITH KAONS ARE MEASURED TO BE SMALL, SO ALL HADRONS  
R31 ARE ASSUMED TO BE PIONS.  
R31 E 692 (0.32) (0.05) BACINO 78 DLCO E+E- ECM=3.1-7.4GEV  
R31 E 33 (0.18) (0.065) JAROS 78 SMAG E+E- ECM > 6 GEV  
R31 35 0.24 0.06 BRANDELIX 80 TASS E+E- ECM=30 GEV  
R31 186 0.150 0.020 BEHREND 82 CELL E+E- ECM=32-36.8GEV ECM  
R31 152 (0.14) 0.02 BLOCKERS 82 SMK2 E+E- ECM = 29 GEV  
R31 A 660 0.148 0.017 ATHARA 84 TPC E+E- ECM=29 GEV  
R31 C 0.098 0.130 0.004 AKERLOF 85 HRS E+E- ECM=29 GEV  
R31 F 367 0.153 0.017 0.019 ALTHOFF 85 TASS E+E- ECM=34.5 GEV  
R31 J 0.136 0.009 BARTEL 85 JADE E+E- ECM=34.6 GEV  
R31 G 0.122 0.041 BERGER 85 PLUT E+E- ECM=34.6 GEV  
R31 B 0.133 0.007 FERNADE 85 MAC E+E- ECM=29 GEV  
R31 E LOW ENERGY EXPERIMENTS ARE NOT IN AVERAGE OR FIT BECAUSE THE  
R31 E SYSTEMATIC ERRORS IN BACKGROUND SUBTRACTION ARE JUDGED TO BE LARGE.  
R31 A STATISTICAL AND SYSTEMATIC ERRORS ARE 0.009 AND 0.015.  
R31 C STATISTICAL AND SYSTEMATIC ERRORS ARE 0.002 AND 0.005.  
R31 F STATISTICAL AND SYSTEMATIC ERRORS ARE 0.011 AND 0.013, 0.016.  
R31 J STATISTICAL AND SYSTEMATIC ERRORS ARE 0.005 AND 0.008.  
R31 G NOT INDEPENDENT OF BERGER 85 R1, R2, R42, AND R43.  
R31 G STATISTICAL AND SYSTEMATIC ERRORS ARE 0.013 AND 0.039.  
R31 B STATISTICAL AND SYSTEMATIC ERRORS ARE 0.003 AND 0.007.  
R31 R31 AVG 0.1334 0.0030 AVERAGE  
R31 FIT 0.1337 0.0030 FROM FIT

$\tau^\pm \rightarrow (5\pi^\pm (\geq 0 \pi^0) \nu_\tau)/\text{total}$  (P20)

R33 10 (0.06) OR LESS CL=.95 BRANDELIX 80 TASS E+E- ECM=30GEV  
R33 2 (0.005) OR LESS CL=.95 BEHREND 82 CELL E+E- ECM = 29 GEV  
R33 4 (0.003) OR LESS CL=.90 ATHARA 84 TPC E+E- ECM=29 GEV  
R33 1 (0.009) OR LESS CL=.95 BEHREND 84 CELL E+E- ECM=14,22 GEV  
R33 0 (0.007) OR LESS CL=.95 ALTHOFF 85 TASS E+E- ECM=34.5 GEV  
R33 J 0.003 0.002 BARTEL 85 JADE E+E- ECM=34.6 GEV  
R33 10 0.0013 0.0004 BELTRAMI 85 HRS E+E- ECM=29 GEV  
R33 C 4 0.0016 0.0009 BURCHAT 85 SMK2 E+E- ECM=29 GEV  
R33 2 (0.0017) OR LESS CL=.95 FERNADE 85 MAC E+E- ECM=29 GEV  
R33 J STATISTICAL AND SYSTEMATIC ERRORS ARE 0.001 AND 0.002.  
R33 C STATISTICAL AND SYSTEMATIC ERRORS ARE 0.0008 AND 0.0004.  
R33 R33 AVG 0.00147 0.00036 AVERAGE  
R33 FIT 0.00147 0.00036 FROM FIT

$\tau^\pm \rightarrow (K^\pm (2 \text{ charged}) (\geq 0 \text{ neutrals}) \nu_\tau)/\text{total}$  (P28)

R34 0.006 OR LESS CL=.90 ATHARA 84 TPC E+E- ECM=29 GEV

$\tau^\pm \rightarrow (3 \text{ hadrons}^\pm \nu_\tau)/\text{total}$  (P17)(P17-P18)

R35 (NOT INDEPENDENT OF R28 AND R31 VALUES.)  
R35 103 0.37 0.35 0.20 ALTHOFF 85 TASS E+E- ECM=34.5 GEV  
R35 B 0.61 0.06 FERNADE 85 MAC E+E- ECM=29 GEV  
R35 B STATISTICAL AND SYSTEMATIC ERRORS ARE 0.03 AND 0.05.  
R35 R35 AVG 0.599 0.059 AVERAGE

$\tau^\pm \rightarrow (3\pi^\pm K^0 (\geq 0 \gamma) \nu_\tau)/\text{total}$  (P29)

R37 0.0027 OR LESS CL=.90 BELTRAMI 85 HRS E+E- ECM=29 GEV

$\tau^\pm \rightarrow (5\pi^\pm \nu_\tau)/\text{total}$  (P3)

R38 B 5 0.00067 0.00030 BELTRAMI 85 HRS E+E- ECM=29 GEV  
R38 B THE ERROR QUOTED IS STATISTICAL ONLY.

$\tau^\pm \rightarrow (5\pi^\pm \pi^0 \nu_\tau)/\text{total}$  (P4)

R39 B 5 0.00067 0.00030 BELTRAMI 85 HRS E+E- ECM=29 GEV  
R39 B THE ERROR QUOTED IS STATISTICAL ONLY.

$\tau^\pm \rightarrow (K^\pm K^\mp \pi^\pm \nu_\tau)/\text{total}$  (P5)

R40 M 9 0.0022 0.0017 0.0011 MILLS 85 DLCO E+E- ECM=29 GEV  
R40 M ERROR CORRELATED WITH MILLS 85 (K PI PI0 NU) VALUE. EXCLUDES  
R40 M 23% SYSTEMATIC ERROR.

$\tau^\pm \rightarrow (K^\pm \pi^\pm \pi^\mp (\geq 0 \pi^0) \nu_\tau)/\text{total}$  (P6)

R41 M 9 0.0022 0.0016 0.0013 MILLS 85 DLCO E+E- ECM=29 GEV  
R41 M ERROR CORRELATED WITH MILLS 85 (K K PI PI0 NU) VALUE. EXCLUDES 23%  
R41 M SYSTEMATIC ERROR.

$\tau^\pm \rightarrow (\text{hadron}^\pm (\geq 1 \pi^0) \nu_\tau)/\text{total}$  (P30)

R42 G 0.427 0.035 BERGER 85 PLUT E+E- ECM=34.6 GEV  
R42 G STATISTICAL AND SYSTEMATIC ERRORS ARE 0.020 AND 0.029.  
R42 R42 FIT 0.408 0.013 FROM FIT

$\tau^\pm \rightarrow (\text{hadron}^\pm \nu_\tau)/\text{total}$  (P12+P7)

R43 G 0.130 0.045 BERGER 85 PLUT E+E- ECM=34.6 GEV  
R43 G STATISTICAL AND SYSTEMATIC ERRORS ARE 0.020 AND 0.040.  
R43 FIT 0.108 0.011 FROM FIT

$\tau^\pm \rightarrow (\mu^\pm \text{ charged particles} + e^\pm \text{ charged particles})/\text{total}$  (P33+P34)

R52 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R52 B 0.04 OR LESS CL=.90 BURMESTZ 77 PLUT E+E- 4-5 GEV ECM  
R52 B ASSUMES SAME MU, E MOM. SPEC. AS (MU E + NOTHING DETECTED).

$\tau^\pm \rightarrow (\mu^\pm \gamma)/\text{total}$  (P31)

R53 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R53 5.5E-4 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (e^\pm \gamma)/\text{total}$  (P32)

R54 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R54 6.4E-4 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (\mu^\pm \mu^\pm \mu^\mp)/\text{total}$  (P35)

R55 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R55 4.9E-4 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (e^\pm \mu^\pm \mu^\mp)/\text{total}$  (P36)

R56 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R56 3.3E-4 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (\mu^\pm e^\pm e^\mp)/\text{total}$  (P37)

R57 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R57 4.4E-4 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (e^\pm e^\pm e^\mp)/\text{total}$  (P38)

R58 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R58 4.0E-4 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (\mu^\pm \pi^0)/\text{total}$  (P39)

R59 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R59 8.2E-4 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (e^\pm \pi^0)/\text{total}$  (P40)

R60 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R60 2.1E-3 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (\mu^\pm K^0)/\text{total}$  (P41)

R61 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R61 1.0E-3 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (e^\pm K^0)/\text{total}$  (P42)

R62 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R62 1.3E-3 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (\mu^\pm \rho^0)/\text{total}$  (P43)

R63 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R63 4.4E-4 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

$\tau^\pm \rightarrow (e^\pm \rho^0)/\text{total}$  (P44)

R64 TEST OF LEPTON FAMILY NUMBER CONSERVATION  
R64 3.7E-4 OR LESS CL=.90 HAYES 82 SMK2 E+E- 3.8-6.8GEV ECM

## $\tau$ DECAY PARAMETERS

$\rho$  (MICHEL) PARAMETER (V-A theory predicts  $\rho=0.75$ )

RHO 594 0.72 0.15 BACINO2 79 DLCO E+E- ECM=3.5-7.4GEV  
RHO A1426 0.71 0.095 BEHREND 85 CLEO E+E- NEAR UPS(4S)  
RHO A STAT. ERROR (0.09) AND SYST. ERROR (0.03) ADDED IN QUADRATURE.  
RHO A USED 699 (727) EVENTS WITH AN ELECTRON (MUON).  
RHO RHO AVG 0.713 0.080 AVERAGE

## $\tau/\mu$ CHARGED COUPLING-CONSTANT RATIO

CC C 0.94 0.15 0.13 ALTHOFF 84 TASS E+E- ECM=43 GEV  
CC C STATISTICAL AND SYSTEMATIC ERRORS ARE +0.12 -0.09 AND 0.09.

## REFERENCES FOR $\tau$

PERL 75 PRL 35 1489 -ABRAMS, BOYARSKI, BREIDENBACH + (LBL,SLAC)  
PERL 76 PL 63B 466 +FELDMAN, ABRAMS, ALAM, BOYARSKI + (SLAC-LBL)

BARBARO- 77 PRL 39 1058 BARBARO-GALLIERI + (LBL-NWES-SLAC-HAWA)  
BRANDELI 77 PL 70B 125 BRANDELIX + (AACH+DESY+HAMB+MPI+TOKY)  
BURMESTZ 77 PL 68B 297 BURMESTER, CRIGEE + (DESY+HAMB+SIEG+WUPG)  
BURMESTZ 77 PL 68B 301 BURMESTER, CRIGEE + (DESY+HAMB+SIEG+WUPG)  
CAVALLI- 77 LNG 20 337 CAVALLI-SFORZA, GOGGI (PAVI-PRIN-LUM)  
ALEXAND 77 PL 70B 487 =FELDMAN, ABRAMS, ALAM, BOYARSKI, (SLAC-LBL)

ALEXAND 78 PL 78B 162 ALEXANDER + (DESY+AACH+HAMB+SIEG+WUPG)  
BACINO 78 PRL 41 13 +FERGUSON, NODULMAN + (UCLA-SLAC+UCI+STON) J  
ALSO 78 TOKYO CONF. P.249 J. KIRZ (19TH INTL. CONF. ON HEP) (STON)  
ALSO 80 PL 96B 214 ZHOLENTZ, KURADAZE, LELCHUK, MISHNEV + (NOVO)  
BARTEL 78 PL 77B 331 +DITTMANN, DITNKER, OLSSON, ONEILL + (DESY+HEF)

BRANDELI 78 PL 73B 109 BRANDELIX + (AACH+DESY+HAMB+ MPI+TOKY) J  
HEILE 78 NP B138 189 =PERL, ABRAMS, ALAM, BOYARSKI + (SLAC+LBL)  
JAROS 78 PRL 40 1120 +ABRAMS, ALAM + (SLAC+LBL-NWES-HAWA)  
SMITH 78 PR D18 1 +FORD, MORSE, MANN, RESVANIS + (COLO.PENN-WISC)

ALLES 79 LNC 25 404 M. ALLES  
BACINO1 79 PRL 42 6 +FERGUSON, NODULMAN + (UCLA+SLAC+UCI+STON) J  
BACINO2 79 PRL 42 749 +FERGUSON, NODULMAN + (UCLA+SLAC+UCI+STON)

For notation, see key on page 91.

Stable Particle Full Listings

SEARCHES FOR MASSIVE NU'S & LEPTON MIXING

Table listing various experiments and their results, including BLOCKER, BRANDOLI, WAGNER, ZHOLENTZ, BERGER, DORFAN, BEHREND, JAROS, AIHARA, ALTHOFF, BEHREND, MILLS, AKERLOF, ALTHOFF, ASH, BALTRUSA, BARTEL, BEHREND, BELTRAMI, BERGER, BURCHART, FERNANDEZ, MILLS, AZIMOV, FELDMAN, FLUGGE, KIRKBY, PERL, etc.

SEARCHES FOR MASSIVE NEUTRINOS AND LEPTON MIXING

OMITTED FROM SUMMARY TABLE

SEE THE NOTE ON NEUTRINOS BY R.E. SHROCK IN THE ELECTRON NEUTRINO SECTION NEAR THE BEGINNING OF THESE DATA LISTINGS.

SEARCHES FOR INDIRECT EFFECTS OF NEUTRINO MASSES AND LEPTON MIXING ARE LISTED HERE. DIRECT SEARCHES FOR MASSES OF DOMINANTLY COUPLED NEUTRINOS ARE LISTED IN THE APPROPRIATE SECTION ON NU(MU), NU(E), OR NU(TAU). RESULTS OF THESE INDIRECT SEARCHES ARE CORRELATED UPPER BOUNDS ON MIXING MATRIX COEFFICIENTS (U, A, J) VERSUS NEUTRINO MASS. THESE RESULTS ARE DIVIDED INTO THREE SECTIONS--

- (A) BOUNDS FROM PARTICLE AND NUCLEAR DECAYS
(B) BOUNDS FROM NEUTRINO REACTIONS
(C) SEARCHES FOR NEUTRINOLESS DOUBLE BETA DECAY

THE SITUATION CAN BE SUMMARIZED AS FOLLOWS. CURRENT EXPERIMENTS YIELD NULL RESULTS, I.E. ONLY UPPER BOUNDS, WITH TWO EXCEPTIONS:

- BOTH OF WHICH ARE IN STRONG CONFLICT WITH SUBSEQUENT EXPERIMENTS:
(1) CAVAIGNAC 84, REPORT EVIDENCE FOR NEUTRINO OSCILLATIONS OF THE TYPE NUBAR(E) -> NUBAR(E), FROM REACTOR NEUTRINO EXPERIMENT AT BUGEY. HOWEVER, ZAJEC 85 RULES OUT ALMOST ALL OF THE ALLOWED REGION IN SIN^2(2\*THETA)\*\*2 AND DELTA(M\*\*2) CLAIMED BY CAVAIGNAC 84.
(2) SIMPSON 85 REPORTS OBSERVATION OF ADMIXED HEAVY NEUTRINO IN TRITIUM BETA DECAY WITH M(NU(J))=17.1 KEV, CABSCU(1, J)\*\*2=0.05+-0.01. HOWEVER, THE CLAIM OF SIMPSON 85 IS STRONGLY REFUTED BY ALTZITZGOLU 85, OHI 85, AND APALIKOV 85.

SEE ALSO BOUNDS ON MASSES OF NEUTRINOS ABOVE.

(A). BOUNDS FROM PARTICLE AND NUCLEAR DECAYS

LIMITS ON |Uij|^2 AS FUNCTION OF MASS(vj)

Table with columns for Uij, Application of kink and peak search test to existing data, and M(NU(J)) values.

SEARCHES FOR DECAYS OF MASSIVE NU

Table listing searches for decays of massive neutrinos by BERGSMAS, SHROCK, GRONAU, COOPER-SARKER.

Uij B BERGSMAS 83 ALSO QUOTE LIMITS ON CABSCU(1, J)\*\*2 WHERE THE INDEX 3 REFERS TO THE MASS EIGENSTATE DOMINANTLY COUPLED TO THE TAU. THOSE LIMITS WERE BASED ON ASSUMPTIONS ABOUT THE D/S MASS AND D/S --> B TAU NUTAU BRANCHING RATIO WHICH ARE NO LONGER VALID. SEE COOPER-SARKER 85.
Uij C COOPER-SARKER 85 ALSO GIVE LIMITS BASED ON MODEL-DEPENDENT ASSUMPTIONS FOR NUTAU FLUX. WE DO NOT LIST THESE.
Uij C NOTE THAT FOR THIS BOUND TO BE NON-TRIVIAL, J IS NOT EQUAL TO 3, I.E. NU(J) CANNOT BE THE DOMINANT MASS EIGENSTATE IN NUTAU SINCE M(NU(3))>70 MEV (ALBRECHT 85). ALSO, OF COURSE, J IS NOT EQUAL TO 1 OR 2, SO A FOURTH GENERATION WOULD BE REQUIRED FOR THIS BOUND TO BE NONTRIVIAL.

KINK SEARCH IN NUCLEAR BETA DECAY

Table listing kink search results in nuclear beta decay for various experiments like SCHRECKEN, SIMPSON, ALTZITZGOLU, etc.

DATA FROM SULFUR-35 BETA DECAY.

Uij S DATA FROM TRITIUM BETA DECAY CONTRADICTED BY ALTZITZGOLU 85. Uij S APALIKOV 85, MARKKY 85, AND OHI 85. ALSO COMMENT IN HARTON 85 Uij S AND KALBEFLEISCH 85. Uij B THIS LIMIT WAS TAKEN FROM THE FIG. 3 OF APALIKOV 85; THE TEXT GIVES A MORE RESTRICTIVE LIMIT OF 1.7E-3 AT CL=90.

LIMITS ON |Uij|^2 AS FUNCTION OF MASS(vj)

Table listing application of peak search test to existing data for various experiments like SHROCK, SHROCK 81, etc.

NEW EXPERIMENTS TO APPLY PEAK SEARCH TEST.

Table listing new experiments to apply peak search test, including ABEA, CALAPRICE, ASANO, HAYANO, MINEHART, etc.

PEAK SEARCH IN MUON-CAPTURE

Table listing peak search in muon-capture limits on |Uij|^2 as function of mass(NU(J)) by DEUTSCH 83, etc.

SEARCHES FOR DECAYS OF MASSIVE NU

Table listing searches for decays of massive neutrinos by COOPER-SARKER, BERGSMAS, COOPER-SARKER 85, etc.

LIMITS ON |Uij|^2

Table listing limits on |Uij|^2 where A=1,2,3 FROM RHO PARAM. IN MU DECAY by SHROCK 81, etc.

LIMITS ON |Uij\*Uij|^2 AS FUNCTION OF MASS(vj)

Table listing limits on |Uij\*Uij|^2 as function of mass(NU(J)) by BERGSMAS, SHROCK 81, etc.

(B). BOUNDS FROM NU REACTIONS

SOLAR NU EXPERIMENTS

SOLAR NU FLUX (units, SNU)

Table listing solar neutrino flux measurements by BERGSMAS, BANCALL, MEAS, etc.

SEE ALSO THE REVIEW BY BANCALL 82 AND THE ANALYSIS BY EHRICH 82.

Stable Particle Full Listings  
SEARCHES FOR MASSIVE ν's & LEPTON MIXING

DEEP MINE EXPERIMENTS

R= (MEASURED FLUX OF νμ)/(EXPECTED FLUX OF νμ)
DU 0.62 0.17 CROUCH 78 CASE WEST/UCI
DU 0.95 0.22 BOLIEV 81 BAKSAN
DU LIMITS ON DELTA(M\*\*2) FOR SIN(2\*THETA)=1 (E\*\*2)
DU L EXCLUDE (2.2-11.2)E-5 CL=.90 LOSECCO 85 CNTR 1MB, FLUX-INDEP.

REACTOR ν EXPERIMENTS

EVENTS (OBSERVED/EXPECTED) FROM REACTOR ν EXPERIMENTS.
RD B (0.89) (0.15) BOEHM 80 ANU(E) P -> E- N
RD A R 0.38 0.21 REINES 80 SEE NOTE A
RD A R 0.40 0.22 REINES 80 SEE NOTE A

νe + ν̄e

Δ(m²) FOR sin²(2θ)=1 (units eV²)
RD1 1.3E-1 OR LESS BELENKII 83 ANU(E) P -> E- N
RD1 G 1.6E-2 OR LESS CL=.90 GABATHULE 84 ANU(E) P -> E- N
RD1 7. E-2 OR LESS CL=.90 AFONIN 85 ANU(E) P -> E- N

Δ(m²) FOR GIVEN sin²(2θ) (units eV²)
RD2 C 0.2 0.1 CAVAINA 84 ANU(E) P -> E- N
RD2 C SIN(2\*THETA)\*\*2=0.25--0.1.
RD2 C THESE ARE FROM BEST FIT TO DATA; SEE CAVAINAC 84 FOR PLOT OF
RD2 C ALLOWED REGIONS IN THESE VARIABLES. THESE DATA FROM BUGEY REACTOR.

sin²(2θ) FOR "LARGE" Δ(m²)
RS1 H 0.4 OR LESS BELENKII 83 ANU(E) P -> E- N
RS1 G 0.16 OR LESS CL=.90 GABATHULE 84 ANU(E) P -> E- N
RS1 0.34 OR LESS CL=.90 AFONIN 85 ANU(E) P -> E- N
RS1 J 0.19 OR LESS CL=.90 ZACEK 85 ANU(E) P -> E- N

ACCELERATOR EXPERIMENTS

BOUNDS ON Δ(m²) VS. sin²(2θ)
WHERE DELTA(M\*\*2) IS MAGNITUDE OF (MASS(NU(I))\*\*2 - MASS(NU(J))\*\*2)
AND THETA IS THE MIXING ANGLE FOR THE SIMPLIFYING ASSUMPTION OF
MIXING BETWEEN TWO NEUTRINO FAMILIES ONLY.

EACH EXPERIMENTAL RESULT IS A PLOT GIVING ALLOWED AND EXCLUDED
REGIONS AS FUNCTIONS OF DELTA(M\*\*2) AND SIN(2\*THETA)\*\*2. WE QUOTE
TWO REPRESENTATIVE LIMITS FROM EACH PLOT --
1) DELTA(M\*\*2) FOR SIN(2\*THETA)\*\*2=1.
2) SIN(2\*THETA)\*\*2 FOR "LARGE" DELTA(M\*\*2), I.E. SUFFICIENTLY
LARGE DELTA(M\*\*2) THAT THE DETECTOR WOULD MEASURE ONLY AN
EFFECT AVERAGED OVER MANY OSCILLATIONS.

νμ + ν̄μ

Δ(m²) FOR sin²(2θ)=1 (units eV²)
D1 1.2 OR LESS CL=.95 BELLOTTI 76 HLBC GGM CERN PS
D1 1.2 OR LESS CL=.95 BLIETSCHA 78 HLBC GGM CERN PS
D1 1.7 OR LESS CL=.90 ARMENISE 81 HLBC GGM CERN PS

sin²(2θ) FOR "LARGE" Δ(m²)
S1 1. E-2 OR LESS CL=.95 BELLOTTI 76 HLBC GGM CERN PS
S1 4. E-3 OR LESS CL=.95 BLIETSCHA 78 HLBC GGM CERN PS
S1 1. E-2 OR LESS CL=.90 ARMENISE 81 HLBC GGM CERN PS

ν̄μ → ν̄e

Δ(m²) FOR sin²(2θ)=1 (units eV²)
D2 1. OR LESS CL=.95 BLIETSCHA 78 HLBC GGM CERN PS
D2 0.91 OR LESS CL=.90 NEMETHY 81 CNTR LAMPF
D2 2.4 OR LESS CL=.90 TAYLOR 83 HLBC 15 FT FNAL

sin²(2θ) FOR "LARGE" Δ(m²)
S2 4. E-3 OR LESS CL=.95 BLIETSCHA 78 HLBC GGM CERN PS
S2 0.2 OR LESS CL=.90 NEMETHY 81 CNTR LAMPF
S2 1.3E-2 OR LESS CL=.90 TAYLOR 83 HLBC 15 FT FNAL

νμ → ντ

Δ(m²) FOR sin²(2θ)=1 (units eV²)
D3 4.6 OR LESS CL=.90 ARMENISE 81 HLBC GGM CERN SPS
D3 3. OR LESS CL=.90 BAKER 81 HLBC 15FT FNAL
D3 6. OR LESS CL=.90 ERRIQUEZ 81 HLBC BEBC CERN SPS

sin²(2θ) FOR "LARGE" Δ(m²)
S3 1.7E-2 OR LESS CL=.90 ARMENISE 81 HLBC GGM CERN SPS
S3 6. E-2 OR LESS CL=.90 BAKER 81 HLBC BEBC CERN SPS
S3 5. E-2 OR LESS CL=.90 ERRIQUEZ 81 HLBC BEBC CERN SPS

ν̄μ → ν̄τ

Δ(m²) FOR sin²(2θ)=1 (units eV²)
D4 2.2 OR LESS CL=.90 ASRATYAN 81 HLBC FNAL
D4 7.4 OR LESS CL=.90 TAYLOR 83 HLBC 15 FT FNAL

sin²(2θ) FOR "LARGE" Δ(m²)
S4 4.4E-2 OR LESS CL=.90 ASRATYAN 81 HLBC FNAL
S4 8.8E-2 OR LESS CL=.90 TAYLOR 83 HLBC 15 FT FNAL

νμ + ν̄μ

Δ(m²) FOR sin²(2θ)=1 (units eV²)
D5 THESE EXPERIMENTS ALSO ALLOW SUFFICIENTLY LARGE DELTA(M\*\*2)
D5 8.0 OR LESS CL=.90 BELIKOV 83 CNTR OR > 1000.
D5 0.29 OR LESS CL=.90 BERGSM 84 CNTR OR > 22.
D5 0.23 OR LESS CL=.90 DYDAK 84 CNTR OR > 100.
D5 10. OR LESS CL=.90 STOCKDALE 84 CNTR OR > 1250.
D5 7. OR LESS CL=.90 BELIKOV 85 CNTR SERPUKHOV
D5 8.0 OR LESS CL=.90 STOCKDALE 85 CNTR OR > 1250.

sin²(2θ) AS FUNCTION OF Δ(m²)
S5 A 0.1 OR LESS CL=.90 BELIKOV 83 CNTR SERPUKHOV
S5 B 0.27 OR LESS CL=.90 BERGSM 84 CNTR CHARM CERN PS
S5 C 0.1 OR LESS CL=.90 DYDAK 84 CNTR CERN PS
S5 D 0.02 OR LESS CL=.90 STOCKDALE 84 CNTR FNAL
S5 E 0.07 OR LESS CL=.90 BELIKOV 85 CNTR SERPUKHOV
S5 F 0.02 OR LESS CL=.90 STOCKDALE 85 CNTR FNAL
S5 A BOUND HOLDS FOR DELTA(M\*\*2)=20--1000 EV\*\*2
S5 B THIS BOUND APPLIES FOR DELTA(M\*\*2)=0.7-9. EV\*\*2. LESS STRINGENT
S5 B BOUNDS APPLY FOR OTHER DELTA(M\*\*2); THESE ARE NONTRIVIAL FOR
S5 B 0.28 < DELTA(M\*\*2) < 22 EV\*\*2.
S5 C THIS BOUND APPLIES FOR DELTA(M\*\*2)=1.-10. EV\*\*2. LESS STRINGENT
S5 C BOUNDS APPLY FOR OTHER DELTA(M\*\*2); THESE ARE NONTRIVIAL FOR
S5 C 0.23 < DELTA(M\*\*2) < 90 EV\*\*2.
S5 D THIS BOUND APPLIES FOR DELTA(M\*\*2)=110 EV\*\*2. LESS STRINGENT
S5 D BOUNDS APPLY FOR OTHER DELTA(M\*\*2); THESE ARE NONTRIVIAL FOR
S5 D 10 < DELTA(M\*\*2) < 1250 EV\*\*2.
S5 E THIS BOUND APPLIES FOR A WIDE RANGE OF DELTA(M\*\*2) > 7 EV\*\*2. FOR
S5 E SOME VALUES OF DELTA(M\*\*2), THE VALUE IS LESS STRINGENT; THE LEAST
S5 E RESTRICTIVE, NONTRIVIAL BOUND OCCURS APPROXIMATELY AT DELTA(M\*\*2)=
S5 E 300 EV\*\*2 WHERE SIN(2\*THETA)\*\*2 < 0.13 AT CL=.90.
S5 F THIS BOUND APPLIES FOR DELTA(M\*\*2)=100 EV\*\*2. LESS STRINGENT
S5 F BOUNDS APPLY FOR OTHER DELTA(M\*\*2); THESE ARE NONTRIVIAL FOR
S5 F 8 < DELTA(M\*\*2) < 1250 EV\*\*2.

νe + ν̄e

Δ(m²) FOR sin²(2θ)=1 (units eV²)
D6 8. OR LESS CL=.90 BAKER 81 HLBC 15FT FNAL
D6 56. OR LESS CL=.90 DENEN 81 HLBC BEBC CERN SPS
D6 10. OR LESS CL=.90 ERRIQUEZ 81 HLBC BEBC CERN SPS
D6 2.3 TO 8 EXCLUDED CL=.90 NEMETHY 81 CNTR LAMPF

sin²(2θ) FOR "LARGE" Δ(m²)
S6 0.6 OR LESS CL=.90 BAKER 81 HLBC 15FT FNAL
S6 0.3 OR LESS CL=.90 DENEN 81 HLBC BEBC CERN SPS
S6 7.E-2 OR LESS CL=.90 ERRIQUEZ 81 HLBC BEBC CERN SPS

ν̄e → ν̄τ

sin²(2θ) FOR "LARGE" Δ(m²)
S7 F 0.7 OR LESS CL=.90 FRITZE 80 HYBR BEBC CERN SPS
S7 F AUTHORS GIVE P(NU(E)-->NU(TAU))>.35, EQUIVALENT TO ABOVE LIMIT.

ν̄μ + ν̄μ

Δ(m²) FOR sin²(2θ)=1 (units eV²)
D8 7. OR LESS CL=.90 STOCKDALE 85 CNTR OR > 1200

sin²(2θ) AS FUNCTION OF Δ(m²)
S8 G 0.02 OR LESS CL=.90 STOCKDALE 85 CNTR FNAL
S8 G THIS BOUND APPLIES FOR DELTA(M\*\*2) BETWEEN 190 AND 320 OR=530 EV\*\*2
S8 G LESS STRINGENT BOUNDS APPLY FOR OTHER DELTA(M\*\*2); THESE ARE NON-
S8 G TRIVIAL FOR 7 < DELTA(M\*\*2) < 1200 EV\*\*2.

For notation, see key on page 91.

SEARCHES FOR MASSIVE ν's & LEPTON MIXING, LIMITS ON NUMBER OF LIGHT ν TYPES

Stable Particle Full Listings

νμ → (ν̄e)L

THIS IS A LIMIT ON LEPTON FAMILY-NUMBER VIOLATION AND TOTAL LEPTON-NUMBER VIOLATION. (NUBAR(E))-L DENOTES A HYPOTHETICAL LEFT-HANDED ANTI-ELECTRON NEUTRINO. THE SOUND IS QUOTED IN TERMS OF DELTA (Δm²), SIN(2θ), AND ALPHA, WHERE ALPHA DENOTES THE FRACTIONAL ADMIXTURE OF (ν+A) CHARGED CURRENT.

α Δ(m²) FOR sin²(2θ)=1 (units eV²)

AD1 C 7. E-1 OR LESS CL-.90 COOPER 82 HLCB
AD1 C EXISTING BOUNDS ON ν+A CURRENTS REQUIRE ALPHA SMALL - SEE COOPER82.

α² sin²(2θ) FOR "LARGE" Δ(m²)

AS1 C 1. E-3 OR LESS CL-.90 COOPER 82 HLCB
AS1 C EXISTING BOUNDS ON ν+A CURRENTS REQUIRE ALPHA SMALL - SEE COOPER82.

νe → (ν̄e)L

SEE NOTE FOR NU(MU) --> (NUBAR(E))-L LIMIT

α Δ(m²) FOR sin²(2θ)=1 (units eV²)

AD2 C 7. OR LESS CL-.90 COOPER 82 HLCB
AD2 C EXISTING BOUNDS ON ν-A CURRENTS REQUIRE ALPHA SMALL - SEE COOPER82.

α² sin²(2θ) FOR "LARGE" Δ(m²)

AS2 C 5. E-2 OR LESS CL-.90 COOPER 82 HLCB
AS2 C EXISTING BOUNDS ON ν-A CURRENTS REQUIRE ALPHA SMALL - SEE COOPER82.

(C). SEARCHES FOR NEUTRINOLESS DOUBLE β DECAY

THE DECAY (Z,A) --> (Z-2,A) + e- + e-, I.E. NEUTRINOLESS DOUBLE BETA DECAY, ISOLATES TOTAL LEPTON NUMBER BY TWO UNITS. IT IS FORBIDDEN IF NEUTRINOS ARE DIRAC PARTICLES BUT CAN OCCUR IF NEUTRINOS ARE MAJORANA PARTICLES AND (A) THEY ARE MASSIVE OR (B) THEY HAVE NON (V-A) COUPLINGS. PRIMAKOFF 81, ROSEN 81, AND HAXTON 83 DISCUSS CORRELATED BOUNDS ON MNM AND RIGHT-HANDED COUPLINGS. FURTHER THEORETICAL DISCUSSIONS INCLUDE DOI 81, HAXTON 82, AND GROTZ 83.

MW MNM, THE EFFECTIVE WEIGHTED SUM OF NEUTRINO MASSES CONTRIBUTING TO NEUTRINOLESS DOUBLE BETA DECAY (UNITS EV)
MW MNM = CABS(SUM FROM I TO N OF UC(I,J)\*\*2\*(NUC(I,J))), WHERE N=NUMBER OF NEUTRINO GENERATIONS, AND NU(I,J) IS A MAJORANA NEUTRINO. NOTE THAT UC(I,J)\*\*2, NOT CABS(CU(I,J))\*\*2, OCCURS IN THE SUM; THE POSSIBILITY OF CANCELLATIONS HAS BEEN STRESSED IN WOLFENSTEIN 81.
MW 10 OR LESS CL-.68 AVIGNONE 83 CNTR +THY GE-76
MW 22 OR LESS CL-.90 BELLOTTI 83 CNTR +THY1 GE-76
MW 8.3 OR LESS CL-.68 BELLOTTI 83 CNTR +THY2 GE-76
MW 3.6 OR LESS CL-.95 KIRSTEN 83 SPEC +THY TE-128/TE-130
MW 5.8 OR LESS CL-.68 BELLOTTI 84 CNTR +THY GE-76
MW 22. OR LESS CL-.68 FORSTER 84 CNTR +THY GE-76
MW 6. OR LESS CL-.68 CALDWELL 85 CNTR +THY GE-76
MW N LIMITS ARE OBTAINED FROM ANALYSIS OF DATA USING THEORETICAL B CALCULATIONS BY DOI-83 (=THY1) AND ROSEN 81 (=THY2).
MW E SEE TABLE 1 OF BELLOTTI 84 FOR THEIR ASSESSMENT OF PREVIOUS BOUNDS.
MW E HALF LIFE FOR 0- -> 0+ TRANSITION >> E22 YR (CL-.90), >1.2E23 YR (CL-.68)
MW C USES RESULTS OF HAXTON 81,82. AUTHORS STATE THAT BOUND COULD BE "TWO OR THREE TIMES LARGER." HALF LIFE FOR 0- -> 0+ TRANSITION > 5E22 YR (CL-.68)
MW N SEE HUBERT 85 FOR LIFETIME LIMITS ON NEUTRINOLESS DOUBLE BETA DECAY
MW N OF GE-76 TO EXCITED STATES OF GE-76.

ETA LIMITS ON LEPTON-NUMBER VIOLATING (ν+A) CURRENT ADMIXTURE
ETA ETA IS DEFINED AS THE FRACTIONAL ADMIXTURE OF (ν+A)
ETA CHARGED CURRENT, RELATIVE TO (ν-A) IN ELECTRON-TYPE LEPTON SECTOR.
ETA 2.4E-5 OR LESS AVIGNONE 83 CNTR +THY GE-76
ETA 4. E-5 OR LESS CL-.68 BELLOTTI 83 CNTR +THY1 GE-76
ETA 1.5E-5 OR LESS CL-.68 BELLOTTI 83 CNTR +THY2 GE-76
ETA 2.4E-5 OR LESS CL-.95 KIRSTEN 83 SPEC +THY TE-128/TE-130
ETA C 0.8E-5 OR LESS CL-.68 BELLOTTI 84 CNTR +THY GE-76
ETA C 0.6E-5 OR LESS CL-.68 BELLOTTI 84 CNTR +THY GE-76
ETA C 1.4E-5 OR LESS CL-.68 CALDWELL 85 CNTR +THY GE-76
ETA C 0.9E-5 OR LESS CL-.68 CALDWELL 85 CNTR +THY GE-76
ETA B LIMITS ARE OBTAINED FROM ANALYSIS OF DATA USING THEORETICAL B CALCULATIONS BY DOI-83 (=THY1) AND ROSEN 81 (=THY2).
ETA C TWO BOUNDS GIVEN, DEPENDING ON TYPES OF CHIRALITY MIXING. SEE REFS. 1

REFERENCES FOR BOUNDS ON ν MASS, MIXING

BELLOTTI 76 LNC 17 553 +CAVALLI,FIORINI,ROLLIER (MILA)
BLIETSCH 78 NP B133 205 BLIETSCHAU-(AACH+LH3+CERN+EPOL+MILA+ORSA)+
CROUCH 78 PR D18 2239 +LANDECKER,LATHROP,REINES- (CASE+UCI+MIT)

BALTAY 81 NU 81 CONF HAWAII C-BALTAY (COLU)
BOLIV 81 SUNP 34 787 +BUTKEVICH,ZAKIDYSHYEV,MAKOEV,MIKHEEV-(CNR)
CALAPRICE 81 PL 1068 175 CALAPRICE,SCHREIBER,SCHNEIDER (PRIN+IND)
DAVIS 81 BNL NU WKSP B.T.CLEVELAND,R.DAVIS JR.,J.K.ROWLEY (BNL)
DEFEN 81 PL 988 310 +(AACH+BONN+CERN+ATEN+LOIC+OXF+SACL)
DOI 81 PL 1039 219 +KOTANI,NISHIURA,OKUDA,TAKASUGI (OSAK)
DOI 81 PTP 66 1739 DOI,KOTANI,NISHIURA,OKUDA,TAKASUGI (OSAK)
DOI 81 PTP 66 1765 DOI,KOTANI,NISHIURA,OKUDA,TAKASUGI (OSAK)
ERRIQUEZ 81 PL 1029 73 +NATALI-(BARI+BIEM+ISH+EPOL+RHEL+SACL)
HAXTON 81 PRL 47 153 +G.J.STEPHENSON,D.STROTTMAN (PURD,LASL)
KWON 81 PR D24 1097 +BOEHM,HANN,HENRIKSON+ (CIT+GREN-NFUK)
NEMETHY 81 PR D23 262 +(YALE+LBL+LASL+MIT+SACL+SIN+CNRC+BERN)
PRIMAKOFF 81 ARNS 31 145 H.PRIMAKOFF,S.P.ROSEN (CIT)
ROSEN 81 NU 81 CONF HAWAII S.P.ROSEN (PURD)
SHROCK 81 PR D24 1232 D.BRYMAN,C.PICCITOTTO (TRIU,VICT)
SHROCK 81 PR D24 1275 R.E.SHROCK (STON)
SILVERMAN 81 PRL 46 467 D.SILVERMAN,A.SONE (UCI+UCLA)
SIMPSON 81 PR D25 2360 J.SIMPSON (STON)
USHIDA 81 PRL 47 1694 (AICH-FNAL+KOE+SEOU-MGCI+NAGO+OSU+OKAY+)
WOLFENST 81 PL 1078 77 L.WOLFENSTEIN (CARN)
BACWAL 82 RMP 54 767 +HUEBNER,LUBOW+ (IAS-LANL+HPC+YALE+UCLA)
COOPER 82 PL 1128 97 +GUY,MICHETTE,TYNDEL,VENUS (RL)
EHRlich 82 PR D25 2282 R.EHRlich (GNLS)
FILLIPPO 82 APJ 253 393 B.W.FILIPPONE,D.N.SCHRAMM (ANL+EFI)
HAXTON 82 PR D25 2360 +STEPHENSON,D.STROTTMAN (LANL+PURD)
HAYANO 82 PRL 49 1305 +TANIGUCHI,YAMANAKA+ (TOKY+KEK-TSUK)
VUILLEUM 82 PL 1148 298 VUILLEUMIER,BOEHM,EGGER+ (CIT-SIN-MUNI)
AVIGNONE 83 PRL 50 721 +BRODZINSKI,BROWN,EVANS,HENSLY+ (SCUC-PURD)
BELENKII 83 JETPL 58 493 +DOBRYNIN,ZEMLYAKOV,MIKHELYAN- (KIEP)
BELLOTTI 83 JETPL 58 661 +VOLKOV,KOCHETKOV,MUKHIN,SVIRIDOV- (SER)
BELLOTTI 83 PL 1218 72 +FIORINI,LIGUORI,PULLIA,SARRACINO- (MILA)
BERGMA 83 PL 1288 361 CHARM COLLAB. (ANIK+CERN+HAMB+ITEP+INFN)
BRYMAN 83 PRL 50 1546 +DUBOIS,NUMAO,OLANIYI,OLIN- (TRIU-CNRC)
DEUTSCH 83 PR D27 1644 BRYMAN,DUBOIS,NUMAO,OLANIYI,OLIN- (TRIU-CNRC)
DOI 83 PTP 69 602 J.P.DEUTSCH,M.LEBRUN,R.PRIEELS (LVN)
GRONAU 83 PR D28 2762 KOTANI,NISHIURA,TAKASUGI (OSAK+KYOT)
GROTZ 83 JPG 9 1169 GRONAU (HAIF)
HAXTON 83 CNP 11 41 W.C.HAXTON (LASL+PURD)
KIRSTEN 83 PRL 50 474 +KIRSTEN,H.RICHTER,E.JESSBERGER (MPIH)
KIRSTEN 83 ZPHY 16 189 T.KIRSTEN,H.RICHTER,E.JESSBERGER (MPIH)
SCHRECKE 83 PL 1298 265 SCHRECKENBACH,COLVIN+ (GREN+ILLG)
TAYLOR 83 PR D28 2705 +CENCE,HARRIS,JONES+ (HAWA+LBL+FNAL)
BALLAGH 84 PR D30 2271 +BINGHAM,+ (UCB+LBL+FNAL+HAMA+WASH+MISC)
BELLOTTI 84 PL 1468 450 +CREMONESI,FIORINI,LIGUORI,PULLIA+ (MILA)
BERGMA 84 PL 1428 103 CHAR C. (AMT+CERN+HAMB+ITEP+INFN)
CAVAIGNAC 84 PL 1488 387 CAVAIGNAC,HOUMMADA,KOANG- (LLG-LAPP)
DYDAK 84 PL 1348 281 +FELDMAN,GUYOT-, (CERN+DORT+HEID+SACL+WARS)
FORSTER 84 PL 1388 301 +KWON,MARKY,BOEHM,HENRIKSON (CIT)
GABATHULER 84 PL 1388 449 GABATHULER,BOEHM- (CIT-SIN-MUNI)
MINEHART 84 PRL 50 804 +ZIOCK,MARSHALL,STEPHENS,DAUM- (VIRG+STON)
STOCKDAL 84 PRL 52 1384 STOCKDALE,BODEK+, (ROCH+CHIC+COLU+FNAL)
AFONIN 85 JETPL 41 435 AFONIN,BOROVOI,DOBRYNIN+ (KIEP)
ALPHESS 85 PR D31 2732 AFONIN,BOGATOV,BOROVOI,DOBRYNIN+ (KIEP)
ALPHESS 85 PR D31 2732 +ARNOSON-, (BNL+BROOK+KEK+OSAK+PENN+STON)
ALPHESS 85 PRL 55 799 ALPHESS,ALPHESS,ALPHESS,DEWEY+, (PRIN)
APALIKOV 85 JETPL 42 289 +BORIS,GOLTYVIN,LAPTEIN,LUBIMOV+ (ITEP)
BELLIKOV 85 SUNP 41 569 +VOLKOV,KOCHETKOV,MUKHIN+ (SER)
BLUMER 85 PL 1618 407 BLUMER,KLEJNKNECHT (DORT)
CALDWELL 85 PRL 54 281 +EISEBERG,GRUMM,HALE,WITHERELL+ (UCSB+LBL)
COOPER-S 85 PL 1608 207 COOPER-SARKAR,HAYWOOD+CERN-LOIC-DFX+SACL)
HAXTON 85 PRL 55 807 W.HAXTON (LASL)
HUBERT 85 NC 85A 19 +LECCIA,DASSIE,MENNRATH- (BORD-ZARA)
KALBFLEI 85 PRL 55 2225 KALBFLEI,SCHMIDT,MILTON (OKLA)
LOSECCO 85 PRL 54 2299 +BIONTA+BLEWITT+BRATTON- (UCI+MICH+BNL)
MARKY 85 PR D32 2215 MARKY,BOEHM (CIT)
OH 85 PL 1608 322 OHI,NAKAJIMA,TAMURA- (TOKY-INUS-KEK)
SIMPSON 85 PRL 54 1891 SIMPSON (GUEL)
STOCKDAL 85 ZPHY C27 53 STOCKDALE,BODEK+, (ROCH+CHIC+COLU+FNAL)
ZACEK 85 PL 1648 193 ZACEK,ZACEK,BOEHM- (MUNI+CIT+SIN)

LIMITS ON NUMBER OF LIGHT NEUTRINO TYPES

OMITTED FROM SUMMARY TABLE

The neutrinos referred to in this section are those of the standard SU(2)×U(1) electroweak model. Light neutrinos are those with M(ν) << M(Z⁰). The limits are on the number of neutrino families or species. See also cosmological limits in the "ν Bounds from Astrophysics and Cosmology" section.

In the subsection on "Limits from pp Colliders," the results assume that there are only three families of quarks and three families of charged leptons light enough to contribute to W or Z decay. The results were derived from

Nν = [ΓZ(measured) - ΓZ(3-family theory)]/Cν + 3.

The term "3" above is for νe, νμ, and ντ; Cν is approximately 0.18 GeV; and the Γ's are measured in GeV. For

# Stable Particle Full Listings

## LIMITS ON NUMBER OF LIGHT $\nu$ TYPES, $\nu$ BOUNDS FROM ASTROPHYSICS & COSMOLOGY

the results reported here,  $\Gamma_Z$  (measured) is not a directly measured number, but rather an inferred number based on measured cross sections times branching fractions:

$$\Gamma_Z = \Gamma_W \frac{\Gamma(Z \rightarrow e^+e^-)}{\Gamma(W \rightarrow e\nu)} \frac{\sigma_Z}{\sigma_W} \left[ \frac{\sigma_W \cdot B(W \rightarrow e\nu)}{\sigma_Z \cdot B(Z \rightarrow e^+e^-)} \right]$$

For each result,  $\Gamma_W$  and  $\frac{\Gamma(Z \rightarrow e^+e^-)}{\Gamma(W \rightarrow e\nu)}$  are calculated from the standard model with three families, while  $\sigma_Z/\sigma_W$  is calculated from QCD (most uncertainties in QCD are thought to cancel in this ratio). Only  $\frac{\sigma_W \cdot B(W \rightarrow e\nu)}{\sigma_Z \cdot B(Z \rightarrow e^+e^-)}$  is a measured quantity. The errors quoted include the uncertainties from each theoretical and experimental quantity except  $\Gamma_W$ .

### LIMITS FROM $p\bar{p}$ COLLIDERS

NP NUMBER OF NU TYPES INCLUDING NU(E), NU(MU), NU(TAU).  
 NP SEE ABOVE NOTE FOR METHOD OF DERIVATION AND CRUCIAL ASSUMPTIONS.  
 NP A 7.3 OR LESS CL=.90 APPEL 86 UA2 EDM=546+630 GEV  
 NP B (10.) OR LESS CL=.90 ARNISON 86 UA1 EDM=546+630 GEV  
 NP A ASSUME MASS(T-QUARK)=40 GEV, C(NU)=0.177 GEV, GAMMA(W)=2.65 GEV  
 NP A AND GAMMA(Z, 3FAMILY THEORY)=2.72 GEV. APPEL 86 REPORTED THEIR  
 NP A LIMIT AS 5.6--1.7 OR LESS AND WE CHOSE THE UPPER VALUE.  
 NP B ASSUME MASS(T-QUARK)=40 GEV, C(NU)=0.182 GEV, GAMMA(W)=2.82 GEV  
 NP B AND GAMMA(Z, 3FAMILY THEORY)=2.83 GEV.

### REFERENCES FOR LIMITS ON NUMBER OF LIGHT $\nu$ TYPES

APPEL 86 ZPHY C TO BE PUB UA2 COLLAB. (BERN-CERN-BOHR-HEID-ORSA-PAVI-)  
 ARNISON 86 PL 9 TO BE PUB UA1 COLLAB. (AACH-ANIK+LAPP-BIRM-CERN-HARV-)

## NEUTRINO BOUNDS FROM ASTROPHYSICS AND COSMOLOGY

### OMITTED FROM SUMMARY TABLE

SEE THE NOTE ON NEUTRINOS BY R.E. SHROCK IN THE ELECTRON NEUTRINO SECTION NEAR THE BEGINNING OF THESE CARD LISTINGS. FOR INFORMATION ON NEUTRINOS DERIVED FROM MORE CONVENTIONAL (TERRESTRIAL) EXPERIMENTS, SEE THE NU(E), NU(MU), NU(TAU), AND HEAVY-NU SECTIONS ABOVE.

### NOTE ON $\nu$ MASS LIMITS

These limits apply to  $m_{tot}$  given by

$$m_{tot} = \sum_{j=1,n} \left( \frac{g_{\nu_j}}{2} \right) m_{\nu_j}$$

where  $n$  is the number of neutrino species and  $g_{\nu_j}$  is the number of independent components in the neutrino field:  $g_{\nu_j} = 4$  for Dirac neutrinos;  $g_{\nu_j} = 2$  for chiral Majorana neutrinos.

### $\nu$ MASS (eV)

M LIMIT ON TOTAL NU MASS, SUMMED OVER ALL SPIN COMPONENTS,  
 M OF EFFECTIVELY STABLE NEUTRINOS (I.E. THOSE WITH MEAN LIVES ROUGHLY  
 M EQUAL TO OR GREATER THAN THE AGE OF THE UNIVERSE) (UNITS EV)  
 M A 132. OR LESS COWSIK 72 COSM  
 M A 80. OR LESS LEE 77 COSM  
 M 100. OR LESS OLIVE 81 COSM  
 M A COWSIK 72 AND LEE 77 HAVE BEEN GENERALIZED TO APPLY TO M(TOT) AS  
 M A DEFINED ABOVE, WHERE ONE ALLOWS EITHER DIRAC OR MAJORANA NEUTRINOS.  
 M A THESE PAPERS ASSUMED DIRAC NEUTRINOS.  
 M FOR OTHER LIMITS, SEE SATO 77, VYSOTSKY 77, DICUS 78, HUT 79,  
 M ZELDOVICH 80, FREESE 84, AND SCHRAMM 84.  
 ML ASTROPHYSICAL AND COSMOLOGICAL LIMITS ON NU MASSES. (UNITS EV)  
 ML ANALYSES OF MASS/LIGHT RATIOS AND DYNAMICS OF GALAXIES AND  
 ML CLUSTERS ARE CONSISTENT WITH PRESENCE OF DARK MATTER. TREMAINE 79  
 ML STATED THAT THE DARK MATTER COULD NOT BE MUON OR ELECTRON NEUTRINOS  
 ML OF NONZERO REST MASS OR ANY NEUTRAL LEPTON LESS MASSIVE THAN 1 MEV.  
 ML AUTHORS BOND 81, DAVIS 81, AND SCHRAMM 81 CLAIMED THAT THIS DARK  
 ML MATTER COULD CONSISTENTLY BE ASCRIBED TO MASSIVE NEUTRINOS  
 ML WITH SUFFICIENTLY LONG LIFETIMES. SUBSEQUENT ANALYSES HAVE  
 ML CHALLENGED THIS CLAIM; SEE PRIMACK 83. FOR DEGENERATE NEUTRINOS,  
 ML SEE FREESE 83.  
 ML ORDER 30 TREMAINE 79 COSM ISOTHERMAL  
 ML BOND 81 COSM ADIABATIC  
 ML 50-100 DAVIS 81 COSM ADIA.+DECAYING NUS  
 ML 4-20 SCHRAMM 81 COSM ISOTHERMAL  
 ML NO CONSISTENT VALUE PRIMACK 83 COSM  
 S 1000 OR LESS SARKAR 84 COSM  
 M 22 OR MORE MADSEN 85 COSM ASSUME ISOTROPY  
 M 18 OR MORE MADSEN 85 COSM SOME ANISOTROPY  
 ML S SARKAR 84 RESULTS APPLY TO UNSTABLE NU(TAU) OR OTHER DIRAC NEUTRINO  
 ML S WITH SIGNIFICANT RADIATIVE DECAY MODES. THEY ASSUME THAT STANDARD  
 ML S COSM MODEL OF PRIMORDIAL NUCLEOSYNTHESIS IS NOT DISRUPTED BY  
 ML S NU(TAU) DECAY.  
 ML M MADSEN 85 ASSUME STABLE NEUTRINO (NU(E)) TO BE DARK MATTER IN  
 ML M GALAXY HALOS (ONLY DATA ON BIG GALAXIES).  
 ML M SECOND MADSEN 85 VALUE IS ALLOWING MAX REASONABLE ANISOTROPY.  
 ML M NOTE HOWEVER THAT IN MADSEN 84 THE ADDITIONAL ASSUMPTION THAT NU(E)  
 ML M IS THE DARK MATTER IN DWARF AS WELL AS BIG GALAXY HALOS (ASSUMING  
 ML M ISOTROPY) LEADS TO THE LIMIT 127 EV OR MORE IN CONFLICT WITH THE  
 ML M EXPERIMENTAL UPPER LIMIT OF 46 EV. FURTHERMORE THE ADDITIONAL  
 ML M ASSUMPTION THAT THE DISTRIBUTION OF DARK MATTER IS AN ISOTHERMAL  
 ML M SPHERE LEADS TO THE LIMIT (LIM 83) OF 500 EV OR MORE.  
 MSW LIMITS ON MASSES OF STABLE RIGHT-HANDED NU  
 MSW (WITH NECESSARILY SUPPRESSED INTERACTION STRENGTHS)  
 MSW O OLIVE 82  
 MSW O ALLOWED VALUES OF MASS ARE STRONGLY CORRELATED WITH INTERACTION  
 MSW O STRENGTH. SEE FIG. 1 OF OLIVE 82.

### $\nu$ RADIATIVE MEAN LIFE VERSUS MASS

T COWSIK 77 COSM  
 T DICUS 77 COSM  
 T GOLDMAN 77 COSM  
 T FALK 78 COSM  
 T STRONGLY CORRELATED COWSIK 79 COSM  
 T LIMITS. SEE REFERENCES GOLDMAN 79 COSM  
 T DERULJULA 80 COSM  
 T STECKER 80 COSM  
 T HENRY 81 COSM  
 T KIMBLE 81 COSM  
 T REPHALI 81 COSM  
 T TURNER 81 COSM  
 T KRAUSS 83 COSM  
 T BINETRUY 84 COSM

### NUMBER OF LIGHT TWO-COMPONENT $\nu$ TYPES (LIGHT MEANS < ABOUT 1 MEV)

N NUMBER COUPLING WITH FULL WEAK STRENGTH  
 N SHIVARTSMA 69 COSM  
 N (7) OR LESS STEIGMAN 77 COSM  
 N (4) OR LESS YANG 79 COSM  
 S CRITICISM OF BOUND STECKER2 80 COSM  
 N MAYBE NO FIRM BOUND OLIVE 81 COSM  
 N (4) OR LESS TURNER 81 COSM  
 N CRITICISM OF BOUND RANA 82 COSM  
 N (4) OR LESS SCHRAMM 82 COSM  
 N (FROM 10 TO 100) OR LESS ELLIS 83 ASTROPHYS MODEL-DEP  
 N (4) OR LESS YANG 84 COSM  
 S SEE HOWEVER OLIVE2 81 CRITIQUE AND STECKER 81 REPLY.  
 N UNCERTAINTIES AND CRITICISMS COME FROM DIFFERING ESTIMATES OF LOWER  
 N LIMIT ON BARYON DENSITY OF THE UNIVERSE AND UPPER LIMIT ON THE  
 N PRIMORDIAL HELIUM-4 ABUNDANCE. SEE ALSO BERNSTEIN 82.  
 NSW NUMBER COUPLING WITH LESS THAN FULL WEAK STRENGTH  
 NSW A (20) OR LESS STEIGMAN 79 COSM  
 NSW A LIMIT VARIES WITH STRENGTH OF COUPLING.

### MAGNETIC MOMENT OF SUFFICIENTLY LIGHT $\nu$ (UNITS EV/GAUSS)

MM S (4.9E-19)OR LESS SUTHERLAND 76 COSM FOR MU(NU)=10 KEV.  
 MM S USES SUTHERLAND 76 EQ.3 WITH F=1/3 FROM THEIR TABLE AS MODIFIED TO  
 MM S APPLY AS A LIMIT ON ANY ONE NEUTRINO SPECIES INDIVIDUALLY.

### REFERENCES FOR $\nu$ BOUNDS FROM ASTRO. AND COSM.

SHVARTSM 69 JETPL 9 184 V. F. SHVARTSMAN (MOSU)  
 COWSIK 72 PRL 29 669 R. COWSIK, J. MC CLELLAND (UCB)  
 SUTHERLAND 76 PR D13 2700 SUTHERLAND, NG, FLOWERS, - (PENN-COLU-NYU)

For notation, see key on page 91.

## Stable Particle Full Listings

### $\nu$ BOUNDS FROM ASTROPHYSICS & COSMOLOGY, HEAVY LEPTON SEARCHES

COWSIK 77 PRL 39 784	R. COWSIK (MPIM+TIFR)
ALSO 79 COWSIK	
DICUS 77 PRL 39 168	D. A. DICUS, E. W. KOLB, V. L. TEPLITZ (TEXA+VPI)
GOLDMAN 77 PR D16 2256	T. GOLDMAN, G. J. STEPHENSON (LASL)
LEE 77 PRL 39 165	B. W. LEE, S. WEINBERG (FNAL+STAN)
SATO 77 PTP 58 1775	K. SATO, M. KOBAYASHI (KOT)
STEIGMAN 77 PL 668 202	G. STEIGMAN, D. SCHRAMM, J. GUNN (YALE, CHIC, CIT)
VYSOTSKY 77 JETPL 26 188	VYSOTSKY, DOLGOV, ZELDOVICH (ITEP)
DICUS 78 PR D17 1529	+KOLB, TEPLITZ, WAGONER (TEXA+VPI+STAN)
FALK 78 PL 79B 511	S. FALK, D. SCHRAMM (CHIC)
ALSO 78 APJ 223 1015	GUNN, LEE (CIT+CAMB+FNAL+CHIC+YALE)
COWSIK 79 PR D19 2219	R. COWSIK (TIFR)
GOLDMAN 79 PR D19 2215	T. GOLDMAN, G. J. STEPHENSON (LASL)
HUT 79 PL 87B 144	P. HUT, K. A. OLIVE (AMSTERDAM+EFI)
STEIGMAN 79 PRL 43 239	G. STEIGMAN, K. OLIVE, D. SCHRAMM (BART+EFI)
TREHARNE 79 PRL 42 407	S. TREHARNE, J. E. GUNN (CIT+CAMB+CAW)
YANG 79 APJ 227 697	YANG, SCHRAMM, STEIGMAN, ROOD (CHIC+YALE+VIRG)
ALSO 79 STEIGMAN	FOOTNOTE 4
DERUJULA 80 PRL 45 942	A. DE RUJULA, S. L. GLASHOW (MIT+HARV)
STECKER 80 PRL 45 1460	F. W. STECKER (NASA)
ALSO 81 NU 81 CONF HAWAII	F. W. STECKER (PROC. V. 1, P. 124) (NASA)
STECKER2 80 PRL 44 1237	F. W. STECKER (NASA)
ZELDOVIC 80 SJNP 31 664	ZELDOVICH, KLIPIN, KHLOPOV, CHECHETKIN
BOND 81 NU 81 CONF HAWAII	J. R. BOND, A. S. SZALAY (UCB+CHIC)
DAVIS 81 APJ 250 423	M. DAVIS, M. LECAR, C. PRYOR, E. WITTEN (HARV+PRIN)
HENRY 81 PRL 47 618	R. C. HENRY, P. D. FELDMAN (JHU)
KIMBLE 81 PRL 46 80	R. KIMBLE, S. BOWYER, P. JAKOBSEN (UCB)
OLIVE 81 APJ 246 557	+SCHRAMM, STEIGMAN, TURNER, YANG (CHIC+BART)
OLIVE2 81 PRL 46 516	K. A. OLIVE, M. S. TURNER (CHIC+UCSB)
REPHAELE 81 PL 106B 73	Y. REPHAELE, A. S. SZALAY (EFI)
SCHRAMM 81 APJ 243 1	D. N. SCHRAMM, G. STEIGMAN (UCSB+CHIC)
STECKER 81 PRL 46 517	F. W. STECKER (NASA)
TURNER 81 NU 81 CONF HAWAII	M. S. TURNER (UCSB+CHIC)
BERNSTEIN 82 PRL 48 774	J. BERNSTEIN (STEV)
OLIVE 82 PR D25 213	K. A. OLIVE, M. S. TURNER (CHIC+UCSB)
RANA 82 PRL 48 209	H. C. RANA (TIFR)
SCHRAMM 82 PRSL A307 43	D. N. SCHRAMM (CHIC)
ELLIS 83 NP B223 256	J. ELLIS, K. A. OLIVE (CERN)
FREISE 83 PR D27 1689	K. FREISE, E. W. KOLB, M. S. TURNER (CHIC+UCSB)
KRAUSS 83 PL 128B 37	L. M. KRAUSS (HARV)
LIN 83 APJ 266 L21	LIN, FABER (UCSC)
PRIMACK 83 PHILADELPHIA	J. PRIMACK (4TH WKSHIP. ON GRAND UNIF.) (UCSC)
ALSO 82 NAT 299 37	BLUMENTHAL, PAELS, PRIMACK (UCSC+ROCK)
BINETRUY 84 PL 134B 174	P. BINETRUY, G. GIRARDI, P. SALATI (LAPP)
FREISE 84 NP B233 167	+SCHRAMM (CHIC+FNAL)
MAIDEN 84 APJ 282 11	J. MAIDEN, R. EPSTEIN (AARH-LANL)
SARKAR 84 PL 148B 347	S. SARKAR, A. W. COOPER (OXF-CERN)
SCHRAMM 84 PL 141B 337	SCHRAMM, STEIGMAN (FNAL+STAN)
YANG 84 APJ (TO BE PUBL)	+TURNER, STEIGMAN, SCHRAMM, OLIVE (CHIC+BART)
MAIDEN 85 PRL 54 2720	J. MAIDEN, R. EPSTEIN (AARH-LANL)

$$\left. \begin{aligned} L^- &\rightarrow \nu_L e^- \bar{\nu}_e \\ L^- &\rightarrow \nu_L \mu^- \bar{\nu}_\mu \\ L^- &\rightarrow \nu_L \text{ hadrons} \end{aligned} \right\} \text{ are allowed.}$$

There could be an increasing mass sequence of such pairs. It is frequently assumed that the neutrinos are massless.

Decay rates are assumed calculable from conventional weak interaction theory. For an  $L^-$  mass between 1 and 3 GeV, the branching fraction to each of the two leptonic modes above should be roughly 10 to 20%. For an  $L^-$  mass above 1 GeV, the mean life should be  $\leq 10^{-12}$  second.

*Paraleptons* ( $E^+$ ,  $E^0$ ) and ( $M^+$ ,  $M^0$ ). These pairs have the same lepton numbers as the opposite-charge ordinary leptons, i.e.,  $e^-$  and  $\mu^-$ , respectively. Radiative decays are again forbidden and decays similar to those allowed for  $L^-$  are allowed here, e.g.,

$$M^+ \rightarrow \nu_\mu e^+ \nu_e$$

or

$$M^+ \rightarrow \nu_\mu \mu^+ \nu_\mu.$$

However, the lightest member is not stable as is the case for sequential leptons, so that bizarre decay schemes such as (assuming  $m_{E^0} < m_{E^+}$ )

$$E^+ \rightarrow E^0 \mu^+ \nu_\mu$$

$$\downarrow$$

$$e^- e^+ \nu_e$$

are allowed.

Heavy leptons of this type were proposed (before the discovery of the  $Z^0$  boson) in unified gauge theories of weak and electromagnetic interactions to cancel unphysical high energy behavior in such processes as  $e^+ e^- \rightarrow W^+ W^-$ .<sup>3</sup>

*Ortholeptons* ( $F^-$  and  $N^-$ ). These have the same lepton numbers as  $e^-$  and  $\mu^-$ , respectively. They may or may not have associated neutral leptons. Radiative decays are allowed in addition to weak modes similar to those of sequential leptons. The radiative mode can dominate or can be relatively unimportant depending on the model.<sup>4</sup> Decays such as

$$F^- \rightarrow e^- + \text{hadrons}$$

are also allowed.

## HEAVY LEPTON SEARCHES

OMITTED FROM SUMMARY TABLE

Data on the  $\tau^\pm$  are listed in a separate section above, following the  $e$  and  $\mu$  listings.

The following section contains information on searches for heavy leptons of other types and searches for the  $\tau^\pm$  in collisions other than  $e^+ e^-$ .

Several types of heavy leptons (that is, non-strongly-interacting fermions other than  $e$  and  $\mu$ ) have been proposed. In the Full Listings we distinguish four types.<sup>1,2</sup> Each has a corresponding antiparticle with opposite charge and lepton number. For convenience we omit writing the antiparticles in the following descriptions. The four types are:

*Sequential leptons* ( $L^-, \nu_L$ ). Such a pair is assumed to have its own separately strictly conserved lepton number  $n_L = +1$ . This means that the radiative decays

$$\left. \begin{aligned} L^- &\rightarrow e^- \gamma \\ L^- &\rightarrow \mu^- \gamma \end{aligned} \right\} \text{ are forbidden,}$$

while the weak decays (assuming  $m_L$  sufficiently large)

# Stable Particle Full Listings

## HEAVY LEPTON SEARCHES

*Long-lived penetrating particles.* Heavy leptons could have long mean lives under certain circumstances. For example, if  $m_{\nu_{\tau}} > m_{L^-}$ , then  $L^-$ , the sequential lepton, is completely stable since its lepton number is conserved.

*Experimental results.* The results are summarized in the Full Listings below. Mass limits for sequential leptons are listed in subsection MS, while all other types are listed together in subsection M.

The Full Listings also contain cross-section upper limits reported as results of unsuccessful searches. We no longer list cross sections for anomalous  $\mu$  events in  $e^+e^-$  collisions. These cross sections are consistent with coming from  $e^+e^- \rightarrow \tau^+\tau^-$  where the  $\tau^\pm$  is assumed to be a spin-1/2 Dirac point particle with a mass about 1785 MeV.

### References

1. M.L. Perl and P. Rapidis, SLAC-PUB-1496 (October 1974).
2. C.H. Llewellyn Smith, Invited paper presented at the Royal Society Meeting on New Particles and New Quantum Numbers, 11 March 1976, Oxford Ref. 33/76.
3. J.D. Bjorken and C.H. Llewellyn Smith, Phys. Rev. **D7**, 887 (1973).
4. F. Wilczek and A. Zee, Nucl. Phys. **B106**, 461 (1976).

SEE PERL 81 FOR A REVIEW

PROPERTIES OF THE TAU-(1785) HEAVY LEPTON AND ITS ASSOCIATED NEUTRINO ARE LISTED SEPARATELY ABOVE FOLLOWING THE E AND MU LISTINGS. THE FOLLOWING SECTION CONTAINS INFORMATION ON SEARCHES FOR HEAVY LEPTONS OF OTHER TYPES AND SEARCHES FOR TAU- IN COLLISIONS OTHER THAN E-E- WE LIST MASS LIMITS AND CROSS SECTION UPPER LIMITS REPORTED AS NEGATIVE SEARCH RESULTS. WE NO LONGER LIST CROSS SECTIONS FOR THE ESTABLISHED PROCESS E+ E- -> TAU+ TAU- AS WAS DONE IN OUR 1977 SUPPLEMENT. SEARCHES FOR FRACTIONALLY CHARGED HEAVY LEPTONS ARE INCLUDED IN THE FREE QUARK SEARCH SECTION.

### HEAVY LEPTON MASS LIMITS

LIMITS APPLY ONLY TO HEAVY LEPTON TYPE GIVEN IN COMMENT AT RIGHT ON DATA LISTING. SEE REVIEW ABOVE FOR DESCRIPTION OF TYPES. L,E,M,F,N STAND FOR SEQUENTIAL LEPTON, PARA-ELECTRON, PARA-MUON, ORTHO-ELECTRON, ORTHO-MUON RESPECTIVELY.

SEQUENTIAL HEAVY LEPTON MASS LIMITS (GEV)		
MS A	(13.) OR MORE	AZIMOV 80 --
MS B	(16.) OR MORE CL=.95	BARBER 80 CNTR --
MS C	NONE ABOVE TO 14.5GEV CL=.95	BERGER 81 PLUT --
MS D	(15.5) OR MORE CL=.95	BRANDELIX 81 TASS --
MS E	(14.) OR MORE CL=.95	ADEVA 82 MRKJ +-
MS F	(18.) OR MORE CL=.95	ADEVA 83 MRKJ +-
MS G	(22.5) OR MORE CL=.95	ADEVA 83 JADE +- ECM=40-47 GEV

MS A AZIMOV 80 ESTIMATED PROBABILITIES FOR M-N TYPE EVENTS IN E, E- -> L+ L- DEDUCTING SEMI-HADRONIC DECAY MULTIPLICITIES OF L FROM E+ E- ANNIHILATION DATA AT WCM=(2/3)\*ML. OBTAINED ABOVE LIMIT COMPARING A THESE WITH E+E- DATA (BRANDELIX 80, PL 92B 199).

MS B BARBER 80 LOOKED FOR E+ E- -> L+ L-, L->NU(L)-X WITH MARK-J AT B DESY-PETRA.

MS C BERGER 81 IS DESY DORIS AND PETRA EXPT. LOOKING FOR E+E- -> L+L-.

MS D BRANDELIX 81 IS DESY PETRA EXPT. LOOKING FOR E+E- -> L+L-.

MS E ADEVA 83 LOOKED FOR MUON OPPOSITE AGAINST A HADRON JET.

MS F BARTEL 83 LIMIT IS FROM PETRA E+E- EXP WITH AVERAGE WCM=34.2 GEV.

MS G ADEVA 85 ANALYZE ONE-ISOLATED-MUON DATA AND SENSITIVE TO TAU < 10 NANOSEC. ASSUME BR(LEPTON)=0.50.

HEAVY LEPTON MASS LIMITS (GEV)		
M A	1.0 OR MORE	BEHREND 65 SPEC - ORTHOELECTRON(F)
M B	NONE BETWEEN 0.12 AND 0.57	BETOURNE 65 SPEC - ORTHOELECTRON(F)
M C	NONE BETWEEN 0.3 AND 0.7	BUDNITZ 66 SPEC - ORTHOELECTRON(F)
M D	NONE BETWEEN 0.2 AND 0.92	BARNA 68 CNTR - LONG-LIVED
M E	NONE BETWEEN 0.97 AND 1.03	BARNA 68 CNTR - LONG-LIVED
M D	NONE BETWEEN 0.1 AND 1.5	BOLEY 68 SPEC - ORTHOELECTRON(F)
M F	NONE BETWEEN 0.2 AND 0.6	LIBERMAN 69 OSPK - ORTHOMUON(N)
M G	0.490 OR MORE	ROTHE 69 RVUE
M H	NONE BETWEEN 0.26 AND 1.32	LICHTENST 70 SPEC - ORTHOELECTRON(F)
M I	20 (0.424) (0.013) (0.002)	RAMM 70 HLBC 0
M J	22 (0.431) (0.004)	RAMM 71 HLBC - ORTHOMUON(N)
M K	0 0.1 OR MORE	ANSORGE 73 HBC - LONG-LIVED
M L	0 0.6 OR MORE	BACCI 73 ELEC +- ORTHOELECTRON(F)
M K	0 2.2 OR MORE	BACCI 73 ELEC +- ORTHOELECTRON(F)
M L	0 2.0 OR MORE CL=.90	BARISH 73 ASPK + PARAMUON (M)
M M	0 1.4 OR MORE CL=.95	BERNARDIN 73 ASPK -- ANY NON-RAD TYPE
M N	0 1.0 OR MORE CL=.95	BERNARDIN 73 ASPK -- ANY NON-RAD TYPE
M N	NONE BETWEEN 0.55 AND 4.5	BUSHNIN 73 CNTR - LONG-LIVED
M O	2.4 OR MORE CL=.90	EICHTEN 73 HLBC + PARAMUON (M)
M P	7.8 OR MORE CL=.95	HANSON 73 WIRE ORTHOELECTRON(F)
M Q	1.8 OR MORE CL=.90	ASRATYAN 74 HLBC - ORTHOMUON (N)
M R	8.4 OR MORE CL=.90	BARISH 74 SPEC + PARAMUON (M)
M S	NONE BETWEEN 0 AND 2.0	GITTLESON 74 SPEC ORTHOMUON (N)
M T	0 1.15 OR MORE CL=.95	ORITO 74 ASPK -- ANY NON-RAD TYPE
M U	NONE BETWEEN 0.25 AND 2.3	BACCI 77 SPEC -- ORTHOELECTRON(F)
M V	1.2 OR MORE	MEYER 77 SMAG 0 NEUTRAL
M W	10.3 OR MORE CL=.98	ASRATYAN 78 -- ORTHOMUON (N)
M X	0 7.5 OR MORE	CNOPS 78 HLBC - ORTHOMUON (N)
M Y	0 9.0 OR MORE	CNOPS 78 CNPS +- PARAMUON (M)
M Z	10.0 OR MORE	ERRIQUEZ 78 BEBC
M Z	12. OR MORE CL=.90	HOLDER 78 CNTR + PARAMUON (N)
M 1	NONE 1 GEV TO 9 GEV CL=.90	CLARK 81 SPEC 0 PARAMUON(MOBAR)
M 1	NONE 1 GEV TO 9 GEV CL=.95	CLARK 81 SPEC +-
M 2	NONE BETWEEN 0.6 AND 3.3	HAYES 82 SMK2 -- ORTHOMUON (N)
M 2	NONE BETWEEN 0.5 AND 3.3	HAYES 82 SMK2 -- ORTHOELECTRON(F)
M 3	24.5 OR MORE CL=.95	BARTEL 83 JADE 0 PARAELECTRON(E0)
M 3	22.5 OR MORE CL=.95	BARTEL 83 JADE 0 PARAELECTRON(E0)

M A BEHREND 65 IS DESY EXPT. LOOKS FOR E P -> F P, F -> E GAMMA.

M A THIS MASS LIMIT CORRESPONDS TO A LIMIT ON LAMBDA\*\*2 OF 6.25\*10\*\*+4.

M B BETOURNE 65 IS ORSAY EXPT. LOOKS FOR E P -> F P. MASS OF .12 CORRESPONDS TO COUPLING CONSTANT LAMBDA\*\*2 GT .0016, MASS OF .57 TO LAMBDA\*\*2 GT .22.

M C BUDNITZ 66 IS CEA EXPT. LOOKS FOR E P -> F P.

M D BARNA 68 IS SLAC PHOTOPRODUCTION EXPT.

M E BOLEY 68 IS CEA EXPT. LOOKS FOR E P -> F P. MASS OF .1 CORRESPONDS TO COUPLING CONSTANT LAMBDA\*\*2 GT 3\*10\*\*+4, MASS LIMIT OF 1.3 TO E LAMBDA\*\*2 GT .01.

M F LIBERMAN 69 IS A BNL EXPT MEASURING MUON BREMSSTRAHLUNG.

M G ROTHE 69 EXAMINES PREVIOUS DATA ON MU PAIR PROD AND PI AND K DECAYS

M H LICHTENSTEIN 70 IS CORNELL EXPT MEASURING E BREMSSTRAHLUNG.

M H MASS LIMIT DEPENDS ON COUPLING CONSTANT. FIRST VALUE ABOVE IS FOR LAMBDA\*\*2 GT .17, SECOND IS FOR LAMBDA\*\*2 GT .42.

M I RAMM 70 FINDS PEAK IN MU PI COMBINED MASS PRODUCED BY NEUTRINO INTERACTIONS. HE ALSO CLAIMS EVIDENCE FOR THIS IN KOMU3 DECAYS IN HBC WHERE PI MU COMBINED MASS PEAKS IN SAME REGION. CLARK 72 FINDS NO EVIDENCE FOR PI MU PEAK IN HIGH STATISTICS KL3 EXPT.

M I RAMM 71 SEES PEAK IN MU GAMMA COMBINED MASS PRODUCED BY NEUTRINOS.

M J ANSORGE 73 LOOKS FOR ELECTRON PAIR PROD AND ELECTRON-LIKE BREMS.

M K BACCI 73 IS FRASCATI E+E- EXPT. LOOKS FOR F -> E GAMMA.

M K MASS LIMIT DEPENDS ON COUPLING CONSTANT LAMBDA FOR THIS DECAY.

M K FIRST VALUE ABOVE IS FOR LAMBDA\*\*2 GT 9\*10\*\*+5, 2ND IS FOR LAMBDA\*\*2 GT 10\*\*+3.

M L BARISH 73 IS FNAL 50,145 GEV NU EXPT. LOOKS FOR (NU NUCLEON -> L+ M+ ANYTHING). ASSUMES (M+ -> MU+ NU NU) WITH BR=.3.

M M BERNARDINI 73 IS FRASCATI E+E- EXPT. FIRST VALUE ASSUMES UNIVERSAL COUPLING TO ORDINARY LEPTONS. SECOND VALUE ALSO ASSUMES COUPLING TO HADRONS.

M N BUSHNIN 73 IS SERPUKOV 70 GEV P EXPT. MASSES ASSUME MEAN LIFE ABOVE 7E-10 AND 3E-8 RESPECTIVELY. CALCULATED FROM CROSS SEC(DC BELOW N AND 30 GEV NUON PAIR PRODUCTION DATA.

M O EICHTEN 73 IS CERN 1-10GEV NU EXPT. LOOKS FOR M+ PRODUCED IN O NU NUCL -> M- HADRONS ASSUMING 15 PERCENT DECAY TO E+ NU NU.

M P HANSON 73 LOOK FOR DEVIATIONS FROM QED IN E+ E- ->2GAMMA. THEY MEASURE THE PRODUCT OF THE F MASS \* THE COUPLING CONSTANT LAMBDA, WHICH IS THE VALUE QUOTED ABOVE.

M Q ASRATYAN 74 USES EICHTEN 73 DATA ON NU NUCL -> E- HADRONS AND Q NUBAR NUCL -> E+ HADRONS TO SET LIMITS ON ORTHOMUON PRODUCTION.

M R BARISH 74 IS FNAL 50,135 GEV NU EXPT. LOOKS FOR (NU NUCLEON -> R+ M+ ANYTHING). ASSUMES (M+ -> MU+ NU NU) WITH BR=.3.

M S GITTLESON 74 IS MU P -> P ORTHOMUON SEARCH. COUPLING CONSTANT LAMBDA\*\*2 IS <.01 FOR MASS UP TO .7 GEV, LIMIT ON LAMBDA\*\*2 RISES TO <.1 FOR MASS OF 2.0 GEV.

M T ORITO 74 LOOKED FOR H+- PAIRS GIVING MU-E PAIRS. MASS LIMIT REFERS TO ANY NON-RADIATIVE TYPE HEAVY LEPTON -- L, E, M, F, N.

M T COUPLING TO HADRON ASSUMED FROM THEORETICAL MODELS.

M U BACCI 77 IS SAME TYPE AS BACCI 73. LOWER MASS LIMIT CORRESPONDS TO LAMBDA\*\*2 LIMIT OF 4\*10\*\*+5, UPPER VALUE IS FOR LAMBDA\*\*2 LIMIT OF 1.5\*10\*\*+3.

M V MEYER 77 LOOKS FOR NARROW NEUTRAL RESONANCE (INCE P1)AND(MU P1) CHANNELS PRODUCED BY E+E- AT 6.8 GEV (ECM). ASSUMED TO BE DECAY PRODUCT OF THE TAU. SEE SECTION NE BELOW.

M W ASRATYAN 78 ANALYZES DEPENDENCE OF N.C./C.C. ON ENERGY OF ASSOC. HADRONS. USES DATA OF HOLDER 77 (PL 72B, 254)-NU(MU) INTERACTIONS AT CERN-SPS.

M X CNOPS 78 IS FNAL EXPT LOOKING FOR NU(MU) NE -> L+(-), FOLLOWED BY L-(-) -> E(-) NU NU.

M Y ERRIQUEZ 78 IS CERN SPS EXPT. LOOKS FOR NU(MU) NUCLEON -> MU- E+ X. FINDS CS FOR PRODUCING HYP LEPT -> E- <.7\*10\*\*+3 %C.C. CS.

M Z HOLDER 78 IS A CERN NU EXPT LOOKING FOR NU(MU) NUCLEON -> MU+ ANY THING. ASSUMES M+ -> MU+ 2NU(MU) WITH BR=0.2.

For notation, see key on page 91.

Stable Particle Full Listings HEAVY LEPTON SEARCHES

M 1 CLARK 81 IS FNAL EXP WITH 209 GEV MUONS. BOUNDS APPLY TO MU WHICH COUPLES WITH FULL WEAK STRENGTH TO MUON. SEE ALSO SECTION MU.

HEAVY LEPTON PROD. CROSS SECTION (μN) (cm²)

MU SEE ALSO SECT 'MU' IN CHARM SEARCHES AND OTHER NEW PARTICLE SEARCHES MU A 1.22E-34 OR LESS LEBRITTON 80 SPEC MU->MU+ MU- NU

HEAVY LEPTON PROD. DIFFERENCE CROSS SEC. (p N) (cm²/sr-GeV)

DC A 0 1.6E-37 OR LESS CL=.90 GOLDOVIN 72 CNTR-70GEV P, SERPUKHOV DC B 0 4. E-38 OR LESS CL=.90 BUSHNIN 73 CNTR-70GEV P, SERPUKHOV

PRODUCTION OF HEAVY LEPTON IN BEAM DUMP

BD A LOSECCO 81 AT BNL AGS SET LIMIT FOR CS(PROD)\*CS(INT) RATIO OF BD A SLOW (BETA<0.89) HEAVY LEPTONS TO PROMPT NU'S AS 2.2E-2(CL=.90).

INVARIANT HEAVY LEPTON PROD. CROSS SEC. (p N) (cm²/GeV²)

IC A 0 5.4E-39 OR LESS CL=.90 CRONIN 74 SPEC - M=1-6.8 GEV IC B 0 6.4E-35 OR LESS CL=.90 BINTINGER 75 SPEC -- M=1-5 GEV

HEAVY LEPTON PROD. CROSS SEC. (σ(HEAVY LEPTON)/σ(π))

RP1 A 0 7. E-12 OR LESS CL=.95 BUSSIERE 80 CNTR Q=-1 M=4-4.5 GEV RP1 A 0 2.5E-12 OR LESS CL=.95 BUSSIERE 80 CNTR Q=-2 M=5-7.5 GEV

NEUTRAL HEAVY LEPTON PROD. CROSS SECTION (cm²)

CN A 5 (1. E-37 OR MORE) KRISHNASWAMY 75 CNTR +/- M=2-5 GEV CN B 0 BENVENUTI 75 SPEC 0

NEUTRAL HEAVY LEPTON PRODUCED IN ν INTERACTIONS

N A 1 POSSIBLY SEEN BARANOV 77 HBCB 0 SERPUKHOV N A 2 POSSIBLY SEEN BARANOV 79 HBCB 0 SERPUKHOV

NEUTRAL HEAVY LEPTON PROD. CROSS SEC. (p N) (cm²)

CP A 0 1. E-29 OR LESS FAISSNER 76 HBCB 0 CP B 0 2.8E-35 OR LESS CL=.90 BECHIS 78 SPEC 0

NEUTRAL HEAVY LEPTON PROD. CROSS SEC. (e+e-) (10⁻⁵ nb)

NE CS\*BR(TAU-> NEW NEUTRAL LEPTON)\*BR(NEUTRAL LEPTON-> E OR MU P1) NE A 450 OR LESS CL=.90 MEYER 77 SMAG FOR M(L)=0.5 GEV

LIMIT ON ντ PRODUCTION IN BEAM DUMP EXPERIMENT

TNU A FRITZE 80 HYBR TNU A FRITZE 80 IS CFBN SPS EXP WITH BECC. NC/CC RATIO CORRESPONDS TO TNU A R=(PROMPT-NU(TAU)-INDUCED EVENTS)/(ALL PROMPT-NU EVENTS) <0.1.

LIMITS ON EXCITED e AND μ (e\* AND μ\*)

EXC A NONE ABOVE 58 GEV CL=.95 ADEVA 82 MRJK E\* PROD IN E+E- EXC B NONE BELOW 10 GEV CL=.95 ADEVA 82 MRJK E\* PROD IN E+E-

REFERENCES FOR HEAVY LEPTON SEARCHES

BEHREND 65 PRL 15 900 -BRASSE, ENGLER, GANSSAUGE- (DESY-KARL) BETHOURNE 65 PL 17 70 +NGUYEN NGOC PEREZ Y JORBA- (ORSA)





For notation, see key on page 91.

Stable Particle Full Listings

$\pi^\pm, \pi^0$

Table with columns for particle type, count, and rate. Includes entries for  $\pi^\pm \rightarrow (\pi^0 e^\pm \nu)/total$  (units  $10^{-8}$ ) and  $\pi^\pm \rightarrow (e^\pm \nu \gamma)/total$  (units  $10^{-8}$ ).

Table with columns for particle type, count, and rate. Includes entries for  $\pi^\pm \rightarrow (e^\pm \nu e^\pm e^-)/total$  (units  $10^{-8}$ ) and  $\pi^+ \rightarrow (\mu^+ \bar{\nu}_e)/total$  (units  $10^{-3}$ ).

Table with columns for particle type, count, and rate. Includes entries for  $\pi^\pm \rightarrow (\mu^+ \nu_e)/total$  (units  $10^{-3}$ ) and  $\pi^\pm$  POLARIZATION OF EMITTED  $\mu^\pm$ .

Table with columns for particle type, count, and rate. Includes entries for  $\pi^+$  POLARIZATION OF EMITTED  $\mu^+$ .

Table with columns for particle type, count, and rate. Includes entries for  $\pi^+$  POLARIZATION OF EMITTED  $\mu^+$ .

POL PI+- INTO MU+- NU
TESTS LORENTZ STRUCTURE OF LEPTONIC CHARGED WEAK INTERACTIONS
POL (0.9959)OR MORE CL=.90 FETSCHER 84 RVUE

REFERENCES FOR  $\pi^\pm$

List of references for  $\pi^\pm$  decays, including authors like CROWE, BARKAS, and ANDERSON, and their respective publications.

PAPERS NOT REFERRED TO IN DATA LISTINGS

List of papers not referred to in data listings, including authors like CARTWRIGHT and MERRISON.

Table with columns for author, journal, and volume/page. Includes entries for DEPOMMIE, WILKIN, and BRYMAN.

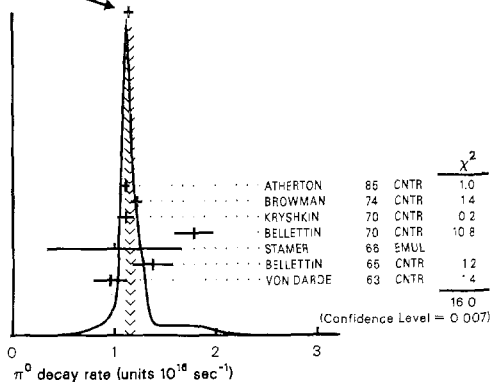
$\pi^0$

$I^G(J^{PC}) = 1^-(0^{+-})$

Table with columns for particle type, count, and rate. Includes entries for  $\pi^\pm \rightarrow \pi^0$  MASS DIFFERENCE (MeV) and  $\pi^0$  MEAN LIFE (units  $10^{-16}$  sec).

Table with columns for particle type, count, and rate. Includes entries for  $\pi^0$  MEAN LIFE (units  $10^{-16}$  sec) and  $\pi^0$  PARTIAL DECAY MODES.

WEIGHTED AVERAGE
1.160 = 0.049 (ERROR SCALED BY 1.8)



$\pi^0$  PARTIAL DECAY MODES

Table with columns for decay mode, count, and rate. Includes entries for  $\pi^0 \rightarrow 2\gamma$ ,  $\pi^0 \rightarrow e^+e^-\gamma$ ,  $\pi^0 \rightarrow 4e$ ,  $\pi^0 \rightarrow 3\gamma$ ,  $\pi^0 \rightarrow 4\gamma$ ,  $\pi^0 \rightarrow e^+e^-$ ,  $\pi^0 \rightarrow 2\nu$ , and  $\pi^0 \rightarrow \mu^+e^- + \mu^-e^+$ .

$\pi^0$  BRANCHING RATIOS

Table with columns for particle type, count, and rate. Includes entries for  $\pi^0 \rightarrow (\gamma e^+ e^-)/(2\gamma)$  (%) and  $\pi^0$  BRANCHING RATIOS.

# Stable Particle Full Listings

$\pi^0, \eta$

$\pi^0 \rightarrow (3\gamma)/total$  (units  $10^{-6}$ ) (P4)  
 R2 FORBIDDEN BY C INVARIANCE.  
 R2 D 0 (4.9) OR LESS CL=.90 DUCLOS 65 CNTR  
 R2 D (4.9) OR LESS CL=.90 KUTIN 65 CNTR  
 R2 D THESE EXPTS. GIVE BR(GAMMA/2GAMMA)<5.0\*10\*\*6.  
 R2 0 (1.5) OR LESS CL=.90 AUERBAC1 78 CNTR  
 R2 0 0.38 OR LESS CL=.90 HIGHLAND 80 CNTR

$\pi^0 \rightarrow (e^+ e^- e^- e^-)/(2\gamma)$  (units  $10^{-5}$ ) (P3)/(P1)  
 R3 N 146 (3.18) (0.30) SAMIOS 62 HBC SEE NOTE N BELOW  
 R3 N 3.28 THEORET. CALC. MIYAZAKI 75 QUANTUM ELECT.

$\pi^0 \rightarrow (4\gamma)/total$  (units  $10^{-5}$ ) (P5)  
 R4 A 0 (6.0) OR LESS CL=.90 ABRAMS 75 ASPK  
 R4 A ABRAMS 75 GIVES BR(4GAMMA/2GAMMA)<6.1\*10\*\*5.  
 R4 0 (3.8) OR LESS CL=.90 AUERBAC2 78 CNTR  
 R4 0 0.44 OR LESS CL=.90 AUERBACH 80 CNTR

$\pi^0 \rightarrow (e^+ e^-)/total$  (units  $10^{-6}$ ) (P6)/(P1)  
 R5 D (2.0) OR LESS CL=.90 DAVIES 74 RVUE  
 R5 (4.0) OR LESS CL=.90 SCHACHER 77 STRC P1 - P -> P10 N  
 R5 B (0.22) 0.240 0.110 FISCHER2 78 SPRK K- EXPT. CL=.90  
 R5 58 (0.18) (0.06) MISCHKE 82 SPEC REPL.BY FRANK 83  
 R5 59 0.17 0.07 FRANK 83 SPEC P1 - P -> P10 N  
 R5 D DAVIES 74 EXTRACTS THIS INFORMATION FROM BLOCH 75 K EXPERIMENT.  
 R5 AVG 0.177 0.065 AVERAGE

$\pi^0 \rightarrow (2\nu)/total$  (units  $10^{-5}$ ) (P7)  
 R6 0 2.4 OR LESS CL=.90 HERCZEG 81 RVUE

$\pi^0 \rightarrow (\mu^+ e^- + \mu^- e^+)/total$  (P8)  
 R7 FORBIDDEN BY LEPTON FAMILY NUMBER CONSERVATION  
 R7 7E-B OR LESS BRYMAN 82 RVUE K- -> P1+ MU E  
 R7 (14E-8) OR LESS HERCZEG 84 RVUE K- -> P1+ MU E  
 R7 (2E-15) OR LESS HERCZEG 84 THEO MU- -> E- CONV

## $\pi^0$ ELECTROMAGNETIC FORM FACTOR

THE AMPLITUDE FOR THE PROCESS  $\pi^0 \rightarrow e^+ e^- \gamma$  CONTAINS A FORM FACTOR  $GAMMA(X**2)$  AT THE  $(\pi^0 \gamma)$  VERTEX WHERE  $X=MASS(E-E-)/MASS(\pi^0)$ . THE PARAMETER A IN THE LINEAR EXPANSION  $GAMMA(X**2)=1+A*(X**2)$  IS LISTED BELOW.

## LINEAR COEFFICIENT OF $\pi^0$ ELECTROMAGNETIC FORM FACTOR

A (-0.15) (0.10) KOBRAK 61 HBC NO RAD. CORR.  
 A 3071 (-0.24) (0.16) SAMIOS 61 HBC NO RAD. CORR.  
 A 2200 (-0.01) (0.11) DEVONS 69 OSPK NO RAD. CORR.  
 A F 30K +0.10 0.03 FISCHER1 78 SPEC RAD. CORR.  
 A T TUPPER 83 THEO FISCHER1 78 DATA  
 A F ERROR STATISTICAL ONLY. RESULT WITHOUT RAD. CORR. = +0.05 - 0.03.  
 A T TUPPER 83 IS THEORETICAL ANALYSIS OF FISCHER1 78 INCLUDING TWO-  
 A T PHOTON EXCHANGE. THEIR ESTIMATE OF THE IMPACT OF THESE CORRECTIONS  
 A T IS THAT THE MODIFIED NUMBER WOULD BE 0.12+0.05-0.04.

## REFERENCES FOR $\pi^0$

PANOFSKY 51 PR 81 565 W K N PANOFSKY, R L AAMODT, J HADLEY (LRL)  
 CHINOWSK 54 PR 93 586 W CHINOWSKY, J STEINBERGER (COLUMBIA)  
 CASSELS 59 PRL 74 92 CASSELS, JONES, MURPHY, O'NEILL (LIVERPOOL)  
 HADDOCK 59 PRL 3 478 HADDOCK, ABASHIAN, CROWE, CZIRI (LRL)  
 HILLMAN 59 NC 14 887 HILLMAN, MIDDELKOOP, YAMAGATA, ZAVATTINI (CERN)  
 BUDAEOV 60 JETP 11 755 BUDAEOV, VIKTOR, DZHELEPOV, ERMOLOV (JINR)  
 JOSEPH 60 NC 16 997 D W JOSEPH (EFI)  
 SAMIOS 60 NC 18 154 N P SAMIOS (COLUMBIA)  
 GLASSER 61 PR 123 1014 R C GLASSER, N SEEMAN, B STILLER (LRL)  
 KOBRAK 61 NC 20 1115 H KOBRAK (EFI)  
 SAMIOS 61 PR 121 275 N P SAMIOS (COLUMBIA+BNL)  
 SAMIOS 62 PR 126 1844 SAMIOS, PLANO, PRODELL (COLUMBIA+BNL)  
 TIETGE 62 PR 127 1324 J TIETGE, W PUESCHEL (MAX PLANCK INST)  
 CZIRI 63 PR 130 341 JOHN B CZIRI (LRL)  
 KOLLER 63 NC 27 1405 E L KOLLER, S TAYLOR, T HUETTER (STEVENS)  
 ALSO 66 STAMER  
 PETRUKHI 63 SIENA CONF 20 V I PETRUKHIN, YU D PROKOSHKIN (JINR)  
 VON DARO 63 PL 4 51 VON DAROEL, DEKKERS, MERMOD, VAN PUTTEN (CERN)  
 SHWE 64 PR 136B 1839 H SHWE, F M SMITH, W H BARKAS (LRL)  
 BELLETTINI 65 NC 40 A 1139 BELLETTINI, BEMPORAD, BRACCINI (PISA+FIRENZE)  
 DUCLOS 65 PL 19 253 DUCLOS, FREYTAG, HEINTZE (CERN+HEIDELBERG)  
 EVANS 65 PR 139 B 982 D A EVANS (OXFORD)  
 KUTIN 65 JETP LETT 2 243 KUTIN, PETRUKHIN, PROKOSHKIN (JINR)  
 STAMER 66 PR 151 1108 STAMER, TAYLOR, KOLLER, HUETTER (STEVENS)  
 VASILEVS 66 PL 23 281 VASILEVSKY, VISHNYAKOV, DONAITSEV (DUBNA)  
 DEVONS 69 PRL 184 1356 NEWENTH, MISSIM, SAGATI, DI CAPUA (COLL+ROMA)  
 BELLETTINI 70 NC 66A 243 BELLETTINI, BEMPORAD, LUBELSMEY (PISA+BONN)  
 KRYSHKIN 70 JETP 30 1037 \*STERLIGOV, YUSVIA (TOMSK POLYTECH. INST.)  
 ABRAMS 73 PL 45B 66 -CARROLL, KYCIA, LI, MICHAEL, MCKEET + (BNL)  
 MIYAZAKI 73 PR D8 2051 T. MIYAZAKI, E. TAKASUGI (TOKY)  
 BROWMAN 74 PRL 33 1400 +DEWIRE, EITTELKAM, HANSON (CORN+BING)  
 DAVIES 74 NC 26A 324 -GUZ, ZIYA (BIRM-RHEL+SHMP)  
 SCHACHER 77 LNC 20 177 -CZAPEK, HAHN, MARTI (CERN)  
 AUERBAC1 78 PRL 41 275 AUERBACH, HIGHLAND, JOHNSON, + (TEMP+LASL)  
 AUERBAC2 78 PL 72B 353 AUERBACH, HIGHLAND, JOHNSON, + (TEMP+LASL)  
 FISCHER1 78 PL 73B 359 +EXTERMANN, GUISSAN, MERMOD, - (GEVA+SACL)  
 FISCHER2 78 PL 73B 364 +EXTERMANN, GUISSAN, MERMOD, MOREL+ (GEVA+SACL)  
 AUERBACH 80 PL 90B 317 +HAIK, HIGHLAND, MCFARLANE, MACKE+ (TEMP+LASL)  
 HIGHLAND 80 PRL 44 628 +AUERBACH, HAIK, MCFARLANE, MACKE+ (TEMP+LASL)  
 HERCZEG 81 PL 100B 347 P. HERCZEG, C. M. HOFFMAN (LANL)  
 SCHARDT 81 PR D23 629 +FRANK, HOFFMANN, MISCHKE, MOIR + (ARZS+LANL)  
 BRYMAN 82 PR D29 2538 D. BRYMAN (TRIU)  
 MISCHKE 82 PRL 48 1153 +FRANK, HOFFMAN, MOIR, SARRACINO + (LANL+ARZS)  
 FRANK 83 PR D28 423 +HOFFMAN, MISCHKE, MOIR+ (LANL+ARZS)  
 TUPPER 83 PR D28 2908 +GROSE, SAMUEL (OXSU)  
 HERCZEG 84 PR D29 1954 +HOFFMAN (LANL)  
 ATHERTON 85 PL 158B 81 +BOVET, COET, + (CERN+ISU+LUND+LPTP+EFI)

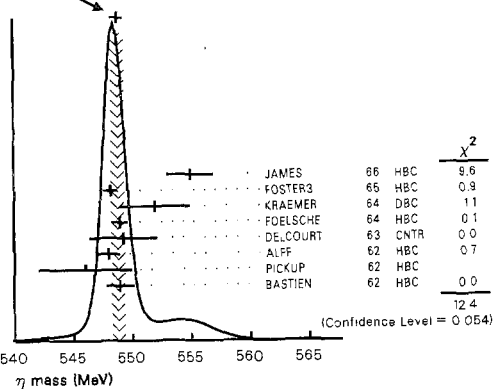
$\eta$

$$J^G(J^{PC}) = 0^+(0^{-+})$$

## $\eta$ MASS (MeV)

M	53	549.0	1.2	BASTIEN	62	HBC
M	35	546.0	4.0	PICKUP	62	HBC
M	91	548.0	1.0	ALFF	62	HBC
M	148	549.3	2.9	DEL COURT	63	CNTR
M	325	552.0	3.0	FOELSCHKE	64	HBC
M	5	548.2	0.65	KRAEMER	64	DBC
M	250	555.0	2.0	FOSTER3	65	HBC
M	1	548.2	0.56	JAMES	66	HBC
M	AVG	548.82	0.56	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4) (SEE IDEOGRAM BELOW)		

WEIGHTED AVERAGE  
 $548.82 \pm 0.58$  (ERROR SCALED BY 1.4)



## $\eta$ WIDTH

WM ETA WIDTH DETERMINED FROM MASS SPECTRUM (UNITS MEV)  
 WM 91 (10.0) OR LESS ALFF 62 HBC  
 WM 148 (10.0) OR LESS FOELSCHKE 64 HBC  
 WM 31 (12.0) OR LESS JAMES 66 HBC  
 WM (4.0) OR LESS BALTAY 66 DBC  
 WM (0.9) OR LESS JONES 66 CNTR

W ETA WIDTH DETERMINED FROM DECAY RATE (UNITS KEV)  
 W THIS IS THE PARTIAL DECAY RATE (W1) FOR THE MODE ( $\eta \rightarrow 2\gamma$ )  
 W DIVIDED BY THE FITTED BRANCHING FRACTION (P1) FOR THAT MODE.  
 W FIT 1.05 0.15 FROM FIT

## $\eta$ PARTIAL DECAY MODES

	DECAY MASSES
P1	$\eta \rightarrow 2\gamma$ 0+ 0
P2	$\eta \rightarrow 3\pi^0$ 135+ 135+ 135
P3	$\eta \rightarrow \pi^+ \pi^- \pi^0$ 140+ 140+ 135
P4	$\eta \rightarrow \pi^+ \pi^- \gamma$ 140+ 140+ 0
P5	$\eta \rightarrow e^+ e^- \pi^0$ 135+ 511+ 511
P6	$\eta \rightarrow e^+ e^- \pi^+ \pi^-$ 140+ 140+ 511+ 511
P7	$\eta \rightarrow \pi^0 2\gamma$ 135+ 0+ 0
P8	$\eta \rightarrow e^+ e^- \gamma$ 511+ 511+ 0
P10	$\eta \rightarrow \pi^+ \pi^- \pi^0 \gamma$ 140+ 140+ 135+ 0
P11	$\eta \rightarrow \pi^+ \pi^- 2\gamma$ 140+ 140+ 0+ 0
P12	$\eta \rightarrow \mu^+ \mu^-$ 106+ 106
P13	$\eta \rightarrow \mu^+ \mu^- \gamma$ 106+ 106+ 0
P14	$\eta \rightarrow \mu^+ \mu^- \pi^0$ 106+ 106+ 135
P15	$\eta \rightarrow \pi^+ \pi^-$ 140+ 140
P16	$\eta \rightarrow e^+ e^-$ 511+ 511
P17	$\eta \rightarrow \mu^+ \mu^- \pi^0 \gamma$ 106+ 106+ 135+ 0
P18	$\eta \rightarrow 3\gamma$ 0+ 0+ 0

## FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions,  $P_i$ , as follows: The diagonal elements are  $P_i \delta P_i$ , where  $\delta P_i = \sqrt{(\delta P_i)^2}$ , while the off-diagonal elements are the normalized correlation coefficients  $(\delta P_i \delta P_j) / (P_i P_j)$ . For the definitions of the individual  $P_i$ , see the listings above; only those  $P_i$  appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

For notation, see key on page 91.

# Stable Particle Full Listings

$\eta$

	P 1	P 2	P 3	P 4	P 7	P 8
P 1	.3894+-0.0640					
P 2	.458	.3190+-0.0035				
P 3	-.856	-.809	.2368+-0.0053			
P 4	-.745	-.704	.798	.0491+-0.0013		
P 7	.026	.021	-.045	-.040	.00078+-0.00012	
P 8	-.110	-.105	-.055	-.054	-.006	.0050+-0.0012

### FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space; i.e.  $G_i = \Gamma_i = \Gamma_{total} P_i$ , in appropriate units. In analogy to the matrix above, the diagonal elements are  $G_i = \delta G_i$ , where  $\delta G_i = \sqrt{(\delta G_i \delta G_i)}$ , while the off-diagonal elements are the normalized correlation coefficients  $(\delta G_i \delta G_j) / (\delta G_i \delta G_j)$ . Note that, because of the error in  $\Gamma_{total}$ , the errors and correlations here are not directly derivable from those above.

	G 1	G 2	G 3	G 4	G 7	G 8
G 1	.4082+-0.0595					
G 2	.996	.334+-0.0489				
G 3	.978	.978	.2482+-0.0369			
G 4	.974	.974	.998	.0515+-0.0077		
G 7	.682	.682	.672	.469	.00081+-0.00017	
G 8	.525	.525	.521	.518	.361	.0052+-0.0014

### NOTE ON THE $\eta$ WIDTH TO TWO PHOTONS

(by S. Cooper, SLAC)

Measurements of  $\Gamma(\eta \rightarrow \gamma\gamma)$  via Primakoff effect and via two-photon production disagree, as shown in the table below. The discrepancy is not yet understood.

$\Gamma(\eta \rightarrow \gamma\gamma)$ (keV)	Technique	Reference
$(1.00 \pm 0.22)^*$	Primakoff	BEMPORAD 67
$0.324 \pm 0.046$	Primakoff	BROWMAN 74
$0.56 \pm 0.16$	two-photon	WEINSTEIN 83 (Crystal Ball-SPEAR)
$0.53 \pm 0.04 \pm 0.04$	two-photon	BARTEL 85 (JADE)
$0.64 \pm 0.14 \pm 0.13$	two-photon	TPC/Two Gamma <sup>1</sup>
$0.58 \pm 0.02 \pm 0.06$ (preliminary)	two-photon	Crystal Ball-DORIS <sup>2</sup>
-----		
$0.56 \pm 0.04$	two-photon	average

\*corrected for  $B(\eta \rightarrow \gamma\gamma) = 0.38 \pm 0.01$

The Primakoff effect is production of  $\eta$ 's in the Coulomb field of a nucleus. Background comes from  $\eta$ 's produced in its hadronic field. They are separated by fitting the  $\eta$  angular distribution, which is calculated to be different for the two effects. The systematic error of this calculation may have been underestimated. BROWMAN 74 has better angular resolution, higher statistics, and a larger range of beam energies than BEMPORAD 67. Both the Primakoff and two-photon techniques can be "calibrated" at the  $\pi^0$ . BROWMAN 74's measurement  $\Gamma(\pi^0 \rightarrow \gamma\gamma) = 8.0 \pm 0.4^3$  agrees with  $7.3 \pm 0.2 \pm 0.1$  from the very accurate  $\pi^0$  lifetime measurement of Atherton et al.<sup>4</sup>

The two-photon calculation to get  $\Gamma(\eta \rightarrow \gamma\gamma)$  from the observed rate of  $e^+e^- \rightarrow e^+e^-\eta$  is pure QED and noncontroversial. These experiments are in good agreement with each other. A preliminary two-photon measurement  $\Gamma(\pi^0 \rightarrow \gamma\gamma) = 7.9 \pm 1.4 \pm 1.6^2$  agrees with Atherton et al., but the error is large.

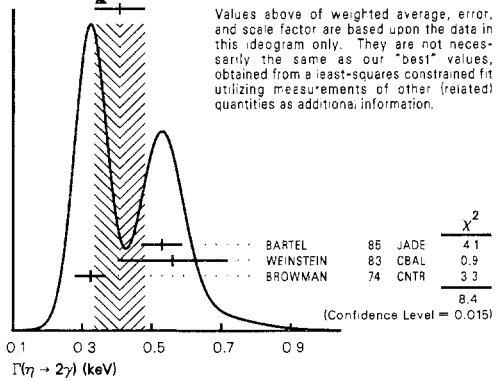
### References

- H. Aihara et al. (TPC/Two Gamma), UCLA-85-013 (1985).
- K. Wacker (Crystal Ball), presented to *International Europhysics Conference on High Energy Physics* (Bari, Italy, 1985).
- A. Browman et al., Phys. Rev. Lett. **33**, 1400 (1974).
- H.W. Atherton et al., Phys. Lett. **158B**, 81 (1985).

### $\eta$ DECAY RATES

$\eta \rightarrow 2\gamma$	(keV)				(G1)
W1 B	(1.00)	(0.22)	BEMPORAD 67 CNTR	PRIMAKOFF EFFECT	
W1	0.324	0.046	BROWMAN 74 CNTR	PRIMAKOFF EFFECT	
W1 J	0.56	0.16	WEINSTEIN 83 CBAL	E+ E- -> E+ E- ETA	
W1 J	0.53	0.06	BARTEL 85 JADE E+	E- -> E+ E- ETA	
W1 B	BEMPORAD 67 GIVES W1=1.21+-0.26 KEV ASSUMING THAT W1/TOTAL=0.314.				
W1 B	BEMPORAD PRIVATE COMMUNICATION GIVES MORE GENERAL RESULT AS				
W1 B	W1*W1/TOTAL=0.38+-0.083. WE EVALUATE THIS USING W1/TOTAL=0.38+-0.01.				
W1 B	NOT INCLUDED IN AVERAGE BECAUSE THE UNCERTAINTY RESULTING FROM THE				
W1 B	SEPARATION OF THE COULOMB AND NUCLEAR AMPLITUDES HAS APPARENTLY				
W1 B	BEEN UNDERESTIMATED.				
W1 J	STAT ERROR (0.04) AND SYST ERROR (0.04) ADDED IN QUADRATURE.				
W1 AVE	0.408	0.073	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.0)		
W1 FIT	0.408	0.059	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7) (SEE IDEOGRAM BELOW)		

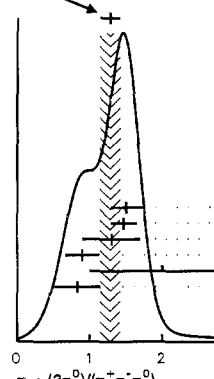
WEIGHTED AVERAGE  
0.408 ± 0.073 (ERROR SCALED BY 2.0)



### $\eta$ BRANCHING RATIOS

$\eta \rightarrow (2\gamma + 3\pi^0 + \pi^0 2\gamma) / (\pi^+ \pi^- \pi^0 + \pi^+ \pi^- \gamma + e^+ e^- \gamma)$		
R1 N	10 (2.5)	(1.0) PICKUP (P1+P2+P7)/(P3+P4+P8)
R1 N	53 (3.20)	(1.26) BASTIEN 62 HBC
R1 N	280 (4.5)	(1.0) JAMES 66 HBC
R1 N	THESE EXPERIMENTS HAVE NOT BEEN USED IN COMPUTING THE AVERAGES	
R1 N	AS THEY WERE UNABLE TO SEPARATE CLEARLY PARTIAL MODES (3) AND (4)	
R1 N	FROM EACH OTHER. THE REPORTED VALUES THUS PROBABLY CONTAIN	
R1 N	SOME (UNKNOWN) FRACTION OF MODE (4).	
R1	2.64	0.23 BALTAY2 67 DBC
R1 FIT	2.438	0.076 FROM FIT

Stable Particle Full Listings

$\eta \rightarrow 2\gamma/(\pi^+\pi^-\pi^0 + \pi^+\pi^-\gamma + e^+e^-\gamma)$		(P1)/(P3+P4+P8)		
R2	0.99	0.48	CRAWFORD 63 HBC	
R2	75	1.51	KENDALL 74 OSPK	
R2 AVG	1.10	0.43	AVERAGE	
R2 FIT	1.339	0.042	FROM FIT	
$\eta \rightarrow (\pi^0 2\gamma)/(2\gamma + 3\pi^0 + \pi^0 2\gamma)$		(P7)/(P1+P2+P7)		
R3	OTHER RESULTS ARE IN SECTION R22.			
R3	(0.04) OR LESS	CL=.90	ABROSIMOV 80 HLBC	
R3	0.0010	0.0002	ALDE 84 CNTR	
R3 FIT	0.00109	0.00017	FROM FIT	
$\eta \rightarrow (\pi^+\pi^-\gamma)/(\pi^+\pi^-\pi^0)$		(P4)/(P3)		
R4	0.14	0.08	FOELSCH 64 HBC	
R4	24	(0.73) (0.25)	PAULI 64 DBC	
R4	0.30	0.06	CRAWFORD 66 HBC	
R4	0.10	0.10	KRAEMER 64 DBC	
R4	0.196	0.041	FOSTER 65 HBC	
R4	0.25	0.035	LITCHEFIEL 67 DBC	
R4	0.28	0.04	BALTAY 67 DBC	
R4	7250	0.201	GORMLEY 70 ASPK	
R4	18K	0.209	THALER 73 ASPK	
R4 AVG	0.2074	0.0037	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)	
R4 FIT	0.2074	0.0033	FROM FIT	
$\eta \rightarrow 3\pi^0/2\gamma$		(P2)/(P1)		
R6	(0.90) OR MORE		CHRETIEN 62 PBC	
R6	(1.25) (0.39)		BACCI 63 CNTR INVERSE BR REPORTED	
R6	0.88	0.16	BALTAY 1 67 DBC	
R6	1.1	0.2	CENCE 67 OSPK	
R6	0.91	0.14	COX 70 HBC	
R6	0.75	0.09	DEVONS 70 OSPK	
R6 AVG	0.842	0.065	AVERAGE	
R6 FIT	0.8193	0.0090	FROM FIT	
$\eta \rightarrow 2\gamma/(\pi^+\pi^-\pi^0)$		(P1)/(P3)		
R7	1.61	0.39	FOSTER 1 65 HBC	
R7	401	1.72	BAGLIN 69 HLBC	
R7 AVG	1.61	0.21	AVERAGE	
R7 FIT	1.644	0.052	FROM FIT	
$\eta \rightarrow (2\gamma + 3\pi^0 + \pi^0 2\gamma)/(\pi^+\pi^-\pi^0)$		(P1+P2+P7)/(P3)		
R8	50	3.6	0.8	KRAEMER 64 DBC
R8		3.8	1.1	PAULI 64 DBC
R8		2.89	0.56	ALFF-STEI 66 HBC
R8	244	5.6	0.6	FLATTE 67 HBC
R8	29	3.4	1.1	AGUILAR-B 72 HBC
R8 B	70	2.83	0.80	BLOODWORTH 72 HBC
R8		2.56	0.89	KENDALL 74 OSPK
R8 B	ERROR INCREASED FROM PUBLISHED VALUE 0.5 BY BLOODWORTH, PRIV. COMM.			
R8 AVG	3.26	0.30	AVERAGE	
R8 FIT	2.995	0.094	FROM FIT	
$\eta \rightarrow (e^+e^-\pi^0)/(\pi^+\pi^-\pi^0)$ (units 10 <sup>-4</sup> )		(P5)/(P3)		
R9	SINGLE PHOTON PROCESS FORBIDDEN BY C-PARITY			
R9	(110.) OR LESS		PRICE 65 HBC	
R9	0 (77.) OR LESS		FOSTER 2 65 HBC	
R9	(2.) OR LESS	CL=.90	BAGLIN 1 67 HLBC	
R9	0 (16.) OR LESS	CL=.90	BILLING 67 HLBC	
R9	1.9	OR LESS	CL=.90 JANE 1 75 OSPK	
$\eta \rightarrow (e^+e^-\pi^+\pi^-)/total$ (units 10 <sup>-2</sup> )		(P6)		
R10	(0.7) OR LESS		RITTENBER 65 HBC	
$\eta \rightarrow (e^+e^-\pi^+\pi^-)/(\pi^+\pi^-\gamma)$		(P6)/(P4)		
R11	1	0.026	0.026	GROSSMAN 66 HBC
$\eta \rightarrow 2\gamma/(2\gamma + 3\pi^0 + \pi^0 2\gamma)$		(P1)/(P1+P2+P7)		
R12	(0.416) (0.044)		DIGIUGNO 66 CNTR ERROR DOUBLED	
R12	(0.44) (0.07)		GRUNHAUS 66 OSPK	
R12	(0.579) (0.052)		FELDMAN 67 OSPK	
R12 T	(0.39) (0.06)		JONES 66 CNTR	
R12 T	THIS RESULT FROM COMBINING CROSS SECTIONS FROM TWO DIFFERENT EXPTS.			
R12	0.59	0.035	BUNIA TOV 67 OSPK	
R12	0.535	0.018	BUTTRAM 70 OSPK	
R12	(0.57) (0.09)		STRIGALSK 71 HLBC	
R12	113	0.60	0.14	KENDALL 74 OSPK
R12	88	0.52	0.09	ABROSIMOV 80 HLBC
R12		0.549	0.004	ALDE 84 CNTR
R12 AVG	0.5489	0.0039	AVERAGE	
R12 FIT	0.5491	0.0027	FROM FIT	
$\eta \rightarrow 3\pi^0/(2\gamma + 3\pi^0 + \pi^0 2\gamma)$		(P2)/(P1+P2+P7)		
R13	(0.209) (0.054)		DIGIUGNO 66 CNTR ERROR DOUBLED	
R13	(0.225) (0.10)		GRUNHAUS 66 OSPK	
R13	(0.177) (0.035)		FELDMAN 67 OSPK	
R13	(0.41) (0.035)		BUNIA TOV 67 OSPK NOT INDEP. OF R12	
R13	0.439	0.024	BUTTRAM 70 OSPK	
R13	(0.32) (0.09)		STRIGALSK 71 HLBC	
R13	75	0.44	0.08	ABROSIMOV 80 HLBC
R13		0.450	0.004	ALDE 84 CNTR
R13 AVG	0.4497	0.0039	AVERAGE	
R13 FIT	0.4498	0.0027	FROM FIT	
$\eta \rightarrow (e^+e^-\pi^0)/total$ (units 10 <sup>-2</sup> )		(P5)		
R15	SINGLE PHOTON PROCESS FORBIDDEN BY C-PARITY			
R15	(0.7) OR LESS		RITTENBER 65 HBC	
R15	(0.084) OR LESS	CL=.90	BAZIN 68 DBC	
R15	0 (0.016) OR LESS	CL=.90	MARTYNOV 76 HLBC	
$\eta \rightarrow (\pi^+\pi^-\pi^0)/(\pi^+\pi^-\pi^0)$ (units 10 <sup>-2</sup> )		(P10)/(P3)		
R17	(7.0) OR LESS		FLATTE 67 HBC	
R17	(0.9) OR LESS		PRICE 67 HBC	
R17	(1.6) OR LESS	CL=.95	BALTAY 2 67 DBC	
R17	(1.7) OR LESS	CL=.90	ARNOLD 68 HLBC	
R17	0	0.24	OR LESS CL=.90 THALER 73 ASPK	
$\eta \rightarrow (\pi^+\pi^-\pi^0)/(\pi^+\pi^-\pi^0)$		(P11)/(P3)		
R18	0.009	OR LESS	PRICE 67 HBC	
R18	(0.016) OR LESS	CL=.95	BALTAY 2 67 DBC	
$\eta \rightarrow 3\pi^0/(\pi^+\pi^-\pi^0)$		(P2)/(P3)		
R19	0.83	0.32	CRAWFORD 63 HBC	
R19	2.0	1.0	FOELSCH 64 HBC	
R19	0.90	0.24	FOSTER 1 65 HBC	
R19	1.3	0.4	BAGLIN 2 67 HLBC	
R19	1.47	0.20	0.17 BULLOCK 68 HLBC	
R19	1.50	0.15	0.29 BAGLIN 69 HLBC	
R19 AVG	1.28	0.14	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)	
R19 FIT	1.347	0.043	FROM FIT (SEE IDEOGRAM BELOW)	
<b>WEIGHTED AVERAGE 1.28 ± 0.14 (ERROR SCALED BY 1.3)</b>				
				
<p>Values above of weighted average, error, and scale factor are based upon the data in this ideogram only. They are not necessarily the same as our "best" values, obtained from a least-squares constrained fit utilizing measurements of other (related) quantities as additional information.</p>				
$\eta \rightarrow (2\gamma + 3\pi^0 + \pi^0 2\gamma)/total$		(P1+P2+P7)		
R21	16K	0.79	0.08	BUNIA TOV 67 OSPK
R21		0.705	0.008	BASILE 71 CNTR MM SPECTROMETER
R21 AVG		0.7058	0.0080	AVERAGE
R21 FIT		0.7091	0.0064	FROM FIT
$\eta \rightarrow (\pi^0 2\gamma)/total$		(P7)		
R22	0	(0.003) OR LESS	CL=.90	DAVYDOV 81 CNTR PI- P-->ETA N
R22	70	0.0095	0.0023	BINON 82 CNTR PI- P-->ETA N
R22 FIT		0.00078	0.00012	FROM FIT
$\eta \rightarrow \mu^+\mu^-/total$ (units 10 <sup>-5</sup> )		(P12)		
R23	0 (2.) OR LESS		CL=.95	WEHMANN 68 OSPK
R23	27	0.65	0.21	DZHEL'YAZ 80 SPEC PI- P-->ETA N
$\eta \rightarrow \mu^+\mu^-\pi^0/total$ (units 10 <sup>-4</sup> )		(P14)		
R24	SINGLE PHOTON PROCESS FORBIDDEN BY C-PARITY			
R24	(5.) OR LESS		WEHMANN 68 OSPK	
R24	0.05	OR LESS	CL=.90	DZHEL'YADI 81 SPEC PI- P-->ETA N
$\eta \rightarrow \mu^+\mu^-/2\gamma$ (units 10 <sup>-5</sup> )		(P12)/(P1)		
R25	(5.9)	(2.2)		HYAMS 69 OSPK
$\eta \rightarrow (\pi^+\pi^-)/total$ (units 10 <sup>-2</sup> )		(P15)		
R27	VIOLATES P AND CP INVARIANCE			
R27	0	0.15	OR LESS	THALER 73 ASPK CON. LEV. NOT GIVEN
$\eta \rightarrow (e^+e^-)/(\pi^+\pi^-\pi^0)$ (units 10 <sup>-2</sup> )		(P8)/(P3)		
R28	J 80	2.1	0.5	JANE 2 75 OSPK
R28	J	VALUE CHANGED BY ERRATUM.		
R28 FIT		2.10	0.50	FROM FIT
$\eta \rightarrow (e^+e^-)/total$ (units 10 <sup>-4</sup> )		(P16)		
R29 D	3.	OR LESS	CL=.90	DAVIES 74 RVUE
R29 D	DAVIES 74 EXTRACTS THIS INFORMATION FROM ESTEN 67.			
$\eta \rightarrow (\mu^+\mu^-)/total$ (units 10 <sup>-4</sup> )		(P13)		
R30	100	(1.5)	(0.75)	BUSHNIN 78 SPEC REPL BY DZHEL'YADI 80
R30	600	3.1	0.4	DZHEL'YADI 80 SPEC PI- P-->ETA N
$\eta \rightarrow (\mu^+\mu^-\pi^0)/total$ (units 10 <sup>-6</sup> )		(P17)		
R31	3.	OR LESS	CL=.90	DZHEL'YADI 81 SPEC PI- P-->ETA N
$\eta \rightarrow (3\gamma)/(2\gamma + 3\pi^0 + \pi^0 2\gamma)$		(P18)/(P1+P2+P7)		
R33	FORBIDDEN BY C INVARIANCE			
R33	7.E-4	OR LESS	CL=.95	ALDE 84 CNTR

For notation, see key on page 91.

## Stable Particle Full Listings

$\eta$

### NOTE ON $\eta$ DECAY PARAMETERS

#### C violation in $\eta$ decays

As a test of possible C violation in electromagnetic interactions, a number of experiments have looked for possible charge asymmetries in the decays  $\eta \rightarrow \pi^+ \pi^- \pi^0$  and  $\eta \rightarrow \pi^+ \pi^- \gamma$ . We list the following parameters:

(a) The left-right asymmetry

$$A = (N^+ - N^-)/(N^+ + N^-),$$

where  $N^\pm$  means the number of events with the  $\pi^\pm$  energy greater than the  $\pi^\mp$  energy in the  $\eta$  rest frame.

(b) The sextant asymmetry

$$A_s = \frac{N_1 + N_3 + N_5 - N_2 - N_4 - N_6}{N_1 + N_2 + N_3 + N_4 + N_5 + N_6}$$

for the decay  $\eta \rightarrow \pi^+ \pi^- \pi^0$ . The numbers refer to sextants of the Dalitz plot (see, for example, Layter et al.<sup>1</sup>).  $A_s$  is sensitive to an  $I = 0$  C-violating asymmetry.

(c) The quadrant asymmetry  $A_q$ , defined in a similar way as  $A_s$ , but with each sector of the Dalitz plot now containing  $\pi/2$  rather than  $\pi/3$  radians.  $A_q$  is sensitive to an  $I = 2$  C-violating final state.

(d) The D-wave contribution to the C-violating amplitude in the decay  $\eta \rightarrow \pi^+ \pi^- \gamma$ . The upper limit for this contribution is measured by the parameter  $\beta$ , defined by

$$dN/d|\cos\theta| \propto \sin^2\theta(1 + \beta \cos^2\theta),$$

where  $\theta$  is the angle between the  $\pi^+$  and the  $\gamma$  in the dipion center of mass. A term proportional to  $\cos^2\theta$  could also be due to P- and F-wave interference.

We list  $A$  for the decay modes  $\eta \rightarrow \pi^+ \pi^- \pi^0$  and  $\eta \rightarrow \pi^+ \pi^- \gamma$ ,  $A_s$  and  $A_q$  for the decay  $\eta \rightarrow \pi^+ \pi^- \pi^0$ , and  $\beta$  for the decay  $\eta \rightarrow \pi^+ \pi^- \gamma$  in the Full Listings below.

#### Dalitz plot for $\eta \rightarrow \pi^+ \pi^- \pi^0$

The Dalitz plot for the decay  $\eta \rightarrow \pi^+ \pi^- \pi^0$  may be fit by the distribution

$$|M(x,y)|^2 \propto 1 + ay + by^2 + cx + dx^2 + exy.$$

Here,

$$x = \sqrt{3}(T_+ - T_-)/Q, \quad y = (3T_0/Q) - 1,$$

$T_+$ ,  $T_-$ , and  $T_0$  are the kinetic energies of the  $\pi^+$ ,  $\pi^-$ , and  $\pi^0$  in the  $\eta$  rest system, and  $Q = m_\eta - m_{\pi^+} - m_{\pi^-} - m_{\pi^0}$ . The coefficient of the term linear in  $x$  is

sensitive to C-violation due to an  $I = 0$  or  $I = 2$  final state. We list papers presenting determinations of the parameters  $a$ ,  $b$ ,  $c$ , and  $d$  in the section DP below. However, we do not tabulate values of these parameters because the assumptions made by different authors are not compatible and do not allow comparison of the numerical values.

#### Dalitz plot for $\eta \rightarrow \pi^0 \pi^0 \pi^0$

The Dalitz plot for the decay  $\eta \rightarrow \pi^0 \pi^0 \pi^0$  may be fit to the expression

$$|M|^2 \propto 1 + 2\alpha z,$$

where

$$z = \frac{2}{3} \sum_{i=1}^3 [3(m_\eta - 3m_\pi)^{-1}(E_i - \frac{1}{3}m_\eta)]^2 = \rho^2/\rho_{\max}^2.$$

Here  $E_i$  is the energy of the  $i^{\text{th}}$  pion in the  $\eta$  rest frame, and  $\rho$  is the distance to the center of the Dalitz plot. We list the parameter  $\alpha$  in section A0 below.

#### Reference

1. J.G. Layter et al., Phys. Rev. Lett. **29**, 316 (1972).

#### $\eta$ C-NONCONSERVING DECAY PARAMETERS

$\pi^+ \pi^- \pi^0$ LEFT-RIGHT ASYMMETRY PARAMETER (units $10^{-2}$ )					
A1	1351	7.2	2.8	BALTAY	66 DBC
A1	1300	5.8	3.4	CLPWY	66 HBC
A1	10665	(0.3)	(1.0)	CNOPS	66 OSPK REPL BY MULLER 69
A1	705	-6.1	4.0	LARRIBE	66 HBC
A1	G36800	(1.5)	(0.5)	GORMLEY3	68 ASPK
A1	10709	0.3	1.1	MULLER	69 OSPK
A1	1138	-4	3	CARPENTER	70 HBC
A1	349	3.2	5.4	DANBURG	70 DBC
A1	220K	-0.05	0.22	LAYTER	72 ASPK
A1	165K	0.28	0.26	JANE1	74 OSPK
A1	G	GORMLEY3 68 ASYMMETRY PROBABLY DUE TO UNMEASURED (E X B) SPK. CH.			
A1	G	EFFECTS. NEW EXPTS. WITH (E X B) CONTROLS DONT OBSERVE ASYMMETRY.			
A1	AVG	0.12	0.17	AVERAGE	

$\pi^+ \pi^- \gamma$ LEFT-RIGHT ASYMMETRY PARAMETER (units $10^{-2}$ )						
A2	33	-2.	17.	DRAWFORD	66 HBC	
A2	M	1620	1.5	2.5	LITCHFIEL	67 DBC
A2	7257	1.22	1.56	MULLER	69 OSPK	
A2	36K	0.5	0.6	GORMLEY	70 ASPK	
A2	35K	1.2	0.6	THALER	72 ASPK	
A2	M	MULLER 69 IS SENSITIVE ONLY TO UPPER .4 OF GAMMA-RAY SPECTRUM.				
A2	AVG	0.88	0.40	AVERAGE		

$\pi^+ \pi^- \pi^0$ SEXTANT ASYMMETRY PARAMETER (units $10^{-2}$ )					
AS	1300	6.8	3.3	CLPWY	66 HBC
AS	705	-2.4	4.0	LARRIBE	66 HBC
AS	37K	0.5	0.5	GORMLEY3	68 WIRE
AS	220K	0.10	0.22	LAYTER	72 ASPK
AS	165K	0.20	0.25	JANE1	74 OSPK
AS	AVG	0.19	0.16	AVERAGE	

$\pi^+ \pi^- \pi^0$ QUADRANT ASYMMETRY PARAMETER (units $10^{-2}$ )					
AQ	220K	-0.07	0.22	LAYTER	72 ASPK
AQ	165K	-0.30	0.25	JANE1	74 OSPK
AQ	AVG	-0.17	0.17	AVERAGE	

#### $\beta$ FOR $\eta \rightarrow \pi^+ \pi^- \gamma$ . SENSITIVE TO D-WAVE CONTRIBUTION.

BET	DN/DCOS THETA = SIN**2 THETA * (1 + BETA * COS**2 THETA)					
BET	7250	-0.060	0.065	GORMLEY	70 WIRE	
BET	L	0.12	0.06	THALER	72 ASPK	
BET	L	35K	0.11	0.11	JANE1	74 OSPK
BET	L	ALTHOUGH DONT BELIEVE THIS TO INDICATE D-WAVE BECAUSE DEPENDENCE OF BETA ON GAMMA ENERGY INCONSISTENT WITH THEOR. PREDICTION.				
BET	L	COS**2 DEPENDENCE MAY ALSO COME FROM P AND F-WAVE INTERFERENCE.				
BET	AVG	0.047	0.062	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)		
BET				(SEE IDEOGRAM BELOW)		

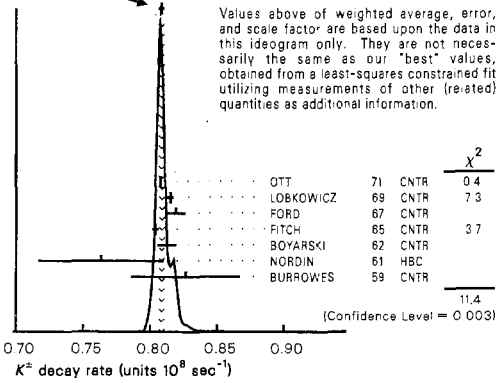


For notation, see key on page 91.

# Stable Particle Full Listings

$K^\pm$

WEIGHTED AVERAGE  
0.8084 ± 0.0021 (ERROR SCALED BY 2.4)



### FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions,  $P_i$ , as follows: The diagonal elements are  $P_i \pm \delta P_i$ , where  $\delta P_i = \sqrt{(\delta P_i^2)}$ , while the off-diagonal elements are the normalized correlation coefficients  $(\delta P_i \delta P_j) / (\delta P_i \delta P_j)$ . For the definitions of the individual  $P_i$ , see the listings above; only those  $P_i$  appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3	P 4	P 5	P 6
P 1	.6351±.0016					
P 2	-.720	-.2117±.0015				
P 3	-.175	-.021	-.0559±.0003			
P 4	-.158	.050	.223	.0173±.0005		
P 5	-.291	-.252	-.165	-.357	.0318±.0010	
P 6	-.340	-.145	.141	.004	.221	.0482±.0005

### FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space; i.e.,  $G_i \equiv \Gamma_i = \Gamma_{total} P_i$ , in appropriate units. In analogy to the matrix above, the diagonal elements are  $G_i \pm \delta G_i$ , where  $\delta G_i = \sqrt{(\delta G_i^2)}$ , while the off-diagonal elements are the normalized correlation coefficients  $(\delta G_i \delta G_j) / (\delta G_i \delta G_j)$ . Note that, because of the error in  $\Gamma_{total}$ , the errors and correlations here are not directly derivable from those above.

	G 1	G 2	G 3	G 4	G 5	G 6
G 1	.5334±.0017					
G 2	-.317	-.1711±.0013				
G 3	-.086	.012	-.0452±.0002			
G 4	-.090	.063	.226	.0140±.0004		
G 5	-.181	-.219	-.173	-.353	.0257±.0008	
G 6	-.157	-.087	.143	.009	.227	.0390±.0004

### ( $K^+ - K^-$ )/AVG., MEAN LIFE DIFFERENCE (%)

THIS QUANTITY IS A MEASURE OF CPT INVARIANCE IN W.1.

DT	0.47	0.30	FORD	67 CNTR
DT	0.090	0.078	LOBKOWICZ	69 CNTR
DT				
DT	AVG	0.114	0.093	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)

### $K^\pm$ PARTIAL DECAY MODES

	DECAY MODES	DECAY MASSES
P1	$K^\pm \rightarrow \mu^\pm \nu$	106+ 0
P2	$K^\pm \rightarrow \pi^\pm \pi^0$	140+ 135
P3	$K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$	140- 140+ 140
P4	$K^\pm \rightarrow \pi^\pm 2\pi^0$	140+ 135+ 135
P5	$K^\pm \rightarrow \pi^0 \mu^\pm \nu$ (called $K_{\mu 3}$ )	135+ 106+ 0
P6	$K^\pm \rightarrow \pi^0 e^\pm \nu$ (called $K_{e 3}$ )	135+ 511+ 0
P7	$K^+ \rightarrow \pi^+ \pi^- e^+ \nu$ (called $K_{e 4}$ )	140+ 140+ 511+ 0
P8	$K^+ \rightarrow \pi^+ \pi^+ e^- \nu$	140+ 140+ 511+ 0
P9	$K^+ \rightarrow \pi^+ \pi^- \mu^+ \nu$ (called $K_{\mu 4}$ )	140+ 140+ 106+ 0
P10	$K^+ \rightarrow \pi^+ \pi^+ \mu^- \nu$	140+ 140+ 106+ 0
P11	$K^\pm \rightarrow e^\pm \nu$	511+ 0
P12	$K^\pm \rightarrow \mu^\pm \nu \gamma$	106+ 0+ 0
P13	$K^\pm \rightarrow \pi^\pm \pi^0 \gamma$	140+ 135+ 0
P14	$K^\pm \rightarrow \pi^\pm \pi^+ \pi^- \gamma$	140+ 140+ 140+ 0
P15	$K^\pm \rightarrow \pi^\pm e^+ e^-$	140+ 511+ 511
P16	$K^\pm \rightarrow \pi^\pm \mu^+ \mu^-$	140+ 106+ 106
P17	$K^\pm \rightarrow \pi^\pm \gamma \gamma$	140+ 0+ 0
P18	$K^\pm \rightarrow \pi^0 \mu^\pm \nu \gamma$	135+ 511+ 0+ 0
P19	$K^\pm \rightarrow \pi^0 e^\pm \nu \gamma$	140+ 511+ 511
P20	$K^\pm \rightarrow \pi^\pm \nu \nu$	140+ 0+ 0
P21	$K^\pm \rightarrow e^\pm \nu \gamma$	511+ 0+ 0
P22	$K^\pm \rightarrow \pi^\pm \gamma$	140+ 0
P23	$K^\pm \rightarrow \pi^\pm 3\gamma$	140+ 0+ 0+ 0
P24	$K^\pm \rightarrow \pi^0 \pi^0 e^\pm \nu$	135+ 135+ 511+ 0
P25	$K^+ \rightarrow \pi^- e^+ \mu^+$	140+ 511+ 106
P26	$K^+ \rightarrow \pi^+ e^+ \mu^-$	140+ 511+ 106
P27	$K^\pm \rightarrow \mu^\pm \nu \nu$	106+ 0+ 0+ 0
P28	$K^\pm \rightarrow \pi^0 \mu^\pm \nu \gamma$	135+ 106+ 0+ 0
P29	$K^+ \rightarrow \pi^+ \mu^+ e^-$	140+ 106+ 511
P30	$K^\pm \rightarrow \mu^\pm \nu e^+ e^-$	106+ 0+ 511+ 511
P31	$K^\pm \rightarrow \mu^\pm \nu e^\pm e^\pm$	106+ 0+ 511+ 511
P32	$K^\pm \rightarrow \nu e^\pm e^+ e^-$	0+ 511+ 511+ 511
P33	$K^\pm \rightarrow e^\pm \nu \nu$	511+ 0+ 0+ 0
P34	$K^+ \rightarrow \mu^+ \nu_e$	106+ 0
P35	$K^+ \rightarrow \mu^+ \bar{\nu}_e$	106+ 0
P36	$K^+ \rightarrow \pi^0 e^+ \bar{\nu}_e$	135+ 511+ 0

### $K^\pm$ CONSTRAINED FIT

OVERALL FIT OF MEAN LIFE, WIDTHS AND BRANCHING RATIOS USES 59 DATA POINTS TO DETERMINE SIX QUANTITIES. OVERALL FIT HAS  $\chi^2_{min} = 78.0$ . MAIN CONTRIBUTION (13.2) COMES FROM R19 OF HAIDT 71 (WE SEE NO REASON TO REJECT THIS EXPERIMENT AT THIS TIME)

### $K^\pm$ DECAY RATES

$K^\pm \rightarrow \mu^\pm \nu$ (units $10^6 \text{ sec}^{-1}$ )		(G1)
W1	51.2 0.8 FORD 67 CNTR --	
W1	FIT 51.34 0.17 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)	
$K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ (units $10^6 \text{ sec}^{-1}$ )		(G3)
W2 F	(4.496) (0.030) FORD 67 CNTR +- SEE NOTE F	
W2 F 3.2M	(4.529) (0.032) FORD 70 ASPK SEE NOTE F	
W2 F	4.511 0.024 FORD 70 ASPK SEE NOTE F	
W2 F	THE LAST IS THE COMBINED RESULT OF FORD 67 AND FORD 70	
W2	FIT 4.518 0.023 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)	

### ( $K^+ - K^-$ )/AVG., DECAY RATE DIFFERENCE (%)

$K^\pm \rightarrow \mu^\pm \nu$ RATE DIFFERENCE (%)		((G1+)-(G1-))/G1
D1	TEST OF CPT CONSERVATION	
D1	-0.54 0.41 FORD 67 CNTR	
$K^\pm \rightarrow \pi^\pm \pi^+ \pi^-$ RATE DIFFERENCE (%)		((G3+)-(G3-))/G3
D2	TEST OF CP CONSERVATION	
D2	-0.50 0.90 FLETCHER 67 OSPK	
D2 F	(-0.04) (0.21) FORD 67 CNTR SEE NOTE F	
D2 F 3.2M	(0.10) (0.14) FORD 70 ASPK SEE NOTE F	
D2 F	0.08 0.12 FORD 70 ASPK SEE NOTE F	
D2 S	(-0.02) (0.16) SMITH 73 ASPK +-	
D2 F	SECOND FORD 70 VALUE IS FIRST FORD 70 COMBINED WITH FORD 67.	
D2 S	SMITH 73 VALUE OF D2 IS DERIVED FROM SMITH 73 VALUE OF D3.	
D2	AVG 0.07 0.12 AVERAGE	
$K^\pm \rightarrow \pi^\pm 2\pi^0$ RATE DIFFERENCE (%)		((G4+)-(G4-))/G4
D3	TEST OF CP CONSERVATION	
D3	1802 -1.1 1.8 HERZO 69 OSPK	
D3	0.08 0.58 SMITH 73 ASPK +-	
D3	AVG -0.03 0.55 AVERAGE	
$K^\pm \rightarrow \pi^\pm \pi^0$ RATE DIFFERENCE (%)		((G2+)-(G2-))/G2
D4	TEST OF CPT CONSERVATION	
D4	0.8 1.2 HERZO 69 OSPK	
$K^\pm \rightarrow \pi^\pm \pi^0 \gamma$ RATE DIFFERENCE (%)		((G13+)-(G13-))/G13
D5	TEST OF CP CONSERVATION	
D5	24 0.0 24.0 EDWARDS 72 OSPK PI KE 58-90 MEV	
D5	4000 1.0 4.0 ABRAMS 73 ASPK +- PI KE 51-100 MEV	
D5	2461 0.8 5.8 SMITH 76 WIRE +- PI+KE 55-90 MEV	
D5	AVG 0.9 3.3 AVERAGE	

### $K^\pm$ BRANCHING RATIOS

$K^\pm \rightarrow (\mu^\pm \nu)/\text{total}$ (units $10^{-2}$ )		(P1)
R1	0 (58.3) (3.0) BIRGE 56 EMUL +	
R1	0 (56.2) (2.6) ALEXANDER 57 EMUL -	
R1	0 OLD EXPERIMENTS NOT INCLUDED IN AVERAGING	
R1	62K 63.24 0.44 CHIANG 72 OSPK + 1.84 GEV/C K-	
R1	FIT 63.51 0.16 FROM FIT	



# Stable Particle Full Listings

$K^\pm$

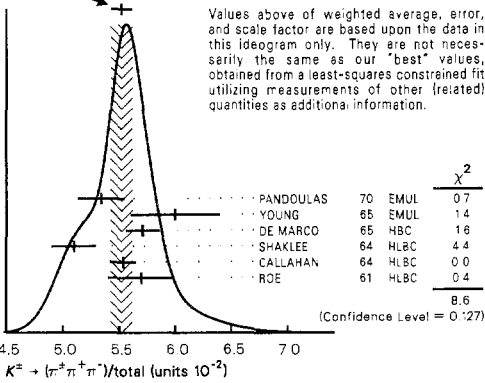
$K^\pm \rightarrow (\pi^\pm \pi^0)/\text{total (units } 10^{-2})$  (P2)

R2 0	(27.7)	(2.7)	BIRGE	56 EMUL +
R2 0	(23.2)	(2.2)	ALEXANDER	57 EMUL +
R2 0	EARLIER EXPERIMENTS NOT AVERAGED			
R2	(21.0)	(0.6)	CALLAHAN	65 HLBC SEE R17
R2	(21.6)	(0.6)	TRILLING	65 RVUE
R2	16K	21.18	CHIANG	72 OSPK + 1.84 GEV/C K+
R2	FIT 21.17 ± 0.15 FROM FIT			

$K^\pm \rightarrow (\pi^\pm \pi^+ \pi^-)/\text{total (units } 10^{-2})$  (P3)

R3 0	(5.6)	(0.4)	BIRGE	56 EMUL +	
R3 0	(6.8)	(0.4)	ALEXANDER	57 EMUL +	
R3 0	(5.2)	(0.3)	TAYLOR	59 EMUL +	
R3 0	EARLIER EXPERIMENTS NOT AVERAGED				
R3	5.7	0.3	ROE	61 HLBC +	
R3	2332	5.54	0.12	CALLAHAN	64 HLBC +
R3	540	5.71	0.2	SHAKLEE	64 HLBC +
R3		5.71	0.15	DE MARCO	65 HBC
R3	44	6.0	0.4	YOUNG	65 EMUL +
R3 P	693	5.34	0.21	PANDOULAS	70 EMUL +
R3 C	2330	(5.56)	(0.20)	CHIANG	72 OSPK + 1.84 GEV/C K+
R3 C	THIS VALUE IS NOT INDEPENDENT OF CHIANG R1,R2,R4,R5, AND R6				
R3 P	INCLUDES EVENTS OF TAYLOR 59.				
R3					
R3 AVG	5.521	0.098	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)		
R3 FIT	5.589	0.030	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1) (SEE IDEOGRAM BELOW)		

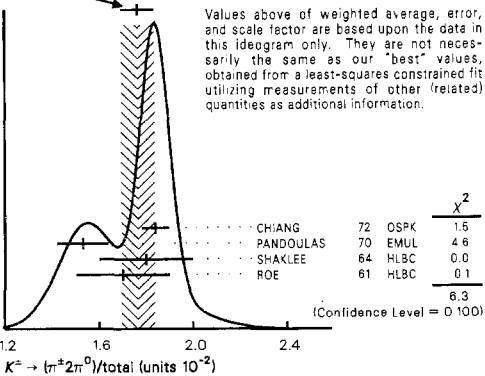
WEIGHTED AVERAGE  
5.521 ± 0.098 (ERROR SCALED BY 1.3)



$K^\pm \rightarrow (\pi^\pm 2\pi^0)/\text{total (units } 10^{-2})$  (P4)

R4 0	(2.1)	(0.5)	BIRGE	56 EMUL +	
R4 0	(2.2)	(0.4)	ALEXANDER	57 EMUL +	
R4 0	(1.5)	(0.2)	TAYLOR	59 EMUL +	
R4 0	EARLIER EXPERIMENTS NOT AVERAGED				
R4	1.7	0.2	ROE	61 HLBC +	
R4	108	1.8	0.2	SHAKLEE	64 HLBC +
R4 P	198	1.53	0.11	PANDOULAS	70 EMUL +
R4	1307	1.84	0.06	CHIANG	72 OSPK + 1.84 GEV/C K+
R4 P	INCLUDES EVENTS OF TAYLOR 59.				
R4					
R4 AVG	1.767	0.071	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)		
R4 FIT	1.733	0.046	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.4) (SEE IDEOGRAM BELOW)		

WEIGHTED AVERAGE  
1.767 ± 0.071 (ERROR SCALED BY 1.4)



$K^\pm \rightarrow (\pi^0 \mu^\pm \nu)/\text{total (units } 10^{-2})$  (P5)

R5 0	(2.8)	(1.0)	BIRGE	56 EMUL +	
R5 0	(5.9)	(1.3)	ALEXANDER	57 EMUL +	
R5 0	(2.8)	(0.4)	TAYLOR	59 EMUL +	
R5 0	EARLIER EXPERIMENTS NOT AVERAGED				
R5	2345	3.33	0.16	CHIANG	72 OSPK + 1.84 GEV/C K+
R5	FIT 3.180 ± 0.095 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.9)				

$K^\pm \rightarrow (\pi^0 e^\pm \nu)/\text{total (units } 10^{-2})$  (P6)

R6 0	(3.2)	(1.3)	BIRGE	56 EMUL +	
R6 0	(5.1)	(0.3)	ALEXANDER	57 EMUL +	
R6 0	EARLIER EXPERIMENTS NOT AVERAGED				
R6	5.0	0.5	ROE	61 HLBC +	
R6	429	4.7	0.3	SHAKLEE	64 HLBC +
R6	3516	4.06	0.10	CHIANG	72 OSPK + 1.84 GEV/C K-
R6					
R6 AVG	4.849	0.093	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)		
R6 FIT	4.819	0.052	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		

$K^\pm \rightarrow (\pi^\pm \pi^0 + \mu^\pm \pi^0 \nu)/\text{total (units } 10^{-2})$  (P2+P5)

R7	WE COMBINE THESE TWO MODES FOR EXPTS MEASURING THEM IN XENON BC BECAUSE OF DIFFICULTIES OF SEPARATING THEM THERE				
R7	23.4	1.1	ROE	61 HLBC +	
R7	886	25.4	0.9	SHAKLEE	64 HLBC +
R7					
R7 AVG	26.60	0.98	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)		
R7 FIT	24.35	0.15	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		

$K^+ \rightarrow (\pi^+ \pi^+ e^- \nu)/\text{total (units } 10^{-7})$  (P8)

R8	TEST OF DELTA-S = DELTA-Q RULE				
R8	(20.0)	OR LESS	CL=.95	BIRGE	65 FBC +
R8	(6.9)	OR LESS	CL=.95	ELY	69 HLBC +
R8	(9.0)	OR LESS	CL=.95	SCHWEINBE	71 HLBC +

$K^+ \rightarrow (\pi^+ \pi^- \mu^+ \nu)/\text{total (units } 10^{-5})$  (P9)

R9	1	(0.77)	(0.54)	CLINE	65 FBC +
----	---	--------	--------	-------	----------

$K^+ \rightarrow (\pi^+ \pi^+ \mu^- \nu)/\text{total (units } 10^{-6})$  (P10)

R10	TEST OF DELTA-S = DELTA-Q RULE					
R10	0	3.0	OR LESS	CL=.95	BIRGE	65 FBC +

$K^\pm \rightarrow (e^\pm \nu)/\text{total (units } 10^{-5})$  (P11)

R11	(160.0)	OR LESS	CL=.95	BORREANI	64 HBC +	
R11	4	(2.1)	(1.8)	(1.3)	BOWEN	67 OSPK +

$K^\pm \rightarrow (\pi^\pm \gamma \gamma)/\text{total (units } 10^{-4})$  (P17)

ALL VALUES GIVEN HERE ASSUME A PHASE SPACE PION ENERGY SPECTRUM

R12	(-0.1)	(0.6)	CHEN	68 OSPK + T(P1) 60-90 MEV		
R12	0	(0.5)	OR LESS	CL=.90	KLEMS	71 OSPK + T(P1)GT 117 MEV
R12	0	(0.35)	OR LESS	CL=.90	LJUNG	73 HLBC + 6-102, 114-127 MEV
R12	0	(-0.42)	(0.52)	ABRAMS	77 SPEC + T(P1)LT 92 MEV	
R12	0	0.084	OR LESS	CL=.90	ASANO	82 CNTR + T(P1)-117-127 MEV

$K^\pm \rightarrow (\pi^\pm \pi^0 \gamma)/\text{total (units } 10^{-4})$  (P13)

R13 0	18	(2.2)	(0.7)	CLINE	64 FBC + P1+ KE 55-80 MEV	
R13 0	(1.9)	OR LESS	CL=.90	EMMERSON	69 OSPK P1+ KE 55-80 MEV	
R13 M	0	(1.0)	OR LESS	MALTSEV	70 HLBC + P1+ KE LT 55 MEV	
R13 A2100	2.71	0.19	BRAXMS	72 ASPK -- P1+ KE 55-90 MEV		
R13 0	24	(2.4)	(0.8)	EDWARDS	72 OSPK P1+ KE 58-90 MEV	
R13 L	(1.5)	(1.1)	(0.6)	LJUNG	73 HLBC + P1+ KE 55-80 MEV	
R13 L	(2.6)	(1.5)	(1.1)	LJUNG	73 HLBC + P1+ KE 55-90 MEV	
R13 OL	17	(4.8)	(3.7)	(2.1)	LJUNG	73 HLBC + P1+ KE 55-102 MEV
R13	2461	2.87	0.52	SMITH	76 WIRE -- P1+ KE 55-90 MEV	
R13 0	ONLY HIGH STATISTICS EXPERIMENTS ARE AVERAGED.					
R13 M	MALTSEV 70 SELECTS LOW P1- ENERGY TO ENHANCE DIRECT EMISSION CONTR.					
R13 L	THE LJUNG 73 VALUES ARE NOT INDEPENDENT.					
R13 A	ABRAMS 72 OBSERVES DIRECT EMISSION BR. RATIO OF (1.56 ± 0.35) × 10 <sup>-5</sup>					
R13 A	± 0.5 × 10 <sup>-5</sup> ADDNL. SYST. ERROR AND INNER BREMSSTRAHLUNG BR. RATIO					
R13 A	OF (2.55 ± 0.18) × 10 <sup>-4</sup> . WE QUOTE THE SUM OF THESE BR. RATIOS.					
R13						
R13 AVG	2.75	0.16	AVERAGE			

$K^\pm \rightarrow (\pi^\pm \pi^+ \pi^- \gamma)/\text{total (units } 10^{-4})$  (P14)

R14	1.0	0.4	STAMER	65 EMUL + EGAM ET 11MEV
-----	-----	-----	--------	-------------------------

$K^\pm \rightarrow (\pi^\pm e^+ e^-)/\text{total (units } 10^{-6})$  (P15)

R15	TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY COMBINED					
R15	FIRST ORDER WEAK AND E.M. INTERACTIONS.					
R15	1	(2.45)	OR LESS	CL=.90	CAMERINI	64 FBC +
R15	(4.4)	OR LESS	CL=.90	BISI	67 DBC +	
R15 C	(0.4)	OR LESS	CL=.90	CLINE1	67 FBC +	
R15 C	(0.88)	OR LESS	CL=.90	CLINE2	67 FBC +	
R15	(32.0)	OR LESS	CL=.90	BEIER	72 OSPK --	
R15	(1.7)	OR LESS	CL=.90	CENCE	74 ASPK + THREE TRACK EVTS	
R15	(0.27)	OR LESS	CL=.90	CENCE	74 ASPK + TWO TRACK EVENTS	
R15 C	CLINE2 REPLACES CLINE1. CLINE1 IS NOT FOR CL=.90.					

$K^\pm \rightarrow (\pi^\pm \mu^+ \mu^-)/\text{total (units } 10^{-6})$  (P16)

R16	TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY COMBINED				
R16	FIRST ORDER WEAK AND E.M. INTERACTIONS.				
R16	(3.0)	OR LESS	CL=.90	CAMERINI	65 FBC +
R16	2.4	OR LESS	CL=.90	BISI	67 DBC +

$K^\pm \rightarrow (\pi^\pm \pi^0)/(\pi^\pm \pi^+ \pi^-)$  (P2)/(P3)

R17	134	3.24	0.34	YOUNG	65 EMUL +
R17	1045	3.96	0.15	CALLAHAN	66 FBC +
R17					
R17 AVG	3.84	0.27	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)		
R17 FIT	3.787	0.034	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		

$K^\pm \rightarrow (\pi^\pm 2\pi^0)/(\pi^\pm \pi^+ \pi^-)$  (P4)/(P3)

R18	2027	0.303	0.009	BISI	65 H-HL +
R18	17	0.393	0.099	YOUNG	65 EMUL +
R18					
R18 AVG	0.3037	0.0090	AVERAGE		
R18 FIT	0.3100	0.0079	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)		



## Stable Particle Full Listings

 $K^\pm$  $K^+ \rightarrow (\pi^+ \pi^+ e^- \nu)/(\pi^+ \pi^- e^+ \nu)$  (units  $10^{-4}$ ) (P8)/(P7)

R37 TEST OF DELTA-S = DELTA-Q RULE  
 R37 0 (130.) OR LESS CL=.95 BOURQUIN 71 ASPK  
 R37 B 3 3.6 OR LESS CL=.95 BLOCH 76 SPEC  
 R37 B CORRESPONDS TO 3E10-4 AT CL=.90.

 $K^\pm \rightarrow (\pi^0 \pi^0 e^\pm \nu)/(\pi^0 e^\pm \nu)$  (units  $10^{-4}$ ) (P24)/(P6)

R38 0 (37.0) OR LESS CL=.90 ROMANO 71 HLBC +  
 R38 2 3.8 5.0 1.2 LJUNG 73 HLBC +

 $K^+ \rightarrow (\pi^- e^+ \mu^+)/\text{total}$  (units  $10^{-8}$ ) (P25)

R39 K- INTO (PI+ E- MU-)/TOTAL IS ALSO INCLUDED HERE  
 R39 (2.8) OR LESS CL=.90 BEIER 72 OSPK --

 $K^+ \rightarrow (\pi^+ e^+ \mu^-)/\text{total}$  (units  $10^{-8}$ ) (P26)

R40 K- INTO (PI- E- MU-)/TOTAL IS ALSO INCLUDED HERE  
 R40 (1.4) OR LESS CL=.90 BEIER 72 OSPK --

 $K^\pm \rightarrow (\mu^\pm 3\nu)/\text{total}$  (units  $10^{-6}$ ) (P27)

R41 P 0 6.0 OR LESS CL=.90 PANG 73 CNTR +  
 R41 P PANG 73 ASSUMES MU SPECTRUM FROM NU-NU INTERACTION OF BARDIN 70.

 $K^\pm \rightarrow (\pi^0 \mu^\pm \nu \gamma)/\text{total}$  (units  $10^{-5}$ ) (P28)

R42 0 6.1 OR LESS CL=.90 LJUNG 73 HLBC + EGAM GT 30 MEV

 $K^\pm \rightarrow (\pi^0 e^\pm \nu)/(\pi^\pm \pi^0)$  (P6)/(P7)

R43 L 786 0.221 0.012 LUCAS2 73 HBC - DALITZ PRS ONLY  
 R43 L LUCAS 73 GIVES N(E3)=786--3.PCT, N(PI2)=3564--3.1PCT. WE DIVIDE.  
 R43  
 R43 FIT 0.2277 0.0031 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)

 $K^\pm \rightarrow (\pi^\pm 2\pi^0)/(\pi^\pm \pi^0)$  (P4)/(P2)

R44 L 574 0.081 0.005 LUCAS2 73 HBC - DALITZ PRS ONLY  
 R44 L LUCAS 73 GIVES N(PI 2PI0)=574--5.9 PCT, N(PI2)=3564--3.1 PCT.  
 R44 L WE QUOTE 0.5\*N(PI 2PI0)/N(PI2) WHERE 0.5 IS BECAUSE ONLY DALITZ  
 R44 L PAIR PI0'S WERE USED.  
 R44  
 R44 FIT 0.0819 0.0022 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)

 $K^\pm \rightarrow (\mu^\pm \nu \gamma)/\text{total}$  (units  $10^{-3}$ ) (P12)

R45 SEE ALSO SECTIONS R60, R61, AND R62 BELOW.  
 R45 12 (5.8) (3.5) WEISSENBE 74 STRC + E-GAMMA ST 9 MEV  
 R45 A 5.4 0.3 AKIBA 85 SPEC P(MU)=231.5MEV/C  
 R45 A ASSUMES MU-E UNIVERSALITY AND USES CONSTRAINTS FROM  $K^+ \rightarrow e^- \mu^+ \nu$  GAMMA.

 $K^\pm \rightarrow (\pi^\pm e^+ e^-)/(\pi^+ \pi^- e^\pm \nu)$  (units  $10^{-3}$ ) (P15)/(P7)

R46 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALL LED BY COMBINED  
 R46 FIRST ORDER WEAK AND E.M. INTERACTIONS.  
 R46 B 41 7.0 1.3 BLOCH 75 S--  
 R46 B BLOCH 75 QUOTES THIS RESULT MULTIPLIED BY OUR 1974 KE4 BR.FRAC.

 $K^\pm \rightarrow (e^\pm \nu \gamma)/(e^\pm \nu)$  (P21)/(P11)

R47 STRUCTURE DEPENDENT PART WITH + GAMMA HELICITY.  
 R47 H 56 (1.05) (0.25) (0.30) HEARD1 75 SPEC + P(E) 236 TO 247  
 R47 H THIS VALUE IS INCLUDED IN THE SECOND HEINTZE 79 VALUE IN SEC.R54  
 R47 H BELOW.

 $K^+ \rightarrow (\pi^\pm \mu^\pm e^+)/(\pi^+ \pi^- e^+ \nu)$  (units  $10^{-4}$ ) (P25+P26)/(P7)

R48 TEST OF LEPTON FAMILY NUMBER OR TOTAL LEPTON NUMBER CONSERVATION.  
 R48 D 0 1.9 OR LESS CL=.90 DIAMANTBE 76 SPEC +  
 R48 D DIAMANTBE 76 QUOTES THIS RESULT TIMES OUR 1975 KE4 BR. RATIO.

 $K^+ \rightarrow (\pi^+ \mu^+ e^-)/(\pi^+ \pi^- e^+ \nu)$  (units  $10^{-4}$ ) (P29)/(P7)

R49 TEST OF LEPTON FAMILY NUMBER CONSERVATION.  
 R49 D 0 1.3 OR LESS CL=.90 DIAMANTBE 76 SPEC +  
 R49 D DIAMANTBE 76 QUOTES THIS RESULT TIMES OUR 1975 KE4 BR RATIO.

 $K^\pm \rightarrow (\mu^\pm \nu e^+ e^-)/(\pi^+ \pi^- e^\pm \nu)$  (units  $10^{-3}$ ) (P30)/(P7)

R50 D 14 (3.3) (0.9) DIAMANTBE 76 SPEC + M(EE) GT 140  
 R50 D 14 (2.33) (0.42) DIAMANTBE 76 SPEC + EXTRAPOLATED BR  
 R50 D DIAMANTBE 76 QUOTES THESE RESULTS TIMES OUR 1975 KE4 BR RATIO.  
 R50 D THE SECOND DIAMANTBE 76 VALUE IS THE FIRST VALUE EXTRAPOLATED TO 0  
 R50 D TO INCLUDE LOW MASS E PAIRS.

 $K^+ \rightarrow (\pi^- e^+ e^+)/(\pi^+ \pi^- e^+ \nu)$  (units  $10^{-4}$ ) (P19)/(P7)

R51 TEST OF TOTAL LEPTON NUMBER CONSERVATION  
 R51 D 0 2.5 OR LESS CL=.90 DIAMANTBE 76 SPEC +  
 R51 D DIAMANTBE 76 QUOTES THIS RESULT TIMES OUR 1975 BR RATIO.

 $K^+ \rightarrow (\mu^- \nu e^+ e^+)/(\pi^+ \pi^- e^+ \nu)$  (units  $10^{-3}$ ) (P31)/(P7)

R52 TEST OF LEPTON FAMILY NUMBER CONSERVATION.  
 R52 D 0 0.5 OR LESS CL=.90 DIAMANTBE 76 SPEC +  
 R52 D DIAMANTBE 76 QUOTES THIS RESULT TIMES OUR 1975 KE4 BR RATIO.

 $K^+ \rightarrow (\nu e^+ e^+ e^-)/(\pi^+ \pi^- e^+ \nu)$  (units  $10^{-2}$ ) (P32)/(P7)

R53 4 0.54 0.54 0.27 DIAMANTBE 76 SPEC +  
 R53 D DIAMANTBE 76 QUOTES THIS RESULT TIMES OUR 1975 KE4 BR RATIO.

 $K^\pm \rightarrow (e^\pm \nu \gamma)/(\mu^\pm \nu)$  (units  $10^{-5}$ ) (P21)/(P11)

R54 STRUCTURE DEPENDENT PART WITH + GAMMA HELICITY  
 R54 H 51 (2.33) (0.42) HEINTZE 79 SPEC +  
 R54 H 107 2.40 0.36 HEINTZE 79 SPEC +  
 R54 H SECOND HEINTZE 79 RESULT IS FIRST COMBINED WITH HEARD1 75 RESULT  
 R54 H FROM SECTION R47 ABOVE.

 $K^\pm \rightarrow (e^\pm \nu \gamma)/\text{total}$  (units  $10^{-4}$ ) (P21)

R55 STRUCTURE DEPENDENT PART WITH - GAMMA HELICITY  
 R55 H 1.6 OR LESS CL=.90 HEINTZE 79 SPEC +  
 R55 H IMPLIES (AXIAL VEC./VECTOR) AMPL. RATIO OUTSIDE RANGE -1.8 TO -.54.

 $K^\pm \rightarrow (e^\pm \nu \nu \bar{\nu})/(e^\pm \nu)$  (P33)/(P11)

R56 0 3.8 OR LESS CL=.90 HEINTZE 79 SPEC -

 $K^+ \rightarrow (\mu^+ \nu_e)/\text{total}$  (P34)

R57 FORBIDDEN BY LEPTON FAMILY NUMBER CONSERVATION.  
 R57 0 0.004 OR LESS CL=.90 LYONS 81 HLBC 200GEV K+ N.B. BEAM  
 R57 (0.012) OR LESS CL=.90 COOPER 82 HLBC WIDEBAND NU BEAM

 $K^+ \rightarrow \mu^+ \bar{\nu}_e/\text{total}$  (units  $10^{-3}$ ) (P35)

R58 FORBIDDEN BY TOTAL LEPTON NUMBER CONSERVATION.  
 R58 3.3 OR LESS CL=.90 COOPER 82 HLBC WIDEBAND NU BEAM

 $K^+ \rightarrow \pi^0 e^+ \bar{\nu}_e/\text{total}$  (P36)

R59 FORBIDDEN BY TOTAL LEPTON NUMBER CONSERVATION.  
 R59 0.003 OR LESS CL=.90 COOPER 82 HLBC WIDEBAND NU BEAM

 $K^\pm \rightarrow (\mu^\pm \nu \gamma)/\text{total}$  (units  $10^{-5}$ ) (P12)

R60 STRUCTURE DEPENDENT PART WITH + GAMMA HELICITY (SD- TERM).  
 R60 BRANCHING RATIOS R60, R61, AND R62 ARE COMPONENTS OF R45 ABOVE.  
 R60 5.0 OR LESS CL=.90 AKIBA 85 SPEC

 $K^\pm \rightarrow (\mu^\pm \nu \gamma)/\text{total}$  (units  $10^{-5}$ ) (P12)

R61 INTERFERENCE TERM BETWEEN INTERNAL BREMSSTRAHLUNG AND SD- TERM.  
 R61 2.7 OR LESS CL=.90 AKIBA 85 SPEC

 $K^\pm \rightarrow (\mu^\pm \nu \gamma)/\text{total}$  (units  $10^{-4}$ ) (P12)

R62 SUM OF STRUCTURE DEPENDENT PART WITH - GAMMA HELICITY (SD- TERM)  
 AND INTERFERENCE TERM BETWEEN INTERNAL BREMSSTRAHLUNG AND SD- TERM.  
 R62 A 2.6 OR LESS CL=.90 AKIBA 85 SPEC  
 R62 A ASSUMES MU-E UNIVERSALITY AND USES CONSTRAINTS FROM  $K^+ \rightarrow e^- \mu^+ \nu$  GAMMA.

 $K^\pm$  LONGITUDINAL POLARIZATION OF EMITTED  $\mu^\pm$ 

POL K+- INTO MU+- NU  
 POL TESTS FOR RIGHT-HANDED CURRENTS IN STRANGENESS-CHANGING DECAY.  
 POL -0.970 0.047 HAYANO 84 CNTR  
 POL -1.0 0.1 CUTTS 69 SPRK  
 POL -0.96 0.12 COMBES 57 CNTR  
 POL AVE -0.974 0.040 AVERAGE

NOTE ON DALITZ PLOT PARAMETERS FOR  $K \rightarrow 3\pi$  DECAYS

The Dalitz plot distribution for  $K^\pm \rightarrow \pi^\pm \pi^\pm \pi^\mp$ ,  $K^\pm \rightarrow \pi^0 \pi^0 \pi^\pm$ , and  $K_L^0 \rightarrow \pi^+ \pi^- \pi^0$  can be parametrized by a series expansion such as that introduced by Weinberg.<sup>1</sup> We use the form

$$|M|^2 \propto 1 + g \frac{(s_3 - s_0)}{m_{\pi^+}^2} + h \left[ \frac{s_3 - s_0}{m_{\pi^+}^2} \right]^2 + j \frac{(s_2 - s_1)}{m_{\pi^+}^2} + k \left[ \frac{s_2 - s_1}{m_{\pi^+}^2} \right]^2 + \dots, \quad (1)$$

where  $m_{\pi^+}^2$  has been introduced to make the coefficients  $g$ ,  $h$ ,  $j$ , and  $k$  dimensionless, and

$$s_i = (P_K - P_i)^2 = (m_K - m_i)^2 - 2m_K T_i, \quad i = 1, 2, 3,$$

$$s_0 = \frac{1}{3} \sum_i s_i = \frac{1}{3} (m_K^2 + m_1^2 + m_2^2 + m_3^2).$$

Here the  $P_i$  are four-vectors,  $m_i$  and  $T_i$  are the mass and kinetic energy of the  $i^{\text{th}}$  pion, and the index 3 is used for the odd pion.

The coefficient  $g$  is a measure of the slope in the variable  $s_3$  (or  $T_3$ ) of the Dalitz plot, while  $h$  and  $k$  measure the quadratic dependence on  $s_3$  and  $(s_2 - s_1)$ , respectively. The coefficient  $j$  is related to the asymmetry of the plot and must be zero if  $CP$  invariance holds. Note also that if  $CP$  is good,  $g$ ,  $h$ , and  $k$  must be the same for  $K^+ \rightarrow \pi^+ \pi^+ \pi^-$  as for  $K^- \rightarrow \pi^- \pi^- \pi^+$ .

For notation, see key on page 91.

# Stable Particle Full Listings

$K^\pm$

Since different experiments use different forms for  $|M|^2$ , in order to compare the experiments we have converted to  $g, h, j,$  and  $k$  whatever coefficients have been measured. For details of this conversion and discussion of the data, see the April 1982 version of this note.<sup>2</sup>

See also the review of Devlin and Dickey,<sup>3</sup> which contains an analysis of  $K \rightarrow 2\pi$  and  $K \rightarrow 3\pi$  data in terms of transition amplitudes with appropriate energy dependence.

## References

1. S. Weinberg, Phys. Rev. Lett. **4**, 87 (1960).
2. Particle Data Group, Phys. Lett. **111B**, 69 (1982).
3. T.J. Devlin and J.O. Dickey, Rev. Mod. Phys. **51**, 237 (1979).

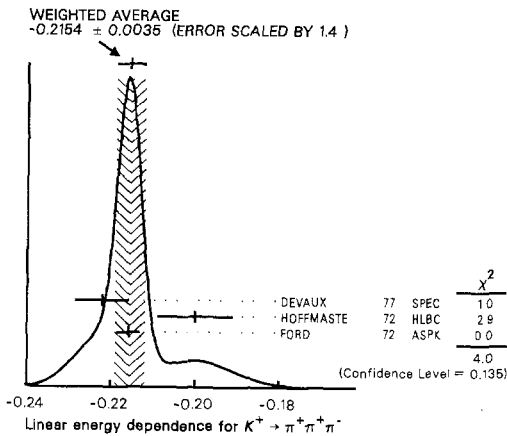
### $K^\pm$ ENERGY DEPENDENCE OF DALITZ PLOT

MATRIX ELEMENT SQUARED =  $1 + G^*U - H^*U^*2 + K^*v^*2$   
 WHERE  $U=(S3-S0)/(MP1^*2)$  AND  $V=(S1-S2)/(MP1^*2)$

**LINEAR COEFFICIENT  $g_T^+$  FOR  $K^+ \rightarrow \pi^+\pi^+\pi^-$**   
 SOME EXPTS USE DALITZ VARIABLES X AND Y. WE GIVE AY=COEFF OF Y TERM AT RIGHT. SEE MINI-REVIEW ABOVE.  

GT+ZL 5428	(-0.223)	(0.024)	ZINCHENKO 67 HBC	+ AY=0.28--0.03
GT+L 9994	(-0.218)	(0.016)	BUTLER 68 HBC	+ AY=0.277--0.020
GT+ G17898	(-0.196)	(0.012)	GRAUMAN 70 HLBC	+ AY=0.228--0.030
GT+ 750K	-0.2157	0.0028	FORD 72 ASPK	+ AY=0.2754--0.0035
GT+H 39819	-0.200	0.009	HOFFMASTE 72 HLBC	+
GT+ 225K	-0.2221	0.0065	DEVAUX 77 SPEC	+ AY=0.2814--0.0082

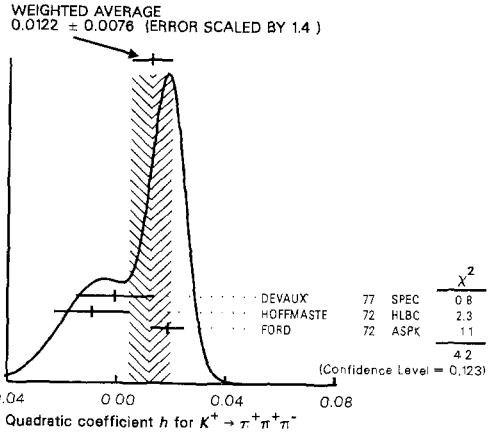
 EXPTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE.  
 GT+ L ALSO INCLUDES OBC EVENTS  
 GT+ G EMULS. DATA ADDED - ALL EVENTS INCLUDED BY HOFFMASTE 72  
 GT+H HOFFMASTE 72 INCLUDES GRAUMAN 70 DATA.  
 GT-  
 GT+ AVG -0.2154 0.0035 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)  
 (SEE IDEOGRAM BELOW)



**QUADRATIC COEFFICIENT  $h$  FOR  $K^+ \rightarrow \pi^+\pi^+\pi^-$**   

HT- 750K	0.0187	0.0062	FORD	72	ASPK	+
HT+ 39819	-0.009	0.014	HOFFMASTE	72	HLBC	+
HT+ 225K	-0.0006	0.0143	DEVAUX	77	SPEC	-
HT+ AVG	-0.0122	0.0076	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)			

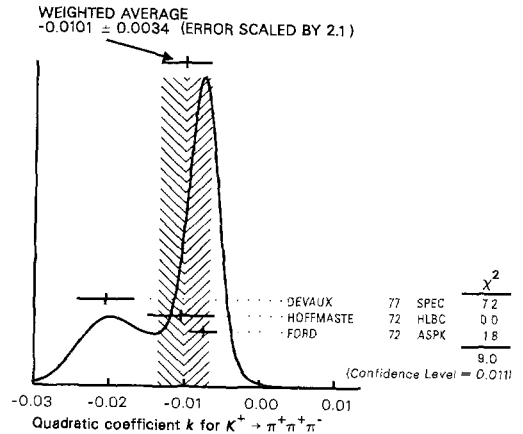
 (SEE IDEOGRAM BELOW)



**QUADRATIC COEFFICIENT  $k$  FOR  $K^+ \rightarrow \pi^+\pi^+\pi^-$**   

KT- 750K	-0.0075	0.0019	FORD	72	ASPK	+
KT+ 39819	-0.0105	0.0065	HOFFMASTE	72	HLBC	-
KT+ 225K	-0.0205	0.0039	DEVAUX	77	SPEC	-
KT+ AVG	-0.0101	0.0034	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.1)			

 (SEE IDEOGRAM BELOW)



**LINEAR COEFFICIENT  $g_T^-$  FOR  $K^- \rightarrow \pi^-\pi^-\pi^+$**   
 GT- FOR DEFINITION OF AY SEE NOTE IN SECTION GT- ABOVE.  

GT- F 1347	(-0.220)	(0.035)	FERRI-LUIZ 61 HBC	- AY=0.28--0.045
GT-ML 5778	(-0.190)	(0.023)	MOSCOSO 68 HBC	- AY=0.242--0.029
GT- 50919	-0.193	0.010	MAST 69 HBC	- AY=0.244 +--0.013
GT- 750K	-0.2186	0.0028	FORD 72 ASPK	- AY=0.2770--0.0035
GT- Q 81K	(-0.199)	(0.008)	LUCAST 73 HBC	- AY=0.252--0.011

 NO RADIATIVE CORRECTIONS INCLUDED.  
 GT- L ALSO INCLUDES OBC EVENTS.  
 GT- Q QUADRATIC DEPENDENCE IS REQUIRED BY KL EXPTS. FOR COMPARISON WE  
 GT- Q AVERAGE ONLY THOSE K- EXPERIMENTS WHICH QUOTE QUADRATIC FIT VALUES.  
 GT-  
 GT- AVG -0.2167 0.0066 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.5)

**QUADRATIC COEFFICIENT  $h$  FOR  $K^- \rightarrow \pi^-\pi^-\pi^+$**   

HT- 50919	-0.001	0.012	MAST	69	HBC	-
HT- 750K	0.0125	0.0062	FORD	72	ASPK	-
HT- AVG	0.0097	0.0055	AVERAGE			

**QUADRATIC COEFFICIENT  $k$  FOR  $K^- \rightarrow \pi^-\pi^-\pi^+$**   

KT- 50919	-0.014	0.012	MAST	69	HBC	-
KT- 750K	-0.0083	0.0019	FORD	72	ASPK	-
KT- AVG	-0.0084	0.0019	AVERAGE			

**$(g_T^+ - g_T^-)/(g_T^+ + g_T^-)$  IN PERCENT**  
 DG A NON-ZERO VALUE FOR THIS QUANTITY INDICATES CP VIOLATION  

DG 3.2M	-0.70	0.53	FORD	70	ASPK
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# Stable Particle Full Listings

$K^\pm$

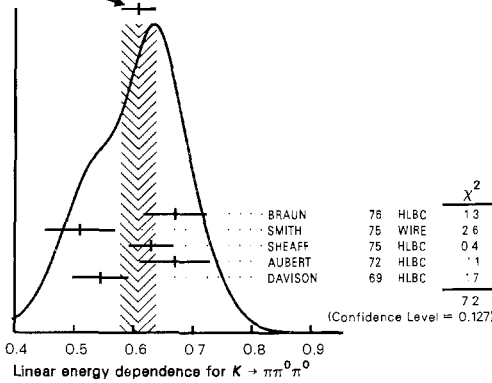
## LINEAR COEFFICIENT $g$ FOR $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ .

UNLESS OTHERWISE STATED, ALL EXPTS INCLUDE TERMS QUADRATIC IN  $(S_3 - S_0)/(M\pi^{**2})$ . SEE MINI-REVIEW ABOVE.

GTP K 1792	(0.48)	(0.04)	KALMUS	64	HLBC +
GTP K 1874	(0.586)	(0.098)	BISI	65	HLBC + ALSO HBC
GTP C 8	0.544	0.048	DAVIDSON	69	HLBC + ALSO EMUL
GTP L 198	(1.57)	(0.102)	PANDOLAS	70	EMUL +
GTP 1365	J.67	0.06	AUBERT	72	HLBC +
GTP K 574	(0.484)	(0.084)	LUCAS2	73	HBC - DALITZ PRS ONLY
GTP 5635	0.630	0.038	SHEAFF	75	HLBC -
GTP 27K	0.510	0.060	SMITH	75	WIRE -
GTP L 4639	(0.806)	(0.220)	BERTRAND	76	EMUL +
GTP 3263	0.670	0.054	BRAUN	76	HLBC +

GTP K AUTHORS GIVE LINEAR FIT ONLY.  
GTP L EXPERIMENTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE.  
GTP AVG 0.607 ± 0.030 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)  
(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
0.607 ± 0.030 (ERROR SCALED BY 1.3)



## QUADRATIC COEFFICIENT $h$ FOR $K^\pm \rightarrow \pi^\pm \pi^0 \pi^0$ .

SEE MINI-REVIEW ABOVE.

HTP 4048	0.026	0.050	DAVIDSON	69	HLBC + ALSO EMUL
HTP L 198	(0.018)	(0.124)	PANDOLAS	70	EMUL +
HTP 1365	-0.01	0.08	AUBERT	72	HLBC +
HTP 5635	0.041	0.030	SHEAFF	75	HLBC +
HTP 27K	0.009	0.040	SMITH	75	WIRE +
HTP L 4639	(0.164)	(0.121)	BERTRAND	76	EMUL +
HTP 3263	0.152	0.082	BRAUN	76	HLBC +

HTP L EXPERIMENTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE.  
HTP AVG 0.054 ± 0.020 AVERAGE

## NOTE ON $K_{e3}^\pm$ AND $K_{e3}^0$ FORM FACTORS

Assuming that only the vector current contributes to  $K \rightarrow \pi e \nu$  decays, we write the matrix element as

$$M \propto f_+(t) [(P_K + P_\pi)_\mu \bar{e} \gamma_\mu (1 + \gamma_5) \nu] + f_-(t) [m_e \bar{e} (1 + \gamma_5) \nu], \quad (1)$$

where  $P_K$  and  $P_\pi$  are the four-momenta of the  $K$  and  $\pi$  mesons,  $m_e$  is the lepton mass, and  $f_+$  and  $f_-$  are dimensionless form factors which can depend only on  $t = (P_K - P_\pi)^2$ , the square of the four-momentum transfer to the leptons. If time-reversal invariance holds,  $f_+$  and  $f_-$  are relatively real.  $K_{\mu 3}$  experiments measure  $f_+$  and  $f_-$ , while  $K_{e3}$  experiments are sensitive only to  $f_+$  because the small electron mass makes the  $f_-$  term negligible.

(a)  $K_{\mu 3}$  experiments. Analyses of  $K_{\mu 3}$  data frequently assume a linear dependence of  $f_+$  and  $f_-$  on  $t$ , i.e.,

$$f_\pm(t) = f_\pm(0) [1 + \lambda_\pm (t/m_\pi^2)]. \quad (2)$$

Most  $K_{\mu 3}$  data are adequately described by Eq. (2) for  $f_+$  and a constant  $f_-$  (i.e.,  $\lambda_- = 0$ ). There are two equivalent parametrizations commonly used in these analyses:

(1)  $\lambda_+$ ,  $\xi(0)$  parametrization. Analyses of  $K_{\mu 3}$  data often introduce the ratio of the two form factors

$$\xi(t) = f_-(t)/f_+(t).$$

The  $K_{\mu 3}$  decay distribution is then described by the two parameters  $\lambda_+$  and  $\xi(0)$  (assuming time reversal invariance and  $\lambda_- = 0$ ). These parameters can be determined by three different methods:

*Method A.* By studying the Dalitz plot or the pion spectrum of  $K_{\mu 3}$  decay. The Dalitz plot density is (see, e.g., Chounet et al.<sup>1</sup>):

$$\rho(E_\pi, E_\mu) \propto f_+^2(t) [A + B\xi(t) + C\xi(t)^2],$$

where

$$A = m_K^2 (2E_\mu E_\nu - m_K E'_\pi) + m_\mu^2 (\frac{1}{4} E'_\pi - E_\nu),$$

$$B = m_\mu^2 (E_\nu - \frac{1}{2} E'_\pi),$$

$$C = \frac{1}{4} m_\mu^2 E'_\pi,$$

$$E'_\pi = E_\pi^{\max} - E_\pi = (m_K^2 + m_\pi^2 - m_\mu^2)/2m_K - E_\pi$$

Here  $E_\pi$ ,  $E_\mu$ , and  $E_\nu$  are, respectively, the pion, muon, and neutrino energies in the kaon center of mass. The density  $\rho$  is fit to the data to determine the values of  $\lambda_+$ ,  $\xi(0)$ , and their correlation.

*Method B.* By measuring the  $K_{\mu 3}/K_{e3}$  branching ratio and comparing it with the theoretical ratio (see, e.g., Fearing et al.<sup>2</sup>) as given in terms of  $\lambda_+$  and  $\xi(0)$ , assuming  $\mu$ -e universality:

$$\Gamma(K_{\mu 3}^\pm)/\Gamma(K_{e3}^\pm) = 0.6457 + 1.4115\lambda_+ + 0.1264\xi(0) + 0.0192\xi(0)^2 + 0.0080\lambda_+\xi(0),$$

$$\Gamma(K_{\mu 3}^0)/\Gamma(K_{e3}^0) = 0.6452 + 1.3162\lambda_+ + 0.1264\xi(0) + 0.0186\xi(0)^2 + 0.0064\lambda_+\xi(0).$$

This cannot determine  $\lambda_+$  and  $\xi(0)$  simultaneously but simply fixes a relationship between them.

*Method C.* By measuring the muon polarization in  $K_{\mu 3}$  decay. In the rest frame of the  $K$ , the  $\mu$  is expected to be polarized in the direction  $\mathbf{A}$  with  $\mathbf{P} = \mathbf{A}/|\mathbf{A}|$ , where  $\mathbf{A}$  is given (Cabibbo and Maksymowicz<sup>3</sup>) by

For notation, see key on page 91.

## Stable Particle Full Listings

$K^\pm$

$$\begin{aligned} \mathbf{A} = & a_1(\xi)\mathbf{p}_\mu \\ & - a_2(\xi) \left[ \frac{\mathbf{p}_\mu}{m_\mu} \left( m_K - E_\pi + \frac{\mathbf{p}_\pi \cdot \mathbf{p}_\mu}{|\mathbf{p}_\mu|^2} (E_\mu - m_\mu) \right) + \mathbf{p}_\pi \right] \\ & + m_K \text{Im}\xi(t)(\mathbf{p}_\pi \times \mathbf{p}_\mu). \end{aligned}$$

If time-reversal invariance holds,  $\xi$  is real, and thus there is no polarization perpendicular to the  $K$ -decay plane. Polarization experiments measure the weighted average of  $\xi(t)$  over the  $t$  range of the experiment, where the weighting accounts for the variation with  $t$  of the sensitivity to  $\xi(t)$ .

(2)  $\lambda_+$ ,  $\lambda_0$  parametrization. Most of the more recent  $K_{\mu 3}$  analyses have parametrized in terms of the form factors  $f_+$  and  $f_0$  which are associated with vector and scalar exchange, respectively, to the lepton pair.  $f_0$  is related to  $f_+$  and  $f_-$  by

$$f_0(t) = f_+(t) + [t/(m_K^2 - m_\pi^2)]f_-(t).$$

Here  $f_0(0)$  must equal  $f_+(0)$  unless  $f_-(t)$  diverges at  $t = 0$ . The earlier assumption that  $f_+$  is linear in  $t$  and  $f_-$  is constant leads to  $f_0$  linear in  $t$ :

$$f_0(t) = f_0(0)[1 + \lambda_0(t/m_\pi^2)].$$

With the assumption that  $f_0(0) = f_+(0)$ , the two parametrizations,  $(\lambda_+, \xi(0))$  and  $(\lambda_+, \lambda_0)$  are equivalent as long as correlation information is retained.  $(\lambda_+, \lambda_0)$  correlations tend to be less strong than  $(\lambda_+, \xi(0))$  correlations.

The experimental results for  $\xi(0)$  and its correlation with  $\lambda_+$  are listed in the  $K^\pm$  and  $K_L^0$  sections of the Stable Particle Full Listings in section XIA, XIB, or XIC depending on whether method A, B, or C discussed above was used. The corresponding values of  $\lambda_+$  are listed in subsection L+M.

Because recent experiments tend to use the  $(\lambda_+, \lambda_0)$  parametrization, we include a subsection L0 for  $\lambda_0$  results. Wherever possible we have converted  $\xi(0)$  results into  $\lambda_0$  results and vice versa.

See the 1982 version of this note<sup>4</sup> for additional discussion of the  $K_{\mu 3}^0$  parameters, correlations, and conversion between parametrizations, and also for a comparison of the experimental results.

(b)  $K_{e3}$  experiments. Analysis of  $K_{e3}$  data is simpler than that of  $K_{\mu 3}$  because the second term of the matrix element assuming a pure vector current [Eq. (1) above]

can be neglected. Here  $f_+$  is usually assumed to be linear in  $t$ , and the linear coefficient  $\lambda_+$  of Eq. (2) is determined.

If we remove the assumption of a pure vector current, then the matrix element for the decay, in addition to the terms in Eq. (2), would contain

$$\begin{aligned} & + 2m_K f_S \bar{\ell}(1 + \gamma_5)\nu \\ & + (2f_T/m_K)(P_K)_\lambda (P_\pi)_\mu \bar{\ell}\sigma_{\lambda\mu}(1 + \gamma_5)\nu, \end{aligned}$$

where  $f_S$  is the scalar form factor, and  $f_T$  is the tensor form factor. In the case of the  $K_{e3}$  decays where the  $f_-$  term can be neglected, experiments have yielded limits on  $|f_S/f_+|$  and  $|f_T/f_+|$ .

The  $K_{e3}$  results for  $\lambda_+$ ,  $|f_S/f_+|$ , and  $|f_T/f_+|$  are listed in the subsections L+M, FS, and FT, respectively, of the  $K^\pm$  and  $K_L^0$  sections of the Stable Particle Full Listings.

### References

1. L.M. Chounet, J.M. Gaillard, and M.K. Gaillard, Phys. Rep. **4C**, 199 (1972).
2. H.W. Fearing, E. Fischbach, and J. Smith, Phys. Rev. **D2**, 542 (1970).
3. N. Cabibbo and A. Maksymowicz, Phys. Lett. **9**, 352 (1964).
4. Particle Data Group, Phys. Lett. **111B**, 73 (1982).

### $K^\pm$ FORM FACTORS

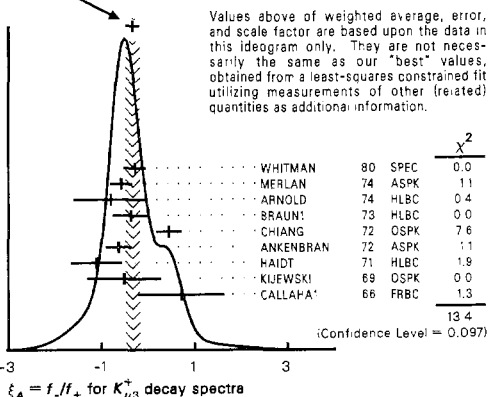
IN THE FORM FACTOR COMMENTS, THE FOLLOWING ABBREVIATIONS ARE USED.  
F+ AND F- ARE FORM FACTORS FOR THE VECTOR MATRIX ELEMENT.  
FS AND FT REFER TO THE SCALAR AND TENSOR TERM.  
FO = (F+) - (F-)\*T/(M\*\*2-MPI\*\*2)  
L+, L-, AND L0 ARE THE LINEAR EXPANSION COEFFS. OF F+, F- AND FO.  
L+ REFERS TO THE KMUS VALUE EXCEPT IN THE KE3 SECTIONS.  
DXI/DL IS THE CORRELATION BETWEEN XI(0) AND L+ IN KMUS.  
DL0/DL+ IS THE CORRELATION BETWEEN L0 AND L+ IN KMUS.  
T = MOMENTUM TRANSFER TO THE PI IN UNITS OF MPI\*\*2.  
DP = DALITZ PLOT ANALYSIS  
PI = PI SPECTRUM ANALYSIS  
MU = MU SPECTRUM ANALYSIS  
POL = MU POLARIZATION ANALYSIS  
BR = KMUS/KE3 BRANCHING RATIO ANALYSIS  
E = POSITRON OR ELECTRON SPECTRUM ANALYSIS  
RC = RADIATIVE CORRECTIONS

$\xi_A = f_-/f_+$ (DETERMINED FROM SPECTRA)									
XIA	76	(+1.8)	(0.6)	BROWN	62	XEBC + DP-BR, L+=0			
XIA	87	(+0.7)	(0.5)	GIACOMELLI	64	EMUL + MU+BR RVUE, L+=0			
XIA	J	(-0.08)	(0.7)	JENSEN	64	XEBC + DR-BR(KMUS,KE3)			
XIA	2648	(0.0)	(1.1)	(0.9)	CALLAHA1	66	FRBC + MU, L+=0, T UNKN		
XIA	C	444	+0.72	0.93	CALLAHA1	66	FRBC + PI, DXI/DL=-17		
XIA	78	(-0.5)	(0.9)	EISLER	68	HLBC + PI, L+=0, NO DX/DL			
XIA	K2041	0.5	0.8	KIJEWSKI	69	OSPK + PI, DXI/DL=-26			
XIA	H3240	-1.1	0.56	HAIDT	71	HLBC - DP, DXI/DL=-29			
XIA	A4025	-0.62	0.28	ANKENBRAND	72	ASPK + PI, DXI/DL=-12			
XIA	B3480	+0.45	0.28	CHIANG	72	OSPK + DP, DXI/DL=-15			
XIA	D1897	-0.36	0.40	BRAUN1	73	HLBC + DP, DXI/DL=-19			
XIA	M 490	-0.8	0.8	ARNOLD	74	HLBC + DP, DXI/DL=-20			
XIA	M6527	-0.57	0.24	MERLAN	74	ASPK - DP, DXI/DL=-9			
XIA	3973	-0.27	0.25	WHITMAN	80	SPEC - DP, DXI/DL=-17			
XIA	AVG	-0.32	0.15	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)					
XIA	FIT	-0.35	0.15	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.6)					
(SEE IDEOGRAM BELOW)									
FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.									
XIA	J	JENSEN 64 GIVES L+M=L-E=-.020+-027. DXI/DL UNKNOWN. INCLUDES							
XIA	J	SHAKLEE 64 XIB(KMUS/KE3).							
XIA	C	CALLAHAN 66 TABLE 1 (PI ANAL) GIVES DXI/DL=(.72-.05)/(0-.04)=-17.							
XIA	C	ERROR RAISED FROM .80 TO AGREE WITH DXI0=.37 FOR FIXED L+.							
XIA	K	KIJEWSKI 69 FIG. 17 WAS USED TO OBTAIN DXI/DL AND ERRORS.							
XIA	H	HAIDT 71 TABLE 8 (DP ANAL) GIVES DXI/DL=(-1.1+0.5)/(0.00-.029)=-29.							
XIA	H	ERROR RAISED FROM .50 TO AGREE WITH DXI0=.20 FOR FIXED L+.							
XIA	A	ANKENBRAND 72 FIG. 3 WAS USED TO OBTAIN DXI/DL.							
XIA	B	CHIANG 72 FIG. 10 WAS USED TO OBTAIN DXI/DL.							
XIA	B	FIT HAD L=L+ BUT WOULD NOT CHANGE FOR L=0. L-PONDROM, PRIV.COM.74							
XIA	D	BRAUN1 73 GIVES XI(T)=-.34+-0.20, DXI(T)/DL=-14 FOR L+=.027, T=6.6.							
XIA	D	WE CALCULATE ABOVE XI(0) AND DXI(0)/DL+ FOR THEIR L+=.025+-0.017.							
XIA	M	ARNOLD 74 FIG. 4 WAS USED TO OBTAIN XIA AND DXI/DL.							
XIA	M	MERLAN 74 FIG.5 WAS USED TO OBTAIN DXI/DL.							

# Stable Particle Full Listings

K±

WEIGHTED AVERAGE  
-0.32 ± 0.15 (ERROR SCALED BY 1.3)



## λ₀ (LINEAR ENERGY DEPENDENCE OF f₀ IN Kμ₃ DECAY)

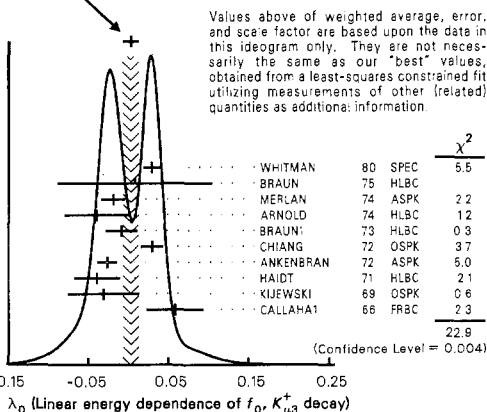
WHEREVER POSSIBLE, WE HAVE CONVERTED THE ABOVE VALUES OF XI(0) INTO VALUES OF LO USING THE ASSOCIATED L+M AND DXI/DL.

Author	Value	χ²
CALLAHA	66 FRBC + PI, DLO/DL = -0.37	
68 HLBC + POL, DLO/DL = -0.60		
CUTTS	69 OSPK + POL, DLO/DL = -0.69	
KLJEWSKI	69 OSPK + PI, DLO/DL = -1.10	
HAIDT	71 HLBC + DP, DLO/DL = -1.34	
ANKENBRAN	72 ASPK + PI, DLO/DL = +0.05	
CHIANG	72 OSPK + DP, DLO/DL = -0.21	
BRAUN	73 HLBC + DP, DLO/DL = -0.53	
ARNOLD	74 HLBC + DP, DLO/DL = -0.62	
BRAUN	74 HLBC + KMUS/KE3 VS. T	
MERLAN	74 ASPK + DP, DLO/DL = +0.27	
BRAUN	75 HLBC + DP, DLO/DL = -0.92	
H 55K	+0.019 (0.010)	
HEINTZE	77 SPEC + BR, DLO/DL = +0.03	
WHITMAN	80 SPEC + DP, DLO/DL = -0.37	
<b>L0 AVG</b>	<b>0.0032 ± 0.0099</b>	<b>AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.7)</b>
<b>L0 FIT</b>	<b>-0.004</b>	<b>0.007 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.6)</b>

(SEE IDEOGRAM BELOW)

LO FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.  
 LO W LO CALCULATED BY US FROM XI(0), L+M, AND DXI/DL.  
 LO L LO VALUE IS FOR L+M=0.03 CALCULATED BY US FROM XI(0) AND DXI/DL.  
 LO D THIS VALUE AND ERROR ARE TAKEN FROM BRAUN 75 BUT CORRESPOND TO THE  
 LO D BRAUN 73 L+M RESULT. DLO/DL IS FROM BRAUN 73 DXI/DL IN KIA ABV.  
 LO B BRAUN 74 IS A COMBINED KMUS-KE3 RESULT. IT IS NOT INDEPENDENT OF  
 LO B BRAUN 73 (KMUS) AND BRAUN 73 (KE3) FORM FACTOR RESULTS.  
 LO M MERLAN 74 LO AND DLO/DL WERE CALCULATED BY US FROM KIA, L-M, AND  
 LO H DXI/DL. THEIR FIG.6 GIVES L0 = -0.025 ± 0.012 AND NO DLO/DL.  
 LO H HEINTZE 77 USES L+M = 0.029 ± 0.003. DLO/DL ESTIMATED BY US.

WEIGHTED AVERAGE  
0.0032 ± 0.0099 (ERROR SCALED BY 1.7)



## ξ\_B = f\_-/f\_+ (DETERMINED FROM Kμ₃/KE₃)

THE KMUS/KE3 BRANCHING RATIO FIXES A RELATIONSHIP BETWEEN XI(0) AND L+. WE QUOTE THE AUTHORS XI(0) AND ASSIGNED L+ BUT DO NOT AVERAGE BECAUSE THE L+ VALUES DIFFER. THE FIT RESULT AND SCALE FACTOR GIVEN IN THE ABOVE NOTE ON KL3 FORM FACTORS ARE NOT OBTAINED FROM THESE XI(0) VALUES. INSTEAD THEY ARE OBTAINED DIRECTLY FROM THE FITTED KMUS/KE3 RATIO (R29), WITH THE EXCEPTION OF HEINTZE 77.

Author	Value	χ²
SHAKLEE	64 XEBC + BR, L+ = 0	
BISI	1 65 HBC + BR, L+ = 0	
CUTTS	65 OSPK + BR, L+ = 0	
CALLAHA	66 FRBC + BR, L+ = 0	
AUERBACH	67 OSPK + BR, L+ = 0	
BOTTERL1	68 ASPK + BR, L+ = -0.23 ± 0.08	
EICHTEN	68 HLBC + BR, SEE NOTE E	
GARLAND	68 OSPK + BR, L+ = 0	
ZELLER	69 ASPK + BR, L+ = -0.23	
BOTTERL	70 OSPK + BR, L+ = -0.45 ± 0.15	
HAIDT	71 HLBC + BR, L+ = -0.28 ± 0.16	
CHIANG	72 OSPK + BR, L+ = -0.3 ± 0.10	
HEINTZE	77 CNTR + BR, L+ = -0.29	
<b>L+ FIT</b>	<b>-0.35 ± 0.15</b>	<b>FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.6)</b>

FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.  
 B BOTTERL 70 IS REEVALUATION OF BOTTERL 2 68 WITH DIFFERENT L+.  
 E EICHTEN 68 HAS L+ = -0.23 ± 0.08, T = 4. INSTEAD OF L+ REPL. BY HAIDT 71.  
 H H CALCULATED BY US FROM LO AND L+ GIVEN BELOW.

## ξ\_C = f\_-/f\_+ (DETERMINED FROM μ POLARIZATION IN Kμ₃)

THE MU POLARIZATION IS A MEASURE OF XI(T), NO ASSUMPTIONS ON L+ NECESSARY, T (WEIGHTED BY SENSITIVITY TO XI(0)) SHOULD BE SPECIFIED. IN L+, XI(0) PARAMETERIZATION THIS IS XI(0) FOR L+ = 0. DXI/DL XI(T).

Author	Value	χ²
BORREANI	65 HLBC + POLARIZATION	
CUTTS	65 OSPK + LONG. POL.	
CALLAHA	66 FRBC + TOTAL POL.	
CALLAHA	66 FRBC + LONG. POL.	
BETTELS	68 HLBC + TOTAL POL. T = 4.9	
CUTTS	69 OSPK + TOTAL POL. T = 4.0	
MERLAN	74 ASPK + POL, DXI/DL = -1.7	
BRAUN	75 HLBC + POL. T = 4.2	
<b>L+ AVG</b>	<b>-0.95 ± 0.21</b>	<b>AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)</b>
<b>L+ FIT</b>	<b>-0.35 ± 0.15</b>	<b>FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.6)</b>

FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.  
 T VALUE NOT GIVEN.  
 B BETTELS 68 DXI/DL XI(T) = -1.0 ± 0.9 ± 4.9.  
 C CUTTS 69 T = 4.0 WAS CALCULATED FROM FIG.8. DXI/DL XI(T) = -0.95 ± 4.9 ± 3.8.  
 M MERLAN 74 POLARIZATION RESULT (FIG.5) NOT POSSIBLE. SEE DISCUSSION  
 M OF POLARIZATION EXPERIMENTS IN NOTE ON KL3 FORM FACTORS ABOVE.  
 D BRAUN 75 DXI/DL XI(T) = -2.5 ± 4.2 ± 1.0.

## IMAGINARY PART OF ξ (TEST OF T REVERSAL)

Author	Value	χ²
CALLAHA	66 FRBC + MU	
CALLAHA	66 FRBC + TOTAL POL.	
CALLAHA	66 FRBC + LONG. POL.	
BETTELS	68 HLBC + TOTAL POL.	
CUTTS	69 OSPK + TOTAL POL. FIG.7	
CAMPBELL	81 CNTR + POL.	
BLATT	83 CNTR POLARIZATION	
COMBINED	RESULT OF MORSE 80 (KMUS) AND CAMPBELL 81 (K+MUS).	
<b>IMX AVG</b>	<b>-0.017 ± 0.025</b>	<b>AVERAGE</b>

## λ₊ (LINEAR ENERGY DEPENDENCE OF f₊ IN Kμ₃ DECAY)

SEE ALSO THE CORRESPONDING ENTRIES AND FOOTNOTES IN SECTIONS XIA, XIC, AND LO.

Author	Value	χ²
CALLAHA	66 FRBC + PI	
KLJEWSKI	69 OSPK + PI	
HAIDT	71 HLBC + DP	
ANKENBRAN	72 ASPK + PI	
CHIANG	72 OSPK + DP	
BRAUN	73 HLBC + DP	
ARNOLD	74 HLBC + DP	
MERLAN	74 ASPK + DP	
WHITMAN	80 SPEC + DP	
<b>L+M AVG</b>	<b>0.0272 ± 0.0072</b>	<b>AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)</b>
<b>L+M FIT</b>	<b>0.033 ± 0.008</b>	<b>FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.6)</b>

FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.  
 A ANKENBRANDT 72 L+ FROM FIG.3 TO MATCH DXI/DL. TEXT GIVES .024 ± 0.022

## λ₊ (LINEAR ENERGY DEPENDENCE OF f₊ IN Kμ₃ DECAY)

FOR RAD. COR. OF KE3 DP SEE GINSBURG 67 AND BECHERRAWY 70.

Author	Value	χ²
BROWN	62 XEBC + PI, NO RC	
JENSEN	64 XEBC + PI, NO RC	
BORREANI	64 HBC + E+, NO RC	
BELLOTT2	67 FBC + DP, USES RC	
IMLAY	67 OSPK + DP, NO RC	
BOTTERL1	68 ASPK + E+, USES RC	
EISLER	68 HLBC + PI, USES RC	
BOTTERL	70 OSPK + E+, USES RC	
STEINER	71 HLBC + DP, USES RC	
CHIANG	72 OSPK + DP, NO RC	
BRAUN2	73 HLBC + DP, NO RC	
BRAUN	74 HLBC + KMUS/KE3 VS. T	
BRAUN2	73 STATES THAT RC OF GINSBURG 67 WOULD LOWER L+ BY .002 BUT	
BRAUN	74 IS A COMBINED KMUS-KE3 RESULT. IT IS NOT INDEPENDENT OF	
BRAUN	73 (KMUS) AND BRAUN2 73 (KE3) FORM FACTOR RESULTS.	
<b>L+M AVG</b>	<b>0.0285 ± 0.0043</b>	<b>AVERAGE</b>

## |f₊/f₊| FOR Kμ₃ DECAY

RATIO OF SCALAR TO F+ COUPLINGS

Author	Value	χ²
BELLOTT2	67 HLBC	
KALMUS	67 HLBC +	
BOTTERL1	68 ASPK	
STEINER	71 HLBC + L+, FS, FT, PHI FIT	
CHIANG	72 OSPK +	
BRAUN	75 HLBC +	
<b>FS AVG</b>	<b>0.125 ± 0.044</b>	<b>AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)</b>

## |f\_T/f₊| FOR Kμ₃ DECAY

RATIO OF TENSOR TO F+ COUPLINGS

Author	Value	χ²
BELLOTT2	67 HLBC	
KALMUS	67 HLBC +	
BOTTERL1	68 ASPK	
STEINER	71 HLBC + L+, FS, FT, PHI FIT	
CHIANG	72 OSPK +	
BRAUN	75 HLBC +	
<b>FT AVG</b>	<b>0.22 ± 0.14</b>	<b>AVERAGE</b>

## f\_T/f₊ FOR Kμ₃ DECAY

RATIO OF TENSOR TO F+ COUPLINGS

Author	Value	χ²
BRAUN	75 HLBC	

For notation, see key on page 91.

Stable Particle Full Listings

K±

DECAY FORM FACTORS FOR K± → π+π-e±ν

Table with 2 columns: Key (KE4) and Content (GIVEN IN THE FOLLOWING PAPERS)

Table with 2 columns: Key (BASILE) and Content (71 ASPK +, 73 OSPK +, 77 SPEC +)

REFERENCES FOR K±

Main list of references on the left page, including entries for Birge, Callahan, Cameron, Greiner, Jensen, Kalmus, etc.

Continuation of references on the right page, including entries for Basile, Bourquin, Haidt, Klems, Kunselma, Romano, Schweinb, Steiner, etc.

Continuation of references on the right page, including entries for R. Kunselman, Brannon, Clark, Edwards, Ford, Hoffmast, etc.



# Stable Particle Full Listings

$K^0, K_S^0$

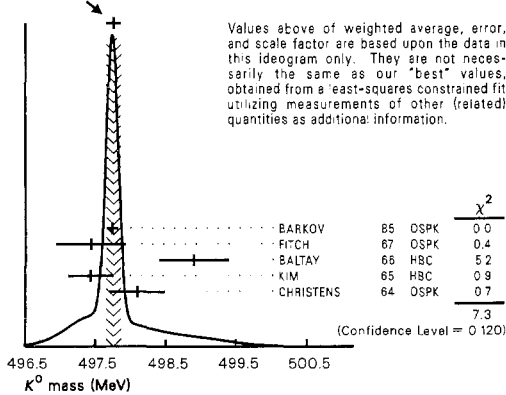
$K^0$

$$I(J^P) = \frac{1}{2}(0^-)$$

### $K^0$ MASS (MeV)

M	498.1	0.4	CHRISTENS	64 OSPK	
M	2223	497.44	0.33	KIM	65 HBC KD FROM PBAR P
M	4500	498.9	0.5	BALTAY	66 HBC KD FROM PBAR P
M	497.44	0.50	FITCH	67 OSPK	
M	780	497.742	0.085	BARKOV	85 OSPK E+ E- TO KL KS
M	AVG	497.76	0.11	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)	
M	FIT	497.718	0.069	FROM FIT (SEE IDEOGRAM BELOW)	

WEIGHTED AVERAGE  
497.76 ± 0.11 (ERROR SCALED BY 1.4)



### $K^0 - K^\pm$ MASS DIFFERENCE (MeV)

D	3.9	0.6	ROSENFELD	59 HBC	-	
D	5.4	1.1	CRAWFORD	59 HBC	-	
D	9	3.90	0.25	BURNSTEIN	65 HBC	-
D	7	3.71	0.35	KIM	65 HBC	K- P TO K0 N
D	417	3.95	0.21	HILL	68 HBC	K+0 TO KOPP
D	AVG	3.92	0.14	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)		
D	FIT	4.051	0.070	FROM FIT		

### REFERENCES FOR $K^0$

CRAWFORD	59 PRL	2	112	CRAWFORD, CRESTI, GOOD, STEVENSON, TICHQ	(LRL)
ROSENFEL	59 PRL	2	110	A H ROSENFELD, F SOLWITZ, R D TRIPP	(LNL)
CHRISTEN	64 PRL	13	138	CHRISTENSON, CROWIN, FITCH, TURLEY	(PRINCETON)
BURNSTEIN	65 PR	138	B 895	R A BURNSTEIN, H A RUBIN	(MARYLAND)
KIM	65 PR	140	B 1334	J K KIM, L KIRSCH, D MILLER	(COLUMBIA)
BALTAY	66 PR	142	932	BALTAY, SANDWEISS, STONEHILL	(YALE-BNL)
FITCH	67 PR	164	1711	FITCH, ROTH, RUSS, VERNON	(PRINCETON)
HILL	68 PR	168	1534	HILL, ROBINSON, SAKITT, CANTER	(BNL, CARNEGIE)
BARKOV	85 JETPL	42	159	-BLINOV, VASSERMAN	(NOVO)

$K_S^0$

$$I(J^P) = \frac{1}{2}(0^-) \quad \text{SHORT-LIVED } K^0$$

### $K_S^0$ MEAN LIFE (units $10^{-10}$ sec)

T	KS MEAN LIFE (PRE-1971 EXPERIMENTS)						
T	0	90	(1.07)	(0.13)	BOLDT	58 CC	
T	0	512	0.94	0.05	0.05	CRAWFORD	59 HBC
T	0	63	(1.09)	(0.18)	(0.15)	BOWEN	60 CC
T	0	OLD EXPTS WITH LOW STATISTICS	NOT INCLUDED IN AVERAGE.				
T	378	0.94	0.05	0.05	BERTANZA	62 HBC	
T	503	0.87	0.05		CHRISTEN	63 HBC	
T	545	0.86	0.04		KREISLER	64 OSPK	
T	572	0.866	0.016		ALFF-STEI	66 OSPK	
T	572	0.90	0.06	0.05	AUERSACH	66 OSPK	
T	4500	0.92	0.04		BALTAY	66 HBC	
T	B	(0.904)	(0.024)		BOTT-BODE	66 OSPK	
T	5000	0.843	0.013		KIRSCH	66 HBC	
T	19994	0.856	0.008		DONALD	68 HBC	
T	H	20000	0.872	0.009		HILL	68 HBC
T	AVG	0.8641	0.0065	0.0065	AVERAGE (ERROR INCL. SCALE FACTOR OF 1.3)		

T	KS MEAN LIFE (POST-1971 EXPERIMENTS)					
T	THESE VALUES ARE USED TO DETERMINE THE STABLE PARTICLE SUMM. TABLE					
T	VALUES OF THE KS MEAN LIFE AND RATES.					
T	H	50K	0.8958	0.0045	SKJEGGESTAD	72 HBC
T	F	2173	(0.867)	(0.024)	FACKLER	73 OSPK
T	G	6M	0.8957	0.0048	GEWERTIGER	74 ASPK
T	C		0.8913	0.0032	CARITHERS	75 SPEC
T	A	26K	0.881	0.009	ARONSON	76 SPEC
T			(0.905)	(0.007)	ARONSON	82 SPEC
T	AVG		0.8923	0.0022	0.0022 AVERAGE	
T	FIT		0.8923	0.0022	FROM FIT	

COMMENTS  
HILL 68 HAS BEEN CHANGED BY THE AUTHORS FROM THE PUBLISHED VALUE (0.865±0.009) BECAUSE OF A CORRECTION IN THE SHIFT DUE TO ETA--. SKJEGGESTAD 72 AND HILL 68 GIVE DETAILED DISCUSSIONS OF SYSTEMATICS ENCOUNTERED IN THIS TYPE OF EXPERIMENT.  
B KS MEAN LIFE NOT THE PRIMARY QUANTITY MEASURED IN THIS EXPT.  
FACKLER 73 DOES NOT INCLUDE SYSTEMATIC ERRORS.  
CARITHERS 75 VALUE IS FOR KL-KS MASS DIFFERENCE DM=5348±.0021. C THE DM DEPENDENCE OF THE TOTAL DECAY RATE (INVERSE MEAN LIFE) IS GAMMA(KS)=((1.122±.004)+16\*(DM-5348)/DM)\*10\*\*10 /SEC.  
C VALUE WOULD NOT CHANGE WITH OUR CURRENT DM=5349±.0022.  
A ARONSON 82 FIND THAT KS MEAN LIFE MAY DEPEND ON THE KAON ENERGY.

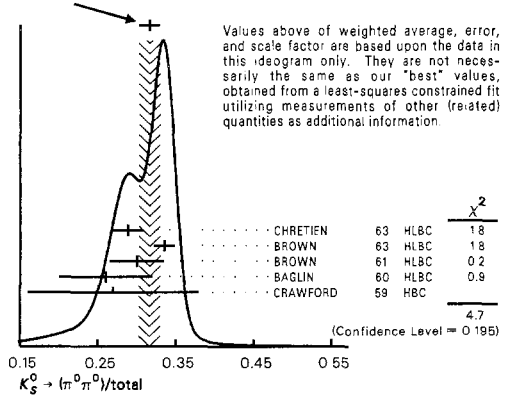
### $K_S^0$ PARTIAL DECAY MODES

P1	$K_S^0 \rightarrow \pi^+ \pi^-$	DECAY MASSES
P2	$K_S^0 \rightarrow \pi^0 \pi^0$	140+ 140
P3	$K_S^0 \rightarrow \mu^+ \mu^-$	135+ 135
P4	$K_S^0 \rightarrow e^+ e^-$	106+ 106
P5	$K_S^0 \rightarrow \pi^+ \pi^- \gamma$	140+ 140+ 0
P6	$K_S^0 \rightarrow \gamma \gamma$	0+ 0
P7	$K_S^0 \rightarrow 3\pi^0$	135+ 135+ 135
P8	$K_S^0 \rightarrow \pi^+ \pi^- \pi^0$	140+ 140+ 135

### $K_S^0$ BRANCHING RATIOS

$K_S^0 \rightarrow (\pi^+ \pi^-)/total$				(P1)		
R1	0.68	0.04	CRAWFORD	59 HBC		
R1	0.70	0.08	COLUMBIA	60 HBC		
R1	U	(0.740)	(0.024)	ANDERSON	62 HBC	
R1	U	3447	0.670	0.010	DOYLE	69 HBC
R1	U	ANDERSON RESULT NOT PUBLISHED, EVENTS ADDED TO DOYLE SAMPLE			PI- P TO LAM K0	
R1	AVG	0.6710	0.0096	AVERAGE		
R1	FIT	0.6861	0.0024	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)		
$K_S^0 \rightarrow (\pi^0 \pi^0)/total$				(P2)		
R2	0.27	0.11	CRAWFORD	59 HBC		
R2	0.26	0.06	BAGLIN	60 HBC		
R2	0.30	0.035	BROWN	61 HBC		
R2	1066	0.335	0.014	BROWN	63 HBC	
R2	198	0.288	0.021	CHRETIEN	63 HBC	
R2	AVG	0.316	0.014	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)		
R2	FIT	0.3139	0.0024	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1) (SEE IDEOGRAM BELOW)		

WEIGHTED AVERAGE  
0.316 ± 0.014 (ERROR SCALED BY 1.3)



(Confidence Level = 0.120)

(Confidence Level = 0.195)

For notation, see key on page 91.

Stable Particle Full Listings

K<sub>S</sub><sup>0</sup>

Table with columns for K<sub>S</sub><sup>0</sup> → (π<sup>+</sup> π<sup>-</sup>)(π<sup>0</sup> π<sup>0</sup>) (P1)/(P2) and rows R3 N 267 (2.12) (0.17) BOZOKI 69 HLBC, R3 G 3016 (2.285) (0.055) GOBBI 69 OSPK, R3 3700 2.10 0.06 MORFIN 69 HLBC, R3 G 7944 2.282 0.045 MOFFETT 70 OSPK, R3 B 6150 2.22 0.095 BALTAY 71 HBC K-P TO KO +NEUTRALS, R3 A 3068 2.22 0.10 ALITTI 72 HBC K-P TO PI+ P KO, R3 6380 2.22 0.08 MORSE 72 DBC K-N TO KOP, R3 N 701 2.10 0.11 NAGY 72 HLBC K-N TO KOP, R3 4799 2.16 0.08 HILL 73 DBC K-D TO KO P P, R3 16K 2.169 0.094 COWELL 74 OSPK PI- P TO LAM KO, R3 1315 2.11 0.09 EVERHART 76 WIRE PI- P TO LAM KO, R3 N NAGY 72 IS A FINAL RESULT WHICH INCLUDES BOZOKI 69, R3 G MOFFETT 70 IS A FINAL RESULT WHICH INCLUDES GOBBI 69, R3 B THE DIRECTLY MEASURED QUANTITY IS KS TO PI+PI-/ALL KOBAR=.345+-0.005, R3 A THE DIRECTLY MEASURED QUANTITY IS KS TO PI+PI-/ALL KO=.345+-0.005, R3 AVG 2.197 0.026 AVERAGE FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1), R3 FIT 2.186 0.025

Table with columns for K<sub>S</sub><sup>0</sup> → (μ<sup>+</sup> μ<sup>-</sup>)/(π<sup>+</sup> π<sup>-</sup>) (units 10<sup>-5</sup>) (P3)/(P1) and rows R4 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT, ALLOWED BY FIRST ORDER WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION., R4 (10.0) OR LESS CL=.90 BOTT-DEBE 67 OSPK, R4 (20.0) OR LESS CL=.90 BOHM 69 OSPK, R4 (1.07) OR LESS CL=.90 HYAMS 69 OSPK, R4 S (32.6) OR LESS CL=.90 STUTZKE 69 OSPK, R4 0.047 OR LESS CL=.90 GJESDAL 73 ASPK, R4 S VALUE CALCULATED BY US, USING 2.3 INSTEAD OF 1 EVENT, 90 PERC.CL.

Table with columns for K<sub>S</sub><sup>0</sup> → (π<sup>+</sup> π<sup>-</sup> γ)/(π<sup>+</sup> π<sup>-</sup>) (units 10<sup>-3</sup>) (P5)/(P1) and rows R5 27 NO RATIO GIVEN BELLOTTI 66 HBC PG GT 50 MEV/C, R5 10 3.3 1.2 WEBBER 70 HBC PG GT 50 MEV/C, R5 B 2.8 0.6 BURGUN 73 HBC PG GT 50 MEV/C, R5 C 29 (3.0) (0.6) BOBISUT 74 HLBC PG GT 40 MEV/C, R5 T 2.68 0.15 TAUREG 76 SPEC PG GT 50 MEV/C, R5 B BURGUN 73 ESTIMATES THAT DIRECT EMISSION CONTRIBUTION IS .3+-0.6, R5 C BOBISUT 74 NOT INCLUDED IN AVERAGE BECAUSE P(GAMMA) CUT DIFFERS., R5 C ESTIMATES DIRECT EMISSION CONTRIBUTION TO BE 0.5 OR LESS, CL=.95., R5 T TAUREG 76 FIND DIRECT EMISSION CONTRIB LT .06, CL=.90., R5 AVG 2.70 0.14 AVERAGE

Table with columns for K<sub>S</sub><sup>0</sup> → (e<sup>+</sup> e<sup>-</sup>)/(π<sup>+</sup> π<sup>-</sup>) (units 10<sup>-5</sup>) (P4)/(P1) and rows R6 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT, ALLOWED BY FIRST ORDER WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION., R6 50.0 OR LESS CL=.90 BOHM 69 OSPK

Table with columns for K<sub>S</sub><sup>0</sup> → 2γ/total (units 10<sup>-3</sup>) (P6) and rows R7 R 0 (21.0) OR LESS CL=.90 BANNER 69 OSPK, R7 R 0 (2.2) OR LESS CL=.90 REPELLIN 71 OSPK, R7 R 0 (0.71) OR LESS CL=.90 BANNER 72 OSPK, R7 R 0 (2.0) OR LESS CL=.90 MORSE 72 DBC, R7 R 0 0.4 OR LESS CL=.90 BARMIN 73 HLBC, R7 R THESE LIMITS ARE FOR MAXIMUM INTERFERENCE IN KS-KL TO 2GAMMAS

Table with columns for K<sub>S</sub><sup>0</sup> → (π<sup>+</sup> π<sup>-</sup> π<sup>0</sup>)/total (units 10<sup>-4</sup>) (P8) and rows R8 (0.85) OR LESS CL=.90 METCALF 72 ASPK, R8 0.49 OR LESS CL=.90 BARMIN 85 HLBC K+ 850 MEV

Table with columns for K<sub>S</sub><sup>0</sup> → (π<sup>0</sup> π<sup>0</sup> π<sup>0</sup>)/total (units 10<sup>-4</sup>) (P7) and rows R9 (4.3) OR LESS CL=.90 BARMIN1 73 HLBC, R9 0.37 OR LESS CL=.90 BARMIN 83 HLBC

NOTE ON CP VIOLATION IN K<sub>S</sub><sup>0</sup> → 3π

For K<sub>S</sub><sup>0</sup> → 3π, the quantities which measure CP violation are the ratios of amplitudes

η<sub>+-0</sub> = A<sub>S</sub>(K<sub>S</sub><sup>0</sup> → π<sup>+</sup> π<sup>-</sup> π<sup>0</sup>) / A<sub>L</sub>(K<sub>L</sub> → π<sup>+</sup> π<sup>-</sup> π<sup>0</sup>)

η<sub>000</sub> = A<sub>S</sub>(K<sub>S</sub><sup>0</sup> → π<sup>0</sup> π<sup>0</sup> π<sup>0</sup>) / A<sub>L</sub>(K<sub>L</sub> → π<sup>0</sup> π<sup>0</sup> π<sup>0</sup>)

If one assumes that CPT invariance holds and that there are no transitions to I = 3 states, then Re(η<sub>+-0</sub>) and Re(η<sub>000</sub>) can be neglected, and CP violation would be observed as nonzero values of Im(η<sub>+-0</sub>) and Im(η<sub>000</sub>). Sections ET+ and ET0 list the relative rates

(Im η<sub>+-0</sub>)<sup>2</sup> = Γ(K<sub>S</sub><sup>0</sup> → π<sup>+</sup> π<sup>-</sup> π<sup>0</sup>) / Γ(K<sub>L</sub> → π<sup>+</sup> π<sup>-</sup> π<sup>0</sup>)

(Im η<sub>000</sub>)<sup>2</sup> = Γ(K<sub>S</sub><sup>0</sup> → π<sup>0</sup> π<sup>0</sup> π<sup>0</sup>) / Γ(K<sub>L</sub> → π<sup>0</sup> π<sup>0</sup> π<sup>0</sup>)

obtained under the above assumptions.

In the above expressions the three pions are restricted to the dominant symmetric I = 1 state, a CP = -1 state which couples to K<sub>S</sub> only if CP is violated. The decay K<sub>S</sub><sup>0</sup> → π<sup>+</sup> π<sup>-</sup> π<sup>0</sup> also has CP-allowed amplitudes to I = 0 and I = 2 states of the three pions. The angular momenta in these states cannot be S wave so they are strongly suppressed by centrifugal barrier effects, and, for the I = 2 state, by the ΔI = 1/2 rule as well. For comparison with the CP-violating rate, we list in section RSQ the CP-conserving rate relative to K<sub>L</sub> → π<sup>+</sup> π<sup>-</sup> π<sup>0</sup> decay. The CP-conserving limit is large and thus determines the branching ratio limit Γ(K<sub>S</sub><sup>0</sup> → π<sup>+</sup> π<sup>-</sup> π<sup>0</sup>)/total given in section R8.

CP VIOLATION PARAMETERS IN K<sub>S</sub><sup>0</sup> DECAY

Table with columns for Im(η<sub>+-0</sub>)<sup>2</sup> and rows ET+ WHERE ETA+-0 = A(KS → PI+ PI- PI0, CP VIOL.)/A(KL → PI+ PI- PI0), ET+ CPT ASSUMED VALID - (1.E RE(ETA+-0)=0), ET+ 18 (3.8) OR LESS CL=.90 ANDERSON 65 HBC INCL. IN WEBBER 70, ET+ (0.45) OR LESS CL=.90 BEHR 66 HLBC, ET+ 71 (0.8) OR LESS CL=.90 WEBBER 70 HBC, ET+ 99 (1.2) OR LESS CL=.90 CHO 71 DBC, ET+ 98 (1.0) OR LESS CL=.90 JAMES 71 HBC INCL. IN JAMES 72, ET+ M 50 (1.2) OR LESS CL=.95 MEISNER 71 HBC CL=.9 NOT AVAIL., ET+ 180 (0.66) OR LESS CL=.90 JAMES 72 HBC, ET+ 99 (1.2) OR LESS CL=.90 JONES 72 OSPK, ET+ 384 0.12 OR LESS CL=.90 METCALF 72 ASPK, ET+ 148 (0.71) OR LESS CL=.90 MALLARY 73 OSPK RE(A)=-.05--.17, ET+ 192 (1.2) OR LESS CL=.90 BALDCEOL 75 HLBC, ET+ B 601 (0.23) OR LESS CL=.90 BARMIN 85 HLBC K+ 850 MEV, ET+ M THESE AUTHORS FIND RE(A)=.275+-0.65, ABOVE VALUE AT RE(A)=0, ET+ B BARMIN 85 FIND RE(ETA+-0)=(0.05+-0.17) AND IM(ETA+-0)=(0.15+-0.33), Im(η<sub>000</sub>)<sup>2</sup>, ETO WHERE ETA000 = A(KS INTO 3PI0, CP VIOLATING)/A(KL INTO 3PI0), ETO SEE COMMENTS IN SECTION ET+ ABOVE., ETO THIS LIMIT DETERMINES BRANCHING RATIO R9 ABOVE., ETO 22 (1.2) OR LESS CL=.90 BARMIN1 73 HLBC, ETO G (0.28) OR LESS CL=.90 GJESDAL 74 SPEC INDIRECT MEAS., ETO B 632 0.12 OR LESS CL=.90 BARMIN 85 HLBC, ETO G GJESDAL 74 USES K2PI, KMU3 AND KE3 DECAY RESULTS, UNITARITY, AND ETO G CPT. CALCULATES ABS(ETA000)=.26+-0.20. WE CONVERT TO UPPER LIMIT., ETO B BARMIN 85 FIND RE(ETA000)=(0.08+-0.18) AND IM(ETA000)=(0.05+-0.27), ETO B ASSUMING CPT INVARIANCE THEY OBTAIN THE LIMIT QUOTED ABOVE., REFERENCES FOR K<sub>S</sub><sup>0</sup>, BOLDT 58 PRL 1 150 E BOLDT, D O CALDWELL, Y PAL (MIT), CRAWFORD 59 PRL 2 266 CRAWFORD, CRESTI, DOUGLASS, GOOD, TICHO + (LRL), BAGLIN 60 NC 18 1043 BAGLIN, BLOCH, BRISSON, HENNESSY + (EPOL), BOWEN 60 PR 119 2030 BOWEN, HARDY, REYNOLDS, SUN, MOORE + (PRIN+BNL), COLUMBIA 60 ROCH CONF 727 M SCHWARTZ + (COLUMBIA), BROWN 61 NC 19 1155 BROWN, BRYANT, BURNSTEIN, GLASER, KADYK. (MICH), ANDERSON 62 CERN CONF 836 J A ANDERSON, F S CRAWFORD + (LRL), BERTANZA 62 PREPRINT 0105 BERTANZA, CONNOLLY, CULWICK, EISLER + (BNL), UNPUBLISHED, BUT RECERTIFIED BY AUTHORS, AUGUST 66., CHRETIEN 63 PR 131 2208 CHRETIEN + (BRANDEIS+BROWN-HARVARD- MIT), BROWN 63 PR 130 769 BROWN, KADYK, TRILLING, ROE - (LRL-MICH), KREISLER 64 PR 136 B 1074 M KREISLER, O OVERSETH, J CRONIN (PRINCETON), CRAWFORD, GOLDEN, STERN, BINFORD + (LRL+MISC), ALFF-STE 66 PL 21 595 ALFF-STEINBERGER, HEUER, KLEINKNECHT + (CERN), AUERBACH 66 PR 149 1052 AUERBACH, DOBBS, LANDE, MANN, SCIULLI + (PENN), ALSO 65 AUERBACH, BALTAY 66 PR 142 932 BALTAY, SANDWEISS, STONEHILL + (CYALE+BNL), BEHR 66 PR 22 540 BEHR, BRISSON, PETHIAU + (EPOL, MILA, PADO, ORSAY), BELLOTTI 66 NC 45A 737 +PULLIA, BALDO-DEOLIN + (MILAN-PADUA), BOTT-BOD 66 PL 23 277 BOTT-BODENHAUSEN, DE BOUARD + (CERN), KIRSCH 66 PR 147 939 L KIRSCH, P SCHMIDT (COLUMBIA), BOTT-BOD 67 PL 24B 194 BOTT-BODENHAUSEN, DE BOUARD, CASSEL + (CERN), DONALD 68 PL 27B 58 DONALD, EDWARDS, NISAR + (LIVP, CERN, IJMP, CDEF), HILL 68 PR 171 1418 HILL, ROBINSON, SAKITT + (BNL, CARNEGIE), BANNER 69 PR 188 2033 +CRONIN, LIU, PILCHER (PRINCETON), BOHM 69 THESIS A. BOHM (AACH), SOZOKI 69 PL 30B 498 +FENYVES, GOMBOSI, NAGY, SURANYI + (BUDAPEST), DOYLE 69 UCRL 18159-THESIS J.C. DOYLE (LRL), GOBBI 69 PRL 22 682 GOBBI, GREEN, HAKEL, MOFFETT, ROSEN + (ROCHESTER), HYAMS 69 PL 29B 521 +KOGH, POTTER, VON LINDERN, LORENZ + (CERN+MP+IM), MORFIN 69 PRL 23 660 MORFIN, SINCLAIR (MICH), STUTZKE 69 PR 177 2009 +ABASHIAN, JONES, MANTSCH, OKS, SMITH (ILLINOIS), MOFFETT 70 BAPS 15 512 +GOBBI, GREEN, HAKEL, ROSEN (ROCHESTER), WEBBER 70 WEBER 70 PR 01 1967 +SOLMITSZ, CRAWFORD, ALSTON-GARNJOIST (LRL), ALSO 69 UCRL 19226 THESIS B R WEBBER (LRL)

Stable Particle Full Listings

K\_S^0, K\_L^0

Table with 3 columns: Author, Particle, Reference. Includes entries like BALTAY 71 PRL 27 1678, ALSO 71 NEVIS-187 THESIS, CHO 71 PR D3 1557, etc.

Table with 3 columns: Author, Particle, Reference. Includes entries like BARMIN1 73 PL 468 465, SARMIN2 73 PL 478 465, BURGUN 73 PL 468 481, etc.

Table with 3 columns: Author, Particle, Reference. Includes entries like BOBISUT 74 LNC 11 646, COMELL 74 PR D10 2085, GEWENIGER 74 PL 483 487, etc.

Table with 3 columns: Author, Particle, Reference. Includes entries like BALDOCEO 75 NC 25A 688, CARITHERS 75 PRL 34 1244, ARONSON 76 NC 32A 236, etc.

Table with 3 columns: Author, Particle, Reference. Includes entries like BJERGE 60 ROCH CONF 601, MULLER 60 PR 4 418, FITCH 61 NC 22 1160, etc.

Table with 3 columns: Author, Particle, Reference. Includes entries like CRAWFORD 62 CERN CONF 827, AUERBACH 65 PRL 14 192, TRILLING 65 UCLR 16473, etc.

K\_L^0

I(J^P) = 1/2(0^-) LONG-LIVED K^0

K\_S^0 - K\_L^0 MASS DIFFERENCE WE GIVE (KL-KS MASS DIFFERENCE / HBAR) IN UNITS OF 10\*\*10 SEC-1

Large table of experimental results for K\_S^0 - K\_L^0 mass difference. Columns include author, value in parentheses, and reference. Includes notes on errors and systematic effects.

K\_L^0 MEAN LIFE (units 10^-8 sec)

Table showing mean life measurements for K\_L^0. Columns: Particle, Value, Error, Reference. Includes entries like KL MEAN LIFE, ASSUMED DS-L00 AND DELTA, etc.

K\_L^0 PARTIAL DECAY MODES

Table of partial decay modes for K\_L^0. Columns: Mode, Branching Ratio, Reference. Includes modes like K\_L^0 -> 3pi^0, K\_L^0 -> pi+ pi- pi^0, etc.

K^0 CONSTRAINED FIT

OVERALL FIT OF MEAN LIFE, WIDTHS AND BRANCHING RATIOS USES 65 DATA POINTS TO DETERMINE SIX QUANTITIES. OVERALL FIT HAS CHI-SQUARED=69.9

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P\_i, as follows: The diagonal elements are P\_i +/- delta P\_i, where delta P\_i = sqrt(delta P\_i^2) / (delta P\_i / P\_i), while the off-diagonal elements are the normalized correlation coefficients (delta P\_i delta P\_j) / (delta P\_i delta P\_j). For the definitions of the individual P\_i, see the listings above; only those P\_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table showing fitted branching fractions P1 through P11 and their correlations.

FITTED PARTIAL DECAY MODE RATES

The matrix below is the branching fraction matrix above, transformed into rate space, i.e., G\_i = 1/Gamma\_total P\_i, in appropriate units. In analogy to the matrix above, the diagonal elements are G\_i +/- delta G\_i, where delta G\_i = sqrt(delta G\_i^2), while the off-diagonal elements are the normalized correlation coefficients (delta G\_i delta G\_j) / (delta G\_i delta G\_j). Note that because of the error in Gamma\_total, the errors and correlations here are not directly derivable from those above.

Table showing fitted decay rates G1 through G11 and their correlations.

K\_L^0 DECAY RATES

Table showing decay rates for K\_L^0 -> pi^0 pi^0 pi^0. Columns: Mode, Rate, Reference. Includes entries like K\_L^0 -> pi^0 pi^0 pi^0 (units 10^6 sec^-1), W1 54 5.22 1.03 0.84 BEHR, etc.

For notation, see key on page 91.

# Stable Particle Full Listings

$K^0_L$

$K^0_L \rightarrow \pi^+ \pi^- \pi^0$  (units  $10^6 \text{ sec}^{-1}$ ) (G2)

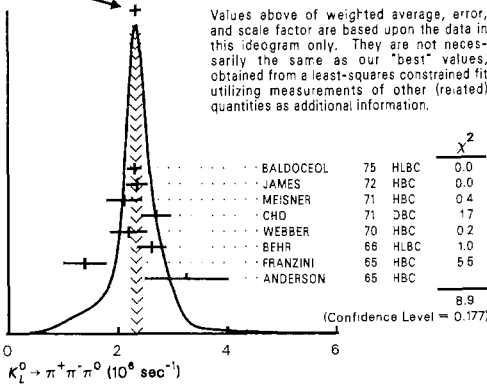
W2	18	3.26	0.77	ANDERSON	65	HBC	
W2	14	1.4	0.4	FRANZINI	65	HBC	
W2	136	2.62	0.28	BEHR	66	HLBC	
W2	53	2.20	0.35	WEBBER	70	HBC	ASSUMES CP
W2	99	2.71	0.28	CHO	71	DBC	ASSUMES CP
W2	98	(2.5)	(0.3)	JAMES	71	HBC	ASSUMES CP
W2	50	2.12	0.33	WEISNER	71	HBC	ASSUMES CP
W2	180	2.35	0.20	JAMES	72	HBC	ASSUMES CP
W2	192	2.32	0.13	BALDOCEOL	75	HLBC	ASSUMES CP

IN THE OVERALL FIT THIS RATE IS WELL DETERMINED BY THE MEAN LIFE AN THE BRANCHING RATIO R2. FOR THIS REASON THE DISCREPANCY BETWEEN THE W2 MEASUREMENTS DOES NOT AFFECT THE SCALE FACTOR OF THE OVERALL FIT W2 JAMES 72 IS A FINAL MEASUREMENT AND INCLUDES JAMES 71.

W2	AVG	2.34	0.11	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.2)
W2	FIT	2.392	0.040	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.2)

(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
2.34 ± 0.11 (ERROR SCALED BY 1.2)



$K^0_L \rightarrow \pi \mu \nu$  (units  $10^6 \text{ sec}^{-1}$ ) (G3)

W6	19	4.54	1.24	1.08	LOWYS	67	HLBC
W6	FIT	5.231	0.087	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.3)		

## $K^0_L$ BRANCHING RATIOS

$K^0_L \rightarrow (\pi^0 \pi^0 \pi^0)/(\pi \ell \nu + \pi^+ \pi^- \pi^0)$  (P1)/(P2+P3+P4)

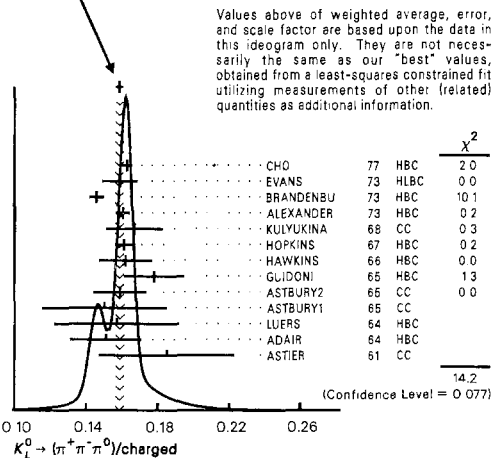
R1	L=E OR MU.						
R1	24	0.24	0.08		ANIKINA	64	CC
R1	549	0.251	0.014		BUDAGOV	68	HLBC
R1	444	0.277	0.021		BUDAGOV	68	HLBC
R1	29	0.31	0.07	0.06	KULYUKINA	68	CC
R1	AVG	0.260	0.011	AVERAGE			
R1	FIT	0.274	0.016	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.8)		

$K^0_L \rightarrow (\pi^+ \pi^- \pi^0)/(\pi \ell \nu + \pi^+ \pi^- \pi^0)$  (P2)/(P2+P3+P4)

R2	L=E OR MU.						
R2	59	0.185	0.038		ASTIER	61	CC
R2	79	0.151	0.020		ADAIR	64	HBC
R2	75	0.157	0.03	0.04	LUERS	64	HBC
R2	66	0.15	0.03	0.04	ASTBURY1	65	CC
R2	326	0.159	0.015		ASTBURY2	65	CC
R2	566	0.178	0.017		GUIDONI	65	HBC
R2	1729	(0.144)	(0.004)		HOPKINS	65	HBC
R2	126	0.162	0.015		HAWKINS	66	HBC
R2		0.161	0.005		HOPKINS	67	HBC
R2	1402	0.167	0.016		KULYUKINA	68	CC
R2	1590	0.1605	0.0038		ALEXANDER	73	HBC
R2	3200	0.146	0.004		BRANDENBU	73	HBC
R2	558	0.159	0.010		EVANS	73	HBC
R2	6499	0.163	0.003		CHO	77	HLBC
R2	AVG	0.1587	0.0024	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.3)		
R2	FIT	0.1585	0.0019	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.2)		

(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
0.1587 ± 0.0024 (ERROR SCALED BY 1.3)



$K^0_L \rightarrow \pi \ell \nu$  (units  $10^6 \text{ sec}^{-1}$ ) (G4)

W3	620	7.52	0.85	0.72	AUBERT	65	HLBC	DS=DQ, CP ASSUMED
W3		7.81	0.56		CHAN	71	HBC	
W3	AVG	7.71	0.46	AVERAGE				
W3	FIT	7.47	0.11	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.3)			

$K^0_L \rightarrow \pi \ell \nu + \pi^+ \pi^- \pi^0$  (units  $10^6 \text{ sec}^{-1}$ ) (G2+G3+G4)

KL INTO CHARGED (UNITS  $10^{**6} \text{ SEC}^{-1}$ ) (G2+G3+G4)

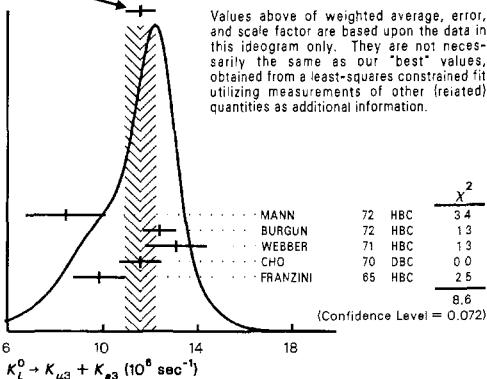
W4	L=E OR MU.						
W4	98	15.1	1.9	AUERBACH	66	OSPK	
W4	FIT	15.09	0.21	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.4)		

$K^0_L \rightarrow \pi \ell \nu$  (units  $10^6 \text{ sec}^{-1}$ ) (G3+G4)

W5	D 109	9.85	1.15	1.05	FRANZINI	65	HBC	K=N TO KO P
W5	C 335	(10.3)	(0.8)		HILL	67	DBC	K=N TO KOP
W5	D 393	11.6	0.9		CHO	70	DBC	K=N TO KOP
W5	D 252	13.1	1.3		WEBBER	71	HBC	K=P TO KOBAR N
W5	D 410	12.4	0.7		BURGUN	72	HBC	K=P TO KOPPI+
W5	D 126	8.47	1.69		MANN	72	HBC	K=P TO KOBAR N
W5	C CHO 70	INCLUDES EVENTS OF HILL 67						
W5	D	ASSUMES DS=DQ RULE						
W5	AVG	11.60	0.65	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.5)			
W5	FIT	12.70	0.18	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.4)			

(SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
11.60 ± 0.65 (ERROR SCALED BY 1.5)



$K^0_L \rightarrow (\pi \mu \nu)/(\pi \ell \nu + \pi^+ \pi^- \pi^0)$  (P3)/(P2+P3+P4)

R3	L=E OR MU.						
R3	C 251	(0.356)	(0.07)		LUERS	64	HBC
R3	C 172	(0.39)	(0.08)	(0.10)	ASTBURY1	65	CC
R3	C 330	(0.335)	(0.055)		KULYUKINA	68	CC
R3	C THIS MODE NOT MEASURED INDEPENDENTLY FROM R2 AND R4						
R3	FIT	0.3465	0.0028	FROM FIT			

$K^0_L \rightarrow (\pi \ell \nu)/(\pi \ell \nu + \pi^+ \pi^- \pi^0)$  (P4)/(P2+P3+P4)

R4	L=E OR MU.						
R4	24	0.46	0.11		NEAGU	61	CC
R4	153	0.487	0.05		LUERS	64	HBC
R4	202	0.46	0.08	0.10	ASTBURY1	65	CC
R4	500	0.498	0.052		KULYUKINA	68	CC
R4	AVG	0.485	0.032	AVERAGE			
R4	FIT	0.4950	0.0029	FROM FIT			

$K^0_L \rightarrow (\pi \ell \nu)/(\pi \ell \nu + (\pi \mu \nu))$  (P4)/(P3+P4)

R5	320	0.415	0.120	ASTIER	61	CC	
R5	FIT	0.5882	0.0032	FROM FIT			

$K^0_L \rightarrow (\pi^+ \pi^- \pi^0)/total$  (P2)

R6							
R6	FIT	0.1240	0.0020	FROM FIT			

$K^0_L \rightarrow (\pi \ell \nu)/total$  (P3+P4)

R7							
R7	FIT	0.6584	0.0089	FROM FIT			

# Stable Particle Full Listings

$K_L^0$

$K_L^0 \rightarrow (2\gamma)/total (units 10^{-4})$  (P9)  
 RB C 32 (1.3) (0.6) CRIEGEE 66 OSPK  
 RB R 33 (7.4) (1.6) TODOROFF 67 OSPK REPL. CRIEGEE66  
 RB R 90 5.5 1.1 KUNZ 68 OSPK NORM. TO SPI(C+N)  
 RB R 23 4.5 1.0 ENSTROM 71 OSPK KL 1.5-9 GEV/C  
 RB R 4.5 1.0 REPELLIN 71 OSPK  
 RB B 4.54 0.84 BANNER2 72 OSPK  
 RB B THIS VALUE USES  $(E00/E+)*2-1.05+-0.14$  IN GENERAL, S13R8 =  
 (4.32+-0.55)/(10\*\*4)\*(E00/E+)\*2.  
 RB B ASSUMES REGEN AMPL IN COPPER AT 22 MV. TO EVALUATE  
 FOR A GIVEN REGEN AMPL AND ERROR, MULTIPLY BY (REGEN AMPL/22MB)\*\*2  
 RB C CRIEGEE 66 REPLACED BY TODOROFF 67  
 RB K CROWN1 67 REPLACED BY KUNZ 68.  
 RB R R21 BELOW GIVES (4.82+-0.52)E-4. COMBINED AVG (4.85+-0.37)E-4.  
 RB AVG 4.89 0.54 AVERAGE

$K_L^0 \rightarrow (\pi^+ \pi^-)/(\pi \ell \nu + \pi^+ \pi^- \pi^0) (units 10^{-3})$  (P5)/(P2+P3+P4)  
 R9 L=E OR MU.  
 R9 O 45 (2.0) (0.4) CHRISTENS 64 OSPK ETA+- = 1.95+-0.20  
 R9 O 54 (2.08) (0.35) GALBRAITH 65 OSPK ETA+- = 1.99+-0.16  
 R9 O (1.93) (0.26) BASILLE 66 OSPK ETA+- = 1.92+-0.13  
 R9 O (0.99) (0.080) BOTT-BODE 66 OSPK ETA+- = 1.95+-0.04  
 R9 M 4200 (2.60) (0.07) MESSNER 73 ASPK ETA+- = 2.23+-0.05  
 R9 O OLD EXPERIMENTS EXCLUDED FROM FIT. SEE SUBSECTION E+- BELOW FOR  
 R9 O AVERAGE ETA+- OF THESE EXPERIMENT AND FOR NOTE ON DISCREPANCY.  
 R9 M FROM SAME DATA AS R27 MESSNER 73, BUT WITH DIFFERENT NORMALIZATION.  
 R9 FIT 2.598 0.053 FROM FIT

$K_L^0 \rightarrow (\pi \mu \nu)/(\pi \ell \nu)$  (P3)/(P4)  
 R10 0 0.81 0.19 ADAIR 64 HBC  
 R10 0 0.82 0.10 DEBOUARD 67 OSPK  
 R10 273 0.7 0.2 HAWKINS 67 HBC  
 R10 0 0.81 0.08 HOPKINS 67 HBC  
 R10 770 0.71 0.05 BUDAGOV 68 HLCB  
 R10 K (0.67) (0.13) KULYUKINA 68 CC  
 R10 B 569 (0.71) (0.04) BEILLIERE 69 HLCB  
 R10 1309 (0.648) (0.030) EVANS 69 HLCB REPL. BY EVANS 73  
 R10 2548 0.68 0.08 BASILLE 70 OSPK  
 R10 6700 0.74 0.04 BRANDENBU 73 HBC  
 R10 1309 0.662 0.030 EVANS 73 HLCB  
 R10 10K 0.662 0.037 WILLIAMS 74 ASPK  
 R10 35K 0.702 0.011 CHO 80 HBC  
 R10 K KULYUKINA 68: R10 IS NOT MEASURED INDEPENDENTLY FROM R2 AND R4.  
 R10 B BEILLIERE 69 IS A SCANNING EXPT USING SAME EXPOSURE AS BUDAGOV 68  
 R10 AVG 0.7001 0.0093 AVERAGE  
 R10 FIT 0.7001 0.0092 FROM FIT

$K_L^0 \rightarrow (\mu^+ \mu^-)/(\pi \ell \nu + \pi^+ \pi^- \pi^0) (units 10^{-6})$  (P6)/(P2+P3+P4)  
 R11 L=E OR MU.  
 R11 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R11 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R11 (100.0) OR LESS ANIKINA 65 CC  
 R11 (250.0) OR LESS CL=90 ALFF-STEI 66 OSPK  
 R11 (2.0) OR LESS CL=90 BOTT-BODE 67 OSPK  
 R11 (35.0) OR LESS CL=90 FITCH 67 OSPK

$K_L^0 \rightarrow (\pi^+ \pi^- \gamma)/total (units 10^{-3})$  (P10)  
 R12 (15.0) OR LESS ANIKINA 65 CC  
 R12 0 (5.0) OR LESS BELLOTTI 66 HLCB GAM KE 40-130 MV  
 R12 1 (3.0) OR LESS NEFKENS 66 OSPK GAM KE 120 MEV  
 R12 (0.4) OR LESS CL=90 THATCHER 68 OSPK GAM KE 20-170 MV  
 R12 (3.2) OR LESS CL=90 BOBISUT 74 HLCB GAM KE GT 40 MEV  
 R12 D 24 (0.062) (0.021) DONALDS1 74 SPEC  
 R12 H (0.46) OR LESS CL=90 WOOD 74 SPEC  
 R12 H 516 (0.052) (0.016) CARROLL2 80 SPEC +-OGAM KE GT 20 MEV  
 R12 J 546 (0.0289) (0.0028) CARROLL2 80 SPEC +-0  
 R12 K1062 0.0441 0.0032 CARROLL2 80 SPEC +-OGAM KE GT 20 MEV  
 R12 D USES  $\pi^+ \rightarrow \pi^0 \pi^+ \pi^-$  FOR ALL KL DECAYS = 0.126  
 R12 H INTERNAL BREMSSTRAHLUNG COMPONENT ONLY.  
 R12 J DIRECT GAMMA EMISSION COMPONENT ONLY.  
 R12 K BOTH COMPONENTS. USES KL TO  $\pi^+ \pi^- \pi^0$  FOR ALL KL DECAYS = 0.1239.

$K_L^0 \rightarrow (e^+ e^-)/(\pi \ell \nu + \pi^+ \pi^- \pi^0) (units 10^{-6})$  (P7)/(P2+P3+P4)  
 R13 L=E OR MU.  
 R13 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R13 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R13 (100.0) OR LESS ANIKINA 65 CC  
 R13 (200.0) OR LESS CL=90 ALFF-STEI 66 OSPK  
 R13 (23.0) OR LESS CL=90 BOTT-BODE 67 OSPK

$K_L^0 \rightarrow (e \mu)/(\pi \ell \nu + \pi^+ \pi^- \pi^0) (units 10^{-4})$  (P8)/(P2+P3+P4)  
 R14 L=E OR MU.  
 R14 TEST OF LEPTON FAMILY NUMBER CONSERVATION.  
 R14 (10.0) OR LESS ANIKINA 65 CC  
 R14 (1.0) OR LESS CL=90 CARPENTER 66 OSPK  
 R14 (0.1) OR LESS CL=90 BOTT-BODE 67 OSPK  
 R14 0.08 OR LESS CL=90 FITCH 67 OSPK

$K_L^0 \rightarrow (\pi^- e^+ \nu)/(\pi^+ e^- \nu)$   
 R15 D 97 (0.90) (0.18) NEAGU 61 CC  
 R15 D (1.01) (0.16) LUERS 64 HBC  
 R15 D 894 (0.99) (0.025) KULYUKINA 66 CC  
 R15 O 1539 (1.06) (0.05) VERHEY 66 OSPK  
 R15 O LOW PRECISION EXPTS NOT AVERAGED. FOR MORE PRECISE VALUE,  
 R15 O SEE S13A2 IN THE CP VIOLATION SECTION BELOW.

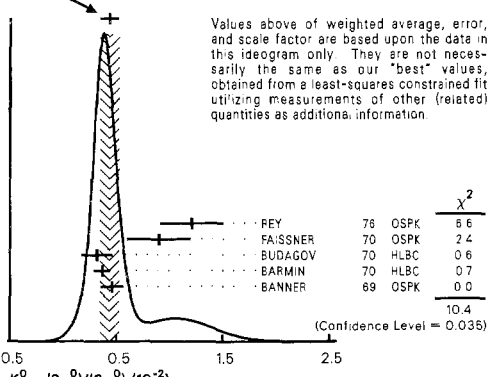
$K_L^0 \rightarrow (\pi^- \mu^+ \nu)/(\pi^+ \mu^- \nu)$   
 R16 IM 1.0081 0.0027 DORFAN 67 OSPK  
 R16 SEE ALSO S13A2 AND S13AL IN THE CP VIOLATION SECTION BELOW.

$K_L^0 \rightarrow (\pi^0 \pi^0)/total (units 10^{-3})$  (P11)  
 R17 C 7 (1.2) (1.5) CRIEGEE 66 OSPK  
 R17 C CRIEGEE EXPT NOT DESIGNED TO MEASURE 2PI0 DECAY MODE  
 R17 B 189 (2.5) (0.8) GAILLARD 69 OSPK E00=3.6+-0.6  
 R17 G LATEST RESULT OF THIS EXPERIMENT GIVEN BY FAISSNER 70 R19  
 R17 FIT 0.94 0.19 FROM FIT

$K_L^0 \rightarrow (3\pi^0)/(\pi^+ \pi^- \pi^0)$  (P1)/(P2)  
 R18 188 2.0 0.6 ALEKSANYA 64 FBC  
 R18 1010 1.80 0.15 BUDAGOV 68 HLCB  
 R18 883 (1.65) (0.07) BARMIN2 72 HLCB ERROR STAT. ONLY  
 R18 R 1.81 0.13 AVERAGE  
 R18 FIT 1.73 0.10 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7)

$K_L^0 \rightarrow (2\pi^0)/(3\pi^0) (units 10^{-2})$  (P11)/(P1)  
 R19 C 109 (1.89) (0.31) CRONIN 1 67 OSPK ETA00=4.9+-0.5  
 R19 C (1.36) (0.18) CRONIN 2 67 OSPK ETA00=3.92+-0.3  
 R19 C CRONIN2 IS FURTHER ANALYSIS OF CRONIN1, NOW BOTH WITHDRAWN  
 R19 NO EVENTS SEEN BARTLETT 68 OSPK SEE E00 BELOW  
 R19 BANNER 69 OSPK ETA00=2.2+-0.3  
 R19 R 133 (1.31) (0.31) CENCE 69 OSPK ETA00=3.7+-0.5  
 R19 29 0.37 0.08 BARMIN 70 HLCB ETA00=2.02+-0.23  
 R19 30 0.32 0.15 BUDAGOV 70 HLCB ETA00=1.9+-0.5  
 R19 F 172 0.90 0.30 FAISSNER 70 OSPK ETA00=3.2+-0.5  
 R19 R 150 1.21 0.30 REY 76 OSPK ETA00=3.8+-0.5  
 R19 F FAISSNER 70 CONTAINS SAME 2PI0 EVENTS AS GAILLARD 69 R17  
 R19 R CENCE 69 EVENTS ARE INCLUDED IN REY 76.  
 R19 AVG 0.437 0.092 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)  
 R19 FIT 0.437 0.085 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.5)  
 (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
 $0.437 \pm 0.092$  (ERROR SCALED BY 1.6)



$K_L^0 \rightarrow (\pi^+ \pi^-)/(\pi \ell \nu) (units 10^{-3})$  (P5)/(P3+P4)  
 R20 O 309 (2.31) (0.23) DEBOUARD 67 OSPK ETA--=2.00+-0.09  
 R20 D 525 (2.35) (0.19) FITCH 67 OSPK ETA--=1.94+-0.08  
 R20 2703 3.04 0.14 DEVOE 77 SPEC ETA--=2.25+-0.05  
 R20 1687 3.13 0.14 COUPLAR 85 SPEC ETA--=2.28+-0.06  
 R20 O OLD EXPERIMENTS EXCLUDED FROM FIT. SEE SUBSECTION E+- BELOW FOR  
 R20 O AVERAGE ETA-- OF THESE EXPERIMENTS AND FOR NOTE ON DISCREPANCY.  
 R20 AVG 3.085 0.099 AVERAGE  
 R20 FIT 3.087 0.066 FROM FIT

$K_L^0 \rightarrow (2\gamma)/(3\pi^0) (units 10^{-3})$  (P9)/(P1)  
 R21 16 2.5 0.7 ARNOLD 68 HLCB VACUUM DECAY  
 R21 115 2.24 0.28 BANNER 69 OSPK  
 R21 28 2.13 0.43 BARMIN 71 HLCB  
 R21 AVG 2.24 0.22 AVERAGE

$K_L^0 \rightarrow (\mu^+ \mu^-)/(\pi^+ \pi^-) (units 10^{-6})$  (P6)/(P5)  
 R22 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R22 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R22 O (140.) OR LESS CL=90 FOETH 69 SPEC  
 R22 O (18.) OR LESS CL=90 DARRIULAT 70 SPEC  
 R22 A O (1.53) OR LESS CL=90 CLARK 71 SPEC  
 R22 C 9 5.8 2.3 1.5 CARITHERS 73 SPEC  
 R22 F 3 4.2 5.1 2.6 FUKUSHIMA 76 SPEC  
 R22 15 4.0 1.4 0.9 SHONET 79 SPEC  
 R22 A CLARK 71 LIMIT RAISED FROM 1.2 E-06 BY FIELD 74 REANALYSIS.  
 R22 A NOT IN AGREEMENT WITH SUBSEQUENT EXPTS. SO NOT AVERAGED.  
 R22 C CARITHERS 73 ERRORS ARE AT CL=0.68. M.CARITHERS, PRIV.COMM. 1979.  
 R22 F FUKUSHIMA 76 ERRORS ARE AT CL=90 PERCENT.  
 R22 AVG 4.47 0.95 AVERAGE

$K_L^0 \rightarrow (e^+ e^-)/(\pi^+ \pi^-) (units 10^{-5})$  (P7)/(P5)  
 R23 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R23 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R23 O 10.0 OR LESS CL=90 FOETH 69 ASPK  
 R23 A (0.10) OR LESS CL=90 CLARK 71 ASPK  
 R23 A POSSIBLE (BUT UNKNOWN) SYSTEMATIC ERRORS. SEE NOTE A IN R22 ABOVE.

$K_L^0 \rightarrow (e \mu)/(\pi^+ \pi^-) (units 10^{-5})$  (P8)/(P5)  
 R24 L=E OR MU.  
 R24 A (0.10) OR LESS CL=90 CLARK 71 ASPK  
 R24 A POSSIBLE (BUT UNKNOWN) SYSTEMATIC ERRORS. SEE NOTE A IN R22 ABOVE.

$K_L^0 \rightarrow (\pi e \nu \gamma)/(K_{e3}) (units 10^{-2})$  (P12)/(P3)  
 R25 10 3.3 2.0 PEACH 71 HLCB GAM KE GT 15 MEV

$K_L^0 \rightarrow (\pi^0 2\gamma)/(3\pi^0) (units 10^{-3})$  (P13)/(P1)  
 R26 0 1.1 OR LESS CL=90 BANNER 69 OSPK

For notation, see key on page 91.

# Stable Particle Full Listings

$K_L^0$

$K_L^0 \rightarrow (\pi^+ \pi^-)(\pi^+ \pi^- \pi^0)$  (units  $10^{-2}$ ) (P5)/(P2)  
 R27 4200 1.64 0.04 MESSNER 73 ASPK ETA-- = 2.23  
 R27 FIT 1.639 0.033 FROM FIT

$K_L^0 \rightarrow (e^+ e^- \gamma)$ /total (units  $10^{-5}$ ) (P14)  
 R28 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R28 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R28 B 0 (2.7) OR LESS CL=.90 BARMINI 72 HLBC  
 R28 C 4 (1.74) OR LESS CL=.87 CARROLLI 80 SPEC +-0  
 R28 D USES KL TO  $\pi^0$  TO  $\pi^0$ /TOTAL=0.214  
 R28 C USES KL TO  $\pi^+$   $\pi^-$   $\pi^0$ /ALL KL DECAYS = 0.1239 .

$K_L^0 \rightarrow (\mu^+ \mu^- \gamma)$ /total (units  $10^{-6}$ ) (P15)  
 R29 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R29 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R29 D (7.81) OR LESS CL=.90 DONALDSON 74 SPEC  
 R29 C 1 0.28 0.28 CARROLLI 80 SPEC +-0  
 R29 D USES KL TO  $\pi^+$   $\pi^-$   $\pi^0$ /ALL KL DECAYS = 0.126 .  
 R29 C USES KL TO  $\pi^+$   $\pi^-$   $\pi^0$ /ALL KL DECAYS = 0.1239 .

$K_L^0 \rightarrow (\mu^+ \mu^- \pi^0)$ /total (units  $10^{-5}$ ) (P16)  
 R30 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R30 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R30 D (5.66) OR LESS CL=.90 DONALDSON 74 SPEC  
 R30 C 0 0.12 OR LESS CL=.90 CARROLLI 80 SPEC  
 R30 D USES KL TO  $\pi^+$   $\pi^-$   $\pi^0$ /ALL KL DECAYS = 0.126 .  
 R30 C USES KL TO  $\pi^+$   $\pi^-$   $\pi^0$ /ALL KL DECAYS = 0.1239 .

$K_L^0 \rightarrow (\pi^+ \pi^- e^+ e^-)$ /total (units  $10^{-6}$ ) (P17)  
 R31 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R31 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R31 D (8.81) OR LESS CL=.90 DONALDSON 76 SPEC  
 R31 C 0 2.5 OR LESS CL=.90 BALATS 83 SPEC  
 R31 D USES KL TO  $\pi^+$   $\pi^-$   $\pi^0$ /ALL KL DECAYS = 0.126 .

$K_L^0 \rightarrow (\pi^0 \pi^\pm e^\mp \nu)$ /total (units  $10^{-3}$ ) (P18)  
 R32 (2.2) OR LESS CL=.90 DONALDSON 74 SPEC  
 R32 16 0.062 0.020 CARROLLI 80 SPEC  
 R32 D DONALDSON 74 USES KL TO  $\pi^+$   $\pi^-$   $\pi^0$ /ALL KL DECAYS = 0.126 .

$K_L^0 \rightarrow ((\pi \mu \text{ atom})\nu)/(\pi \mu \nu)$  ( $10^{-7}$ ) (P19)/(P3)  
 R33 18 SEEN 0.41 COOMBS 76 WIRE  
 R33 155 3.88 0.41 ARONSON 82 SPEC

$K_L^0 \rightarrow (e^+ e^- \pi^0)$ /total (units  $10^{-6}$ ) (P20)  
 R34 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R34 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R34 C 0 2.3 OR LESS CL=.90 CARROLLI 80 SPEC  
 R34 C USES KL TO  $\pi^+$   $\pi^-$   $\pi^0$ /ALL KL DECAYS = 0.1239 .

$K_L^0 \rightarrow (\mu^+ \mu^- e^+ e^-)$ /total (units  $10^{-6}$ ) (P21)  
 R35 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R35 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R35 4.9 OR LESS CL=.90 BALATS 83 SPEC

$K_L^0 \rightarrow (e^+ e^- e^+ e^-)$ /total (units  $10^{-6}$ ) (P22)  
 R36 TEST FOR DELTA(S) = 1 WEAK NEUTRAL CURRENT. ALLOWED BY FIRST ORDER  
 R36 WEAK INTERACTION COMBINED WITH ELECTROMAGNETIC INTERACTION.  
 R36 2.6 OR LESS CL=.90 BALATS 83 SPEC

### $K_L^0$ ENERGY DEPENDENCE OF DALITZ PLOT

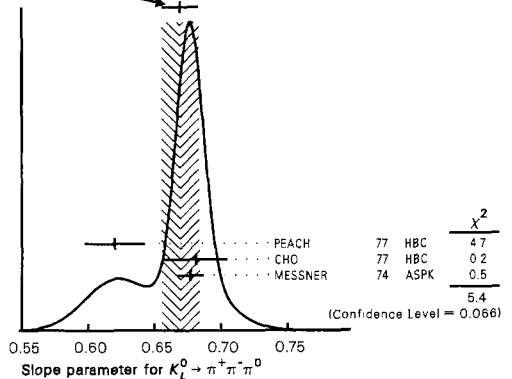
FOR DISCUSSION, SEE NOTE ON SLOPE PARAMETERS IN THE K-- SECTION OF THE FULL LISTINGS ABOVE.

MATRIX ELEMENT SQUARED =  $1 + G_{\mu 0} - H_{\mu 0} + J_{\mu 0} + K_{\mu 0} + V_{\mu 0}$   
 WHERE  $U = (S3-S0)/(MPl**2)$  AND  $V = (S1-S2)/(MPl**2)$

LINEAR COEFFICIENT  $g$  FOR  $|M(K_L^0 \rightarrow \pi^+ \pi^- \pi^0)|^2$

GTO Q	79	(0.55)	(0.23)	ADAIR	64 HBC	AV=-7.6 +- 1.7
GTO Q	77	(0.51)	(0.20)	LUEERS	64 HBC	AV=-7.3 +- 1.6
GTO Q	66	(0.32)	(0.13)	ASTBURY1	65 CC	AV=-5.5 +- 1.5
GTO Q	310	(0.51)	(0.09)	ASTBURY2	65 CC	AV=-7.3 +- 0.8
GTO Q	280	(0.64)	(0.17)	ANIKINA	66 CC	AV=-8.2 +- 0.8
GTO Q	126	(0.70)	(0.12)	HAWKINS	66 HBC	AV=-8.6 +- 0.7
GTO Q	1350	(0.649)	(0.044)	HOPKINS	67 HBC	AT=-0.294 +- .018
GTO Q	1198	(0.428)	(0.055)	NEFKENS	67 OSPK	AU=-0.204 +- .025
GTO Q	2446	(0.400)	(0.045)	BASTILEZ	68 OSPK	AT=-0.188 +- .020
GTO Q	29K	(0.650)	(0.012)	ALBROW	70 OSPK	AY=-0.858 +- .015
GTO Q	36K	(0.505)	(0.022)	BUCHANAN	70 SPEC	AU=-0.278 +- .010
GTO Q	4400	(0.664)	(0.056)	SMITH	70 OSPK	AT=-0.306 +- 0.024
GTO Q	180	(0.50)	(0.11)	JAMES	72 HBC	
GTO Q	1486	(0.681)	(0.043)	KRENZ	72 HLBC	AT=-0.277 +- .018
GTO Q	384	(0.688)	(0.074)	METCALF	72 ASPK	AT=-0.31 +- .03
GTO Q	3200	(0.73)	(0.04)	BRANDENBU	73 HBC	
GTO Q	20K	(0.619)	(0.027)	BISI	74 ASPK	AT=-0.282 +- .011
GTO Q	509K	0.677	0.010	MESSNER	74 ASPK	AY=-0.917 +- .013
GTO Q	192	(0.69)	(0.07)	BALDOCELO	75 HLBC	
GTO Q	56K	(0.590)	(0.022)	BUCHANAN	75 SPEC	AU=-0.277 +- .010
GTO	H6499	0.681	0.024	CHO	77 HBC	
GTO	4709	0.620	0.023	PEACH	77 HBC	
GTO Q	QUADRATIC DEPENDENCE REQUIRED BY SOME EXPERIMENTS (SEE SECTIONS GTO Q AND KTO BELOW). CORRELATIONS PREVENT US FROM AVERAGING RESULTS					
GTO Q	OF FITS NOT INCLUDING C, H, AND K TERMS.					
GTO B	BUCHANAN 70 RESULT REVISED BY BUCHANAN 75 TO INCLUDE RADIATIVE COR.					
GTO B	AND TO USE MORE RELIABLE KL MOM.SPECT. OF 2ND EXPT.(HAD SAME BEAM).					
GTO C	81S1 74 VALUE COMES FROM QUADRATIC FIT WITH QUAD. TERM CONSISTENT					
GTO C	WITH ZERO. GTO ERROR IS THUS LARGER THAN IF LINEAR FIT WERE USED.					
GTO						
GTO AVG	0.670	0.014	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)	(SEE IDEOGRAM BELOW)		

WEIGHTED AVERAGE  
 0.670 ± 0.014 (ERROR SCALED BY 1.6)



### QUAD. COEFFICIENT $h$ FOR $|M(K_L^0 \rightarrow \pi^+ \pi^- \pi^0)|^2$

HTO Q	29K	(-0.011)	(0.018)	ALBROW	70 ASPK
HTO Q	4400	(0.0490)	(0.032)	SMITH	70 OSPK
HTO	509K	0.079	0.007	MESSNER	74 ASPK
HTO	6499	0.095	0.032	CHO	77 HBC
HTO	4709	0.048	0.026	PEACH	77 HBC
HTO	SEE NOTES IN SECTION GTO ABOVE.				
HTO AVG	0.0786	0.0067	AVERAGE		

### QUAD. COEFFICIENT $k$ FOR $|M(K_L^0 \rightarrow \pi^+ \pi^- \pi^0)|^2$

KTO	509K	0.0097	0.0018	MESSNER	74 ASPK
KTO	4709	-0.008	0.012	PEACH	77 HBC
KTO	AVERAGE				
KTO AVG	0.0098	0.0018	AVERAGE		

### LINEAR COEFFICIENT $j$ FOR $K_L^0 \rightarrow \pi^+ \pi^- \pi^0$ CP VIOLATING TERM

JTR LISTED IN CP VIOLATION SECTION BELOW.

### $K_L^0$ FORM FACTORS

FOR DISCUSSION, SEE NOTE ON FORM FACTORS IN THE K-- SECTION OF THE FULL LISTINGS ABOVE.

IN THE FORM FACTOR COMMENTS, THE FOLLOWING ABBREVIATIONS ARE USED.  
 F+ AND F- ARE FORM FACTORS FOR THE VECTOR MATRIX ELEMENT.  
 FS AND FT REFER TO THE SCALAR AND TENSOR TERM.  
 F+(CF+) = (F-)\*T/(MK\*\*2-MPl\*\*2)  
 L+, L- AND LO ARE THE LINEAR EXPANSION COEFFS. OF F+, F- AND FO.  
 L+ REFERS TO THE KMUS VALUE EXCEPT IN THE KE3 SECTIONS.  
 DXI/DL IS THE CORRELATION BETWEEN XI(0) AND L+ IN KMUS.  
 DLO/DL IS THE CORRELATION BETWEEN LO AND L+ IN KMUS.  
 T = MOMENTUM TRANSFER TO THE PI IN UNITS OF MPl\*\*2.  
 DP = DALITZ PLOT ANALYSIS  
 PI = PI SPECTRUM ANALYSIS  
 MU = MU SPECTRUM ANALYSIS  
 POL = MU POLARIZATION ANALYSIS  
 BR = KMUS/KE3 BRANCHING RATIO ANALYSIS  
 E = POSITRON OR ELECTRON SPECTRUM ANALYSIS  
 RC = RADIATIVE CORRECTIONS

### $\xi_A = f_-/f_+$ (DETERMINED FROM SPECTRA)

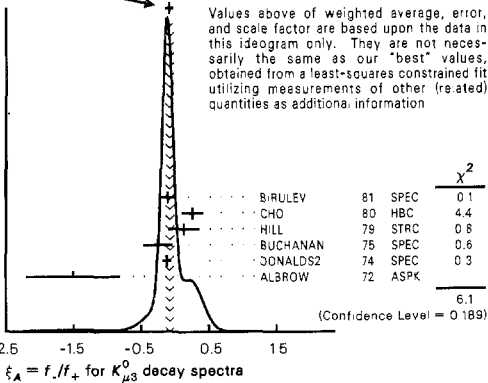
XIA	A1341	+1.2	(0.8)	CARPENTER	66 OSPK	DP, DXI/DL=-18
XIA	B 3140	(-3.9)	(0.4)	BASTILE	70 OSPK	DP, INDEP OF L+
XIA	C 16K	(-0.68)	(0.12)	CHIEN	70 ASPK	DP, DXI/DL=-26
XIA	D9086	-1.5	0.7	ALBROW	72 ASPK	DP, DXI/DL=-28
XIA	C 16K	(+0.50)	(0.61)	DALLY	72 ASPK	DP, DXI/DL UNKN.
XIA	E1385	-1.00	(0.45)	PEACH	73 HLBC	DP, DXI/DL=-20
XIA	F1.6M	-0.11	0.07	DONALDSON	74 SPEC	DP, DXI/DL=-17
XIA	G 32K	-0.25	0.22	BUCHANAN	75 SPEC	DP, DXI/DL=-5.9
XIA	H 16K	+0.13	0.23	HILL	79 STRC	DP, DXI/DL=-20
XIA	I 14K	+0.26	0.16	CHO	80 HBC	DP, DXI/DL=-13
XIA	J150K	-0.10	0.09	BIRULEV	81 SPEC	DP, DXI/DL=-12
XIA	AVG	-0.074	0.061	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)		
XIA	FIT	-0.11	0.09	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 2.3)	(SEE IDEOGRAM BELOW)	

XIA FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.  
 XIA A CARPENTER 66 XI(0) IS FOR L=0. DXI/DL IS FROM FIG. 9.  
 XIA B BASTILE 70 IS INCOMPATIBLE WITH ALL OTHER RESULTS. AUTHORS SUGGEST XIA B THAT EFFICIENCY ESTIMATES MIGHT BE RESPONSIBLE.  
 XIA C CHIEN 70 ERRORS ARE STATISTICAL ONLY. DXI/DL FROM FIG. 4.  
 XIA C DALLY 72 IS A REANALYSIS OF CHIEN 70. THE DALLY 72 RESULT IS NOT COMPATIBLE WITH ASSUMPTION L=0 SO NOT INCLUDED IN OUR FIT.  
 XIA C THE NON-ZERO L- VALUE AND THE RELATIVELY LARGE L+ VALUE FOUND BY XIA C DALLY 72 COME MAINLY FROM A SINGLE LOW T BIN (FIGS.1,2).  
 XIA C THE (F- XI) CORRELATION WAS IGNORED.  
 XIA C WE ESTIMATE FROM FIG. 2 THAT FIXING L=0 WOULD GIVE XI(0)=-1.4+-0.3 AND WOULD ADD 10 TO CHI SQUARED. DXI/DL IS NOT GIVEN.  
 XIA D ALBROW 72 FIT HAS L- FREE, GETS L=-.050+-0.060 OR LAM+15+-17-.11.  
 XIA E PEACH 73 GIVES XI(0)=-.95+-0.45 FOR L=+1=-.025. THE ABOVE VALUE IS FOR L=0. K.PEACH, PRIVATE COMMUNICATION(1974).  
 XIA F DONALDSON 74 GIVES XI=-.11+-0.02 NOT INCLUDING SYSTEMATICS. ABOVE XIA F ERROR AND DXI/DL WERE CALCULATED BY US FROM LO AND L+ ERRORS (WHICH INCLUDE SYSTEMATICS) AND DLO/DL+.  
 XIA G BUCHANAN 75 IS CALCULATED BY US FROM LO, L- AND DLO/DL+ BECAUSE G THEIR APPENDIX A VALUE -.20+-22 ASSUMES XI(T) CONSTANT, I.E. L=-L+.  
 XIA H HILL 79 AND CHO 80 CALCULATED BY US FROM LO, L+, AND DLO/DL+.  
 XIA I BIRULEV 81 ERROR, DXI/DL CALC. BY US FROM LO, L-, DLO/DL+ USED.

# Stable Particle Full Listings

$K^0_L$

WEIGHTED AVERAGE  
-0.074 ± 0.061 (ERROR SCALED BY 1.2)



$\xi_B = f_-/f_+$  (DETERMINED FROM  $K_{\mu 3}/K_{e 3}$ )

X1B THE KMUS/KE3 BRANCHING RATIO FIXES A RELATIONSHIP BETWEEN XI(0) AND L+. WE QUOTE THE AUTHORS XI(0) AND ASSOCIATED L+ BUT DO NOT AVERAGE BECAUSE THE L+ VALUES DIFFER. THE FIT RESULT AND SCALE FACTOR GIVEN IN THE NOTE ON KL3 FORM FACTORS IN THE K+- SECTION OF THIS PAPER ARE NOT OBTAINED FROM THESE XIB VALUES. INSTEAD THEY ARE OBTAINED DIRECTLY FROM THE AUTHORS KMUS/KE3 BRANCHING RATIO VIA THE FITTED KMUS/KE3 RATIO (R10).

XIB 389 (+1.1) (-1.1)	ADAIR	64 HBC	BR, L+=0	
XIB (+0.66) (-0.9)	(1.3)	LUERS	64 HBC	BR, L+=0
XIB (-0.2) (-0.8)	(1.2)	KULYUKINA	68 CC	BR, L+=0
XIB 569 (-0.45) (-0.28)		BELLIERE	69 HLBC	BR, L+=0
XIB E 1309 (-0.22) (-0.30)		EVANS	69 HLBC	
XIB 3548 (-0.5) (-0.5)		BASILE	70 OSPK	BR, L+=.02
XIB 6700 (-0.5) (-0.4)		BRANDENBU	73 HBC	BR, L+=.019+-.013
XIB E1509 (-0.08) (-0.25)		EVANS	73 HLBC	BR, L+=.02

XIB FIT -0.11 ± 0.09 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 2.3)

XIB FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.

XIB E EVANS 73 REPLACES EVANS 69.

$\xi_C = f_-/f_+$  (DETERMINED FROM  $\mu$  POLARIZATION IN  $K_{\mu 3}$ )

X1C THE MU POLARIZATION IS A MEASURE OF XI(T). NO ASSUMPTIONS ON L- NECESSARY, T (WEIGHTED BY SENSITIVITY TO XI0) SHOULD BE SPECIFIED. IN L+, XI(0) PARAMETERIZATION THIS IS XI(0) FOR L=0. DXI/DL+T.

X1C FOR RAD. CORR. TO MUON POLARIZATION IN KMUS, SEE EINSBERG 73.

X1C T 2608 (-1.2) (-0.5)	AUERBACH	66 OSPK	POLARIZATION		
X1C T 638 (-1.6) (-0.5)	ABRAMS	68 OSPK	POLARIZATION		
X1C L -1.81	0.50	0.26	LONGO	69 CNTR	POL, T=3.3
X1C S2.2M -0.385	0.105	SANDWEISS	73 CNTR	POL, DXI/DL+-6	
X1C H207K +0.178	0.105	CLARK	77 SPEC	POL, DXI/DL+-6.8	

X1C AVG -0.17 ± 0.28 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 3.9)

X1C FIT -0.11 ± 0.09 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 2.3)

(SEE IDEOGRAM BELOW)

X1C FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.

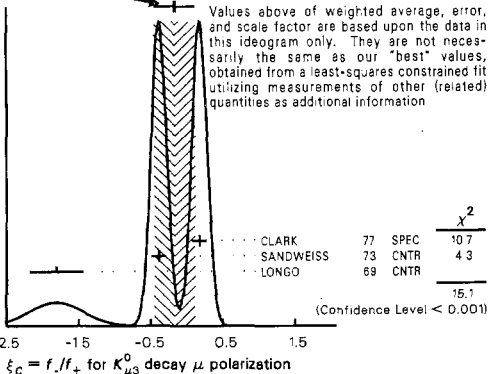
X1C T T VALUE NOT GIVEN.

X1C L LONGO 69 T=3.3 CALC. FROM DXI/DL=-6.0 (TABLE 1) DIVIDED BY XI+-1.81

X1C S SANDWEISS 73 IS FOR L+=0 AND T=0.

X1C H CLARK 77 T=3.80, DXI/DL+XI(T)\*T=.178\*3.80+-.68.

WEIGHTED AVERAGE  
-0.17 ± 0.28 (ERROR SCALED BY 3.9)



IMAGINARY PART OF  $\xi$  (TEST OF T REVERSAL)

X1I	-0.2	0.6	ABRAMS	68 OSPK	POLARIZATION
X1I	-0.02	0.08	LONGO	69 CNTR	POL, T=3.3
X1I 2.2M	-0.060	0.045	SANDWEISS	73 CNTR	POL, T=0
X1I S2.2M	-0.085	0.064	SANDWEISS	73 CNTR	POL, T=0
X1I C207K	0.35	0.30	CLARK	77 SPEC	POL, T=0
X1I	(0.012)	(0.026)	SCHMIDT	79 CNTR	REPL. BY MORSE 80
X1I	0.009	0.030	MORSE	80 CNTR	POLARIZATION
X1I S	SANDWEISS 73 VALUE CORRECTED FROM VALUE QUOTED IN THEIR PAPER DUE TO NEW VALUE OF RE(XI). SEE FTNOTE 4 OF SCHMIDT 79.				
X1I C	CLARK 77 VALUE HAS ADDITIONAL XI0 DEPENDENCE +0.21*RE(XI0).				

X1I AVG -0.020 ± 0.022 AVERAGE

$\lambda_+$  (LINEAR ENERGY DEPENDENCE OF  $f_+$  IN  $K_{\mu 3}$  DECAY)

SEE ALSO THE CORRESPONDING ENTRIES AND NOTES IN SECTION XIA AND LO. FOR RAD. CORR. OF KMUS DP SEE GINSBURG 70 AND BECHERRAY 70.

L+M A9086	0.085	0.015	ALBROW	72 ASPK	DP
L+M 16K (0.11)	(0.04)		DALLY	72 ASPK	DP
L+M 82K (0.046)	(0.008)		ALBRECHT	74 WIRE	REPL. BY BIRULEV 81
L+M 1.6M	0.030	0.003	DONALDS2	74 SPEC	DP
L+M 32K	0.046	0.030	BUCHANAN	75 SPEC	DP
L+M 129K (0.0337)	(0.0033)		DZHRDZHA	77 SPEC	REPL. BY BIRULEV 81
L+M 16K	0.028	0.011	HILL	79 STRC	DP
L+M 14K	0.028	0.010	CHO	80 HBC	DP
L+M 150K	0.0427	0.0044	BIRULEV	81 SPEC	DP

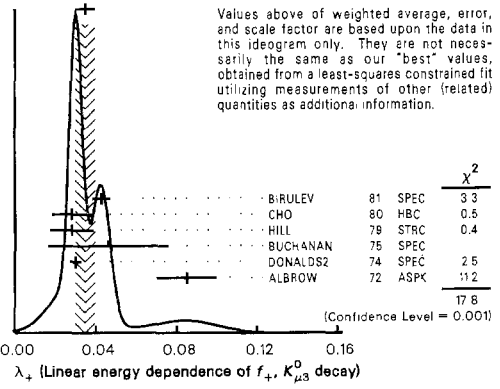
L+M AVG 0.0347 ± 0.0049 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.1)

L+M FIT 0.034 ± 0.005 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 2.3)

(SEE IDEOGRAM BELOW)

L+M FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.

WEIGHTED AVERAGE  
0.0347 ± 0.0049 (ERROR SCALED BY 2.1)



$\lambda_0$  (LINEAR ENERGY DEPENDENCE OF  $f_0$  IN  $K_{\mu 3}$  DECAY)

WHEREVER POSSIBLE, WE HAVE CONVERTED THE ABOVE VALUES OF XI(0) INTO VALUES OF LO USING THE ASSOCIATED L+M AND DXI/DL.

LO L 1371 +0.08	(0.07)	CARPENTER	66 OSPK	DP, DLO/DL+-0.54
LO L -0.140	(0.043)	(0.022)LONGO	69 CNTR	POL, DLO/DL+-0.49
LO B 3140 (-0.333)	(0.034)	BASILE	70 OSPK	DP, DLO/DL+-1.1
LO A 9086 -0.043	0.052	ALBROW	72 ASPK	DP, DLO/DL+-1.39
LO C 16K (-0.067)	(0.227)	DALLY	72 ASPK	DP, DLO/DL UNKN.
LO R 6700 (+0.06)	(0.03)	BRANDENBU	73 HBC	BR, L+=.019+-.013
LO P 1385 -0.060	(0.038)	PEACH	73 HLBC	DP, DLO/DL+-0.71
LO L 2.2M -0.018	(0.009)	SANDWEISS	73 CNTR	POL, DLO/DL+-0.49
LO E 82K (+0.024)	(0.011)	ALBRECHT	74 WIRE	REPL. BY BIRULEV 81
LO F 1.6M +0.019	0.004	DONALDS2	74 SPEC	DP, DLO/DL+-0.47
LO F 32K -0.025	0.019	BUCHANAN	75 SPEC	DP, DLO/DL+-0.5
LO L 207K -0.047	(0.009)	CLARK	77 SPEC	POL, DLO/DL+-1.06
LO L 47K (+0.0485)	(0.076)	DZHRDZHA	77 SPEC	REPL. BY BIRULEV 81
LO L 16K +0.039	0.010	HILL	79 STRC	DP, DLO/DL+-0.67
LO L 14K +0.050	0.008	CHO	80 HBC	DP, DLO/DL+-0.11
LO G 14K (0.041)	(0.008)	CHO	80 HBC	BR, L+=0.028
LO H 150K 0.0341	0.0057	BIRULEV	81 SPEC	DP, DLO/DL+-0

LO AVG 0.0279 ± 0.0057 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)

LO FIT 0.025 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 2.3)

(SEE IDEOGRAM BELOW)

LO FIT DISCUSSED IN NOTE ON KL3 FORM FACTORS IN 1982 EDITION.

LO L LO VALUE IS FOR L+=0.3 CALCULATED BY US FROM X10 AND DXI/DL.

LO B BASILE 70 LO IS FOR L+=0. CALCULATED BY US FROM X1A WITH DXI/DL=0.

LO B BASILE 70 IS INCOMPATIBLE WITH ALL OTHER RESULTS. AUTHORS SUGGEST THAT EFFICIENCY ESTIMATES MIGHT BE RESPONSIBLE.

LO A ALBROW 72 LO IS CALCULATED BY US FROM X1A, L+ AND DXI/DL. THEY GIVE LO = -.043+-.039 FOR L+=0. WE USE OUR LARGER CALCULATED ERROR.

LO C DALLY 72 GIVES FO=-1.20+-.35, LO=-.080+-.272, LOPRIME=-.006+-.045, BUT WITH A DIFFERENT DEFINITION OF LO. OUR QUOTED LO IS HIS LO/FO.

LO C WE CANNOT CALCULATE TRUE LO ERROR WITHOUT HIS (LO,FO) CORRELATIONS.

LO C SEE ALSO NOTE C IN SECTION XIA.

LO P PEACH 73 ASSUMES L+=0.025. CALCULATED BY US FROM X10 AND DXI/DL+.

LO R FIT FOR LO DOES NOT INCLUDE THIS VALUE BUT INSTEAD INCLUDES THE KMUS/KE3 RESULT FROM THIS EXPERIMENT.

LO R DONALDS2 74 DLO/DL+ OBTAINED FROM FIG. 18.

LO F BUCHANAN 75 VALUE IS FROM THEIR APPENDIX A AND USES ONLY KMUS DATA.

LO C DLO/DL+ WAS OBTAINED BY PRIVATE COMMUNICATION, C. BUCHANAN, 1976.

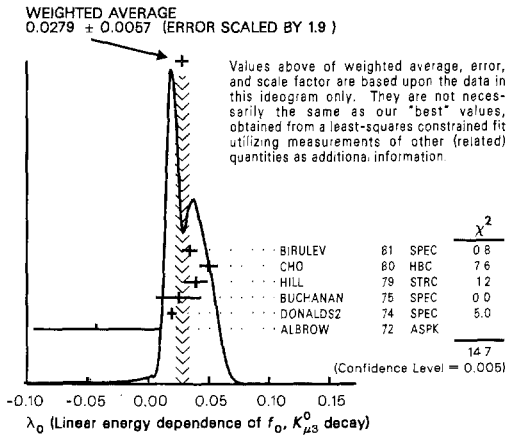
LO G CHO 80 BR RESULT NOT INDEPENDENT OF THEIR DP RESULT.

LO H BIRULEV 81 GIVES DLO/DL+-1.5, GIVING AN UNREASONABLY NARROW ERROR ELLIPSE WHICH DOMINATES ALL OTHER RESULTS. WE USE DLO/DL+=0.

For notation, see key on page 91.

# Stable Particle Full Listings

$K_L^0$



$|f_T/f_+|$  FOR  $K_{e3}$  DECAY

FT	RATIO OF TENSOR TO F+ COUPLINGS	CL	EXPERIMENT
FT	(1.0) OR LESS	CL=.68	KULYUKINA 67 CC
FT	5600 (1.0) OR LESS	CL=.95	ALBROW 73 ASPK
FT	25K 0.23 OR LESS	CL=.68	BLUMENTHA 75 SPEC
FT	48K (0.34) OR LESS	CL=.68	BIRULEV 76 SPEC SEE ALSO BIRULEV 81
FT	18K (0.40) OR LESS	CL=.95	HILL 78 STRC

$|f_T/f_+|$  FOR  $K_{\mu 3}$  DECAY

FTM	RATIO OF TENSOR TO F+ COUPLINGS	BIRULEV	81 SPEC
FTM	0.12 0.12		

## NOTE ON CP VIOLATION IN $K_L^0$ DECAY

We list the parameters which measure CP violation in  $K_L^0$  decays and compare them with superweak model predictions.

### Parameters

There are two different  $K_L^0$  decays in which CP violation has been observed (for details, see Kleinknecht<sup>1</sup>).

(a) Asymmetry in the  $K_L \rightarrow \pi^\mp \ell^\pm \nu$  decays. The quantity measured and compiled here is

$$\delta = \frac{\Gamma(K_L \rightarrow \pi^- \ell^+ \nu) - \Gamma(K_L \rightarrow \pi^+ \ell^- \nu)}{\Gamma(K_L \rightarrow \pi^- \ell^+ \nu) + \Gamma(K_L \rightarrow \pi^+ \ell^- \nu)}$$

This asymmetry violates CP invariance. If CPT is good, for a pure  $K_L^0$  beam,  $\delta$  can be written as

$$\delta = 2[(1 - |x|^2)/(1 + |x|^2)] \text{Re } \epsilon,$$

where x is defined below in the "Note on the  $\Delta S = \Delta Q$  Rule in  $K^0$  Decay," and  $\epsilon$  is the parameter of the expansion

$$|K_L\rangle = [(1 + \epsilon)|K\rangle - (1 - \epsilon)|\bar{K}\rangle]/[2(1 + |\epsilon|^2)]^{1/2}, \quad (1a)$$

$$|K_S\rangle = [(1 + \epsilon)|K\rangle + (1 - \epsilon)|\bar{K}\rangle]/[2(1 + |\epsilon|^2)]^{1/2}. \quad (1b)$$

We list  $\delta$  separately for  $K_L^0 \rightarrow \pi\mu\nu$  and  $K_L^0 \rightarrow \pi e\nu$  in sections A1 and A2 respectively, and list the combined values in section AL.

(b)  $K_L \rightarrow 2\pi$  decay. The relevant parameters are

$$\eta_{+-} = A(K_L \rightarrow \pi^+ \pi^-)/A(K_S \rightarrow \pi^+ \pi^-)$$

$$= |\eta_{+-}| \exp(i\phi_{+-}),$$

$$\eta_{00} = A(K_L \rightarrow \pi^0 \pi^0)/A(K_S \rightarrow \pi^0 \pi^0)$$

$$= |\eta_{00}| \exp(i\phi_{00}),$$

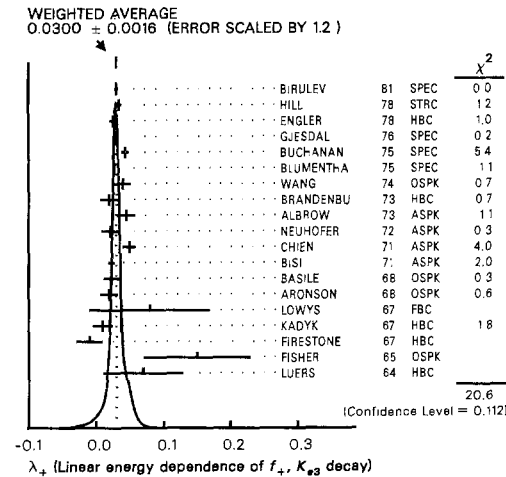
$\epsilon$ , defined in Eqs. (1) above, and

$$\epsilon' = \frac{1}{2}i\sqrt{2} \exp[i(\delta_2 - \delta_0)] \text{Im}(A_2/A_0).$$

Here the decay amplitudes to  $\pi\pi$  states with definite

### $\lambda_+$ (LINEAR ENERGY DEPENDENCE OF $f_+$ IN $K_{e3}$ DECAY)

L+E	FOR RAD. COR.	OF KE3 DP	SEE GINSBURG 67 AND BECHERRAWY 70.	EXPERIMENT	81 SPEC	$\chi^2$
L+E	153	+0.07	0.06	LUERS	64	HBC DP, NO RC
L+E	577	-0.15	0.08	FISHER	65	OSPK DP, NO RC
L+E	762	-0.01	0.02	FIRESTONE	67	HBC DP, NO RC
L+E	531	+0.01	0.015	KADYK	67	HBC E, P1, NO RC
L+E	240	+0.08	0.10	LOWYS	67	FBC P1
L+E	1000	0.02	0.015	ARONSON	68	OSPK P1
L+E	4800	+0.023	0.012	BASTIE	68	OSPK DP, NO RC
L+E	42K	0.023	0.005	BISI	71	ASPK DP
L+E	16K	0.05	0.01	CHIEN	71	ASPK DP, NO RC
L+E	1910	0.022	0.014	NEUHOFER	72	ASPK P1
L+E	5600	0.045	0.014	ALBROW	73	ASPK DP
L+E	1871	0.019	0.013	BRANDENBU	73	HBC P1 TRANSV.
L+E	2171	0.040	0.012	WANG	74	OSPK DP
L+E	25K	0.0270	0.0028	BLUMENTHA	75	SPEC DP
L+E	24K	0.044	0.006	BUCHANAN	75	SPEC DP
L+E	48K	(0.032)	(0.0042)	BIRULEV	76	SPEC REPL. BY BIRULEV 81
L+E	500K	0.0312	0.0025	GJESDAL	76	SPEC DP
L+E	E 12K	0.025	0.005	ENGLER	78	HBC DP
L+E	18K	0.0348	0.0044	HILL	78	STRC DP
L+E	26K	(0.0286)	(0.0049)	BIRULEV	79	SPEC REPL. BY BIRULEV 81
L+E	E 19K	(0.029)	(0.005)	CHO	80	HBC DP
L+E	74K	0.0506	0.0054	BIRULEV	81	SPEC DP
L+E	E	ENGLER 78 USES UNIQUE KE3 SUBSET OF CHO 80 EVENTS AND IS LESS SUBJECT TO SYSTEMATIC EFFECTS.				
L+E	AVG	0.0300	0.0016	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2) (SEE IDEOGRAM BELOW)		



$|f_S/f_+|$  FOR  $K_{e3}$  DECAY

FS	RATIO OF SCALAR TO F+ COUPLINGS	CL	EXPERIMENT
FS	(0.15) OR LESS	CL=.68	KULYUKINA 67 CC
FS	5600 (0.19) OR LESS	CL=.95	ALBROW 73 ASPK
FS	25K 0.04 OR LESS	CL=.68	BLUMENTHA 75 SPEC
FS	48K (0.07) OR LESS	CL=.68	BIRULEV 76 SPEC SEE ALSO BIRULEV 81
FS	18K (0.095) OR LESS	CL=.95	HILL 78 STRC



## Stable Particle Full Listings

$K_L^0$

isospin  $I$  are given by

$$\langle I=0 | T | K \rangle = \exp(i\delta_0)A_0,$$

$$\langle I=2 | T | K \rangle = \exp(i\delta_2)A_2.$$

where the  $\delta_I$  are the  $\pi\pi$  scattering phase shifts at the  $K$  mass. Wu and Yang<sup>2</sup> derived the relationships

$$\eta_{+-} = \epsilon + \epsilon', \quad \eta_{00} = \epsilon - 2\epsilon',$$

assuming  $CPT$  invariance, using a phase convention in which  $A_0$  is real, and neglecting small corrections of order  $\text{Re}A_2/\text{Re}A_0$ .

Measurements of  $|\eta_{+-}|$ ,  $|\eta_{00}|^2$ ,  $\phi_{+-}$ , and  $\phi_{00}$  are listed in sections E+-, EOS, F+-, and F00. The FIT values given in these sections come from constrained fits which include the  $|\eta_{00}|/|\eta_{+-}|$  and  $\phi_{00} - \phi_{+-}$  measurements from sections ER and DF, respectively.

One can determine the values of  $|\epsilon|$  and  $|\epsilon'|$  using theoretical input on their phases. The phase of  $\epsilon$  is determined from unitarity and  $CPT$  invariance to be within a few degrees of the "superweak value" of  $44^\circ$  (for a recent analysis, see V. Barmin et al.<sup>3</sup>). The phase of  $\epsilon'$  is given by  $\delta_2 - \delta_0 + \pi/2$ , which equals  $48 \pm 8^\circ$  from an analysis of  $\pi\pi$  phase shifts.<sup>4</sup> Thus  $\epsilon$  and  $\epsilon'$  have the same phase within  $10$  to  $15^\circ$ . Therefore to a good approximation

$$\frac{\epsilon'}{\epsilon} = \frac{1}{6} \left\{ |\eta_{+-}/\eta_{00}|^2 - 1 \right\} = (-3.3 \pm 4.2) \times 10^{-3}.$$

The value of  $\phi_{00} - \phi_{+-}$  is not used in this analysis, since taken literally it represents a violation of  $CPT$  invariance.<sup>3</sup>

**Superweak model predictions for  $|\eta_{00}/\eta_{+-}|$ ,  $\phi_{+-}$ , and  $\text{Re } \epsilon$**

The superweak model<sup>5</sup> predicts that<sup>6</sup>

$$|\eta_{00}/\eta_{+-}| = 1,$$

$$\phi_{+-} = \phi_{00} = \tan^{-1} \left( \frac{2\Delta m \tau_S}{\hbar} \right),$$

and

$$\text{Re } \epsilon = |\eta_{+-}| \left[ 1 + \left( \frac{2\Delta m \tau_S}{\hbar} \right)^2 \right]^{-1/2}$$

The latter two expressions and the values of the  $K_S^0 - K_L^0$  mass difference  $\Delta m = (0.5349 \pm 0.0022) \times 10^{10} \hbar \text{ sec}^{-1}$ , the  $K_S^0$  mean life  $\tau_S = (0.8923 \pm 0.0022) \times 10^{-10} \text{ sec}$ , and the magnitude of the  $(K_L^0 \rightarrow \pi^+\pi^-)/(K_S^0 \rightarrow \pi^+\pi^-)$  amplitude ratio  $|\eta_{+-}| = (2.275 \pm 0.021) \times 10^{-3}$ , all from the current edition,

result in the predictions that

$$\phi_{+-} = \phi_{00} = (43.67 \pm 0.14)^\circ$$

and

$$\text{Re } \epsilon = (1.646 \pm 0.016) \times 10^{-3}.$$

The above predictions can be compared with the experimental values

$$|\eta_{00}/\eta_{+-}| = 1.010 \pm 0.013,$$

$$\phi_{+-} = (44.6 \pm 1.2)^\circ,$$

$$\phi_{00} = (54.5 \pm 5.3)^\circ,$$

$$\text{Re } \epsilon = (1.621 \pm 0.088) \times 10^{-3},$$

where  $\text{Re } \epsilon$  has been computed using the relation

$$\text{Re } \epsilon = \frac{\delta}{2} \left( \frac{|1-x|^2}{1-|x|^2} \right),$$

and our current values of the charge asymmetry parameter for leptonic  $K_L^0$  decay  $\delta = (0.330 \pm 0.012)\%$  and the  $\Delta S = -\Delta Q$  amplitude  $(\text{Re}x, \text{Im}x) = (0.009 \pm 0.020, -0.004 \pm 0.026)$ .

The superweak predictions are within one standard deviation of the data except for the measured value of  $\phi_{00}$ , which is two standard deviations above the prediction. This results primarily from the CHRISTENSON179 measurement  $\phi_{00} = (55.7 \pm 5.8)^\circ$ .

### Searches for $CP$ violation in $K_L \rightarrow 3\pi$

As was discussed in the note on  $K \rightarrow 3\pi$  decay in the  $K^\pm$  section of the Full Listings, the Dalitz plot distribution for this decay contains a charge asymmetry term with coefficient  $j$ , the presence of which would indicate  $CP$  violation. Experimenters have used several forms for this  $CP$ -violation term. As described in the "Note on Slope Parameters for  $K \rightarrow 3\pi$  Decays" in the 1982 edition of this Review,<sup>7</sup> we have converted all results to coefficient  $j$  and have listed the results in section JT0 below. The coefficient  $j$  is consistent with zero, i.e., absence of  $CP$  violation.

### References

1. K. Kleinknecht, Ann. Rev. Nucl. Sci. **26**, 1 (1976).
2. T.T. Wu and C.N. Yang, Phys. Rev. Lett. **13**, 380 (1964).
3. V. Barmin et al., Nucl. Phys. **B247**, 293 (1984).
4. T.T. Devlin and J.O. Dickey, Rev. Mod. Phys. **51**, 237 (1979).
5. L. Wolfenstein, Phys. Rev. Lett. **13**, 562 (1964).

For notation, see key on page 91.

# Stable Particle Full Listings

$K_L^0$

- 6. T.D. Lee and L. Wolfenstein, Phys. Rev. **138B**, 1490 (1965).
- 7. Particle Data Group, Phys. Lett. **111B**, 69 (1982).

## CP VIOLATION PARAMETERS IN $K_L^0$ DECAYS

### CHARGE ASYM. $\pi^+\pi^-\pi^0$ DECAYS

CP VIOL. COEFFICIENT  $j$  FOR  $|\mu(K_L^0 \rightarrow \pi^+\pi^-\pi^0)|^2$ .

JTO DEFINED AT BEGINNING OF SECTION GTO ABOVE. SEE ALSO NOTE ON SLOPE PARAMETERS IN K+- SECTION AND NOTE ON CP VIOLATION IN KL DECAY ABOVE.

JTO	DECAY ABOVE	0.001	0.004	BLANPIED 68
JTO	3M	0.0013	0.0009	SCRIBANO 70
JTO	4400	0.0	0.017	SMITH 70 OSPK
JTO	6499	0.001	0.011	CHO 77
JTO	4709	-0.001	0.003	PEACH 77
JTO	AVG	0.00110	0.00084	AVERAGE

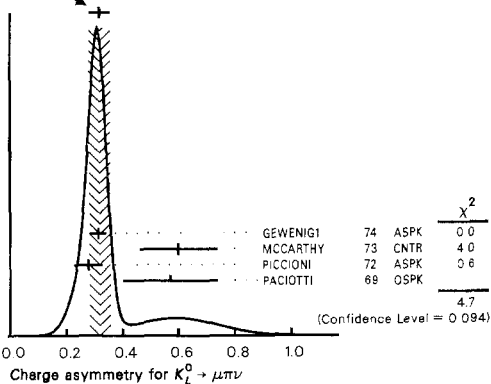
### CHARGE ASYM. IN LEPTONIC DECAYS

SUCH ASYMMETRY VIOLATES CP. IT IS RELATED TO REAL(EPSILON).

$K_L^0 \rightarrow ((\mu^+\pi^-\nu) - (\mu^-\pi^+\nu)) / ((\mu^+\pi^-\nu) + (\mu^-\pi^+\nu))$  (%)

JTO	DECAY ABOVE	0.001	0.004	BLANPIED 68
A1 D	1M	0.403	0.134	DORFAN 67 OSPK DERIVED FROM R16
A1	1M	0.57	0.17	PACIOTTI 69 OSPK
A1	7.7M	0.278	0.051	PICCIONI 72 ASPK
A1	4.1M	0.60	0.14	MCCARTHY 73 CNTR
A1	15M	0.313	0.029	GEWENIG1 74 ASPK
A1 D	PACIOTTI 69 IS A REANALYSIS OF DORFAN 67 AND IS CORRECTED FOR MU- MU+ RANGE DIFFERENCE IN MC CARTHY 72.			
A1	AVG	0.319	0.038	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5) (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
0.319 ± 0.038 (ERROR SCALED BY 1.5)



$K_L^0 \rightarrow ((e^+\pi^-\nu) - (e^-\pi^+\nu)) / ((e^+\pi^-\nu) + (e^-\pi^+\nu))$  (%)

JTO	DECAY ABOVE	0.001	0.004	BLANPIED 68
A2 B	10M	0.224	0.036	BENNETT 67 CNTR
A2 B	10M	0.246	0.059	SAL 69 CNTR
A2	10M	0.346	0.033	MARK 70 CNTR
A2	600K	0.36	0.18	ASHFORD 72 ASPK
A2	40M	0.318	0.038	FITCH 73 ASPK
A2	34M	0.341	0.018	GEWENIG1 74 ASPK
A2 B	SAL 69 IS A REANALYSIS OF BENNETT 67			
A2	AVG	0.333	0.014	AVERAGE

$K_L^0 \rightarrow ((\pi^+\pi^-\nu) - (\pi^-\pi^+\nu)) / ((\pi^+\pi^-\nu) + (\pi^-\pi^+\nu))$  (%)

JTO	DECAY ABOVE	0.001	0.004	BLANPIED 68
AL	10M	0.246	0.059	SAL 69 CNTR KE3
AL	1M	0.57	0.17	PACIOTTI 69 OSPK KMU3
AL	10M	0.346	0.033	MARK 70 CNTR KE3
AL	600K	0.36	0.18	ASHFORD 72 ASPK KE3
AL	7.7M	0.278	0.051	PICCIONI 72 ASPK KMU3
AL	40M	0.318	0.038	FITCH 73 ASPK KE3
AL	4.1M	0.60	0.14	MCCARTHY 73 CNTR KMU3
AL	33M	0.333	0.050	WILLIAMS 73 ASPK KMU3+KE3
AL	15M	0.313	0.029	GEWENIG1 74 ASPK KMU3
AL	34M	0.341	0.018	GEWENIG1 74 ASPK KE3
AL	AVG	0.330	0.012	AVERAGE

## PARAMETERS FOR $K_L^0 \rightarrow 2\pi$ DECAY

ETA+- = A(KL INTO PI+PI-)/A(KS INTO PI+PI-)  
ETA00 = A(KL INTO PI0PI0)/A(KS INTO PI0PI0)

THE FITTED VALUES OF ETA+- AND ETA00 GIVEN BELOW ARE THE RESULTS OF A FIT TO ETA+-, ETA00 AND ETA00/ETA+- RESULTS. THE VALUES LISTED BELOW WHICH ARE NOT PARENTHEZIZED AND BEAR THE FOOTNOTE X DO NOT ENTER THE FIT AS SHOWN. THESE EXPERIMENTS GIVE BRANCHING RATIOS AND ENTER THE FIT VIA THE QUANTITY ACTUALLY MEASURED -- BRANCHING RATIOS R9, R20 AND R27 (ETA+-) AND R17 AND R19 (ETA00). THESE BRANCHING RATIOS ARE COMBINED WITH CURRENT NORMALIZATIONS AND CURRENT KL AND KS MEAN LIVES TO OBTAIN PI PI RATES. THE ETA+- AND ETA00 VALUES OBTAINED FROM THESE RATES ARE ENTERED BELOW WITH THE NAME 'GKL/GKS'.

$|\eta_{00}|^2 = |A(K_L^0 \rightarrow 2\pi^0)/A(K_S^0 \rightarrow 2\pi^0)|^2$  (UNITS 10\*\*--6)

EOS X	0	(-2.)	(7.0)	BARTLETT	68 OSPK
EOS X	57	(4.9)	(1.2)	BANNER	69 OSPK
EOS XR	133	(14.1)	(3.4)	CENCE	69 OSPK
EOS XF	180	(13.)	(4.)	GAILLARD	69 OSPK
EOS X	29	(4.0B)	(0.9)	BARMIN	70 HBC
EOS X	30	(3.6I)	(1.9)	BUDAGOV	70 HBC
EOS C	8.7	3.7		CHOLLET	70 OSPK CU REG., 4GAMMAS
EOS XF	172	(9.9)	(3.4)	FAISSNER	70 OSPK
EOS C	56	7.4	2.0	WOLF	71 OSK
EOS XR	150	(14.1)	(5.4)	REY	76 OSPK
EOS	5.43	0.84		CHRISTE1	79 ASPK
EOS X	5.1	1.0		GKL/GKS	86 RVUE BR SCALE FACTOR=1.5

EOS X SEE NOTE ABOVE REGARDING FITTED VALUES OF ETA+- AND ETA00.  
EOR R CENCE 69 EVENTS ARE INCLUDED IN REY 76.  
EOS F FAISSNER 70 CONTAINS SAME 2PI0 EVENTS AS GAILLARD 69  
EOS C CHOLLET 70 GIVES ETA00=(1.25+-0.24)\*(REGEN AMPL, 2GEV/C CU)/10000MB  
EOS C WOLFF 71 GIVES ETA00=(1.13+-0.12)\*(REGEN AMPL, 2GEV/C CU)/10000MB  
EOS C WE COMPUTE BOTH ETA00\*2 VALUES FOR (REGEN AMPL, 2GEV/C CU)=24+-2MB.  
EOS C THIS REGEN AMPL RESULTS FROM AVERAGING OVER FAISSNER 69,  
EOS C EXTRAPOLATED USING OPTICAL MODEL CALCULATIONS OF BOHM ET AL.  
EOS C PL 27B 69 AND THE DATA OF BALATS 71. (FROM H. FAISSNER,  
EOS C PRIVATE COMMUNICATION)  
EOS C FAISSNER 70 CONTAINS SAME 2PI0 EVENTS AS GAILLARD 69

$|\eta_{+-}| = |A(K_L^0 \rightarrow \pi^+\pi^-)/A(K_S^0 \rightarrow \pi^+\pi^-)|$  (units 10<sup>-3</sup>)

THE FIT VALUE ABOVE FOR ETA00\*2 CORRESPONDS TO ETA00=2.299+-0.036

$|\eta_{+-}| = |A(K_L^0 \rightarrow \pi^+\pi^-)/A(K_S^0 \rightarrow \pi^+\pi^-)|$  (units 10<sup>-3</sup>)

E-- X	45	(1.95)	(0.20)	CHRISTENS	64 OSPK
E-- X	54	(1.99)	(0.16)	GALBRAITH	65 OSPK
E-- X	1	(1.92)	(0.13)	BASILE	66 OSPK
E-- X	1	(1.95)	(0.04)	BOTT-BODE	66 OSPK
E-- X	1	(2.00)	(0.09)	DEBOUARD	67 OSPK
E-- X	1	(1.94)	(0.08)	FITCH	67 OSPK
E-- AX	1	(1.95)	(0.03)	GKL/GKS	71 RVUE EXPTS. BEFORE 71

E-- A AVERAGE OF ABOVE EXPERIMENTS. THESE ARE EXCLUDED FROM THE GKL/GKS,  
E-- A AVERAGE, AND FIT VALUES BELOW SINCE THEY DO NOT AGREE WITH MORE  
E-- X RECENT PRECISE AND IN PRINCIPLE SUPERIOR EXPERIMENTS.  
E-- X 4200 (2.23) (0.05) MESSNER 75 ASPK  
E-- X 2703 (2.25) (0.05) GEWENIG2 74 ASPK  
E-- X 2703 (2.25) (0.05) DEVOE 77 SPEC  
E-- Z 2.27 (0.12) CHRISTE2 79 ASPK  
E-- Z (2.09) (0.02) ARONSON2 82 SPEC E=30-110 GEV  
E-- X C1687 2.28 0.06 COUPL 85 SPEC P(K)=70 GEV/C  
E-- X 2.258 0.027 GKL/GKS 86 RVUE BR EXP. AFTER 71  
E-- Z ARONSON 82 FITS THAT ETA+- MAY DEPEND ON THE KAON ENERGY.  
E-- X SEE NOTE ABOVE REGARDING FITTED VALUES OF ETA+- AND ETA00.  
E-- C COUPL 85 CONCLUDES: NO ENERGY DEPENDENCE OF E+-, BECAUSE THEIR  
E-- C VALUE IS CONSISTENT WITH ABOVE VALUES WHICH OCCUR AT LOWER ENERGIES

$\eta_{00}/\eta_{+-}$

E-- AVG	2.274	0.020	AVERAGE
E-- FIT	2.275	0.021	FROM FIT

ER APPROX EQUAL TO 1-3\*EPSILON/|EPSILON|.  
ER 124 1.03 0.07 BANNER1 72 OSPK  
ER 167 1.00 0.06 HOLDER 72 ASPK  
ER C (1.00) (0.09) CHRISTE1 79 ASPK  
ER B5152 1.014 0.017 BERNSTEIN 85 SPEC  
ER 1122 0.995 0.025 BLACK 85 SPEC  
ER C NOT INDEPENDENT OF E+- AND EOS VALUES WHICH ARE INCLUDED IN FIT.  
ER B STAT ERROR (0.016) AND SYST ERROR (0.007) ADDED IN QUADRATURE.  
ER AVG 1.008 0.013 AVERAGE  
ER FIT 1.010 0.013 FROM FIT

$\phi \rightarrow$  PHASE OF  $\eta_{+-}$  (degrees)

F-- THE DEPENDENCE OF THE PHASE ON THE KL-KS MASS DIFFERENCE IS GIVEN FOR EACH EXPERIMENT IN THE COMMENTS BELOW, WHERE DM IS (MASS DIFF./HBAR) IN UNITS 10\*\*10 SEC-1. WE HAVE EVALUATED THESE MASS DEPENDENCES USING OUR APRIL 1982 VALUE, DM=0.5349+-0.0022 TO OBTAIN THE VALUES AND AVERAGE QUOTED BELOW. WE ALSO GIVE THE REGENERATOR PHASE FR IN THE COMMENTS BELOW.

F-- O	(45.0)	(50.0)	FITCH	65 OSPK	BE REGEN
F-- O	(30.0)	(45.0)	FIRESTONE	66 HBC	
F-- O	(70.0)	(21.0)	BOTT-BODE	67 OSPK	C REGEN
F-- O	(25.0)	(35.0)	HISCHE	67 OSPK	CU REGEN

F-- OLD EXPERIMENTS WITH LARGE ERRORS NOT INCLUDED IN AVERAGE.  
F-- N (51.0) (11.0) BENNETT2 68 CNTR CU REG. USES  
F-- C 34.2 10.0 BENNETT 69 CNTR CU REGEN  
F-- B 45.3 12.0 BOHM 69 OSPK VACUUM REGEN  
F-- F 45.2 7.4 FAISSNER 69 ASPK CU REGEN  
F-- J 40.6 4.2 JENSEN 70 ASPK VACUUM REGEN  
F-- D 37.2 12.0 BALATS 71 OSPK CU REGEN  
F-- P 36.2 6.1 CARNELIE 72 ASPK CU REGEN  
F-- G 46.5 1.6 GEWENIG2 74 ASPK VACUUM REGEN  
F-- H 45.5 2.8 CARITHERS 75 SPEC C REGEN  
F-- A 41.7 3.5 CHRISTE2 79 ASPK CU REGEN  
F-- A (35.3) (3.9) ARONSON2 82 SPEC E=30-110 GEV

F-- AVG 44.6 1.2 AVERAGE  
F-- FIT 44.6 1.2 FROM FIT





# Stable Particle Full Listings

$K^0_L, D^\pm$

GINSBERG 67 PR 162 1570	EDWARD S GINSBERG (U. MASS BOSTON)
RUBBIA 67 PL 248 531	C. RUBBIA, J. STEINBERGER (CERN+COLU)
ALSO 1 66 PL 20 207	T. BECHERRAY (ROCH)
ALSO 2 66 PL 21 595	ALFF-STEINBERGER, HEUER, KLEINKNECHT+ (CERN)
ALSO 3 66 PL 23 167	C. RUBBIA, J. STEINBERGER (CERN+COLU)
SCHMIDT 67 NEVIS 160(THEISIS) P. SCHMIDT	(COLUMBIA)
CRONIN 68 VIENNA CONF P.281	CRONIN, RAPORTEURS TALK (PRINCETON)
BECHERRA 70 PR D1 1452	T. BECHERRAY (ROCH)
GINSBERG 70 PR D1 229	E S GINSBERG (IIT HAIFA)
HEUSER 70 LNC 3 449	-AUBERT, PASCAUD, VIALLE (ORSAY)
GINSBERG 73 PR D8 3887	E S GINSBERG, J SMITH (MIT-STON)
KLEINKNE 76 ARNS 26 1	K. KLEINKNECHT (DOBT)

## FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions,  $P_i$ , as follows: The diagonal elements are  $P_i = \delta P_i$ , where  $\delta P_i = \sqrt{(\delta P_i \delta P_i)}$ , while the off-diagonal elements are the normalized correlation coefficients  $(\delta P_i \delta P_j) / (\delta P_i \delta P_j)$ . For the definitions of the individual  $P_i$ , see the listings above; only those  $P_i$  appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

P 1	P 17	P 19
.1144+-0112	E+E-4.03, 4.41 ECM	
P 17	-.105	-.0436+-0155
P 19	-.639	-.832
		.8220+-0200

## CHARMED MESONS

$D^\pm$

$$J(P) = \frac{1}{2}(0^-)$$

### $D^\pm$ MASS (MeV)

M	50(187.6)	(15.)	PERUZZI 76 SMAG +- K+-P1-PI-
M	(187.4)	(5.)	GOLDBER 77 SMAG +- DD, D+ RECOIL SPC
M	(186.8, 3)	(0.9)	PERUZZI 77 SMAG +- E+E- 3.77GEV ECM
M	(187.4)	(11.)	PICCOLO 77 SMAG +- E+E-4.03, 4.41 ECM
M	(186.8, 4)	(0.5)	SCHINDLER 81 SMK2 +- E+E- 3.77GEV ECM
M	1869.4	0.6	TRILLING 81 RVUE +- E+E- 3.77GEV ECM
M	6 1860.	16.	ADAMOVICH 84 EMUL PHOTOPRODUCTION
M	1863.	4.	DERRICK 84 HRS E+E- 29GEV ECM
M	P	PERUZZI 77 AND SCHINDLER 81 ERRORS DO NOT INCLUDE THE 0.13 PERCENT	
M	P	UNCERTAINTY IN THE ABSOLUTE SPEAR ENERGY CALIBRATION. TRILLING 81	
M	P	USES THE HIGH PRECISION J/PSI AND PSI(3685) MEASUREMENTS OF	
M	P	ZHOLENTZ 80 TO DETERMINE THIS UNCERTAINTY AND COMBINES THE	
M	P	PERUZZI 77 AND SCHINDLER 81 RESULTS TO OBTAIN THE VALUE QUOTED.	
M	AVG	1869.25	0.59
		AVERAGE	

### $D^\pm$ MEAN LIFE (units $10^{-13}$ sec)

T	4	(8.)	OR LESS	CL-.90	ARMENISE 79 HYBR NU P --> DIMUONS +	
T	4	2.5	2.2	1.1	ALLASIA 80 EMUL NU WIDEBAND	
T	A	1	(10.4)	(3.9)	(2.9)	BACINO 80 DLCO E+E- 3.77 GEV ECM
T	C	1	(2.2)	(2.3)	(1.1)	BALLAGH 81 HYBR FNAL 15FT, NU HE-H2
T	E	70	2.5	3.1	1.9	ALBINI 82 SILI CERN GAM S1
T	T	15	8.4	3.5	2.2	AGUILAR 83 HYBR PI- P, P P
T	T	7	6.3	5.0	2.7	BADERTSCH 83 HYBR CERN PI- NUCL.
T	T	11	11.5	7.5	3.5	USHIDA 83 EMUL REPL, USHIDA 80
T	H	21	7.2	2.3	2.0	ABE 84 HYBR SLAC GAM P 20GEV
T	G	12	3.91	2.35	1.25	ADAMOVICH 84 EMUL PHOTOPRODUCTION
T	T	28	10.6	3.6	2.4	BAILEY 85 SILI + PI-BE 200 GEV
T	A					USES THEORETICAL RATE D TO (K E NU)=1.4*10**11 SEC**1
T	C					BALLAGH 81 VALUE QUOTED HERE ASSUMES THAT ALL DILEPTON EVENTS
T	C					CONTAIN DO OR D+, EACH WITH EQUAL NUMBERS OF SEMILEPTONIC DECAYS.
T	E					ALBINI 82 ASSUMES D MOMENTUM IS 1/2 BEAM MOMENTUM.
T	H					SOME EVENTS MAY BE D/S+ AND 5 EVENTS COULD BE LAMBDA/C+.
T	G					ESTIMATE SYSTEMATIC ERROR LESS THAN +1.2 -1.3
T	AVG	9.2	1.3	1.0	AVERAGE	

### $D^\pm$ PARTIAL DECAY MODES

			DECAY MASSES
P1	$D^+ \rightarrow K^- \pi^+ \pi^+$		494+ 140- 140
P2	$D^+ \rightarrow \bar{K}^0 \pi^+$		498+ 140
P3	$D^+ \rightarrow \pi^+ \pi^+ \pi^-$		140+ 140+ 140
P4	$D^+ \rightarrow \pi^+ K^+ K^-$		140+ 494+ 494
P5	$D^+ \rightarrow K^+ \pi^+ \pi^-$		494+ 140+ 140
P6	$D^+ \rightarrow e^+ \nu_e$		.511+ 0
P7	$D^+ \rightarrow e^+$ anything		
P8	$D^+ \rightarrow K^-$ anything		
P9	$D^+ \rightarrow \bar{K}^0$ anything + $K^0$ anything		
P10	$D^+ \rightarrow K^+$ anything		
P11	$D^+ \rightarrow \bar{K}^0(892)^0 \pi^+$		892+ 140
P12	$D^+ \rightarrow \bar{K}^0 \pi^+ \pi^0$		498- 140+ 135
P13	$D^+ \rightarrow \bar{K}^0 \pi^+ \pi^+ \pi^-$		498+ 140+ 140+ 140
P14	$D^+ \rightarrow K^- \pi^+ \pi^+ \pi^+ \pi^-$		494+ 140+ 140+ 140+
P15	$D^+ \rightarrow \pi^+ \pi^0$		140+ 135
P16	$D^+ \rightarrow \bar{K}^0 K^+$		498+ 494
P17	$D^+ \rightarrow K^- \pi^+ \pi^+ \pi^0$		494+ 140+ 140+ 135
P18	$D^+ \rightarrow \bar{K}^0 \rho^+$		498+ 769
P19	$D^+ \rightarrow$ unfitted modes (keeps fit program happy)		
P20	$D^+ \rightarrow \mu^+ \nu_\mu$		106+ 0
P21	$D^+ \rightarrow \eta$ anything		
P22	$D^+ \rightarrow \mu^+$ anything		
P23	$D^+ \rightarrow \mu^+ \mu^-$ anything		
P24	$D^+ \rightarrow \phi \pi^+$		140+1020
P25	$D^+ \rightarrow \bar{K}^0(892)^0 K^+$		892- 494
P26	$D^+ \rightarrow \pi^+ K^+ K^-$ (non-resonant)		140+ 494+ 494

D- MODES ARE CHARGE CONJUGATES OF THE ABOVE MODES

## NOTE ON CHARMED PARTICLE BRANCHING FRACTIONS

The determination of  $D^0$  and  $D^+$  hadronic branching fractions has been performed by three distinct techniques. The earliest measurements (PERUZZI 77, SCHARRE 78, and SCHINDLER 81) come from  $e^+e^-$  collisions at the  $\psi(3770)$  resonance. The rate at which  $D$  mesons are observed in a given channel is divided by the height of the resonance (assumed to decay wholly to  $D\bar{D}$  pairs) to obtain branching fractions. The height of the resonance is determined from a fine energy scan. The second technique (employed by AGUILAR 84 and AGUILAR2 84), in  $\pi p$  and  $pp$  scattering, relies on a determination of branching fractions through the identification of specific 3- and 4-body final states relative to all 3- and 4-charged-prong vertices. These are normalized by the absolute topological branching fractions of  $D$  mesons determined from tagged  $e^+e^-$  data. The third technique (BALTRUSAITIS 86), also coming from  $e^+e^-$  collisions at the  $\psi(3770)$  resonance, compares the number of fully reconstructed  $D\bar{D}$  pairs with the number of single  $D$  mesons reconstructed. This provides an absolute measurement which does not rely on the resonance height and assumptions about its decay to charmed particles.

The third technique is the most direct way of measuring the  $D$  meson branching fractions, relying solely on kinematics. The results indicate significantly higher branching fractions than the first technique, and are consistent with the second. For consistency in normalization, the results of SCHINDLER 81 have been rescaled to reflect the cross section for charm production determined in BALTRUSAITIS 86. Earlier  $e^+e^-$  results are excluded from the averages unless they are in ratios independent of this scale. The technique is justified, as the production rates ( $\sigma_D \cdot B$ ) in  $e^+e^-$  experiments at  $E_{cm} = 3.77$  GeV are generally in agreement, implying that the discrepancy in branching fractions lies either in the direct resonance-height determination or in the assumptions about the resonance decay. Other

For notation, see key on page 91.

Stable Particle Full Listings

D±

experiments measuring D branching fractions by relative measurements have been rescaled (when possible) to the measurement of BALTRUSAITIS 86.

D± BRANCHING RATIOS

Table of D± branching ratios including sections like (PB), (P10), (P9), (P7), (P22), (P23), (P24), (P15), (P16), (P2), (P1), (P5), (P17), (P1). Contains data for various decay channels such as (K- anything)/total, (K+ anything)/total, (K0 anything + K0+ anything)/total, (eta anything)/(total D+ and D0), (mu+ anything)/(total charm), (e+ nu\_e)/(D+ -> e+ anything + D0 -> e+ anything), (mu+ nu\_mu)/total, (mu+ mu- anything)/(total charm), (e+ anything)/(total D+ and D0), (K- (892)0 K+)/total, (K0 K+)/total, (K- (892)0 pi+)/total.

Table of D± branching ratios including sections like (P25), (P4), (P3), (P1), (P12), (P26), (P13), (P17), (P14), (P15)/(P2), (P16)/(P2), (P2)/(P1), (P24)/(P1), (P25)/(P1), (P3)/(P1), (P26)/(P1), (P5)/(P1), (P17)/(P1). Contains data for various decay channels such as (K- (892)0 K+)/total, (phi pi+)/total, (pi+ pi+ pi-)/total, (K- pi+ pi+)/total, (K0 pi+ pi+)/total, (K- K+ pi+ non-resonant)/total, (K0 pi+ pi+ pi-)/total, (K- pi+ pi+ pi0)/total, (K- pi+ pi+ pi+ pi-)/total, (pi+ pi0)/(K0 pi+), (K0 K+)/total, (K0 pi+)/total, (K- (892)0 K+)/total, (K- (892)0 K+)/total, (pi+ pi+ pi-)/(K- pi+ pi+), (K+ K- non-resonant)/(K- pi+ pi+), (K+ pi+ pi-)/(K- pi+ pi+), (K- pi+ pi+ pi0)/(K- pi+ pi+).



For notation, see key on page 91.

Stable Particle Full Listings D<sup>0</sup>

P36 D<sup>0</sup> -> K<sup>0</sup>K<sup>+</sup>π<sup>-</sup> + K<sup>0</sup>K<sup>0</sup>π<sup>-</sup> (non-resonant) 498. 494+ 140
P37 D<sup>0</sup> -> K\*(892)<sup>0</sup>K<sup>0</sup> + K\*(892)<sup>0</sup>K<sup>+</sup> 892+ 498
P38 D<sup>0</sup> -> K\*(892)<sup>+</sup>K<sup>+</sup> + K\*(892)<sup>+</sup>K<sup>-</sup> 892+ 494

DOBAR MODES ARE CHARGE CONJUGATES OF ABOVE MODES

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P<sub>i</sub>, as follows: The diagonal elements are P<sub>i</sub>±δP<sub>i</sub>, where δP<sub>i</sub> = √(δP<sub>i</sub>P<sub>i</sub>), while the off-diagonal elements are the normalized correlation coefficients (δP<sub>i</sub>δP<sub>j</sub>)/δP<sub>i</sub>δP<sub>j</sub>. For the definitions of the individual P<sub>i</sub>, see the listings above; only those P<sub>i</sub> appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Table with 5 columns: P 1, P 2, P 3, P 5, P 7, P 19. Values range from -0.384 to 0.119.

D<sup>0</sup> BRANCHING RATIOS

SEE NOTE IN D<sup>0</sup> SECTION CONCERNING REVISIONS TO HADRONIC BRANCHING FRACTIONS.

D<sup>0</sup> -> (K<sup>-</sup> anything)/total (P11)
R1 19 0.35 0.10 VUILLEMIN 78 SMAG E+E- 3.772 GEV ECM
R1 121 0.55 0.11 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R1 AVG 0.440 0.100 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)

D<sup>0</sup> -> (K<sup>+</sup> anything)/total (P12)
R2 25 0.98 0.03 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM

D<sup>0</sup> -> (K<sup>0</sup> anything + K<sup>2\*</sup> anything)/total (P13)
R3 6 0.57 0.26 VUILLEMIN 78 SMAG E+E- 3.772 GEV ECM
R3 13 0.29 0.11 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R3 AVG 0.33 0.10 AVERAGE

D<sup>0</sup> -> (e<sup>+</sup> anything)/total (P10)
R5 0 (0.04) OR LESS CL=.95 BACINO 80 DLCO E+E- 3.77 GEV ECM
R5 12 0.055 0.037 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R5 3 0.051 0.048 AGUILAR 83 HYBR P1-P PP P
R5 A 137 0.075 0.012 BALTRUSA 86 SMK3 E+E- 3.77GEV ECM
R5 A THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.011 AND 0.004.
R5 AVG 0.070 0.011 AVERAGE

D<sup>0</sup> -> (μ<sup>+</sup> μ<sup>-</sup>)/total (P28)
R7 3.4E-4 OR LESS CL=.90 AUBERT 85 EMC 0 DEEP INEL MU-NUC

D<sup>0</sup> -> (μ<sup>-</sup> anything, via D<sup>0</sup>)(μ<sup>+</sup> any + μ<sup>-</sup> any) (P26)/(P26+P27)
R8 THIS IS A DO-DOBAR MIXING LIMIT.
R8 0.044 OR LESS CL=.90 BODEK 82 SPEC P1-P, FE-->D0

D<sup>0</sup> -> (K<sup>+</sup> π<sup>-</sup>, via D<sup>0</sup>)(K<sup>-</sup> π<sup>+</sup> + K<sup>+</sup> π<sup>-</sup>) (P6)/(P1-P6)
R9 THIS IS A DO-DOBAR MIXING LIMIT.
R9 (0.16) OR LESS CL=.90 FELDMAN 77 SMAG D+ TO DO P1+
R9 (0.18) OR LESS CL=.90 GOLDBABER 77 SMAG
R9 (0.11) OR LESS CL=.90 AVERY 80 SPEC GAMMA NUC -->D+
R9 A (0.23) OR LESS CL=.90 ALTHOFF 84 TASS E+E- 34.4 GEV
R9 0.081 OR LESS CL=.90 YAMAMOTO 85 DLCO 0 E+E- 29 GEV ECM
R9 2 (0.11) OR LESS CL=.90 ALBRECHT 85 ARG 0 E+E- 10 GEV ECM
R9 A TECHNIQUE USES K<sup>-</sup> P1 DECAY TO DERIVE LIMIT
R9 A ALSO DERIVE RATIO OF CHARM TO AVERAGE ALPHA-STRONG OF
R9 A 1.00+-0.20+-0.20 AT 34.4 GEV USING K<sup>+</sup> PION MODES

D<sup>0</sup> -> (π<sup>+</sup> π<sup>-</sup>)/total (P5)
R10 A 39 (0.0018) (0.0007) BALTRUSA 86 SMK3 E+E- 3.77GEV ECM
R10 A STATISTICAL AND SYSTEMATIC ERRORS ARE 0.0006 AND 0.0004.
R10 A NOT INDEPENDENT OF BALTRUSA2 85 R26.

D<sup>0</sup> -> (K<sup>+</sup> K<sup>-</sup>)/total (P7)
R11 A 118 (0.0068) (0.0013) BALTRUSA 86 SMK3 E+E- 3.77GEV ECM
R11 A STATISTICAL AND SYSTEMATIC ERRORS ARE 0.0011 AND 0.0007.
R11 A NOT INDEPENDENT OF BALTRUSA2 85 R27.

D<sup>0</sup> -> (K<sup>-</sup> π<sup>+</sup>)/total (P1)
R12A S 130 0.054 0.012 PERUZZI 77 SMAG E+E- 3.77GEV ECM
R12B S 263 0.054 0.008 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R12 C 930 0.056 0.005 BALTRUSA 86 SMK3 E+E- 3.77GEV ECM
R12 S CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM BALTRUSA 86.
R12A THE UNCORRECTED VALUE WAS 0.022+-0.006.
R12B THE UNCORRECTED VALUE WAS 0.030+-0.006.
R12 C THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.004 AND 0.003.
R12 AVG 0.0555 0.0040 AVERAGE
R12 FIT 0.0540 0.0037 FROM FIT

D<sup>0</sup> -> (K<sup>0</sup> π<sup>0</sup>)/total (P9)
R13 (0.06) OR LESS CL=.90 SCHARRE 78 SMAG E+E- 3.77 GEV
R13 S 8 0.040 0.018 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R13 S THIS ENTRY WAS CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM
R13 S BALTRUSA 86. THE UNCORRECTED VALUE WAS 0.022+-0.011.

D<sup>0</sup> -> (K<sup>0</sup> ρ<sup>0</sup>)/total (P17)
R14 S 1 0.0013 0.0094 0.0013 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R14 S THIS ENTRY WAS CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM
R14 S BALTRUSA 86. THE UNCORRECTED VALUE WAS 0.0010,0.006-0.001.

D<sup>0</sup> -> (K<sup>-</sup> ρ<sup>+</sup>)/total (P16)

R15 S 31 0.129 0.051 0.053 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R15 A 5 0.075 0.053 0.041 SUMMERS 84 TPS 0 PHOTOPRODUCTION
R15 S THIS ENTRY WAS CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM
R15 S BALTRUSA 86. THE UNCORRECTED VALUE WAS 0.072+-0.030-0.031.
R15 A THIS ENTRY WAS CORRECTED BY US USING B(K<sup>-</sup> P1+)=0.056 FROM BALTRUSA
R15 A 86. THE UNCORRECTED VALUE WAS 0.032+-0.023.
R15 AVG 0.099 0.035 AVERAGE

D<sup>0</sup> -> (K\*(892)<sup>-</sup> π<sup>+</sup>)/total (P14)

R16 S 25 0.070 0.027 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R16 A 2 0.021 0.033 0.021 SUMMERS 84 TPS 0 PHOTOPRODUCTION
R16 S THIS ENTRY WAS CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM
R16 S BALTRUSA 86. THE UNCORRECTED VALUE WAS 0.034+-0.014.
R16 A THIS ENTRY WAS CORRECTED BY US USING B(K<sup>-</sup> P1+)=0.056 FROM BALTRUSA
R16 A 86. THE UNCORRECTED VALUE WAS 0.034+-0.039.
R16 AVG 0.071 0.025 AVERAGE

D<sup>0</sup> -> (K\*(892)<sup>0</sup> π<sup>0</sup>)/total (P15)

R17 S 4 0.025 0.040 0.025 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R17 A 2 0.021 0.033 0.021 SUMMERS 84 TPS 0 PHOTOPRODUCTION
R17 S THIS ENTRY WAS CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM
R17 S BALTRUSA 86. THE UNCORRECTED VALUE WAS 0.014+-0.023-0.014.
R17 A THIS ENTRY WAS CORRECTED BY US USING B(K<sup>-</sup> P1+)=0.056 FROM BALTRUSA
R17 A 86. THE UNCORRECTED VALUE WAS 0.009+-0.014.
R17 AVG 0.023 0.021 AVERAGE

D<sup>0</sup> -> (π<sup>+</sup> π<sup>-</sup> π<sup>0</sup>)/total (P29)

R18 A 10 0.011 0.004 BALTRUS2 85 SMK3 E+E- 3.77 GEV
R18 A ALL EVENTS CONSISTENT WITH RHOD P10. STATISTICAL AND SYSTEMATIC
R18 A ERRORS ARE 0.004 AND 0.002.

D<sup>0</sup> -> (K<sup>-</sup> π<sup>+</sup> π<sup>0</sup>)/total (P8)

R19 7 (0.12) (0.06) SCHARRE 78 SMAG E+E- 3.77 GEV
R19 S 37 0.152 0.054 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R19 B 93 0.175 0.018 BALTRUSA 86 SMK3 E+E- 3.77GEV ECM
R19 S THIS ENTRY WAS CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM
R19 S BALTRUSA 86. THE UNCORRECTED VALUE WAS 0.085+-0.032.
R19 S THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.013 AND 0.013.
R19 AVG 0.173 0.017 AVERAGE

D<sup>0</sup> -> (K<sup>-</sup> π<sup>+</sup> π<sup>0</sup> non-resonant)/total (P32)

R20 S 0.043 OR LESS CL=.90 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R20 A 21 (0.12) (0.065) SUMMERS 84 TPS 0 PHOTOPRODUCTION
R20 S THIS ENTRY WAS CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM
R20 A THIS ENTRY WAS CORRECTED BY US USING B(K<sup>-</sup> P1+)=0.056 FROM BALTRUSA
R20 A 86. THE UNCORRECTED VALUE WAS 0.052+-0.029.

D<sup>0</sup> -> (K<sup>0</sup> π<sup>+</sup> π<sup>-</sup>)/total (P3)

R21A S 28 0.103 0.029 PERUZZI 77 SMAG E+E- 3.77GEV ECM
R21B S 32 0.067 0.019 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R21 S CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM BALTRUSA 86.
R21A THE UNCORRECTED VALUE WAS 0.040+-0.013.
R21B THE UNCORRECTED VALUE WAS 0.038+-0.012.
R21 AVG 0.078 0.017 AVERAGE
R21 FIT 0.085 0.014 FROM FIT

D<sup>0</sup> -> (K<sup>0</sup> π<sup>+</sup> π<sup>-</sup> non-resonant)/total (P33)

R22 S 10 0.020 0.017 0.016 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R22 S CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM BALTRUSA 86.
R22 S THE UNCORRECTED VALUE WAS 0.011+-0.009.

D<sup>0</sup> -> (π<sup>+</sup> π<sup>-</sup> π<sup>+</sup> π<sup>-</sup>)/total (P18)

R23 B 9 0.015 0.006 BALTRUS2 85 SMK3 E+E- 3.77 GEV
R23 B STATISTICAL AND SYSTEMATIC ERRORS ARE 0.006 AND 0.002.

D<sup>0</sup> -> (K<sup>-</sup> π<sup>+</sup> π<sup>+</sup> π<sup>-</sup>)/total (P2)

R24A S 44 0.080 0.024 PERUZZI 77 SMAG E+E- 3.77GEV ECM
R24B S 185 0.152 0.030 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM
R24 6 0.10 0.04 AGUILAR 84 HYBR P1-P PP 360 GEV
R24 8 0.071 0.025 AGUILAR2 84 HYBR 0 P1-P 360 GEV
R24 C 992 0.118 0.014 BALTRUSA 86 SMK3 E+E- 3.77GEV ECM
R24 S CORRECTED BY US USING THE SIGMA D0=4.48 NB FROM BALTRUSA 86.
R24A THE UNCORRECTED VALUE WAS 0.032+-0.011.
R24B THE UNCORRECTED VALUE WAS 0.085+-0.021.
R24 C THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.009 AND 0.011.
R24 AVG 0.107 0.012 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)
R24 FIT 0.1093 0.0096 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1) (SEE IDEOGRAM BELOW)

D<sup>0</sup> -> (K<sup>-</sup> π<sup>+</sup> π<sup>0</sup> π<sup>0</sup>)/total (P20)

R25 1 SEEN ADEVA 81 HYBR P1-P -->D0 DOBAR

D<sup>0</sup> -> (π<sup>+</sup> π<sup>-</sup>)/(K<sup>-</sup> π<sup>+</sup>) (P5)/(P1)

R26 (0.07) OR LESS CL=.90 PICCOLO 77 SMAG E+E- 4.03 GEV ECM
R26 B 39 0.033 0.015 ABRAMS 79 SMK2 E+E- 3.77GEV ECM
R26 B 39 0.033 0.015 BALTRUS2 85 SMK3 E+E- 3.77 GEV
R26 B STATISTICAL AND SYSTEMATIC ERRORS ARE 0.010 AND 0.006.
R26 AVG 0.0330 0.0094 AVERAGE
R26 FIT 0.0330 0.0094 FROM FIT

D<sup>0</sup> -> (K<sup>+</sup> K<sup>-</sup>)/(K<sup>-</sup> π<sup>+</sup>) (P7)/(P1)

R27 (0.07) OR LESS CL=.90 PICCOLO 77 SMAG E+E- 4.03GEV ECM
R27 0.113 0.030 ABRAMS 79 SMK2 E+E- 3.77GEV ECM
R27 B 118 0.122 0.022 BALTRUS2 85 SMK3 0 E+E- 3.77 GEV
R27 B THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.018 AND 0.012.
R27 AVG 0.119 0.018 AVERAGE
R27 FIT 0.119 0.018 FROM FIT

D<sup>0</sup> -> (K<sup>0</sup> π<sup>+</sup> π<sup>-</sup>)/(K<sup>-</sup> π<sup>+</sup>) (P3)/(P1)

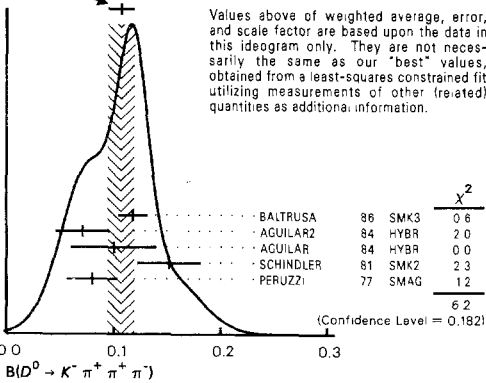
R28 116 2.8 1.0 PICCOLO 77 SMAG E+E-4.03,4.41ECM
R28 35 1.7 0.8 AVERY 80 SPEC GAMMA NUC-->D+
R28 AVG 2.13 0.62 AVERAGE
R28 FIT 1.57 0.28 FROM FIT



# Stable Particle Full Listings

$D^0, D_s^\pm$

WEIGHTED AVERAGE  
0.107 ± 0.012 (ERROR SCALED BY 1.2)



$D^0 \rightarrow (K^- \pi^+ \pi^+ \pi^-)/(K^- \pi^+ \pi^-)$  (P2)/(P1)  
 R29 P 2.14 2.2 0.8 PICCOLO 77 SMAG E-E-4.03,4.41ECM  
 R29 E 10 2.0 1.0 BAILEY 83 SPEC P1 BE --> D0  
 R29 E 2.17 0.36 ALBRECHT 85 ARG E+E- 10 GEV ECM  
 R29 P THIS CHANNEL DOMINATED BY  $K^- \pi^+ \rho^0$  (85-15 PERCENT).  
 R29 P  $K^- \pi^+ \rho^0$  AND  $K^- \rho^0$  CONSISTENT WITH 0,  $K^- \rho^0$  FRAC IS 0.1-0.1.  
 R29 E STATISTICAL AND SYSTEMATIC ERRORS ARE 0.28 AND 0.25.  
 R29 E NOT INDEPENDENT OF  $(K^- \pi^+ \pi^-)$ /TOTAL.  
 R29 AVG 2.16 0.31 AVERAGE  
 R29 FIT 2.02 0.20 FROM FIT

$D^0 \rightarrow (K^- \pi^+ (1320)^+)/(K^- \pi^+ \pi^+ \pi^-)$  (P25)/(P2)  
 R31 FOLLOWED BY DECAY  $A_2^+ \rightarrow \pi^+ \pi^+ \rho^0$  (B.R.=0.35)  
 R31 0.06 OR LESS PICCOLO 77 SMAG E-E-4.03,4.41ECM

$D^0 \rightarrow (K^- \pi^+ \rho^0)/(K^- \pi^+ \pi^+ \pi^-)$  (P25)/(P2)  
 R32 180 0.85 0.11 0.22 PICCOLO 77 SMAG E+E-4.03,4.41ECM  
 R32 2 (0.2) (0.2) CL=.90 BAILEY 83 SPEC P1 BE --> D0

$D^0 \rightarrow (\bar{K}^0(892)^0 \pi^+ \pi^-)/(K^- \pi^+ \pi^+ \pi^-)$  (P24)/(P2)  
 R33 FOLLOWED BY DECAY  $K^* \text{BAR}(892)^0 \rightarrow K^- \pi^+$  (B.R.=0.67)  
 R33 P 0 (0.10) (0.2) (0.0) PICCOLO 77 SMAG E+E-4.03,4.41ECM  
 R33 0 0.18 OR LESS CL=.90 BAILEY 83 SPEC P1 BE --> D0  
 R33 P CORRESPONDS TO 0.33 AT CL=.90.

$D^0 \rightarrow (\bar{K}^0 \phi)/(K^- \pi^+ \pi^-)$  (P30)/(P3)  
 R34 26 0.186 0.052 ALBRECHT 85 ARG E+E- 10 GEV ECM

$D^0 \rightarrow (\bar{K}^0 K^+ K^-)/(K^- \pi^+ \pi^-)$  (P31)/(P3)  
 R35 A 52 0.185 0.055 ALBRECHT 85 ARG E+E- 10 GEV ECM  
 R35 A RESONANT CONTRIBUTIONS TO KOBAR  $K^+ K^-$  ARE NOT DISTINGUISHED  
 R35 A (KOBAR PH1 IS INCLUDED).

$D^0 \rightarrow (\bar{K}^0(892)^0 \rho^0)/(K^- \pi^+ \pi^+ \pi^-)$  (P22)/(P2)  
 R36 FOLLOWED BY DECAY  $K^* \text{BAR}(892)^0 \rightarrow K^- \pi^+$  (B.R.=0.67)  
 R36 20 0.10 0.11 0.10 PICCOLO 77 SMAG E+E-4.03,4.41ECM  
 R36 5 (0.5) (0.2) CL=.90 BAILEY 83 SPEC P1 BE --> D0

$D^0 \rightarrow (\pi^- \pi^- \pi^+ \pi^+)/(K^- \pi^+ \pi^+ \pi^-)$  (P18)/(P2)  
 R37 0.21 OR LESS CL=.90 SCHINDLER 81 SMK2 E+E- 3.771 GEV ECM

$D^0 \rightarrow (\bar{K}^0 \phi)/\text{total}$  (P30)  
 R38 B 0.011 0.008 0.005 BALTRUS2 86 SMK3 E+E- 3.77 GEV  
 R38 B STAT ERROR (+.007-.005) AND SYST ERROR (-.004-.002) ADDED IN QUADR.

$D^0 \rightarrow (\bar{K}^0 K^+ K^- \text{ non-resonant})/\text{total}$  (P34)  
 R39 B 0.011 0.005 0.004 BALTRUS2 86 SMK3 E+E- 3.77 GEV  
 R39 B EXCLUDES CONTRIBUTIONS FROM D0 --> KOBAR PH1.  
 R39 B STAT ERROR (-.004-.003) AND SYST ERROR (-.003-.002) ADDED IN QUADR.

$D^0 \rightarrow (\bar{K}^0 K^0)/\text{total}$  (P35)  
 R40 0.006 OR LESS CL=.90 BALTRUS2 86 SMK3 E+E- 3.77 GEV

$D^0 \rightarrow (K^0 K^- \pi^+ + \bar{K}^0 K^+ \pi^- \text{ non-resonant})/\text{total}$  (P36)  
 R41 B 0.016 OR LESS CL=.90 BALTRUS2 86 SMK3 E+E- 3.77 GEV  
 R41 B EXCLUDES CONTRIBUTIONS FROM D0 -->  $K^*(892)^0$ .

$D^0 \rightarrow (\bar{K}^*(892)^0 K^0 + K^*(892)^0 \bar{K}^0)/\text{total}$  (P37)  
 R42 0.0073 OR LESS CL=.90 BALTRUS2 86 SMK3 E+E- 3.77 GEV

$D^0 \rightarrow (K^*(892)^- K^+ + K^*(892)^+ K^-)/\text{total}$  (P38)  
 R43 0.011 0.005 BALTRUS2 86 SMK3 E+E- 3.77 GEV

## REFERENCES FOR $D^0$

GOLDHABE 76 PRL 37 255	GOLDHABER,PIERRE,ABRAMS,ALAM+ (LBL-SLAC)
FELDMAN 77 PRL 38 1313	+PERUZZI,PICCOLO,ABRAMS,ALAM (SLAC-LBL)
GOLDHABE 77 PRL 698 503	GOLDHABER,WISS,ABRAMS,ALAM+ (LBL+SLAC)
PERUZZI 77 PRL 39 1301	+PICCOLO,FELDMAN+ (SLAC+LBL+WES+HAWA)
PICCOLO 77 PRL 708 260	+PERUZZI,LUTH,NGUYEN,WISS,ABRAMS+(SLAC-LBL)
BALTAY 78 PRL 41 73	+CAROUMBALIS,FRENCH,HIBBS,HYLTON+(COLU+BNL)
SCHARRE 78 PRL 40 74	+BARBARO-GALTIERI+ (SLAC-LBL+WES+HAWA)
VUILLEMI 78 PRL 41 1149	VUILLEMIN,FELDMAN+ (LBL-SLAC+WES+HAWA)
ABRAMS 79 PRL 43 481	+ALAM,BLOCKER,BOYARSKI+ (SLAC-LBL)
ARMENISE 79 PRL 868 115	+ERRIQUEZ+ (BARI+CERN+EPOL=MILA+ORSA)
ATIYA 79 PRL 43 414	+HOLMES,KNAPP,LEE+ (COLU+ILL+FNAL)

ALASTIA 80 NP 8176 13	(ANKA+LBN+CERN+DUUC+LOUC+KEYN+PISA+ROMA+)
ASTON 80 PL 948 113	+BONN+CERN+EPOL+GLAS+LANC+MCHS+LALO+LNP+)
AVERY 80 PRL 44 1309	+WISS,BUTLER,GLADDING+ (ILL+FNAL+COLU)
BACINO 80 PRL 45 329	+FERGUSON+ (UCLA+SLAC+STAN+UCI+STON)
ZHOLENTZ 80 PL 968 214	+KURDADZE,LELCHUK,MISHNEV,NIKITIN+ (NOVO)
ALSD 81 YAD.PHYS.34 1471	ZHOLENTZ+ (NOVO)
ADEVA 81 PL 1028 285	+ASULAR-BENITEZ+ (LEBC-EHS COLLABORATION)
BALLAGH 81 PR D24 7	+BINGHAM+ (LBL+UCB+FNAL+HAMA+WASH+WISC)
ALSD 80 PL 898 423	BALLAGH+ (LBL+UCB+FNAL+HAMA+WASH+WISC)
FIORINO 81 LNC 30 166	FIORINO+ (PHOTON-EMUL, OMEGA-PHOTON COLLS.)
FUCHI 81 LNC 31 199	+HOSHINO,MITSUSHI+ (NAGO-AICHI+TOKY+YOKO)
SCHINDLER 81 PR D24 78	SCHINDLER,ALAM,BOYARSKI+ (SLAC+LBL)
BODEK 82 PL 1138 82	+BREEDON,COLEMAN+ (ROCH+CIT+CHIC+FNAL+STAN)
USHIDA 82 PRL 48 844	(AICH+FNAL+KOBE+SEOU+MCGI+MAGO+OSU+OKAY+)
AGUILAR 85 PL 1228 312	LEBC-EHS COLLAB., AGUILAR-BENITEZ+ (CERN+)
BADERTSCH 85 PL 1238 471	BADERTSCHER,HAHN,HUGENTOBLER+ (BERN+MPI)
BAILEY 83 PL 1328 237	ACCMOR COLLAB., +BARDSLEY,BECKER+ (CERN+)
ABE 84 PR D30 1	SLAC HYBRID FACILITY PHOTON COL. (BIRM+BROW)
ADAMOVIĆ 84 PL 1408 123	WASB C., ADAMOVIĆ+ (BGNA+CERN+FIRZ+GENO+)
AGUILAR 84 PL 1358 237	LEBC-EHS C. (ANIK+BRUX+CERN+MADR+MONS+NIJM+)
AGUILAR2 84 PL 1468 266	LEBC-EHS C. (ANIK+BRUX+CERN+MADR+MONS+NIJM+)
ALTHOFF 84 PL 1388 317	TASSO C. (ACGI+DOMN+DESY+HAMB+LOIC+OXF+RL+)
DERRICK 84 PRL 53 1971	HRS C., DERRICK+(ANL+IND+MICH+PURD+LBL+SLAC)
SUMMERS 84 PRL 52 410	+ (UCSB+CARL+COLO+FNAL+TNTO+OKLA+CNRC)
VELTON 84 PRL 52 2019	MARK III COLLABORATION (SLAC+LBL+HARV)
ALBRECHT 85 PL 1588 235	ARGUS C. (DESY+DORT+HEID+IPPC+KANS+LUND+)
ALBRECHT2 85 PL 1588 525	ARGUS C. (DESY+DORT+HEID+IPPC+KANS+LUND+)
AUBERT 85 PL 1558 461	EMC COLL. (CERN+DESY+FREI+KIEL+LANC+LAPP+)
BAILEY 85 ZPHY C28 357	ABCCMR COLLAB. (ANIK+BRIS+CERN+CRAC+RL+MPI)
ALBRECHT 85 PRL 54 1976	MARK III C., BALTRUSAITIS+ (CIT+USC+ILL+)
BALTRUS2 85 PRL 55 150	MARK III C., BALTRUSAITIS+ (SLAC-CIT+USC+)
YAMAMOTO 85 PRL 54 522	DELCO COLLAB. (CIT+SLAC+STAN)
YAMAMOTO2 85 PR 052 2901	DELCO COLLAB. YAMAMOTO+ (CIT+SLAC+STAN)
BALTRUSA 86 SLAC-PUB-3861	MARK III C., BALTRUSAITIS+ (SLAC-CIT+USC+)
BALTRUS2 86 SLAC-PUB-3858	MARK III C., BALTRUSAITIS+ (SLAC-CIT+USC+)

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA LISTINGS

NGUYEN 77 PRL 39 262 -WISS,ABRAMS,ALAM,BOYARSKI+ (LBL+SLAC)J

REVIEWS  
 BARBARO 78 LBL-8537 A.BARBARO-GALTIERI (ERICE 1978) (LBL)  
 WOJCIK 78 SLAC-PUB-2232 S.WOJCIK (SLAC SUMMER INST. 1978) (SLAC)  
 KIRKBY 79 SLAC-PUB-2619 J.KIRKBY (LEPTON CONF. BATAVIA, 1979) (SLAC)  
 TRILLING 81 PRL 75 57 G.H.TRILLING (LBL+USC)J

## CHARMED STRANGE MESON

$D_s^\pm$  was  $F^\pm$

$$I(J^P) = 0(0^-)$$

QUANTUM NUMBERS NOT MEASURED. VALUES ARE ASSIGNED HERE ASSUMING CHARMED-STRANGE GROUND STATE  $D_s^0$  MESON. CHEN 83 OBSERVATIONS ARE CONSISTENT WITH  $J_{d0}$ .

### NOTE ON THE $D_s$ MESON (by W. Toki, SLAC)

The mass of the  $D_s(\bar{c}s)$  meson (formerly called the  $F$  meson) is not as controversial as it was two years ago. Recent results from ARGUS, HRS, and ACCMOR confirmed the decay  $D_s^\pm \rightarrow \phi \pi^\pm$ , first seen by CLEO. The table summarizes the published masses and branching fractions (B) to  $\phi \pi^\pm$ .

Expt.	Mass (MeV)	Mode	Production ( $D_s^\pm$ )	B $\rightarrow \phi \pi^\pm$
CLEO <sup>1</sup>	1970 $\pm 5 \pm 5$	$\phi \pi^\pm$	$e^+ e^-$	4.4 $\pm 1.1\%$
TASSO <sup>2</sup>	1975 $\pm 9 \pm 10$	$\phi \pi^\pm$	$e^+ e^-$	13 $\pm 5\%$
ARGUS <sup>3</sup>	1972.8 $\pm 3.1 \pm 3.0$	$\phi \pi^\pm$	$e^+ e^-$	3.2 $\pm 0.7 \pm 0.5\%$
ARGUS	1975.7 $\pm 4.7 \pm 3.0$	$\phi 3 \pi^\pm$	$e^+ e^-$	
HRS <sup>4</sup>	1963 $\pm 3 \pm 3$	$\phi \pi^\pm$	$e^+ e^-$	3.3 $\pm 1\%$
ACCMOR <sup>5</sup>	1975 $\pm 4$	$K^+ K^- \pi^\pm$	Hadron + Be	

For notation, see key on page 91.

# Stable Particle Full Listings

$D_s^\pm$

In addition, ACCMOR has published a  $D_s^\pm$  lifetime based on 3 events of  $(3.2^{+3.0}_{-1.3}) \times 10^{-13}$  sec, and HRS has submitted for publication a lifetime based on 17 events measured at 1964 MeV of  $(3.5^{+2.4}_{-1.8} \pm 0.9) \times 10^{-13}$  sec.<sup>6</sup>

The TPC and ARGUS collaborations have published evidence for the  $D_s^*$ . TPC observed  $D_s^{*\pm} \rightarrow \gamma K^+ K^- \pi^\pm$  with a  $D_s^* - D_s$  mass difference of  $139.5 \pm 8.3 \pm 9.7$  MeV.<sup>7</sup> ARGUS observed  $D_s^{*\pm} \rightarrow \gamma \phi \pi^\pm$  with a  $D_s^* - D_s$  mass difference of  $144 \pm 9 \pm 7$  MeV.<sup>8</sup> In both measurements the  $D_s^*$  was seen only when the  $\gamma$  was associated with a  $K^+ K^- \pi^\pm$  mass near 1970 MeV.

In  $\pi^- p$  collisions the LEBC-EHS collaboration has published upper limits for  $D_s^\pm \rightarrow \phi \pi^\pm$ .<sup>9</sup> No cross section has been given by ACCMOR, so a comparison is not possible.

The evidence for the  $D_s^\pm \rightarrow \eta \pi^\pm$  decay mode disagrees with the  $\phi \pi^\pm$  evidence. The original DASP measurement of  $D_s^\pm \rightarrow \eta \pi^\pm$  had a mass of  $2030 \pm 60$  MeV.<sup>10,11</sup> Although this mass is consistent with the 1970 MeV mass, some results are in conflict with CRYSTAL BALL measurements. DASP observed a large threshold in inclusive  $\eta$  production slightly above charm threshold, whereas CRYSTAL BALL measured no increase.<sup>12</sup> Also the DASP cross section  $\sigma(e^+e^- \rightarrow D_s D_s^*) \times B(D_s^\pm \rightarrow \eta \pi^\pm)$  was  $0.41 \pm 0.18$  nb, and the CRYSTAL BALL upper limit to this mode was 0.09–0.29 nb.<sup>13</sup> In photoproduction, the Omega group has published evidence for  $D_s \rightarrow \eta \pi^\pm, \eta 3\pi^\pm, \eta 3\pi^\pm$ , and  $\phi \rho^\pm$  at  $2020 \pm 10$  MeV and also has set upper limits for  $D_s \rightarrow \phi \pi^\pm$ .<sup>14–16</sup> The upper limit on  $\phi \pi^\pm$  is 35 times smaller than the combination of the other modes they have observed.

The recent evidence indicates that the  $D_s$  mass is 1970 MeV. The weighted mean of the masses in the table (using statistical errors only) is  $1970 \pm 3$  MeV.

The branching fractions for  $D_s^\pm \rightarrow \phi \pi$  have been calculated using a crude estimate of  $\sigma_{D_s D_s^*}$  by counting quarks, assuming an  $s\bar{s}$  sea of 15% and scaling the hadron rate. (See the CLEO and TASSO publications.) The TASSO result is  $2\sigma$  higher than the other results, which are in reasonable agreement with each other. The weighted mean of these estimated branching fractions is  $3.6 \pm 0.3\%$ .

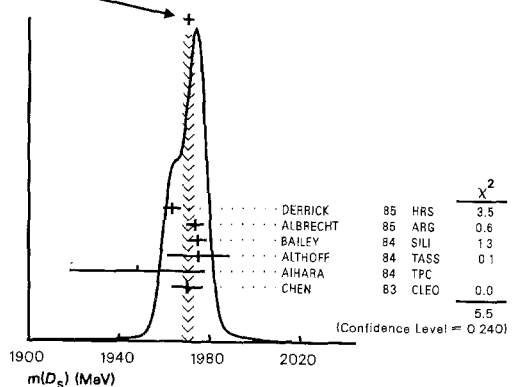
The TASSO and HRS groups have published  $x$  ( $2E / \sqrt{s}$ ) distributions. The TASSO distribution is peaked at  $x = 0.5$ , and the HRS distribution is softer. Neither result is statistically compelling.

## References

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$D_s^\pm$ MASS (MeV)		
M	4(2030.) (60.)	BRANDELIK 77 DASP -- IN BRANDELIK 79
M	6(2030.) (60.)	BRANDELIK 79 DASP +- E+E- ECM=4.62GEV
M	1(2017.) (25.)	AMMAR 80 HYBR + NU WIDEBAND
M	1(2026.) (56.)	USHIDA 80 EMUL - FNAL NU WIDEBAND
M	1(2089.) (121.)	USHIDA 80 EMUL + FNAL NU WIDEBAND
M	A 460(2020.) (22.)	ASTON 81 OMEG ++ GAMMA P
M	M 30(2049.) (15.)	ASTON2 81 OMEG +- GAMMA P
M	M (1970.) (10.)	ARGUS 83 ARG PRELIMINARY
M	C 17(2017.) (13.)	ATKINSON 83 OMEG +- GAMMA P
M	S 104 1970. 7.	CHEN 83 CLED -- E+E- ECM=10.5GEV
M	B 65 1948. 30.	AIHARA 84 TPC + E+E- 29 GEV ECM
M	S 49 1975. 14.	ALTHOFF 84 TASS +- E+E- ECM14-25GEV
M	S 5 1975.0 4.0	BALLBY 84 SILLI HAD-BE--> PHI PI+X
M	L 163 1973.6 4.0	ALBRECHT 85 ARG E+E- 10 GEV ECM
M	D 30 1963. 4.	DERRICK 85 HRS + E+E- 29 GEV ECM
M	A	ERROR QUOTED BY ASTON 81 IS 10 MEV STAT AND <20 MEV SYST.
M	A	AVERAGE OF THREE MODES LISTED IN SECTIONS R2, R3, AND R4 BELOW.
M	C	ATKINSON 83 MASS ERROR INCLUDES SYSTEMATIC UNCERTAINTIES.
M	S	STATISTICAL AND SYSTEMATIC ERRORS COMBINED IN QUADRATURE.
M	B	STAT. AND SYST. ERRORS COMBINED IN QUADR. FROM PEAK IN K+ K- PI
M	D	STATISTICAL AND SYSTEMATIC ERRORS ADDED IN QUADRATURE.
M	L	STATISTICAL AND SYSTEMATIC ERRORS ARE 2.6 AND 3.0.
M	M	
M	AVG	1970.5 +- 2.5 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2) (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
1970.5 ± 2.5 (ERROR SCALED BY 1.2)



# Stable Particle Full Listings

$D_s^\pm, B^\pm, B^0$

## $D_s^\pm$ MEAN LIFE (units $10^{-13}$ sec)

T	2	(2.24)	(2.78)	(1.05)	USHIDA	80	EMUL	NU	WIDEBAND
T	1	(1.4)			AMMAR	80	MYR	NU	WIDEBAND
T	4	1.9	1.3	0.7	USHIDA	83	EMUL	REPL	USHIDA 80
T	3	3.2	3.0	1.3	BAILEY	84	SILLI	HAD+BE-->	PHI PI+X
T	B	17	3.5	2.6	JUNG	86	HRS	+	E+E- -> PHI PI+X
T	B	STAT. ERROR (-2.4 -1.8)		2.0	AND SYST. ERROR (0.9) ADDED IN QUADR.				
T	AVG	2.80	1.57	0.74	AVERAGE				

## $D_s^\pm$ PARTIAL DECAY MODES

	DECAY MODE	DECAY MASSES
P1	$D_s^\pm \rightarrow \eta \pi^\pm$	549- 140
P2	$D_s^\pm \rightarrow \eta$ anything	549+ 0
P3	$D_s^\pm \rightarrow \eta \pi^\pm \pi^+ \pi^-$	549- 140+ 140+ 140
P4	$D_s^\pm \rightarrow \eta' \pi^\pm \pi^+ \pi^-$	958+ 140+ 140- 140
P5	$D_s^\pm \rightarrow \phi \rho^\pm$	1020- 769
P6	$D_s^\pm \rightarrow \phi \pi^\pm$	1020- 140
P7	$D_s^\pm \rightarrow \mu^\pm \nu$	106+ 0
P8	$D_s^\pm \rightarrow \phi \pi^\pm \pi^+ \pi^-$	1020+ 140+ 140+ 140

## $D_s^\pm$ BRANCHING RATIOS

$D_s^\pm \rightarrow (\eta \pi^\pm)/(\eta$ anything)	(P1)/(P2)
R1 A 6 (0.09) (0.06)	BRANDELIK 79 DASP E-E- ECM=4.42GEV
R1 A	DENOMINATOR INCONSISTENT WITH PARTRIDGE 81 (CRYSTAL BALL)
$D_s^\pm \rightarrow \eta \pi^\pm$	(P1)
R2 40 +- 9 EVENTS SEEN	ASTON 81 OMEG GAMMA P
R2 17 -- 6 EVENTS SEEN	ATKINSON 83 OMEG GAMMA P
$D_s^\pm \rightarrow \eta \pi^\pm \pi^+ \pi^-$	(P3)
R3 360 +- 90 EVENTS SEEN	ASTON 81 OMEG GAMMA P
$D_s^\pm \rightarrow \eta' \pi^\pm \pi^+ \pi^-$	(P4)
R4 60 +- 20 EVENTS SEEN	ASTON 81 OMEG GAMMA P
$D_s^\pm \rightarrow \phi \rho^\pm$	(P5)
R5 85 +- 26 EVENTS SEEN	ASTON2 81 OMEG GAMMA P
$D_s^\pm \rightarrow (\phi \pi^\pm)/total$	(P6)
R6 (SEEN)	ARGUS 83 ARG PRELIMINARY
R6 A 104 (0.044)	CHEN 83 CLEO +- E-E- ECM=10.5GEV
R6 A 49 (0.13) (0.05) (0.08)	ALTHOFF 84 TASS +- E-E- ECM=14-25GEV
R6 100 EVENTS SEEN	ALBRECHT 85 ARG E-E- 10 GEV ECM
R6 B 30 (0.033) (0.011)	DERRICK 85 HRS + E-E- 29 GEV ECM
R6 A	BOTH VALUES BASED ON SAME CRUDE ESTIMATE OF D/S+- PRODUCTION LEVEL.
R6 A	ALTHOFF 84 ERRORS ARE STATISTICAL AND SYSTEMATIC COMBINED IN
R6 A	QUADRATURE WITH ADDITIONAL NEGATIVE ERROR FOR D/S+- FROM PRIMARY B.
R6 B	SAME ASSUMPTIONS AS (A) APPLY; STATISTICAL ERRORS ONLY.
$D_s^\pm \rightarrow (\mu^\pm \nu)/total$	(P7)
R7 A 0 (0.03) OR LESS	AUBERT 83 SPEC MU- FE, 250 GEV
R7 A	AUBERT 83 OBTAIN THIS LIMIT ASSUMING THAT D/S+- PRODUCTION RATE IS 20 PERCENT OF TOTAL CHARM PRODUCTION RATE.
$D_s^\pm \rightarrow \phi \pi^\pm \pi^+ \pi^-$	(P8)
R8 62 EVENTS SEEN	ALBRECHT 85 ARG E-E- 10 GEV ECM
$D_s^\pm \rightarrow (\phi \pi^\pm \pi^+ \pi^-)/(\phi \pi^\pm)$	(P8/P6)
R9 A 1.11 0.46	ALBRECHT 85 ARG E-E- 10 GEV ECM
R9 A	STATISTICAL AND SYSTEMATIC ERRORS ARE 0.37 AND 0.28.

## REFERENCES FOR $D_s^\pm$

BRANDELIK 77 PL 708 132	BRANDELIK - (AACH+DESY+HAMB+MPIM+TOKY)
BRANDELIK 79 PL 808 412	BRANDELIK+ (AACH+DESY+HAMB+MPIM+TOKY)
AMMAR 80 PL 945 118	+ (KANS+FNAL+SERP+ITEP+CRAC+JINR+WASH+)
USHIDA 80 PRL 43 1053	(AICH+FNAL+KOE+SEU+MCGI+NAGO+OSU+OKAY+)
ASTON 81 PL 1008 91	(BONN+CERN+EPOL+GLAS+LANC+MCHS+LALO+LPP+)
ASTON2 81 NP 819 205	(BONN+CERN+EPOL+GLAS+LANC+MCHS+LALO+LPP+)
PARTRIDGE 81 PRL 47 760	PARTRIDGE+PECK+ (CIT+HARV+PRIN+STAN+SLAC)
ARGUS 83 CERN COUR. 23 423	ARGUS COLLABORATION (PRELIMINARY)
ATKINSON 83 ZPHY 177 1	(BONN+CERN+GLAS+LANC+MCHS+LPP+RL+SHEP+)
AUBERT 83 NP B213 31	EUROPEAN MUON COLLAB. (CERN+DESY+FREI+)
CHEN 83 PRL 51 634	+ (SYRA+VAND+HARV+OSU+CORN+ITHA+ROCH+RUTG+)
USHIDA 83 PRL 51 2362	(AICH+FNAL+KOE+SEU+MCGI+NAGO+OSU+OKAY+)
AIHARA 84 PRL 53 2465	TPC COLL. (LBL+UCLA+UCR+JHU+YALE+TOKY+MASA)
ALTHOFF 84 PL 136B 130	TASSO C. (AACH+BONN+DESY+HAMB+LOIC+OXF+RL+)
BAILEY 84 PL 139B 320	ACCMOR C. (ANIK+BRIS+CERN+CRAC+MPIM+RL)
ALBRECHT 85 PL 153B 343	ARGUS C. (DESY+DORT+HEID+IPPC+KANS+LUND+)
DERRICK 85 PRL 54 2568	HRS COLL. (ANL+IND+MICH+PURD+LBL+SLAC)
JUNG 86 PRL (TO BE PUBL.)	HRS C., -ABACHI (IND+ANL+MICH+PURD+LBL)
REVIEWS	
TRILLING 81 PRPL 75 57	G.H. TRILLING (LBL+UCB)

## BOTTOM MESONS

$B^\pm$

$$I(J^P) = \frac{1}{2}(0^-)$$

QUANTUM NUMBERS NOT MEASURED. VALUES SHOWN ARE QUARK MODEL PREDICTIONS. SEE ALSO THE LISTING FOR THE B (FOLLOWING THE ENTRY FOR THE B0) FOR MEASUREMENTS WHICH DO NOT IDENTIFY THE CHARGE STATE.

## $B^\pm$ MASS (MeV)

M A 6 (5270.8)	(3.0)	BEHREND 83 CLEO +- D*- PI+ PI+ + CC
M G 5271.2	3.0	GILES 84 CLEO E+E- UPSILON 4S
M A	STATISTICAL (2.3 MEV) AND SYSTEMATICAL (2.0 MEV) ERRORS COMBINED.	
M G	STATISTICAL AND SYSTEMATIC ERRORS ARE 2.2 AND 2.0.	

## $B^\pm$ PARTIAL DECAY MODES

	DECAY MODE	DECAY MASSES
P1	$B^+ \rightarrow \bar{D}^0 \pi^+$	1865+ 140
P2	$B^+ \rightarrow D^*(2010)^- \pi^+ \pi^+$	2007+ 140+ 140
P3	$B^+ \rightarrow J/\psi(3097) K^+$	3097+ 494
P4	$B^+ \rightarrow \rho^0 \pi^+$	769+ 140

B- MODES ARE CHARGE CONJUGATES OF THE ABOVE MODES.

## $B^\pm$ BRANCHING RATIOS

$B^+ \rightarrow (\bar{D}^0 \pi^+)/total$	(P1)
R1 CB 2 (0.023) (0.023)	BEHREND 83 CLEO E+ E-, UPSILON(4S)
R1 GB	GILES 84 CLEO E+E- UPSILON 4S
R1 G	STATISTICAL AND SYSTEMATIC ERRORS ARE 0.004 AND 0.003.
R1 B	CORRECTED BY US USING OUR NEW VALUE (0.054) FOR D0--> K- PI+.
R1 B	UNCORRECTED VALUES WERE A FACTOR 1.8 LARGER.
R1 B	INCLUDES CONTAMINATION FROM (D*0 PI- PI-) AND (D*- PI+).
R1 C	ADDED STATISTICAL AND SYSTEMATIC ERRORS IN QUADRATURE.
$B^+ \rightarrow (D^*(2010)^- \pi^+ \pi^+)/total$	(P2)
R2 C 6 0.027 0.017	BEHREND 83 CLEO E+ E-, UPSILON(4S)
R2 G	ADDED STATISTICAL AND SYSTEMATIC ERRORS IN QUADRATURE.
R2 G	CORRECTED BY US USING OUR NEW VALUE (0.054) FOR D0--> K- PI+.
R2 G	UNCORRECTED VALUES WERE A FACTOR 1.8 LARGER.
$B^+ \rightarrow (J/\psi(3097) K^+)/total$	(P3)
R3 1 0.0026 OR LESS CL-.90	GILES 84 CLEO + E+E- UPSILON 4S
$B^+ \rightarrow (\rho^0 \pi^+)/total$	(P4)
R4 0 0.0006 OR LESS CL-.90	GILES 84 CLEO - E+E- UPSILON 4S

## REFERENCES FOR $B^\pm$

BEHREND 83 PRL 50 881	+ (ROCH+RUTG+SYRA+VAND+CORN+ITHA+HARV+OSU)
GILES 84 PR D30 2279	CLEO C. (HARV+OSU+ROCH+RUTG+SYRA+VAND+CORN+)

$B^0$

$$I(J^P) = \frac{1}{2}(0^-)$$

QUANTUM NUMBERS NOT MEASURED. VALUES SHOWN ARE QUARK MODEL PREDICTIONS. SEE ALSO THE LISTING FOR THE B (FOLLOWING THIS ENTRY) FOR MEASUREMENTS WHICH DO NOT IDENTIFY THE CHARGE STATE.

## $B^0$ MASS (MeV)

M A 5 (5274.2)	(2.8)	BEHREND 83 CLEO 0 D*- PI+ + CC
M G 5275.2	2.8	GILES 84 CLEO E+E- UPSILON 4S
M A	STATISTICAL (1.9 MEV) AND SYSTEMATICAL (2.0 MEV) ERRORS COMBINED.	
M G	STATISTICAL AND SYSTEMATIC ERRORS ARE 1.9 AND 2.0.	

## $B^0 - B^+$ MASS DIFFERENCE (MeV)

DM A (3.4) (3.6)	BEHREND 83 CLEO E+E-, UPSILON(4S)
DM G 4.0 3.4	GILES 84 CLEO E+E- UPSILON 4S
DM A	STATISTICAL (3.0) AND SYSTEMATICAL (2.0) ERRORS COMBINED.
DM G	STATISTICAL AND SYSTEMATIC ERRORS ARE 2.7 AND 2.0.

For notation, see key on page 91.

# Stable Particle Full Listings

$B^0, B$

## $B^0$ PARTIAL DECAY MODES

	DECAY MASSES
P1 $B^0 \rightarrow \bar{D}^0 \pi^+ \pi^-$	1865- 140+ 140
P2 $B^0 \rightarrow D^*(2010)^- \pi^+$	2007+ 140
P3 $B^0 \rightarrow D^*(2010)^- \rho^+$	2007- 769
P4 $B^0 \rightarrow J/\psi(3097) K^+ \pi^-$	3097+ 494+ 140
P5 $B^0 \rightarrow \pi^+ \pi^-$	140- 140
P6 $B^0 \rightarrow e^+ e^-$	.511+.511
P7 $B^0 \rightarrow \mu^+ \mu^-$	106+ 106
P8 $B^0 \rightarrow (e^+ \mu^-) + (e^- \mu^+)$	.511- 106
P9 $B^0 \rightarrow \bar{B}^0$ via mixing	

BOBAR MODES ARE CHARGE CONJUGATES OF THE ABOVE MODES.

## B MEAN LIFE (units $10^{-13}$ sec)

T	14.	OR LESS	CL=.95	BARTEL	82 JADE E+E-,AVG ECM 34 GEV	
T	A	18.	7.2	FERNAND2	83 MAC E+E- AT ECM=29 GEV	
T	B	12.0	5.4	LOCKYER	83 SMK2 E+E- AT ECM=29 GEV	
T	C	18.3	5.3	ALTHOFF2	84 TASS 0 E+E- 30-46.8 GEV	
T	D	46	11.6	4.4	KLEM	84 DLCO E+E- AT ECM=29 GEV
T	E	2	POSS.DIFF.BETW. B+ AND B0	ALBANESE	85 MYR 350GEV P1-P EMUL	
T	A	THE STATISTICAL AND SYSTEMATIC ERRORS ARE 6 AND 4.				
T	B	THE STATISTICAL ERRORS ARE +4.5 AND -3.6, THE SYSTEMATIC ERROR IS				
T	B	3.0. THE LIFETIME IS AN AVERAGE OVER BOTTOM PARTICLES PRODUCED.				
T	C	STATISTICAL AND SYSTEMATIC ERRORS ARE +3.8 -3.7 AND +3.7 -3.4				
T	D	STATISTICAL ERRORS +3.7 -3.4 SYSTEMATIC ERRORS +2.3.				
T	E	THE MEAN FLIGHT TIME FOR THE ONE B0 WAS 5E-13 SEC WHILE THE ONE B-				
T	E	WAS 0.8E-13 SEC.				
T	AVG	14.2	2.7	AVERAGE		

## B PARTIAL DECAY MODES

- P1  $B \rightarrow e \nu$  hadrons
- P2  $B \rightarrow \mu \nu$  hadrons
- P3  $B \rightarrow e^+ e^-$  anything
- P4  $B \rightarrow \mu^+ \mu^-$  anything
- P5  $B \rightarrow K$  anything
- P6  $B \rightarrow J/\psi(3097)$  anything
- P7  $B \rightarrow D^0$  anything
- P8  $B \rightarrow \rho$  anything
- P9  $B \rightarrow \Lambda$  anything
- P10  $B \rightarrow e \nu$  hadrons (noncharm)
- P11  $B \rightarrow D^*(2010)^+$  anything
- P12  $B \rightarrow D^0 \pi^+, D^- \pi^+, D^*(2010)^0 \pi^+,$  or  $D^*(2010)^- \pi^+$

## $B^0$ BRANCHING RATIOS

- $B^0 \rightarrow (\bar{D}^0 \pi^+ \pi^-)/total$  (P1)
  - R1 G 5 0.07 0.05 BEHREND 83 CLEO E+E-, UPSILON(4S)
  - R1 G GILES 84 ADDED STATISTICAL AND SYSTEMATIC ERRORS IN QUADRATURE.
  - R1 G CORRECTED BY US USING OUR NEW VALUE (0.054) FOR  $D^0 \rightarrow K^- \pi^+$ .
  - R1 G UNCORRECTED VALUES WERE A FACTOR 1.8 LARGER.
  - R1 G INCLUDES CONTAMINATION FROM ( $D^0 \pi^+ \pi^-$ ).
- $B^0 \rightarrow (D^*(2010)^- \pi^+)/total$  (P2)
  - R2 B 5 (0.014) (0.011) BEHREND 83 CLEO 0 E+E-, UPSILON(4S)
  - R2 G 4 0.017 0.007 GILES 84 CLEO E+E- UPSILON 4S
  - R2 B ADDED STATISTICAL AND SYSTEMATIC ERRORS IN QUADRATURE.
  - R2 B CORRECTED BY US USING OUR NEW VALUE (0.054) FOR  $D^0 \rightarrow K^- \pi^+$ .
  - R2 B UNCORRECTED VALUES WERE A FACTOR 1.8 LARGER.
  - R2 B ASSUMES  $B(D^+ \rightarrow D^0 \pi^+) = 0.6$
  - R2 G STATISTICAL AND SYSTEMATIC ERRORS ARE 0.005 AND 0.005. ASSUMES
  - R2 G  $B(D^+ \rightarrow D^0 \pi^+) = 0.6$ . ALSO ASSUMES CHARGED/NEUTRAL B RATIO FROM
  - R2 G UPS(4S) IS 6 TO 4. QUOTED NUMBER IS CORRECTION DESCRIBED IN
  - R2 G CHEN 85.
- $B^0 \rightarrow (D^*(2010)^- \rho^+)/total$  (P3)
  - R3 C 19 0.081 0.066 0.038 CHEN 85 CLEO 0 E+E-, UPSILON(4S)
  - R3 C USES  $B(D^+ \rightarrow D^0 \pi^+) = 0.6$  AND  $B(UPS(4S) \rightarrow B^0 \text{--} \text{BAR}) = 0.4$ . STATISTICAL
  - R3 C AND SYSTEMATIC ERRORS COMBINED IN QUADRATURE.
- $B^0 \rightarrow (J/\psi(3097) K^+ \pi^-)/total$  (P4)
  - R4 2 0.0063 OR LESS CL=.90 GILES 84 CLEO E+E- UPSILON 4S
- $B^0 \rightarrow (\pi^+ \pi^-)/total$  (P5)
  - R5 4 0.0005 OR LESS CL=.90 GILES 84 CLEO 0 E+E- UPSILON 4S
- $B^0 \rightarrow (e^+ e^-)/total$  (P6)
  - R6 TEST FOR DELTA(B) = 1 WEAK NEUTRAL CURRENT.
  - R6 0.0003 OR LESS CL=.90 GILES 84 CLEO 0 E+E- UPSILON 4S
- $B^0 \rightarrow (\mu^+ \mu^-)/total$  (P7)
  - R7 TEST FOR DELTA(B) = 1 WEAK NEUTRAL CURRENT.
  - R7 0.0002 OR LESS CL=.90 GILES 84 CLEO E+E- UPSILON 4S
- $B^0 \rightarrow ((e^+ \mu^-) + (e^- \mu^+))/total$  (P8)
  - R8 TEST OF LEPTON FAMILY NUMBER CONSERVATION.
  - R8 0.0003 OR LESS CL=.90 GILES 84 CLEO E+E- UPSILON 4S
- $B^0 \rightarrow (\bar{B}^0 \text{ via mixing})/total$  (P9)
  - R9 A (0.30) OR LESS CL=.90 AVERY 84 CLEO E+E- UPSILON 4S
  - R9 B 0.12 OR LESS CL=.90 SCHAAD 85 SMK2 0 E+E- ECM=29 GEV
  - R9 A SAME SIGM OILDETON EVENTS. LIMIT ASSUMES SEMILEPTONIC BR
  - R9 A FOR B+ AND B0 EQUAL. IF  $B^0/B^+$  RATIO < 0.58, NO LIMIT EXISTS.
  - R9 B LIMIT IS AVERAGE PROB. FOR HADRON CONTAINING B QUARK
  - R9 B TO PRODUCE A POSITIVE LEPTON. CM ENERGY HIGH ENOUGH TO ALLOW
  - R9 B PRODUCTION OF B-STRANGE HADRONS.

## REFERENCES FOR $B^0$

BEHREND 83 PRL 50 881 (ROCH+RUTG+SYRA+VAND+CORN+ITHA+HARV+OSU)  
 AVERY 84 PRL 53 1309 CLEO C.(CORN+ITHA+HARV+OSU+ROCH+RUTG+SYRA+)  
 GILES 84 PR D30 2279 CLEO C.(HARV+OSU+ROCH+RUTG+SYRA+VAND+CORN+)  
 CHEN 85 PR D31 2386 CLEO C.(SYRA+VAND+ALBA+CORN+HARV+ITHA+ROCH+)  
 SCHAAD 85 PL 1608 188 MARK II C., NELSON,ABRAMS+ (SLAC+LBL+HARV)

**B**

$$I(J^P) = \frac{1}{2}(0^-)$$

QUANTUM NUMBERS NOT MEASURED. VALUES SHOWN ARE QUARK MODEL  
 PREDICTIONS. THIS ENTRY LISTS MEASUREMENTS OF B MESON  
 PARAMETERS FOR WHICH THE CHARGE STATES ARE NOT SEPARATED.  
 MEASUREMENTS IN WHICH THE CHARGE STATE IS CLEARLY IDENTIFIED  
 ARE LISTED IN THE PRECEDING B+- AND B0 ENTRIES.

## B MASS (MeV)

M	A	5180 to 5290	ANDREWS	80 CLEO UPSIL(4S) THRESHOLD	
M	A	5289 OR LESS	FINOCCHIA	80 CUSB UPSIL(4S) THRESHOLD	
M	AB	5263 to 5278	SCHAMBERG	82 CUSB UPSIL(4S) THRESHOLD	
M	G	5273.0 2.4	GILES	84 CLEO E+E- UPSILON 4S	
M	A	REDETERMINATION OF THE CESR ENERGY SCALE CHANGED THE ANDREWS 80,			
M	A	FINOCCHIA 80, AND SCHAMBERGER 82 LIMITS FROM THEIR PUBLISHED			
M	A	VALUES TO THOSE GIVEN ABOVE.			
M	AB	SCHAMBERGER 82 DEDUCED ABOVE LIMITS FROM NON-OBSERVATION OF 50 MEV			
M	A	GAMMAS FROM $B^+ \rightarrow B \text{ GAMMA}$ AND UPSIL(4S) $\rightarrow B \text{ GAMMA}$ .			
M	G	STATISTICAL AND SYSTEMATIC ERRORS ARE 1.3 AND 2.0.			

## B BRANCHING RATIOS

- $B \rightarrow (e \nu \text{ hadrons})/total$  (P1)
  - R1 D 0.152 0.016 KLOPFENST 83 CUSB DIRECT E AT UPS(4S)
  - R1 E (0.156) (0.027) NELSON 83 SMK2 E+E- ECM=29 GEV
  - R1 G 0.111 0.052 ALTHOFF3 84 TASS E+E- 34-66GEV ECM
  - R1 F 0.120 0.009 CHEN 84 CLEO DIR E UPS(4S)
  - R1 H (0.146) (0.028) KOOP 84 DLCO E+E- ECM=29 GEV
  - R1 I (0.110) (0.021) AIHARA 85 TPC E+E- ECM=29 GEV
  - R1 J (0.149) (0.022) (0.019)PAL 86 DLCO E+E- 29 GEV ECM
  - R1 D STATISTICAL AND SYSTEMATIC ERRORS ARE 0.008 AND 0.014.
  - R1 E THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.021 AND 0.017.
  - R1 F THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.007 AND 0.005
  - R1 G THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.037 AND 0.040.
  - R1 H THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.018 AND 0.010.
  - R1 I THIS MEASUREMENT SUPERSEDES KOOP 84.
  - R1 ONLY THE EXPERIMENTS AT UPS(4S) ARE USED IN THE AVERAGE.
  - R1 AVG 0.1226 0.0078 AVERAGE
- $B \rightarrow (\mu \nu \text{ hadrons})/total$  (P2)
  - R2 A (0.105) (0.020) ADEVA2 83 MRKJ E+E- ECM=33-38.5GEV
  - R2 B (0.155) (0.054) (0.029) FERNAND1 83 MAC E+E- AT ECM=29 GEV
  - R2 D (0.117) (0.030) ALTHOFF 84 TASS E+E- ECM=34.5 GEV
  - R2 C 0.108 0.012 CHEN 84 CLEO DIR MU UPS(4S)
  - R2 E 0.112 0.013 LEVMAN 84 CUSB DIR MU UPS(4S)
  - R2 F (0.114) (0.031) BARTEL 85 JADE E+E- ECM=34.6 GEV
  - R2 A THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.015 AND 0.015.
  - R2 D THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.028 AND 0.010.
  - R2 C THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.006 AND 0.01.
  - R2 E THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.009 AND 0.01.
  - R2 F THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.018 AND 0.025.
  - R2 THE AVERAGE OF THE FOUR HIGH-ENERGY RESULTS IS 0.115+0.014.
  - R2 THESE EXPERIMENTS PRODUCE OTHER BOTTOM PARTICLES IN ADDITION TO
  - R2 THE B MESON.
  - R2 AVG 0.1098 0.0088 AVERAGE
- $B \rightarrow (e^+ e^- \text{ anything})/total$  (P3)
  - R3 TEST FOR DELTA(B) = 1 WEAK NEUTRAL CURRENT.
  - R3 0.05 OR LESS CL=.90 BEBEK 81 CLEO E+E- AT UPSIL(4S)
- $B \rightarrow (\mu^+ \mu^- \text{ anything})/total$  (P4)
  - R4 TEST FOR DELTA(B) = 1 WEAK NEUTRAL CURRENT.
  - R4 (0.017) OR LESS CL=.90 CHADWICK 81 CLEO E+E- AT UPSIL(4S)
  - R4 0.007 OR LESS CL=.95 ADEVA1 83 MRKJ E+E- 30-MCM=38 GEV
  - R4 0.07 OR LESS CL=.95 BARTEL 83 JADE E+E- 33-MCM=36.76GEV
  - R4 (0.02) OR LESS CL=.95 ALTHOFF 84 TASS E+E-, ECM=34.5 GEV
- $B \rightarrow (\text{dilepton anything})/total$  (P3+P4)
  - R5 TEST FOR DELTA(B) = 1 WEAK NEUTRAL CURRENT.
  - R5 (0.008) OR LESS CL=.90 MATTEUZZI 83 SMK2 E+E- AT ECM=29 GEV
  - R5 A 0.0062 OR LESS CL=.90 AVERY 84 CLEO E+E- UPSILON 4S
  - R5 A DETERMINE RATIO OF B+ TO B0 SEMILEPTONIC DECAYS TO BE IN THE RANGE
  - R5 A 0.25 TO 2.9
- $B \rightarrow (K \text{ anything})/total$  (P5)
  - R6 C SEEN BRODY 82 CLEO KAONS AT UPSIL(4S)
  - R6 D SEEN GIANNINI 82 CUSB KAONS AT UPSIL(4S)
  - R6 C ASSUMING UPSILON(4S)  $\rightarrow B \text{ BBAR}$ , A TOTAL OF 3.38+-0.34+-0.68 KAONS
  - R6 C PER UPSILON(4S) DECAY IS FOUND (THE SECOND ERROR IS SYSTEMATIC). IN
  - R6 C THE CONTEXT OF THE STANDARD B-DECAY MODEL, THIS LEADS TO A VALUE
  - R6 C FOR (B-QUARK  $\rightarrow$  C-QUARK)/(B-QUARK  $\rightarrow$  ALL) OF 1.09+-0.35+-0.13.
  - R6 D GIANNINI 82 AT CESR-CUSB OBSERVED 1.58+- .35 KO PER HADRONIC EVENT
  - R6 D MUCH HIGHER THAN 0.82+-0.10 BELOW THRESHOLD. CONSISTENT WITH
  - R6 D PREDOMINANT B- $\rightarrow$ C X DECAY.

# Stable Particle Full Listings

*B, p*

**B → (J/ψ(3097) anything)/total** (P6)  
 R7 (0.049)OR LESS CL=90 MATTEUZZI 83 SMK2 E+ E- AT ECM-29 GEV  
 R7 A 7 0.014 0.006 0.005 ALBRECHT 85 ARG E+E- NEAR UPS(4S)  
 R7 C 46 0.011 0.003 HAAS 85 CLEO E+E- NEAR UPS(4S)  
 R7 A STATISTICAL AND SYSTEMATIC ERRORS WERE ADDED IN QUADRATURE.  
 R7 A ALBRECHT 85 ALSO REPORT A CL=90 LIMIT OF 0.007 FOR B→ J/PSI + X  
 R7 A WHERE M(X) < 1 GEV.  
 R7 C THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.0021 AND 0.0023.  
 R7 C DIMUON AND DIELECTRON EVENTS USED.  
 R7  
 R7 AVG 0.0117 0.0026 AVERAGE

**B → (D<sup>0</sup> anything)/total** (P7)  
 R8 A 0.8 0.28 GREEN 83 CLEO E+ E- AT UPS1L(4S)  
 R8 A OBSERVED 0.8+0.2(STAT.)+0.2(SYST.) DO PER B DECAY.

**B → (p anything)/total** (P8)  
 R9 A 0.036 OR MORE ALAM 83 CLEO PROTONS AT UPS1(4S)  
 R9 A THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.006 AND 0.004. VALUES  
 R9 A ARE FOR (BR(B→p X)+BR(B→pBAR X))/2. DATA ARE CONSISTENT WITH  
 R9 A EQUAL YIELDS OF BARYONS AND ANTI-BARYONS. USING ASSUMED YIELDS  
 R9 A BELOW CUT, BR(B→pX)+0.03 NOT INCLUDING PROTONS FROM LAM.DECAYS.

**B → (Λ anything)/total** (P9)  
 R10 A 0.022 OR MORE ALAM 83 CLEO LAMBDA S AT UPS1(4S)  
 R10 A THE STATISTICAL AND SYSTEMATIC ERRORS ARE 0.007 AND 0.004. VALUES  
 R10 A ARE FOR (BR(LAMBDA X)+BR(LAMBDA BAR X))/2. DATA ARE CONSISTENT WITH  
 R10 A EQUAL YIELDS OF BARYONS AND ANTI-BARYONS. USING ASSUMED YIELDS  
 R10 A BELOW CUT, BR(B→LAMBDA X) 0.03.

**B → e ν hadrons (noncharm)/(e ν hadrons)** (P10/P1)  
 R11 (0.055)OR LESS CL=90 KLOPFENST 83 CUSB DIRECT E AT UPS(4S)  
 R11 0.04 OR LESS CL=90 CHEN 84 CLEO E+E-DIR E UPS4S

**B → (D<sup>+</sup>(2010)<sup>+</sup> anything)/total** (P11)  
 R12 A 510 0.23 0.09 CSORNA 85 CLEO E+E- UPS1LON(4S)  
 R12 A ASSUMES UPS(4S)→B BBAR, AND V-A MOMENTUM SPECTRUM TO  
 R12 A EXTRAPOLATE BELOW 1 GEV/C D\* MOMENTUM. STAT AND SYST. ERRORS  
 R12 A COMBINED IN QUADR.

**B → (D<sup>0</sup>π<sup>+</sup>, D<sup>-</sup>π<sup>+</sup>, D<sup>+</sup>(2010)<sup>0</sup>π<sup>+</sup>, D<sup>+</sup>(2010)<sup>-</sup>π<sup>+</sup>)/total** (P12)  
 R13 G 0.020 0.008 GILES 84 CLEO E+E- UPS1LON 4S  
 R13 G STATISTICAL AND SYSTEMATIC ERRORS ARE 0.006 AND 0.005.  
 R13 G NO DEPENDENCE ON D; USED FAST-PI MOMENTUM.

**REFERENCES FOR BOTTOM MESON B**

ANDREWS 80 PRL 45 219	- (CORN+HARV+ITHA+SYRA+ROCH+RUTG+VAND)
FINOCCHI 80 PRL 45 222	FINOCCHIARO, GIANNINI, - (STON+COLU+LSU)
BEBEK 81 PRL 46 84	+ (HARV+ITHA+SYRA+ROCH+RUTG+VAND+CORN)
CHADWICK 81 PRL 46 88	+GANGI+(ROCH+RUTG+SYRA+VAND-CORN+HARV+ITHA)
BARTEL 82 PL 1148 71	JADE C. (DESY+HAMB+HEID+LANC+MCHS+RHEL+TOKY)
BRODY 82 PRL 48 1070	CHEN, +(SYRA+VAND+CORN+ITHA+HARV+ROCH+RUTG)
GIANNINI 82 NP B206 1	+FINOCCHIARO, FRANZINI- (STON+COLU+LSU+MPI)
SCHAMBER 82 PR D26 720	SCHAMBERGER+ (STON+COLU+LSU+MPI)
ADEVA1 83 PRL 50 799	MARK J. C. (AACH+DESY+MIT+MADR+AIKO+BHEP+CIT)
ADEVA2 83 PRL 51 443	(VAND+HARV+OHIO+CORN+ITHA+ROCH+RUTG+SYRA)
ALAM 83 PRL 51 1143	JADE C. (DESY+HAMB+HEID+LANC+MCHS+RHEL+TOKY)
BARTEL 83 PL 132B 241	+MAC. C. (COLO+FRAS+HOUS+NEAS+SLAC+UTAH+WISC)
FERNAND1 83 PRL 50 2054	+ (RUTG+SYRA+VAND+HARV+OHIO+CORN+ITHA+ROCH)
FERNAND2 83 PRL 51 1022	KLOPFENSTEIN, CUSB C. (STON+COLU+CORN+LSU+)
GREEN 83 PRL 51 347	+JAROS, NELSON, ABRAMS, - (SLAC+LBL+UCB+HARV)
KLOPFENS 83 PL 130B 444	MATTEUZZI, ABRAMS, AMIDEI+ (SLAC+LBL+UCB+HARV)
LODYKER 83 PRL 51 1316	+BLONDEL, TRILLING+ (LBL+UCB+SLAC+HARV)
MATTEUZZ 83 PL 129B 141	TASSO C. (AACH+BNON+DESY+HAMB+LOIC+OXF+RL+)
NELSON 83 PRL 50 1542	TASSO C. (AACH+BNON+DESY+BRIS+HAMB+LOIC+)
ALTHOFF 84 ZPHY C22 219	TASSO C. (AACH+BNON+DESY+HAMB+LOIC+OXF+)
ALTHOFF2 84 PL 149B 524	CLEO C. (CORN+ITHA+HARV+OSU+ROCH+RUTG+SYRA+)
ALTHOFF3 84 PL 146B 443	CLEO C. (SYRA+VAND+CORN+ITHA+HARV+OSU+ROCH+)
AVERY 84 PRL 53 1309	CLEO C. (HARV+OSU+ROCH+RUTG+SYRA+VAND+CORN+)
CHEN 84 PRL 52 1084	DELCO COLLABORATION (CIT+SLAC+STAN)
GILES 84 PR D30 2279	DELCO COLLABORATION (CIT+SLAC+STAN)
KLEM 84 PRL 53 1875	CUSB C. (LSU+MPI+STON+COLU+CORN)
KOOP 84 PRL 52 970	TPC COLL. (LBL+UCI+UCR+JHU+YALE+TOKY+KASA)
LEWMAN 84 PL 141B 271	WAT5 EXPER. (BARI+CEEN+DUUC+LOUC+NAGO+INFN+)
AIHARA 85 ZPHY C27 39	ARGUS C. (DESY+DORT+HEID+IPPC+KANS+LUND+)
ALBARECHT 85 PL 162B 395	JADE C. (DESY+HAMB+HEID+LANC+MCHS+RL+TOKY+)
BARTEL 85 PL 163B 277	CLEO C. (SYRA+VAND+ALBA+CORN+ITHA+ROCH+)
CSORNA 85 PRL 54 1894	CLEO C. (OHIO+ROCH+RUTG+SYRA+VAND+ALBA+)
HAAS 85 PRL 55 1248	DELCO COLLABORATION (CIT+SLAC+STAN)
PAL 86 CALT-68-1283	

		p MASS (MeV)			
M1	938.3	0.5	BAMBERGER	70	CNTR
M1	938.179	0.058	HU	75	CNTR
M1	938.229	0.049	ROBERSON	77	CNTR
M1	938.30	0.13	ROBERTS	78	CNTR
M1	AVG	938.216	0.036	AVERAGE	

**NOTE ON PROTON MEAN LIFE LIMITS**

(by M. Goldhaber, Brookhaven National Laboratory, and F. Reines, University of California, Irvine)

Current ideas on the unification of the weak, electromagnetic, and strong forces suggest that baryon number might not be strictly conserved, so that the proton could decay. In the Particle Properties Summary Tables there are nearly thirty particles listed with a mass smaller than that of the proton (if we count both particles and antiparticles and different members of multiplets separately). Ten of these particles are fermions and the remainder bosons. There are then a great many possible two-body decay modes of the proton and an even larger number of three-body, etc., decay modes which satisfy charge, energy, momentum, and angular momentum conservation. Each decay mode has to contain at least one fermion to satisfy angular momentum conservation.

The "decay signature" distributions as well as the backgrounds depend on detector characteristics (the material from which the detector is made, the method of detection, timing information, time resolution, etc.). The background, due chiefly to atmospheric neutrinos, depends also on the geomagnetic latitude and on the phase of the solar cycle with which the magnetic field of the sun is associated. The depth-dependent cosmic ray background is due to cosmic ray muons and their progeny. For each possible proton decay signature there is a finite probability of a background event with a similar signature, where the probability depends on the detector characteristics.

The simplest grand unified theory, minimal SU(5), predicts  $e^+ \pi^0$  to be the predominant proton decay mode. The IMB lower limit on the partial mean life for this mode,  $2 \times 10^{32}$  years, is at least a factor of 10 higher than predicted by minimal SU(5) theory.

See also the reviews by Reines and Schultz, Surveys in High Energy Physics **1**, 89 (1980); Goldhaber, Langacker, and Slansky, Science **210**, 851 (1980); Goldhaber and Sulak, Comments Nucl. Part. Phys. **10**, 215 (1981); D.H. Perkins, Ann. Rev. Nucl. Part. Sci. **34**, 1 (1984); and J.M. LoSecco, Comments Nucl. Part. Phys. **15**, 23 (1985).

**NUCLEONS**

**p**

$$I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$$

**p MASS (MeV)**

M	(938.256)	(0.005)	COHEN	65	RVE
M	(938.272)	(0.0052)	TAYLOR	79	RVE
M	938.2796	0.0027	COHEN	73	RVE

USING NEW E/H



# Stable Particle Full Listings

p, n

<b>p → μ<sup>+</sup>μ<sup>+</sup>μ<sup>-</sup> (10<sup>30</sup> years)</b>				(P17)
T17	1(.7)	> 44.	CL=.90 BLEWITT 85 CNTR	P (FREE PROTON)
T17	1(.9)	>190.	CL=.90 BLEWITT 85 CNTR	P
T17 B	1	> 2.1	CL=.90 BATTISTON 82 CNTR	P
T17 B	WE HAVE CONVERTED 1 POSS EVENT TO .90 CL LIMIT.			
<b>N → ν K*(892) (10<sup>30</sup> years)</b>				(P18)
T18	10(16)	>5.8	CL=.90 BLEWITT 85 CNTR	P (FREE PROTON)
T18	7(6)	>9.6	CL=.90 BLEWITT 85 CNTR	P
T18	1(4)	>.7.	CL=.90 PARK 85 CNTR	N
T18 B	1	>2.1	CL=.90 BATTISTON 82 CNTR	P
T18 B	WE HAVE CONVERTED 1 POSS EVENT TO .90 CL LIMIT.			
<b>N → e<sup>+</sup>π<sup>0</sup> anything (10<sup>30</sup> years)</b>				(P19)
T19	0	>0.6	CL=.90 LEARNED 79 CNTR	P,N
<b>N → ν anything (years)</b>				(P20)
T20 L	0	>2.E26	CL=.90 LEARNED 79 CNTR	P,N
T20 L	ANYTHING = PI, RHO, K, ETC.			
<b>n → 3ν (years)</b>				(P21)
T21	0	>5.E26	CL=.90 LEARNED 79 CNTR	N
<b>p → e<sup>+</sup>η (10<sup>30</sup> years)</b>				(P22)
T22	5(6.5)	> 64.	CL=.90 BLEWITT 85 CNTR	P (FREE PROTON)
T22	5(4.7)	>200.	CL=.90 BLEWITT 85 CNTR	P
T22 C	2	> 1.2	CL=.90 CHERRY 81 CNTR	P
T22 C	WE HAVE CONVERTED 2 POSS EVENTS TO .90 CL LIMIT.			
<b>p → μ<sup>+</sup>η (10<sup>30</sup> years)</b>				(P23)
T23	6(6)	>17.	CL=.90 BLEWITT 85 CNTR	P (FREE PROTON)
T23	7(8)	>46.	CL=.90 BLEWITT 85 CNTR	P
<b>n → νη (10<sup>30</sup> years)</b>				(P24)
T24	4(3)	>18.	CL=.90 PARK 85 CNTR	N
T24 C	2	> 0.6	CL=.90 CHERRY 81 CNTR	N
T24 C	WE HAVE CONVERTED 2 POSS EVENTS TO .90 CL LIMIT.			
<b>N → νρ (10<sup>30</sup> years)</b>				(P25)
T25	6(7)	> 4.1	CL=.90 BLEWITT 85 CNTR	P (FREE PROTON)
T25	6(5)	> 8.4	CL=.90 BLEWITT 85 CNTR	P
T25	7(3)	> 2.	CL=.90 PARK 85 CNTR	N
T25 C	2	> 0.9	CL=.90 CHERRY 81 CNTR	P
T25 C	2	> 0.6	CL=.90 CHERRY 81 CNTR	N
T25 C	WE HAVE CONVERTED 2 POSS EVENTS TO .90 CL LIMIT.			
<b>n → νω (10<sup>30</sup> years)</b>				(P26)
T26	1(2)	>16.	CL=.90 PARK 85 CNTR	N
T26 C	2	> 2.0	CL=.90 CHERRY 81 CNTR	N
T26 C	WE HAVE CONVERTED 2 POSS EVENTS TO .90 CL LIMIT.			
<b>n → e<sup>+</sup>e<sup>-</sup>ν (10<sup>30</sup> years)</b>				(P27)
T27	4(3)	>26.	CL=.90 PARK 85 CNTR	N
<b>n → μ<sup>+</sup>μ<sup>-</sup>ν (10<sup>30</sup> years)</b>				(P28)
T28	4(7)	>19.	CL=.90 PARK 85 CNTR	N
<b>n → e<sup>-</sup>π (10<sup>30</sup> years)</b>				(P29)
T29	2(4)	>25.	CL=.90 PARK 85 CNTR	N
<b>n → μ<sup>-</sup>π (10<sup>30</sup> years)</b>				(P30)
T30	2(3)	>27.	CL=.90 PARK 85 CNTR	N
<b>n → e<sup>-</sup>ρ (10<sup>30</sup> years)</b>				(P31)
T31	5(3)	>12.	CL=.90 PARK 85 CNTR	N
<b>n → μ<sup>-</sup>ρ (10<sup>30</sup> years)</b>				(P32)
T32	2(2)	>.9.	CL=.90 PARK 85 CNTR	N
<b>N → 2 BODIES, ν-FREE (10<sup>30</sup> years)</b>				(P6+P7+P8+P9)
T33	0	>1.3	CL=.90 ALEXSEEV 81 CNTR	P,N

**p MAGNETIC MOMENT (e/2m<sub>p</sub>)**

MM	(2.79276) (0.00002)	COHEN 65 RVUE
MM	(2.792782 0.000017)	TAYLOR 69 RVUE
MM A	2.7928456 .0000011	COHEN 73 RVUE USING NEW E/H
MM A	COHEN 73 HAVE MADE A CORRECTION (MULTIPLYING BY 1.0000256) OF THE	
MM A	ORIGINAL DATA FOR DIAMAGNETIC SCREENING OF THE PROTON IN THE WATER	
MM A	MOLECULE. MAMYRIN 83 REPORT THAT ACCOUNTING FOR NEW ATOMIC MASS	
MM A	VALUES (AND OTHER CORRECTIONS) WILL REDUCE THE QUOTED VALUE BY	
MM A	1.000 000537 (WHICH WOULD GIVE 2.7928446 AFTER THE WATER	
MM A	CORRECTION).	

**n MAGNETIC MOMENT (e/2m<sub>p</sub>)**

MM1 O	(-1.8)	(1.2)	BUTTON 62 CNTR
MM1 R	(-2.83)	(0.10)	FOX 72 CNTR
MM1 R	(-2.819)	(0.056)	ROBERTS 74 CNTR
MM1	-2.791	0.021	HU 75 CNTR
MM1 R	-2.817	0.048	ROBERTS 78 CNTR
MM1 O	OLD EXPERIMENT WITH LARGE ERROR. NOT AVERAGED.		
MM1 R	ROBERTS 74 IS REANALYSIS OF FOX 72 DATA. REPLACES OLD FOX VALUE.		
MM1 R	ROBERTS 78 IS A REANALYSIS OF ROBERTS 74.		
MM1			
MM1 AVG	-2.795	0.019	AVERAGE

**p ELECTRIC DIPOLE MOMENT (units 10<sup>-23</sup> e-cm)**

FORBIDDEN BY BOTH T INVARIANCE AND P INVARIANCE

EDM	1G (700.) (900.)	HARRISON 69 MSR
EDM	(5000.) OR LESS	KHRIPLOVI 76
EDM A	(130.) (200.)	WILKENING 84
EDM A	400. 1400.	WILKENING 84
EDM	(400.) OR LESS	DZUBA 85 THEO USE EDM OF 129XE
EDM A	FIRST WILKENING 84 NUMBER INCLUDES A FINITE SIZE EFFECT AND A	
EDM A	MAGNETIC EFFECT; SECOND NUMBER IS MORE CAUTIOUS AND EXCLUDES FINITE	
EDM A	SIZE EFFECT WHICH RELIES ON UNCERTAIN NUCLEAR INTEGRALS.	

**p - e CHARGE DIFFERENCE (units e)**

SEE ALSO SECTION Q IN THE FULL LISTINGS FOR NEUTRONS BELOW.

DQ	D	1.0E-21 OR LESS	DYLLA 73	NEUTRALITY OF SF6
DQ	D	(0.8E-21) OR LESS	MARINELLI 84 RVUE	
DQ	D	ASSUMES THAT Q(NEUTRON)=Q(PROTON)-Q(E-). SEE DYLLA 73 FOR A		
DQ	D	SUMMARY OF EXPERIMENTS ON THE NEUTRALITY OF MATTER.		

**REFERENCES FOR p**

BLEWITT 85 PRL 55 2114 + (UCI-MICH+BNL+CIT+CLEV+HAWA+LOUC+WARS)  
 DZUBA 85 PL 1548 53 DZUBA, FLAMBAUM, SILVESTROV (NOVO)  
 PARK 85 PRL 54 22 - (UCI-MICH+BNL+CIT+CLEV+HAWA+LOUC+WARS)  
 MARINELLI 84 PL 137B 439 MARINELLI, MORPURGO (GENO)  
 WILKENING 84 PR A29 425 WILKENING, RAMSEY, LARSON (HARV+VIRG)  
 BARTLETT 83 PRL 50 651 +COURANT, HELLER, JOYE, MARSHAK. (MINN-ANL)  
 BATTISTON 83 PL 133B 454 BATTISTON, BELLOTTI+ (FRAS-MILA+LCGT-CERN)  
 MAMRYN 83 JETP 57 1152 +ARUEV, ALEKSEENKO (IOFF)  
 BATTISTON 82 PL 118B 461 BATTISTON, BELLOTTI+ (FRAS-MILA+LCGT-CERN)  
 KRISHNAS 82 PL 115B 349 KRISHNASHAMY, MENON, MORDAL. (TIFR+OSKC+TOKY)  
 ALEXSEEV 81 JETPL 33 651 +BAKATANOV, BUTKEVICH, VOEVODSKII - (LENI)  
 CHERRY 81 PRL 47 1507 +DEAKYNE, LANDE, LEE, STEINBERG + (PENN-BNL)  
 COWSIK 80 PR D22 2204 R. COWSIK, V. S. NARASIMHAN (TIFR)  
 BELL 79 PL 86B 215 +CALVETTI, CARRON, CHANEY, CITTOLIN+ (CERN)  
 GOLDEN 79 PRL 43 1196 +HORAN, MAUGER, BADHMAR, LACY+ (NASA-PSL)  
 LEARNED 79 PRL 43 907 +REINES, SONI (UCI)  
 BREGMAN 78 PL 78B 174 +CALVETTI, CARRON, CITTOLIN, HAUER, HERR+ (CERN)  
 GANGULI 78 PL 74B 130 +MALHOTRA, RACHAVAN, SUBRAMANIAN + (TIFR)  
 ROBERTS 78 PR D17 358 B. L. ROBERTS (WILL+RHEL)  
 EVANS 77 SCIENCE 197 989 +STEINBERG (BNL+PENN)  
 ROBERSON 77 PR C16 1945 -KING, KUNSELMAN+ (NYOM+CIT+CARN+VPI+WILL)  
 KHRIPLOV 76 JETP 44 25 I. S. KHRIPLOVICH (NIC. PHYS. INST., SIBERIA)  
 HU 75 NP A254 403 +ASANO, CHEN, CHENG, DUGAN+ (COLU+YALE)  
 ROBERTS 74 PRL 33 1181 +COX, ECKHAUSE+ (WILL+VPI+CARN+NYOM+CIT-BNL)  
 COHEN 75 PR D12 1232 ROBERTS, COX + (WILL+VPI+CARN+NYOM+CIT-BNL)  
 DYLLA 73 PR A7 1224 H. F. DYLLA, J. G. KING (MIT)  
 FOX 72 PRL 29 193 +BARNES, EISENSTEIN+ (BNL+CARN+VPI+WILL+NYOM)  
 BAMBERGE 70 PL 35B 233 BAMBERGER, LYNN, PIEKARZ+ (MPIH+CERN+KARL)  
 DIX 70 THESIS CASE F. E. DIX (CASE)  
 HARRISON 69 PRL 22 1263 HARRISON, SANDARS, WRIGHT (CLARENDON OXFORD)  
 TAYLOR 69 RMP 41 375 +PARKER, LANGENBERG (PRIN+UCI+PENN)  
 GURR 67 PR 15B 1521 GURR, KRÖPP, REINES, MEYER (CASE, JONAHNBURG)  
 COHEN 65 RMP 37 537 +DUMOND (W. AMER. AVIATION SCIENCE CENT., CIT)  
 BUTTON 62 PR 127 1297 J. BUTTON, B. MAGLIC (LBL)  
 FLEROV 58 SOV PHYS DOK 3 79 FLEROV, KLOCHKOV, SKOBKIN, TERENTEV (USSR)

QUANTUM NUMBER DETERMINATIONS NOT REFERRED TO IN THE DATA LISTINGS  
 KALOGERO 76 PRL 37 1037 KALOGEROPOULOS, CHIU, SUDARSHAN (SYRA-TEXA)P  
 FRANKLIN 77 PR D16 910 JERROLD FRANKLIN (HAIF)P

$$\bar{n} \quad I(J^P) = \frac{1}{2} \left( \frac{1}{2} \right)^4$$

**n MASS (MeV)**

M T	(939.5527) (0.0052)	TAYLOR 69 RVUE	USING NEW E/H
M T	939.5731 0.0027	COHEN 73 RVUE	
M T	(939.5730) (0.0026)	GREENWOOD 80 CNTR	N P --> D GAMMA
M T	THESE DETERMINATIONS OF NEUTRON MASS NOT INDEPENDENT OF		
M T	NEUTRON-PROTON MASS DIFFERENCE MEASUREMENTS BELOW.		

**n - p MASS DIFFERENCE (MeV)**

D M	(1.29344) (0.00007)	MATTAUCH 65 RVUE
D M	(1.293429) (0.000036)	COHEN 73 RVUE
D M	1.293330 0.000040	VYLOV 78 CNTR
D M	1.293322 0.000017	GREENWOOD 80 CNTR
D M	WE HAVE CONVERTED MEASURED NEUTRON-HYDROGEN MASS DIFFERENCE TO	
D M	NEUTRON-PROTON MASS DIFFERENCE USING CURRENT VALUE OF ELECTRON MASS	
D M	AND A HYDROGEN BINDING ENERGY OF 13.6 EV.	
D		
D AVG	1.293323 0.000016	AVERAGE

For notation, see key on page 91.

# Stable Particle Full Listings

n

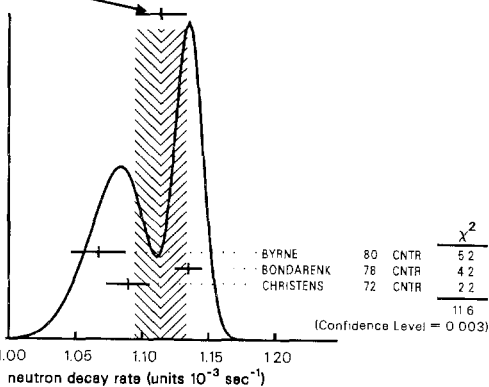
## n MEAN LIFE (units 10<sup>3</sup> sec)

THE ORIGIN OF THE DISCREPANCIES BETWEEN MEASUREMENTS OF THE NEUTRON MEAN LIFE IS NOT KNOWN. WE AVERAGE THE HIGHEST PRECISION MODERN EXPERIMENTS AND INCLUDE THE RESULTING SCALE FACTOR IN THE ERROR, AS USUAL. THE IDEOGRAM SHOWS THE INVERSE MEAN LIFE, THE QUANTITY ACTUALLY USED IN CALCULATING OUR AVERAGE VALUE.

LIFETIMES FOR SOUND NEUTRONS ARE GIVEN IN THE PROTON MEAN LIVES SECTION

T	(1.013)	(0.026)	SOSNOVSKI	59 CNTR	
T	(0.935)	(0.014)	CHRISTENS	67 CNTR REPL BY CHRISTENS72	
T	0.918	0.014	CHRISTENS	72 CNTR	
T	0.881	0.008	BONDARENK	78 CNTR	
T	0.937	0.018	BYRNE	80 CNTR	
T	(0.875)	(0.095)	KOSVINTSE	80 CNTR	
T	(0.902)	(0.010)	WILKINSON	82 RVUE	INFERRED VALUE
T	A INCLUDES CORRECTION FOR RECOIL PROTON SCATTERING SEE BONDARENKO 82.				
T	B WILKINSON 82 VALUE INFERRED FROM N DECAY CORRELATIONS.				
T	AVG	0.898	0.016	0.016	AVERAGE (ERROR INCL. SCALE FACTOR OF 2.4) (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
1.114 ± 0.019 (ERROR SCALED BY 2.4)



## n MAGNETIC MOMENT (NUCLEAR MAGNETONS, 938.2 MeV)

MM	(-1.913148)	0.0000066	COHEN	56 RVUE	
MM	(-1.91304211)	0.0000088	GREENE	77 MRS REPL. BY GREENE 79	
MM	(-1.91304184)	0.0000088	GREENE	79 MRS	
MM	-1.91304308	0.0000054	GREENE	82 MRS	REPLACES GREENE 79
MM	A GREENE 82 MEASURES THE NEUTRON MAGNETIC MOMENT IN BOHR MAGNETONS				
MM	A TO BE 1.041 875 64E-31-D.000 000 26E-3. THE VALUE QUOTED ABOVE IS				
MM	A OBTAINED BY MULTIPLYING BY M(PROTON)/M(ELECTRON)=1836.153000.				
MM	A A MODIFIED VALUE WOULD BE OBTAINED USING A NEW UNPUBLISHED VALUE OF				
MM	A 1836.152701				

## n ELECTRIC DIPOLE MOMENT (units 10<sup>-23</sup> e-cm)

FORBIDDEN BY BOTH T INVARIANCE AND P INVARIANCE

EDM	M	(-20.)	(30.)	MILLER	67 MRS	
EDM	M	+24.	39.	SHULL	67 CNTR	
EDM	M	(30.)	OR LESS	DRESS	68 MRS	ABSOLUTE VALUE
EDM	M	(5.)	OR LESS	BAIRD	69 MRS	INCLUDED IN DRESS73
EDM	M	- 2.	39.	APOSTOLES	70 MRS	
EDM	M	0.32	0.75	DRESS	73 MRS	< 10** -23 (CL=.80)
EDM	M	0.04	0.15	DRESS	77 MRS	< 3 E-24 (CL=.90)
EDM	A	0.040	0.075	ALTAREV	79 MRS	< 1.6E-24 (CL=.90)
EDM	A	0.021	0.024	ALTAREV	81 MRS	< 6 E-25 (CL=.90)
EDM	M	DRESS 68 INCLUDES DATA OF MILLER 67.				
EDM	A	ALTAREV 79 AND 81 USE ULTRACOLD NEUTRONS.				
EDM	AVG	0.023	0.023	AVERAGE		

## n CHARGE

SEE ALSO SECTION DQ IN THE FULL LISTINGS FOR PROTONS ABOVE

Q	NEUTRON CHARGE (UNITS 10** -20 E)			
Q	(-1.5)	(2.2)	CL=.90	GAHLER 82 CNTR REACTOR N BEAM

## LIMIT ON n- $\bar{n}$ OSCILLATIONS

NAN MEAN TIME FOR N-ANTI N TRANSITION IN VACUUM (UNITS SEC.)  
 NAN TEST OF BARYON CONSERVATION.  
 NAN LIMITS ARE DERIVED FROM EXPERIMENTAL LIMITS ON DELTA-B = 2 NUCLEAR  
 NAN DECAY PROCESSES, USING THEORETICAL ASSUMPTIONS FOR NUCLEAR PHYSICS  
 NAN EFFECTS. SEE ALSO THE THEORETICAL ANALYSIS OF DOVER, GAL, AND  
 NAN RICHARDS, PR D27, 1090 (1983). SEE THE REVIEWS OF H.L. ANDERSON,  
 NAN PROC. LASL CONF. ON NUCLEAR PARTICLE PHYSICS AT ENERGIES UP TO 31  
 NAN GEV, 1981.

NAN	(1.E8)	OR MORE	CL=.90	CHETYRKIN	81 THEO	
NAN	(5.E7)	OR MORE	CL=.90	COWSIK	81 THEO	
NAN	(3.E7)	OR MORE		ALBERICO	82 THEO	
NAN	(2.E7)	OR MORE		CHERRY	83 CNTR-THY	
NAN	(3.E7)	OR MORE		BATTISTON	84 CNTR	
NAN	(2.7E7)	TO 1.E8	OR MORE	JONES	84 CNTR+THY	
NAN	(1.E6)	OR MORE	CL=.90	FIDECARO	85 CNTR	FREE REACTOR N
NAN	8.E7	OR MORE	CL=.90	PARK	85 CNTR-THY	

## n PARTIAL DECAY MODES

P1	n	→ p e <sup>-</sup> $\bar{\nu}_e$	DECAY MASSES	
P2	n	→ p $\nu_e$ $\bar{\nu}_e$	DECAY MASSES	
			938.-.511+ 0	
			938. 0- 0	

## n BRANCHING RATIOS

n	→ (p $\nu_e$ $\bar{\nu}_e$ )/(p e <sup>-</sup> $\bar{\nu}_e$ )	(P2)/(P1)
R1	FORBIDDEN BY CHARGE CONSERVATION	
R1	(3. E-19) OR LESS	NORMAN 79 CNTR RB87-->SR87M+NEUTRL
R1	9. E-24 OR LESS	BARABANOV 80 CNTR GA71-->GE71 + ANY
R1	(7.9E-21) OR LESS	VAIDYA 83 CNTR RB87-->SR87M+NEUTRL
R1	(9.7E-18) OR LESS	CL=.90 ROY 83 CNTR
R1	R CD13 --> IN115M + NEUTRALS	

## NOTE ON BARYON DECAY PARAMETERS

### A/V ratio for baryon leptonic decays

Consider the decay

$$B_i \rightarrow B_f + \ell + \nu$$

Assuming V, A theory, neglecting "induced" scalar, "induced" pseudoscalar, and axial weak-magnetism terms, and neglecting the q<sup>2</sup> dependence of the form factors, the baryon part of the matrix element for these decays may be written as<sup>1</sup>

$$\bar{B}_f [\gamma_\lambda (g_V - g_A \gamma_5) + (g_W/m_{B_i}) \sigma^{\lambda\nu} q_\nu] B_i$$

Here B<sub>i</sub> and  $\bar{B}_f$  are spinors which represent initial and final baryons, g<sub>A</sub> and g<sub>V</sub> are the axial and vector coupling constants, g<sub>W</sub> is the weak magnetism coupling constant, and q<sub>ν</sub> is the sum of the lepton momenta.

The Pauli representation is used for the γ matrices. The ratio g<sub>A</sub>/g<sub>V</sub> may be written as

$$g_A/g_V = |g_A/g_V| \exp(i\phi),$$

where φ is 0 plus nπ if time reversal holds.<sup>2</sup>

Experiments on the leptonic decays of baryons other than the neutron have generally assumed φ to be either 0 or π, and have thus measured the magnitude and sign of g<sub>A</sub>/g<sub>V</sub>. In studying neutron beta decay, however, experiments have been sensitive enough to measure φ more precisely, and we include the phase angle in our Listings for this case. It is consistent with time-reversal invariance, and by using the above definition of the matrix element with the Pauli representations, the value of g<sub>A</sub>/g<sub>V</sub> in neutron beta decay is negative.



## Stable Particle Full Listings

*n*

Due to statistical limitations, the weak magnetism form factor  $g_W$  is usually assumed from CVC and SU(3), so that usually only  $g_A$  and  $g_V$  are determined experimentally. This determination is accomplished in a variety of ways:

(a) The lepton-neutrino angular correlation provides a measure of the absolute value of  $g_A/g_V$  (for relevant formulas, see, e.g., Albright<sup>3</sup>).

(b) The up-down asymmetry of the lepton from polarized baryon decays provides a measure of  $g_A/g_V$  with its sign (for relevant formulas, see, e.g., Albright<sup>3</sup>).

(c) The lepton spectrum, given enough statistics, provides a measure of  $g_A/g_V$  with its sign (for relevant formulas, see, e.g., Bender et al.<sup>4</sup>). The lepton spectrum also provides a measure of  $g_W/g_A$  if the CVC-SU(3) assumption is relaxed.

(d) The polarization of the decay baryon, from polarized or unpolarized initial baryon, also provides  $g_A/g_V$  with its sign (for formulas, see, e.g., Willis and Thompson<sup>5</sup>).

(e) The presence of a triple correlation term proportional to

$$\sigma_{B_i} \cdot (\mathbf{p}_e \times \mathbf{p}_\nu),$$

where the initial baryon is polarized or

$$\sigma_{B_f} \cdot (\mathbf{p}_e \times \mathbf{p}_\nu),$$

where the polarization of the decay baryon is observed provides a measure of the deviation of  $\phi$  from 0 or  $\pi$ , and is thus a test of time-reversal invariance (see, e.g., Willis and Thompson<sup>5</sup>).

We compile the ratio  $g_A/g_V$  with its sign, for those decays for which it has been measured.

All the coupling constants and decay rates for baryon leptonic decays are related by Cabibbo's theory,<sup>6</sup> extended to six quarks (and three mixing angles) by Kobayashi and Maskawa.<sup>7</sup> A discussion of the Kobayashi-Maskawa mixing matrix is given in the Miscellaneous Section of this Review.

### Asymmetry parameters in nonleptonic hyperon decays

The transition matrix for hyperon decay may be written as

$$M = s + p(\boldsymbol{\sigma} \cdot \mathbf{q}), \quad (1)$$

where  $s$  and  $p$  are the parity-changing and the parity-conserving amplitudes, respectively;  $\boldsymbol{\sigma}$  is the Pauli spin operator, and  $\mathbf{q}$  is a unit vector along the direction of the decay baryon in the hyperon rest frame.

The asymmetry parameters are defined by the relations

$$\alpha = 2 \operatorname{Re}(s^*p)/(|s|^2 + |p|^2),$$

$$\beta = 2 \operatorname{Im}(s^*p)/(|s|^2 + |p|^2),$$

$$\gamma = (|s|^2 - |p|^2)/(|s|^2 + |p|^2).$$

With the transition matrix  $M$  given by Eq. (1) above, the angular distribution of the decay baryon, in the hyperon rest system, is of the form

$$I = 1 + \alpha \mathbf{P}_Y \cdot \mathbf{q},$$

where  $\mathbf{P}_Y = \langle Y | \boldsymbol{\sigma} | Y \rangle$  is the hyperon polarization. In the notation of Lee and Yang,<sup>8</sup> the polarization  $P_B$  of the decay baryons is

$$\mathbf{P}_B = \frac{(\alpha + \mathbf{P}_Y \cdot \mathbf{q})\mathbf{q} + \beta(\mathbf{P}_Y \times \mathbf{q}) + \gamma\mathbf{q} \times (\mathbf{P}_Y \times \mathbf{q})}{1 + \alpha \mathbf{P}_Y \cdot \mathbf{q}},$$

where  $\mathbf{P}_B$  is defined in that rest system of the baryon obtained by a Lorentz transformation along  $\mathbf{q}$  from the hyperon rest system in which  $\mathbf{q}$  and  $\mathbf{P}_Y$  are defined. Note that  $\alpha$  is the helicity of the decay baryon for unpolarized hyperons.

The three parameters  $\alpha$ ,  $\beta$ , and  $\gamma$  satisfy the relation

$$\alpha^2 + \beta^2 + \gamma^2 = 1.$$

It is then convenient to describe hyperon nonleptonic decays in terms of the two independent parameters  $\alpha$  and the angle  $\phi$  defined by

$$\beta = (1 - \alpha^2)^{1/2} \sin \phi,$$

$$\gamma = (1 - \alpha^2)^{1/2} \cos \phi,$$

which has a more nearly Gaussian distribution of measurement error than  $\beta$  or  $\gamma$ . Evidently

$$-\frac{1}{2}\pi \leq \phi \leq \frac{1}{2}\pi \quad \text{for } \gamma > 0,$$

$$+\frac{1}{2}\pi \leq \phi \leq \frac{3}{2}\pi \quad \text{for } \gamma < 0.$$

In discussing time-reversal invariance, the quantity of interest is  $\Delta$ , defined by

$$\alpha = 2|s||p| \cos \Delta / (|s|^2 + |p|^2),$$

$$\beta = -2|s||p| \sin \Delta / (|s|^2 + |p|^2);$$

that is,  $\Delta$  is the phase angle of  $s$  relative to  $p$ . Evidently

$$-\frac{1}{2}\pi \leq \Delta \leq \frac{1}{2}\pi \quad \text{for } \alpha > 0,$$

$$+\frac{1}{2}\pi \leq \Delta \leq \frac{3}{2}\pi \quad \text{for } \alpha < 0.$$

For notation, see key on page 91.

# Stable Particle Full Listings

Under the assumption of time-reversal invariance, the angle  $\Delta$  must satisfy the relation

$$\Delta = \delta_s - \delta_p,$$

modulo  $\pi$ , where  $\delta_s$  and  $\delta_p$  are the pion-baryon scattering phase shifts at the appropriate energy and for the appropriate isospin state. For  $\Lambda$  decay, assuming the validity of the  $|\Delta I| = 1/2$  rule,

$$\Delta = \delta_s - \delta_p = (7.0 \pm 1.0) \text{ deg.}^9$$

In the Stable Particle Full Listings we give  $\alpha$  and  $\phi$  for each decay since they are the most closely related to the experiments and are essentially uncorrelated. Whenever necessary we have changed the signs of the reported values, so as to agree with our conventions. In the Stable Particle Summary Table we give  $\alpha$ ,  $\phi$ , and  $\Delta$  with errors; and for convenience we also give the central value of  $\gamma$ , without an error.

### References

- M.L. Goldberger and S.B. Treiman, Phys. Rev. **11**, 354 (1958).
- J.D. Jackson, S.B. Treiman, and H.W. Wyld Jr., Phys. Rev. **106**, 517 (1957).
- C.H. Albright, Phys. Rev. **115**, 750 (1959).
- I. Bender, V. Linke, and H.J. Rothe, Z. Physik **212**, 190 (1968).
- W. Willis and J. Thompson, in *Advances in Particle Physics*, eds. R.L. Cool and R.E. Marshak (Wiley, New York), Vol. 1, p. 295 (1968).
- N. Cabibbo, Phys. Rev. Lett. **10**, 531 (1963).
- M. Kobayashi and T. Maskawa, Progr. Theor. Phys. **49**, 652 (1973).
- T.D. Lee and C.N. Yang, Phys. Rev. **108**, 1615 (1957). Note that this paper contains a misprint. The minus sign in the definition of  $\beta$  should be replaced by a 2. In addition, our unit vector  $\mathbf{q}$  is the direction of the baryon, whereas their unit vector  $\mathbf{p}$  is the direction of the pion.
- This value for  $\delta_s - \delta_p$  is derived from the phase-shift analyses by R. Ayed, CEA-N-192, Saclay thesis (1976). The error is our estimation of the uncertainty allowing for possible correlations.

### DECAY PARAMETERS FOR $n \rightarrow p e^- \bar{\nu}$

$g_A/g_V$	(see note above for sign convention)		
AV C	(-1.250) (0.044)	CONFORTO 67 RVUE	SEE NOTE C BELOW
AV EP	(-1.25) (0.01)	CHRISTENS 67 CNTR	N DECAY FT VALUE
AV P	(-1.22) (0.08)	GRIGOREV 68 CNTR	E-MU ANG CORREL
AV P	(-1.26) (0.02)	CHRISTENS 70 CNTR	PE, NEUT SPIN CORREL
AV EP	(-1.27) (0.025)	EROZOLIMS 71 CNTR	REPL. BY EROZOLIMS79
AV EP	(-1.259) (0.011)	CHRISTENS 72 CNTR	N DEC. + FT VALUE
AV P	(-1.263) (0.016)	KROPP 74 RVUE	N DECAY ALONE
AV P	-1.250 0.009	KROPP 74 RVUE	N DEC. + FT VALUE
AV EP	(-1.250) (0.036)	DOBROZEM 75 CNTR	REPL. BY STRATOWA 78
AV K	-1.253 0.021	KROHN 75 CNTR	PE, NEUT SPIN CORREL
AV	(-1.263) (0.015)	EROZOLIMS 77 CNTR	REPL. BY EROZOLIMS79
AV E	-1.259 0.017	STRATOWA 78 CNTR	PROTON RECOL SPECT
AV E	-1.261 0.012	EROZOLIMS 79 CNTR	PE, NEUT SPIN CORREL
AV	-1.226 0.042	MOSTOVOY 83 RVUE	

AV C CON-ORTO 67 COMBINES FREE NEUTRON DATA TO 1967. REPL. BY KROPP 74. THESE EXPERIMENTS MEASURE THE ABSOLUTE VALUE OF  $g_A/g_V$  ONLY.

AV E KROPP 74 VALUE OBTAINED BY FITTING ALL DATA THROUGH 1972.

AV K KROHN 75 PAPER GIVES -1.258--0.015 INCLUDING EVENTS OF CHRISTENS 70. AV K THE VALUE QUOTED ABOVE IS DERIVED FROM HIS A, BASED ON NEW EXPT ONLY

AV AVG 1.2539 0.0063 AVERAGE

PHASE ANGLE OF $g_A$ RELATIVE TO $g_V$ (degrees)			
F	C	TIME REVERSAL INVARIANCE WOULD REQUIRE THIS TO BE 0 OR 180 DEGREES	
F	P	(175.3) (10.)	BURGY 60 CNTR POLAR. NEUTRONS
F	P	(198.3) (27.)	CLARK 60 CNTR POLAR. NEUTRONS
F	C	(176.1) (6.4)	CONFORTO 67 RVUE
F	P	(181.3) (1.3)	EROZOLIMS 70 CNTR POLAR. NEUTRON
F	P	180.35 0.43	EROZOLIMS 74 CNTR POLAR. NEUTRONS
F	P	181.1 1.3	KROPP 74 RVUE N DECAY
F	P	180.14 0.22	STEINBERG 74 CNTR POLAR. NEUTRONS
F	P	179.71 0.39	EROZOLIMS 78 CNTR POLAR. NEUTRONS
F	C	CONFORTO 67 COMBINES FREE NEUTRON DATA TO 1967. REPL. BY KROPP 74.	
F	P	KROPP 74 VALUE OBTAINED BY FITTING ALL DATA THROUGH 1972.	
F	AVG	180.11 0.17	AVERAGE

### TRIPLE CORRELATION COEFFICIENT

D1 MEASURES COMPONENT OF NEUTRON SPIN PERPENDICULAR TO THE DECAY PLANE IN BETA DECAY. SHOULD BE ZERO IF T-INVARIANCE NOT VIOLATED. SEE NOTE ON BARYON DECAY PARAMETERS ABOVE.			
D1		-0.01 0.01	EROZOLIMS 70 CNTR POLAR. NEUTRONS
D1	E	0.0027 0.0050	EROZOLIMS 74 CNTR POLAR. NEUTRONS
D1	E	-0.0011 0.0017	STEINBERG 74 CNTR POLAR. NEUTRONS
D1	E	+0.0022 0.0030	EROZOLIMS 78 CNTR POLAR. NEUTRONS
D1	E	EROZOLIMS78 SAYS ASYMMETRIC PROTON LOSSES AND NON-UNIFORM BEAM	
D1	E	POLARIZATION MAY GIVE SYSTEMATIC ERROR UP TO 0.005, THUS INCREASING	
D1	E	THE EROZOLIMS78 74 ERROR TO 0.005. STEINBERG 74 76 ESTIMATES	
D1	E	THESE SYSTEMATIC ERRORS TO BE INSIGNIFICANT IN THEIR EXPERIMENT.	
D1	AVG	-0.0007 0.0014	AVERAGE

### REFERENCES FOR $n$

COHEN 56 PR 104 283	V W COHEN, CORNGOLD, RAMSEY (BNL-HARVARD)
SOSNOVSK 59 JETP 9 717	SOSNOVSKII, SPIVAK, PROKOFEV + (IAE MOSCOW)
BUREY 60 PR 120 1829	+KROHN, MOVSEY, RINGO (ANL-CHIC)
CLARK 60 CJP 38 693	-ROBSON
MATTAUCH 65 NP 67 1	+THEIELE, MAPSTRA (MAX PLANCK INST. CHEM.)
CHRISTEN 67 PL 268 11	CHRISTENSEN, NIELSON, BARNSEN, BROWN+ (RISO)
CONFORTO 67 APAP 22 15	G. CONFORTO (CERN)
MILLER 67 PRL 19 381	+DRESS, BAIRD, RAMSEY (ORNL+HARV)
SHULL 67 PRL 19 384	C.G. SHULL, R. NATHANS (MIT-BNL)
DRESS 68 PR 170 1200	+BAIRD, MILLER, RAMSEY (ORNL+HARV)
GRIGOREV 68 SJNP 6 239	+GRISHIN, VLADIMIRSKII, NIKOLAEVSKII + (ITEP)
BAIRD 69 PR 179 1285	+MILLER, DRESS, RAMSEY (ORNL, HARV)
TAYLOR 69 RMP 41 375	+PARKER, LANGENBERG (PRIN-UCI+PENN)
APOSTOLE 70 RRP 15 343	APOSTOLESCU, IONESCU, IONESCU-BUJOR + (BUCH)
CHRISTEN 70 PR C1 1693	CHRISTENSEN, KROHN, RINGO (ANL)
EROZOLIM 70 SJNP 11 583	EROZOLIMSKY, BONDARENKO, + (KIAE)
ALSO PL 278 557	EROZOLIMSKY, BONDARENKO +
EROZOLIM 71 JETPL 13 252	EROZOLIMSKII, BONDARENKO + (KIAE)
CHRISTEN 72 PR D5 1628	CHRISTENSEN, NIELSON, BARNSEN, BROWN+ (RISO)
COHEN 73 J. PHYS. CHEM. REF. DATA 2, P. 663	E.R. COHEN, B.N. TAYLOR (ORNL+HARV)
DRESS 73 PR D7 3147	DRESS, MILLER, RAMSEY (ORNL+HARV)
KROPP 74 ZPHY 267 129	A KROPP, H PAUL (LINZ)
ALSO 70 NP A154 160	H PAUL (VIEN)
EROZOLIM 74 JETPL 20 345	EROZOLIMSKII, MOSTOVOI, FEDUNIN, FRANK+ (KIAE)
STEINBERG 74 PRL 33 41	STEINBERG, LIAU, VIGNON, HUGHES (YALE+GREEN)
ALSO 76 PR D25 2469	STEINBERG, LIAU, VIGNON, HUGHES (YALE+GREEN)
DOBROZEM 75 PR D11 510	DOBROZEMSKY, KERSCHBAUM, MORAW, PAUL (SEIB)
KROHN 75 PL 558 175	KROHN, RINGO (ANL)
DRESS 77 PR D15 9	+MILLER, PENDLEBURY, PERRIN+ (ORNL+GREEN+HARV)
EROZOLIM 77 JETPL 23 663	EROZOLIMSKII, FRANK, MOSTOVOI+ (KIAE)
GREENE 77 PL 718 297	+RAMSEY, MAMPE+ (HARV-ILLG+SUSS-ORNL-CENG)
EROZOLIM 78 SJNP 28 48	EROZOLIMSKII, MOSTOVOI, FEDUNIN, FRANK+ (KIAE)
STRATOWA 78 PR D18 3970	+DOBROZEMSKY, WEINZLER (SEIB)
BONDAREN 78 JETPL 28 303	BONDARENKO, KURGUZOV, PROKOFEV+ (KIAE)
ALSO 82 SMOLENICE CONF.	L.N. BONDARENKO (KIAE)
VILOY 78 SJNP 28 585	+GROMOV, IVANOV, OSIPENKO, FROLOV (JINR)
ALTAREV 79 JETPL 29 730	+BORISOV, BRANDIN, EGOROV, EZHOV, IVANOVA-(LENI)
EROZOLIM 79 SJNP 30 356	EROZOLIMSKII, FRANK, MOSTOVOI+ (KIAE)
GREENE 79 PR D20 2139	+RAMSEY, MAMPE+ (HARV-ILLG+SUSS+ORNL+CENG)
NORMAN 79 PRL 43 1226	E.B. NORMAN, A.G. SEAMSTER (WASH)
BARABANO 80 JETPL 32 359	BARABANOV, VERETENKIN, GAVRIN + (LENI)
BYRNE 80 PL 928 274	+MORSE, SMITH, SHAIKH, GREEN, GREENE (SUSS-RL)
GREENWOOD 80 PR C21 498	R.C. GREENWOOD, R.C. GREEN (INEL-BNL)
KOSVINTS 80 JETPL 31 236	KOSVINTSEV, KUSHNIR, MOROZOV, TEREKHOV (JINR)
ALTAREV 81 PL 1028 13	+BORISOV, BOROVIKOVA, BRANDIN, EGOROV + (LENI)
CHEZYRKI 81 PL 998 358	CHEZYRKIN, KAZARNOVSKY, KUZMIN- (INRM)
COWSIK 81 PL 1018 237	COWSIK, NUSSINOV (UMD)
ALBERICO 82 PL 1148 266	+BOTTINO, MOLINARI (CERN+TORI)
BAHLER 82 PR D25 2887	+KALUS, MAMPE (HARV-ILLG)
GREENE 82 METROLOGIA 18 93	+RAMSEY+ (YALE+HARV+ILLG+SUSS+ORNL+CENG)
WILKINSO 82 NP A377 474	D.H. WILKINSON (SUSS-BNL)
CHERRY 83 PRL 50 1354	+LAMDE, LEE, STEINBERG, CLEVELAND (PENN-BNL)
MOSTOVOY 83 JETPL 37 196	YU.A. MOSTOVOI (KIAE)
ROY 83 PR D28 1770	+VAIDYA, EPHRAIM, DATAR, BHATTI+ (TIFR)
VAIDYA 83 PR D27 486	+ROY, EPHRAIM, DATAR, BHATTACHERJEE (TIFR)
BATTISTO 84 PL 1338 454	BATTISTONI, BELLOTTI+ (FRAS+MILA+LDGT+CERN)
JONES 84 PRL 52 1191	IMB COLLAB (UCI-MICH-BNL-CIT+CLEV+HAWA+LOUC)
FIDECARO 85 PL 1568 122	+LANCERIT+ (CERN+ILLG+PADO+PAL+SUSS)
PARK 85 NP B252 261	IMB COLLAB (UCI-MICH-BNL-CIT+CLEV+HAWA+LOUC)

PAPERS NOT REFERRED TO IN DATA LISTINGS

JACKSON 57 PR 106 517	JACKSON, TREIMAN, WYLD (PRINCETON)
COHEN 65 RMP 37 537	+DUMOND (N. AMER. AVIATION SCIENCE CENT., CIT)
C P SHALL 64 PL 19 201	C P SHALL (ALABAMA)
BYRNE 82 RPP 45 115	D.I. BYRNE (SUSS)
FRANK 82 SPU 25 280	A.J. FRANK (KIAE)
WILKINSO 82 NP A377 474	WILKINSON (SUSS)

# Stable Particle Full Listings

A

## STRANGENESS -1 BARYONS

$\Lambda$

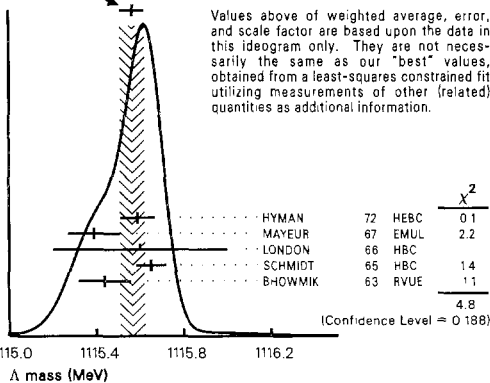
$$I(J^P) = 0(\frac{1}{2}^+)$$

### A MASS (MeV)

SINCE OUR FINAL VALUES FOR THE SIGMA AND LAMBDA MASSES COME FROM DOING AN OVERALL FIT TO ALL MEASURED MASSES AND MASS DIFFERENCES, WE HAVE USED THE UNCORRELATED MEASUREMENTS FROM SCHMIDT 65 RATHER THAN THE ONES COMING FROM THE OVERALL FIT REPORTED IN THAT PAPER. SINCE THERE SEEMS TO BE NO CONVINCING ARGUMENT AS TO WHY ONE SHOULD IGNORE DATA USING RANGE MEASUREMENTS, WE HAVE INCLUDED HERE VALUES DEPENDING ON PROTON AND PION RANGES. THE SCHMIDT 65 MASSES HAVE BEEN REEVALUATED USING OUR APRIL 1973 PROTON AND K<sup>-</sup> AND PI<sup>-</sup> MASSES. P. SCHMIDT, PRIVATE COMMUNICATION, (1974).

M	1115.44	0.12	BHOWMIK	63 RVUE +	SEE NOTE L BELOW
M	L				ABOVE LAMBDA MASS HAS BEEN RAISED 35 KEV TO ACCOUNT FOR 46 KEV INCREASE IN PROTON MASS AND 11 KEV DECREASE IN PI <sup>-</sup> MASS.
M	S	635(1115.86)	(0.09)	BALTAY	65 HBC ERROR IS STATIS.
M	488	1115.65	0.07	SCHMIDT	65 HBC SEE NOTE N
M	S	1147(1115.74)	(0.04)	CHIEN	66 HBC 6.9 PBAR P
M	S	972(1115.69)	(0.05)	CHIEN	66 HBC 6.9 PBAR PANTIL
M	1115.6	0.4	LONDON	66 HBC	
M	(1116.0)	(0.2)	BADIER	67 HBC	2.4 PBAR P, LLBAR
M	195	1115.39	0.12	MAYEUR	67 EMUL
M	B	1524(1115.52)	(0.05)	BOHM	70 EMUL
M	935	1115.59	0.08	HYMAN	72 HBC
M	B	AVERAGE OF VERY INCONSISTENT DATA. ERROR STATISTICAL ONLY. AUTHORS			
M	B	DETECT SYSTEMATIC EFFECT OF ABOUT .15 MEV, WHICH THEY ATTRIBUTE			
M	B	TO ERROR IN RANGE-ENERGY RELATIONS, IN REGION BETA=0.6-0.7.			
M	B	THIS EFFECT, IF CONFIRMED, WOULD AFFECT VERY LITTLE THE VALUES OF			
M	B	BHOWMIK 63 AND MAYEUR 67.			
M	S	ERROR PURELY STATISTICAL.			
M	AVG	1115.566	0.056	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.3)
M	FIT	1115.596	0.046	FROM FIT	(ERROR INCLUDES SCALE FACTOR OF 1.2)

### WEIGHTED AVERAGE 1115.566 ± 0.056 (ERROR SCALED BY 1.3)



### A - A MASS DIFFERENCE (MeV)

DM	TEST OF CPT		CHIEN	66 HBC	6.9 PBAR P
DM	0.05	0.06	BADIER	67 HBC	2.4 PBAR P
DM	0.29	0.15			
DM	AVG	0.083	AVERAGE	(ERROR INCLUDES SCALE FACTOR OF 1.5)	

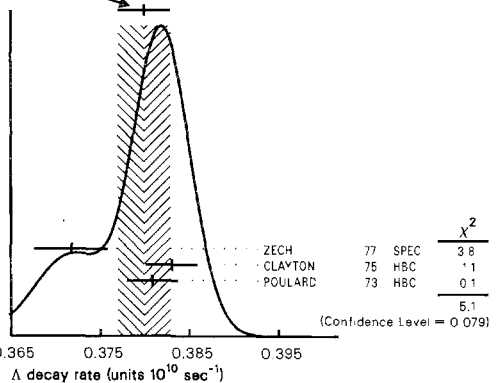
### A MEAN LIFE (units 10<sup>-10</sup> sec)

T	0	188	(2.63)	(0.21)	(0.21)	BOLDT	58 CC
T	0	825	(2.72)	(0.16)	(0.16)	CRAWFORD	59 HBC
T	0	140	(2.72)	(0.29)	(0.27)	BOWEN	60 CC
T	0	186	(2.60)	(0.28)	(0.20)	CHANG	62 HBC
T	0	799	(2.69)	(0.11)	(0.11)	HUMPHREY	62 HBC
T	0	2239	(2.36)	(0.06)	(0.06)	BLICK	63 HBC
T	0	706	(2.76)	(0.20)		CHRETIEN	63 HLBC
T	0	794	(2.59)	(0.09)		HUBBARD	64 HBC
T	0	2260	(2.31)	(0.10)		KREISLER	64 OSPK
T	0	1378	(2.59)	(0.07)		SCHWARTZ	64 HBC
T	0	635	(2.51)	(0.16)		BALTAY	65 HBC
T	0	2534	(2.6)	(0.1)		HILL	65 OSPK
T	0	916	(2.35)	(0.09)		BURAN	66 HLBC
T	S	1147	(2.50)	(0.14)		CHIEN	66 HBC
T	S	972	(2.70)	(0.20)		CHIEN	66 HBC
T	0	2213	(2.452)	(0.056)	(0.054)	VENIGELMANN	67 HBC
T	0	585	(2.68)	(0.13)	(0.11)	AUERBACH	67 DSPK

T	0	(2.44)	(0.15)	BADIER	67 HBC	2.4 PBAR P
T	0	(2.55)	(0.15)	BADIER	67 HBC	2.4 PBAR P, ANTIL
T	0	8342	(2.535)	(0.035)	GRIMM	68 HBC
T	0	2600	(2.47)	(0.08)	HEPP	68 HBC
T	0	1059	(2.39)	(0.10)	DEMIDOV	70 HLBC PI-P, 3.86 GEV/C
T	0	4572	(2.54)	(0.04)	BALTAY	71 HBC K-P AT REST
T	0	6582	(2.69)	(0.05)	ALTHOFF2	73 OSPK PI-N TO K+LAMBDA
T	36K	2.626	0.020	POULLARD	73 HBC K-P, KMOM .4702.3	
T	34K	2.611	0.020	CLAYTON	75 HBC K-P, KMOM .96-1.4	
T	53K	2.69	0.03	ZECH	77 SPEC NEUTRAL HYP. BEAM	
T	0	OLD LOWER STATISTICS EXPERIMENTS NOT INCLUDED IN AVERAGE.				
T	S	ERROR PURELY STATISTICAL.				
T	AVG	2.632	0.020	0.020	AVERAGE (ERROR INCL. SCALE FACTOR OF 1.6)	

(SEE IDEOGRAM BELOW)

### WEIGHTED AVERAGE 0.3799 ± 0.0029 (ERROR SCALED BY 1.6)



### (A - A)/AVG., MEAN LIFE DIFFERENCE

DT	TEST OF CPT		BADIER	67 HBC	2.4 PBAR P
DT	0.044	0.085			

### A MAGNETIC MOMENT (MAGNETONS, 938.26 MeV)

MM	-1.5	0.5	COOL	62 OSPK		
MM	0.0	0.6	KERNAN	63 CC		
MM	8553	-1.39	0.72	ANDERSON	64 HBC	
MM	151	-0.5	0.28	CHARRIERE	65 EMUL	
MM	49	(-0.67)	(0.31)	(0.37)	BARKOV	71 EMUL PRELIM. RESULT
MM	1300	-0.66	0.07	DAHLJENSE	71 EMUL	MAG FIELD=200KG
MM	3868	-0.73	0.18	HILL	71 OSPK	
MM	57	-0.65	0.28	BARKOV	72 EMUL	INCLUDES BARKOV 0.7
MM	1.2M	-0.57	0.05	BUNCE	76 SPEC	
MM	350K	-0.59	0.07	HELLER	77 SPEC	
MM	3M	-0.6138	0.0047	SCHACHING	78 SPEC	
MM	200K	-0.606	0.015	COX	81 SPEC	
MM	AVG	-0.6130	0.0044	AVERAGE		

### A ELECTRIC DIPOLE MOMENT (units 10<sup>-14</sup> e-cm) NONZERO VALUE IMPLIES VIOLATION OF T AND P

EDM	(5.0)	OR LESS	CL=.95	GIBSON	66 EMUL	
EDM	B	(1.0)	OR LESS	CL=.95	BARONI	71 EMUL
EDM	P	0.015	OR LESS	CL=.95	PONDROM	81 SPEC
EDM	B	BARONI MEASURES (-5.9--2.9)*10**15 E CM				
EDM	P	PONDROM 81 MEASURE (-3.0--7.4)*10**17 E CM				

### A PARTIAL DECAY MODES

P1	$\Lambda \rightarrow p \pi^-$	938+ 140
P2	$\Lambda \rightarrow n \pi^0$	940+ 135
P3	$\Lambda \rightarrow p \mu^- \nu$	938+ 106- 0
P4	$\Lambda \rightarrow p e^- \nu$	938- 511+ 0
P5	$\Lambda \rightarrow p \pi^- \gamma$	938- 140+ 0

### A BRANCHING RATIOS

$\Lambda \rightarrow (p \pi^-) / ((p \pi^-) + (n \pi^0))$		(P1)/(P1+P2)		
R1	0.627	0.031		
R1	0.65	0.05		
R1	U	(0.685)	(0.017)	
R1	903	0.643	0.016	
R1	U	6736	0.635	0.007
R1	4572	0.646	0.008	
R1	U	ANDERSON RESULT NOT PUBLISHED, EVENTS ADDED TO DOYLE SAMPLE.		
R1	AVG	0.6399	0.0049	
R1	FIT	0.6419	0.0049	

FROM FIT

For notation, see key on page 91.

Stable Particle Full Listings

A

Table with columns for particle type, values, and authors. Includes entries for (p pi^0)/(p pi^-) + (n pi^0) and (P2)/(P1+P2).

Table with columns for particle type, values, and authors. Includes entries for (p e^- nu)/total (units 10^-3) and (P4)/(P1+P2).

Table with columns for particle type, values, and authors. Includes entries for (p mu^- nu)/total (units 10^-4) and (P3)/(P1+P2).

Table with columns for particle type, values, and authors. Includes entries for (p e^- nu)/(p pi^-) (units 10^-3) and (P4)/(P1).

Table with columns for particle type, values, and authors. Includes entries for (p pi^+ gamma)/(p pi^-) (units 10^-3) and (P5)/(P1).

DELTA DECAY PARAMETERS

SEE NOTE ON BARYON DECAY PARAMETERS IN NEUTRON SECTION ABOVE.

Table with columns for alpha\_-(lambda -> pi^- p), values, and authors. Includes entries for lambda = 1156, 10130, M 2529, etc.

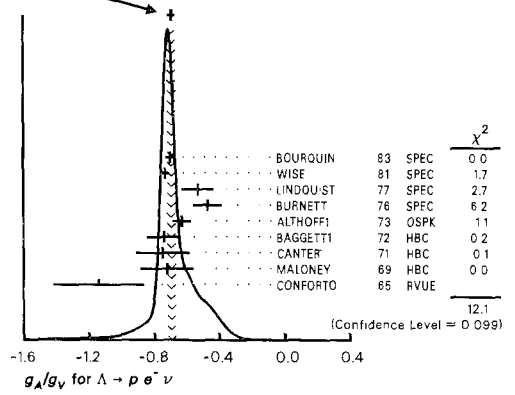
Table with columns for alpha\_0/alpha\_- FOR DELTA (lambda -> pi^0 n / lambda -> pi^- p), values, and authors. Includes entries for lambda = 1.10, 4.760, etc.

Table with columns for alpha(lambda)/alpha(alpha), values, and authors. Includes entries for lambda = 10k, 10130, etc.

Table with columns for phi ANGLE (degrees), values, and authors. Includes entries for lambda = 1156, 10130, etc.

Table with columns for g\_A/g\_V FOR DELTA -> p e^- nu, values, and authors. Includes entries for lambda = 1.14, 102, etc.

WEIGHTED AVERAGE -0.694 +/- 0.025 (ERROR SCALED BY 1.3)



REFERENCES FOR DELTA

A list of references for the DELTA baryon, including authors like EISLER, BOLDT, CRAWFORD, BAGLIN, BOWEN, CORK, COLUMBIA, HUMPHREY, ALSTON, AUBERT, CHANG, COOL, GOOD, HUMPHREY, ALSTON, BHOWMIK, BLOCK, BROWN, CHRETIEN, CRONIN, ELY, KERNAN, ANDERSON, BAGLIN, HUBBARD, KERNAN, KREISLER, LIND, RONNE, SCHWARTZ, BAGLIN, BALTAY, BARLOW, CHARRIERE, CONFORTO, CONFORTO, ELY, HILL, SCHMIDT, BERGE, BURAN, CHEN, ENGELMANN, GIBSON, LONDON, AUERBACH, BAJOIE, DOYLE, MALONEY, BOHM, DENIDOV, OLSEN, ALTHOFF1, ALTHOFF2, BALYAY, BARKOV, BARONI, CANTER, CANTER1, DAHLJENS, HILL, LINDQUIST, BAGGETT1, BAGGETT2, BAGGETT3, BARKOV, CLELAND, HYMAN, EISLER, PLANO, SAMIOS, SCHWARTZ, E BOLDT, D O CALDWELL, Y PAL, CRAWFORD, CRESTI, DOUGLASS, GOOD, BAGLIN, BLOCH, BRISSON, HENNESSY, BOWEN, HARDY, REYNOLDS, SUN, CORK, KERTH, WENZEL, CRONIN, M SCHWARTZ, HUMPHREY, KIRZ, ROSENFELD, RHEE, ANDERSON, CRAWFORD, GOLDEN, LLOYD, AUBERT, BRISSON, HENNESSY, SIX, CHUEN CHUEN CHANG, COOL, HILL, MARSHALL, M L GOOD, V G LIND, W E HUMPHREY, R R ROSS, ALSTON, KIRZ, NEUFELD, SOLMITS, WOHLMUT, B BHOWMIK, D P GOYAL, BLOCK, GESSARDI, RATTI, NWES-BGNA, SYRA, ORNL, BROWN, KADYK, TRILLING, ROE, CHRETIEN, CROUCH, W J CRONIN, O E OVERSETH, ELY, GIDAL, KALMUS, OSWALD, POWELL, KERNAN, MOVEY, WARSHAW, WATTENBERG, J A ANDERSON, F S CRAWFORD, BAGLIN, BINGHAM, HUBBARD, BERGE, KALB, FLEISCH, SHAFER, KERNAN, POWELL, SANDLER, M N KREISLER, D OVERSETH, J CRONIN, LIND, BINFORD, GOOD, STERN, RONNE, COERN+EPOL+LOUC+UNIV.BERGEN, JOSEPH ADAM SCHWARTZ, BAGLIN, BALTAY, SANDWEISS, CULWICK, KOPP, BARLOW, BLAIR, CONFORTO, CHARRIERE, GIBSON, CONFORTO, CONFORTO, ELY, HILL, L I, JENKINS, KYCIA, RUDERMAN, P SCHMIDT, BERGE, CABIBBO, BURAN, EVINSDON, SKJEFFEGSTAD, TOFTE, LACH, SANDWEISS, TAIT, YEH, OREN, ENGELMANN, FILTHUTH, ALEXANDER, W M GIBSON, K GREEN, LONDON, RAU, GOLDBERG, LICHTMAN, AUERBACH, BOWEN, DOBBS, LANDE, MANN, BONNET, BRIANDET, SADOULET, CLELAND, BIEMLEIN, CONFORTO, MAYEUR, E. TOMPA, J WICKENS, O E OVERSETH, R F ROTH, H.-J. GRIMM, V. HEPP, H. SCHLEICH, MERRILL, SHAFER, +BERGE, HUBBARD, MERRILL, MILLER, J.C. DOYLE, MALONEY, DECHEN, KRECKER, +BERL+BRUX+DUUC+LOUC+WARS, +KIRILLOV+UGRYUMOV, PONOSOV, PROTASOV, +PONDROM, HANDLER, LIMON, SMITH, ALTHOFF1, ALTHOFF2, BALYAY, BARKOV, BARONI, CANTER, CANTER1, DAHLJENS, HILL, LINDQUIST, BAGGETT1, BAGGETT2, BAGGETT3, BARKOV, CLELAND, HYMAN, +BAGGETT, EISELE, FILTHUTH, FREHSE, +BAGGETT, EISELE, FILTHUTH, FREHSE, HEPP, +BUREVICH, MAKARINA, MARTYNYANOV, +CONFORTO, EATON, GERBER, +BUNNELL, BERRICK, FIELDS, KATZ,

# Stable Particle Full Listings

## $\Lambda, \Sigma^+$

ALTHOFF 73 PL 438 237 ALTHOFF 75 NP 866 29 KATZ 75 THESIS (MARYLAND) POULARD 73 PL 468 135	+BROWN, FREYTAG, HEARD, HEINTZE- (CERN+HEID) +BROWN, FREYTAG, HEARD, HEINTZE- (CERN-HEID) C.N. KATZ (UMD) +GIVERNAUD, BORG (SACL)
ASTBURY 75 NP 899 30 CLAYTON 75 NP 895 130 BUNCE 76 PRL 36 1113 BURNETT 76 NC 34A 14	+GALLIVAN, JAFAR + (LOIC-CERN+ETH+SACL) +BACON, BUTTERWORTH, WATERS + (LOIC+RHEL) +HANDLER, MARCH, MARTIN + (WISC+MICH+RUTE) +INNES, HASEK, HAUNE, MILLER, RUDERMAN+ (UCSC)
HELLER 77 PL 688 480 LINDQUIST 77 PR D16 2104 ALSO 76 PG 2 L211 ZECH 77 NP 8124 413 SCHACHIN 78 PRL 41 1348	+OVERSETH, BUNCE, DYDAK + (MICH+WISC+HEID) LINDQUIST, SWALLOW, SUMNER + (EFI+OSU+ANL) LINDQUIST, SWALLOW, SUMNER+ (EFI+WISC+OSU+ANL) +DYDAK, NAVARRIA+ (STEG+CERN+DORT+HEID) SCHACHINGER, BUNCE, COX + (MICH, RUTG+WISC)
WISE 80 PL 918 165 COX 81 PRL 46 877 PONDROM 81 PR D23 814 WISE 81 PL 988 123 BOURQUIN 83 ZPHY C21 1 CHAUVAUT 85 PL 1638 273	-JENSEN, KREISLER, LOMANNO, POSTER- (MASA+BNL) +DWORKIN (MICH+WISC+RUTG+MINN+BNL) +HANDLER, SHEAFF, COX + (WISC+MICH+RUTG+MINN) +JENSEN, KREISLER, LOMANNO, POSTER+ (MASA+BNL) +BROWN+ (BRIS+GEVA+HEID-LALD+RL+STRB) +ERHAN, HAYES, + (CERN+UDCF+UCLA+SACL)

PAPERS NOT REFERRED TO IN DATA LISTINGS

ARMENTER 62 CERN CONF 236 BALTAY 62 CERN CONF 233 BERGE 63 THESIS (BERKELEY)	ARMENTEROS+ (CERN+EPOL+LOIC+BIRM+CEN+SACLAY) BALTAY, FOWLER, SANDWEISS, CULWICK+ (YALE+BNL) J PETER BERGE (LRL)
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$\Sigma^+$

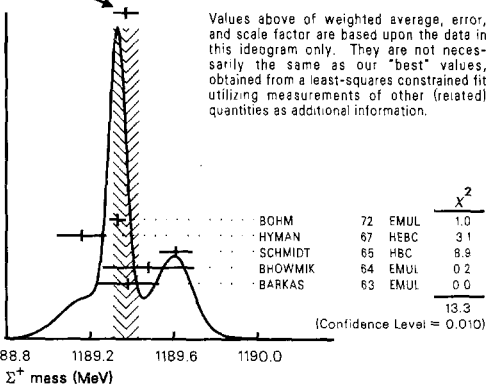
$$I(J^P) = 1(\frac{1}{2}^+)$$

### $\Sigma^+$ MASS (MeV)

SEE NOTE PRECEDING LAMBDA MASS LISTINGS

M 144 1189.58 0.15 M 58 1189.48 0.22 M S ABOVE SIGMA- MASSES HAVE BEEN RAISED 30 KEV TO ACCOUNT FOR 46 KEV M S INCREASE IN PROTON MASS AND 21 KEV DECREASE IN PION MASS M 4205 1189.61 0.08 M 1189.16 0.12 M B 607 1189.33 0.04 M B BOHM 72 UPDATED WITH PDG AVER. 73 K-, PI- AND P10 MASSES. M 1189.371 0.060 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8) M FIT 1189.565 0.057 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.8) (SEE IDEOGRAM BELOW)	BARKAS 63 EMUL + SEE NOTE S BELOW BHOWMIK 64 EMUL + SEE NOTE S BELOW SCHMIDT 65 HBC SEE NOTE N HYMAN 67 HBC BOHM 72 EMUL BOHM 72 UPDATED WITH PDG AVER. 73 K-, PI- AND P10 MASSES. BARKAS 63 EMUL + SEE NOTE S BELOW BHOWMIK 64 EMUL + SEE NOTE S BELOW SCHMIDT 65 HBC SEE NOTE N HYMAN 67 HBC BOHM 72 EMUL BOHM 72 UPDATED WITH PDG AVER. 73 K-, PI- AND P10 MASSES.
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WEIGHTED AVERAGE  
1189.371 ± 0.060 (ERROR SCALED BY 1.8)



Values above of weighted average, error, and scale factor are based upon the data in this ideogram only. They are not necessarily the same as our "best" values, obtained from a least-squares constrained fit utilizing measurements of other (related) quantities as additional information.

### $\Sigma^+$ MEAN LIFE (units 10<sup>-10</sup> sec)

T 127 0.98 0.16 T 41 0.82 0.34 T 117 0.85 0.14 T 54 0.80 0.10 T 23 0.76 0.22 T 49 0.75 0.13 T 140 0.82 0.10 T 192 0.749 0.056 T 456 0.765 0.04 T 203 0.84 0.12 T 181 0.84 0.09 T 900 0.76 0.03 T C 1300 0.83 0.052 T S 125 (0.86) (0.15) T S 117 (1.10) (0.24) T 10664 0.803 0.008 T 20K 0.795 0.010 T 526 0.85 0.14 T 5719 0.807 0.013 T 30K 0.798 0.005 T S CHANGE ERROR 0.018 RAISED BY US. SEE 1970 EDITION, RMP 42, 123(1970) T S ERROR PURELY STATISTICAL T AVG 0.7997 0.0036 0.0036 AVERAGE	GLASER 58 RVUE PUSCHEL 60 EMUL EVANS 60 EMUL FREDEN 60 EMUL KAPLON 60 EMUL CHIESA 61 EMUL BERTHELOT 61 HLBC BARKAS 61 EMUL GRAND 62 HBC HUMPHREY 62 HBC BHOWMIK 64 EMUL BALTAY 65 HBC CARAYAN 65 HBC CHANG 66 HBC CHIEN 66 HBC CHIEN 66 HBC + 6.9 PBAR P COOK 66 OSPK BARLOUTAU 69 HBC K-P -4.1.2 GEV/C EISELE 70 HBC K-P AT REST BAKKER 71 HBC K-M TO SIG+ 2PI- CONFORTO 76 HBC K-P 1-1.4 GEV/C MARRAFFIN 80 HBC K-P TO SIG+ PI-
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### $\Sigma^+$ MAGNETIC MOMENT (MAGNETONS, 938.26 MeV)

MM 381 1.5 1.1 MM 52 3.5 1.5 MM 53 5.0 1.2 MM 69 3.5 1.2 MM 29333 2.1 1.0 MM 955 2.67 0.97 MM 2651 2.7 0.9 MM S 8503 (2.95) (0.31) MM S 14K 2.30 0.14 MM 44K 2.38 0.02 MM S SETTLES 79 INCLUDES DOBLE 77 DATA. MM AVG 2.379 0.020 AVERAGE	COOK 66 OSPK KOTELCHUC 67 EMUL K-P AT 1.15BEV/C SULLIVAN 67 EMUL PHOTOPRODUCTION COMBE 68 EMUL MAST 68 HBC K-P AT .4 GEV/C ALLEY 71 OSPK 1.28 GEV/C PI+P SAHA 73 HLBC K-P .25TO .55GEV/C DOBLE 77 HBC K-P .46 GEV/C SETTLES 79 HBC K-P .42TO .50GEV/C ANKENBRAN 83 CNTR 210GEV HYP. BEAM
--	---

### $\Sigma^+$ PARTIAL DECAY MODES

P1 $\Sigma^+ \rightarrow p \pi^0$ P2 $\Sigma^+ \rightarrow n \pi^+$ P3 $\Sigma^+ \rightarrow n \pi^+ \gamma$ P4 $\Sigma^+ \rightarrow \Lambda e^+ \nu$ P5 $\Sigma^+ \rightarrow p \gamma$ P6 $\Sigma^+ \rightarrow n \mu^+ \nu$ P7 $\Sigma^+ \rightarrow n e^+ \nu$ P8 $\Sigma^+ \rightarrow p e^+ e^-$	DECAY MASSES 938+ 135 940+ 140 940+ 140+ 0 1116+.511+ 0 938+ 0 940+ 106+ 0 940+ 511+ 0 938+.511+.511
--	--

### $\Sigma^+$ BRANCHING RATIOS

$\Sigma^+ \rightarrow (n \pi^+)/(N \pi)$ R1 308 0.490 0.024 R1 534 0.46 0.02 R1 1331 0.488 0.010 R1 537 0.484 0.015 R1 1861 0.488 0.008 R1 M 10K 0.4828 0.0036 R1 M MARRAFFINO 80 GIVES BR TO (P P10)/ALL. WE QUOTE 1-BR. R1 AVG 0.4836 0.0030 AVERAGE	HUMPHREY 62 HBC CHANG 66 HBC BARLOUTAU 69 HBC K-P -4.1.2 GEV/C TOVES 71 EMUL NOMAK 78 HBC BRITISH 1.5M (TST) MARRAFFIN 80 HBC K-P 420-500MEV/C WE QUOTE 1-BR.
$\Sigma^+ \rightarrow (n \pi^+ \gamma)/(\pi^+ n)$ (units 10 <sup>-3</sup> ) R2 29 (1.8) ABOUT R2 180 (0.27) (0.05) R2 180 0.93 10 R2 P1+ MOMENTUM CUTS DIFFER, NOT AVERAGED. LATEST VALUE USED IN SUMMARY TABLE.	BAZINZ 65 HBC PI- LT 116 MEV/C ANG 69 HBC PI+ LT 110 MEV/C EBENHOH 73 HBC PI- LT 150 MEV/C
$\Sigma^+ \rightarrow (\Lambda e^+ \nu)/total$ (units 10 <sup>-5</sup> ) R3 W 4 (3.3) (1.7) R3 6 2.0 0.8 R3 5 1.6 0.7 R3 10 2.9 1.0 R3 AVG 2.02 0.47 AVERAGE	WILLIS 64 HBC STOP K- EISELE1 69 BARASH 67 HBC STOP K- BALTAY 69 HBC STOP K- EISELE1 69 HBC STOP K-
$\Sigma^+ \rightarrow (p \gamma)/(p \pi^0)$ (units 10 <sup>-3</sup> ) R4 1 (0.068) OR LESS R4 24 0.37 0.08 R4 6 (0.17) R4 45 0.21 0.03 R4 31 0.276 0.051 R4 46 0.211 0.058 R4 155 0.246 0.030 R4 AVG 0.236 0.019 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)	CARRARA 64 HBC BAZINZ 65 HBC QUARENZI 65 EMUL ANG 69 HBC STOP K- GERSHWIN 69 HBC MANZ 80 HBC K-P --> SIGMA- PI- BIAGI 85 CNTR INCLUDES HYP. BEAM
$\Sigma^+ \rightarrow (n e^+ \nu)/(n \pi^+)$ (units 10 <sup>-5</sup> ) R5 TEST OF DELTA-S = DELTA-Q RULE R5 E 0 (16220) EFFECTIVE DENOM. R5 E 0 (2720) EFFECTIVE DENOM. R5 E 0 1 (9690) EFFECTIVE DENOM. R5 U 0 (32406) EFFECTIVE DENOM. R5 U 0 (80400) EFFECTIVE DENOM. R5 U 0 1 (30000) EFFECTIVE DENOM. R5 O OLDER LOWER STATISTICS EXPTS. NOT INCLUDED IN AVERAGE. R5 U 0 105000 EFFECTIVE DENOM. R5 U A 0 111000 EFFECTIVE DENOM. R5 U A EFFECTIVE DENOM. CALCULATED BY US R5 E EFFECTIVE DENOM. TAKEN FROM EISELE 67 R5 E A EISELE2 69 REPLACED BY EBENHOH 74. R5 AVG 1.1 OR LESS. CL=.90 OUR AVERAGE (2.3 EVTS)/(EFF.DENOM.SUM) R5 NUMBER OF EVENTS INCREASED TO 2.5 FOR 90% CONFIDENCE LEVEL	COURANT 64 HBC SEE NOTE E MURPHY 64 HBC SEE NOTE E NAUENBERG 64 HBC SEE NOTE E BIERMAN 68 HBC EISELE2 69 HBC NORTON 69 HBC SECHIZORN 73 HBC STOP K- EBENHOH 74 HBC STOP K- CALCULATED BY US TAKEN FROM EISELE 67 OUR AVERAGE (2.3 EVTS)/(EFF.DENOM.SUM) NUMBER OF EVENTS INCREASED TO 2.5 FOR 90% CONFIDENCE LEVEL
$\Sigma^+ \rightarrow (n \mu^+ \nu)/(\pi^+ n)$ (units 10 <sup>-5</sup> ) R6 TEST OF DELTA-S = DELTA-Q RULE R6 1 (120) ANALYZED EVENTS R6 E 0 10150 EFFECTIVE DENOM. R6 E 0 1710 EFFECTIVE DENOM. R6 U 2 62000 EFFECTIVE DENOM. R6 E 0 33800 EFFECTIVE DENOM. R6 U EFFECTIVE DENOM. CALCULATED BY US R6 AVG 6.2 OR LESS. CL=.90 OUR AVERAGE (6.7 EVTS)/(EFF.DENOM.SUM) R6 NUMBER OF EVENTS INCREASED TO 6.7 FOR 90% CONFIDENCE LEVEL	BALTIERI 62 EMUL NO RATIO QUOTED COURANT 64 HBC SEE NOTE E NAUENBERG 64 HBC SEE NOTE E EISELE2 69 HBC BAGGETT 69 HBC CALCULATED BY US OUR AVERAGE (6.7 EVTS)/(EFF.DENOM.SUM) NUMBER OF EVENTS INCREASED TO 6.7 FOR 90% CONFIDENCE LEVEL
$(\Sigma^+ \rightarrow n e^+ \nu)/(\Sigma^+ \rightarrow n e^- \nu)$ R7 TEST OF DELTA-S = DELTA-Q RULE R7 0 (0.35) OR LESS R7 1 (0.08) OR LESS R7 AVG 0.043 OR LESS. CL=.90 OUR AVERAGE USING R5 AND R6	BAGGETT 67 HBC NORTON 69 HBC
$\Sigma^+ \rightarrow (p e^+ e^-)/total$ (units 10 <sup>-6</sup> ) R8 A ANG 69 FOUND 3 E+ E- EVENTS IN AGREEMENT WITH GAMMA CONVERSION OF PROTON GAMMA DECAY - LIMIT GIVEN HERE IS FOR NEUTRAL CURRENT	ANG 69 HBC STOP K-

For notation, see key on page 91.

Stable Particle Full Listings

$\Sigma^+, \Sigma^-$

Table with columns for particle type (R9, R9), parameters (0.06, 0.045, 0.03), and references (EISELE2, 69 HBC). Includes sub-sections for  $\Sigma^+$  and  $\Sigma^-$  decay parameters.

$\Sigma^+$  DECAY PARAMETERS

SEE NOTE ON BARYON DECAY PARAMETERS IN NEUTRON SECTION ABOVE.

Table for  $\alpha_+$  FOR  $\Sigma^+ \rightarrow \pi^+ n$  with columns for  $\alpha_+$ ,  $\alpha_0$ , and various experimental data points.

Table for  $\alpha_+$  FOR  $\Sigma^+ \rightarrow \pi^+ n$  with columns for  $\alpha_+$ ,  $\alpha_0$ , and various experimental data points.

Table for  $\alpha_0$  FOR  $\Sigma^+ \rightarrow \pi^0 p$  with columns for  $\alpha_0$ ,  $\alpha_+$ , and various experimental data points.

Table for  $\phi_+$  ANGLE  $(\Sigma^+ \rightarrow \pi^+ n)$  (degrees) with columns for  $\phi_+$ ,  $\alpha_+$ , and various experimental data points.

Table for  $\alpha_\gamma$  FOR  $\Sigma^+ \rightarrow p \gamma$  with columns for  $\alpha_\gamma$ ,  $\alpha_+$ , and various experimental data points.

Table for  $\phi_0$  ANGLE  $(\Sigma^+ \rightarrow \pi^0 p)$  (degrees) with columns for  $\phi_0$ ,  $\alpha_+$ , and various experimental data points.

REFERENCES FOR  $\Sigma^+$

List of references for  $\Sigma^+$  decays, including authors like CORK, ERTH, WENZEL, and institutions like CERN, BNL, and various universities.

Main table listing experimental data for  $\Sigma^+$  and  $\Sigma^-$  particles, including columns for particle type, experiment name, date, and references.

$\Sigma^-$   $(J^P) = (1/2^+)$

$\Sigma^-$  MASS (MeV)

Table showing mass measurements for  $\Sigma^-$  particles, including columns for mass (M), neutron number (N), and references.

$\Sigma^- - \Sigma^+$  MASS DIFFERENCE (MeV)

Table showing mass differences between  $\Sigma^-$  and  $\Sigma^+$ , including columns for mass difference (D), mass (M), and references.

$\Sigma^- - \Lambda$  MASS DIFFERENCE (MeV)

Table showing mass differences between  $\Sigma^-$  and  $\Lambda$ , including columns for mass difference (DL), mass (M), and references.

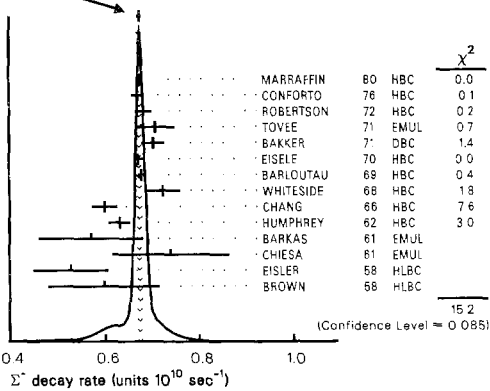
# Stable Particle Full Listings

$\Sigma^-$

## $\Sigma^-$ MEAN LIFE (units $10^{-10}$ sec)

T	1.67	0.40	0.28	BROWN	58	HLBC	
T	1.89	0.33	0.25	EISLER	58	HLBC	
T	45	1.35	0.32	0.17	CHIESA	61	EMUL
T	41	1.75	0.39	0.30	BARKAS	61	EMUL
T	1208	1.58	0.06	0.06	HUMPHREY	62	HBC
T	3267	1.666	0.075		CHANG	66	HBC
T	61	(2.08)	(0.22)		CHIEN	66	HBC
T	64	(1.46)	(0.31)		CHIEN	66	HBC
T	506	1.38	0.07		WHITESIDE	68	HBC
T	10253	1.472	0.016		BARLOUTAU	69	HBC
T	0.1M	1.485	0.022		EISELE	70	HBC
T	1383	1.42	0.05		BAKKER	71	DBC
T		1.41	0.09	0.08	TOVEE	71	EMUL
T	2400	1.463	0.039		ROBERTSON	72	HBC
T	8437	1.49	0.03		CONFORTO	76	HBC
T	16k	1.480	0.014		MARRAFFIN	80	HBC
T	CHANG ERROR 0.018 RAISED BY US. SEE 1970 EDITION, RMP 42,125(1970)						
T	ERROR PURELY STATISTICAL.						
T	S						
T	AVG	1.482	0.011	0.011	AVERAGE (ERROR INCL. SCALE FACTOR OF 1.3)		
T	(SEE IDEOGRAM BELOW)						

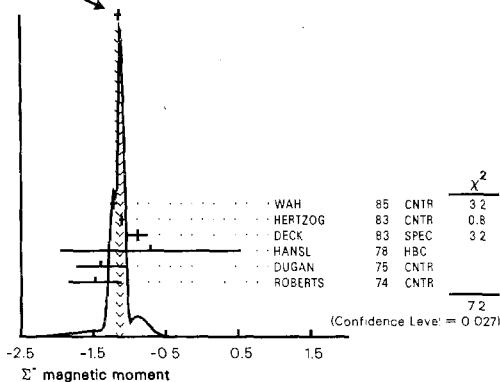
WEIGHTED AVERAGE  
0.6747 ± 0.0050 (ERROR SCALED BY 1.3)



## $\Sigma^-$ MAGNETIC MOMENT (MAGNETONS, 938.26 MeV)

MM	R	BTWN -1.6 AND -0.8	FOX	73	CNTR	SIG-ATOM FINE ST	
MM	R	-1.48	0.37	ROBERTS	74	CNTR	
MM	D	-1.40	0.41	0.28	DUGAN	75	CNTR
MM	D	(0.65)	(0.28)	(0.40)	DUGAN	75	CNTR
MM		28k	-0.71	1.25	HANSL	78	HBC
MM		516k	-0.89	0.14	DECK	83	SPEC
MM			-1.11	0.033	HERTZOG	83	CNTR
MM			-1.23	0.05	WAN	85	CNTR
MM	R	ROBERTS 74 INCLUDES DATA FROM FOX 73.					
MM	D	DUGAN 75 NEGATIVE VALUE AVERAGED SINCE IT AGREES WITH ROBERTS 74.					
MM	W	STATISTICAL AND SYSTEMATIC ERRORS ARE EACH 0.03.					
MM	AVG	-1.141	0.051	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.9)			
				(SEE IDEOGRAM BELOW)			

WEIGHTED AVERAGE  
-1.141 ± 0.051 (ERROR SCALED BY 1.9)



## $\Sigma^-$ PARTIAL DECAY MODES

P1	$\Sigma^- \rightarrow n\pi^-$	DECAY MASSES
P2	$\Sigma^- \rightarrow n\pi^-\gamma$	940+ 140
P3	$\Sigma^- \rightarrow n\mu^-\nu$	940+ 140+ 0
P4	$\Sigma^- \rightarrow n e^-\nu$	940+ 106+ 0
P5	$\Sigma^- \rightarrow \Lambda e^-\nu$	940+ 511+ 0

## $\Sigma^-$ BRANCHING RATIOS

$\Sigma^- \rightarrow (n\mu^-\nu)/(n\pi^-)$ (units $10^{-3}$ )					(P3)/(P1)			
R1	22	0.66	0.15	COURANT	64	HBC		
R1	11	0.56	0.20	BAZIN	65	HBC		
R1	72	0.43	0.09	BAGGETT	69	HBC		
R1	72	0.43	0.06	ANG 1	69	HBC		
R1	13	0.38	0.11	COLE	71	HBC		
R1	AVG	0.447	0.043	AVERAGE				
$\Sigma^- \rightarrow (\Lambda e^-\nu)/(n\pi^-)$ (units $10^{-3}$ )					(P4)/(P1)			
R2	9	(1.0)	(0.4)	(0.3)	MURPHY	64	HLBC	
R2	16	(1.37)	(0.34)		NAUENBERG	64	HBC	
R2	16	(1.15)	(0.4)		MILLER	64	FBC	
R2	51	(1.4)	(0.3)		COURANT	64	HBC	
R2	180	1.11	0.09		BIESMAN	68	HBC	
R2	331	(1.02)	(0.08)		ANG 1	69	HBC	
R2	57	0.97	0.15		COLE	71	HBC	
R2	A	452	0.46	0.07	0.13	SECHIZON	73	HBC
R2	A	601	1.09	0.06	0.08	ESENHOH	74	HBC
R2		2847	0.96	0.05		BOURQUIN 1	83	SPEC
R2	A	ADDITIONAL NEGATIVE SYSTEMATIC ERROR INCLUDED FOR INTERNAL						
R2	A	RADIATIVE CORRECTIONS AND LATEST FORM FACTORS. SEE BOURQUIN 83.						
R2	AVG	1.022	0.034	AVERAGE				
$\Sigma^- \rightarrow (\Lambda e^-\nu)/(n\pi^-)$ (units $10^{-4}$ )					(P5)/(P1)			
R3	11	0.75	0.28	COURANT	64	HBC		
R3	35	0.66	0.12	BARASH	67	HBC		
R3	31	0.69	0.12	EISELE 1	69	HBC		
R3	31	0.52	0.09	BALTAY	69	HBC		
R3	H	122	(0.60)	(0.11)	HERBERT	78	ASPK	
R3	H	114	0.63	0.11	THOMPSON	80	ASPK	
R3	B	1620	0.561	0.031	BOURQUIN 2	83	SPEC	
R3	H	HERBERT 78 REPLACED BY THOMPSON 80.						
R3	B	VALUE IS FROM BOURQUIN 2 83. INCLUDES RAD. CORR. AND NEW ACCEPTANCE.						
R3	AVG	0.574	0.027	AVERAGE				
$\Sigma^- \rightarrow (n\pi^-\gamma)/(n\pi^-)$ (units $10^{-3}$ )					(P2)/(P1)			
R4		(1.1) APPROXIM.		BAZIN	65	HBC		
R4	23	(0.10)	(0.02)	ANG 2	69	HBC		
R4	292	0.46	0.06	ESENHOH	73	HBC		
R4		PI+ MOMENTUM CUTS DIFFER, NOT AVERAGED. LATEST VALUE USED IN SUMMARY TABLE.						

## $\Sigma^-$ DECAY PARAMETERS

SEE NOTE ON BARYON DECAY PARAMETERS IN NEUTRON SECTION ABOVE.

$\alpha_-(\Sigma^-)$							
A-	0	(-0.16)	(0.21)	TRIPP	62	HBC	
A-	0	6500	(-0.10)	(0.043)	BANGERTER	66	HBC
A-	0	6068	(-0.104)	(0.04)	BERLEY	67	HBC
A-	5	1000	-0.071	0.012	BANGERTER	69	HBC
A-	B	5978	(-0.134)	(0.034)	BERLEY	70	HBC
A-	6	0000	(-0.067)	0.011	BOGERT	70	HBC
A-	8	28k	-0.062	0.024	HANSL	78	HBC
A-	O OLD RESULTS. HAVE BEEN REPLACED.						
A-	B BERLEY 70 REPLACED BY BOGERT 70.						
A-	AVG	-0.0681	0.0077	AVERAGE			

$\phi$ ANGLE (degrees) (tan $\phi = \beta/\gamma$ )							
F-	0	1006	(+22.)	(30.)	BERLEY	67	HBC
F-		1385	14.	19.	BANGERTI	69	HBC
F-	C	1092	+5.	23.	BERLEY	70	HBC
F-	C CHANGED FROM -5 TO +5 TO AGREE WITH SIGN CONVENTION						
F-	O OLD RESULTS. HAVE BEEN REPLACED.						
F-	AVG	10.3	14.6	AVERAGE			

## NOTE ON $\Sigma^- \rightarrow \Lambda e^-\nu$

(by J.A. Thompson, University of Pittsburgh)

The decay  $\Sigma^- \rightarrow \Lambda e^-\nu$  is of special interest because its form is predicted by the strong form of CVC and is not sensitive to the current octet assumptions or SU(3) structure constants which enter into Cabibbo's predictions for the other hyperon decays. For  $\Delta S = 0$  transitions, the weak interaction vector current is related to the electromagnetic current through a multiplicative constant, set by neutron beta decay, and an isospin rotation.

For notation, see key on page 91.

# Stable Particle Full Listings

Σ<sup>-</sup>

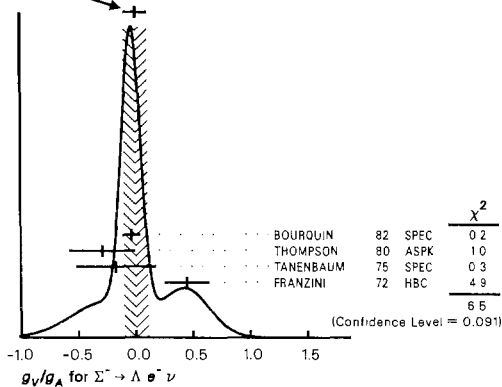
The decay  $\Sigma^0 \rightarrow \Lambda \gamma$  (the isospin-rotation analogue of  $\Sigma^- \rightarrow \Lambda e^- \nu$ ) is mediated predominantly through the magnetic interaction, assuming there are no inhomogeneities in the  $\Sigma^0$ ,  $\Lambda$  charge distributions. Thus we expect the  $g_{WM}$  term,

$$g_{WM} \sim \frac{\mu_{\Sigma\Lambda}}{\sqrt{2}} \sim -\frac{\sqrt{3}}{2} \mu_n \text{ [by SU(3)]},$$

to dominate the vector part of the weak current. The strong CVC predictions are thus:  $g_V/g_A = 0$  and  $g_{WM} \sim 1.6$ .

**$g_V/g_A$  FOR  $\Sigma^- \rightarrow \Lambda e^- \nu$**   
 AV (FOR SIGN CONVENTION, SEE NOTE ON BARYON DECAY PARAMETERS IN NEUTRON SECTION ABOVE.)  
 AV PREDICTED TO BE ZERO BY CONSERVED VECTOR CURRENT THEORY.  
 AV VALUES AVERAGED ASSUME CVC-SU3 WEAK MAGNETISM TERM.  
 AV FB 45 (0.31) (0.30) BARASH 67 HBC  
 AV FS 51 (0.7) (0.4) BALTAY 69 HBC USING SIG--  
 AV FS 81 (-0.22) (0.28) EISELE1 69 HBC  
 AV FS 186 0.45 0.20 FRANZINI 72 HBC USING SIG--  
 AV 55 -0.17 0.35 TANENBAUM 75 SPEC BNL HYPERON BEAM  
 AV 114 -0.29 0.29 THOMPSON 80 ASPK BNL HYPERON BEAM  
 AV S 1620 -0.034 0.080 BOURQUIN 82 SPEC SPS HYPERON BEAM  
 AV B BARASH 67 MEASURED ABSOLUTE VALUE.  
 AV S SIGN CHANGED TO AGREE WITH OUR CONVENTION.  
 AV F FRANZINI 72 INCLUDES EVENTS OF BARASH 67, EISELE1 69, BALTAY 69.  
 AV AVG 0.01 0.10 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5) (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
 $0.01 \pm 0.10$  (ERROR SCALED BY 1.5)

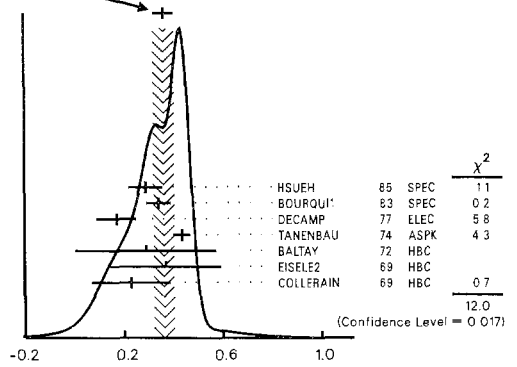


**$g_{WM}/g_A$  FOR  $\Sigma^- \rightarrow \Lambda e^- \nu$**   
 WM VALUES QUOTED ASSUME THE CVC PREDICTION  $g_V=0$ .  
 WM 186 2.4 2.1 FRANZINI 72 HBC USING SIG--  
 WM 55 3.5 4.5 TANENBAUM 75 SPEC BNL HYPERON BEAM  
 WM 114 1.75 3.5 THOMPSON 80 ASPK BNL HYPERON BEAM  
 WM AVG 2.4 1.7 AVERAGE

**$g_A/g_V$  FOR  $\Sigma^- \rightarrow \Lambda e^- \nu$**   
 AV1 (FOR SIGN CONVENTION, SEE NOTE ON BARYON DECAY PARAMETERS IN NEUTRON SECTION ABOVE.)  
 AV1 (0.05) (0.23) (0.32) GERSHWIN 68 HBC REPLACED BY GER.69  
 AV1 57 0.19 (0.20) 0.17 GERSHWIN 69 HBC POLARIZED SIGMAS  
 AV1 61 +0.19 (0.20) 0.17 GERSHWIN 69 HBC POLARIZED SIGMAS  
 AV1 63 -0.33 (0.30) 0.85 BOBERT 70 HBC K-P AT 400 MEV/C  
 AV1 43 -0.2 (0.52) 1.5 ELLIS 72 ASPK POLARIZED SIGMAS  
 AV1 S 193 +0.15 OR LESS CL-.95 KELLER 82 SPEC POLARIZED SIGMAS  
 AV1B S4456 +0.33 BOURQUIN 83 SPEC SPS HYPERON BEAM  
 AV1H S 25K +0.29 0.07 HSUEH 85 SPEC 250 GEV/C SIGMA-  
 AV1 S SIGN CHANGED TO AGREE WITH OUR CONVENTION.  
 AV1B THE VALUE +0.33 IS PREFERRED OVER -0.33 BY AT LEAST 2.6 STD. DEV.  
 AV1B INCLUDING SYSTEMATIC ERRORS.  
 AV1H FROM MEASUREMENT OF ELECTRON ASYMMETRY  $A_E = -0.53 \pm 0.14$

**$|g_A/g_V|$  FOR  $\Sigma^- \rightarrow \Lambda e^- \nu$**   
 AV2 49 0.23 0.16 COLLERAIN 69 HBC NEUTRON SCATTER  
 AV2 33 0.37 0.26 0.19 EISELE2 69 HBC NEUTRON SCATTER  
 AV2 36 0.29 0.28 0.29 BALTAY 72 HBC NEUTRON SCATTER  
 AV2 3507 0.435 0.035 TANENBAUM 74 ASPK  
 AV2 519 0.17 0.07 DECAMP 77 ELEC H.E. HYPERON BEAM  
 AV2 B4456 0.34 0.05 BOURQUIN 83 SPEC SPS HYPERON BEAM  
 AV2H S 25K 0.29 0.07 HSUEH 85 SPEC 250 GEV/C SIGMA-  
 AV2 B POSITIVE SIGN OF  $g_A/g_V$  FAVORED BY AT LEAST 2.6 STANDARD DEVIATIONS.  
 AV2H FROM MEASUREMENT OF ELECTRON ASYMMETRY  $A_E = -0.53 \pm 0.14$   
 AV2 AVG 0.362 0.043 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.7) (SEE IDEOGRAM BELOW)

WEIGHTED AVERAGE  
 $0.362 \pm 0.043$  (ERROR SCALED BY 1.7)



## REFERENCES FOR $\Sigma^-$

BROWN 58 CERN CONF 270	BROWN, GLASER, GRAVES, PERL, CRONIN + (MICH)
EISLER 58 NC SER10 10 150	EISLER, BASSI, CONVERSI + (COLU, BNL, BGNA, PISA)
BARKAS 61 PR 124 1209	BARKAS, DYER, MASON, NICKOLS, SMITH (LRL)
CHIESA 61 NC 19 1171	A M CHIESA, B QUASSIATI, G RINAUDO (TURIN)
HUMPHREY 62 PR 127 1505	W E HUMPHREY, R R ROSS (LRL)
TRIPP 62 PR 9 66	R O TRIPP, M WATSON, M FERRO-LUZZI (LRL)
BARKAS 63 PRL 11 26	W H BARKAS, J N DYER, H H HECKMAN (LRL)
BURNSTEI 64 PRL 13 66	BURNSTEIN, DAY, KEHOE, SECHI ZORN, SNOW (JMD)
COURANT 64 PR 136 B 1791	COURANT, FILTHUTH + (CERN+HEID+UMD+NRL+BNL)
MILLER 64 PL 11 262	MILLER, STANNARD, BEZAGUET + (LOUC, EPOL+BERG)
MURPHY 64 PR 134 B 188	C THORNTON MURPHY (WISCONSIN)
NAUENBER 64 PRL 12 679	NAUENBERG, SCHMIDT, MARATECK + (COLU+RUTG+PRIN)
BAZIN 65 PR 140 B 1358	BAZIN, PLANO, SCHMIDT + (PRIN-RUTG+COLU)
DOSCH 65 PL 14 239	DOSCH, ENGELMANN, FILTHUTH, HEPP, KLUGE + (HEID)
ALSO 66 PR 151 1081	CHUNG YUN CHANG (COLUMBIA)
SCHMIDT 65 PR 140 B 1328	P SCHMIDT (COLUMBIA)
BANGERTER 66 PRL 17 495	BANGERTER, GALTIERI, BERGE, MURRAY + (LRL)
CHANG 66 PR 151 1081	CHUNG YUN CHANG (COLUMBIA)
CHIEN 66 PR 152 1171	+LACH, SANDWEISS, TAFT, YEH, OREN + (YALE+BNL)
BARASH 67 PRL 19 181	BARASH, DAY, GLASSER, KEHOE, KNOP + (MARYLAND)
BERLEY 67 PRL 19 979	BERLEY, HERTZBACH, KOFLER + (BNL, MASA, YALE)
BIERMAN 68 PRL 20 1459	BIERMAN, KOUNOSU, NAUENBERG + (COLUMBIA)
GERSHWIN 68 PRL 20 1270	GERSHWIN, ALSTON-GARNJOST, BANGERTER + (LRL)
HEPP 68 ZPHY 214 71	V. HEPP, H. SCHLEICH (HEIDELBERG)
WHITESID 68 NC 54A 537	H. WHITESIDE, J. GOLLUB (OBERLIN)
ANG 1 69 ZPHY 223 103	ANG, EISELE, ENGELMANN, FILTHUTH + (HEID)
ANG 2 69 ZPHY 228 151	+EBENHORN, EISELE, ENGELMANN, FILTHUTH + (HEID)
BAGGETT 69 PRL 23 249	BAGGETT, KEHOE, SNOW (UNIV MARYLAND)
BALTAY 69 PRL 22 615	BALTAY, FRANZINI, NEWMAN, NORTON + (COLU, STON)
BANGTERE 69 UCRL-19244	ROGER ODELL BANGERTER (THEIS) (LRL)
BANGERT1 69 PR 187 1821	BANGERTER, GARNJOST, GALTIERI, GERSHWIN + (LRL)
BARLOUTA 69 NP 814 153	BARLOUTAUD, BELLEFON, GRANET + (SACL-CERN-HEID)
COLLRAI 69 PRL 23 198	COLLERAINE, DAY, GLASSER, KNOP (UNIV MARYLAND)
EISELE1 69 ZPHY 221 1	+ENGELMANN, FILTHUTH, FOHLISCH, HEPP + (HEID)
EISELE2 69 ZPHY 223 487	EISELE, ENGELMANN, FILTHUTH, FOHLISCH + (HEID)
GERSHWIN 69 UCRL-19246	LAWRENCE KENNETH GERSHWIN (THEIS) (LRL)
BERLEY 70 PR D1 2015	+YAMIN, HERTZBACH, KOFLER + (BNL, MASA, YALE)
ROBERT 70 PR D2 6	+LUCAS, TAFT, WILLIS, BERLEY + (BNL, MASA, YALE)
EISELE 70 ZPHY 238 372	+FILTHUTH, HEPP, PRESSER, ZECH (HEIDELBERG)
BAKKE 71 LNC 1 37	+SABRE COLLAB. (ZEEM+SACL+BGNA+REHO+EPOL)
COLE 71 PR D4 631	+LEE-FRANZINI, LOVELESS, BALTAY + (STON, COLU)
ALSO 69 NEVIS-175 THESIS	HERBERT NORTON
TOVAY 71 NP B33 493	LOUC, BELGRADE, BERL, BRUX, DUBLIN, WARS COLLAB
BALTAY 72 PR D5 1569	+FEINMAN, FRANZINI, NEWMAN, YEH + (COLU+STON)
BOHM 72 NP B48 1	BERLIN+BELGRADE-BRUX-DUBLIN+LOUC+WARSAW
ELLIS 72 NP B39 77	OXF+AERE-RHEL+LOOM+LYON+NWES+ITEP COLLABOR
FRANZINI 72 PR D6 2417	COLUMBIA+HEIDELBERG+MARYLAND+STONY BROOK
ROBERTSD 72 THESIS	R.M. ROBERTSON (IIT)
EBENHON 73 ZPHY 264 413	+EISELE, FILTHUTH, HEPP, LEITNER, THOUW + (HEID)
FOX 73 PRL 31 1084	+LAM, BARNES, EISENSTEIN + (BNL, VPI+WILL+HYOM)
SCHIZOR 73 PR D8 12	B. SECHI-ZORN, G. SNOW (JMD)
EBENHON 74 ZPHY 266 367	+EISELE, ENGELMANN, FILTHUTH, HEPP + (HEID)
ROBERTS 74 PRL 32 1265	WILL+VPI+CARN+WYOM+CIT COLLABORATION
ALSO 74 PRL 33 122	ERRATUM TO ROBERTS 74
ALSO 75 PR D12 1232	ROBERTS, COX + (WILL+VPI+CARN+WYOM+CIT+BNL)
TANENBAU 74 PRL 33 175	TANENBAUM, HUNGERBUUEHLER + (YALE+FNAL+BNL)
ALSO 75 TANENBAUM	
DUGAN 75 NP A254 396	+ASANO, CHEN, CHENG, HU, LIDOFSKY + (COLU+YALE)
TANENBAU 75 PR D12 1871	TANENBAUM, HUNGERBUUEHLER + (YALE+FNAL+BNL)
CONFORTO 76 NP 8105 189	+GOPAL, KALMUS, LITCHFIELD, ROSS + (RHEL+LOIC)
DECAMP 77 PL 668 495	+BADIER, BLAND, CHOLLET, GAILLARD, (CALO-EPOL)
HANSU 78 NP 8132 45	+MANZ, MATT, REUCROFT, SETTLES + (MP+M+VAND)
HERBERT 78 PRL 40 1250	+CLELAND, COOPER, DRIS, ENGLS + (PITT+BNL)
MARRAFFI 80 PR D21 2501	MARRAFFINO, REUCROFT, ROOS, WATERS + (VAND+MP+IN)
THOMPSON 80 PR D21 25	+CLELAND, COOPER, DRIS, ENGLS + (PITT+BNL)
BOURQUIN 82 ZPHY C12 307	+BROWN + (BRIS-GEVA-HEID+LALO-RL+STRB)
ALSO 83 BOURQUIN	
KELLER 82 PRL 48 971	+LESNIK, ROMANOWSKI, KEIG + (OSU+CHIC+ANL)
BOURQUIN 83 ZPHY C21 17	BOURQUIN + (BRIS+GEVA-HEID+LALO+RL+STRB)
BOURQUIN 83 ZPHY C21 27	BOURQUIN + (BRIS+GEVA-HEID+LALO+RL+STRB)





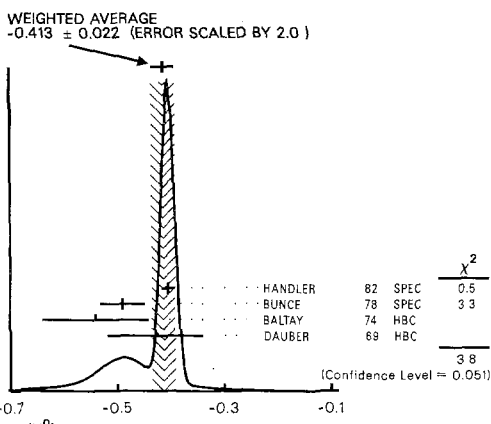




For notation, see key on page 91.

# Stable Particle Full Listings

$\Xi^0, \Omega^-$



$\phi$ ANGLE (degrees)	( $\tan \phi = \beta/\gamma$ )	Experiment	Count	Chi-squared	
146	-8.	BERGE	66	HBC	SEE NOTE D BELOW
739	38.	DAUBER	69	HBC	SEE NOTE A BELOW
652	16.0	BALTAY	74	HBC	1.75 GEV/C K-

USED ALPHA LAMBDA = 0.647 -- 0.020.  
ERRORS MULTIPLIED BY 1.2 DUE TO APPROXIMATIONS USED FOR XI  
POLARIZATION. (SEE DAUBER 69 FOR DETAILED DISCUSSION)

### REFERENCES FOR $\Xi^0$

ALVAREZ 59 PRL 2 215	ALVAREZ, EBERHARD, GOOD, GRAZIANO, TICHO+ (LRL)
JAUNEAU 63 SIENA CONF 1 1	JAUNEAU+ (EPOL+CERN+LOUC-RHEL+BERGEN)
ALSO 63 PL 4 49	JAUNEAU+ (EPOL+CERN+LOUC-RHEL+BERGEN)
TICHO 63 BNL CONF 410	HAROLD K TICHO (UCLA)
CARMONY 64 PRL 12 482	CARMONY, PJERROU, SCHLEIN, SLATER, STORK+(UCLA)
HUBBARD 64 PR 135 B 183	HUBBARD, BERGE, KALBFLEISCH, SHAFER+ (LRL)
PJERROU 65 PRL 14 275	+ CONLEIN, SLATER, SMITH, STORK, TICHO (UCLA)
PJERROU 65 THESIS	G M PJERROU (UCLA)
BERGE 66 PR 147 945	BERGE, EBERHARD, HUBBARD, MERRILL+ (LRL)
HUBBARD 66 UCL 11510	J RICHARD HUBBARD (THESTS, BERKELEY) (LRL)
LONDON 66 PR 143 1034	LONDON, RAU, GOLDBERG, LICHTMAN+(BNL+SYRACUSE)
PALMER 68 PL 268 323	PALMER, RADDUJIC, RAU, RICHARDSON+ (BNL, SYR)
DAUBER 69 PR 179 1262	+BERGE, HUBBARD, MERRILL, MILLER (LRL)
MAYEUR 72 NP 847 333	+VAN BINST, WILQUET+ (BRUX+CERN-TUFT-LOUC)
ALSO 75 NP 853 268	ERRATUM TO MAYEUR 72
WILQUET 72 PL 428 372	+FLIAGNE, SUX, KNIGHT+ (BRUX-CERN-TUFT-LOUC)
BALTAY 74 PR D9 49	+BRIDGEWATER, COOPER, GERSHWIN+ (COLU-BING)J
YEH 74 PR D10 3545	+GAIAGALAS, SMITH, ZENOLE, BALTAY+ (BING-COLU)
GEMENIGE 75 PL 378 193	GEMENIGE, BJESDAL, PRESSER+ (CERN+HEID)
ZECH 77 NP 8124 413	+DYDAK, NAVARRIA+ (SIEG+CERN-DORT+HEID)
BUNCE 78 PR D18 633	+HANDLER, MARCH, MARTIN+ (WISC-MICH+RUTG)
SUNCE 79 PL 868 384	+OVERSETH, COX, DWORKIN+ (BNL-MICH-RUTG+WISC)
COX 81 PRL 46 877	+DWORKIN+ (MICH+WISC-RUTG+MINN+BNL)
HANDLER 82 PR D25 639	+GROBEL, PONDROM+ (WISC-MICH-MINN+RUTG)

## STRANGENESS -3 BARYON

$\Omega^-$

$$I(J^P) = 0(\frac{3}{2}^+)$$

QUANTUM NUMBERS ASSIGNED FROM SU3

### $\Omega^-$ MASS (MeV)

Experiment	Count	Mass (MeV)	Chi-squared
EISENBERG	54	EMUL	
FRY1	55	EMUL	
FRY2	55	EMUL	
ABRAMS	64	HBC	INTO XI- P10
PALMER	68	HBC	K-P 4.6, 5. GEV/C
SCHULTZ	68	HBC	K-P 5.5 GEV/C
SCOTTER	68	HBC	K-P 6. GEV/C
SPEITH	69	HBC	K-P 10. GEV/C
ABCLV	73	HBC	K-P 10. GEV/C
DIBIANCA	75	DBC	4.9 GEV/C K-D
BAUBILLIE	78	HBC	8.25 GEV/C K-P
HEMINGWAY	78	HBC	4.2 GEV/C K-P
HARTOUNI	85	SPEC	80-280 GEV KL C

AVG 1672.43 0.32 AVERAGE

THE UNAMBIGUOUS DISCOVERY OF OMEGA- IN BOTH PRODUCTION AND DECAY MODES OCCURRED IN BARNES 1 64.  
EISENBERG 54 MASS CALCULATED FOR DECAY IN FLIGHT. ALVAREZ 73 HAS SHOWN THAT THE OMEGA INTERACTED WITH AG NUCLEUS TO GIVE K- XI AG. BOTH FRY EVENTS IDENTIFIED AS OMEGA- BY ALVAREZ 73.  
FRY MASSES ASSUME DECAY TO LAMBDA K- AT REST. DECAY FROM ATOMIC ORBIT COULD DOPPLER SHIFT THE K- ENERGY AND RESULTING OMEGA- MASS BY SEVERAL MEV FOR FRY 2. THIS SHIFT IS NEGLIGIBLE FOR FRY 1 BECAUSE THE OMEGA DECAY IS APPROXIMATELY PERPENDICULAR TO ITS ORBITAL VELOCITY, AS IS KNOWN BECAUSE THE LAMBDA STRIKES THE NUCLEUS (L. ALVAREZ, PRIVATE COMM. 1973). WE HAVE CALCULATED THE ERROR ASSUMING THAT ORBITAL N IS 4 OR LARGER.  
ABCLV VALUE INCLUDES THE SPEITH 69 EVENTS. EXCLUDED FROM AVERAGE.  
SEE NOTE D IN THE OMEGA- MEAN LIFE SECTION BELOW.  
DIBIANCA 75 GIVES MASS FOR EACH EVENT. WE QUOTE AVERAGE.

### $\Omega^-$ MASS (MeV)

Experiment	Count	Mass (MeV)	Chi-squared	Notes
FIRESTONE	71	HBC	1.0	12 GEV/C K+D
HARTOUNI	85	SPEC	1.0	80-280 GEV KL C
AVG	1672.55	0.71	AVERAGE	

### $\Omega^-$ MEAN LIFE (units $10^{-10}$ sec)

Experiment	Count	Mean Life	Chi-squared	Notes
BAUBILLIE	78	HBC	0.80	8.25 GEV/C K-P
DEUTSCHMA	78	HBC	0.12	10, 16 GEV/C K-P
HEMINGWAY	78	HBC	0.15	4.2 GEV/C K-P
BOURQUIN	84	SPEC	0.75	CERN SPS HYPERON BM
DEUTSCHMANN	78	EXCLUDED	0.823	FROM AVERAGE BECAUSE OF SIGNIFICANT DISAGREEMENT WITH OTHER RECENT EXPERIMENTS, POSSIBLY DUE TO XI- CONTAMINATION.
AVG	0.822	0.013	0.013	AVERAGE

### $\Omega^-$ PARTIAL DECAY MODES

Mode	Decay Masses
$\Omega^- \rightarrow \Lambda K^-$	1116+ 494
$\Omega^- \rightarrow \Xi^0 \pi^-$	1315+ 140
$\Omega^- \rightarrow \Xi^- \pi^0$	1321+ 135
$\Omega^- \rightarrow \Lambda \pi^-$	1116+ 140
$\Omega^- \rightarrow \Xi^- \gamma$	1321+ 0
$\Omega^- \rightarrow \Xi(1530)^0 \pi^-$	1533+ 140
$\Omega^- \rightarrow \Xi^0 e^- \nu$	1315+ .511+ 0
$\Omega^- \rightarrow \Xi^- \pi^+ \pi^-$	1321+ 140+ 140

### $\Omega^-$ BRANCHING RATIOS

Mode	Branching Ratio	Notes
$\Omega^- \rightarrow \Lambda K^-$	0.678	0.007
$\Omega^- \rightarrow \Xi^0 \pi^-$	0.236	0.007
$\Omega^- \rightarrow \Xi^- \pi^0$	0.086	0.004
$\Omega^- \rightarrow (\Lambda \pi^-)/total$ (units $10^{-3}$ )	0.19	OR LESS CL=.90
$\Omega^- \rightarrow (\Xi^- \gamma)/total$ (units $10^{-3}$ )	2.2	OR LESS CL=.90
$\Omega^- \rightarrow (\Xi(1530)^0 \pi^-)/total$ (units $10^{-3}$ )	0.64	0.51
$\Omega^- \rightarrow (\Xi^0 e^- \nu)/total$ (units $10^{-2}$ )	0.56	0.28
$\Omega^- \rightarrow (\Xi^- \pi^+ \pi^-)/total$ (units $10^{-4}$ )	4.3	3.4

### $\Omega^-$ DECAY PARAMETERS

SEE NOTE ON BARYON DECAY PARAMETERS IN NEUTRON SECTION ABOVE.

Experiment	Count	Mass (MeV)	Chi-squared	Notes
KOCHER	74	HBC	0.36	10 GEV/C K-P
BAUBILLIE	78	HBC	0.50	8.25 GEV/C K-P
HEMINGWAY	78	HBC	0.50	4.2 GEV/C K-P
BOURQUIN	84	SPEC	0.028	CERN SPS HYPERON BM
AVG	1672.43	0.32	AVERAGE	

Experiment	Count	Mass (MeV)	Chi-squared	Notes
BOURQUIN	84	SPEC	0.09	0.14
AVG	614	0.05	0.21	



For notation, see key on page 91.

# Stable Particle Full Listings

$\Lambda_c^+$ ,  $\Xi_c^+$ ,  $\Omega_c^0$ ,  $\Lambda_b^0$

$\Lambda_c^+ \rightarrow (e^+ \text{ anything})/\text{total}$  (P12)

R17	0.045	0.017	VELLA	82	SMK2 E+ E-	4.5-6.8	GEV
-----	-------	-------	-------	----	------------	---------	-----

$\Lambda_c^+ \rightarrow (p e^+ \text{ anything})/\text{total}$  (P13)

R18 M	0.018	0.009	VELLA	82	SMK2 E+ E-	4.5-6.8	GEV
R18 M	THIS INCLUDES PROTONS FROM LAMBDA DECAY.						

$\Lambda_c^+ \rightarrow (\Lambda e^+ \text{ anything})/\text{total}$  (P14)

R19	1	(0.022)OR LESS	CL-.90	BALLAGH	81	HYBR NU WE-H2	IN 15-FT
R19 N	0.011	0.008	VELLA	82	SMK2 E+ E-	4.5-6.8	GEV
R19 N	THIS INCLUDES LAMBDA'S FROM SIGMA0 DECAY.						

$\Lambda_c^+ \rightarrow (\Lambda e^+ \text{ anything})/(\Lambda \text{ anything})$  (P14)/(P16)

R20 S	0.027	0.017	SON	82	DBC NU D	IN FNAL 15-FT	
R20 V	0.065	0.041	VELLA	82	SMK2 E+ E-	4.5-6.8	GEV
R20 S	SON 82 USES OWN DATA AND MU- E+ LAMBDA EVENTS OF MURTAGH 79.						
R20 V	THIS VALUE DEDUCED BY SON 82 FROM DATA OF VELLA 82.						
R20							
R20 AVG	0.033	0.016	AVERAGE				

### REFERENCES FOR $\Lambda_c^+$

CAZZOLI 75 PRL 34 1125	+CNOPS, CONNOLLY, LOUITT, MURTAGH, + (BNL)
KNAPP 76 PRL 37 882	+LEE, LEUNG, SMITH + (COLU+HAWA+ILL+FNAL)
BARISH 77 PR D15 1	+DERRICK, DOMBECK, MUSGRAVE (ANL+PURD)
ANGELINI 79 PL 848 150	(ANKA+L1BH+CERN+DUUC+LOUC+KEYN+PISA+ROMA+)
SALTAY 79 PRL 42 1721	+CAROUBALIS, FRENCH, HIBBS, + (COLU+BNL)
CNOPS 79 PRL 42 197	+CONNOLLY, KAHN, KIRK, MURTAGH, PALMER+ (BNL)
DRIJARD 79 PL 858 452	+FISCHER+ (CERN+CDEF+DORT+HEID+LAPP+WARS)
GIBONI 79 PL 858 437	+DIBITONTO+ (AACH+CERN+HARV+MUNI+NWES+UCR)
KERNAN 79 LEPTON CONF. FNAL A	KERNAN (UCR)
LOCKMAN 79 PL 858 443	+MEYER, RANDER, SCHLEIN, WEBB+ (UCLA+SACL)
MURTAGH 79 FERMILAB. SYMP. 277	M. J. MURTAGH (FNAL)
ABRAMS 80 PRL 44 10	+ALAM, BLOCKER, BOYARSKI, + (SLAC+LBL)
ALLASIA 80 NP B176 13	(ANKA+L1BH+CERN+DUUC+LOUC+KEYN+PISA+ROMA+)
CALICCHI 80 PL 938 521	+ (BARI+BIRM+BRUX+CERN+EPOL+RHEL+SACL+LOUC)
KITAGAKI 80 PRL 45 955	+TANAKA, YUTA, ABE, + (TOHO+IIT+UMD+STON+TUFT)
WEISS 80 TORONTO CONF 319	J M WEISS (SLAC)
BALLAGH 81 PR D24 7	+BINGHAM + (LBL+UCB+FNAL+HAWA+WASH+WISC)
BASILE 81 NC 62A 14	-CARA ROMEO + (CERN+BGNA+PGIA+FRAS)
IRION 81 PL 998 495	+SEEBRUNNER, + (AACH+CERN+HARV+MUNI+NWES)
RUSSELL 81 PRL 46 799	+AVERY, BUTLER, GLADDING + (ILL+FNAL+COLU)
BOSSETTI 82 PL 1098 234	+GRASSLER, + (AACH+BNL+CERN+MPIM+DFP)
KITAGAKI 82 PRL 48 299	+TANAKA, YUTA, ABE, + (TOHO+IIT+UMD+STON+TUFT)
SON 82 PRL 49 1128	+SNOW, CHANG, KUMORI + (UMD+IIT+STON+TOHO+TUFT)
VELLA 82 PRL 48 1515	+TRILLING, ABRAMS, ALAM, + (SLAC+LBL+UCB)
SON 83 PR D28 2129	+SNOW, CHANG, KUMORI + (UMD+IIT+STON+TOHO+TUFT)
USHIDA 83 PRL 51 2362	(AICH+FNAL+KOE+SEUL+MCGI-NAGO+OSU+OKAY+)
ADAMOVIĆ 84 PL 140B 119	ADAMOVIĆ, WA58 C. (BGNA+CERN+FRZ+GENO+)
ALEEV 84 ZPHY C23 333	BIS-2 COLLAB(BERL+JINR+LEBD+MSU+PRAG+SOFI)
BOGWICK 85 PRL 55 923	CLEO C. (HARV+OHIO+ROCH+RUTG+SYRA+VAND+)

THEORY AND REVIEW

DERUJULA 75 PR D12 147	+GEORGI, GLASHOW (HARV)
GAISSER 76 PR D14 3153	T. K. GAISSER, F. HALZEN (BART+WISC)
LEE 77 PR D15 157	+QUIGG, ROSNER (FNAL)
MULLER 79 CERN/EP 79-148	F. MULLER (CARGESE LEC. 1979)
DIBITONTO 81 MADISON CONF.	D. DIBITONTO (CERN)
TRILLING 81 PRPL 75 57	G H TRILLING (LBL)

$\Xi_c^+$  was  $A^+$

$$I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$$

A NARROW SIGNAL (WIDTH COMPATIBLE WITH THE 23-MEV RESOLUTION) INTERPRETED AS A STABLE CHARMED STRANGE BARYON (QUARK CONTENT CSU). THIS INTERPRETATION NEEDS CONFIRMATION.

BIAG12 85 LOOK FOR THE ISOSPIN PARTNER  $\Xi_1^0$  IN LAMBDA K-  $\pi^+$  AND OTHER CHANNELS WITHOUT SUCCESS.

### $\Xi_c^+$ MASS (MeV)

M	82 2460.	25.	BJAGI	83	SPEC	SIG- BE-->XI/C+ X
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### $\Xi_c^+$ MEAN LIFE (units $10^{-13}$ sec)

T	53	4.8	2.9	1.8	BIAG11	85	SPEC	SIG- BE-->XI/C+ X
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### $\Xi_c^+$ PARTIAL DECAY MODES

P1	$\Xi_c^+ \rightarrow \Lambda K^- \pi^+ \pi^+$	DECAY MASSES	1116+ 494+ 140- 140
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### $\Xi_c^+$ BRANCHING RATIOS

$\Xi_c^+ \rightarrow \Lambda K^- \pi^+ \pi^+$  (P1)

R1 A	82	SEEN	BIAGI	83	SPEC	SIG- BE-->XI/C+ X
R1 A	BIAG12 85 LOOK FOR BUT DO NOT SEE THE $\Xi_1^0$ IN P K- K0BAR $\pi^+$ (BRANCHING FRACTION < 0.08 WITH 90 PERCENT CL), P 2K- 2P1- (<0.03, R1 A 90 PERCENT CL), OMEGA- K+ $\pi^+$ , LAMBDA K <sup>0</sup> $\pi^+$ , AND SIGMA(1385)+ K- $\pi^+$ .					

### REFERENCES FOR $\Xi_c^+$

BIAGI	83	PL 122B 455	+ (BRIS+CERN+GEVA+HEID+LAUS+LOQM+MELB+RL)
BIAG11	85	PL 150B 230	+ (BRIS+CERN+GEVA+HEID+LAUS+LOQM+MELB+RAL)
BIAG12	85	ZPHY C28 175	+ (BRIS+CERN+GEVA+HEID+LAUS+LOQM+MELB+RAL)

$\Omega_c^0$  was  $T^0$

$$I(J^P) = ( )$$

A CLUSTER OF THREE  $\Xi_1^-$  K-  $\pi^+$   $\pi^+$  EVENTS. THE OMEGA/CO -  $\Xi_1^0$  MASS DIFFERENCE IS 280-10 MEV. THE INTERPRETATION AS BEING THE OMEGA/CO NEEDS CONFIRMATION.

OMITTED FROM SUMMARY TABLE

### $\Omega_c^0$ MASS (MeV)

M	3 2740.	20.	BIAGI	85	SPEC	SIG- BE-->
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### REFERENCES FOR $\Omega_c^0$

BIAGI	85	ZPHY C28 175	+ (BRIS+CERN+GEVA+HEID+LAUS+LOQM+MELB+RAL)
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## BOTTOM BARYON

$\Lambda_b^0$

$$I(J^P) = 0(\frac{1}{2}^+)$$

QUANTUM NUMBERS NOT ESTABLISHED, TAKEN FROM QUARK MODEL. THE CLAIM BY BASILE 81 IS HOTLY DISPUTED BY DRIJARD 82. BASILE 82 IS THE REPLY, AND DRIJARD1 82 IS THE REPLY TO THAT.

OMITTED FROM SUMMARY TABLE

### $\Lambda_b^0$ MASS (MeV)

M	5425.0	175.0	75.0	BASILE	81	SFM	0 P P	62	GEV	ECM
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### $\Lambda_b^0$ PARTIAL DECAY MODES

P1	$\Lambda_b^0 \rightarrow p D^0 \pi^-$	DECAY MASSES	938+1865- 140
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### $\Lambda_b^0$ BRANCHING RATIOS

$\Lambda_b^0 \rightarrow (p D^0 \pi^-)/\text{total}$  (P1)

R1	SEEN	BASILE	81	SFM	DO TO K- $\pi^+$
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### REFERENCES FOR $\Lambda_b^0$

BASILE	81	LNC 31 97	+BONVICINI, CARA ROMEO+(CERN-BGNA+FRAS+PGIA)
BASILE	82	NC 68A 289	+BONVICINI, CARA ROMEO+ (CERN-BGNA+FRAS)
DRIJARD	82	PL 108B 361	+FISCHER,+ (CERN+CDEF+DORT+HEID+LAPP+WARS)
DRIJARD1	82	CERN/EP 82-31	+FISCHER,+ (CERN+CDEF+DORT+HEID+LAPP+WARS)



For notation, see key on page 91.

Stable Particle Full Listings
FREE QUARK SEARCHES

C I - BOUND TO NUCLEI.
M FOR X-SECT READ FRACTION OF FRAGMENTS.
G - FOR X-SECTION READ X-SECT(Q-Q)/X-SECT(MU-MU)
F - 3E-5 < LIFETIME < 1E-3 S.
C E - X-SECTION CM2/GEV2
D - INCLUDES BOTT 72 RESULTS.
C C - HADRONIC OR LEPTONIC QUARKS.
C B - ASSUMES ISOTROPIC CM PRODUCTION.
C A - CROSS SECTION INFERRED FROM FLUX.

QUARK DIFFERENTIAL PROD. CROSS SECT. --- ACCELER. SEARCHES

Table with columns: QUARK, X-SECT, CHARGE, MASS, REFERENCE, DET, BEAM. Rows include experimental results for various quarks like u, d, s, c, b.

QUARK FLUX --- ACCELERATOR SEARCHES

Table with columns: QUARK, FLUX, CHARGE, MASS, REFERENCE, DET, BEAM. Rows include flux measurements for various quarks.

FR E - FOR FLUX READ QUARK PROD. X-SECTION RATIO TO CSE+D-MU+MU
FR D - QUARK LIFETIMES > 1E-8 S.
FR C - ONE CANDIDATE M < 17 GEV.
FR B - HADRONIC QUARK.
FR A - LEPTONIC QUARK.
FR \* - QUARK FLUX PER CHARGED PARTICLE.

QUARK FLUX --- COSMIC RAY SEARCHES

Table with columns: QUARK, FLUX, CHARGE, MASS, REFERENCE, DET, SHIELDING. Rows include cosmic ray search results for various quarks.

F - LIFETIME > 10\*\*8 S; CHARGE --, 0, .68, .42; AND MASS
F D - 4, 4, AND 20 GEV, RESPECTIVELY.
F E - ALSO 1/4 AND 1/6E CHARGES.
F D - NO EVENTS IN SUBSEQUENT EXPTS.
F C - LEPTONIC QUARKS.
F B - PROMPT AIR SHOWER SEARCH.
F A - TIME DELAYED AIR SHOWER SEARCH.
F \* - ALTITUDE IN KM. ALL OTHERS SEA LEVEL.

QUARK DENSITY --- MATTER SEARCHES

Table with columns: QUARK, DENSITY, CHARGE, MASS, REFERENCE, MATERIAL/METHOD. Rows include matter search results for various quarks.

RHO B - ALSO SET LIMITS FOR Q=+1/6E.
RHO A - LIMIT INFERRED BY JONES 77 (RMP 69, 717).

REFERENCES FOR QUARK SEARCHES

List of references for quark searches, including authors, journals, and years.





For notation, see key on page 91.

Stable Particle Full Listings
MAGNETIC MONOPOLE SEARCHES, AXION SEARCHES

F F - USED DKMPR MECHANISM AND PENNING EFFECT.
F D - ASSUMES MONOPOLE ATTACHES FERMION NUCLEUS.
F D - ANOMALOUS LONG-RANGE ALPHA TRACKS.
F C CATALYSIS OF NUCLEON DECAY.
F B - REEVALUATES PARKER 70 LIMIT FOR GUT MONOPOLES.
F A - ALVAREZ 75, FLEISCHER 75, FRIEDLANDER 75, ROSS 76 EXPLAIN AS FRAGMENTING NUCLEUS. EBERHARD 75 DISCUSSES CONFLICT WITH OTHER EXP'TS. HAGSTROM 77 REINTERPRETS AS ANTINEUTRONS. PRICE 78 REASSESSSES.

FLEISCHS 69 PR 184 1398
ALSO 70 JAP 41 958
CARTHERS 66 PR 149 1070
AMALDI 63 NC 28 773
GOTO 63 PR 132 387
PETUKHOV 63 NP 49 87
PURCELL 63 PR 129 2326
FIDECARD 61 NC 22 657
BRADNER 59 PR 114 603
MALKUS 51 PR 83 899
FLEISCHER, PRICE, WOODS
MART, JACOBS, PRICE, SCHWARTZ, WOODS (GESC)
CARTHERS, STEFANSKI, ADAIR (VALF-BNL)
BARONI, MANFREDINI, BRADNER + (ROMA-UCSD-BRAN)
KOLM, FORD (TOKY-MIT-BRAN)
YAKIMENKO (LEBD)
COLLINS, FUJII, HORNOSTEL, TURKOT (HARV-BNL)
FINOCCHIARO, GIACOMELLI (CERN)
ISBEL (LSBL)
MALKUS (CMB)

MONOPOLE DENSITY -- MATTER SEARCHES

Table with columns: D, MONOPOLE EVENTS, DENSITY, CHARGE, REFERENCE, DET, MATERIAL. Contains search results for iron ore, sun, earth heat, moon rock, moon wake, manganesite, meteorite, etc.

CRAYEN 85 FERMILAB-85/13
RUZICKA 80 JINR 2-80-850
GIACOMEL 84 RNC 7 12 1
GIACOMELLI (BGNA)

AXION SEARCHES

OMITTED FROM SUMMARY TABLE

VARIOUS AXION SEARCHES IN PRODUCTION OR DECAY

AX (VALUES QUOTED ARE AXION PRODUCTION RATIO TO PID PROD CROSS SEC.)
AX FOR THEORY AND REVIEW SEE WEINBERG PR 40, 223 (1978); WILCZEK
AX PRL 40, 279 (1978); DONNELLY PR D18, 1607 (1978); BARSHAY ET AL.,
PRL 46, 1361 (1981); BAROSO-MUKHOPADHYAY PL 106, 91 (1981); AND
PECEI, PROC. INTL. NEUTRINO CONF., MAUI (1981), ED. CENCE (HAWAII).
AX FOR INVISIBLE-AXION THEORY SEE E.G. WISE, GEORGI AND GLASHOW
AX PRL 47, 402 (1981); AND FOR VERY HEAVY AXION THEORY SEE E.G. TYE
PRL 47, 1035 (1981). FOR K --> PI AD SEE ALSO E.G., GOLDMAN
AND HOFFMAN, PRL 40, 220 (1978), FRERE ET AL., PL 103B, 129 (1981).

REFERENCES FOR MAGNETIC MONOPOLE SEARCHES

ARAFUNE 85 PR D32 2586
BERMON 85 PR 15 1850
BRACCI 85 PR B258 726
ALSO 85 LNC 42 123
CAPLIN 85 NAT 317 234
ERISU 85 NP 11 885
PARK 85 NP B252 61
FRYBERGE 84 PR D29 1524
HARVEY 84 NP B236 255
INCANDIEL 84 PR 55 2067
KAJINO 84 PRL 52 909
KAJINO2 84 JPE 10 447
KAWAGOE 84 LNC 41 315
KRISHNAS 84 PL 1428 99
LISS 84 PR D30 884
PRICE1 84 PRL 52 1265
PRICE2 84 PL 140B 112
TARLE 84 PRL 52 909
ANDERSON 83 PR D28 2308
ARAFUNE 83 PL 1338 380
AUBERT 83 PL 120B 465
BARWICK 83 PR D28 2336
BARTELT 83 PRL 50 655
BATTISTONI 83 PL 1338 454
BONARELLI 83 PL 126B 157
BOSETTI 83 PL 1338 245
CARRERA 83 PRL 51 1933
DOKE 83 PL 129B 370
ERREDE 83 PRL 51 243
FREES 83 PRL 50 655
GROOM 83 PRL 50 655
MASHIMO 83 PL 128B 327
MIKHAILO 83 PL 130B 351
MUSSET 83 PL 126B 157
REPHAELI 83 PL 121B 115
SCHATTEN 83 PR D27 1525
ALEXEYEV 82 LNC 35 413
BONARELLI 82 PL 112B 100
CARRERA 82 PRL 48 1378
DIMOPOUL 82 PL 119B 320
DELL 82 NP B209 44
KINOSHIT 82 PRL 52 909
KOLB 82 PRL 49 1373
MASHIMO 82 JPSJ 51 3067
TURNER 82 PR D26 1296
BARTLETT 81 PR D24 612
BONNARDEAU 81 PR D23 323
KINOSHIT 81 PR D24 1707
ULLMAN 81 PRL 47 289
CARRIGAN 80 NAT 288 348
BRODERICK 79 PR D19 1046
BARTLETT 78 PR D18 2553
CARRIGAN 78 PR D17 1754
HOFFMANN 78 LNC 42 1183
PRICE 78 PR D18 1382
HAGSTROM 77 PRL 38 729
CARRIGAN 76 PR D17 1823
DELL 76 LNC 15 269
ROSS 76 LBL-4665
STEVENS 76 PR D14 2207
ZRELOV 76 JETP 36 1506
ALVAREZ 75 LBL-4260
BURKE 75 PL 60B 113
CARRERA 75 PHD THESIS
CARRIGAN 75 NP 810 279
ALSO 71 PR D3 56
EBERHARD 75 LBL-4289
EBERHARD 75 PR D11 3099
FLEISCHER 75 PR D12 3512
FRIEDLANDER 75 PRL 35 1167
GIACOMEL 75 NC 28A 21
PRICE 75 PRL 35 487
CARRIGAN 74 PR D10 2567
CARRIGAN 73 PR D8 3717
ROSS 73 PR D8 698
ALSO 71 PR D4 3260
ROSS 70 SI 167 701
BARTLETT 72 PR D6 1817
GUREVICH 72 PL 38B 549
ALSO 72 JETP 34 917
ALSO 71 SI 319 3940
FLEISCHER 71 PR D4 24
KOLM 71 PR D4 1285
PARKER 70 APJ 160 383
SCHATTEN 70 PR D1 2245
FLEISCHER 69 PR 177 2029
FLEISCHER 69 PR 184 1393

AXION PROD. IN HADR. COLLISIONS + BEAM DUMP EXPTS.

AXP LIMITS ARE FOR C(SAO)/CS(PI0).
AXP 1. E-8 OR LESS CL=90 ALBRAN 78 HYBR BEAM DUMP
AXP 6. E-9 OR LESS CL=95 ASRATYAN 78 CALO BEAM DUMP
AXP A 5.4E-14 OR LESS CL=90 BELLOTTI 78 HLBC FOR MASS=1.5 MEV
AXP A 4.1E-9 OR LESS CL=90 BELLOTTI 78 HLBC FOR MASS=1 MEV
AXP A 1.5E-8 OR LESS CL=90 BELLOTTI 78 HLBC BEAM DUMP
AXP B 1. E-8 OR LESS CL=90 BOSETTI 78 HYBR BEAM DUMP
AXP C DONNELLY 78
AXP D HANSLI 78 WIRE BEAM DUMP
AXP E MICELMAC 78
AXP F VYSOTSSKII 78
AXP G BECHIS 79 CNTR
AXP H COTEUS 79 OSKP BEAM DUMP
AXP I DISHAW 79 CALO 600 GEV P P
AXP J 1. E-8 OR LESS CL=90 FAISSNER 80 OSKP BEAM DUMP, AO--E-E-
AXP J 1. E-8 OR LESS CL=90 JACQUES 80 HLBC 28 GEV PROTONS
AXP J 1. E-14 OR LESS CL=90 JACQUES 80 HLBC BEAM DUMP
AXP K SOUKAS 80 CALO 285V P B-DUMP
AXP L 12 FAISSNER1 81 OSKP CERN PS NU WIDEBAN
AXP M 15 FAISSNER2 81 OSKP BEAM DUMP, AO--26AM
AXP N 8 KIM 81 OSKP 26 GEV P NUC--AO X
AXP O FAISSNER 82 RVUE CF FAISSNER2 81
AXP P 24 FAISSNER 83 OSKP BEAM DUMP, AO--26AM
AXP Q FAISSNER2 83 RVUE LAMP BEAM DUMP
AXP R FRANK 83 RVUE LAMP BEAM DUMP
AXP S HOFFMAN 83 CNTR PL--N(AO--E-E-)
AXP T 0 2. E-11 OR LESS CL=90 BERGSM 85 CHRM CERN BEAM DUMP
AXP T 0 1. E-13 OR LESS CL=90 BERGSM 85 CHRM CERN BEAM DUMP
AXP A BELLOTTI 78 FIRST VALUE COMES FROM SEARCH FOR A--E--E-, SECOND
AXP A VALUE COMES FROM SEARCH FOR A--2GAMMA, ASSUMING MASS=2 MASS(E-).
AXP A FOR ANY MASS SATISFYING THIS, LIMIT IS ABOVE VALUE (MASS\*\*4).
AXP A THIRD VALUE COMES FROM PL 60B 401 AND QUOTES CS(PROD)\*CS(INTER-
AXP A ACTION)=10\*\*67 CM\*\*4.
AXP B BOSETTI 78 QUOTES CS(PROD)\*CS(INTERACT)> 2.E-67 CM\*\*4
AXP C DONNELLY 78 EXAMINES DATA FROM REACTOR NEUTRON DECAYS OF REINES 76
AXP C AND GURR 74 AS WELL AS SLAC BEAM DUMP EXPT. EVIDENCE IS NEGATIVE.
AXP D MICELMACHER 78 FINDS NO EVIDENCE OF AXION EXISTENCE IN REACTOR
AXP D EXPTS OF REINES 76 AND GURR 74. (SEE REF UNDER DONNELLY 78 BELOW).
AXP E VYSOTSSKII 78 DERIVED LOWER LIMIT FOR THE AXION MASS. 25 KEV FROM
AXP E LUMINOUSITY OF THE SUN AND 200 KEV FROM RED SUPERGIANTS.
AXP F BECHIS 79 LOOKED FOR THE AXION PRODUCTION IN LOW ENERGY ELECTRON
AXP F BREMSSTRAHLUNG AND THE SUBSEQUENT DECAY INTO EITHER 2GAMMAS OR
AXP F E--E-, NO SIGNAL FOUND. C.L.=0.90 LIMITS FOR MODEL PARAMETERS(S)
AXP F ARE GIVEN.
AXP G COTEUS 79 IS A BEAM DUMP EXPERIMENT AT BNL.
AXP H DISHAW 79 IS A CALORIMETRIC EXPERIMENT AND LOOKS FOR LOW ENERGY
AXP H TAIL OF ENERGY DISTRIBUTIONS DUE TO ENERGY LOST TO WEAKLY
AXP H INTERACTING PARTICLES.
AXP I FAISSNER 81 IS SIN BEAM DUMP EXPT WITH 590 MEV PROTONS LOOKING FOR
AXP I AO-->E+ E- DECAY. ASSUMING AO/PI0=5E-7, OBTAINED DECAY RATE LIMIT
AXP I 20/(AO MASS) MEV/SEC (CL=90) WHICH IS ABOUT 10\*\*7 BELOW THEORY
AXP I AND INTERPRETED AS UPPER LIMIT TO MASS(AO) < 2\* MASS(E-).
AXP J JACQUES 80 IS A BNL BEAM DUMP EXPT - FIRST LIMIT ABOVE COMES FROM
AXP J NON-OBSERVATION OF EXCESS NC-TYPE EVENTS (CS(PROD)\*CS(INTERACT) <
AXP J 7.E-68 CM\*\*4, CL=0.90) - SECOND LIMIT IS FROM NON-OBSERVATION OF
AXP J AXION DECAYS INTO TWO GAMMAS OR E-E-, AND FOR AXION MASS A FEW MEV.
AXP K SOUKAS 80 AT BNL OBSERVED NO EXCESS OF NC-TYPE EVENTS IN BEAM DUMP.
AXP L FAISSNER1 81 SEE EXCESS MU E EVENTS. SUGGEST AXION INTERACTIONS.
AXP M FAISSNER2 81 IS SIN 90MEV NEUTRON BEAM DUMP. OBSERVED 14.5+-5.0 EVS
AXP M OF 2GAMMA DECAY OF LONG-LIVED NEUTRAL PENETRATING PARTICLE WITH
AXP M N(2GAMMA) < APPROX. 1 MEV. AXION INTERPR. WITH ETA-AO MIXING GIVES
AXP M M(AO)=(250+-25)KEV, TAUC(2GAMMA)=(7.3+-3.7)E-8 SEC FROM ABOVE RATE.
AXP M SEE CRITICAL REMARKS BELOW IN COMMENTS OF FLEISCHER 82, FAISSNER 83,
AXP M FAISSNER2 83, FRANK 83, AND BERGSM 85, ALSO SEE IN THE NEXT
AXP M SUBSECTION ALEKSEEV 82, CAVAINAG 83, AND ANANEV 85.



For notation, see key on page 91.

# Stable Particle Full Listings

## SUPERSYMMETRIC PARTICLE SEARCHES

SE A BRANDELIK 82 LIMIT IS FROM NO ENHANCEMENT IN EVENTS WITH THETA NEAR  
SE A 90 DEG(LIGHT SCALAR) AND WITH LARGE ACOPLANARITY (HEAVY SCALAR).

SE B GLADNEY 83 LOOKED FOR LARGE PT E FROM SINGLY PRODUCED S-ELECTRONS.

SE C BARTEL 84 MAKE NON-STANDARD ASSUMPTION THAT PHOTINO DECAYS TO  
SE C GOLDSTINO +GAMMA. THEY LOOK FOR 26AM EVTS FROM PHOTINO PAIR PROD.  
SE C FOR SUPERSYM BREAKING PARAM D=(100 GEV)\*\*2, THEY FIND AT CL=.95  
SE C M(S-ELECTRON) > 80 GEV. LIMIT IS ALSO APPLICABLE IF THE PHOTINO  
SE C DECAYS RADIATIVELY WITHIN THE DETECTOR.

SE D BARTEL 84 LIMIT IS FROM NEUTRALINO SEARCH ASSUMING THAT  
SE D MASS(NEUTRALINO) < 30 GEV. USING ACOPLANARITY EVENTS WITH MISSING PT.  
SE D UNDER THESE CONDITIONS, M(S-ELECTRON) > 50 GEV AT CL=.90.

SE E FERNANDEZ 84 ANALYZED SINGLE ELECTRON EVENTS FROM SINGLY PRODUCED  
SE E S-ELECTRON. ENERGY DISTRIBUTION IS CONSISTENT WITH E+GAMMA BKGD.

SE F BARTEL 85 CONSIDER SINGLE AND PAIR S-ELECTRON PRODUCTION AS WELL AS  
SE F PHOTINO PAIR+GAMMA. FIRST LIMIT IS FOR MASSLESS PHOTINO AND SECOND  
SE F IS FOR M(PHOTINO) < 15 GEV. BOTH ASSUME EQUAL MASS FOR LEFT AND  
SE F RIGHT S-ELECTRON. FOR M(S(E)) >> M(S(E,L)) AND MASSLESS PHOTINO,  
SE F LIMIT IS 21.8 GEV (SEE ALSO THEIR FIG.2).

SE G FERNANDEZ 85 ANALYZE SINGLE GAMMA EVENT EXPECTED FROM E-E ->>  
SE G PHOTINO PAIR+GAMMA ABOVE LIMIT IS FOR M(PHOTINO)=0 AND EQUAL  
SE G MASS LEFT AND RIGHT S-E. (SEE ALSO FIG.2). FOR M(RIGHT) >>  
SE G M(LEFT), LIMIT IS 30 GEV.

SE H BARTHA 86 SEARCH FOR ANOMALOUS PHOTONS FROM E-E ->>  
SE H (PHOTINO PHOTINO PHOTON) WHERE PHOTINOS ARE ASSUMED TO LEAVE  
SE H APPARATUS UNDETECTED. FOR M(RIGHT) >> M(LEFT) AND M(PHOTINO)=0,  
SE H LIMIT IS 42 GEV AT CL=.90.

### SCALAR $\mu$ MASS LIMIT

SMU NONE 3GEV TO 15GEV CL=.95 BARBER 80 MRKJ E-E-  
SMU A NONE 3.3GEV TO 16GEV CL=.95 BEHREND 82 CELL E-E-  
SMU A NONE BELOW 16.4GEV CL=.95 BRANDELIK 82 TASS E-E-  
SMU B NONE BELOW 13.8GEV CL=.95 FERNANDEZ 83 MKC E-E-  
SMU B NONE BELOW 20 GEV CL=.95 ADEVA 85 MRKJ E-E- ECM=40-47 GEV  
SMU C NONE BELOW 21. GEV CL=.95 BARTEL 85 JADE E-E- ECM=32-47 GEV  
SMU C NONE BELOW 20.9GEV CL=.95 BARTEL 85 JADE E-E- STABLE S-MU

SMU A BRANDELIK 82 LIMIT IS FROM NO ENHANCEMENT IN EVENTS WITH THETA NEAR  
SMU A 90 DEG(LIGHT SCALAR) AND WITH LARGE ACOPLANARITY (HEAVY SCALAR).

SMU B FERNANDEZ 83 AND ADEVA 85 OBSERVED NO EXCESS ACOPLANAR MU+MU-  
SMU B EVENTS.

SMU C BARTEL 85 SECOND LIMIT IS FOR STABLE S-MU FROM NON-OBSERVATION OF  
SMU C SMU-PAIR PROD. FIRST LIMIT IS FOR S-MU ->> MU + STABLE-PHOTINO  
SMU C FROM NON-OBSERVATION OF ACOPLANAR MU PAIR WITH LARGE MISSING PT  
SMU C AND APPLIES IF M(PHOTINO) < 15 GEV. FIRST LIMIT ASSUMES M(LEFT)=  
SMU C M(RIGHT); IF M(LEFT) >> M(RIGHT) THE LIMIT IS 20.3 GEV.

### SCALAR $\tau$ MASS LIMIT

STA SCALAR-TAU->>TAU PHOTINO LOOKS IDENTICAL TO CHGD-HIGGS->>TAU NU  
STA AS IN STABLE PARTICLE SEARCH SECTION. WERE TAKEN FROM RESULTS  
STA QUOTED FOR CHGD-HIGGS, THE LIMITS BELOW CORRESPOND TO  
STA BR(TAU NU)=1.0 LINE IN MASS-LIMIT GRAPHS FOR PAIR OF  
STA CGHD-HIGGS->>(TAU NU) (TAU NU).

STA 14GEV OR MORE CL=.95 ADEVA 82 MRKJ  
STA NONE 46EV TO 14GEV CL=.95 BARTEL 82 JADE  
STA A NONE 66EV TO 15.3GEV CL=.95 BEHREND 82 CELL  
STA A NONE M(TAU) TO 5.8GEV CL=.95 BEHREND 82 CELL  
STA NONE M(TAU) TO 9.9GEV CL=.90 BLOCKER 82 SMKZ  
STA B NONE BELOW 17 GEV CL=.95 ADEVA 85 MRKJ E-E- ECM=40-47 GEV

STA A BEHREND 82 FIRST LIMIT FOR SCALAR TAU IS FROM PT-CUT TAU-PAIR  
STA A ANALYSIS, SECOND LIMIT IS FROM NO EXCESS TAU PAIR EVENTS.

STA B NO EXCESS ACOPLANAR MU AND HADRONIC JET.

### SCALAR QUARK MASS LIMIT

SGK A NONE BELOW 3 GEV NAPPI 82 RVUE Q-2/3 SQUARK ONLY  
SGK B 1300GEV OR LESS HINCHLIFF 82 RVUE M(GLUINO)=2 GEV  
SGK B 3100GEV OR LESS HINCHLIFF 82 RVUE M(GLUINO)=4 GEV  
SGK C NONE BELOW 100 GEV BERGSM 83 RVUE FOR M(GLUINO)=2 GEV  
SGK D NONE BELOW 40 GEV BARGER 84 RVUE P PBAR ECM=540 GEV  
SGK E NONE BELOW 40 GEV ELLIS2 84 RVUE M(GLNO) >> M(SQK)  
SGK F NONE BELOW 25 GEV BAER 85 RVUE P PBAR  
SGK G NONE BELOW 70 GEV REYA 85 RVUE M(GLNO) >> M(SQK)  
SGK G NONE BELOW 70 GEV CL=.90 BARNETT 86 RVUE FOR M(GLNO)= 80 GEV  
SGK G NONE BELOW 60 GEV CL=.90 BARNETT 86 RVUE FOR M(GLNO)= 80 GEV

SGK A NAPPI 82 LIMIT APPLIES TO CHARGE 2/3 SQUARK. NO LIMIT FOUND FOR  
SGK A CHARGE 1/3. LIMIT FROM P-WAVE SQUARK-SQUARK BOUND STATE NON-  
SGK A OBSERVATION IN E-E ANNIHILATION.

SGK B ABSTRACTED FROM DISCUSSION OF BEAM CONTAMINATION EXPER AT FNAL  
SGK B IN HINCHLIFF 82. ASSUMES GLUINO->> QUARK QUARK PHOTINO.

SGK C REANALYSIS OF CERN-SPS BEAM DUMP DATA. SEE THEIR FIG. 1.

SGK D BARGER 84 USE MONOJET + MISSING-PT EVENTS OBSERVED IN UA1  
SGK D EXPERIMENT AT ECM=546 GEV AS UPPER LIMIT DUE TO SQUARK ->>  
SGK D QUARK+PHOTINO ASSUMING M(GLUINO)=2\*M(SQUARK).

SGK E ELLIS2 84 AND REYA 85 USE UA1 MONOJET+MISSING-PT EVENTS AT ECM=546  
SGK E GEV. ASSUME SQUARK ->> QUARK+PHOTINO.

SGK F BAER 85 EVALUATE THE CONTRIBUTION OF SQUARK->>QUARK +Z-INO, W-INO.  
SGK F NO CHANGE IS NEEDED OF LIMIT ON M(GLUINO), M(SQUARK) DERIVED  
SGK F ASSUMING THE DOMINANCE OF SQUARK->> QUARK+PHOTINO.

SGK G BARNETT 86 USE THE UA1 MONOJET+MISSING-PT EVENTS AT 650 GEV.  
SGK G THEY CONSIDER ALL SUPERSYMMETRIC PROCESSES CONSISTENT WITH  
SGK G M(GLUINO) > M(PHOTINO) AND NEGLECT DECAYS TO W-INOS AND Z-INOS.  
SGK G THEY CONCLUDE UA1 MONOJETS CANNOT ORIGINATE IN GLUINO OR  
SGK G SCALAR QUARK PRODUCTION.

### GLUINO MASS LIMIT

GNO THERE IS AN ONGOING CONTROVERSY (REFLECTED IN THESE LISTINGS) ABOUT  
GNO WHETHER VERY LIGHT GLUINOS ARE RULED OUT. THESE PAPERS SOMETIMES  
GNO MAKE DIFFERENT ASSUMPTIONS AND USE DIFFERENT CALCULATIONAL  
GNO TECHNIQUES.

GNO A NONE BELOW 2.3GEV KANE 82 RVUE BEAM DUMP  
GNO B NONE BELOW 2 GEV BERGSM 83 RVUE FOR M(S-QK)<100 GEV  
GNO C BALL 84 CALO IF PHOTINO DECAYS  
GNO D FARRAR 84 RVUE  
GNO E NONE BELOW 3.76EV CL=.68 LLEWELLYN 84 RVUE FOR M(S-QK)=30 GEV  
GNO F NONE BELOW 60 GEV ALLAN 85 RVUE FOR M(SQ)=20-35 GEV  
GNO G NONE BETWEEN 15-40 GEV ALTARELLI 85 RVUE FOR M(SQ)=100 GEV  
GNO H BAER 85 RVUE P PBAR  
GNO I NONE BETWEEN 0.5-4 GEV COOPER-SA 85 8DMP FOR M(SQ)< 65 GEV  
GNO I NONE BETWEEN 0.5-3 GEV COOPER-SA 85 8DMP FOR M(SQ)=150 GEV  
GNO I NONE BETWEEN 0.5-2 GEV COOPER-SA 85 8DMP FOR M(SQ)=300 GEV  
GNO J NONE BETWEEN 2-4 GEV DAWSON 85 RVUE TAU > E-7 SEC  
GNO J NONE BETWEEN 1-2.5 GEV DAWSON 85 RVUE FOR M(SQ)=100 GEV  
GNO K DERUJULA 85 RVUE FOR M(SQ)=80 GEV  
GNO L NONE BELOW 40 GEV ELLIS 85 RVUE P PBAR ECM=546 GEV  
GNO M NONE BETWEEN .5-4.1GEV CL=.90 FARRAR 85 RVUE FNAL BEAM DUMP  
GNO N NONE BELOW 1 GEV GOLDMAN 85 RVUE GLUONIUM  
GNO P NONE BELOW 1-2 GEV HABER 85 RVUE  
GNO Q NONE BELOW 25 GEV REYA 85 RVUE M(SQK) >> M(GLNO)  
GNO R BARGER 86 RVUE FOR M(SQ)= 100 GEV  
GNO S NONE BETWEEN 3-62 GEV CL=.90 BARNETT 86 RVUE FOR M(SQ)= 75 GEV  
GNO S NONE BETWEEN 3-50 GEV CL=.90 BARNETT 86 RVUE FOR M(SQ)= 75 GEV

GNO A KANE 82 INFERED ABOVE GLUINO MASS LIMIT FROM RETROACTIVE ANALYSIS  
GNO A OF HADRONIC COLLISION AND BEAM DUMP EXPERIMENTS. LIMITS VALID IF  
GNO A GLUINO DECAYS INSIDE DETECTOR.

GNO B BERGSM 83 IS REANALYSIS OF CERN-SPS BEAM-DUMP DATA. SEE THEIR  
GNO B FIG. 1.

GNO C BALL 84 IS FNAL BEAM DUMP EXPERIMENT. OBSERVED NO INTERACTIONS OF  
GNO C PHOTINO IN THE CALORIMETER, WHERE PHOTINOS ARE EXPECTED TO COME  
GNO C FROM PAIR PROD. SEARCH FOR LONG LIFETIME PHOTINO  
GNO C INTERACTING IN CALORIMETER 5CM FROM TARGET. LIMIT IS FOR M(S-QUARK)  
GNO C =40 GEV AND PROD CROSS SECTION PROPORTIONAL TO A\*\*0.72. BALL 84  
GNO C FIND NO GLUINO ALLOWED BELOW 4.1 GEV AT CL=.90. THEIR FIG.1 SHOWS  
GNO C DEPENDENCE ON M(S-QUARK) AND A. SEE ALSO KANE 82 (PHYS. LETT.  
GNO C 112B, 227).

GNO D FARRAR 84 ARGUES THAT M(GLUINO) < 100 MEV IS NOT RULED OUT IF THE  
GNO D LIGHTEST R-HADRONS ARE LONG-LIVED. A LONG LIFETIME WOULD OCCUR IF  
GNO D R-HADRONS ARE LIGHTER THAN PHOTINOS OR IF M(SQUARK) > 100 GEV.  
GNO E FROM BEBC DATA.

GNO F ALLAN 85 LIMIT IS FROM CONSIDERING Q-GLUON->> SQ-GLUINO FOLLOWED BY  
GNO F SQ->> Q-PHOTINO AND GLUINO->> Q+SQ->> Q+QBAR-PHOTINO ASSUMING  
GNO F M(GLUINO)=M(SQUARK). LIMIT IS FOR SQUARK MASS BETWEEN 20 GEV TO 35  
GNO F GEV, AND USED THE UA1 MONOJET DATA AT ECM=546 GEV.

GNO G ALTARELLI 85 SAY THAT M(GLUINO)<15 GEV CANNOT BE EXCLUDED BY UA1  
GNO G MONOJET DATA AT 546 GEV DUE TO UNCERTAINTY IN FRAGMENTATION OF  
GNO G LIGHT GLUINOS. THEY ANALYZE GLUINO-GLUINO PRODUCTION.

GNO H BAER 85 EVALUATE THE CONTRIBUTION OF SQUARK->>QUARK +Z-INO, W-INO.  
GNO H NO CHANGE IS NEEDED OF LIMIT ON M(GLUINO), M(SQUARK) DERIVED  
GNO H ASSUMING THE DOMINANCE OF SQUARK->> QUARK+PHOTINO.

GNO I COOPER-SARKAR 85 IS BEBC BEAM-DUMP. GLUINOS DECAYING IN DUMP WOULD  
GNO I YIELD PHOTINOS IN THE DETECTOR GIVING NEUTRAL-CURRENT-LIKE  
GNO I INTERACTIONS. FOR M(SQUARK) > 330 GEV, NO LIMIT IS SET.

GNO J DAWSON 85 FIRST LIMIT FROM NEUTRAL PARTICLE SEARCH. SECOND LIMIT  
GNO J BASED ON FNAL BEAM DUMP EXPERIMENT.

GNO K DERUJULA 85 SAY THAT M(GLUINO)=3 GEV CANNOT BE EXCLUDED BY UA1  
GNO K MONOJET DATA AT 546 GEV DUE TO UNCERTAINTY IN FRAGMENTATION OF  
GNO K LIGHT GLUINOS. THEY ANALYZE GLUINO-GLUINO PRODUCTION.

GNO L ELLIS 85 USE UA1 MONOJET+MISSING-PT EVENTS AT ECM=546 GEV AS UPPER  
GNO L LIMIT FOR GLUINO PAIR PRODUCTION WITH GLUINO ->> Q QBAR PHOTINO.  
GNO L GLUINO WITH MASS AROUND 3 GEV TO 20 GEV SEEMS EXCLUDED EVEN AFTER  
GNO L THE EFFECT OF VARIOUS REFINEMENTS ARE INCLUDED.

GNO M FARRAR 85 POINTS OUT THAT BALL 84 ANALYSIS APPLIES ONLY IF THE  
GNO M GLUINOS DECAY BEFORE INTERACTING, I.E. M(SQUARK)>50\*M(GLUINO)\*\*1.5.  
GNO M FARRAR 85 FINDS M(GLUINO)<0.5 NOT EXCLUDED FOR M(SQ)=30-1000 GEV  
GNO M AND M(GLUINO)<1.0 NOT EXCLUDED FOR M(SQ)=100-500 GEV BY BALL 84  
GNO M EXPERIMENT.

GNO N GOLDMAN 85 USE NON-OBSERVATION OF A PSEUDOSCALAR GLUINO  
GNO N BOUND STATE IN RADIATIVE PSI DECAY.

GNO P HABER 85 IS BASED ON SURVEY OF ALL PREVIOUS SEARCHES SENSITIVE TO  
GNO P LOW MASS GLUINOS. LIMIT MAKES ASSUMPTIONS REGARDING THE LIFETIME  
GNO P AND ELECTRIC CHARGE OF THE LIGHTEST SUPERSYMMETRIC PARTICLE.

GNO Q REYA 85 USE UA1 MONOJET+MISSING-PT EVENTS AT ECM=546 GEV. ASSUME  
GNO Q GLUINO ->> QUARK+QUARK-PHOTINO.

GNO R BARGER 86 SAY THAT M(GLUINO)=3-5 GEV CANNOT BE EXCLUDED BY UA1  
GNO R MONOJET DATA AT 546 GEV DUE TO UNCERTAINTY IN FRAGMENTATION OF  
GNO R LIGHT GLUINOS. THEY ANALYZE GLUINO-GLUINO PRODUCTION. THEY  
GNO R ALSO ANALYZE THE GLUINO DISTRIBUTION FUNCTION AND ITS ROLE IN  
GNO R GLUINO PRODUCTION.

GNO S BARNETT 86 USES INCREASED STATISTICS OF UA1 MONOJET RESULTS AT  
GNO S ECM=630 GEV PLUS A HIGHER-ORDER PROCESS TO RULE OUT LIGHT GLUINOS  
GNO S WITH M(GLUINO)=3-5 GEV.

GNO S BARNETT 86 CONSIDER ALL SUPERSYMMETRIC PROCESSES CONSISTENT WITH  
GNO S M(GLUINO) > M(PHOTINO) AND NEGLECT DECAYS TO W-INOS AND Z-INOS.  
GNO S THEY CONCLUDE UA1 MONOJETS CANNOT ORIGINATE IN GLUINO OR  
GNO S SCALAR QUARK PRODUCTION.





# Stable Particle Full Listings

## OTHER STABLE PARTICLE SEARCHES

### OBS. OF EVENTS WITH LARGE MISSING $p_T$ IN COLL. EXPTS

MPT A 7 EVENTS ARNISON 84 UA1 P AP, ECM=540 GEV  
MPT B 4 EVENTS BAGNAIA 84 UA2 P AP, ECM=540 GEV  
MPT A ARNISON 84 OBSERVED 7 EVENTS WITH MISSING  $p_T > 40$  GEV, 5 ASSOCIATED  
MPT A WITH A NARROW HADRONIC JET, 2 WITH NEUTRAL EM CLUSTER. EXCEPT ONE  
MPT A OF THE LATTER EVENTS, NONE OF THEM CAN BE INTERPRETED BY  $W \rightarrow e \nu$   
MPT A AND ARNISON 84 SAYS THEY ARE UNLIKELY TO BE STANDARD BACKGROUND.  
MPT A THESE HAVE BEEN REFERRED TO AS MONOJET EVENTS. PRELIMINARY UA1  
MPT A RESULTS (REPORTED AT 1985 CONFERENCES) WITH A LARGER INTEGRATED  
MPT A LUMINOSITY INDICATE THAT MANY OR POSSIBLY ALL OF THE MONOJET  
MPT A EVENTS ARE DUE TO THE SUM OF SEVERAL DIFFERENT STANDARD MODEL  
MPT A BACKGROUNDS.  
MPT B BAGNAIA 84 OBSERVED 4 EVENTS WITH MISSING  $p_T > 25$  GEV, WITH JET(S)  
MPT B OF  $p_T > 30$  GEV, AND WITH ELECTRON OF  $p_T > 15$  GEV. PRELIMINARY UA2  
MPT B RESULTS (REPORTED AT 1985 CONFERENCES) WITH A LARGER INTEGRATED  
MPT B LUMINOSITY CONTAIN NO ADDITIONAL EVENTS.

### HIGHLY IONIZING PARTICLE FLUX (number/m<sup>2</sup>-year)

IDN 0 0.4 OR LESS CL-.95 KINOSHITA 81 PLAS Z/BETA 30-100

### TACHYON FLUX IN COSMIC RAYS (number/cm<sup>2</sup>-sec- $\sigma$ )

TCF SEE SMITH 77 FOR A REVIEW OF EARLIER COSMIC RAY AND ACCELERATOR  
EXPERIMENTS.  
TCF A NONE PRESCOTT 76 CNTR  
TCF B NONE SMITH 77 CNTR  
TCF C 0 2.3E-10 OR LESS CL-.95 BHAT 79 CNTR  
TCF D ? 2.4E-9 OR LESS CL-.90 MARINI 82 CNTR V/C>1.2  
TCF A PRESCOTT 76 REANALYZED CLAY AND CROUCH(C.C.C.)74 DATA(NATURE 248,28).  
TCF A FOUND APPARATUS EFFECT, CORRECTION FOR WHICH MUCH REDUCES THE STAT.  
TCF A SIGNIFICANCE OF POSITIVE C.G. RESULT. ALSO PERFORMED TWO NEW EXPTS  
TCF A ONE USING C.C.C. APPARATUS, ANOTHER WITH NEW APPARATUS. SET UPPER LIMIT  
TCF A AT CL-.95 OF ABOUT 30 TACHYONS PER SHOWER WITH AVERAGE SIZE  $N=6E+5$ .  
TCF B SMITH 77 ANALYZED MORE THAN 20000 SHOWERS(223 DAYS) WITH  $E=1E+14$   
TCF B EV SCANNING 290 E-6 SEC PERIOD BEFORE EACH SHOWER. OBSERVED EXCESS  
TCF B 46--40 EVS DOES NOT CONSTITUTE STATISTICALLY SIGNIFICANT EVIDENCE.  
TCF C BHAT 79 IS AT OOTACAMUND(2200M ABOVE SEA). NO SIGNAL IN 3621 HOURS.  
TCF D MARINI 82 IS TOF MEASUREMENT USING PEP-CNTR AT SEA LEVEL.

TCM TACHYON SEARCHES IN  $E-E$  ANNIHILATION  
TCM A 0 1. E-6 OR LESS CL-.90 PEREPILIT 77 CNTR  $V^*(EQ) < 1$   
TCM A 0 1. E-5 OR LESS CL-.90 PEREPILIT 77 CNTR  $1 < V^*(EQ) < 15$   
TCM A PEREPILIT 77 IS MICHELSON-TYPE EXPT FOR PAIR-PRODUCED TACHYONS  
TCM A IN  $E+E$  ANNIHILATION ( $E+$  FROM CU ISOTOPE). ABOVE LIMITS ARE FOR  
TCM A CS( $E+E \rightarrow$  TACHYON PAIR)/CS( $E+E \rightarrow$  2 $\gamma$ ), AND  $V^*(EQ)$  IS TACHYON  
TCM A VELOCITIES TIMES EARTH EQUATOR COMPONENT OF VELOCITY OF PREFERRED  
TCM A REFERENCE FRAME.

TCO SEARCHES FOR TACHYONIC DECAY (LOWER LIMIT FOR MEAN LIFE IN YEARS)  
TCO A SEE LJUBICIC 75 FIG. 1 FOR REVIEW OF EARLIER EXPERIMENTS.  
TCO A 4.8E-15 OR MORE LJUBICIC 75 ELEC M(TACHYON) $> 1.1$  KEV  
TCO A LJUBICIC 75 USED LEAD OXIDE CATHODE AND ELECTRON MULTIPLIER LOOKING  
TCO A FOR IONIZATION DUE TO TACHYONIC DECAY (SPONTANEOUS ACQUISITION OF  
TCO A ENERGY OF BOUND-STATE  $E+$ , SENSITIVE TO PROPER TACHYON MASS $> 1$  KEV.  
TCO A ABOVE LIMIT IS OBTAINED FROM OBSERVED  $E-$  EMISSION RATE 3/HOUR.

### FLAVOR-CHANGING AXIONLIKE PARTICLE (FAMILON=FA)

FCA A DICUS 83 RVUE COSMOLOGY  
FCA A THE PRIMORDIAL HEAVY NEUTRINO MUST DECAY INTO NU AND FA EARLY SO  
FCA A THAT THE RED-SHIFTED DECAY PRODUCTS ARE BELOW CRITICAL DENSITY, SEE  
FCA A THEIR TABLE. IN ADDITION,  $K \rightarrow \pi F_A$  AND  $\mu \rightarrow e F_A$  ARE UNSEEN. COM-  
FCA A BINGING THESE EXCLUDES M(HEAVY NU) BETWEEN 5E-5 TO 5E-4 MEV(MU DECAY)  
FCA A AND M(HEAVY NU) BETWEEN 5E-5 TO 0.1 MEV(K-DECAY).

### PROD. OF NEW PENETRATING NON-NU LIKE STATES

BD A IN BEAM DUMP  
BD A LOSECCO 81 CALO 28 GEV PROTONS  
BD A NO EXCESS N.C. EVS LEADS TO CSP\*CS\*ACCEPTANCE=2.26E-71 CM\*\*4/NUC\*\*2  
BD A (CL-.90) FOR LIGHT NEUTRALS. ACC. DEPENDS ON MODELS (0.1 TO 4.E-4).

### HEAVY PARTICLE PROD. CROSS SECTION IN $e^+e^-$

EE (RATIO TO CS( $E+E \rightarrow$  MU-MU)). SEE ALSO IN QUARK SEARCH AND MAGNETIC  
EE MONOPOLE SEARCHES.  
EE A 0 5.0 E-2 OR LESS CL-.90 BARTEL 80 JADE Q=(3,4,5)/3 2-12GEV  
EE B 0 1.6 E-2 OR LESS CL-.95 KINOSHITA 82 PLAS Q=3-180, M<14.5GEV  
EE A BARTEL 80 IS DESY PETRA EXPT WITH WCM=27-35 GEV. ABOVE LIMIT IS FOR  
EE A INCLUSIVE PAIR PROD AND RANGES BETWEEN 1.E-1 - 1.E-2 DEPENDING ON  
EE A MASS AND PROD MOMENTUM DIST. (SEE THEIR FIGS.9,10,11).  
EE B KINOSHITA 82 IS SLAC PEP EXPT AT WCM=29 GEV USING LEXAN AND CR-39  
EE B PLASTIC SHEETS SENSITIVE TO HIGHLY IONIZING PARTICLES.

### HEAVY PARTICLE PROD. CROSS SECTION (cm<sup>2</sup>)

CH A 0 1. E-31 OR LESS LEIPUNER 73 CNTR  $\rightarrow$  M=3-11 GEV  
CH B 0.3-1.3E-31 OR LESS CARROLL 78 SPEC M=2.-2.5 GEV  
CH A LEIPUNER 73 IS AN NAL 300 GEV P EXPT. WOULD HAVE DETECTED PARTICLES  
CH A WITH LIFETIME GREATER THAN 200 NSEC.  
CH B CARROLL 78 LOOK FOR NEUTRAL, S=2 DIHYPERON RESONANCE IN  
CH B P P  $\rightarrow$  2 $\pi$  X. CS VARIES WITHIN ABOVE LIMITS OVER MASS RANGE AND  
CH B PLAS=5.1-5.9 GEV/C.

### HEAVY PARTICLE PROD. CROSS SECTION (cm<sup>2</sup>/N)

CS A 0 2.5E-35 OR LESS GUSTAFSON 76 CNTR 0 TAU GT 10\*\*7  
CS A GUSTAFSON 76 IS A 300 GEV FNAL EXPT LOOKING FOR HEAVY (M GT 2 GEV)  
CS A LONGLIVED NEUTRAL HADRONS IN THE M4 NEUTRAL BEAM. THE ABOVE TYPICAL  
CS A VALUE IS FOR M=3 GEV AND ASSUMES AN INTERACTION CROSS SECTION OF  
CS A 1 MB. VALUES AS A FUNCTION OF MASS AND INTERACTION CROSS SECTION  
CS A ARE GIVEN IN FIG. 2.

### HEAVY PARTICLE PROD. DIFFERENTIAL CROSS SECTION (cm<sup>2</sup>/sr-GeV)

D A 0 1.5E-36 OR LESS DORFAN 65 CNTR 8E TARGET M=3-7GEV  
D A 0 3.0E-36 OR LESS DORFAN 65 CNTR FE TARGET M=5-7GEV  
D B 0 2.4E-35 OR LESS CL-.90 BINON 69 CNTR Q=- M=1-1.8 GEV  
D B 0 2.4E-35 OR LESS CL-.90 ANTIPOV1 71 CNTR Q=- M=1.2-1.7, 2.1-4  
D C 0 1.2E-35 OR LESS CL-.90 ANTIPOV2 71 CNTR Q=- M=2.2-2.8  
D D 0 5.8E-34 OR LESS CL-.90 ALPER 75 SPEC  $\rightarrow$  M=1.5-24 GEV  
D E 0 1. E-31 OR LESS CL-.90 APPEL 74 CNTR  $\rightarrow$  M=3.2-7.2 GEV  
D F 0 2.2E-33 OR LESS CL-.90 ALBROW 75 SPEC Q=-1 M=4-15 GEV  
D F 0 1.1E-33 OR LESS CL-.90 ALBROW 75 SPEC Q=-2 M=6-27 GEV  
D G 0 8. E-35 OR LESS CL-.90 JOVANOVIC 75 CNTR  $\rightarrow$  M=15-26 GEV  
D G 0 1.5E-34 OR LESS CL-.90 JOVANOVIC 75 CNTR Q=-2, M=3-10 GEV  
D G 0 6. E-35 OR LESS CL-.90 JOVANOVIC 75 CNTR Q=-2, M=10-26 GEV  
D H 0 2.6E-36 OR LESS CL-.90 BALDIN 76 CNTR Q=-1, M=2.1-9.4 GEV  
D A DORFAN 65 IS A 30 GEV/C P EXPT AT BNL. UNITS ARE PER GEV MOMENTUM  
D A PER NUCLEUS.  
D B ANTIPOV1 71 LIMIT INFERRED FROM FLUX RATIO. 70 GEV P EXPERIMENT.  
D C ANTIPOV2 71 IS FROM SAME 70 GEV P EXP. AS ANTIPOV1 71 AND BINON 69.  
D D ALPER 73 IS CERN ISR 26+26 GEV P+P EXPT.  $P>.9$  GEV,  $.2 < BETA < .65$ .  
D E APPEL 74 IS NAL 300 GEV P+W EXPERIMENT. STUDIES FORWARD PRODUCTION  
D E OF HEAVY (UP TO 24 GEV) CHARGED PARTICLES WITH MOMENTA 24-200GEV(-)  
D E AND 40-150GEV(+CHG). ABOVE TYPICAL VALUE IS FOR 75 GEV AND IS  
D E PER GEV MOMENTUM PER NUCLEON.  
D F ALBROW 75 IS A CERN ISR EXPT WITH ECM=53 GEV. THETA=40 MR. SEE  
D F FIG. 5 FOR MASS RANGES UP TO 35 GEV.  
D G JOVANOVIC 75 IS A CERN ISR 26+26 AND 15-15 GEV P-P EXPERIMENT.  
D G FIG. 4 COVERS RANGES  $Q=1/3$  TO 2 AND M=3 TO 26 GEV.  
D G VALUE IS PER GEV MOMENTUM.  
D H BALDIN 76 IS A 70 GEV SERP EXP. VALUE IS PER AL NUCLEUS AT  
D H THETA=0. FOR OTHER CHARGES IN RANGE -0.5 TO -3.0, CL-.90 LIMIT IS  
D H (2.6E-36)/ABS(CHARGE) FOR MASS RANGE (2.1 TO 9.4GEV)\*ABS(CHARGE).  
D H ASSUMES STABLE PARTICLE INTERACTING WITH MATTER AS DO ANTI-PROTONS.

### LONGLIVED HEAVY PARTICLE INVARIANT C.S. (cm<sup>2</sup>/GeV<sup>2</sup>/N)

ICH A 0 1.1E-37 OR LESS CL-.90 CUTTS 78 CNTR MASS=4-10 GEV  
ICH B 0 3.0E-37 OR LESS CL-.90 VIDAL 78 CNTR MASS=4.5-6 GEV  
ICH C 0 6. E-33 OR LESS CL-.90 ARMITAGE 79 SPEC M=1.87 GEV  
ICH C 0 1.5E-33 OR LESS CL-.90 ARMITAGE 79 SPEC M=1.5-3.0 GEV  
ICH D 0 BOZZOLI 79 CNTR Q=- (2/3, 1/4, 3/2)  
ICH E 0 2.5E-36 OR LESS CL-.90 THRON 85 CNTR Q=-1 M=4-12 GEV  
ICH E 1. 1E-35 OR LESS CL-.90 THRON 85 CNTR Q=-1 M=4-12 GEV  
ICH A CUTTS 78 IS P BE EXPT AT FNAL SENSITIVE TO PARTICLES OF TAU=5E-8SEC  
ICH A VALUE IS FOR  $-3 < X < 0$  AND  $PT=0.175$ .  
ICH B VIDAL 78 IS FNAL 400 GEV PROTON EXPT. VALUE IS FOR  $X=0$  AND  $PT=0$ .  
ICH B PUTS LIFETIME LIMIT OF  $< 5 \times 10^{-8}$  SEC ON PARTICLE IN THIS MASS RANGE  
ICH C ARMITAGE 79 IS CERN-ISR EXPT AT ECM=53 GEV. VALUE IS FOR  $X=0.1$  AND  
ICH C  $PT=0.15$ . OBSERVED PARTICLES AT M=1.87 GEV ARE FOUND ALL CONSISTENT  
ICH C WITH BEING ANTI-DEUTERONS.  
ICH D BOZZOLI 79 IS CERN-SPS 200 GEV P N EXPERIMENT. LOOKS FOR PARTICLE  
ICH D WITH TAU LARGER THAN  $10^{**}-8$  SEC. SEE THEIR FIG. 11-18 FOR PRODUCTION  
ICH D CROSS SECTION UPPER LIMITS VS MASS.  
ICH E THRON 85 IS FNAL 400 GEV PROTON EXP. MASS DETERMINED FROM MEASURED  
ICH E VELOCITY AND MOMENTUM. LIMITS ARE FOR TAU=  $3 \times 10^{**}-9$  SEC.

### LONGLIVED HEAVY PARTICLE PRODUCTION

RPI (CS(HEAVY PARTICLE)/CS(PION)) KD  
RPI A 0 BUSSIÈRE 80 CNTR Q=- (2/3, 1/4, 3/2)  
RPI A BUSSIÈRE 80 IS CERN-SPS EXPT WITH 200-240 GEV PROTONS ON BE AND AL  
RPI A TARGET. SEE THEIR FIG. 6-7 FOR CS RATIO VS MASS.

### PROD. AND CAPT. OF LONG-LIVED MASSIVE PARTICLES (cm<sup>2</sup>)

CA A 0 0.1-9E-36 OR LESS FRANKEL 74 CNTR TAU=1 TO 1000 HRS  
CA B 0 1.4-9E-36 OR LESS FRANKEL 75 CNTR TAU=50 MS TO 10 HRS  
CA C 0 2-20E-34 OR LESS ALEKSEEV 76 ELEC TAU=100 MS TO 1 DAY  
CA C 0 0.2-8E-34 OR LESS ALEKSEEV 76 ELEC TAU=5 MS TO 1 DAY  
CA A FRANKEL 74 LOOKS FOR PARTICLES PRODUCED IN THICK AL TARGETS BY  
CA A 300-400 GEV/C PROTONS.  
CA B FRANKEL 75 IS EXTENSION OF FRANKEL 74.  
CA C ALEKSEEV(1,2) 76 ARE 61-70 GEV P SERP EXPT. CS IS PER PB NUCLEUS.





# Meson Full Listings

$\pi^\pm, \pi^0, \eta, \rho(770)$

## S=0, C=0, B=0 MESON STATES

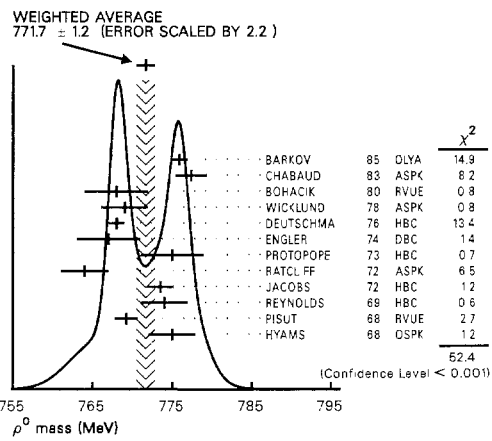
$\pi^\pm$	$I^G(J^PC) = 1^-(0^-)$
	SEE STABLE PARTICLE FULL LISTINGS
$\pi^0$	$I^G(J^PC) = 1^-(0^{++})$
	SEE STABLE PARTICLE FULL LISTINGS
$\eta$	$I^G(J^PC) = 0^-(0^{--})$
	SEE STABLE PARTICLE FULL LISTINGS
$\rho(770)$	$I^G(J^PC) = 1^+(1^{--})$
	OUR LATEST MINIREVIEW ON THIS PARTICLE CAN BE FOUND IN THE 1984 EDITION.

### $\rho(770)$ MASS (MeV)

WE NO LONGER LIST S-WAVE BREIT-WIGNER FITS, PBAR P DATA WITH HIGH COMBINATORIAL BACKGROUND, AND INSIGNIFICANT OR DOUBTFUL DATA.

CHARGED ONLY		NEUTRAL ONLY	
M2 R (760.0)	(9.0)	M2 P 4000	765. 5.0
M2 R (768.0)	(5.0)	M2 P 765.0	10.0
M2 R (765.0)	(5.0)	M2 P 775. 5.0	
M2 R (760.0)	(5.0)	M2 P 140K	767.7 1.9
M2 R (765.0)	(5.0)	M2 P 1930	767.0 4.0
M2 R 2775 (753.5)	(10.5)	M2 P 2430	770.0 2.7
M2 R (758.0)	(10.0)	M2 P 767.6	2.7
M2 R (749.0)	(3.0)	M2 AVG	766.1 1.3
M2 R (768.0)	(5.0)		
M2 R (773.0)	(2.0)		
M2 Z 900	767. 6.		
M2 A 9650	766.8 1.5		
M2 X 6500	766. 7.		
M2 AVG	766.8 1.4		
AVERAGE			
M2 R 300 (760.0)	(10.0)	M2 R 500 (770.0)	(10.0)
M2 R (760.0)	(5.0)	M2 R (770.0)	(5.0)
M2 R (775.0)	(5.0)	M2 R 4207 (758.0)	(7.5)
M2 R (770.0)	(5.0)	M2 R (765.0)	(8.0)
M2 R 4207 (758.0)	(7.5)	M2 R (760.0)	(5.0)
M2 R (765.0)	(8.0)	M2 R (768.0)	(2.0)
M2 R (770.0)	(5.0)	M2 R (761.0)	(3.0)
M2 R (775.0)	(2.0)	M2 R (770.0)	(4.0)
M2 R (768.0)	(2.0)	M2 R (775.0)	(2.0)
M2 H (778.0)	(2.0)	M2 R (768.4)	(2.4)
M2 H (770.0)	(9.0)	M2 H (778.0)	(2.0)
M2 C (775.0)	(2.0)	M2 H (770.0)	(9.0)
M2 D (776.3)	(0.4)	M2 C (775.0)	(2.0)
M2 G (776.1)	(2.6)	M2 D (776.3)	(0.4)
M2 CH (769.5)	(0.7)	M2 G (776.1)	(2.6)
M2 B (770.0)	(2.0)	M2 CH (769.5)	(0.7)
M2 D 2250	775.0 3.0	M2 B (770.0)	(2.0)
M2 A 13300	769.2 1.5	M2 D 2250	775.0 3.0
M2 C 1700	774.0 3.0	M2 A 13300	769.2 1.5
M2 Z 11200	773.5 1.7	M2 C 1700	774.0 3.0
M2 6800	764.0 3.0	M2 Z 11200	773.5 1.7
M2 C 32000	775.0 4.0	M2 6800	764.0 3.0
M2 4100	767. 4.	M2 C 32000	775.0 4.0
M2 76000	768.0 1.0	M2 4100	767. 4.
M2 X 769.0	3.0	M2 76000	768.0 1.0
M2 C H 768.0	4.0	M2 X 769.0	3.0
M2 G 777.4	2.0	M2 C H 768.0	4.0
M2 K 775.9	1.1	M2 G 777.4	2.0
M2 AVG	771.7 1.2	M2 K 775.9	1.1
AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.2) (SEE IDEOGRAM BELOW)			

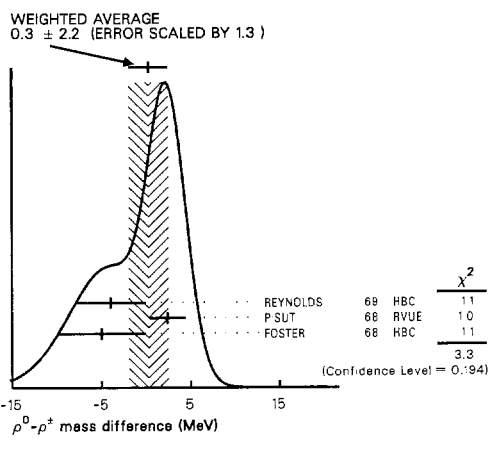
-----NOTES-----  
M A FROM FIT OF 3-PARAMETER RELATIVISTIC P-WAVE BREIT WIGNER TO TOTAL MASS DISTRIBUTION.  
M B HEYN 80 INCLUDES ALL SPACELIKE AND TIMELIKE F(PI) VALUES UNTIL 1978.  
M C FROM POLE EXTRAPOLATION  
M D ENERGY-DEPENDENT ANALYSIS OF BATON 70, HYAMS 73, PROTOPOESCU 73 PHASE SHIFTS.  
M E FROM FIT OF 3-PARAMETER RELATIVISTIC BREIT-WIGNER TO HELICITY ZERO PART OF P-WAVE INTENSITY. CHABAUD 83 AND BECKER 79 INCLUDE DATA OF GRAYER 74.  
M F FROM PHASE SHIFT ANALYSIS OF GRAYER 74 DATA.  
M G FROM PHOTOPRODUCTION.  
M H INCLUDED IN PISUT 68 RVUE  
M I PHASE SHIFT ANALYSIS. SYSTEMATIC ERRORS ADDED CORRESPONDING TO SPREAD OF DIFFERENT FITS.  
M J MASS ERRORS ENLARGED BY US TO WIDTH/SQRT(M), SEE K\* TYPED NOTE FROM THE GOUNARIS-SAKURAI PARAMETRIZATION OF THE PION FORM FACTOR



### $\rho^0 - \rho^\pm$ MASS DIFFERENCE (MeV)

D A 3600	-5. 5.	FOSTER	68 HBC	+0 PBAR P AT REST
D 22950	2.4 2.1	PISUT	68 RVUE	PI N TO RHO N
D A 3000	-4.0 4.0	REYNOLDS	69 HBC	-0 2.26 PI- P
D AVG	0.3 2.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3) (SEE IDEOGRAM BELOW)		

D A FROM QUOTED MASSES OF CHARGED AND NEUTRAL MODES



### $\rho(770)$ RANGE PARAMETER (GeV<sup>-1</sup>)

R	5.3	0.9	0.7	CHABAUD	83 ASPK	0 17 PI-P POLARIZ
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For notation, see key on page 91.

Meson Full Listings

ρ(770)

ρ(770) WIDTH (MeV)

Table listing meson widths in MeV. Columns include W, W2, W3, W4, W5, W6, W7, W8, W9, W10, W11, W12, W13, W14, W15, W16, W17, W18, W19, W20, W21, W22, W23, W24, W25, W26, W27, W28, W29, W30, W31, W32, W33, W34, W35, W36, W37, W38, W39, W40, W41, W42, W43, W44, W45, W46, W47, W48, W49, W50, W51, W52, W53, W54, W55, W56, W57, W58, W59, W60, W61, W62, W63, W64, W65, W66, W67, W68, W69, W70, W71, W72, W73, W74, W75, W76, W77, W78, W79, W80, W81, W82, W83, W84, W85, W86, W87, W88, W89, W90, W91, W92, W93, W94, W95, W96, W97, W98, W99, W100. Includes sub-sections for CHARGED ONLY, NEUTRAL ONLY, and AVERAGE.

ρ(770) BRANCHING RATIOS

Table listing branching ratios. Columns include R1, R2, R3, R4, R5, R6, R7, R8, R9, R10, R11, R12, R13, R14, R15, R16, R17, R18, R19, R20, R21, R22, R23, R24, R25, R26, R27, R28, R29, R30, R31, R32, R33, R34, R35, R36, R37, R38, R39, R40, R41, R42, R43, R44, R45, R46, R47, R48, R49, R50, R51, R52, R53, R54, R55, R56, R57, R58, R59, R60, R61, R62, R63, R64, R65, R66, R67, R68, R69, R70, R71, R72, R73, R74, R75, R76, R77, R78, R79, R80, R81, R82, R83, R84, R85, R86, R87, R88, R89, R90, R91, R92, R93, R94, R95, R96, R97, R98, R99, R100. Includes sub-sections for ρ → 4π/2π, ρ → (e+e-)/(π+π-), ρ → (π η)/(2π), ρ → (μ+μ-)/(π+π-), ρ → (π+π-π0)/(π+π-), and ρ → (η γ)/total.

ρ(770) PARTIAL DECAY MODES

Table listing partial decay modes. Columns include P1, P2, P3, P4, P5, P6, P7, P8 and decay masses.

ρ(770) PARTIAL WIDTHS (keV)

Table listing partial widths in keV. Columns include ρ → (π γ), W3, W5, W6, W7, W8, W9, W10, W11, W12, W13, W14, W15, W16, W17, W18, W19, W20, W21, W22, W23, W24, W25, W26, W27, W28, W29, W30, W31, W32, W33, W34, W35, W36, W37, W38, W39, W40, W41, W42, W43, W44, W45, W46, W47, W48, W49, W50, W51, W52, W53, W54, W55, W56, W57, W58, W59, W60, W61, W62, W63, W64, W65, W66, W67, W68, W69, W70, W71, W72, W73, W74, W75, W76, W77, W78, W79, W80, W81, W82, W83, W84, W85, W86, W87, W88, W89, W90, W91, W92, W93, W94, W95, W96, W97, W98, W99, W100.

REFERENCES FOR ρ(770)

List of references for ρ(770) studies, including authors, journals, and years.







For notation, see key on page 91.

# Meson Full Listings

$\eta'(958)$

<p><math>\eta' \rightarrow (\pi^+ \pi^- \gamma + \text{neutrals})/\text{total}</math> (P2C-P5)</p> <p>R4 EXCLUDING (P1- P1- ETA(NEUTRAL DECAY))/TOTAL R4 42 0.045 0.029 RITTENBER 69 HBC 1.7-2.7 K-P</p> <p>R4 FIT 0.0907 0.0069 FROM FIT</p>	<p><math>\eta' \rightarrow (\pi^+ \pi^- \gamma)/(\pi^+ \pi^- \eta(\text{neutral decay}))</math> (P3)/(P1N)</p> <p>R27 473 0.92 0.14 DANBURG 73 HBC 2.2 K-P, LAM X0 R27 192 1.11 0.18 JACOBS 73 HBC 2.9 K-P, LAM X0</p> <p>R27 AVG 0.99 0.11 AVERAGE R27 FIT 0.992 0.075 FROM FIT</p>
<p><math>\eta' \rightarrow \text{neutrals}/\text{total}</math> (P2N+P4)</p> <p>R5 123 0.189 0.026 RITTENBER 69 HBC 1.7-2.7 K-P R5 535 0.185 0.022 BASILE1 71 CNTR 1.6 PI- P, N X0</p> <p>R5 AVG 0.187 0.017 AVERAGE R5 FIT 0.182 0.014 FROM FIT</p>	<p><math>\eta' \rightarrow (2\gamma)/(\pi^0 \pi^0 \eta(\text{neutral decay}))</math> (P4)/(P2N)</p> <p>R28 16 0.188 0.058 APEL 72 OSPK 3.8 PI- P, N X0</p> <p>R28 FIT 0.115 0.015 FROM FIT</p>
<p><math>\eta' \rightarrow (\pi^+ \pi^- \gamma (\text{including } \rho \gamma))/\text{total}</math> (P3)</p> <p>R6 35 0.34 0.09 BADIER 65 HBC 3.0 K-P R6 20 0.2 0.1 LONDON 66 HBC 2.2 K-P R6 298 0.329 0.033 RITTENBER 69 HBC 1.7-2.7 K-P</p> <p>R6 AVG 0.319 0.030 AVERAGE R6 FIT 0.300 0.016 FROM FIT</p>	<p><math>\eta' \rightarrow (\gamma \mu^+ \mu^-)/(2\gamma)(\text{units } 10^{-3})</math> (P20)/(P4)</p> <p>R29 33 4.9 1.2 VIKTOROV 80 CNTR 25,33 PI-P, 2M U G</p>
<p><math>\eta' \rightarrow (\pi^+ \pi^- \gamma (\text{including } \rho \gamma))/(\pi \pi \eta)</math> (P3)/(P1+P2)</p> <p>R7 0.31 0.15 DAVIS 68 HBC 5.5 K- P</p> <p>R7 FIT 0.459 0.035 FROM FIT</p>	<p><math>\eta' \rightarrow (\eta \mu^+ \mu^-)/\text{total} (\text{units } 10^{-5})</math> (P21)</p> <p>R30 (1.5) OR LESS CL=.90 DZHELAD 81 CNTR 30 PI-P, ETAP N</p>
<p><math>\eta' \rightarrow (\pi^0 e^+ e^-)/\text{total}</math> (P16)</p> <p>R8 (0.013) OR LESS RITTENBER 65 HBC 2.7 K-P</p>	<p><math>\eta' \rightarrow (\pi^0 \mu^+ \mu^-)/\text{total} (\text{units } 10^{-5})</math> (P22)</p> <p>R31 (6.0) OR LESS CL=.90 DZHELAD 81 CNTR 30 PI-P, ETAP N</p>
<p><math>\eta' \rightarrow (\eta e^+ e^-)/\text{total}</math> (P17)</p> <p>R9 (0.011) OR LESS RITTENBER 65 HBC 2.7 K-P</p>	<p><math>\eta' \rightarrow (3\pi^0)/(\pi^0 \pi^0 \eta)</math> (P6)/(P2)</p> <p>R32 0.0075 0.0018 BINON 84 SPEC 30-40 PI-P, G GAM</p> <p>R32 FIT 0.0075 0.0018 FROM FIT</p>
<p><math>\eta' \rightarrow (\pi^0 \rho^0)/\text{total}</math> (P18)</p> <p>R10 (0.04) OR LESS RITTENBER 65 HBC 2.7 K-P</p>	<p><b><math>\eta'(958) C</math> - NONCONSERVING DECAY PARAMETER</b></p> <p>SEE THE NOTE ON ETA DECAY PARAMETERS IN THE STABLE PARTICLE FULL LISTINGS FOR DEFINITION OF THIS PARAMETER</p>
<p><math>\eta' \rightarrow (\pi^+ \pi^- e^+ e^-)/\text{total}</math> (P10)</p> <p>R12 (0.006) OR LESS RITTENBER 65 HBC 2.7 K-P</p>	<p>A DECAY ASYMMETRY PARAMETER FOR PI+ PI- GAMMA RITTENBER 65 HBC 2.1-2.7 K-P A 103 .00 .10 KALBFLEI 75 HBC 2.2 K-P A 295 -.069 .078 GRIGORIA 75 STRC 2.1 PI-P A AVG -.0001 0.049 AVERAGE</p>
<p><math>\eta' \rightarrow (2\pi)/\text{total}</math> (P11)</p> <p>R13 (0.07) OR LESS LONDON 66 HBC COMPILATION</p>	<p><b>REFERENCES FOR <math>\eta'(958)</math></b></p>
<p><math>\eta' \rightarrow (3\pi)/\text{total}</math> (P12)</p> <p>R14 (0.07) OR LESS LONDON 66 HBC COMPILATION</p>	<p>DAUBER 64 PRL 13 449 DAUBER, SLATER, SMITH, STORK, TICH0 (UCLA)JP ALSO 64 DUBNA CONF 1 418 GOLDBERG 64 PRL 12 546 +GUNDZIK, LICHTMAN, CONNOLLY, HART, + (SYR+BNL) GOLDBERG 64 PRL 13 249 +GUNDZIK, LEITNER, CONNOLLY, HART, + (SYR+BNL) KALBFLEI 64 PRL 12 527 KALBFLEISCH, ALVAREZ, BARBARO-GALTIERI, + (LRL)JP KALBFLEI 64 PRL 13 349 G.R. KALBFLEISCH, O. DAHL, A. RITTENBERG (LRL)JP</p>
<p><math>\eta' \rightarrow (4\pi)/\text{total}</math> (P13)</p> <p>R15 (0.01) OR LESS LONDON 66 HBC COMPILATION</p>	<p>BADIER 65 PL 17 337 BADIER, DEMOULIN, BARLOUTAUD+(EPOL-SACL-AMST) KIENZLE 65 PL 19 438 KIENZLE, MAGLIC, LEVRAT, LEFEBVRES + (CERN) RITTENBER 65 PRL 15 556 RITTENBERG, KALBFLEISCH (LRL+BNL) TRILLING 65 PL 19 427 +BROWN, GOLDBERG, KADYK, SCANIO (LRL)</p>
<p><math>\eta' \rightarrow (6\pi)/\text{total}</math> (P15)</p> <p>R16 (0.01) OR LESS LONDON 66 HBC COMPILATION</p>	<p>COHN 66 PL 21 347 COHN, MCCULLOCH, BUGG, CONDO (ORNL-TEN+UCSD) LONDON 66 PR 143 1034 LONDON, RAU, SAMIOS, GOLDBERG + (BNL-SYRACUSE)EIJ MARTIN 66 PL 22,352 MARTIN, CRITTENDEN, SCHROEDER (INDIANA UI)</p>
<p><math>\eta' \rightarrow (\omega \gamma)/(\pi^+ \pi^- \eta)</math> (P5)/(P1)</p> <p>R17 68 0.068 0.013 ZANFINDO 77 ASPK 8.4 PI-P</p> <p>R17 FIT 0.065 0.013 FROM FIT</p>	<p>BARBARO-GALTIERI, MATISON, RITTENBERG- (LRL)I=0 BARLOUTAUD + (SACLAY-AMST-BGNA+REHO-EPOL)I=0 BOHLER, DALPIAZ, MASSAMA. (CERN+BGNA-STRB) DAVIS 68 PL 27 B 532 +AMAR, MOTT, DAGAN, DERRICK, FIELDS (NWES-ANL)</p>
<p><math>\eta' \rightarrow (\pi^+ \pi^- \gamma (\text{including } \rho \gamma))/(\pi \pi \eta + \omega \gamma)</math> (P3)/(P1+P2+P5)</p> <p>R18 0.25 0.14 DAUBER 64 HBC 1.95 K-P</p> <p>R18 FIT 0.441 0.034 FROM FIT</p>	<p>DUFEY 69 PL 29 B 605 +GOBBI, POUCHON, CNOPS + (ETH-CERN+SACL)JP MOTT 69 PR 177 1066 +AMMAR, DAVIS, KROPAC, SLATE, DAGAN + (NWES-ANL) RITTENBER 69 UCRL-18863 ALAN RITTENBERG (THESES) (LRL)I=0</p>
<p><math>\eta' \rightarrow (2\gamma)/\text{total}</math> (P4)</p> <p>R19 31 0.020 0.008 0.006 HARVEY 71 OSPK 3.65 PI- P, N X0 R19 68 0.0171 0.0033 DALPIAZ 72 CNTR 1.6 PI- P, N X0 R19 0.025 0.007 DUANE 74 MMS PI-P, N MM R19 6000 0.018 0.002 APEL 79 CNTR 15-60 PI- P</p> <p>R19 AVG 0.0183 0.0016 AVERAGE R19 FIT 0.0185 0.0016 FROM FIT</p>	<p>AGUILAR 70 PRL 25 1635 AGUILAR-BENITEZ, BASSANO, SAMIOS, BARNES-(BNL) BENSINGER 70 PL 35 B 505 BENSINGER, ERWIN, THOMPSON, W.D. WALKER (MISC)</p>
<p><math>\eta' \rightarrow (\pi^+ \pi^-)/\text{total}</math> (P11)</p> <p>R20 (0.02) OR LESS RITTENBER 69 HBC 1.7-2.7 K-P R20 (0.08) OR LESS CL=.95 DANBURG 73 HBC 2.2 K-P, LAM X0</p>	<p>BARBARDIN 71 PR 40 2711 BARBARDIN-OTWINOWSKA, HOFMOKL, MICHEJDA (MARS) BASILE1 71 NC 3 A 371 +BOLLINI, DALPIAZ, FRABETTI, + (CERN+BGNA+STRB) BASILE2 71 PR 32 425 +BOLLINI, DALPIAZ, FRABETTI, + (CERN+BGNA+STRB) HARVEY 71 PRL 27 883 +MAGUIE, PETERSON, RHODES, + (MILN-MICH) OGIEVETS 71 PL 35 B 69 OGIEVETSKY, TYBOR, ZASLAVSKY (DUBNA)</p>
<p><math>\eta' \rightarrow (\pi^+ \pi^- \pi^0)/\text{total}</math> (P12)</p> <p>R21 (0.05) OR LESS RITTENBER 69 HBC 1.7-2.7 K-P R21 (0.09) OR LESS CL=.95 DANBURG 73 HBC 2.2 K-P, LAM X0</p>	<p>AGUILAR 72 PR D 6 29 AGUILAR-BENITEZ, CHUNG, FISHER, SAMIOS (BNL) APEL 72 PL 40 B 680 +AUSLANDER, MULLER, BERTOLUCCI, + (KARL-MICH+LBL) BINNIE 72 PL 39 B 275 +CAMILLERI, DUANE, GARbutt, BURTON, (LOIC+SHMP) BLOODWORTH 72 NP B 39 525 BLOODWORTH, JACKSON, PRENTICE, YOON (TORONTO) DALPIAZ 72 PL 42 B 377 +FRABETTI, MASSAMA, NAVARRA, ZICHICH (CERN) RAEER 72 PR D 6 3059 +ABOLINS, DAHL, DANBURG, DAVIES, HOCH, + (LBL)</p>
<p><math>\eta' \rightarrow (\pi^+ \pi^+ \pi^- \pi^-)/\text{total}</math> (P13)</p> <p>R22 (0.01) OR LESS RITTENBER 69 HBC 1.7-2.7 K-P R22 (0.01) OR LESS CL=.95 DANBURG 73 HBC 2.2 K-P, LAM X0</p>	<p>DANBURG 73 PR D 8 3744 +KALBFLEISCH, BORENSTEIN, CHAPMAN, + (BNL-MICH) JACOBS 73 PR D 8 18 +CHANG, GAUTHIER, + (BRAN+UMD+SYR+TUFT) JP KALBFLEI 73 PRL 31 333 KALBFLEISCH, CHAPMAN, + (BNL-MICH+LBL) JP</p>
<p><math>\eta' \rightarrow (\pi^+ \pi^+ \pi^- \pi^- \pi^0)/\text{total}</math> (P14)</p> <p>R23 (0.01) OR LESS RITTENBER 69 HBC 1.7-2.7 K-P</p>	<p>BALTAY 74 PR D9 2999 +COHEN, CSORNA, HABIBI, KALELKAR, + (COLU-BING) JP DELANE 74 PR 32 425 +BINNIE, CAMILLERI, CARR, DEBENHAM, (LOR+SHAW) GAULT 74 NC 24 A 259 +JONES, SCARDON, THEMS (DURH+LOIC-ARIZ) KALBFLEI 74 PR D10 916 G.R. KALBFLEISCH (BNL)</p>
<p><math>\eta' \rightarrow (\pi^+ \pi^+ \pi^- \pi^- + \text{neutrals})/\text{total}</math> (P15+...)</p> <p>R24 (0.01) OR LESS RITTENBER 69 HBC 1.7-2.7 K-P</p>	<p>GRIGORIAN 75 NP 891 232 GRIGORIAN, LAJAGE, MELLEMA, RUDNICK, + (UCLA) KALBFLEISCH 75 PR D11 987 KALBFLEISCH, STRAND, CHAPMAN (BNL-MICH)</p>
<p><math>\eta' \rightarrow (\rho^0 \gamma)/(\pi^+ \pi^- \gamma)</math> (P7)/(P3)</p> <p>R25 0.94 0.20 AGUILAR 70 HBC 3.9-4.6 K-P R25 473 1.15 0.10 DANBURG 73 HBC 2.2 K-P, LAM X0 R25 473 (0.95) OR MORE CL=.95 DANBURG 73 HBC 2.2 K-P, LAM X0 R25 137 1.01 0.15 JACOBS 73 HBC 2.9 K-P, LAM X0</p> <p>R25 AVG 1.082 0.077 AVERAGE</p>	<p>CERRADA 77 NP B 126 189 CERRADA 77 NP B 126 189 DELAGUILL 77 PR D14 2833 DELAGUILL AND W.G. DONNELLY (BARCELONA) GESSAROL 77 NP B 126 382 GESSAROLI, + (BGNA+FRIZ+GENO+MILA+OXF+PAVI) LEDNICKY 77 E2-10521, 22, 23 LEDNICKY (JINR) ZANFINDO 77 PRL 38 930 +BROCKMAN, DANKOWYCH, + (CARL-MCGI+OHIO+TNTO)</p>
<p><b>E EQUIVALENT STATEMENTS</b></p>	<p>ABRAMS 79 PRL 43 477 +ALAM, BLOCKER, BOYARSKI, + (SLAC-LBL) APEL 79 PL 83 B 131 +AUGENSTEIN, BERTOLUCCI (KARL-PISA-SERP+WIEN) BINNIE 79 PL 83 B 141 +CARR, DEBENHAM, JONES, KARAMI, KEYNE, + (LOIC) DZHELAD 79 PL B 88 379 DZHELADIN, GOLOVKIN, GRIZUK, KACHANOV + (SERP) VIKTOROV 80 SJNP 32 520 +GOLOVKIN, DZHELADIN, ZAITSEV, MUKHIN, + (NOVO) DZHELAD 81 PL 105 B 239 DZHELADIN, GOLOVKIN, KONSTANTINOV, + (SERP)</p>
<p><math>\eta' \rightarrow (\pi^0 \pi^0 \eta (3\pi^0 \text{ decay}))/\text{total}</math> (P2N(3P10))</p> <p>R26 4 0.11 0.06 BENSINGER 70 NBC 2.2 P1+ P</p> <p>R26 FIT 0.0680 0.0062 FROM FIT</p>	<p>BARTEL 82 PL 113 B 190 +CORDS + (DESY+HAMB+HEID-LANC+MCHS+RHEL-TOKY) BICKERSTAFF 82 ZPHY C 16 171 BICKERSTAFF, MCKELLAR (MELB)</p>
	<p>BEHRND 83 PL 125 B 518 +D-AGOSTINI - (DESY+KARL+MPI+LALO+LNP+SACL) ALSO 82 PL 114 B 378 BEHRND - (DESY+KARL+MPI+LALO+LNP+SACL) JENNY 83 PR D 27 1031 +BURKE, TELNOV, ABRAMS, BLOCKER - (SLAC-LBL)</p>

# Meson Full Listings

$\eta'(958), f_0(975)$

ALTHOFF 84 PL 147 B 487 (TASSO COLLABORATION)  
 BERGER 84 PL 142 B 125 (AACH-BERG+DESY+GLAS+HAMB+UMD+SIEG+TELA+)  
 BINON 84 PL 140 B 264 +DONSKOV, DUTELL,+ (SERP+BELG+LAPP+CERN)

SENS 85 SLAC-PUB-3754 PHYSICS IN COLLISION V CONFERENCE (SLAC)

$f_0(975)$   
was  $S(975)$

$$I^G(J^{PC}) = 0^{+}(0^{++})$$

FORMERLY CALLED  $S^*$   
 UNDER THIS ENTRY WE LIST PARAMETERS OF THE POLE IN THE ISOSCALAR S-WAVE.

FOR EARLY WORK USING BREIT-WIGNER OR SCATTERING LENGTH PARAMETRIZATION IN FITS TO THE K KBAR MASS SPECTRUM, SEE REFERENCE SECTION AND OUR 1972 EDITION.

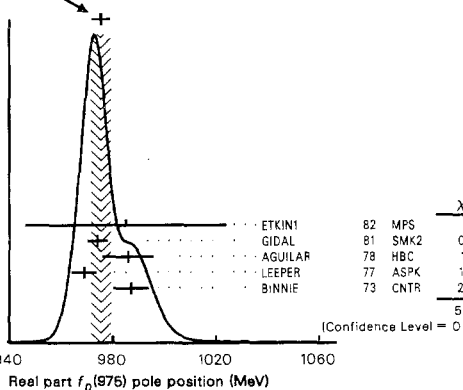
## $f_0(975)$ MASS OR REAL PART OF POLE POSITION (MeV)

M1	POLE POSITION DETERMINATIONS	PROPOPE 73 HBC	Pi+ P
M1 R	(997.) (6.)	ESTABROOK 73 ASPK	17 P1-P, P1-P1-N
M1 A	(997.) (6.)	GRAY 73 ASPK	17 P1-P, P1+P1-N
M1 R	(1012.) (20.)	HYAMS 73 ASPK	17 P1-P, N P1+P1-
M1 A	(986.) (5.)	FUJII 75 RVUE	17 P1-P, P1+P1-N
M1 AD	(998.0) (7.0)	BOHACIK 80 RVUE	
M1 E	(988.) (7.0)	IRVING 81 RVUE	
M1 F	(966.) (7.0)	IRVING 81 RVUE	
M2	987. 7.	BINNIE 73 CNTR	P1- P, MM N
M2 C	969.0 5.0	LEEPER 77 ASPK	2-2.4 P1-P
M2	986. 10.	AGUILAR 78 HBC	.7 PBAR P, KS KS
M2	(986.) (7.0)	MARTIN 79 RVUE	
M2 G	974.0 4.0	GIDAL 81 SMK2	J/PSI DECAY
M2	985.0 39.0	ETKIN1 82 MPS	25 P1-P, 2K0S N
M2 AVG	975.4 3.7	TORNQVIST 82 RVUE	

M3 MASS DETERMINATIONS (REAL PART OF MASS MATRIX EIGENVALUE)  
 M3 B (975.) ACHASOV 80 RVUE  
 M3 B (985.) TORNAVIST 82 RVUE

M A FROM SINGLE CHANNEL FIT TO HYAMS 73 DATA.  
 M B COUPLED CHANNEL ANALYSIS WITH FINITE WIDTH CORRECTIONS.  
 M C FROM COUPLED CHANNEL FIT TO HYAMS 73 AND PROPOPE 73 DATA.  
 M D POLE POSITIONS FROM ALMOST MODEL-INDEPENDENT PARAMETRIZATION.  
 M E FROM COUPLED CHANNEL ANALYSIS OF PI N --> PI PI N OR K KBAR N DATA.  
 M F SIMILAR TO (E), BUT OMIT PI N --> PI PI N MOMENTS, AND INCLUDE K-P --> PI+ PI- Y DATA.  
 M G INCLUDED IN AGUILAR 78 FIT  
 M H ETKIN1 82 QUOTES ERRORS +9/-39 MEV. WE USE +39 MEV IN THE AVERAGE.

WEIGHTED AVERAGE  
 975.4 ± 3.7 (ERROR SCALED BY 1.4)



## $f_0(975)$ WIDTH OR IMAG PART OF POLE POSITION (MeV)

M1	POLE POSITION DETERMINATIONS (CORRESPONDS TO HALF-WIDTH, NOT FULL WIDTH)	PROPOPE 73 HBC	Pi+ P
M1 R	(27.) (8.)	ESTABROOK 73 ASPK	17 P1-P, P1+P1-N
M1 A	(5.) (5.)	GRAY 73 ASPK	17 P1-P, P1+P1-N
M1 R	(16.) (5.)	HYAMS 73 ASPK	17 P1-P, N P1+P1-
M1 A	(19.) (3.)	FUJII 75 RVUE	17 P1-P, P1+P1-N
M1 AD	(19.0) (6.0)	BOHACIK 80 RVUE	
M1 E	(8.) (7.0)	IRVING 81 RVUE	
M1 F	(24.) (7.0)	IRVING 81 RVUE	
M2	24. 7.	BINNIE 73 CNTR	P1- P, MM N
M2 C	15.0 4.0	LEEPER 77 ASPK	2-2.4 P1-P
M2	50. 40.	AGUILAR 78 HBC	.7 PBAR P, KS KS
M2	(7.) (7.0)	MARTIN 79 RVUE	
M2	14.0 5.0	GIDAL 81 SMK2	J/PSI DECAY
M2	60.0 141.0	ETKIN1 82 MPS	25 P1-P, 2K0S N
M2 AVG	16.4 2.8	TORNQVIST 82 RVUE	

W3 FULL WIDTH DETERMINATIONS (FROM IMAG PART OF MASS MATRIX EIGENVALUE)  
 W3 B 70 TO 300 ACHASOV 80 RVUE  
 W3 B (400.) APPROX. TORNAVIST 82 RVUE

W A FROM SINGLE CHANNEL FIT TO HYAMS 73 DATA.  
 W B COUPLED CHANNEL ANALYSIS WITH FINITE WIDTH CORRECTIONS.  
 W C FROM COUPLED CHANNEL FIT TO HYAMS 73 AND PROPOPE 73 DATA.  
 W D POLE POSITIONS FROM ALMOST MODEL-INDEPENDENT PARAMETRIZATION.  
 W E FROM COUPLED CHANNEL ANALYSIS OF PI N --> PI PI N OR K KBAR N DATA.  
 W F SIMILAR TO (E), BUT OMIT PI N --> PI PI N MOMENTS, AND INCLUDE K-P --> PI+ PI- Y DATA.  
 W R INCLUDED IN AGUILAR 78 FIT

## $f_0(975)$ PARTIAL DECAY MODES

P1	$f_0(975) \rightarrow K\bar{K}$	DECAY MASSES
P2	$f_0(975) \rightarrow \pi\pi$	498-498
P3	$f_0(975) \rightarrow \eta\eta$	140-140
		549-549

## $f_0(975)$ BRANCHING RATIOS

$f_0(975) \rightarrow (\pi\pi)/total$	HYAMS 75 ASPK	17.2 P1-P, P1-P1-N
R1 (0.71)	WETZEL 78 OSK	8.9 P1-P, KS KS N
R1 0.78 0.03	CASON 78 STRC	7.1 P1-P, KS KS N
R1 0.81 0.09	LOVERRE 80 HBC	4. P1-P, K K N
R1 0.67 0.09		
R1 AVG 0.776 0.026	AVERAGE	

## REFERENCES FOR $f_0(975)$

WANG 61 JETP 13 323 WANG TSU-TSENG, VEKSLER, VRANA,+ (JINR)  
 BIGI 62 CERN CONF 247 A BIGI, S BRANDT, R CARRARA + (CERN)  
 BINGHAM 62 CERN CONF 240 H B BINGHAM, M BLOCH - (EPOL-CERN)  
 ERWIN 62 PRL 9 34 ERWIN, HOYER, MARCH, WALKER, WANGLER (WISC+BNL)

BALTAY 64 DUBNA CONF 1 409 BALTAY, LACH, CRENNELL, OREN, STUMP + (YALE-BNL)  
 BARMIN 64 DUBNA CONF 1 433 BARMIN, DOLGOLENKO, YEROFEEV, KRESTNI+ (ITEP)

CRENNELL 66 PRL 16 1025 CRENNELL, KALBFLEISCH, LAI, SCARR, SCHU- (BNL)  
 HESS 66 PRL 17 1109 +DAHL+HARDY+KIRZ-MILLER (LRL)

BARLOW 67 NC 50A 701 +LILLESTOL+MONTANET+ (CERN+CDEF+IRAD+LIVP)  
 BEUSCH 67 PL 25 B 357 +FISCHER, GOBBI, ASTBURY+ (ETH-CERN)  
 DAHL 67 PR 163 1377 +HARDY+HESS+KIRZ-MILLER (LRL)

ALITTI 68 PRL 21 1705 +BARNES, CRENNELL, FLAMINIO, GOLDBERG,+ (BNL)  
 LAI 68 PHILAD. CONF. P. 303 KWAN WU LAI (BNL)  
 PHELAN 68 THESIS JAMES J. PHELAN (ANL+ST. LOUIS UNIV)  
 ALSO 68 PRL 21 316 HOANG, EARTLY, PHELAN, ROBERTS + (ANL+CHIC-NDAM)

AGUILAR 69 PL 29 B 241 M. AGUILAR-BENITEZ, J. BARLOW,+ (CERN+CDEF)  
 ALSO 69 NP B 14 195 M. AGUILAR-BENITEZ, J. BARLOW,+ (CERN+CDEF)  
 HOANG 69 NC 61 A 325 T. F. HOANG (ANL)  
 HOANG 69 PR 184 1363 +EARTLY, PHELAN, ROBERTS,+ (ANL+ILLC)

BADIER 70 NP B 22 512 +BONNET, DREVILLON, BAUBILLIER,+ (EPOL+IPNP)  
 BATON 70 PL 33 B 528 +LAURENS, REIGNIER (SACLAY)  
 BEUSCH 70 PHLA. CONF. P. 185 W. BEUSCH (ETH-CERN)  
 HYAMS 70 PHLA. CONF. P. 41 +KOECH, BEUSCH,+ (CERN+MPI+ETH-LOIC+HAW)  
 ALSO 70 NP B 22 189 HYAMS, KOCH, PÖTTER, VON LINDERN,+ (CERN+MPI)  
 OH 70 PR D 1 2494 +GARFINKEL, MORSE, WALKER, PRENTICE (WISC+TNT)

ALSTON-G 71 PL 36 B 152 ALSTON-GARNJOST, BARBARO-GALTIERI,+ (LBL)

BASDEVANT 72 PL 41 B 178 BASDEVANT, FROGGATT, PETERSEN (CERN)  
 DAMERI 72 NC 9 A 1 +BORZATTA, GOUSSU,+ (GENO-MILA-SACL)  
 DUBOC 72 NP B 46 429 +GOLDBERG, MAKOWSKI, DONALD,+ (LBNL+LIVP)  
 FLATTE 72 PL 38 B 232 +ALSTON-GARNJOST, BARBARO-GALTIERI,+ (LBL)  
 GRAY 72 PHIL. CONF. PROC. 5 HYAMS, JONES, SCHLEIN, BLUM, DIETL+ (CERN-MPI)  
 WILLIAMS 72 PR D 6 3178 P. K. WILLIAMS (FSU)

BINNIE 73 PRL 31 1534 +CARR, DEBENHAM, DUANE, GARBUTT,+ (LOIC+SHMP)  
 DIAMOND 73 PR D 7 1977 +BINKLEY,+ (WISC+DUKE+COLO+TNTO-OHIO)  
 ESTABROOK 73 TALLAHASSEE ESTABROOKS, MARTIN, GRAY, HYAMS+ (CERN-MPI)  
 FUJII 73 NC 13 A 311 Y. FUJII, M. KATO (TOKYO)  
 GRAY 73 TALLAHASSEE +HYAMS, JONES, BLUM, DIETL, KOCH+ (CERN-MPI)  
 HYAMS 73 NP B 64 134 +JONES, WEILHAMMER, BLUM, DIETL+ (CERN-MPI)  
 OCHS 73 THESIS W. OCHS (MPI)  
 PROPOPE 73 PR D 7 1280 PROPOPEUCO, GARNJOST, GALTIERI, FLATTE+(LBL)

GRAY 74 NP B 75 189 +HYAMS, JONES, BLUM, DIETL, KOCH+ (CERN-MPI)  
 GRAY 74 NP B 76 375 +HYAMS, JONES, BLUM, DIETL (CERN-MPI)  
 MORGAN 74 PL 518 71 D. MORGAN (RHEL)

FUJII 75 NP 885 179 Y. FUJII, M. FUKUGITA (TOKYO)  
 HYAMS 75 NP 100 205 +JONES, WEILHAMMER, BLUM, DIETL+ (CERN-MPI)  
 MORGAN 75 ARGONNE CONF. 45 D. MORGAN (RHEL)  
 PAWLICKI 75 PR D12 631 +AYRES, DIEBOLD, GREENE, KRAMER, WICKLUND (ANL)

BRANDEB 76 NP B 104 413 +CARNEGIE, CASHMORE, DAVIER, LASINSKI,+ (SLAC)  
 BUTTRAM 76 PR D 13 1153 +CRAWLEY, DUKE, LAMB, LEEPER, PETERSON (ISU)  
 CERRADA 76 PL 62 B 353 +GONZALEZ-ARROYO, RUBIO, YNDURAIN (CERN+MADR)  
 LATTE 76 PL 63 B 228 S. M. FLATTE (CERN)  
 WETZEL 76 NP B 115 208 +REUDENRICH, BEUSCH,+ (ETH-CERN+LOIC)  
 WILKINS 76 PR D 13 1831 +ALBRIGHT, S+V HAGOPIAN, LANNUTTI (FSU)

FROGGATT 77 NP B 129 89 +PETERSEN (GLASSOW-COPENHAGEN)  
 LEEPER 77 PR D 16 2054 +BUTTRAM, CRAWLEY, DUKE, LAMB, PETERSON (ISU)  
 MARTIN 77 NP B 121 514 +OZMUTLU, SOUHRES (DURHAM)  
 PAWLICKI 77 PR D 15 3196 +AYRES, COHEN, DIEBOLD, KRAMER, WICKLUND (ANL) J

AGUILAR 78 NP B 140 73 +CERRADA,+ (MADRID+BOMBAY+CERN-PARIS)  
 BALAND 78 NP B 140 220 +GRARD, JOHNSON,+ (MONS-BELG-CERN+LOIC-LALO)  
 CASON 78 PRL 41 271 +BAUMBAUGH, BISHOP, BISWAS, KENNEY,+ (NDAM+ANL)

For notation, see key on page 91.

Meson Full Listings
f0(975), a0(980)

Table listing meson production experiments with columns for author, year, experiment, and meson type.

a0(980) PARTIAL DECAY MODES table showing decay channels and branching ratios.

a0(980) BRANCHING RATIOS table providing detailed branching ratios for various decay modes.

a0(980) was delta(980)

I(G^PC) = 1^-(0^+)

OUR LATEST MINIREVIEW ON THIS PARTICLE CAN BE FOUND IN THE 1984 EDITION.

a0(980) MASS (MeV)

Table of a0(980) mass measurements from various experiments.

Systematic error and energy calibration added. Coupled channel analysis with finite width corrections.

a0(980) WIDTH (MeV)

Table of a0(980) width measurements from various experiments.

Coupled channel analysis with finite width corrections. Error in the paper is wrongly quoted at one point.

REFERENCES FOR a0(980)

Extensive list of references for the a0(980) particle, including experiment names and publication details.







# Meson Full Listings

$h_1(1190)$ ,  $b_1(1235)$

$h_1(1190)$   
was  $H(1190)$

$$I^G(J^{PC}) = 0^-(1^{--})$$

$h_1(1190)$ MASS (MeV)			
M	C	1190. 60.	DANKOWYCH 81 SPEC
M	T	(1175.0) APPROX	TORNQVIST 82 RVUE
M	C	USES THE MODEL OF BOWLER 75	
M	T	FROM A UNITARIZED QUARK MODEL CALCULATION	

$h_1(1190)$ WIDTH (MeV)			
W	C	320. 50.	DANKOWYCH 81 SPEC
W	T	(365.0) APPROX	TORNQVIST 82 RVUE
W	C	USES THE MODEL OF BOWLER 75	
W	T	FROM A UNITARIZED QUARK MODEL CALCULATION	

$h_1(1190)$ PARTIAL DECAY MODES			
P1	$h_1(1190) \rightarrow \rho \pi$	DECAY MASSES	769- 135

$h_1(1190)$ BRANCHING RATIOS			
$h_1(1190) \rightarrow (\rho \pi)/\text{total}$			(P1)
R1	SEEN	DANKOWYCH 81 SPEC	8 PI P,3 PI N
R1	SEEN	ATKINSON 84 OMEG	20-70GAM P,3 PI

REFERENCES FOR $h_1(1190)$			
BOWLER 75 NP B97 227	-GAME,ATTCHISON,DAINTON	(OXF+DARE)	
DANKOWYCH 81 PRL 46 580	+BROCKMAN,EDWARDS+(TNT0-BNL-CARL+MCGI+OHIO)		
TORNQVIST 82 NP B 203 268	TORNQVIST	(HELS)	
ATKINSON 84 NP B 231 15	ATKINSON+ (BONN+CERN+GLAS+LANC+MCHS+LPNP+)		

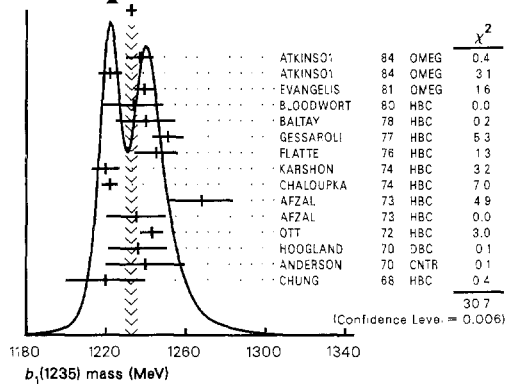
$b_1(1235)$   
was  $B(1235)$

$$I^G(J^{PC}) = 1^+(1^{+-})$$

$b_1(1235)$ MASS (MeV)			
M	W	(1228.) (5.)	FRENKIEL 72 HBC -- 0. PBAR PI,5 PI
M	T	360(1208.0) (18.0)	GAVILLET 78 HBC + 4.2 K-P,BACKWARD
M	C	(1243.0) APPROX	TORNQVIST 82 RVUE
M	C	(1215.) (5.)	ATKINSON 84 OMEG 020-70GAM P
M	C	(1271.) (11.)	COLLICK 84 SPEC + 200 PI+Z,Z PIOME
M	C	1220. 20.	CHUNG 68 HBC - 3.2,4.2 PI- P
M	C	1240.0 20.0	ANDERSON 70 CNTR 0 5-18 GAMMA P
M	C	1236.0 15.0	HOOGLAND 70 DBC - 3.0 K- D
M	O	1163 1243. 6.	OTT 72 HBC + 7.1 PI+ P
M	C	1235. 15.	AFZAL 75 HBC + 11.7 PI+ P
M	C	1268. 16.	AFZAL 75 HBC - 11.2 PI- P
M	C	1400 1222. 4.	CHALOUKPA 74 HBC - 3.9 PI-P
M	C	600 1220. 7.	KARSHON 74 HBC + 4.9 PI+P
M	C	890 1245.0 11.0	FLATTE 76 HBC - 4.2 K-P,PI-OMEGA
M	C	450 1251.0 8.0	GESSAROLI 77 HBC - 11 PI-P,PI- OME
M	C	225 1240.0 15.0	BALTAY 78 HBC + 15 PI+P,P 4PI
M	C	105 1234.0 15.0	BLOODWORT 80 HBC - 8.2 K- P
M	C	1239. 5.	EVANGELIS 81 OMEG - 12 PI-P,OME PI P
M	C	1222. 5.	ATKINSO1 84 OMEG - 25-55GAM P,OM PI
M	C	1237. 7.	ATKINSO1 84 OMEG 025-55GAM P,OM PI
M	AVG	1232.6 3.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5) (SEE IDEOGRAM BELOW)

M O FROM FIT OF THE MASS SPECTRUM  
M W FIT REQUIRES AN ADDITIONAL JP=1- RESONANCE  
M W AT 1256 MEV, WIDTH 129 MEV.  
M T FROM A UNITARIZED QUARK MODEL CALCULATION

WEIGHTED AVERAGE  
1232.6 ± 3.0 (ERROR SCALED BY 1.5)



$b_1(1235)$ WIDTH (MeV)			
W	W	(126.) (10.)	FRENKIEL 72 HBC -- 0. PBAR PI,5 PI
W	T	360 (163.0) (50.8)	GAVILLET 78 HBC + 4.2 K-P,BACKWARD
W	C	(118.0) APPROX	TORNQVIST 82 RVUE
W	C	(231.) (14.)	ATKINSON 84 OMEG 020-70GAM P
W	C	(232.) (29.)	COLLICK 84 SPEC + 200 PI+Z,Z PIOME
W	C	150. 20.	CHUNG 68 HBC - 3.2,4.2 PI- P
W	C	132.0 20.0	HOOGLAND 70 DBC - 3.0 K- D
W	O	1163 134. 23.	OTT 72 HBC + 7.1 PI+ P
W	C	120. 50.	AFZAL 73 HBC + 11.7 PI+ P
W	C	135. 50.	AFZAL 75 HBC - 11.2 PI- P
W	C	1400 135. 20.	CHALOUKPA 74 HBC - 3.9 PI-P
W	C	600 156. 22.	KARSHON 74 HBC + 4.9 PI+P
W	C	890 182.0 45.0	FLATTE 76 HBC - 4.2 K-P,PI-OMEGA
W	C	450 155.0 32.0	GESSAROLI 77 HBC - 11 PI-P,PI- OME
W	C	225 170.0 50.0	BALTAY 78 HBC + 15 PI+P,P 4PI
W	C	105 150.0 50.0	BLOODWORT 80 HBC - 8.2 K- P
W	C	170. 15.	EVANGELIS 81 OMEG - 12 PI-P,OME PI P
W	AVG	150.0 7.3	AVERAGE

W O FROM FIT OF THE MASS SPECTRUM  
W T SEE NOTE UNDER THE MASS ABOVE.  
W T FROM A UNITARIZED QUARK MODEL CALCULATION

$b_1(1235)$ PARTIAL DECAY MODES			
P1	$b_1(1235) \rightarrow \omega \pi$	DECAY MASSES	783+ 140
P2	$b_1(1235) \rightarrow 2\pi^+ 2\pi^-$		140+ 140+ 140+ 140
P3	$b_1(1235) \rightarrow K \bar{K}$		494+ 494
P4	$b_1(1235) \rightarrow \pi \pi$		140+ 140
P5	$b_1(1235) \rightarrow \pi \phi$		135+1020
P6	$b_1(1235) \rightarrow \eta \pi$ (FORBIDDEN BY G)		549+ 140
P7	$b_1(1235) \rightarrow K \bar{K} \pi$		494+ 494+ 140
P8	$b_1(1235) \rightarrow \eta \rho$		549+ 769

$b_1(1235)$ PARTIAL WIDTHS (keV)			
$b_1(1235) \rightarrow \pi^\pm \gamma$			(63)
W3	230.0	60.0	COLLICK 84 SPEC + 200 PI+Z,Z PIOME

$b_1(1235)$ D-wave/S-wave RATIO IN $\omega \pi$ DECAY			
DS	0.3	0.1	CHALOUKPA 74 HBC - 3.9-7.5 PI-P
DS	600 0.35	0.25	KARSHON 74 HBC + 4.9 PI+P
DS	0.21	0.08	CHUNG 75 HBC + 7.1 PI+P
DS	0.4	0.1	0.1 GESSAROLI 77 HBC - 11 PI-P,PI- OME
DS	0.235	0.047	ATKINSO3 84 OMEG 20-70 GAM P
DS	AVG	0.260 0.035	AVERAGE

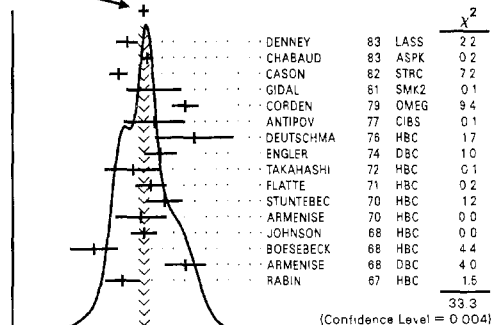


Meson Full Listings

$\rho(1250)$ ,  $f_2(1270)$

Table of references for  $\rho(1250)$  and  $f_2(1270)$ , listing authors, publication details, and values.

WEIGHTED AVERAGE 176.2 ± 5.3 (ERROR SCALED BY 1.5)



f2(1270) was f(1270)

I^G(J^PC) = 0^-(2^+)

f2(1270) MASS (MeV)

Table of f2(1270) mass measurements from various experiments, including authors, years, values, and errors.

AVG 1273.8 ± 1.5 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)

Notes regarding data inclusion and error analysis for the mass measurements.

f2(1270) PARTIAL WIDTHS (keV)

Table of partial widths for f2(1270) decays into various meson pairs, listing authors, widths, and errors.

Notes regarding data inclusion and error analysis for the partial widths.

f2(1270) WIDTH (MeV)

Table of f2(1270) total width measurements from various experiments.

AVG 176.2 ± 5.3 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)

Notes regarding data inclusion and error analysis for the width measurements.

f2(1270) BRANCHING RATIOS

Table of branching ratios for f2(1270) decays into  $2\pi^+ 2\pi^-$  and  $(\pi^+ \pi^-)$ .

Table of branching ratios for f2(1270) decays into  $(\pi^+ \pi^- 2\pi^0)$ .

Table of branching ratios for f2(1270) decays into  $(K \bar{K})$  and other channels.

f2(1270) PARTIAL DECAY MODES

Table of partial decay modes for f2(1270), listing decay modes and their masses.

Notes regarding data inclusion and error analysis for the branching ratios.

Table of branching ratios for f2(1270) decays into  $(K^0 \bar{K}^0 \pi^+ + C.C.)$ .

Table of branching ratios for f2(1270) decays into  $(\eta \pi \pi)$ .





For notation, see key on page 91.

Meson Full Listings

f<sub>1</sub>(1285)

Table with columns M, R, T, W, X, Y, Z containing experimental data for f<sub>1</sub>(1285). Includes entries like BROMBERG 80 SPEC, DIONISI 80 HBC, etc.

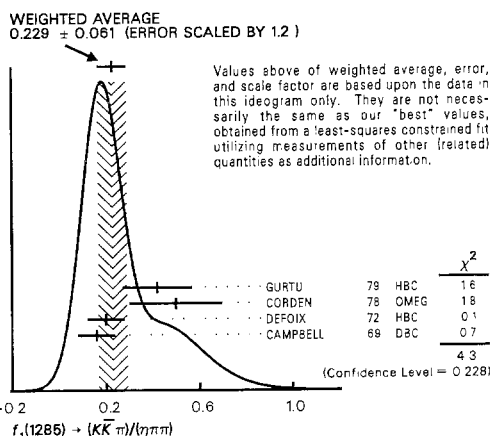


Table titled 'f1(1285) WIDTH (MeV)' listing various experiments and their widths. Includes a summary row: 'AVG 1285.4 1.3 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.1)'. Also includes notes about phase shift analysis and unitarized quark model calculation.

Table titled 'f1(1285) PARTIAL DECAY MODES' listing decay channels and corresponding masses. Includes modes like P1 f1(1285) -> K K pi, P2 f1(1285) -> pi pi rho, etc.

Table showing various decay rates and branching ratios for f1(1285). Includes entries like f1(1285) -> (K0(980) pi)/(eta pi pi) and f1(1285) -> (2 pi+ 2 pi- (including rho pi pi))/(eta pi+ pi-).

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions. The diagonal elements are Pi\_i = delta Pi\_i, where delta Pi\_i = sqrt(delta Pi\_i delta Pi\_i), while the off-diagonal elements are the normalized correlation coefficients (delta Pi\_i delta Pi\_j) / (delta Pi\_i delta Pi\_j). For the definitions of the individual Pi\_i see the listings above; only those Pi\_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Correlation matrix for branching fractions. Matrix elements include values like 0.266, 0.086, 0.086, 0.2931, 0.4883, 0.086, 0.6152, 0.9341, 0.4000, 0.0711.

f1(1285) BRANCHING RATIOS

THE f1(1285) BRANCHING RATIOS FIT IS MADE WITH THE ASSUMPTION THAT THE f1(1285) INTO 4PI DECAY IS ALWAYS VIA DECAY INTO I=1 PI PI PAIRS (E.G., RHO PI PI).

Table listing branching ratios for f1(1285) into various partial decay modes. Includes modes like (pi pi rho)/(K K pi) and (K K pi)/(eta pi pi) with associated ratios and errors.

REFERENCES FOR f1(1285)

List of references for f1(1285) studies, including authors and journal information. Examples include BARLOW, ADAMSON, CHUNG, DALH, HESS, HARDY, KIRZ, etc.











# Meson Full Listings

$f_2(1410), f_1(1420)$

## $f_2(1410)$ PARTIAL DECAY MODES

P1	$f_2(1410) \rightarrow K_S^0 K_3^0$	DECAY MASSES 498+ 498
<b>REFERENCES FOR <math>f_2(1410)</math></b>		
BEUSCH	67 PL 25 B 357	+FISCHER, BOBBI, ASTBURY+ (ETH+CERN)
BEUSCH	70 EXPERIM. MESON SPECTROSCOPY, P. 185, COLUMBIA U. PRESS (ETH+CERN)	
BALOSHIN	76 SJNP 24 297	+BOLDNKIN, VLADIMIRSKI J, GRIGORJEV, + (ITEP)
WETZEL	76 NP B 115 208	+FREUDENREICH, BEUSCH, + (ETH+CERN+LOIC)
POLYCHRO	79 PR D 19 1317	POLYCHRONAKOS, CASON, BISHOP+ (NDAM+ANL)
CHABAUD	81 APP B 12 575	+NICZYPORUK, BECKER- (CERN+CRAC+MPIM)JP
ETKIN	82 PR D 25 1786	+FOLEY, LAI, LINDENBAUM+ (BNL+CUNY+TUFT+VAND)
CHABAUD	83 NP B 223 1	+GORLICH, CERRADA+ (CERN+CRAC+MPIM)JP
PALKA	83 THESIS 1230/PH	H. PALKA (CRAC) JP
DAUM	84 ZPHY C 23 339	+HERTZBERGER+ (AMST+CERN+CRAC+MPIM+OXF+RHEL) JP
TURNAU	84 ZPHY C 25 299	J. TURNAU (CRAC)

$f_1(1420)$   
was  $E(1420)$

$$I^G(J^{PC}) = 0^{++}(1^{++})$$

### NOTE ON $f_1(1420)$ AND $\eta(1440)$

For this edition we continue to split the data on  $f_1(1420)/\eta(1440)$  into two entries according to the proposed  $J^{PC}$  assignments.

The  $J^{PC} = 1^{++}$  state (DIONISI 80, ARMSTRONG 84), now called  $f_1(1420)$ , appears to have a dominant decay mode into the  $K^*(892)\bar{K}$  system.

The state with  $J^{PC} = 0^{-+}$  is now called  $\eta(1440)$ . Under this entry we group the results obtained in the early  $\bar{p}p$  annihilation experiment at rest (BAILLON 67,83), in the study of the radiative decay of the  $J/\psi$  resonance (SCHARRE 80,81, EDWARDS 82, RICHMAN 85), and in a hadroproduction experiment (CHUNG 85). The  $\eta(1440)$  is largely coupled to the  $a_0(980)\pi$  decay channel, although the lack of a signal in the  $\eta\pi\pi$  system (EDWARDS 83) is a source of concern.

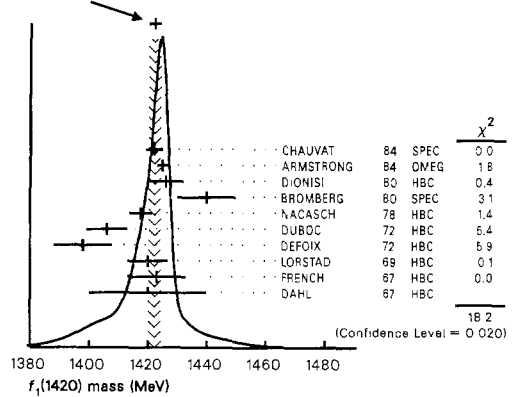
The CHUNG 85 study of the  $K\bar{K}\pi$  system is based on more than ten times the statistics of DIONISI 80. The  $J^{PC} = 0^{-+}$  wave dominates the 1420-MeV mass region with little evidence for a  $1^{++}$  resonance. The CHUNG 85 results may suggest that the  $f_1(1420)$  is in fact the  $\eta(1440)$ .

### $f_1(1420)$ MASS (MeV)

M	1420.	20.	DAHL	67 HBC	1.6-4.2 PI- P
M	1423.0	10.0	FRENCH	67 HBC	3-4 PBAR P
M	310 1420.	7.	LORSTAD	69 HBC	.7PB P, 4, 5-BODY
M	170 1398.	10.	DEFOIX	72 HBC	0.7 PBAR P, 7 PI
M	280 1406.	7.	DUBOC	72 HBC	1.2 PBAR P, 2K4PI
M	1417.5	4.	NACASCH	78 HBC	.7, .76 PBAR P
M	1440.0	10.0	BROMBERG	80 SPEC	100 PI-P, 2KPIX
M	221 1426.0	6.0	DIONISI	80 HBC	4. PI-P, K K PI N
M	(1423.0) APPROX		TORNQVIST	82 RVUE	
M	1520 1425.0	2.0	ARMSTRONG	84 OMEG	85 PI-P, K KBARPI
M	1422.0	3.0	CHAUVAUT	84 SPEC	ISR 31.5 PP
M	AVG	1422.3	2.1	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5) (SEE IDEOGRAM BELOW)	

M A MASS ERROR INCREASED TO ACCOUNT FOR A0(980) MASS CUT  
M A UNCERTAINTIES  
M T FROM A UNITARIZED QUARK MODEL CALCULATION

WEIGHTED AVERAGE  
1422.3 ± 2.1 (ERROR SCALED BY 1.5)



### $f_1(1420)$ WIDTH (MeV)

W	60.0	20.0	DAHL	67 HBC	1.6-4.2 PI- P
W	45.	20.	FRENCH	67 HBC	3-4 PBAR P
W	310 60.	20.	LORSTAD	69 HBC	.7PB P, 4, 5-BODY
W	170 50.	10.	DEFOIX	72 HBC	0.7 PBAR P, 7 PI
W	280 50.	12.	DUBOC	72 HBC	1.2 PBAR P, 2K4PI
W	53.	20.0	NACASCH	78 HBC	.7, .76 PBAR P
W	62.0	14.0	BROMBERG	80 SPEC	100 PI-P, 2KPIX
W	221 40.0	15.0	DIONISI	80 HBC	4. PI-P, K K PI N
W	1520 62.0	5.0	ARMSTRONG	84 OMEG	85 PI-P, K KBARPI
W	47.0	10.0	CHAUVAUT	84 SPEC	ISR 31.5 PP
W	AVG	55.9	3.4	AVERAGE	

### $f_1(1420)$ PARTIAL DECAY MODES

P1	$f_1(1420) \rightarrow K \bar{K}^*(892)$	DECAY MASSES 498+ 892
P2	$f_1(1420) \rightarrow K \bar{K} \pi$	498+ 498+ 140
P3	$f_1(1420) \rightarrow \pi \pi \rho$	140+ 140+ 769
P4	$f_1(1420) \rightarrow a_0(980) \pi$	983+ 140
P5	$f_1(1420) \rightarrow \eta \pi \pi$	549+ 140+ 140
P6	$f_1(1420) \rightarrow 4\pi$	140+ 140+ 140+ 140
P7	$f_1(1420) \rightarrow \gamma \gamma$	0+ 0

### $f_1(1420)$ $\Gamma(I)^*\Gamma(\gamma\gamma)/\Gamma(\text{total})$ (keV)

$\Gamma(K \bar{K} \pi)^*\Gamma(\gamma\gamma)/\Gamma(\text{total})$	$\Gamma(\gamma\gamma)/\Gamma(\text{total})$
W7 (8.0) OR LESS	CL=0.95 JENNI 83 SMK2 GAM GAM, KKBAR PI

### $f_1(1420)$ BRANCHING RATIOS

$f_1(1420) \rightarrow (\bar{K}K^*(892) + \text{C.C.})/(K \bar{K} \pi)$	(P1)/(P2)
R1 0.76 0.06	BROMBERG 80 SPEC 100 PI-P, 2KPIX
R1 0.86 0.12	DIONISI 80 HBC 4. PI-P, K K PI N
R1 AVERAGE MEANINGLESS	
$f_1(1420) \rightarrow (\pi \pi \rho)/(K \bar{K} \pi)$	(P3)/(P2)
R2 (2.0) OR LESS	DAHL 67 HBC 0 1.6-4.2 PI- P
R2 (0.3) OR LESS CL=.95	CORDEN 78 OMEG 12-15PI-P
$f_1(1420) \rightarrow (\eta \pi \pi)/(K \bar{K} \pi)$	(P5)/(P2)
R3 (1.5) OR LESS CL=.95	FOSTER 68 HBC -0.0 PBAR P
R3 1.5 0.8	DEFOIX 72 HBC 0.7 PBAR P
R3 (0.5) OR LESS CL=.95	CORDEN 78 OMEG 12-15PI-P
$f_1(1420) \rightarrow (a_0(980) \pi)/(\eta \pi \pi)$	(P4)/(P5)
R4 0.4 0.2	DEFOIX 72 HBC 0.7 PBAR P, 7 PI
R4 NOT SEEN IN EITHER MODE	CORDEN 78 OMEG 12-15PI-P
$f_1(1420) \rightarrow (4\pi)/(\bar{K}K^*(892) + \text{C.C.})$	(P6)/(P1)
R5 (0.90) OR LESS CL=.95	DIONISI 80 HBC 4. PI-P, K K PI N
$f_1(1420) \rightarrow (K \bar{K} \pi)/(a_0(980) \pi) + \bar{K}K^*(892) + \text{C.C.}$	(P2)/(P1+P4)
R6 C 0.65 0.27	DIONISI 80 HBC 4. PI-P
R6 C CALCULATED USING	
R6 C (A0(980) INTO K KB)/(A0(980) INTO ETA PI) = 0.24 ± 0.07	
$f_1(1420) \rightarrow (a_0(980) \pi)/(\bar{K}K^*(892))$	(P4)/(P1)
R7 (0.04) OR LESS CL=.68	ARMSTRONG 84 OMEG 85 PI-P, KKBAR PI













Meson Full Listings

π<sub>2</sub>(1680), φ(1680), ρ<sub>3</sub>(1690)

Table listing meson production experiments with columns for experiment name, particle ID, and production method. Includes experiments like ARMENISE 69, BRANDENB 70, BEKETOV 71, etc.

φ(1680) PARTIAL DECAY MODES

Table showing decay modes for φ(1680) such as φ(1680) → ωππ, φ(1680) → 3π, φ(1680) → K<sup>±</sup>K<sup>0</sup>, etc., with associated decay masses.

φ(1680) BRANCHING RATIOS

Table of branching ratios for φ(1680) decays, including theoretical predictions and experimental data points like R1, R2, R3, R4.

THIS COMBINATION OF A PARTIAL WIDTH WITH THE PARTIAL WIDTH INTO E-E- AND WITH THE TOTAL WIDTH IS OBTAINED FROM THE INTEGRATED CROSS SECTION INTO CHANNEL (1) IN E-E- ANNIHILATION.

Table showing integrated cross sections and branching ratios for φ(1680) decays, with columns for Γ(ωππ)Γ(e<sup>+</sup>e<sup>-</sup>)/Γ(total) and Γ(K<sup>±</sup>K<sup>0</sup>)/Γ(total).

φ(1680)

I<sup>G</sup>(J<sup>PC</sup>) = 0<sup>-</sup>(1<sup>-</sup>-)

FORMERLY CALLED PHIPRIME. FIRST IDENTIFIED USING DALITZ PLOT ANALYSIS OF E+E- INTO K<sup>+</sup>K<sup>0</sup>(892) (BIZOT 80, DELCOURT 81) WE LIST BELOW CANDIDATES FOR THE OMEGA AND PHI RADIAL EXCITATIONS UNTIL THEY ARE WELL ESTABLISHED.

φ(1680) MASS (MeV)

Table of φ(1680) mass measurements from various experiments like M K KBAR MODE, M B, M C, M D, etc.

PIONS MODE

Table of pions mode mass measurements for φ(1680) from experiments like COSME, ESPOSITO, CORDIER, etc.

M B JP NOT UNAMBIGUOUSLY 1- FROM GLOBAL FIT OF RHO, OMEGA, PHI AND THEIR RADIAL EXCITATIONS TO CHANNELS OMEGA PI-PI-, K-K-, KS KL, KS K-S PI-. ASSUME MASS 1570 MEV AND WIDTH 510 MEV FOR RHO RAD. EXIT., MASS 1570 AND WIDTH 500 MEV FOR OMEGA RADIAL EXCITATION.

φ(1680) WIDTH (MeV)

Table of φ(1680) width measurements from experiments like M K KBAR MODE, W B, W C, W D, etc.

PIONS MODE

Table of pions mode width measurements for φ(1680) from experiments like COSME, ESPOSITO, CORDIER, etc.

M B JP NOT UNAMBIGUOUSLY 1- SEE NOTE C UNDER MASS. W D FIT TO ONE CHANNEL ONLY, NEGLECTING INTERFERENCE WITH OMEGA, W D RHO(1600). W A MAY BE PHI OR OMEGA RADIAL EXCITATION. INTERPRETATION COMPLICATED.

REFERENCES FOR φ(1680)

Bibliographic references for φ(1680) including experiments like COSME, ASTON, BIZOT, ESPOSITO, etc.

ρ<sub>3</sub>(1690) was g(1690)

I<sup>G</sup>(J<sup>PC</sup>) = 1<sup>-</sup>(3<sup>-</sup>-)

ρ<sub>3</sub>(1690) MASS (MeV)

WE ONLY INCLUDE HIGH STATISTICS EXPERIMENTS IN THE AVERAGE FOR THE ZPI AND KKBAR MODES.

ZPI MODE

Table of ρ<sub>3</sub>(1690) mass measurements from experiments like M ZPI MODE, M B, M C, M D, M E, etc.

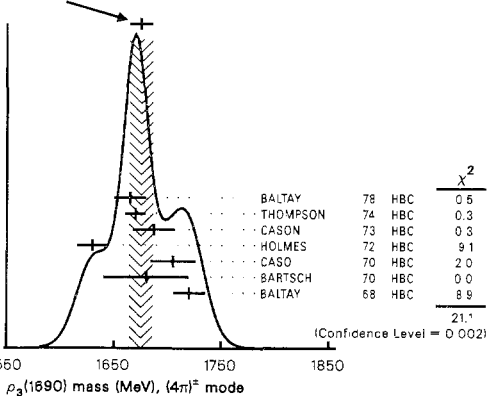
For notation, see key on page 91.

Meson Full Listings  
 $\rho_3(1690)$

K KBAR + K KBAR PI MODE
1690.0 16.0 ADERHOLZ 69 HBC + 8 PI+ P, KKBARPI
...
AVG 1690.9 2.6 AVERAGE

FROM RHO- RHOD MODE, NOT INDEPENDENT OF (A), (C)
FROM (A2)- P1D MODE, NOT INDEPENDENT OF (A), (C)
FROM (A2)D PI- MODE, NOT INDEPENDENT OF (A), (B)
FROM (RHO-- RHOD) MODE

WEIGHTED AVERAGE
1675.2 ± 11.1 (ERROR SCALED BY 1.9)



OMEGA PI MODE
1654. 24. BARNHAM 70 HBC + 10 K+ P, OMEGA PI
...
AVG 1680.1 6.6 AVERAGE

rho\_3(1690) WIDTH (MeV)

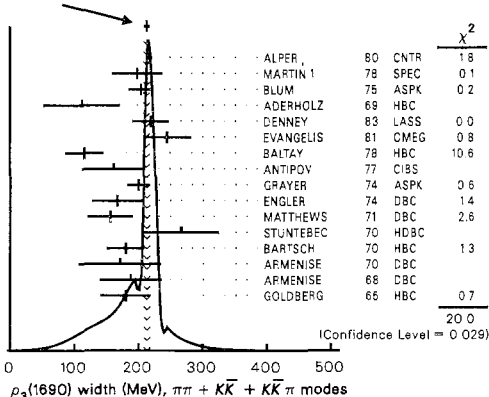
WE ONLY INCLUDE HIGH STATISTICS EXPERIMENTS IN THE AVERAGE FOR THE 2PI AND KKBAR MODES.

2PI MODE
180.0 40.0 GOLDBERG 65 HBC 0 6 PI+D, 8 PI-P
...
220. 29. DENNEY 83 LASS 10 PI-W /PI+ P

FROM PHASE-SHIFT ANALYSIS
ERROR TAKES ACCOUNT OF SPREAD OF DIFFERENT PHASE-SHIFT SOLUTIONS
USES SAME DATA AS HYAMS 75 AND BECKER 79
FROM A PHASE SHIFT SOLUTION CONTAINING A F2(1525) WIDTH
TWO TIMES LARGER THAN THE K KBAR RESULT.
WIDTH ERRORS ENLARGED BY US TO 4\*WIDTH/SQRT(N), SEE K\*(892) NOTE

K KBAR + K KBAR PI MODE
112.0 60.0 ADERHOLZ 69 HBC + 8 PI+ P, KKBARPI
...
AVG 213.6 5.1 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)

WEIGHTED AVERAGE
213.6 ± 5.1 (ERROR SCALED BY 1.4)



(4PI)-- MODE
100. 35. BALTAY 68 HBC + 7, 8.5 PI+ P
...
AVG 117.8 12.9 AVERAGE

OMEGA PI MODE
130. 73. 43. BARNHAM 70 HBC + 10 K+ P, OMEGA PI
...
AVG 112.9 20.6 AVERAGE

rho\_3(1690) PARTIAL DECAY MODES

P1 rho\_3(1690) -> pi pi 140+ 140
P2 rho\_3(1690) -> 4pi (including pi^0 S) 140+ 140+ 140+ 140
P3 rho\_3(1690) -> K K pi 498+ 498+ 140
P4 rho\_3(1690) -> K K 498+ 498
P5 rho\_3(1690) -> pi pi rho (excluding 2rho and a2 pi) 140+ 140+ 769
P6 rho\_3(1690) -> a2(1320) pi 1318+ 140
P7 rho\_3(1690) -> omega pi 140+ 783
P8 rho\_3(1690) -> 2rho 769+ 769
P9 rho\_3(1690) -> phi pi 1020+ 140
P10 rho\_3(1690) -> eta pi 549+ 140
P11 rho\_3(1690) -> pi +/- pi +/- pi^0 140+ 140+ 140+ 140

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions, P\_i, as follows: The diagonal elements are P\_i +/- delta P\_i, where delta P\_i = sqrt(delta P\_i delta P\_i), while the off-diagonal elements are the normalized correlation coefficients (delta P\_i delta P\_j) / (delta P\_i delta P\_j). For the definitions of the individual P\_i, see the listings above; only those P\_i appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.



For notation, see key on page 91.

# Meson Full Listings

$f_2(1720)$ ,  $f_0(1730)$ ,  $\pi(1770)$ ,  $f_2(1810)$

$f_2(1720)$   
was  $\theta(1690)$

$$I^G(J^{PC}) = 0^+(2^{++})$$

NAMED THETA BY EDWARDS 82  
SEEN IN J/PSI INTO GAMMA F2(1720), THEREFORE C=+.  
F2(1720) DECAYS INTO 2ETA, THEREFORE IG=0-.  
JP=2- IS PREFERRED OVER 0-, HIGHER SPINS NOT  
STUDIED.  
MASS AND WIDTH DETERMINATION COMPLICATED BY OVERLAP WITH F2\*(1525)  
IN MASS SPECTRA. POSSIBLE CONNECTION OF THIS STATE WITH STRUCTURE  
SEEN IN J/PSI TO GAMMA RHO RHO AND IN J/PSI TO GAMMA ETA PI PI  
IS UNCLEAR (SEE BURKE 82, HITLIN 83). RECENT RESULTS (BALTRUSAITIS 85)  
INDICATE THAT THE RHO RHO ENHANCEMENT IS JP=0-, HENCE UNRELATED TO THE  
F2(1720).

### $f_2(1720)$ MASS (MeV)

M	E	(1640.)	(50.)	EDWARDS	82	CBAL	0	J/PSI, GAM	2ETA	
M		1708.0	30.0	FRANKLIN	82	SMK2	E+E-, GAM	K+ K-		
M		1670.	50.	BLOOM	83	CBAL	0	J/PSI, GAM	2ETA	
M		1720.	10.	PERRIER	84	SMK3	J/PSI, GAM	K+ K-		
M		1713.	15.	PERRIER	84	SMK3	J/PSI, GAM	PI-PI-		
M	AVG	1716.0	7.9	AVERAGE						
M	E	FROM FIT NEGLECTING NEARBY F2*(1525). REPLACED BY BLOOM 83.								

### $f_2(1720)$ WIDTH (MeV)

W	E	(228.)	(100.)	(70.)	EDWARDS	82	CBAL	0	J/PSI, GAM	2ETA
W		156.0	60.0		FRANKLIN	82	SMK2	E+E-, GAM	K+ K-	
W		160.	80.		BLOOM	83	CBAL	0	J/PSI, GAM	2ETA
W		130.	20.		PERRIER	84	SMK3	J/PSI, GAM	K+ K-	
W	AVG	134.1	18.5	AVERAGE						
W	E	FROM FIT NEGLECTING NEARBY F2*(1525). REPLACED BY BLOOM 83.								

### $f_2(1720)$ PARTIAL DECAY MODES

P1	$f_2(1720) \rightarrow \eta \eta$	549+ 549	DECAY MASSES
P2	$f_2(1720) \rightarrow K \bar{K}$	498+ 498	
P3	$f_2(1720) \rightarrow \rho \rho$	769+ 769	
P4	$f_2(1720) \rightarrow \omega \omega$	783+ 783	
P5	$f_2(1720) \rightarrow \pi \pi$	140+ 140	

### REFERENCES FOR $f_2(1720)$

ALTHOFF 82 ZPHY C 16 13 -BOERNER, BURKHARDT+ (TASSO COLLABORATION)  
BARNES 82 PL 116 B 365 T. BARNES AND P. E. CLOSE (RHEL)  
BARNES 82 NP B 198 360 -CLOSE, MONAGHAN (RHEL+OXF)  
BURKE 82 PRL 49 632 -TRILLING, ABRAMS, ALAM, BLOCKER,+ (LBL-SLAC)  
EDWARDS 82 PRL 48 458 -PARTIDGE, PECK,+ (CIT+HARV-PRIN-STAN-SLAC)  
FRANKLIN 82 SLAC-254 M. E. B. FRANKLIN (SLAC)  
TANIMOTO 82 PL 116 B 198 M. TANIMOTO (BIEL)  
ALTHOFF 83 PL 121 B 216 -BRANDELIC, BOERNER+ (TASSO COLLABORATION)  
BARNETT 83 PL 120 B 455 -BLOCKUS, BURKA, CHIEN, CHRISTIAN+ (CHU)  
BLOOM 83 ARNS 35 143 BLOOM, PECK (SLAC-CIT)  
HITLIN 83 CORNELL CONF. 746 D. HITLIN (CIT)  
PERRIER 84 SLAC-PUB-3436 J. PERRIER (UCSC-SLAC)  
BALTRUSA 85 SLAC-PUB-3682 BALTRUSAITIS, (CIT-UCSC-ILL-SLAC-WASH)

$f_0(1730)$   
was  $S(1730)$

$$I^G(J^{PC}) = 0^+(0^{-+})$$

NAMED S\* PRIME BY ETKIN 82  
SEEN IN PHASE SHIFT ANALYSIS OF K0S K0S SYSTEM.  
NEEDS CONFIRMATION.

OMITTED FROM  
SUMMARY TABLE

### $f_0(1730)$ MASS (MeV)

M	A	S	1730.0	22.0	ETKIN2	82	MPS	0	23	PI-P, KSKS	N
M			1742.0	15.0	WILLIAMS	84	MPSF	200PI-N, KSKS	X		
M	AVERAGE MEANINGLESS										
M	A	FROM AN AMPLITUDE ANALYSIS OF THE K0S K0S SYSTEM.									
M	S	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.									

### $f_0(1730)$ WIDTH (MeV)

W	A	200.0	156.0	9.0	ETKIN1	82	MPS	0	23	PI-P, KSKS	N
W		57.0	38.0		WILLIAMS	84	MPSF	200PI-N, KSKS	X		
W	AVG	82.0	54.3	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)							
W	A	FROM AN AMPLITUDE ANALYSIS OF THE K0S K0S SYSTEM.									

### $f_0(1730)$ PARTIAL DECAY MODES

P1	$f_0(1730) \rightarrow K \bar{K}$	DECAY MASSES
		498+ 498

### REFERENCES FOR $f_0(1730)$

ETKIN1 82 PR D 25 1786 +FOLEY, LAI, LINDENBAUM+ (BNL-CUNY+TUFT-VAND)JP  
ETKIN2 82 PR D 25 2446 +FOLEY, LAI, LINDENBAUM+ (BNL-CUNY+TUFT-VAND)  
WILLIAMS 84 PR D 30 877 +DIAMOND+(VAND+NDAM+TUFT+ARIZ+FNAL+FSU+VPI)

$\pi(1770)$

$$I^G(J^{PC}) = 1^-(0^{-+})$$

OMITTED FROM  
SUMMARY TABLE  
SEEN IN PARTIAL WAVE ANALYSIS OF THE DIFFRACTIVELY  
PRODUCED 3PI SYSTEM.  
NEEDS CONFIRMATION.

### $\pi(1770)$ MASS (MeV)

M	1100	1770.	30.	BELLINI	82	SPEC	-	40	PI-A, 3PI	A
---	------	-------	-----	---------	----	------	---	----	-----------	---

### $\pi(1770)$ WIDTH (MeV)

W	1100	310.	50.	BELLINI	82	SPEC	-	40	PI-A, 3PI	A
---	------	------	-----	---------	----	------	---	----	-----------	---

### $\pi(1770)$ PARTIAL DECAY MODES

P1	$\pi(1770) \rightarrow f_0(1300) \pi$	1300+ 140	DECAY MASSES
P2	$\pi(1770) \rightarrow \rho \pi$	769+ 140	

### $\pi(1770)$ BRANCHING RATIOS

$\pi(1770) \rightarrow (f_0(1300) \pi)/total$	(P1)			
R1	DOMINANT	BELLINI 82 SPEC - 40	PI-A, 3PI	A
$\pi(1770) \rightarrow (\rho \pi)/total$	(P2)			
R2	NOT SEEN	BELLINI 82 SPEC - 40	PI-A, 3PI	A

### REFERENCES FOR $\pi(1770)$

BELLINI 82 PRL 48 1697 -FRABETTI, IVANSHIN, LITKIN+ (MILA-BGNA-JINR)

$f_2(1810)$   
was  $f(1810)$

$$I^G(J^{PC}) = 0^+(2^{++})$$

FORMERLY CALLED X(1850).  
FROM AN AMPLITUDE ANALYSIS OF THE K-K- SYSTEM SEEN  
IN PI- P INTO K- K- N AT 10 GEV/C. NOT CONFIRMED  
BY ETKIN 82. SEEN ALSO IN PI-PI- TO 2PI0 AMPLITUDE  
ANALYSIS (CASON 82), BUT NOT SEEN IN THE PARTIAL  
WAVE ANALYSIS OF THE PI-PI- SYSTEM.  
NEEDS CONFIRMATION.

OMITTED FROM  
SUMMARY TABLE

### $f_2(1810)$ MASS (MeV)

M	A	1857.0	35.0	24.0	COSTA	80	OMEG	0	10	PI-P, K+ K- N	
M		1799.0	15.0		CASON	82	STRC	0	8	PI-P, PI+2PI0	
M	AVERAGE MEANINGLESS (SCALE FACTOR = 1.8)										
M	A	ERROR INCREASED BY SPREAD OF TWO SOLUTIONS.									

### $f_2(1810)$ WIDTH (MeV)

W	A	185.0	102.0	135.0	COSTA	80	OMEG	0	10	PI-P, K+ K- N	
W		280.0	42.0	35.0	CASON	82	STRC	0	8	PI+P, PI+2PI0	
W	AVERAGE MEANINGLESS										
W	A	ERROR INCREASED BY SPREAD OF TWO SOLUTIONS.									

### $f_2(1810)$ PARTIAL DECAY MODES

P1	$f_2(1810) \rightarrow K^+ K^-$	494+ 494	DECAY MASSES
P2	$f_2(1810) \rightarrow \pi \pi$	140+ 140	







# Meson Full Listings

$a_3(2050)$ ,  $\pi_2(2100)$ ,  $\rho(2150)$

## $a_3(2050)$ BRANCHING RATIOS

$a_3(2050) \rightarrow (\rho_3(1690) \pi)/(3\pi)$  (P2)/(P1)  
 R1 DOMINANT KALELKAR 75 HBC - 15 P1-P, P 3P1

### REFERENCES FOR $a_3(2050)$

HUSON 68 PL 28 B 208	+LUBATTI, BELLINI, BINGHAM, + (ORSA-MILA-LBL)
BEMPORAD 71 NP B 33 397	-DUFÉY, COOLING, + (CERN+ETH+LOIC-MILA)
CLAYTON 72 NP B 47 81	-MASON, MUIRHEAD, RICOPOULOS, - (LIVP-PATR)
BASTIEN 73 UPPSALA CONF. 73	+DUNN, HARRIS, LUBATTI, BINGHAM, + (SEAT-UCB)
ORSEN 74 NP B71 189	-COOPER, FIELDS, RHINES, WHITMORE, - (ANL-OXF)
DEUTSCHM 75 NP B99 397	DEUTSCHMANN, + (ABBCCCH COLLABORATION)
KALELKAR 75 THESIS (NEVIS 207)	M.S. KALELKAR (COLU)
ANTIPOV 77 NP B 119 45	+BUSNELLO, DAMGAARD, KIENZLE - (CERN-SERP)
BALTAY 77 PRL 39 591	+CAUTIS, KALELKAR (COLUMBIA)JP
HARRIS 81 ZPHYS C 9 275	+DUNN, LUBATTI, MORIYASU, PODOLSKY, + (SEAT-UCB)
CAUTIS 77 THESIS NEVIS 221	C.V. CAUTIS (COLUMBIA)JP
BALTAY 78 PR D 17 52	-CAUTIS, COHEN, CSORNA, KALELKAR - (COLU-BING)

$\pi_2(2100)$   
was  $A(2100)$

$$I^G(J^{PC}) = 1^-(2^{--})$$

FORMERLY CALLED P1.  
 SEEN IN THE (RHO P1), (FO(1300) P1), AND (F2(1270) P1) JP = 2- WAVES OF THE DIFFRACTIVELY PRODUCED 3P1 SYSTEM. NEEDS CONFIRMATION.

OMITTED FROM SUMMARY TABLE

### $\pi_2(2100)$ MASS (MeV)

M L 2100. 150. DAUM 81 CNTR 63,94 P1 - P, 3P1
M L FROM A TWO RESONANCE FIT TO FOUR 2-0- WAVES.

### $\pi_2(2100)$ WIDTH (MeV)

W L 651. 50. DAUM 81 CNTR 63,94 P1 - P, 3P1
W L FROM A TWO RESONANCE FIT TO FOUR 2-0- WAVES.

### $\pi_2(2100)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$\pi_2(2100) \rightarrow 3\pi$	140- 140- 140
P2	$\pi_2(2100) \rightarrow \rho \pi$	769- 140
P3	$\pi_2(2100) \rightarrow f_2(1270) \pi$	1274- 140
P4	$\pi_2(2100) \rightarrow f_0(1300) \pi$	1300- 140

### $\pi_2(2100)$ BRANCHING RATIOS

$\pi_2(2100) \rightarrow (\rho \pi)/(3\pi)$ (P2)/(P1)	R1 L 0.19 0.05 DAUM 81 CNTR 63,94 P1 - P
$\pi_2(2100) \rightarrow (f_2(1270) \pi)/(3\pi)$ (P3)/(P1)	R2 L 0.36 0.09 DAUM 81 CNTR 63,94 P1 - P
$\pi_2(2100) \rightarrow (f_0(1300) \pi)/(3\pi)$ (P4)/(P1)	R3 L 0.45 0.07 DAUM 81 CNTR 63,94 P1 - P
<b>D-wave/S-wave RATIO FOR <math>\pi_2(2100) \rightarrow f_2(1270) \pi</math></b>	
4 L 0.39 0.23 DAUM 81 CNTR 63,94 P1 - P	
R L FROM A TWO RESONANCE FIT TO FOUR 2-0- WAVES.	

### REFERENCES FOR $\pi_2(2100)$

DAUM 81 NP B 182 269	+HERTZBERGER-(AMST-CERN+CRAC-MPII-OMX-RHEL)
----------------------	---

$\rho(2150)$

$$I^G(J^{PC}) = 1^-(1^{--})$$

OMITTED FROM SUMMARY TABLE

THIS ENTRY WAS PREVIOUSLY CALLED T1(2190). CONTAINS ONLY RESULTS FROM FORMATION EXPERIMENTS, FOR PRODUCTION EXPERIMENTS SEE THE NBAR N(1200-3600) ENTRY. SEE ALSO F2(2150), RHO3(2250), F4(2300), RHO5(2350)

OUR LATEST MINIREVIEW ON THIS PARTICLE CAN BE FOUND IN THE 1984 EDITION.

### $\rho(2150)$ MASS (MeV)

M	PBAR P INTO PI P1		
M	P (2100.0) APPROX.	MARTIN A 80 RVUE	
M	P (2170.0) APPROX.	MARTIN B 80 RVUE	
M	S CHANNEL NUCLEON ANTINUCLEON		
M	B 2190. 10. ABRAMS 70 CNTR S CHANNEL PBAR N		
M	I 2193. 2. ALPSPECTOR 73 CNTR S CHANNEL PBAR P		
M	E I 2155.0 15.0 COUPLAND 77 CNTR 0.7-2.4PB-P, PB-P		
M	I (2190.0) APPROX. CUTTS 78 CNTR .97-3. PB P, NB N		
M	AVERAGE MEANINGLESS		
M	B I=1, JP=1- FROM SIMULTANEOUS ANALYSIS OF P PB --> P1-P1- AND P10 P10		
M	B SEEN AS BUMP IN 1-1 STATE. SEE ALSO COOPER 68.		
M	P PEASLEE 75 CONFIRM PBAR P RESULTS OF ABRAMS 70, NO NARROW STRUCTURE		
M	E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.		
M	I ISOSPINS 0 AND 1 NOT SEPARATED		

### $\rho(2150)$ WIDTH (MeV)

W	PBAR P INTO PI P1		
W	P (200.0) APPROX.	MARTIN A 80 RVUE	
W	P (250.0) APPROX.	MARTIN B 80 RVUE	
W	S CHANNEL NUCLEON ANTINUCLEON		
W	B (85.) APPROX. ABRAMS 70 CNTR S CHANNEL PBAR N		
W	J 98. 8. ALPSPECTOR 73 CNTR S CHANNEL PBAR P		
W	E I 135.0 75.0 COUPLAND 77 CNTR 0.7-2.4PB-P, PB-P		
W	P I=1, JP=1- FROM SIMULTANEOUS ANALYSIS OF P PB --> P1-P1- AND P10 P10		
W	AVERAGE MEANINGLESS		
W	B SEE NOTE B ABOVE.		
W	E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.		
W	I ISOSPINS 0 AND 1 NOT SEPARATED		

### REFERENCES FOR $\rho(2150)$

ABRAMS 67 PRL 18 1209	+COOL, GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL)
COOPER 68 PRL 20 1059	-HYMAN, MANNER, MUSGRAVE, VOYVODIC (ANL)
BRICMAN 69 PL 29 B 451	+FERRO-LUZZI, BIZARD, - (CERN-CAEN-SACL)
ABRAMS 70 PR D 1 1917	-COOL, GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL)
BACON 71 NP B 32 66	-BUTTERWORTH, MILLER, PHELAN, + (RHEL-LIVP)
FIELDS 71 PRL 27 1749	+COOPER, RHINES, ALLISON (ANL-OXF)
YOH 71 PRL 26 922	+BARISH, CAROLL, LOBKOVICZ - (CIT+BNL-ROCH)
ALEXANDER 72 NP B 45 29	ALEXANDER, BAR-NIR, BEVARY, DAGAN, + (TELA)
DONALD 72 PL 40 B 586	+GALLETTY, EDWARDS, DE BILLY, + (LIVP-LPNP)
ALPSPECTO 73 PRL 30 511	ALPSPECTOR, COHEN, CVIJANOVICH, - (RUTG-UPN)J
BACON 73 PR D 7 577	+BUTTERWORTH, (RHEL-LIVP)
BETTINI 73 NC 15 A 563	+GARNJOST, BIGI, + (PADO-LBL-PISA-TORI)
DONALD 73 NP B 61 333	+EDWARDS, GIBBINS, BRIAND, DUBOC, + (LIVP-LPNP)
NICHOLSO 73 PR D 7 2572	NICHOLSON, DELORNE, CARROLL, + (CIT-ROCH+BNL)
BERTANZA 74 NC 23A 209	+BIGI, CASALI, LARICCIA, + (PISA-PADO-TORI)
HYAMS 74 NP B 73 202	+JONES, WEILHAMMER, BLUM, + (CERN-MPII)
DONNACHIE 75 NC 26 A 317	A. DONNACHIE, P. R. THOMAS (MANCHESTER)
EISENHAN 75 NP B 96 109	EISENHANDLER, GIBSON, + (LOQM+LIVP-DARE-RHEL)
HANDLER 75 NP B101 35	+JACQUES, JONES, PANDOULAS, - (RUTG-STEV-ALBA)
HUESMAN 75 NC 25A 91	+GARNJOST, ROSS, + (LBL-PADO-PISA-TORI)
PEASLEE 75 PL 57B 189	+DEMARZO, GUERRIERO, + (CANS+BARIL-BROU-MIT)
GAY 76 NC 31 A 593	+JEANNERET, BOGDANSKI, + (NEUC+LAUS+LIVP-LPNP)
ZEMANY 76 NP B 103 557	+MING MA, MOUNTZ, SMITH (MSU)
CARTER 1 77 PL 67 B 117	-COUPLAND, EISENHANDLER, ASTBURY, + (LOQM+RHEL)JP
CARTER 2 77 PL 67 B 122	A. A. CARTER (LOQM)JP
CARTER 3 77 NP B 127 202	-COUPLAND, ATKINSON, ARNISON, (LOQM-DARE-RHEL)
COUPLAND 77 PL 71 B 460	-EISENHANDLER, GIBSON, ASTBURY, - (LOQM+RHEL)
JONES 77 NP B 119 476	M. D. JONES, R. J. PLANO (RUTG)
MONTANET 77 BOSTON CONF. 260	L. MONTANET (CERN)
CARTER 1 78 NP B 132 176	A. A. CARTER (LOQM)JP
CARTER 2 78 NP B 141 467	A. A. CARTER (LOQM)
CUTTS 78 PR D 17 16	+GOOD, GRANNIS, GREEN, LEE, PITTMAN + (STON-WISC)
MARTIN 79 PL 86 B 93	A. D. MARTIN, M. R. PENNINGTON (DURH)
MARTIN A 80 NP B 169 216	A. D. MARTIN, M. R. PENNINGTON (DURH)JP
MARTIN B 80 NP B 176 355	B. R. MARTIN, J. MORGAN (LOUC+RHEL)JP



# Meson Full Listings

$\rho_3(2250)$ ,  $f_4(2300)$ ,  $\rho_5(2350)$

M J I=1, JP=3- FROM AMPLITUDE ANALYSIS.  
 M K I=0, 1, JP=3- FROM BARRELET ZERO ANALYSIS.  
 M P I=1, JP=3- FROM SIMULTANEOUS ANALYSIS OF P PB --> PI-PI+ AND P10 P10

S CHANNEL NUCLEON ANTINUCLEON

M B 2190. 10. ABRAMS 70 CNTR S CHANNEL PBAR N  
 M I 2193. 2. ALSPECTOR 73 CNTR S CHANNEL PBAR P  
 M E I 2155.0 15.0 COUPLAND 77 CNTR 0 .7-2.4PB-P,PB-P  
 M I (2190.0) APPROX. CUTTS 78 CNTR .97-3. PB P,N,N

M AVERAGE MEANINGLESS

M B SEEN AS BUMP IN I=1 STATE. SEE ALSO COOPER 68.  
 M B PEASLEE 75 CONFIRM PBAR P RESULTS OF ABRAMS 70, NO NARROW STRUCTURE  
 M E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.  
 M I ISOSPINS 0 AND 1 NOT SEPARATED

### $\rho_3(2250)$ WIDTH (MeV)

W PBAR P INTO PI PI OR K KB

W J (200.0) APPROX. CARTER 1 77 CNTR 0 .7-2.4PB P,PIPI  
 W K (150.0) APPROX. CARTER 2 78 CNTR 0 .7-2.4PB P,K-K+  
 W P (200.0) APPROX. MARTIN A 80 RVUE  
 W P (250.0) APPROX. MARTIN B 80 RVUE

M J I=1, JP=3- FROM AMPLITUDE ANALYSIS.  
 M K I=0, 1, JP=3- FROM BARRELET ZERO ANALYSIS.  
 M P I=1, JP=3- FROM SIMULTANEOUS ANALYSIS OF P PB --> PI-PI+ AND P10 P10

S CHANNEL NUCLEON ANTINUCLEON

W B (85.) APPROX. ABRAMS 70 CNTR S CHANNEL PBAR N  
 W I 98. 8. ALSPECTOR 73 CNTR S CHANNEL PBAR P  
 W E I 135.0 75.0 COUPLAND 77 CNTR 0 .7-2.4PB-P,PB-P

W AVERAGE MEANINGLESS

W B SEE NOTE B ABOVE.  
 W E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.  
 W I ISOSPINS 0 AND 1 NOT SEPARATED

### REFERENCES FOR $\rho_3(2250)$

ABRAMS 67 PRL 18 1209 +COOL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)  
 COOPER 68 PRL 20 1059 +HYMAN,WANNER,MUSGRAVE,VOJVODIC (ANL)  
 ABRAMS 70 PR D 1 1917 +COOL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)  
 FIELDS 71 PRL 27 1749 +COOPER,RHINES,ALLISON (ANL+OXF)  
 YOH 71 PRL 26 922 +BARISH,CAROLL,LOBKOVICZ+ (CIT-BNL-ROCH)

ALSPECTOR 73 PRL 30 511 ALSPECTOR,COHEN,CVIJANOVICH,+ (RUTG+UPNJ)  
 BETTINI 73 NC 15 A 563 -GARNJOST,BIGI,+ (PADO-LBL-PISA-TORI)  
 DONNACHI 73 LNC 7 285 A.DONNACHIE,P.R.THOMAS (MANCHESTER)  
 NICHOLSO 73 PR D 7 2572 NICHOLSON,DELORME,CARROLL,+ (CIT-ROCH-BNL)

BERTANZA 74 NC 23A 209 +BIGI,CASALI,LARICCIA,+ (PISA-PADO-TORI)  
 ZEMANY 76 NP B 103 537 +MING MA,MUNTZ,SMITH (MSU)

CARTER 1 77 PL 67 B 117 +COUPLAND,EISENHANDLER,ASTBURY,+ (LOQM+RHEL)JP  
 CARTER 2 77 PL 67 B 122 A.A.CARTER (LOQM)JP  
 CARTER 3 77 NP B 127 202 +COUPLAND,ATKINSON,ARNISON+(LOQM+DARE+RHEL)  
 COUPLAND 77 PL 71 B 460 -EISENHANDLER,GIBSON,ASTBURY,+ (LOQM+RHEL)  
 MONTANET 77 BOSTON CONF. 260 L.MONTANET (CERN)

CARTER 1 78 NP B 132 176 A.A.CARTER (LOQM)JP  
 CARTER 2 78 NP B 141 467 A.A.CARTER (LOQM)  
 CUTTS 78 PR D 17 16 +GOOD,GRANNIS,GREEN,LEE,PITTMAN+(STON+WISC)

MARTIN 79 PL 86 B 93 A.D. MARTIN,M.R. PENNINGTON (DURH)

MARTIN A 80 NP B 169 216 A.D. MARTIN,M.R. PENNINGTON (DURH)JP  
 MARTIN B 80 NP B 176 355 B.R.MARTIN,D.MORGAN (LOUC+RHEL)JP

$f_4(2300)$   
was  $\epsilon(2300)$

$$I^G(J^{PC}) = 0^{+}(4^{++})$$

THIS ENTRY WAS PREVIOUSLY CALLED U(2350). CONTAINS ONLY RESULTS FROM FORMATION EXPERIMENTS, FOR PRODUCTION EXPERIMENTS SEE THE NBARN(1200-3600) ENTRY. SEE ALSO RHO(2150), F2(2150), RHO5(2250), RHO5(2350).

### $f_4(2300)$ MASS (MeV)

M PBAR P INTO PI PI OR KB K

M J (2310.0) APPROX. CARTER 1 77 CNTR 0 .7-2.4PB P,PIPI  
 M K (2340.0) APPROX. CARTER 2 78 CNTR 0 .7-2.4PB P,K-K+  
 M J (2330.0) APPROX. DULUDEZ 78 OSPK 1.-2.PB P,P10P10  
 M P (2300.0) APPROX. MARTIN A 80 RVUE  
 M P (2500.0) APPROX. MARTIN B 80 RVUE

M J I=0, JP=4+ FROM AMPLITUDE ANALYSIS.  
 M K I=0, 1, JP=4+ FROM BARRELET ZERO ANALYSIS.  
 M P I=0, 1, JP=4+ FROM SIMULTANEOUS ANALYSIS OF P PB --> PI-PI+ AND P10 P10

S CHANNEL PBAR P OR NBAR N

M A 2375. 10. ABRAMS 70 CNTR S CHANNEL NBAR N  
 M I (2359.) (2.) ALSPECTOR 73 CNTR S CHANNEL PBAR P  
 M E I (2345.) (15.0) COUPLAND 77 CNTR 0 .7-2.4PB-P,PB-P  
 M I (2380.0) APPROX. CUTTS 78 CNTR .97-3. PB P,N,N

M E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.  
 M I ISOSPINS 0 AND 1 NOT SEPARATED

### $f_4(2300)$ WIDTH (MeV)

W PBAR P INTO PI PI OR KB K

W J (210.0) APPROX. CARTER 1 77 CNTR 0 .7-2.4PB P,PIPI  
 W K (150.0) APPROX. CARTER 2 78 CNTR 0 .7-2.4PB P,K-K+  
 W P (200.) APPROX. MARTIN A 80 RVUE

M J I=0, JP=4+ FROM AMPLITUDE ANALYSIS.  
 M K I=0, 1, JP=4+ FROM BARRELET ZERO ANALYSIS.  
 M P I=0, 1, JP=4+ FROM SIMULTANEOUS ANAL. OF P PBAR --> PI-PI+ AND P10 P10

W S CHANNEL PBAR P OR NBAR N

W I (190.) APPROX. ABRAMS 70 CNTR S CHANNEL NBAR N  
 W E I (155.0) (15.0) (8.) COUPLAND 77 CNTR 0 .7-2.4PB-P,PB-P  
 W I (155.0) (150.0) (65.0) COUPLAND 77 CNTR 0 .7-2.4PB-P,PB-P

W E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.  
 W I ISOSPINS 0 AND 1 NOT SEPARATED

### REFERENCES FOR $f_4(2300)$

BRICMAN 69 PL 29 B 451 +FERRO-LUZZI,BIZARD,+ (CERN+CAEN-SACL)  
 ABRAMS 70 PR D 1 1917 -COOL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)  
 FIELDS 71 PRL 27 1749 -COOPER,RHINES,ALLISON (ANL+OXF)  
 YOH 71 PRL 26 922 -BARISH,CAROLL,LOBKOVICZ+ (CIT-BNL-ROCH)

EASTMAN 72 NP B 51 29 +MING MA,OH,PARKER,SMITH,SPRAKKA (MSU)

ALSPECTOR 73 PRL 30 511 ALSPECTOR,COHEN,CVIJANOVICH,+ (RUTG+UPNJ)  
 DONNACHI 73 LNC 7 285 A.DONNACHIE,P.R.THOMAS (MANCHESTER)  
 NICHOLSO 73 PR D 7 2572 NICHOLSON,DELORME,CARROLL,+ (CIT-ROCH+BNL)

HYAMS 74 NP B 73 202 +JONES,WELHAMMER,BLUM,+ (CERN+MPIM)  
 MING MA 74 NP B68 214 +MOUNTZ,ZEMANY,SMITH (MICH)

DONNACHI 75 NC 26 A 317 A.DONNACHIE,P.R.THOMAS (MANCHESTER)  
 EISENHAN 75 NP B 96 109 EISENHANDLER,GIBSON,+ (LOQM+LIVP+DARE+RHEL)

CARTER 1 77 PL 67 B 117 +COUPLAND,EISENHANDLER,ASTBURY,+ (LOQM+RHEL)JP  
 CARTER 2 77 PL 67 B 122 A.A.CARTER (LOQM)JP  
 CARTER 3 77 NP B 127 202 +COUPLAND,ATKINSON,ARNISON+(LOQM+DARE+RHEL)  
 COUPLAND 77 PL 71 B 460 -EISENHANDLER,GIBSON,ASTBURY,+ (LOQM+RHEL)  
 MONTANET 77 BOSTON CONF. 260 L.MONTANET (CERN)

CARTER 1 78 NP B 132 176 A.A.CARTER (LOQM)JP  
 CARTER 2 78 NP B 141 467 A.A.CARTER (LOQM)  
 CUTTS 78 PR D 17 16 +GOOD,GRANNIS,GREEN,LEE,PITTMAN+(STON+WISC)  
 DULUDE1 78 PL 79 B 329 -LANDOU,MASSIMO,PEASLEE+ (BROW+MIT+BARI)JP  
 DULUDE2 78 PL 79 B 335 +LANDOU,MASSIMO,PEASLEE+ (BROW+MIT+BARI)JP

MARTIN 79 PL 86 B 93 A.D. MARTIN,M.R. PENNINGTON (DURH)

BOWCOCK 80 LNC 28 21 J.E.BOWCOCK,D.C.HODGSON (BIRM)  
 MARTIN A 80 NP B 169 216 A.D. MARTIN,M.R. PENNINGTON (DURH)JP  
 MARTIN B 80 NP B 176 355 B.R.MARTIN,D.MORGAN (LOUC+RHEL)JP

$\rho_5(2350)$   
was  $\rho(2350)$

$$I^G(J^{PC}) = 1^{+}(5^{-})$$

THIS ENTRY WAS PREVIOUSLY CALLED U(2400). CONTAINS ONLY RESULTS FROM FORMATION EXPERIMENTS, FOR PRODUCTION EXPERIMENTS SEE THE NBARN(1200-3600) ENTRY. SEE ALSO RHO(2150), F2(2150), RHO5(2250), F4(2300).

### $\rho_5(2350)$ MASS (MeV)

M PBAR P INTO PI PI OR KB K

M J (2480.0) APPROX. CARTER 1 77 CNTR 0 .7-2.4PB P,PIPI  
 M K (2500.0) APPROX. CARTER 2 78 CNTR 0 .7-2.4PB P,K-K+  
 M P (2250.0) APPROX. MARTIN A 80 RVUE  
 M P (2300.0) APPROX. MARTIN B 80 RVUE

M J I=1, JP=5- FROM AMPLITUDE ANALYSIS.  
 M K I=0, 1, JP=5- FROM BARRELET ZERO ANALYSIS.  
 M P I=1, JP=5- FROM SIMULTANEOUS ANALYSIS OF P PB --> PI-PI+ AND P10 P10

S CHANNEL NUCLEON ANTINUCLEON

M A 2350. 10. ABRAMS 70 CNTR S CHANNEL NBAR N  
 M N (2360.0) (25.0) OH 70 HDBC -OPBAR(P,N),K\*K2P1  
 M I (2359.0) (2.) ALSPECTOR 73 CNTR S CHANNEL PBAR P  
 M E I (2345.0) (15.0) COUPLAND 77 CNTR 0 .7-2.4PB-P,PB-P  
 M I (2380.0) APPROX. CUTTS 78 CNTR .97-3. PB P,N,N

M A FOR I=1 NBAR N  
 M E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.  
 M I ISOSPINS 0 AND 1 NOT SEPARATED  
 M N NO EVIDENCE FOR THIS BUMP SEEN IN THE PBAR P DATA OF CHAPMAN 71  
 M N NARROW STATE NOT CONFIRMED BY OH 73 WITH MORE DATA.

### $\rho_5(2350)$ WIDTH (MeV)

W PBAR P INTO PI PI OR KB K

W J (210.0) APPROX. CARTER 1 77 CNTR 0 .7-2.4PB P,PIPI  
 W K (150.0) APPROX. CARTER 2 78 CNTR 0 .7-2.4PB P,K-K+  
 W P (300.0) APPROX. MARTIN A 80 RVUE  
 W P (250.0) APPROX. MARTIN B 80 RVUE

M J I=1, JP=5- FROM AMPLITUDE ANALYSIS.  
 M K I=0, 1, JP=5- FROM BARRELET ZERO ANALYSIS.  
 M P I=1, JP=5- FROM SIMULTANEOUS ANALYSIS OF P PB --> PI-PI+ AND P10 P10

For notation, see key on page 91.

# Meson Full Listings

$\rho_5(2350)$ ,  $a_6(2450)$ ,  $f_6(2510)$ ,  $e^+e^-(1100-2200)$ ,  $\bar{N}N(1200-3600)$

S CHANNEL NUCLEON ANTINUCLEON  
 W W (140.) APPROX. ABRAMS 67 CNTR S CHANNEL PBAR N  
 W N (60.0) OR LESS OH 70 HOBC -OPBAR(P,N),K\*KZI  
 W I (165.) (18.) (8.) ALSPECTOR 73 CNTR S CHANNEL PBAR P  
 W EI (135.0) (150.0) (65.0) COUPLAND 77 CNTR 0 .7-2.4PB-P,PB-P  
 W E FROM A FIT TO THE TOTAL ELASTIC CROSS SECTION.  
 W I ISOSPINS 0 AND 1 NOT SEPARATED  
 W N NO EVIDENCE FOR THIS BUMP SEEN IN THE PBAR P DATA OF CHAPMAN 71  
 W N NARROW STATE NOT CONFIRMED BY OH 73 WITH MORE DATA.

## $f_6(2510)$ WIDTH (MeV)

W 240.0 60.0 BINON 84 SPEC 0 23 PI-P,N 2PI0

## $f_6(2510)$ PARTIAL DECAY MODES

P1  $f_6(2510) \rightarrow \pi\pi$  DECAY MASSES  
 135+ 135

## REFERENCES FOR $\rho_5(2350)$

ABRAMS 67 PRL 18 1209 +COOL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)  
 BRICHAN 69 PL 29 B 451 +FERRO-LUZZI,BIZARD,+ (CERN-CAEN+SACL)  
 CASO 69 LNC 3 707 +CONTE,BENZ,+ (GENO+DESY+HAMB-MILA+SACL)  
 ABRAMS 70 PR D 1 1917 +COOL,GIACOMELLI,KYCIA,LEONTIC,LI,+ (BNL)  
 OH 70 PRL 24 1257 +PARKER,EASTMAN,SMITH,SPRAFKA,MA (MSU)  
 CHAPMAN 71 PR D 4 1275 +GREEN,LVS,MURPHY,RING,+ (MICH)  
 FIELDS 71 PRL 27 1749 +COOPER,RHINES,ALLISON (ANL-ORF)  
 YOH 71 PRL 26 922 +BARISH,CARROLL,LOBKOVICZ+ (CIT+BNL+ROCH)  
 EASTMAN 72 NP B 51 29 +MING MA,OH,PARKER,SMITH,SPRAFKA (MSU)  
 MING MA 72 NP B 51 77 +EASTMAN,OH,PARKER,SMITH,SPRAFKA (MSU)  
 OH 72 NP B 51 57 +EASTMAN,MING MA,PARKER,SMITH,+ (MSU)  
 ALSPECTOR 73 PRL 30 511 ALSPECTOR,COHEN,CVJANOVICH,+ (RUTG-UPN)  
 NICHOLSO 73 PR D 7 2572 NICHOLSON,DELORME,CARROLL,+ (CIT+ROCH-BNL)  
 HYAMS 74 NP B 73 202 +JONES,WEILHAMMER,BLUM,+ (CERN+MPIM)  
 MING MA 74 NP B 68 214 +MOUNT,ZEMANY,SMITH (MICH)  
 DONNACHI 75 NC 26 A 317 A.DONNACHIE,P.R.THOMAS (MANCHESTER)  
 EISENHAN 75 NP B 96 109 EISENHANDLER,GIBSON,+ (LOQM-LIVP+DARE+RHEL)  
 CARTER 1 77 PL 67 B 117 +COUPLAND,EISENHANDLER,ASTBURY,+ (LOQM+RHEL)JP  
 CARTER 2 77 PL 67 B 122 A.A.CARTER (LOQM)JP  
 CARTER 3 77 NP B 127 202 +COUPLAND,ATKINSON,ARNISON-(LOQM+DARE+RHEL)  
 COUPLAND 77 PL 71 B 460 +EISENHANDLER,GIBSON,ASTBURY,+ (LOQM+RHEL)  
 MONTANET 77 BOSTON CONF. 260 L.MONTANET (CERN)  
 CARTER 1 78 NP B 132 176 A.A.CARTER (LOQM)JP  
 CARTER 2 78 NP B 141 467 A.A.CARTER (LOQM)  
 CUTTS 78 PR D 17 16 +GOOD,GRANNIS,GREEN,LEE,PITTMAN-(STON-WISC)  
 MARTIN 79 PL 86 B 93 A.D. MARTIN,M.R. PENNINGTON (DURH)  
 BOWDOCK 80 LNC 28 21 J.E. BOWDOCK,D.C. HODGSON (BIRM)  
 MARTIN A 80 NP B 169 216 A.D. MARTIN,M.R. PENNINGTON (DURH)JP  
 MARTIN B 80 NP B 176 355 B.R. MARTIN,D. MORGAN (LOUC-RHEL)JP

## REFERENCES FOR $f_6(2510)$

BINON 84 LNC 39 41 +DONSKOV,DUTEIL,GOUANERE+ (SERP-BELG-LAPP) JP

## $e^+e^-(1100-2200)$ $I^G(J^{PC}) = ?(1^{--})$

OMITTED FROM SUMMARY TABLE THIS ENTRY CONTAINS NON-STRANGE VECTOR MESONS COUPLED TO  $e^+e^-$  (PHOTON) BETWEEN PHI AND J/PSI MASS REGION. SEE ALSO RHO(1250), RHO(1600), AND PHI(1680).

## $e^+e^-(1100-2200)$ MASSES AND WIDTHS (MeV)

M (1097.0) (16.0) (19.0) BARTALUCC 79 OSPK 7 GAM P,E+ E- P  
 W (31.0) (24.0) (20.0) BARTALUCC 79 OSPK 7 GAM P,E+ E- P  
 W (1830.0) APPROX. PETERSON 78 SPEC GAM P,K- K- P  
 W (120.0) APPROX. PETERSON 78 SPEC GAM P,K- K- P  
 M C (2130.) APPROX. ESPOSITO 78 FRAM E+E-,K\*(892)+..  
 W C (30.) APPROX. ESPOSITO 78 FRAM E+E-,K\*(892)+..  
 M A (1820.) APPROX. SPINETTI 79 RVUE E+E-,4 PI+- 2GAM  
 W A (30.) APPROX. SPINETTI 79 RVUE E+E-,4 PI+- 2GAM  
 A INTEGRATED CROSS-SECTION OF BACCI 77, BARBIELINI 77, ESPOSITO 77.  
 C NOT SEEN BY DELCOURT 79.

## REFERENCES FOR $e^+e^-(1100-2200)$

BACCI 75 PL 58 B 481 +BIDOLI,PENSO,STELLA,BALDINI,+ (ROMA+FRAS)  
 BACCI 76 PL 64 B 356 +BIDOLI,PENSO,STELLA,BALDINI,+ (ROMA+FRAS)  
 BACCI 77 PL B 68 393 +DE ZORZI,PENSO,STELLA,BALDINI,+ (ROMA+FRAS)  
 BARBIELI 77 PL B 68 397 BARBIELINI, BARLETTA,+ (FRAS+NAPL+PISA+SANI)  
 BARTALUCC 77 NC A 39 374 BARTALUCCI, BERTOLUCCI, BRADASCHIA (DEST+FRAS)  
 ESPOSITO 77 PL B 68 389 +FELICETTI, MARINI,+ (FRAS+NAPL+PADO+ROMA)  
 AMBROSIO 78 PL 80 B 141 +CERRITO,BEMPORAD,BROSCO,+ (NAPL+PISA+ROMA)  
 BALDINI 78 PL 78 B 167 +BATTISTONI,CAPON,BACCI,DEZORZI-(FRAS+ROMA)  
 ESPOSITO 78 LNC 22 305 ESPOSITO,FELICETTI+ (FRAS+NAPL+PADO+ROMA)  
 ESPOSITO 78 LNC 23 604 ESPOSITO,FELICETTI+ (FRAS+NAPL+PADO+ROMA)  
 PETERSON 78 PR D 18 3955 +DIXON,EHRlich,GALIK,LARSON+ (CORN+HARV)  
 BARTALUCC 79 NC 49 A 207 BARTALUCCI,BASINI,BERTOLUCCI+ (DEST+FRAS)  
 DELCOURT 79 BATAVIA CONF. 499 +BERTRAND,BISELLO,BIZOT,BUON,CORDIER-(LALO)  
 ESPOSITO 79 LNC 25 5 +MARINI,PALLOTTA+ (FRAS+UMD+PADO+ROMA)  
 SPINETTI 79 BATAVIA CONF. 506 M. SPINETTI (FRAS)  
 BALDINI 81 LNC 30 337 +BATTISTONI,CAPON,BACCI,DEZORZI+(FRAS+ROMA)

## $a_6(2450)$ was $\delta(2450)$

$I^G(J^{PC}) = 1^-(6^{--})$

SEEN IN PARTIAL WAVE ANALYSIS OF THE K KBAR SYSTEM. NEEDS CONFIRMATION.

OMITTED FROM SUMMARY TABLE

## $a_6(2450)$ MASS (MeV)

M C 2450. 130. CLELAND 82 SPEC +- 50 PI P,KS K+-P  
 M C FROM AN AMPLITUDE ANALYSIS

## $a_6(2450)$ WIDTH (MeV)

W C 400. 250. CLELAND 82 SPEC +- 50 PI P,KS K+-P  
 M C FROM AN AMPLITUDE ANALYSIS

## $a_6(2450)$ PARTIAL DECAY MODES

P1  $a_6(2450) \rightarrow K\bar{K}$  DECAY MASSES  
 498- 498

## REFERENCES FOR $a_6(2450)$

CLELAND 82 NP B 208 228 +DELFOSSÉ,DORSAZ,GLOOR(DURH+GEVA-LAUS-PITT)

## $f_6(2510)$ was $r(2510)$

$I^G(J^{PC}) = 0^+(6^{++})$

SEEN IN  $\pi^0\pi^0$ . NEEDS CONFIRMATION.

OMITTED FROM SUMMARY TABLE

## $f_6(2510)$ MASS (MeV)

M 2510.0 30.0 BINON 84 SPEC 0 38 PI-P,N 2PI0

## $\bar{N}N(1200-3600)$

OMITTED FROM SUMMARY TABLE THIS ENTRY CONTAINS VARIOUS HIGH MASS, NON-STRANGE STRUCTURES COUPLED TO THE BARYON-ANTIBARYON SYSTEM AS WELL AS QUASI-NUCLEAR BOUND STATES BELOW THRESHOLD.

SEE ALSO X(1935), RHO(2150), F2(2150), RHO3(2250), F4(2300), RHOS(2350). EVIDENCE FOR STRUCTURES COUPLED TO THE ANTIHYPERON NUCLEON (OR C.C.) SYSTEM IS LISTED UNDER K(2200).

## $\bar{N}N(1200-3600)$ MASSES AND WIDTHS (MeV)

M W G 1210. 5.0 RICHTER 83 CNTR 0 STOPPED PBAR S  
 M W G (1395.) PAVLOPOUL 78 CNTR STOPPED PBAR S  
 M 1637.1 5.6 7.3 ADIELS 84 CNTR PBAR HELIUM I  
 M W G 1638. 3.0 RICHTER 83 CNTR 0 STOPPED PBAR S  
 M 1644.0 5.6 7.3 ADIELS 84 CNTR PBAR HELIUM I  
 M W G (1646.) PAVLOPOUL 78 CNTR STOPPED PBAR S  
 M W G (1684.) PAVLOPOUL 78 CNTR STOPPED PBAR S  
 M 1687.1 5.0 4.3 ADIELS 84 CNTR PBAR HELIUM I



For notation, see key on page 91.

Meson Full Listings

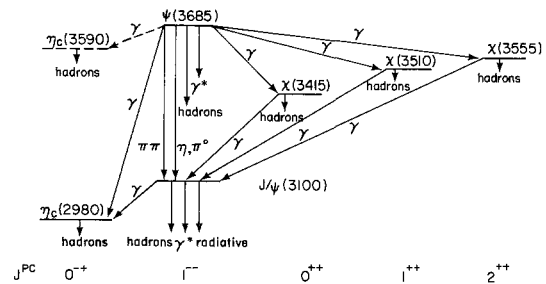
X(1900-3600), CHARMONIUM, η<sub>c</sub>(2980)

X(1900-3600) MASSES AND WIDTHS (MeV)

Table listing mesons with columns for mass, width, and experimental references. Includes entries for ALDE, THOMPSON, BOESEBECK, CHLIAPNIK, CASO, KRAMER, TAKAHASHI, AJINENKO, BALTHAY, BAUD, ANDERSON, DENNEY, SABAU, YOST, and BAUD.

Table listing mesons with columns for mass, width, and experimental references. Includes entries for BLANAR, CLINE, AJINENKO, ASTON, BARTH, DENNEY, ATKINSON, ALDE, and DOVER.

THE CHARMONIUM SYSTEM



The current state of knowledge of the charmonium system and transitions, as interpreted by the charmonium model. Uncertain states and transitions are indicated by dashed lines. The notation γ\* refers to decay processes involving intermediate virtual photons, including decays to e+e- and μ+μ-.

ηc(2980) I<sup>G</sup>(J<sup>PC</sup>) = 0<sup>+</sup>(0<sup>-+</sup>) OBSERVED IN THE INCLUSIVE GAMMA SPECTRUM GENERATED FROM PSI(3685) DECAY, THEREFORE C++, FROM THE 4P1 DECAY G++, THEREFORE I=0. FROM ANGULAR DISTRIBUTION IN J/PSI TO ETA/C, ETA/C TO PHI PHI, J<sup>PC</sup>=0- (BALTRUSAITIS 84).

ηc(2980) MASS (MeV) table listing masses and widths for different experimental groups: HIMEL, PARTRIDGE, BALTRUSAITIS, BALTRUSAITIS, GAISER, and AVERAGE.

REFERENCES FOR X(1900-3600)

References for X(1900-3600) listing authors and publications: CLAYTON, BOESEBECK, ANDERSON, BAUD, KRAMER, SABAU, YOST, TAKAHASHI, THOMPSON, BALTHAY, KALELKAR, and KEMP.

ηc(2980) WIDTH (MeV) table listing widths and experimental references: HIMEL, PARTRIDGE, BALTRUSAITIS, BALTRUSAITIS, GAISER, and AVERAGE.







# Meson Full Listings

## $J/\psi(3097)$

<b><math>J/\psi(3097) \rightarrow (K^+ K^-)/total (units 10^{-4})</math></b> (P13)	<b><math>J/\psi(3097) \rightarrow (\phi \eta)/total (units 10^{-3})</math></b> (P37)
R13 2 2.0 1.6 VANNUCCI 77 SMAG E=E-,K-K-	R37 5 1.0 0.6 VANNUCCI 77 SMAG E=E-
R13 6 2.2 0.9 BRANDELIK 79 DASP E=E-,K+K-	R37 0.69 0.11 BALTRUS2 85 SMK3 K+K-NEUTRALS
R13S 110 2.39 0.33 BALTRUSAS 85 SMK3 E=E-,K+K-	R37 0.64 0.17 BALTRUS2 85 SMK3 E+E-,K+K-ETA
R13S SYSTEMATIC ERROR ADDED QUADRATICALLY BY US	R37 0.61 0.16 BALTRUS2 85 SMK3 E+E-,K+K-SPI
R13 R37	R37
R13 AVG 2.35 0.30 AVERAGE	R37 AVG 0.665 0.079 AVERAGE
<b><math>J/\psi(3097) \rightarrow (K_S^0 \bar{K}_S^0)/total (units 10^{-4})</math></b> (P14)	<b><math>J/\psi(3097) \rightarrow (\phi \eta')/total (units 10^{-3})</math></b> (P38)
R14C (0.052)OR LESS CL=0.90 BALTRUSAI 85 SMK3 E=E-,HADRONS	R38 (1.3) OR LESS CL=0.90 VANNUCCI 77 SMAG E=E-
R14C FORBIDDEN BY CP.	R38 0.39 0.12 BALTRUS2 85 SMK3 K+K-RHO GAMMA
<b><math>J/\psi(3097) \rightarrow (K \bar{K} \pi)/total (units 10^{-4})</math></b> (P15)	R38 0.365 0.064 BALTRUS2 85 SMK3 K+K-ETA P1+P1-
R15 126 78.0 21.0 VANNUCCI 77 SMAG E=E-,KS K+- P1-	R38 AVG 0.371 0.056 AVERAGE
R15 25 55.2 12.0 FRANKLIN 83 SMK2 E=E-,K- K- P10	
R15 R38	
R15 AVG 60.8 10.2 AVERAGE	<b><math>J/\psi(3097) \rightarrow (\phi F_2(1270))/total (units 10^{-4})</math></b> (P39)
<b><math>J/\psi(3097) \rightarrow (\pi^+ \pi^- K^+ K^-)/total</math></b> (P16)	R39 (3.7) OR LESS CL=0.90 VANNUCCI 77 SMAG E+E-
R16 205 0.0072 0.0023 VANNUCCI 77 SMAG E=E-	<b><math>J/\psi(3097) \rightarrow (\phi F_2'(1525))/total (units 10^{-4})</math></b> (P40)
<b><math>J/\psi(3097) \rightarrow (2(\pi^+ \pi^-) K^+ K^-)/total</math></b> (P17)	R40B 6 8.0 5.0 VANNUCCI 77 SMAG E+E-
R17 30 0.0031 0.0013 VANNUCCI 77 SMAG E=E-	R40B 46 3.4 1.3 GIDAL 81 SMK2 E+E-
<b><math>J/\psi(3097) \rightarrow (\pi^+ \pi^- \pi^0 K^+ K^-)/total</math></b> (P18)	R40B ASSUMES P2'(1525) INTO K KBAR IS 100 PER CENT.
R18 309 0.012 0.003 VANNUCCI 77 SMAG E+E-	R40 5.7 1.3 AVERAGE
<b><math>J/\psi(3097) \rightarrow 2(K^+ K^-)/total</math></b> (P19)	<b><math>J/\psi(3097) \rightarrow (\phi F_0(975))/total (units 10^{-4})</math></b> (P41)
R19 0.0007 0.0003 VANNUCCI 77 SMAG E=E-	R41 50 2.6 0.6 GIDAL 81 SMK2 E+E-
<b><math>J/\psi(3097) \rightarrow (\rho \pi)/total</math></b> (P20)	<b><math>J/\psi(3097) \rightarrow (\pi^\pm a_2)/total</math></b> (P42)
R20 543 0.010 0.002 BARTEL 1 76 CNTR E=E-	R42 (0.0043)OR LESS CL=0.90 BRAUNSWH 76 DASP E+E-
R20 153 0.013 0.003 JEAN-MARI 76 SMAG E=E-	<b><math>J/\psi(3097) \rightarrow (\rho a_2(1320))/total</math></b> (P43)
R20 183 0.016 0.004 ALEXANDER 78 PLUT E+E-	R43 36 0.0084 0.0045 VANNUCCI 77 SMAG E=E-
R20 0.0133 0.0021 BRANDELIK 78 DASP E+E-,PI-PI-GAMMA	<b><math>J/\psi(3097) \rightarrow (K \bar{K}^*(892) + C.C.)/total (units 10^{-4})</math></b> (P44)
R20 150 0.013 0.003 FRANKLIN 83 SMK2 E+E-,HADRONS	R44 39 82. 24. BRAUNSWH 76 DASP E=E-,K+- K*BAR--
R20 0.0133 0.0015 BALTRUS2 85 SMK3 E+E-,PI0PI+PI-	R44 48 64. 12. VANNUCCI 77 SMAG E+E-,K+ K*BAR--
R20 AVG 0.01270 0.00051 AVERAGE	R44 45 54. 12. VANNUCCI 77 SMAG E+E-,K0 K*BAR0
<b><math>J/\psi(3097) \rightarrow (\rho^0 \pi^0)/(\rho \pi)</math></b> (P21)/P(20)	R44 24 32.8 11.4 FRANKLIN 83 SMK2 E+E-,K+- K*BAR--
R21 0.39 0.11 BARTEL 1 76 CNTR E+E-	R44A 100. 13. BALTRUS2 85 SMK3 E+E-,K+-K*+
R21 0.37 0.09 JEAN-MARI 76 SMAG E=E-	R44B 120. 15. BALTRUS2 85 SMK3 E+E-,K-K*-
R21 0.35 0.08 ALEXANDER 78 PLUT E+E-	R44C 78. 15. BALTRUS2 85 SMK3 E+E-,K0 K*0
R21 0.32 0.08 BRANDELIK 78 DASP E+E-,PI-PI-GAMMA	R44A FROM K*+ K- - C.C. TO FINAL STATE K+- KS P1- AND ISOSPIN INV.
R21 (0.36) (0.03) SCHARRE 79 SMAG E=E-	R44B FROM K*+ K- - C.C. TO FINAL STATE K- K- P10 AND ISOSPIN INV.
R21 R21	R44C FROM K*0 K0BAR + C.C TO FINAL STATE K+- KS P1- AND ISOSPIN INV.
R21 AVG 0.352 0.044 AVERAGE	R44 AVG 74.6 9.3 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8)
<b><math>J/\psi(3097) \rightarrow (\rho \eta)/total (units 10^{-3})</math></b> (P22)	<b><math>J/\psi(3097) \rightarrow (K \bar{K}^*(1430))/total</math></b> (P45)
R22 0.18 0.04 BALTRUS2 85 SMK3 E=E-,PI+PI-ETA	R45 (0.006)OR LESS CL=0.90 BRAUNSWH 76 DASP E+E-,K+-K2*BAR-
<b><math>J/\psi(3097) \rightarrow (\rho \eta')/total (units 10^{-3})</math></b> (P23)	R45 (0.004) OR LESS CL=0.90 VANNUCCI 77 SMAG E+E-,K0 K2*BAR0
R23 (0.1) OR LESS BALTRUS2 85 SMK3 E+E-,4PI ETA	<b><math>J/\psi(3097) \rightarrow (K^*(892)^0 \bar{K}^*(892)^0)/total</math></b> (P46)
<b><math>J/\psi(3097) \rightarrow (\omega \pi^+ \pi^-)/total (units 10^{-3})</math></b> (P24)	R46 (0.0005)OR LESS CL=0.90 VANNUCCI 77 SMAG E=E-
R24 215 7.8 1.6 BURMESTER 77 PLUT E+E-	<b><math>J/\psi(3097) \rightarrow (K_1^*(1430)^0 \bar{K}_1^*(1430)^0)/total</math></b> (P47)
R24 348 6.8 1.9 VANNUCCI 77 SMAG E+E-	R47 (0.0029)OR LESS CL=0.90 VANNUCCI 77 SMAG E+E-
R24 AVG 7.4 1.2 AVERAGE	<b><math>J/\psi(3097) \rightarrow (K^*(892)^0 \bar{K}_2^*(1430)^0 + C.C.)/total</math></b> (P48)
<b><math>J/\psi(3097) \rightarrow (\omega \pi \pi)/(2(\pi^+ \pi^-) \pi^0)</math></b> (P24)/(P9)	R48 40 0.0067 0.0026 VANNUCCI 77 SMAG E+E-
R25K J (0.3) JEAN-MARI 76 SMAG E=E-	<b><math>J/\psi(3097) \rightarrow (b_1(1235)^\pm \pi^\mp)/total</math></b> (P49)
R25 J FINAL STATE 2(P1+ PI-)PI0	R49 87 0.0029 0.0007 BURMESTER 77 PLUT E+E-
R25K UNDER THE ASSUMPTION THAT PI PI IS ISOSPIN 0	<b><math>J/\psi(3097) \rightarrow (\bar{p} p)/total (units 10^{-3})</math></b> (P50)
<b><math>J/\psi(3097) \rightarrow (\omega 2\pi^+ 2\pi^-)/total</math></b> (P26)	R50 2.0 0.5 BESCH 78 BONA E+E-
R26 140 0.0085 0.0034 VANNUCCI 77 SMAG E=E-	R50 A 351 2.2 0.2 PERUZZI 78 SMAG E+E-
<b><math>J/\psi(3097) \rightarrow (\omega K \bar{K})/total</math></b> (P27)	R50 133 2.5 0.4 BRANDELIK 79 DASP E=E-
R27 22 0.0016 0.0010 FELDMAN 77 SMAG E+E-	R50S 1420 2.16 0.17 EATON 84 SMK2 E=E-,HADRONS GAM
<b><math>J/\psi(3097) \rightarrow (\omega F_2(1270))/total</math></b> (P28)	R50S SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.
R28 81 0.0019 0.0008 VANNUCCI 77 SMAG E=E-	R50 A ASSUMING ANGULAR DISTRIBUTION (1.+COS(THETA)**2)
R28 70 0.0040 0.0016 BURMESTER 77 PLUT E=E-	R50 AVG 2.20 0.12 AVERAGE
R28 R28	<b><math>J/\psi(3097) \rightarrow (\bar{p} p)/(\mu^+ \mu^-)</math></b> (P50)/(P2)
R28 AVG 0.00232 0.00084 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)	R51 A 20 .051 .02 WIK 75 PLUT E+E-
<b><math>J/\psi(3097) \rightarrow (\omega F_2'(1525))/total (units 10^{-4})</math></b> (P29)	R51 A ASSUMING ANGULAR DISTRIBUTION (1.-COS(THETA)**2)
R29 (1.6) OR LESS CL=0.90 VANNUCCI 77 SMAG E+E-	<b><math>J/\psi(3097) \rightarrow (p \bar{p} \pi^0)/total (units 10^{-3})</math></b> (P52)
<b><math>J/\psi(3097) \rightarrow (\omega \eta)/total (units 10^{-3})</math></b> (P30)	R52 109 1.00 0.15 PERUZZI 78 SMAG E=E-,P PB
R30 1.9 0.4 BALTRUS2 85 SMK3 E=E-,3PI ETA	R52S 1.4 0.4 BRANDELIK 79 DASP E=E-
<b><math>J/\psi(3097) \rightarrow (\omega \eta')/total (units 10^{-3})</math></b> (P31)	R52S 685 1.13 0.13 EATON 84 SMK2 E=E-,HADRONS GAM
R31 0.39 0.13 BALTRUS2 85 SMK3 3PI RHO GAMMA	R52S SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.
R31 0.43 0.20 0.23 BALTRUS2 85 SMK3 PI0PI-PI-ETA2P1	R52 R52
R31 AVG 0.40 0.11 AVERAGE	R52 AVG 1.093 0.095 AVERAGE
<b><math>J/\psi(3097) \rightarrow (\omega \pi^0)/total (units 10^{-3})</math></b> (P32)	<b><math>J/\psi(3097) \rightarrow (p \bar{p} \pi^-)/total (units 10^{-3})</math></b> (P53)
R32 0.67 0.13 BALTRUS2 85 SMK3 PI0PI-PI- PI0	R53 194 2.16 0.29 PERUZZI 78 SMAG E=E-,P P1-
<b><math>J/\psi(3097) \rightarrow (\phi \pi^0)/total (units 10^{-3})</math></b> (P33)	R53B 204 2.04 0.27 PERUZZI 78 SMAG E+E-,P P1+
R33 (0.013) OR LESS BALTRUS2 85 SMK3 E+E-, K-K-PI0	R53S 32 1.7 0.7 BESCH 81 BONA E=E-
<b><math>J/\psi(3097) \rightarrow (\phi \pi^+ \pi^-)/total</math></b> (P34)	R53B 8 1.6 1.2 BESCH 81 BONA E=E-
R34 23 0.0021 0.0009 FELDMAN 77 SMAG E+E-	R53S 1288 2.02 0.24 EATON 84 SMK2 E+E-,P P1-
<b><math>J/\psi(3097) \rightarrow (2(\pi^+ \pi^-))/total</math></b> (P35)	R53S B1197 1.93 0.24 EATON 84 SMK2 E+E-,P P1+
R35 (0.0015)OR LESS CL=0.90 VANNUCCI 77 SMAG E=E-	R53S SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.
<b><math>J/\psi(3097) \rightarrow (\phi K \bar{K})/total</math></b> (P36)	R53B FROM ANTI-CHANNEL (PBAR N P1+)
R36 14 0.0018 0.0008 FELDMAN 77 SMAG E=E-	R53 R53
	R53 AVG 2.01 0.13 AVERAGE
	<b><math>J/\psi(3097) \rightarrow (p \bar{p} \pi^+ \pi^-)/total (units 10^{-3})</math></b> (P54)
	R54 533 5.5 0.6 PERUZZI 78 SMAG E+E-,P PB 1-2P1
	R54 48 3.8 1.6 BESCH 81 BONA E=E-
	R54S 1435 6.46 0.46 EATON 84 SMK2 E+E-,HADRONS GAM
	R54S SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.
	R54 R54
	R54 AVG 5.99 0.48 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)







# Meson Full Listings

$\chi_2(3555)$ ,  $\eta_c(3590)$

$\chi_2(3555)$   
was  $\chi(3555)$

$$I^G(J^{PC}) = 0^+(2^{-+})$$

OBSERVED IN RADIATIVE DECAY OF  $\Psi(3685)$  INTO  $\chi_{12}(3555)$  GAMMA. THEREFORE C++, THE OBSERVED DECAY INTO  $4P_1$  AND  $6P_1$  IMPLY  $G_{++}$ , THUS  $I=0$ .  
J=0 IS EXCLUDED BY THE ANGULAR DISTRIBUTION IN THE HADRONIC DECAYS. JP ABNORMAL EXCLUDED BY  $P_1, P_1-$  AND  $K^*, K^-$  DECAYS.  
JP=2- PREFERRED (FELDMAN 77, OREGLIA 82).

### $\chi_2(3555)$ MASS (MeV)

M	(3550.0)	(10.0)	TRILLING	76	SMAG	E-E-, HADRONS	GAM	
M	4(3543.0)	(10.0)	WHITAKER	76	SMAG	E-E-, J/ $\Psi$	2 GAM	
M	D 360 3563.0	7.0	BIDDICK	77	CNTR	E-E-, MONOCHR.	GAM	
M	D 3553.0	4.0	BARTEL	78	CNTR	E-E-, J/ $\Psi$	2 GAM	
M	D M 3552.0	6.0	TANENBAUM	78	SMAG	E-E-		
M	M 15 3551.0	11.0	BRANDEL2	79	DASP	E-E-, J/ $\Psi$	26AM	
M	D 69 3557.7	1.5	HIMEL	80	SMK2	E-E-, J/ $\Psi$	2 GAM	
M	P 66 3553.4	2.2	LEMOIGNE	82	GOLI	190 PI-BE, GAM2MU		
M	E 3555.9	0.7	OREGLIA	82	CBAL	E-E-, J/ $\Psi$	2 GAM	
M	F 3556.8	0.6	BAGLIN	85	SPEC	PBAR P, E- E-		
M	AVG 3556.31	0.42	AVERAGE					

M D MASS VALUE SHIFTED BY US BY AMOUNT APPROPRIATE FOR  
M D  $\Psi(3685)$  MASS=3686 AND  $\Psi(3097)$  MASS=3097.  
M E ASSUMING  $\Psi(3685)$  MASS=3686 AND  $\Psi(3097)$  MASS=3097.  
M F SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.  
M M FROM A SIMULTANEOUS FIT TO RADIATIVE AND HADRONIC DECAY CHANNELS  
M M SYSTEMATIC ERROR ADDED QUADRATICALLY BY US  
M P J/ $\Psi$  MASS CONSTRAINED TO 3097

### $\chi_2(3555)$ WIDTH (MeV)

W	2.9	1.8	1.1	BAGLIN	85	SPEC	PBAR P, E- E-
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### $\chi_2(3555)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$\chi_2(3555) \rightarrow \pi^+ \pi^-$	140- 140
P2	$\chi_2(3555) \rightarrow K^+ K^-$	494- 494
P3	$\chi_2(3555) \rightarrow 2(\pi^+ \pi^-)$	140- 140- 140- 140
P4	$\chi_2(3555) \rightarrow 3(\pi^+ \pi^-)$	
P5	$\chi_2(3555) \rightarrow \pi^+ \pi^- K^+ K^-$	140- 140- 494- 494
P6	$\chi_2(3555) \rightarrow J/\psi(3097) \gamma$	3097- 0
P7	$\chi_2(3555) \rightarrow \gamma \gamma$	0- 0
P8	$\chi_2(3555) \rightarrow \pi^+ \pi^- \rho \bar{\rho}$	140- 140- 938- 938
P9	$\chi_2(3555) \rightarrow \rho^0 \pi^+ \pi^-$	769- 140- 140
P10	$\chi_2(3555) \rightarrow K^*(892)^0 K^+ K^-$	892- 494- 140
P11	$\chi_2(3555) \rightarrow \rho \bar{\rho}$	938- 938
P12	$\chi_2(3555) \rightarrow J/\psi(3097) \pi^+ \pi^- \pi^0$	3097- 140- 140- 135

### $\chi_2(3555)$ BRANCHING RATIOS

$\chi_2(3555) \rightarrow (\gamma \gamma)/\text{total}$	(P7)
R1 T (0.0006) OR LESS CL=0.90	YAMADA 77 DASP E- E-, 3 GAMMA
$\chi_2(3555) \rightarrow 2(\pi^+ \pi^-)/\text{total}$	(P3)
R2 T 0.022 0.005	TANENBAUM 78 SMAG $\Psi(3685)$ TO GAM CHI
$\chi_2(3555) \rightarrow (\pi^+ \pi^- K^+ K^-)/\text{total}$	(P5)
R3 T 0.019 0.005	TANENBAUM 78 SMAG $\Psi(3685)$ TO GAM CHI
$\chi_2(3555) \rightarrow 3(\pi^+ \pi^-)/\text{total}$	(P4)
R4 T 0.011 0.008	TANENBAUM 78 SMAG $\Psi(3685)$ TO GAM CHI
$\chi_2(3555) \rightarrow (\pi^+ \pi^- + K^+ K^-)/\text{total}$	(P1-P2)
R5 T 0.0025 0.0010	TANENBAUM 78 SMAG $\Psi(3685)$ TO GAM CHI
$\chi_2(3555) \rightarrow (\pi^+ \pi^- \rho \bar{\rho})/\text{total}$	(P8)
R6 T 0.0033 0.0013	TANENBAUM 78 SMAG $\Psi(3685)$ TO GAM CHI
$\chi_2(3555) \rightarrow (J/\psi(3097) \gamma)/\text{total}$	(P6)
R7 T (0.28) (0.13)	BIDDICK 77 CNTR $\Psi(3685)$ TO GAM CHI
R7 T 0.13 0.03	BARTEL 78 CNTR $\Psi(3685)$ TO GAM CHI
R7 T 0.11 0.13	SPITZER 78 PLUT $\Psi(3685)$ TO GAM CHI
R7 T 0.13 0.08	TANENBAUM 78 SMAG $\Psi(3685)$ TO GAM CHI
R7 T 0.18 0.05	BRANDEL2 79 DASP $\Psi(3685)$ TO GAM CHI
R7 T 0.14 0.04	HIMEL 80 SMK2 $\Psi(3685)$ TO GAM CHI
R7 T 479 0.162 0.028	OREGLIA 82 CBAL $\Psi(3685)$ TO GAM CHI
R7 AVG 0.148 0.017	AVERAGE
$\chi_2(3555) \rightarrow (\rho^0 \pi^+ \pi^-)/\text{total}$	(P9)
R8 T 0.0067 0.0040	TANENBAUM 78 SMAG $\Psi(3685)$ TO GAM CHI
$\chi_2(3555) \rightarrow (K^*(892)^0 K^+ K^-)/\text{total}$	(P10)
R9 T 0.0047 0.0028	TANENBAUM 78 SMAG $\Psi(3685)$ TO GAM CHI

$\chi_2(3555) \rightarrow (\pi^+ \pi^-)/\text{total (units } 10^{-3})$	(P1)
R10 T 4 1.9 1.0	BRANDEL1 79 DASP $\Psi(3685)$ TO GAM CHI
$\chi_2(3555) \rightarrow (K^+ K^-)/\text{total (units } 10^{-3})$	(P2)
R11 T 2 1.5 1.1	BRANDEL1 79 DASP $\Psi(3685)$ TO GAM CHI
$\chi_2(3555) \rightarrow (p \bar{p})/\text{total (units } 10^{-3})$	(P11)
R12 T (0.9) OR LESS CL=0.90	BRANDEL2 79 DASP $\Psi(3685)$ TO GAM CHI
R12 7.6 4.0 2.0	BAGLIN 85 SPEC PBAR P, E- E-
$\chi_2(3555) \rightarrow (J/\psi(3097) \pi^+ \pi^- \pi^0)/\text{total}$	(P12)
R13 (0.015) OR LESS CL=0.90	BARATE 81 SPEC 190 PI-BE, 2P12MU
R T ESTIMATED USING $\Psi(3685)$ TO GAMMA $\chi_{12}(3555)$ TOTAL=0.078	
R T THE ERRORS DO NOT CONTAIN THE UNCERTAINTY IN THE $\Psi(3685)$ DECAY.	

### REFERENCES FOR $\chi_2(3555)$

FELDMAN 75 PRL 35 821	+JEAN-MARIE, SADDULET, VANNUCCI, + (LBL-SLAC)
ALSO 75 PRL 35 1189 (ERRATA)	
TANENBAUM 75 PRL 35 1323	TANENBAUM, WHITAKER, ABRAMS, + (LBL-SLAC)
TRILLING 76 STANFORD SYMP. 437	G. H. TRILLING (LBL)
WHITAKER 76 PRL 37 1596	-TANENBAUM, ABRAMS, ALAM, BOYARSKI, + (SLAC-LBL)
BIDDICK 77 PRL 38 1324	+BURNETT (UCSD+UMD+PAVI+PRIN+SLAC+STAN)
FELDMAN 77 PL 33 C 285	+PERL (LBL-SLAC)
YAMADA 77 HAMB. CONF. P. 69	YAMADA (DESY-TOKY)
BARTEL 78 PL 79 B 492	DITTMANN, DUINKER, OLSOON, O'NEILL, + (DESY+HEID)
SPITZER 78 KYOTO SUM. INST. 47	H. SPITZER (HAMB)
TANENBAUM 78 PR D 17 1751	TANENBAUM, ALAM, BOYARSKI, + (SLAC-LBL)
ALSO 82 PRIVATE COMM.	G.H. TRILLING (LBL+UCB)
BRANDEL1 79 ZPHY C 1 233	BRANDEL1, CORDS, -(AACH+DESY+HAMB+MPI+TOKY)
BRANDEL2 79 NP B 160 426	BRANDEL1, CORDS, -(AACH+DESY+HAMB+MPI+TOKY)
KIRK 79 PRL 42 619	+GOODMAN, ALVERSON, +(FNL+HARV+ILL+OXF+TUFT)
HIMEL 80 PRL 44 920	+ABRAMS, ALAM, BLOCKER, + (LBL-SLAC)
ALSO 82 PRIVATE COMM.	G.H. TRILLING (LBL+UCB)
BARATE 81 PR D 24 2994	+ASTBURY, MCEWEN, + (SACL+LOIC-SHMP-CERN+IND)
LEMOIGNE 82 PL 113 B 509	+BARATE, ASTBURY, MCEWEN, -(SACL+LOIC-SHMP+IND)
OREGLIA 82 PR D 25 2259	+PARTRIDGE, BLOOM, -(SLAC+CIT+HARV+PRIN+STAN)
ALSO 82 PRIVATE COMM.	M. OREGLIA (EFI)
BARATE 83 PL 121 B 449	+BAREVRE, ASTBURY, MCEWEN (SACL+LOIC-SHMP-IND)
BAGLIN 85 CERN-PRE-85	(LAPP-CERN+GENO-LYON-OSLO-ROMA+STRB-TORI)

$\eta_c(3590)$

$$I^G(J^{PC}) = ?^?(?^+)$$

OMITTED FROM SUMMARY TABLE OUR LATEST MINIREVIEW ON THIS PARTICLE CAN BE FOUND IN THE 1984 EDITION. NEEDS CONFIRMATION.

### $\eta_c(3590)$ MASS (MeV)

M	A	3594.0	5.0	EDWARDS	82	CBAL	E+E-, GAM INCL
M	A	ASSUMING MASS OF $\Psi(3685) = 3686$ MEV.					

### $\eta_c(3590)$ WIDTH (MeV)

W	(8.0)	OR LESS	CL=.95	EDWARDS	82	CBAL	E+E-, GAM INCL
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### $\eta_c(3590)$ PARTIAL DECAY MODES

P1	$\eta_c(3590) \rightarrow \text{hadrons}$	DECAY MASSES
----	---	--------------

### $\eta_c(3590)$ BRANCHING RATIOS

$\eta_c(3590) \rightarrow \text{hadrons}$	(P1)
R1 T SEEN	EDWARDS 82 CBAL E+E-, GAM INCL

### REFERENCES FOR $\eta_c(3590)$

BARTEL 78 PL 79 B 492	-DITTMANN, DUINKER, OLSOON, + (DESY+HEID)
PORTER 81 SLAC SUM. CONF. 355	-EDWARDS, - (CIT+HARV+PRIN+STAN+SLAC)
EDWARDS 82 PRL 48 70	+PARTRIDGE, PECK, + (CIT+HARV+PRIN+STAN+SLAC)
OREGLIA 82 PR D 25 2259	+PARTRIDGE, BLOOM, -(SLAC+CIT+HARV+PRIN+STAN)

For notation, see key on page 91.

# Meson Full Listings

## $\psi(3685)$

**$\psi(3685)$**

$$J^{G}(J^{PC}) = 0^{--}(1^{--})$$

### $\psi(3685)$ MASS (MeV)

WE USE INDEPENDENT MEASUREMENTS OF THE  $J/\Psi(3097)$  MASS, THE  $\Psi(3685)$  MASS, AND THE MASS DIFFERENCE TO PERFORM A CONSTRAINED FIT.

M S	3680.3	37.	CRIEGEE	75 PLUT	E+E-
M R	(3684.)	(5.)	LUTH	75 SMAG	E+E-
M	3684.	9.	PREPOST	75 SPEC	21. GAMMA D
M F	140(3683.0)	(6.0)	LEMOIGNE	79 SOL1	0 150 P1-BE,2MU
M	3686.	3.	BRANDEL1	79 DASP	E+ E-
M	413 3686.00	0.10	ZHOLENTZ	80 OLYA	E+E- COLL.BEAMS
M	AVG	3686.000			AVERAGE
M	FIT	3686.00			FROM FIT

M F FROM A SIMULTANEOUS FIT TO E+, E-, MU-, MU- AND HADRONIC CHANNELS  
M R ASSUMING  $G(E+E-) = G(\mu^+\mu^-)$   
M R REDUNDANT WITH DATA IN MASS DIFFERENCE BELOW  
M S ERROR OF ABOUT 1 PER CENT FROM THE UNCERTAINTY IN CALIBRATION OF THE BEAM ENERGY.

### $\psi(3685) - J/\Psi(3097)$ MASS DIFFERENCE (MeV)

DM	588.7	.8	LUTH	75 SMAG	E+E-
DM R	(589.07)	(0.13)	ZHOLENTZ	80 OLYA	E+E-
DM	589.7	1.2	LEMOIGNE	82 GOL1	190 P1-BE,2MU
DM	589.01	0.67			AVERAGE
DM	FIT	589.06			FROM FIT
DM R	REDUNDANT WITH DATA IN MASS ABOVE				

### $\psi(3685)$ WIDTH (keV)

W	228.	56.	LUTH	75 SMAG	E+E-
W F	202.	57.	BRANDEL1	79 DASP	E+ E-
W F	FROM A SIMULTANEOUS FIT TO E+, E-, MU+ MU- AND HADRONIC CHANNELS				
W F	ASSUMING $G(E+ E-) = G(\mu^+ \mu^-)$				
W	AVG	215.2			39.9 AVERAGE

### $\psi(3685)$ PARTIAL DECAY MODES

P	DECAY MODE	DECAY MASSES
P1	$\psi(3685) \rightarrow e^+ e^-$	.511+.511
P2	$\psi(3685) \rightarrow \mu^+ \mu^-$	106+ 106
P3	$\psi(3685) \rightarrow \text{hadrons}$	
P4	$\psi(3685) \rightarrow \text{virtual } \gamma \rightarrow \text{hadrons}$	
P	DECAYS INTO $J/\Psi(3097) + \text{ANYTHING}$	
P	-----	
P11	$\psi(3685) \rightarrow J/\Psi(3097) + \text{anything}$	
P12	$\psi(3685) \rightarrow J/\Psi(3097) + \text{neutrals}$	
P13	$\psi(3685) \rightarrow J/\Psi(3097) \pi^+ \pi^-$	3097+ 140. 140
P14	$\psi(3685) \rightarrow J/\Psi(3097) \pi^0 \pi^0$	3097+ 135. 135
P15	$\psi(3685) \rightarrow J/\Psi(3097) \eta$	3097+ 549
P16	$\psi(3685) \rightarrow J/\Psi(3097) \gamma \gamma$	3097+ 0+ 0
P17	$\psi(3685) \rightarrow J/\Psi(3097) \pi^0$	3097. 135
P17	SMALL -- NOT USED IN FIT	
P	HADRONIC DECAYS	
P	-----	
P21	$\psi(3685) \rightarrow \pi^+ \pi^-$	140+ 140
P22	$\psi(3685) \rightarrow \rho \pi$	769+ 140
P23	$\psi(3685) \rightarrow K^+ K^-$	494+ 494
P24	$\psi(3685) \rightarrow 2(\pi^+ \pi^-)$	140+ 140+ 140+ 140
P25	$\psi(3685) \rightarrow 2(\pi^+ \pi^-) \pi^0$	140+ 140+ 140+ 140
P26	$\psi(3685) \rightarrow \pi^+ \pi^- K^+ K^-$	140+ 140+ 494+ 494
P27	$\psi(3685) \rightarrow \bar{\rho} \rho$	938+ 938
P28	$\psi(3685) \rightarrow \Delta \Delta$	1116+ 1116
P29	$\psi(3685) \rightarrow \Xi \Xi$	1321+ 1321
P31	$\psi(3685) \rightarrow \pi^+ \pi^- \rho \bar{\rho}$	140+ 140+ 938+ 938
P32	$\psi(3685) \rightarrow 3(\pi^+ \pi^-)$	769+ 140+ 140
P33	$\psi(3685) \rightarrow \rho^0 \pi^+ \pi^-$	892+ 494+ 140
P34	$\psi(3685) \rightarrow K^*(892)^0 K^\pm \pi^\mp$	938+ 938- 135
P35	$\psi(3685) \rightarrow \bar{\rho} \rho \pi^0$	140+ 140+ 135
P36	$\psi(3685) \rightarrow \pi^+ \pi^- \pi^0$	494+ 494+ 135
P37	$\psi(3685) \rightarrow 3(\pi^+ \pi^-) \pi^0$	494+ 494+ 135
P38	$\psi(3685) \rightarrow K^+ K^- \pi^0$	494+ 494+ 135
P39	$\psi(3685) \rightarrow K^\pm K^*(892)^\mp$	494+ 892

### RADIATIVE DECAYS

P	RADIATIVE DECAYS		
P	-----		
P51	$\psi(3685) \rightarrow \gamma \gamma$	0+	0
P52	$\psi(3685) \rightarrow \pi^0 \gamma$	135+	0
P53	$\psi(3685) \rightarrow \eta \gamma$	549+	0
P54	$\psi(3685) \rightarrow \eta' \gamma$	958+	0
P56	$\psi(3685) \rightarrow \chi_0(3415) \gamma$	3415+	0
P58	$\psi(3685) \rightarrow \chi_1(3510) \gamma$	3510+	0
P59	$\psi(3685) \rightarrow \chi_2(3555) \gamma$	3556+	0
P60	$\psi(3685) \rightarrow \chi_1(3510) + \text{anything}$		
P61	$\psi(3685) \rightarrow \eta_c(2980) \gamma$	2981+	0
P62	$\psi(3685) \rightarrow \eta(1440) \gamma$	1440+	0
P63	$\psi(3685) \rightarrow \eta_c(3590) \gamma$	3594+	0

### FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions,  $P_i$ , as follows: The diagonal elements are the  $P_i = \delta P_i$ , where  $\delta P_i = \sqrt{(\delta P_i / P_i)^2 + (\delta P_i / P_i)^2}$ , while the off-diagonal elements are the normalized correlation coefficients  $(\delta P_i \delta P_j) / (\delta P_i \delta P_j)$ . For the definitions of the individual  $P_i$ , see the listings above; only those  $P_i$  appearing in the matrix are assumed in the matrix to be nonzero and are thus constrained to add to 1.

$J/\Psi \rightarrow e^+ e^-$	$J/\Psi \rightarrow \mu^+ \mu^-$	$J/\Psi \rightarrow \pi^+ \pi^-$	$J/\Psi \rightarrow \eta$	$J/\Psi \rightarrow \text{OTHER}$	NON- $J/\Psi$		
		-.4106	-.0821				
		.8343	.2176	-.0522			
		.0155	.0130	.0265	-.0036		
		.4864	.0643	-.0298	.0373	-.0464	
		-.9819	-.8192	-.0273	-.5935	.3081	-.1510

### $\psi(3685)$ PARTIAL WIDTHS

$\psi(3685) \rightarrow e^+ e^-$ (keV)	2.1	.3	LUTH	75 SMAG	E+E-	(61)
W1 F	2.1	0.3	BRANDEL1	79 DASP	E+ E-	
W1 F	FROM A SIMULTANEOUS FIT TO E+, E-, MU+ MU- AND HADRONIC CHANNELS					
W1 F	ASSUMING $G(E+ E-) = G(\mu^+ \mu^-)$					
W1	AVG	2.05			0.21 AVERAGE	

$\psi(3685) \rightarrow \text{hadrons}$ (keV)	224.	56.	LUTH	75 SMAG	E+E-	(63)	
W3							
$\psi(3685) \rightarrow \gamma \gamma$ (eV)	43.	OR LESS	CL=0.90	BRANDEL1	79 DASP	E+ E-	(651)
W51							

### $\psi(3685)$ BRANCHING RATIOS

$\psi(3685) \rightarrow (e^+ e^-)/\text{total}$	.0088	.0013	FERDMAN	77 RVUE	E+E-	(P1)
R1 L	FROM AN OVERALL FIT ASSUMING EQUAL PARTIAL WIDTHS FOR (E+E-)					
R1 L	AND (MU+MU-), FOR A MEASUREMENT OF THE RATIO SEE THE ENTRY R4 BELOW					
R1 L	INCLUDES LUTH 75, HILGER 75, BURMESTER 77					
$\psi(3685) \rightarrow (\mu^+ \mu^-)/\text{total}$	.0077	.0017	HILGER	75 SPEC	E+E-	(P2)
R2 H	RE-STATIED BY US USING (J/Psi(3097)+ANYTHING)/TOTAL = 0.55					
$\psi(3685) \rightarrow (\text{hadrons})/\text{total}$	.981	.003	LUTH	75 SMAG	E+E-	(P3)
R3 P	INCLUDES CASCADE DECAY INTO $J/\Psi(3097)$					
$\psi(3685) \rightarrow (\mu^+ \mu^-)/(e^+ e^-)$	.89	.16	BOYARSKI	75 SMAG	E+E-	(P2)/(P1)
$\psi(3685) \rightarrow (\gamma \rightarrow \text{hadrons})/\text{total}$	.029	.004	LUTH	75 SMAG	E+E-	(P4)
R5 C	INCLUDED IN R3					
R	DECAYS INTO $J/\Psi(3097) + \text{ANYTHING}$					
R	-----					
$\psi(3685) \rightarrow (J/\Psi(3097) + \text{anything})/\text{total}$	.57	.08	ABRAMS	75 SMAG	E+E-	(P11)
R10	.51	0.12	BRANDEL1	79 DASP	E+ E-	
R10						
R10 AVG	.552	0.067			AVERAGE	
R10 FIT	.69	0.15			FROM FIT (ERROR INCLUDES SCALE FACTOR OF 3.1)	
$\psi(3685) \rightarrow (J/\Psi(3097) + \text{neutrals})/(J/\Psi(3097) + \text{anything})$	.41	.02	TAMENBAUM	76 SMAG	E+E-	(P12)/(P11)
R11						
R11 FIT	.395	0.027			FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.4)	
$\psi(3685) \rightarrow (J/\Psi(3097) \pi \pi)/\text{total}$	0.48	0.06	ABRAMS1	75 SMAG	E+E-, JPS1 P1+P1-	(P13+P14)
R12	.51	0.087	ABRAMS1	75 SMAG	E+E-, JPS1 2P10	
R12	.54	0.09	WIK	75 DASP	E+E-, JPS1 P1+P1-	
R12	.54	0.18	WIK	75 DASP	E+E-, JPS1 2P10	
R12						
R12 AVG	.503	0.042			AVERAGE	
R12 FIT	.411	0.082			FROM FIT (ERROR INCLUDES SCALE FACTOR OF 2.7)	

# Meson Full Listings

## $\psi(3685)$

$\psi(3685) \rightarrow (J/\psi(3097) \pi^0 \pi^0)/(J/\psi(3097) \pi^+ \pi^-)$  (P14)/(P13)

R14 H	(.64)	(.15)	HILGER 75 SPEC	E+E-
R14	0.53	0.06	TANENBAUM 76 SMAG	E+E-

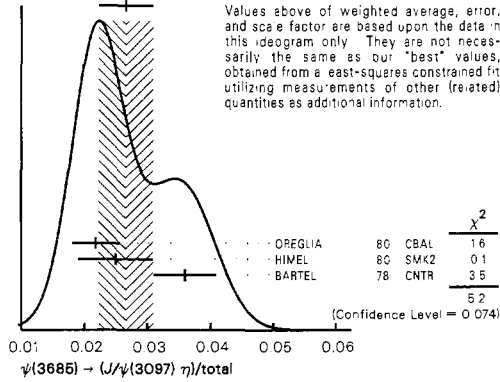
R14 H IGNORING THE (J/PSI ETA) AND (J/PSI GAMMA GAMMA) DECAYS

R14 FIT	0.530	0.070	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
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$\psi(3685) \rightarrow (J/\psi(3097) \eta)/\text{total}$  (P15)

R15 S	44	(.043)	(.008)	TANENBAUM 76 SMAG	E-E-
R15	164	0.036	0.005	BARTEL 78 CNTR	E+E-
R15 S	17	(0.035)	(0.009)	BRANDEL 79 DASP	E+E-, PSI 2GAM
R15	166	0.025	0.006	HIMEL 80 SMK2	E+E-
R15 D	386	0.0218	0.0038	OREGLIA 80 CBAL	E+E-, PSI 2GAM
R15 AVG		0.0266	0.0044	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6)	
R15 FIT		0.0265	0.0036	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.3)	
				(SEE IDEOGRAM BELOW)	

WEIGHTED AVERAGE  
0.0266 ± 0.0044 (ERROR SCALED BY 1.6)



Values above of weighted average, error, and scale factor are based upon the data in this ideogram only. They are not necessarily the same as our "best" values, obtained from a least-squares constrained fit utilizing measurements of other (related) quantities as additional information.

$\psi(3685) \rightarrow (J/\psi(3097) \pi^0)/\text{total}$  (P17)

R16	7	0.0015	0.0006	HIMEL 80 SMK2	E+E-
R16 D	23	0.0009	0.0003	OREGLIA 80 CBAL	E+E-, PSI 2GAM
R16 AVG		0.00102	0.00027	AVERAGE	

R HADRONIC DECAYS

$\psi(3685) \rightarrow (\pi^+ \pi^-)/\text{total (units } 10^{-4})$  (P21)

R20	(0.5)	OR LESS	CL=0.90	FELDMAN 77 SMAG	E+E-
R20	0.8			BRANDEL 79 DASP	E+E-

$\psi(3685) \rightarrow (2\pi^+ \pi^- \pi^0)/\text{total}$  (P25)

R22	42	0.0035	0.0015	ABRAMS 75 SMAG	E+E-, HADRONS
R22		0.005	0.0008	FRANKLIN 83 SMK2	E+E-, HADRONS
R22 AVG		0.00311	0.00071	AVERAGE	

$\psi(3685) \rightarrow (K^+ K^-)/\text{total (units } 10^{-4})$  (P23)

R23	(0.5)	OR LESS	CL=0.90	FELDMAN 77 SMAG	E+E-
R23	1.0			BRANDEL 79 DASP	E+E-

$\psi(3685) \rightarrow (\pi^+ \pi^- K^+ K^-)/\text{total}$  (P26)

R24K	0.0016	0.0004	TANENBAUM 78 SMAG	E+E-
------	--------	--------	-------------------	------

R24K ASSUMING ENTIRELY STRONG DECAY

$\psi(3685) \rightarrow (\bar{p} p)/\text{total (units } 10^{-4})$  (P27)

R25	4	2.5	0.7	FELDMAN 77 SMAG	E+E-
R25		1.4	0.8	BRANDEL 79 DASP	E+E-
R25 AVG		1.91	0.53	AVERAGE	

$\psi(3685) \rightarrow (\rho \pi)/\text{total (units } 10^{-4})$  (P22)

R26 N	(10.)	OR LESS	CL=0.90	ABRAMS 75 SMAG	E+E-	
R26	(10.)	OR LESS	CL=0.90	BARTEL 76 CNTR	E+E-	
R26	1	(0.83)	OR LESS	CL=0.90	FRANKLIN 83 SMK2	E+E-, HADRONS
R26 N	FINAL STATE RHO0 P10					

$\psi(3685) \rightarrow 2(\pi^+ \pi^-)/\text{total}$  (P24)

R27	0.00045	0.0001	TANENBAUM 78 SMAG	E+E-
-----	---------	--------	-------------------	------

$\psi(3685) \rightarrow (\Delta \bar{\Delta})/\text{total}$  (P28)

R28	(0.0004)	OR LESS	CL=0.90	FELDMAN 77 SMAG	E+E-
-----	----------	---------	---------	-----------------	------

$\psi(3685) \rightarrow (\Xi^- \bar{\Xi}^+)/\text{total}$  (P29)

R29	(0.0002)		FELDMAN 77 SMAG	E+E-
-----	----------	--	-----------------	------

$\psi(3685) \rightarrow (\pi^+ \pi^- \bar{p} p)/\text{total (units } 10^{-3})$  (P31)

R31S	0.8	0.2	TANENBAUM 78 SMAG	E+E-
------	-----	-----	-------------------	------

R31S ASSUMING ENTIRELY STRONG DECAY

$\psi(3685) \rightarrow 3(\pi^+ \pi^-)/\text{total (units } 10^{-3})$  (P32)

R32S	0.15	0.1	TANENBAUM 78 SMAG	E+E-
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$\psi(3685) \rightarrow (\rho^0 \pi^+ \pi^-)/\text{total (units } 10^{-3})$  (P33)

R33	0.42	0.15	TANENBAUM 78 SMAG	E+E-
-----	------	------	-------------------	------

$\psi(3685) \rightarrow (K^*(892)^0 K^\pm \pi^\mp)/\text{total (units } 10^{-3})$  (P34)

R34	0.67	0.25	TANENBAUM 78 SMAG	E+E-
-----	------	------	-------------------	------

$\psi(3685) \rightarrow (p \bar{p} \pi^0)/\text{total (units } 10^{-4})$  (P35)

R35	9	1.4	0.5	FRANKLIN 83 SMK2	E+E-, HADRONS
-----	---	-----	-----	------------------	---------------

$\psi(3685) \rightarrow (\pi^+ \pi^- \pi^0)/\text{total (units } 10^{-5})$  (P36)

R36	4	8.5	4.6	FRANKLIN 83 SMK2	E+E-, HADRONS
-----	---	-----	-----	------------------	---------------

$\psi(3685) \rightarrow (3\pi^+ \pi^- \pi^0)/\text{total}$  (P37)

R37	6	.0035	.0016	FRANKLIN 83 SMK2	E+E-, HADRONS
-----	---	-------	-------	------------------	---------------

$\psi(3685) \rightarrow (K^+ K^- \pi^0)/\text{total (units } 10^{-5})$  (P38)

R38	1	(2.96)	OR LESS	CL=0.90	FRANKLIN 83 SMK2	E+E-, HADRONS
-----	---	--------	---------	---------	------------------	---------------

$\psi(3685) \rightarrow (K^\pm K^*(892)^\mp)/\text{total (units } 10^{-5})$  (P39)

R39	0	(1.79)	OR LESS	CL=0.90	FRANKLIN 83 SMK2	E+E-, HADRONS
-----	---	--------	---------	---------	------------------	---------------

R RADIATIVE DECAYS

$\psi(3685) \rightarrow (\pi^0 \gamma)/\text{total}$  (P52)

R42 U	(.0054)	OR LESS	CL=.95	LIBERMAN 75 SPEC	E+E-
R42	(.01)	OR LESS	CL=.90	WIK 75 DASP	E+E-

$\psi(3685) \rightarrow (\eta \gamma)/\text{total (units } 10^{-2})$  (P53)

R43	(0.02)	OR LESS	CL=0.90	YAMADA 77 DASP	E+E-, 3 GAMMA
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$\psi(3685) \rightarrow (\eta' \gamma)/\text{total (units } 10^{-2})$  (P54)

R44C	(0.11)	OR LESS	CL=0.90	BARTEL 76 CNTR	E+E-
R44 R	(0.6)	OR LESS	CL=0.90	BRAUNSCHW 77 DASP	E+E-

$\psi(3685) \rightarrow (X_{0(3415)} \gamma)/\text{total (units } 10^{-2})$  (P56)

R55 A	7.5	2.6	WHITAKER 76 SMAG	E+E-
R55 A	7.2	2.3	BIDDICK 77 CNTR	E+E-, MONOCHR. GAM
R55 D	9.9	0.9	GAISER 85 CBAL	E+E-, MONOCHR. GAM
R55 AVG	9.35	0.80	AVERAGE	

$\psi(3685) \rightarrow (X_{1(3510)} \gamma)/\text{total (units } 10^{-2})$  (P58)

R58 B	7.1	1.9	BIDDICK 77 CNTR	E+E-, MONOCHR. GAM
R58 D	9.0	0.9	GAISER 85 CBAL	E+E-, MONOCHR. GAM
R58 AVG	8.65	0.81	AVERAGE	

$\psi(3685) \rightarrow (X_{2(3555)} \gamma)/\text{total (units } 10^{-2})$  (P59)

R59 B	7.0	2.0	BIDDICK 77 CNTR	E+E-, MONOCHR. GAM
R59 F	8.0	0.9	GAISER 85 CBAL	E+E-, MONOCHR. GAM
R59 AVG	7.83	0.82	AVERAGE	

$\psi(3685) \rightarrow (\eta_c(2980) \gamma)/\text{total (units } 10^{-2})$  (P61)

R60	0.28	0.06	GAISER 85 CBAL	E+E-, MONOCHR. GAM
-----	------	------	----------------	--------------------

$\psi(3685) \rightarrow (\eta(1440) \gamma)/\text{total (units } 10^{-3})$  (P62)

R61 E	(0.12)	OR LESS	CL=.90	SCHARRE 80 SMAG	E+E-
-------	--------	---------	--------	-----------------	------

$\psi(3685) \rightarrow (\eta_c(3590) \gamma)/\text{total (units } 10^{-2})$  (P63)

R62	(0.2)	TO 1.3	CL=.95	EDWARDS 82 CBAL	E+E-, MONOCHR. GAM
-----	-------	--------	--------	-----------------	--------------------

R A ANGULAR DISTRIBUTION (1-COS\*2) ASSUMED  
R B VALID FOR ISOTROPIC DISTRIBUTION OF THE PHOTON  
R C THE VALUE IS NORMALIZED TO THE BRANCHING RATIO FOR PSI(3685)  
R D INTO (J/PSI(3097) ETA)/TOTAL  
R E SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.  
R F INCLUDES UNKNOWN BRANCHING FRACTION ETA(1440) INTO K KBAR PI.  
R G ANGULAR DISTRIBUTION (1-0.052\*cos\*2) ASSUMED  
R H ANGULAR DISTRIBUTION (1-0.189\*cos\*2) ASSUMED  
R I RE-STAT'D BY US USING (MU-MU-J)/TOTAL = .0077  
R J RE-STAT'D BY US USING TOTAL DECAY WIDTH 228 KEV.  
R S LOW STATISTICS DATA REMOVED FROM AVERAGE.

$$\psi(3685) \quad \Gamma(I^+ \Gamma^-(e^+ e^-))/\Gamma(\text{total}) \quad (\text{keV})$$

THIS COMBINATION OF A PARTIAL WIDTH WITH THE PARTIAL WIDTH INTO E+E- AND WITH THE TOTAL WIDTH IS OBTAINED FROM THE INTEGRATED CROSS-SECTION INTO CHANNEL (1) IN THE E+E- ANNIHILATION. WE ONLY LIST DATA NOT HAVING BEEN USED TO DETERMINE THE PARTIAL WIDTH G(1) OR THE BRANCHING RATIO G(1)/TOTAL.

$\Gamma(\text{hadronic}) \cdot \Gamma(e^+ e^-)/\Gamma(\text{total})$

G3	2.2	.4	ABRAMS 75 SMAG	E+E-
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### REFERENCES FOR $\psi(3685)$

ABRAMS 74 PRL 33 1453	-BRIGGS, AUGUSTIN, BOYARSKI+	(LBL+SLAC)
ABRAMS 75 STANFORD SYMP. 25	G. S. ABRAMS	(LBL)
ABRAMS 75 PRL 34 1181	+BRIGGS, CHINOVSKY, FRIEDBERG, +	(LBL+SLAC)
AUBERT 75 PRL 35 1624	+BECKER, BIGGS, BURGER, GLENN+	(MIT+BNL)
BOYARSKI 75 PALERMO CONF. 54	+BRIDENBACH, BULOS, ABRAMS, BRIGGS, (SLAC+LBL)	
CAMERINI 75 PRL 35 483	+LEARNED, PREPOST, ASH, ANDERSON, +	(WISC+SLAC)
CRIGGEE 75 PL 55B 489	+BEINE, FRANK, HORLITZ, KRECHLOK- (DESY)	
DASPE 75 PL 57B 407	BRUNSCHWIG, KONIGS+, (AACH-DESY+MPI+M. TOKYO)	
FELDMAN 75 PRL 35 821	+JEAN-MARIE, SAOULET, VANNUCCI, +	(LBL-SLAC)
GRECO 75 PL 56B 367	+PANCHERI-SRIVASTAVA, SRIVASTAVA (FRAS)	

For notation, see key on page 91.

# Meson Full Listings

$\psi(3685)$ ,  $\psi(3770)$ ,  $\psi(4030)$

JACKSON 75 NIM 128 13	J. D. JACKSON, D. SCHARRE (LBL)
HILGER 75 PRL 35 625	+BERON, FORD, HOFSTADTER, HOWELL, + (STAN-PENN)
LIBERMAN 75 STANFORD SYMP. 55	A. D. LIBERMAN (STANFORD)
LUTH 75 PRL 35 1124	+BOYARSKI, LYNCH, BREIDENBACH, + (SLAC+LBL) JPC
PREPOST 75 STANFORD SYMP. 241	R. PREPOST (WISCONSIN)
SIMPSON 75 PRL 35 699	+BERON, FORD, HILGER, HOFSTADTER, + (STAN-PENN)
WIJK 75 STANFORD SYMP. 69	B. N. WIJK (DESY)
BARTEL 76 PL 64 B 483	+DUINKER, OLSSON, STEFFEN, HEINTZE - (DESY-HEID)
SNYDER 76 PRL 36 1415	+HOM, LEDERMAN, APPEL, KAPLAN - (COLU-FNAL-STON)
TANENBAU 76 PRL 36 402	TANENBAUM, ABRAMS, BOYARSKI, BULOS, - (SLAC-LBL) JG
WHITAKER 76 PRL 37 1596	+TANENBAUM, ABRAMS, ALAM, BOYARSKI, - (SLAC+LBL)
BIDDICK 77 PRL 38 1324	+BURNETT, - (UCSD-UMD-PAVI-PRIN-SLAC-STAN)
BRAUNSCH 77 PL 67 B 249	BRAUNSCHEWIG, - (AACH-DESY-HAMB-MPI-M+TOKY)
BURMESTER 77 PL 66 B 395	BURMESTER, CRIGEE, - (DESY-HAMB-SIEG+WUPP)
FELDMAN 77 PL 33 C 285	+PERL (LBL+SLAC)
YAMADA 77 HAMB. CONF. P. 69	YAMADA (DESY+TOKY)
BARTEL 78 PL 79 B 492	DITTMANN, DUINKER, OLSSON, O'NEILL, - (DESY+HEID)
TANENBAU 78 PR D 17 1731	TANENBAUM, ALAM, BOYARSKI, + (SLAC-LBL)
BRANDEL 79 ZPHY C 1 233	BRANDELIK, CORDS, + (AACH-DESY-HAMB-MPI-M+TOKY)
BRANDEL 79 NP B 160 426	BRANDELIK, CORDS, + (AACH-DESY-HAMB-MPI-M+TOKY)
LEMOIGNE 79 FERMILAB CONF. 524	+ABOLINS, BARATE, + (SACL+LOIC-SHMP-IND)
HIMEL 80 PRL 44 920	+ABRAMS, ALAM, BLOCKER, + (LBL+SLAC)
ORegLIA 80 PRL 45 959	+PARTRIDGE - (SLAC+CIT+HARV+PRIN+STAN)
PARTRIDGE 80 PRL 45 1150	PARTRIDGE, PECK, - (CIT+HARV+PRIN+STAN-SLAC)
SCHARRE 80 PL 97 B 329	+TRILLING, ABRAMS, ALAM, BLOCKER + (SLAC-LBL)
ZHOLENTZ 80 PL 96 B 214	+KURDADZE, LELI CHUK, MISHNEV, NIKITIN - (NOVO)
ALSO 81 YAD. PHYS. 34 1471	ZHOLENTZ ET AL. (NOVO)
BARATE 81 PR D 24 2994	+ASTBURY, MCEWEN, + (SACL+LOIC-SHMP-CERN+IND)
EDWARDS 82 PRL 48 70	+PARTRIDGE, PECK, - (CIT+HARV+PRIN+STAN+SLAC)
LEMOIGNE 82 PL 113 B 509	+BARATE, ASTBURY, MCEWEN + (SACL+LOIC-SHMP-IND)
BARATE 83 PL 121 B 449	+BAREYRE, ASTBURY, MCEWEN (SACL+LOIC-SHMP-IND)
FRANKLIN 83 PRL 51 963	+FRANKLIN, FELDMAN, ABRAMS, ALAM - (LBL+SLAC)
FRANKLI 83 SLAC-254 THESIS	M. E. B. FRANKLIN (STAN)
GAISER 85 SLAC-PUB-2899	+BLOOD, BULOS - (CIT+HARV+PRIN+STAN-SLAC)

## REFERENCES FOR $\psi(3770)$

PERUZZI 77 PRL 39 1301	+PICCOLO, FELDMAN, PERL, - (SLAC, LBL, NWES-HAWA)
RAPIDIS 77 PRL 39 526	+GOBBI, LUKE, PERL, - (STAN+SLAC+LBL-NWES-HAWA)
BACINO 78 PRL 40 671	+BAUMGARTEN, BIRKWOOD, - (SLAC-STAN-UCLA-UCI)
SCHINDLER 80 PR D 21 2716	SCHINDLER, SIEGRIST, ALAM, BOYARSKI - (SLAC+LBL)

## $\psi(4030)$

$$I^G(J^{PC}) = ?(1^{--})$$

### NOTE ON $\psi(4030)$ AND $\psi(4160)$

Although these peaks are clearly separated from each other in the DASP experiment (BRANDELIK 78, 79) and are confirmed with much fewer statistics by PLUTO (BURMESTER 77), the existence of two resonances is in doubt. The SPEAR MARK-I data (SIEGRIST 82) show only one broad structure between 3.9 and 4.1 GeV, upon which there may be a sharp rise at 4.0 GeV. Part of the discrepancy between the experiments may come from the radiative corrections which tend to enhance any structure seen. A deviation from a naive Breit-Wigner-like structure is expected at 4.02 GeV, where the  $D^*\bar{D}^*$  threshold opens up.

It is difficult to accommodate within a conventional charmonium model two resonances with large  $e^+e^-$  coupling in this region. A  $3S$  state is expected near 4.10 GeV (RICHARDSON 79), while the slightly heavier  $2D$  state should couple only very weakly to  $e^+e^-$ , even when relativistic or coupled-channel effects are included (HEIKKILA 84). Therefore an identification of  $\psi(4160)$  as the  $2D$  state is hardly possible. Thus, if there really are two peaks, one must have a new type of state, e.g., a "hybrid"  $c\bar{c}g$  state (ONO 84). New data in this region would be most welcome.

## $\psi(3770)$

$$I^G(J^{PC}) = ?(1^{--})$$

### $\psi(3770)$ MASS (MeV)

M	E	3772.0	(6.0)	RAPIDIS	77	SMAG	0	E+ E-
M	E	3770.	(6.0)	BACINO	78	DLCO	0	E+ E-
M	E	3764.0	(5.0)	SCHINDLER	80	SMAG	E+ E-	
M E ERRORS INCLUDE SYSTEMATIC COMMON TO ALL EXPERIMENTS								
M	M	3769.9	2.5	FROM PSI(3685) MASS AND MASS DIFFERENCE BELOW				

### $\psi(3770) - \psi(3685)$ MASS DIFFERENCE (MeV)

DM	S	88.0	3.0	RAPIDIS	77	SMAG	E+ E-
DM	S	86.0	2.0	BACINO	78	DLCO	E+ E-
DM	S	80.0	2.0	SCHINDLER	80	SMAG	E+ E-
DM	S	85.9	2.4	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8)			

### $\psi(3770)$ WIDTH (MeV)

W		28.0	5.0	RAPIDIS	77	SMAG	0	E+ E-
W		24.0	5.0	BACINO	78	DLCO	0	E+ E-
W		24.0	5.0	SCHINDLER	80	SMAG	E+ E-	
W	AVG	25.3	2.9	AVERAGE				

### $\psi(3770)$ PARTIAL DECAY MODES

				DECAY MASSES	
P1	$\psi(3770) \rightarrow e^+e^-$			.511+.511	
P2	$\psi(3770) \rightarrow D\bar{D}$			1869+1869	

### $\psi(3770)$ PARTIAL WIDTHS (keV)

$\psi(3770) \rightarrow e^+e^-$ (G1)								
W1	R	0.37	0.09	RAPIDIS	77	SMAG	0	E+ E-
W1	R	0.18	0.06	BACINO	78	DLCO	0	E+ E-
W1	R	0.276	0.050	SCHINDLER	80	SMAG	E+ E-	
W1 R SEE ALSO R2 BELOW								
W1	AVG	0.257	0.046	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)				

### $\psi(3770)$ BRANCHING RATIOS

$\psi(3770) \rightarrow (D\bar{D})/\text{total}$ (P2)							
R1	DOMINANT		PERUZZI	77	SMAG	E+ E-, D DBAR	
$\psi(3770) \rightarrow (e^+e^-)/\text{total (units } 10^{-5})$ (P1)							
R2	1.3	0.2	RAPIDIS	77	SMAG	0	E+ E-

### $\psi(4030)$ MASS (MeV)

M		4028.0	2.5	GOLDHABER	77	SMAG	E+ E-
M		4040.0	10.0	BRANDELIK	78	DASP	E+ E-
M	AVG	4028.7	2.8	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2)			

### $\psi(4030)$ WIDTH (MeV)

W	52.0	10.0	BRANDELIK	78	DASP	E+ E-
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### $\psi(4030)$ PARTIAL DECAY MODES

				DECAY MASSES	
P1	$\psi(4030) \rightarrow D\bar{D}$			1869+1869	
P2	$\psi(4030) \rightarrow D^*(2010) \bar{D} + C.C.$			2007+1865	
P3	$\psi(4030) \rightarrow D^*(2010) \bar{D}^*(2010)$			2007+2007	
P4	$\psi(4030) \rightarrow J/\psi(3097) + \text{hadrons}$				
P5	$\psi(4030) \rightarrow e^+e^-$			.511+.511	
P6	$\psi(4030) \rightarrow \mu^+\mu^-$			106, 106	

### $\psi(4030)$ PARTIAL WIDTHS (keV)

$\psi(4030) \rightarrow e^+e^-$ (G5)								
W5	0.75	0.15	BRANDELIK	78	DASP	E+ E-		



# Meson Full Listings

$\psi(4030), \psi(4160), \psi(4415), \text{BOTTOMONIUM}$

## $\psi(4030)$ BRANCHING RATIOS

$\psi(4030) \rightarrow (D^0 \bar{D}^0)/(D^*(2010)^0 \bar{D}^0 + C.C.)$ (P1)/(P2)			
R1	P	0.05 0.03	GOLDBER 77 SMAG 0 E+ E-
$\psi(4030) \rightarrow (D^*(2010)^0 \bar{D}^*(2010)^0)/(D^*(2010)^0 \bar{D}^0 + C.C.)$ (P3)/(P2)			
R2	P	32.0 12.0	GOLDBER 77 SMAG 0 E+ E-
R P PHASE-SPACE FACTOR (P**3) EXPLICITLY REMOVED.			
$\psi(4030) \rightarrow (J/\psi(3097) + \text{hadrons})/\text{total}$ (P4)			
R3		LOOKED FOR	BURMESTER 77 PLUT E+E-
$\psi(4030) \rightarrow (e^+ e^-)/\text{total (units } 10^{-5})$ (P5)			
R4		(1.0) APPROX.	FELDMAN 77 SMAG E+ E-

## REFERENCES FOR $\psi(4030)$

AUGUSTIN 75 PRL 34 764	-BOYARSKI, ABRAMS, BRIGGS+	(SLAC-LBL)
BACCI 75 PL 588 481	-BIDDLI, PENSO, STELLA, .	(ROMA-FRAS)
BOYARSKI 75 PRL 34 762	-BREIDENBACH, ABRAMS, BRIGGS, +	(SLAC-LBL)
ESPOSITO 75 PL 588 478	-FELICETTI, PERUZZI, -	(FRAS-NAPL-PADO-ROMA)
PERUZZI 76 PRL 37 569	-PICCOLO, FELDMAN, NGUYEN, WISS, -	(SLAC-LBL)
BURMESTE 77 PL 66 B 395	+CRIEGEE, DEHNE+	(DESY-HAMB-SIEG-WUPP)
GOLDBER 77 PL 69 B 503	GOLDBER, WISS, ABRAMS, ALAM, LUTH, +	(LBL-SLAC)
FELDMAN 77 PL 33 C 285	-PERL	(LBL-SLAC)
LUTH 77 PL 70 B 120	-PIERRE, ABRAMS, ALAM, BOYARSKI, +	(LBL-SLAC)
BRANDELI 78 PL 76 B 361	BRANDELIK, CORDS, -	(AACH+DESY-HAMB-MPIM-TOKY)
ALSO 79 ZPHY C 1 233	BRANDELIK, CORDS, -	(AACH+DESY-HAMB-MPIM-TOKY)
KIRKBY 79 FERMILAB SYMP. 107	J. KIRKBY RAPPORTEUR	(SLAC)
RICHARDS 79 PL 82 B 272	J.L. RICHARDSON	(SLAC)
SIEGRIST 82 PR D 26 969	-SCHWITTERS, ALAM, CHINOWSKY, +	(SLAC-LBL)
HEIKKILA 84 PR D 29 110	+TORNVIST, ONO	(HELS-TOKY)
ONO 84 ZPHY C 26 307	SEIJI ONO	(ORSA)

## $\psi(4160)$

$$I^G(J^{PC}) = ?(1^{--})$$

SEE NOTE UNDER  $\psi(4030)$

## $\psi(4160)$ MASS (MeV)

M	4159.0	20.0	BRANDELIK 78 DASP E-E-
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## $\psi(4160)$ WIDTH (MeV)

W	78.0	20.0	BRANDELIK 78 DASP E-E-
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## $\psi(4160)$ PARTIAL DECAY MODES

P1	$\psi(4160) \rightarrow e^+ e^-$	DECAY MASSES .511+.511
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## $\psi(4160)$ PARTIAL WIDTHS (keV)

$\psi(4160) \rightarrow e^+ e^-$ (G1)	W1	0.77	0.23	BRANDELIK 78 DASP E+ E-
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## REFERENCES FOR $\psi(4160)$

BURMESTE 77 PL 66 B 395	+CRIEGEE, DEHNE+	(DESY-HAMB-SIEG-WUPP)
BRANDELI 78 PL 76 B 361	BRANDELIK, CORDS+	(AACH-DESY-HAMB-MPIM-TOKY)
KIRKBY 79 FERMILAB SYMP. 107	J. KIRKBY RAPPORTEUR	(SLAC)
ONO 84 ZPHY C 26 307	SEIJI ONO	(ORSA)

## $\psi(4415)$

$$I^G(J^{PC}) = ?(1^{--})$$

## $\psi(4415)$ MASS (MeV)

M	4414.	7.	SIEGRIST 76 SMAG E+E-	
M	(4400.)	APPROX.	KNIES 77 PLUT 0 E+E-, MU+ MU-	
M	4417.0	10.0	BRANDELIK 78 DASP E+E-	
M	AVG	4415.0	5.7	AVERAGE

## $\psi(4415)$ WIDTH (MeV)

W	33.	10.	SIEGRIST 76 SMAG E-E-	
W	66.0	15.0	BRANDELIK 78 DASP E+E-	
W	AVG	43.2	15.2	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.8)

## $\psi(4415)$ PARTIAL DECAY MODES

P1	$\psi(4415) \rightarrow e^+ e^-$	DECAY MASSES .511+.511
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## $\psi(4415)$ PARTIAL WIDTHS (keV)

$\psi(4415) \rightarrow e^+ e^-$ (G1)	W1	0.44	0.14	SIEGRIST 76 SMAG E- E-	
	W1	0.49	0.13	BRANDELIK 78 DASP E- E-	
	W1	AVG	0.467	0.095	AVERAGE

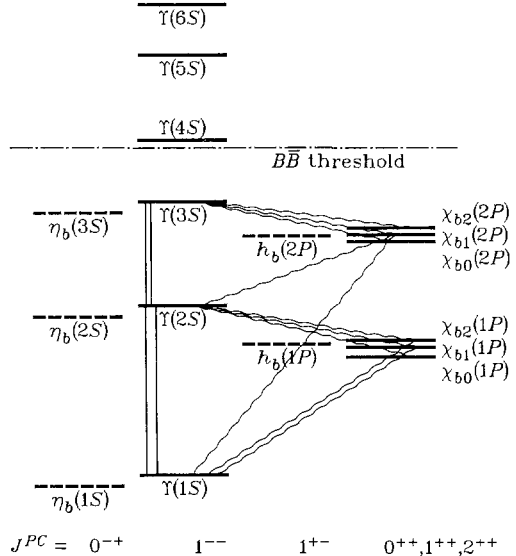
## $\psi(4415)$ BRANCHING RATIOS

$\psi(4415) \rightarrow \text{hadrons}/\text{total}$ (G1)	R2	DOMINANT	SIEGRIST 76 SMAG E+E-
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## REFERENCES FOR $\psi(4415)$

SIEGRIST 76 PRL 36 700	+ABRAMS, BOYARSKI, BREIDENBACH, -	(LBL-SLAC)
BURMESTE 77 PL 66 B 395	+CRIEGEE, DEHNE-	(DESY-HAMB-SIEG-WUPP)
KNIES 77 HAMBURG SYMP. 93	G. KNIES HAMBURG TALK ON PLUTO COLLAB. (DESY)	
LUTH 77 PL 70 B 120	+PIERRE, ABRAMS, ALAM, BOYARSKI, +	(LBL-SLAC)
BRANDELI 78 PL 76 B 361	BRANDELIK, CORDS-	(AACH-DESY-HAMB-MPIM-TOKY)

# THE BOTTOMONIUM SYSTEM



The level scheme of the  $b\bar{b}$  states with the names as adopted in this issue of the Review of Particle Properties. Singlet states are called  $\eta_b$  and  $h_b$ , triplet states  $T$  and  $\chi_{bJ}$ . In parentheses it is sufficient to give the radial quantum number and the orbital angular momentum to specify the states with all their quantum numbers. E.g.,  $h_b(2P)$  means  $2^1P_1$  with  $n = 2, L = 1, S = 0, J = 1, PC = +-.$  If found,  $D$ -wave states would be called  $\eta_b(nD), h_b(nD), T_J(nD),$  and  $\chi_{bJ}(nD),$  with  $J = 1, 2, 3$  and  $n = 1, 2, 3, 4, \dots$ . The figure also shows the observed hadronic transitions as solid lines and the observed radiative transitions as photonic lines.

For notation, see key on page 91.

## Meson Full Listings

### BOTTOMONIUM, $\Upsilon(9460)$

#### NOTE ON WIDTH DETERMINATIONS OF THE $\Upsilon$ STATES

As is the case for  $J/\psi(3097)$  and  $\psi(3685)$ , the full widths of the bound  $b\bar{b}$  states  $\Upsilon(1S)$ ,  $\Upsilon(2S)$ , and  $\Upsilon(3S)$  are not directly measurable, since they are much smaller than the energy resolution of the  $e^+e^-$  storage rings where these states are produced. The common indirect method to determine  $\Gamma$  starts from

$$\Gamma = \Gamma_{\ell\ell}/B_{\ell\ell}, \quad (1)$$

where  $\Gamma_{\ell\ell}$  is one leptonic partial width and  $B_{\ell\ell}$  is the corresponding branching fraction ( $\ell = e, \mu, \text{ or } \tau$ ). One then assumes  $e\text{-}\mu\text{-}\tau$  universality and uses

$$\Gamma_{\ell\ell} = \Gamma_{ee}, \quad (2)$$

$$B_{\ell\ell} = \text{average of } B_{ee}, B_{\mu\mu}, \text{ and } B_{\tau\tau}.$$

The electronic partial width  $\Gamma_{ee}$  is also not directly measurable at  $e^+e^-$  storage rings, only the combination  $\Gamma_{ee}\Gamma_{\text{had}}/\Gamma$ , where  $\Gamma_{\text{had}}$  is the hadronic partial width and

$$\Gamma_{\text{had}} + 3\Gamma_{ee} = \Gamma. \quad (3)$$

This combination is obtained experimentally from the energy-integrated hadronic cross section

$$\int_{\text{resonance}} \sigma(e^+e^- \rightarrow \Upsilon \rightarrow \text{hadrons}) dE = \frac{6\pi}{M^2} \frac{\Gamma_{ee}\Gamma_{\text{had}}}{\Gamma} \times C_r, \quad (4)$$

where  $M$  is the  $\Upsilon$  mass and  $C_r$  is the radiative correction. The knowledge of the  $B_{\ell\ell}$  then allows one to extract  $\Gamma_{ee}$  using Eq. (3):

$$\Gamma_{ee} = \frac{\Gamma_{ee}\Gamma_{\text{had}}/\Gamma}{(1 - 3B_{\ell\ell})}. \quad (5)$$

The Listings give experimental results only on  $B_{ee}$ ,  $B_{\mu\mu}$ ,  $B_{\tau\tau}$ , and  $\Gamma_{ee}$ . The  $\Gamma_{ee}$  entries are checked (within statistical errors) or re-evaluated for the proper  $(1 - 3B_{\ell\ell})$  correction in Eq. (5). The full widths of  $\Upsilon(1S)$ ,  $\Upsilon(2S)$ , and  $\Upsilon(3S)$  are then calculated from the averages for  $\Gamma_{ee}$  and  $B_{\ell\ell}$  using Eqs. (1) and (2). We no longer list  $\Gamma$  evaluations of individual experiments. Also note that the procedure adopted here relies only on experimental results and not on model assumptions like  $\Gamma_{ggg}[\Upsilon(2S)]/\Gamma_{ee}[\Upsilon(2S)] = \Gamma_{ggg}[\Upsilon(1S)]/\Gamma_{ee}[\Upsilon(1S)]$ , where  $\Gamma_{ggg}$  is the partial width for three-gluon decay. This assumption was made in, e.g., MAGERAS 81, NICZYPORUKZ 81, and in former editions of this Review.

$\Upsilon(9460)$   
or  $\Upsilon(1S)$

$$J^{G}(J^{PC}) = ?(1^{--})$$

#### $\Upsilon(9460)$ MASS (MeV)

M	Q	9460.6	0.4	ARTAMONO	84	REDE	E+E	-HADRONS
M	Q	9459.97	0.13	MAC KAY	84	REDE	E+E	-HADRONS
M	M	AVG	9460.03	0.19	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.5)			
M	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.						

#### $\Upsilon(9460)$ WIDTH (keV)

W	B	43.1	3.1					
W	B	OUR EVAL. FROM W2,R1,R2,R3 BELOW ASSUMING E-MU-TAU UNIVERSALITY						

#### $\Upsilon(9460)$ PARTIAL DECAY MODES

		DECAY MASSES	
P1	$\Upsilon(9460) \rightarrow \mu^+\mu^-$	106+	106
P2	$\Upsilon(9460) \rightarrow e^+e^-$	.511+	.511
P3	$\Upsilon(9460) \rightarrow \tau^+\tau^-$	1784+	1784
P	HADRONIC DECAYS		
P11	$\Upsilon(9460) \rightarrow \rho\pi$	769+	140
P12	$\Upsilon(9460) \rightarrow J/\psi(3097) + \text{anything}$		

#### $\Upsilon(9460)$ PARTIAL WIDTHS (keV)

$\Upsilon(9460) \rightarrow e^+e^-$								(G2)
W2	Q	1.33	0.14	BERGER	79	PLUT	E+E	-HADRONS
W2	Q	1.08	0.25	BOCK	80	CNTR	E+E	-HADRONS
W2	Q	1.23	0.16	ALBRECHT	82	DASP	E+E	-HADRONS
W2	Q	1.13	0.12	NICZYPOR	82	LENA	E+E	-HADRONS
W2	Q	1.15	0.11	TUTS	83	CUSB	E+E	-HADRONS
W2	Q	1.30	0.09	GILES	84	CLEO	E+E	-HADRONS
W2	AVG	1.224	0.050	AVERAGE				
W2	Q	SYSTEMATIC ERRORS ADDED QUADRATICALLY BY US.						

#### $\Upsilon(9460)$ BRANCHING RATIOS

$\Upsilon(9460) \rightarrow (\mu^+\mu^-)/\text{total}$								(P1)	
R1	Q	0.022	0.020	BERGER	79	PLUT	E+E	-MU+MU-	
R1	Q	0.014	0.034	0.014	BOCK	80	CNTR	E+E	-MU+MU-
R1	Q	0.032	0.013	ALBRECHT	82	DASP	E+E	-MU+MU-	
R1	Q	0.038	0.015	NICZYPOR	82	LENA	E+E	-MU+MU-	
R1	Q	0.027	0.004	ANDREWS	83	CLEO	E+E	-MU+MU-	
R1	Q	0.027	0.004	TUTS	83	CUSB	E+E	-MU+MU-	
R1	Q	0.029	0.004	BESSION	84	CLEO	Y2S	-PI+PI-ZMU	
R1	AVG	0.0278	0.0022	AVERAGE					
$\Upsilon(9460) \rightarrow (e^+e^-)/\text{total}$								(P2)	
R2	Q	0.051	0.030	BERGER	80	PLUT	E+E	-E+E-	
R2	Q	0.028	0.005	ALBRECHT	84	ARG	Y2S	-PI+PI-E+E-	
R2	Q	0.028	0.004	BESSION	84	CLEO	Y2S	-PI+PI-E+E-	
R2	AVG	0.0282	0.0031	AVERAGE					
$\Upsilon(9460) \rightarrow (\tau^+\tau^-)/\text{total}$								(P3)	
R3	Q	0.034	0.006	GILES	83	CLEO	E-E	-TAU-TAU-	
R3	A	0.0307	0.0051	ALBRECHT	85	ARG	Y2S	-PI+PI-2TAU	
R3	AVG	0.0321	0.0039	AVERAGE					
R3	A	USING BR(Y1S->EE)=BR(Y1S->MMU)=0.029, NOT USED FOR TOTAL							
R3	A	WIDTH EVALUATION.							
$\Upsilon(9460) \rightarrow (\rho\pi)/\text{total}$								(P11)	
R11	Q	(.0021)OR	LESS	CL=0.90	NICZYPOR	83	LENA		
$\Upsilon(9460) \rightarrow (J/\psi(3097) + \text{anything})/\text{total}$								(P12)	
R12	Q	(.02)	OR	LESS	CL=0.90	NICZYPOR	83	LENA	
R	Q	SYSTEMATIC ERRORS ADDED QUADRATICALLY BY US.							

#### REFERENCES FOR $\Upsilon(9460)$

COBB	77	PL	72	B	273	+IWATA,FABJAN,GOLDBERG+(BNL-CERN-SYRA-YALE)
HERB	77	PRL	39	252	+HOM,LEDERMAN,APPEL,ETO,-(COLU+PNAL-STON)	
INNES	77	PRL	39	1240	+APPEL,BROWN,HERB,HOM,FISK+(COLU+PNAL-STON)	
BERGER	78	PL	76	B	243	+ALEXANDER,DAUM,-(AACH+DESY-HAMB-SIEG-WUPG)
BIENLEIN	78	PL	78	B	360	+GLAWE,BOCK,BLANAR,+ (DESY+HAMB-HEID-MPI-M)
DARDEN	78	PL	76	B	246	+C.W.DARDEN+(DASPR-COLLAB.)
GARELICK	78	PR	D	18	945	+GAUTHIER,HICKS,OLIVER,+ (NEAS+WASH-TUFT)
KAPLAN	78	PRL	40	435	+APPEL,HERB,HOM,LEDERMAN,+ (STON+PNAL-COLU)	
YOH	78	PRL	41	684	+HERB,HOM,LEDERMAN,UENO,- (COLU+PNAL-STON)	

# Meson Full Listings

$\Upsilon(9460)$ ,  $\chi_{b0}(9860)$ ,  $\chi_{b1}(9895)$ ,  $\chi_{b2}(9915)$

ANGELIS 79 PL 87 B 398	+BESCH, BLUMENFELD,+ (CERN-COLU+OXF+ROCK)
BADIER 79 PL 86 B 98	+BOUCROT, BURGUN+ (SACL-CERN-CDF+EPOL-LALO)
BERGER 79 ZPHY C 1 343	+ALEXANDER+ (AACH-DESY+HAMB-SIEG-WUPG)
DARDEN 79 PL 80 B 419	C.W.DARDEN- (DASP2 COLLAB.)
ALBRECHT 80 PL 93 B 500	H. ALBRECHT- (DASP2 COLLAB.)
ANDREWS 80 PRL 44 1108	+ (CORN-HARV-ITHA+LEMO+ROCH+RUTG-SYRA-VAND)
BERGER 80 PL 93 B 497	+LACKAS, RAUPACH,+ (AACH+DESY+HAMB-SIEG-WUPP)
BOCK 80 ZPHY C 6 125	+BLANAR, BLUM, BIENLEIN+(HEID-MPI+DESY-HAMB)
BOHRINGER 80 PRL 44 1111	BOHRINGER, COSTANTINI, FINOCCHIARO+(COLU+STON)
KOURKOU 80 PL 91 B 481	KOURKOUMLIS+(ATHU+NTUA-BNL-CERN-SYRA+YALE)
MAGERAS 81 PRL 46 1115	-BOHRINGER, FINOCCHIARO+(COLU+STON+LSU+MPI+M)
MUELLER 81 PRL 46 1181	- (RUTG+SYRA+LEMO+VAND+CORN+ITHA+HARV+ROCH)
NICZYPOR 81 PRL 46 92	NICZYPORUK, CHEN, VOGEL, WEGENER+(LENA COLLAB)
ALBRECHT 82 PL 116 B 383	H. ALBRECHT+ (DASP2 COLLAB.)
ARTAMONO 82 PL 116 B 225	+BARU, BLINDV, BONDAR, BUKIN, GROSHEV- (NOVO)
NICZYPOR 82 ZPHY C 15 299	NICZYPORUK, FOLGER, BIENLEIN+ (LENA COLLAB)
ANDREWS 83 PRL 50 807	- (CORN+ITHA+HARV+OSU+ROCH+RUTG+SYRA+VAND)
GILES 83 PRL 50 877	- (HARV+OSU+ROCH+RUTG+SYRA+VAND+CORN+ITHA)
NICZYPOR 83 ZPHY C 17 197	- (CRAC+ERLA+DESY-NIJM-PITT+SACL+TELA+WURZ)
TUTS 83 CORNELL CONF. P. 248	P.M. TUTS+(CUSB COLLAB.)
ALBRECHT 84 PL 134 B 137	H. ALBRECHT- (ARGUS COLLAB.)
ARTAMONO 84 PL 137 B 272	ARTAMONOV+BARU+BLINDV+BONDAR+ (NDVOSI+IRSK)
BESSON 84 PR D 30 1433	+GREEN+HICKS+NAMJOSHI+SANNAS+(CLEO COLLAB.)
GILES 84 PR D 29 1285	+HASSARD+HEMPSTEAD+KINOSHITA+(CLEO COLLAB.)
MAC KAY 84 PR D 29 2483	+CUSB COLL.+ (HARV-CORN-COLU+LSU+MPI+ALBA)
ALBRECHT 85 PL 154 B 452	H. ALBRECHT+ (ARGUS COLLAB.)

$\chi_{b0}(9860)$   
or  $\chi_{b0}(1P)$

$$I^G(J^{PC}) = ?(0 \text{ preferred}^{++})$$

OBSERVED IN RADIATIVE DECAY OF THE UPSILON(10023), THEREFORE C=+. BRANCHING RATIO REQUIRES E1 TRANSITION, M1 IS STRONGLY DISFAVOURD, THEREFORE P=+.

### $\chi_{b0}(9860)$ MASS (MeV)

M U Q	(9872.8)	(5.1)	KLOPFENS 83 CUSB	Y2S->GAMMA X
M U Q	9864.1	7.1	HAAS1 84 CLEO	Y2S->CONV.GAM X
M U Q	9858.3	3.1	NERNST 85 CBAL	Y2S->GAMMA X
M U Q	9860.0	1.5	ALBRECH2 85 ARG	Y2S->CONV.GAM X
M	9859.8	1.3	AVERAGE	
M U	FROM GAMMA ENERGY BELOW, ASSUMING UPSILON(10023) MASS = 10023.4 MEV			
M Q	SYSTEMATIC ERRORS ADDED QUADRATICALLY BY US			

### $\gamma$ ENERGY IN $\Upsilon(10023)$ DECAY (MeV)

DM Q	(149.4)	(5.1)	KLOPFENS 83 CUSB	Y2S->GAMMA X
DM Q	158.0	7.1	HAAS1 84 CLEO	Y2S->CONV.GAM X
DM Q	163.8	3.1	NERNST 85 CBAL	Y2S->GAMMA X
DM Q	162.1	1.5	ALBRECH2 85 ARG	Y2S->CONV.GAM X
DM	162.3	1.3	AVERAGE	
DM Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.			

### $\chi_{b0}(9860)$ PARTIAL DECAY MODES

P1	$\chi_{b0}(9860) \rightarrow \Upsilon(9460) \gamma$	9460, 0
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### $\chi_{b0}(9860)$ BRANCHING RATIOS

$\chi_{b0}(9860) \rightarrow (\Upsilon(9460) \gamma)/\text{total}$	(P1)	
R1	(0.11) OR LESS CL=.90 PAUSS 83 CUSB	Y2S->2 GAM L+L-
R1	(0.06) OR LESS CL=.90 WALK 85 CBAL	Y2S->2 GAM L+L-

### REFERENCES FOR $\chi_{b0}(9860)$

KLOPFENS 83 PRL 51 160	KLOPFENSTEIN+HORSTKOTTE+ (CUSB COLLAB.)
PAUSS 83 PL 130 B 439	+DIETL,EIGEN+ (MPIM+COLU+CORN+LSU+STON)
HAAS1 84 PRL 52 799	+JENSEN+KAGAN+KASS-BEHREND+ (CLEO COLLAB.)
ALBRECH2 85 PL 160 B 331	H. ALBRECHT- (ARGUS COLLAB.)
NERNST 85 PRL 54 2195	-ANTREASYAN+ASCHMAN+ (CRYSTAL BALL COLLAB.)
WALK 85 SLAC-PUB-3820	+ZSCHORSCH+ (CRYSTAL BALL COLLAB.)

$\chi_{b1}(9895)$   
or  $\chi_{b1}(1P)$

$$I^G(J^{PC}) = ?(1 \text{ preferred}^{++})$$

OBSERVED IN RADIATIVE DECAY OF THE UPSILON(10023), THEREFORE C=+. BRANCHING RATIO REQUIRES E1 TRANSITION, M1 IS STRONGLY DISFAVOURD, THEREFORE P=+.

### $\chi_{b1}(9895)$ MASS (MeV)

M U Q	9894.4	3.0	KLOPFENS 83 CUSB	Y2S->GAMMA X
M U Q	9892.0	3.0	PAUSS 83 CUSB	Y2S->2 GAM L+L-
M U Q	9893.6	1.3	HAAS1 84 CLEO	Y2S->CONV.GAM X
M U Q	9892.0	2.5	NERNST 85 CBAL	Y2S->GAMMA X
M U Q	9890.8	1.1	ALBRECH2 85 ARG	Y2S->CONV.GAM X
M U Q	9890.8	1.6	WALK 85 CBAL	Y2S->2 GAM L+L-
M AVG	9891.89	0.68	AVERAGE	
M U	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US			
M Q	FROM GAMMA ENERGY BELOW ASSUMING UPSILON(10023) MASS = 10023.4 MEV			

### $\gamma$ ENERGY IN $\Upsilon(10023)$ DECAY (MeV)

DM Q	128.1	3.0	KLOPFENS 83 CUSB	Y2S->GAMMA X
DM Q	130.6	3.0	PAUSS 83 CUSB	Y2S->2 GAM L+L-
DM Q	129.0	1.3	HAAS1 84 CLEO	Y2S->CONV.GAM X
DM Q	130.6	2.5	NERNST 85 CBAL	Y2S->GAMMA X
DM Q	131.7	1.1	ALBRECH2 85 ARG	Y2S->CONV.GAM X
DM Q	131.7	1.6	WALK 85 CBAL	Y2S->2 GAM L+L-
DM AVG	130.65	0.68	AVERAGE	
DM Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.			

### $\chi_{b1}(9895)$ PARTIAL DECAY MODES

P1	$\chi_{b1}(9895) \rightarrow \Upsilon(9460) \gamma$	9460, 0
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### $\chi_{b1}(9895)$ BRANCHING RATIOS

$\chi_{b1}(9895) \rightarrow (\Upsilon(9460) \gamma)/\text{total}$	(P1)		
R1	0.47 0.18	KLOPFENS 83 CUSB	Y2S->2 GAM L+L-
R1	0.32 0.09	WALK 85 CBAL	Y2S->2 GAM L+L-
R1 AVG	0.350 0.080	AVERAGE	
R1 Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US		

### REFERENCES FOR $\chi_{b1}(9895)$

KLOPFENS 83 PRL 51 160	KLOPFENSTEIN+HORSTKOTTE+ (CUSB COLLAB.)
PAUSS 83 PL 130 B 439	+DIETL,EIGEN+ (MPIM+COLU+CORN+LSU+STON)
HAAS1 84 PRL 52 799	+JENSEN+KAGAN+KASS-BEHREND+ (CLEO COLLAB.)
ALBRECH2 85 PL 160 B 331	H. ALBRECHT- (ARGUS COLLAB.)
NERNST 85 PRL 54 2195	-ANTREASYAN+ASCHMAN+ (CRYSTAL BALL COLLAB.)
WALK 85 SLAC-PUB-3820	+ZSCHORSCH+ (CRYSTAL BALL COLLAB.)

$\chi_{b2}(9915)$   
or  $\chi_{b2}(1P)$

$$I^G(J^{PC}) = ?(2 \text{ preferred}^{++})$$

OBSERVED IN RADIATIVE DECAY OF THE UPSILON(10023), THEREFORE C=+. BRANCHING RATIO REQUIRES E1 TRANSITION, M1 IS STRONGLY DISFAVOURD, THEREFORE P=+.

### $\chi_{b2}(9915)$ MASS (MeV)

M U Q	9914.6	2.0	KLOPFENS 83 CUSB	Y2S->GAMMA X
M U Q	9914.0	4.0	PAUSS 83 CUSB	Y2S->2 GAM L+L-
M U Q	9913.3	1.2	HAAS1 84 CLEO	Y2S->CONV.GAM X
M U Q	9912.4	2.3	NERNST 85 CBAL	Y2S->GAMMA X
M U Q	9912.2	1.0	ALBRECH2 85 ARG	Y2S->CONV.GAM X
M U Q	9915.8	1.7	WALK 85 CBAL	Y2S->2 GAM L+L-
M AVG	9913.29	0.63	AVERAGE	
M U	SYSTEMATIC ERRORS ADDED QUADRATICALLY BY US			
M Q	FROM GAMMA ENERGY BELOW ASSUMING UPSILON(10023) MASS = 10023.4 MEV			



# Meson Full Listings

$\chi_{b0}(10235)$ ,  $\chi_{b1}(10255)$ ,  $\chi_{b2}(10270)$ ,  $\Upsilon(10355)$

$\chi_{b0}(10235)$   
or  $\chi_{b0}(2P)$

$$I^G(J^{PC}) = ?(0 \text{ preferred}^{+-})$$

OBSERVED IN RADIATIVE DECAY OF THE UPSILON(10355), THEREFORE C=+. BRANCHING RATIO REQUIRES E1 TRANSITION, M1 IS STRONGLY DISFAVOURER, THEREFORE P=+. NEEDS CONFIRMATION.

OMITTED FROM SUMMARY TABLE

### $\chi_{b0}(10235)$ MASS (GeV)

M	U	10.2327	.0050					
M	U	Q						
M	U	AVG						

FROM GAMMA ENERGY BELOW ASSUMING UPSILON(10355) MASS = 10355.5 MEV

### $\gamma$ ENERGY IN $\Upsilon(10355)$ DECAY (MeV)

DM	Q	122.1	5.0	TUTS	83	CUSB	E+E-->2 GAM L+L-
DM	Q						E+E-->GAMMA X

SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.

### $\chi_{b0}(10235)$ PARTIAL DECAY MODES

			DECAY MASSES
P1	$\chi_{b0}(10235) \rightarrow \Upsilon(9460) \gamma$		9460+ 0
P2	$\chi_{b0}(10235) \rightarrow \Upsilon(10023) \gamma$		10023+ 0

### REFERENCES FOR $\chi_{b0}(10235)$

EIGEN	82	PRL	49	1616	+BOHRINGER, HERB. (MPIM+COLU+CORN+STON+LSU)
HAN	82	PRL	49	1612	+HORSTKOTTE, IMLAY. (COLU+STON+CORN+LSU+MPIM)
TUTS	83	CORNELL CONF.	P.284		P.M.TUTS(CUSB COLLAB.)

$\chi_{b1}(10255)$   
or  $\chi_{b1}(2P)$

$$I^G(J^{PC}) = ?(1 \text{ preferred}^{+-})$$

OBSERVED IN RADIATIVE DECAY OF THE UPSILON(10355), THEREFORE C=+. BRANCHING RATIO REQUIRES E1 TRANSITION, M1 IS STRONGLY DISFAVOURER, THEREFORE P=+.

### $\chi_{b1}(10255)$ MASS (GeV)

M	U	10.2560	.0020	EIGEN	82	CUSB	E+E-->2 GAM L+L-
M	U	Q		TUTS	83	CUSB	E+E-->GAMMA X
M	U	AVG					

FROM GAMMA ENERGY BELOW ASSUMING UPSILON(10355) MASS = 10355.5 MEV. SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.

### $\gamma$ ENERGY IN $\Upsilon(10355)$ DECAY (MeV)

DM	Q	99.0	2.0	EIGEN	82	CUSB	E+E-->2 GAM L+L-
DM	Q	101.4	3.0	TUTS	83	CUSB	E+E-->GAMMA X
DM	Q	AVG					

SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.

### $\chi_{b1}(10255)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$\chi_{b1}(10255) \rightarrow \Upsilon(9460) \gamma$	9460+ 0
P2	$\chi_{b1}(10255) \rightarrow \Upsilon(10023) \gamma$	10023+ 0

### REFERENCES FOR $\chi_{b1}(10255)$

EIGEN	82	PRL	49	1616	+BOHRINGER, HERB. (MPIM+COLU+CORN+STON+LSU)
HAN	82	PRL	49	1612	+HORSTKOTTE, IMLAY. (COLU+STON+CORN+LSU+MPIM)
TUTS	83	CORNELL CONF.	P.284		P.M.TUTS(CUSB COLLAB.)

$\chi_{b2}(10270)$   
or  $\chi_{b2}(2P)$

$$I^G(J^{PC}) = ?(2 \text{ preferred}^{+-})$$

OBSERVED IN RADIATIVE DECAY OF THE UPSILON(10355), THEREFORE C=+. BRANCHING RATIO REQUIRES E1 TRANSITION, M1 IS STRONGLY DISFAVOURER, THEREFORE P=+.

### $\chi_{b2}(10270)$ MASS (GeV)

M	U	10.2712	.0030	EIGEN	82	CUSB	E+E-->2 GAM L+L-
M	U	Q		TUTS	83	CUSB	E+E-->GAMMA X
M	U	AVG					

FROM GAMMA ENERGY BELOW ASSUMING UPSILON(10355) MASS = 10355.5 MEV. SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.

### $\gamma$ ENERGY IN $\Upsilon(10355)$ DECAY (MeV)

DM	Q	84.0	3.0	EIGEN	82	CUSB	E+E-->2 GAM L+L-
DM	Q	84.2	2.0	TUTS	83	CUSB	E+E-->GAMMA X
DM	Q	AVG					

SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.

### $\chi_{b2}(10270)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$\chi_{b2}(10270) \rightarrow \Upsilon(9460) \gamma$	9460+ 0
P2	$\chi_{b2}(10270) \rightarrow \Upsilon(10023) \gamma$	10023+ 0

### REFERENCES FOR $\chi_{b2}(10270)$

EIGEN	82	PRL	49	1616	+BOHRINGER, HERB. (MPIM+COLU+CORN+STON+LSU)
HAN	82	PRL	49	1612	+HORSTKOTTE, IMLAY. (COLU+STON+CORN+LSU+MPIM)
TUTS	83	CORNELL CONF.	P.284		P.M.TUTS(CUSB COLLAB.)

$\Upsilon(10355)$   
or  $\Upsilon(3S)$

$$I^G(J^{PC}) = ?(1^{--})$$

### $\Upsilon(10355)$ MASS (GeV)

M	10.3555	0.0005	ARTAMOND	84	REDE	E+E-->HADRONS
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### $\Upsilon(10355)$ WIDTH (keV)

W	B	12.0	10.0	4.0
W	B	OUR EVAL.	FROM W2,R1	BELOW ASSUMING E-MU-UNIVERSALRTY

### $\Upsilon(10355)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$\Upsilon(10355) \rightarrow \mu^+ \mu^-$	106+ 106
P2	$\Upsilon(10355) \rightarrow e^+ e^-$	.511-.511
P3	$\Upsilon(10355) \rightarrow \pi^+ \pi^- \Upsilon(9460)$	140+ 140+9460
P4	$\Upsilon(10355) \rightarrow \pi^+ \pi^- \Upsilon(10023)$	140+ 140+10023
P5	$\Upsilon(10355) \rightarrow \chi_{b2}(10270) \gamma$	10271+ 0
P6	$\Upsilon(10355) \rightarrow \chi_{b1}(10255) \gamma$	10254+ 0
P7	$\Upsilon(10355) \rightarrow \chi_{b0}(10235) \gamma$	10233+ 0

### $\Upsilon(10355)$ PARTIAL WIDTHS (keV)

$\Upsilon(10355) \rightarrow (e^+ e^-)$					(62)		
W2	Q	0.39	0.04	TUTS	83	CUSB	E+E-->HADRONS
W2	Q	0.42	0.05	GILES	84	CLEO	E+E-->HADRONS
W2	Q	AVG					

SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.

For notation, see key on page 91.

# Meson Full Listings

$\Upsilon(10355), \Upsilon(10575), \Upsilon(10860), \Upsilon(11020)$

### $\Upsilon(10355)$ BRANCHING RATIOS

$\Upsilon(10355) \rightarrow (\mu^+ \mu^-)/\text{total}$							
R1	Q	0.033	0.015	ANDREWS 83 CLEO	E+E-	--MU, MU-	(P1)
$\Upsilon(10355) \rightarrow (\Upsilon(9460) \pi^+ \pi^-)/\text{total}$							
R3	Q	0.22	0.049	0.010	GREEN 82 CLEO	Y3S-->PI+PI-L-L-	(P3)
R3		26	0.039	0.013	MAGERAS 82 CUSB	Y3S-->PI+PI-L-L-	
R3	AVG	0.0453	0.0079	AVERAGE			
$\Upsilon(10355) \rightarrow (\Upsilon(10023) \pi^+ \pi^-)/\text{total}$							
R4		5	0.031	0.020	MAGERAS 82 CUSB	Y3S-->PI+PI-L-L-	(P4)
$\Upsilon(10355) \rightarrow (X_{b2}(10270) \gamma)/\text{total}$							
R5			0.127	0.041	TUTS 83 CUSB	E+E-	--GAMMA X (P5)
$\Upsilon(10355) \rightarrow (X_{b1}(10255) \gamma)/\text{total}$							
R6			0.156	0.042	TUTS 83 CUSB	E+E-	--GAMMA X (P6)
$\Upsilon(10355) \rightarrow (X_{b0}(10235) \gamma)/\text{total}$							
R7			0.076	0.035	TUTS 83 CUSB	E+E-	--GAMMA X (P7)
R	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US					

### REFERENCES FOR $\Upsilon(10355)$

COBB 77 PL 72 B 273	+IWATA, FABIAN, GOLDBERG-(BNL+CERN+SYRA+YALE)
HERB 77 PRL 39 252	+HOM, LEDERMAN, APPEL, ITO, -(COLU+FNAL+STON)
INNES 77 PRL 39 1240	+APPEL, BROWN, HERB, HOM, FISK-(COLU+FNAL+STON)
KAPLAN 78 PRL 40 435	+APPEL, HERB, HOM, LEDERMAN, +(STON+FNAL+COLU)
YOH 78 PRL 41 684	+HERB, HOM, LEDERMAN, UENO, -(COLU+FNAL+STON)
UENO 79 PRL 42 486	+BROWN, HERB, HOM, FISK, ITO, -(FNAL+COLU+STON)
ANDREWS 80 PRL 44 1108	+ (CORN+HARV+ITHA+LEMO+ROCH+RUTG+SYRA+VAND)
BOHRINGER 80 PRL 44 1111	BOHRINGER, COSTANTINI, FINOCCHIARO(COLU+STON)
GREEN 82 PRL 49 617	+ (RUTG+SYRA+VAND+CORN+ITHACA+HARV+OSU+ROCH)
HAN 82 PRL 49 1612	+HORSTKOTTE, IMLAY-(COLU+STON+CORN+LSU+MPIM)
MAGERAS 82 PL 118 B 453	+HERB, IMLAY, -(COLU+CORN+LSU+MPIM+STON)
PETERSON 82 PL 114 B 277	+GIANNINI, LEE-FRANZINI-(COLU+STON+LSU+MPIM)
ANDREWS 83 PRL 50 807	+ (CORN+ITHA+HARV+OSU+ROCH+RUTG+SYRA+VAND)
TUTS 83 CORNELL CONF. P. 284	P.M. TUTS(CUSB COLLAB.)
ARTAMONOV 84 PL 137 B 272	ARTAMONOV+BARU+BLINOV+BUNDAR-(NOVOSIBIRSK)
GILES 84 PR D 29 1285	-HASSARD-HEMPSTEAD-KINOSHITA-(CLEO COLLAB.)

$\Upsilon(10575)$   
or  $\Upsilon(4S)$

$$J^G(J^{PC}) = ?(1^{--})$$

### $\Upsilon(10575)$ MASS (GeV)

M	A	Q	10.5775	0.0058	BESSON 85 CLEO	E+E-	--HADRONS
M	B	Q	10.5774	0.0058	LOVELOCK 85 CUSB	E+E-	--HADRONS
M	AVG		10.5775	0.0041	AVERAGE		
M	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US					
M	A	SYS.ERR.=0.004 INCREASED FOR AVERAGING PURPOSE, SEE PROCEDURES C2					
M	B	NO SYSTEMATIC ERROR GIVEN, WE ASSUME IT EQUAL TO BESSON 85					

### $\Upsilon(10575)$ WIDTH (MeV)

W	Q	20.0	4.5	BESSON 85 CLEO	E+E-	--HADRONS	
W		25.	2.5	LOVELOCK 85 CUSB	E+E-	--HADRONS	
W	AVG	23.8	2.2	AVERAGE			
W	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.					

### $\Upsilon(10575)$ PARTIAL DECAY MODES

P1	$\Upsilon(10575) \rightarrow e^+ e^-$	DECAY MASSES	.511+.511
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### $\Upsilon(10575)$ PARTIAL WIDTHS (keV)

$\Upsilon(10575) \rightarrow (e^+ e^-)$				(61)			
W1	Q	0.192	0.039	BESSON 85 CLEO	E+E-	--HADRONS	
W1		0.283	0.037	LOVELOCK 85 CUSB	E+E-	--HADRONS	
W1	AVG	0.240	0.045	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.7)			
W1	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US					

### REFERENCES FOR $\Upsilon(10575)$

ANDREWS 80 PRL 45 219	+ (CORN+HARV+ITHA+LEMO+ROCH+RUTG+SYRA+VAND)
FINOCCHI 80 PRL 45 222	FINOCCHIARO, GIANNINI, BOHRINGER, +(COLU+STON)
BESSON 85 PRL 54 381	+GREEN+NAMJOSHI+SANNES+SKUBIC+(CLEO COLLAB)
LOVELOCK 85 PRL 54 377	+HORSTKOTTE+KLOPFENSTEIN+(CUSB COLLAB)

$\Upsilon(10860)$   
or  $\Upsilon(5S)$

$$J^G(J^{PC}) = ?(1^{--})$$

### $\Upsilon(10860)$ MASS (GeV)

M	Q	10.868	0.008	BESSON 85 CLEO	E+E-	--HADRONS	
M		10.845	0.020	LOVELOCK 85 CUSB	E+E-	--HADRONS	
M	AVG	10.8648	0.0079	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)			
M	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US					

### $\Upsilon(10860)$ WIDTH (MeV)

W	Q	112.0	29.0	BESSON 85 CLEO	E+E-	--HADRONS	
W		110.0	15.0	LOVELOCK 85 CUSB	E+E-	--HADRONS	
W	AVG	110.4	13.3	AVERAGE			
W	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.					

### $\Upsilon(10860)$ PARTIAL DECAY MODES

P1	$\Upsilon(10860) \rightarrow e^+ e^-$	DECAY MASSES	.511+.511
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### $\Upsilon(10860)$ PARTIAL WIDTHS (keV)

$\Upsilon(10860) \rightarrow (e^+ e^-)$				(61)			
W1	Q	0.22	0.09	BESSON 85 CLEO	E+E-	--HADRONS	
W1		0.365	0.070	LOVELOCK 85 CUSB	E+E-	--HADRONS	
W1	AVG	0.310	0.070	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3)			
W1	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US					

### REFERENCES FOR $\Upsilon(10860)$

BESSON 85 PRL 54 381	+GREEN+NAMJOSHI+SANNES+SKUBIC+(CLEO COLLAB)
LOVELOCK 85 PRL 54 377	+HORSTKOTTE+KLOPFENSTEIN+(CUSB COLLAB)

$\Upsilon(11020)$   
or  $\Upsilon(6S)$

$$J^G(J^{PC}) = ?(1^{--})$$

### $\Upsilon(11020)$ MASS (GeV)

M	Q	11.019	0.009	BESSON 85 CLEO	E+E-	--HADRONS	
M		11.020	0.030	LOVELOCK 85 CUSB	E+E-	--HADRONS	
M	AVG	11.0191	0.0086	AVERAGE			
M	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US					

### $\Upsilon(11020)$ WIDTH (MeV)

W	Q	61.0	26.0	BESSON 85 CLEO	E+E-	--HADRONS	
W		90.0	20.0	LOVELOCK 85 CUSB	E+E-	--HADRONS	
W	AVG	79.2	15.9	AVERAGE			
W	Q	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US.					

### $\Upsilon(11020)$ PARTIAL DECAY MODES

P1	$\Upsilon(11020) \rightarrow e^+ e^-$	DECAY MASSES	.511+.511
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# Meson Full Listings

$\Upsilon(11020), K^\pm, K^0, K^*(892)$

### $\Upsilon(11020)$ PARTIAL WIDTHS (keV)

$\Upsilon(11020) \rightarrow (e^+ e^-)$  (G1)

W1	Q	0.095	0.046	BESSON	85	CLEO	E+E- -->HADRONS
W1		0.156	0.040	LOVELOCK	85	CUSB	E+E- -->HADRONS
W1	AVG	0.130	0.030	AVERAGE			

W1 Q SYSTEMATIC ERROR ADDED QUADRATICALLY BY US

### REFERENCES FOR $\Upsilon(11020)$

BESSON 85 PRL 54 381 -GREEN+NAIJOSHI+SANNES+SKUBIC+(CLEO COLLAB)  
 LOVELOCK 85 PRL 54 377 -HORSTKOTTE+KLOPFENSTEIN+ (CUSB COLLAB)

$S = \pm 1, C = 0, B = 0$  MESON STATES

$K^\pm$

$$I(J^P) = \frac{1}{2}(0^-)$$

SEE STABLE PARTICLE FULL LISTINGS

$K^0$

$$I(J^P) = \frac{1}{2}(0^-)$$

SEE STABLE PARTICLE FULL LISTINGS

$K^*(892)$

$$I(J^P) = \frac{1}{2}(1^-)$$

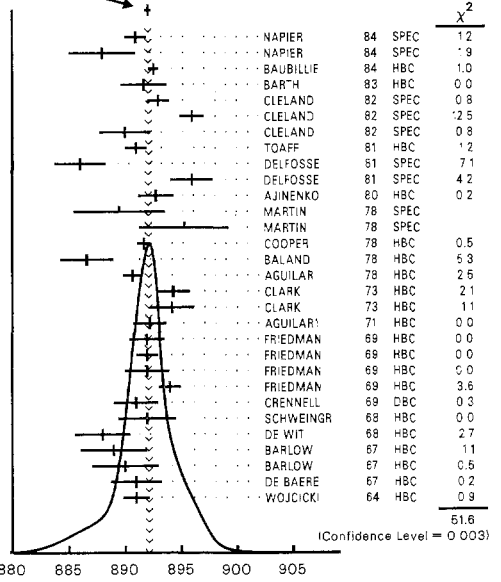
### $K^*(892)$ MASS (MeV)

CHARGED ONLY. THIS IS WHAT APPEARS ON MESON SUMMARY TABLE							
M W	1700	891.0	1.2	WOJCICKI	64	HBC	- 1.7 K-P(K0 P1-)
M D	620	891.	2.5	DE BAERE	67	HBC	+ 3.5 K-P (K0 P1+)
M	720	890.	3.0	BARLOW	67	HBC	-- 1.2 PBAR(K0 P1)
M	600	889.	3.0	BARLOW	67	HBC	-- 1.2 PBAR(K P1)
M D	540	888.	2.5	DE WIT	68	HBC	- 3.0 K-P
M D	341	892.0	2.6	SCHWEINGR	68	HBC	- 5.5 K-P(K0 P1-)
M	1000	891.0	2.0	CRENNELL	69	DBC	- 5.9 K-N (K0 P1-)
M	2886	894.	1.0	FRIEDMAN	69	HBC	- 2.1 K-P(K0 P1-)
M	728	892.	2.	FRIEDMAN	69	HBC	- 2.45 K-P(K0 P1-)
M	3229	892.	1.0	FRIEDMAN	69	HBC	- 2.6 K-P(K0 P1-)
M D	1027	892.	1.6	FRIEDMAN	69	HBC	- 2.7 K-P(K0 P1-)
M	4404	892.2	1.5	AGUILARI	71	HBC	- 3.9,4.6 K-P
M D	765	894.2	2.0	CLARK	73	HBC	- 3.13 K-P(K0 P1-)
M W	01150	894.3	1.5	CLARK	73	HBC	- 3.5 K-P, P P1-K0
M I	9000	(891.9)	(0.7)	PALER	75	HBC	- 14.3 K-P, K- X+
M	1800	890.7	0.9	AGUILAR	78	HBC	-- .76 PB P, K KS P1
M	1225	886.6	2.4	BALAND	78	HBC	-- 12 PB P, INCLUSIV
M X	6706	891.7	0.6	COOPER	78	HBC	-- .76 PB P, INCLUSIV
M X		895.3	4.0	MARTIN	78	SPEC	+ 10 K+-P, KS P1 P
M X		889.5	4.1	MARTIN	78	SPEC	+ 10 K+-P, KS P1 P
M		892.8	1.6	AJINENKO	80	HBC	- 32 K-P
M	380	896.0	1.9	DELFOSSÉ	81	SPEC	- K-P, K+- P10 P
M	187	886.0	2.3	DELFOSSÉ	81	SPEC	- K-P, K+- P10 P
M	4100	891.0	1.0	TOAFF	81	HBC	- 6.5 K-P, K0 P1-P
M W D	800	890.0	2.3	CLELAND	82	SPEC	+ 30 K+-P, KS P1-P
M W D	3200	896.0	1.1	CLELAND	82	SPEC	+ 50 K+-P, KS P1-P
M W D	3500	895.0	1.0	CLELAND	82	SPEC	- 50 K+-P, KS P1-P
M	3700	891.7	2.1	BARTH	83	HBC	+ 70 K+-P, K0 P1+ X
M	5840	892.6	0.5	BAUBILLIE	84	HBC	- 8.25 K-P, K0 P1-P
M		888.0	3.0	NAPIER	84	SPEC	- 200 P1-P, 2K0 X
M		891.0	1.0	NAPIER	84	SPEC	- 200 P1-P, 2K0 X
M AVG		892.11	0.32	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4) (SEE IDEOGRAM BELOW)			

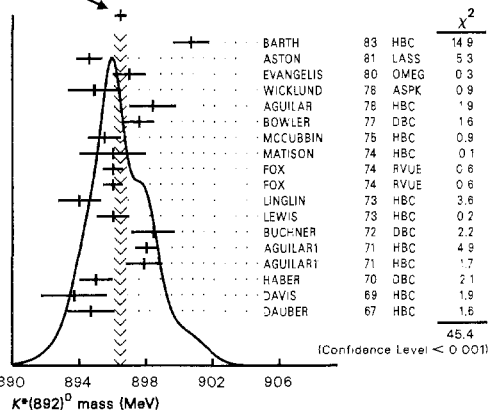
NEUTRAL ONLY.							
M	91040	894.7	1.4	DAUSER	67	HBC	02.0 K-P(K+P1-)
M	10K	893.7	2.0	DAVIS	69	HBC	0 12. K+P(K+P1-)
M W	4300	895.0	1.0	HABER	70	DBC	0 3. K-N (K+P1+)
M D	2934	897.9	1.1	AGUILARI	71	HBC	0 3.9,4.6 K- P
M	D5362	898.0	(0.7)	AGUILAR	71	HBC	0 3.9,4.6 K- P
M	01700	898.4	1.3	BUCHNER	72	DBC	04.6 K+ N, K+ P1-
M	3186	896.0	1.0	LEWIS	73	HBC	02.1-2.7 K+P
M C		894.0	1.3	LINGLIN	73	HBC	02-13 K+P, K+P1-
M W	10K	895.0	0.6	FOX	74	RVUE	0 2 K+P, K+P1-P
M		896.0	0.6	FOX	74	RVUE	0 2 K+N, K+P1-P
M C	896.		2.	MATISON	74	HBC	012 K+P, K+P1-
M	3600	895.5	1.0	MCCUBBIN	75	HBC	0 3.6 K+K, K+P1-N
M I	2200	(892.7)	(0.7)	PALER	75	HBC	014.3 K-P, *D X0
M		897.6	0.9	BOWLER	77	DBC	05.4 K+D, K+P1-P
M	1180	898.4	1.4	AGUILAR	78	HBC	0 .76 PB P, K KS P1
M P		(894.9)	(0.3)	ESTABROOK	78	ASPK	0 15 K+-P, K+P1+-
M C		(892.8)	(1.3)	WICKLUND	78	ASPK	0 3,4,6 P1+-PN
M	28K	897.	1.	LANG	79	RVUE	0
M		897.6	1.	EVANGELIS	80	DMEG	0 10 P1-P
M	5900	900.7	1.1	ASTON	81	LASS	0 11 K-P, K+ P1- X
M		891.0	1.0	BARTH	83	HBC	0 70 K+-P, K- P1- X
M AVG		896.45	0.37	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.6) (SEE IDEOGRAM BELOW)			

M C FROM POLE EXTRAPOLATION.  
 M D MASS ERRORS ENLARGED BY US TO GAMMA/SQRT(N). SEE TYPED NOTE.  
 M I INCLUSIVE REACTION. COMPLICATED BACKGROUND AND PHASE-SPACE EFFECTS  
 M P FROM PHASE SHIFT ANALYSIS OF 155000 EVENTS.  
 M W NUMBER OF EVENTS IN PEAK REEVALUATED BY US  
 M X SYSTEMATIC ERROR ADDED

WEIGHTED AVERAGE  
 892.11 ± 0.32 (ERROR SCALED BY 1.4)



WEIGHTED AVERAGE  
 896.45 ± 0.37 (ERROR SCALED BY 1.6)



### NOTE ON $K^*(892)$ MASSES AND MASS DIFFERENCES

Unrealistically small errors are reported by some experiments. We use simple "realistic" tests for the minimum errors on the determination of mass and width from a sample of  $N$  events:















For notation, see key on page 91.

# Meson Full Listings

$K_2(1580)$ ,  $K_2(1770)$

## $K_2(1580)$ WIDTH (MeV)

W	(110.)	APPROX.	OTTER	79	-	10,14,16	K-P
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## $K_2(1580)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$K_2(1580) \rightarrow K^*(892) \pi$	892+ 140
P2	$K_2(1580) \rightarrow K_2^*(1430) \pi$	1425+ 140

## $K_2(1580)$ BRANCHING RATIOS

$K_2(1580) \rightarrow (K^*(892) \pi)/total$					(P1)
W1	SEEN	OTTER	79	HBC	- 10,14,16 K-P
$K_2(1580) \rightarrow (K_2^*(1430) \pi)/total$					(P2)
W2	POSSIBLY SEEN	OTTER	79	HBC	- 10,14,16 K-P

## REFERENCES FOR $K_2(1580)$

OTTER	79	NP B 147 1	-	RUDDOLPH, +	(AACH+BERL+CERN-LOIC+WIEN)JP
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$K_2(1770)$   
was  $L(1770)$

$$I(J^P) = \frac{1}{2}(2^-)$$

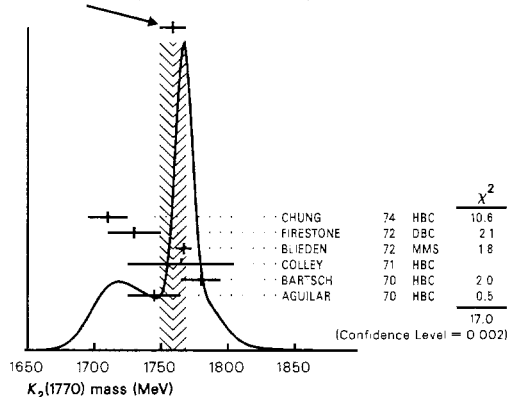
OUR LATEST MINIREVIEW ON THIS PARTICLE CAN BE FOUND IN THE 1984 EDITION.

## $K_2(1770)$ MASS (MeV)

M	1745.0	20.0	AGUILAR	70	HBC	-	4.6	K-P
M	1780.0	15.0	BARTSCH	70	HBC	-	10.1	K-P
M	(1760.0)	(15.0)	LUDLAM	70	HBC	-	12.6	K-P
M X	1765.0	40.0	COLLEY	71	HBC	+	10.0	K-P, K 2PI D
M	(1740.0)		DENERGI	71	DBC	-	12.6	K-D, K 2PI D
M	1767.	6.	BLIEDEN	72	MMS	-	11-16.	K-P
M P	306 1730.	20.	FIRESTONE	72	DBC	-	12.	K+D
M	60 1710.	15.	CHUNG	74	HBC	-	7.3	K-P, K-OMEGA P
M	(1820.)	APPROX.	DAUM	81	CNTR	-	63	K-P, K 2PI P
M	(1750.)	APPROX.	ARMSTRONG	83	OMEG	-	18.5	K-P, 3K P
M	AVG	1758.9	10.0					

M P PRODUCED IN CONJUNCTION WITH EXCITED DEUTERON.  
M X SYSTEMATIC ERRORS ADDED CORRESP. TO SPREAD OF DIFFERENT FITS.

WEIGHTED AVERAGE  
1758.9 ± 10.0 (ERROR SCALED BY 2.1)

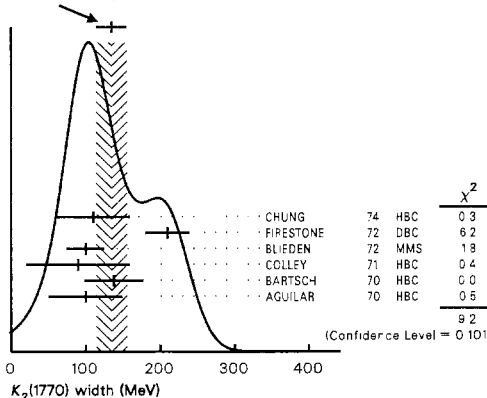


## $K_2(1770)$ WIDTH (MeV)

W	100.0	50.0	AGUILAR	70	HBC	-	4.6	K-P
W	138.0	40.0	BARTSCH	70	HBC	-	10.1	K-P
W	(50.0)	(40.0)	LUDLAM	70	HBC	-	12.6	K-P
W X	90.	70.	COLLEY	71	HBC	+	10.0	K-P, K 2PI D
W	(130.0)		DENERGI	71	DBC	-	12.6	K-D, K 2PI D
W	100.	26.	BLIEDEN	72	MMS	-	11-16.	K-P
W P	306 210.	50.	FIRESTONE	72	DBC	-	12.	K+D
W	60 110.		CHUNG	74	HBC	-	7.3	K-P, K-OMEGA P
W	(200.)	APPROX.	DAUM	81	CNTR	-	63	K-P, K 2PI P
W	(220.)	APPROX.	ARMSTRONG	83	OMEG	-	18.5	K-P, 3K P
W	AVG	135.1	20.9					

W P PRODUCED IN CONJUNCTION WITH EXCITED DEUTERON  
W X SYSTEMATIC ERRORS ADDED CORRESP. TO SPREAD OF DIFFERENT FITS.

WEIGHTED AVERAGE  
135.1 ± 20.9 (ERROR SCALED BY 1.4)



## $K_2(1770)$ PARTIAL DECAY MODES

	DECAY MASSES	
P1	$K_2(1770) \rightarrow K \pi \pi$	498+ 135+ 135
P2	$K_2(1770) \rightarrow K_2^*(1430) \pi$	135+1425
P3	$K_2(1770) \rightarrow K \pi \pi \pi$	498+ 135+ 135+ 135
P4	$K_2(1770) \rightarrow K^*(892) \pi$	892- 135
P5	$K_2(1770) \rightarrow K^*(892) \rho$	892- 769
P6	$K_2(1770) \rightarrow K^*(892) \omega$	892+ 783
P7	$K_2(1770) \rightarrow K^*(892) \pi \pi$	892+ 135+ 135
P8	$K_2(1770) \rightarrow K \omega$	498- 783
P9	$K_2(1770) \rightarrow K f_2(1270)$	498-1274
P10	$K_2(1770) \rightarrow K \phi$	494+1020

## $K_2(1770)$ BRANCHING RATIOS

$K_2(1770) \rightarrow (K_2^*(1430) \pi) / (K \pi \pi)$	(P2)/(P1)
R1	(K2*(1430) INTO K PI)
R1	(1.0)
R1	0.2
R1	(1.0) OR LESS
R1	(1.0) OR LESS
R1 P	(1.0) APPROX.
R1	(0.6) APPROX.

$K_2(1770) \rightarrow (K \omega)/total$	(P8)
R2	SEEN
R2	SEEN

$K_2(1770) \rightarrow (K^*(892) \pi)/(K \pi \pi)$	(P4)/(P1)
R3	(0.24) APPROX.

$K_2(1770) \rightarrow (K f_2(1270))/(K \pi \pi)$	(P9)/(P1)
R4	(F2(1270) INTO PI PI)
R4	(0.16) APPROX.

$K_2(1770) \rightarrow (K \phi)/total$	(P10)
R5	SEEN



For notation, see key on page 91.

# Meson Full Listings

$K^*(1790)$ ,  $K(1830)$ ,  $K_4^*(2060)$

## $K^*(1790)$

$$I(J^P) = \frac{1}{2}(1^-)$$

OMITTED FROM SUMMARY TABLE

SEEN IN PARTIAL WAVE ANALYSIS OF THE KO PI+PI- OR K- PI+ SYSTEMS.

### $K^*(1790)$ MASS (MeV)

M	1650 APPROX.	ESTABROOK 78 ASPK	0 13K+-P,K+-PI+-N
M	1800. 70.	ETKIN 80 MPS	0 6 K-P,KO PI-PI-N
M	1700 APPROX.	ASTON 81 LASS	0 11 K-P,K- PI- N
M D	1786.0 9.0	ASTON 84 LASS	0 11K-P,KO ZPI N
M	1786.2 8.9	AVERAGE	
M D	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US		

### $K^*(1790)$ WIDTH (MeV)

W	250-300 APPROX.	ESTABROOK 78 ASPK	0 13K+-P,K+-PI+-N
W	370. 30.	ETKIN 80 MPS	0 6 K-P,KO PI-PI-N
W	(200) APPROX.	ASTON 81 LASS	0 11 K-P,K- PI+ N
W D	195.0 22.0	ASTON 84 LASS	0 11K-P,KO ZPI N
W	186.3 17.7	AVERAGE	
W D	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US		

### $K^*(1790)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$K^*(1790) \rightarrow K \pi$	494- 140
P2	$K^*(1790) \rightarrow K^*(892) \pi$	892- 140
P3	$K^*(1790) \rightarrow K \rho$	494- 769

### $K^*(1790)$ BRANCHING RATIOS

$K^*(1790) \rightarrow (K \rho)/(K^*(892) \pi)$	(P3)/(P2)
R1	3.4 1.3 ASTON 84 LASS 0 11 K-P,KO ZPI N
$K^*(1790) \rightarrow (K \pi)/(K^*(892) \pi)$	(P1)/(P2)
R2	2.8 1.1 ASTON 84 LASS 0 11 K-P,KO ZPI N
$K^*(1790) \rightarrow (K \rho)/(K \pi)$	(P3)/(P1)
R3	1.2 0.4 ASTON 84 LASS 0 11 K-P,KO ZPI N

### REFERENCES FOR $K^*(1790)$

ESTABROOK 78 NP B 133 490	ESTABROOKS, CARNEGIE, + (MONT+CARL+DURH+SLAC) JP
ETKIN 80 PR D 22 42	-FOLEY, LINDENBAUM, KRAMER, + (BNL+CUNY) JP
ASTON 81 PL 106 B 235	-CARNEGIE, DUNWOODIE, DURKIN+(SLAC+CARL+OTTA) JP
ASTON 84 PL 149 B 258	-CARNEGIE, DUNWOODIE, DURKIN+(SLAC+CARL+OTTA) JP

## $K(1830)$

$$I(J^P) = \frac{1}{2}(0^-)$$

OMITTED FROM SUMMARY TABLE

SEEN IN PARTIAL WAVE ANALYSIS OF K- PHI SYSTEM. NEEDS CONFIRMATION.

### $K(1830)$ MASS (MeV)

M	(1830.0) APPROX	ARMSTRONG 83 OMEG -	18.5 K-P, 3K P
W	(250.0) APPROX	ARMSTRONG 83 OMEG -	18.5 K-P, 3K P

### $K(1830)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$K(1830) \rightarrow K \phi$	494+1020

### REFERENCES FOR $K(1830)$

ARMSTRON 83 NP B 221 1	ARMSTRONG- (BARI+BIRM-CERN-MILA+LPNP+PAVI) JP
------------------------	---

## $K_4^*(2060)$ was $K^*(2060)$

$$I(J^P) = \frac{1}{2}(4^+)$$

### $K_4^*(2060)$ MASS (MeV)

M	488 2115. 46.	CARMONY 77 HBC	0 9 K-D,K- PIONS
M C	(2092.) (21.)	ASTON 1 81 LASS	011.K-P,K- PI- N
M D	2070. 100. 40.	ASTON 2 81 LASS	011.K-P,K- PI- N
M	650 2088. 20.	BAUBILLIE 82 HBC	- 8.25 K-P,KS PI-P
M W B	400 2039. 10.	CLELAND 82 SPEC	-- 50 K-P,KS PI--P
M	AVERAGE MEANINGLESS (SCALE FACTOR = 1.8)		
M B	FROM A FIT TO 8 MOMENTS.		
M W	NUMBER OF EVENTS EVALUATED BY US.		
M C	FROM A FIT TO Y(5,0), Y(7,0) AND Y(8,0) MOMENTS.		
M D	FROM ENERGY INDEPENDENT PWA.		

### $K_4^*(2060)$ WIDTH (MeV)

W	300. 200.	CARMONY 77 HBC	0 9 K-D,K- PIONS
W C	(205.) (70.) (55.)	ASTON 1 81 LASS	011.K-P,K- PI- N
W D	240. 500. 100.	ASTON 2 81 LASS	011.K-P,K- PI- N
W	650 170. 100. 50.	BAUBILLIE 82 HBC	- 8.25 K-P,KS PI-P
W B	400 189. 35.	CLELAND 82 SPEC	+- 50 K-P,KS PI--P
W	AVERAGE MEANINGLESS		
W B	FROM A FIT TO 8 MOMENTS.		
W W	NUMBER OF EVENTS EVALUATED BY US.		
W C	FROM A FIT TO Y(5,0), Y(7,0) AND Y(8,0) MOMENTS.		
W D	FROM ENERGY INDEPENDENT PWA.		

### $K_4^*(2060)$ PARTIAL DECAY MODES

		DECAY MASSES
P1	$K_4^*(2060) \rightarrow K \pi$	494+ 140
P2	$K_4^*(2060) \rightarrow K^*(892) \pi \pi$	892+ 140- 140
P3	$K_4^*(2060) \rightarrow \rho K \pi$	769+ 498- 140
P4	$K_4^*(2060) \rightarrow \omega K \pi$	783+ 498- 140
P5	$K_4^*(2060) \rightarrow K^*(892) \pi \pi \pi$	892+ 140- 140- 135

### $K_4^*(2060)$ BRANCHING RATIOS

$K_4^*(2060) \rightarrow (K \pi)/total$	(P1)
R1	0.07 0.01 ASTON 2 81 LASS 0 11 K-P,K- PI- N
$K_4^*(2060) \rightarrow (K^*(892) \pi \pi)/total$	(P2)
R2	SEEN BAUBILLIE 82 HBC - 8.25K-P,KS 3PI P
$K_4^*(2060) \rightarrow (\rho K \pi)/total$	(P3)
R3	SEEN BAUBILLIE 82 HBC - 8.25K-P,KS 3PI P
$K_4^*(2060) \rightarrow (\omega K \pi)/total$	(P4)
R4	SEEN BAUBILLIE 82 HBC - 8.25K-P,KS 3PI P
$K_4^*(2060) \rightarrow (K^*(892) \pi \pi \pi)/total$	(P5)
R5	POSSIBLY SEEN BAUBILLIE 82 HBC - 8.25K-P,KS 3PI P

### REFERENCES FOR $K_4^*(2060)$

CARMONY 71 PRL 27 1160	+CORDS, CLOPP, ERWIN, MEIERE, + (PURD+UCD-IND)
CARMONY 77 PR D 16 1251	+CLOPP, LANDER, MEIERE, YEN, + (PURD+UCD+IUPUI)
BROMBERG 80 PR D 22 1513	+HAGGERTY, AERAMS, DZIERBA(CIT-FNAL+ILL-IND)
CLELAND 80 PL 97B 465	+DORSAZ, MARTIN, NEF, + (PITT-GEVA-LAUS+DURH) JP
ASTON 1 81 PL 99 B 502	+DUNWOODIE, DURKIN, FIEGUTH+ (SLAC-CARL+OTTA) JP
ASTON 2 81 PL 106 B 235	+CARNEGIE, DUNWOODIE, DURKIN+(SLAC-CARL+OTTA) JP
BAUBILLIE 82 PL 118 B 447	BAUBILLIER, BURNS, (BIRM-CERN-GLAS+MSU+LPNP)
CLELAND 82 NP B 208 189	+DELFOSE, DORSAZ, GLOOR(DURH-GEVA-LAUS+PITT)



## Meson Full Listings

 $K_2(2250)$ ,  $K_3(2320)$ ,  $K_4(2500)$ ,  $D^\pm$ ,  $D^0$ ,  $D^*(2010)^\pm$  $K_2(2250)$   
was  $K(2250)$ 

$$I(J^P) = \frac{1}{2}(2^-)$$

OMITTED FROM  
SUMMARY TABLEFORMERLY CALLED  $K^*$ .  
THIS ENTRY CONTAINS VARIOUS PEAKS IN STRANGE MESON  
SYSTEMS REPORTED IN THE 2100-2300 MEV REGION AS  
WELL AS ENHANCEMENTS SEEN IN ANTIHYPERON NUCLEON  
SYSTEM, EITHER IN THE MASS SPECTRA OR IN THE JP=2-  
WAVE. $K_2(2250)$  MASS (MeV)

M	20(2240.)	(20.)	LISSAUER 70 HBC	9. K, P
M	(2200.)	APPROX.	SLATTERY 71 RVUE	8-13 K- P
M	37(2147.)	(4.)	CHLIAPNIK 79 HBC	+ K-P TO LAM-BAR P
M	Q 2235.	50.	BAUBILLIE 81 HBC	- 8. K-P, LAM PBAR
M	Q 2260.	20.	CLELAND 81 SPEC	-- 50 K+P, LAM PBAR
M	Q 2200.0	40.0	ARMSTRONG 83 OMEG	- 18 K-P, LAM PBAR
M	AVG	2246.5	16.8	AVERAGE

M C COMPILATION OF (ANTIHYPERON) MASS IN K+ P 8.-13. GEV/C  
M Q JP=2- FROM MOMENTS ANALYSIS. $K_2(2250)$  WIDTH (MeV)

W	20	(80.)	(20.)	LISSAUER 70 HBC	9. K, P
W	C	(200.)	APPROX.	SLATTERY 71 RVUE	8-13 K- P
W	Q	37	(40.)	APPROX.	CHLIAPNIK 79 HBC + K+P TO LAM-BAR P
W	Q	(200.)	APPROX.	BAUBILLIE 81 HBC	- 8. K-P, LAM PBAR
W	Q	210.	30.	CLELAND 81 SPEC	-- 50 K+P, LAM PBAR
W	Q	150.0	30.0	ARMSTRONG 83 OMEG	- 18 K-P, LAM PBAR
W	AVG	180.0	50.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.4)	

M C COMPILATION OF (ANTIHYPERON) MASS IN K, P 8.-13. GEV/C  
M Q JP=2- FROM MOMENTS ANALYSIS. $K_2(2250)$  PARTIAL DECAY MODES

		DECAY MASSES
P1	$K_2(2250) \rightarrow K \pi \pi$	498+ 135- 135
P2	$K_2(2250) \rightarrow \Delta \bar{p}$	1116+ 938

REFERENCES FOR  $K_2(2250)$ 

ALEXANDE 68 PRL 20 755	ALEXANDER, FIRESTONE, GOLDBABER, SHEN	(LRL)
LISSAUER 70 NP B 18 491	-ALEXANDER, FIRESTONE, GOLDBABER	(LBL)
SLATTERY 71 UR 875-332(PREP)	P. SLATTERY, A REVIEW OF STRANGE MESONS (ROCH)	
CHLIAPNI 79 NP B 158 253	CHLIAPNIKOV, GERDYUKOV	(CERN-BELG-MONS)
BAUBILLI 81 NP B 183 1	BAUBILLIER, +	(BIRM-CERN-GLAS-MSU-LPNP) JP
CLELAND 81 NP B 184 1	-NEF, MARTIN, +	(PITT+GEVA-LAUS-DURH) JP
ARMSTRON 83 NP B 227 365	ARMSTRONG+ (BARI-BIRM-CERN+MILA-LPNP+PAVI)	

 $K_3(2320)$   
was  $K(2320)$ 

$$I(J^P) = \frac{1}{2}(3^+)$$

OMITTED FROM  
SUMMARY TABLETHIS ENTRY CONTAINS ENHANCEMENTS SEEN IN THE  
JP=3- WAVE OF THE ANTIHYPERON NUCLEON SYSTEM $K_3(2320)$  MASS (MeV)

M	P	2320.0	30.0	CLELAND 81 SPEC	-- 50 K+P, LAM PBAR
M	P	2330.0	40.0	ARMSTRONG 83 OMEG	- 18 K-P, LAM PBAR
M	AVG	2323.6	24.0	AVERAGE	

M P JP=3+, FROM MOMENTS ANALYSIS

 $K_3(2320)$  WIDTH (MeV)

W	P	(250.0)	APPROX.	CLELAND 81 SPEC	-- 50 K-P, LAM PBAR
W	P	150.0	30.0	ARMSTRONG 83 OMEG	- 18 K-P, LAM PBAR
W	P	JP=3-	FROM MOMENTS ANALYSIS		

 $K_3(2320)$  PARTIAL DECAY MODES

		DECAY MASSES
P1	$K_3(2320) \rightarrow \Delta \bar{p}$	1116- 938

REFERENCES FOR  $K_3(2320)$ 

CLELAND 81 NP B 184 1	-NEF, MARTIN, +	(PITT+GEVA-LAUS-DURH)
ARMSTRON 83 NP B 227 365	ARMSTRONG-	(BARI-BIRM-CERN+MILA-LPNP+PAVI)

 $K_4(2500)$   
was  $K(2500)$ 

$$I(J^P) = \frac{1}{2}(4^-)$$

OMITTED FROM  
SUMMARY TABLETHIS ENTRY CONTAINS ENHANCEMENTS SEEN IN THE  
JP=4- WAVE OF THE ANTIHYPERON NUCLEON SYSTEM $K_4(2500)$  MASS (MeV)

M	R	2490.0	20.0	CLELAND 81 SPEC	-- 50 K+P, LAM PBAR
M	R	JP=4-	FROM MOMENTS ANALYSIS		

 $K_4(2500)$  WIDTH (MeV)

W	R	(250.0)	APPROX.	CLELAND 81 SPEC	-- 50 K+P, LAM PBAR
W	R	JP=4-	FROM MOMENTS ANALYSIS		

REFERENCES FOR  $K_4(2500)$ 

CLELAND 81 NP B 184 1	-NEF, MARTIN, -	(PITT+GEVA-LAUS+DURH)
-----------------------	-----------------	-----------------------

 $C = \pm 1$ ,  $B = 0$  MESON STATES $D^\pm$ 

$$I(J^P) = \frac{1}{2}(0^-)$$

SEE STABLE PARTICLE FULL LISTINGS

 $D^0$ 

$$I(J^P) = \frac{1}{2}(0^-)$$

SEE STABLE PARTICLE FULL LISTINGS

 $D^*(2010)^\pm$ 

$$I(J^P) = \frac{1}{2}(1^-)$$

 $D^*(2010)^\pm$  MASS (MeV)

M	G	(2008.)	(3.)	GOLDBABER 77 SMAG	-- E-E-
M	P	(2008.6)	(1.0)	PERUZZI 77 SMAG	-- E-E-
M	MASS	2010.1	0.7	FROM DO MASS (TRILLING 81 RVUE) AND MASS DIFFERENCE BELOW	

M G FROM SIMULTANEOUS FIT TO  $D^*(2010)^\pm$ ,  $D^*(2010)_0$ ,  $D^+$ , AND  $D_0$ , NOT  
M G INDEPENDENT OF FELDMAN 77 MASS DIFFERENCE BELOW.  
M P PERUZZI 77 MASS NOT INDEPENDENT OF FELDMAN 77 MASS DIFFERENCE  
M P BELOW AND PERUZZI 77 DO MASS VALUE. $D^*(2010)^\pm - D^0$  MASS DIFFERENCE (MeV)

DM	30	145.3	0.5	FELDMAN 77 SMAG	$D^+ \rightarrow D_0 \pi^-$
DM	2	145.2	0.6	BLIETSCHA 79 BEBC	NEUTRINO P
DM	(145.5)	APPROX.	AVERY 80 SPEC	GAMMA A	
DM	60	145.5	0.3	FITCH 81 SPEC	$\pi^- \rightarrow \pi^0$
DM	14	145.5	0.5	YELTON 82 SMK2	$29 E \rightarrow E, K \rightarrow \pi^+$
DM	16	145.8	1.5	AHLEN 83 HRS	$D^+ \rightarrow D_0 \pi^+$
DM	12	145.1	1.8	BAILEY 83 SPEC	$D^+ \rightarrow D_0 \pi^+$
DM	28	145.5	0.3	BAILEY 83 SPEC	$D^+ \rightarrow D_0 \pi^+$
DM	14	145.1	0.5	BAILEY 83 SPEC	$D^+ \rightarrow D_0 \pi^+$
DM	D	145.46	0.08	ALBRECHT 85 ARG	$D^+ \rightarrow D_0 \pi^+$
DM	AVG	145.451	0.072	AVERAGE	

DM D SYSTEMATIC ERROR ADDED QUADRATICALLY BY US

For notation, see key on page 91.

# Meson Full Listings

$D^*(2010)^{\pm}, D^*(2010)^0, D^*(2420)^0$

**$D^*(2010)^{\pm} - D^*(2010)^0$  MASS DIFFERENCE (MeV)**

EM	P	2.6	1.8	PERUZZI	77	SMAG	--	E+E-
EM	P	NOT INDEPENDENT OF FELDMAN 77 MASS DIFFERENCE ABOVE, PERUZZI 77						
EM	P	DO MASS, AND GOLDBER 77 $D^*(2010)^0$ MASS.						
EM								
EM	DMASS	2.9	1.3	FROM $(D^{*+})-(D^0)$ AND $(D^{*0})-(D^0)$				
EM				MASS DIFFERENCES				

**$D^*(2010)^{\pm}$  WIDTH (MeV)**

W	30	(2.0)	OR LESS	CL=.90	FELDMAN	77	SMAG	D*+ TO D0 P1+
W		(2.2)	OR LESS		YELTON	82	SMK2	29 E+E-,K-PI+P1-

**$D^*(2010)^{\pm}$  PARTIAL DECAY MODES**

		DECAY MASSES	
P1	$D^*(2010)^+ \rightarrow D^0 \pi^+$	1865-	140
P2	$D^*(2010)^+ \rightarrow D^+ \gamma$	1869-	0
P3	$D^*(2010)^+ \rightarrow D^+ \pi^0$	1869+	135
P	$D^*(2010)^0$ MODES ARE CHARGE CONJUGATES OF ABOVE MODES		

**FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS**

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions,  $P_i$ , as follows: The diagonal elements are  $P_i = \delta P_i$ , where  $\delta P_i = \sqrt{(\delta P_i)^2 / P_i^2}$ , while the off-diagonal elements are the normalized correlation coefficients  $(\delta P_i \delta P_j) / (P_i P_j)$ . For the definitions of the individual  $P_i$ , see the listings above; only those  $P_i$  appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3
P 1	.4892+-0832		
P 2	-.7652	-.1708+-1087	
P 3	.0000	-.6438	.3400+-0700

**$D^*(2010)^{\pm}$  BRANCHING RATIOS**

<b><math>D^*(2010)^+ \rightarrow (D^0 \pi^+)/total</math></b>		(P1)
R1 G	0.6 0.15	GOLDBER 77 SMAG + E+E-
R1 R	0.44 0.10	COLES 82 SMK2 E+ E-
R1 G	ASSUMING THAT ISOSPIN IS CONSERVED IN THE DECAY	
R1 R		
R1 AVG	0.489 0.083	AVERAGE
R1 FIT	0.489 0.083	FROM FIT
<b><math>D^*(2010)^+ \rightarrow (D^+ \gamma)/total</math></b>		(P2)
R2 C	(0.22) (0.12)	COLES 82 SMK2 E+ E-
R2 C	NOT INDEPENDENT OF R1 AND R3 MEASUREMENT.	
R2 R		
R2 FIT	0.17 0.11	FROM FIT
<b><math>D^*(2010)^+ \rightarrow (D^+ \pi^0)/total</math></b>		(P3)
R3 R	0.34 0.07	COLES 82 SMK2 E+ E-
R3 R		
R3 FIT	0.340 0.070	FROM FIT

**REFERENCES FOR  $D^*(2010)^{\pm}$**

PERUZZI 76 PRL 37 569	+PICCOLO,FELDMAN,NGUYEN,WISS,+ (SLAC-LBL)
FELDMAN 77 PRL 38 1313	+PERUZZI,PICCOLO,ABRAMS,ALAM- (SLAC+LBL)
PERUZZI 77 PRL 39 1301	+PICCOLO,FELDMAN,PERL,+(SLAC,LBL,NWES-HAWA)
GOLDBER 77 PL 69 B 503	+WISS,ABRAMS,ALAM,BOVARSKI,+ (LBL+SLAC)
BLIETSCH 79 PL 86 B 108	BLIETSCHAU,- (AACH+BDNN-CERN-MPIM+OXF)
AVERY 80 PRL 44 1309	+WISS,BINKLEY,ATIYA,+ (ILL+FNAL+COLU)
FITCH 81 PRL 46 761	+DEVAUX,CAVAGLIA,MAY,+ (PRIN+SACL+TORI+BNL)
TRILLING 81 PRPL 75 57	G.H. TRILLING (LBL+UCB)
BEBEK 82 PRL 49 610	+ (HARV+OSU+ROCH+RUTG+SYRA+VAND+CORN+ITHACA)
COLES 82 PR D26 2190	+ABRAMS,BLOCKER,BLONDEL+ (LBL+SLAC)
YELTON 82 PRL 49 430	+FELDMAN,GOLDBER,+ (SLAC-LBL+UCB+HARV)
AHLEN 83 PRL 51 1147	+AKERLOF+ (ANL+IND+LBL+MICH+PURD+SLAC)
ALTHOFF 83 PL 126 B 493	+FISCHER,BURKHARDT+ (TASSO COLLABORATION)
BATILEY 83 PL 132 B 230	+BARDSLEY+ (AMST+BRIS-CERN+CRAC+MPIM-RHEL)
ALBRECHT 85 PL 150 B 235	+DRESCHER,HELLER+ (ARGUS COLLABORATION)

**$D^*(2010)^0$**

$$I(J^P) = \frac{1}{2}(1^-)$$

J CONSISTENT WITH 1, VALUE 0 RULED OUT (NGUYEN 77).

**$D^*(2010)^0$  MASS (MeV)**

M	G	(2006.)	(1.5)	GOLDBER 77 SMAG	E+E-
M	G	FROM SIMULTANEOUS FIT TO $D^*(2010)^{\pm}, D^*(2010)^0, D^+$ , AND $D^0$ .			
M					
M	MASS	2007.2	2.1	FROM DO MASS (TRILLING 81 RVUE) AND MASS DIFFERENCE BELOW	

**$D^*(2010)^0 - D^0$  MASS DIFFERENCE (MeV)**

DM	G	142.7	1.7	GOLDBER 77 SMAG	0 E+E-
DM	G	142.2	2.0	SADROZIN 80 CBAL	0 D*0 TO D0 P10
DM	G	FROM SIMULTANEOUS FIT TO $D^*(2010)^{\pm}, D^*(2010)^0, D^+$ , AND $D^0$ .			
DM					
DM	AVG	142.5	1.3	AVERAGE	

**$D^*(2010)^0$  WIDTH (MeV)**

W	(5.)	OR LESS	GOLDBER 76 SMAG	E+E- TO D*0+
---	------	---------	-----------------	--------------

**$D^*(2010)^0$  PARTIAL DECAY MODES**

		DECAY MASSES	
P1	$D^*(2010)^0 \rightarrow D^0 \pi^0$	1865+	135
P2	$D^*(2010)^0 \rightarrow D^0 \gamma$	1865+	0
P	$D^*(2010)^0$ MODES ARE CHARGE CONJUGATES OF ABOVE MODES		

**$D^*(2010)^0$  BRANCHING RATIOS**

<b><math>D^*(2010)^0 \rightarrow (D^0 \gamma)/(D^0 \pi^0 + D^0 \gamma)</math></b>		(P2)/(P1-P2)
R1 G	0.45 0.15	GOLDBER 77 SMAG E+F-
R1 R	0.47 0.12	COLES 82 SMK2 E+ E-
R1 R	0.53 0.13	BARTEL 85 JADE E+ E-,HADRONS
R1 G	WE QUOTE THE NORMAL FIT VALUE FROM TABLE 1. THE ISO-SPIN	
R1 G	CONSTRAINED FIT IS NOW KNOWN TO GIVE A DO GAMMA FRACTION WHICH IS	
R1 G	TOO LARGE. SEE DETAILS IN FOOTNOTE 21 OF FELDMAN 77 REVIEW.	
R1 R		
R1 AVG	0.485 0.076	AVERAGE

**REFERENCES FOR  $D^*(2010)^0$**

GOLDBER 76 PRL 37 255	GOLDBER,PIERRE,ABRAMS,ALAM,- (LBL+SLAC)
GOLDBER 76 SLAC CONF. 379	G.GOLDBER (AVAIL. AS LBL-5534) (LBL+SLAC)
GOLDBER 77 PL 69 B 503	GOLDBER,ABRAMS,ALAM- (LBL-SLAC)
FELDMAN 77 BANFF SUM. INST 75	G.J.FELDMAN (SLAC)
NGUYEN 77 PRL 39 262	+WISS,ABRAMS,ALAM,BOVARSKI,- (LBL+SLAC)
SADROZIN 80 MADISON CONF. 681	SADROZINSKI,+ (PRIN-CIT.HARV+SLAC+STAN)
TRILLING 81 PRPL 75 57	G.H. TRILLING (LBL+UCB)
COLES 82 PR D26 2190	+ABRAMS,BLOCKER,BLONDEL+ (LBL+SLAC)
BARTEL 85 PL 1618 197	+DIETRICH,AMBRUS+ (JADE COLLABORATION)

**$D^*(2420)^0$**

$$I(J^P) = \frac{1}{2}(?)$$

OMITTED FROM SUMMARY TABLE

SEEN IN  $D^*(2010)^0$  P1- JP-0+ RULED OUT

**$D^*(2420)^0$  MASS (MeV)**

M	2420.	6.	ALBRECHT 85 ARG	E+E-,D*+ PI-X
---	-------	----	-----------------	---------------

**$D^*(2420)^0$  WIDTH (MeV)**

W	70.	21.	ALBRECHT 85 ARG	E+E-,D*+ PI-X
---	-----	-----	-----------------	---------------

**$D^*(2420)^0$  PARTIAL DECAY MODES**

		DECAY MASSES	
P1	$D^*(2420)^0 \rightarrow D^*(2010)^0 \pi^-$	2010+	140

# Meson Full Listings

$D^*(2420)^0$ ,  $D_s^\pm$ ,  $D_s^*(2110)$ ,  $B^\pm$ ,  $B^0$ ,  $B^*(5325)$ , EXOTICS

## $D^*(2420)^0$ BRANCHING RATIOS

$D^*(2420)^0 \rightarrow (D^*(2010)^+ \pi^-)/\text{total}$  (P1)  
 R1 SEEN ALBRECHT 85 ARG E+E-,D\*, P1- X

## REFERENCES FOR $D^*(2420)^0$

ALBRECHT 85 DESY 85-119 (ARGUS COLLABORATION)

$D_s^\pm$   
was  $F^\pm$

$I(J^P) = 0(0^-)$

SEE STABLE PARTICLE FULL LISTINGS

$D_s^*(2110)$   
was  $F^*(2140)$

$I(J^P) = ?(?)$

OMITTED FROM SUMMARY TABLE

## $D_s^*(2110)$ MASS (MeV)

M D 2113. 8.  
 M D OUR EVALUATION FROM D/S MASS= 1972 MEV AND DM AVERAGE BELOW

## $D_s^*(2110)^+$ - $D_s^+$ MASS DIFFERENCE (MeV)

DM	110.	46.	BRANDELIK 79 DASP	+ E+E-,D/S GAMMA
DM D 60	139.5	12.8	AIHARA 84 TPC	+ -E+E-,HADRONS
DM D	144.0	11.4	ALBRECHT 84 ARG	E+E-,D/S GAMMA
DM B	143.0	18.0	ASRATTAN 85 HLSC	FNAL 15FT,NU-H2
DM AVG	141.3	7.6	AVERAGE	
M D	SYSTEMATIC ERROR ADDED QUADRATICALLY BY US			

## $D_s^*(2110)$ PARTIAL DECAY MODES

P1  $D_s^*(2110) \rightarrow D_s \gamma$  DECAY MASSES 1971, 0

## $D_s^*(2110)$ BRANCHING RATIOS

$D_s^*(2110) \rightarrow (D_s \gamma)/\text{total}$  (P1)  
 R1 SEEN BY ALL FOUR GROUPS QUOTED ABOVE FOR DM

## REFERENCES FOR $D_s^*(2110)$

BRANDELI 77 PL 70 B 132	BRANDELIK,CORDS,+ (AACH+DESY+HAMB+MPIM-TOKY)
BRANDELI 78 PL 76 B 361	BRANDELIK,CORDS,+ (AACH+DESY+HAMB+MPIM-TOKY)
BRANDELI 79 PL 80 B 412	BRANDELIK,CORDS,+ (AACH+DESY+HAMB+MPIM-TOKY)
AIHARA 84 PRL 53 2455	+ALSTON- (LBL+UCI+UCR+JHU+MASA+TOKY+YALE)
ALBRECHT 84 PL 146 B 111	+ (DESY+DORI+HEID+IPPC-KANS+LUND+ITEP-USCC)
ASRATTAN 85 PL 156 B 441	+FFDOTOV,AMMOISOV,BURTOVOY- (ITEP+SERP)

## $B = \pm 1$ MESON STATES

$B^\pm$

$I(J^P) = ?(?)$

SEE STABLE PARTICLE FULL LISTINGS

$B^0$

$I(J^P) = ?(?)$

SEE STABLE PARTICLE FULL LISTINGS

$B^*(5325)$

$I(J^P) = ?(?)$

OMITTED FROM SUMMARY TABLE

## $B^*(5325)$ MASS (MeV)

M D 1400 5325.0 5.0 HAN 85 CUSB 0 E+E-,GAMMA E X  
 M D FROM B MASS 5272.3--2.5 MEV AND MASS DIFFERENCE BELOW

## $B^*(5325)$ - B MASS DIFFERENCE (MeV)

DM D 1400 52.0 4.5 HAN 85 CUSB 0 E+E-,GAMMA E X  
 M D SYSTEMATIC ERROR ADDED QUADRATICALLY BY US

## REFERENCES FOR $B^*(5325)$

HAN 85 PRL 55 36 +KLOPFENSTEIN,MAGERAS+ (COLU+LSU+MPIM-STON)

## EXOTIC MESON STATES

EXOTICS

OMITTED FROM SUMMARY TABLE

THE PURPOSE OF THIS ENTRY IS TO PROVIDE A LIST OF REFERENCES FOR EXOTIC MESON SEARCHES (SEE THE SECTION ON THE NONRELATIVISTIC QUARK MODEL IN THE MISCELLANEOUS SECTION OF THIS REVIEW), AS WELL AS THEORETICALLY BASED SUGGESTIONS FOR EXPERIMENTS. NOTE THAT LIPKIN 73 PROPOSES EXPERIMENTS WHICH ARE CONCLUSIVE EVEN IF NEGATIVE RESULTS ARE OBTAINED.

## REFERENCES FOR EXOTICS

ROSENFEL 68 PHILA.CONF.P.455	A.H.ROSENFELD (LRL)
DODD 69 PR 177 1991	+JOLDERSM4, PALMER, SAMIOS (BNL)
CHO 70 PL 32 B 409	+DERRICK,JOHNSON,MUSGRAVE,+ (ANL+NWES-KANS)
GIACOMEL 70 PL 33 B 373	G.GIACOMELLI + (BGNA-SACL+AMST+REHO-EPOL)
LYS 70 PR D 2 2525	J.LYS+ (MICH)
ROSNER 70 EXP.MESON SPECTROSCOPY,ED. C.BALTAY AND A.H.ROSENFELD,P.499	
BUHL 72 NP B 37 421	+CLINE,TERRELL (WISCONSIN)
COHEN 73 NP B 53 1	+FERBEL,SLATTERY,WERNER (ROCHESTER)
DURUSOY 73 PL 45 B 517	-BAUBILLIER,GEORGE,ARMENISE,+ (LPN+BAR1)
ALAM 74 PL 53B 207	+BRABSON,GALLOWAY,+ (IND+PURD+SLAC+VAND)
COHEN 74 BOSTON	D.COHEN REVIEW TALK (COLU)
OREN 74 NP B71 189	+COOPER,FIELDS,RHINES,WHITMORE,+ (ANL+OXF)
BALTAY 75 PL 57B 293	+CAUTIS,COHEN,KALELKAR,PISELLO,-(COLU-BING)
DAVIS 75 NP B96 426	+AMMAR,KROPAC,YARGER,+ (KANS+CCAC-ANL)
BRUNDIR 76 PL 64 B 107	BRUNDIERS,BRUN,FLURI,+ (FREIBURG+SACL+ETH)
BOUCROT 77 NP B 121 251	+NAVACH,RIVET,+ (LALO-CERN+CDEF-EPOL)
HOOGLAND 77 NP B 126 109	+GRAYER,HYAMS,BLUM,DITL,+ (AMST+CERN+MPIM)
HOOGLAND 77 NP B 126 109	+GRAYER,HYAMS,BLUM,DITL,+ (AMST+CERN+MPIM)
MOSER 77 NP B 129 28	F.L.MOSER (EFI)
ALAM 78 PRL 40 1685	+BAGGETT,BAGLIN,BONAMY+(IND+PURD-SLAC-VAND)
ARMSTRON 78 PL 77 B 447	ARMSTRONG,FRAME,HUGHES,BIENLEIN+(GLAS+DESY)
LEMOIGNE 79 BATAVIA CONF.524	+ABOLINS,BARATE- (SACL+LOIC+SHP+IND)
KOOIJMAN 80 PRL 45 316	+ARENTON,AYRES,DIEBOLD,MAY- (ANL+EFI)
AGUILAR 81 ZPHY C 6 109	-ALBAJAR,SJOGREN,+ (CERN+CDEF+MADR-STOH)
APEL 81 NP B 193 269	-AUGENSTEIN,BERTOLUCCI,DONSKOV,+ (SERP-CERN)
BIONTA 81 PRL 46 970	-CARROLL,EDELSTEIN,+ (BNL+CERN+FNAL-SNAS)
EVANGELI 81 NP B 178 197	EVANGELISTA+(BAR+BOON+CERN+DARE+LVP+MILA)
FRAME 81 PL 107 B 301	+HUGHES,COLLEY,ARMSTRONG,+ (GLAS+BIRM+CERN)
IRVING 81 NP B 193 1	-LOVERRE,AGUILAR,+ (CERN+CDEF+MADR+STOH)
DOVER 84 PL 146 B 103	C.B.DOVER (ORSA)
JENKINS 84 PR D 30 1409	-DIAMOND,KIRSCH,+ (FSU+BRAN-BNL-CINC+SMAS)
KITAZOE 84 ZPHY C 24 143	+WADA,KABURAGI,KAWAGUCHI,MORII,+ (KOBE-MIT)

SUGGESTIONS FOR SEARCHES  
 ROSNER 68 PRL 21 950,1468 J.L.ROSNER (TEL-AVIV)

ROSNER 70 EXP.MESON SPECTROSCOPY,ED. C.BALTAY AND A.H.ROSENFELD,P.499	
FAIMAN 73 PL 43 B 307	D.FAIMAN,G.GOLDHABER,Y.ZARMI (CERN)
LIPKIN 73 PR D 7 2262	H.J.LIPKIN (ARGONNE-FNAL)
HOLMGREN 78 PL 77 B 304	+PENNINGTON (STDN-CERN)
ARENTON 82 PR D 25 224	+AYRES,DIEBOLD,MAY,SWALLOW- (ANL+ILL)
ACHASOV 85 ZPHY C 27 99	+DEVYANIN,SHESTAKOV (NOVO)

For notation, see key on page 91.

## Baryon Full Listings

$N$ 's and  $\Delta$ 's

### NOTE ON $N$ AND $\Delta$ RESONANCES

#### I. Introduction

(by G. Höhler, University of Karlsruhe)

The excited states of the nucleon have been studied in a large number of formation and production experiments. Production experiments are not suitable for accurate determination of resonance parameters but will be essential in searching for the many predicted nucleon resonances that decouple from the  $\pi N$  channel.<sup>1</sup>

The masses, widths, and elasticities of the  $N$  and  $\Delta$  resonances in the Baryon Summary Table come almost entirely from partial-wave analyses of  $\pi N$  total, elastic, and charge-exchange scattering data (see Sec. II, below). Similar methods have been used to get the  $N\eta$ ,  $\Delta K$ , and  $\Sigma K$  branching fractions. Other branching fractions come from isobar-model analyses of  $\pi N \rightarrow N\pi\pi$  data (Sec. III). Finally, some  $N\gamma$  branching fractions have been determined from photoproduction experiments (Sec. IV).

Table 1 lists all the  $N$  and  $\Delta$  entries in the Baryon Listings and gives our evaluation of the status of each, both overall and channel by channel. Only the established resonances (overall status 3 or 4 stars) appear in the Baryon Summary Table. A resonance is considered to be well established only if it has been seen in at least two independent analyses and if its partial wave does not behave erratically or have large errors. Good reason for a cautious attitude is the fact that some recent data<sup>2,3</sup> differ appreciably from earlier data and from predictions of the analyses.

The Baryon Listings give, in addition to the usual Breit-Wigner parameters, the locations and the residues of the poles of the resonant partial waves on the second sheet of the complex energy plane as obtained from  $\pi N$  partial-wave analyses and from the isobar-model analyses of  $\pi N \rightarrow N\pi\pi$ .

The Listings are much shortened by the omission of many now-obsolete results, nearly all which were published before 1975. There also used to be separate entries for bumps seen in production experiments — bumps with masses in the 1440-MeV region, the 1520-MeV region, etc. — but these have been removed. All the omitted material may be found in our 1982 edition.<sup>4</sup>

There are two recent extensive reviews of nucleon resonances.<sup>5,6</sup>

Further progress in understanding the  $N$  and  $\Delta$  resonances depends on investigations of three different types.

Table 1. The status of the  $N$  and  $\Delta$  resonances. Only those with an overall status of \*\*\* or \*\*\*\* are included in the main Baryon Summary Table.

Particle	$L_{2J-2J}$	Overall status	Status as seen in —						
			$N\pi$	$N\eta$	$\Delta K$	$\Sigma K$	$\Delta\pi$	$N\rho$	$N\gamma$
$N(939)$	$P_{11}$	****							
$N(1440)$	$P_{11}$	****	****	*			***	*	***
$N(1520)$	$D_{13}$	****	****	*			****	****	****
$N(1535)$	$S_{11}$	****	****	****			*	**	***
$N(1540)$	$P_{11}$	*					*	*	*
$N(1540)$	$P_{13}$	****	****	*	***	**	***	*	***
$N(1650)$	$D_{11}$	****	****	*	*		****	*	***
$N(1675)$	$D_{15}$	****	****	*	*		****	****	****
$N(1680)$	$F_{15}$	****	****	*	*		****	****	****
$N(1700)$	$D_{15}$	***	***	*	**	*	**	*	**
$N(1710)$	$P_{11}$	****	****	*	**	*	**	*	***
$N(1720)$	$P_{13}$	****	****	*	**	*	*	*	*
$N(1960)$	?	*							
$N(1990)$	$F_{17}$	**	**	*	*	*			*
$N(2000)$	$F_{15}$	**	**	*	*	*			*
$N(2080)$	$D_{13}$	**	**	*	*	*			*
$N(2090)$	$S_{11}$	*	*						
$N(2100)$	$P_{11}$	*	*						
$N(2190)$	$G_{17}$	****	****	*	*	*			*
$N(2200)$	$D_{15}$	**	**	*	*	*			*
$N(2220)$	$H_{19}$	****	****	*					*
$N(2250)$	$G_{19}$	****	****	*					*
$N(2600)$	$I_{11}$	***	***						*
$N(2700)$	$K_{113}$	**	**						*
$N(\sim 3000)$									
$\Delta(1232)$	$P_{33}$	****	****	F					****
$\Delta(1550)$	$P_{31}$	*		o			*	*	*
$\Delta(1600)$	$P_{31}$	**	**	r			**	*	**
$\Delta(1620)$	$S_{31}$	****	****	b			****	****	***
$\Delta(1700)$	$D_{33}$	****	****	i		*	***	**	***
$\Delta(1900)$	$S_{31}$	****	****	d		*	**	*	*
$\Delta(1905)$	$F_{35}$	****	****	d		*	**	*	***
$\Delta(1910)$	$P_{31}$	****	****	e		*	*	*	*
$\Delta(1920)$	$P_{33}$	***	***	n		*	*	*	*
$\Delta(1930)$	$D_{35}$	***	***	F		*			*
$\Delta(1940)$	$D_{33}$	*	*	o		*			*
$\Delta(1950)$	$F_{37}$	****	****	r		*	***	*	***
$\Delta(2000)$	$F_{35}$	**	**	b			**		**
$\Delta(2150)$	$S_{31}$	*	*	i					*
$\Delta(2200)$	$G_{37}$	*	*	d					*
$\Delta(2300)$	$H_{39}$	**	**	e					*
$\Delta(2350)$	$D_{35}$	*	*	n					*
$\Delta(2390)$	$F_{37}$	*	*						*
$\Delta(2400)$	$G_{37}$	**	**						*
$\Delta(2420)$	$H_{39}$	****	****						*
$\Delta(2750)$	$I_{313}$	**	**						*
$\Delta(2950)$	$K_{315}$	**	**						*
$\Delta(\sim 3000)$									

\*\*\*\* Good, clear, and unmistakable.  
 \*\*\* Good, but in need of clarification or not absolutely certain.  
 \*\* Not established; needs confirmation.  
 \* Evidence weak; could disappear.

(1) *New accurate data*: Much new data is coming from groups working at LAMPF,<sup>7</sup> and there is also some new accurate data from the Leningrad group.<sup>8</sup> These groups are also preparing to measure spin-rotation data, the first such in the resonance region. Very unfortunately, however, none of this work extends above a mass of about 1500 MeV, and to our knowledge there are no plans anywhere for new measurements at higher masses. The recently published results of an older high-statistics measurement of  $\pi^+p$  backward differential cross sections from 1.3 to 2.5 GeV/ $c^3$  disagree significantly with previous high-statistics experiments and thus also with predictions from partial-wave analyses.<sup>5</sup>

## Baryon Full Listings

$N$ 's and  $\Delta$ 's

(2) *New partial-wave analyses*: Existing solutions will need to be adjusted to get a good fit to the new data.<sup>9</sup> However, this is only the first step of an iterative procedure. A unique and *reliable* result can be obtained only if, in addition, rather strong analyticity constraints are imposed on the amplitudes. The CMU-LBL and Karlsruhe groups have computer programs for this purpose, but new fits has not yet been made due to lack of manpower.

(3) *New theoretical investigations*: The relation between the resonance parameters determined from partial-wave analyses and the quantities derived from lattice calculations and various models needs clarifying. Until this difficult problem is solved, the uncertainties involved in comparing theoretical predictions with the parameters listed in our Table should be kept much in mind.

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## II. Two-body partial-wave analyses and determination of resonance parameters

(by G. Höhler, University of Karlsruhe)

*$\pi N$  partial-wave analysis*: Even if  $\pi N \rightarrow \pi N$  scattering data were measured with infinite accuracy, it would not be possible in the inelastic region to determine a unique set of partial waves from the data alone. It is essential to add theoretical constraints, and unitarity, analyticity, and isospin invariance are chosen in

order to avoid the biases that a specific model or parametrization might introduce.

Atkinson et al.,<sup>1</sup> continuing earlier work, investigated to what extent amplitudes are restricted by just unitarity if the  $d\sigma/d\Omega$  and  $P$  angular distributions for  $\pi^+ p$  elastic scattering are given at one energy with very high precision. They found a variety of solutions differing from one another substantially in some of the lower partial waves and strongly in the tail of high partial waves. They concluded that cutting off the partial-wave expansion sharply (which was done in many early and in some recent analyses<sup>2,3</sup>) is not justified.

In QCD, isospin is not exactly conserved in strong interactions because the masses of the up and down quarks are different. The only well-established experimental evidence for a violation is in the  $\Delta(1232)$  region, where one is expected because the  $\Delta^{++}$  and  $\Delta^0$  masses are different. Other reported violations turned out to be caused by errors in the data or the analysis.<sup>4</sup>

The uniqueness problem remains even if one includes data for all three reactions plus unitarity and isospin invariance; it is still necessary to add analyticity constraints. Many analyses have used as input predictions for the forward amplitudes, which follow from total-cross-section data, the optical theorem, and forward dispersion relations, but this is still not nearly enough.

Constraints based on Mandelstam's 2-variable analyticity have so far been used successfully only in the CMU-LBL<sup>5</sup> and Karlsruhe-Helsinki<sup>6,7</sup> analyses. In both, long tails of high partial waves were admitted, but only some global results for these waves should be taken seriously, not the value of a particular high wave. The resonance masses, widths, and elasticities in the Baryon Summary Table are mainly determined by these two analyses, whose partial-wave amplitudes are shown in Fig. 1. More detailed figures and speed plots may be found in Ref. 8.

Results from other recent analyses should be considered preliminary as long as the compatibility with analyticity constraints and the effect of the neglect of higher partial waves have not been investigated (see Sec. 2.1 in Ref. 7).

Substantial progress may be expected when final results of several experiments now in progress are available.<sup>9</sup> The analysis will be simplified and improved if predictions for the tail of high partial waves, based on new evaluations of the nearby parts of the Mandelstam double spectral function,<sup>10</sup> are used. R. Koch<sup>11</sup> included these predictions in getting a smooth interpola-

For notation, see key on page 91.

## Baryon Full Listings

N's and  $\Delta$ 's

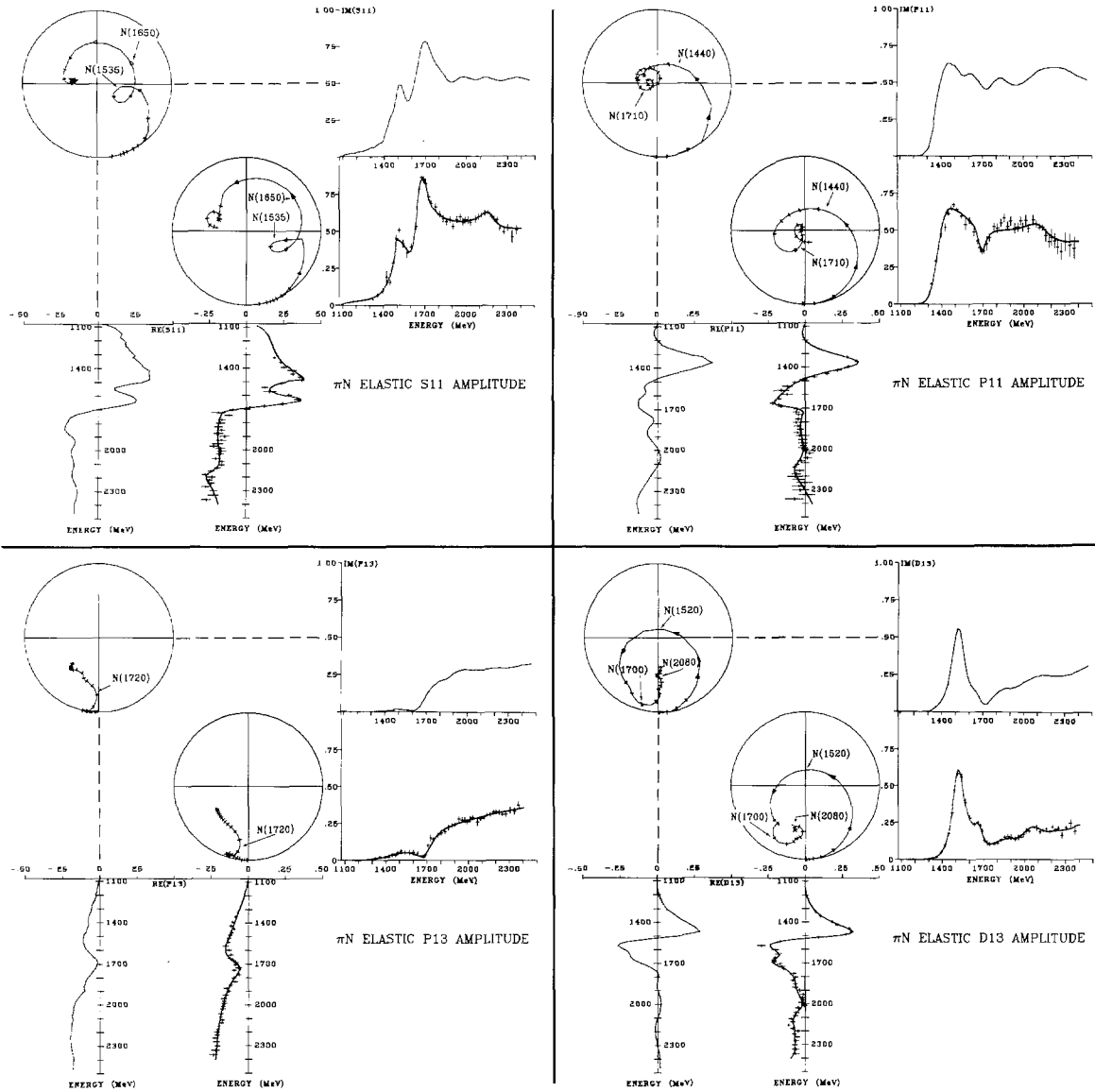


Fig. 1(a). The  $L_{2J,2J} = S_{11}, P_{11}, P_{13},$  and  $D_{13}$  partial-wave amplitudes for  $\pi N$  elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

# Baryon Full Listings

$N$ 's and  $\Delta$ 's

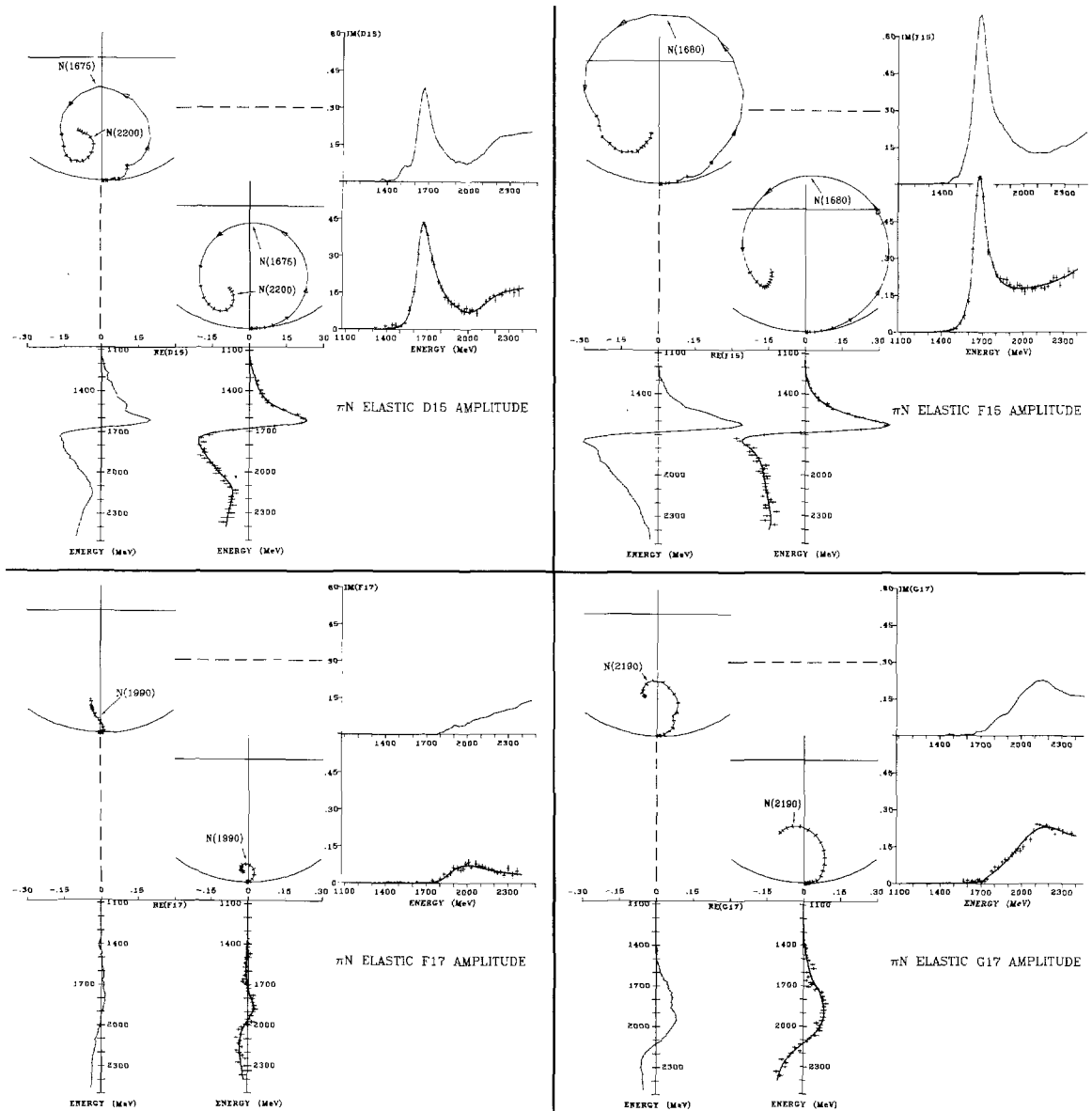


Fig. 1(b). The  $L_{2J,2J} = D_{15}, F_{15}, F_{17},$  and  $G_{17}$  partial-wave amplitudes for  $\pi N$  elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

For notation, see key on page 91.

## Baryon Full Listings

*N*'s and  $\Delta$ 's

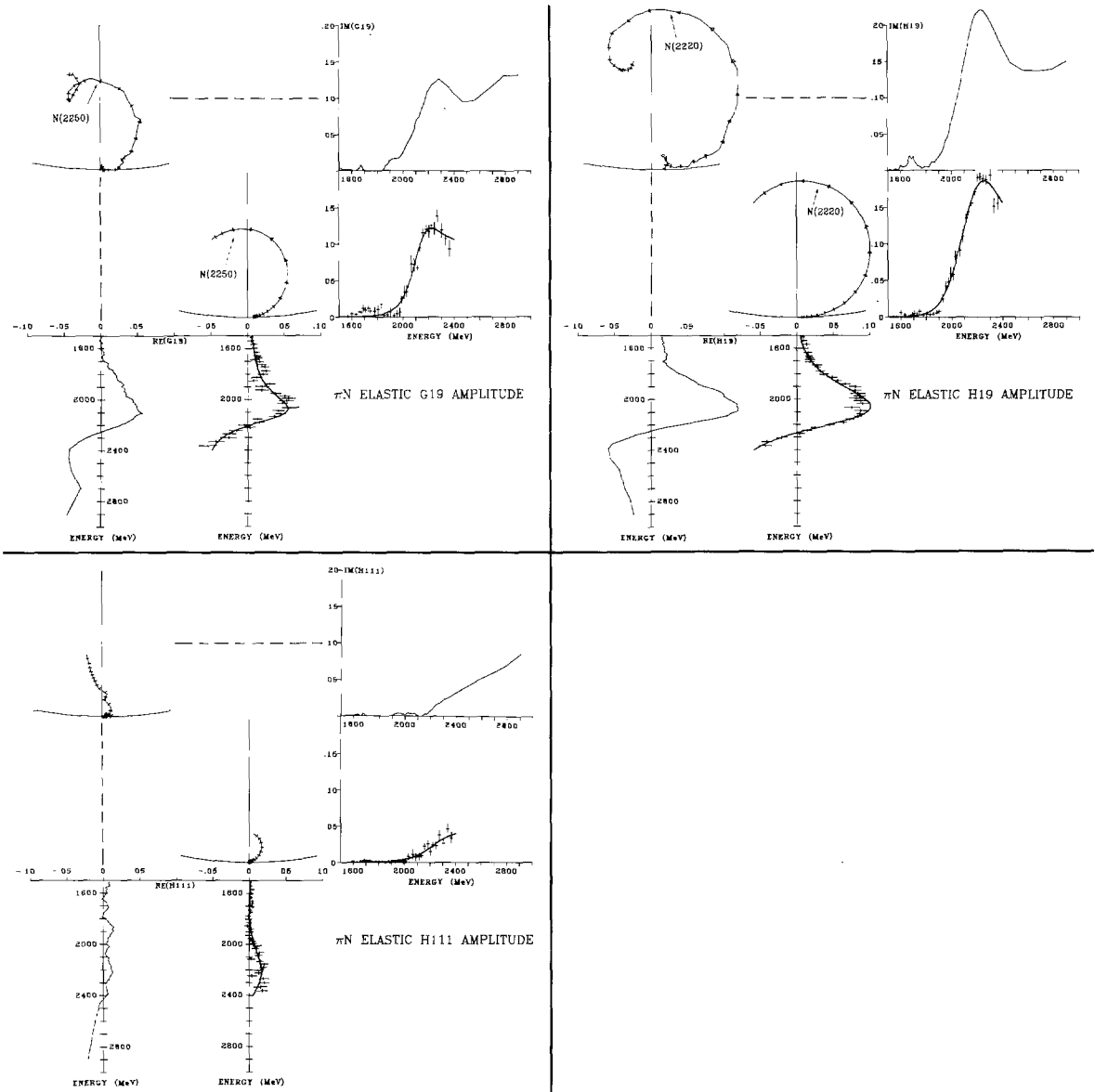


Fig. 1(c). The  $L_{2J,2J} = G_{19}$ ,  $H_{19}$ , and  $H_{111}$  partial-wave amplitudes for  $\pi N$  elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).



# Baryon Full Listings

$N$ 's and  $\Delta$ 's

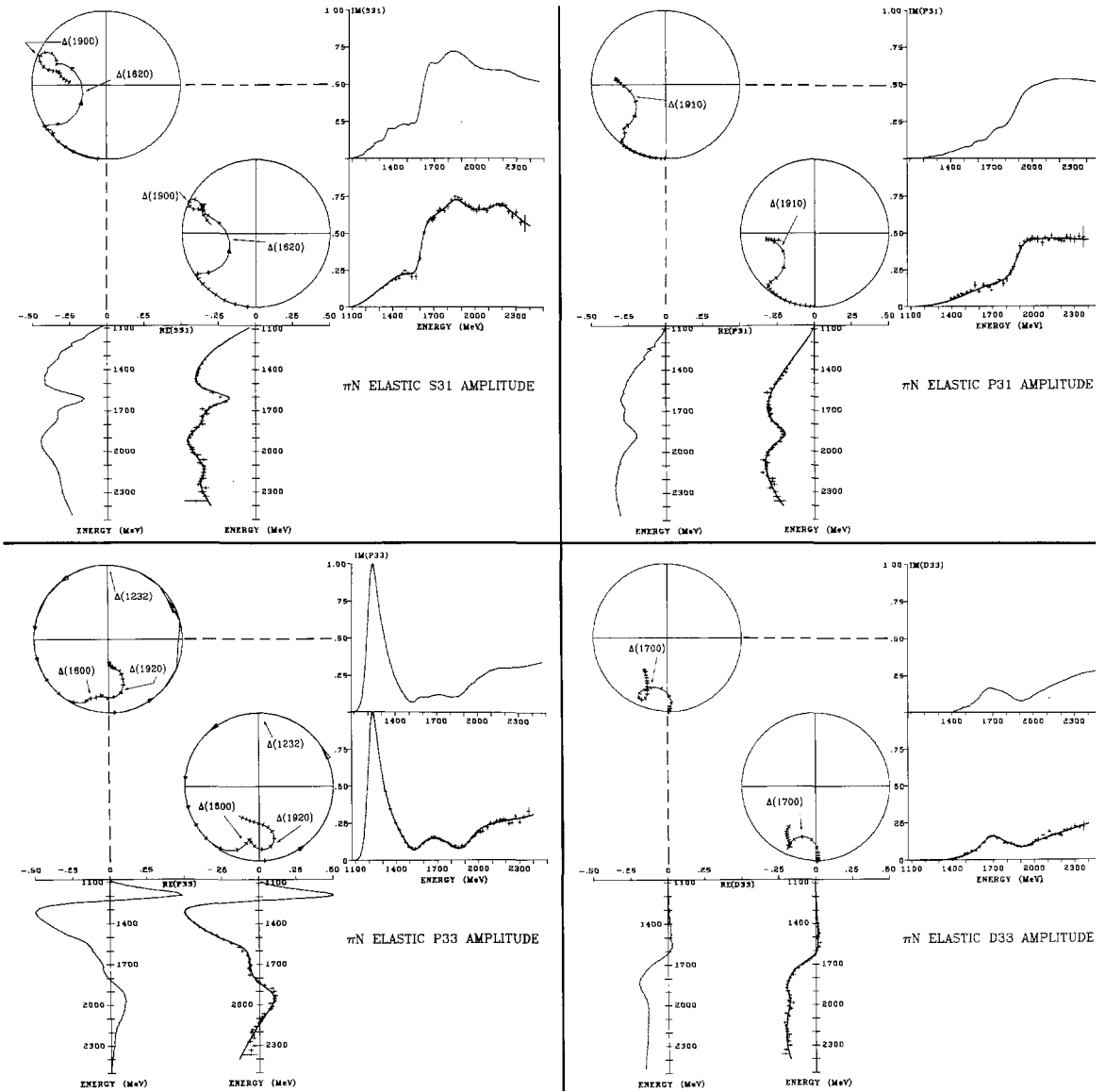


Fig. 1(d). The  $L_{2J,2J} = S_{31}, P_{31}, P_{33}$ , and  $D_{33}$  partial-wave amplitudes for  $\pi N$  elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

For notation, see key on page 91.

## Baryon Full Listings

*N*'s and  $\Delta$ 's

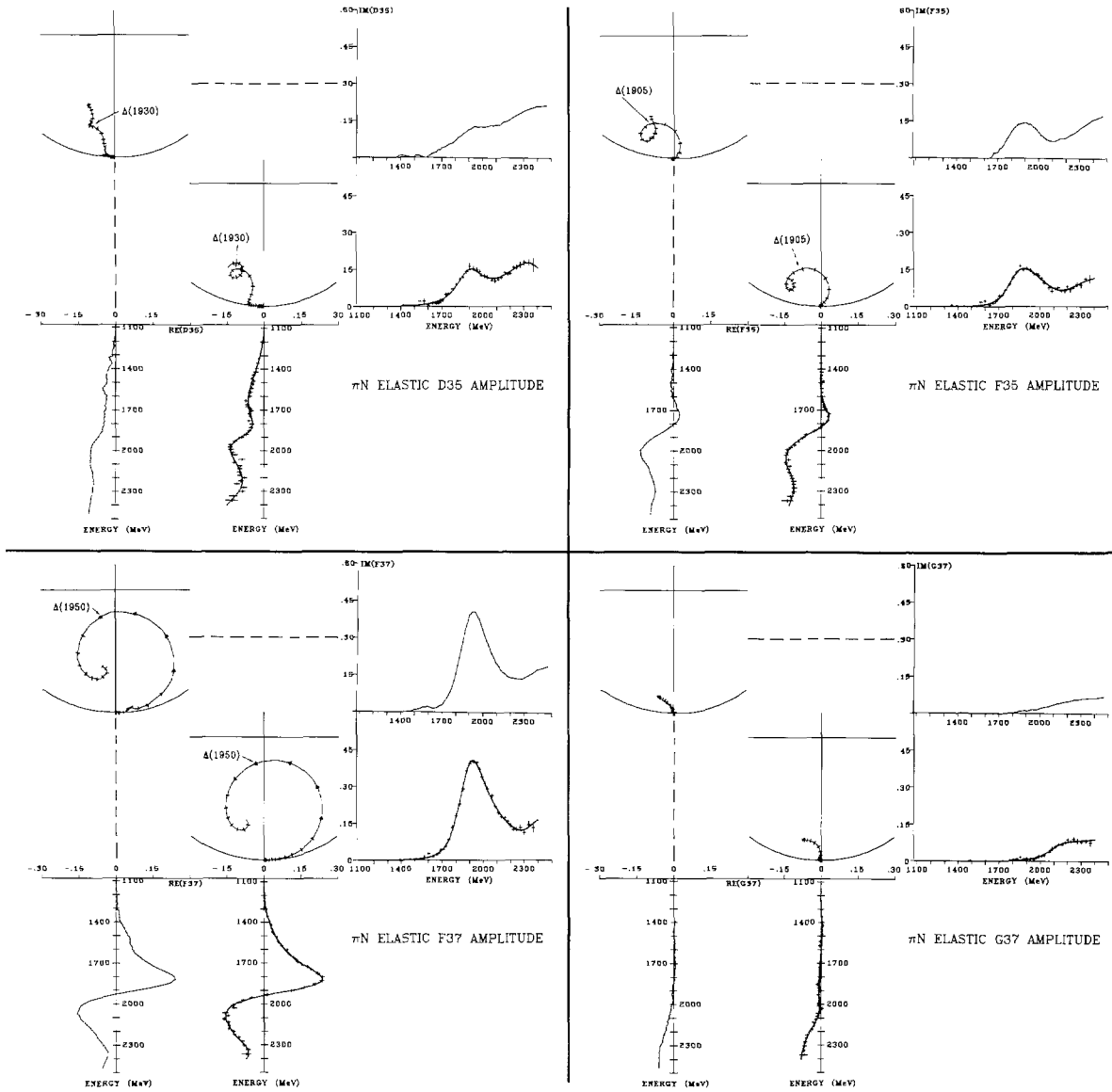


Fig. 1(e). The  $L_{2J;2J} = D_{35}, F_{35}, F_{37},$  and  $G_{37}$  partial-wave amplitudes for  $\pi N$  elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

# Baryon Full Listings

$N$ 's and  $\Delta$ 's

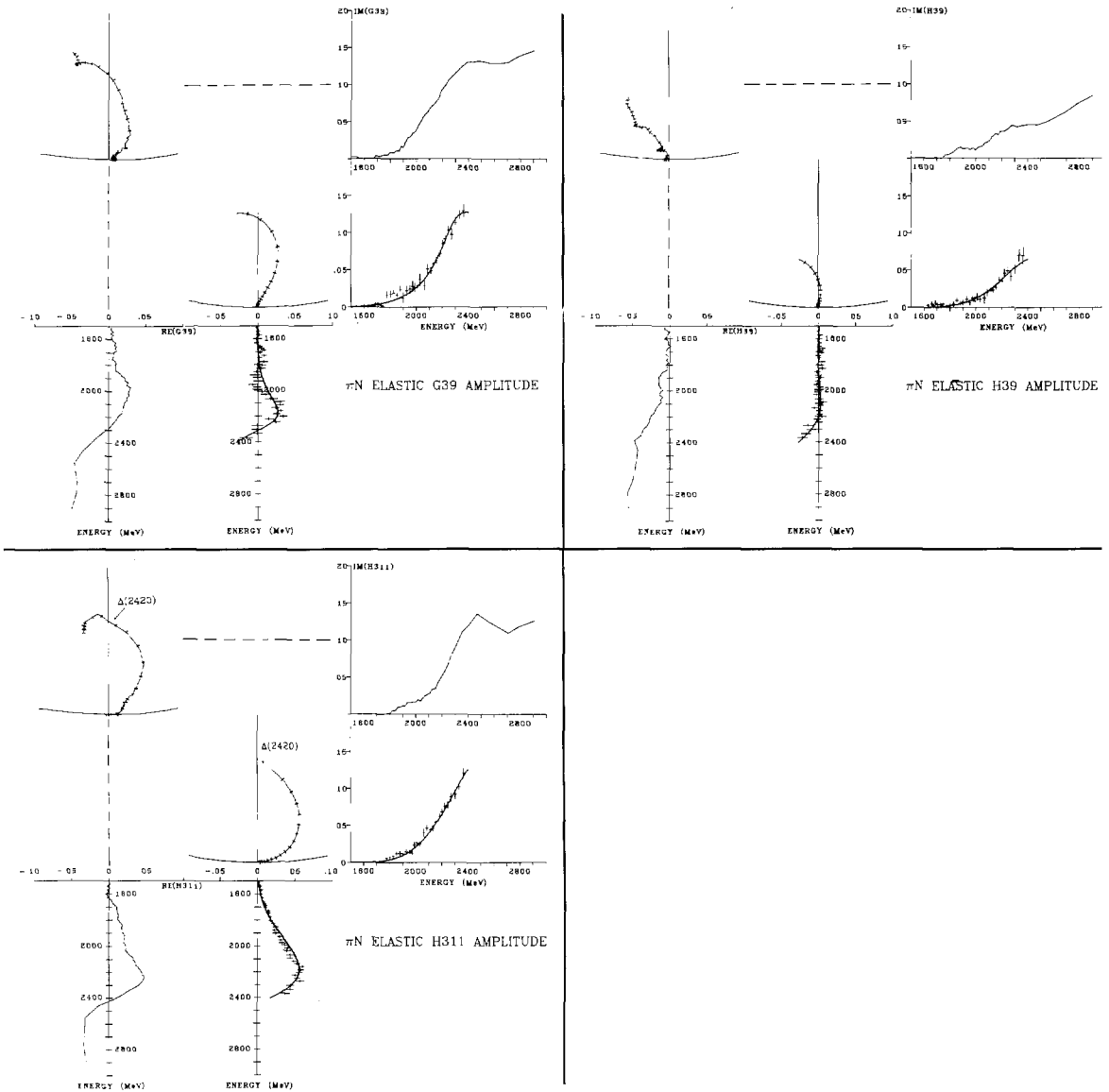


Fig. 1(f). The  $L_{2J,2J} = G_{39}, H_{39},$  and  $H_{311}$  partial-wave amplitudes for  $\pi N$  elastic scattering. The upper plot for each amplitude is from HOEHLER 79 and the lower one is from CUTKOSKY 80. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots (in the projections of the CUTKOSKY 80 amplitudes, the "data points" are results of energy-independent fits, and the curves are from an energy-dependent fit to join them).

## Baryon Full Listings

$N$ 's and  $\Delta$ 's

For notation, see key on page 91.

tion of the earlier Karlsruhe solution for  $D$  and higher partial waves constrained by the condition that the partial-wave dispersion relations be satisfied. Koch's solution improves the information on the shape of the resonance structures and has been used for a determination of the pole parameters.<sup>12</sup>

In a second paper,<sup>13</sup> Koch presented real parts of partial waves up to 500 MeV/ $c$  from a projection of fixed- $t$  dispersion relations (an exact version of the approach of Ref. 14). Agreement with the results of the quite different application of analyticity constraints in the first paper is in general good, which justifies the use of Mandelstam analyticity in partial-wave analysis. A smooth solution, taking the  $S$  and  $P$  waves from Ref. 13 and the  $D$  and higher waves from Ref. 11, is being prepared. The solution will be checked by calculating the zero trajectories and should be a good starting solution in the analysis of new data.

The data are still too poor for a good test of predictions for the highest resonances (masses  $> 2.2$  GeV). Evidence for resonances in this range has been reported by Koch<sup>6</sup> and by Hendry.<sup>15</sup> See also Ref. 7.

**Determination of resonance parameters:** Since a dynamical theory of  $\pi N$  scattering does not yet exist, the resonance parameters are not uniquely defined. One can fit a partial-wave amplitude with a phenomenological ansatz consisting of a generalized Breit-Wigner form plus a background term, and most of the earlier analyses, including the first CMU-LBL analysis and the KH 78 analysis,<sup>6</sup> used a prescription of this type. A more sophisticated multichannel coupled resonance scheme was used in the more recent work of the CMU-LBL group.<sup>5</sup> The parameters listed in the Baryon Summary Table were obtained using these methods.

A difficulty that becomes more and more important as the energy increases is that some "background terms" such as diffraction and  $\rho$ -exchange make contributions to the partial waves that resemble highly inelastic resonances (see Sec. 2.4.1.1 in Ref. 7). The energy dependences are different, but at high energies the speed with which an amplitude traverses the complex plane cannot be accurately determined due to insufficient data. Furthermore, it is a dynamical question whether this background is part of the resonance mechanism.

If the resonances are ordered according to the shapes of their Argand plots, one finds a continuous transition from textbook-type resonances to tiny wiggles superimposed on a huge background. The Baryon Summary Table lists all objects that have a "resonance-like" shape of the Argand diagram and a maximum of the speed. It

is up to the reader to decide which of these objects are "resonances" in the framework of his or her model.

The above discussion shows that a comparison of the resonance masses listed in the Baryon Summary Table with predictions from lattice calculations or from quark-shell, bag, Skyrmion, or other models is quite uncertain, especially where mass differences are small, since the models cannot yet describe the scattering process and take into account the background.

A cautionary example has been given by Blankleider and Walker,<sup>16</sup> who used separable potentials for the reactions  $\pi N \rightarrow \pi N$ ,  $\pi N \rightarrow \pi \Delta$ , and  $\pi \Delta \rightarrow \pi \Delta$  and got excellent fits to the  $P_{11}$ ,  $D_{13}$ ,  $D_{15}$ , and  $S_{31}$  waves without explicitly introducing the lowest 4-star resonances in these waves. Since various parameters have been fitted, the authors do not claim that their mechanism explains the resonances. But the opening of an inelastic channel contributes to the shape of a resonant amplitude, as may also be seen from the partial wave dispersion relation (or from the discussion by Ball and Frazer<sup>17</sup>). This effect is disregarded in the usual extraction of resonance parameters from models.

The Baryon Listings contain a second set of resonance parameters, the locations and residues of the resonance poles on the second sheet of the  $s$  plane. These may be determined in a (more or less) model-independent way. Table 2 summarizes some of the recent results.<sup>2(b),5,12</sup> Note, however, that Fonda et al.<sup>18</sup> were able to fit the resonant  $P_{33}$  amplitude *without a pole*. A theoretical assumption that excludes parametrizations of this type is needed.

Remarkably, there exist families of resonances in which the splittings of the pole positions are comparable with the errors; i.e., degeneracy is not excluded.<sup>7</sup> For example, all six isospin-1/2 partial waves from  $S_{11}$  to  $F_{15}$  have a well-established resonance with a pole near  $\sqrt{s} = (1665-60i)$  MeV, and at least six of the seven possible isospin-3/2 resonances from  $S_{31}$  to  $F_{37}$  have a pole near  $(1880-120i)$  MeV.

We have not included in the Listings the zeros of the partial-wave amplitudes given in Ref. 2(b) because a zero in the neighborhood of a resonance pole gives only information on the background. However, zero trajectories of the invariant and transversity amplitudes may be of fundamental importance (see Sect. 2.4.3 in Ref. 7).

**Inelastic 2-body reactions:** Partial-wave analyses of the inelastic 2-body reactions  $\pi N \rightarrow N\eta$ ,  $\Delta K$ , and  $\Sigma K$  are similar to  $\pi N \rightarrow \pi N$  analyses. However, since the data are far less complete and accurate, energy-dependent parametrizations must be used.

# Baryon Full Listings

$N$ 's and  $\Delta$ 's

Table 2. Recent determinations of pole parameters of 3- and 4-star  $N$  and  $\Delta$  resonances. Cutkosky et al.<sup>5</sup> and Arndt et al.<sup>2(b)</sup> have taken into account inelastic channels in the isobar approximation. In general, a resonance has a pole in several sheets of the energy plane. The parameters here are of the pole reached most directly from the physical region. In certain cases, this condition is ambiguous because a strong inelastic channel ( $\Delta\pi$ ,  $N\eta$ ,  $N\rho$ , etc.) opens within the width of the resonance. If theoreticians working on quark or bag models use the parameters, they should keep in mind that their models do not include the effects caused by the opening of inelastic channels.<sup>16</sup> In particular, they should not try to describe the two poles given by Arndt et al.<sup>2(b)</sup> for the  $N(1440) P_{11}$ . A splitting of this resonance was claimed earlier from an analysis of the elastic data alone, but this was not tenable (see the remark in Sec. 2.1.8 of Ref. 7). The pole parameters determined by M. Sararu<sup>12</sup> for a selection of resonances follow from Koch's smoothed version<sup>11</sup> of the Karlsruhe solution without taking into account data for inelastic scattering.

Resonance	Pole position (MeV)		Residue		Ref. <sup>†</sup>
	Re $W$	$-2 \times \text{Im} W$	$ r $ (MeV)	$\theta$ ( $^\circ$ )	
$N(1440) P_{11}$	1375 $\pm$ 30	180 $\pm$ 40	52 $\pm$ 5	-100 $\pm$ 35	C
	1355	200	62	-108	A
	1416	156	118	-4	
$N(1520) D_{13}$	1510 $\pm$ 5	114 $\pm$ 10	35 $\pm$ 2	-12 $\pm$ 5	C
	1508	124	40	-9	A
$N(1535) S_{11}$	1510 $\pm$ 50	260 $\pm$ 80	120 $\pm$ 40	+15 $\pm$ 45	C
	1464	150	40	-44	A
$N(1650) S_{11}$	1640 $\pm$ 20	150 $\pm$ 30	60 $\pm$ 10	-75 $\pm$ 25	C
	1656	108	34	-54	A
$N(1675) D_{15}$	1660 $\pm$ 10	140 $\pm$ 10	31 $\pm$ 5	-30 $\pm$ 10	C
	1658	136	32	-20	A
$N(1680) F_{15}$	1667 $\pm$ 5	110 $\pm$ 10	34 $\pm$ 2	-25 $\pm$ 5	C
	1668	110	33	-18	A
	1671	122	25	-20	S
$N(1700) D_{13}$	1660 $\pm$ 30	90 $\pm$ 40	6 $\pm$ 3	0 $\pm$ 50	C
	1676	48	2	+43	A
$N(1710) P_{11}$	1690 $\pm$ 20	80 $\pm$ 20	8 $\pm$ 2	+175 $\pm$ 35	C
	(not seen)				A
$N(1720) P_{13}$	1680 $\pm$ 30	120 $\pm$ 40	8 $\pm$ 2	-160 $\pm$ 30	C
	1690	66	3.7	-138	A
	1670	188	8	-127	S
$N(2190) G_{17}$	2100 $\pm$ 50	400 $\pm$ 160	25 $\pm$ 10	-30 $\pm$ 50	C
	2056	580	40	-18	S
$N(2220) H_{19}$	2160 $\pm$ 80	480 $\pm$ 100	45 $\pm$ 20	-45 $\pm$ 25	C
	2130	340	19	-47	S
$N(2250) G_{19}$	2150 $\pm$ 50	360 $\pm$ 100	20 $\pm$ 6	-50 $\pm$ 20	C
$N(2600) I_{111}$	2589	460			S
$\Delta(1232) P_{33}$	1210 $\pm$ 1	100 $\pm$ 2	53 $\pm$ 2	-47 $\pm$ 1	C
	1211	102	56	-30	A
	1209	100			S
$\Delta(1620) S_{31}$	1600 $\pm$ 15	120 $\pm$ 20	15 $\pm$ 2	-110 $\pm$ 20	C
	1592	108	13	-117	A
$\Delta(1700) D_{33}$	1675 $\pm$ 25	220 $\pm$ 40	13 $\pm$ 3	-20 $\pm$ 25	C
	1674	336	32	-24	A
	1680	226	14	+34	S
$\Delta(1900) S_{31}$	1870 $\pm$ 40	180 $\pm$ 50	10 $\pm$ 3	+20 $\pm$ 40	C
$\Delta(1905) F_{35}$	1830 $\pm$ 40	280 $\pm$ 60	25 $\pm$ 8	-50 $\pm$ 20	C
	1872	228	23	-13	A
	1850	220	10	-11	S
$\Delta(1910) P_{31}$	1880 $\pm$ 30	200 $\pm$ 40	20 $\pm$ 4	-90 $\pm$ 30	C
	1883	392	27	-89	S
$\Delta(1920) P_{33}$	1900 $\pm$ 80	300 $\pm$ 100	24 $\pm$ 4	-150 $\pm$ 30	C
	(not seen)				A
$\Delta(1930) D_{35}$	1890 $\pm$ 50	260 $\pm$ 60	18 $\pm$ 6	-20 $\pm$ 40	C
$\Delta(1950) F_{37}$	1890 $\pm$ 15	260 $\pm$ 40	50 $\pm$ 7	-33 $\pm$ 8	C
	1864	216	50	-20	A
	1890	242	32	-22	S
$\Delta(2420) H_{311}$	2360 $\pm$ 100	420 $\pm$ 100	18 $\pm$ 6	-30 $\pm$ 40	C

<sup>†</sup>C = Cutkosky et al.,<sup>5</sup> A = Arndt et al.,<sup>2(b)</sup> and S = Sararu.<sup>12</sup>

The best results, which give resonance masses and widths as well as couplings, follow from the  $\pi^- p \rightarrow \Delta K^0$  data of the Rutherford group.<sup>19</sup> In one analysis, the nonresonant and high waves were represented by a reggeized  $K^*$  exchange term.<sup>19,20</sup> In another analysis, a Lagrangian model was used for the long-range forces.<sup>21</sup> In general, agreement with the  $\pi N \rightarrow \pi N$  analyses is good, but there are discrepancies for the  $P_{11} N(1710)$  and  $D_{15} N(1675)$  widths and for the  $D_{15} N(2200)$  mass.

In an analysis of the less accurate  $\pi^- p \rightarrow n\eta$  data,<sup>22</sup> the partial waves were parametrized as Breit-Wigner resonances without background. The resonance spectrum was assumed and the data were used to determine the  $n\eta$  couplings. For resonances with relatively large couplings, the masses and widths were varied in a second step.

The results derived from the bubble-chamber data for  $\pi^+ p \rightarrow \Sigma^+ K^+$ <sup>23</sup> have large uncertainties. Values of the resonance masses were assumed and Breit-Wigner forms and an empirical ansatz for the background were used for partial waves up to  $F$  waves (the  $G$  waves are probably not negligible at 1.7 GeV/c). The recent addition of precise data from 1820 to 2350 MeV<sup>24</sup> has allowed an improved analysis.<sup>25</sup> The solution found is unique. Above 2 GeV, all the resonances with two or more stars are seen, but none of the 1-star states is supported.

In a recent note,<sup>26</sup> Isgur has pointed out that distortions of resonance couplings can occur in cases such as  $\pi N \rightarrow \Delta \rightarrow \Sigma K$  if the threshold is just below the resonance mass.

A possible new resonance, an  $N(1960)$ , has been detected in the final state  $\Sigma(1385)^- K^+$  in neutron carbon reactions.<sup>27</sup> If confirmed, the small width (27  $\pm$  15 MeV) indicates an exotic nature.

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## Baryon Full Listings

$N$ 's and  $\Delta$ 's

For notation, see key on page 91.

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### III. The $\pi N \rightarrow N\pi\pi$ reaction

(by D.M. Manley, Lawrence Livermore National Laboratory)

Dalitz plots reveal that the  $\pi N \rightarrow N\pi\pi$  reaction is dominated by the formation of  $N$  and  $\Delta$  resonances which decay into quasi-2-body channels. This observation prompts the analysis of this reaction by isobar models in which the partial-wave amplitudes are represented by a coherent sum over quasi-2-body channels, such as  $\pi N \rightarrow \Delta(1232)\pi$ ,  $\pi N \rightarrow N\rho$ ,  $\pi N \rightarrow N(\pi\pi)_S$ , and  $\pi N \rightarrow N(1440)\pi$ , where  $(\pi\pi)_S$  is the strong isospin-0,  $S$ -wave  $\pi\pi$  interaction. The resulting fitted amplitudes contain valuable information about the couplings of the  $N$  and  $\Delta$  resonances to the quasi-2-body channels.

While the analyses determine the relative phases of the quasi-2-body amplitudes, their overall phase is arbitrary. In this edition, all couplings are expressed according to the "baryon-first" convention,<sup>1</sup> with the overall phase determined by choosing the coupling sign of the  $S_{31}\Delta(1620)$  to the  $\Delta\pi$  channel to be negative. Experimentally, this coupling is strong in all analyses performed thus far. Furthermore, quark-model calculations for this coupling are simplified by the absence of other low-mass  $S_{31}$  resonances. Further details of the isobar-model formalism and a definition of the couplings are given in our 1982 edition.<sup>2</sup>

The Baryon Summary Table includes branching fractions for the  $N\pi\pi$  quasi-2-body channels. For this edition, these branching fractions were re-evaluated (a number of the branching fractions have been changed considerably) based upon the results of five energy-independent partial-wave analyses (IPWA), one of which is new. The Listings give the results from these five. Further details of the four older analyses are also discussed in our 1982 edition.

LONGACRE 75 and 78 (LBL-SLAC)<sup>3</sup> are resonance analyses of the partial-wave solution of Herndon et al.,<sup>4</sup> which was obtained from an isobar-model analysis of 170,000  $\pi^-p \rightarrow n\pi^-\pi^+$ ,  $\pi^-p \rightarrow p\pi^-\pi^0$ , and  $\pi^+p \rightarrow p\pi^+\pi^0$  events between 1300 and 1990 MeV c.m. energy. The  $\Delta(1232)\pi$ ,  $N\rho$ , and  $N(\pi\pi)_S$  channels were included. The couplings and  $T$ -matrix poles of nine  $N$  and five  $\Delta$  resonances are given in the Listings.

LONGACRE 77 (Saclay)<sup>5</sup> is a similar analysis of 91,000 events between 1360 and 1760 MeV. This work, unlike that of Herndon et al.,<sup>4</sup> included a nonzero range in the centrifugal barrier term that enters into the isobar-model parametrization. The importance of a

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nonzero range has been discussed by Dolbeau et al.<sup>1</sup> Several waves (especially  $N\rho$  waves) that were small in the LBL-SLAC solution<sup>4</sup> contribute significantly in this analysis, mainly because of the choice of a nonzero-range barrier. The Listings include couplings and pole positions for ten  $N$  and six  $\Delta$  resonances, including a  $P_{13}N(1540)$  and a  $P_{31}\Delta(1550)$ , which this analysis suggested for the first time.

NOVOSELLER 78 (Cal Tech)<sup>6</sup> is an analysis between 1630 and 1990 MeV based in part on the analysis of Herndon et al.<sup>4</sup> Starting with the LBL-SLAC solution as input, a solution was generated without the unitarity constraints imposed by Herndon et al.<sup>4</sup> Another solution was generated in a similar manner by including in the fit a high-partial-wave background derived from single-pion exchange. The most important contribution from this background is in the  $D$ -wave  $N\rho$  channel of the  $G_{17}$  wave. It was concluded that the contribution from the  $N\pi\pi$  channel, with the two pions in the isospin-2  $S$  wave, is small below 1970 MeV. It also was concluded that the solution of Herndon et al.<sup>4</sup> may be questionable for the  $P_{11}$ ,  $P_{13}$ , and  $F_{35}$  waves above 1700 MeV because of inconsistencies in the solutions obtained under different fitting conditions.

BARNHAM 80 (Imperial College)<sup>7</sup> is an analysis of 44,000  $\pi^+p \rightarrow p\pi^+\pi^0$  and  $\pi^+p \rightarrow n\pi^+\pi^+$  events between 1400 and 1700 MeV. The analysis included  $\Delta(1232)\pi$ ,  $N\rho$ , and  $N(1440)\pi$  channels, and used a zero-range barrier. Evidence was found for the  $P_{31}\Delta(1550)$  seen by LONGACRE 77, and for significant decay of the  $P_{33}\Delta(1600)$  into  $N(1440)\pi$ , a mode first considered by this analysis. The Listings include couplings for four  $\Delta$  resonances.

MANLEY 84 (VPI&SU)<sup>8</sup> is a major new analysis of 241,000  $\pi^-p \rightarrow n\pi^-\pi^+$ ,  $\pi^-p \rightarrow p\pi^-\pi^0$ ,  $\pi^+p \rightarrow p\pi^+\pi^0$ , and  $\pi^+p \rightarrow n\pi^+\pi^+$  events between 1320 and 1930 MeV. The main channels considered were  $\Delta(1232)\pi$ ,  $N\rho$ ,  $N(\pi\pi)_S$ , and  $N(1440)\pi$ . Also investigated, however, were the channels  $N(1520)\pi$ ,  $N(1535)\pi$ , and  $N(1675)\pi$ . Actual resonance parameters have not been published, so the Listings include only the signs and sizes (large or small) of the couplings. The analysis used a nonzero-range centrifugal barrier factor, and for the first time investigated the importance of partial waves with  $L > 3$  ( $G$  waves).

The analysis found evidence for a large  $D$ -wave  $N\rho$  decay in the  $G_{17}$  wave, possibly associated with the  $G_{17}N(2190)$ . Unlike earlier analyses, essentially all inelasticity in the  $P_{31}$  wave was accounted for by decay into  $N(1440)\pi$ . A large  $N(1440)\pi$  decay for the

$P_{33}\Delta(1600)$  was found, as did BARNHAM 80. No evidence was found for the  $P_{13}N(1540)$  and  $P_{31}\Delta(1550)$  claimed by LONGACRE 77 and BARNHAM 80. In agreement with LONGACRE 75, no evidence was found for an  $N\pi\pi$  decay of the  $D_{35}\Delta(1930)$ , even though  $\pi N$  elastic analyses predict a large inelastic decay for this resonance. Also in agreement with LONGACRE 75, a large  $F$ -wave  $\Delta(1232)\pi$  decay was found for the  $F_{37}\Delta(1950)$ , although the resonant phases determined by the two analyses differ by about  $60^\circ$ . It was pointed out<sup>8</sup> that this conflict could be resolved if the large  $P$ -wave  $N\rho$  decay observed by LONGACRE 75 and MANLEY 84 in the  $F_{35}$  wave were associated with a second resonance, an  $F_{35}\Delta(2000)$ , above the  $F_{35}\Delta(1905)$ . This higher mass  $F_{35}$  resonance is of interest because it is predicted by the Isgur-Karl quark model;<sup>9</sup> failure to observe it in  $\pi N$  elastic analyses is not surprising since it should couple weakly to the  $\pi N$  channel.

This analysis also found tentative evidence for several new resonances, including a  $P_{13}$  and an  $F_{15}$  at about 1850 MeV, and an  $S_{11}$ , a  $P_{11}$ , and a  $D_{13}$  at about 1900 MeV. (Since no actual masses or other parameters have been published, these resonances are not included in the Listings.)

Table 3 gives a compilation, adapted in part from MANLEY 84, of the signs of the  $\pi N \rightarrow N\pi\pi$  couplings. Only those decay channels found to be significant in at least one analysis are included.

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For notation, see key on page 91.

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Table 3. The signs of the  $\pi N \rightarrow N\pi\pi$  couplings according to four analyses. Letter subscripts give the final-state orbital angular momenta. The  $\rho$  states are labeled by twice the total final-state spin,  $2S$ . A "0" indicates a negligibly small coupling. A question mark indicates the sign is indeterminate or questionable because the resonant amplitude is rotated more than  $70^\circ$  away from the imaginary axis. An asterisk at the right indicates the coupling sign is considered to be well established by Manley 84.

Resonance	Decay channel	Long-acre 75,78 <sup>3</sup>	Long-acre 77 <sup>5</sup>	Barnham 80 <sup>7</sup>	Manley 84 <sup>8</sup>
$N(1440) P_{11}$	$(\Delta\pi)_P$	+	+	+	+
	$(N\rho_1)_P$	-	-?	-	0
	$(N\rho_3)_P$	0	+	0	0
	$N(\pi\pi)_S$	-?	-	-	+
$N(1520) D_{13}$	$(\Delta\pi)_S$	-	-	-	-
	$(\Delta\pi)_D$	-	-	-	-
	$(N\rho_3)_S$	-	-?	-	-
	$N(\pi\pi)_S$	-	-?	0	0
$N(1535) S_{11}$	$(\Delta\pi)_D$	+?	0	0	0
	$(N\rho_1)_S$	-	-?	-	-
	$N(\pi\pi)_S$	+	+	+	+
$N(1540) P_{13}$	$(\Delta\pi)_P$	-	-?	0	0
	$(N\rho_1)_P$	-	+	0	0
$N(1650) S_{11}$	$(\Delta\pi)_D$	-	+	+	+
	$(N\rho_1)_S$	-?	+?	-	-
	$(N\rho_3)_D$	0	+	+	+?
	$N(\pi\pi)_S$	+	0	+	+
$N(1675) D_{15}$	$(\Delta\pi)_D$	+	+	+	+
	$(N\rho_3)_D$	0	-	-	-
	$N(\pi\pi)_S$	0	+	0	0
$N(1680) F_{15}$	$(\Delta\pi)_P$	-	-	-	-
	$(\Delta\pi)_F$	+	+	+	+
	$(N\rho_3)_P$	-	-	-	-
	$(N\rho_3)_F$	0	-	-	-
$N(1700) D_{13}$	$(N\rho_3)_F$	+	+	+	+
	$(\Delta\pi)_S$	-	0	-?	-?
	$(\Delta\pi)_D$	+	-?	+	+
	$(N\rho_3)_S$	+	-?	0	0
$N(1710) P_{11}$	$N(\pi\pi)_S$	+	0	+	+
	$(\Delta\pi)_P$	+	-	-	-
	$(N\rho_1)_P$	-	+	+	+
	$(N\rho_3)_P$	0	+	0	0
$N(1720) P_{13}$	$N(\pi\pi)_S$	-?	-	-?	-?
	$(\Delta\pi)_P$	0	-	0	0
	$(N\rho_1)_P$	+	-	0	0
	$(N\rho_3)_P$	0	+	0	0
$\Delta(1550) P_{31}$	$N(\pi\pi)_S$	0	-	0	0
	$(\Delta\pi)_P$	-	-	?	0
	$(N\rho_3)_P$	-	-	?	0
	$(\Delta\pi)_F$	+	+	+	+
$\Delta(1600) P_{33}$	$(\Delta\pi)_F$	0	-?	?	0
	$(N\rho_1)_P$	0	+	?	0
	$(N\rho_3)_P$	0	+	?	0
	$\dagger(N^*\pi)_P$	0	+	?	0
	$\dagger(N^*\pi)_P$	0	+	?	0
$\Delta(1620) S_{31}$	$(\Delta\pi)_D$	-	-	-	-
	$(N\rho_1)_S$	+	+	+	+
	$(N\rho_3)_D$	0	-	?	-
$\Delta(1700) D_{33}$	$(\Delta\pi)_S$	+	+	+	+
	$(\Delta\pi)_D$	+?	+	?	+
	$(N\rho_1)_D$	0	0	+	0
	$(N\rho_3)_S$	-	-?	?	+
$\Delta(1905) F_{35}$	$(\Delta\pi)_P$	0	-	+	+
	$(\Delta\pi)_F$	+	-	+	+
$\Delta(1910) P_{31}$	$(\Delta\pi)_P$	-	+	0	0
	$(N\rho_3)_P$	-	+	0	0
	$\dagger(N^*\pi)_P$	-	+	+	+
$\Delta(1950) F_{37}$	$(\Delta\pi)_F$	+	-	+	+
	$(N\rho_3)_F$	+?	-	0	0
$\Delta(2000) F_{35}$	$(N\rho_3)_P$	+	-	+	+

$\dagger$ The  $N^*$  here is the  $P_{11} N(1440)$ .

### IV. Photoproduction and Compton scattering

(by R.L. Crawford, University of Glasgow)

The  $\gamma N$  couplings of the  $N$  and  $\Delta$  resonances are obtained from partial-wave analyses of single-pion photoproduction on protons and neutrons and of Compton scattering on protons. The definitions of the couplings and a discussion of the methods used in the analyses may be found in our 1976 edition.<sup>1</sup> The large amount of pion photoproduction data, including many measurements from single and double polarization experiments, has permitted accurate measurement of many of the couplings and determination of the signs of many more. Recently, photon couplings in the second and third resonance regions have also been obtained from accurate measurements of proton Compton-scattering differential cross sections; the results, with a few exceptions, agree with those from photoproduction. All these analyses rely heavily upon  $\pi N \rightarrow \pi N$  analyses for information on the existence, masses, and widths of the resonances. The only photoproduction analyses that quote resonance masses and widths as well as couplings are BERENDS 75, BERENDS 77, BARBOUR 78, and CRAWFORD 80. These results are of interest since they give access to the charge +1 states of the resonances. In particular, the mass of the  $\Delta(1232)^+$  seems to be as well determined as those of the  $\Delta^{++}$  and  $\Delta^0$  seen in elastic  $\pi N$  scattering.

There are three main methods of analysis.

(a) *The simple isobar model*: The simple isobar model is an energy-dependent partial-wave analysis (DPWA) in which the partial waves are parametrized as Breit-Wigner resonances plus background. The model is relatively simple but is flexible enough to give good fits. However, there are possible problems about the uniqueness of the solutions obtained, and it is not clear how the form of the parametrization distorts the solutions.

The Listings give the results of the isobar analyses of TAKEDA 80 and BRATASHEVSKIJ 80 (photoproduction), and of ISHII 80 and WADA 84 (Compton scattering). The two photoproduction analyses use small data sets. The large-scale isobar analysis of METCALF 74 is now obsolete and has been omitted.

(b) *Fixed- $t$  dispersion relations (FTDR)*: Here only the imaginary parts of the production amplitudes are parametrized, and the real parts are calculated from them using fixed- $t$  dispersion relations. The resonance dominance of the imaginary parts permits a relatively simple parametrization scheme, and there are fewer



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ambiguity problems than in the isobar model. However, it is less flexible than the isobar model and gives poorer fits, and it can only be used properly in a large-scale analysis over a wide energy range.

The Listings give the results from the FTDR analyses of AZNAURYAN 77, BARBOUR 78, ARAI 80, CRAWFORD 80, FUJII 81, and AWAJI 81. NOELLE 78 is a hybrid analysis using FTDR in a coupled-channel isobar calculation.

(c) *Energy-independent analyses (IPWA)*: These evaluate the partial-wave amplitudes by fitting at essentially single energies. At low energies, Watson's theorem is used to fix the complex phases of many of the partial waves. This allows a unique solution but becomes difficult above the first resonance region due to the onset of inelasticity. BERENDS 77 is the only true IPWA that has been extended into the second resonance region. CRAWFORD 83 is an IPWA that depends on the CRAWFORD 80 FTDR analysis to give stable solutions, and is thus not independent of the energy-dependent analysis.

*New data in the Listings*: Reflecting the decline in interest in baryon physics, there is only one new analysis in the Listings. WADA 84 is an isobar-model analysis of Compton scattering on protons using data on differential cross sections in the second and third resonance regions. Most of the couplings agree well with previously determined values from photoproduction. The exceptions are  $A_{1/2}^p$  for the  $P_{11}N(1440)$  and the  $P_{33}\Delta(1600)$ , which disagree badly with other measurements, particularly with those from photoproduction. However, these results should not yet be taken seriously since the quality of the photoproduction data is much higher and restricts the values of the couplings more strongly than do the Compton scattering data. Also, WADA 84 does not fit the couplings of all the resonances in the energy range of the analysis. Many of the more important couplings are not varied but are given values from the photoproduction analysis of AWAJI 81. Fitting all the couplings might give results for them all that are compatible with the photoproduction results.

The Baryon Summary Table now gives the  $N\gamma$  branching fractions for the resonances whose couplings are considered to have an unambiguous sign. These branching fractions have been calculated from the couplings given by the most reliable analyses in the Listings using the partial width

$$\Gamma_\gamma = \frac{k^2}{\pi} \frac{2M_N}{M_R(2J+1)} \left[ |A_{1/2}|^2 + |A_{3/2}|^2 \right],$$

where  $M_N$  and  $M_R$  are the masses of the nucleon and resonance,  $J$  is the resonance spin, and  $k$  is the c.m. decay momentum.

*Resonance couplings in the Listings*: The Listings omit a number of analyses that now must be considered to be obsolete. These are ROSSI 73, HEMMI1 73, HEMMI2 73, BENEVENTANO 74, METCALF 74, KRIVETS 75 (all isobar model), MOORHOUSE 73, DEVENISH 73, KNIES 74, MOORHOUSE 74, DEVENISH2 74, CRAWFORD 75, and BARBOUR 76 (all FTDR). They may all be found in our 1982 edition.<sup>7</sup>

The errors quoted for the couplings in the Listings are calculated in different ways in the different analyses. They are therefore not comparable and should be used with care. FELLER 76, AZNAURYAN 77, and ARAI 80 obtain errors using the sensitivity of the "best possible"  $\chi^2$  to the value of each coupling. BARBOUR 78, CRAWFORD 80, and CRAWFORD 83 attempt to assess the effects of the systematic errors caused by the different possible parametrization schemes, especially those for the background, and quote errors that reflect this. AWAJI 81 gives errors which also include a contribution from the uncertainty in the  $\pi N$  elasticity used to calculate the couplings from the partial waves.

In general, it is likely that the systematic differences between the analyses caused by the different parametrization schemes are more indicative than are the individual errors quoted in each analysis.

Table 4 gives a compilation of the couplings from BARBOUR 78, ARAI 80, CRAWFORD 80, FUJII 81, AWAJI 81, and CRAWFORD 83. The errors given are a combination of the statistical errors quoted in the analyses and the systematic differences between them. The Table is unchanged from the previous edition except that two values are quoted for  $A_{1/2}$  of the  $S_{31}\Delta(1620)$  to take account of the large spread in values obtained for this coupling. This can be attributed to imaginary background in the partial wave, which is treated differently, or ignored, in the different analyses. The second value uses only the Glasgow analyses (BARBOUR 78, CRAWFORD 80, and CRAWFORD 83). The Glasgow FTDR analyses obtain stable and acceptable values for the mass and width of this resonance, and it is thus reasonable to infer that the coupling obtained is accurate.

For notation, see key on page 91.

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Table 4. A compilation of measured  $\gamma N$  decay couplings. Sources are given in the text.

(a) Proton-target couplings			
Resonance	Helicity	Couplings ( $\text{GeV}^{-1/2} \times 10^{-3}$ )	Status
$N(1440) P_{11}$	1/2	$-69 \pm 7$	good
$N(1520) D_{13}$	1/2	$-22 \pm 10$	good
	3/2	$+167 \pm 10$	good
$N(1535) S_{11}$	1/2	$+73 \pm 14$	good
$N(1650) S_{11}$	1/2	$+48 \pm 16$	good
$N(1675) D_{15}$	1/2	$+19 \pm 12$	good, nonzero
	3/2	$+19 \pm 12$	good, nonzero
$N(1680) F_{15}$	1/2	$-17 \pm 10$	good, nonzero
	3/2	$+127 \pm 12$	good
$N(1700) D_{13}$	1/2	$-22 \pm 13$	good, small
	3/2	$0 \pm 19$	fair, small
$N(1710) P_{11}$	1/2	$+5 \pm 16$	fair, small
$N(1720) P_{13}$	1/2	$+52 \pm 39$	poor
	3/2	$-35 \pm 24$	fair
$N(1990) F_{17}$	1/2	$+24 \pm 30$	poor
	3/2	$+31 \pm 55$	bad
$\Delta(1232) P_{33}$	1/2	$-141 \pm 5$	very good
	3/2	$-258 \pm 11$	very good
$\Delta(1550) P_{31}$	1/2	$+16 \pm 16$	doubtful
$\Delta(1600) P_{33}$	1/2	$-20 \pm 29$	poor, small
	3/2	$+1 \pm 22$	fair, small
$\Delta(1620) S_{31}$	1/2	$+19 \pm 16$	fair
	1/2	$+30 \pm 10$	good -- see text
$\Delta(1700) D_{33}$	1/2	$+116 \pm 17$	good
	3/2	$+77 \pm 28$	fair
$\Delta(1900) S_{31}$	1/2	$+10 \pm ?$	?
$\Delta(1905) F_{35}$	1/2	$+27 \pm 13$	good
	3/2	$-47 \pm 19$	fair
$\Delta(1910) P_{31}$	1/2	$-12 \pm 30$	poor
$\Delta(1920) P_{33}$	1/2	$+40 \pm ?$	?
	3/2	$+23 \pm ?$	?
$\Delta(1930) D_{35}$	1/2	$-30 \pm 40$	poor
	3/2	$-10 \pm 35$	poor
$\Delta(1950) F_{37}$	1/2	$-73 \pm 14$	good
	3/2	$-90 \pm 13$	good
(b) Neutron-target couplings			
Resonance	Helicity	Couplings ( $\text{GeV}^{-1/2} \times 10^{-3}$ )	Status
$N(1440) P_{11}$	1/2	$+37 \pm 19$	fair
$N(1520) D_{13}$	1/2	$-65 \pm 13$	good
	3/2	$-144 \pm 14$	good
$N(1535) S_{11}$	1/2	$-76 \pm 32$	fair
$N(1650) S_{11}$	1/2	$-17 \pm 37$	poor
$N(1675) D_{15}$	1/2	$-47 \pm 23$	fair
	3/2	$-69 \pm 19$	fair
$N(1680) F_{15}$	1/2	$+31 \pm 13$	good
	3/2	$-30 \pm 14$	good
$N(1700) D_{13}$	1/2	$0 \pm 56$	bad
	3/2	$-2 \pm 44$	bad
$N(1710) P_{11}$	1/2	$-5 \pm 23$	fair, small
$N(1720) P_{13}$	1/2	$-2 \pm 26$	fair, small
	3/2	$-43 \pm 94$	very bad
$N(1990) F_{17}$	1/2	$-49 \pm 45$	poor
	3/2	$-122 \pm 55$	poor

**The  $E/M$  ratio of the  $P_{33}\Delta(1232)$ :** There is continuing interest in the  $E_2/M_1$  ratio of the  $\Delta(1232)$ . In the simplest form of the quark model, it is predicted to be zero,<sup>2</sup> but in more recent models this is not the case,<sup>3</sup> an effect attributed to tensor quark-quark forces giving  $D$ -wave components in the wave functions. Tanabe and Ohta<sup>4</sup> have evaluated the ratio using a phenomenological model, derived from the work of Olsson,<sup>5</sup> which includes nonresonant background not usually included when extracting the  $\Delta \rightarrow N\gamma$  couplings. Effectively, this work defines renormalized couplings that are evaluated by fitting the model to the partial waves from the IPWA of BERENDS 75. The couplings obtained are  $A_{1/2} = -0.093 \pm 0.005 \text{ GeV}^{-1/2}$  and  $A_{3/2} = -0.142 \pm 0.009 \text{ GeV}^{-1/2}$ . These are not included in the Listings since it is not clear that they are the same quantities as are given there. The  $E_2/M_1$  ratio obtained is  $+0.037 \pm 0.004$ . A more recent, unpublished analysis, also based on the approach of Olsson and using IPWA of the first resonance region, has been carried out by Davidson, Mukhopadhyay, and Wittman.<sup>6</sup> The values obtained are  $A_{1/2} = -0.125 \pm 0.003 \text{ GeV}^{-1/2}$ ,  $A_{3/2} = -0.229 \pm 0.006 \text{ GeV}^{-1/2}$ , and  $E_2/M_1 = -0.015 \pm 0.002$ .

It should be noted that the errors in the compiled values of  $A_{1/2}$  and  $A_{3/2}$  in Table 4 do not accurately reflect the precision with which  $E_2/M_1$  is given in the Listings, since the errors contain components, such as in overall normalization, that do not contribute to this ratio. Calculation of  $E_2/M_1$  from the individual entries and then combining them gives  $-0.013 \pm 0.005$ .

### References for section IV

1. Particle Data Group, Rev. Mod. Phys. **48**, S157 (1976).
2. C.M. Becchi and G. Morpurgo, Phys. Lett. **17**, 352 (1965).
3. N. Isgur, G. Karl, and R. Koniuk, Phys. Rev. **D25**, 2394 (1983); G. Kälbermann and J.M. Eisenberg, Phys. Rev. **D28**, 71 (1983); J. Dey and M. Dey, Phys. Lett. **138B**, 200 (1984); and D. Drechsel and M.M. Giannini, Phys. Lett. **143B**, 329 (1984).
4. H. Tanabe and K. Ohta, Phys. Rev. **C31**, 1876 (1985).
5. M.G. Olsson, Nucl. Phys. **B78**, 55 (1974).
6. R. Davidson, N.C. Mukhopadhyay, and R. Wittman, Preprint, Rensselaer Polytechnic Inst., New York (1985).
7. Particle Data Group, Phys. Lett. **111B** (1982).

# Baryon Full Listings

$N$ 's and  $\Delta$ 's,  $p$ ,  $n$ ,  $N(1440)$

## V. Electroproduction

The excitation of the  $N$  and  $\Delta$  resonances has been investigated using pion and  $\eta$  electroproduction data. However, the level of activity in this field is presently low, and the interested reader is referred to our 1982 edition<sup>1</sup> and to an extensive review.<sup>2</sup>

Since 1982, additional information has been obtained from  $\pi^+$  electroproduction about the switching in dominance of the helicity 3/2 and 1/2 amplitudes of the  $D_{13} N(1520)$  and  $F_{15} N(1680)$  as  $Q^2$  moves from the photoproduction limit.<sup>3</sup> In  $\eta$  electroproduction,<sup>4</sup> the  $Q^2$  dependence of the  $S_{11} N(1535)$  cross section has been examined for  $Q^2$  up to 3 GeV<sup>2</sup>. It is found that the slow decrease of production of the  $S_{11}$  compared to that of the  $D_{13} N(1520)$  continues to these values of  $Q^2$ . Both these features of electroproduction are already well known and have been described in the reviews mentioned above.

## References for section V

1. Particle Data Group, Phys. Lett. **111B** (1982).
2. F. Foster and G. Hughes, Rep. Prog. Phys. **46**, 1445 (1983).
3. H. Breuker et al., Zeit. Physik **C13**, 113 (1982).
4. F.W. Brasse et al., Zeit. Physik **C22**, 33 (1984).

## VI. Production experiments

Partial-wave analyses of course separate partial waves, whereas a peak in a cross section or an invariant mass distribution usually cannot be disentangled from background and analyzed for its quantum numbers; and more than one resonance may be contributing to the peak. We used to have separate entries in the Baryon Listings for bumps seen in production experiments in the 1440-MeV region, the 1520-MeV region, etc., but these have been removed. They last appeared in our 1982 edition.<sup>1</sup>

## Reference for section VI

1. Particle Data Group, Phys. Lett. **111B** (1982).

## $N$ BARYONS ( $S=0$ , $I=1/2$ )

$P$

$$I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$$

SEE STABLE PARTICLES.

$n$

$$I(J^P) = \frac{1}{2}(\frac{1}{2}^-)$$

SEE STABLE PARTICLES.

$N(1440) P_{11}$

$$I(J^P) = \frac{1}{2}(\frac{1}{2}^-) \quad \text{Status: ****}$$

MOST OF THE RESULTS PUBLISHED BEFORE 1975 ARE NOW OBSOLETE AND HAVE BEEN OMITTED. THEY MAY BE FOUND IN OUR 1982 EDITION (PHYSICS LETTERS 111B).

IN ADDITION, RESULTS IN THIS REGION FROM PRODUCTION EXPERIMENTS, WHICH USED TO BE LISTED SEPARATELY AS THE NEXT ENTRY, HAVE BEEN ENTIRELY REMOVED. THEY TOO MAY BE FOUND IN OUR 1982 EDITION.

### $N(1440)$ MASS (MeV)

M	A	1415. OR 1390.	LONGACRE	75 IPWA	PI N TO 2P1 N
M	A	THE 2 SETS OF PARAMETERS ARE FROM METHODS 1 AND 2 OF LONGACRE 75.			
M		(1460.0)	BERENDS	77 IPWA	PI-N PHOTOPROD.
M	B	(1380.0)	LONGACRE	77 IPWA	PI N TO 2P1 N
M	B	ALL LONGACRE77 PARAMETERS ARE FROM SOLUTION 52, EXCEPT FOR THE POLE POSITION WHICH IS FROM SOLUTIONS S1 AND C1.			
M		(1417.0)	BARBOUR	78 DPWA	PI-N PHOTOPROD.
M	C	(1472.0)	BAKER	79 DPWA	0 PI- P TO ETA N
M		(1450.0) (30.0)	CUTKOSKY	79 IPWA	PI N TO P1 N
M		1410.0 12.0	HOEHLER	79 IPWA	PI N TO P1 N
M		(1411.0)	CRAWFORD	80 DPWA	PI N PHOTOPROD.
M		1440.0 30.0	CUTKOSKY	80 IPWA	PI N TO P1 N

### $N(1440)$ WIDTH (MeV)

W	A	180. OR 200.	LONGACRE	75 IPWA	PI N TO 2P1 N
W		(279.0)	BERENDS	77 IPWA	PI-N PHOTOPROD.
W	B	(200.0)	LONGACRE	77 IPWA	PI N TO 2P1 N
W		(331.0)	BARBOUR	78 DPWA	PI-N PHOTOPROD.
W	C	(113.0)	BAKER	79 DPWA	0 PI- P TO ETA N
W		(370.0) (80.0)	CUTKOSKY	79 IPWA	PI N TO P1 N
W		135.0 10.0	HOEHLER	79 IPWA	PI N TO P1 N
W		(334.0)	CRAWFORD	80 DPWA	PI N PHOTOPROD.
W		340.0 70.0	CUTKOSKY	80 IPWA	PI N TO P1 N

### $N(1440)$ REAL PART OF POLE POSITION (MeV)

RE	B	1360.0 OR 1333.0	LONGACRE	77 IPWA	PI N TO 2P1 N
RE		1381.0 OR 1379.0	LONGACRE	78 IPWA	PI N TO 2P1 N
RE		(1369.0)	CUTKOSKY	79 IPWA	PI N TO P1 N
RE	F	1375.0 30.0	CUTKOSKY	80 IPWA	PI N TO P1 N
RE	F	(1359.0)	ARNDT	85 DPWA	PI N TO P1 N
RE	F	ARNDT 84 FIND A SECOND P11 POLE AT (1410, -80) MEV.			

### $N(1440)$ -2\*IMAG PART OF POLE POSITION (MeV)

IM	B	167.0 OR 234.0	LONGACRE	77 IPWA	PI N TO 2P1 N
IM		209.0 OR 210.0	LONGACRE	78 IPWA	PI N TO 2P1 N
IM		(178.0)	CUTKOSKY	79 IPWA	PI N TO P1 N
IM		180.0 40.0	CUTKOSKY	80 IPWA	PI N TO P1 N
IM	F	(200.0)	ARNDT	85 DPWA	PI N TO P1 N

### $N(1440)$ REAL PART OF ELASTIC POLE RESIDUE (MeV)

RER		(-9.0)	CUTKOSKY	79 IPWA	PI N TO P1 N
RER		-9.0 31.0	CUTKOSKY	80 IPWA	PI N TO P1 N

### $N(1440)$ IMAG PART OF ELASTIC POLE RESIDUE (MeV)

IMR		(-48.0)	CUTKOSKY	79 IPWA	PI N TO P1 N
IMR		-51.0 7.0	CUTKOSKY	80 IPWA	PI N TO P1 N

### $N(1440)$ ABSOLUTE VALUE OF POLE RESIDUE (MeV)

ABS		52.0 5.0	CUTKOSKY	80 IPWA	PI N TO P1 N
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For notation, see key on page 91.

Baryon Full Listings
N(1650), N(1675)

N(1650) -> n gamma, helicity = 1/2 (GeV^-1/2)
Table with columns for author, values, and experimental conditions.

N(1675) REAL PART OF ELASTIC POLE RESIDUE (MeV)
Table with columns for author, values, and experimental conditions.

N(1675) IMAG PART OF ELASTIC POLE RESIDUE (MeV)
Table with columns for author, values, and experimental conditions.

REFERENCES FOR N(1650)

FOR EARLY REFERENCES, SEE PHYSICS LETTERS 1118 (1982).

Table listing references for N(1650) with columns for author, year, and journal.

N(1675) PARTIAL DECAY MODES

Table listing partial decay modes for N(1675) with columns for mode, author, and decay masses.

N(1675) BRANCHING RATIOS

Table listing branching ratios for N(1675) with columns for mode, author, and values.

N(1675) D15

I(J^P) = 1/2(1/2^-) Status: \*\*\*\*

MOST OF THE RESULTS PUBLISHED BEFORE 1975 ARE NOW OBSOLETE AND HAVE BEEN OMITTED. THEY MAY BE FOUND IN OUR 1982 EDITION (PHYSICS LETTERS 1118).

IN ADDITION, RESULTS IN THIS REGION FROM PRODUCTION EXPERIMENTS, WHICH USED TO BE LISTED SEPARATELY IN AN ENTRY FOLLOWING THE N(1700), HAVE BEEN ENTIRELY REMOVED. THEY TOO MAY BE FOUND IN OUR 1982 EDITION.

N(1675) MASS (MeV)

Table listing mass measurements for N(1675) with columns for author, value, and experimental conditions.

N(1675) WIDTH (MeV)

Table listing width measurements for N(1675) with columns for author, value, and experimental conditions.

N(1675) REAL PART OF POLE POSITION (MeV)

Table listing real part of pole position for N(1675) with columns for author, value, and experimental conditions.

N(1675) -2\*IMAG PART OF POLE POSITION (MeV)

Table listing imaginary part of pole position for N(1675) with columns for author, value, and experimental conditions.

N(1675) PHOTON DECAY AMPLITUDES (GeV^-1/2)

FOR DEFINITION OF GAMMA-NUCLEON DECAY AMPLITUDES, SEE MINI-REVIEW PRECEDING THE BARYON LISTINGS.

Table listing photon decay amplitudes for N(1675) with columns for mode, author, and values.

















For notation, see key on page 91.

# Baryon Full Listings

$N(2090)$ ,  $N(2100)$ ,  $N(2190)$

### $N(2090) - 2^*IMAG PART OF POLE POSITION (MeV)$

IM	139.0 OR 131.0	LONGACRE 78 IPWA	PI N TO 2PI N
IM	350.0 100.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $N(2090) REAL PART OF ELASTIC POLE RESIDUE (MeV)$

RER	40.0 20.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $N(2090) IMAG PART OF ELASTIC POLE RESIDUE (MeV)$

IMR	0.0 60.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $N(2090) PARTIAL DECAY MODES$

			DECAY MASSES
P1	$N(2090) \rightarrow N\pi$		938- 140
P2	$N(2090) \rightarrow \Delta K$		1116- 498

### $N(2090) BRANCHING RATIOS$

$N(2090) \rightarrow (N\pi)/total$		(P1)	
R1	0.09 0.05	HOEHLER 79 IPWA	PI N TO PI N
R1	0.18 0.08	CUTKOSKY 80 IPWA	PI N TO PI N
$N(2090) in N\pi \rightarrow \Delta K$		SORT (P1*P2)	
R2	NOT SEEN	SAXON 80 DPWA	0 PI- P TO K LAM

### REFERENCES FOR $N(2090)$

LONGACRE 78 PR D17 1795 (LBL-SLAC)  
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL. 12-1 (KARL)IJP  
 ALSO 80 TORONTO CONF 3 KAISER, KOCH, PIETARINEN (KARL)IJP  
 CUTKOSKY 80 TORONTO CONF 19 R KOCH (KARL)IJP  
 SAXON 80 NP 5162 522 -FORSYTH, BARCOCK, KELLY, HENDRICK (CARN+LBL)IJP  
 -BAKER, BELL, BLISSSETT, BLOODWORTH+(RHEL+BRIS)IJP

### $N(2100) P_{11}$

$I(J^P) = \frac{1}{2}(\frac{1}{2}^-)$  Status: \*

OMITTED FROM SUMMARY TABLE

### $N(2100) MASS (MeV)$

M	2050.0 20.0	HOEHLER 79 IPWA	PI N TO PI N
M	2125.0 75.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $N(2100) WIDTH (MeV)$

W	200.0 30.0	HOEHLER 79 IPWA	PI N TO PI N
W	260.0 100.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $N(2100) REAL PART OF POLE POSITION (MeV)$

RE	2120.0 40.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $N(2100) - 2^*IMAG PART OF POLE POSITION (MeV)$

IM	240.0 80.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $N(2100) REAL PART OF ELASTIC POLE RESIDUE (MeV)$

RER	11.0 7.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $N(2100) IMAG PART OF ELASTIC POLE RESIDUE (MeV)$

IMR	8.0 6.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $N(2100) PARTIAL DECAY MODES$

		DECAY MASSES
P1	$N(2100) \rightarrow N\pi$	938- 140

### $N(2100) BRANCHING RATIOS$

$N(2100) \rightarrow (N\pi)/total$		(P1)	
R1	0.10 0.04	HOEHLER 79 IPWA	PI N TO PI N
R1	0.12 0.03	CUTKOSKY 80 IPWA	PI N TO PI N

### REFERENCES FOR $N(2100)$

HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL. 12-1 (KARL)IJP  
 KOCH 80 TORONTO CONF 3 -KAISER, KOCH, PIETARINEN (KARL)IJP  
 CUTKOSKY 80 TORONTO CONF 19 +FORSYTH, BARCOCK, KELLY, HENDRICK (CARN+LBL)IJP

### $N(2190) G_{17}$

$I(J^P) = \frac{1}{2}(\frac{7}{2}^-)$  Status: \*\*\*\*

MOST OF THE RESULTS PUBLISHED BEFORE 1975 ARE NOW OBSOLETE AND HAVE BEEN OMITTED. THEY MAY BE FOUND IN OUR 1982 EDITION (PHYSICS LETTERS 1118).

### $N(2190) MASS (MeV)$

M	(2117.0)	BARBOUR 78 DPWA	PI-N PHOTOPROD.
M	2140.0 40.0	HENDRY 78 MPWA	PI N TO PI N
M	(2140.0)	BAKER 79 DPWA	0 PI- P TO ETA N
M	(2150.0) (100.0)	CUTKOSKY 79 IPWA	PI N TO PI N
M	2140.0 12.0	HOEHLER 79 IPWA	PI N TO PI N
M	(2098.0)	CRAWFORD 80 DPWA	PI N PHOTOPROD.
M	2200.0 70.0	CUTKOSKY 80 IPWA	PI N TO PI N
M	(2180.0)	SAXON 80 DPWA	0 PI- P TO K LAM

### $N(2190) WIDTH (MeV)$

W	(220.0)	BARBOUR 78 DPWA	PI-N PHOTOPROD.
W	270.0 50.0	HENDRY 78 MPWA	PI N TO PI N
W	(319.0)	BAKER 79 DPWA	0 PI- P TO ETA N
W	(300.0) (100.0)	CUTKOSKY 79 IPWA	PI N TO PI N
W	390.0 30.0	HOEHLER 79 IPWA	PI N TO PI N
W	(238.0)	CRAWFORD 80 DPWA	PI N PHOTOPROD.
W	500.0 150.0	CUTKOSKY 80 IPWA	PI N TO PI N
W	(80.0)	SAXON 80 DPWA	0 PI- P TO K LAM

### $N(2190) REAL PART OF POLE POSITION (MeV)$

RE	(2111.0)	CUTKOSKY 79 IPWA	PI N TO PI N
RE	2100.0 50.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $N(2190) - 2^*IMAG PART OF POLE POSITION (MeV)$

IM	(308.0)	CUTKOSKY 79 IPWA	PI N TO PI N
IM	400.0 160.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $N(2190) REAL PART OF ELASTIC POLE RESIDUE (MeV)$

RER	(24.0)	CUTKOSKY 79 IPWA	PI N TO PI N
RER	22.0 14.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $N(2190) IMAG PART OF ELASTIC POLE RESIDUE (MeV)$

IMR	(-12.0)	CUTKOSKY 79 IPWA	PI N TO PI N
IMR	-13.0 20.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $N(2190) PARTIAL DECAY MODES$

		DECAY MASSES
P1	$N(2190) \rightarrow N\pi$	938+ 140
P2	$N(2190) \rightarrow N\eta$	940+ 549
P3	$N(2190) \rightarrow \Delta K$	1116+ 494
P4	$N(2190) \rightarrow \Sigma K$	1189+ 494
P5	$N(2190) \rightarrow N\pi\pi$	938+ 140+ 140
P6	$N(2190) \rightarrow N\rho, S=3/2, D\text{-wave}$	938+ 769
P7	$N(2190) \rightarrow p\gamma, helicity=1/2$	938+ 0
P8	$N(2190) \rightarrow p\gamma, helicity=3/2$	938+ 0
P9	$N(2190) \rightarrow n\gamma, helicity=1/2$	940- 0
P10	$N(2190) \rightarrow n\gamma, helicity=3/2$	940- 0

### $N(2190) BRANCHING RATIOS$

$N(2190) \rightarrow (N\pi)/total$		(P1)	
R1	0.16 0.04	HENDRY 78 MPWA	PI N TO PI N
R1	(0.16) (0.07)	CUTKOSKY 79 IPWA	PI N TO PI N
R1	0.14 0.02	HOEHLER 79 IPWA	PI N TO PI N
R1	0.12 0.06	CUTKOSKY 80 IPWA	PI N TO PI N







# Baryon Full Listings

$N(2700)$ ,  $N(\sim 3000)$ ,  $\Delta(1232)$

## $N(2700)$ PARTIAL DECAY MODES

	DECAY MASSES	
$P1$	$N(2700) \rightarrow N\pi$	938+ 140

## $N(2700)$ BRANCHING RATIOS

$N(2700) \rightarrow (N\pi)/total$				(P1)
R1	0.07	0.02	HENDRY 78 MPWA	PI N TO PI N
R1'	0.04	0.01	HOEHLER 79 IPWA	PI N TO PI N

## REFERENCES FOR $N(2700)$

HENDRY 78 PRL 41 222 A W HENDRY (IND+LBL)JIP  
 -- THE ANALYSIS AND RESULTS ARE DISCUSSED MORE FULLY IN HENDRY 81.  
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1  
 +KAISER, KOCH, PIETARINEN (KARL)JIP  
 ALSO 80 TORONTO CONF 3 R KOCH (KARL)JIP  
 HENDRY 81 ANP 136 1 A W HENDRY (IND)

## $\sim 3000$ MeV REGION - $I=1/2$ FORMATION EXPERIMENTS

WE LIST HERE MISCELLANEOUS HIGH-MASS CANDIDATES FOR ISOSPIN-1/2 RESONANCES FOUND IN PARTIAL-WAVE ANALYSES. SO FAR, NO ANALYSIS OF THIS REGION HAS USED ALL THE AVAILABLE DATA OR INCORPORATED ANALYTICITY CONSTRAINTS.

OUR 1982 EDITION ALSO HAD AN  $N(3030)$ , AN  $N(3245)$ , AN  $N(3690)$ , AND AN  $N(3755)$ . NOTHING HAS BEEN HEARD FROM THEM SINCE THE 1960'S, AND UNDER THE AUTHORITY GRANTED UNTO US BY THE STATUTE OF LIMITATIONS WE DECLARE THEM TO BE DEAD. THE LAST THREE WERE NARROW PEAKS SEEN IN PRODUCTION EXPERIMENTS. THE  $N(3030)$  WAS DEDUCED FROM TOTAL-CROSS-SECTION AND 180-DEG-ELASTIC-CROSS-SECTION MEASUREMENTS; PLACED IN THE BARYON SUMMARY TABLE IN THE ANYTHING-GOES 1960'S, IT REMAINED THERE DUE TO INATTENTION UNTIL THE 1984 EDITION.

## $N(\sim 3000)$ MASS (MeV)

M	3500.0	200.0	HENDRY 78 MPWA	PI N L115
M	3800.0	200.0	HENDRY 78 MPWA	PI N M117
M	4100.0	200.0	HENDRY 78 MPWA	PI N M119
M	(2600.0)		KOCH 80 IPWA	PI N D13
M	(3100.0)		KOCH 80 IPWA	PI N L115
M	(3500.0)		KOCH 80 IPWA	PI N M117
M	3500.0	TO 4000.0	KOCH 80 IPWA	PI N M119

## $N(\sim 3000)$ WIDTH (MeV)

W	1500.0	200.0	HENDRY 78 MPWA	PI N L115
W	1600.0	200.0	HENDRY 78 MPWA	PI N M117
W	1900.0	300.0	HENDRY 78 MPWA	PI N M119

## $N(\sim 3000)$ PARTIAL DECAY MODES

	DECAY MASSES	
$P1$	$N(\sim 3000) \rightarrow N\pi$	938+ 140

## $N(\sim 3000)$ BRANCHING RATIOS

$N(\sim 3000) \rightarrow (N\pi)/total$				(P1)
R1	0.055	0.02	HENDRY 78 MPWA	PI N L115
R1	0.040	0.015	HENDRY 78 MPWA	PI N M117
R1	0.030	0.015	HENDRY 78 MPWA	PI N M119

## REFERENCES FOR $N(\sim 3000)$

HENDRY 78 PRL 41 222 A W HENDRY (IND+LBL)JIP  
 -- THE ANALYSIS AND RESULTS ARE DISCUSSED MORE FULLY IN HENDRY 81.  
 KOCH 80 TORONTO CONF 3 R KOCH (KARL)JIP  
 HENDRY 81 ANP 136 1 A W HENDRY (IND)JIP

## $\Delta$ BARYONS ( $S=0, I=3/2$ )

### $\Delta(1232) P_{33}$

$$I(J^P) = \frac{3}{2}(\frac{3}{2}^+)$$

Status: \*\*\*\*

MOST OF THE RESULTS PUBLISHED BEFORE 1977 ARE NOW OBSOLETE AND HAVE BEEN OMITTED. THEY MAY BE FOUND IN OUR 1982 EDITION (PHYSICS LETTERS 111B).

IN ADDITION, RESULTS IN THIS REGION FROM PRODUCTION EXPERIMENTS, WHICH USED TO BE LISTED SEPARATELY AS THE NEXT ENTRY, HAVE BEEN ENTIRELY REMOVED. THEY TOO MAY BE FOUND IN OUR 1982 EDITION.

## $\Delta(1232)$ MASS (MeV)

M	1233.0	2.0	HOEHLER 79 IPWA	PI N TO PI N
	1232.0	3.0	CUTKOSKY 80 IPWA	PI N TO PI N
M++	1231.1	0.2	PEDRONI 78 DPWA ++	PI N 70-370 MEV
M++	1230.9	0.3	KOCH 80 IPWA ++	PI N TO PI N
M++	1230.6	0.2	ZIDELL 80 DPWA ++	PI N 0-350 MEV
M+	(1231.8)		BERENDS 75 IPWA +	GAM P TO PI NUC
M+	(1231.2)		BARBOUR 78 DPWA	PI-N PHOTOPROD.
M+	1234.9	1.4	MIROSHNIC 79 +	FIT PHOTOPROD.
M+	(1231.6)		CRAWFORD 80 DPWA	PI N PHOTOPROD.
M0	1233.8	0.2	PEDRONI 78 DPWA 0	PI N 70-370 MEV
M0	1233.6	0.5	KOCH 80 IPWA 0	PI N TO PI N
M0	1232.5	0.3	ZIDELL 80 DPWA 0	PI N 0-350 MEV

## $\Delta(1232)$ WIDTH (MeV)

W	116.0	5.0	HOEHLER 79 IPWA	PI N TO PI N
	120.0	5.0	CUTKOSKY 80 IPWA	PI N TO PI N
W++	111.3	0.5	PEDRONI 78 DPWA ++	PI N 70-370 MEV
W++	111.0	1.0	KOCH 80 IPWA ++	PI N TO PI N
W++	113.2	0.3	ZIDELL 80 DPWA ++	PI N 0-350 MEV
W+	(111.0)		BARBOUR 78 DPWA	PI-N PHOTOPROD.
W+	151.1	2.4	MIROSHNIC 79 +	FIT PHOTOPROD.
W+	(111.2)		CRAWFORD 80 DPWA	PI N PHOTOPROD.
W0	117.9	0.9	PEDRONI 78 DPWA 0	PI N 70-370 MEV
W0	113.0	1.5	KOCH 80 IPWA 0	PI N TO PI N
W0	121.3	0.4	ZIDELL 80 DPWA 0	PI N 0-350 MEV

## $\Delta(1232)^0 - \Delta(1232)^{++}$ MASS DIFFERENCE (MeV)

D A (2.7) (0.3) PEDRONI 78 DPWA PI N 70-370 MEV  
 D A REDUNDANT WITH DATA IN MASS LISTING.  
 D A USING PI--D AS WELL, PEDRONI 78 DETERMINE (M- - M++)/(M0 - M++)/3 =  
 D A 4.6+-0.2 MEV.

## $\Delta(1232)^0 - \Delta(1232)^{++}$ WIDTH DIFFERENCE (MeV)

WD A (6.6) (1.0) PEDRONI 78 DPWA PI N 70-370 MEV  
 WD A REDUNDANT WITH DATA IN MASS LISTING.

## $\Delta(1232)$ REAL PART OF POLE POSITION (MeV)

RE	1210.0	1.0	CUTKOSKY 80 IPWA	PI N TO PI N
RE	(1210.)		ARNDT 85 DPWA	PI N TO PI N
R++ B	1209.6	0.5	VASAN 76	++ FIT CARTER 73
R++ B	FROM FITS TO COULOMB-BARRIER-CORRECTED CARTER 73			PHASE SHIFT.
R++ C	1210.5	TO 1210.8	VASAN 76	++ FIT CARTER 73
R++ C	FROM FITS TO CARTER 73 NUCLEAR PHASE SHIFT WITHOUT COULOMB BARRIER			
R++ C	CORRECTIONS.			
R++ D	(1210.4)	(0.17)	ZIDELL 78	++ FIT ZIDELL 78
R++ D	FIT TO ZIDELL 78 NUCLEAR PHASE SHIFT WITHOUT COULOMB			
R++ D	BARRIER CORRECTIONS.			
R++ E	1210.70	0.16	ZIDELL 80 DPWA ++	PI N 0-350 MEV
R++ E	FIT TO ZIDELL 80 NUCLEAR PHASE SHIFTS. THE ACCURACY CLAIMED ON THE			
R++ E	REAL PART IS CONSIDERABLY BETTER THAN IS ALLOWED BY UNCERTAINTIES			
R++ E	IN THE BEAM MOMENTUM.			
RE+	1208.0	2.0	CAMPBELL 76	- FIT PHOTOPROD.
RE+	1206.9+-0.9	TO 1210.5+-1.8	MIROSHNIC 79	+ FIT PHOTOPROD.
RE0 B	1210.75	0.6	VASAN 76	0 FIT CARTER 73
RE0 C	(1210.2)		VASAN 76	0 FIT CARTER 73
RE0 D	(1209.5)	(0.41)	ZIDELL 78	0 FIT ZIDELL 78
RE0 E	1210.50	0.36	ZIDELL 80 DPWA	0 PI N 0-350 MEV

## $\Delta(1232)$ -IMAG PART OF POLE POSITION (MeV)

IM	50.0	1.0	CUTKOSKY 80 IPWA	PI N TO PI N
IM	(50.)		ARNDT 85 DPWA	PI N TO PI N
I++ B	50.4	0.5	VASAN 76	++ FIT CARTER 73
I++ B	49.9	TO 50.0	VASAN 76	++ FIT CARTER 73
I++ D	(49.745)	(0.14)	ZIDELL 78	++ FIT ZIDELL 78
I++ E	49.61	0.12	ZIDELL 80 DPWA ++	PI N 0-350 MEV





















# Baryon Full Listings

$\Delta(2200)$ ,  $\Delta(2300)$ ,  $\Delta(2350)$

### $\Delta(2200)$ REAL PART OF POLE POSITION (MeV)

RE	(2094.0)		CUTKOSKY 79 IPWA	PI N TO PI N
RE	2100.0	50.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $\Delta(2200)$ -2\*IMAG PART OF POLE POSITION (MeV)

IM	(294.0)		CUTKOSKY 79 IPWA	PI N TO PI N
IM	340.0	80.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $\Delta(2200)$ REAL PART OF ELASTIC POLE RESIDUE (MeV)

RER	(2.0)		CUTKOSKY 79 IPWA	PI N TO PI N
RER	3.0	5.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $\Delta(2200)$ IMAG PART OF ELASTIC POLE RESIDUE (MeV)

IMR	(-7.0)		CUTKOSKY 79 IPWA	PI N TO PI N
IMR	-8.0	3.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $\Delta(2200)$ PARTIAL DECAY MODES

P1	$\Delta(2200) \rightarrow N\pi$	DECAY MASSES
		938- 140

### $\Delta(2200)$ BRANCHING RATIOS

$\Delta(2200) \rightarrow (N\pi)/total$ (P1)				
R1	0.09	0.02	HENDRY 78 MPWA	PI N TO PI N
R1	(0.05)		CUTKOSKY 79 IPWA	PI N TO PI N
R1	0.05	0.02	HOEHLER 79 IPWA	PI N TO PI N
R1	0.06	0.02	CUTKOSKY 80 IPWA	PI N TO PI N

### REFERENCES FOR $\Delta(2200)$

HENDRY 78 PRL 41 222 A W HENDRY (IND+LBL)JP  
 -- THE ANALYSIS AND RESULTS ARE DISCUSSED MORE FULLY IN HENDRY 81.  
 CUTKOSKY 79 PR D20 2839 +FORSYTH,HENDRICK,KELLY (CARN+LBL)JP  
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1 (KARL)JP  
 +KAISER,KOCH,PIETARINEN  
 ALSO 80 TORONTO CONF 3 R KOCH (KARL)JP  
 CUTKOSKY 80 TORONTO CONF 19 +FORSYTH,BABCOCK,KELLY,HENDRICK (CARN+LBL)JP  
 HENDRY 81 ANP 136 1 A W HENDRY (IND)

## $\Delta(2300) H_{39}$

$I(J^P) = \frac{3}{2}(\frac{9}{2}^-)$  Status: \*\*

OMITTED FROM SUMMARY TABLE

### $\Delta(2300)$ MASS (MeV)

M	2450.0	100.0	HENDRY 78 MPWA	PI N TO PI N
M	2217.0	80.0	HOEHLER 79 IPWA	PI N TO PI N
M	(2204.5)	(3.4)	CHEW 80 BPWA --	PI-P TO PI-P
M	2400.0	125.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $\Delta(2300)$ WIDTH (MeV)

W	500.0	200.0	HENDRY 78 MPWA	PI N TO PI N
W	300.0	100.0	HOEHLER 79 IPWA	PI N TO PI N
W	(32.3)	(1.0)	CHEW 80 BPWA ++	PI-P TO PI+P
W	425.0	150.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $\Delta(2300)$ REAL PART OF POLE POSITION (MeV)

RE	2370.0	80.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $\Delta(2300)$ -2\*IMAG PART OF POLE POSITION (MeV)

IM	420.0	160.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $\Delta(2300)$ REAL PART OF ELASTIC POLE RESIDUE (MeV)

RER	9.0	4.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $\Delta(2300)$ IMAG PART OF ELASTIC POLE RESIDUE (MeV)

IMR	-3.0	5.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $\Delta(2300)$ PARTIAL DECAY MODES

P1	$\Delta(2300) \rightarrow N\pi$	DECAY MASSES
		938- 140

### $\Delta(2300)$ BRANCHING RATIOS

$\Delta(2300) \rightarrow (N\pi)/total$ (P1)				
R1	0.08	0.02	HENDRY 78 MPWA	PI N TO PI N
R1	0.03	0.02	HOEHLER 79 IPWA	PI N TO PI N
R1	(0.05)		CHEW 80 BPWA ++	PI+P TO PI+P
R1	0.06	0.02	CUTKOSKY 80 IPWA	PI N TO PI N

### REFERENCES FOR $\Delta(2300)$

HENDRY 78 PRL 41 222 A W HENDRY (IND+LBL)JP  
 -- THE ANALYSIS AND RESULTS ARE DISCUSSED MORE FULLY IN HENDRY 81.  
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1 (KARL)JP  
 +KAISER,KOCH,PIETARINEN  
 ALSO 80 TORONTO CONF 3 R KOCH (KARL)JP  
 CHEW 80 TORONTO CONF 123 D M CHEW (LBL)JP  
 CUTKOSKY 80 TORONTO CONF 19 +FORSYTH,BABCOCK,KELLY,HENDRICK (CARN+LBL)JP  
 HENDRY 81 ANP 136 1 A W HENDRY (IND)

## $\Delta(2350) D_{35}$

$I(J^P) = \frac{3}{2}(\frac{5}{2}^-)$  Status: \*

OMITTED FROM SUMMARY TABLE

### $\Delta(2350)$ MASS (MeV)

M	2305.0	26.0	HOEHLER 79 IPWA	PI N TO PI N
M	2400.0	125.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $\Delta(2350)$ WIDTH (MeV)

W	300.0	70.0	HOEHLER 79 IPWA	PI N TO PI N
W	400.0	150.0	CUTKOSKY 80 IPWA	PI N TO PI N

### $\Delta(2350)$ REAL PART OF POLE POSITION (MeV)

RE	2400.0	125.0	CUTKOSKY 80 IPWA	PI N TO PI N
----	--------	-------	------------------	--------------

### $\Delta(2350)$ -2\*IMAG PART OF POLE POSITION (MeV)

IM	400.0	150.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $\Delta(2350)$ REAL PART OF ELASTIC POLE RESIDUE (MeV)

RER	5.0	17.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $\Delta(2350)$ IMAG PART OF ELASTIC POLE RESIDUE (MeV)

IMR	-14.0	10.0	CUTKOSKY 80 IPWA	PI N TO PI N
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### $\Delta(2350)$ PARTIAL DECAY MODES

P1	$\Delta(2350) \rightarrow N\pi$	DECAY MASSES
		938- 140

### $\Delta(2350)$ BRANCHING RATIOS

$\Delta(2350) \rightarrow (N\pi)/total$ (P1)				
R1	0.04	0.02	HOEHLER 79 IPWA	PI N TO PI N
R1	0.20	0.10	CUTKOSKY 80 IPWA	PI N TO PI N

### REFERENCES FOR $\Delta(2350)$

HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1 (KARL)JP  
 +KAISER,KOCH,PIETARINEN  
 ALSO 80 TORONTO CONF 3 R KOCH (KARL)JP  
 CUTKOSKY 80 TORONTO CONF 19 +FORSYTH,BABCOCK,KELLY,HENDRICK (CARN+LBL)JP



# Baryon Full Listings

$\Delta(2420)$ ,  $\Delta(2750)$ ,  $\Delta(2950)$ ,  $\Delta(\sim 3000)$

### $\Delta(2420)$ BRANCHING RATIOS

$\Delta(2420) \rightarrow (N \pi)/total$		(P1)	
R1	0.11 0.02	HENDRY	78 MPWA PI N TO PI N
R1	0.08 0.015	HOEHLER	79 IPWA PI N TO PI N
R1	(0.22)	CHEW	80 BPWA + PI+P TO PI+P
R1	0.08 0.03	CUTKOSKY	80 IPWA PI N TO PI N

### REFERENCES FOR $\Delta(2420)$

FOR EARLY REFERENCES, SEE PHYSICS LETTERS 111B (1982).

HENDRY 78 PRL 41 222 A W HENDRY (IND-LBL)JJP  
 -- THE ANALYSIS AND RESULTS ARE DISCUSSED MORE FULLY IN HENDRY 81.  
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1  
 +KAISER,KOCH,PIETARINEN (KARL)JJP  
 ALSO 80 TORONTO CONF 3 R KOCH (KARL)JJP  
 CHEW 80 TORONTO CONF 123 D M CHEW (LBL)JJP  
 CUTKOSKY 80 TORONTO CONF 19 -FORSYTH, BARCOCK,KELLY,HENDRICK (CARN-LBL)JJP  
 HENDRY 81 ANP 136 1 A W HENDRY (IND)

**$\Delta(2750)$   $I_{313}$**

$I(J^P) = \frac{3}{2}(\frac{13}{2}^-)$  Status: \*\*

OMITTED FROM SUMMARY TABLE

### $\Delta(2750)$ MASS (MeV)

M	2650.0	100.0	HENDRY	78 MPWA	PI N TO PI N
M	2794.0	80.0	HOEHLER	79 IPWA	PI N TO PI N

### $\Delta(2750)$ WIDTH (MeV)

W	500.0	100.0	HENDRY	78 MPWA	PI N TO PI N
W	350.0	100.0	HOEHLER	79 IPWA	PI N TO PI N

### $\Delta(2750)$ PARTIAL DECAY MODES

$\Delta(2750) \rightarrow N \pi$		DECAY MASSES	
P1		938-	140

### $\Delta(2750)$ BRANCHING RATIOS

$\Delta(2750) \rightarrow (N \pi)/total$		(P1)	
R1	0.05 0.01	HENDRY	78 MPWA PI N TO PI N
R1	0.04 0.015	HOEHLER	79 IPWA PI N TO PI N

### REFERENCES FOR $\Delta(2750)$

HENDRY 78 PRL 41 222 A W HENDRY (IND-LBL)JJP  
 -- THE ANALYSIS AND RESULTS ARE DISCUSSED MORE FULLY IN HENDRY 81.  
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1  
 +KAISER,KOCH,PIETARINEN (KARL)JJP  
 ALSO 80 TORONTO CONF 3 R KOCH (KARL)JJP  
 HENDRY 81 ANP 136 1 A W HENDRY (IND)

**$\Delta(2950)$   $K_{315}$**

$I(J^P) = \frac{3}{2}(\frac{15}{2}^-)$  Status: \*\*

OMITTED FROM SUMMARY TABLE

### $\Delta(2950)$ MASS (MeV)

M	2850.0	100.0	HENDRY	78 MPWA	PI N TO PI N
M	2990.0	100.0	HOEHLER	79 IPWA	PI N TO PI N

### $\Delta(2950)$ WIDTH (MeV)

W	700.0	200.0	HENDRY	78 MPWA	PI N TO PI N
W	350.0	100.0	HOEHLER	79 IPWA	PI N TO PI N

### $\Delta(2950)$ PARTIAL DECAY MODES

$\Delta(2950) \rightarrow N \pi$		DECAY MASSES	
P1		938-	140

### $\Delta(2950)$ BRANCHING RATIOS

$\Delta(2950) \rightarrow (N \pi)/total$		(P1)	
R1	0.03 0.01	HENDRY	78 MPWA PI N TO PI N
R1	0.04 0.02	HOEHLER	79 IPWA PI N TO PI N

### REFERENCES FOR $\Delta(2950)$

HENDRY 78 PRL 41 222 A W HENDRY (IND-LBL)JJP  
 -- THE ANALYSIS AND RESULTS ARE DISCUSSED MORE FULLY IN HENDRY 81.  
 HOEHLER 79 HANDBOOK OF PI-N SCATTERING, PHYSIK DATEN VOL.12-1  
 +KAISER,KOCH,PIETARINEN (KARL)JJP  
 ALSO 80 TORONTO CONF 3 R KOCH (KARL)JJP  
 HENDRY 81 ANP 136 1 A W HENDRY (IND)

### $\sim 3000$ MeV REGION - $I=3/2$ FORMATION EXPERIMENTS

WE LIST HERE MISCELLANEOUS HIGH-MASS CANDIDATES FOR ISOSPIN-3/2 RESONANCES FOUND IN PARTIAL-WAVE ANALYSES. SO FAR, NO ANALYSIS OF THIS REGION HAS USED ALL THE AVAILABLE DATA OR INCORPORATED ANALYTICITY CONSTRAINTS.

OUR 1982 EDITION ALSO HAD A DELTA(2850) AND A DELTA(3230). NOTHING HAS BEEN HEARD FROM THEM IN 10 YEARS, AND UNDER THE AUTHORITY GRANTED UNTO US BY THE STATUTE OF LIMITATIONS, WE DECLARE THEM TO BE DEAD. THE EVIDENCE FOR THEM WAS DEDUCED FROM TOTAL-CROSS-SECTION AND 180-DEG-ELASTIC-CROSS-SECTION MEASUREMENTS. PLACED IN THE BARYON SUMMARY TABLE IN THE ANYTHING-GOES 1960'S, THEY REMAINED THERE DUE TO INATTENTION UNTIL THE 1984 EDITION.

### $\Delta(\sim 3000)$ MASS (MeV)

M	2850.0	150.0	HENDRY	78 MPWA	PI N 1311
M	3200.0	200.0	HENDRY	78 MPWA	PI N K313
M	3300.0	200.0	HENDRY	78 MPWA	PI N L317
M	3700.0	200.0	HENDRY	78 MPWA	PI N M319
M	4100.0	300.0	HENDRY	78 MPWA	PI N N321
M	(3300.0)		KOCH	80 IPWA	PI N L317
M	(3500.0)		KOCH	80 IPWA	PI N M315
M	IN ADDITION, KOCH 80 REPORT SOME EVIDENCE FOR AN S31 DELTA(2700) AND A P33 DELTA(2800).				

### $\Delta(\sim 3000)$ WIDTH (MeV)

W	700.0	200.0	HENDRY	78 MPWA	PI N 1311
W	1000.0	300.0	HENDRY	78 MPWA	PI N K313
W	1100.0	300.0	HENDRY	78 MPWA	PI N L317
W	1300.0	400.0	HENDRY	78 MPWA	PI N M319
W	1600.0	500.0	HENDRY	78 MPWA	PI N N321

### $\Delta(\sim 3000)$ PARTIAL DECAY MODES

$\Delta(\sim 3000) \rightarrow N \pi$		DECAY MASSES	
P1		938-	140

### $\Delta(\sim 3000)$ BRANCHING RATIOS

$\Delta(\sim 3000) \rightarrow (N \pi)/total$		(P1)	
R1	0.06 0.02	HENDRY	78 MPWA PI N 1311
R1	0.045 0.02	HENDRY	78 MPWA PI N K313
R1	0.05 0.01	HENDRY	78 MPWA PI N L317
R1	0.025 0.01	HENDRY	78 MPWA PI N M319
R1	0.018 0.01	HENDRY	78 MPWA PI N N321

### REFERENCES FOR $\Delta(\sim 3000)$

HENDRY 78 PRL 41 222 A W HENDRY (IND-LBL)JJP  
 -- THE ANALYSIS AND RESULTS ARE DISCUSSED MORE FULLY IN HENDRY 81.  
 KOCH 80 TORONTO CONF 3 R KOCH (KARL)JJP  
 HENDRY 81 ANP 136 1 A W HENDRY (IND)







For notation, see key on page 91.

# Baryon Full Listings

$Z_1(2150)$ ,  $Z_1(2500)$ ,  $\Lambda$ 's and  $\Sigma$ 's

### REFERENCES FOR $Z_1(2150)$

ABRAMS 70 PR D1 1917 +COOL,GIACOMELLI,KYCIA,LEONTIC,LI + (BNL)  
 ALSO 67 PRL 19 257 ABRAMS,COOL,GIACOMELLI,KYCIA,LEONTIC- (BNL)  
 ARNDT 85 PR D31 2250 ARNDT,ROPER (VPI)

**$Z_1(2500)$  BUMPS**  $I(J^P) = 1( )$  Status: \*

OMITTED FROM SUMMARY TABLE

A SMALL BUMP IN THE TOTAL CROSS SECTION AT 2.7 GEV/C.

### $Z_1(2500)$ MASS (MeV)

M 2500.0 20.0 ABRAMS 70 CNTR ++ K+P TOTAL

### $Z_1(2500)$ WIDTH (MeV)

W (160.0) ABRAMS 70 CNTR -- K+P TOTAL

### $Z_1(2500)$ PARTIAL DECAY MODES

P1  $Z_1(2500) \rightarrow NK$  DECAY MASSES 938+ 494

### $Z_1(2500)$ BRANCHING RATIOS

$Z_1(2500) \rightarrow (NK)/total$  (P1)  
 R1 J IS NOT KNOWN, THE FOLLOWING IS (J-1/2)\*P1  
 R1 (0.03) ABRAMS 70 CNTR ++ K+P TOTAL

### REFERENCES FOR $Z_1(2500)$

ABRAMS 70 PR D1 1917 +COOL,GIACOMELLI,KYCIA,LEONTIC,LI + (BNL)  
 ALSO 67 PRL 19 257 ABRAMS,COOL,GIACOMELLI,KYCIA,LEONTIC- (BNL)

Table 1. The status of the  $\Lambda$  and  $\Sigma$  resonances. Only those with an overall status of \*\*\* or \*\*\*\* are included in the main Baryon Table.

Particle	$L_f, 2J$	Overall status	Status as seen in --			
			$N\bar{K}$	$\Lambda\pi$	$\Sigma\pi$	Other channels
$\Lambda(1116)$	$P_{01}$	****				$N\pi$ (weakly)
$\Lambda(1405)$	$S_{01}$	****	****	F	****	
$\Lambda(1520)$	$D_{03}$	****	****	o	****	$\Lambda\pi\pi, \Lambda\gamma$
$\Lambda(1600)$	$P_{01}$	***	***	r	**	
$\Lambda(1670)$	$S_{01}$	****	****	b	****	$\Lambda\eta$
$\Lambda(1690)$	$D_{03}$	****	****	i	****	$\Lambda\pi\pi, \Sigma\pi\pi$
$\Lambda(1800)$	$S_{01}$	***	***	d	**	$N\bar{K}^*, \Sigma(1385)\pi$
$\Lambda(1800)$	$P_{01}$	***	***	d	**	$N\bar{K}^*$
$\Lambda(1820)$	$F_{05}$	****	****	e	****	$\Sigma(1385)\pi$
$\Lambda(1830)$	$D_{05}$	****	****	n	****	$\Sigma(1385)\pi$
$\Lambda(1890)$	$F_{03}$	****	****	F	**	$N\bar{K}^*, \Sigma(1385)\pi$
$\Lambda(2000)$		*	*	o	*	$\Lambda\omega, N\bar{K}^*$
$\Lambda(2020)$	$F_{07}$	*	*	r	*	
$\Lambda(2100)$	$G_{07}$	****	****	b	***	$\Lambda\omega, N\bar{K}^*$
$\Lambda(2110)$	$F_{05}$	***	**	i	*	$\Lambda\omega, N\bar{K}^*$
$\Lambda(2325)$	$D_{03}$	*	*	d	*	$\Lambda\omega$
$\Lambda(2350)$		***	***	d	*	
$\Lambda(2585)$		**	**	e	n	
$\Sigma(1193)$	$P_{11}$	****				$N\pi$ (weakly)
$\Sigma(1385)$	$P_{13}$	****	****	****	****	
$\Sigma(1480)$		*	*	*	*	
$\Sigma(1560)$		**	**	**	**	
$\Sigma(1580)$	$D_{13}$	**	*	*	*	
$\Sigma(1620)$	$S_{11}$	**	**	*	*	
$\Sigma(1660)$	$P_{11}$	****	****	*	**	
$\Sigma(1670)$	$D_{13}$	****	****	****	****	several others
$\Sigma(1690)$		**	*	**	*	$\Lambda\pi\pi$
$\Sigma(1750)$	$S_{11}$	***	***	**	*	$\Sigma\eta$
$\Sigma(1770)$	$P_{11}$	*				
$\Sigma(1775)$	$D_{15}$	****	****	****	****	several others
$\Sigma(1840)$	$P_{13}$	*	*	**	*	
$\Sigma(1880)$	$P_{11}$	**	**	**	**	$N\bar{K}^*$
$\Sigma(1915)$	$F_{15}$	****	****	****	****	$\Sigma(1385)\pi$
$\Sigma(1940)$	$D_{13}$	****	*	****	**	quasi-2-body
$\Sigma(2000)$	$S_{13}$	*	*	*	*	$N\bar{K}^*, \Lambda(1520)\pi$
$\Sigma(2030)$	$F_{17}$	****	****	****	**	several others
$\Sigma(2070)$	$F_{17}$	*	*	*	*	
$\Sigma(2080)$	$P_{15}$	**	*	**	*	
$\Sigma(2100)$	$P_{13}$	*	*	*	*	
$\Sigma(2250)$	$G_{17}$	****	***	*	*	
$\Sigma(2455)$		**	*	*	*	
$\Sigma(2620)$		**	*	*	*	
$\Sigma(3000)$		*	*	*	*	
$\Sigma(3170)$		*	*	*	*	multi-body

\*\*\*\* Good, clear, and unmistakable.  
 \*\*\* Good, but in need of clarification or not absolutely certain.  
 \*\* Not established; needs confirmation.  
 \* Evidence weak; could disappear.

## NOTE ON $\Lambda$ AND $\Sigma$ RESONANCES

### I. Introduction

There is actually a new result or two on  $\Lambda$  and  $\Sigma$  resonances for this edition. There are new results on the  $\Sigma(1385)$  and  $\Lambda(1405)$ , leftovers from the old 4.2 GeV/c  $K^-p$  experiment.<sup>1</sup> There are two modest additions to the set of scattering data, mentioned in Sec. II below. There is some recent work on the nature of the low-energy  $\bar{K}N$  system bearing on the interpretation of the  $\Lambda(1405)$ , and this difficult object is also being got at by studying radiative annihilation of the  $K^-$  hydrogen atom (see the Listings for references). Generally, however, the field remains at a standstill. It can only be revived if one of the lower energy accelerators becomes a kaon factory.

Table 1 is an attempt to evaluate the status, both overall and channel by channel, of each  $\Lambda$  and  $\Sigma$  resonance in the full Baryon Listings; the evaluations are of course partly subjective. A blank indicates there is no evidence at all; either the relevant couplings are small or the resonance does not really exist. The main Baryon

Summary Table includes only the established resonances (overall status 3 or 4 stars). A number of the 1- and 2-star entries may eventually disappear, but there are certainly many resonances yet to be discovered underlying the established ones.

None of the  $\Lambda$ 's and  $\Sigma$ 's proposed in the last decade couple strongly to the main 2-body decay channels  $N\bar{K}$ ,  $\Lambda\pi$ , and  $\Sigma\pi$ , and thus they seldom appear in cross sections or invariant mass distributions. However, when the reactions  $\bar{K}N \rightarrow \bar{K}N$ ,  $\bar{K}N \rightarrow \Lambda\pi$ , and  $\bar{K}N \rightarrow \Sigma\pi$  are analyzed, some of the partial-wave amplitudes traverse small, more-or-less resonance-like circles. The question in each case is: Is this really a resonance, or is it an idle meander? Is the effect even real, or is it the result of imperfect data and analysis?

## Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

What follows is the review, slightly revised, from our 1984 edition: it summarizes "recent" progress and problems. (For another brief overview, see Tripp.<sup>2</sup>) In the Listings, some obsolete results, nearly all from before 1975, have been removed. This has been done only for the established  $\Lambda$ 's and  $\Sigma$ 's, where the addition of much improved data to partial-wave analyses has really made obsolete the older results. Where little new has been learned in the last decade, or where the situation is uncertain, nothing has been removed.

### II. Formation experiments

(by G.P. Gopal, Rutherford Appleton Laboratory)

Partial-wave analyses have been made mainly for the  $N\bar{K}$ ,  $\Lambda\pi$ , and  $\Sigma\pi$  channels, but there are also a few results for the  $\Xi K$ ,  $\Lambda\omega$ , and some quasi-2-body channels. Early analyses usually covered only the range of a single bubble chamber experiment. Although the amplitudes from analyses in neighboring mass ranges often did not join smoothly, they did give fairly reliable information about the strongly coupled resonances. More recent analyses have used the Breit-Wigner forms of the dominant resonances as input to provide constraints in determining the overall amplitudes and thus in learning about the less prominent resonances. Besides covering wider ranges, some of the more ambitious of the analyses at the lower energies have treated several channels simultaneously, so that unitarity constraints are automatically satisfied and only a single mass and width is obtained for each resonance.

In the mid and late 1970's, much new data became available. Results from several large  $K^-p$  bubble chamber experiments were published,<sup>3-6</sup> and other bubble chamber experiments studied  $K^-n$  reactions<sup>7</sup> and  $K_L^0 p$  reactions.<sup>8</sup> Counter experiments measured the  $K^-p \rightarrow \bar{K}^0 n$  total and differential cross sections at low energies,<sup>9</sup> the  $K^-p$  polarizations down to 1630 MeV for the first time,<sup>10</sup> the  $K^-p$  polarizations from 1700 to 1900 MeV with an order of magnitude increase in statistics,<sup>11</sup> the  $K^-n$  elastic angular distributions from 1600 to 1800 MeV<sup>12</sup> and from 1900 to 2300 MeV,<sup>13</sup> and the  $180^\circ K^-p$  and  $0^\circ \Sigma^- \pi^+$  differential cross sections from 1550 to 1900 MeV.<sup>14</sup>

More recently, there have been new measurements of  $K^-n$  elastic scattering between 1600 and 1740 MeV.<sup>15</sup> Also, new total and differential cross-section data on  $K^-p$ ,  $\bar{K}^0 n$ ,  $\Sigma^\pm \pi^\mp$ , and  $\Lambda\pi^0$  between 1437 and 1486 MeV have become available.<sup>16</sup> They clearly show the onset of  $P$ -wave amplitudes by 1450 MeV, which brings

into question analyses of low energy data that assumed only  $S$  waves were significant. Finally, there are new  $\Sigma^\pm \pi^\mp$  differential cross-section and polarization distributions in a region where data were sparse, from 1650 to 1715 MeV.<sup>17</sup>

We now compare the more recent analyses with each other and with the data. Some of the data have yet to be incorporated into any analysis.

**The  $N\bar{K}$  channel:** The most recent analysis<sup>18</sup> is an update of the old Rutherford Lab-Imperial College (RLIC 77) analysis.<sup>19</sup> As before, it is a conventional energy-dependent analysis with the added constraint that the masses and widths of the resonances had to be consistent with those determined in the inelastic channels analyzed previously —  $\Lambda\pi$ ,  $\Sigma\pi$ ,  $\Lambda(1520)\pi$ ,  $\Sigma(1385)\pi$ , and  $N\bar{K}^*$  (892). The analysis also goes closer to threshold, covering 1470 to 2170 MeV. It does not include the data from a number of the more recent experiments mentioned above. As before, angular distributions (a total of 5110 data points) were fit directly. The new amplitudes differ little from the RLIC 77 amplitudes. However, the  $K^-n$  data removed some of the uncertainties in the  $\Sigma$  resonances.

The LBL-Mt. Holyoke-CERN analysis<sup>20</sup> covers the narrower range of 1500 to 1940 MeV and also includes most of the new data. It is an energy-dependent analysis using a unitary background parametrized in terms of scattering lengths. The cusp effects at the  $\Lambda\eta$  and  $\Sigma\eta$  thresholds are included by introducing a square-root singularity in the energy variation of the widths of the appropriate resonances. This group's own high-statistics charge-exchange data<sup>9</sup> (which do not agree with bubble chamber measurements) all but kill the less well-established resonances.

The University College, London (UCL)  $K$ -matrix energy-dependent analysis<sup>21</sup> covers from 1540 to 2000 MeV. The  $N\bar{K}$  amplitudes are consistent with those of the other analyses over most of this range. However, at the low end there are major differences, due to the absence of constraints from the  $\Lambda(1520)$ , which lies just outside the range covered. The  $K^-n$  angular distributions and  $K^-p$  polarization measurements are not fit very well.

The above analyses, all below 2200 MeV, are complemented by the College de France-Saclay (CdF-S) energy-dependent analysis<sup>6</sup> covering from 2070 to 2440 MeV. Besides the conventional polynomial parametrization of the background amplitudes, also tried is a parametrization using constraints imposed by the duality hypothesis (that  $s$ -channel backgrounds come

## Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

For notation, see key on page 91.

exclusively from the  $t$ -channel Pomeron exchange amplitude). With 30 fewer free parameters, the results are consistent with the conventional approach.

**The  $\Sigma\pi$  channel:** There is very little agreement, particularly about the lower partial waves, between the two multichannel analyses.<sup>19,21</sup> The low-energy  $K_L^0 p \rightarrow \Sigma^0 \pi^+$  data<sup>8</sup> are better explained by the RLIC 77 amplitudes than by the UCL amplitudes. At the high end, there is good continuity between the RLIC 77 amplitudes and those from the single-channel analysis of the CdF-S collaboration<sup>6</sup> covering from 2070 to 2440 MeV. The  $\Lambda(1520)$  and  $\Lambda(2110)$  resonances, which lie outside the range covered by the UCL analysis, clearly provide strong constraints.

**The  $\Lambda\pi$  channel:** This isospin-1 channel has been the subject of many energy-dependent and -independent analyses (for example, RLIC 77,<sup>19</sup> UCL,<sup>21</sup> Baillon-Litchfield,<sup>22</sup> de Bellefon-Berthon,<sup>23</sup> and Van Horn<sup>24</sup>). However, even the widespread use of the method of Barrelet zeroes has not helped to resolve the  $\Sigma$  spectrum — probably because most  $\Sigma$  resonances simply do not couple strongly to the  $N\bar{K}$  initial state.

**Quasi-2-body channels:** The RLIC group has made energy-dependent analyses of the  $\Lambda(1520)\pi$ ,  $\Sigma(1385)\pi$ , and  $N\bar{K}^*(892)$  channels over the widest ranges for which data are available. The data were extracted from the appropriate 3-particle final states by making 4-

variable fits to an incoherent superposition of quasi-2-body final states and 3-particle Lorentz-invariant phase space. The quality of the fits suggests a maximum model-dependent systematic uncertainty of 10%. The  $\Lambda\omega$  channel has been analyzed from threshold to 2440 MeV by the CdF-S collaboration.<sup>6</sup>

**Sign conventions for resonance couplings:** In terms of the isospin-0 and -1 elastic scattering amplitudes  $A_0$  and  $A_1$ , the amplitude for  $K^- p \rightarrow \bar{K}^0 n$  scattering is  $\pm(A_1 - A_0)/2$ , where the sign depends on conventions used in conjunction with the Clebsch-Gordan coefficients (such as, is the baryon or the meson the “first” particle). If this reaction is partial-wave analyzed and if the overall phase is chosen so that, say, the  $D_{15} \Sigma(1775)$  amplitude at resonance points along the positive imaginary axis (points “up”), then any  $\Sigma$  at resonance will point “up” and any  $\Lambda$  at resonance will point “down” (along the negative imaginary axis). Thus the phase at resonance determines the isospin. The above ignores background amplitudes in the resonating partial waves.

That is the basic idea. In a similar but somewhat more complicated way, the phases of the  $\bar{K}N \rightarrow \Lambda\pi$  and  $\bar{K}N \rightarrow \Sigma\pi$  amplitudes for a resonating wave help determine the SU(3) multiplet to which the resonance belongs. Again, a convention has to be adopted for some overall arbitrary phases: which way is “up”? Our convention is that of Levi-Setti<sup>25</sup> and is shown in Fig. 1, which also compares experimental results with

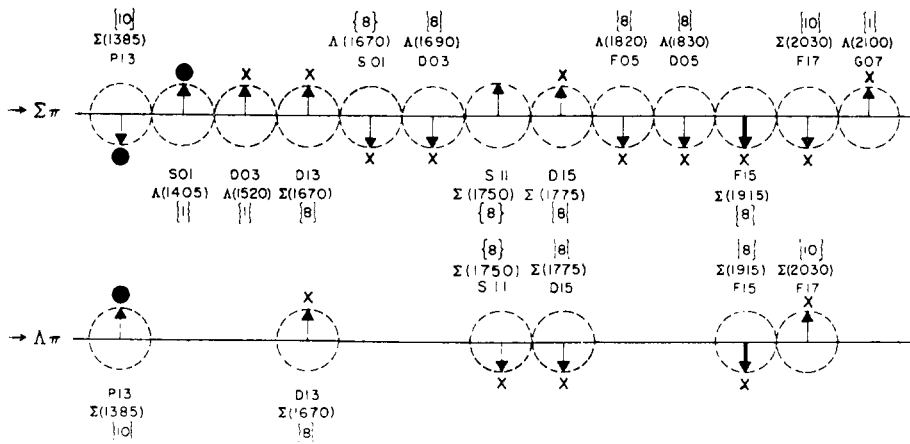


Fig. 1. The signs of the imaginary parts of resonating amplitudes in the  $\bar{K}N \rightarrow \Lambda\pi$  and  $\Sigma\pi$  channels. The signs of the  $\Sigma(1385)$  and  $\Lambda(1405)$ , marked with a  $\bullet$ , are set by convention, and then the others are determined relative to them. The signs required by the SU(3) assignments of the resonances are shown with an arrow, and the experimentally determined signs are shown with an X.

## Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

theoretical predictions for the signs of several resonances. In the Listings, a + or - sign in front of a measurement of an inelastic resonance coupling indicates the sign (the *absence* of a sign means that the sign is not determined, *not* that it is positive). For more details, see Appendix II of our 1982 edition.<sup>26</sup>

**Argand plots:** Figure 2 shows some representative Argand plots of partial-wave amplitudes. For the  $N\bar{K}$  channel we show the amplitudes from RLIC 77<sup>19</sup> and from LBL-Mt. Holyoke-CERN,<sup>20</sup> and for the  $\Lambda\pi$  and  $\Sigma\pi$  channels we show those from RLIC 77<sup>19</sup> and from UCL.<sup>21</sup>

**Errors on masses and widths:** The errors quoted on resonance parameters from partial-wave analyses are often only statistical, and the parameters can change by more than these errors when a different parametrization of the waves is used. Furthermore, the different analyses use more or less the same data, so it is not really appropriate to treat the different determinations of the resonance parameters as independent or to average them together. In any case, the spread of the masses, widths, and branching fractions from the different analyses is certainly a better indication of the uncertainties than are the quoted errors. In the Baryon Summary Table, usually a range reflecting the spread of the values is given rather than a particular value with error.

For three states, the  $\Lambda(1520)$ , the  $\Lambda(1820)$ , and the  $\Sigma(1775)$ , there is enough information to make an overall fit to the various branching fractions. It is then necessary to use the quoted errors, but the errors obtained from the fit should not be taken seriously.

### III. Production experiments

Partial-wave analyses of course separate partial waves, whereas a peak in a cross section or an invariant mass distribution usually cannot be disentangled from background and analyzed for its quantum numbers; and more than one resonance may be contributing to the peak. Results from partial-wave analyses and from production experiments are generally kept separate in the Listings, and in the Baryon Summary Table results from production experiments are used only for the low-mass states. The  $\Sigma(1385)$  and  $\Lambda(1405)$  of course lie below the  $\bar{K}N$  threshold and everything about them comes from production experiments; and production and formation experiments agree quite well in the case of  $\Lambda(1520)$  and results have been combined. There is some disagreement between production and formation experi-

ments in the 1600-1700-MeV region: see the Note on the  $\Sigma(1670)$ .

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For notation, see key on page 91.

## Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

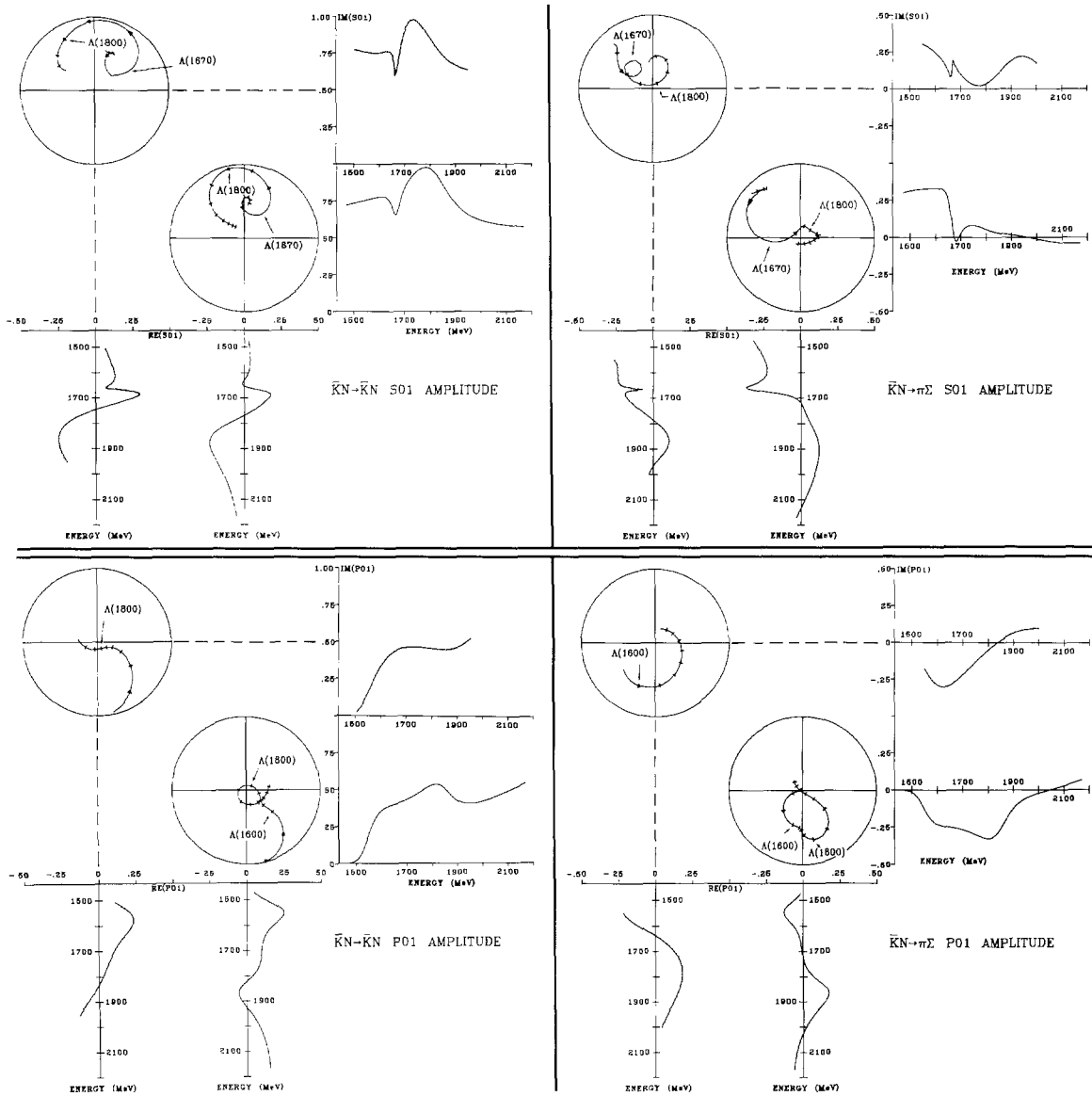


Fig. 2(a). The  $L_{I-2J} = S_{01}$  and  $P_{01}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plots for the elastic amplitudes are from ALSTON 78, and the upper plots for the  $\Sigma\pi$  amplitudes are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions [the  $S_{01}$   $\Lambda(1405)$  is of course below threshold and is not shown]. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

# Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

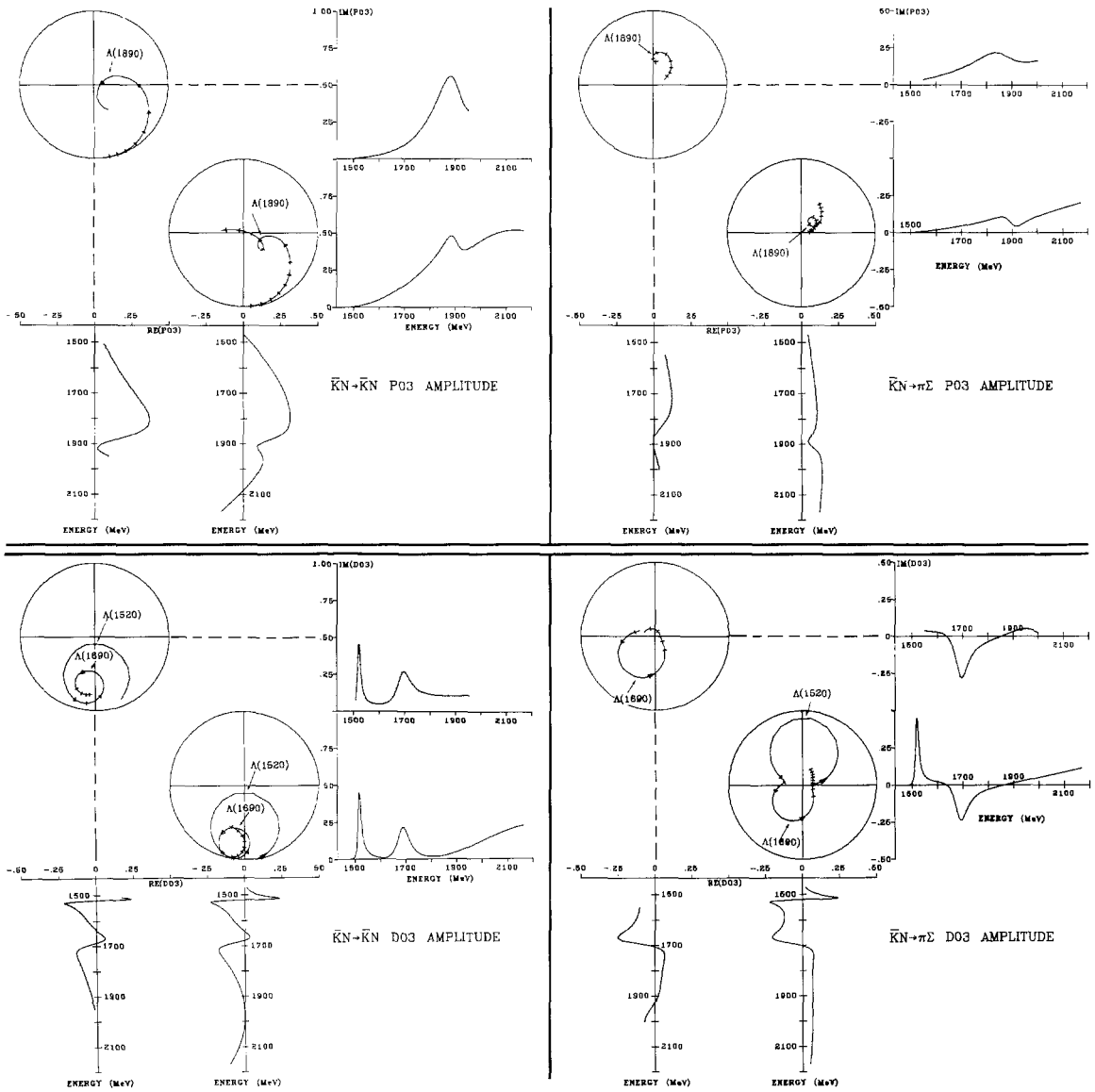


Fig. 2(b). The  $L_{I,2J} = P_{03}$  and  $D_{03}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plots for the elastic amplitudes are from ALSTON 78, and the upper plots for the  $\Sigma\pi$  amplitudes are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

For notation, see key on page 91.

### Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

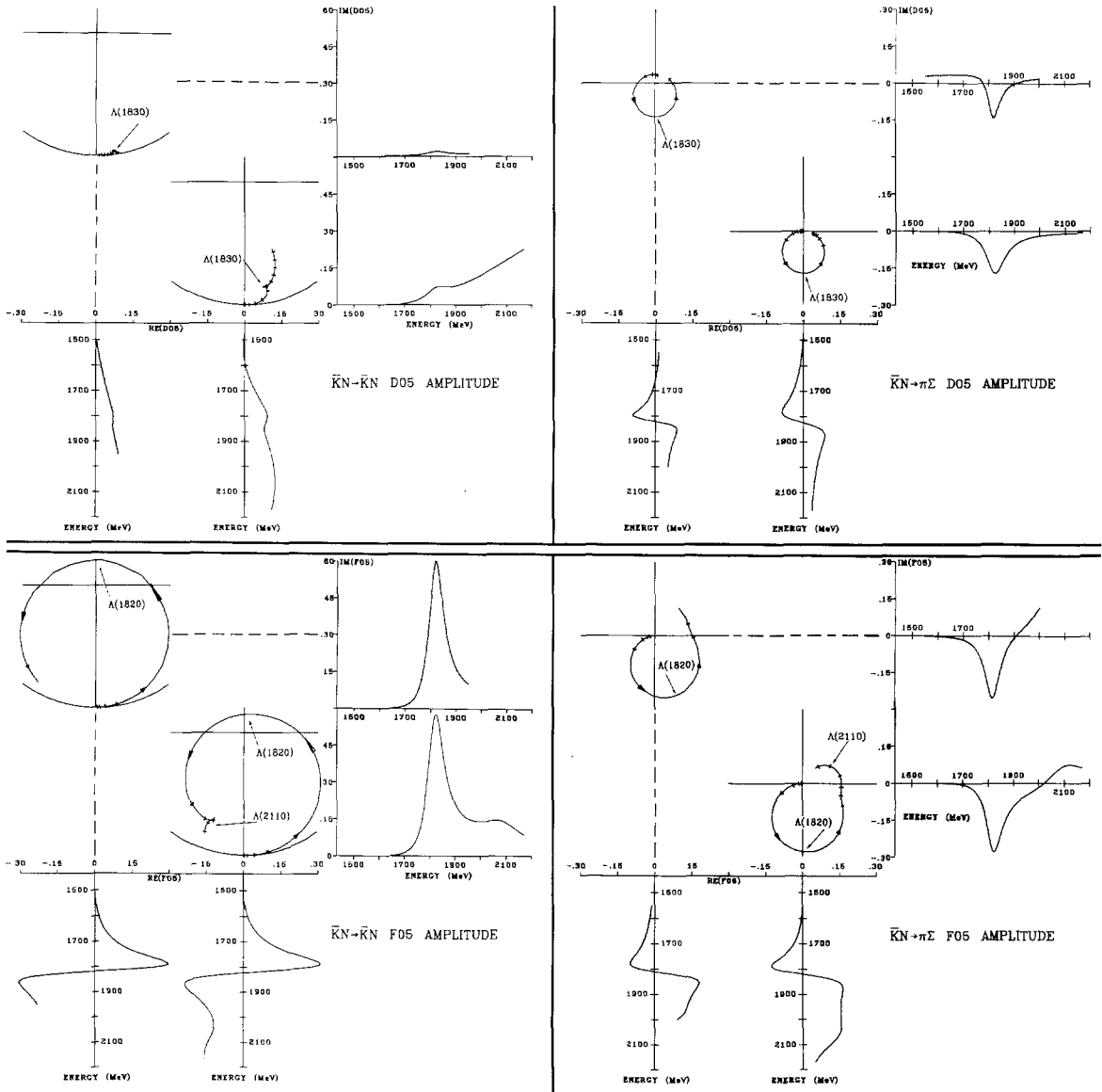


Fig. 2(c). The  $L_{J,2J} = D_{05}$  and  $F_{05}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plots for the elastic amplitudes are from ALSTON 78, and the upper plots for the  $\Sigma\pi$  amplitudes are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.



# Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

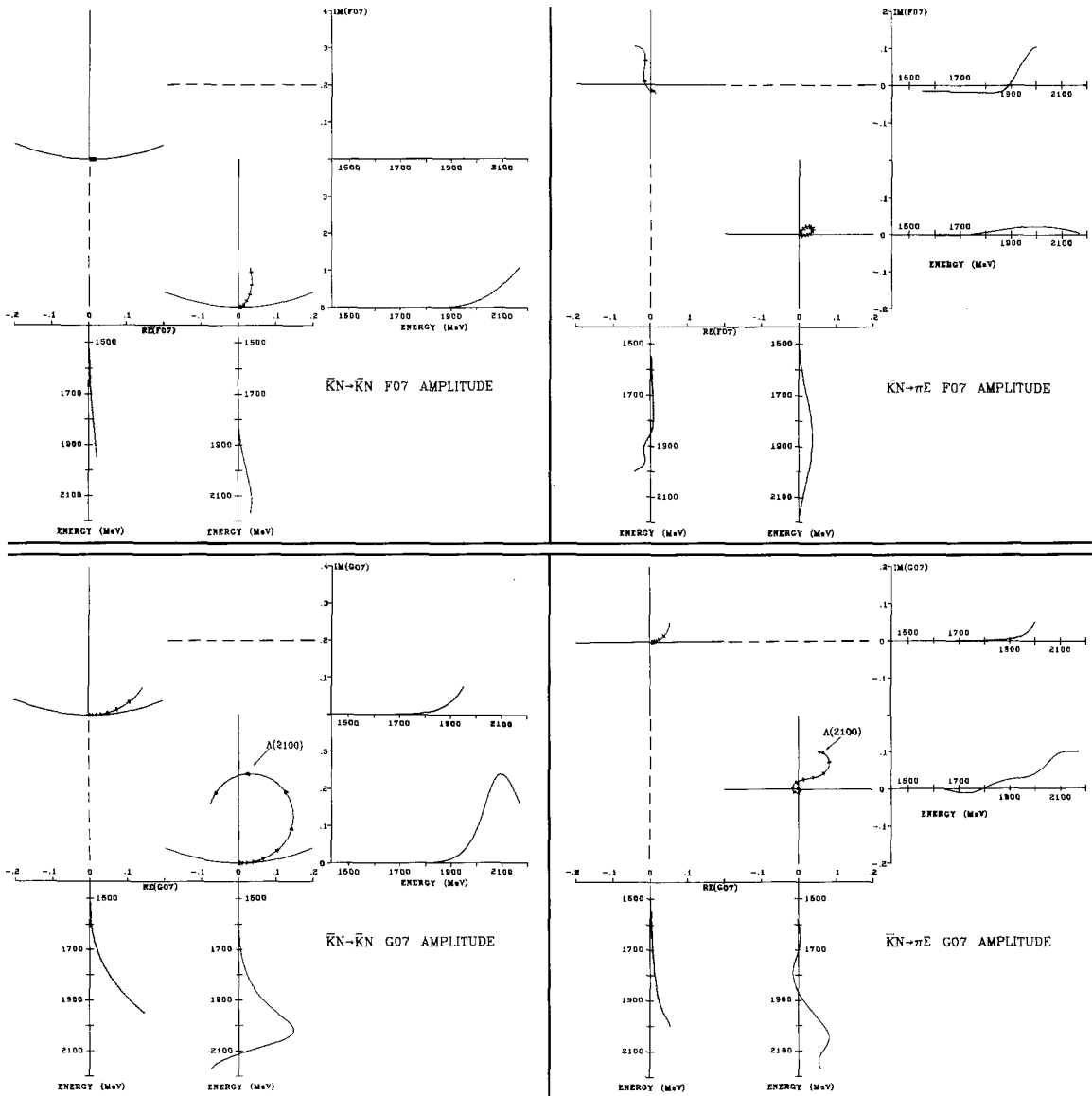


Fig. 2(d). The  $L_{1,2J} = F_{07}$  and  $G_{07}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plots for the elastic amplitudes are from ALSTON 78, and the upper plots for the  $\Sigma\pi$  amplitudes are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

For notation, see key on page 91.

## Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

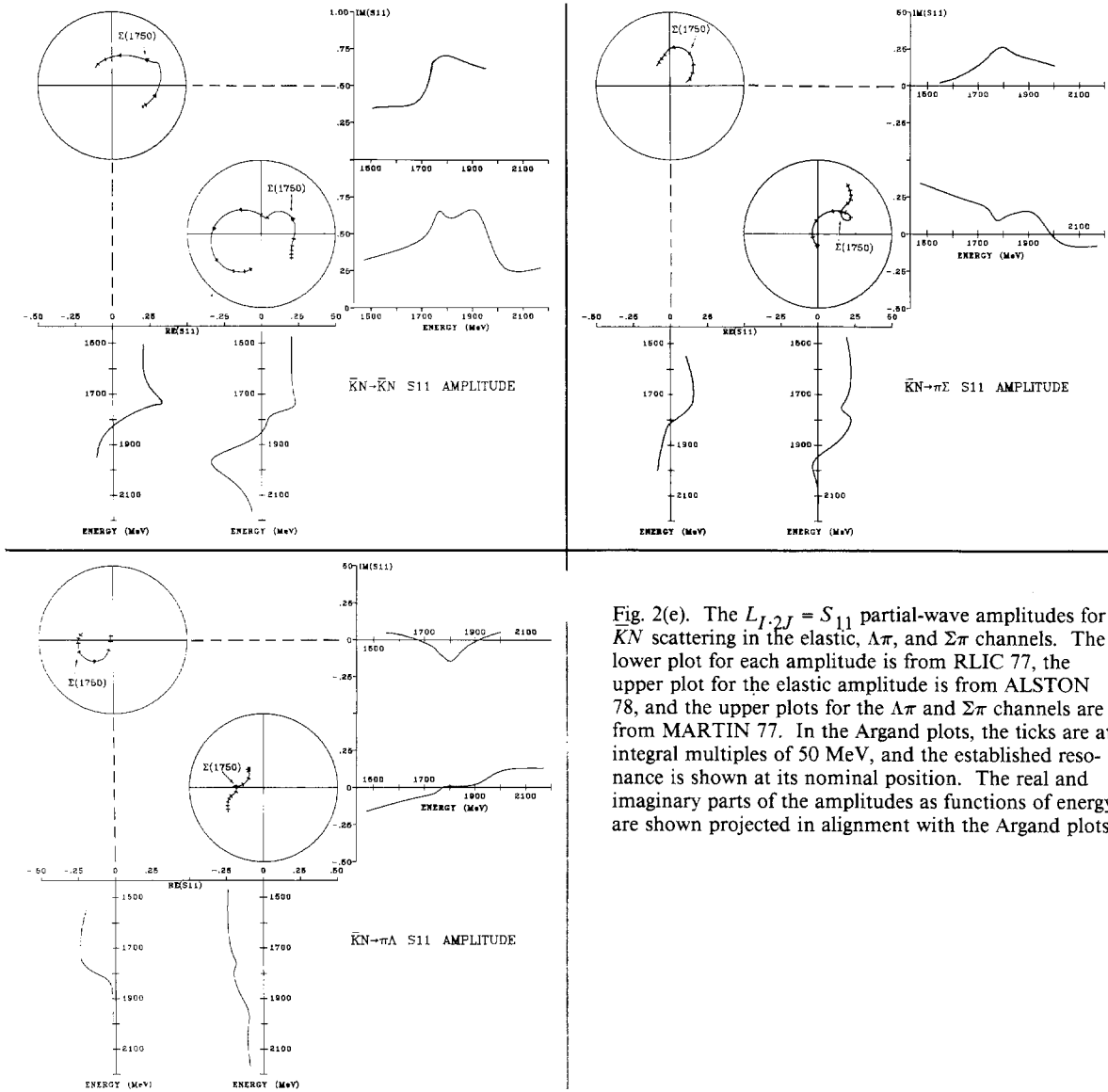


Fig. 2(e). The  $L_{I,2J} = S_{11}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic,  $\Lambda\pi$ , and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the  $\Lambda\pi$  and  $\Sigma\pi$  channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

# Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

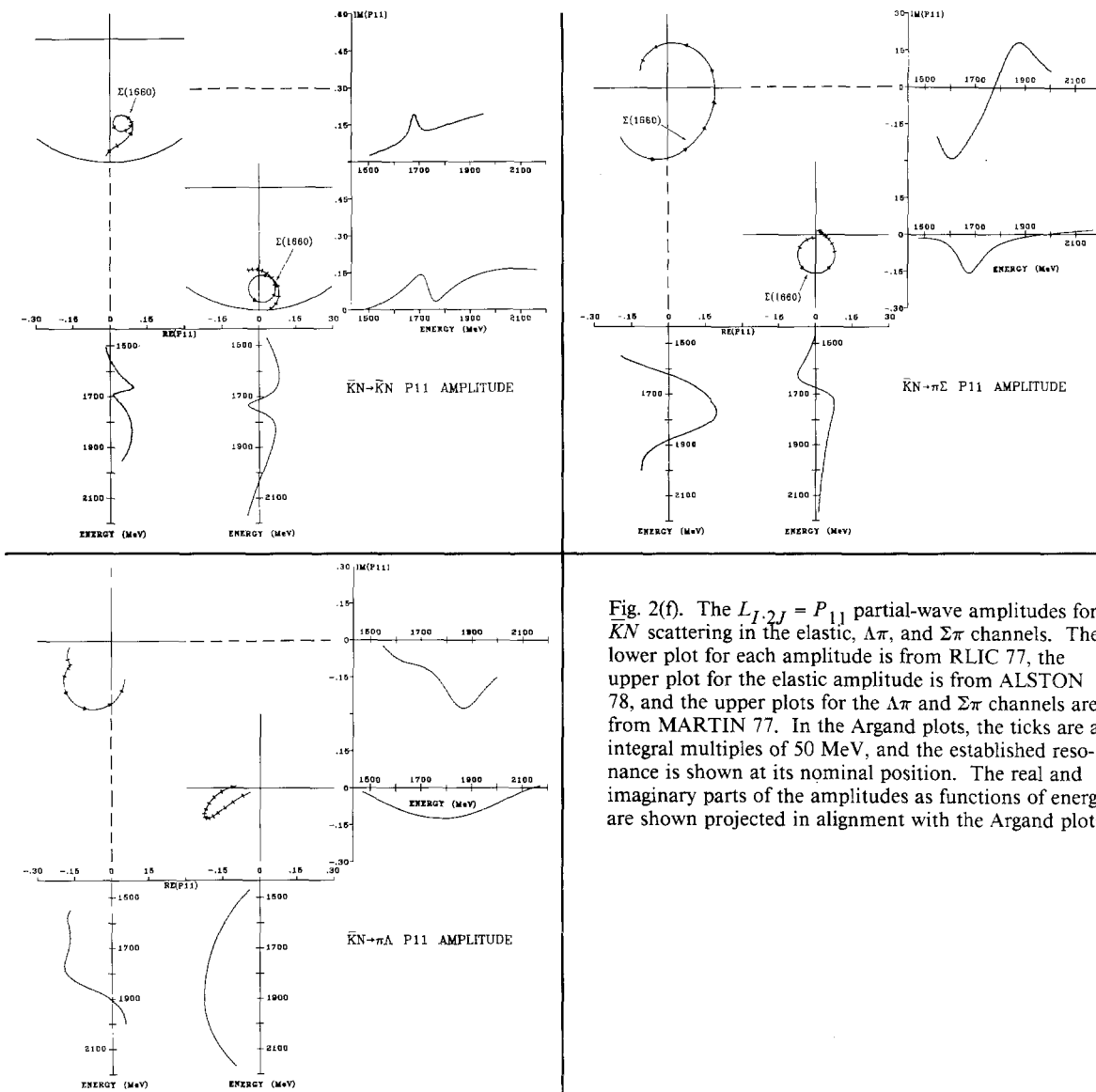


Fig. 2(f). The  $L_{I,2J} = P_{11}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic,  $\Delta\pi$ , and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the  $\Delta\pi$  and  $\Sigma\pi$  channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

For notation, see key on page 91.

## Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

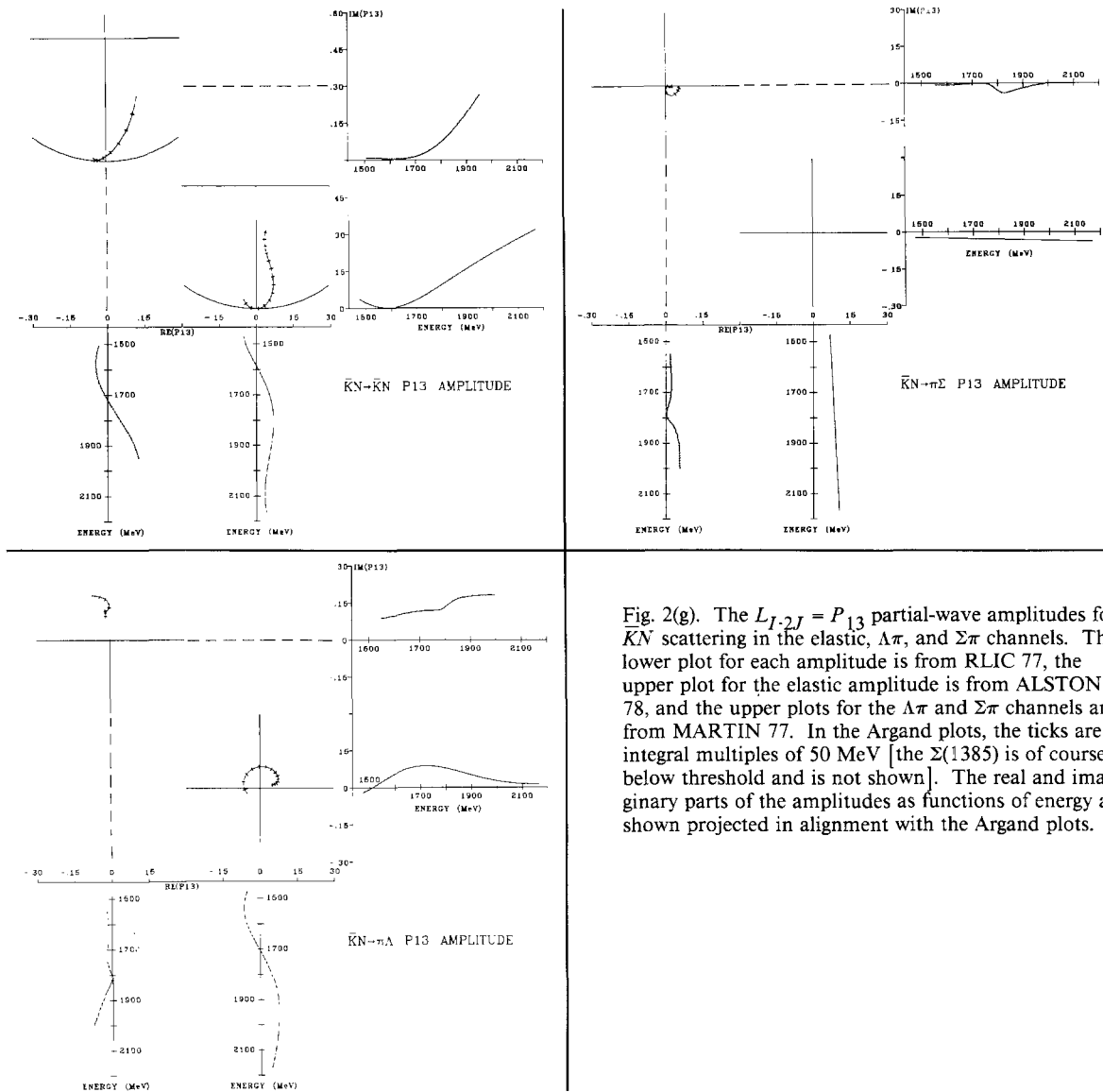


Fig. 2(g). The  $L_{J-2J} = P_{13}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic,  $\Lambda\pi$ , and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the  $\Lambda\pi$  and  $\Sigma\pi$  channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV [the  $\Sigma(1385)$  is of course below threshold and is not shown]. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

# Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

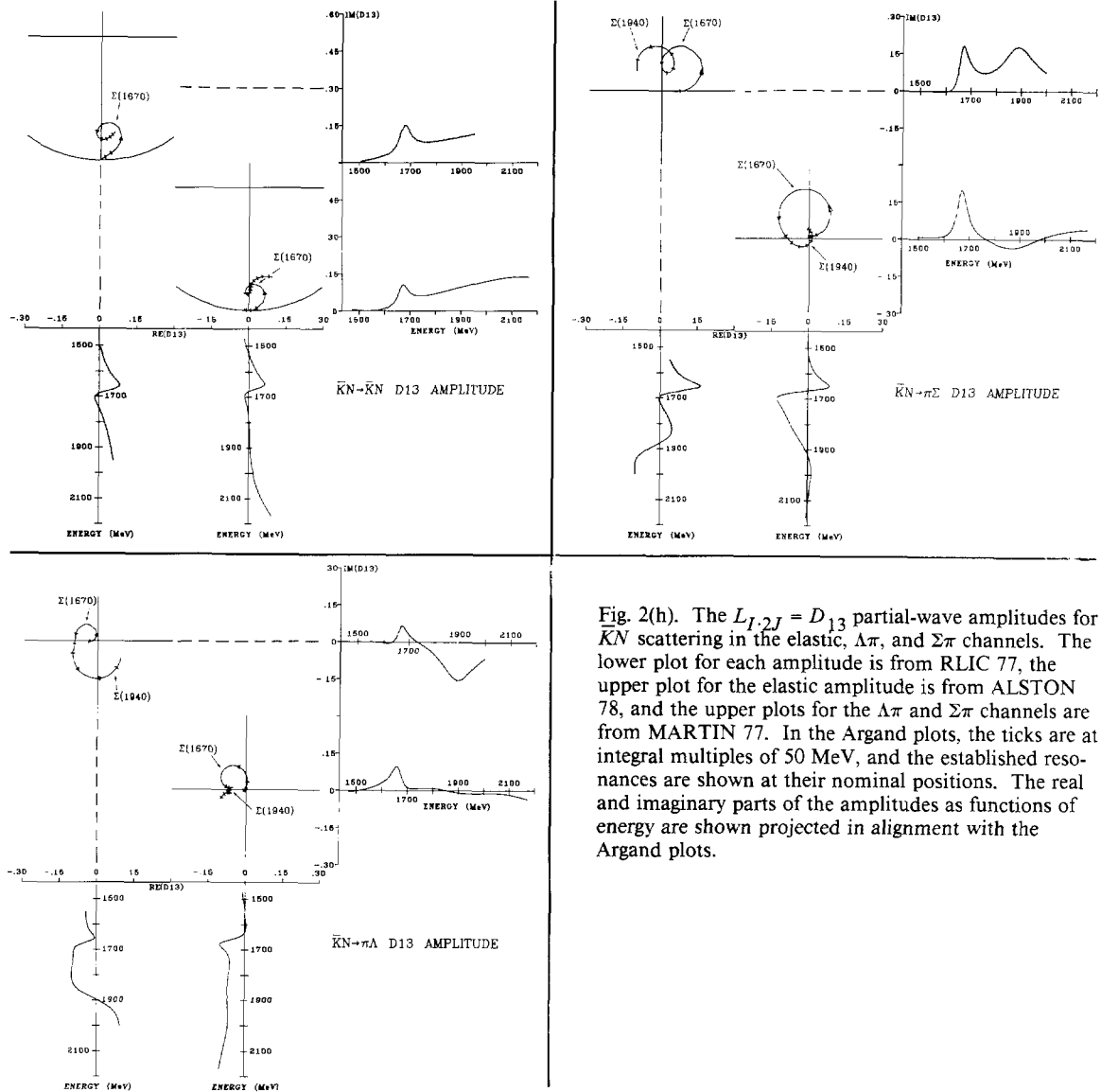


Fig. 2(h). The  $L_{I,2J} = D_{13}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic,  $\Delta\pi$ , and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the  $\Delta\pi$  and  $\Sigma\pi$  channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonances are shown at their nominal positions. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

## Baryon Full Listings

 $\Lambda$ 's and  $\Sigma$ 's

For notation, see key on page 91.

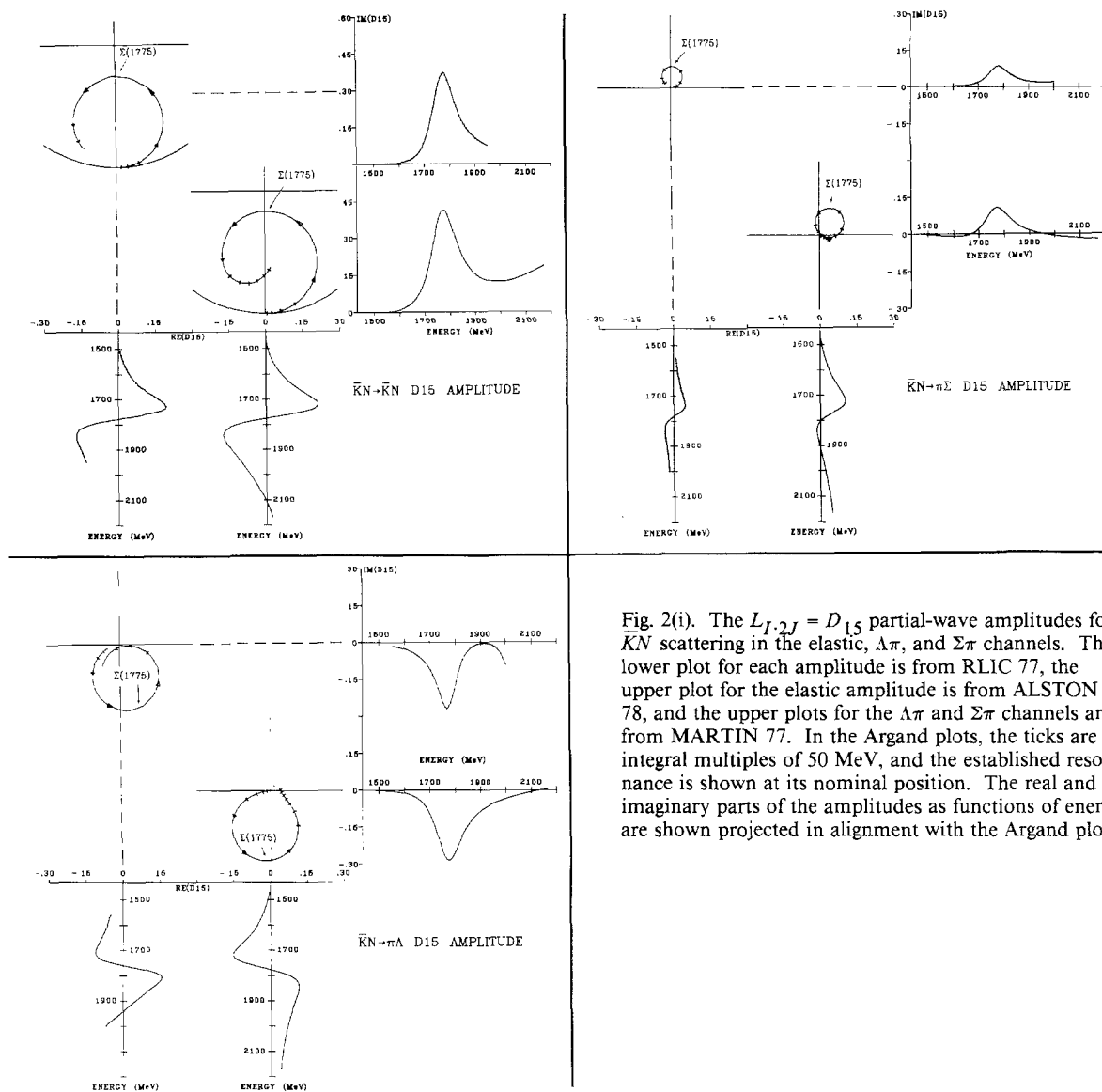


Fig. 2(i). The  $L_I \cdot 2J = D_{15}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic,  $\Lambda\pi$ , and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the  $\Lambda\pi$  and  $\Sigma\pi$  channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

# Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

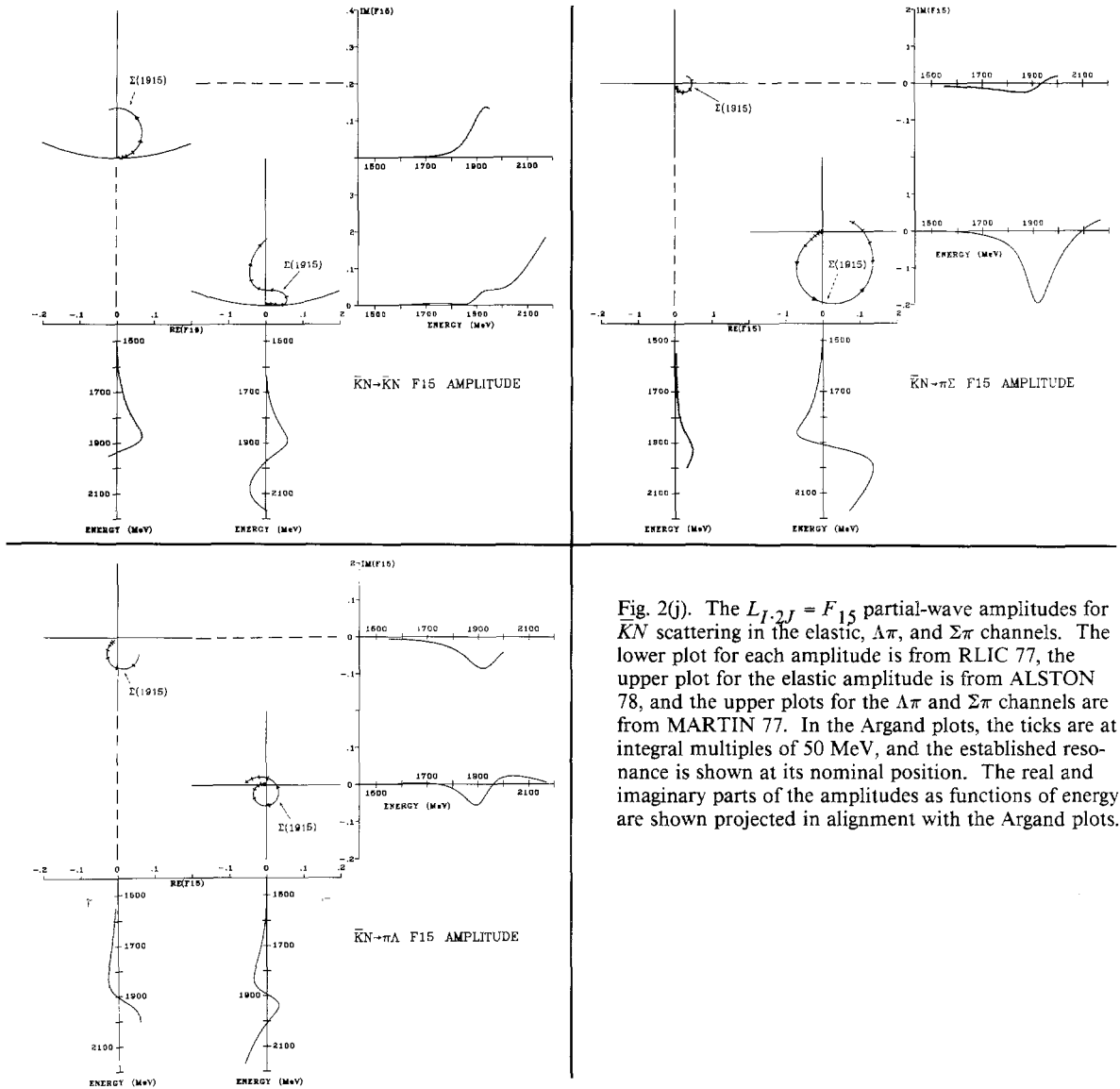


Fig. 2(j). The  $L_{J,2J} = F_{15}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic,  $\Lambda\pi$ , and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the  $\Lambda\pi$  and  $\Sigma\pi$  channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.

For notation, see key on page 91.

## Baryon Full Listings

$\Lambda$ 's and  $\Sigma$ 's

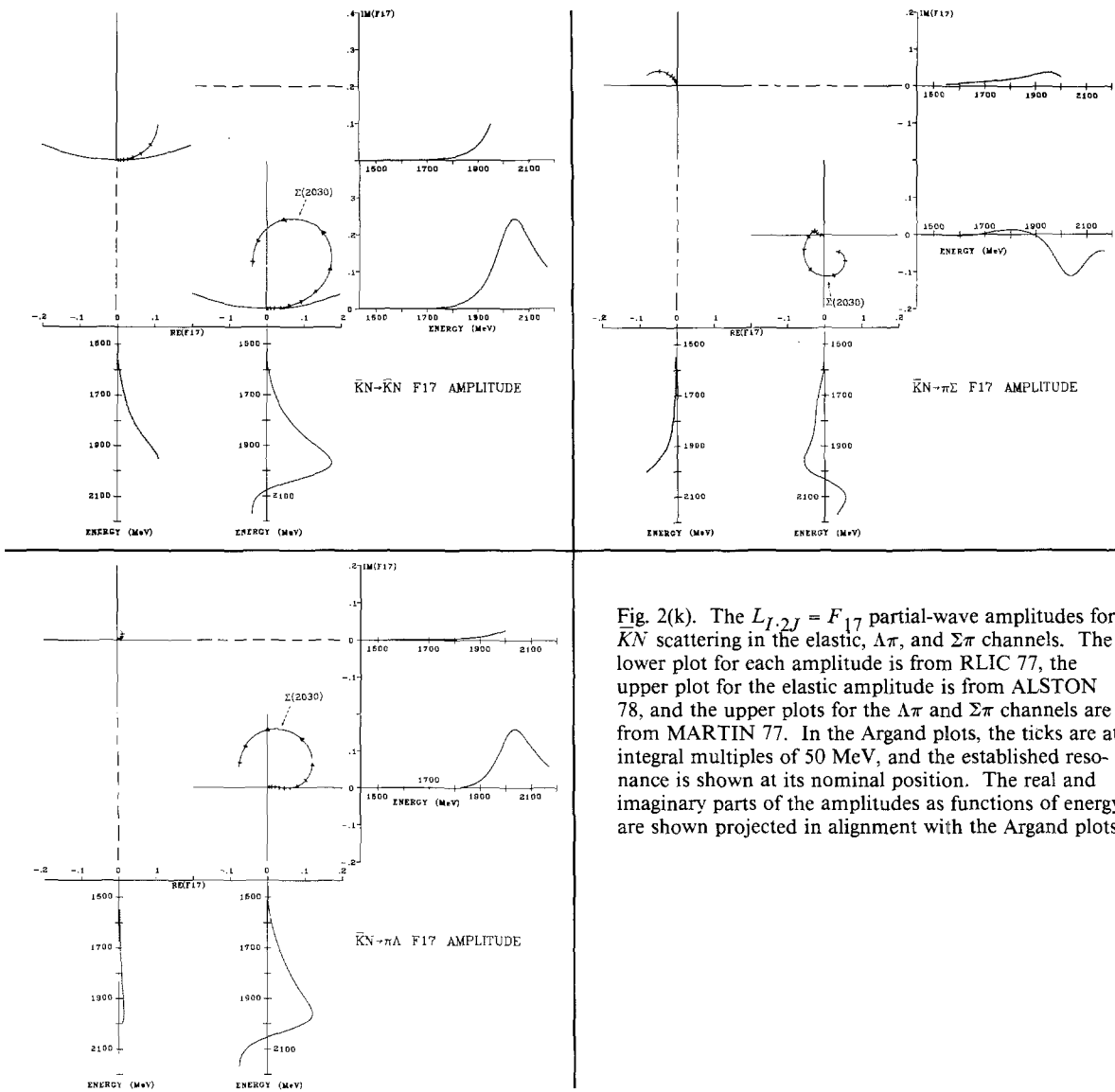


Fig. 2(k). The  $L_{I,2J} = F_{17}$  partial-wave amplitudes for  $\bar{K}N$  scattering in the elastic,  $\Delta\pi$ , and  $\Sigma\pi$  channels. The lower plot for each amplitude is from RLIC 77, the upper plot for the elastic amplitude is from ALSTON 78, and the upper plots for the  $\Delta\pi$  and  $\Sigma\pi$  channels are from MARTIN 77. In the Argand plots, the ticks are at integral multiples of 50 MeV, and the established resonance is shown at its nominal position. The real and imaginary parts of the amplitudes as functions of energy are shown projected in alignment with the Argand plots.





For notation, see key on page 91.

Baryon Full Listings
Lambda(1520), Lambda(1600)

Lambda(1520) PARTIAL DECAY MODES

Table with columns for decay mode, branching fraction, and decay masses. Includes modes like Lambda(1520) -> N K-bar, Lambda(1520) -> Sigma pi, Lambda(1520) -> Delta pi pi, etc.

FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions. P1 as follows. The diagonal elements are P1, P1, while the off-diagonal elements are the normalized correlation coefficients (P1, P1)/(P1, P1). For the definitions of the individual Pi, see the listings above; only those Pi appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

Correlation matrix table with columns P1 through P6 and rows P1 through P6.

Lambda(1520) BRANCHING RATIOS

Table showing branching ratios for Lambda(1520) -> (Sigma pi)/(N K-bar) and Lambda(1520) -> (Delta pi pi)/(N K-bar) with associated error and fit information.

Table showing branching ratios for Lambda(1520) -> (Delta pi pi)/(N K-bar) with associated error and fit information.

Table showing branching ratios for Lambda(1520) -> (Sigma pi)/(Delta pi pi) with associated error and fit information.

Table showing branching ratios for Lambda(1520) -> (Lambda gamma)/total with associated error and fit information.

Table showing branching ratios for Lambda(1520) -> (Sigma 0 gamma)/total with associated error and fit information.

Table showing branching ratios for Lambda(1520) -> (N K-bar)/total with associated error and fit information.

Table showing branching ratios for Lambda(1520) -> (Sigma pi)/total with associated error and fit information.

Table showing branching ratios for Lambda(1520) -> (Sigma pi pi)/total with associated error and fit information.

Lambda(1520) -> (Sigma(1385) pi) -> Lambda pi pi / Delta pi pi

Table listing various decay channels and their branching fractions for Lambda(1520) -> (Sigma(1385) pi) -> Lambda pi pi / Delta pi pi.

Table showing the total branching fraction for Lambda(1520) -> (Sigma(1385) pi) with associated error and fit information.

Table showing the total branching fraction for Lambda(1520) -> (Lambda pi pi) with associated error and fit information.

Table showing the total branching fraction for Lambda(1520) -> (Lambda pi pi) S-wave with associated error and fit information.

REFERENCES FOR Lambda(1520)

List of references for Lambda(1520) decay studies, including authors like Ferro-Luzzi, R D Tripp, M B Watson, etc.

Lambda(1600) P01

I(J^P) = 0(1/2^+)

Status: \*\*\*

SEE THE NOTE FOR THE LAMBDA(1600), JP=1/2+, P01. SOMEWHERE IN THIS REGION THERE IS PROBABLY ONE, AND PERHAPS TWO, P01 STATES.

Lambda(1600) MASS (MeV)

Table listing various decay channels and their branching fractions for Lambda(1600) MASS.

Lambda(1600) WIDTH (MeV)

Table listing various decay channels and their branching fractions for Lambda(1600) WIDTH.

Lambda(1600) PARTIAL DECAY MODES

Table showing partial decay modes for Lambda(1600) with branching fractions.

Table showing partial decay modes for Lambda(1600) with branching fractions.



For notation, see key on page 91.

Baryon Full Listings

Λ(1690), Λ(1800)

Λ(1690) BRANCHING RATIOS

THE SUM OF ALL THE QUOTED BRANCHING RATIOS IS MORE THAN 1.0. THE TWO-BODY RATIOS ARE FROM PARTIAL-WAVE ANALYSES, AND THUS PROBABLY ARE MORE RELIABLE THAN THE THREE-BODY RATIOS, WHICH ARE DETERMINED FROM BUMPS IN CROSS SECTIONS. OF THE LATTER, THE SIGMA PI PI BUMP LOOKS MORE SIGNIFICANT (THE ERROR GIVEN FOR THE LAMBDA PI PI RATIO LOOKS UNREASONABLY SMALL). HARDLY ANY OF THE SIGMA PI PI DECAY CAN BE VIA SIGMA(1385), FOR THEN NINE TIMES AS MUCH LAMBDA PI PI DECAY WOULD BE REQUIRED.

Table with columns for decay modes (e.g., Σπ, Δππ, Σππ, Δη, Σ(1385)π, S-wave) and branching ratios from various experiments (e.g., MARTIN, ALSTON, GOPAL).

REFERENCES FOR Λ(1690)

List of references for Λ(1690) including authors like ARMENTEROS, LANGBEIN, BARTLEY, etc., and their respective publications.

Λ(1800) PARTIAL DECAY MODES

Table listing partial decay modes for Λ(1800) such as Σπ, Σ(1385)π, NK\*(892), and NK\*(892), along with their masses.

Λ(1800) BRANCHING RATIOS

Table with columns for decay modes (e.g., Σπ, Σ(1385)π, NK\*(892), D3-wave) and branching ratios from various experiments (e.g., BRICMAN, KIM, LANGBEIN).

REFERENCES FOR Λ(1800)

List of references for Λ(1800) including authors like BRICMAN, KIM, LANGBEIN, MARTIN, etc., and their respective publications.

Λ(1800) P01

I(J^P) = 0(1/2^+) Status: \*\*\*

THE EVIDENCE FOR THIS STATE IS SOMEWHAT CONFUSED. IT WAS FIRST SUGGESTED IN A PARTIAL-WAVE ANALYSIS OF N KBAR DATA BY THE BEHAVIOUR OF THE P01 AMPLITUDE WHEN IT WAS PARAMETRIZED AS A TWO-STRAIGHT-LINE BACKGROUND (ARMENTEROS 68).

ALMOST ALL THE RECENT ANALYSES CONTAIN A P01 STATE, AND SOMETIMES TWO, BUT THE MASSES, WIDTHS, AND BRANCHING RATIOS VARY GREATLY. SEE ALSO THE LAMBDA(1600) P01 LISTING.

Λ(1800) S01

I(J^P) = 0(1/2^-) Status: \*\*\*

THE S01 AMPLITUDE SHOWS A RATHER CLEAR SECOND RESONANCE BEHAVIOR IN THE 1700-1900 MEV REGION. THERE ARE MAJOR DISAGREEMENTS ABOUT THE MASS, WIDTH, AND COUPLINGS.

Λ(1800) MASS (MeV)

Table listing mass values for Λ(1800) from various experiments (e.g., BRICMAN, KIM, LANGBEIN, MARTIN).

Λ(1800) WIDTH (MeV)

Table listing width values for Λ(1800) from various experiments (e.g., BRICMAN, KIM, LANGBEIN, MARTIN).

Λ(1800) MASS (MeV)

Table listing mass values for Λ(1800) from various experiments (e.g., ARMENTEROS, LANGBEIN, MARTIN).

Λ(1800) WIDTH (MeV)

Table listing width values for Λ(1800) from various experiments (e.g., ARMENTEROS, BAILEY, LANGBEIN, MARTIN).





# Baryon Full Listings

$\Lambda(1890)$ ,  $\Lambda(2000)$ ,  $\Lambda(2020)$ ,  $\Lambda(2100)$

ALBROW 71 NP B29 413 +ANDERSON,BOSNIAKOVIC,DAUM,ERNZ,+ (CERN)  
 CONFORTO 71 NP B54 411 +LEVI SETTI,LASINSKI, OBERLACK+ (EFI-HEID)IJP  
 KIM 71 PRL 27 356 J K KIM (HARV)IJP  
 ALSO 70 DUKE 161 J K KIM (HARV)IJP  
 LANGBEIN 72 NP B47 477 +WAGNER (MPIM)IJP  
 LEA 73 NP B56 77 +MARTIN,MOORHOUSE- (RHEL-LOUC+GLAS+AARH)IJP  
 HEMINGWA 75 NP B91 12 HEMINGWAY,EADES,HARMSEN+ (CERN+HEID-MPIM)IJP  
 NAKKASYA 75 NP B93 85 A NAKKASYAN (CERN)IJP  
 BACCARI 77 NC 41A 96 +POULARD,REVEL,TALLINI+ (SACL-CDEF)IJP  
 MARTIN 77 NP B127 349 MARTIN,PIDCOCK,MOORHOUSE (LOUC-GLAS)IJP  
 ALSO 77 NP B126 266 MARTIN,PIDCOCK (LOUC)  
 ALSO 77 NP B126 285 MARTIN,PIDCOCK (LOUC)IJP  
 RLIC 77 NP B119 362 GOPAL,ROSS,VAN HORN,MCPHERSON+ (LOIC-RHEL)IJP  
 ALSTON 78 PR D18 182 +KENNEY,POLLARD,ROSS+ (LBL+MTHO-CERN)IJP  
 ALSO 77 PRL 38 1007 ALSTON-GARN,JOST,KENNEY,+ (LBL+MTHO-CERN)IJP  
 CAMERON 78 NP B143 189 +FRANEK,GOPAL,BACON,BUTTERWORTH-(RHEL-LOIC)IJP  
 CAMERON2 78 NP B146 327 +FRANEK,GOPAL,KALMUS,MCPHERSON+-(RHEL-LOIC)IJP  
 GOPAL 80 TORONTO CONF 159 G P GOPAL (RHEL)IJP

$\Lambda(2020) F_{07}$

$I(J^P) = 0(\frac{7}{2}^-)$  Status: \*

OMITTED FROM SUMMARY TABLE

EFFECTS IN THIS PARTIAL WAVE HAVE BEEN OBSERVED AT DIFFERENT ENERGIES IN TWO CHANNELS. IN LITCHFIELD 71, NEED FOR THE STATE RESTS SOLELY ON POSSIBLY INCONSISTENT POLARIZATION MEASUREMENT AT 1.784 GEV/C. HEMINGWAY 75 ANALYSIS OF N KBAR DOES NOT REQUIRE THIS STATE. RLIC 77 DO NOT NEED IT IN EITHER N KBAR OR SIGMA P1 WITH NEW K- NEUTRON ANGULAR DISTRIBUTIONS INCLUDED, DECLAIS 77 SEE THIS STATE. HOWEVER, THIS AND OTHER NEW DATA ARE INCLUDED IN GOPAL 80 AND THIS STATE IS NOT REQUIRED. BACCARI 77 WEAKLY SUPPORTS THIS STATE.

$\Lambda(2000)$

$I(J^P) = 0(\quad)$  Status: \*

OMITTED FROM SUMMARY TABLE

WE LIST HERE ALL THE AMBIGUOUS RESONANCE POSSIBILITIES WITH A MASS AROUND 2 GEV. THE PROPOSED QUANTUM NUMBERS ARE D3 (GALTIERI 70 IN SIGMA P1), D3F5, P3+05, OR P1+03 (BRANDSTETTER 72 IN LAMBDA OMEGA), AND S1 (CAMERON2 78 IN N K\*). THE FIRST TWO OF THE ABOVE ANALYSES SHOULD NOW BE CONSIDERED OBSOLETE.

**$\Lambda(2020)$  MASS (MeV)**

M	(2020.0)	(20.0)	GALTIERI	70 DPWA	0 K-P TO SIGMA P1
M	(2100.0)	(30.0)	LITCHFIE	71 DPWA	K-P TO KBAR N
M	(2140.0)		BACCARI	77 DPWA	0 K-P TO LAM. OMG.
M	(2117.0)		DECLAIS	77 DPWA	KBAR N TO KBAR N

**$\Lambda(2020)$  WIDTH (MeV)**

W	(160.0)	(30.0)	GALTIERI	70 DPWA	0 K-P TO SIGMA P1
W	(120.0)	(30.0)	LITCHFIE	71 DPWA	K-P TO KBAR N
W	(128.0)		BACCARI	77 DPWA	0 K-P TO LAM. OMG.
W	(167.0)		DECLAIS	77 DPWA	KBAR N TO KBAR N

**$\Lambda(2000)$  MASS (MeV)**

M	(2010.0)	(30.0)	GALTIERI	70 DPWA	0 K-P TO SIGMA P1
M	A	1935. TO 1971.	BRANDSTETTER	72 DPWA	0 K-P TO LAM. OMG.
M	A	1951. TO 2034.	BRANDSTETTER	72 DPWA	0 K-P TO LAM. OMG.
M	A	PARAMETERS QUOTED ARE RANGES FROM THREE BEST FITS, THE LOWER			
M	A	(HIGHER) MASS STATE PROBABLY HAS J.L.3/2(5/2).			
M		2030.0	30.0	CAMERON2	78 DPWA K-P TO K*(892) N

**$\Lambda(2020)$  PARTIAL DECAY MODES**

P1	$\Lambda(2020) \rightarrow N \bar{K}$	DECAY MASSES	938+ 494
P2	$\Lambda(2020) \rightarrow \Sigma \pi$		1189+ 140
P3	$\Lambda(2020) \rightarrow \Lambda \omega$		1116+ 783

**$\Lambda(2000)$  WIDTH (MeV)**

W	(130.0)	(50.0)	GALTIERI	70 DPWA	0 K-P TO SIGMA P1
W	A	180. TO 240. (LWR. MASS)	BRANDSTETTER	72 DPWA	0 K-P TO LAM. OMG.
W	A	73. TO 154. (HGR. MASS)	BRANDSTETTER	72 DPWA	0 K-P TO LAM. OMG.
W		125.0	25.0	CAMERON2	78 DPWA K-P TO K*(892) N

**$\Lambda(2020)$  BRANCHING RATIOS**

$\Lambda(2020) \rightarrow (N \bar{K})/\text{total}$					
R1	(0.05)	(0.02)	LITCHFIE	71 DPWA	K-P TO KBAR N (P1)
R1	(0.05)		DECLAIS	77 DPWA	KBAR N TO KBAR N
$\Lambda(2020) \text{ in } N \bar{K} \rightarrow \Sigma \pi$					
R2	(-0.15)	(0.02)	GALTIERI	70 DPWA	0 K-P TO SIGMA P1
$\Lambda(2020) \text{ in } N \bar{K} \rightarrow \Lambda \omega$					
R3	LESS THAN 0.05		BACCARI	77 DPWA	0 K-P TO LAM. OMG.

**$\Lambda(2000)$  PARTIAL DECAY MODES**

P1	$\Lambda(2000) \rightarrow N \bar{K}$	DECAY MASSES	938+ 494
P2	$\Lambda(2000) \rightarrow \Sigma \pi$		1189+ 140
P3	$\Lambda(2000) \rightarrow \Lambda \omega$		1116+ 783
P4	$\Lambda(2000) \rightarrow N \bar{K}^*(892), S1\text{-wave}$		940+ 892
P5	$\Lambda(2000) \rightarrow N \bar{K}^*(892), D3\text{-wave}$		940+ 892

**$\Lambda(2000)$  BRANCHING RATIOS**

$\Lambda(2000) \text{ in } N \bar{K} \rightarrow \Sigma \pi$					
R1	(-0.20)	(0.04)	GALTIERI	70 DPWA	0 K-P TO SIGMA P1
$\Lambda(2000) \text{ in } N \bar{K} \rightarrow \Lambda \omega$					
R2	A	0.17 TO 0.25 (LWR.)	BRANDSTETTER	72 DPWA	0 K-P TO LAM. OMG.
R2	A	0.04 TO 0.15 (HGR.)	BRANDSTETTER	72 DPWA	0 K-P TO LAM. OMG.
$\Lambda(2000) \text{ in } N \bar{K} \rightarrow N \bar{K}^*(892), S1\text{-wave}$					
R3	B	-0.12	0.03	CAMERON2	78 DPWA K-P TO K*N
R3	B	THE SIGN HERE IS CHANGED TO BE IN ACCORD WITH THE BARYON-FIRST			
R3	B	CONVENTION.			
$\Lambda(2000) \text{ in } N \bar{K} \rightarrow N \bar{K}^*(892), D3\text{-wave}$					
R4		-0.09	0.03	CAMERON2	78 DPWA K-P TO K*N

**REFERENCES FOR  $\Lambda(2000)$**

GALTIERI 70 DUKE CONF 173 A BARBARO-GALTIERI (LRL)IJP  
 BRANDSTETTER 72 NP B39 13 BRANDSTETTER,BUTTERWORTH,+ (RHEL-CDEF+SACL)IJP  
 CAMERON2 78 NP B146 327 +FRANEK,GOPAL,KALMUS,MCPHERSON+ (RHEL+LOIC)IJP  
 PAPERS NOT REFERRED TO IN DATA LISTINGS  
 NAKKASYA 75 NP B93 85 A NAKKASYAN (CERN)IJP

$\Lambda(2100) G_{07}$

$I(J^P) = 0(\frac{7}{2}^-)$  Status: \*\*\*\*

FOR MOST RESULTS PUBLISHED BEFORE 1973 (THEY ARE NOW OBSOLETE), SEE OUR 1982 EDITION (PHYSICS LETTERS 111B). ALL THE REFERENCES HAVE BEEN RETAINED.

THIS ENTRY ONLY INCLUDES RESULTS FROM PARTIAL-WAVE ANALYSES. PARAMETERS OF PEAKS SEEN IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS AROUND 2100 MEV ARE GIVEN IN A SEPARATE ENTRY BELOW.

**$\Lambda(2100)$  MASS (MeV)**

M	2115.0	(10.0)	KANE	74 DPWA	K-P TO PI SIG
M	2105.0	(10.0)	HEMINGWAY	75 DPWA	0 K-P TO KBAR N
M	A	2110. OR 2085.	NAKKASYA	75 DPWA	0 K-P TO LAM. OMG.
M	A	QUOTED PARAMETERS CORRESPOND TO THE TWO BEST SOLUTIONS FOUND.			
M	A	EACH HAS THE LAMBDA(2100) AND ONE ADDITIONAL RESONANCE (P3 OR F5).			
M	(2094.0)		BACCARI	77 DPWA	0 K-P TO LAM. OMG.
M	(2094.0)		DECLAIS	77 DPWA	KBAR N TO KBAR N
M	2110.0	(10.0)	RLIC	77 DPWA	KBAR N MULTICHNL
M	2106.0	(30.0)	BELLEFDON	78 DPWA	0 KBAR N TO KBAR N
M	2104.0	(10.0)	GOPAL	80 DPWA	KBAR N ELASTIC





Baryon Full Listings

$\Lambda(2110), \Lambda(2325), \Lambda(2350), \Lambda(2585)$

$\Lambda(2110)$  PARTIAL DECAY MODES

Table with columns: Mode (P1, P2, P3, P4, P5), Decay, and Decay Masses (938-494, 1189+140, 1116+783, 1385+140, 940-892)

$\Lambda(2110)$  BRANCHING RATIOS

$\Lambda(2110)$  in  $N\bar{K} \rightarrow \Sigma\pi$  (SQRT(P1\*P2))

$\Lambda(2110) \rightarrow (N\bar{K})/\text{total}$  (P1)

$\Lambda(2110)$  in  $N\bar{K} \rightarrow \Lambda\omega$  (SQRT(P1\*P3))

$\Lambda(2110)$  in  $N\bar{K} \rightarrow \Sigma(1385)\pi, P\text{-wave}$  (SQRT(P1\*P4))

$\Lambda(2110)$  in  $N\bar{K} \rightarrow N\bar{K}^*(892), F1\text{-wave}$  (SQRT(P1\*P5))

REFERENCES FOR  $\Lambda(2110)$

BERTHON 70 NP 824 417 ... KANE 72 PR 05 1583 ... GOPAL 80 TORONTO CONF 159 6 P GOPAL

$\Lambda(2325) D_{03}$

$I(J^P) = 0(\frac{3}{2}^-)$  Status: \*

OMITTED FROM SUMMARY TABLE

BACCARI 77 FIND THIS STATE WITH JP EITHER 3/2- OR 5/2- IN A DPWA OF K-P TO LAMBDA OMEGA FROM 2070 TO 2436 MEV...

$\Lambda(2325)$  MASS (MeV)

Table with columns: Mode, Mass, and Reference (BACCARI 77 DPWA, BELLEFON 78 DPWA)

$\Lambda(2325)$  WIDTH (MeV)

Table with columns: Mode, Width, and Reference (BACCARI 77 IPWA, BELLEFON 78 DPWA)

$\Lambda(2325)$  PARTIAL DECAY MODES

Table with columns: Mode (P1, P2), Decay, and Decay Masses (938-494, 1116+783)

$\Lambda(2325)$  BRANCHING RATIOS

$\Lambda(2325)$  in  $N\bar{K} \rightarrow \Lambda\omega$  (SORT(P1\*P2))

$\Lambda(2325) \rightarrow (N\bar{K})/\text{total}$  (P1)

REFERENCES FOR  $\Lambda(2325)$

BACCARI 77 NC 41A 96 ... BELLEFON 78 NC 42A 403

$\Lambda(2350)$  BUMPS

$I(J^P) = 0(\frac{3}{2}^+)$  Status: \*\*\*

DAUM 68 FAVORS JP=7/2- OR 9/2-. BRICMAN 70 FAVORS 9/2-. LASINSKI 71 SUGGESTS THREE STATES IN THIS REGION USING A POMERON - RESONANCES MODEL...

$\Lambda(2350)$  MASS (MeV) (PROD. EXP.)

Table with columns: Mode, Mass, Reference (BUGG 68 CNTR, BRICMAN 70 CNTR, COOL 70 CNTR)

$\Lambda(2350)$  WIDTH (MeV) (PROD. EXP.)

Table with columns: Mode, Width, Reference (BUGG 68 CNTR, BRICMAN 70 CNTR, COOL 70 CNTR)

$\Lambda(2350)$  PARTIAL DECAY MODES (PROD. EXP.)

Table with columns: Mode (P1, P2, P3), Decay, and Decay Masses

$\Lambda(2350)$  BRANCHING RATIOS (PROD. EXP.)

$\Lambda(2350) \rightarrow (N\bar{K})/\text{total}$  (P1)

$\Lambda(2350)$  in  $N\bar{K} \rightarrow \Sigma\pi$  (SQRT(P1\*P2))

$\Lambda(2350)$  in  $N\bar{K} \rightarrow \Lambda\omega$  (SQRT(P1\*P3))

$\Lambda(2350) \rightarrow (N\bar{K})/\text{total}$  (J+1/2)\*(P1)

REFERENCES FOR  $\Lambda(2350)$  (PROD. EXP.)

BUGG 68 PR 169 1466 ... BRICMAN 70 PR D1 1887 ... LASINSKI 71 NP 829 125 ... BELLEFON 78 NC 28A 289

$\Lambda(2585)$  BUMPS

$I(J^P) = 0(\ )$  Status: \*\*

OMITTED FROM SUMMARY TABLE

$\Lambda(2585)$  MASS (MeV) (PROD. EXP.)

Table with columns: Mode, Mass, Reference (ABRAMS 70 CNTR, LU 70 CNTR)





For notation, see key on page 91.

Baryon Full Listings

Σ(1385), Σ(1480), Σ(1560)

Σ(1385) → (Δ γ)/total (P3)
R2 1 (0.17) (0.17) MEISNER 72 HBC 0 1 EVENT ONLY
Σ(1385) in N K̄ → Δ π SQRT(P1\*P4)
R3 S -0.586 0.319 DEVENISH 74 0 FIXED T DISP REL
R3 S EXTRAPOLATION OF PARAMETRIZED AMPLITUDE BELOW THRESHOLD.

Σ(1385) → (Δ γ)/(Δ π) (P3)/(P1)
R4 (0.06) OR LESS CL=90 COLAS 75 HLBC 0 K-P 575-970 MEV

Σ(1385) → (Σ γ)/(Δ π) (P5)/(P1)
R5 (0.05) OR LESS CL=90 COLAS 75 HLBC 0 K-P 575-970 MEV

REFERENCES FOR Σ(1385)

ALSTON 60 PRL 5 520
BASTIEN 61 PRL 6 702
BERGE 61 PRL 6 557
DAHIL 61 PRL 6 142
ELY 61 PRL 7 461
MARTIN 61 PRL 6 283
ALSTON 62 CERN CONF 311
COLLEY 62 PR 128 1930
CURTIS 63 PR 152 1771
COOPER 64 PL 8 365
HUME 64 UCRL-11291 THESIS
ALSTON 69 PR 180 1824
ARMENTEROS 65 PL 19 75
BALTAY 65 PR 140 81027
MUSGRAVE 65 NC 35 735
SMITH 65 THESIS (UCLA)
BIRMINGHAM 66 PR 151 1148
LONDON 66 PR 143 1034
SIEGEL 67 UCRL 18041 THESIS
AGUILAR 70 PRL 25 58
ATHERTON 71 NP 829 477
COLLEY 71 NP 831 611
AGUILAR 72 PR 06 29
MEISNER 72 NC 124 62
AMMANN 73 PR 07 1345
MASTZ 73 PR 07 3212
ALSTON 73 PR 07 5
HABIBI 73 NEVIS 199(THESIS)
ALSTON 73 PURD73, PG. 387
THOMAS 73 NP B56 15
BERTHON 74 NC 214 146
BORENSTEIN 74 NP 09 3006
DEVENISH 74 NP B81 330
LICHTENBERG 74 PR D10 3865
ATHERTON 75 NC 254 1
BARDADIN 75 NP 898 418
COLAS 75 NP 891 253
BARREIRO 77 NP 8126 319
HOLMGREN 77 NP B119 261
ALSTON 78 PR D18 182
CAMERON 78 NP B143 189
DIONISI 78 PL 788 154
BANEJEE 79 ZPHY C3 1
BAUBILLIER 79 NP B148 18
CAUTIS 79 NP B156 507
SUGAHARA 79 NP B156 237
BAKER 80 NP B166 207
AGUILAR 81 AFIS 477 144
BAUBILLIER 84 ZPHY C23 213

MALAMUD 64 PL 10 145
SHAFFER 66 PR 134 B1372
HUNGERBUHLER 74 PR D10 205
WALTER 79 ZPHY C3 89
AGUILAR 80 ZPHY C6 109
BALAND 80 ZPHY C3 187
RIACAT 8 2PHY C9 305

Σ(1480) BUMPS

I(J^P) = I( ) Status: \*

OMITTED FROM SUMMARY TABLE

PEAKS ARE SEEN IN LAMBDA PI AND SIGMA PI SPECTRA IN THE REACTION PI+P TO K+ PI+ Y AT 1.7 GEV/C. ALSO THE Y POLARIZATION OSCILLATES IN THE SAME REGION.

SEE MILLER 70 FOR A DISCUSSION OF THIS STATE. HE SUGGESTS A POSSIBLE ALTERNATE EXPLANATION IN TERMS OF A REFLECTION OF N(1675) DECAY TO LAMBDA K. HOWEVER, SUCH AN EXPLANATION FOR THE K- SIGMA- PIO CHANNEL SEEMS UNLIKELY (SEE PAN 70) IN TERMS OF KNOWN DELTA(1650) DECAY INTO SIGMA K. IN ADDITION SUCH REFLECTIONS WOULD ALSO HAVE TO ACCOUNT FOR THE OSCILLATION OF THE Y POLARIZATION IN THE 1480 MASS REGION.

HANSON 71, WITH FEWER DATA THAN PAN 70, CAN NEITHER CONFIRM NOR DENY THE EXISTENCE OF THIS STATE. MASTZ 75 SEES NO STRUCTURE IN THIS MASS REGION IN K- P TO LAMBDA PIO.

ENGELEN 81 PERFORM A MULTI-CHANNEL ANALYSIS OF K-P -> KO PI- P AT 4.2 GEV/C. THEY OBSERVE A 3.5 STD. DEV SIGNAL AT 1480 MEV IN P KBAR WHICH CANNOT BE EXPLAINED AS A REFLECTION OF ANY COMPETING CHANNEL.

Σ(1480) MASS (MeV) (PROD. EXP.)

Table with 4 columns: M, MASS (MeV), PROD. EXP., REFERENCE. Rows include MALAMUD (1479.0, 10.0, PAN, 70 HBC + PI+P TO K PI LAM), SHAFFER (1465.0, 15.0, PAN, 70 HBC + PI+P TO K PI SIG), HUNGERBUHLER (1485.0, 10.0, CLINE, 73 MPWA K-D TO LM PI-P), and ENGELEN (120(1480.0), 10.0, ENGELEN, 80 HBC + K-P TO KO PI- P).

Σ(1480) WIDTH (MeV) (PROD. EXP.)

Table with 4 columns: W, WIDTH (MeV), PROD. EXP., REFERENCE. Rows include PAN (31.0, 15.0, PAN, 70 HBC + PI+P TO K PI LAM), ENGELEN (40.0, 20.0, PAN, 70 HBC + PI+P TO K PI SIG), and ENGELEN (120, 80.0, ENGELEN, 80 HBC + K-P TO KO PI- P).

Σ(1480) PARTIAL DECAY MODES (PROD. EXP.)

Table with 3 columns: P, DECAY MODE, DECAY MASSES. Rows include P1 Σ(1480) -> N K̄ (958+ 494), P2 Σ(1480) -> Δ π (1116+ 140), and P3 Σ(1480) -> Σ π (1189+ 140).

Σ(1480) BRANCHING RATIOS (PROD. EXP.)

Table with 3 columns: R, BRANCHING RATIO, PROD. EXP. Rows include R1 Σ(1480) -> (Σ π)/(Δ π) (0.82, 0.51, PAN, 70 HBC +), R2 Σ(1480) -> (p K̄^0)/(Δ π) (0.36, 0.25, PAN, 70 HBC +), and R3 Σ(1480) -> (N K̄)/total (SMALL, CLINE, 73 MPWA K-D TO LM PI- P).

REFERENCES FOR Σ(1480) (PROD. EXP.)

PAN 70 PR D2, 49
CLINE 73 LNC 6 205
ENGELEN 80 NP B167 61
+FORMAN, K.O., HAGOPIAN, SELOVE (PENNY)
CLINE, LAUMANN, MAPP (WISC) IJP
+HEINEN, KITTEL, METZGER-(CNJM-AMST+CERN-OXF)

PAPERS NOT REFERRED TO IN DATA LISTINGS

YU-LI PA 69 PRL 23 806
YU-LI PA 69 PRL 23 808
HANSON 71 PR D4 1296
MAST 75 PR D11 3078
YU-LI PAN, F L FORMAN (PENNY) I
YU-LI PAN, F L FORMAN (PENNY) I
D H MILLER (REVIEW TALK) (PURD)
+KALMUS, LOUIE (LBL) I
+ALSTON-GARNJOST, BANGERTER (LBL)

Σ(1560) BUMPS

I(J^P) = I( ) Status: \*\*

OMITTED FROM SUMMARY TABLE

THIS ENTRY LISTS PEAKS REPORTED IN MASS SPECTRA AROUND 1560 MEV WITHOUT IMPLYING THAT THEY ARE NECESSARILY RELATED.

DIONISI 78 OBSERVE A 6 STD. DEVIATION ENHANCEMENT AT 1553 MEV IN THE CHARGED (LAMBDA/SIGMA PI) MASS SPECTRA FROM K-P -> LAMBDA/SIGMA PI K KBAR AT 4.2 GEV/C. IN A CERN ISR EXPERIMENT, LOCKMAN 78 REPORT A NARROW 6 STD. DEVIATION ENHANCEMENT AT 1572 MEV IN THE LAMBDA PI-/PI- SYSTEMS FROM THE REACTION P P -> LAMBDA PI- PI- + ANYTHING AT C.M. ENERGIES OF 53 AND 62 GEV.

THESE ENHANCEMENTS ARE UNLIKELY TO BE ASSOCIATED WITH THE SIGMA(1580) (WHICH HAS NOT BEEN CONFIRMED BY SEVERAL RECENT EXPERIMENTS -- SEE THE LISTINGS BELOW).

CARROLL 76 OBSERVE A BUMP AT 1550 MEV (AS WELL AS AT 1580 MEV) IN THE K-N 1:1 TOTAL CROSS SECTION, BUT UNCERTAINTIES IN CROSS SECTION MEASUREMENTS OUTSIDE THE MASS RANGE OF THE EXPERIMENT PRECLUDE ESTIMATING ITS SIGNIFICANCE.

SEE ALSO MEADOWS 80 FOR A REVIEW OF THIS STATE.

Σ(1560) MASS (MeV) (PROD. EXP.)

Table with 4 columns: M, MASS (MeV), PROD. EXP., REFERENCE. Rows include DIONISI (121 1553.0, 7.0, DIONISI, 78 HBC +- K-P TO Y\* K KBAR) and LOCKMAN (40 1572.0, 4.0, LOCKMAN, 78 SPEC +- PP TO L PI PI X).

Σ(1560) WIDTH (MeV) (PROD. EXP.)

Table with 4 columns: W, WIDTH (MeV), PROD. EXP., REFERENCE. Rows include DIONISI (121 79.0, 30.0, DIONISI, 78 HBC +- K-P TO Y\* K KBAR), LOCKMAN (40 44.0, 6.0, LOCKMAN, 78 SPEC +- PP TO L PI PI X), and CARROLL (A OBSERVED WIDTH CONSISTENT WITH EXPERIMENTAL RESOLUTION).

Σ(1560) PARTIAL DECAY MODES (PROD. EXP.)

Table with 3 columns: P, DECAY MODE, DECAY MASSES. Rows include P1 Σ(1560) -> Δ π (1116+ 140) and P2 Σ(1560) -> Σ π (1189+ 140).

Σ(1560) BRANCHING RATIOS (PROD. EXP.)

Table with 3 columns: R, BRANCHING RATIO, PROD. EXP. Rows include R1 Σ(1560) -> Σ π/(Σ π + Δ π) (0.35, 0.12, DIONISI, 78 HBC +- K-P TO Y\* K KBAR) and R2 Σ(1560) -> (Δ π)/total (SEEN, LOCKMAN, 78 SPEC +- PP TO L PI PI X).



For notation, see key on page 91.

## Baryon Full Listings

$\Sigma(1620)$ ,  $\Sigma(1660)$ ,  $\Sigma(1670)$

### $\Sigma(1620)$ BRANCHING RATIOS (PROD. EXP.)

$\Sigma(1620) \rightarrow (\Lambda \pi \pi)/(\Lambda \pi)$	(P4)/(P3)
R1 14 (2.5) APPROX	BLUMENFELD 69 HBC +
$\Sigma(1620) \rightarrow (N \bar{K})/(\Lambda \pi)$	(P1)/(P2)
R2 (0.0) (0.1)	CRENNELL 68 DBC +
R2 0.4 0.4	AMMANN 70 DBC K-P 4.5 GEV/C
$\Sigma(1620) \rightarrow (\Lambda \pi)/\text{total}$	(P2)
R3 LARGE	CRENNELL 68 DBC --
$\Sigma(1620) \rightarrow (\Sigma(1385) \pi)/(\Lambda \pi)$	(P3)/(P2)
R4 (0.2) (0.1)	CRENNELL 68 DBC --
R4 (0.3) OR LESS CL=.95	AMMANN 70 DBC K-P 4.5 GEV/C
$\Sigma(1620) \rightarrow (\Sigma \pi)/(\Lambda \pi)$	(P5)/(P2)
R5 (1.1)(.95 PC UPPER LIMIT)	AMMANN 70 DBC K-N 4.5 GEV/C
$\Sigma(1620) \rightarrow (\Lambda(1405) \pi)/(\Lambda \pi)$	(P6)/(P2)
R6 0.7 0.4	AMMANN 70 DBC K-P 4.5 GEV/C

### REFERENCES FOR $\Sigma(1620)$ (PROD. EXP.)

CRENNELL 68 PRL 21 648	-DELANEY, FLAMINIO, KARSHON, + (BNL-CUNY) I
BLUMENFELD 69 PL 298 58	BLUMENFELD, KALBFLEISCH (BNL) I
CRENNELL 69 LUND PAPER 183	+KARSHON, LAI, ONEIL, SCARR, - (BNL-CUNY) I
RESULTS ARE QUOTED IN LEVI SETTI 69.	
AMMANN 70 PRL 24 327	- GARFINKEL, CARMONY, GUTAY, + (PURD-IND)
ALSO 73 PR D7 1345	AMMANN, CARMONY, GARFINKEL, (PURD-IUPUI)
	PAPERS NOT REFERRED TO IN DATA LISTINGS
ARMENTER 68 NP 88 183	ARMENTEROS, BAILLON + (CERN-HEID-SACL)
LEVISETTI 69 LUND CONF	R LEVI SETTI (RAPPORTEUR) (EPJ)
TRIPP 69 UCRL 19361	R D TRIPP (LRL)
ARMENTER 70 DUKE 123	ARMENTEROS, BAILLON + (CERN-HEID-SACL)
MILLER 70 DUKE 229	D H MILLER (REVIEW TALK) (PURD)
SABRE 70 NP 816 201	SABRE COLLAB. (SACL-AMST-BGNA-REHO-EPOL)
HUNGERBU 74 PR D10 2051	HUNGERBUHLER, MAJKA, + (YALE-FNAL-BNL-PITT)

$\Sigma(1660) P_{11}$

$I(J^P) = 1(\frac{1}{2}^-)$  Status: \*\*\*

FOR RESULTS PUBLISHED BEFORE 1974 (THEY ARE NOW OBSOLETE), SEE OUR 1982 EDITION (PHYSICS LETTERS 1118). ALL THE REFERENCES HAVE BEEN RETAINED.

### $\Sigma(1660)$ MASS (MeV)

M	1670.0 (20.0)	KANE 74 DPWA K-P TO PI SIG
M A	(1660.0) (30.0)	BAILLON 75 IPWA KBAR N TO LAM PI
M A	FROM SOLUTION 1 OF BAILLON 75,	NOT PRESENT IN SOLUTION 2
M B	(1671.0) (2.0)	PONTE 75 DPWA 0 K- P TO LAM PI
M B	FROM SOLUTION 2 OF PONTE 75,	NOT PRESENT IN SOLUTION 1
M	1668.0 (25.0)	VANHORN 75 DPWA 0 K- P TO LAM PI
M C	1565. OR 1597.	MARTIN 77 DPWA KBAR N MULTICHNL
M C	THE TWO ENTRIES FOR MARTIN 77	CORRESPOND TO EXTRACTION OF RESONANCE
M C	PARAMETERS FROM THE T-MATRIX POLE AND FROM A 3-4 FIT, RESPECTIVELY.	
M	1676.0 (15.0)	RLIC 77 DPWA KBAR N MULTICHNL
M	1679.0 (10.0)	ALSTON 78 DPWA KBAR N ELASTIC
M	1670.0 (10.0)	GOPAL 80 DPWA KBAR N ELASTIC
M D	1665.1 (11.2)	KOTISO 85 DPWA K-P TO SIGMA PI
M D	BUT THE EVIDENCE IS WEAK.	

### $\Sigma(1660)$ WIDTH (MeV)

W	250.0 (110.0)	KANE 74 DPWA K-P TO PI SIG
W A	(80.0) (40.0)	BAILLON 75 IPWA KBAR N TO LAM PI
W B	(81.0) (10.0)	PONTE 75 DPWA 0 K- P TO LAM PI
W C	230.0 (165.0) (60.0)	VANHORN 75 DPWA 0 K- P TO LAM PI
W	202. OR 217.	MARTIN 77 DPWA KBAR N MULTICHNL
W	120.0 (20.0)	RLIC 77 DPWA KBAR N MULTICHNL
W	38.0 (10.0)	ALSTON 78 DPWA KBAR N ELASTIC
W	152.0 (20.0)	GOPAL 80 DPWA KBAR N ELASTIC
W D	81.5 (22.2)	KOTISO 85 DPWA K-P TO SIGMA PI

### $\Sigma(1660)$ PARTIAL DECAY MODES

P1	$\Sigma(1660) \rightarrow N \bar{K}$	DECAY MASSES
P2	$\Sigma(1660) \rightarrow \Sigma \pi$	958+ 494
P3	$\Sigma(1660) \rightarrow \Lambda \pi$	1189+ 140
		1116+ 140

### $\Sigma(1660)$ BRANCHING RATIOS

$\Sigma(1660)$ in $N \bar{K} \rightarrow \Sigma \pi$	SQRT(P1*P2)
R1 -0.11 (0.01)	KANE 74 DPWA K-P TO PI SIG
R1 NOT SEEN	HEPP2 76 DPWA 0 K- NUC TO SIG PI
R1 C -0.34 OR -0.37	MARTIN 77 DPWA KBAR N MULTICHNL
R1 -0.16 (0.03)	RLIC 77 DPWA KBAR N MULTICHNL
R1 D -0.13 (0.04)	KOTISO 85 DPWA K-P TO SIGMA PI

### $\Sigma(1660) \rightarrow (N \bar{K})/\text{total}$

R2 C 0.27 OR 0.29	MARTIN 77 DPWA KBAR N MULTICHNL
R2 E LESS THAN 0.04	RLIC 77 DPWA KBAR N MULTICHNL
R2 (N KBAR)/TOTAL FROM RLC 77 IS SUPERSEDED BY GOPAL 80.	
R2 0.10 (0.05)	ALSTON 78 DPWA KBAR N ELASTIC
R2 0.12 (0.03)	GOPAL 80 DPWA KBAR N ELASTIC

### $\Sigma(1660)$ in $N \bar{K} \rightarrow \Lambda \pi$

R3 A (-0.04) (0.02)	BAILLON 75 IPWA KBAR N TO LAM PI
R3 B (-0.16) (0.01)	PONTE 75 DPWA 0 K- P TO LAM PI
R3 0.12 (0.12) (0.04)	VANHORN 75 DPWA 0 K- P TO LAM PI
R3 C -0.10 OR -0.11	MARTIN 77 DPWA KBAR N MULTICHNL
R3 LESS THAN 0.04	RLIC 77 DPWA KBAR N MULTICHNL

### REFERENCES FOR $\Sigma(1660)$

ARMENTER 70 DUKE 123	ARMENTEROS, BAILLON, + (CERN-HEID)IJP
KIM 71 PRL 27 356	J K KIM (HARV)IJP
ALSO 70 DUKE 161	J K KIM (HARV)IJP
HART 73 PURDUE CONF. 311	-RICE, BACASTOW, FUNG, + (TENN-UCR-MASA-BUFF)IJP
LEA 73 NP 856 77	-MARTIN, MOORHOUSE + (RHEL-LOUC-GLAS-AARH)IJP
KANE 74 LBL-2452	D K KANE (LBL)IJP
BAILLON 75 NP 819 39	B BAILLON, P J LITCHFIELD (CERN-RHEL)IJP
PONTE 75 PR D12 2597	-HERTZBACH, BUTTON-SHAFER + (MASA-TENN)UCR)IJP
VANHORN 75 NP 887 145	A J VAN HORN (LBL)IJP
ALSO 75 NP 887 157	A J VAN HORN (LBL)IJP
HEPP2 76 PL 658 487	-BRAUN, GRIMM, STROBELE, THOL + (CERN-HEID-IMP)IJP
MARTIN 77 NP 8127 349	MARTIN, PIDCOCK, MOORHOUSE (LOUC-GLAS)IJP
ALSO 77 NP 8126 266	MARTIN, PIDCOCK (LOUC)IJP
ALSO 77 NP 8126 285	MARTIN, PIDCOCK (LOUC)IJP
RLIC 77 NP 8119 362	GOPAL, ROSS, VAN HORN, MCPHERSON + (LOIC-RHEL)IJP
ALSTON 78 PR D18 182	+KENNEY, POLLARD, ROSS + (LBL-MTHO-CERN)IJP
ALSO 77 PRL 38 1007	ALSTON-GARNJOST, KENNEY, + (LBL-MTHO-CERN)IJP
GOPAL 80 TORONTO CONF 159	G P GOPAL (RHEL)IJP
KOTISO 85 NP 4433 619	-SAI, YAMAMOTO, KOFLER (TOKY-MASA)

### NOTE ON THE $\Sigma(1670)$

**Production experiments:** The measured  $\Sigma\pi/\Sigma\pi\pi$  branching ratio for produced  $\Sigma(1670)$ 's is strongly dependent on momentum transfer. This was first discovered by EBERHARD 69, who suggested that there exist two  $\Sigma$  resonances with the same mass and quantum numbers: one with a large  $\Sigma\pi\pi$  [mainly  $\Lambda(1405)\pi$ ] decay mode produced peripherally, and the other with a large  $\Sigma\pi$  decay mode produced at larger angles. These results were confirmed by AGUILAR-BENITEZ 70, ASPELL 74, ESTES 74, and TIMMERMANS 76. The most likely quantum numbers for both the  $\Sigma\pi$  and the  $\Lambda(1405)\pi$  states are  $D_{13}$ . There is also possibly a third  $\Sigma$ , the  $\Sigma(1690)$  in the Listings, the main evidence for which is a large  $\Lambda\pi/\Sigma\pi\pi$  branching ratio. These topics have been reviewed by EBERHARD 73 and by MILLER 70.

**Formation experiments:** Two states are also observed near this mass in formation experiments. One of these, the  $D_{13} \Sigma(1670)$ , has the same quantum numbers as those observed in production and has a large  $\Sigma\pi/\Sigma\pi\pi$  branching ratio. It may well be the  $\Sigma(1670)$  produced at larger angles (see TIMMERMANS 76). The other state, the  $P_{11} \Sigma(1660)$ , has different quantum numbers from those seen in production, and its  $\Sigma\pi/\Sigma\pi\pi$  branching ratio is unknown. Thus its relation to the produced  $\Sigma(1670)$ 's remains obscure.













For notation, see key on page 91.

Baryon Full Listings

Σ(1915), Σ(1940)

Σ(1915) PARTIAL DECAY MODES

Table with columns P1, P2, P3, P4, P5 and decay modes like Σ(1915) → N K-bar, Σ(1915) → Δ π, Σ(1915) → Σ π, Σ(1915) → Σ(1385) π, P-wave, Σ(1915) → Σ(1385) π, F-wave.

Σ(1915) BRANCHING RATIOS

Table with columns R1, R2, R3, R4, R5 and branching ratios for Σ(1915) → (N K-bar)/total, Σ(1915) in N K-bar → Δ π, Σ(1915) in N K-bar → Σ π.

Table with columns R2, R3, R4, R5 and branching ratios for Σ(1915) in N K-bar → Δ π, Σ(1915) in N K-bar → Σ π.

Table with columns R3, R4, R5 and branching ratios for Σ(1915) in N K-bar → Σ π.

Table with columns R4 and branching ratios for Σ(1915) in N K-bar → Σ(1385) π, P-wave.

Table with columns R5 and branching ratios for Σ(1915) in N K-bar → Σ(1385) π, F-wave.

REFERENCES FOR Σ(1915)

SMART 66 PRL 17 556 W M SMART, A KERNAN, G E KALMUS, R P ELY (LRL)JJP
ARMENTER 67 PL 249 198 ARMENTEROS, FERRO-LUZZI+ (CERN-HEID+SACL)
ARMENTERI 67 NP 83 592 ARMENTEROS, FERRO-LUZZI+ (CERN-HEID+SACL)
CONFORTO 68 NP 88 265 -HARSEN, LASINSKI, + (CHIC+HEJD)
SMART 68 PR 169 1330 W M SMART (LRL)JJP
BERTHON 70 NP 82 1674 -RANGAN, VRANA, + (CDFE+RHEL+SACL)JJP
BERTHON 70 NP 824 417 -VRANA, BUTTERWORTH, + (CDFE+RHEL+SACL)JJP
BRICMAN 70 PL 338 511 -FERRO-LUZZI, LAGNAUX (CERN)
COX 70 NP 819 61 +ISLAM, COLLEY, + (BIRM+EDIN+GLAS+LOIC)JJP
KANE 72 PR 5 1573 A BARBARO-GALTIERI (LRL)JJP
LITCFHIE 70 NP 822 269 P J LITCFHIE (RHEL)JJP
CONFORTO 71 NP 834 41 -LEVI SETTI, LASINSKI, OBERLACK++ (EFI-HEID)JJP
LITCFHIE 71 NP 830 125 LITCFHIE, D, ...+LEBQUOY, ... (RHEL+CDEF+SACL)JJP
KANE 72 PR 5 1573 A BARBARO-GALTIERI (LRL)JJP
DEVENISH 74 NP 881 330 DEVENISH, FROGGATT, MARTIN (DESY+NORD+LOUC)
KANE 74 LBL-2452 D F KANE (LBL)JJP
BAILLON 75 NP 894 39 P J BAILLON, P J LITCFHIE (CERN+RHEL)JJP
HEMINGWA 75 NP 891 12 HEMINGWAY, GADES, HARMSEN+ (CERN+HEID+MIM)JJP
VANHORN 75 NP 887 145 A J VAN HORN (LBL)JJP
ALSO 75 NP 887 157 A J VAN HORN (LBL)JJP
BELLEFON 76 NP 8109 129 DE BELLEFON, BERTHON (CDFE)JJP
CORDEN 76 NP 8104 382 +COX, DARTNELL, KENYON, ONEALE, SUMOROK+ (BIRM)JJP
CORDEN 77 NP 8125 61 +COX, KENYON, ONEALE, STUBBS, SUMOROK+ (BIRM)JJP
DECLAIS 77 CERN 77-16 +DUCHON, LOUVEL, PATRY, SEGUINOT+ (CAEN+CERN)JJP
MARTIN 77 NP 8127 349 MARTIN, PIDCOCK, MOORHOUSE (LOUC+BLAS)JJP
ALSO 77 NP 8126 266 MARTIN, PIDCOCK (LOUC)
ALSO 77 NP 8126 285 MARTIN, PIDCOCK (LOUC)JJP
RLIC 77 NP 8119 362 GOPAL, ROSS, VAN HORN, MCPHERSON+ (LOIC+RHEL)JJP
ALSTON 78 PR D18 182 +KENNEY, POLLARD, ROSS+ (LBL+MTHO+CERN)JJP
ALSO 77 PRL 38 1007 ALSTON-GARNJUST, KENNEY, + (LBL+MTHO+CERN)JJP
CAMERON 78 NP B143 189 +FRANEK, GOPAL, BACON, BUTTERWORTH+ (RHEL+LOIC)JJP
GOPAL 80 TORONTO CONF 159 6 P GOPAL (RHEL)JJP

1915 MeV REGION - PRODUCTION AND σTOTAL EXPPTS

I(J^P) = I( )

SEE THE NOTES TO THE SIGMA(1915) AND SIGMA(1940), WHICH IMMEDIATELY PRECEDE AND FOLLOW THIS ENTRY. HERE WE LIST ONLY PARAMETERS OF PEAKS SEEN IN CROSS SECTIONS AND INVARIANT-MASS DISTRIBUTIONS. THE CROSS SECTION PEAKS ARE ALMOST CERTAINLY ASSOCIATED WITH THE F15 SIGMA(1915) SEEN IN PARTIAL-WAVE ANALYSES. THE INVARIANT-MASS PEAKS SEEM MORE LIKELY TO BE ASSOCIATED WITH THE D13 SIGMA(1940).

Σ(1915) MASS (MeV) (PROD. EXP.)

Table with columns M, B, M and cross-section peaks for Σ(1915).

Table with columns M, B, M and invariant-mass-distribution peaks for Σ(1915).

Σ(1915) WIDTH (MeV) (PROD. EXP.)

Table with columns W, B, W and cross-section peaks for Σ(1915).

Σ(1915) PARTIAL DECAY MODES (PROD. EXP.)

Table with columns P1, P2, P3, P4 and decay masses for Σ(1915).

Σ(1915) BRANCHING RATIOS (PROD. EXP.)

Table with columns R1, R2, R3, R4, R5 and branching ratios for Σ(1915).

Table with columns R2 and branching ratios for Σ(1915) → (N K-bar)/(Σ π).

Table with columns R3 and branching ratios for Σ(1915) → (Δ π)/(Σ π).

Table with columns R4 and branching ratios for Σ(1915) → (Σ K)/total.

REFERENCES FOR Σ(1915) (PROD. EXP.)

BOCK 65 PL 17 166 +COOPER, FRENCH, KINSON, + (CERN+SACL) I
BOCK 66 PRL 16 1228 +GIACOMELLI, KYCIA, LEONTIC, LI, LUNDBY, + (BNL) I
BUGG 68 PR 168 1466 +GILMORE, KNIGHT, DAVIES- (BIRM+CAVE+RHEL)I
PRIMER 68 PRL 20 610 +GOLDBERG, JAEGER, BARNES, DORNAN + (SYRA+BNL)
SUPERSEDED BY BARNES 69 AND AGUILAR-BENITEZ 70.
BARNES 69 PRL 22 479 +FLAMINIO, MONTANET, SAMIOS + (BNL+SYRA)
AGUILAR 70 PRL 25 58 AGUILAR-BENITEZ, BARNES, + (BNL+SYRA)
BRICMAN 70 PL 318 152 +FERRO LUZZI, PERREAU, + (CERN+CAEN+SACL)
COOL 70 PR D1 1887 +GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL) I
DADO 72 PRL 29 1695 +BIRMAN, GOLDBERG, WEISS (IATF)JP
BRIEFEL 77 PR D16 2706 +GOUREVITCH, CHANG+ (BRAN+UMD+SYRA+TUFT)
FERRER 81 NP B178 373 +TREILLE, RIVET, VOLTE+ (CERN+CDEF+EPO+LALO)

Σ(1940) D13

I(J^P) = I(3/2^-) Status: \*\*\*

FOR RESULTS PUBLISHED BEFORE 1974 (THEY ARE NOW OBSOLETE), SEE OUR 1982 EDITION (PHYSICS LETTERS 111B). ALL THE REFERENCES HAVE BEEN RETAINED.

SOME, NOT ALL, PARTIAL-WAVE ANALYSES SUGGEST A STATE IN THIS REGION. IT IS PERHAPS ASSOCIATED WITH THE BUMPS SEEN IN PRODUCTION EXPERIMENTS NEAR THIS MASS (SEE THE PRECEEDING ENTRY). THIS STATE IS NOT REQUIRED IN K- NEUTRON TO (PI SIGMA)- ANALYSIS OF GOPAL 77. KBAR N ANALYSIS (GOPAL 80) WITH K- NEUTRON ELASTIC DATA DOES NOT REQUIRE THIS STATE.

Σ(1940) MASS (MeV)

Table with columns M, B, M and invariant-mass (MeV) for Σ(1940).







For notation, see key on page 91.

Baryon Full Listings

Σ(2250), Σ(2455), Σ(2620), Σ(3000)

Σ(2250) PARTIAL DECAY MODES (PROD. EXP.)

Table with columns P1, Σ(2250) → N K-bar, DECAY MASSES, and values like 938- 494, 1116- 135, 1189- 140.

Σ(2250) BRANCHING RATIOS (PROD. EXP.)

Table with columns R1, Σ(2250) → (N K-bar)/total, Σ(2250) in N K-bar → Δ π, Σ(2250) in N K-bar → Σ π, Σ(2250) → (N K-bar)/(Σ π), Σ(2250) → (Δ π)/(Σ π), Σ(2250) in N K-bar → Σ(1530)0 K0, Σ(2250) → (N K-bar)/total.

REFERENCES FOR Σ(2250) (PROD. EXP.)

BLANPIED 65 PRL 14 741 +GREENBERG, HUGHES, KITCHING, + (YALE-CEA)
BOCK 65 PL 17 166 +COOPER, FRENCH, KINSON, + (CERN+SACL)
BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (RHCL+BIRM+CAVE) I

Σ(2455) BUMPS

I(JP) = 1( ) Status: \*\*

OMITTED FROM SUMMARY TABLE

THERE IS ALSO SOME SLIGHT EVIDENCE FOR Y\* STATES IN THIS MASS REGION FROM THE REACTION GAMMA + P TO K + MISSING MASS -- SEE GREENBERG 68.

Σ(2455) MASS (MeV) (PROD. EXP.)

Table with columns M, Σ(2455) MASS (MeV), DECAY MASSES, and values like 2455.0, 7.0, 68 CNTR, 70 CNTR, 938- 494.

Σ(2455) WIDTH (MeV) (PROD. EXP.)

Table with columns W, Σ(2455) WIDTH (MeV), DECAY MASSES, and values like 100.0, 20.0, 68 CNTR, 70 CNTR, 938- 494.

Σ(2455) PARTIAL DECAY MODES (PROD. EXP.)

Table with columns P1, Σ(2455) → N K-bar, DECAY MASSES, and values like 938+ 494.

Σ(2455) BRANCHING RATIOS (PROD. EXP.)

Table with columns R1, Σ(2455) → (N K-bar)/total, Σ(2455) in N K-bar → Δ π, Σ(2455) in N K-bar → Σ π, Σ(2455) → (N K-bar)/(Σ π).

REFERENCES FOR Σ(2455) (PROD. EXP.)

BUGG 68 PR 168 1466 +GILMORE, KNIGHT, + (RHCL+BIRM+CAVE) I
ABRAMS 70 PR D1 1917 +COOL, GIACOMELLI, KYCIA, LEONTIC, + (BNL) I

Σ(2620) BUMPS

I(JP) = 1( ) Status: \*\*

OMITTED FROM SUMMARY TABLE

Σ(2620) MASS (MeV) (PROD. EXP.)

Table with columns M, Σ(2620) MASS (MeV), DECAY MASSES, and values like 2620.0, 15.0, 70 CNTR, 75 DBC, 2542.0, 22.0.

Σ(2620) WIDTH (MeV) (PROD. EXP.)

Table with columns W, Σ(2620) WIDTH (MeV), DECAY MASSES, and values like 175.0, 221.0, 81.0, 70 CNTR, 75 DBC.

Σ(2620) PARTIAL DECAY MODES (PROD. EXP.)

Table with columns P1, Σ(2620) → N K-bar, DECAY MASSES, and values like 938- 494.

Σ(2620) BRANCHING RATIOS (PROD. EXP.)

Table with columns R1, Σ(2620) → (N K-bar)/total, Σ(2620) in N K-bar → Δ π, Σ(2620) in N K-bar → Σ π, Σ(2620) → (N K-bar)/(Σ π).

REFERENCES FOR Σ(2620) (PROD. EXP.)

ABRAMS 67 PRL 19 678 +COOL, GIACOMELLI, KYCIA, LEONTIC, LI, + (BNL)
SUPERSEDED BY ABRAMS 70.
ABRAMS 70 PR D1 1917 +COOL, GIACOMELLI, KYCIA, LEONTIC, + (BNL) I

Σ(3000) BUMPS

I(JP) = 1( ) Status: \*

OMITTED FROM SUMMARY TABLE

ENHANCEMENT IN LAMBDA PI AND KBAR N INVARIANT MASS SPECTRA AND IN MISSING MASS OF NEUTRALS RECOLLING AGAINST K0.

Σ(3000) MASS (MeV) (PROD. EXP.)

Table with columns M, Σ(3000) MASS (MeV), DECAY MASSES, and values like 3000.0, 66 HBC, 0 PI-P, 7.91 GEV/C.

Σ(3000) PARTIAL DECAY MODES (PROD. EXP.)

Table with columns P1, Σ(3000) → N K-bar, DECAY MASSES, and values like 938- 494, 1116- 140.



# Baryon Full Listings

$\Sigma(3000)$ ,  $\Sigma(3170)$ ,  $\Xi$ 's

## REFERENCES FOR $\Sigma(3000)$ (PROD. EXP.)

EHRlich 66 PR 152 1194 R EHRlich, M SELOVE, H YUTA (PENN+BNL) 1

**$\Sigma(3170)$  BUMPS**

 $I(J^P) = 1( )$  Status: \*
 

OMITTED FROM SUMMARY TABLE

SEEN BY AMIRZADEH 79 AS A NARROW 6.5 STD. DEV. ENHANCEMENT IN THE REACTION  $K^-p \rightarrow Y^* \pi^-$  USING DATA FROM TWO INDEPENDENT HIGH STATISTICS BUBBLE CHAMBER EXPERIMENTS AT 8.25 AND 6.5 GEV/C. THE DOMINANT DECAY MODES ARE INTO MULTI-BODY, MULTI-STRANGE FINAL STATES AND THE PRODUCTION IS VIA  $I=3/2$  BARYON EXCHANGE.  $I=1$  IS FAVORED.

NOT SEEN IN A  $K^-p$  EXPERIMENT IN LASS AT 11 GEV/C (ASTON 85).

### $\Sigma(3170)$ MASS (MeV) (PROD. EXP.)

M 35 3170.0 5.0 AMIRZAD 79 HBC + K-P TO Y\* PI-

### $\Sigma(3170)$ WIDTH (MeV) (PROD. EXP.)

W A 35 (20.0) OR LESS AMIRZAD 79 HBC + K-P TO Y\* PI-  
 W A OBSERVED WIDTH CONSISTENT WITH EXPERIMENTAL RESOLUTION.

### $\Sigma(3170)$ PARTIAL DECAY MODES (PROD. EXP.)

	DECAY MASSES
P1	$\Sigma(3170) \rightarrow \Lambda \bar{K} \bar{K} + \pi^+$ s
P2	$\Sigma(3170) \rightarrow \Sigma \bar{K} \bar{K} + \pi^+$ s
P3	$\Sigma(3170) \rightarrow \Xi K + \pi^+$ s

### $\Sigma(3170)$ BRANCHING RATIOS (PROD. EXP.)

$\Sigma(3170) \rightarrow (\Lambda \bar{K} \bar{K} + \pi^+)$ /total	(P1)
R1 SEEN	AMIRZAD 79 HBC + K-P TO Y* PI-
$\Sigma(3170) \rightarrow (\Sigma \bar{K} \bar{K} + \pi^+)$ /total	(P2)
R2 SEEN	AMIRZAD 79 HBC + K-P TO Y* PI-
$\Sigma(3170) \rightarrow (\Xi K + \pi^+)$ /total	(P3)
R3 SEEN	AMIRZAD 79 HBC + K-P TO Y* PI-

### REFERENCES FOR $\Sigma(3170)$ (PROD. EXP.)

AMIRZAD 79 PL 89B 125 AMIRZADEH+ (BIRM-CERN+GLAS+MSU+LBNL-CAMB+J)  
 ALSO 80 TORONTO CONF. 263 J B KINSON+ (BIRM-CERN+GLAS+MSU+LBNL-CAMB+J)  
 ASTON 85 PR 032 2270 + LASS COLLABORATION (SLAC+CARL-CNRC+CINC)

## NOTE ON $\Xi$ RESONANCES

Not much is known about  $\Xi$  resonances. This is because (1) they can only be produced as a part of a final state, and so the analysis is more complicated than if direct formation were possible, (2) they are produced with small cross sections (typically a few  $\mu\text{b}$ ), and (3) the final states are topologically complicated and difficult to study with electronic techniques. Thus our early information about  $\Xi$  resonances came entirely from bubble chamber experiments, where the numbers of events are small, and the best information is still from bubble chamber experiments.

The accompanying table gives our evaluation of the present status of the  $\Xi$  resonances. Until fairly recently only the  $\Xi(1530)$  was really well established. However, the late 1970's saw a major improvement with results from the large  $K^-p$  bubble chamber experiment at 4.2

Table 1. The status of the  $\Xi$  resonances. Only those with an overall status of \*\*\* or \*\*\*\* are included in the Baryon Summary Table.

Particle	$L_{21-2J}$	Overall status	Status as seen in --				
			$\Xi\pi$	$\Lambda\bar{K}$	$\Sigma\bar{K}$	$\Xi(1530)\pi$	Other channels
$\Xi(1318)$	$P_{11}$	****					
$\Xi(1530)$	$P_{13}$	****	****				
$\Xi(1630)$		*	*				
$\Xi(1680)$		**		*	**		
$\Xi(1820)$	13	***	*	***	**	***	
$\Xi(1940)$		**	**			**	
$\Xi(2030)$	1	***		**	***		
$\Xi(2120)$		*		*			3-body decays
$\Xi(2250)$		**					3-body decays
$\Xi(2370)$	1	**					3-body decays
$\Xi(2500)$		*		*	*		3-body decays

\*\*\*\* Good, clear, and unmistakable.  
 \*\*\* Good, but in need of clarification or not absolutely certain.  
 \*\* Not established; needs confirmation.  
 \* Evidence weak; could disappear.

GeV/c (GAY 76, HEMINGWAY 77). The  $\Xi(1820)$  and  $\Xi(2030)$  were firmly established as narrow states (widths of about 20 MeV), and the spin of the  $\Xi(1820)$  was found to be  $3/2$  (TEODORO 78).

Since then, however, not much has changed, although there is some improved evidence for the  $\Xi(2250)$  and the  $\Xi(2370)$ . There is probably at least one other  $\Xi$  in the 1850-2000-MeV region, and there are indications of several others above 2000 MeV. Indeed, there should be many  $\Xi$  resonances below 2500 MeV, and the broad (and not completely established)  $\Xi(1940)$  could well be a mixture of several of them. For now we are forced to group together differing observations and await new results. The differences among experiments are shown in ideograms in the Listings; the wait for new results is almost certain to be a long one.

In the last few years, results from experiments using electronic methods have appeared. BIAGI 81 used the CERN hyperon beam to study inclusive  $\Lambda\bar{K}$  and  $\Xi\pi$  mass spectra from 102 and 135 GeV/c  $\Xi^-$  incident on hydrogen and deuterium. They saw a large  $\Xi(1820)$  signal in  $\Lambda\bar{K}$  as well as a peak at about 1700 MeV, which might be associated with the threshold enhancement seen by DIONISI 78. The  $\Xi(1940)$  appears as a broad bump in the  $\Xi\pi$  mass spectrum, and there is a very clean  $\Xi(1530)$  signal. And Brookhaven multiparticle spectrometer measurements of  $K^-p \rightarrow K^+$  anything at 5 GeV/c (JENKINS 83) have seen all the established  $\Xi$  resonances and also the less well-established  $\Xi(2250)$ ,  $\Xi(2370)$ , and  $\Xi(2500)$ .

For a detailed review, see Meadows.<sup>1</sup>

## References

1. B.T. Meadows, in *Proceedings of the IV<sup>th</sup> International Conference on Baryon Resonances* (Toronto, 1980), ed. N. Isgur, p. 283.

For notation, see key on page 91.

# Baryon Full Listings

$\Xi^-, \Xi^0, \Xi(1530)$

## $\Xi$ BARYONS ( $S = -2, I = 1/2$ )

$\Xi^-$

$$I(J^P) = \frac{1}{2}(\frac{1}{2})$$

SEE STABLE PARTICLES.

$\Xi^0$

$$I(J^P) = \frac{1}{2}(\frac{1}{2})$$

SEE STABLE PARTICLES.

$\Xi(1530) P_{13}$

$$I(J^P) = \frac{1}{2}(\frac{3}{2}^+)$$
 Status: \*\*\*\*

THIS IS THE ONLY XI RESONANCE WHOSE PROPERTIES ARE ALL REASONABLY WELL KNOWN. SPIN-PARITY 3/2+ IS FAVORED BY THE DATA.  
WE DO NOT USE DETERMINATIONS OF THE MASS AND WIDTH HERE UNLESS THEY ARE ACCOMPANIED BY SOME DISCUSSION OF SYSTEMATICS AND RESOLUTION.

### $\Xi(1530)$ MASS (MeV)

M MIXED CHARGES		BERTANZA 62 HBC	-0 K-P 2.3 GEV/C
M 20(1535.0)		PJERROU 62 HBC	-0 K-P 1.8 GEV/C
M 55(1529.0)	(5.0)	BADIER 64 HBC	-0 K-P 3 GEV/C
M (1532.0)	(2.0)		
M- NEGATIVE CHARGE ONLY			
M- 38 1535.7	3.2	LONDON 66 HBC	- K-P 2.24 GEV/C
M- 354(1534.7)	(1.1)	BALTAY 72 HBC	- K-P 1.75 GEV
M- 185 1536.2	1.6	KIRSCH 72 HBC	- K-P 2.87GEV/C
M- 1535.3	2.0	ROSS2 73 HBC	- XI KBAR PI (PI)
M- 48(1540.0)	(3.0)	BERTHON 74 HBC	- QUASI 2 BODY CS
M- 1534.5	1.2	BELLEFOZ 75 HBC	- K-P TO XI- K PI
M- AVG 1535.18	0.84	SIXEL 79 HBC	0 INCL. K-P 10 GEV
M- FIT 1534.97	0.63	SIXEL 79 HBC	0 INCL. K-P 16 GEV
		W C	EXPERIMENTAL RESOLUTION OF 15 MEV NOT UNFOLDED.
		W D	2700 (12.8) (1.0) BAUBILLIE 81 HBC
		W D	FIT TO INCLUSIVE SPECTRUM. RESOLUTION (5 MEV) NOT UNFOLDED.
		W D	AVG 9.14 0.48 AVERAGE
MO NEUTRAL CHARGE ONLY			
MO 76 1528.7	1.1	LONDON 66 HBC	0 K-P 2.24 GEV/C
MO 59 1531.4	0.8	BADIER 72 HBC	0 K-P AT 3.95GEV/C
MO 1262 1532.0	0.4	BALTAY 72 HBC	0 K-P 1.75 GEV
MO 324 1531.3	0.6	BORENSTEI 72 HBC	0 K-P 2.26GEV/C
MO 286 1532.3	0.7	KIRSCH 72 HBC	0 K-P 2.87GEV/C
MO 1533.0	1.0	ROSS2 73 HBC	XI KBAR PI (PI)
MO 97(1533.6)	(1.4)	BERTHON 74 HBC	0 K-P AT 8.25 GEV
MO 1532.2	0.7	BELLEFOZ 75 HBC	0 K-P TO XI- K PI
MO 80(1527.0)	(6.0)	SIXEL 79 HBC	0 INCL. K-P 10 GEV
MO 100(1535.0)	(4.0)	SIXEL 79 HBC	0 INCL. K-P 16 GEV
MO A 2700(1532.1)	(0.6)	BAUBILLIE 81 HBC	0 K-P AT 8.25 GEV
MO A FIT TO INCLUSIVE SPECTRUM. RESOLUTION (5 MEV) NOT UNFOLDED.		BIAGI 81 SPEC	- HYPERON BEAM
MO 450(1530.0)	(1.0)	ASTON 85 LASS	K-P 11 GEV/C
MO 1244(1532.1)	(0.4)		
MO AVG 1531.78	0.34		
MO FIT 1531.80	0.31		

### $\Xi(1530)^- - \Xi(1530)^0$ MASS DIFFERENCE (MeV)

D	5.7	3.0	PJERROU 65 HBC	-0 1.8-1.95 GEV/C
D B	(7.0)	(4.0)	LONDON 66 HBC	-0 2.24 GEV/C
D	2.0	3.2	MERRILL 66 HBC	-0 1.7-2.7 GEV/C
D	2.7	1.0	BALTAY 72 HBC	-0 K-P 1.75 GEV
D B	(3.9)	(1.8)	KIRSCH 72 HBC	-0 K-P 2.87 GEV/C
D B	REDUNDANT WITH DATA IN MASS LISTING.			
D	2.92	0.91	AVERAGE	
D FIT	3.17	0.64	FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.0)	

### $\Xi(1530)$ WIDTH (MeV)

W MIXED CHARGES	20 (35.0)	OR LESS	BERTANZA 62 HBC	-0 K-P 2.3 GEV/C
W- NEGATIVE CHARGE ONLY				
W-	7.8	3.5	7.8 BALTAY 72 HBC	- K-P 1.75 GEV
W-	16.2	4.6	KIRSCH 72 HBC	- XI- PI0, XI0 PI-
W-	8.3	3.6	ROSS2 73 HBC	- XI KBAR PI (PI)
W-	9.6	2.8	BELLEFOZ 75 HBC	- K-P TO XI- K PI
W-	AVG 10.1	1.9	AVERAGE	
WO NEUTRAL CHARGE ONLY				
WO	7.0	2.0	SCHLEIN 63 HBC	0 1.8, 1.95 GEV/C
WO	7.0	7.0	BERGE 66 HBC	0 1.5-1.7 GEV/C
WO	8.5	3.5	LONDON 66 HBC	0 2.24 GEV/C
WO	11.0	2.0	BADIER 72 HBC	0 K-P AT 3.95GEV/C
WO	9.0	0.7	BALTAY 72 HBC	0 K-P 1.75 GEV
WO	8.4	1.4	BORENSTEI 72 HBC	0 XI- PI- MODE
WO	11.0	1.8	KIRSCH 72 HBC	0 XI- PI-
WO	9.1	2.4	ROSS2 73 HBC	0 XI KBAR PI (PI)
WO	9.5	1.2	BELLEFOZ 75 HBC	0 K-P TO XI- K PI
WO	80 (19.0)	(6.0)	SIXEL 79 HBC	0 INCL. K-P 10 GEV
WO	100 (14.0)	(5.0)	SIXEL 79 HBC	0 INCL. K-P 16 GEV
WO	EXPERIMENTAL RESOLUTION OF 15 MEV NOT UNFOLDED.			
WO	2700 (12.8)	(1.0)	BAUBILLIE 81 HBC	0 K-P AT 8.25 GEV
WO	FIT TO INCLUSIVE SPECTRUM. RESOLUTION (5 MEV) NOT UNFOLDED.			
WO	AVG 9.14	0.48	AVERAGE	

### $\Xi(1530)$ REAL PART OF POLE POSITION (MeV)

RE0	1531.6	0.4	LICHTENB 74	0 EXTRAP HABIBI73
RE-	1534.4	1.1	LICHTENB 74	- EXTRAP HABIBI73

### $\Xi(1530)$ IMAGINARY PART OF POLE POSITION (MeV)

IM0	4.45	0.35	LICHTENB 74	0 EXTRAP HABIBI73
IM-	3.9	1.75	3.9 LICHTENB 74	- EXTRAP HABIBI73

### $\Xi(1530)$ PARTIAL DECAY MODES

P1	$\Xi(1530) \rightarrow \Xi \pi$	DECAY MASSES
P2	$\Xi(1530) \rightarrow \Xi \gamma$	1321+ 140
		1321+ 0

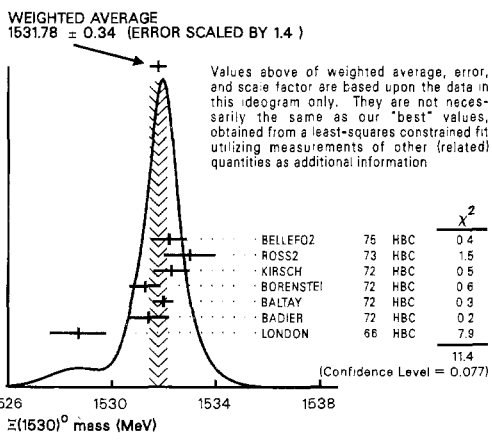
### $\Xi(1530)$ BRANCHING RATIOS

$\Xi(1530) \rightarrow (\Xi \gamma)/total$	
R1	(0.04) OR LESS CL.-90 KALBFLEI 75 HBC - K-P AT 2.18 GEV

### REFERENCES FOR $\Xi(1530)$

BERTANZA 62 PRL 9 180 +BRISSON, CONNOLLY, GOLDBERG, GRAY, + (BNL+SYRA) IJ  
 PJERROU 62 PRL 9 114 +PROMSE, SCHLEIN, SLATER, STORK, TICHO (UCLA) I  
 SCHLEIN 63 PRL 11 167 +CARMONY, PJERROU, SLATER, STORK, TICHO (UCLA) IJP  
 BADIER 64 DUBNA I 593 +BENDULIN, GOLDBERG, + (EPOL+SACL+AMST) I  
 PJERROU 65 PRL 14 275 +SCHLEIN, SLATER, SMITH, STORK, TICHO (UCLA)  
 BERGE 66 PR 147 945 +EBERHARD, HUBBARD, MERRILL, B-SHAFER, + (LRL) I  
 LONDON 66 PR 143 1034 +RAU, SAMIOS, YAMAMOTO, GOLDBERG, + (BNL+SYRA) IJ  
 MERRILL 66 UCRL-16455 THESIS D W MERRILL (LRL) JP  
 BADIER 72 NP 537 429 +BARRETT, CHARLTON, VIDEAU (EPOL)  
 BALTAY 72 PL 42B 129 +BRIDGEWATER, COOPER, GERSHWIN, + (COLU+BING)  
 HABI (COLU)  
 ALSO 73 NEVIS 199 THESIS  
 BORENSTEIN 72 PR D5 1559 BORENSTEIN, DANBURG, KALBFLEISCH++ (BNL+MICH) I  
 KIRSCH 72 NP 840 349 SCHWIDT, CHANG, HEINGHAY+(BRAN+UMD+SYRA+TUFT) I  
 ROSS2 73 PURDUE CONF. 355 ROSS, LLOYD, RADJOJIC (OXF)  
 BERTHON 74 NC 21A 146 +BERNTHON, TRISTRAM, + (CDEF+RHEL+SACL+STRB)  
 LICHTENB 74 PR D10 3865 D B LICHTENBERG (IND)  
 ALSO 74 PRIV. COMM. D B LICHTENBERG (IND)  
 BELLEFOZ 75 NC 28A 289 DE BELLEFON, BERTHON, BILLOIR+ (CDEF+SACL)  
 KALBFLEI 75 PR D11 987 KALBFLEISCH, STRAND, CHAPMAN (BNL+MICH)  
 SIXEL 79 NP 8159 125 +BOTTCHEER, KLEIN+ (AACH+BERL+CERN+LOIC+VIEN)  
 BAUBILLI 81 PR 192 1 +BAUBILLIER, + (BRM+CERN+GLAS+MSU+LMP)  
 BIAGI 81 ZPHY C9 305 + (BRIS+CAMB+GEVA+HEID+LAUS+LOQM+RHEL)  
 ASTON 85 PR D32 2270 + LASS COLLABORATION (SLAC+CARL+CNCR+IINC)

PAPERS NOT REFERRED TO IN DATA LISTINGS  
 SHAFER 66 PR 142 883 BUTTON-SHAFER, LINDSEY, MURRAY, SMITH (LRL) JP  
 HUNGERBUHLER, MAJKA, + (YALE+PAL-BNL+PITT)  
 BRIEFEL 75 PR D12 1859 +GOUREVITCH, KIRSCH+ (BRAN+UMD+SYRA+TUFT)  
 BRIEFEL 77 PR D16 2706 +GOUREVITCH, CHANG+ (BRAN+UMD+SYRA+TUFT)  
 MAZZUCATO 81 NP 8178 1 MAZZUCATO, PENNING+ (AMST+CERN-NIJM-OXF)



# Baryon Full Listings

$\Xi(1630)$ ,  $\Xi(1680)$ ,  $\Xi(1820)$

## $\Xi(1630)$

$$I(J^P) = \frac{1}{2}(\quad) \quad \text{Status: } *$$

OMITTED FROM SUMMARY TABLE

SEEN ONLY IN THE XI P1 CHANNEL.

BARTSCH 69 SEE A SMALL, BROAD ENHANCEMENT NEAR 1650 MEV - IT IS NOT CLEAR THAT IT IS THE SAME PHENOMENON AS BRIEFEL 77, WHO FIND CS=2.6--0.9 MICROBARN AT 2.87 GEV/C INCIDENT K- MOMENTUM.

BORNSTEIN 72 SEE NO EFFECT IN THIS REGION. THEY FIND CS=2 MICROBARN AT 2.18 GEV/C.

ROSS 72 ARGUE THAT THE EFFECT THEY SEE IS NOT THE SAME AS THAT SEEN BY BRIEFEL 77 (WHOSE PRELIMINARY RESULTS WERE REPORTED IN BMST 70), AND FIND CS=2+-1 MICROBARN AT 3.5 GEV/C.

BELLEFON 75 FIND A CS OF AROUND 10 MICROBARN NEAR 2 GEV/C, BUT LESS THAN 3 MICROBARN AROUND 2.3 GEV/C.

NOT SEEN BY HASSALL 81 IN A HIGH STATISTICS BUBBLE CHAMBER EXPERIMENT (46 EVENTS/MICROBARN) AT 6.5 GEV/C.

### $\Xi(1630)$ MASS (MeV)

M	29	1606.0	6.0	ROSS	72	HBC	0	K-P	AT	3.1-3.7
M	34	1633.0	12.0	BELLEF02	75	HBC	0	K-P	TO	XI- K PI
M	31	1624.0	3.0	BRIEFEL	77	HBC	0	K-P	2.87	GEV/C

### $\Xi(1630)$ WIDTH (MeV)

W	29	21.0	7.0	ROSS	72	HBC	0	XI-PI+	K*(0.890)	
W	34	40.0	15.0	BELLEF02	75	HBC	0	K-P	TO XI- K PI	
W A	31	(22.5)		BRIEFEL	77	HBC	0	K-P	2.87 GEV/C	
W A	GOODNESS OF FIT INSENSITIVE TO VALUES BETWEEN 15 AND 30 MEV.									

### $\Xi(1630)$ PARTIAL DECAY MODES

P1	$\Xi(1630) \rightarrow \Xi \pi$	DECAY MASSES
		1321- 140
	SEEN IN K- P TO XI- PI+ K0 AND XI- P10 K+.	

### REFERENCES FOR $\Xi(1630)$

ROSS 72 PL 388 177 +BURAN, LLOYD, MULVEY, RADOJICIC (OXF) I  
 BELLEF02 75 NC 28A 289 DE BELLEFON, BERTHON, BILLOIR. (CDFE+SACL)  
 BRIEFEL 77 PR D16 2706 +GOURVITCH, CHANG+ (BRAN+UMD-SYRA+TUFT)  
 ALSO 70 DUKE CONF. 317 BMST (BRAN+UMD-SYRA+TUFT)

PAPERS NOT REFERRED TO IN DATA LISTINGS

APSELL 69 PRL 23 884 + (BRAN+UMD-SYRA+TUFT)  
 SUPERSEDED BY BMST 70.  
 BARTSCH 69 PL 288 439 + (AACH+BERL+CERN-LOIC+VIEV)  
 KALBLEI 70 DUKE CONF 331 + (KBL) I  
 BORNSTEIN 72 PR D5 1559 G R KALBFLEISCH  
 SCHMIDT 73 PURDUE CONF. 363 BORNSTEIN, DANBURG, KALBFLEISCH++ (BNL+MICH) I  
 HUNGERBU 74 PR D10 2051 SCHMIDT (BRAN)  
 BRIEFEL 75 PR D12 1859 HUNGERBUHLER, MAJKA, + (YALE+FNAL+BNL+PITT)  
 HASSALL 81 NP B189 397 +GOURVITCH, KIRSCH+ (BRAN+UMD-SYRA+TUFT)  
 +ANSORGE, CARTER, NEALE, RUSHBROOKE+(CAMB+MSU)

## $\Xi(1680)$

$$I(J^P) = \frac{1}{2}(\quad) \quad \text{Status: } **$$

OMITTED FROM SUMMARY TABLE

SEEN BY DIONISI 78 AS A THRESHOLD ENHANCEMENT IN BOTH THE NEUTRAL AND NEGATIVELY CHARGED SIGMA K<sub>BAR</sub> MASS SPECTRA FROM THE REACTIONS K-P --> (SIGMA K<sub>BAR</sub>) K PI AT 4.2 GEV/C. THE DATA FROM THE SIGMA K<sub>BAR</sub> CHANNELS ALONE CANNOT DISTINGUISH BETWEEN A RESONANCE INTERPRETATION AND A LARGE SCATTERING LENGTH.

WEAKER EVIDENCE FOR AN ENHANCEMENT AT THE SAME MASS IS SEEN IN THE CORRESPONDING LAMBDA K<sub>BAR</sub> CHANNELS AND A COUPLED CHANNEL ANALYSIS YIELDS RESULTS CONSISTENT WITH A NEW XI.

THE HYPERON BEAM EXPERIMENT OF BIAGI 81 OBSERVE AN ENHANCEMENT AT 1700 MEV IN THE DIFFRACTIVELY PRODUCED LAMBDA K- SYSTEM. A PEAK IS ALSO OBSERVED IN THE LAMBDA K0 MASS SPECTRUM AT 1660 MEV WHICH IS CONSISTENT WITH A RESONANCE OF MASS 1720 MEV DECAYING INTO SIGMA<sup>0</sup> K-, WITH THE GAMMA FROM THE SIGMA<sup>0</sup> DECAY NOT DETECTED.

### $\Xi(1680)$ MASS (MeV)

M0	NEUTRAL CHARGE									
M0	A	175(1699.0)	(5.0)	DIONISI	78	HBC	0	K-P	AT	4.2 GEV/C
M0	A	FROM FIT TO SIGMA+ K- SPECTRUM								
M0	B	183(1684.0)	(5.0)	DIONISI	78	HBC	0	K-P	AT	4.2 GEV/C
M0	B	FROM COUPLED CHANNEL ANALYSIS OF SIGMA- K- AND LAMBDA K0 SPECTRA								
M-	NEGATIVE CHARGE									
M-	C	45(1694.0)	(6.0)	DIONISI	78	HBC	-	K-P	AT	4.2 GEV/C
M-	C	FROM COUPLED CHANNEL ANALYSIS OF SIGMA K- AND LAMBDA K- SPECTRA								
M-	D	150(1700.0)	(10.0)	BIAGI	81	SPEC	-	HYPERON BEAM		
M-	D	FIT TO INCLUSIVE SPECTRUM FROM XI-N --> LAM K- X								

### $\Xi(1680)$ WIDTH (MeV)

W0	NEUTRAL CHARGE									
W0	A	175 (44.0)	(23.0)	DIONISI	78	HBC	0	K-P	AT	4.2 GEV/C
W0	B	185 (20.0)	(4.0)	DIONISI	78	HBC	0	K-P	AT	4.2 GEV/C
W-	NEGATIVE CHARGE									
W-	C	45 (26.0)	(6.0)	DIONISI	78	HBC	-	K-P	AT	4.2 GEV/C
W-	D	150 (47.0)	(14.0)	BIAGI	81	SPEC	-	HYPERON BEAM		

### $\Xi(1680)$ PARTIAL DECAY MODES

P1	$\Xi(1680) \rightarrow \Sigma \bar{K}$	DECAY MASSES
		1192- 498
P2	$\Xi(1680) \rightarrow \Lambda \bar{K}$	
		1116- 498
P3	$\Xi(1680) \rightarrow \Xi \pi$	
		1315- 135
P4	$\Xi(1680) \rightarrow \Xi(1530) \pi$	
		1533- 135
P5	$\Xi(1680) \rightarrow \Xi \pi \pi$ (including $\Xi(1530) \pi$ )	
		1315- 135, 135

### $\Xi(1680)$ BRANCHING RATIOS

$\Xi(1680) \rightarrow (\Sigma \bar{K})/(\Lambda \bar{K})$	(P1)/(P2)
R1 E (2.7) (0.9)	DIONISI 78 HBC 0 K-P AT 4.2 GEV/C
R1 E NEUTRAL CHARGE	
R1 F (3.1) (1.4)	DIONISI 78 HBC - K-P AT 4.2 GEV/C
R1 F NEGATIVE CHARGE	
$\Xi(1680) \rightarrow (\Xi \pi)/(\Sigma \bar{K})$	(P3)/(P1)
R2 (0.09) OR LESS	DIONISI 78 HBC 0 K-P AT 4.2 GEV/C
$\Xi(1680) \rightarrow (\Xi^- \pi^+ \pi^0)/(\Sigma \bar{K})$	(P5)/(P1)
R3 (0.04) OR LESS	DIONISI 78 HBC 0 K-P AT 4.2 GEV/C
$\Xi(1680) \rightarrow (\Xi^- \pi^+ \pi^-)/(\Sigma \bar{K})$	(P5)/(P1)
R4 (0.03) OR LESS	DIONISI 78 HBC - K-P AT 4.2 GEV/C
$\Xi(1680) \rightarrow (\Xi(1530) \pi)/(\Sigma \bar{K})$	(P4)/(P1)
R5 (0.06) OR LESS	DIONISI 78 HBC - K-P AT 4.2 GEV/C

### REFERENCES FOR $\Xi(1680)$

DIONISI 78 PL 80B 145 +DIAZ, ARMENTEROS- (CERN-AMST-NIJM-0XF) I, JP  
 BIAGI 81 ZPHY C9 305 + (BRIS+CAMB-GEVA+HETD-LAUS+LODM+RHEL)

## $\Xi(1820)$

$$I(J^P) = \frac{1}{2}(\frac{3}{2}) \quad \text{Status: } ***$$

WE LIST HERE EVERYTHING REPORTED IN THE MASS RANGE 1750-1875 MEV.

The clearest evidence for the  $\Xi(1820)$  comes from GAY 76, who saw an 8-standard-deviation peak in  $\Lambda K^-$  and smaller signals in  $\Xi(1530)\pi$  and  $\Sigma \bar{K}$ . The peak is narrow ( $\Gamma = 21 \pm 7$  MeV), whereas earlier (and much smaller) experiments found widths of up to 100 MeV (see the Listings below). A spin-parity analysis of the GAY 76 data, but with more events (TEODORO 78), favors spin 3/2 but cannot make a parity discrimination.

BIAGI 81 used the CERN hyperon beam to study  $\Xi^-$  interactions in hydrogen and deuterium. The diffractively produced  $\Lambda K^-$  system has a broad peak ( $\Gamma = 72 \pm 20$  MeV) at 1830 MeV on top of a substantial background. There is also a smaller peak in the inclusive  $\Lambda K_S^0$  spectrum.

Neither GAY 76 nor BIAGI 81 saw a peak in the  $\Xi \pi$  channel. It is possible that  $\Xi \pi$  peaks seen in this region by some of the lower momentum experiments were at least partly due to the  $\Xi(1940)$ , with a shape distorted by the limited phase space available (SMITH 65). Furthermore, some of the early experiments were forced to add several different channels together to overcome poor statistics (CRENNELL 70, BADIER 71).

For notation, see key on page 91.

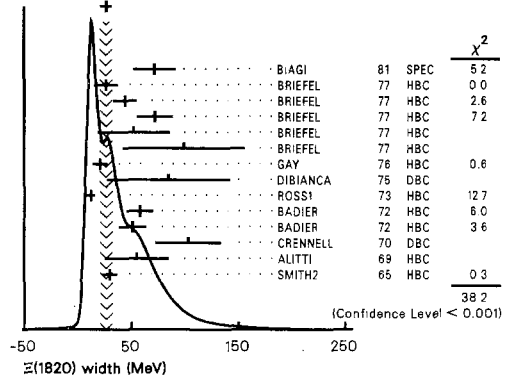
# Baryon Full Listings

$\Xi(1820)$

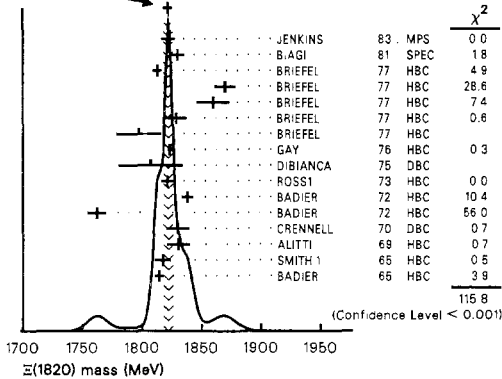
## $\Xi(1820)$ MASS (MeV)

M	(1770.0)		HALSTEINS	63	FBC	-0	K-FR	3.5	GEV/C		
M	30	1814.0	4.0	BADIER	65	HBC	0	LAMBDA	KOBAR		
M	29	1817.0	7.0	SMITH	1	65	HBC	-0	LAMBDA	KBAR	
M	40	1830.0	10.0	ALITTI	69	HBC	-	LAM,	SIG	KBAR	
M	A	25	1830.0	10.0	CRENNELL	70	DBC	-0	3.6,	3.9	GEV/C
M	B	FROM FIT TO INCLUSIVE XI PI, XI PI PI AND LAMBDA K- SPECTRA		CRENNELL	70	DBC	-0	3.6,	3.9	GEV/C	
M	(1826.0)	(12.0)									
M	B	FROM FIT TO INCLUSIVE XI PI AND XI PI PI SPECTRA ONLY									
M	C	28	1762.0	8.0	BADIER	72	HBC	-0	XI PI, XI2PI, K Y		
M	C	38	1836.0	5.0	BADIER	72	HBC	-0	XI PI, XI2PI, K Y		
M	C	BADIER 72 ADDS ALL CHANNELS AND DIVIDES PEAK IN LOWER AND HIGHER MASS REGIONS. THE DATA CAN ALSO BE FITTED WITH A SINGLE BREIT-WIGNER OF MASS 1800 AND WIDTH 150 MEV.									
M	D	30	1821.0	5.0	ROSS1	73	HBC	-0	LAMBDA	K-/KBARO	
M	D	LESS SIGNIFICANT ENHANCEMENTS SEEN IN XI(1530) PI (M=1825, W=100) AND SIGMA KBAR (M=1810+-9, W=16+-11).									
M	D	1807.0	27.0	DIBIANCA	75	DBC	-0	XI 2PI, XI* PI			
M	130	1825.0	2.0	GAY	76	HBC	-	K- P	AT 4.2	GEV	
M	64	1797.0	19.0	BRIEFEL	77	HBC	0	XI PI (2.87	K-P)		
M	78	1829.0	9.0	BRIEFEL	77	HBC	-0	XI(1530) PI			
M	39	1860.0	14.0	BRIEFEL	77	HBC	-	SIGMA- KOBAR			
M	44	1870.0	9.0	BRIEFEL	77	HBC	0	LAMBDA	KOBAR		
M	57	1813.0	4.0	BRIEFEL	77	HBC	-	LAMBDA	K-		
M	E	300	1830.0	6.0	BIAGI	81	SPEC	-	HYPERON	BEAM	
M	E	FIT TO INCLUSIVE SPECTRUM FROM XI-N --> LAM K- X		JENKINS	83	MPS	-	K- P	TO K+ MM		
M	1826.0	6.0									
M	AVG	1821.9	3.9	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 3.0) (SEE IDEOGRAM BELOW)							

## WEIGHTED AVERAGE 26.3 ± 5.8 (ERROR SCALED BY 2.2)



## WEIGHTED AVERAGE 1821.9 ± 3.9 (ERROR SCALED BY 3.0)



## FITTED PARTIAL DECAY MODE BRANCHING FRACTIONS

The matrix below is derived from the error matrix for the fitted partial decay mode branching fractions,  $P_i$ , as follows: The diagonal elements are  $P_i = \delta P_i$ , where  $\delta P_i = \sqrt{(\delta P_i^2)}$ , while the off-diagonal elements are the normalized correlation coefficients  $(\delta P_i \delta P_j) / (\delta P_i \delta P_j)$ . For the definitions of the individual  $P_i$  see the listings above; only those  $P_i$  appearing in the matrix are assumed in the fit to be nonzero and are thus constrained to add to 1.

	P 1	P 2	P 3	P 4
P 1	.4974+-0.0871			
P 2	-.761	.1889+-0.0529		
P 3	.133	-.508	.1463+-0.0477	
P 4	-.822	.581	-.504	.1674+-0.0647

## $\Xi(1820)$ BRANCHING RATIOS

$\Xi(1820) \rightarrow (\Delta \bar{K})/\text{total}$	(P1)
R1	0.30 0.15 ALITTI 69 HBC - K-P 3.9-5.0 GEV
R1 FIT	0.497 0.087 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.8)
$\Xi(1820) \rightarrow (\Xi \pi)/\text{total}$	(P2)
R2	0.10 0.10 ALITTI 69 HBC - K-P 3.9-5.0 GEV
R2 FIT	0.189 0.053 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.2)
$\Xi(1820) \rightarrow (\Xi \pi)/(\Delta \bar{K})$	(P2)/(P1)
R21	0.20 0.20 BADIER 65 HBC 0 K-P AT 3 GEV
R21	(0.56) OR LESS CL=.95 GAY 76 HBC - K- P AT 4.2 GEV
R21 FIT	0.38 0.16 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.5)
$\Xi(1820) \rightarrow (\Xi \pi)/(\Xi(1530) \pi)$	(P2)/(P4)
R22	1.5 0.6 0.4 APSELL 70 HBC 0 K-P AT 2.87 GEV
R22 FIT	1.13 0.36 FROM FIT
$\Xi(1820) \rightarrow (\Sigma \bar{K})/\text{total}$	(P3)
R3	(0.02) OR LESS TRIPP 67 RVUE - K-P 3.9-5.0 GEV
R3	0.30 0.15 ALITTI 69 HBC - K-P 3.9-5.0 GEV
R3 FIT	0.146 0.048 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
$\Xi(1820) \rightarrow (\Sigma \bar{K})/(\Delta \bar{K})$	(P3)/(P1)
R31	0.24 0.10 GAY 76 HBC - K- P AT 4.2 GEV
R31 FIT	0.29 0.10 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.1)
$\Xi(1820) \rightarrow (\Xi(1530) \pi)/\text{total}$	(P4)
R4	0.30 0.15 ALITTI 69 HBC - K-P 3.9-5.0 GEV
R4 F	(0.25) OR LESS DAUBER 69 HBC K-P 2.7 GEV/C
R4 G	USES IN PART THE SAME DATA AS SMITH 65
R4	NOT SEEN HASSALL 81 HBC K-P 6.5 GEV/C
R4	SEEN ASTON 85 LASS K-P 11 GEV/C
R4 G	INCLUDING XI PI PI
R4 FIT	0.167 0.065 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.6)
$\Xi(1820) \rightarrow (\Xi(1530) \pi)/(\Delta \bar{K})$	(P4)/(P1)
R41	0.26 0.13 SMITH1 65 HBC -0 K-P 2.45-2.70GEV
R41	1.0 0.3 GAY 76 HBC - K- P AT 4.2 GEV
R41	
R41 AVG	0.38 0.27 AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.3)
R41 FIT	0.34 0.18 FROM FIT (ERROR INCLUDES SCALE FACTOR OF 1.7)

## $\Xi(1820)$ WIDTH (MeV)

W	(80.0)	OR LESS	HALSTEINS	63	FBC	-0	K-FR	3.5	GEV/C		
W	(12.0)	(4.0)	BADIER	65	HBC	0	LAMBDA	KOBAR			
W	30.0	7.0	SMITH2	65	HBC	-0	LAMBDA	KBAR			
W	55.0	40.0	20.0	ALITTI	69	HBC	-	LAM,	SIG	KBAR	
W	A	103.0	38.0	24.0	CRENNELL	70	DBC	-0	3.6,	3.9	GEV/C
W	B	(48.0)	(36.0)	(19.0)	CRENNELL	70	DBC	-0	3.6,	3.9	GEV/C
W	C	51.0	13.0	BADIER	72	HBC	-0	LOWER	MASS		
W	D	30	12.0	4.0	BADIER	72	HBC	-0	HIGHER	MASS	
W	85.0	58.0	ROSS1	73	HBC	-0	LAMBDA	K-/KBARO			
W	130	21.0	7.0	DIBIANCA	75	DBC	-0	XI 2PI, XI* PI			
W	74	99.0	57.0	GAY	76	HBC	-	K- P	AT 4.2	GEV	
W	68	52.0	34.0	BRIEFEL	77	HBC	0	XI PI (2.87	K-P)		
W	39	72.0	17.0	BRIEFEL	77	HBC	-	SIGMA- KOBAR			
W	44	44.0	11.0	BRIEFEL	77	HBC	0	LAMBDA	KOBAR		
W	57	26.0	11.0	BRIEFEL	77	HBC	-	LAMBDA	K-		
W	E	300	72.0	20.0	BIAGI	81	SPEC	-	HYPERON	BEAM	
W	AVG	26.3	5.8	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 2.2) (SEE IDEOGRAM BELOW)							

## $\Xi(1820)$ PARTIAL DECAY MODES

P1	$\Xi(1820) \rightarrow \Delta \bar{K}$	DECAY MASSES	1116+ 498
P2	$\Xi(1820) \rightarrow \Xi \pi$		1321+ 140
P3	$\Xi(1820) \rightarrow \Sigma \bar{K}$		1197+ 498
P4	$\Xi(1820) \rightarrow \Xi(1530) \pi$		1533+ 140
P5	$\Xi(1820) \rightarrow \Xi \pi \pi$ (excluding $\Xi(1530) \pi$ )		1321- 140+ 140

# Baryon Full Listings

## $\Xi(1820), \Xi(1940)$

$\Xi(1820) \rightarrow (\Xi \pi \pi) / (\Lambda \bar{K})$  (P5)/(P1)  
 R51 (0.1) OR MORE SMITH1 65 HBC -0 K-P 2.45-2.70GEV

$\Xi(1820) \rightarrow (\Xi \pi \pi) / (\Xi(1530) \pi)$  (P5)/(P4)  
 R52 H (0.3) APSELL 70 HBC 0 K-P AT 2.87 GEV  
 R52 H OR LESS UPPER LIMIT FOR THE 5-BODY DECAY  
 R52 CONSISTENT WITH ZERO GAY 76 HBC - K-P AT 4.2 GEV

$\Xi(1820) \rightarrow (\Xi \pi \pi \text{ including } \Xi(1530) \pi) / (\Lambda \bar{K})$   
 R53 (0.14) OR LESS BADIER 65 HBC 0 1 STD.DEV.LIMIT  
 R53 I FOR THE DECAY MODE (XI- PI+ PI0) ONLY

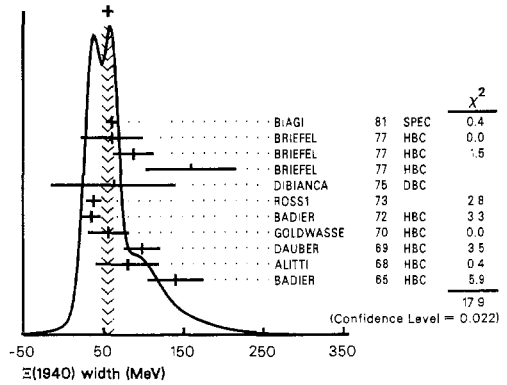
### $\Xi(1940)$ WIDTH (MeV)

W	35	140.0	35.0	BADIER	65 HBC	0 XI- PI+
W	27	80.0	40.0	ALITTI	68 HBC	0 XI- PI+
W	66	90.0	23.0	DAUBER	69 HBC	- XI PI
W	21	56.0	26.0	GOLDWASSE	70 HBC	- XI PI
W	29	35.0	11.0	BADIER	72 HBC	XI PI, XI2PI, K Y
W		38.0	10.0	ROSS1	73	(XI PI)-
W		65.0	78.0	DIBIANCA	75 DBC	XI PI
W	139	159.0	57.0	BRIEFEL	77 HBC	0 XI-PI+(2.87 K-P)
W	44	87.0	26.0	BRIEFEL	77 HBC	- XI0P1-(2.87 K-P)
W	56	60.0	39.0	BRIEFEL	77 HBC	-0 XI(1530) PI
W	A 150	60.0	8.0	BIAGI	81 SPEC	0 HYPERON BEAM

(SEE IDEOGRAM BELOW)

### REFERENCES FOR $\Xi(1820)$

HALSTEIN 63 SIENA CONF 173	HALSTEINSLID, + (BERG+CERN+EPOL+RHEL+LOUC) I
BADIER 65 PL 16 171	+DEMOUNIN, GOLDBERG, + (EPOL+SACL+AMST) I
SMITH1 65 PRL 14 25	+LINDSEY, BUTTON-SHAFER, MURRAY (LRL) I, JP
SMITH2 65 ATHENS CONF 251	G A SMITH, J S LINDSEY (LRL)
TRIPP 67 NP 83 10	+ LEITH, + (LRL+SLAC+CERN+HEID+SACL)
USES DATA OF SMITH1.	
ALITTI 69 PRL 22 79	+BARNES, FLAMINIO, METZGER, + (BNL+SYRA) I
DAUBER 69 PR 179 1262	+BERGE, HUBBARD, MERRILL, MULLER (LRL)
APSELL 70 PRL 24 777	+ (BRAN+UMD+SYRA+TUFT) I
CRENNELL 70 PR D1 847	+KARSHON, LAI, ONEALL, SCARR, SCHUMANN(BNL)
BADIER 72 NP 837 429	+BARRELET, CHARLTON, VIDEAU (EPOL)
ROSS1 73 PURDUE CONF. 345	ROSS, LLOYD, RADOJICIC (OXF)
DIBIANCA 75 NP 898 137	DIBIANCA, ENDORF (CARN)
GAY 76 PL 628 477	+ARMENTEROS, BERGE, GAVILLET-(AMST+CERN+NIJM) I, J
BRIEFEL 77 PR D16 2706	+GOUREVITCH, CHANG + (BRAN+UMD+SYRA+TUFT)
ALSO 70 DUKE CONF. 317	BMST + (BRAN+UMD+SYRA+TUFT)
BIAGI 81 ZPHY C9 305	+ (BRIS+CAMB+GEVA+HEID+LAUS+LOOM+RHEL)
HASSALL 81 NP B189 397	+ANSORGE, CARTER, NEALE, RUSHBROOKE+(CAMB+MSU)
JEWKINS 83 PRL 51 951	+ALBRIGHT, DIAMOND, +(FSU+BRAN+LBU+CINC+SMAS)
ASTON 85 PR D32 2270	+ LASS COLLABORATION (SLAC+CARL+CNRC+CINC)
PAPERS NOT REFERRED TO IN DATA LISTINGS	
SMITH 64 PRL 13 61	+LINDSEY, MURRAY, BUTTON-SHAFER+ (LRL) I, JP
MERRILL 68 PR 167 1202	D W MERRILL, J BUTTON-SHAFER (LRL)
APSELL 69 PRL 23 884	+ (BRAN+UMD+SYRA+TUFT)
SUPERSEDED BY BRIEFEL 77.	
SCHMIDT 73 PURDUE CONF. 363	SCHMIDT (BRAN)
BRIEFEL 75 PR D12 1859	+GOUREVITCH, KIRSCH+ (BRAN-UMD-SYRA+TUFT)
TEODORO 78 PL 77B 451	+DIAZ, DIONISI, BLOKZIJL+(AMST+CERN+NIJM+OXF) JP



## $\Xi(1940)$

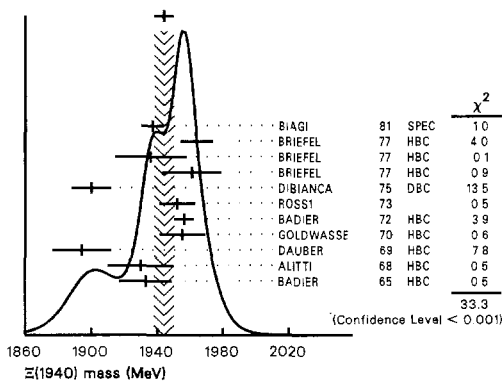
$I(J^P) = \frac{1}{2} ( )$  Status: \*\*

OMITTED FROM SUMMARY TABLE

WE LIST UNDER XI(1940) EVERYTHING REPORTED IN THE MASS RANGE 1875-2000 MEV.

### $\Xi(1940)$ MASS (MeV)

M	35	1933.0	16.0	BADIER	65 HBC	0 XI- PI+
M	27	1930.0	20.0	ALITTI	68 HBC	0 XI- PI-
M	66	1894.0	18.0	DAUBER	69 HBC	- XI PI
M	21	1885.0	14.0	GOLDWASSE	70 HBC	- XI PI
M	29	1956.0	6.0	BADIER	72 HBC	XI PI, XI2PI, K Y
M	25	1952.0	11.0	ROSS1	73	(XI PI)-
M		1900.0	12.0	DIBIANCA	75 DBC	XI PI
M	139	1961.0	18.0	BRIEFEL	77 HBC	0 XI-PI+(2.87 K-P)
M	44	1936.0	22.0	BRIEFEL	77 HBC	- XI0P1-(2.87 K-P)
M	56	1964.0	10.0	BRIEFEL	77 HBC	-0 XI(1530) PI
M	A 150	1937.0	7.0	BIAGI	81 SPEC	0 HYPERON BEAM
M	A FIT TO INCLUSIVE SPECTRUM FROM XI- --> XI- PI- X (SEE IDEOGRAM BELOW)					



### $\Xi(1940)$ PARTIAL DECAY MODES

P1	$\Xi(1940) \rightarrow \Xi \pi$	1321+ 140
P2	$\Xi(1940) \rightarrow \Xi(1530) \pi$	1533+ 140
P3	$\Xi(1940) \rightarrow \Xi \pi \pi$ (excluding $\Xi(1530) \pi$ )	1321+ 140+ 140
P4	$\Xi(1940) \rightarrow \Xi^0 \pi^-$	1315+ 140
P5	$\Xi(1940) \rightarrow \Xi^- \pi^0$	1321+ 135
P6	$\Xi(1940) \rightarrow \Sigma \bar{K}$	1197+ 498

### $\Xi(1940)$ BRANCHING RATIOS

THE XI(1940) IS SEEN MAINLY IN XI PI AND SOME IN XI(1530) PI. IT HAS BEEN LOOKED FOR IN OTHER CHANNELS BUT ONLY OBSERVED BY HASSALL 81 WHO SEE A SIGMA EFFECT IN SIGMA KBAR.

$\Xi(1940) \rightarrow (\Xi \pi) / (\Xi(1530) \pi)$ (P1)/(P2)	R1	2.8	0.7	0.6	APSELL	70 HBC	0
$\Xi(1940) \rightarrow (\Xi \pi \pi) / (\Xi(1530) \pi)$ (P3)/(P2)	R2	0.0	0.3		APSELL	70 HBC	0
$\Xi(1940) \rightarrow (\Xi^0 \pi^-) / (\Xi^- \pi^0)$ (P4)/(P5)	R3	25	2.6	6.0	1.6	ROSS1	73 (XI PI)-
$\Xi(1940) \rightarrow (\Sigma \bar{K}) / \text{total}$ (P6)	R4	17	POSSIBLY SEEN		HASSALL	81 HBC	- K-P AT 6.5 GEV/C

### REFERENCES FOR $\Xi(1940)$

BADIER 65 PL 16 171	+DEMOUNIN, GOLDBERG, + (EPOL+SACL+AMST) I
ALITTI 68 PRL 21 1119	+FLAMINIO, METZGER, RADOJICIC, + (BNL+SYRA) I
DAUBER 69 PR 179 1262	+BERGE, HUBBARD, MERRILL, MULLER (LRL) I
APSELL 70 PRL 24 777	+ (BRAN+UMD+SYRA+TUFT) I
GOLDWASS 70 PR D1 1960	E L GOLDWASSER, P F SCHULTZ (ILLINOIS)
BADIER 72 NP 837 429	+BARRELET, CHARLTON, VIDEAU (EPOL)
ROSS1 73 PURDUE CONF. 345	ROSS, LLOYD, RADOJICIC (OXF)
DIBIANCA 75 NP 898 137	DIBIANCA, ENDORF (CARN)
BRIEFEL 77 PR D16 2706	+GOUREVITCH, CHANG + (BRAN+UMD+SYRA+TUFT)
ALSO 70 DUKE CONF. 317	BMST + (BRAN+UMD+SYRA+TUFT)
BIAGI 81 ZPHY C9 305	+ (BRIS+CAMB+GEVA+HEID+LAUS+LOOM+RHEL)
HASSALL 81 NP B189 397	+ANSORGE, CARTER, NEALE, RUSHBROOKE+(CAMB+MSU)
PAPERS NOT REFERRED TO IN DATA LISTINGS	
APSELL 69 PRL 23 884	+ (BRAN+UMD+SYRA+TUFT)
SUPERSEDED BY BMST 70.	
SCHMIDT 73 PURDUE CONF. 363	SCHMIDT (BRAN)
BRIEFEL 75 PR D12 1859	+GOUREVITCH, KIRSCH+ (BRAN+UMD+SYRA+TUFT)

For notation, see key on page 91.

# Baryon Full Listings

$\Xi(2030)$ ,  $\Xi(2120)$

## $\Xi(2030)$

$I(J^P) = \frac{1}{2}(\frac{5}{2})$  Status: \*\*\*

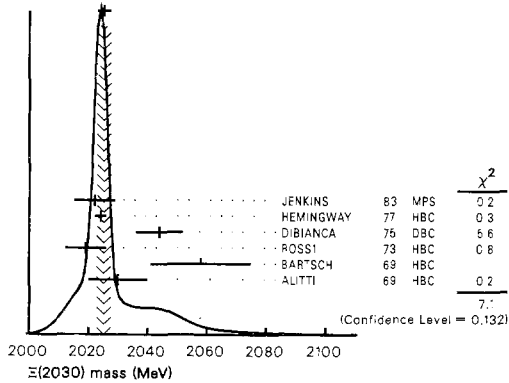
THE EVIDENCE FOR THIS STATE HAS BEEN MUCH IMPROVED BY HEMINGWAY 77, WHO SEE AN 8 STD. DEV. ENHANCEMENT IN SIGMA K $\bar{A}$ R AND A WEAKER COUPLING TO LAMBDA K $\bar{A}$ R. ALITTI 68 AND HEMINGWAY 77 OBSERVE NO SIGNALS IN THE XI PI PI (OR XI(1530) PI) CHANNEL, IN CONTRAST TO DIBIANCA 75, THE DECAY INTO LAMBDA/SIGMA K $\bar{A}$ R PI REPORTED BY BARTSCH 69 IS ALSO NOT CONFIRMED BY HEMINGWAY 77.

A MOMENTS ANALYSIS OF THE HEMINGWAY 77 DATA INDICATES THAT THE SPIN IS GREATER THAN OR EQUAL TO 5/2 AT A LEVEL OF 3 STD. DEVIATIONS.

### $\Xi(2030)$ MASS (MeV)

M	42	2030.0	10.0	ALITTI	69	HBC	-	K-P	3.9-5	GEV/C
M	40	2058.0	17.0	BARTSCH	69	HBC	-	K-P	10	GEV/C
M	15	2019.0	7.0	ROSSI	73	HBC	-	SIGMA	K $\bar{A}$ R	
M	2044.0	8.0	DIBIANCA	75	DBC	-	0	XI	2PI, XI* PI	
M	200	2024.0	2.0	HEMINGWAY	77	HBC	-	K-P	AT 4.2	GEV
M	2022.0	7.0	JENKINS	83	MPS	-	K-P	TO K+ MM		
M	AVG	2025.1	2.4	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.3) (SEE IDEOGRAM BELOW)						

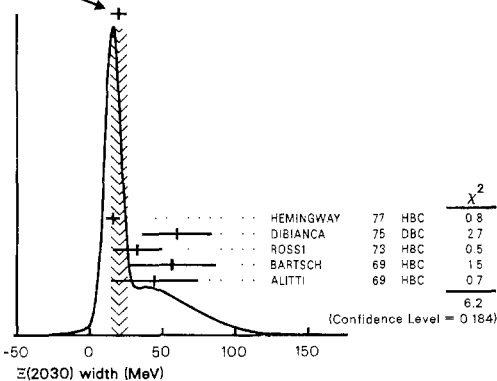
WEIGHTED AVERAGE  
2025.1 ± 2.4 (ERROR SCALED BY 1.3)



### $\Xi(2030)$ WIDTH (MeV)

W	45.0	40.0	20.0	ALITTI	69	HBC	-	K-P	3.9-5	GEV/C
W	57.0	30.0		BARTSCH	69	HBC	-	K-P	10	GEV/C
W	15	33.0	17.0	ROSSI	73	HBC	-	SIGMA	K $\bar{A}$ R	
W	60.0	24.0		DIBIANCA	75	DBC	-	0	XI	2PI, XI* PI
W	200	16.0	5.0	HEMINGWAY	77	HBC	-	K-P	AT 4.2	GEV
W	AVG	20.5	5.7	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.2) (SEE IDEOGRAM BELOW)						

WEIGHTED AVERAGE  
20.5 ± 5.7 (ERROR SCALED BY 1.2)



## $\Xi(2030)$ PARTIAL DECAY MODES

P1	$\Xi(2030) \rightarrow \Xi \pi$	1321+ 140	DECAY MASSES
P2	$\Xi(2030) \rightarrow \Lambda \bar{K}$	1116+ 498	
P3	$\Xi(2030) \rightarrow \Sigma \bar{K}$	1197+ 498	
P4	$\Xi(2030) \rightarrow \Xi(1530) \pi$	1533+ 140	
P5	$\Xi(2030) \rightarrow \Xi \pi \pi$ (excluding $\Xi(1530) \pi$ )	1321+ 140+ 140	
P6	$\Xi(2030) \rightarrow \Lambda \bar{K} \pi$	1116+ 498+ 140	
P7	$\Xi(2030) \rightarrow \Sigma \bar{K} \pi$	1189+ 498+ 140	

## $\Xi(2030)$ BRANCHING RATIOS

$\Xi(2030) \rightarrow (\Xi \pi)/(\text{modes P1 to P4})$	(P1)/(P1+P2+P3+P4)
R1 (0.30) OR LESS	ALITTI 69 HBC - 1 STD DEV LIMIT
$\Xi(2030) \rightarrow (\Xi \pi)/(\Sigma \bar{K})$	(P1)/(P3)
R11 (0.19) OR LESS CL=.95	HEMINGWAY 77 HBC - K-P AT 4.2 GEV
$\Xi(2030) \rightarrow (\Lambda \bar{K})/(\text{modes P1 to P4})$	(P2)/(P1+P2+P3+P4)
R2 0.25 0.15	ALITTI 69 HBC - K-P 3.9-5 GEV/C
$\Xi(2030) \rightarrow (\Lambda \bar{K})/(\Sigma \bar{K})$	(P2)/(P3)
R21 0.22 0.09	HEMINGWAY 77 HBC - K-P AT 4.2 GEV
$\Xi(2030) \rightarrow (\Sigma \bar{K})/(\text{modes P1 to P4})$	(P3)/(P1+P2+P3+P4)
R3 0.75 0.20	ALITTI 69 HBC - K-P 3.9-5 GEV/C
$\Xi(2030) \rightarrow (\Xi(1530) \pi)/(\text{modes P1 to P4})$	(P4)/(P1+P2+P3+P4)
R4 (0.15) OR LESS	ALITTI 69 HBC - 1 STD DEV LIMIT
$\Xi(2030) \rightarrow (\Xi \pi \pi \text{ including } \Xi(1530) \pi)/(\Sigma \bar{K})$	(P4+P5)/(P3)
R41 A (0.11) OR LESS CL=.95	HEMINGWAY 77 HBC - K-P AT 4.2 GEV
R41 A FOR THE DECAY MODE (XI- PI+ PI-) ONLY	
$\Xi(2030) \rightarrow (\Lambda \bar{K} \pi)/\text{total}$	(P6)
R6 SEEN	BARTSCH 69 HBC - K-P AT 10 GEV
$\Xi(2030) \rightarrow (\Lambda \bar{K} \pi)/(\Sigma \bar{K})$	(P6)/(P3)
R61 (0.32) OR LESS CL=.95	HEMINGWAY 77 HBC - K-P AT 4.2 GEV
$\Xi(2030) \rightarrow (\Sigma \bar{K} \pi)/\text{total}$	(P7)
R7 SEEN	BARTSCH 69 HBC - K-P AT 10 GEV
$\Xi(2030) \rightarrow (\Sigma \bar{K} \pi)/(\Sigma \bar{K})$	(P7)/(P3)
R71 B (0.04) OR LESS CL=.95	HEMINGWAY 77 HBC - K-P AT 4.2 GEV
R71 B FOR THE DECAY MODE (SIGMA-- K- PI-) ONLY	

## REFERENCES FOR $\Xi(2030)$

ALITTI 69 PRL 22 79 +BARNES,FLAMINIO,METZGER, + (BNL-SYRA) I  
 BARTSCH 69 PL 288 439 + (AACH+BERL+CERN-LOIC-VIEN)  
 ROSSI 73 PURDUE CONF. 345 ROSS,LLOYD,RADJOJIC (OXF)

DIBIANCA 75 NP 898 137 DIBIANCA,ENDORF (CERN)  
 HEMINGWAY 77 PL 688 197 HEMINGWAY,ARMENTEROS+ (AMST+CERN+NIJ+OXF) IJ  
 ALSO 76 PL 628 477 GAY,ARMENTEROS,BERGE- (AMST+CERN+NIJ)  
 JENKINS 83 PRL 51 951 +ALBRIGHT,DIAMOND,-(FSU-BRAN+LBL+CINC+SMAS)

## $\Xi(2120)$

$I(J^P) = \frac{1}{2}(\ )$  Status: \*

OMITTED FROM SUMMARY TABLE

THIS EFFECT IS REPORTED IN GAY 76 AS A FOUR STANDARD DEVIATION ENHANCEMENT IN LAMBDA K-. AN ANALYSIS OF THE SAME DATA BY HEMINGWAY 77, BUT WITH ADDITIONAL STATISTICS, POINTS OUT THAT THE SIGNIFICANCE OF THE ENHANCEMENT IS GREATLY REDUCED IF A RESTRICTIVE FOUR-MOMENTUM CUT (U-CUT) IS MADE. THIS SUGGESTS AN ANOMALOUS PRODUCTION MECHANISM IF THE STATE IS GENUINE.

CHLIAPNIKOV 79 REPORT A BUMP OF 18 EVENTS AT 2137 MEV IN AN INCLUSIVE LAMBDA $\bar{A}$ R K+ SPECTRUM FROM K+P INTERACTIONS AT 32 GEV/C. THE K- ARE NOT UNIQUELY IDENTIFIED. BUMPS WITH LOWER NUMBERS OF EVENTS ARE ALSO REPORTED AT 2240, 2830, AND 2540 MEV.

### $\Xi(2120)$ MASS (MeV)

M	2123.0	7.0	GAY	76	HBC	-	K-P	AT 4.2	GEV
M	18(2137.0)	(4.0)	CHLIAPNIKOV	79	HBC	+	LAMBDA	BAR K-	

### $\Xi(2120)$ WIDTH (MeV)

W	25.0	12.0	GAY	76	HBC	-	K-P	AT 4.2	GEV
W	18	(20.0) OR LESS	CHLIAPNIKOV	79	HBC	+	LAMBDA	BAR K+	

# Baryon Full Listings

$\Xi(2120)$ ,  $\Xi(2250)$ ,  $\Xi(2370)$ ,  $\Xi(2500)$

## $\Xi(2120)$ PARTIAL DECAY MODES

P1	$\Xi(2120) \rightarrow \Lambda \bar{K}$	DECAY MASSES
		1116+ 498

## $\Xi(2120)$ BRANCHING RATIOS

$\Xi(2120) \rightarrow (\Lambda \bar{K})/\text{total}$	(P1)
R1 SEEN	GAY 76 HBC - K-P AT 4.2 GEV

## REFERENCES FOR $\Xi(2120)$

GAY 76 PL 628 477 + ARMENTEROS, BERGE, GAVILLET, (AMST+CERN+NIJM) I  
 HEMINGWAY 77 PL 688 197 HEMINGWAY, ARMENTEROS, (AMST+CERN+NIJM+OXF)  
 CHLIAPNI 79 NP B158 253 CHLIAPNIKOV, GERDYUKOV, (SERP+BELG+MONS)

## $\Xi(2250)$

$$I(J^P) = \frac{1}{2} ( ) \quad \text{Status: **}$$

OMITTED FROM SUMMARY TABLE

THE EVIDENCE FOR THIS STATE IS MIXED. BARTSCH 69 SEE A BUMP OF NOT MUCH STATISTICAL SIGNIFICANCE IN LAMBDA-KBAR-PI, SIGMA-KBAR-PI, AND XI-PI-PI MASS SPECTRA. GOLDWASSER 70 SEE A NARROWER BUMP IN XI-PI-PI AT A HIGHER MASS. NOT SEEN BY HASSALL 81 WITH 45 EVENTS/MICROBARN AT 6.5 GEV/C. SEEN BY JENKINS 83.

## $\Xi(2250)$ MASS (MeV)

M	35 2244.0	52.0	BARTSCH 69 HBC	K-P 10 GEV/C
M	18 2295.0	15.0	GOLDWASSE 70 HBC	K-P 5.5 GEV/C
M	2274.0	5.0	JENKINS 83 MPS	K-P TO K+ MM

## $\Xi(2250)$ WIDTH (MeV)

W	130.0	80.0	BARTSCH 69 HBC	K-P 5.5 GEV/C
W	LESS THAN	30.0	GOLDWASSE 70 HBC	K-P 5.5 GEV/C

## $\Xi(2250)$ PARTIAL DECAY MODES

P1	$\Xi(2250) \rightarrow \Xi \pi \pi$	DECAY MASSES
P2	$\Xi(2250) \rightarrow \Lambda \bar{K} \pi$	1321+ 140- 140
P3	$\Xi(2250) \rightarrow \Sigma \bar{K} \pi$	1116+ 498- 140
		1197+ 498- 140

## REFERENCES FOR $\Xi(2250)$

BARTSCH 69 PL 288 439 - (AACH-BERL+CERN-LOIC+VIEN)  
 GOLDWASS 70 PR D1 1960 E L GOLDWASSER, P F SCHULTZ (ILL)  
 JENKINS 83 PRL 51 951 +ALBRIGHT, DIAMOND, +(FSU+BRAN-LBL-CINC-SMAS)

PAPERS NOT REFERRED TO IN DATA LISTINGS

HASSALL 81 NP B189 397 +ANSORGE, CARTER, NEALE, RUSHBROOKE+(CAMB+MSU)

## $\Xi(2370)$

$$I(J^P) = \frac{1}{2} ( ) \quad \text{Status: **}$$

OMITTED FROM SUMMARY TABLE

SEEN BY AMIRZADEH 80 AND HASSALL 81 IN THE CHARGED AND NEUTRAL LAMBDA/SIGMA KBAR PI MASS SPECTRA FROM THE REACTIONS K-P  $\rightarrow$  XI(1530) K AND XI(1530) K PI. AMIRZADEH 80 ALSO OBSERVE A SMALL EFFECT AT THE SAME MASS IN THE OMEGA- K MASS SPECTRUM. KINSON 80 RE-ANALYSE THE DATA OF AMIRZADEH 80 BUT WITH 50 PER CENT MORE STATISTICS.

## $\Xi(2370)$ MASS (MeV)

M	2392.0	27.0	DIBIANCA 75 DBC	XI 2P1
M	94 2373.0	8.0	AMIRZAD 80 HBC	-0 K-P AT 8.25 GEV
M	50(2370.0)		HASSALL 81 HBC	-0 K-P AT 6.5 GEV/C
M	2356.0	10.0	JENKINS 83 MPS	K-P TO K+ MM
M	AVG 2367.7	7.0	AVERAGE (ERROR INCLUDES SCALE FACTOR OF 1.1)	

## $\Xi(2370)$ WIDTH (MeV)

W	75.0	69.0	DIBIANCA 75 DBC	XI 2P1
W	94 80.0	25.0	AMIRZAD 80 HBC	-0 K-P AT 8.25 GEV
W	50 (80.0)		HASSALL 81 HBC	-0 K-P AT 6.5 GEV/C
W	AVG 79.4	23.5	AVERAGE	

## $\Xi(2370)$ PARTIAL DECAY MODES

P1	$\Xi(2370) \rightarrow \Lambda \bar{K} \pi$ (including P4+P6)	DECAY MASSES
P2	$\Xi(2370) \rightarrow \Sigma \bar{K} \pi$ (including P5+P6) <td>1116+ 498- 140</td>	1116+ 498- 140
P3	$\Xi(2370) \rightarrow \Omega^- K$	1197+ 498- 140
P4	$\Xi(2370) \rightarrow \Lambda \bar{K}^*(892)$	1672+ 498
P5	$\Xi(2370) \rightarrow \Lambda \bar{K}^*(892)$	1116+ 892
P6	$\Xi(2370) \rightarrow \Sigma \bar{K}^*(892)$	1197+ 892
		1385+ 498

## $\Xi(2370)$ BRANCHING RATIOS

$\Xi(2370) \rightarrow (\Lambda \bar{K} \pi)/\text{total}$	(P1)
R1 SEEN	AMIRZAD 80 HBC -0 K-P AT 8.25 GEV
$\Xi(2370) \rightarrow (\Sigma \bar{K} \pi)/\text{total}$	(P2)
R2 SEEN	AMIRZAD 80 HBC -0 K-P AT 8.25 GEV
$\Xi(2370) \rightarrow (\Lambda/\Sigma \bar{K}^* \pi)/\text{total}$	(P1+P2)
R3 50 SEEN	HASSALL 81 HBC -0 K-P AT 6.5 GEV/C
$\Xi(2370) \rightarrow (\Omega^- K)/\text{total}$	(P5)
R4 0.09 0.04	KINSON 80 HBC - K-P AT 8.25 GEV
$\Xi(2370) \rightarrow (\Lambda/\Sigma \bar{K}^*(892))/\text{total}$	(P4+P5)
R5 0.22 0.13	KINSON 80 HBC - K-P AT 8.25 GEV
$\Xi(2370) \rightarrow (\Sigma(1385) \bar{K})/\text{total}$	(P6)
R6 0.12 0.08	KINSON 80 HBC - K-P AT 8.25 GEV

## REFERENCES FOR $\Xi(2370)$

DIBIANCA 75 NP 898 137 DIBIANCA, ENDORF (CERN)  
 AMIRZAD 80 PL 908 324 AMIRZADEH, (BIRM-CERN+GLAS+MSU-LPMP) I  
 KINSON 80 TORONTO CONF. 263 J B KINSON, (BIRM-CERN+GLAS+MSU-LPMP) I  
 HASSALL 81 NP B189 397 +ANSORGE, CARTER, NEALE, RUSHBROOKE+(CAMB+MSU)  
 JENKINS 83 PRL 51 951 +ALBRIGHT, DIAMOND, -(FSU+BRAN+LBL+CINC+SMAS)

## $\Xi(2500)$

$$I(J^P) = \frac{1}{2} ( ) \quad \text{Status: *}$$

OMITTED FROM SUMMARY TABLE

THE ALITTI 69 PEAK MIGHT BE INSTEAD THE XI(2370) OR MIGHT BE NEITHER THE XI(2370) NOR THE XI(2500).

## $\Xi(2500)$ MASS (MeV)

M	30 2430.0	20.0	ALITTI 69 HBC	K-P 4.6-5 GEV/C
M	45 2500.0	10.0	BARTSCH 69 HBC	-0 K-P 10 GEV/C
M	2505.0	10.0	JENKINS 83 MPS	K-P TO K+ MM

## $\Xi(2500)$ WIDTH (MeV)

W	150.0	60.0	40.0	ALITTI 69 HBC
W	59.0	27.0		BARTSCH 69 HBC -0

## $\Xi(2500)$ PARTIAL DECAY MODES

P1	$\Xi(2500) \rightarrow \Xi \pi$	DECAY MASSES
P2	$\Xi(2500) \rightarrow \Lambda \bar{K}$	1321+ 140
P3	$\Xi(2500) \rightarrow \Sigma \bar{K}$	1116+ 498
P4	$\Xi(2500) \rightarrow \Xi(1530) \pi$	1197+ 498
P5	$\Xi(2500) \rightarrow \Lambda$ (or $\Sigma$ ) $\bar{K} \pi$	1533+ 140
P6	$\Xi(2500) \rightarrow \Xi \pi \pi$	1116+ 498- 140
		1321+ 140- 140

## $\Xi(2500)$ BRANCHING RATIOS

$\Xi(2500) \rightarrow (\Xi \pi)/(\text{modes P1 to P4})$	(P1)/(P1+P2+P3+P4)
R1 (0.5) OR LESS	ALITTI 69 HBC 1 STD DEV LIMIT
$\Xi(2500) \rightarrow (\Lambda \bar{K})/(\text{modes P1 to P4})$	(P2)/(P1+P2+P3+P4)
R2 0.5 0.2	ALITTI 69 HBC -
$\Xi(2500) \rightarrow (\Sigma \bar{K})/(\text{modes P1 to P4})$	(P3)/(P1+P2+P3+P4)
R3 0.5 0.2	ALITTI 69 HBC -
$\Xi(2500) \rightarrow (\Xi(1530) \pi)/(\text{modes P1 to P4})$	(P4)/(P1+P2+P3+P4)
R4 (0.2) OR LESS	ALITTI 69 HBC 1 STD DEV LIMIT
$\Xi(2500) \rightarrow (\Lambda$ (or $\Sigma$ ) $\bar{K} \pi)/\text{total}$	(P5)
R5 SEEN	BARTSCH 69 HBC -0
$\Xi(2500) \rightarrow (\Xi \pi \pi)/\text{total}$	(P6)
R6 SEEN	BARTSCH 69 HBC -0

For notation, see key on page 91.

## Baryon Full Listings

$\Xi(2500)$ ,  $\Omega^-$ ,  $\Lambda_c^+$ ,  $\Sigma_c(2450)$ ,  $\Xi_c^+$ ,  $\Omega_c^0$ ,  $\Lambda_b^0$ , DIBARYONS

### REFERENCES FOR $\Xi(2500)$

ALITTI 69 PRL 22 79 +BARNES,FLAMINIO,METZGER, + (BNL-SYRA) I  
BARTSCH 69 PL 288 439 + (AACH-BERL+CERN-LOIC-VJEN)  
JENKINS 83 PRL 51 951 +ALBRIGHT,DIAMOND,-(FSU-BRAN+LBL+CINC+SMAS)

### $\Omega$ BARYONS ( $S=-3$ , $I=0$ )

$\Omega^-$

$I(J^P) = 0( )$  Status: \*\*\*\*

SEE STABLE PARTICLES.

### CHARMED BARYONS

$\Lambda_c^+$

$I(J^P) = 0( )$

SEE STABLE PARTICLES.

$\Sigma_c(2450)$

$I(J^P) = 1( )$  Status: \*\*

OMITTED FROM SUMMARY TABLE

THE SIGMA/C DECAYS TO LAMBDA/C P1, AND THE SCHISM IN MASSES HERE REFLECTS THAT IN MEASUREMENTS OF THE LAMBDA/C MASS (THE HIGHER MASSES ARE PRESENTLY FAVORED). THE IMPRESSIVE AGREEMENT ON THE SIGMA/C-LAMBDA/C MASS DIFFERENCE STRONGLY INDICATES THIS TO BE THE CASE, RATHER THAN THAT TWO STATES (THE SIGMA/C AND  $\Sigma_c^*$ ) ARE BEING OBSERVED.

THIS PARTICLE IS AT ABOUT THE 2-AND-1/2-STAR LEVEL. A DEFINITIVE EXPERIMENT IS NEEDED.

### $\Sigma_c(2450)$ MASS (MeV)

M	1	2426.0	12.0	CAZZOLI	75	HBC	++	NU	P	IN	BNL	7-FT
M	9	(2460.0)		KNAPP	76	SPEC	-	GAMMA	BE			
M	1	(2439.0)	OR MORE	BARTSH	77	DBC	--	NU	D	IN	12-FT	
M	6	2425.0	10.0	BALTAY	79	HLBC	--	NU	NE-H	IN	15-FT	
M	1	2457.0	4.0	CALICCHIO	80	HBC	+	NU	P	IN	BEBC-TST	
M	1	2454.0	5.0	BOSETTI	82	HBC	++	NU	P	IN	BEBC	
M	1	(2480.0)		ADAMOVIĆ	84	EMUL	--	GAMMA	NUC--	OMEGA		

### $\Sigma_c(2450) - \Lambda_c^+$ MASS DIFFERENCE (MeV)

D	1	166.0	15.0	CAZZOLI	75	HBC	++	NU	P	IN	BNL	7-FT
D <td>6</td> <td>168.0</td> <td>3.0</td> <td>BALTAY</td> <td>79</td> <td>HLBC</td> <td>++</td> <td>NU</td> <td>NE-H</td> <td>IN</td> <td>15-FT</td> <td></td>	6	168.0	3.0	BALTAY	79	HLBC	++	NU	NE-H	IN	15-FT	
D <td>1</td> <td>168.0</td> <td>3.0</td> <td>CALICCHIO</td> <td>80</td> <td>HBC</td> <td>+</td> <td>NU</td> <td>P</td> <td>IN</td> <td>BEBC-TST</td> <td></td>	1	168.0	3.0	CALICCHIO	80	HBC	+	NU	P	IN	BEBC-TST	
D <td>1</td> <td>166.0</td> <td>1.0</td> <td>BOSETTI</td> <td>82</td> <td>HBC</td> <td>++</td> <td>NU</td> <td>P</td> <td>IN</td> <td>BEBC</td> <td></td>	1	166.0	1.0	BOSETTI	82	HBC	++	NU	P	IN	BEBC	

### $\Sigma_c(2450)$ PARTIAL DECAY MODES

P1	$\Sigma_c(2450) \rightarrow \Lambda_c^+ \pi$	DECAY MASSES
		2282 + 140

### REFERENCES FOR $\Sigma_c(2450)$

CAZZOLI 75 PRL 34 1125 -OMPS,CONNOLLY,LOUITTIT,MURTAGH, - (BNL)  
KNAPP 76 PRL 37 882 -LEE,LEUNG,SMITH, - (COLU+HAWA-ILL+FNAL)  
BARTSH 77 PR D15 1 +DERRICK,DOMBECK,MUSGRAVE, - (ANL+PURD)  
BALTAY 79 PRL 42 1721 -CAROUBALIS,FRENCH,HIBBS, - (COLU+BNL)I  
CALICCHIO 80 PL 959 521 + (BARL-BIRM+BRUX+CERN+EPOL+RHEL+SACL+LOU)  
BOSETTI 82 PL 1098 234 +GRASSLER, + (AACH+SOHN+CERN-MPI-M-ORF)  
ADAMOVIĆ 84 PL 1408 119 ADAMOVIĆ+ (BGNA+CERN-FIRZ+GENO+MADR+LEBD+)

#### THEORY AND REVIEW

DERUJULA 75 PR D12 147 +GEORGI,GLASHOW (HARV)  
LEE 77 PR D15 157 +QUIGG,ROSNER (FNAL)  
TRILLING 81 PRPL 75 57 G H TRILLING (LBL)

$\Xi_c^+$   
was  $A^+$

$I(J^P) =$

OMITTED FROM SUMMARY TABLE

SEE STABLE PARTICLES.

$\Omega_c^0$   
was  $T^0$

$I(J^P) =$

OMITTED FROM SUMMARY TABLE

SEE STABLE PARTICLES.

### BOTTOM (BEAUTY) BARYON

$\Lambda_b^0$

$I(J^P) =$

OMITTED FROM SUMMARY TABLE

SEE STABLE PARTICLES.

### NOTE ON DIBARYON RESONANCES

(by L.D. Roper, Virginia Polytechnic Institute and State University)

Probably the first modern theoretical discussion of dibaryon resonances was by Oakes.<sup>1</sup> The first experimental hints were in a  $\Delta p$  invariant mass distribution by Dahl et al.,<sup>2</sup> and in the  $^1D_2$  state in a  $pp$  partial-wave analysis by Arndt.<sup>3</sup> [The notation is  $(2S+1)L_J$ , where  $S$ ,  $L$ , and  $J$  are the spin, orbital, and total angular momenta. The Pauli principle restricts two nucleons to the following states:

$$I=0: (^3S_1, ^3D_1), ^1P_1, ^3D_2, (^3D_3, ^3G_3), ^1F_3, ^3G_4$$

$$I=1: ^1S_0, ^3P_0, ^3P_1, (^3P_2, ^3F_2), ^1D_2, ^3F_3, \dots$$

Here the states that couple together (same  $J^P$ ) are grouped together in parentheses. Similarly, only certain states are allowed for  $\Lambda\Lambda$ , etc.]

Interest in dibaryons rose dramatically in 1977 when strong energy dependence was unexpectedly observed in  $pp$  polarization experiments at Argonne.<sup>4</sup> At about the same time, Hoshizaki claimed to have found dibaryon resonances in a  $pp$  partial-wave analysis,<sup>5</sup> and Jaffe gave a detailed theoretical treatment of multi-quark states.<sup>6</sup> There is now a vast literature on dibaryon resonances.



## Baryon Full Listings

### DIBARYONS

However, there is still disagreement about what “dibaryon resonances” are. There is little doubt there are distinct structures in  $NN$  partial-wave amplitudes that look very much like ordinary highly inelastic resonances, such as are seen in  $\pi N$  scattering. The question is whether these structures are caused by resonance poles in the complex energy plane or by some other structure of the scattering amplitude.

One of the arguments about dibaryon resonances is whether they are calculable in terms of quark theory or should instead be calculated using some hadron interaction theory without reference to the underlying quarks. Both approaches have had successes and failures, so possibly both will make contributions to unraveling the mysteries of dibaryon resonances.

The idea that dibaryon resonances are “pseudo-resonances”<sup>7</sup> has taken some new turns. The idea is that box diagrams (e.g., involving  $N\Delta$  in  $NN$  scattering) create resonance-like loops in the Argand diagram without resonance poles actually existing. The problem with believing this is whether poles would be created when one unitarizes the box-diagram calculations in order to calculate physical scattering amplitudes. Kloet and Tjon<sup>8</sup> have recently shown that a model exists in which, indeed, this is the case. However, resonance hunters should definitely report pole positions rather than looping Argand diagrams in the future. All who suggest that the  $NN$   $^1D_2$  and  $^3F_3$  resonance-like structures are resonances or are instead due to some other dynamics must take their case to the world collection of  $NN$  scattering data in the form of a detailed partial-wave analysis, and should report the existence or nonexistence of resonance poles.

Closely related to the work described above is that of Ueda,<sup>9</sup> in which Faddeev  $\pi NN$  dynamics is fitted to the  $NN$  partial-wave amplitudes. Although the fit is not good, the approximately correct structure is present. The interesting point is that poles do occur on the “resonance” sheets in the complex energy plane. Ueda claims this work is important because many claims for resonance poles assume that the poles exist, but the Faddeev approach does not make such prior assumptions.

VerWest<sup>10</sup> recently reported separable potential model fits to the  $^1D_2$  and  $^3F_3$   $NN$  amplitudes and claims some solutions had no resonance poles, but Kloet and Tjon<sup>11</sup> have shown that these potential model fits all do, indeed, have resonance poles.

Dinucleon resonances also communicate with the  $\gamma d$  and  $\pi d$  channels. There is not much  $\gamma d$  data, and the multipole analysis does not tell much about which

dibaryons might be involved. In the  $\pi d$  channel, uncertainties abound, and partial-wave analyses yield poor fits compared to  $NN$  analyses. In addition to  $NN$  analyses, results from analyses of  $\pi d \rightarrow \pi d$ ,  $\pi d \rightarrow \pi pn$ ,  $pp \rightarrow \pi d$ , and  $\gamma d \rightarrow pn$  scattering are listed below. Most of these strongly indicate the existence of dibaryon resonances in the  $^1D_2$  and  $^3F_3$   $NN$  states, and some indicate possible resonances in the  $^1S_0$ ,  $^3S_1$ ,  $^3P_1$ ,  $^3P_2$ ,  $^3D_3$ ,  $^1F_3$ ,  $^1G_4$ , and other states.

Since our last edition, many more papers have been published about dibaryon resonances. There are eleven new references for dinucleons and one for strange dibaryons giving values for the resonance-pole (or Breit-Wigner) parameters. Many of the new references are fits to  $\pi d$  partial-wave amplitudes, in which the couplings to dibaryons appear to be larger than in  $NN$  amplitudes.

Some notable papers in which narrow (width  $< 20$  MeV) dibaryon production peaks are seen are:

- (1) Siemiarczuk et al.<sup>12</sup> for the  $dp \rightarrow (np)p$  and  $dp \rightarrow p\pi X$  reactions;
- (2) Bairamov et al.<sup>13</sup> for  $pC \rightarrow npX$ ;
- (3) Zelinski et al.<sup>14</sup> for  $^4\text{He}p \rightarrow dppn$ ; and
- (4) Tatischeff et al.<sup>15</sup> for  $p^3\text{He} \rightarrow dX$  and  $^3\text{He}p \rightarrow dX$  reactions.

However, Katayama et al.<sup>16</sup> claims to have done an experiment similar to that of Siemiarczuk et al. and sees no peaks; and Combes et al.<sup>17</sup> sees no narrow dibaryon resonances in  $dd \rightarrow dX$ . Tatischeff<sup>15</sup> suggests that the reason some experiments do not see these narrow peaks is that they appear as small bumps on a large background and may not be observable except at the maxima of the production cross sections.

Since our last edition, only one paper has appeared giving data on strange dibaryon states. It appears that the strangeness  $-1$  dominant resonance is in the  $^3S_1$  state, an SU(3) partner of the deuteron. Dalitz,<sup>18</sup> in an excellent review, concludes that the  $S = -1$   $^3S_1$  resonance pole probably exists. However, May et al.<sup>19</sup> sees no enhancements in the  $\Sigma^+n$ ,  $\Sigma^0p$ , or  $\Lambda p$  invariant mass spectra in  $K^-d \rightarrow \pi^-X$ ; and Arenton et al.<sup>20</sup> sees none in the  $\Lambda p$  spectrum in  $pp \rightarrow \Lambda pK^+$ .

In the Listings below, we separate the determinations of pole positions and Breit-Wigner parameters. To be a resonance, the pole must occur on the lower half of the second sheet for the elastic channel; it may be a bound state or resonance for inelastic channels.

In summary, this reviewer feels that the evidence, both experimental and theoretical, for the  $^1D_2$  and  $^3F_3$  dinucleon resonances is now very strong. The theoretical calculations almost all now agree that resonance





For notation, see key on page 91.

Baryon Full Listings
DIBARYONS

Table with 4 columns: Particle ID, Energy (MeV), Width (MeV), and Reference/Notes. Includes entries like GREIN 82 NNF 3S1 OR 3D3, AKEMOTO 83 PIDC Q=0, etc.

Table with 4 columns: Particle ID, Energy (MeV), Width (MeV), and Reference/Notes. Includes entries like BERETVAS, NIELO, SPINKA+, BRISSAUD, DIDELEZ, PERRIN+, etc.

S = -1 DIBARYON

BARYON NUMBER 2, STRANGENESS -1 STATES

IN THIS SECTION WE USE THE FOLLOWING ABBREVIATIONS FOR TYPES OF ANALYSES--

Legend table: BB BARYON-BARYON SCATTERING COMBINED AMPLITUDE ANALYSIS, LNIM LAMBDA-N INVARIANT MASS, LPIM LAMBDA-P INVARIANT MASS, etc.

N(2130) I=1/2, 3S1

Status: \*\*

OMITTED FROM SUMMARY TABLE

B=2, S=0 MISCELL. -- BREIT-WIGNER WIDTH (MeV)

Table with 4 columns: Particle ID, Energy (MeV), Width (MeV), and Reference/Notes. Includes entries like ALADASHVILI 76 DB PP INVARIANT MASS, GREIN 80 NNF 3S1 OR 3D3, etc.

B=2, S=0 MISCELL. -- BREIT-WIGNER ELASTICITY

Table with 4 columns: Particle ID, Energy (MeV), Width (MeV), and Reference/Notes. Includes entries like HASHIMOTO 80 NN 1F3 ASSUMED BCKGRND, LOCHER 83 PID 164, etc.

REFERENCES FOR B=2, S=0 STATES

ALADASHVILI, ELAGOLEV+ (JINR+WARS+WINR), GREIN, P KROLL (KARL+WUPP), N HOSHIZAKI (KYOT), etc.

B=2, S=-1 -- BREIT-WIGNER MASS (MeV)

BREIT-WIGNER MASS APPROXIMATELY EQUALS RE (POLE POSITION).

Table with 4 columns: Particle ID, Energy (MeV), Width (MeV), and Reference/Notes. Includes entries like COHN 64 LNIM Q=0, CLINE 68 LPIM 3S1 Q=1, etc.

B=2, S=-1 -- BREIT-WIGNER WIDTH (MeV)

BREIT-WIGNER WIDTH APPROXIMATELY EQUALS 2 TIMES IM (POLE POSITION).

Table with 4 columns: Particle ID, Energy (MeV), Width (MeV), and Reference/Notes. Includes entries like COHN 64 LPIM Q=0, CLINE 68 LPIM 3S1 Q=1, etc.

# Baryon Full Listings

## DIBARYONS

### B=2, S=-1 --- Re(POLE POSITION) (MeV)

Re(POLE POSITION) APPROXIMATELY EQUALS BREIT-WIGNER MASS.

RE	J	(2132.0)	NAGELS	79	BB	3S1	Q=1	
RE	J	(2137.0)	NAGELS	79	BB	3S1	Q=0	
RE	B	(2129.0)	DOSCH	80	LPIM	3S1	Q=1	
RE	J	(2127.0)	TAKAHASHI	80	BB	3S1	Q=0,1,2	
RE	J	(2148.0)	TAKAHASHI	80	BB	1P	Q=0,1,2	
RE	J	NAGELS 79 REPORTS POLE POSITION FOR TWO DIFFERENT 3S1 CHARGE STATES.						

### B=2, S=-1 --- Im(POLE POSITION) (MeV)

Im(POLE POSITION) APPROXIMATELY EQUALS ONE-HALF BREIT-WIGNER WIDTH.

IM	J	(2.4)	NAGELS	79	BB	3S1	Q=1
IM	J	(2.6)	NAGELS	79	BB	3S1	Q=0
IM	B	(3.0)	DOSCH	80	LPIM	3S1	Q=1
IM	J	(9.0)	TAKAHASHI	80	BB	3S1	Q=0,1,2
IM	J	(8.0)	TAKAHASHI	80	BB	1P	Q=0,1,2

### REFERENCES FOR B=2, S=-1 STATES

COHN	64	PRL	22	668	H O COHN, K H BHATT, W M BUGG	(ORNL-TENN)
CLINE	68	PRL	20	1452	D CLINE, R LAUMANN, J MAPP	(WISC)
ALEXANDE	69	PRL	22	483	ALEXANDER, HALL, JEW, KALMUS, KERNAN	(LBL-UCR)
JAIN	69	PR	167	1816	P L JAIN	(SLAC)
TAN	69	PRL	23	395	T H TAN	(BUFF)
EASTWOOD	71	PR	D3	2603	=FRY, HEATHCOTE, ISLAN+ (BIRM+EDIN+GLAS+LOIC)	(FSU)
SIMS	71	PR	D3	1162	=O'NEAL, ALBRIGHT, BRUCKER, LANUTTI	(JINR)
SHAHBAZI	73	NP	B53	19	B SHAHBAZIAN, A TIMONINA	(JINR)
SODHI	75	NP	B97	403	A SODHI, D GOYAL	(DELH)
BRAUN	77	NP	B124	45	=GRIMM, HEPP, STROEBELE, THOEL-	(HEID+MPIM)
GOYAL	78	PR	D18	948	D GOYAL, A SODHI	(DELH)
NAGELS	79	PR	D20	1633	M NAGELS, T RIJKEN, J DESWART	(NIJM)
DOSCH	80	ZPHY	C3	249	H DOSCH, I STAWATESCU	(HEID)
GOYAL	80	PTP	64	700	D GOYAL, J MISRA	(DELH)
TAKAHASHI	80	NP	A336	347	TAKAHASHI, IWAMURA, KIMURA, KUME	(TOKY)
SHAHBAZI	82	NP	A374	73C	SHAHBAZIAN, TEMNIKOV, TIMONINA	(JINR)
PIGOT	85	NP	B249	172	=DE BRION, CAILLET, CHEZE-	(ROMA-SACL+VAND)
PAPERS NOT REFERRED TO IN DATA LISTINGS						
PIROUE	64	PL	11	164	P A PIROUE	(PRIN)
ALEXANDE	68	PR	173	1452	=KASHORN, SHAPIRA+	(REHO+HEID)
BUNNEL	70	PR	D2	98	=DERRICK, FIELDS, HYMAN, KEYES	(NWES-ANL)
GOYAL	71	PR	D5	1259	D P GOYAL	(DELH)
KADYK	71	NP	B27	13	=ALEXANDER, CHAN, GAPOSCHKIN, TRILLING	(LBL)
DOSCH	78	PR	D18	4071	H G DOSCH, V HEPP	(HEID)
MIZUNO	79	PTP	62	1691	T MIZUNO	(TOKY)
ROSEN	79	NO	B49	217	=VANDERVELDE-WILQUET, WICKENS-	(LOUC+BRUX)
D'AGOSTI	81	PL	1048	330	(ROMA-SACL+VAND)	
KIMURA	81	PTP	65	649	M KIMURA, Y IWAMURA, Y TAKAHASHI	(TOKY)
TOKER	81	NP	A362	405	G TOKER, A GAL, J EISENBERG	(HEBR+TELA)
AERTS	85	NP	B253	116	A T M AERTS, C B DOVER	(CERN-IPN)

## S=-2 DIBARYON

OMITTED FROM SUMMARY TABLE

BARYON NUMBER 2, STRANGENESS -2 STATES

IN THIS SECTION WE USE THE FOLLOWING ABBREVIATIONS FOR MEASURED QUANTITIES--

LLIM	LAMBDA-LAMBDA INVARIANT MASS
LLPI	LAMBDA-LAMBDA-PI INVARIANT MASS
XPIM	XI-P INVARIANT MASS

### B=2, S=-2 --- MASS (MeV)

M	A	(2367.0)	(4.0)	BEILLIERE 72	LLIM	Q=0	GAUSSIAN FIT
M	B	(2365.3)	(9.6)	SHAHBAZIA 73	LLIM	Q=0	
M	C	(2480.0)		GOYAL 80	XPIM	Q=0	
M	A	K-D TO XI-P Q=0.					
M	B	N P TO LAMBDA LAMBDA X AND PI- P TO LAMBDA LAMBDA X FOR P IN C12.					
M	C	GOYAL 80 ALSO SEES A SHOULDER AT 2360 MEV.					

### B=2, S=-2 --- WIDTH (MeV)

W	A	(15.0)	(4.0)	BEILLIERE 72	LLIM	Q=0	GAUSSIAN FIT
W	B	(47.0)	(15.7)	SHAHBAZIA 73	LLIM	Q=0	

### REFERENCES FOR B=2, S=-2 STATES

BEILLIERE	72	PL	398	671	BEILLIERE, MAYEUR+	(BRUX-CERN-TUFT-LOUC)
SHAHBAZI	73	NP	B53	19	B SHAHBAZIAN, A TIMONINA	(JINR)
GOYAL	80	PR	D21	607	D GOYAL, J MISRA, A SODHI	(DELH)
SHAHBAZI	82	NP	A374	73C	SHAHBAZIAN, TEMNIKOV, TIMONINA	(JINR)
PAPERS NOT REFERRED TO IN DATA LISTINGS						
CARROLL	78	PRL	41	777	+CHIANG, JOHNSON, KYCIA, KI-	(BNL-PRIN)
D'AGOSTI	82	NP	B209	1	(INFN+SACL+VAND-CERN)	
AERTS	85	NP	B253	116	A T M AERTS, C B DOVER	(CERN-IPN)
WALCHER	85	NP	A434	343C	T WALCHER	(MPIH)

## ACCESSING AND USING PARTICLE PHYSICS DATABASES

A number of publicly accessible computer databases containing particle physics information now exist at various institutions. Some of these databases are for literature searching, allowing the user to locate papers of interest, while others contain actual numerical data. The following discussion gives some idea of what is available and how to get started accessing and using these databases. The two locations covered are SLAC and Rutherford Appleton Laboratory (RAL).

### The SLAC Particle Physics Databases

The databases of interest at SLAC are: (1) HEP, a literature-searching guide for all particle physics journal articles, preprints, reports, theses, etc., indexed by the standard bibliographic quantities; (2) DATAGUIDE, an adjunct to HEP, which indexes papers containing experimental data by accelerator, detector, beam momentum, reactions and particles studied; (3) PARTICLES (formerly RPP), containing the Full Listings from this Review of Particle Properties, indexed by particle and particle property; (4) REACTIONS, containing numerical data (e.g., cross sections, polarizations, etc.) on reactions; and (5) EXPERIMENTS, a guide to current and past particle physics experiments, indexed similarly to the HEP and DATAGUIDE databases.

All these databases are managed by the SPIRES database management system, which runs interactively under VM/CMS on SLAC's IBM 3081 computer. To enter SPIRES, once you are logged onto the computer, key in SPIRES. You can then obtain information about the database you are interested in by typing in, say, EXPLAIN PARTICLES. To actually access the database, enter, for example, SELECT PARTICLES. You may then find out what terms are available for searching on by keying in SHOW INDICES. To see the form of the contents of a particular index, say the PP (particle property) index of the PARTICLES database, key in BROWSE PP; this will give you an idea of what kinds of expressions appear in this index, and thus will suggest what form you should use in your search. Then to do an actual search for information, say for the RPP Full Listings on the  $\eta$  meson mass, you would key in a command like FIND PP ETA MASS, followed by the command TYPE; this would print out the Listings for the  $\eta$  mass. At any time, you may get help by typing in such commands as EXPLAIN EXPLAIN, EXPLAIN SHOW INDICES, EXPLAIN BROWSE, EXPLAIN FIND, EXPLAIN TYPE, etc. When you are finished searching, key in EXIT, which gets you out of SPIRES.

Anyone who has an account on the SLAC computing system can access these databases online. If you do not have an account and cannot find anyone who does (at main laboratories, ask at the library), please contact SLAC directly. More information on how to access and search the databases can be found in the report "A User's Guide to Particle Physics Computer-Searchable Databases on the SLAC-SPIRES System," LBL-19173, available from the Particle Data Group, Bldg. 50 — Room 308, Lawrence Berkeley Laboratory, Berkeley, CA 94720, USA. An extensive wall poster, "A Guide to VM SPIRES," is available from the Library, SLAC, P.O. Box 4349, Stanford, CA 94305, USA. You may also contact Alan Rittenberg at LBL (CMS-id AXRVX, tel. 415-486-4723, or 451-4723 on FTS), or Louise Addis at SLAC (CMS-id ADDIS, tel. 415-854-3300, ext. 2411).

### The Durham-RAL Particle Physics Databases

These databases contain compilations of current and past experimental particle physics data (e.g., reaction cross sections), and are available for interactive searching under CMS on both the Rutherford Appleton and CERN central computers. The topics included are: (1) two-body (and quasi-two-body) reactions; (2) hadron and photon one- and two-particle inclusive distributions; (3) lepton-produced inclusive data (i.e., deep inelastic scattering, structure functions, etc.); and (4) data from  $e^+e^-$  annihilations. The databases also contain complete bibliographic information on these and other related topics, status information of current particle physics experiments, and the Full Listings from this Review of Particle Properties. To insure that the databases are up to date, experimentalists are urged to send their data to the compilers immediately on completion.

The databases can easily be used by anyone having network access to the RAL computer (PSS address 23422351919169, then RLIB) or to CERN; a guest account (PDG password HEPDATA) is available at RAL for those who do not have their own CMS account. An EXEC file, HEPDATA, on the UDISK gives direct access to the databases, and contains an extensive built-in HELP facility to assist the unfamiliar user. Data are retrieved using a simple keyword-based search, and can be displayed in either tabular or (at RAL) graphical form.

For more information, or a guide to the service, please contact Mike Whalley at Durham University, England (CMS-id MRW; tel. 0385-64971, ext. 591), or Dick Roberts at RAL (CMS-id RGR; tel. 0235-21900, ext. 5259).

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† Omitted from this edition; see listed page number in Phys. Lett. **111B** (1982).

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