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## INVESTIGATION OF THE EXTERNAL FLOW ANALYSIS FOR DENSITY MEASUREMENTS AT HIGH ALTITUDE

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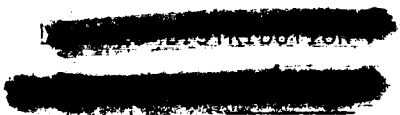
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Grant NSG-1630  
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National Aeronautics and  
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## I. INTRODUCTION

Accurate experimental determination or verification of aerodynamic force coefficients ( $C_D$  and/or  $C_L$ ) requires accurate simultaneous measurements of the forces (or accelerations) and the dynamic pressure ( $q = 1/2\rho U^2$ ). Comparison with theoretical predictions requires independent knowledge of the density ( $\rho$ ) and velocity ( $U$ ) to establish the proper values of the non-dimensional parameters such as Reynolds number ( $Re$ ) and Mach number ( $M$ ). These parameters in principle require independent measurement of temperature and measurement or inference of viscosity. Under hypersonic conditions during the early phases of re-entry the Mach number becomes a secondary parameter while the relevant Reynolds number is based on viscosity within the gas layer near the vehicle and is only weakly dependent on free stream temperature. During the earliest part of re-entry, independent knowledge of density is necessary to establish the degree of rarefaction generally measured by the Knudsen number  $Kn = \lambda/L$  where  $\lambda$  is a relevant mean free path and  $L$  the characteristic physical dimension (either vehicle size for the overall flowfield or entrance dimensions for the local behavior at the instrument).

The Shuttle Upper Atmosphere Mass Spectrometer (SUMS) Experiment<sup>(1)</sup> is designed to provide independent measurement of  $q = 1/2\rho U^2$  within the high altitude range. When combined with information of vehicle velocity, it will provide independent determination of upper atmosphere density and

coupled to accelerometer data will give the aerodynamic force coefficients within a regime difficult to simulate on the ground. The experiment is primarily intended to provide information between about 80 Km and 140 Km where rarefaction effects on the force coefficients are most important for a vehicle of the size of the Space Shuttle. It is also a regime where information on the atmosphere is relatively sparse as it lies below the altitude traversed by satellites and above that regularly assessed by ground launched meteorological vehicles. The interpretation of the measurements, however, requires an adequate understanding of the flowfield around the Space Shuttle within the vicinity of the SUMS experiment in order to provide the proper data reduction procedure and an assessment of the accuracy of the results.

At sufficiently low altitudes (below about 80 Km for the Space Shuttle), conventional pitot probe measurements can provide the dynamic pressure with straightforward data reduction and relatively minor corrections. At sufficiently high altitude (above about 150 Km), free molecular theory can be used to infer free stream conditions from surface measurements. The forces, however, are small and of little interest while the measurements require instruments of high sensitivity and are therefore difficult. In the intermediate range of altitudes where SUMS is designed to provide data, the typical molecular mean free path is of the same order as the characteristic vehicle dimensions. Figure 1 shows the

variation with altitude of the free stream Knudsen number ( $Kn_{\infty} = \lambda_{\infty}/D$ ) based on the free stream mean free path  $\lambda_{\infty}$  and the Shuttle diameter  $D$  at the location of the SUMS orifice just ahead of the wheel well. Note that  $Kn_{\infty} = .01$  at about 87 Km and  $Kn_{\infty} = 10.0$  at about 136 Km. Intermolecular collisions can, therefore, neither be neglected (free molecular theory) nor represented by the resultant transport properties (continuum theory) over the major portion of the SUMS measurement regime. The gas properties at the entrance to the instrument are, therefore, dependent on a flowfield that can only be determined on the basis of a "molecular" theory.

In addition to the above "external flow problem", needed to establish properties near the surface of the vehicle, the entrance region of the instrument is typically either smaller or comparable to a local mean free path. In such circumstances, the connection between the gas properties at some distance into the internal plumbing and those at the vehicle surface, can be very sensitive to the velocity and angular distribution of the incoming molecules. This requires both a high degree of detail from the "external flow" and a local analysis that must assess the molecular behavior at the instrument entrance. We shall refer to this as the "entrance problem". Figure 1 also shows an approximate band of Knudsen numbers for the entrance region of the SUMS experiment.  $Kn_s$  is based on the mean free path  $\lambda_s$  at the vehicle surface and the orifice diameter ( $d_o = .235$  cm) with surface properties fitted between free molecular results at high altitude and continuum Newtonian values at low altitude.  $Kn_c$  is a similar Knudsen number based on

conditions behind the "entrance" tube with  $\lambda_c$  estimated on free molecular results using Hughes and deLeeuw theory<sup>(2)</sup> at high altitude, with continuum constant pressure results applied at low altitude. Note that conditions within the entrance tube range from clearly free molecular behavior above about 110 Km to transitional behavior near 80 Km with fully continuum results only approached at the lowest altitude of interest.

The subsequent connection between the properties immediately behind the entrance tube and the measurements at the mass spectrometer shall be referred to as the "internal problem". The analytical procedures for calculating pressure profiles through the internal plumbing are well established and will be further verified by instrument calibration<sup>(1)</sup>.

During the preliminary phase of NASA Grant NSG 1630 (July 1979 to November 1979), the feasibility of examining the "external flow problem" for the Space Shuttle nose region within the relevant altitude range was established. A previously developed Direct Simulation Monte Carlo Computer Code<sup>(3,4,5)</sup> was found to be suitable as the starting point for this geometry and altitude range. Preliminary results were obtained at 87, 95, 105, 115 Km altitudes.

During the subsequent grant periods (November 1979 to September 1982) improvements in the modelling of the geometry and the molecular interactions have been incorporated in the external flow computer code. A number of runs at altitudes of 87, 95, 105, and 115 Km have been made to obtain a range of the relevant parameters and to provide input information

at the SUMS entrance location. The "entrance problem" has been examined both by using published information (2,6,7) and a previously developed Monte Carlo code for internal geometries (6). Because of the combination of entrance geometry (very long tube) and the range of local Knudsen number over the altitudes considered, a totally new "entrance" computer code had to be developed. This code provides the connection between the flux information at the orifice entrance obtained from the external code and the local gas properties behind the entrance tube where the gas is in equilibrium with the "cold" walls of the internal plumbing. This new code has only been exercised to a limited extent, but preliminary results relating the pressure within the tube behind the tile to the free stream dynamic pressure have been obtained. This information coupled with an appropriate calibration of the mass spectrometer provides the basis for a viable data reduction procedure of the SUMS experiment.

Section IIA contains a brief description of the operation of the EXTERNAL computer code (the detailed code is attached in Appendix A). Section IIB describes the issues associated with geometric modelling of the shuttle nose region and the modelling of intermolecular collisions including rotational energy exchange and a preliminary analysis of the vibrational excitation and dissociation effects. Section IIC discusses the selection of the trial runs and presents the major results.

Section IIIA contains a brief description of both the

original version and the modified present code (INTERNAL) for the entrance problem (Appendix C contains the code listing). Section IIIB contains a disucssion justifying the selection of geometric, collisional and surface modelling parameters used for the trial runs. Section IIIC presents the preliminary results and discusses the major effects.

Section IV presents the conclusions that can be drawn from the present study, provides a preliminary estimate of the data reduction procedure and suggests future work.



## II. EXTERNAL FLOWFIELD

The physical properties of the gas monitored by the SUMS instrument are not those of the free stream but are altered both by the intermolecular interactions in the external flowfield and by the combination of intermolecular and surface interactions within the entrance orifice and tubing leading to the instrument. External flowfield effects can be summarized in terms of the relation between local "stagnation" pressure and the free stream dynamic pressure ( $q = 1/2\rho U^2$ ) at sufficiently low altitudes. Within the "transition" regime of interest (80-140 Km) the gas entering the orifice is neither in equilibrium with the surface nor simply related to the free stream. The only currently available technique for describing the flowfield within this regime and providing sufficiently detailed data on the physical state of the gas at the surface is the Direct Simulation Monte Carlo Code.

### A. DIRECT SIMULATION MONTE CARLO COMPUTER CODE

"Monte Carlo" is the technique of using a simulated situation and random numbers to generate solutions from which information for the real case is then deduced statistically. The Monte Carlo approach ranges from being a strictly mathematical technique for evaluating the complicated multi-dimensional Boltzmann collision integral to a complete simulation of a number of molecules, with randomness

introduced only in the initial conditions. A modification of this latter approach is the one used in the present development. It consists of simultaneously following a large number of particles which yields, to some degree, a "direct simulation" of the processes taking place. Because there are finite limits on computer storage space, a modification to the direct simulation technique was developed by G.A. Bird (Ref. 5) wherein the real gas is simulated by several thousand "sample" particles populated into cells of the sample space considered. For collision calculations, all the particles in one cell are used as a representation of the local distribution function from which collision pairs are chosen at random, but in proportion to their collision probability based on the real gas. This preceding discussion applies to a general program incorporating the direct simulation procedure. A specific computer program for the generalized three-dimensional program for axisymmetric bodies in a hypersonic multi-fluid flow is described below.

The program (EXTERNAL) conducts numerical experiments with a model multi-component gas. The real gas is simulated by several thousand molecules which may be thought of as a representative sample of the many billions of molecules in the corresponding real gas. The positions and velocity components of the simulated molecules are stored in the computer and typical collisions are computed among them as a time parameter is advanced. Since the flow is at an angle of attack to the body, three position coordinates, three velocity components and appropriate internal energy levels

must be stored for each simulated molecule.

The computation of collisions starts at zero time with the molecules moving along the flow axis at the required freestream Mach number. The body is inserted into this flow at the zero time and the desired steady flow is obtained as the large-time solution of the resulting unsteady flow.

The free-stream flow vector lies in the x-y plane. The simulated region is bounded by the x-y plane as a plane of symmetry, an outer cylindrical boundary (the axis of the cylinder is the x axis), and two planes parallel to the y-z plane. These boundaries must be set sufficiently far from the body to eliminate interference. The simulated region is divided into a number of cells which are sufficiently small for the expected change in flow properties across the cells to be small.

The first step is to generate the initial, or zero time, configuration of molecules. The molecules are distributed over the simulated region and the velocity components assigned are appropriate to a gas in Maxwellian equilibrium and moving at the required Mach number. The body is then inserted into the flow and the molecules are allowed to move and collide among themselves. The move and collide processes are uncoupled by computing a number of collisions appropriate to a time interval  $\Delta t_m$  equivalent to a small fraction of the mean time between collisions, and then moving the molecules through distances appropriate to  $\Delta t_m$  and their instantaneous velocities. The distortion produced in the molecular paths by this approximation is small as long as  $\Delta t_m$  is small

compared with the mean time between collisions, and smaller than the typical transit time of a molecule through a cell.

Since the change in flow properties over the width of one cell is assumed small, the molecules in a cell at any instant may be regarded as a sample of the molecules at the location of the cell. The relative location of the various molecules within the cell can then be disregarded. A pair of molecules is chosen at random from those within the cell under consideration and is retained or rejected in such a way that the probability of retention is proportional to the relative collision probability for the appropriate interaction law. When a pair is retained, a typical collision is computed between the two molecules and the new velocity components and internal energies are stored in place of the old ones.

In general, the relative number ratio of the species of molecules in the multi-component gas will differ from unity, requiring the computation of different types of collisions. There is, therefore, one time counter for each type of collision in each cell. For each collision, the correct time counter is advanced for the cell by an amount appropriate to the collision parameters. The probability of collision, and therefore the time advancement per collision, is made proportional to the number of molecules in the cell, and the relative velocity and cross-sections of the colliding molecules. Collisions are computed in each cell until all the time counters have advanced through at least a time  $\Delta t_m$ . When this procedure has been carried out for all cells, the

overall experiment time is advanced by  $\Delta t_m$  and the molecules are moved through appropriate distances.

The set of molecules in each cell changes as the molecules are moved and appropriate conditions must be applied at the boundaries of the region being simulated. Every boundary is treated as a source of molecules with velocity components representative of molecules moving in thermal equilibrium at the appropriate fraction of the free-stream Mach number. (The fractional Mach number is determined by the cosine of the angle between the local boundary normal and the flow direction.) Any molecule which moves outward across any boundary is regarded as being lost and is removed from the sample. The plane of symmetry (the x-y plane) is regarded as a specularly reflecting surface. Interactions with the body are also computed. The body consists of a number of conic sections rotated about the axis of symmetry. Each section must be separately specified according to the coefficients of the defining equation, a procedure to be described later in the report. For the purpose of computing the momentum and energy transfers to the surface, each region of the body can be subdivided into smaller sections. Within these smaller sections, the following three parameters must be specified: wall temperature/gas temperature, energy accommodation coefficient for each species, and tangential (momentum) accommodation coefficient for each species.

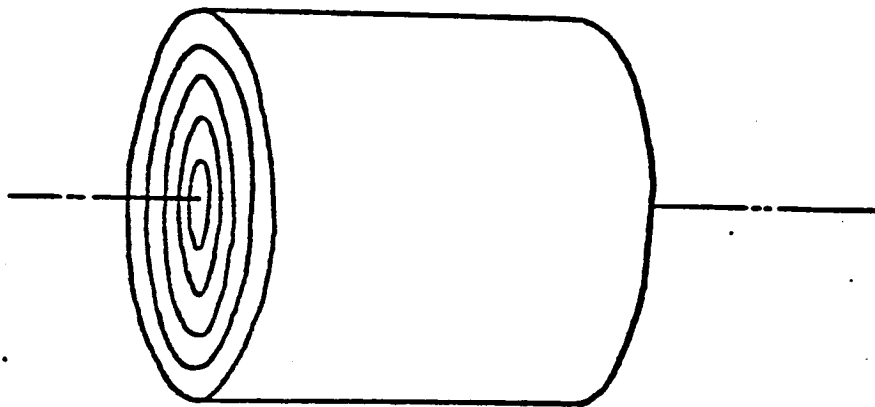
After the flow has settled down to a steady state, the number flux, momentum and energy transfers to the surface are

accumulated and used to compute the aerodynamic data. The time required to establish steady flow is usually assumed to be close to the time required for the free-stream flow to traverse several body lengths. The overall number flux, drag, and heat transfer coefficients are determined, along with their distribution along the surface.

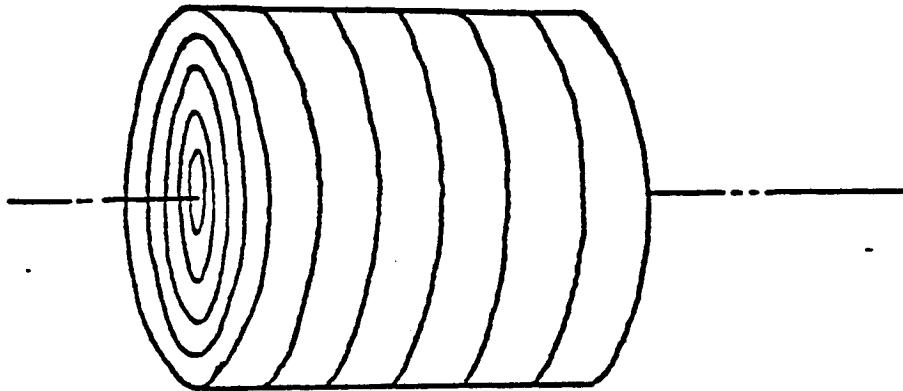
In addition, it is possible to generate data on the body surface which can be used as input to the companion program, INTERNAL, described separately. (INTERNAL) computes the flow field regime inside an axially symmetric cavity, such as might be used for a spacecraft-borne sensor. The input data needed for this computation includes the molecular distribution functions present at the orifice to the cavity, the orifice being on the surface of the spacecraft.) This data consists of velocity and internal energy samples in three coordinate directions for all species of the mixture.

Flow field properties are also computed. Instantaneous values are sampled at appropriate time intervals and these are time-averaged for greater accuracy. Number densities, velocities and temperatures are printed for each cell.

The numerical experiment takes place in a cylindrical block of space whose axis is coincident with the axis of revolution of the conic surfaces comprising the test body. This space is subdivided into cells in which the flow field properties of the experiment can be monitored. Cylindrical surfaces concentric with the axis partition the space into nested cylindrical volumes, as shown:

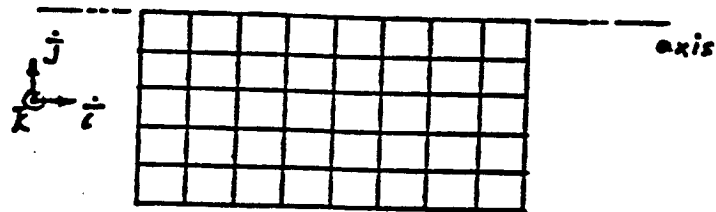


Planes parallel to the end-faces of the cylinders divide the cylinders into a stack of nested rings.



Finally, planes perpendicular to the preceding planes and passing through the axis, called radial planes, divide up the rings in the azimuthal direction, producing cells which are wedge-shaped pieces of rings. The geometry is more easily visualized if one considers a trace of the cell configuration in a radial plane. A typical planar trace is shown:





The axis shown is the axis of symmetry of the sample space (and of the test body). This axis is considered to be the x axis. When this direction must be specified in vector algebra computations, a unit vector  $\bar{i}$  is assumed in the x direction.

Assume that the planar trace shown above is bounded by the x-axis as described and by the -y axis. This plane is at  $0^\circ$  azimuth angle and is called the zero plane. It is the plane normally depicted when describing the sample space.

A unit vector  $\bar{j}$  points in the +y direction. The z direction points out of the paper, and in this direction is the unit vector  $\bar{k}$ , which is given by  $\bar{i} \times \bar{j}$ .

The flow velocity is in the direction  $\bar{i} \cos \alpha + \bar{j} \sin \alpha$ , where  $\alpha$  is the defined angle of attack between  $0^\circ$  and  $90^\circ$ . Since the flow is thus in the xy plane, there is symmetry in the z direction. That is, any condition in effect at +z is also in effect at -z. Thus azimuth angles need be specified only from  $0^\circ$  to  $180^\circ$  where  $0^\circ$  is in the -y direction,  $90^\circ$  is in the +z direction, and  $180^\circ$  is in the +y direction.

Now the fact is that the gas density in the vicinity of the stagnation point of the body can become many tens of times higher than the density far from the body. It is thus desirable to use small cells in this region while the cells are larger in the regions of relatively low density. This partitioning of cell sizes is accomplished in two ways.

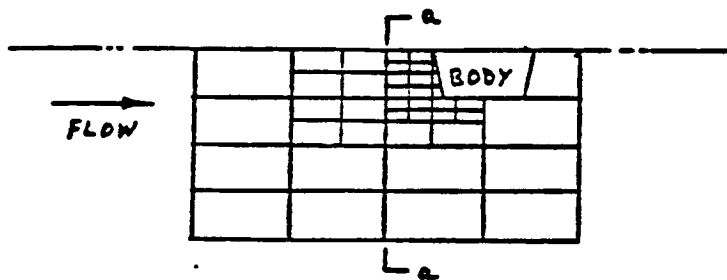
First, the rings can be divided azimuthally into two different sizes. This is done by specifying an angle, called THETAZ, and the number of wedge divisions both below and above this theta plane, called NWEDGE<sub>1</sub> and NWEDGE<sub>2</sub>. ("Below the theta plane" means azimuth angles between 0° and THETAZ, and "above the theta plane" means azimuth angles between THETAZ and 180°.)

Second, the cells below the theta plane can be subdivided in the axial and radial directions down to second and then third level cells. In the 0° radial plane representation of the sample space, this would appear as large rectangles being sub-divided into small rectangles.

In this way, the sample space geometry can be tailored to the configuration of the test body angle of attack to the flow. The following examples are presented to clarify the above statement. Assume for all cases that the test body in question is a short cylinder. As explained in the section TEST BODY, the cylinder cannot have flat end faces, so the ends are cones with apex angles of about 175°.

- a) For 0° angle of attack, the flow impinges directly on the left face of the cylinder.

It is thus desirable to have, if possible, constant azimuthal angles since there is no angles since there is no azimuthal assymetry. One possible configuration is therefore: THETAZ = 180° NWEDGE<sub>1</sub> = 6, NWEDGE<sub>2</sub> = 0 (producing 30° wedges), and the axial and radial directions can be subdivided any convenient way, producing a zero radial plane that looks like:



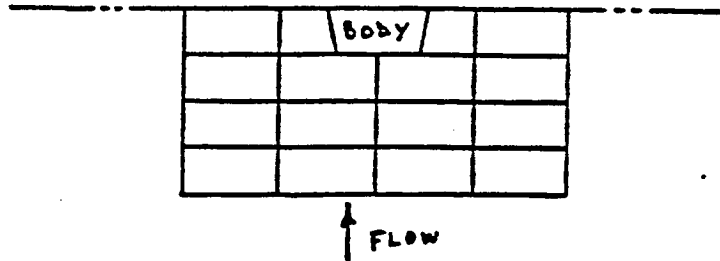
In this type of geometry, all radial planes are the same as the zero plane.

b) For 90° angle of attack, the flow impinges on

the curved cylindrical surface of the cylinder at the bottom. It is thus necessary to have small azimuthal wedges on the lower portion of the cylinder at and near the stagnation point, while larger wedges will suffice on the upper portions of the body (in the wake of the flow).

For instance, an acceptable set of parameters is:  $\text{THETAZ} = 60^\circ$ ,  $\text{NWEDGE}_1 = 3$ ,  $\text{NWEDGE}_2 = 3$  (producing  $20^\circ$  and  $40^\circ$  wedges), and again the axial and radial directions can be subdivided in any convenient way. Any radial plane up to  $60^\circ$  looks like the radial plane in example (a), with the body and smaller cells centered about a-a, while any radial plane between  $60^\circ$  and  $180^\circ$  looks like:

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- c) For an angle of attack between  $0^\circ$  and  $90^\circ$ , the configuration looks generally like that of example (b)). In this case, however, the theta plane generally should be at an azimuth angle which is relatively low (near  $60^\circ$ ) for a high angle of attack ( $45^\circ - 90^\circ$ ) and relatively high (near  $120^\circ$ ) for a low angle of attack.

Because the test body is located on the axis of the cylindrical sample space, for each particle that interacts directly with the body, many more do not. In the interest of minimizing the program running time necessary to permit a statistically sufficient number of particles to strike the body, the computation makes use of zonal weighting factors. That is to say, each particle in the sample-space in reality represents one or more particles, the actual number depending upon the weighting factor of the zone in which the particle currently exists.

Up to five cylindrical boundaries are selected across which the zonal weighting factors change. These boundaries

are specified in terms of the number of first level cells between the axis and the boundary. The change in the zonal weighting factor across each boundary can be given by:

$$LF_n = \frac{LD_{n+1} + LD_n}{n-1 \prod_{j=0}^{n-1} LF_j} \quad n=1,2,\dots,5$$

where the LD values are the number of first level cells between the axis and the cylindrical boundaries; and  $LF_0=1$ ,  $LD_6=LD_5=LD_K=NH$  were  $K \gg K_{last}$ .

This equation is the result of having the zonal weighting factors defined in such a way that they are equal to the ratio of sample space volume in the zone above the weighting-factor boundary to the sample space volume in the zone below the weighting factor boundaries. The importance of the region near the axis can be emphasized by choosing LF values larger than those given above.

The sample space is populated with a distribution of gas molecules. Each molecule is assigned a velocity, a rotational

energy and a position, such that the sample space is uniformly (albeit in a random manner) filled with a gas in thermal equilibrium and flowing at the required Mach number. Each molecule is assigned a number corresponding to its zonal weighting factor. The molecule thus represents in actuality a number of molecules (including itself) equal to the zonal weighting factor. While the molecule moves in such a way as to stay within the given zone, its weighting factor does not change. If it crosses a weighting factor boundary while moving in toward the axis, a number of molecules is added to the distribution. The number of molecules added is equal to:

$$\frac{\text{old weighting factor}}{\text{new weighting factor}} - 1$$

For instance, assume that above a boundary, the zonal weighting factor is 6, and below the boundary it is 3. Hence  $6/3 - 1$ , or 1, molecule must be added to the distribution when the molecule crosses the boundary. This is clearer if one considers that above the boundary, the molecule represents a total of 6 molecules. When the molecule drops below the boundary, it can only represent 3 molecules, so another molecule must be added to the distribution to represent the other 3 molecules.

The added molecule(s) is (are) assigned the same position and velocity components as the original molecule. While this does not approximate true kinetic theory at first, in practice the positions and velocities are soon randomized

by collision processes.

On the other hand, if a molecule crosses a weighting factor boundary while moving away from the axis of the sample space, there is a probability that it must be dropped from the molecular distribution. The probability is given by

$$1 - \frac{\text{old weighting factor}}{\text{new weighting factor}}$$

The random number generator is used to generate a histogram of disappearing molecules to match the actual probability of disappearance. The whole idea behind using weighting factors is to increase test body - flow field interaction in a given running time. Thus, when body surface quantities are accumulated (like flux, energy, momentum, heat), they are accumulated in terms of the weighted number of molecules striking the body. This is particularly important if the body exists in two or more weighting zones, so that surface quantities in different zones can be correctly compared.

The program is set up to handle a test body consisting of a sequence of connected conic sections rotated about a common axis of symmetry. Some typical surfaces which can be considered without modification are sections of spheres, cones, cylinders, ellipsoids, hyperboloids, and paraboloids. A disc perpendicular to the symmetry axis cannot be handled without modification, since it represents a multi-valued function in  $r$ . Cones with very large apex angles ( $\sim 180^\circ$ ) are



used in place of discs. The procedure used to specify body surfaces is to generate a form of the standard equation of the surface in question, substitute into this equation the actual values of the constants, and reduce the standard form to the following type of equation:

$$Ax^2 + Br^2 + Cx + D = 0$$

where A, B, C, and D are numerical values

These coefficients are used as input data to the program. The program requires additional parameters for each conic surface. Each surface can be axially divided into several segments for the purpose of accumulating the body surface parameters like flux, heat transfer, etc. The x-coordinate on the right side of the segment is required as data. (The segments are divided azimuthally by the same radial planes that partition the sample-space.) In addition, the temperature, energy accommodation coefficients and tangential accommodation coefficients for each species are required for each segment. The temperature of a surface is normalized by the free-stream temperature.

The tangential accommodation coefficient specifies the fractional part of incoming tangential momentum that is lost to the surface. For the reflection model used in this program, this also represents the fractional part of the molecules colliding with the body surface whose collisions are diffuse. The remainder collide specularly.

The energy accommodation coefficient specifies the ratio

of the net molecular energy flux absorbed by the body to that energy which would be absorbed if all re-emission were appropriate to equilibrium at the surface temperature of the body. This coefficient together with the tangential accommodation coefficient and the surface temperature determines the effective temperature of the diffuse component of the re-emitted molecules in the reflection model used in this program.

It is also possible to collect velocity samples of the colliding molecules on a restricted number of surface segments to generate distributions for the internal flow program. For this purpose, molecules that collide with the body surface are considered in two classes. The class called "UNCOLLIDED" consists of "free-stream" molecules. That is, these molecules have not previously collided with the body surface or with any other molecules other than "free-stream" molecules. The other class, called "COLLIDED", consists of molecules which have previously collided with either the body surface or with other "COLLIDED" molecules.

Appendix A contains the full listing of the computer code in Fortran IV. The main program consists of dimensioning statements coupled to a fairly detailed description of the input cards (using the NAMELIST input). The main operating program is called RUN which in turn calls the appropriate subroutines. Figure 2 shows a schematic of the flow chart for the operation of the Program EXTERNAL.

## B. Modelling of the SUMS Problem

### 1. Geometric Modelling

The computer code described above requires an axisymmetric geometry of the body, although the flow vector can be at an arbitrary direction. The Space Shuttle nose geometry in the vicinity of the SUMS orifice has to be modelled by an equivalent body of revolution. A paraboloid of revolution around an axis inclined at about  $8^\circ$  with respect to the shuttle axis models the lower surface cross sections of the body both in the symmetry plane and in the transverse direction reasonably well. Figure 3 shows a sketch of the actual and modelled Space Shuttle nose geometry. This model will be called the parabola model, and the flow direction of  $32^\circ$  with respect to the axis will be used to represent the  $40^\circ$  angle of attack of the Space Shuttle.

Both for the purpose of benchmarking the test runs against continuum results and to achieve better resolution an alternative axisymmetric flow model was examined. This model consists of a hyperbola rotated about the flow velocity vector passing through the stagnation point. Figure 4 shows a typical cell geometry for this model. This is also the model used by Professor Clark Lewis for his continuum calculations. As will be seen from the results in the next section the latter (axisymmetric) model is questionable as a representation of the flow in the stagnation region of the

Space Shuttle, at least at higher altitudes.

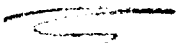
## 2. Modelling of Cross Sections

In order to avoid uncertainties associated with the choice of hard sphere cross-section to best model the "typical" collision, an energy dependent inverse power law cross-section collision code has been incorporated in the program. The power law exponent and reference cross-section is chosen to provide the best fit to the viscosity temperature dependence over the range between the wall and stagnation conditions which represents the energy range of the important collisions. Figure 5 shows that viscosity can be matched to within a few percent over the relevant range with a single choice of exponent and reference conditions

Using the exponent of  $N=.552$  and a reference temperature of  $1000^{\circ}\text{K}$  four axisymmetric cases of a hyperbola at  $0^{\circ}$  have been run to simulate 87KM, 95KM, 105KM and 115KM altitudes. Since in these cases rotational energy was not included these represent a fictitious monatomic gas of  $\gamma=5/3$ . Figure 6 shows "smoothed" temperature profiles normalized to the free stream temperature. Note that at the two lower altitudes a relatively "flat" Rankine Hugoniot (R.H.) region exists, while at 105KM R.H. conditions are barely reached and at 115KM the temperature peaks at about 70% of R.H. temperature. The shock layer in all cases is, however, very much thicker than the inviscid result based on nose radius  $R_N$ . Note, however, that the hyperbola is very blunt and it is not clear that the nose radius is the appropriate dimension or that this models properly the flow about the Space Shuttle Nose.

### 3. Rotational Energy Exchange

On the basis of previous experience in modelling rotational energy exchange in Reference 4 the model of Larsen and Borgnakke (9) was chosen as appropriate for the blunt geometry under investigation. The External Flow Program was therefore re-coded to include an arbitrary number of internal energy modes but with a single relaxation time related to the parameter  $\phi$  which ranges from  $\phi=0$  representing no exchange to  $\phi=1$  simulating maximum available exchange at each collision. An indication of the effect at 115KM is shown in Figure 7 and 8 for the most rapid energy exchange ( $\phi=1$ ). Note that, as expected, the peak temperature is lower for the  $\gamma=1.4$  case because some of the energy goes into rotation. Also note that significant non-equilibrium exists between translation and rotation even at this maximum energy exchange rate corresponding to  $\phi=1$ . The effect on surface properties shown in Figure 8 is virtually non-existent for the shear, but noticeable on the heat flux and pressure.



#### 4. Preliminary Attempts at Comparison with Continuum Results

Some preliminary comparisons of both the monatomic ( $\gamma = 1.667$ ) and diatomic ( $\gamma = 1.4$ ) runs for the axisymmetric hyperbola were made with continuum results provided by Prof. C.H. Lewis based on continuum theory (10,11). Among the surface properties only the pressure distribution agreed reasonably well. Since this is the property least sensitive to detailed flow field behavior the agreement is not a very sensitive test. A typical comparison is shown on Figure 8. A comparison of heat flux is not shown on this figure as it is not the same scale at this high altitude. The continuum result of Prof. Lewis gives the stagnation point heat transfer coefficient  $C_H = Q/1/2\rho U^3 = 2.8$  which is physically unrealistic. At 105KM the continuum heat transfer coefficient  $C_H$  is below one but still appreciably above the Monte Carlo results. At 95KM the Monte Carlo results give heat transfer that is almost twice as large as Prof. Lewis's result for  $\gamma = 1.4$ .

The most significant discrepancy between the continuum and Monte Carlo results arises in the shock layer thickness. While the definition is somewhat arbitrary, a comparison of the subsonic region in the vicinity of the stagnation point does give a good indication of the extent of influence of the "downstream" portions of the body. Figure 9 shows the Monte Carlo results for the  $M=1$  line at 115KM together with the shock line from C.H. Lewis. This discrepancy led to a whole re-examination of the modelling of the nose geometry.

#### *5. Re-examination of Geometric Modelling*

Since the subsonic region appears to extend to the "shoulder" wherever the hyperbola is terminated as shown in Figure 9, the applicability of the hyperbola model becomes suspect, at least for the Monte Carlo calculation where upstream influence cannot be eliminated. In order to further assess the effect of the geometric model on the results a comparison of the hyperbola at  $0^\circ$  to the parabola at  $32^\circ$  are presented in Figures 10 and 11. It can be seen from Figure 10 that the subsonic layer in the vicinity of the SUMS orifice is significantly different in the two cases. Figure 11 shows that while the pressure is not dramatically affected by the model the heat flux and shear are significantly altered. On the basis of these results it was concluded that while the runs for the hyperbola may indicate trends they are neither representative of the SUMS region on the Space Shuttle nor good candidates for benchmarking with the

continuum results. All subsequent runs were therefore made using the geometric model of a parabola at  $32^\circ$  angle of attack.

#### 6. *Inelastic vibrational Excitation and Dissociations*

Because of the high energies of collision between free-stream and reflected molecules, inelastic collisions (certainly vibration and dissociation) are in principle possible. At the highest altitude the number of such collisions is expected to be insignificant because the cross-sections are low enough so that a typical molecule will reach the body with a negligible probability of a previous vibrationally exciting or dissociating collision. At the lower altitudes 87 to 95Km the number of collisions suffered by a typical molecule before reaching the body surface is measured in the tens or hundreds and therefore inelastic cross-sections that are only a few hundredths of the elastic and rotational ones may produce significant effects. The detail needed to properly model these collisions in the Monte Carlo Programs far exceeds the available experimental information, which primarily gives overall rate constants. An investigation of the best combination of analytical and empirical information was initiated early within the grant period. A theoretical attempt to couple low energy vibrational excitation experimental information to the highly non-equilibrium high energy collisions through theoretical work was initiated. Appendix B contains a Master's thesis



presenting the formalism and giving the initial results of a theoretical formalism. Based on those results, coupled to some limited experimental data, a method for determining the probabilities of specific outcomes in individual collisions in the Monte Carlo Programs can be developed. The increase in computing times, however, may make the feasibility of using such a code, for anything but benchmarking, prohibitive.

### C. MAJOR RESULTS FOR EXTERNAL FLOW

The primary effort during the grant period up to February 1981 was spent in developing the code for external flow and establishing the appropriate geometry to model the region of the Space Shuttle in the vicinity of the SUMS experiment. The major results to that date are presented in the renewal proposal for the period February 1, 1981 to July 30, 1982 which also served as a progress report on the previous grant period. (12) We will only summarize those results here and update them on the basis of additional external flow computations.

A representation of the shuttle nose geometry as a paraboloid of revolution around an axis  $8^\circ$  from the actual shuttle axis (Figure 3) was found to adequately model the windward side of the shuttle in the vicinity of the SUMS orifice. Computations for both this model at an effective angle of attack of  $32^\circ$  and an alternative axisymmetric flow about hyperboloid model centered on an axis through the nominal stagnation point, demonstrated that the paraboloid model at angle of attack is necessary to adequately model conditions near the SUMS orifice. The typical body and computation cell geometry is shown in Figure 12.

A reasonable indication of the flowfield can be obtained by examining density contours or Mach number contours about the body. Figure 13 shows the sonic lines and the  $M = 5$  lines indicating approximately the outer extent of the "shock" layer at 95 Km and 115 Km, within the plane formed by

the velocity vector and the axis of the modelled paraboloid. Note the greater "shock" layer thickness at the higher altitude. Also note the fact that even at 95 Km, the shock layer is a large fraction of the local body dimension such as nose radius. The "shock" thickness is a major portion of the entire "shock layer" casting doubt on any calculation incorporating a thin "shock" assumption. Figure 14 shows some density contours on the "windward" side of the body at 95 Km altitude as well as the sonic line. Note the rather constant density rise towards the body with no discernible separation between "shock" and "boundary layer." Also shown are estimates of the stagnation streamline and another streamline along which the velocity only goes slightly subsonic. Note the rather gradual turning along the latter streamline and the rather diffuse nature of the shock layer even at this altitude where the nominal free stream Knudsen number is 0.04. In order to give a better picture of the three-dimensional aspect of the flowfield, a sketch of one-half of the paraboloid and some contour plots of  $M = 1$ . and  $M = 5$ . lines are shown on Figure 15.

Information in Figures 13 through 15 gives some indication of the nature of the flowfield in the vicinity of the SUMS orifice. The ultimate objective, however, is to establish properties at the vehicle surface. Figure 16 shows the variation of two surface fluxes (normal pressure and heat flux) at 95 Km along the four cross-planes shown in Figure 15. The normal pressure  $p_s$  normalized by  $1/2\rho U^2$  ( $C_p = p_s / 1/2\rho U^2$ ) is shown on the right side of the figure while a heat

transfer coefficient ( $C_H = Q/1/2\rho U^3$ ) is shown on the left. Figure 17 shows the variation versus angle in the cross-plane nearest the one containing the SUMS orifice, and compares the results to some theoretical and semi-empirical predictions. The pressure coefficient, as expected, lies approximately between the free molecular ( $C_p$  (F.M.)) value and the modified Newtonian ( $C_p$  (NEWT)) value, up to about  $90^\circ$ . (Note that at  $90^\circ$  around the axis the local angle ( $\beta$ ) between the velocity vector and the surface normal is approximately  $64^\circ$  at this cross-plane.) The heat transfer coefficient lies substantially below the free molecular value. It is also compared to a semi-empirical extrapolation of experimental results of stagnation heat transfer on hemispheres presented in reference 13. Direct comparison is clearly dubious due to the substantially different body geometry and the implicit assumption of totally local behavior contained in the  $\cos\beta$  variation with local angle of the surface to the velocity vector. Some additional degree of uncertainty is contained in the choice of "body size" that is used to evaluate the correlation parameter  $K_R^2$ .

The ultimate objective of the external flow calculations is to provide information on the properties of the gas entering the SUMS orifice. It is the potential non-equilibrium nature of the entering distribution of molecules that is responsible for the difference between the surface "pressure" and the gas pressure within the internal plumbing around the mass spectrometer. The non-equilibrium aspect is most commonly represented by a temperature jump and a

velocity slip within the continuum formalism. Under the highly rarefied conditions at the upper end of the altitude range of interest, even that description may be inadequate because of the highly non-Boltzmann distribution of the molecules arriving at the surface. Figure 18 shows the distribution function of the flux of molecules arriving at the SUMS orifice at 95 Km and 115 Km obtained from the external program. The flux distribution is plotted both versus molecular velocities normal to the local surface and tangential within the plane formed by the free stream velocity vector and the local normal. Note that at 95 Km one could fit the distribution by a Maxwellian with some temperature different than the body temperature and a slip velocity which is comparable to the free stream mean molecular speed. At 115 Km the distribution is clearly composed of two components, the free stream molecules arriving unperturbed at the surface, and the collided molecules having a broader distribution possibly representable by another Maxwellian. Clearly this potential bimodal character of the incoming molecules must be recognized in the evaluation of the "entrance" problem at the SUMS entrance.

### III. Entrance Problem

As a companion to the initial version of the external code (EXT) an internal code (INT) was developed (8). The objective of this code was to determine local properties and surface fluxes in a cavity connected by a tube to the exterior surface. The input molecular distribution is obtained from the surface flux information provided by the external code. While the code is in principle general to allow a wide variety of geometric configurations it is optimized for a short tube-large cavity geometry. Early attempts to apply this code directly to the SUMS inlet geometry resulted in very long running times with little assurance that a steady state solution had been achieved. A total recoding of the entrance problem was therefore implemented resulting in a code (INTERNAL) that allows greater flexibility in handling the long tube geometry of SUMS as well as incorporates all of the changes in intermolecular collisions, numbers of species and rotational energy exchange developed for the code EXTERNAL.

#### A. Entrance Computer Code INTERNAL

The purpose of the program described in this section is to determine the fluid field inside a cavity which consists of a connected sequence of conic sections rotated about an axis of symmetry. The cavity is considered in two parts: a main cavity on whose interior surface the sensor will be

located, and an inlet tube whose orifice is presumed to be at the exterior surface of a spacecraft. A detailed code listing is given in Appendix C.

The input data for this program includes a molecular distribution function which is obtained from an external-flow run. Some of the other data refers to input parameters of the external-flow run. Thus it is seen that the pair of programs can be run as a set, computing the conditions inside a cavity which exist for a given set of conditions in the undisturbed free-stream flow. The programs have not yet been directly coupled, although they are written with this intent.

The program was constructed by turning an external-flow program inside out. In doing this, the basic molecule/body collision mechanism is preserved with only minor changes, while the molecule/molecule collision mechanism are not changed in any way.

A general description of the way in which the internal-flow program conducts the experiment is unnecessary since the description in the preceding section generally applies here. Any important differences will be described as they occur.

The numerical experiment takes place in a cylindrical block of space quite similar to that used in external flow. However, the inlet region must be cylindrical in shape while the cavity region can be defined by conic sections.

There is no subdivision of 1st level cells into smaller sized cells. Also, the number of azimuthal wedges is specified for a full  $360^\circ$  (since there is no plane of symmetry determined by an input flow direction) and is the

same for both regions. No weighting factors are used as the desired surface fluxes on the walls generally occur near the outer edges of the sample space.

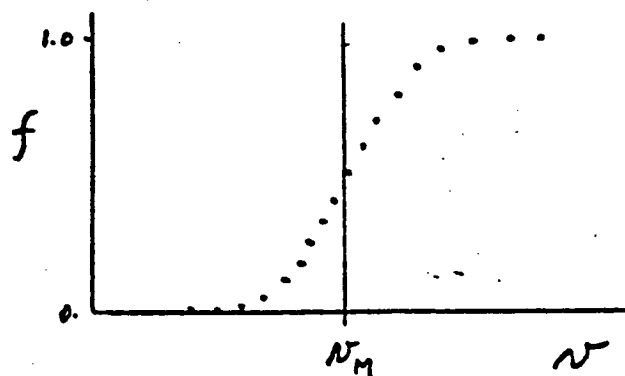
The orientation of the sample space is determined by the geometry of the particular body surface segment in external flow which is used as the orifice area of internal flow. The body surface normal of external flow becomes the x-axis (axis of symmetry) of internal flow; the direction given by the cross product of the x-axis of external flow and the body surface normal becomes the z-axis of internal flow; and the direction given by the cross product of the above cross product and the body surface normal becomes the y-axis.

The zero radial plane is the positive x-y plane, and azimuth angles range from  $0^\circ$  to  $360^\circ$ , with the positive z-axis at  $90^\circ$ .

The sample space is initially populated with gas molecules in a manner similar to that used in external flow. In this program, however, the gas molecules are initially in thermal with the walls locally. The selection of the density profile initially in the tube indeed poses a problem and is very critical in minimizing running time. The code therefore allows an arbitrary initial distribution specified by the user. The choices and procedures for selecting them are described in the next section. The key objective is to provide a distribution that is consistent with the input flux and will not have to be severely altered in magnitude to arrive at the steady state.



The molecular input distributions are generated from data produced by an external-flow run. This data consists of velocity samples of molecules impinging on the external body surface segment which is considered to be the orifice area of the inlet region. The form of the velocity samples is a series of horizontal S-shaped curves, one for each molecule type, in each of three directions, for both UNCOLLIDED (free-stream) and COLLIDED molecules. Each s-curve gives the fraction of molecules (impinging on the body surface segment) with a velocity  $\leq$  the given velocity. A typical curve for UNCOLLIDED molecules is given [ $v_M$  is the velocity corresponding to the center velocity of the distribution as computed in the external flow program]:



The information from these curves is used directly to generate the molecular distribution of the molecules entering

the inlet.

The population (pressure) inside the cavity region is selected initially as an input. The computer simulation proceeds until a steady profile inside the tube entrance region is generated. Since this does not necessarily require zero net flux, a series of runs with different cavity pressures is generated and a cross plot of flux versus pressure is the output for a single external flow input condition. The equilibrium solution if desired can be obtained from the zero flux point.

#### B. Scope of "Entrance Problem" for SUMS experiment

Within the free molecular regime the "entrance problem" has been examined by Hughes and deLeeuw<sup>(2)</sup>. In order to indicate the potential magnitude of the problem, Figure 20 shows the variation of both the surface pressure ( $p_s$ ) and the chamber pressure ( $p_c$ ) versus angle of attack under free molecular conditions at infinite speed ratio and a very long tube (conditions approached at high altitude during Space Shuttle re-entry). For the SUMS location, the local angle of attack  $\beta$  is approximately  $28^\circ$  giving a possible difference of a factor of 10 between the surface pressure ( $p_s$ ) and the internal chamber ( $p_c$ ) that directly affects the mass spectrometer reading. While this theoretical result is expected to be accurate at 140 Km and above, at lower altitudes the local flux distribution will have both a directed and a rather diffuse component (see Figure 18), and

also the effect of internal collisions within the tube will begin to play a role (see Figure 1).

Because the tube length to diameter ratio is very large (37) the time constant and therefore the computing time to reach equilibrium is very long. For this reason no attempt is made to simulate the problem all the way to the condition where the chamber pressure behind the tube is at the correct value to nearly balance the net flux. Instead a series of runs with different assumed "cavity" pressures are performed and the zero net flux (or a given small value for the dynamic condition) can be selected to interpolate the correct "cavity" pressure. This procedure if it proves generally successful can of course be automated within the code.

#### C. Preliminary Results of INTERNAL Code

Preliminary results for the SUMS entrance tube geometry are presented as a couple points on Figure 21 giving to the ratio  $p_c/p_s$  versus altitude.

Also shown are some theoretical curves for full continuum (Continuum), for continuum external flow but free molecular flow through the tube (No slip,  $Kn_s \rightarrow \infty$ ), and for fully free molecular flow (Free Molecule). In addition, some points using a best fit Maxwellian to the external flow flux and free molecular internal flow from reference 2, are shown at 95 Km and 115 Km. These results are preliminary, but they do indicate the magnitude of the correction and the fact that no simple use of the currently available results can cover

the entire range of altitudes of interest.

#### IV. DISCUSSION

If we couple information from INTERNAL (Figure 21) to the results of the external flow computations (summarized as a plot of  $p_s/1/2\rho U^2$  versus altitude in Figure 22), we can produce a preliminary estimate of the data reduction curve that could be coupled to the calibration of the mass spectrometer to deduce the  $q = 1/2\rho U^2$  during Space Shuttle re-entry. Figure 23 is a cross-plot of the  $q/p_c$  versus  $p_c$  obtained from Figures 21 and 22 with the 87 and 105 Km results only estimated on the basis of interpolation of Figure 21. The establishment of such a data reduction curve for the nominal re-entry conditions, together with associated error bars as well sensitivities to wall temperature and angle of attack is necessary for the proper interpretation of data to be obtained by the SUMS instrument.

The data measured by the mass spectrometer in the SUMS experiment essentially provides collector currents of charged species of different masses. Calibration of the instrument can relate these to the overall pressure and composition at the entrance to the instrument being calibrated. Since the actual environment of the flux of molecules to be encountered under flight conditions cannot be simulated the calibration is performed with the incident flux essentially in equilibrium with the instrument outer walls (room temperature and no flow). The data reduction procedure must therefore relate the effective environment in the ground test simulation to the desired dynamic pressure  $q$  under the

flowing non-equilibrium condition and through the calibration to the instrument measurements.

As discussed above, capability now exists for calculating the surface flux distribution of molecules entering an opening at the SUMS location. The computational procedures for connecting that information to the pressure immediately behind the entrance tube has also been developed, although not fully exercised over the entire range of parameters applicable to the SUMS experiment. That pressure, in turn, can be directly related to the ground test environment used to produce the calibration curves for the mass spectrometer.

The primary objective of future work must be to provide a data reduction procedure that relates the spectrometer reading to the free stream dynamic pressure ( $q = 1/2 \rho U^2$ ). With currently available procedures a relation between the pressure  $p_c$  and  $q$  such as the preliminary one shown on Figure 23 can be obtained using the best available information on the flight parameters such as velocity, angle of attack, tile temperature, surface conditions, molecular collision parameters, etc. This relation must then be combined with the calibration curve where the instrument readings are related either to  $p_c$  directly, or to a calibration pressure which can be related to  $p_c$  by conventional means. A single plot of  $q$  versus total measured collector current can thus be obtained from the combination of these results.

The relation between the dynamic pressure  $q$  and the calibration pressure  $p_c$  depends on many parameters of the

problem. Some of these, such as flight velocity, angle of attack, and tile surface temperature are expected to vary only slightly from their nominal values. Since the actual measured values of these quantities will be available on each individual flight, corrections to the data reduction relation should be evaluated in the form of sensitivity coefficients for small changes from the nominal. Studies to determine the effects of these parameters are necessary to establish the significant sensitivity coefficients that must be incorporated into the data reduction scheme.

The parameters such as surface accommodation, surface recombination, and free stream composition can also affect the results. In addition, modeling simplifications of both the geometry of the problem and the molecular collision phenomena can alter the quantitative value of the relation between  $q$  and  $p_c$ . Because of the unavailability of any in-flight measurements that could lead to an evaluation of these parameters, bounds on the uncertainties they produce should be studied. Sensitivity of the data reduction relation to the most significant of these can then produce bounds on uncertainties on the dynamic pressure  $q$ , due to reasonable variations. The final goal of future work is an algorithm for the evaluation of the dynamic pressure  $q$ , from the calibration pressure  $p_c$ , together with error bounds, due to the uncertainties associated with the external flow and the entrance problems.

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FIGURE 1. RELEVANT KNUDSEN NUMBERS VERSUS ALTITUDE

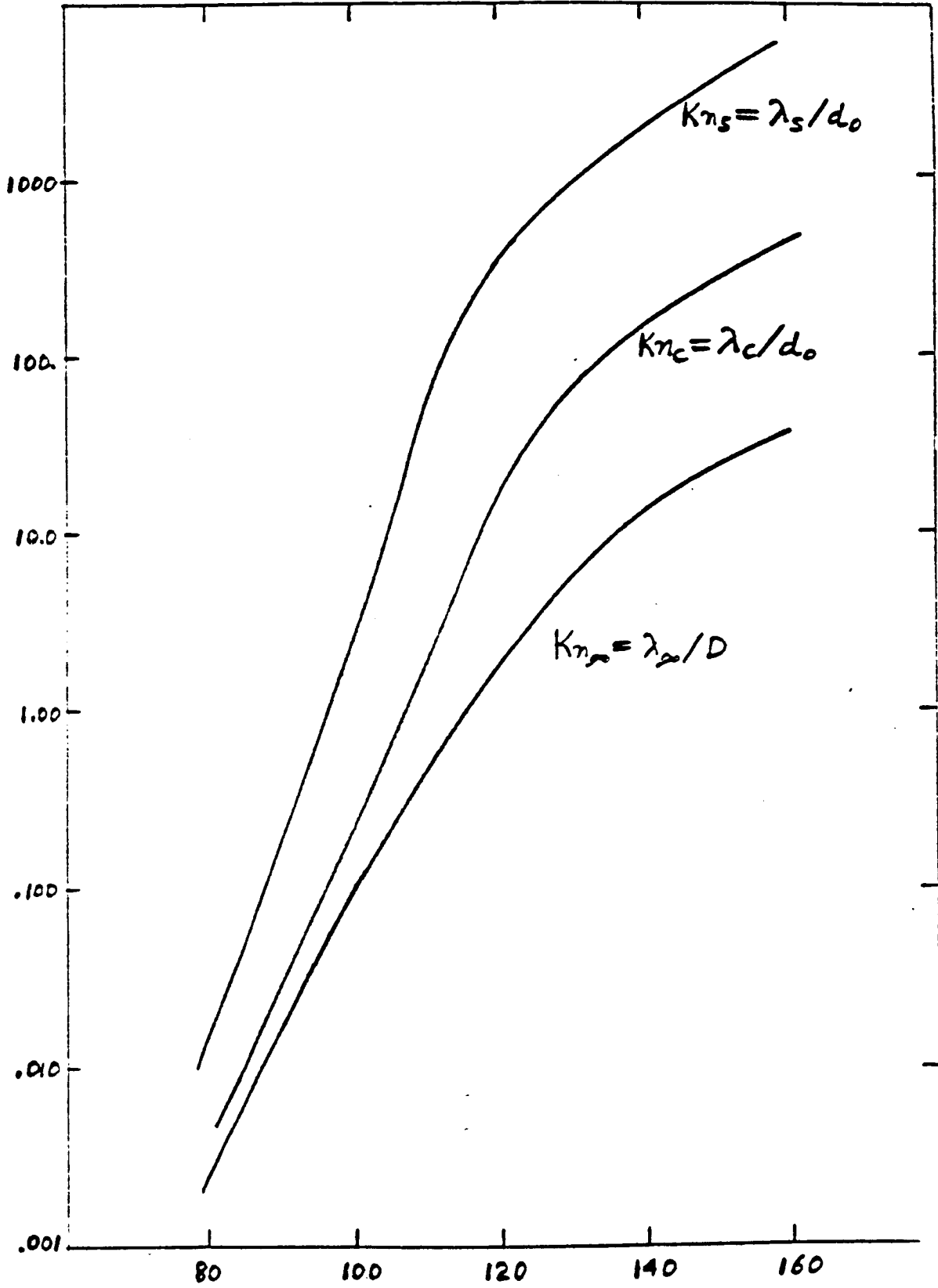


FIGURE 2. SYSTEM FLOW-CHART

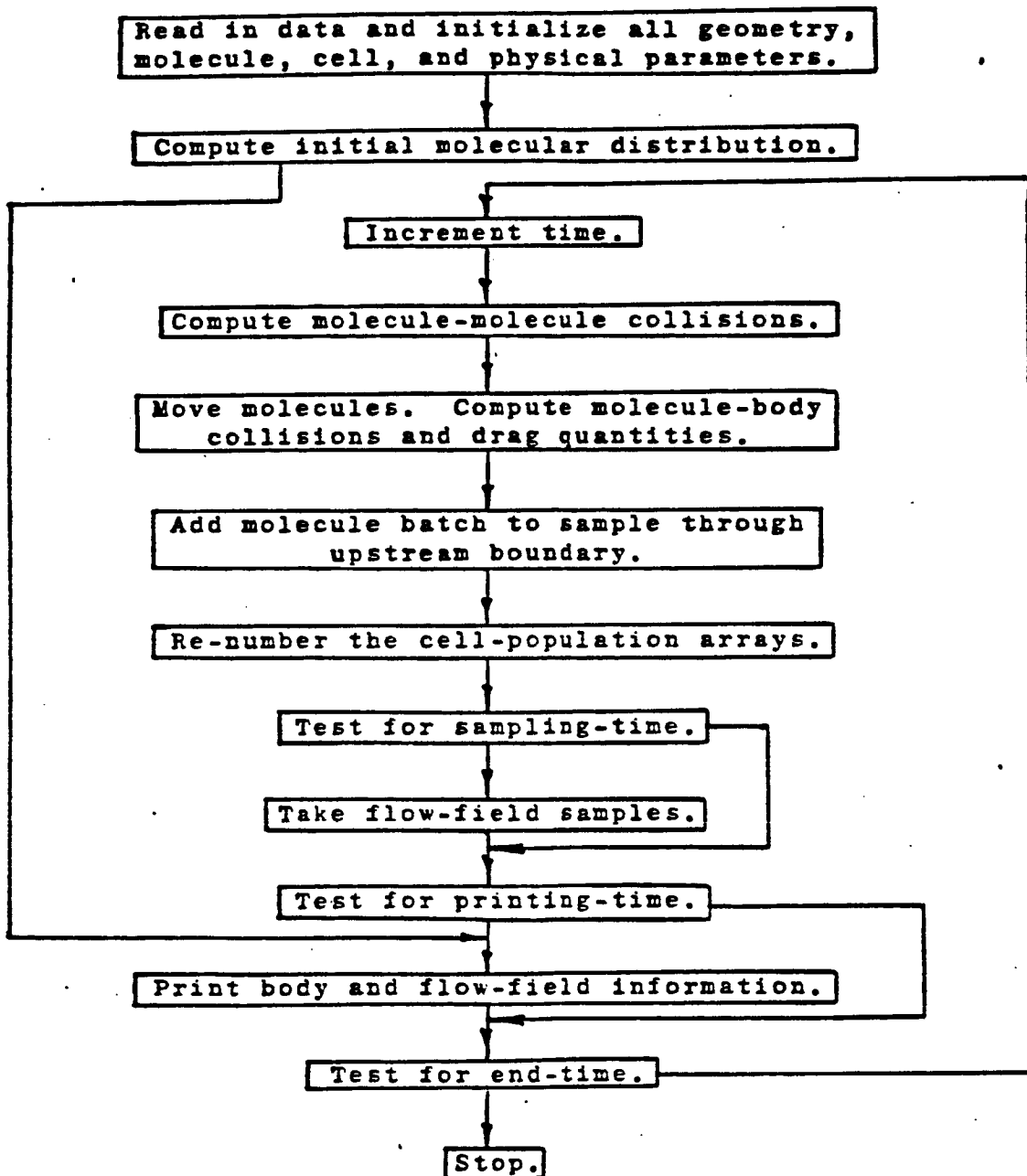
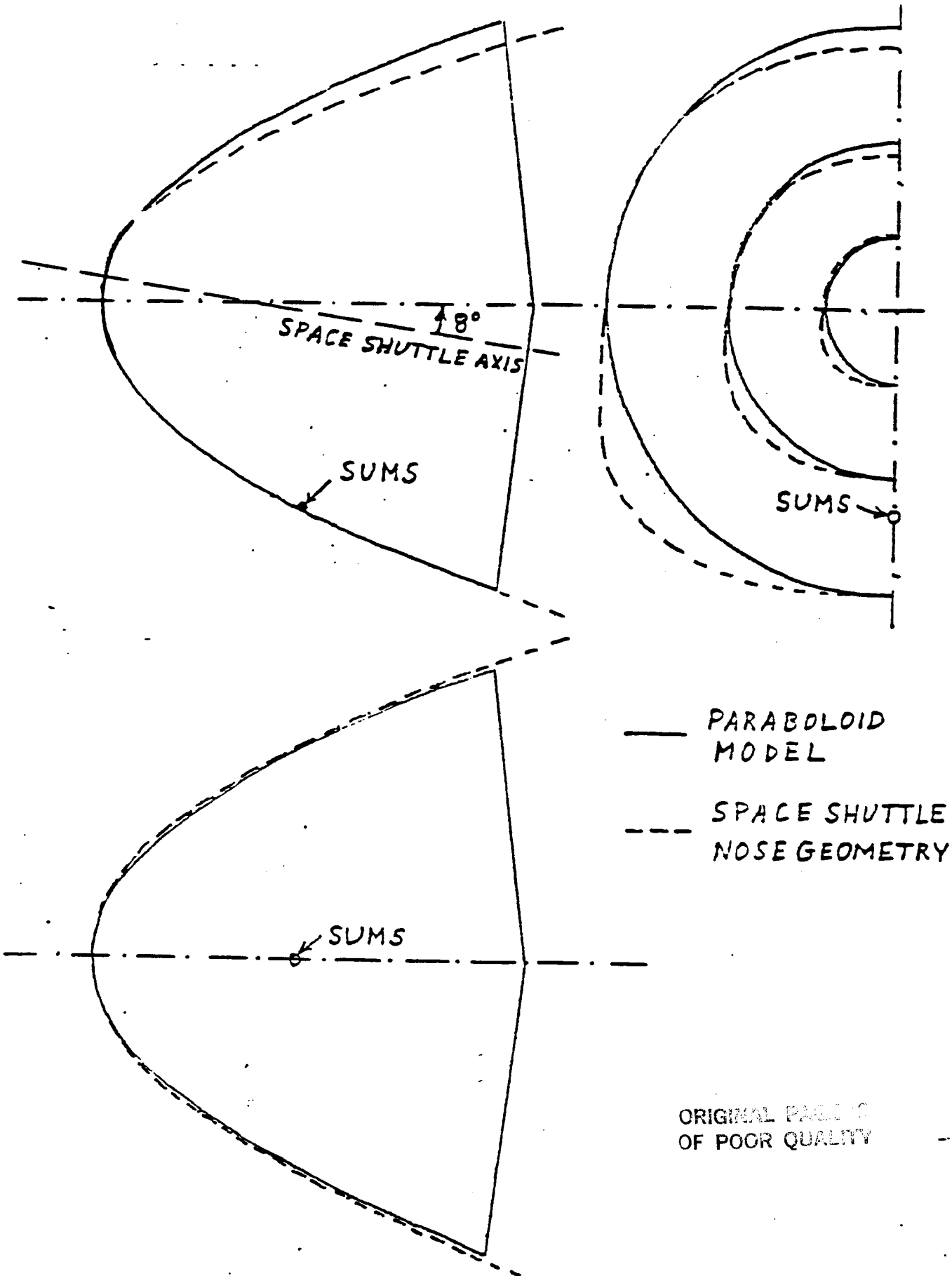


FIGURE 3. MODELING OF SPACE SHUTTLE NOSE (PARABOLA)



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FIGURE 4. TYPICAL CELL CONFIGURATION FOR HYPERBOLA

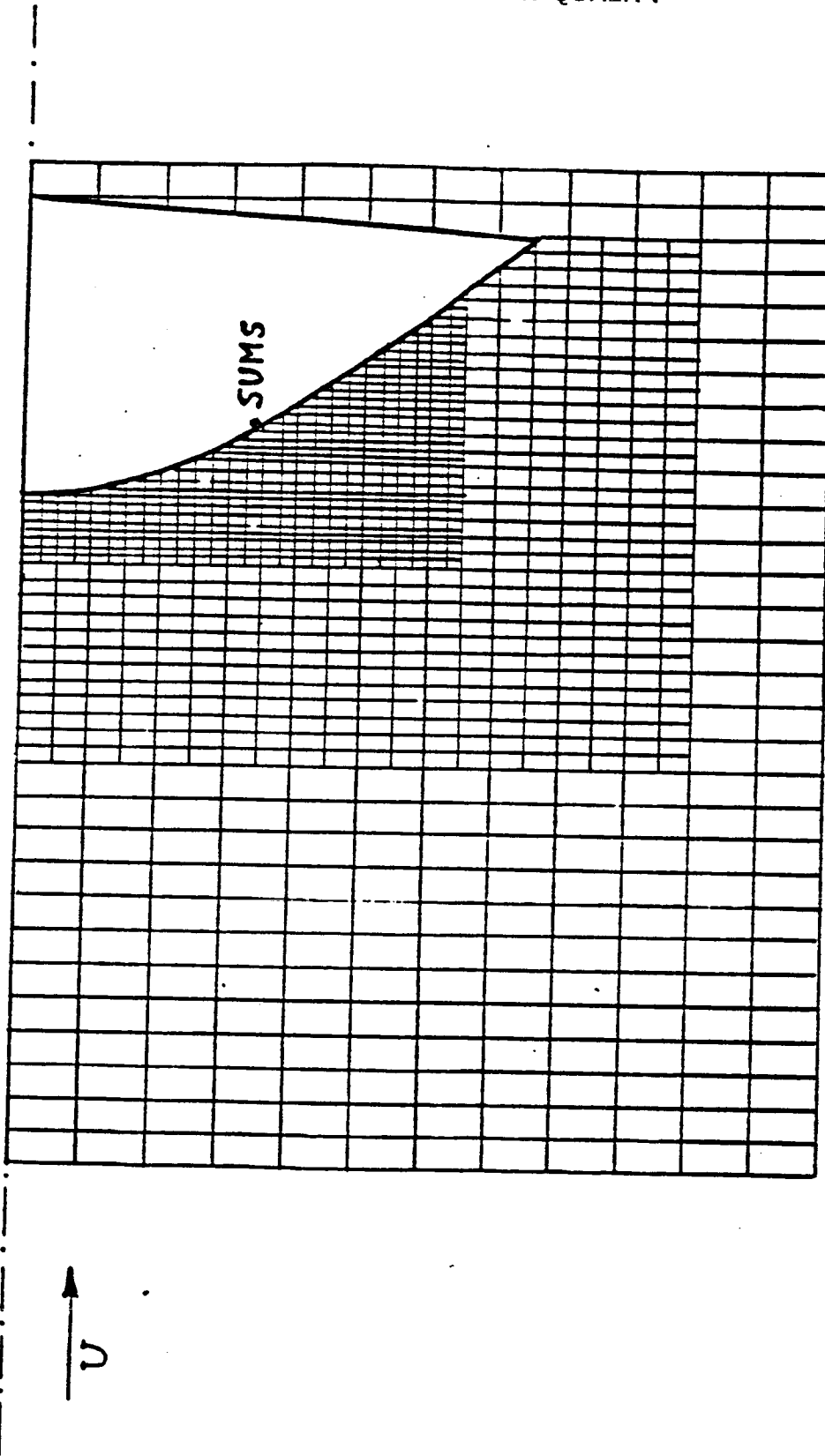


FIGURE 5. MODELLING OF VISCOSITY-TEMPERATURE DEPENDENCE

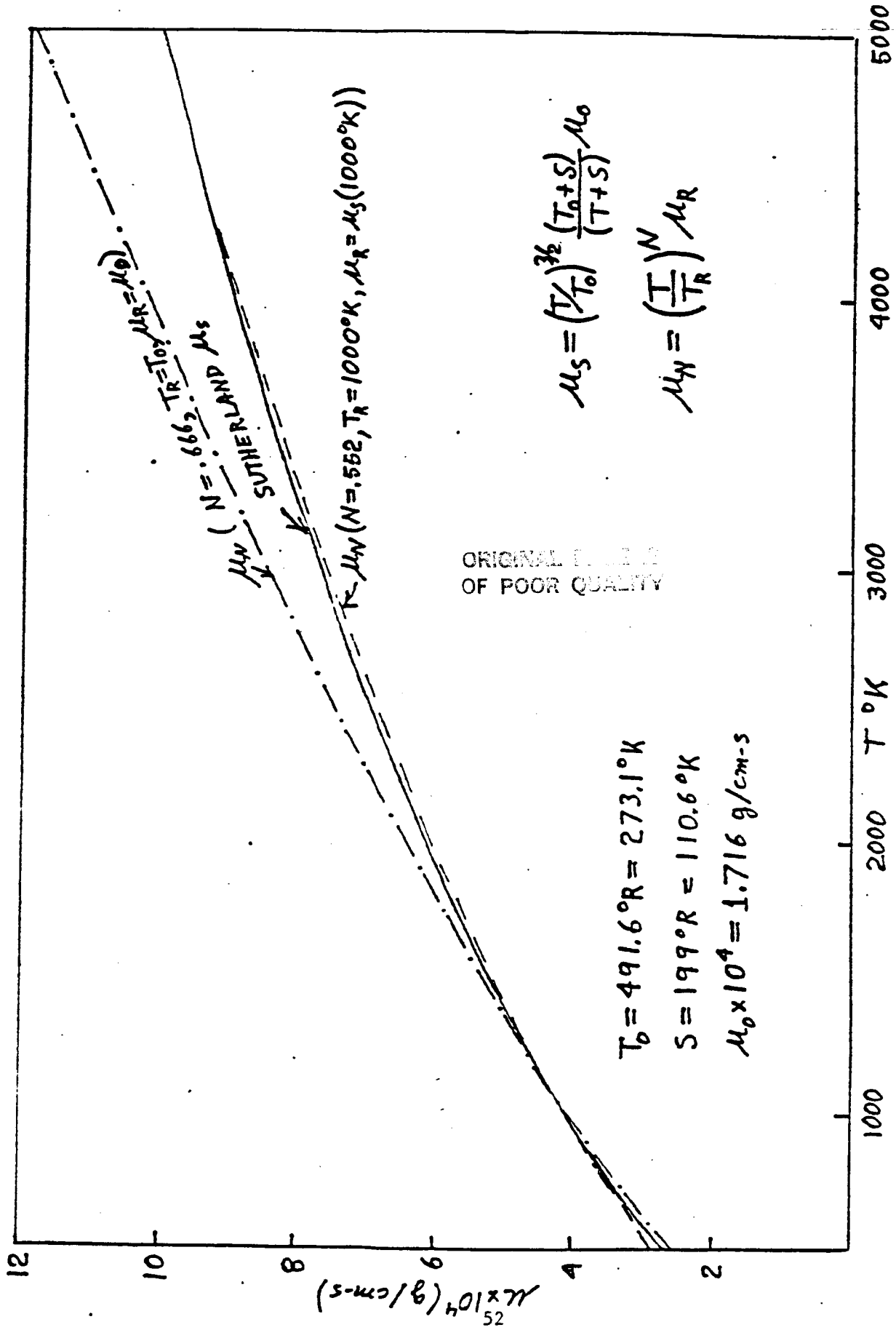
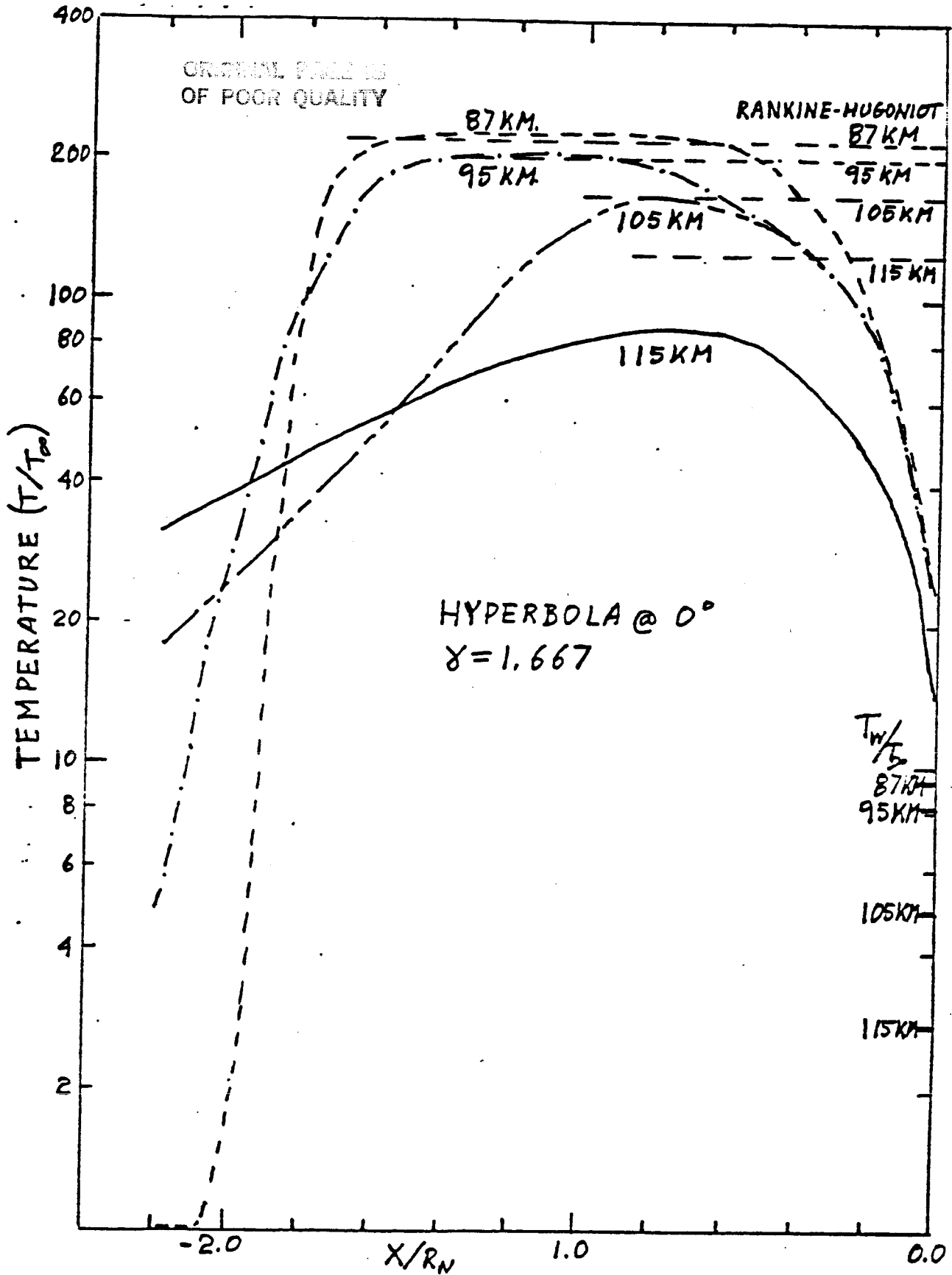


FIGURE 6. TEMPERATURE DISTRIBUTIONS (EFFECT OF ALTITUDE)



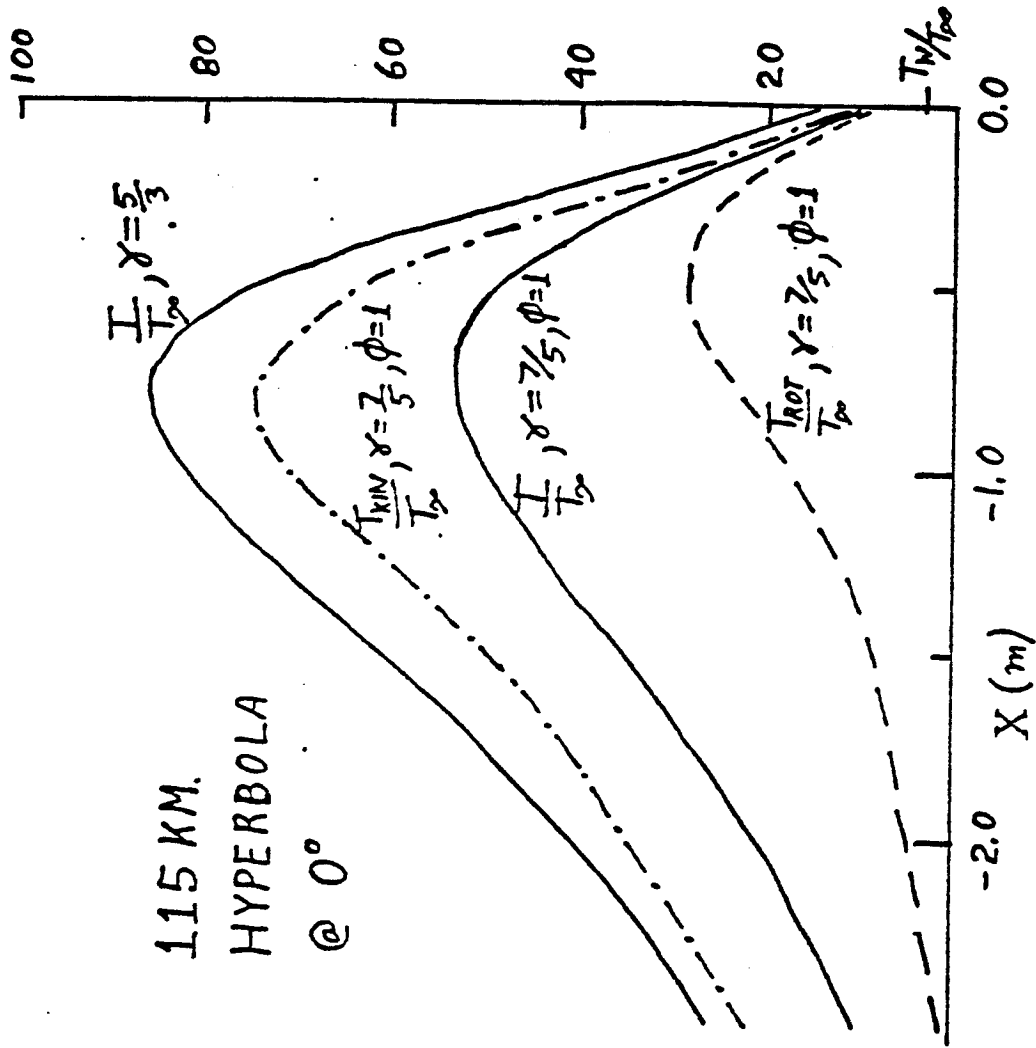
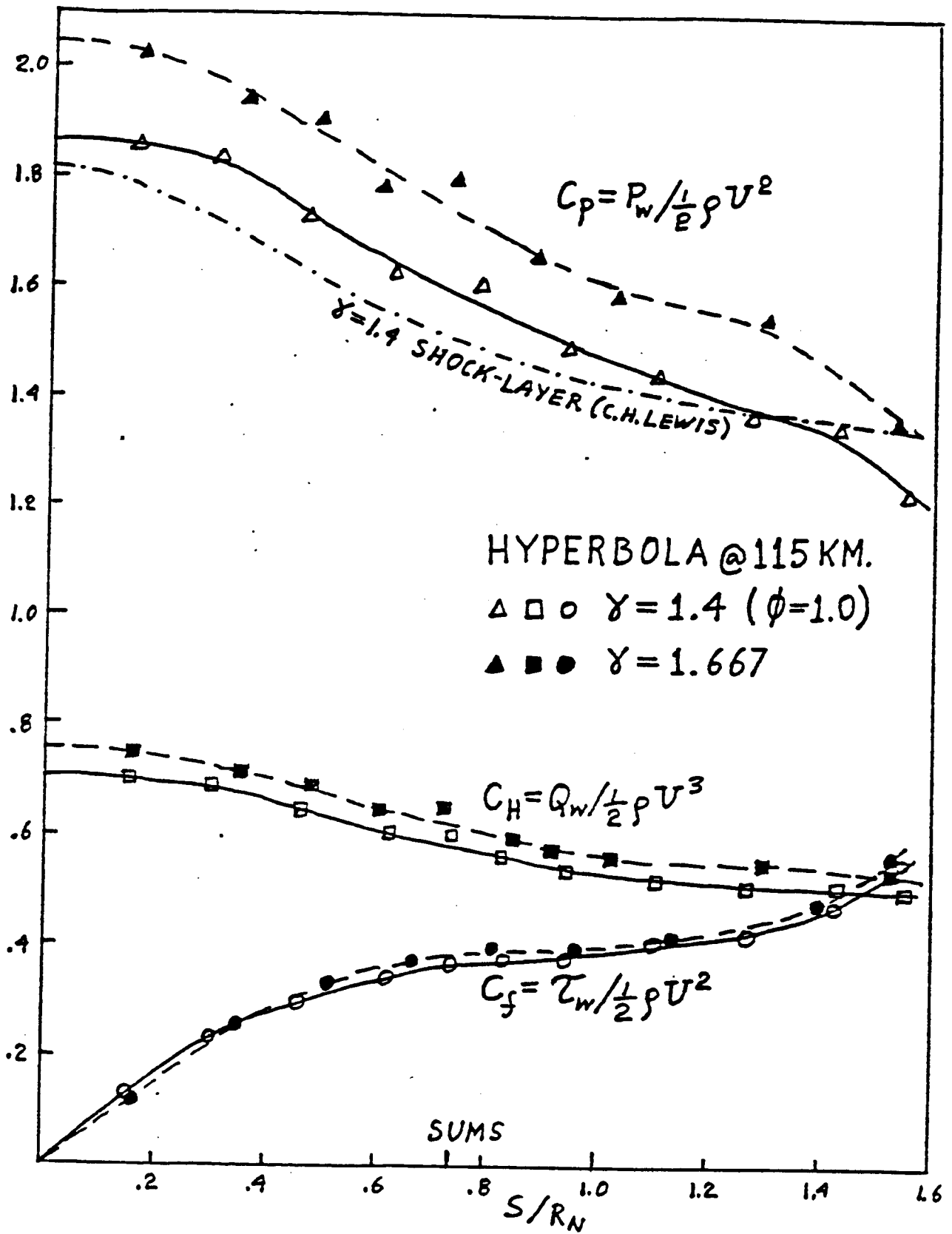


FIGURE 7. TEMPERATURE DISTRIBUTION ALONG STAGNATION STREAMLINE (EFFECT OF  $\gamma$ )



FIGURE 8. SURFACE FLUXES (EFFECT OF  $\gamma$ )



Normalized Distance from Stagnation Point

FIGURE 9. SONIC LINES (EFFECT OF BODY SIZE)

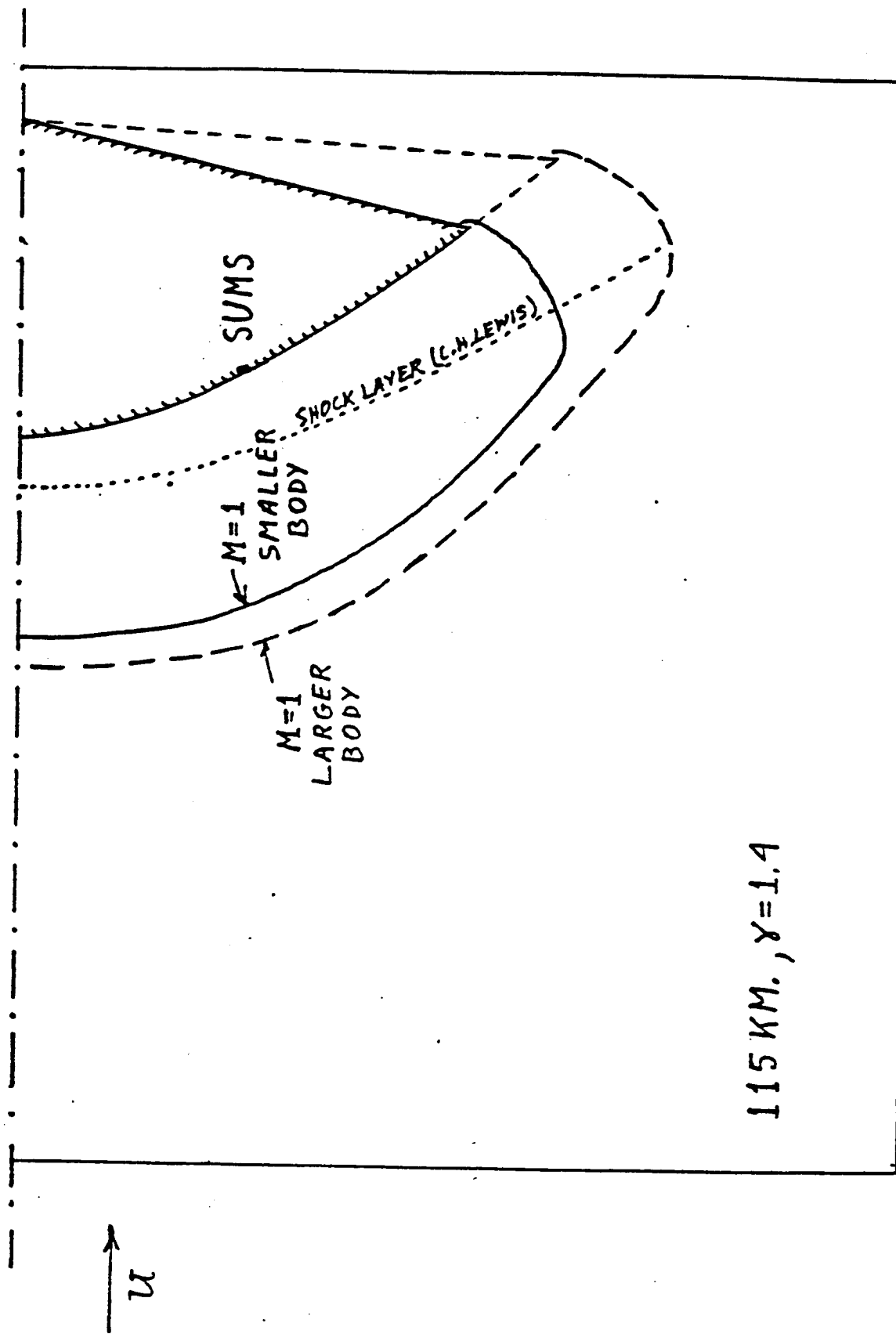


FIGURE 10. SONIC LINES (EFFECT OF GEOMETRIC MODEL)

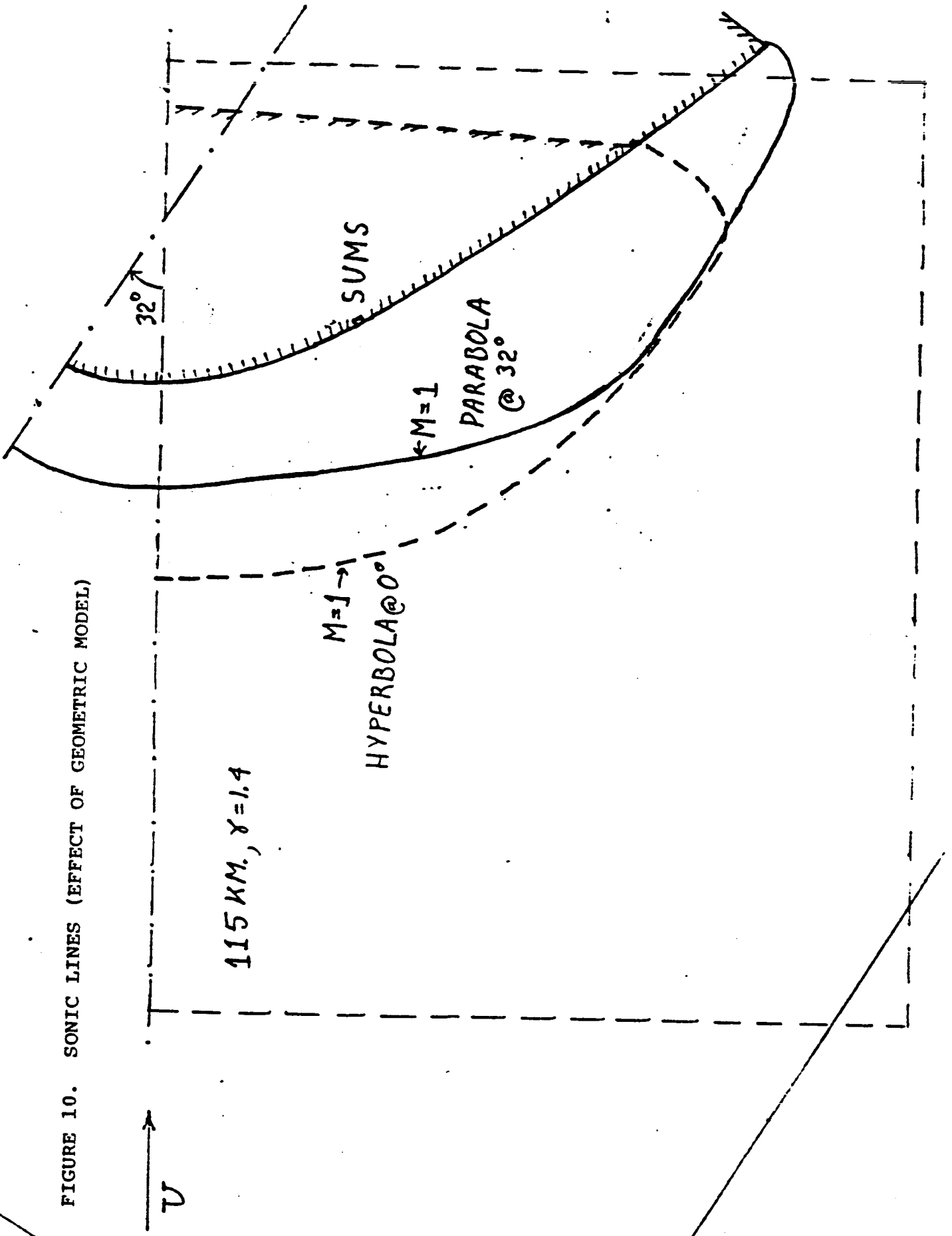
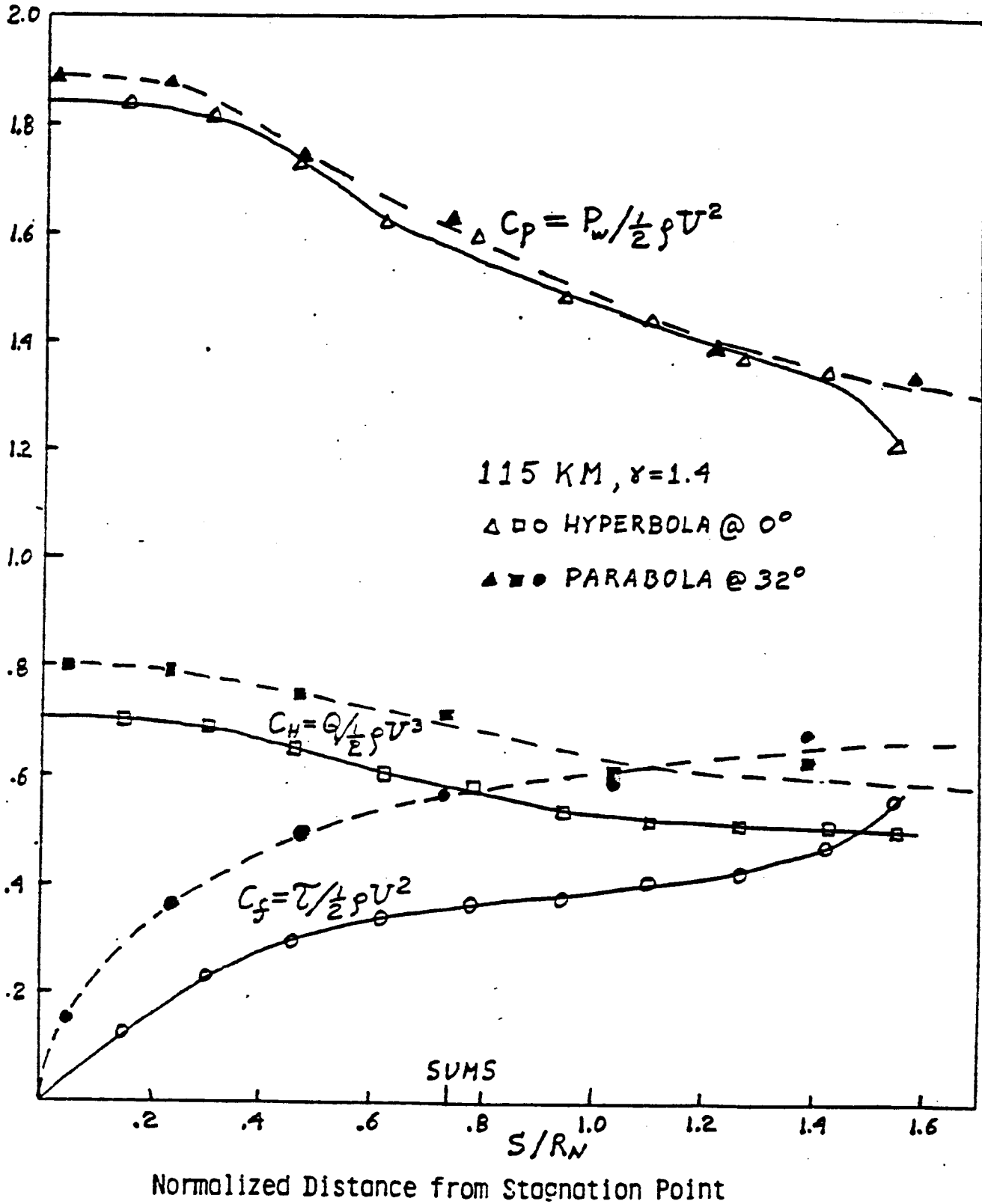


FIGURE 11. SURFACE FLUXES (EFFECT OF GEOMETRIC MODEL)



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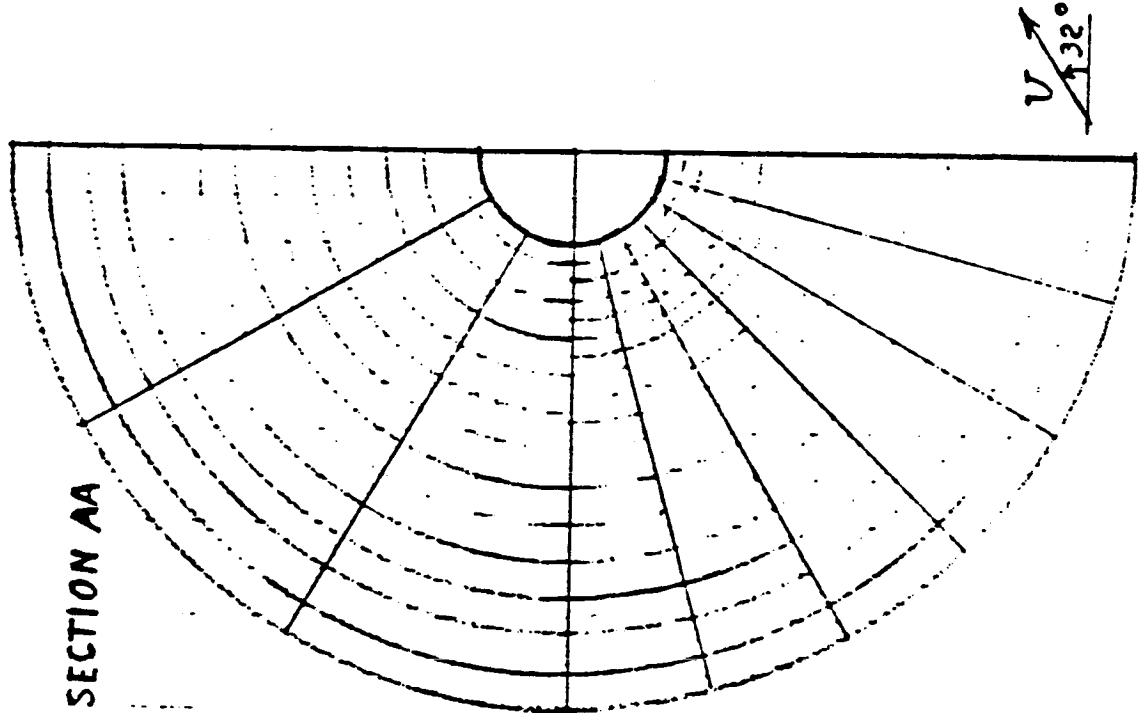
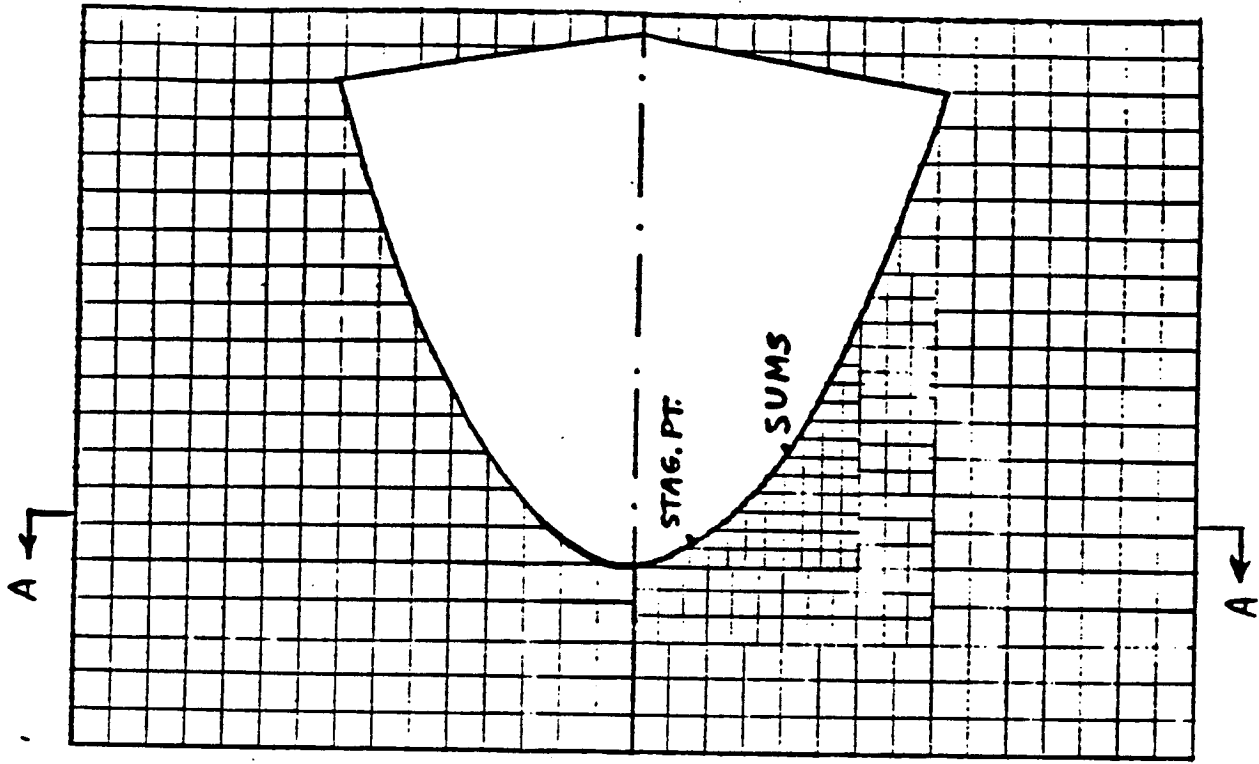


FIGURE 12. TYPICAL CELL CONFIGURATION (95 Km RUN)

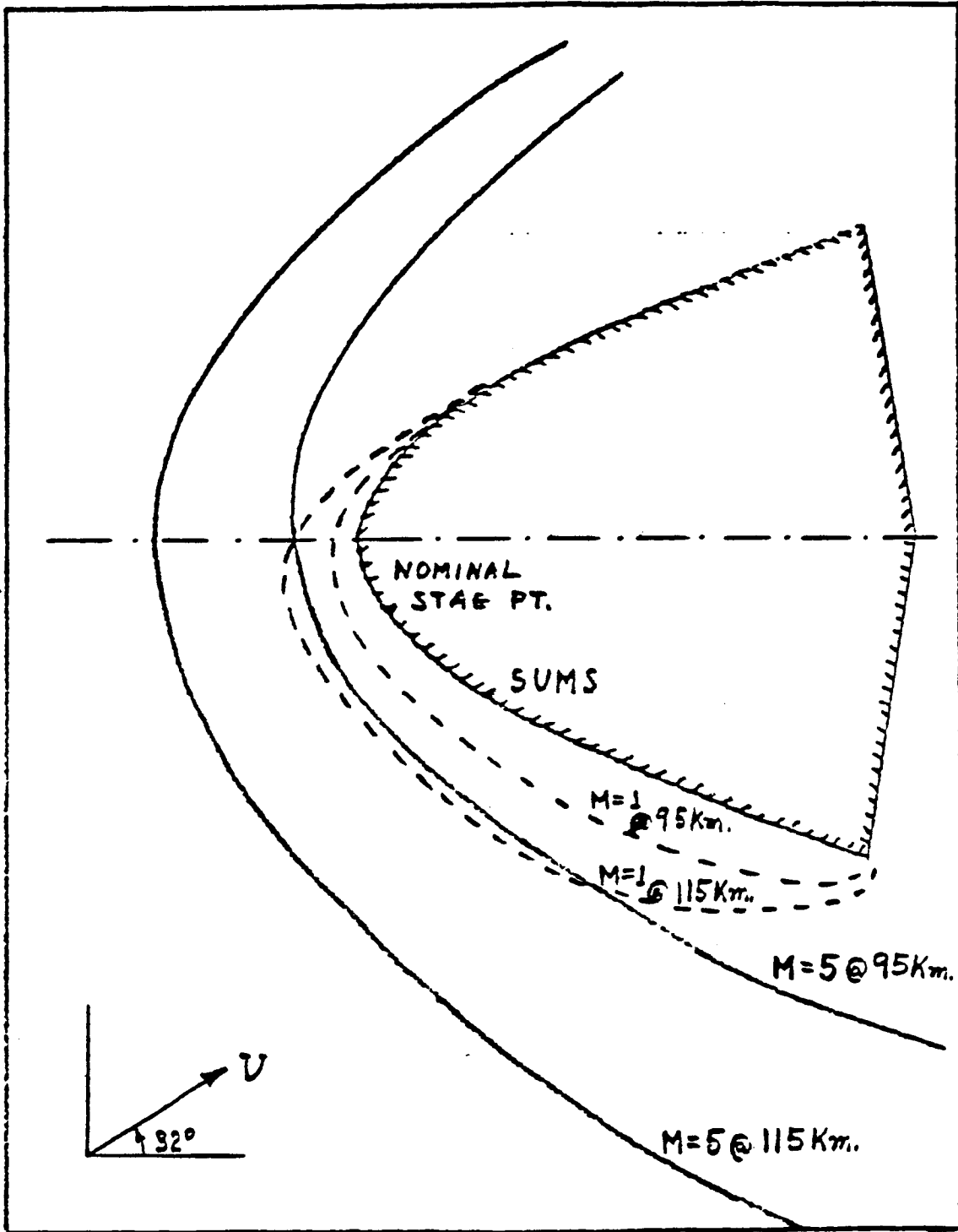


FIGURE 13. MACH NUMBER CONTOURS IN SYMMETRY PLANE AT 95 Km and 115 Km

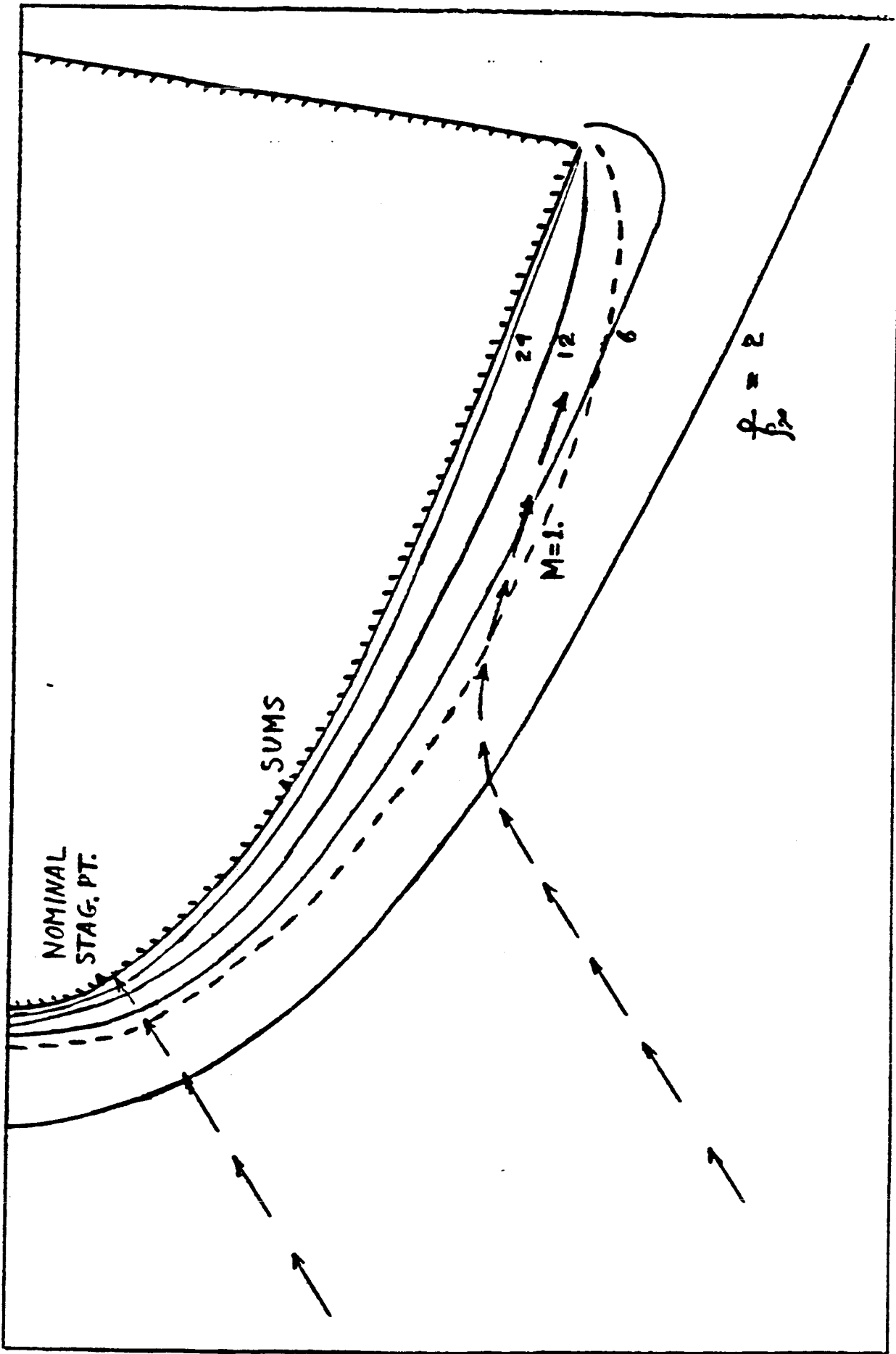
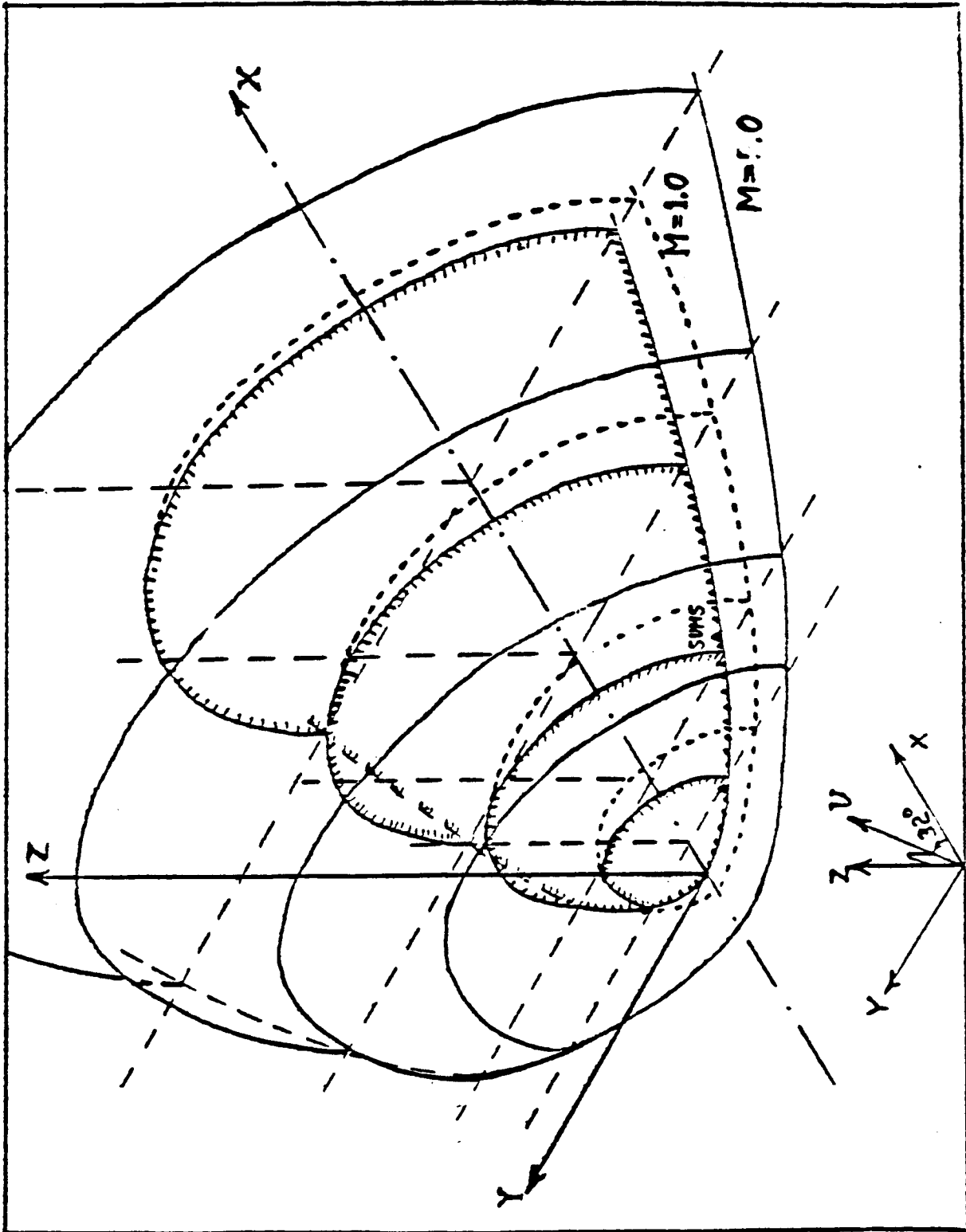


FIGURE 14. DENSITY CONTOURS ON WINDWARD SIDE AT 95 KM ALTITUDE

FIGURE 4.3. APPROXIMATE ORIGINAL SKETCH OF MACH NUMBER CONTOURS  
AT 95 Km ALTITUDE





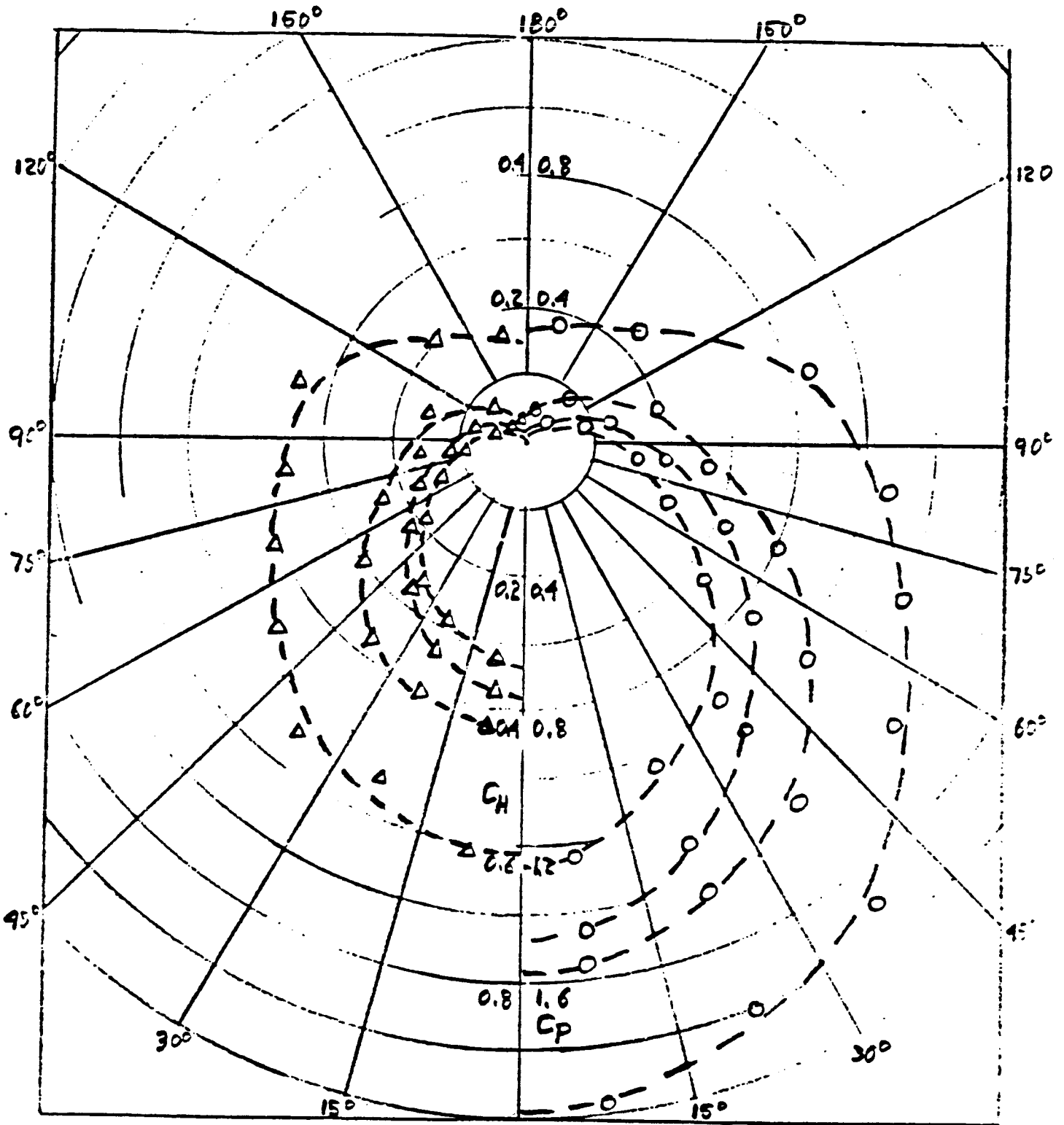


FIGURE 16. SURFACE PROPERTIES; PRESSURE COEFFICIENT  $C_p = p_s/1/2\rho U^2$  AND HEAT TRANSFER COEFFICIENT  $C_H = Q/1/2\rho U^3$  IN FOUR CROSS-PLANES AT 95 Km ALTITUDE

FIGURE 17. SURFACE PROPERTIES; PRESSURE  $C_p$  AND HEAT TRANSFER COEFFICIENT  $C_H$  VERSUS ANGLE IN CROSS-PLANE  $x = 75$  cm (NEAR THE SUMS ENTRANCE) AT 95 Km ALTITUDE

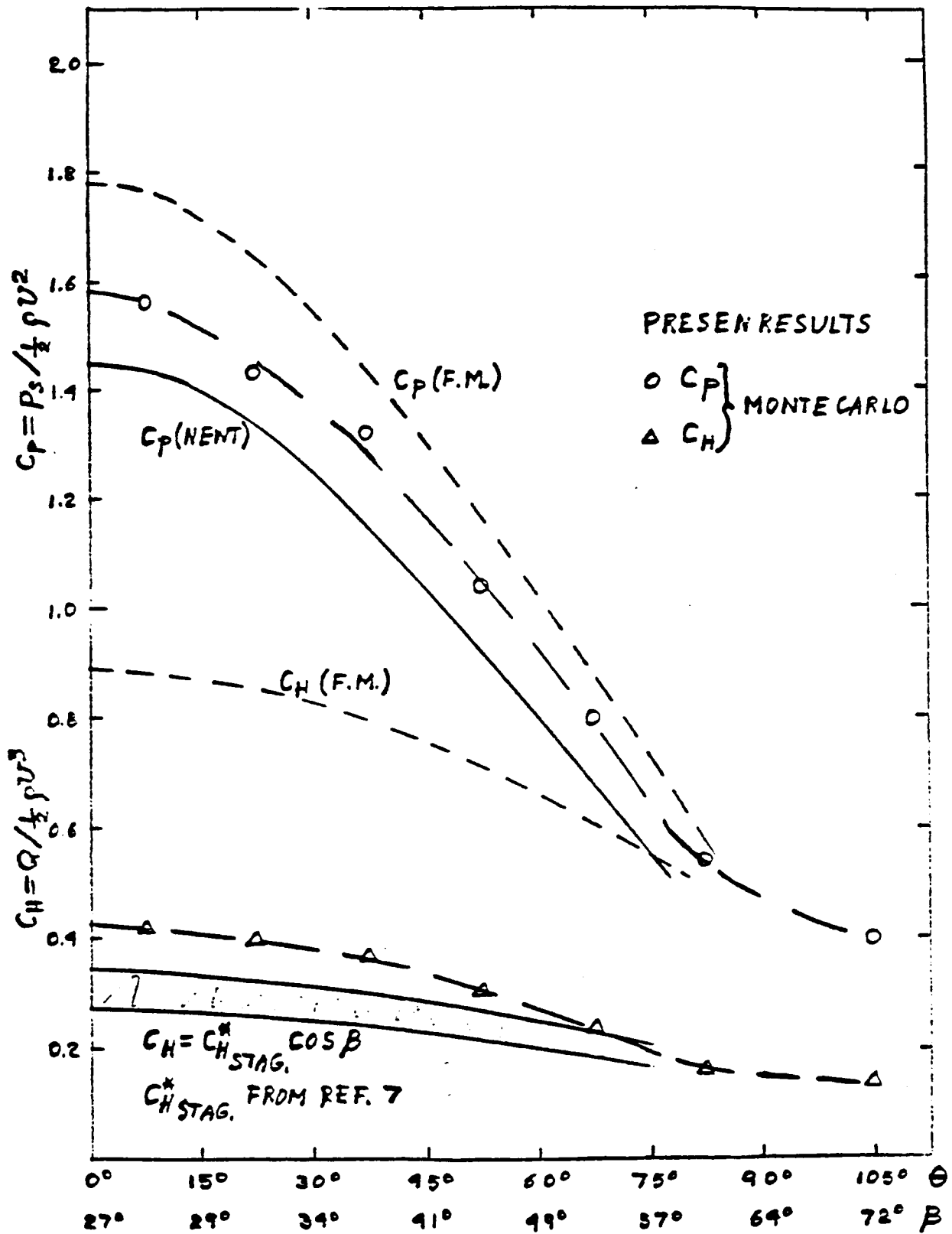


FIGURE 18. FLUX DISTRIBUTION FUNCTION AT SUMS INLET

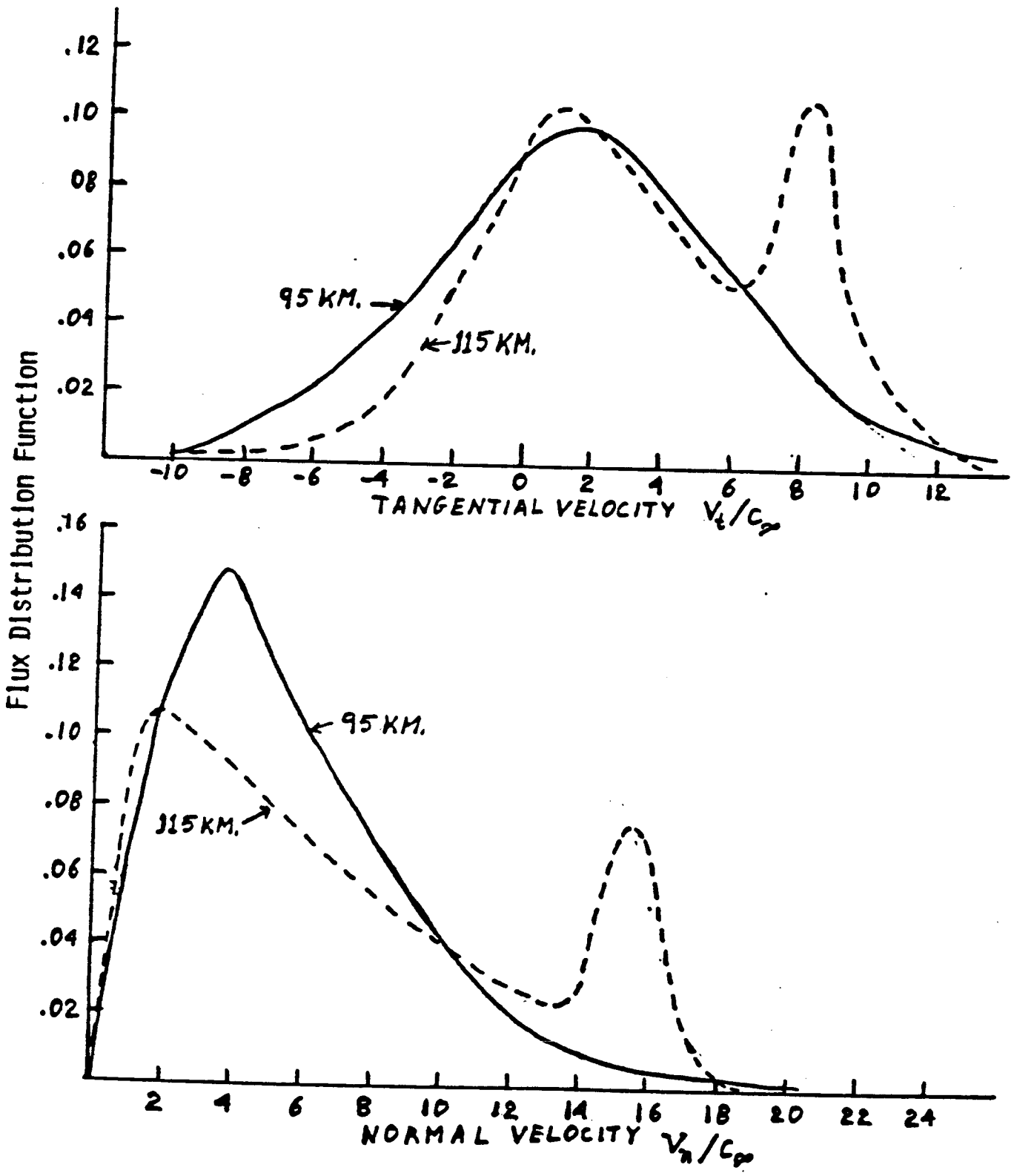


FIGURE 19. SUMS INLET GEOMETRY

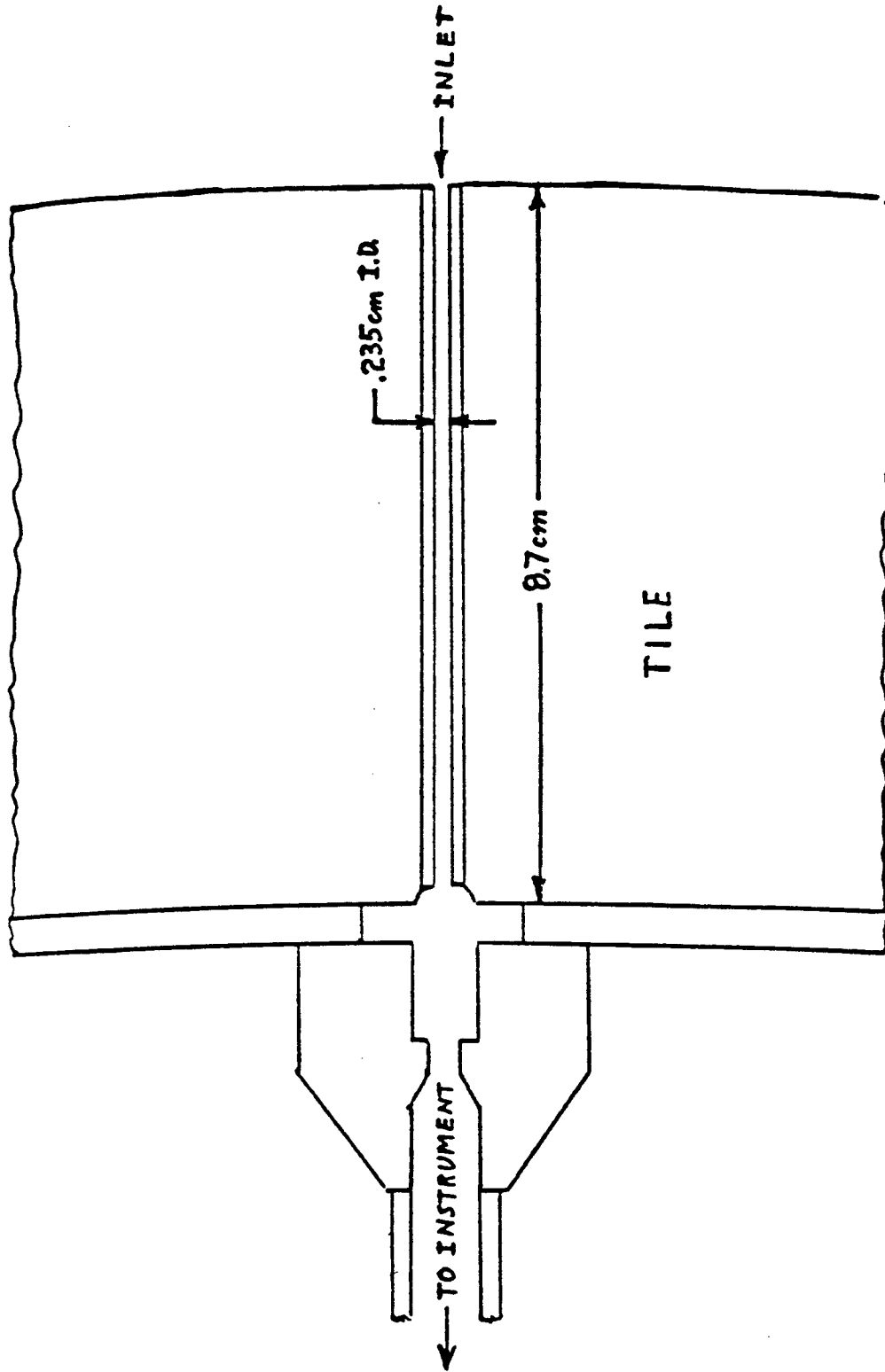


FIGURE 20. FREE MOLECULAR PROBE RESPONSE ( $S \rightarrow \infty$ ,  $d/L \rightarrow 0$ )

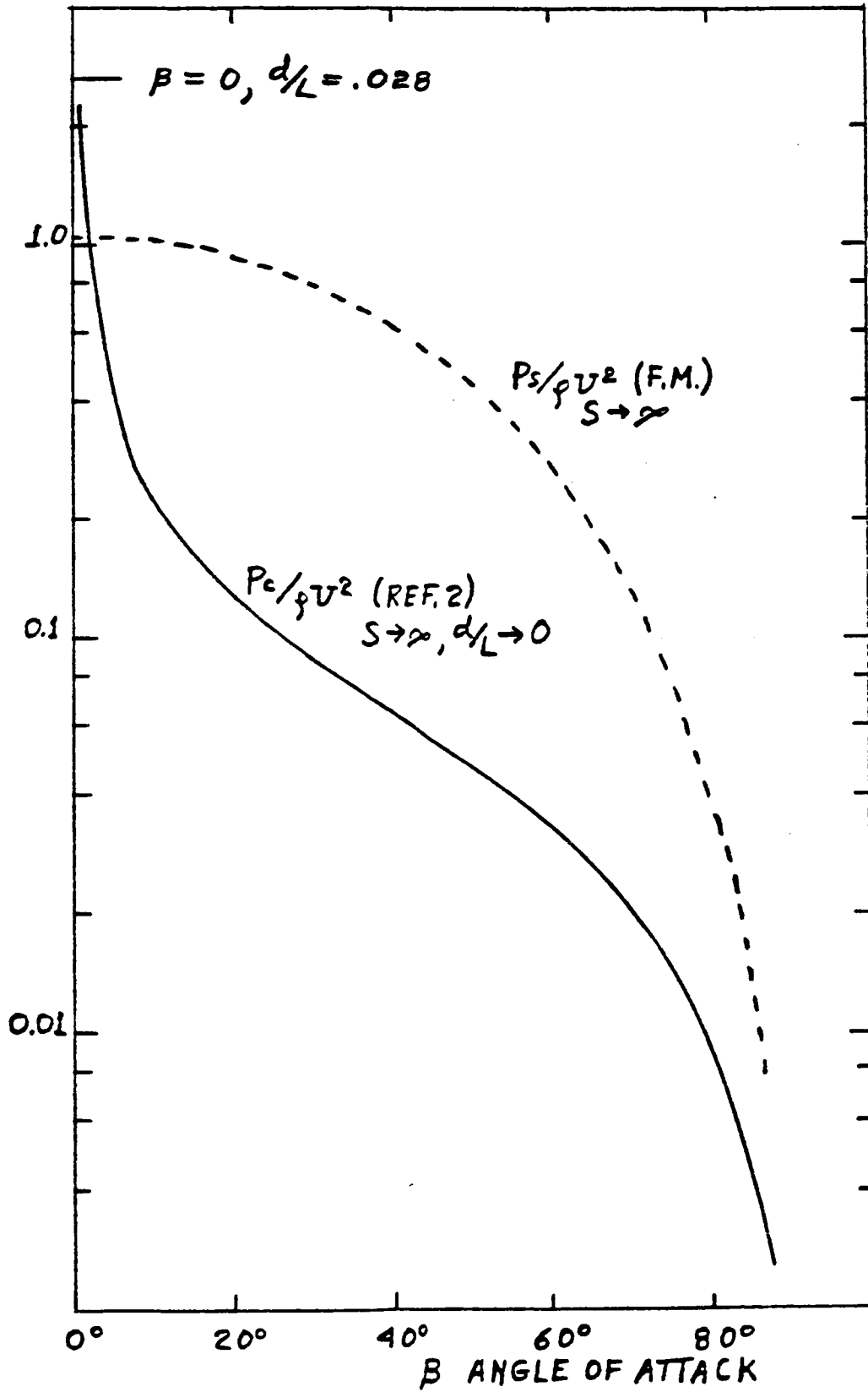


FIGURE 21. PRESSURE CORRECTION FACTOR ( $P_c/P_s$ )

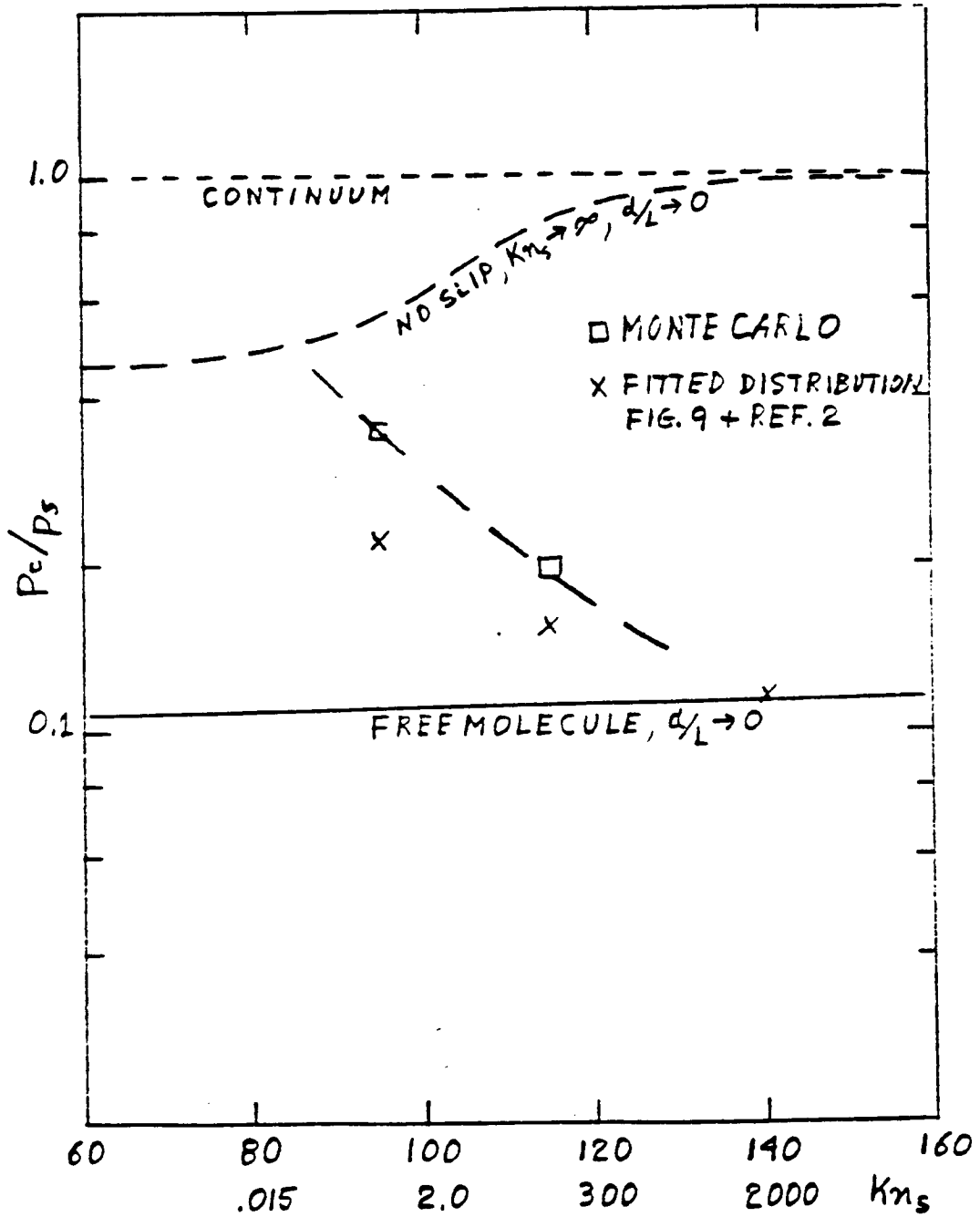


FIGURE 22. SURFACE PRESSURE AT SUMS INLET  
VERSUS ALTITUDE

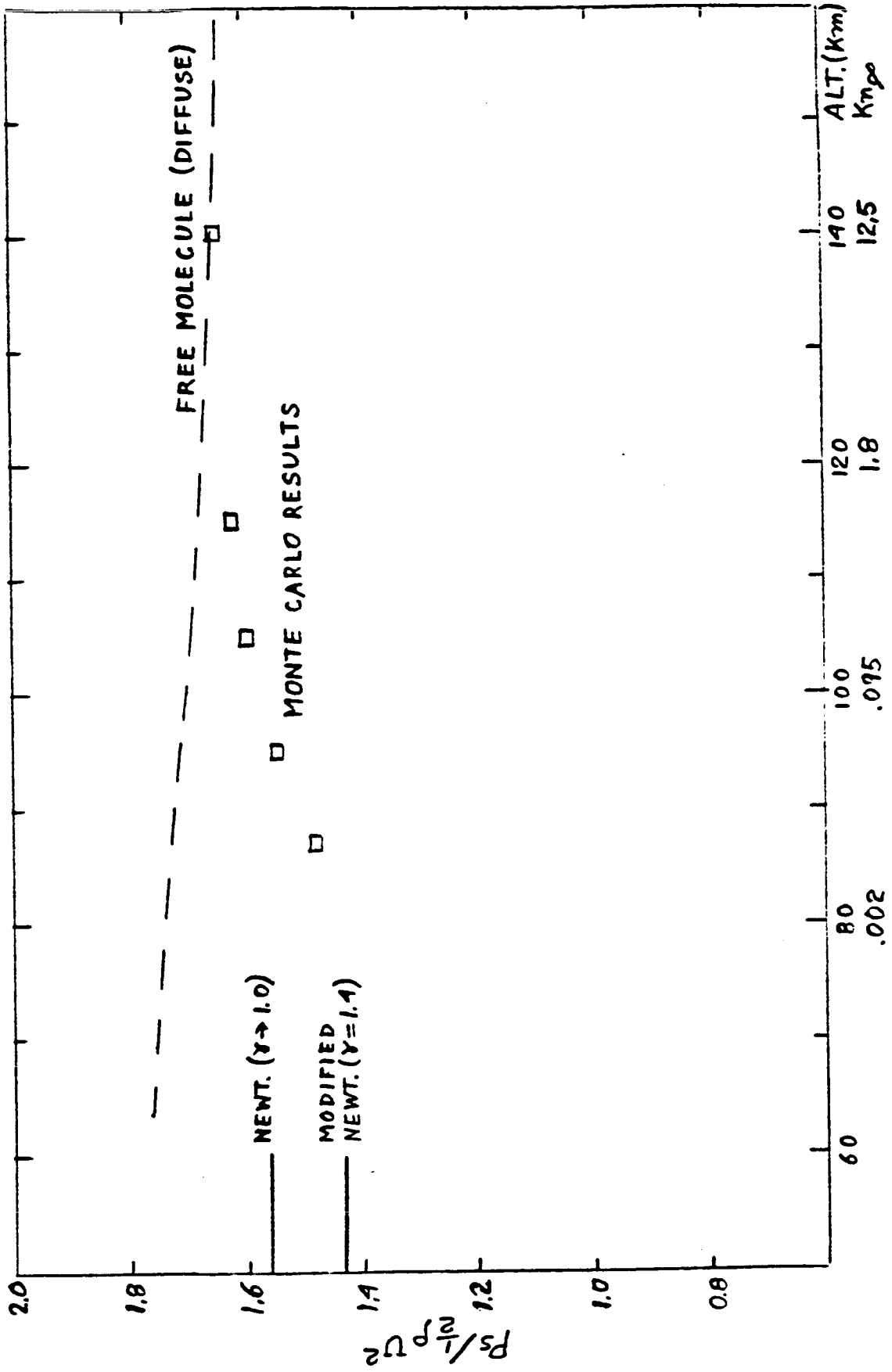
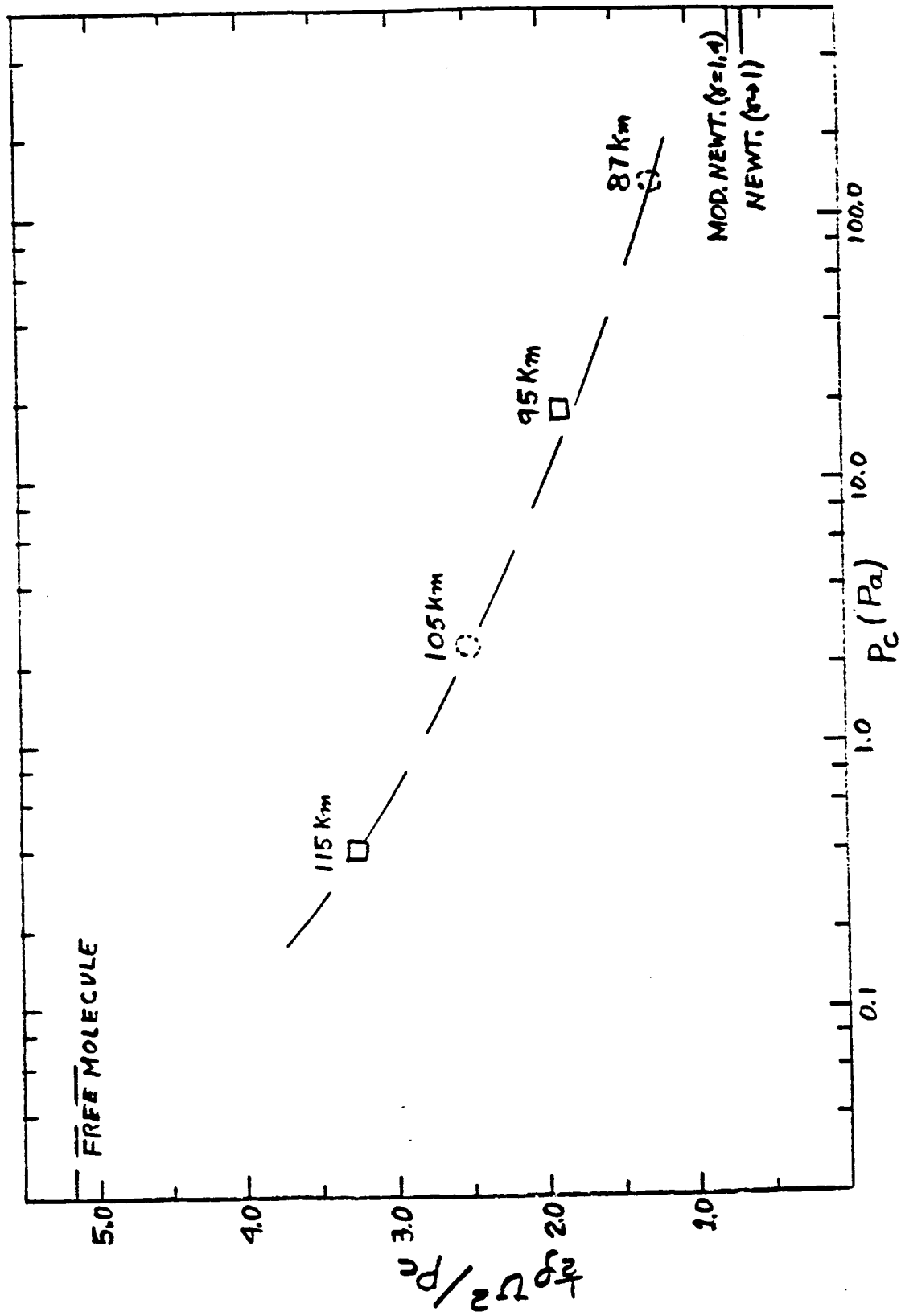


FIGURE 23. PRELIMINARY ESTIMATE OF DATA REDUCTION  
RELATION  $q/p_c$  VERSUS  $p_c$





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```
// JOB GKB 0367425.GKBSPACE N=EXTCOMP REG=500 T=.8 P=150 C=0
// *FORMAT PR, DDNAME=SYMSG, DEST=VM370
// EXEC PORTXCL, PARM.FORT='XREP'
// PORT.SYSIN DD *
```

MAIN PROGRAM FOR MONTE CARLO 3-D EXTERNAL FLOW CALCULATIONS  
OBJECTIVE OF THIS MAIN PROGRAM IS TO SET THE DIMENSIONS  
MAIN RUNNING PROGRAM IS \*\*\* RUN \*\*\*

FOLLOWING TWO CARDS HAVE TO BE ELIMINATED FOR NON IBM MACHINES  
\*\*\*\*\*  
INTEGER\*2 LB, NBM, NBN, NB, NBF, NBT, NBS, NUMCEL  
INTEGER\*2 LM, LPF, LCOL, LKW  
\*\*\*\*\*

THE NEXT CARD IS ASSOCIATED WITH PRINCETON RANDOM NUMBER GENERATOR  
\*\*\*\*\*  
COMMON/RANCOM/NRAN(4)  
\*\*\*\*\*

THE FOLLOWING DIMENSION STATEMENTS SET THE MAJOR ARRAY DIMENSIONS  
AND MUST BE CONSISTENT WITH THE FOLLOWING DATA CARD -----

NSP=NUMBER OF SPECIES - EXAMPLE BELOW NSP=1  
NMB=NUMBER OF BOXES WITHOUT COUNTING SUBDIVISIONS OR THOSE  
OCCUPIED BY THE BODY - EXAMPLE BELOW NMB=2600  
NMC=NUMBER OF FINAL CELLS - EXAMPLE BELOW NMC=1500  
NMP=MAX NUMBER OF MOLECULES OF EACH SPECIES ALLOWED IN PROGRAM.  
IF EXCEEDED, PROGRAM EITHER FAILS OR RESTARTS AT BEGINNING  
WITH NUMBER REDUCED BY 10% - EXAMPLE BELOW NMP=20000  
NPB=MAXIMUM NUMBER IN EACH CELL - EXAMPLE NPB=100

DIMENSION DBA(1, 1500), NB(1, 1500), NBF(1, 1500), NBT(1, 1500), LKW(1500)  
DIMENSION TMP(1, 1500), TRPA(1, 1500), YV(1, 1500), XVA(1, 1500)  
DIMENSION YV(1, 1500), YVA(1, 1500), ZV(1, 1500), ZVA(1, 1500), DR(1, 1500)  
DIMENSION TRP(1, 1500), TRPA(1, 1500), NBS(1, 1500), NBM(1, 1500)  
DIMENSION NBN(1500), T(1, 1, 1500)  
DIMENSION LB(20000), LM(1, 20000), ER(1, 20000)  
DIMENSION LPF(1, 20000), PAU(1, 20000), PAV(1, 20000), PAW(1, 20000)  
DIMENSION PAX(1, 20000), PAY(1, 20000), PAZ(1, 20000), LCOL(1, 20000)  
DIMENSION PNB(2600), XC(2600), YC(2600), ZC(2600), NUMCEL(2600)  
DATA NSP/1/, NMB/2600/, NMC/1500/, NMP/20000/, NPB/100/

2 FORMAT(/17X, 'NORMAL TERMINATION OF THE PROGRAM')  
NAMELIST/DIM/NSP, NMB, NMC, NMP, NPB, NRAN

INITIALIZATION OF RANDOM NUMBER GENERATOR - PRINCETON ROUTINE

NRAN(1)=0  
NRAN(2)=0  
NRAN(3)=0  
NRAN(4)=0

PRINTOUT OF MAJOR ARRAY DIMENSIONS USED ABOVE

FILE: GKTEXT DECK A

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WRITE(6,DIM)

CALL OF MAIN OPERATING PROGRAM WHICH REQUIRES INPUTS:  
 &CONTRL,&STIMES,&FLOREF,&MOLEC,&SHAPES,&GEOM,&COUPLE  
 THESE INPUTS ARE ALL CURRENTLY IN THE NAMELIST FORMAT  
 AND MAY HAVE TO BE CHANGED IF THAT CONVENTION IS NOT AVAILABLE  
 BRIEF DESCRIPTION OF THE PARAMETERS FOLLOWS

&CONTRL - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
NAME	8 BLANKS	ANY ALPHANUMERIC NAME UP TO 8 CHARACTERS
TITLE	24 BLNKS	ANY ALPHANUMERIC TITLE UP TO 24 CHARACTERS
PERCNT	.001	ACCURACY IN INTEGRATION PROCEDURES
ICOPI	1	NUMBER OF ADDITIONAL COPIES OF OUTPUT
DUMP	.TRUE.	IF TRUE WILL CAUSE SYSTEM DUMP FOR ANY OF 12 PROGRAMMER DESIGNED ERROR HALTS.
DEBUG (1)	.FALSE.	IF TRUE WILL PRINT MESSAGE WHEN CELL POP. EXCEEDS MNB
DEBUG (2)	.FALSE.	IF TRUE WILL PRINT CPU TIME AROUND EACH PART OF LOOP
DEBUG (3)	.TRUE.	IF TRUE WILL PRINT CPU TIME REMAINING AT END OF LOOP
NEW	.FALSE.	IF TRUE - NEW RUN - IF FALSE - RESTART OF RUN
SAVE	.FALSE.	IF TRUE - SNAPSHOT SAVED ON TAPE(9) FOR RESTART
REDO	.FALSE.	IF TRUE PROGRAM WILL AUTOMATICALLY RESTART WITH 90% OF TOTAL IF TOTAL CELL POPULATION EXCEEDS MNB

&STIMES - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
DTM	- - -	REAL NUMBER - FRACTION OF MEAN FREE TIME PER CYCLE
ITS	- - -	INTEGER - NUMBER OF CYCLES PER SAMPLE
ITP	- - -	INTEGER - NUMBER OF CYCLES BETWEEN PRINTOUTS
TST	- - -	INTEGER - ESTIMATE OF NUMBER OF CYCLES TO STEADY STATE
TLIM	- - -	INTEGER - TOTAL NUMBER OF CYCLES TO END OF RUN - WILL TERMINATE SOONER IF CPU TIME IS TO BE EXCEEDED

&FLOREF - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
INN	- - -	INITIAL NUMBER OF MOLECULES INN<MNB< OR = MNB
MNB	- - -	MAXIMUM NUMBER OF MOLECULES PER SPECIES
MNB	- - -	MAXIMUM NUMBER PER CELL - DIAGNOSTIC ONLY
MSP	- - -	NUMBER OF MOLECULAR SPECIES (MAX. IS 3)
MET	0	IF 0 - DATA IS IN SI (METRIC) UNITS IF >0 - DATA IS IN ENGLISH UNITS
U	- - -	FLOW VELOCITY (M/SEC) OR (FT/SEC)
ANGLE	- - -	ANGLE OF ATTACK (DEGREES)
RNU	0.0	ARRAY GIVING MOLE FRACTIONS OF SPECIES IN FREE STREAM
RMA	0.0	ARRAY GIVING MOLECULAR WEIGHTS OF SPECIES ABOVE
TF	- - -	FREE STREAM TEMPERATURE (K OR R)
DENP	- - -	FREE STREAM NUMBER DENSITY (NUM/M**3 OR NUM/PT**3)

&MOLEC - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
TRF	- - -	REFERENCE TEMPERATURE FOR MOLECULAR DATA
DIR	0.0	CROSS-SECTIONS AT REFERENCE TEMP. (MSP/MSP)
ETA	0.0	PARAMETERS IN DIFFUSION AND VISCOSITY LAW (MSP/MSP)
PHI	0.0	PARAMETERS FOR ROTATIONAL RELAXATION (MSP/MSP)
CHI	0.0	ROTATIONAL DEGREE OF FREEDOM PARAMETER (NROT/2 - 1)
ACR	.001	ACCURACY IN MOLECULAR COLLISION CALCULATIONS

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:ESHAPES - ND+1 OCCURRENCES WHERE ND=NUMBER OF BODY SEGMENTS (NEW RUN)
:PARAMETER DEFAULT DEFINITION
:           FIRST OCCURRENCE
:BODY(1) 0.0 STARTING POINT OF BODY FROM FRONT OF CELLS (M OR FT)
:BODY(I) I>1 - - NEED NOT BE SPECIFIED
:           SUBSEQUENT OCCURRENCES (ND)
:BODY(1) - - - X COORDINATE FROM FRONT OF BODY OF THE DOWNSTREAM
:           EDGE OF THE CURRENT BODY SEGMENT
:BODY(2) - - - TEMPERATURE OF THIS BODY SEGMENT
:BODY(3) - - - SURFACE AREA/TOTAL CROSS-SECTIONAL AREA FOR SEGMENT.
:           IF 0.0 PROGRAM WILL COMPUTE THIS QUANTITY
:BODY(4) - - - SWITCH - IF 0.0 THIS SEGMENT'S EQ. WILL APPEAR LATER
:           IF >0.0 THE EQ. OF THIS AND PRECEDING SEGMENTS IS
:           GIVEN BY BODY(6+2*MSP) TO BODY(9+2*MSP)
:BODY(5) - - - SWITCH - IF NOT 0.0 THIS IS THE LAST SHAPES CARD
:BODY(I) I EVEN ALPHA - ENERGY ACCOMODATION COEFFICIENT FOR SPECIES
:BODY(J) J ODD SIGMA - TANGENTIAL ACCOMODATION COEFF. FOR SPECIES
:           I AND J < (6+2*MSP)
:BODY(6+2*MSP) ORIGIN OF COORDINATES WITH RESPECT TO BODY START
:           FOR THE EQUATION OF THIS BODY SECTION
:BODY(7+2*MSP) - COEFFICIENTS A,B,C IN THE EQUATION
:BODY(9+2*MSP) R**2+A*X**2+B*X+C=0.0 FOR THIS BODY SECTION

:EGEOM - ONE OCCURRENCE (NEW RUN ONLY)
:PARAMETER DEFAULT DEFINITION
:NWEDGE - - - TWO INTEGERS GIVING THE NUMBER OF WEDGES BELOW
:           AND ABOVE THE ANGLE THETAZ
:THETAZ - - - ANGLE FROM NEG. Y AXIS DIVIDING DIFF. WEDGE SIZES
:RMB - - - MAX. BODY RADIUS - IF 0. WILL BE COMPUTED BY PROGRAM
:BW - - - WIDTH (DEL X) OF FIRST LEVEL CELLS (M OR FT)
:BH - - - HEIGHT (DEL R) OF FIRST LEVEL CELLS (M OR FT)
:NW - - - NUMBER OF FIRST LEVEL CELLS IN X DIRECTION
:NH - - - NUMBER OF FIRST LEVEL CELLS IN RADIAL DIRECTION
:NL - - - NUMBER OF LEVELS OF CELLS
:NPA - - - 1 NUMBER OF FIRST LEVEL CELLS AHEAD OF LEVEL 2
:NCA - - - 0 NUMBER OF FIRST LEVEL CELLS SUBDIVIDED INTO SECOND
:           LEVEL CELLS ALONG THE X DIRECTION
:NHA - - - 0 AS ABOVE BUT IN RADIAL DIRECTION
:MW - - - 0 NUMBER OF SECOND LEVEL CELLS IN THE X DIPECTION
:MR - - - 0 NUMBER OF SECOND LEVEL CELLS IN THE RADIAL DIRECTION
:NFB - - - 0 NUMBER OF SECOND LEVEL CELLS AHEAD OF LEVEL 3
:NCB - - - 0 NUMBER OF SECOND LEVEL CELLS SUBDIVIDED INTO THIRD
:           LEVEL CELLS ALONG THE X DIRECTION
:NHB - - - 0 AS ABOVE BUT IN RADIAL DIRECTION
:LW - - - 0 NUMBER OF THIRD LEVEL CELLS IN X DIRECTION
:LH - - - 0 NUMBER OF THIRD LEVEL CELLS IN RADIAL DIRECTION
:LD - - - NUMBER OF FIRST LEVEL CELLS FROM AXIS IN R DIRECTION
:           FOR WEIGHTING FACTOR BOUNDARIES (5 INTEGERS)
:LF - - - WEIGHTING FACTOR RATIOS AT BOUNDARIES ABOVE (5 INTEGERS)

:ECOUPLE - ONE OCCURRENCE (NEW RUN ONLY) - DISTRIBUTION FUNCTION
:PARAMETER DEFAULT DEFINITION
:NS - - - THE NUMBER OF BODY SEGMENTS FOR ACCUMULATING VELOCITY
:           DISTRIBUTION FUNCTION INFORMATION

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C MS      - - -   ARRAY(NS) OF AXIAL SEGMENT NUMBERS
C IVS     - - -   ARRAY(NS) OF AZIMUTHAL WEDGE NUMBERS
C VEL     3.,3.,4. THERMAL VELOCITY SPREAD FOR THE UNCOLLIDED MOLECULES
C MJ      20      NUMBER OF DEL V REGIONS FOR SAMPLING VELOCITY SPACE
C SL      - - -   ARRAY GIVING LOWER BOUND ON THE VELOCITY SAMPLE OF
C          - - -   COLLIDED MOLECULES (MSP X NS X 3)
C DELS    - - -   THE RANGE (SL<V<SL+DELS) FOR SAMPLE OF COLLIDED
C          - - -   MOLECULES (MSP X NS X 3)
IF DISTRIBUTION FUNCTION INFORMATION IS NOT DESIRED USE:
&COUPLE NS=0 &END
  
```

A SAMPLE INPUT DECK IS GIVEN BELOW:

```

&CONTRL NAME='SHUT', 'FILE ', TITLE='HYPE', 'EBOL', 'A AT', ' 95K', 'M M', 'ON. ',
  DEBUG=.F., .T., .T., NEW=.T., SAVE=.T., ICOPY=0, REDO=.T. &END
&TIMES DTM=.025, ITS=5, ITP=1000, TST=400, TLM=1000 &END
&FLOREP INM=4500, MNM=20000, MNB=100, MSP=1, MET=0, U=7485.9, ANGLE=0.0, RND=1., 2*0.,
  RMA=28.94, 0., 0., TP=195.51, DENP=2.52E+19 &END
&MOLEC TRF=1000, DIR=3.5E-19, ETA=.104, PHI=0.0, CHI=-1., ACR=.001 &END
&SHAPES BODY=1.00 &END
&SHAPES BODY=.0173, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.0672, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.1444, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.2432, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.3579, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.4842, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.6192, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.7405, 1590., 3*0.0, 2*1.0 &END
&SHAPES BODY=.7500, 1590., 0., 1., 0., 2*1., 0., -1.423278, -2.286, 0. &END
&SHAPES BODY=.9000, 1590., 0.0, 4*1.0, 0.9000, -111.82, 2*0.0 &END
&GEOM KWEDGE=1, 0, THETAZ=180., RMB=0.0, BW=.05, BH=.1, NW=40, NH=24, NL=2,
  NPA=15, NCA=20, NHA=16, NW=40, MH=32, LD=1, 2, 4, 8, 12, LP=3, 2, 2, 2, 1 &END
&COUPLE RS=3, MS=1, 5, 8, IVS=1, 1, 1, VEL=2.5, 2.5, 2.0, MJ=20, SL=9*1., 18*-9.,
  DELS=27*20. &END
  
```

```

CALL RUN(NSP, RMB, NMC, NMP, NPB, DBA, NB, NBP, NBT, TMP, TMPA, XV, YVA, YV,
  1YVA, ZV, ZVA, T, DB, FNB, XC, YC, ZC, NUMCEL, LM, LPP, PAU, PAV, PAW, PAX, PAY,
  2PAZ, LCOL, TRP, TRPA, ER, LKW, NBS, LB, NBM, NBN)
WRITE(6, 2)
STOP
END
  
```

```

SUBROUTINE RUN(NSP, RMB, NMC, NMP, NPB, DBA, NB, NBP, NBT, TMP, TMPA, XV, YVA,
  1 YV, YVA, ZV, ZVA, T, DB, FNB, IC, YC, ZC, NUMCEL, LM, LPP, PAU, PAV, PAW, PAX,
  2PAY, PAZ, LCOL, TRP, TRPA, ER, LKW, NBS, LB, NBM, NBN)
MAIN RUNNING PROGRAM ** RUN *** CALLS ALL OTHER SUBROUTINES
  
```

```

*****
INTEGER*2 LM, LPP, LCOL, LKW
INTEGER*2 LB, NBM, NBN, NB, NBP, NBT, NBS, NUMCEL
INTEGER PRT, SAMP, TST, TLM, TIME, Q
LOGICAL DUMP, DEBUG(3), SAVE, NEW, REDO
  
```

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REAL INTGRL,LAM,HU,NU,JAY,KAY  
DIMENSION BTA(3),C1(3),C2(3),C3(3),C7(3),C8(3),DPA(3),FL(3)  
DIMENSION DELANG(2),PDR(3),HTI(3),HTR(3),JNT(3),KHM(3),NM(3),SR(3)  
DIMENSION NAME(2),TITLE(6),NWEDGE(2),LD(5),LP(5),LWF(6),RLD(6)  
DIMENSION RNU(3),RMA(3),WTH(3),CHI(3),DIR(3,3),DAH(3,3),PHI(3,3)  
DIMENSION ETA(3,3),CNB(3,3),CNG(3),CNG(3),CN(3,3,3),CM(3,3,3)  
DIMENSION CTI(3,3),CTR(3,3),CNI(3,3),CNR(3,3),LEV(3),SN(3),ST(3)  
DIMENSION D1(3),D2(3),D3(3),D4(3),BODY(15),DBG1(3,3),LIMIT(10)  
DIMENSION VL(3,3),DELV(3,3),SSA(2,3),SSB(2,3),VEL(3),COEPP(4,9)  
DIMENSION XLIM(12),NTCP(3,3),NS(3),IWS(3),TANGN(3),NCOL(3,3)  
DIMENSION PV(3,3,2,20,3),NTCV(3,3,2,20,3),SL(3,3,3),DELS(3,3,3)  
DIMENSION XCB(18),XS(18),YCB(18),TB(18),ALPHA(3,18),SIGMA(3,18)  
DIMENSION NTS(3,18,12),NTSP(3,18,12),UTL(3,18,12),UTT(3,18,12)  
DIMENSION VTS(3,18,12),HTSI(3,18,12),HTS(3,18,12)  
DIMENSION UTLI(3,18,12),UTTI(3,18,12),VTSI(3,18,12)  
DIMENSION ENT(2,3,6,12),REM(2,3,6,12),ENTS(3,12),REMS(3,12)  
DIMENSION PTH(3,12),THETA(12),DTH(12)  
DIMENSION LB(NMP),NBN(NMC),NBN(NSP,NMC),LS(NSP,NMP)  
DIMENSION ER(NSP,NMP),TRP(NSP,NMC),TRPA(NSP,NMC)  
DIMENSION DBA(NSP,NMC),NB(NSP,NMC),NBP(NSP,NMC),NBT(NSP,NMC)  
DIMENSION THP(NSP,NMC),THPA(NSP,NMC),YV(NSP,NMC),YVA(NSP,NMC)  
DIMENSION YV(NSP,NMC),YVA(NSP,NMC),ZV(NSP,NMC),ZVA(NSP,NMC)  
DIMENSION T(NSP,NSP,NMC),DB(NSP,NMC),LKW(NMC),NBS(NSP,NMC)  
DIMENSION PNB(NMB),YC(NMB),YC(NMB),ZC(NMB),NUHCEL(NMB)  
  
DIMENSION LPP(NSP,NMP),PAU(NSP,NMP),PAV(NSP,NMP),PAW(NSP,NMP)  
DIMENSION PAX(NSP,NMP),PAY(NSP,NMP),PAZ(NSP,NMP),LCOL(NSP,NMP)  
  
EXTERNAL PNCTN,PRCTN  
  
COMMON /RANCOM/NRAN(4),KAWLS  
COMMON /FIRST/NL,NW,NH,MW,MH,LW,LE,EXA,NXB,NCA,NCB,NFA,NPB,NHA,NHBRUNO480  
COMMON /SECON/BW,BH,BWB,BHB,BWC,BHC,XLB,XLC RUNO490  
COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN RUNO500  
COMMON /PORTH/NBY,RY,RY,DUMP,C9,BWPM,LL RUNO510  
COMMON /PIPTH/ND,TIME,DTM,TI,ITS,ITP,TST,TLIM,RMA,RND,DIR RUNO520  
COMMON /SIXTH/RMB,IXTAET,JNM,MNM,MNB,NEW,SAVE,PERCENT,NSR,TR RUNO530  
COMMON /SVNTH/LAM,HU,NU,MT,N,J,I,Y,Z,TOSE RUNO540  
COMMON/EIGHT/DENP,U,TP,ANGLE,TRP,CHI,PHI,ETA,WTH,DAH,VELR,XREP  
NAMELIST/CONTRL/NAME,TITLE,PERCENT,ICOPI,DUMP,DEBUG,NEW,SAVE,REDO  
NAMELIST/TIMES/DTM,ITS,ITP,TST,TLIM RUNO580  
NAMELIST/FLOREP/INM,MNM,MNB,MSP,MET,U,ANGLE,RNU,RMA,TP,DENP  
NAMELIST/MOLEC/TRP,DIR,ETA,PHI,CHI,ACR  
NAMELIST/SHAPES/BODY RUNO570  
NAMELIST/GEOM/NWEDGE,THETAZ,RMB,BW,BH,NW,NH,NL,NFA,ECA,NHA,MW,MH,N  
HB,LW,LE,LD,LP,NPB,NCB  
NAMELIST/COUPLE/NS,MS,IWS,VEL,NJ,SL,DELS RUNO590  
DATA IC/0/,ICOPI/1/  
DATA DBG1/' GAS', AT ', '110 ', 'FLOW', AT ', '130 ', ' RUN', AT ', RUNO630  
1'303 '// RUNO640  
DATA LIMIT/12,9,18,500,3600,70,900,3,20,3/  
DATA TITLE/' // RUNO660  
DATA NAME/' // RUNO670  
DATA CPC/0.0/,CPH/0.0/,CPB/0.0/,CPJ/15.0/  
  
RUNO360  
RUNO440  
RUNO450  
RUNO460  
RUNO480  
RUNO490  
RUNO500  
RUNO510  
RUNO520  
RUNO530  
RUNO540  
RUNO580  
RUNO570  
RUNO590  
RUNO630  
RUNO640  
RUNO660  
RUNO670  
RUNO680

```

*****RUN0690
*****RUN0700
*****RUN0710
*****RUN0720
*****RUN0730
*****RUN0740
*****RUN0750
*****RUN0760
*****RUN0770
*****
          FORMATS
*****
1  FORMAT (1H1)
2  FORMAT (1H1/17X,'RARIIFIED SUPERSONIC FLOW OF BINARY GAS',T74,'I')
3  FORMAT ('+',103X,'COPY ',I2)
4  FORMAT (/17X,'FLOW THROUGH ALL THE BOUNDARIES'//)
5  FORMAT (3X,3I4,6E18.6/)
6  FORMAT (7X,2I4,      E18.6,72X,E18.6)
30 FORMAT ('1TIME =',F6.3,60X,'RANDOM NUMBER GENERATOR HAS BEEN CALLED
1  ',I10,' TIMES')
31 FORMAT (' CPU TIME LEFT- ',F8.3)
32 FORMAT (7X,'-MOLECULES-'/3X,3I6)
33 FORMAT (' TIME = ',F8.3,5X,'COLLISION LOOP=',F8.3,5X,'MOVE LOOP = '
1  ,F8.3,5X,'TOTAL TIME = ',F8.3/21X,'2ND MOVE LOOP = ',F8.3,5X,
2  'CLEANUP LOOP=',F8.3,4X,'PARTICLE NUMBERS = ',4I6)
34 FORMAT (9X,'-MOLECULAR COLLISIONS-'/3(3I14/))
35 FORMAT (2X,'-COLLISIONS WITH SURFACE-'/3I,3I8)
36 FORMAT (' MAXIMUM NUMBER OF MOLECULES SO FAR- ',I6//)
38 FORMAT (' EXCESS MOLECULES OCCURRED IN ',3A4)
40 FORMAT (' SOMETHING IS WRONG WITH BOX NUMBERING IN RUN '/5I5,2E17.
17,3I5,2E17.7,5I5/2E17.7)
42 FORMAT (' SOMETHING WRONG IN COMPUTING XSTART'/1X,8E16.8)
44 FORMAT (' NB ('.I2,'.',I4,') POPULATION EXCEEDED ',I3,' IN MAIN AT TRUN0930
1  TIME = ',F7.3)
50 FORMAT (///' SNAP SAVED ON TAPE')
*****
          CPA=ELTIME (0)
          CALL NOUNDF
          LIMIT (4)=NMC
          LIMIT (5)=NMP
          LIMIT (6)=NBP
          LIMIT (7)=NBB
          LIMIT (10)=NSP
          KAWLS=0
          PI=3.141593
          PIROOT=SQRT (PI)
          NET=0
          LARGE=0
          NL=1
          NPA=0
          NCA=0
          NRA=0
          NW=0
          NB=0
          NPB=0
          ECB=0
          NHB=0
          LW=0
          LH=0
*****
          RUN0800
          RUN0810
          RUN0880
          RUN0890
          RUN0900
          RUN0910
          RUN0920
          RUN0930
          RUN0940
          RUN0950
          RUN0960
          RUN0970
          RUN0980
          RUN1030
          RUN1040
          RUN1060
          RUN1080
          RUN1090
          RUN1100
          RUN1110
          RUN1120
          RUN1130
          RUN1140
          RUN1150
          RUN1160
          RUN1170
          RUN1180

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ILE: GKBEXT DECK A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

```

MJ=20
DUMP=.TRUE.
DEBUG(1)=.FALSE.
DEBUG(2)=.FALSE.
DEBUG(3)=.TRUE.
SAVE=.FALSE.
NEW=.TRUE.
REDO=.FALSE.
PERCNT=.001
ACR=.001
DO 58 I=1,15
58 BODY(I)=0.0
DO 60 I=1,3
RNU(I)=0.0
RMA(I)=0.0
CHI(I)=0.0
DO 59 J=1,18
ALPHA(I,J)=1.0
59 SIGMA(I,J)=1.0
DO 60 K=1,3
ETA(I,K)=0.0
PHI(I,K)=0.0
60 DIR(I,K)=0.0
VEL(1)=3.
VEL(2)=3.
VEL(3)=4.
WRITE(6,1)
READ(5,CONTRL)
WRITE(6,CONTRL)
IF(NEW) GO TO 103
REWIND 9
READ(9) DENP,U,XREP,TRF, KAWLS,NL,NW,NR,MW,MH,LW,LH,NXA,NXB,NCA
1 ,NCB,NPA,NPB,NHA,NHB,BW,BB,BWB,BHB,BWC,BHC,XLB,XLC,PI,NREG
2 ,S,SINANG,COSANG,AKN,NBX,RE,IR,ND,TIME,DTM,TI,ITS,ITP,TST
3 ,TLIM,RMA,RNU,DIR,XSTART,JNE,MNM,MNB,TR,BZC,CN7,DRP,PCF
4 ,PNA,HTP,INH,ITYPE,JTYPE,MJ,NAV,NMAX,NS,NWEDG,PRT,SAMP
5 ,BTA,C1,C2,C3,C7,C8,DAM,DPA,PL,DELANG,PDN,HTI,HTR,JNT,KNM
6 ,NM,WTH,C4,VRM,NCOL,LD,LP,LWP,RLD,CTI,CTR,CNI,CNE,LEV,SN
7 ,ST,D1,D2,D3,D4,SSA,SSB,MS,NSP,NMB,NMC,NMP,NPB,NRAN,VELR
8 ,IWS,TANGN,XLIM,COEPP,XCB,IS,YCB,TB,ALPHA,SIGMA,NTS,NTSF
9 ,UTL,UTT,VTS,HTS,HTSI,ENT,REM,ENTS,REMS,PTH,THETA,DTM,THPAR
A ,DBA,NB,NBF,NBT,TMP,YV,YVA,YV,ZV,ZVA,T,DE,PNB,IC,YC,ZCRUN
B ,NUMCEL,PAU,PAV,PAW,PAX,PAY,PAZ,FV,NTCV,NTCP,LPP,LCOL,LM
C,ETA,PHI,CHI,CN,CM,CNG,CMG,CN8,TRP,TRPA,THETAZ,NWEDGE,MSP,ANGLE,TF
D,UTLI,UTTI,VTSI,ER,RMB,LKW,NBS,LB,NBM,NBN
REWIND 9
DTMO=DTM
READ(5,TIMES)
WRITE(6,TIMES)
IF(DTM.EQ.DTMO) GO TO 100
AIME=TIME*DTMO
TIME=AIME/DTM+0.1
DO 99 J=1,NSP
DO 99 L=1,NWEDG
ENTS(J,L)=ENTS(J,L)*DTM/DTMO

```

RUN1190

RUN1230  
RUN1260  
RUN1240  
RUN1250

RUN1300  
RUN1310  
RUN1320  
RUN1330  
RUN1340

RUN1360  
RUN1370

RUN1380

RUN1390

RUN1400

RUN1410

RUN1420

RUN1430

RUN1440

RUN1460

RUN1470

RUN1480

RUN1490

RUN1500

```

DO 98 K=1,6
DO 98 I=1,2
ENT (I,J,K,L)=ENT (I,J,K,L)*DTM/DTMO
98 CONTINUE
99 CONTINUE
100 IF (TI.GT.0.0) TST=TI/DTM
WRITE (6,2)
WRITE (6,4)
WRITE (6,5) (( (I,J,L, (ENT (I,J,K,.,.,K=1,6), L=1, NWEDG), J=1, MSP), I=1,2)
WRITE (6,6) (( (J,L, ENTS (J,L), PTR (J,L), L=1, NWEDG), J=1, MSP)
WRITE (6,2)
CALL PRINTA (THETAZ, NWEDGE, TITLE, NAME, XCB, YCB, TB, ALPHA, SIGMA, LD, LP, RUN1540
1 XLIM, COEFP, LIMIT, MSP)
CALL PRINTB (PNA, MSP, PNB, LEV, LWP, NM, RLD, XLIM, XC, YC, ZC, NB, NUMCEL, LKW
1, MSP)
GO TO 280
103 READ (5, TIMES)
WRITE (6, TIMES)
READ (5, FLOREF)
WRITE (6, FLOREF)
READ (5, MOLEC)
WRITE (6, MOLEC)
IF (MSP.GT.LIMIT (10)) CALL DIAG (10, LIMIT (10), MSP)
CHIM=0.0
RMR=0.0
DMR=0.0
DO 115 M=1, MSP
RMR=RMR+RMA (M)*RNU (M)
CHIM=CHIM+CHI (M)*RNU (M)
DO 115 K=1, MSP
115 DMR=DMR+RNU (M)*RNU (K)*DIR (M,K)*(TRP/TF)**(ETA (M,K)/2.)
XREF=1./ (DENP*DMR*1.414214)
VELR=SQRT (16628.64*TF/RMR)
IF (MET.NE.0) VELR=SQRT (99437.92*TF/RMR)
THR=XREF/VELR
S=U/VELR
XLIM (1)=0.
TR=0.
YR=0.
NREG=0
ND=0
READ (5, SHAPES)
WRITE (6, SHAPES)
XO=BODY (1)/XREF
104 READ (5, SHAPES)
WRITE (6, SHAPES)
ND=ND+1
IF (ND.GT.LIMIT (3)) CALL DIAG (3, LIMIT (3), ND)
XCB (ND)=BODY (1)/XREF+XO
TB (ND)=BODY (2)/TF
YCB (ND)=BODY (3)
DO 1104 M=1, MSP
ALPHA (M, ND)=BODY (4+2*M)
1104 SIGMA (M, ND)=BODY (5+2*M)
IF (TB (ND).GT.TR) TR=TB (ND)

```

RUN1590

RUN1610

RUN1620

RUN1630

RUN1640

RUN1650

RUN1670

RUN1680

RUN1690

RUN1770



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FILE: GKBEYX DECK A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

```

IF (YCB (ND) .GT. YR) YR=YCB (ND)
IF (BODY (4) .EQ. 0.0) GO TO 104
NREG=NREG+1
X LIM (NREG+2)=XCB (ND)
IT=BODY (6+2*NREG) /XREP+XO
A=BODY (7+2*NREG)
B=BODY (8+2*NREG) /XREP
C=BODY (9+2*NREG) /XREP**2
COEFP (1, NREG)=A
COEFP (2, NREG)=1.0
COEFP (3, NREG)=B-2.*A*YT
COEFP (4, NREG)=A*YT**2-B*YT+C
IF (BODY (5) .EQ. 0.0) GO TO 104
IF (NREG.GT.LIMIT (2)) CALL DIAG (2, LIMIT (2), NREG)
NSTEP=NREG+2
A=COEFP (1, 1)
B=COEFP (3, 1)
C=COEFP (4, 1)
DISC=B*B-4.*A*C
IF (DISC.LT.0.) DISC=0.
DISC=SQRT (DISC)
IF (A.NE.0.) GO TO 108
XSTART=-C/B
GO TO 109
108 X1=.5/A* (-B+DISC)
X2=.5/A* (-B-DISC)
XSTART=AMIN1 (X1, X2)
XMAX=AMAX1 (X1, X2)
IF (XMAX.LT.XLIM (3)) XSTART=XMAX
IF (XSTART.GT.0.) GO TO 109
WRITE (6, 42) XSTART, A, COEFP (2, 1), B, C, DISC, X1, X2
IF (DUMP) CALL ABEND (1)
STOP
109 X LIM (2)=XSTART
A=COEFP (1, NREG)
B=COEFP (3, NREG)
C=COEFP (4, NREG)
DISC=B*B-4.*A*C
IF (DISC.LT.0.) DISC=0.
DISC=SQRT (DISC)
IF (A.NE.0.) GO TO 111
X LIM (NSTEP)=-C/B
GO TO 112
111 X1=.5/A* (-B+DISC)
X2=.5/A* (-B-DISC)
X LIM (NSTEP)=AMAX1 (X1, X2)
XMIN=AMIN1 (X1, X2)
IF (XMIN.GT.XLIM (NSTEP-1)) X LIM (NSTEP)=XMIN
IF (X LIM (NSTEP) .GT. X LIM (NSTEP-1)) GO TO 112
WRITE (6, 42) X LIM (NSTEP), A, COEFP (2, NREG), B, C, DISC, X1, X2
IF (DUMP) CALL ABEND (1)
STOP
112 AKN=1./ (X LIM (NSTEP)-XSTART)
XCB (ND)=X LIM (NSTEP)
DO 260 N=1,3

```

RUN1780

RUN1800

RUN1860

RUN1870

RUN1880

RUN1890

RUN1900

RUN1910

RUN1920

RUN1930

RUN1940

RUN1950

RUN1960

RUN1970

RUN1980

RUN1990

RUN2000

RUN2010

RUN2020

RUN2030

RUN2040

RUN2050

RUN2060

RUN2070

RUN2080

RUN2090

RUN2100

RUN2110

RUN2120

RUN2130

RUN2140

RUN2150

RUN2160

RUN2170

RUN2180

```

MS(N)=0
IWS(N)=0
DO 259 M=1,3
NCOL(N,M)=0
DO 258 K=1,3
SL(N,M,K)=0.0
258 DELS(N,M,K)=0.0
259 CONTINUE
260 CONTINUE
READ(5,GEOM)
READ(5,COUPLE)
WRITE(6,GEOM)
WRITE(6,COUPLE)
BW=BW/XREP
BH=BH/XREP
RMB=RMB/XREP
IF(RMB.GT.0.) GO TO 264
DO 262 K=1,NREG
XBEG=XLIM(K+1)
XEND=ILIM(K+2)
A=COEPP(1,K)
B=COEPP(2,K)
C=COEPP(3,K)
D=COEPP(4,K)
REND=SQRT(ABS((A*XEND**2+C*XEND+D)/B))
IPEAK=0.
IF(A.NE.0.) IPEAK=-.5*C/A
IF((XPEAK.LE.XBEG).OR.(XPEAK.GE.XEND)) GO TO 261
RTEMP=SQRT(ABS((A*IPEAK**2+C*IPEAK+D)/B))
IF(RTEMP.GT.REND) REND=RTEMP
261 IF(REND.GT.RMB) RMB=REND
262 CONTINUE
264 CONTINUE
NWEDG=NWEDGE(1)+NWEDGE(2)
IF(NWEDGE(2).EQ.0) THETAZ=180.0
IF(SAVE) REWIND 9
IF(NWEDG.GT.LIMIT(1)) CALL DIAG(1,LIMIT(1),NWEDG)
IF(MM.GT.LIMIT(5)) CALL DIAG(5,LIMIT(5),MM)
IF(MNB.GT.LIMIT(6)) CALL DIAG(6,LIMIT(6),MNB)
IF(NS.GT.LIMIT(8)) CALL DIAG(8,LIMIT(8),NS)
IF((NS.NE.0).AND.(NJ.GT.LIMIT(9))) CALL DIAG(9,LIMIT(9),NJ)
JNN=INN
DELANG(1)=THETAZ/NWEDGE(1)
DELANG(2)=0.0
IF(NWEDGE(2).NE.0) DELANG(2)=(180.-THETAZ)/NWEDGE(2)
SINANG=SIN(ANGLE/180.*PI)
COSANG=COS(ANGLE/180.*PI)
XR=BW*NR
XLIM(NSTEP+1)=XR
RH=BH*NR
VOL=PI*RH*RH*XR
NXA=NW*NH*NWEDG
NXB=HW*HB*NWEDGE(1)
NXC=LW*LB*NWEDGE(1)
IXA=NW*NH*NWEDGE(1)

```

RUN2200

RUN2220

RUN2230

RUN2240

RUN2250

RUN2290

RUN2300

RUN2320

RUN2330

RUN2380

RUN2390

RUN2400

RUN2420

RUN2430

RUN2440

RUN2450

```

NBX=NXA+NXB+NXC
IP(NBX.GT.LIMIT(7)) CALL DIAG(7,LIMIT(7),NBX)
BR=SQRT(TR)
DO 113 N=1,5
113 RLD(N)=BH*LD(N)
RLD(6)=RM
LWP(1)=1
RWPM=RLD(1)
B=RWPM*RWPM
C=B
DO 114 N=2,6
A=RLD(N)*RLD(N)
LWP(N)=LWP(N-1)*LP(N-1)
B=B+(A-C)/LWP(N)
D=RLD(N)/LWP(N)
IF(D.GT.RWPM) RWPM=D
114 C=A
INM=INM*RM*RM/B
DDN=INM/VOL
DO 140 MT=1,MSP
WTM(MT)=RMA(MT)/RMR
FDN(MT)=RNU(MT)*DDN
DPA(MT)=RNU(MT)
BTA(MT)=SQRT(WTM(MT))
SR(MT)=S*BTA(MT)
SN(MT)=SR(MT)*COSANG
ST(MT)=SR(MT)*SINANG
DO 117 K=1,MSP
DAM(K,MT)=DIR(K,MT)*(TRF/TF)**(ETA(K,MT)/2.)/DMR
CH8(K,MT)=DDN/DAM(K,MT)*1.414214
BT=AMIN1(BTA(K),BTA(MT))
VR1=S+3.*(1.+SQRT(TR))/BT
VR2=3.*SQRT((1.+2.*S**2/(5.+CHIM))*(1./WTM(K)+1./WTM(MT)))
CH(K,MT,1)=AMAX1(VR1,VR2)
CN(K,MT,1)=RAND(0)*CH(K,MT,1)
DF=PHI(K,MT)*(CHI(K)+CHI(MT)+2.)-1
DS=PHI(K,MT)*(2.-.5*ETA(K,MT))-1.0
DO 917 N=2,3
IPN=ACR**AMIN1(DF,DS)
IF((DF.GT.0.).AND.(DS.GT.0.)) XPN=(DF/(DF+DS)**DF*(DS/(DF+DS)**DS)
XPN=ACR**AMAX1(DF,DS)
IF((DF.LT.0.).AND.(DS.LT.0.)) IPN=(DF/(DF+DS)**DF*(DS/(DF+DS)**DS)
CH(K,MT,N)=IPN-XPN
CN(K,MT,N)=RAND(0)*CH(K,MT,N)
DF=CHI(K)
DS=CHI(MT)
917 CONTINUE
117 CONTINUE
ARG=SN(MT)
DO 119 NT=1,2
D=ERRF(ARG)
TEMPA=EXP(-SN(MT)*SN(MT))/3.544908+0.5*ARG*D
TEMPA=TEMPA/BTA(MT)
TEMPC=0
DO 118 N=1,6

```

```

RUN2460
RUN2470
RUN2480
RUN2490
RUN2500
RUN2510
RUN2520
RUN2530
RUN2540
RUN2550
RUN2560
RUN2570
RUN2580
RUN2590
RUN2600
RUN2610
RUN2620
RUN2630

```

```

RUN2730
RUN2740
RUN2750

```

```

RUN2760
RUN2770
RUN2780
RUN2790
RUN2800
RUN2810
RUN2820

```

```

TEMPB=RLD(N)*RLD(N)
TEMPD=INH*TEMPA*DTH/(XR*RM*RM)*(TEMPB-TE MPC)/LWF(N)
DO 116 K=1,NWEDG
DELTH=DELANG(1)
IF(K.GT.NWEDGE(1)) DELTH=DELANG(2)
ENT(NT,MT,N,K)=TEMPD*DELTH/180.
116 REN(NT,MT,N,K)=0.
TE MPC=TEMPB
118 CONTINUE
SSR=ARG*ARG
SSA(NT,MT)=ARG+SQRT(SSR+2.)
SSB(NT,MT)=SSA(NT,MT)*(.25*SSA(NT,MT)-ARG)
119 ARG=-ARG
CHT=CHI(MT)
IF(CHT.GT.0.) CHG(MT)=CHT**CHT*EXP(-CHT)
IF(CHT.EQ.0.) CHG(MT)=1.0
IF(CHT.LT.0.) CHG(MT)=ACR**CHT*EXP(-ACR)
CMG(MT)=RAND(0)*CMG(MT)
ANG=0.
N=0
ARG=ST(MT)
TEMPD=INH/PI*DTH/RM*2./(BTA(MT)*LWF(6))
DO 135 M=1,2
I=NWEDGE(M)
IF(I.EQ.0) GO TO 135
DO 134 K=1,I
N=N+1
AA=ANG
BB=DELANG(M)
CC=AA+BB
ANG=ANG+BB
THETA(N)=AA
DTH(N)=BB
AA=AA*PI/180.
CC=CC*PI/180.
ARGA=ARG*COS(AA)
ARGC=ARG*COS(CC)
ENTS(MT,N)=0.
REMS(MT,N)=0.
PTH(MT,N)=0.
TE PCA=0.
IF(ARGA.LE.-10.) GO TO 120
TE PCA=FNCTM(ARGA,PIROOT,L,COEFF)
PTH(MT,N)=2.*PIROOT*TE PCA
120 TEMPB=0.
IF(ARGC.LE.-10.) GO TO 125
TEMPB=FNCTM(ARGC,PIROOT,L,COEFF)
125 SUM1=TE PCA+TEMPB
TE MPC=.5*SUM1*(CC-AA)*TEMPD
IF(TE MPC.LT.1.E-06) GO TO 134
CALL SIMPSN(ARGA,ARGC,0,INTGRL,PERCNT,COEFF,PIROOT,SUM1,FNCTM)
ENTS(MT,N)=INTGRL*TEMPD*(CC-AA)
134 CONTINUE
135 CONTINUE
140 CONTINUE

```

```

RUN2830
RUN2840
RUN2850
RUN2860
RUN2870
RUN2880
RUN2890
RUN2900
RUN2910
RUN2920
RUN2930
RUN2940
RUN2950
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RUN2970
RUN2960
RUN2980
RUN2990
.
RUN3010
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RUN3020
RUN3030
RUN3040
RUN3050
RUN3060
RUN3070
RUN3080
RUN3090
RUN3100
RUN3110
RUN3120
RUN3130
RUN3140
RUN3150
RUN3160
RUN3170
RUN3180
RUN3190
RUN3200
RUN3210
RUN3220
RUN3230
RUN3240
RUN3250
RUN3260
RUN3270
RUN3280
RUN3290
.
RUN3330

```

```

XS (1) = .5 * (XCB (1) - XSTART) * AKN
DO 155 N=2, ND
155 IS (N) = (.5 * (XCB (N) + XCB (N-1)) - XSTART) * AKN
IF (NS.EQ.0) GO TO 160
DO 159 I=1, NS
N=NS (I)
X=XS (N) / AKN + XSTART
J=0
157 J=J+1
IF (X.GT.XLIM (J+2)) GO TO 157
CALL HEIGHT (X, Y, J, COEFF, 3)
TANGN (I) = (IWS (I) - .5) * DELANG (1)
IF (IWS (I) * DELANG (1) .GT. THETAZ) TANGN (I) = THETAZ + (IWS (I) - NWEDGE (1) -
15) * DELANG (2)
Z=Y * SIN (TANGN (I) * PI / 180.)
Y=-Y * COS (TANGN (I) * PI / 180.)
CALL NORMAL (EYE, JAY, KAY, ONE, COEFF)
SNN=-S * (COSANG * EYE + SINANG * JAY)
ST1=S * (COSANG * ONE - EYE * SINANG * JAY / ONE)
ST2=-S * SINANG * KAY / ONE
DO 159 MT=1, MSP
VL (MT, 1) = AMAX1 (0., SNN - VEL (1) / BTA (MT))
VL (MT, 2) = ST1 - VEL (2) / BTA (MT)
VL (MT, 3) = ST2 - VEL (3) / BTA (MT)
DELV (MT, 1) = SNN + VEL (1) / BTA (MT) - VL (MT, 1)
DELV (MT, 2) = 2. * VEL (2) / BTA (MT)
DELV (MT, 3) = 2. * VEL (3) / BTA (MT)
AMJ=MJ-1
DO 159 K=1, 3
DO 159 J=1, MJ
PV (MT, I, 1, J, K) = VL (MT, K) + (J-1) / AMJ * DELV (MT, K)
PV (MT, I, 2, J, K) = SL (MT, I, K) + (J-1) / AMJ * DELS (MT, I, K)
159 CONTINUE
160 CONTINUE
IF (YR.GT.0.) GO TO 169
A=XSTART
B=XCB (1)
L=1
SUM1=FNCTN (A, PIROOT, L, COEFF) + FNCTN (B, PIROOT, L, COEFF)
CALL SIMPSN (A, B, L, INTGRL, PERCNT, COEFF, PIROOT, SUM1, FNCTN)
INTGRL=INTGRL / (ABS (COEFF (2, L)) * RMB * RMB) * (B-A)
YCB (1) = INTGRL
DO 168 N=2, ND
A=XCB (N-1)
B=XCB (N)
IF (XLIM (L+2) .GE. B) GO TO 167
L=L+1
167 SUM1=FNCTN (A, PIROOT, L, COEFF) + FNCTN (B, PIROOT, L, COEFF)
CALL SIMPSN (A, B, L, INTGRL, PERCNT, COEFF, PIROOT, SUM1, FNCTN)
INTGRL=INTGRL / (ABS (COEFF (2, L)) * RMB * RMB) * (B-A)
168 YCB (N) = INTGRL
169 LV=2
IF (NWEDGE (2) .EQ. 0) LV=1
CALL CELL (THETAZ, LV, BW, BH, NW, NH, 0., 0, 0, DELANG, NWEDGE, IC, IC, ZC, PNB)
IF (NL.LT.2) GO TO 170

```

RUN3370  
 RUN3380  
 RUN3390  
 RUN3400  
 RUN3410  
 RUN3420  
 RUN3430  
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 RUN3790  
 RUN3800  
 RUN3810  
 RUN3820  
 RUN3830  
 RUN3840  
 RUN3850  
 RUN3860  
 RUN3870  
 RUN3890

CALL ZERO (NW, NH, NFA, NCA, NHA, 0, NWEDGE (1), PNB)	RUN3900
XLB=BW*NFA	RUN3910
BWB=BW*NCA/MW	RUN3920
BBB=BH*NHA/MH	RUN3930
CALL CELL (THETAZ, 1, BWB, BBB, MW, MH, XLB, NXA, 0, DELANG, NWEDGE, XC, YC, ZC, PNB)	RUN3940
1 PNB)	RUN3950
IF (NL.LT.3) GO TO 170	RUN3960
CALL ZERO (MW, MH, NFB, NCB, NHB, NIA, NWEDGE (1), PNB)	RUN3970
XLC=XLB+BWB*NFB	RUN3980
BWC=BWB*NCB/LW	RUN3990
BHC=BHB*NHB/LE	RUN4000
CALL CELL (THETAZ, 1, BWC, BHC, LW, LH, XLC, NXA, NXB, DELANG, NWEDGE, XC, YC, ZC, PNB)	RUN4010
1C, PNB)	RUN4020
170 CALL SBTRECT (1, IXA, NIA, BW, BH, DELANG, XC, YC, PNB, XLIM, COEFF)	RUN4030
LEV (1) =NBX+1	
LEV (2) =NBY+1	
IF (NL.LT.2) GO TO 190	RUN4040
NI=NXA+1	RUN4050
NP=NIA+NXB	RUN4060
LEV (1) =NI	RUN4070
CALL SBTRECT (NI, NP, NP, BWB, BBB, DELANG, XC, YC, PNB, XLIM, COEFF)	RUN4080
IF (NL.LT.3) GO TO 190	RUN4090
NI=NP+1	RUN4100
NP=NBX	RUN4110
LEV (2) =NI	RUN4120
CALL SBTRECT (NI, NP, NP, BWC, BHC, DELANG, XC, YC, PNB, XLIM, COEFF)	RUN4130
190 PNA=0.0	RUN4140
H=0	RUN4150
DO 210 N=1, NBX	RUN4160
NUMCEL (N) =0	RUN4170
IF (PNB (N) .LE.0.) GO TO 210	RUN4180
M=M+1	RUN4190
NUMCEL (N) =M	RUN4200
PNA=PNA+PNB (N)	RUN4210
DYC=BH/2.	
IF (N.GT.NIA) DYC=BHB/2.	
IF (N.GT.NXA+NXB) DYC=BHC/2.	
YTC=YC (N) +DYC	
DO 200 LA=1, 6	
IF (YTC.LE.RLD (LA)) GO TO 201	
200 CONTINUE	
201 LKW (M) =LWP (LA)	
210 CONTINUE	
NPX=M	RUN4220
IF (NPX.GT.LIMIT (4)) CALL DIAG (4, LIMIT (4), NPX)	RUN4230
220 TIME=0	RUN4240
LARGE=0	RUN4250
SAHP=0	RUN4260
PRT=0	RUN4270
NAV=0	
TIME=0.	RUN4290
TI=1.	RUN4300
NHAX=0	RUN4310
DO 250 NT=1, 3	RUN4380
C1 (NT) =RAND (0)	RUN4410

C2 (MT) =RAND(0)	RUN4420
C3 (MT) =RAND(0)	RUN4430
C7 (MT) =RAND(0)	RUN4440
C8 (MT) =RAND(0)	RUN4450
D1 (MT) =RAND(0)	RUN4460
D2 (MT) =RAND(0)	RUN4470
D3 (MT) =RAND(0)	RUN4480
D4 (MT) =RAND(0)	RUN4490
FL (MT) =0.	RUN4500
HTI (MT) =0.	RUN4510
HTR (MT) =0.	RUN4520
JNT (MT) =0	RUN4530
NM (MT) =0	RUN4540
DO 230 N=1,3	RUN4550
CTI (MT,N) =0.	RUN4560
CTR (MT,N) =0.	RUN4570
CNI (MT,N) =0.	RUN4580
230 CNR (MT,N) =0.	RUN4590
DO 240 N=1,KD	RUN4600
DO 240 K=1,NWEDG	RUN4610
NTS (MT,N,K) =0	RUN4620
NTSP (MT,N,K) =0	RUN4630
HTSI (MT,N,K) =0.	RUN4640
UTLI (MT,N,K) =0.	
UTTI (MT,N,K) =0.	
VTSI (MT,N,K) =0.	
UTL (MT,N,K) =0.	RUN4650
UTT (MT,N,K) =0.	RUN4660
VTS (MT,N,K) =0.	RUN4670
240 HTS (MT,N,K) =0.	RUN4680
DO 250 I=1,3	
NTCP (MT,I) =0	RUN4820
DO 250 L=1,2	RUN4830
DO 250 K=1,3	RUN4840
DO 250 J=1,MJ	RUN4850
HTCV (MT,I,L,J,K) =0	RUN4860
250 CONTINUE	RUN4870
DO 245 N=1,NPX	RUN4690
DO 245 MT=1,NSP	
NB (MT,N) =0	RUN4700
NBP (MT,N) =0	RUN4710
NBS (MT,N) =0	
NBT (MT,N) =0	
DBA (MT,N) =0.	RUN4720
XVA (MT,N) =0.	RUN4730
YVA (MT,N) =0.	RUN4740
ZVA (MT,N) =0.	RUN4750
TRPA (MT,N) =0.	RUN4760
TRPA (MT,N) =0.0	RUN4790
DO 245 NN=1,NSP	
T (MT,NN,N) =0.0	
FND=DDN	
245 CONTINUE	RUN4800
DRF=2./ (FND*S*S*RMB*RMB*PI)	RUN4930
PCF=1./ (FND*S*RMB*RMB*PI)	RUN4940

```

HTF=.5*DRP/S
C9=RAND(0)*RWFN
LL=PNA/VOL*INH
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) ((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)
WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)
WRITE(6,2)
CALL PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,LD,LP,
1XLM,COEPP,LIMIT,MSP)
CALL GAS(NWEDG,THETAZ,DELANG,NWEDGE,BTA,C1,DPA,NE,RLD,LWF,FNB,DB,
1B,NBP,LPP,PAU,PAV,PAW,PAX,PAY,PAZ,XLM,COEPP,LM,LIMIT(4),LIMIT(6),
2LARGE,MNM,MNB,DEBUG(1),LCOL,NUMCEL,MSP,ER,CHI,CNG,CMG,NSP,LE,NBM,
3NBN)
CPUTYM=TFIND(0)
IF(LARGE.NE.0) GO TO 345
DO 265 I=1,MSP
265 IF(NM(I).GT.NMAX) NMAX=NM(I)
CALL PRINTB(PNA,MSP,PNB,LEV,LWF,NE,RLD,XLM,IC,YC,ZC,NB,NUMCEL,LKW
1,NSP)
IF(DEBUG(2)) WRITE(6,1)
CALL ACCUM(NMC,NPB,FNB,NE,PAU,PAV,PAW,ER,TRP,TRP,IV,YV,ZV,LM,MSP,
1NSP,LPP,NBP,NBM)
CPA=ELTIME(0)
CPI=CPA
GO TO 340
280 TIME=TIME+1
285 LARGE=0
CPI=ELTIME(0)
AIME=TIME*DTM
IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
PRT=PRT+1
SAMP=SAMP+1
CALL COLIDE(CN,CE,WTM,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,LM,MSP,
1LIMIT(4),LIMIT(6),NUMCEL,ETA,PHI,CHI,CN8,NSP,LPP,LKW,NBP,NBM)
KNM(1)=0
KNM(2)=0
KNM(3)=0
CPC=ELTIME(0)
IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
CALL MOVE(0,AKN,EJ,NS,NWEDG,THETAZ,XSTART,LIMIT(3),LIMIT(1),LIMIT
1(8),LIMIT(9),DELANG,NWEDGE,BTA,C2,C3,DPA,PL,HTI,HTR,JNT,KNM,NE,XCB
2,XLM,ES,IWS,NTCP,NTCV,FV,CTI,CTB,CNI,CNR,ALPHA,SIGMA,COEPP,HTS,HTR
3SI,HTS,NTSP,UTL,UTT,VTS,PAU,PAV,PAW,PAX,PAY,PAZ,LPP,LCOL,TB,MSP,ER
4,CHI,CNG,CMG,NSP,UTLI,UTTI,VTSI)
KNM(1)=NM(1)
KNM(2)=NM(2)
KNM(3)=NM(3)
CPM=ELTIME(0)
IF(DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
CALL FLOW(NWEDG,MNM,LARGE,BTA,C1,C7,C8,D1,D2,D3,D4,DTH,NE,SN,ST,TR
1ETA,LWF,RLD,PTH,ENTS,REMS,SSA,SSB,PAU,PAV,PAW,PAX,PAY,PAZ,LPP,ENT,
2REN,LCOL,MSP,ER,CHI,CNG,CMG,NSP)
IF(LARGE.NE.0) GO TO 345
CPB=ELTIME(0)

```

RUN4960  
 RUN4970

RUN5000  
 RUN5010

RUN5030  
 RUN5040

RUN5060

RUN5130  
 RUN5140

RUN5170

RUN5180  
 RUN5190

RUN5220  
 RUN5230

RUN5280  
 RUN5290

RUN5330



```

IF (DEBUG (1)) WRITE (6, 33) AIME, CPC, CPM, CPI, CPB, CPA, (NM(I), I=1, 3), NMAX
CALL MOVE (1, AKN, SJ, NS, NWEDG, THETAZ, XSTART, LIMIT (3), LIMIT (1), LIMIT
1 (8), LIMIT (9), DELANG, NWEDGE, BTA, C2, C3, DPA, FL, HTI, HTR, JNT, KNH, NH, XCBRUN5350
2, XLIM, HS, IWS, NTCF, NTCV, PV, CTI, CTR, CNI, CNR, ALPHA, SIGMA, COEPP, HTS, HTRUN5360
3SI, NTS, NTSF, UTL, UTT, VTS, PAU, PAV, PAW, PAX, PAY, PAZ, LPF, LCOL, TB, MSP, ER
4, CHI, CNG, CHG, NSP, UTLI, UTTI, VTSI)
DO 330 MT=1, MSP
CPB=CPB+ELTIME (0)
IF (DEBUG (1)) WRITE (6, 33) AIME, CPC, CPE, CPI, CPB, CPA, (NM(I), I=1, 3), NMAX
M=0
DO 290 N=1, NBX
IF (NUMCEL(N).EQ.0) GO TO 290
M=M+1
NB (MT, M) =0
NBF (MT, M) =0
290 CONTINUE
NG=NM (MT)
N=0
295 N=N+1
IF (N.GT.NG) GO TO 310
X=PAX (MT, N)
Y=PAY (MT, N)
Z=PAZ (MT, N)
R=SQRT (Y*Y+Z*Z)
ARG=Y/R
TANG=180. * (1.-ARCCOS (ARG)/PI)
IWDGE=TANG/DELANG (1)
IF ((IWDGE.GE.NWEDGE (1)).AND. (DELANG (2).NE.0.)) IWDGE=(TANG-THETAZ)
1/DELANG (2)+NWEDGE (1)
IF (IWDGE.LT.0) IWDGE=0
IF (IWDGE.GE.NWEDG) IWDGE=NWEDG-1
L=X/BW+1.
IF (L.GT.NW) L=NW
M=R/BH
IF (M.GE.NH) M=NH-1
K=(IWDGE*NH+M)*NW+L
IF (K.LE.NXA) GO TO 296
WRITE (6, 40) L, M, K, MT, N, X, Y, NWEDGE, NWEDG, TANG, Z, IWDGE, NH, NW
IF (DUMP) CALL ABEND (4)
STOP
296 KW=0
IF (NL.EQ.1) GO TO 300
IF (IWDGE.GE.NWEDGE (1)) GO TO 300
IF (FNB (K).GT.0.) GO TO 300
L=(X-XLB)/BWB+1.
IF (L.GT.MW) L=MW
M=R/BHB
IF (M.GE.MH) M=MH-1
K=(IWDGE*MH+M)*MW+L+NIA
IF (K.LE.NXA+NXB) GO TO 297
WRITE (6, 40) L, M, K, MT, N, X, Y, NWEDGE, NWEDG, TANG, Z, IWDGE, MH, MW, NIA
IF (DUMP) CALL ABEND (5)
STOP
297 IF (NL.EQ.2) GO TO 300
IF (FNB (K).GT.0.) GO TO 300

```

RUN5410  
 RUN5420  
 RUN5440  
 RUN5450  
 RUN5460  
 RUN5470  
 RUN5480  
 RUN5490  
 RUN5500  
 RUN5510  
 RUN5520  
 RUN5530  
 RUN5540  
 RUN5550  
 RUN5560  
 RUN5570  
 RUN5580  
 RUN5600  
 RUN5610  
 RUN5620  
 RUN5630  
 RUN5640  
 RUN5650  
 RUN5660  
 RUN5670  
 RUN5680  
 RUN5690  
 RUN5700  
 RUN5710  
 RUN5730  
 RUN5740  
 RUN5750  
 RUN5760  
 RUN5770  
 RUN5780  
 RUN5790  
 RUN5800  
 RUN5810  
 RUN5820  
 RUN5830  
 RUN5840

	L=(X-XLC)/BWC+1.	RUN5850
	IF (L.GT.LW) L=LW	RUN5860
	M=R/BWC	RUN5870
	IF (M.GE.LH) M=LH-1	RUN5880
	K=(INDGE*LH+M)*LW+L+NIA+HIB	RUN5890
	IF (K.LE.NBX) GO TO 300	RUN5900
	WRITE (6,40) L, M, K, MT, N, X, Y, NWEDGE, NWEDG, TANG, Z, INDGE, LH, LW, NIA, HIB	RUN5910
	IF (DUMP) CALL ABEND (6)	RUN5920
	STOP	RUN5930
200	LA=0	RUN5940
	IF (NUMCEL(K).EQ.0) GO TO 306	
301	LA=LA+1	RUN5960
	IF (R.GT.RLD(LA)) GO TO 301	RUN5970
	KW=LWF(LA)	RUN5980
	KKW=LPF(MT,N)	RUN5990
	IF (KW.EQ.KKW) GO TO 305	RUN6000
	IF (KW.LT.KKW) GO TO 302	RUN6010
	M=KW/KKW	RUN6020
	A=RAND(0)	RUN6030
	B=M	RUN6040
	B=1./B	RUN6050
	IF (A.LT.B) GO TO 305	RUN6060
306	PAX(MT,N)=PAX(MT,NG)	RUN6070
	PAY(MT,N)=PAY(MT,NG)	RUN6080
	PAZ(MT,N)=PAZ(MT,NG)	RUN6090
	PAU(MT,N)=PAU(MT,NG)	RUN6100
	PAV(MT,N)=PAV(MT,NG)	RUN6110
	PAW(MT,N)=PAW(MT,NG)	RUN6120
	ER(MT,N)=ER(MT,NG)	
	LPP(MT,N)=LPP(MT,NG)	RUN6130
	LCOL(MT,N)=LCOL(MT,NG)	RUN6140
	N=N-1	RUN6150
	NH(MT)=NH(MT)-1	RUN6160
	NG=NH(MT)	RUN6170
	GO TO 295	RUN6180
302	M=KKW/KW-1	RUN6190
	IF ((NH(MT)+M).LE.MNE) GO TO 307	
	LARGE=3	
	GO TO 345	RUN6230
207	CONTINUE	RUN6240
	DO 304 L=1,M	
	NH(MT)=NH(MT)+1	RUN6200
	NG=NH(MT)	RUN6210
	PAX(MT,NG)=PAX(MT,N)	RUN6260
	PAY(MT,NG)=PAY(MT,N)	RUN6270
	PAZ(MT,NG)=PAZ(MT,N)	RUN6280
	PAU(MT,NG)=PAU(MT,N)	RUN6290
	PAV(MT,NG)=PAV(MT,N)	RUN6300
	PAW(MT,NG)=PAW(MT,N)	RUN6310
	ER(MT,NG)=ER(MT,N)	
	LCOL(MT,NG)=LCOL(MT,N)	RUN6320
204	LPP(MT,NG)=KW	RUN6330
305	LPP(MT,N)=KW	RUN6340
	Q=NUMCEL(K)	RUN6350
	J=NB(MT,Q)+1	RUN6360

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IF (J.LE.MNB) GO TO 308
IF (DEBUG(1)) WRITE(6,44) MT,K,MNB,AIME
308 NB(MT,Q)=J
NBF(MT,Q)=NBF(MT,Q)+KW
LB(N)=Q
GO TO 295
310 CONTINUE
NBH(MT,1)=0
DO 320 N=1,NBX
M=NUMCEL(N)
IF (M.EQ.0) GO TO 320
A=NBF(MT,M)
DB(MT,M)=A*DPA(MT)/PNB(N)
NBH(MT,M+1)=NBH(MT,M)+NB(MT,M)
NBN(M)=NBH(MT,M)
320 CONTINUE
IF (NM(MT).GT.NMAX) NMAX=NM(MT)
DO 325 N=1,NG
Q=LB(N)
NBN(Q)=NBN(Q)+1
NA=NBN(Q)
325 LM(MT,NA)=N
330 CONTINUE
IF (SAMP.LT.ITS) GO TO 335
CALL ACCUM(NMC,NPB,PNB,NB,PAU,PAV,PAW,ER,TEP,TRP,IV,YV,ZV,LM,MSP,
1NSP,LPF,NBF,NEM)
SAMP=0
IF (TIME.LE.TST) GO TO 335
CALL AVRGE(PFB,DB,DEA,NB,NET,IV,YV,ZV,XVA,YVA,ZVA,TEP,TMPA,TRP,TRP
1A,MSP,NSP,NBP,NBS)
NAV=NAV+1
335 CPA=ELTIME(0)
CPI=CPC+CPM+CPB+CPA
CPJ=2.*CPI+5.
340 CPUTYM=TFIND(0)
IF (DEBUG(2)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
IF ((TIME.GE.TLIM).OR.(CPUTYM.LE.CPJ)) GO TO 345
IF (PRT.LT.ITP) GO TO 280
PRT=0
345 WRITE(6,30) AIME,KAWLS
IF (DEBUG(3)) WRITE(6,31) CPUTYM
WRITE(6,32) (NM(I),I=1,3)
WRITE(6,34) ((NCOL(I,J),J=1,3),I=1,3)
WRITE(6,35) (JNT(I),I=1,3)
IF (LARGE.NE.0) GO TO 360
WRITE(6,36) NMAX
IF (.NOT.SAVE) GO TO 355
IF (PRT.NE.0.AND.CPUTYM.GT.CPJ.AND.TIME.LT.TLIM) GO TO 355
WRITE(9) DENP,U,XREP,TRP,KAWLS,NL,NF,NE,EW,MH,LW,LH,NXA,NXB,NCA
1 ,NCB,NPA,NPB,NHA,NHB,BW,BH,BWB,BHB,BWC,BHC,XLB,XLC,PI,NREG
2 ,S,SINANG,COSANG,AKN,NBX,RM,IR,ND,TIME,DTM,TL,ITS,ITP,TST
3 ,TLIM,REA,BNU,DIR,XSTART,JNM,MNM,MNB,TR,BZC,CN7,DEF,PCF
4 ,PNA,HTP,INM,ITYPE,JTYPE,MJ,NAV,NMAX,NS,NWEDG,PRT,SAMP
5 ,BTA,C1,C2,C3,C7,C8,DAM,DPA,PL,DELANG,PDN,HTI,HTR,JNT,KNM
6 ,NH,WTH,C4,VRM,NCOL,LD,LP,LWP,RLD,CTI,CTR,CNI,CNR,LEV,SN

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7      ,ST,D1,D2,D3,D4,SSA,SSB,SS,NSP,NMB,NMC,NMP,NPB,NRAN,VZLR
8      ,IWS,TANGN,XLIM,COEPP,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,NTSP  RUN7000
9      ,UTL,UTT,VTS,HTS,HTSI,ENT,REM,ENTS,REMS,PTH,THETA,DTH,TMPARUN7010
A      ,DBA,NB,NBP,NBT,TMP,XV,YVA,YV,YVA,ZV,ZVA,T,DB,FNB,XC,YC,ZCRUN7020
B      ,NUMCEL,PAU,PAV,PAW,PAX,PAY,PAZ,PV,NTCV,NTCP,LPP,LCOL,LM  RUN7030
C,ETA,PHI,CHI,CN,CH,CNG,CMG,CNS,TRP,TRPA,THETAZ,NWEDGE,MSP,ANGLE,TF
D,UTLI,UTTI,VTSI,ER,RMB,LKW,NBS,LB,NBM,NBN
REWIND 9
WRITE(6,50)
355 CONTINUE
IF(TIME.LE.TST) GO TO 350
DT=AIME-TI
CALL PRINT1(DT,COSANG,SINANG,RMA,RNU,DRP,PCP,HTP,PL,HTI,HTE,CTI,
1CTR,CHI,CNR)
CALL PRINT2(AKN,XSTART,DT,RNU,RMA,DPP,PCP,HTP,UTLI,UTTI,VTSI,HTSI,
1DELANG,NWEDGE,XS,XCB,YCB,HTS,NTS,NTSP,UTL,UTT,VTS,LIMIT(3),
2LIMIT(1),MSP)
IF(NS.NE.0) CALL PRINT3(      MSP,MJ,NS,NWEDG,LIMIT(3),LIMIT(1),
1LIMIT(8),LIMIT(9),RMA,XS,IWS,SS,TANGN,NTSP,NTCP,NTCV,PV)
CALL PRINT4(MSP,CHI,RNU,NSP,TRPA,NUMCEL,PDN,WTH,DBA,NBS,TMPA,XVA,
1YVA,ZVA,1,NBT,XC,YC,ZC,LEV,LKW)
GO TO 353
350 CONTINUE
CALL PRINT4(MSP,CHI,RNU,NSP,TRP,NUMCEL,FDN,WTH,DB,NB,TMP,XV,YV,ZV,
10,NBP,XC,YC,ZC,LEV,LKW)
353 IF(DEBUG(2)) WRITE(6,1)
IF((TIME.LT.TLIM).AND.(CPUTM.GT.CPJ)) GO TO 280
IF(IC.EQ.ICOPY) RETURN
IC=IC+1
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) (((I,J,L),(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)
WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)
WRITE(6,2)
WRITE(6,3) IC
CALL PRINTA(THETAZ,NWEDGE,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,LD,LP,
1XLIM,COEPP,LIMIT,MSP)
CALL PRINTB(PNA,MSP,FNB,LEV,LWP,NS,RLD,XLIM,IC,YC,ZC,NB,NUMCEL,LKW
1,MSP)
SAVE=.FALSE.
GO TO 345
360 WRITE(6,38) (DBG1(I,LARGE),I=1,3)
IF(REDO) GO TO 364
IF(DUMP) CALL ABEND(9)
STOP
364 CONTINUE
IF(NEW) GO TO 365
READ(9) DENP,U,XREF,TRP,      KAWLS,NL,NW,NH,MW,MH,LW,LE,NXA,NIB,NCA
1      ,NCB,NFA,NPB,NHA,NHB,SW,BH,BWB,BHB,BWC,BHC,XLB,XLC,PI,NREG
2      ,S,SINANG,COSANG,AKN,NBI,RM,IR,ND,TIME,DTH,TI,ITS,ITP,TST
3      ,TLIM,RMA,RNU,DIR,XSTART,JNM,MNE,MNB,TR,BZC,CN7,DRP,PCP
4      ,PNA,HTP,INH,ITYPE,JTYPE,MJ,NAV,NHAX,NS,NWEDG,PRT,SAMP
5      ,BTA,C1,C2,C3,C7,C8,DAH,DPA,PL,DELANG,FDN,HTI,HTE,JNT,KHM
6      ,NE,WTH,C4,VRE,NCOL,LD,LP,LWP,RLD,CTI,CTR,CNI,CNR,LEV,SM
7      ,ST,D1,D2,D3,D4,SSA,SSB,SS,NSP,NMB,NMC,NMP,NPB,NRAN,VZLR

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RUN7040  
RUN7050

RUN6770

RUN6860

RUN7080

RUN7120

RUN7130

RUN7140

RUN7180

RUN7190

RUN7200

RUN7220

RUN7230

```

8      ,IWS,TANGH,XLIH,COEPP,ICB,XS,YCB,TB,ALPHA,SIGMA,NTS,NTSP
9      ,UTL,UTT,VTS,HTS,HTSI,ENT,REM,ENTS,REMS,PTH,THETA,DTH,THPA
A      ,DBA,NB,NBF,NBT,THP,XV,YVA,YV,YVA,ZV,ZVA,T,DB,PNB,IC,YC,ZC
B      ,NUMCEL,PAU,PAV,PAW,PAY,PAZ,PV,HTCV,NTCP,LPP,LCOL,LM
C,ETA,PHI,CHI,CH,CM,CNG,CMG,CNS,TRP,TRPA,THETAZ,NWEDGE,HSP,ANGLE,TF
D,UTLI,UTTI,VTSI,ER,RMB,LKW,NBS,LB,NBM,NBN

```

REWIND 9

365 JNM=9\*JNM/10

ANN=INN

INN=9\*ANN/10

DDN=.9\*DDN

DRF=DRF/.9

PCF=PCF/.9

HTF=HTF/.9

DO 370 MM=1,MSP

FDN(MM)=FDN(MM)\*INN/ANN

DO 366 KK=1,MSP

366 CNB(KK,MM)=CNB(KK,MM)\*.9

DO 370 LT=1,NWEDG

ENTS(MM,LT)=ENTS(MM,LT)\*INN/ANN

REMS(MM,LT)=0.0

DO 370 NK=1,2

DO 370 NJ=1,6

ENT(NK,MM,NJ,LT)=ENT(NK,MM,NJ,LT)\*INN/ANN

370 REM(NK,MM,NJ,LT)=0.0

IF(NEW) GO TO 220

TST=TIME+TST

TI=-1.

PRT=ITP

WRITE(6,2)

WRITE(6,4)

WRITE(6,5) (((I,J,L,(ENT(I,J,K,L),K=1,6),L=1,NWEDG),J=1,MSP),I=1,2)

WRITE(6,6) ((J,L,ENTS(J,L),PTH(J,L),L=1,NWEDG),J=1,MSP)

WRITE(6,2)

IF((LARGE.EQ.2).OR.(LARGE.EQ.3)) GO TO 280

REDO=.FALSE.

GO TO 360

END

RUN7260

RUN7410

SUBROUTINE DIAG(N,ITEST,NUM)

REAL\*8 PARAM(10) /' NWEDGE','

NREG','

ND','

NPX','

1 MNH','

MNB','

NBX','

NS','

MJ','

MSP' /

DIAG010

DIAG040

DIAG050

DIAG060

DIAG070

DIAG080

DIAG090

DIAG100

DIAG110

DIAG120

DIAG130

DIAG140

DIAG150

DIAG160

---

 FORMATS

32 FORMAT(9X,'ENT,REM,ENTS,REMS,PTH,THETA,DTH')

42 FORMAT(///5X,43H ARRAY DIMENSIONS ARE ABOUT TO BE VIOLATED./)

44 FORMAT(5X,18H MAXIMUM VALUE IS 15,20H, WHEREAS YOU INPUT 15,3H

1A8,1H)

56 FORMAT(/5X,78H IF YOU DESIRE TO USE THIS VALUE, THE FOLLOWING ARRADIAG130

1YS MUST BE RE-DIMENSIONED./)

62 FORMAT(9X,'HTS,HTSI,NTS,NTSP,UTL,UTT,VTS')

64 FORMAT(9X,'XLIH,COEPP'//11X,'NOTE THAT THE ILM ARRAY MUST BE DIMEDIAG160

```

1 INSIONED TO 3 MORE THAN THE COEFF ARRAY.')
66 FORMAT (9X, 'XCB, XS, YCB, TB, ALPHA, SIGMA')
68 FORMAT (9X, 'DBA, NB, NBP, NBT, TMP, TMPA, XV, IVA, YV, YVA, ZV, ZVA, T, DB')
70 FORMAT (9X, 'LPF, PAU, PAV, PAW, PAX, PAY, PAZ, LCOL')
72 FORMAT (8X, 3H LM)
74 FORMAT (//5X, 76H IF YOU CHANGE THE ARRAY DIMENSIONS, ALSO CHANGE THDIAG220
1E 'LIMIT' DATA STATEMENT.)
75 FORMAT (9X, 'ALL ARRAYS ASSOCIATED WITH SPECIES')
76 FORMAT (9X, 'PNE, XC, YC, ZC, NUMCEL')
78 FORMAT (9X, 'PV, NTCV, NTCF, MS, IWS, SL, DELS, TANGN')
80 FORMAT (9X, 'PV, NTCV')

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```

WRITE (6, 42)
WRITE (6, 44) ITEST, NUM, PARAM (N)
WRITE (6, 56)
GO TO (1, 2, 3, 4, 5, 6, 7, 8, 9, 10), N
1 WRITE (6, 62)
WRITE (6, 32)
GO TO 11
2 WRITE (6, 64)
GO TO 11
3 WRITE (6, 66)
WRITE (6, 62)
GO TO 11
4 WRITE (6, 68)
GO TO 11
5 WRITE (6, 70)
GO TO 11
6 WRITE (6, 72)
GO TO 11
7 WRITE (6, 76)
GO TO 11
8 WRITE (6, 78)
GO TO 11
9 WRITE (6, 80)
GO TO 11
10 WRITE (6, 75)
11 WRITE (6, 74)
STOP
END

```

```

SUBROUTINE PRINTA (THETAZ, NWEDGE, TITLE, NAME, XCB, YCB, TB, ALPHA, SIGMA, PRA0010
1 LD, LP, XLIM, COEFF, LIMIT, MSP)
INTEGER TST, TLIM, TIME
LOGICAL SAVE, NEW
DIMENSION NWEDGE (2), LIMIT (1), TITLE (6), NAME (2), XCB (1), YCB (1), TB (1)
DIMENSION ALPHA (3, 1), SIGMA (3, 1), LD (1), LP (1), XLIM (1), COEFF (4, 1)
DIMENSION RNU (3), RMA (3), CHI (3), DIR (3, 3), PHI (3, 3), ETA (3, 3)
DIMENSION WTM (3), DAM (3, 3), VELS (3), XSP (3)
COMMON /FIRST/NL, NW, NH, HW, HE, LW, LH, NIA, NIB, NCA, NCB, NPA, NPB, NHA, NEB
COMMON /SECND/BW, BE
COMMON /THIRD/PI, NREG, S, SINANG, COSANG, AKN
COMMON /FIFTH/ND, TIEZ, DTM, TI, ITS, ITP, TST, TLIM, RMA, RNU, DIR

```

COMMON /SIXTH/RMB, XSTART, JNH, MNH, MNB, NEW, SAVE, PERCNT, NSR, TR PRA0110  
 COMMON/EIGHT/DENF, U, TP, ANGLE, TRF, CHI, PHI, ETA, WTM, DAM, VELR, XREF PRA0120  
 DATA NOT/'NOT '/ PRA0130

## FORMATS

PRA0140  
 PRA0150  
 PRA0160  
 PRA0170

1 FORMAT(16X,40('-'),T74,'I'//9X,'3-D',I2,'-FLUID PROGRAM - ')  
 2 FORMAT('+',31X,A4) PRA0190  
 3 FORMAT('+',35X, 'A RESTART OF A PREVIOUS RUN',T74,'I'/12X,2 PRA0200  
 1A4,' - ',6A4,' - ',I2,' REGIONS',T74,'I',16(/T74,'I')) PRA0210  
 4 FORMAT(7X,'FRONT OF BODY =',E12.4,' XSTART MAX HEIGHT =',E12.4,'  
 1RMB',T74,'I'/7X,'X-LIMIT',T37,'BODY COEFFICIENTS',T74,'I')  
 6 FORMAT(5F14.6,3X,'I') PRA0240  
 10 FORMAT(1X,72('-')) PRA0250  
 12 FORMAT(/14X,'PARAMETERS OF SEGMENTS FOR BODY COLLISIONS',T96,'I'/  
 18X,'Y-COORD. TEMP. ALPHA1 ALPHA2 ALPHA3 SIGMA1 SIGMA2  
 2 SIGMA3 AREAS',T96,'I')  
 14 FORMAT(4X,E12.4,7F9.4,E12.4,T96,'I')  
 16 FORMAT(23X,'WEIGHTING FACTORS'/1X,10I6,T96,'I')  
 17 FORMAT(///25X,'ARRAY STORAGE USED'/5X,I6,' \*\*',10I6,T96,'I')  
 18 FORMAT(1H1/17X,'LENGTH OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BW'  
 A,T76,'I'  
 1/17X,'HEIGHT OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BH',T76,'I'  
 2/16X,'NUMBER OF L1 CELLS ALONG FLOW AXIS =',I13,' NW',T76,'I'  
 3/17X,'NUMBER OF L1 CELLS IN RADIAL DIR. =',I13,' NH',T76,'I'  
 4/21X,'NUMBER OF LEVELS OF CELL SIZE =',I13,' NL',T76,'I')  
 20 FORMAT(11X,'NUMBER OF L1 CELLS IN FRONT OF L2 CELLS =',I13,' NFA'  
 1,T76,'I'/15X,'NUMBER OF AXIAL SUBDIVIDED L1 CELLS =',I13,' NCA',T  
 276,'I'/14X,'NUMBER OF RADIAL SUBDIVIDED L1 CELLS =',I13,' NHA',T7  
 36,'I'/16X,'NUMBER OF L2 CELLS ALONG FLOW AXIS =',I13,' MW',T76,'I'  
 4'/17X,'NUMBER OF L2 CELLS IN RADIAL DIR. =',I13,' MH',T76,'I')  
 22 FORMAT(11X,'NUMBER OF L2 CELLS IN FRONT OF L3 CELLS =',I13,' NFB'  
 1,T76,'I'/15X,'NUMBER OF AXIAL SUBDIVIDED L2 CELLS =',I13,' NCB',T  
 276,'I'/14X,'NUMBER OF RADIAL SUBDIVIDED L2 CELLS =',I13,' NHB',T7  
 36,'I'/16X,'NUMBER OF L3 CELLS ALONG FLOW AXIS =',I13,' LW',T76,'I'  
 4'/17X,'NUMBER OF L3 CELLS IN RADIAL DIR. =',I13,' LH',T76,'I')  
 23 FORMAT(3X,'NUMBER OF AZIMUTHAL WEDGES IN LOWER ',I3,' DEGREES =',IPRA0470  
 113,' NWEDGE1 I'/3X,'NUMBER OF AZIMUTHAL WEDGES IN UPPER ',I3,' DE  
 2GREES =',I13,' NWEDGE2 I'/)  
 24 FORMAT(16X,'BASIC TIME INTERVAL FOR COLLISIONS =',E13.4,' DTM  
 1 I'/8X,'TIME INTERVAL FOR SAMPLING FLOW FIELD INFO =',E13.4,' DTS  
 2 I'/24X,'TIME INTERVAL FOR PRINTING =',E13.4,' DTP I'/9X,  
 3'TIME TO STEADY-STATE CONDITIONS (ASSUMED) =',E13.4,' TST I'/  
 419X,'TIME AT WHICH RUN IS TERMINATED =',E13.4,' TTM I'/)  
 26 FORMAT(9X,'INITIAL NUMBER OF MOLECULES - EITHER TYPE =',I13,' INH  
 1 I'/9X,'MAXIMUM NUMBER OF MOLECULES - EITHER TYPE =',I13,' MN  
 2H I'/1X,'MAX NUMBER OF MOLECULES IN ANY CELL - EITHER TYPE =',PRA0570.  
 3I13,' MNB I')  
 27 FORMAT(/22X,'VELOCITY OF FREE STREAM FLOW =',E13.4,' U',T76,'I'/1  
 19X,'SPEED RATIO OF FREE STREAM FLOW =',E13.4,' S',T76,'I'/19X,'MAC  
 AH NUMBER OF FREE STREAM FLOW =',E13.4,' N',T76,'I'/19X,'SPECIFIC H  
 BEAT RATIO (CALCULATED) =',E13.4,' GAMMA',T76,'I'/ 35X,'ANG  
 2LE OF ATTACK =',F13.4,' ANGLE I'/16X,'NUMBER DENSITY OF FREE ST

```

3REAN FLOW =',E13.4,' N',T76,'I'/19X,'TEMPERATURE OF FREE STREAM FL
4OW =',F13.4,' TP',T76,'I'/16X,'MOLE FRACTIONS OF FREE STREAM FLOW
5=' ,3E13.4,' RNU I'/16X,'MOLECULAR WEIGHTS OF SPECIES ABOVE =',3F13
6.4,' RMA I')
28 FORMAT(/10X,'REFERENCE TEMPERATURE FOR MOLECULAR DATA =',F12.4,'
1TRP',T90,'I'/14X,'CROSS-SECTION',26X,'TEMP EXPONENT',T90,'I'/3(3X,
23E12.4,3X,3F12.6,T90,'I'/) / 5X,'CHI/2-1',11X,'ROTATIONAL PARAMETER
3 PHI',T90,'I'/3(F12.4,5X,3F12.6,T90,'I'/)
29 FORMAT(9X,'DATA SAVED ON TAPE 9')
30 FORMAT(/31X,'REF MOLECULAR SPEED =',E13.4,' VELR',T76,'I'/20X,'SP
1ECIES FREE STREAM MOLECULAR SPEEDS',T76,'I'/14X,3E16.6,T76,'I'/26X
2,'REFERENCE MEAN FREE PATH =',E13.4,' KREF',T76,'I'/26X,'SPECIES N
3EAN FREE PATHS',T76,'I'/14X,3E16.6,T76,'I'/11X,'LONGITUDINAL KNUDS
4EN NUMBER (CALCULATED) =',E13.4,' AKN',T76,'I'/13X,'TRANSVERSE KNUD
5SEN NUMBER (CALCULATED) =',E13.4,' AKT',T76,'I')

```

PRA0670

PRA0680

PRA0690

PRA0700

PRA0710

PRA0720

PRA0730

```

IARRAY=708+LIMIT(3)*(32+56*LIMIT(1))+20*LIMIT(2)+LIMIT(4)*(120+4*1
LIMIT(6))+56*LIMIT(5)+LIMIT(8)*(68+96*LIMIT(9))+20*LIMIT(7)+224*LI
2IT(1)
WRITE(6,1) MSP
IF(NEW) WRITE(6,2) NOT
WRITE(6,3) NAME,TITLE,NREG
WRITE(6,4) XSTART,RMB
DO 100 I=1,NREG
100 WRITE(6,6) XLIM(I+2),(COEFF(J,I),J=1,4)
WRITE(6,10)
WRITE(6,12)
DO 110 I=1,ND
110 WRITE(6,14) YCB(I),TB(I),(ALPHA(J,I),J=1,3),(SIGMA(J,I),J=1,3),YCB
1(I)
WRITE(6,10)
WRITE(6,16) (LD(N),LP(N),N=1,5)
WRITE(6,17) IARRAY,(LIMIT(I),I=1,10)
WRITE(6,18) BW,BH,NW,NE,EL
IF(NL.GT.1) WRITE(6,20) NPA,NCA,NHA,MW,MH
IF(NL.GT.2) WRITE(6,22) NPB,NCB,NHB,LW,LE
IETAZ=THETAZ
JETAZ=180-IETAZ
WRITE(6,23) IETAZ,NWEDGE(1),JETAZ,NWEDGE(2)
DTS=DTM*ITS
DTP=DTM*ITP
AST=DTM*TST
ALIM=DTM*TLIM
CHT=0.0
DO 120 J=1,MSP
120 CHT=CHT+CHI(J)*RNU(J)
GAMMA=(7.+2.*CHT)/(5.+2.*CHT)
AM=S*SQRT(2./GAMMA)
WRITE(6,24) DTM,DTS,DTP,AST,ALIM
WRITE(6,26) JNH,MNH,MNB
WRITE(6,27) U,S,AM,GAMMA,ANGLE,DENP,TP,(RNU(I),I=1,3),(RMA(I),I=1,3
1)
WRITE(6,28) TRP,((DIR(I,K),K=1,3),(ETA(I,K),K=1,3),I=1,3),(CHI(I),

```

PRA0750

PRA0760

PEA0770

PRA0780

PRA0790

PRA0800

PRA0810

PRA0820

PRA0850

PRA0860

PRA0880

PRA0890

PEA0900

PRA0910

PRA0920

PRA0930

PRA0940

PRA0950

PRA0960

PRA0970

PRA0980



```

1 (PHI (I,K),K=1,3),I=1,3)
  DO 210 I=1,3
  VELS (I)=0.0
210 XSP (I)=0.0
  DO 220 J=1,MSP
  VELS (J)=VELR/SQRT (WTH (J))
  XT=0.0
  DO 215 M=1,MSP
215 XT=XT+RNU (M)*DAM (J,M)*SQRT (1.+WTH (J)/WTH (M))
220 XSP (J)=1.414214*XREP/XT
  AKT=1./RMB
  WRITE (6,30) VELR, (VELS (I),I=1,3),XREP, (XSP (I),I=1,3),AKN,AKT
  IF (SAVE) WRITE (6,29)
  RETURN
  END

```

PRA1040  
PRA1050  
PRA1060

```

SUBROUTINE PRINTB (PNA,MSP,PNB,LEV,LWP,NH,RLD,XLIM,XC,YC,ZC,NB,
1 NUMCEL,LKW,N)
  INTEGER*2 LKW,NB,NUMCEL
  DIMENSION PNB (1),LEV (1),LWP (1),NH (1),RLD (1),XLIM (1),XC (1)
  DIMENSION YC (1),ZC (1),NB (N,1),NUMCEL (1),LKW (1)
  COMMON /FIRST/HL
  COMMON /THIRD/PI,NREG,S,SINANG,COSANG,AKN
  COMMON /FOUR/NBX

```

PRB0050  
PRB0060  
PRB0070  
PRB0080

```

1 FORMAT (1H1)
2 PCMAT (2X,'-----CELL GEOMETRY-----'
1-----'/2X,'BOX LEVEL' POSITION OF CENTER VOLUME
2EIGHTING POPULATION'/2X,'NUM.',12X,'X',7X,'Y THETA',12X,'FACTO
3R ',15X,' CELL#')
3 FORMAT (1X,I4,I5,3X,2P8.3,F7.1,E12.3,2X,I2,4X,3I5,3X,I4)
4 PCMAT (2X,'-----TOTALS-----',E12.4,8X,3I5)

```

PRB0090  
PRB0100  
PRB0110

```

WRITE (6,1)
WRITE (6,2)
DO 200 I=1,NBX
  IF (NUMCEL (I).EQ.0) GO TO 200
  X=(XC (I)-XLIM (2))*AKN
  Y=YC (I)*AKN
  LEVEL=1
  IF (NL.LT.2) GO TO 120
  IF (I.LT.LEV (1)) GO TO 120
  LEVEL=2
  IF (NL.LT.3) GO TO 120
  IF (I.LT.LEV (2)) GO TO 120
  LEVEL=3
120 CONTINUE
  J=NUMCEL (I)
  M1=NB (1,J)
  M2=0
  IF (MSP.GE.2) M2=NB (2,J)
  M3=0
  IF (MSP.GE.3) M3=NB (3,J)
140 WRITE (6,3) I,LEVEL,X,Y,ZC (I),PNB (I),LKW (J),M1,M2,M3,J
200 CONTINUE
  NM2=0
  IF (MSP.GE.2) NM2=NM (2)

```

PRB0150  
PRB0170  
PRB0180  
PRB0190  
PRB0200  
PRB0210  
PRB0220  
PRB0230  
PRB0240  
PRB0250  
PRB0260  
PRB0320  
PRB0370

```

NM3=0
IF (NHP.GE.3) NM3=NM(3)
WRITE (6,4) PNA, NM(1), NM2, NM3
RETURN
END

```

PRB0390  
PRB0400

```

SUBROUTINE SIMPSN(A,B,L,INTGRL,PERCNT,COEPP,PIROOT,SUM1,PUN)
REAL INTGRL
DIMENSION COEPP(4,1)

```

SIMP010  
SIMP020  
SIMP030  
SIMP040

-----

THE PURPOSE OF THIS SUBROUTINE IS TO PERFORM A SIMPSON'S RULE  
INTEGRATION.

SIMP050  
SIMP060  
SIMP070  
SIMP080

FORMAT

SIMP090  
SIMP100  
SIMP110

4 FORMAT(/30H TOO MANY ITERATIONS. TEST IS E15.7,14H, INTEGRAL IS E15.7,14H IF THE RANGE E15.7,4H TO E15.7)

SIMP120  
SIMP130  
SIMP140

```

PREV=0.0
N=4
N1=3
G=.5*(B-A)+A
SUME=FUR(G,PIROOT,L,COEPP)
SUMO=0.0
K=0

```

SIMP150  
SIMP160  
SIMP170

235 SUME=SUME+SUMO

SIMP180

N1=N1+K

SIMP190

K=N1+1

SIMP200

IF (K.LT.5000) GO TO 237

SIMP210

WRITE (6,4) TEST,PREV,A,B

SIMP220

STOP

SIMP230

237 SUMO=0.0

SIMP240

DO 240 I=1,N1,2

SIMP250

G=(I\*(B-A))/K+A

SIMP260

240 SUMO=SUMO+PUN(G,PIROOT,L,COEPP)

SIMP270

INTGRL=(SUM1+4.\*SUMO+2.\*SUME)/(3.\*K)

SIMP280

TEST=ABS(2.-4.\*PREV/(INTGRL+PREV))

SIMP290

PREV=INTGRL

SIMP300

IF (TEST.GT.PERCNT) GO TO 235

SIMP310

RETURN

SIMP320

END

SIMP330

SIMP340

SIMP350

SIMP360

SIMP370

```

FUNCTION FNCTN(X,PIROOT,L,COEPP)
DIMENSION COEPP(4,1)

```

FNCTN010

FNCTN020

FNCTN030

-----

THE PURPOSE OF THIS FUNCTION IS TO EVALUATE THE INTEGRAND USED IN  
THE SIMPSON-S RULE INTEGRATION ROUTINE.

FNCTN040

FNCTN050

FNCTN060

A=COEPP(1,L)

FNCTN070

B=COEPP(2,L)

FNCTN080

C=COEPP(3,L)

FNCTN090

D=COEPP(4,L)

FNCTN100

AA=4.\*(A-B)\*(A\*X+C)\*X+C\*C-4.\*B\*D

FNCTN110

```

IF (AA.LT.0.) AA=0.
FNCTN=SQRT (AA)
RETURN
END
FNCN120
FNCN130
FNCN140
FNCN150

SUBROUTINE HEIGHT (X,R,L,COEFF,I)
DIMENSION COEFF (4,1),DBG2 (2,3)
DATA DBG2/'SBTR','CT ','GAS ',' ','RUN',' '/
-----
THE PURPOSE OF THIS SUBROUTINE IS TO COMPUTE THE R-COORDINATE OF
THE BODY CURVE AT ANY GIVEN X-COORDINATE, IN THE X-R PLANE.
-----
2 FORMAT (/' MESSAGE ',I2,E17.8,3X,2A4)
ARG=- ((COEFF (1,L)*X+COEFF (3,L)) *I+COEFF (4,L))/COEFF (2,L)
IF (ARG.GE.0.) GO TO 100
WRITE (6,2) L,X, (DBG2 (M,I),M=1,2)
ARG=0.
100 R=SQRT (ARG)
RETURN
END
HGHT010
HGHT020
HGHT030
HGHT040
HGHT050
HGHT060
HGHT070
HGHT080
HGHT090
HGHT100
HGHT110
HGHT120
HGHT130
HGHT140
HGHT150

FUNCTION FNCTM (ARG,PIROOT,L,COEFF)
DIMENSION COEFF (4,1)
D=0.
IF (ABS (ARG).LT.10.) D=EXP (-ARG*ARG)/PIROOT
E=0.
IF (ARG.GT.-10.) E=ARG*ERRF (ARG)
TEM=-.5*(D+E)
FNCTM=TEM
RETURN
END
FNCM010
FNCM020
FNCM030
FNCM040
FNCM050
FNCM060
FNCM070
FNCM080
FNCM090
FNCM100

SUBROUTINE CELL (TH,LEV,A,B,K,KH,XO,I,J,DELANG,NWEDGE,XC,YC,ZC,PNB)
DIMENSION DELANG (1),NWEDGE (1),XC (1),YC (1),ZC (1),PNB (1)
COMMON /THIRD/PI
-----
THE PURPOSE OF THIS SUBROUTINE IS TO
1. COMPUTE THE VOLUME OF EACH CELL (ALL 3 POSSIBLE LEVELS)
AND STORE THE RESULT IN THE ARRAY CALLED 'PNB'.
2. COMPUTE THE X, R, AND THETA COORDINATES OF THE CENTER OF
EACH CELL (ALL 3 POSSIBLE LEVELS) AND STORE THE RESULTS IN
ARRAYS CALLED 'XC', 'YC', AND 'ZC'.
-----
INDEX=I+J
ZO=0.
DO 120 MT=1,LEV
ICNT=NWEDGE (MT)
ANGLE=DELANG (MT)
FACTOR=ANGLE/180.*PI*B*B*A
Z=ZO-.5*ANGLE
DO 110 L=1,ICNT
Z=Z+ANGLE
Y=-.5*B
CELL010
CELL020
CELL030
CELL040
CELL050
CELL060
CELL070
CELL080
CELL090
CELL100
CELL110
CELL120
CELL130
CELL140
CELL150
CELL160
CELL170
CELL180
CELL190
CELL200
CELL210
CELL220
CELL230

```

DO 110 N=1,KH	CELL240
X=YO-.5*A	CELL250
Y=Y+B	CELL260
DO 110 N=1,K	CELL270
X=X+A	CELL280
INDEX=INDEX+1	CELL290
KC(INDEX)=X	CELL300
YC(INDEX)=Y	CELL310
ZC(INDEX)=Z	CELL320
110 PNB(INDEX)=FACTOR*(2*N-1)	CELL330
ZO=TH	CELL340
120 CONTINUE	CELL350
RETURN	CELL360
END	CELL370
-----	
SUBROUTINE ZERO(NWIDE, NHI, NBEG, NLONG, NUP, NAREA, ICNT, PNB)	ZERO010
DIMENSION PNB(1)	ZERO020
-----	
THIS SUBROUTINE SETS THE SIZES TO ZERO OF THOSE CELLS WHICH ARE	ZERO030
TO BE SUBDIVIDED INTO SMALLER CELLS.	ZERO040
-----	
NGO=NBEG+1	ZERO050
NSTOP=NBEG+NLONG	ZERO060
DO 110 N=NGO, NSTOP	ZERO070
DO 110 M=1, NUP	ZERO080
DO 110 L=1, ICNT	ZERO090
INDEX=NWIDE*(NHI*(L-1)+M-1)+N+NAREA	ZERO100
110 PNB(INDEX)=0.0	ZERO110
RETURN	ZERO120
END	ZERO130
-----	
SUBROUTINE SBTRCT(NGO, NTEMP, NSTOP, BWIDTH, BHITE, DELANG, XC, YC, PNB, XLSBCT	SBCT010
1IM, COEPP)	SBCT020
DIMENSION DELANG(1), XC(1), YC(1), PNB(1), XLIM(1), COEPP(4, 1)	SBCT030
COMMON /THIRD/PI, NREG	SBCT040
-----	
THIS SUBROUTINE SUBTRACTS FROM EACH CELL SIZE ('PNB' ARRAY) THAT	SBCT050
PORTION OCCUPIED BY THE BODY.	SBCT060
-----	
FACTOR=PI*DELANG(1)/180.	SBCT070
DO 150 N=NGO, NSTOP	SBCT080
IF(PNB(N).LE.0.) GO TO 150	SBCT090
DFNB=.005*PNB(N)	SBCT100
SLICE=.01*BWIDTH	SBCT110
IF(N.GT.NTEMP) FACTOR=PI*DELANG(2)/180.	SBCT120
X=XC(N)-.5*(BWIDTH+SLICE)	SBCT130
YBOT=YC(N)-.5*BHITE	SBCT140
YTOP=YBOT+BHITE	SBCT150
DO 130 S=1, 100	SBCT160
X=X+SLICE	SBCT170
IF(X.LE.XLIM(2)) GO TO 130	SBCT180
	SBCT190
	SBCT200
	SBCT210
	SBCT220

```

DO 120 L=1,NREG
IF (X.LT.XLIM(L+2)) GO TO 125
120 CONTINUE
GO TO 130
125 CALL HEIGHT(X,YBODY,L,COEFF,1)
IF (YBODY.LE.YBOT) GO TO 130
YTEMP=YTOP
IF (YBODY.LT.YTOP) YTEMP=YBODY
FNB(N)=FNB(N)-SLICE*(YTEMP*YTEMP-YBOT*YBOT)*FACTOR
130 CONTINUE
IF (FNB(N)-LT.DFNB) FNB(N)=0.
150 CONTINUE
RETURN
END

```

SBCT230  
 SBCT240  
 SBCT250  
 SBCT260  
 SBCT270  
 SBCT280  
 SBCT290  
 SBCT300  
 SBCT310  
 SBCT320  
 SBCT330  
 SBCT340  
 SBCT350  
 SBCT360

```

SUBROUTINE IMPACT(RM,G1,G2,G3,ET,EI,PHI,CHI,ETA,XM,CIM)
COMMON/THIRD/PI
IF (PHI.EQ.0.) GO TO 20
IF (CHI.EQ.0.) GO TO 20
DF=PHI*CHI-1.
DS=PHI*(2.-.5*ETA)-1.
E=ET+EI
10 X=RAND(0)
IF (X.EQ.0.0) GO TO 10
XT=X**DF*(1.-X)**DS
IF (XT.GT.XM) GO TO 15
CIM=CIM+XT
IF (CIM.LT.XM) GO TO 10
CIM=CIM-XM
15 ET=(1.-PHI)*ET+(1.-X)*PHI*E
EI=(1.-PHI)*EI+X*PHI*E
20 GP=SQRT(ET/RM)
EP=2.*PI*RAND(0)
CSX=2.*RAND(0)-1.
SSX=SQRT(1.-CSX**2)
G1=GP*CSX
G2=GP*SSX*COS(EP)
G3=GP*SSX*SIN(EP)
RETURN
END

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SUBROUTINE GAS(NWEDG,THETAZ,DELANG,NWEDGE,BTA,C1,DPA,NM,RLD,LWF,PGAS0010
1B,DB,NB,NBF,LPP,PAU,PAV,PAW,PAX,PAY,PAZ,XLIM,COEFF,LM,I2,I3,LARGE,GAS0020
2MNE,ENB,DEBUG1,LCOL,NUMCEL,IP,ER,CHI,CNG,CMG,I,LB,NBN,NBF)
INTEGER*2 LM(I,1),LPP(I,1),LCOL(I,1),LB(1),NBN(I,1),NBN(1)
INTEGER*2 NB,NBF,NUMCEL
INTEGER Q
LOGICAL DUMP,DEBUG1
DIMENSION DELANG(2),NWEDGE(2),NUMCEL(1)
DIMENSION BTA(1),C1(1),DPA(1),NM(1),RLD(1),LWF(1),FNB(1),CHI(1)
DIMENSION DB(I,1),NB(I,1),NBF(I,1),PAU(I,1),PAV(I,1),PAW(I,1)
DIMENSION PAX(I,1),PAY(I,1),PAZ(I,1),ER(I,1),COEFF(4,1),XLIM(1)
DIMENSION CNG(1),CMG(1)
COMMON /FIRST/NL,NW,NH,NW,MB,LB,LH,NXA,NXB
COMMON /SECND/BW,BH,BWB,BHB,BWC,BHC,XLB,XLC

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GAS0050  
 GAS0060  
 GAS0070  
  
 GAS0110  
 GAS0120

\*\*\*\*\* / TIME / P1, P2, P3, P4, P5, P6, P7, P8, P9, P10, P11, P12, P13, P14, P15, P16, P17, P18, P19, P20, P21, P22, P23, P24, P25, P26, P27, P28, P29, P30, P31, P32, P33, P34, P35, P36, P37, P38, P39, P40, P41, P42, P43, P44, P45, P46, P47, P48, P49, P50, P51, P52, P53, P54, P55, P56, P57, P58, P59, P60, P61, P62, P63, P64, P65, P66, P67, P68, P69, P70, P71, P72, P73, P74, P75, P76, P77, P78, P79, P80, P81, P82, P83, P84, P85, P86, P87, P88, P89, P90, P91, P92, P93, P94, P95, P96, P97, P98, P99, P100, P101, P102, P103, P104, P105, P106, P107, P108, P109, P110, P111, P112, P113, P114, P115, P116, P117, P118, P119, P120, P121, P122, P123, P124, P125, P126, P127, P128, P129, P130, P131, P132, P133, P134, P135, P136, P137, P138, P139, P140, P141, P142, P143, P144, P145, P146, P147, P148, P149, P150, P151, P152, P153, P154, P155, P156, P157, P158, P159, P160, P161, P162, P163, P164, P165, P166, P167, P168, P169, P170, P171, P172, P173, P174, P175, P176, P177, P178, P179, P180, P181, 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C9=C9+R/A	GAS0580
IF (C9.LT.RWPM) GO TO 140	GAS0590
C9=C9-RWPM	GAS0600
IF (X.LE.XLIM(2)) GO TO 159	GAS0610
DO 152 L=1,NREG	GAS0620
IF (X.LT.XLIM(L+2)) GO TO 154	GAS0630
152 CONTINUE	GAS0640
GO TO 159	GAS0650
154 CALL HEIGHT(X,YBODY,L,COEFF,2)	GAS0660
IF (R.LT.YBODY) GO TO 130	GAS0670
159 PAX(MT,N)=X	GAS0680
D=PI*RAND(0)	GAS0690
PAY(MT,N)=R*COS(D)	GAS0700
PAZ(MT,N)=R*SIN(D)	GAS0710
TANG=180.*(1.-D/PI)	GAS0720
IWDGE=TANG/DELANG(1)	GAS0730
IF ((IWDGE.GE.NWEDGE(1)).AND.(DELANG(2).NE.0.)) IWDGE=(TANG-THETAZ)	
1/DELANG(2)+NWEDGE(1)	GAS0750
IF (IWDGE.GE.NWEDG) IWDGE=NWEDG-1	GAS0760
L=X/BW+1.	GAS0770
IF (L.GT.NW) L=NW	GAS0780
M=R/BH	GAS0790
IF (M.GE.NH) M=NH-1	GAS0800
K=(IWDGE*NH+M)*NW+L	GAS0810
IF (K.LE.NXA) GO TO 160	GAS0820
WRITE(6,2) L,M,K,MT,N,DELANG,NWEDGE,NWEDG,D,TANG,IWDGE,NH,NW	GAS0830
IF (DUMP) CALL ABEND(11)	GAS0840
STOP	GAS0850
160 IF (NL.EQ.1) GO TO 162	
IF (IWDGE.GE.NWEDGE(1)) GO TO 162	GAS0870
IF (PNB(K).GT.0.) GO TO 162	GAS0880
L=(X-XLB)/BWB+1.	GAS0890
IF (L.GT.MW) L=MW	GAS0900
M=R/BHB	GAS0910
IF (M.GE.MH) M=MH-1	GAS0920
K=(IWDGE*MH+M)*MW+L+NXA	GAS0930
IF (K.LE.NYA+NXB) GO TO 161	GAS0940
WRITE(6,2) L,M,K,MT,N,DELANG,NWEDGE,NWEDG,D,TANG,IWDGE,MH,MW,NYA	GAS0950
IF (DUMP) CALL ABEND(12)	GAS0960
STOP	GAS0970
161 IF (NL.EQ.2) GO TO 162	GAS0990
L=(X-ILC)/BWC+1.	GAS0980
IF (PNB(K).GT.0.) GO TO 162	GAS1000
IF (L.GT.LW) L=LW	GAS1010
M=R/BHC	GAS1020
IF (M.GE.LH) M=LH-1	GAS1030
K=(IWDGE*LH+M)*LW+L+NYA+NXB	GAS1040
IF (K.LE.NBX) GO TO 164	GAS1050
WRITE(6,2) L,M,K,MT,N,DELANG,NWEDGE,NWEDG,D,TANG,IWDGE,LH,LW,NYA,NXB	GAS1060
1B	GAS1070
IF (DUMP) CALL ABEND(13)	GAS1080
STOP	GAS1090
162 IF (PNB(K).GT.0.) GO TO 164	GAS1100
WRITE(6,3) L,M,K,MT,N,PNB(K)	GAS1110
IF (DUMP) CALL ABEND(14)	

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STOP
164 Q=NUMCEL(K)
IF (Q.GT.0) GO TO 165
WRITE (6,2) Q,L,H,K,H,DELANG,NWEDGE,NWEDG,D,TANG,INDGE,LE,LW,NXA,NYB
1,X,R
IF (DUMP) CALL ABEND(15)
STOP
165 J=NB(MT,Q)+1
KW=LWP(LA)
LPP(MT,N)=KW
LCOL(MT,N)=0
LLC=LLC+KW
IF (J.LE.MNB) GO TO 166
IF (DEBUG1) WRITE (6,4) MT,Q,MNB
GO TO 167
166 NB(MT,Q)=J
LB(N)=Q
NBP(MT,Q)=NBP(MT,Q)+KW
167 IF (LLC.LT.LL) GO TO 110
NM(MT)=N
NBN(MT,1)=0
DO 170 Q=1,NBX
N=NUMCEL(Q)
IF (N.EQ.0) GO TO 170
A=NBP(MT,N)
DB(MT,N)=A*DPA(MT)/PNE(Q)
NBN(MT,N+1)=NBN(MT,N)+NB(MT,N)
NBN(N)=NBN(MT,N)
170 CONTINUE
NG=NS(MT)
DO 175 N=1,NG
Q=LB(N)
NBN(Q)=NBN(Q)+1
NA=NEB(Q)
175 LB(MT,NA)=N
180 CONTINUE
RETURN
190 LARGE=1
RETURN
END

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GAS1120  
GAS1130  
GAS1140  
GAS1150  
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GAS1180  
GAS1190  
GAS1200  
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GAS1420  
GAS1430

SUBROUTINE FLOW(NWEDG,MNB,LARGE,BTA,C1,C7,C8,D1,D2,D3,D4,DTH,NH,SHFLOO010  
1,ST,THETA,LWP,RLD,PTH,ENTS,REMS,SSA,SSB,PAU,PAV,PAW,PAI,PAY,PAZ,LPPFLOO020  
2F,ENT,REN,LCOL,IP,ER,CHI,CNG,CHG,I)

INTEGER\*2 LPP,LCOL
DIMENSION BTA(1),D1(1),D2(1),D3(1),D4(1),DTH(1),NH(1),SW(1),ST(1)
DIMENSION C1(1),C7(1),C8(1),LWP(1),RLD(1),THETA(1),SSA(2,1)
DIMENSION SSB(2,1),PAU(I,1),PAV(I,1),PAW(I,1),PAI(I,1),PAY(I,1)
DIMENSION PAZ(I,1),LPP(I,1),ENT(2,3,6,1),REN(2,3,6,1),LCOL(I,1)
DIMENSION ENTS(3,1),REMS(3,1),PTH(3,1),ER(I,1),CHI(1),CNG(1),
1CHG(1)

COMMON /THIRD/PI
COMMON /FORTH/NBY,RM,XR

FLOO100  
FLOO110  
FLOO120  
FLOO130

THE PURPOSE OF THIS SUBROUTINE IS TO ADD A NEW BATCH OF MOLECULES



TO THE SAMPLE THROUGH THE UPSTREAM BOUNDARY.	FLO0140
-----	FLO0150
DO 370 NT=1,IP	FLO0170
ARG=SN(MT)	FLO0180
XGO=0.	FLO0190
E=1.	FLO0200
DO 180 NT=1,2	FLO0210
SM=AHXI(0.,ARG-4.)	FLO0220
SSM=AHXI(0.,ARG)	FLO0230
TEMPC=0.	FLO0240
DO 170 LA=1,6	FLO0250
TEMPB=RLD(LA)*RLD(LA)	FLO0260
AY = TEMPB-TEMPC	FLO0270
C = TEMPC	FLO0280
TEMPB=TEMPB	FLO0290
DO 170 K=1,NREDG	FLO0300
AM=ENT(MT,MT,LA,K)+REM(MT,MT,LA,K)	FLO0310
H=AM	FLO0320
AMH=H	FLO0330
REM(MT,MT,LA,K)=AM-AMH	FLO0340
IF(H.EQ.0) GO TO 170	FLO0350
DY =AY/AMH	FLO0360
DO 160 N=1,H	
IF(NB(MT).GE.MNH) GO TO 380	FLO0370
NH(MT)=NH(MT)+1	FLO0390
NMX=NH(MT)	FLO0400
R = SQRT(C + DY*(N+RAND(0) - 1.))	FLO0410
D=(THETA(K)+RAND(0)*DTH(K))*PI/180.	FLO0420
PAY(MT,NMX)=R*COS(D)	FLO0430
PAZ(MT,NMX)=R*SIN(D)	FLO0440
LPP(MT,NMX)=LWP(LA)	FLO0450
LCOL(MT,NMX)=0	FLO0460
130 V=SM+RAND(0)*(SSM+4.-SM)	FLO0470
C1(MT)=C1(MT)+2.*V*EXP(SSB(MT,MT)+2.*ARG*V-V*V)/SSA(MT,MT)	FLO0480
IF(C1(MT).LT.1.) GO TO 130	FLO0490
C1(MT)=C1(MT)-1.	FLO0500
PAU(MT,NMX)=E*V/BTA(MT)	FLO0510
140 V=8.*RAND(0)-4.	FLO0520
C7(MT)=C7(MT)+EXP(-V*V)	FLO0530
IF(C7(MT).LT.1.) GO TO 140	FLO0540
C7(MT)=C7(MT)-1.	FLO0550
PAV(MT,NMX)=(V+ST(MT))/BTA(MT)	FLO0560
150 V=8.*RAND(0)-4.	FLO0570
C8(MT)=C8(MT)+EXP(-V*V)	FLO0580
IF(C8(MT).LT.1.) GO TO 150	FLO0590
C8(MT)=C8(MT)-1.	FLO0600
PAW(MT,NMX)=V/BTA(MT)	
ER(MT,NMX)=0.0	
IF(CHI(MT).LE.-1.) GO TO 160	
125 X=9.*RAND(0)	
IF(X.LE.0.0) GO TO 125	
XT=X**CHI(MT)*EXP(-X)	
IF(XT.GE.CMG(MT)) GO TO 126	
CNG(MT)=CNG(MT)+XT	
IF(CNG(MT).LT.CMG(MT)) GO TO 125	

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CNG(MT)=CNG(MT)-CNG(MT)
126 ER(MT,NMX)=X
160 PAX(MT,NMX)=XGO
170 CONTINUE
    ARG=-ARG
    XGO=XR
    E=-1.
180 CONTINUE
    DO 370 K=1,NWEDG
    AM=ENTS(MT,K)+REMS(MT,K)
    S=AM
    AMH=S
    REMS(MT,K)=AM-AMH
    IF(M.EQ.0) GO TO 370
    DX=XR/AMH
    DO 365 N=1,M
    IF(NH(MT).GE.MNH) GO TO 380
    NH(MT)=NH(MT)+1
    NMX=NH(MT)
    PAX(MT,NMX)=(N-1.+RAND(0))*DX
330 TH=(THETA(K)+RAND(0)*DTH(K))*PI/180.
    A=COS(TH)
    B=SIN(TH)
    SM=ST(MT)*A
    C=0.
    IF(ABS(SM).LT.10.) C=EXP(-SM*SM)
    D=0.
    IF(SM.GT.-10.) D=SQRT(PI)*SM*ERRF(SM)
    D1(MT)=D1(MT)+(C+D)/PTH(MT,K)
    IF(D1(MT).LT.1.) GO TO 330
    D1(MT)=D1(MT)-1.
    PAY(MT,NMX)=-RM*A
    PAZ(MT,NMX)=RM*B
    VNH=.5*SM+SQRT(.25*SM*SM+.5)
    VN=AMAX1(0.,SM-4.)
340 V=VN+RAND(0)*(SM+4.-VN)
    D2(MT)=D2(MT)+V*EXP(VNH*(VNH-2.*SM)-V*(V-2.*SM))/VNH
    IF(D2(MT).LT.1.) GO TO 340
    D2(MT)=D2(MT)-1.
    VN=V
350 V=8.*RAND(0)-4.
    D3(MT)=D3(MT)+EXP(-V*V)
    IF(D3(MT).LT.1.) GO TO 350
    D3(MT)=D3(MT)-1.
    VT1=SM(MT)+V
360 V=8.*RAND(0)-4.
    D4(MT)=D4(MT)+EXP(-V*V)
    IF(D4(MT).LT.1.) GO TO 360
    D4(MT)=D4(MT)-1.
    VT2=ST(MT)*B+V
    PAU(MT,NMX)=VT1/BTA(MT)
    PAV(MT,NMX)=(VN*A+VT2*B)/BTA(MT)
    PAW(MT,NMX)=(-VN*B+VT2*A)/BTA(MT)
    LCOL(MT,NMX)=0
    ER(MT,NMX)=0.0

```

```

FLO0610
FLO0620
FLO0630
FLO0640
FLO0650
FLO0660
FLO0670
FLO0680
FLO0690
FLO0700
FLO0710
FLO0720
FLO0730
FLO0740
FLO0750
FLO0770
FLO0780
FLO0790
FLO0800
FLO0810
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FLO0950
FLO0960
FLO0970
FLO0980
FLO0990
FLO1000
FLO1010
FLO1020
FLO1030
FLO1040
FLO1050
FLO1060
FLO1070
FLO1080
FLO1090
FLO1100
FLO1110
FLO1120

```

```

IF (CHI (MT) .LE. -1.) GO TO 365
225 X=9.*RAND(0)
IF (X.EQ.0.0) GO TO 225
IT=X**CHI (MT) *EXP (-X)
IF (IT.GE.CHG (MT)) GO TO 226
CNG (MT)=CNG (MT)+IT
IF (CNG (MT) .LT. CHG (MT)) GO TO 225
CNG (MT)=CNG (MT)-CHG (MT)
226 ER (MT, NMI) =X
365 LPP (MT, NMI) =LWP (6)
370 CONTINUE
RETURN
380 LARGE=2
RETURN
END

```

FLO1130  
FLO1140  
FLO1150  
FLO1160  
FLO1170  
FLO1180

```

SUBROUTINE COLIDE (CN, CM, WTM, DB, DBA, NB, NCOL, LCOL, PAU, PAV, PAW, ER, T,
1LM, MT, I2, I3, NUMCEL, ETA, PHI, CHI, CN8, NP, LPP, LKW, NBP, NBM)
INTEGER TIME
INTEGER*2 LM (NP, 1), LCOL (NP, 1), LPP (NP, 1), LKW (1)
INTEGER*2 NBM, NB, NBP, NUMCEL
DIMENSION CN (3, 3, 1), CM (3, 3, 1), WTM (1), DB (NP, 1), DBA (NP, 1), NB (NP, 1)
DIMENSION NCOL (3, 1), T (NP, NP, 1), NUMCEL (1), ETA (3, 1), PHI (3, 1), CHI (1)
DIMENSION PAU (NP, 1), PAV (NP, 1), PAW (NP, 1), ER (NP, 1), CN8 (3, 1), WA (2)
DIMENSION NBP (NP, 1), NBM (NP, 1)
COMMON /FOUR/NBX
COMMON /FIFTH/ND, TIME, DTM
1 FORMAT (' TIME = ', P9.4, ' COLL. TIMES = ', 2F9.4, ' NUMBERS = ', 2I5/)
2 FORMAT (' COLIDE REACHED LINE 160 IN BOX NUMBER = ', I5, ' AT CPU TIM
1E = ', P9.4, ' VR= ', E12.4, ' REL VEL G = ', 3E12.4, ' EI = ', E12.4)
3 FORMAT (' COLIDE REACHED LINE 165 IN BOX NUMBER = ', I5, ' AT CPU TIM
1E = ', P9.4, ' VR= ', E12.4, ' REL VEL G = ', 3E12.4, ' EI = ', E12.4)

```

COL0030  
COL0080  
COL0090

-----

THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE ELAPSED TIMES IN COL0110  
CELLS BY AN AMOUNT APPROXIMATELY EQUAL TO THE PRE-SELECTED COLLISCOL0120  
TIME. THERE ARE FOUR TIMES FOR EACH CELL, SAVED IN AN AREA CALLECOL0130  
'T', CORRESPONDING TO THE FOUR TYPES OF MOLECULAR COLLISIONS WHICCOL0140  
CAN OCCUR. TO ADVANCE THE VARIOUS TIMES, AN APPROPRIATE NUMBER OFCOL0150  
THE CORRESPONDING MOLECULAR COLLISIONS IS COMPUTED. THE ACTUAL COL0160  
MOLECULES TO COLLIDE ARE SELECTED AT RANDOM, AND THEIR VELOCITY VCOL0170  
DIRECTIONS AFTER COLLISION ARE SELECTED AT RANDOM. COL0180

-----

```

AIMP=DTM*TIME
DO 240 MTA=1, MT
DO 230 MTB=1, MTA
D = WTM (MTA) + WTM (MTB)
WA (MTA) =WTM (MTA) /D
WB (MTB) =WTM (MTB) /D
RM=WTM (MTA) *WTM (MTB) /D
CHT=CHI (MTA) +CHI (MTB) +2.0
PHT=PHI (MTA, MTB)
ETT=ETA (MTA, MTB)
DO 220 M=1, NBX
N=NUMCEL (M)
IF (N.LE.0) GO TO 220

```

COL0190

```

AKW=LKW(N)
IF (T(MTA,MTB,N).LT.AIME) GO TO 100
IF (T(MTB,MTA,N).GE.AIME) GO TO 220
100 NA=NB(MTA,N)*NB(MTB,N)
IF (MTA.EQ.MTB) NA=(NA-NB(MTA,N))/2.
IF (NA.LT.1) GO TO 220
KS=0
120 KC=0
CPUT=ELTIME(0)
KS=KS+1
IF (KS.GT.NA) GO TO 220
130 KC=KC+1
IF (KC.GT.NA) GO TO 220
135 I=NB(MTA,N)*RAND(0)+1+NBM(MTA,N)
IF (I.GT.NBM(MTA,N+1)) I=NBM(MTA,N+1)
J=LJ(MTA,I)
CR=LPP(MTA,J)/AKW
IF (CR.GT.0.99) GO TO 140
IF (RAND(0).GT.CR) GO TO 135
140 K=NB(MTB,N)*RAND(0)+1+NBM(MTB,N)
IF (K.GT.NBM(MTB,N+1)) K=NBM(MTB,N+1)
IF (MTA.EQ.MTB.AND.I.EQ.K) GO TO 140
L=LJ(MTB,K)
CR=LPP(MTB,L)/AKW
IF (CR.GT.0.99) GO TO 145
IF (RAND(0).GT.CR) GO TO 140
145 CONTINUE
GM1=WA(MTA)*PAU(MTA,J)+WA(MTB)*PAU(MTB,L)
GM2=WA(MTA)*PAV(MTA,J)+WA(MTB)*PAV(MTB,L)
GM3=WA(MTA)*PAW(MTA,J)+WA(MTB)*PAW(MTB,L)
G1=PAU(MTA,J)-PAU(MTB,L)
G2=PAV(MTA,J)-PAV(MTB,L)
G3=PAW(MTA,J)-PAW(MTB,L)
GS=G1**2+G2**2+G3**2
IF (GS.LT.1.0E-8) GO TO 130
ET=RM*GS
EI=ER(MTA,J)+ER(MTB,L)
VR=GS**(.5-ZTT/2.)
IF (VR.GE.CH(MTA,MTB,1)) GO TO 160
CN(MTA,MTB,1)=CN(MTA,MTB,1)+VR
IF (CN(MTA,MTB,1).LT.CH(MTA,MTB,1)) GO TO 130
CN(MTA,MTB,1)=CN(MTA,MTB,1)-CN(MTA,MTB,1)
160 CONTINUE
CPUT=ELTIME(0)
IF (N.EQ.1196) WRITE(6,2) N,CPUT,VR,G1,G2,G3,EI
CALL IMPACT(RM,G1,G2,G3,ET,EI,PHT,CST,ETT,CN(MTA,MTB,2),CN(MTA,MTB
1,2))
165 CONTINUE
CPUT=ELTIME(0)
IF (N.EQ.1196) WRITE(6,3) N,CPUT,VR,G1,G2,G3,EI
IF (PHT.EQ.0.) GO TO 175
X1=0.0
IF (CHI(MTA).EQ.-1.) GO TO 175
X1=1.0
IF (CHI(MTB).EQ.-1.) GO TO 175

```

```

170 X1=Rand(0)
    IF ((CHI(MTA).EQ.0.).AND.(CHI(MTB).EQ.0.)) GO TO 175
    XT=X1**CHI(MTA)*(1.-X1)**CHI(MTB)
    IF (XT.GT.CH(MTA,MTB,3)) GO TO 175
    CN(MTA,MTB,3)=CN(MTA,MTB,3)+XT
    IF (CN(MTA,MTB,3).LT.CH(MTA,MTB,3)) GO TO 170
    CN(MTA,MTB,3)=CN(MTA,MTB,3)-CH(MTA,MTB,3)
175 CONTINUE
    C=DBA(MTA,N)
    D=DBA(MTB,N)
    IF (C.EQ.0.0) C=DB(MTA,N)
    IF (D.EQ.0.0) D=DB(MTB,N)
    IF (T(MTA,MTB,N).GE.AIME) GO TO 180
    PAU(MTA,J)=GM1+WA(MTB)*G1
    PAV(MTA,J)=GM2+WA(MTB)*G2
    PAW(MTA,J)=GM3+WA(MTB)*G3
    IF (PHT.GT.0.) ER(MTA,J)=X1*EI
    LCOL(MTA,J)=1
    NCOL(MTA,MTB)=NCOL(MTA,MTB)+1
    T(MTA,MTB,N)=T(MTA,MTB,N)+CNS(MTA,MTB)*LPP(MTA,J)/NBP(MTA,N)/D/VR
    IF (MTA.EQ.MTB) GO TO 190
180 IF (T(MTB,MTA,N).GE.AIME) GO TO 210
190 PAU(MTB,L)=GM1-WA(MTA)*G1
    PAV(MTB,L)=GM2-WA(MTA)*G2
    PAW(MTB,L)=GM3-WA(MTA)*G3
    IF (PHT.GT.0.) ER(MTB,L)=(1.-X1)*EI
    LCOL(MTB,L)=1
    NCOL(MTB,MTA)=NCOL(MTB,MTA)+1
    T(MTB,MTA,N)=T(MTB,MTA,N)+CNS(MTB,MTA)*LPP(MTB,L)/NBP(MTB,N)/C/VR
210 CONTINUE
    IF (M.EQ.1196) WRITE(6,1) AIME,T(MTA,MTB,N),T(MTB,MTA,N),NBP(MTA,N)
    1,NBP(MTB,N)
    IF (T(MTA,MTB,N).LT.AIME.OR.T(MTB,MTA,N).LT.AIME) GO TO 120
220 CONTINUE
230 CONTINUE
240 CONTINUE
    RETURN
    END

```

COL0810

COL0920

```

SUBROUTINE MOVE(KSWCH,AKN,MJ,NS,NWEDG,THETAZ,XSTART,I2,I3,I4,I5,DEMOV0010
1LANG,NWEDGE,BTA,C2,C3,DPA,PL,HTI,HTR,JNT,KNN,NH,ICB,XLIM,MS,INS,NTNOV0020
2CP,NTCV,PV,CTI,CTR,CNI,CNR,ALPHA,SIGMA,COEFP,HTS,HTSI,NTS,NTSP,UTLNOV0030
3,UTT,VTS,PAU,PAV,PAW,PAX,PAY,PAZ,LPP,LCOL,TB,IP,ER,CHI,CNG,CMG,I,
4UTLI,UTTI,VTSI)
    INTEGER*2 LPP(I,1),LCOL(I,1)
    INTEGER SWTCH,TIME
    LOGICAL DUMP
    REAL LAM,MU,NU
    DIMENSION DELANG(1),NWEDGE(1),BTA(1),C2(1),C3(1),PL(1),HTI(1)
    DIMENSION HTR(1),TB(1),ICB(1),ALPHA(3,1),SIGMA(3,1),COEFP(4,1)
    DIMENSION PAU(I,1),PAV(I,1),PAW(I,1),CTI(3,1),CTR(3,1)
    DIMENSION CNI(3,1),CNR(3,1),DPA(1),JNT(1),XLIM(1),KNN(1),NH(1)
    DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3),NTSP(3,I2,I3)
    DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)
    DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3),NTCP(3,I4)

```

HOV0060  
HOV0070  
HOV0080  
HOV0090

```

DIMENSION PAX(I,1),PAY(I,1),PAZ(I,1),IWS(1),MS(1)
DIMENSION NTCV(3,I4,2,I5,3),FV(3,I4,2,I5,3)
DIMENSION ER(I,1),CHG(1),CHG(1),CHI(1)
COMMON /THIRD/PI,NREG
COMMON /FOURTH/NSI,RE,XR,DUMP
COMMON /FIFTH/ND,TIME,DTH
COMMON /SYNTH/LAM,NU,NU,MT,N,J,XI,YI,ZI,TUSE
LABELIST/CHECK/TIME,X,Y,Z,DX,DY,DZ,TLEFT,RADS,RMS,XR
-----
THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE SPATIAL POSITION
OF ALL THE MOLECULES BY AN AMOUNT APPROPRIATE TO THEIR CURRENT
LOCITIES AND THE PRE-SELECTED COLLISION TIME.
-----
2 FORMAT(27H SOMETHING IS WRONG IN MOVE/3E20.7,4I7,E20.7)
NAREA=NREG+3
RMS=RM**2
DO 150 MT=1,IP
N=KNM(MT)
10 N=N+1
TLEFT=DTH
IF(KSWCH.EQ.1) TLEFT=TLEFT*RAND(0)
IF(N.GT.NM(MT)) GO TO 150
15 LAM=PAU(MT,N)
IF(LAM.EQ.0.) LAM=.0000001
NU=PAV(MT,N)
NU=PAW(MT,N)
XI=PAX(MT,N)
YI=PAY(MT,N)
ZI=PAZ(MT,N)
DX=TLEFT*LAM
DY=TLEFT*NU
DZ=TLEFT*NU
X=XI+DX
Y=YI+DY
Z=ZI+DZ
RADS=Y**2+Z**2
IF((RADS.GT.2.*RMS).OR.(ABS(X).GT.2.*XR)) WRITE(6,CHECK)
IF(RADS.GT.RMS) GO TO 100
RAD=SQRT(RMS)
KEY=ABS(LAM)/LAM+.5
DO 60 L=1,NAREA
IF(XI-XLIM(L)) 65,55,60
55 J=L+KEY-2
GO TO 70
60 CONTINUE
WRITE(6,2) DTH,XI,LAM,N,KEY,NAREA,L,XLIM(L)
IF(DUMP) CALL ABEND(16)
STOP
65 J=L-2
70 K=J+KEY+1
IF((K.EQ.0).OR.(K.EQ.NAREA+1)) GO TO 100
TUSE=(XLIM(K)-XI)/LAM
XTEMP=XLIM(K)
IF(TUSE.LE.TLEFT) GO TO 75
TUSE=TLEFT

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```

XTEMP=XI+TUSE*LAM
75 IF ((J.EQ.0).OR.(J.EQ.NREG+1)) GO TO 85
CALL INTERS(AKN,NJ,NS,NWEDG,SWTCH,THETAZ,XSTART,I2,I3,I4,I5,DELANG,NOV0650
1,NWEDGE,BTA,C2,C3,DFA,FL,HTI,HTR,JNT,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,NOV0660
2SIGMA,COEPP,HTS,HTSI,NTS,NTSP,UTL,UTT,VTS,HS,IWS,NTCF,NTCV,FV,PAU,NOV0670
3PAV,PAW,LPP,LCOL,IP,ER,CHI,CNG,CHG,I,UTLI,UTTI,VTSI)
IP(SWTCH.EQ.1) GO TO 90
85 XI=XTEMP
YI=YI+TUSE*NU
ZI=ZI+TUSE*NU
90 PAX(MT,N)=XI
PAY(MT,N)=YI
PAZ(MT,N)=ABS(ZI)
PAW(MT,N)=ABS(ZI)/ZI*PAW(MT,N)
TLEPT=TLEPT-TUSE
IF(TLEPT.GT.0.) GO TO 15
GO TO 10
100 NZ=NM(MT)
PAX(MT,N)=PAX(MT,NZ)
PAY(MT,N)=PAY(MT,NZ)
PAZ(MT,N)=PAZ(MT,NZ)
PAU(MT,N)=PAU(MT,NZ)
PAV(MT,N)=PAV(MT,NZ)
PAW(MT,N)=PAW(MT,NZ)
ER(MT,N)=ER(MT,NZ)
LPP(MT,N)=LPP(MT,NZ)
LCOL(MT,N)=LCOL(MT,NZ)
N=N-1
NM(MT)=NM(MT)-1
GO TO 10
150 CONTINUE
RETURN
END

```

```

SUBROUTINE ACCUM(I2,I3,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,
1IP,I,LPP,NBP,NBM)
INTEGER*2 LM(I,1),LPP(I,1),NBM(I,1)
INTEGER*2 NB,NBP
DIMENSION FNB(1),NB(I,1),PAU(I,1),PAV(I,1),PAW(I,1),TMP(I,1)
DIMENSION XV(I,1),YV(I,1),ZV(I,1),ER(I,1),TRP(I,1),NBP(I,1)
COMMON /FORTH/NBX

```

-----  
 THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE TEMPERATURES,  
 VELOCITIES, AND DENSITIES IN VARIOUS ARRAYS FOR DETERMINING THE  
 AVERAGE FLOW FIELD PROPERTIES AFTER STEADY-STATE HAS BEEN REACHED  
 -----

```

N=0
DO 180 K=1,NBX
IF(FNB(K).LE.0.) GO TO 180
N=N+1
DO 110 MT=1,IP
IV(MT,N)=0.0
YV(MT,N)=0.0
ZV(MT,N)=0.0
TMP(MT,N)=0.0

```

TRP(MT,N)=0.0	ACUM200
TTX=0.	ACUM210
TTY=0.	ACUM220
TTZ=0.	
TTR=0.0	ACUM230
N=NB(MT,N)	ACUM240
IP(M.LT.1) GO TO 110	ACUM250
U=0.	ACUM260
V=0.	ACUM270
W=0.	ACUM280
DO 100 L=1,N	
NA=NBE(MT,N)+L	
J=LE(MT,NA)	
AKW=LPF(MT,J)	
PU=PAU(MT,J)	ACUM300
PV=PAV(MT,J)	ACUM310
PW=PAW(MT,J)	ACUM320
U=U+PU*AKW	
V=V+PV*AKW	
W=W+PW*AKW	
TTR=TTR+ER(MT,J)*AKW	
TTX=TTX+PU*PU*AKW	
TTY=TTY+PV*PV*AKW	
20 TTZ=TTZ+PW*PW*AKW	
S=NBP(MT,N)	ACUM390
XV(MT,N)=U/S	ACUM400
YV(MT,N)=V/S	ACUM410
ZV(MT,N)=W/S	ACUM420
TMP(MT,N)=(TTX+TTY+TTZ)/S	
TRP(MT,N)=TTR/S	ACUM430
10 CONTINUE	ACUM440
30 CONTINUE	ACUM450
RETURN	ACUM460
END	

SUBROUTINE AVRGE(FNB,DB,DBA,NB,NBT,XV,YV,ZV,XVA,YVA,ZVA,TMP,TMPA,  
 1TRP,TRPA,IP,I,NBP,NBS)  
 INTEGER\*2 NB,NBT,NBP,NBS  
 DIMENSION FNB(1),DB(1,1),DBA(1,1),NB(1,1),NBT(1,1),TMP(1,1)  
 DIMENSION TMPA(1,1),XV(1,1),XVA(1,1),YV(1,1),YVA(1,1),ZV(1,1)  
 DIMENSION ZVA(1,1),TRP(1,1),TRPA(1,1),NBP(1,1),NBS(1,1)  
 COMMON /FORTH/NBX

---

THE PURPOSE OF THIS SUBROUTINE IS TO COMPUTE THE AVERAGE FLOW  
 FIELD PROPERTIES.

---

N=0	AVG0050
DO 110 M=1,NBX	AVG0060
IF(FNB(M).LE.0.) GO TO 110	AVG0070
N=N+1	AVG0080
DO 100 MT=1,IP	AVG0090
A=NBT(MT,N)	AVG0100
B=NBP(MT,N)	AVG0110
C=A+B	AVG0120
NBT(MT,N)=C	AVG0130
	AVG0150
	AVG0170
	AVG0180



```

NBS(MT,N)=NBS(MT,N)+NB(MT,N)
IF(C.LT.1.) GO TO 100
DBA(MT,N)=(DBA(MT,N)*A+DB(MT,N)*B)/C
IVA(MT,N)=(IVA(MT,N)*A+IV(MT,N)*B)/C
YVA(MT,N)=(YVA(MT,N)*A+YV(MT,N)*B)/C
ZVA(MT,N)=(ZVA(MT,N)*A+ZV(MT,N)*B)/C
TMPA(MT,N)=(TMPA(MT,N)*A+TMP(MT,N)*B)/C
TRPA(MT,N)=(TRPA(MT,N)*A+TRP(MT,N)*B)/C
100 CONTINUE
110 CONTINUE
RETURN
END

```

AVG0190  
AVG0200  
AVG0210  
AVG0220  
AVG0230  
AVG0240  
AVG0250  
AVG0260  
AVG0270  
AVG0280

```

SUBROUTINE DRAG(AKW,NJ,NS,NWEDG,THETAZ,ISTART,I2,I3,I4,I5,DELANG,
1NWEDGE,BTA,C2,C3,DFA,FL,HTI,HTR,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA
2,COEPP,HTS,HTSI,NTS,NTSP,UTL,UTT,VTS,MS,IWS,NTCP,NTCV,FV,PAU,PAV,P
3AW,LPP,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
INTEGER*2 LPP,LCOL
INTEGER TIME,TST
REAL LAM,MU,NU,JAY,KAY
DIMENSION DELANG(1),NWEDGE(1),BTA(1),C2(1),C3(1),FL(1),HTI(1)
DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1),COEPP(4,1)
DIMENSION LPP(1,1),PAU(1,1),PAV(1,1),PAW(1,1),CTI(3,1),CTR(3,1)
DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3),NTSP(3,I2,I3)
DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3),LCOL(1,1)
DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)
DIMENSION CNI(3,1),CNR(3,1),DPA(1),IWS(1),MS(1),NTCP(3,I4)
DIMENSION NTCV(3,I4,2,I5,3),FV(3,I4,2,I5,3)
DIMENSION ER(1,1),CNG(1),CMG(1),CHI(1)
COMMON /THIRD/PI
COMMON /FIFTH/ND,TIME,DTM,TI,ITS,ITP,TST
COMMON /SYNTH/LAM,MU,NU,MT,N,J,XCL,YCL,ZCL

```

DRG0040  
DRG0050  
DRG0060  
DRG0070  
DRG0140  
DRG0150  
DRG0160  
DRG0170

---

```

THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE THE DRAG AND HEAT
TRANSFER INCREMENTS ON THE BODY CONTRIBUTED BY EACH MOLECULE WHICH
COLLIDES WITH THE BODY. IN ADDITION, EACH MOLECULE WHICH COLLIDES
WITH THE BODY IS ASSIGNED AN APPROPRIATE NEW VELOCITY (OF REFLECTED
WHICH IS USED TO CONTINUE ITS SPATIAL TRANSLATION (IN SUBROUTINE

```

DRG0180  
DRG0190  
DRG0200  
DRG0210  
DRG0220  
DRG0230

---

```

CALL NORMAL(EYE,JAY,KAY,ONE,COEPP)
AKW=LPP(MT,N)
RAD=SQRT(YCL*YCL+ZCL*ZCL)
ARG=YCL/RAD
TANG=180.*(1.-ARCCOS(ARG)/PI)
IWDG=TANG/DELANG(1)+1.
IF((IWDG.GT.NWEDGE(1)).AND.(DELANG(2).NE.0.)) IWDG=(TANG-THETAZ)
1/DELANG(2)+NWEDGE(1)+1
IF(IWDG.LT.1) IWDG=1
IF(IWDG.GT.NWEDG) IWDG=NWEDG
D=(LAM*LAM+MU*MU+NU*NU)*AKW
G=ER(MT,N)*AKW
H=G
DO 100 H=1,ND
IF(XCL.LT.XCB(3)) GO TO 110
100 CONTINUE

```

DRG0240  
DRG0250  
DRG0260  
DRG0270  
DRG0280  
DRG0290  
DRG0310  
DRG0320  
DRG0330  
DRG0340  
DRG0350

```

110 UI=LAM                                DRG0360
    WI=(NU*JAY-HU*KAY)/ONE                DRG0370
    VID=LAM*EYE+HU*JAY+NU*KAY            DRG0380
    UID=LAM*ONE-EYE*(HU*JAY+NU*KAY)/ONE  DRG0390
    E=Rand(0)                              DRG0400
    IF (E.LT.SIGMA(MT,M)) GO TO 115        DRG0410
    VRD=-VID                               DRG0420
    URD=UID                                 DRG0430
    WR=WI                                   DRG0440
    GO TO 125                               DRG0450
115 V=4.*RAND(0)                           DRG0460
    C2(MT)=C2(MT)+.544331*V*V*V*EXP(1.5-V*V) DRG0470
    IF (C2(MT).LT.1.) GO TO 115           DRG0480
    C2(MT)=C2(MT)-1.                       DRG0490
    IF (NTSP(MT,M,IWDG).NE.0) GO TO 117   DRG0500
    ATR=ALPHA(MT,M)*TB(M)/SIGMA(MT,M)     DRG0510
    GO TO 118                               DRG0520
117 ATR=ALPHA(MT,M)*TB(M)/SIGMA(MT,M)+(1.-ALPHA(MT,M)/SIGMA(MT,M))*HTSDRG0530
    I(MT,M,IWDG)/NTSP(MT,M,IWDG)/(3.+CHI(MT))
118 ABR=SQRT(ATR)                           DRG0550
    V=V*ABR/BTA(MT)                        DRG0560
120 A=Rand(0)                              DRG0570
    C3(MT)=C3(MT)+A                         DRG0580
    IF (C3(MT).LT.1.) GO TO 120           DRG0590
    C3(MT)=C3(MT)-1.                       DRG0600
    B=SQRT(1.-A*A)                         DRG0610
    C=2.*PI*RAND(0)                        DRG0620
    VRD=V*A                                 DRG0630
    URD=V*B*COS(C)                         DRG0640
    WR=V*B*SIN(C)                          DRG0650
    IF (CHI(MT).EQ.-1.) GO TO 125
122 X=9.*RAND(0)                           DRG0660
    IF (X.EQ.0.0) GO TO 122               DRG0670
    XTEMP=1.0                              DRG0680
    IF (CHI(MT).NE.0.0) XTEMP=X**CHI(MT)  DRG0690
    CNG(MT)=CNG(MT)+XTEMP*EXP(-X)
    IF (CNG(MT).LT.CNG(MT)) GO TO 122
124 CONTINUE
    CNG(MT)=CNG(MT)-CNG(MT)
    IF (CNG(MT).GE.CNG(MT)) GO TO 124
    ER(MT,M)=X*ATR
    H=ER(MT,M)*AKW
125 UR=EYE*VRD+ONE*URD                     DRG0660
    PAU(MT,M)=UR                           DRG0670
    PAV(MT,M)=JAY*VRD-(KAY*WR+EYE*JAY*URD)/ONE DRG0680
    PAW(MT,M)=KAY*VRD+(JAY*WR-EYE*KAY*URD)/ONE DRG0690
    IF (TIME.LE.TST) RETURN
    IF (TI.GT.0.) GO TO 130                DRG0710
    TI=TST*DTH
130 IMZ=(ICL-XSTART)*AKN                   DRG0730
    IMZ=RAD*AKN                             DRG0740
    B=(URD*URD+VED*VRD+WR*WR)*AKW         DRG0750
    UTI=UID*UID+WI*WI                       DRG0760
    UYI=-(UID*EYE*JAY+WI*KAY)/ONE         DRG0770
    UYR=-(URD*EYE*JAY+WR*KAY)/ONE        DRG0780

```

```

PL(MT)=PL(MT)+AKW*DPA(MT)
HTI(MT)=HTI(MT)+D+G
HTR(MT)=HTR(MT)-B-H
CTI(MT,1)=CTI(MT,1)+UID*ONE*AKW
CTI(MT,2)=CTI(MT,2)+UYI*AKW
CTI(MT,3)=CTI(MT,3)+(XMZ*UYI-YMZ*UID*ONE)*AKW
CNI(MT,1)=CNI(MT,1)+VID*EYE*AKW
CNI(MT,2)=CNI(MT,2)+VID*JAY*AKW
CNI(MT,3)=CNI(MT,3)+(XMZ*JAY-YMZ*EYE)*VID*AKW
CTR(MT,1)=CTR(MT,1)-URD*ONE*AKW
CTR(MT,2)=CTR(MT,2)-UYR*AKW
CTR(MT,3)=CTR(MT,3)-(XMZ*UYR-YMZ*URD*ONE)*AKW
CNR(MT,1)=CNR(MT,1)-VRD*EYE*AKW
CNR(MT,2)=CNR(MT,2)-VRD*JAY*AKW
CNR(MT,3)=CNR(MT,3)-(XMZ*JAY-YMZ*EYE)*VRD*AKW
NTS(MT,M,IWDG)=NTS(MT,M,IWDG)+1
NTSF(MT,M,IWDG)=NTSF(MT,M,IWDG)+AKW
UTLI(MT,M,IWDG)=UTLI(MT,M,IWDG)+UID*AKW
UTL(MT,M,IWDG)=UTL(MT,M,IWDG)+(UID-URD)*AKW
UTTI(MT,M,IWDG)=UTTI(MT,M,IWDG)+WI*AKW
UTT(MT,M,IWDG)=UTT(MT,M,IWDG)+(WI-WR)*AKW
VTSI(MT,M,IWDG)=VTSI(MT,M,IWDG)-VID*AKW
VTS(MT,M,IWDG)=VTS(MT,M,IWDG)+(VRD-VID)*AKW
HTSI(MT,M,IWDG)=HTSI(MT,M,IWDG)+D+G
HTS(MT,M,IWDG)=HTS(MT,M,IWDG)+D-B+G-H
IF(NS.EQ.0) RETURN
NC=(2*LCOL(MT,N))/(1+LCOL(MT,N))+1
DO 160 L=1,NS
MTEST=MS(L)
IWT=IWS(L)
IF((M.NE.MTEST).OR.(IWDG.NE.IWT)) GO TO 160
IF(NC.EQ.1) NTCF(MT,L)=NTCF(MT,L)+AKW
DO 145 JJ=1,MJ
IP(ABS(VID).GT.PV(MT,L,NC,JJ,1)) GO TO 135
NTCV(MT,L,NC,JJ,1)=NTCV(MT,L,NC,JJ,1)+AKW
135 IP(UID.GT.PV(MT,L,NC,JJ,2)) GO TO 140
NTCV(MT,L,NC,JJ,2)=NTCV(MT,L,NC,JJ,2)+AKW
140 IP(WI*ABS(ZCL)/ZCL.GT.PV(MT,L,NC,JJ,3)) GO TO 145
NTCV(MT,L,NC,JJ,3)=NTCV(MT,L,NC,JJ,3)+AKW
145 CONTINUE
160 CONTINUE
RETURN
END

```

DRG0790  
 DRG0820  
 DRG0830  
 DRG0840  
 DRG0850  
 DRG0860  
 DRG0870  
 DRG0880  
 DRG0890  
 DRG0900  
 DRG0910  
 DRG0920  
 DRG0930  
 DRG0940  
 DRG0950  
 DRG0960  
 DRG0970  
 DRG0980  
 DRG1010  
 DRG1020  
 DRG1060  
 DRG1080  
 DRG1150  
 DRG1160  
 DRG1170  
 DRG1180

```

SUBROUTINE INTERS(AKW,MJ,NS,NWEDG,SWTCH,THETAZ,XSTART,I2,I3,I4,I5,INT0010
1DELANG,NWEDGE,BTA,C2,C3,DPA,FL,HTI,HTR,JNT,TB,XCB,CTI,CTR,CNI,CNR,INT0020
2ALPHA,SIGMA,COEFF,HTS,HTSI,NTS,NTSF,UTL,UTT,VTS,MS,IWS,NTCF,NTCV,PINT0030
3V,PAU,PAV,PAW,LPP,LCOL,IP,ER,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
INTEGER*2 LPP,LCOL
INTEGER SWTCH
REAL LAN,MU,NU,LINEAR
DIMENSION DELANG(1),NWEDGE(1),BTA(1),C2(1),C3(1),FL(1),HTI(1)
DIMENSION HTR(1),TB(1),XCB(1),ALPHA(3,1),SIGMA(3,1),COEFF(4,1)
DIMENSION LPP(I,1),PAU(I,1),PAV(I,1),PAW(I,1),CTI(3,1),CTR(3,1)
DIMENSION HTS(3,I2,I3),HTSI(3,I2,I3),NTS(3,I2,I3),NTSF(3,I2,I3)

```

INT0050  
 INT0060  
 INT0070  
 INT0080

ORIGINAL FILE IS  
OF POOR QUALITY

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

FILE: GKBEY DECK A

DIMENSION UTL(3,I2,I3),UTT(3,I2,I3),VTS(3,I2,I3),LCOL(I,1)  
DIMENSION UTLI(3,I2,I3),UTTI(3,I2,I3),VTSI(3,I2,I3)  
DIMENSION CNI(3,1),CNR(3,1),DPA(1),IWS(1),MS(1),NTCP(3,I4)  
DIMENSION JNT(1),NTCV(3,I4,2,I5,3),FV(3,I4,2,I5,3)  
DIMENSION ER(I,1),CNG(1),CHG(1),CHI(1)  
COMMON /SVNTH/LAM,NU,NU,HT,N,J,XI,YI,ZI,TUSE

-----INTO150  
THE PURPOSE OF THIS SUBROUTINE IS TO DETERMINE FOR EACH MOLECULARINTO160  
TRAJECTORY IF THERE IS AN INTERSECTION OF THE TRAJECTORY WITH THEINTO170  
BODY SURFACE.-----INTO180  
INTO190

-----INTO200  
SWTCH=0INTO210  
A=COEPP(1,J)INTO220  
B=COEPP(2,J)INTO230  
C=COEPP(3,J)INTO240  
D=COEPP(4,J)INTO250  
ONE=A\*LAMINTO260  
TWO=B\*NUINTO270  
TRE=B\*NUINTO280  
PUR=.5\*C\*LAMINTO290  
SQUARE=ONE\*LAM+TWO\*NU+TRE\*NUINTO300  
LINEAR=ONE\*XI+TWO\*YI+TRE\*ZI+PURINTO310  
CONST=(A\*XI+C)\*XI+B\*(YI\*YI+ZI\*ZI)+DINTO320  
IF(SQUARE.EQ.0.) GO TO 150INTO330  
DISCR=LINEAR\*LINEAR-SQUARE\*CONSTINTO340  
IF(DISCR.LT.0.) RETURNINTO350  
SDISC = SQRT(DISCR)INTO360  
TYME = (-LINEAR-SDISC)/SQUAREINTO370  
GO TO 250INTO380  
150 IF(LINEAR.EQ.0.) RETURNINTO390  
TYME=-.5\*CONST/LINEARINTO400  
250 IF(TYME.GT.TUSE) RETURNINTO410  
IF(TYME.LE.0.) RETURNINTO420  
JNT(HT)=JNT(HT)+1INTO430  
XI=XI+LAM\*TYMEINTO440  
YI=YI+NU\*TYMEINTO450  
ZI=ZI+NU\*TYMEINTO460  
CALL DRAG(AKW,NJ,NS,NWEDG,THETAZ,ISTART,I2,I3,I4,I5,DELANG,NWEDGE,INTO470  
1BTA,C2,C3,DPA,PL,HTI,HTB,TE,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA,COEPP,INTO480  
2HTS,HTSI,NTS,NTSP,UTL,UTT,VTS,NS,IWS,NTCP,NTCV,FV,PAU,PAV,PAW,LPP,INTO490  
3LCOL,IP,ER,CHI,CNG,CHG,I,UTLI,UTTI,VTSI)  
LCOL(HT,N)=2INTO510  
TUSE=TYMEINTO520  
SWTCH=1INTO530  
RETURNINTO540  
ENDINTO550

NORM010  
SUBROUTINE NORMAL(EYE,JAY,KAY,ONE,COEPP)NORM020  
REAL LAM,NU,NU,JAY,KAYNORM030  
DIMENSION COEPP(4,1)NORM040  
COMMON /SVNTH/LAM,NU,NU,HT,N,J,XCL,YCL,ZCLNORM050  
DPDX=2.\*COEPP(1,J)\*XCL+COEPP(3,J)NORM060  
DPDY=2.\*COEPP(2,J)\*YCLNORM070  
DPDZ=2.\*COEPP(4,J)\*ZCLNORM080  
DENOM=SQRT(DPDX\*DPDX+DPDY\*DPDY+DPDZ\*DPDZ)

FILE: GKBEXT DECK A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

```

EYE=DFDX/DENOM
JAY=DFDY/DENOM
KAY=DFDZ/DENOM
ONE=SQRT(JAY*JAY+KAY*KAY)
RETURN
END
NORM090
NORM100
NORM110
NORM120
NORM130
NORM140

FUNCTION ARCCOS(ARG)
COMMON /THIRD/PI
IF (ARG) 30, 10, 20
10 A=-.5*PI
GO TO 40
20 A=ATAN(SQRT(1.-ARG*ARG)/ARG)
GO TO 40
30 A=PI+ATAN(SQRT(1.-ARG*ARG)/ARG)
40 ARCCOS=A
RETURN
END
ARCS010
ARCS020
ARCS030
ARCS040
ARCS050
ARCS060
ARCS070
ARCS080
ARCS090
ARCS100
ARCS110

FUNCTION ERRF(SS)
ERRF=ERFC(-SS)
RETURN
END
EERF010
ERRF050
ERRF060

SUBROUTINE PRINT1(DT,COSANG,SINANG,RMA,RNU,DRP,PCP,HTP,PL,HTI,HTR,
1CTI,CTR,CNI,CNR)
DIMENSION DD(3),WD(2,5),PP(4,4),QQ(4,4),RR(4,4),SS(4,4),TT(4,4)
DIMENSION UU(4,4),P1(4,4),Q1(4,4),R1(4,4),PA(4),PB(4),PC(4)
DIMENSION FL(1),HTI(1),HTR(1),CTI(3,1),CTR(3,1),CNI(3,1),CNR(3,1)
DIMENSION RMA(1),RNU(1)
DATA WD/'X-PO','RCE ','Y-PO','RCE ','Z-NO','MENT','DRAG',' ','LPT10070
1IPT',' '/
PT10080
-----
THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE GROSS SURFACE PT10100
COEFFICIENTS OF THE BODY.
PT10110
-----
PT10120
PT10130
PT10140
PT10150
PT10160
PT10170
PT10180

FORMATS

1 FORMAT(//1X,50('*'),' GROSS SURFACE COEFFICIENTS ',50('*')/' MOLEC
1ULAR WEIGHT',12X,P8.3,3(19X,P8.3)/25X,
2 'INC. REP. TOT. INC. REP. TOT. INC.
3REP. TOT. INC. REP. TOT.')
```

```

RMR=0.0
DO 50 MT=1,3
DD (MT)=RMA (MT)*RNU (MT)*DRP/DT
50 RMR=RMR+RMA (MT)*RNU (MT)
WRITE (6, 1) (RMA (MT), MT=1,3), RMR
PF=PL (1)*PCF/DT
QF=PL (2)*PCF/DT
RF=PL (3)*PCF/DT
SF=PF+QF+RF
WRITE (6, 10) PF, QF, RF, SF
DO 200 I=1,3
PP (4, I)=0.0
QQ (4, I)=0.0
RR (4, I)=0.0
SS (4, I)=0.0
TT (4, I)=0.0
UU (4, I)=0.0
P1 (4, I)=0.0
Q1 (4, I)=0.0
R1 (4, I)=0.0
DO 150 MT=1,3
PP (MT, I)=CTI (MT, I)*DD (MT)/RMR
QQ (MT, I)=CTR (MT, I)*DD (MT)/RMR
SS (MT, I)=CNI (MT, I)*DD (MT)/RMR
TT (MT, I)=CNR (MT, I)*DD (MT)/RMR
P1 (MT, I)=PP (MT, I)+SS (MT, I)
Q1 (MT, I)=QQ (MT, I)+TT (MT, I)
RR (MT, I)=PP (MT, I)+QQ (MT, I)
UU (MT, I)=SS (MT, I)+TT (MT, I)
R1 (MT, I)=P1 (MT, I)+Q1 (MT, I)
PP (4, I)=PP (4, I)+PP (MT, I)
QQ (4, I)=QQ (4, I)+QQ (MT, I)
RR (4, I)=RR (4, I)+RR (MT, I)
SS (4, I)=SS (4, I)+SS (MT, I)
TT (4, I)=TT (4, I)+TT (MT, I)
UU (4, I)=UU (4, I)+UU (MT, I)
P1 (4, I)=P1 (4, I)+P1 (MT, I)
Q1 (4, I)=Q1 (4, I)+Q1 (MT, I)
R1 (4, I)=R1 (4, I)+R1 (MT, I)
150 CONTINUE
WRITE (6, 12) (WD (J, I), J=1, 2), (PP (K, I), QQ (K, I), RR (K, I), K=1, 4)
WRITE (6, 14) (SS (K, I), TT (K, I), UU (K, I), K=1, 4)
WRITE (6, 16) (P1 (K, I), Q1 (K, I), R1 (K, I), K=1, 4)
200 CONTINUE
AA=COSANG
BB=SINANG
DO 300 I=4,5
DO 250 K=1,4
PP (K, 4)=AA*PP (K, 1)+BB*PP (K, 2)
QQ (K, 4)=AA*QQ (K, 1)+BB*QQ (K, 2)
RR (K, 4)=AA*RR (K, 1)+BB*RR (K, 2)
SS (K, 4)=AA*SS (K, 1)+BB*SS (K, 2)
TT (K, 4)=AA*TT (K, 1)+BB*TT (K, 2)
UU (K, 4)=AA*UU (K, 1)+BB*UU (K, 2)
P1 (K, 4)=AA*P1 (K, 1)+BB*P1 (K, 2)

```

PT10400

PT10460

PT10470

PT10480

PT10490

PT10630

PT10640

PT10650

PT10660

PT10680

PT10690

PT10700

PT10710

PT10720

PT10730

PT10740

```

250 Q1(K,4)=AA*Q1(K,1)+BB*Q1(K,2)
R1(K,4)=AA*R1(K,1)+BB*R1(K,2)
WRITE(6,12) (WD(J,I),J=1,2), (PP(K,4),QQ(K,4),RR(K,4),K=1,4)
WRITE(6,14) (SS(K,4),TT(K,4),UU(K,4),K=1,4)
WRITE(6,16) (P1(K,4),Q1(K,4),R1(K,4),K=1,4)
AA=-SINANG
BB=COSANG
300 CONTINUE
HD=HTF/DT
PA(4)=0.0
PB(4)=0.0
PC(4)=0.0
DO 400 MT=1,3
PA(MT)=HTI(MT)*RMA(MT)*RNU(MT)*HD/RMR
PB(MT)=HTR(MT)*RMA(MT)*RNU(MT)*HD/RMR
PC(MT)=PA(MT)+PB(MT)
PA(4)=PA(4)+PA(MT)
PB(4)=PB(4)+PB(MT)
PC(4)=PC(4)+PC(MT)
400 CONTINUE
WRITE(6,18) (PA(I),PB(I),PC(I),I=1,4)
RETURN
END

```

PT10750  
PT10760PT10800  
PT10810  
PT10820  
PT10830PT10950  
PT10960

```

SUBROUTINE PRINT2(AKN,XSTART,DT,RNU,RMA,DRF,FCF,HTF,UTLI,UTTI,VTSI
1,HTSI,DELANG,NWEDGE,XS,XCB,YCB,HTS,NTS,NTSP,UTL,UTT,VTS,I2,I3,IP)
DIMENSION RMA(1),RNU(1),DELANG(1),NWEDGE(1),XS(1),XCB(1),YCB(1)
DIMENSION HTS(3,I2,I3),NTS(3,I2,I3),NTSP(3,I2,I3),UTL(3,I2,I3)
DIMENSION UTT(3,I2,I3),VTS(3,I2,I3),UTLI(3,I2,I3),UTTI(3,I2,I3)
DIMENSION VTSI(3,I2,I3),HTSI(3,I2,I3)
COMMON /FIFTH/ND

```

PT20060

PT20070

```

-----
THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE DISTRIBUTION
ON SURFACE OF THE SURFACE COEFFICIENTS
-----

```

PT20080

PT20090

PT20100

PT20110

PT20120

PT20130

PT20140

PT20150

PT20160

## FORMATS

```

8 FORMAT(/1X,45('*'),' DISTRIBUTION ON SURFACE ',45('*')/71X,'INC.
1 TOT. INC. TOT. INC. TOT. '/11X,' SEGMENT GEOMETRY',
214X,'MOL. HOLE SAMP NUM. SKIN SKIN PRES- PRES-
3 HEAT HEAT'/' NO. CENTER DELI CENTER DELANG',4X,'WGHT. P
4RACT.',10X,
5 'PLUX FRCTN FRCTN SURE SURE TRNSP TRNSP')
10 FORMAT(1X,I3,F8.3,F7.3,F9.3,F8.3,1X,2F8.4,I6,7F8.4)
11 FORMAT(37X,F8.4,' 1.0000',I6,7F8.4)

```

PT20230

\*\*\*\*\* PT20240

PT20250

PT20260

```

RMR=0.0
DO 50 MT=1,IP

```

```

50 RMR=RMR+RMA(MT)*RNU(MT)
WRITE(6,8)
I=0
DO 110 N=1,ND
DTY=DT*YCB(N)/180.
P=XS(N)
Q=2.*(ICB(N)-ISTART)*AKN-IS(N)
ANGLE=0.
R=0.
J=0
DO 105 L=1,2
R=R+.5*ANGLE
ANGLE=DELANG(L)
R=R-.5*ANGLE
ICNT=NWEDGE(L)
IF(ICNT.EQ.0) GO TO 105
PMLT=PCF/(DTY*ANGLE)
QMLT=DRF/(DTY*ANGLE)
SMLT=HTF/(DTY*ANGLE)
DO 100 K=1,ICNT
R=R+ANGLE
I=I+1
J=J+1
M3=0
P3=0.0
Q3=0.0
Q4=0.0
R3=0.0
R4=0.0
S3=0.0
S4=0.0
DO 90 MT=1,IP
M1=NTS(MT,N,J)
M3=M3+M1
P1=NTSF(MT,N,J)*PMLT*RNU(MT)
P3=P3+P1
Q1=SQRT(UTLI(MT,N,J)**2+UTTI(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
Q2=SQRT(UTL(MT,N,J)**2+UTT(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
Q3=Q3+Q1
Q4=Q4+Q2
R1=VTSI(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
R2=VTS(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
S1=HTSI(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR
S2=HTS(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR
R3=R3+R1
R4=R4+R2
S3=S3+S1
S4=S4+S2
90 WRITE(6,10)I,P,Q,R,ANGLE,RMA(MT),RNU(MT),M1,P1,Q1,Q2,R1,R2,S1,S2
WRITE(6,11)RMR,M3,P3,Q3,Q4,R3,R4,S3,S4
100 CONTINUE
105 CONTINUE
110 CONTINUE
RETURN
END

```

PT20280  
PT20320  
PT20330  
PT20340  
PT20350  
PT20360  
PT20370  
PT20380  
  
PT20400  
  
PT20420  
  
PT20440  
PT20450  
PT20460  
PT20470  
PT20480  
PT20490  
PT20500

PT20660

PT20670

PT20680

PT20690



```

SUBROUTINE PRINT3(      IP,NJ,NS,NWEDG,I2,I3,I4,I5, RMA,XS,IWS,
1MS,TANGN,NTSP,NTCF,NTCV,FV)      PT30020
DIMENSION RMA(1),XS(1),IWS(1),MS(1),TANGN(1),NTSP(3,I2,I3)
DIMENSION NTCF(3,I4), NTCV(3,I4,2,I5,3),FV(3,I4,2,I5,3),QUO(3,3) PT3

```

```

-----
                FORMATS
2 FORMAT (//1X,40('*'),' MOMENTS OF INCIDENT DISTRIBUTION FUNCTIONS ' PT30050
1,40('*'))      PT30060
6 FORMAT (/21X,I5,' UNCOLLIDED MOLECULES',P8.4,27X,I5,' COLLIDED MOLE PT30070
1CULES',P8.4/) PT30080
8 FORMAT (I5,4X,A3,12(1X,P9.4))      PT30090
10 FORMAT (12X,12(1X,P9.4))          PT30100
12 FORMAT (1H )                      PT30110

```

```

*****
WRITE(6,2)
DO 155 I=1,NS
MR=MS(I)
ITT=IWS(I)
N=(MR-1)*NWEDG+ITT
DO 150 MT=1,IP
A=0.
B=1.
IC=NTCF(MT,I)
ID=NTSP(MT,MR,ITT)-IC
E=NTSP(MT,MR,ITT)
IF(E.LE.0.) GO TO 110
A=IC/E
B=1.-A
110 WRITE(6,6) IC,A,ID,B
E=NTCF(MT,I)
DO 121 NC=1,2
DO 120 K=1,3
QUO(NC,K)=0.
IF(E.EQ.0.) GO TO 120
QUO(NC,K)=NTCV(MT,I,NC,1,K)/E
120 CONTINUE
E=NTSP(MT,MR,ITT)-NTCF(MT,I)
121 CONTINUE
WRITE(6,8) N, RMA(MT),FV(MT,I,1,1,1),QUO(1,1),FV(MT,I,1,1,2),QUO
1(1,2),FV(MT,I,1,1,3),QUO(1,3),FV(MT,I,2,1,1),QUO(2,1),FV(MT,I,2,1, PT30220
22),QUO(2,2),FV(MT,I,2,1,3),QUO(2,3) PT30230
IF(NJ.EQ.1) GO TO 150 PT30240
DO 140 J=2,NJ
E=NTCF(MT,I)
DO 131 NC=1,2
DO 130 K=1,3
QUO(NC,K)=0.
IF(E.EQ.0.) GO TO 130

```

```

      QUO(NC,K)=NTCV(MT,I,NC,J,K)/E
130 CONTINUE
      E=NTSP(MT,MR,ITT)-NTCP(MT,I)
131 CONTINUE
140 WRITE(6,10) FV(MT,I,1,J,1),QUO(1,1),FV(MT,I,1,J,2),QUO(1,2),FV(MT,I,1,J,3),QUO(1,3),
      FV(MT,I,2,J,1),QUO(2,1),FV(MT,I,2,J,2),QUO(2,2),FV(MT,I,2,J,3),QUO(2,3)
150 WRITE(6,12)
155 CONTINUE
      RETURN
      END

```

PT30640  
PT30650  
PT30660  
PT30670  
PT30680  
PT30690  
PT30700  
PT30710  
PT30720  
PT30730

```

SUBROUTINE PRINT4(MSP,CHI,RNU,I,TRP,NUMCEL,PDN,WTS,DB,NS,TMP,YV,
1YV,ZV,KS,NB,IC,YC,ZC,LEV,LKW)
INTEGER*2 LKW(1)
INTEGER*2 NB,NUMCEL,NS
DIMENSION PDN(1),RNU(1),CHI(1),WTS(1),NUMCEL(1),TMP(I,1),TRP(I,1)
DIMENSION DB(I,1),NB(I,1),YV(I,1),ZV(I,1),DBT(3),NS(I,1)
DIMENSION IC(1),YC(1),ZC(1),LEV(1)
COMMON /PORTH/NBI

```

PT50050  
PT50060  
PT50070  
PT50080

-----  
THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE INSTANTANEOUS  
FLOW-FIELD PROPERTIES.  
-----

PT50090  
PT50100  
PT50110  
PT50120  
PT50130  
PT50140  
PT50150  
PT50160  
PT50170

#### FORMATS

```

1 FORMAT(//1X,45(' '), ' INSTANTANEOUS FLOW FIELD INFORMATION ',45(' '
1))
2 FORMAT(//2X, 'LEVEL=',I3,3X, 'WEIGHTING FACTOR (MAX)=' ,I3,3X, 'WEDGE AN
1GLE =' ,F7.2, ' DEGREES',3X, ' RADIAL POSITION =' ,E11.3/2X, 'BOX# X PO
2SITION SAMP DENSITY MACH NO I VEL. Y VEL. Z VEL. T(KIN) T(RO
3T) TEMP.',14X, 'MOLE FRACTIONS')
3 FORMAT(//1X,46(' '), ' ACCUMULATED FLOW FIELD INFORMATION ',46(' '
1))
4 FORMAT(1X,I4,E11.3,I6,8F8.3,3X,3E11.3)

```

PT50220  
PT50230  
PT50240  
PT50250

```

*****
      IF(KS.EQ.0) WRITE(6,1)
      IF(KS.NE.0) WRITE(6,3)
      DO 40 MT=1,3
40 DBT(MT)=0.0
      FDA=0.
      CHT=0.
      DO 50 MT=1,MSP
      CHT=CHT+CHI(MT)*RNU(MT)
50 FDA=FDA+PDN(MT)*WTS(MT)
      YCT=0.0
      ZCT=0.0
      LEVEL=1

```

PT50290

```

DO 110 M=1,NBX
N=NUMCEL(M)
IF(N.LE.0) GO TO 110
IF((ZC(M).EQ.ZCT).AND.(YC(M).EQ.YCT)) GO TO 52
ZCT=ZC(M)
YCT=YC(M)
IF(M.GE.LEV(1)) LEVEL=2
IF(M.GE.LEV(2)) LEVEL=3
WRITE(6,2) LEVEL,LKW(N),ZCT,YCT
52 ICT=XC(M)
NSAMP=0
DBA=0.
XVM=0.
YVM=0.
ZVM=0.
TMPM=0.
TRPM=0.
E=0.
F=0.
DO 100 MT=1,MSP
NSAMP=NSAMP+NS(MT,N)
XVM=XVM+XV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
YVM=YVM+YV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
ZVM=ZVM+ZV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
DBA=DBA+DB(MT,N)*WTM(MT)
TMPM=TMPM+TMP(MT)*RNU(MT)*NB(MT,N)
TRPM=TRPM+TRP(MT)*RNU(MT)*NB(MT,N)
E=E+WTM(MT)*RNU(MT)*NB(MT,N)
100 F=F+RNU(MT)*NB(MT,N)
DBA=DBA/FDA
IF(E.EQ.0.0) GO TO 55
XVM=XVM/E
YVM=YVM/E
ZVM=ZVM/E
VS=XVM**2+YVM**2+ZVM**2
TMPM=TMPM/E-VS
TRPM=TRPM/F
55 CONTINUE
TTM=(TMPM+TRPM)/(2.5+CHT)
TMPM=TMPM/1.5
IF(CHT.NE.-1.) TRPM=TRPM/(1.+CHT)
AMS=SQRT(VS)
IF(TTM.GT.0.) AMS=SQRT((5.+2.*CHT)*VS/(TTM*(3.5+CHT)))
CCZ=COS(ZCT/57.29578)
SCZ=SQRT(1.-CCZ**2)
RVM=ZVM*SCZ-YVM*CCZ
TVM=YVM*SCZ+ZVM*CCZ
DO 60 MT=1,MSP
DBT(MT)=RNU(MT)*NB(MT,N)
IF(F.NE.0.) DBT(MT)=DBT(MT)/F
60 CONTINUE
WRITE(6,4) M,XCT,NSAMP,DBA,AMS,XVM,RVM,TVM,TMPM,TRPM,TTM,(DBT(J),
1J=1,3)
110 CONTINUE
RETURN

```

PT50470  
PT50480

FILE: GKBEIT DECK A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

```
END
/LKED.SYSLMOD DD DSN=U.GKBSPACE.ROTATION(FORTH),
/ DISP=OLD,UNIT=3350,VOL=SER=RES101,SPACE=
/LKED.SYSIN DD *
  INCLUDE SYSLMOD(FORTH)
  ENTRY MAIN
/ EXEC COMPRESS,DSN='U.GKBSPACE.ROTATION',RLSE=RLSE
```

PT50490

APPENDIX B

MASTER'S THESIS OF Y.P. TSAI

COLLISION INDUCED VIBRATIONAL TRANSITION  
PROBABILITIES IN DIATOMIC MOLECULES

## ABSTRACT

In order to improve the Monte-Carlo Direct Simulation calculations for hypersonic flow in the transition regime, we have to incorporate the effects due to vibrational non-equilibrium and potential dissociation for a diatomic gas. The state-to-state transition probabilities are desired. In this paper, we model the diatomic molecule as a harmonic oscillator which collides with another molecule collinearly. Two different methods have been developed, the first one is the semi-classical treatment and the second one is the fully quantum mechanical approach. The interaction potential between two molecules is assumed to be the Lennard-Jones 12-6 interaction law which is a small perturbation to the colliding system. Some numerical results of the state-to-state transition probabilities and comparisons are presented in Chapter 4. Discussions, which present the important aspects of this kind of problem for further study, are made in Chapter 5.

## Chapter 1 Introduction

The characteristic flow in a highly rarefied gas is called "free molecular flow". In this regime the mean free path is large compared to the characteristic dimensions of an aerodynamic body in the flow; and molecules that impinge on the body, and are then reemitted from it will, in general, be far away from the body before they strike another molecule. The characteristic flow in a moderately rarefied gas is called "slip flow". The flow regime intermediate between slip and molecular flow is known as the "transition flow regime". It corresponds to densities for which the mean free path has the same general order of magnitude as the characteristic dimension of the flow field. There is a dimensionless parameter called the Knudsen number  $Kn$ , introduced to serve as a criterion for determining the relative importance of these rarefaction effects.

where 
$$Kn = \frac{\lambda}{d} \sim \frac{M}{Re}$$

- and
- $\lambda$  = molecular mean free path
  - $d$  = characteristic dimension of vehicle
  - $M$  = Mach Number
  - $Re$  = Reynolds Number

Free molecular flow is usually defined as that flow for which  $Kn > 10$ . Slip flow is characterized by a Knudsen number of a few per cent,  $0.01 < Kn < 0.1$ , and the intermediate transition regime corresponds to Knudsen number in the range  $0.1 < Kn < 10$ .



These values are, of course, arbitrary. Since the Knudsen number is defined as the ratio of particle mean free path to body size, it therefore increases with altitude for a fixed size body travelling through the atmosphere.

Generally, for a body of the order of a meter at altitudes 150 Kilometers above the earth's surface, the mean free path of the particles is much larger than vehicle size (here the free stream mean free path  $\lambda_{\infty} \approx 40\text{m}$ ). Hence the particle-particle collision process in the vicinity of the vehicle need not be considered and the relation of measurement on the surface to atmospheric properties is explained in a relatively straightforward manner through free molecular theory. Below 90Km, (here the free stream mean free path  $\lambda_{\infty} \approx 2.5\text{ cm}$ ), the Knudsen number is so small that the fluid can be treated as a continuum with limited influence of transport properties and slip-flow boundary conditions. However, between these two zones, i.e. in the lower region of the thermosphere, is what we called the "transition flow regime" in which neither of the limiting theories is applicable.

Flight at very high altitudes often involves extremely high velocities and resulting high gas temperature. At velocities which correspond to effective temperature of the order of a few thousand degrees Kelvin, the so called "real gas" effects associated with vibration, dissociation, and ionization of the gas molecules can begin to be of importance. We are interested in atmospheric entry of a

vehicle such as the space shuttle. The Mach number is generally above 20 in passing through "transition flow regime" and the flow is hypersonic there. At such high speeds, any significant number of collisions between incoming particles and those reflected from the body can produce extreme changes in the environment near the surface of the vehicle. Measurements on the surface are thus strongly affected making inference of conditions in the ambient atmosphere extremely difficult. For the purpose of data interpretation, previous calculation based on direct simulation Monte-Carlo computer technique have been developed by G.A. Bird(1). By using the Monte-Carlo method, the real gas molecules are replaced by their statistical models, and the motion of one or more of the chosen particles is traced by the computer. In the "Bird" method the real gas molecules are simulated by several thousand modeled molecules, rigid spheres in the simplest version(1). Theoretical calculations of the heat transfer and aerodynamic characteristics of a body submerged in the transition flow regime may be carried out in this way. Modifications involving more realistic interaction laws have been also carried out. Molecules, however, contain internal structure, which is important primarily for its effect on the energy content of the flowing gas. Being composed of nuclei and electrons that have motion relative to the center of mass of the molecule, the molecules can possess rotational and vibrational as well as electronic

internal states. Therefore, translational energy is not necessarily conserved in all collisions. The hypersonic flow past the sharp leading edge of a flat plate incorporating the effects of rotational non-equilibrium for a diatomic gas was studied by D.I. Pullin, J.K. Harvey, and G.K. Bienkowski(2), and subsequently applied to other blunt body problems(3,4). A different model was used to incorporate the same effect in other works such as references 5 and 6. At the present stage, we want to improve the Monte-Carlo calculations by including the effects due to the vibrational degree of freedom of diatomic molecules.

It is convenient to divide vibrational energy exchange into two cases: the V-V process, in which the total vibrational quantum number of the system is unchanged, and the V-T process in which energy is exchanged between translation and vibration without conserving the vibrational quantum. For harmonic oscillators in the V-V process, the amount of vibrational energy lost by one molecule is gained by the other and no vibrational-translational energy transfer occurs. Considerable interest has been shown in the details of inelastic molecular collisions. The treatment of scattering between particles with internal structure is capable of producing differential and total cross section for state-to-state transitions. Our purpose is to calculate scattering cross sections for different transitions with known initial and final states as a function of collision energy or initial relative velocity. We emphasize here that

we need state-to-state transition cross sections rather than some overall "rates of transition" (rate coefficients) which are averaged over an equilibrium velocity distribution for the relative translational motion (e.g. Maxwellian distribution). The reasons are twofold:

1. We are dealing with a hypersonic flow system in a highly non-equilibrium state with a non-Maxwellian distribution of relative velocities of collision. The relative contributions of different energy molecules to the overall rates may therefore be drastically different than in the equilibrium state. This effect is accentuated by the steep rise of cross-sections with energy coupled to the generally decreasing magnitude of the distribution function with energy. This unifies that small changes in the fraction of molecules with high energies due to the non-equilibrium aspect of the flow can have extreme effects on inelastic processes without corresponding effects on mean properties such as density or fluid momentum.
2. The Direct Simulation Monte Carlo Code (DSMCC) consists of tracing a set of "test" molecules through a designated volume surrounding the body. The velocities and internal states of the simulated molecules are altered on the basis of collisions computed (as determined by local

collision probabilities) at fixed time steps. The positions are then advanced to new values on the basis of motion through the time step increments. This detailed computation of individual molecular collisions requires, in principle, state to state transition probabilities in order to incorporate the inelastic energy exchange into the computation of molecular velocities and internal states after collision. While the results we desire must ultimately be consistent with overall measured rates the level of detail necessary within the program is well beyond the level of availability of experimental data.

Early theoretical studies of vibrational, rotational, and translational energy transfer in collisions were based on approximate analytical solutions to the quantum mechanical and classical equations of motion. The method of Zener (7), later to become known as the distorted wave method, and the Born (8) approximation are leading examples of approximate solutions to quantum mechanical collision problems based on first order perturbation theory. A more detailed literature review on previous work in the field of vibrational collisions is given in the next Chapter.

Instead of doing an approximate treatment of a three-dimensional realistic system, in this work, we do an exact numerical treatment of a simpler model one-dimensional system

which has some important features in common with the real one. A calculation of transition probabilities for vibrational-vibrational-translational energy transfer in a collision of two diatomic molecules is to be presented. This simple collision model is approximate, utilizing a collinear collision of harmonic oscillators with an exponential repulsion between center atoms and no chemical reaction between the molecules. It may be argued that the configuration allowing the most efficient transfer of energy between translation and vibration is that in which the atoms are collinear. Collinear or head-on collisions make the most significant contribution to the transition probability. The averaged probability is equal to the probability of excitation in a head-on collision times a "steric factor" smaller than unity which takes account of unfavorable trajectories. For homonuclear molecules it is usually taken as  $\frac{1}{3}$  (the average of  $\cos^2\theta$  taken over a sphere). A more detailed theory for the steric factor of linear molecules has been propounded by Herzfeld (9). It is our belief that an accurate treatment of a collinear model is of more worth than an approximate result for the three-dimensional problem. The latter approach frequently contains errors which are difficult to estimate. The general magnitudes and trends of the transition probability obtained by this restricted treatment can show us qualitatively, or semi-quantitatively, some characteristic features of the problem.

Two different methods for calculating transition

probabilities are discussed in this paper. One is a semi-classical approach, and the other one is a fully quantum mechanical treatment. In the semi-classical calculation,, it is assumed that the vibrational amplitude of the harmonic oscillators are small and so the molecular oscillations do not greatly affect the external classical collision trajectory. The trajectory can be calculated from the classical equation of motion. The classical trajectory is assumed to define a time-dependent perturbation potential for the colliding system and quantum theory is used to derive the transition probability. Essentially, the Schrodinger equation is solved subject to certain initial conditions according to time-dependent perturbation theory. A detailed discussion about this theory can be found in the book Quantum Mechanics by Schiff (10).

In the quantum mechanical calculation of transition probability, the wavefunction of the whole system is expanded in terms of the complete set of eigenfunctions of vibrational states of the diatomic molecule. With the aid of the orthonormality property of these eigenfunctions, a set of coupled second order differential equations is obtained for the translational wavefunctions. The transition probability is given by the solution to this set of equations in the asymptotic region, subject to appropriate boundary condition. Several different numerical methods for solving the set of coupled equations have been developed by Diestler and Mckoy (11).

Riley and Kuppermann (12), and Gutshick et al (13). Solving a problem of quantum scattering between two diatomic molecules is then reduced to the task of finding a good numerical scheme for integrating a system of coupled differential equations accurately. In this work, we use IBM IMSL ROUTINE DGEAR to solve for the scattered wavefunction in the asymptotic region and then calculate the transition probability.

In Chapter 2, a brief review of previous semiclassical and quantum mechanical methods on vibrational collisions will be given so that we can identify the new points in the current work in Chapter 3. Chapter 3 is devoted to a discussion of the general theory. We derive some equations and expressions there which make a numerical algorithm feasible. In Chapter 4, theoretical results of state-to-state transition probability specifically for  $N_2-N_2$  collisions at different relative velocities are presented. Comparison is made with published results. Finally, we discuss some important problems related to inelastic molecular collisions which deserve further study because they make extension to a more realistic treatment of molecular scattering possible.



## Chapter 2 Literature Review

In 1931, Oldenberg (14) discussed molecular collision processes qualitatively to show the persistence of the rotational and vibrational motion. Zener (7) was the first to give a detailed mathematical treatment for collisions in which molecular vibrations are excited or de-excited. He restricted himself to collinear collisions between a diatomic molecule and an atom. His theory was based on the distorted wave method which includes direct transition from the initial state to the final state and assumes that the probability of transition is small. It is a perturbation method and cannot treat strongly coupled system. Takayanagi (15) then extended Zener's one-dimensional treatments to three-dimensional collisions. In order to save computational labor, the modified wave number approximation was introduced. Meanwhile, Schwartz, Slawsky and Herzfeld (16) gave a mathematical formulation for the vibrational transitions based on the distorted wave approximation due to Jackson and Mott (17), simplified by the modified wave number approximation due to Takayanagi, in diatom-diatom collisions. Their formulation is referred to as the SSH theory now and is a quantum mechanical result. For purpose of future reference, we describe SSH theory in more detail.

Consider the head-on collision between two diatomic molecules AB and CD (assume harmonic oscillators). For exponential intermolecular interaction between nearest

atoms B and C, one has,

$$\begin{aligned} V &= V_0 \exp(-\alpha r_{BC}) \\ &= V_0 \exp[-\alpha(R - \lambda_1 r_1 - \lambda_2 r_2)] \end{aligned} \quad (2-1)$$

where R is the distance between centers of mass of two molecules,  $r_1$  and  $r_2$  are vibrational coordinates,  $\frac{1}{\alpha}$  is the range of the potential and  $V_0$  is a constant.

$$\lambda_1 = \frac{m_A}{m_A + m_B}, \quad \lambda_2 = \frac{m_D}{m_C + m_D}$$

$m_i$  is the mass of the  $i$ -th atom. Solving the Schrodinger equation, the transition probability is given in the closed analytic form:

$$P(n_1 n_2 \rightarrow n_1' n_2') = \frac{\pi^2}{4} |V_1(n_1' n_1)|^2 |V_2(n_2' n_2)|^2 \left\{ \frac{\Delta q^2}{\cosh \pi q' - \cosh \pi q} \right\}^2 \cdot \sinh \pi q' \cdot \sinh \pi q \quad (2-2)$$

where

$$V_1(n_1' n_1) = \int_{-\infty}^{+\infty} z_1(n_1', r) \exp[\alpha \lambda_1 (r - r_{e1})] z_1(n_1, r) dr \quad (2-3)$$

$$V_2(n_2' n_2) = \int_{-\infty}^{+\infty} z_2(n_2', r) \exp[\alpha \lambda_2 (r - r_{e2})] z_2(n_2, r) dr \quad (2-4)$$

$z_i(n_j, r)$  is the vibrational wave function for harmonic oscillator "i" in quantum state  $n_j$ ;  $r_{e1}$  and  $r_{e2}$  are equilibrium separations for AB and CD respectively, and

$$q = q_{n_1 n_2} = \frac{2kn_1 n_2}{\alpha} \quad (2-5)$$

$$q' = q_{n_1' n_2'} = \frac{2kn_1' n_2'}{\alpha} \quad (2-6)$$

$$\Delta q^2 = q'^2 - q^2 = - \left( \frac{8\mu}{h^2 \alpha^2} \right) [\epsilon_1(n_1') + \epsilon_2(n_2') - \epsilon_1(n_1) - \epsilon_2(n_2)] \quad (2-7)$$

The wave number of relative motion before and after collision are  $k_{n_1 n_2}$  and  $k_{n_1' n_2'}$ , which satisfy the energy conservation law:

$$\epsilon_1(n_1) + \epsilon_2(n_2) + \frac{h^2 k_{n_1 n_2}^2}{2\mu} = \epsilon_1(n_1') + \epsilon_2(n_2') + \frac{h^2 k_{n_1' n_2'}^2}{2\mu} \quad (2-8)$$

$\mu$  is the reduced mass of the whole system;  $\epsilon_i(n)$  is the energy of molecule "i" in the n-th vibration state. For the special case, V-V transition, in which  $n_1 + n_2 = n_1' + n_2'$  and  $\Delta q = 0$ . Applying L'Hospital's rule, we get:

$$\lim_{q' \rightarrow q} \frac{q'^2 - q^2}{\cosh \pi q' - \cosh \pi q} = \frac{2q}{\pi \sinh \pi q}$$

Then the transition probability for V-V process becomes

$$P(n_1 n_2 \rightarrow n_1' n_2') = |V_1(n_1' n_1)|^2 |V_2(n_2' n_2)|^2 q^2 \quad (2-9)$$

SSH theory has been most widely used for quantitative comparison with experimental measurements of vibrational relaxation, but the coupling between rotation has been ignored.

Zelechow, Rapp and Sharp (ZRS) (18) have developed a semi-classical method for calculating transition probabilities for V-V and V-T energy transfer in a collision of two diatomic molecules. Their basic

assumption are:

- (1) The perturbation potential is linearized in the oscillator coordinates.
- (2) The collision velocity is not too high (e.g. the upper limit for  $N_2-N_2$  collision is 10Km/sec) and the collision induced time-varying force constant  $k'(t)$  is small compared to  $k$ , the characteristic force constant of the molecule.

Under these two conditions, Kerner (19) method can be applied to solve the Schrodinger equation and closed form analytical results are obtained. However, this approach restricts itself to the transitions of processes of symmetric type only. The general formula is:



where  $n, m$  are vibration quantum numbers before collision and  $n', m'$  are that after collision. The collision is symmetric in the sense that the two B atoms are in the center.

The development of the high speed electronic computer has made it possible to solve the collision problem by direct numerical techniques. T.E. Sharp and D. Rapp (20) have calculated the vibrational transition probabilities for collisions between a diatomic molecule and an atom. In their semi-classical treatment, an N-state approximation method is used, in which the total wave function is expanded in terms of N eigenfunctions of stationary states of the system including the initial, final and all

energetically intervening states. A Runge-Kutta single-step integration method is employed in the computation program. Generally, the value of  $N$  needed in expanding the total wave function increases with collision velocity. An "exact" solution for any transition probability  $P_{j \rightarrow k}$  is reached when the addition of more states to the computation results in no significant change in  $P_{j \rightarrow k}$ . We extend this method to collisions between two diatomic molecules (Chapter 3-A).

In quantum mechanical treatment of collisions between two diatomic molecules AB and CD, taking B and C as the inner atoms of the system, the total wave function is expanded in terms of normalized vibrational wave functions  $Z_{AB}(n_1, r_1)$  and  $Z_{CD}(n_2, r_2)$ . That is:

$$\psi = \sum_{n_1} \sum_{n_2} f_{n_1 n_2}(R) Z_{AB}(n_1, r_1) Z_{CD}(n_2, r_2)$$

Inevitably, we have to solve a system of coupled differential equations, which are equivalent to the Schrodinger equation, of the following form (detailed discussions will be given in Chapter 3-C).

$$\frac{\hbar^2}{2\mu} \frac{d^2}{dR^2} + k_{n_1' n_2'}^2 f_{n_1 n_2}(R) = \sum_{n_1} \sum_{n_2} \langle n_1' n_2' | V | n_1 n_2 \rangle f_{n_1 n_2}(R) \quad (2-10)$$

where  $\langle n_1' n_2' | V | n_1 n_2 \rangle$  is the matrix element. In principle if we can obtain the solution to equation (2-10) with the asymptotic form

$$\begin{aligned}
f_{n_1 n_2}'(R) &\rightarrow 0 & R \rightarrow -\infty \\
f_{n_1 n_2}'(R) &\rightarrow \delta_{n_1 n_1}' \delta_{n_2 n_2}' \exp(-ik_{n_1 n_2}' R) \\
&\quad + A_{n_1 n_2}' ; n_1 n_2 \exp(+ik_{n_1 n_2}' R) & R \rightarrow +\infty
\end{aligned} \tag{2-11}$$

the probability per collision for transition  $(n_1, n_2) \rightarrow (n_1', n_2')$  will be given by:

$$P(n_1 n_2 \rightarrow n_1' n_2') = \frac{k_{n_1' n_2'}}{k_{n_1 n_2}} |A_{n_1' n_2}' ; n_1 n_2|^2 \tag{2-12}$$

A number of numerical methods (7-9) have been proposed for solving the systems of equations (2-10). However, due to the rapidly oscillating wavelike solutions to the Schrodinger equation, the numerical technique is not straightforward. Riley (12) developed the initial-value technique with periodic "reorthogonalization". Gadschick et al. (13), on the other hand introduced a technique of integration using Dirichlet boundary condition and simple one-step Euler integration. A new method for constructing wave function for bound states and scattering has been proposed by Roy G. Gordon (21), perhaps this procedure can save much computer time. Our quantum mechanical treatment of this molecule scattering problem is similar to the method due to Riley and Kuppermann (12). It is relatively simple and straightforward, but in our procedures, the virtual states (energetically inaccessible) are not included in the total wavefunction expansion.

## Chapter 3 Theory

### A. General Formalism

The collision model is shown in Fig. 3-1.

Fig. 3-1 Collision Coordinates

This figure is the collinear collision configuration between two diatoms AB and CD. Assuming that CD is the target, and AB is the incident projectile from right. The laboratory coordinates of A, B, C and D are  $x_A$ ,  $x_B$ ,  $x_C$  and  $x_D$ ; their masses are  $m_A$ ,  $m_B$ ,  $m_C$  and  $m_D$  respectively. Let  $V_{AB}$  and  $V_{CD}$  be the binding potential of molecules AB and CD. The short range interaction is assumed to be a sum of interatomic interactions,

$$V_{INT} = V_{AC}(x_A - x_C) + V_{AD}(x_A - x_D) + V_{BD}(x_B - x_D) + V_{BC}(x_B - x_C)$$

The interatomic potentials are exponentially decreasing functions, so that for the collinear configuration under consideration, only the term  $V_{BC}(x_B - x_C)$  is important and thus

$$V_{INT} \approx V_{BC}(x_B - x_C)$$

The Schrodinger equation for the system is:

$$\left\{ -\frac{\hbar^2}{2M_A} \frac{\partial^2}{\partial x_A^2} - \frac{\hbar^2}{2M_B} \frac{\partial^2}{\partial x_B^2} - \frac{\hbar^2}{2M_C} \frac{\partial^2}{\partial x_C^2} - \frac{\hbar^2}{2M_D} \frac{\partial^2}{\partial x_D^2} + V_{AB}(x_A - x_B) + V_{CD}(x_C - x_D) \right. \\ \left. + V_{INT}(x_B - x_C) \right\} \psi(x_A, x_B, x_C, x_D) = E_{TOT} \psi(x_A, x_B, x_C, x_D) \quad (3-1)$$

We designate the distance between the centers of mass of two molecules as  $R$ . In molecule  $AB$  the distance between the atoms is  $x$ ; in  $CD$  it is  $y$ .  $x$  and  $y$  are internal coordinates. i.e.

$$R = \frac{m_A x_A + m_B x_B}{m_A + m_B} - \frac{m_C x_C + m_D x_D}{m_C + m_D}$$

$$x = x_A - x_B$$

$$y = x_C - x_D$$

Let  $R_{CM} = \frac{m_A x_A + m_B x_B + m_C x_C + m_D x_D}{m_A + m_B + m_C + m_D}$  which is the coordinate of

center of mass of the whole system. In terms of the new coordinates  $(x, y, R, R_{CM})$  the Schrodinger equation becomes:

$$\left\{ -\frac{\hbar^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{\hbar^2}{2\mu_{CD}} \frac{\partial^2}{\partial y^2} - \frac{\hbar^2}{2\mu} \frac{\partial^2}{\partial R^2} - \frac{\hbar^2}{2M} \frac{\partial^2}{\partial R_{CM}^2} + V_{AB}(x) + V_{CD}(y) \right. \quad (3-2)$$

$$\left. + V_{INT}(R - \gamma_{AB}x - \gamma_{CD}y) \right\} \psi(x, y, R, R_{CM}) = E_{TOT} \psi(x, y, R, R_{CM})$$

where  $M = m_A + m_B + m_C + m_D =$  total mass of the system.

$$\mu = \frac{(m_A + m_B)(m_C + m_D)}{M} = \text{reduced mass of the system}$$



$$\mu_{AB} = \frac{m_A m_B}{m_A + m_B} \quad = \text{reduced mass of the molecule AB}$$

$$\mu_{CD} = \frac{m_C m_D}{m_C + m_D} \quad = \text{reduced mass of the molecule CD}$$

$$\gamma_{AB} = \frac{m_A}{m_A + m_B}, \quad \gamma_{CD} = \frac{m_D}{m_C + m_D}$$

Since there is no external force applied to the system, the center of mass of the system moves like a free particle and its motion can be described by a plane wave  $\frac{h^2 k_{CM}^2}{2M} e^{i k_{CM} R_{CM}}$ , and the energy of the center of mass,  $T_{CM} = \frac{h^2 k_{CM}^2}{2M}$ , is also a constant of motion. This does not affect the energy transfer and need not be considered further. We can remove the  $R_{CM}$ -dependent part of the wave function  $(x, y, R, R_{CM})$  by separation of variables.

Let

$$\Psi(x, y, R, R_{CM}) = \psi(x, y, R) e^{i k_{CM} R_{CM}}, \quad E_{TOT} = E + T_{CM}$$

Substituting these into equation (3-2), we arrive at a Schrodinger equation concerning the internal coordinates  $x, y$  and the relative motion  $R$  of the two colliding molecules as follows:

$$\left\{ -\frac{\hbar^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{\hbar^2}{2\mu_{CD}} \frac{\partial^2}{\partial y^2} - \frac{\hbar^2}{2\mu} \frac{\partial^2}{\partial R^2} + V_{AB}(x) + V_{CD}(y) + V_{INT}(R - \gamma_{AB}x - \gamma_{CD}y) \right\}$$

$$\psi(x, y, R) = E\psi(x, y, R)$$

(3-3)

Define

$$X = x - x_{eq}$$

$$Y = y - y_{eq}$$

where  $x_{eq}$  and  $y_{eq}$  are equilibrium separations of AB and CD. Then  $x$  and  $y$  are displacements from equilibrium of each oscillator. Now we introduce harmonic bonds (intramolecular potential) into AB and CD with force constants  $k_{AB}$  and  $k_{CD}$ , hence

$$V_{AB} = 1/2 k_{AB} x^2,$$

$$V_{CD} = 1/2 k_{CD} y^2,$$

A conventional representation of the intermolecular potential energy curve is given by the Lennard-Jones 12-6 equation. Since the elementary models for energy transfer are based on exponential potential, the exponential function  $V_{INT}(x_B - x_C) - \epsilon e^{-(x_B - x_C)/L}$  must be fitted to the Lennard-Jones potential (Appendix 1), where  $L$  is a parameter characterizing the range of the interaction. Landau and Teller (22) assumed that only the short range repulsive part of the intermolecular potential is steep enough to influence energy transfer, so that the long-range attractive potential  $\epsilon$  can be neglected. The molecular interaction is then assumed to be an exponential repulsion between atoms C and B.

Let

$v$  = initial relative velocity

$E = 1/2 \mu v_0^2$  = initial relative kinetic energy

$$\bar{R} = R - R_T$$

where  $R_T$  is the distance at the classical turning point. The potential energy in equation (3-3) may be expressed as

(16):

$$V_{\text{INT}}(X, Y, \bar{R}) = E_0 \exp \left[ \frac{-1}{L} (\bar{R} - \gamma_{AB} X - \gamma_{CD} Y) \right] \quad (3-4)$$

Equation (3-3) can be written in the form:

$$\left\{ \frac{-h^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{h^2}{2\mu_{CD}} \frac{\partial^2}{\partial Y^2} - \frac{h^2}{2\mu} \frac{\partial^2}{\partial \bar{R}^2} + \frac{1}{2} k_{AB} X^2 + \frac{1}{2} k_{CD} Y^2 \right. \\ \left. + E_0 \exp \left[ \frac{-1}{L} (\bar{R} - \gamma_{AB} X - \gamma_{CD} Y) \right] \right\} \psi(X, Y, \bar{R}) = E \psi(X, Y, \bar{R}) \quad (3-5A)$$

or

$$\left\{ - \frac{h^2}{2\mu_{AB}} \frac{\partial^2}{\partial x^2} - \frac{h^2}{2\mu_{CD}} \frac{\partial^2}{\partial Y^2} + \frac{1}{2} k_{AB} X^2 + \frac{1}{2} k_{CD} Y^2 + E_0 \exp \left( \frac{-1}{L} (\bar{R} - \gamma_{AB} X - \gamma_{CD} Y) \right) \right\}$$

$$\psi(x, Y, \bar{R}) = ih \frac{\partial}{\partial t} \psi(x, Y, \bar{R}) \quad (3-5B)$$

We shall solve the time-independent Schrodinger equation (3-5A) by purely quantum mechanical method and the time-dependent Schrodinger equation (3-5B) by semi-classical method.

#### B. Semi-classical Calculation

Since the deBroglie wavelength of the relative motion of two molecules is usually very small compared with atomic dimensions (e.g. for  $N_2-N_2$  collisions at velocity 5Km/sec, the deBroglie wavelength is of the order  $10^{-15}$  cm, however the dimension of  $N_2$  molecules is of the order of a few Å), it is a fairly good approximation to use the classical trajectory for relative motion. The classical equations of motion are:

$$\mu \frac{d^2 \bar{R}}{dt^2} = - \frac{\partial}{\partial \bar{R}} V_{INT}(X, Y, \bar{R}) \quad (3-6A)$$

$$\mu_{AB} \frac{d^2 X}{dt^2} = - k_{AB} X - \frac{\partial}{\partial X} V_{INT}(X, Y, \bar{R}) \quad (3-6B)$$

$$\mu_{CD} \frac{d^2 Y}{dt^2} = - k_{CD} Y - \frac{\partial}{\partial Y} V_{INT}(X, Y, \bar{R}) \quad (3-6C)$$

Generally, the incident energy is much larger than the change in the vibrational energy, or in other words, due to small transition probabilities, only a small fraction of the translational energy is transferred to vibrational energy. We may then assume that during the collision, the vibrational amplitudes of the oscillators are not driven to large values, that means

$$\begin{aligned} X &\ll L, \\ Y &\ll L. \end{aligned} \quad (3-7)$$

In a series calculation given by Wolfberg and Kelley (23), we can see that conditions (3-7) are justifiable. Wolfberg and Kelley have calculated the energy transfer for collisions involving two harmonic oscillators via an exponential collision with  $L=0.22\text{\AA}$ . Other parameters and data are:  $m_A=m_B=m_C=m_D=12\text{a.m.u.}$ , angular frequency  $=2.3 \times 10^{14} \text{sec}^{-1}$  (cf: for  $N_2-N_2$  collisions  $m_A=m_B=m_C=m_D=14 \text{a.m.u.}$ ,  $=4.45 \times 10^{14} \text{sec}^{-1}$ ), the initial energy  $E_0 = 5.078 \text{ ev}$  (corresponds to  $v_0 = 9 \text{ Km/sec}$ ), then the vibrational energy transferred to each diatomic molecule is  $\Delta E_{CD} = 1.78 \times 10^{-3} \text{ ev}$ . Obviously, both  $\Delta E_{AB}$  and  $\Delta E_{CD}$  are much less than  $E_0$ .

The vibrational amplitude never exceeds  $0.007 \text{ \AA}$ , and since  $L \approx 0.22 \text{ \AA}$  the conditions (3-7) are fairly well satisfied. At low velocities (still high enough so that the deBroglie wavelength is much less than the atomic dimension), the energy transfer becomes smaller because the oscillator can readjust adiabatically to the perturbation caused by the incident particle. Conditions (3-7) are satisfied even better. Under conditions (3-7), the molecular oscillations do not greatly affect the external classical collision trajectory. Therefore, one can neglect the motion in X and Y in treating the motion in R. Equations (3-6A), (3-6B), and (3-6C) may then be replaced by:

$$\mu \frac{d^2 \bar{R}}{dt^2} = - \frac{\partial}{\partial \bar{R}} V_{\text{INT}}(\bar{R}, X=0, Y=0) \quad (3-8)$$

Solving equations (3-8) with  $V_{\text{INT}}$  given by (3-4), one finds that the trajectory  $R(t)$  satisfies the relation

$$\exp\left(-\frac{\bar{R}(t)}{L}\right) = \text{sech}^2\left(\frac{u_0 t}{2L}\right) \quad (3-9)$$

$R(t)$  is then inserted into the interaction potential function  $V_{\text{INT}}(X, Y, \bar{R})$  to obtain  $V_{\text{INT}}(X, Y, t)$ . In this semi-classical treatment,  $V_{\text{INT}}(X, Y, t)$  is used as a transition inducing perturbation acting upon a quantum mechanical harmonic oscillator. Finally, from equation (3-5B) we get the time-dependent Schrodinger equation to be

solved numerically by first order time-dependent perturbation theory.

$$\left\{ -\frac{\hbar^2}{2\mu_{AB}} \frac{\partial^2}{\partial X^2} + \frac{1}{2} k_{AB} X^2 - \frac{\hbar^2}{2\mu_{CD}} \frac{\partial^2}{\partial Y^2} + \frac{1}{2} k_{CD} Y^2 + E_0 \exp\left(\frac{\gamma_{AB} X + \gamma_{CD} Y}{L}\right) \right. \quad (3-10)$$

$$\left. \operatorname{sech}^2\left(\frac{u_C t}{2L}\right) \right\} \psi(X, Y, t) = i\hbar \frac{\partial}{\partial t} \psi(X, Y, t)$$

There is one more comment about the assumptions of this semi-classical model. In principle, if the deBroglie wavelength associated with the relative motion is much less than the atomic dimensions, it is a good approximation to use classical mechanics for solving the relative trajectory. As just mentioned before, if the criteria (3-7) are fulfilled, in other words, each of the oscillation amplitudes of the two colliding diatomic molecules is very small compared to the characteristic range of the interaction potential, a further simplification can be made in solving for  $\bar{R}(t)$ , the relative trajectory. These procedures correspond to the "approximate" classical method, because the trajectory  $\bar{R}(t)$  is determined with  $X$  and  $Y$  set equal to zero. As a result, these calculations do not include the conservation of energy, and  $E$  is assumed to remain as the energy in coordinate  $\bar{R}$ , regardless of how much excitation occurs in the oscillators. For more accurate semi-classical calculations, one must solve equations (3-6A), (3-6B), and (3-6C) for "exact" classical trajectory  $\bar{R}(t)$  to be used as a time-dependent perturbation. Prior to the work of Wolfsberg and

Kelley (23), it has been thought that the criteria (3-7) are automatically satisfied for all low-velocity collisions. Actually, this is not the case Wolfsberg and Kelley proved that the approximate classical method should be limited to collisions between a light particle and a heavy oscillator.

For our case of  $N_2-N_2$  collision, Wolfsberg and Kelley's requirement is satisfied. The energy transferred to each diatomic species is very small compared to the kinetic energy of relative motion. Both molecules are negligibly distorted during collision approach and the molecular oscillations never deviate substantially from their equilibrium configurations. The semi-classical approach is a good approximation even within the high velocity range in which we are interested. The reason is that as  $v_0$  increases, the deBroglie wavelength characterizing the relative motion in  $R$  is small compared to the distance over which the interaction potential varies significantly in  $R$ , i.e.  $L.k_{n1n2} \gg 1$ . This implies that both  $q$  and  $q'$  defined by equations (2-5) and (2-6) are much greater than unity. In this limit,  $q \gg 1$  and  $q' \gg 1$ , the quantum mechanical results of equation (2-2) are reduced to the classical results of Landau and Teller (22).

Let  $\omega_{AB}$  and  $\omega_{CD}$  be the angular frequencies for the AB and CD molecules, their isolated Hamiltonians are  $H^{(AB)}(X)$  and  $H^{(CD)}(Y)$  respectively, then

$$\omega_{AB} = \frac{k_{AB}}{\mu_{AB}}$$

$$\omega_{CD} = \frac{k_{CD}}{\mu_{CD}}$$

$$H^{(AB)}(X) = -\frac{\hbar^2}{2\mu_{AB}} \frac{\partial^2}{\partial X^2} + \frac{1}{2} \mu_{AB} \omega_{AB}^2 X^2$$

$$H^{(CD)}(Y) = -\frac{\hbar^2}{2\mu_{CD}} \frac{\partial^2}{\partial Y^2} + \frac{1}{2} \mu_{CD} \omega_{CD}^2 Y^2$$

Suppose the individual eigenfunctions of  $H^{(AB)}(X)$  and  $H^{(CD)}(Y)$  are  $\phi_n^{(AB)}(X)$  and  $\phi_j^{(CD)}(Y)$ , we have:

$$H^{(AB)}(X) \phi_n^{(AB)}(X) = (n + \frac{1}{2}) \hbar \omega_{AB} \phi_n^{(AB)}(X) \quad n = 0, 1, 2, \dots$$

$$H^{(CD)}(Y) \phi_j^{(CD)}(Y) = (j + \frac{1}{2}) \hbar \omega_{CD} \phi_j^{(CD)}(Y) \quad j = 0, 1, 2, \dots$$

Set

$$\phi_{nj}(X, Y) = \phi_n^{(AB)}(X) \phi_j^{(CD)}(Y) \quad (3-11)$$

$$W_{ij} = (n + \frac{1}{2}) \hbar \omega_{AB} + (j + \frac{1}{2}) \hbar \omega_{CD} \quad (3-12)$$

$$H_0(X, Y) = H^{(AB)}(X) + H^{(CD)}(Y) \quad (3-13)$$

$H_0(X, Y)$  is the unperturbed Hamiltonian of the system. It is obvious that  $H_0(X, Y) \phi_{nj}(X, Y) = W_{nj} \phi_{nj}(X, Y)$  (3-14)  $W_{ij}$  are eigenvalues of the unperturbed Hamiltonian  $H_0(X, Y)$ . The solution to the equation (3-10), i.e. the total wavefunction for the system of two oscillators  $\psi(X, Y, t)$  can be expanded in terms of the individual harmonic oscillator wavefunctions  $\phi_n^{(AB)}(X)$  and  $\phi_j^{(CD)}(Y)$ .



$$\psi(X, Y, t) = \sum_n \sum_j a_{nj}(t) \phi_n^{(AB)}(X) \phi_j^{(CD)}(Y) e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \quad (3-15)$$

where the expansion coefficients  $a_{nj}$  depend on time. If the oscillators AB and CD are initially in state N and J respectively, then the initial conditions are:

$$n = 0, 1, 2, \dots$$

$$a_{nj}(-\infty) = \delta_{nN} \delta_{jJ}$$

(Appendix 2)

$$j = 0, 1, 2, \dots$$

The probability of the system ending up with AB in state Q and CD in state K is

$$P_{NJ \rightarrow QK} = |a_{QK}(+\infty)|^2$$

We are now in the position to solve for the expansion coefficients  $a(t)$ . Substituting equation (3-15) into Schrodinger equation (3-10), and using the relations (3-11)  $\rightarrow$  (3-14), we find:

$$\begin{aligned} & (i\hbar) \sum_{nj} \frac{d}{dt} a_{nj}(t) \phi_n^{(AB)}(X) \phi_j^{(CD)}(Y) e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \\ & = E_0 \exp\left(\frac{\gamma_{AB}X + \gamma_{CD}Y}{L}\right) \operatorname{sech}^2\left(\frac{u_0 t}{2L}\right) \sum_{nj} a_{nj}(t) \phi_n^{(AB)}(X) \phi_j^{(CD)}(Y) \\ & \quad e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \end{aligned} \quad (3-16)$$

Equation (3-16) is multiplied by  $\phi_{n'}^{(AB)*}(X) \phi_{j'}^{(CD)*}(Y)$  on both side, where  $\phi_{n'}^{(AB)*}$  is the complex conjugate function of  $\phi_{n'}^{(AB)}$  and  $\phi_{j'}^{(CD)*}(Y)$  is the complex conjugate of  $\phi_{j'}^{(CD)}(Y)$ , and then integrated. Equation (3-16) then becomes:

$$(ih) \sum_{nj} \frac{d}{dt} a_{nj}(t) \delta_{nn'} \delta_{jj'} e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t}$$

$$= E_0 \operatorname{sech}^2\left(\frac{u_0 t}{2L}\right) \sum_{nj} a_{nj}(t) \left[ \int_{-\infty}^{+\infty} \phi_{n'}^{(AB)*}(X) e^{\gamma_{AB}X} \phi_n^{(AB)}(X) dX \right] \quad (3-17)$$

$$\left[ \int_{-\infty}^{+\infty} \phi_j^{(CD)*}(Y) e^{\gamma_{CD}Y} \phi_j^{(CD)}(Y) dY \right] e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t}$$

where use has been made of the orthonormal property for the set of harmonic oscillator energy eigenfunctions  $\phi_n^{(AB)}(X)$  and  $\phi_j^{(CD)}(Y)$ .

$$\int_{-\infty}^{+\infty} \phi_{n'}^{(AB)*}(X) \phi_n^{(AB)}(X) dX = \delta_{nn'}$$

$$\int_{-\infty}^{+\infty} \phi_{j'}^{(CD)*}(Y) \phi_j^{(CD)}(Y) dY = \delta_{jj'}$$

Let 
$$U_{n'n} = \int_{-\infty}^{+\infty} \phi_{n'}^{(AB)*}(X) e^{\gamma_{AB}X} \phi_n^{(AB)}(X) dX$$

$$V_{j'j} = \int_{-\infty}^{+\infty} \phi_{j'}^{(CD)*}(Y) e^{\gamma_{CD}Y} \phi_j^{(CD)}(Y) dY$$

$U_{n'n}$  and  $V_{j'j}$  are the coupling terms between states, called matrix elements. Equations (3-17) can be written now as:

$$\frac{d}{dt} a_{n'j'}(t) = \left(\frac{E_0}{i\hbar}\right) \operatorname{sech}^2\left(\frac{u_0 t}{2L}\right) \sum_{nj} a_{nj}(t) U_{n'n} V_{j'j} e^{-i(n+\frac{1}{2})\omega_{AB}t} e^{-i(j+\frac{1}{2})\omega_{CD}t} \quad (3-18)$$

Equation (3-18) is a set of coupled first order differential equations subject to certain initial conditions. Actually, this set of differential equations is equivalent to the Schrodinger equation (3-10). Next, let

$$a_{nj}(t) = A_{nj}(t) + iB_{nj}(t) \quad (3-19)$$

where  $A_{nj}(t)$  is the real part of  $a_{nj}(t)$ , and  $B_{nj}(t)$  is the imaginary part of  $a_{nj}(t)$ . Both  $A_{nj}(t)$  and  $B_{nj}(t)$  are real functions of time  $t$ . We also know that  $U_{m'n}$ ,  $V_{jj}$  are real, and

$$e^{i(n+\frac{1}{2})\omega_{AB}t} e^{i(j+\frac{1}{2})\omega_{CD}t} = \cos [(n+\frac{1}{2})\omega_{AB}t + (j+\frac{1}{2})\omega_{CD}t] + i \sin [(n+\frac{1}{2})\omega_{AB}t + (j+\frac{1}{2})\omega_{CD}t] \quad (3-20)$$

Inserting equations (3-19) and (3-20) into equation (3-18), with some algebra, the real part and imaginary part on each side of the equation must be equal separately. So we obtain a set of coupled differential equations for A's and B's:

$$\frac{d}{dt} A_{n',j'}(t) = \frac{E_0}{h} \operatorname{sech}^2 \left( \frac{u_0 t}{2L} \right) \sum_{nj} \{ A_{nj}(t) \sin [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] + B_{nj}(t) \cos [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \} U_{n'n} V_{j'j} \quad (3-21)$$

$$\frac{d}{dt} B_{n',j'}(t) = \frac{E_0}{h} \operatorname{sech}^2 \left( \frac{v_0 t}{2L} \right) \sum_{nj} \{ B_{nj}(t) \sin [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] - A_{nj}(t) \cos [(n'-n)\omega_{AB}t + (j'-j)\omega_{CD}t] \} U_{n'n} V_{n'j'}$$

$A_{rs}(+\infty)$  and  $B_{rs}(+\infty)$  are to be found, hence

$$|a_{rs}(+\infty)|^2 = |A_{rs}(+\infty)|^2 + |B_{rs}(+\infty)|^2,$$

the desired transition probability.

In Chapter 2, we mentioned that in the work of ZRS (18), they treated collinear collisions between two

diatomic molecules with symmetric configuration only. Therefore,  $\gamma_{AB}$  is equal to  $\gamma_{CD}$  i.e.:

$$\gamma_{AB} = \gamma_{CD} = \gamma$$

The intermolecular potential function in equation (3-10),  $E_0 \text{sech}^2 \left( \frac{v_0 t}{2L} \right) \exp \left( \frac{\gamma_{AB} X + \gamma_{CD} Y}{L} \right)$  is expanded in a Taylor series to the second order of X and Y;

$$E_0 \text{sech}^2 \left( \frac{v_0 t}{2L} \right) \exp \left[ \frac{\gamma}{L} (x+Y) \right] = E_0 \text{sech}^2 \left( \frac{v_0 t}{2L} \right) \left\{ 1 + \frac{\gamma}{L} (X+Y) + \frac{1}{2} \left( \frac{\gamma}{L} \right)^2 (x+Y)^2 \right\} \quad (3-22)$$

neglecting high order terms in X,Y. Equation (3-10), with interaction potential given by equation (3-22), can be solved analytically by applying Kerner (19) method. Our semi-classical approach differs from the method of ZRS in that we do not expand the potential function  $E_0 \text{sech}^2 \left( \frac{v_0 t}{2L} \right) \exp \left( \frac{\gamma_{AB} X + \gamma_{CD} Y}{L} \right)$  into a Taylor series. We make a direct numerical integration of the coupled equations (3-21), which are equivalent to equation (3-10). It is appropriate to carry out classical calculations in which the approximate equation (3-8) is used, but the potential expansion (3-22) is not carried out. Numerical comparison between these two methods for  $N_2-N_2$  collision will be given in Chapter 4.

### C. Quantum Mechanical Calculation

To clarify the presentation of the theory, we shall use a scaled Schrodinger equation. Define (13):

$$y^* = \left( \frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} y \quad (3-23A)$$

$$x^* = \left( \frac{\mu_{AB} k_{AB}}{h^2} \right)^{\frac{1}{4}} x \quad (3-23B)$$

Inserting equations (3-23A) and (3-23B) into equation (3-5A), it becomes:

$$\begin{aligned} & \left\{ \frac{1}{2} h\omega_{CD} \left( -\frac{\partial^2}{\partial y^{*2}} + y^{*2} \right) + \frac{1}{2} h\omega_{AB} \left( -\frac{\partial^2}{\partial x^{*2}} + x^{*2} \right) - \frac{h^2}{2\mu} \frac{\partial^2}{\partial \bar{R}^2} \right. \\ & \left. + E_0 \exp \left[ \frac{-1}{L} \gamma_{CD} \left( \frac{h^2}{\mu_{CD} k_{CD}} \right)^{\frac{1}{4}} \left( \frac{1}{\gamma_{CD}} \left( \frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} \bar{R} - y^* - \frac{\gamma_{AB}}{\gamma_{CD}} \left( \frac{\mu_{CD} k_{CD}}{\mu_{AB} k_{AB}} \right)^{\frac{1}{4}} x^* \right) \right] \right\} \\ & \psi(x^*, y^*, \bar{R}) = E\psi(x^*, y^*, \bar{R}) \quad (3-24A) \end{aligned}$$

Dividing both sides by  $1/2h_{CD}$ , equation (3-24) looks simpler. If we define:

$$\begin{aligned} \omega &= \frac{\omega_{AB}}{\omega_{CD}}, \quad E^* = \frac{1}{2} \frac{E}{h\omega_{CD}}, \quad r^* = \frac{1}{\gamma_{CD}} \left( \frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{4}} \bar{R} \\ L^* &= \frac{1}{L} \gamma_{CD} \left( \frac{h^2}{\mu_{CD} k_{CD}} \right)^{\frac{1}{4}} = \frac{1}{L} \frac{m_D}{m_C + m_D} \left( \frac{h^2}{\mu_{CD} k_{CD}} \right)^{\frac{1}{4}} \\ \frac{1}{m} &= \frac{h^2}{\mu\omega_{CD}} \left( \frac{\mu_{CD} k_{CD}}{h^2} \right)^{\frac{1}{2}} \frac{1}{\gamma_{CD}} = \frac{\mu_{CD}}{\mu} \left( \frac{m_C + m_D}{m_D} \right)^2 = \frac{m_C M}{(m_A + m_B) m_D} \\ \beta &= \frac{\gamma_{AB}}{\gamma_{CD}} \left( \frac{k_{CD} \mu_{CD}}{k_{AB} \mu_{AB}} \right)^{\frac{1}{4}} = \frac{m_C}{m_B} \left( \frac{\mu_{AB}}{\mu_{CD}} \right)^{\frac{1}{2}} \frac{1}{\omega^{1/2}} \end{aligned}$$

The new equation (3-24) is dimensionless, namely;

$$\left\{ \left( -\frac{\partial^2}{\partial y^{*2}} + y^{*2} \right) + \left( -\frac{\partial^2}{\partial x^{*2}} + x^{*2} \right) - \frac{1}{\bar{m}} \frac{\partial^2}{\partial j^{*2}} + \left( \frac{E_0}{\frac{1}{2}h\omega_{CD}} \right) \exp[-L^*(r^*-y^*-\beta x^*)] \right\} \psi(x^*, y^*, r^*) = E^* \psi(x^*, y^*, r^*) \quad (3-24B)$$

Since the constant coefficient  $\frac{E_0}{\frac{1}{2}h\omega_{CD}}$  can be absorbed into the argument of the exponential function  $\exp(-L^*(r^*-y^*-\beta x^*))$ , and the operator  $\frac{\partial}{\partial r^{*2}}$  is invariant under the transformation

$$r^* \longrightarrow r^* + \text{any constant,}$$

we can choose  $\frac{E_0}{\frac{1}{2}h\omega_{CD}} = 1$  with no loss of generality. The notation  $x^*$ ,  $y^*$  and  $r^*$  refer to dimensionless variables, we now drop the \*\*\* for convenience. Hence the scaled Schrodinger equation is:

$$\left\{ \left( -\frac{\partial^2}{\partial y^2} + y^2 \right) + \omega \left( -\frac{\partial^2}{\partial x^2} + x^2 \right) - \frac{1}{\bar{m}} \frac{\partial^2}{\partial r} + \exp[-L^*(r-y-\beta x)] \right\} \psi(x, y, r) = E^* \psi(x, y, r) \quad (3-25)$$

The system is specified by the five dimensionless parameters. They are  $\omega$ ,  $\bar{m}$ ,  $\beta$ ,  $L^*$  and  $E^*$ ,  $E^*$  is energy of the system in term of ground state energy of molecule CD. The dimensionless form of the unperturbated Hamiltonian  $H_0$  is:

$$H_0(x, y) = \left( -\frac{\partial^2}{\partial y^2} + y^2 \right) + \omega \left( -\frac{\partial^2}{\partial x^2} + x^2 \right)$$

with eigenfunctions  $\phi_{ij}(x, y)$  and eigenvalues  $W_{ij}$ , then

$$H_0(x, y) \phi_{ij}(x, y) = W_{ij} \phi_{ij}(x, y) \quad \begin{array}{l} i = 0, 1, 2, \dots \\ j = 0, 1, 2, \dots \end{array}$$

where  $\phi_{ij}(x,y) = \phi_i(y)\phi_j(x)$

$$W_{ij} = (2i+1) + (2j+1) \omega$$

$\phi_{ij}(x,y)$  is the product of the individual harmonic oscillator wavefunction which indicates that molecule AB with internal coordinate  $x$  is the vibrational state  $j$ , and molecule CD with internal coordinate  $y$  is in the state  $i$ . Let the system be in a particular initial state  $(n_0, m_0)$ . We can expand the total stationary scattering wavefunction  $\psi_{n_0 m_0}(x,y,r)$ , i.e. the solution to equation (3-25) in terms of  $\phi_{nm}(x,y)$  because they form a complete set.

$$\psi_{n_0 m_0}(x,y,r) = \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) \quad (3-26)$$

Substituting equation (3-26) into equation (3-25) get:

$$\begin{aligned} & \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} \{ W_{nm} f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) - \frac{1}{m} \phi_n(y) \phi_m(x) \frac{d^2}{dr^2} f_{nm, n_0 m_0}(r) \\ & + \exp[-L(r-y-\beta x)] f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) \} \\ & = E^* \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{nm, n_0 m_0}(r) \phi_n(y) \phi_m(x) \end{aligned} \quad (3-27)$$

where we have used the relation

$$H_0(x,y) \phi_{ij}(x,y) = W_{ij} \phi_{ij}(x,y)$$

When equation (3-27) is multiplied by  $\phi_n^*(y)\phi_m^*(x)$  on both sides and integrated over  $x$  and  $y$ , we have:

$$\begin{aligned} & \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} \{ W_{nm} f_{nm, n_0 m_0}(r) \langle \phi_{n',m'}(x,y) \rangle - \frac{1}{m} \frac{d^2}{dy^2} f_{nm, n_0 m_0}(r) \langle \phi_{n',m'}(x,y) | \phi_{nm}(x,y) \rangle \\ & + f_{n_0 m_0, nm}(r) \langle \phi_{n',m'}(x,y) | \exp[-L(r-y-\beta x)] | \phi_{nm}(x,y) \rangle \} \quad (3-28) \\ & = E^* \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} f_{nm, n_0 m_0}(r) \langle \phi_{n',m'}(x,y) | \phi_{nm}(x,y) \rangle \end{aligned}$$

$$\begin{aligned} \text{where } \langle \phi_{n'm'}(x,y) | \phi_{nm}(x,y) \rangle &\equiv \int_{-\infty}^{+\infty} \phi_{n'}(y) \phi_n(y) dy \int_{-\infty}^{+\infty} \phi_{m'}(x) \phi_m(x) dx \\ &= \delta_{n'm'} \delta_{m'm} \end{aligned} \quad (3-29)$$

$$\begin{aligned} &\langle \phi_{n'm'}(x,y) | \exp [-L^*(r-y-\beta x)] | \phi_{nm}(x,y) \rangle \\ &= \exp (-L^* r) \int_{-\infty}^{+\infty} \phi_{n'}(y) \exp (L^* y) \phi_n(y) dy \int_{-\infty}^{+\infty} \phi_{m'}(x) \exp (L^* \beta x) \phi_m(x) dx \end{aligned}$$

$$\text{Defining: } V_{n'm',nm}(r) = \bar{m} \langle \phi_{n'm'}(x,y) | \exp [-L^*(r-y-\beta x)] | \phi_{nm}(x,y) \rangle \quad (3-30)$$

and using equation (3-29) and (3-30), equation (3-28)

becomes

$$\begin{aligned} \bar{m} W_{n'm',n_0m_0}(r) - \frac{d^2}{dr^2} f_{n'm',n_0m_0}(r) + \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} V_{n'm',nm}(r) f_{nm,n_0m_0}(r) \\ = \bar{m} E^* f_{n'm',n_0m_0}(r) \end{aligned}$$

Rearranging this equation, we get

$$\begin{aligned} \frac{d^2}{dr^2} f_{n'm',n_0m_0}(r) = \sum_{n=0}^{N-1} \sum_{m=0}^{M-1} V_{n'm',nm}(r) f_{nm,n_0m_0}(r) \\ - \bar{m} (E^* - W_{n'm'}) f_{n'm',n_0m_0}(r) \end{aligned} \quad (3-31)$$

Define

$$\begin{aligned} k_{ij} &= \sqrt{\bar{m} (E^* - W_{ij})} \\ &= \sqrt{\bar{m} (E - (2i+1) - (2j+1))} \end{aligned} \quad (3-32)$$

N and M are the number of states of CD and AB included in the expansion. We introduce (9)

$$i = n + m.N \quad n = 0, 1, 2, \dots, N-1 \quad m = 0, 1, \dots, M-1$$

$$j = n + m.N \quad n = 0, 1, 2, \dots, N-1 \quad m = 0, 1, \dots, M-1$$

$$k = n + m.N \quad n = 0, 1, 2, \dots, N-1 \quad m = 0, 1, \dots, M-1$$

to indicate the states of the system. For example,  $i = (n,m)$ ,  $j = (n,m)$  and  $k = (n,m)$  etc. By incorporating the



definition of  $k$ , equation (3-32), we obtain a system of coupled second order ordinary differential equations in matrix form which is equivalent to the Schrodinger equation (3-32).

$$\frac{d^2}{dr^2} \vec{F}(r) = (\vec{V} - \vec{K}^2) \vec{F} \quad (3-33)$$

where  $(\vec{F}(r))_{ij} = (\vec{F}(r))_{i(n,m)j(n'm')} = f_{nm,n'm'}(r)$

$$(\vec{V}(r))_{ij} = (\vec{V}(r))_{i(n,m)j(n'm')} = V_{nm,n'm'}(r)$$

and  $(\vec{K})_{ij} = (\vec{K})_{i(n,m)j(n'm')} = k_{nm} \delta_{mn'} \delta_{mm'}$

In the asymptotic region where  $r$  is very large,  $V(r)$  tends to zero, so integrating equation (3-33), we get

$$\lim_{r \rightarrow \infty} \vec{F}(r) = e^{-i\vec{K}r} \vec{G} + e^{i\vec{K}r} \vec{J} \quad (3-34)$$

Equation (3-34) is the asymptotic form of  $\vec{F}(r)$  at large  $r$ . In principle if  $\vec{G}$  and  $\vec{J}$  are determined, then the transition probability from state  $i = i(n_0, m_0)$  to state  $j = j(n, m)$  is given by:

$$P_{ij} = |(\vec{J} \vec{G}^{-1})_{ji}| \frac{k_j}{k_i} \quad (3-35)$$

where  $\vec{G}^{-1}$  is the inverse matrix of  $\vec{G}$ . However, it is rather difficult to find matrices  $\vec{G}$  and  $\vec{J}$  in a straightforward manner. We go to the following alternative way:

Set 
$$\frac{d}{dr} \vec{F}(r) = \vec{E}(r) \quad (3-36)$$

$$\frac{d}{dr} \vec{E}(r) = (\vec{V}(r) - \vec{K}^2) \vec{F}(r)$$

Letting  $r_0$  be some point in the asymptotic regime, the asymptotic form of  $\vec{F}(r)$  may be also written as

$$\lim_{r \rightarrow \infty} \vec{F}(r) = e^{-i\vec{K}(r-r_0)} \vec{G} + e^{i\vec{K}(r-r_0)} \vec{J} \quad (3-37)$$

and then,

$$\lim_{r \rightarrow \infty} \vec{E}(r) = \lim_{r \rightarrow \infty} \frac{d\vec{F}}{dr} = -i\vec{K}e^{-i\vec{K}(r-r_0)} \vec{G} + i\vec{K}e^{i\vec{K}(r-r_0)} \vec{J}.$$

So, as the point  $r = r_0$  in the asymptotic region, it is obvious that:

$$\vec{F}(r_0) = \vec{G} + \vec{J} \quad (3-38A)$$

$$\vec{E}(r_0) = -i\vec{K}\vec{G} + i\vec{K}\vec{J} \quad (3-38B)$$

These two relations will be used later. From equation (3-37) we have:

$$\lim_{r \rightarrow \infty} \vec{F}(r) \vec{G}^{-1} e^{-i\vec{K}r_0} = e^{-i\vec{K}r} + e^{i\vec{K}r} (e^{-i\vec{K}r_0} \vec{J} \vec{G}^{-1} e^{-i\vec{K}r_0}),$$

then the transition probabilities are:

$$P_{ij} = |(e^{-i\vec{K}r_0} \vec{J} \vec{G}^{-1} e^{-i\vec{K}r_0})_{ji}|^2 \frac{k_j}{k_i} \quad (3-39)$$

Since  $k_j$  is real for open channels, that is to say the incoming particle has sufficient energy to excite the bound particle to any of its lowest  $N$  eigenstates, matrix  $\vec{K}$  is real too, and these are the only observed ones. In equation (3-39) since  $\vec{K}$  is real, we end up with

$$P_{ij} = (\vec{J} \vec{G}^{-1})_{ji} \frac{k_j}{k_i}$$

This result is exactly the same as the probabilities obtained based on the asymptotic form of  $\vec{F}(r)$  in equation (3-34). We can find  $\vec{G}$  and  $\vec{J}$  in terms of  $\vec{F}(r_0)$ , and  $\vec{E}(r_0)$  with no difficulty, since  $\vec{G}$  and  $\vec{J}$  are related to  $\vec{F}(r_0)$ ,  $\vec{E}(r_0)$  by equation (3-38A) and (3-38B). Solving for  $\vec{G}$  and

we obtain:

$$2\vec{G} = \vec{F}(r_0) + i\vec{K}^{-1}\vec{E}(r_0)$$

$$2\vec{J} = \vec{F}(r_0) - i\vec{K}^{-1}\vec{E}(r_0)$$

Defining

$$\vec{S} = \text{Re } \vec{S} + i \cdot \text{Im } \vec{S} = \vec{J}\vec{G}^{-1}.$$

where  $\text{Re } \vec{S}$  is the real part of  $\vec{S}$  and  $\text{Im } \vec{S}$  is the imaginary part of  $\vec{S}$ , (both are real) then

$$|\vec{J}\vec{G}^{-1}|^2 = |\vec{S}|^2 = (\text{Re } \vec{S})^2 + (\text{Im } \vec{S})^2$$

We must find  $\text{Re } \vec{S}$  and  $\text{Im } \vec{S}$  in terms of  $\vec{F}(r_0)$ ,  $\vec{E}(r_0)$  and  $\vec{K}$ .

We know that (19) given three matrices  $\vec{A}$ ,  $\vec{X}$ , and  $\vec{Y}$  such that

$$\vec{A} = \vec{X} + i\vec{Y}$$

where  $\vec{X}$  and  $\vec{Y}$  are real, if

$$\begin{pmatrix} X & Y \\ -Y & X \end{pmatrix}^{-1} = \begin{pmatrix} Z & W \\ -W & Z \end{pmatrix}$$

then the inverse matrix of A is

$$\vec{A}^{-1} = \vec{Z} + i\vec{W}.$$

In our case  $\vec{F}(r_0)$ ,  $\vec{E}(r_0)$  and  $\vec{K}$  are all real, we may define

$$\vec{D} = \vec{K}^{-1}\vec{E}(r_0)$$

so that equation (3-40) becomes:

$$2\vec{G} = \vec{F}(r_0) + i\vec{D}$$

$$2\vec{J} = \vec{F}(r_0) - i\vec{D}$$

Here we want to find the inverse of  $\vec{G}$ , then set:

$$\begin{pmatrix} \vec{F}(r_0) & \vec{D} \\ -\vec{D} & \vec{F}(r_0) \end{pmatrix}^{-1} = \begin{pmatrix} \vec{Z} & \vec{W} \\ -\vec{W} & \vec{Z} \end{pmatrix} \quad (3-41)$$

and solve for  $\vec{Z}$ ,  $\vec{W}$  in terms of  $\vec{F}(r_0)$  and  $\vec{D}$ . Equation (3-41) means that

$$\begin{pmatrix} \vec{F}(r_0) & \vec{D} \\ -\vec{D} & \vec{F}(r_0) \end{pmatrix} \begin{pmatrix} \vec{Z} & \vec{W} \\ -\vec{W} & \vec{Z} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Or equivalently

$$\vec{F}(r_0)\vec{Z} - \vec{D}\vec{W} = 1$$

$$\vec{F}(r_0)\vec{W} + \vec{D}\vec{Z} = 0$$

From the second equation of (3-42),

$$\vec{W} = -(\vec{F}(r_0))_{-1} \vec{D} \vec{Z}$$

Substituting  $\vec{W}$  into the first equation of (3-42), get:

$$(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D}) \vec{Z} = 1$$

hence 
$$\vec{Z} = (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}$$

and 
$$\vec{W} = -(\vec{F}(r_0))^{-1} \vec{D} (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}$$

in turn, 
$$\vec{G} = 2((\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} (\vec{F}(r_0))^{-1} \vec{D} (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1})$$

The derived result is clear now:

$$\begin{aligned} \vec{J} \vec{G}^{-1} &= (\vec{F}(r_0) - i \vec{D}) ((\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} (\vec{F}(r_0))^{-1} \\ &\quad \vec{D} (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1}) = \text{Re } \vec{S} + i \text{Im } \vec{S} \end{aligned} \quad (3-43)$$

Since the real part and imaginary part on each side of equation (3-43) are equal to each other respectively, we have:

$$\begin{aligned} \text{Re } \vec{S} &= (\vec{F}(r_0) - \vec{D}(\vec{F}(r_0))^{-1} \vec{D}) (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} \\ \text{Im } \vec{S} &= - (2\vec{D}) (\vec{F}(r_0) + \vec{D}(\vec{F}(r_0))^{-1} \vec{D})^{-1} \end{aligned} \quad (3-44)$$

where  $\vec{D} = \vec{K}^{-1} \vec{E}(r_0)$ .

The transition probability from state  $i = (n_0, m_0)$  to state  $j = (n, m)$  is:

$$P_{i(n_0, m_0) \rightarrow j(n, m)} = ((\text{Re } \vec{S})_{ji}^2 + (\text{Im } \vec{S})_{ji}^2) \frac{k_j}{k_i}$$

where  $\text{Re } \vec{S}$  and  $\text{Im } \vec{S}$  are given by equation (3-44). This completes the basic principle in calculating transition probability for one-dimensional scattering problem within the quantum mechanical approach. Discussions of boundary conditions in integrating equation (3-36) will be made in the next Chapter.

The original Schrodinger equation has rapidly oscillating wavelike solutions which are difficult to represent numerically. The integration of equation (3-33)

is numerically unstable, unless special algorithms are used. Secrest and Johnson, ( ) in their exact quantum mechanical treatment of the one-dimensional scattering problem, convert the coupled differential equations into equivalent integral equations. The integral involved is then replaced by a quadrature sum. The resulting matrix equation is then solved indirectly by numerical method to obtain the transition probabilities. Chan et al. (24), propose a different numerical approach to this problem. It involves converting the set of coupled second-order equations for the translational wavefunctions into first-order equations in matrix form and then solving it by an exponential method developed by W. Magnus (25). The idea was first conceived by Light et al. (26). The method we just discussed for quantum mechanical calculation of transition probabilities is a direct integration of the state expanded Schrodinger equation. This treatment is similar to the method due to Riley and Kupermann (12). It is relatively simple and straightforward, but in our procedures, the virtual states are not included in the total wavefunction expansion. Roy G. Gordon (21) developed another method for integrating coupled differential equations arising in bound state and scattering problems in quantum mechanics. The wavefunctions are constructed in piecewise analytic form, to any prescribed accuracy. The chief advantage of this method is that it avoids searching

for the correct initial derivatives of the wavefunction.  
It is claimed to be numerically very stable.

## Chapter 4 Numerical Results

### A. Semi-classical Results

For the specific  $N^2-N^2$  molecular collision, the two molecules are identical, then  $\omega_0 = \omega_{AB} = \omega_{CD} = 4.45 \times 10^{-14}$  sec,  $\gamma_{AB} = \gamma_{CD}$ , and  $m_A = m_B = m_C = m_D = 14.0$  a.m.u... This implies that the matrix elements  $U_{n'n}$  and  $V_{n'n}$  are equal. The matrix elements  $U_{n'n}$  are given in Table 4-1.

As mentioned in Chapter 3-B, we may set:

$$\begin{aligned} k &= j + (n-1) J, & j &= 1, 2, 3, \dots \\ & & n &= 1, 2, 3, \dots \end{aligned}$$

where  $J$  is an integer which is the number of states of molecules CD and AB included in the expansion of the total wavefunction  $\psi$ . The integer  $k$  is used to represent the state  $(j,n)$  which means that molecule CD is in the state  $j$  and molecule AB is in the state  $n$ . For example, if we choose  $J = 4$ , there are  $4 \times 4 = 16$  states involved in the expansion of total wavefunction  $\psi$  in terms of the individual harmonic oscillator wavefunctions. In other words, there are  $16 \times 2 = 32$  coupled first order differential equations to be solved in equation (3-21). (for  $J$  greater than 4, the extension is straightforward). In general,  $J=N$ , let  $K=k+N.N$ , where  $k=1,2,3,\dots(N.N)$ . In this way,  $A_{nj}$  and  $B_{nj}$  can be designated as:

$$\begin{aligned} Y(k) &= A_{nj} & k &= 1, 2, \dots (N.N) \\ Y(K) &= B_{nj} & K &= (N.N)+1, (N.N)+2, \dots 2(N.N) \end{aligned}$$

Y is then a vector of length  $2(N.N)$ . This is the suitable form for doing numerical integration of equation (3-21). In our computer program, the IBM IMSL ROUTINE DGEAR is called. On input,  $Y(1), Y(2), \dots, Y(2N*N)$  supply initial values which are initial conditions for the system. One of the arguments in the subroutine DGEAR, TOL, must be chosen suitably. Otherwise the computer time is unnecessarily long. This parameter TOL, is an estimate of the local truncation error. In a series test calculations, we choose  $N=11$ ,  $V_0 = 8$  km/sec and initial state = (1,2), for three different values of TOL,  $10^{-7}$ ,  $10^{-8}$ , and  $10^{-9}$ . We obtain the data as shown in Table 4-2. We then choose  $TOL = 10^{-8}$ . The initial value of the step size H, is chosen small enough at the beginning of integration so that it can pass the error test (based on TOL). In the subsequent procedures H is adjusted by the routine itself, but changing in the step size always satisfies the error test. The number of states used in the total wavefunction expansion,  $N*N$  plays a very important role in integration of equation (3-21).

In principle, we have to increase N until the final transition probabilities converge to values independent of N. Table 4-3 to Table 4-5 show the transition probability as a function of N for initial state (1,2) at low, medium, and high initial relative velocity; i.e.  $v_0 = 3$  Km/sec, 6 Km/sec, and 9 Km/sec respectively. From these tables, it is obvious that for high value of velocity  $v_0$ , we need more



states in the expansion of total wavefunction  $\psi$ . Generally speaking, for  $v_0$  less than 6 Km/sec,  $N = 7$  i.e.  $7 \times 7 = 49$  states expansion is good enough for initial states (1,1), (1,2), (2,2), (3,1), and (3,2). For  $v_0 = 7$  Km/sec,  $N$  should be no less than 9; and for  $v_0 = 8$  Km/sec,  $N$  should be no less than 11. For  $v_0 = 9$  Km/sec,  $N$  must be larger than 12. We also find that the value of  $N$  depends on the initial state. For example, for initial state (4,1) at velocity 6 Km/sec, only  $N = 9$ , i.e. 81 states expansion makes the transition probability converge. The integration limits are adjusted until the constraint equation (4-1) is satisfied: (for suitable  $N$ )

$$\left| \sum_{\text{final state}} P_{\text{initial state} \rightarrow \text{final state}}^{-1.0} \right| < \delta \quad (4-1)$$

where  $\delta$  is an arbitrary small number, we choose  $\delta = 10^{-4}$  here. The integration limits depend on the initial relative velocity  $v_0$ , e.g. for  $v_0 = 3$  Km/sec, the lower limit  $T = -1.6 \times 10^{-13}$  and the upper limit  $TEND = 1.6 \times 10^{-13}$ ; for  $v_0 = 6$  Km/sec,  $T = -1.0 \times 10^{-13}$ , and  $TEND = 1.0 \times 10^{-14}$ ; for  $v_0 = 9$  Km/sec,  $T = -5.0 \times 10^{-14}$ , and  $TEND = 5.0 \times 10^{-14}$ . Condition (4-1) serves as a useful criterion on numerical calculations.

The numerical results of transition probabilities for different initial and final states as a function of  $v_0$  are obtained by the semi-classical method and shown in Tables

4-6 to 4-10. Figure 4-1 through Figure 4-5 plot the transition probabilities as a function of  $v_0$  for five initial states (1,1), (1,2), (2,2), (3,1) and (3,2). A scaling relationship can be very useful for both the theoretical and experimental analysis of molecular scattering problem. With this relationship, it is easier to write the direct Monte Carlo computation program and save computer storage space. For a systematic study of the scaling relationship, we need more data. It is too expensive to be done at this time. In Figure 4-6a and 4-6b, we plot only the scaling relationship for V-T processes (1,i) (1,i+1) at different  $v_0$ . Since the collision is symmetric, there is nothing new in the results by changing the initial state (i,j) to (j,i). We examine Figure 4-1 through Figure 4-5 and find that at low energies the probability of transferring a given number of quanta by a V-V process is much greater than the probability of converting them by a V-T process into translational energy. For V-T process, the probability increases rapidly with increasing collision energies in the low energy regime. However, the probability of V-V transfer rises less sharply with increasing collision energy. For V-V transitions involving two quanta jumps such as (3,1) (1,3) and (3,2) (1,4), the transition probability is less than that of V-V process (1,2) (2,1), which involve only one quantum jump. Generally, at low energies, the transition probabilities are very small and multiple quantum transitions are assumed

primarily due to stepwise transitions via single collision  $j \rightarrow j+1 \rightarrow \dots \rightarrow K-1 \rightarrow K$ . At high collision energies the direct transition  $j \rightarrow K$  has significant contribution to the transition probability of multiple quantum jump. For V-V-T transfer to an adjacent level, processes involving transfer of a single quantum (such as  $(2,1) \rightarrow (3,1)$  and  $(2,1) \rightarrow (2,1)$ ) are much more probable than processes of several quanta such as  $(1,2) \rightarrow (3,1)$ . Likewise, transitions  $(3,2) \rightarrow (4,2)$ ,  $(3,2) \rightarrow (2,2)$  and  $(3,2) \rightarrow (3,1)$  are much more probable than the transition  $(3,2) \rightarrow (2,1)$ . In Figure 4-6a and 4-6b we notice that the general trend of the scaling relationship for transitions

$$(1,i) \rightarrow (1,i+1),$$

seems to be a weak  $v_0$ -dependent function. We need more data for further analysis.

A useful check on numerical results is provided by time-reversal invariance (which leads to the principle of detailed balance). Stated classically, the principle implies that a system executes its motion in reverse if time is allowed to run backward. In quantum scattering processes this means that  $P_{ij} = P_{ji}$ , i.e., the probability of a transition for state  $i$  to state  $j$  is equal to that for transition for state  $j$  to  $i$ . For example, check table 4-8 and table 4-9, we have  $P_{(2,2) \rightarrow (3,1)} = 0.845 \times 10^{-3}$ ,  $P_{(3,1) \rightarrow (2,2)} = 0.843 \times 10^{-3}$ , this gives

$$\frac{|P_{(3,1) \rightarrow (2,2)} - P_{(2,2) \rightarrow (3,1)}|}{P_{(3,1) \rightarrow (2,2)}} = 0.237\%, \text{ at}$$

$v_0 = 4 \text{ Km/sec}$

At  $v_0 = 6 \text{ Km/sec}$ , both  $P_{(3,1) \rightarrow (2,2)}$  and  $P_{(2,2) \rightarrow (3,1)}$  are equal to 0.121. The principle of time reversal invariance is satisfied quite well.

## B. Quantum Mechanical Results and Comparison

Referring to Chapter 3-C, it is clear that the numerical procedures for calculating transition probability by the quantum mechanical method are as follows:

(1) Integrate equation (3-36) and solve for  $\vec{F}(r_0)$  and  $\vec{E}(r_0)$ .

(2) Form  $\vec{D}$  and the expression  $(\vec{F}(r_0) + \vec{D}(\vec{F}(r_0)))^{-1} \vec{D}^{-1}$ .

(3) Construct  $\text{Re } \vec{S}$  and  $\text{Im } \vec{S}$ .

(4)  $P_{ij} = ((\text{Re } \vec{S})^2 + (\text{Im } \vec{S})^2)_{ji} \frac{k_j}{k_i}$ .

For  $N^2-N^2$  collisions the parameters of the system are  $\omega = 0.113$  (corresponding to  $L = 0.2 \text{ \AA}$ ),  $\omega = 1.0$ ,  $\beta = 1.0$ , and  $\bar{m} = 0.5$ . The total energy  $E$  can be assigned a suitable value which corresponds to some value of  $v_0$ . Having all these parameters, the IBM IMSL routine DGEAR is called to integrate equation (3-36) and find  $\vec{F}(r_0)$ ,  $\vec{E}(r_0)$ . There are four factors in this problem that can affect the numerical integration.

(1) Integration error.

(2) Number of states retained in the state expansions.

(3) Starting point of integration  $r_s$ .

(4) End point of integration  $r_0$ .

The local truncation error is controlled by TOL, which is one of the arguments of the subroutine DGEAR. We choose  $\text{TOL} = 1.0 \times 10^{-8}$  here. The starting point is chosen as the point  $r = r_s$  at which the largest diagonal element of  $\vec{V}(r_0)$

is equal to twice of the total energy there. The starting point is just beyond the classical turning point and in the classical forbidden region. Now, it is appropriate to discuss the initial conditions for this system. We set  $N = M = 2$ , so that there are four states involved in the total wavefunction expansion (for high energy collisions we need larger  $N$  and  $M$  to get more accurate results). The initial states of the molecule AB and CD are  $n_0$  and  $m_0$  definitely at the beginning, where  $n_0 = 1, 2$ ;  $m_0 = 1, 2$ . Therefore the initial value of the matrix  $\vec{F}(r)$  is a unit matrix:

$$\vec{F}_0^{\dagger} = \begin{matrix} & \begin{matrix} 1.0 & 0 & 0 & 0 \end{matrix} \\ \begin{matrix} 0 \\ 0 \\ 0 \\ 0 \end{matrix} & \begin{matrix} 1.0 & 0 & 0 & 0 \\ 0 & 1.0 & 0 & 0 \\ 0 & 0 & 1.0 & 0 \\ 0 & 0 & 0 & 1.0 \end{matrix} \end{matrix}$$

Since the point  $r = r_s$  is in the classical forbidden region, the diagonal element of  $\vec{V}(r_s)$  is much larger than the element of  $K$ , equation (3-33) becomes:

$$\frac{d^2 \vec{F}}{dr^2} = \vec{V}(r_s) \cdot \vec{F}$$

The diagonal element of  $\vec{F}$ ,  $f$ , satisfies the equation

$$\frac{d^2 f}{dr^2} = m e^{-\alpha r_s} \cdot f(r) \quad \text{as } r \rightarrow -\infty$$

The asymptotic solution is then  $\lim_{r \rightarrow \infty} f(r) = e^{\sqrt{m e^{-\alpha r_s} r}}$ . Recall that

$E = \frac{d\vec{F}}{dr}$ , this gives the initial value of matrix  $\vec{E}$  at starting point  $r_s$  as:

$$\vec{E}_0 = \begin{matrix} & g & 0 & 0 & 0 \\ & 0 & g & 0 & 0 \\ & 0 & 0 & g & 0 \\ & 0 & 0 & 0 & g \end{matrix}$$

where  $g = \sqrt{me^{-r_0}}/s$ . The stopping point  $r_0$  is chosen as the point where the diagonal element of  $V(r_0)$  are less than  $\frac{1.0}{5000}$  of the total energy. In principle, we have to increase the total number of states  $N, M$ , until the probabilities do not change significantly. We keep  $N=2$ , and  $M=2$  in this work and calculate transition probabilities among states (1,1), (1,2) and (2,2). Obviously, this four-state expansion is good for low energy collisions only. Table 4-11 and Table 4-12 display the quantum transition probabilities as a function of  $V_0$  for initial states (2,1) and (1,1). In Table 4-13 through Table 4-16, we list the probabilities for transitions (1,2)  $\rightarrow$  (1,1), (1,2)  $\rightarrow$  (2,2), (1,2)  $\rightarrow$  (2,1) and (1,1)  $\rightarrow$  (1,2) by the four different methods. Since the computer expense is prohibitively large for fully quantum mechanical method, we do not have enough data for plotting purpose. Here, we show that the method just explored does work. The plots of these tables are thus given by Figure 4-7, Figure 4-8, Figure 4-9 and Figure 4-10 respectively.

We check these Figures (4-7) through (4-10) and find that for the three V-T transitions the probability obtained by the semi-classical method is almost one order of magnitude smaller than that obtained from ZRS analytic

method. For V-1 process (1,2)  $\rightarrow$  (2,1), the semi-classical and ZRS results are very close to each other, however, the ZRS results are always slightly larger than the semi-classical results. The SSH theory is only good for low energy collisions, if the energy is too high the probability is greater than unity. This theory breaks down there. Since we are interested in high energy collision processes of two molecules. We concentrate on the semi-classical method. We believe the semi-classical treatment can supply a reasonable estimate in calculating transition probabilities. This is very helpful because the semi-classical method can save much computer time and the numerical algorithm is relatively simple.



## Chapter 5 Discussions

In this final chapter, we discuss some important problems requiring further study for vibrational energy transfer.

### A. Three-Dimensional Collisions

One of the assumptions which we have made in the collinear molecular collision model is that the target molecule is struck in the direction of its axis. To avoid this assumption in the collinear treatment, we have to average over the relative orientation of the molecule at the proper stage of calculations. However, the period of rotation is usually comparable with the duration of the collision, there is no simple way to take the average. A constant steric factor is generally used.

Since the rotational energy spacing is much smaller than the vibrational spacing, appreciable rotational scattering occurs over a range of molecule-molecule separations that is considerably longer than that for which vibrational transitions are important. The coupling between rotational and translation is usually strong too, so that the rotational state generally changes before the vibrational transition occurs. When a vibrational transition takes place the corresponding energy change will appear in either translational motion, or rotational motion of molecules, or both. It is obvious that if we calculate the vibrational transition probabilities,

effects of rotational motion have to be considered. In the collinear treatment however, we have assumed that the simultaneous rotational and vibrational transitions are not important and the impact parameter is zero. The realistic three-dimensional analyses that take rotational transitions into account should include the correction to the vibrational transition probability that results from the finite size of the rotational energy spacing in future work. Also, the incident particle is described by a plane wave which contains the partial waves of different orbital angular momentum (the one-dimensional model corresponding to an s-wave scattering problem). Usually, many partial waves have to be considered and this makes the problem very difficult.

#### B. Effect of Anharmonicity

It has been found by experience that the potential energy function of actual diatomic molecules can be represented quite accurately by a simple analytical function called Morse potential, which contains three adjustable parameters. If Morse potentials are used to describe the intramolecular forces, the diagonal matrix elements of the interaction potential which enter into the quantum theory of vibrational energy transfer are approximately but not identically equal. In the calculation given by F.H. Mies, (27) the consideration was restricted to the head-on collision between a structureless incident particle and a diatomic molecule. The transition

probability is found to decrease markedly when the ratio of the diagonal elements of the initial and final oscillator states is allowed to deviate even slightly from one. The deviation in turn, increases with the anharmonicity of the molecular vibrations, and an anharmonic correction factor of the order of  $10^{-1}$  to  $10^{-2}$  should be applied to the generally used probability expression for atom-molecule collision. There must exist a correction factor of this kind for molecule-molecule collision.

### C. Interaction Potential

Choice of a potential function to be used in calculating the transition probability is a very important task since it affects the results considerably. In the theory of inelastic molecular collisions, the scattering potential to be adopted should be simple enough to make the calculations feasible as long as the essential features of the physics of collision is not lost. This requirement is relaxed if we deal with numerical calculations. The chosen interaction potential for some pair of molecules must be relatively accurate and can be used to represent the real situation. If the intermolecular interaction is strongly orientation dependent, as in the polar gas, the molecules may take a particular orientation during the encounter. For this problem of preferential orientation, a somewhat different treatment is required.

#### D. Exact Classical Trajectory

In our semi-classical approximation, the classical equations are first solved to obtain the relative motion of the molecules as a function of the time. The time-dependent Schrodinger equation for the internal motion under the external perturbation is then solved to obtain the probabilities of various transitions. However, the occurrence of inelastic processes are not taken into account in solving the classical equation of motion in that the effective intermolecular potential and the effective translational energy depend on the internal state. If the incident energy is much greater than the internal energy, the influence of inelastic process on the relative motion is unimportant. It is a good approximation to ignore the internal state in calculating classical trajectory. For more rigorous calculations, an exact classical trajectory must be found in which the energy conservation law is satisfied. For high energy collisions, the semi-classical treatment is, however, a fairly good approximation. It requires less computational effort and saves much computer time.

## Appendix 1 Exponential Interaction Potential

A conventional representation of the intermolecular potential energy curve is given by the Lennard-Jones 12-6 power law (28);

$$V(r) = 4\epsilon \left( \left( \frac{\sigma}{r} \right)^{12} - \left( \frac{\sigma}{r} \right)^6 \right) \quad (\text{A1-1})$$

where  $V(r)$  is the potential energy at separation  $r$ , and  $r$  is the distance between atom B and C. This is shown graphically in Figure A1-1.

is the depth of the potential well at intermolecular distance  $r_m$ , where the repulsive force  $\left( \frac{\sigma}{r} \right)^{12}$  takes over the long range attractive force  $\left( \frac{\sigma}{r} \right)^6$ , and  $V(r_m)$  is the minimum of the potential function  $V(r)$ .  $\epsilon$  is the

separation at zero energy, when  $V(r) = 0$ , sometimes loosely called the "collision diameter". The exponential function

$$V_{\text{INT}}(r) = \text{constant} \cdot e^{-r/L} - \epsilon \quad (\text{A1-2})$$

must be fitted to the Lennard-Jones potential  $V(r)$ , equation (A1-1). Here, the choice is made that the magnitudes and slopes of the potentials are set equal at  $r = r_c$ .  $r_c$  is the minimum value of  $r$ . These two potentials are illustrated in Figure A1-2.

Figure A1-2. The exponential potential  $V_{\text{INT}}(r)$  fitted to the Lennard Jones potential  $V(r)$ . The magnitudes and slopes of the two potentials are set equal to each other at the classical turning point  $r = r_c$ .

We deduce an approximate formula:

$$L \approx \frac{\sigma}{17.5}$$

For  $N_2$  molecule,  $\sigma = 3.749 \text{ \AA}$ , so  $L = 0.21 \text{ \AA}$ .

## Appendix 2 The Choice of Initial Time Reference

### Coordinate in Quantum Scattering Process

Consider the general scattering in one-dimensional space, A flux of incoming particles with mean momentum  $p_0$  are incident from left and scattered by an arbitrary potential distribution  $V(x)$  as shown in Figure A2-1, where  $V(x)$  is finite and  $V(x) \rightarrow 0$  as  $x \rightarrow \pm\infty$

Figure A2-1. Particles scattered by an arbitrary potential.

For large and negative  $x$ , the wave packet with mean momentum  $p_0$  can be superposed as:

$$\psi(x,t) = \int_{-\infty}^{+\infty} dp \exp(-\alpha(p-p_0)^2) \exp\left(\frac{ipx}{\hbar}\right) \exp\left(-i \frac{p^2}{2m\hbar} t\right) \quad (\text{A2-1})$$

$$+ \int_{-\infty}^{+\infty} dp R(p) \exp(-\alpha(p-p_0)^2) \exp\left(-\frac{ipx}{\hbar}\right) \exp\left(-i \frac{p^2}{2m\hbar} t\right)$$

where  $R(p)$  is the reflection coefficient which is a constant over a region  $\Delta p \sim \frac{1}{\sqrt{\alpha}}$ . For large and positive  $x$ , the transmitted wave packet is:



$$\psi(x,t) = \int_{-\infty}^{+\infty} dp T(p) \exp(-\alpha(p-p_0)^2) \exp\left(\frac{ipx}{h}\right) \exp\left(-i \frac{p^2 t}{2mh}\right) \quad (\text{A2-2})$$

where  $T(p)$  is the transmission coefficient which is constant over a region  $\Delta p \sim \frac{1}{\sqrt{\alpha}}$ . Since for large  $|x|$  and  $|t|$ , the term  $\exp\left(i\left(\frac{px}{h} - \frac{p^2 t}{2mh}\right)\right)$  in equation (A2-1) and (A2-2) is a very rapidly varying function of momentum  $p$ , the integrals are essentially zero unless  $p$ ,  $x$ , and  $t$  are such that the stationary phase conditions are satisfied:

$$p = p_0$$

$$\frac{\partial}{\partial p} \left( \frac{px}{h} - \frac{p^2 t}{2mh} \right) = 0 \quad (\text{A2-3})$$

From equation (A2-3) we get:

$$x = \frac{p_0}{m} t \quad (\text{A2-4})$$

where  $\frac{p_0}{m}$  is the classical velocity. Equation (A2-4) gives the result that  $t \ll 0$ , if  $x \ll 0$ . Let's check the second term in the right hand side of equation (A2-1). Since  $R(p)$  is approximately a constant within a width  $\Delta p \sim \frac{1}{\sqrt{\alpha}}$  centered at  $p = p_0$ , then,

$$\int_{-\infty}^{+\infty} dp R(p) \exp(-(p-p_0)^2) \exp\left(-i \frac{px}{h}\right) \exp\left(-i \frac{p^2 t}{2mh}\right)$$

$$= R(p_0) \int_{-\infty}^{+\infty} dp \exp(-(p-p_0)^2) \exp\left(-i \left(\frac{px}{h} + \frac{p^2 t}{2mh}\right)\right)$$

Stationary phase conditions require that:

$$p = p_0$$

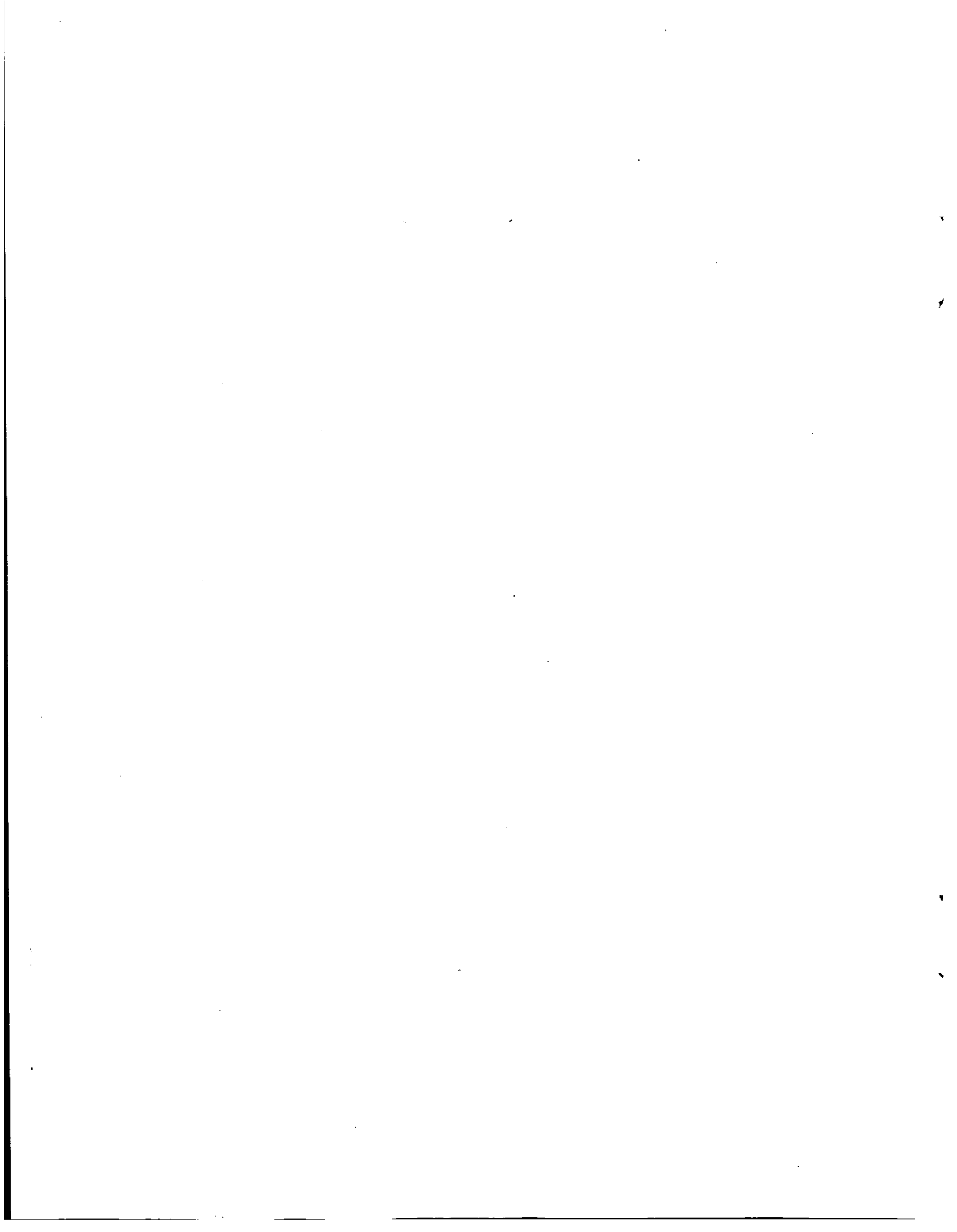
and

$$\frac{\partial}{\partial p} \left( \frac{px}{h} + \frac{p^2 t}{2mh} \right) = 0 \quad (\text{A2-5})$$

Equation (A2-5) implies that  $x = \frac{-p_0 t}{m}$ , this means that reflection occurs only when  $t > 0$  because the reflected wave exists only at  $x \longrightarrow -\infty$ . Combining the discussions just made, we conclude that the incident particles hit the potential at  $x = 0$  and  $t = 0$ , and the initial conditions of the system are described at  $t \longrightarrow -\infty$ , the final conditions are the states at  $t \longrightarrow +\infty$ .

## REFERENCES

1. G.A. Bird, "Shock Wave Structure in a Rigid Sphere Gas"  
Academic Press, Vol. 1, P. 216, New York 1965.
2. D.I. Pullin, J.K. Harvey, and G.K. Bienkowski  
"Hypersonic Leading Edge Flow of a Diatomic Gas by the  
Direct Simulation Method".



APPENDIX C

COMPUTER CODE INTERNAL

ORIGINAL FILE IS  
OF POOR QUALITY

FILE: GKBINT AUG82 A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

```
// JOB GKB 0367425.GKBSPACE N=WATRUN2 REG=560 T=1.0 P=100 C=0
// EXEC WATPIV
//WATPIV.PT09P001 DD DISP=OLD,DSN=U.GKBSPACE.RM115
//WATPIV.SYSIN DD DATA
$JOB      BIEKOWSKI,T=59,P=100,NOLIST
C        MAIN PROGRAM FOR MONTE CARLO 3-D ENTRANCE PROBLEM CALCULATIONS
C        OBJECTIVE OF THIS MAIN PROGRAM IS TO SET THE DIMENSIONS
C        MAIN RUNNING PROGRAM IS *** RUN ***
C
C        FOLLOWING TWO CARDS HAVE TO BE ELIMINATED FOR NON IBM MACHINES
C*****
C        INTEGER*2 LB,NBM,NBN,NB,NBT
C        INTEGER*2 LM,LCOL
C*****
C
C        THE NEXT CARD IS ASSOCIATED WITH PRINCETON RANDOM NUMBER GENERATOR
C*****
C        COMMON/RANCOM,/NRAN(4)
C*****
C
C        THE FOLLOWING DIMENSION STATEMENTS SET THE MAJOR ARRAY DIMENSIONS
C        AND MUST BE CONSISTENT WITH THE FOLLOWING DATA CARD -----
C
C        NSP=NUMBER OF SPECIES - EXAMPLE BELOW NSP=1
C        NJV=NUMBER OF SUBDIVISIONS OF INPUT DISTRIBUTION FUNCTION
C             EXAMPLE BELOW NJV=22
C        NMC=NUMBER OF FINAL CELLS - EXAMPLE BELOW NMC=150
C        NMP=MAX NUMBER OF MOLECULES OF EACH SPECIES ALLOWED IN PROGRAM.
C             IF EXCEEDED, PROGRAM EITHER FAILS OR RESTARTS AT BEGINNING
C             WITH NUMBER REDUCED BY 10% - EXAMPLE BELOW NMP=5000
C        NPB=MAXIMUM NUMBER IN EACH CELL - EXAMPLE NPB=150
C
C        DIMENSION DBA(1,150),NB(1,150),NBT(1,150)
C        DIMENSION TMP(1,150),TMPA(1,150),XV(1,150),YVA(1,150)
C        DIMENSION YV(1,150),YVA(1,150),ZY(1,150),ZVA(1,150),DB(1,150)
C        DIMENSION TRP(1,150),TRPA(1,150),NBM(1,150)
C        DIMENSION NBN(150),T(1,1,150)
C        DIMENSION LB(5000),LM(1,5000),ER(1,5000)
C        DIMENSION PAU(1,5000),PAV(1,5000),PAW(1,5000)
C        DIMENSION PAX(1,5000),PAY(1,5000),PAZ(1,5000),LCOL(1,5000)
C        DIMENSION PNB(150),IC(150),YC(150),ZC(150)
C        DIMENSION VEL(22,4,1),PFV(22,4,1)
C        DATA NSP/1/,NJV/22/,NMC/150/,NMP/5000/,NPB/150/
C
C
C        2 FORMAT(/17X,'NORMAL TERMINATION OF THE PROGRAM')
C        NAMELIST/DIM/NSP,NJV,NMC,NMP,NPB,NRAN
C
C        INITIALIZATION OF RANDOM NUMBER GENERATOR - PRINCETON ROUTINE
C
C        NRAN(1)=0
C        NRAN(2)=0
C        NRAN(3)=0
C        NRAN(4)=0
```

## PRINTOUT OF MAJOR ARRAY DIMENSIONS USED ABOVE

WRITE(6,DIM)

CALL OF MAIN OPERATING PROGRAM WHICH REQUIRES INPUTS:  
 &CONTRL,&TIMES,&FLOREP,&MOLEC,&SHAPES,&GEON,&INCUPL,&INOUT  
 THESE INPUTS ARE ALL CURRENTLY IN THE NAMELIST FORMAT  
 AND MAY HAVE TO BE CHANGED IF THAT CONVENTION IS NOT AVAILABLE  
 BRIEF DESCRIPTION OF THE PARAMETERS FOLLOWS

&amp;CONTRL - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
NAME	8 BLANKS	ANY ALPHANUMERIC NAME UP TO 8 CHARACTERS
TITLE	24 BLANKS	ANY ALPHANUMERIC TITLE UP TO 24 CHARACTERS
PERCNT	.001	ACCURACY IN INTEGRATION PROCEDURES
ICOPY	1	NUMBER OF ADDITIONAL COPIES OF OUTPUT
DUMP	.TRUE.	IF TRUE WILL CAUSE SYSTEM DUMP FOR ANY OF 12 PROGRAMMER DESIGNED ERROR HALTS.
DEBUG (1)	.FALSE.	IF TRUE WILL PRINT MESSAGE WHEN CELL POP. EXCEEDS MNB
DEBUG (2)	.FALSE.	IF TRUE WILL PRINT CPU TIME AROUND EACH PART OF LOOP
DEBUG (3)	.TRUE.	IF TRUE WILL PRINT CPU TIME REMAINING AT END OF LOOP
NEW	.TRUE.	IF TRUE - NEW RUN - IF FALSE - RESTART OF RUN
SAVE	.FALSE.	IF TRUE - SNAPSHOT SAVED ON TAPE(9) FOR RESTART
REDO	.FALSE.	IF TRUE PROGRAM WILL AUTOMATICALLY RESTART WITH 90% OF TOTAL IF TOTAL CELL POPULATION EXCEEDS MNM

&amp;TIMES - ONE OCCURRENCE (NEW OR RESTART)

PARAMETER	DEFAULT	DEFINITION OR EXPLANATION
DTM	- - -	REAL NUMBER - FRACTION OF MEAN FREE TIME PER CYCLE
ITS	- - -	INTEGER - NUMBER OF CYCLES PER SAMPLE
ITP	- - -	INTEGER - NUMBER OF CYCLES BETWEEN PRINTOUTS
TST	- - -	INTEGER - ESTIMATE OF NUMBER OF CYCLES TO STEADY STATE
TLIM	- - -	INTEGER - TOTAL NUMBER OF CYCLES TO END OF RUN - WILL TERMINATE SCORER IF CPU TIME IS TO BE EXCEEDED

&amp;FLOREP - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
LLM	- - -	INITIAL NUMBER OF MOLECULES $LLM < MNM < OR = NNP$
MNM	- - -	MAXIMUM NUMBER OF MOLECULES PER SPECIES
MNB	- - -	MAXIMUM NUMBER PER CELL - DIAGNOSTIC ONLY
MSP	- - -	NUMBER OF MOLECULAR SPECIES (MAX. IS 3)
NET	0	IF 0 - DATA IS IN SI (METRIC) UNITS IF > 0 - DATA IS IN ENGLISH UNITS
U	- - -	FLOW VELOCITY (M/SEC) OR (FT/SEC)
ANGLE	- - -	ANGLE OF ATTACK (DEGREES)
RNU	0.0	ARRAY GIVING MOLE FRACTIONS OF SPECIES IN FREE STREAM
RMA	0.0	ARRAY GIVING MOLECULAR WEIGHTS OF SPECIES ABOVE
TF	- - -	FREE STREAM TEMPERATURE (K OR R)
DENF	- - -	FREE STREAM NUMBER DENSITY (NUM/M**3 OR NUM/FT**3)

&amp;MOLEC - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
TRF	- - -	REFERENCE TEMPERATURE FOR MOLECULAR DATA
DIR	0.0	CROSS-SECTIONS AT REFERENCE TEMP. (MSP1MSP)
ETA	0.0	PARAMETERS IN DIFFUSION AND VISCOSITY LAW (MSP1MSP)
PHI	0.0	PARAMETERS FOR ROTATIONAL RELAXATION (MSP1MSP)

CHI 0.0 ROTATIONAL DEGREE OF FREEDOM PARAMETER (NROT/2 - 1)  
 ACR .001 ACCURACY IN MOLECULAR COLLISION CALCULATIONS

ESHAPES - ND+1 OCCURRENCES WHERE ND=NUMBER OF BODY SEGMENTS (NEW RUN)

PARAMETER	DEFAULT	DEFINITION
		FIRST OCCURRENCE
BODY(I) I>3	--	NEED NOT BE SPECIFIED
BODY(1)	--	START OF BODY (ISTART) IN ARBITRARY COORDINATE
BODY(2)	0.0	TEMPERATURE AT FRONT OF TUBE IN K OR R
BODY(3)	0.0	DIAMETER OF TUBE IN METERS OR FT.
		SUBSEQUENT OCCURRENCES (ND)
BODY(1)	--	X COORDINATE FROM FRONT OF BODY OF THE DOWNSTREAM EDGE OF THE CURRENT BODY SEGMENT
BODY(2)	--	TEMPERATURE AT THE BACK OF THIS BODY SEGMENT
BODY(3)	--	SWITCH - IF NOT 0.0 THIS IS THE LAST SHAPES CARD
BODY(I) I EVEN		ALPHA - ENERGY ACCOMODATION COEFFICIENT FOR SPECIES
BODY(J) J ODD		SIGMA - TANGENTIAL ACCOMODATION COEFF. FOR SPECIES
		I AND J < (4+2*MSP)

EGEOM - ONE OCCURRENCE (NEW RUN ONLY)

PARAMETER	DEFAULT	DEFINITION
NWEDG	--	INTEGER GIVING THE NUMBER OF WEDGES WITHIN 180 DEGREES
NW	--	NUMBER OF FIRST LEVEL CELLS IN X DIRECTION
NH	--	NUMBER OF FIRST LEVEL CELLS IN RADIAL DIRECTION

EINCOPL - ONE OCCURRENCE (NEW RUN ONLY) - INPUT DISTRIBUTION

PARAMETER	DEFAULT	DEFINITION
FLUXIN	1.0	FLUX INPUT IN TERMS OF FREE STREAM FLUX - ONE NUMBER FOR EACH SPECIES
PCOL	0.0	FRACTION OF ARRIVING MOLECULES THAT HAVE PREVIOUSLY COLLIDED
RHP	--	RATIO OF "CAVITY" PRESSURE TO THE EFFECTIVE PRESSURE OF THE INCOMING STREAM AT ENTRANCE
JV	--	THE NUMBER OF VELOCITY INTERVALS FOR DISTRIBUTION FUNCTION INFORMATION
KMX	4	NUMBER OF COMPONENTS OF DISTRIBUTION KMX=3 IF NO POTATIONAL ENERGY (CHI=-1) KMX=4 IF ROTATIONAL ENERGY IS INCLUDED (CHI>-1)

FINOUT - MSP\*KMX OCCURRENCES

PARAMETER	DEFAULT	DEFINITION
		MT DESIGNATES SPECIES
		K=1 DESIGNATES NORMAL VELOCITY
		K=2 DESIGNATES TANGENTIAL VELOCITY IN FLOW DIRECTION
		K=3 DESIGNATES TRANSVERSE TANGENTIAL VELOCITY
		K=4 DESIGNATES ROTATIONAL ENERGY
VARG(J)	--	VELOCITY BOUNDARIES FOR DISTRIBUTION FUNCTION 1<J<JV VELOCITY BCUNDARIES
CURV(J)	--	PROBABILITY OF VELOCITY (OR ROTATIONAL ENERGY) INCIDENT AT ENTRANCE BEING BELOW VARG(J)

A SAMPLE INPUT DECK IS GIVEN BELOW:



```

&CONTROL NAME='INTE','BNAL',TITLE=' PAR','ABOL','A AT',' 95K','M H','OH. ',
  DEBUG=.F.,.F.,.T.,NEW=.T.,SAVE=.F.,ICOPY=0,REDO=.T. &END
&TIMES DTN=.010,ITS=5,ITP=1000,TST=400,TLIN=1000 &END
&FLOREP LLM=2000,MNH=5000,MNB=150,MSP=1,NET=0,U=7485.9,ANGLE=28.,RNU=1.,2*0.,
  RMA=28.94,0.,0.,TP=195.51,DENF=2.52E+19 &END
&MOLEC TRP=1000,DIR=3.5E-19,ETA=.104,PHI=0.0,CHI=-1.,ACR=.001 &END
&SHAPES BODY=0.0,1000.,.00235 &END
&SHAPES BODY=.0025,555.,0.0,2*1.0 &END
&SHAPES BODY=.0050,345.,0.0,2*1.0 &END
&SHAPES BODY=.0100,300.,0.0,2*1.0 &END
&SHAPES BODY=.0200,300.,0.0,2*1.0 &END
&SHAPES BODY=.0300,300.,0.0,2*1.0 &END
&SHAPES BODY=.0400,300.,0.0,2*1.0 &END
&SHAPES BODY=.0500,300.,0.0,2*1.0 &END
&SHAPES BODY=.0600,300.,0.0,2*1.0 &END
&SHAPES BODY=.0700,300.,0.0,2*1.0 &END
&SHAPES BODY=.0800,300.,0.0,2*1.0 &END
&SHAPES BODY=.0870,300.,1.0,2*1.0 &END
&GGEOM MWEDG=2,NW=20,NH=3, &END

```

```

&INCUPL FLUXIN=2.1429,PCOL=1.0,RMP=0.0,JV=22,KMI=3 &END
&INOUT VARG=0.,1.,2.,3.,4.,5.,6.,7.,8.,9.,10.,11.,12.,13.,14.,15.,16.,17.,18.,
  19.,20.,21.,CURV=0.0,.070,.170,.282,.369,.459,.537,.599,.656,.710,.750,.785,
  .815,.845,.872,.900,.922,.951,.975,.988,.996,1.00, &END.
&INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
  11.,13.,15.,17.,19.,CURV=4*0.,.003,.013,.036,.084,.149,.250,.406,.611,.762,
  .871,.932,.962,.984,.995,.999,3*1.0, &END
&INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
  11.,13.,15.,17.,19.,CURV=4*0.0,.003,.013,.036,.084,.149,.250,.406,.611,.762,
  .871,.932,.962,.984,.995,.999,3*1.0, &END

```

```

CALL RUN(NSP,NJV,NMC,NMP,NPB,DBA,NB,NBT,TMP,TMPA,IV,XVA,IV,
  1YVA,ZV,ZVA,T,DB,FNB,IC,YC,ZC,LM,PAU,PAV,PAW,PAX,PAY,
  2PAZ,LCOL,TRP,TRPA,ER,LB,NBM,NBN,VEL,PPV)
WRITE(6,2)
STOP
END

```

```

SUBROUTINE RUN(NSP,NJV,NMC,NMP,NPB,DBA,NB,NBT,TMP,TMPA,IV,XVA,
  1 YV,YVA,ZV,ZVA,T,DB,FNB,IC,YC,ZC,LM,PAU,PAV,PAW,PAX,PAY,PAZ,
  2LCOL,TRP,TRPA,ER,LB,NBM,NBN,VEL,PPV)
  MAIN RUNNING PROGRAM ** RUN *** CALLS ALL OTHER SUBROUTINES

```

\*\*\*\*\*

```

INTEGER*2 LM,LCOL
INTEGER*2 LB,NBM,NBN,NB,NBT
INTEGER PRT,SAMP,TST,TLIN,TIME,Q
LOGICAL DUMP,DEBUG(3),SAVE,NEW,REDO
REAL INTGRL,LAM,MU,NU,JAY,KAY
DIMENSION BTA(3),C1(3),C2(3),C3(3),C7(3),C8(3),DPA(3),FL(3)
DIMENSION FDN(3),HTI(3),HTR(3),JNT(3),KNH(3),NH(3),SR(3)
DIMENSION NAME(2),TITLE(6)
DIMENSION RNU(3),RMA(3),WTH(3),CHI(3),DIR(3,3),DAN(3,3),PHI(3,3)
DIMENSION ETA(3,3),CNS(3,3),CNG(3),CHG(3),CH(3,3,3),CH(3,3,3)

```



FILE: GKBINT AUG82 1

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

```

2 FORMAT (1H1/17X, 'RARIIFIED SUPERSONIC FLOW OF BINARY GAS', T74, 'I') RUN0760
3 FORMAT ('+', 103X, 'COPY ', I2) RUN0770
4 FORMAT (/17X, 'FLOW THROUGH ALL THE BOUNDARIES'/5X, 'MT', 5X, 'MASS MOL
1E PR.', 10X, 'PCOL. FLUXIN FLUXES (ENT)')
5 FORMAT (3X, I4, F9.2, F9.4, 5X, 4F10.4)
25 FORMAT (/5X, 'PRESSURE RATIO (INSIDE/ENTRANCE) - EITHER TYPE =', F13.5
1, ' RMP', T76, 'I'/6X, 'DENSITY RATIO (INSIDE/ENTRANCE) - EITHER TYPE
2=', F13.5, ' RMN', T76, 'I'/9X, 'FLUX RATIO (INSIDE/ENTRANCE) - EITHER
3TYPE =', F13.5, ' RMP', T76, 'I'/)
26 FORMAT (/10X, 'FLUX RATIOS FOR SPECIES RMA =', F7.2/2X, 'BOUNDARY',
1 ' INFLUX REFLUX NET FLOX NET FLUX/RHO*U'/2X, 'ENTRANCE',
24F10.4/4X, 'CAVITY', 4F10.4/)
30 FORMAT (' TIME =', F6.3, 60X, 'RANDOM NUMBER GENERATOR HAS BEEN CALLED
1 ', I10, ' TIMES') RUN0800
31 FORMAT (' CPU TIME LEFT- ', F8.3) RUN0810
32 FORMAT (7X, '-MOLECULES-' /3X, 3I6)
33 FORMAT (' TIME = ', F8.3, 5X, 'COLLISION LOOP=', F8.3, 5X, 'MOVE LOOP = '
1, F8.3, 5X, 'TOTAL TIME = ', F8.3/2X, '2ND MOVE LOOP =', F8.3, 5X,
2 'CLEANUP LOOP=', F8.3, 4X, 'PARTICLE NUMBERS = ', 4I6)
34 FORMAT (9X, '-MOLECULAR COLLISIONS-' /3 (3I14/))
35 FORMAT (2X, '-COLLISIONS WITH SURFACE-' /3X, 3I8)
36 FORMAT (' MAXIMUM NUMBER OF MOLECULES SO FAR- ', I6//) RUN0880
38 FORMAT (' EXCESS MOLECLES OCCURRED IN ', 3I4) RUN0890
40 FORMAT (/ ' SOMETHING IS WRONG WITH BOX NUMBERING IN RUN '/9I5, 5E14. RUN0900
15)
44 FORMAT (' NB (' , I2, ', ', I4, ') POPULATION EXCEEDED ', I3, ' IN MAIN AT TRUNC930
TIME = ', F7.3) RUN0940
50 FORMAT (/// ' SNAP SAVED ON TAPE') RUN0950
: RUN0960
: ***** RUN0970
: ***** RUN0980

```

```

CPA=ELTIME (0)
CALL NOUNDF
CALL TRAPS (0, 1, 1000000, 1, 1)
LIMIT (4)=NMC
LIMIT (5)=NMP
LIMIT (6)=NPB
LIMIT (7)=NJY
LIMIT (10)=RSP
KAVLS=0
PI=3.141593 RUN1030
PIROOT=SQRT (PI) RUN1040
SET=0
LARGE=0 RUN1060
NL=1 RUN1080
DUMP=.TRUE. RUN1190
DEBUG (1)=.FALSE.
DEBUG (2)=.FALSE.
DEBUG (3)=.TRUE.
SAVE=.FALSE. RUN1230
NEW=.TRUE. RUN1260
REDO=.FALSE. RUN1240
PERCMT=.001 RUN1250
ACR=.001
DO 58 I=1, 15

```



```

58 BODY(I)=0.0
   DO 60 I=1,3
   RNU(I)=0.0
   RMA(I)=0.0
   CHI(I)=0.0
   FLUXIN(I)=0.0
   PCOL(I)=0.0
   LL(I)=0
   DO 59 J=1,18
   ALPHA(I,J)=1.0
59 SIGMA(I,J)=1.0
   DO 60 K=1,3
   ETA(I,K)=0.0
   PHI(I,K)=0.0
60 DIR(I,K)=0.0
   WRITE(6,1)
   READ(5,CONTRL)
   WRITE(6,CONTRL)
   IF(NEW) GO TO 103
   REWIND 9
   READ(9) DENP,U,XREP,TRP,KAWLS,NL,NW,NH,BW,BH,HREG,XLB,XLC,PI,ND,
2     S,SINANG,COSANG,AKN,AKT,NBX,BM,XR,TIME,DTM,TI,ITS,ITP,TST,
3     TLIH,RMA,RNU,DIR,XSTART,MNE,MNB,TR,BZC,CN7,DRP,PCP,FNA,
4     HTF,INH,LLM,NAV,NEAX,NWEDG,PRT,SAMP,AKM1,AKM2,AKT1,AKT2,
5     BTA,C1,C2,C3,C7,C8,DAM,DPA,PL,DELANG,PDN,HTI,HTR,JNT,KNS,
6     NH,WTH,C4,VRM,NCOL,CTI,CTR,CNI,CNR,SM,
7     ST,D1,D2,D3,D4,NRAN,VELR,RMP,RMN,RMP,IPLUX,FLUXIN,
8     XLIM,COEFP,XCB,XS,YCB,TB,ALPHA,SIGMA,NTS,
9     UTL,UTT,VTS,HTS,HTSI,ENT,REM,TMPA,
A     DBA,HB,NBT,TMP,XV,XVA,YV,YVA,ZV,ZVA,T,DB,PNB,XC,YC,ZC,
B     PAU,PAV,PAW,PAX,PAY,PAZ,LCOL,LM,
C     ETA,PHI,CHI,CN,CH,CNG,CHG,CN8,TRP,TRPA,MSP,ANGLE,TF,
D     UTLI,UTTI,VTSI,ER,RMB,LB,NBM,NBN,VEL,PPV,PCOL,JV
   DTMO=DTM
   READ(5,TIMES)
   WRITE(6,TIMES)
   IF(DTM.EQ.DTMO) GO TO 100
   AIME=TIME*DTMO
   TIME=AIME/DTM+0.1
   DO 99 J=1,NSP
   DO 99 L=1,2
   ENT(J,L)=ENT(J,L)*DTM/DTMO
99 CONTINUE
100 IF(TI.GT.0.0) TST=TI/DTM
   WRITE(6,2)
   WRITE(6,4)
   WRITE(6,5) (MT,RMA(MT),RNU(MT),PCOL(MT),FLUXIN(MT),(ENT(MT,K),K=1,2
1) ,MT=1,MSP)
   WRITE(6,2)
   CALL PRINTA(NWEDG,TITLE,NAME,XCB,YCB,TB,ALPHA,SIGMA,XLIM,
1COEFP,LIMIT,MSP)
   CALL PRINTB(PNA,MSP,PNB,NM,XLIM,XC,YC,ZC,NB,NSP)
   GO TO 280
103 READ(5,TIMES)
   WRITE(6,TIMES)

```

RUN1330  
 RUN1340  
 RUN1360  
 RUN1370

RUN1540  
 RUN1590

C-3

```

READ (5, FLOREP)
WRITE (6, FLOREP)
READ (5, MOLEC)
WRITE (6, MOLEC)
IF (MSP.GT.LIMIT(10)) CALL DIA (10, LIMIT(10), MSP)
CHIM=0.0
RMR=0.0
DMR=0.0
ETT=0.0
DO 105 M=1, MSP
RMR=RMR+RMA (M) *RNU (M)
CHIM=CHIM+CHI (M) *RNU (M)
105 CONTINUE
DO 115 K=1, MSP
DO 115 M=1, MSP
ETT=ETT+RNU (M) *RNU (K) *ETA (M, K)
SRMA=SQRT (.5*RME* (1./RMA (M) +1./RMA (K)))
115 DMR=DMR+RNU (M) *RNU (K) *DIR (M, K) * (TRP/TF) ** (ETA (M, K)/2.) *SRMA
XREP=1./ (DENP*DMR*1.414214)
VELR=SQRT (16628.64*TF/RMR)
IF (MET.EE.0) VELR=SQRT (99437.92*TF/RMR)
TMR=XREP/VELR
S=U/VELR
NREG=1
ND=1
READ (5, SHAPES)
WRITE (6, SHAPES)
XCB (1)=BODY (1) /XREP
TB (1)=BODY (2) /TF
XLIM (1)=XCB (1)
XSTART=XLIM (1)
RMB=.5*BODY (3) /XREP
TR=TB (1)
104 READ (5, SHAPES)
WRITE (6, SHAPES)
ND=ND+1
IF (ND.GT.LIMIT (3)) CALL DIAG (3, LIMIT (3), ND)
XCB (ND)=BODY (1) /XREP
TB (ND)=BODY (2) /TF
DO 1104 M=1, MSP
ALPHA (M, ND)=BODY (2+2*M)
1104 SIGMA (M, ND)=BODY (3+2*M)
IF (TB (ND).GT.TR) TR=TB (ND)
IF (BODY (3).EQ.0.0) GO TO 104
NSTEP=NREG+1
ILIM (NSTEP)=XCB (ND)
COEFF (1)=0.0
COEFF (2)=1.0
COEFF (3)=0.0
COEFF (4)=-RMB**2
IR=XLIM (NSTEP)-XSTART
AKN=1./IR
AKT=.5/RMB
RHFP1=TB (1) ** (.5+ETT/2.) / (2.*PIROOT*S)
RHFP2=TB (ND) ** (.5+ETT/2.) / (2.*PIROOT*S)

```

RUN1640  
 RUN1650

RUN1670  
 RUN1680  
 RUN1690

RUN1770  
 RUN1800

```

AKN1=AKN*EMFP1
AKN2=AKN*EMFP2
AKT1=AKT*EMFP1
AKT2=AKT*EMFP2
DO 260 N=1,3
DO 260 M=1,3
NCOL(N,M)=0
260 CONTINUE
READ(5,GEOM)
READ(5,INCUPL)
WRITE(6,GEOM)
WRITE(6,INCUPL)
DO 116 J=1,42
VARG(J)=0.0
116 CURV(J)=0.0
RMN=RSP*TB(1)/TB(ND)
RMP=SQRT(RMN*RSP)
RM=RMB
BW=XR/NW
BH=RM/NH
IF(NWEDG.GT.LIMIT(1)) CALL DIAG(1,LIMIT(1),NWEDG)
IF(MNB.GT.LIMIT(5)) CALL DIAG(5,LIMIT(5),MNB)
IF(MNB.GT.LIMIT(6)) CALL DIAG(6,LIMIT(6),MNB)
DELANG=180./NWEDG
SINANG=SIN(ANGLE/180.*PI)
COSANG=COS(ANGLE/180.*PI)
VOL=PI*RM*RM*XR
NBX=NW*NH*NWEDG
IF(NBX.GT.LIMIT(4)) CALL DIAG(4,LIMIT(4),NBX)
IF(JV.NE.LIMIT(7)) CALL DIAG(7,LIMIT(7),JV)
BR=SQRT(TR)
SRMX=0.0
DO 916 MT=1,MSP
WTH(MT)=RMA(MT)/RMR
BTA(MT)=SQRT(WTH(MT))
SR(MT)=S*BTA(MT)
SRT=SR(MT)*PLUXIN(MT)/RNU(MT)
IF(SRT.GT.SRMX) SRMX=SRT
916 CONTINUE
INM=LLM*SQRT(TB(1))/PIROCT/SRMX/(1.+RMN)
DDN=INM/VOL
DO 140 MT=1,MSP
FDN(MT)=RNU(MT)*DDN
DPA(MT)=RNU(MT)
SN(MT)=SR(MT)*COSANG
ST(MT)=SR(MT)*SINANG
DO 117 K=1,MSP
DAM(K,MT)=DIR(K,MT)*(TRP/TF)**(ETA(K,MT)/2.)/DMR
CNB(K,MT)=DDN/DAM(K,MT)*1.414214
BT=AMIN1(BTA(K),BTA(MT))
VR1=S+3.*(1.+SQRT(TR))/BT
VR2=3.*SQRT((1.+2.*S**2/(5.+CHIN))*(1./WTH(K)+1./WTH(MT)))
CH(K,MT,1)=AMAX1(VR1,VR2)
CN(K,MT,1)=RAND(0)*CH(K,MT,1)
DF=PHI(K,MT)*(CHI(K)+CHI(MT)+2.)-1

```

RUN2200

RUN2230

RUN2240

RUN2320

RUN2330

RUN2420

RUN2470

RUN2480

RUN2740

RUN2750

```

DS=PHI (K,MT) *(2.-.5*ETA (K,MT) )-1.0
DO 917 N=2,3
XPH=ACR**AMIN1 (DF,DS)
IF ((DF.GT.0.) .AND. (DS.GT.0.)) XPH=(DF/(DF+DS) ) **DF* (DS/ (DF+DS) ) **DS
XPN=ACR**AMAX1 (DF,DS)
IF ((DF.LT.0.) .AND. (DS.LT.0.)) IPN=(DF/(DF+DS) ) **DF* (DS/ (DF+DS) ) **DS
CM (K,MT,N) =XPH-XPN
CN (K,MT,N) =RAND (0) *CM (K,MT,N)
DF=CHI (K)
DS=CHI (MT)
917 CONTINUE
117 CONTINUE
DO 118 K=1,KMI
READ (5,INOUT)
WRITE (6,INOUT)
DO 118 J=1,JV
VEL (J,K,MT) =VARG (J)
PPV (J,K,MT) =CURV (J)
118 CONTINUE
ENT (MT,1) =INH*S*DTM*AKN*FLUXIN (MT) /RND (MT)
ENT (MT,2) =RNN*ENT (MT,1) *SQRT (TB (ND) /TB (1) )
REN (MT,1) =0.0
REN (MT,2) =0.0
LL (MT) =INH*PIROOT* (1.+REN) *SR (MT) /SQRT (TB (1) ) *FLUXIN (MT) /RND (MT)
CHT=CHI (MT)
IF (CHT.GT.0.) CMG (MT) =CHT**CHT*EXP (-CHT)
IF (CHT.EQ.0.) CMG (MT) =1.0
IF (CHT.LT.0.) CMG (MT) =ACR**CHT*EXP (-ACR)
CMG (MT) =RAND (0) *CMG (MT)
140 CONTINUE
XS (1) =0.0
DO 155 N=2,ND
155 XS (N) = (.5* (XCB (N) +ICB (N-1) ) -XSTART) *AKN
YCB (1) =0.0
DO 160 N=2,ND
160 YCB (N) =2. * (YCB (N) -ICB (N-1) ) /RNB
CALL CELL (BW, BR, NW, NE, XSTART, DELANG, NWEDG, IC, YC, ZC, PNB)
FNA=0.0
DO 210 N=1,NBX
?10 FNA=FNA+FNB (N)
NPX=NBX
220 TIME=0
LARGE=0
SAMP=0
PRT=0
NAV=0
AIME=0.
TI=0.0
DT=DTM
EMAX=0
DO 250 MT=1,3
C1 (MT) =RAND (0)
C2 (MT) =RAND (0)
C3 (MT) =RAND (0)
C7 (MT) =RAND (0)

```

RUN3330

RUN3380

RUN3390

RUN4250

RUN4260

RUN4270

RUN4290

RUN4300

RUN4310

RUN4380

RUN4410

RUN4420

RUN4430

RUN4440

FILE: GKBINT AUG82 A

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C8 (MT)=RAND(0)	RUN4450
D1 (MT)=RAND(0)	RUN4460
D2 (MT)=RAND(0)	RUN4470
D3 (MT)=RAND(0)	RUN4480
D4 (MT)=RAND(0)	RUN4490
FL (MT)=0.	RUN4500
HTI (MT)=0.	RUN4510
HTR (MT)=0.	RUN4520
JNT (MT)=0	RUN4530
NM (MT)=0	RUN4540
IPLUX (MT,1)=0	
IPLUX (MT,2)=0	
DO 230 N=1,3	RUN4550
CTI (MT,N)=0.	RUN4560
CTR (MT,N)=0.	RUN4570
CNI (MT,N)=0.	RUN4580
230 CNR (MT,N)=0.	RUN4590
DO 240 N=1,ND	RUN4600
DO 240 K=1,NWEDG	RUN4610
NTS (MT,N,K)=0	RUN4620
HTSI (MT,N,K)=0.	RUN4640
UTLI (MT,N,K)=0.	
UTTI (MT,N,K)=0.	
VTSI (MT,N,K)=0.	
UTL (MT,N,K)=0.	RUN4650
UTT (MT,N,K)=0.	RUN4660
VTS (MT,N,K)=0.	RUN4670
240 HIS (MT,N,K)=0.	RUN4680
DO 245 N=1,NPX	RUN4690
NB (MT,N)=0	RUN4700
NBT (MT,N)=0	RUN4720
DBA (MT,N)=0.	RUN4730
XVA (MT,N)=0.	RUN4740
YVA (MT,N)=0.	RUN4750
ZVA (MT,N)=0.	RUN4760
TMPA (MT,N)=0.	RUN4790
TRPA (MT,N)=0.0	
DO 245 NN=1,NSP	
T (MT,NN,N)=0.0	
245 CONTINUE	RUN4800
250 CONTINUE	
PND=DDN	
DRP=2./ (PKD*S*S*RMB*RMB*PI)	RUN4930
PCP=1./ (PND*S*RMB*RMB*PI)	RUN4940
HTP=.5*DRP/S	
WRITE (6,2)	
WRITE (6,4)	
WRITE (6,5) (MT,RMA (MT),RNU (MT),PCOL (MT),FLUXIN (MT), (ENT (MT,K),K=1,2	
1), MT=1, MSP)	
WRITE (6,2)	RUN5000
CALL PRINTA (NWEDG, TITLE, NAME, XCB, YCB, TB, ALPHA, SIGMA, XLIM,	
1COEPP, LIMIT, MSP)	
CALL GAS (NWEDG, DELANG, ND, BTA, C1, DPA, HE, FNB, DB, NB, NBE, NBN,	
1PAU, PAV, PAW, PAX, PAY, PAZ, XLIM, COEPP, LB, LIMIT (4), LIMIT (6), XCB, TB,	
2LARGE, MNR, MNB, DEBUG (1), LCOL, MSP, ER, CHI, CHG, CHG, MSP, LB)	



```

CPUTYM=TPIND(0)
IF (LARGE.NE.0) GO TO 345
DO 265 I=1,MSP
265 IF (NM(I).GT.NMAX) NMAX=NM(I)
CALL PRINTB(FNA,MSP,FNB,NM,XLIM,XC,YC,ZC,NB,NSP)
IF (DEBUG(2)) WRITE(6,1)
CALL ACCUM(NMC,NPB,FNB,NB,PAU,PAV,PAW,ER,TRP,TRP,XV,YV,ZV,LM,MSP,
1NSP,NBM)
CPA=ELTIME(0)
CPI=CPA
GO TO 340
280 TIME=TIME+1
IF (TIME.NE.TST+1) GO TO 285
TI=TST*DTM
DO 282 MT=1,MSP
IFLUX(MT,1)=0
282 IFLUX(MT,2)=0
285 LARGE=0
CPI=ELTIME(0)
AIME=TIME*DTM
DT=AIME-TI
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
PRT=PRT+1
SAMP=SAMP+1
CALL COLIDE(CM,CM,WTM,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,LM,MSP,
1LIMIT(4),LIMIT(6),ETA,PHI,CHI,CNB,NSP,NBM)
KNM(1)=0
KNM(2)=0
KNM(3)=0
CPC=ELTIME(0)
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
CALL MOVE(0,AKN,NWEDG,XSTART,LIMIT(3),LIMIT(1),LIMIT(8),LIMIT(9),
1DELANG,BTA,C2,C3,DFA,PL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,CTI,CTR,
2CNI,CNR,ALPHA,SIGMA,COEPP,HTS,HTSI,NTS,UTL,UTT,VTS,PAU,PAV,PAW,
3PAX,PAY,PAZ,LCOL,TB,MSP,ER,CHI,CNG,CHG,NSP,UTLI,UTTI,VTSL,IFLUX)
KNM(1)=NM(1)
KNM(2)=NM(2)
KNM(3)=NM(3)
CPM=ELTIME(0)
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
TBI=TB(ND)
CALL FLOW(NWEDG,NEM,LARGE,BTA,C1,C7,C8,ENT,REN,LCOL,MSP,NM,SN,ST,
1TBI,PAU,PAV,PAW,PAX,PAY,PAZ,ER,CHI,CNG,CHG,NSP,JV,FCOL,VEL,PFV)
IF (LARGE.NE.0) GO TO 345
CPB=ELTIME(0)
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
CALL MOVE(1,AKN,NWEDG,ISTART,LIMIT(3),LIMIT(1),LIMIT(8),LIMIT(9),
1DELANG,BTA,C2,C3,DFA,PL,HTI,HTR,JNT,KNM,NM,XCB,XLIM,CTI,CTR,
2CNI,CNR,ALPHA,SIGMA,COEPP,HTS,HTSI,NTS,UTL,UTT,VTS,PAU,PAV,PAW,
3PAX,PAY,PAZ,LCOL,TB,MSP,ER,CHI,CNG,CHG,NSP,UTLI,UTTI,VTSL,IFLUX)
CPB=CPB+ELTIME(0)
IF (DEBUG(1)) WRITE(6,33) AIME,CPC,CPM,CPI,CPB,CPA,(NM(I),I=1,3),NMAX
DO 330 MT=1,MSP
DO 290 N=1,NBI
NB(MT,N)=0

```

RUN5060

RUN5130  
RUN5140

RUN5170

RUN5180  
RUN5190

RUN5220  
RUN5230

RUN5280  
RUN5290

RUN5330

RUN5420  
RUN5450

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290	CONTINUE	
	NG=NH(MT)	RUN5470
	N=0	RUN5480
295	N=N+1	RUN5490
	IF(N.GT.NG) GO TO 310	RUN5500
	X=PAY(MT,N)	RUN5510
	Y=PAY(MT,N)	RUN5520
	Z=PAZ(MT,N)	RUN5530
	R=SQRT(Y*Y+Z*Z)	RUN5540
	TANG=180.*ATAN2(Z,-Y)/PI	RUN5550
	IWDGE=TANG/DELANG+1	
	IF(IWDGE.LT.1) IWDGE=1	
	IF(IWDGE.GT.NWEDG) IWDGE=NWEDG	
	L=X/BW	
	IF(L.GE.NW) L=NW-1	RUN5610
	M=R/BH	RUN5620
	IF(M.GE.NH) M=NH-1	RUN5630
	K=NWEDG*(L*NH+M)+IWDGE	RUN5640
	IF(K.LE.NBX) GO TO 305	
	WRITE(6,40)L,M,IWDGE,NW,NH,NWEDG,K,MT,N,X,Y,Z,R,TANG	RUN5660
	IF(DUMP) CALL ABEND(4)	
	STOP	RUN5680
		RUN5690
305	J=NB(MT,K)+1	
	IF(J.LE.MNB) GO TO 308	
	IF(DEBUG(1)) WRITE(6,44) MT,K,MNB,AIME	
308	NB(MT,K)=J	RUN6400
	LB(N)=K	
	GO TO 295	RUN6430
310	CONTINUE	
	NB5(MT,1)=0	
	DO 320 M=1,NB5	RUN6450
	A=NB(MT,M)	RUN6480
	DB(MT,M)=A*DFA(MT)/PNE(M)	RUN6490
	NBM(MT,M+1)=NB5(MT,M)+NB(MT,M)	
	NBN(M)=NB5(MT,M)	
320	CONTINUE	RUN6500
	IF(N5(MT).GT.NMAX) NMAX=N5(MT)	RUN6510
	DO 325 N=1,NG	
	Q=LB(N)	
	NBN(Q)=NBN(Q)+1	
	NA=NBN(Q)	
325	LM(MT,NA)=N	
330	CONTINUE	RUN6520
	IF(SAMP.LT.ITS) GO TO 335	
	CALL ACCUM(NMC,NPB,PNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,HSP,	
	1NSP,NBM)	
	SAMP=0	RUN6570
	IF(TIME.LE.TST) GO TO 335	
	CALL AVRGE(PNB,DB,DBA,NB,NBT,XV,YV,ZV,XVA,YVA,ZVA,TMP,TMPA,TRP,TRP	
	1A,HSP,NSP)	
	NAV=NAV+1	RUN6600
335	CPA=ELTIME(0)	
	CPI=CPC+CPN+CPB+CPA	
	CPJ=2.*CPI+5.	
340	CPUTYM=TFIND(0)	

```

IF (DEBUG (3)) WRITE (6, 33) AIME, CPC, CPM, CPI, CPB, CPA, (NM (I), I=1, 3), NMAX
IF (TIME.EQ.0) GO TO 355
IF ((TIME.GE.TLIM).OR. (CPUTYM.LE.CPJ)) GO TO 345
IF (PRT.LT.ITP) GO TO 280
PRT=0
345 WRITE (6, 30) AIME, KAWLS
IF (DEBUG (3)) WRITE (6, 31) CPUTYM
WRITE (6, 32) (NM (I), I=1, 3)
WRITE (6, 34) ((NCOL (I, J), J=1, 3), I=1, 3)
WRITE (6, 35) (JNT (I), I=1, 3)
IF (LARGE.NE.0) GO TO 360
WRITE (6, 36) NMAX
IF (.NOT.SAVE) GO TO 355
IF (PRT.EE.0.AND.CPUTYM.GT.CPJ.AND.TIME.LT.TLIM) GO TO 355
REWIND 9
WRITE (9) DENP, U, XREF, TRP, KAWLS, NL, NW, NH, BW, BH, NREG, XLB, XLC, PI, ND,
2 S, SINANG, COSANG, AKN, AKT, NBI, RM, XR, TIME, DTM, TI, ITS, ITP, TST,
3 TLIM, RMA, RNU, DIR, XSTART, RNM, MNB, TR, BZC, CN7, DRP, PCP, PNA,
4 HTP, INM, LLM, NAV, NMAX, NWEDG, PRT, SAMP, AKN1, AKN2, AKT1, AKT2,
5 BTA, C1, C2, C3, C7, C8, DAM, DPA, PL, DELANG, PDN, HTI, HTR, JNT, KNE,
6 NB, WTM, C4, VEM, NCOL, CTI, CTE, CNI, CNR, SN,
7 ST, D1, D2, D3, D4, NRAM, VELR, RMP, RMN, RMP, IFLOX, FLUXIN,
8 ILM, COEPP, XCB, XS, YCB, TB, ALPHA, SIGMA, NTS,
9 UTL, UTT, VTS, HTS, HTSI, ENT, BEH, TMPA,
A DBA, NB, NBT, TMP, IV, YVA, YV, YVA, ZV, ZVA, T, DB, PNB, XC, YC, ZC,
B PAU, PAV, PAW, PAX, PAY, PAZ, LCOL, LM,
C ETA, PHI, CHI, CN, CM, CNG, CHG, CNB, TRP, TRPA, MSP, ANGLE, TP,
D UTLI, UTTI, VTSI, ER, RMB, LB, NEM, EBN, VEL, PVV, PCOL, JV
WRITE (6, 50)
355 CONTINUE
WRITE (6, 25) REP, RNM, RMP
DO 356 MT=1, MSP
PIN1=ENT (MT, 1)
PIN2=ENT (MT, 2)
RF1=IFLOX (MT, 1) *DTM/DT
RF2=IFLOX (MT, 2) *DTM/DT
RNF1=1.-RF1/PIN1
RNF2=(PIN2-RF2)/PIN1
RPS1=RNF1*FLOXIN (MT) /RNU (MT)
RPS2=RNF2*FLOXIN (MT) /RNU (MT)
356 WRITE (6, 26) BMA (MT), PIN1, RF1, RNF1, RPS1, PIN2, RF2, RNF2, RPS2
IF (TIME.EQ.0) GO TO 280
IF (TIME.LE.TST) GO TO 350
CALL PRINT1 (DT, COSANG, SINANG, RMA, RNU, DRP, PCP, HTP, PL, HTI, HTR, CTI,
1CTE, CNI, CNR)
CALL PRINT2 (AKN, XSTART, DT, RNU, BMA, DRP, PCP, HTP, UTLI, UTTI, VTSI, HTSI,
1DELANG, NWEDG, XS, XCB, YCB, HTS, NTS, UTL, UTT, VTS, LIMIT (3), LIMIT (1), MSP)
CALL PRINT4 (MSP, CHI, RNU, MSP, TRPA, PDN, WTM, DBA, NET, TMPA, IVA,
1YVA, ZVA, 1, NBT, XC, YC, ZC)
GO TO 353
350 CONTINUE
CALL PRINT4 (MSP, CHI, RNU, MSP, TRP, PDN, HTS, DB, NB, TMP, IV, YV, ZV,
10, NB, XC, YC, ZC)
353 IF (DEBUG (2)) WRITE (6, 1)
IF ((TIME.LT.TLIM).AND. (CPUTYM.GT.CPJ)) GO TO 280

```

RUN6650  
 RUN6660  
 RUN6670

RUN6740  
 RUN6750  
 RUN6900

RUN7050

RUN6860

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IF (IC.EQ.ICOPY) RETURN
IC=IC+1
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) (MT,RMA(MT),RNU(MT),PCOL(MT),FLUXIN(MT), (ENT(MT,K),K=1,2
1),MT=1,MSP)
WRITE(6,2)
WRITE(6,3) IC
CALL PRINTA(NWEDG,TITLE,NAME,ICB, YCB, TB, ALPHA, SIGMA, XLIM,
1COEFF, LIMIT, MSP)
CALL PRINTB(PNA, MSP, PNB, NM, XLIM, IC, YC, ZC, NB, NSP)
SAVE=.FALSE.
GO TO 345
360 WRITE(6,38) (DBG1(I,LARGE),I=1,3)
IF ((REDO).AND. (TIME.LE.TST)) GO TO 364
IF (DUMP) CALL ABEND(9)
STOP
364 CONTINUE
IF (NEW) GO TO 365
REWIND 9
READ(9) DENF, U, XREP, TRP, KAWLS, KL, NW, NE, BW, BH, NREG, XLB, XLC, PI, ND,
2 S, SINANG, COSANG, AKN, AKT, NBI, RM, XR, TIME, DTE, TI, ITS, ITP, TST,
3 TLIH, RMA, RNU, DIR, XSTART, MNH, MNB, TB, BZC, CN7, DRP, PCF, PNA,
4 HTP, INH, LLH, NAV, NMAX, NWEDG, PRT, SAMP, AKN1, AKN2, AKT1, AKT2,
5 BTA, C1, C2, C3, C7, C8, DAM, DPA, PL, DELANG, PDN, HTI, HTR, JNT, KMM,
6 NH, WTS, C4, VRH, NCOL, CTI, CTR, CNI, CNR, SN,
7 ST, D1, D2, D3, D4, NRAM, VELR, REP, RMN, RMP, IFLUX, FLUXIN,
8 XLIM, COEFF, ICB, XS, YCB, TB, ALPHA, SIGMA, NTS,
9 UTL, UTT, VTS, HTS, HTSI, ENT, REM, TRPA,
A DBA, NB, NBT, TMP, XV, IVA, YV, YVA, ZV, ZVA, T, DB, PNB, IC, YC, ZC,
B PAU, PAV, PAW, PAX, PAY, PAZ, LCOL, LH,
C ETA, PHI, CHI, CN, CH, CNG, CHG, CN8, TRP, TRPA, MSP, ANGLE, TP,
D UTLI, UTTI, VTSI, ER, RMB, LB, NEM, NBN, VEL, PPV, PCOL, JV
365 ANH=INH
INH=9*ANH/10
DDH=.9*DDN
DRP=DRP/.9
PCF=PCF/.9
HTP=HTP/.9
DO 370 MM=1, MSP
PDN(MM)=PDN(MM)*INH/ANH
LL(MM)=9*LL(MM)/10
DO 366 KK=1, MSP
366 CN8(KK,MM)=CN8(KK,MM)*.9
DO 370 NK=1, 2
ENT(MM,NK)=ENT(MM,NK)*INH/ANH
370 REM(MM,NK)=0.0
IF (NEW) GO TO 220
TST=TIME+TST
TI=0.
PRT=ITP
WRITE(6,2)
WRITE(6,4)
WRITE(6,5) (MT,RMA(MT),RNU(MT),PCOL(MT),FLUXIN(MT), (ENT(MT,K),K=1,2
1),MT=1,MSP)

```

RUN7080

RUN7120

RUN7130

RUN7180

RUN7190

RUN7200

RUN7220

RUN7230

RUN7260

```

WRITE (6,2)
IF ((LARGE.EQ.2).OR.(LARGE.EQ.3)) GO TO 280
REDO=.FALSE.
GO TO 360
END

```

RUN7410

```

SUBROUTINE DIAG(N, ITEST, NUM)
REAL*8 PARAM(10) /' NWEDGE', ' NREG', ' ND', ' NPX', '
1 HNS', ' ENB', ' NBX', ' NS', ' NJ', ' MSP' /

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DIAG010

DIAG040

DIAG050

DIAG060

DIAG070

DIAG080

DIAG090

DIAG100

---

**FORMATS**

```

32 FORMAT(9X, 'ENT, REM, ENTS, REMS, PTH, THETA, DTH')
42 FORMAT(///5X, 43H ARRAY DIMENSIONS ARE ABOUT TO BE VIOLATED./)
44 FORMAT(5X, 17H MAXIMUM VALUE IS, 15, 19H, WHEREAS YOU INPUT, 15, 3H (,
1A8, 1H))
56 FORMAT(/5X, 78H IF YOU DESIRE TO USE THIS VALUE, THE FOLLOWING ARE
1YS MUST BE RE-DIMENSIONED./)
62 FORMAT(9X, 'HTS, HTSI, NTS, NTSP, UTL, UTT, VTS')
64 FORMAT(9X, 'XLIM, COEPP' // 11X, 'NOTE THAT THE XLIM ARRAY MUST BE DIMEN
1SIONED TO 3 MORE THAN THE COEPP ARRAY.')
66 FORMAT(9X, 'XCB, XS, YCB, TB, ALPHA, SIGMA')
68 FORMAT(9X, 'DBA, NB, NBT, TMP, TPA, IV, XVA, YV, YVA, ZV, ZVA, T, DB')
70 FORMAT(9X, 'PAV, PAV, PAW, PAX, PAY, PAZ, LCOL')
72 FORMAT(8X, 3H LM)
74 FORMAT(///5X, 76H IF YOU CHANGE THE ARRAY DIMENSIONS, ALSO CHANGE TH
1E 'LIMIT' DATA STATEMENT.)
75 FORMAT(9X, 'ALL ARRAYS ASSOCIATED WITH SPECIES')
76 FORMAT(9X, 'PNB, XC, YC, ZC')
78 FORMAT(9X, 'PV, NTCV, NTCF, MS, IWS, SL, DELS, TANGN')
80 FORMAT(9X, 'VEL, PFV')

```

DIAG120

DIAG130

DIAG140

DIAG150

DIAG160

DIAG170

DIAG180

DIAG190

DIAG210

DIAG220

DIAG230

DIAG240

DIAG250

DIAG260

DIAG270

DIAG280

DIAG290

DIAG300

DIAG310

DIAG320

```

WRITE (6,42)
WRITE (6,44) ITEST, NUM, PARAM(N)
WRITE (6,56)
GO TO (1,2,3,4,5,6,7,8,9,10), N
1 WRITE (6,62)
WRITE (6,32)
GO TO 11
2 WRITE (6,64)
GO TO 11
3 WRITE (6,66)
WRITE (6,62)
GO TO 11
4 WRITE (6,68)
WRITE (6,76)
GO TO 11
5 WRITE (6,70)
GO TO 11
6 WRITE (6,72)
GO TO 11

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DIAG340

DIAG350

DIAG370

DIAG390

DIAG400

DIAG420

DIAG440

DIAG460

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7 WRITE(6,80) DIAG480  
GO TO 11  
8 WRITE(6,78) DIAG500  
GO TO 11  
9 WRITE(6,80) DIAG520  
GO TO 11  
10 WRITE(6,75)  
11 WRITE(6,74)  
STOP DIAG540  
END DIAG550

SUBROUTINE PRINTA(NWEDG, TITLE, NAME, XCB, YCB, TB, ALPHA, SIGMA, XLIM,  
1 COEFF, LIMIT, MSP)  
INTEGER TST, TLIM, TIME PRA0030  
LOGICAL SAVE, NEW PRA0040  
DIMENSION LIMIT(1), TITLE(6), NAME(2), XCB(1), YCB(1), TB(1)  
DIMENSION ALPHA(3,1), SIGMA(3,1), XLIM(1), COEFF(4)  
DIMENSION RNU(3), RMA(3), CHI(3), DIR(3,3), PHI(3,3), ETA(3,3)  
DIMENSION WTM(3), DAM(3,3), VEL5(3), XSP(3)  
COMMON /FIRST/NL, NW, NH  
COMMON /SECND/BW, BH, RMP, RMN, RMP  
COMMON /THIRD/PI, NREG, S, SINANG, COSANG, AKN, AKT, AKN1, AKN2, AKT1, AKT2  
COMMON /FOURTH/NBX, RM, XR, DUMP, C9, LL(3), LLM  
COMMON /FIFTH/ND, TIME, DTM, TI, ITS, ITP, TST, TLIM, RMA, RNU, DIR PRA0100  
COMMON /SIXTH/RMB, XSTART, INN, MNN, MNB, NEW, SAVE, PERCENT, NSR, TR  
COMMON/EIGHT/DENP, U, TP, ANGLE, TRP, CHI, PHI, ETA, WTM, DAM, VELR, XREF  
DATA NOT/'NOT '/ PRA0120  
PRA0130  
PRA0140  
PRA0150  
PRA0160  
PRA0170

FORMATS

1 FORMAT(16X,40('-'),T74,'I'//9X,'3-D',I2,'-FLUID PROGRAM - ') PRA0190  
2 FORMAT('+',31X,A4) PRA0200  
3 FORMAT('+',35X, 'A RESTART OF A PREVIOUS RUN',T74,'I'/12X,2PRA0210  
1A4,' - ',6A4,' - ',I2,' REGIONS',T74,'I',16(/T74,'I'))  
4 FORMAT(7X,'FRONT OF BODY =',E12.4,' XSTART MAX HEIGHT =',E12.4,  
1RMB',T74,'I'/7X,'BODY TEMPERATURES =',F12.2,' T PR.STRM.',F12.2,  
2' T ENTR.',F12.2,' T CAVITY'/7X,'X-LIMIT',T37,'BODY COEFFICIENTS',  
3T74,'I')  
6 FORMAT(5F14.6,3X,'I') PRA0240  
10 FORMAT(1X,72('-')) PRA0250  
12 FORMAT(//14X,'PARAMETERS OF SEGMENTS FOR BODY COLLISIONS',T96,'I'/  
18X,'X-COORD. TEMP. ALPHA1 ALPHA2 ALPHA3 SIGMA1 SIGMA2  
2 SIGMA3 AREAS',T96,'I')  
14 FORMAT(4X,E12.4,7P9.4,E12.4,T96,'I')  
17 FORMAT(///25X,'ARRAY STORAGE USED'/5X,I6,' \*',10I6,T96,'I')  
18 FORMAT(1H1/17X,'LENGTH OF CELL IN MEAN-FREE-PATHS. = ',F12.4,' BW'  
A,T76,'I'  
1/17X,'HEIGHT OF CELL IN MEAN-FREE-PATHS = ',F12.4,' BH',T76,'I'  
2/16X,'NUMBER OF L1 CELLS ALONG FLOW AXIS = ',I13,' NW',T76,'I'  
3/17X,'NUMBER OF L1 CELLS IN RADIAL DIR. = ',I13,' NH',T76,'I'  
4/21X,'NUMBER OF LEVELS OF CELL SIZE = ',I13,' NL',T76,'I')  
23 FORMAT(3X,'NUMBER OF AZIMUTHAL WEDGES WITHIN ',I3,' DEGREES = ',IPRA0470  
113,' NWEDG I'/)

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24 FORMAT(16X,'BASIC TIME INTERVAL FOR COLLISIONS =',E13.4,' DTM
1 I'/8X,'TIME INTERVAL FOR SAMPLING FLOW FIELD INPO =',E13.4,' DTS
2 I'/24X,'TIME INTERVAL FOR PRINTING =',E13.4,' DTP I'/9X,
3'TIME TO STEADY-STATE CONDITIONS (ASSUMED) =',E13.4,' TST I'/
419X,'TIME AT WHICH RUN IS TERMINATED =',E13.4,' TTIM I'/)
25 FORMAT(5X,'PRESSURE RATIO (INSIDE/ENTRANCE) - EITHER TYPE =',F13.5
1,' RMP',T76,'I'/6X,'DENSITY RATIO (INSIDE/ENTRANCE) - EITHER TYPE
2=',F13.5,' RDN',T76,'I'/9X,'FLUX RATIO (INSIDE/ENTRANCE) - EITHER
3TYPE =',F13.5,' RFP',T76,'I')
26 FORMAT(5X,'FREE STREAM NUMBER OF MOLECULES - EITHER TYPE =',I13,'
A INH I'/9X,'INITIAL NUMBER OF MOLECULES - MAXIMUM =',I13,
1' LLN I'/9X,'MAXIMUM NUMBER OF MOLECULES - EITHER TYPE =',I13
2,' MNB I'/1X,'MAX NUMBER OF MOLECULES IN ANY CELL - EITHER TY
3PE =',I13,' MNB I')
27 FORMAT(/22X,'VELOCITY OF FREE STREAM FLOW =',E13.4,' U',T76,'I'/1
19X,'SPEED RATIO OF FREE STREAM FLOW =',E13.4,' S',T76,'I'/19X,'MAC
AH NUMBER OF FREE STREAM FLOW =',E13.4,' M',T76,'I'/19X,'SPECIFIC H
BEAT RATIO (CALCULATED) =',E13.4,' GAMMA',T76,'I'/
35X,'ANG
2LE OF ATTACK =',F13.4,' ANGLE I'/16X,'NUMBER DENSITY OF FREE ST
3REAM FLOW =',E13.4,' N',T76,'I'/19X,'TEMPERATURE OF FREE STREAM FL
4OW =',F13.4,' TP',T76,'I'/16X,'MOLE FRACTIONS OF FREE STREAM FLOW
5=',3E13.4,' RND I'/16X,'MOLECULAR WEIGHTS OF SPECIES ABOVE =',3F13
6.4,' RRA I'/18X,'INITIAL NUMBERS OF SPECIES ABOVE =',3I13,' LL I'
7)
28 FORMAT(/10X,'REFERENCE TEMPERATURE FOR MOLECULAR DATA =',F13.4,'
1TRP',T90,'I'/14X,'CROSS-SECTION',26X,'TEMP EXPONENT',T90,'I'/3(3X,
23E12.4,3X,3F12.6,T90,'I'/)/ 5X,'CHI/2-1',11X,'ROTATIONAL PARAMETER
3 PHI',T90,'I'/3(F12.4,5X,3F12.6,T90,'I'/))
29 FORMAT(/9X,'DATA SAVED ON TAPE 9')
30 FORMAT( 31X,'REP MOLECULAR SPEED =',E13.4,' VELR',T76,'I'/20X,'SP
1ECIES FREE STREAM MOLECULAR SPEEDS',T76,'I'/14X,3E16.6,T76,'I'/26X
2,'REFERENCE MEAN FREE PATH =',E13.4,' XREP',T76,'I'/26X,'SPECIES M
3EAN FREE PATHS',T76,'I'/14X,3E16.6,T76,'I'/11X,'LONGITUDINAL KNUDS
4EN NUMBER (FREE STRM.) =',E13.4,' AKN',T76,'I'/13X,'TRANSVERSE KNUD
5SEN NUMBER (FREE STRM.) =',E13.4,' AKT',T76,'I'/11X,'LONGITUDINAL K
6NUDSEN NUMBER ( ENTRANCE ) =',E13.4,' AKN1',T76,'I'/13X,'TRANSVERSE
7 KNUDSEN NUMBER ( ENTRANCE ) =',E13.4,' AKT1',T76,'I'/11X,'LONGITUD
8INAL KNUDSEN NUMBER ( CAVITY ) =',E13.4,' AKN2',T76,'I'/13X,'TRAN
9VERSE KNUDSEN NUMBER ( CAVITY ) =',E13.4,' AKT2',T76,'I')

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PRA0680  
PRA0690  
PRA0700  
PRA0710  
PRA0720  
PRA0730

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IARRAY=708+LIMIT(3)*(32+56*LIMIT(1))+20*LIMIT(2)+LIMIT(4)*(120+4*LPRA0710
LIMIT(6))+56*LIMIT(5)+LIMIT(8)*(68+96*LIMIT(9))+20*LIMIT(7)+224*LIMIT(1)
2IT(1)
WRITE(6,1) MSP
IF(NEW) WRITE(6,2) NOT
WRITE(6,3) NAME,TITLE,NREG
WRITE(6,4) ISTART,RMB
DO 100 I=1,NREG
100 WRITE(6,6) XLIM(I+1), (COEFF(J),J=1,4)
WRITE(6,10)
WRITE(6,12)
DO 110 I=1,ND
110 WRITE(6,14) XCB(I),TB(I), (ALPHA(J,I),J=1,3), (SIGMA(J,I),J=1,3),YCB

```

PRA0750  
PRA0760  
PRA0770  
PRA0780  
PRA0800  
PRA0810  
PRA0820

FILE: GKBINT AUG82 A

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```

1(I)
WRITE(6,10) PRA0850
WRITE(6,17) IARRAY, (LIMIT(I), I=1,10) PRA0880
WRITE(6,18) BW, BH, NW, NE, NL
IETAZ=180.
WRITE(6,23) IETAZ, NWEDG PRA0940
DTS=DTM*ITS PRA0950
DTP=DTM*ITP PRA0960
AST=DTM*TST PRA0970
ALIM=DTM*TLIM
CHT=0.0
DO 120 J=1, MSP
120 CHT=CHT+CHI(J)*RNU(J)
GAMMA=(7.+2.*CHT)/(5.+2.*CHT)
AN=S*SQRT(2./GAMMA) PRA0980
WRITE(6,24) DTM, DTS, DTP, AST, ALIM
WRITE(6,26) INM, LLM, MNM, MNB
WRITE(6,25) RMP, RMN, RMP
WRITE(6,27) U, S, AN, GAMMA, ANGLE, DENF, TF, (RNU(I), I=1,3), (RMA(I), I=1,3)
1), (LL(I), I=1,3)
WRITE(6,28) TRF, ((DIR(I,K), K=1,3), (ETA(I,K), K=1,3), I=1,3), (CHI(I),
1(PHI(I,K), K=1,3), I=1,3)
DO 210 I=1,3
VELS(I)=0.0
210 XSP(I)=0.0
DO 220 J=1, MSP
VELS(J)=VELR/SQRT(WTM(J))
XT=0.0
DO 215 M=1, MSP
215 XT=XT+RNU(M)*DAM(J,M)*SQRT(1.+WTM(J)/WTM(M))
220 XSP(J)=1.414214*XREF/XT
WRITE(6,30) VELR, (VELS(I), I=1,3), XREF, (XSP(I), I=1,3), AKN, AKT,
1AKN1, AKT1, AKN2, AKT2
IF(SAVE) WRITE(6,29) PRA1040
RETURN PRA1050
END PRA1060

SUBROUTINE PRINTB(PNA, MSP, PNB, NE, XLIM, XC, YC, ZC, NB, N)
INTEGER*2 NB
DIMENSION PNB(1), NE(1), XLIM(1), XC(1)
DIMENSION YC(1), ZC(1), NB(N,1)
COMMON /FIRST/NL, NW, NH
COMMON /THIRD/PI, NREG, S, SINANG, COSANG, AKN, AKT
COMMON /PORTH/NBX PRA0070
1 FORMAT(1H1) PRA0080
2 FORMAT(2X, '-----CELL GEOMETRY-----' PRA0090
1-----'/2X, 'BOX LEVEL POSITION OF CENTER VOLUME PRA0100
2INITIAL POPULATION'/2X, 'NUM.', 12X, 'X', 7X, 'Y THETA', 12X, ' TOT
3AL ', ' EACH SPECIES ', ' CELL')
3 FORMAT(1X, I4, I5, 3X, 2F8.3, F7.1, E12.3, 2X, I4, 2X, 3I5, 3X, I4)
4 FORMAT(2X, '-----TOTALS-----', E12.4, 8X, 3I5) PRA0150
WRITE(6,1)
WRITE(6,2)
DO 200 I=1, NBX PRA0170
X=(XC(I)-XLIM(1))*AKN PRA0180

```



```

Y=YC(I)*AKT*2.0
IX=X*NW+1
IY=Y*NH+1
M1=NB(1,I)
M2=0
IF(MSP.GE.2) M2=NB(2,I)
M3=0
IF(MSP.GE.3) M3=NB(3,I)
MM=M1+M2+M3
140 WRITE(6,3) IX,IY,X,Y,ZC(I),PNB(I),MM,M1,M2,M3,I
200 CONTINUE
NM2=0
IF(MSP.GE.2) NM2=MM(2)
NM3=0
IF(MSP.GE.3) NM3=MM(3)
WRITE(6,4) PNA,MM(1),NM2,NM3
RETURN
END

```

PRB0370  
PRB0390  
PRB0400

```

SUBROUTINE CELL(A,B,KW,KH,XO,DELANG,NWEDG,XC,YC,ZC,PNB)
DIMENSION XC(1),YC(1),ZC(1),PNB(1)
COMMON /THIRD/PI

```

-----

THE PURPOSE OF THIS SUBROUTINE IS TO

1. COMPUTE THE VOLUME OF EACH CELL (ALL 3 POSSIBLE LEVELS) AND STORE THE RESULT IN THE ARRAY CALLED 'PNB'.
2. COMPUTE THE X, Y, AND THETA COORDINATES OF THE CENTER OF EACH CELL (ALL 3 POSSIBLE LEVELS) AND STORE THE RESULTS IN ARRAYS CALLED 'XC', 'YC', AND 'ZC'.

CELL030  
CELL040  
CELL050  
CELL060  
CELL070  
CELL080  
CELL090  
CELL100  
CELL110  
CELL120  
CELL130

```

I=0
X=XO-0.5*A
FACTOR=DELANG*PI*B*B*A/180.
DO 110 K=1,KW
X=X+A
Y=-.5*B
DO 110 L=1,KH
Y=Y+B
Z=-.5*DELANG
DO 110 M=1,NWEDG
Z=Z+DELANG
I=I+1
XC(I)=X
YC(I)=Y
ZC(I)=Z
110 PNB(I)=FACTOR*(2*L-1)
RETURN
END

```

CELL360  
CELL370

```

SUBROUTINE IMPACT(RM,G1,G2,G3,ET,EI,PHI,CHI,ETA,XM,CIN)
COMMON/THIRD/PI
IF(PHI.EQ.0.) GO TO 20

```

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```

IF (CHI.EQ.0.) GO TO 20
DP=PHI*CHI-1.
DS=PHI*(2.-.5*ETA)-1.
E=ET+EI
10 X=Rand(0)
IF (X.EQ.0.0) GO TO 10
XT=X**DP*(1.-X)**DS
IF (XT.GT.IX) GO TO 15
CIN=CIN+XT.
IF (CIN.LT.IX) GO TO 10
CIN=CIN-IX
15 ET=(1.-PHI)*ET+(1.-X)*PHI*E
EI=(1.-PHI)*EI+X*PHI*E
20 GP=SQRT(ET/RM)
EP=2.*PI*Rand(0)
CSX=2.*Rand(0)-1.
SSX=SQRT(1.-CSX**2)
G1=GP*CSX
G2=GP*SSX*cos(EP)
G3=GP*SSX*sin(EP)
RETURN
END

```

```

SUBROUTINE GAS (NWEDG, DELANG, ND, BTA, C1, DPA, NM, FNB, DB, NB, NBM, NBN,
1PAU, PAV, PAW, PAX, PAY, PAZ, XLIM, COEPP, LM, I2, I3, XCB, TB, LARGE,
2MMH, MNB, DEBUG1, LCOL, IP, ER, CHI, CNG, CMG, I, LB)
INTEGER*2 LM(I, 1), LCOL(I, 1), LB(1), NBM(I, 1), NBN(1)
INTEGER*2 NB
LOGICAL DUMP, DEBUG1
DIMENSION BTA(1), C1(1), DPA(1), NM(1), FNB(1), CHI(1)
DIMENSION DB(I, 1), NB(I, 1), PAU(I, 1), PAV(I, 1), PAW(I, 1)
DIMENSION PAX(I, 1), PAY(I, 1), PAZ(I, 1), ER(I, 1), COEPP(4), XLIM(1)
DIMENSION CNG(1), CMG(1), XCB(1), TB(1)
COMMON /FIRST/NL, NW, NE
COMMON /SECND/BW, BH, BHP, RMN, RMP
COMMON /THIRD/PI, NREG, S, SINANG, COSANG, AKN, AKT
COMMON /FOURTH/NBX, RM, XR, DUMP, C9, LL(3)

```

GAS0060

-----GAS0150

THE PURPOSE OF THIS SUBROUTINE IS TO

1. COMPUTE THE INITIAL VELOCITY OF EACH MOLECULE.  
THE VELOCITY ARRAYS ARE 'PAU', 'PAV', AND 'PAW'.
2. COMPUTE THE INITIAL POSITION OF EACH MOLECULE.  
THE POSITION ARRAYS ARE 'PAX', 'PAY', AND 'PAZ'.
3. CREATE AN ARRAY WHICH STORES THE CELL POPULATIONS-  
'NB' FOR THE ACTUAL POPULATIONS
4. CREATE A CROSS-REFERENCING ARRAY (WHOSE CONSTRUCTION IS  
INDICATED BY A COMMENT CARD) CALLED 'LM'.
5. COMPUTE AN ARRAY WHICH STORES THE NUMBER DENSITY IN EACH  
CELL

GAS0160  
GAS0170  
GAS0180  
GAS0190  
GAS0200  
GAS0210  
GAS0220  
GAS0230  
GAS0250  
GAS0260  
GAS0270  
GAS0280  
GAS0290  
-----GAS0300

```

2 FORMAT(/' SOMETHING IS WRONG WITH BOX NUMBERING IN GAS'/9I5,5E14.5
1//)
3 FORMAT(/' SOMETHING WRONG IN CELL VOLUMES IN GAS'/5I,5I5,2E14.5)

```

```

4  FORMAT(' NB(',I2,',',I4,') POPULATION EXCEEDED ',I3,' IN GAS')      GAS0340
   BP=(1.-RNN)
   CP=1.-RNN**2
   DO 180 MT=1,IP
   N=0
110  N=N+1
      IF(N.GT.LL(MT)) GO TO 180
      IF(N.GT.RNN) GO TO 190
120  P=.001+.998*RAND(0)
      IP(BP.NE.0.0) P=(1.-SQRT(1.-CP*P))/BP
      X=XLIS(1)+P*IR
      R=RN*SQRT(RAND(0))
      D=PI*RAND(0)
      PAY(MT,N)=X
      PAZ(MT,N)=R*COS(D)
      PAZ(MT,N)=R*SIN(D)
      DO 126 J=2,ND
      IF(X.LE.ICB(J)) GO TO 128
126  CONTINUE
      WRITE(6,3) J,ND,MT,LL(MT),N,X,XCB(J)
      IF(DUMP) CALL ABEND(14)
      STOP
128  TL=TB(J-1)+(TB(J)-TB(J-1))*(X-ICB(J-1))/(ICB(J)-ICB(J-1))
130  V=4.*RAND(0)
      VV=V*V
      C1(MT)=C1(MT)+VV*EXP(1.-VV)
      IF(C1(MT).LT.1.) GO TO 130
      C1(MT)=C1(MT)-1.
      A=1.-2.*RAND(0)
      B=SQRT(1.-A*A)
      C=2.*PI*RAND(0)
      V=V/BTA(MT)*SQRT(TL)
      PAU(MT,N)=V*A
      PAV(MT,N)=V*B*COS(C)
      PAW(MT,N)=V*B*SIN(C)
      EI=0.0
      IF(CHI(MT).LE.-1.) GO TO 136
135  EI=9.*RAND(0)
      IF(EI.EQ.0.0) GO TO 135
      XT=EI**CHI(MT)*EXP(-EI)
      IF(XT.GE.CMG(MT)) GO TO 136
      CNG(MT)=CMG(MT)+XT
      IF(CNG(MT).LT.CMG(MT)) GO TO 135
      CNG(MT)=CMG(MT)-CNG(MT)
136  ER(MT,N)=EI*TL
      TANG=180.*(1.-D/PI)
      IWDGE=TANG/DELANG+1
      IF(IWDGE.GT.NWEDG) IWDGE=NWEDG
      L=X/BR
      IF(L.GE.NW) L=NW-1
      H=R/BR
      IF(H.GE.NH) H=NH-1
      K=NWEDG*(L*NH+H)+IWDGE
      IF(K.LE.HBX) GO TO 165
      WRITE(6,2) L,N,IWDGE,NW,NH,NWEDG,K,MT,N,X,Y,Z,R,TANG

```

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```

IF (DUMP) CALL ABEND(11)
STOP
165 J=NB (MT, K) +1
LCOL (MT, N)=0
IF (J.LE.HNB) GO TO 166
IF (DEBUG1) WRITE (6,4) MT,K,HNB
166 NB (MT, K) =J
LB (N)=K
167 IF (N.LT.LL (MT)) GO TO 110
NM (MT) =N
NBM (MT, 1)=0
DO 170 N=1, NBI
A=NB (MT, N)
DB (MT, N) =A*DFA (MT) /PNB (N)
NBM (MT, N+1) =NBM (MT, N) +NB (MT, N)
NBN (N) =NBM (MT, N)
170 CONTINUE
NG=NM (MT)
DO 175 N=1, NG
NQ=LB (N)
NBN (NQ) =NBN (NQ) +1
NA=NBN (NQ)
175 LH (MT, NA) =N
180 CONTINUE
RETURN
190 LARGE=1
RETURN
END

```

GAS0830  
GAS0840  
GAS1190  
GAS1240  
GAS1250  
GAS1270  
GAS1310  
GAS1330  
GAS1360  
GAS1370  
GAS1380  
GAS1390  
GAS1400  
GAS1410  
GAS1420  
GAS1430

SUBROUTINE FLOW (NWEDG, HNB, LARGE, BTA, C1, C7, C8, ENT, REM, LCOL, IP, NM,  
1SN, ST, TBI, PAU, PAV, PAW, PAX, PAY, PAZ, ER, CHI, CNG, CMG, I, JV, FCOL, VEL,  
2PFV)

```

INTEGER*2 LCOL
DIMENSION BTA (1), NM (1), SN (1), ST (1)
DIMENSION C1 (1), C7 (1), C8 (1), CMG (1), FCOL (1), VELK (4)
DIMENSION PAU (I, 1), PAV (I, 1), PAW (I, 1), PAX (I, 1), PAY (I, 1), PAZ (I, 1)
DIMENSION ENT (3, 1), REM (3, 1), LCOL (I, 1), ER (I, 1), CHI (1), CNG (1)
DIMENSION VEL (JV, 4, 1), PFV (JV, 4, 1)
COMMON /THIRD/PI
COMMON /PORTH/NBX, RM, IR

```

FLO0100  
FLO0110  
FLO0120  
FLO0130  
FLO0140  
FLO0150

-----  
THE PURPOSE OF THIS SUBROUTINE IS TO ADD A NEW BATCH OF MOLECULES  
TO THE SAMPLE THROUGH THE UPSTREAM BOUNDARY.  
-----

FLO0180  
FLO0190

```

DO 370 MT=1, IP
XGO=0.
E=1.
FRAC=FCOL (MT)
ARG=SN (MT)
STT=ST (MT)
TV=1./BTA (MT)
TR=1.
DO 180 NT=1, 2
VM=(SQRT (ARG**2+2.) +ARG) /2.
SE=AMAX1 (0.0, VM-4.)

```

FLO0200

```

SMH=VM+4.-SM
AM=ENT(MT,NT)+RES(MT,NT)
M=AM
RES(MT,NT)=AM-M
IF(M.EQ.0) GO TO 170
DO 160 N=1,M
IF(NH(MT).GE.SMH) GO TO 380
NH(MT)=NH(MT)+1
NMN=NH(MT)
R=RE*SQRT(RAND(0))
D=PI*RAND(0)
PAY(MT,NMX)=R*COS(D)
PAZ(MT,NMX)=R*SIN(D)
LCOL(MT,NMX)=0
IF(PFAC.EQ.0.0) GO TO 130
PF=RAND(0)
IF(PF.GT.PFAC) GO TO 130
KMX=3
VELK(4)=0.0
IF(CHI(MT).GT.-1) KMX=4
DO 110 K=1,KMX
P=RAND(0)
DO 102 J=2,JV
IF(PPV(J,K,MT).GT.P) GO TO 105
102 CONTINUE
VELK(K)=VEL(JV,K,MT)
GO TO 110
105 VELK(K)=VEL(J-1,K,MT)+(P-PPV(J-1,K,MT))*(VEL(J,K,MT)-VEL(J-1,K,
1MT))/ (PPV(J,K,MT)-PPV(J-1,K,MT))
110 CONTINUE
PAU(MT,NMX)=VELK(1)
PAV(MT,NMX)=VELK(2)
PAW(MT,NMX)=VELK(3)
ER(MT,NMX)=VELK(4)
GO TO 160
130 V=SM+RAND(0)*SMH
C1(MT)=C1(MT)+V*EXP(VH**2-V**2+2.*ARG*(V-VH))/VM
IF(C1(MT).LT.1.) GO TO 130
C1(MT)=C1(MT)-1.
PAU(MT,NMX)=E*V*TV
140 V=8.*RAND(0)-4.
C7(MT)=C7(MT)+EXP(-V*V)
IF(C7(MT).LT.1.) GO TO 140
C7(MT)=C7(MT)-1.
PAV(MT,NMX)=STT+V*TV
150 V=8.*RAND(0)-4.
C8(MT)=C8(MT)+EXP(-V*V)
IF(C8(MT).LT.1.) GO TO 150
C8(MT)=C8(MT)-1.
PAW(MT,NMX)=V*TV
X=0.0
IF(CHI(MT).LE.-1.) GO TO 156
155 X=9.*RAND(0)
IF(X.LE.0.0) GO TO 155
XT=X**CHI(MT)*EXP(-X)

```

FLO0340

FLO0360

FLO0370

FLO0390

FLO0420

FLO0430

FLO0480

FLO0490

FLO0510

FLO0520

FLO0530

FLO0540

FLO0560

FLO0570

FLO0580

FLO0590

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```
IF (XT.GE.CMG(MT)) GO TO 156
CMG(MT)=CMG(MT)+IT
IF (CMG(MT).LT.CMG(MT)) GO TO 155
CMG(MT)=CMG(MT)-CMG(MT)
156 ER(MT,NMX)=X*TR
160 PAX(MT,NMX)=XGO
170 CONTINUE
ARG=0.0
XGO=XR
E=-1.
FRAC=0.0
STT=0.0
TV=SQRT(TBI)/BTA(MT)
TR=TBI
180 CONTINUE
370 CONTINUE
RETURN
380 LARGE=2
RETURN
END
```

FLO0610  
FLO0620

FLO0640  
FLO0650

FLO0660  
FLO1140  
FLO1150  
FLO1160  
FLO1170  
FLO1180

```
SUBROUTINE COLIDE(CF,CM,WTH,DB,DBA,NB,NCOL,LCOL,PAU,PAV,PAW,ER,T,
1LM,MT,I2,I3,ETA,PHI,CHI,CN8,NP,NBM)
INTEGER TIME
INTEGER*2 LM(NP,1), LCOL(NP,1)
INTEGER*2 NBM,NB
DIMENSION CN(3,3,1),CM(3,3,1),WTH(1),DB(NP,1),DBA(NP,1),NB(NP,1)
DIMENSION NBE(NP,1),NCOL(3,1),T(NP,NP,1),ETA(3,1),PHI(3,1),CHI(1)
DIMENSION PAU(NP,1),PAV(NP,1),PAW(NP,1),ER(NP,1),CN8(3,1),WA(2)
COMMON /PORTH/NBX
COMMON /PIPTH/ND,TIME,DTH
```

COL0030

COL0080  
COL0090  
COL0100

-----  
THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE ELAPSED TIMES IN  
CELLS BY AN AMOUNT APPROXIMATELY EQUAL TO THE PRE-SELECTED COLLIS  
TIME. THERE ARE FOUR TYPES FOR EACH CELL, SAVED IN AN ARRAY CALLE  
'T', CORRESPONDING TO THE FOUR TYPES OF MOLECULAR COLLISIONS WHI  
CAN OCCUR. TO ADVANCE THE VARIOUS TIMES, AN APPROPRIATE NUMBER OF  
THE CORRESPONDING MOLECULAR COLLISIONS IS COMPUTED. THE ACTUAL  
MOLECULES TO COLLIDE ARE SELECTED AT RANDOM, AND THEIR VELOCITY  
DIRECTIONS AFTER COLLISION ARE SELECTED AT RANDOM.  
-----

```
AIMC=DTH*TIME
DO 240 MTA=1,MT
DO 230 MTB=1,MTA
D = WTH(MTA) + WTH(MTB)
WA(MTA)=WTH(MTA)/D
WA(MTB)=WTH(MTB)/D
RM=WTH(MTA)*WTH(MTB)/D
CHT=CHI(MTA)+CHI(MTB)+2.0
PHT=PHI(MTA,MTB)
ETT=ETA(MTA,MTB)
DO 220 N=1,NB
IF (T(MTA,MTB,N).LT.AIME) GO TO 100
IF (T(MTB,MTA,N).GE.AIME) GO TO 220
100 NA=NB(MTA,N)*NB(MTB,N)
```

```

IF (MTA.EQ.MTB) NA=(NA-NB(MTA,N))/2.
IF (NA.LT.1) GO TO 220
KS=0
120 KC=0
   CPUT=ELTIME(0)
   KS=KS+1
   IF (KS.GT.NA) GO TO 220
130 KC=KC+1
   IF (KC.GT.NA) GO TO 220
135 I=NB(MTA,N)*RAND(0)+1+NBM(MTA,N)
   IF (I.GT.NBM(MTA,N+1)) I=NBM(MTA,N+1)
   J=LM(MTA,I)
140 K=NB(MTB,N)*RAND(0)+1+NBM(MTB,N)
   IF (K.GT.NBM(MTB,N+1)) K=NBM(MTB,N+1)
   IF (MTA.EQ.MTB.AND.I.EQ.K) GO TO 140
   L=LM(MTB,K)
   GM1=WA(MTA)*PAU(MTA,J)+WA(MTB)*PAU(MTB,L)
   GM2=WA(MTA)*PAV(MTA,J)+WA(MTB)*PAV(MTB,L)
   GM3=WA(MTA)*PAW(MTA,J)+WA(MTB)*PAW(MTB,L)
   G1=PAU(MTA,J)-PAU(MTB,L)
   G2=PAV(MTA,J)-PAV(MTB,L)
   G3=PAW(MTA,J)-PAW(MTB,L)
   GS=G1**2+G2**2+G3**2
   IF (GS.LT.1.0E-8) GO TO 130
   ET=RM*GS
   EI=ER(MTA,J)+ER(MTB,L)
   VR=GS*(.5-ETT/2.)
   IF (VR.GE.CN(MTA,MTB,1)) GO TO 160
   CN(MTA,MTB,1)=CN(MTA,MTB,1)+VR
   IF (CN(MTA,MTB,1).LT.CN(MTA,MTB,1)) GO TO 130
   CN(MTA,MTB,1)=CN(MTA,MTB,1)-CN(MTA,MTB,1)
160 CONTINUE
   CPUT=ZLTIME(0)
   CALL IMPACT(RM,G1,G2,G3,ET,EI,PHT,CHT,ETT,CN(MTA,MTB,2),CN(MTA,MTB
1,2))
165 CONTINUE
   CPUT=ELTIME(0)
   IF (PHT.EQ.0.) GO TO 175
   X1=0.0
   IF (CHI(MTA).EQ.-1.) GO TO 175
   X1=1.0
   IF (CHI(MTB).EQ.-1.) GO TO 175
170 X1=RAND(0)
   IF ((CHI(MTA).EQ.0.).AND.(CHI(MTB).EQ.0.)) GO TO 175
   XT=X1**CHI(MTA)*(1.-X1)**CHI(MTB)
   IF (XT.GT.CN(MTA,MTB,3)) GO TO 175
   CN(MTA,MTB,3)=CN(MTA,MTB,3)+XT
   IF (CN(MTA,MTB,3).LT.CN(MTA,MTB,3)) GO TO 170
   CN(MTA,MTB,3)=CN(MTA,MTB,3)-CN(MTA,MTB,3)
175 CONTINUE
   C=DBA(MTA,N)
   D=DBA(MTB,N)
   IF (C.EQ.0.0) C=DB(MTA,N)
   IF (D.EQ.0.0) D=DB(MTB,N)
   IF (T(MTA,MTB,N).GE.AIME) GO TO 180

```

```

PAU (MTA, J) = GM1 + WA (MTB) * G1
PAV (MTA, J) = GM2 + WA (MTB) * G2
PAW (MTA, J) = GM3 + WA (MTB) * G3
IF (PHT.GT. 0.) ER (MTA, J) = I1 * EI
LCOL (MTA, J) = 1 + LCOL (MTA, J)
NCOL (MTA, MTB) = NCOL (MTA, MTB) + 1
T (MTA, MTB, N) = T (MTA, MTB, N) + CNB (MTA, MTB) / NB (MTA, N) / D / VR
IF (MTA.EQ. MTB) GO TO 190
180 IF (T (MTB, MTA, N) .GE. AIME) GO TO 210
190 PAU (MTB, L) = GM1 - WA (MTA) * G1
PAV (MTB, L) = GM2 - WA (MTA) * G2
PAW (MTB, L) = GM3 - WA (MTA) * G3
IF (PHT.GT. 0.) ER (MTB, L) = (1. - I1) * EI
LCOL (MTB, L) = 1 + LCOL (MTB, L)
NCOL (MTB, MTA) = NCOL (MTB, MTA) + 1
T (MTB, MTA, N) = T (MTB, MTA, N) + CNB (MTB, MTA) / NB (MTB, N) / C / VR
210 CONTINUE
IF (T (MTA, MTB, N) .LT. AIME .OR. T (MTB, MTA, N) .LT. AIME) GO TO 120
220 CONTINUE
230 CONTINUE
240 CONTINUE
RETURN
END

```

COL0810

COL0920

SUBROUTINE MOVE (KSWCH, AKN, NWEDG, XSTART, I2, I3, I4, I5, DELANG,  
1BTA, C2, C3, DFA, FL, HTI, HTR, JNT, KNE, NM, XCB, XLIM, CTI, CTR, CNI,  
2CNR, ALPHA, SIGMA, COEPP, HTS, HTSI, NTS, UTL, UTT, VTS, PAU, PAV, PAW, PAX,  
3PAY, PAZ, LCOL, TB, IP, ER, CHI, CNG, CHG, I, UTLLI, UTTI, VTSI, IFLUX)

INTEGER\*2 LCOL (I, 1)

INTEGER TST, TIME

LOGICAL DUMP

REAL LAM, MU, NU

DIMENSION BTA (1), C2 (1), C3 (1), FL (1), HTI (1)

DIMENSION HTR (1), TB (1), XCB (1), ALPHA (3, 1), SIGMA (3, 1), COEPP (4)

DIMENSION PAU (I, 1), PAV (I, 1), PAW (I, 1), CTI (3, 1), CTR (3, 1)

DIMENSION CNI (3, 1), CNR (3, 1), DFA (1), JNT (1), XLIM (1), KNE (1), NM (1)

DIMENSION HTS (3, I2, I3), HTSI (3, I2, I3), NTS (3, I2, I3)

DIMENSION UTLLI (3, I2, I3), UTTI (3, I2, I3), VTSI (3, I2, I3)

DIMENSION UTL (3, I2, I3), UTT (3, I2, I3), VTS (3, I2, I3)

DIMENSION PAX (I, 1), PAY (I, 1), PAZ (I, 1)

DIMENSION ER (I, 1), CNG (1), CHG (1), CHI (1), IFLUX (3, 2)

COMMON /THIRD/PI, NREG

COMMON /FORTH/NBX, RM, XR, DUMP

COMMON /FIFTH/ND, TIME, DTM, TI, ITS, ITP, TST

COMMON /SVNTH/LAM, MU, NU, NT, N, J, XI, YI, ZI, TUSE

NAMelist/CHECK/TIME, X, Y, Z, DX, DY, DZ, TLEFT, RADS, RMS, XR

MOV0070

MOV0080

MOV0170

MOV0180

MOV0200

----- 210  
THE PURPOSE OF THIS SUBROUTINE IS TO ADVANCE THE SPATIAL POSITION 220  
OF ALL THE MOLECULES BY AN AMOUNT APPROPRIATE TO THEIR CURRENT VE 230  
LOCITIES AND THE PRE-SELECTED COLLISION TIME. 240  
----- 250

NAREA = NREG + 1

RMS = RM \*\* 2

DO 150 MT = 1, IP

N = KNE (MT)

MOV0270

MOV0290



```

10 H=N+1
    TLEPT=DTE
    IF (KSWCB.EQ.1) TLEPT=TLEPT*RAND(0)
    IF (N.GT.NH(MT)) GO TO 150
15 LAM=PAU(MT,N)
    NU=PAV(MT,N)
    NU=PAW(MT,N)
    XI=PAY(MT,N)
    YI=PAY(MT,N)
    ZI=PAZ(MT,N)
    DX=TLEPT*LAM
    DY=TLEPT*NU
    DZ=TLEPT*NU
    X=XI+DX
    Y=YI+DY
    Z=ZI+DZ
    RADS=Y**2+Z**2
    TUSE=TLEPT
    KS=0
    IF (RADS.GT.RMS) CALL INTERS(X,Y,Z,RMS,KS)
    IF (X.GT.XLIM(1)) GO TO 18
    IFLUX(MT,1)=IFLUX(MT,1)+1
    GO TO 100
18 IF (X.LT.XLIM(NAREA)) GO TO 19
    IFLUX(MT,2)=IFLUX(MT,2)+1
    GO TO 100
19 CONTINUE
    IF (KS.GT.0) CALL DRAG(AKN,NWEDG,XSTART,I2,I3,I4,I5,DELANG,JNI,BTA,
1C2,C3,DFR,PL,HTI,HTB,TB,XCB,CTI,CTR,CNI,CNR,ALPHA,SIGMA,EM,HTS,
2HTSI,NTS,UTL,UTT,VTS,LCOL,IP,EE,CHI,CNG,CMG,I,UTLI,UTTI,VTSI)
    IF (Z.GT.0.0) GO TO 20
    Z=-Z
    NU=-NU
20 CONTINUE
    PAX(MT,N)=X
    PAY(MT,N)=Y
    PAZ(MT,N)=Z
    PAU(MT,N)=LAM
    PAV(MT,N)=NU
    PAW(MT,N)=NU
    IF ((X.GT.0.0).AND.(X.LT.XR).AND.((Y**2+Z**2).LT.RMS)) GO TO 25
    WRITE(6,CHECK)
    GO TO 100
25 CONTINUE
    TLEPT=TLEPT-TUSE
    IF (TLEPT.GT.0.0) GO TO 15
    GO TO 10
100 NZ=NH(MT)
    PAX(MT,N)=PAX(MT,NZ)
    PAY(MT,N)=PAY(MT,NZ)
    PAZ(MT,N)=PAZ(MT,NZ)
    PAU(MT,N)=PAU(MT,NZ)
    PAV(MT,N)=PAV(MT,NZ)
    PAW(MT,N)=PAW(MT,NZ)
    ER(MT,N)=ER(MT,NZ)

```

```

HOV0300
HOV0310
HOV0320
HOV0330
HOV0340
HOV0360
HOV0370
HOV0380
HOV0390
HOV0400
HOV0410
HOV0420
HOV0430
HOV0440
HOV0450
HOV0460

```

```

HOV0820
HOV0830
HOV0840
HOV0850
HOV0860
HOV0870
HOV0880

```

```

LCOL(MT,N)=LCOL(MT,NZ)-
N=N-1
NM(MT)=NM(MT)-1
GO TO 10
150 CONTINUE
RETURN
END

```

MOV0900  
 MOV0910  
 MOV0920  
 MOV0930  
 MOV0940  
 MOV0950  
 MOV0960

```

SUBROUTINE ACCUM(I2,I3,FNB,NB,PAU,PAV,PAW,ER,TMP,TRP,XV,YV,ZV,LM,
1IP,I,NBM)
INTEGER*2 LM(I,1),NBM(I,1)
INTEGER*2 NB
DIMENSION FNB(1),NB(I,1),PAU(I,1),PAV(I,1),PAW(I,1),TMP(I,1)
DIMENSION XV(I,1),YV(I,1),ZV(I,1),ER(I,1),TRP(I,1)
COMMON /PORTH/NBX

```

---

THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE TEMPERATURES, VELOCITIES, AND DENSITIES IN VARIOUS ARRAYS FOR DETERMINING THE AVERAGE FLOW FIELD PROPERTIES AFTER STEADY-STATE HAS BEEN REACHED

---

```

DO 180 N=1,NBX
DO 110 MT=1,IP
XV(MT,N)=0.0
YV(MT,N)=0.0
ZV(MT,N)=0.0
TMP(MT,N)=0.
TRP(MT,N)=0.0
TTX=0.
TTY=0.
TTZ=0.
TTR=0.0
M=NB(MT,N)
IF (M.LT.1) GO TO 110
U=0.
V=0.
W=0.
DO 100 L=1,M
NA=NBM(MT,N)+L
J=LM(MT,NA)
PU=PAU(MT,J)
PV=PAV(MT,J)
PW=PAW(MT,J)
U=U+PU
V=V+PV
W=W+PW
TTR=TTR+ER(MT,J)
TTX=TTX+PU*PU
TTY=TTY+PV*PV
100 TTZ=TTZ+PW*PW
M=NB(MT,N)
XV(MT,N)=U/M
YV(MT,N)=V/M
ZV(MT,N)=W/M
TMP(MT,N)=(TTX+TTY+TTZ)/M
TRP(MT,N)=TTR/M

```

ACUM050  
 ACUM060  
 ACUM070  
 ACUM080  
 ACUM090  
 ACUM100  
 ACUM120  
 ACUM160  
 ACUM170  
 ACUM180  
 ACUM190  
 ACUM200  
 ACUM210  
 ACUM220  
 ACUM230  
 ACUM240  
 ACUM250  
 ACUM260  
 ACUM270  
 ACUM280  
 ACUM300  
 ACUM310  
 ACUM320  
 ACUM390  
 ACUM400  
 ACUM410  
 ACUM420

```

110 CONTINUE ACUM430
180 CONTINUE ACUM440
RETURN ACUM450
END ACUM460

```

```

SUBROUTINE AVRGE (PNB, DB, DBA, NB, NBT, XV, YV, ZV, XVA, YVA, ZVA, TMP, TRPA,
1TRP, TRPA, IP, I)
INTEGER*2 NB, NBT
DIMENSION PNB (1), DB (1, 1), DBA (1, 1), NB (1, 1), NBT (1, 1), TMP (1, 1)
DIMENSION TRPA (1, 1), XV (1, 1), XVA (1, 1), YV (1, 1), YVA (1, 1), ZV (1, 1)
DIMENSION ZVA (1, 1), TRP (1, 1), TRPA (1, 1)
COMMON /PORTH/NBX

```

AVG0050

AVG0060

```

-----
THE PURPOSE OF THIS SUBROUTINE IS TO COMPUTE THE AVERAGE FLOW
FIELD PROPERTIES.
-----

```

AVG0070

AVG0080

AVG0090

AVG0110

```

DO 110 N=1, NBX
DO 100 MT=1, IP
A=NBT (MT, N)
B=NB (MT, N)
C=A+B
NBT (MT, N)=C
IF (C.LT. 1.) GO TO 100
DBA (MT, N) = (DBA (MT, N) *A+DB (MT, N) *B) /C
XVA (MT, N) = (XVA (MT, N) *A+XV (MT, N) *B) /C
YVA (MT, N) = (YVA (MT, N) *A+YV (MT, N) *B) /C
ZVA (MT, N) = (ZVA (MT, N) *A+ZV (MT, N) *B) /C
TMPA (MT, N) = (TMPA (MT, N) *A+TMP (MT, N) *B) /C
TRPA (MT, N) = (TRPA (MT, N) *A+TRP (MT, N) *B) /C
100 CONTINUE
110 CONTINUE
RETURN
END

```

AVG0150

AVG0170

AVG0180

AVG0190

AVG0200

AVG0210

AVG0220

AVG0230

AVG0240

AVG0250

AVG0260

AVG0270

AVG0280

```

SUBROUTINE DRAG (AKN, NWEDG, ISTART, I2, I3, I4, I5, DELANG, JNT, BTA,
1C2, C3, DPA, PL, HTI, HTR, TB, XCB, CTI, CTR, CNI, CNR, ALPHA, SIGMA, BM, HTS,
2HTSI, NTS, JTL, UTT, VTS, LCOL, IP, ER, CHI, CNG, CHG, I, UTLI, UTTI, VTSI)
INTEGER*2 LCOL
INTEGER TIME, TST
REAL LAM, NU, NU, JAY, KAY
DIMENSION BTA (1), C2 (1), C3 (1), PL (1), HTI (1), CTI (3, 1), CTR (3, 1)
DIMENSION HTR (1), TB (1), XCB (1), ALPHA (3, 1), SIGMA (3, 1)
DIMENSION HTS (3, I2, I3), HTSI (3, I2, I3), NTS (3, I2, I3)
DIMENSION UTL (3, I2, I3), UTT (3, I2, I3), VTS (3, I2, I3), LCOL (1, 1)
DIMENSION UTLI (3, I2, I3), UTTI (3, I2, I3), VTSI (3, I2, I3)
DIMENSION CNI (3, 1), CNR (3, 1), DPA (1), JNT (1)
DIMENSION ER (1, 1), CNG (1), CHG (1), CHI (1)
COMMON /THIRD/PI
COMMON /FIFTH/BD, TIME, DTM, TI, ITS, ITP, TST
COMMON /SVNTH/LAM, NU, NU, MT, N, J, ICL, YCL, ZCL

```

DRG0040

DRG0050

DRG0060

DRG0140

DRG0150

DRG0160

```

-----
THE PURPOSE OF THIS SUBROUTINE IS TO ACCUMULATE THE DRAG AND HEAT
TRANSFER INCREMENTS ON THE BODY CONTRIBUTED BY EACH MOLECULE WHICH
COLLIDES WITH THE BODY. IN ADDITION, EACH MOLECULE WHICH COLLIDES
WITH THE BODY IS ASSIGNED AN APPROPRIATE NEW VELOCITY (OF REFLECTION)

```

170

180

190

200

210

WHICH IS USED TO CONTINUE ITS SPATIAL TRANSLATION (IN SUBROUTINE DRAG)

230

```

CALL NORMAL(JAY,KAY,RH)
JNT(MT)=JNT(MT)+1
TANG=180.*ATAN2(ZCL,-YCL)/PI
IWDG=TANG/DELANG+1
IF(IWDG.GT.NWEDG) IWDG=NWEDG DRG0310
D=(LAM*LAM+MU*MU+NU*NU) DRG0320
G=ER(MT,N)
H=G
DO 100 M=2,ND DRG0330
IF(XCL.LT.XCB(M)) GO TO 110 DRG0340
100 CONTINUE DRG0350
110 TBX=TB(M-1)+(TB(M)-TB(M-1))*(XCL-XCB(M-1))/(XCB(M)-XCB(M-1))
WI=(MU*JAY-NU*KAY) DRG0370
VID=MU*JAY+NU*KAY
UID=LAM
E=RAND(0) DRG0400
IF(E.LT.SIGMA(MT,M)) GO TO 115 DRG0410
VRD=-VID DRG0420
URD=UID DRG0430
WR=WI DRG0440
GO TO 125 DRG0450
115 V=4.*RAND(0) DRG0460
C2(MT)=C2(MT)+.544331*V*V*V*EXP(1.5-V*V) DRG0470
IF(C2(MT).LT.1.) GO TO 115 DRG0480
C2(MT)=C2(MT)-1. DRG0490
IF(NTS(MT,M,IWDG).NE.0) GO TO 117 DRG0500
ATR=ALPHA(MT,M)*TBX/SIGMA(MT,M) DRG0510
GO TO 118 DRG0520
117 ATR=ALPHA(MT,M)*TBX/SIGMA(MT,M)+(1.-ALPHA(MT,M)/SIGMA(MT,M))*HTS DRG0530
II(MT,M,IWDG)/NTS(MT,M,IWDG)/(3.+CHI(MT))
118 ABR=SQRT(ATR) DRG0550
V=V*ABR/BTA(MT) DRG0560
AA=RAND(0) DRG0570
A=SQRT(AA)
B=SQRT(1.-AA)
C=2.*PI*EAND(0)
VRD=V*A
URD=V*B*COS(C)
WR=V*B*SIN(C)
IF(CHI(MT).EQ.-1.) GO TO 125
122 X=9.*RAND(0)
IF(X.EQ.0.0) GO TO 122
XTEMP=1.0
IF(CHI(MT).NE.0.0) XTEMP=X**CHI(MT)
CNG(MT)=CNG(MT)+XTEMP*EXP(-X)
IF(CNG(MT).LT.CMG(MT)) GO TO 122
124 CONTINUE
CNG(MT)=CNG(MT)-CMG(MT)
IF(CNG(MT).GE.CMG(MT)) GO TO 124
ER(MT,N)=X*ATR
H=ER(MT,N)
125 UR=URD
LAM=UR

```

```

MU=JAY*VRD-KAY*WR
NU=KAY*VRD+JAY*WR
IF (TIME.LE.TST) RETURN
XMZ= (XCL-XSTART) *AKN
YMZ= RM*AKN
B= (URD*URD+VRD*VRD+WR*WR)
UTI=UID*UID+WI*WI
UYI=-WI*KAY
UYR=-WR*KAY
FL (MT) =FL (MT) +DPA (MT)
HTI (MT) =HTI (MT) +D+G
HTR (MT) =HTR (MT) -B-H
CTI (MT, 1) =CTI (MT, 1) +UID
CTI (MT, 2) =CTI (MT, 2) +UYI
CTI (MT, 3) =CTI (MT, 3) + (XMZ*UYI-YMZ*UID)
CNI (MT, 2) =CNI (MT, 2) +VID*JAY
CNI (MT, 3) =CNI (MT, 3) +XMZ*JAY*VID
CTR (MT, 1) =CTR (MT, 1) -URD
CTR (MT, 2) =CTR (MT, 2) -UYR
CTR (MT, 3) =CTR (MT, 3) - (XMZ*UYR-YMZ*URD)
CNR (MT, 2) =CNR (MT, 2) -VRD*JAY
CNR (MT, 3) =CNR (MT, 3) -XMZ*JAY*VRD
NTS (MT, M, IWDG) =NTS (MT, M, IWDG) +1
UTLI (MT, M, IWDG) =UTLI (MT, M, IWDG) +UID
UTL (MT, M, IWDG) =UTL (MT, M, IWDG) + (UID-URD)
UTTI (MT, M, IWDG) =UTTI (MT, M, IWDG) + WI
UTT (MT, M, IWDG) =UTT (MT, M, IWDG) + (WI-WR)
VTSI (MT, M, IWDG) =VTSI (MT, M, IWDG) -VID
VTS (MT, M, IWDG) =VTS (MT, M, IWDG) + (VRD-VID)
HTSI (MT, M, IWDG) =HTSI (MT, M, IWDG) +D+G
HTS (MT, M, IWDG) =HTS (MT, M, IWDG) +D-B+G-H
RETURN
END

```

DRG0730  
DRG0740  
DRG0750  
DRG0760  
  
DRG0790  
  
DRG0820  
DRG0830  
DRG0840  
DRG0860  
  
DRG0880  
DRG0890  
DRG0900  
DRG0920  
  
DRG0940  
  
DRG0960  
  
DRG0970  
  
DRG0980  
  
DRG1170  
DRG1180

```

SUBROUTINE INTERS (X, Y, Z, RMS, KS)
REAL LAM, MU, NU
COMMON /SYNTH/LAM, MU, NU, MT, N, J, XI, YI, ZI, TUSE
1 FORMAT (// ' SOMETHING IS WRONG IN INTERS' /6E15.6/4E15.6)
2 FORMAT (// ' TROUBLE IN INTERS - TTIME =', E15.6/6E15.6/4E15.6)
A=MU**2+NU**2
B= (YI*MU+ZI*NU) /A
C= (RMS-YI**2-ZI**2) /A
IF (C.LT.0.0) C=0.0
DISCR=B**2+C
IF (DISCR.GE.0.0) GO TO 10
WRITE (6, 1) XI, YI, ZI, MU, NU, RMS, A, B, C, DISCR
STOP
10 TTIME=SQRT (DISCR) -B
IF (C.EQ.0.0) TTIME=0.0
IF ((TTIME.LE.TUSE) .AND. (TTIME.GE.0.0)) GO TO 11
WRITE (6, 2) TTIME, XI, YI, ZI, MU, NU, RMS, A, B, C, DISCR
IF (TTIME.GT.TUSE) TTIME=TUSE
IF (TTIME.LT.0.0) TTIME=0.0
11 KS=1
TTIME=TTIME

```

ILE: GKBINT AUG82 A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

```

XI=XI+LAM*TYME
YI=YI+MU*TYME
ZI=ZI+NU*TYME
IF (ZI.GT.0.0) GO TO 20
ZI=-ZI
NU=-NU
20 CONTINUE
YI=.9999*YI
ZI=.9999*ZI
X=XI
Y=YI
Z=ZI
RETURN
END

```

```

SUBROUTINE NORMAL(JAY,KAY,RM)
REAL LAM,MU,NU,JAY,KAY
COMMON /SVNTH/LAM,MU,NU,HT,N,J,XCL,YCL,ZCL
JAY=-YCL/RM
KAY=-ZCL/RM
RETURN
END

```

NORM020  
NORM040  
  
NORM130  
NORM140

```

SUBROUTINE PRINT1(DT,COSANG,SINANG,RMA,RNU,DRP,PCP,HTF,PL,HTI,HTR,
1CTI,CTR,CNI,CNR)
DIMENSION DD(3),WD(2,5),PP(4,4),QQ(4,4),RR(4,4),SS(4,4),TT(4,4)
DIMENSION UU(4,4),P1(4,4),Q1(4,4),R1(4,4),PA(4),PB(4),PC(4)
DIMENSION PL(1),HTI(1),HTR(1),CTI(3,1),CTR(3,1),CNI(3,1),CNR(3,1)
DIMENSION RMA(1),RNU(1)
DATA WD/'X-PO','RCE ','Y-PO','RCE ','Z-MO','MENT','DRAG',' ','LPT10070
1IPT',' '/

```

-----  
THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE GROSS SURFACE  
COEFFICIENTS OF THE BODY.  
-----

FORMATS

```

1 FORMAT(//1X,50('*'),' GROSS SURFACE COEFFICIENTS ',50('*')/' MOLEC
1ULAR WEIGHT',12X,P8.3,3(19X,P8.3)/25X,
2 'INC. REP. TOT. INC. REP. TOT. INC.
3REP. TOT. INC. REP. TOT. ')
10 FORMAT(' NUMBER FLUX ',4(P9.3,18X))
12 FORMAT(1X,2A4,2X,' SHEAR ',4(3F8.3,3X))
14 FORMAT(11X,' PRESSURE ',4(3F8.3,3X))
16 FORMAT(11X,' TOTAL ',4(3F8.3,3X)/)
18 FORMAT(' HEAT TRANSFER',7X,4(3F8.3,3X)/)

```

PT10070  
PT10080  
PT10090  
PT10100  
PT10110  
PT10120  
PT10130  
PT10140  
PT10150  
PT10160  
PT10170  
PT10180  
  
\*\*\*\*\*  
PT10280  
PT10290  
PT10300  
PT10310

RMR=0.0

```

DO 50 MT=1,3
DD(MT)=RMA(MT)*RNU(MT)*DRF/DT
50 RMR=RMB+RMA(MT)*RNU(MT)
WRITE(6,1) (RMA(MT),MT=1,3),RMR
PF=PL(1)*PCF/DT
QP=PL(2)*PCF/DT
RP=PL(3)*PCF/DT
SF=PF+QP+RP
WRITE(6,10) PF,QP,RP,SF
DO 200 I=1,3
PP(4,I)=0.0
QQ(4,I)=0.0
RR(4,I)=0.0
SS(4,I)=0.0
TT(4,I)=0.0
UU(4,I)=0.0
P1(4,I)=0.0
Q1(4,I)=0.0
R1(4,I)=0.0
DO 150 MT=1,3
PP(MT,I)=CTI(MT,I)*DD(MT)/RMR
QQ(MT,I)=CTR(MT,I)*DD(MT)/RMR
SS(MT,I)=CNI(MT,I)*DD(MT)/RMR
TT(MT,I)=CHR(MT,I)*DD(MT)/RMR
P1(MT,I)=PP(MT,I)+SS(MT,I)
Q1(MT,I)=QQ(MT,I)+TT(MT,I)
RR(MT,I)=PP(MT,I)+QQ(MT,I)
UU(MT,I)=SS(MT,I)+TT(MT,I)
R1(MT,I)=P1(MT,I)+Q1(MT,I)
PP(4,I)=PP(4,I)+PP(MT,I)
QQ(4,I)=QQ(4,I)+QQ(MT,I)
RR(4,I)=RR(4,I)+RR(MT,I)
SS(4,I)=SS(4,I)+SS(MT,I)
TT(4,I)=TT(4,I)+TT(MT,I)
UU(4,I)=UU(4,I)+UU(MT,I)
P1(4,I)=P1(4,I)+P1(MT,I)
Q1(4,I)=Q1(4,I)+Q1(MT,I)
R1(4,I)=R1(4,I)+R1(MT,I)
150 CONTINUE
WRITE(6,12) (WD(J,I),J=1,2),(PP(K,I),QQ(K,I),RR(K,I),K=1,4)
WRITE(6,14) (SS(K,I),TT(K,I),UU(K,I),K=1,4)
WRITE(6,16) (P1(K,I),Q1(K,I),R1(K,I),K=1,4)
200 CONTINUE
AA=COSANG
BB=SINANG
DO 300 I=4,5
DO 250 K=1,4
PP(K,4)=AA*PP(K,1)+BB*PP(K,2)
QQ(K,4)=AA*QQ(K,1)+BB*QQ(K,2)
RR(K,4)=AA*RR(K,1)+BB*RR(K,2)
SS(K,4)=AA*SS(K,1)+BB*SS(K,2)
TT(K,4)=AA*TT(K,1)+BB*TT(K,2)
UU(K,4)=AA*UU(K,1)+BB*UU(K,2)
P1(K,4)=AA*P1(K,1)+BB*P1(K,2)
Q1(K,4)=AA*Q1(K,1)+BB*Q1(K,2)

```

PT10400

PT10460  
PT10470  
PT10480  
PT10490

PT10630  
PT10640  
PT10650  
PT10660

PT10680  
PT10690  
PT10700  
PT10710  
PT10720  
PT10730  
PT10740  
PT10750

```

250 R1(K,4)=AA*R1(K,1)+BB*R1(K,2)
WRITE(6,12) (WD(J,I),J=1,2), (PP(K,4),QQ(K,4),RR(K,4),K=1,4)
WRITE(6,14) (SS(K,4),TT(K,4),UU(K,4),K=1,4)
WRITE(6,16) (P1(K,4),Q1(K,4),R1(K,4),K=1,4)
AA=-SINANG
BB=COSANG
300 CONTINUE
HD=HTF/DT
PA(4)=0.0
PB(4)=0.0
PC(4)=0.0
DO 400 MT=1,3
PA(MT)=HTI(MT)*RMA(MT)*RNU(MT)*HD/RMR
PB(MT)=HTR(MT)*RMA(MT)*RNU(MT)*HD/RMR
PC(MT)=PA(MT)+PB(MT)
PA(4)=PA(4)+PA(MT)
PB(4)=PB(4)+PB(MT)
PC(4)=PC(4)+PC(MT)
400 CONTINUE
WRITE(6,18) (PA(I),PB(I),PC(I),I=1,4)
RETURN
END

```

PT10760

PT10800

PT10810

PT10820

PT10830

PT10950

PT10960

```

SUBROUTINE PRINT2(AKN,XSTART,DT,RNU,RMA,DRP,PCF,HTF,UTLI,UTTI,VTSI
1,HTSI,DELANG,NWEDG,XS,XCB,YCB,HTS,NTS,UTL,UTT,VTS,I2,I3,IP)
DIMENSION RMA(1),RNU(1),XS(1),XCB(1),YCB(1)
DIMENSION HTS(3,I2,I3),NTS(3,I2,I3),UTL(3,I2,I3)
DIMENSION UTT(3,I2,I3),VTS(3,I2,I3),UTLI(3,I2,I3),UTTI(3,I2,I3)
DIMENSION VTSI(3,I2,I3),HTSI(3,I2,I3)
COMMON /FIFTH/ND

```

PT20060

PT20070

THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE DISTRIBUTION ON SURFACE OF THE SURFACE COEFFICIENTS

PT20080

PT20090

PT20100

PT20110

PT20120

PT20130

PT20140

PT20150

PT20160

FORMATS

```

8 FORMAT(/1X,45('*'), ' DISTRIBUTION ON SURFACE ',45('*')/1X, 'INC.
1 TOT. INC. TOT. INC. TOT. '/11X, 'SEGMENT GEOMETRY',
214X, 'MOL. MOLE SAMP NUM. SKIN SKIN PRES- PRES-
3 HEAT HEAT'/' NO. CENTER DELI CENTER DELANG',4X, 'WGHT. P
4RACT.',10X,
5 'FLUX PRCTN PRCTN SURE SURE TRNSF TRNSP')
10 FORMAT(1X,I3,F8.3,F7.3,F9.3,F8.3,1X,2F8.4,I6,7F8.4)
11 FORMAT(37X,F8.4, ' 1.0000',I6,7F8.4)

```

PT20230

PT20240

PT20250

PT20260

```

RMR=0.0
DO 50 MT=1,IP
50 RMR=RMR+RMA(MT)*RNU(MT)

```



```

WRITE(6,8)
I=0
DO 110 N=2,ND
DTY=DT*YCB(N)/180.
P=IS(N)
Q=2.*(ICB(N)-I*START)*AKH-XS(N)
J=0
R=-.5*DELANG
PHLT=PCP/(DTY*DELANG)
QMLT=DRP/(DTY*DELANG)
SMLT=HTP/(DTY*DELANG)
DO 100 K=1,NWEDG
R=R+DELANG
I=I+1
J=J+1
M3=0
P3=0.0
Q3=0.0
Q4=0.0
E3=0.0
R4=0.0
S3=0.0
S4=0.0
DO 90 MT=1,IP
M1=NTS(MT,N,J)
M3=M3+M1
P1=NTS(MT,N,J)*PHLT*RNU(MT)
P3=P3+P1
Q1=SQRT(UTLI(MT,N,J)**2+UTTI(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
Q2=SQRT(UTL(MT,N,J)**2+UTT(MT,N,J)**2)*RNU(MT)*RMA(MT)*QMLT/RMR
Q3=Q3+Q1
Q4=Q4+Q2
R1=VTSI(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
R2=VTS(MT,N,J)*RNU(MT)*RMA(MT)*QMLT/RMR
S1=HTSI(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR
S2=HTS(MT,N,J)*RNU(MT)*RMA(MT)*SMLT/RMR
R3=R3+R1
R4=R4+R2
S3=S3+S1
S4=S4+S2
90 WRITE(6,10)I,P,Q,R,DELANG,RMA(MT),RNU(MT),M1,P1,Q1,Q2,R1,R2,S1,S2
WRITE(6,11)RMR,M3,P3,Q3,Q4,R3,R4,S3,S4
100 CONTINUE
110 CONTINUE
RETURN
END

```

PT20280  
PT20320  
PT20330  
PT20340  
PT20350  
PT20380

PT20490  
PT20500

PT20660  
PT20670  
PT20680  
PT20690

```

SUBROUTINE PRINT4(MSP,CHI,RNU,I,TRP,FDN,WTH,DB,NS,TMP,YV,
1YV,ZV,KS,NB,IC,YC,ZC)
INTEGER*2 NB,NS
DIMENSION FDN(1),RNU(1),CHI(1),WTH(1),TMP(I,1),TRP(I,1)
DIMENSION DB(I,1),NB(I,1),YV(I,1),ZV(I,1),DBT(3),NS(I,1)
DIMENSION IC(1),YC(1),ZC(1)
COMMON /SECND/BW,BH,BTP,RMR,BFP
COMMON /FORTH/NBX,RA,XR

```

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FILE: GKBINT AUG82 1

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

-----  
THE PURPOSE OF THIS SUBROUTINE IS TO PRINT OUT THE INSTANTANEOUS  
FLOW-FIELD PROPERTIES.  
-----

FORMATS

1 FORMAT (//1X,45('\*'), ' INSTANTANEOUS FLOW FIELD INFORMATION ',45('\*'  
1')) PT50060  
2 FORMAT (/2X, 'HORIZONTAL NUMBER=', I3,3X, 'VERTICAL NUMBER=', I3,3X, 'HO  
RIZONTAL POSITION=', F8.5,3X, ' RADIAL POSITION =', F8.5/2X, ' BOX\*  
ZANGLE SAMP DENSITY EACH NO X VEL. Y VEL. Z VEL. T(KIN) T(RO  
3T) TEMP.', 14X, 'MOLE FRACTIONS') PT50070  
3 FORMAT (//1X,46('\*'), ' ACCUMULATED FLOW FIELD INFORMATION ',46('\*'  
1)) PT50080  
4 FORMAT (1X, I4, E11.3, I6, 8F8.3, 3X, 3E11.3) PT50090  
PT50100  
PT50110  
PT50120  
PT50130  
PT50140  
PT50150  
PT50160  
PT50170

\*\*\*\*\* PT50220  
\*\*\*\*\* PT50230  
\*\*\*\*\* PT50240  
\*\*\*\*\* PT50250

```

IF (KS.EQ.0) WRITE (6,1)
IF (KS.NE.0) WRITE (6,3)
DO 40 MT=1,3
40 DBT(MT)=0.0
   FDA=0.
   CHT=0.
   DO 50 MT=1,MSP
     CHT=CHT+CHI(MT)*RNU(MT)
50 FDA=FDA+PDN(MT)*WTH(MT)
   YCT=0.0
   XCT=0.0
   DO 110 N=1,NBX
     IF ((XC(N).EQ.XCT).AND.(YC(N).EQ.YCT)) GO TO 52
     ICT=XC(N)
     YCT=YC(N)
     XXT=XCT/XR
     YYT=YCT/RM
     IXC=XCT/BW + 1
     IYC=YCT/BH + 1
     WRITE (6,2) IXC,IYC,XXT,YYT
52 ZCT=ZC(N)
   NSAMP=0
   DBA=0.
   XVM=0.
   YVM=0.
   ZVM=0.
   TRPM=0.
   E=0.
   F=0.
   DO 100 MT=1,MSP
     NSAMP=NSAMP+NS(MT,N)

```

PT50290

FILE: GKBINT AUG82 A

PRINCETON UNIVERSITY TIME-SHARING SYSTEM

```

XVM=XVM+XV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
YVM=YVM+YV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
ZVM=ZVM+ZV(MT,N)*RNU(MT)*WTM(MT)*NB(MT,N)
CBA=DBA+DB(MT,N)*WTM(MT)
TMPM=TMPM+WTM(MT)*RNU(MT)*TMP(MT,N)*NB(MT,N)
TRPM=TRPM+RNU(MT)*NB(MT,N)*TRP(MT,N)
E=E+WTM(MT)*RNU(MT)*NB(MT,N)
100 F=F+RNU(MT)*NB(MT,N)
    DBA=DBA/PDA
    IF(E.EQ.0.0) GO TO 55
    XVM=XVM/E
    YVM=YVM/E
    ZVM=ZVM/E
    VS=XVM**2+YVM**2+ZVM**2
    TMPM=TMPM/E-VS
    TRPM=TRPM/F
55 CONTINUE
    TTM=(TMPM+TRPM)/(2.5+CHT)
    TMPM=TMPM/1.5
    IF(CHT.NE.-1.) TRPM=TRPM/(1.+CHT)
    AMS=SQRT(VS)
    IF(TTM.GT.0.) AMS=SQRT((5.+2.*CHT)*VS/(TTM*(3.5+CHT)))
    CCZ=COS(ZCT/57.29578)
    SCZ=SQRT(1.-CCZ**2)
    RVM=ZVM*SCZ-YVM*CCZ
    TVM=YVM*SCZ+ZVM*CCZ
    DO 60 MT=1,NSP
    DBT(MT)=RNU(MT)*NB(MT,N)
    IF(F.NE.0.) DBT(MT)=DBT(MT)/F
60 CONTINUE
    WRITE(6,4) N,ZCT,NSAMP,DBA,AMS,XVM,YVM,TVM,TMPM,TRPM,TTM,(DBT(J),
1J=1,3)
110 CONTINUE
    RETURN
    END

```

PT50470  
PT50480  
PT50490

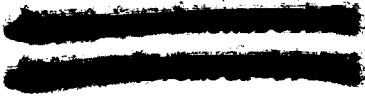
```

ENTRY
CONTROL NAME='INTE','ENAL',TITLE=' PAR','ABOL','A AT',' 95K','M M','ON. ',
DEBUG=.F.,.F.,.T.,NEW=.F.,SAVE=.T.,ICOPY=0,REDO=.T. &END
TIMES DTM=.005,ITS=6,ITP=6,TST=2,TLIM=12 &END
FLOREF LLM=2000,ENH=5000,MNB=150,NSP=1,NET=0,U=7485.9,ANGLE=28.,RNU=1.,2*0.,
RMA=28.94,0.,0.,TF=195.51,DENF=2.52E+19 &END
MOLEC TRP=1000,DIR=3.5E-19,ETA=.104,PHI=0.0,CHI=-1.,ACR=.001 &END
SHAPES BODY=0.0,1000.,.00235 &END
SHAPES BODY=.0025,555.,0.0,2*1.0 &END
SHAPES BODY=.0050,345.,0.0,2*1.0 &END
SHAPES BODY=.0100,300.,0.0,2*1.0 &END
SHAPES BODY=.0200,300.,0.0,2*1.0 &END
SHAPES BODY=.0300,300.,0.0,2*1.0 &END
SHAPES BODY=.0400,300.,0.0,2*1.0 &END
SHAPES BODY=.0500,300.,0.0,2*1.0 &END
SHAPES BODY=.0600,300.,0.0,2*1.0 &END
SHAPES BODY=.0700,300.,0.0,2*1.0 &END
SHAPES BODY=.0800,300.,0.0,2*1.0 &END
SHAPES BODY=.0870,300.,1.0,2*1.0 &END
GEOM BWEDG=2,NW=20,NH=3, &END

```

```
;INCOPL FLUXIN=2.1429,PCOL=1.0,RMP=0.3,JV=22,KMX=3 &END
;INOUT VARG=0.,1.,2.,3.,4.,5.,6.,7.,8.,9.,10.,11.,12.,13.,14.,15.,16.,17.,18.,
19.,20.,21.,CURV=0.0,.070,.170,.282,.369,.459,.537,.599,.656,.710,.750,.785,
.815,.845,.872,.900,.922,.951,.975,.988,.996,1.00, &END
;INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
11.,13.,15.,17.,19.,CURV=4*0.,.003,.013,.036,.084,.149,.250,.406,.611,.762,
.871,.932,.962,.984,.995,.999,3*1.0, &END
;INOUT VARG=-20.,-19.,-17.,-15.,-13.,-11.,-9.,-7.,-5.,-3.,-1.,1.,3.,5.,7.,9.,
11.,13.,15.,17.,19.,CURV=4*0.0,.003,.013,.036,.084,.149,.250,.406,.611,.762,
.871,.932,.962,.984,.995,.999,3*1.0, &END
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16. Abstract <p>This report presents the results of analysis performed on the external flow around the shuttle orbiter nose region at the Shuttle Upper Atmosphere Mass Spectrometer (SUMS) inlet orifice. The purpose of the analysis is to quantitatively characterize the flow conditions to facilitate SUMS flight data reduction and subsequent determination of orbiter aerodynamic force coefficients in the hypersonic rarefied flow regime. Experimental determination of aerodynamic force coefficients requires accurate simultaneous measurement of forces (or acceleration) and dynamic pressure along with independent knowledge of density and velocity. SUMS provides independent measurement of dynamic pressure; however, it does so indirectly and requires knowledge of the relationship between measured orifice conditions and the dynamic pressure which can only be determined on the basis of molecule or theory for a winged configuration.</p> <p>Monte Carlo direct simulation computer codes were developed for both the flow field solution at the orifice and for the internal orifice flow. These codes were used to study issues associated with geometric modeling of the orbiter nose geometry and the modeling of intermolecular collisions including rotational energy exchange and a preliminary analysis of vibrational excitation and dissociation effects. Data obtained from preliminary simulation runs are presented. Conclusions from these results are included along with recommendations for a SUMS flight data reduction procedure and future work to generate the data required for that procedure.</p>					
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