# Application of Stiffened Cylinder Analysis to ATP Interior Noise Studies 

E.G.Wilby and J.F.Wilby<br>Bolt Beranek and Newman Inc.<br>Canoga Park, CA 01303

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Space Administration

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### 1.0 INTRODUCTION

The transmission of propeller noise into an airplane fuselage has a direct influence on the design of General Aviation and Advanced Turboprop (ATP) aircraft. For this reason NASA has undertaken several analytical or experimental studies which consider different aspects of the overall problem. In one analytical study [1-4] a method was developed to predict sound levels in a stiffened cylindrical fuselage when the exterior of the fuselage was exposed to a propeller noise field. The analytical model was compared with experimental results obtained from a small test cylinder. Under a separate investigation [5], noise transmission measurements were made in the laboratory using the fuselage of a Fairchild Metro II airplane. This experimental investigation provided validation data for ATP noise control studies.

The present study has two objectives. The first is to adapt the stiffened cylinder analytical model [1-4] to the test conditions associated with the Metro II experiment [5] and to compare predicted and measured results for the noise reduction provided by the fuselage structure and treatment. The second objective is to extend the analytical model to include turbulent boundary layer excitation so that comparative noise reduction predictions can be performed for different types of excitation.

The first part of this report (Section 2) describes the analytical model [1] and the changes made to it in order to accomplish the objectives of the study. Section 3 presents an outline of the Metro II test with emphasis being placed on the information relevant to the current study. Analytical representations of the Metro II test structure and the test excitation field are given in Sections 4 and 5. Then the predicted and measured noise reductions for the test fuselage are compared in Section 6. The effect of type of excitation on the noise reduction is discussed in Section 7 and final conclusions are presented in Section 8.

### 2.0 ANALYTICAL MODEL

### 2.1 Outline

An analytical model for aircraft interior noise prediction was developed under a program sponsored by NASA Langley Research Center, (NASA Contract NAS1-15782), and is described in [1-4]. The model calculates space-average sound pressure levels inside a cylindrical fuselage with a floor and sidewall treatment, when the exterior pressure field is generated by a propeller. In addition the model derives the noise reduction for an exterior reverberant (diffuse) acoustic field.

For this report, the analytical model has been extended to calculate the noise reduction associated with turbulent boundary layer excitation. In addition, the sidewall treatment has been modified so that it may consist of 1 to 4 trim elements. Modifications have also been made to the noise transmitted from the cabin floor to the interior, to allow variation of the floor treatment transmission loss.

### 2.2 Tone Transmission

The band-limited, space-average mean-square pressure in the interior of the fuselage, for harmonic $H$ at frequency $\omega_{H}$, is given by Equations (3), (8) and (10) in [1]. If $^{2} \tau_{\mathrm{ML}}^{\mathrm{H}}$ is written in terms of its component parts, $\tau_{M L}^{H}=\left(\frac{2 \rho I c_{O I}}{\omega_{H} m}\right)^{2}$, then

$$
\begin{align*}
&\left\langle p_{i}^{2}\right\rangle \\
& H, t=\frac{\rho_{I}^{2} c_{O I}^{4}}{V^{2}} \cdot \omega_{H}^{4} \cdot \sum_{n} \frac{\varepsilon_{n}}{\omega_{n}^{4}} \cdot \sum_{r} \frac{A^{2} \bar{f}^{2}(n, r)}{M^{2} \omega_{r}^{4}} \cdot \Psi_{G}(r, H) \\
& x \frac{1}{\left[\left(1-\frac{\omega_{H}^{2}}{\omega_{n}^{2}}\right)^{2}+\eta_{n}^{2}\right]\left[\left(1-\frac{\omega_{H}^{2}}{\omega_{r}^{2}}\right)^{2}+\left(\eta_{r}^{\prime}+\eta_{r}^{\prime \prime}\right)^{2}\right]} \tag{1}
\end{align*}
$$

where $A=$ interior (cylinder + floor) surface area associated with interior coupling factor $\bar{f}^{\prime}(n, r)$
$=L\left[L_{p}+2 a\left(\pi-\theta_{0}\right)\right]$,
and $\rho_{I}$ and $c_{o I}$ are the interior density and speed of sound, respectively. Other symbols are defined in Appendix A.

The interior coupling factor, $\bar{f}^{\prime}(n, r)$, between a fuselage structural mode and a cabin acoustic mode is defined in Equation (52) Ref.[1] for a bare cabin floor. The $\bar{f}^{\prime}(n, r)$ includes a sidewall trim transmission coefficient $\tau_{t}$, and, in a similar way, a transmission coefficient, $\tau_{F}$, could be included for the treatment on the floor.* If the floor treatment transmission loss is known and is defined as

Floor treatment transmission loss $=-10 \log \left(\tau_{F}\right) d B$

Equation (52) Ref.[1] becomes

$$
\begin{align*}
& \bar{f}^{\prime}(n, r)=\bar{f}^{\prime}(q i, r)=\frac{I}{\bar{L}} \int_{0}^{L} \cos \frac{q z}{L} \sin \frac{M z}{L} d z  \tag{3}\\
& x \frac{1}{L_{p}+2 a\left(\pi-\theta_{0}\right)}\left[\sqrt{\tau_{t}^{H}} \int_{\theta_{0}}^{2 \pi-\theta_{0}} a \psi_{S}^{r}(\theta) \phi_{i}(\theta) d \theta-\sqrt{T_{F}^{H}} \int_{-L_{p} / 2}^{L_{p} / 2} \psi_{p}^{r}(x) \phi_{i}(x) d x\right]
\end{align*}
$$

The sidewall trim is assumed to cover all the curved surface of the fuselage above the floor, and the trim transmission coefficient $\tau_{t}$ is defined in Appendix $A[1]$.

The analytical model divides the frequency range into "low" and "high" regimes. At low frequencies acoustic coupling between structural and cavity modes is calculated on a mode-by-mode basis. As frequency increases the number of acoustic modes in a given frequency band becomes very large and the coupling between structural and acoustic modes is calculated on a band-average basis (as is done in the statistical energy analysis method).

[^0]An estimate of the change-over frequency from low freqency to high frequency regimes is based on the volume of the cavity. The change-over frequency is the center frequency of the one-third octave band in which the following empirical values lies

$$
\mathrm{f}_{\mathrm{V}}=\frac{4.1 \mathrm{c}_{\mathrm{oI}}}{\sqrt[3]{V}} \mathrm{~Hz}
$$

where $c_{O I}$ is the speed of sound ( $\mathrm{m} / \mathrm{s}$ ) in the cavity of volume $\mathrm{V}\left(\mathrm{m}^{3}\right)$.

For high frequencies, the expected value of the space average mean-square interior pressure for harmonic $H$ is given by equation (16) in [1]. Without using the approximation for $g_{r}$ developed in Section 3.2 of [1], this becomes
$E\left[\left\langle p_{i}^{2}\right\rangle{ }_{s, t}^{H}\right]=\frac{\rho_{I}^{2} c_{O I}}{2 \pi V} \cdot \frac{\omega_{H}^{3}}{\bar{n}_{n}} \sum_{r \in \Delta \omega} \frac{\Psi_{G}(r, H)}{M_{r}^{2} \omega_{r}^{4}\left[\left(1-\frac{\omega_{H}^{2}}{\omega_{r}^{2}}\right)^{2}+\eta_{r}^{\prime 2}\right]} \cdot\left[A^{2} j_{r}^{2}\left(\omega_{H}^{r e v}\right)\right]$ (4) (4)
The bandwidth $\Delta \omega$, containing $\omega_{H}$, should be wide enough to ensure smoothness and may be selected to include a sufficiently large number of modes for computational accuracy. The notation resw implies that $\omega_{r}$ also lies within the band $\Delta \omega$.

It is now assumed that the modal displacement of the surface which forms the inner boundary of the sidewall trim or floor covering can be given approximately by the function $\sqrt{\tau(\bar{x})} \psi^{r}(\bar{x})$ where $\tau(\bar{x})$ is a general transmission coefficient

Then the interior joint acceptance (i.e., the joint acceptance coupling the fuselage structure to the interior sound field) can be defined for the interior cavity above the floor as

$$
\begin{aligned}
j_{r}^{2}(\omega)= & \frac{I}{A^{2}} \iint_{\text {interior }} \sqrt{\tau(\bar{x}) \tau\left(\bar{x}^{\prime}\right)} C\left(\bar{x} \mid \bar{x}^{\prime} ; \omega\right) \psi^{r}(\bar{x}) \psi^{r} d \bar{x} d \bar{x}^{\prime} \\
\simeq & \frac{1}{A^{2}}\left\{\iint_{\text {fliselage above }} \int_{\text {floor }} \tau_{\mathrm{t}} C\left(\bar{x} \mid \bar{x}^{\prime} ; w\right) \psi^{r}(\bar{x}) \psi^{r}\left(\bar{x}^{\prime}\right) d \bar{x} \bar{d}^{\prime}\right. \\
& \left.+\iint_{f l \infty r} \tau_{F} C\left(\bar{x} \mid \bar{x}^{\prime} ; \omega\right) \psi^{r}(\bar{x}) \psi^{r}\left(\bar{x}^{\prime}\right) d \bar{x} d^{\prime} \bar{x}^{\prime}\right\}
\end{aligned}
$$

This approximation assumes that the cross terms between the floor and the fuselage above the floor can be neglected. For a reverberant field, the correlation function is of the form

$$
\begin{equation*}
C\left(\bar{x} \mid \bar{x}^{\prime} ; \omega\right)=\frac{\sin k\left(\bar{x}-\bar{x}^{\prime}\right)}{\left.k(\bar{x}-\bar{x})^{\prime}\right)} \tag{5}
\end{equation*}
$$

The neglect of the cross terms is therefore reasonable at high frequencies, but is only approximate in the mid-frequency range.

Calculation of the joint acceptance of the fuselage above the floor is considerably more complex than that for the whole fuselage, because the integration from floor-to-floor does not cover full modal wavelengths. However some simplifying approximations are possible. In Sections 5.4 and 5.5 of Ref.[6], it was shown that for reverberant field excitation, at high frequencies, $A\left\langle j_{r}^{2}(\underset{\omega}{\omega})^{r}\right\rangle$ is essentially independent of area. Thus

$$
\begin{equation*}
\left[\bar{A}_{\underset{r}{2}(\omega)}^{\substack{r e v \\ \text { whole } \\ \text { cylinder }}} \underset{ }{ } \simeq A_{t} j_{\substack{\text { cylinder } \\ \text { above floor }}}\right. \tag{6}
\end{equation*}
$$

where $\bar{A}=$ fuselage surface area

$$
\begin{equation*}
=\text { L. } 2 \pi а \tag{7}
\end{equation*}
$$

$A_{t}=$ fuselage surface area above floor

$$
=\text { Transmitting area of fuselage with trim }
$$

$$
\begin{equation*}
=L .2 a\left(\pi-\theta_{0}\right) \tag{8}
\end{equation*}
$$

This gives

$$
\left[A^{2} j_{r}^{2}\left(\omega_{H}^{r e v}\right)\right]_{\text {interior }}=\tau_{t} A_{t} \bar{A}\left[j_{r}^{\text {rev }} j_{\substack{2  \tag{9}\\
\text { fuselage } \\
\text { interior }}}+\tau_{F} A_{f}^{2}\left[\begin{array}{c}
\text { rev } \\
j_{r}^{2}\left(\omega_{H}\right)
\end{array}\right]_{\text {floor }}\right.
$$

The reverberant field joint acceptance for the fuselage is defined in Equations 62-73, Ref.[1], and for the floor is defined ir Section 2.7 of this report. The value of the speed of sound, $c_{0 I}$, for the interior volume should be used.

### 2.3 Noise Reduction due to Reverberant Field Excitation

For a reverberant field, the exterior mean square pressure $\left\langle p_{e}^{2}\right\rangle_{s, t}$ is related to the mean square blocked pressure $\left\langle p_{b 1}^{2}\right\rangle_{s, t}$ incident on the fuselage by

$$
\left\langle p_{e}^{2}\right\rangle_{s, t}=\left\langle p_{b 1}^{2}\right\rangle_{s, t} / 2
$$

The noise reduction of the fuselage for the one-third octave band at center frequency $\omega$ is given by Equation 18 in Ref.[1]. Expanding the expression for $\tau_{M L}$ gives

$$
\begin{align*}
\frac{\left\langle p_{i}^{2}\right\rangle s, t}{\left\langle p_{e}^{2}\right\rangle s, t} & =\frac{2 \rho_{I}^{2} c_{o I}^{4} \bar{A}^{2}}{c_{\omega} V^{2}} \cdot \sum_{n} \varepsilon_{n} \sum_{r} \frac{\left[j_{r}^{2 r e v}(\omega)\right]_{e x t} A^{2} \bar{f}^{2}(n, r)}{M_{r}^{2} D_{n r}} \\
& x\left\{\left(\frac{c_{r}-c_{n}}{4}\right) \ln _{n}+\left(\frac{2 c_{n}\left(b_{r}-b_{n}\right)-b_{n}\left(c_{r}-c_{n}\right)}{4 n_{n} \omega_{n}^{2}}\right) \arctan n\right.  \tag{10}\\
& \left.+\left(\frac{c_{n}-c_{r}}{4}\right) \ln n_{r}+\left(\frac{2 c_{r}\left(b_{n}-b_{r}\right)-b_{r}\left(c_{n}-c_{r}\right)}{4\left(\eta_{r}^{\prime}+n_{r}^{\prime \prime}\right) \omega_{r}^{2}}\right) \arctan _{r}\right\}
\end{align*}
$$

where $\bar{A}$ is the fuselage surface area and $A$ is the interior (cylinder + floor) coupling area.

For $n$ or $r=j$,

$$
\begin{aligned}
& \ln j=\ln \left\{\frac{\left|\left(1+c_{\omega} / 2\right)^{4} \omega^{4}+b_{j}\left(1+c_{\omega} / 2\right)^{2} \omega^{2}+c_{j}\right|}{\left|\left(1-c_{\omega} / 2\right)^{4} \omega^{4}+b_{j}\left(1-c_{\omega} / 2\right)^{2} \omega^{2}+c_{j}\right|}\right\}, \\
& \arctan j=\tan ^{-1}\left[\frac{\left(2+c_{\omega}\right)^{2} \omega^{2}-4 \omega_{j}^{2}}{4 n_{j} \omega_{j}^{2}}\right]-\tan ^{-1}\left[\frac{\left(2-c_{\omega}\right)^{2} \omega^{2}-4 \omega_{j}^{2}}{4 n_{j} \omega_{j}^{2}}\right]
\end{aligned}
$$

where when $j=r, \eta_{j}$ above $=\eta_{r}^{\prime}+\eta_{r}^{\prime}$, and when $j=n_{i}, \eta_{j}=\eta_{n}$. Also,

$$
\begin{aligned}
& D_{n r}=\left(c_{r}-c_{n}\right)^{2}+\left(b_{n}-b_{r}\right)\left(b_{n} c_{r}-b_{r} c_{n}\right), \\
& b_{n}=-2 \omega_{n}^{2} ; \quad b_{r}=-2 \omega_{r}^{2}, \\
& c_{n}=\omega_{n}^{4}\left(1+\eta_{n}^{2}\right) \quad ; \quad c_{r}=\omega_{r}^{4}\left[1+\left(\eta_{r}^{\prime}+\eta_{r}^{\prime \prime}\right)^{2}\right]
\end{aligned}
$$

The expression for $\bar{f}^{\prime}(n, r)$, including the floor transmission loss is given in Equation (3), and the expression for the exterior reverberant field joint acceptance $\left[j_{r}^{2}(\omega)\right]$ ext is given in Equations 62-73, Ref.[1] using the speed of sound, coE, for the exterior pressure field.

For high frequencies, when the acoustic modal density is high, the model was developed along the lines of $[4,7,8]$. The power absorbed on the interior wall of the fuselage is given by Equation 22 of Reference [1], assuming that the response is resonant acoustic. Structural modes resonant below the frequency band ( $r<\Delta \omega$ ) and resonant in the frequency band ( $r \varepsilon \Delta \omega$ ) are included. The noise reduction at the one-third octave band frequency $\omega$ is given by

$$
\begin{align*}
& \frac{\left\langle p_{e}^{2}\right\rangle_{S, t}}{\left\langle p_{i}\right\rangle_{S, t}}=\left\{\frac{\alpha S}{4 \rho_{I}{ }^{c} O I}+\frac{\rho_{I} c_{O E}}{\rho_{E} c_{O I}} \cdot \pi n_{r}\left\langle\frac{A^{2}}{M_{r}} \cdot j_{r}^{2}(\omega)\right\rangle_{i n t}^{r e v}\right\rangle_{r \Delta \omega} \\
& \left.+\frac{\rho_{E}}{2 \pi c_{o E}} \sum_{r<\Delta \omega}\left(\frac{A^{2}}{M_{r}} \cdot j_{r}^{2 r(\omega)}\right)_{\text {int }} \cdot\left(\frac{\bar{A}^{2}}{M_{r}} \cdot j_{r}^{2(\omega)}\right)_{\text {ext }}^{\text {rev }}\right\} \\
& \div\left\{\frac{\rho_{I} \omega n_{r}}{2 c_{o I} \bar{n}_{r}}\left\langle\left(\frac{\bar{A}^{2}}{M_{r}} \cdot j_{r}^{2}(\omega)\right)_{\underset{r \varepsilon \Delta \omega}{ }(\omega)}^{\operatorname{ext}_{r}} \cdot\left(\frac{A^{2}}{M_{r}} \cdot j_{r}^{2(\omega)}\right)_{i n t}^{r e v}\right\rangle\right. \\
& \left.+\frac{\rho I}{2 \pi c_{o I}} \sum_{r<\Delta \omega}\left(\frac{A^{2}}{\bar{M}_{r}} \cdot j_{r}^{2 r(\omega)}\right)_{\text {int }}^{r e v}\left(\frac{\bar{A}^{2}}{\bar{M}_{r}} \cdot j_{r}^{2(\omega)}\right)_{e x t}^{2(\omega)}\right\} \tag{11}
\end{align*}
$$

where $\left(A^{2} j_{r}^{r e v}(\omega)\right)_{\text {int }}$ is given by Equation (9) and $\left(j_{r}^{2}(\underset{\sim}{e v})\right)_{\text {ext }}$ is given by Equations 62-73 in Reference [1] using the exterior speed of sound $c_{O E}$. The exterior air density is $\rho_{E}$.
The band average absorption coefficient $\alpha$ is associated with the absorbing surface area $S$.

### 2.4 Noise Reduction due to Boundary Layer or Progressive Wave Excitation Fields

The development of the noise reduction due to an exterior field such as boundary layer or progressive wave excitation is very similar to that for the reverberant field excitation. In this
case, the blocked pressure $\left\langle p_{b l}^{2}\right\rangle$ is used, rather than the exterior pressure $\left\langle p_{e}^{2}\right\rangle$, and the exterior field joint acceptance $\left[j_{r}^{2}(\stackrel{e}{\omega})\right]$ ext must be used in place of the reverberant field joint acceptance in Equation (10). This gives

$$
\begin{align*}
& \frac{\left\langle p_{i}^{2}\right\rangle_{s, t}}{\left\langle p_{b 1}^{2}\right\rangle_{s, t}}=\frac{\rho_{I}^{2} c_{o I}^{4} \bar{A}^{2}}{c_{\omega} V^{2}} \cdot \sum_{n} \varepsilon_{n} \sum_{r} \frac{\left[j_{r}^{2}(\omega)\right]_{\text {ext }} A^{2} \bar{f},^{2}(n, r)}{M_{r}^{2} D_{n r}} \\
& x\left\{\left(\frac{c_{r}-c_{n}}{4}\right) \ln n_{n}+\left(\frac{2 c_{n}\left(b_{r}-b_{n}\right)-b_{n}\left(c_{r}-c_{n}\right)}{4 n_{n} \omega_{n}^{2}}\right) \arctan n\right.  \tag{12}\\
& \left.+\left(\frac{c_{n}^{-c} r}{4}\right) \ln n_{r}+\left(\frac{2 c_{r}\left(b_{n}-b_{r}\right)-b_{r}\left(c_{n}-c_{r}\right)}{4\left(\eta_{r}^{\prime}+\eta_{r}^{\prime \prime}\right) \omega_{r}^{2}}\right) \arctan r\right\},
\end{align*}
$$

where the functions $b, c, D$, In are defined in Equation (10). The exterior field joint acceptance is defined in Section 2.8 for boundary layer and progressive wave excitation.

Again for high frequencies, an expression similar to Equation (10) is developed, and the noise reduction is given by

$$
\begin{align*}
& \frac{\left\langle p_{b I}^{2}\right\rangle_{S, t}}{\left\langle p_{i}^{2}\right\rangle_{S, t}}=\left\{\frac{\alpha S}{4 \rho_{I} c_{O I}}+\frac{\rho_{I} c_{O E}}{\rho_{E} c_{O I}} \cdot \pi n_{r}\left\langle\frac{A^{2}}{M_{r}} \cdot j_{r}^{2}(\omega)\right\rangle_{i n t}^{\text {rev }}{ }_{r \in \Delta \omega}\right. \\
& \left.+\frac{\rho_{E}}{2 \pi c_{o E}} \sum_{r<\Delta \omega}\left(\frac{A^{2}}{M_{r}} \cdot j_{r}^{2 r e v}\right)_{\text {int }} \cdot\left(\frac{\bar{A}^{2}}{M_{r}} \cdot j_{r}^{2 r e v}(\omega)\right)_{e x t}\right\} \\
& \div\left\{\frac{\rho_{I} \omega n_{r}}{4 c_{o I} \bar{n}_{r}}\left\langle\left(\frac{\bar{A}^{\prime}}{M_{r}} \cdot j_{r}^{2}(\omega)\right)_{r \varepsilon \Delta \omega} \cdot\left(\frac{A^{2}}{M_{r}} \cdot j_{r}^{2}(\omega)\right)_{i n t}^{r e v}\right)_{r \varepsilon \Delta \omega}\right\rangle \\
& \left.+\frac{\rho I}{4 \pi c_{o I}} \sum_{r<\Delta \omega}\left(\frac{A^{2}}{M_{r}} \cdot j_{r}^{2(\omega)}\right)_{\text {int }}^{r e v}\left(\frac{\bar{A}^{2}}{M_{r}} \cdot j_{r}^{2}(\omega)\right)_{e x t}^{e f}\right\} \tag{13}
\end{align*}
$$

where $\left(A^{2} j_{r}\left(\begin{array}{r}\text { ev }\end{array}\right)\right.$ int is given by Equation (9), and $\left(j_{r}^{2}(\omega)\right)$ ext is given in Section 2.8 .

### 2.5 Four Element Sidewall

### 2.5.1 Sidewall Transfer Matrix

In Appendix A, Ref.[1], the sidewall transfer matrix is developed for a single trim element consisting of insulation and a limp mass. For present purposes the sidewall may be represented by up to 4 elements, each element being made up of a layer of insulation (or air) lined with a limp, dissipative mass.

The transfer matrix across all the elements is given by

$$
\begin{aligned}
& =\left[\begin{array}{l}
a_{11} a_{12} \\
a_{21} a_{22}
\end{array}\right]\left\{\begin{array}{l}
p_{1}^{i} \\
w_{1}
\end{array}\right\}
\end{aligned}
$$

where 1 is the element in contact with the skin, $n$ is the finishing element in the cabin, and $p$ and $w$ are the pressure and displacement at the points shown in Figure 1. The coefficients all etc. are defined in Ref.[1], and are dependent on the acoustical properties of the materials used and the mass and loss factor of the limp mass lining. The complex wave impedance, $W$, and the propagation constant, $\gamma$, of the insulation are required for the coefficients $a_{11}$ etc., where

$$
\begin{equation*}
\gamma=\alpha-12 \pi / \lambda_{m} \tag{15}
\end{equation*}
$$

The amplitude and phase of $W$, the attenuation $\alpha$ in $d B / m$ and the wavelength $\lambda_{\mathrm{m}}$ may be calculated using [9, p. 258 on].


FIGURE 1. SIDEWALL TRIM: INSULATION AND LINING

### 2.5.2 Sidewall Stiffness

The sidewall transfer matrix method developed in [1] assumed that the inner lining was a limp mass, and did not originally include the stiffness of an inner wall. The transfer matrix across a stiff lining may be represented by

$$
\left\{\begin{array}{l}
p_{2}  \tag{16}\\
w_{2}
\end{array}\right\}=\left[\begin{array}{cc}
1 & m_{t}\left(\omega^{2}-\omega_{t}^{2}\right)+i \omega^{2} n_{t} m_{t} \\
0 & 1
\end{array}\right]\left\{\begin{array}{l}
p_{1} \\
w_{1}
\end{array}\right\}
$$

$$
\text { where } \begin{aligned}
m_{t} & =\text { mass of lining } \\
\omega_{t} & =\text { fundamental resonance frequency of lining, and } \\
\eta_{t} & =\text { loss factor associated with lining. }
\end{aligned}
$$

This approximation assumes that the lining is locally reacting as before.

The inner wall resonance frequencies are calculated for a curved panel extending from floor to floor. Because of the strong curvature, the lowest frequencies would typically have many half-wavelengths in the circumferential direction from floor to floor and one half-wavelength in the axial direction.

The inclusion of the fundamental wall resonance frequency in the transfer matrix reduces the effective mass of the inner wall and hence raises the double-wall resonance frequency slightly. However, the effect on the transmission loss of the sidewall is small. Thus inner wall stiffness is not included in the final analysis.

## 2. 6 Loss Factor for Structure with Trim

The analytical model [1] introduces a total structural loss factor, $\eta_{r}^{\prime}$, when a trim is present on the sidewall. This loss
factor replaces the structural loss factor, $\eta_{r}$, which represents dissipation in the structure when there is no trim. According to the analysis in Ref.[1],

$$
\begin{equation*}
{\eta_{r}^{\prime}}^{2}=\frac{\left|C_{W}\right|^{2}}{m^{2} \omega^{4}}-\frac{2 C_{W}^{I}}{m} \frac{\eta_{r}}{\omega_{r}^{2}}+\eta_{r}^{2} \tag{17}
\end{equation*}
$$

where $C_{W}=C_{W}^{R}+i C_{W}^{I}$, and $m$ is the average surface mass of the structure.

When predictions were compared with data from scale model tests performed at NASA Langley (see Appendix E of [l]) it was found that the computed values of $\eta_{r}^{\prime}$ were much greater than measured values at low frequencies (Fig.E-17 of [1]). Consequently an upper limit of $\eta_{r}^{\prime}=0.15$ was introduced into the analytical model.

Further analysis by L.D. Pope [private communication] indicated that, since the trim was present only on the sidewalls, and not on the floor, the parameter $m$ in Eq. (17) above should exclude the floor. He replaced $m$ in $E q .(17)$ by

$$
\begin{equation*}
\bar{m}=m \frac{M_{r}}{M_{r}^{E}} \tag{18}
\end{equation*}
$$

where $M_{r}$ is the generalized mass for mode $r$ of the total structure (shell plus floor) and $M \underset{r}{E}$ is the corresponding generalized mass for the region of the structure covered by trim. The analytical model has been modified to incorporate Eq. (18).

## 2. 7 Joint Acceptance for Fuselage Floor with Reverberant Excitation

The structural modes of the fuselage with floor are derived in Appendix $D$, Ref.[1]. The symmetric mode shape of mode $r$ for the floor is assumed to be a finite series of the form

$$
\begin{equation*}
\psi^{r}(z, x)=\sin \frac{M \pi z}{L} \sum_{n=0}^{n^{*}} c_{M n}^{p r} \cos \frac{n \pi x}{L_{p}} \tag{19}
\end{equation*}
$$

and the antisymmetric mode is

$$
\begin{equation*}
\psi^{r}(z, x)=\sin \frac{M \pi z}{L} \sum_{n=1}^{n *} C_{M n}^{p r} \sin \frac{n \pi x}{L_{p}} \tag{20}
\end{equation*}
$$

where $z$ is the axial coordinate measured from the forward end of the fuselage and $x$ is measured horizontally from the floor center in the lateral direction. The width of the floor $L_{p}$ is given by

$$
L_{\mathrm{p}}=2 \mathrm{a} \sin \theta_{0}
$$

where $\theta_{0}$ is the floor angle measured from the bottom centerline. The coefficients $\mathrm{C}_{\mathrm{Mr}}^{\mathrm{pr}}$ are the generalized coordinates for mode $r$ obtained using Appendix $D, \operatorname{Ref.[1].}$

The reverberant field joint acceptance for the floor is given by

$$
\begin{equation*}
j_{r}^{2}(\omega)=\frac{1}{L^{2} L_{p}^{2}} \iint C\left(\bar{x} \mid \bar{x}^{\prime} ; \omega\right) \psi^{r}(\bar{x}) \psi^{r}(\bar{x}) d \bar{x} d \bar{x}, \tag{21}
\end{equation*}
$$

where $C(\bar{x} \mid \bar{x} ; \omega)$ is assumed separable in the axial and transverse directions and is given by

$$
\begin{equation*}
C(\bar{x} \mid \bar{x} \prime ; \omega)=\frac{\sin k\left(z-z^{\prime}\right)}{k\left(z-z^{\prime}\right)} \cdot \frac{\sin k\left(x-x^{\prime}\right)}{k\left(x-x^{\prime}\right)} \tag{22}
\end{equation*}
$$

where

$$
k=\frac{\omega}{c_{O I}}
$$

The joint acceptance may be written in the form

$$
\begin{equation*}
j_{r}^{2}(\omega)=j_{M}^{2}(\omega) \sum_{n_{1}}^{\text {rev }} \sum_{n_{2}} C_{M n_{1}}^{p r} c_{M n_{2}}^{p r} j_{n_{1} n_{2}}^{2}(\omega) \tag{23}
\end{equation*}
$$

The axial component of the joint acceptance is given by [10] as

$$
\begin{equation*}
j_{M}^{2}(\omega)=I_{1}(M)+I_{2}(M)+I_{3}(M) \tag{24}
\end{equation*}
$$

where

$$
\begin{aligned}
& I_{1}(M)=\frac{1}{2 \pi M k L}\{\operatorname{Cin}(k L+M \pi)-\operatorname{Cin}|M \pi-k L|\} \\
& I_{2}(M)=\frac{1}{2 k L}\{S i(k L+M \pi)-S i(M \pi-k L)\} \\
& I_{3}(M)=\frac{1-(-1)^{M} \cos k L}{(M \pi)^{2}-(k L)^{2}}
\end{aligned}
$$

Si and Cin are the sine and cosine integrals [11].
The lateral component of the joint acceptance $\left.j_{n_{2} n_{2}}^{( }\right)$fev for symmetric modes is

$$
\begin{equation*}
j_{n_{1} n_{2}}^{2}(\omega)=\frac{1}{L_{p}^{2}} \iint_{-L_{p / 2}}^{L_{p / 2}} \frac{\sin k\left(x_{1}-x_{2}\right)}{k\left(x_{1}-x_{2}\right)} \cdot \cos \frac{n_{1} \pi x_{1}}{L_{p}} \cdot \cos \frac{n_{2} \pi x_{2}}{L_{p}} \cdot d x_{1} d x_{2} \tag{25}
\end{equation*}
$$

and for antisymmetric modes is
$j_{n_{1} n_{2}}^{2}(\omega)=\frac{1}{L_{p}^{2}} \iint_{-L_{p / 2}}^{L_{p / 2}} \frac{\sin k\left(x_{1}-x_{2}\right)}{k\left(x_{1}-x_{2}\right)} \cdot \sin \frac{n_{1} \pi x_{1}}{L_{p}} \cdot \sin \frac{n_{2} \pi x_{2}}{L_{p}} \cdot d x_{1} d x_{2}(26)$
It is necessary to evaluate the cross-terms of the lateral joint acceptance, when $n_{1} \neq \mathrm{n}_{2}$, since the individual component mode shapes for the floor are not necessarily orthogonal to one another.

Evaluating the integral in Equation (25) for symmetric modes gives

$$
\begin{align*}
& j_{0, O}^{2}(\omega)=\frac{2}{k L_{p}} \operatorname{Si}\left(k I_{p}\right)-\frac{2}{k^{2} L_{p}^{2}}\left(1-\cos k L_{p}\right) \\
& j_{n, n}^{2}(\omega)=\frac{1}{2 k I_{p}}\left[S i\left(n \pi+k I_{p}\right)-S i\left(n \pi-k L_{p}\right)\right] \\
& -\frac{\cos n \pi}{2 k L_{p} n \pi}\left[\operatorname{cin}\left(n \pi+k L_{p}\right)-\operatorname{Cin}\left(n \pi-k L_{p}\right)\right]  \tag{27}\\
& +\frac{1-\cos n \pi \cos k L_{p}}{n^{2} \pi^{2}-k^{2} L_{p}^{2}} \\
& j_{n_{1} n_{2}}^{2}(\omega)=\frac{1}{\pi k L_{p}\left(n_{1}^{2}-n_{2}^{2}\right)} \cdot\left\{A\left(S i\left(n_{1} \pi+k L_{p}\right)-S i\left(n_{1} \pi-k L_{p}\right)\right)\right. \\
& -B\left(\operatorname{Cin}\left(n_{1} \pi+k I_{p}\right)-\operatorname{Cin}\left(n_{1} \pi-k I_{p}\right)\right) \\
& +A\left(S i\left(n_{2} \pi+k I_{p}\right)-S i\left(n_{2} \pi-k L_{p}\right)\right) \\
& \left.+c\left(\operatorname{cin}\left(n_{2} \pi+k L_{p}\right)-\operatorname{cin}\left(n_{2} \pi-k L_{p}\right)\right)\right\}
\end{align*}
$$

and Equation (26) for antisymmetric modes gives

$$
\begin{aligned}
j_{n n}^{2}(\omega) & =\frac{I}{2 k L_{p}}\left[S i\left(n \pi+k I_{p}\right)-S i\left(n \pi-k I_{p}\right)\right] \\
& +\frac{\cos n \pi}{2 k L_{p} n \pi}\left[\operatorname{Cin}\left(n \pi+k L_{p}\right)-\operatorname{Cin}\left(n \pi-k L_{p}\right)\right] \\
& +\frac{1-\cos n \pi \operatorname{coskI_{p}}}{n^{2} \pi^{2}-k^{2} I_{p}^{2}}
\end{aligned}
$$

$$
\begin{align*}
j_{n_{1} n_{2}}^{2}(\omega)=\frac{1}{\pi k L_{p}\left(n_{1}^{2}-n_{2}^{2}\right)}\{ & -D\left(\operatorname{Si}\left(n_{1} \pi+k L_{p}\right)-S i\left(n_{1} \pi-k L_{p}\right)\right) \\
& -C\left(\operatorname{Cin}\left(n_{1} \pi+k L_{p}\right)-\operatorname{Cin}\left(n_{1} \pi-k L_{p}\right)\right)  \tag{28}\\
& -D\left(\operatorname{Si}\left(n_{2} \pi+k L_{p}\right)-\operatorname{Si}\left(n_{2} \pi-k L_{p}\right)\right) \\
& \left.+B\left(\operatorname{Cin}\left(n_{2} \pi+k L_{p}\right)-\operatorname{Cin}\left(n_{2} \pi-k L_{p}\right)\right)\right\}
\end{align*}
$$

where

$$
\begin{aligned}
& A=n_{1} \sin \frac{\pi n_{1}}{2} \cos \frac{\pi n_{2}}{2}-n_{2} \cos \frac{\pi n_{1}}{2} \sin \frac{\pi n_{2}}{2} \\
& B=n_{1} \cos \frac{\pi n_{1}}{2} \cos \frac{\pi n_{2}}{2}+n_{2} \sin \frac{\pi n_{1}}{2} \sin \frac{\pi n_{2}}{2} \\
& C=n_{1} \sin \frac{\pi n_{1}}{2} \sin \frac{\pi n_{2}}{2}+n_{2} \cos \frac{\pi n_{1}}{2} \cos \frac{\pi n_{2}}{2} \\
& D=n_{1} \cos \frac{\pi n_{1}}{2} \sin \frac{\pi n_{2}}{2}-n_{2} \sin \frac{\pi n_{1}}{2} \cos \frac{\pi n_{2}}{2}
\end{aligned}
$$

Using Equations (24), (27) and (28) in Equation (23) gives the floor joint acceptance for reverberant field excitation.

## 2. 8 Joint acceptance for fuselage with exterior excitation

 fieldThe joint acceptance, describing the coupling between the exterior excitation field and the fuselage structure is defined by

$$
\begin{equation*}
\left.\left.j_{r}^{2}(\omega)=\frac{I}{\bar{A}^{2}} \iint C\left(\bar{x} \mid \bar{x}^{\prime} ; \omega\right) \psi^{r}\right) \bar{x}\right) \psi^{r}\left(\bar{x}{ }^{\prime}\right) d \bar{x} d \bar{x}{ }^{\prime} \tag{29}
\end{equation*}
$$

The correlation function is assumed separable in the axial and circumferential directions and is of the form

$$
\begin{aligned}
C\left(\bar{x} \mid \bar{x}^{\prime} ; w\right) & =C_{z}\left(z-z^{\prime}, w\right) C_{y}\left(a\left(\theta-\theta^{\prime}\right), \omega\right) \\
& =C_{z}\left(z-z^{\prime}, w\right) C_{y}\left(y-y^{\prime}, w\right)
\end{aligned}
$$

where $z$ is the axial coordinate, $\theta$ is the angular coordinate relative to the fuselage bottom centerline and $y=a \theta$ is the distance around the circumference of the fuselage.

For a reverberant (diffuse) excitation field, the correlation functions are

$$
\begin{align*}
& C_{z}\left(z-z^{\prime}, \omega\right)=\frac{\sin k\left(z-z^{\prime}\right)}{k\left(z-z^{\prime}\right)} \text { where } k=\frac{\omega}{C_{O E}}  \tag{30}\\
& C_{y}\left(y-y^{\prime}, \omega\right)=\frac{\sin k\left(y-y^{\prime}\right)}{k\left(y-y^{\prime}\right)} \tag{31}
\end{align*}
$$

For boundary layer excitation, the correlation functions for the axial and circumferential directions are given by Cockburn and Jolly [12] as

$$
\begin{align*}
& C_{z}\left(z-z^{\prime}, \omega\right)=\exp \left[-\left\{\left(0.1 \frac{\omega}{U_{c}}\right)^{2}+\left(\frac{0.034}{\delta^{*}}\right)^{2}\right\}^{\frac{1}{2}}\left|z-z^{\prime}\right|\right] \cos \frac{\omega}{U_{c}}\left(z-z^{\prime}\right)  \tag{32}\\
& C_{y}\left(y-y^{\prime}, \omega\right)=\exp \left[-\left\{\left(0.72 \frac{\omega}{U_{c}}\right)^{2}+\left(\frac{0.244}{\delta^{*}}\right)^{2}\right\}^{\frac{1}{2}}\left|y-y^{\prime}\right|\right] \tag{33}
\end{align*}
$$

where $U_{C}=$ axial trace velocity

$$
\delta^{*}=\text { boundary layer displacement thickness }
$$

The preceding analytical formulation can be readily modified to represent progressive wave excitation where the waves propagate in the axial direction. Although this excitation is not applied in the present investigation the analysis is included here for completeness. An analytical representation of this form was used to describe the pressure field over the Aft Cargo Carrier of the Space Shuttle [13], where the source is random and extends over a large volume to the rear of the vehicle.

In the axial direction, the correlation function is

$$
\begin{equation*}
C_{z}\left(z-z^{\prime}, \omega\right)=\exp \left[-C \frac{\omega}{\left.z \frac{U_{z}}{}\left|z-z^{\prime}\right|\right] \cos \frac{\omega}{U_{z}}\left(z-z^{\prime}\right), ~(z)}\right. \tag{34}
\end{equation*}
$$

where $U_{Z}=$ axial trace velocity

$$
c_{x}=\text { axial decay parameter }
$$

In the circumferential direction the field is assumed reverberant with a correlation function as in Eq. (31).

The structural modes of the fuselage are derived in Appendix $D$ Ref.[1]. The symmetric mode shape of mode $r$ for the fuselage wall is assumed to be a finite series of the form

$$
\begin{equation*}
\psi^{r}(z, \theta)=\sin \frac{M \pi z}{L} \sum_{n=u}^{n^{*}} C_{M n}^{S r}(-I)^{n} \cos n \theta \tag{35}
\end{equation*}
$$

and the antisymmetric mode is

$$
\begin{equation*}
\psi^{r}(z, \theta)=-\sin \frac{M \pi z}{L} \sum_{n=1}^{n} C_{M n}^{s r}(-1)^{n} \sin n \theta \tag{36}
\end{equation*}
$$

$$
\begin{aligned}
\text { where } M= & \text { number of longitudinal half-wavelengths for mode } r \\
n= & \text { number of circumferential wavelengths in the fuselage } \\
& \text { shell } \\
n^{*}= & \text { maximum number of circumferential wavelengths used } \\
& \text { to represent mode } r \\
C_{M n}^{S r}= & \text { fuselage generalized coordinate for mode } r \text { associated } \\
& \text { with } n
\end{aligned}
$$

The joint acceptances for the axial and circumferential directions may be calculated independently for mode $r$ using

$$
\begin{equation*}
j_{r}^{2}(\omega)=j_{M}^{2}(\omega) \cdot \sum_{n=0}^{n^{*}}\left(C_{M n}^{s r}\right)^{2} j_{n}^{2}(\omega) \tag{37}
\end{equation*}
$$

where the axial joint acceptance is given by
$j_{M}^{2}(\omega)=\frac{1}{L^{2}} \int_{0}^{I} \int_{0}^{L} C_{z}\left(z_{1}-z_{2}, \omega\right) \sin \frac{M \pi z_{1}}{L} \cdot \sin \frac{M \pi z_{2}}{L} \cdot d z_{1} d z_{2}$
and the circumferential joint acceptance terms are

$$
j_{n}^{2}(\omega)=\frac{1}{(2 \pi a)^{2}} \int_{0}^{2 \pi a} \int_{0}^{2 \pi a} c_{y}\left(y_{1}-y_{2}, \omega\right)\left\{\begin{array}{c}
\cos \frac{n y_{1}}{a} \cos \frac{n y_{2}}{a}  \tag{39}\\
\sin \frac{n y_{1}}{a} \frac{\sin \frac{n y_{2}}{a}}{a}
\end{array}\right\} d y_{1} d y_{2}
$$

Terms in $\cos \frac{n y}{a}$ and $\sin \frac{n y}{a}$ refer to the symmetric and antisymmetric modes of the fuselage respectively. Unlike the floor joint acceptance in Section 2.7 , cross terms, $j_{n_{1} n_{2}}^{2}(\omega)$, are equal and opposite in sign to $j_{n_{2} n_{1}}^{2}(\omega)$, since $C_{y}\left(y_{1}-y_{2}, w\right)$ is an even function, and hence do not appear in Eq. (37).

### 2.8.1 Axial joint acceptance for reverberant field excitation

The axial joint acceptance is given by [10] as

$$
j_{M}^{2 r e v}(\omega)=I_{1}(M)+I_{2}(M)+I_{3}(M)
$$

where

$$
\begin{align*}
& I_{1}(M)=\frac{1}{2 \pi M k I}\{\operatorname{Cin}(k L+M \pi)-\operatorname{Cin} \mid M \pi-k L \quad\} \\
& I_{2}(M)=\frac{1}{2 k L}\{S i(k I+M \pi)-S i(M \pi-k L)\}  \tag{40}\\
& I_{3}(M)=\frac{1-(-I)^{M} \cos k L}{(M \pi)^{2}-(k L)^{2}}
\end{align*}
$$

Si and Cin are the sine and cosine integrals [II] and $k=\omega / c_{o E}$ is the acoustic wavenumber.

### 2.8.2 Axial joint acceptance for boundary layer or progressive wave excitation

The correlation function from Eqs.(32) and (34) may be written in the general form

$$
\begin{equation*}
C_{z}\left(z_{1}-z_{2}, \omega\right)=e^{-A\left|z_{1}-z_{2}\right|} \cos \frac{\omega}{U}\left(z_{1}-z_{2}\right) \tag{41}
\end{equation*}
$$

where

$$
U=\text { convection or trace velocity }
$$

and for boundary layer excitation

$$
A=\left\{\left(0.1 \frac{\omega}{U}\right)^{2}+\left(\frac{0.034}{\delta^{*}}\right)^{2}\right\}^{\frac{1}{2}}
$$

For progressive wave excitation

$$
A=c_{z \bar{U}}
$$

The axial joint acceptance is given by [14] as
where

$$
\left.\begin{array}{rl}
j_{M}^{2}(\omega) & =\frac{2}{(M \pi)^{2} \Delta_{M}^{2}}\left[p _ { M } \left\{1-(-1)^{M} e^{-A L} \cos \frac{\omega}{U} L\right.\right.
\end{array}\right\}
$$

$$
\begin{aligned}
& \Delta_{M}=\left[I+\left(\frac{\omega L}{U M \pi}\right)^{2}+\left(\frac{A L}{M \pi}\right)^{2}\right]^{2}-4\left[\frac{\omega L}{U M \pi}\right]^{2} \\
& p_{M}=\left[1-\left(\frac{\omega L}{U M \pi}\right)^{2}+\left(\frac{A L}{M \pi}\right)^{2}\right]^{2}-4\left[\frac{\omega L}{U M \pi}\right]^{2}\left[\frac{A L}{M \pi}\right]^{2}
\end{aligned}
$$

$$
\begin{aligned}
& q_{M}=\left[\frac{\omega L}{U M \pi}\right]\left[\frac{A L}{M \pi}\right]\left[1-\left(\frac{\omega L}{U M \pi}\right)^{2}+\left(\frac{A L}{M \pi}\right)^{2}\right] \\
& r_{M}=\left[\frac{A L}{M \pi}\right]\left[1+\left(\frac{\omega L}{U N \pi}\right)^{2}+\left(\frac{A L}{M \pi}\right)^{2}\right]
\end{aligned}
$$

### 2.8.3 Circumferential joint acceptance for reverberant field joint acceptance

The circumferential joint acceptance is given by [10]. For symmetric modes

$$
\begin{align*}
& j_{0}^{2}(\omega)=\frac{S i(2 \pi k a)}{\pi k a}-\frac{(1-\cos 2 \pi k a)}{2(\pi k a)^{2}} \text { for } n=0 \\
& j_{n}^{2}(\omega)=I_{2}(n)+I_{3}(n)+I_{1}(n) \text { for } n \neq 0 \tag{43}
\end{align*}
$$

And for antisymmetric modes

$$
\begin{equation*}
j_{n}^{2}(w)=I_{2}(n)+I_{3}(n)+I_{1}(n) \tag{44}
\end{equation*}
$$

where

$$
\begin{aligned}
& I_{1}(n)=\frac{1}{4 \pi n \cdot 2 \pi k a}\{\operatorname{Cin}[2 \pi(n+k a)]-\operatorname{Cin}[2 \pi(n-k a)]\} \\
& I_{2}(n)=\frac{1}{4 \pi k a}\{\operatorname{Si}[2 \pi(n+k a)]-\operatorname{Si}[2 \pi(n-k a)]\} \\
& I_{3}(n)=\frac{1-\cos 2 \pi k a}{(2 \pi n)^{2}-(2 \pi k a)^{2}}
\end{aligned}
$$

2.8.4 Circumferential joint acceptance for boundary layer excitation

The correlation function from Eq.(33) may be written as

$$
\begin{equation*}
c_{y}\left(y_{2}-y_{2}, w\right)=e^{-B\left|y_{1}-y_{2}\right|} \tag{45}
\end{equation*}
$$

where

$$
B=\left\{\left(0.072 \frac{\omega}{U_{C}}\right)^{2}+\left(\frac{0.244}{\delta^{*}}\right)^{2}\right\}^{\frac{1}{2}}
$$

Let

$$
\delta=2 \pi a B
$$

Expanding the techniques used in [10], the circumferential joint acceptances, for symmetric modes, are given by

$$
\begin{align*}
& \left.j_{0}^{2}(\omega)=\frac{2}{\delta}-\frac{2 \delta^{2}\left(1-e^{-\delta}\right.}{\delta^{2}}\right) \\
& j_{n}^{2}(\omega)=\frac{\delta}{(2 \pi n)^{2}+\delta^{2}}-\frac{2 \delta^{2}\left(1-e^{-\delta}\right)}{\left[(2 \pi n)^{2}+\delta^{2}\right]^{2}} \tag{46}
\end{align*}
$$

and for antisymmetric modes

$$
\begin{equation*}
j_{n}^{2}(\omega)=\frac{\delta}{(2 \pi n)^{2}+\delta^{2}}+\frac{2(2 \pi n)^{2}\left(1-e^{-\delta}\right)}{\left[(2 \pi n)^{2}+\delta^{2}\right]^{2}} \tag{47}
\end{equation*}
$$

### 3.0 SUMMARY OF METRO II TEST

The noise transmission test conducted by Prydz et al [5] on a Fairchild Metro II fuselage was performed in an anechoic chamber using an electropneumatic acoustic source. Aspects of the test relevant to the present study are given here for ready reference.

### 3.1 Test Fuselage Structure

The test structure consisted of a section of a Fairchild Metro II fuselage with a special floor installed about one foot above the structural floor level of a Metro II in airline service. The special floor could be rigidly attached by brackets to the frames along the full length of the fuselage structure or mounted on air mounts which were located on the structural floor of the fuselage. Cross-sections through the test fuselage are shown in Figure 2 for the attached and isolated floor configurations.

The fuselage shell is of conventional aluminum skin-stringerframe construction. Basic structural properties of the shell are listed in Table 1. Structural loss factors were measured by Prydz et al [5] for the baseline (bare) fuselage; empirical values are given in Table 2. As part of the test program, the surface density of the structure was increased by the addition of a sheet of iron-oxide vinyl with a surface density of $4.88 \mathrm{~kg} / \mathrm{m}^{2}$. Measured loss factors for the structure with added vinyl are given in Table 3.


FIGURE 2. CROSS-SECTION THROUGH METRO II FUSELAGE SHOWING FLOOR INSTALLATION FOR TEST [5]

TABLE 1 - OUTER WALL STRUCTURAL PROPERTIES OF
METRO II TEST FUSELAGE

| Item | Units |  | Values |  |
| :---: | :---: | :---: | :---: | :---: |
|  | SI | English | SI | English |
| Cylindrical structure |  |  |  |  |
| Diameter | m | (ft) | 1.68 | (5.5) |
| Length | m | (ft) | 9.02 | (29.6) |
| Skin thickness | cm | (in.) | 0.102 | (0.040) |
| Suriace density of skin | $\mathrm{kg} / \mathrm{m}^{2}$ | (pst) | 2.78 | (0.57) |
| Surface density skin plus stiffeners | $\mathrm{kg} / \mathrm{m}^{2}$ | (pst) | 4.59 | (0.94) |
| Frames |  |  |  |  |
| Spacing | cm | (in.) | 38.1 | (15) |
| Depth | cm | (in.) | 5.08 | (2) |
| Cross sectional area | $\mathrm{cm}^{2}$ | (in ${ }^{2}$ ) | 0.884 | (0.137) |
| Area centroid (re: skin q) | cm | (in.) | 2.835 | (1.116) |
| 2nd area moment (re: skin q. | cm ${ }^{4}$ | (in ${ }^{4}$ ) | 9.303 | (0.2235) |
| Torsion constant | $\mathrm{cm}^{4}$ | $\left(i^{4}\right)$ | 0.0030 | $\left(0.73 \times 10^{-4}\right)$ |
| Stringers |  |  |  |  |
| Spacing | cm | (in.) | 18.3 | (7.2) |
| Depth | cm | (in.) | 2.223 | (0.875) |
| Cross sectional area | $\mathrm{cm}^{2}$ | $\left(\mathrm{in}^{2}\right)$ | 0.716 | (0.111) |
| Area centroid (re: skin q) | cm | (in.) | 0.772 | (0.304) |
| 2nd area moment (re: skin qu) | $\mathrm{cm}^{4}$ | $\left(\mathrm{in}^{4}\right)$ | 0.11904 | (0.00286) |
| Torsion constant | cm ${ }^{4}$ | (in ${ }^{4}$ ) | 0.00450 | $\left(0.108 \times 10^{-3}\right)$ |


| Dual Panel Grid |  | Circumferential Grid |  |
| :---: | :---: | :---: | :---: |
| Frequency ( $\mathrm{H}_{2}$ ) | Loss Factor (\%) | Frequency ( Hz ) | Loss Factor (\%) |
| 129.39 | 2.18 | 40.00 |  |
| 158.69 | 1.37 | 83.01 |  |
| 258.79 | 1.97 | 83.01 | 6.79 |
| 285.64 | 0.58 | 146.48 | 4.74 |
| 307.62 | 1.39 | 222.17 | 2.45 |
| 341.80 | 1.66 | 266.11 | 2.27 |
| 356.45 | 1.66 | 297.85 | 1.99 |
| 383.30 | 1.31 | 405.27 | 3.79 |
| 446.78 | 1.28 | 424.80 | 2.91 |
| 446.78 493.16 | 2.19 | 446.78 | 1.94 |
| 493.16 | 2.25 | 493.16 | 2.31 |
| 561.52 | 0.97 | 556.64 | 1.74 |
| 686.04 | 0.93 | 607.91 | 1.36 |
| 778.81 | 1.72 | 671.39 | 0.50 |
| 800.78 | 2.45 | 708.01 | 1.32 |
| 808.10 | 0.78 | 756.84 | 0.83 |
| 849.61 | 1.08 | 781.25 | 0.91 |
| 891.11 | 0.80 | 908.20 | 1.32 |
| 910.64 | 1.87 | 1037.60 | 1.00 |
| 1037.60 | 1.23 | 1137.70 |  |
| 1069.34 | 1.84 |  | 0.91 |
| 1110.84 | 1.63 |  |  |
| 1164.55 | 0.87 |  |  |

TABLE 3 - DAMPING VALUES FOR FUSELAGE STRUCTURE WITH VINYL SHEET ON EXTERIOR

| Single Panel Grid |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Frequency <br> $(\mathrm{Hz})$ | Loss Factor <br> $(\%)$ | Frequency <br> $(\mathrm{Hz})$ | Loss Factor <br> $(\%)$ |  |  |  |
| 65.92 | 8.33 | 825.20 | 4.80 |  |  |  |
| 97.66 | 11.85 | 834.96 | 1.91 |  |  |  |
| 275.88 | 3.79 | 852.05 | 1.70 |  |  |  |
| 305.18 | 5.15 | 866.70 | 1.52 |  |  |  |
| 429.69 | 3.47 | 888.67 | 0.63 |  |  |  |
| 456.54 | 7.20 | 947.27 | 4.96 |  |  |  |
| 571.29 | 2.60 | 957.03 | 5.41 |  |  |  |
| 676.27 | 2.71 | 1101.07 | 0.82 |  |  |  |
| 773.93 | 3.73 | 1164.55 | 2.00 |  |  |  |
| 810.55 | 0.62 | 1191.41 | 9.29 |  |  |  |

TABLE 4 - INTERIOR ABSORPTION COEFFICIENTS

| Frequency <br> $(\mathrm{Hz})$ | Decay Method |  | Reference Sound <br> Source Method |
| :---: | :---: | :---: | :---: |
|  | Bare Interior | Added Interior Absorption | Added Absorption |
| 200 | 0.14 | 0.22 | 0.35 |
| 250 | 0.13 | 0.39 | 0.32 |
| 315 | 0.14 | 0.43 | 0.44 |
| 400 | 0.10 | 0.32 | 0.29 |
| 500 | 0.10 | 0.32 | 0.51 |
| 630 | 0.14 | 0.34 | 0.34 |
| 800 | 0.10 | 0.33 | 0.27 |
| 1000 | 0.10 | 0.31 | 0.28 |
| 1250 | 0.09 | 0.34 | 0.28 |
| 1600 | 0.10 | 0.38 | 0.26 |
| 2000 | 0.08 | 0.39 | 0.26 |
| AVG | 0.11 | 0.34 | 0.33 |

### 3.2 Test Floor Structure

The cabin floor used in the Metro II tests consisted of plywood panels ( $1.9 \mathrm{~cm}, 0.75$ inch thick) on steel I-beams. Crosssections through parts of the floor are shown in Fig. 3. The 7.6 cm ( 3 x .17 inch, $1.98 \mathrm{lb} / \mathrm{ft}$ ) deep beams supported the floor panels in the transverse direction on a typical spacing of 1.2 m ( 48 inches). The two 15.2 cm ( 6 x .23 inch, $4.3 \mathrm{lb} / \mathrm{ft}$ ) deep beams were placed in the longitudinal direction, near to the edge of the floor panels and the air mounts (see Fig.2(b)).

Rigid connection between the floor and fuselage shell was achieved by bolting the plywood floor panels to brackets on the frames along the fuselage. For the "floating" floor configuration the transverse I-beams rested on air mounts with low resonance frequencies.

### 3.3 Excitation

Two excitation acoustic fields were used during the Metro II tests [5] -- (a) random, broadband and (b) deterministic, discrete frequency. The latter excitation is of interest here. The sound levels were generated using an electropneumatic source coupled to a horn which was pointed towards the fuselage (Fig. 4). The deterministic field had strong spatial characteristics which were intended to simulate those of a propeller noise field. Longitudinal and circumferential spatial distributions of sound pressure level are shown in Fig. 5. In the case of the longitudinal distribution, the sound levels are free-field values, whereas for the circumferential direction the sound levels were measured on the test cylinder and include reflection effects.

The test section of the fuselage exposed to high intensity acoustic excitation is 3.66 m (144 in) long, with the noise source at the center of the test section (Sta. 377.3).


FIGURE 3. DETAILS OF FLOOR IN METRO II TEST FUSELAGE [5]


FIGURE 4. TEST SET-UP FOR ACOUSTIC TESTS ON METRO II [5]
(a) Longitudinal Distribution of Free Field Sound Levels

(b) Circumferential Distribution of Sound Levels on Fuselage


FIGURE 5. SPATIAL DISTRIBUTION OF SOUND LEVELS FOR TESTS ON METRO II [5]

### 3.4 Interior Treatment

The interior treatment of the test fuselage consisted of a double-panel sidewall with a fiberglass blanket and airspace in between. Sections through the sidewall are shown in Figure 6. The fiberglass had a fiber diameter of about 2 microns and a flow resistance of approximately 257 mks rayls/cm.

The baseline interior trim panel consisted of an aluminum panel with a thickness of 0.089 cm . Additional material, in the form of one or two sheets of vinyl impregnated with iron oxide, was added to the trim panel for noise reduction parametric studies.

The acoustic absorption within the fuselage was increased by the installation of open-pore Scott foam placed on the floor near the sidewall, along the full length of the cabin. Interior absorption coefficients measured in the test fuselage with and without the foam material installed, are given in Fig. 7 and Table 4. These are normalized to the curved surface transmitting area of $28.9 \mathrm{~m}^{2}\left(311 \mathrm{ft}^{2}\right)$.

### 3.5 Test Configurations

Nine test configurations were investigated by Prydz et al [5]. These configurations are listed in Table 5. The first six configurations are associated with a floating floor (suspended on air mounts) and a floating trim (supported by the floor). Since the analytical model assumes that there is no structureborne path between the fuselage structure and the trim panel, it forms a reasonably good representation of the floating trim.


FIGURE 6. SECTIONS THROUGH SIDEWALL TREATMENT FOR TESTS ON METRO II [5]


FIGURE 7. ABSORPTION COEFFICIENTS MEASURED IN TREATED METRO II FUSELAGE [5]

TABLE 5. TEST CONFIGURATIONS

| Configuration No. | Surface Density, $\mathrm{kg} / \mathrm{m}^{2}(\mathrm{psf})$ |  |  |  | Floor Trim Configuration |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\underset{\sigma_{1}}{\text { Outer }}$ |  | $\underset{\sigma_{2}}{\text { Inner }} \text { Wall }$ |  |  |  |
| 1 | 4.590* | (0.94) | 2.25** | (0.46) | Floating | Floating |
| 2 | 4.59 | (0.94) | 7.13 | (1.46) | Floating | Floating |
| 3 | 4.59 | (0.94) | 12.01 | (2.46) | Floating | Floating |
| 4 | 9.47 | (1.94) | 2.25 | (0.46) | Floating | Floating |
| 5 | 9.47 | (1.94) | 7.13 | (1.46) | Floating | Floating |
| 6 | 9.47 | (1.94) | 12.01 | (2.46) | Floating | Floating |
| 7 | 9.47 | (1.94) | 7.13 | (1.46) | Attached | Disconn. |
| 8 | 9.47 | (1.94) | 7.13 | (1.46) | Disconn. | Attached |
| 9 | 9.47 | (1.94) | 7.13 | (1.46) | Attached | Attached |

* Baseline (bare) outer wall surface density.
** Baseline (bare) inner wall surface density.


### 4.0 ANALYTICAL REPRESENTATION OF METRO II TEST STRUCTURE

### 4.1 Fuselage Shell

The fuselage shell was modeled as a cylinder of uniform thickness, representing the skin and 'smeared-out' frame and stringer areas. Additional bending stiffnesses representing 'smearedout' frame and stringer bending stiffnesses were included.

Using the structural data in Table 1 , the baseline configuration was:

Surface density of skin and stiffeners $=4.59 \mathrm{~kg} / \mathrm{m} 2$ ( 0.94 psf ) Actual skin thickness $=0.1016 \mathrm{~cm}$ ( 0.04 in )
Equivalent (skin \& stiffener) thickness $=0.1640 \mathrm{~cm}(0.064 \mathrm{in})$ Additional stiffener bending rigidity $=$

$$
\left[\frac{\text { EI }}{\text { Spacing }}\right]_{\text {Stiffener }}+\operatorname{D}_{\substack{\text { actual } \\ \text { Skin }}}-D_{\substack{\text { equivalent } \\ \text { skin }}}
$$

Additional frame bending rigidity $=1.7655 \times 10^{4} \mathrm{~N} . \mathrm{m}$ Additional stringer bending rigidity $=4.4822 \times 102 \mathrm{~N} . \mathrm{m}$

For configurations where vinyl was added to the outer wall, the only change made to the analytical model was to the mass, the total surface density becoming $9.47 \mathrm{~kg} / \mathrm{m}^{2}$ ( 1.94 psf ). The skin thickness and stiffnesses were assumed to be unchanged.

### 4.2 Floor Structure

### 4.2.1 Attached Floor

The attached, or fixed, floor was modeled as a uniform plate connected to the fuselage along its longitudinal sides. The floor could either be hinged along the attachment line or
rigidly connected so that there was no relative rotation between fuselage and floor along the attachment line. Since the brackets connecting the floor to the frames (see Figures 2(a)) occur at approximately 38.1 cm ( 15 in ) intervals, the floor was assumed to be hinged along the attachment line.

The plywood floor, 1.905 cm thick ( 0.75 in), was assumed to have the following characteristics:

$$
\begin{aligned}
& \text { Density }=544.6 \mathrm{~kg} / \mathrm{m} 3(34 \mathrm{lb} / \mathrm{ft} 3) \\
& \text { Young's Modulus } E=1.103 \times 10^{10} \mathrm{~N} / \mathrm{m}^{2}\left(1.6 \times 10^{6} \mathrm{Ib} / \mathrm{in}\right) \\
& \text { Poisson's ratio }=0 \\
& \text { Typical moment of inertia for } 0.30 \mathrm{~m}(12-\text { inch }) \text { wide plywood } \\
& \text { panel }=8.782 \times 10^{-8} \mathrm{~m}^{4}(0.211 \mathrm{in} 4)
\end{aligned}
$$

The transverse floor support beams had an average spacing of 1.224 m ( 48.2 in ) and were $7.6-\mathrm{cm}$ I-beams ( 3 x .17 in , $1.98 \mathrm{lb} / \mathrm{ft}$ ). Two $15.2-\mathrm{cm} \mathrm{I}$ beams ( $6 \mathrm{x} .23 \mathrm{in}, 4.3 \mathrm{lb} / \mathrm{ft}$ ) supported the floor longitudinally along the edges, 0.508 m (20 in.) from the floor centerline. The floor had a width of 158 cm ( 62.2 in ) and was 28 cm ( 11 in .) below the fuselage centerline; the radius from the center of the fuselage to the floor/fuselage junction made an angle of $70.5^{\circ}$ to the vertical.

The stiffeners were again assumed 'smeared-out' over the floor giving, for the fixed floor,

$$
\text { Surface density of floor }+ \text { stiffeners }=20.879 \mathrm{~kg} / \mathrm{m}^{2}
$$

$$
\text { Actual floor thickness }=1.905 \mathrm{~cm}(0.75 \mathrm{in})
$$

Equivalent (floor + stiffener) thickness $=4.012 \mathrm{~cm}(1.58 \mathrm{in})$

Additional transverse bending rigidity $=1.0081 \times 105 \mathrm{~N} . \mathrm{m}$ Additional longitudinal bending rigidity $=1.5790 \times 10^{6} \mathrm{~N} . \mathrm{m}$

### 4.2.2 Floating Floor Representation

Although the plywood floor was located 28 cm (ll in) below the fuselage centerline, the floating floor was supported on air mounts attached to the structural floor (Figure 2). It was assumed therefore that the floating floor connection to the fuselage occurred 68.6 cm ( 27 in ) below the fuselage centerline and that the floor had a width of 96.4 cm ( 37.95 in ); the radius to the floor/fuselage junction made an angle of $35.1^{\circ}$ to the vertical.

The same total mass was used as for the fixed floor, thus increasing the surface density of the equivalent narrower floating floor. No information was available on the structural floor of the fuselage, but it was assumed to be at least as stiff as the fixed floor. [Photographs in Figs. 4 and 6 of Ref.[5] show closely spaced circumferential frames and longitudinal stiffness in the structural floor]. The same beam stiffnesses were taken as for the attached floor, with a spacing of 38.1 cm ( 15 in) for the transverse beams and a spacing of 50.8 cm (20 in) for the longitudinal beams. The parameters used were

$$
\text { Surface density }=34.328 \mathrm{~kg} / \mathrm{m}^{2}
$$

Equivalent plywood floor thickness $=4.915 \mathrm{~cm}(1.93 \mathrm{in})$

Additional transverse bending rigidity $=3.9536 \times 105 \mathrm{~N} . \mathrm{m}$

Additional longitudinal bending rigidity $=3.1613 \times 10^{6} \mathrm{~N} \cdot \mathrm{~m}$

### 4.3 Structural Modes

The structural modes of the fuselage with the floating floor were calculated with and without the vinyl added to the outer skin. The modes are calculated using displacement series for both the shell and floor displacements [1], allowing 14 terms for the shell and 5 terms for the floor. The fuselage was assumed to be freely supported at its ends, allowing axial displacements.

The number of axial half-wavelengths, $M$, was varied from 1 to 10 , and the first 20 symmetric and 20 antisymmetric modes were calculated for each value of M, giving a total of 400 modes. Tables 6, 7 and 8 list the first 16 modes for the floating floor models, without and with add-on vinyl on the outer skin and for the fixed floor with add-on vinyl. These tables list only the 5 largest coefficients associated with the circumferential shell displacement series and the 3 largest for the transverse floor displacement series, which are used in calculating the fuselage response.

Figure 8 shows the first four structural mode shapes for the floating floor configuration with add-on vinyl, for both fuselage and floor. At these frequencies, the heavy floor is acting as a rigid mass.

### 4.4 Measured Structural Modes

Some mode shapes were measured [5] using circumferential grids and rectangular panel grids. It was not possible to compare the mode shapes measured by the panel grid with calculations, but a comparison was made for the mode measured by the circumferential grid shown in Figure 9.

| mulue | Fktou | Mune | Siskl |  |  | Plate |  |  | generaliled mass TOTAL SHELL |  | $(K G)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Nu | ( $H C$ ) | TYPL | n | * | CMid | M | $N$ | CAN |  |  | PLATE W |
| 1 | 37.05 | SYMA | 1 | 1 | $.02<00$ | 1 | 0 | . 76213 | 207.28017 | 35.24797144 .70867 |  |
|  |  |  | 1 | 2 | -0.47230 | 1 | 1 | .03214 |  |  |  |
|  |  |  | 1 | 3 | -.10472 | 1 | 2 | . 00303 |  |  |  |
|  |  |  | 1 | 4 | .00200 |  |  |  |  |  |  |
|  |  |  | 1 | 3 | -.vくく17 |  |  |  |  |  |  |
| 2 | 63.31 | SYMM | 1 | 2 | .35320 | 1 | 1 | -. 02700 | 34.71538 | 23.86738 | .21442 |
|  |  |  | 1 | 1 | . $317 \%$ | 1 | 0 | -. 01470 |  |  |  |
|  |  |  | 1 | 3 | .16779 | 1 | 3 | $.000 \pm 1$ |  |  |  |
|  |  |  | 1 | 4 | -.04700 |  |  |  |  |  |  |
|  |  |  | 1 | 3 | .010s3 |  |  |  |  |  |  |
| 3 | 08.20 | SYMM | 2 | 2 | -08675 | 2 | $\begin{aligned} & 0-.71010 \\ & 1-.06464 \\ & 2-.00063 \end{aligned}$ |  | 142.84940 | 45.91989 | 84.23744 |
|  |  |  | 2 | 3 | . 17508 | 2 |  |  |  |  |  |  |
|  |  |  | 2 | 4 | -.1103 | 2 |  |  |  |  |  |  |
|  |  |  | 2 | 1 | -. $070<5$ |  |  |  |  |  |  |  |
|  |  |  | 2 | 5 | .04137 |  |  |  |  |  |  |  |
| 4 | 77.50 | ANTI | 1 | 2 | -.50757 |  | $\begin{aligned} & 1 \\ & 2 \\ & 3 \end{aligned}$ | $\begin{aligned} & -.09205 \\ & .00506 \\ & .00048 \end{aligned}$ | 40.52281 | 27.50887 | . 57172 |
|  |  |  | 1 | 1 | -. 33307 | $1$ |  |  |  |  |  |
|  |  |  | 1 | 1 | -. 26474 |  |  |  |  |  |  |
|  |  |  | 1 | 5 | .02030 |  |  |  |  |  |  |
|  |  |  | 1 | 0 | -.02110 |  |  |  |  |  |  |
| 5 | 110.47 | ANTI | 2 | 2 | -. 52041 | 2 | $\begin{array}{r} 1-.01106 \\ 2-.00514 \\ 4-.00003 \end{array}$ |  | 32.80658 | 26.47796 | . 00365 |
|  |  |  | 2 | 3 | -. 43466 | 2 |  |  |  |  |  |  |
|  |  |  | 2 | 1 | -. 11784 | 2 |  |  |  |  |  |  |
|  |  |  | 2 | 5 | $.038 \pm 0$ |  |  |  |  |  |  |  |
|  |  |  | 2 |  | -.03737 |  |  |  |  |  |  |  |
| 6 | 111.10 | SYAM | 3 | 2 | . 83544 | 3 | 0 | -. 77703 | 160.32208 | 40.94124107 .73928 |  |
|  |  |  | 3 | 4 | -. 10 d 70 | 1 | 1 | -. 11255 |  |  |  |  |
|  |  |  | 3 | 1 | -.073.5 | 3 | 2 | -.02188 |  |  |  |  |
|  |  |  | 3 | $b$ | -00011 |  |  |  |  |  |  |  |
|  |  |  | 3 |  | $-.01009$ |  |  |  |  |  |  |  |
| 7 | 118.19 | ANTI | 1 | 3 | . 52148 | 1 | 1 |  | 54.48301 | 29.77090 | .65431 |
|  |  |  | 1 | 1 | -.43<17 | 1 | 2 | -. $\operatorname{cc} 447$ |  |  |  |
|  |  |  | 1 | 2 | - 2 J b 7 | 1 | 3 | . 00075 |  |  |  |
|  |  |  | 1 | 4 | .07784 |  |  |  |  |  |  |
|  |  |  | 1 | 5 | -.04400 |  |  |  |  |  |  |
| 8 | 128.86 | SYMM | 2 | 3 | -35402 |  | 0 |  | 29.49626 | 16.93045 | 6.79309 |
|  |  |  | 2 | 2 | - 3ubut, | 2 | 1 | $. C 4413$ |  |  |  |
|  |  |  | c | 1 | - 213>0 | 2 | 2 | .00878 |  |  |  |
|  |  |  | $L$ | 4 | .11512 |  |  |  |  |  |  |
|  |  |  | 2 |  | -. 04304 |  |  |  |  |  |  |

## TABLE 6. STRUCTURAL MODES FOR FLOATING FLOOR CONFIGURATION (Continued)

| Mujut | FREL | muni. | SHELL |  |  | plate |  |  | GENERALILED MASS |  | $\begin{aligned} & \text { IKG } \\ & \text { PLATE } \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Nu | (HZ) | TYPL | 4 | $N$ | CMis | M | N | CMN | TGTAL | SHELL W |  |
| 4 | 148.74 | A WTI | 3 | 3 | . 55243 | 3 | 1 | -. 10557 | 32.46110 | 20.60793 | .87220 |
|  |  |  | 3 | 2 | .41645 | 3 | 2 | -.0c303 |  |  |  |
|  |  |  | 3 | 4 | . 10416 | 3 | 3 | .00077 |  |  |  |
|  |  |  | 3 | 5 | -.03238 |  |  |  |  |  |  |
|  |  |  | 3 | 1 | . 02200 |  |  |  |  |  |  |
| 10 | 150.06 | SYMM | 1 | 3 | .60104 | 111 | 102 | .14830 <br> .09829 <br> .00678 | 43.01100 | 32.41422 | 5.91548 |
|  |  |  | 1 | 4 | -35581 |  |  |  |  |  |  |
|  |  |  | 1 | 1 | -. 13273 |  |  |  |  |  |  |
|  |  |  | 1 | 5 | -.08200 |  |  |  |  |  |  |
|  |  |  | 1 | 8 | -.01033 |  |  |  |  |  |  |
| 11 | 155.91 | ANTI | 1 | 3 | -. 07068 |  |  | $\begin{array}{r} .41814 \\ -.00223 \\ .00110 \end{array}$ | 263.86270 | 51.8278113 .07471 |  |
|  |  |  | 1 | 2 | . 31225 |  |  |  |  |  |  |  |  |
|  |  |  | 1 | 4 | -. 27636 |  |  |  |  |  |  |  |  |
|  |  |  | 1 | 1 | -. 0.07620 |  |  |  |  |  |  |  |  |
|  |  |  | 1 | 5 | . 02808 |  |  |  |  |  |  |  |  |
| 12 | 158.01 | S YMM | 3 | 3 | . 53727 | 333 | 102 | .08483 <br> .08109 <br> .00543 | 27.99343 | 22.04181 | 2.85620 |
|  |  |  | 3 | 2 | . 23616 |  |  |  |  |  |  |
|  |  |  | 3 | 4 | . 2320 B |  |  |  |  |  |  |
|  |  |  | 3 | 5 | -.05570 |  |  |  |  |  |  |
|  |  |  | 3 | 1 | . 03986 |  |  |  |  |  |  |
| 13 | 158.09 | SYAM | 4 | 2 | -.498d8 |  |  | $\begin{aligned} & .70448 \\ & .17600 \\ & .03333 \end{aligned}$ | 137.90877 | $32 \cdot 43707100 \cdot 30573$ |  |
|  |  |  | 4 | 3 | .44505 |  |  |  |  |  |  |  |  |
|  |  |  | 4 | 4 | . 32343 |  |  |  |  |  |  |  |  |
|  |  |  | 4 | 5 | -. 12543 |  |  |  |  |  |  |  |  |
|  |  |  | 4 | 1 | .08467 |  |  |  |  |  |  |  |  |
| 14 | 167. 67 | ANTI | 2 | 3 | .50433 | 2 1 -.41638 <br> 2 2 .00339 <br> 2 3 .00288 |  |  | 48.57781 | 25.81028 | 12.72050 |
|  |  |  | 2 | 1 | -. 34443 |  |  |  |  |  |  |  |  |
|  |  |  | 2 | 2 | -. 23570 |  |  |  |  |  |  |  |  |
|  |  |  | 2 | 4 | . .19515 |  |  |  |  |  |  |  |  |
|  |  |  | 2 |  | -. 02677 |  |  |  |  |  |  |  |  |
| 13 | 169.25 | SYMM | 2 | 3 | 3. 33443 | 22 |  | $\begin{aligned} & 1-13740 \\ & 0-\operatorname{csu} 0 \\ & 2-00437 \end{aligned}$ | 45.46452 | 33.25248 | . 43477 |
|  |  |  | 2 | 4 | -.37920 |  | 0 |  |  |  |  |
|  |  |  | 2 | 1 | -.37336 |  |  |  |  |  |  |
|  |  |  | 2 | 2 | - -14345 |  |  |  |  |  |  |
|  |  |  | 2 |  | -.05063 |  |  |  |  |  |  |
| 10 | 162.58 | ANYI | 1 |  | 3-.55240 |  | 1 | . 447u0 | 163.14356 | 27.31395 | 72.33242 |
|  |  |  | 1 | 2 | - 3 5 ¢ b |  |  | $-.01273$ |  |  |  |
|  |  |  | 1 | 4 | 4-201ヵn |  |  | -. UUく06 |  |  |  |
|  |  |  | 1 |  | $1-.11081$ |  |  |  |  |  |  |
|  |  |  |  |  | -.03416 |  |  |  |  |  |  |

TABLE 7. STRUCTURAL MODES FOR FLOATING FLOOR CONFIGURATION WITH ADD-ON VINYL


TABLE 7. STRUCTURAL MODES FOR FLOATING FLOOR CONFIGURATION WITH ADD-ON VINYI (Continued)

CALAC METKU. DIAMETER L.G甘M, LENOTH 9M, FLOATING FLOUR MUUEL + ADO-ON VINYL

| Mude | FKEZ | MOUE | SHELL |  |  | PLATE |  |  | GENERALILED MASS TOTAL SHELL |  | $\begin{aligned} & \text { (KG) } \\ & \text { PLATE } \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| NU | (H2) | TYPE | H | N | CMN | M | N | CMN |  |  |  |
| 4 | 104.13 | ANTI | 3 | 3 | . 54147 | 3 | 1 | -. 08738 | 64.99790 | 54.51492 | .59896 |
|  |  |  | 3 | 2 | . 42153 | 3 | 2 | -. 00263 |  |  |  |
|  |  |  | 3 | 4 | . 10854 | 3 | 3 | .00092 |  |  |  |
|  |  |  | 3 | 5 | -. 03203 |  |  |  |  |  |  |
|  |  |  | 3 | 1 | . 02447 |  |  |  |  |  |  |
| 10 | 106.96 | SYMM | 1 | 3 | .67031 | 1 | 0 | .15933 | 80.72721 | 69.77585 | 8.24122 |
|  |  |  | 1 | 4 | . 38571 | 1 | 1 | .11365 |  |  |  |
|  |  |  | 1 | 1 | -. 11950 | 1 |  | $.00548$ |  |  |  |
|  |  |  | 1 | 5 | -. 077790 |  |  |  |  |  |  |
|  |  |  | 1 | 2 | -. 02692 |  |  |  |  |  |  |
| 11 | 113.13 | SYMM | 3 | 3 | . 60286 | 3 | 0 | . 21069 | 70.82922 | 53.95066 | 10.59583 |
|  |  |  | 3 | 4 | . 31543 | 3 | 1 | . 08481 |  |  |  |
|  |  |  | 3 | 2 | . 09800 | 3 | 2 | .00622 |  |  |  |
|  |  |  | 3 | 5 | -. 06578 |  |  |  |  |  |  |
|  |  |  | 3 | 1 | .04207 |  |  |  |  |  |  |
| 12 | 118.98 | SYMM | 2 | 3 | -. 50907 | 2 | 0 | .10587 | 95.55709 | 67.98480 | . 46419 |
|  |  |  | 2 | 1 | . 40924 | 2 | 1 | -. 08537 |  |  |  |
|  |  |  | 2 | 4 | -. 37481 | 2 | 2 | . 00434 |  |  |  |
|  |  |  | 2 | 2 | $.18624$ |  |  |  |  |  |  |
|  |  |  | 2 | 5 | $.04216$ |  |  |  |  |  |  |
| 13 | 124.05 | ANTI | 2 | 3 | $-.46676$ |  |  |  | 75.85460 | 50.06625 | 6.90991 |
|  |  |  | 2 | 1 | $.32713$ | 2 | 2 | -. 00573 |  |  |  |
|  |  |  | 2 | 4 | -. 25978 | 2 | 3 | -. 00256 |  |  |  |
|  |  |  | 2 | 2 | . 22371 |  |  |  |  |  |  |
|  |  |  | 2 | 6 | .03674 |  |  |  |  |  |  |
| 14 | 230.09 | SYMM | 4 | 3 | . 0.7465 | 4 | 0 | . 32579 | 102.31225 | 71.81968 | 22.52389 |
|  |  |  | 4 | 4 | $.41778$ | 4 | 1 | $.09602$ |  |  |  |
|  |  |  | 4 | 5 | -. 066966 | 4 |  | $.00749$ |  |  |  |
|  |  |  | 4 | 2 | -. 04954 |  |  |  |  |  |  |
|  |  |  | 4 | 1 | . 02212 |  |  |  |  |  |  |
| 15 | 133.98 | ANTI | 4 | 3 | . 56383 | 4 | 1 | -. 11020 | 59.57925 | 51.03835 | . 89649 |
|  |  |  | 4 | 2 | . 29937 | 4 | 3 | . 00118 |  |  |  |
|  |  |  | 4 | 4 | .21383 | 4 | 2 | .00051 |  |  |  |
|  |  |  | 4 | 6 | $-.01329$ |  |  |  |  |  |  |
|  |  |  | 4 | 7 | $.01160$ |  |  |  |  |  |  |
| 16 | 134.72 | ANTI | 1 | 3 | -. 54529 | 1 | 1 | . 34732 | 199.24715 | 96.71058 | 8.53262 |
|  |  |  | 1 | 4 | -. 50007 | 1 | 2 | -. 00997 |  |  |  |
|  |  |  | 1 | 2 | . 43134 | 1 |  | -. 00246 |  |  |  |
|  |  |  | 1 | 1 | $-.23126$ |  |  |  |  |  |  |
|  |  |  | 1 |  | -. 08406 |  |  |  |  |  |  |

TABLE 8. STRUCTURAL MODES FOR FIXED FLOOR CONFIGURATION WITH ADD-OiN VIivyL


TABLE 8．STRUCTURAL MODES FOR FIXED FLOOR CONFIGURATION WITH ADD－ON VINYL（Continued）

| Mude | FKEG | mude | Stiell |  |  | plate |  |  | GENEKALILEO MASS |  | （KG） |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Nu | （HI） | TYPE | M | N | CMN | M | N | CPN | YOTAL | SHELL | plate m |
| $y$ | 77.00 | SYMM | 3 | 3 | －． 00007 | 3 | 1 | － 99943 | 74.43382 | .00000 | 74.43382 |
|  |  |  | 3 | 2 | －．U0v90 | 3 | 0 | －CCUO6 |  |  |  |
|  |  |  | 3 | 4 | ．U0001 | 3 | 2 | －CCOO1 |  |  |  |
|  |  |  | 3 | 1 | －U00．j |  |  |  |  |  |  |
|  |  |  | 3 | 6 | －．víojo |  |  |  |  |  |  |
| 10 | 41.15 | SYAM | 3 | 3 | ．00873 | $\begin{aligned} & 3 \\ & 3 \\ & 3 \end{aligned}$ | $\begin{aligned} & 0 \\ & 2 \\ & 4 \end{aligned}$ | $\begin{array}{r} -.51053 \\ -.13454 \\ .00777 \end{array}$ | 125．19129 | 71.26489 | 40.11496 |
|  |  |  | 3 | 2 | －5uつつ4 |  |  |  |  |  |  |
|  |  |  | 3 | 4 | －．U0500 |  |  |  |  |  |  |
|  |  |  | 3 | 1 | －．02409 |  |  |  |  |  |  |
|  |  |  | 3 | 6 | －0し4y8 |  |  |  |  |  |  |
| 12 | 95.44 | SYMA | 2 | 1 | ． 34023 | $\begin{aligned} & 2 \\ & 2 \\ & 2 \end{aligned}$ | $\begin{aligned} & 0 \\ & 2 \\ & 4 \end{aligned}$ | $\begin{array}{r} .28171 \\ \cdot 15249 \\ -. C 817 \end{array}$ | 76.52660 | 33.74691 | 23.40918 |
|  |  |  | 2 | 2 | ． 10020 |  |  |  |  |  |  |
|  |  |  | 2 | 3 | －＜ 4714 |  |  |  |  |  |  |
|  |  |  | 2 | 4 | $.07544$ |  |  |  |  |  |  |
|  |  |  | 2 | 4 | $\text { -. } 00689$ |  |  |  |  |  |  |
| 12 | 98.52 | ANTI | 3 | 2 | －．5574．3 | $\begin{aligned} & 3 \\ & 3 \\ & 3 \end{aligned}$ |  | $\begin{array}{r} -.42017 \\ -.00035 \\ .00081 \end{array}$ | 95.70690 | 68．54229 | 13.14369 |
|  |  |  | 3 | 3 | ． 54305 |  |  |  |  |  |  |
|  |  |  | 3 | 4 | ． 00168 |  |  |  |  |  |  |
|  |  |  | 3 |  | －．00836 |  |  |  |  |  |  |
|  |  | － | 3 |  | －． 00004 |  |  |  |  |  |  |
| 23 | 99.61 | ANTI | 2 | 3 | .61603 |  | $\begin{aligned} & 1 \\ & 3 \\ & 5 \end{aligned}$ | $\begin{array}{r} -.05609 \\ =.00144 \\ .00018 \end{array}$ | 64.33322 | 53.75372 | .23428 |
|  |  |  | 2 | 4 | ．1401J | 2 |  |  |  |  |  |
|  |  |  | 2 | 1 | .14105 |  |  |  |  |  |  |
|  |  |  | 2 | 2 | .14126 |  |  |  |  |  |  |
|  |  |  | 2 | 6 | －． 02345 |  |  |  |  |  |  |
| 14 | 107.51 | ANTI | 1 | 3 | －．01021 | 1 | $\begin{aligned} & 1-.01054 \\ & 3-.02909 \\ & 5 \quad .00362 \end{aligned}$ |  | 342．56．263 | 59.14559 | 48.96392 |
|  |  |  | 1 | 1 | ． 32645 | 1 |  |  |  |  |  |  |
|  |  |  | 1 | 4 | －．14174 | 1 |  |  |  |  |  |  |
|  |  |  | 1 |  | $.0<445$ |  |  |  |  |  |  |  |
|  |  |  | 1 |  | －．01033 |  |  |  |  |  |  |  |
| 15 | 107.79 | SYMM | 4 |  | －．ucuus |  |  |  | 74.43382 | .00000 | 74．43382 |
|  |  |  | 4 | 4 | －．U0Uu3 | 4 | C | $\text { - } \mathrm{CcOO} 4$ |  |  |  |
|  |  |  | 4 | 4 | －U0002 | 4 |  | － CCOOL |  |  |  |
|  |  |  | 4 | 1 | ．0c001 |  |  |  |  |  |  |
|  |  |  | 4 |  | －．00000 |  |  |  |  |  |  |
| 10 | 215.48 | ANTI | 1 | 4 | ． 55130 | 1 | 1 | －． 55164 | 120.37357 | 74.54285 | 9.22010 |
|  |  |  | 1 | 3 | ． 42255 | 1 | 3 | －． 01540 |  |  |  |
|  |  |  | 1 | 1 | －． 31304 | 1 | 5．CO140 |  |  |  |  |
|  |  |  | 1 |  | ． 20104 |  |  |  |  |  |  |  |
|  |  |  | 1 |  | －．0450\％ |  |  |  |  |  |  |  |

CALAC METRO. UIAMETEK L. GBM, LENGTH GM, FLUATIHG FLUIJK MODEL + ADDGBN VIIML

```
MUDE NUMBER 1
    FREQUENCY= 32.1<U3 SYMMETKIC MIIUES M = l
```



FIGURE 8. STRUCTURAL MODE SHAPES: FLOATING FLOOR CONFIGURATION WITH ADD-ON VINYL

CALAC METKU. UIAMETER L.BdM, LENGTH YM. FLJGTING FLUUK MUOEL + AUU-ON VINYL

```
MJUE NUMBEK }
FREOUENCY:44.1579 SYMMETRIC MOOES M= 1
```



FIGURE 8. CONTINUED



FIGURE 8. CONTINUED


FIGURE 9. MEASURED CIRCUMFERENTIAL MODE SHAPE AT 83 Hz FOR BARE FUSELAGE STRUCTURE [5]

This mode was assumed to be symmetric so that the displacements normal to the fuselage could be averaged for the 2 sides and then normalized to a maximum value of 1.0 . Figure 10 compares the measured mode shape with that calculated for the fuselage shell only (no floor), for $M=1$ and $n=3$. Figures 11, 12 and 13 compare the measured mode shape and frequency with corresponding calculated values for the floating floor, rigid joint model, and the attached floor, rigid and hinged joint models. The comparison suggests that the floating floor model puts too large a constraint on the fuselage below the floor line. However, because of the uncertainty regarding the identification of the order of the measured mode, it is difficult to make any definitive conclusions. The uncertainty in measured mode order is one factor influencing the agreement between predicted and measured resonance frequencies. It is interesting to note that the measured frequency ( 40 Hz ) of the lowest order mode of the structure is close to the predicted values of 35 Hz for the fixed floor and 37 Hz for the floating floor.

### 4.5 Structural Loss Factors

Measured structural loss factors for the baseline (bare) test structure are listed in Table 2. The data show large variations from frequency to frequency, and different values were obtained using different test methods. Thus, some averaging, smoothing and interpolation was performed in one-third octave frequency bands in order to obtain input data for the model. The resulting average loss factors are given in Table 9. Also, the average loss factor curve is shown superimposed on the test data in Fig. 14. The average loss factors for the increased density test structure are also shown in Table 9.

a) Measured 83 Hz
b) Calculated 120 Hz

FIGURE 10. COMPARISON OF MEASURED MODE SHAPE WITH FUSELAGE SHELL MODEL (NO FLOOR)


FIGURE 11. COMPARISON OF MEASURED MODE SHAPE WITH FLOATING FLOOR MODEL (RIGID JOINT)

a) Measured 83 Hz
b) Calculated 150 Hz

FIGURE 11. CONTINUED

a) Measured 83 Hz
b) Calculated 88 Hz

FIGURE 12. COMPARISON OF MEASURED MODE SHAPE WITH ATTACHED FLOOR MODEL (RIGID JOINT)

a) Measured 83 Hz
B) Calculated, 88 Hz

FIGURE 13. COMPARISON OF MEASURED MODE SHAPE WITH ATTACHED FLOOR MODEL (HINGED JOINT)

TABLE 9

AVERAGE STRUCTURAL LOSS FACTORS FOR FUSELAGE

| Frequency (Hz) | Average Loss Factor |  |
| :---: | :---: | :---: |
|  | Baseline | Increased Density |
|  |  |  |
| 60 | 0.0771 | 0.0833 |
| 80 | 0.0725 | 0.0833 |
| 100 | 0.0679 | 0.1000 |
| 125 | 0.0390 | 0.1185 |
| 160 | 0.0218 | 0.0920 |
| 200 | 0.0306 | 0.0670 |
| 250 | 0.0245 | 0.0510 |
| 315 | 0.0212 | 0.0379 |
| 400 | 0.0141 | 0.0515 |
| 630 | 0.0224 | 0.0347 |
| 800 | 0.0210 | 0.0720 |
| 1000 | 0.0102 | 0.0266 |
| 1250 | 0.0122 | 0.0213 |
|  | 0.0148 | 0.0373 |
|  |  |  |
|  |  |  |

$\triangle$ MEASURED- PANEL GRID

- MEASURED- CIRCUMFERENTIAL GRID
- averrge $1 / 3$ gb mersured values


FIGURE 14. STRUCTURAL LOSS FACTORS FOR BASELINE STRUCTURE

### 4.6 Floor Structural Modes used in the Analytical Model

The floating floor configuration introduces an inconsistency in the analytical representation of the coupling between floor and cavity modes. The acoustic modes in the cavity are determined by the physical presence of the plywood floor, which is located 28 cm (11 inches) below the fuselage centerline. In contrast the weight of the plywood floor is supported on fuselage structure which is below the test floor line. Thus, for structural modeling purposes, the floor line has been taken at a distance of 68.6 cm ( 27 inches) below the fuselage centerline. The inconsistency arises because the model calculates the coupling factor between the structural and acoustic modes and, therefore, the physical dimensions for the two sets of modes must correspond.

To overcome this inconsistency, it was assumed that the floor mode shape, calculated for the floating floor at a lower position on the fuselage structure, could be applied to the actual plywood floor. The transferred mode shapes were assumed to have the same frequencies and maximum deflections, but the corresponding modal wavelengths were increased in direct proportion to the increases in the effective width of the floor. The structural/acoustic coupling factor was then calculated for the assumed mode shapes for the plywood floor and for the fuselage shell above the plywood floor only. The actual generalized mass for the floating floor was used in calculating the response.

### 4.7 Sidewall Treatment

The analytical representation for the sidewall treatment is discussed in Section 2.5, and the required parameters required for the insulation are identified in Eq. (15). Values for these parameters used in the analysis of the Metro II test fuselage are given in Table 10 for the fiberglass material and the airgap.

trim insulation typeAIR GAP

| FREQUENCY | ALPHA <br> DZ | LAMBDA | MUD(W) | PHASE(W) |
| ---: | ---: | ---: | ---: | ---: |
|  |  |  | MKS RAYLS | DEG |

The sidewall model was not used to calculate the acoustic loss factors of the modes in the cavity, but it was used to calculate the transmission loss through the sidewall and the effective damping of the sidewall on the structure [1, Appendix A].

The mass of the interior trim lining was changed for the different sidewall configurations, as shown in Table 11. These masses can be used, in a simplified analysis, to estimate sidewall resonance frequencies based on a mass-spring-mass model,

$$
f_{d w}=\frac{1}{2 \pi} \sqrt{\frac{\rho c^{2}\left(m_{1}+m_{2}\right)}{d m_{1} m_{2}}}
$$

where $d$ is the distance between the panels of surface density $\mathrm{m}_{1}$ and $\mathrm{m}_{2}$. The resulting frequencies are listed in Table 11. The mechanics of the analytical model for sidewall transmission are, however, somewhat different in that the model utilizes the trim mass in the trim.transfer matrix (Eq. (16)) and, thence, in the trim transmission coefficient (Eq. (A.6)[1]) and the damping of the outer panel (Eq. (17)).

The loss factor associated with the lining was taken as 0.50 in all cases. This value might be high for the sidewall configuration without add-on vinyl but, as the mass-spring-mass double wall resonances lie below 200 Hz , the effect of lining loss factor will not be too important.

The lowest resonance frequencies of the inner trim panel were calculated for the test section, as a 3.66 m long (144 in) curved panel extending from floor to floor. Because of the strong curvature, the lowest frequencies were found to be about 50 Hz for the modes with one half-wavelength in the axial direction and 10 to 15 half-wavelengths from floor to floor. When the effect of the fundamental resonance frequency of the wall
was included in the transfer matrix, the effective mass of the inner wall was reduced and the double-wall resonance frequency was raised slightly. However, above 200 Hz , the effect on the transmission loss of the sidewall was small; thus the inner wall stiffness was not included in the final computations.

TABLE 11. SIDEWALL RESONANCE FREQUENCIES

| Configuration <br> No. | Surface density, $\mathrm{kg} / \mathrm{m}^{2}$ |  | Double-Wall <br> Resonance, Hz |
| :---: | :---: | :---: | :---: |
|  |  |  |  |
| 1 | 4.59 | 2.25 | 193.5 |
| 2 | 4.59 | 7.13 | 142.3 |
| 3 | 4.59 | 12.01 | 130.4 |
| 4 | 9.47 | 2.25 | 176.3 |
| 5 | 9.47 | 7.13 | 12.01 |

### 4.8 Cavity Model

### 4.8.1 Cavity Modes

Resonance frequencies and mode shapes for the acoustic modes inside the fuselage were computed using the finite difference method developed for the analytical model. Examples of the predicted modes are shown in Figure 15 for modes with zero axial half-wavelengths ( $q=0$ ) in the cavity for the first 12 circumferential modes ( $i=0$ to 11).

A total of 400 acoustic modes were included in the analysis, for $q=0$ to 19 axially, and $i=0$ to 19 circumferentially, where

```
q = no of axial half-wavelengths
i = assigned order of 2-dimensional modal pattern
        in cylinder cross section.
```

Because of the length of the cavity, $9.02 \mathrm{~m}(29.6 \mathrm{ft})$, the 400 acoustic modes cover only the frequency range 19 Hz to 670 Hz . Calculations made using the low frequency modal model (individual acoustic modes) will therefore not be valid above the 500 Hz one-third octave band.

### 4.8.2 Acoustic Loss Factors

Acoustic absorption coefficients were measured by Prydz et al [5] for the bare and treated interiors of the test fuselage. Two methods were tried, one being the use of a standard sound source of known acoustic output power and the other being the sound level decay method. The decay method was used for the bare interior, and both methods were applied to the treated cabin. Resulting data are presented in Table 4 and Fig. 7 where it is seen that the two methods can, at some frequencies, give quite different values.

|  | B＊ | $1^{\circ}$ | $5{ }^{\circ}$ | $\varepsilon \cdot$ | $1 \cdot$ | 0．－ | 1－－ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $1 \cdot$ | $1 *$ | $9 \cdot$ | ＊＊ | $2 \cdot$ | 1 | 10－ | 1：－ |
| ＊＊ | ＊＊ | E＊ | $2 \cdot$ | 0 ＊ | $1 \cdot$ | で－ | 2＊－ |
| $1 *$ | $0 *$ | 0＊－ | 1＊－ | 2＊－ | 2＊－ | 2＊－ | 20－ |
| c－ | $\varepsilon \cdot-$ | ＊－ | ＊＊－ | ＋＊－ | ${ }^{\circ}-$ | 2＊－ | 2＊－ |
|  | 9＊－ | 9＊－ | 9＊－ | ＊－ | 2＊－ | 1＊－ | 0．－ |
|  | 2－ | 2－ | 9＊－ | c＊－ | O－ | 2 | E |
|  |  | 9．－ | ＊－－ | $1 *-$ | $2 *$ | s＊ | $9 \cdot$ |
|  |  |  | c＊－ | $1 *$ | ${ }^{\circ}$ | $0 *$ | $6^{\circ}$ |
|  |  |  |  |  | 2＊ | $0^{\circ}$ | $0 \cdot 1$ |
|  |  | 入วขแ | $K_{S}$ | $\mathrm{zH}^{\text {H }}$ | 282 | S ap | O |
|  | 2－ | 10－ | $9 \cdot-$ | $9{ }^{\circ}-$ | 5＊－ | s＊－ | 5＊－ |
| 2－－ | 2－－ | $9 \cdot-$ | 9＊－ |  | 5＊－ | 5＊－ | ＊＊－ |
| 9＊－ | 9＊－ | s＊－ | 5＊－ | ＊＊ | ¢＊－ | E＊－ | ［＊－ |
| 5＊－ | ＊＊－ | ＊＊ | ＊－ | で－ | 2＊－ | 10 | 1＊－ |
| E＊－ | 2＊－ | 2＊－ | 1＂－ | $0 \%$ | $1 *$ | ＊＊ | $1 \cdot$ |
|  | 1－－ | $0 \cdot$ | ＊＊ | 2＊ | E | ＊＊ | $\stackrel{*}{ }$ |
|  | $2 \times$ | $2^{*}$ | ¢＊ | ¢＊ | s＊ | $9{ }^{\circ}$ | $9 \bullet$ |
|  |  | $\cdots$ | $5{ }^{\circ}$ | $9 *$ | ${ }^{\circ}$ | e＊ | $日^{*}$ |
|  |  |  | $9 *$ | $2^{\circ}$ | $6^{*}$ | $0 \cdot$ | $0^{\circ} 1$ |
|  |  |  |  |  | $6{ }^{\circ}$ | $0 \cdot 1$ | $0 \cdot 1$ |
|  |  | 17 วur | $\kappa_{S}$ | 2 H | 191 | 乙 əp | W |




|  | $0 \cdot 1$ | $6{ }^{\circ}$ | $日^{*}$ | $1 *$ | c ${ }^{\text {c }}$ | $2^{*}$ | $0^{\circ}-$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.1 | 0.1 | $6^{*}$ | ＊${ }^{\text {－}}$ | $9 \cdot$ | 5 | 2＊ | 0－－ |
| $0 \cdot 1$ | $0 \cdot 1$ | $6 \cdot$ | A | $9 *$ | $\cdots$ | 2 | 0－－ |
| $0 \cdot 1$ | $0 \cdot 1$ | $6^{\circ}$ | ค． | $9 \cdot$ | $* *$ | $2 \cdot$ | 0－－ |
| $0 *$ | $6^{*}$ | $6^{\circ}$ | $\iota^{\circ}$ | $9 \cdot$ | ＊＊ | $2 \cdot$ | 0＊－ |
|  | t＊ | ${ }^{*}$ | $2^{\bullet}$ | $9 \cdot$ | $\cdots$ | $2 \cdot$ | $0 \cdot$ |
|  | $8 \cdot$ | 2＊ | $1 \cdot$ | $5 \cdot$ | $\cdots$ | 7＊ | 0＊－ |
|  |  | ${ }^{\circ}$ | $9{ }^{-}$ | s• | ¢ | 2 | 0＊－ |
|  |  |  | $5^{\circ}$ | $\stackrel{\bullet}{ }$ | ${ }^{*}$ | ${ }^{*}$ | $0 \cdot$ |
|  |  |  |  |  | $\varepsilon \cdot$ | $1{ }^{*}$ | 00 |


|  | 5＊－ | \＆${ }^{-}$ | $0^{\circ}-$ | $\varepsilon^{\circ}$ | $2^{\circ}$ | $6 *$ | $0 \cdot 1$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 9＊－ | 90－ | ＊＊＊ | 1＊－ | $\varepsilon^{\circ}$ | 9 • | 6. | $0 \cdot 1$ |
| （＊－ | 9＊－ | ＊＊－ | 1\％－ | 20 | s＊ | $8 \cdot$ | $6{ }^{\circ}$ |
| 8＊－ | 1＊－ | 5＊－ | 20－ | $1 *$ | ＊＊ | $9 \cdot$ | 2＊ |
| 8＇－ | 10－ | 9－－ | \＆${ }^{-}$ | 0＊－ | $\varepsilon^{\bullet}$ | $5 \cdot$ | $5{ }^{\circ}$ |
|  | 2＊－ | －＊－ | ＊＊ | 1＊－ | $2 *$ | $\varepsilon$＊ | $4 *$ |
|  | 9＊－ | 5＊－ | ＋＊－ | 1＊－ | $1 *$ | 2• | $6^{*}$ |
|  |  | 5．－ | \＆＊－ | 2＊－ | 0 ＊ | $1 \cdot$ | $2{ }^{\circ}$ |
|  |  |  | ［＊－ | 1＊－ | 0 ＊ | $1{ }^{\text {－}}$ | $1 *$ |
|  |  |  |  |  | $0^{\circ}$ | $1{ }^{\circ}$ | 1 ＊ |


$0.100 .10 .100 .1 \quad 0.100 .10 .1$ $\begin{array}{llllllll}0.1 & 0.1 & 0.1 & 0.1 & 0.1 & 0.1 & 0.1 & 0.1\end{array}$ $0.10 .10 .1 \quad 0.1 \quad 0.1 \quad 0.1 \quad 0.1 \quad 0.1$ 0.10 .10 .10 .10 .10 .10 .10 .1 0.100 .10 .10 .10 .100 .100 .100 .1
$0.10 .1 \quad 0.1 \quad 0.1 \quad 0.1 \quad 0.180 .1$
$0.10 .1 \quad 0.1 \quad 0.1 \quad 0.120 .100 .1$ $0.10 .1 \quad 0.1 \quad 0.1 \quad 0.1 \quad 0.1$ $\begin{array}{lllll}0.1 & 0.1 & 0.1 & 0.1 & 0.1\end{array}$
$0.1 \quad 0.1 \quad 0.1$



Data from Table 4 and Fig. 7 were combined to provide a mean absorption coefficient spectrum over the frequency range of interest. The absorption coefficient $\alpha$ was then converted to the associated acoustic loss factor $\eta_{n}$ by means of the relationship

$$
\eta_{n}=\frac{\alpha c_{0} s}{4 \omega V}
$$

where $V$ is the volume of the cavity $(14.08 \mathrm{~m} 3)$ and $S$ the transmitting surface area ( $28.9 \mathrm{~m}^{2}$ ). Values for the mean absorption coefficient and acoustic loss factor are given in Table l2. These empirical loss factor data were used as input to the computation process instead of using the analytical model itself to calculate acoustic loss factors on the basis of the dynamic characteristics of the trim panels. This approach of inserting empirical acoustic loss factor data as data input was used because much of the acoustic dissipation within the test fuselage was provided by the foam placed along the edges of the floor, rather than by a uniform distribution over the sidewall treatment. The analytical model was not designed to cater for the test situation.

It should be noted that the value of area $S$ used to compute $\alpha$ was not the total surface area in the cavity. The value was associated with the curved, transmitting area since this was the area used by Prydz et al [5] to calculate average values for the absorption coefficient $\alpha$.

TABLE 12.
ACOUSTIC ABSORPTION COEFFICIENTS AND LOSS FACTORS FOR THE INTERIOR OF THE TREATED TEST FUSELAGE

| Frequency (Hz) | Absorption Coefficient | Loss Factor |
| :---: | :--- | :--- |
|  |  |  |
| 60 | 0.005 | 0.00280 |
| 80 | 0.01 | 0.00445 |
| 100 | 0.015 | 0.00525 |
| 125 | 0.02 | 0.00560 |
| 160 | 0.03 | 0.00672 |
| 200 | 0.08 | 0.01401 |
| 250 | 0.285 | 0.03992 |
| 315 | 0.355 | 0.03978 |
| 400 | 0.435 | 0.03868 |
| 500 | 0.305 | 0.02136 |
| 630 | 0.415 | 0.02325 |
| 800 | 0.34 | 0.01512 |
| 1000 | 0.30 | 0.01050 |
| 1250 | 0.295 | 0.00826 |
| 1600 | 0.31 | 0.00695 |
| 2000 | 0.32 | 0.00560 |
|  |  |  |

### 5.0 ANALYTICAL REPRESENTATION OF METRO II TEST EXCITATION

The analytical model requires that the exterior acoustic pressure field be described at a series of points on a grid covering the sidewall of the cylinder. The description of the field is given in terms of the free field pressure amplitude and phase for a particular frequency of interest, the model having been developed for harmonics of a propeller blade passage frequency. The grid points ( $k, \ell$ ) used for defining the pressure field are shown in Figure 16 , the grid spacing being 0.1463 m ( 5.76 in ). In the case of the Metro II tests, the point $(8,1)$ was 1.3 m from the electropneumatic noise source; for application in this model, the pressure field was described over a region of the fuselage approximately 1 meter forward and aft of the noise source.

Free field test data were provided for the Metro II along the longitudinal axis of the grid point array and "blocked" pressures along the circumferential direction. These data are shown in Figure 5. Since. the model requires that the acoustic pressures be free field, the circumferential pressure distribution was converted to an equivalent free-field pattern using the inverse of the equation contained in the computer program (see Eq. (43) of [1]).

$$
\begin{equation*}
p_{b l k}=\left\{10\left[0.3-0.000224 \mathrm{e}^{0.08 \gamma}\right]\right\} p_{\text {free }} \tag{48}
\end{equation*}
$$

where $\gamma$ is the "incidence angle" in degrees, which in this case is the angle between a line connecting the center of the horn to the grid point on the fuselage and the normal to the surface at that point, (see Figure 17). The calculated reflection effects in the vertical plane containing the noise source axis are given in Table 13. Furthermore, the three distribution patterns shown in Fig. 5(b) for three frequencies were normalized relative to peak values, and an average spatial distribution was calculated. This was then combined with the reflection effect to give an estimated circumferential free field distribution in Table 13.


$\begin{aligned} \bar{x}_{m}=(a, \theta, z)= & \underset{ }{\text { Location of Free }} \\ & \text { Point Prediction } \\ \left(r_{p}, \phi, z_{p}\right) & =\text { Propeller Position }\end{aligned}$

FIGURE 17. PROPELLER AND FUSELAGE SURFACE POINT GEOMETRY

## TABLE 13 <br> CIRCUMFERENTIAL VARIATION OF FREE FIELD PRESSURE

| Grid Point | Reflection <br> effect <br> $(\mathrm{dB})$ | Average Measured <br> spatial variation <br> (dB re max) | Estimated <br> free field <br> variation <br> (dB re max) |  |
| :---: | :---: | :---: | :---: | :---: |
| $k$ | $\ell$ | 0 | -18.9 | -12.9 |
| 8 | 10 | 0 | -17.1 | -11.1 |
| 8 | 9 | 0 | -14.8 | $-10.0 *$ |
| 8 | 8 | 2.6 | -12.9 | -9.5 |
| 8 | 7 | 4.6 | -9.9 | -8.5 |
| 8 | 6 | 5.5 | -6.8 | -6.3 |
| 8 | 5 | 5.8 | -4.0 | -3.8 |
| 8 | 4 | 5.9 | -1.3 | -1.2 |
| 8 | 3 | 6.0 | -0.2 | -0.2 |
| 8 | 2 | 6.0 | -0.0 | 0 |
| 8 | 1 |  |  |  |

TABLE 14
AXIAL VARIATION OF FREE FIELD PRESSURE

| Grid Point |  | Axial Distance <br> From Source Plane <br> $(\mathrm{m})$ | Free Field <br> Variation <br> (dB re max) |
| :---: | :---: | :---: | :---: |
| $k$ | $\ell$ | 0 | 0 |
| 8 | 1 | .146 | -.45 |
| 9 | 1 | .293 | -1.30 |
| 10 | 1 | .439 | -2.65 |
| 11 | 1 | .585 | -4.30 |
| 12 | 1 | .731 | -6.40 |
| 13 | 1 | .878 | -8.35 |
| 14 | 1 | 1.024 | -10.70 |
| 15 | 1 | 1.170 | -12.30 |
| 16 | 1 |  |  |

The longitudinal free field distribution from Fig. 5(a) is given in Table 14. It was assumed that the maximum pressure at any longitudinal station was defined by Fig. 5(a), along the grid line $(\ell=1)$. It was also assumed that the circumferential free field distribution could be applied at any longitudinal station, so that the pressure amplitude at all grid points could be calculated. A maximum sound pressure level of 134 dB was assumed for the array as a whole. The actual value of the maximum level was not critical to the analysis because only noise reduction was being calculated. The 134 dB level was chosen as a typical value used in the test program.

Phase angle data were not available from the Metro II test. Consequently, an analytical model had to be constructed for relative phase angle at each frequency of interest. This was accomplished by assuming that the sound field consisted of spherical waves originating in the horn. The distance from horn to the grid fuselage structure was calculated for each grid point and converted to relative phase using the appropriate values for the speed of sound and frequency. In this case the relative phase angle is the same for similar grid points above and below the grid longitudinal centerline; the two phase angles would have different values for propeller noise because of propeller rotation.

The "blocked" pressure amplitude and phase calculated for the frequencies 383,766 and 1149 Hz are shown in Table 15. It can be seen that the amplitude falls off 11 - 12 dB over the axial distance covered by the grid points.
PROPELLER HARMO:NIC 1 AT $3 H 3.0112$


PRESSURE AMPLITUDE, IHR RE 20 MICRO PA

| 19 | 180.00 | 110.4 | 112.8 | 114.7 | $11 \% .8$ | 118.5 | 11\%.R | 120.6 | 121.1 | 120.6 | 119.8 | 118.5 | 116.B | 114.7 | 112.8 | 110.4 | 108.8 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 18 | 170.00 | 112.2 | 114.5 | 119.5 | $11 \% .6$ | 12 U .3 | 121.6 | 122.4 | 122.9 | 122.4 | 121.6 | 120.3 | 118.6 | 116.5 | 114.5 | 112.2 | 110.6 |
| 17 | 16\%.00 | .113 .1 | 115.7 | 117.5 | 117.7 | 121.4 | 122. 1 | 121.5 | 124.0 | 123.5 | 122.7 | 121.4 | 119.7 | 117.6 | 115.7 | 113.3 | 111.7 |
| 16 | 159.00 | 116.2 | 118.6 | 120.6 | 122.7 | 1<4.4 | 125.8 | $1<6.6$ | 127.1 | 126.6 | 125.8 | 124.4 | 122.7 | 120.6 | 118.6 | 116.2 | 114.5 |
| 15 | 140.00 | 112.1 | 121.5 | 123.5 | 123.7 | 127.4 | 12H.H | $1<4.6$ | 130.1 | 129.6 | 128.8 | 127.4 | 125.7 | 123.5 | 121.5 | 114.1 | 117.4 |
| 14 | 130.00 | 122.2 | 124.6 | 125.6 | 128.8 | 130.5 | 131.8 | 132.7 | 133.2 | 132.7 | 131.8 | 130.5 | 128.8 | 126.6 | 124.6 | 122.2 | 120.5 |
| 13 | 120.00 | 125.1 | 127.5 | 129.5 | 131.7 | 133.3 | 134.7 | 135.6 | 136.0 | 135.6 | 134.7 | 133.3 | 131.7 | 129.5 | 127.5 | 125.1 | 123.5 |
| 12 | 110.00 | 127.? | 130.3 | 132.3 | 134.4 | 136.L | 131.4 | 138.3 | $13 \mathrm{H}$. ? | 138.3 | 137.4 | 136.1 | 134.4 | 132.3 | 130.3 | 127.9 | 126.3 |
| 11 | 100.00 | 127.0 | 131.4 | 137.3 | 135.5 | 177.1 | 173.5 | 134.3 | 139.8 | 139.3 | 138.5 | 137.1 | 135.5 | 133.3 | 131.4 | 129.0 | 127.3 |
| 10 | 10.00 | 129.2 | 111.6 | 133.6 | 135.7 | 137.3 | 138.7 | 134.5 | 140.0 | 139.5 | $13 \mathrm{H.7}$ | 137.3 | 135.7 | 133.6 | 131.6 | 129.2 | 127.6 |
| 9 | 80.00 70.00 | 129.0 | 131.4 | 133.3 | 135.5 | 137.1 | 13 H .5 | $139 \cdot 3$ | 139.8 | 139.3 | 138.5 | 137.1 | 135.5 | 133.3 | 131.4 | 129.0 | 127.3 |
| 7 | 70.00 60.00 | 127.9 125.1 | 130.3 127.5 | 132.3 124.5 | 134.4 131.7 | 136.1 133.3 | 137.4 134.7 | 138.3 135.6 | 138.7 | 138.3 | 137.4 | 136.1 | 134.4 | 132.3 | 130.3 | 127.9 | 126.3 |
| 6 | 60.00 50.00 | 125.1 122.2 | 127.5 124.5 | 129.5 126.6 | 131.7 128.8 | 133.3 130.5 | 134.7 131.8 | 135.6 132.7 | 136.0 133.2 | 135.6 | 134.7 | 133.3 | 131.7 | 129.5 | 127.5 | 125.1 | 123.5 |
| 5 | 40.00 | 119.1 | 121.5 | 123.5 | 125.7 | 127.4 |  | 129.6 | 130.1 | 129.6 | 128.8 | 127.4 | 125.7 | 123.5 | 121.5 | 119. | 120.5 117.4 |
| 4 | 30.00 | 115.2 | 118.6 | 120.6 | 122.7 | 124.4 | 125.8 | 126.6 | 127.1 | 126.6 | 125.8 | 124.4 | 122.7 | 120.6 | 118.6 | 116.2 | 114.4 |
| 3 | 20.00 | 113.3 | 115.7 | 117.6 | 119.7 | 121.4 | 122.7 | 123.5 | 124.0 | 123.5 | 122.7 | 121.4 | 119.7 | 117.6 | 115.7 | 113.3 | 111.7 |
| 2 | 10.00 | 112.2 | 114.5 | 116.5 | 118.6 | 120.3 | 121.6 | 122.4 | 122.9 | 122.4 | 121.6 | 120.3 | 118.6 | 116.5 | 114.5 | 112.2 | 110.6 |
| 1 | 0.00 | 110.4 | 112.8 | 114.7 | 116.8 | 118.5 | 114.8 | 120.6 | 121.1 | 120.6 | 119.8 | 118.5 | 116.8 | 114.7 | 112.8 | 110.4 | 108.8 |


| 19 | 180.00 | 128.2 | 105.8 | 86.3 | 70.1 | 37.3 | 48.1 | 42.5 | 40.6 | 42.5 | 40.1 | 57.3 | 70.1 | 86.3 | 105.8 |  | 6 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 18 | 170.00 | 77.2 | 53.5 | 32.7 | 15.8 | 2.3 | 352.4 | 34 t. 4 | 344.4 | 346.4 | 352.4 | 2.2 | 15.8 | 32.9 | 53.5 | 77.2 | 103.9 |
| 17 | 160.00 | 25.0 | 359.9 | 33H.0 | 114.9 | 305.3 | 214. 8 | 288.4 | 286.2 | 288.4 | 244.8 | 305.3 | 319.6 | 338.0 | 359.9 | 25.0 | 53.1 |
| 16 | 150.00 | 333.1 | 303.4 | 2日5.1 | 263.6 | 24.4 .0 | 236.7 | 229.9 | 227.6 | 229.9 | 236.7 | 248.0 | 263.6 | 283.1 | 306.4 | 333.1 | 2.9 |
| 15 | 141.100 | 283.3 | 254.9 | 23).1 | 203. 1 | 192.3 | 1411 | 172.7 | 170.2 | 172.7 | 180.1 | 192.3 | 209.1 | 230.1 | 254.9 | 283.3 | 314.4 |
| 14 | 130.00 | 237.7 | 207.6 | 141.0 | 154.5 | 140.5 | 127.4 | 119.3 | 116.5 | 119.3 | 127.4 | 140.5 | 158.5 | 181.0 | 207.6 | 237.7 | 271.1 |
| 13 | 170.00 | 198.6 | 136.8 | 139.6 | 114.6 | 45.3 | 81.2 | 72.5 | 69.6 | 12.5 | 81.2 | 95.3 | 114.6 | 138.6 | 166.8 | 198.6 | 233.7 |
| 12 | 110.00 | 168.4 | 135.1 | 103.5 | 90.2 | 59.8 | 44.7 | 35.5 | 32.4 | 35.5 | 44.7 | 59.8 | 80.2 | 105.5 | 135.1 | 168.4 | 205.0 |
| 11 | 100.00 | 144.2 | 114.9 | 84.3 | 58.1 | 36.9 | 21.2 | 11.6 | 8.4 | 11.6 | 21.2 | 36.9 | 58.1 | 84.3 | 114.9 | 149.2 | 186.8 |
| 10 | 90.00 | 142.7 | 108.0 | 77.0 | 50.5 | 24.0 | 13.1 | 3.3 | 0.0 | 3.3 | 13.1 | 29.0 | 50.5 | 77.0 | 108.0 | 142.7 | -179.4 |
| 9 | 80.100 | 147.2 | 114.7 | $154 \cdot 3$ | 53.1 | 36.7 | 31.2 | 11.6 | F. 4.4 | 11.6 | $\overline{21.2}$ | 36.9 | 58.1 | 84.3 | 114.9 | 149.2 | $-173.2$ |
| 8 | 70.70 | 168.4 | 135.1 | 105.5 | 80.2 | 54.8 | 44.7 | 35.5 | 32.4 | 35.5 | 44.7 | 59.8 | 80.2 | 105.5 | 135.1 | 168.4 | 155.0 |
| 7 | 60.00 | -161.4 | 163.8 | 138.6 | 114.6 | 75.3 | 81.2 | 72.5 | 67.6 | 72.5 | 81.2 | 95.3 | 114.6 | 138.6 | 166.8 | -161 | -126.3 |
| 6 | 50.00 | $-122.3$ | -152.4 | -179.0 | 158.5 | 140.5 | 127.4 | 119.3 | 116.6 | 119.3 | 127.4 | 140.5 | 158.5 | -179.0 | $-152.4$ | $-122.3$ | -88.9 |
| 5 | 40.00 | -76.7 | -105.1 | $-130.0$ | -150.7 | -167.7 | -179.9 | 172.7 | 170.2 | 172.7 | -179.9 | -167.7 | $-150.9$ | -130.0 | $-105.1$ | -70.7 | -45.1 |
| 4 | 30.00 | $-26.9$ | -53.6 | $-76.9$ | -76.4 | $-112.9$ | $-123.3$ | $-130.1$ | -132.4 | -130.1 | -123.3 | -112.0 | -96.4 | -76.9 | $-53.6$ | -26.9 | 2.9 |
| 3 | 20.00 | <5.0 | -.1 | $-27.0$ | -411. | $-1.4 .7$ | -65. | $-71.6$ | $-73.8$ | -71.6 | -65.2 | -54.7 | $-40.2$ | -22.0 | -. 1 | 25.0 | 53.1 |
| 2 | 10.00 | 77.2 | 53.5 | 32.7 | 15.8 | 2.? | -7.1 | $-13.6$ | $-15.6$ | -13.6 | $-7.6$ | 2.2 | 15.8 | 32.9 | 53.5 | 77.2 | 103.9 |
| 1. | 0.00 | 128.2 | 105.8 | 46.3 | 70.1 | 47.1 | $4 \mathrm{H.1}$ | 42.5 | 40.5 | 42.5 | 48.1 | 57.3 | 70.1 | 86.3 | 105.8 | 128.2 | 153.6 |

CIRCUM. AXIAL LOCATIUN
 PRESSURE AMPLITUOE, OB RE. 20 MICRO PA

| 19 | 180.00 | 110.4 | 112.8 | 114.7 | 115.8 | 118.5 | 119.8 | 120.6 | 121.1 | 120.6 | 119.8 | 118.5 | 116.8 | 114.7 | 112.8 | 110.4 | 108.8 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 18 | 170.00 | 112.2 | 114.5 | 116.5 | 118.6 | $120 \cdot 3$ | 121.6 | 122.4 | 122.9 | 122.4 | 121.6 | 120.3 | 118.6 | 116.5 | 114.5. | 112.2 | 110.6 |
| 17 | 160.00 | 111.1 | 115.1 | 117.6 | 114.7 | 121.4 | 122.7 | 123.5 | 124.0 | 123.5 | 122.7 | 121.4 | 114.7 | 117.6 | 115.7 | 115.3 | 111.7 |
| 16 | 150.00 | 115.2 | 118.6 | 120.6 | 122.7 | 124.4 | 123.8 | 126.6 | 127.1 | 126.6 | 125.8 | 124.4 | 122.7 | 120.6 | 118.6 | 116.2 | 114.5 |
| 15 | 140.00 | 117.1 | $1<1.5$ | 123.5 | 125.7 | 121.4 | 128.8 | 124.6 | 130.1 | $1 \angle 7.6$ | 12 B .8 | 127.4 | 125.7 | 123.5 | 121.5 | 114.1 | 117.4 |
| 14 | 130.00 | 122.2 | 124.6 | 126.6 | 124.8 | 130.5 | 131.8 | 132.7 | 133.2 | 112.7 | 131.8 | 130.5 | 128.8 | 126.6 | 124.6 | 122.2 | 120.5 |
| 13 | 120.00 | 125.1 | 127.5 | 129.5 | 131.7 | 133.3 | 134.7 | 135.6 | 136.0 | 135.6 | 134.7 | 133.3 | 131.7 | 129.5 | 127.5 | 125.1 | 123.5 |
| 12 | 110.00 | 127.9 | 130.3 | 132.3 | 114.4 | 136.1 | 137.4 | 138.3 | 138.7 | 138.3 | 137.4 | 136.1 | 134.4 | 132.3 | 130.3 | 127.9 | 126.3 |
| 11 | 100.00 | 127.0 | 131.4 | 133.3 | 115.3 | 137.1 | 134.5 | 134.3 | 137.8 | 139.3 | 138.5 | 137.1 | 135.5 | 133.3 | 131.4 | 129.0 | 127.3 |
| 10 | 90.00 | 129.2 | 131.6 | 133.6 | 135.7 | 137.3 | 138.7 | 139.5 | 140.0 | 139.5 | 138.7 | 137.3 | 135.7 | 133.6 | 131.6 | 129.2 | 127.6 |
| 9 | 80.00 | 129.0 | 131.4 | 133.3 | 135.5 | 137.1 | 138.5 | 139.3 | 139.8 | 139.3 | 138.5 | 137.1 | 135.5 | 133.3 | 131.4 | 129.0 | 127.3 |
| 8 | 70.00 | 127.9 | 130.3 | 132.3 | 134.4 | 136.1 | 137.4 | 138.3 | 134.7 | 138.3 | 137.4 | 136.1 | 134.4 | 132.3 | 130.3 | 127.9 | 126.3 |
| 7 | 60.00 | 125.1 | 127.5 | 179.5 | 131.7 | 133.3 | 134.7 | 135.6 | 136.0 | 135.6 | 134.7 | 133.3 | 131.7 | 129.5 | 127.5 | 125.1 | 123.5 |
| 6 | 50.00 | 122.2 | 124.6 | 126.6 | 128.8 | 130.5 | 131.8 | 132.7 | 133.2 | 132.7 | 131.8 | 130.5 | 128.8 | 126.6 | 124.6 | 122.2 | 120.5 |
| 5 | 40.00 | 119.1 | $1<1.5$ | 123.5 | 123.7 | 127.4 | 124.8 | 129.6 | 130.1 | 129.6 | 128.8 | 127.4 | 125.7 | 123.5 | 121.5 | 119.1 | 117.4 |
| 4 | 30.00 | 116.2 | 118.6 | 120.6 | 122.7 | 124.4 | 125.8 | 126.6 | 127.1 | 126.6 | 125.8 | 124.4 | 122.7 | 120.6 | 118.6 | 116.2 | 114.5 |
| 3 | 20.00 | 113.3 | 115.7 | 117.6 | 119.7 | 121.4 | 122.7 | 123.5 | 124.0 | 123.5 | 122.7 | 121.4 | 119.7 | 117.6 | 115.7 | 113.3 | 111.7 |
| 2 | 10.00 | 112.2 | 114.5 | 116.5 | 118.6 | 120.3 | 121.6 | 122.4 | 122.3 | 122.4 | 121.6 | 120.3 | 118.6 | 116.5 | 114.5 | 112.2 | 110.6 |
| 1 | 0.00 | 110.4 | 112.8 | 114.7 | 116.H | 114.5 | 119.8 | 120.6 | 121.1 | 120.6 | 119.8 | 118.5 | 116.8 | 114.7 | 112.8 | 110.4 | 108.8 |
| PHASE (DEGREES) |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |  |


| 19 | 190.00 | 256.5 | 211.5 | 172.6 | 140.2 | 114.7 | 96.2 | 85.0 | 81.3 | 85.0 | 96.2 | 114.7 | 140.2 | 172.6 | 211.5 | 256.5 | 307.2 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 18 | 170.00 | 154.4 | 107.0 | 135.9 | 31.6 | 4.4 | 344.8 | 332.9 | 328.1 | 332.9 | 344.8 | 4.4 | 31.6 | 65.9 | 107.0 | 154.4 | 207.7 |
| 17 | 180.00 | 50.0 | 359.7 | 316.1 | 279.5 | 250.5 | 227.5 | 216.8 | 212.5 | 216.8 | 229.5 | 250.5 | 279.5 | 316.1 | 359.7 | 50.0 | 106.2 |
| 16 | 150.00 | 306.2 | 252.8 | 206.2 | 167.1 | 136.0 | 113.4 | 97.7 | 75.1 | 99.7 | 113.4 | 136.0 | 167.1 | 206.2 | $252.8{ }^{\circ}$ | 306.2 | 5.8 |
| 15 | 140.00 | 206.t | 14\%.4 | 100.1 | 5\%.2 | 24.1 | . 3 | 345.5 | 340.5 | 345.3 | - 3 | 24.7 | 58.2 | 100.1 | 149.8 | 206.6 | 269.7 |
| 14 | 130.00 | 115.4 | 55.1 | 2.0 | 317.1 | 281.1 | 254.7 | 238.7 | 233.3 | 238.7 | 254.7 | 281.1 | 317.1 | 2.0 | 55.1 | 115.4 | 182.2 |
| 13 | 120.00 | 37.3 | 333.5 | 277.2 | 229.2 | 190.7 | 162.4 | 145.0 | 139.2 | 145.0. | 162.4 | 190.7 | 229.2 | 277.2 | 333.5 | 37.3 | 107.5 |
| 12 | 110.00 | 336.18 | 270.2 | 211.0 | 160.4 | 117.6 | 89.5 | 71.0 | 64.8 | 11.0 | 89.5 | 119.6 | 160.4 | 211.0 | 270.2 | 336.8 | 50.0 |
| 11 | 100.00 | 298.5 | 229.8 | 164.7 | 116.3 | 73.8 | 42.4 | 23.2 | 16.7 | 23.2 | 42.4 | 73.8 | 116.3 | 168.7 | 229.8 | 298.5 | 13.6 |
| 10 | 90.00 | -14.7 | -144.1 | 154.1 | 101.0 | 513.9 | 2t. 1 | 6.6 | 0.0 | 6.5 | 26.1 | 58.0 | 101.0 | 154.1 | -144.1 | -74.7 | 1.1 |
| 9 | 30.00 | -61.3 | -130.2 | 168.7 | 116.3 | 13.8 | 42.4 | 23.2 | 16.7 | 23.2 | 42.4 | 73.8 | 116.3 | 168.7 | -130.2 | -61.5 | 13.6 |
| 8 | 70.00 | $-23.2$ | -83.8 | -149.0 | 150.4 | 117.6 | 87.5 | 71.0 | 64.8 | 11.0 | 89.5 | 114.6 | 160.4 | -144.0 | -89.8 | $-23.2$ | 50.0 |
| 7 | 60.00 | 17.3 | $-25.5$ | $-32.8$ | $-139.6$ | $-169.1$ | 162.4 | 145.0 | 139.2 | 145.0 | 152.4 | -169.3 | -130.8 | -82.8 | $-26.5$ | 37.3 | 107.5 |
| 6 | 30.00 | 115.4 | 55.1 | 2.0 | $-42.9$ | -14.4 | -105.3 | -121.3 | $-126.7$ | -121.3 | -105.3 | -78.9 | -42.4 | 2.0 | 55.1 | 115.4 | $-177.8$ |
| 5 | 40.00 | -153.4 | 149.4 | 100.1 | 3H.2 | 24.7 | . 3 | -14.5 | $-17.5$ | $-14.5$ | - 3 | 24.7 | 58.2 | 100.1 | 149.8 | -153.4 | $-40.3$ |
| 4 | 39.00 | $-33.4$ | -107. 2 | $-153.8$ | 16.7 .1 | 136.0 | 113.4 | 97.7 | 95.1 | 4.7 .7 | 113.4 | 136.0 | 167.1 | -153.8 | -107.2 | -53.8 | 5.8 |
| 3 | 29.011 | 30.0 | -. 3 | -43.7 | -130.3 | -10'.) | -131.5 | $-143.2$ | -147.3 | $-143.2$ | -130.5 | $-109.5$ | -80.5 | -43.9 | -. 3 | 50.0 | 106.2 |
| 2 | 10.00 | 154.4 | 107.0 | 55.7 | 31.6 | 4.4 | $-15.2$ | -27.1 | -31.1 | -27.1 | -15.2 | 4.4 | 11.6 | 65.9 | 107.0 | 154.4 | -152.3 |
| 1 | 0.00 | -103.5 | $-147.5$ | 172.5 | 141).2 | 114.7 | 9', 2 | 85.0 | H1.3 | 85.0 | 96.2 | 114.7 | 140.2 | 172.6 | $-148.5$ | -103.3 | -32.8 |

TABLE 15. CONTINUED


PRISSUKE AMPLITUDE: DA RE 20 MICRO PA


TABIE 15. CONTINUED

### 6.0 PREDICTED AND MEASURED NOISE REDUCTIONS FOR METRO II TEST

Nine test conditions were considered by Prydz et al [5] in the Metro II experimental program. These configurations are listed in Table 5. The first seven configurations are associated with either a floating or a disconnected trim panel. Since the analytical model assumes that there is no structureborne path between the fuselage structure and the trim panel, it forms a reasonably good representation of the floating trim. Thus, comparisons between measured and predicted noise reductions have been restricted to Configurations 1 through 7 .

The computations were performed at six discrete frequencies. Three of these frequencies were the same as those used in the Metro II tests [5], i.e., 383, 766 , and 1149 Hz . The frequencies were computed as harmonics of a fundamental frequency of 191.5 Hz . Following the criterion given in Section 2.2 , the low frequency form of the analytical model was applied at frequencies below 630 Hz and the high frequency form at frequencies above 630 Hz .

The peak exterior sound pressure level was assumed to be 134 dB at each frequency. This value represents a free-field sound pressure level and is consistent with the Metro II data [5]. Space-average sound pressure levels were calculated for the interior of the test fuselage and noise reduction computed by taking the difference between the exterior free field level of 134 dB and the interior space-average level [5]. The relationship between the free-field and blocked pressures is given by Eq. (48). The relationship corresponds to empirical data of Magliozzi [15].

The analytical model described herein and in Ref.[1] considers the fuselage cylindrical shell and cabin floor as a single unit. Thus the modeling of the floor is important to the computation
of noise reduction. The floor used in the test fuselage, described in Sections 3.2 and 4.0 , is of unusual construction and installation relative to normal aircraft design. Thus, the appropriate analytical modeling of the floor has been of concern in the present study.

In summary, the test floor is a heavy, stiff structure which is either mounted on vibration isolators or attached to the fuselage frames by means of brackets. Certain assumptions introduced in the modeling are discussed in Section 4.0. In addition, the floor treatment transmission coefficient, $\tau_{F}$, described in Section 2.2 can be used as a means of representing vibration isolation provided by the air mounts supporting the floating floor.

The influence of the floor treatment transmission coefficient can be seen in Figures 18 through 20 which compare predicted noise reductions for Configurations 1 through 6. The noise reductions were computed for four values of the floor treatment transmission coefficient, of which transmission losses of 0,20 and 400 dB are shown in the figures. The transmission loss of $0 d B$ represents the case where the untreated floor is rigidly attached to the fuselage shell, and the transmission loss of 400 dB represents, essentially, a situation where there is no acoustic transmission via the floor structure. The intermediate, 20 dB , case is a possible model for the floating floor of the test structure.

Inspection of Figures 18 through 20 shows that the influence of the floor can dominate the noise transmission into the fuselage interior if zero transmission loss is assumed for the floor treatment (or mounts). Thus, in Figure 18, increases in the sidewall treatment transmission loss have little or no effect on the noise reduction for the fuselage as a whole. As the assumed transmission loss through the floor installation is increased,
(a) Baseline Fuselage Structure

(b) Heavy Fuselage Structure


FIGURE 18. PREDICTED NOISE REDUCTIONS FOR METRO II TEST FUSELAGE EXPOSED TO ELECTROPNEUMATIC EXCITATION (Floor Treatment Transmission Loss $=0 \mathrm{~dB}$ )


FIGURE 19. PREDICTED NOISE REDUCTIONS FOR METRO II TEST FUSELAGE EXPOSED TO ELECTROPNEUMATIC EXCITATION
(Floor Treatment Transmission Loss $=20 \mathrm{~dB}$ )


FIGURE 20. PREDICTED NOISE REDUCTIONS FOR METRO II TEST FUSELAGE EXPOSED TO ELECTROPNEUMATIC EXCITATION
(Floor Treatment Transmission Loss $=400 \mathrm{~dB}$ )
the sidewall treatment has greater influence. The full benefit of the sidewall treatment is achieved only when the floor treatment transmission loss is of similar magnitude. The predicted transmission loss provided by the sidewall treatment alone is represented by the trim factor plotted in Figure 21. The trim factor is defined as $10 \log \left(1 / \tau_{t}\right)$ where $\tau_{t}$ is the trim transmission coefficient discussed in Section 2.2 and in Appendix A of Ref.[1]. Figure 21 contains predicted trim factor spectra for the three sidewall treatments investigated in the Metro II tests. A comparison of Figures 20 and 21 shows that the relative effect of the sidewall treatments can be seen when a high value ( 400 dB ) is assumed for the floor treatment transmission loss. (It should be noted that the trim factor is not the only parameter describing the effect of the sidewall treatment on noise transmission into the cabin. The analytical model also allows the treatment to influence the response of the outer structure through the structural loss factor $\eta_{r}$, Eq. (17)).

The noise reduction spectra shown in Figures 18 through 20 can be superimposed on the corresponding measured noise reductions [5] for Configurations 1 through 6. The comparison is shown in Figure 22. The immediate observation is that, by choosing appropriate values for the floor treatment transmission loss, the predicted noise reductions can be made to envelope the measured values. The closest agreement between predicted and measured noise reductions for the Configurations (1 through 3) with the baseline outer wall is achleved when the floor treatment transmission loss is about 20 dB . A somewhat higher value is appropriate for Configurations 4 through 6 for the heavy outer wall. This difference is reasonable since the addition of iron-oxide vinyl to the fuselage shell changes the predicted structural modes and fuselage response.


FIGURE 21. PREDICTED TRIM FACTOR (Sidewall Treatment Transmission Loss) FOR METRO \|I TEST TREATMENTS


FIGURE 22. COMPARISON OF MEASURED AND PREDICTED NOISE REDUCTIONS FOR METRO \|I TEST FUSELAGE AND ELECTROPNEUMATIC EXCITATION (Different Assumed Floor Treatment Transmission Losses)


FIGURE 22. CONTINUED


FIGURE 22. CONTINUED

The predicted noise reduction spectra assume that the floor treatment transmission loss is independent of frequency. This is a simplification; a better fit between measured and predicted results could be achieved by making the transmission loss frequency dependent. However, such a detailed fit, although physically reasonable is considered to be outside the objective of the present comparison.

The data in Figure 22 generally show better agreement between measured and predicted noise reductions for Configurations 1 through 3, with the baseline outer wall, than for Configurations 4 through 6 for the heavy outer wall. One possible reason for this difference lies in the analytical representation for the fuselage shell. In practice, the heavy wall was achieved by bonding a sheet of iron-oxide vinyl to the fuselage skin; the analytical model assumed simply that the skin panel had an increased skin density.

Fuselage test Configuration 7 is similar to Configuration 5 with respect to the fuselage structure and sidewall treatment, except that the floor is attached directly to the fuselage shell rather than being supported on air mounts. Thus the structural modeling is different for the two configurations. Predicted noise reduction spectra, associated with Configuration 7 , are shown in Figure 23 for different assumed values for the floor treatment transmission loss. The predicted results are compared with a single measured spectrum [5]. (Ref. 5 does not present a range of measured noise reductions). In this case it appears that closest agreement between measured and predicted noise reductions occurs for a floor treatment transmission loss of 10 dB rather than the 20 dB or higher value appropriate to the floating floor configurations. This reduction in effective floor treatment transmission loss is consistent with test data (Figure 76 of [5]), reproduced in Figure 24, which show that the vibration levels measured on the attached floor are higher than


FIGURE 23. MEASURED AND PREDICTED NOISE REDUCTIONS FOR METRO II TEST CONFIGURATION WITH ATTACHED FLOOR (Configuration \#7)


FIGURE 24. COMPARISON OF FLOOR VIBRATION LEVELS WITH AND WITHOUT THE FLOOR ATTACHED TO THE FUSELAGE FRAMES [5]


FIGURE 25. MEASURED [5] AND PREDICTED INTERIOR SOUND LEVELS FOR DIFFERENT FLOOR MOUNTINGS
those on the floating floor. Thus, if the floor treatment transmission loss is interpreted as a representation of the isolation provided by the attachment method, it can take into account the attenuation provided by the air mounts.

Figure 25 compares predicted noise reductions for Configuration 5, assuming a floor treatment transmission loss of 20 dB , and Configuration 7 , with a transmission loss of 10 dB . The figure also shows measured noise reduction spectra for the two configurations. The agreement is generally good, but it could be improved by selection of frequency-dependent transmission losses for the floor treatment parameter.

The uncertainty regarding the participation of the test floor arises because of problems encountered in the analytical modeling of the floor structure. These problems have been discussed earlier in this report but are worth repeating. First, the test floor structure did not form an integral part of the fuselage structure as it would in a production airplane and as it is modeled in the analysis. Secondy, when the floor is mounted on the air mounts, the structural mode shapes of the fuselage and floor have to be determined by modeling the test floor. However, the coupling between the modes of the floor and the cabin volume has to be modeled with the test floor represented as a vibrating partition in its actual location. Finally the air mounts will influence the vibration levels of the test floor, but such an effect is not included in the analytical model except through the floor treatment transmission coefficient. For these reasons, the test configuration does not provide a very good basis for evaluating the analytical model.

### 7.0 EFFECT OF EXCITATION CHARACTERISTICS

The original analytical model [1-4] utilized in this study was constructed to represent propeller noise (deterministic) and reverberant field (random) excitations. Then, as described in Section 1 , the model has been extended to include other excitations. The discussion in Section 6 is concerned with only one type of excitation, that associated with an electropneumatic source with directivity designed to represent a high-speed propeller. In the present section four other exterior pressure fields will be considered for the Metro II fuselage, with emphasis being placed on test Configuration 1 . These pressure fields are:
(a) General aviation airplane propeller
(b) Turbulent boundary layer on a general aviation airplane fuselage
(c) Turbulent boundary layer on an ATP airplane fuselage
(d) Reverberant pressure field.

### 7.1 General Aviation Propeller

The propeller noise field has to be described for the analytical model in terms of amplitude and phase over the fuselage gridpoint array shown in Figure 16. Since the prediction of such a pressure field distribution can be time-consuming, use was made of an existing distribution computed for a general aviation type propeller [1]. The pressure field was then modified to fit the dimensions of the Metro II fuselage.

Scaling of the propeller noise field was performed on the basis of cylinder or fuselage radius, while maintaining the same propeller tip speed. The reference cylinder [1] had a diameter of $1.12 \mathrm{~m}(3.67 \mathrm{ft})$ and the propeller had a diameter of 0.76 m (2.49 ft). Thus, scaling to Metro II dimensions, with a fuselage diameter of $1.67 \mathrm{~m}(5.5 \mathrm{ft})$, the propeller diameter becomes 1.26 m (4.1 ft). The clearance between the fuselage and propeller tip is $0.13 \mathrm{~m}(0.41 \mathrm{ft})$ or ten percent of the propeller diameter. In order to maintain the same tip speed, the propeller rpm was taken to be 2424; the blade passage frequency is then 121.2 Hz for a three-bladed propeller.

The amplitude and phase components of the propeller pressure field are listed in Table 16 for the three-lowest-order harmonics. The presentation follows the same format as that given in Table 15 for the electropneumatic excitation. Now, however, the phase angle in the lower quadrant of the fuselage differs from that in the upper quadrant to take into account the rotation of the propeller pressure field. In the case of the electropneumatic source the phase angle distribution is symmetric about the grid axes.

Noise reductions predicted for the propeller excitation are shown in Figure 26 for fuselage test Configuration 1 and four assumed values of the floor treatment transmission loss. The noise reduction spectra show a broad peak centered at the fifth harmonic ( 606 Hz ). This peak seems to be associated with reduced excitation efficiency of the fuselage structure.

### 7.2 Turbulent Boundary Layer

Turbulent boundary layer pressure fields were predicted for two flight conditions typical of general aviation (GA) and advanced turboprop (ATP) airplane cruise conditions. These conditions are listed in Table 17. The conditions were selected mainly to explore the effect of airspeed on the noise transmission.

PROPELLER HARMONIC 1 AT 121.2 Hz


PRUPELLER HARMONIC 2 AT 242.4 HZ

| CIRCUM. |  |  |  |  |  |  | AXIAL | LOCATI |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| LOCATION | $K=1$ | 2 | 3 | 4 | 5 | 6 | 7 | B | 9 | 10 | 11 | 12 | 13 | 14 | 13 | 16 |
| L THETA | 2=2.439 | 2.586 | 2.732 | 2.878 | 3.025 | 3.171 | 3.317 | 3.464 | 3.610 | 3.756 | 3.402 | 4.049 | 4.195 | 4. 141 | 4.488 | 4.634 |

PRESSURE AMPLITUDE, DB RE $\angle 0$ MICRU PA

| 19 | 180.00 | 75.5 | 77.6 | 79.3 | 81.0 | 82.5 | 83.7 | 84.6 | 85.3 | 85.6 | 85.6 | 45. 1 | 64.3 | 81.1 | 81.8 | Bi).0 | 70.1 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 18 | 170.00 | 75.7 | 78.2 | 80.1 | 82.1 | 83.6 | 84.9 | 85.8 | 8t. 4 | 86.6 | 86.5 |  | 34.4 | 81.4 | 61.8 | 14.8 | 76.t |
| 17 | 160.00 | 75.9 | 78.8 | 81.2 | 83.5 | 85.3 | 86.6 | 87.5 | 88.0 | 88. 1 | -7. 7 | 86.7 | H5.5 | 43.0 | 61.7 | 74.3 | 70.7 |
| 16 | 150.00 | 75.9 | 79.4 | 82.3 | 85.2 | 87.5 | 89.1 | 89.8 | 70.1 | 40.0 | (3). 4 | 88. 2 | 86.3 | 43.8 | H1.4 | 74.4 | 75.4 |
| 15 | 140.00 | 78.1 | 82.4 | 86. 2 | 90.0 | 93.2 | 95.3 | 96.0 | 75.9 | 95.6 | 44.7 | 92.8 | 90.0 | 86.6 | 83.4 | 74.7 | 76.0 |
| 14 | 130.00 | 78.9 | 84.2 | 88.9 | 93.9 | 98.3 | 101.3 | 102.3 | 101.5 | 101.1 | 99.8 | 46.4 | 92.9 | 38.4 | 64.3 | 74.8 | 75.4 |
| 13 | 120.00 | 78.2 | 84.5 | 90.3 | 96.8 | 102.9 | 107.6 | 109.3 | 107.5 | 107.5 | 105.5 | 100.8 | 95.0 | 89.0 | 83.8 | $7 \mathrm{H}$. | 73.2 |
| 12 | 110.00 | 76.6 | 83.7 | 90.5 | 78.5 | 106.4 | 113.2 | 116.7 | 113.7 | 114.7 | 110.8 | 103.9 | 16.2 | 88.7 | $8<.4$ | 76.0 | 70.3 |
| 11 | 100.00 | 74.5 | 82.4 | 90.0 | 94.1 | 100.5 | 117.3 | 123.3 | 114.2 | 121.2 | 114.8 | 105.8 | 96.6 | 117.0 | 80.7 | 71.6 | 67.3 |
| 10 | 90.00 | 74.1 | 82.1 | 89.9 | 99.1 | 108.8 | 118.0 | 124.5 | 120.3 | 122.4 | 115.5 | 100.1 | 76.6 | d7.6 | 80.4 | 73.2 | 66.8 |
| 9 | 80.00 | 74.5 | 82.4 | 90.0 | 99.1 | 108.5 | 117.3 | 123.3 | 119.2 | 121.2 | 114.8 | 105.8 | 46.6 | H7.8 | 60.7 | 73.6 | 67.1 |
| 8 | 70.00 | 76.6 | 83.7 | 90.5 | 98.5 | 106.4 | 113.2 | 116.7 | 113.7 | 114.7 | 110.8 | 103.4 | 96.2 | 88.7 | H2.4 | 76.0 | 70.3 |
| 7 | 60.00 | 78.2 | 84.5 | 90.3 | 96.8 | 102.9 | 107.6 | 109.3 | 107.5 | 107.5 | 105.5 | 100.H | 45.0 | 84.0 | 83.8 | 78.3 | 73.2 |
| 6 | 50.00 | 78.9 | 84.2 | 88.9 | 93.9 | 98.3 | 101.3 | 102.3 | 101.5 | 101.1 | 49.8 | 76.9 | 72.4 | 88.4 | 84.3 | 74.8 | 75.4 |
| 5 | 40.00 | 78.1 | 82.4 | 86.2 | 90.0 | 93.2 | 95.3 | 96.0 | 95.9 | 95.6 | 94.7 | 92.8 | 70.0 | d0.6 | 63.4 | 74.7 | 76.0 |
| 4 | 30.00 | 75.9 | 79.4 | 82.3 | 85.2 | 87.5 | 89.1 | 89.8 | 90.1 | 90.0 | 89.4 | 88. 2 | H6. 3 | d1.0 | 81.4 | 78.4 | 75.4 |
| 3 | 20.00 | 75.9 | 78.8 | 81.2 | 83.5 | 85.3 | 86.6 | 87.5 | 84.0 | $8 \mathrm{H.1}$ | 87.8 | 86.9 | 63.5 | 83.6 | 81.7 | 74.3 | 76.1 |
| 2 | 10.00 | 75.7 | 78.2 | 80.1 | 82.1 | 83.6 | 84.9 | 85.8 | 86.4 | 86.6 | 86.5 | 85.9 | 44.8 | 83.4 | 81.8 | 74.8 | 77.6 |
| 1 | 0.00 | 75.5 | 77.6 | 74.3 | 81.0 | 82.5 | 83.7 | 84.6 | 85.3 | 85.6 | BS.6 | 85.1 | 84.3 | 83.1 | 81.8 | 80.0 | 7H.1 |


| 19 | 180.00 | 179.0 | 168.6 | 161.6 | 156.6 | 154.7 | 155.7 | 159.4 | 164.9 | 171.6 | 178.9 | -173.2 | -164.6 | -13. 0 | $-145.5$ | -133.5 | -120.2 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 18 | 170.00 | 148.1 | 137.2 | 129.9 | 124.6 | 122.8 | 124.4 | 129.2 | 136.2 | 144.2 | 152.5 | 1.60 .4 | 169.6 | 17\%.1 | -171.6 | $-154.8$ | $-146.5$ |
| 17 | 160.00 | 115.5 | 104.2 | 96.5 | 90.8 | 88.9 | 91.0 | 97.3 | 106.7 | 117.0 | 126.8 | 135.5 | 144.0 | 153.0 | 161.8 | 173.1 | $-174.0$ |
| 16 | 150.00 | 92.3 | 80.7 | 72.8 | 66.8 | 64.6 | 67.0 | 75.0 | 87.9 | 101.9 | 113.5 | 122.3 | 124.9 | 137.0 | 143.3 | 154.0 | 167.7 |
| 15 | 140.00 | 78.2 | 66.7 | 58.8 | 52.7 | 50.1 | 52.4 | 61.9 | 79.9 | 99.7 | 113.3 | 121.4 | 127.0 | 13L. 3 | 13 J .5 | 147.1 | 150.0 |
| 14 | 130.00 | 75.4 | 64.2 | 56.6 | 50.5 | 47.5 | 49.1 | 59.3 | 84.5 | 113.5 | 128.4 | 134.2 | 136.9 | 134.5 | 143.2 | 149.7 | 154.0 |
| 13 | 120.00 | 91.2 | 80.7 | 73.5 | 67.6 | 64.2 | 64.4 | 73.5 | 109.2 | 151.5 | 164.5 | 166.8 | 166.3 | 163.9 | 167.0 | 171.0 | 178.3 |
| 12 | 110.00 | 121.1 | 111.3 | 104.5 | 98.9 | 95.1 | 93.8 | 99.9 | 147.4 | -157.9 | -150.1 | $-150.8$ | -153.3 | -155.8 | $-150.6$ | -134.6 | -14\%.2 |
| 11 | 100.00 | 165.4 | 156.1 | 149.7 | 144.2 | 140.2 | 137.8 | 140.4 | -160.4 | -94. 1 | -96.3 | -98.7 | $-102.3$ | -103.t | -107.8 | -101.4 | -103. 2 |
| 10 | 90.00 | -128.9 | -138.1 | -144.4 | -149.8 | -153.9 | -156.5 | -154.5 | -92.8 | -31.0 | -29.1 | $-31.7$ | -35.5 | -34.1 | -41.2 | -41.0 | -37.1 |
| 9 | 80.00 | -65.7 | -75.0 | -81.4 | -86.9 | -90.9 | -93.3 | -90.7 | -31.5 | 29.8 | 32.6 | 30.2 | 26.6 | 23.1 | <1.1 | 21.5 | 23.0 |
| 8 | 70.00 | -3.7 | $-13.5$ | -20.2 | -25.9 | -29.7 | -31.0 | -24.9 | 22.7 | 77.3 | 85.1 | 84.5 | 81.9 | 74.4 | 76.t | 80.7 | 86.0 |
| 7 | 60.00 | 40.9 | 30.4 | 23.2 | 17.3 | 13.9 | 14.1 | 23.2 | 58.7 | 101.2 | 114.2 | 116.5 | 116.0 | 117.5 | 116.1 | 120.7 | 120.0 |
| 6 | 50.00 | 69.9 | 58.8 | 51.1 | 45.1 | 42.1 | 43.7 | 53.9 | 79.1 | 108.0 | 123.0 | 128.7 | 131.4 | 134.1 | 137.1 | 144.3 | 153.5 |
| 5 | 40.00 | 94.5 | 83.0 | 75.1 | 68.9 | 66.4 | 68.7 | 78.2 | 96.1 | 115.9 | 129.6 | 117.6 | 143.2 | 14 d. 7 | 134.0 | 103.4 | 114.3 |
| 4 | 30.00 | 113.2 | 101.6 | 93.6 | 87.7 | 85.5 | 87.9 | 95.9 | 108.8 | 122.8 | 134.4 | 143.2 | 150.8 | 1) (\%) | 160.2 | 176.5 | -171.4 |
| 3 | 20.00 | 128.7 | 117.3 | 104.6 | 104.0 | 102.0 | 104.1 | 110.5 | 119.8 | 130.2 | 139.9 | 148.7 | 157.2 | 1ヵt.1 | 1) ${ }^{1}$ | -173.H | $-1611.4$ |
| 2 | 10.00 | 144.6 | 133.7 | 126.4 | 121.1 | 119.2 | 120.8 | 125.6 | 132.5 | 140.7 | 148.4 | 157.3 | 168.0 | 17\%.) | -17\%.1 | -163.3 | -1;0.1 |
| 1 | 0.00 | 152.1 | 141.6 | 134.7 | 129.7 | 127.7 | 128.8 | 132.5 | 138.0 | 144.7 | 152.0 | 159.t | 1,H.5 | 178.1 | $-172.4$ | -160.4 | $-141.1$ |



| 19 | 180.00 | 154.3 | 135.4 | 121.8 | 110.2 | 102.9 | 99.6 | 100.2 | 103.9 | 110.1 | 118.2 | 128.0 | 139.6 | 153.3 | 167.3 | -174.7 | 4 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 18 | 170.00 | 109.4 | 90.0 | 76.0 | 64.1 | 56.8 | 54.0 | 55.5 | 60.6 | 68.0 | 77.0 | 87.3 | 97.1 | 112.7 | 126.6 | 144.4 | 164.6 |
| 17 | 160.00 | 61.8 | 41.9 | 27.7 | 15.7 | 8.5 | 6.2 | 9.1 | 16.3 | 26.0 | 36.3 | 46.7 | 58.3 | 71.2 | 84.6 | 101.9 | 121.8 |
| 16 | 150.00 | 27.9 | 8.2 | -5.8 | -17.5 | -24.4 | -26.0 | -21.5 | -11.1 | 2.0 | 14.0 | 24.2 | 34.2 | 45.4 | 37.4 | 13.6 | 92.6 |
| 15 | 1.40 .00 | 7.6 | -11.3 | -24.5 | -35.3 | -41.6 | -42.7 | -36.5 | -21.4 | -2.4 | 11.1 | 20.4 | 27.2 | 35.2 | 44. | 58.8 | h. 3 |
| 14 | 130.00 | 4.2 | -13.4 | -25.3 | -34.9 | -40.5 | -41.4 | -34.2 | -12.1 | 17.1 | 32.7 | 38.0 | 40.3 | 43.8 | 49.H | 60.7 | 7h. 3 |
| 13 | 120.00 | 29.5 | 13.4 | 2.9 | -5.4 | -10.6 | -12.1 | -5.5 | 26.6 | 71.0 | 84.3 | 35.3 | 83.2 | 81.7 | 83.6 | 40,5 | 102.7 |
| 12 | 110.00 | 76.8 | 61.6 | 52.0 | 44.4 | 34.2 | 36.5 | 40.5 | 83.7 | 142.0 | 149.2 | 147.0 | 142.4 | 137.9 | 136.1 | 138.6 | 147.0 |
| 11 | 100.00 | 143.6 | 130.3 | 121.5 | 114.2 | 108.8 | 105.1 | 106.1 | 160.4 | -134.2 | -132.4 | -136.3 | -141.7 | -147.3 | -151.3 | -151.3 | -144.4 |
| 10 | 90.00 | -121.6 | -133.1 | -140.6 | -147.1 | -152.2 | -155.9 | -155.5 | -98.3 | -32.8 | -31.9 | $-35.8$ | -41.0 | $-46.1$ | $-48.7$ | -46.5 | -37.1 |
| 9 | 80.00 | -23.0 | -36.4 | $-45.1$ | -52.4 | -57.4 | -61.5 | -60.6 | -6. 2 | 59.2 | 61.0 | 57.1 | 51.7 | 45.7 | 42.1 | 42.1 | $4 \% .0$ |
| 8 | 70.00 | 69.7 | 54.5 | 44.9 | 37.3 | 32.0 | 29.3 | 33.3 | 76.6 | 134.8 | 142.0 | 139.9 | 135.3 | 130.8 | 128.9 | 131.5 | 13.0. |
| 7 | 60.00 | 134.0 | 117.9 | 107.4 | 99.1 | 94.0 | 92.4 | 99.0 | 131.1 | 175.5 | -171.2 | $-170.2$ | -172.3 | -173.6 | -171.4 | -165.0 | -152.8 |
| 6 | 50.00 | 176.0 | 158.4 | 146.5 | 136.9 | 131.3 | 130.4 | 137.7 | 159.7 | -171.0 | $-155.4$ | -150.1 | -147.8 | -144.4 | -13H.4 | -127.4 | -111.7 |
| 5 | 40.00 | -148.0 | -166.9 | 179.9 | 169.1 | 162.7 | 161.7 | 167.9 | $-177.0$ | -158.0 | -143.9 | -135.2 | -128.4 | -120.4 | -110.7 | -36.8 | -74.1 |
| 4 | 30.00 | -120.8 | -140.5 | -154.5 | -166.2 | -173.1 | $-174.7$ | -170.1 | -159.8 | $-146.7$ | -134.7 | -124.4 | -114.5 | $-103.2$ | -71.3 | -75.1 | $-36.0$ |
| 3 | 20.00 | -98.5 | -118.3 | -132.6 | -144.5 | -151.8 | -154.0 | -151.1 | -143.7 | -134.1 | -124.0 | -113.4 | -102.0 | -44.0 | -15. 1 | $-38.4$ | -30.5 |
| 2 | 10.00 | -75.4 | -45.4 | -109.4 | -121.2 | -128.6 | -131.4 | -129.9 | -124.8 | -117.3 | -108.3 | -18.0 | -36.3 | -72.7 | -28.3 | -41.0 | $-20.8$ |
| 1 | 0.00 | -66.1 | -84.9 | -34.6 | -110.1 | $-117.5$ | $-120.1$ | $-120.2$ | $-116.4$ | $-110.3$ | $-102.2$ | -42.4 | - HO 0.1 | -67.1 | -33.1 | -35.1 | -14.3 |



FIGURE 26. PREDICTED NOISE REDUCTIONS FOR TYPICAL G.A. PROPELLER NOISE EXCITATION ( METRO II TEST FUSELAGE CONFIGURATION \#1)

TABLE 17

FLIGHT CONDITIONS ASSOCIATED WITH BOUNDARY LAYER ESTIMATES

| Parameter | GA Airplane | ATP Airplane |
| :---: | :---: | :---: |
| Airspeed m/s (ft/sec) | 95 (416) | 182 (796) |
| Altitude m (ft) | 3650 (12000) | 9150 (30000) |
| Distance from airplane nose to point of interest on fuselage m (ft) | 6.1 (20) | 16.9 (55.5) |
| Boundary layer thickness mm (in) | 70 (2.8) | 174 (6.8) |
| Boundary layer displacement thickness mm (in) | 8.8 (0.35) | 21.7 (0.86) |

Parameters required for the analytical model are the pressure field convection velocity and the boundary layer displacement thickness (see Eqs.(41) and (45)). The convection velocity was taken to be $80 \%$ of the flight airspeed.

Since the turbulent boundary layer pressure field is broadband and random in character, use is made of the analytical model feature which predicts noise reduction in one-third octave frequency bands rather than the deterministic, tonal transmission calculation used for propeller noise and the electropneumatic source. The relevant part of the analytical model is described in Section 2. Predicted noise reductions associated with the GA boundary layer pressure field are shown in Figure 27 for fuselage test Configuration 1 and four assumed values of the floor treatment transmission loss. The results show a general increase of noise reduction with frequency, except for a spectral trough at 125 Hz associated with sidewall resonance. In this case noise reduction refers to the surface or blocked pressure on the exterior of the fuselage rather than the free field sound pressure used for the propeller and electropneumatic noise sources discussed in preceding sections.

### 7.3 Comparison of Predicted Noise Reductions

Noise reductions predicted for boundary layer, GA propeller noise and simulated propeller (electropneumatic) excitations are compared in Figure 28. Two values of. the floor treatment transmission loss ( 0 dB and 20 dB ) are considered. Again, it should be remembered that the reference sound level for propeller and electropneumatic noise sources is the free field value, whereas the surface pressure is used as the reference for the boundary layer cases. If surface (blocked) pressure were used in all cases the noise reductions for propeller and electropneumatic noise excitation should be increased by 6 dB .


FIGURE 27. PREDICTED NOISE REDUCTIONS FOR TYPICAL GENERAL AVIATION AIRPLANE BOUNDARY LAYER EXCITATION (METRO II TEST FUSELAGE CONFIGURATION \#1)


FIGURE 28. COMPARISON OF PREDICTED NOISE REDUCTIONS FOR DIFFERENT EXCITATION PRESSURE FIELDS (METRO II TEST FUSELAGE CONFIGURATION 1)

Comparing first the results for GA propeller and simulated propeller (electropneumatic) excitations it is seen that the noise reductions are higher for the propeller than for the electropneumatic source. Two factors, amplitude and phase of the excitation pressure field, are involved. Both of these factors determine the response of the fuselage structure. Pressure phase differences for the two excitations have been discussed in Section 7.1; the pressure amplitude differences are associated with the spatial distribution, particularly in the longitudinal direction. A representative comparison of the longitudinal distribution of harmonic sound levels is given in Figure 29 where it is seen that the propeller noise pressure amplitude decreases more rapidly with distance than does the electropneumatic source sound field. The latter sound field was designed to simulate propfan pressure levels (which it does successfully [16]), so that Figure 29 essentially compares GA and ATP propeller noise levels.

The different excitation characteristics of the excitation fields will result in different structural responses and, hence, noise transmission into the fuselage. In this case, the more rapid the fall-off of the pressure amplitude, the higher the noise reduction. It is interesting to note that Prydz et al [16] have used an associated argument regarding average and peak exterior sound levels to increase predicted noise reductions by 7 dB . (The predictions in that case [16] were computed initially under the assumption of spatially uniform pressure amplitude and then adjusted to account for the spatially-varying amplitude).

Now compare noise reductions for acoustic (propeller and electropneumatic) and aerodynamic (turbulent boundary layer) excitations. It is seen in Figure 28 that the predicted noise reductions associated with boundary layer excitation lie between those for propeller and electropneumatic sources, even when the $6 d B$


FIGURE 29. TYPICAL LONGITUDINAL DISTRIBUTIONS OF EXCITATION PRESSURE AMPLITUDE FOR G.A.PROPELLER AND SIMULATED PROPFAN (ELECTROPNEUMATIC) NOISE SOURCES
adjustment from free field to blocked pressure excitation is applied. As before, excitation pressure amplitude and phase are playing important roles. The boundary layer has a uniformlydistributed pressure amplitude, in contrast to the spatiallyvarying amplitudes for the propeller and electropneumatic sources. Also the phase characteristics are quite different, the boundary layer pressure field being convected in the longitudinal directions and the propeller pressure field circumferentially. As the airplane airspeed and pressure field convection velocity increase, the predicted noise reduction decreases and approaches values associated with the electropneumatic source.

### 7.4 Reverberant Field Excitation

Noise reductions have been predicted for Configuration 1 of the Metro II test fuselage when the excitation was a reverberant acoustic field. The resulting noise reduction spectra are compared in Figure 30 for different values of the floor treatment transmission loss. Figure 30 also contains a measured noise reduction spectrum [private communication from R.A. Prydz] measured on the test fuselage when placed in a reverberation chamber. It is seen that the measured and predicted spectra have similar shapes, although the analytical model underpredicts the noise reduction at lower frequencies. It is also observed that the noise reductions predicted for reverberant excitation are lower than those predicted for the other excitations considered in Figure 28. A comparison of predicted noise reductions for reverberant acoustic and boundary layer aerodynamic excitations is contained in Figure 31, where the floor treatment transmission loss is assumed to be zero.


FIGURE 30. COMPARISON OF MEASURED AND PREDICTED NOISE REDUCTION FOR REVERBERANT FIELD EXCITATION (METRO II TEST FUSELAGE CONFIGURATION 1)


FIGURE 31. COMPARISON OF PREDICTED NOISE REDUCTIONS FOR REVERBERANT ACOUSTIC AND TURBULENT BOUNDARY LAYER EXCITATIONS (METRO \|I TEST FUSELAGE CONFIGURATION \#1, FLOOR TREATMENT TRANSMISSION LOSS $=0 \mathrm{DB}$ )

### 8.0 DISCUSSION AND CONCLUSIONS

The analytical model developed in Ref.[1] for the prediction of sound pressure levels in propeller-driven aircraft provides a representation of the response of the combined fuselage shell and cabin floor structure. This inclusion of the cabin floor as an integral part of the analytical model is a significant departure from earlier analytical models. As a result, a certain amount of experience has to be gained in the application of the model to practical structures.

In the comparisons with data from the Metro II noise transmission tests performed by Prydz et al, it became apparent that certain changes were required to the way in which the analytical model [1] represented the floor structure. The analytical model was modified to incorporate the changes but unfortunately the particular floor configuration chosen for the Metro II tests posed a significant problem to the analytical representation. The floor in the test fuselage did not constitute a structural member of the fuselage, rather it was a secondary item (albeit heavy and stiff) introduced to provide a cavity which was geometrically similar to that of a wide-body airplane. The test floor rested on the structural floor of the Metro II but it did so via air mounts which would provide some vibration isolation.

Since the analytical model could not handle the test floor structure without a major change to the structural model, an alternative approach was adopted whereby a "floor treatment transmission coefficient" was introduced to represent, among other things, the vibration isolation provided by the air mounts. The coefficient could, equally well, represent attenuation provided by any floor coverings, just as the trim transmission coefficient provides an estimate of the attenuation provided by the sidewall trim.

Parametric variations of the floor treatment transmission coefficient allowed the predicted noise reductions to stradde the measured values. Appropriate values of the coefficient could be selected to achieve good agreement between measured and predicted results. While this may appear to be an indirect approach it can be justified because of the unusual floor configuration in the test fuselage. Furthermore the values of the floor treatment transmission loss that provide good agreement between measured and predicted noise reductions are consistent with the measured reductions in floor vibration attributed to the use of air mounts (as shown in Figure 24).

It should be noted also that the analytical model assumes that the sound-absorbing material is distributed uniformly over the sidewall surface. In the actual test the sidewall treatments were highly reflective and additional absorptive material was placed on the floor near the sidewall along the full length of the cabin. Since the analytical model would not account for this ad hoc distribution of absorptive material, the appropriate empirical information was introduced into the computations.

Because of these problems in representing the floor installation and absorptive material, the validation of the analytical model was not conclusive. However, the results do show that the analytical model predicts noise reductions which are consistent with the measured data.

The capability of the analytical model has been extended to include turbulent boundary layer pressure fluctuations as an exterior excitation field. By the same token the model can represent a convected acoustic pressure field. The model already had the capability of describing noise transmission from a reverberant acoustic field. In all these cases the excitation is broadband and random in character rather than tonal and deterministic as is the case with propeller noise or simulated propeller (electropneumatic) noise.

The ability of the analytical model to predict noise reductions for the fuselage structure exposed to various exterior pressure fields has been used to compare propeller and simulated propeller excitation, turbulent boundary layer and reverberant acoustic fields. While the results are probably dependent to some extent on the fuselage structure chosen (Metro II test Configuration \#l in Ref.[5]), the comparisons show certain trends of interest.

Propeller noise (or simulated propeller noise) has the characteristic that the sound pressure level changes markedly with distance away from the plane of rotation of the propeller, but the coherence maintains high values for, at least, the lower order harmonics [17]. In contrast, the boundary layer pressure field has a sound pressure level which is essentially independent of location but a coherence function which decays with increasing separation distance.

It can be shown (see, for example, Eq. (55) of [6]) that the amplitude spatial decay and coherence spatial decay play similar roles in the calculation of panel joint acceptance. Consequently, it might be expected that the noise reductions associated with propeller noise and boundary layer excitation would be similar. This appears to be the case in Figure 28, particularly with respect to the simulated propeller (where the sound pressure level decays relatively slowly with distance) and the higher speed (ATP airplane) boundary layer. As the amplitude spatial decay rate (GA propeller) or the coherence spatial decay rate (GA boundary layer) increases, the predicted noise reduction increases.

Obviously, the results presented in this report represent a limited number of special cases. A more extensive parametric study would define the roles of the different excitations more clearly. Alternatively, calculations should be made for specific airplane configurations of interest.

In conclusion, the analytical model developed in Ref.[1] and extended in this study predicts noise reductions which are consistent with values measured in the Metro II test [5]. However certain configurations peculiar to the test set-up make a definitive validation difficult. Application of the analytical model to different types of excitation show that there can be differences in the predicted noise reductions. For example the predicted noise reductions asssociated with turbulent boundary layer excitation were greater than those for simulated propfan excitation. However, before general conclusions are drawn, calculations should be performed for different fuselage structures with special attention placed on the modeling of the floor.

One further point to be noted is that, because of the unusual distribution of sound-absorbing material in the test cabin, empirical rather than analytical means had to be used to describe the absorption in the cabin. Consequently, it was not possible to make a complete evaluation of the analytical model for the sidewall treatment.

Strictly speaking, the validation of the analytical model has been performed for only one fuselage structure and one excitation. Thus, confidence in the analytical model when applied to other structures and excitation is dependent to some extent on a qualitative judgment of the accuracy of the extrapolation of the model to other conditions. The role played by the floor in a conventional fuselage configuration is of particular interest since it was the modeling of the test floor that posed a major problem in the present study. In spite of these difficulties it seems justifiable to apply the analytical model to airplane configurations such as those discussed in Reference 6 in order to obtain alternative estimates of the treatment weights required to achieve the cabin noise goal. In this manner it should be possible to obtain an indication of the confidence limits for the predictions by the use of different analytical approaches.

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## APPENDIX A

## List of Symbols

This appendix contains the list of symbols used in the equations in Reference [1]. The equations in this report rely heavily on [1] for their development, and the symbols have not been redefined in the text since the reference will receive wide distribution. The reader should refer to Reference [1] for a detailed presentation of the development of the basic analytical model.

## LIST OF SYMBOLS

| A | Cylinder surface area |
| :---: | :---: |
| or |  |
| A | ```Interior (cylinder & floor) surface area, used in conjunction with structural/acoustic coupling function f'(n,r)``` |
| $A^{\prime}$ | Transmitting area of cylinder without trim |
| $\bar{A}$ | Exterior cylinder surface area |
| $A_{t}$ | Transmitting area of cylinder with trim |
| $\left(A+A^{\prime}\right)$ | Total transmitting area of cylinder, floor to floor |
| $A_{H}^{m}$ | Amplitude of Fourier component of blocked propeller pressure signature at propeller harmonic $H$ and grid location $m \equiv(k, l)$, see Eq. (43) of [I] |
| $\begin{gathered} \overline{\mathrm{A}}_{\mathrm{H}}^{\mathrm{m}} \end{gathered}$ | Amplitude of Fourier component of free field propeller pressure at harmonic $H$ and grid location $m \equiv(k, \ell)$ |
| a | Radius of cylinder |
| $\mathrm{a}_{\mathrm{H}}^{\mathrm{m}}$ | Fourier series coeffficient of propeller pressure, for harmonic $H$ at location $m$; see Eqs.(27),(29)[1] |
| $a_{0}^{m / 2}$ | Mean propeller pressure amplitude at location m (defined in Eq. (28) of [1]) |
| $\left.\arctan _{n} \arctan _{r}\right\}$ | Functions defined in Section 2.3, Eq. (10), et seq |
| B | Number of propeller blades |
| BPF | Propeller blade passage frequency ( Hz ) ; BPF $=\frac{N B}{60}$ |

## LIST OF SYMBOLS

(Continued)
$b_{n}, b_{r} \quad$ Functions defined in Section 2.3, Eq. (10) et seq
$\begin{array}{ll}b_{H}^{m} & \begin{array}{l}\text { Fourier series coefficient of } \\ \text { for harmonic } H \text { at location } m ;\end{array} \\ C(j) & \text { Function defined in Eq. (76)[1] }\end{array}$
$\operatorname{Cin}(z) \quad$ Cosine integral; $\operatorname{Cin}(z)=\int_{0}^{z}(1-\cos t) d t / t$
$C_{\mathrm{Mn}}^{\mathrm{pr}} \mathrm{C}_{\mathrm{IM}}^{\mathrm{sr}} \quad$ Floor and shell generalized coordinates for structural mode $r \equiv(M, N)$; see Eqs. (46)-(49)[I]
$C_{p_{b l}}\left(\bar{x} \mid \bar{x}^{\prime} ; \omega\right)$ Cospectral density function of the blocked exterior pressure field
$C_{p}\left(\bar{x} \mid \bar{x}^{\prime} ; \omega\right)$ Cospectral density function of the exterior pressure field
$C_{W}, C_{W}^{R}, C_{W}^{I} \quad$ Trim parameter, derived from the trim transfer matrix, Eq. (A. 8 ); $C_{w}=C_{w}^{R}+i C_{w}^{I}$
$C_{x}(\xi, \omega) C_{y}(\zeta, \omega)$ Cospectral density functions of the exterior pressure field in the axial and transverse directions respectively; see Eq. (60) of [1]
$c_{0} \quad$ Speed of sound in air
$c_{n}, c_{r} \quad$ Functions defined in Section 2.3, Eq.(IO) et seq
$c_{\omega} \quad$ Constant percentage bandwidth parameter, where $\Delta \omega=c_{\omega} \omega\left[c_{\omega}=0.232\right.$ for one-third octave bands]

## LIST OF SYMBOLS

## (Continued)

| $\mathrm{D}_{\mathrm{nr}}$ | Function defined in Section 2.3, Eq. (10) et seq |
| :---: | :---: |
| $E[$ ] | Expected value of a function |
| $f^{\prime}(n, r)$ | Interior structural/acoustic coupling factor; see Eq. (45) |
| $\begin{aligned} \bar{f}^{\prime}(n, r)=\bar{f}^{\prime}(q i, r) & \text { Interior structural/acoustic coupling factor } \\ & \text { including effect of trim factor } \tau_{t} \text { and floor } \\ & \text { treatment factor } \tau_{F} \text {, see Eq. (3) } \end{aligned}$ |  |
| $f^{\prime}$ | Frequency of propeller lst harmonic; $f_{I}=1 / T_{0}=$ BPF |
| $f_{H}$ | Frequency of propeller harmonic $\mathrm{H} ; \mathrm{f}_{\mathrm{H}}=\mathrm{Ff}_{1}$ |
| $\mathrm{f}_{\mathrm{qm}}$ | Acoustic/structural coupling factor in axial direction; see Eq.(57) of [1] |
| $g_{n}, g_{r}$ | Functions defined in Section 3.2, page 3-11 [1] |
| H | Propeller harmonic order, used as superscript to denote functions evaluated at frequency $\omega_{H}$ |
| i | Acoustic mode number counter for fuselage crosssection modes, associated with mode $n \equiv(q, i)$ |
| $\mathrm{I}_{1} \mathrm{I}_{2} \mathrm{I}_{3}$ | Integrals defined in Equations (63) and (70)-(73)[1] |
| j | Circumferential location on fuselage wall, $\theta_{j}$, a boundary point at which the acoustic eigenvector is evaluated (see Fig.C-2 [1]) |

## LIST OF SYMBOLS

(Continued)

| $j_{\text {Fl }}^{2}(\omega)$ | Structural joint acceptance function in axial direction |
| :---: | :---: |
| $j_{i J}^{2}(\omega)$ | Structural joint acceptance function in circumferential direction |
| $j_{r}^{2}(\omega)$ | $\begin{aligned} & \text { Structural joint acceptance in axial and } \\ & \text { circumferential directions; } j_{r}^{2}(\omega) \equiv j_{M N}^{2}(\omega) \\ & =j_{M}^{2}(\omega) j_{N}^{2}(\omega) ; \text { see Eq. }(2 I) \end{aligned}$ |
| $j_{r}^{2 r e v}(\omega)$ | Joint acceptance for reverberant/diffuse excitation |
| $\left\langle j^{2}(\omega)\right\rangle_{r}^{r e v}$ | Joint acceptance for reverberant excitation <br> averaged over structural modes resonant in band $\Delta \omega$ |
| $k$ | Acoustic wave number, $k=2 \pi / \lambda$ |
| k | Axial non-dimensional coordinate for grid point; see Figure 16 |
| L | Fuselage structure length |
| $L_{p}$ | Floor width (wall to wall) |
| $\ell$ | Circumferential non-dimensional coordinate for grid point; see Figure 16 |
| $\ln _{n}, \ln n_{r}$ | Functions defined in Section 2.3, Eq. (10) et seq |
| M | Number of axial half-wavelengths for structural mode $r \equiv(M, N)$ |

## LIST OF SYMBOLS <br> (Continued)

| $M_{r}$ | Generalized modal mass, for structure mode r |
| :---: | :---: |
| $m \equiv(k, \ell)$ | Grid point on surface of cylinder used for propeller noise predictions; see Figure 16 |
| or |  |
| m | Average surface mass/unit area of cylinder |
| N | Structural mode counter, associated with mode $r \equiv(M, N)$ |
| or |  |
| N | Propeller rpm |
| $N_{n}, N_{r}$ | Number of acoustic modes or structural modes in frequency band $\Delta \omega$ |
| n | Symbolizes acoustic mode $n \equiv(q, i)$ |
| or |  |
| n | Number of circumferential wavelengths (or transverse half-wavelengths) in fuselage shell (or floor); see Eqs. (35-36) |
| n* | Number of terms in displacement series for fuselage shell (or floor) |
| $\mathrm{n}_{\mathrm{b}}, \mathrm{n}_{\mathrm{b}}^{\prime}$ | Number of boundary points on the fuselage shell (or floor) at which the acoustic eigenvectors are defined |
| $n_{n}$ | Modal density of acoustic modes |
| $n_{r}$ | Modal density of structural modes |

## LIST OF SYMBOLS

(Continued)

| $P\left(\omega_{n}\right)$ | Probability distribution function for $\omega_{n}$ in $\Delta \omega$ |
| :---: | :---: |
| $p\left(\omega_{n}\right)$ | Probability density function for $\omega_{n}$ in $\Delta \omega$ |
| $p_{b l}(\bar{x}, t)$ | Exterior pressure over the blocked (immobile) fuselage |
| $\left\langle\mathrm{p}_{\mathrm{b} 1}^{2}\right\rangle$ | Band-limited mean square blocked pressure |
| $\left\langle\mathrm{p}_{\dot{j}}^{2}(\bar{\xi}, \omega)\right\rangle_{t}$ | Interior mean square pressure at location $\bar{\xi}$ |
| $\left\langle p_{i}^{2}\right\rangle_{s, t}$ | Space-averaged band-limited mean square interior pressure |
| $\left\langle\mathrm{p}_{\mathrm{e}}^{2}\right\rangle_{s, t}$ | Space-averaged band-limited mean square exterior pressure for a reverberant field |
| $\left\langle p_{n}^{2}\right\rangle_{s, t}$ | Space-averaged band-limited mean square modal pressure, for nth mode in interior volume $V$ |
| $Q_{H}$ | Function defined in Section 3,2, Eq. (12) [1] |
| q | Number of axial half-wavelengths for acoustic mode $n \equiv(q, i)$ |
| $r$ | Symbolizes structural mode $r \equiv$ (M.N) |
| $r_{p}$ | Radial distance from center of fuselage cylinder to the axis of rotation of the propeller. |

## LIST OF SYMBOLS

(Continued)

| $\mathrm{R}_{\mathrm{p}_{\mathrm{b}}}\left(\mathrm{x} \mid \mathrm{x}^{\prime} ; \omega\right)$ | Average cross correlation of the exterior blocked pressure over the fuselage |
| :---: | :---: |
| S | Absorbing surface area of fuselage sidewall |
| $S_{e}$ | Absorbing surface area on each end surface (bulkhead) |
| $S_{p}(\omega)$ | Power spectral density of exterior pressure |
| $\mathrm{s}_{\mathrm{p}_{\mathrm{b}}}(\omega)$ | Power spectral density of exterior blocked pressure |
| $\mathrm{S}_{\mathrm{pbI}}\left(\overline{\mathrm{x}} \mid \overline{\mathrm{x}}^{\prime} ; w\right)$ | Cross spectral density of exterior blocked pressure |
| Si(z) |  |
| t | time |
| T | Period of rotation of propeller; $T=60 / \mathrm{N}$ |
| $\mathrm{T}_{0}, \mathrm{~T}_{1}$ | Period of propeller noise signature; $T_{0}=T_{1}=$ $(\mathrm{BPF})^{-1}=\mathrm{T} / \mathrm{E}$ |
| u | In-plane axial displacement of cylinder wall (or floor) |
| v | Volume of cavity |
| v | Circumferential (or transverse) displacement of cylinder wall (or floor) |
| $W_{\text {diss }}$ | Power dissipated on the cabin walls |

## LIST OF SYMBOLS

(Continued)

| $W_{\text {in }}$ | Net power inflow |
| :---: | :---: |
| $W_{r a d}^{\text {int }}(\omega)$ | Spectral density of power radiated by structure into interior acoustic space |
| $\mathrm{w}_{\mathrm{abs}}^{\operatorname{int}}(\omega)$ | Spectral density of power absorbed on inner wall of the space from interior acoustic field |
| w | Cylinder wall (or floor) normal displacement |
| x | Transverse coordinate; see Figure 19 |
| $\overline{\mathrm{x}}$ | Location on exterior surface of fuselage |
| $\bar{x}_{\mathrm{m}}$ | Location of grid point on exterior surface of fuselage |
| $x_{1} x_{2} x_{3}$ | Local coordinate systems used for PROPFAN propeller noise prediction; see Figure E-6 |
| y | Vertical coordinate, relative to fuselage centerline; see Figure 17 |
| $z$ | Axial coordinate, relative to forward end of the fuselage structure (of length L); see Figure 16 |
| $z_{k}$ | Axial coordinate for grid point $k$, see Figure 16 |
| $z_{p}$ | Location of propeller relative to the forward end of the fuselage structure (of length $L$ ); see Figure 17 |


|  | LIST OF SYMBOLS <br> (Continued) |
| :---: | :---: |
| $\bar{\alpha}$ | Band average absorption coefficient |
| $\gamma$ | Incidence angle between propeller and location $\bar{x}_{\mathrm{m}}$; see Figure 17 |
| $\Delta=\frac{\pi a}{18}$ | Grid spacing for propeller noise predictions |
| $\Delta A=\Delta^{2}$ | Area associated with each grid point; see Figure 16 |
| $\Delta \omega$ (radians/sec) Frequency band of width $\Delta \omega=c_{\omega} \omega$ <br> $n<\Delta \omega$ symbolizes modes resonant below band <br> $n \varepsilon \Delta \omega$ symbolizes modes resonant inside band <br> $n>\Delta \omega$ symbolizes modes resonant above band |  |
| $\delta()$ | Delta function |
| $\varepsilon_{n}=\mathrm{V} / \int_{\mathrm{V}} \phi_{\mathrm{n}}^{2} \overline{\mathrm{~V}}$ ( Acoustic mode normalization factor |  |
| $\varepsilon_{q}$ | Acoustic mode normalizing factor in axial direction (see Eq.C.II [1]) |
| $\zeta$ | Transverse coordinate; see Section 3.6 [I] |
| $n_{n}$ | Acoustic mode loss factor |
| $n_{r}$ | Structural mode loss factor |
| $\eta_{r}^{\prime}$ | Structural loss factor, including damping due to trim; Eq.(17) |
| $\eta_{r}^{\prime \prime}$ | Internal radiation loss factor, due to closely coupled structural and acoustic modes; Eq.(83) |

## LIST OF SYMBOLS

(Continued)

| $\bar{n}_{n}$ | Average one-third octave band acoustic mode loss factor |
| :---: | :---: |
| $\bar{\eta}_{r}=\eta_{r}^{\prime}+\eta_{r}^{\prime}$ | +next Average one-third octave band structural mode loss factor |
| $\eta_{\mathrm{r}}^{\text {struc }}$ | Average one-third octave band structural loss factor |
| $n_{r}^{\mathrm{rad}}$ | Average one-third octave band radiation loss factor |
| $\eta_{\text {rad }}^{\text {int }}, \eta_{\text {rad }}^{\text {ext }}$ | Average one-third octave band internal and external radiation loss factors defined after Eq. (19) [1] |
| $\theta$ | Angular coordinate, relative to fuselage bottom centerline; see Figure 17 |
| $\theta_{2}$ | Angular coordinate for grid location ( $k, \ell$ ) |
| $\theta_{0}$ | Angle at which fuselage shell/floor joint is located |
| $\theta_{j}$ | Angle $\theta$ for point $j$ on fuselage wall, a boundary point for the acoustic eigenvectors |
| $\theta^{\prime} j^{\theta} 2 j$ | Angles defining mid-points between boundary point $j$ and adjacent boundary points |
| $\xi$ | Axial coordinate; see Section 3.6[1] |
| or |  |
| $\xi$ | Conductance for trim on end surface of cylinder interior |

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## LIST OF SYMBOLS

(Continued)


## LIST OF SYMBOLS

(Continued)

| $\phi_{\mathrm{n}}(\bar{\xi}) \equiv \phi^{\prime}$ | $\bar{\xi})$ Mode shape, or eigenfunction, of the nth mode of the cavity at location $\bar{\xi}$ |
| :---: | :---: |
| $\phi_{i}\left(\theta_{j}\right)$ | Mode shape of ith acoustic mode of the fuselage cross-section evaluated on the fuselage wall at location $j$, angle $\theta_{j}$ |
| $\Psi_{G}(\mathrm{r}, \mathrm{H})$ | Generalized modal forcing function due to propeller noise, mode $r$ at propeller harmonic $H$; see Section 3.4 [I] |
| $\psi^{r}(\bar{x})$ | Mode shape, or eigenfunction, of the rth mode of the structure, at location $\bar{x}$ |
| $\psi_{p}^{r}(z, x)$ | Floor displacement in structure mode $r$ |
| $\psi_{S}^{r}(z, \theta)$ | Fuselage shell displacement in structure mode $r$ |
| $\omega$ | Angular frequency (rads/sec) |
| ${ }^{\omega} \mathrm{H}$ | Angular frequency of propeller harmonic H |
| $\omega_{n}$ | Acoustic mode resonance angular frequency |
| $\omega_{r}$ | Structure mode resonance angular frequency |
| $\left\rangle_{s, t}\right.$ | Band-limited, space-averaged and time-averaged value |


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[^0]:    * $\tau_{F}$ can be used to represent vibration isolation mounts supporting the floor.

