

JAN 6 1964

UNCLASSIFIED

MASTER

CLASSIFICATION LEVEL
(S, C OR U)



ATOMICS INTERNATIONAL
A Division of North American Aviation Inc.

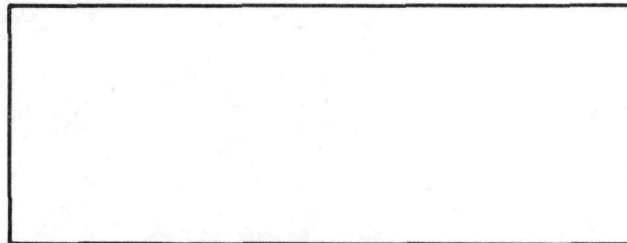
NAA-SR-MEMO COVER SHEET

REPORT TITLE

AUTHOR

NAA-SR-MEMO 10421

(This Document Contains 87 Pages.
This is Copy of Series)



CLASSIFICATION TYPE
(RD OR DI)

NAA-SR-MEMOs are working papers and may be expanded, modified, or withdrawn at any time, and are intended for internal use only.

THIS REPORT MAY NOT BE PUBLISHED WITHOUT THE APPROVAL OF THE PATENT BRANCH, AEC.

LEGAL NOTICE

This report was prepared as an account of Government sponsored work. Neither the United States, nor the Commission, nor any person acting on behalf of the Commission:

A. Makes any warranty or representation, express or implied, with respect to the accuracy, completeness, or usefulness of the information contained in this report, or that the use of any information, apparatus, method, or process disclosed in this report may not infringe privately owned rights; or

B. Assumes any liabilities with respect to the use of, or for damages resulting from the use of information, apparatus, method, or process disclosed in this report.

As used in the above, "person acting on behalf of the Commission" includes any employee or contractor of the Commission, or employee of such contractor, to the extent that such employee or contractor of the Commission, or employee of such contractor prepares, disseminates, or provides access to, any information pursuant to his employment or contract with the Commission, or his employment with such contractor.

UNCLASSIFIED

CLASSIFICATION LEVEL
(S, C OR U)

DO NOT REMOVE THIS SHEET

DISCLAIMER

This report was prepared as an account of work sponsored by an agency of the United States Government. Neither the United States Government nor any agency Thereof, nor any of their employees, makes any warranty, express or implied, or assumes any legal liability or responsibility for the accuracy, completeness, or usefulness of any information, apparatus, product, or process disclosed, or represents that its use would not infringe privately owned rights. Reference herein to any specific commercial product, process, or service by trade name, trademark, manufacturer, or otherwise does not necessarily constitute or imply its endorsement, recommendation, or favoring by the United States Government or any agency thereof. The views and opinions of authors expressed herein do not necessarily state or reflect those of the United States Government or any agency thereof.

DISCLAIMER

Portions of this document may be illegible in electronic image products. Images are produced from the best available original document.

7

ATOMICS INTERNATIONAL <small>A Division of North American Aviation, Inc.</small>			NAA SR TDR NO. 10421	APPROVALS
TECHNICAL DATA RECORD			PAGE 1 OF 87	
AUTHOR R. H. Norman <i>R.H. Norman</i> D. G. Oliver <i>D.G. Oliver</i> W. R. Castle <i>wrcastle</i>	DEPT. & GROUP NO. 727-75 727-75 727-75	DATE 9-9-64	GO NO. 7611	<i>K. E. Buttrey</i>
TITLE A Method of Calculating Annular Void Reactivity Effects in Cylindrical Reactors		S/A NO. 2321	TWR	
		SECURITY CLASSIFICATION		
		<small>(CHECK ONE BOX ONLY)</small>	<small>(CHECK ONE BOX ONLY)</small>	
PROGRAM Aerospace Safety		SUBACCOUNT TITLE Critical Configuration Analysis	UNCL. <input checked="" type="checkbox"/>	RESTRICTED DATA <input type="checkbox"/>
DISTRIBUTION			CONF. <input type="checkbox"/>	DEFENSE INFO. <input type="checkbox"/>
5 D. Arnold - 727 2 E. Ash - 726 2 R. Balent - 720 2 R. Brehm - 791 2 J. Brunings - 726 5 K. Buttrey - 727 1 G. Calkins - 735 5 D. Cockeram - 720 2 R. Courson - 720 5 R. Cranor - 727 5 R. Cummings - 727 5 W. Davis - 727 5 R. Detterman - 727 2 H. Dieckamp - 720 5 R. Elliott - 727 2 W. Engle - 791 5 E. Faelten - 727 5 S. Fields - 727 5 R. Gimera - 727 2 R. Hart - 726 5 W. Henoch - 724 5 C. Johnson - 727 5 R. Johnson - 727 2 A. Martin - 720 5 L. Maki - 727 5 D. McGoff - 727 2 J. Miller - 791 5 L. Mims - 727 5 L. Moss - 727 5 A. Piccot - 727 2 W. Roberts - 791 5 E. Rosell - 727 5 J. Rohlfs - 727 5 V. Sanders - 727 5 W. Scott - 724 5 R. Simms - 727 5 R. Tobler - 727 5 J. Susnir - 727 5 R. Wallerstedt - 722 5 E. Weisner - 720 5 A. Weitzberg - 727 5 R. Williams - 720 2 R. Wilson - 722	STATEMENT OF PROBLEM Develop a method of calculating void sleeve reactivity worths.	AUTHORIZED CLASSIFIER SIGNATURE <i>K. E. Buttrey</i>	DATE 9/28/64	
		ABSTRACT: A problem of increasing concern in reactor neutronic analysis is the criticality calculation of a reactor containing a void gap. This report presents the development of a method of calculating annular void sleeve reactivity worths in SNAP reactors. A code was developed to compute a geometrical escape probability and absorption cross section (DB ²) for the void region. This DB ² is derived from S _n transport theory and characterizes axial leakage from the void in reactivity calculations with the DTK S _n transport code. Calculations of void sleeve reactivity worths by this method for the water immersed S8ER core compare very well with measurements.		
Nuclear & Dynamics Unit (8)				

A Method of Calculating Annular Void
Reactivity Effects in Cylindrical Reactors

<u>Contents</u>	<u>Page</u>
1. INTRODUCTION	3
2. CALCULATION OF ESCAPE PROBABILITIES	
a. Escape Probability Model	4
b. Derivation of Equations	4
c. Angular Distribution	15
d. Solution of Equations	15
e. Isotropic Escape Probabilities	16
3. CALCULATION OF DB^2	
a. S_n Transport Theory Model	19
b. Derivation and Solution of DB^2 Equation	19
4. COMPARISON OF REACTIVITY CALCULATIONS WITH MEASUREMENTS	24
5. CONCLUSIONS AND FURTHER WORK	26
REFERENCES	27
APPENDIX - The FOSDICK Code	
a. Description	28
b. Code Equations	29
c. Program Listing	33
d. Sample Problem	57

1. INTRODUCTION

One of the areas of much difficulty and increasing concern to reactor neutronic analysts is the treatment of void regions where neutron streaming occurs. In heterogeneous lattices, methods have been developed through the use of collision probabilities in integral transport theory. Benoist⁽¹⁾ uses integral transport theory to obtain radial and transverse diffusion coefficients in terms of "directed collision probabilities", considering the lattice as a whole. Leslie's method⁽²⁾ of weighting diffusion coefficients is restricted to a single lattice cell and considers a "tilted source" in the moderator regions. These techniques represent analytical closed-form solutions to the problem of calculating L^2 , but do not apply to our problem in that we are considering small reactors (i.e., B^2 is large), with a homogeneous mixture of fuel and moderator. The method presented in this report is based on purely geometrical considerations of a cylindrical annular void and is extended to the calculation of reactivity effects by characterizing the transverse leakage from the void by a macroscopic absorption cross section in an S_n transport theory criticality code. Although D and B^2 are undefined in a void, the product DB^2 may be used as an absorption cross section because the one-dimensional transport theory does not distinguish between a neutron lost by absorption and a neutron lost by transverse leakage. At the present time this method is divided into two parts: (1) the geometrical probability that a neutron leaving the core is travelling in the right direction to escape out the ends of an annular void between the core and reflector, and (2) the calculation of an effective absorption cross section (i.e., DB^2) to be used in a one-dimensional S_n transport theory criticality calculation. A comparison is made between void reactivity worth calculations and measurements from water immersion experiments of the S8ER core with surrounding void sleeves.

2. CALCULATION OF ESCAPE PROBABILITIES

a. Escape Probability Model

Considering a single neutron at any point along the radial surface of a cylindrical core surrounded by an annular void gap, we are primarily concerned with whether or not the neutron will escape from the gap in the transverse direction, i.e., the escape probability of a distribution of neutrons on the core boundary. We are not concerned with neutrons after they reach the reflector; what happens there is treated by the S_n transport theory. There is another escape probability of interest, however, namely that for a neutron on the inner radial surface of the reflector. Again, those neutrons, having left the reflector, which do not escape out the ends of the gap are treated by the S_n transport theory within the core. The problem is to define an escape probability which would account for transverse leakage from the gap for neutrons on both boundaries of the void. The present model assumes the escape probability for a neutron on the core boundary and the reflector boundary to be the same. Both core neutron and reflector neutron escape probabilities are derived and were calculated, but only the core neutron escape probability is used in the derivation of DB^2 for reactivity calculations. Several factors that affect this assumption are: the greater importance of reflector neutrons, the lower density of reflector neutrons, and the larger escape probability for reflector neutrons. The assumption, in effect, is saying that these factors nullify one another. Neutron "importance" considerations have been neglected to the extent that a spatial fission neutron density distribution is used to weight the escape probability along the length of the core, rather than an "importance" distribution (e.g., the product of flux times adjoint flux).

b. Derivation of Equations

The escape probability of a neutron on the boundary of a void gap is simply defined as the probability of the neutron travelling in the direction of the open ends of the gap. The number of neutrons travelling in directions $d\omega$ about ω is simply: $N(\phi, \theta) d\omega$, where $N(\phi, \theta)$ is the number of neutrons in solid angle $d\omega = \sin \phi d\theta d\phi$.

The differential probability of neutrons leaving the surface in angles $d\omega$ about ω is then

$$dP = \frac{N(\varphi, \theta) d\omega}{\int_0^{2\pi} N(\varphi, \theta) d\omega} \quad (1a)$$

Considering one quadrant of the sphere in Figure 1, with origin of the neutron at a distance z_0 from the top of the core, φ is defined as the angle between the z -axis and the flight path of the neutron, and θ is defined as the angle (in the x - y plane) between the x -axis and the flight path of the neutron. The escape geometry is symmetrical about $\theta = 0$ (the x -axis) such that only the quadrant defined by $\varphi = 0$ to π radians, and $\theta = 0$ to $\pi/2$ radians need be considered. Looking at the upper octant of this quadrant in detail we see a line of natural division in the latitude φ_1 . This is the projection on the sphere of the point on the rim of the reflector nearest to the core neutron (origin of the flight path). The latitude φ_2 defines the most distant point on the rim of the reflector. From $\varphi = 0$ to φ_1 , φ is constant for any angle θ , therefore the probability of escape through this region is simply the integral of the differential probability:

$$P_1 = \frac{\int_0^{\varphi_1} \int_0^{\pi/2} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi}{\int_0^{\pi} \int_0^{\pi/2} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi} \quad (1b)$$

To simplify the writing of escape probabilities henceforth, let us define the total number of neutrons emerging from the surface by symbol A ,

$$A \equiv \int_0^{\pi} \int_0^{\pi/2} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi \quad (1c)$$

Between φ_1 , and φ_2 , due to the different location of the origin of the void boundary radii and the origin of the neutron flight path (see Figures 1 and 2) there arises an area with one boundary on which θ

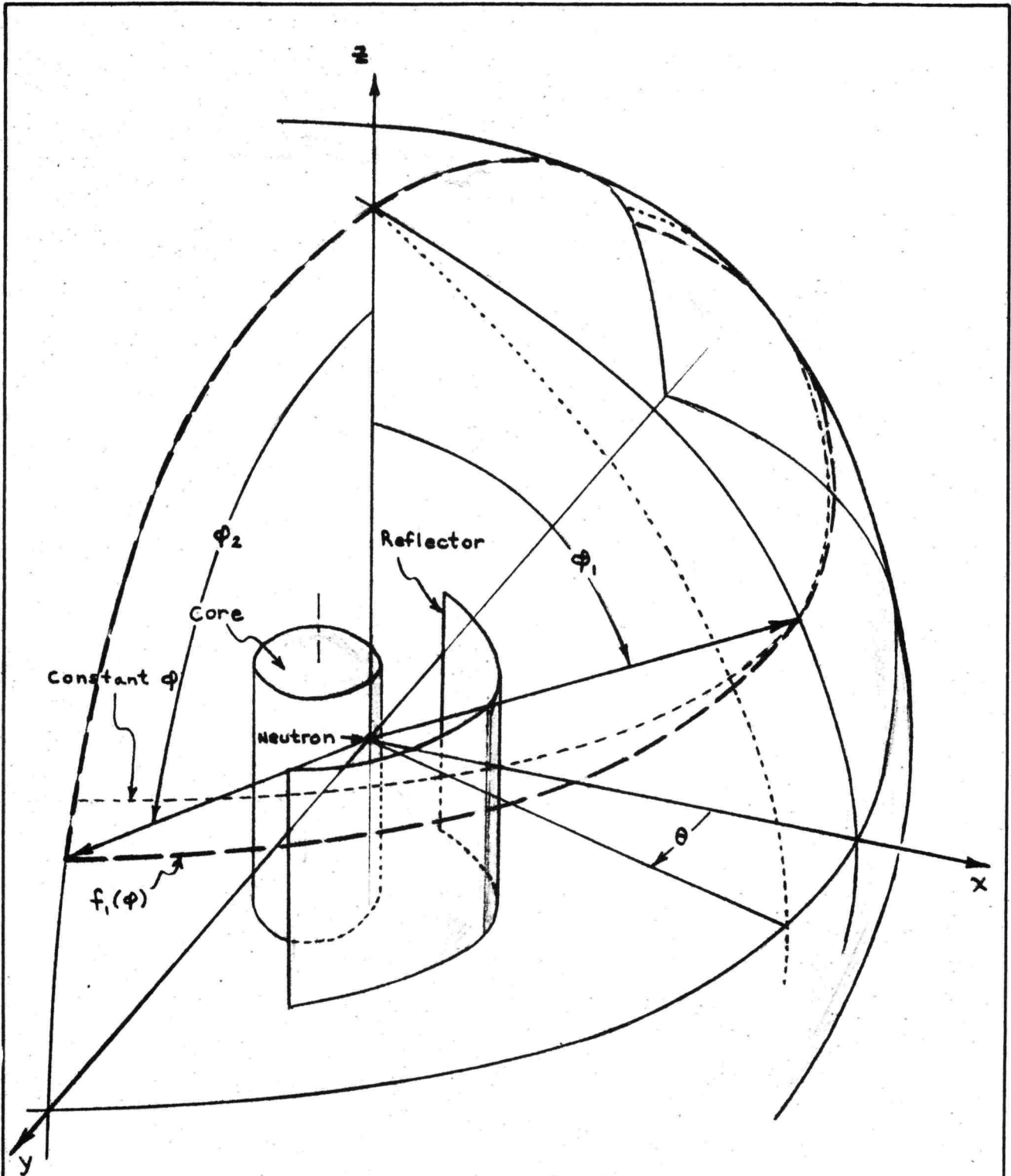


FIGURE 1. Illustration of the Probability of a Neutron Emitted from the Core Escaping from an Annular Void Gap.

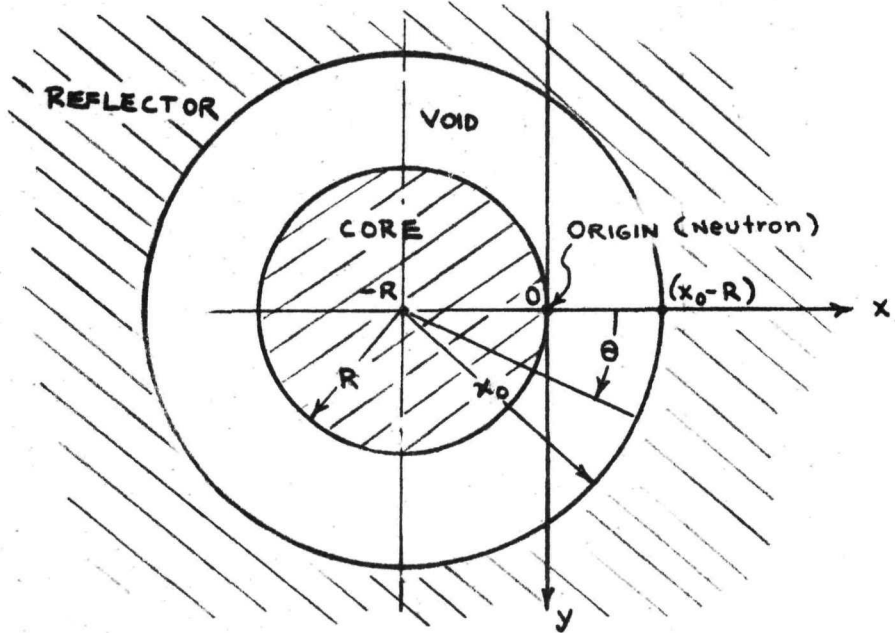


FIGURE 2. Top View of Void Gap.

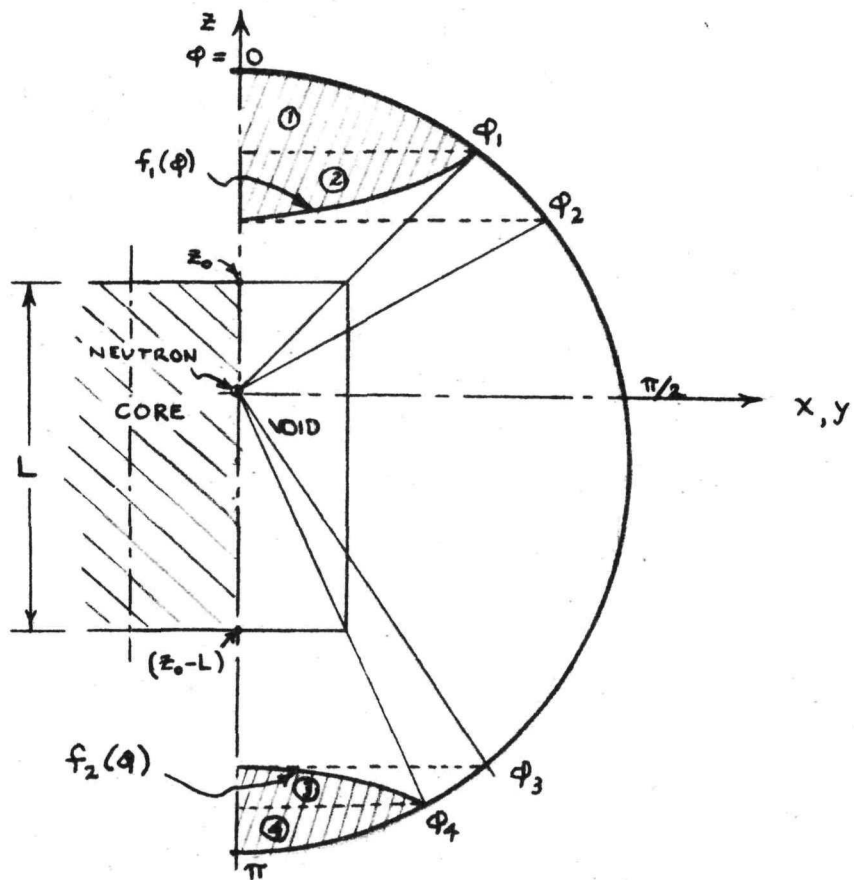


FIGURE 3. Side View of the Limits of Integration over ϕ .

varies with ϕ . To derive this relationship, $\theta = f_1(\phi)$, we begin by considering a cylindrical core of radius R within a reflector of radius x_0 (see Figure 2). With the distance from the origin of the neutron flight path to the top of the core given by z_0 , and a core length of L , the equations of the cylinder of radius x_0 with origin at $-R$ are

$$(x + R)^2 + y^2 = x_0^2 \quad (2a)$$

$$z = z_0 \quad (2b)$$

Transforming to spherical coordinates (ρ, θ, ϕ) by

$$\left. \begin{aligned} x &= \rho \cos \theta \sin \phi \\ y &= \rho \sin \theta \sin \phi \\ z &= \rho \cos \phi \end{aligned} \right\} \quad (2c)$$

equation (2a) becomes

$$\rho^2 \cos^2 \theta \sin^2 \phi + 2R\rho \cos \theta \sin \phi + \rho^2 \sin^2 \theta \sin^2 \phi = x_0^2 - R^2 \quad (2d)$$

and equation (2b) becomes

$$\rho = \frac{z_0}{\cos \phi} \quad (2e)$$

Putting (2e) into (2d), and solving for θ ,

$$\theta = f_1(\phi) = \cos^{-1} \left[\left(\frac{x_0^2 - R^2}{2Rz_0} \right) \cot \phi - \frac{z_0}{2R} \tan \phi \right] \quad (3a)$$

The second probability in the upper octant is therefore

$$P_2 = \frac{1}{A} \int_{\phi_1}^{\phi_2} \int_{f_1(\phi)}^{\pi/2} \sin \phi N(\phi, \theta) d\theta d\phi \quad (4)$$

Correspondingly, in the lower octant we have

$$P_3 = \frac{1}{A} \int_{\varphi_3}^{\varphi_4} \int_{f_2(\varphi)}^{\pi/2} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi \quad (5)$$

and

$$P_4 = \frac{1}{A} \int_{\varphi_4}^{\pi} \int_0^{\pi/2} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi \quad (6)$$

where, in the lower octant, $f_2(\varphi)$ is obtained by simply replacing z_0 in the equation (3a) with $(z_0 - L)$,

$$f_2(\varphi) = \cos^{-1} \left[\left(\frac{x_0^2 - R^2}{2R(z_0 - L)} \right) \cot \varphi - \frac{(z_0 - L)}{2R} \tan \varphi \right] \quad (3b)$$

In the present model the angular distribution $N(\varphi, \theta)$ is assumed separable in angles φ and θ , and only the φ angular dependence is accounted for,

$$N(\varphi, \theta) = a \cdot N'(\varphi) \quad (7)$$

where "a" is a constant. (The θ angular dependence remains to be investigated.) Then the total escape probability for a neutron leaving the core at a given point z_0 , is

$$\begin{aligned}
 P_c(z_0) = & \frac{1}{2} \int_0^{\theta_1} \sin \theta \cdot N(\theta) \cdot d\theta \\
 & + \frac{1}{\pi} \int_{\theta_1}^{\theta_2} \sin \theta \cdot \left\{ \frac{\pi}{2} - \cos^{-1} \left[\left(\frac{x_0^2 - R^2}{2Rz_0} \right) \cot \theta - \frac{z_0}{2R} \tan \theta \right] \right\} \cdot N(\theta) d\theta \\
 & + \frac{1}{\pi} \int_{\theta_3}^{\theta_4} \sin \theta \cdot \left\{ \frac{\pi}{2} - \cos^{-1} \left[\left(\frac{x_0^2 - R^2}{2R(z_0 - L)} \right) \cot \theta - \frac{(z_0 - L)}{2R} \tan \theta \right] \right\} \cdot N(\theta) d\theta \\
 & + \frac{1}{2} \int_{\theta_4}^{\pi} \sin \theta \cdot N(\theta) \cdot d\theta
 \end{aligned} \tag{8}$$

Where the normalized angular distribution, $N(\theta)$, is defined by

$$N(\theta) \equiv \frac{a \cdot N'(\theta)}{A} \tag{9a}$$

and the limits of integration are (see Figure 3)

$$\left. \begin{aligned} \varphi_1 &= \tan^{-1} \left(\frac{x_0 - R}{z_0} \right) \\ \varphi_2 &= \tan^{-1} \left(\frac{\sqrt{x_0^2 - R^2}}{z_0} \right) \\ \varphi_3 &= \tan^{-1} \left(\frac{\sqrt{x_0^2 - R^2}}{z_0 - L} \right) \\ \varphi_4 &= \tan^{-1} \left(\frac{x_0 - R}{z_0 - L} \right) \end{aligned} \right\} \quad (9b)$$

For a neutron leaving the reflector at a given point z_0 , the escape probability, $P_R(z_0)$, is the sum of six probabilities. Figure 4 shows the three areas of consideration in the upper octant (there are correspondingly three more in the lower octant). The sum of probabilities for a neutron leaving the reflector can be written

$$\begin{aligned} P_R(z_0) = \frac{1}{A} & \left[\int_0^{\varphi_1} \int_0^{\theta_0} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi \right. \\ & + \int_{\varphi_1}^{\varphi_2} \int_{f_3(\varphi)}^{\theta_0} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi \\ & + \int_0^{\varphi_5} \int_{\theta_0}^{f_4(\varphi)} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi \\ & + \int_{\varphi_6}^{\pi} \int_{\theta_0}^{f_5(\varphi)} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi \\ & + \int_{\varphi_3}^{\varphi_4} \int_{f_6(\varphi)}^{\theta_0} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi \\ & \left. + \int_{\varphi_4}^{\pi} \int_0^{\theta_0} \sin \varphi \cdot N(\varphi, \theta) d\theta d\varphi \right] \end{aligned} \quad (10)$$

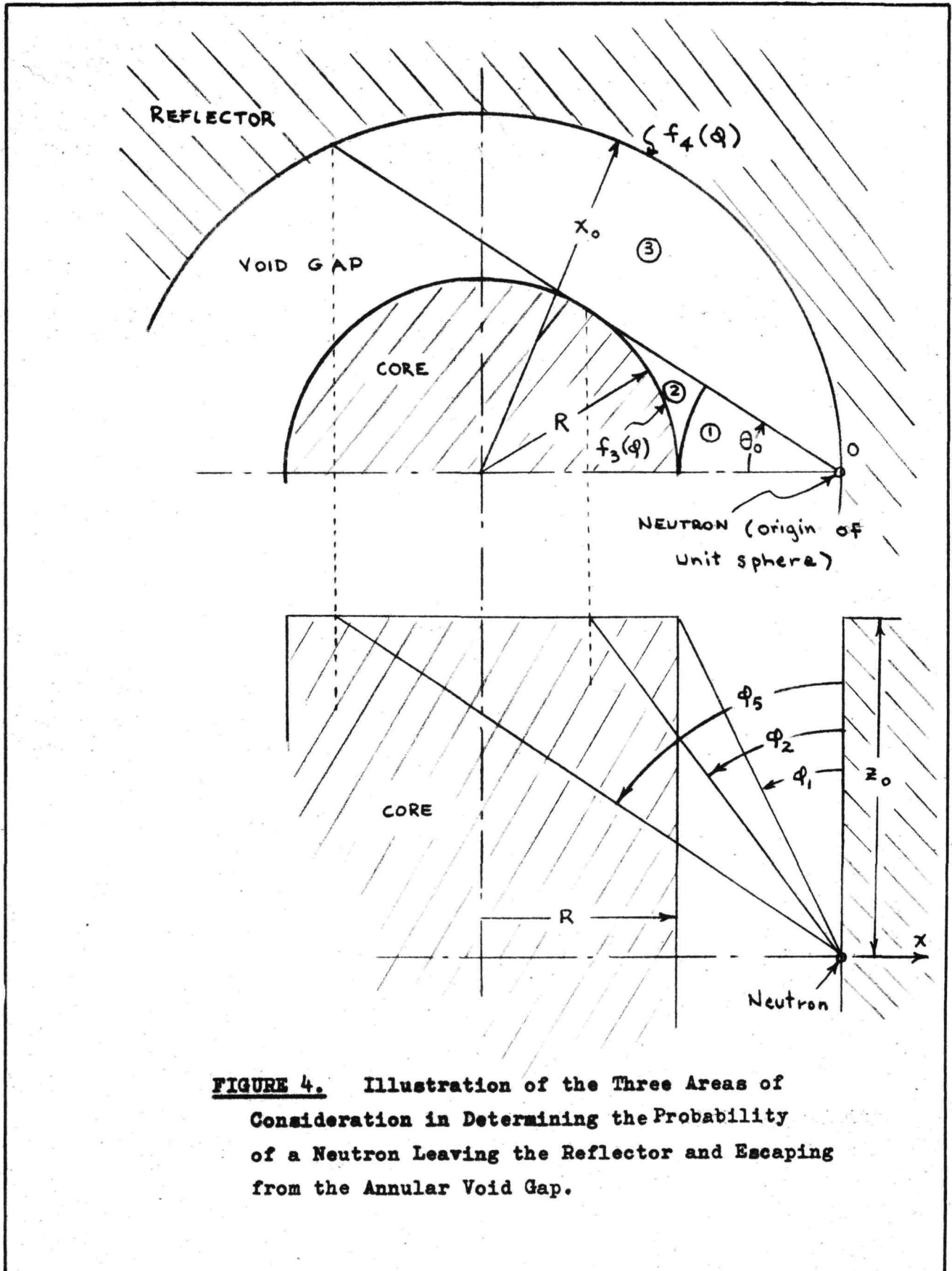


FIGURE 4. Illustration of the Three Areas of Consideration in Determining the Probability of a Neutron Leaving the Reflector and Escaping from the Annular Void Gap.

where

$$\left. \begin{aligned} \theta_0 &= \tan^{-1} \left(\frac{R}{\sqrt{x_0^2 - R^2}} \right) \\ \varphi_5 &= \tan^{-1} \left(\frac{2\sqrt{x_0^2 - R^2}}{z_0} \right) \\ \varphi_6 &= \tan^{-1} \left(\frac{2\sqrt{x_0^2 - R^2}}{z_0 - L} \right) \end{aligned} \right\} \quad (11)$$

and $\varphi_1, \varphi_2, \varphi_3, \varphi_4$ are given by equations (9b). To derive the functions $f_3(\varphi)$ and $f_6(\varphi)$, we write down the equations of a cylindrical core of radius R with origin at x_0 (the location of the origin of the neutron flight path in the void)

$$(x - x_0)^2 + y^2 = R^2 \quad (12a)$$

$$z = z_0 \quad (12b)$$

and transform to spherical coordinates by equations (2c), solving for θ ,

$$\theta = f_3(\varphi) = \cos^{-1} \left[\frac{x_0^2 - R^2 + z_0^2 \tan^2 \varphi}{2 z_0 x_0 \tan \varphi} \right] \quad (13)$$

and

$$f_6(\varphi) = \cos^{-1} \left[\frac{x_0^2 - R^2 + (z_0 - L)^2 \tan^2 \varphi}{2 (z_0 - L) x_0 \tan \varphi} \right] \quad (14)$$

The functions $f_4(\varphi)$ and $f_5(\varphi)$ are derived from the equations of a cylindrical reflector of radius x_0 with origin at x_0 ,

$$(x - x_0)^2 + y^2 = x_0^2 \quad (15a)$$

$$z = z_0 \quad (15b)$$

Transforming to spherical coordinates by equations (2c), and again solving for θ ,

$$\theta = f_4(\varphi) = \cos^{-1} \left[\frac{z_0 \tan \varphi}{2x_0} \right] \quad (16)$$

and

$$f_5(\varphi) = \cos^{-1} \left[\frac{(z_0 - L) \tan \varphi}{2x_0} \right] \quad (17)$$

Again, accounting for the φ angular dependence only, by equation (7), the total escape probability for a neutron leaving the reflector at point z_0 is

$$\begin{aligned} P_R(z_0) = & \frac{\theta_0}{\pi} \int_0^{\varphi_1} \sin \varphi \cdot N(\varphi) \cdot d\varphi + \frac{1}{\pi} \int_{\varphi_1}^{\varphi_2} \sin \varphi \cdot \left\{ \theta_0 - \right. \\ & \left. \cos^{-1} \left[\frac{x_0^2 - R^2 + z_0^2 \tan^2 \varphi}{2z_0 x_0 \tan \varphi} \right] \right\} \cdot N(\varphi) \cdot d\varphi + \frac{1}{\pi} \int_0^{\varphi_5} \sin \varphi \cdot \left\{ \cos^{-1} \left[\frac{z_0 \tan \varphi}{2x_0} \right] - \theta_0 \right\} \cdot N(\varphi) \cdot d\varphi \\ & + \frac{1}{\pi} \int_{\varphi_6}^{\pi} \sin \varphi \cdot \left\{ \cos^{-1} \left[\frac{(z_0 - L) \tan \varphi}{2x_0} \right] - \theta_0 \right\} \cdot N(\varphi) \cdot d\varphi + \frac{1}{\pi} \int_{\varphi_3}^{\varphi_4} \sin \varphi \cdot \left\{ \theta_0 - \right. \\ & \left. \cos^{-1} \left[\frac{x_0^2 - R^2 + (z_0 - L)^2 \tan^2 \varphi}{2(z_0 - L) x_0 \tan \varphi} \right] \right\} \cdot N(\varphi) \cdot d\varphi + \frac{\theta_0}{\pi} \int_{\varphi_4}^{\pi} \sin \varphi \cdot N(\varphi) \cdot d\varphi \quad (18) \end{aligned}$$

where the angular distribution normalization, the limits of integration, and Θ_0 are given by equations (9a), (9b), and (11).

The average escape probability over the entire core length $z_0 = 0$ to L , for a gap thickness ($x_0 - R$), is given by

$$\bar{P}_c = \frac{\int_0^L P_c(z_0) \cdot F(z_0) dz_0}{\int_0^L F(z_0) dz_0} \quad (19a)$$

where $P_c(z_0)$ and $F(z_0)$ are the escape probability and fission neutron density along the length of the core.

c. Angular Distribution

The angular distribution function, $N(\varphi)$, is obtained from a high order S_n transport calculation* of the reactor configuration under consideration. The angular fluxes are summed over all energy groups for each discrete value of φ and normalized by equation (9a). It was found necessary to use at least 12th order in the S_n approximation to obtain a suitable angular distribution, S_{12} yielding 6 discrete φ angles in the forward direction between 0 and $\pi/2$ radians. The energy dependence of the angular distribution was found to be quite weak at the boundaries of the void gap, on the basis of S_4 calculations with 16 energy groups from 10 Mev to .0092 ev. Although energy dependence has been neglected in the present model, it would be simple, in principle, to calculate energy-dependent escape probabilities. The angular distribution at the void boundaries also depends on the gap thickness, but in the present model a single representative distribution is used for all gaps.

d. Solution of Equations

Equation (8) has been programmed in FORTRAN II language for the IBM 7094 computer. The integrations are performed by the 5th order Legendre-Gauss quadrature method.⁽³⁾ The first term in equation (8), for example, is computed as:

$$\frac{1}{2} \int_0^{\varphi_1} \sin \varphi \cdot N(\varphi) \cdot d\varphi = \frac{1}{2} \left\{ \frac{\varphi_1 - 0}{2} \sum_{k=1}^5 \sin(\varphi_k) N(\varphi_k) w_k \right\},$$

* More detail is found in Section 3.

where
$$\varphi_k = \left(\frac{\varphi_1 + 0}{2} \right) + \left(\frac{\varphi_1 - 0}{2} \right) \cdot \xi_k$$

ξ_k are the zeros of Legendre polynomials, and w_k are the weights (Christoffel numbers). The angular weighting factor, $N(\varphi_k)$, used in any particular numerical integration by the code is obtained from a linear interpolation in $\sin \varphi_k$ of normalized angular fluxes. The end points of the distribution are linear (in $\sin \varphi_k$) extrapolations of the normalized angular fluxes. The integrals in equation (19a) are simply programmed as summations,

$$\bar{P}_c = \frac{\int_0^L P_c(z_0) F(z_0) dz_0}{\int_0^L F(z_0) dz_0} \approx \frac{\sum_{i=1}^m P_{ci} F_i}{\sum_{i=1}^m F_i} \quad (19b)$$

where m is the number of equal interval mesh points. Equations (8) and (19b) are a part of the computer program designated FOSDICK; the details of the entire program appear in the Appendix.

e. Isotropic Escape Probabilities

A special case of the escape probability as already derived arises if the angular distribution of neutrons on the boundary of the void is isotropic, i.e., all directions of travel are equally probable. In this case the probability of a neutron escaping from the gap is the area projected onto the surface of a unit (area) sphere of the space between the core and reflector as seen by a neutron on the surface of the core or reflector. The equations (8) and (18) are simply modified by taking $N(\varphi) = 1$. Although an isotropic distribution of neutrons at the void boundary is unrealistic, the isotropic escape probability was calculated for the sake of comparison. Figure 5 shows the effect

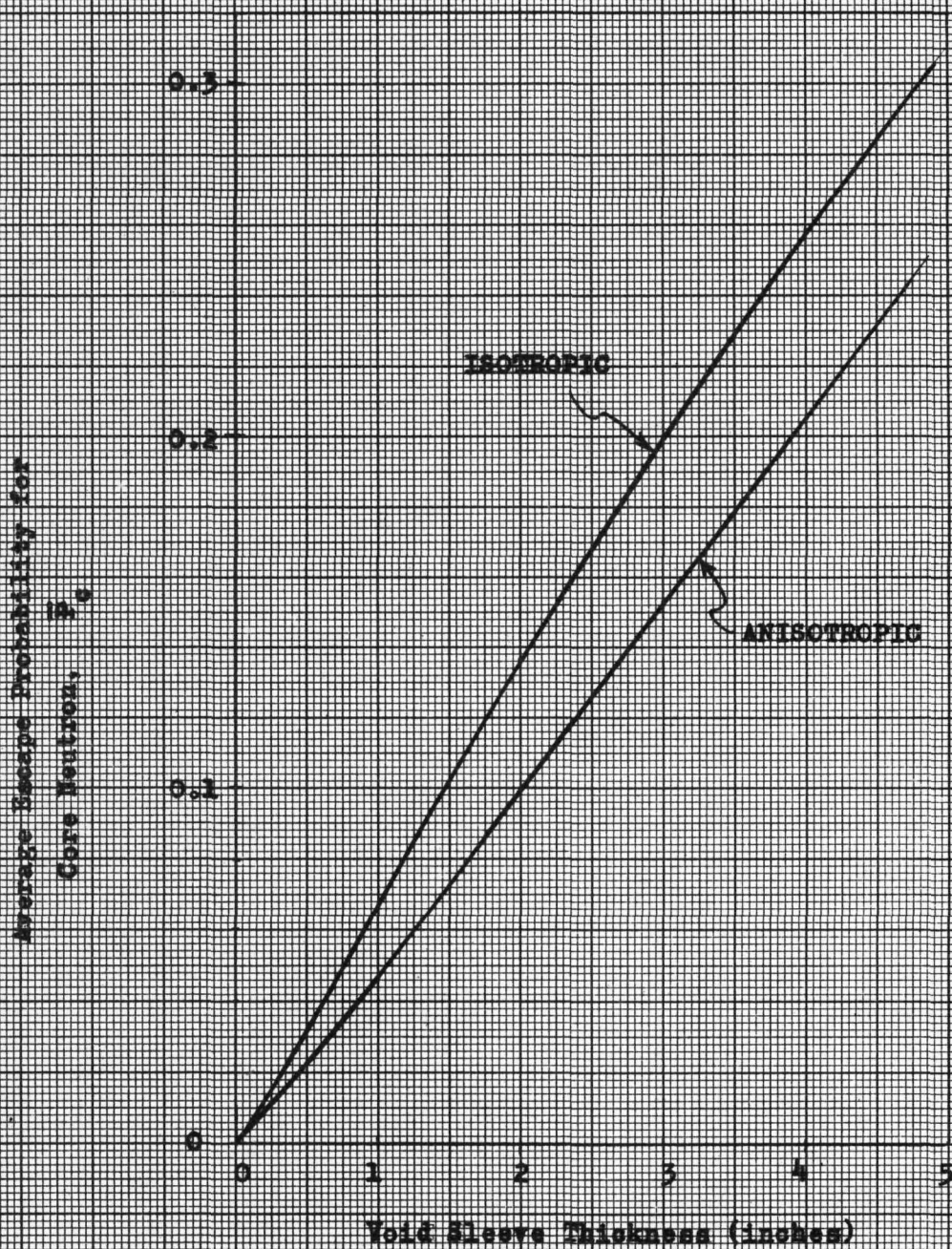


FIGURE 5. Comparison of Anisotropic and Isotropic Escape Probability for S₁ER Model.

ATOMICS INTERNATIONAL
A Division of North American Aviation, Inc.

NO. NAA SR TDR 10421
DATE 9-9-64
PAGE 18 OF 87

of angular distribution in the S8ER void-sleeve model, where an S_{12} calculated angular distribution for the 3-inch void thickness was used for all cases. The anisotropic curve is significantly lower than the isotropic curve showing the effect of the strong peaking of angular neutron flux at a core/void boundary in the forward direction.

[Equation (18) was programmed for $N(\varphi) = 1$, and showed the escape probability for a neutron on the reflector surface to be significantly greater than that for neutrons on the core surface. The present model does not consider P_R , however. It remains to obtain an "effective" escape probability for both core and reflector neutrons.]

3. CALCULATION OF DB²

a. S_n Transport Theory Model

The one-dimensional DTK S_n transport code ⁽⁴⁾ is used to calculate angular fluxes at the void boundaries. The discrete directions are determined by the order of the S_n approximation and the geometry, S₁₂ having 21 discrete directions per spherical octant in cylindrical geometry. The S₄ approximation has 3 discrete directions per spherical octant in cylindrical geometry. This means there are 6 values of ϕ from 0 to $\pi/2$ radians in S₁₂, and only 2 in S₄. To account for transverse leakage from the void gap a macroscopic absorption cross section (or DB²) is assigned to it. The DTK code is limited in that it does not use angular-dependent absorption cross sections, therefore an average DB² is calculated to represent the total leakage in all directions from the gap. An energy-dependent DB² could be calculated but, as previously mentioned, the angular distribution at the core/void interface is not strongly energy-dependent.

b. Derivation and Solution of DB² Equation

The escape probability for neutrons at the void boundary of surface area A_i is defined as the total number of neutrons per second escaping from the gap divided by the total number of neutrons per second entering the gap at A_i,

$$\bar{P}_{ci} \equiv \frac{J_i^+ A_i - J_{i+1}^+ A_{i+1}}{J_i^+ A_i} \quad (22)$$

where "i+1" denotes the reflector boundary of the void gap. J_i⁺ is the neutron flow in the forward directions at the core boundary of the void gap, and is given in the discrete S_n approximation by

$$J_i^+ \equiv \sum_{m; \mu_m > 0} w_m \mu_m N_{m,i} \quad (23)$$

where ω_m is the solid angle through which the neutrons travel in the m^{th} direction, μ_m is the m^{th} direction cosine, and $N_{m,i}$ is the angular neutron flux in direction m at the core boundary of the void gap. Considering only the escape probability in direction m we write

$$\bar{P}_{c,m,i} = \frac{\mu_m N_{m,i} A_i - \mu_m N_{m,i+1} A_{i+1}}{\mu_m N_{m,i} A_i} \quad (24)$$

This requires an expression for $N_{m,i+1}$ in terms of $N_{m,i}$ and σ_i , consequently we turn to the one-dimensional, steady-state, Boltzmann neutron transport conservation equation. The "diamond"-central difference representation of this equation, as formulated by Lee⁽⁵⁾, may be written

$$\frac{\mu_m}{V_i} \cdot (A_{i+1} N_{m,i+1} - A_i N_{m,i}) + \frac{1}{2V_i} (A_{i+1} - A_i) \cdot [(\delta + \beta)_{m+1} N_{m+1} - (\delta - \beta)_{m+1} N_m] + \sigma_i N'_{m,i} = S_{m,i} \quad (25a)$$

where "i" is a spatial index, "m" is an angular index, and the surface and volume elements in cylindrical geometry are given by

$$A_i = 2\pi r_i \quad ; \quad V_i = \pi (r_{i+1} + r_i) \cdot (r_{i+1} - r_i) \quad (25b)$$

σ_i is the macroscopic cross section, $S_{m,i}$ is the source, δ and β are angular coefficients determined by the integration and differentiation properties implicit in the finite-difference form of the transport conservation equations, and $N'_{m,i}$ is given by the "diamond" difference scheme as

$$2 N'_{m,i} = N_{m,i} + N_{m,i+1} \quad (25c)$$

The DB^2 (i.e., σ_i) is defined such that the conservation condition for radial transport is correct when transverse leakage from the void is included; the angular transport being calculated by the code independent of σ_i . In a void gap the source $S_{m,i}$ is identically zero,

then equation (25a) may be written

$$\frac{\mu_m}{V_i} (A_{i+1} N_{m,i+1} - A_i N_{m,i}) + \sigma_i N'_{m,i} = 0 \quad (26)$$

To get the desired expression for $N_{m,i+1}$, we insert equation (25c) into the conservation equation (26), and obtain

$$N_{m,i+1} = N_{m,i} \cdot \left(\frac{2\mu_m A_i - \sigma_i V_i}{2\mu_m A_{i+1} + \sigma_i V_i} \right) \quad (27)$$

Putting equation (27) back into equation (24), multiplying the angular fluxes by their respective solid angles, and summing over all forward angles yields the total escape probability

$$\bar{P}_{ci} = \frac{\sum_{m; \mu_m > 0} \omega_m \mu_m N_{m,i} \left[A_i - A_{i+1} \left(\frac{2\mu_m A_i - \sigma_i V_i}{2\mu_m A_{i+1} + \sigma_i V_i} \right) \right]}{\sum_{m; \mu_m > 0} \omega_m \mu_m N_{m,i} A_i} \quad (28)$$

or, by definition (23), equation (28) may be shortened to

$$\bar{P}_{ci} = 1 - \frac{A_{i+1}}{J_i^+ A_i} \sum_{m; \mu_m > 0} J_{m,i}^+ \left(\frac{2\mu_m A_i - \sigma_i V_i}{2\mu_m A_{i+1} + \sigma_i V_i} \right) \quad (29)$$

where

$$J_{m,i}^+ = \omega_m \mu_m N_{m,i} \quad ; \quad \mu_m > 0$$

Equation (29) is the DB² equation; it is solved for σ_i by the second-order Newton iteration method (or "Newton-Raphson algorithm")⁽⁶⁾ which consists in drawing successive tangents to the $F(\sigma_i)$ curve. $F(\sigma_i)$ is defined by simply rearranging equation (29),

$$F(\sigma_i) \equiv Q_i - R_i \cdot \sum_{m; \mu_m > 0} J_{m,i}^+ \left(\frac{2\mu_m A_i - \sigma_i V_i}{2\mu_m A_{i+1} + \sigma_i V_i} \right) \quad (30)$$

ATOMICS INTERNATIONAL
 A Division of North American Aviation, Inc.

where $Q_i \equiv (1 - \bar{P}_{ci})$

$$R_i \equiv \frac{A_{i+1}}{J_i^+ A_i}$$

Then its first derivative is

$$F'(\sigma_i) = \frac{\partial F(\sigma_i)}{\partial \sigma_i} = 2 R_i V_i (A_{i+1} + A_i) \sum_{m; \mu_m > 0} \frac{J_{mi}^+ \mu_m}{(2\mu_m A_{i+1} + \sigma_i V_i)^2} \quad (31)$$

Consider the i^{th} gap whose transverse leakage probability is characterized by σ_i , so we may henceforth drop the subscript "i". The iteration scheme begins with an initial guess for σ (eg. $\sigma_0 = 0$). The first calculated value of σ is then

$$\sigma_1 = \sigma_0 - \frac{F(\sigma_0)}{F'(\sigma_0)} \quad (32a)$$

or, more generally, after j iterations,

$$\sigma_{j+1} = \sigma_j - \frac{F(\sigma_j)}{F'(\sigma_j)} \quad (32b)$$

Convergence is satisfied when

$$\frac{\sigma_{j+1} - \sigma_j}{\sigma_{j+1}} = \frac{-F(\sigma_j)}{\sigma_{j+1} F'(\sigma_j)} \leq \epsilon \quad (33)$$

where ϵ is some specified convergence criterion. The value of σ obtained in this calculation is precisely the macroscopic absorption cross section for the void region. (In the DTK code this quantity must be entered as both an absorption and a transport cross section.) An estimate of the absorption cross section (DB^2) should be made for the high order S_n calculation of angular distribution. A good first guess based on the isotropic escape probability, $\bar{P}_{c iso}$, (calculated with FOSDICK by inputting $N(\varphi) = 1$) can be derived from the DB^2 equation (29).

ATOMICS INTERNATIONAL

A Division of North American Aviation, Inc.

NO. NAA SR TDR 10421

DATE 9-9-64

PAGE 23 OF 87

Assuming isotropic flux at the void boundary, $J_i^+ = \frac{\phi_i}{4}$, and $2\mu_m A_{i+1} > \sigma_i V_i$, one obtains from equation (29)

$$DB^2 = C \cdot \frac{R}{X^2 - R^2} \cdot \bar{P}_{c iso} \quad (34)$$

where R and X are radii (in cm.) of the core and reflector, respectively; and C is a constant approximately correcting for assumptions applied to equation (29). If the angular distribution is slightly less forward-peaked than a sine distribution, C varies from .6 to .7 for gaps of 10 to 100 mils, and .9 to 1.3 for gaps of 1 to 4 inches surrounding a typical SNAP core.

4. COMPARISON OF REACTIVITY CALCULATIONS WITH MEASUREMENTS

Reactivity calculations have been made utilizing the DTK S_n transport code ⁽⁴⁾, with 16-group zirconium-hydride spectrum weighted cross section library, for the SNAP 8 Experimental Reactor (S8ER) in a water environment. The calculational model is presented in Addendum I of this report. ⁽⁷⁾ For the S8ER water flooded and reflected core the reactivity effect of void sleeves of 1, 3, and 4 inches were calculated and a comparison with measurements from critical assemblies ⁽⁸⁾ is presented in Table I, and plotted in Figure 6.

Table I

S8ER Void Sleeve Worth Calculations and Measurements

Void Gap Width (inches)	S_4 Calculations $\Delta K_{eff} / .008$ (\$)	Measured Reactivity Worth (\$)
1	-4.02	-3.90 [±] .36
3	-8.44	-8.14 [±] .54
4	-9.75	-9.44 [±] .61

A modification of the reported measurements was made to account for the reactivity effect of the water on the ends of the void gap, since the calculations assumed open ends. This modification was estimated on the basis of results from another critical experiment. ⁽²⁾

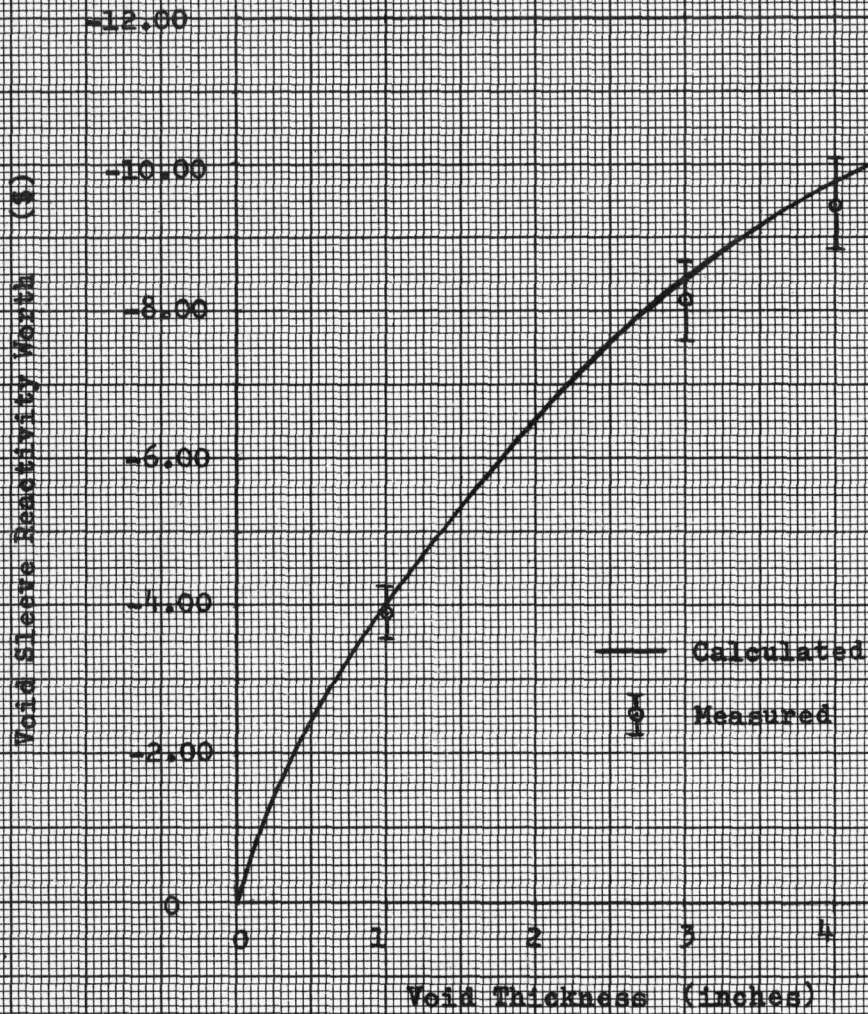


FIGURE 6. 38ER Void Sleeve Worth Calculations and Measurements.

5. CONCLUSIONS AND FURTHER WORK

a. A method has been developed to calculate escape probabilities for neutrons entering an annular void, accounting for anisotropic angular distributions and axial spatial distributions of neutrons on the surface of a cylindrical core. The escape probability is used to calculate a DB^2 which simulates axial leakage from the void in reactivity calculations with the DTK S_n transport code. The FOSDICK code has been written to perform the escape probability and DB^2 calculations. The present method of calculating DB^2 assumes the same escape probability for reflector neutrons as for core neutrons. An attempt to redefine the escape probability for the void in terms of both P_c and P_R should be made.

b. In the calculation of angular distribution only the ϕ angular dependence was considered. The inclusion of the θ angular dependence would not cause serious difficulty, and it may be a necessary factor in considering small gaps.

c. The FOSDICK code in its present form could be quite easily modified to read in angular fluxes and S_n constants directly from the output tape of a DTK calculation.

d. A comparison of reactivity calculations, using the FOSDICK code for DB^2 and the DTK code for reactivity, with measurements of void sleeve reactivity worths in SCA-4B water immersion experiments (S8ER core) shows very close agreement.

REFERENCES

1. P. Benoist, A.E.R.E. TRANS-842, translated by I.S. Grant (1958).
2. D. C. Leslie, Reactor Science and Technology, 16, 1-11 (1962).
3. R. E. Bellman, R. E. Kalaba, and M. C. Prestrud, Invariant Imbedding and Radiative Transfer in Slabs of Finite Thickness, pp. 11-13. American Elsevier Publishing Co., Inc., New York, 1963.
4. W. W. Engle, "The DTK/DTF II S_n Transport Codes", private communication.
5. C. E. Lee, "The Discrete S_n Approximation to Transport Theory", LA-2595, Los Alamos Scientific Laboratory (March 9, 1962).
6. L. Lapidus, Digital Computation for Chemical Engineers, pp. 288-292. McGraw-Hill Book Co., Inc., New York, 1962.
7. R. H. Norman, D. G. Oliver, and W. R. Castle, "A Method of Calculating Annular Void Reactivity Effects in Cylindrical Reactors - S8ER Computational Model", NAA-SR-TDR-10421, Addendum I, C.R.D. (September 9, 1964).
8. S. Yee, "Reactivity Worth of Poison and Void Sleeves", Reactor Test Procedure No. SCA-4B-12, January 11, 1962; Part VI. Results and Conclusions, June 13, 1963.
9. L. I. Moss, "Reactivity Worths of Poison Sleeves - Water in Core", Reactor Test Procedure No. SCA-4B-7, April 3, 1962; Part VI. Results and Conclusions, November 8, 1962.
10. D. J. McGoff and E. M. Faelten, private communication (February, 1964).

APPENDIX - The FOSDICK Code

This Appendix is a presentation of the FORTRAN II program, designated FOSDICK, written to compute escape probabilities and DB^2 to be used in the DTK/DTF II S_n Transport codes for reactivity worth calculations of cylindrical annular voids.

a. Description

The present form of the FOSDICK code consists of a main program of input-output and the following four subroutines:

(1) ANGLE - takes data from DTK/DTF II output (which was read into the main program from any S_n transport calculation up through the S_{16} approximation) and sets up an angular distribution of neutron flux as a function of sine of the angle (ϕ) measured from the z -axis. This subroutine is no more than a manipulation of the input angular fluxes (which have been summed over all energy groups at the core/void boundary), the S_n direction cosines, and the S_n point weights. The angular fluxes are (solid angle) weighted over each discrete (average) ϕ angle. The sine of ϕ is taken directly from the direction cosine. The end points of the distribution ($\phi = 0, \pi$) are obtained from linear extrapolation in $\sin \phi$. The peak of the distribution ($\phi = \pi/2$) is computed to be one-half the value between a linear extrapolation in $\sin \phi$ and a flat peak determined by a straight line connecting the angular fluxes at the nearest points on each side of $\pi/2$. This subroutine also prints the (unnormalized) angular distribution in the output.

(2) NORM - normalizes the angular distribution (obtained from ANGLE) for direct use in the escape probability numerical integration scheme, and prints the normalization factor in the output.

(3) GPESC - computes the geometrical global probability of escape for each axial position specified, and averages it over an axial fission density distribution which was read into the main program. The

escape probability is computed as four separate escape probability integrals, an anisotropic weighting factor being obtained by linear interpolation (in $\sin \phi$) between the discrete values in the normalized angular distribution. The average escape probability is computed for each gap thickness specified, and each of the four integrals as well as their sum are printed out by axial position and gap thickness.

(4) DBSQ - computes DB^2 for the void gap using the average escape probability from GPESC and the angular fluxes, S_n direction cosines, and S_n point weights from DTK/DTF II output which were read into the main program. The angular fluxes used here should be of the same order as the reactivity calculation, usually S_4 . The DB^2 is calculated by S_n theory, the transcendental equation being solved by the second-order Newton iteration method. DB^2 is computed and printed out for each gap thickness specified.

b. Code Equations

The bulk of the code equations are found in the subroutines GPESC and DBSQ, the derivation of the major equations being presented in detail in the body of this report. The following is a presentation of all the equations used in the code leading to the computation of escape probabilities and DB^2 .

The gap thicknesses considered in a given calculation are determined by

$$\left. \begin{aligned} x_i &= \left(\frac{x_{max} - x_{min}}{N} \right) \cdot i + x_{min} \\ t_i &= x_i - R \end{aligned} \right\} \quad (34)$$

where i = index denoting radial position,

x_i = radius of reflector medium for the i^{th} gap,

x_{max} = radius of reflector medium for largest gap considered,

x_{min} = radius of reflector medium for smallest gap considered,

or usually, the core radius,

N = number of gaps considered,

t_i = gap thickness of i^{th} gap,

R = core radius.

The axial positions along half the length of the core are determined by

$$z_j = \left(\frac{L}{2}\right) \cdot \frac{j}{M} \quad (35)$$

where L = core length,

z_j = j^{th} axial position,

M = number of axial positions considered,

j = index denoting axial position.

The four escape probabilities for a neutron on the surface of the core are given in equation (8). The limits of integration are given by equations (9b). The integrals in equation (8) are evaluated by the 5th order Legendre-Gauss quadrature method as:

$$P_{1ij} = \frac{q_{1ij}}{4} \sum_{k=1}^5 \sin \varphi_k \cdot N_1(\varphi_k) \cdot w_k \quad \left. \vphantom{P_{1ij}} \right\} \quad (36a)$$

where $\varphi_k = \frac{q_{1ij}}{2} \cdot (1 + u_k)$

u_k, w_k are the Legendre zeros and Christoffel weights, respectively.

$$P_{2ij} = \frac{(q_{2ij} - q_{1ij})}{2\pi} \sum_{k=1}^5 \sin \varphi_k \cdot N_2(\varphi_k) \cdot \left[\frac{\pi}{2} - f_{1ij}(\varphi_k)\right] \cdot w_k \quad \left. \vphantom{P_{2ij}} \right\} \quad (36b)$$

where $\varphi_k = \left(\frac{q_{2ij} + q_{1ij}}{2}\right) + \left(\frac{q_{2ij} - q_{1ij}}{2}\right) \cdot u_k$

$$P_{3ij} = \frac{(q_{4ij} - q_{3ij})}{2\pi} \sum_{k=1}^5 \sin \varphi_k \cdot N_3(\varphi_k) \cdot \left[\frac{\pi}{2} - f_{2ij}(\varphi_k)\right] \cdot w_k \quad \left. \vphantom{P_{3ij}} \right\} \quad (36c)$$

where $\varphi_k = \left(\frac{q_{4ij} + q_{3ij}}{2}\right) + \left(\frac{q_{4ij} - q_{3ij}}{2}\right) \cdot u_k$

$$P_{4ij} = \frac{(\pi - q_{4ij})}{4} \sum_{k=1}^5 \sin \varphi_k \cdot N_4(\varphi_k) \cdot w_k \quad \left. \vphantom{P_{4ij}} \right\} \quad (36d)$$

where $\varphi_k = \left(\frac{\pi + q_{4ij}}{2}\right) + \left(\frac{\pi - q_{4ij}}{2}\right) \cdot u_k$

and

$$f_{1ij}(\varphi_k) = \cos^{-1} \left[\left(\frac{x_i^2 - R^2}{2Rz_j} \right) \cot \varphi_k - \frac{z_j}{2R} \tan \varphi_k \right] \quad (36e)$$

$$f_{2ij}(\varphi_k) = \cos^{-1} \left[\left(\frac{x_i^2 - R^2}{2R(z_j - L)} \right) \cot \varphi_k - \frac{(z_j - L)}{2R} \tan \varphi_k \right]$$

The anisotropic weighting factors $N_1(\varphi_k)$, etc., are determined from the discrete angular distribution, Φ_ℓ (as a function of $\sin \varphi_\ell$), by a linear interpolation scheme. For example,

$$N_1(\varphi_k) = \Phi_\ell - \left(\frac{\sin \varphi_\ell - \sin \varphi_k}{\sin \varphi_\ell - \sin \varphi_{\ell-1}} \right) \cdot (\Phi_\ell - \Phi_{\ell-1}) \quad (37)$$

where Φ_ℓ is the normalized angular flux at angle φ_ℓ in the computed angular distribution; $\Phi_{\ell-1}$ is at the previous (smaller) angle, $\varphi_{\ell-1}$; $N_1(\varphi_k)$ is the normalized angular flux at angle φ_k , as determined by the quadrature integration. The angular distribution normalization factor (in NORM) is computed as

$$\text{"Normalization factor"} = \frac{\pi}{4} \sum_{k=1}^5 \sin \varphi_k \cdot N_m(\varphi_k) \cdot w_k \quad (38)$$

where $\varphi_k = \frac{\pi}{2}(1 + u_k)$, and $N_m(\varphi_k)$ is obtained by equation (37). The average escape probability is then computed by

$$\bar{P}_{ci} = \frac{\sum_{j=1}^m P_{cij} \cdot \phi_j}{\sum_{j=1}^m \phi_j} \quad (39)$$

where $P_{cij} = P_{1ij} + P_{2ij} + P_{3ij} + P_{4ij}$
 ϕ_j = axial fission neutron density as a function of z_j
(alternately, an "importance" distribution).

The DB^2 equation and its solution by the Newton iteration method give rise to three summations, appearing ⁱⁿ equations (25), (30), and (31).

Considering a given gap thickness t_i , (therefore dropping the subscript "i") the sums are computed for the j^{th} iteration as:

$$S_1 = \sum_{m=1}^M w_m \mu_m N_m \quad (40)$$

$$S_{2j} = \sum_{m=1}^M w_m \mu_m N_m \left(\frac{2\mu_m A_1 - \sigma_j V_1}{2\mu_m A_2 + \sigma_j V_1} \right) \quad (41)$$

$$S_{3j} = \sum_{m=1}^M \frac{w_m N_m \mu_m^2}{(2\mu_m A_2 + \sigma_j V_1)^2} \quad (42)$$

where N_m = angular neutron flux

w_m = S_n solid-angle weight for flux in the m^{th} direction

μ_m = S_n m^{th} direction cosine

$M = \frac{n(n+2)}{8}$, (n = order of S_n)

A_1 = inner surface area of annular void (cm^2)

A_2 = outer surface area of annular void (cm^2)

V_1 = volume of annular void (cm^3)

σ_j = DB^2 for the j^{th} iteration (cm^{-1})

Then the iteration scheme is comprised of the equations (30) through (33) written for the code as:

$$\left. \begin{aligned} Q &= 1 - \bar{P}_c \\ R &= A_2 / (A_1 \cdot S_1) \\ F_j &= Q - R \cdot S_{2j} \\ F'_j &= 2R \cdot V_1 \cdot (A_1 + A_2) \cdot S_{3j} \end{aligned} \right\} \quad (43)$$

After j iterations DB^2 is computed as

$$\sigma_{j+1} = \sigma_j - \frac{F_j}{F'_j} \quad (44)$$

and the convergence is

$$\epsilon = - \frac{F_j}{\sigma_{j+1} F'_j} \quad (45)$$

ATOMICS INTERNATIONAL
A Division of North American Aviation, Inc.

NO. NAA SR TDR 10421
DATE 9-9-64
PAGE 33 OF 87

c. Program Listing

The FOSDICK code is presented in the following listing. For further information and program deck see R. H. Norman or W. R. Castle.

*7911-99

(001)

001 020 25 R H NORMAN-

1054001 A72 727-753

00-00-00

LISTING (COMPILATION)

BEGIN COMPILATION 29

8 TAPES 00-00-00

08/17/24

```

C FØSDICK - CODE FOR CALCULATION OF TRANSVERSE LEAKAGE PROBABILITIES00000010
C FROM ANNULAR VOIDS IN CYLINDRICAL REACTORS 00000020
C MAIN PROGRAM . . . . 00000030
  DIMENSION X(20),Z(100),P1(20,100),P2(20,100),P3(20,100),PRØB(20,1000000100
10),U(5),W(5),T(20),P4(20,100),PHI(100),AVG(20),PMS(20),AFX(20),AFX00000200
25(20),PHX(40),PM7(40),WØ(40),PHY(40),SIG(20),CØN(20),AFLUX(3),S4M700000201
3(3),S4WØ(3) 00000202
51 READ INPUT TAPE 5,1,M,(PHI(J),J=1,M) 00000400
  1 FØRMAT(I12/(6E12.8)) 00000600
52 READ INPUT TAPE 5,57,NSN,MM,MAX,EPS,(PHX(MI),MI=1,MM),(PM7(MI),MI=00000610
  11,MM),(WØ(MI),MI=1,MM) 00000611
57 FØRMAT(3I12,E12.8/(6E12.8)) 00000620
50 READ INPUT TAPE 5,55,R,XMAX,XMIN,CØRL,N 00000700
55 FØRMAT(4F12.8,I12) 00000710
100 READ INPUT TAPE 5,101,(AFLUX(I),I=1,3),(S4M7(I),I=1,3),(S4WØ(I),I=00000750
  11,3) 00000751
101 FØRMAT(6E12.8) 00000760
  WRITE OUTPUT TAPE 6,6,R,XMAX,XMIN,CØRL,N,M 00000800
  6 FØRMAT(1HØ 2X, 16HTHE FØSDICK CODE///2X,22H* * * INPUT DATA 00000900
  1* * //2X,19HCYLINDER RADIUS, R=F8 5//2X,31HRADIUS OF MEDIUM (MAX 00000901
  2), X-MAX=F8.5//2X,31HRADIUS OF MEDIUM (MIN ), X-MIN=F8 5//2X,19HCY00000902
  3LINDER LENGTH, L=F8.5//2X,25HNUMBER OF X INTERVALS, N=12//2X,26HNU00000903
  4MBER OF AXIAL POINTS, M=13) 00000904
  WRITE OUTPUT TAPE 6,59,NSN 00001600
59 FØRMAT(//2X,14HØRDER OF SN = I2) 00001610
  WRITE OUTPUT TAPE 6,58,EPS,MAX,(PHX(MI),PM7(MI),WØ(MI),MI=1,MM) 00001700
58 FØRMAT(///2X,24HCØNVERGENCE CRITERIØN = F9 6,4X,22HMAX NØ ITERATØ0001800
  1IØNS = 13//4X,12HANGULAR FLUX5X,17HDIRECTION CØSINES5X,13HPØINT WE00001801
  2IGHTS/(4X,E12.6,8X,F9.6,F20 6)) 00001802
  WRITE OUTPUT TAPE 6,102,(AFLUX(I),S4M7(I),S4WØ(I),I=1,3) 00001850
102 FØRMAT(//2X,31HS-4 FLUXES FOR DBSQ CALCULATION/4X,12HANGULAR FLUX500001860
  1X,17HDIRECTION CØSINES5X,13HPØINT WEIGHTS/(4X,E12 6,8X,F9.6,F20.6)00001861
  2) 00001862
  U(1)=-.9061798 00001900
  U(2)=-.5384693 00001901
  U(3)=0.0 00001902
  U(4)=.5384693 00001903
  U(5)=.9061798 00001904
  W(1)=.236927 00001905

```

W(2)=.478629	00001906
W(3)=.568889	00001907
W(4)=.478629	00001908
W(5)=.236927	00001909
CALL ANGLE (PM7,PHX,W0,PMS,AFX,NSN,L)	00002000
CALL NORM (U,W,PMS,AFX,L,AFX5,AFX4)	00003010
CALL GPESC (R,XMAX,XMIN,C0RL,N,M,U,W,PHI,L,PMS,AFX5,X,T,AVG,Z,P1,P00004220	00004221
12,P3,PR0B,P4,AFX4)	00004221
CALL DBSQ (R,MM,N,X,AVG,EPS,MAX,SIGMA,C0NV,SIG,C0N,AFLUX,S4M7,S4W000009420	00009430
1)	00012650
DO 308 I=1,N	00012700
WRITE OUTPUT TAPE 6,7,T(I),AVG(I),C0N(I),SIG(I)	00012800
7 FORMAT(//2X,6HF0R T=F7.4,6X,44HAVERAGE ESCAPE PROBABILITY, PC-ANIS00012801	00012801
10TR0PIC =E12.5//22X,14HC0NVERGENCE = F9.7,6X,15HSIGMA (DBSQ) = F9.00012802	00012802
27//)	00013000
WRITE OUTPUT TAPE 6,9	00013100
9 FORMAT(1H ,8X,1HZ,13X,2HP1,15X,2HP2,15X,2HP3,15X,2HP4,15X,2HPC)	00013300
DO 308 J=1,M	00013400
WRITE OUTPUT TAPE 6,10,Z(J),P1(I,J),P2(I,J),P3(I,J),P4(I,J),PR0B(I	00013401
1,J)	00013600
10 FORMAT(F13.5,5E17.5)	00013700
308 CONTINUE	00013800
IF DIVIDE CHECK 44,45	00013900
44 CALL DUMP	00014000
45 GO TO 50	
END(1,0,0,0,0,0,1,0,0,1,0,0,0,0,0)	

STORAGE NOT USED BY PROGRAM

DEC OCT
11014 25406

DEC OCT
32561 77461

STORAGE LOCATIONS FOR VARIABLES APPEARING IN DIMENSION AND EQUIVALENCE STATEMENTS

DEC OCT
AFLUX 10483 24363
CON 10503 24407
P4 2474 04652
PM7 10643 24623
S4W0 10477 24355
W0 10603 24553

DEC OCT
AFX5 10703 24717
P1 10474 24352
PHI 10863 25157
PMS 10743 24767
SIG 10523 24433
W 10888 25210

DEC OCT
AFX 10723 24743
P2 8474 20432
PHX 10683 24673
PR0B 4474 10572
T 10883 25203
X 11013 25405

DEC OCT
AVG 10763 25013
P3 6474 14512
PHY 10563 24503
S4M7 10480 24360
U 10893 25215
Z 10993 25361

STORAGE LOCATIONS FOR VARIABLES NOT APPEARING IN COMMON, DIMENSION, OR EQUIVALENCE STATEMENTS

DEC OCT
AFX4 474 00732
L 470 00726
N 466 00722
XMAX 462 00716

DEC OCT
CONV 473 00731
MAX 469 00725
NSN 465 00721
XMIN 461 00715

DEC OCT
CORL 472 00730
MM 468 00724
R 464 00720

DEC OCT
EPS 471 00727
M 467 00723
SIGMA 463 00717

SYMBOLS AND LOCATIONS FOR SOURCE PROGRAM FORMAT STATEMENTS

EFN LOC
8)1 1 00712
8)A 10 00471
8)1R 59 00617

EFN LOC
8)6 6 00677
8)1N 55 00703
8)35 101 00701

EFN LOC
8)7 7 00531
8)1P 57 00707
8)36 102 00557

EFN LOC
8)9 9 00503
8)1Q 58 00612

LOCATIONS FOR OTHER SYMBOLS NOT APPEARING IN SOURCE PROGRAM

DEC OCT
2) 295 00447
C)G2 459 00713

DEC OCT
3) 299 00453
C)100 460 00714

DEC OCT
4) 32767 77777
D)20K 241 00361

DEC OCT
6) 305 00461
D)30K 240 00360

LOCATIONS OF NAMES IN TRANSFER VECTOR

	DEC	OCT		DEC	OCT		DEC	OCT		DEC	OCT
ANGLE	5	00005	DBSQ	8	00010	DUMP	9	00011	GPESC	7	00007
NORM	6	00006	(FIL)	4	00004	(FPT)	0	00000	(RTN)	2	00002
(STH)	3	00003	(TSH)	1	00001						

ENTRY POINTS TO SUBROUTINES NOT OUTPUT FROM LIBRARY

ANGLE	DBSQ	DUMP	GPESC	NORM	(FIL)	(FPT)	(RTN)
(STH)	(TSH)						

EXTERNAL FORMULA NUMBERS WITH CORRESPONDING INTERNAL FORMULA NUMBERS AND OCTAL LOCATIONS

EFN	IFN	L0C	EFN	IFN	L0C	EFN	IFN	L0C	EFN	IFN	L0C
51	14	00017	52	20	00034	50	32	00075	100	34	00113
308	81	00431	44	83	00443	45	84	00446			

BEGIN COMPILATION 29

8 TAPES 01-42-42

08/15/64

```

SUBROUTINE ANGLE (PM7,PHX,W0,PMS,AFX,NSN,L)          00002001
DIMENSION PM7(40),PHX(40),W0(40),PMS(20),AFX(20),SUM1(10),SUM2(10) 00002002
1,AFX0(20),PMU(20)                                00002003
200 MM=(NSN*(NSN+2))/8                             00002010
201 L=NSN+3                                         00002012
202 LM=(L+1)/2                                      00002014
203 LLM=LM-1                                        00002016
204 PMU(1)=1.0000                                  00002018
205 PMU(LM)=0.0                                     00002020
206 PMU(L)=-1.0000                                 00002022
DO 208 LA=2,LLM                                    00002024
207 MI=MM+2-LA                                     00002026
208 PMU(LA)=PM7(MI)                                00002028
209 LE=(L-3)/2                                     00002030
DO 210 LA=1,LE                                     00002032
100 LAL=LA+(L+3)/2-1                               00002033
101 MAM=MM-LE+LA                                   00002034
210 PMU(LAL)=-PM7(MAM)                             00002035
DO 212 N=1,10                                      00002036
SUM1(N)=0.0                                        00002037
212 SUM2(N)=0.0                                    00002038
213 NS=NSN/2                                       00002040
214 J=1                                             00002042
215 K=0                                             00002044
DO 222 N=1,NS                                     00002046
216 J=J+N-1                                        00002048
217 K=K+N                                          00002050
DO 219 MI=J,K                                     00002052
218 SUM1(N)=SUM1(N)+PHX(MI)*W0(MI)                00002054
219 SUM2(N)=SUM2(N)+W0(MI)                        00002056
220 AFX0(N)=SUM1(N)/SUM2(N)                        00002058
221 AFX(N+1)=AFX0(N)                              00002060
102 M=L-N                                          00002061
222 AFX(M)=AFX(N+1)                               00002062
223 XPX=(QATANF(SQRTF(1.-(PM7(MM))**2),PM7(MM)))/(QATANF(SQRTF(1.-(PM700002064
1(MM-1))**2),PM7(MM-1))-QATANF(SQRTF(1.-(PM7(MM))**2),PM7(MM))) 00002066
224 AFX(1)=AFX(2)-XPX*(AFX(3)-AFX(2))            00002068
225 AFX(L)=AFX(1)                                 00002070
103 MA1=MM-NS+1                                   00002071

```

```
104 MA2=MM-NS+2 00002072
226 ACM1=CATANF(SQRTF(1.-(PM7(MA1))**2),PM7(MA1)) 00002073
227 ACM2=CATANF(SQRTF(1.-(PM7(MA2))**2),PM7(MA2)) 00002074
228 AFXL0=AFX(LLM) 00002076
229 AFXHI=AFX(LLM)+(AFX(LLM)-AFX(LLM-1))*(1.5707963-ACM1)/(ACM1-ACM2) 00002078
230 GAMMA=0.5 00002080
231 AFX(LM)=AFXL0+GAMMA*(AFXHI-AFXL0) 00002082
D0 233 LA=1,L 00002083
233 PMS(LA)=SINF(CATANF(SQRTF(1.-(PMU(LA))**2),PMU(LA))) 00002084
WRITE OUTPUT TAPE 6,232,(PMS(LA),AFX(LA),LA=1,L) 00002086
232 FORMAT(///2X,18H* * * OUTPUT * * *//2X,21HANGULAR DISTRIBUTION,5X,00002087
14HSINE9X,4HFLUX//(25X,F8.5,4X,E13.6)) 00002088
RETURN 00002090
END(1,0,0,0,0,0,1,0,0,1,0,0,0,0,0)
```

STORAGE NOT USED BY PROGRAM

DEC	OCT	DEC	OCT
509	00775	32561	77461

STORAGE LOCATIONS FOR VARIABLES APPEARING IN DIMENSION AND EQUIVALENCE STATEMENTS

DEC	OCT	DEC	OCT	DEC	OCT	DEC	OCT
AFXO	488 00750	PMU	468 00724	SUM1	508 00774	SUM2	498 00762

STORAGE LOCATIONS FOR VARIABLES NOT APPEARING IN COMMON, DIMENSION, OR EQUIVALENCE STATEMENTS

DEC	OCT	DEC	OCT	DEC	OCT	DEC	OCT
ACM1	448 00700	ACM2	447 00677	AFXHI	446 00676	AFXLO	445 00675
GAMMA	444 00674	J	443 00673	K	442 00672	LAL	441 00671
LA	440 00670	LE	439 00667	LLM	438 00666	LM	437 00665
MA1	436 00664	MA2	435 00663	MAM	434 00662	MI	433 00661
MM	432 00660	M	431 00657	N	430 00656	NS	429 00655
XPX	428 00654						

SYMBOLS AND LOCATIONS FOR SOURCE PROGRAM FORMAT STATEMENTS

EFN	LCC	EFN	LCC	EFN	LCC	EFN	LCC
8178	232 00633						

LOCATIONS FOR OTHER SYMBOLS NOT APPEARING IN SOURCE PROGRAM

DEC	OCT	DEC	OCT	DEC	OCT	DEC	OCT
1)	415 00637	2)	379 00573	3)	385 00601	6)	389 00605
9)	412 00634	C)G1	418 00642	C)G2	419 00643	C)G4	420 00644
C)G5	421 00645	C)G6	422 00646	C)G7	423 00647	C)G8	424 00650
C)G9	425 00651	C)GA	426 00652	C)GC	427 00653	D)107	196 00304

LOCATIONS OF NAMES IN TRANSFER VECTOR

DEC	OCT	DEC	OCT	DEC	OCT	DEC	OCT
GATAN	1 00001	SIN	2 00002	SQRT	0 00000	(FIL)	4 00004
(STH)	3 00003						

ENTRY POINTS TO SUBROUTINES NOT OUTPUT FROM LIBRARY

CATAN SIN SQRT (FIL) (STH)

EXTERNAL FORMULA NUMBERS WITH CORRESPONDING INTERNAL FORMULA NUMBERS AND OCTAL LOCATIONS

EFN	IFN	L0C	EFN	IFN	L0C	EFN	IFN	L0C	EFN	IFN	L0C
200	5	00123	201	6	00140	202	7	00145	203	8	00156
204	9	00163	205	10	00165	206	11	00167	207	13	00176
208	14	00204	209	15	00211	100	17	00225	101	18	00242
210	19	00250	212	22	00260	213	23	00264	214	24	00272
215	25	00274	216	27	00305	217	28	00311	218	30	00317
219	31	00323	220	32	00330	221	33	00333	102	34	00335
222	35	00341	223	36	00347	224	37	00403	225	38	00413
103	39	00416	104	40	00424	226	41	00431	227	42	00451
228	43	00472	229	44	00474	230	45	00510	231	46	00512
233	48	00525									

BEGIN COMPILATION 29

8 TAPES 01-43-02

08/15/64


```
SUBROUTINE NORM (U,W,PMS,AFX,L,AFX5,AFX4)          00003011
DIMENSION U(5),W(5),PMS(20),AFX(20),AFX5(20)      00003012
AFX4=0.0                                           00003098
56 DO 40 K=1,5                                     00003100
46 A11=(3.1415926/2.)*(1.+U(K))                   00003200
DO 49 LA=1,L                                       00003300
48 IF(PMS(LA)-SINF(A11))49,54,54                 00003400
49 CONTINUE                                        00003500
54 A12 = (PMS(LA)-SINF(A11))/(PMS(LA)-PMS(LA-1)) 00003600
55 AFX6 = AFX(LA)-A12*(AFX(LA)-AFX(LA-1))        00003700
40 AFX4 = (3.14159265/4.)*SINF(A11)*AFX6*W(K) + AFX4 00004000
DO 42 LA=1,L                                       00004100
42 AFX5(LA) = AFX(LA)/AFX4                        00004200
WRITE OUTPUT TAPE 6,43,AFX4                       00004205
43 FORMAT(//2X,44+ANGULAR DISTRIBUTION NORMALIZATION FACTOR = E12.6) 00004206
RETURN                                             00004210
END(1,0,0,0,0,0,C,1,C,0,1,0,0,C,C,0)
```

STORAGE NOT USED BY PROGRAM

DEC OCT
173 00255

DEC OCT
32561 77461

STORAGE LOCATIONS FOR VARIABLES NOT APPEARING IN COMMON, DIMENSION, OR EQUIVALENCE STATEMENTS

DEC OCT
A11 172 00254

DEC OCT
A12 171 00253

DEC OCT
AFX6 170 00252

DEC OCT
LA 169 00251

SYMBOLS AND LOCATIONS FOR SOURCE PROGRAM FORMAT STATEMENTS

EFN LOC
811B 43 00242

EFN LOC

EFN LOC

EFN LOC

LOCATIONS FOR OTHER SYMBOLS NOT APPEARING IN SOURCE PROGRAM

DEC OCT
1) 165 00245
9) 163 00243
E)3 70 00106

DEC OCT
2) 139 00213
C)G0 167 00247
E)6 80 00120

DEC OCT
3) 141 00215
C)G1 168 00250

DEC OCT
6) 147 00223
C)407 81 00121

LOCATIONS OF NAMES IN TRANSFER VECTOR

DEC OCT
SIN 0 00000

DEC OCT
(FIL) 2 00002

DEC OCT
(STH) 1 00001

DEC OCT

ENTRY POINTS TO SUBROUTINES NOT OUTPUT FROM LIBRARY

SIN (FIL) (STH)

EXTERNAL FORMULA NUMBERS WITH CORRESPONDING INTERNAL FORMULA NUMBERS AND OCTAL LOCATIONS

EFN IFN LOC
56 6 00060
54 11 00122

EFN IFN LOC
46 -7 00061
55 12 00137

EFN IFN LOC
48 9 00073
40 13 00147

EFN IFN LOC
49 10 00115
42 15 00174

BEGIN COMPILATION 29

8 TAPES 01-43-19

08/15/64

```

SUBROUTINE GPESC (R,XMAX,XMIN,CORL,N,M,U,W,PHI,L,PMS,AFX5,X,T,AVG,00004230
1Z,P1,P2,P3,PRGB,P4,AFX4) 00004231
DIMENSION U(5),W(5),PHI(100),PMS(20),AFX5(20),X(20),T(20),AVG(20),00004233
1Z(100),P1(20,100),P2(20,100),P3(20,100),PRGB(20,100),P4(20,100) 00004234
67 AN=FLCATF(N) 00004250
DO 2 I=1,N 00004300
68 AI=FLCATF(I) 00004350
69 X(I)=((XMAX-XMIN)/AN)*AI+XMIN 00004400
2 T(I)=X(I)-R 00004500
DEN=0.0 00004600
70 AM=FLCATF(M) 00004650
DO 3 J=1,M 00004700
71 AJ=FLCATF(J) 00004750
DEN=PHI(J)+DEN 00004800
3 Z(J)=(CORL/2.)*AJ/AM 00004900
DO 60 I=1,20 00004908
DO 60 J=1,100 00004909
P1(I,J)=0.0 00004910
P2(I,J)=0.0 00004911
P3(I,J)=0.0 00004912
60 P4(I,J)=0.0 00004915
UNP=0.0 00004950
DO 53 I=1,N 00005000
DO 51 J=1,M 00005100
11 A1=QATANF((X(I)-R),Z(J)) 00005200
12 A2=QATANF(SQRTF(X(I)**2-R**2),Z(J)) 00005300
13 A3=QATANF(SQRTF(X(I)**2-R**2),(Z(J)-CORL)) 00005400
14 A4=QATANF((X(I)-R),(Z(J)-CORL)) 00005500
A7=3.14159265 00005600
DO 90 K=1,5 00005700
23 A6=(A1/2.)*(1.+U(K)) 00005800
DO 36 LA=1,L 00005810
25 IF(PMS(LA)-SINF(A6))36,26,26 00005820
36 CONTINUE 00005830
26 A8 = (PMS(LA)-SINF(A6))/(PMS(LA)-PMS(LA-1)) 00005840
27 AFX1 = AFX5(LA)-A8*(AFX5(LA)-AFX5(LA-1)) 00005900
15 P1(I,J)=(A1/4.)*SINF(A6)*AFX1*W(K)+P1(I,J) 00006000
17 A5=((A2+A1)/2.)+((A2-A1)/2.)*U(K) 00006600
DO 37 LA=1,L 00006610

```

```

29 IF(PMS(LA)-SINF(A5))37,30,30                                00006620
37 CONTINUE                                                    00006630
30 A9 = (PMS(LA)-SINF(A5))/(PMS(LA)-PMS(LA-1))                00006640
31 AFX2 = AFX5(LA)-A9*(AFX5(LA)-AFX5(LA-1))                    00007100
18 C1=((1.-(X(I)/R)**2)*COSF(A5)/SINF(A5)+((Z(J)/R)**2)*SINF(A5)/COSF
1(A5))/(2.*Z(J)/(-R))                                          00007400
21 P2(I,J)=((A2-A1)/6.2831853)*SINF(A5)*AFX2*(1.5707963-(CATANF(SQRTF
1(1.-C1**2),C1)))*W(K)+P2(I,J)                                00007600
19 B5=((A4+A3)/2.)+(A4-A3)/2.*U(K)                              00007800
DO 38 LA=1,L                                                  00007810
33 IF(PMS(LA)-SINF(B5))38,34,34                                00007820
38 CONTINUE                                                    00007830
34 A10 = (PMS(LA)-SINF(B5))/(PMS(LA)-PMS(LA-1))               00007840
35 AFX3 = AFX5(LA)-A10*(AFX5(LA)-AFX5(LA-1))                  00008300
20 D1=((1.-(X(I)/R)**2)*COSF(B5)/SINF(B5)+(((Z(J)-C0RL)/R)**2)*SINF(B
15)/COSF(B5))/(2.*(C0RL-Z(J))/R)                              00008600
4 P3(I,J) = ((A4-A3)/6.2831853)*SINF(B5)*AFX3*(1.5707963-CATANF(SQRT
1F(1.-D1**2),D1))*W(K) + P3(I,J)                              00008800
88 B4 = ((A7+A4)/2.) + ((A7-A4)/2.)*U(K)                      00008900
DO 39 LA=1,L                                                  00008910
40 IF(PMS(LA)-SINF(B4))39,41,41                                00008920
39 CONTINUE                                                    00008930
41 A13 = (PMS(LA)-SINF(B4))/(PMS(LA)-PMS(LA-1))               00008940
89 AFX6 = AFX5(LA)-A13*(AFX5(LA)-AFX5(LA-1))                  00008950
90 P4(I,J) = ((A7-A4)/4.)*SINF(B4)*AFX6*W(K)+P4(I,J)         00008960
5 PR0B(I,J) = P1(I,J) + P2(I,J) + P3(I,J) + P4(I,J)         00009000
IF(PR0B(I,J))55,51,50                                         00009010
55 PR0B(I,J)=0.0                                               00009020
GO TO 51                                                       00009030
50 IF(PR0B(I,J)-1.0)51,51,56                                    00009040
56 WRITE OUTPUT TAPE 6,80,PR0B(I,J)                            00009050
80 FORMAT(/3X,28+ERROR, ESCAPE PR0BABILITY = F9.6//)         00009060
51 UNM=PHI(J)*PR0B(I,J)+UNM                                    00009200
52 AVG(I)=UNM/DEN                                              00009300
53 UNM=0.0                                                      00009400
RETURN                                                         00009410
END(1,C,0,C,0,C,1,C,0,1,0,C,C,0)

```

STORAGE NOT USED BY PROGRAM

DEC OCT
871 01547

DEC OCT
32561 77461

STORAGE LOCATIONS FOR VARIABLES NOT APPEARING IN COMMON, DIMENSION, OR EQUIVALENCE STATEMENTS

DEC	OCT	DEC	OCT	DEC	OCT	DEC	OCT
A10	870 01546	A13	869 01545	A1	868 01544	A2	867 01543
A3	866 01542	A4	865 01541	A5	864 01540	A6	863 01537
A7	862 01536	A8	861 01535	A9	860 01534	AFX1	859 01533
AFX2	858 01532	AFX3	857 01531	AFX6	856 01530	AI	855 01527
AJ	854 01526	AM	853 01525	AN	852 01524	B4	851 01523
B5	850 01522	C1	849 01521	D1	848 01520	CEN	847 01517
I	846 01516	J	845 01515	LA	844 01514	UNM	843 01513

SYMBOLS AND LOCATIONS FOR SOURCE PROGRAM FORMAT STATEMENTS

EFN	LOC	EFN	LOC	EFN	LOC	EFN	LOC
812G	80 01475						

LOCATIONS FOR OTHER SYMBOLS NOT APPEARING IN SOURCE PROGRAM

DEC	OCT	DEC	OCT	DEC	OCT	DEC	OCT
1)	832 01500	2)	805 01445	3)	809 01451	6)	816 01460
9)	830 01476	C)G0	838 01506	C)G1	839 01507	C)G2	840 01510
C)G3	841 01511	C)100	842 01512	D)20B	331 00513	D)218	779 01413
D)4CA	266 00412	D)40F	364 00554	D)40N	438 00666	D)40P	564 01064
D)40T	576 01100	D)413	711 01307	D)50P	563 01063	D)50T	575 01077
D)60A	265 00411	E)D	353 00541	E)G	363 00553	E)M	437 00665
E)S	574 01076	E)V	700 01274	E)12	710 01306	E)14	752 01360

LOCATIONS OF NAMES IN TRANSFER VECTOR

DEC	OCT	DEC	OCT	DEC	OCT	DEC	OCT
CCS	3 00003	QATAN	0 00000	SIN	2 00002	SCRT	1 00001
(FIL)	5 00005	(STH)	4 00004				

ENTRY POINTS TO SUBROUTINES NOT OUTPUT FROM LIBRARY

CGS CATAN SIN SQRT (FIL) (STH)

EXTERNAL FORMULA NUMBERS WITH CORRESPONDING INTERNAL FORMULA NUMBERS AND OCTAL LOCATIONS

EFN	IFN	L0C	EFN	IFN	L0C	EFN	IFN	L0C	EFN	IFN	L0C
67	5	00263	68	7	00275	69	8	00302	2	9	00310
70	11	00320	71	13	00332	3	15	00342	60	21	00362
11	25	00413	12	26	00424	13	27	00446	14	28	00473
23	31	00514	25	33	00526	36	34	00550	26	35	00555
27	36	00572	15	37	00602	17	38	00622	29	40	00641
37	41	00662	30	42	00667	31	43	00704	18	44	00714
21	45	00764	19	46	01031	33	48	01050	38	49	01073
34	50	01101	35	51	01116	20	52	01126	4	53	01175
88	54	01243	40	56	01261	39	57	01303	41	58	01310
89	59	01325	90	60	01335	5	61	01361	55	63	01371
50	65	01374	56	66	01401	51	68	01414	52	69	01426
53	70	01432									

BEGIN COMPILATION 29

8 TAPES 01-43-45

08/15/64

NAA SR TDR 10421
9-9-64
Page 53 of 87


```

SUBROUTINE DBSQ (R,MM,N,X,AVG,EPS,MAX,SIGMA,C0NV,SIG,C0N,AFLUX,S4M00009421
17,S4W0) 00009422
DIMENSION X(20),AVG(20),SIG(20),C0N(20),AFLUX(3),S4M7(3),S4W0(3) 00009425
100 R1=2.540005*R 00009500
    D0 8 J=1,N 00010000
    SIGMA=0.0 00010054
    NITT=0 00010056
105 R2 = 2.540005*X(J) 00010100
102 T = X(J)-R 00010150
101 A1 = 6.2831853*R1 00010200
103 A2 = 6.2831853*R2 00010210
104 V1 = 3.14159265*(R2+R1)*(R2-R1) 00010220
107 C = 1.-AVG(J) 00010230
106 NITT=NITT+1 00010600
    SUM1 = 0.0 00010700
    SUM2 = 0.0 00010710
    SUM3 = 0.0 00010750
    D0 113 I=1,3 00010800
113 SUM1 = S4W0(I)*S4M7(I)*AFLUX(I)+SUM1 00010900
    D0 127 I=1,3 00011200
127 SUM2 = S4W0(I)*S4M7(I)*AFLUX(I)*((2.*S4M7(I)*A1-SIGMA*V1)/(2.*S4M700011300
1(I)*A2+SIGMA*V1))+SUM2 00011301
    D0 114 I=1,3 00011400
114 SUM3 = S4W0(I)*AFLUX(I)*(S4M7(I)**2)/(2.*S4M7(I)*A2+SIGMA*V1)**2+S00011410
    IUM3 00011411
128 RR = A2/(A1*SUM1) 00011500
115 F = C-RR*SUM2 00011550
116 FPRIME = 2.*RR*V1*(A1+A2)*SUM3 00011600
117 SIGMA = SIGMA-F/FPRIME 00011700
118 C0NV = -F/(FPRIME*SIGMA) 00011800
119 IF(ABSF(C0NV)-EPS)120,120,122 00011900
122 IF(NITT-MAX)106,106,123 00012000
123 WRITE OUTPUT TAPE 6,124,T,C0NV,SIGMA 00012100
124 FORMAT(///2X,3HT =F8.4,6X,53HMAXIMUM NUMBER OF ITERATIONS EXCEEDED00012400
1, CONVERGENCE = F9.7/2X,8HSIGMA = F9.7) 00012401
120 SIG(J) = SIGMA 00012500
121 C0N(J) = C0NV 00012600
    8 CONTINUE 00012610
    RETURN 00012620
END(1,0,0,C,0,C,1,C,0,1,0,C,C,C,0)

```

STORAGE NOT USED BY PROGRAM

DEC OCT
296 00450

DEC OCT
32561 77461

STORAGE LOCATIONS FOR VARIABLES NOT APPEARING IN COMMON, DIMENSION, OR EQUIVALENCE STATEMENTS

DEC OCT
A1 295 00447
NITT 291 00443
RR 287 00437
T 283 00433

DEC OCT
A2 294 00446
Q 290 00442
SUM1 286 00436
VI 282 00432

DEC OCT
FPRIME 293 00445
R1 289 00441
SUM2 285 00435

DEC OCT
F 292 00444
R2 288 00440
SUM3 284 00434

SYMBOLS AND LOCATIONS FOR SOURCE PROGRAM FORMAT STATEMENTS

EFN LOC
8)3S 124 00423

EFN LOC

EFN LOC

EFN LOC

LOCATIONS FOR OTHER SYMBOLS NOT APPEARING IN SOURCE PROGRAM

DEC OCT
1) 277 00425
9) 276 00424

DEC OCT
2) 243 00363
C)GC 281 00431

DEC OCT
3) 247 00367
E)2 103 00147

DEC OCT
6) 253 00375

LOCATIONS OF NAMES IN TRANSFER VECTOR

DEC OCT
(FIL) 1 00001

DEC OCT
(STH) 0 00000

DEC OCT

DEC OCT

ENTRY POINTS TO SUBROUTINES NOT OUTPUT FROM LIBRARY

(FIL) (STH)

EXTERNAL FORMULA NUMBERS WITH CORRESPONDING INTERNAL FORMULA NUMBERS AND OCTAL LOCATIONS

EFN IFN LOC
100 5 00102
103 12 00125
113 20 00162

EFN IFN LOC
105 9 00114
104 13 00130
127 22 00173

EFN IFN LOC
102 10 00117
107 14 00143
114 24 00224

EFN IFN LOC
101 11 00122
106 15 00150
128 25 00253

1199

08/15/64 PAGE 3

115 26 00261
119 30 00320
121 35 00353

116 27 00266
122 31 00326
8 36 00355

117 28 00302
123 32 00333

118 29 00310
120 34 00351

d. Sample Problem

INPUT - The following sample input data is for the calculation of escape probabilities and DB^2 for 1 to 4 inch void sleeves between the S8ER core and surrounding water, and for 20 to 100 mil gaps between the SNAP 10A core and reflector. The angular fluxes for the angular distribution were calculated by the DTK S_{12} approximation for the S8ER water flooded core surrounded by 3-inch void sleeve and 6-inch water reflector. The angular fluxes for the DB^2 calculation were calculated in the S8ER model by the DTK S_4 approximation for a 3-inch sleeve, and in the SNAP 10A model by the DTF II S_4 approximation for a 50 mil gap. The axial fission density distribution was calculated by the DTK S_4 approximation for a SNAP 10A-type reactor. ⁽¹⁰⁾ FOSDICK is presently set up to run multiple cases where the axial fission density distribution (not printed out) and the angular flux data for the escape probability calculation is common to all cases. Only the geometry of the void gap and the angular flux data for the DB^2 calculation need be respecified in successive cases. All data are entered in E12.8 format except the integer data which is placed to the extreme right of the data field.

FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 1 of 8 JOB NO. _____

NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1	*	
13	D A T A	
25	
37	
49	73.....80	
611	
1	6.0	Number of axial points, $M \leq 100$
13	
25	
37	
49	73.....80	
612	
1	. 8 2 1 6	Axial fission density distribution (equal intervals over half core) arbitrarily normalized. Enter M values starting from end of core to center. Format (6E12.8)
13	. 9 7 9 5	
25	1 . 1 0 7 2	
37	1 . 2 2 2 4	
49	1 . 3 3 0 0	
61	1 . 4 3 2 63
1	1 . 5 3 1 5	
13	1 . 6 2 7 8	
25	1 . 7 2 2 0	
37	1 . 8 1 4 3	
49	1 . 9 0 5 0	73.....80
61	1 . 9 9 4 24

FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 2 of 8 JOB NO. _____

	NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1	2 . 0 8 2 0		
13	2 . 1 6 8 5		
25	2 . 2 5 3 5		
37	2 . 3 3 7 1		
49	2 . 4 1 9 3		73: 80:
61	2 . 5 0 0 1	5	
1	2 . 5 7 9 5		
13	2 . 6 5 7 4		
25	2 . 7 3 3 9		
37	2 . 8 0 8 8		
49	2 . 8 8 2 1		73: 80:
61	2 . 9 5 4 0	6	
1	3 . 0 2 4 2		
13	3 . 0 9 2 8		
25	3 . 1 5 9 8		
37	3 . 2 2 5 1		
49	3 . 2 8 8 7		73: 80:
61	3 . 3 5 0 6	7	
1	3 . 4 1 0 8		
13	3 . 4 6 9 2		
25	3 . 5 2 5 9		
37	3 . 5 8 0 7		
49	3 . 6 3 3 7		73: 80:
61	3 . 6 8 4 9	8	

FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 3 of 8 JOB NO. _____

	NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1	.3 . . 7 3 4 2		
13	3 . 7 8 1 6		
25	3 . 8 2 7 1		
37	3 . . 8 7 0 6		
49	3 . 9 1 2 2		73 80
61	3 . 9 5 1 9		9
1	3 . 9 8 9 6		
13	4 . 0 2 5 2		
25	4 . 0 5 8 9		
37	4 . 0 9 0 5		
49	4 . 1 2 0 1		73 80
61	4 . 1 4 7 6		1 0
1	4 . 1 7 3 0		
13	4 . 1 9 6 4		
25	4 . 2 1 7 6		
37	4 . 2 3 6 8		
49	4 . 2 5 3 8		73 80
61	4 . 2 6 8 7		1 1
1	4 . 2 8 1 5		
13	4 . 2 9 2 2		
25	4 . 3 0 0 7		
37	4 . 3 0 7 0		
49	4 . 3 1 1 3		73 80
61	4 . 3 1 3 3		1 2

FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 4 of 8 JOB NO. _____

NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1	1 2	Order of S_n , $NSN \leq 16$
13	2 1	Number of Angular Fluxes, $MM = NSN(NSN+2)/8$
25	2 5	Maximum No. of iterations in DBSQ, MAX
37	. 0 0 0 0 1	Convergence criterion on DBSQ, EPS; Format (E12.8)
49	73 80	
61	1 3	
1	5 0 7 1 3 8 - 0 2	Angular Fluxes from DTK/DTF II calculation.
13	3 6 9 6 6 9 - 0 2	These values are group-summed, forward-angle fluxes at
25	1 2 0 6 2 9 - 0 1	the core/void boundary.
37	4 2 7 1 3 8 - 0 2	Enter MM values corresponding to appropriate direction
49	5 0 6 8 8 9 - 0 2	cosines.
61	1 7 4 5 4 2 - 0 1	Format (6E12.8)
1	4 6 5 5 0 7 - 0 2	
13	3 7 8 2 0 7 - 0 2	
25	8 5 2 1 0 0 - 0 2	
37	2 0 8 1 3 3 - 0 1	
49	4 8 9 0 0 1 - 0 2	
61	3 7 9 9 8 9 - 0 2	
1	4 9 7 7 0 8 - 0 2	
13	1 1 8 7 4 8 - 0 1	
25	2 2 9 6 0 7 - 0 1	
37	5 2 4 0 7 3 - 0 2	
49	4 2 6 6 0 9 - 0 2	
61	4 4 6 8 5 3 - 0 2	

FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 5 of 8 JOB NO. _____

	NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1	8 1 5 2 5 5 - 0 2		
13	1 7 5 8 7 8 - 0 1		
25	2 5 1 4 2 2 - 0 1		
37	. 1 7 4 0 8		S_n Direction Cosines, forward directions from
49	. 1 7 4 0 8	73 80	DTK/DTF II data. Enter MM values (in this order only).
61	. 4 6 0 5 7		
		1 7	
1	. 1 7 4 0 8		
13	. 4 6 0 5 7		
25	. 6 2 7 6 5		
37	. 1 7 4 0 8		
49	. 4 6 0 5 7	73 80	
61	. 6 2 7 6 5		
		1 8	
1	. 7 5 8 7 9		
13	. 1 7 4 0 8		
25	. 4 6 0 5 7		
37	. 6 2 7 6 5		
49	. 7 5 8 7 9	73 80	
61	. 8 7 0 3 9		
		1 9	
1	. 1 7 4 0 8		
13	. 4 6 0 5 7		
25	. 6 2 7 6 5		
37	. 7 5 8 7 9		
49	. 8 7 0 3 9	73 80	
61	. 9 6 9 2 2		
		2 0	

FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 6 of 8 JOB NO. _____

	NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1	. 0 3 1 8 8		S_n Point Weights from DTK/DTF II data.
13	. 0 2 8 4 5		Enter MM values corresponding to direction cosines.
25	. 0 2 8 4 5		
37	. 0 2 2 6 9		
49	. 0 1 7 3 7	73	80
61	. 0 2 2 6 9		2 1
1	. 0 2 2 6 9		
13	. 0 1 5 1 4		
25	. 0 1 5 1 4		
37	. 0 2 2 6 9		
49	. 0 2 8 4 5	73	80
61	. 0 1 7 3 7		2 2
1	. 0 1 5 1 4		
13	. 0 1 7 3 7		
25	. 0 2 8 4 5		
37	. 0 3 1 8 8		
49	. 0 2 8 4 5	73	80
61	. 0 2 2 6 9		2 3
1	. 0 2 2 6 9		
13	. 0 2 8 4 5		
25	. 0 3 1 8 8		
37			
49		73	80
61			2 4

FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 7 of 8 JOB NO. _____

NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH
1		Core radius, R (inches)
13		Maximum reflector radius, XMAX (inches)
25		Minimum reflector radius, XMIN (inches)
37		Core length, L (inches)
49	4 73 80	Number of gaps, $N \leq 20$ (each using same axial and angular distributions input)
61	2 5	
1		} S_4 Angular Fluxes, (3 values)
13		
25		
37		} S_4 Direction Cosines, (3 values)
49	73 80	
61	2 6	
1		} S_4 Point Weights (3 values)
13		
25		
37		
49	73 80	
61	2 7	
1		R Second Case
13		XMAX
25		XMIN
37		L
49	5 73 80	N
61	2 8	

FORTRAN FIXED 10 DIGIT DECIMAL DATA

DECK NO. _____ PROGRAMMER _____ DATE _____ PAGE 8 of 8 JOB NO. _____

NUMBER	IDENTIFICATION	DESCRIPTION DO NOT KEY PUNCH	
1		} S_4 Angular Fluxes	
13			. 0 3 5 0 1 3 2 3 .
25			. 0 3 0 9 1 0 4 0
37		} S_4 Direction Cosines	
49			. 3 3 3 3 3 3 3
61			. 8 8 1 9 1 7 1 0
	2 9		
1		} S_4 Point Weights	
13			. 1 6 6 6 6 6 6 7
25			. 1 6 6 6 6 6 6 7
37			
49		7 3	8 0
61		3 0	
1			
13			
25			
37			
49		7 3	8 0
61			
1			
13			
25			
37			
49		7 3	8 0
61			

ATOMICS INTERNATIONAL

A Division of North American Aviation, Inc.

OUTPUT - The following sample output should be quite self-explanatory. When running multiple cases, each new case begins with "THE FOSDICK CODE"; all input data, whether common between cases or not, is printed out for each case. All spatial dimensions are in inches, and SIGMA (DBSQ) has units of cm^{-1} . The execution time for this problem was 32 seconds on the IBM 7094 computer.

* DATA
BEGIN LOADING 25

01-44-03 08/15/64

00000010

NAME	LOAD	ENTRY	NAME	LOAD	ENTRY	NAME	LOAD	ENTRY	NAME	LOAD	ENTRY
000000	00200	00212	ANGLE	25606	25617	NORM	26603	26612	GPESC	27060	27072
DBSQ	30627	30635									

SUBROUTINES LOADED FROM LIBRARY.

(FPT) 31277 31306	.FPTRP*31277*31306	DUMP 31501 31527	PDUMP *31501*31517
XDUMP 31501 31511	(KEYS) 31501 31507	DDUMP *31501*31523	(KIKL) 31501 32116
EXIT 32503 32511	(TES) 32550 32554	(CDIN) 32550 32567	(LQUT) 32550 32572
(CDXC)*32550*32571	(TAPU) 32550 32561	(CTES) 32550 32555	(TCRT)*32550*32557
(IDCC)*32550*32562	(WDCT) 32550 32566	(CTJB)*32550*32570	(WFT) 32550 32576
(FTC) 32550 32643	(NUSE)*32550*32573	(DATE)*32550*32560	(TIME)*32550*32556
(JOBX)*32550*32563	(BOXN)*32550*32564	(LSEC)*32550*32565	(BACAL 32550 32574
(FIBBX 32550 32575	SIN 32646 32666	SINE *32646*32666	SINI *32646*32666
SIND *32646*32653	COS 32646 32660	COSINE*32646*32660	CCS1 *32646*32660
COSD *32646*32656	ECSIN *32646*33054	EPSIN 33061 33061	SCRT 33074 33101
ECSQRT*33074*33177	EPSQRT 33203 33203	ATN1 *33205*33311	ATAN *33205*33311
ARCTAN*33205*33311	ATAND *33205*33212	QATAN 33205 33220	QATAND*33205*33222
ECQTAN*33205*33442	EPQTAN 33446 33446	ESCORT 33455 33471	ESCSWT*33455*34011
PAKUP) 33455 33673	ERRG0 *33455*34014	(ESOUT 34023 34025	(ESCFI 34023 34026
(SESC) 34027 34027	(ESC0) 34030 34030	TRAPEX 34031 34032	(IGS) 34033 34037
(RDS) 34033 34145	(WRS) 34033 34146	(BSR) 34033 34147	(WEF) 34033 34150
(REW) 34033 34151	(ETT) 34033 34156	(RCH) 34033 34160	(TEF) 34033 34161
(TC0) 34033 34162	(TRC) 34033 34163	(RUN) 34033 34152	(BTT) *34033*34157
(STC) *34033*34164	(LCH) *34033*34165	(BSF) *34033*34153	(SDH) *34033*34154
(SDL) *34033*34155	(UBCD)*34033*34166	(IGU) 34243 34246	(TSH) 34267 34275
(STH) 34317 34334	(CNTV)*34317*34370	(CNTL)*34317*34376	(WER) 34457 34472
(WTC) 34457 34522	(RER) 34600 34612	(RDC) 34600 34647	(IGH) 34745 34753
(FIL) 34745 36641	(RTN) 34745 36652	PRINT) 37026 37026	(SCRTV 37313 37314
(YTOP)*37313*37316	(YREG)*37313*37317	(CTAPE 37313 37315	

LOWEST COMMON LOC. (77461)8 HIGHEST PROGRAM LOC.(37317)8 STORAGE NOT USED BY PROGRAM. (40141)8
BEGIN EXECUTION 25 01-44-12

THE FOSDICK CODE

* * * INPUT DATA * * *

CYLINDER RADIUS, R= 4.71417

RADIUS OF MEDIUM (MAX.), X-MAX= 8.71417

RADIUS OF MEDIUM (MIN.), X-MIN= 4.71417

CYLINDER LENGTH, L=14.00000

NUMBER OF X INTERVALS, N= 4

NUMBER OF AXIAL POINTS, M= 60

ORDER OF SN = 12

CONVERGENCE CRITERION = 0.000010 MAX. NO. ITERATIONS = 25

ANGULAR FLUX	DIRECTION COSINES	POINT WEIGHTS
0.507138E-02	0.174080	0.031880
0.369669E-02	0.174080	0.028450
0.120629E-01	0.460570	0.028450
0.427138E-02	0.174080	0.022690
0.506889E-02	0.460570	0.017370
0.174542E-01	0.627650	0.022690
0.465507E-02	0.174080	0.022690
0.378207E-02	0.460570	0.015140
0.852100E-02	0.627650	0.015140
0.208133E-01	0.758790	0.022690
0.489001E-02	0.174080	0.028450
0.379989E-02	0.460570	0.017370
0.497708E-02	0.627650	0.015140
0.118748E-01	0.758790	0.017370
0.229607E-01	0.870390	0.028450
0.524073E-02	0.174080	0.031880
0.426609E-02	0.460570	0.028450
0.446853E-02	0.627650	0.022690

0.815255E-02	0.758790	0.022690
0.175878E-01	0.870390	0.028450
0.251422E-01	0.969220	0.031880

S-4 FLUXES FOR DBSQ CALCULATION		
ANGULAR FLUX	DIRECTION COSINES	POINT WEIGHTS
0.218511E-01	0.333333	0.166667
0.168993E-01	0.333333	0.166667
0.300621E-01	0.881917	0.166667

* * * OUTPUT * * *

ANGULAR DISTRIBUTION,	SINE	FLUX
-0.		0.244553E-02
0.24620		0.507138E-02
0.49236		0.787979E-02
0.65134		0.925896E-02
0.77850		0.100997E-01
0.88762		0.106759E-01
0.98473		0.113028E-01
1.00000		0.114835E-01
0.98473		0.113028E-01
0.88762		0.106759E-01
0.77850		0.100997E-01
0.65134		0.925896E-02
0.49236		0.787979E-02
0.24620		0.507138E-02
0.		0.244553E-02

ANGULAR DISTRIBUTION NORMALIZATION FACTOR = 0.100049E-01

FOR T= 1.0000 AVERAGE ESCAPE PROBABILITY, PC-ANISOTROPIC = 0.46621E-01

CONVERGENCE = 0.000002 SIGMA (DBSQ) = 0.0107013

Z	P1	P2	P3	P4	PC
0.11667	0.42880E-00	0.18084E-01	0.85452E-03	0.38179E-03	0.44812E-00
0.23333	0.36584E-00	0.34464E-01	0.87121E-03	0.38883E-03	0.40156E-00

0.35000	0.30850E-00	0.48145E-C1	0.88838E-C3	0.39607E-C3	0.35793E-00
0.46667	0.25814E-00	0.58536E-C1	0.90604E-C3	0.40350E-C3	0.31799E-00
0.58333	0.21514E-00	0.65621E-C1	0.92422E-C3	0.41115E-C3	0.28209E-00
0.70000	0.17882E-00	0.69834E-C1	0.94294E-C3	0.41901E-C3	0.25002E-00
0.81667	0.14853E-00	0.71672E-C1	0.96221E-C3	0.42710E-C3	0.22160E-00
0.93333	0.12366E-00	0.71634E-C1	0.98205E-C3	0.43543E-C3	0.19671E-00
1.05000	0.10358E-00	0.70228E-C1	0.10025E-C2	0.44399E-C3	0.17526E-00
1.16667	0.87214E-01	0.67860E-C1	0.10236E-C2	0.45281E-C3	0.15655E-00
1.28333	0.73857E-01	0.64891E-C1	0.10453E-C2	0.46190E-C3	0.14026E-00
1.40000	0.62782E-01	0.61574E-C1	0.10677E-C2	0.47125E-C3	0.12590E-00
1.51667	0.53757E-01	0.58092E-C1	0.10908E-C2	0.48089E-C3	0.11342E-00
1.63333	0.46356E-01	0.54515E-C1	0.11146E-C2	0.49083E-C3	0.10248E-00
1.75000	0.40175E-01	0.51042E-C1	0.11392E-C2	0.50108E-C3	0.92857E-01
1.86667	0.35021E-01	0.47723E-C1	0.11646E-C2	0.51165E-C3	0.84420E-01
1.98333	0.30726E-01	0.44523E-C1	0.11908E-C2	0.52255E-C3	0.76963E-01
2.10000	0.27113E-01	0.41508E-C1	0.12179E-C2	0.53380E-C3	0.70373E-01
2.21667	0.24054E-01	0.38675E-C1	0.12458E-C2	0.54541E-C3	0.64521E-01
2.33333	0.21448E-01	0.36031E-C1	0.12747E-C2	0.55741E-C3	0.59312E-01
2.45000	0.19215E-01	0.33574E-C1	0.13046E-C2	0.56980E-C3	0.54664E-01
2.56667	0.17291E-01	0.31263E-C1	0.13355E-C2	0.58260E-C3	0.50472E-01
2.68333	0.15624E-01	0.29138E-C1	0.13675E-C2	0.59584E-C3	0.46725E-01
2.80000	0.14172E-01	0.27184E-01	0.14006E-C2	0.60952E-C3	0.43367E-01
2.91667	0.12903E-01	0.25388E-C1	0.14349E-C2	0.62368E-C3	0.40350E-01
3.03333	0.11789E-01	0.23733E-C1	0.14703E-C2	0.63833E-C3	0.37631E-01
3.15000	0.10810E-01	0.22204E-C1	0.15071E-C2	0.65350E-C3	0.35175E-01
3.26667	0.99414E-02	0.20798E-C1	0.15452E-C2	0.66921E-C3	0.32954E-01
3.38333	0.91685E-02	0.19504E-C1	0.15847E-C2	0.68548E-C3	0.30942E-01
3.50000	0.84783E-02	0.18305E-C1	0.16257E-C2	0.70235E-C3	0.29112E-01
3.61667	0.78597E-02	0.17184E-C1	0.16683E-C2	0.71984E-C3	0.27432E-01
3.73333	0.73035E-02	0.16149E-01	0.17125E-02	0.73799E-C3	0.25903E-01
3.85000	0.68035E-02	0.15194E-01	0.17584E-02	0.75682E-C3	0.24513E-01
3.96667	0.63514E-02	0.14311E-01	0.18061E-02	0.77638E-C3	0.23245E-01
4.08333	0.59411E-02	0.13495E-01	0.18557E-02	0.79669E-C3	0.22088E-01
4.20000	0.55679E-02	0.12738E-01	0.19072E-02	0.81779E-C3	0.21031E-01
4.31667	0.52275E-02	0.12038E-01	0.19608E-02	0.83974E-C3	0.20066E-01
4.43333	0.49163E-02	0.11389E-01	0.20166E-02	0.86257E-C3	0.19184E-01
4.55000	0.46312E-02	0.10786E-01	0.20747E-02	0.88633E-C3	0.18378E-01
4.66667	0.43693E-02	0.10224E-01	0.21352E-02	0.91107E-C3	0.17640E-01
4.78333	0.41284E-02	0.96969E-C2	0.21982E-02	0.93685E-C3	0.16960E-01
4.90000	0.39062E-02	0.92055E-02	0.22639E-C2	0.96372E-C3	0.16339E-01
5.01667	0.37010E-02	0.87471E-02	0.23325E-C2	0.99174E-C3	0.15772E-01
5.13333	0.35110E-02	0.83189E-02	0.24040E-02	0.10210E-02	0.15255E-01
5.25000	0.33349E-02	0.79186E-02	0.24787E-02	0.10515E-02	0.14784E-01
5.36667	0.31714E-02	0.75440E-02	0.25568E-C2	0.10834E-02	0.14356E-01

5.48333	0.30192E-02	0.71932E-02	0.26384E-02	0.11168E-02	0.13968E-01
5.60000	0.28775E-02	0.68638E-02	0.27237E-02	0.11517E-02	0.13617E-01
5.71667	0.27452E-02	0.65547E-02	0.28130E-02	0.11882E-02	0.13301E-01
5.83333	0.26217E-02	0.62648E-02	0.29065E-02	0.12265E-02	0.13019E-01
5.95000	0.25061E-02	0.59930E-02	0.30049E-02	0.12666E-02	0.12771E-01
6.06667	0.23978E-02	0.57372E-02	0.31080E-02	0.13087E-02	0.12552E-01
6.18333	0.22962E-02	0.54962E-02	0.32162E-02	0.13528E-02	0.12361E-01
6.30000	0.22008E-02	0.52689E-02	0.33298E-02	0.13993E-02	0.12199E-01
6.41667	0.21111E-02	0.50544E-02	0.34491E-02	0.14480E-02	0.12063E-01
6.53333	0.20266E-02	0.48519E-02	0.35745E-02	0.14994E-02	0.11952E-01
6.65000	0.19470E-02	0.46604E-02	0.37063E-02	0.15534E-02	0.11867E-01
6.76667	0.18719E-02	0.44793E-02	0.38451E-02	0.16104E-02	0.11807E-01
6.88333	0.18010E-02	0.43078E-02	0.39913E-02	0.16705E-02	0.11770E-01
7.00000	0.17339E-02	0.41453E-02	0.41453E-02	0.17339E-02	0.11758E-01

FOR T= 2.0000

AVERAGE ESCAPE PROBABILITY, PC-ANISOTROPIC = 0.99786E-01

CONVERGENCE = 0.000001

SIGMA (DBSQ) = 0.0114974

Z	F1	P2	P3	P4	PC
0.11667	0.46188E-00	0.76847E-02	0.24610E-02	0.17670E-02	0.47379E-00
0.23333	0.42880E-00	0.15159E-01	0.25096E-02	0.18010E-02	0.44827E-00
0.35000	0.39680E-00	0.22228E-01	0.25596E-02	0.18359E-02	0.42342E-00
0.46667	0.36584E-00	0.28779E-01	0.26109E-02	0.18719E-02	0.39910E-00
0.58333	0.33634E-00	0.34729E-01	0.26638E-02	0.19089E-02	0.37564E-00
0.70000	0.30850E-00	0.39983E-01	0.27182E-02	0.19470E-02	0.35315E-00
0.81667	0.28246E-00	0.44494E-01	0.27742E-02	0.19862E-02	0.33171E-00
0.93333	0.25814E-00	0.48247E-01	0.28320E-02	0.20266E-02	0.31125E-00
1.05000	0.23569E-00	0.51258E-01	0.28915E-02	0.20682E-02	0.29191E-00
1.16667	0.21514E-00	0.53590E-01	0.29529E-02	0.21111E-02	0.27379E-00
1.28333	0.19615E-00	0.55302E-01	0.30161E-02	0.21552E-02	0.25662E-00
1.40000	0.17882E-00	0.56452E-01	0.30812E-02	0.22008E-02	0.24055E-00
1.51667	0.16309E-00	0.57090E-01	0.31482E-02	0.22477E-02	0.22558E-00
1.63333	0.14853E-00	0.57308E-01	0.32174E-02	0.22962E-02	0.21136E-00
1.75000	0.13544E-00	0.57121E-01	0.32887E-02	0.23462E-02	0.19819E-00
1.86667	0.12366E-00	0.56606E-01	0.33623E-02	0.23978E-02	0.18603E-00
1.98333	0.11309E-00	0.55811E-01	0.34381E-02	0.24511E-02	0.17479E-00
2.10000	0.10358E-00	0.54790E-01	0.35164E-02	0.25061E-02	0.16439E-00
2.21667	0.94997E-01	0.53605E-01	0.35972E-02	0.25629E-02	0.15476E-00
2.33333	0.87214E-01	0.52297E-01	0.36807E-02	0.26217E-02	0.14581E-00
2.45000	0.80214E-01	0.50886E-01	0.37668E-02	0.26824E-02	0.13755E-00
2.56667	0.73857E-01	0.49355E-01	0.38558E-02	0.27452E-02	0.12981E-00
2.68333	0.68032E-01	0.47791E-01	0.39478E-02	0.28102E-02	0.12258E-00

2.80000	0.62782E-01	0.46211E-01	0.40428E-02	0.28775E-02	0.11591E-00
2.91667	0.58043E-01	0.44603E-01	0.41410E-02	0.29471E-02	0.10973E-00
3.03333	0.53757E-01	0.42991E-01	0.42427E-02	0.30192E-02	0.10401E-00
3.15000	0.49877E-01	0.41405E-01	0.43478E-02	0.30939E-02	0.98723E-01
3.26667	0.46356E-01	0.39848E-01	0.44565E-02	0.31714E-02	0.93832E-01
3.38333	0.43135E-01	0.38332E-01	0.45691E-02	0.32516E-02	0.89287E-01
3.50000	0.40175E-01	0.36861E-01	0.46856E-02	0.33349E-02	0.85056E-01
3.61667	0.37478E-01	0.35439E-01	0.48063E-02	0.34213E-02	0.81144E-01
3.73333	0.35021E-01	0.34042E-01	0.49313E-02	0.35110E-02	0.77505E-01
3.85000	0.32779E-01	0.32669E-01	0.50615E-02	0.36042E-02	0.74114E-01
3.96667	0.30726E-01	0.31352E-01	0.51965E-02	0.37010E-02	0.70976E-01
4.08333	0.28843E-01	0.30091E-01	0.53365E-02	0.38016E-02	0.68072E-01
4.20000	0.27113E-01	0.28883E-01	0.54817E-02	0.39062E-02	0.65385E-01
4.31667	0.25521E-01	0.27729E-01	0.56324E-02	0.40151E-02	0.62898E-01
4.43333	0.24054E-01	0.26627E-01	0.57889E-02	0.41284E-02	0.60598E-01
4.55000	0.22700E-01	0.25549E-01	0.59513E-02	0.42464E-02	0.58447E-01
4.66667	0.21448E-01	0.24517E-01	0.61201E-02	0.43693E-02	0.56455E-01
4.78333	0.20289E-01	0.23529E-01	0.62954E-02	0.44975E-02	0.54612E-01
4.90000	0.19215E-01	0.22588E-01	0.64777E-02	0.46312E-02	0.52912E-01
5.01667	0.18218E-01	0.21692E-01	0.66673E-02	0.47707E-02	0.51348E-01
5.13333	0.17291E-01	0.20838E-01	0.68644E-02	0.49163E-02	0.49910E-01
5.25000	0.16428E-01	0.20026E-01	0.70696E-02	0.50685E-02	0.48592E-01
5.36667	0.15624E-01	0.19252E-01	0.72831E-02	0.52275E-02	0.47386E-01
5.48333	0.14873E-01	0.18514E-01	0.75055E-02	0.53938E-02	0.46287E-01
5.60000	0.14172E-01	0.17811E-01	0.77372E-02	0.55679E-02	0.45289E-01
5.71667	0.13517E-01	0.17142E-01	0.79789E-02	0.57501E-02	0.44388E-01
5.83333	0.12903E-01	0.16494E-01	0.82313E-02	0.59411E-02	0.43569E-01
5.95000	0.12327E-01	0.15875E-01	0.84944E-02	0.61413E-02	0.42838E-01
6.06667	0.11789E-01	0.15285E-01	0.87688E-02	0.63514E-02	0.42194E-01
6.18333	0.11285E-01	0.14722E-01	0.90552E-02	0.65719E-02	0.41634E-01
6.30000	0.10810E-01	0.14186E-01	0.93542E-02	0.68035E-02	0.41154E-01
6.41667	0.10363E-01	0.13675E-01	0.96665E-02	0.70471E-02	0.40751E-01
6.53333	0.99414E-02	0.13187E-01	0.99928E-02	0.73035E-02	0.40425E-01
6.65000	0.95439E-02	0.12722E-01	0.10334E-01	0.75743E-02	0.40174E-01
6.76667	0.91685E-02	0.12278E-01	0.10691E-01	0.78597E-02	0.39997E-01
6.88333	0.88138E-02	0.11854E-01	0.11062E-01	0.81606E-02	0.39891E-01
7.00000	0.84783E-02	0.11449E-01	0.11449E-01	0.84783E-02	0.39855E-01

FOR T= 3.0000

AVERAGE ESCAPE PROBABILITY, PC-ANISOTROPIC = 0.15318E-00

CONVERGENCE = 0.0000001

SIGMA (DBSQ) = 0.0116863

Z

P1

P2

P3

P4

PC

0.11667	0.47302E-00	0.44624E-02	0.45043E-02	0.44542E-02	0.48644E-00
0.23333	0.45077E-00	0.88671E-02	0.45919E-02	0.45414E-02	0.46877E-00
0.35000	0.42880E-00	0.13159E-01	0.46819E-02	0.46312E-02	0.45127E-00
0.46667	0.40740E-00	0.17285E-01	0.47743E-02	0.47235E-02	0.43419E-00
0.58333	0.38633E-00	0.21205E-01	0.48692E-02	0.48185E-02	0.41722E-00
0.70000	0.36584E-00	0.24913E-01	0.49668E-02	0.49163E-02	0.40063E-00
0.81667	0.34599E-00	0.28380E-01	0.50670E-02	0.50170E-02	0.38446E-00
0.93333	0.32686E-00	0.31575E-01	0.51701E-02	0.51207E-02	0.36873E-00
1.05000	0.30850E-00	0.34480E-01	0.52760E-02	0.52275E-02	0.35348E-00
1.16667	0.29093E-00	0.37085E-01	0.53849E-02	0.53375E-02	0.33874E-00
1.28333	0.27419E-00	0.39386E-01	0.54969E-02	0.54510E-02	0.32452E-00
1.40000	0.25814E-00	0.41384E-01	0.56122E-02	0.55679E-02	0.31071E-00
1.51667	0.24296E-00	0.43086E-01	0.57307E-02	0.56884E-02	0.29747E-00
1.63333	0.22863E-00	0.44508E-01	0.58526E-02	0.58128E-02	0.28480E-00
1.75000	0.21514E-00	0.45672E-01	0.59785E-02	0.59411E-02	0.27273E-00
1.86667	0.20231E-00	0.46598E-01	0.61081E-02	0.60735E-02	0.26109E-00
1.98333	0.19018E-00	0.47296E-01	0.62416E-02	0.62102E-02	0.24993E-00
2.10000	0.17882E-00	0.47778E-01	0.63790E-02	0.63514E-02	0.23933E-00
2.21667	0.16819E-00	0.48072E-01	0.65206E-02	0.64972E-02	0.22928E-00
2.33333	0.15807E-00	0.48192E-01	0.66664E-02	0.66478E-02	0.21958E-00
2.45000	0.14853E-00	0.48153E-01	0.68167E-02	0.68035E-02	0.21031E-00
2.56667	0.13965E-00	0.47957E-01	0.69715E-02	0.69645E-02	0.20154E-00
2.68333	0.13137E-00	0.47631E-01	0.71312E-02	0.71310E-02	0.19326E-00
2.80000	0.12366E-00	0.47198E-01	0.72957E-02	0.73035E-02	0.18546E-00
2.91667	0.11649E-00	0.46646E-01	0.74655E-02	0.74825E-02	0.17808E-00
3.03333	0.10980E-00	0.46011E-01	0.76405E-02	0.76678E-02	0.17112E-00
3.15000	0.10358E-00	0.45310E-01	0.78211E-02	0.78597E-02	0.16457E-00
3.26667	0.97781E-01	0.44552E-01	0.80074E-02	0.80585E-02	0.15840E-00
3.38333	0.92310E-01	0.43749E-01	0.81996E-02	0.82646E-02	0.15252E-00
3.50000	0.87214E-01	0.42897E-01	0.83981E-02	0.84783E-02	0.14699E-00
3.61667	0.82466E-01	0.42013E-01	0.86028E-02	0.86999E-02	0.14178E-00
3.73333	0.78039E-01	0.41085E-01	0.88142E-02	0.89299E-02	0.13687E-00
3.85000	0.73857E-01	0.40145E-01	0.90325E-02	0.91685E-02	0.13220E-00
3.96667	0.69907E-01	0.39198E-01	0.92580E-02	0.94164E-02	0.12778E-00
4.08333	0.66222E-01	0.38250E-01	0.94909E-02	0.96738E-02	0.12364E-00
4.20000	0.62782E-01	0.37283E-01	0.97317E-02	0.99414E-02	0.11974E-00
4.31667	0.59569E-01	0.36310E-01	0.99805E-02	0.10220E-01	0.11608E-00
4.43333	0.56566E-01	0.35348E-01	0.10238E-01	0.10509E-01	0.11266E-00
4.55000	0.53757E-01	0.34400E-01	0.10504E-01	0.10810E-01	0.10947E-00
4.66667	0.51128E-01	0.33468E-01	0.10778E-01	0.11123E-01	0.10650E-00
4.78333	0.48665E-01	0.32553E-01	0.11060E-01	0.11449E-01	0.10373E-00
4.90000	0.46356E-01	0.31656E-01	0.11352E-01	0.11789E-01	0.10115E-00
5.01667	0.44185E-01	0.30780E-01	0.11653E-01	0.12144E-01	0.98762E-01
5.13333	0.42117E-01	0.29910E-01	0.11966E-01	0.12515E-01	0.96508E-01

5.25000	0.40175E-01	0.29061E-01	0.12289E-01	0.12903E-01	0.94427E-01
5.36667	0.38349E-01	0.28233E-01	0.12624E-01	0.13308E-01	0.92514E-01
5.48333	0.36633E-01	0.27429E-01	0.12970E-01	0.13730E-01	0.90762E-01
5.60000	0.35021E-01	0.26634E-01	0.13328E-01	0.14172E-01	0.89156E-01
5.71667	0.33504E-01	0.25861E-01	0.13700E-01	0.14634E-01	0.87699E-01
5.83333	0.32075E-01	0.25110E-01	0.14084E-01	0.15118E-01	0.86387E-01
5.95000	0.30726E-01	0.24383E-01	0.14482E-01	0.15624E-01	0.85216E-01
6.06667	0.29453E-01	0.23679E-01	0.14895E-01	0.16154E-01	0.84181E-01
6.18333	0.28250E-01	0.22995E-01	0.15322E-01	0.16709E-01	0.83276E-01
6.30000	0.27113E-01	0.22331E-01	0.15756E-01	0.17291E-01	0.82491E-01
6.41667	0.26037E-01	0.21682E-01	0.16206E-01	0.17901E-01	0.81827E-01
6.53333	0.25019E-01	0.21045E-01	0.16671E-01	0.18542E-01	0.81277E-01
6.65000	0.24054E-01	0.20429E-01	0.17153E-01	0.19215E-01	0.80852E-01
6.76667	0.23140E-01	0.19835E-01	0.17653E-01	0.19922E-01	0.80549E-01
6.88333	0.22272E-01	0.19261E-01	0.18170E-01	0.20666E-01	0.80368E-01
7.00000	0.21448E-01	0.18706E-01	0.18706E-01	0.21448E-01	0.80308E-01

FOR T= 4.0000

AVERAGE ESCAPE PROBABILITY, PC-ANISOTROPIC = C.20473E-00

CONVERGENCE = 0.000001

SIGMA (DBSQ) = C.0115432

Z	P1	P2	P3	P4	PC
0.11667	0.47860E-00	0.29723E-02	0.67769E-02	0.86437E-02	0.49699E-00
0.23333	0.46188E-00	0.59221E-02	0.69033E-02	0.88138E-02	0.48351E-00
0.35000	0.44523E-00	0.88273E-02	0.70329E-02	0.89887E-02	0.47008E-00
0.46667	0.42880E-00	0.11667E-01	0.71657E-02	0.91685E-02	0.45680E-00
0.58333	0.41272E-00	0.14421E-01	0.73019E-02	0.93535E-02	0.44379E-00
0.70000	0.39680E-00	0.17070E-01	0.74415E-02	0.95439E-02	0.43086E-00
0.81667	0.38115E-00	0.19613E-01	0.75847E-02	0.97398E-02	0.41809E-00
0.93333	0.36584E-00	0.22044E-01	0.77316E-02	0.99414E-02	0.40555E-00
1.05000	0.35089E-00	0.24351E-01	0.78822E-02	0.10149E-01	0.39327E-00
1.16667	0.33634E-00	0.26521E-01	0.80368E-02	0.10363E-01	0.38126E-00
1.28333	0.32220E-00	0.28545E-01	0.81954E-02	0.10583E-01	0.36952E-00
1.40000	0.30850E-00	0.30418E-01	0.83581E-02	0.10810E-01	0.35809E-00
1.51667	0.29525E-00	0.32136E-01	0.85250E-02	0.11044E-01	0.34695E-00
1.63333	0.28246E-00	0.33699E-01	0.86962E-02	0.11285E-01	0.33614E-00
1.75000	0.27012E-00	0.35105E-01	0.88720E-02	0.11533E-01	0.32563E-00
1.86667	0.25814E-00	0.36358E-01	0.90525E-02	0.11789E-01	0.31534E-00
1.98333	0.24668E-00	0.37460E-01	0.92379E-02	0.12054E-01	0.30543E-00
2.10000	0.23569E-00	0.38416E-01	0.94281E-02	0.12327E-01	0.29586E-00
2.21667	0.22518E-00	0.39240E-01	0.96236E-02	0.12610E-01	0.28665E-00
2.33333	0.21514E-00	0.39932E-01	0.98243E-02	0.12903E-01	0.27779E-00
2.45000	0.20547E-00	0.40510E-01	0.10030E-01	0.13205E-01	0.26922E-00

2.56667	0.19615E-00	0.40969E-C1	0.10242E-01	0.13517E-01	0.26088E-00
2.68333	0.18727E-00	0.41316E-01	0.10458E-01	0.13839E-01	0.25288E-00
2.80000	0.17882E-00	0.41562E-01	0.10679E-01	0.14172E-01	0.24523E-00
2.91667	0.17078E-00	0.41715E-01	0.10907E-01	0.14517E-01	0.23792E-00
3.03333	0.16309E-00	0.41778E-01	0.11141E-01	0.14873E-01	0.23089E-00
3.15000	0.15562E-00	0.41758E-01	0.11382E-01	0.15242E-01	0.22400E-00
3.26667	0.14853E-00	0.41662E-01	0.11629E-01	0.15624E-01	0.21745E-00
3.38333	0.14181E-00	0.41491E-01	0.11883E-01	0.16019E-01	0.21120E-00
3.50000	0.13544E-00	0.41252E-01	0.12144E-01	0.16428E-01	0.20526E-00
3.61667	0.12939E-00	0.40956E-01	0.12413E-01	0.16852E-01	0.19961E-00
3.73333	0.12366E-00	0.40602E-01	0.12689E-01	0.17291E-01	0.19424E-00
3.85000	0.11823E-00	0.40194E-01	0.12973E-C1	0.17746E-01	0.18915E-00
3.96667	0.11309E-00	0.39747E-01	0.13264E-01	0.18218E-01	0.18432E-00
4.08333	0.10821E-00	0.39266E-01	0.13564E-01	0.18707E-01	0.17974E-00
4.20000	0.10358E-00	0.38756E-01	0.13873E-01	0.19215E-01	0.17543E-00
4.31667	0.99197E-01	0.38220E-01	0.14184E-01	0.19742E-01	0.17134E-00
4.43333	0.94997E-01	0.37653E-01	0.14504E-01	0.20289E-01	0.16744E-00
4.55000	0.91002E-01	0.37068E-01	0.14832E-01	0.20858E-01	0.16376E-00
4.66667	0.87214E-01	0.36469E-01	0.15170E-01	0.21448E-01	0.16030E-00
4.78333	0.83622E-01	0.35859E-01	0.15518E-01	0.22062E-01	0.15706E-00
4.90000	0.80214E-01	0.35225E-01	0.15875E-01	0.22700E-01	0.15401E-00
5.01667	0.76979E-01	0.34585E-01	0.16242E-01	0.23364E-01	0.15117E-00
5.13333	0.73857E-01	0.33942E-01	0.16619E-01	0.24054E-01	0.14847E-00
5.25000	0.70869E-01	0.33294E-01	0.17007E-01	0.24773E-01	0.14594E-00
5.36667	0.68032E-01	0.32648E-01	0.17405E-01	0.25521E-01	0.14361E-00
5.48333	0.65340E-01	0.31984E-01	0.17815E-01	0.26301E-01	0.14144E-00
5.60000	0.62782E-01	0.31325E-01	0.18236E-01	0.27113E-01	0.13946E-00
5.71667	0.60352E-01	0.30673E-01	0.18669E-01	0.27960E-01	0.13765E-00
5.83333	0.58043E-01	0.30027E-01	0.19104E-01	0.28843E-01	0.13602E-00
5.95000	0.55847E-01	0.29389E-01	0.19548E-01	0.29765E-01	0.13455E-00
6.06667	0.53757E-01	0.28760E-01	0.20003E-01	0.30726E-01	0.13325E-00
6.18333	0.51769E-01	0.28140E-01	0.20471E-01	0.31730E-01	0.13211E-00
6.30000	0.49877E-01	0.27530E-01	0.20950E-01	0.32779E-01	0.13114E-00
6.41667	0.48074E-01	0.26925E-01	0.21442E-01	0.33875E-01	0.13032E-00
6.53333	0.46356E-01	0.26325E-01	0.21946E-01	0.35021E-01	0.12965E-00
6.65000	0.44718E-01	0.25738E-01	0.22459E-01	0.36220E-01	0.12914E-00
6.76667	0.43135E-01	0.25162E-01	0.22979E-01	0.37478E-01	0.12875E-00
6.88333	0.41620E-01	0.24598E-01	0.23507E-01	0.38795E-01	0.12852E-00
7.00000	0.40175E-01	0.24046E-01	0.24046E-01	0.40175E-01	0.12844E-00

THE FOSDICK CODE

* * * INPUT DATA * * *

CYLINDER RADIUS, R= 4.50000

RADIUS OF MEDIUM (MAX.), X-MAX= 4.60000

RADIUS OF MEDIUM (MIN.), X-MIN= 4.50000

CYLINDER LENGTH, L=12.25000

NUMBER OF X INTERVALS, N= 5

NUMBER OF AXIAL POINTS, M= 60

ORDER OF SN = 12

CONVERGENCE CRITERION = 0.000010 MAX. NO. ITERATIONS = 25

ANGULAR FLUX	DIRECTION COSINES	POINT WEIGHTS
0.507138E-02	0.174080	0.031880
0.369669E-02	0.174080	0.028450
0.120629E-01	0.460570	0.028450
0.427138E-02	0.174080	0.022690
0.506889E-02	0.460570	0.017370
0.174542E-01	0.627650	0.022690
0.465507E-02	0.174080	0.022690
0.378207E-02	0.460570	0.015140
0.852100E-02	0.627650	0.015140
0.208133E-01	0.758790	0.022690
0.489001E-02	0.174080	0.028450
0.379989E-02	0.460570	0.017370
0.497708E-02	0.627650	0.015140
0.118748E-01	0.758790	0.017370
0.229607E-01	0.870390	0.028450
0.524073E-02	0.174080	0.031880
0.426609E-02	0.460570	0.028450
0.446853E-02	0.627650	0.022690
0.815255E-02	0.758790	0.022690
0.175878E-01	0.870390	0.028450
0.251422E-01	0.969220	0.031880

S-4 FLUXES FOR DBSQ CALCULATION		
ANGULAR FLUX	DIRECTION COSINES	POINT WEIGHTS
0.350132E-01	0.333333	0.166667
0.309104E-01	0.333333	0.166667
0.420314E-01	0.881917	0.166667

* * * OUTPUT * * *

ANGULAR DISTRIBUTION,	SINE	FLUX
	-0.	0.244553E-02
	0.24620	0.507138E-02
	0.49236	0.787979E-02
	0.65134	0.925896E-02
	0.77850	0.100997E-01
	0.88762	0.106759E-01
	0.98473	0.113028E-01
	1.00000	0.114835E-01
	0.98473	0.113028E-01
	0.88762	0.106759E-01
	0.77850	0.100997E-01
	0.65134	0.925896E-02
	0.49236	0.787979E-02
	0.24620	0.507138E-02
	0.	0.244553E-02

ANGULAR DISTRIBUTION NORMALIZATION FACTOR = 0.100049E-01

T = 0.0200 MAXIMUM NUMBER OF ITERATIONS EXCEEDED, CONVERGENCE = 0.0000153
 SIGMA = 0.0052440

FOR T= 0.0200 AVERAGE ESCAPE PROBABILITY, PC-ANISOTROPIC = 0.48695E-03
 CONVERGENCE = 0.0000153 SIGMA (DBSQ) = 0.0052440

Z	F1	P2	P3	F4	PC
0.10208	0.35571E-02	0.37278E-01	0.29807E-05	0.16644E-06	0.40839E-01
0.20417	0.74732E-03	0.15204E-01	0.30329E-05	0.16928E-06	0.15955E-01

0.30625	0.30902E-03	0.79167E-02	0.30865E-05	C.17219E-C6	0.82290E-02
0.40833	0.16715E-03	0.46321E-02	0.31414E-05	C.17518E-C6	0.48025E-02
0.51042	0.10438E-03	0.29536E-02	0.31979E-05	C.17824E-C6	0.30613E-02
0.61250	0.71280E-04	0.20030E-02	0.32559E-05	C.18139E-C6	0.20777E-02
0.71458	0.51732E-04	0.14262E-02	0.33155E-05	0.18462E-C6	0.14814E-02
0.81667	0.39241E-04	0.10567E-02	0.33768E-05	0.18795E-C6	0.10995E-02
0.91875	0.30779E-04	0.80992E-03	0.34398E-05	0.19135E-C6	0.84433E-03
1.02083	0.24784E-04	0.63768E-03	0.35045E-05	0.19486E-C6	0.66616E-03
1.12292	0.20384E-04	0.51337E-03	0.35711E-05	0.19846E-C6	0.53752E-03
1.22500	0.17058E-04	0.42112E-03	0.36396E-05	0.20216E-C6	0.44202E-03
1.32708	0.14485E-04	0.35109E-03	0.37101E-05	0.20597E-C6	0.36949E-03
1.42917	0.12452E-04	0.29683E-03	0.37826E-05	C.20988E-C6	0.31328E-03
1.53125	0.10819E-04	0.25397E-03	0.38573E-05	0.21391E-C6	0.26886E-03
1.63333	0.94875E-05	0.21957E-03	0.39343E-05	0.21805E-C6	0.23321E-03
1.73542	0.83873E-05	0.19158E-03	0.40135E-05	0.22232E-C6	0.20420E-03
1.83750	0.74678E-05	0.16853E-03	0.40952E-05	0.22672E-C6	0.18031E-03
1.93958	0.66917E-05	0.14932E-03	0.41794E-05	0.23124E-C6	0.16043E-03
2.04167	0.60305E-05	0.13317E-03	0.42662E-05	0.23590E-C6	0.14371E-03
2.14375	0.54627E-05	0.11947E-03	0.43557E-05	0.24071E-C6	0.12953E-03
2.24583	0.49714E-05	0.10775E-03	0.44481E-05	0.24565E-C6	0.11741E-03
2.34792	0.45435E-05	0.97648E-04	0.45434E-05	0.25076E-C6	0.10699E-03
2.45000	0.41686E-05	0.88886E-04	0.46419E-05	0.25603E-C6	0.97953E-04
2.55208	0.38382E-05	0.81238E-04	0.47435E-05	0.26147E-C6	0.90082E-04
2.65417	0.35456E-05	0.74526E-04	0.48486E-05	0.26708E-C6	0.83187E-04
2.75625	0.32853E-05	0.68603E-04	0.49572E-05	0.27287E-C6	0.77119E-04
2.85833	0.30525E-05	0.63353E-04	0.50694E-05	0.27885E-C6	0.71753E-04
2.96042	0.28437E-05	0.58677E-04	0.51855E-05	0.28502E-C6	0.66991E-04
3.06250	0.26556E-05	0.54495E-04	0.53056E-05	0.29142E-C6	0.62748E-04
3.16458	0.24855E-05	0.50741E-04	0.54300E-05	0.29802E-C6	0.58955E-04
3.26667	0.23313E-05	0.47359E-04	0.55587E-05	0.30486E-C6	0.55554E-04
3.36875	0.21910E-05	0.44301E-04	0.56921E-05	0.31193E-C6	0.52496E-04
3.47083	0.20630E-05	0.41528E-04	0.58303E-05	0.31924E-C6	0.49741E-04
3.57292	0.19459E-05	0.39006E-04	0.59737E-05	0.32683E-C6	0.47252E-04
3.67500	0.18384E-05	0.36705E-04	0.61223E-05	0.33468E-C6	0.45000E-04
3.77708	0.17397E-05	0.34600E-04	0.62766E-05	0.34283E-C6	0.42959E-04
3.87917	0.16487E-05	0.32670E-04	0.64368E-05	0.35127E-C6	0.41107E-04
3.98125	0.15646E-05	0.30896E-04	0.66031E-05	0.36003E-C6	0.39424E-04
4.08333	0.14868E-05	0.29262E-04	0.67760E-05	0.36911E-C6	0.37894E-04
4.18542	0.14147E-05	0.27754E-04	0.69558E-05	0.37855E-C6	0.36503E-04
4.28750	0.13477E-05	0.26359E-04	0.71428E-05	0.38835E-C6	0.35237E-04
4.38958	0.12853E-05	0.25065E-04	0.73374E-05	0.39855E-C6	0.34087E-04
4.49167	0.12272E-05	0.23864E-04	0.75400E-05	0.40914E-C6	0.33041E-04
4.59375	0.11730E-05	0.22747E-04	0.77512E-05	0.42017E-C6	0.32092E-04
4.69583	0.11222E-05	0.21707E-04	0.79713E-05	0.43164E-C6	0.31232E-04

4.79792	0.10747E-05	0.20735E-04	0.82009E-05	0.44360E-06	0.30455E-04
4.90000	0.10301E-05	0.19828E-04	0.84406E-05	0.45605E-06	0.29754E-04
5.00208	0.98826E-06	0.18978E-04	0.86909E-05	0.46904E-06	0.29126E-04
5.10417	0.94891E-06	0.18181E-04	0.89525E-05	0.48258E-06	0.28565E-04
5.20625	0.91186E-06	0.17434E-04	0.92260E-05	0.49674E-06	0.28069E-04
5.30833	0.87694E-06	0.16731E-04	0.95123E-05	0.51152E-06	0.27632E-04
5.41042	0.84398E-06	0.16070E-04	0.98120E-05	0.52696E-06	0.27253E-04
5.51250	0.81286E-06	0.15447E-04	0.10126E-04	0.54312E-06	0.26930E-04
5.61458	0.78342E-06	0.14860E-04	0.10455E-04	0.56003E-06	0.26659E-04
5.71667	0.75555E-06	0.14305E-04	0.10801E-04	0.57776E-06	0.26440E-04
5.81875	0.72914E-06	0.13781E-04	0.11164E-04	0.59632E-06	0.26271E-04
5.92083	0.70410E-06	0.13285E-04	0.11546E-04	0.61580E-06	0.26151E-04
6.02292	0.68032E-06	0.12815E-04	0.11947E-04	0.63625E-06	0.26079E-04
6.12500	0.65773E-06	0.12370E-04	0.12370E-04	0.65773E-06	0.26055E-04

FOR T= 0.0400

AVERAGE ESCAPE PROBABILITY, PC-ANISOTROPIC = 0.12177E-02

CONVERGENCE = 0.

SIGMA (DBSQ) = 0.0065598

Z	P1	P2	P3	P4	PC
0.10208	0.17516E-01	0.69689E-01	0.84490E-05	0.66889E-06	0.87214E-01
0.20417	0.35571E-02	0.33706E-01	0.85986E-05	0.68034E-06	0.37272E-01
0.30625	0.14173E-02	0.19304E-01	0.87522E-05	0.69207E-06	0.20731E-01
0.40833	0.74732E-03	0.12202E-01	0.89099E-05	0.70412E-06	0.12959E-01
0.51042	0.45846E-03	0.81922E-02	0.90719E-05	0.71647E-06	0.86604E-02
0.61250	0.30902E-03	0.57591E-02	0.92384E-05	0.72916E-06	0.60781E-02
0.71458	0.22207E-03	0.42125E-02	0.94095E-05	0.74219E-06	0.44447E-02
0.81667	0.16715E-03	0.31791E-02	0.95854E-05	0.75557E-06	0.33566E-02
0.91875	0.13030E-03	0.24599E-02	0.97662E-05	0.76930E-06	0.26007E-02
1.02083	0.10438E-03	0.19477E-02	0.99522E-05	0.78344E-06	0.20629E-02
1.12292	0.85484E-04	0.15725E-02	0.10144E-04	0.79795E-06	0.16689E-02
1.22500	0.71280E-04	0.12916E-02	0.10340E-04	0.81286E-06	0.13740E-02
1.32708	0.60339E-04	0.10774E-02	0.10543E-04	0.82821E-06	0.11491E-02
1.42917	0.51732E-04	0.91037E-03	0.10752E-04	0.84400E-06	0.97370E-03
1.53125	0.44842E-04	0.77798E-03	0.10967E-04	0.86025E-06	0.83465E-03
1.63333	0.39241E-04	0.67153E-03	0.11188E-04	0.87694E-06	0.72283E-03
1.73542	0.34626E-04	0.58483E-03	0.11416E-04	0.89416E-06	0.63177E-03
1.83750	0.30779E-04	0.51340E-03	0.11652E-04	0.91187E-06	0.55674E-03
1.93958	0.27539E-04	0.45408E-03	0.11894E-04	0.93012E-06	0.49444E-03
2.04167	0.24784E-04	0.40421E-03	0.12144E-04	0.94892E-06	0.44209E-03
2.14375	0.22423E-04	0.36192E-03	0.12403E-04	0.96830E-06	0.39771E-03
2.24583	0.20384E-04	0.32578E-03	0.12669E-04	0.98826E-06	0.35982E-03
2.34792	0.18610E-04	0.29468E-03	0.12944E-04	0.10089E-05	0.32725E-03

2.45000	0.17058E-04	0.26774E-03	0.13228E-04	0.10301E-05	0.29906E-03
2.55208	0.15693E-04	0.24426E-03	0.13522E-04	0.10521E-05	0.27453E-03
2.65417	0.14485E-04	0.22368E-03	0.13825E-04	0.10747E-05	0.25306E-03
2.75625	0.13411E-04	0.20555E-03	0.14139E-04	0.10981E-05	0.23419E-03
2.85833	0.12452E-04	0.18949E-03	0.14463E-04	0.11222E-05	0.21753E-03
2.96042	0.11593E-04	0.17522E-03	0.14799E-04	0.11472E-05	0.20276E-03
3.06250	0.10819E-04	0.16247E-03	0.15147E-04	0.11730E-05	0.18961E-03
3.16458	0.10121E-04	0.15105E-03	0.15507E-04	0.11996E-05	0.17787E-03
3.26667	0.94875E-05	0.14077E-03	0.15879E-04	0.12272E-05	0.16736E-03
3.36875	0.89120E-05	0.13149E-03	0.16266E-04	0.12558E-05	0.15792E-03
3.47083	0.83873E-05	0.12309E-03	0.16666E-04	0.12854E-05	0.14943E-03
3.57292	0.79075E-05	0.11545E-03	0.17082E-04	0.13160E-05	0.14176E-03
3.67500	0.74678E-05	0.10850E-03	0.17513E-04	0.13477E-05	0.13483E-03
3.77708	0.70638E-05	0.10215E-03	0.17961E-04	0.13806E-05	0.12856E-03
3.87917	0.66917E-05	0.96336E-04	0.18426E-04	0.14147E-05	0.12287E-03
3.98125	0.63482E-05	0.90998E-04	0.18909E-04	0.14501E-05	0.11771E-03
4.08333	0.60305E-05	0.86087E-04	0.19412E-04	0.14868E-05	0.11302E-03
4.18542	0.57361E-05	0.81559E-04	0.19935E-04	0.15250E-05	0.10875E-03
4.28750	0.54627E-05	0.77376E-04	0.20479E-04	0.15646E-05	0.10488E-03
4.38958	0.52084E-05	0.73503E-04	0.21045E-04	0.16058E-05	0.10136E-03
4.49167	0.49714E-05	0.69911E-04	0.21636E-04	0.16487E-05	0.98167E-04
4.59375	0.47502E-05	0.66574E-04	0.22251E-04	0.16933E-05	0.95269E-04
4.69583	0.45435E-05	0.63468E-04	0.22893E-04	0.17397E-05	0.92644E-04
4.79792	0.43500E-05	0.60573E-04	0.23563E-04	0.17881E-05	0.90274E-04
4.90000	0.41686E-05	0.57870E-04	0.24262E-04	0.18385E-05	0.88139E-04
5.00208	0.39983E-05	0.55342E-04	0.24993E-04	0.18911E-05	0.86225E-04
5.10417	0.38382E-05	0.52976E-04	0.25758E-04	0.19459E-05	0.84518E-04
5.20625	0.36876E-05	0.50756E-04	0.26558E-04	0.20032E-05	0.83005E-04
5.30833	0.35456E-05	0.48673E-04	0.27396E-04	0.20630E-05	0.81677E-04
5.41042	0.34117E-05	0.46714E-04	0.28273E-04	0.21256E-05	0.80525E-04
5.51250	0.32853E-05	0.44870E-04	0.29194E-04	0.21910E-05	0.79540E-04
5.61458	0.31657E-05	0.43133E-04	0.30160E-04	0.22595E-05	0.78718E-04
5.71667	0.30525E-05	0.41494E-04	0.31174E-04	0.23313E-05	0.78051E-04
5.81875	0.29454E-05	0.39946E-04	0.32240E-04	0.24066E-05	0.77538E-04
5.92083	0.28437E-05	0.38482E-04	0.33361E-04	0.24856E-05	0.77173E-04
6.02292	0.27472E-05	0.37098E-04	0.34542E-04	0.25685E-05	0.76955E-04
6.12500	0.26556E-05	0.35786E-04	0.35786E-04	0.26556E-05	0.76882E-04

FOR T= 0.0600

AVERAGE ESCAPE PROBABILITY, PC-ANISOTROPIC = 0.20584E-02

CONVERGENCE = 0.000072

SIGMA (DBSQ) = 0.0073959

Z

P1

P2

P3

P4

PC

0.10208	0.42622E-01	0.89637E-01	0.15554E-04	0.15121E-05	0.13228E-00
0.20417	0.90481E-02	0.50729E-01	0.15831E-04	0.15380E-05	0.59794E-01
0.30625	0.35571E-02	0.30988E-01	0.16117E-04	0.15646E-05	0.34563E-01
0.40833	0.18478E-02	0.20446E-01	0.16410E-04	0.15919E-05	0.22312E-01
0.51042	0.11196E-02	0.14251E-01	0.16710E-04	0.16199E-05	0.15389E-01
0.61250	0.74732E-03	0.10316E-01	0.17020E-04	0.16487E-05	0.11082E-01
0.71458	0.53285E-03	0.76835E-02	0.17338E-04	0.16782E-05	0.82354E-02
0.81667	0.39852E-03	0.58883E-02	0.17664E-04	0.17086E-05	0.63062E-02
0.91875	0.30902E-03	0.46183E-02	0.18001E-04	0.17397E-05	0.49470E-02
1.02083	0.24648E-03	0.36909E-02	0.18346E-04	0.17717E-05	0.39575E-02
1.12292	0.20110E-03	0.29982E-02	0.18702E-04	0.18046E-05	0.32198E-02
1.22500	0.16715E-03	0.24732E-02	0.19069E-04	0.18385E-05	0.26612E-02
1.32708	0.14110E-03	0.20673E-02	0.19446E-04	0.18733E-05	0.22297E-02
1.42917	0.12068E-03	0.17489E-02	0.19834E-04	0.19091E-05	0.18913E-02
1.53125	0.10438E-03	0.14956E-02	0.20234E-04	0.19459E-05	0.16222E-02
1.63333	0.91171E-04	0.12919E-02	0.20646E-04	0.19838E-05	0.14057E-02
1.73542	0.80312E-04	0.11254E-02	0.21071E-04	0.20228E-05	0.12288E-02
1.83750	0.71280E-04	0.98793E-03	0.21510E-04	0.20630E-05	0.10828E-02
1.93958	0.63688E-04	0.87326E-03	0.21962E-04	0.21044E-05	0.96101E-03
2.04167	0.57246E-04	0.77675E-03	0.22428E-04	0.21471E-05	0.85858E-03
2.14375	0.51732E-04	0.69488E-03	0.22909E-04	0.21910E-05	0.77171E-03
2.24583	0.46978E-04	0.62489E-03	0.23406E-04	0.22363E-05	0.69751E-03
2.34792	0.42849E-04	0.56478E-03	0.23919E-04	0.22831E-05	0.63383E-03
2.45000	0.39241E-04	0.51272E-03	0.24449E-04	0.23313E-05	0.57874E-03
2.55208	0.36069E-04	0.46735E-03	0.24997E-04	0.23811E-05	0.53079E-03
2.65417	0.33267E-04	0.42758E-03	0.25563E-04	0.24325E-05	0.48885E-03
2.75625	0.30779E-04	0.39257E-03	0.26149E-04	0.24856E-05	0.45198E-03
2.85833	0.28560E-04	0.36159E-03	0.26755E-04	0.25404E-05	0.41944E-03
2.96042	0.26572E-04	0.33405E-03	0.27382E-04	0.25970E-05	0.39060E-03
3.06250	0.24784E-04	0.30948E-03	0.28032E-04	0.26556E-05	0.36495E-03
3.16458	0.23171E-04	0.28747E-03	0.28705E-04	0.27162E-05	0.34206E-03
3.26667	0.21711E-04	0.26768E-03	0.29402E-04	0.27789E-05	0.32157E-03
3.36875	0.20384E-04	0.24982E-03	0.30125E-04	0.28437E-05	0.30318E-03
3.47083	0.19175E-04	0.23367E-03	0.30874E-04	0.29109E-05	0.28663E-03
3.57292	0.18070E-04	0.21900E-03	0.31652E-04	0.29805E-05	0.27170E-03
3.67500	0.17058E-04	0.20565E-03	0.32460E-04	0.30526E-05	0.25822E-03
3.77708	0.16129E-04	0.19346E-03	0.33298E-04	0.31273E-05	0.24601E-03
3.87917	0.15274E-04	0.18231E-03	0.34170E-04	0.32049E-05	0.23495E-03
3.98125	0.14485E-04	0.17208E-03	0.35075E-04	0.32853E-05	0.22492E-03
4.08333	0.13755E-04	0.16267E-03	0.36017E-04	0.33688E-05	0.21581E-03
4.18542	0.13079E-04	0.15401E-03	0.36998E-04	0.34555E-05	0.20754E-03
4.28750	0.12452E-04	0.14601E-03	0.38019E-04	0.35457E-05	0.20002E-03
4.38958	0.11869E-04	0.13860E-03	0.39082E-04	0.36394E-05	0.19320E-03
4.49167	0.11326E-04	0.13175E-03	0.40190E-04	0.37368E-05	0.18700E-03

4.59375	0.10819E-04	0.12538E-03	0.41346E-04	0.38383E-05	0.18138E-03
4.69583	0.10346E-04	0.11945E-03	0.42552E-04	0.39439E-05	0.17629E-03
4.79792	0.99027E-05	0.11393E-03	0.43811E-04	0.40539E-05	0.17170E-03
4.90000	0.94875E-05	0.10879E-03	0.45127E-04	0.41687E-05	0.16757E-03
5.00208	0.90978E-05	0.10397E-03	0.46503E-04	0.42883E-05	0.16386E-03
5.10417	0.87317E-05	0.99471E-04	0.47941E-04	0.44131E-05	0.16056E-03
5.20625	0.83873E-05	0.95252E-04	0.49448E-04	0.45436E-05	0.15763E-03
5.30833	0.80628E-05	0.91293E-04	0.51026E-04	0.46798E-05	0.15506E-03
5.41042	0.77568E-05	0.87573E-04	0.52680E-04	0.48223E-05	0.15283E-03
5.51250	0.74678E-05	0.84074E-04	0.54415E-04	0.49714E-05	0.15093E-03
5.61458	0.71948E-05	0.80778E-04	0.56236E-04	0.51276E-05	0.14934E-03
5.71667	0.69364E-05	0.77671E-04	0.58151E-04	0.52911E-05	0.14805E-03
5.81875	0.66917E-05	0.74738E-04	0.60163E-04	0.54627E-05	0.14706E-03
5.92083	0.64597E-05	0.71967E-04	0.62281E-04	0.56427E-05	0.14635E-03
6.02292	0.62396E-05	0.69346E-04	0.64512E-04	0.58318E-05	0.14593E-03
6.12500	0.60305E-05	0.66864E-04	0.66864E-04	0.60305E-05	0.14579E-03

FOR T= 0.0800

AVERAGE ESCAPE PROBABILITY, PC-ANISOTROPIC = 0.29626E-02

CONVERGENCE = 0.

SIGMA (DBSQ) = 0.0079879

Z	P1	P2	P3	P4	PC
0.10208	0.74629E-01	0.98459E-01	0.23993E-04	0.27009E-05	0.17311E-00
0.20417	0.17516E-01	0.64953E-01	0.24425E-04	0.27473E-05	0.82496E-01
0.30625	0.68934E-02	0.42161E-01	0.24868E-04	0.27949E-05	0.49082E-01
0.40833	0.35571E-02	0.28799E-01	0.25323E-04	0.28437E-05	0.32384E-01
0.51042	0.21385E-02	0.20578E-01	0.25790E-04	0.28939E-05	0.22746E-01
0.61250	0.14173E-02	0.15238E-01	0.26271E-04	0.29454E-05	0.16685E-01
0.71458	0.10044E-02	0.11579E-01	0.26765E-04	0.29983E-05	0.12613E-01
0.81667	0.74732E-03	0.89761E-02	0.27273E-04	0.30526E-05	0.97537E-02
0.91875	0.57695E-03	0.71065E-02	0.27796E-04	0.31084E-05	0.77144E-02
1.02083	0.45846E-03	0.57298E-02	0.28334E-04	0.31657E-05	0.62198E-02
1.12292	0.37284E-03	0.46929E-02	0.28888E-04	0.32247E-05	0.50978E-02
1.22500	0.30902E-03	0.38894E-02	0.29458E-04	0.32853E-05	0.42311E-02
1.32708	0.26022E-03	0.32635E-02	0.30044E-04	0.33476E-05	0.35571E-02
1.42917	0.22207E-03	0.27688E-02	0.30649E-04	0.34118E-05	0.30249E-02
1.53125	0.19171E-03	0.23721E-02	0.31272E-04	0.34777E-05	0.25986E-02
1.63333	0.16715E-03	0.20505E-02	0.31914E-04	0.35457E-05	0.22531E-02
1.73542	0.14701E-03	0.17869E-02	0.32575E-04	0.36156E-05	0.19702E-02
1.83750	0.13030E-03	0.15695E-02	0.33258E-04	0.36876E-05	0.17367E-02
1.93958	0.11627E-03	0.13879E-02	0.33962E-04	0.37618E-05	0.15419E-02
2.04167	0.10438E-03	0.12348E-02	0.34688E-04	0.38383E-05	0.13777E-02
2.14375	0.94229E-04	0.11046E-02	0.35438E-04	0.39171E-05	0.12382E-02

2.24583	0.85484E-04	0.99327E-03	0.36213E-C4	C.39983E-C5	0.11190E-02
2.34792	0.77900E-04	0.89733E-03	C.37C13E-C4	C.40822E-C5	0.10163E-02
2.45000	0.71280E-04	0.81416E-03	C.37840E-C4	C.41687E-C5	0.92745E-03
2.55208	0.65469E-04	0.74165E-03	C.38694E-04	0.42579E-C5	0.85008E-03
2.65417	0.60339E-04	0.67816E-03	0.39578E-04	0.43500E-05	0.78243E-03
2.75625	0.55788E-04	0.62234E-03	0.40492E-04	0.44452E-05	0.72306E-03
2.85833	0.51732E-04	0.57294E-03	0.41438E-04	0.45436E-05	0.67065E-03
2.96042	0.48103E-04	0.52903E-03	0.42417E-04	0.46452E-05	0.62420E-03
3.06250	0.44842E-04	0.48985E-03	0.43431E-04	C.47503E-C5	0.58287E-03
3.16458	0.41902E-04	0.45476E-03	C.44482E-04	C.48590E-C5	0.54600E-03
3.26667	0.39241E-04	0.42322E-C3	C.45571E-C4	0.49714E-C5	0.51300E-03
3.36875	0.36825E-04	0.39477E-C3	0.46701E-C4	0.50879E-05	0.48339E-03
3.47083	0.34626E-04	0.36904E-03	0.47873E-C4	C.52084E-05	0.45675E-03
3.57292	0.32618E-04	0.34569E-03	C.49089E-04	0.53333E-C5	0.43273E-03
3.67500	0.30779E-04	0.32444E-03	C.50352E-C4	0.54627E-C5	0.41103E-03
3.77708	0.29092E-04	0.30505E-03	C.51663E-C4	0.55969E-05	0.39140E-03
3.87917	0.27539E-04	0.28731E-03	0.53027E-04	C.57361E-C5	0.37361E-03
3.98125	0.26107E-04	0.27105E-03	C.54444E-C4	0.58806E-05	0.35748E-03
4.08333	0.24784E-04	0.25610E-03	C.55919E-C4	0.60305E-05	0.34284E-03
4.18542	0.23560E-04	0.24234E-03	C.57454E-04	C.61863E-C5	0.32954E-03
4.28750	0.22423E-04	0.22964E-03	C.59053E-04	0.63482E-05	0.31746E-03
4.38958	0.21367E-04	0.21789E-03	C.60719E-04	0.65166E-05	0.30649E-03
4.49167	0.20384E-04	0.20701E-03	0.62456E-04	0.66917E-05	0.29654E-03
4.59375	0.19467E-04	0.19691E-03	0.64268E-04	C.68740E-C5	0.28752E-03
4.69583	0.18610E-04	0.18752E-03	0.66160E-04	0.70638E-05	0.27935E-03
4.79792	0.17809E-04	0.17878E-03	0.68135E-04	0.72617E-05	0.27198E-03
4.90000	0.17058E-04	0.17062E-03	0.70200E-04	C.74679E-05	0.26535E-03
5.00208	0.16354E-04	0.16300E-03	C.72358E-04	0.76830E-05	0.25940E-03
5.10417	0.15693E-04	0.15588E-03	0.74618E-04	C.79076E-05	0.25410E-03
5.20625	0.15071E-04	0.14921E-03	C.76983E-C4	C.81421E-C5	0.24940E-03
5.30833	0.14485E-04	0.14295E-03	0.79462E-04	0.83873E-05	0.24528E-03
5.41042	0.13932E-04	0.13707E-03	C.82062E-C4	C.86437E-05	0.24170E-03
5.51250	0.13411E-04	0.13154E-03	0.84790E-04	C.89120E-05	0.23865E-03
5.61458	0.12918E-04	0.12633E-03	C.87655E-04	0.91930E-05	0.23610E-03
5.71667	0.12452E-04	0.12143E-03	0.90666E-04	0.94875E-05	0.23404E-03
5.81875	0.12011E-04	0.11680E-03	C.93833E-C4	0.97964E-05	C.23244E-03
5.92083	0.11593E-04	0.11243E-03	C.97168E-C4	0.10121E-04	0.23131E-03
6.02292	0.11196E-04	0.10830E-03	0.10068E-03	0.10461E-04	0.23064E-03
6.12500	0.10819E-04	0.10439E-03	0.10439E-03	C.10819E-04	0.23041E-03

FOR T= 0.1000

AVERAGE ESCAPE PROBABILITY, PC-ANISOTROPIC = C.39065E-02

CONVERGENCE = 0.000038

SIGMA (DBSQ) = C.0084309

Z	P1	P2	P3	P4	PC
0.102C8	0.10821E-00	0.99921E-C1	0.33595E-C4	0.42398E-C5	0.20817E-00
0.20417	0.28843E-01	0.76307E-C1	0.34202E-C4	0.43128E-C5	0.10519E-00
0.30625	0.11533E-01	0.52330E-C1	0.34826E-C4	0.43878E-C5	0.63902E-01
0.40833	0.59411E-02	0.36962E-C1	0.35467E-C4	0.44646E-C5	0.42943E-01
0.51042	0.35571E-02	0.26973E-01	0.36126E-C4	0.45436E-C5	0.30571E-01
0.61250	0.23462E-02	0.20301E-01	0.368C4E-C4	0.46246E-C5	0.22688E-01
0.71458	0.16551E-02	0.15673E-C1	0.375C0E-C4	0.47078E-C5	0.17370E-01
0.81667	0.12265E-02	0.12322E-01	0.38217E-C4	0.47933E-C5	0.13592E-01
0.91875	0.94346E-03	0.98381E-02	0.38954E-C4	0.48811E-C5	0.10825E-01
1.02083	0.74732E-03	0.79833E-C2	0.39712E-C4	0.49714E-C5	0.87753E-02
1.12292	0.6C605E-03	0.65758E-C2	0.40493E-C4	0.50642E-C5	0.72274E-02
1.22500	0.50108E-03	0.54856E-02	0.41296E-C4	0.51597E-C5	0.60332E-02
1.32708	0.42101E-03	0.46235E-C2	0.42124E-C4	0.52578E-C5	0.50919E-02
1.42917	0.35860E-03	0.39343E-C2	0.42977E-C4	0.53588E-C5	0.43412E-02
1.53125	0.3C902E-03	0.33791E-C2	0.43856E-C4	0.54627E-C5	0.37375E-02
1.63333	0.26901E-03	0.29267E-02	0.44762E-C4	0.55696E-C5	0.32461E-02
1.73542	0.23626E-03	0.25541E-C2	0.45696E-C4	0.56798E-C5	0.28417E-02
1.83750	0.20912E-03	0.22446E-C2	0.46659E-C4	0.57933E-C5	0.25062E-02
1.93958	0.18638E-03	0.19852E-C2	0.47654E-C4	0.59101E-C5	0.22252E-02
2.04167	0.16715E-03	0.17667E-02	0.48680E-C4	0.60305E-C5	0.19886E-02
2.14375	0.15074E-03	0.15811E-C2	0.49739E-C4	0.61547E-C5	0.17877E-02
2.24583	0.13662E-03	0.1422C0E-C2	0.50833E-C4	0.62827E-C5	0.16157E-02
2.34792	0.12440E-03	0.12848E-C2	0.51964E-C4	0.64148E-C5	0.14675E-02
2.450C0	0.11373E-03	0.11657E-C2	0.53132E-C4	0.65510E-C5	0.13391E-02
2.55208	0.10438E-03	0.10618E-C2	0.54340E-C4	0.66917E-C5	0.12272E-02
2.65417	0.96138E-04	0.97069E-C3	0.55589E-C4	0.68370E-C5	0.11293E-02
2.75625	0.88830E-04	0.89041E-C3	0.56882E-C4	0.69870E-C5	0.10431E-02
2.85833	0.82323E-04	0.81936E-03	0.58219E-C4	0.71420E-C5	0.96704E-03
2.96042	0.76505E-04	0.75622E-C3	0.59605E-C4	0.73022E-C5	0.89963E-03
3.06250	0.7128C0E-04	0.70001E-C3	0.61040E-C4	0.74679E-C5	0.83979E-03
3.16458	0.66573E-04	0.64965E-C3	0.62527E-C4	0.76392E-C5	0.78638E-03
3.26667	0.62316E-04	0.60437E-03	0.64C68E-C4	0.78166E-C5	0.73858E-03
3.36875	0.58454E-04	0.56355E-C3	0.65667E-C4	0.80002E-C5	0.69567E-03
3.47083	0.54939E-04	0.52661E-03	0.67326E-C4	0.81903E-C5	0.65706E-03
3.57292	0.51732E-04	0.49309E-03	0.69048E-C4	0.83873E-C5	0.62226E-03
3.67500	0.48798E-04	0.46261E-C3	0.70837E-C4	0.85915E-C5	0.59083E-03
3.77708	0.46106E-04	0.43479E-03	0.72696E-C4	0.88032E-C5	0.56240E-03
3.87917	0.43630E-04	0.40936E-03	0.74628E-C4	0.90228E-C5	0.53664E-03
3.98125	0.41348E-04	0.38604E-03	0.76637E-C4	0.92508E-C5	0.51328E-03
4.08333	0.39241E-04	0.36462E-03	0.78728E-C4	0.94875E-C5	0.49208E-03
4.18542	0.37290E-04	0.34489E-03	0.8C905E-C4	0.97335E-C5	0.47282E-03
4.28750	0.35481E-04	0.32669E-C3	0.83173E-C4	0.99891E-C5	0.45534E-03

4.38958	0.33801E-04	0.30987E-03	0.85536E-04	0.10255E-04	0.43946E-03
4.49167	0.32237E-04	0.29429E-03	0.88001E-04	0.10531E-04	0.42506E-03
4.59375	0.30779E-04	0.27983E-03	0.90573E-04	0.10819E-04	0.41200E-03
4.69583	0.29418E-04	0.26639E-03	0.93258E-04	0.11119E-04	0.40019E-03
4.79792	0.28145E-04	0.25388E-03	0.96063E-04	0.11432E-04	0.38952E-03
4.90000	0.26953E-04	0.24222E-03	0.98995E-04	0.11757E-04	0.37992E-03
5.00208	0.25834E-04	0.23133E-03	0.10206E-03	0.12097E-04	0.37132E-03
5.10417	0.24784E-04	0.22114E-03	0.10527E-03	0.12452E-04	0.36365E-03
5.20625	0.23797E-04	0.21160E-03	0.10864E-03	0.12823E-04	0.35688E-03
5.30833	0.22868E-04	0.20266E-03	0.11216E-03	0.13211E-04	0.35090E-03
5.41042	0.21991E-04	0.19426E-03	0.11586E-03	0.13616E-04	0.34573E-03
5.51250	0.21165E-04	0.18637E-03	0.11974E-03	0.14040E-04	0.34131E-03
5.61458	0.20384E-04	0.17894E-03	0.12381E-03	0.14485E-04	0.33763E-03
5.71667	0.19645E-04	0.17195E-03	0.12810E-03	0.14951E-04	0.33464E-03
5.81875	0.18946E-04	0.16534E-03	0.13261E-03	0.15439E-04	0.33234E-03
5.92083	0.18283E-04	0.15911E-03	0.13736E-03	0.15952E-04	0.33070E-03
6.02292	0.17655E-04	0.15322E-03	0.14236E-03	0.16492E-04	0.32973E-03
6.12500	0.17058E-04	0.14764E-03	0.14764E-03	0.17058E-04	0.32940E-03

END OF FILE A2 ,JOB COMPLETED.

R H NCRMAN- 01470017911-99 L55 4 TAPES 01-44-44
**PRINTING= 1209 LINES **PUNCHING= 114 CARDS **CRT=

08/15/64
0-35MM 0-9INCH